

# Neutron star mergers as materials science

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Alford, Bovard, Hanauske, Rezzolla, Schwenzer, [arXiv:1707.09475](#)

Alford, Harutyunyan, Sedrakian, [arXiv:1907.04192](#), [2006.07975](#)

Alford and Harris, [arXiv:1803.00662](#), [arXiv:1907.03795](#)

Alford and Haber, [arXiv:2009.05181](#)



U.S. DEPARTMENT OF  
**ENERGY**

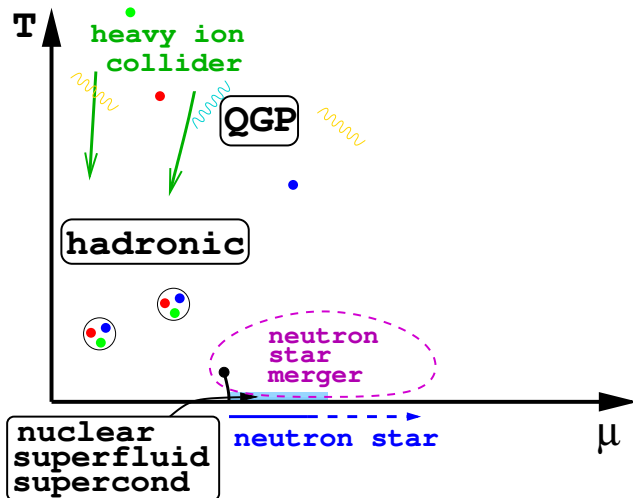
Office of  
Science

# Outline

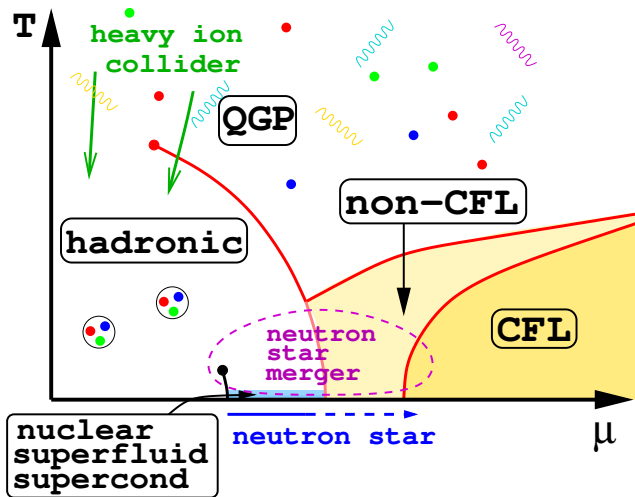
- I Neutron star mergers as a probe of dense matter:  
Different phases with similar equation of state may be distinguishable by their **Transport Properties**
- II **Thermal conductivity** in mergers  
Damping time for temperature inhomogeneities: is it fast enough to affect mergers?
- III **Bulk viscosity** in mergers  
Damping time for density oscillations: is it fast enough to affect mergers?
- IV Summary and Prospects

# I. Mergers as a probe of dense matter

# Observed QCD Phase diagram



# Conjectured QCD Phase diagram



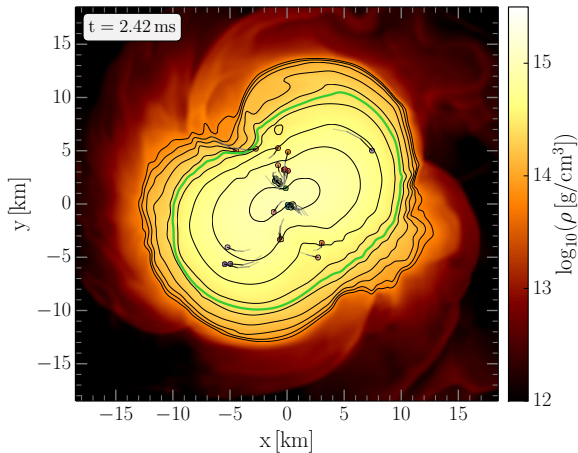
**heavy ion collisions**: deconfinement crossover and chiral critical point

