

Parameterization of amplitudes, finding resonances and unitarity, peculiarities and traps

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Matter To The Deepest, Katowice IX 2019

$\rho(770)$

$$I^G(J^{PC}) = 1^+(1^{--})$$

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$\rho(770)$ MASS

We no longer list S-wave Breit-Wigner fits, or data with high com background.

NEUTRAL ONLY, e^+e^-

VALUE (MeV)	EVTS	DOCUMENT ID	TECN	CHG	COMMENT
775.26 ± 0.25 OUR AVERAGE					
775.02 ± 0.35		1 LEES	12G	BABR	e^+e^-
775.97 ± 0.46 ± 0.70	900k	2 AKHMETSHIN 07			e^+e^-
774.6 ± 0.4 ± 0.5	800k	3.4 ACHASOV	06	SND	e^+e^-
775.65 ± 0.64 ± 0.50	114k	5.6 AKHMETSHIN 04		CMD2	e^+e^-
775.9 ± 0.5 ± 0.5	1.98M	7 ALOISIO	03	KLOE	$1.02 \epsilon^-$
775.8 ± 0.9 ± 0.2	500k	7 ACHASOV	02	SND	$1.02 \epsilon^-$
775.9 ± 1.1		8 BARKOV	85	OLYA	e^+e^-
• • • We do not use the following data for averages, fits, limits, etc. • • •					
775.8 ± 0.5 ± 0.3	1.98M	9 ALOISIO	03	KLOE	$1.02 \epsilon^-$
775.9 ± 0.6 ± 0.5	1.98M	10 ALOISIO	03	KLOE	$1.02 \epsilon^-$
775.0 ± 0.6 ± 1.1	500k	11 ACHASOV	02	SND	$1.02 \epsilon^-$
775.1 ± 0.7 ± 5.3		12 BENAYOUN	98	RVUE	e^+e^-
770.5 ± 1.9 ± 5.1		13 GARDNER	98	RVUE	$0.28-1 \mu^+$
764.1 ± 0.7		14 O'CONNELL	97	RVUE	e^+e^-
757.5 ± 1.5		15 BERNICHA	94	RVUE	e^+e^-
768 ± 1		16 GESHKEN...	89	RVUE	e^+e^-

CHARGED ONLY, τ DECAYS and e^+e^-

VALUE (MeV)	EVTS	DOCUMENT ID	TECN	CHG	COM
775.11 ± 0.34 OUR AVERAGE					
774.6 ± 0.2 ± 0.5	5.4M	17.18 FUJIKAWA	08	BELL	$\pm$
775.5 ± 0.7		18.19 SCHAEAL	05C	ALEP	τ^-
775.5 ± 0.5 ± 0.4	1.98M	7 ALOISIO	03	KLOE	$1.02 \pi^-$
775.1 ± 1.1 ± 0.5	87k	20.21 ANDERSON	00A	CLE2	τ^-
• • • We do not use the following data for averages, fits, limits, etc. • • •					
774.8 ± 0.6 ± 0.4	1.98M	10 ALOISIO	03	KLOE	$-1.02 \pi^-$
776.3 ± 0.6 ± 0.7	1.98M	10 ALOISIO	03	KLOE	$+1.02 \pi^-$
773.9 ± 2.0 ± 0.3		22 SANZ-CILLERO3		RVUE	τ^-
774.5 ± 0.7 ± 1.5	500k	7 ACHASOV	02	SND	$\pm 1.02 \pi^-$
775.1 ± 0.5		23 PICH	01	RVUE	τ^-

MIXED CHARGES, OTHER REACTIONS

VALUE (MeV)	EVTS	DOCUMENT ID	TECN	CHG	COMMENT
763.0 ± 0.3 ± 1.2	600k	24 ABELE	99E	CBAR	$0 \pm 0.0 \bar{p} p \rightarrow \pi^+ \pi^- \pi^0$

CHARGED ONLY, HADROPRODUCED

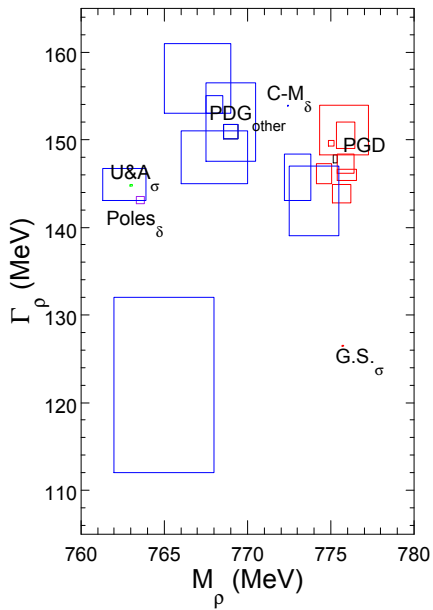
VALUE (MeV)	EVTS	DOCUMENT ID	TECN	CHG	COMMENT
766.5 ± 1.1 OUR AVERAGE					
763.7 ± 3.2		ABELE	97	CBAR	$\bar{p} n \rightarrow \pi^- \pi^0 \pi^0$
768 ± 9		AGUILAR...	91	EHS	400 $p p$
767 ± 3	2935	25 CAPRARO	87	SPEC	$-200 \pi^- \pi^0 \text{Cu} \rightarrow \pi^- \pi^0 \text{Cu}$
761 ± 5	967	25 CAPRARO	87	SPEC	$-200 \pi^- \pi^0 \text{Pb} \rightarrow \pi^- \pi^0 \text{Pb}$
771 ± 4		HUSTON	86	SPEC	$+202 \pi^+ \pi^0 \text{A} \rightarrow \pi^+ \pi^0 \text{A}$
766 ± 7	6500	26 BYERLY	73	OSPK	$-5 \pi^- p$
766.8 ± 1.5	9650	27 PISUT	68	RVUE	$-1.7-3.2 \pi^- p, t < 10$
767 ± 6	900	25 EISNER	67	HBC	$-4.2 \pi^- p, t < 10$

NEUTRAL ONLY, PHOTOPRODUCED

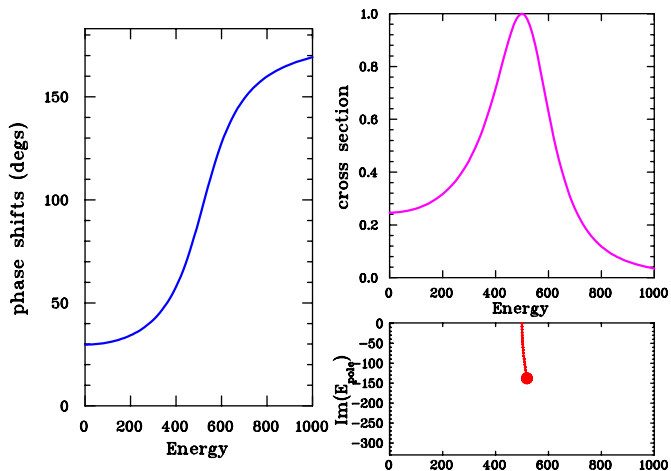
VALUE (MeV)	EVTS	DOCUMENT ID	TECN	COMMENT
769.0 ± 1.0 OUR AVERAGE				
771 ± 2 ± 2	63.5k	28 ABRAMOWICZ12	ZEUS	$e p \rightarrow e \pi^+ \pi^- p$
770 ± 2 ± 1	79k	29 BREITWEG	98B	ZEUS 50-100 γp
767.6 ± 2.7		BARTALUCCI	78	CNTR $\gamma p \rightarrow e^+ e^- p$
775 ± 5		GLADDING	73	CNTR 2.9-4.7 γp
767 ± 4	1930	BALLAM	72	HBC 2.8 γp
770 ± 4	2430	BALLAM	72	HBC 4.7 γp
765 ± 10		ALVENSLEB...	70	CNTR $\gamma A, t < 0.01$
767.7 ± 1.9	140k	BIGGS	70	CNTR $< 4.1 \gamma C \rightarrow \pi^+ \pi^- C$
765 ± 5	4000	ASBURY	67B	CNTR $\gamma + \text{Pb}$
• • • We do not use the following data for averages, fits, limits, etc. • • •				
771 ± 2	79k	30 BREITWEG	98B	ZEUS 50-100 γp

NEUTRAL ONLY, OTHER REACTIONS

VALUE (MeV)	EVTS	DOCUMENT ID	TECN	CHG	COMMENT
769.0 ± 0.9 OUR AVERAGE					
Error includes scale factor of 1.4. See the ideogram below.					
765 ± 6		BERTIN	97C	OBLX	$0.0 \bar{p} p \rightarrow \pi^+ \pi^- \pi^0$
773 ± 1.6		WEIDENAUER	93	ASTE	$\bar{p} p \rightarrow \pi^+ \pi^- \omega$
762.6 ± 2.6		AGUILAR...	91	EHS	400 $p p$
770 ± 2		31 HEYN	81	RVUE	Pion form factor
768 ± 4	32.33	BOHACIK	80	RVUE	0
769 ± 4		26 MARGALIT	80	ASPK	0

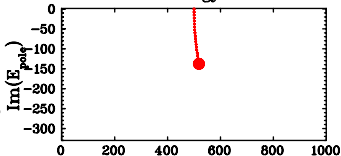
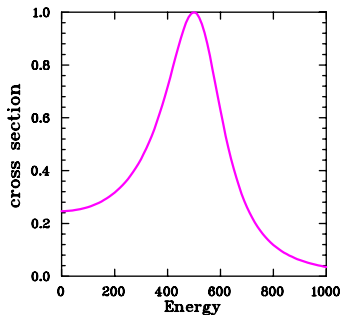
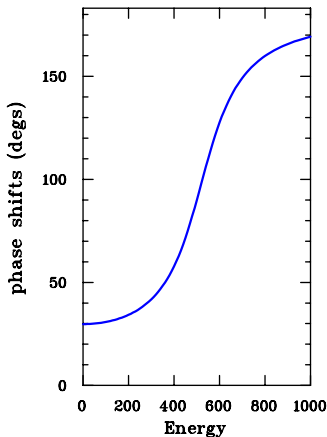


Resonances



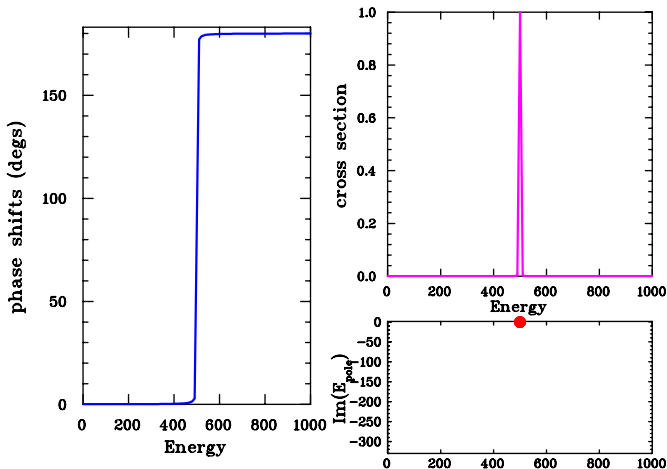
Resonances: Breit-Wigner (BW) approximation:

$$\text{Ampl} = \frac{C}{M_{BW} - E - i\Gamma_{BW}/2}, \quad \delta = \text{ArcTan}\left(\frac{\Gamma/2}{M_{BW} - E}\right), \quad \sigma = \frac{C^2}{(M_{BW} - E)^2 + \Gamma^2/4}$$



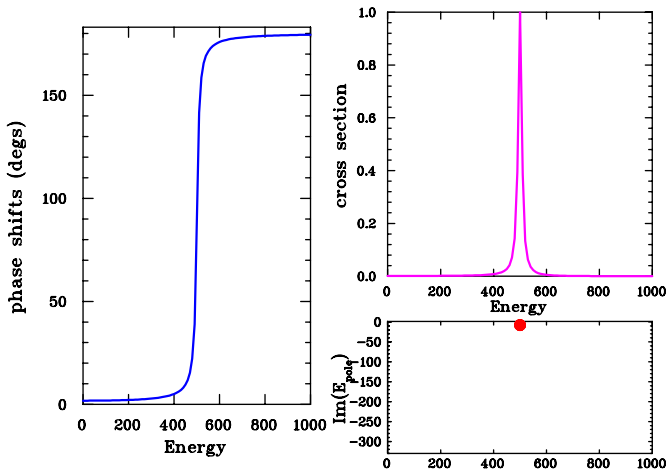
Resonances: exercise - Breit Wigner parameterization

$$\text{Ampl} = \frac{C}{M_{BW} - E - i\Gamma_{BW}/2}, \quad \delta = \text{ArcTan}\left(\frac{\Gamma/2}{M_{BW} - E}\right), \quad \sigma = \frac{C^2}{(M_{BW} - E)^2 + \Gamma^2/4}$$



