

HQ hadronization & hadronic phase in PHSD

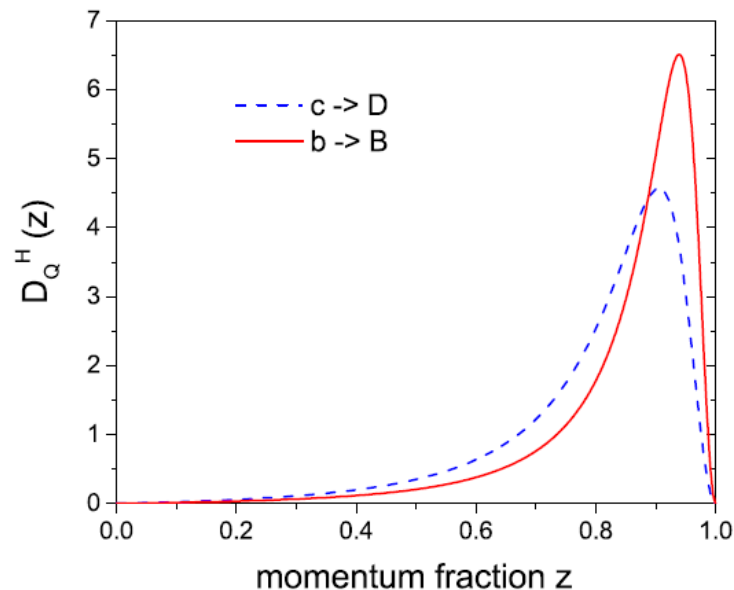
Taesoo Song, Elena Bratkovskaya

1. HQ fragmentation

- transverse momentum
- Peter's fragmentation function

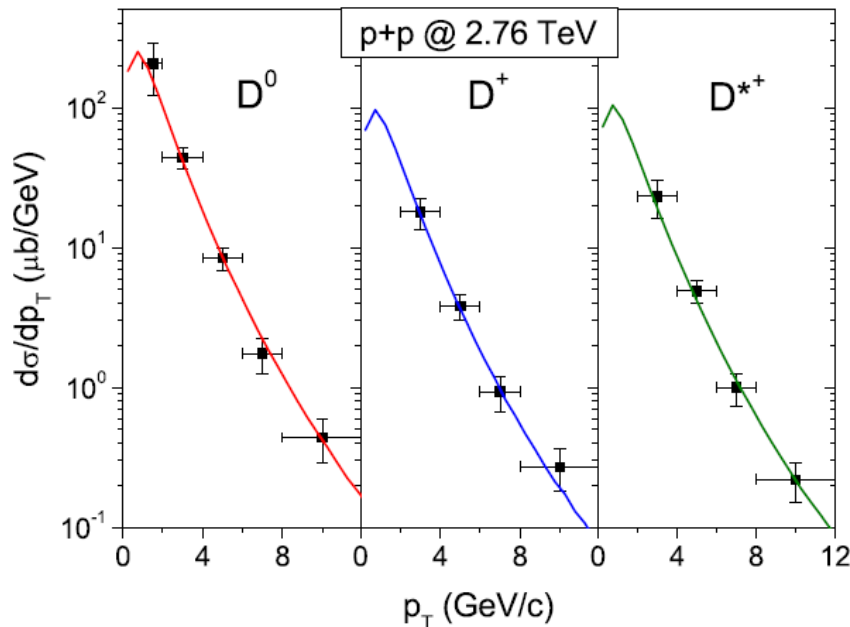
$$D_Q^H(z) \sim \frac{1}{z[1 - 1/z - \epsilon_Q/(1 - z)]^2},$$

