

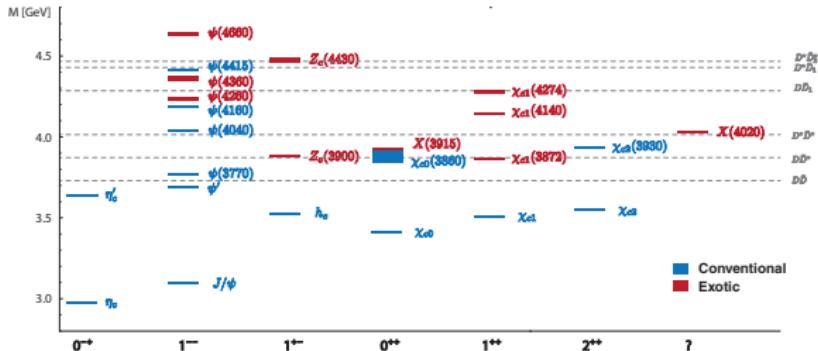


Exotic mesons with functional methods

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LIP & IST Lisboa

Mass in the Standard Model and Consequences of its Emergence
ECT* Trento, Italy
April 22, 2021

Exotic mesons



- Several tetraquark candidates in **charmonium spectrum**: $X(3872)$, $X(3915)$, $Z_c(3900)$,
- Z states cannot be $c\bar{c}$ since they carry charge
- Oldest tetraquark candidates: **light scalar mesons**

Reviews:

- Chen, Chen, Liu, Zhu,
Phys. Rept. 639 (2016), 1601.02092
- Lebed, Mitchell, Swanson
PPNP 93 (2017), 1610.04528
- Esposito, Pilloni, Polosa,
Phys. Rept. 668 (2017), 1611.07920
- Guo, Hanhart, Meißner et al.,
Rev. Mod. Phys. 90 (2018), 1705.00141
- Ali, Lange, Stone,
PPNP 97 (2017), 1706.00610
- Olsen, Skwarnicki, Zieminska,
Rev. Mod. Phys. 90 (2018), 1708.04012
- Liu, Chen, Chen, Liu, Zhu,
PPNP 107 (2019), 1903.11976
- Brambilla, Eidelman, Hanhart et al.,
Phys. Rept. 873 (2020)

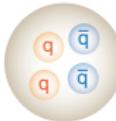
Light exotic mesons



Glueballs?



Hybrid
mesons?



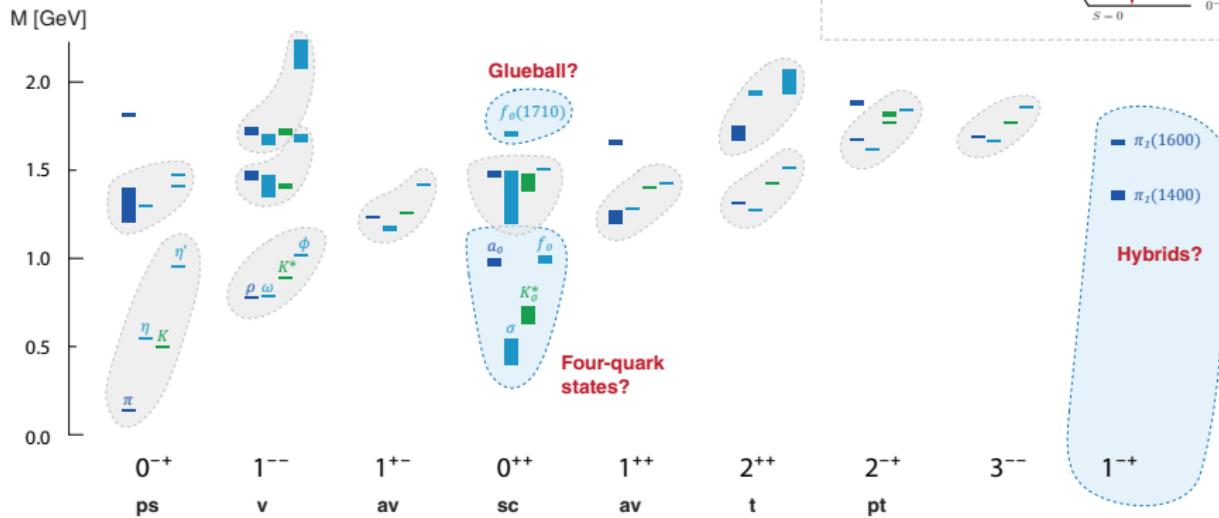
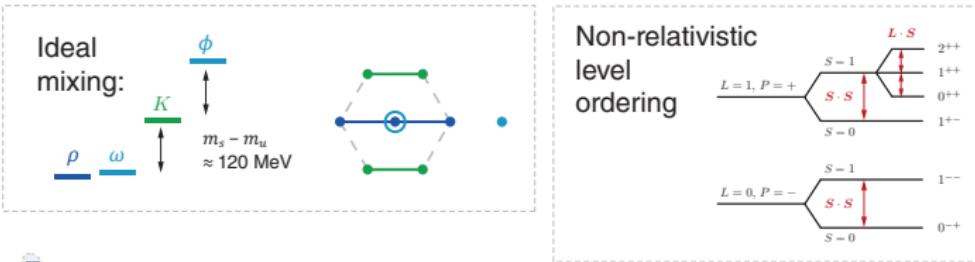
Tetraquarks?

What does “exotic” mean for light hadrons?

- Exotic quantum numbers: $1^{-+}, \dots$
→ **Hybrids?** $\pi_1(1400)?, \pi_1(1600)$
[Meyer, Swanson, Prog. Part. Nucl. Phys. 82 \(2015\), 1502.07276 \[hep-ph\]](#)
[Rodas et al., PRL 122 \(2019\), 1810.04171 \[hep-ph\]](#)
- Overpopulation of multiplets,
unconventional properties
→ **Glueballs?** $f_0(1??0)$
[Ochs, J. Phys. G 40 \(2013\), 1301.5183 \[hep-ph\]](#)
→ **Four-quark states?** $\sigma, \kappa, a_0, f_0, \dots$
[Pelaez, Phys. Rept. 658 \(2016\), 1510.00653 \[hep-ph\]](#)

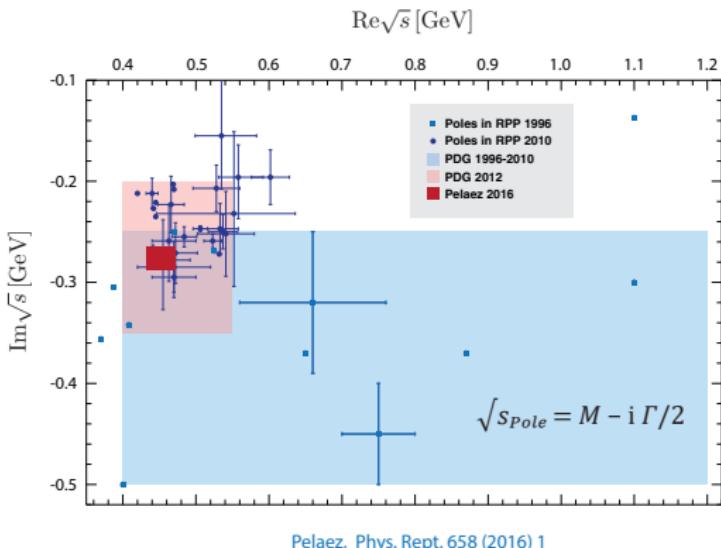
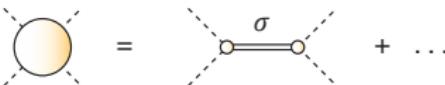
Light exotic mesons

Light meson spectrum
(PDG 2020)



