

# Dalitz Plot Decomposition

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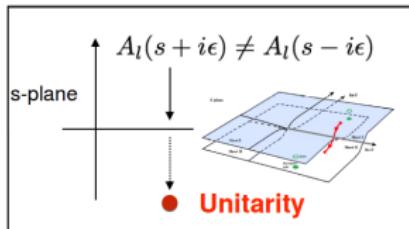
December 18<sup>nd</sup>, 2019



# Non-perturbative hadron physics

obeys general principles of the scattering theory

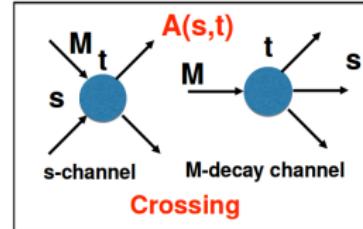
- **Lorentz invariance** = independence of the reference frame, known behavior under boosts and rotations
- **Unitarity** = constraint to imaginary part of scattering amplitude
- **Analyticity** = implementation of relevant, closest singularities
- **Crossing** = decay and scattering regions are analytically connected



$$A(s, t) = \sum_l A_l(s) P_l(z_s)$$

**Analyticity**

$$A_l(s) = \lim_{\epsilon \rightarrow 0} A_l(s + i\epsilon)$$

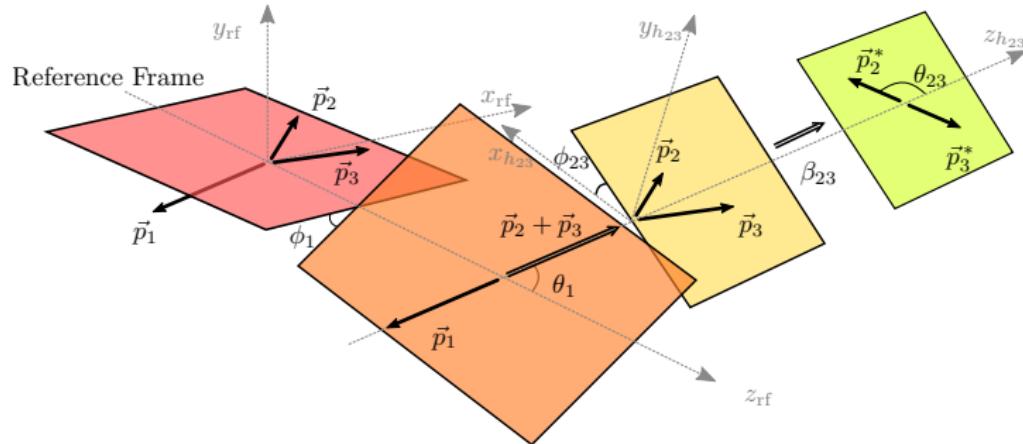


# Introduction

# Decay variables and decay chain

## Three-body kinematics

- depends on 5 variables
- a decay chain - a sequence of decays, (here is two)



$$D^J(\phi_1, \theta_1, 0) \times (m_{23}^2) D^S(\phi_{23}, \theta_{23}, 0)$$

- the full amplitude - sum of many decay chains
- the intensity - amplitude squared, summed over spin orientation

# Conventional helicity approach

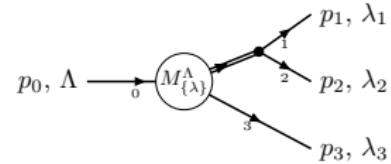
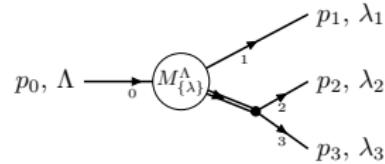
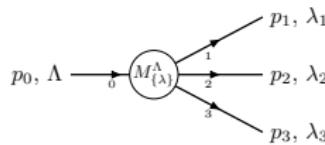
Complicated cases: particles with spin in isobar model [Herndon(1975)], [Hansen (1974)]

# Conventional helicity approach

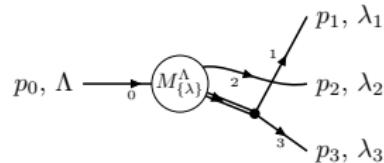
Complicated cases: particles with spin in isobar model [Herndon(1975)], [Hansen (1974)]

$$M_{\{\lambda\}}^{\Lambda} = M_{1,\{\lambda\}}^{\Lambda} + M_{2,\{\lambda\}}^{\Lambda} + M_{3,\{\lambda\}}^{\Lambda}$$

$$\underbrace{M_{\{\lambda\}}^{\Lambda}}_{= H_1 D(\phi_1, \theta_1, 0) D(\phi_{23}, \theta_{23}, 0) W_1(\dots)} + \underbrace{H_3 D(\phi_3, \theta_3, 0) D(\phi_{12}, \theta_{12}, 0) W_3(\dots)}$$



$$+ \underbrace{H_2 D(\phi_2, \theta_2, 0) D(\phi_{31}, \theta_{31}, 0) W_2(\dots)}$$



