Heavy Quark Radiative Energy Loss in T-matrix Approach

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Based on work: Liu and Rapp: JHEP 08 (2020) 168



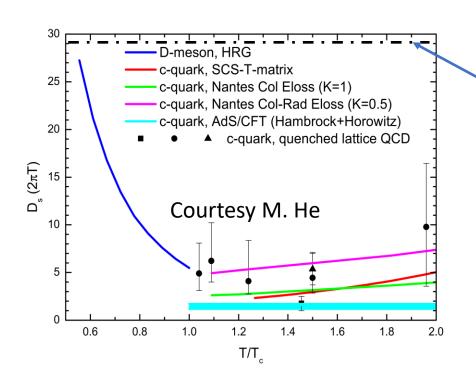


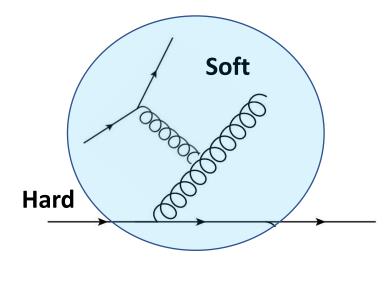
<u>Outline</u>

- 1) Background and Motivation
- 2) Many-Body Approach to Radiative Process
- 3) Study the Non-Perturbative (NP) Effects on Radiative Process
- 4) Conclusion

Non-Perturbative Effects for High Energy Partons?

- A multi-scale problem
- Gluons emitted at soft scale
- Interactions at soft scale, strong!



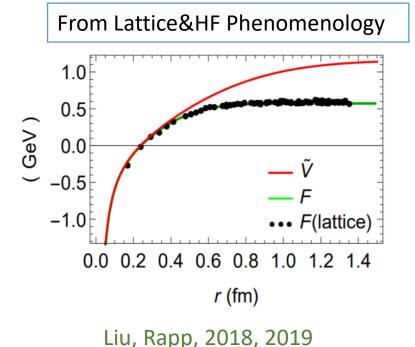


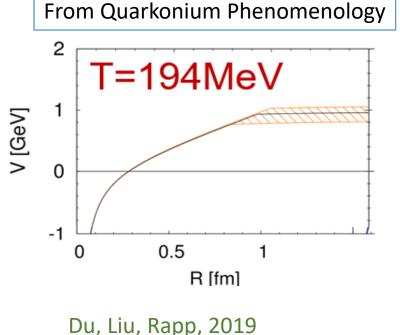
LO pQCD

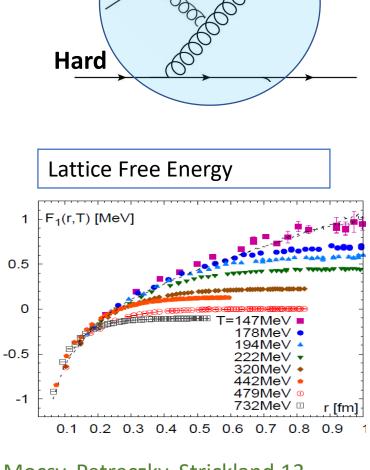
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Not just Coulomb interaction







Soft

Mocsy, Petreczky, Strickland, 13

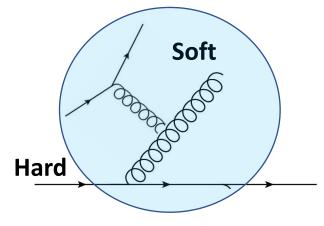
Non-Perturbative Effects for High Energy Partons?

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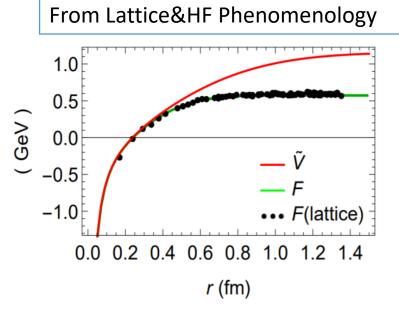
Remnant of confining force,

