

SIGNALS OF COMPOSITE DARK MATTER

Enrico Rinaldi



RIKEN BNL Research Center
RIKEN Nishina Center

SIGNALS OF COMPOSITE DARK MATTER

Enrico Rinaldi



RIKEN BNL Research Center
RIKEN Nishina Center





Lattice Strong Dynamics collaboration

Argonne: Jin, Osborn

Liverpool: Schaich

Bern: Gasbarro

LLNL: Vranas

Boston: Brower,
Rebbi, Howarth

UC Davis: Kiskis

Nvidia: Weinberg

Yale: Appelquist,
Fleming, Cushman

Colorado: Neil,
Witzel

Oregon: Kribs

RIKEN: ER



Lattice Strong Dynamics collaboration

Argonne: Jin, Osborn

Liverpool: Schaich

Bern: Gai

Lattice for BSM workshops

Boston:

Rebbi, H

Nvidia: V

Colorad

Witzel

- University of Colorado Boulder, April 2018
- Boston University, April 2017
- Argonne National Lab, April 2016
- Lawrence Livermore National Lab, April 2015
- Brookhaven National Lab, December 2013
- University of Colorado Boulder, October 2012
- Fermilab, October 2011
- Boston University, November 2009
- Lawrence Livermore National Lab, May 2008

kis

st,

man

RIKEN: ER

NEW STRONG DYNAMICS

Composite Higgs

Composite Dark Matter

NEW STRONG DYNAMICS

Composite Higgs

Composite Dark Matter

New $SU(\mathbf{N}_c)$ gauge sector with \mathbf{N}_f fermions in the \mathbf{N}_r representation of the gauge group

NEW STRONG DYNAMICS

Composite Higgs

Composite Dark Matter

New $SU(\mathbf{N}_c)$ gauge sector with \mathbf{N}_f fermions in the \mathbf{N}_r representation of the gauge group

Most of the theory work is done using EFTs at the SM scale, but we want to use UV complete models

What is composite dark matter?



What is composite dark matter?



- ◆ Dark Matter is a **composite** object

What is composite dark matter?



- ◆ Dark Matter is a **composite** object

e.g. hidden sector
baryon or glueball

What is composite dark matter?



e.g. hidden sector
baryon or glueball

- ◆ Dark Matter is a **composite** object
- ◆ Interesting and complicated internal structure
- ◆ Properties dictated by strong dynamics
- ◆ Self-interactions are natural

What is composite dark matter?



- ◆ Dark Matter is a **composite** object
- ◆ Interesting and complicated internal structure
- ◆ Properties dictated by strong dynamics
- ◆ Self-interactions are natural

e.g. hidden sector
baryon or glueball

Similar to QCD

What is composite dark matter?



- ◆ Dark Matter is a **composite** object
- ◆ Interesting and complicated internal **structure**
- ◆ Properties dictated by **strong dynamics**
- ◆ Self-interactions are natural
- ◆ Composite object is **neutral**
- ◆ Constituents may **interact with Standard Model** particles

e.g. hidden **sector baryon** or **glueball**

Similar to **QCD**

What is composite dark matter?



- ◆ Dark Matter is a **composite** object
- ◆ Interesting and complicated internal **structure**
- ◆ Properties dictated by **strong dynamics**
- ◆ Self-interactions are natural
- ◆ Composite object is **neutral**
- ◆ Constituents may **interact with Standard Model** particles

e.g. hidden **sector baryon** or **glueball**

Similar to **QCD**

Chance to **observe them in experiments** and give the correct **relic abundance**

What is composite dark matter?



◆ Dark Matter is a **composite** object

e.g. hidden sector
baryon or glueball

◆ Interesting and complicated internal
structure

Lattice Field Theory methods

◆ Properties dictated by **strong dynamics**

Similar to **QCD**

◆ **Self-interactions** are natural

◆ Composite object is **neutral**

Chance to **observe them**
in **experiments** and give the
correct **relic abundance**

◆ Constituents may **interact with Standard
Model** particles

“Natural” features



“Natural” features

Stability is a direct
consequence of
accidental **symmetries**



“Natural” features

Stability is a direct consequence of accidental **symmetries**

Neutrality follows naturally from **confinement** into singlet objects wrt. SM charges

“Natural” features

Stability is a direct consequence of accidental **symmetries**

Neutrality follows naturally from **confinement** into singlet objects wrt. SM charges

Small **interactions** with SM particles arise from form factor **suppression** (higher dim. operators)



“Natural” features

Stability is a direct consequence of accidental **symmetries**

Neutrality follows naturally from **confinement** into singlet objects wrt. SM charges

Small **interactions** with SM particles arise from form factor **suppression** (higher dim. operators)

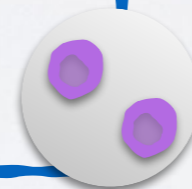
Self-interactions are included due to **strongly coupled** dynamics

COMPOSITE DARK MATTER

COMPOSITE DARK MATTER

★ Pion-like (dark quark-antiquark)

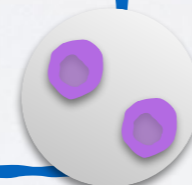
- ◆ pNGB DM [*Hietanen et al., 1308.4130*]
- ◆ Quirky DM [*Kribs et al., 0909.2034*]
- ◆ Ectocolor DM [*Buckley&Neil, 1209.6054*]
- ◆ SIMP [*Hochberg et al., 1411.3727*]
- ◆ Minimal SU(2) [*Francis et al., 1610.10068*]



COMPOSITE DARK MATTER

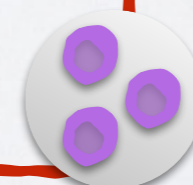
★ Pion-like (dark quark-antiquark)

- ◆ pNGB DM [*Hietanen et al.*, 1308.4130]
- ◆ Quirky DM [*Kribs et al.*, 0909.2034]
- ◆ Ectocolor DM [*Buckley&Neil*, 1209.6054]
- ◆ SIMP [*Hochberg et al.*, 1411.3727]
- ◆ Minimal SU(2) [*Francis et al.*, 1610.10068]



★ Baryon-like (multiple quarks)

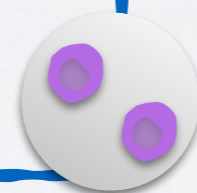
- ◆ “Technibaryons” [*LSD*, 1301.1693]
- ◆ Stealth DM [*LSD*, 1503.04203-1503.04205]
- ◆ One-family TC [*LatKMI*, 1510.07373]
- ◆ Sextet CH [*LatHC*, 1601.03302]



COMPOSITE DARK MATTER

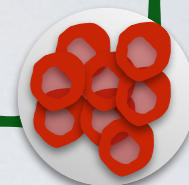
★ Pion-like (dark quark-antiquark)

- ◆ pNGB DM [*Hietanen et al., 1308.4130*]
- ◆ Quirky DM [*Kribs et al., 0909.2034*]
- ◆ Ectocolor DM [*Buckley&Neil, 1209.6054*]
- ◆ SIMP [*Hochberg et al., 1411.3727*]
- ◆ Minimal SU(2) [*Francis et al., 1610.10068*]



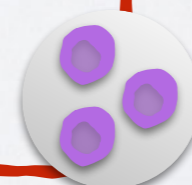
★ Glueball-like (only gluons)

- ◆ SUNonia [*Boddy et al., 1402.3629*]



★ Baryon-like (multiple quarks)

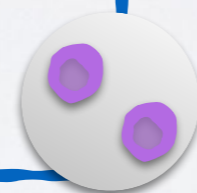
- ◆ “Technibaryons” [*LSD, 1301.1693*]
- ◆ Stealth DM [*LSD, 1503.04203-1503.04205*]
- ◆ One-family TC [*LatKMI, 1510.07373*]
- ◆ Sextet CH [*LatHC, 1601.03302*]



COMPOSITE DARK MATTER

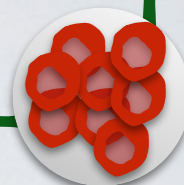
★ Pion-like (dark quark-antiquark)

- ◆ pNGB DM [*Hietanen et al.*, 1308.4130]
- ◆ Quirky DM [*Kribs et al.*, 0909.2034]
- ◆ Ectocolor DM [*Buckley&Neil*, 1209.6054]
- ◆ SIMP [*Hochberg et al.*, 1411.3727]
- ◆ Minimal SU(2) [*Francis et al.*, 1610.10068]



★ Glueball-like (only gluons)

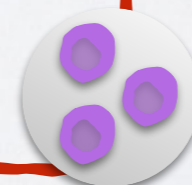
- ◆ SUNonia [*Boddy et al.*, 1402.3629]



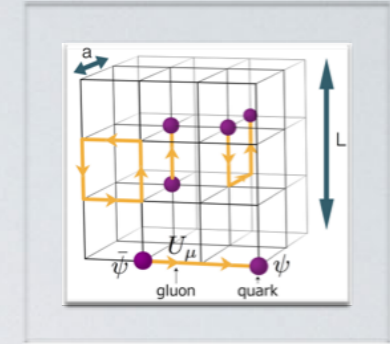
★ Dark Nuclei [*Detmold et al.*, 1406.2276-1406.4116]

★ Baryon-like (multiple quarks)

- ◆ “Technibaryons” [*LSD*, 1301.1693]
- ◆ Stealth DM [*LSD*, 1503.04203-1503.04205]
- ◆ One-family TC [*LatKMI*, 1510.07373]
- ◆ Sextet CH [*LatHC*, 1601.03302]



Importance of lattice simulations



- ★ Lattice simulations are needed to numerically solve strong dynamics
- ★ **Controllable** systematic errors and room for improvement
- ★ Naive dimensional analysis and EFT approaches can miss important non-perturbative contributions
- ★ **EFTs** inspired by QCD might not work when the dynamics is different
- ★ Lattice studies can reliably point out **similarities or differences** as the parameter space (N_c, N_f, N_r) is scanned

“Stealth Dark Matter” model

[LSD collab., Phys. Rev. D92 (2015) 075030]



- ◆ **New strongly-coupled SU(4) gauge sector** “like” QCD with a **plethora of composite states** in the spectrum: all mass scales are technically natural for hadrons
- ◆ New **Dark fermions**: have **dark color** and also have **electroweak charges** ($W/Z, \gamma$)
- ◆ Dark fermions have **electroweak breaking masses** (Yukawa couplings) and **electroweak preserving masses** (from confinement)
- ◆ A global symmetry naturally stabilizes the **dark lightest baryonic** composite states (e.g. dark U(1) “baryon number”) which is a singlet of 4 dark fermions: **spin 0 (!)**

“Stealth Dark Matter” model

[LSD collab., Phys. Rev. D92 (2015) 075030]

- The field content of the model consists in **8 Weyl fermions**
- Dark fermions interact with the SM Higgs and obtain **current/chiral masses**
- Introduce **vector-like masses** for dark fermions that do not break EW symmetry
- Diagonalizing in the mass eigenbasis gives **4 Dirac fermions**
- Assume **custodial SU(2) symmetry** arising when **$u \leftrightarrow d$**

Field	$SU(N)_D$	$(SU(2)_L, Y)$	Q
$F_1 = \begin{pmatrix} F_1^u \\ F_1^d \end{pmatrix}$	\mathbf{N}	$(\mathbf{2}, 0)$	$\begin{pmatrix} +1/2 \\ -1/2 \end{pmatrix}$
$F_2 = \begin{pmatrix} F_2^u \\ F_2^d \end{pmatrix}$	$\overline{\mathbf{N}}$	$(\mathbf{2}, 0)$	$\begin{pmatrix} +1/2 \\ -1/2 \end{pmatrix}$
F_3^u	\mathbf{N}	$(\mathbf{1}, +1/2)$	$+1/2$
F_3^d	\mathbf{N}	$(\mathbf{1}, -1/2)$	$-1/2$
F_4^u	$\overline{\mathbf{N}}$	$(\mathbf{1}, +1/2)$	$+1/2$
F_4^d	$\overline{\mathbf{N}}$	$(\mathbf{1}, -1/2)$	$-1/2$

$$\mathcal{L} \supset + y_{14}^u \epsilon_{ij} F_1^i H^j F_4^d + y_{14}^d F_1 \cdot H^\dagger F_4^u - y_{23}^d \epsilon_{ij} F_2^i H^j F_3^d - y_{23}^u F_2 \cdot H^\dagger F_3^u + h.c.$$

$$\mathcal{L} \supset M_{12} \epsilon_{ij} F_1^i F_2^j - M_{34}^u F_3^u F_4^d + M_{34}^d F_3^d F_4^u + h.c.$$

$$y_{14}^u = y_{14}^d \quad y_{23}^u = y_{23}^d \quad M_{34}^u = M_{34}^d$$

“Stealth Dark Matter” model

[LSD collab., Phys. Rev. D92 (2015) 075030]

EW interactions

- The field content of the model consists in **8 Weyl fermions**
- Dark fermions interact with the SM Higgs and obtain **current/chiral masses**
- Introduce **vector-like masses** for dark fermions that do not break EW symmetry
- Diagonalizing in the mass eigenbasis gives **4 Dirac fermions**
- Assume **custodial SU(2) symmetry** arising when **$u \leftrightarrow d$**

Field	$SU(N)_D$	$(SU(2)_L, Y)$	Q
$F_1 = \begin{pmatrix} F_1^u \\ F_1^d \end{pmatrix}$	\mathbf{N}	$(\mathbf{2}, 0)$	$\begin{pmatrix} +1/2 \\ -1/2 \end{pmatrix}$
$F_2 = \begin{pmatrix} F_2^u \\ F_2^d \end{pmatrix}$	$\overline{\mathbf{N}}$	$(\mathbf{2}, 0)$	$\begin{pmatrix} +1/2 \\ -1/2 \end{pmatrix}$
F_3^u	\mathbf{N}	$(\mathbf{1}, +1/2)$	$+1/2$
F_3^d	\mathbf{N}	$(\mathbf{1}, -1/2)$	$-1/2$
F_4^u	$\overline{\mathbf{N}}$	$(\mathbf{1}, +1/2)$	$+1/2$
F_4^d	$\overline{\mathbf{N}}$	$(\mathbf{1}, -1/2)$	$-1/2$

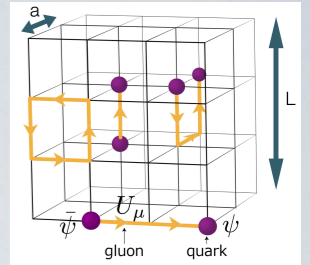
$$\mathcal{L} \supset + y_{14}^u \epsilon_{ij} F_1^i H^j F_4^d + y_{14}^d F_1 \cdot H^\dagger F_4^u - y_{23}^d \epsilon_{ij} F_2^i H^j F_3^d - y_{23}^u F_2 \cdot H^\dagger F_3^u + h.c.$$

$$\mathcal{L} \supset M_{12} \epsilon_{ij} F_1^i F_2^j - M_{34}^u F_3^u F_4^d + M_{34}^d F_3^d F_4^u + h.c.$$

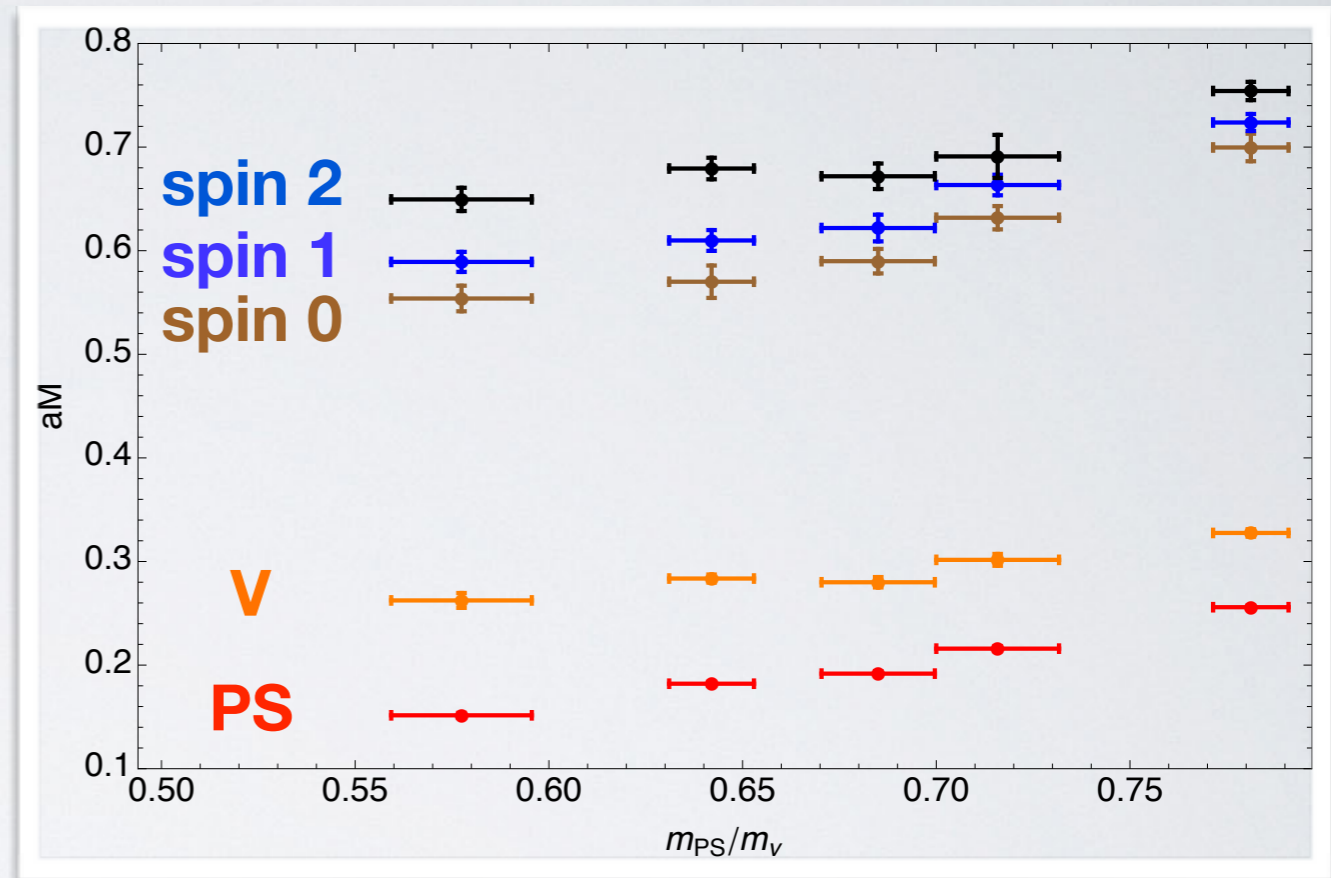
$$y_{14}^u = y_{14}^d \quad y_{23}^u = y_{23}^d \quad M_{34}^u = M_{34}^d$$

Lattice Stealth DM

[KEK-Japan]



- Non-perturbative lattice calculations of the spectrum confirm that **lightest baryon has spin zero**
- The ratio of **pseudoscalar (PS, pion)** to **vector (V, rho)** is used as probe for different dark fermion masses
- The meson-to-baryon mass ratio is a **non-perturbative number** which can only be extracted from lattice simulations

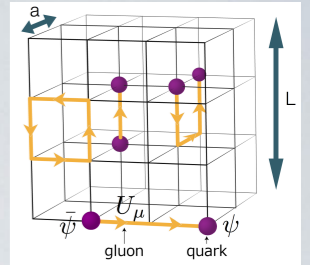


[LSD collab., Phys. Rev. D89 (2014) 094508]