$\epsilon_Q = 0.01$ for charm and 0.004 for bottom

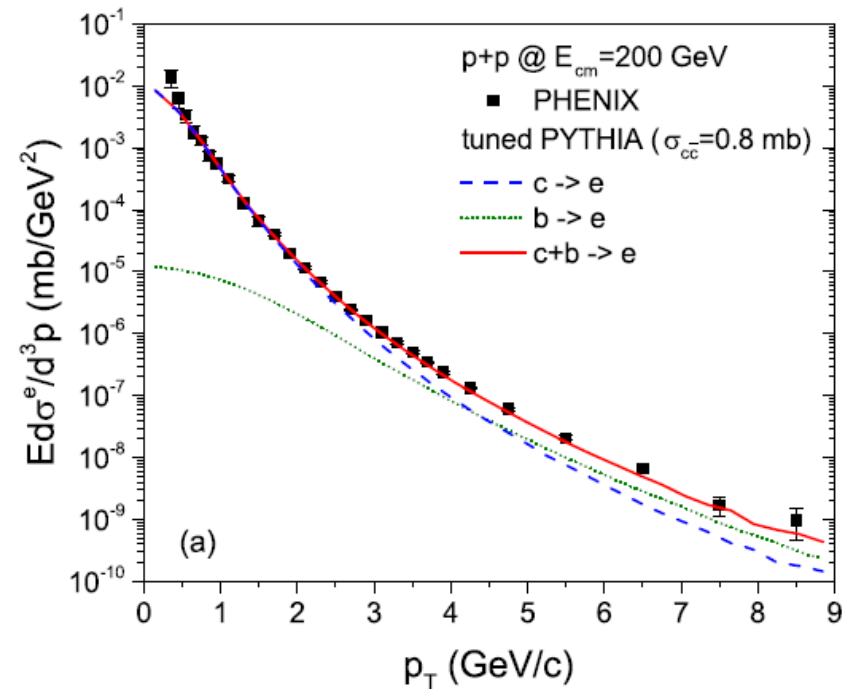


- rapidity
- is assumed to be same after the fragmentation

The chemical fractions of the charm quark decay into D^+ , D^0 , D^{*+} , D^{*0} , D_s , and Λ_c are taken to be 14.9, 15.3, 23.8, 24.3, 10.1, and 8.7% [30,43–45], respectively, and those of the bottom quark decay into B^- , B^0 , B_s^0 , and Λ_b are 39.9, 39.9, 11, and 9.2% [42].



T. Song et al. PRC 93 (2016) 034906



T. Song et al. PRC 96 (2017) 014905

2. HQ coalescence

- Wigner function for S state coalescence probability

$$\Phi(\mathbf{r}, \mathbf{p}) = 8 \exp \left[-\frac{r^2}{\sigma^2} - \sigma^2 p^2 \right],$$

$$\mathbf{r} = \mathbf{r}_1 - \mathbf{r}_2,$$

$$\mathbf{p} = (m_2 \mathbf{p}_1 - m_1 \mathbf{p}_2) / (m_1 + m_2),$$

in center-of-mass frame

$$\begin{aligned} \langle r_M^2 \rangle &= \frac{1}{2} \langle (\mathbf{R} - \mathbf{r}_1)^2 + (\mathbf{R} - \mathbf{r}_2)^2 \rangle \\ &= \frac{1}{2} \frac{m_1^2 + m_2^2}{(m_1 + m_2)^2} \langle r^2 \rangle = \frac{3}{4} \frac{m_1^2 + m_2^2}{(m_1 + m_2)^2} \sigma^2, \end{aligned}$$

multiplied by 4 (=1+3) to include D* and divided by 36(=6*6)

- Wigner function for P state coalescence probability

$$\left(\frac{16}{3} \frac{y_i^2}{\sigma_i^2} - 8 + \frac{16}{3} \sigma_i^2 k_i^2 \right) \exp \left(-\frac{y_i^2}{\sigma_i^2} - k_i^2 \sigma_i^2 \right)$$

multiplied by 8 (=3+5)
to include $J^P = 1^+, 2^+$
and divided by 36(=6*6)

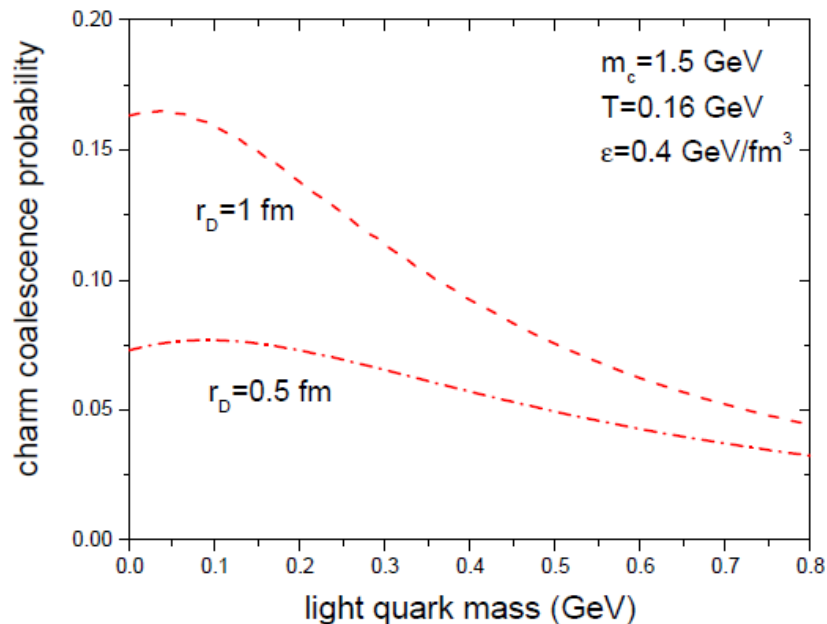
- We assume P states decay to $D(D^*) + \pi$ as soon as hadronized
- D_s is similarly treated

0.9 fm

Model Study: total coalescence probability of thermal charm

Charm can be projected into all possible states (S-, P-, D-states ...) with different coalescence radii $\rightarrow$ too complicated !!