Resummation,

Non-quasiparticle

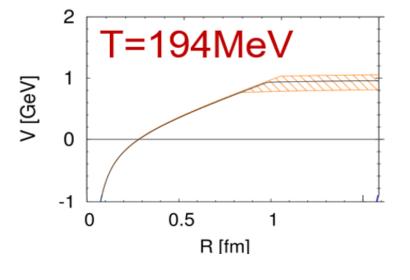


Not just Coulomb interaction



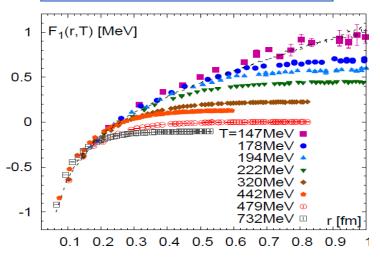
Liu, Rapp, 2018, 2019

From Quarkonium Phenomenology



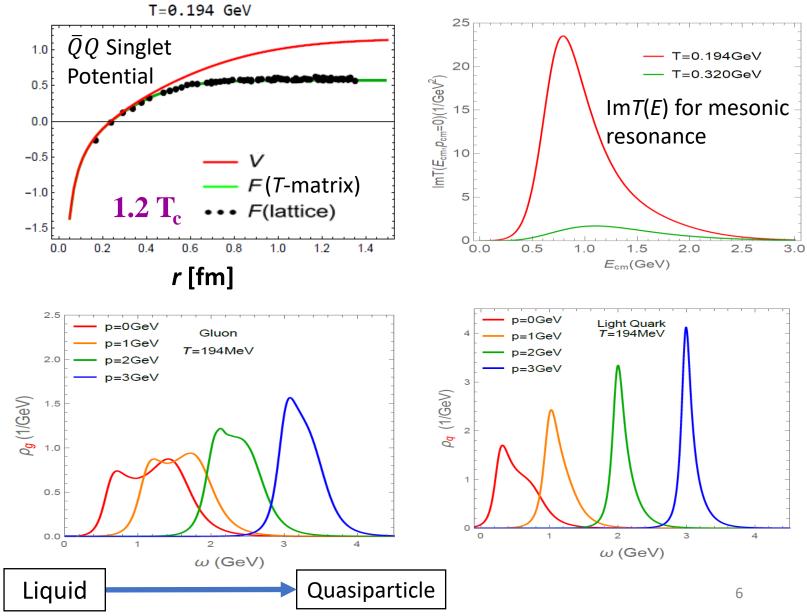
Du, Liu, Rapp, 2019

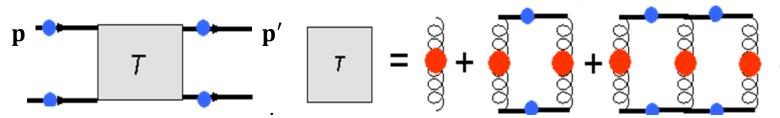
Lattice Free Energy



Mocsy, Petreczky, Strickland,13

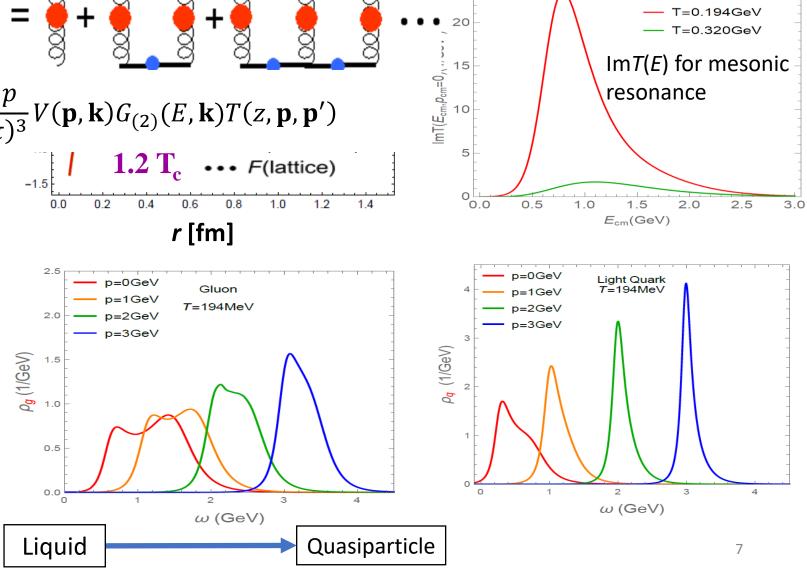
- Remnant of confining force
- Ladder resummationresonances
- Melts low-momentum quasiparticles