**neutron stars**: quark matter core?

**neutron star mergers**: dynamics of warm and dense matter

# Neutron star mergers

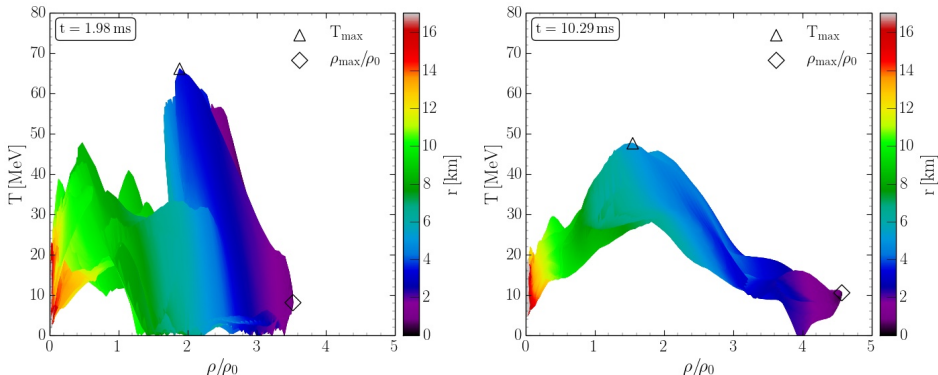
Mergers probe the properties of nuclear/quark matter at high density (up to  $\sim 4n_{\text{sat}}$ ) and temperature (up to  $\sim 80$  MeV)



Which  
**material properties**  
are relevant?

- Equation of state  $\rho(p)$
- And...?

# Nuclear material in a neutron star merger



M. Hanauske, Rezzolla group, Frankfurt

Significant spatial/temporal variation in:

temperature  
fluid flow velocity  
density

so we need to allow for  
thermal conductivity  
shear viscosity  
bulk viscosity

# Role of transport/dissipation in mergers

The important dissipation mechanisms are the ones whose equilibration time is  $\lesssim 20$  ms

## Executive Summary:

### Thermal conductivity:

Might play a role, if

- neutrinos are trapped
- there are short-distance temperature gradients

Shear viscosity: Similar conclusion.

### Bulk viscosity:

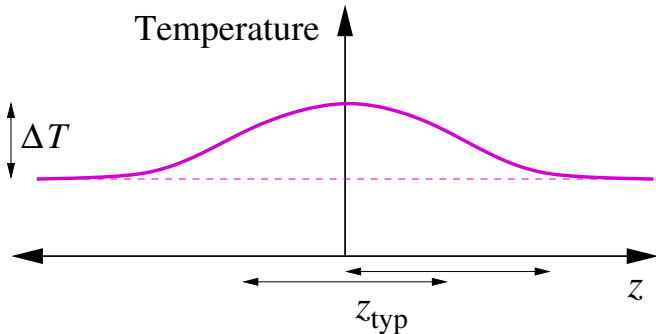
Could damp density oscillations on the same timescale as the merger.

Include bulk viscosity in merger simulations.



## **II. Thermal equilibration under merger conditions**

# Thermal equilibration time



Volume

$$V \sim z_{\text{typ}}^3$$

Surface area

$$A \sim 6z_{\text{typ}}^2$$

Time to equilibrate:  $\tau_{\kappa} = \frac{\text{extra heat in region}}{\text{rate of heat outflow}} = \frac{E_{\text{therm}}}{W_{\text{therm}}}$

Thermal diffusion is important if  $\tau_{\kappa} \lesssim 20 \text{ ms}$

# Estimating thermal equilibration time

Extra heat in region:  $E_{\text{therm}} = c_V V \Delta T \approx c_V z_{\text{typ}}^3 \Delta T$

Rate of heat outflow:  $W_{\text{therm}} = \kappa \frac{dT}{dz} A \approx \kappa \frac{\Delta T}{z_{\text{typ}}} 6z_{\text{typ}}^2$

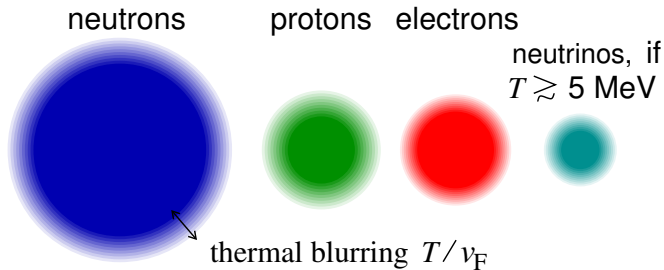
Time to equilibrate:  $\tau_{\kappa} = \frac{E_{\text{therm}}}{W_{\text{therm}}} \approx \frac{c_V z_{\text{typ}}^2}{6\kappa}$

