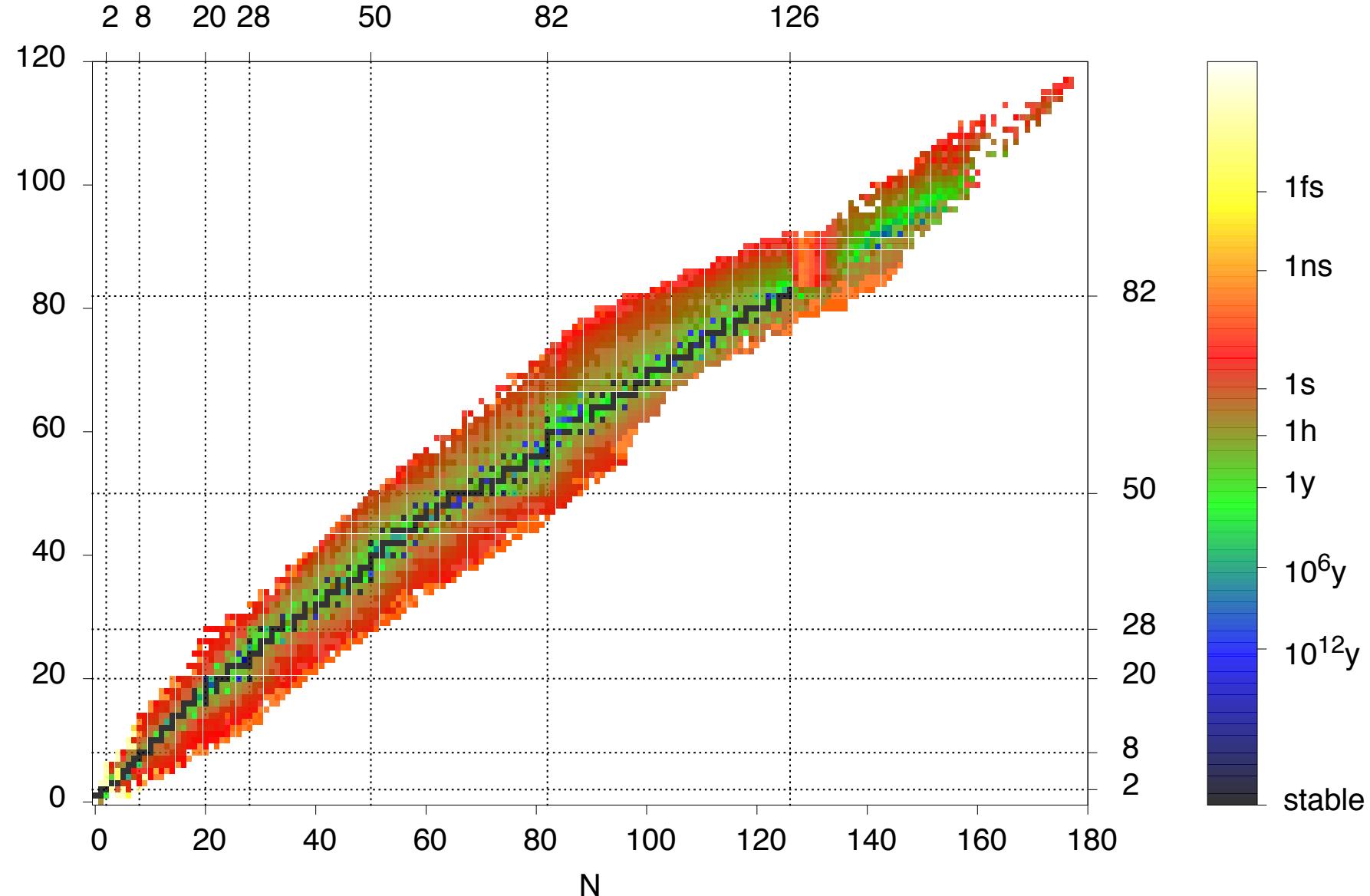


# *Coherent and chaotic dynamics of open quantum systems*

**Alexander Volya**  
Florida State University

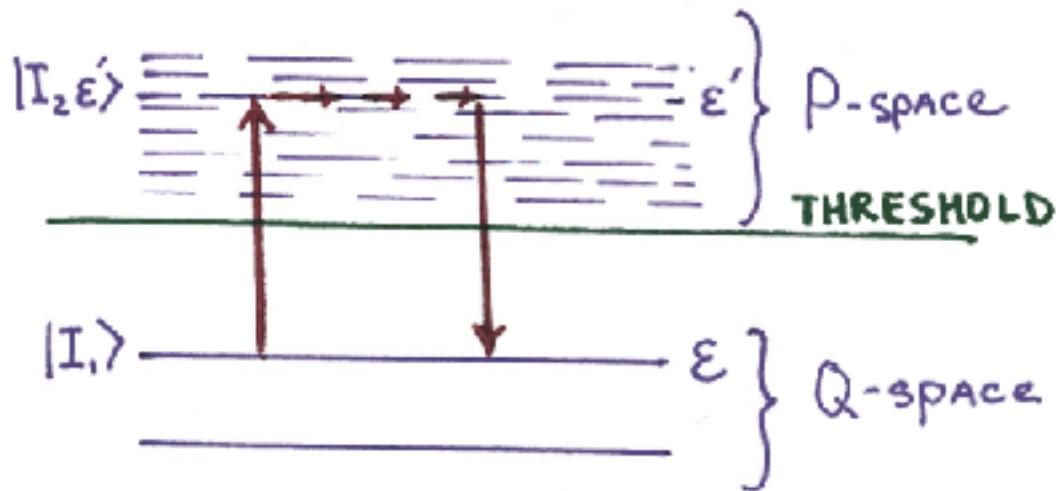
# Atomic nuclei are open quantum systems



# Topics for discussion

- Physics of coupling to continuum
  - Effective Non-Hermitian Hamiltonian formalism
  - Time dependent approach
- Features of open systems
  - Virtual excitations into continuum
  - Resonances and direct decay
  - Superradiance, alignment of structure
- Decay collectivity and intrinsic collectivities.
- Related questions

# Physics of coupling to continuum



## The role of continuum-coupling

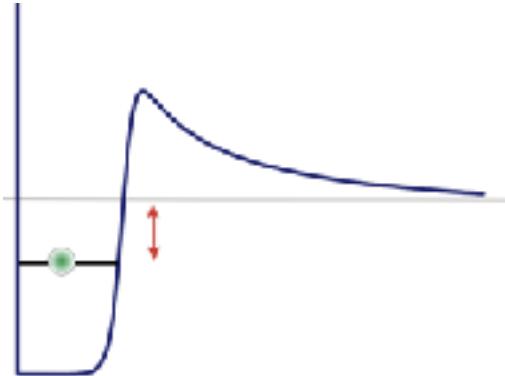
$$H'(\epsilon) = \int_0^{\infty} d\epsilon' A^*(\epsilon') \frac{1}{\epsilon - \epsilon' + i0} A(\epsilon') \quad A(\epsilon') \equiv \langle I_2, \epsilon' | H_{PQ} | I_1 \rangle$$

# Physics of coupling to continuum

$$H'(\epsilon) = \int_0^\infty d\epsilon' \frac{|A(\epsilon')|^2}{\epsilon - \epsilon' + i0}.$$

Integration region involves no poles

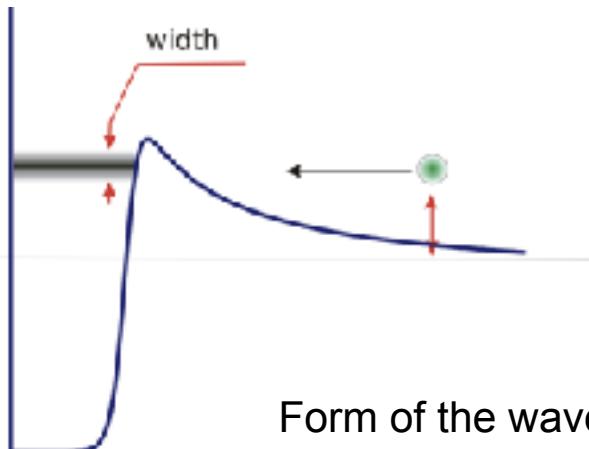
$$H'(\epsilon) = \Delta(\epsilon) \quad \Delta(\epsilon) = \int d\epsilon' \frac{|A(\epsilon')|^2}{\epsilon - \epsilon' + i0}$$



State embedded in the continuum

$$\frac{1}{x \pm i0} = \text{p.v.} \frac{1}{x} \mp i\pi\delta(x)$$

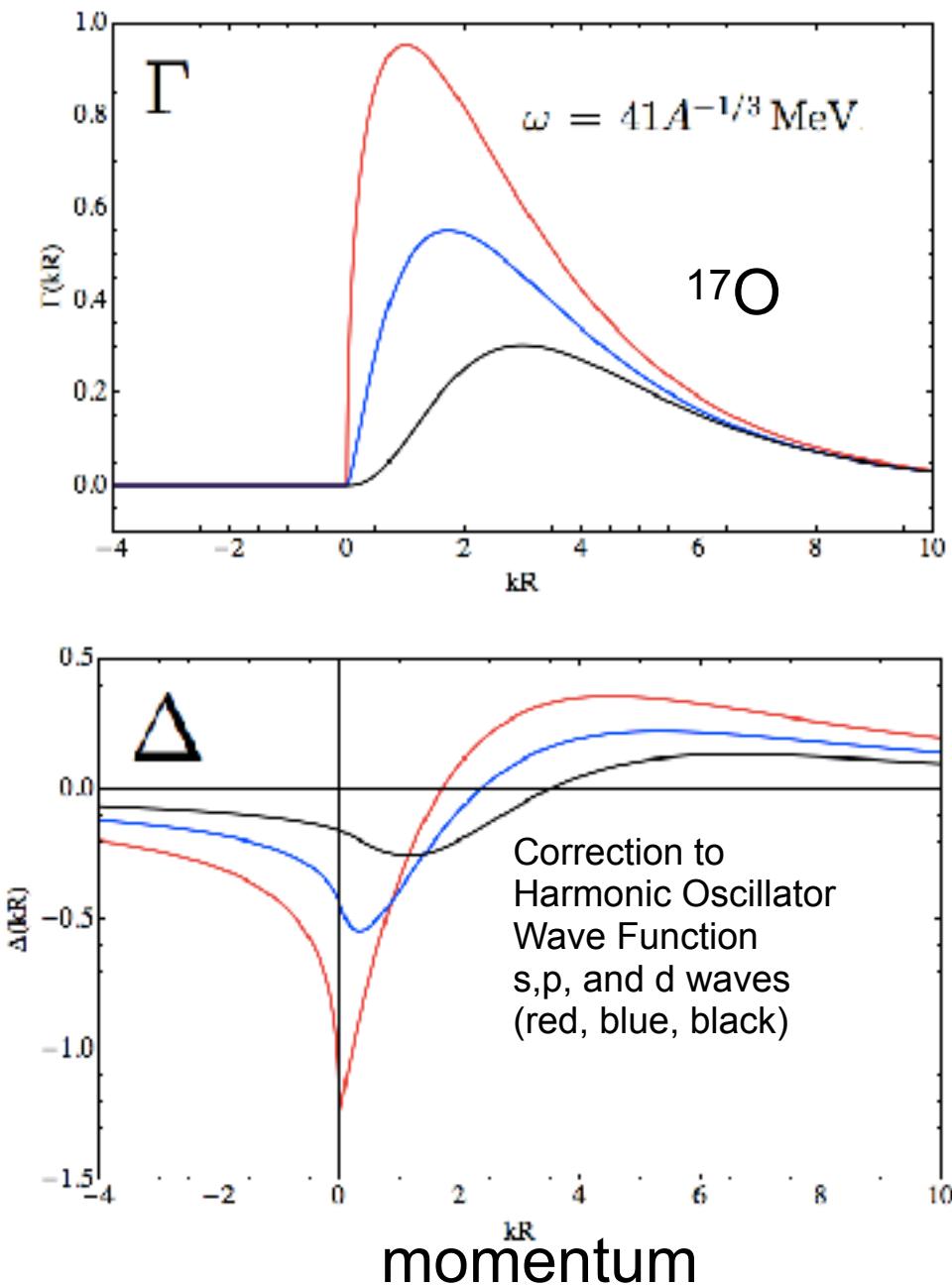
$$H'(\epsilon) = \Delta(\epsilon) - \frac{i}{2}\Gamma(\epsilon) \quad \Gamma(\epsilon) = 2\pi |A(\epsilon)|^2$$



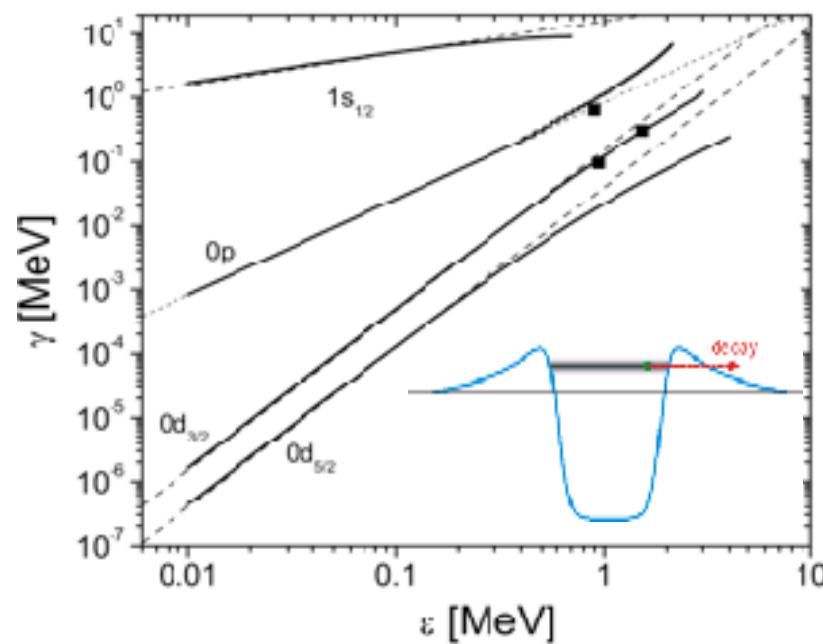
Form of the wave function and probability

$$|\exp(-iEt)|^2 = 1 \rightarrow |\exp(-iEt - \Gamma t/2)|^2 = \exp(-\Gamma t)$$

# Self energy, interaction with continuum

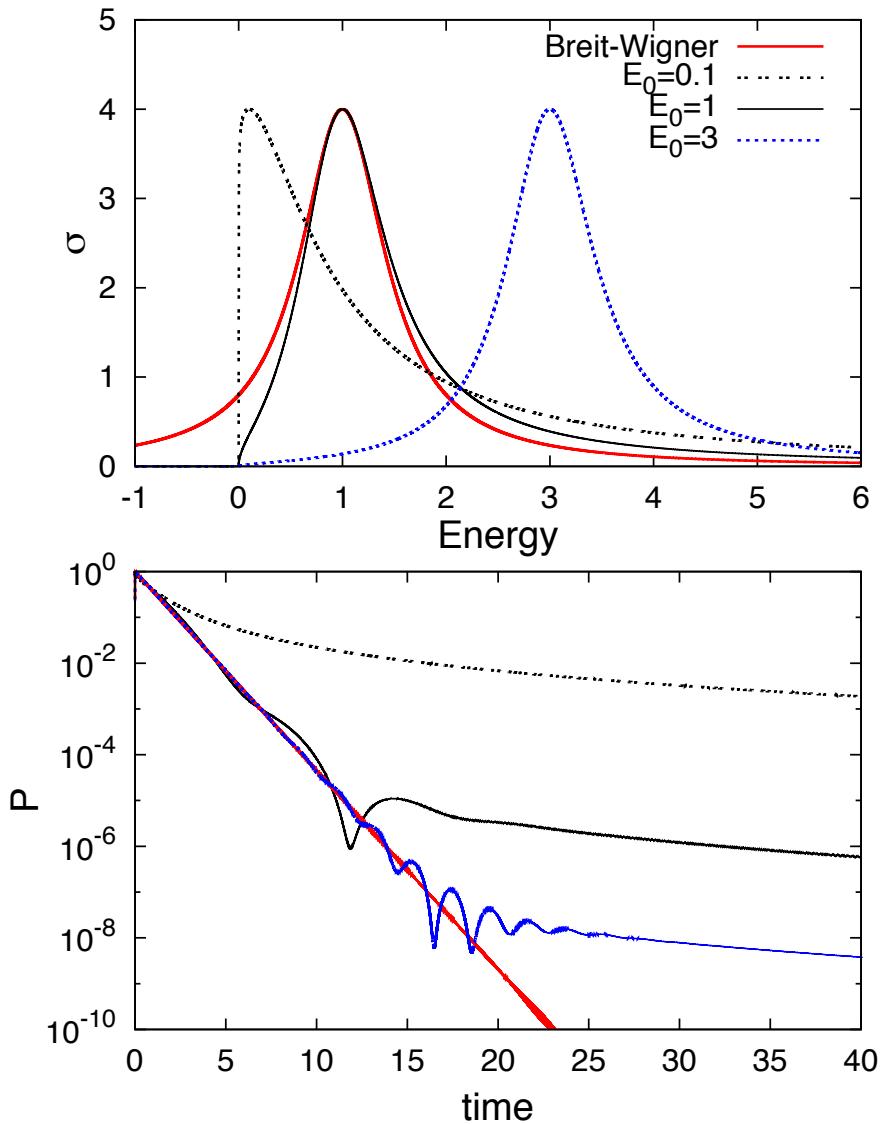


## Width as a function of energy



D. Abrahamsen, A. Volya, and I. Wiedenhoever,  
*Effective R-matrix parameters of the Woods-Saxon nuclear potential*, APS Volume 57, Number 16, section KA 26 (2012).

