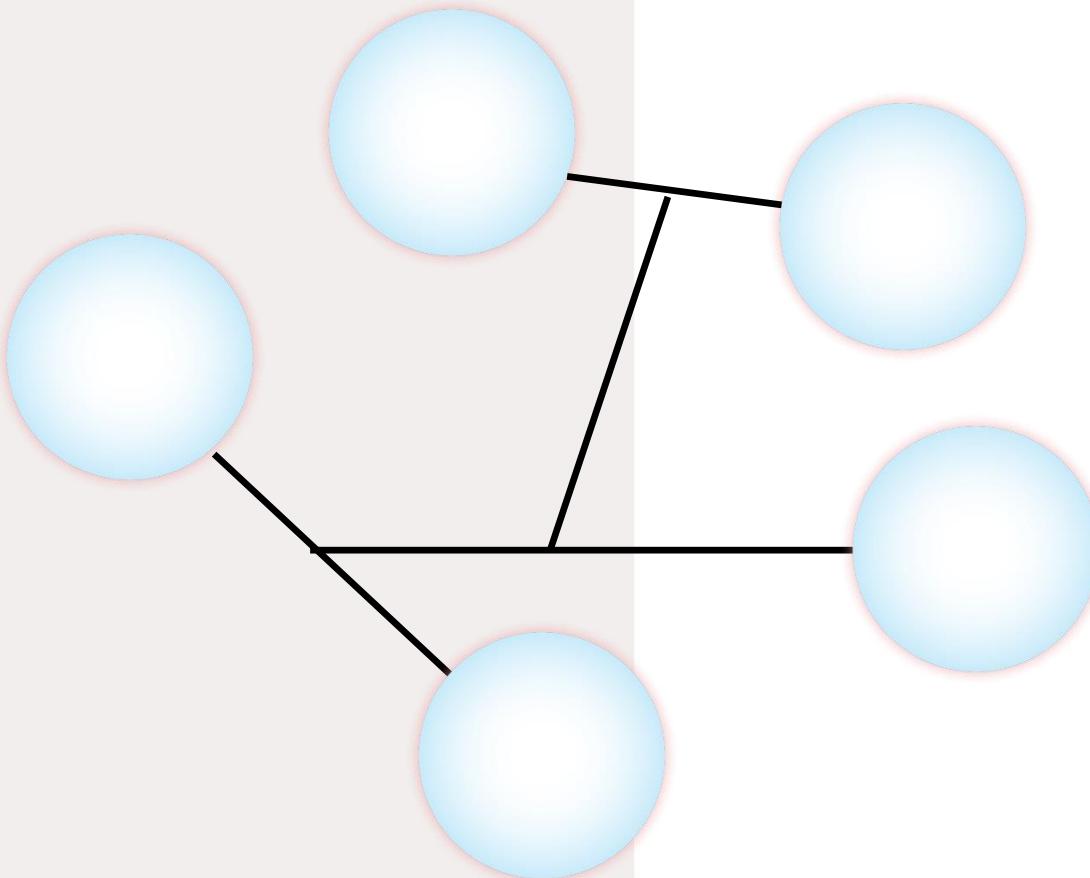


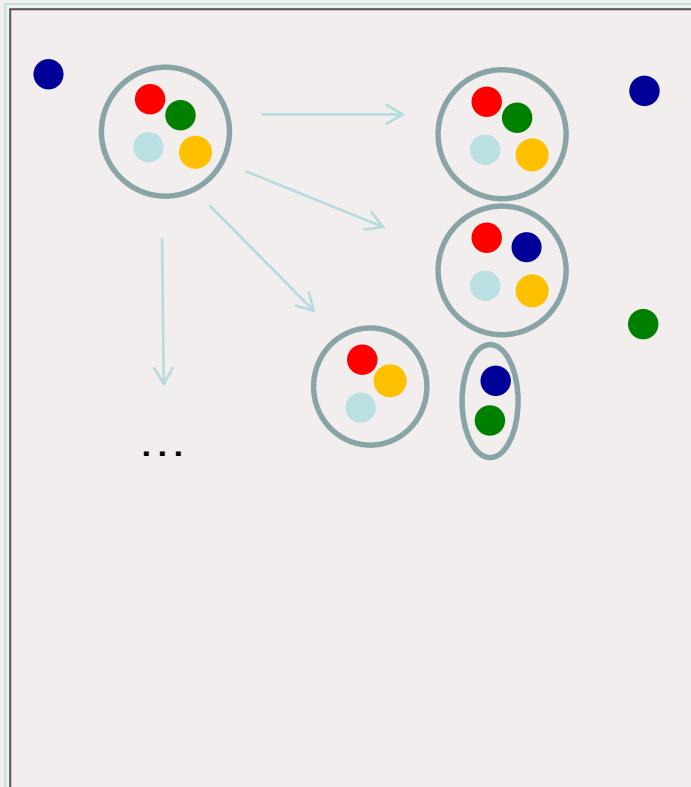
# Description of few-nucleon systems by solving Faddeev-Yakubovsky equations



- Faddeev-Yakubovsky equations
- Some applications
  - $4N$  systems
  - $N-{}^4\text{He}$  elastic scattering
  - $p\text{v}$  in low-energy  $n-{}^4\text{He}$  scattering
  - Resonances in  ${}^5\text{H}$

## Non-relativistic Collisions

- In configuration space wave functions extend to infinity!
- Increasingly complex asymptotic behaviour for  $A > 2$  systems!!



How to take care of the boundary condition?

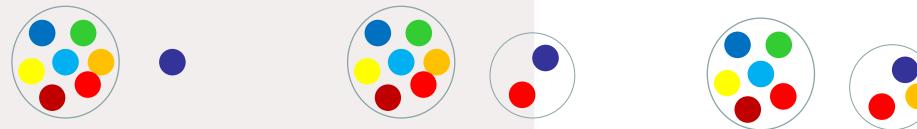
- ✓ Conceptual difficulties to uncouple different particle channel, to constrain asymptotes of the solutions in all directions and thus get unique (physical) solution to the Schrödinger eq.
  - It is ok, as long as there is single particle channel (elastic plus target/projectile excitations)
  - Mathematically ill-conditioned problem when several particle channels are open
- ✓ Faddeev-Yakubovsky equations efficiently separates asymptotes of the binary channels

L. D. Faddeev, Zh. Eksp. Teor. fiz. 39, 1459 (1960). [Sov. Phys. JETP 12, 1014 (1961)].  
O. A. Yakubovsky, Sov. J. Nucl. Phys. 5, 937 (1967).

# Properties of the rigorous scattering eq.

- Should separate all possible scattering channels to incorporate proper asymptotes! Number of binary channels increases  $\sim 2^N$

$$\Psi_N = \sum_{perm} \Psi_{(N-1)(1)} + \sum_{perm} \Psi_{(N-2)(2)} + \sum_{perm} \Psi_{(N-3)(3)} + \dots$$



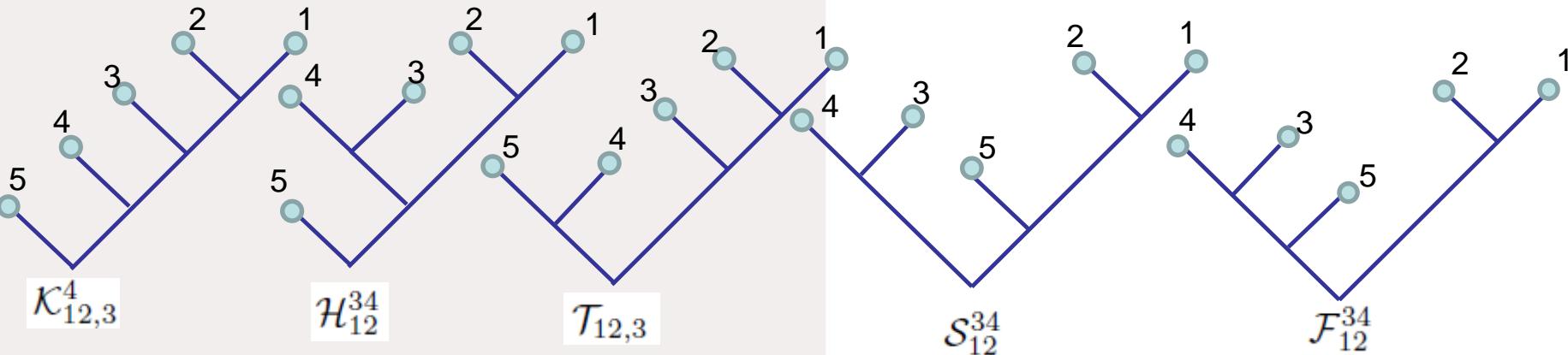
- Should be systematically reducible to smaller subsystems, in order to built proper asymptotic solutions and to be consistent to its subsystems: chain of partitions (tree-like structures to break system in clusters & subclusters)

$$\Psi_{(N-i)(i)} = \left( \Psi_{N-i} \bigcup \Psi_i \right)$$

- FY equations are derived following this pattern, reconnecting different partition chains

*very fast growth of components with N!!*

# Faddeev-Yakubovsky eq



Merits:

- ✓ Handling of symmetries
- ✓ Boundary conditions for binary channels
- ✓ Easy reduction to subsystems
- ✓ 3BF implemented at reasonable price
- ✓ Built for short-ranged interactions.

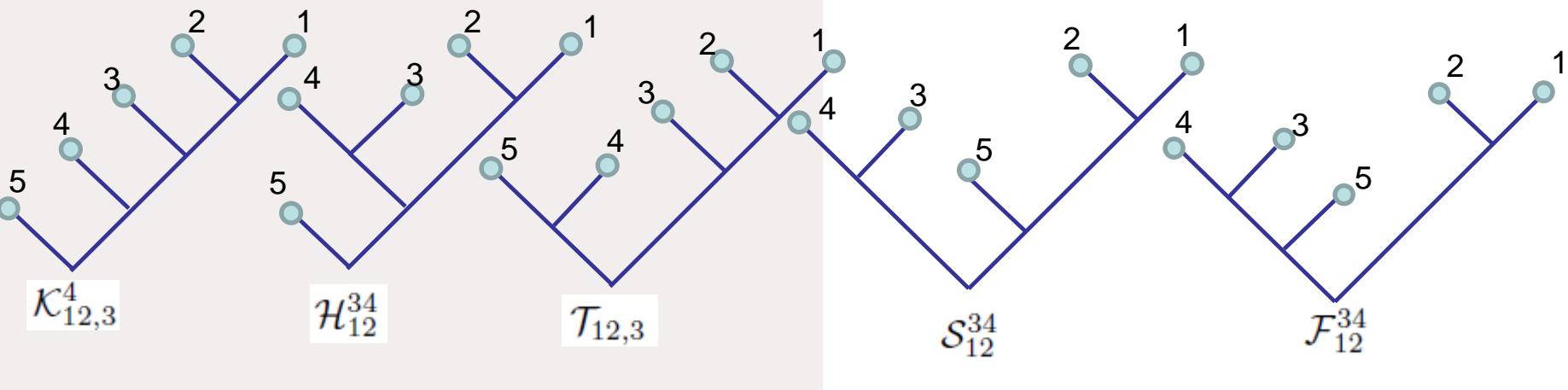
Treatment of Coulomb – true adventure,  
still reasonable for repulsive case.