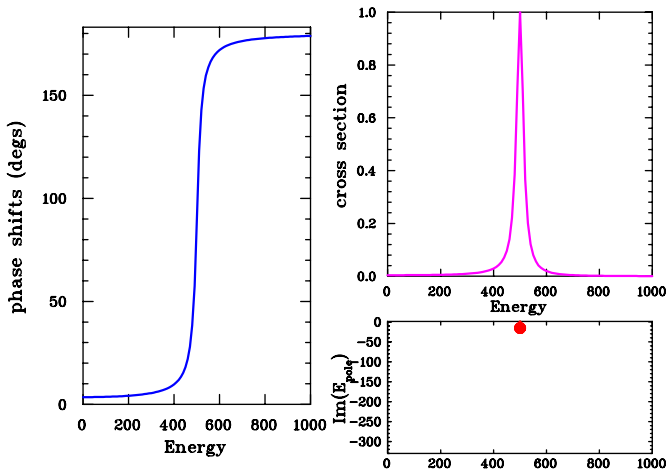
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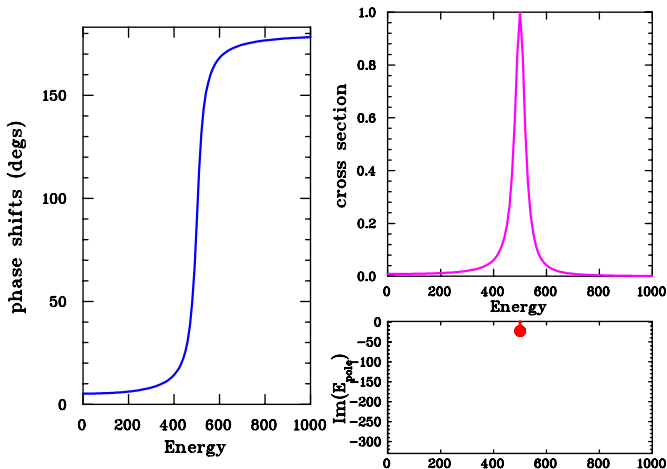
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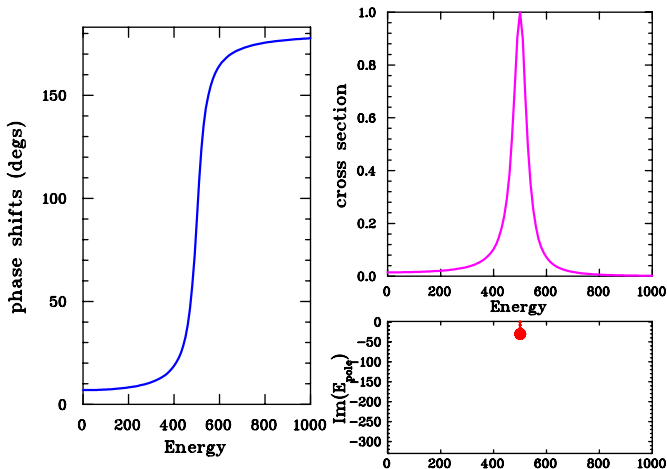
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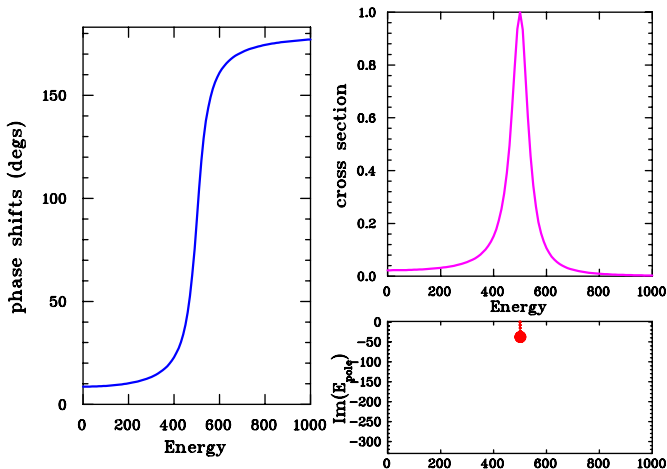
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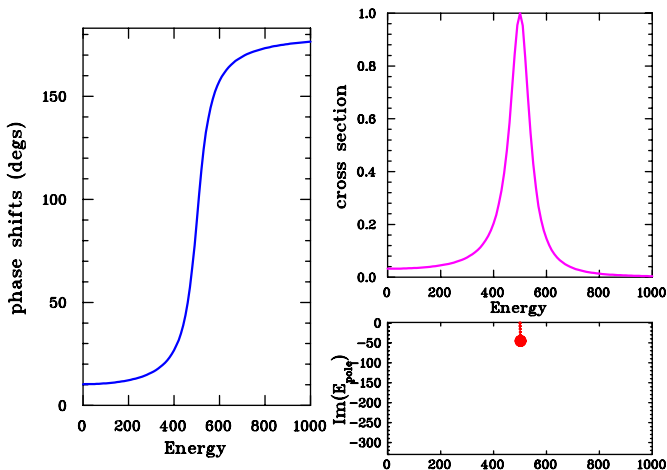
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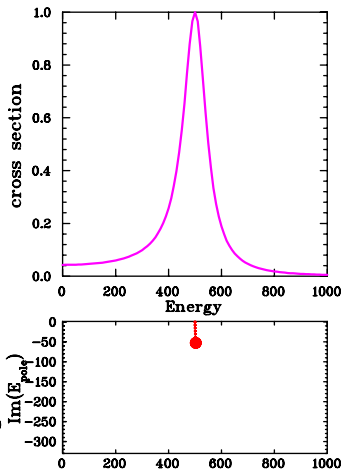
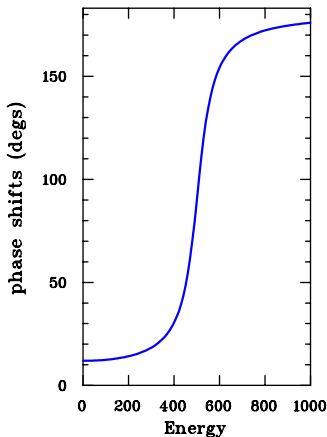
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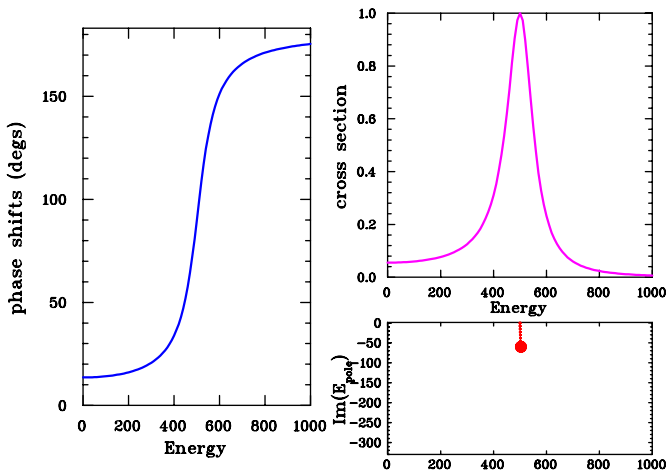
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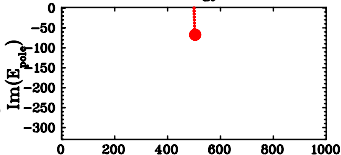
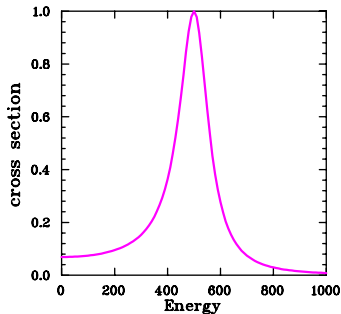
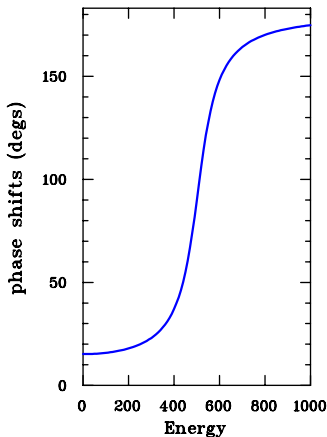
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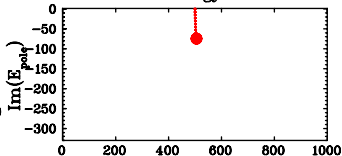
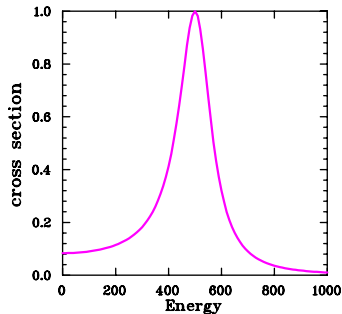
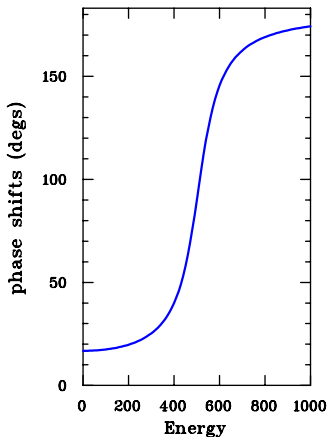
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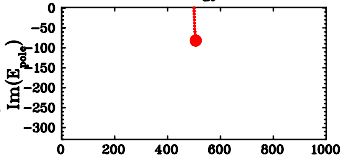
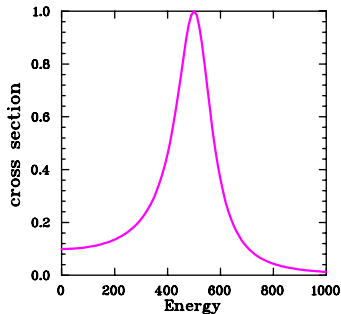
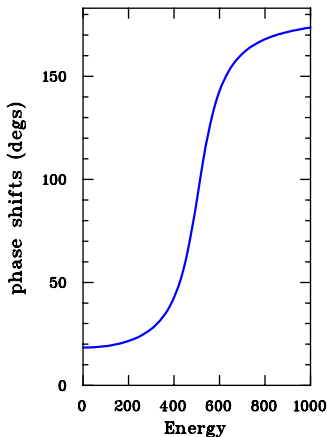
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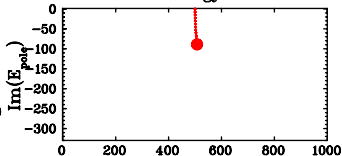
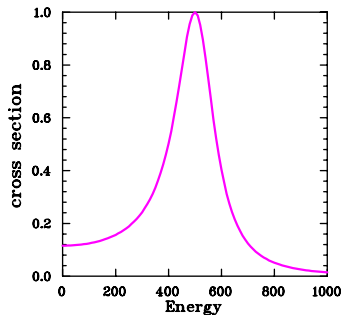
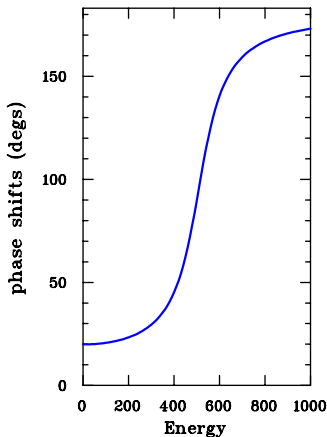
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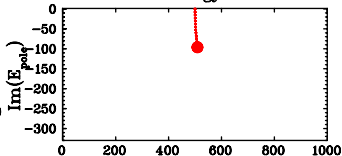
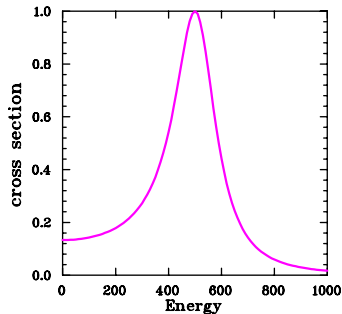
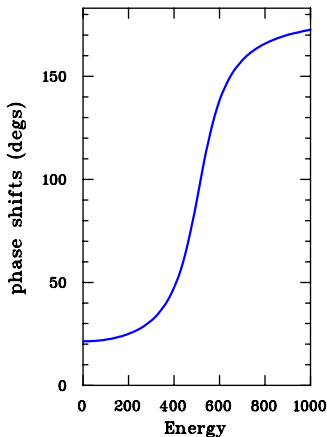
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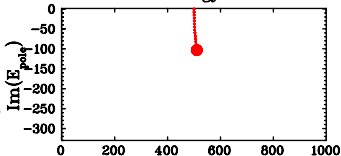
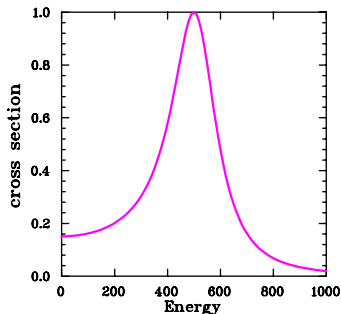
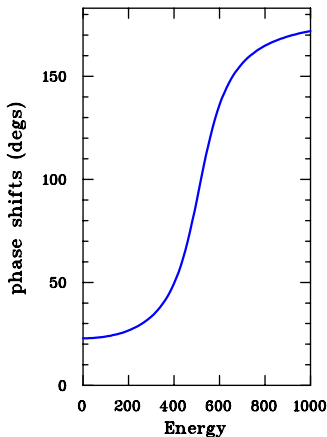
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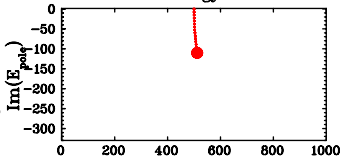
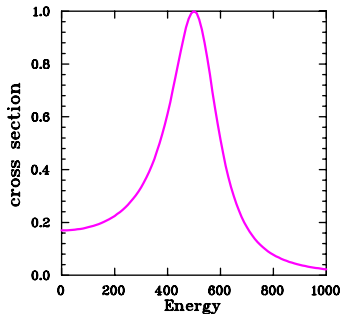
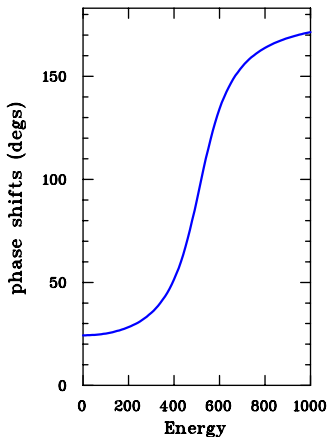
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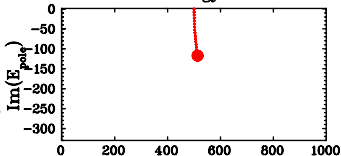
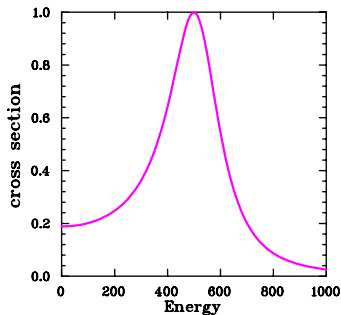
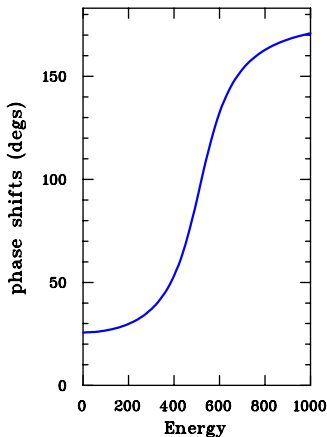
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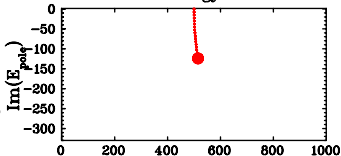
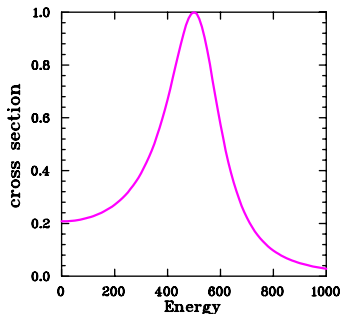
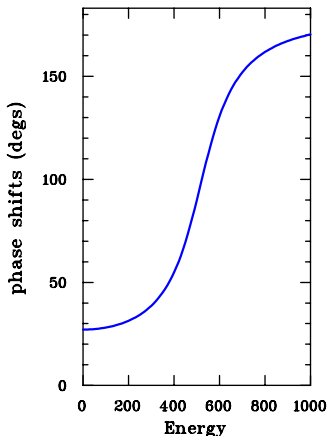
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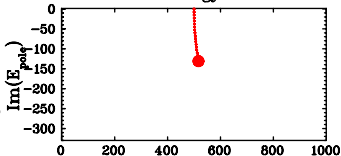
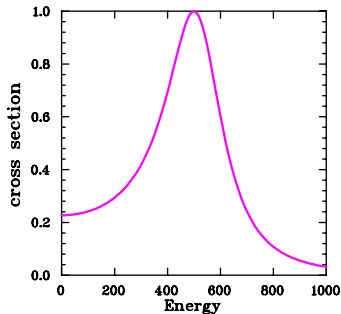
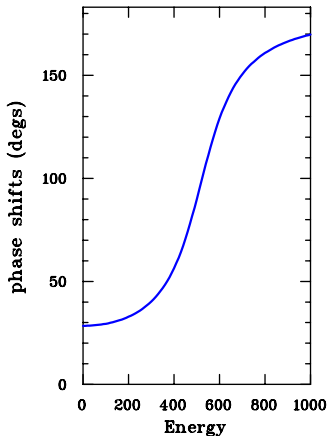
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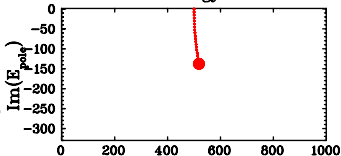
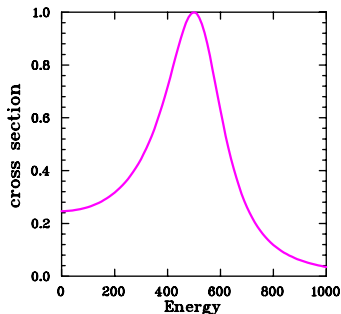
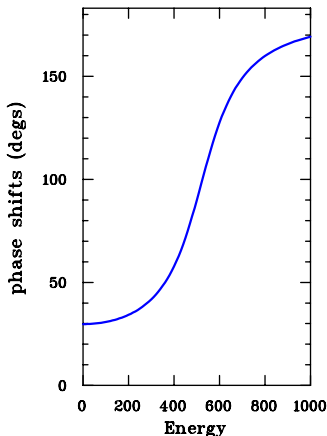
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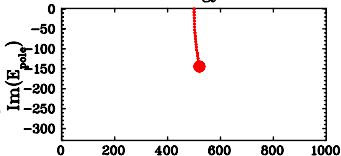
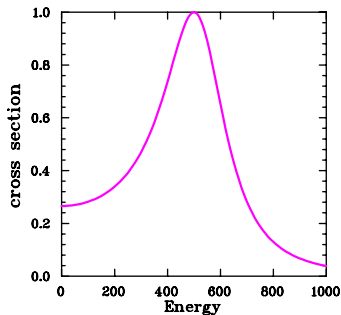
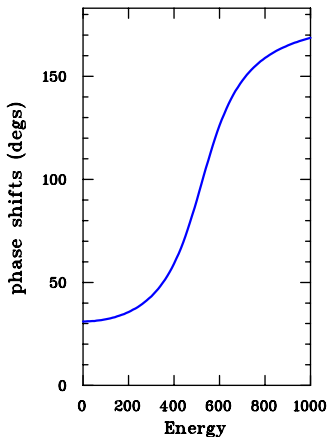
Resonances: exercise - Breit Wigner parameterization

