

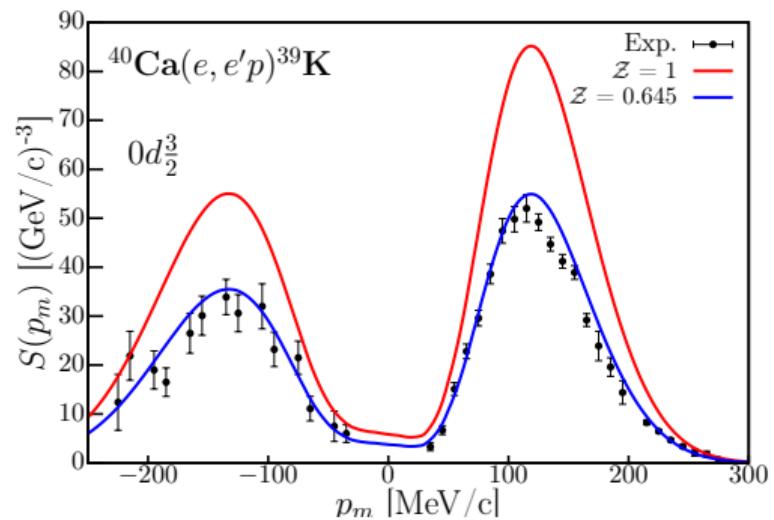
Quenching Spectroscopic Factors

Mack C. Atkinson

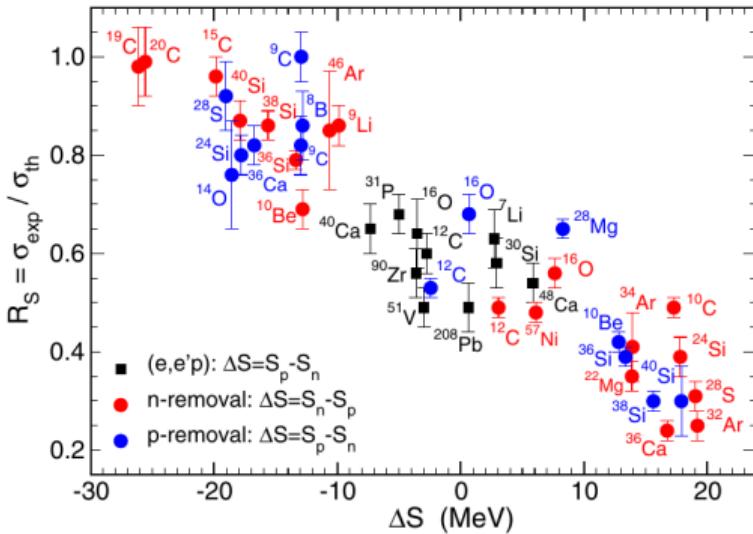
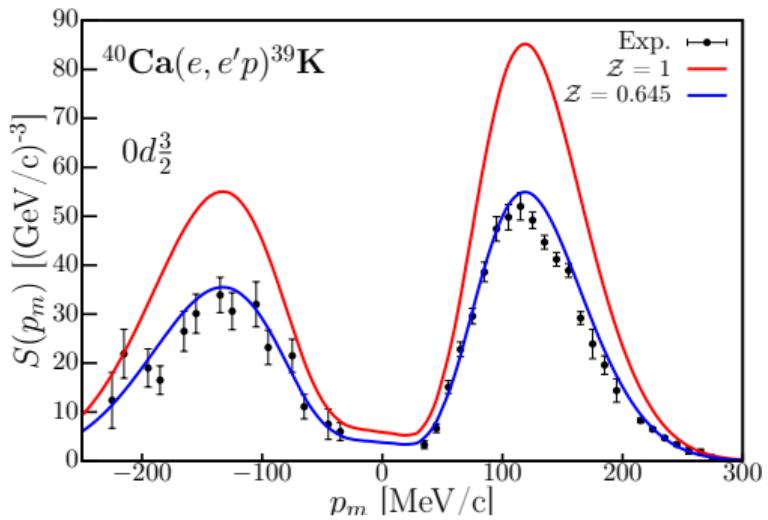
TRIUMF

Nuclear Physics at the Edge of Stability

Motivation

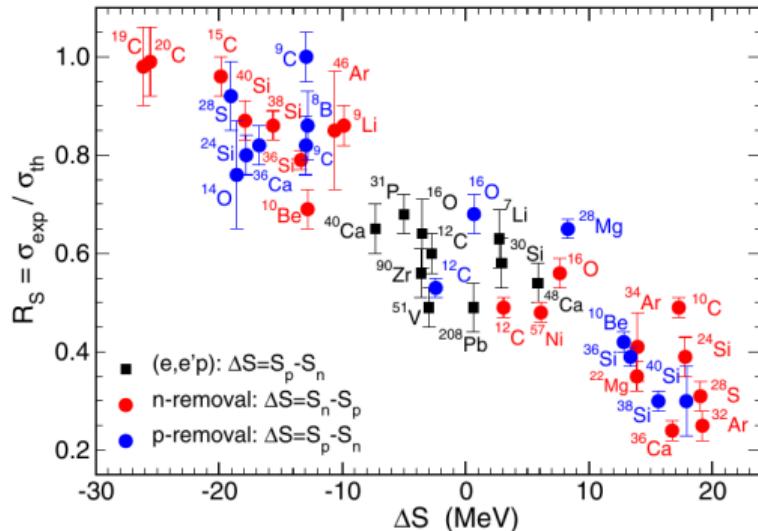
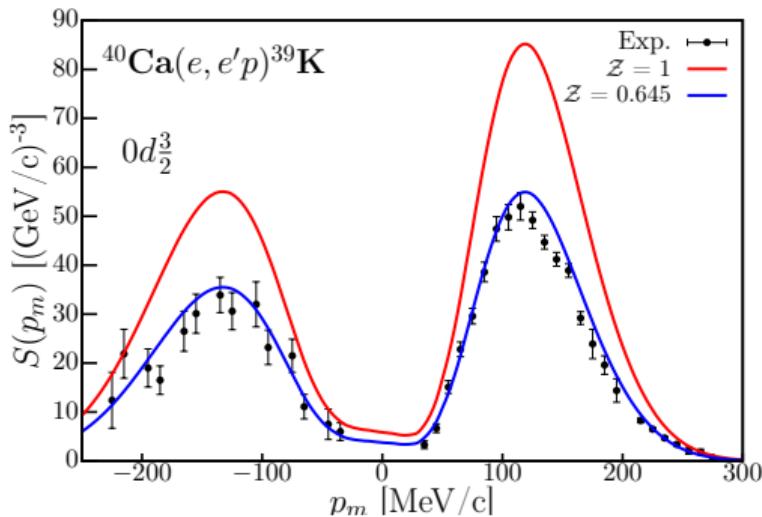


Motivation



Motivation

- Why do spectroscopic factors decrease with increasing neutron asymmetry?



I: The dispersive optical model

Single-Particle Propagator and the Spectral Function

$$G_{\ell j}(r, r'; E) = \sum_m \frac{\langle \Psi_0^A | a_{r\ell j} | \Psi_m^{A+1} \rangle \langle \Psi_m^{A+1} | a_{r'\ell j}^\dagger | \Psi_0^A \rangle}{E - (E_m^{A+1} - E_0^A) + i\eta} \\ + \sum_n \frac{\langle \Psi_0^A | a_{r'\ell j}^\dagger | \Psi_n^{A-1} \rangle \langle \Psi_n^{A-1} | a_{r\ell j} | \Psi_0^A \rangle}{E - (E_0^A - E_n^{A-1}) - i\eta}$$

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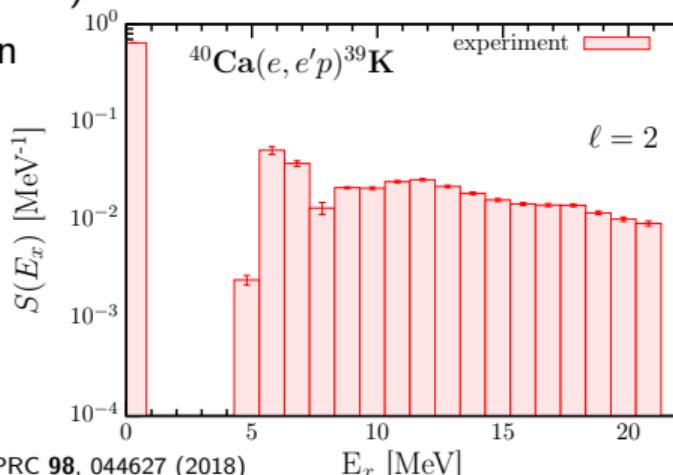
$$S_{\ell j}^h(r; E) = \frac{1}{\pi} \text{Im} G_{\ell j}(r, r; E) \theta(E - (E_0^A - E_0^{A-1}))$$