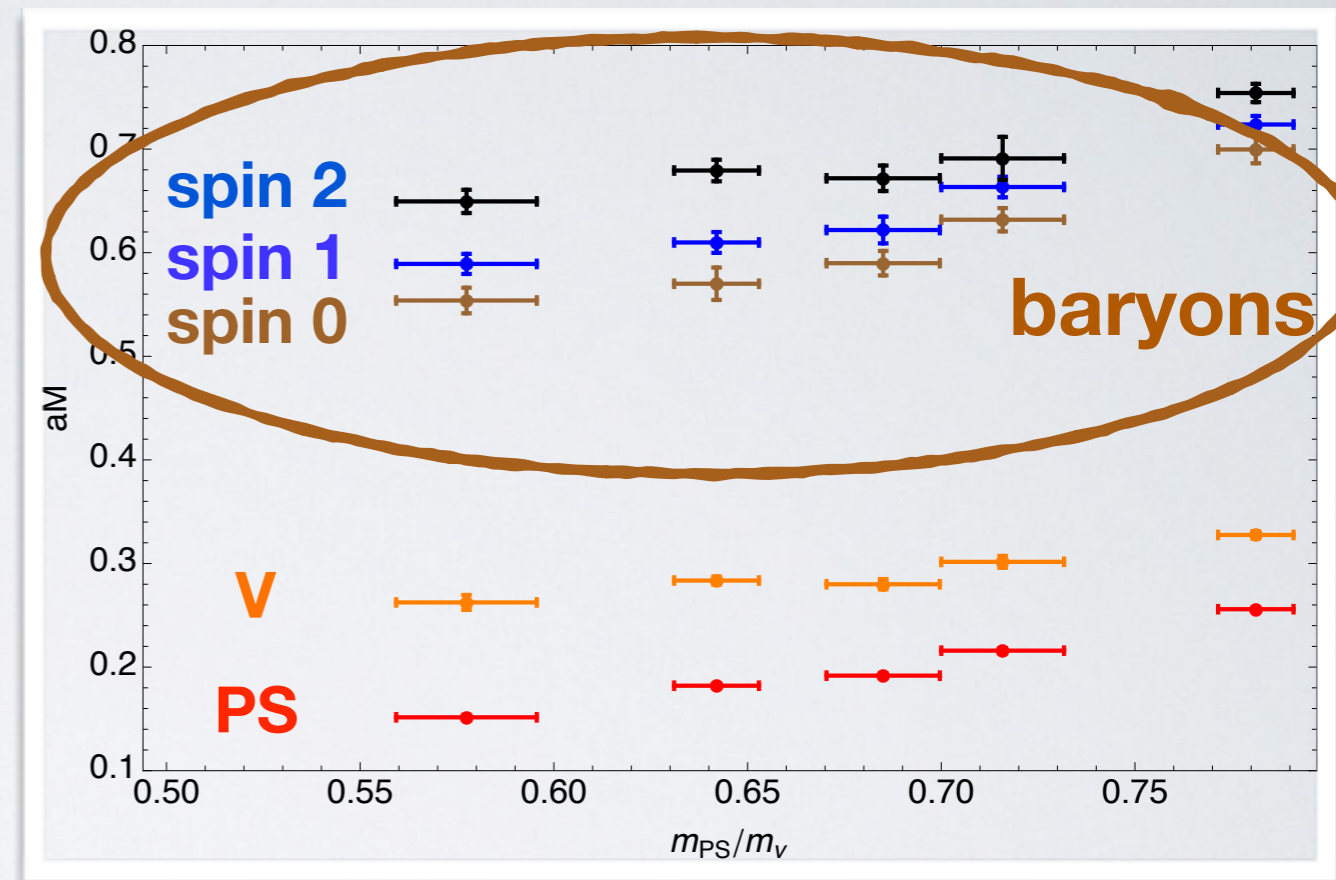
- Lattice discretization and finite volume effects are studied using multiple simulations (similar to what is done in QCD)

Lattice Stealth DM

[KEK-Japan]



- Non-perturbative lattice calculations of the spectrum confirm that **lightest baryon has spin zero**
- The ratio of **pseudoscalar (PS, pion)** to **vector (V, rho)** is used as probe for different dark fermion masses
- The meson-to-baryon mass ratio is a **non-perturbative number** which can only be extracted from lattice simulations

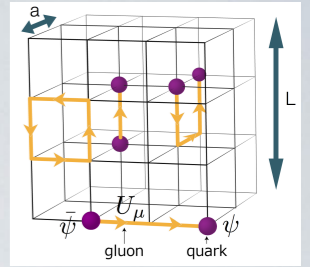


[LSD collab., Phys. Rev. D89 (2014) 094508]

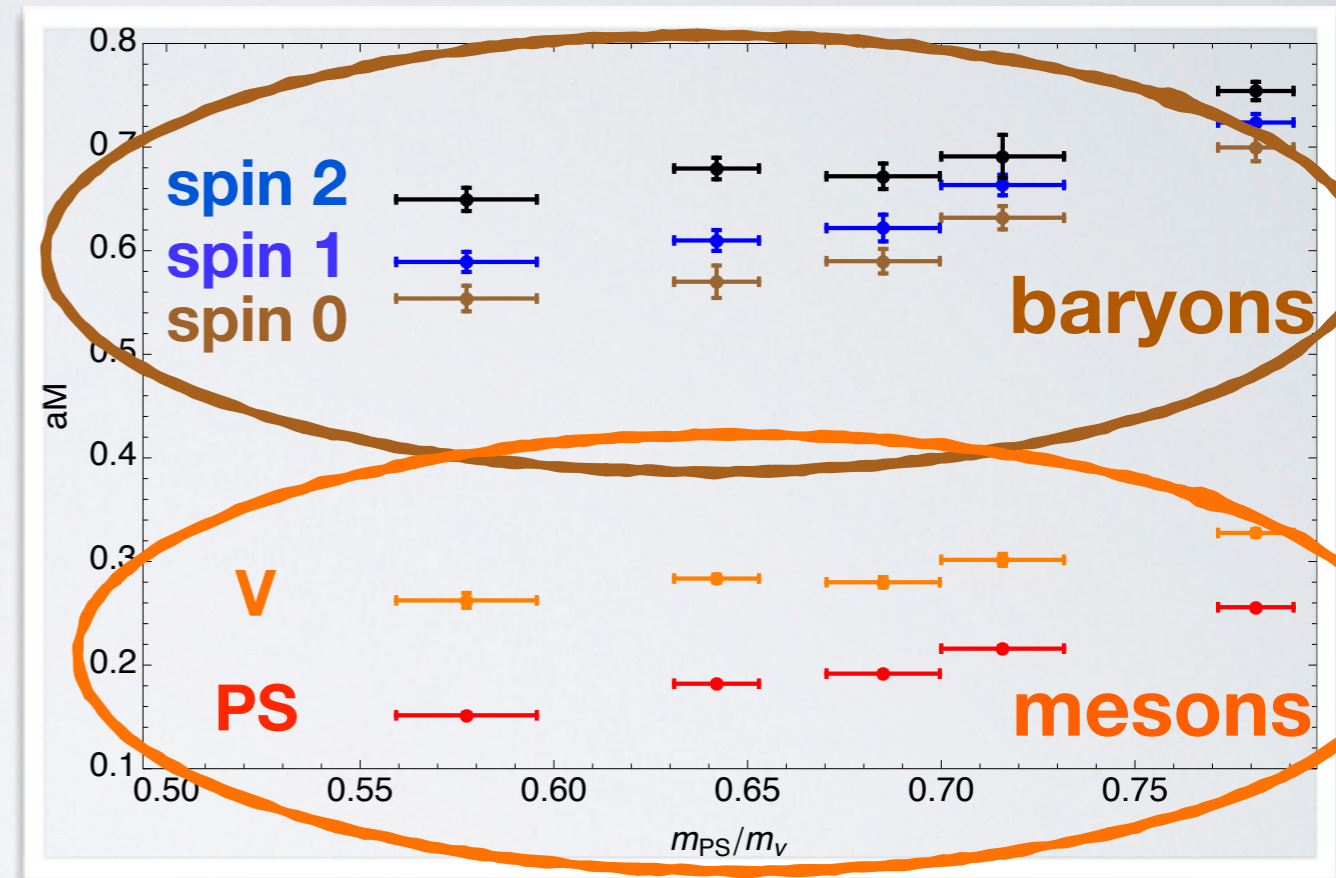
- Lattice discretization and finite volume effects are studied using multiple simulations (similar to what is done in QCD)

Lattice Stealth DM

[KEK-Japan]

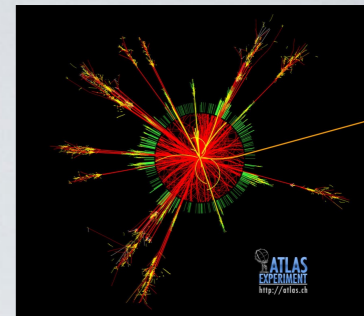


- Non-perturbative lattice calculations of the spectrum confirm that **lightest baryon has spin zero**
- The ratio of **pseudoscalar (PS, pion)** to **vector (V, rho)** is used as probe for different dark fermion masses
- The meson-to-baryon mass ratio is a **non-perturbative number** which can only be extracted from lattice simulations



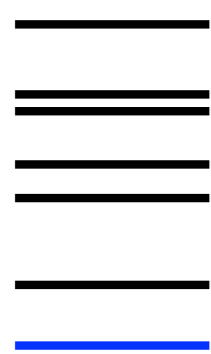
[LSD collab., Phys. Rev. D89 (2014) 094508]

- Lattice discretization and finite volume effects are studied using multiple simulations (similar to what is done in QCD)



Stealth Dark Matter at colliders

SUSY

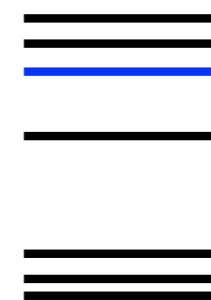


heavier
superpartners

LSP

VS.

Stealth



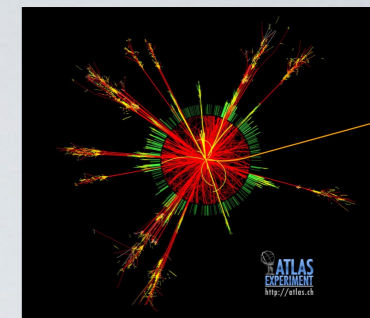
baryon excited
resonances

scalar baryon

ρ

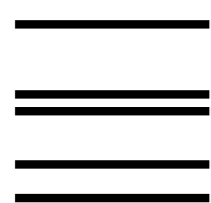
Π_s

Plot by G. Kribs



Stealth Dark Matter at colliders

SUSY



heavier
superpartners



LSP

vs.

Stealth



baryon excited
resonances



ρ



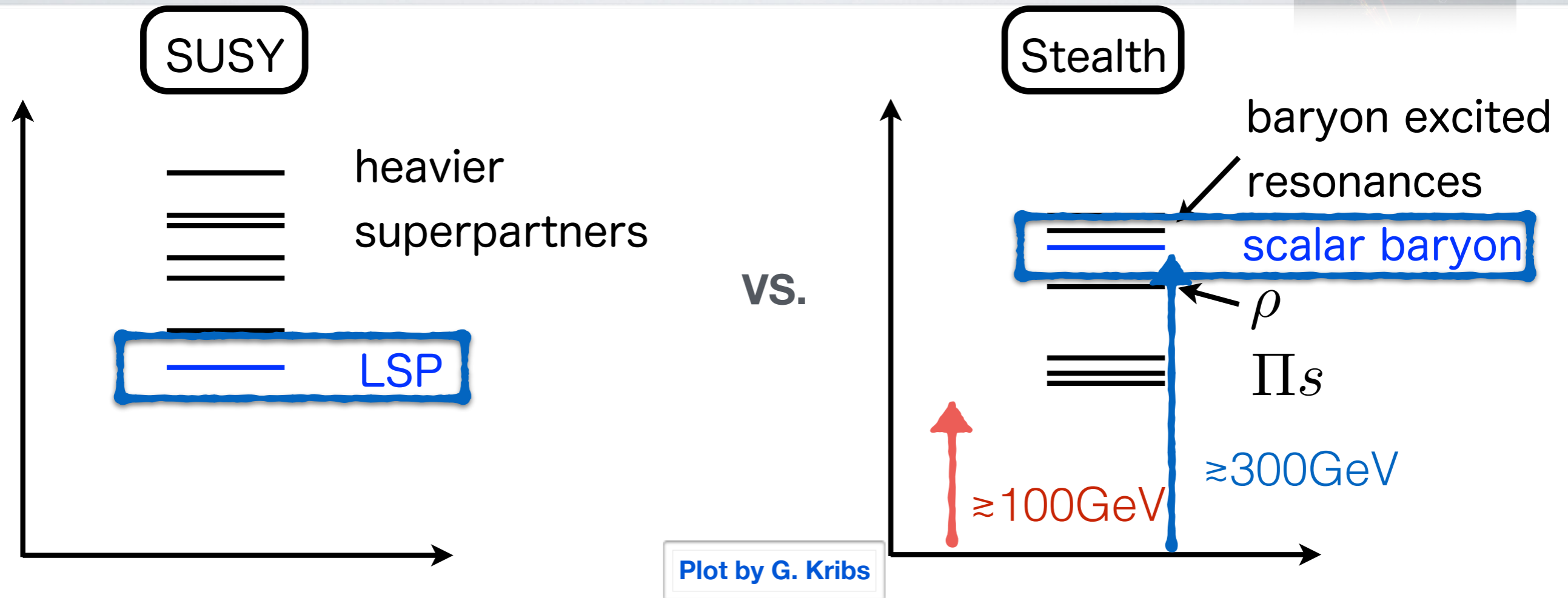
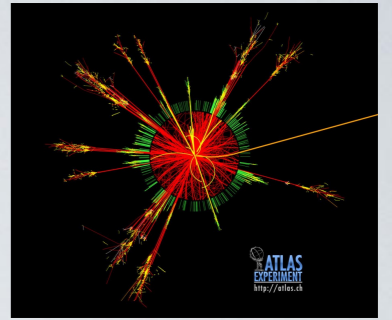
Π_s

Plot by G. Kribs

- ◆ Signatures are not dominated by missing energy: **DM is not the lightest particle!** The interactions are suppressed (form factors)

Stealth Dark Matter at colliders

[ATLAS]

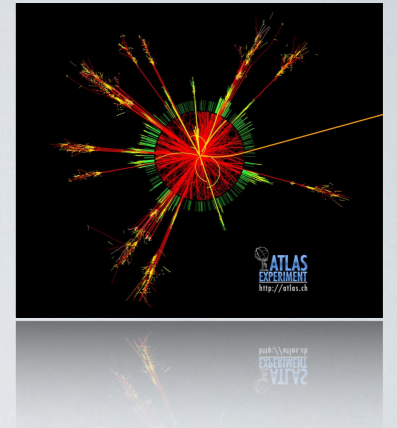


- ◆ Signatures are not dominated by missing energy: **DM is not the lightest particle!** The interactions are suppressed (form factors)
- ◆ Light dark meson production and decay give interesting signatures: **the model can be constrained by LHC/LEP data!** $\Pi \gtrsim 100\text{GeV}$

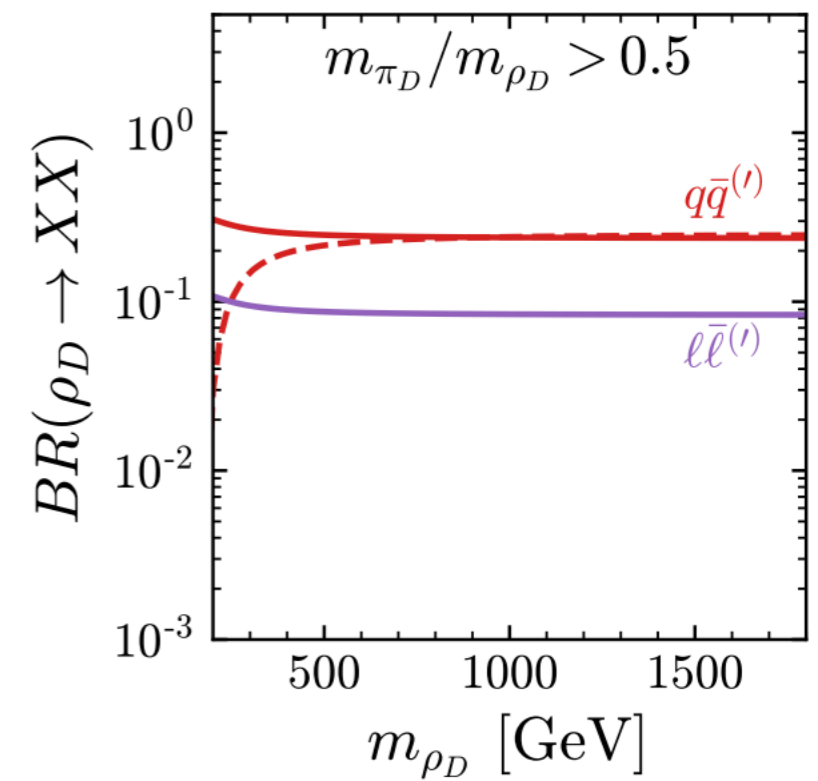
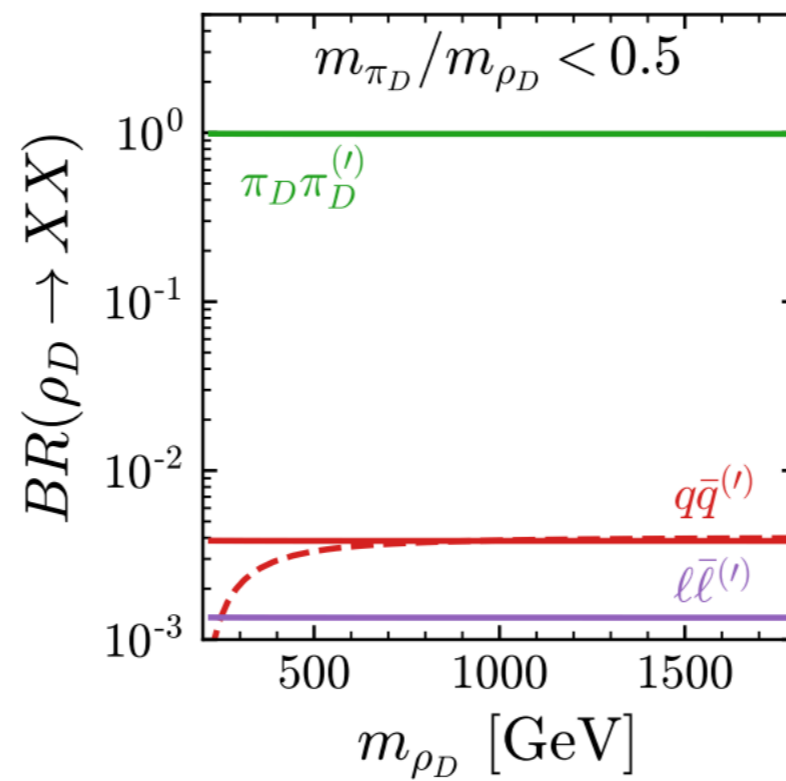
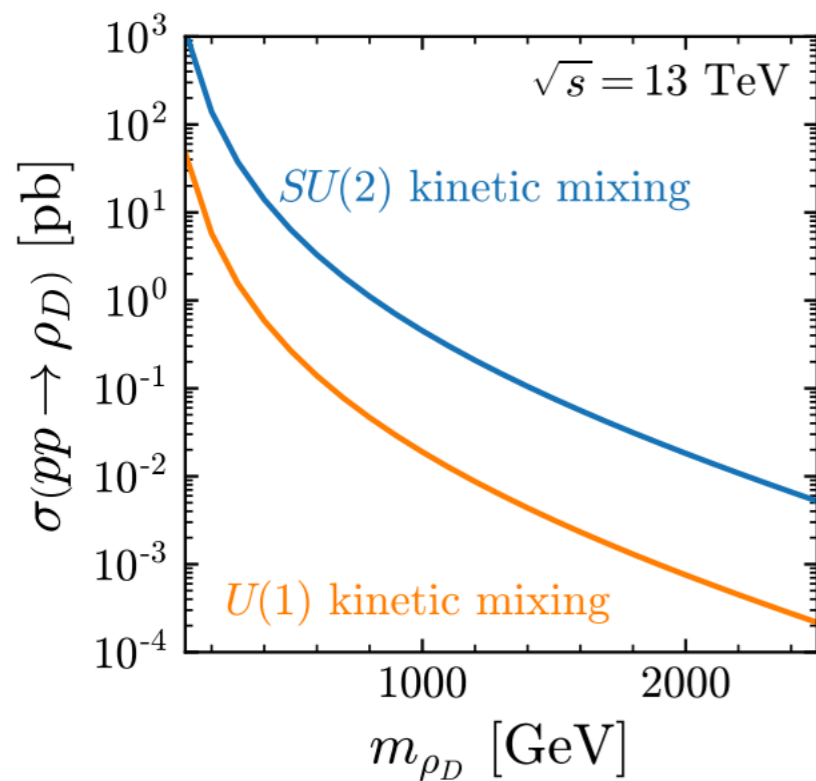
More general bounds on “dark” mesons