$$T(E, \mathbf{p}, \mathbf{p}') = V(\mathbf{p}, \mathbf{p}') + \int \frac{d^3p}{(2\pi)^3} V(\mathbf{p}, \mathbf{k}) G_{(2)}(E, \mathbf{k}) T(z, \mathbf{p}, \mathbf{p}')$$

- Remnant of confining force
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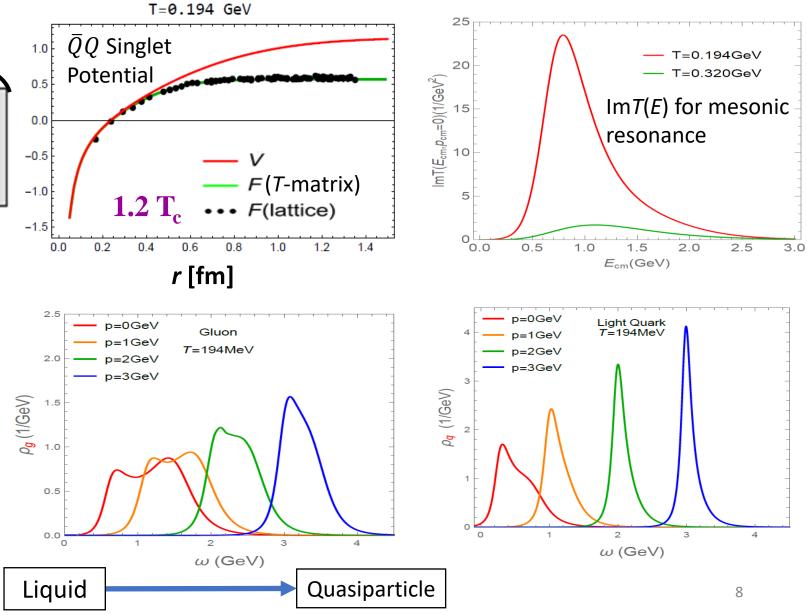


$$\Sigma = \sum_{s,c,f} \int \widetilde{d^4k} T(G) G$$

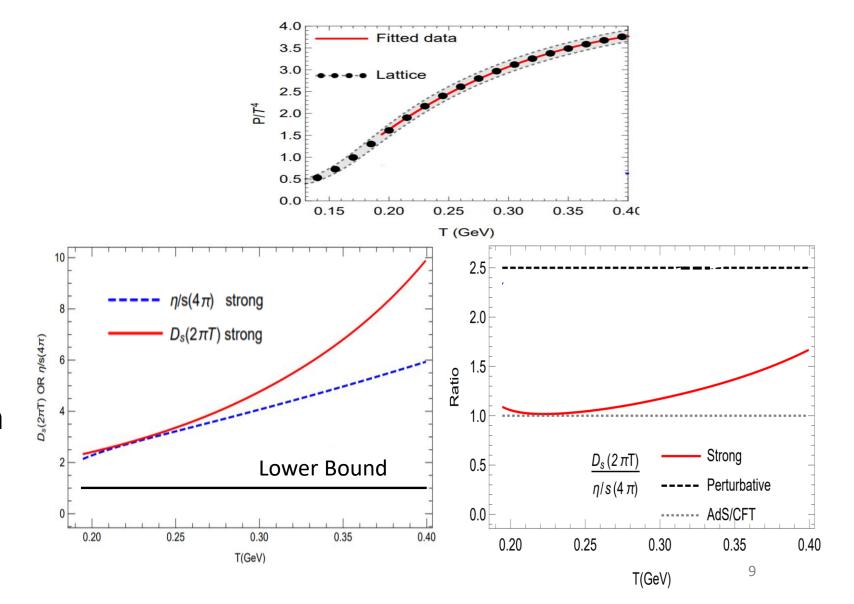
$$\rho(\omega) = \frac{1}{\pi} \operatorname{Im} \frac{-1}{\omega - \varepsilon_p - \Sigma}$$

$$T$$

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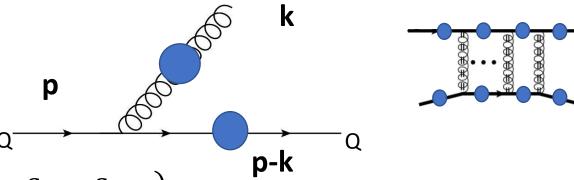