T. Song, G. Coci, arXiv:2104.10987



- Ground state from the coalescence
- Heavier states from the statistical model

$$\frac{\text{number density of all charm hadrons}}{\text{number density of } D^0} = 7.47,$$

at T_c

The total coalescence probability is much lower than 1

There must be correlations in color, momentum, coordinate spaces

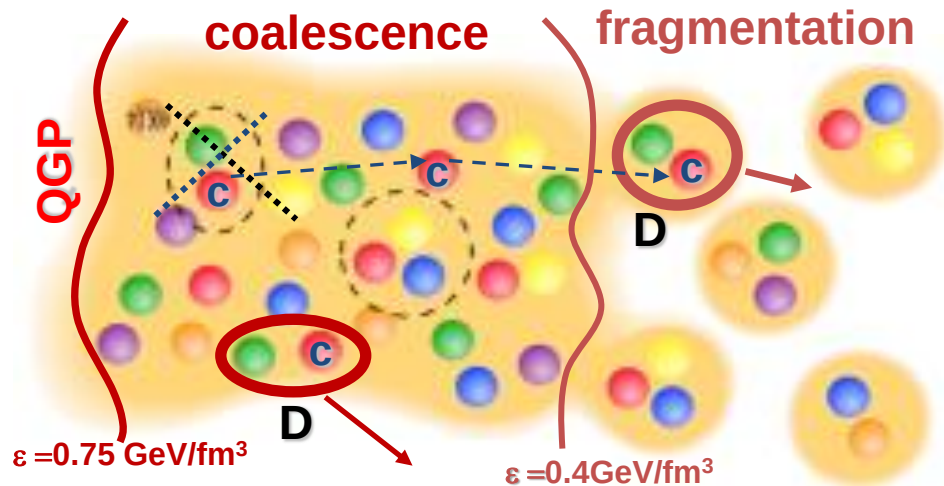
Hadronization of heavy quarks in A+A

□ PHSD: if the local energy density $\varepsilon \rightarrow \varepsilon_c \rightarrow$ hadronization of heavy quarks to hadrons

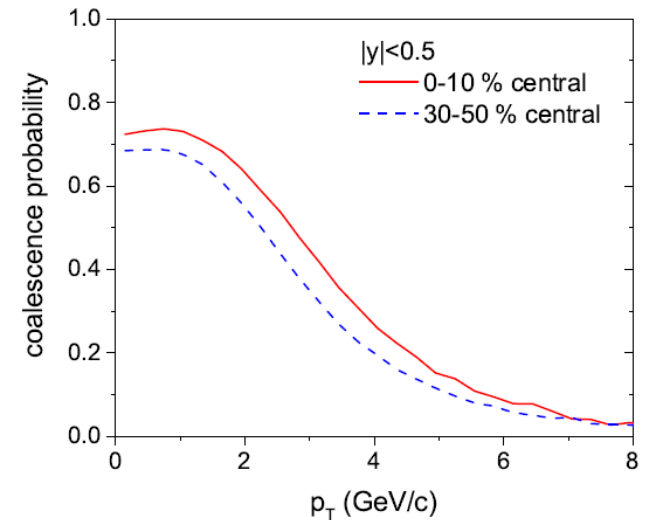
T. Song et al., PRC 93 (2016) 034906

Dynamical hadronization scenario for heavy quarks :

coalescence with $\langle r \rangle = 0.9$ fm & fragmentation
 $0.4 < \varepsilon < 0.75$ GeV/fm³ $\varepsilon < 0.4$ GeV/fm³



Coalescence probability in Pb+Pb at LHC



T. Song et al. PRC 93 (2016) 034906

3. Hadronic scattering

Modelling of D-meson scattering in the hadronic gas

1. D-meson scattering with mesons

Model: effective chiral Lagrangian approach with heavy-quark spin symmetry

L. M. Abreu, D. Cabrera, F. J. Llanes-Estrada, J. M. Torres-Rincon, Annals Phys. 326, 2737 (2011)

Interaction of $D=(D^0, D^+, D^+_s)$ and $D^*=(D^{*0}, D^{*+}, D^{*+}_s)$ with octet $(\pi, K, Kbar, \eta)$:

$$\begin{aligned} \mathcal{L}_{LO} = & \langle \nabla^\mu D \nabla_\mu D^\dagger \rangle - m_D^2 \langle D D^\dagger \rangle - \langle \nabla^\mu D^{*\nu} \nabla_\mu D_\nu^{*\dagger} \rangle \\ & + m_D^2 \langle D^{*\mu} D_\mu^{*\dagger} \rangle + ig \langle D^{*\mu} u_\mu D^\dagger - D u^\mu D_\mu^{*\dagger} \rangle \\ & + \frac{g}{2m_D} \langle D_\mu^* u_\alpha \nabla_\beta D_\nu^{*\dagger} - \nabla_\beta D_\mu^* u_\alpha D_\nu^{*\dagger} \rangle \epsilon^{\mu\nu\alpha\beta} \end{aligned} \quad \text{with}$$

$$\begin{aligned} \nabla_\mu &= \partial_\mu - \frac{1}{2}(u^\dagger \partial_\mu u + u \partial_\mu u^\dagger), \\ u_\mu &= i(u^\dagger \partial_\mu u - u \partial_\mu u^\dagger) \end{aligned}$$