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M.C. Atkinson *et al.*, PRC 98, 044627 (2018)

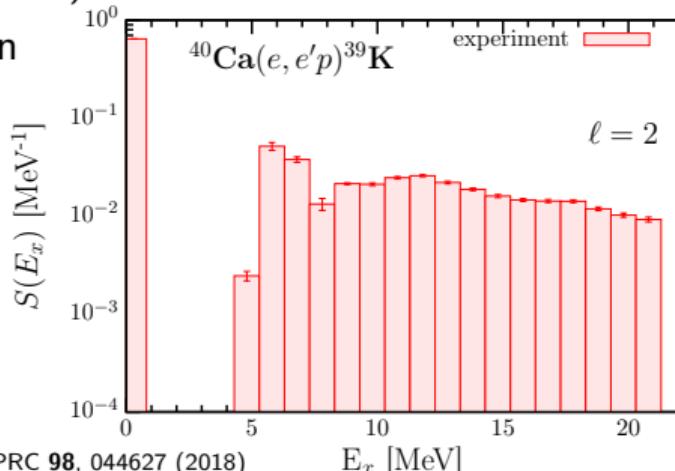
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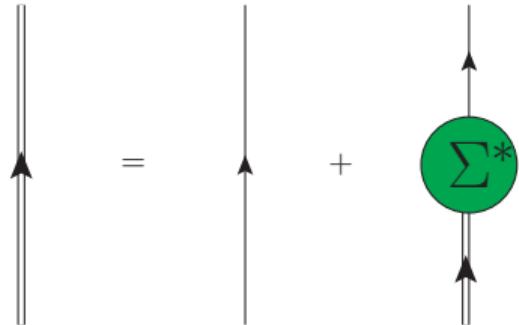
- Close connection with experimental observables



M.C. Atkinson *et al.*, PRC 98, 044627 (2018)

Irreducible self-energy and the Dyson equation

- Perturbative expansion of G leads to the Dyson equation



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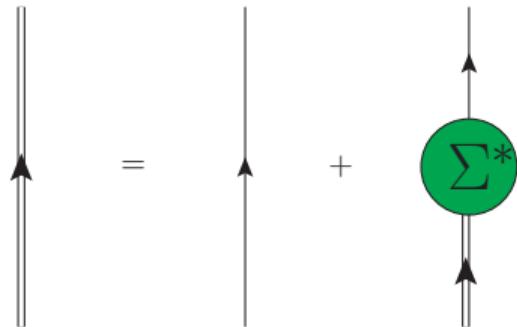
$$\Sigma^*(\mathbf{r}, \mathbf{r}'; E) \rightarrow V_{vol}(\mathbf{r}, \mathbf{r}'; E) + V_{sur}(\mathbf{r}, \mathbf{r}'; E) + V_{so}(\mathbf{r}, \mathbf{r}'; E)$$



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- $\Sigma^*(\mathbf{r}, \mathbf{r}'; E)$ is **nonlocal**, parametrized with β

$$\Sigma^*(\mathbf{r}, \mathbf{r}'; E) = \Sigma^* \left(\frac{\mathbf{r} + \mathbf{r}'}{2}; E \right) e^{\frac{-(\mathbf{r}-\mathbf{r}')^2}{\beta^2}} \pi^{-\frac{3}{2}} \beta^{-3}$$

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Can this also describe negative energy
observables?

Dispersive Optical Model

- The DOM makes use of complex analysis to formulate a consistent self-energy

Dispersive Correction

$$\begin{aligned} Re\Sigma_{\ell j}(r, r'; E) = & Re\Sigma_{\ell j}(r, r'; \epsilon_F) - \frac{1}{\pi}(\epsilon_F - E)\mathcal{P} \int_{\epsilon_T^+}^{\infty} dE' Im\Sigma_{\ell j}(r, r'; E') \left[\frac{1}{E - E'} - \frac{1}{\epsilon_F - E'} \right] \\ & + \frac{1}{\pi}(\epsilon_F - E)\mathcal{P} \int_{-\infty}^{\epsilon_T^-} dE' Im\Sigma_{\ell j}(r, r'; E') \left[\frac{1}{E - E'} - \frac{1}{\epsilon_F - E'} \right] \end{aligned}$$

¹ C. Mahaux, R. Sartor, *Adv. Nucl. Phys.*, **20**, 96 (1991)

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- Dispersion relation constrains self-energy at all energies
- This constraint ensures bound and scattering quantities are simultaneously described

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Schematic of the Fitting Process

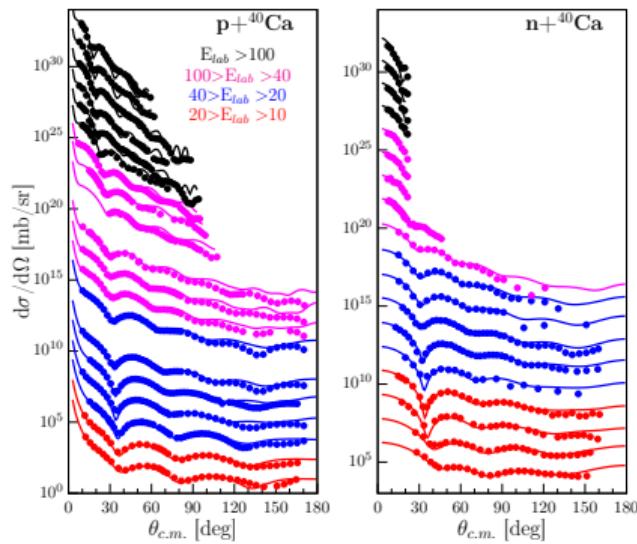


Fitting the Self-energy (^{40}Ca)