[Kribs, Martin, Ostdiek, Tong 1809.10184]

[ATLAS]



- ◆ Dark mesons can be created through Drell-Yan or vector kinetic mixing
- ◆ Dark vector meson can decay to dark “pions” depending on mass ratio
- ◆ Re-cast existing SUSY searches (not optimal)
- ◆ Bounds can be strong depending on model parameters (a few benchmarks)

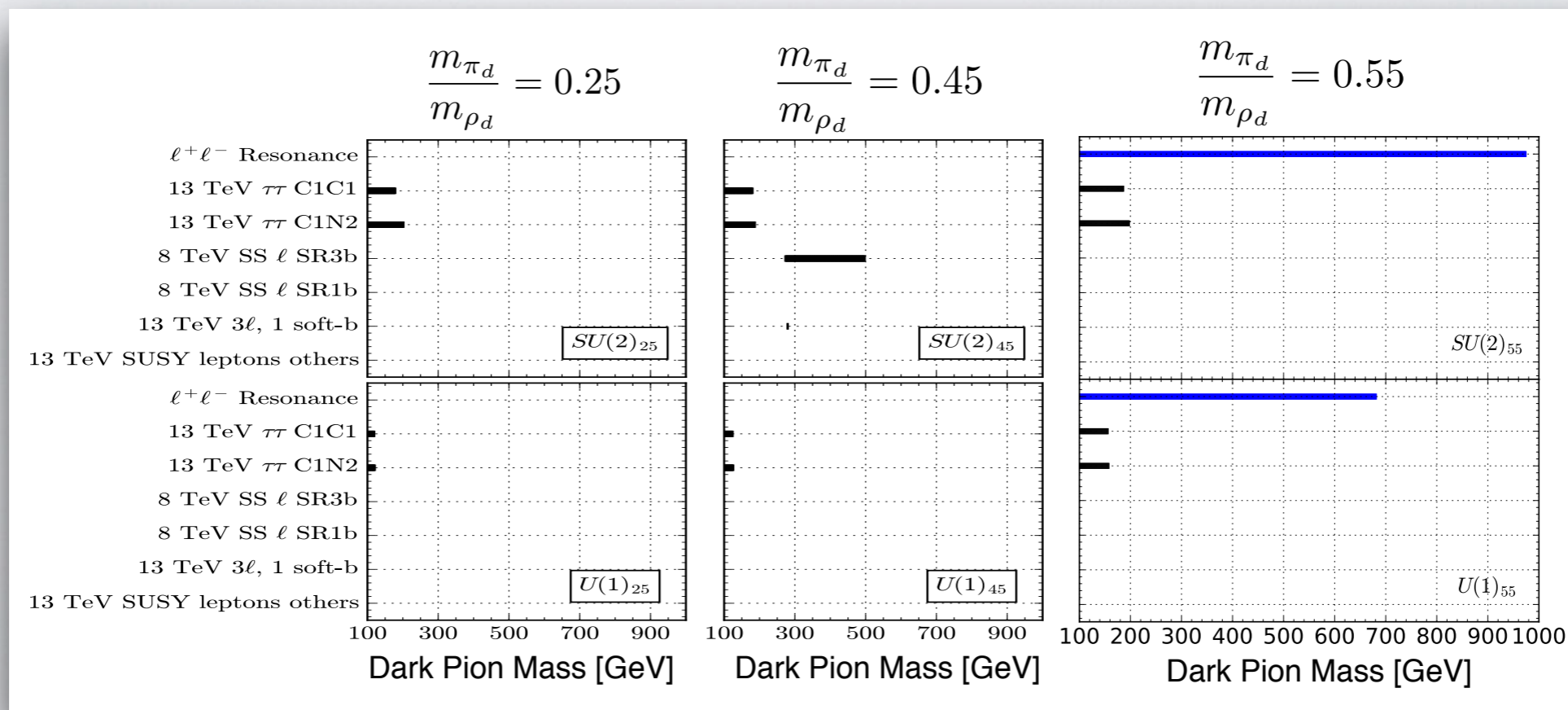


More general bounds on “dark” mesons

[Kribs, Martin, Ostdiek, Tong 1809.10184]

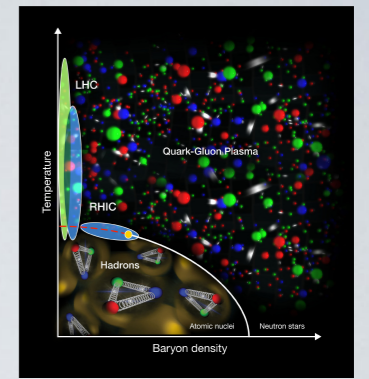


- ◆ Dark mesons can be created through Drell-Yan or vector kinetic mixing
- ◆ Dark vector meson can decay to dark “pions” depending on mass ratio
- ◆ Re-cast existing SUSY searches (not optimal)
- ◆ Bounds can be strong depending on model parameters (a few benchmarks)

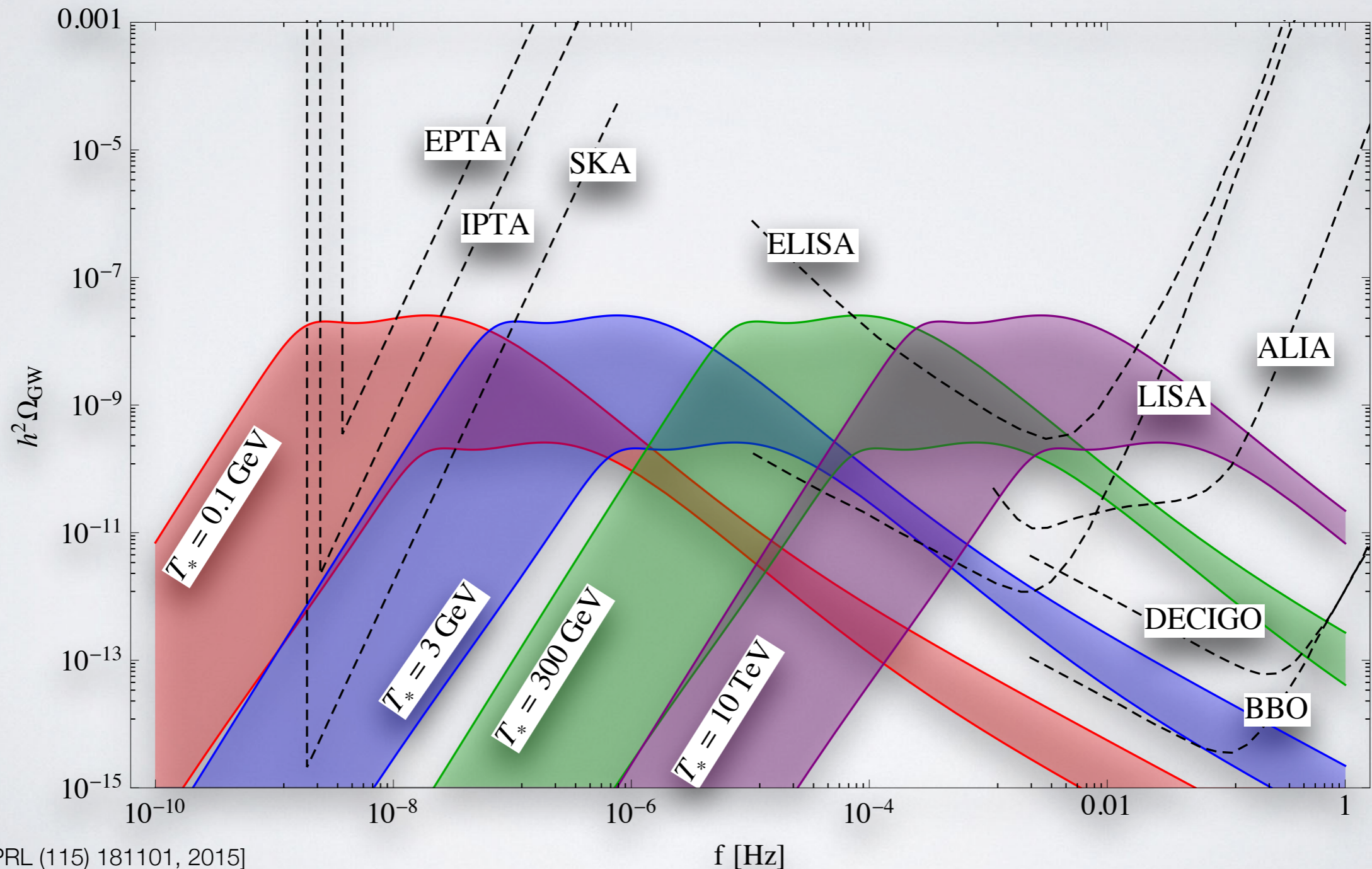


Stealth DM: gravitational wave signatures

[RHIC]

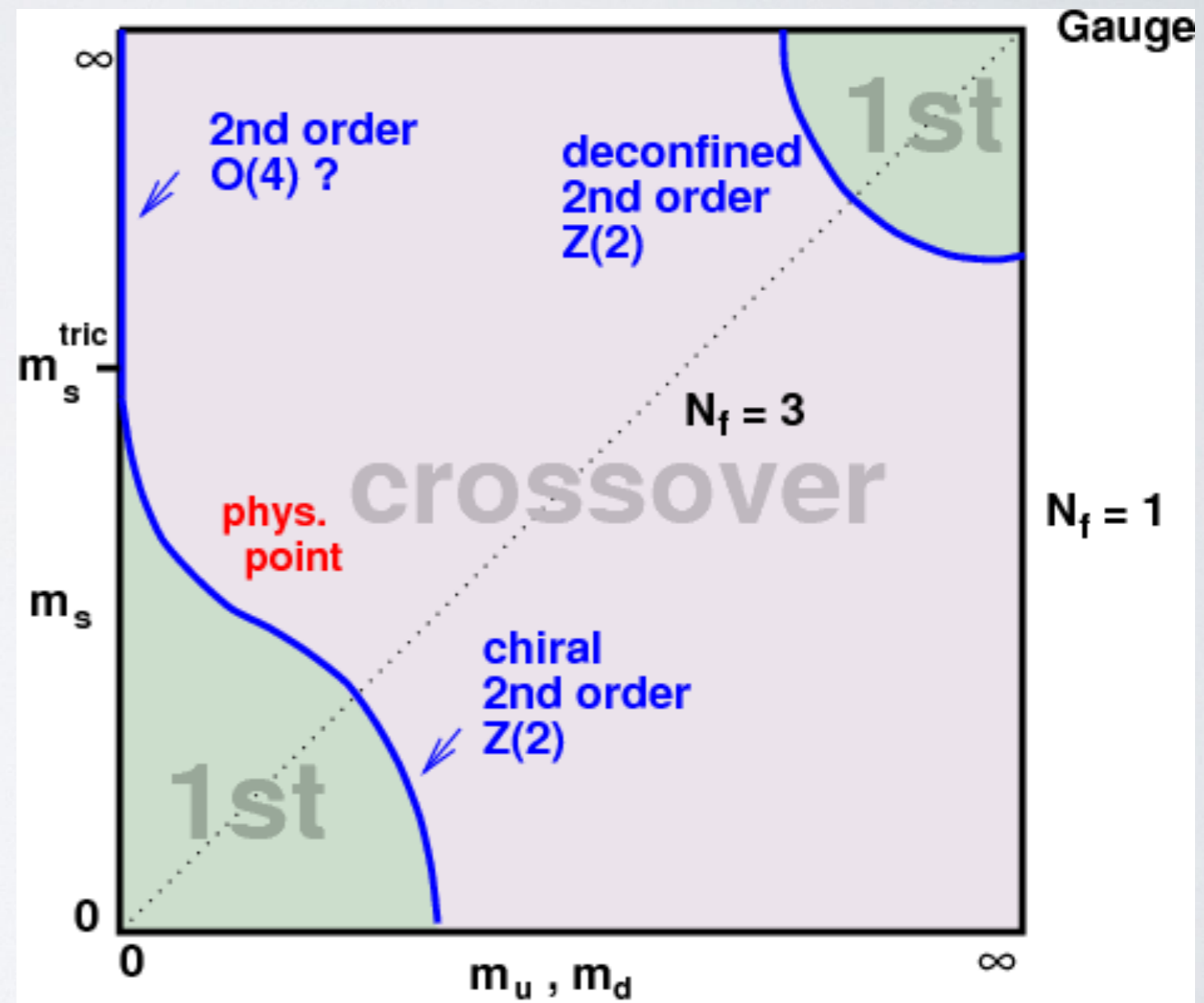


- ◆ Spectrum of GW from a deconfinement first order phase transition in the dark sector, as a function of the dark transition temperature



Phase Transitions in Strongly-coupled Theories

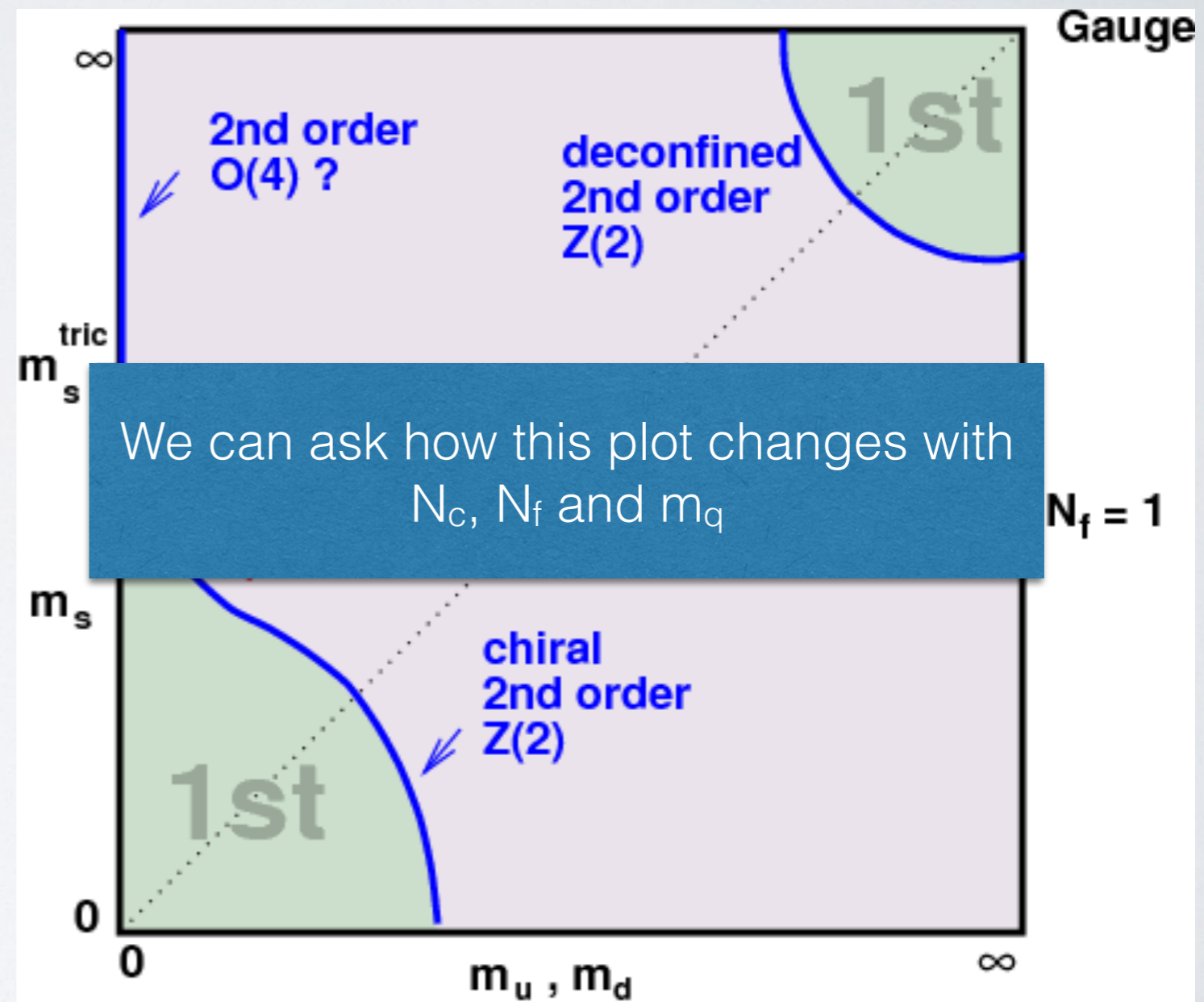
- We know that QCD has a weak cross-over transition between a confined phase and a plasma phase
- We know this is related to the specific values of the u, d and s quark masses
- Our understanding of the QCD phase diagram is heavily based on lattice simulations at finite temperature