# Time-dependent picture



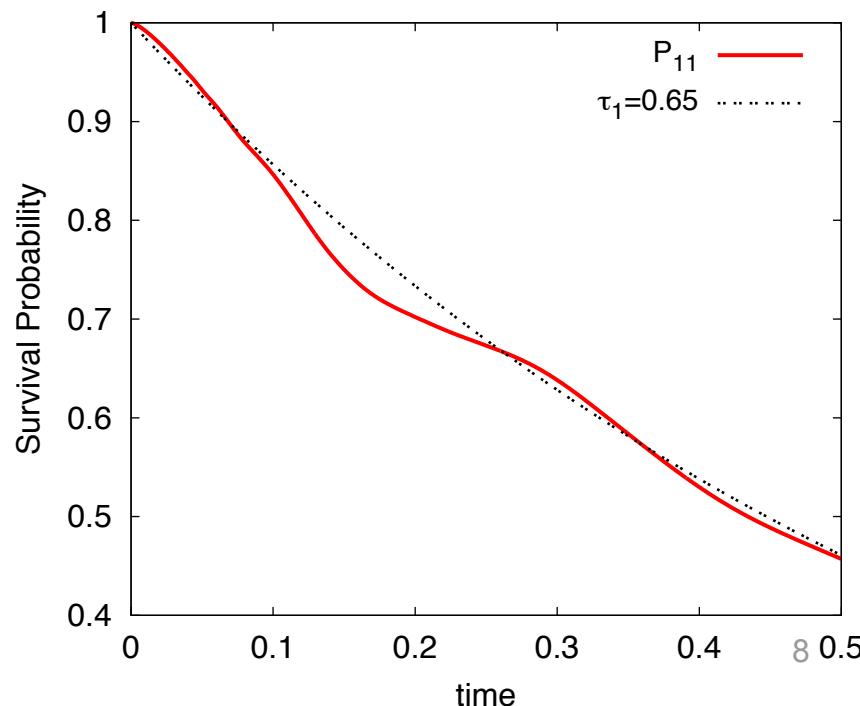
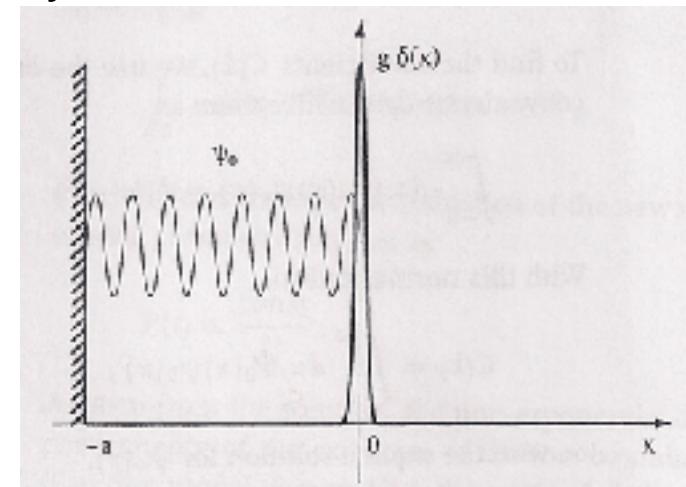
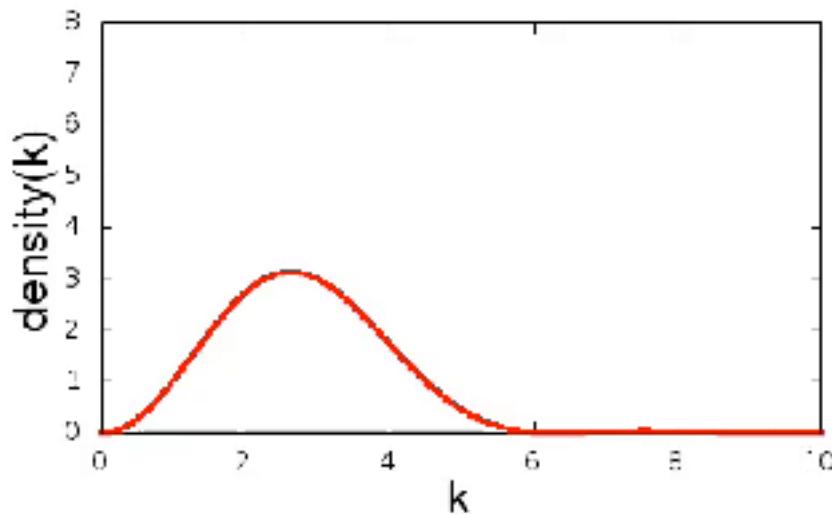
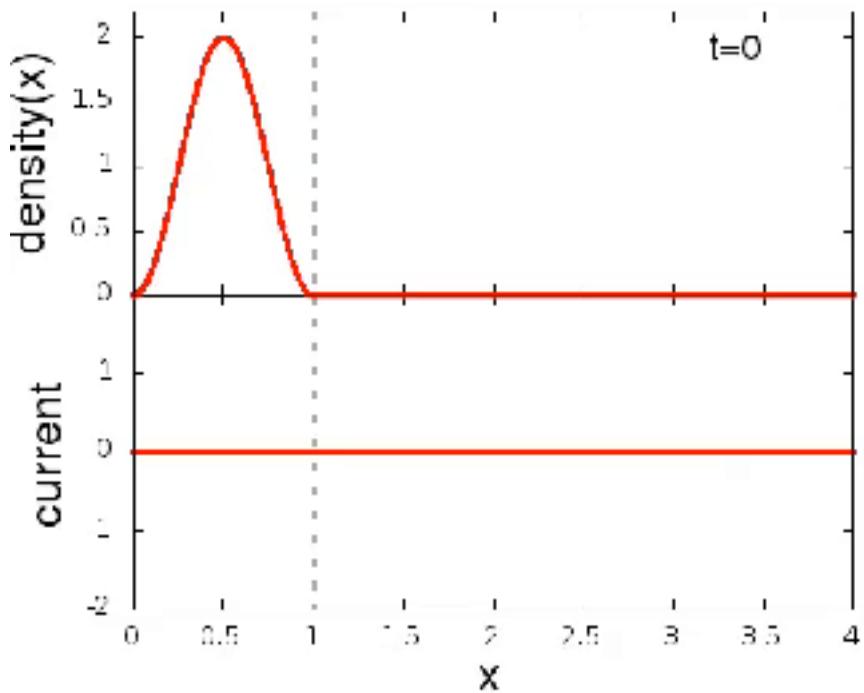
$$G = \frac{1}{E - E_o + i/2 \Gamma(E)}$$

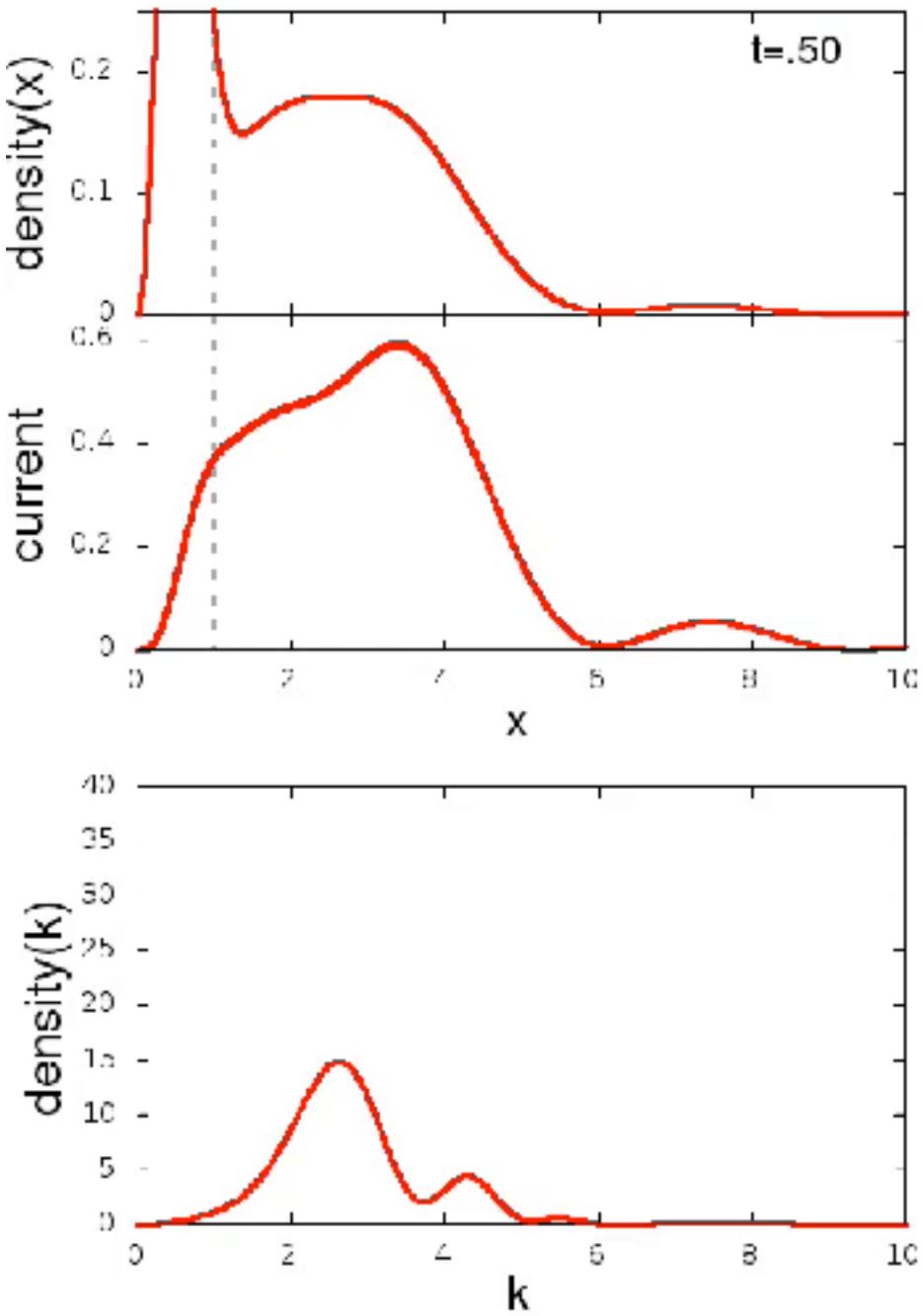
$$\Gamma(E) \propto \sqrt{E}$$

Power-law remote decay rate!

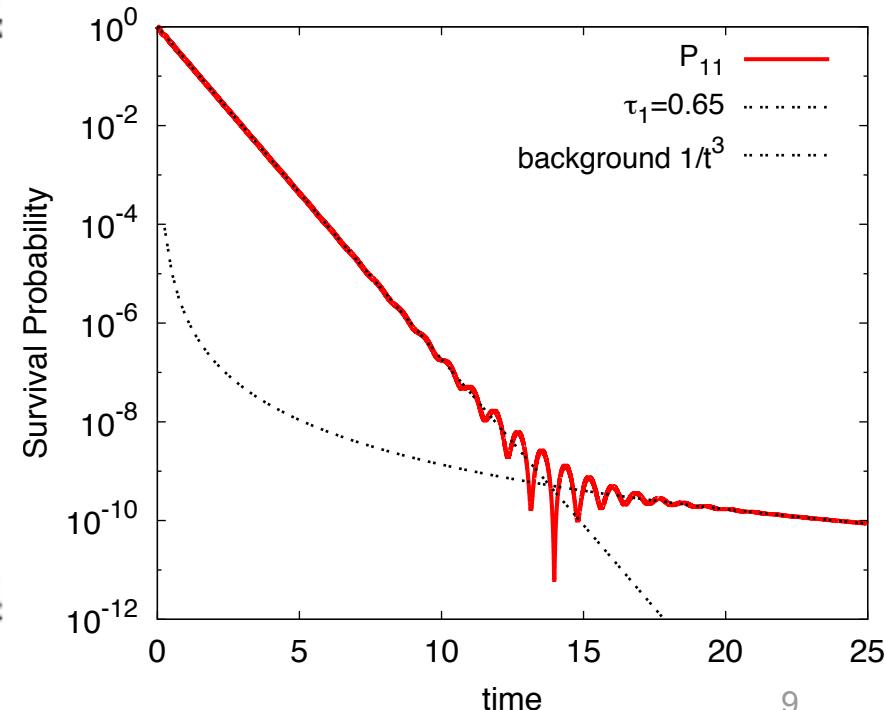
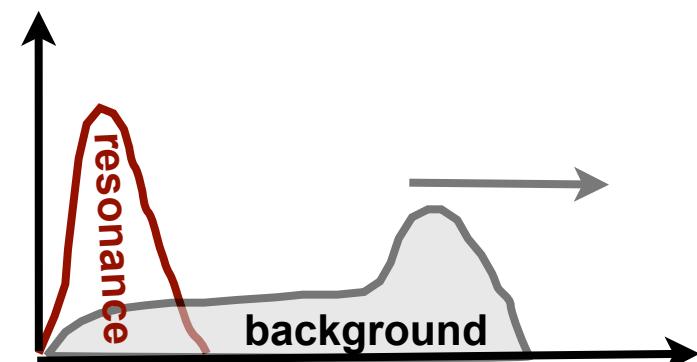
# Time dependence of decay, Winter's model

Winter, Phys. Rev., 123, 1503 1961.

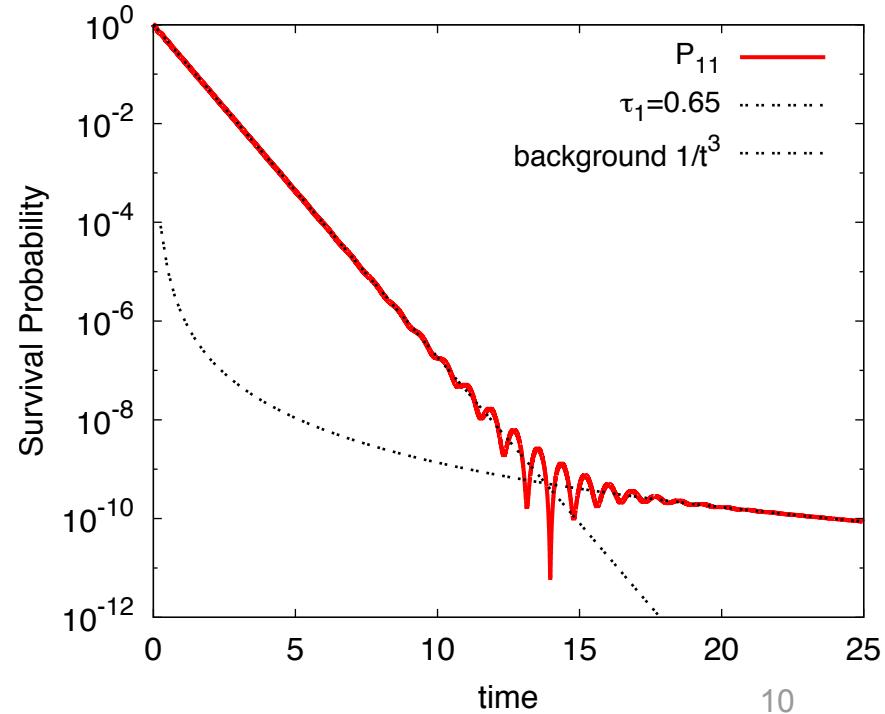
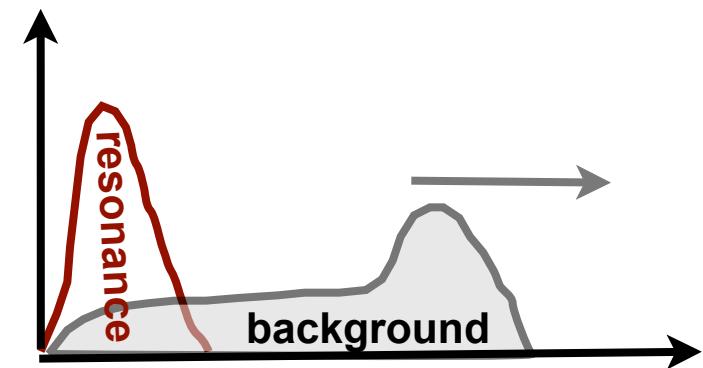
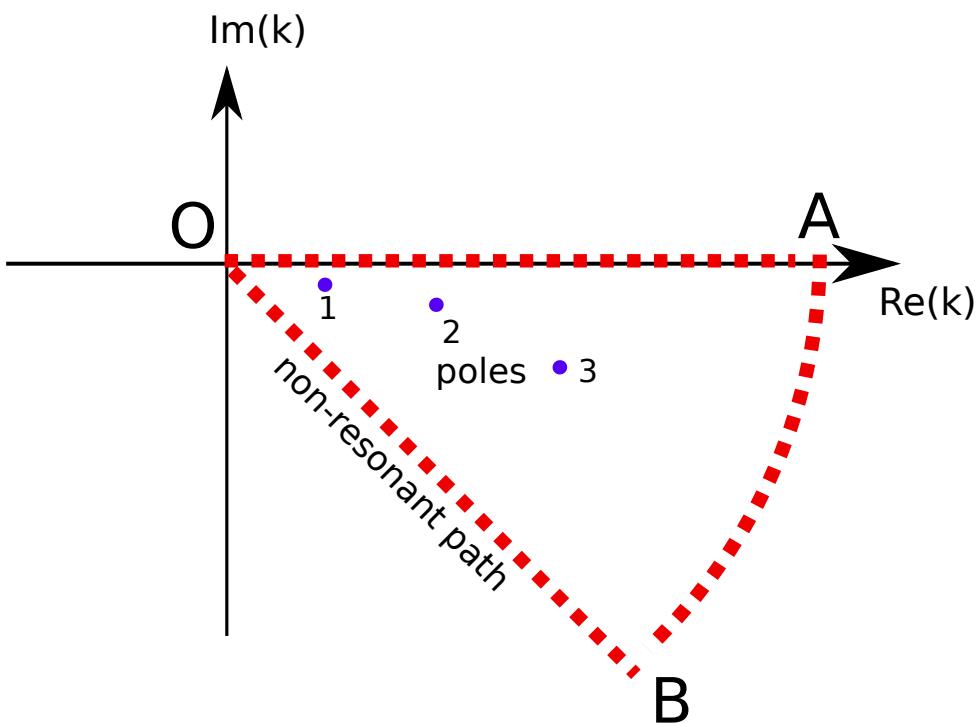