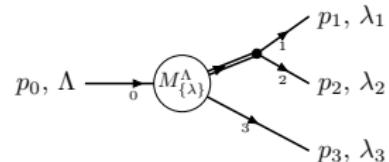
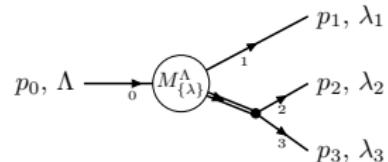
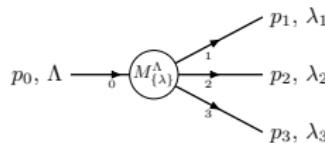
- A special set of angles for every decay chain

# Conventional helicity approach

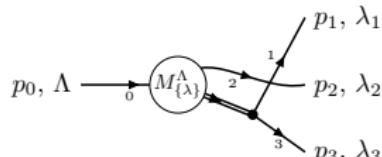
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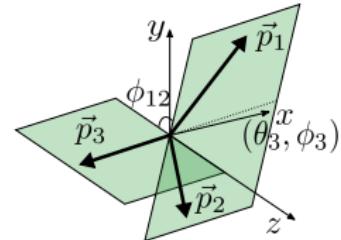
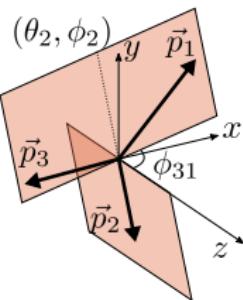
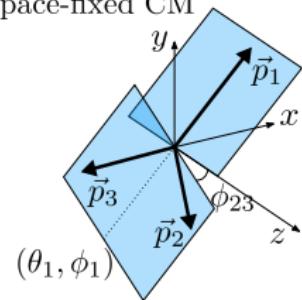


- A special set of angles for every decay chain
- Consistently of quantization direction – **Wigner rotations**

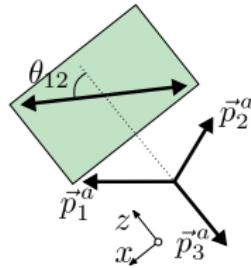
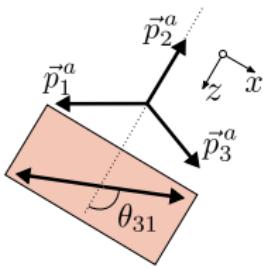
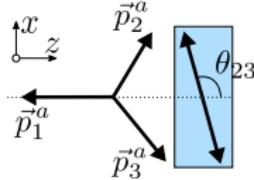
$$W_i(\dots) = D^{j_1}(\tilde{\phi}_1^i, \tilde{\theta}_1^i, 0) D^{j_2}(\tilde{\phi}_2^i, \tilde{\theta}_2^i, 0) D^{j_3}(\tilde{\phi}_3^i, \tilde{\theta}_3^i, 0)$$

# Different angles for decay chains

space-fixed CM



aligned CM



$$\begin{aligned}
 M_{\{\lambda\}} = & H_1 D(\phi_1, \theta_1, 0) D(\phi_{23}, \theta_{23}, 0) W_1(\dots) + H_3 D(\phi_3, \theta_3, 0) D(\phi_{12}, \theta_{12}, 0) W_3(\dots) \\
 & + H_2 D(\phi_2, \theta_2, 0) D(\phi_{31}, \theta_{31}, 0) W_2(\dots)
 \end{aligned}$$

# The Dalitz-Plot decomposition

[MM et al.(JPAC), arXiv:1910.04566]

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Reformulation of the helicity approach

$$p_0, \Lambda \xrightarrow[0]{} M_{\{\lambda\}}^A \xrightarrow[1]{2}{3} p_1, \lambda_1 \quad p_2, \lambda_2 \quad p_3, \lambda_3$$

$$\left. \begin{array}{c} \sigma_3 \\ \sigma_1 \end{array} \right\} = \sum_{\nu} \underbrace{D_{\Lambda\nu}^{J*}(\phi_1, \theta_1, \phi_{23})}_{\text{Decay-plane orientation}} \times \underbrace{O_{\{\lambda\}}^\nu(\{\sigma\})}_{\text{Dalitz-plot function}}$$

Model-independent factorization of the overall rotation:

- Exploits properties of the Lorentz group (orientation – just three Euler angles)
- Dalitz-plot function depends entirely on 2 variables,  $\{\sigma\} \equiv \{\sigma_1, \sigma_2, \sigma_3\}$
- No azimuthal phase factors in  $O_{\{\lambda\}}^\nu$ .