“Columbia” plot

Phase Transitions in Strongly-coupled Theories

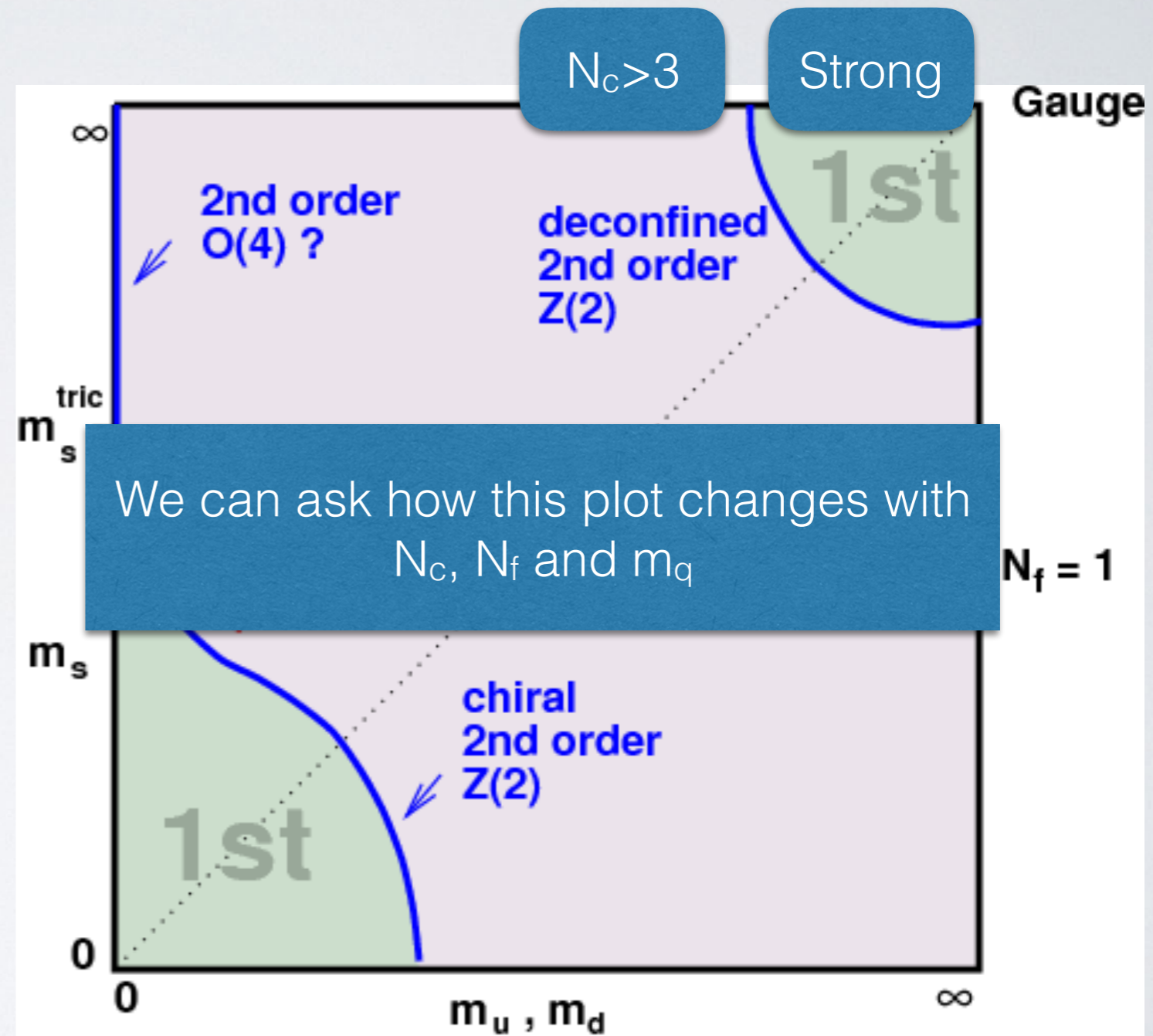
- We know that QCD has a weak cross-over transition between a confined phase and a plasma phase
- We know this is related to the specific values of the u, d and s quark masses
- Our understanding of the QCD phase diagram is heavily based on lattice simulations at finite temperature



“Columbia” plot

Phase Transitions in Strongly-coupled Theories

- We know that QCD has a weak cross-over transition between a confined phase and a plasma phase
- We know this is related to the specific values of the u, d and s quark masses
- Our understanding of the QCD phase diagram is heavily based on lattice simulations at finite temperature

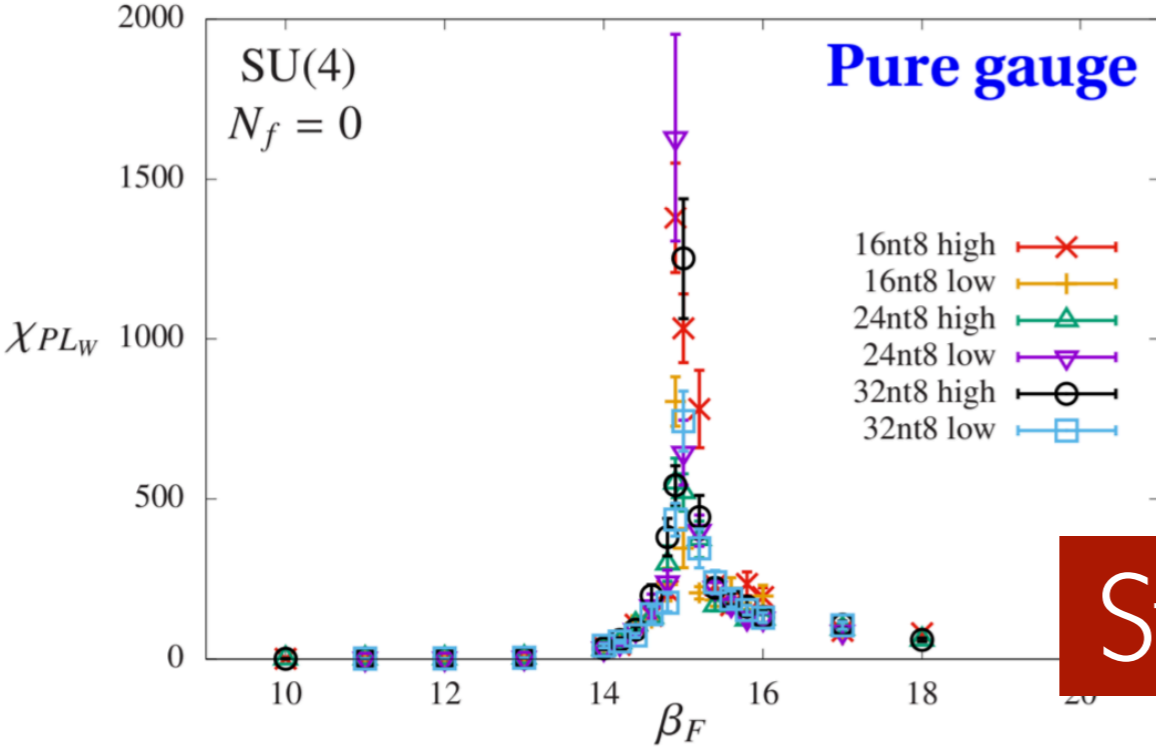
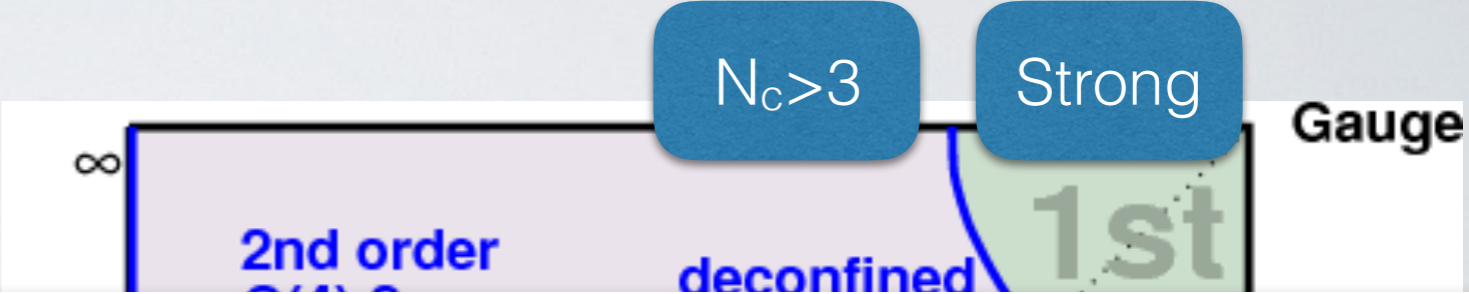


We can ask how this plot changes with N_c , N_f and m_q

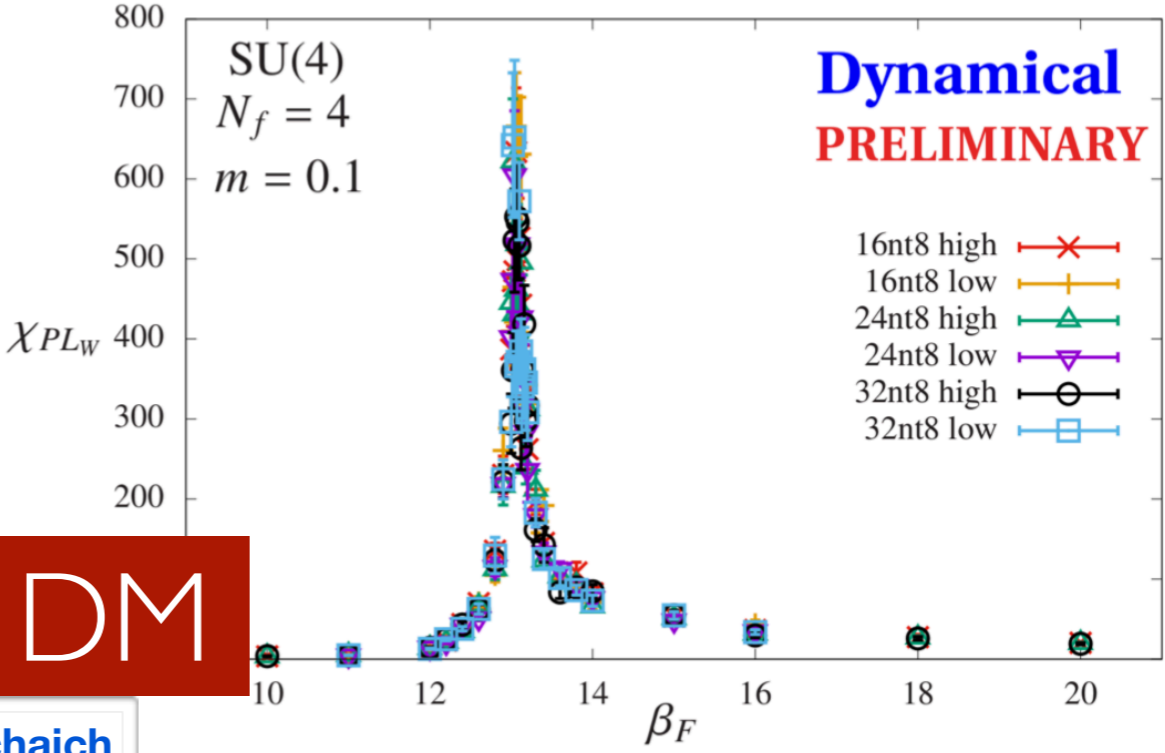
"Columbia" plot

Phase Transitions in Strongly-coupled Theories

- We know that QCD has a weak

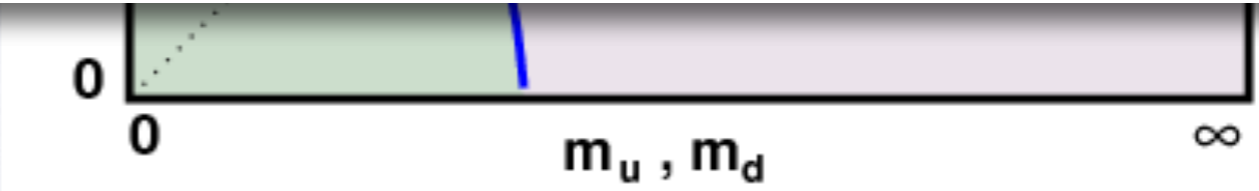


Plot by D. Schaich

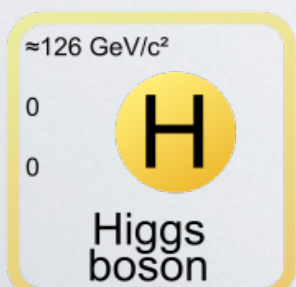
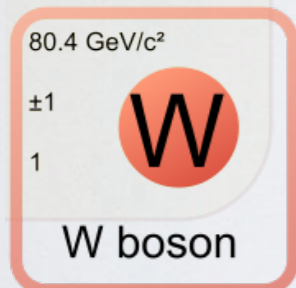
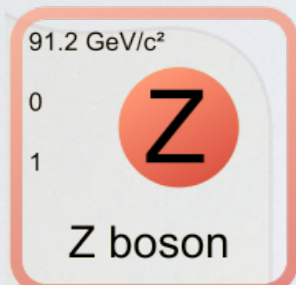


“Columbia” plot

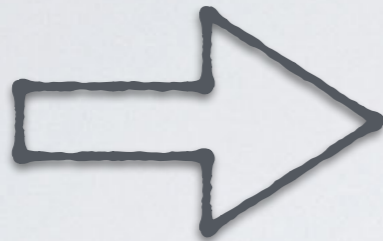
on lattice simulations at finite temperature



The darkness of Composite Dark Matter

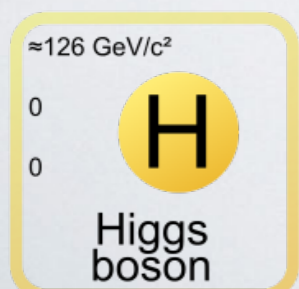
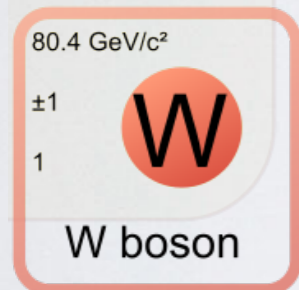
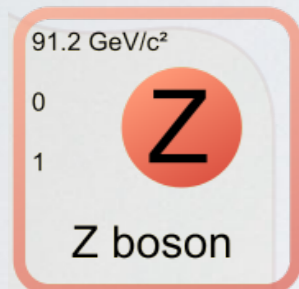


The darkness of Composite Dark Matter

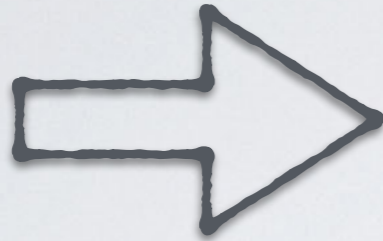


Relevant if the constituents have SM color charges

[Chivukula et al., hep-ph/9210274] [Godbole et al., 1506.01408] [Bay&Osborne, 1506.07110]

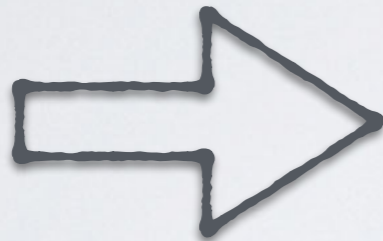


The darkness of Composite Dark Matter



Relevant if the constituents have SM color charges

[Chivukula et al., hep-ph/9210274] [Godbole et al., 1506.01408] [Bay&Osborne, 1506.07110]

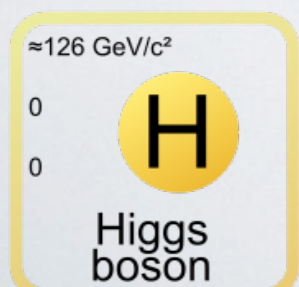
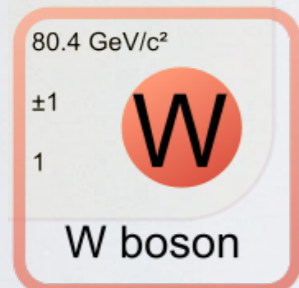
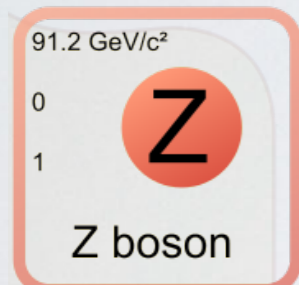


Lowest dimensional operators:

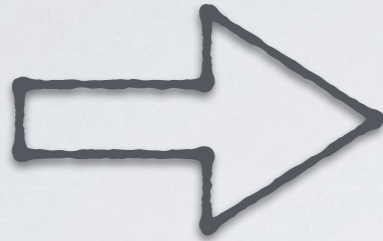
★ magnetic dipole (5)

★ charge radius (6)

★ polarizability (7)

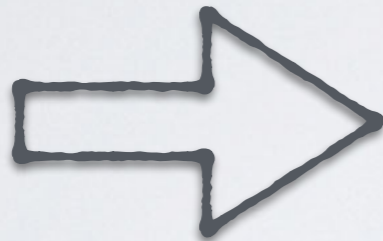


The darkness of Composite Dark Matter



Relevant if the constituents have SM color charges

[Chivukula et al., hep-ph/9210274] [Godbole et al., 1506.01408] [Bay&Osborne, 1506.07110]

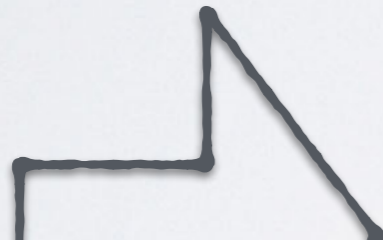
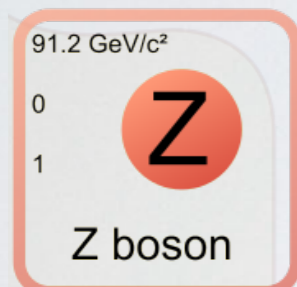


Lowest dimensional operators:

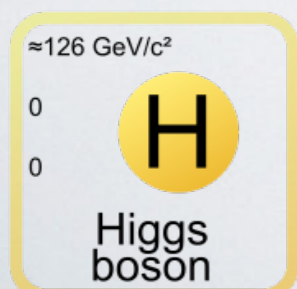
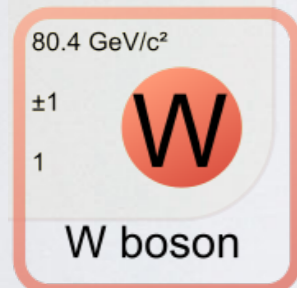
★ magnetic dipole (5)

★ charge radius (6)

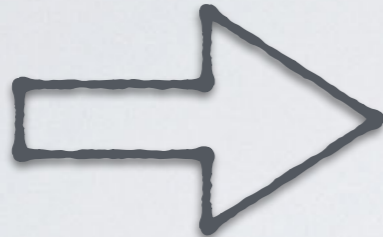
★ polarizability (7)



Same as γ but suppressed due to heavy mass

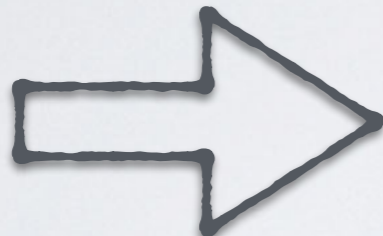


The darkness of Composite Dark Matter



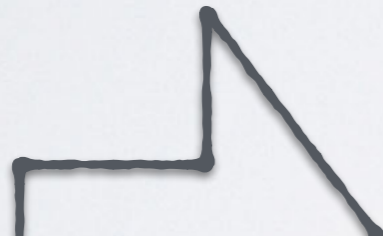
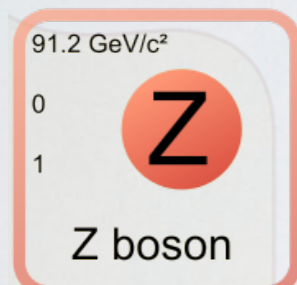
Relevant if the constituents have SM color charges

[Chivukula et al., hep-ph/9210274] [Godbole et al., 1506.01408] [Bay&Osborne, 1506.07110]

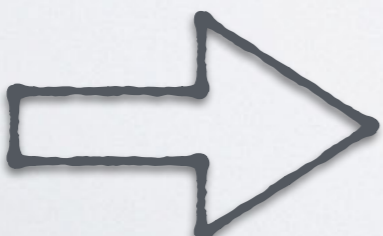
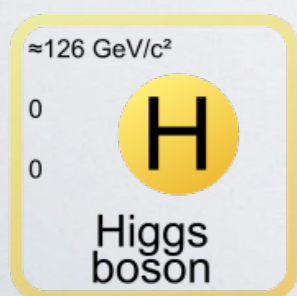
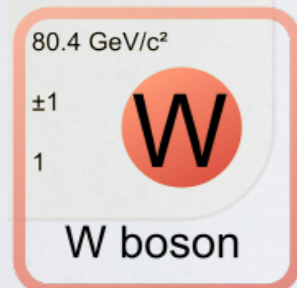


Lowest dimensional operators:

- ★ magnetic dipole (5)
- ★ charge radius (6)
- ★ polarizability (7)

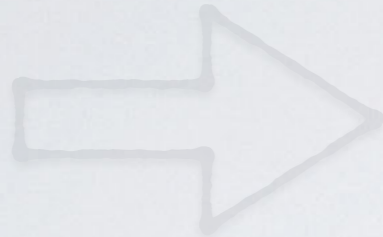


Same as γ but suppressed due to heavy mass



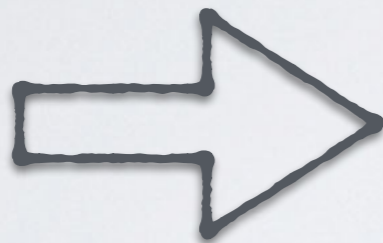
Most relevant interaction if constituents have Yukawa couplings!