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T-matrix Approach for Energy Loss

 Leading skeleton order radiation:



• Main idea:

On Shell
$$\delta(\varepsilon_{p} - \varepsilon_{k} - \varepsilon_{p-k})$$
Off Shell $\int d\omega dv \delta(\varepsilon_{p} - \omega - v) \rho(\omega, k) \rho(v, p - k)$

Momentum transition rate:

$$w(\mathbf{p}, \mathbf{k}) = \int d^{4}\tilde{q}d\omega'(2\pi)^{4} \delta^{(4)} |M_{Q \leftrightarrow Qg}|^{2} \rho_{Q} [1 - n_{Q}] \rho_{g} [1 + n_{g}]$$

$$\sum |u(\bar{p}')\gamma_{\mu}u(p)\epsilon^{\mu}(k)|^{2}$$

Spectral functions

$$\rho(\omega) = \frac{1}{\pi} \operatorname{Im} \frac{-1}{\omega - \varepsilon_{p} - \Sigma} \longrightarrow \begin{cases} \operatorname{Encode} \\ \operatorname{Medium Eff} \end{cases}$$

T-matrix Approach for Energy Loss

 Leading skeleton order radiation: p $-\varepsilon_k - \varepsilon_{n-k}$

• Main idea:

On Shell $\delta(\varepsilon_{p} - \varepsilon_{k} - \varepsilon_{p-k})$ Off Shell $\int d\omega dv \delta(\varepsilon_{p} - \omega - v) \rho(\omega, k) \rho(v, p - k)$

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Spectral functions

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Encode
Medium Effects

Four Cases with Different NP Effects

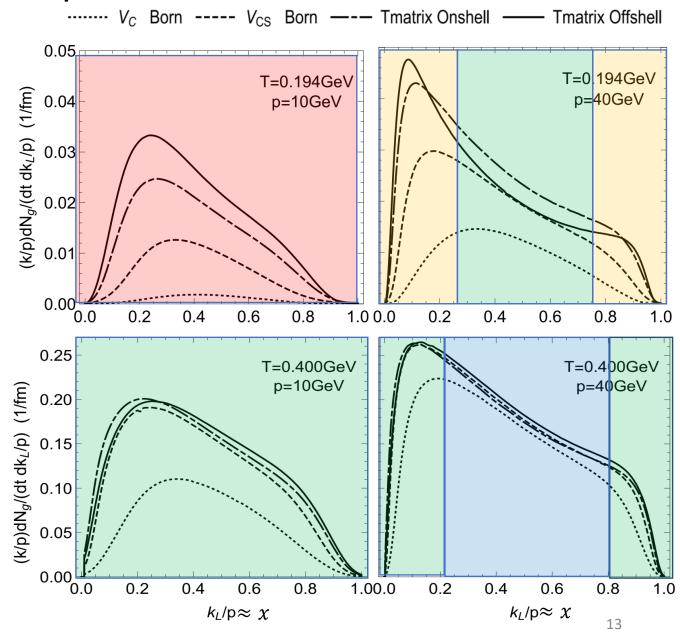
(1) V_C Born	(2) V_{CS} Born	(3) Tmatrix Onshell	(4) Tmatrix Offshell
Coulomb	Coulomb+String	Coulomb+String	Coulomb+String
Leading-order	Leading-order	All order	All order
quasi-particle	quasi-particle	quasi-particle	off-shell spectra
	BOBOBOBO A		
Close to pQCD	Add confining interaction	Add the t-channel resummation	Add off-shell medium partons; our full T-matrix prediction