# The Dalitz-Plot decomposition

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Reformulation of the helicity approach

$$p_0, \Lambda \xrightarrow[0]{} M_{\{\lambda\}}^A \xrightarrow[1]{} p_1, \lambda_1 \quad \xrightarrow[2]{} p_2, \lambda_2 \quad \xrightarrow[3]{} p_3, \lambda_3$$

$$\left. \begin{array}{c} p_1, \lambda_1 \\ p_2, \lambda_2 \\ p_3, \lambda_3 \end{array} \right\} \begin{array}{l} \sigma_3 \\ \sigma_1 \end{array} = \sum_{\nu} \underbrace{D_{\Lambda\nu}^{J*}(\phi_1, \theta_1, \phi_{23})}_{\text{Decay-plane orientation}} \times \underbrace{O_{\{\lambda\}}^{\nu}(\{\sigma\})}_{\text{Dalitz-plot function}}$$

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Gives significant benefits to

- Pentaquark analysis,  $\Lambda_b/\Lambda_c$  polarionation measurements, Baryonic decay chains, . . .

# Dalitz-Plot function

[MM et al.(JPAC), arXiv:1910.04566]

## Master formula $0 \rightarrow 123$ decay with arbitrary spins

$$O_{\{\lambda\}}^{\nu}(\{\sigma\}) = \sum_{(ij)k} \sum_s^{(ij) \rightarrow i,j} \sum_{\tau} \sum_{\{\lambda'\}} n_J n_s d_{\nu, \tau - \lambda'_k}^J(\hat{\theta}_{k(1)}) X_s^{\tau, \lambda'_k; \lambda'_i, \lambda'_j}(\sigma_k) d_{\tau, \lambda'_i - \lambda'_j}^s(\theta_{ij}) \\ \times d_{\lambda'_1, \lambda_1}^{j_1}(\zeta_{k(0)}^1) d_{\lambda'_2, \lambda_2}^{j_2}(\zeta_{k(0)}^2) d_{\lambda'_3, \lambda_3}^{j_3}(\zeta_{k(0)}^3),$$

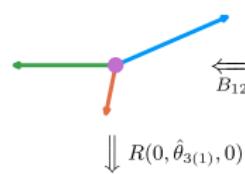
- Three decay chains,  $(ij)k \in \{(12)3, (23)1, (31)2\}$ .
- $\theta_{ij} = \theta_{ij}(\{\sigma\})$  is an isobar decay angle
- $\hat{\theta}_{k(1)} = \hat{\theta}_{k(1)}(\{\sigma\})$  is the particle-0 Wigner angle
- $\zeta_{k(0)}^i = \zeta_{k(0)}^i(\{\sigma\})$  is the particle- $i$  Wigner angle
- $X_s^{\tau, \lambda'_k; \lambda'_i, \lambda'_j}(\sigma_k) \Rightarrow X_s^{LS; l's'}(\sigma_k)$  is the only model-dependent input (lineshape functions)

# Origin of the Wigner angle

[MM et al.(JPAC), arXiv:1910.04566]

I) aligned CM

chain 3



II) isobar-( $ij$ )

$B_{12}$

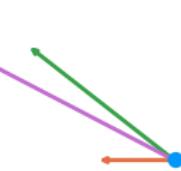
$R_{12}$

III)  $p_1$ -aligned

chain 3



IV)  $p_1$ -rest

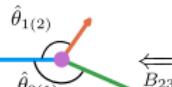


$R(0, \hat{\theta}_{3(1)}, 0)$

$R(0, \zeta_{3(1)}^1, 0)$

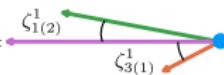
chain 1

$\hat{\theta}_{1(2)}$   
 $\hat{\theta}_{3(1)}$



$B_{23}$

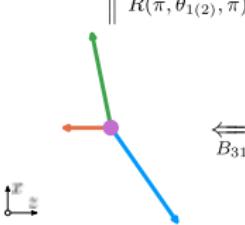
$R_{23}$



$R(\pi, \hat{\theta}_{1(2)}, \pi)$

$R(\pi, \zeta_{1(2)}^1, \pi)$

chain 2



$B_{31}$

$R_{31}$



$\hat{x}$   
 $\hat{z}$

# $\Lambda_c \rightarrow p K \pi$ polarization studies

- proposal for the electromagnetic dipole moments of charmed baryons  
[EPJC 77 (2017) 181, arXiv:1708.08483]
- polarization information for complex decay chains