## Winter's model: Dynamics at remote times

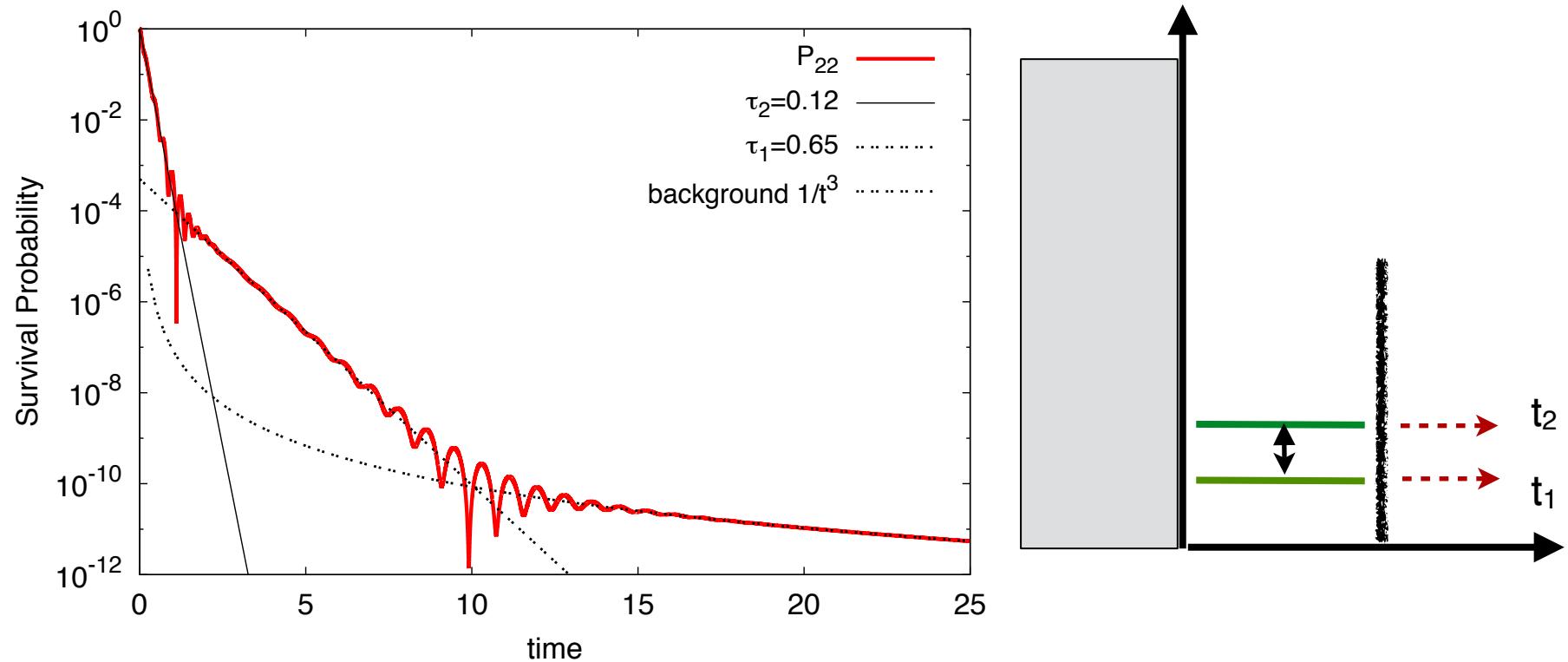


## Winter's model: Dynamics at remote times



# Internal dynamics in decaying system

## Winter's model



# Effective Hamiltonian Formulation

The Hamiltonian in P is:

$$\mathcal{H}(E) = H + \Delta(E) - \frac{i}{2}W(E)$$

Channel-vector:

$$|A^c(E)\rangle = H_{QP}|c; E\rangle$$

Self-energy:

$$\Delta(E) = \frac{1}{2\pi} \int dE' \sum_c \frac{|A^c(E')\rangle\langle A^c(E')|}{E - E'}$$

Irreversible decay to the excluded space:

$$W(E) = \sum_{c(\text{open})} |A^c(E)\rangle\langle A^c(E)|$$

- [1] C. Mahaux and H. Weidenmüller, *Shell-model approach to nuclear reactions*, Amsterdam 1969
- [2] A. Volya and V. Zelevinsky, Phys. Rev. Lett. **94**, 052501 (2005).
- [3] A. Volya, Phys. Rev. C **79**, 044308 (2009).

# Scattering matrix and reactions

$$\mathbf{T}_{cc'}(E) = \langle A^c(E) | \left( \frac{1}{E - \mathcal{H}(E)} \right) | A^{c'}(E) \rangle$$

$$\mathbf{S}_{cc'}(E) = \exp(i\xi_c) \{ \delta_{cc'} - i \mathbf{T}_{cc'}(E) \} \exp(i\xi_{c'})$$

Cross section:

$$\sigma = \frac{\pi}{k'^2} \sum_{cc'} \frac{(2J+1)}{(2s'+1)(2I'+1)} |\mathbf{T}_{cc'}|^2$$

## Additional topics:

- Angular (Blatt-Biedenharn) decomposition
- Coulomb cross sections, Coulomb phase shifts, and interference
- Phase shifts from remote resonances.

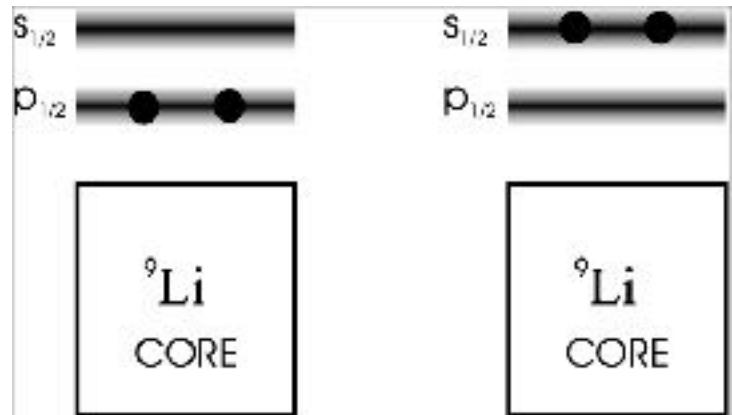
# **Interference between resonances**

# $^{11}\text{Li}$ model

Dynamics of two states coupled to a common decay channel

- Model  $\mathcal{H}$

$$\mathcal{H}(E) = \begin{pmatrix} \epsilon_1 - \frac{i}{2}\gamma_1 & v - \frac{i}{2}A_1 A_2 \\ v - \frac{i}{2}A_1 A_2 & \epsilon_2 - \frac{i}{2}\gamma_2 \end{pmatrix}$$

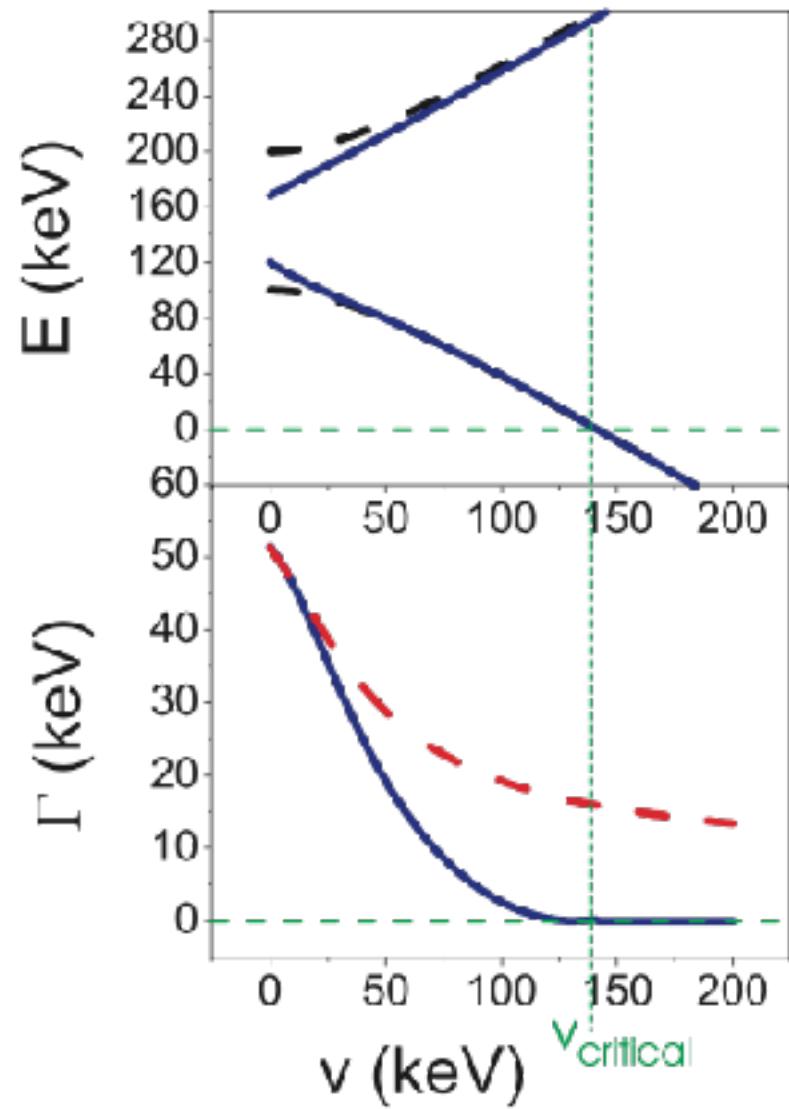
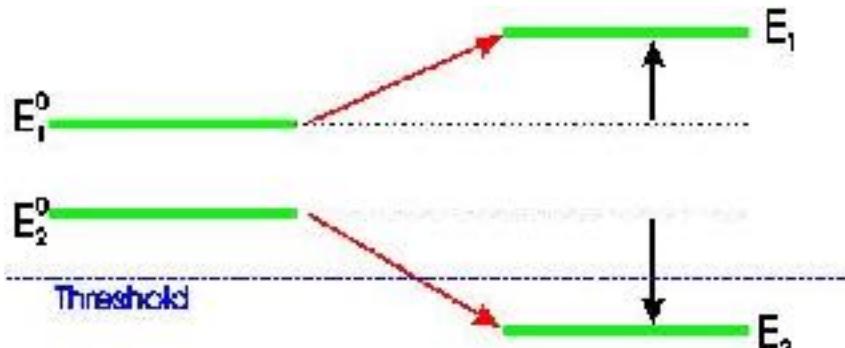


# $^{11}\text{Li}$ model

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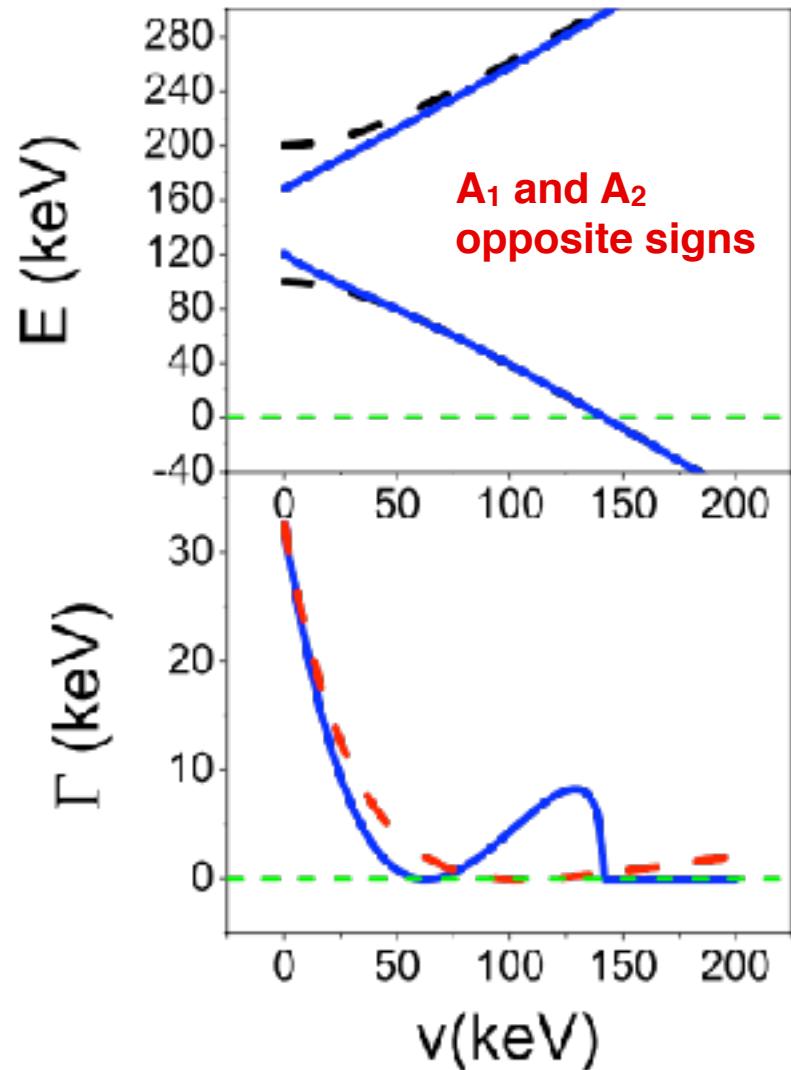
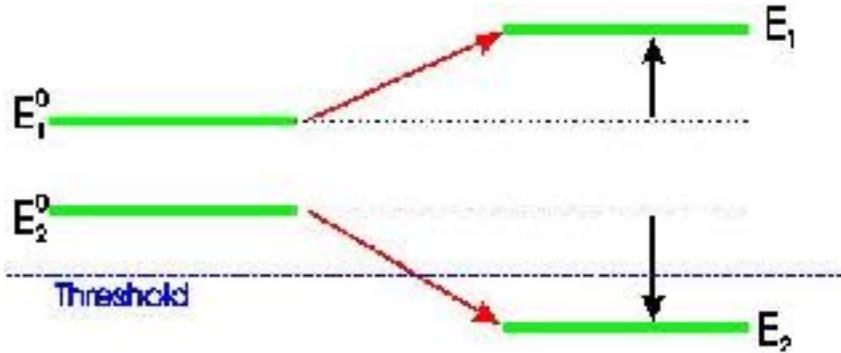