The darkness of Composite Dark Matter



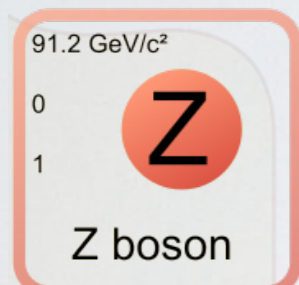
Relevant if the constituents have SM color charges

[Chivukula et al., hep-ph/9210274] [Godbole et al., 1506.01408] [Bay&Osborne, 1506.07110]

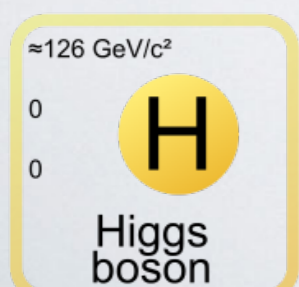
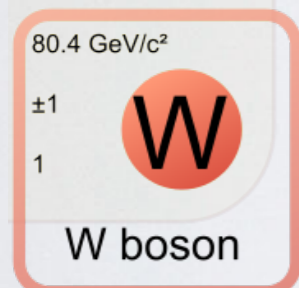


Lowest dimensional operators:

- ★ magnetic dipole (5)
- ★ charge radius (6)
- ★ polarizability (7)



Same as γ but suppressed due to heavy mass



Most relevant interaction if constituents have Yukawa couplings!

LATTICE RESULTS

Template Models

Spectrum

Higgs

Mag. Dip.

Charge r.

Polariz.

SU(2) $N_f=1$ SU(2) $N_f=2$ SU(3) $N_f=2,6$ SU(3) $N_f=8$ SU(3) $N_f=2$ (S)SU(4) $N_f=4$ SO(4) $N_f=2$ (V)SU(N) $N_f=0$ 



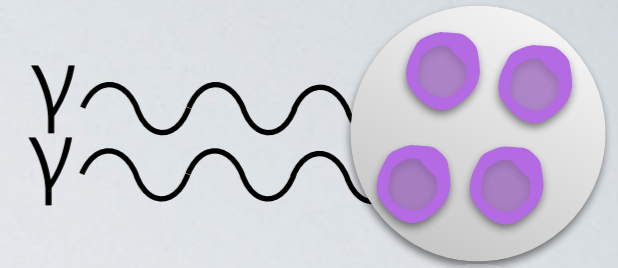
LATTICE RESULTS

Template Models

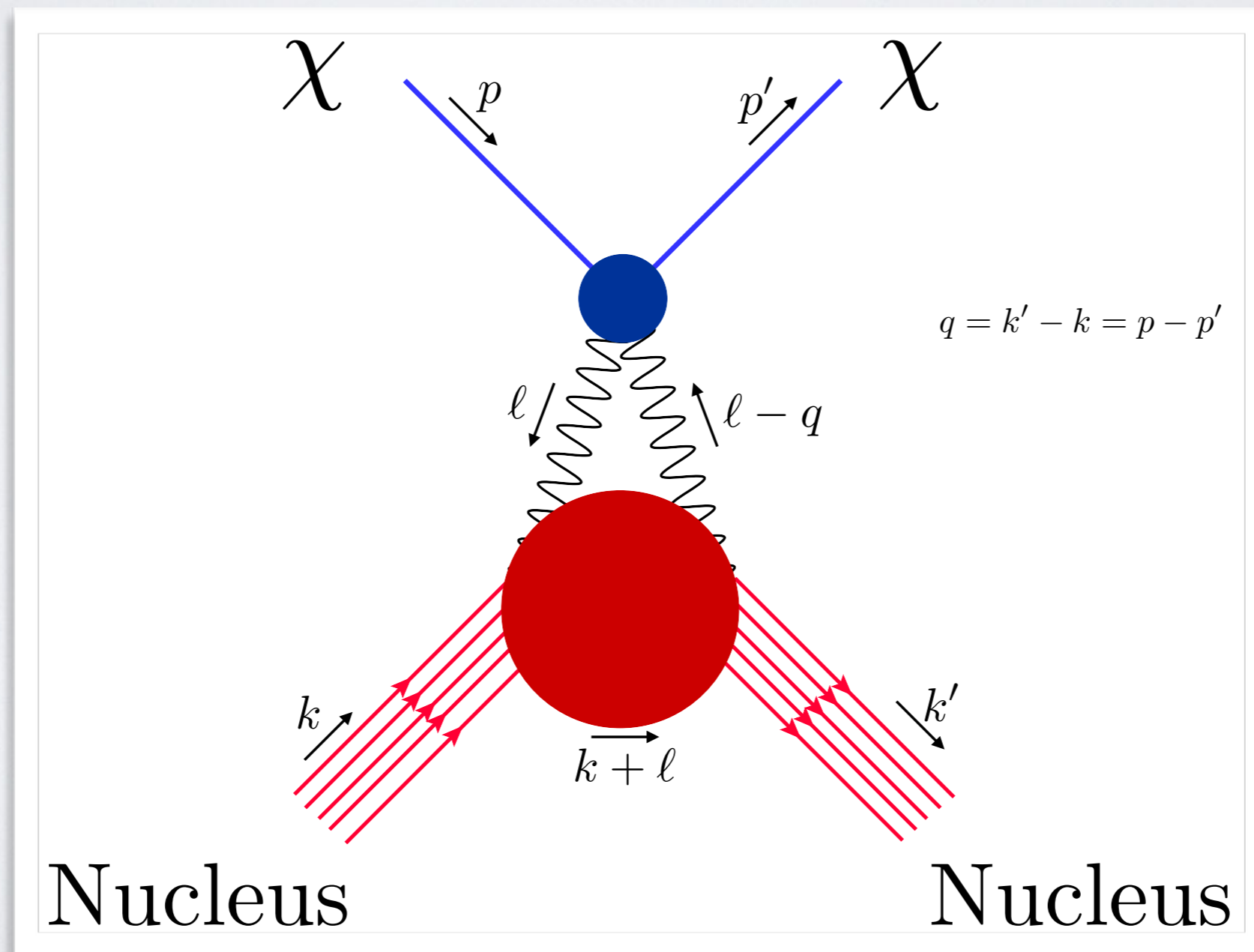
Spectrum Higgs Mag. Dip. Charge r. Polariz.

SU(2) $N_f=1$	★				
SU(2) $N_f=2$	★	★	■ forbidden in pNGB DM	★	★
SU(3) $N_f=2,6$	★		★	★	
SU(3) $N_f=8$	★	★			
SU(3) $N_f=2$ (S)	★				
SU(4) $N_f=4$	★	★	■ forbidden in Stealth DM	■	★
SO(4) $N_f=2$ (V)	★				
SU(N) $N_f=0$	★	■	■ forbidden in SUNonia	■	■

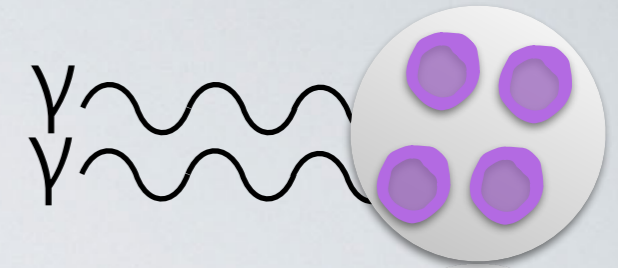
Computing polarizability



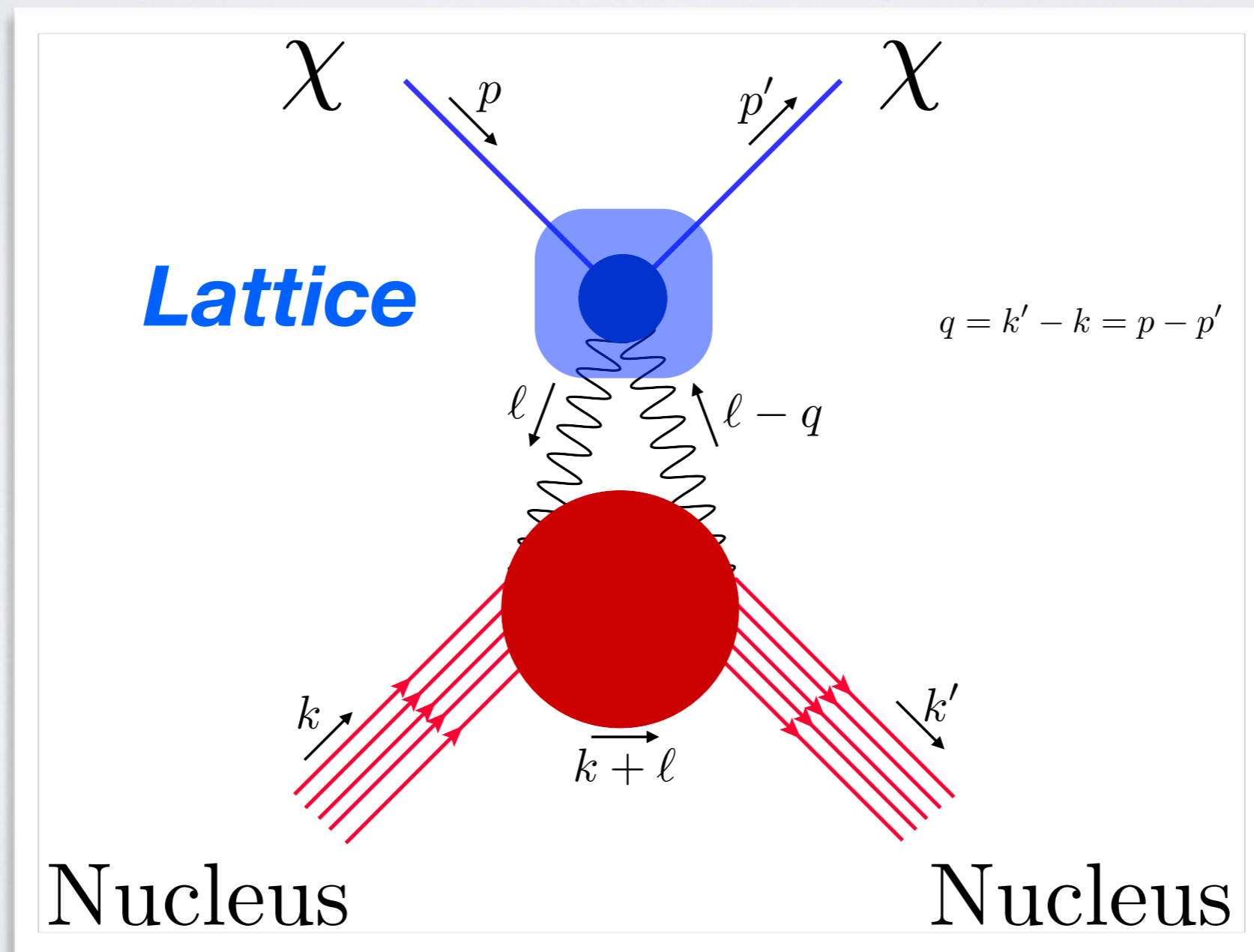
$$\frac{c_F e^2}{m_\chi^3} \chi^* \chi F^{\mu\alpha} F_{\alpha}^{\nu} v_{\mu} v_{\nu}$$



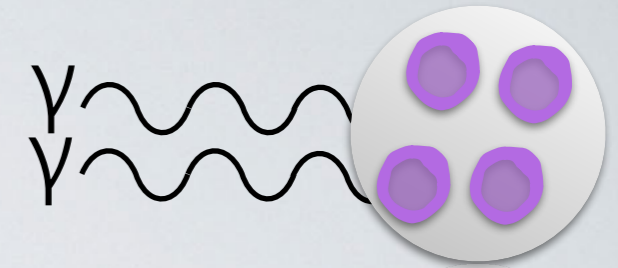
Computing polarizability



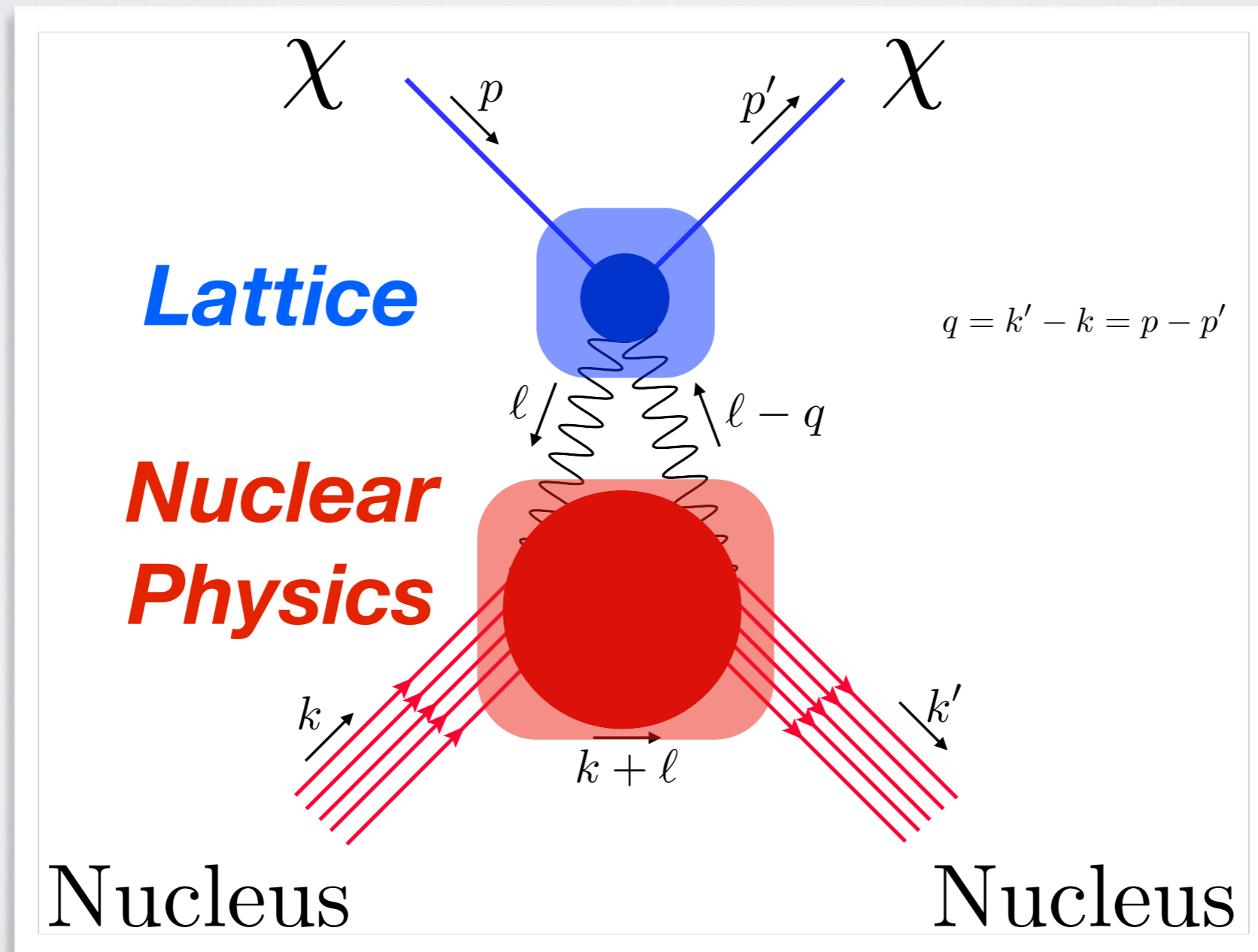
$$\frac{c_F e^2}{m_\chi^3} \chi^* \chi F^{\mu\alpha} F_\alpha^\nu v_\mu v_\nu$$



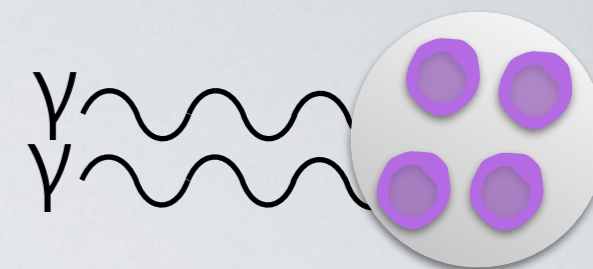
Computing polarizability



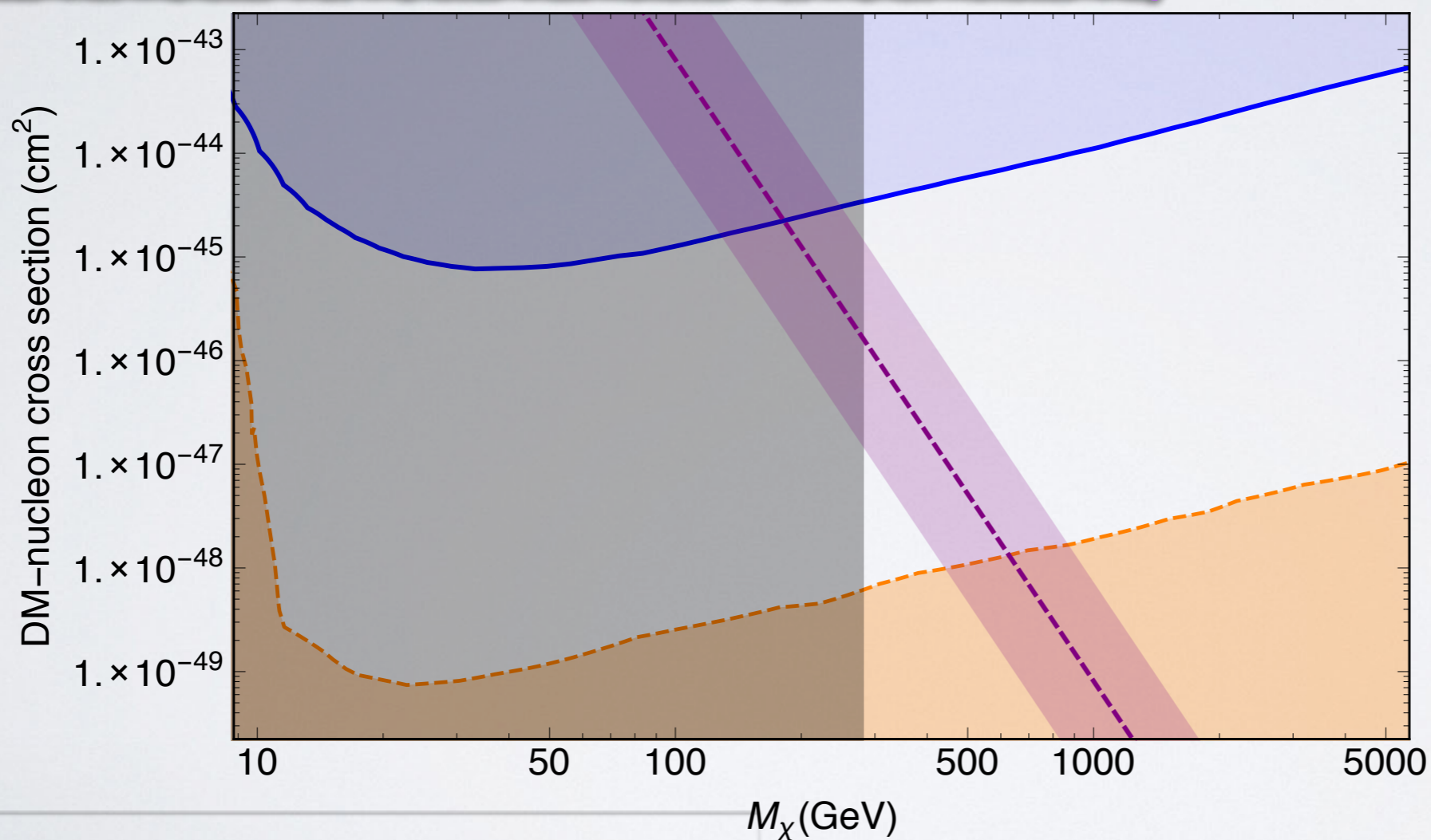
$$\frac{c_F e^2}{m_\chi^3} \chi^* \chi F^{\mu\alpha} F_{\alpha}^{\nu} v_\mu v_\nu$$