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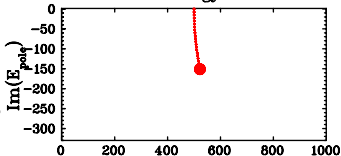
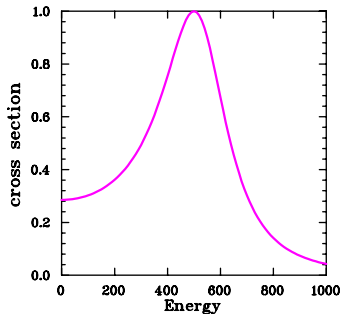
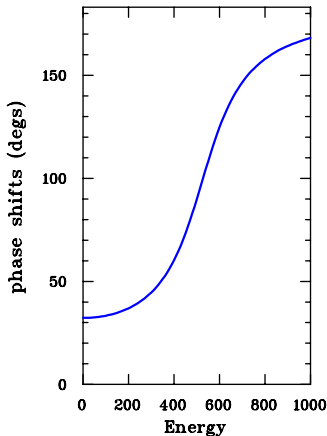
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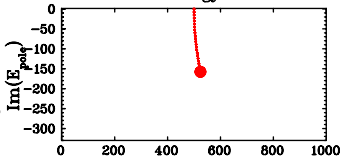
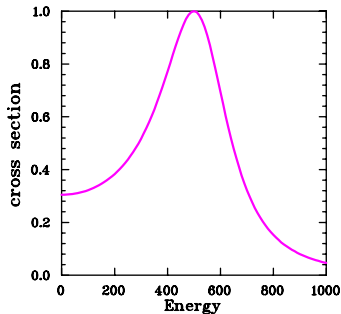
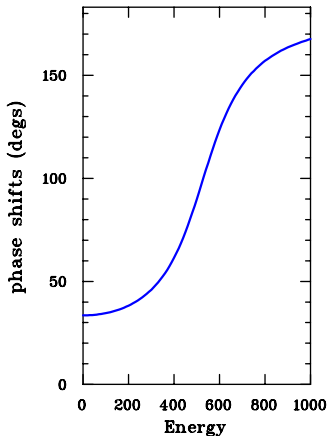
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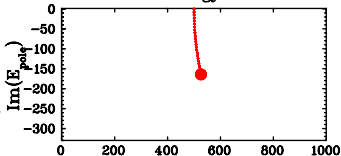
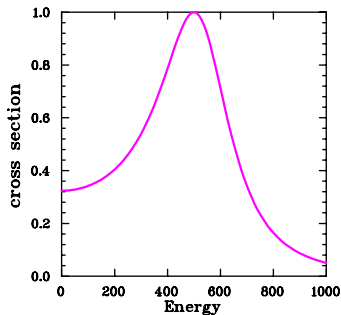
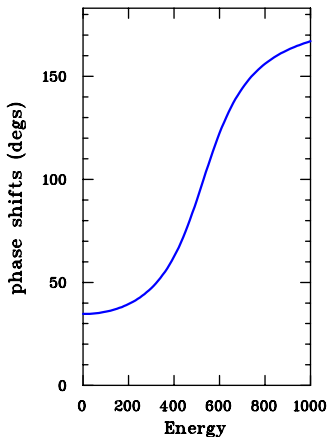
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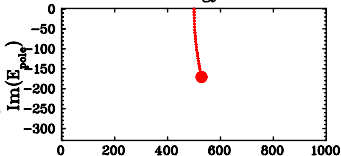
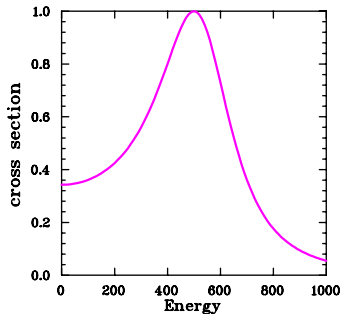
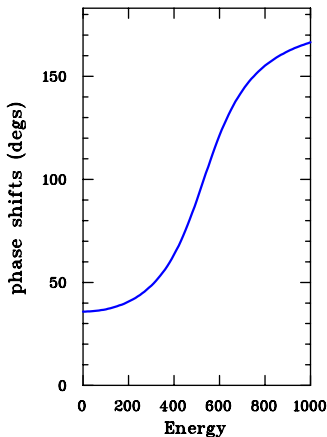
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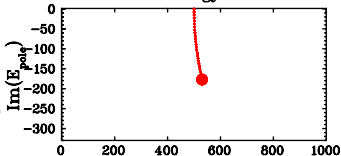
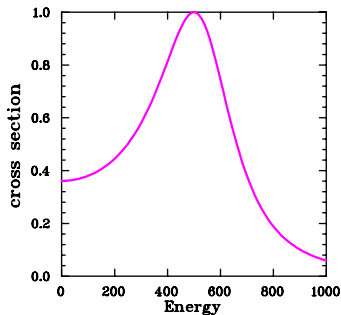
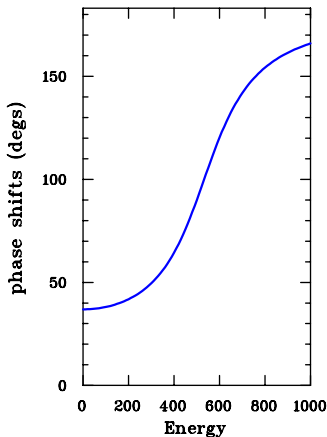
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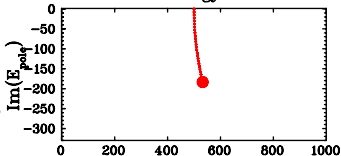
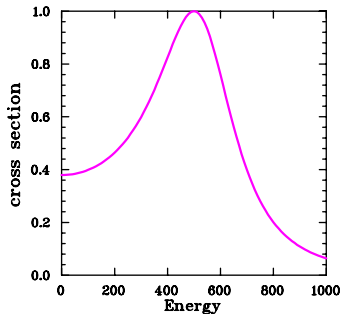
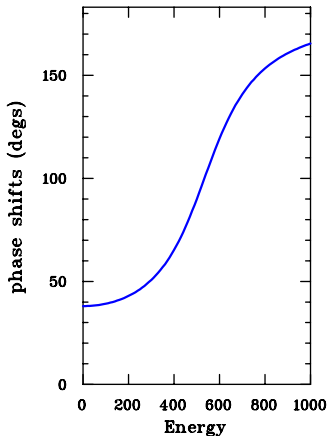
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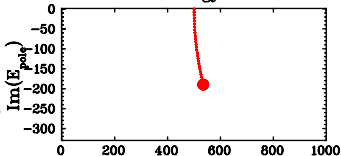
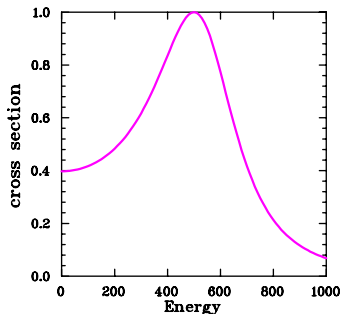
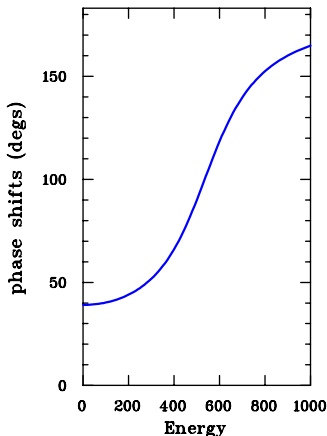
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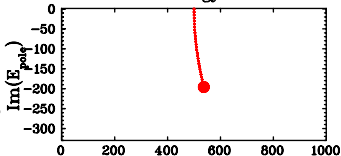
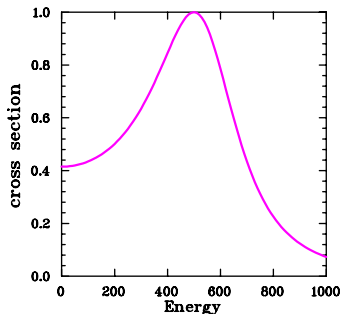
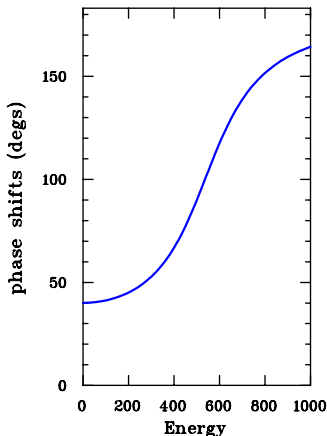
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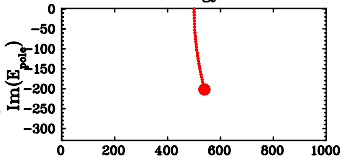
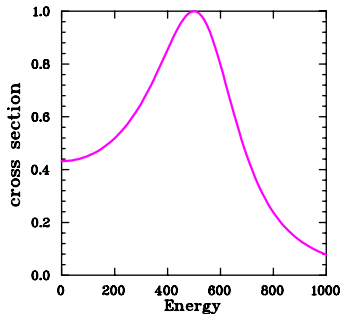
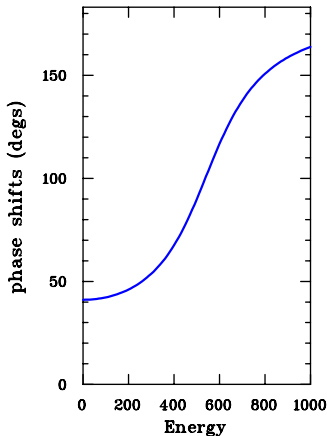
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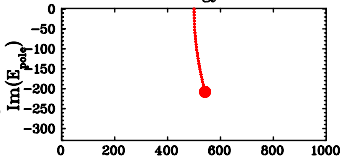
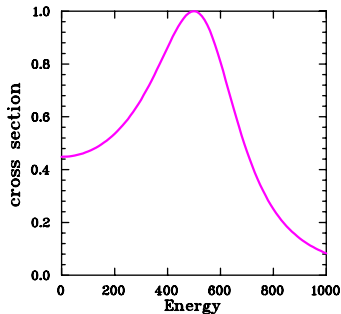
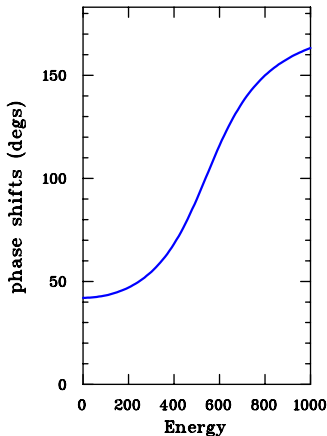
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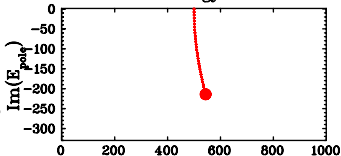
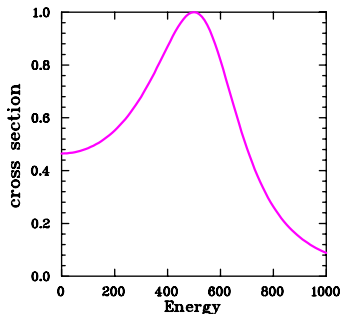
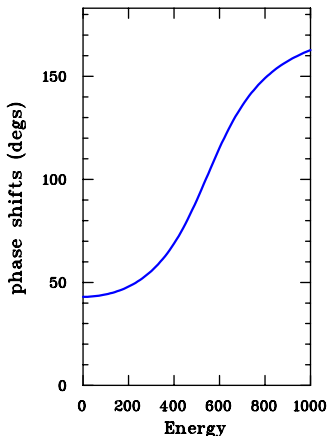
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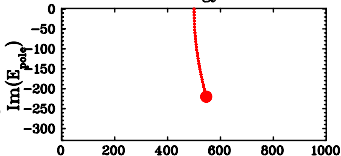
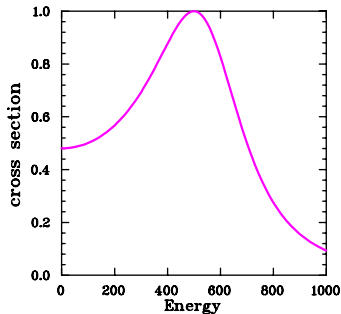
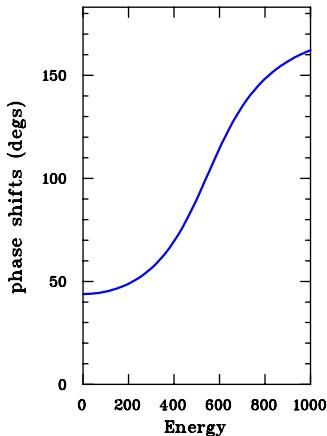
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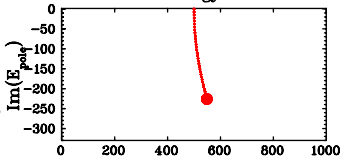
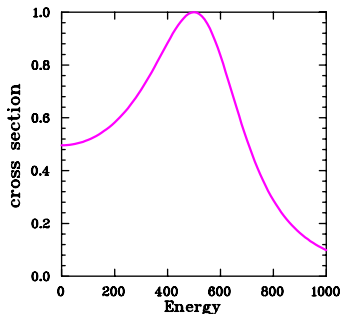
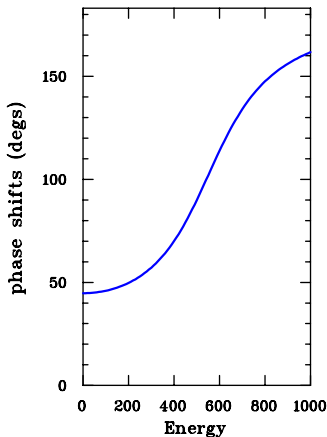
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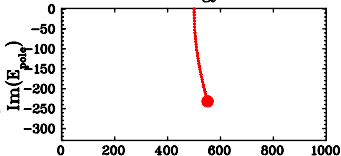
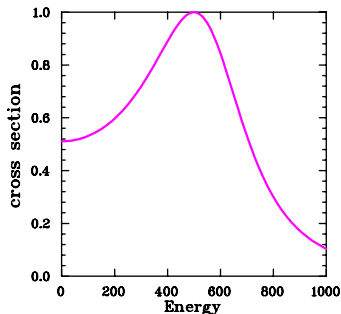
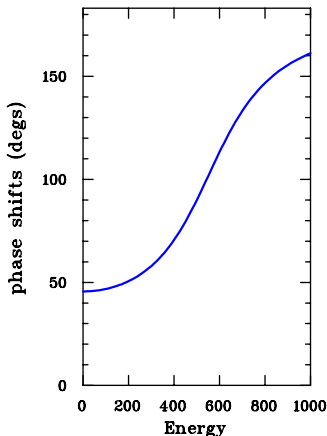
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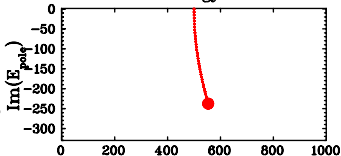
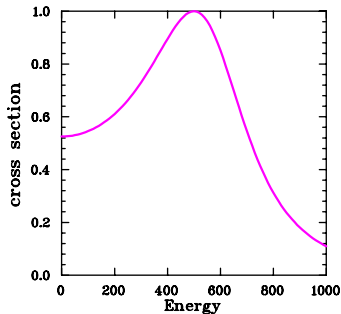
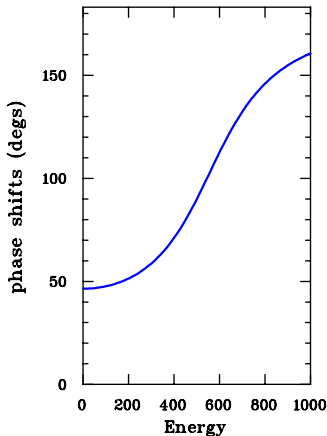
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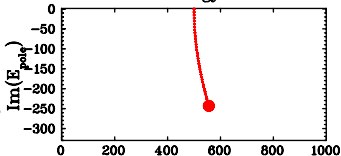
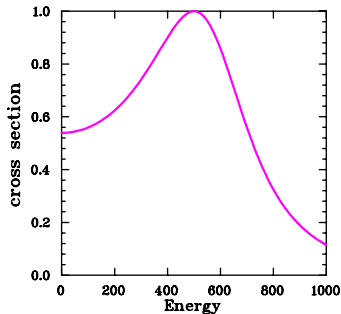
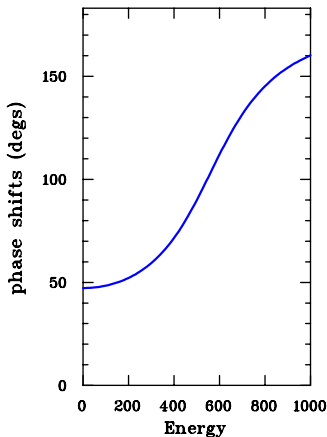
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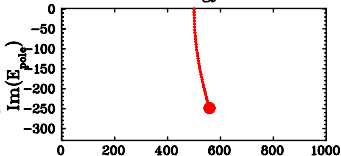
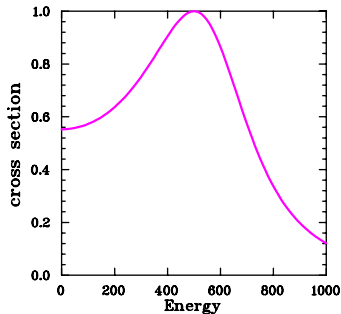
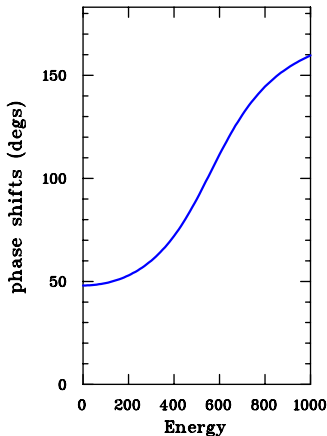
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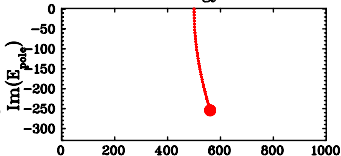
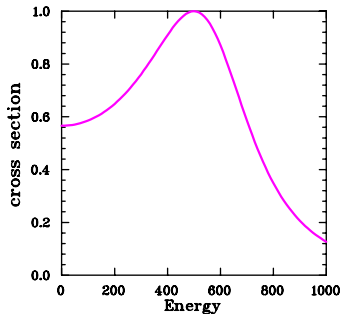
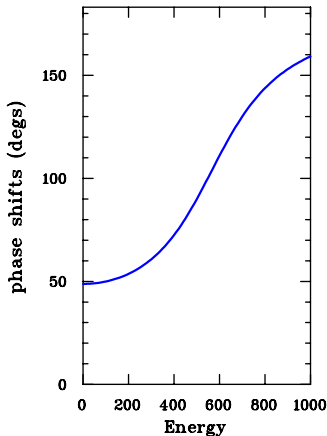
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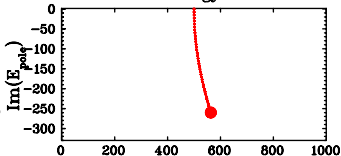
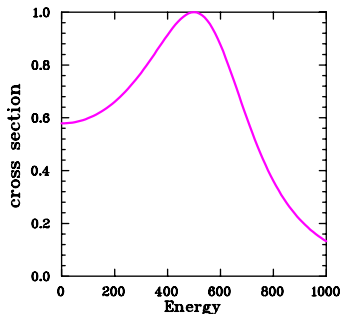
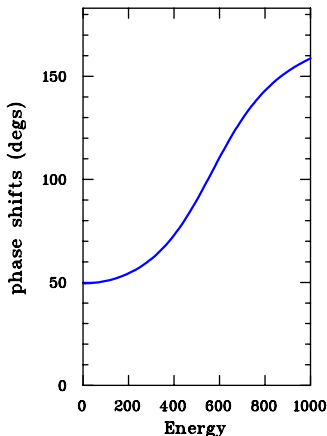
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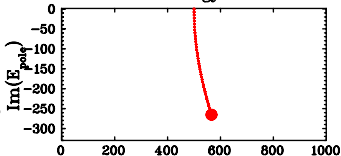
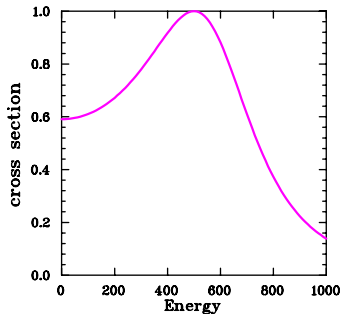
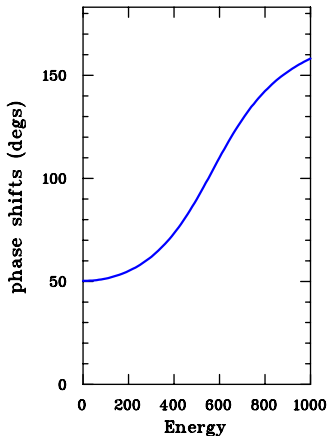
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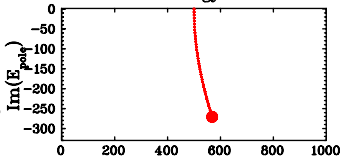
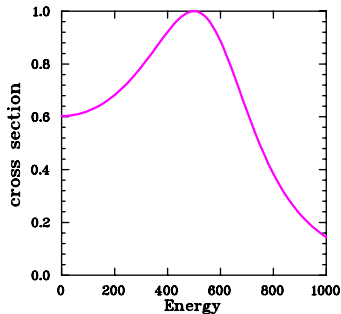
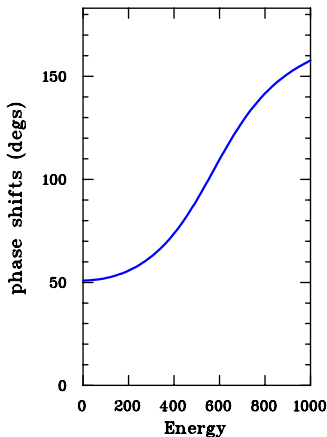
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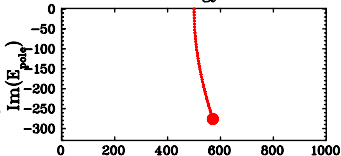
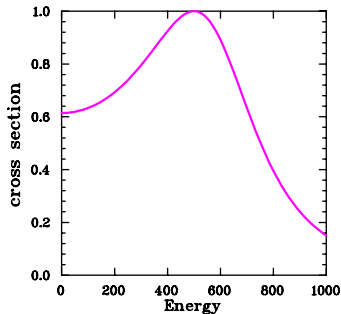
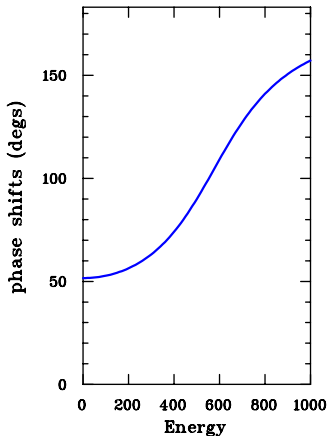
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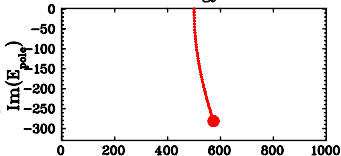
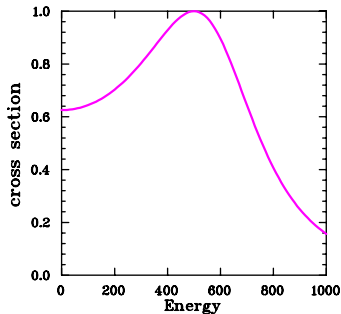
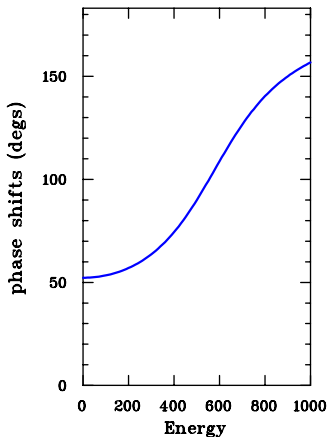
Resonances: exercise - Breit Wigner parameterization

