Hadron tomography and gravitational form factors

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Workshop on Mass in the Standard Model and Consequences of its Emergence

Online, ECT*, Trento, Italy, April 19-23, 2021

https://indico.ectstar.eu/event/97/timetable/#20210419.detailed

April 20, 2021

Contents

References on my GPD-related works

- **Possible GPD studies at hadron accelerator facilities**
  Novel two-to-three hard hadronic processes and possible studies of generalized parton distributions at hadron facilities,
  S. Kumano, M. Strikman, K. Sudoh, Phys. Rev. D 80 (2009) 074003.
- **Possible GPD measurement by exclusive Drell-Yan at J-PARC**
  Accessing proton generalized parton distributions and pion distribution amplitudes with the exclusive pion-induced Drell-Yan process at J-PARC,
  T. Sawada, Wen-Chen Chang, S. Kumano, Jen-Chieh Peng, S. Sawada, K. Tanaka, Phys. Rev. D 93 (2016) 114034.
- **Letter of Intent for the 27th J-PARC PAC meeting, January 16 - 18 January, 2019**
  Studying Generalized Parton Distributions with Exclusive Drell-Yan process at J-PARC,
  JungKeun Ahn *et al*. (S. Kumano 10th author),
  Contact persons: W. C. Chang, H. Noumi, S. Sawada

GPD studies at hadron facilities (J-PARC, NICA, GSI-FAIR, . . .)

- **GPDs for exotic hadrons**
  Tomography of exotic hadrons in high-energy exclusive processes,
  H. Kawamura, S. Kumano, Phys. Rev. D 89 (2014) 054007.
- **Timelike GPDs (GDAs) and KEKB-data analysis for gravitational form factors**
  Hadron tomography by generalized distribution amplitudes in pion-pair production process $\gamma^* \gamma \to \pi^0 \pi^0$ and gravitational form factors for pion,
  S. Kumano, Qin-Tao Song, O. V. Teryaev, Phys. Rev.D 97 (2018) 014020.

GPD studies at e^+e^- facilities (KEKB, ILC, CEPC, . . .)

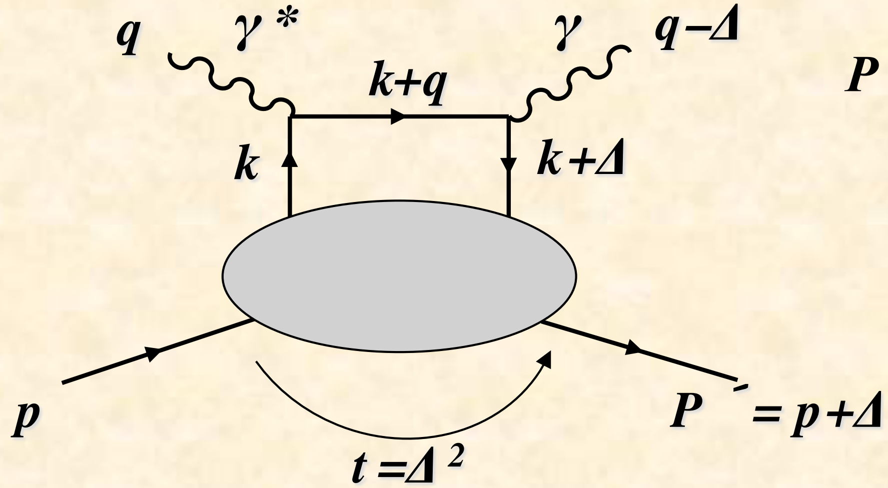
- **High-energy neutrino interactions and GPDs**
  High-energy neutrino-nucleus interactions,
  S. Kumano, EPJ Web Conf. 208 (2019) 07003.
- **Synergies between EIC (Electron-Ion Collider) project and neutrino reactions**
  EIC yellow report, R. Abdul Khalek *et al*. (S. Kumano 150th author;
  Sec. 7.5.2, Neutrino physics by S. Kumano and R. Petti), arXiv:2103.05419.

GPD studies at neutrino facilities (LBNF, . . .)

Introduction:

Origins of nucleon spin and mass

Generalized Parton Distributions (GPDs)



$$P = \frac{p+p'}{2},\quad \Delta = p' - p$$

Bjorken variable $x = \dfrac{Q^2}{2p\cdot q}$

Momentum transfer squared $t = \Delta^2$

Skewdness parameter $\xi = \dfrac{p^+ - p'^+}{p^+ + p'^+} = -\dfrac{\Delta^+}{2P^+}$

GPDs are defined as correlation of off-forward matrix:

$$\int \frac{dz^-}{4\pi} e^{ixP^+z^-} \langle p'|\bar\psi(-z/2)\gamma^+\psi(z/2)|p\rangle\Big|_{z^+=0,\vec z_\perp=0} = \frac{1}{2P^+}\left[H(x,\xi,t)\bar u(p')\gamma^+ u(p) + E(x,\xi,t)\bar u(p')\frac{i\sigma^{+\alpha}\Delta_\alpha}{2M}u(p)\right]$$

$$\int \frac{dz^-}{4\pi} e^{ixP^+z^-} \langle p'|\bar\psi(-z/2)\gamma^+\gamma_5\psi(z/2)|p\rangle\Big|_{z^+=0,\vec z_\perp=0} = \frac{1}{2P^+}\left[\tilde H(x,\xi,t)\bar u(p')\gamma^+\gamma_5 u(p) + \tilde E(x,\xi,t)\bar u(p')\frac{\gamma_5\Delta^+}{2M}u(p)\right]$$

Forward limit: PDFs $H(x,\xi,t)\big|_{\xi=t=0} = f(x),\quad \tilde H(x,\xi,t)\big|_{\xi=t=0} = \Delta f(x),$

First moments: Form factors

Dirac and Pauli form factors F_1, F_2 $\quad \int_{-1}^{1} dx\, H(x,\xi,t) = F_1(t),\quad \int_{-1}^{1} dx\, E(x,\xi,t) = F_2(t)$

Axial and Pseudoscalar form factors G_A, G_P $\quad \int_{-1}^{1} dx\, \tilde H(x,\xi,t) = g_A(t),\quad \int_{-1}^{1} dx\, \tilde E(x,\xi,t) = g_P(t)$

Second moments: Angular momenta

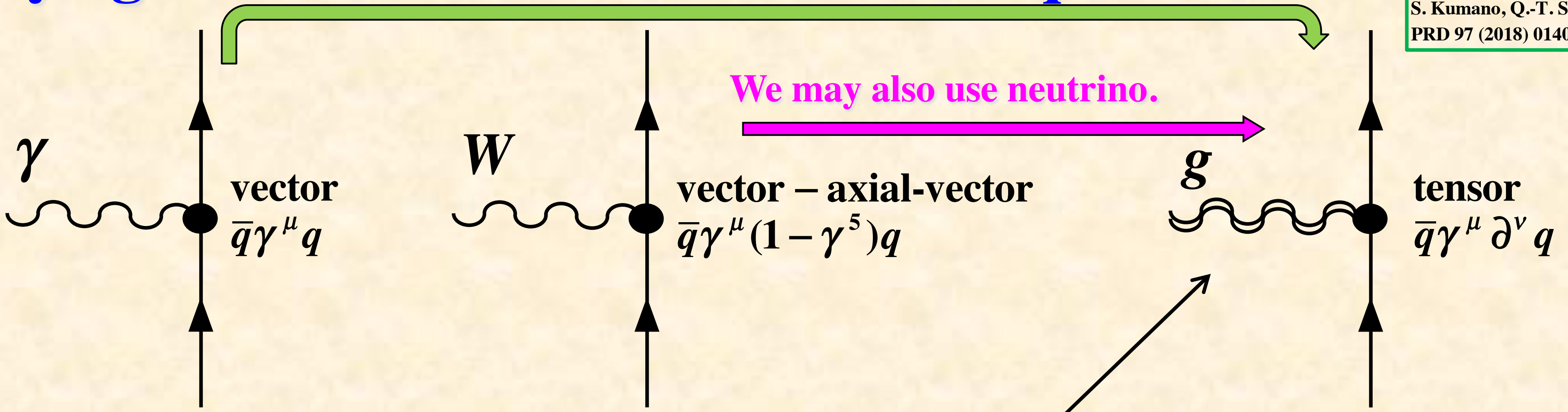
Sum rule: $J_q = \dfrac{1}{2}\int_{-1}^{1} dx\, x\left[H_q(x,\xi,t=0) + E_q(x,\xi,t=0)\right],\quad J_q = \dfrac{1}{2}\Delta q + L_q$

$\Rightarrow$ probe L_q, key quantity to solve the spin puzzle!

Why "gravitational" interactions with quarks

We studied in 2017-2018.

S. Kumano, Q.-T. Song, O. Teryaev, PRD 97 (2018) 014020.



It is possible to probe gravitational sources in the microscopic level without gravitons.

GPDs (Generalized Parton Distributions), GDAs (Generalized Distribution Amplitudes) = timelike GPDs

$$\int \frac{dz^-}{4\pi} e^{ixP^+z^-} \langle p' | \bar{q}(-z/2)\gamma^+ q(z/2) | p \rangle \Big|_{z^+=0,\vec{z}_\perp=0} = \frac{1}{2P^+}\left[H(x,\xi,t)\bar{u}(p')\gamma^+ u(p) + E(x,\xi,t)\bar{u}(p')\frac{i\sigma^{+\alpha}\Delta_\alpha}{2M}u(p)\right]$$

Non-local operator of GPDs/GDAs:

$$(P^+)^n \int dx x^{n-1} \int \frac{dz^-}{2\pi} e^{ixP^+z^-} \left[\bar{q}(-z/2)\gamma^+ q(z/2)\right]_{z^+=0,\vec{z}_\perp=0}$$

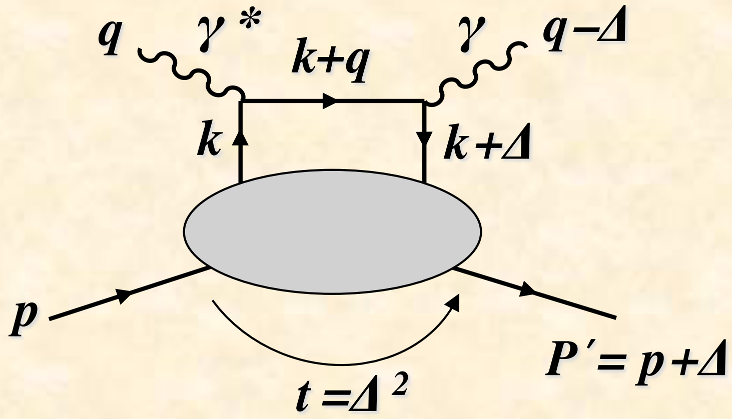
$$= \left(i\frac{\partial}{\partial z^-}\right)^{n-1} \left[\bar{q}(-z/2)\gamma^+ q(z/2)\right]_{z=0}$$

$$= \bar{q}(0)\gamma^+ \left(i\overleftrightarrow{\partial}^+\right)^{n-1} q(0)$$

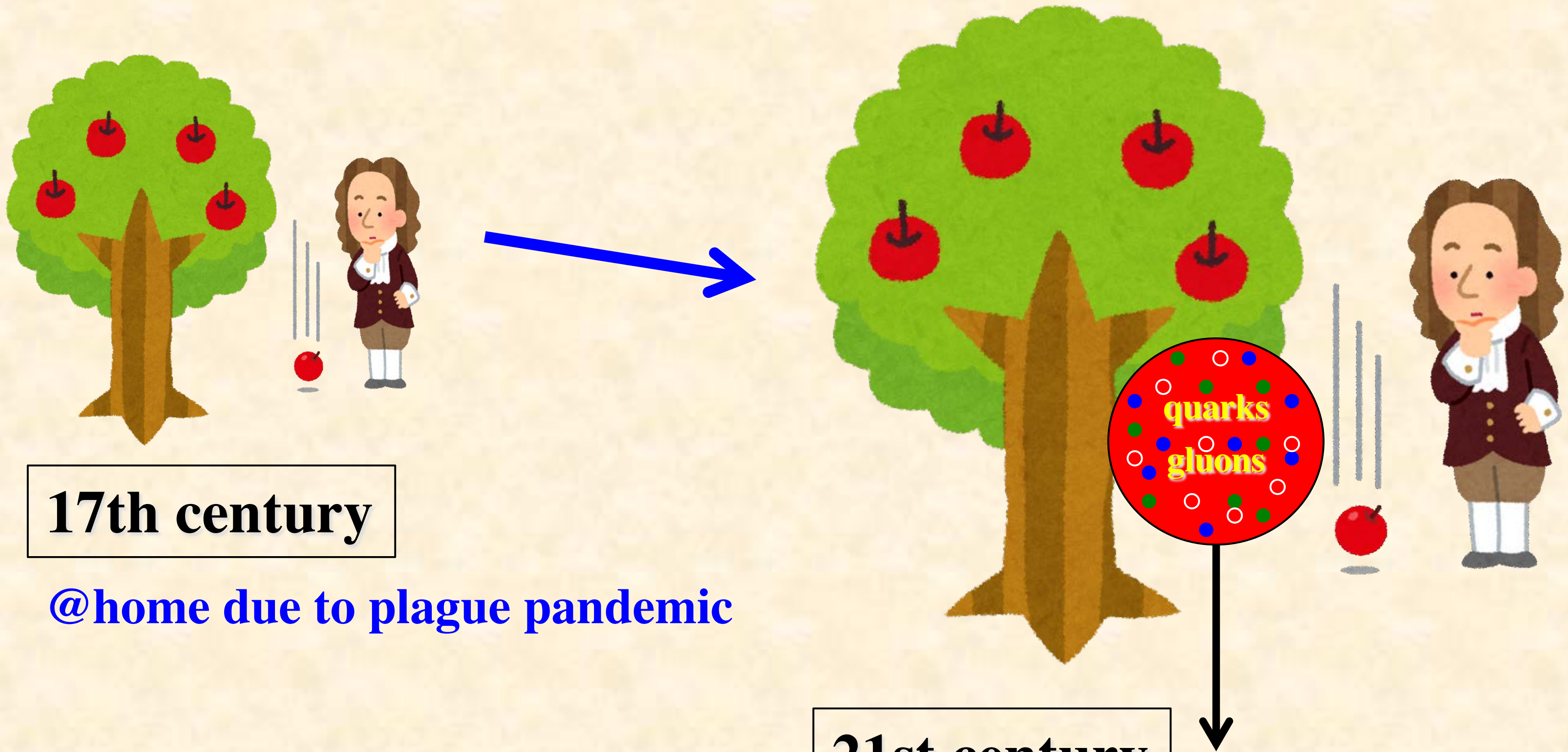
= energy-momentum tensor of a quark for $n = 2$
(electromagnetic for $n = 1$)

= source of gravitiy

Virtual Compton
or (timelike) two-photon process



Time has come to understand the gravitational sources in microscopic (instead of usual macroscopic) world in terms of quark and gluon degrees of freedom.



17th century

@home due to plague pandemic

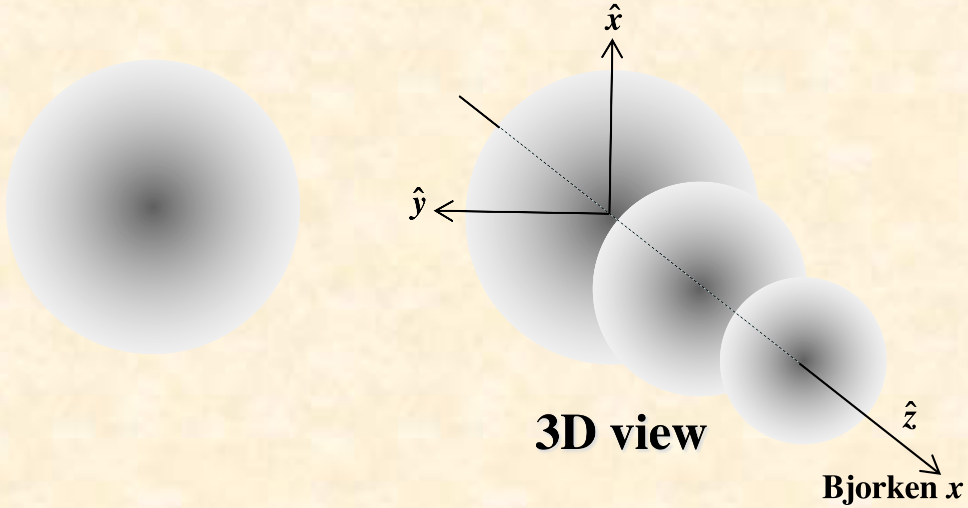
21st century

@home due to coronavirus pandemic

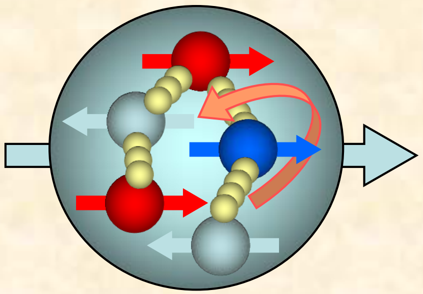
Proton (hadrons) puzzle studies by hadron tomography