- Similar diagrams for rescattering of the outgoing gluons
- Similar medium parton density in all cases (fixed by lattice EoS)

Radiative Power Spectrum

- Power spectrum $\frac{(k/p)dN_g}{dt \ d(k_L/p)} \approx \frac{xdN_g}{dt \ dx}$
 - **Red region**: Confining interaction, Resummation, Off-shell medium
 - **Yellow region :** Confining interaction, Resummation, Off-shek medium
 - **Green region**: Confining interaction, Resummation, Off-shell medium
 - Blue region: Confining interaction, Resummation, Off-shell medium

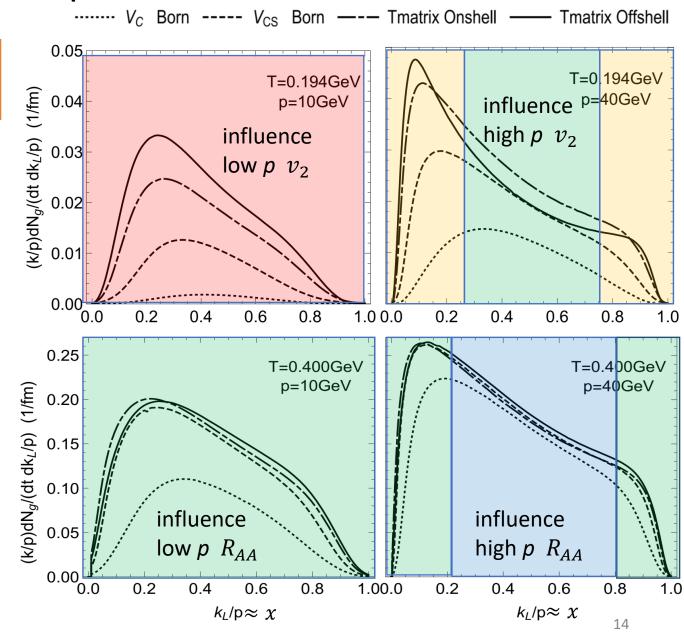
 Perturbative



Radiative Power Spectrum

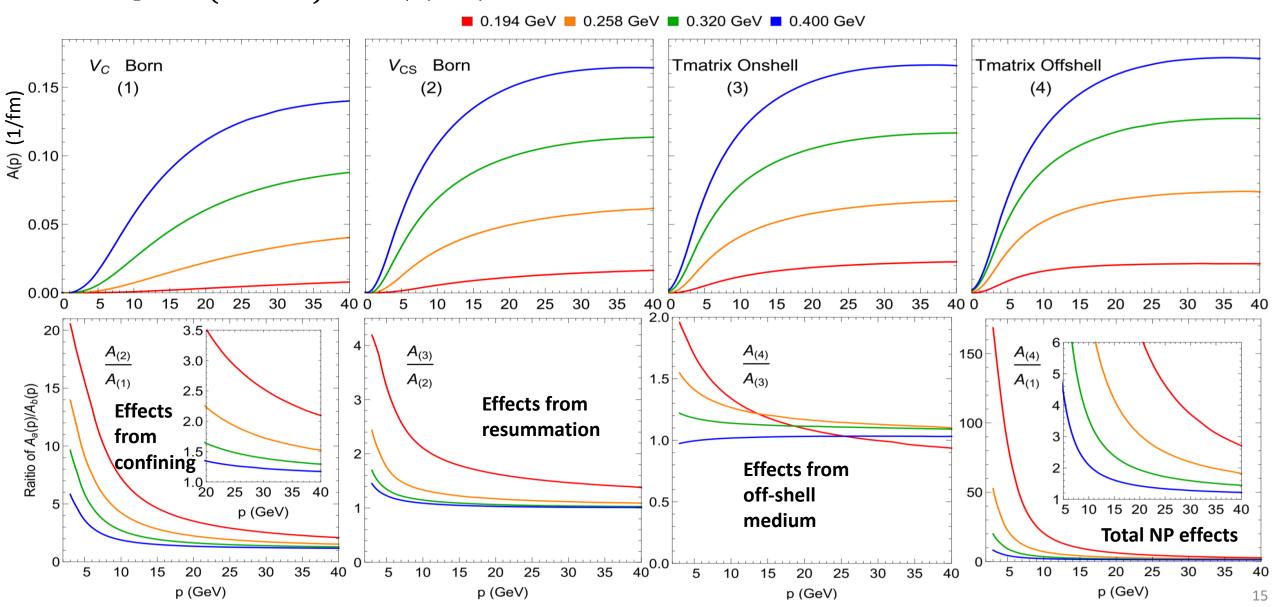
- Adding confining interaction is quite helpful in lots of phase space
- **Red region**: Confining interaction, Resummation, Off-shell medium
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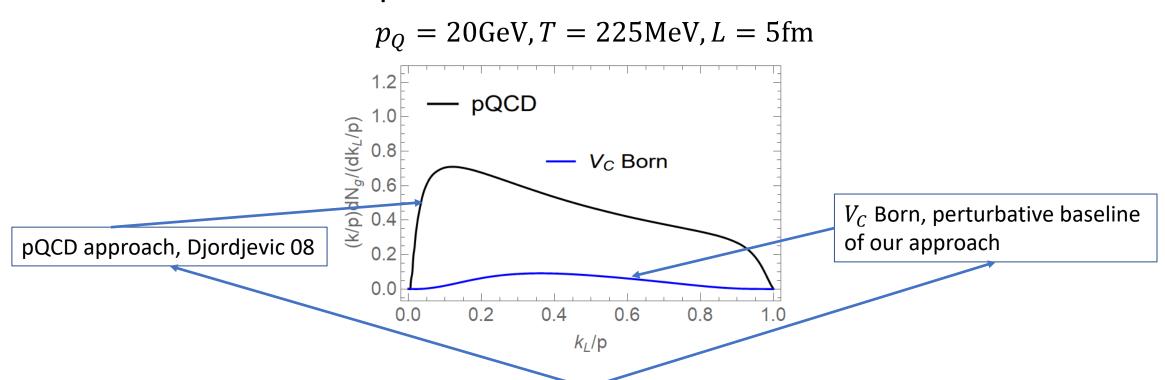
 Perturbative