Lowest bound from EM polarizability



Electric polarizability from lattice simulations

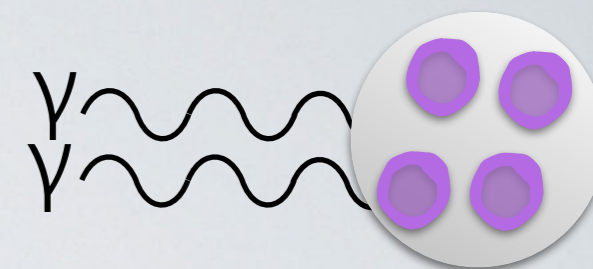


SU(4) $N_f=4$ Stealth DM

[LSD, 1503.04205]

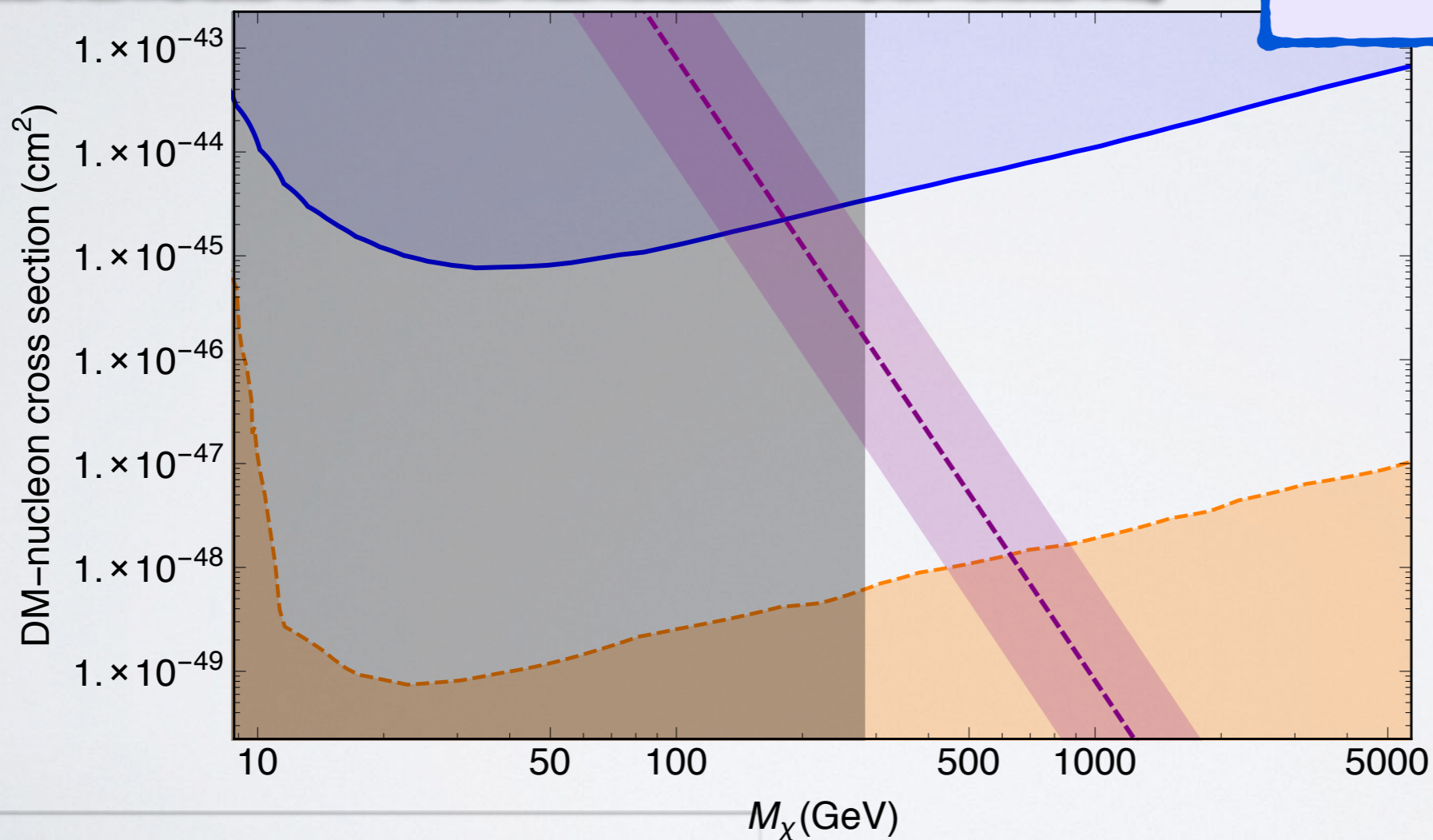
$$\sigma_{\text{nucleon}}(Z, A) = \frac{Z^4}{A^2} \frac{144\pi\alpha^4 \mu_{n\chi}^2 (M_F^A)^2}{m_\chi^6 R^2} [c_F]^2$$

Lowest bound from EM polarizability



Electric polarizability from lattice simulations

LUX exclusion bound for spin-independent cross section

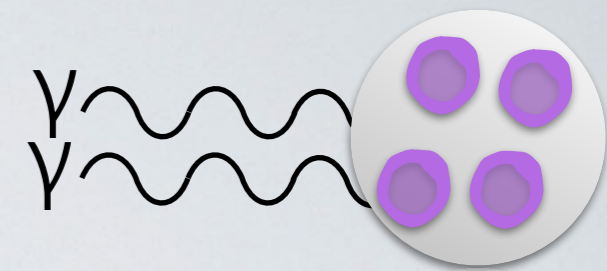


SU(4) $N_f=4$ Stealth DM

[LSD, 1503.04205]

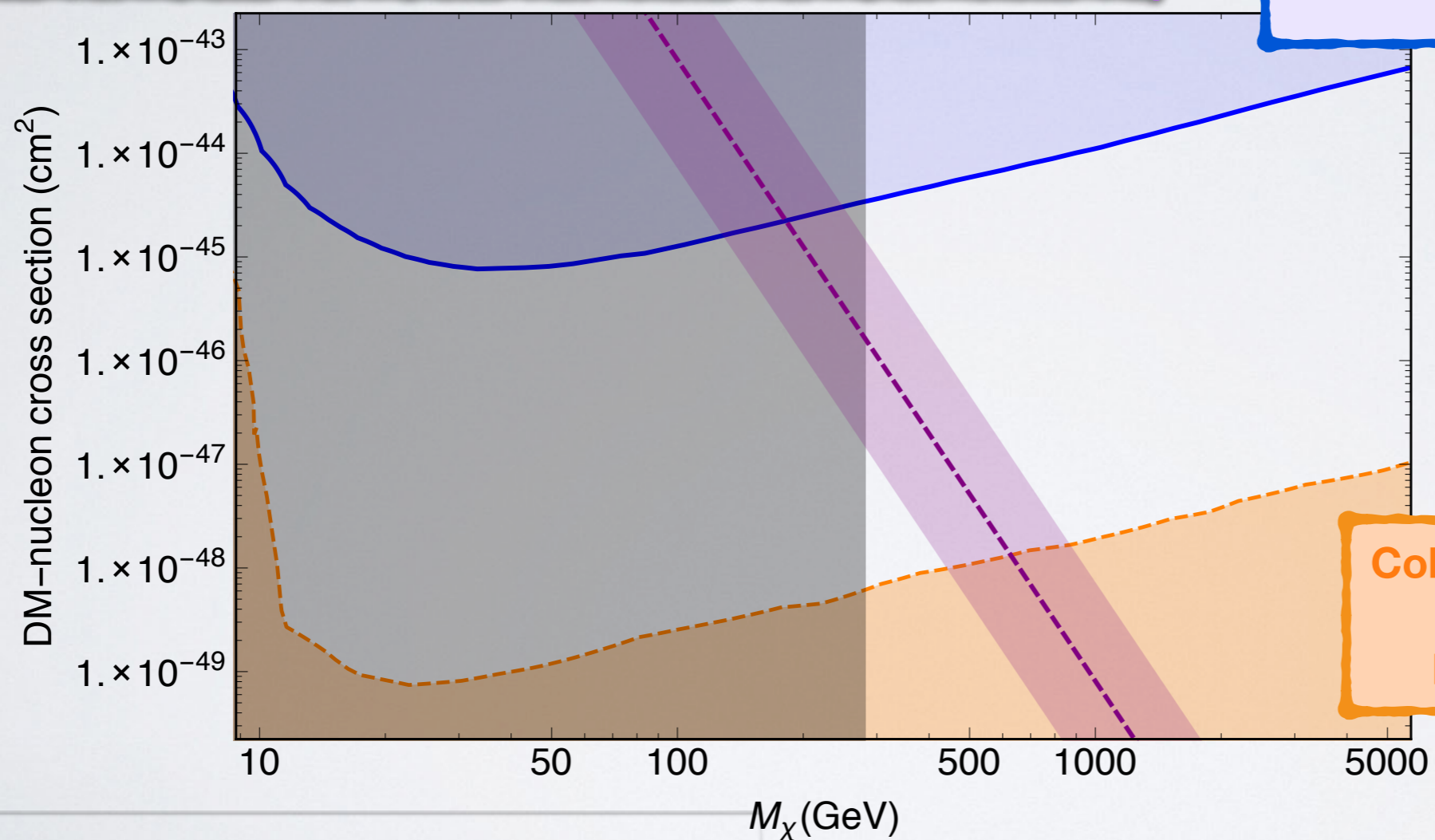
$$\sigma_{\text{nucleon}}(Z, A) = \frac{Z^4}{A^2} \frac{144\pi\alpha^4 \mu_{n\chi}^2 (M_F^A)^2}{m_\chi^6 R^2} [c_F]^2$$

Lowest bound from EM polarizability



Electric polarizability from lattice simulations

LUX exclusion bound for spin-independent cross section

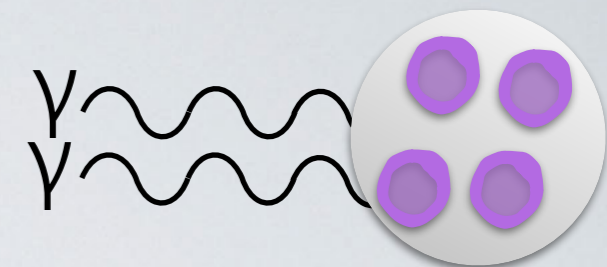


SU(4) $N_f=4$ Stealth DM
[LSD, 1503.04205]

Coherent neutrino scattering background

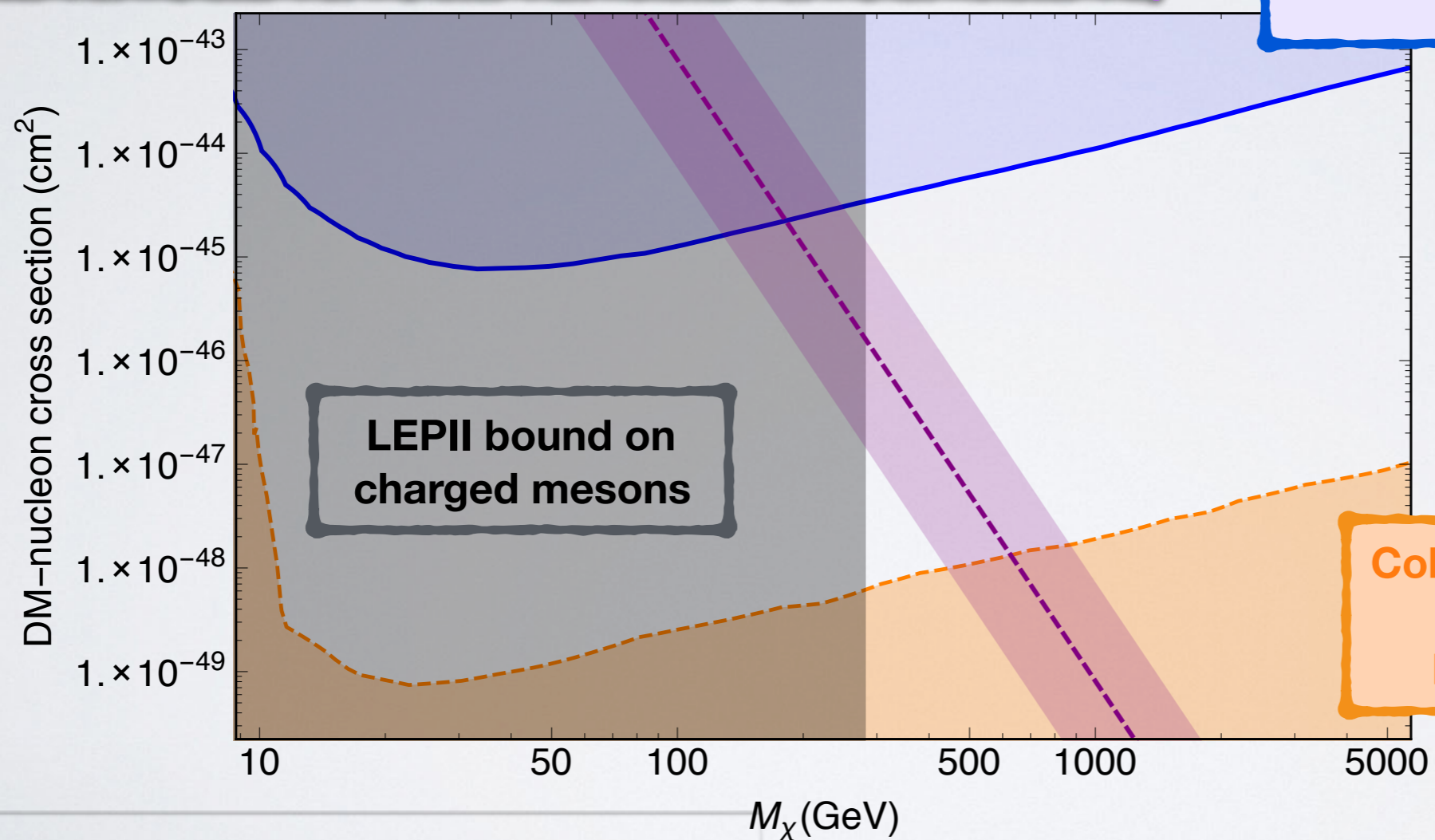
$$\sigma_{\text{nucleon}}(Z, A) = \frac{Z^4}{A^2} \frac{144\pi\alpha^4 \mu_{n\chi}^2 (M_F^A)^2}{m_\chi^6 R^2} [c_F]^2$$

Lowest bound from EM polarizability



Electric polarizability from lattice simulations

LUX exclusion bound for spin-independent cross section

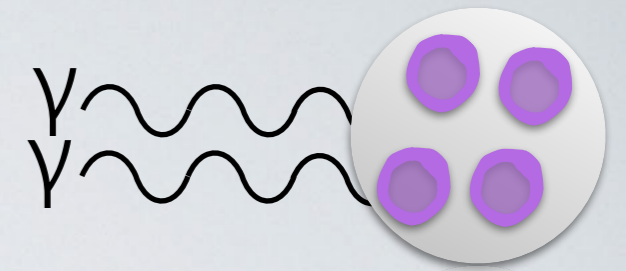


SU(4) $N_f=4$ Stealth DM
[LSD, 1503.04205]

Coherent neutrino scattering background

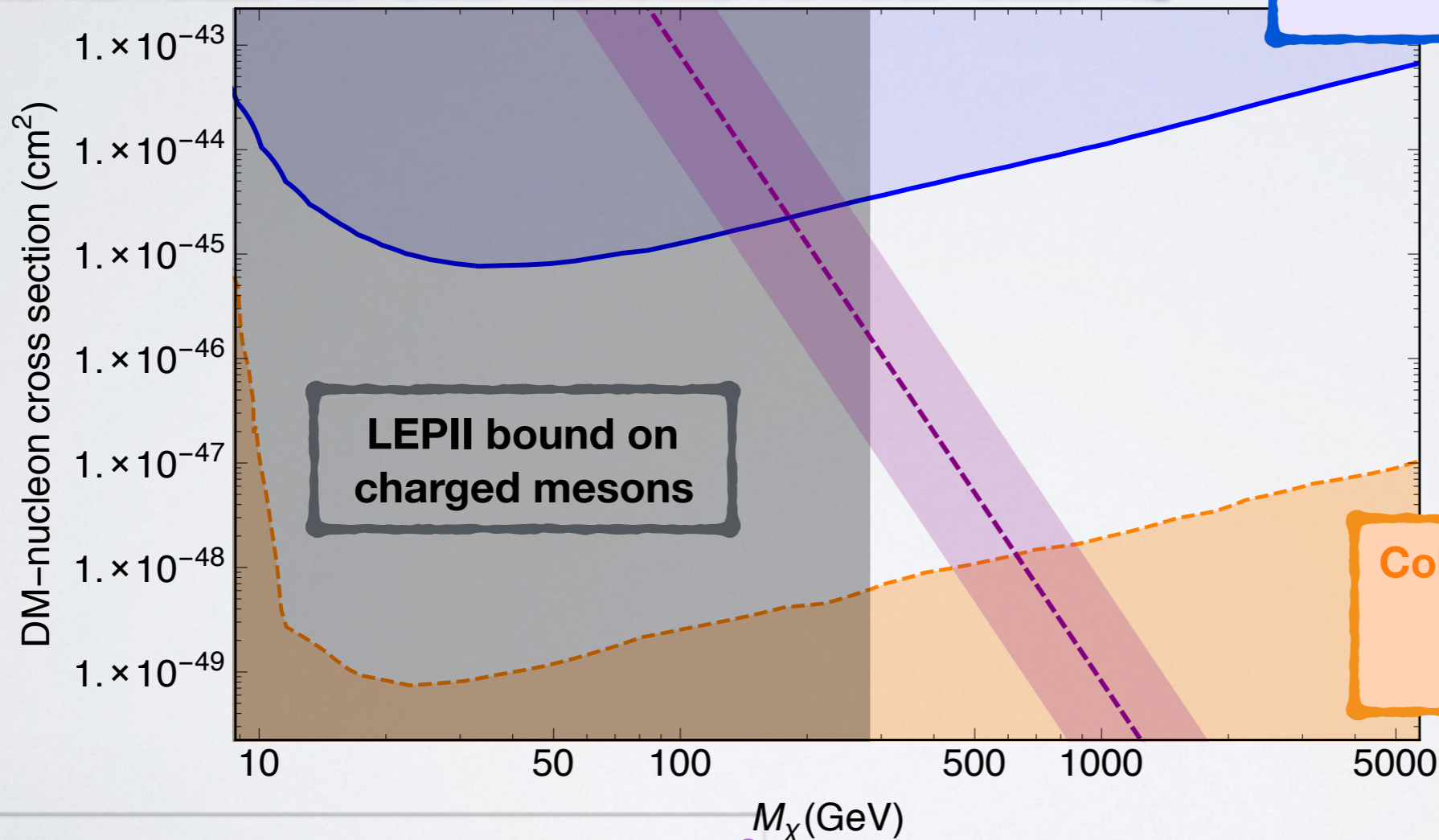
$$\sigma_{\text{nucleon}}(Z, A) = \frac{Z^4}{A^2} \frac{144\pi\alpha^4 \mu_{n\chi}^2 (M_F^A)^2}{m_\chi^6 R^2} [c_F]^2$$