Light scalar mesons

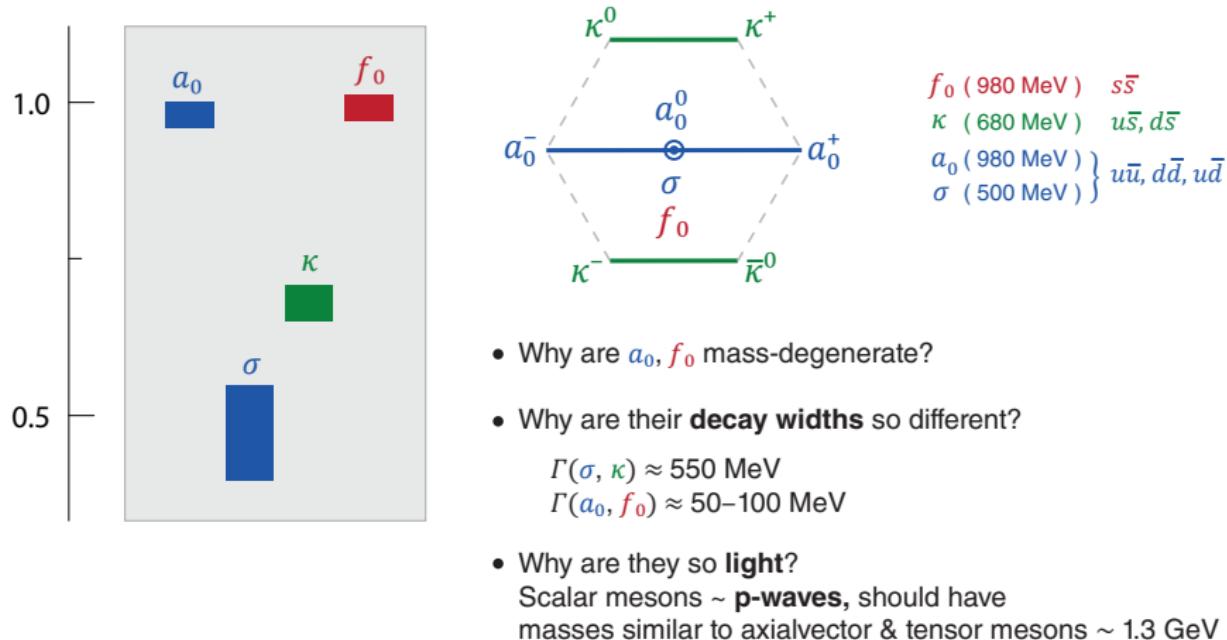
$\sigma/f_0(500)$ is a **resonance** in $\pi\pi$ scattering:



- **PDG 2010:** “ $f_0(600)$ ”
 $\sqrt{s} \sim (400 \dots 1200) - i (250 \dots 500) \text{ MeV}$
- **Dispersive analyses:**
 $\sqrt{s} \sim 450(20) - i 275(10) \text{ MeV}$
 - Caprini, Colangelo, Leutwyler 2006
 - Garcia-Martin, Kaminski, Pelaez, Ruiz de Elvira 2011
 - Moussallam 2011
 - Masjuan, Ruiz de Elvira, Sanz-Cillero 2014
 - Pelaez 2016
- **PDG 2012:** “ $f_0(500)$ ”
 $\sqrt{s} \sim (400 \dots 550) - i (200 \dots 350) \text{ MeV}$
- Pole locations from **lattice QCD**
Briceno, Dudek, Edwards, Wilson, PRL 118 (2017), PRD 97 (2018)

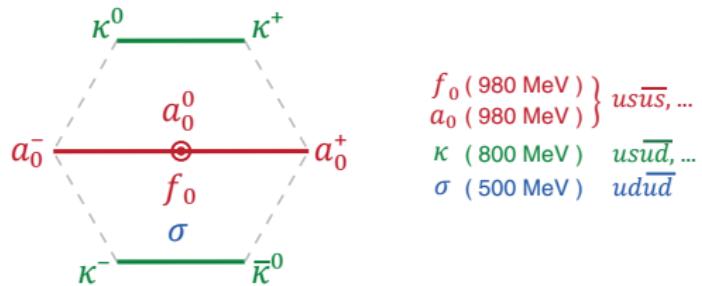
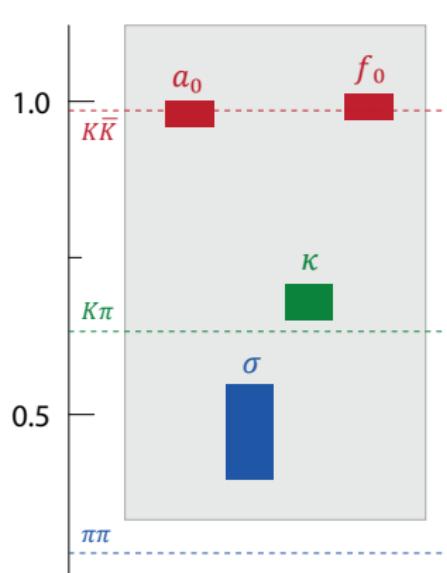
Light scalar mesons

Light scalar (0^{++}) mesons don't fit into the conventional meson spectrum:



Light scalar mesons

What if they were **tetraquarks** (diquark-antidiquark)? Jaffe 1977, Close, Tornqvist 2002, Maiani, Polosa, Riquer 2004



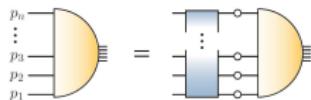
- Explains **mass ordering & decay widths**: f_0 and a_0 couple to $K\bar{K}$, large widths for σ , κ
- Alternative: **meson molecules?**
Weinstein, Isgur 1982, 1990; Close, Isgur, Kumano 1993
- **Non-q-q̄ nature** of σ supported by
dispersive analyses, unitarized ChPT, large N_c ,
extended linear σ model, quark models

Pelaez, Phys. Rept. 658 (2016)



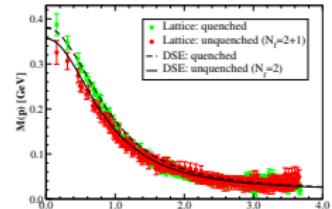
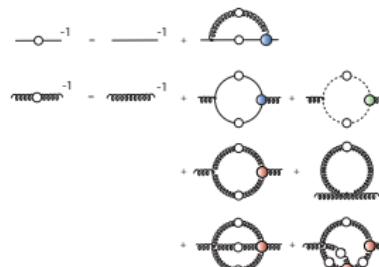
Functional methods

- Hadronic **bound-state equations**
(BSEs, Faddeev eqs, ...)



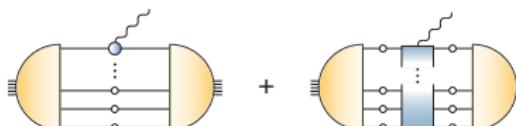
“QFT analogue of Schrödinger eq.”
→ hadron masses & “wave functions”
→ **spectroscopy calculations**

- Ingredients: **QCD's n-point functions**,
Satisfy Dyson-Schwinger equations (**DSEs**):
QCD's quantum eqs. of motion



→ running **quark mass**

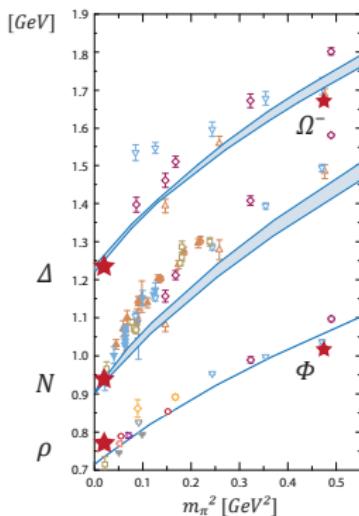
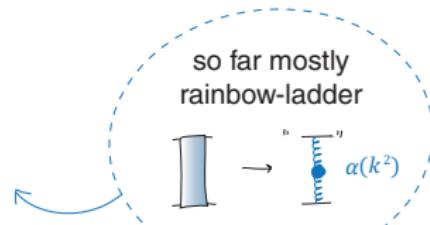
- Structure calculations: form factors, PDFs, GPDs, two-photon processes, ...



Baryons

- Covariant 3-quark Faddeev equation for baryons

GE, Alkofer, Nicmorus, Krassnigg, PRL 104 (2010)

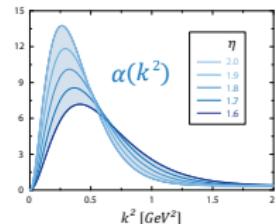


- Octet and decuplet spectra, elastic and transition form factors, ...

GE, Sanchis-Alepuz, Williams, Alkofer, Fischer,
Prog. Part. Nucl. Phys. 91 (2016), 1606.09602

- Heavy baryons

Qin, Roberts, Schmidt, PRD 97 (2018), FBS 60 (2019)



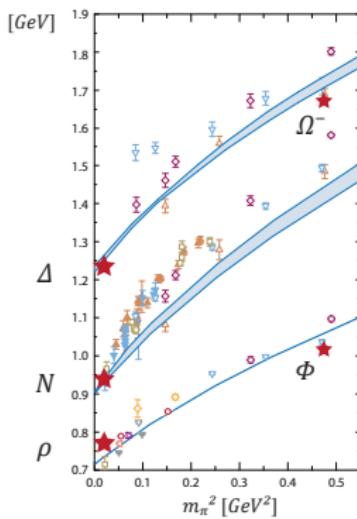
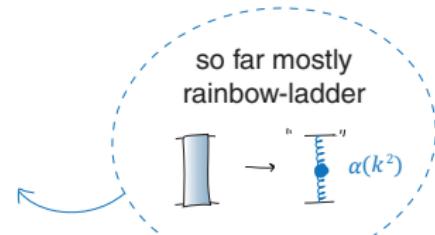
Maris, Tandy, PRC 60 (1999),
Qin et al., PRC 84 (2011)

Scale set by f_π ,
shape parameter
→ bands

Baryons

- Covariant 3-quark Faddeev equation for baryons

GE, Alkofer, Nicmorus, Krassnigg, PRL 104 (2010)



- Octet and decuplet spectra, elastic and transition form factors, ...