We need to know

- specific heat capacity  $c_V$
- thermal conductivity  $\kappa$

# Nuclear material constituents

Fermi  
surfaces:



neutrons:  $\sim 90\%$  of baryons

$$p_{Fn} \sim 350 \text{ MeV}$$

protons:  $\sim 10\%$  of baryons

$$p_{Fp} \sim 150 \text{ MeV}$$

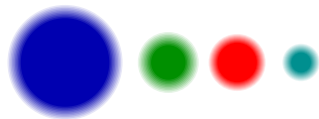
electrons: same density as protons

$$p_{Fe} = p_{Fp}$$

neutrinos: only present if  $\text{mfp} \ll 10 \text{ km}$

i.e. when  $T \gtrsim 5 \text{ MeV}$

# Specific heat capacity

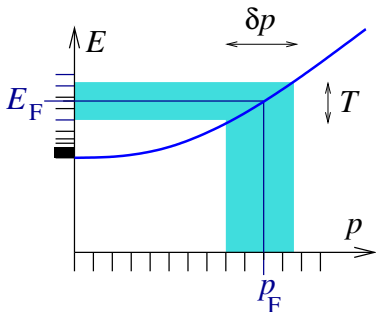


Dominated by neutrons

$c_V \sim$  number of states available  
to carry energy  $\lesssim T$

$\sim$  vol of mom space with states available to carry energy  $\lesssim T$

$\sim p_{Fn}^2 \delta p$



$$\delta p = \frac{T}{v_{Fn}} = T \times \frac{m_n^*}{p_{Fn}}$$

$$c_V \sim p_{Fn}^2 \delta p \sim p_{Fn}^2 \frac{m_n^*}{p_{Fn}} T \sim m_n^* p_{Fn} T$$

(Note: neutron density  $n_n \sim p_{Fn}^3$ )

$$c_V \approx 1.0 m_n^* n_n^{1/3} T$$

# Thermal conductivity

Thermal conductivity  $\kappa \propto n v \lambda$

Dominated by the species with the right combination of

- high density
- weak interactions  $\Rightarrow$  long mean free path (mfp)  $\lambda$

neutrons: high density, but strongly interacting (short mfp) ❌

protons: low density, strongly interacting (short mfp) ❌

electrons: low density, only EM interactions (long mfp) ✓

neutrinos:  $\left\{ \begin{array}{l} T \lesssim 5 \text{ MeV: } \lambda > \text{ size of merged stars, so} \\ \text{they all escape, density} = 0 \\ T \gtrsim 5 \text{ MeV: } \lambda < \text{ size of merged stars,} \\ \text{but still very long mfp!} \end{array} \right.$  ❌



# Neutrino-dominated thermal equilibration

Neutrino-trapped regime,  $T \gtrsim 5 \text{ MeV}$

equilibration time  $\tau_{\kappa} \approx \frac{c_V z_{\text{typ}}^2}{6\kappa} \quad \kappa^{(\nu)} \approx 0.33 \frac{n_\nu^{2/3}}{G_F^2 m_n^{*2} n_e^{1/3} T}$

$$\tau_{\kappa}^{(\nu)} \approx \boxed{700 \text{ ms}} \left( \frac{z_{\text{typ}}}{1 \text{ km}} \right)^2 \left( \frac{T}{10 \text{ MeV}} \right)^2 \left( \frac{\mu_e}{2\mu_\nu} \right)^2 \left( \frac{0.1}{x_p} \right)^{1/3} \left( \frac{m_n^*}{0.8 m_n} \right)^3$$

Neutrino thermal transport may be important if there are thermal gradients on  $\lesssim 0.1 \text{ km}$  scale