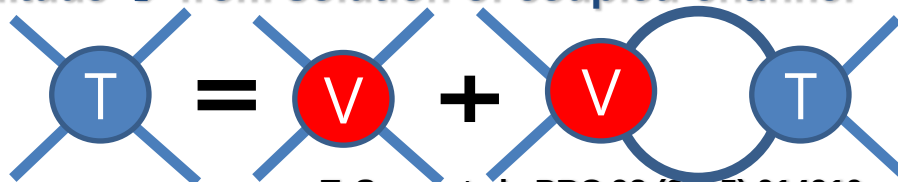
$$U = u^2 = \exp\left(\frac{\sqrt{2}i\Phi}{f}\right) \quad \Phi = \begin{pmatrix} \frac{1}{\sqrt{2}}\pi^0 + \frac{1}{\sqrt{6}}\eta & \pi^+ & K^+ \\ \pi^- & -\frac{1}{\sqrt{2}}\pi^0 + \frac{1}{\sqrt{6}}\eta & K^0 \\ K^- & \bar{K}^0 & -\frac{2}{\sqrt{6}}\eta \end{pmatrix}$$

2. D-meson scattering with baryons

Model: G-matrix approach: interactions of $D=(D^0, D^+, D^+_s)$ and $D^*=(D^{*0}, D^{*+}, D^{*+}_s)$ with nucleon octet $J^P=1/2^+$ and Delta decuplet $J^P=3/2^+$

C. Garcia-Recio, J. Nieves, O. Romanets, L. L. Salcedo, L. Tolos, Phys. Rev. D 87, 074034 (2013)

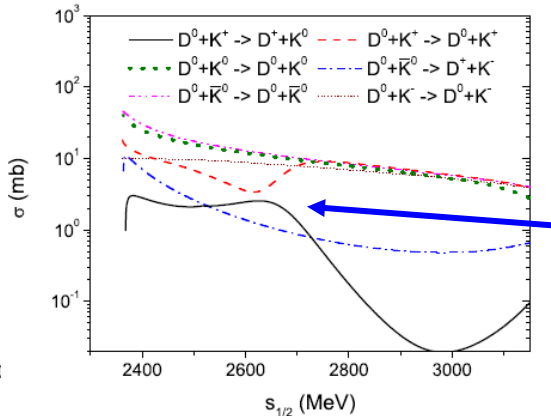
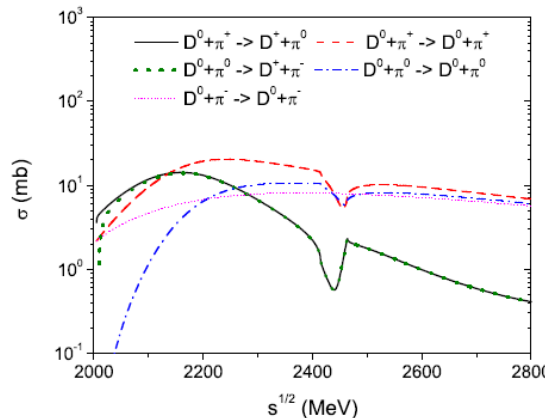
Unitarized scattering amplitude $\rightarrow$ from solution of coupled-channel Bethe-Salpeter equations:



T. Song et al., PRC 92 (2015) 014910, arXiv:1503.03039

D-meson scattering in the hadron gas

1. D-meson scattering with mesons

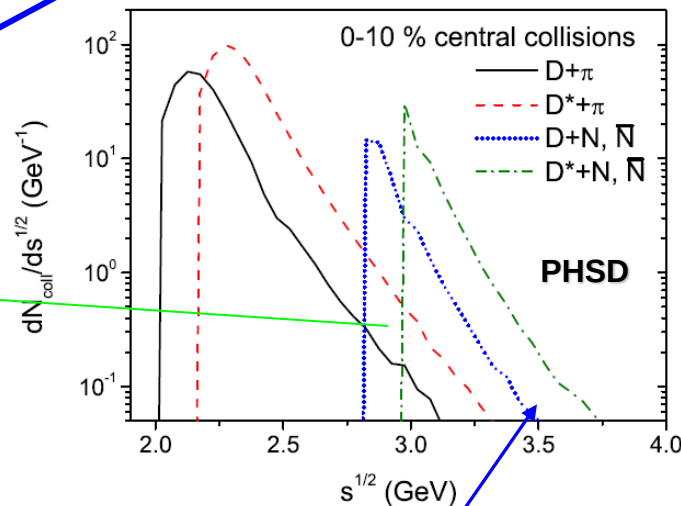
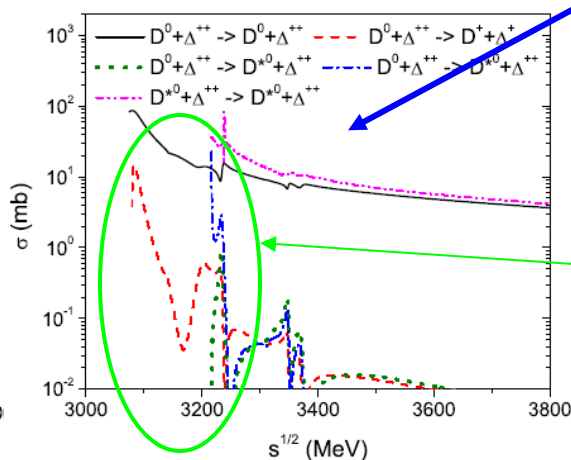
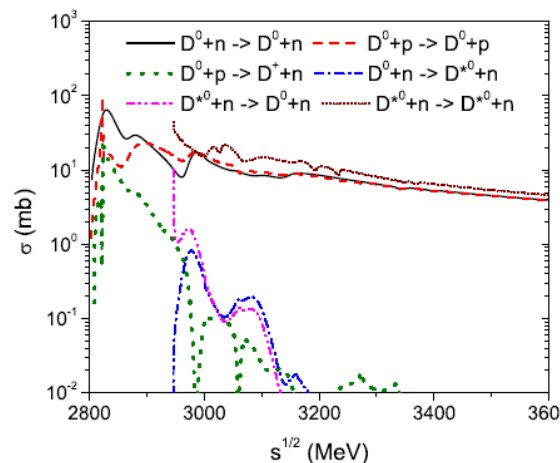