# $^{11}\text{Li}$ model

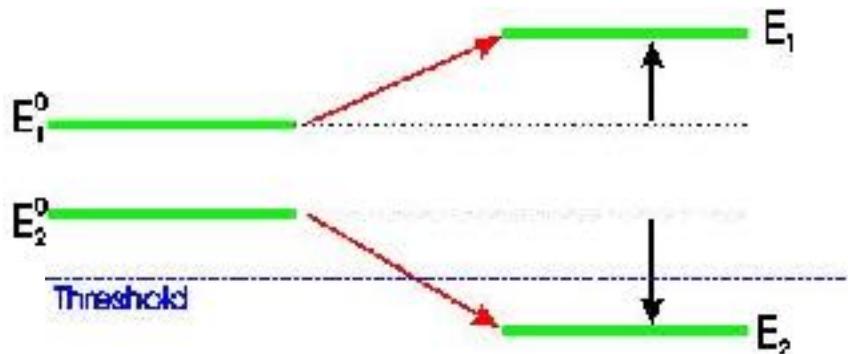
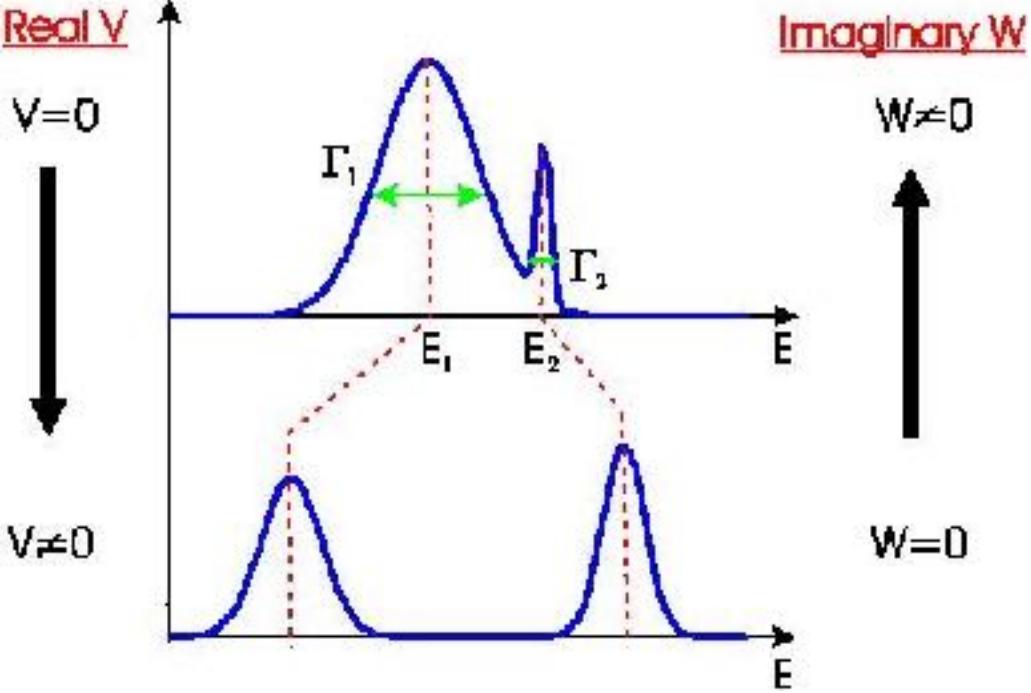
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# Example of interacting resonances

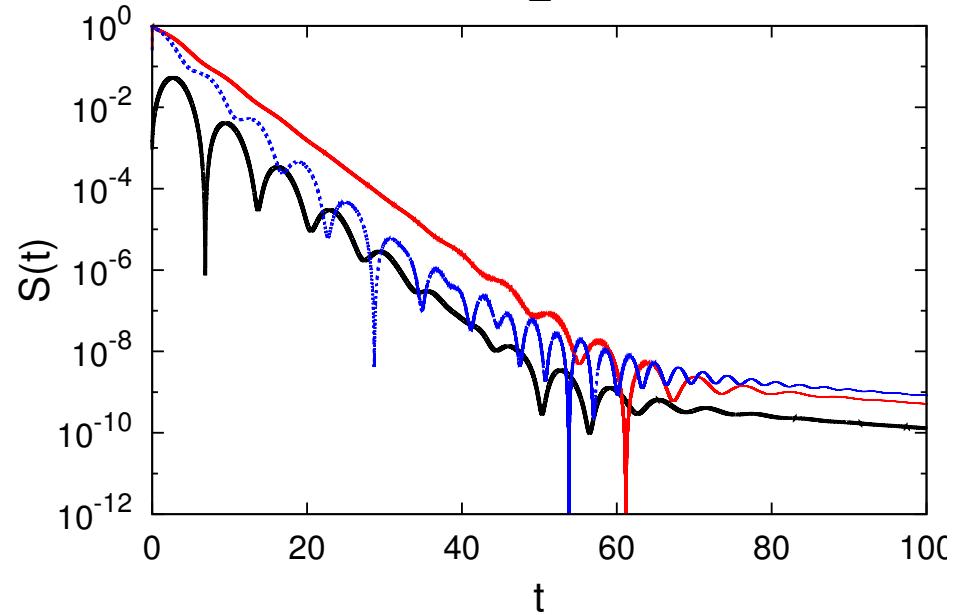
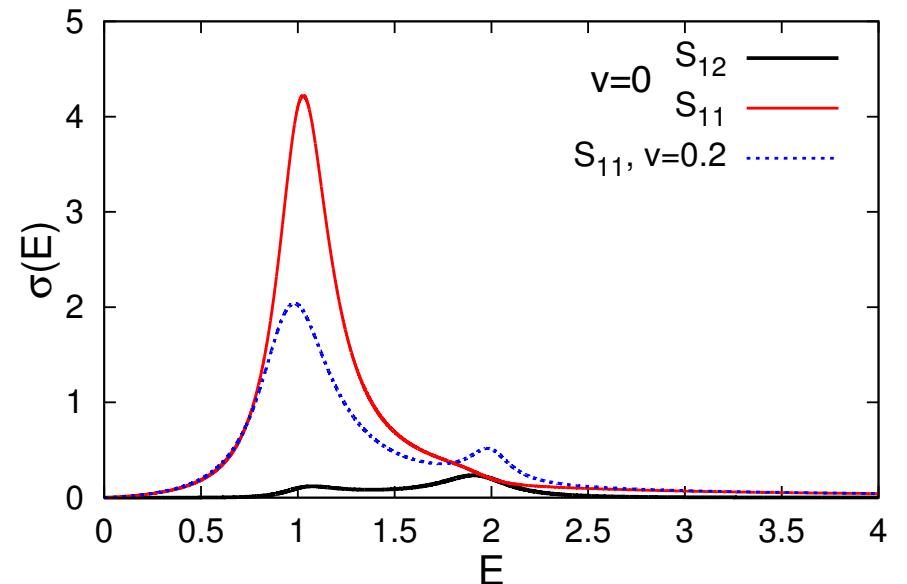
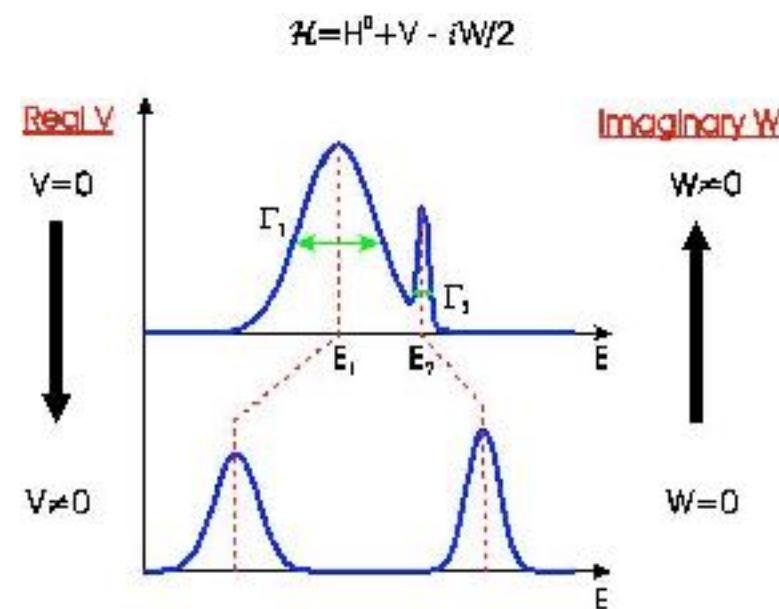


# Two-level system

$$\mathcal{H} = \begin{pmatrix} \epsilon_1 - (i/2)\Gamma_1 & v - (i/2)A_1 A_2 \\ v - (i/2)A_1 A_2 & \epsilon_2 - (i/2)\Gamma_2 \end{pmatrix}$$

$$\Gamma_1 = A_1^2, \quad \Gamma_2 = A_2^2,$$

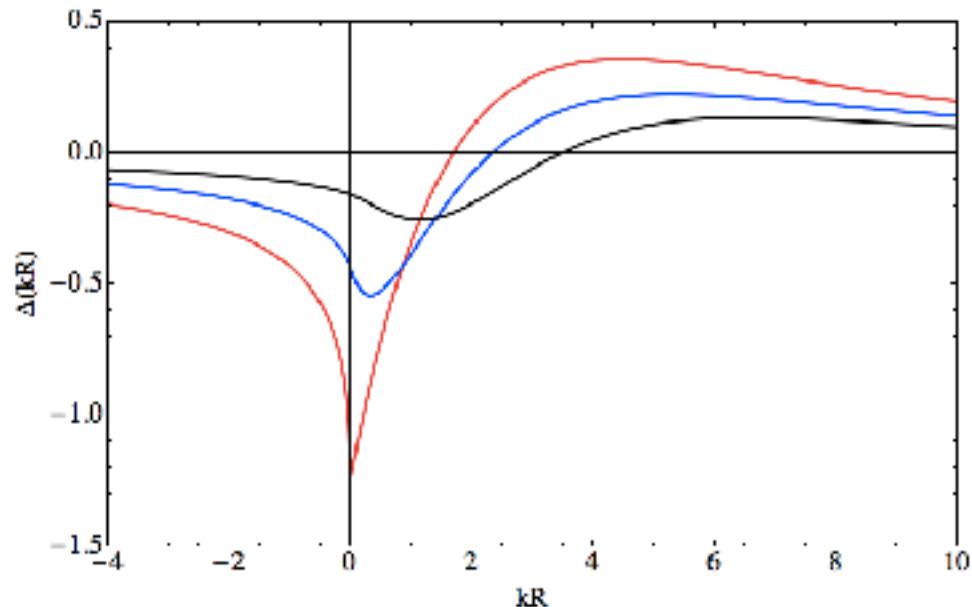
$$S(t) = |\langle \Psi(0) | \Psi(t) \rangle|^2$$



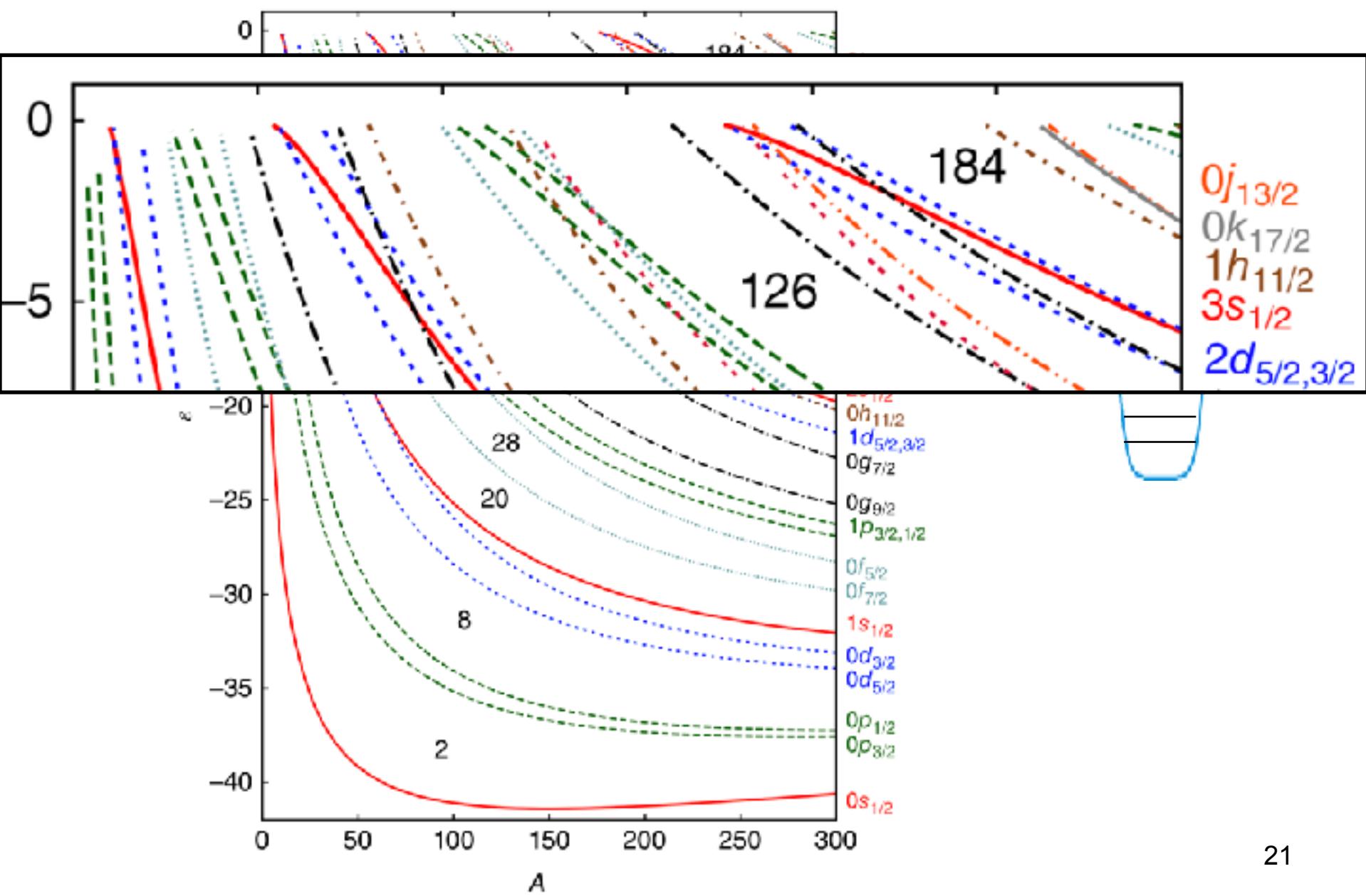
# Virtual excitations

$$H'(\epsilon) = \Delta(\epsilon) - \frac{i}{2}\Gamma(\epsilon)$$

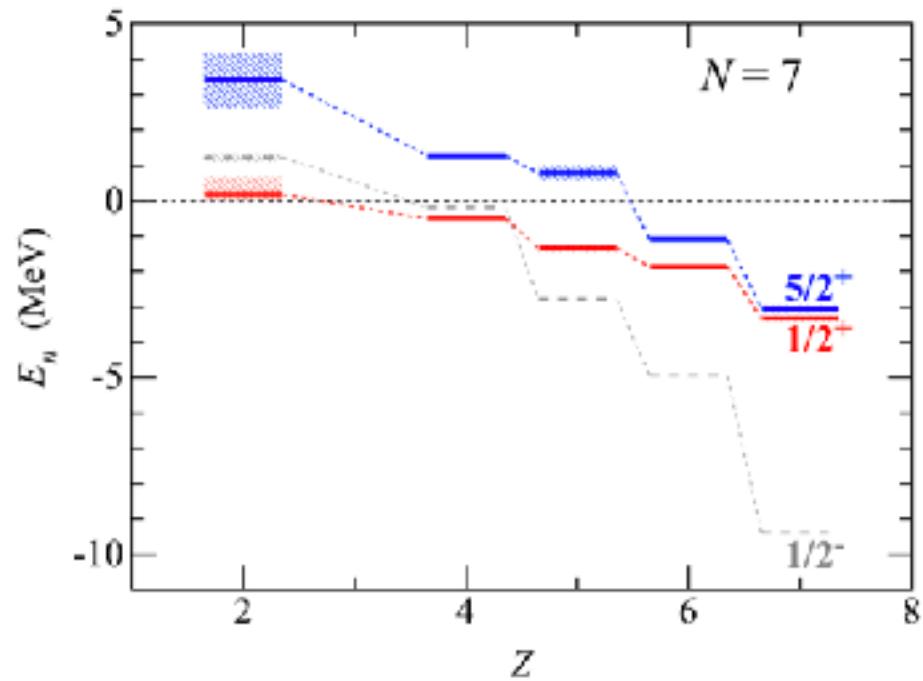
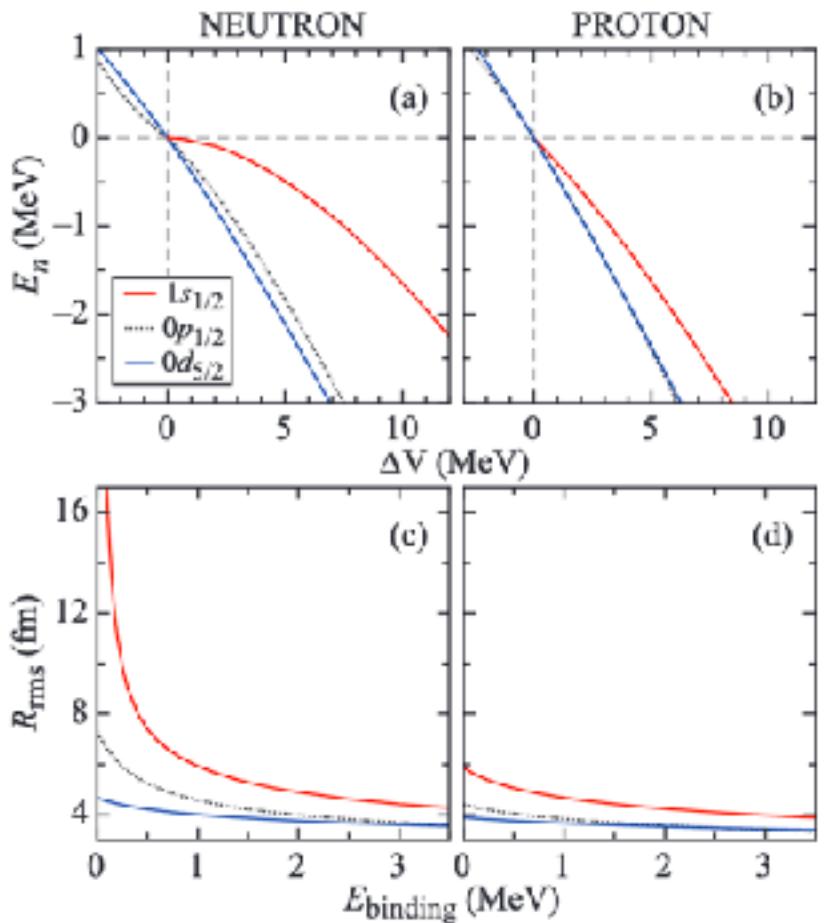
$$\Delta(E) = \frac{1}{2\pi} \oint dE' \sum_c \frac{|A^c(E')\rangle\langle A^c(E')|}{E - E'}$$