- Parameters of self-energy varied to minimize χ^2

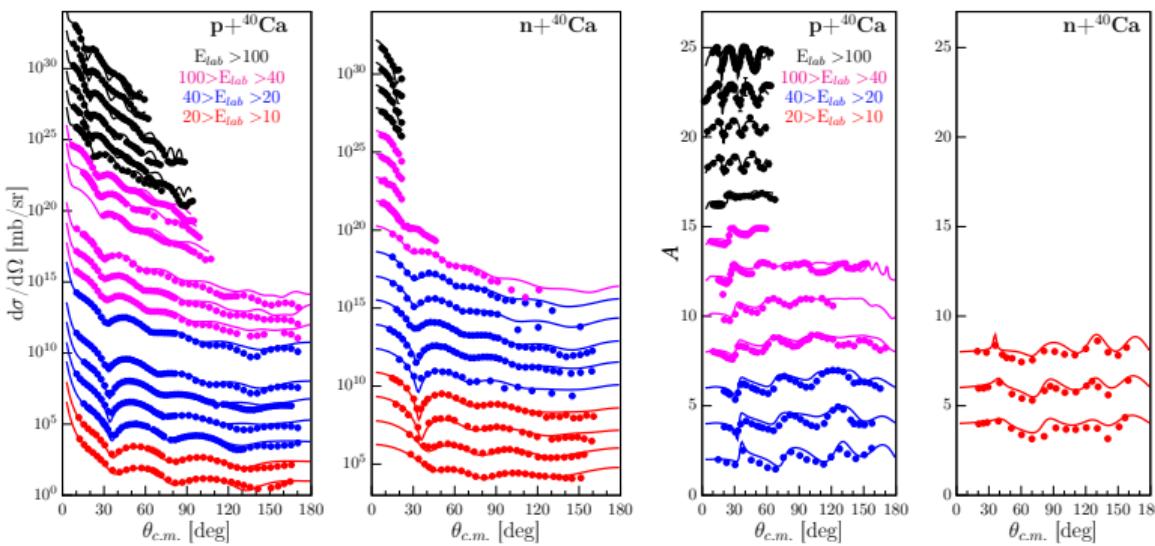
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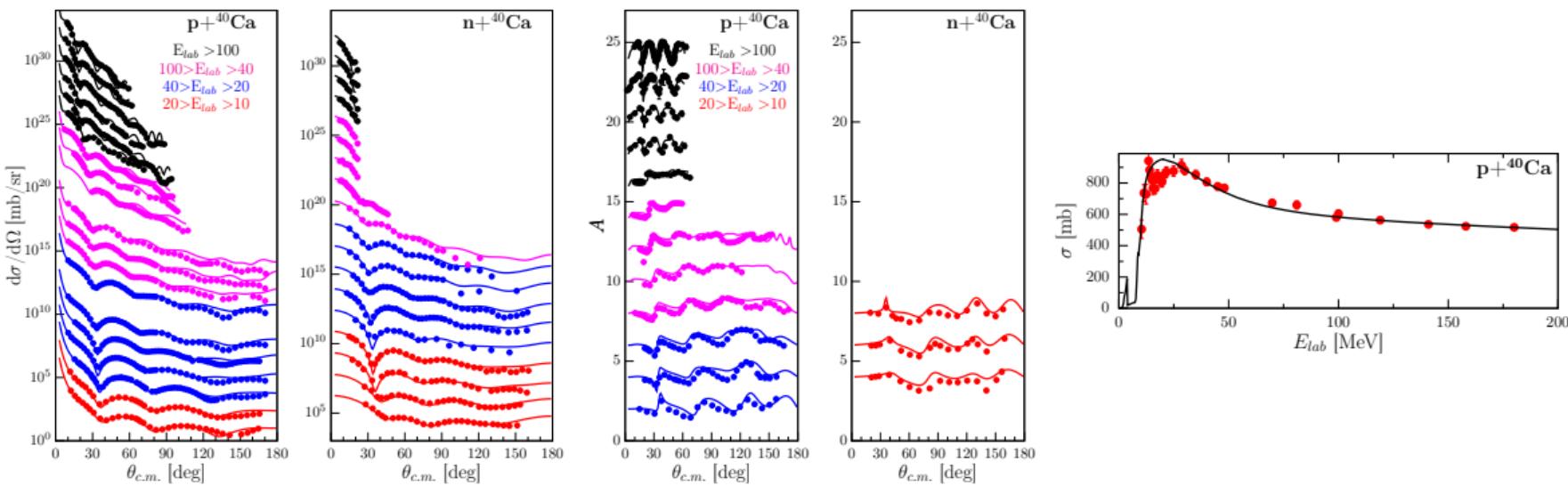
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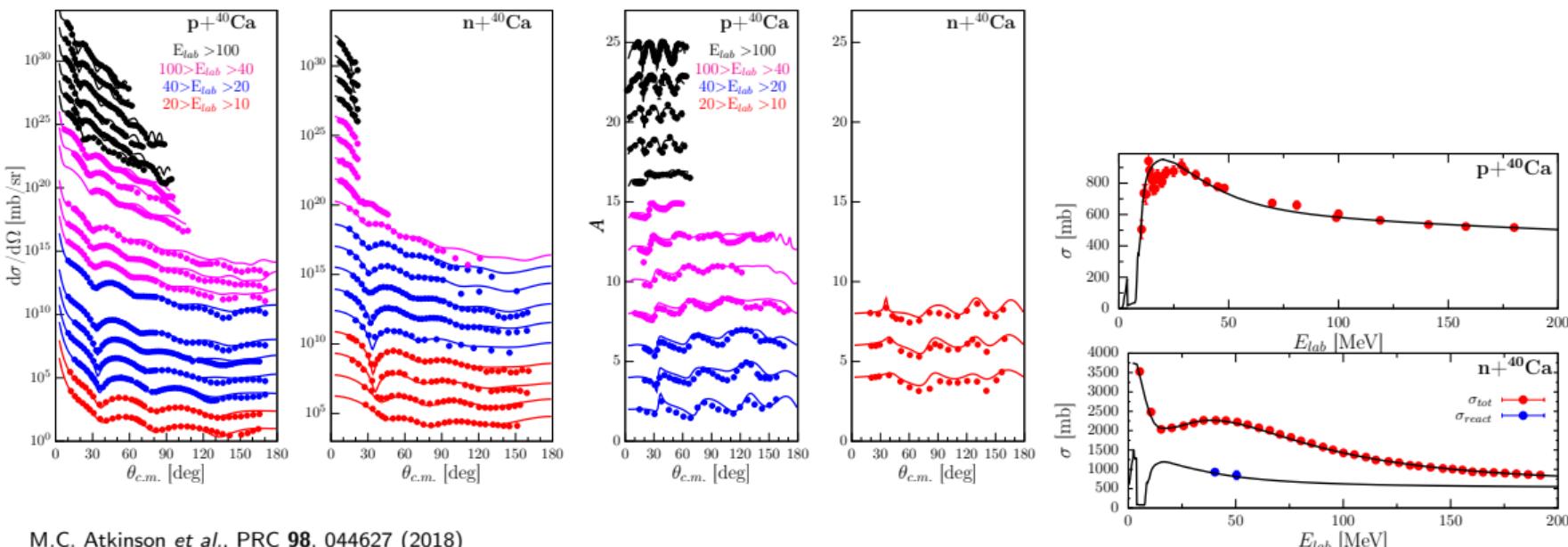
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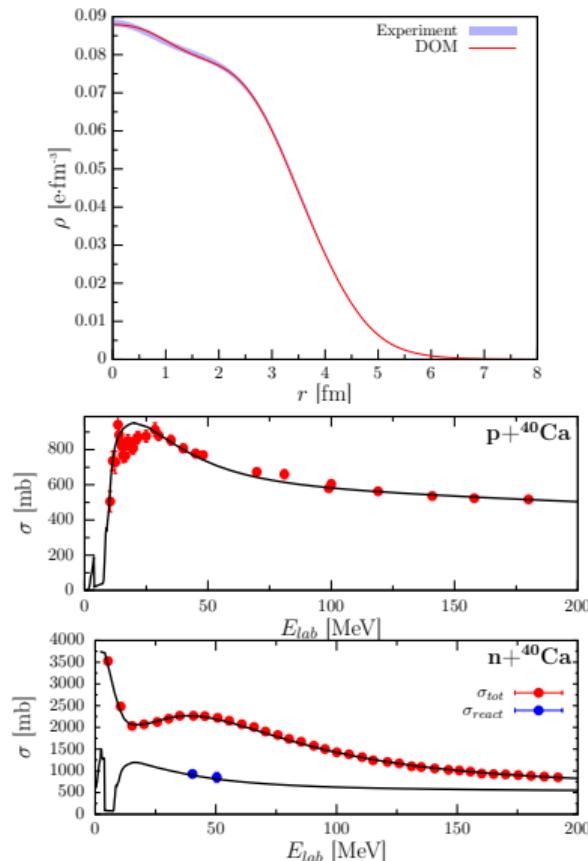
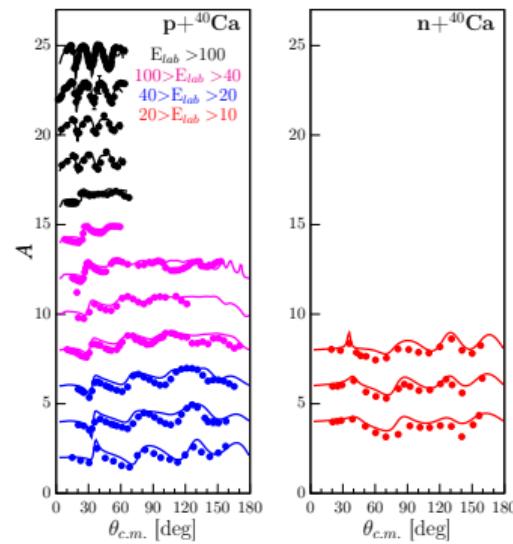
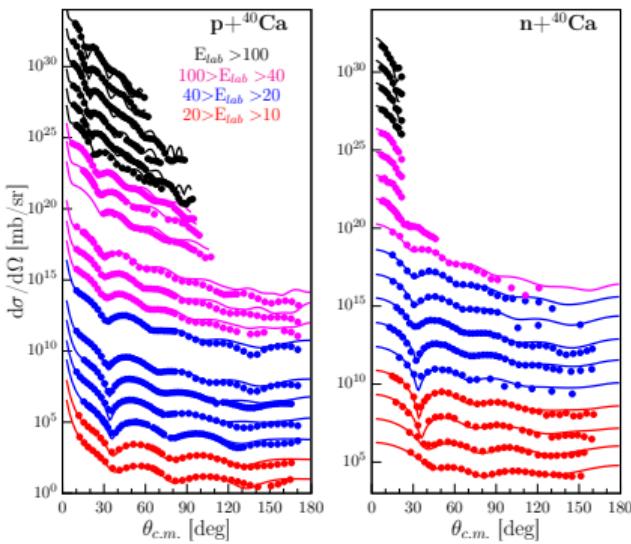
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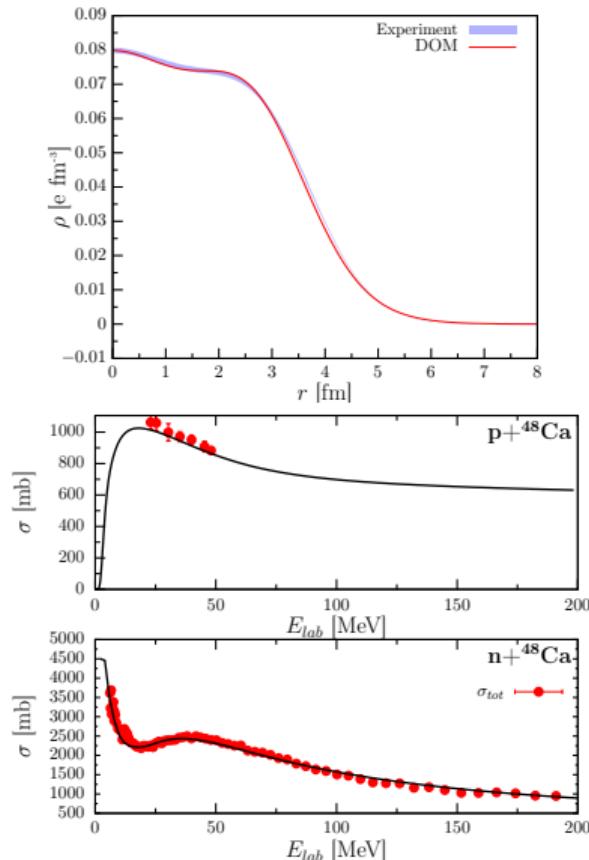
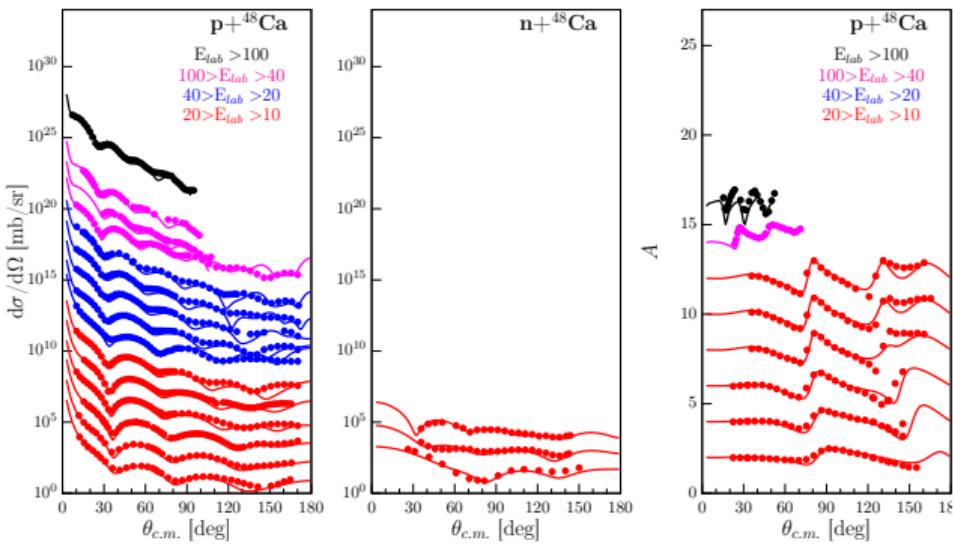
Fitting the Self-energy (^{40}Ca)