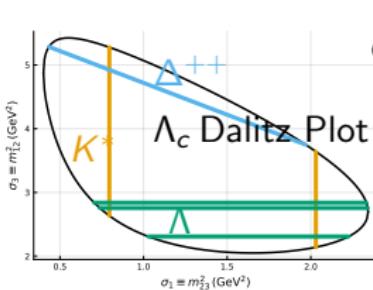
$$M_\lambda^\Lambda = \sum_\nu D_{\Lambda\nu}^{1/2*}(\phi_1, \theta_1, \phi_{23}) O_\lambda^\nu(\sigma_1, \sigma_3),$$

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$$M_\lambda^\Lambda = \sum_\nu D_{\Lambda\nu}^{1/2*}(\phi_1, \theta_1, \phi_{23}) O_\lambda^\nu(\sigma_1, \sigma_3),$$

- Resonances in all channels,  $\Lambda^0$ ,  $K^{*0}$ ,  $\Delta^{++}$
- Possible Triangle Singularity near  $\Lambda\eta$  threshold [Liu, Xiao-Hai et al., PRD100 (2019)]



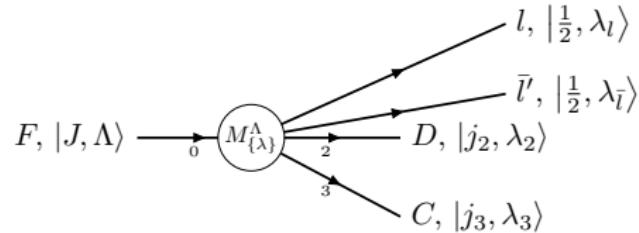
$$\begin{aligned}
 O_\lambda^\nu(\sigma_1, \sigma_3) = & \sum_s \sum_\tau d_{\nu, \tau - \lambda}^{1/2}(0) X_s^{\tau, \lambda}(\sigma_1) d_{\tau, 0}^s(\theta_{23}) \\
 & + \sum_s \sum_{\tau, \lambda'} d_{\nu, \tau}^{1/2}(\hat{\theta}_{2(1)}) X_s^{\tau, \lambda'}(\sigma_2) d_{\tau, -\lambda'}^s(\theta_{31}) d_{\lambda', \lambda}^{1/2}(\tilde{\theta}_{2(1)}^1) \\
 & + \sum_s \sum_{\tau, \lambda'} d_{\nu, \tau}^{1/2}(\hat{\theta}_{3(1)}) X_s^{\tau, \lambda'}(\sigma_3) d_{\tau, \lambda'}^s(\theta_{12}) d_{\lambda', \lambda}^{1/2}(\tilde{\theta}_{3(1)}^1).
 \end{aligned}$$

# Four-body decay with leptons

Four-body decay - 8 degrees of freedom,

$$M_{\lambda_2, \lambda_3, \lambda_I, \lambda_{\bar{I}}}^{\Lambda} \left( \underbrace{\phi_1, \theta_1, \phi_{23}}_{\text{polarization}}, \underbrace{q^2, m_{XD}^2, m_{XC}^2}_{\text{Dalitz-Plot dynamics}}, \underbrace{\phi_I, \theta_I}_{\text{decay to lepton}} \right) =$$

horribly complicated ?

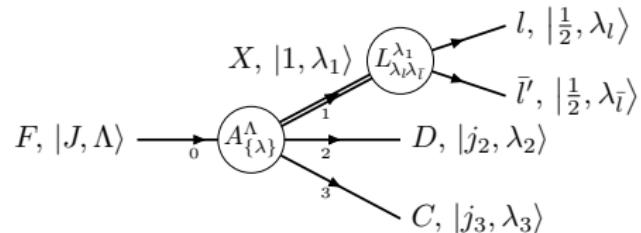


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Model-independent factorization:

$$= \sum_{\nu, \lambda_1} D_{\lambda, \nu}^{J*}(\phi_1, \theta_1, \phi_{23}) O_{\lambda_1, \lambda_2, \lambda_3}^{\nu}(q^2, m_{XD}^2, m_{XC}^2) D_{\lambda_1, \lambda_I - \lambda_{\bar{I}}}^{1*}(\phi_I, \theta_I, 0).$$



# Conclusion

## Dalitz-Plot decomposition

General **S-matrix principles** and **symmetries** are the main reliable tools for dealing with non-perturbative theory

- Requirement of Lorentz Invariance
  - Explicit separation of degrees of freedom
- ⇒ significant simplification of amplitude construction

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Practicalities:

- Julia framework [ThreeBodyDecay.jl](#)
- c++ code (thanks Vincent)



# Thank you for attention

JPAC group:

- Andrew Jakura, Alessandro Pilloni, Vincent Mathieu, Miguel Albaladejo, Cesar Fernandez, Adam Szczepaniak, Lukasz Bibrzycki, Daniel Winney.

LHCb collaborators:

- Sara Mitchell, Marco Pappagallo, Tomasz Skwarnicki