GE, Sanchis-Alepuz, Williams, Alkofer, Fischer,
Prog. Part. Nucl. Phys. 91 (2016), 1606.09602

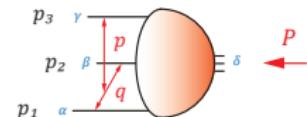
- Heavy baryons

Qin, Roberts, Schmidt, PRD 97 (2018), FBS 60 (2019)

- Keep full structure of baryon's Faddeev amplitude:

$$\Psi_{\alpha\beta\gamma\delta}(p, q, P) = \sum_i f_i(p^2, q^2, p \cdot q, p \cdot P, q \cdot P) \tau_i(p, q, P)_{\alpha\beta\gamma\delta}$$

Lorentz-invariant
dressing functions

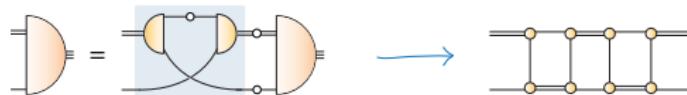


Dirac-Lorentz tensors:
64 (128) for spin 1/2 (3/2)

Diquark correlations

- Quark-diquark (two-body) equation

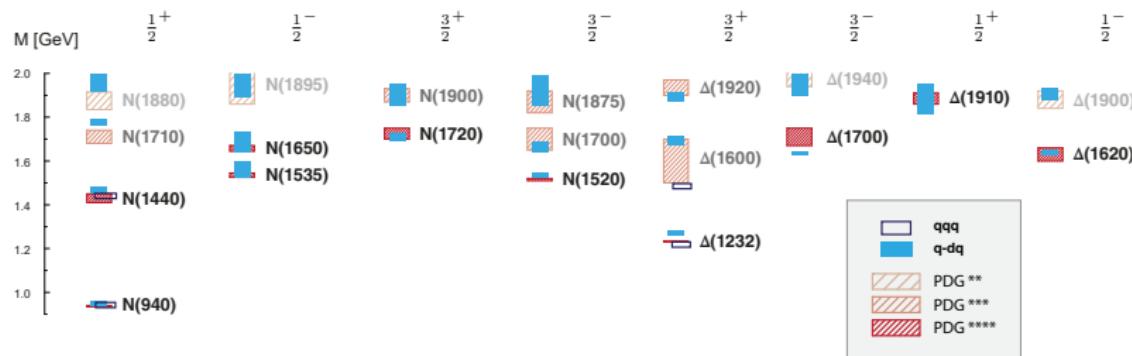
Oettel et al., PRC 58 (1998), GE et al., Ann. Phys. 323 (2008), Cloet et al., FBS 46 (2009), Segovia et al., PRL 115 (2015), Chen et al., PRD 97 (2018)



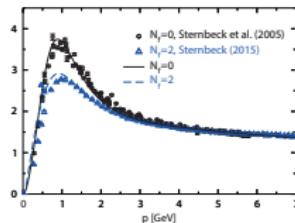
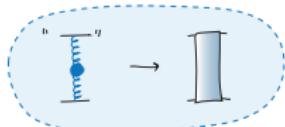
- Three-quark and quark-diquark results very similar

GE, Fischer, Sanchis-Alepuz, PRD 94 (2016)

Diquark clustering in baryons?
Barabanov et al., Prog. Part. Nucl. Phys. 116 (2021)



Towards ab-initio



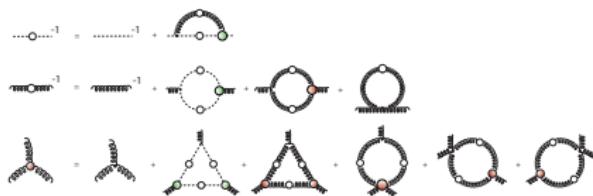
Gluon propagator: DSE vs. lattice

Williams, Fischer, Heupel,
PRD 93 (2016)

See also: Aguilar, De Soto, Ferreira,
Papavassiliou, Rodriguez-Quintero,
Zafeiropoulos, EPJ C 80 (2020)

- Compute higher n-point functions from **DSEs, FRG, lattice QCD**

Binosi, Ibanez, Papavassiliou, JHEP 09 (2014),
GE, Williams, Alkofer, Vujinovic, PRD 89 (2014),
Williams, EPJA 51 (2015), Huber, PRD 101 (2020),
Cyrol, Mitter, Pawłowski, Strodthoff, PRD 97 (2018),
Oliveira, Silva, Skullerud, Sternbeck, PRD 99 (2019), ...

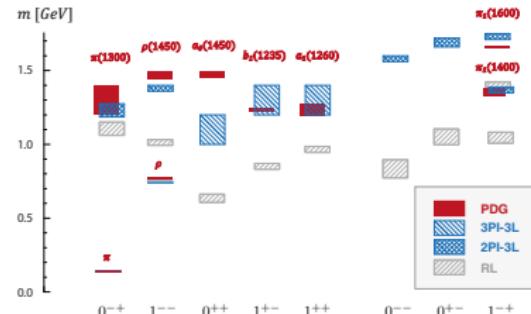


- Kernels from **chiral symmetry** constraints

Chang, Roberts, PRL 103 (2009), Chang, Liu, Roberts, PRL 106 (2001),
Qin, Roberts, 2009.13637

Beyond rainbow-ladder calculations improve **light-meson spectrum**

Williams, Fischer, Heupel, PRD 93 (2016)



Towards ab-initio

- **Glueballs** in pure Yang-Mills theory (no quarks)

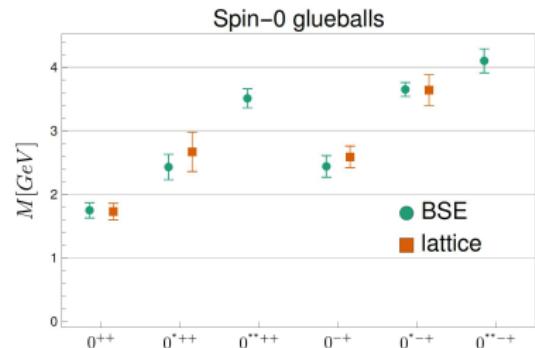
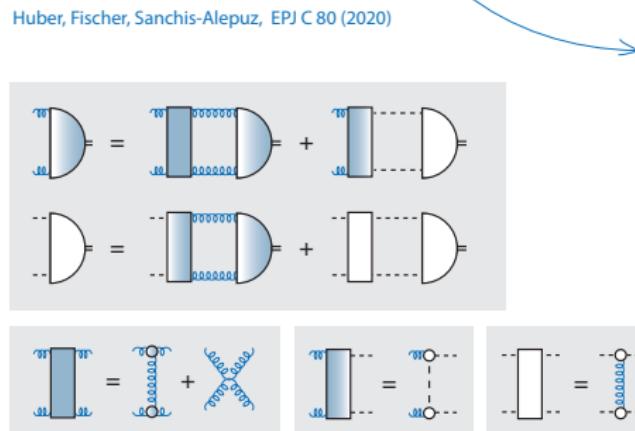
Meyers, Swanson, PRD 87 (2013)

Sanchis-Alepuz, Fischer, Kellermann, Smekal, PRD 92 (2015)

Souza, Ferreira, Aguilar, Papavassiliou, Roberts, Xu, EPJA 56 (2020)

Kernel derived from n-point functions,
no model input

Huber, Fischer, Sanchis-Alepuz, EPJ C 80 (2020)



Lattice:

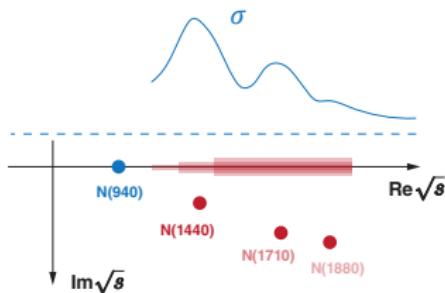
Morningstar, Peardon, PRD 60 (1999)

Chen, Alexandru, Dong, Draper, Horvath, PRD 73 (2006)

Athenodorou, Teper, JHEP 11 (2020)

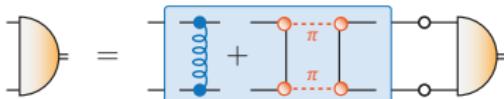
Resonances

- Most hadrons are **resonances** and decay
 \Leftrightarrow poles in complex momentum plane



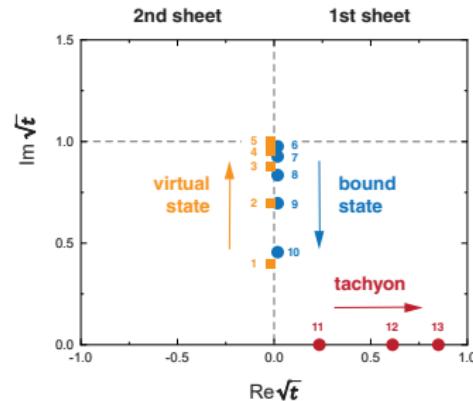
- BSE kernel must be aware of decay channels:
 ρ meson becomes resonance