Lowest bound from EM polarizability



Electric polarizability from lattice simulations

LUX exclusion bound for spin-independent cross section



SU(4) N_f=4 Stealth DM

[LSD, 1503.04205]

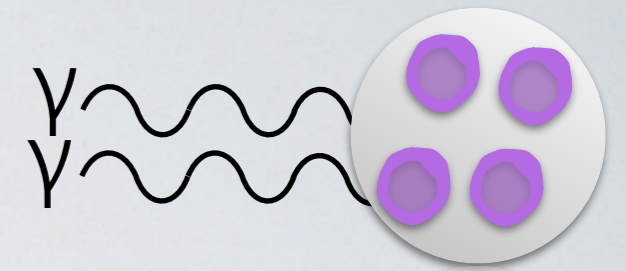
Coherent neutrino scattering background

LEP II bound on charged mesons

$$\sigma_{\text{nucleon}}(Z, A) = \frac{Z^4}{A^2} \frac{144\pi\alpha^4 \mu_{n\chi}^2 (M_F^A)^2}{m_\chi^6 R^2} [c_F]^2$$

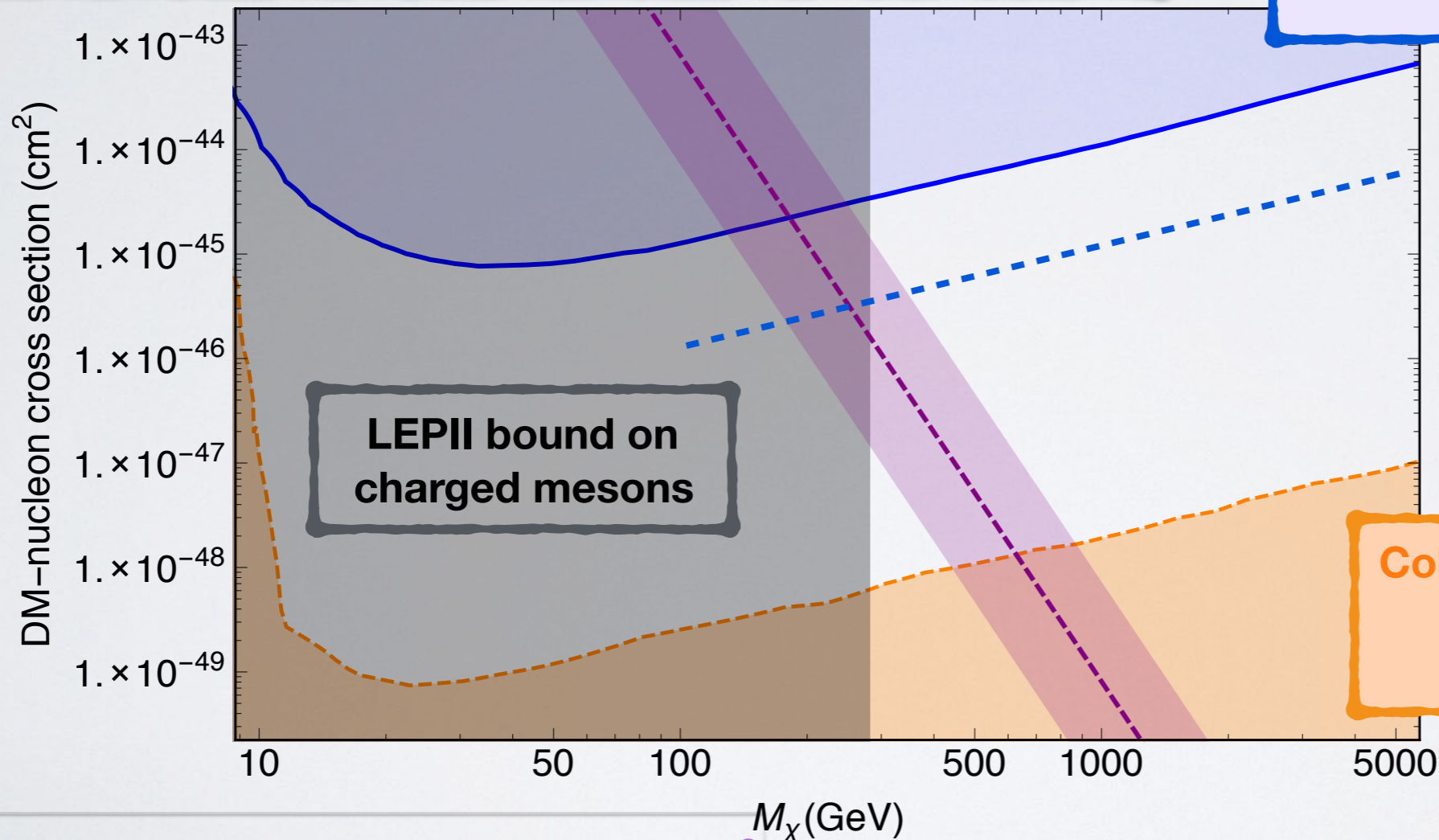
lowest allowed direct detection cross-section for composite dark matter theories with EW charged constituents

Lowest bound from EM polarizability



Electric polarizability from lattice simulations

LUX exclusion bound for spin-independent cross section



[XENON1T 34.2days 1705.06655]

SU(4) N_f=4 Stealth DM

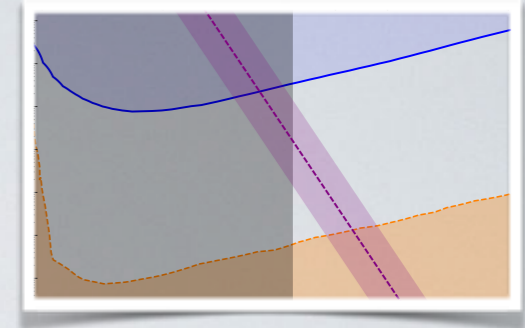
[LSD, 1503.04205]

Coherent neutrino scattering background

$$\sigma_{\text{nucleon}}(Z, A) = \frac{Z^4}{A^2} \frac{144\pi\alpha^4 \mu_{n\chi}^2 (M_F^A)^2}{m_\chi^6 R^2} [C_F]^2$$

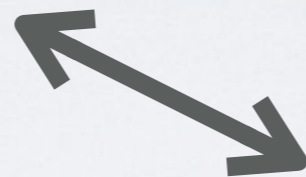
lowest allowed direct detection cross-section for composite dark matter theories with EW charged constituents

Concluding remarks



- ◆ **Composite** dark matter is a viable interesting possibility with rich phenomenology: all scales are natural
- ◆ **Lattice** methods can help in calculating direct detection cross sections, *production rates at colliders*, *self-interaction cross sections* and the *spectrum of gravitational waves*. **Direct phenomenological relevance.**
- ◆ Dark matter constituents can carry electroweak charges and still the stable composites are currently undetectable. **Stealth cross section.**
- ◆ **Lowest bound for composite dark matter models: ~ 300 GeV** (colliders+direct detection+lattice) (can be improved!)

DEEP LEARNING AND PHYSICS



ECT* workshop (3/2020):
Machine Learning for High
Energy Physics, on and off
the Lattice

Deep Learning And Physics DLAP2019



<http://kabuto.phys.sci.osaka-u.ac.jp/~koji/workshop/DLAP2019/>

> Yukawa Institute for Theoretical Physics
> Kyoto, Japan
> 31 Oct - 2 Nov 2019 ■

- > Invited speakers include:
- > Shun-Ichi Amari (RIKEN)
 - > *James Halverson (Northeastern U)
 - > Hong-Ye Hu (UC San Diego)
 - > Gurtej Kanwar (MIT)
 - > Sven Krippendorf (Sommerfeld Center)
 - > Junwei Liu (HKUST)
 - > Zi Yang Meng (CAS)
 - > Youichiro Miyake (Square Enix)
 - > Masayuki Ohzeki (Tohoku U)
 - > Fabian Ruhle (CERN)
 - > Rak-Kyeong Seong (Samsung SDS)
 - > Gary Shiu (U Wisconsin)
 - > Lei Wang (CAS)
 - > Youhei Yamaji (Tokyo U)
 - > Greg Yang (Microsoft Research)
 - > Hajime Yoshino (Osaka U)
 - > *special lecture at Osaka U on 30 Oct ■

> Organizers:

- > Koji Hashimoto (Osaka U)
- > Masatoshi Imada (Toyota RIKEN / Waseda U)
- > Kouji Kashiwa (Fukuoka Institute of Technology)
- > Yuki Nagai (JAEA CCSE)
- > Masayuki Ohzeki (Tohoku U)
- > Enrico Rinaldi (RIKEN & Arithmer Inc.)
- > Akinori Tanaka (RIKEN AIP)
- > Akio Tomiya (RIKEN BNL) ■

This conference is supported by: Yukawa Institute, JSPS/MEXT KAKENHI No.2902, 17H06462, 18K13548, Post-K CDSMI

本研究会は以下の研究費にサポートを受けて開催されます: 新学術領域科研究費「次世代物質探索のための離散幾何学」、
科研費17H06462、科研費18K13548、ポスト「京」プロジェクト「世代の産業を支える新機能デバイス・高性能材料の創成」

COME UN CACCIA INVISIBILE AI RADAR

G+ 1 | Tweet 7 | Share 215

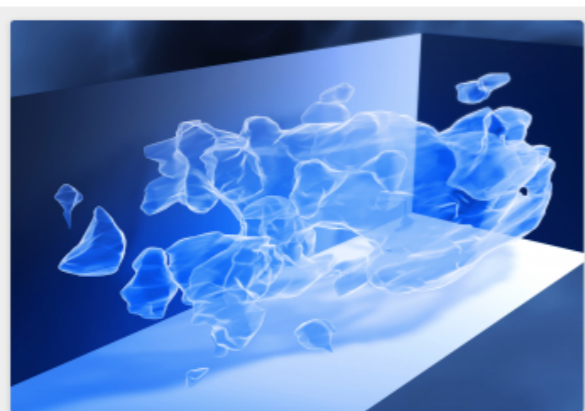
Materia oscura "stealth"

Quark oscuri tenuti insieme da un'interazione forte a sua volta oscura. Ecco come la dark matter riuscirebbe a eludere a ogni tentativo d'incastarla. Enrico Rinaldi (LLNL): «Esiste la possibilità che questo "mondo oscuro", con le sue nuove particelle, possa essere rivelato dagli esperimenti in corso al Large Hadron Collider al CERN di Ginevra»

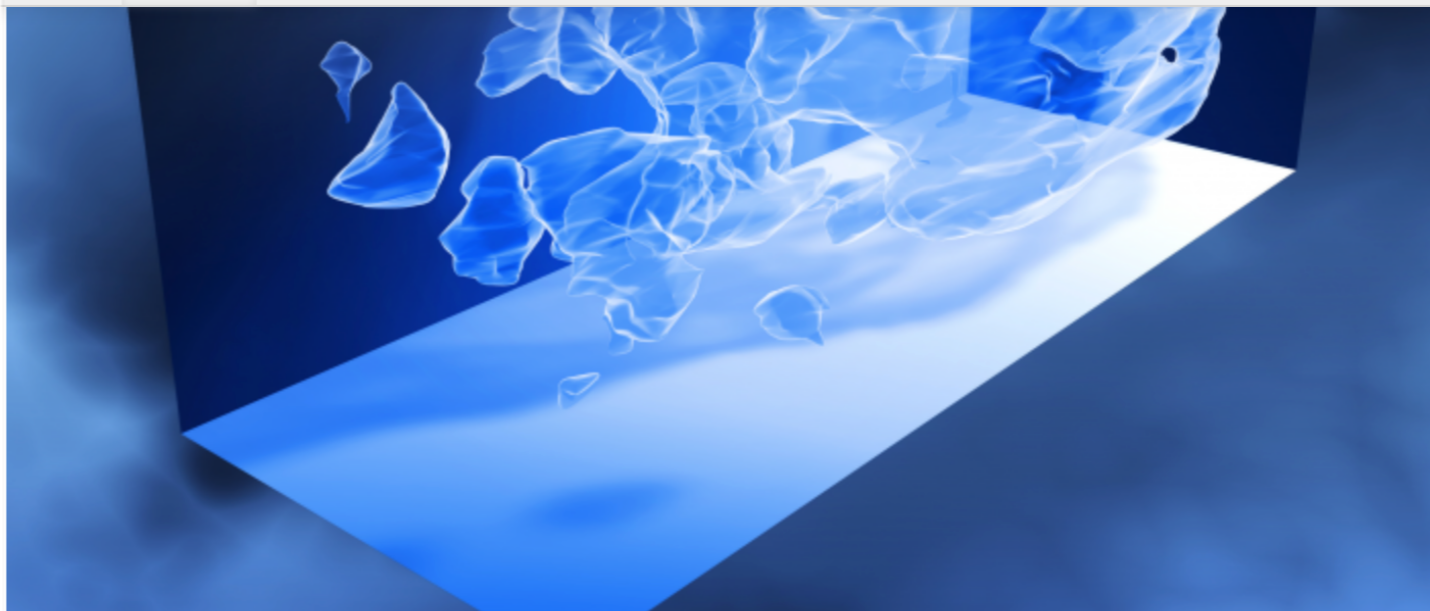
di Marco Malaspina | Segui @malamiao

venerdì 25 settembre 2015 @ 16:15

Stealth come furtiva. Stealth come imprevedibile. Stealth come quei minacciosi aerei da guerra dal profilo sagomato così da essere invisibili ai radar. Da quanto emerge dai calcoli dei fisici dell'LLNL, il Lawrence Livermore National Laboratory californiano, e dai modelli dati in pasto a Vulcan (un supercomputer per il calcolo parallelo in grado di masticare numeri al ritmo dei petaflop), sarebbe questa la natura della materia oscura: *stealthy*, appunto. Per forza non c'è ancora esperimento che sia riuscito a incastrarla.



Mappa 3D della distribuzione su larga scala della materia oscura ricostruita da misure di lente gravitazionale debole utilizzando il telescopio spaziale Hubble



This 3D map illustrates the large-scale distribution of dark matter, reconstructed from measurements of weak gravitational lensing by using the Hubble Space Telescope. (Download Image)

New 'stealth dark matter' theory may explain mystery of the universe's missing mass



Lawrence Livermore National Laboratory (LLNL) scientists have come up with a new theory that may identify why dark matter has evaded direct detection in Earth-based experiments.

Anne M Stark
stark8@llnl.gov
925-422-9799

Detecting Stealth Dark Matter Directly through Electromagnetic Polarizability.

Overview of attention for article published in Physical Review Letters, October 2015



About this score

In the top 5% of all research outputs scored by Altmetric

MORE...

SUMMARY

News | Blogs | Twitter | Facebook | Google+

Title Detecting Stealth Dark Matter Directly through Electromagnetic Polarizability.
Published in Physical Review Letters, October 2015
DOI 10.1103/physrevlett.115.171803
Pubmed ID 26551103
Authors T. Appelquist, E. Berkowitz, R. C. Brower, M. I. Buchoff, G. T. Fleming, X.-Y. Jin, J. Kiskis, G. D...
Abstract We calculate the spin-independent scattering cross section for direct detection that results from...
Abstract We calculate the spin-independent scattering cross section for direct detection that results from...

TWITTER DEMOGRAPHICS

MENDELEY READERS

SCORE IN C...

The data shown below were collected from the profiles of 23 tweeters who shared this research output. Click here to find out how the information was compiled.