1a) cross sections with $m=\rho, \omega, \phi, K^*, \dots$ taken as

$$\sigma(D, D^* + m) = 10 \text{ mb}$$

→ Strong **isospin dependence** and complicated structure due to the resonance coupling of $D+m$, $D+B$ cross sections!

2. D-meson scattering with baryons



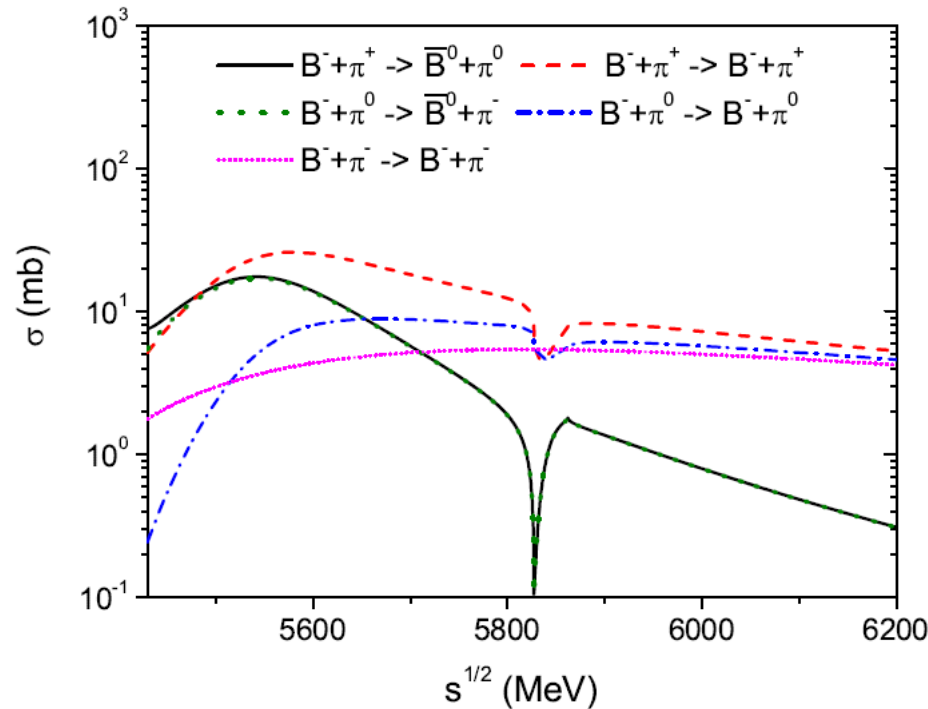
→ Hadronic interactions become ineffective for the energy loss of D, D^* mesons at high transverse momentum (i.e. large $s^{1/2}$)