Williams, PLB 798 (2019), Miramontes, Sanchis-Alepuz, EPJA 55 (2019),
Miramontes, Sanchis-Alepuz, Alkofer, 2102.12541



- Contour deformations** as tool to go beyond thresholds

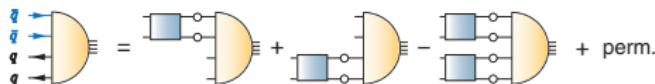
GE, Duarte, Peña, Stadler, PRD 100 (2019)



Four-quark states

- Light scalar mesons (σ, κ, a_0, f_0) as **four-quark states**:

GE, Fischer, Heupel, PLB 753 (2016)



$$\Gamma(p, q, k, P) = \sum_i f_i(p^2, q^2, k^2, \{\omega_j\}, \{\eta_j\}) \tau_i(p, q, k, P) \otimes \text{Color} \otimes \text{Flavor}$$

9 Lorentz invariants:

$$p^2, \quad q^2, \quad k^2, \quad P^2 = -M^2$$

$$\omega_1 = q \cdot k \quad \eta_1 = p \cdot P$$

$$\omega_2 = p \cdot k \quad \eta_2 = q \cdot P$$

$$\omega_3 = p \cdot q \quad \eta_3 = k \cdot P$$

256 Dirac-Lorentz tensors

2 Color tensors:
 $3 \otimes \bar{3}, \quad 6 \otimes \bar{6}$ or
 $1 \otimes 1, \quad 8 \otimes 8$
(Fierz-equivalent)

$$K \psi_i = \lambda_i \psi_i$$

	dim K	memory
Mesons	10^3	20 MB
Baryons	10^8	10^7 GB
Tetraquarks	10^{13}	10^{18} GB

Four-quark states

- Light scalar mesons (σ, κ, a_0, f_0) as **four-quark states**:

GE, Fischer, Heupel, PLB 753 (2016)

$$\begin{array}{c} q \\ \bar{q} \\ q \\ \bar{q} \end{array} \equiv = \begin{array}{c} \text{---} \\ \square \\ \text{---} \end{array} + \begin{array}{c} \text{---} \\ \square \\ \text{---} \end{array} - \begin{array}{c} \text{---} \\ \square \\ \text{---} \end{array} + \text{perm.}$$

$$\Gamma(p, q, k, P) = \sum_i f_i(p^2, q^2, k^2, \{\omega_j\}, \{\eta_j\}) \tau_i(p, q, k, P) \otimes \text{Color} \otimes \text{Flavor}$$

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256 Dirac-Lorentz tensors

2 Color tensors:

$$3 \otimes \overline{3}, \quad 6 \otimes \overline{6} \quad \text{or} \\ 1 \otimes 1, \quad 8 \otimes 8 \\ (\text{Fierz-equivalent})$$

- Group momentum variables into multiplets of **permutation group S4**: can switch off groups of variables without destroying symmetries

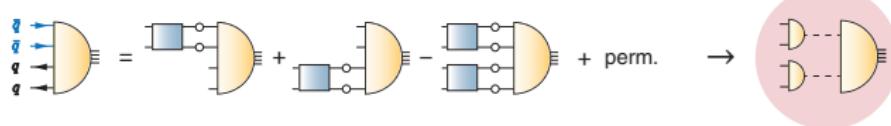
GE, Fischer, Heupel, PRD 92 (2015)

$$f_i(S_0, \nabla, \triangle, \circ)$$

Four-quark states

- Light scalar mesons (σ, κ, a_0, f_0) as **four-quark states**:

GE, Fischer, Heupel, PLB 753 (2016)



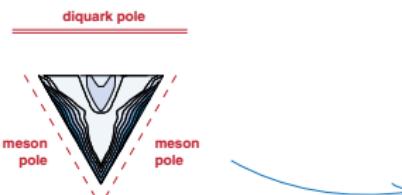
- BSE dynamically generates **meson poles** in BS amplitude:

$$f_i(S_0, \nabla, \Delta, \circ) \rightarrow 1500 \text{ MeV}$$

$$f_i(S_0, \nabla, \Delta, \circ) \rightarrow 1500 \text{ MeV}$$

$$f_i(S_0, \nabla, \Delta, \circ) \rightarrow 1200 \text{ MeV}$$

$$f_i(S_0, \nabla, \Delta, \circ) \rightarrow 350 \text{ MeV !}$$

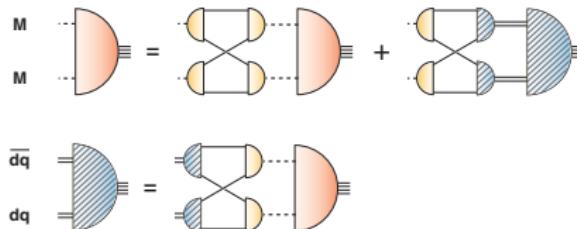


- "Light scalar mesons" look like **meson molecules**, diquark-antidiquark components almost negligible
- Lightness is inherited from pseudoscalar Goldstone bosons!

Four-quark states

Two-body formulation: **meson-meson / diquark-antidiquark**,
follows from four-quark eq. (analogue of quark-diquark for baryons)

Heupel, GE, Fischer, PLB 718 (2012)

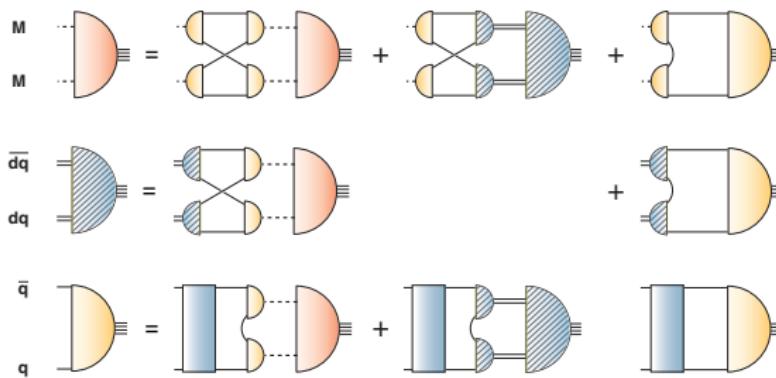


- Interaction by **quark exchange**
- System 'wants' to be **meson-meson-like** (no diagonal dq-dq term)
- Similar results as in 4-quark approach:
 $m_\sigma \sim 400$ MeV, etc.

Four-quark states

Two-body formulation: **meson-meson / diquark-antidiquark**,
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Heupel, GE, Fischer, PLB 718 (2012)



Include mixing with $q\bar{q}$:
 $\pi\pi$ still dominant

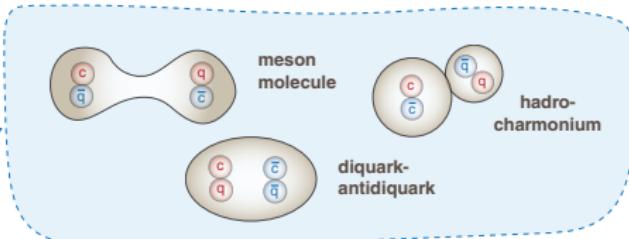
Santowsky, GE, Fischer, Wallbott,
Williams, PRD 102 (2020)

[MeV]	ground state mass	first excitation
$\pi\pi$	416 ± 26	970 ± 130
$\pi\pi + 0^+ 0^+$	416 ± 26	970 ± 130
$q\bar{q}$	667 ± 2	1036 ± 8
$\pi\pi + q\bar{q}$	472 ± 22	1080 ± 280
$\pi\pi + 0^+ 0^+ + q\bar{q}$	456 ± 24	1110 ± 110

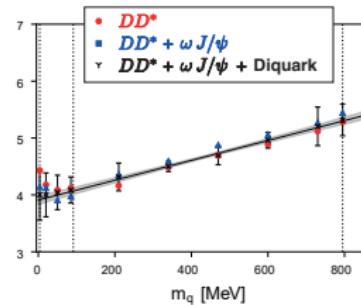
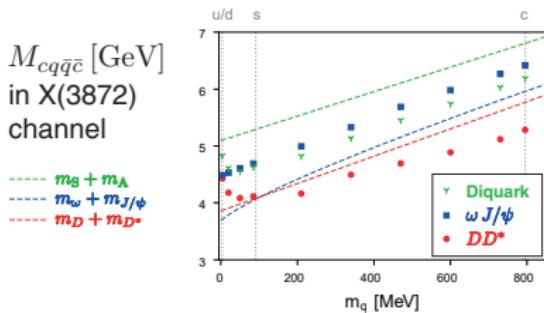
Four-quark states

- Heavy-light four-quark states:
what is their internal decomposition?