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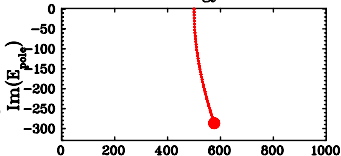
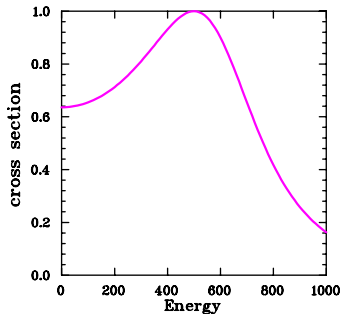
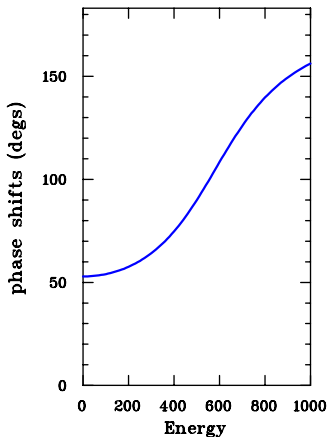
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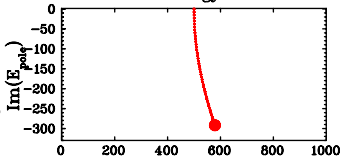
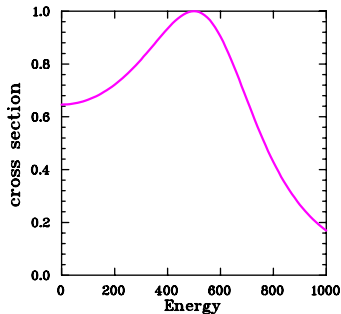
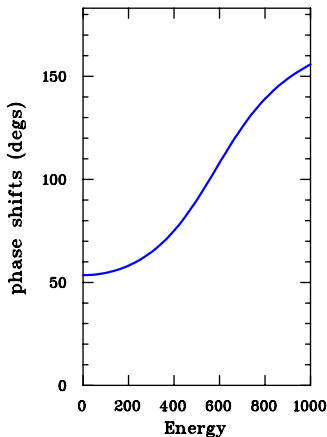
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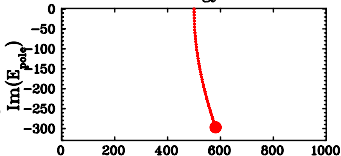
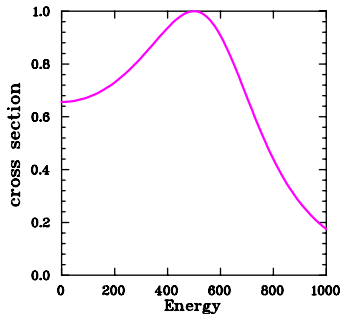
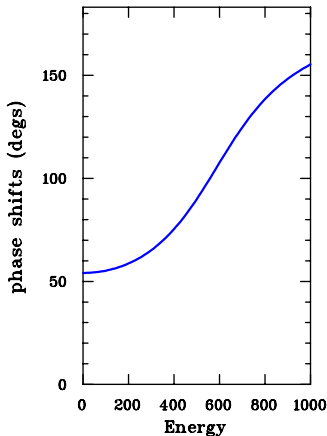
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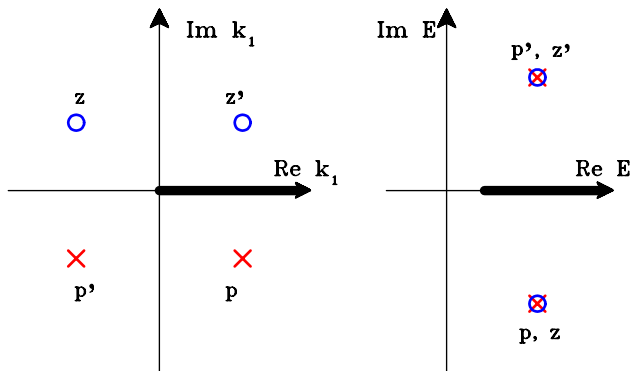
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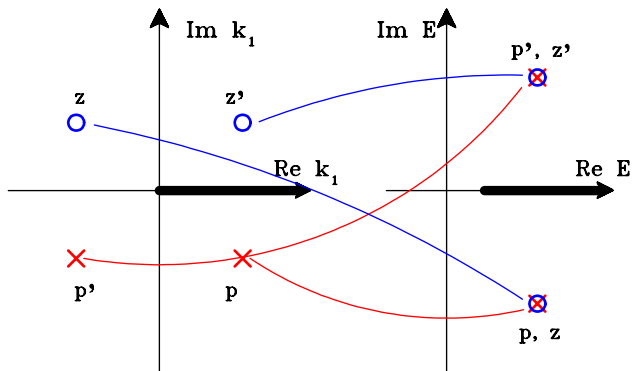
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One channel scattering