Hadron tomography

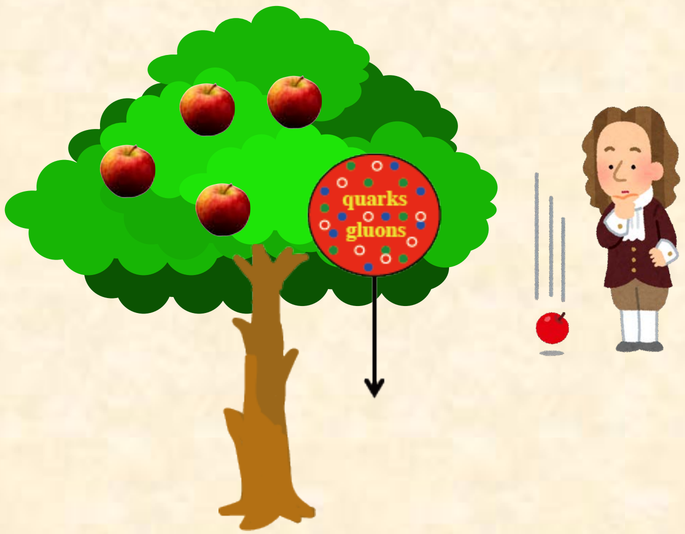
Proton radius puzzle



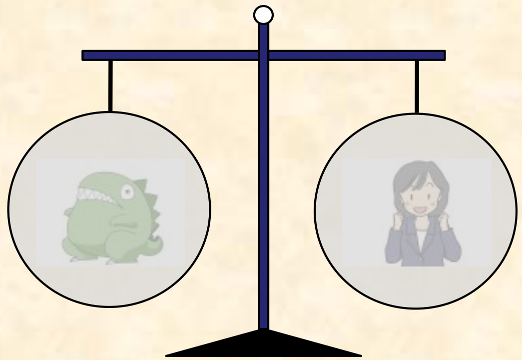
Origin of nucleon spin



Source of gravity (mass)

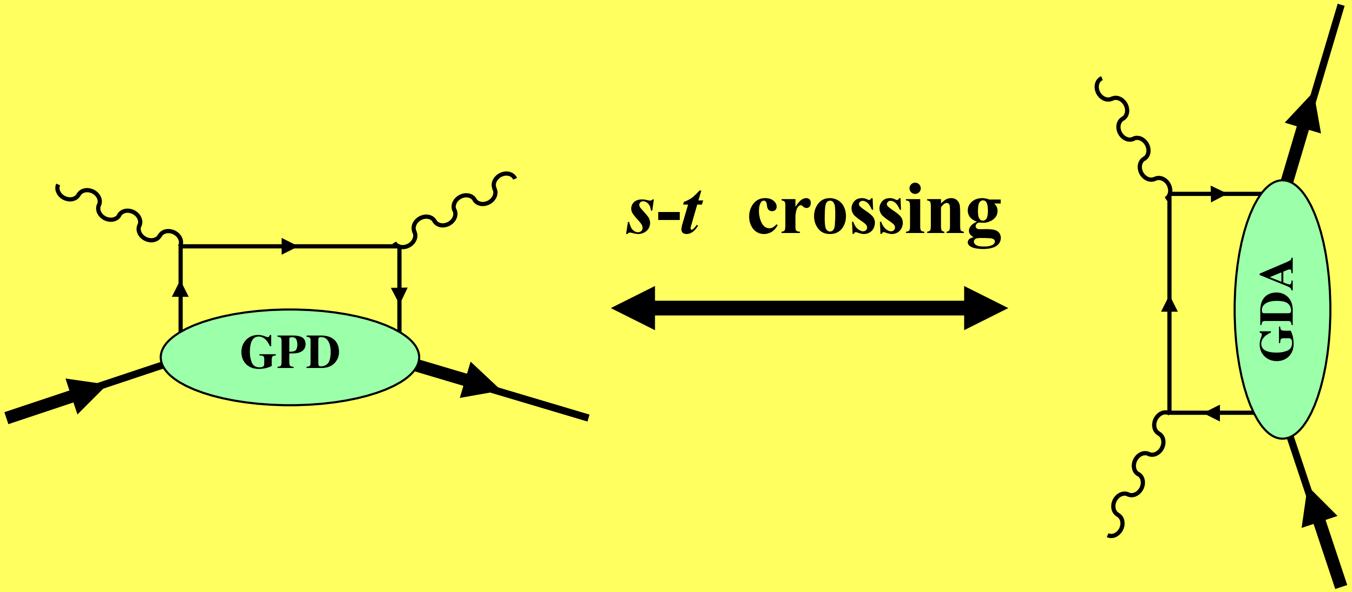


Exotic hadrons



Generalized Distribution Amplitudes (GDAs) and extraction of gravitational form factors from KEKB data

GDA = Timelike GPDs



S. Kumano, Q.-T. Song, O. Teryaev, Phys. Rev. D 97 (2018) 014020.

GPD $H_q^h(x,\xi,t)$ and GDA $\Phi_q^{hh}(z,\zeta,W^2)$

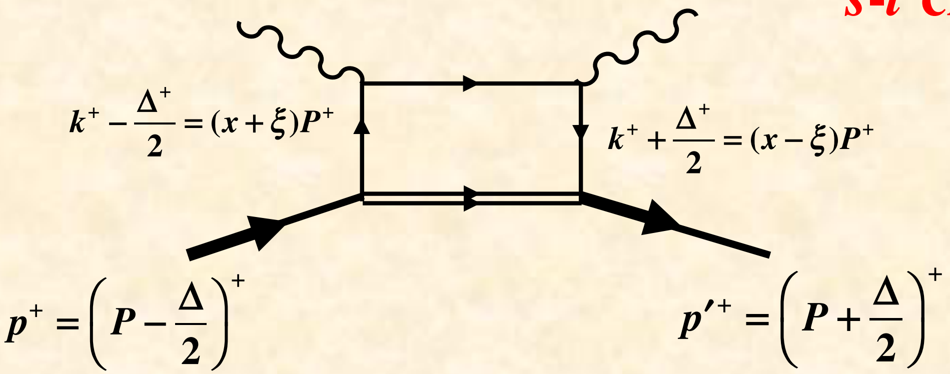
GPD: $H_q(x,\xi,t) = \int \frac{dy^-}{4\pi} e^{ixP^+y^-} \langle h(p') | \bar{\psi}(-y/2)\gamma^+\psi(y/2) | h(p) \rangle \Big|_{y^+=0,\vec{y}_\perp=0}$, $\quad P^+ = \frac{(p+p')^+}{2}$

GDA: $\Phi_q(z,\zeta,s) = \int \frac{dy^-}{2\pi} e^{izP^+y^-} \langle h(p)\bar{h}(p') | \bar{\psi}(-y/2)\gamma^+\psi(y/2) | 0 \rangle \Big|_{y^+=0,\vec{y}_\perp=0}$

DA: $\Phi_q^\pi(z,\zeta,s) = \int \frac{dy^-}{2\pi} e^{izP^+y^-} \langle \pi(p) | \bar{\psi}(-y/2)\gamma^+\gamma_5\psi(y/2) | 0 \rangle \Big|_{y^+=0,\vec{y}_\perp=0}$

$H_q^h(x,\xi,t)$ ⟷ $\Phi_q^{hh}(z,\zeta,W^2)$

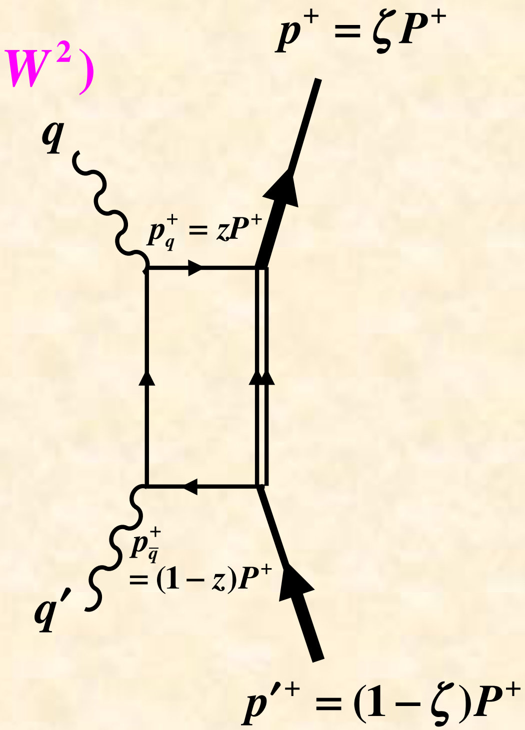
***s*-*t* crossing**



$z \Leftrightarrow \frac{1 - x/\xi}{2}$

$\zeta \Leftrightarrow \frac{1 - 1/\xi}{2}$

$W^2 \Leftrightarrow t$



JLab / COMPASS

KEKB

$P = \frac{p+p'}{2}, \quad \Delta = p' - p$

Bjorken variable: $x = \frac{Q^2}{2p \cdot q}$

Momentum transfer squared: $t = \Delta^2$

Skewdness parameter: $\xi = \frac{p^+ - p'^+}{p^+ + p'^+} = -\frac{\Delta^+}{2P^+}$

Bjorken variable for $\gamma\gamma^*$: $z = \frac{Q^2}{2q \cdot q'}$

Light-cone momentum ratio for a hadron in $h\bar{h}$: $\zeta = \frac{p^+}{P^+} = \frac{1 + \beta\cos\theta}{2}$

Invariant mass of $h\bar{h}$: $W^2 = (p + p')^2$

Cross section for $\gamma^*\gamma \to \pi^0\pi^0$

$$\frac{d\sigma}{d(\cos\theta)} = \frac{1}{16\pi(s+Q^2)}\sqrt{1-\frac{4m_\pi^2}{s}}\,\overline{\sum_{\lambda,\lambda'}}|\mathcal{M}|^2$$

$$\mathcal{M} = \varepsilon_\mu^\lambda(q)\varepsilon_\nu^{\lambda'}(q')T^{\mu\nu} = e^2 A_{\lambda\lambda'},\quad T^{\mu\nu} = i\int d^4\xi\, e^{-i\xi\cdot q}\left\langle \pi(p)\pi(p')\right| T J_{em}^\mu(\xi) J_{em}^\nu(0)\left|0\right\rangle$$

$$A_{\lambda\lambda'} = \frac{1}{e^2}\varepsilon_\mu^\lambda(q)\varepsilon_\nu^{\lambda'}(q')T^{\mu\nu} = -\varepsilon_\mu^\lambda(q)\varepsilon_\nu^{\lambda'}(q')g_T^{\mu\nu}\sum_q \frac{e_q^2}{2}\int_0^1 dz\frac{2z-1}{z(1-z)}\Phi_q^{\pi\pi}(z,\zeta,W^2)$$

GDA: $\Phi_q^{\pi\pi}(z,\zeta,s) = \int \frac{dy^-}{2\pi}e^{izP^+y^-}\left\langle \pi(p)\pi(p')\right|\overline{\psi}(-y/2)\gamma^+\psi(y/2)\left|0\right\rangle\Big|_{y^+=0,\vec{y}_\perp=0}$

$$\frac{d\sigma}{d(\cos\theta)} \simeq \frac{\pi\alpha^2}{4(s+Q^2)}\sqrt{1-\frac{4m_\pi^2}{s}}|A_{++}|^2,\quad A_{++} = \sum_q \frac{e_q^2}{2}\int_0^1 dz\frac{2z-1}{z(1-z)}\Phi_q^{\pi\pi}(z,\zeta,W^2)$$

- **Continuum: GDAs without intermediate-resonance contribution**

$$\Phi_q^{\pi\pi}(z,\zeta,W^2) = N_\pi z^\alpha(1-z)^\alpha(2z-1)\zeta(1-\zeta)F_q^\pi(s)$$

$$F_q^\pi(s) = \frac{1}{\left[1+(s-4m_\pi^2)/\Lambda^2\right]^{n-1}},\quad n=2 \text{ according to constituent counting rule}$$

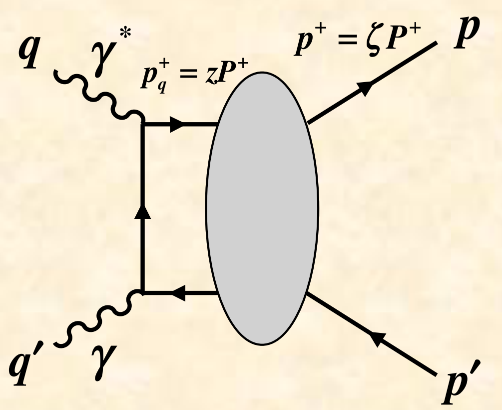
- **Resonances: There exist resonance contributions to the cross section.**

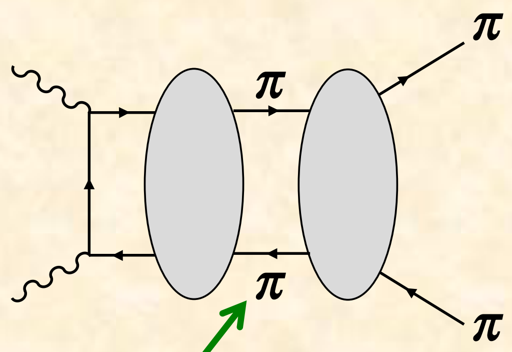
$$\sum_q \Phi_q^{\pi\pi}(z,\zeta,W^2) = 18N_f z^\alpha(1-z)^\alpha(2z-1)\left[\tilde{B}_{10}(W)+\tilde{B}_{12}(W)P_2(\cos\theta)\right]$$

$$P_2(x) = \frac{1}{2}(3x^2-1)$$

$\tilde{B}_{10}(W)$ = resonance $[f_0(500), f_0(980)]$ + continuum

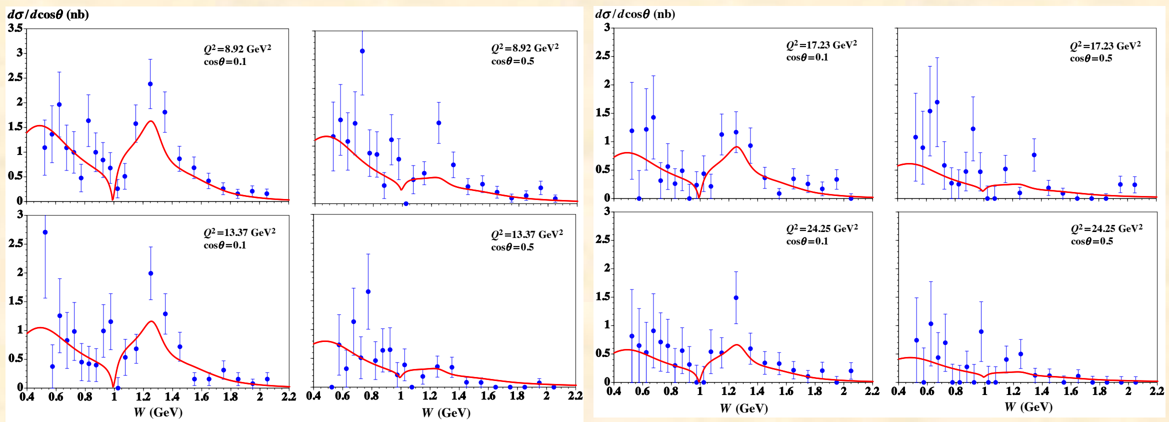
$\tilde{B}_{12}(W)$ = resonance $[f_2(1270)]$ + continuum





Including intermediate resonance contributions

$Q^2 = 8.92,\ 13.37\ \text{GeV}^2$ $\qquad$ $Q^2 = 17.23,\ 24.25\ \text{GeV}^2$



Belle measurements:
M. Masuda *et al.*,
PRD93 (2016) 032003.

Gravitational form factors and radii for pion

$$\int_0^1 dz(2z-1)\Phi_q^{\pi^0\pi^0}(z,\zeta,s) = \frac{2}{(P^+)^2}\left\langle \pi^0(p)\pi^0(p')\middle| T_q^{++}(0)\middle|0\right\rangle$$

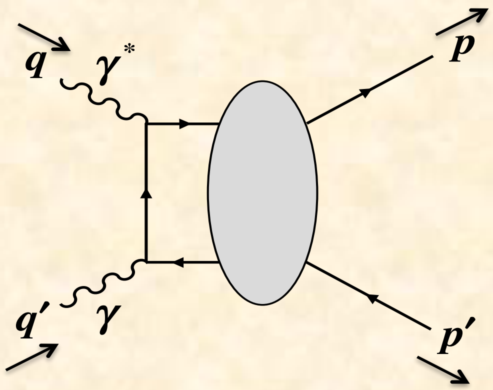
$$\left\langle \pi^0(p)\pi^0(p')\middle| T_q^{\mu\nu}(0)\middle|0\right\rangle = \frac{1}{2}\left[\left(sg^{\mu\nu} - P^\mu P^\nu\right)\Theta_{1,q}(s) + \Delta^\mu\Delta^\nu\Theta_{2,q}(s)\right]$$

$$P = \frac{p+p'}{2}, \qquad \Delta = p' - p$$

See also Hyeon-Dong Son, Hyun-Chul Kim, PRD90 (2014) 111901.