Price

- ✓ Overcomplexity with  $N$

Problem	Number eq. (identical particles)	Number eq. (different particles)
$A=2$	1	1
$A=3$	1	3
$A=4$	2	18
$A=5$	5	180
$A=6$	15	2700
$A=N$	$\text{nint}\left(\frac{2(N-1)!}{(\pi/2)^N}\right)$	$\frac{N! (N-1)!}{2^{N-1}}$

# 5-body Faddeev-Yakubovski eq



$$\mathcal{K}_{12,3}^4(\vec{x}, \vec{y}, \vec{z}, \vec{w}, S, L, T) = \sum_{\alpha_K = (l_{..}, s_{..}, t_{..})} \frac{f_{\alpha_K}(x, y, z, w)}{xyzw} \left[ \left\{ (l_x l_y)_{l_{xy}} (l_z l_w)_{l_{zw}} \right\}_L \{ \dots \}_S \right]_{JM} \{ \dots \}_T$$

## NUMERICAL SOLUTION

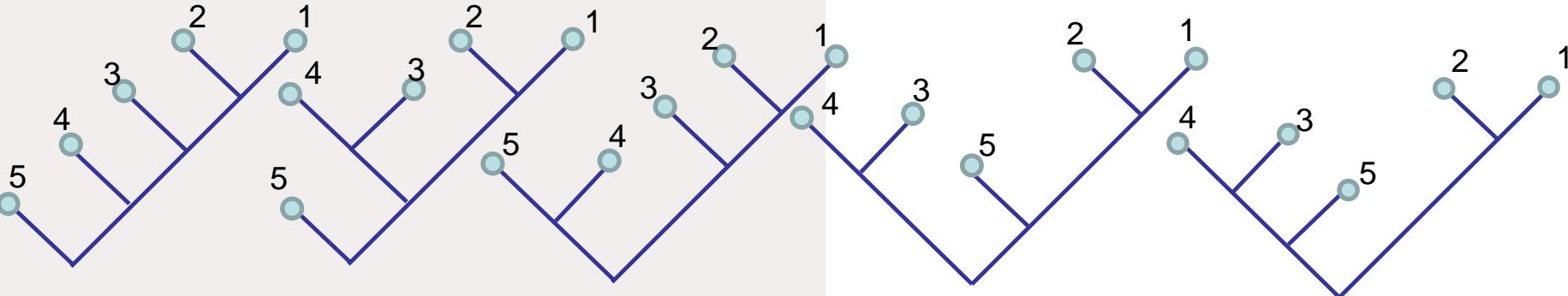
\*R.L., PhD Thesis, Université Joseph Fourier, Grenoble (2003).

- PW decomposition of the components  $K, H, T, S, F$
- Radial parts expanded using Lagrange-mesh method

D. Baye, Physics Reports 565 (2015) 1

- Resulting linear algebra problem solved using iterative methods
- Observables extracted using integral relations

# Numerical costs



Problem	Number eq. (ident particles)	Number eq. (diff. particles)	PW basis.	Radial disc.
2N	1	1	2	$\sim N$
3N	1	3	$\sim 100$	$\sim N^2$
4N	2	18	$\sim 10^4$	$\sim N^3$
5N	5	180	$\sim 10^6$	$\sim N^4$

## NUMERICAL SOLUTION

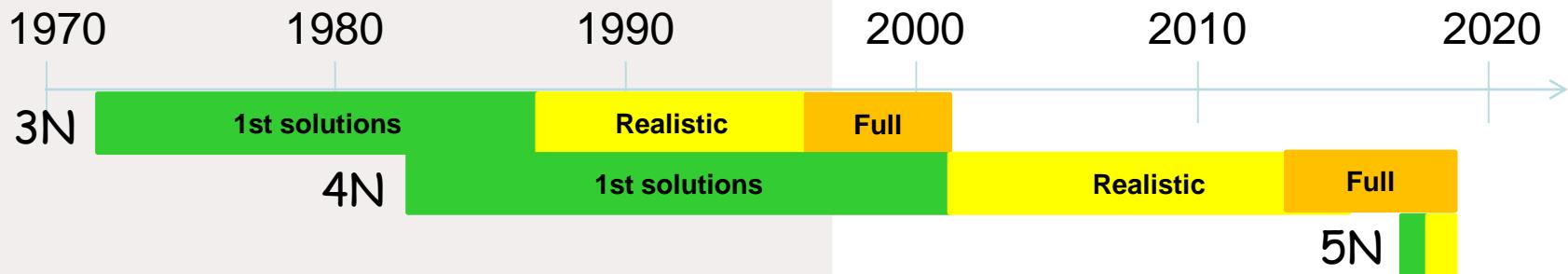
\*R.L., PhD Thesis, Université Joseph Fourier, Grenoble (2003).

- PW decomposition of the components K,H,T,S,F
- Radial parts expanded using Lagrange-mesh method

D. Baye, Physics Reports 565 (2015) 1

- Resulting linear algebra problem solved using iterative methods
- Observables extracted using integral relations

# Short overview of nuclear problems by FY eq's



## 3N-problem (Faddeev eq.)

1<sup>st</sup> solution: A. Laverne and C. Gignoux: Nucl. Phys. A 203 (1973) 597

G. Gignoux, A. Laverne, and S. P. Merkuriev Phys. Rev. Lett. 33 (1974) 1350

Review: W. Glöckle et al., Physics Reports 274 (1996) 107-285

Full: H. Witała et al., Phys. Rev. C 59, 3035 (1999)  
A. Deltuva et al., Phys. Rev. C 72, 054004 (2005)

## 4N-problem (Faddeev eq.)

1<sup>st</sup> solution: S. P. Merkuriev, S.L. Yakovlev, C. Gignoux, Nucl. Phys. A 431 (1984) 125.

Benchmarks: A. Nogga, et al., Phys. Rev. C 65 (2002) 054003 (bound state)

R. Lazauskas et al., Phys. Rev. C 71 (2005) 034004 ( $n^3H$  scattering)

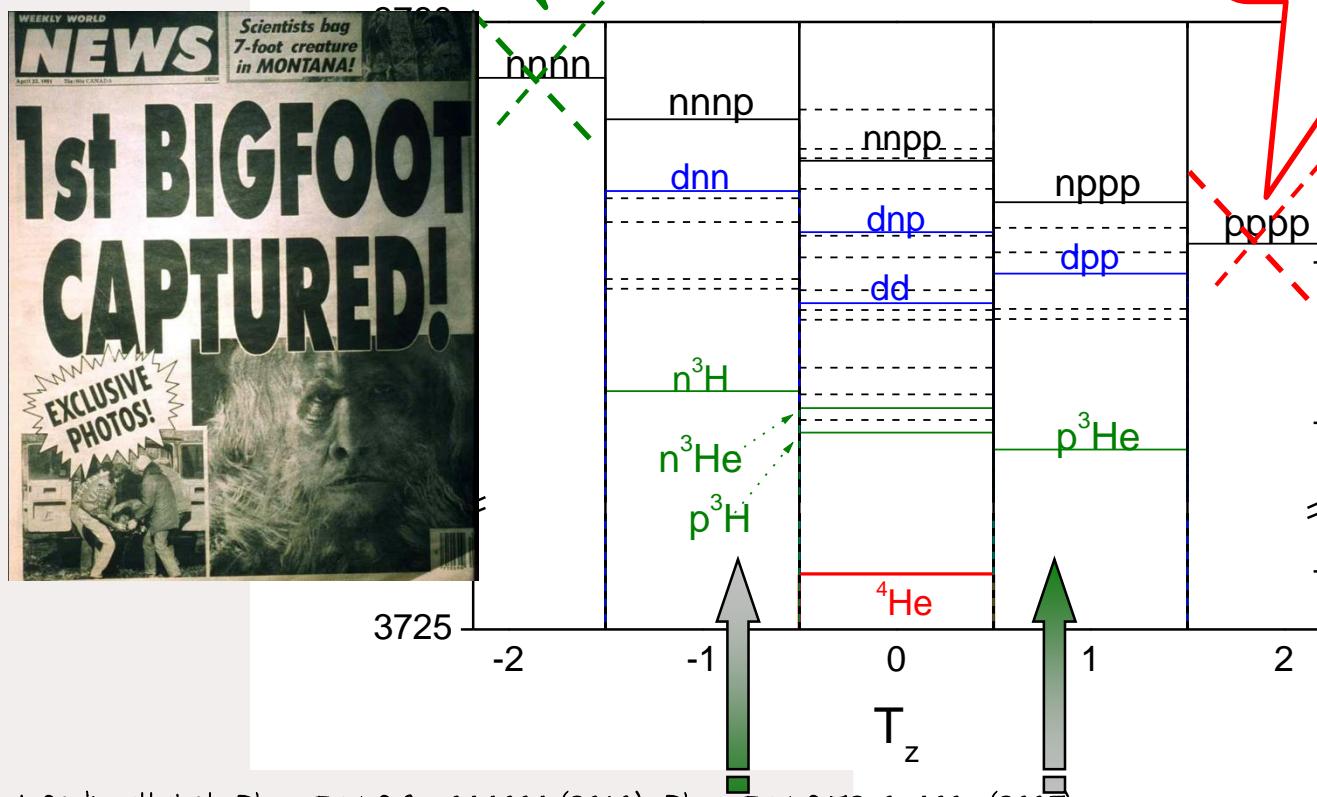
M. Viviani et al., Phys. Rev. C 84 (2011) 054010 ( $p^3He$  scattering)

M. Viviani et al., [arXiv:1610.09140](https://arxiv.org/abs/1610.09140) ( $p^3H, n^3He$  scattering)

Full: A. Deltuva and A. C. Fonseca, Phys. Rev. C 87, 054002(2013), Phys. Rev. C 95, 024003 (2017)  
R. Lazauskas, Phys. Rev. C 91, 041001(R) (2015)