$cq\bar{q}\bar{c}$



- Four-quark BSE: all mix together



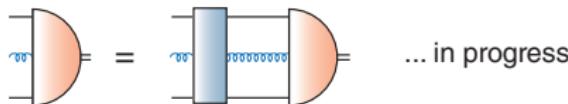
Wallbott, GE, Fischer,
PRD 100 (2019),
PRD 102 (2020)

$cq\bar{q}\bar{c} \rightarrow$ strong meson-meson component: DD^* for $X(3872)$, $Z_c(3900)$:

$cc\bar{q}\bar{q} \rightarrow$ diquarks also play role

Hybrids

- **Three-body equation** (quark, antiquark, gluon)



- **Two-body equation** [quark–gluon]–antiquark
with model ansätze

Xu, Cui, Chang, Papavassiliou, Roberts, EPJA 55 (2019)

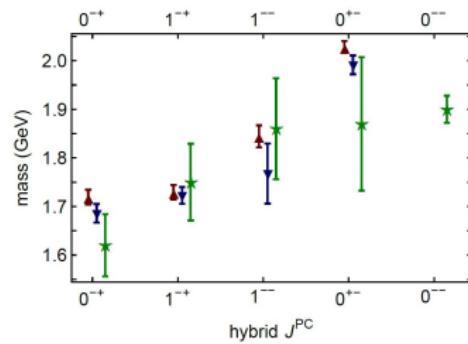
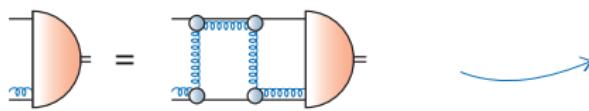
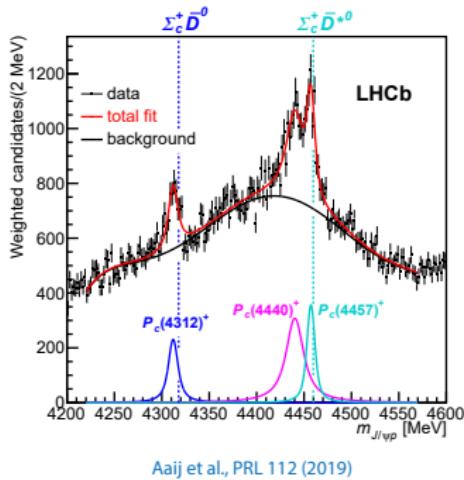


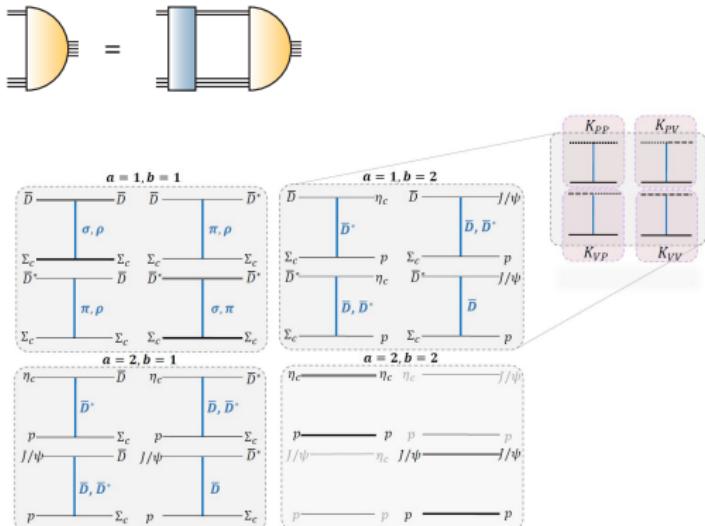
FIG. 2. Comparison between our ACM-improved spectrum (stars, green), Row 2 in Table I, and the rescaled lQCD results in Rows 3 (up-triangles, red) and 4 (down-triangles, blue).

Lattice: Dudek, Edwards, Peardon, Richards, Thomas, PRD 82 (2010)

Pentaquarks?

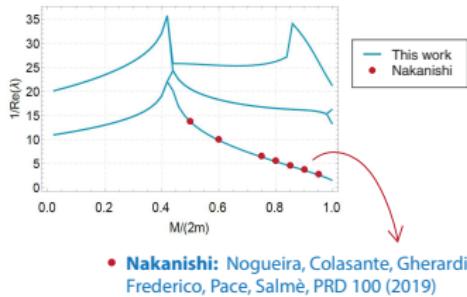


- Meson-baryon equation with hadronic exchanges
GE, Lourenco, Stadler, in preparation



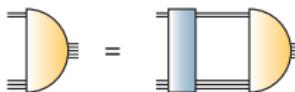
Pentaquarks?

- Many crossings of real and complex conjugate eigenvalues

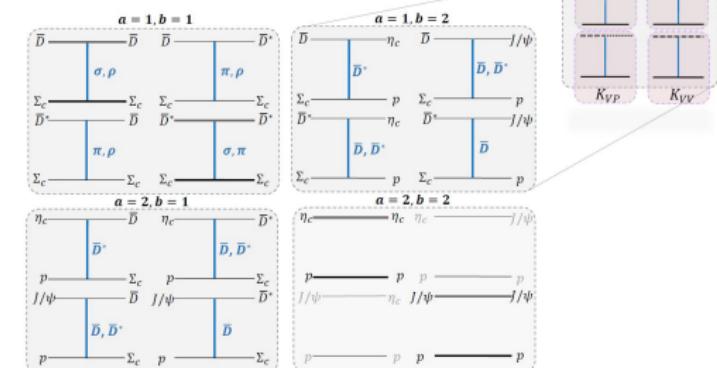
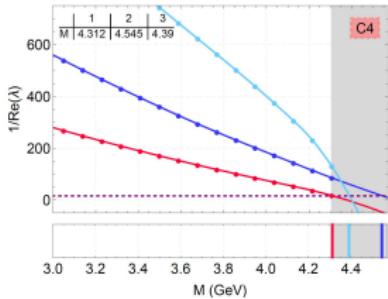


- Meson-baryon equation with hadronic exchanges

GE, Lourenco, Stadler, in preparation



- With spectral reconstruction:



The road ahead

- **Four-quark states:**

GE, Fischer, Heupel, Santowsky, Wallbott, FBS 61 (2020), arXiv:2008.10240

- Towards ab-initio calculations
(include higher n-point functions)
- Resonance properties
- Hybrid mesons
- Exotic baryons ~ **pentaquarks**:
Meson-baryon interactions, admixture of 5q (...) components?
- Exotic dibaryons ~ **hexaquarks**:
from quarks and gluons to nuclei

Thank you!

Backup slides

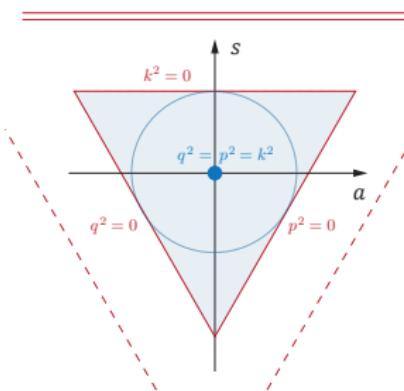
Structure of the amplitude

- **Singlet:** symmetric variable, carries overall scale:

$$S_0 = \frac{1}{4} (p^2 + q^2 + k^2)$$

- **Doublet:** $\mathcal{D}_0 = \frac{1}{4S_0} \begin{bmatrix} \sqrt{3}(q^2 - p^2) \\ p^2 + q^2 - 2k^2 \end{bmatrix}$

Mandelstam triangle,
outside: **meson and diquark poles!**

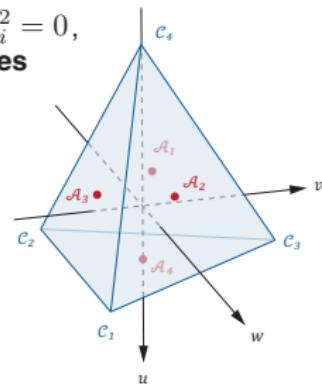


Lorentz invariants can be grouped into
multiplets of the permutation group S_4 :

GE, Fischer, Heupel, PRD 92 (2015)

- **Triplet:** $\mathcal{T}_0 = \frac{1}{4S_0} \begin{bmatrix} 2(\omega_1 + \omega_2 + \omega_3) \\ \sqrt{2}(\omega_1 + \omega_2 - 2\omega_3) \\ \sqrt{6}(\omega_2 - \omega_1) \end{bmatrix}$

tetrahedron bounded by $p_i^2 = 0$,
outside: **quark singularities**



- **Second triplet:**
3dim. sphere

$$\mathcal{T}_1 = \frac{1}{4S_0} \begin{bmatrix} 2(\eta_1 + \eta_2 + \eta_3) \\ \sqrt{2}(\eta_1 + \eta_2 - 2\eta_3) \\ \sqrt{6}(\eta_2 - \eta_1) \end{bmatrix}$$