# Evolution of single particle energies

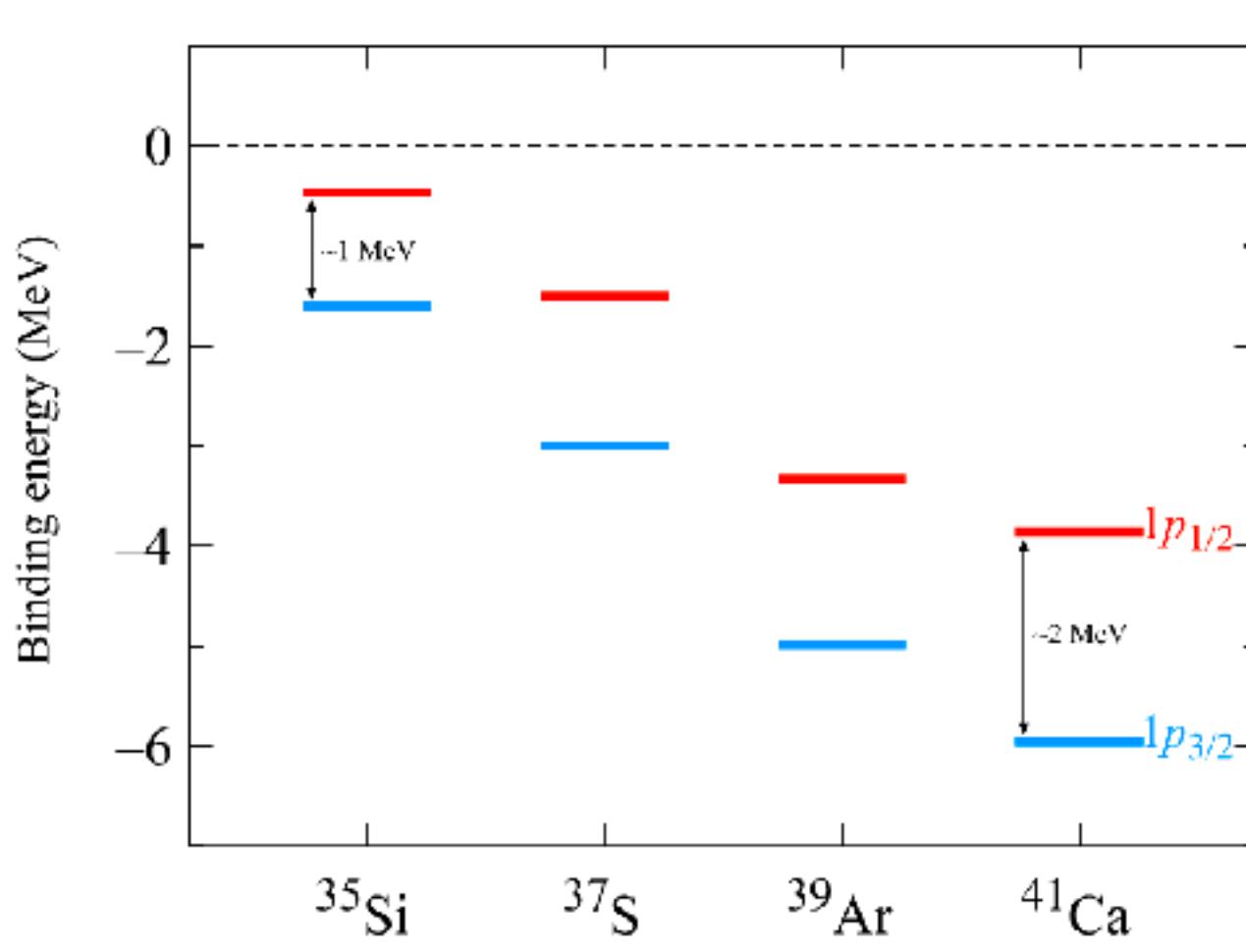


# Effect of weak binding



C. R. Hoffman, B. P. Kay, and J. P. Schiffer Phys. Rev. C 89, 061305(R)  
B. P. Kay, C. R. Hoffman, and A. O. Macchiavelli Phys. Rev. Lett. 119, 182502

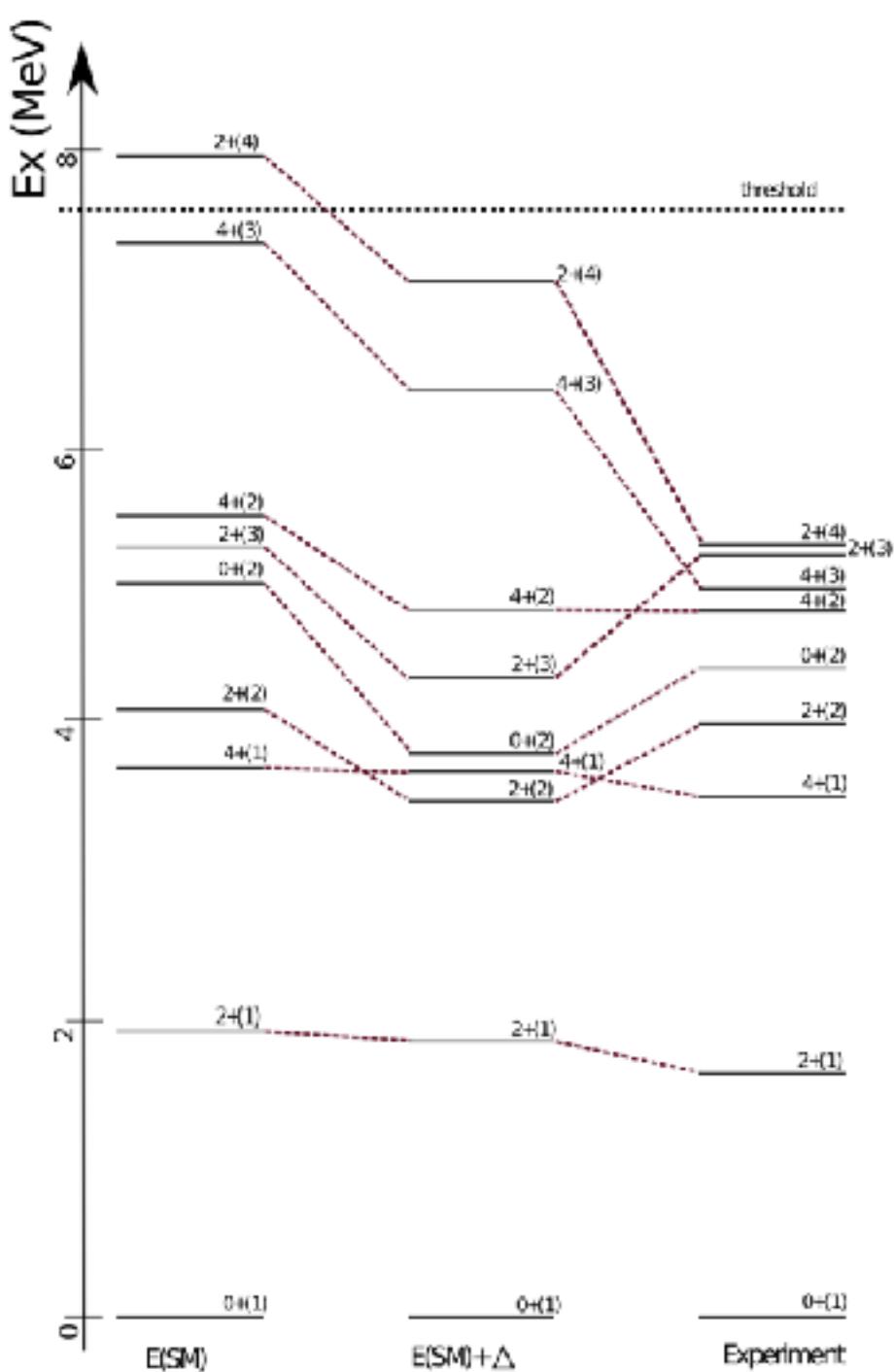
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C. R. Hoffman, B. P. Kay, and J. P. Schiffer Phys. Rev. C 89, 061305(R)  
B. P. Kay, C. R. Hoffman, and A. O. Macchiavelli Phys. Rev. Lett. 119, 182502

# Role of virtual excitations

Spectrum of  $^{20}\text{O}$



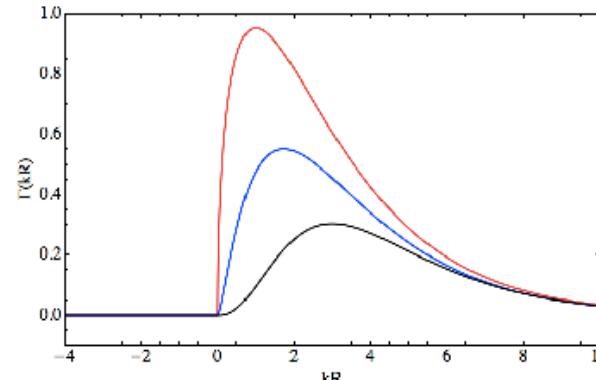
# Superradiance

$$\mathcal{H}(E) = H + \Delta(E) - \frac{i}{2}W(E)$$

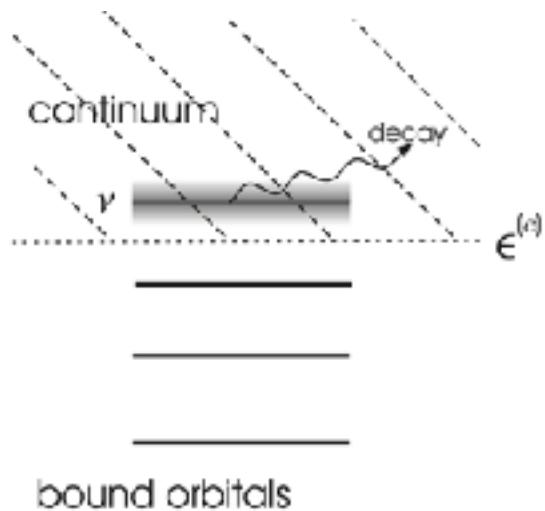
$$W_{12}(E) = 2\pi \sum_{c(\text{open})} A_1^c(E) A_2^{c*}(E) \quad \text{Factorized operator}$$

Factorized form leads to unitarity of scattering matrix

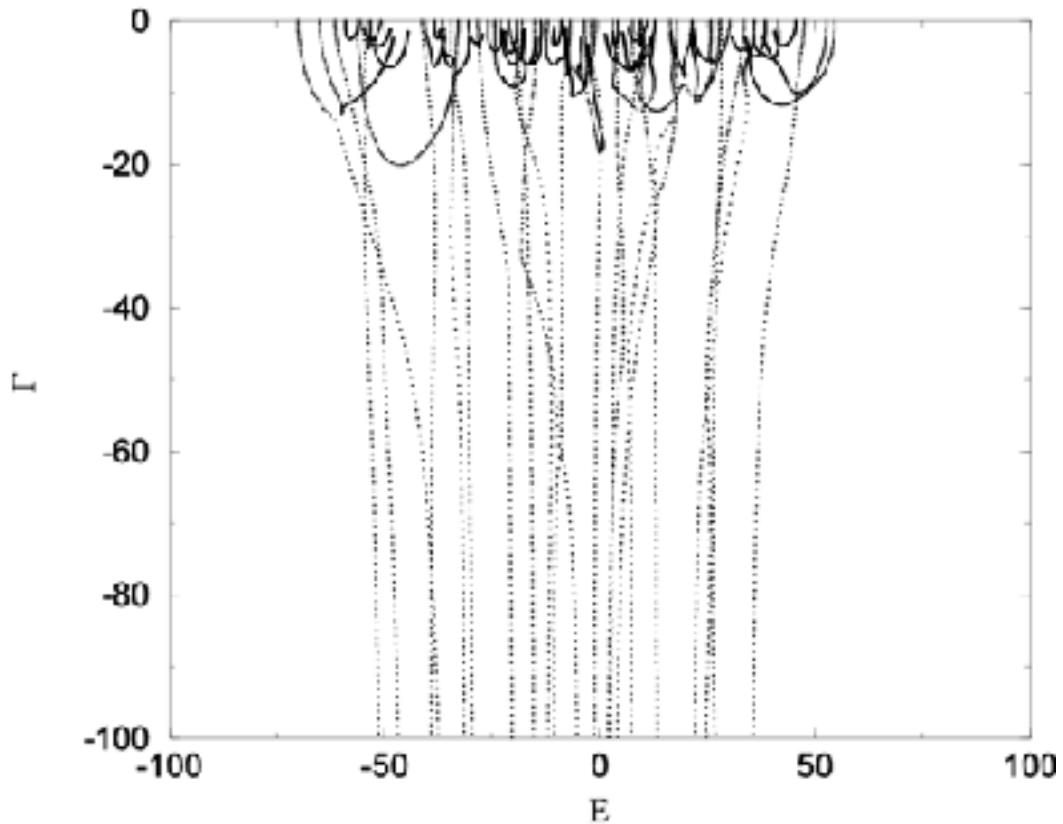
$$T_{cc'}(E) = \langle A^c(E) | \left( \frac{1}{E - \mathcal{H}(E)} \right) | A^{c'}(E) \rangle$$



# Single-particle decay in many-body system



Evolution of complex energies  $E=E-i\Gamma/2$  as a function of  $\gamma$

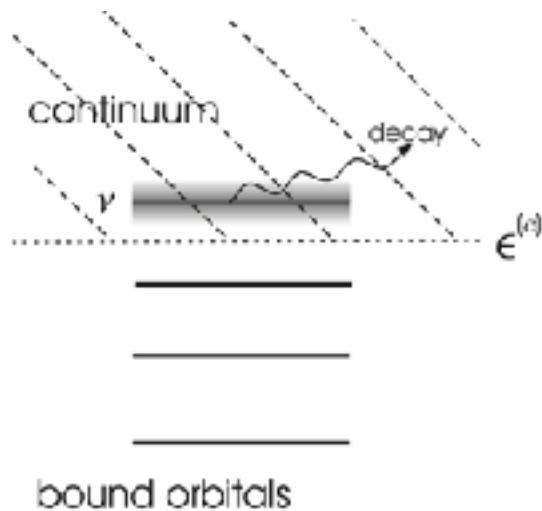


- Assume energy independent  $W$
- Assume one channel  $\gamma=A^2$
- System 8 s.p. levels, 3 particles
- One s.p. level in continuum  $e=\epsilon - i\gamma/2$