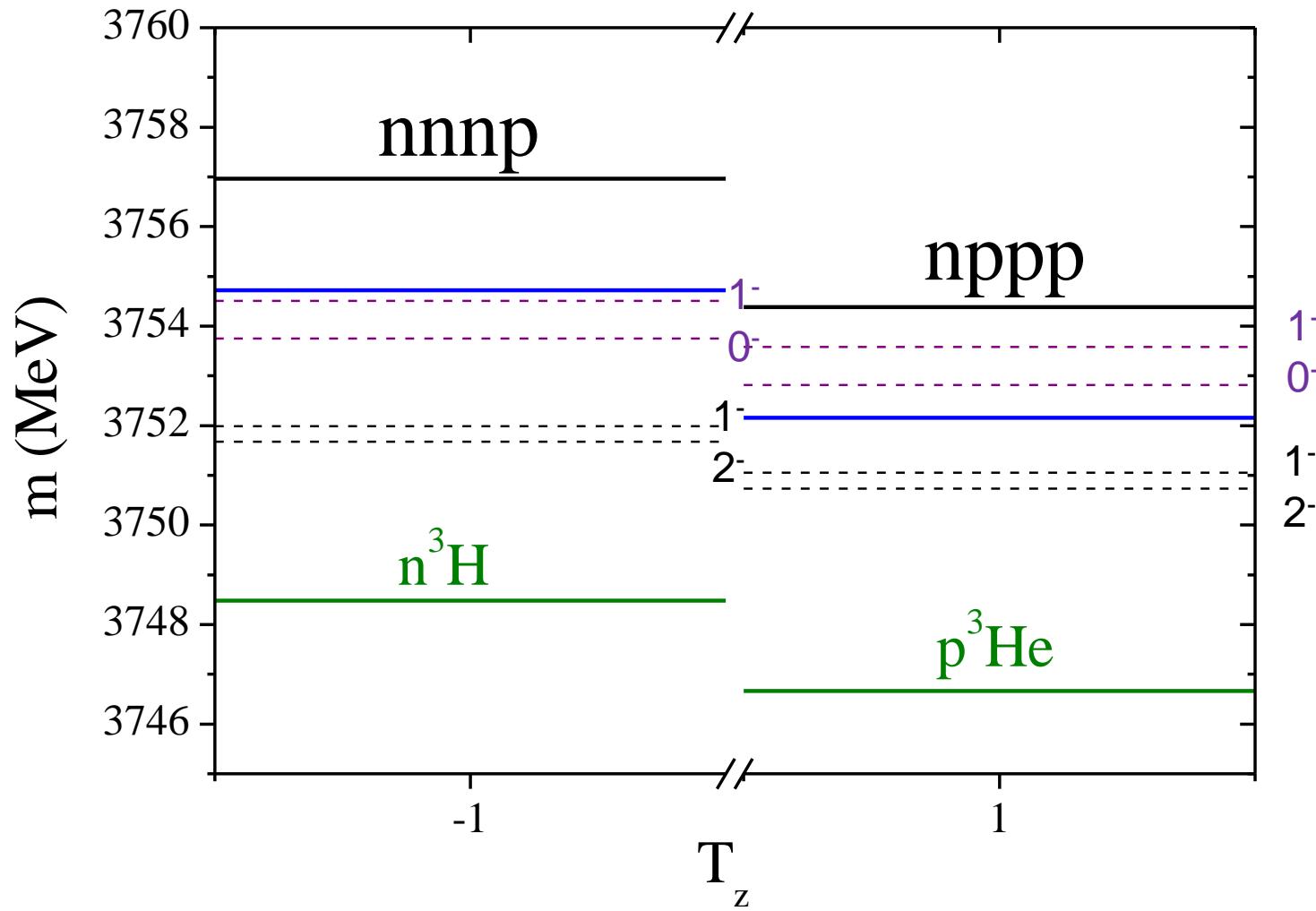
Occasional  
experimental &  
theoretical claims:  
Bigfoot effect?

No claims yet. El  
Dorado for  
“Bigfoot hunters”

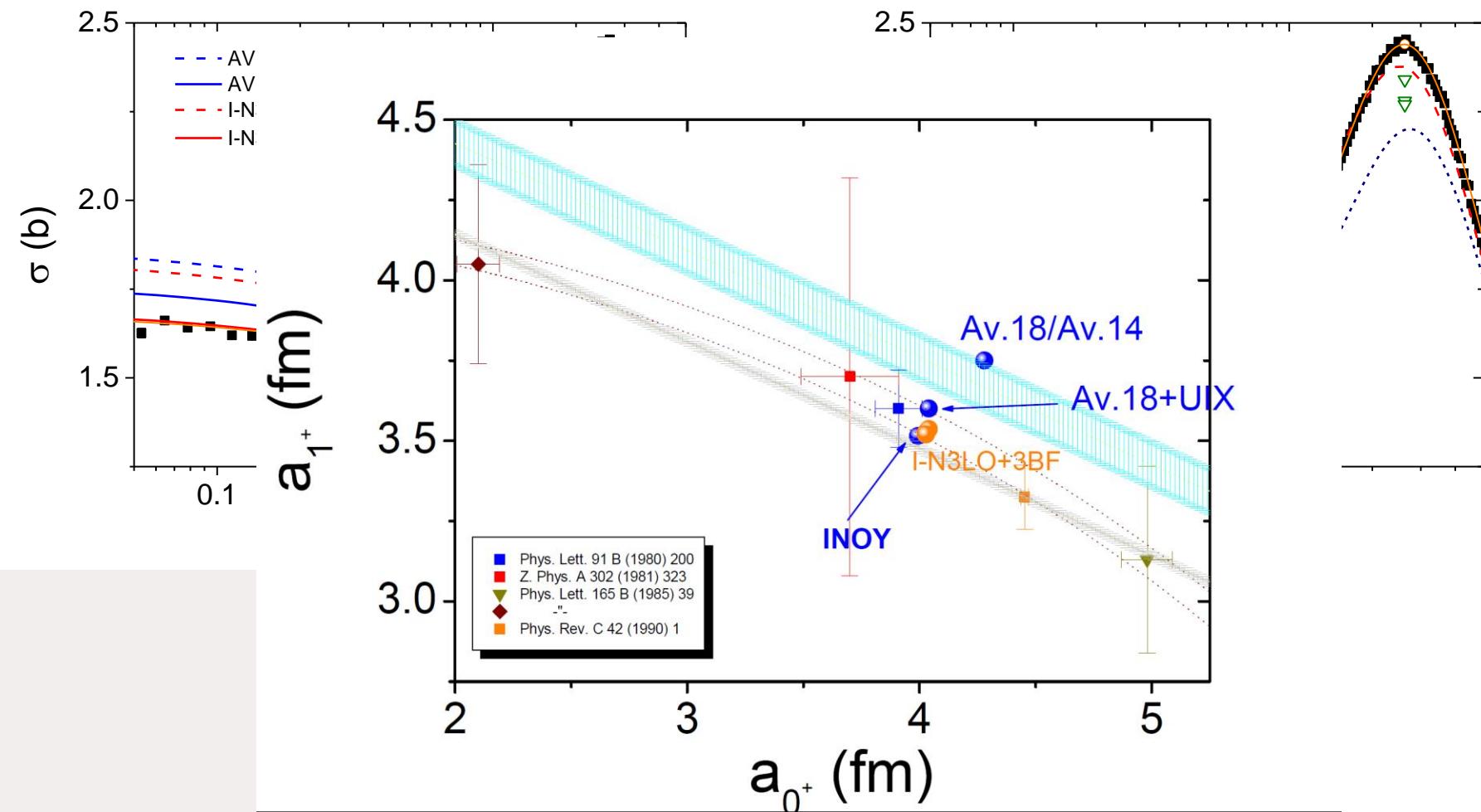


R. L. J. Carbonell et al., Phys. Rev. C 93, 044004 (2016), Phys. Rev. C 72, 034003 (2005)

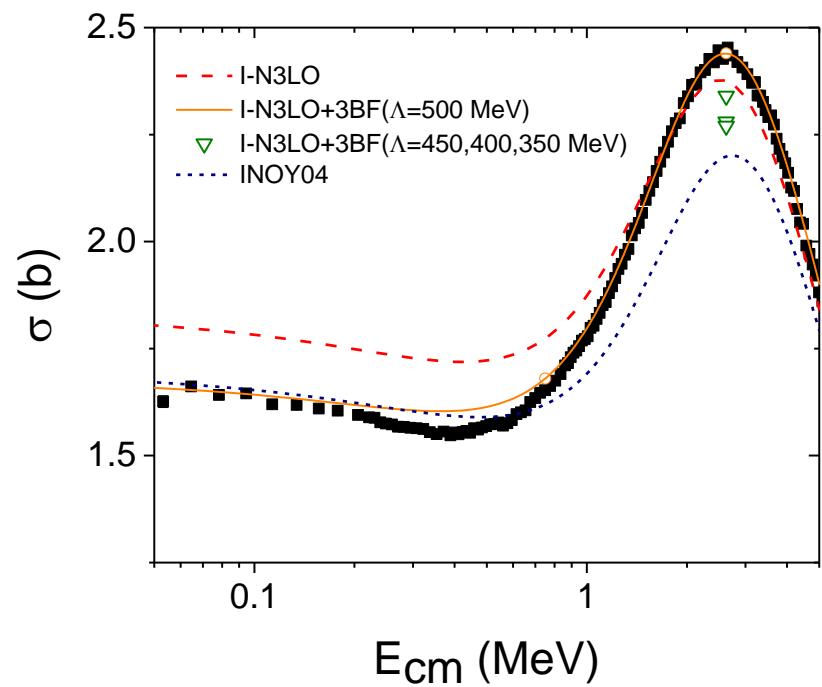
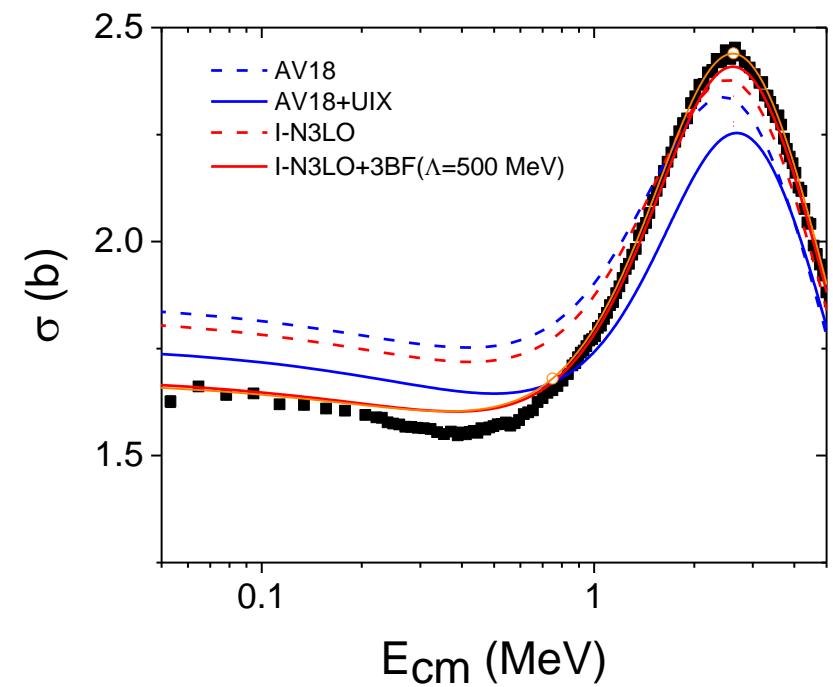
# 4N systems