B-meson scattering in the hadron gas

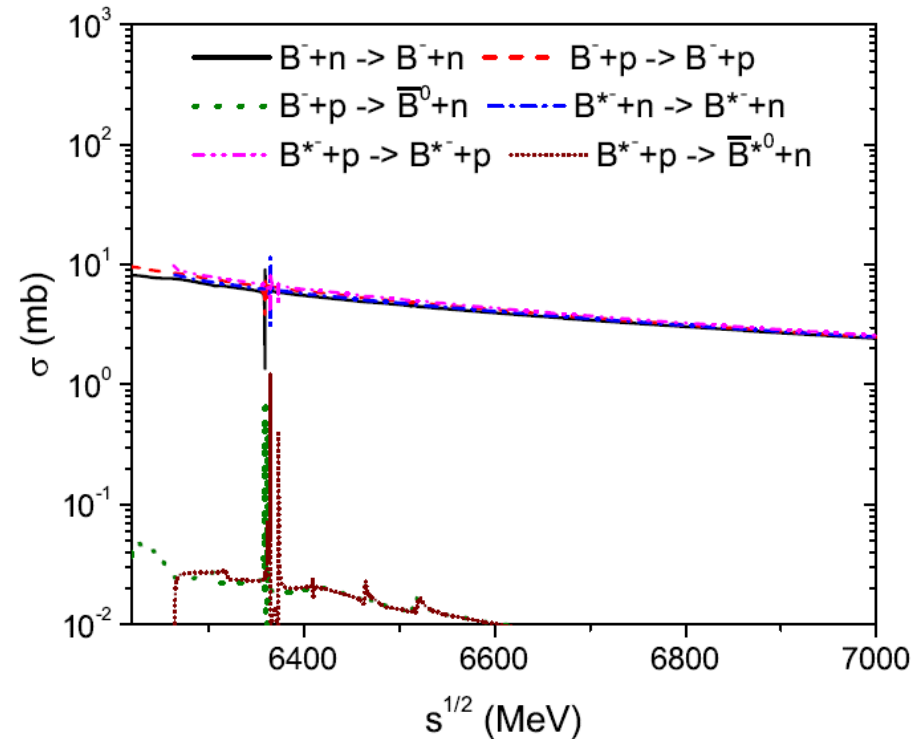
L. Tolos and J. M. Torres-Rincon, Phys. Rev. D 88, 074019 (2013)