$T_q^{\mu\nu}$: energy-momentum tensor for quark

$\Theta_{1,q}$, $\Theta_{2,q}$: gravitational form factos for pion



Analyiss of $\gamma^*\gamma \to \pi^0\pi^0$ cross section

⇒ Generalized distribution amplitudes $\Phi_q^{\pi^0\pi^0}(z,\zeta,s)$

⇒ Timelike gravitational form factors $\Theta_{1,q}(s)$, $\Theta_{2,q}(s)$

⇒ Spacelike gravitational form factors $\Theta_{1,q}(t)$, $\Theta_{2,q}(t)$

⇒ Gravitational radii of pion

Gravitational form factors:

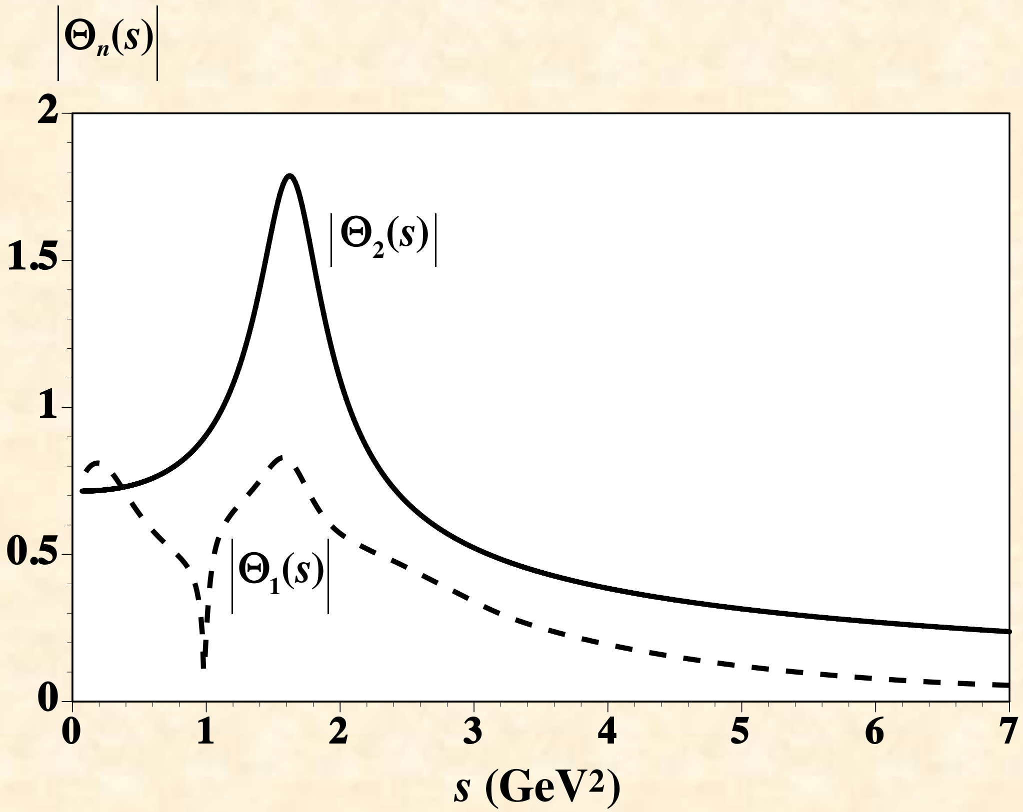
Original definition: H. Pagels, Phys. Rev. 144 (1966) 1250.

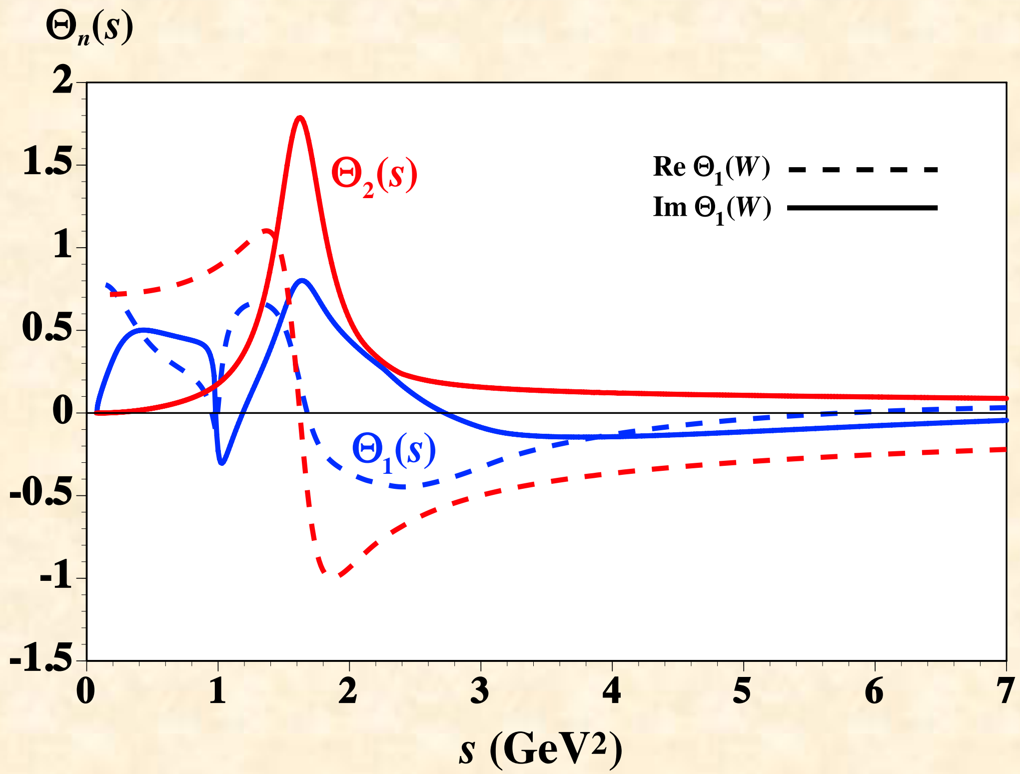
Operator relations: K. Tanaka, Phys. Rev. D 98 (2018) 034009.

Timelike gravitational form factors for pion

$$\left\langle \pi^a(p)\pi^b(p')\left|T_q^{\mu\nu}(0)\right|0\right\rangle = \frac{\delta^{ab}}{2}\left[(sg^{\mu\nu}-P^\mu P^\nu)\Theta_{1(q)}(s)+\Delta^\mu\Delta^\nu\Theta_{2(q)}(s)\right], \quad P=p+p', \quad \Delta=p'-p$$

- $\Theta_{1(q)}(s) = -\frac{3}{10}\tilde{B}_{10}(W^2)+\frac{3}{20}\tilde{B}_{12}(W^2) = -4B_{(q)}(s)$
- $\Theta_{2(q)}(s) = \frac{9}{20\beta^2}\tilde{B}_{12}(W^2) = A_{(q)}(s)$





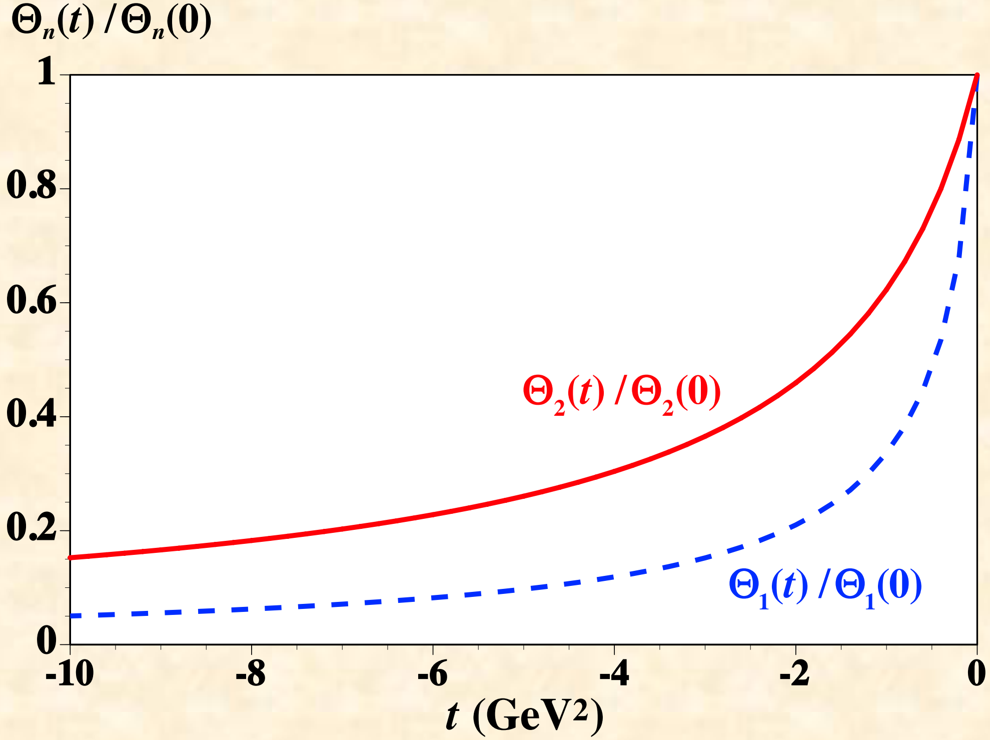
Spacelike gravitational form factors and radii for pion

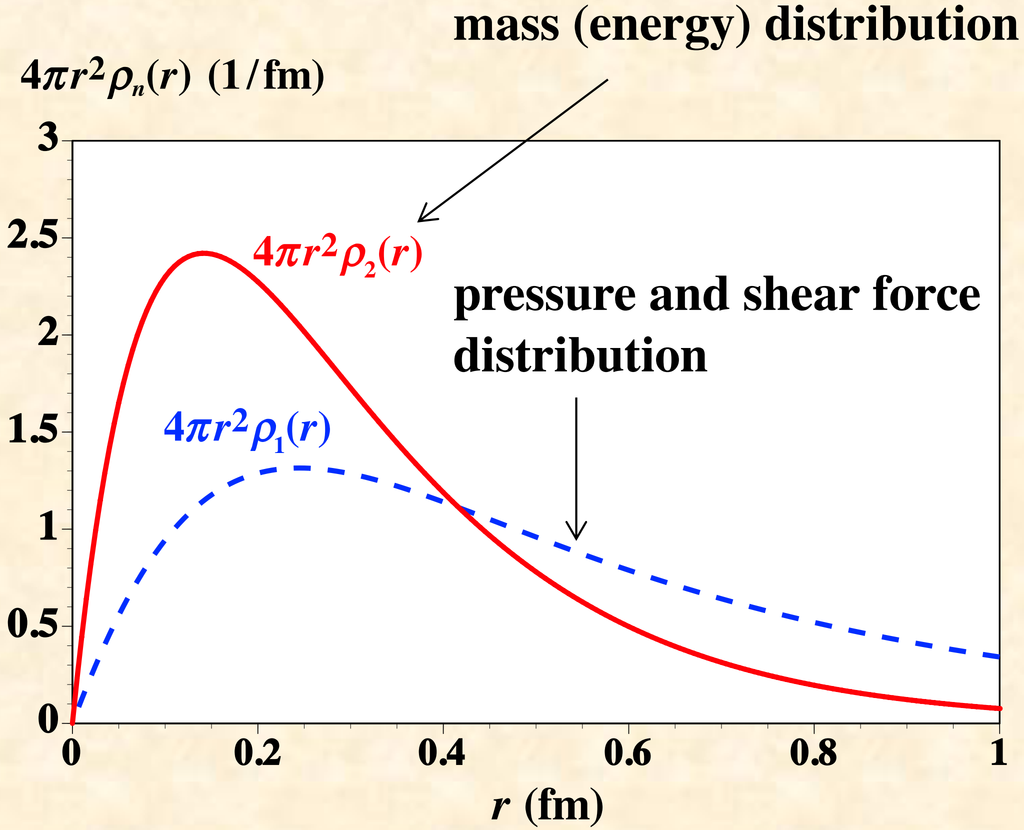
$$F(s)=\Theta_1(s),\ \Theta_1(s),\quad F(t)=\int_{4m_\pi^2}^{\infty} ds\,\frac{\operatorname{Im}F(s)}{\pi(s-t-i\varepsilon)},\quad \rho(r)=\frac{1}{(2\pi)^3}\int d^3q\,e^{-i\vec{q}\cdot\vec{r}}F(q)=\frac{1}{4\pi^2}\frac{1}{r}\int_{4m_\pi^2}^{\infty} ds\,e^{-\sqrt{s}r}\operatorname{Im}F(s)$$

This is the first report on gravitational radii of hadrons from actual experimental measurements.

$\sqrt{\langle r^2\rangle_{\text{mass}}}=0.32\sim0.39$ fm, $\sqrt{\langle r^2\rangle_{\text{mech}}}=0.82\sim0.88$ fm ← First finding on gravitational radius from actual experimental measurements

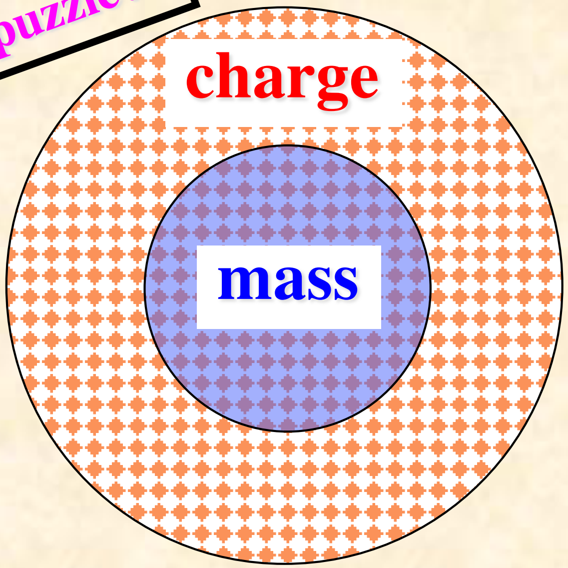
$\Leftrightarrow\ \sqrt{\langle r^2\rangle_{\text{charge}}}=0.672\pm0.008$ fm





Hadron mass radius puzzle?

Hadron-mass radius puzzle?!



For pion

$$\sqrt{\langle r^2 \rangle_{\text{mass}}} = 0.32 \sim 0.39 \text{ fm} \iff \sqrt{\langle r^2 \rangle_{\text{charge}}} = 0.672 \pm 0.008 \text{ fm}$$

S. Kumano, Q.-T. Song, O. Teryaev, PRD 97 (2018) 014020;
Erratum in v3 of arXiv:1711.08088.

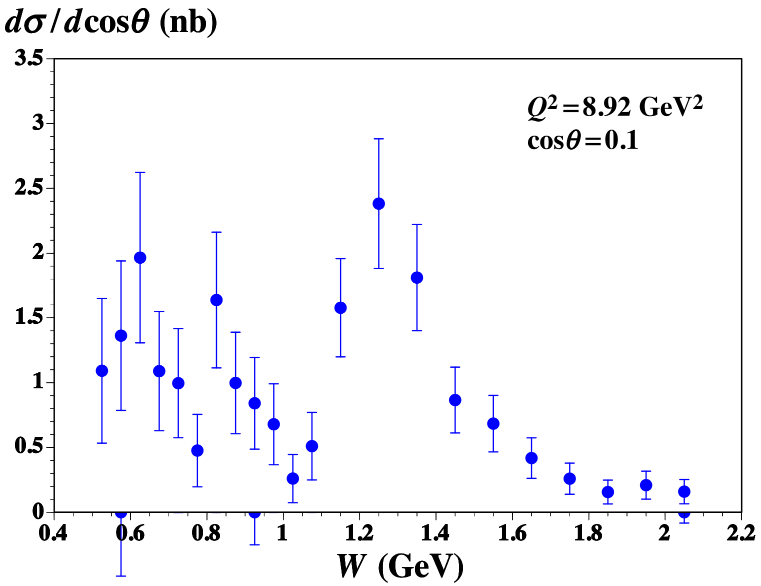
Mass radius seems to be much smaller than the charge radius for pion.

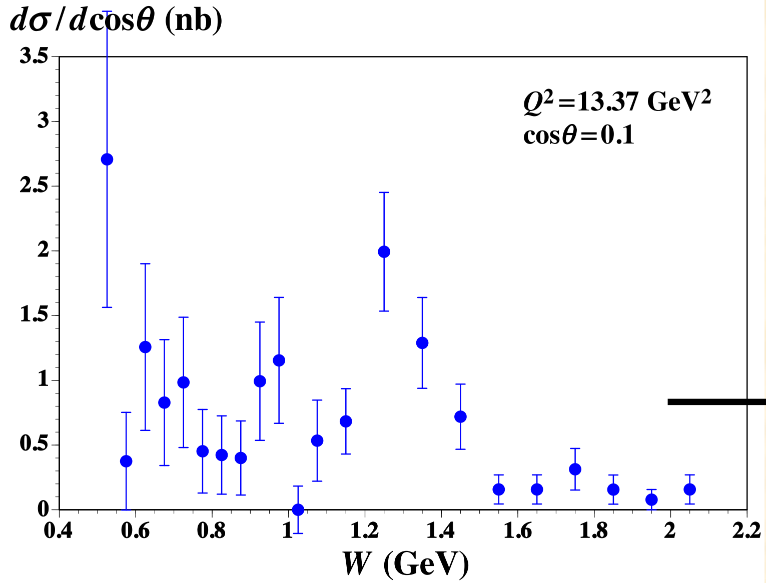
This is the first result on the mass radius from actual measurement,
so further studies are needed to find whether there is acutally a significant difference

Quarks contribute to both charge and mass distributions,
but gluons contribute to only the mass distribution.

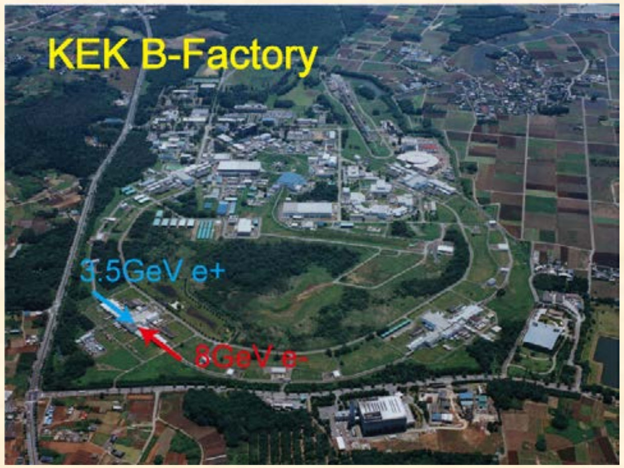
Electric interactions are repulsive (or could be attractive) and
gravitational interactions are always attractive,
so there would be some differences in both radii.
However, the difference of the factor of 2 may not be expected.