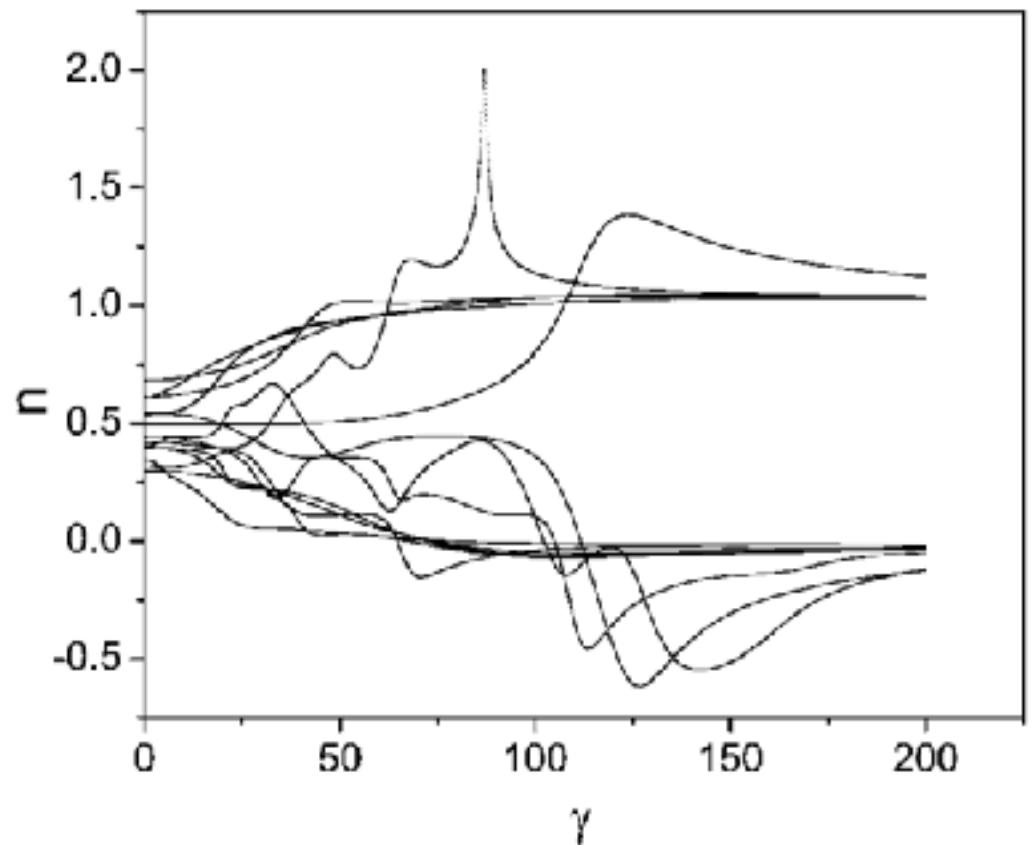
Total states  $8!/(3! 5!) = 56$ ; states that decay fast  $7!/(2! 5!) = 21$

# Single-particle decay in many-body system

Evolution of occupancies  
as a function of  $\gamma$



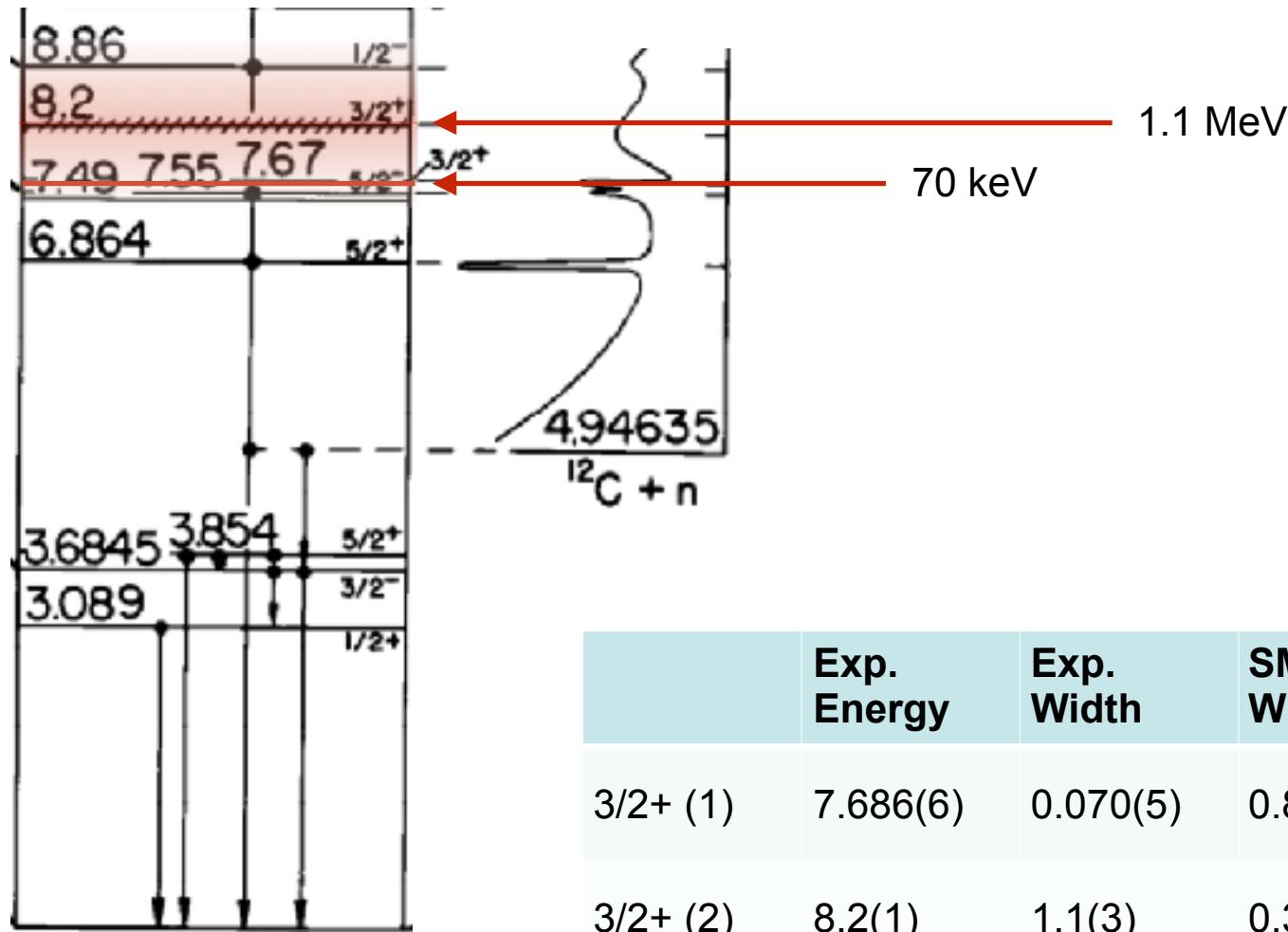
$$n_v(j; \gamma) = \frac{\partial \Gamma_j(\gamma)}{\partial \gamma}$$



# Superradiance

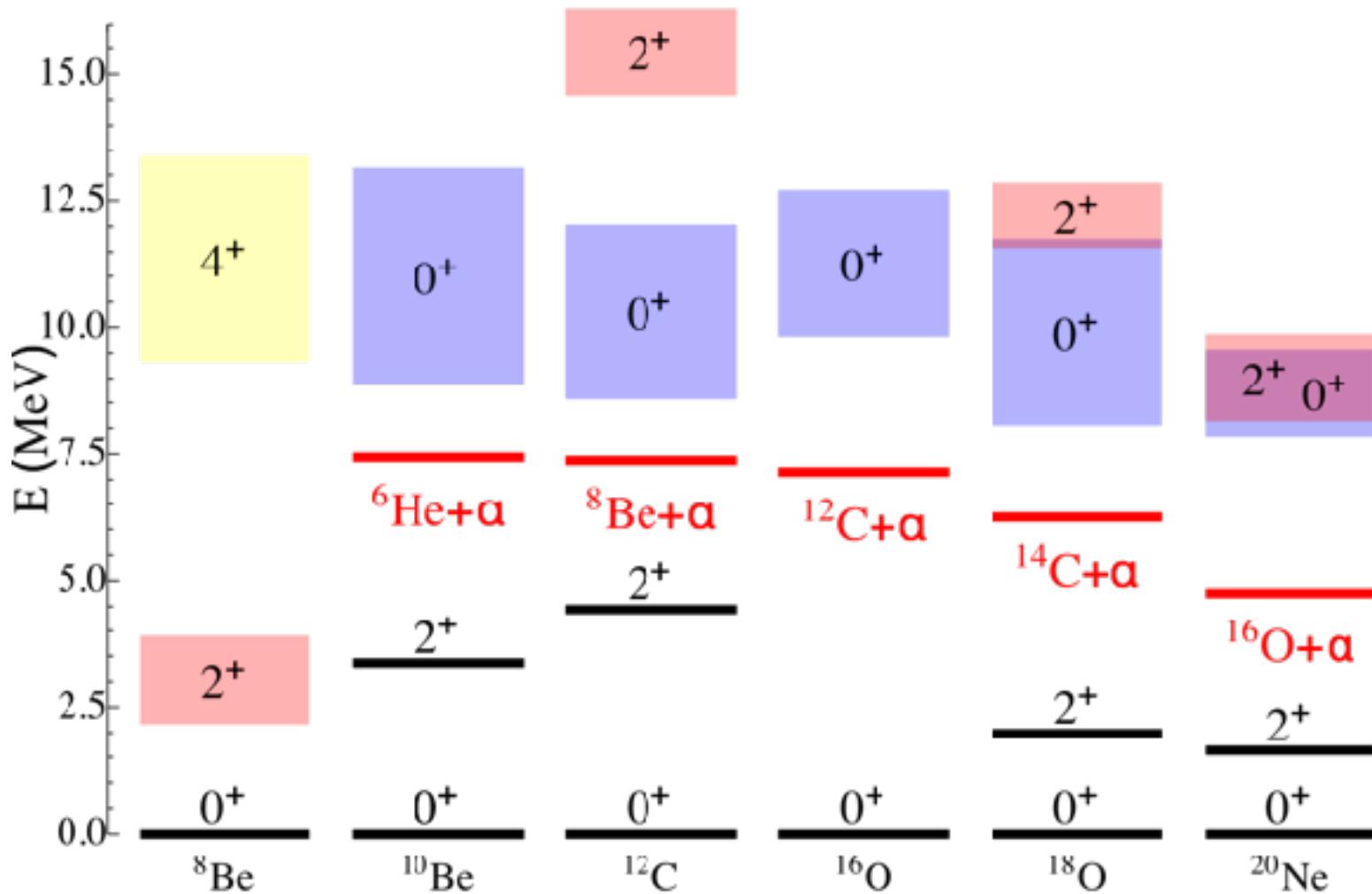
- Factorized form of non-hermitian component consistent with unitarity
- Low operator rank, number of channels versus number of many-body states
- When imaginary part dominates states separate into
  - Superradiant (strongly coupled to continuum)
  - Decoupled from decay
- Internal wave functions are “reoriented” either along or away from decay
- Coupling to decay is a collective phenomenon
- There is a phase transition in many-body dynamics associated with superradiance

# Superradiance in $^{13}\text{C}$



$^{13}\text{C}$

# Searching for clustering states



# Interplay of collectivities

## Definitions

$n$  - labels particle-hole state

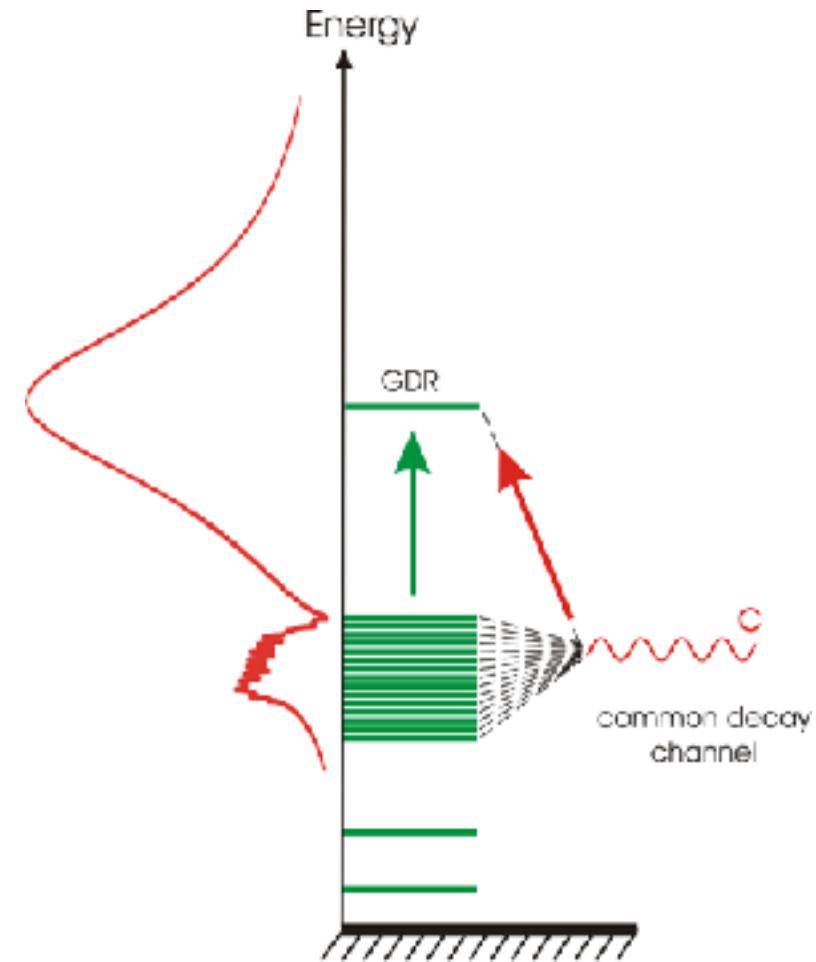
$\varepsilon_n$  – excitation energy of state  $n$

$d_n$  - dipole operator

$A_n$  – decay amplitude of  $n$

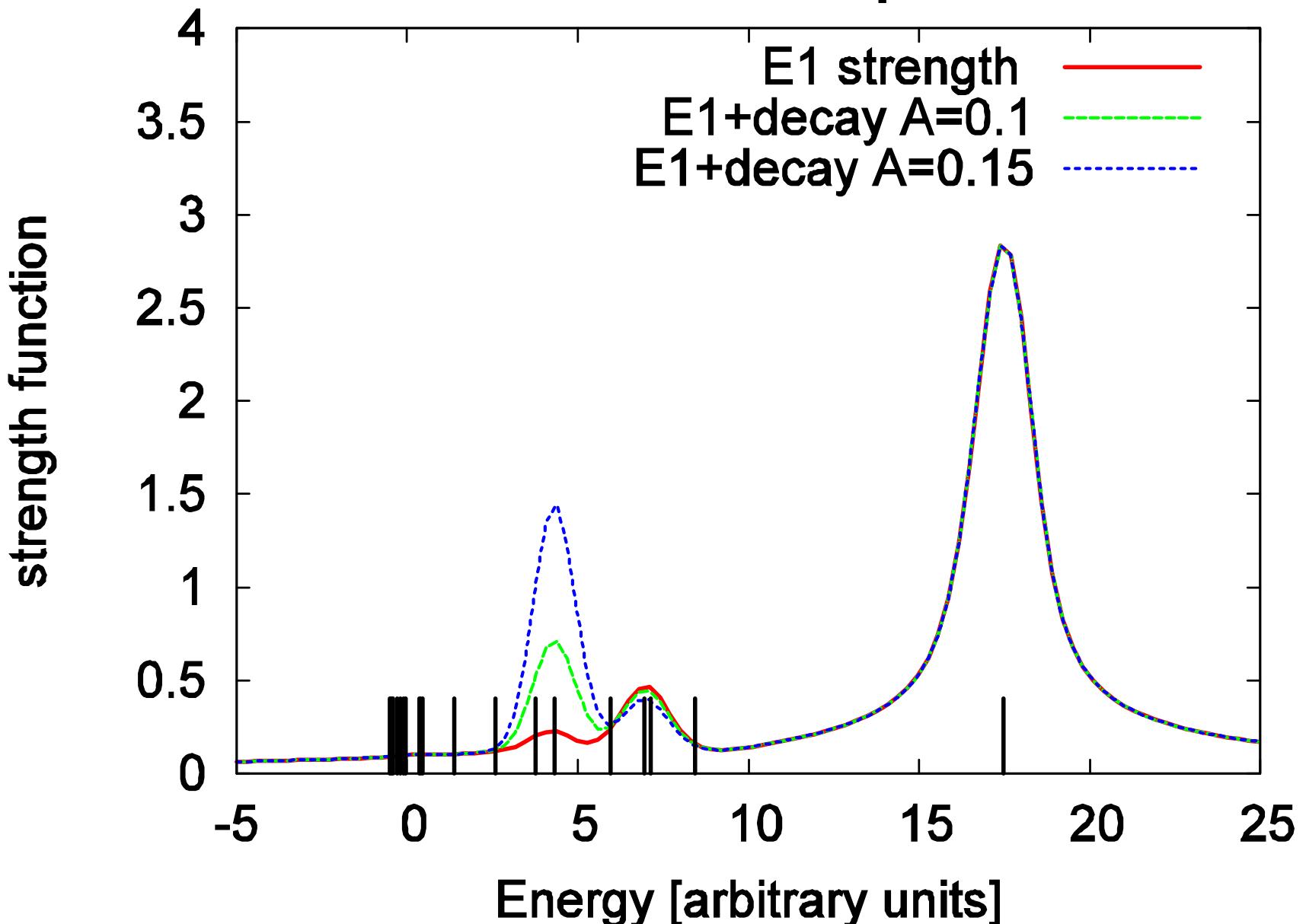
## Model Hamiltonian

$$\mathcal{H}_{nn'} = \epsilon_n \delta_{nn'} + \lambda d_n d_{n'} - \frac{i}{2} A_n A_{n'},$$