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- Reproducing the data means self-energy is found



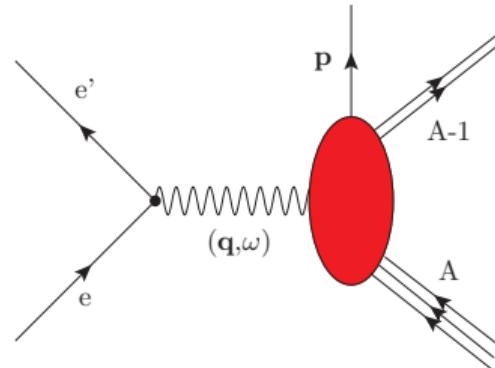
Fitting the Self-energy (^{48}Ca)

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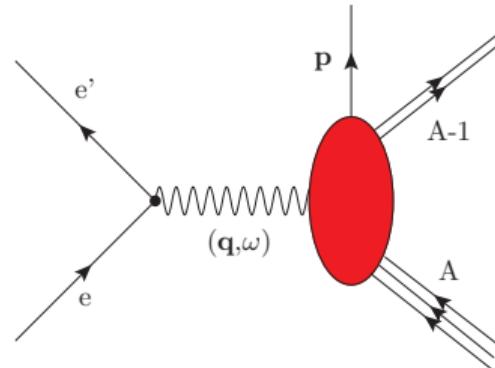
DOM calculation of $^{40}\text{Ca}(e, e'p)^{39}\text{K}$

- DWIA for exclusive reaction (C. Giusti's DWEEPY code)



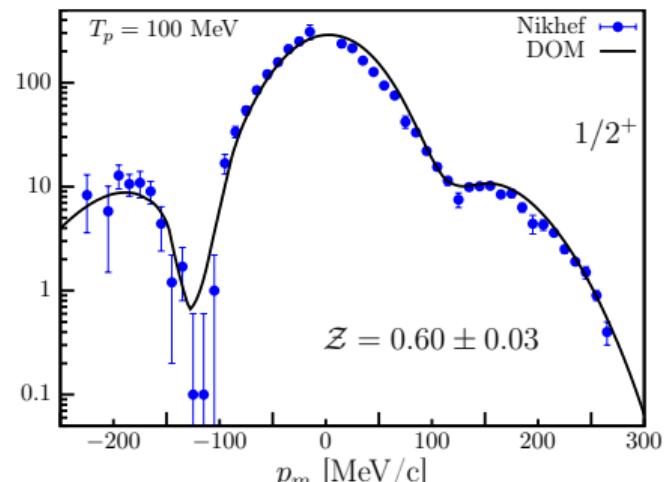
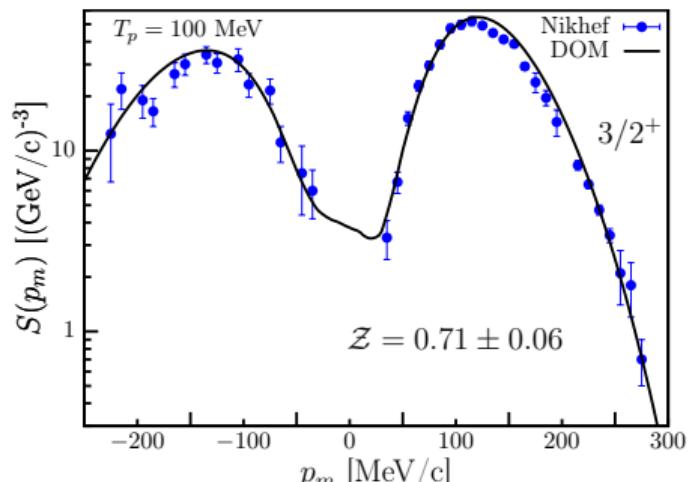
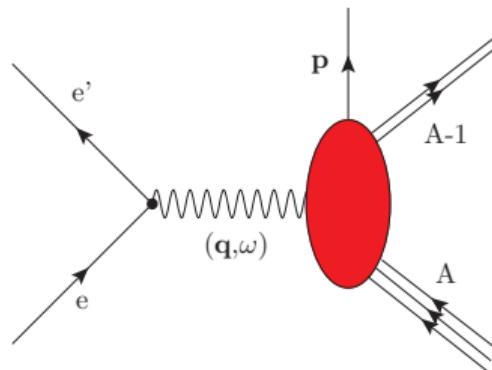
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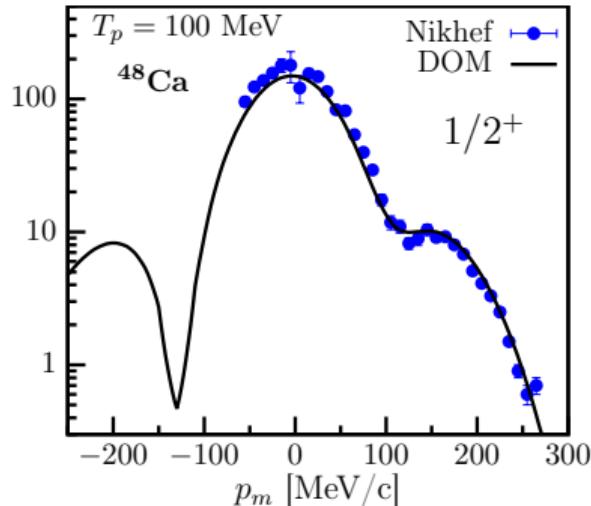
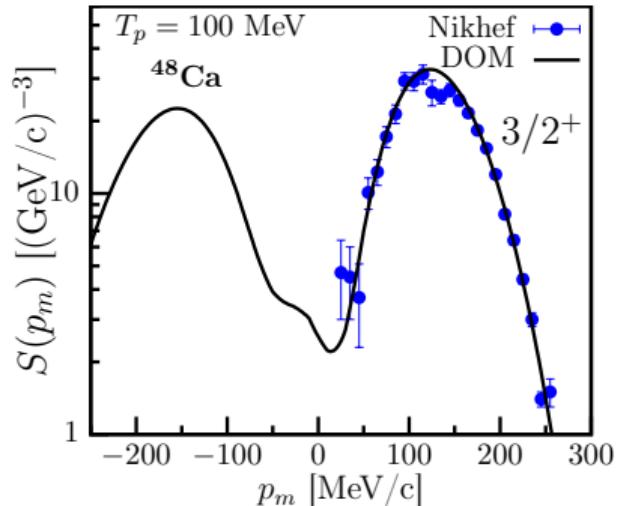
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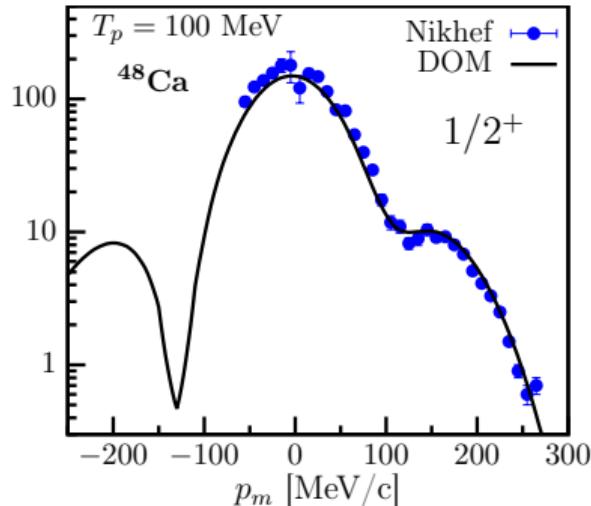
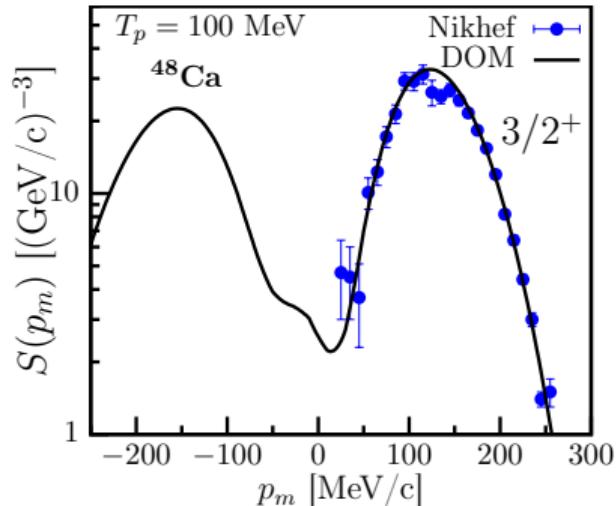
$^{48}\text{Ca}(\text{e},\text{e}'\text{p})^{47}\text{K}$ Momentum Distribution