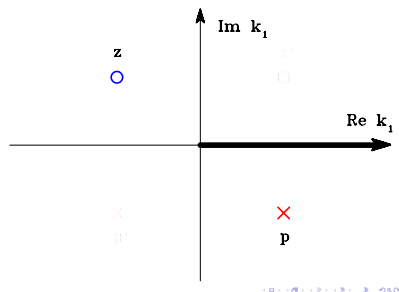
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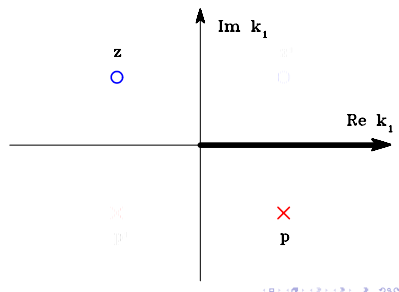
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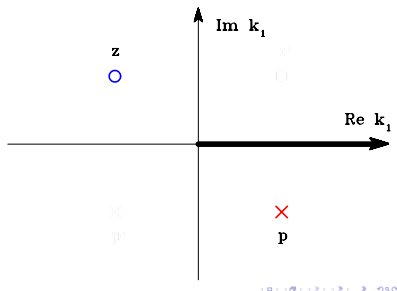
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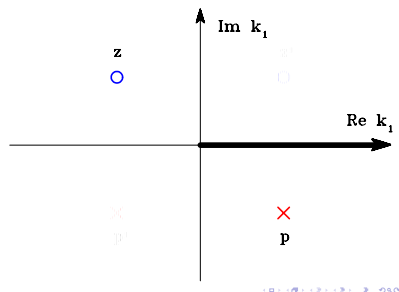
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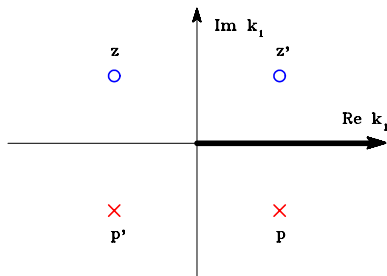
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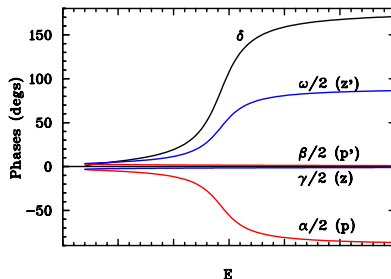
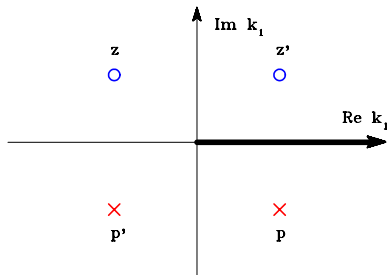
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$$\text{angle} = \text{ArcTan}\left(\frac{-\text{Im}k_j}{k - \text{Re}k_j}\right)$$



Breit Wigner approximation

- ▶ $BW(E) = \frac{\Gamma/2k}{M_{BW}-E-i\Gamma/2}$
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Let's check it for $\rho(770)$:

$M_{BW} = 775.26 \pm 0.25$ MeV (PDG'2016),

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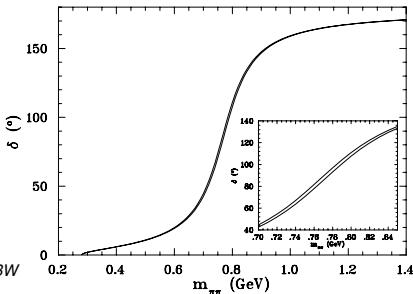
$2\sqrt{a^2 + m^2} < M_{BW}$ by ≈ 9 MeV !!!

Left (upper) line:

A_{BW} fitted to the data

Right (lower) line:

A_S fitted to the data with $2\sqrt{a^2 + m^2} = M_{BW}$



More resonances (but still one channel)

Adding resonances (for simplicity two resonances, both with $S = e^{2i\delta}$):

- **Isobar model:** adding amplitudes (even unitary ones) violates unitarity:

$$T_{1,2} = T_1 + T_2 = \frac{S_1 - 1}{2ik} + \frac{S_2 - 1}{2ik} \rightarrow S_1 + S_2 = e^{2i\delta_1} + e^{2i\delta_2}$$

of course $|S_1 + S_2| \neq 1$,

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For example $S_{1,2} = \frac{(-k - k_1)(-k + k_1^*)(-k - k_2)(-k + k_2^*)}{(k - k_1)(k + k_1^*)(k - k_2)(k + k_2^*)}$

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$$\text{Of course } T_{1,2} = \frac{S_{1,2} - 1}{2ik}$$

- **Sum of K matrices:** $S = 1 + 2iT = (1 + iK)/(1 - iK)$ does not violate unitarity, for example $T_{1,2} = \frac{1}{k} \frac{K_1 + K_2}{1 - iK_1 - iK_2}$

More channels: $k_2 = \pm \sqrt{k_1^2 + m_1^2 - m_2^2}$

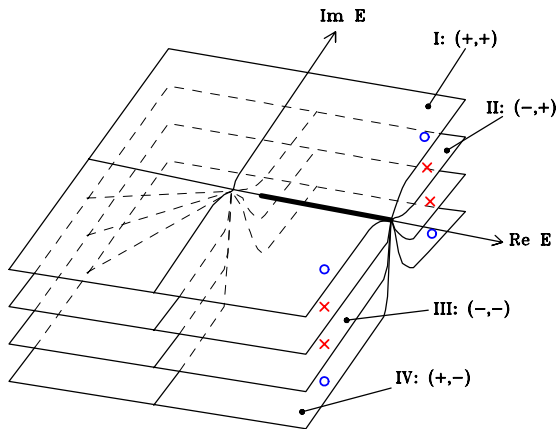
$(\text{Im}(k_1), \text{Im}(k_2))$: $(+,+)$, $(-,+)$

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Multiplication and displacement of S matrix poles

- Multiplication:

1 pole $\longrightarrow 2^{n-1}$ poles due to $(\pm k)^2$ ambiguity and 2^n Riemann sheets

- Displacement:

$$S_{11} = \frac{D_1(-k_1)D_2(k_2)}{D_1(k_1)D_2(k_2)} \text{ in decoupled case}$$

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$$S = \frac{D_1(-k_1)D_2(k_2) + C(-k_1, k_2)}{D_1(k_1)D_2(k_2) + C(k_1, k_2)} \text{ in coupled case}$$