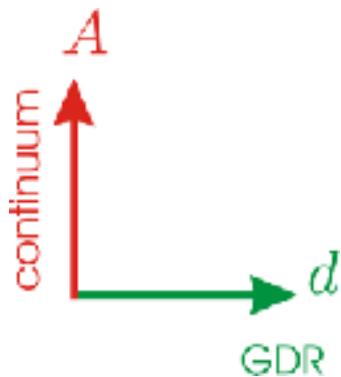


**Everything depends on  
angle between multi dimensional vectors  
 $\mathbf{A}$  and  $\mathbf{d}$**

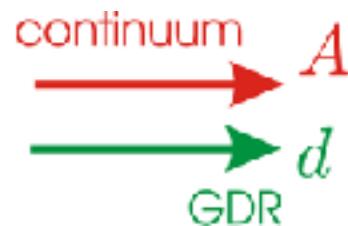
# Model Example



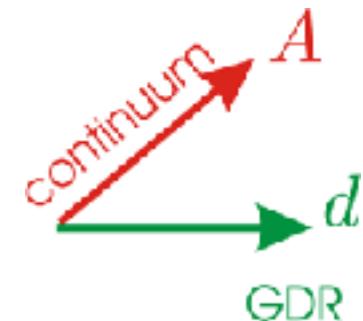
## Pigmy resonance



**Orthogonal:**  
GDR not seen

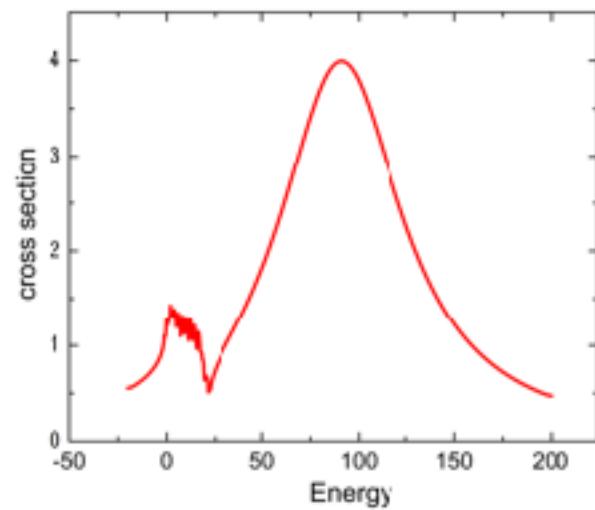
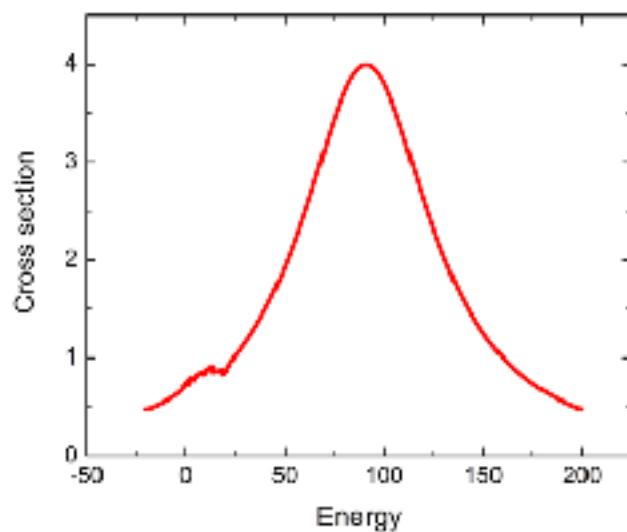


**Parallel:**  
Most effective excitation  
of GDR from continuum

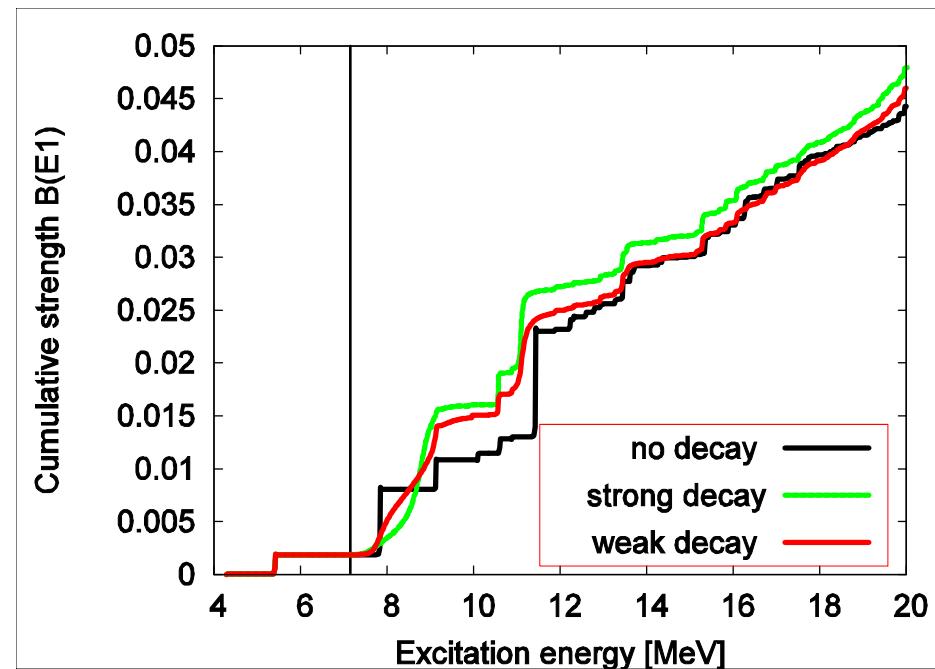
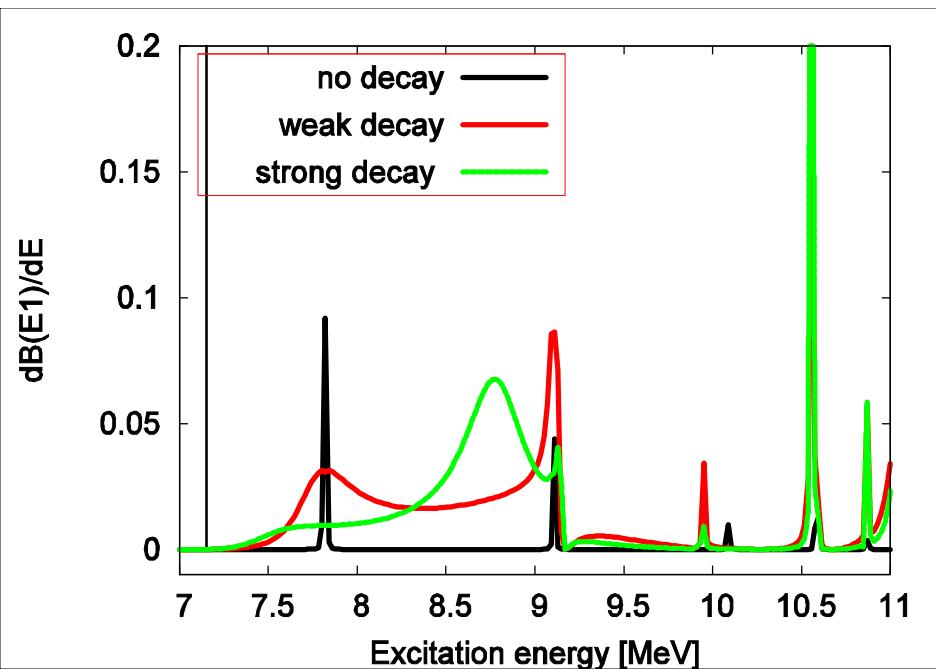


**At angle:**  
excitation of GDR  
and pigmy

A model of 20 equally distant levels is used



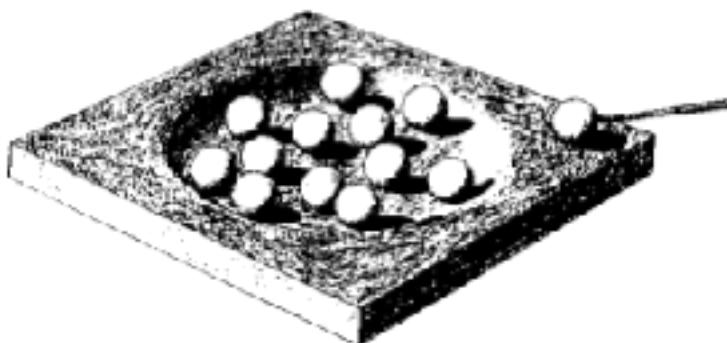
# $B(E1)$ strength in $^{22}\text{O}$



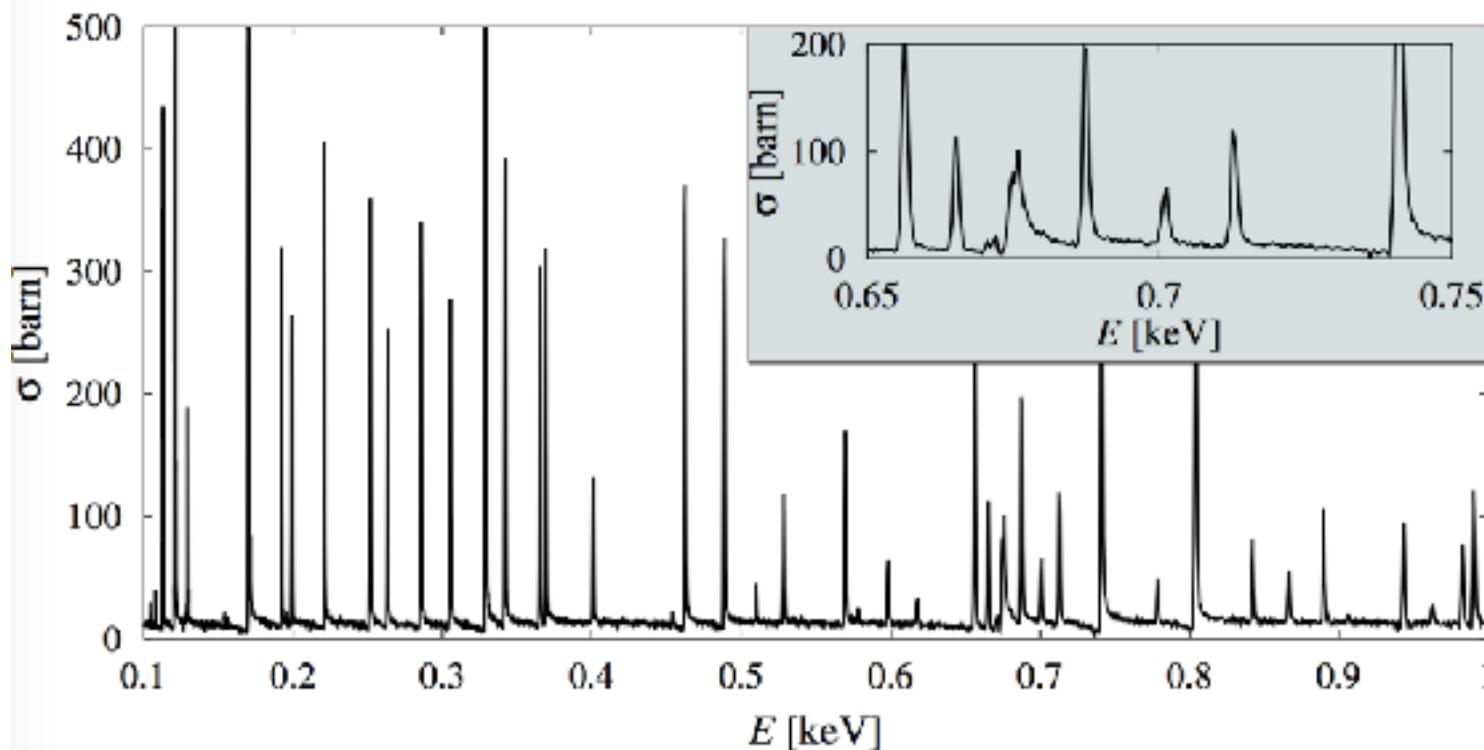
Super-radiance phenomenon,  
strong decay makes leads to width  
distribution

Cumulative strength shows that  
super-radiance increases low-  
lying  
dipole strength

# Distribution of decay widths in a chaotic system



Wooden toy model illustrating Bohr's compound nucleus, from Nature 137, 351 (1936)



## Many-body complexity and reduced widths

$ c\rangle$	Channel-vector (normalized)	Reduced width
$ I\rangle$	Eigenstate	$\gamma_I^c =  \langle I c\rangle ^2$

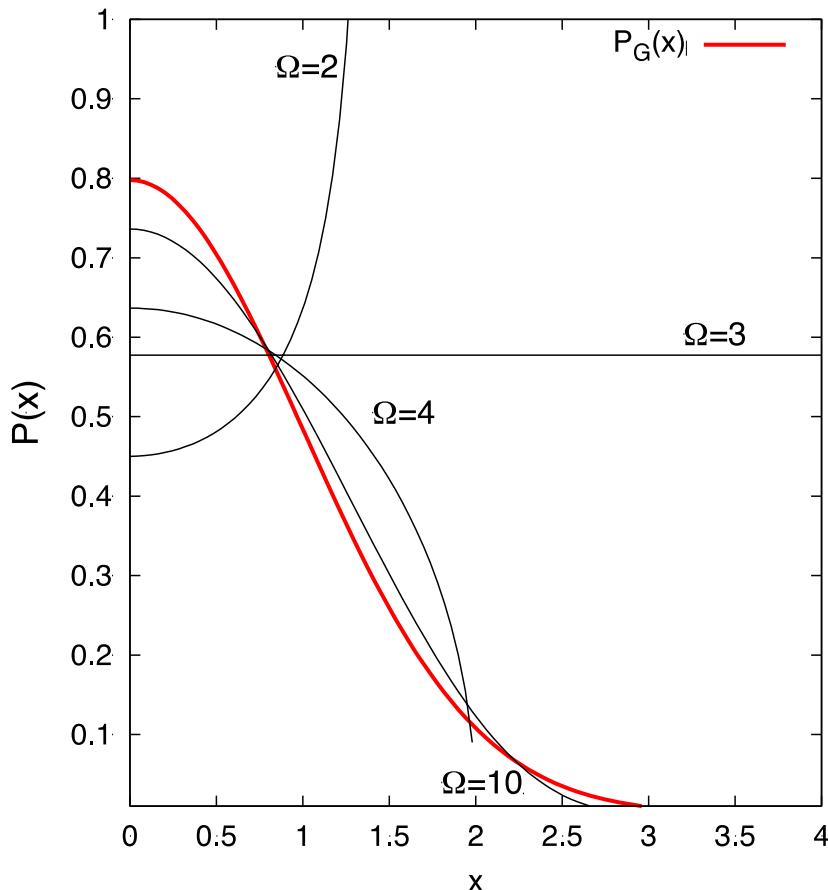
What is the distribution of the reduced width?