Super KEKB





The errors are dominated by statistical errors, and they will be significantly reduced by super-KEKB.

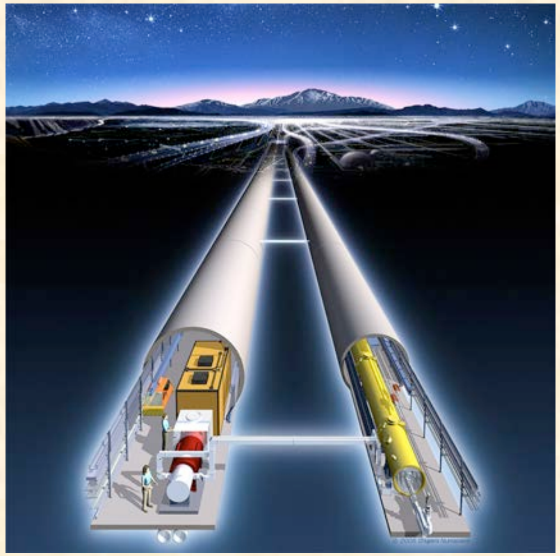


From KEKB to ILC

- Very Large Q^2
- Large W^2

for extracting GDAs

ILC



GSI-FAIR (PANDA)

arXiv:0903.3905 [hep-ex]

FAIR/PANDA/Physics Book i

Physics Performance Report for:

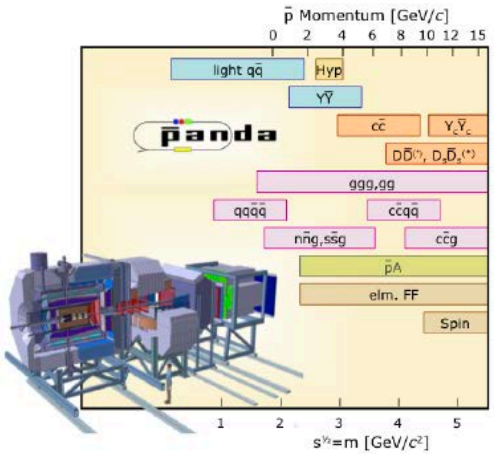
$\overline{\text{P}}$ANDA

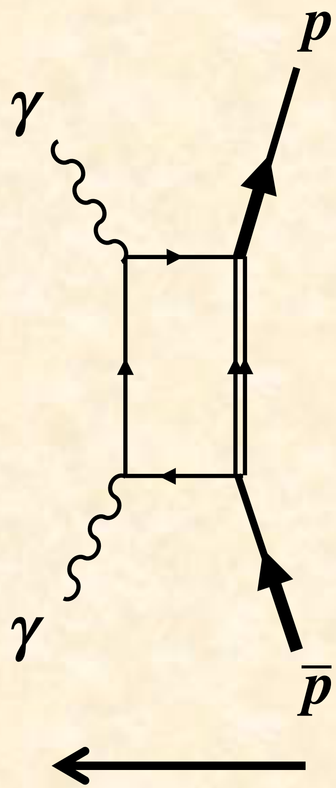
(AntiProton Annihilations at Darmstadt)

Strong Interaction Studies with Antiprotons

$\overline{\text{P}}$ANDA Collaboration

To study fundamental questions of hadron and nuclear physics in interactions of antiprotons with nucleons and nuclei, the universal $\overline{\text{P}}$ANDA detector will be build. Gluonic excitations, the physics of strange and charm quarks and nucleon structure studies will be performed with unprecedented accuracy thereby allowing high-precision tests of the strong interaction. The proposed $\overline{\text{P}}$ANDA detector is a state-of-the-art internal target detector at the HESR at FAIR allowing the detection and identification of neutral and charged particles generated within the relevant angular and energy range.
This report presents a summary of the physics accessible at $\overline{\text{P}}$ANDA and what performance can be expected.





GDAs for the proton!
(also @super-KEKB)

GPDs for exotic hadrons

H. Kawamura and S. Kumano
Phys. Rev. D 89 (2014) 054007.

Constituent counting rule for exotic hadrons:
H. Kawamura, S. Kumano, T. Sekihara, PRD 88 (2013) 034010;
W.-C. Chang, SK, and T. Sekihara, PRD 93 (2016) 034006.

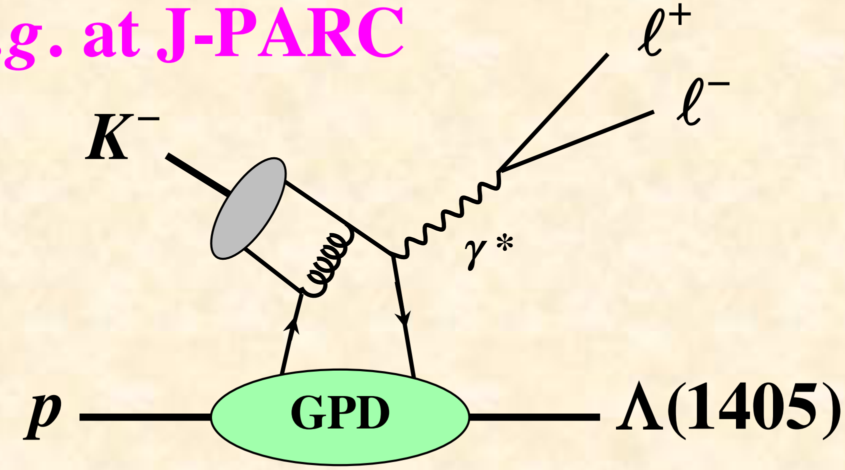
GPDs for exotic hadrons !?

Because stable targets do not exist for exotic hadrons,
it is not possible to measure their GPDs in a usual way.

→ Transition GPDs

or → $s \leftrightarrow t$ crossed qunatity = GDAs at KEKB, Linear Collider

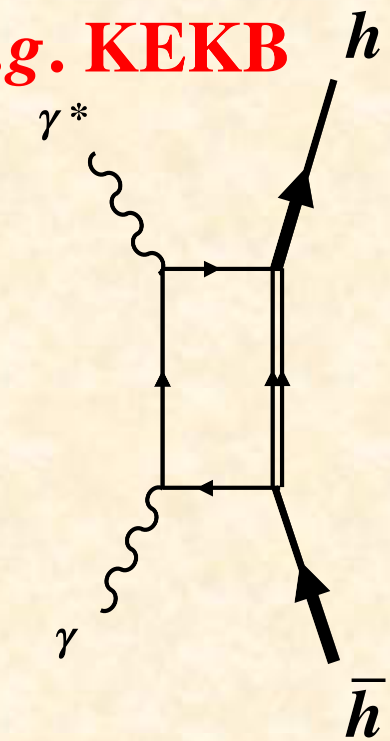
e.g. at J-PARC



$K^-(\bar{u}s) + p(uud) \rightarrow \Lambda_{1405}(uud\bar{u}s) + \gamma^*$

Λ_{1405} = pentaquark ($\bar{K}N$ molecule) candidate

e.g. KEKB



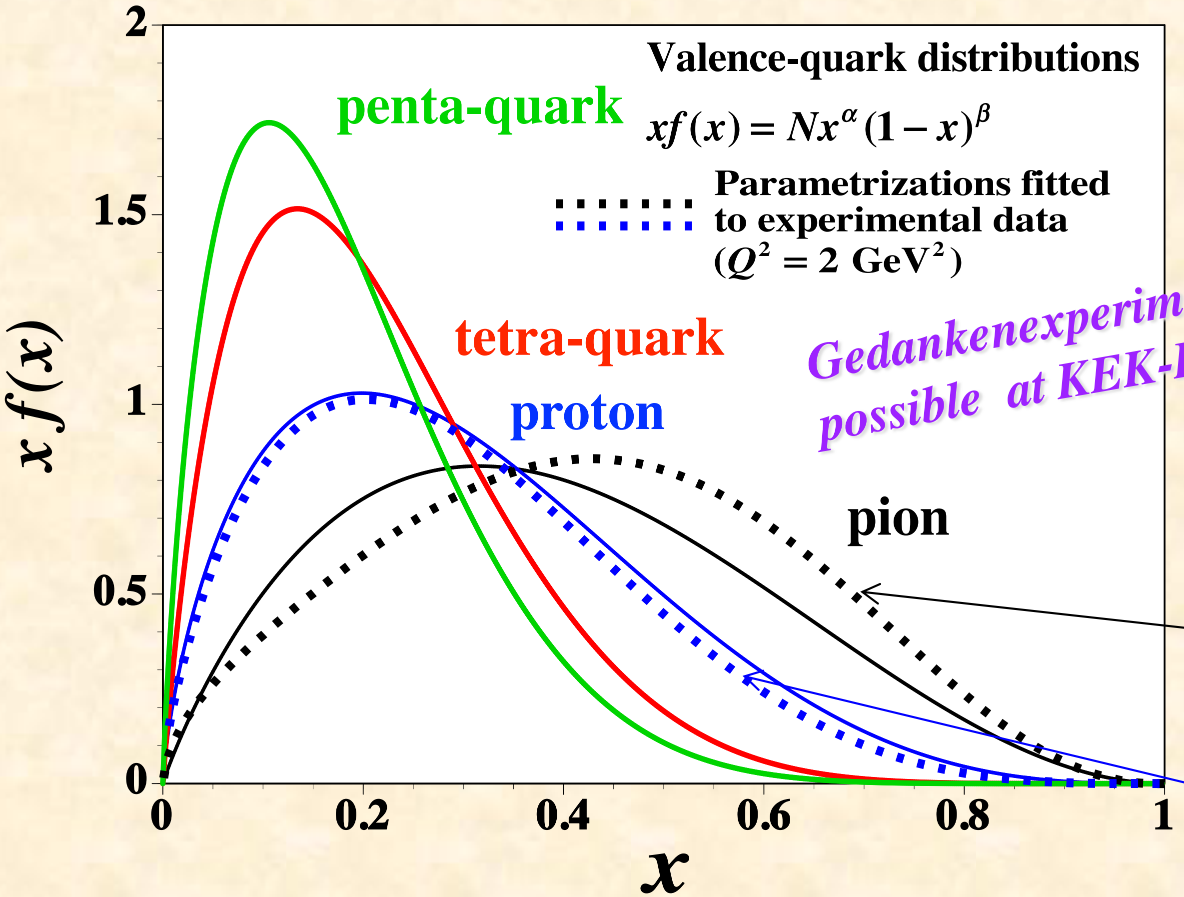
Simple function of GPDs

$$H_q^h(x,t) = f(x)F(t,x)$$

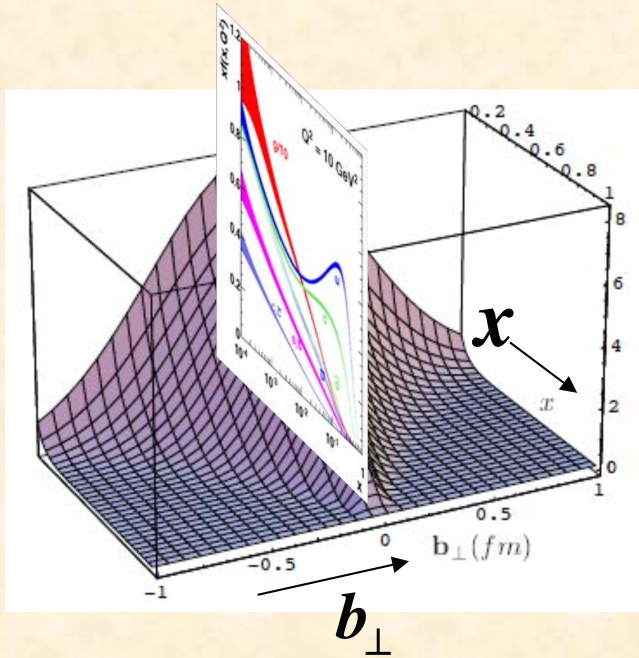
M. Guidal, M.V. Polyakov, A.V. Radyushkin, M. Vanderhaeghen, PRD 72, 054013 (2005).

Longitudinal-momentum distribution (PDF) for valence quarks: $f(x) = q_v(x) = c_n x^{\alpha_n}(1-x)^{\beta_n}$

- Valence-quark number sum rule (charge and baryon numbers): $\int_0^1 dx\, f(x) = n$
- Constituent conting rule at $x \to 1$: $\beta_n = 2n - 3 + 2\Delta S$ (n = number of constituents)
- Momentum carried by quarks $\langle x \rangle_q \simeq \int_0^1 dx\, x f(x)$



Gedankenexperiment, but possible at KEK-B, ILC, ...

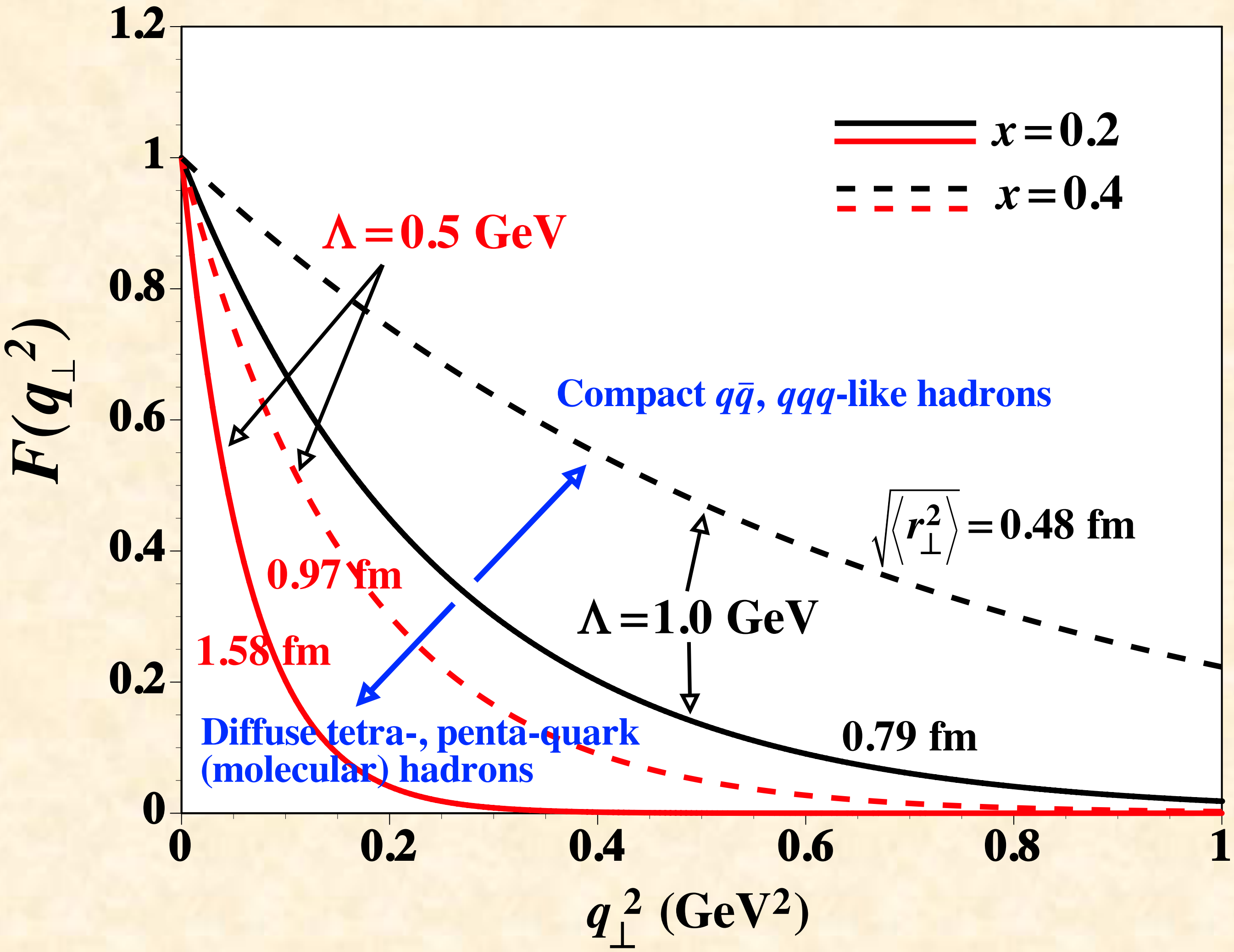


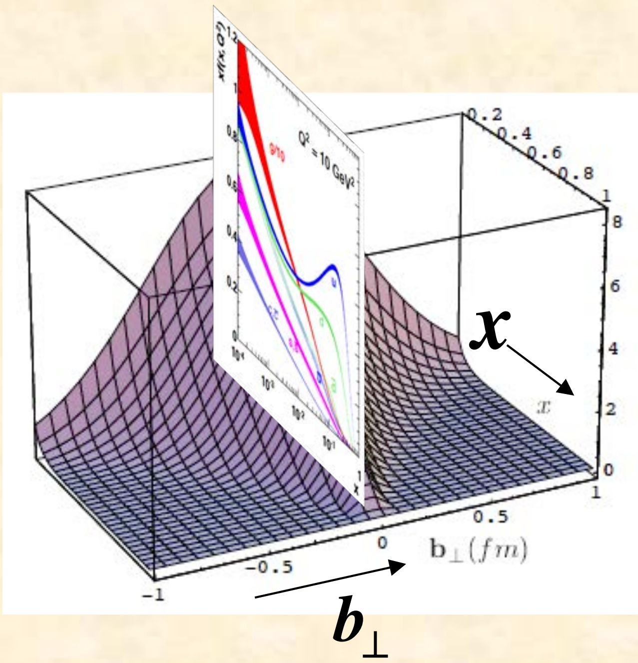
π: M. Aicher, A. Schafer, W. Vogelsang, PRL 105 (2010) 252003.