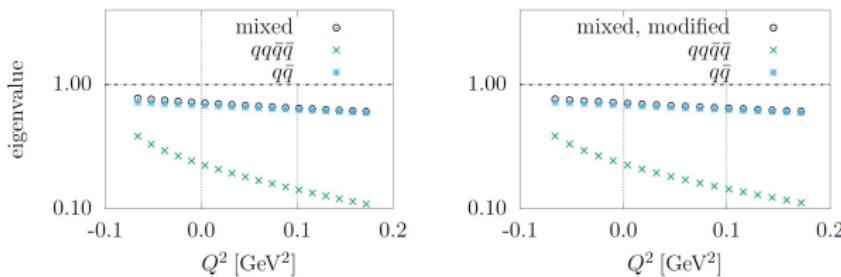
Reply to comment

- Recent comment: BSE kernel is reducible
[Blankleider & Kvinikhidze, 2102.05818](#)
- Irrelevant for homogeneous BSE: same spectrum



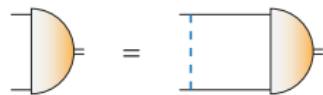
- Results are identical within numerical errors

[Santowsky, GE, Fischer, Wallbott, Williams, 2103.14673](#)



Bound states & resonances

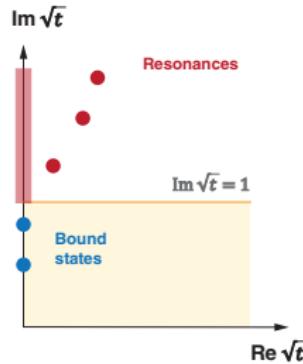
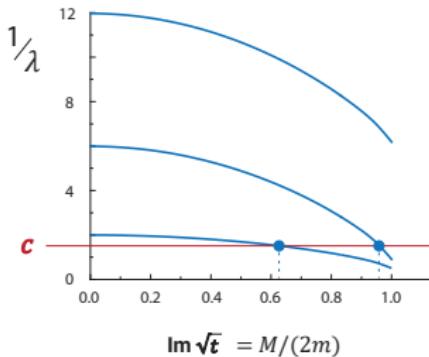
- Homogeneous BSE:



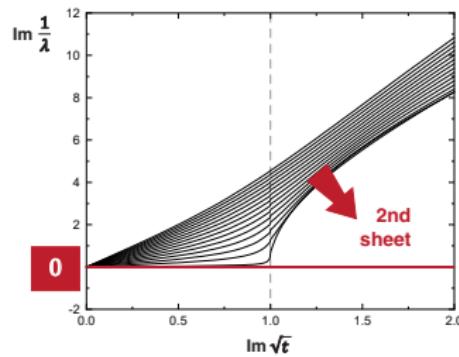
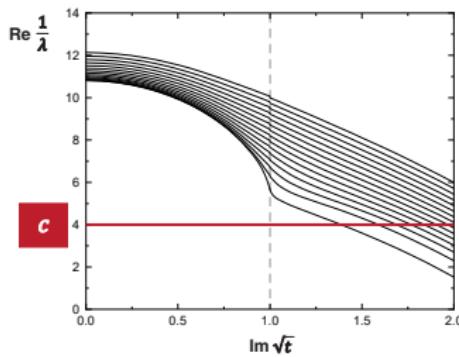
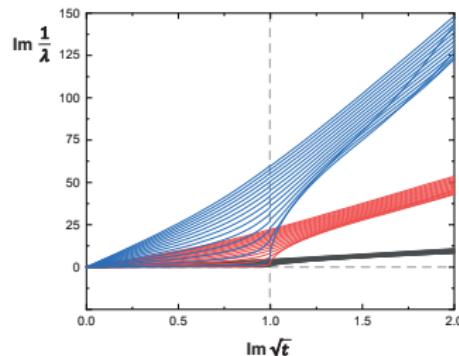
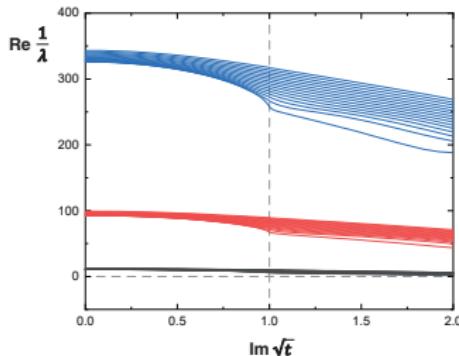
$$\Rightarrow \psi(t) = \mathbf{c} K G_0(t) \psi(t)$$

$$\psi(X, Z, t) = \mathbf{c} \int dx \int dz \ K(X, x, Z, z, t) \ G_0(x, z, t) \ \psi(x, z, t)$$

$$\Rightarrow \frac{1}{\lambda(t)} = \mathbf{c}$$



BSE Eigenvalues



$$\frac{1}{\lambda(t)} = c + 0 \cdot i$$

still valid for
complex poles:
can detect
resonances from
homogeneous BSE

Complex eigenvalues?

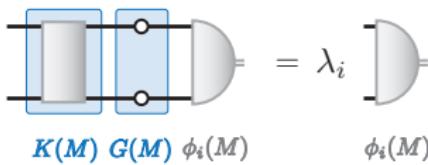
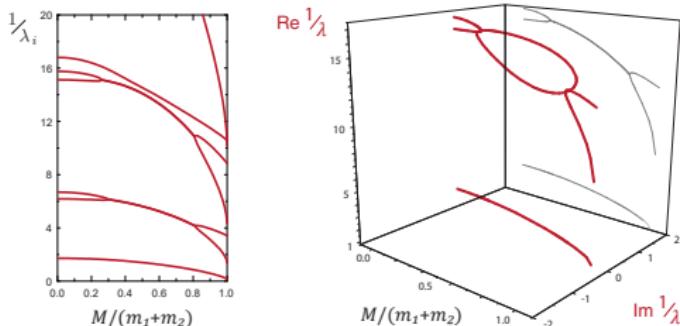
Excited states: some EVs are complex conjugate?

Typical for **unequal-mass** systems, already in Wick-Cutkosky model

Wick 1954, Cutkosky 1954

Connection with “**anomalous**” states?

Ahlig, Alkofer, Ann. Phys. 275 (1999)



If $G = G^\dagger$ and $G > 0$:
Cholesky decomposition $G = L^\dagger L$

$$K L^\dagger L \phi_i = \lambda_i \phi_i \\ (LKL^\dagger)(L\phi_i) = \lambda_i (L\phi_i)$$

⇒ Hermitian problem with same EVs!

K and G are Hermitian (even for unequal masses!) but KG is not

Complex eigenvalues?

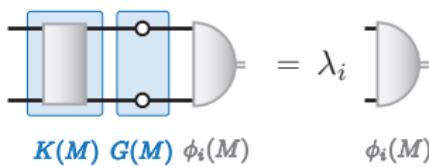
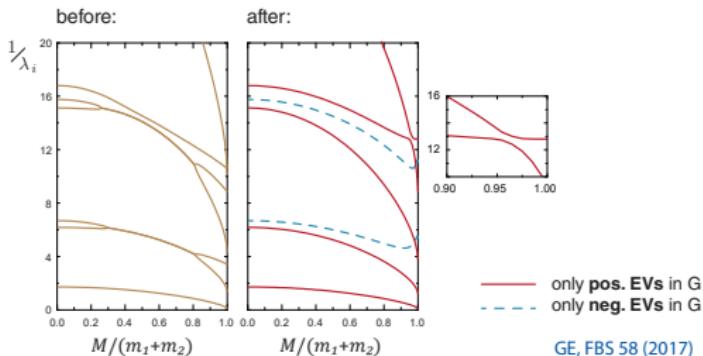
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⇒ Hermitian problem with same EVs!

- ⇒ all EVs strictly **real**
- ⇒ level repulsion
- ⇒ “anomalous states” removed?

Complex eigenvalues?

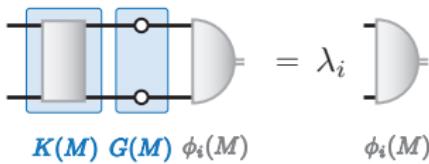
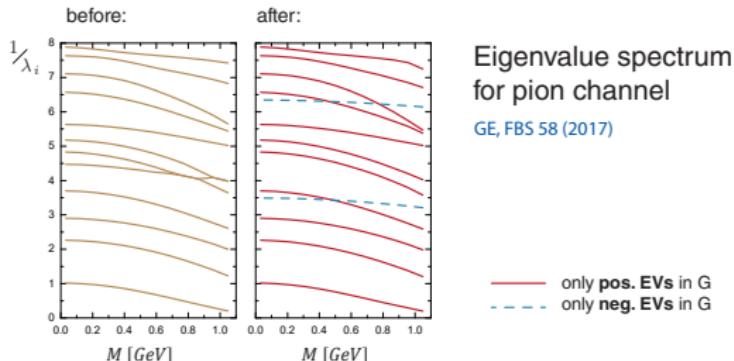
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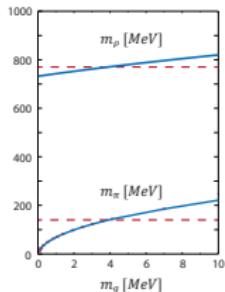
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⇒ Hermitian problem with same EVs!