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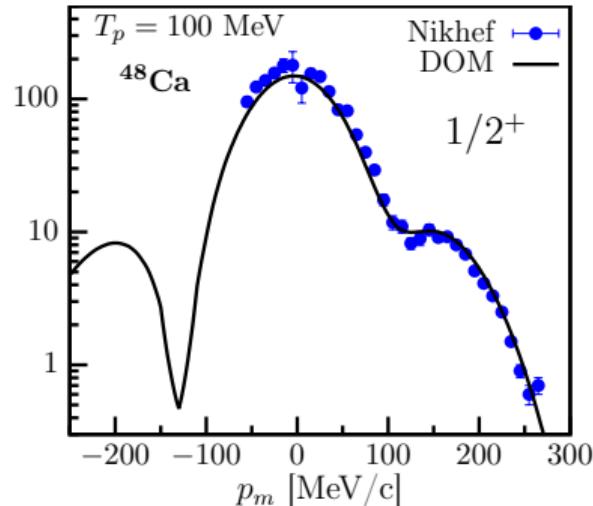
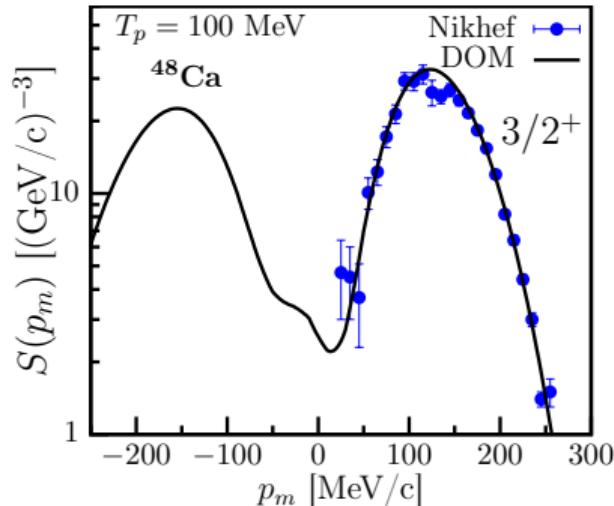


| $\mathcal{Z}_{\ell j}^n$ | $0d\frac{3}{2}$ | $1s\frac{1}{2}$ |
|--------------------------|-----------------|-----------------|
| ^{40}Ca | 0.71 ± 0.04 | 0.60 ± 0.03 |
| ^{48}Ca | 0.58 ± 0.03 | 0.55 ± 0.03 |

M.C. Atkinson and W.H. Dickhoff, Phys. Lett. B, **798**, 135027 (2019)

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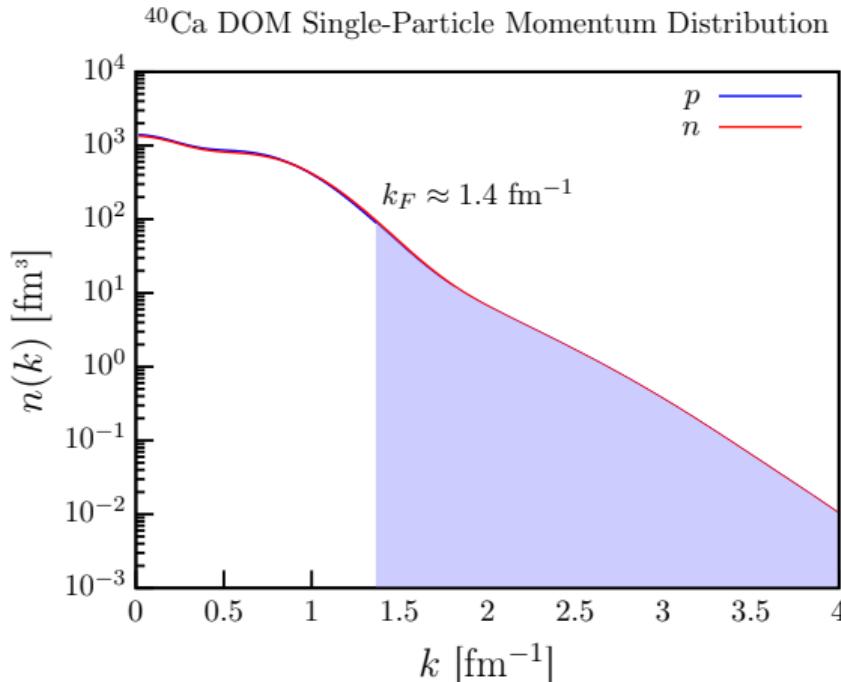
Why are protons more correlated in ^{48}Ca ?