p: A. D. Martin, R. G. Roberts, W. J. Stirling, PLB 636, 259 (2006)

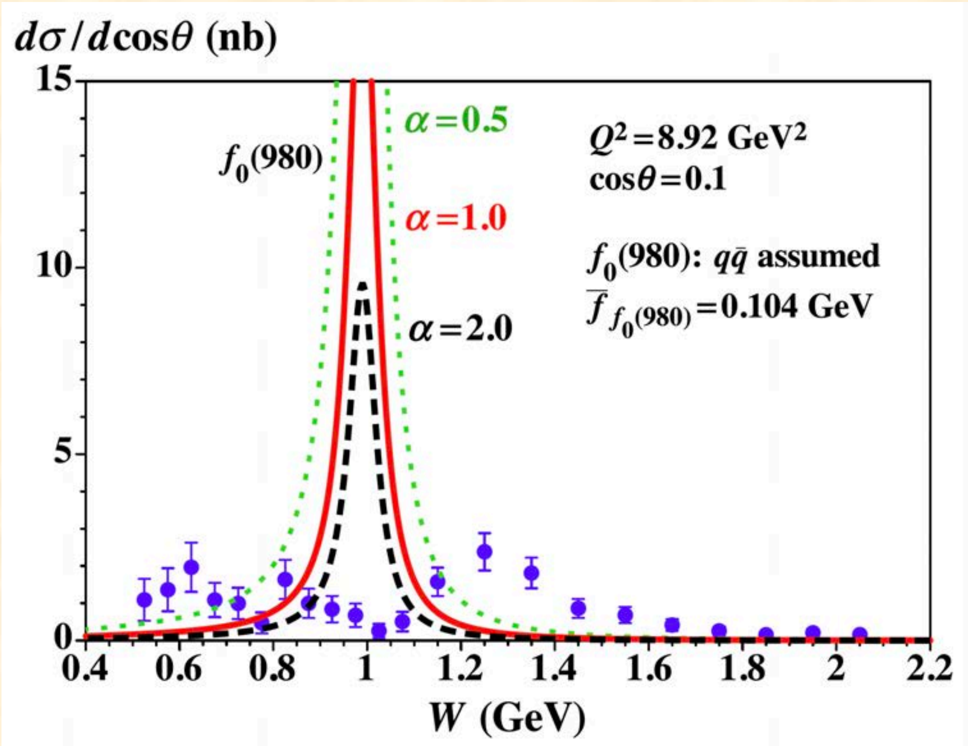
Two-dimensional form factor

$$H_q^h(x,t) = f(x)F(t,x), \quad F(t,x) = e^{(1-x)t/(x\Lambda^2)}, \quad \langle r_\perp^2 \rangle = \frac{4(1-x)}{x\Lambda^2}$$





$f_0(980)$ contribution to $\gamma^*\gamma \to \pi^0\pi^0$



- $f_0(980)$ decay constant is calculated so far by assuming $q\bar{q}$ configuration.

$\to$ not consistent with data

$\to$ $f_0(980)$ is not a $q\bar{q}$ state but likely to be tetra quark or $K\bar{K}$ molecule.

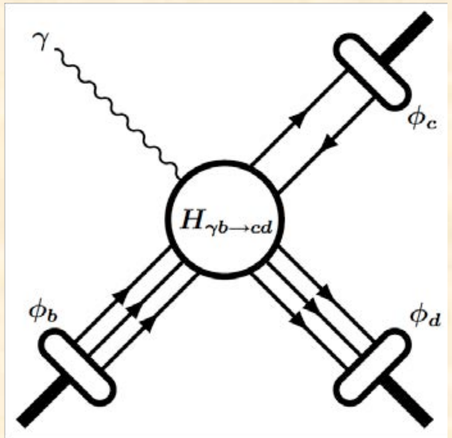
$\to$ $f_0(980)$ is not included in our analysis.

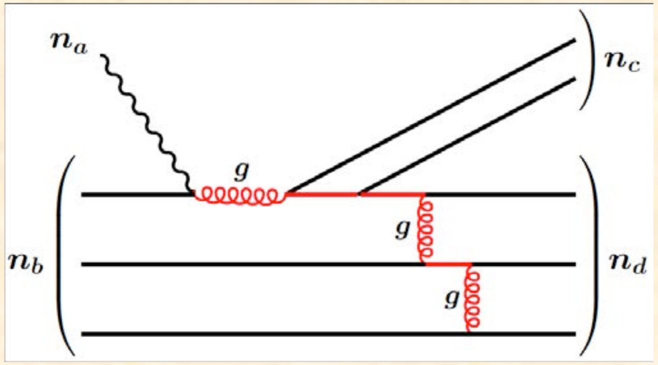
- Resonances: There exist resonance contributions to the cross section.

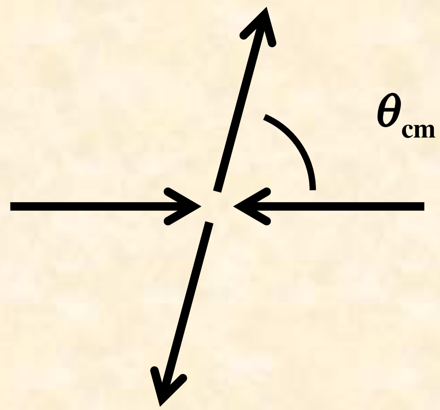
$$\sum_q \Phi_q^{\pi\pi}(z,\zeta,W^2) = 18N_f z^\alpha(1-z)^\alpha(2z-1)\left[\tilde{B}_{10}(W)+\tilde{B}_{12}(W)P_2(\cos\theta)\right]$$

$\gamma^*\gamma\to\pi^0\pi^0$ analysis:
S. Kumano, Q.-T. Song, O. Teryaev,
Phys. Rev. D 97 (2018) 014020.

JLab hyperon productions







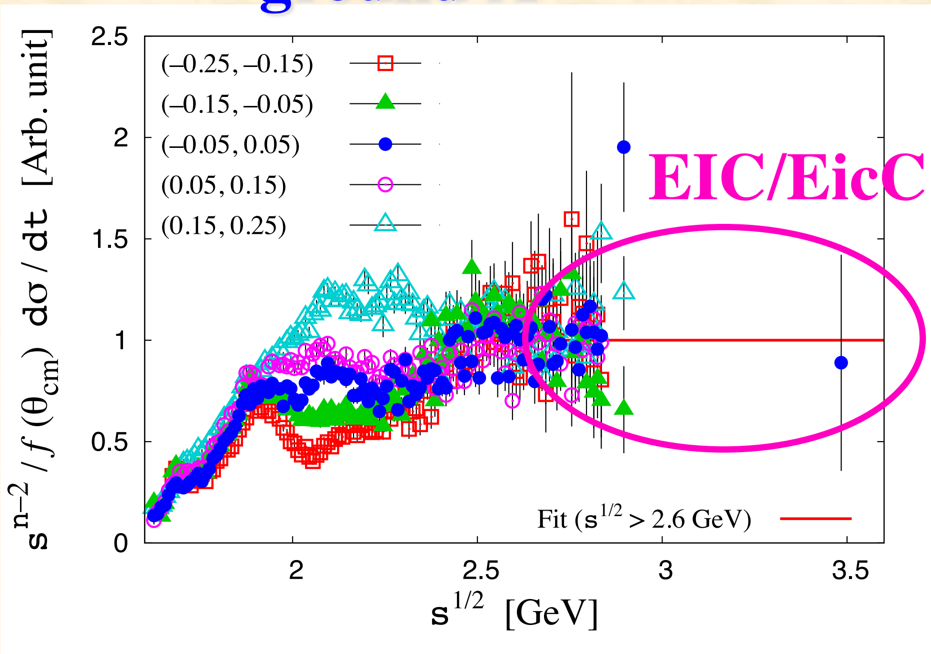
5 bins: $-0.25 < \cos\theta_{cm} < -0.15, \ \cdots, \ 0.15 < \cos\theta_{cm} < 0.25$

4 bins: $-0.20 < \cos\theta_{cm} < -0.10, \ \cdots, \ 0.10 < \cos\theta_{cm} < 0.20$

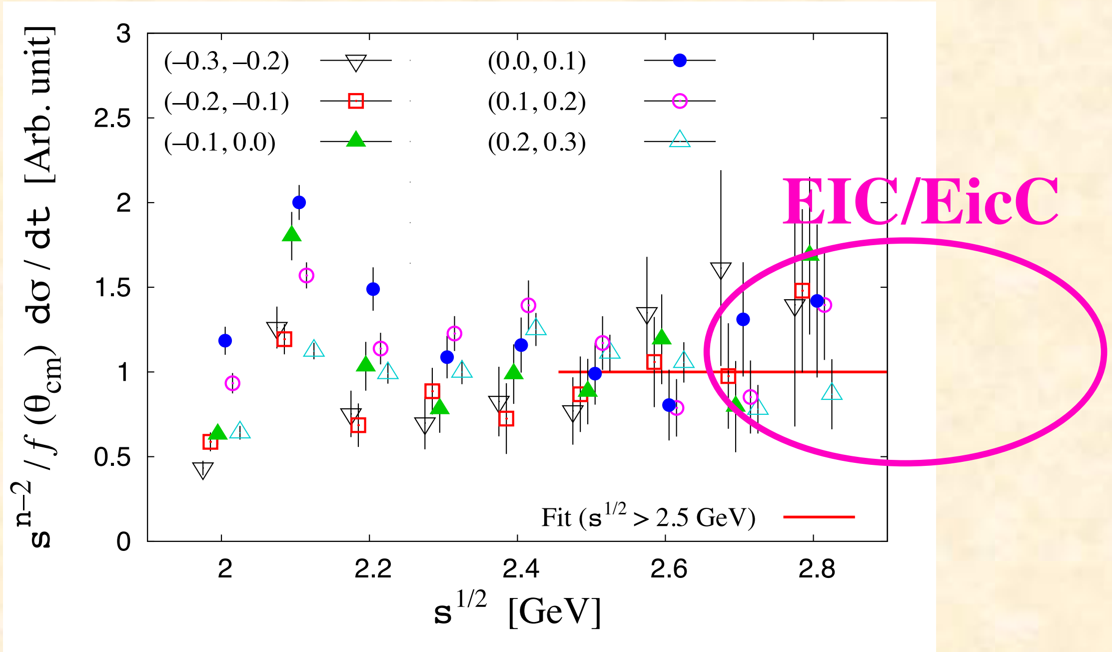
...

1 bin: $-0.05 < \cos\theta_{cm} < +0.05$

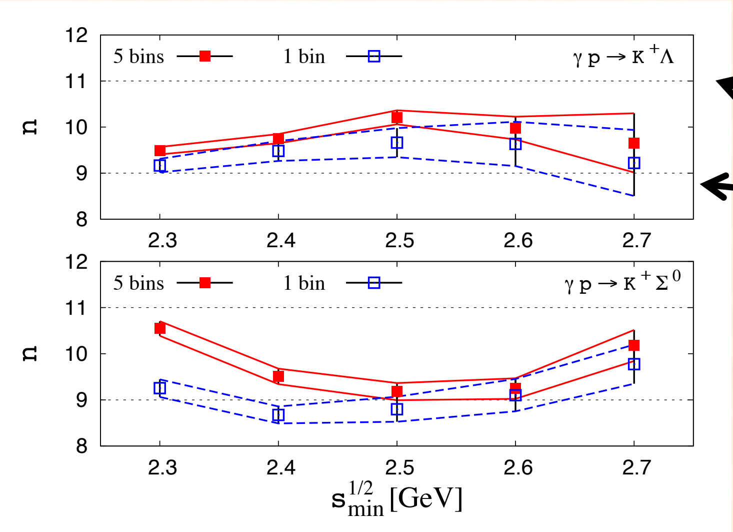
ground Λ



Λ(1405)

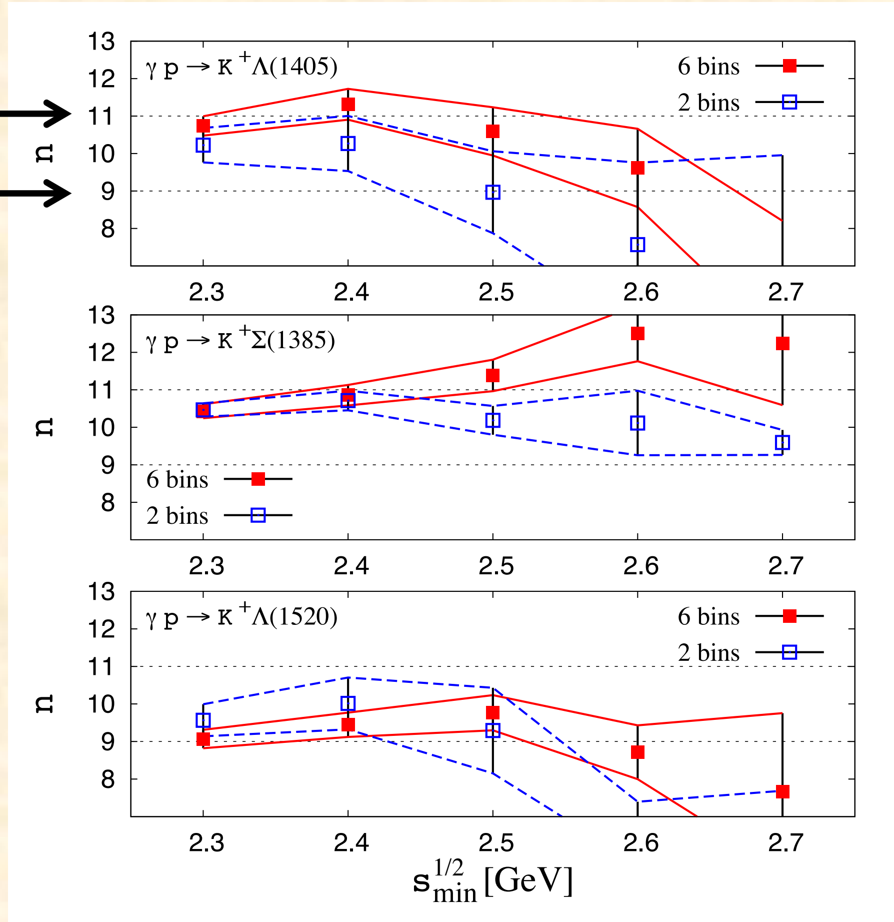


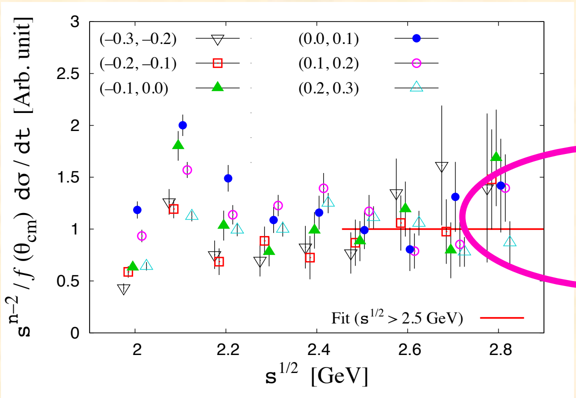
JLab hyperon productions including Λ(1405)



$n_\Lambda = 5$

$n_\Lambda = 3$





Range of
12 GeV JLab!

- Λ. Λ(1520) and Σ seem to be consistent with ordinary baryons with $n = 3$.
- Λ(1405) looks penta-quark at low energies but $n \sim 3$ at high energies???
- Σ(1385) : $n = 5$???