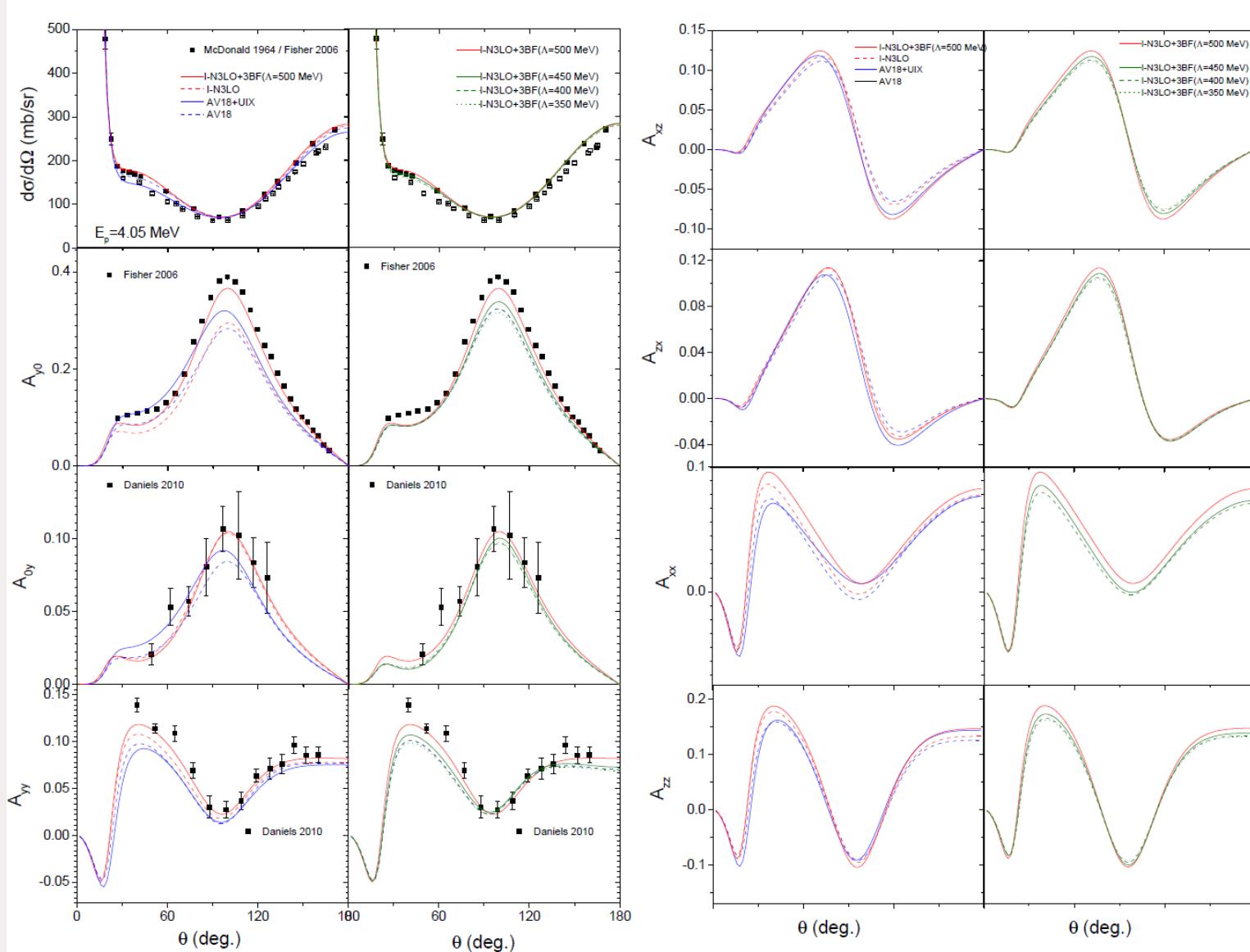
# 4N problem: $n-^3H$ elastic scattering



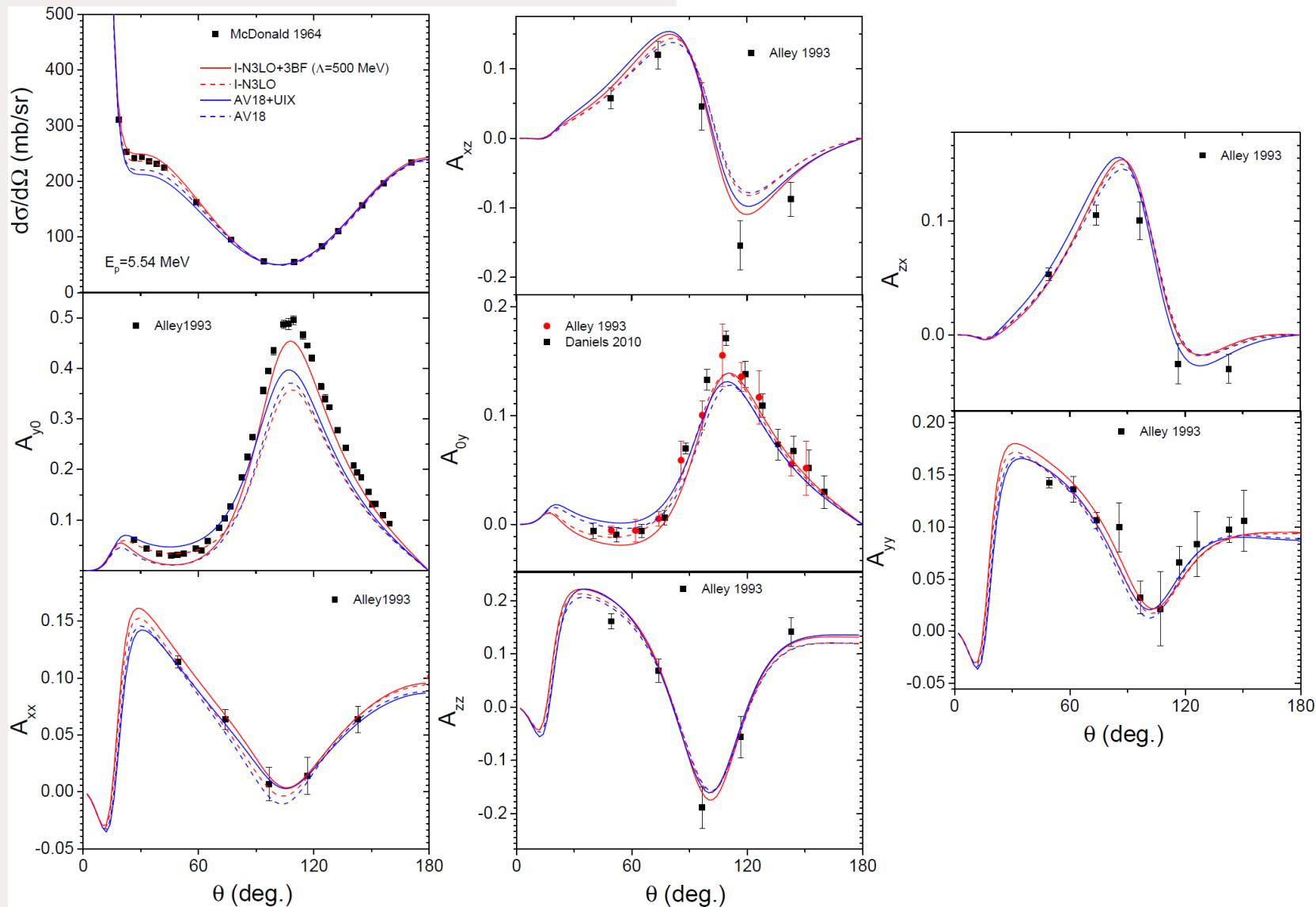
# 4N problem: $n-^3H$ elastic scattering