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Momentum Distributions

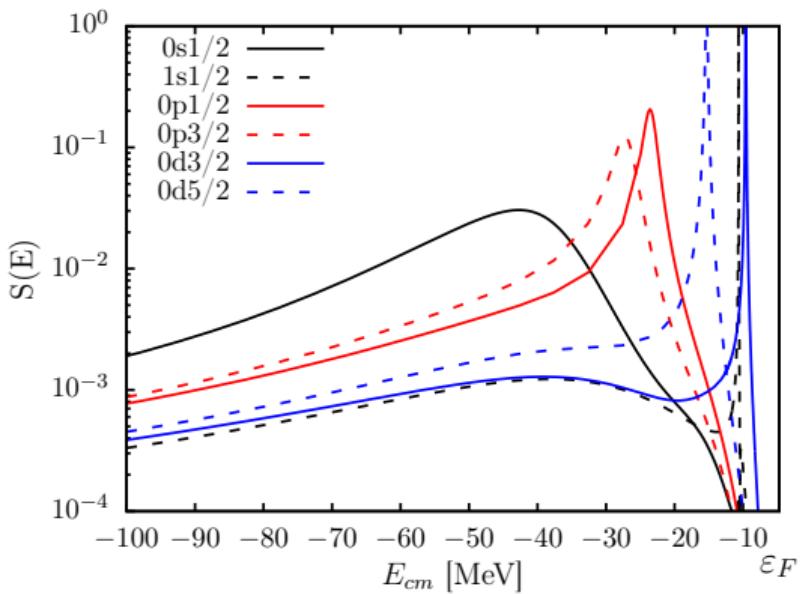
$$n(\mathbf{k}) = \int d^3r \int d^3r' e^{i\mathbf{k}\cdot(\mathbf{r}-\mathbf{r}')} \rho(\mathbf{r}, \mathbf{r}')$$



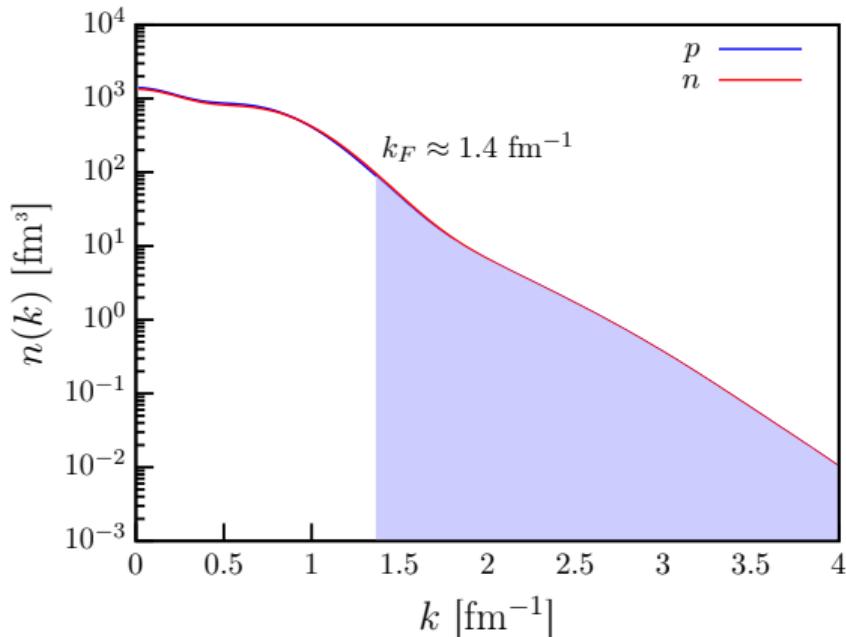
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Proton Spectral Functions in ^{40}Ca

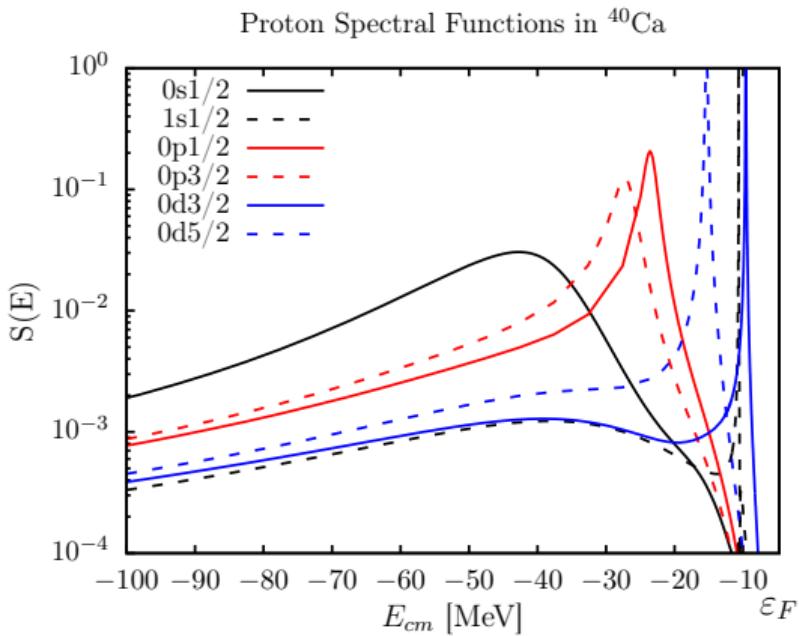


^{40}Ca DOM Single-Particle Momentum Distribution

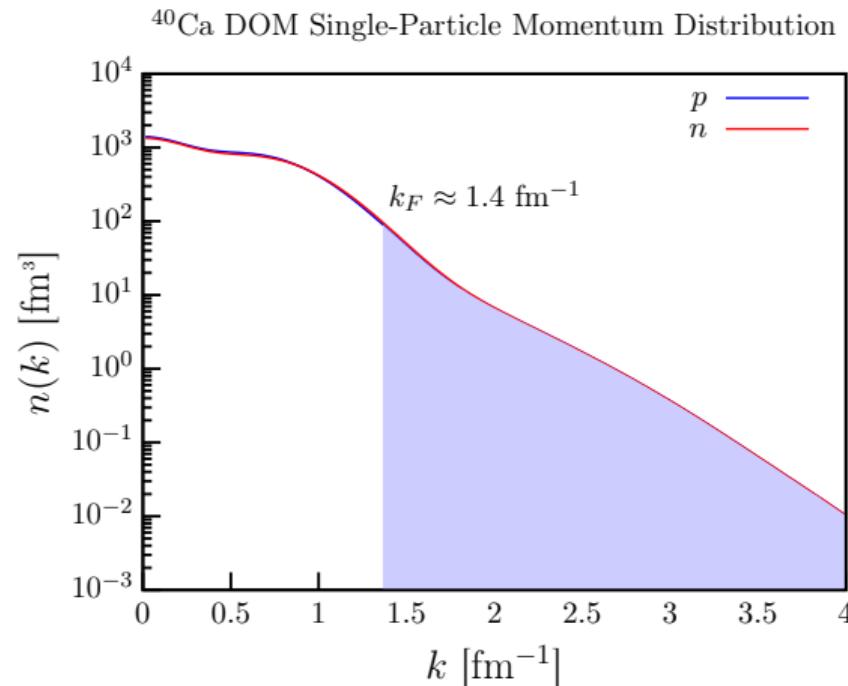


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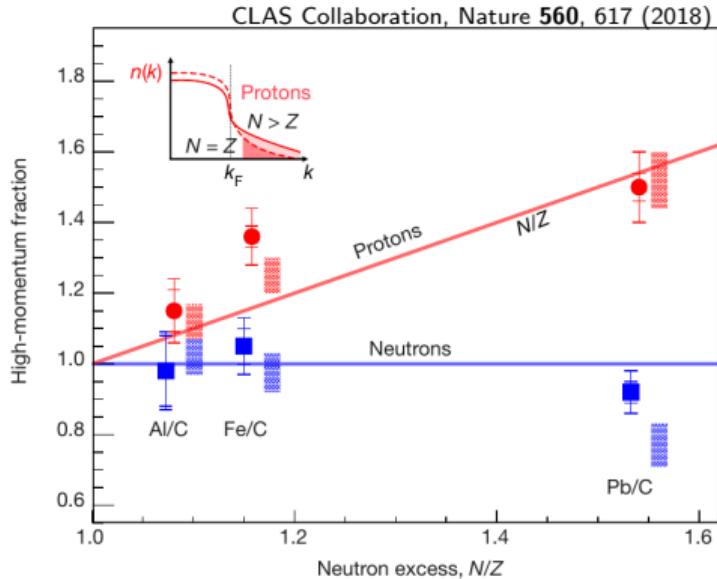
- Short-range correlations (SRC) responsible for this high-momentum content