Radiative Contribution to Drag Coefficients

• A(p) $\approx (E^{-1}dE)/dt$ (1/fm) percentage of energy loss per (fm) time

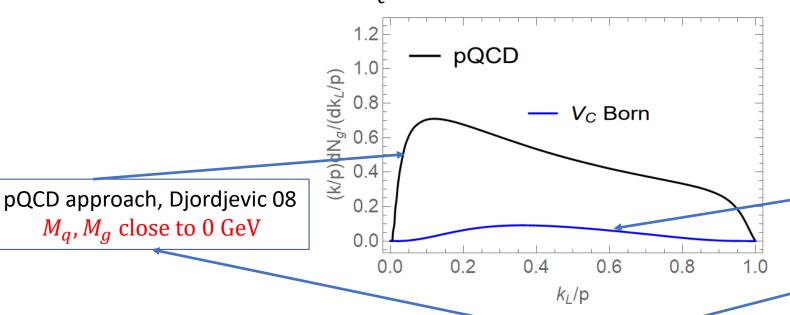




Both are perturbative calculation Why different?

 M_q , M_g close to 0 GeV

$$p_O = 20 \text{GeV}, T = 225 \text{MeV}, L = 5 \text{fm}$$



 V_C Born, perturbative baseline of our approach

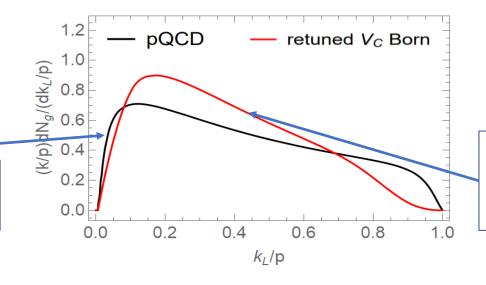
 $M_q \sim 0.5 \text{ GeV}$, $M_q \sim 1 \text{ GeV}$