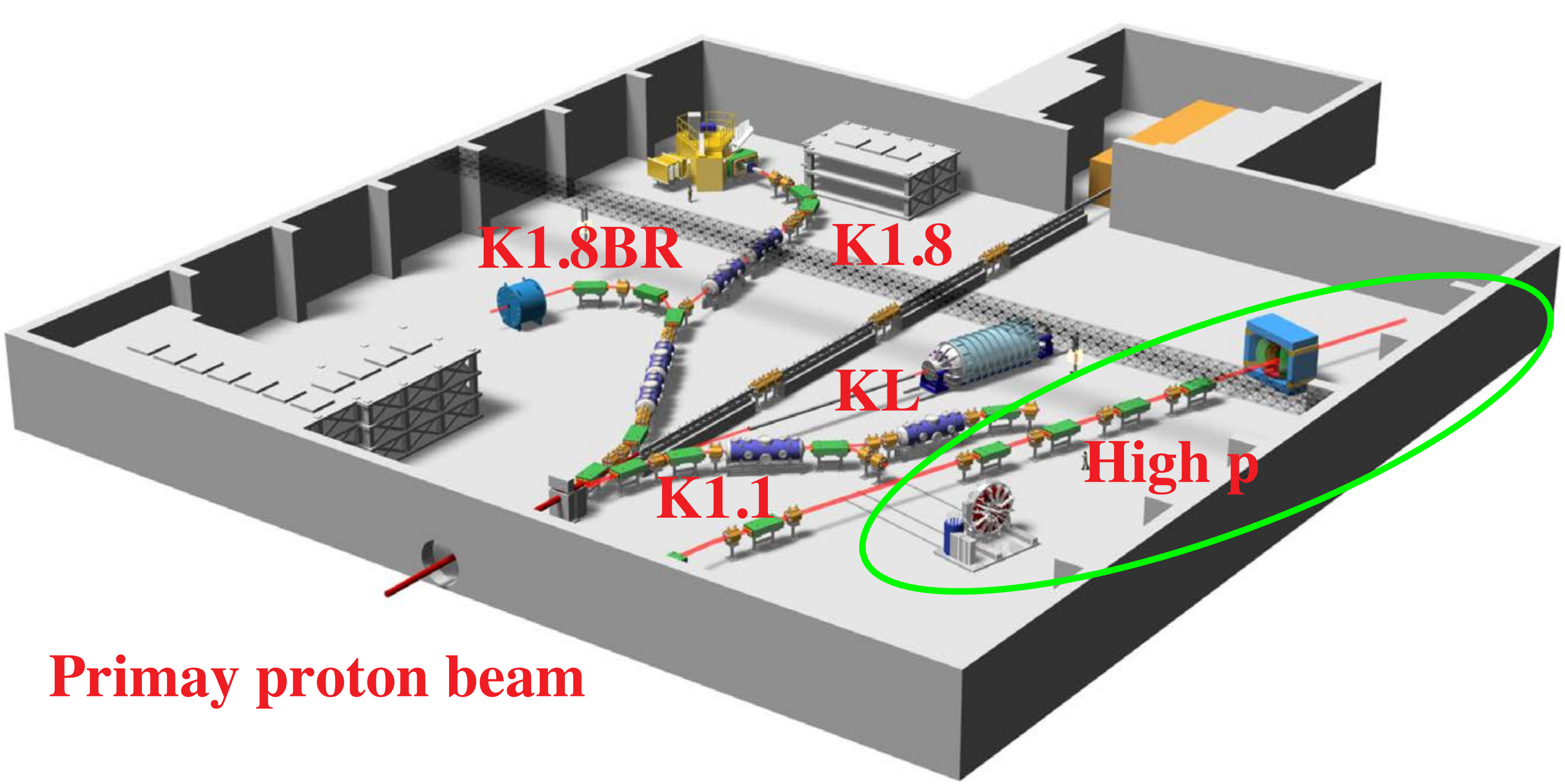
→ In order to clarify the nature of Λ(1405) $\left[qqq,\ \bar{K}N,\ qqqq\bar{q} \right]$,
the JLab 12-GeV experiment plays an important role!

Possible studies on GPDs at hadron accelerator facilities

S. Kumano, M. Strikman, K. Sudoh,
PRD 80 (2009) 074003;
T. Sawada, W.-C. Chang, S. Kumano, J.-C. Peng,
S. Sawada, and K. Tanaka, PRD 93 (2016) 114034.
J-PARC LoI 2019-07, J.-K. Ahn *et al*. (2019).

Hadron facility

Workshops on high-momentum beamline physics,
http://www-conf.kek.jp/hadron1/j-parc-hm-2013/
http://research.kek.jp/group/hadron10/j-parc-hm-2015/.

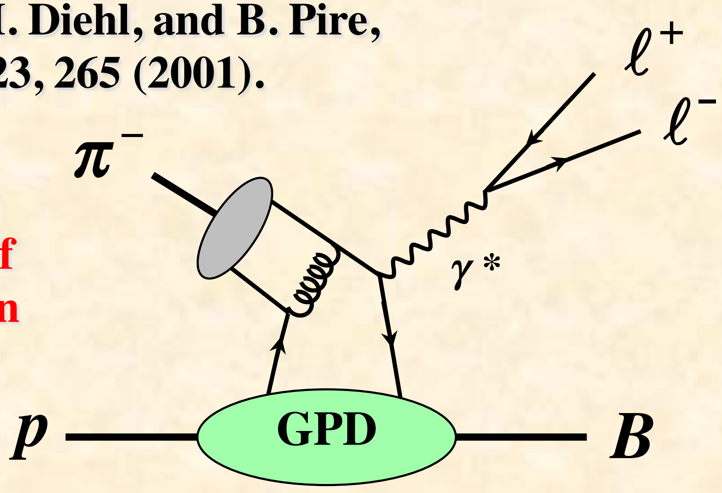


- Proton beam up to 30 GeV
- Unseparated hadron (pion, …) beam up to 15~20 GeV

Toward J-PARC experiments

E. R. Berger, M. Diehl, and B. Pire,
Phys. Lett. B 523, 265 (2001).

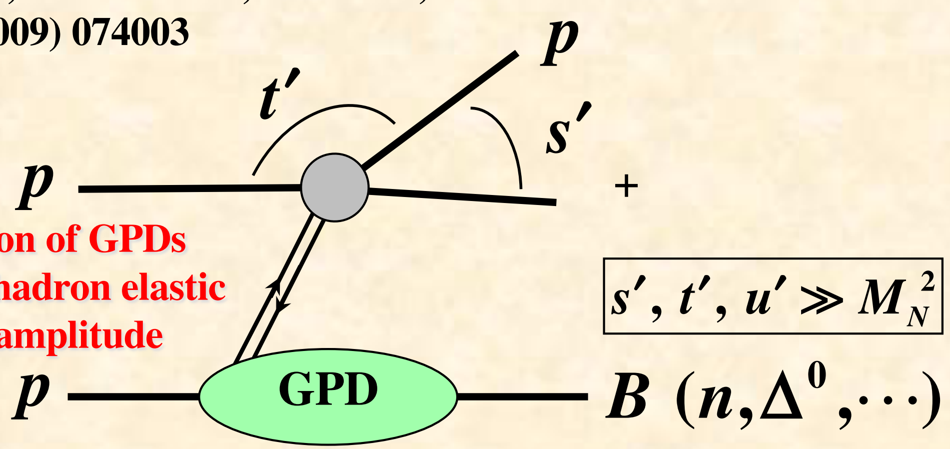
Investigation of GPDs with pion distribution amplitude



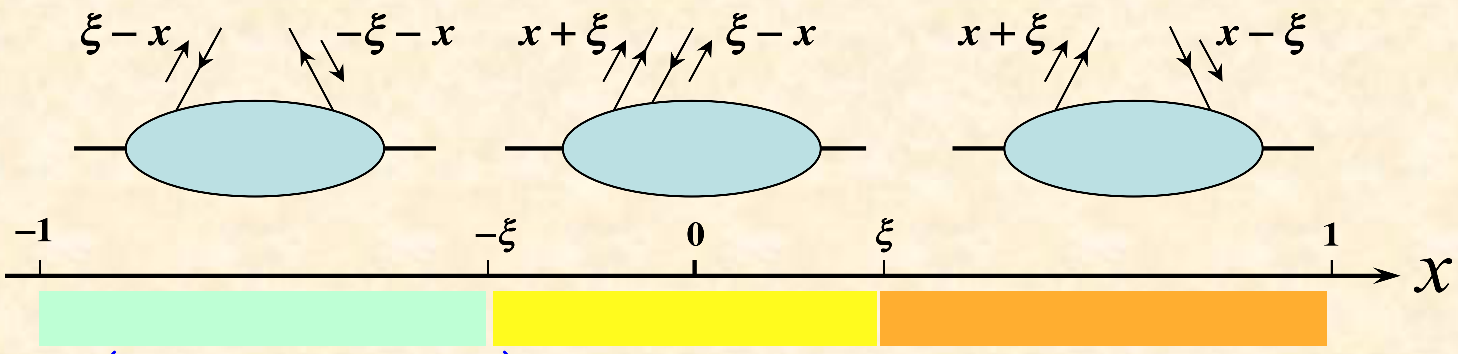
$$\pi^-(\bar{u}d) + p(uud) \rightarrow B(udd) + \gamma^*(\rightarrow \ell^+\ell^-)$$

S. Kumano, M. Strikman, K. Sudoh,
PRD 80 (2009) 074003

Investigation of GPDs with 2→2 hadron elastic scattering amplitude



GPDs in different x regions and GPDs at hadron facilities



$-1 < x < \xi \quad (x + \xi < 0,\ x - \xi < 0)$

$\xi < x < 1 \quad (x + \xi > 0,\ x - \xi > 0)$

$-\xi < x < \xi \quad (x + \xi > 0,\ x - \xi < 0)$

Quark distribution

Emission of quark with momentum fraction $x+\xi$

Absorption of quark with momentum fraction $x-\xi$

q$\bar{\text{q}}$(meson)-like distribution amplitude

Emission of quark with momentum fraction $x+\xi$

Emission of antiquark with momentum fraction $\xi-x$

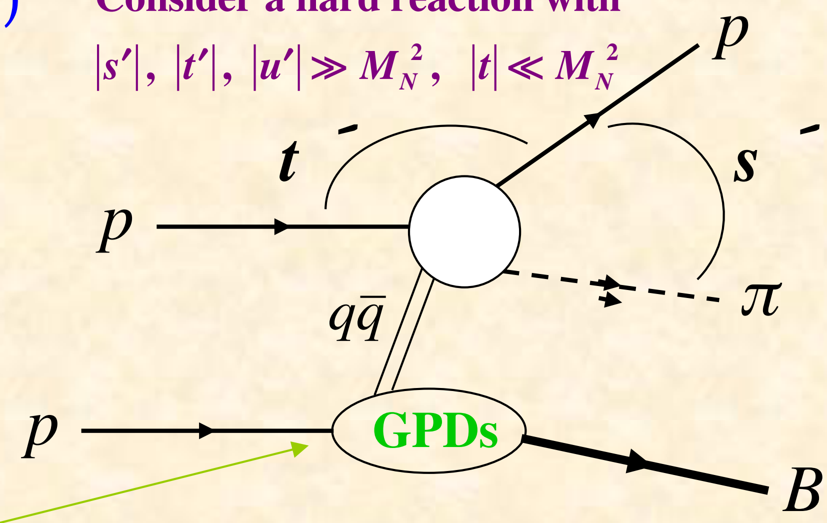
Antiquark distribution

Emission of antiquark with momentum fraction $\xi-x$

Absorption of antiquark with momentum fraction $-\xi-x$

Consider a hard reaction with

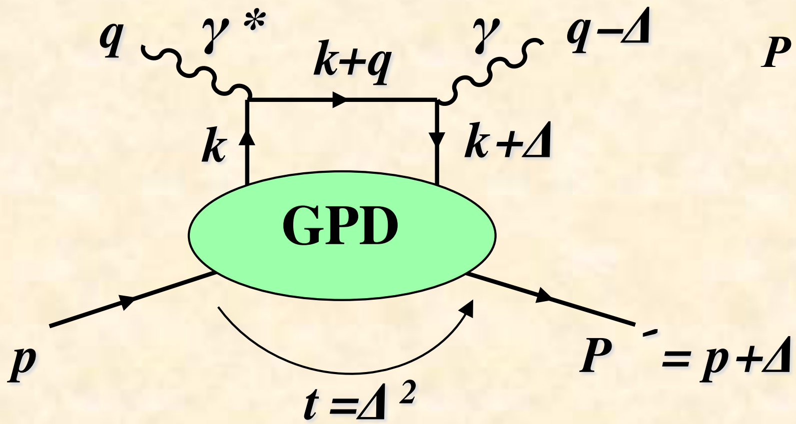
$|s'|,\ |t'|,\ |u'| \gg M_N^2,\ \ |t| \ll M_N^2$



GPDs at J-PARC: S. Kumano, M. Strikman, and K. Sudoh, PRD 80 (2009) 074003.

Efremov-Radyushkin -Brodsky-Lepage (ERBL) region

Generalized Parton Distributions (GPDs)



$$P = \frac{p+p'}{2}, \quad \Delta = p' - p$$

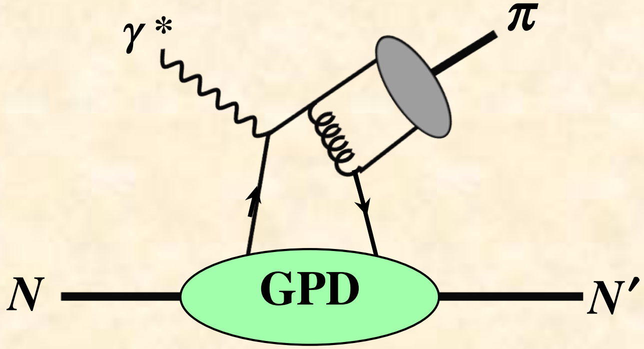
Bjorken variable: $x = \dfrac{Q^2}{2p\cdot q}$

Momentum transfer squared: $t = \Delta^2$

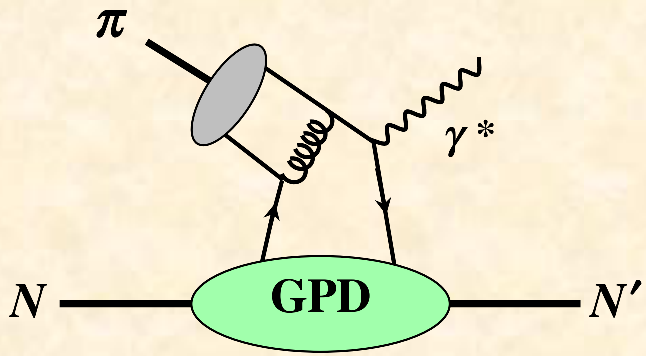
Skewdness parameter: $\xi = \dfrac{p^+ - p'^+}{p^+ + p'^+} = -\dfrac{\Delta^+}{2P^+}$

$$\int \frac{dz^-}{4\pi} e^{ixP^+z^-} \langle p' | \bar\psi(-z/2)\gamma^+\psi(z/2) | p \rangle \Big|_{z^+=0,\vec z_\perp=0} = \frac{1}{2P^+}\left[H(x,\xi,t)\bar u(p')\gamma^+ u(p) + E(x,\xi,t)\bar u(p')\frac{i\sigma^{+\alpha}\Delta_\alpha}{2M}u(p)\right]$$

JLab



J-PARC



$$\int \frac{dz^-}{4\pi} e^{ixP^+z^-} \langle p' | \bar\psi(-z/2)\gamma^+\gamma_5\psi(z/2) | p \rangle \Big|_{z^+=0,\vec z_\perp=0} = \frac{1}{2P^+}\left[\tilde H(x,\xi,t)\bar u(p')\gamma^+\gamma_5 u(p) + \tilde E(x,\xi,t)\bar u(p')\frac{\gamma_5\Delta^+}{2M}u(p)\right]$$

Exclusive Drell-Yan $\pi^- + p \to \mu^+\mu^- + n$ and GPDs

$$\frac{d\sigma_L}{dQ'^2 dt} = \frac{4\pi\alpha^2}{27}\frac{\tau^2}{Q'^2} f_\pi^2 \left[(1-\xi^2)\left|\tilde{H}^{du}(-\xi,\xi,t)\right|^2 - 2\xi^2 \,\mathrm{Re}\left\{\tilde{H}^{du}(-\xi,\xi,t)^* \tilde{E}^{du}(-\xi,\xi,t)\right\} - \xi^2 \frac{t}{4m_N^2}\left|\tilde{E}^{du}(-\xi,\xi,t)\right|^2 \right]$$

$$Q'^2 = q'^2, \;\; t = (p-p')^2, \;\; \tau = \frac{Q'^2}{2p\cdot q_\pi} \simeq \frac{Q'^2}{s-m_\pi^2}$$

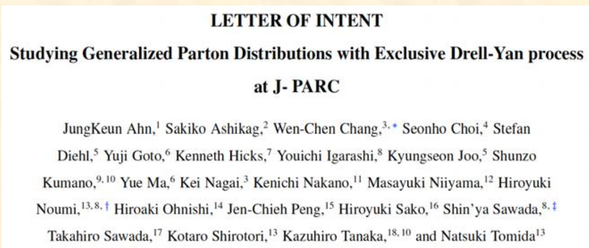
$$\int \frac{dz^-}{4\pi} e^{ixP^+z^-} \langle p(p') | \bar{q}(-z/2)\gamma^+\gamma_5 q(z/2) | p(p) \rangle \Big|_{z^+=0,\vec{z}_\perp=0} = \frac{1}{2P^+}\left[\tilde{H}_p^q(x,\xi,t)\bar{u}(p')\gamma^+\gamma_5 u(p) + \tilde{E}_p^q(x,\xi,t)\bar{u}(p')\frac{\gamma_5\Delta^+}{2M}u(p) \right]$$

$$\int \frac{dz^-}{4\pi} e^{ixP^+z^-} \langle n(p') | \bar{q}_d(-z/2)\gamma^+\gamma_5 q_u(z/2) | p(p) \rangle \Big|_{z^+=0,\vec{z}_\perp=0} = \frac{1}{2P^+}\left[\tilde{H}_{p\to n}^{du}(x,\xi,t)\bar{u}(p')\gamma^+\gamma_5 u(p) + \tilde{E}_{p\to n}^{du}(x,\xi,t)\bar{u}(p')\frac{\gamma_5\Delta^+}{2M}u(p) \right]$$

$$\tilde{H}^{du}(x,\xi,t) = \frac{8}{3}\alpha_s \int_{-1}^{1} dz \frac{\phi_\pi(z)}{1-z^2}\int_{-1}^{1} dx' \left[\frac{e_d}{x-x'-i\varepsilon} - \frac{e_u}{x+x'-i\varepsilon} \right]\left[\tilde{H}^d(x',\xi,t) - \tilde{H}^u(x',\xi,t) \right]$$

$$\tilde{E}^{du}(x,\xi,t) = \frac{8}{3}\alpha_s \int_{-1}^{1} dz \frac{\phi_\pi(z)}{1-z^2}\int_{-1}^{1} dx' \left[\frac{e_d}{x-x'-i\varepsilon} - \frac{e_u}{x+x'-i\varepsilon} \right]\left[\tilde{E}^d(x',\xi,t) - \tilde{E}^u(x',\xi,t) \right]$$