That can lead to  $\tau_{\kappa} \lesssim 20 \text{ ms}$

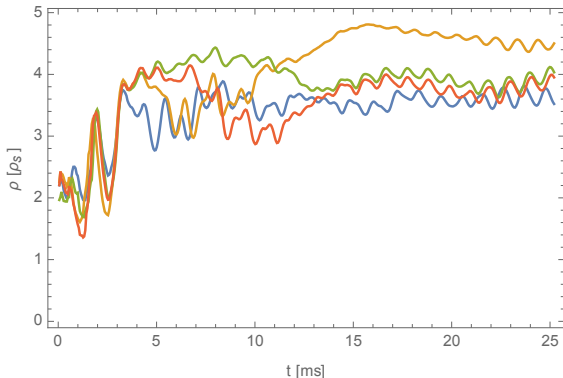
### **III. Damping of density oscillations in mergers**

Equivalently: sound attenuation via bulk viscosity



# Density oscillations in mergers

Density vs time for tracers in merger  
(Bulk viscosity neglected)



Tracers (co-moving fluid elements) show dramatic density oscillations, especially in the first 5 ms.

**Amplitude:** up to 50%

**Period:** 1–2 ms

How long does it take for bulk viscosity to dissipate a sizeable fraction of the energy of a density oscillation?

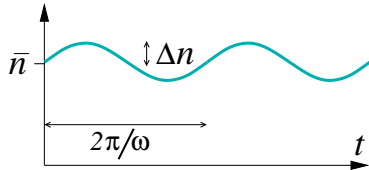
What is the damping time  $\tau_\zeta$ ?

Can we get  $\tau_\zeta \lesssim 20$  ms?

# Density oscillation damping time $\tau_\zeta$

Density oscillation of amplitude  $\Delta n$  at angular freq  $\omega$ :

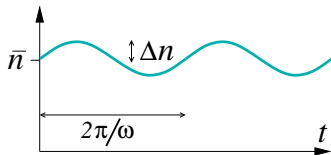
$$n(t) = \bar{n} + \Delta n \cos(\omega t)$$



$$\text{Damping Time: } \tau_\zeta = \frac{\text{energy stored in oscillation}}{\text{rate of energy loss}} = \frac{E_{\text{comp}}}{W_{\text{comp}}}$$

Bulk viscous damping is important if  $\tau_\zeta \lesssim 20$  ms

# Calculating damping time



Energy of density oscillation:  
( $K$  = nuclear incompressibility)

$$E_{\text{comp}} = \frac{K}{18} \bar{n} \left( \frac{\Delta n}{\bar{n}} \right)^2$$

Compression dissipation rate:  
( $\zeta$  = bulk viscosity)

$$W_{\text{comp}} = \zeta \frac{\omega^2}{2} \left( \frac{\Delta n}{\bar{n}} \right)^2$$

Damping Time: $\tau_{\zeta} = \frac{E_{\text{comp}}}{W_{\text{comp}}} = \frac{K \bar{n}}{9 \omega^2 \zeta}$
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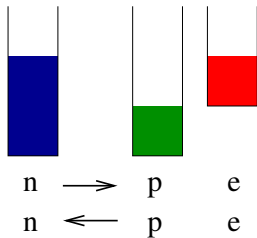
We need:

- nuclear incompressibility  $K$  (from EoS)
- bulk viscosity  $\zeta$  (from beta-equilibration of proton fraction)

# Bulk viscosity and beta equilibration

When you compress nuclear matter, the proton fraction wants to change.

Only **weak interactions** can change proton fraction, and they are rather slow...



	neutrino-transparent ( $T \lesssim 5 \text{ MeV}$ )*	neutrino-trapped ( $T \gtrsim 5 \text{ MeV}$ )*
neutron decay	$n \rightarrow p + e^- + \bar{\nu}_e$	$\nu_e + n \rightarrow p + e^-$
electron capture	$p + e^- \rightarrow n + \nu_e$	$p + e^- \rightarrow n + \nu_e$
	forward $\neq$ backward	$A + B \leftrightarrow C + D$

\* Neutrino transparency is a finite volume effect, which occurs when the neutrino mean free path is greater than the size of the system. Our system is a neutron star,  $R \sim 10 \text{ km}$

# Bulk viscosity: phase lag in system response

Some property of the material (proton fraction) takes time to equilibrate.