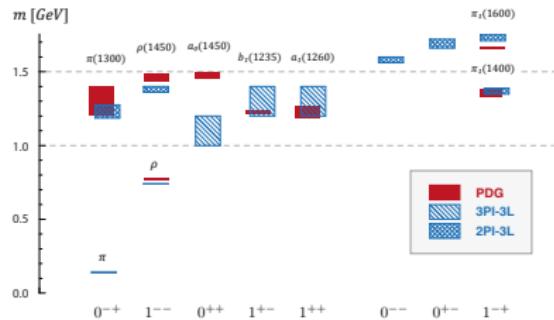
- ⇒ all EVs strictly **real**
- ⇒ level repulsion
- ⇒ “anomalous states” removed?

Mesons

- Pion is **Goldstone boson**: $m_\pi^2 \sim m_q$



- Light meson spectrum beyond rainbow-ladder**

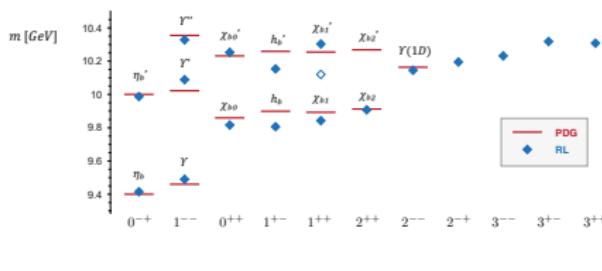


Williams, Fischer, Heupel,
PRD 93 (2016)

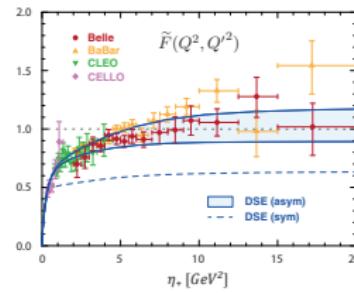
GE, Sanchis-Alepuz, Williams,
Alkofer, Fischer, PPNP 91 (2016)

- Bottomonium spectrum**

Fischer, Kubrak, Williams, EPJ A 51 (2015)



- Pion transition form factor**



GE, Fischer, Weil, Williams,
PLB 774 (2017)

Bound-state equations

Bethe-Salpeter equation for baryons: GE, Sanchis-Alepuz, Williams, Alkofer, Fischer, PPNP 91 (2016), 1606.09602

$$\text{Diagram} = \text{Diagram} + \text{Diagram} + \text{Diagram} + \text{Diagram}$$

Quark-diquark approximation:

$$\text{Diagram} = \text{Diagram}$$

Rainbow-ladder:

$$\text{Diagram}^{-1} = \text{Diagram}^{-1} + \text{Diagram}$$

$$\text{Diagram} = \text{Diagram}$$

Maris, Tandy, PRC 60 (1999),
Qin et al., PRC 84 (2011)

$$\begin{aligned}\text{Diagram} &= \text{Diagram} \\ \text{Diagram}^{-1} &= \text{Diagram} + \text{Diagram}\end{aligned}$$

DSE / BSE / Faddeev landscape

	NJL / contact	q-dq model	DSE (RL)	DSE (bRL)
u/d	N, Δ masses N, Δ em. FFs $N \rightarrow \Delta \gamma$	✓ ✓ ✓	✓ ✓ ✓	✓ ✓ ✓
	N^*, Δ^* masses (+) $N \rightarrow N^* \gamma$	✓ ✓	✓ ✓	✓ ✓
	N^*, Δ^* masses (-) $N \rightarrow N^* \gamma$	✓ ✓	✓ ✓	✓ ✓
s	ground states excited states em. FFs & TFFs	✓ ✓ ✓	✓ ✓ ✓	✓ ✓ ✓
c, b	ground states excited states	✓ ✓	✓ ✓	✓ ✓

✓ ... before 2015

✓ ... after 2015

Cloet, Thomas,
Roberts, Bashir,
Segovia, Chen,
Wilson, Lu, ...

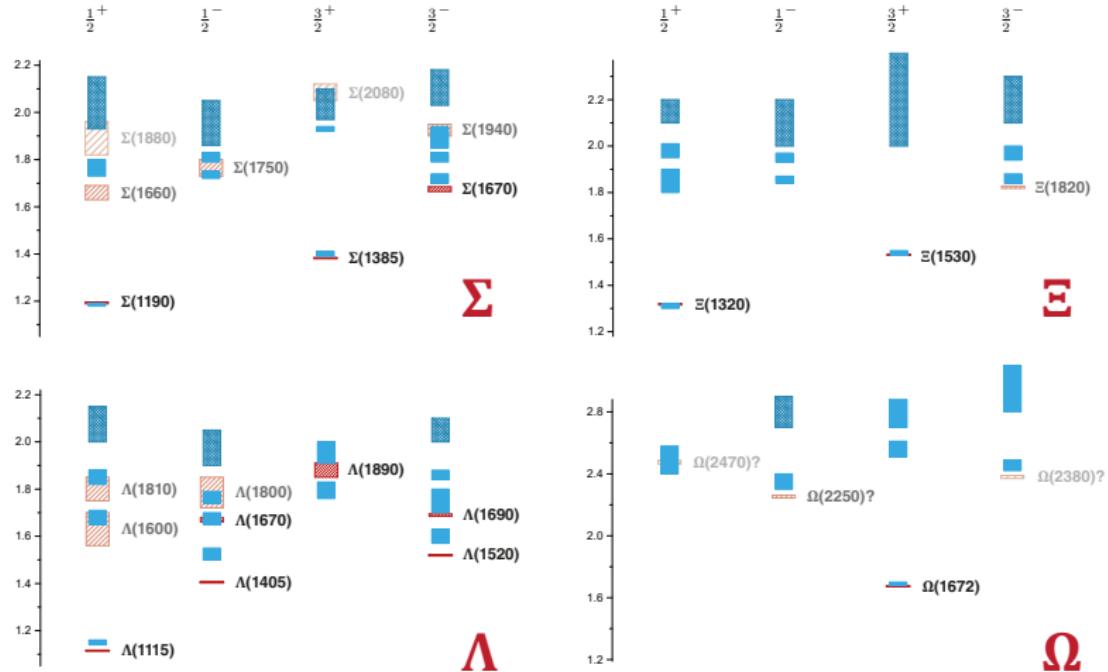
Oettel, Alkofer,
Roberts, Cloet,
Segovia, Chen,
El-Bennich, ...

GE, Alkofer,
Nicmorus,
Sanchis-Alepuz,
Fischer

GE, Sanchis-Alepuz,
Fischer, Alkofer,
Qin, Roberts

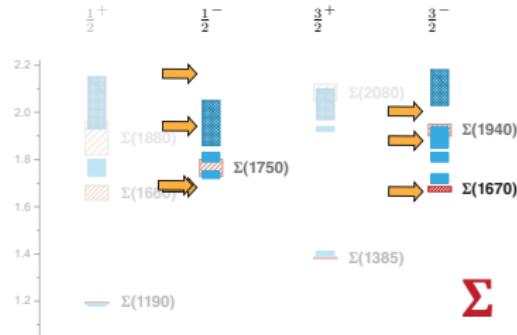
Sanchis-Alepuz,
Williams, Fischer

Strange baryons

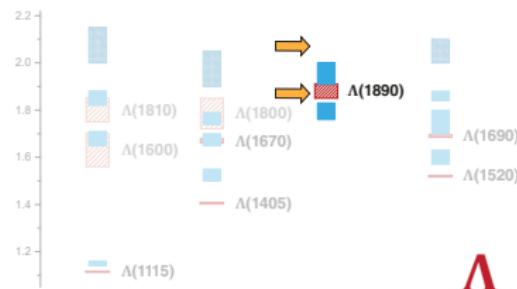


GE, Fischer, FBS 60 (2019), Fischer, GE, PoS Hadron 2017

Strange baryons



New states from Bonn-Gatchina
Sarantsev et al., 1907.13387 [nucl-ex]