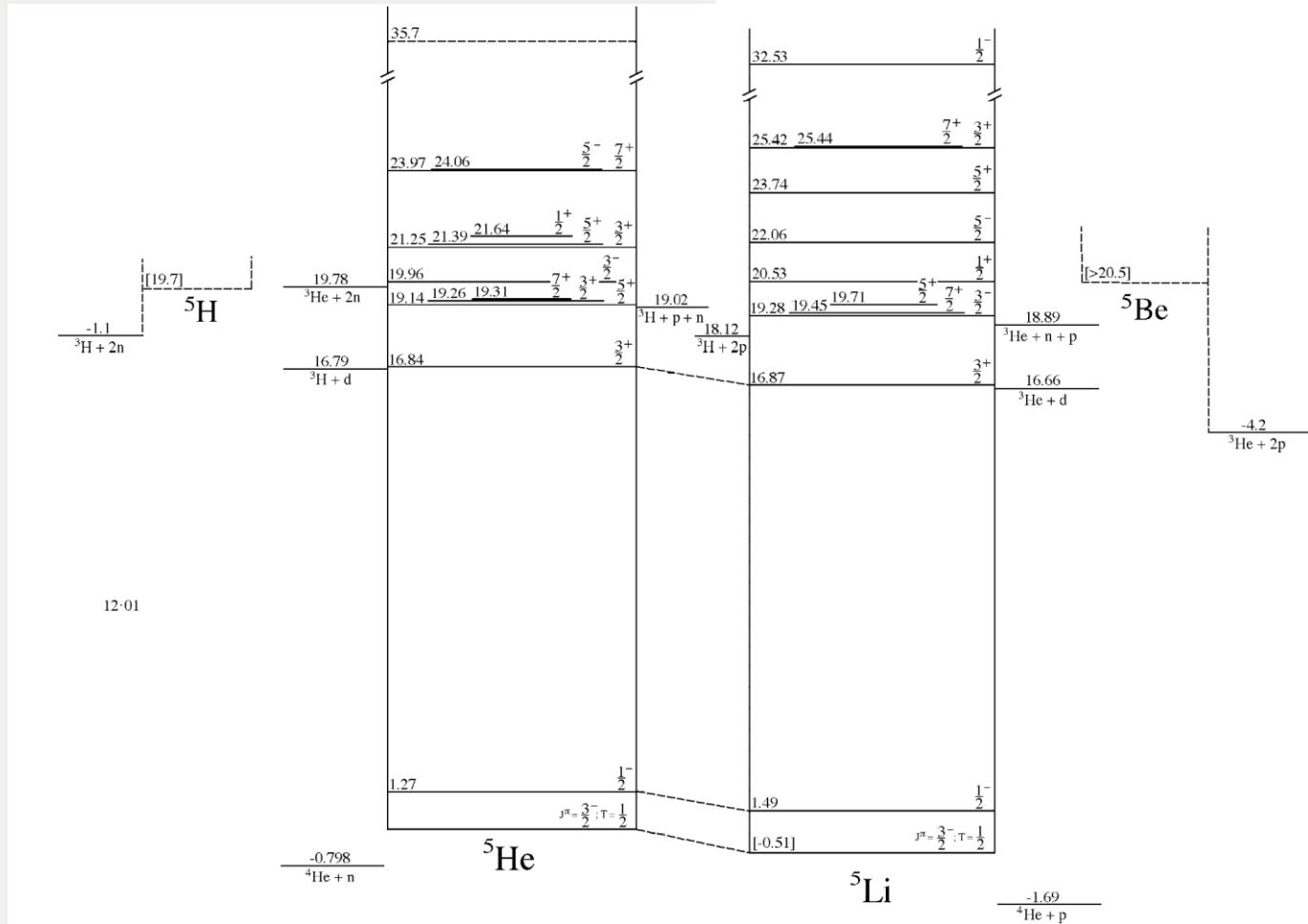
# 4N problem: p- $^3$ He elastic scattering



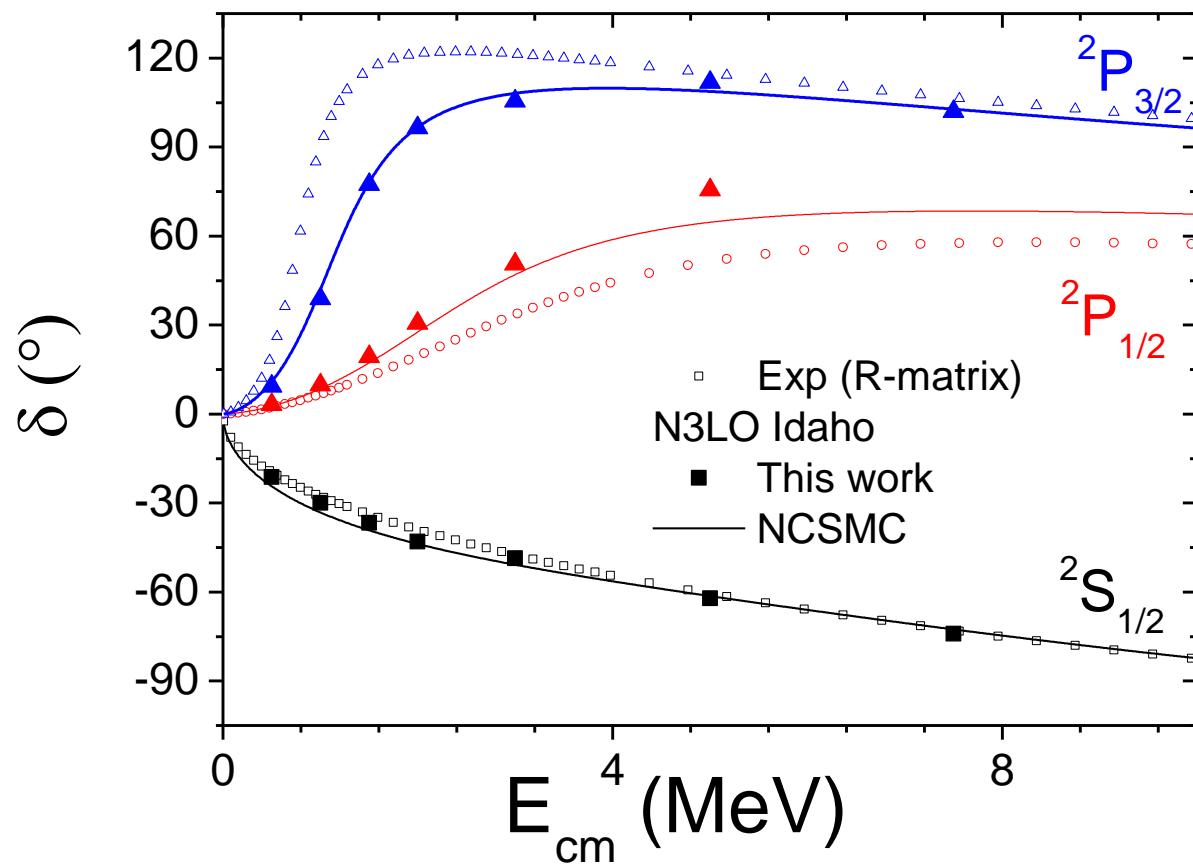
# 4N problem: p- $^3$ He elastic scattering