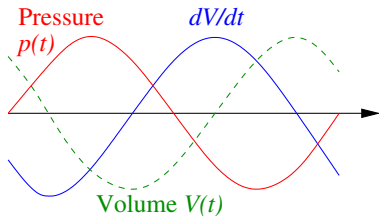
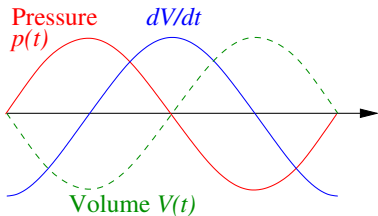
Baryon density  $n$  and hence fluid element volume  $V$  gets out of phase with applied pressure  $p$ :

$$\text{Dissipation} = - \int p dV = - \int p \frac{dV}{dt} dt$$

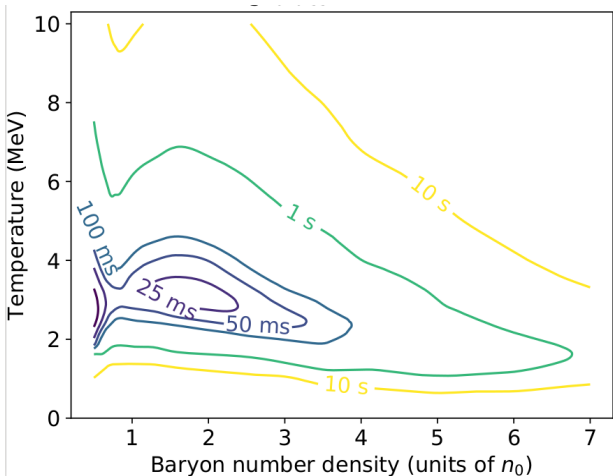
No phase lag.  
Dissipation = 0



Some phase lag.  
Dissipation > 0



# Damping time results ( $\nu$ -transparent)



EoS: IUFSU

$M_{\text{max}}: 1.96 M_{\odot}$

$R_{1.4 M_{\odot}}: 12.8 \text{ km}$

Oscillation freq:

$f = 1 \text{ kHz}$

Fast damping at  $T \sim 3 \text{ MeV}$ ,  $n \lesssim 2n_{\text{sat}}$

## IV. Summary and Prospects

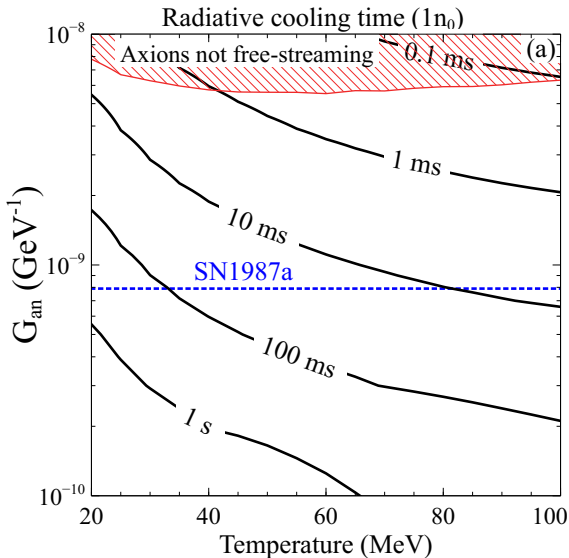
- ▶ Some forms of dissipation are probably physically important for neutron star mergers.
- ▶ **Thermal conductivity** and **shear viscosity** may become significant in the neutrino-trapped regime ( $T \gtrsim 5 \text{ MeV}$ ) if there are fine-scale gradients ( $z \lesssim 100 \text{ m}$ ).
- ▶ **Bulk viscosity** quickly damps density oscillations in neutrino-transparent nuclear matter (at low density and  $T \sim 3 \text{ MeV}$ ).

Prospects:

- ▶ Include bulk viscosity in merger simulations.
- ▶ Calculate bulk viscous damping for other forms of matter: hyperonic, muonic, pion condensed, nuclear pasta, quark matter, etc
- ▶ Beyond Standard Model physics...?

# Cooling by axion emission

Time for a hot region to cool to half its original temperature



Harris, Fortin, Sinha, Alford  
arXiv:2003.09768