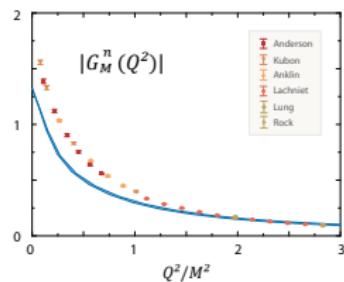
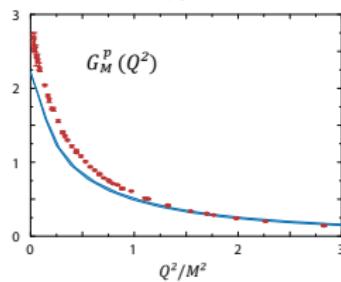
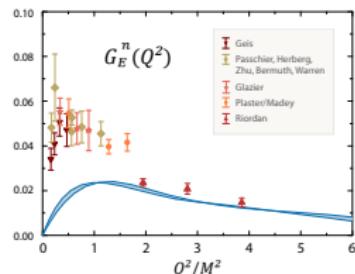
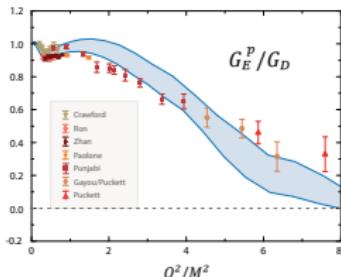
GE, Fischer, FBS 60 (2019), Fischer, GE, PoS Hadron 2017

Baryon form factors

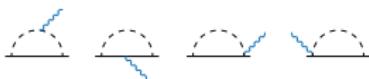
Nucleon em. form factors from three-quark equation

GE, PRD 84 (2011)

$$J^\mu = \text{Diagram 1} + \text{Diagram 2} + \text{Diagram 3} + \text{Diagram 4}$$



- “Quark core without pion cloud”

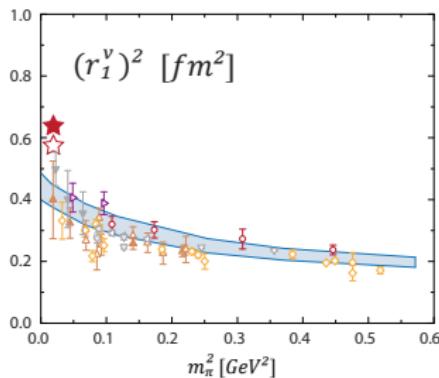


- similar: $N \rightarrow \Delta\gamma$ transition, axial & pseudoscalar FFs, octet & decuplet em. FFs

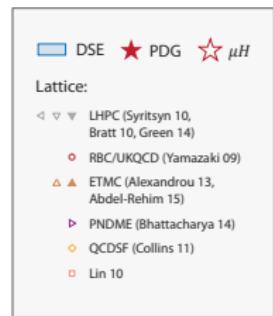
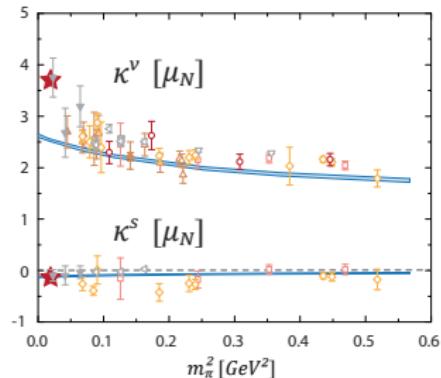
Review: GE, Sanchis-Alepuz, Williams, Fischer, Alkofer, PPNP 91 (2016), 1606.09602

Nucleon em. form factors

Nucleon charge radii:
isovector (p-n) Dirac (F1) radius



Nucleon magnetic moments:
isovector (p-n), isoscalar (p+n)



- Pion-cloud effects missing
(\Rightarrow divergence!), agreement with lattice at larger quark masses.



- But: pion-cloud cancels in $\kappa^s \Leftrightarrow$ quark core

Exp: $\kappa^s = -0.12$

Calc: $\kappa^s = -0.12(1)$

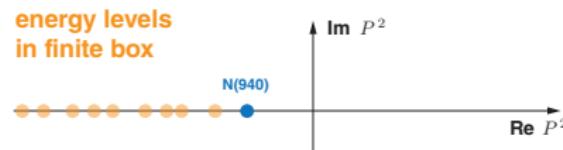


GE, PRD 84 (2011)

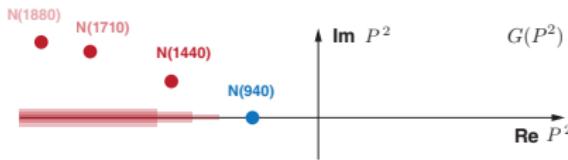
Resonances?

Lattice QCD:

$$\langle \dots \rangle = \int \mathcal{D}[\psi, \bar{\psi}, A] e^{-S[\psi, \bar{\psi}, A]} (\dots)$$



- **Finite volume:**
bound states & scattering states



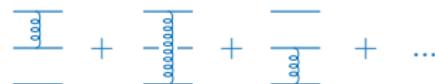
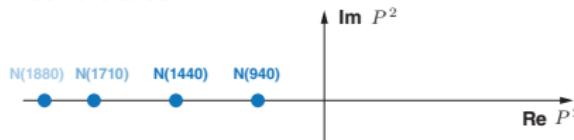
vary volume,
Luescher method

- **Infinite volume:**
Bound states, resonances,
branch cuts

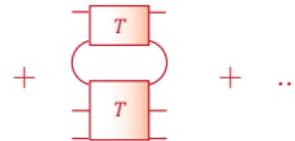
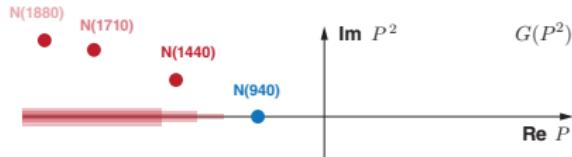
Resonances?

In terms of quarks and gluons?

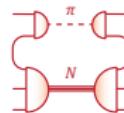
Bound states:



Resonances by meson-baryon interactions:



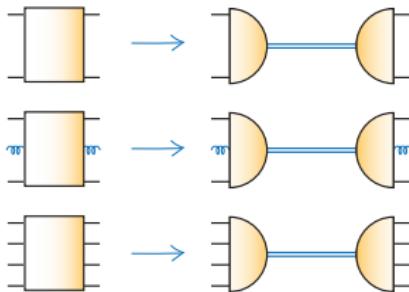
Both **bound states** and **resonances**
must be generated from quark-gluon structure!



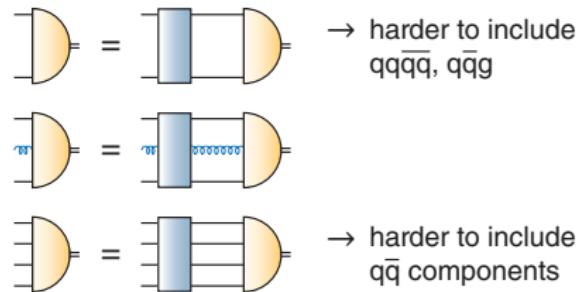
Analogue for $\rho \rightarrow \pi\pi$:
Williams, 1804.11161 [hep-ph],
Miramontes, Sanchis-Alepuz,
1906.06227 [hep-ph]

Some thoughts

- Same spectral representation for **all** correlation functions that produce $q\bar{q}$:



- Without** truncations, BSEs for qq , $qq\bar{q}\bar{q}$, $q\bar{q}g$, ... should produce **same spectrum**



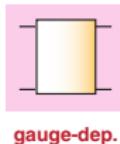
With truncations, $q\bar{q}$ ($qq\bar{q}\bar{q}$, $q\bar{q}g$, ...) BSE should give more reliable spectrum for $q\bar{q}$ ($qq\bar{q}\bar{q}$, $q\bar{q}g$, ...) dominated states

Exotic quantum numbers not excluded in principle (need gluon-rich kernel in $q\bar{q}$ BSE)

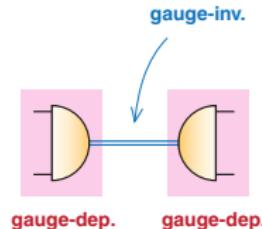
Some thoughts

- Why don't we see exotic quantum numbers in **lattice calculations** with $q\bar{q}$ operators?

$$\langle q_\alpha \bar{q}_\beta q_\gamma \bar{q}_\delta \rangle =$$



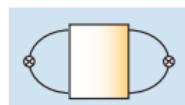
gauge-dep.



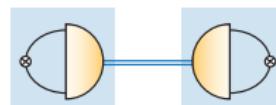
gauge-dep.

gauge-inv.

$$\langle (\bar{q} \Gamma q)(\bar{q} \Gamma q) \rangle =$$



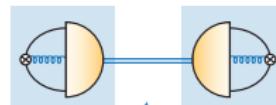
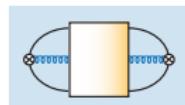
gauge-inv.



gauge-inv.

gauge-inv.

= 0 for exotic
quantum nrs (?)



gauge-inv.

gauge-inv.

≠ 0 for exotic
quantum nrs

Could test with **gauge-fixed**
lattice calculations: same spectrum?