Average width       $\bar{\gamma} = \frac{1}{\Omega} \sum_I \gamma_I^c = \frac{\langle c|c\rangle}{\Omega}$       Amplitude       $x_I = \sqrt{\gamma_I/\bar{\gamma}}$

If any direction in the  $\Omega$ -dimensional Hilbert space is equivalent

$$P(x_{I_1}, \dots, x_{I_\Omega}) \sim \delta \left( \Omega - \sum_I x_I^2 \right)$$

# Why Porter-Thomas Distribution?



For large  $v$  channels

Projection of a randomly oriented vector in  $\Omega$ -dimensional space

$$P(x) = \frac{V_{\Omega-1}}{\sqrt{\Omega} V_\Omega} (1 - x^2/\Omega)^{(\Omega-3)/2}$$

$$V_\Omega = \frac{\Omega \pi^{\Omega/2}}{\Gamma(\Omega/2 + 1)}$$

For large  $\Omega$  this leads to Gaussian

$$P_G(x) = \sqrt{\frac{2}{\pi}} \exp(-x^2/2)$$

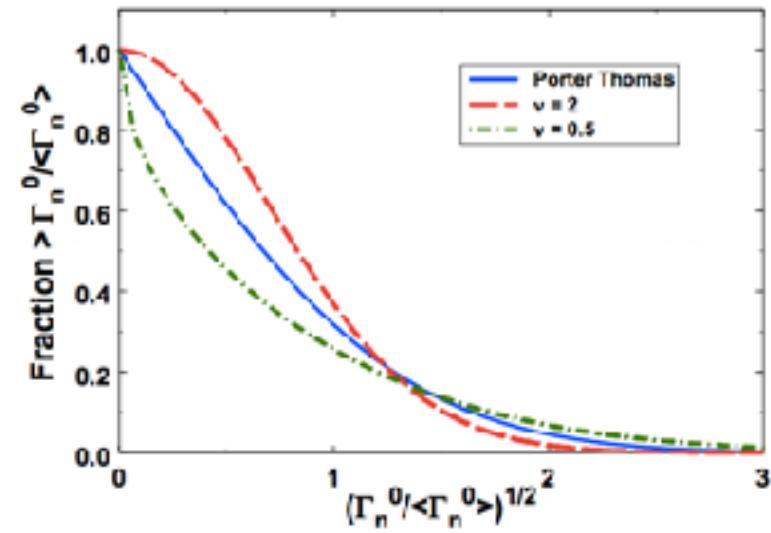
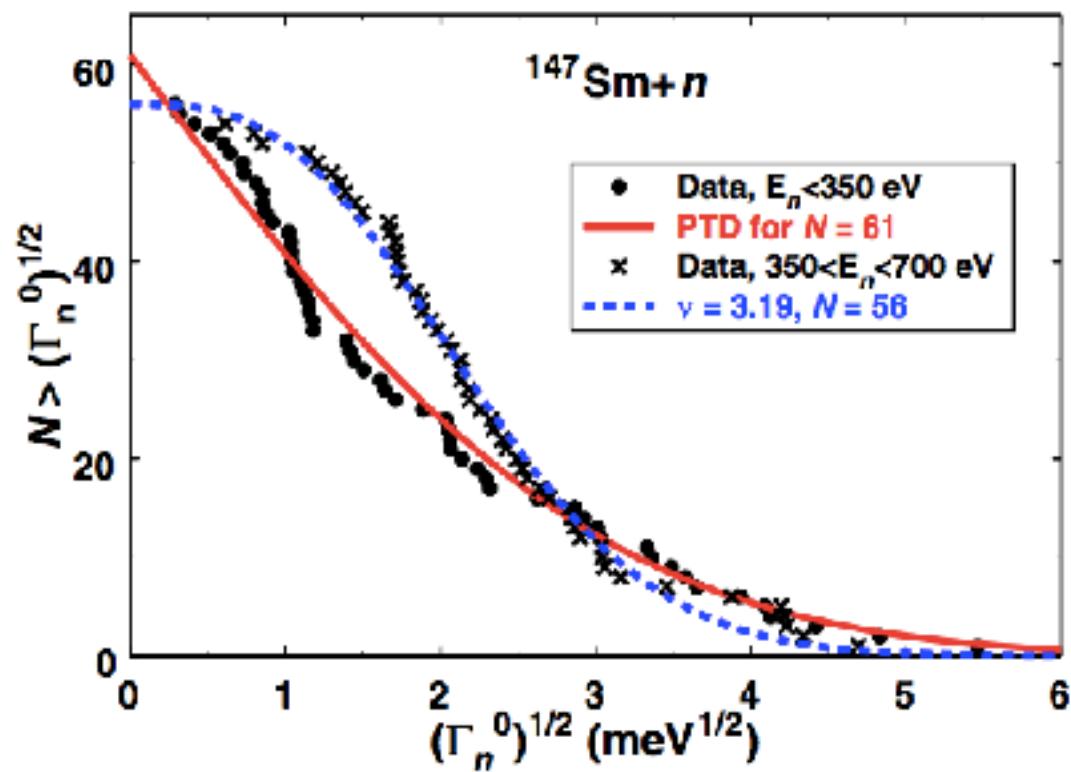
$$P_\nu(\gamma) = \frac{1}{\gamma} \left( \frac{\nu \gamma}{2 \bar{\gamma}} \right)^{\nu/2} \frac{1}{\Gamma(\nu/2)} \exp\left(-\frac{\nu \gamma}{2 \bar{\gamma}}\right)$$

# Nuclear theory nudged? Violation of Porter-Thomas Distribution

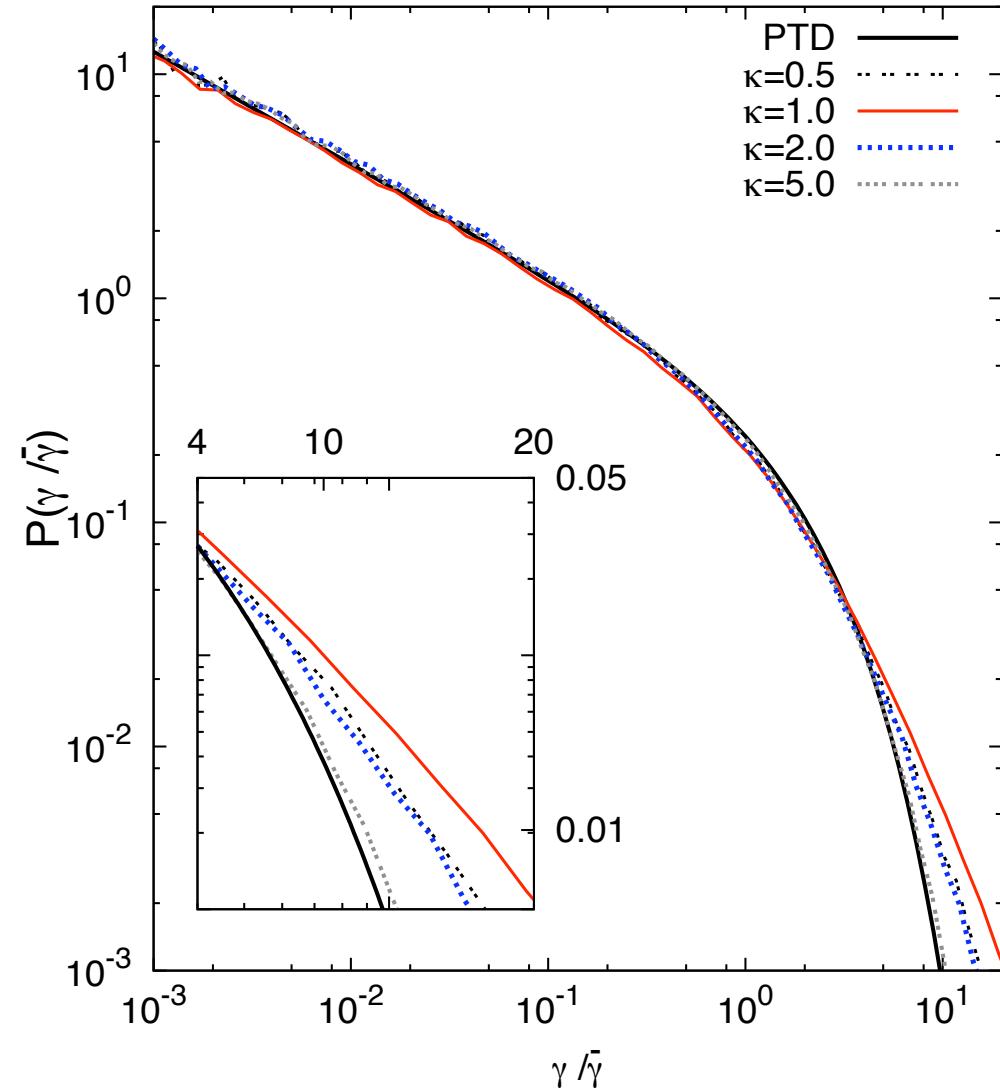
Random matrix theory is rejected with 99.997% probability [Koehler, et. al. Phys. Rev. Lett. 105, 072502 (2010)] In platinum  $\nu = 0.5$

## Implications:

Capture rates, astrophysical reactions, nuclear reactors, critical mass, shielding...



# Superradiance: decay collectivity

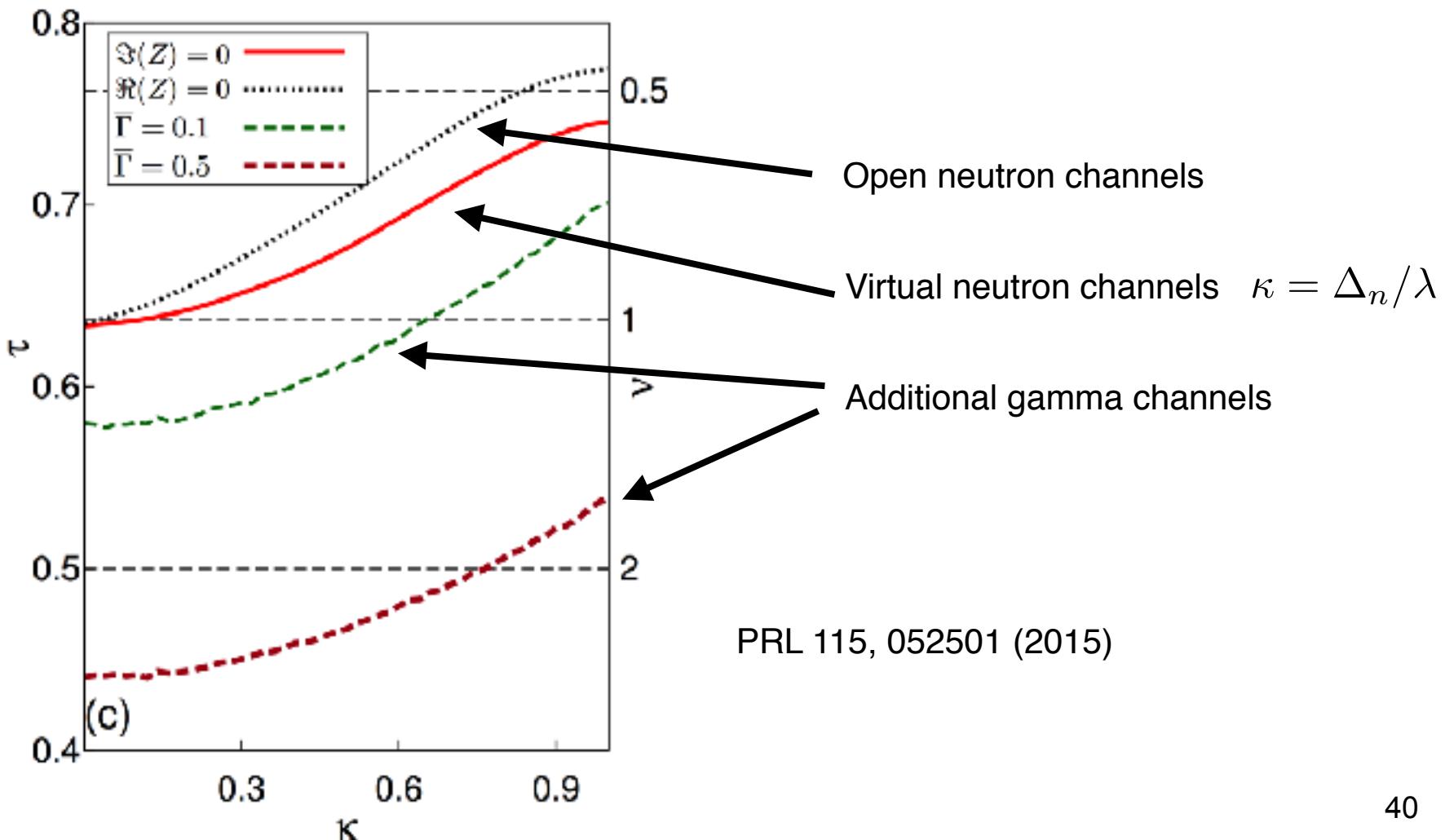


$\Omega=10000$

SR leads to a small number of broad states.

# Virtual excitations as possible explanation

$$\mathcal{H}_{\text{eff}} = H_{\text{GOE}} + \Delta_n + \frac{i}{2} W_n + \frac{i}{2} W_\gamma$$



# Open Question: Dipole-moment and violation of P and T-symmetries

$$\mathbf{d} = \frac{\langle \mathbf{d} \cdot \mathbf{J} \rangle}{J(J+1)} \mathbf{J}$$

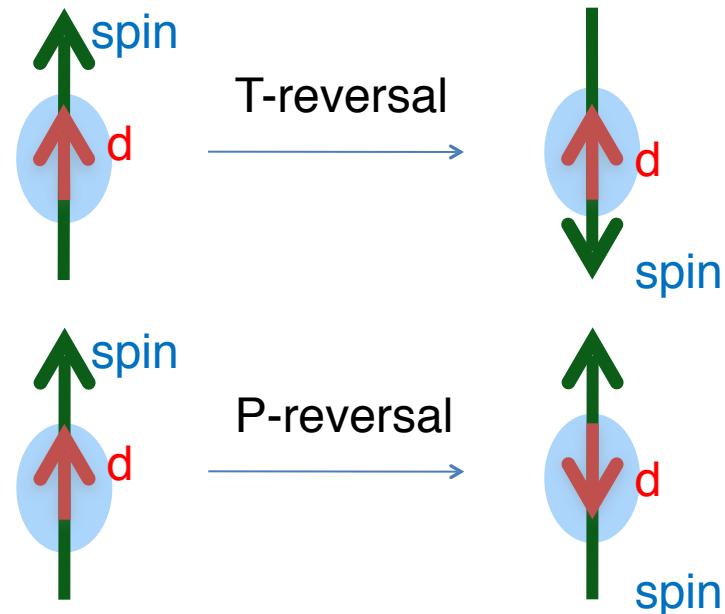
Observation of the dipole moment is an indication of parity and time-reversal violation

## Limit on EDM in electron

Experiment  $10^{-27}$  e cm

Standard model »  $10^{-38}$  e cm

Physics beyond SM »  $10^{-28}$  e cm



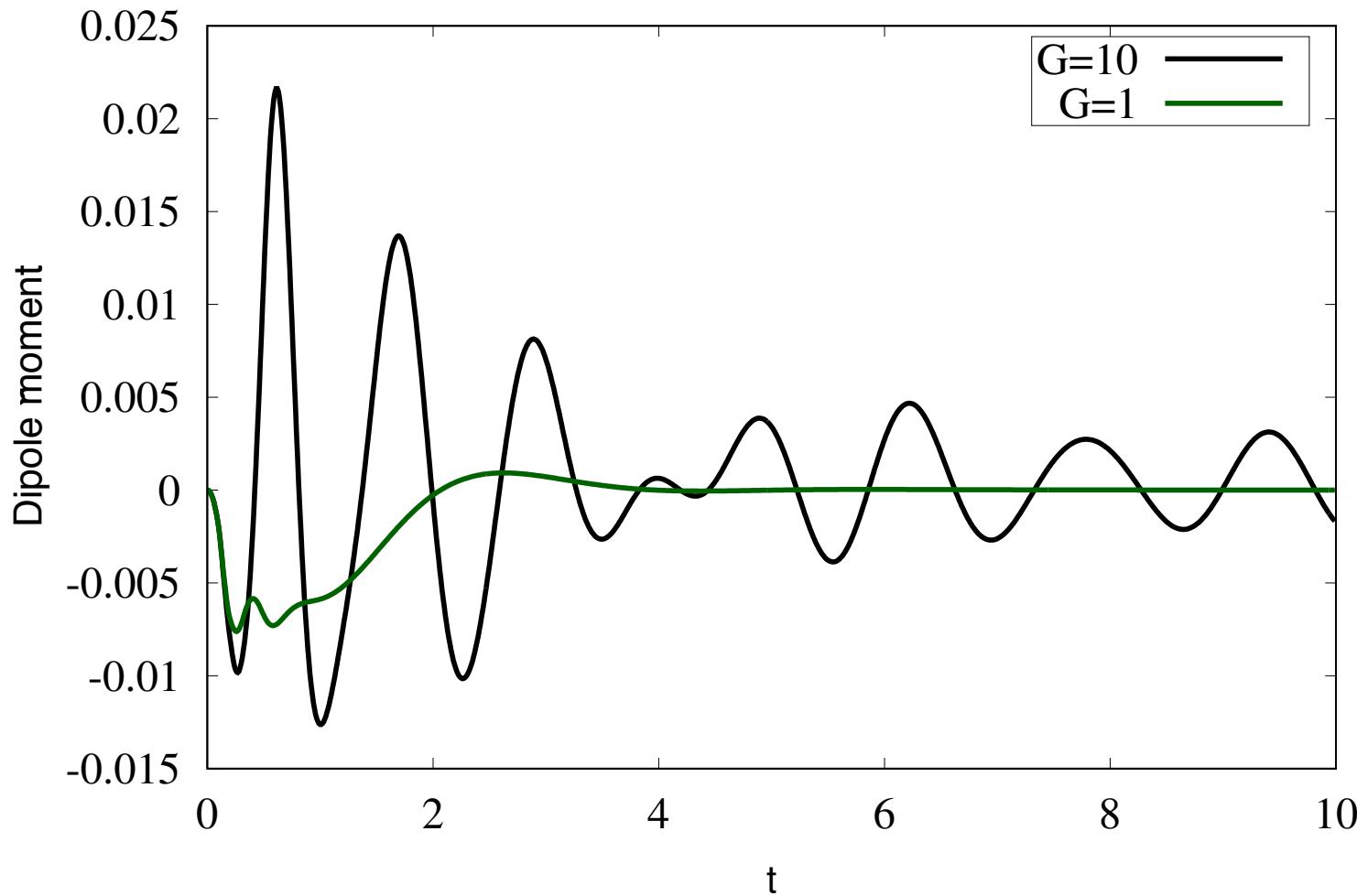
# *Why is this interesting?*

- Sensitive test of CP violation in the standard model
- Baryon asymmetry in the universe.
- Physics beyond standard model.

system	EDM limit	SM
e (electron)	$10^{-27}$	$10^{-40}$
n (neutron)	$3.0 \times 10^{-26}$	$10^{-32}$
$^{225}\text{Ra}$	$1.4 \times 10^{-23}$	$10^{-33}$
$^{199}\text{Hg}$	$7.4 \times 10^{-30}$	$10^{-33}$

# Dipole moment in decaying system

Symmetric Winter's model



Winter's model, with slightly broken parity by coupling to continuum strength

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