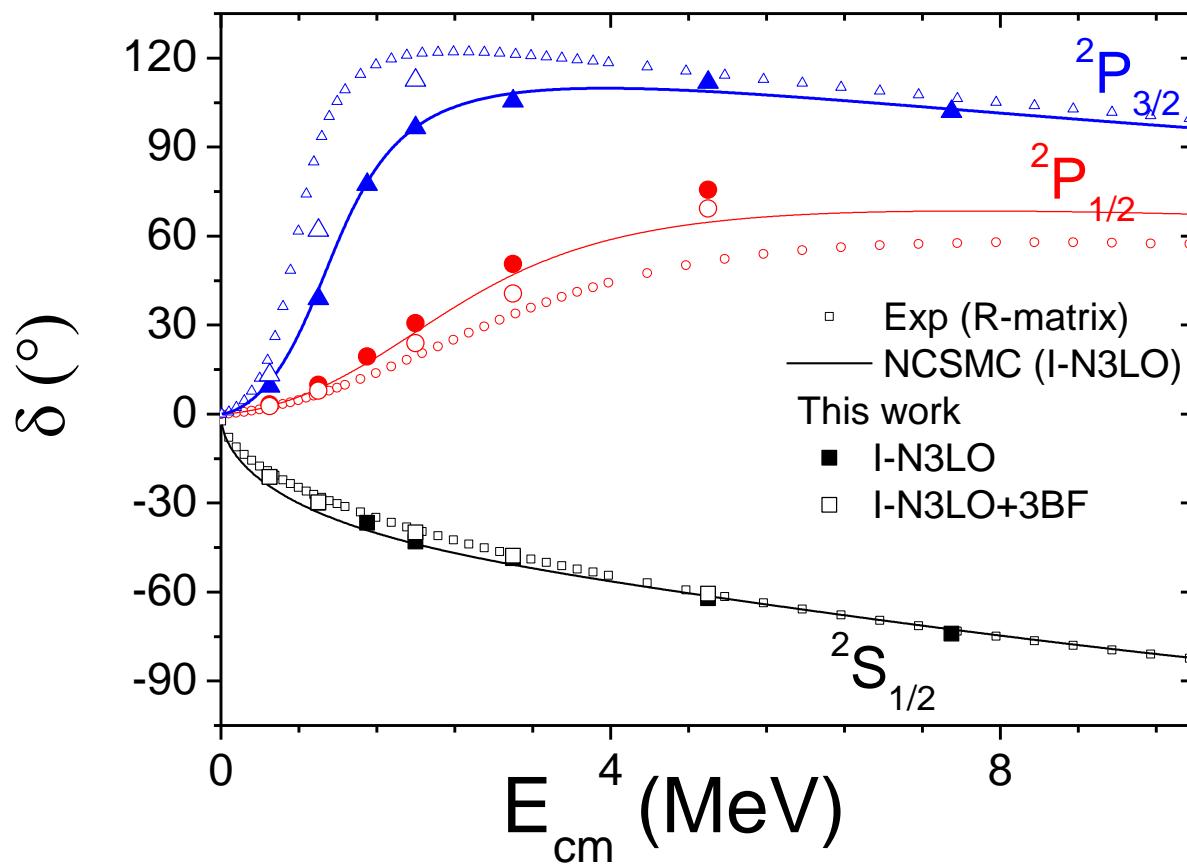
# 5N system



# $n-{}^4He$ scattering

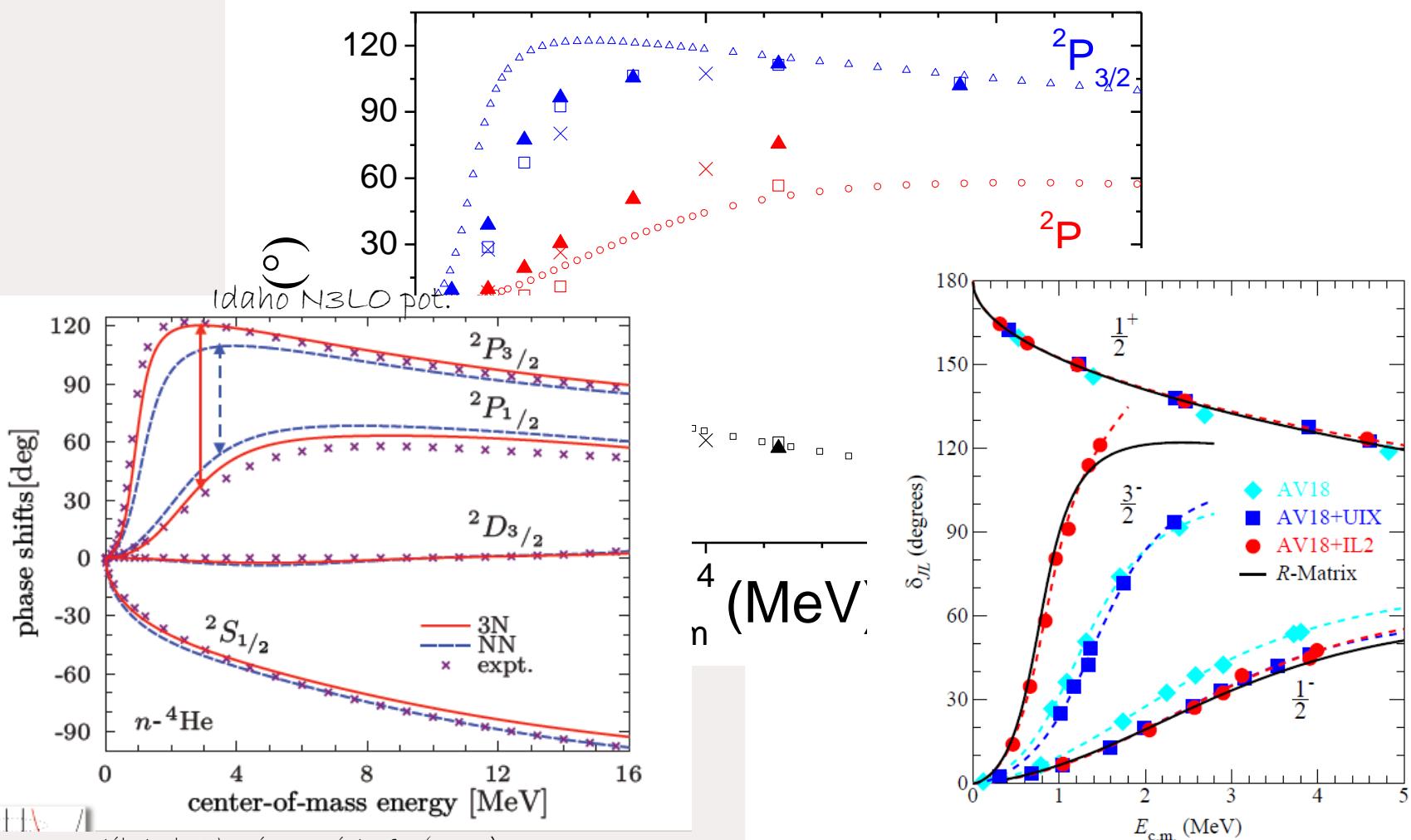


NCSMC: P. Navratil et al., Physica Scripta **91** (2016) 053002



NCSMC: P. Navratil et al., Physica Scripta **91** (2016) 053002

# $n$ - ${}^4\text{He}$ scattering



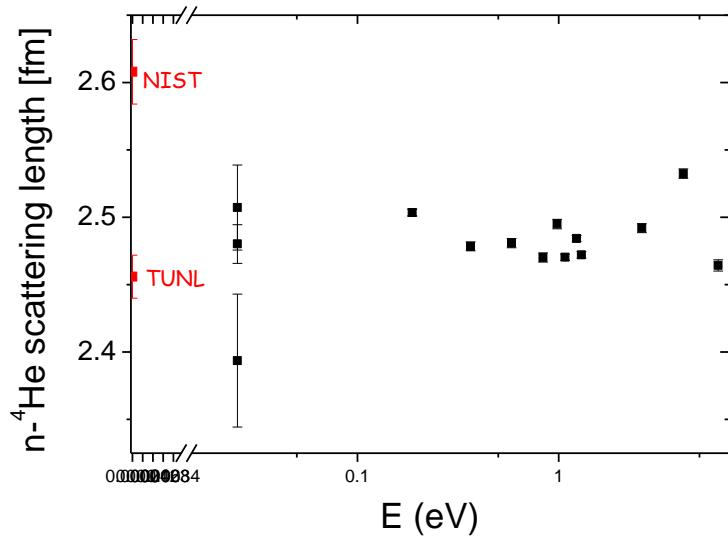
P. Navratil et al., Physica Scripta 91 (2016) 053002

K.M. Nollett et. al., Phys. Rev. Lett. 99:022502, 2007

# Case of little interest: S-wave

Experimental n- ${}^4\text{He}$  scattering length ...

nothing should be as easy to measure...



NIST (Neutron News 3, 1992)

	Coh a (fm)	Inc b (fm)
${}^1\text{H}$	-3.7406(11) -3.79406(11)	25.274(9)
${}^2\text{H}$	6.671(4)	4.04(3)
${}^3\text{H}$	4.792(27)	-1.04(17)
${}^3\text{He}$	$5.74(7)-1.483(2)i$	$-2.5(6)+2.568(3)i$
${}^4\text{He}$	3.26(3)	

TUNL: D.R. Tilley et al., Nucl. Phys. A 708 (2002) 3

NIST: <https://www.ncnr.nist.gov>

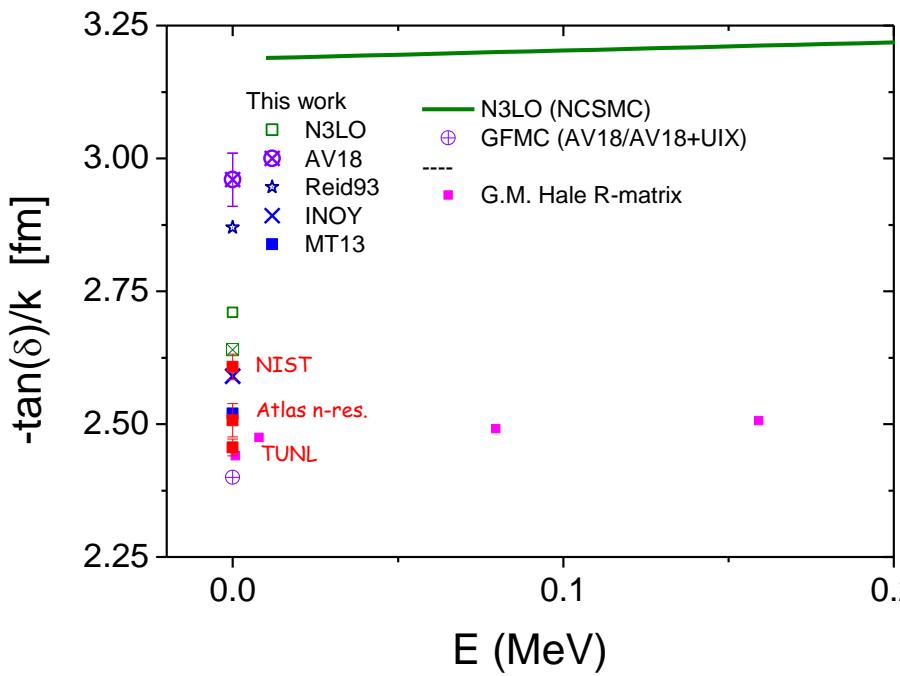
Experimental data:

D.C.Rorer et al., Nucl. Phys. A 133 (1969) 410

S.F.Mughabghab, Atlas of Neutron Resonances (2006)

R.Genin et al., Journal de Physique 24 (1963) 21

# Case of little interest: S-wave



TUNL: D.R. Tilley et al., Nucl. Phys. **A708** (2002) 3

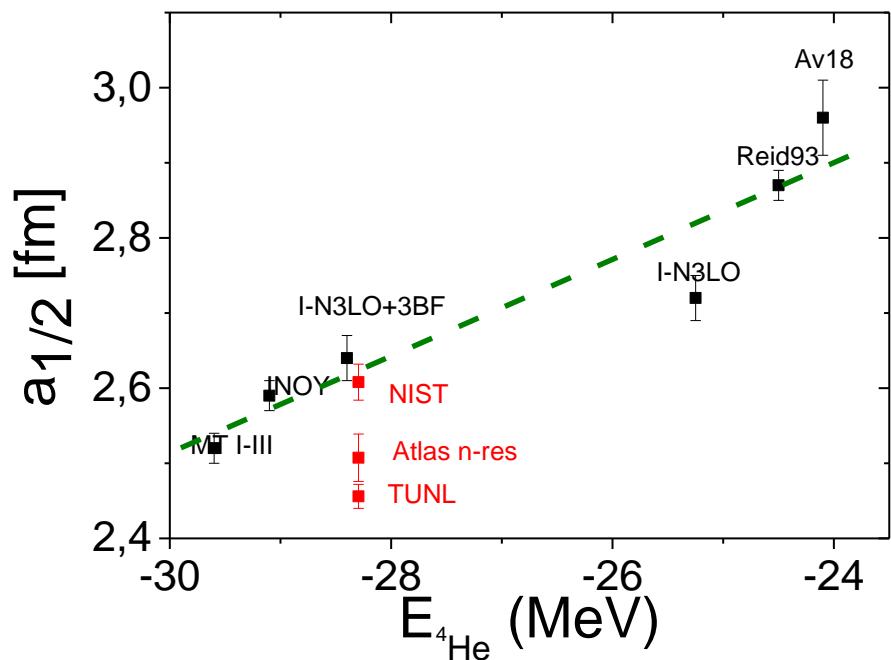
NIST: <https://www.ncnr.nist.gov>

S. Ali PSA: S. Ali et al., Rev. Mod. Phys. **57** (1985) 923

Bang-Gignoux pot: J. Bang, C. Gignoux, Nucl. Phys. A **313** (1979)

NCSMC: P. Navratil et al., Physica Scripta **91** (2016) 053002

GFMC: K.M. Nollett, PRL **99**, 022502 (2007)



# PV violation for $\vec{n}$ - ${}^4\text{He}$

Slow  $\vec{n}$  spin rotation studyt at NIST  
 E. Swanson et al. PRC **100** (2019) 015204

✓ Weak process  $V^{\text{weak}} \ll V^{\text{strong}}$

1<sup>st</sup> order perturbation:

$$R_{f \leftarrow i}^{\text{weak}} \propto \langle \Psi_f^{\text{strong}} | V^{\text{weak}} | \Psi_i^{\text{strong}} \rangle$$

✓ The last expression one may calculate within FY framework, without passing directly to total system's wave function

R. Lazauskas, Y.H. Song, PRC **99** (2019) 054002

Input:  $\frac{V^{\text{strong}}}{V^{\text{weak}}} = 1 - \text{N3LO+3BF}$   
 DDH meson exchange pot.  $(\pi, \rho, \omega, \rho')$ . B. Desplanques et al, Ann.Phys. **124** (1980) 449.

ultracold  $\vec{n}$ - ${}^4\text{He}$  spin rotation angle in  $10^{-7}\text{ rad/m}$ :

$$\frac{d\phi}{dz} = -0.144(1)h_\pi^1 + 0.058(8)h_\omega^0 - 0.402(1)h_\rho^0 + \mathbf{0.0298}h_\omega^1 + \mathbf{0.0296}h_\rho^1 + \mathbf{0.0061}h_\rho^1,$$

$$\frac{d\phi}{dz} = \begin{cases} 3.7 & \text{DDH-best} \\ 3.0 & \text{DZ} \\ 0.8 & \text{FCDH} \\ 12. & \text{large } N_c \end{cases}$$

R. Lazauskas, Y.H. Song, PRC **99** (2019) 054002.

S. Gardner et al., Ann. Rev. Nucl. Part. Sci. **67** (2017) 69

# $^5\text{H}$ resonances: experiment

TABLE I. Summary of experimental results for  $^5\text{H}$ . Resonance energies are given relative to  $^3\text{H} + 2n$ .