T. Sawada, W.-C. Chang, S. Kumano, J.-C. Peng, S. Sawada, and K. Tanaka, PRD93 (2016) 114034.

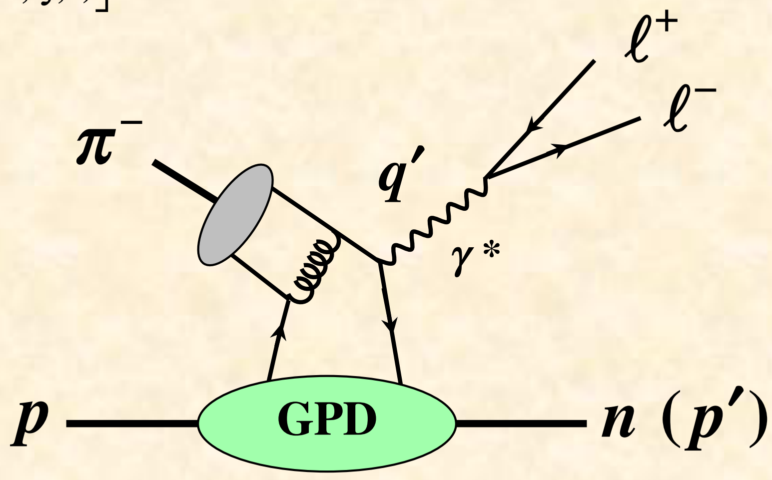


LETTER OF INTENT

Studying Generalized Parton Distributions with Exclusive Drell-Yan process at J- PARC

JungKeun Ahn,[1] Sakiko Ashikag,[2] Wen-Chen Chang,[3,*] Seonho Choi,[4] Stefan Diehl,[5] Yuji Goto,[6] Kenneth Hicks,[7] Youichi Igarashi,[8] Kyungseon Joo,[5] Shunzo Kumano,[9,10] Yue Ma,[6] Kei Nagai,[3] Kenichi Nakano,[11] Masayuki Niiyama,[12] Hiroyuki Noumi,[13,8,†] Hiroaki Ohnishi,[14] Jen-Chieh Peng,[15] Hiroyuki Sako,[16] Shin'ya Sawada,[8,‡] Takahiro Sawada,[17] Kotaro Shirotori,[13] Kazuhiro Tanaka,[18,10] and Natsuki Tomida[13]

LoI for a J-PARC experiment

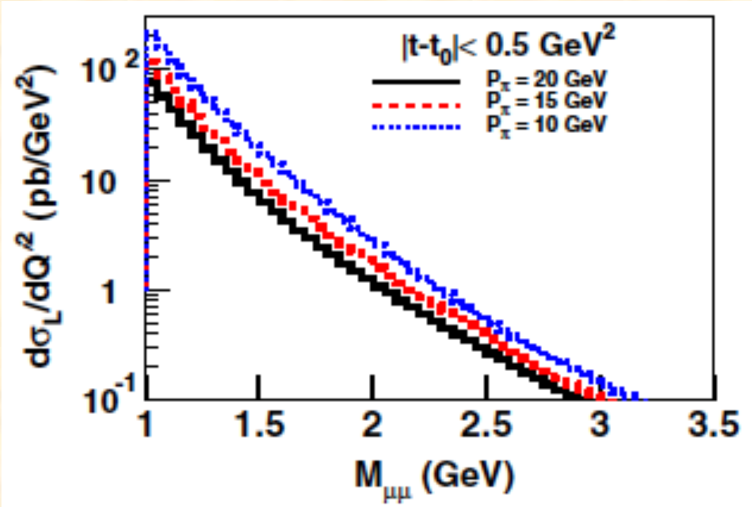


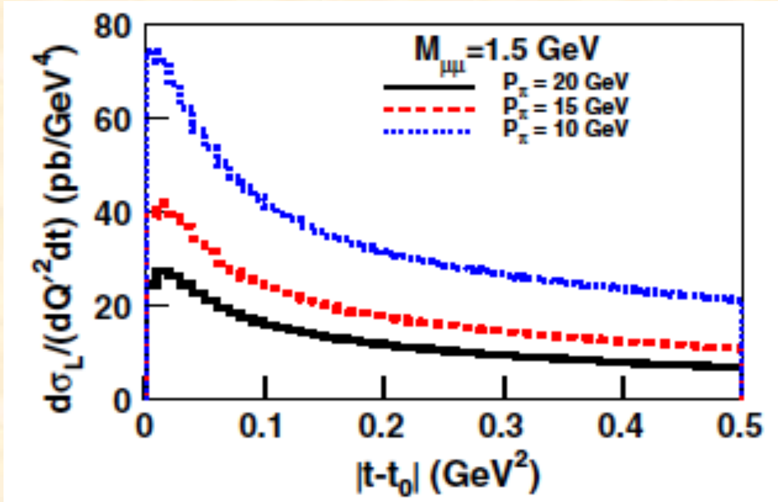
$$\pi^-(\bar{u}d) + p(uud) \to n(udd) + \gamma^*(\to \ell^+\ell^-)$$

Expected Drell-Yan events at J-PARC

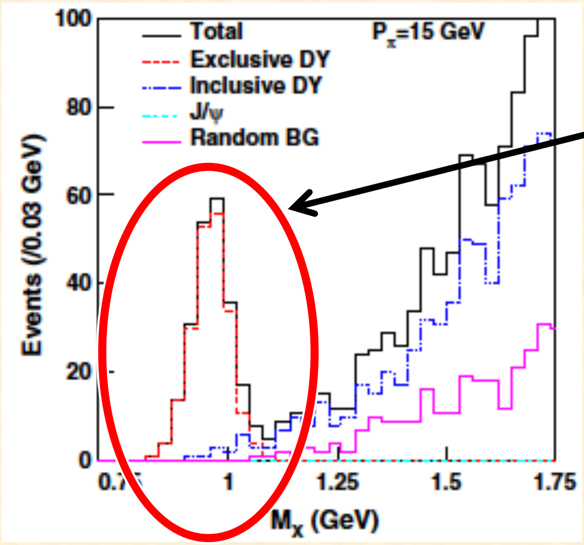
$$Q'^2 = q'^2,\ t = (p - p')^2,\ \tau = \frac{Q'^2}{2p \cdot q_\pi} \simeq \frac{Q'^2}{s - m_N^2}$$

$$\frac{d\sigma_L}{dQ'^2 dt} = \frac{4\pi\alpha^2}{27}\frac{\tau^2}{Q'^2} f_\pi^2 \left[(1-\xi^2)\left|\tilde{H}^{du}(-\xi,\xi,t)\right|^2 - 2\xi^2 \,\mathrm{Re}\left\{\tilde{H}^{du}(-\xi,\xi,t)^* \tilde{E}^{du}(-\xi,\xi,t)\right\} - \xi^2 \frac{t}{4m_N^2}\left|\tilde{E}^{du}(-\xi,\xi,t)\right|^2 \right]$$

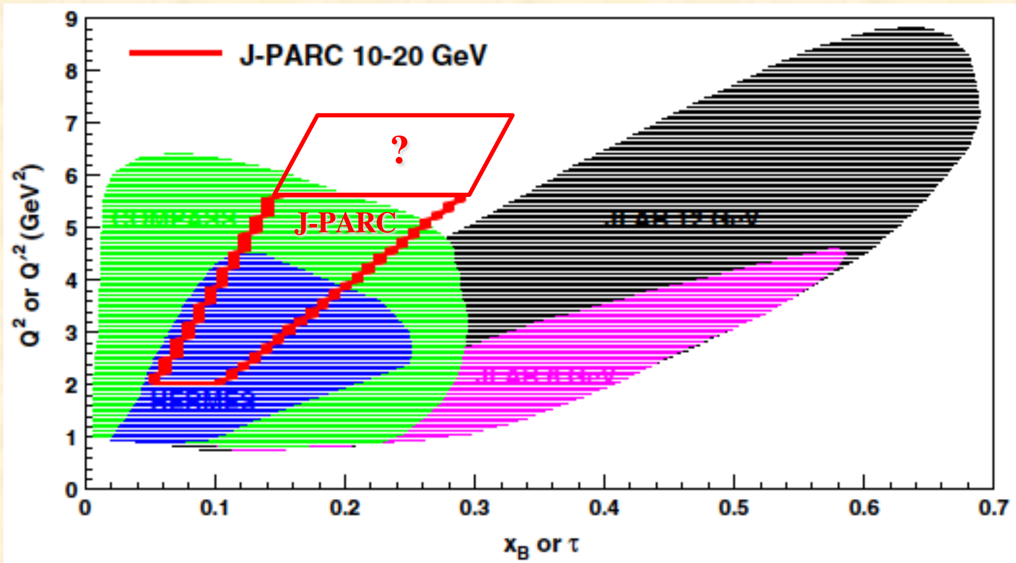




Missing mass



Exclusive Drell-Yan

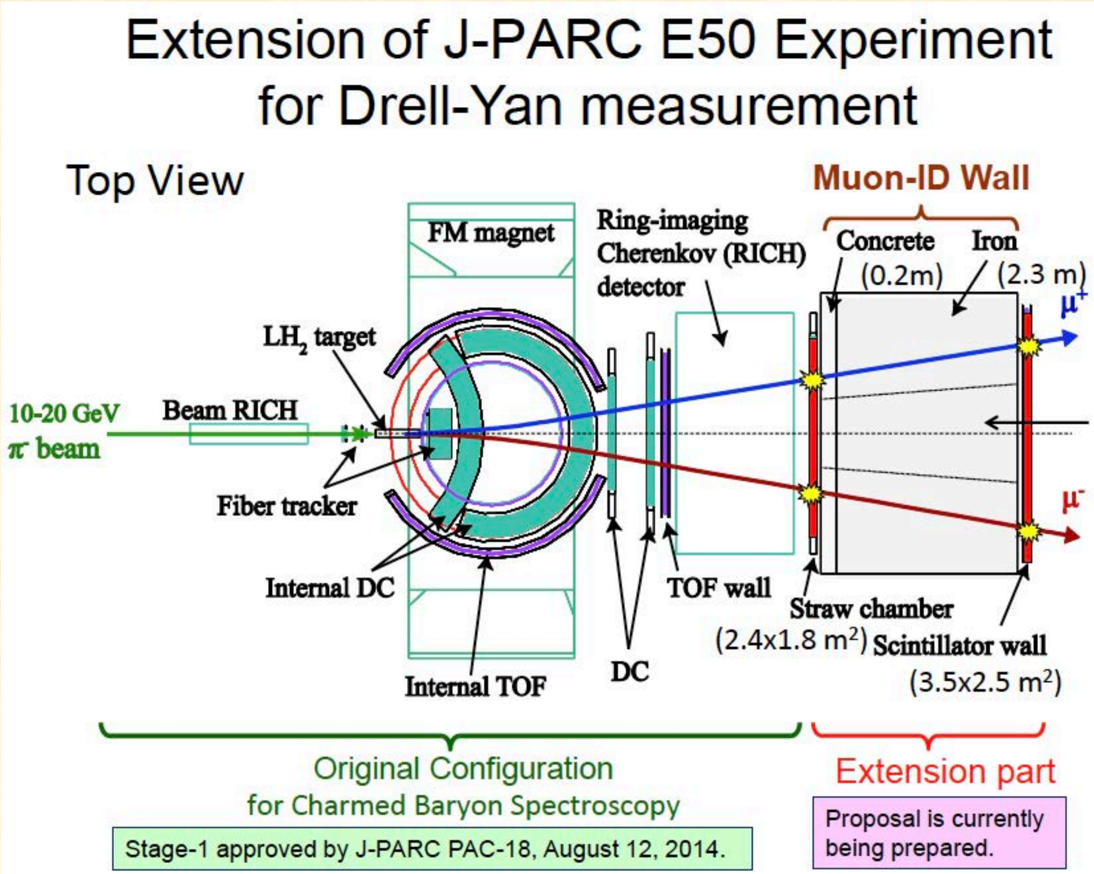


Letter of Intent to join J-PARC-E50 collaboration (Jan. 2019)

LETTER OF INTENT

Studying Generalized Parton Distributions with Exclusive Drell-Yan process at J- PARC

JungKeun Ahn,[1] Sakiko Ashikag,[2] Wen-Chen Chang,[3,*] Seonho Choi,[4] Stefan Diehl,[5] Yuji Goto,[6] Kenneth Hicks,[7] Youichi Igarashi,[8] Kyungseon Joo,[5] Shunzo Kumano,[9,10] Yue Ma,[6] Kei Nagai,[3] Kenichi Nakano,[11] Masayuki Niiyama,[12] Hiroyuki Noumi,[13,8,†] Hiroaki Ohnishi,[14] Jen-Chieh Peng,[15] Hiroyuki Sako,[16] Shin'ya Sawada,[8,‡] Takahiro Sawada,[17] Kotaro Shirotori,[13] Kazuhiro Tanaka,[18,10] and Natsuki Tomida[13]

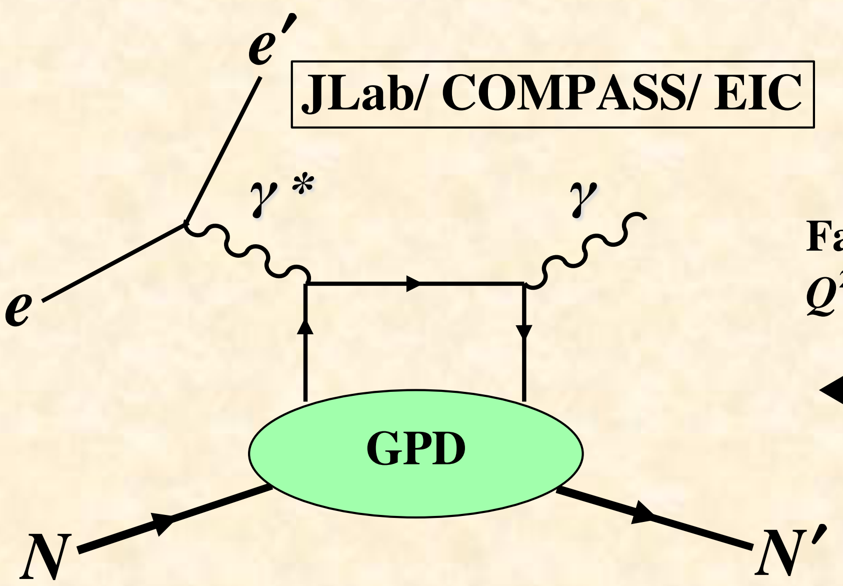


Possible studies on GPDs at neutrino facilities

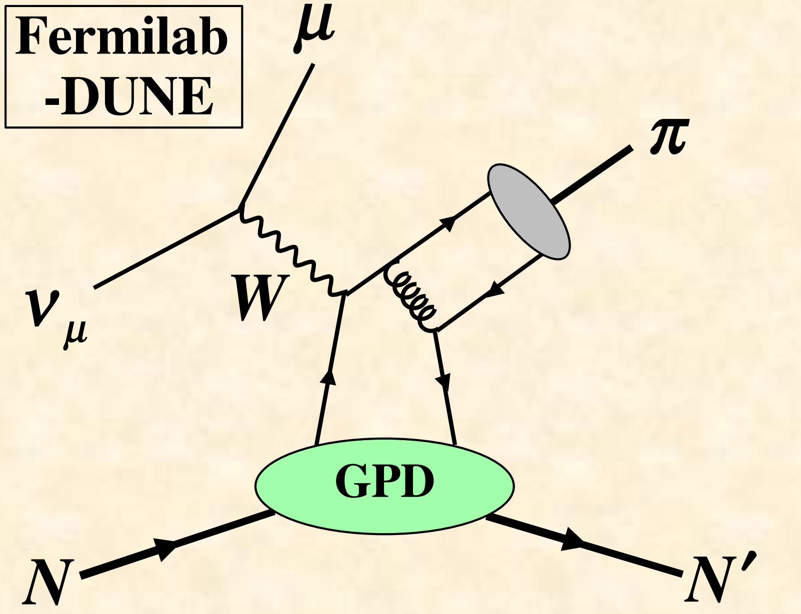
S. Kumano, EPJ Web Conf. 208 (2019) 07003.

EIC yellow report, R. Abdul Khalek *et al*., arXiv:2103.05419
Sec. 7.5.2, Neutrino physics by S. Kumano and R. Petti.