J. M. Torres-Rincon, L. Tolos and O. Romanets, Phys. Rev. D 89, 074042 (2014)

1. B-meson scattering with mesons

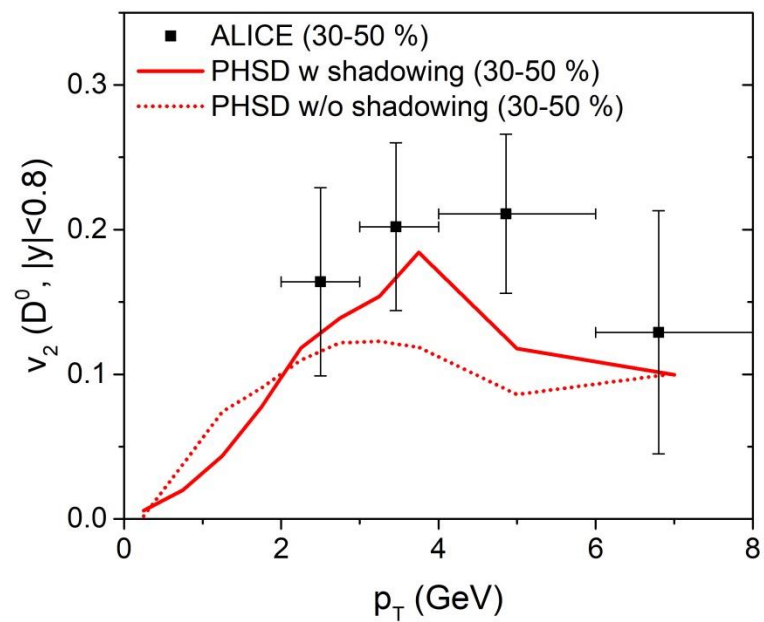
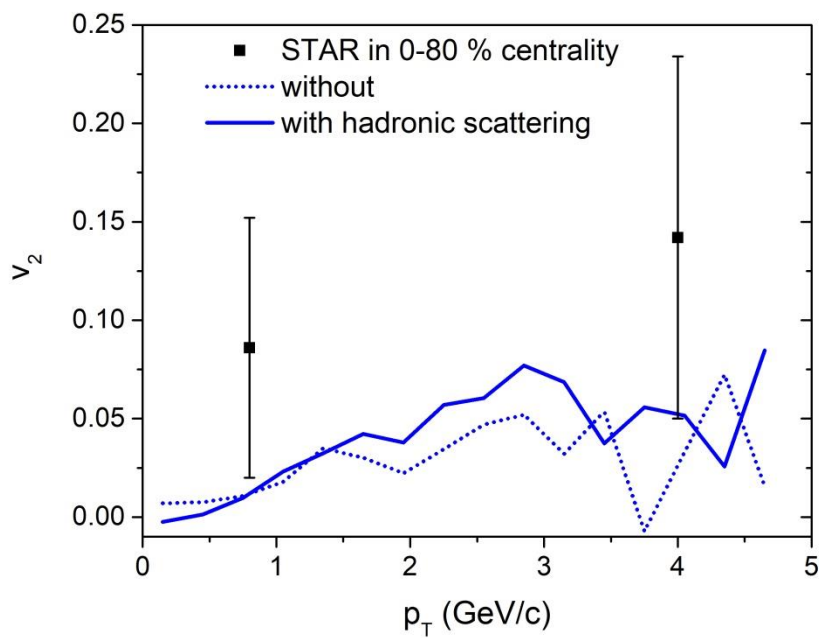
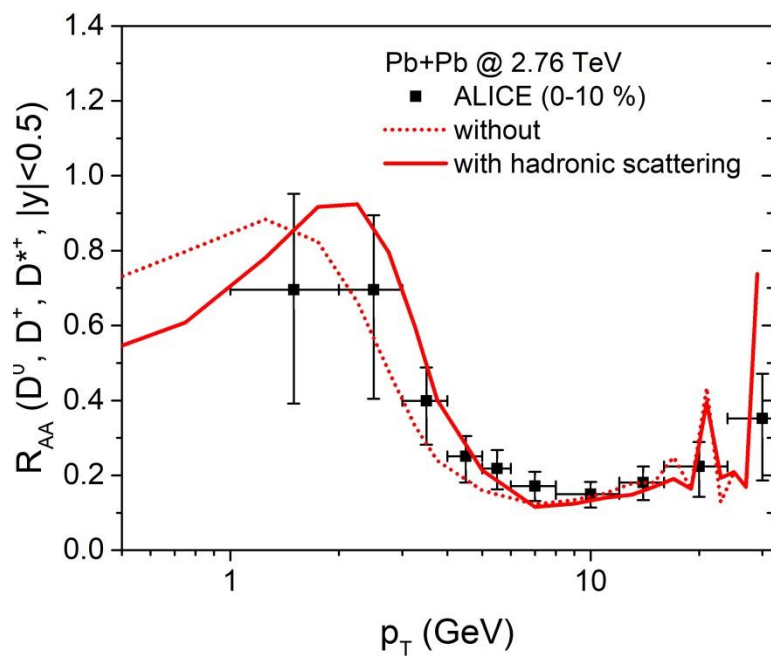
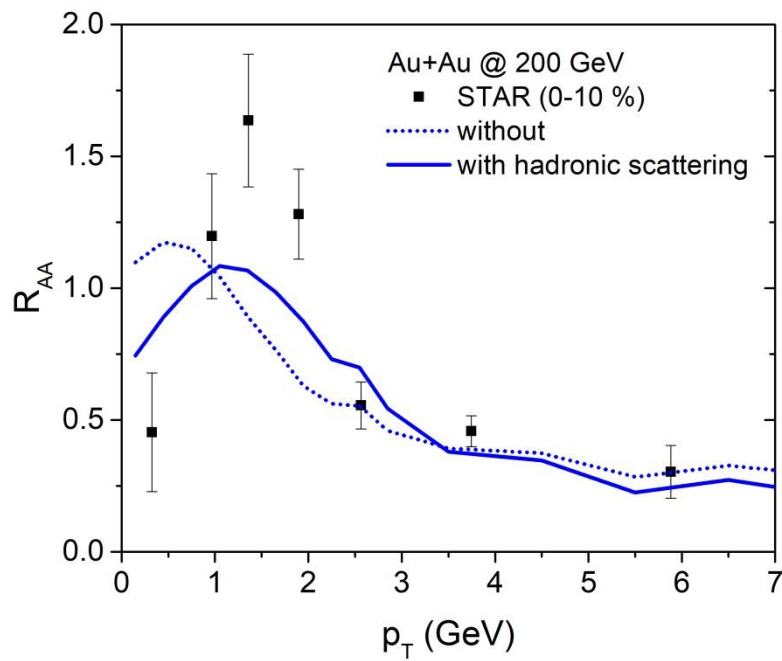


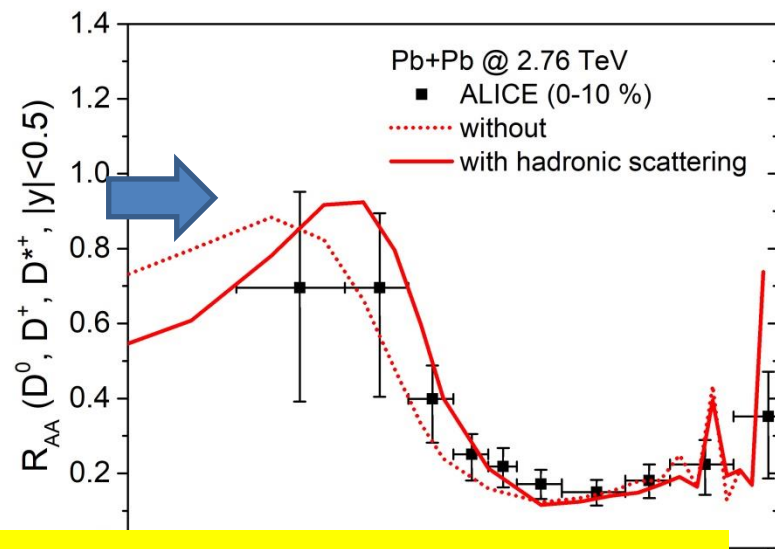
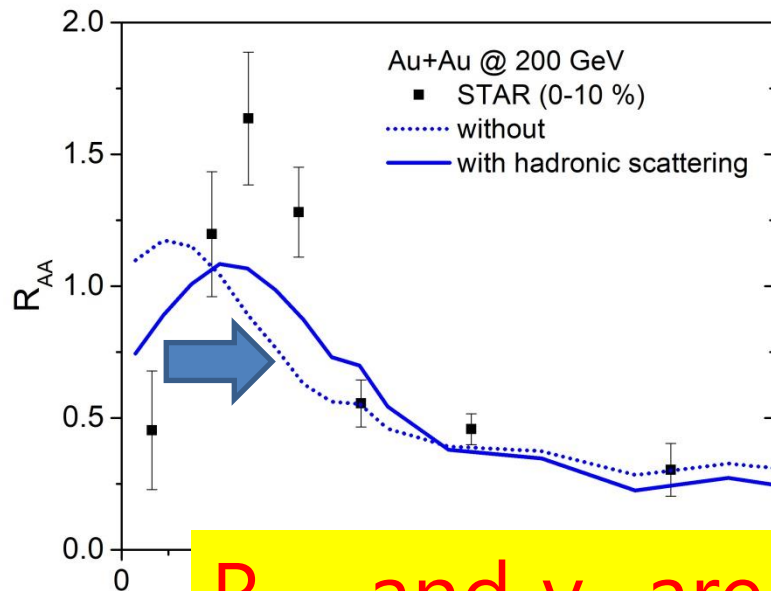
2. B-meson scattering with baryons