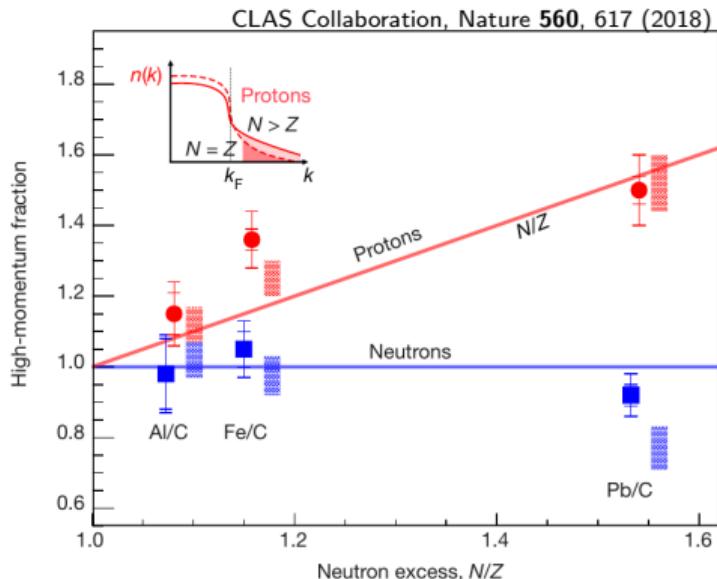
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Asymmetry Dependence of High-Momentum Content

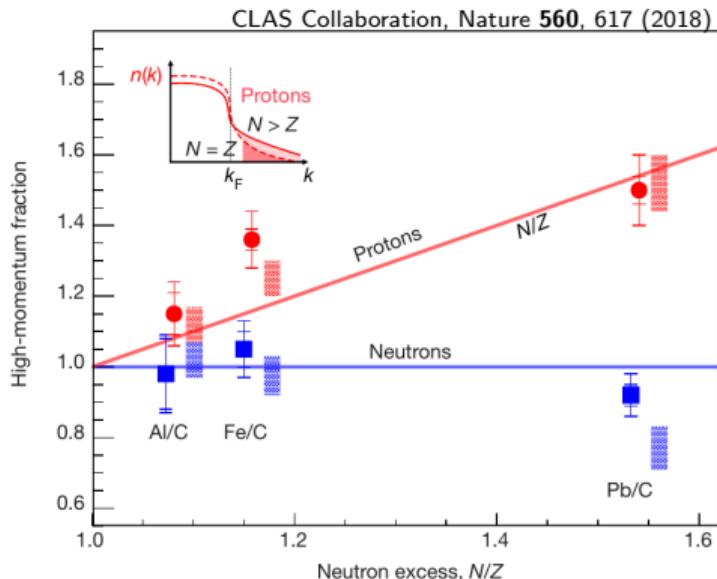


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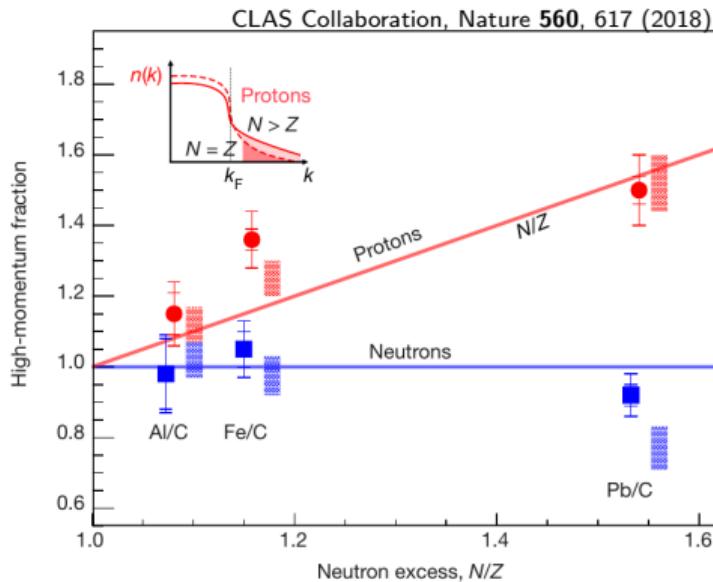
- $n-p$ interaction stronger than $n-n$ or $p-p$

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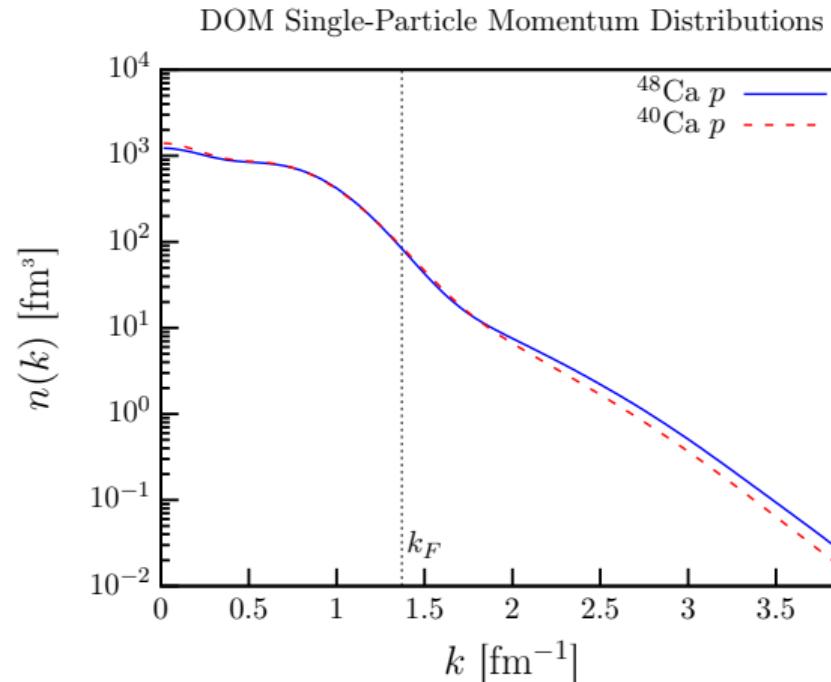
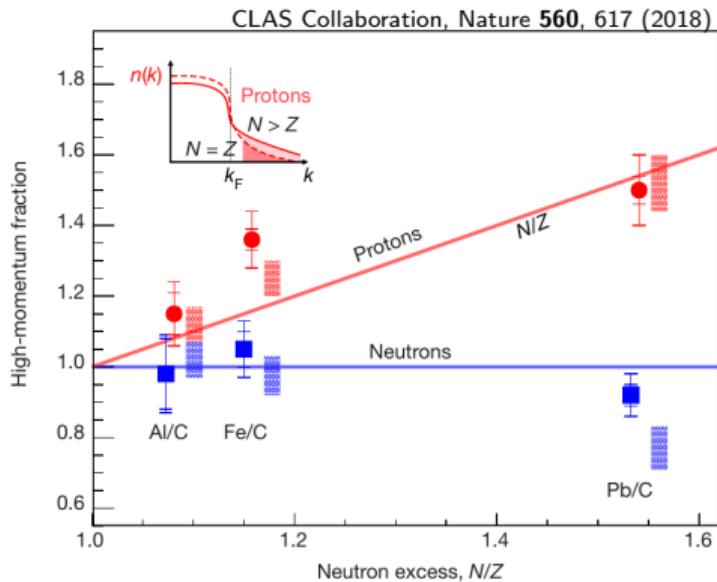
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- More neutrons means more $n-p$ SRC pairs ($n-p$ dominance)

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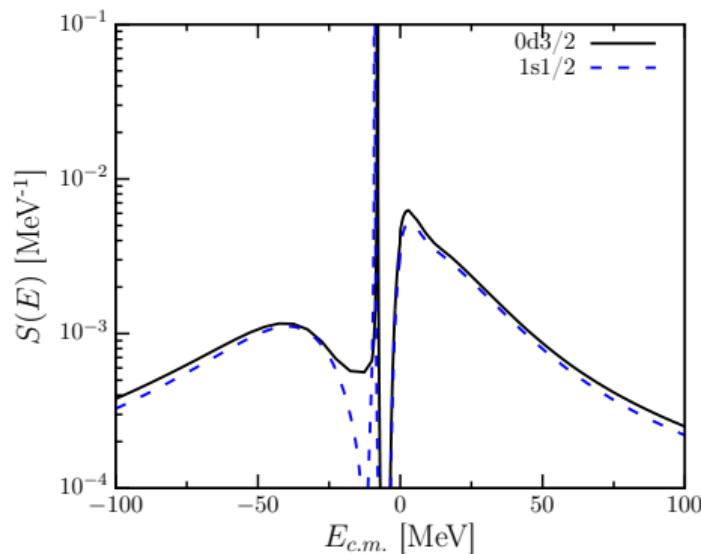
Asymmetry Dependence of High-Momentum Content