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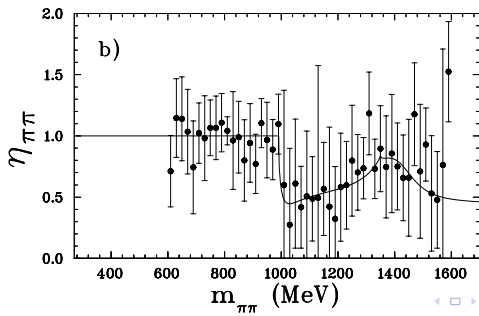
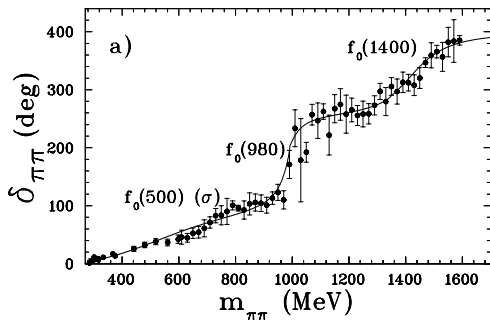
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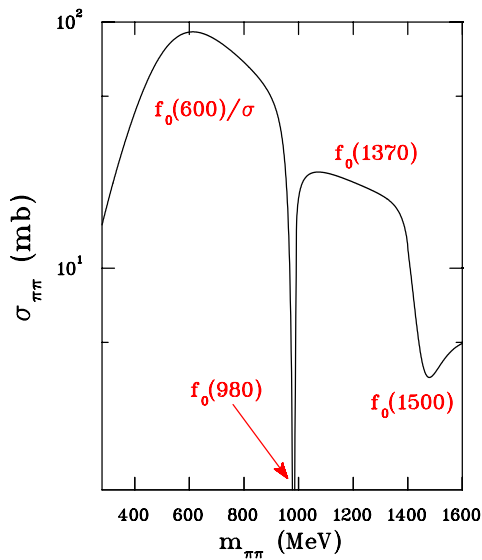
$$S = \begin{pmatrix} \eta e^{2i\delta_1} & i\sqrt{1-\eta^2} e^{i(\delta_1+\delta_2)} \\ i\sqrt{1-\eta^2} e^{i(\delta_1+\delta_2)} & \eta e^{2i\delta_2} \end{pmatrix} = \begin{pmatrix} \frac{D(-k_1, k_2)}{D(k_1, k_2)} & S_{12} \\ S_{21} & \frac{D(k_1, -k_2)}{D(k_1, k_2)} \end{pmatrix}$$

$$\text{where } S_{12}^2 = S_{21}^2 = S_{11}S_{22} - \frac{D(-k_1, -k_2)}{D(k_1, k_2)}$$

Example for two channels: $JI = S0$ wave



Puzzling (J/ψ) S_0 wave $\pi\pi$ cross section



Example for two channels: $JI = S0$ wave

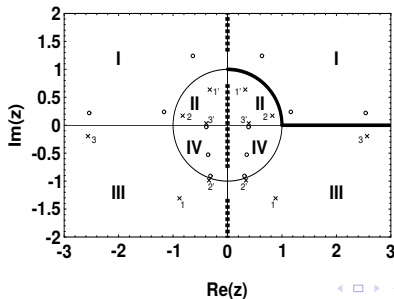
Pole	ReE_{pole} MeV	ImE_{pole} MeV	R. sheet
1	639.6	-323.9	$(-, -) : III$
1'	511.4	-230.6	$(-, +) : II$
2	982.0	-36.9	$(-, +) : II$
2'	432.4	-8.4	$(-, -) : III$
3	1431.7	-79.3	$(-, -) : III$
3'	1394.9	-120.6	$(-, +) : II$

Example for two channels: $J I = S 0$ wave

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$$z = \frac{k_1 + k_2}{\sqrt{m_K^2 - m_\pi^2}}$$

Rysunek 16: Położenie biegunów (krzyże) i zer (kółka) elementu macierzyowego S_{11} macierzy rozpraszania dla dopasowania do zestawu D_{CKM} A. Gruba linia ciągła oznacza obszar fizyczny rozpraszania w kanałach sprzężonych $\pi\pi$ i $K\bar{K}$. Grubą linią przerywaną przedstawione jest położenie cięć funkcji Josta. Cienką linią zaznaczony jest okrąg $|z| = 1$. Numeracja poszczególnych płatów i biegunów została wyjaśniona w tekście.

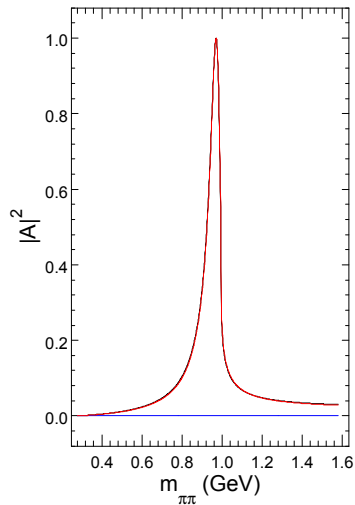
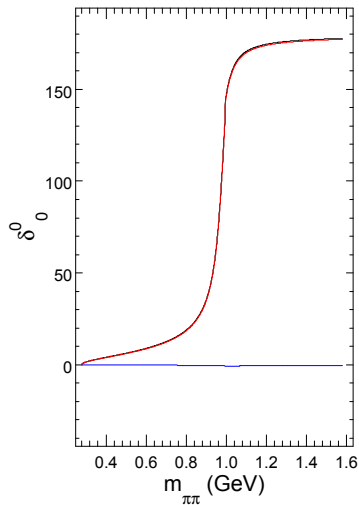


For two channels" TWO POLES 2 and 2' ($f_0(980)$)

2: 982.0-i*36.9

2': 432.4-i*8.4

both 2 and 2'

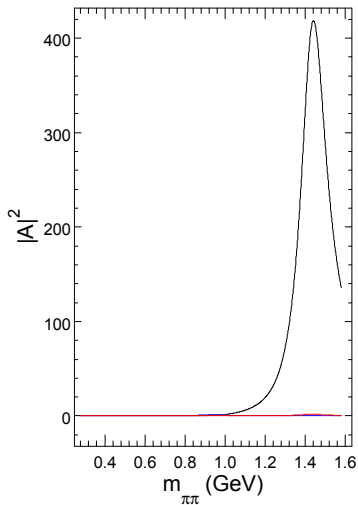
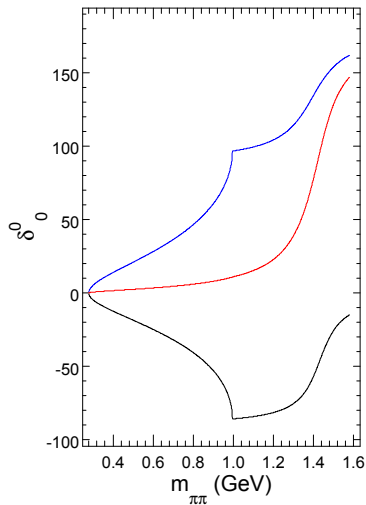


For two channels" TWO POLES 3 and 3' ($f_0(1400)$)

3: 1431.7-i*79.3

3': 1394.9-i*120.6

both 3 and 3'

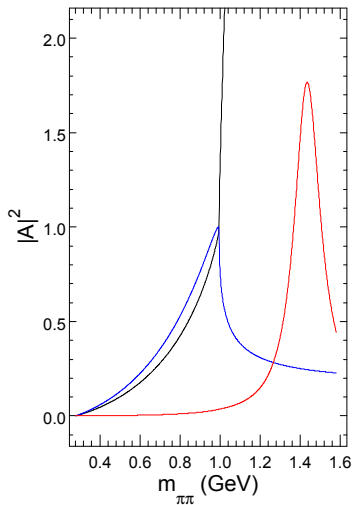
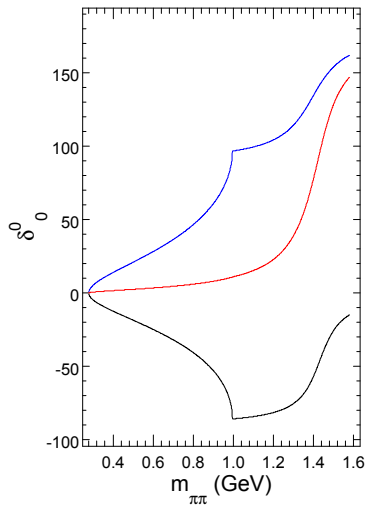


For two channels" TWO POLES 3 and 3' ($f_0(1400)$)

3: 1431.7-i*79.3

3': 1394.9-i*120.6

both 3 and 3'

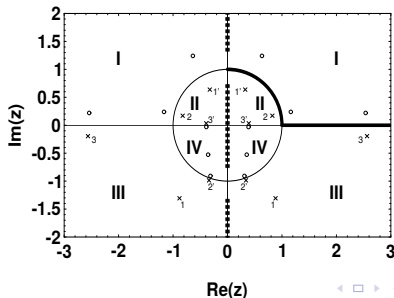


Example for two channels: $J0 = S0$ wave

Pole	$Re E_{pole}$ MeV	$Im E_{pole}$ MeV	R. sheet
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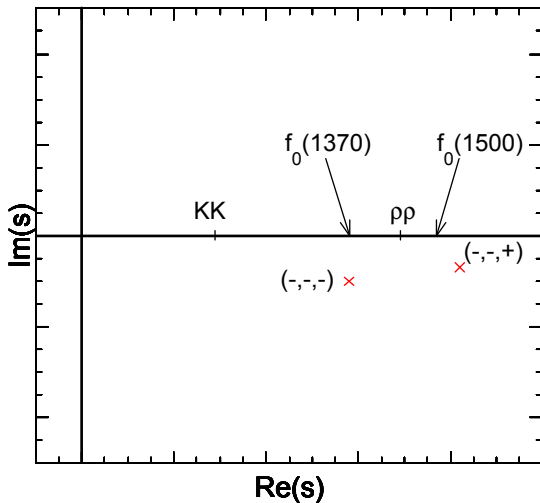
More than 2 channels

- ▶ more poles and more Riemann sheets (2^n)
- ▶ no similar "z" variable

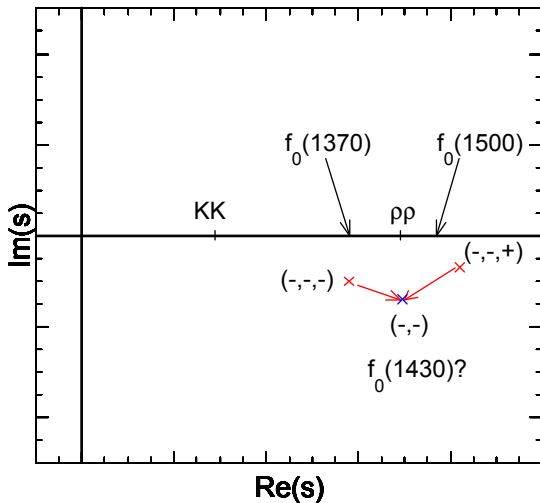
2^n Riemann sheets for n channels

channel	$C = 0$		$C = 1$		sign Imk_π, Imk_K, Imk_3	sheet
	ReE	ImE	ReE	ImE		
$\pi\pi$	658	-607	564	-279	$-, -, -$	VI
			518	-261	$-, +, +$	II
			211	0	$-, +, -$	VII
			532	-315	$-, -, +$	III
			235	0	$+, +, -$	VIII
$\pi\pi$	1346	-275	1405	-74	$-, -, -$	VI
			1445	-116	$-, +, +$	II
			1424	-94	$-, +, -$	VII
			1456	-47	$-, -, +$	III
$K\bar{K}$	881	-498	170	0	$+, -, -$	V
			159	0	$-, -, -$	VI
			418	-10	$-, -, +$	III
			1038	-204	$-, +, -$	VII
			988	-31	$-, +, +$	II
$\sigma\sigma$	118	-2227	4741	-4688	$-, -, -$	VI
			3687	-2875	$-, +, -$	VII
			3626	-3456	$+, -, -$	V
			3533	-579	$+, +, -$	VIII

$f_0(1370)$ and $f_0(1500)$: positions of poles, $C = 1$



$f_0(1370)$ and $f_0(1500)$: positions of poles, $C = 0$



2^n Riemann sheets for n channels

channel	$C = 0$		$C = 1$		sign Imk_π, Imk_K, Imk_3	sheet
	ReE	ImE	ReE	ImE		
$\pi\pi$	658	-607	564	-279	$-, -, -$	VI
			518	-261	$-, +, +$	II
			211	0	$-, +, -$	VII
			532	-315	$-, -, +$	III
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2^n Riemann sheets for n channels

channel	$C = 0$		$C = 1$		sign $\text{Im}k_\pi, \text{Im}k_K, \text{Im}k_3$	sheet
	$\text{Re}E$	$\text{Im}E$	$\text{Re}E$	$\text{Im}E$		
$\pi\pi$	658	-607	564	-279	$-, -, -$	VI
			518	-261	$-, +, +$	II
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$\leftarrow f_0(500)$
 $\leftarrow f_0(1370)?$
 $\leftarrow f_0(1500)?$
 $\leftarrow f_0(980)$

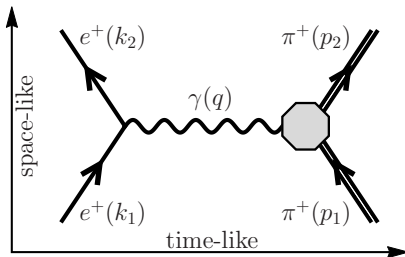
Pion electromagnetic form factor in the P wave

What are the correct $\rho(770)$ meson mass and width values?

PRD 96, 113004 (2017)

Erik Bartoš, Stanislav Dubnička, Andrej Liptaj Anna Zuzana Dubničkova and RK

$$\langle \pi^+(p_2) | J_\pi^\mu(0) | \pi^-(p_1) \rangle = e (p_1 + p_2)^\mu F_\pi(q^2)$$



Gounaris-Sakurai pion EM FF and $\rho(770)$

Gounaris-Sakurai pion electromagnetic form factor at the elastic region.

What are the correct $\rho(770)$ meson mass and width values?