Reference	Reaction	Detected	$E_R$ (MeV)	$\Gamma$ (MeV)	$E_{beam}$ (A MeV)
[17]	$^3\text{H}(t, p)^5\text{H}$	$p$	$\approx 1.8$	$\approx 1.5$	7.42
[18]	$^6\text{He}(p, 2p)^5\text{H}$	$2p$	$1.7 \pm 0.3$	$1.9 \pm 0.4$	36
[19]	$^3\text{H}(t, p)^5\text{H}$	$t, p, n$	$1.8 \pm 0.1$	$< 0.5$	19.2
[21]	$^3\text{H}(t, p)^5\text{H}$	$t, p, n$	$\approx 2$	—	19.2
[22]	$^3\text{H}(t, p)^5\text{H}$	$t, p, n$	$\approx 2$	$\approx 1.3$	19.2
[24]	$^6\text{He}(^{12}\text{C}, X + 2n)^5\text{H}$	$t, 2n$	$\approx 3$	$\approx 6$	240
[25]	$^6\text{He}(d, ^3\text{He})^5\text{H}$	$^3\text{He}, t$	$1.8 \pm 0.1$	$< 0.6$	22
[26]	$^6\text{He}(d, ^3\text{He})^5\text{H}$	$^3\text{He}, t$	$1.8 \pm 0.2$	$1.3 \pm 0.5$	22
[27]	$^6\text{He}(d, ^3\text{He})^5\text{H}$	$^3\text{He}, t$	$1.7 \pm 0.3$	$\approx 2.5$	22
[28]	$^9\text{Be}(\pi^-, pt)^5\text{H}$	$p, t$	$5.2 \pm 0.3$	$5.5 \pm 0.5$	$E_\pi < 30$ MeV
[28]	$^9\text{Be}(\pi^-, dd)^5\text{H}$	$p, t$	$6.1 \pm 0.4$	$4.5 \pm 1.2$	$E_\pi < 30$ MeV

- [17] P. G. Young, Richard H. Stokes, and Gerhard Röder, *Phys. Rev.* **173**, 949 (1968).
- [18] A. A. Korsheninnikov, M. S. Golovkov, A. S. Fomichev, S. I. Sidorchuk, S. Chelnokov, V. A. Gorshkov, D. D. Bogdan Ter-Akopian *et al.*, *Phys. Rev. Lett.* **87**, 092501 (2001).
- [19] M. S. Golovkov, Yu. Ts. Oganessian, D. Fomichev, A. M. Rodin, S. I. Sidorchuk, S. Stepansov, G. M. Ter-Akopian, R. Wolski, *Nature* **566**, 70 (2003).
- [21] M. S. Golovkov, L. V. Grigorenko, A. Krupko, Yu. Ts. Oganessian, A. M. Rodin, R. S. Slepnev, S. V. Stepansov, G. M. Ter-Akopian, *Phys. Rev. Lett.* **93**, 262501 (2004).
- [22] M. S. Golovkov, L. V. Grigorenko, A. Krupko, Yu. Ts. Oganessian, A. M. Rodin, R. S. Slepnev, S. V. Stepansov, G. M. Ter-Akopian, *Phys. Rev. C* **72**, 064612 (2005).

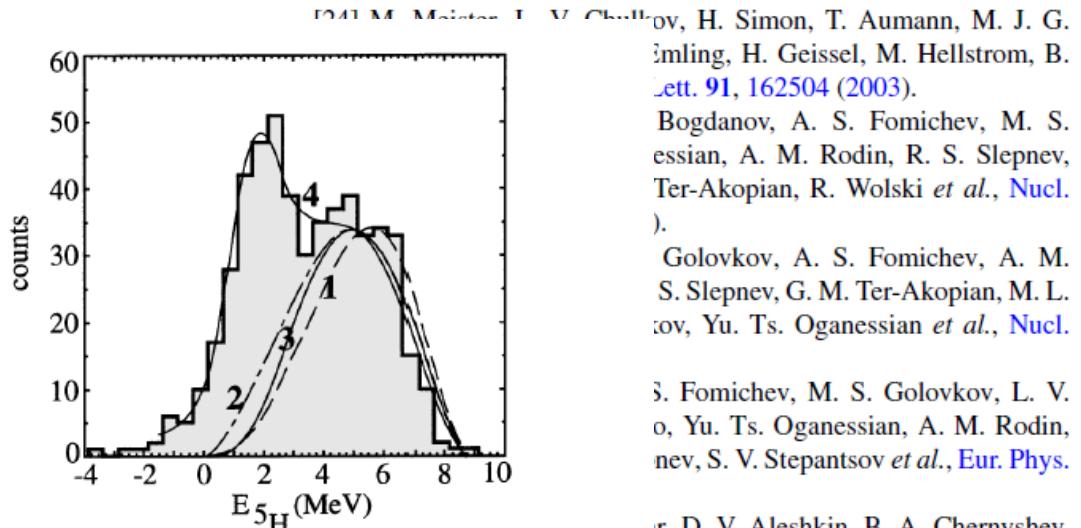


FIG. 3. Spectrum of  $^5\text{H}$  from the reaction  $p(^6\text{He}, ppt)$ . Curves show calculations explained in the text.

# $^5\text{H}$ resonances?

TABLE II. Summary of some theoretical results for  $^5\text{H}$ . Resonance energies are given relative to  $^3\text{H} + 2n$ .

Reference	Method	$E_R$ (MeV)	$\Gamma$ (MeV)
[7]	Cluster, model with source	2–3	4–6
[23]	Three-body cluster	2.5–3	3–4
[31,35]	Cluster, $J$ -matrix, resonating group model	1.39	1.60
[36]	Cluster, complex scaling adiabatic expansion	1.57	1.53
[32]	Cluster, generator coordinate method	$\approx 3$	$\approx 1$ –4
[33]	Cluster, complex scaling	1.59	2.48
[34]	Cluster, analytic coupling in continuum constant	$1.9 \pm 0.2$	$0.6 \pm 0.2$

[7] L. V. Grigorenko, N. K. Timofeyuk, and M. V. Zhukov, *Eur. Phys. J. A* **19**, 187 (2004).

[31] A. V. Nesterov, F. Arickx, J. Broeckhove, and V. S. Vasilevsky, *Phys. Part. Nucl.* **41**, 716 (2010).

[32] P. Descouvemont and A. Kharbach, *Phys. Rev. C* **63**, 027001 (2001).

[33] K. Arai, *Phys. Rev. C* **68**, 034303 (2003).

[34] A. Adachour and P. Descouvemont, *Nucl. Phys. A* **813**, 252 (2008).

[35] J. Broeckhove, F. Arickx, P. Hellinckx, V. S. Vasilevsky, and A. V. Nesterov, *J. Phys. G* **34**, 1955 (2007).

[36] R. de Diego, E. Garrido, D. V. Fedorov, and A. S. Jensen, *Nucl. Phys. A* **786**, 71 (2007).

Predictivity?

$^5\text{H}$  states does not appear naturally

Cluster models, involving approximations for 5-body dynamics

- $^3\text{H}+n+n$  models: without  $n$ -antisymmetrization between the core & valence
- $^3\text{H}+n+n$  models: including  $n$ -antisymmetrization, however by freezing  $^3\text{H}$  core

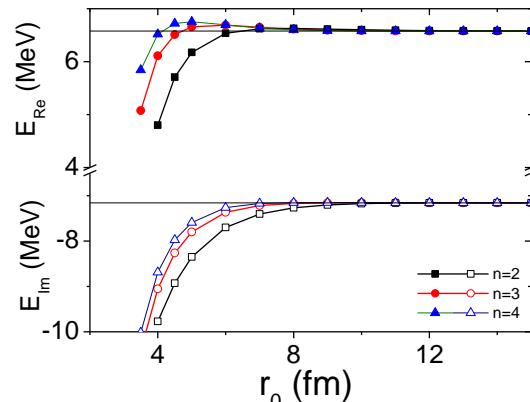
# $^5\text{H}$ resonances?

## How to handle resonances?

- ACCC : **Analytic continuation in the coupling constant** method (V.I. Kukulin et al., « Theory of resonances », Kluwer AP 1989)
  - Artificially bind  $^5\text{H}$  with some additional potential  $V = \lambda V_0$  (we use 5-body pot not to affect  $^3\text{H}$  threshold!!)
  - Study  $B_{^5\text{H}}(\lambda)$  and determine  $\lambda_0$  such that  $B_{^5\text{H}}(\lambda_0) = B_{^3\text{H}}$
  - Smartly extrapolate  $B_{^5\text{H}}(\lambda) = f(\lambda - \lambda_0)$  to determine  $E_{^5\text{H}} = B_{^5\text{H}}(0)$
- « Dirty » smooth exterior complex scaling method (DEXCSM)
 

B. Simon. Phys. Letters A, 71 (1979) 211

  - Choose sharp transformation function, which almost does not affect  $r$  in  $r < r_0$
  - Fix  $r_0$  beyond the physical interaction region
  - Ignore inconsistencies in transformation between different Jacobi bases



2b-example

$$V(l=1) = \frac{1.8}{r} (1438.72e^{-3.11r} - 626.885e^{-1.55r})$$

$$r \rightarrow (1 - f(r))r + f(r)r e^{i\theta} \quad f(r) = \exp\left(-\left(\frac{r_0}{r}\right)^n\right)$$