- $n\text{-}p$ interaction stronger than $n\text{-}n$ or $p\text{-}p$
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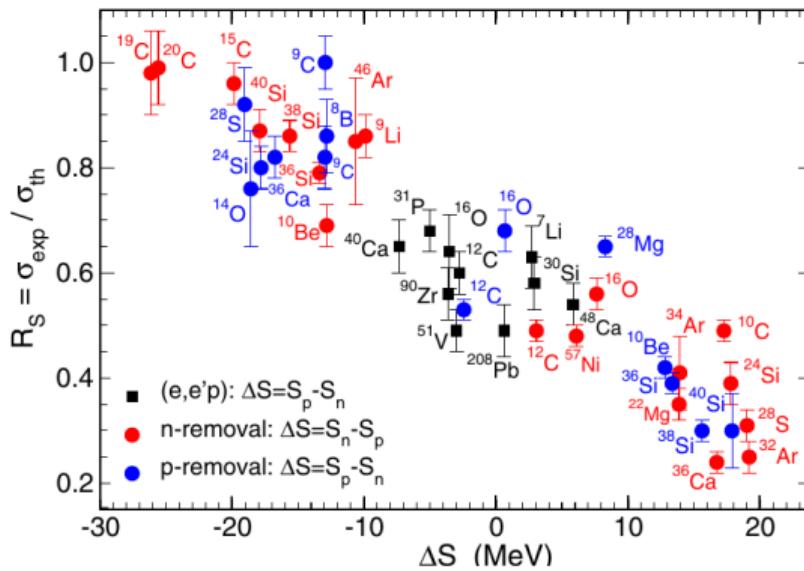
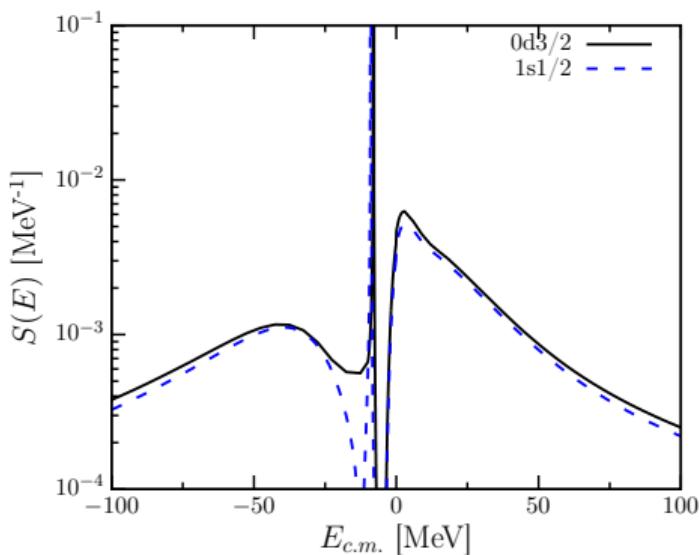
Quenching of Spectroscopic Factors

- Increased proton correlations (SRC) pulls strength from S_F



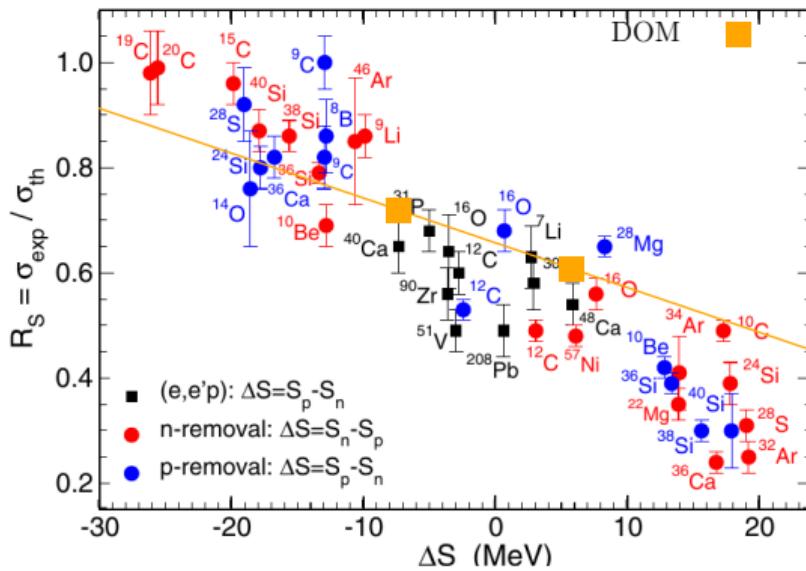
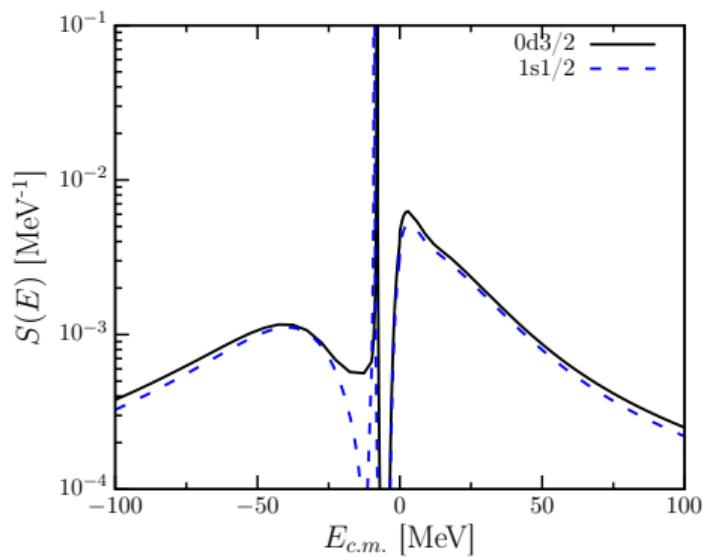
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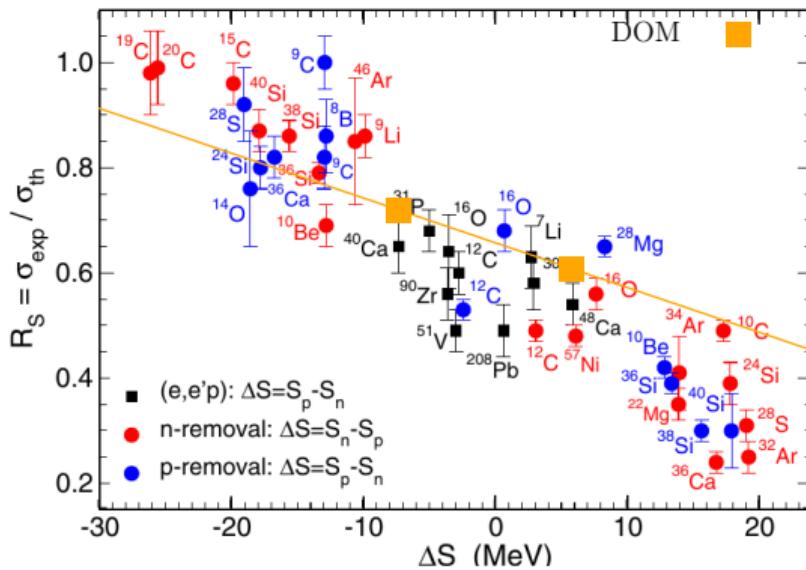
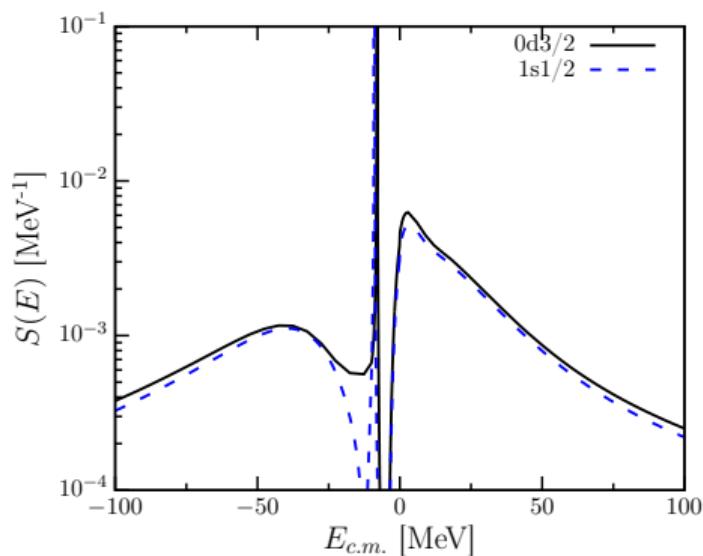
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Quenching is consistent with $n-p$ dominance (SRC)

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- The non-negligible high-momentum content in nuclei is due to SRC correlations
- Proton high-momentum content increases with increasing neutron excess
- The quenching of proton spectroscopic factors is consistent with the np dominance picture

Thanks

- Willem Dickhoff
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