Erik Bartoš, Stanislav Dubnička, Andrej Liptaj Anna Zuzana Dubničková and RK
PRD 96, 113004 (2017)

$$\sigma_{tot}(e^+e^- \rightarrow \pi^+\pi^-) = \frac{\pi\alpha^2(0)}{3s} \beta_\pi^3(s) \left| F_\pi^{EM,l=1}(s) + Re^{i\phi} \frac{m_\omega^2}{m_\omega^2 - s - im_\omega\Gamma_\omega} \right|^2$$

where pion "velocity" $\beta_\pi(s) = \sqrt{\frac{s-4m_\pi^2}{s}}$, R - amplitude for $\rho - \omega$ interference (free parameter), phase $\phi = \text{ArcTan} \frac{m_\rho\Gamma_\rho}{m_\rho^2 - m_\omega^2}$ is the $\delta_1^1 = \delta_\rho$ fixed at $s = m_\omega^2$

Fit to data for $\sigma_{tot}(e^+e^- \rightarrow \pi^+\pi^-)$

- M. Ablikin et al. (BESIII Collaboration), Phys. Lett. B 753, 629 (2016).
- J. P. Lees et al. (BABAR Collaboration), Phys. Rev. D 86, 032013 (2012).

Gounaris-Sakurai pion electromagnetic form factor

G. J. Gounaris and J. J. Sakurai, *Phys. Rev. Lett.* 21, 244 (1968)

Assumption: $\frac{q^3}{\sqrt{s}} \text{Cotg}\delta_1^1(s) = a + bq^2 + h(s)q^2$ where $h(s) = \frac{2q}{\pi\sqrt{s}} \text{Log}\left(\frac{\sqrt{s}+2q}{2m_\pi}\right)$

Then $F_\pi^{GS}(s) = \frac{\sqrt{s}}{q^3} \frac{1}{\text{Cotg}\delta_1^1(s)-i}$ - no direct dependence on $\rho(770)$ parameters,
however taking into account two conditions:

► $\text{Cotg}\delta_1^1(s) \Big|_{s=m_\rho^2} = 0$ and

► $F_\pi^{BW}(s) = \frac{m_\rho^2}{m_\rho^2 - s - im_\rho\Gamma_\rho} \rightarrow \delta_1^1(s) = \text{ArcTan} \frac{m_\rho\Gamma_\rho}{m_\rho^2 - s} \rightarrow \frac{d\delta_1^1(s)}{ds} \Big|_{s=m_\rho^2} = \frac{1}{m_\rho\Gamma_\rho}$

$$a = \frac{4q_\rho^5}{m_\rho^2\Gamma_\rho} + 4q_\rho^4 h'(m_\rho^2)$$

$$b = -\frac{4q_\rho^3}{m_\rho^2\Gamma_\rho} - 4q_\rho^4 h'(m_\rho^2) - h(m_\rho^2)$$

$$F_\pi^{GS}(s) = \frac{m_\rho^2 + m_\rho\Gamma_\rho \left[\frac{3m_\rho^2}{\pi q_\rho^2} \log\left(\frac{m_\rho + 2mq_\rho}{2m_\pi}\right) + \frac{m_\rho}{2\pi q_\rho} - \frac{m_\rho^2}{\pi q_\rho^3} \right]}{(m_\rho^2 - s) + \Gamma_\rho \frac{m_\rho^2}{q_\rho^3} \left[q^2 (h(s) - h(m_\rho^2)) + q_\rho^2 h'(m_\rho^2)(m_\rho^2 - s) \right] - im_\rho\Gamma_\rho \left(\frac{q}{q_\rho} \right)^3 \frac{m_\rho}{\sqrt{s}}}$$

Fit to unified BESIII-BABAR data at the elastic region using G-S model

- ▶ $\chi^2 = 40.6$ pdf
- ▶ $m_\rho = (775.73 \pm 0.10)$ MeV
- ▶ $\Gamma_\rho = (126.51 \pm 0.13)$ MeV
- ▶ PDG'2017:

$$m_\rho = 775.26 \pm 0.25 \text{ MeV (from } e^+e^- \rightarrow \pi^+\pi^-)$$

$$m_\rho = 769.0 \pm 0.9 \text{ MeV (other)}$$

$$\Gamma_\rho = 147.8 \pm 0.9 \text{ MeV (from } e^+e^- \rightarrow \pi^+\pi^-)$$

$$\Gamma_\rho = 150.9 \pm 1.7 \text{ MeV (other)}$$

in Ref. [6].

Now, the $\rho^0(770)$ meson parameters will be determined by an application of the U&A model of the pion EM FF to an optimal description of the same data on the total cross section of the $e^+e^- \rightarrow \pi^+\pi^-$ process.

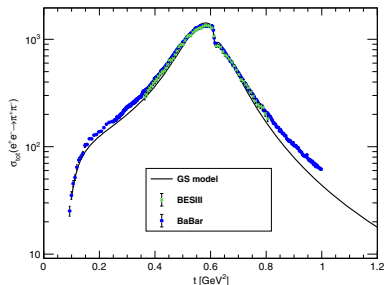


FIG. 2. Optimal description of the unified BESIII-BABAR data on $\sigma_{\text{tot}}(e^+e^- \rightarrow \pi^+\pi^-)$ at the elastic region by the pion EM FF G.-S. model.

Fit to unified BESIII-BABAR data at the elastic region using U&A model

- ▶ $F_{\pi}^{VDM}(s) = \frac{m_{\rho}^2 f_{\rho\pi\pi} / f_{\rho}}{m_{\rho}^2 - s}$
- ▶ $F_{\pi}^{VDM}(q) = \frac{(i-q_{\rho})(i+q_{\rho})}{(q-q_{\rho})(q+q_{\rho})} f_{\rho\pi\pi} / f_{\rho}$
- ▶ $F_{\pi}^{VDM}(q) = \frac{(q-q_Z)(i-q_P)(i-q_{\rho})(i+q_{\rho})}{(q-q_P)(i-q_Z)(q-q_{\rho})(q+q_{\rho})} f_{\rho\pi\pi} / f_{\rho}$
- ▶ $\chi^2 = 1.54 \text{ pdf}$
- ▶ $m_{\rho} = (763.03 \pm 0.14) \text{ MeV}$
- ▶ $\Gamma_{\rho} = (144.8 \pm 0.23) \text{ MeV}$

The optimal description of the recent data [1,2] on the total cross section of the $e^+e^- \rightarrow \pi^+\pi^-$ process at the region of the ρ meson resonance by (21) (see Fig. 3), achieved with $\chi^2/ndf = 1.5443$, gives the ρ meson mass and width

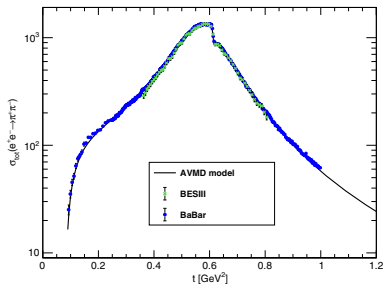


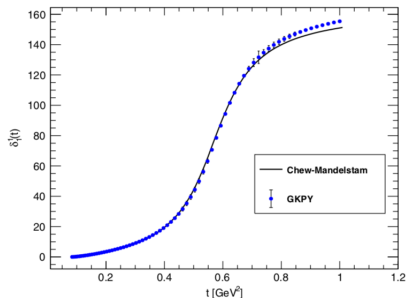
FIG. 3. Optimal description of the unified BESIII-BABAR data on $\sigma_{\text{tot}}(e^+e^- \rightarrow \pi^+\pi^-)$ at the elastic region by the pion EM FF U&A model.

FIG. P-wave effect

113004-4

Fit to δ_1^1 data using Chew-Mandelstam type effective-range formula

- ▶ $\frac{q^3}{\sqrt{s}} \text{Cotg}\delta_1^1(s) = a + bq^2 + h(s)q^2$
where $h(s) = \frac{2q}{\pi\sqrt{s}} \text{Log}\left(\frac{\sqrt{s}+2q}{2m_\pi}\right)$
- ▶ $a = 0.2860 \pm 0.0011 \text{ MeV}^2$,
 $b = -2.7025 \pm 0.0089$
- ▶ $\chi^2 = 2.45 \text{ pdf}$
- ▶ $m_\rho = (772.42 \pm 0.03) \text{ MeV}$
- ▶ $\Gamma_\rho = (153.85 \pm 0.11) \text{ MeV}$



Fit to δ_1^1 data

► $F_{\pi}^{EM,I=1}(s) =$

$$P_n(s) \exp \left[\frac{s}{\pi} \int_4^{\infty} \frac{\delta_1^1(s')}{s'(s' - s)} ds' \right]$$

► $\delta_1^1(q) = \text{ArcTan} \frac{A_3 q^3 + A_5 q^5 + \dots}{1 + A_2 q^2 + A_4 q^4 + \dots}$

► $\equiv \text{Log} \frac{(q - q_2)(q - q_3)(q - q_4)(q - q_5)}{(q - q_2^*)(q - q_3^*)(q - q_4^*)(q - q_5^*)}$

► $\chi^2 = 0.024$ pdf

► $m_{\rho} = (763.56 \pm 0.51) \text{ MeV}$

► $\Gamma_{\rho} = (143.09 \pm 0.82) \text{ MeV}$

$$\delta_1^1(q) = \text{arctg} \frac{A_3 q^3 + A_5 q^5 + \dots}{1 + A_2 q^2 + A_4 q^4 + \dots}, \quad (26)$$

where $A_1 \equiv 0$ in order to secure the threshold behavior of $\delta_1^1(q)$. An optimal description of the GKPY phase shift $\delta_1^1(q)$ data is achieved (see Fig. 5) with $\chi^2/ndf = 0.0244$ and four nonzero coefficients A_2 , A_3 , A_4 , and A_5 .

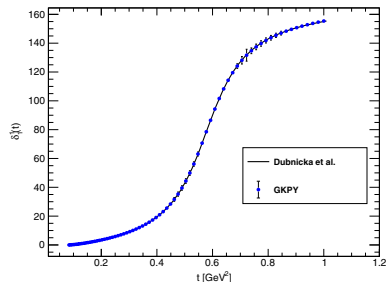
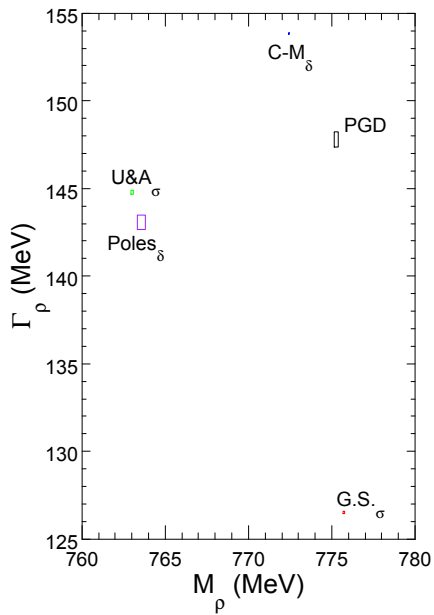
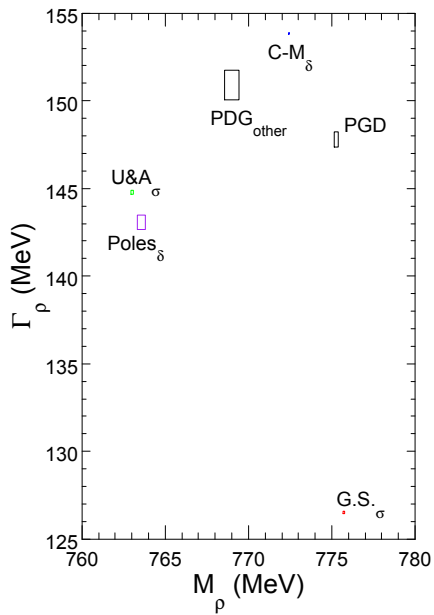
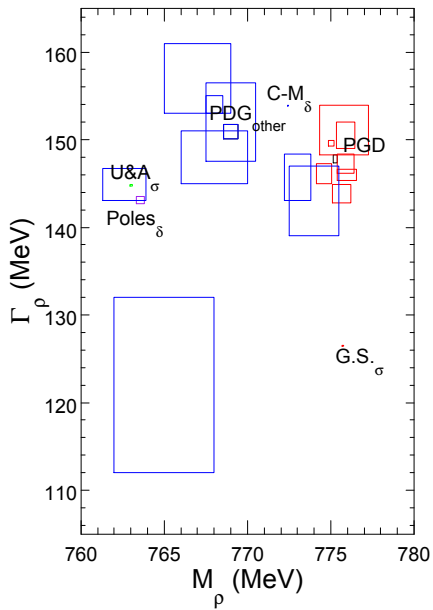


FIG. 5. Optimal description of the most accurate up-to-now P-wave isovector $\pi\pi$ scattering phase shift $\delta_1^1(t)$ data with model-independent parametrization (26).







Generalisation of Gounaris-Sakurai model to excited ρ mesons $\rho(1450)$ and $\rho(1700)$

- ▶ $F_\pi = \frac{1}{1+\beta+\gamma} \left[F_{\rho(770)}^{GS}(s) \left(1 + \delta \frac{s}{m_\omega^2} BW_\omega(s) \right) + \beta F_{\rho(1450)}^{GS}(s) + \gamma F_{\rho(1700)}^{GS}(s) \right]$
- ▶ $\chi^2 = 0.98 \text{ pdf}$

Generalisation of U&A model to excited ρ mesons

- ▶ $F_\pi = \frac{\Pi(q-q_i)}{\Pi(q+q_i^*)}$
- ▶ $\chi^2 = 1.84 \text{ pdf}$

ERIK BARTOŠ *et al.*

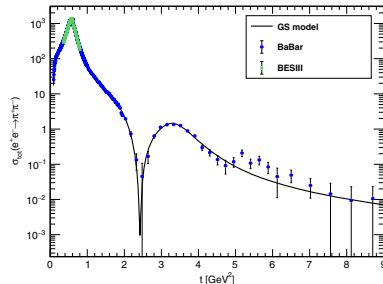


FIG. 6. Optimal description of the unified BESIII-BABAR complete data on $\sigma_{\text{tot}}(e^+e^- \rightarrow \pi^+\pi^-)$ by the generalized pion EM FF G.-S. model.

determined by the original pion EM FF G.-S. model (13) to be valid only at the elastic region.