Masses are different

pQCD approach, Djordjevic 08

 M_q , M_g close to 0 GeV

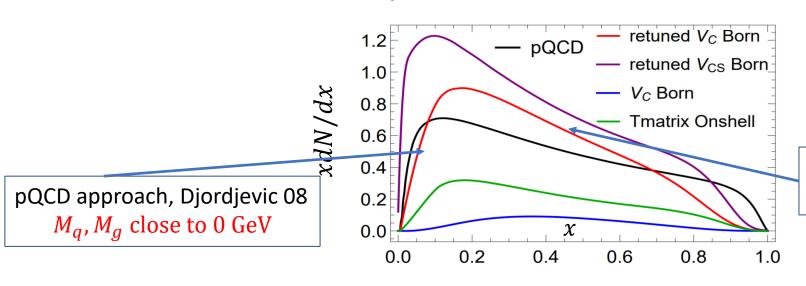
$$p_{Q} = 20 \text{GeV}, T = 225 \text{MeV}, L = 5 \text{fm}$$



Retuned V_C Born, Perturbative baseline of our approach with same mass as pQCD that M_q , M_g close to 0 GeV

 Similar if masses are tune to be the same

$$p_O = 20 \text{GeV}, T = 225 \text{MeV}, L = 5 \text{fm}$$



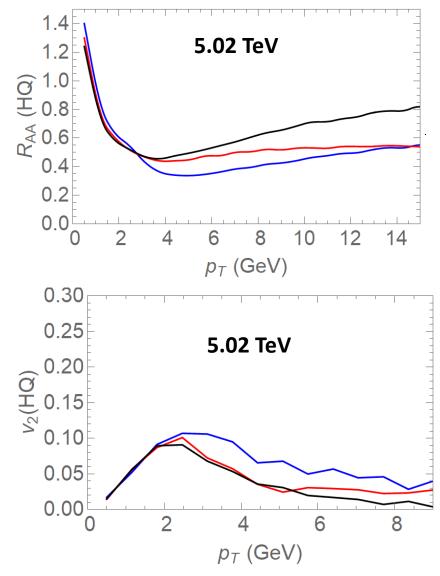
- **Retuned** means pQCD mass
- Otherwise, NP masses

Green line: pQCD interaction + pQCD masses > NP interactions + NP masses

Which picture is more correct for radiative energy loss at intermediate/low momentum?

Apply to Heavy-ion Collisions

- T-matrix elastic only
- T-matrix elastic+ radiation in this work
- T-matrix elastic + High Twist radiation using NP \hat{q} (>5 × \hat{q}_{pQCD} at small T, p) from T-matrix
 - Hybrid, High Twist + NP interactions
 Code used in Cao, Sun, Li, Liu, Xing, Qin PLB, 2020
- Radiation is important for R_{AA} and v_2 at low energy~5 GeV, where NP effects are unavoidable
- Gluon (radiative) Mass (Large(~1GeV)/small) leads to a difference between red and blue
- Could we use these to understand/constrain the NP gluon masses experimentally?

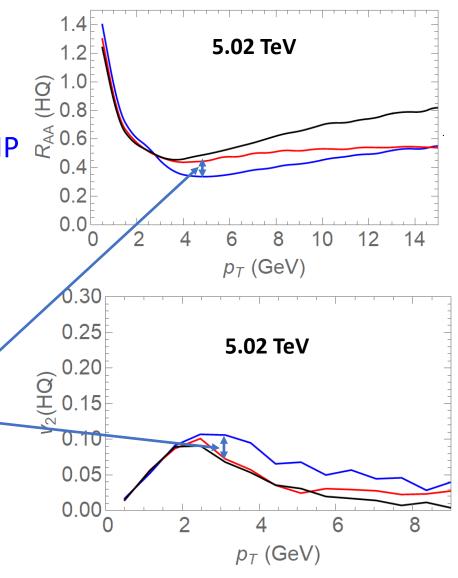


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Conclusion

- Developed many-body approach to study the non-perturbative effects for radiative energy loss
- Non-perturbative effects on radiation are important for R_{AA} and v_2 at small/intermediate p_T (5-10 GeV)
- Adding confining force is probably enough for lots of phase space; resummation and off-shell medium effects are important in low p and low T
- Non-perturbative masses are also important uncertainties for radiative process