# ${}^5\text{H}(\text{J}=1/2^+)$

- $\text{nn}$  interaction described by the MT I-III potential
- auxilliary potential for ACCC

$$V_{5b}(\rho) = \lambda \rho^p \exp(-\rho^2/\rho_0^2).$$

$$\rho^2 = x^2 + y^2 + z^2 + w^2 = 2 \sum_{i=1}^5 r_i^2$$

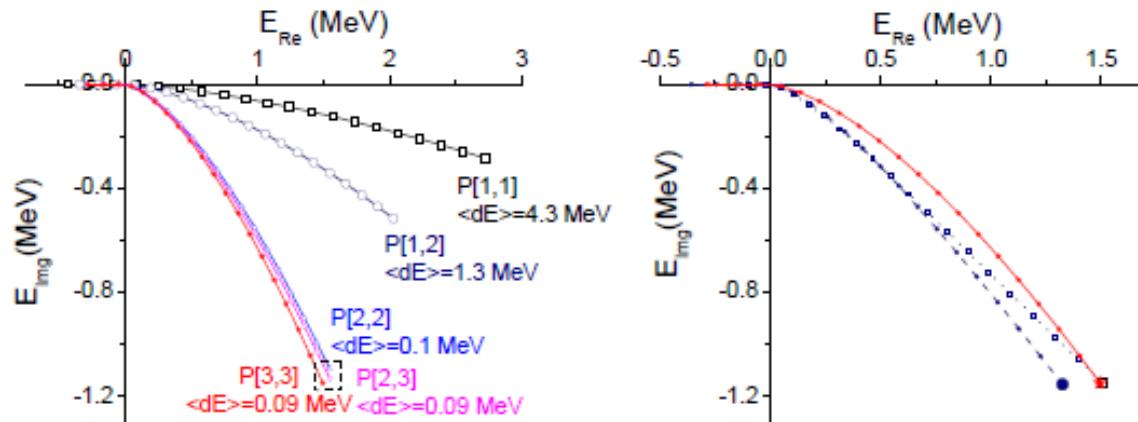


Fig. 3 Resonance trajectories for a  $J^\pi = 1/2^+$  state of  ${}^5\text{H}$  with respect to  ${}^3\text{H}$  threshold. Each trajectory is split by points in 20 intervals of equal step in  $\lambda$ , starting at the position where  ${}^5\text{H}$  nucleus is still weakly bound. The endpoint of the trajectory indicates extrapolated value for the bare NN interaction, corresponding  $\lambda = 0$  case. In the left panel convergence of the results with respect to order of Padé expansion is presented; calculation is based on auxiliary potential defined in eq. (13) with  $\rho_0^2 = 78.4 \text{ fm}^2$  and  $p = 0$ . In the right panel converged results for three different external potentials are presented.

$$\mathcal{E}({}^5\text{H}) - \mathcal{E}({}^3\text{H}) = 1.4(1) - i 1.2(1)$$

$$\mathcal{E}({}^5\text{H}) - \mathcal{E}({}^3\text{H}) = 1.7(2) - i 1.2(1)$$

# $^5\text{H}(\text{J}=1/2^+)$

- nn interaction described by the MT I-III potential

$$\mathbf{J=1/2^+ (L=0^+, S=1/2)}$$

ACCC:

$$E(^5\text{H}) - E(^3\text{H}) = 1.4(1) - i1.2(1)$$

DEXCSM:

$$E(^5\text{H}) - E(^3\text{H}) = 1.6(2) - i1.2(1)$$

$$\mathbf{J=5/2^+ (L=2^+, S=1/2)}$$

DEXCSM:

$$E(^5\text{H}) - E(^3\text{H}) = 2.50(15) - i1.90(15)$$

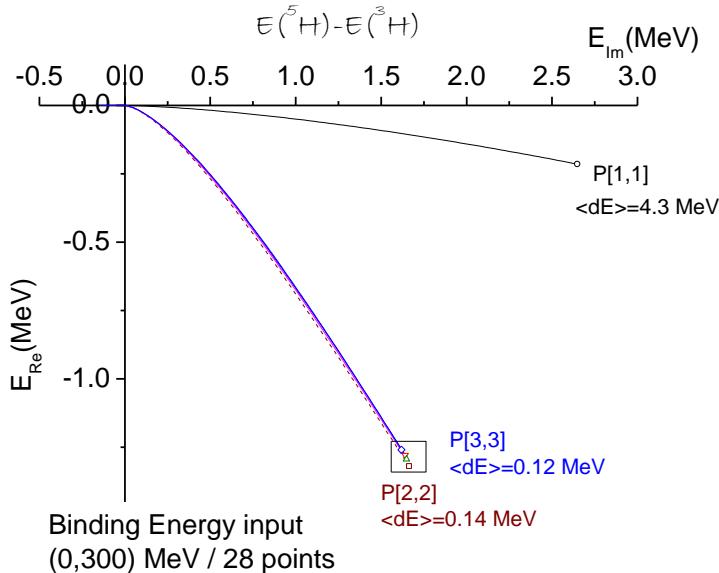
**Negative parity states & ones with  $S=3/2$  are much more broader**

To compare with  $^4\text{H}$  resonances:

$$E(^4\text{H}) - E(^3\text{H}) = \begin{cases} 1.08(1) - i2.04(2) & (S=1, L=1^-) \\ 0.88(3) - i2.20(4) & (S=0, L=1^-) \end{cases}$$

# $^5\text{H}(\text{J}=1/2^+)$

INOY Potential



$$E(\text{H}) - E(\text{H}_3) = 1.65(5) - i1.26(6)$$

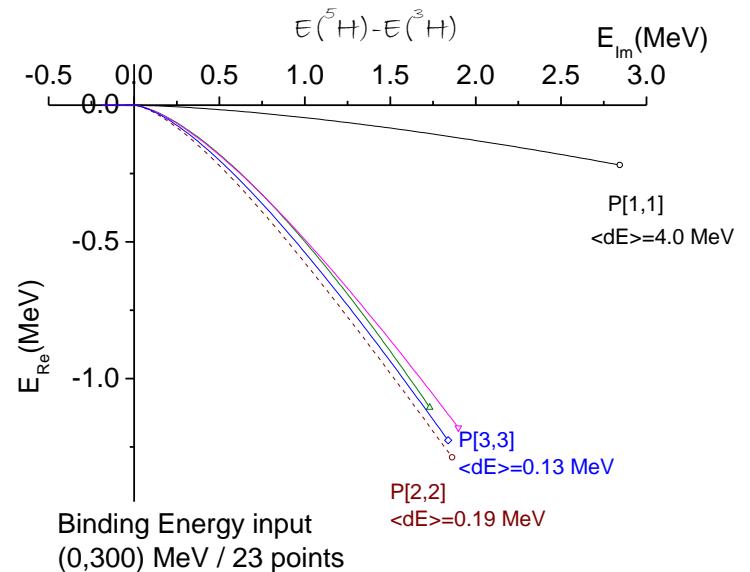
DEXCSM:  $1.8(1) - i1.2(1)$

To compare with  $^4\text{H}$  resonances  $J=2^-$ :

$$E(\text{H}) - E(\text{H}_3) = 1.31(3) - 2.08(2)$$

R. Lazauskas, E. Hiyama, J. Carbonell, Phys. Lett. B 791 (2019) 335

$N_3$ LO Potential



$$E(\text{H}) - E(\text{H}_3) = 1.8(1) - i1.15(15)$$

DEXCSM:  $1.85(10) - i1.20(5)$

$$E(\text{H}) - E(\text{H}_3) = 1.17(3) - 1.99(3)$$

# Conclusion

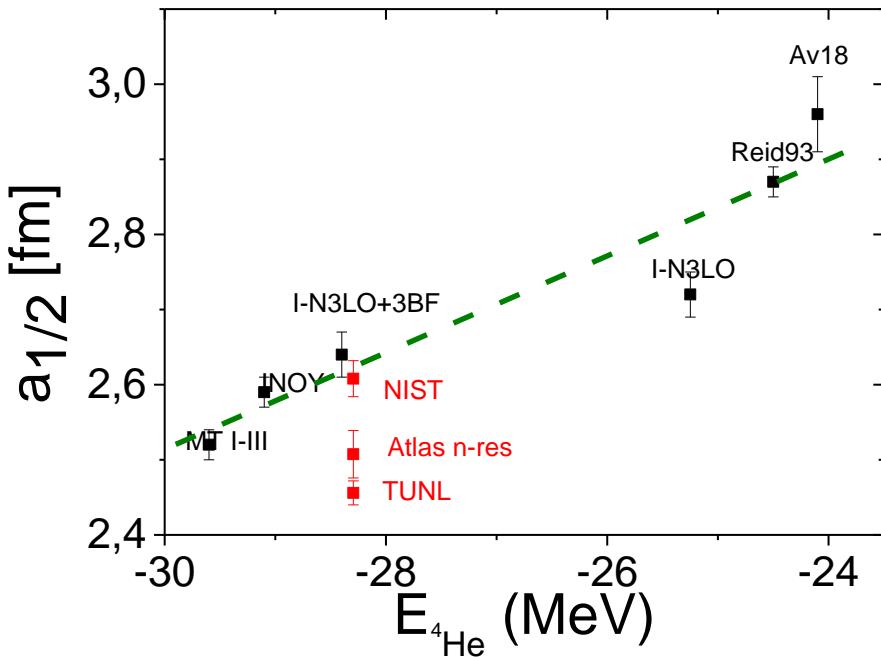
- FY eq. formalism remains reference in few-body scattering calculations. The first solutions of 5-body FY equations are presented.
- Reliable results have been obtained for n- $^4\text{He}$  scattering at low energies using realistic interactions. Satisfactory description is obtained when using Idaho N3LO NN +N2LO NNN interactions.
- The first fully realistic calculation of weak process in 5-nucleon sector is performed.
- Description of the  ${}^5\text{H}$  resonant states have been performed for the first time using fully realistic description and two different methods to calculate resonance positions. Presence of broad resonant states is confirmed!

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# $n$ - ${}^4He$ scattering

Experimental  $n$ - ${}^4He$  scattering length ...

nothing should be as easy to measure...



TUNL: D.R. Tilley et al., Nucl. Phys. **A708** (2002) 3

NIST: <https://www.ncnr.nist.gov>

Experimental data:

D.C. Rorer et al., Nucl. Phys. **A133** (1969) 410

S.F. Mughabghab, Atlas of Neutron Resonances (2006)

R. Genin et al., Journal de Physique **24** (1963) 21

NIST (Neutron News 3, 1992)

	Coh a (fm)	Inc b (fm)
${}^1H$	-3.7406(11) -3.79406(11)	25.274(9)
${}^2H$	6.671(4)	4.04(3)
${}^3H$	4.792(27)	-1.04(17)
${}^3He$	5.74(7)-1.483(2) <i>i</i>	-2.5(6)+2.568(3) <i>i</i>
${}^4He$	3.26(3)	

# $n-{}^4He$ scattering