A totally different situation is in a generalization of the U&A pion EM FF model. Here, the contribution of all three vector mesons is at an equal level. Only now, the effective inelastic threshold, which is left as a free parameter of the model, has to be taken into account explicitly. Therefore, instead of the q variable, the W variable

TABL
of BE
 e^+e^-
Analy

Param

m_ρ
 $m_{\rho'}$
 $m_{\rho''}$
 Γ_ρ
 $\Gamma_{\rho'}$
 $\Gamma_{\rho''}$

χ^2/ndf

$W(t)$

is no
U&A
surfa
Th
of all

Parameter	PDG MeV	G.S. MeV	U&A MeV
m_ρ	775.26 ± 0.25	774.81 ± 0.01	763.88 ± 0.04
$m_{\rho'}$	1465.00 ± 25.00	1497.70 ± 1.07	1326.35 ± 3.46
$m_{\rho''}$	1720.00 ± 20.00	1848.40 ± 0.09	1770.54 ± 5.49
Γ_ρ	149.10 ± 0.80	149.22 ± 0.01	144.28 ± 0.01
$\Gamma_{\rho'}$	400.00 ± 60.00	442.15 ± 0.54	324.13 ± 12.01
$\Gamma_{\rho''}$	250.00 ± 100.00	322.48 ± 0.69	268.98 ± 11.40
χ^2 pdf		0.98	1.84
		14 param.	11 param.

WHAT ARE THE CORRECT $\rho^0(770)$ MESON MASS ...

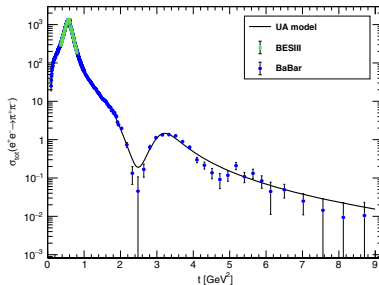


FIG. 7. Optimal description of the unified BESIII-BABAR complete data on $\sigma(e^+e^- \rightarrow \pi^+\pi^-)$ by the generalized pion

meson resonances. For this aim, totally the P-wave isovector the total cross section is exploited.

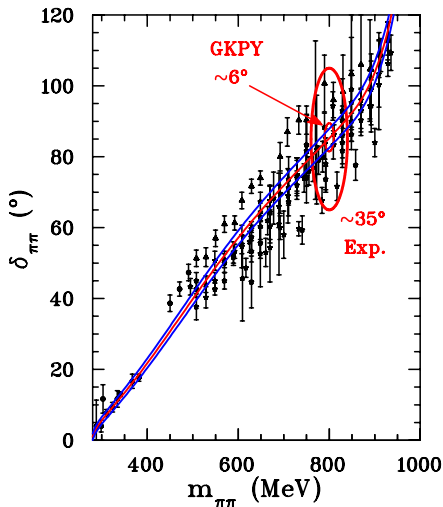
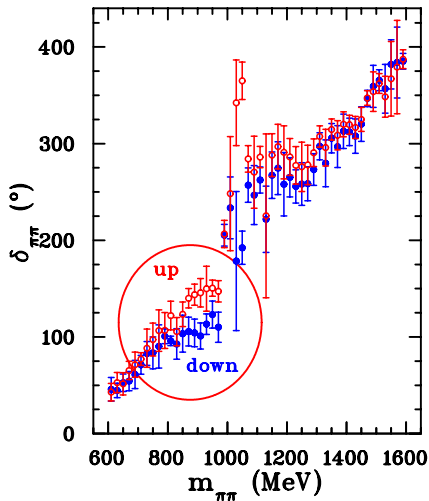
Just by a comparison of the ρ^0 meson parameters to the conclusion, the most likely given.

We conjecture that the $\rho(770)$ and the $\rho(770)$ are listed in Table II, i.e. $m_\rho = 775.26 \pm 0.25$ MeV, $\Gamma_\rho = 144.06 \pm 0.80$ MeV. Considerations in the next section parameters in the

We would like to mention that the $\rho(770)$ is listed in [15, 16] in which

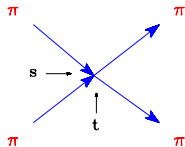
Experimental data for the $\pi\pi$ in the S0 wave (J1)

In PWA (CERN-Munich group'74) $A(s, t) \sim \cos(\theta_S - \theta_P)$



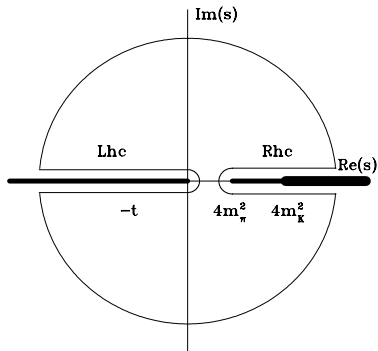
Dispersion relations with imposed crossing symmetry condition for $\pi\pi$ interactions theory $\longleftrightarrow$ experiment

crossing symmetry:



$$\rightarrow \vec{T}_s(s, t) = \hat{C}_{st} \vec{T}_t(t, s)$$

$\vec{T}(s, t)$ + crossing symmetry $\rightarrow$ dispersion relations for $4m_\pi^2 < s < \sim (1150 \text{ MeV})^2$

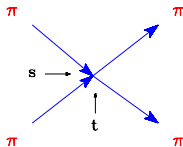


Once subtracted DR:

$$\begin{aligned} \text{Re } \vec{F}(s, t) &= \text{Re } \vec{F}(s_0, t) + \frac{s - s_0}{\pi} \\ &\times \left[\int_{4m_\pi^2}^{\infty} ds' \frac{\text{Im } \vec{F}(s', t)}{(s' - s_0)(s' - s)} \right. \\ &\left. + \int_{-t}^{-\infty} ds' \frac{\text{Im } \vec{F}(s', t)}{(s' - s_0)(s' - s)} \right] \end{aligned}$$

Dispersion relations with imposed crossing symmetry condition for $\pi\pi$ interactions theory $\longleftrightarrow$ experiment

crossing symmetry:



$$\rightarrow \vec{T}_s(s, t) = \hat{C}_{st} \vec{T}_t(t, s)$$

$\vec{T}(s, t)$ + crossing symmetry $\rightarrow$ dispersion relations for $\underline{4m_\pi^2 < s < \sim (1150 \text{ MeV})^2}$

Once subtracted dispersion relations ("GKPY" for the S and P waves):

$$\text{Re } t_\ell^{I(OUT)}(s) = \sum_{I'=0}^2 C_{st}^{II'} a_0^{I'} + \sum_{I'=0}^2 \sum_{\ell'=0}^4 \int_{4m_\pi^2}^{\infty} ds' K_{\ell\ell'}^{II'}(s, s') \text{Im } t_{\ell'}^{I'(IN)}(s')$$

$a_0^{I'}$ - subtraction constant = $\vec{T}_s(s = 4m_\pi^2, t = 0)$ - scattering lengths from only S wave

due to $\text{Re } t_\ell^I(k) = k^{2\ell}(a_\ell^I + b_\ell^I k^2 + O(k^4))$

$$\underline{\text{Re } t_\ell^{I(OUT)}(s) - \text{Re } t_\ell^{I(IN)}(s) \rightarrow 0}$$

GKPY equations and $\pi\pi$ amplitudes

partial waves: JI

experiment

F1

D2

S0

D0

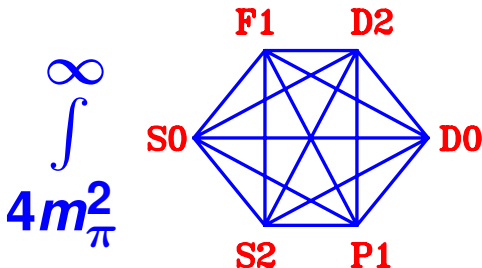
S2

P1

GKPY equations and $\pi\pi$ amplitudes

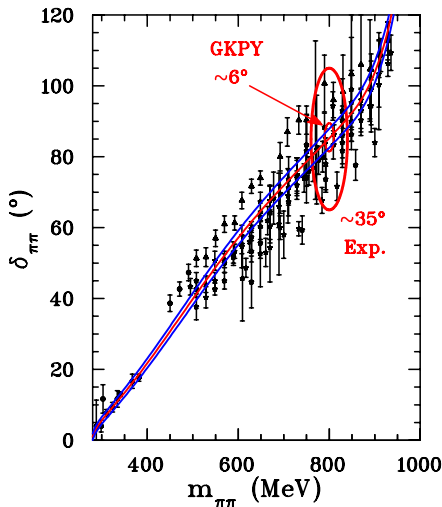
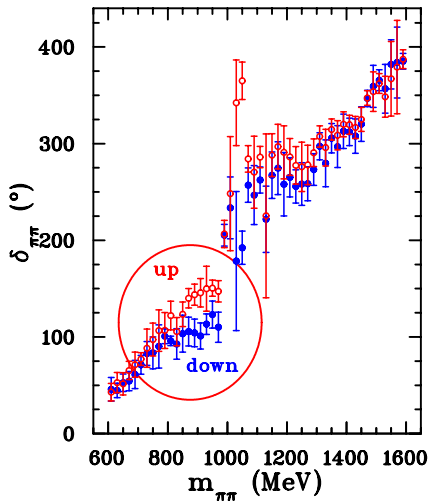
partial waves: Jl

experiment + theory (GKPY)



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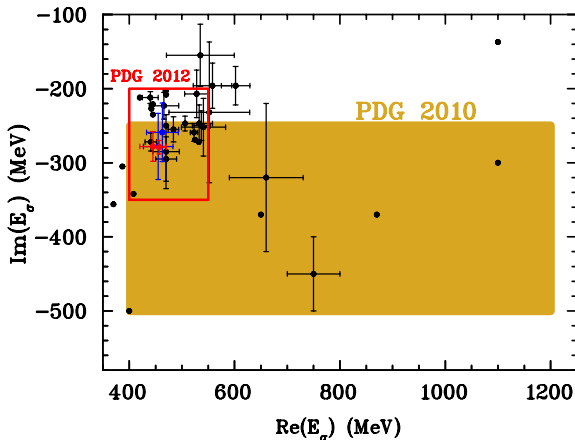
precise determination of $f_0(500)$ (σ) meson and threshold parameters

$f_0(500)$ (σ)

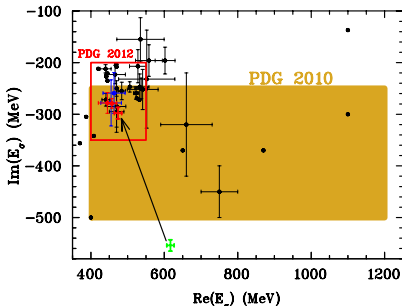
- ▶ PDG 2010:
 $M = 400 - 1200$ MeV
 $\Gamma = 2 \times (250 - 500)$ MeV
- ▶ PDG 2012:
 $M = 400 - 550$ MeV
 $\Gamma = 2 \times (200 - 350)$ MeV
- ▶ GKPY:
 $E_\sigma = 457 \pm 14 - i279^{+11}_{-7}$ MeV

threshold parameters, e.g. a_0^0 :

- ▶ ChPT + Roy eqs (Bern group):
 $0.220 \pm 0.005 m_\pi^{-1}$
- ▶ GKPY:
 $0.220 \pm 0.008 m_\pi^{-1}$



what forces GKPY eqs to pull up-left the sigma pole?



Two things: **trigonometry** and **crossing symmetry algebra** lead to narrower and lighter σ .

Modified $\pi\pi$ amplitude with σ pole PRD 90, 116005 (2014) P. Bydzovský, 1, R. Kamiński, V. Nazari

Nothing more and nothing instead of it is needed.

Conclusions and what to do?

Conclusions:

- ▶ there is significant difference between parameters (especially mass) of a resonance using different methods of parameterisation of amplitudes,
- ▶ amplitudes should be unitary and analytic
- ▶ we should always determine what definition we use,
- ▶ the determination of the most significant poles in amplitudes becomes difficult for $n > 2$,

What to do?

- ▶ never use isobar model!
- ▶ anyway be carefull! = check the 2nd point of conclusions!
- ▶ possible parameterizations: made for wide energy range including many resonances we are interested in,
- ▶ please use only parameterization based on S - or K -matrix approach!
- ▶ possible degrees of freedom: poles themselves! mesons
- ▶ on $\pi\pi$, $K\bar{K}$... amplitudes in all partial waves up to 1.8 GeV ask RK - they are done
- ▶ on other channels ask RK - they may be done

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Conclusions on scalar resonances

Experimental situation in the 0^+0^{++} sector is finally clear (although without new experiments!),

- ▶ $f_0(500)$ (σ) - loosely bound $\pi\pi$ system or/and $2q2\bar{q}$ state,
- ▶ $f_0(980)$ - dynamically generated $2q2\bar{q}$ state or/and $K\bar{K}$ quasi bound state,
- ▶ $f_0(1370)$ - ordinary $q\bar{q}$ state,
- ▶ $f_0(1500)$ - ordinary $q\bar{q}$ state + glueball,
- ▶ $f_0(1720)$ - glueball

Probably $K_0^*(800)$ (κ) is also dynamical effect of coupling of $K_0^*(1430)$