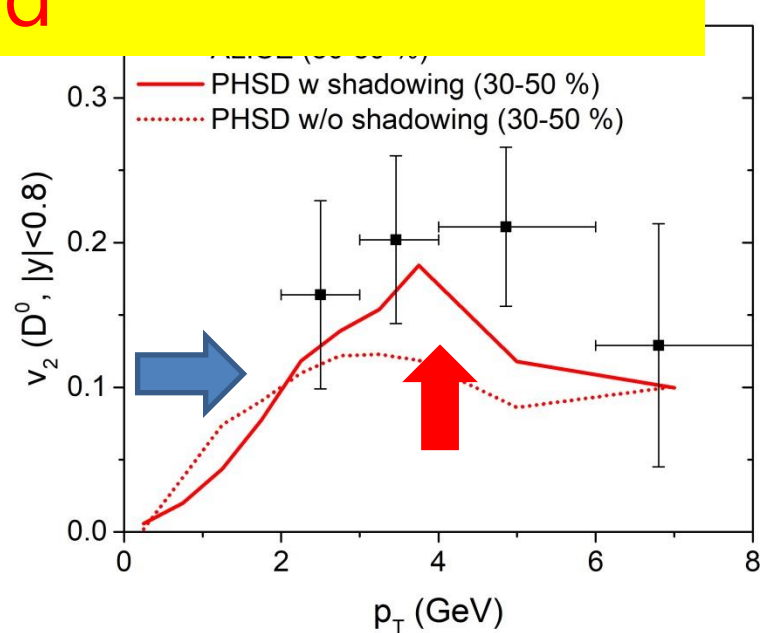
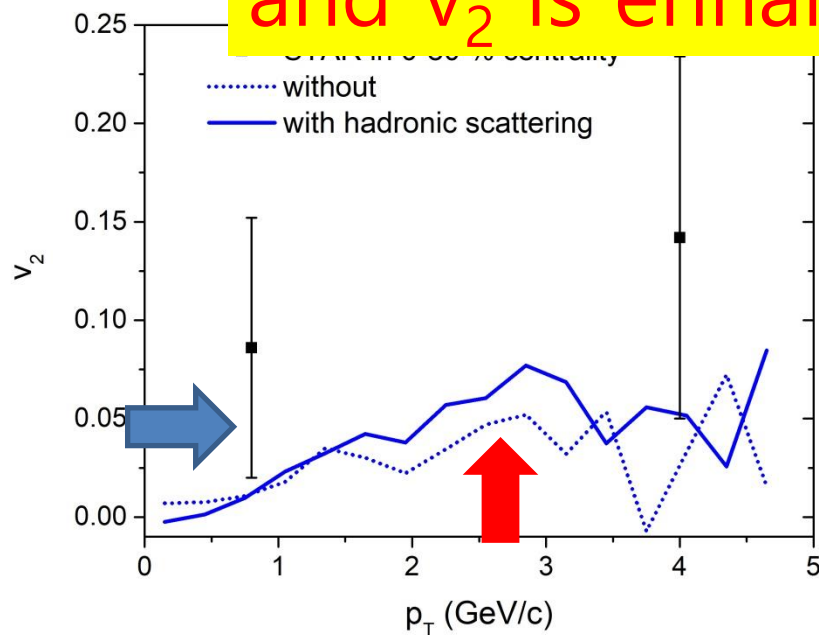
➤ **More than 200 hadronic channels** ➔ implemented in the PHSD

➤ thanks to Juan M. Torres-Rincon!





R_{AA} and v_2 are shifted to higher p_T
and v_2 is enhanced



Thank you for your attention!