Neutrino reactions for gravitational form factors @Fermilab-DUNE (Origins of hadron masses and pressures)

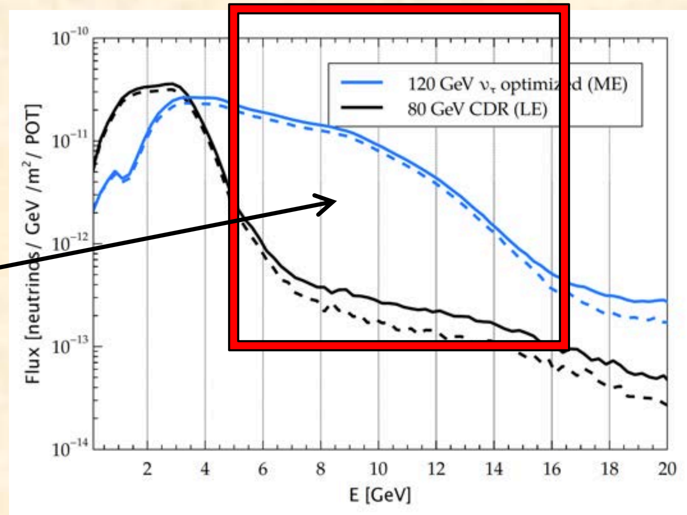


Factorization condition:
$Q^2 \gg |t|, \Lambda^2_{QCD}$



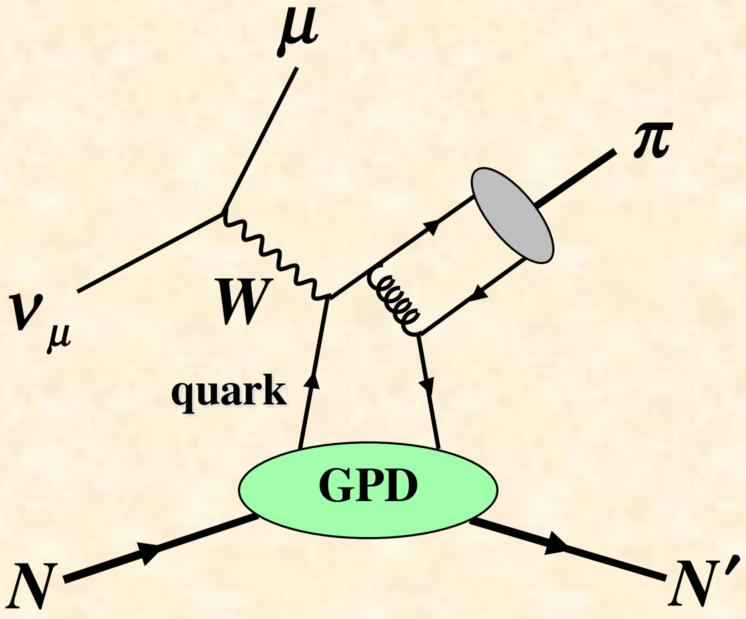
Deep Underground Neutrino Experiment (DUNE)
at Long-Baseline Neutrino Facility (LBNF)

High-energy part of the LBNF ν beam can be used for the GPD studies.



J. Rout *et al*., PRD 102 (2020) 116018

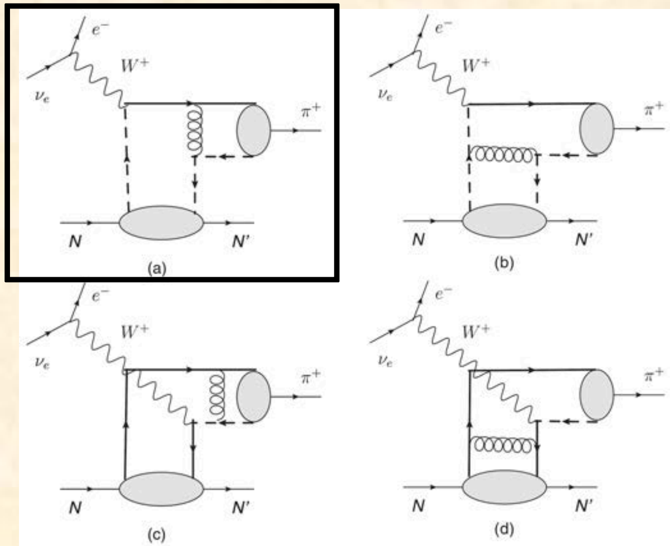
Recent work on pion production in neutrino reaction for GPD studies



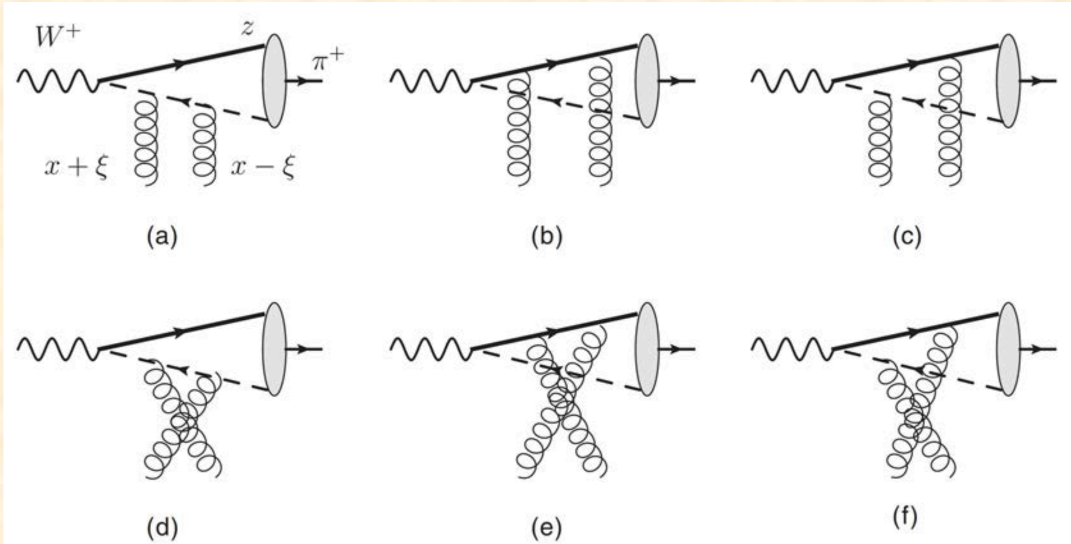
B. Pire, L. Szymanowski, and J. Wagner,
Phys. Rev. D 95, 114029 (2017).

There are several processes to contribute to the pion-production cross section, including the gluon GPD terms.

Quark GPDs



Gluon GPDs



Cross section estimates

proton: $\nu p \to \ell^- \pi^+ p$ **neutron:** $\nu n \to \ell^- \pi^+ n$

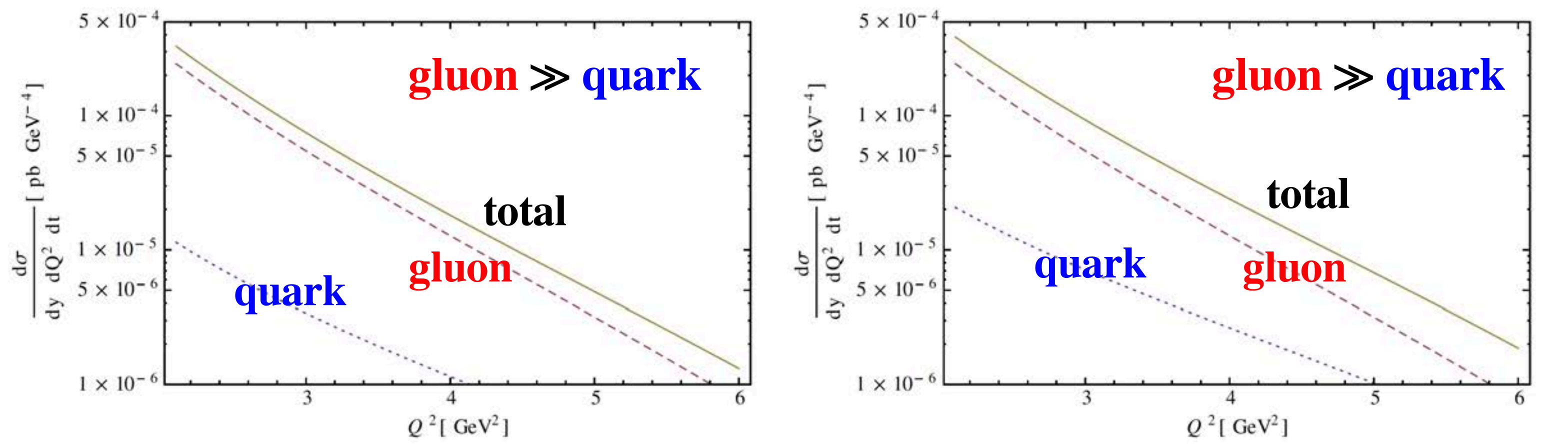


FIG. 3. The Q^2 dependence of the cross section $\frac{d^3\sigma(\nu N \to l^- N \pi^+)}{dy dQ^2 dt}$ (in pb GeV^{-4}) for $y = 0.7$, $\Delta_T = 0$ and $s = 20$ GeV^2, on a proton (left panel) and on a neutron (right panel). The quark contribution (dotted curves) is significantly smaller than the gluon contribution (dashed curves). The solid curves are the sum of the (quark + gluon + interference) contributions.

neutron → proton: $\nu n \to \ell^- \pi^0 p$

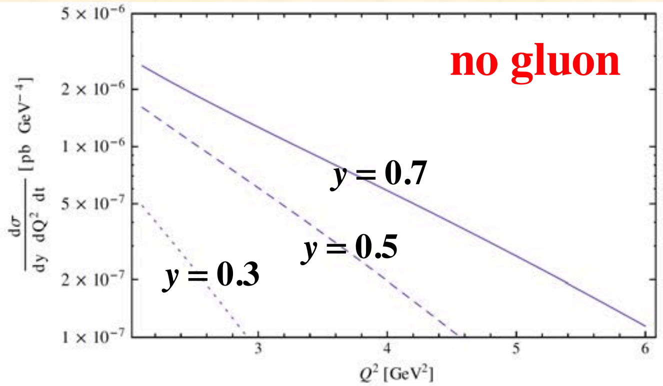
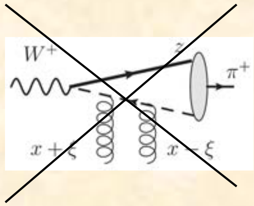


FIG. 6. The Q^2 dependence of the cross section $\frac{d^3\sigma(\nu n \to l^- p \pi^0)}{dy dQ^2 dt}$ (in pb GeV^{-4}) for $\Delta_T = 0$ and $s = 20$ GeV^2. The solid, dashed, and dotted lines correspond to $y = 0.7$, 0.5, and 0.3, respectively. There is no gluon contribution to this amplitude.

Neutrino GPD studies are complementary to the charged-lepton projects.

- **Gluon GPDs could be probed in charged-pion production.**
- **Flavor dependece of quark GPDs could be investigated.**

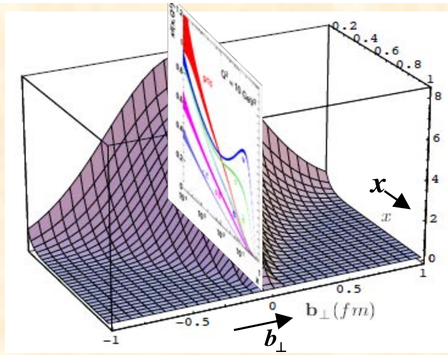


no gluon for π^0

Prospects & Summary

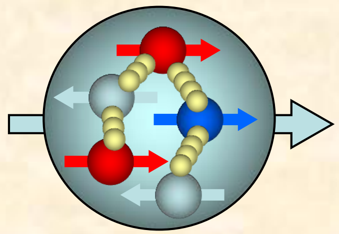
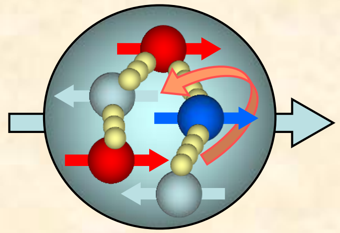
By hadron tomography

3D view of hadrons

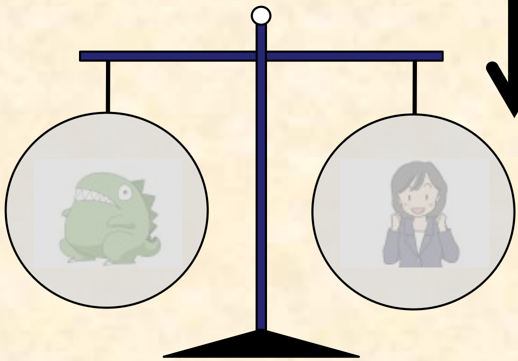


Origin of nucleon spin

By the tomography, we determine

 or  .

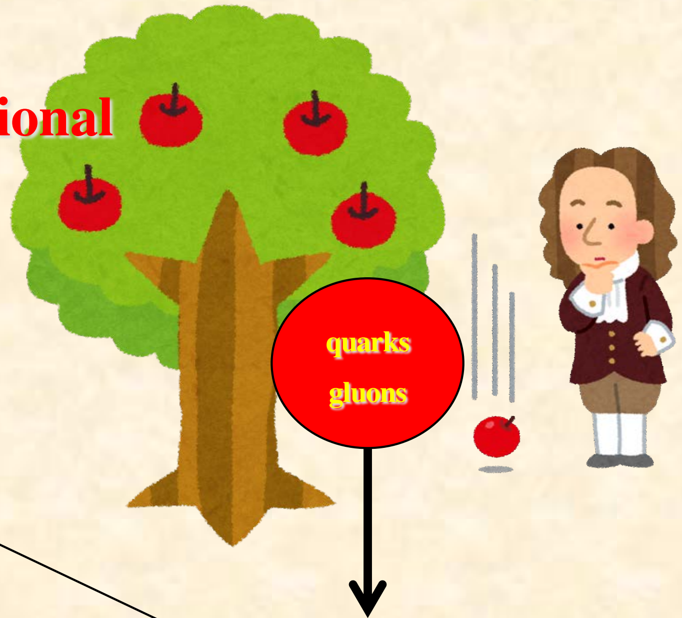
Exotic hadrons



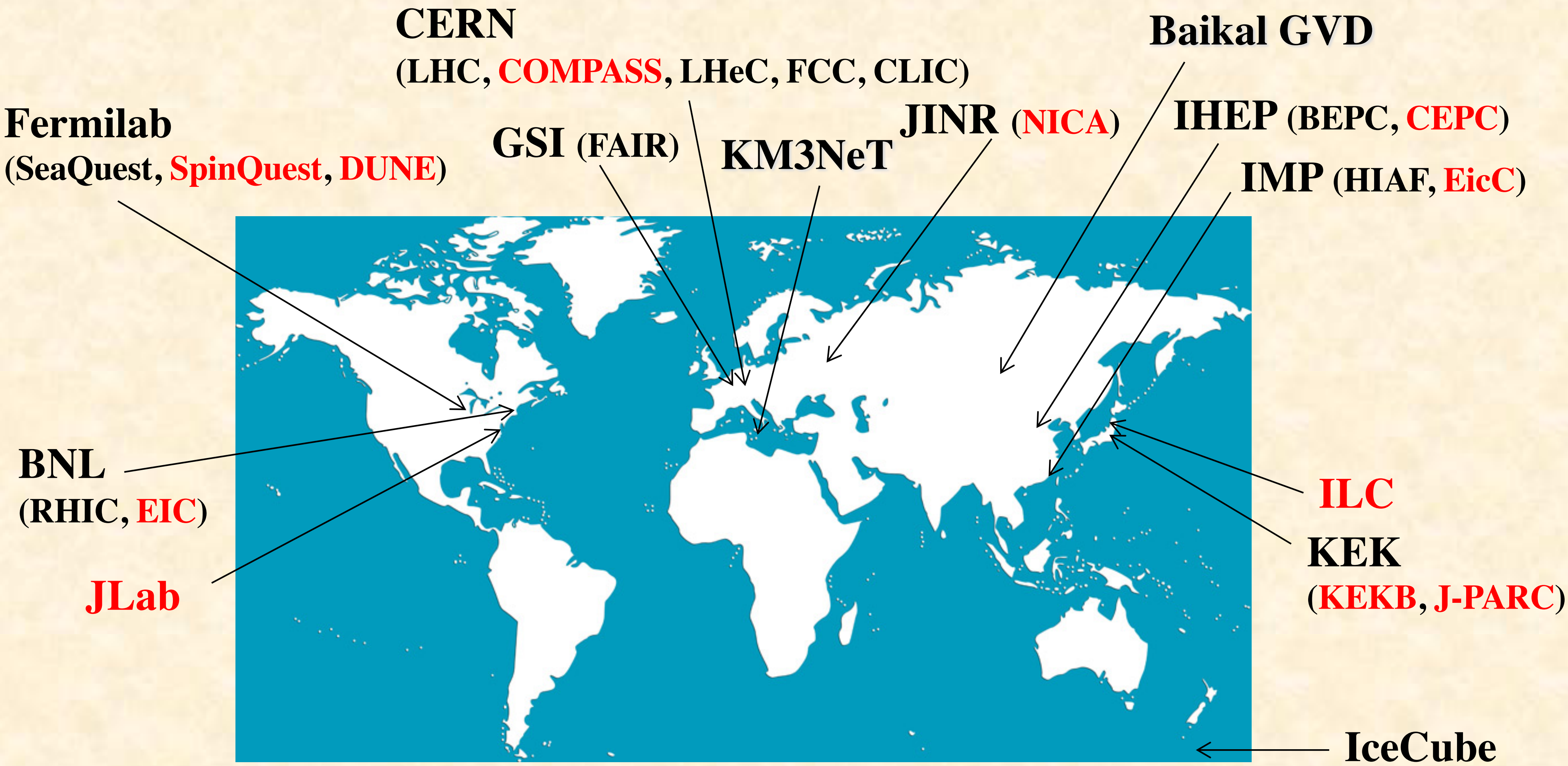
By tomography, we determine  or  .

Origin of gravitational source (mass)

By tomography,
we determine gravitational sources in terms of quarks and gluons.



High-energy hadron physics experiments



Facilities on hadron structure functions on GPDs including future possibilities.

Summary

Hadron-tomography and gravitational form factors

- Puzzle to find the origin of hadron masses and pressures
  in terms of quark and gluon degrees of freedom

- Puzzle to find the origin of nucleon spin

- Exotic hadron candidates could be studied
  in the same tomography method.

- There are world-wide lepton and hadron accelerator facilities
  which has been used and could be used in future for our studies.

Time has come to understand the gravitational sources
and their interactions in microscopic
(instead of usual macroscopic/cosmic)
world in terms of quark and gluon degrees of freedom.

The End

The End