Le Scienze

EDIZIONE ITALIANA DI SCIENTIFIC AMERICAN

LA RIVISTA IN EDIZIONE

Materia oscura
Nuove ipotesi su una misteriosa componente ancora più bizzarra
In edicola dal 1 settembre

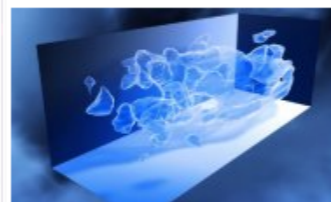
ABBONAMENTI E INFO

ZOOM SU | comportamento | cosmologia | neuroscienze | alimentazione

iflammenco! | Festival Flamenco 5-11 ottobre

28 settembre 2015

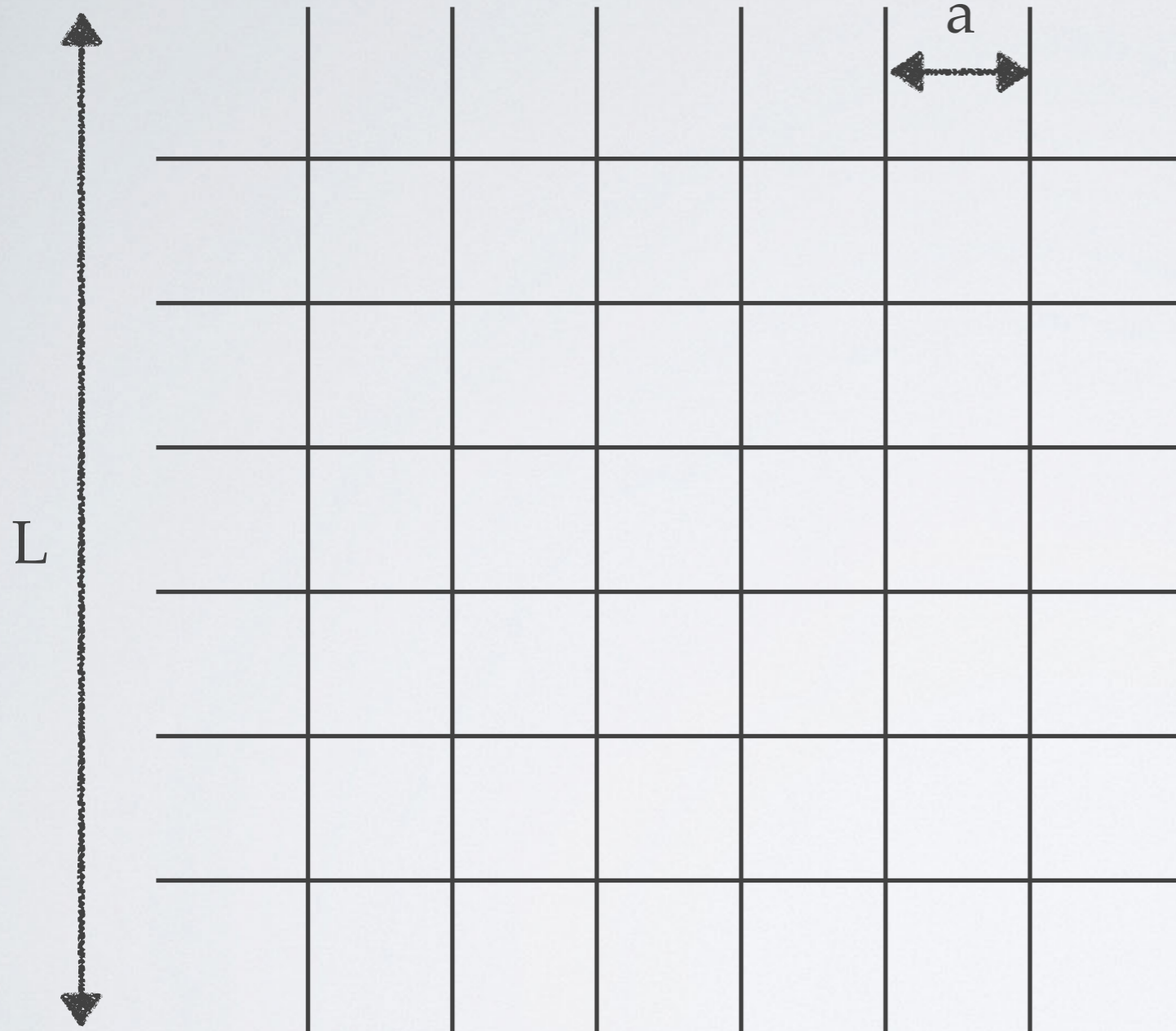
Un nuovo modello per la materia oscura



Cortesia Lawrence Livermore National Laboratory

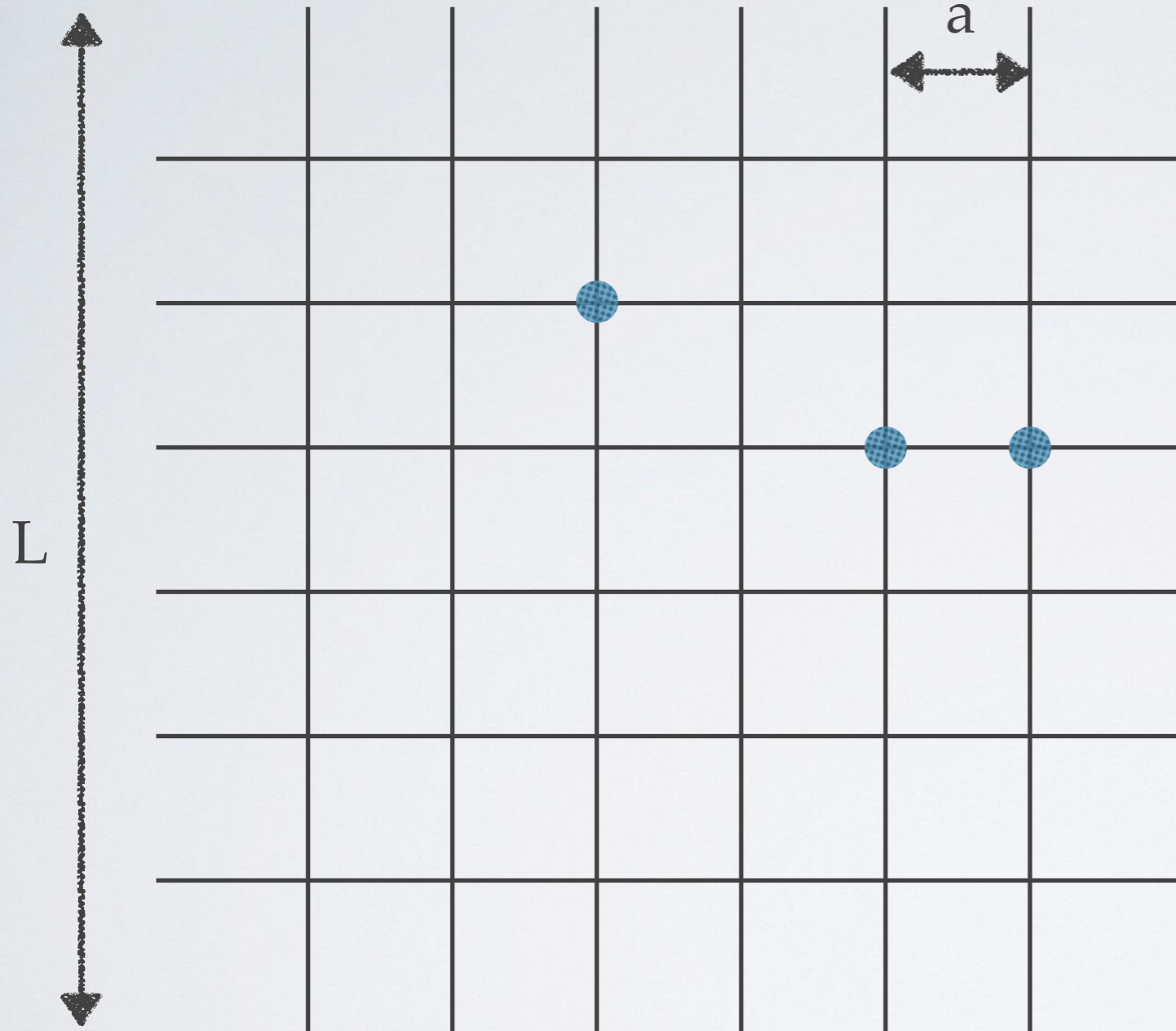
Questa forma misteriosa di materia potrebbe avere una struttura composita come la materia ordinaria, con "quark oscuri" aggregati e tenuti insieme da un analogo della forza che permette ai normali nuclei di rimanere stabili. I componenti di questo tipo di materia oscura, definita stealth matter, potrebbero essere studiati in modo indiretto dal collisore Large Hadron Collider del CERN di Ginevra

Lattice primer



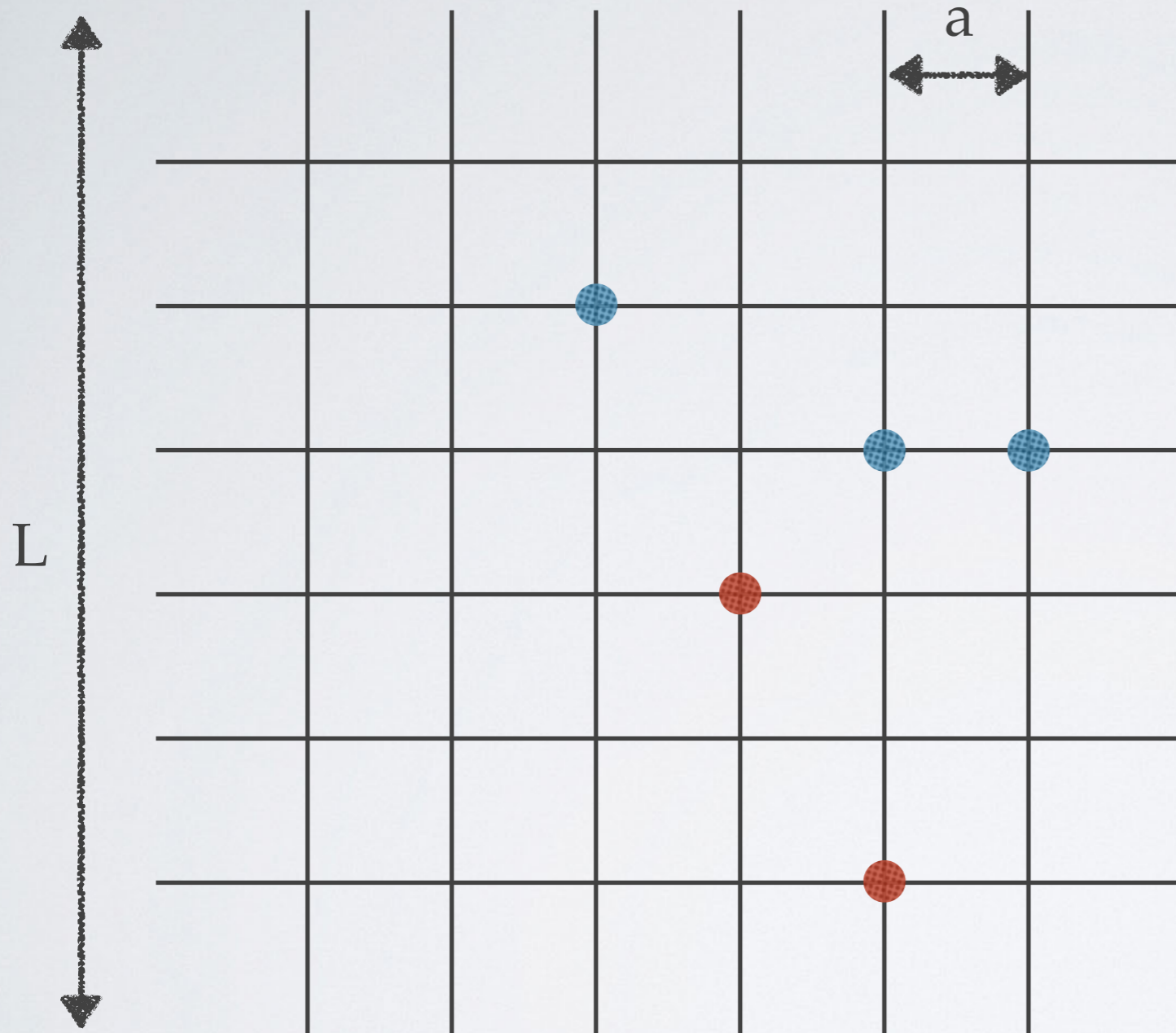
- Discretize space and time
 - lattice spacing “a”
 - lattice size “L”
- Keep all d.o.f. of the theory
 - not a model!
 - no simplifications
- Amenable to numerical methods
 - Monte Carlo sampling
 - use supercomputers
- Precisely quantifiable and improvable errors
 - Systematic
 - Statistical

Lattice primer



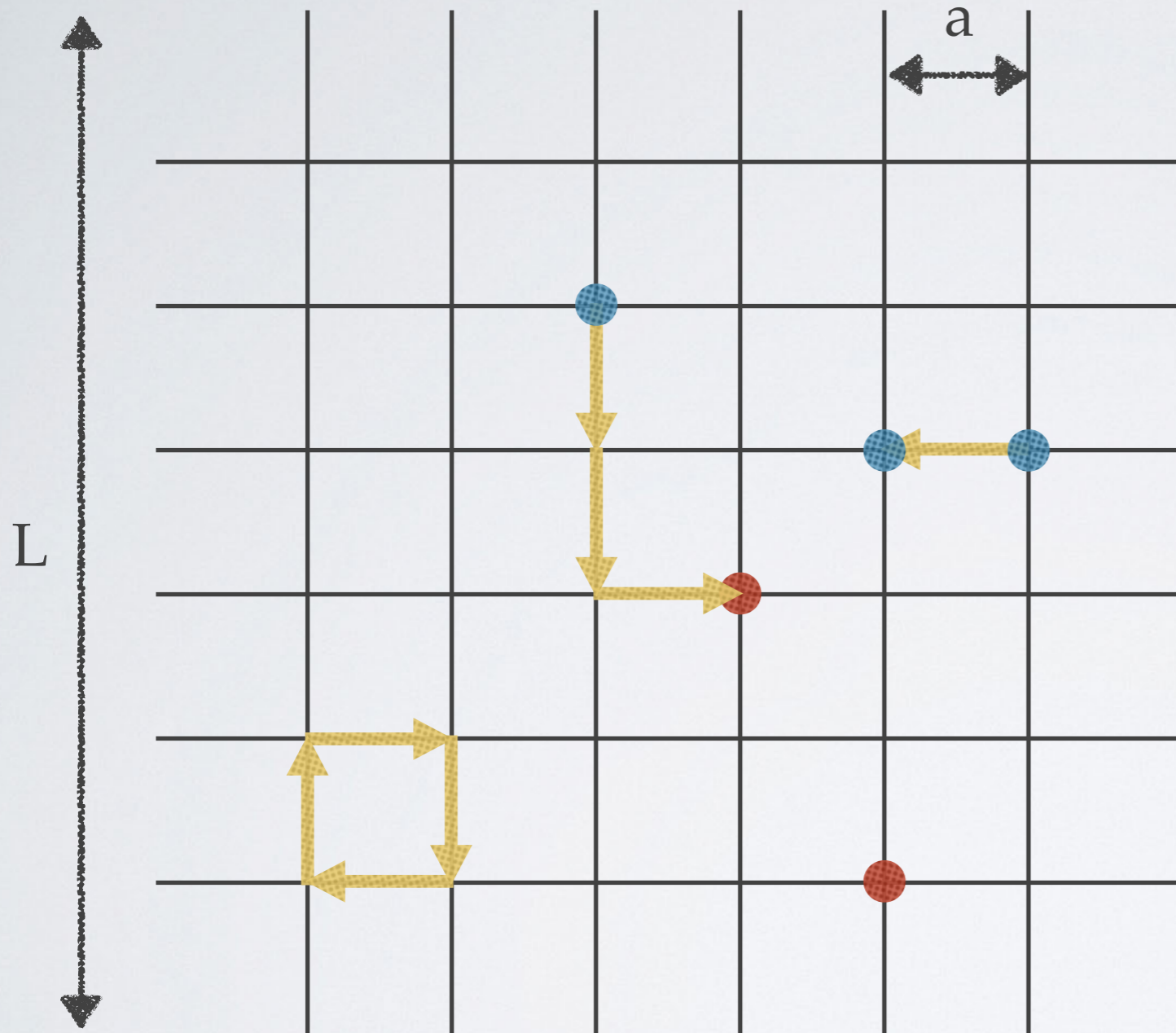
- Discretize space and time
 - lattice spacing “a”
 - lattice size “L”
- Keep all d.o.f. of the theory
 - not a model!
 - no simplifications
- Amenable to numerical methods
 - Monte Carlo sampling
 - use supercomputers
- Precisely quantifiable and improvable errors
 - Systematic
 - Statistical

Lattice primer



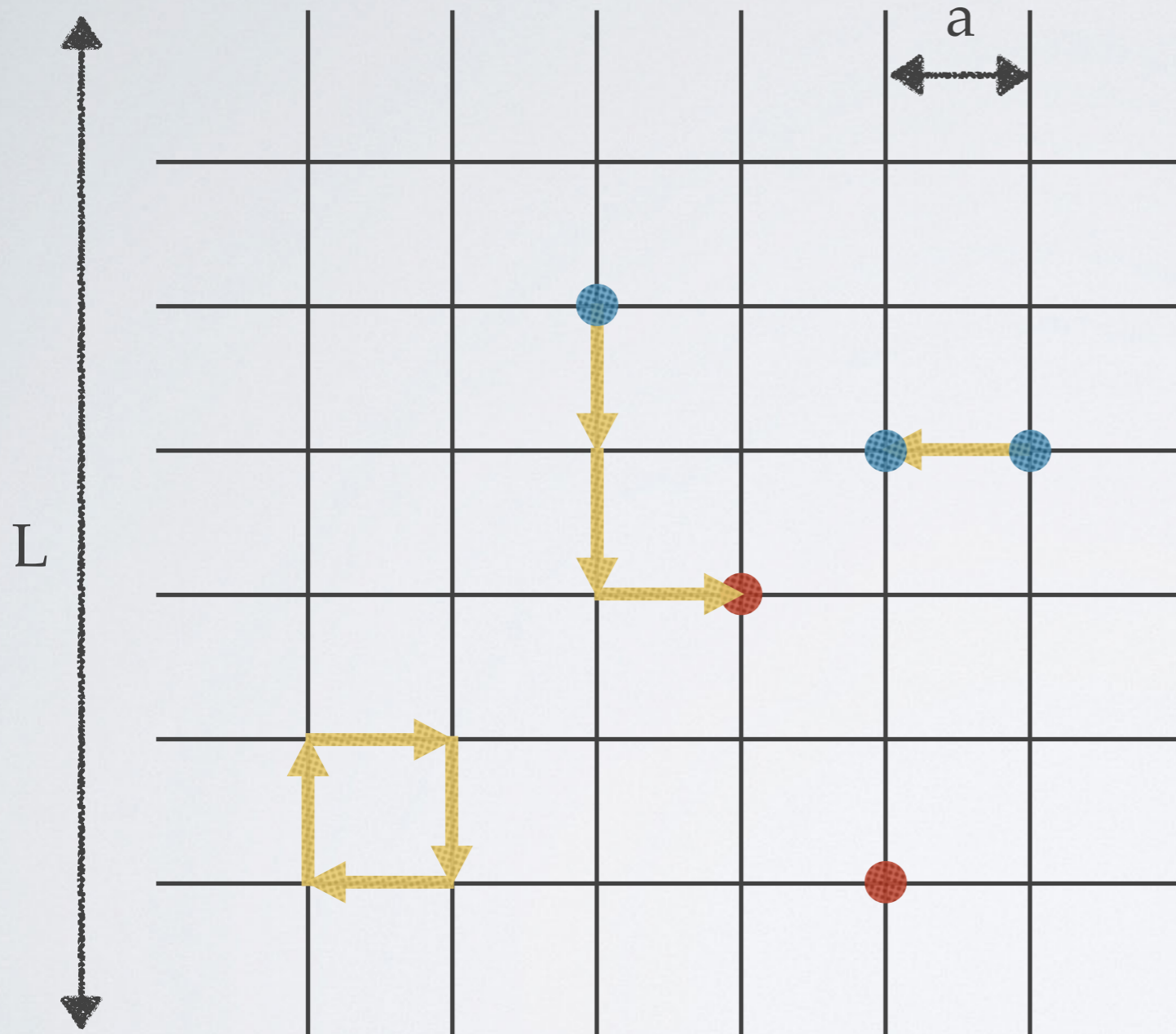
- Discretize space and time
 - lattice spacing “ a ”
 - lattice size “ L ”
- Keep all d.o.f. of the theory
 - not a model!
 - no simplifications
- Amenable to numerical methods
 - Monte Carlo sampling
 - use supercomputers
- Precisely quantifiable and improvable errors
 - Systematic
 - Statistical

Lattice primer



- Discretize space and time
 - lattice spacing “ a ”
 - lattice size “ L ”
- Keep all d.o.f. of the theory
 - not a model!
 - no simplifications
- Amenable to numerical methods
 - Monte Carlo sampling
 - use supercomputers
- Precisely quantifiable and improvable errors
 - Systematic
 - Statistical

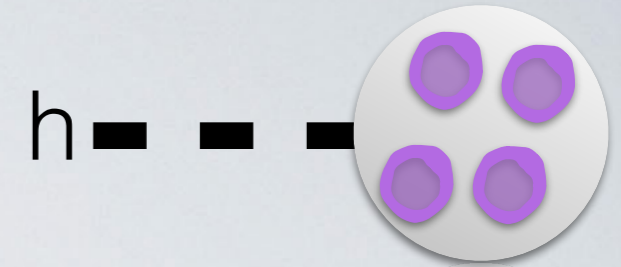
Lattice primer



- Discretize space and time
 - lattice spacing “a”
 - lattice size “L”
- Keep all d.o.f. of the theory
 - not a model!
 - no simplifications
- Amenable to numerical methods
 - Monte Carlo sampling
 - use supercomputers
- Precisely quantifiable and improvable errors
 - Systematic
 - Statistical

N_c **N_f** **N_r** parameters that can be easily changed

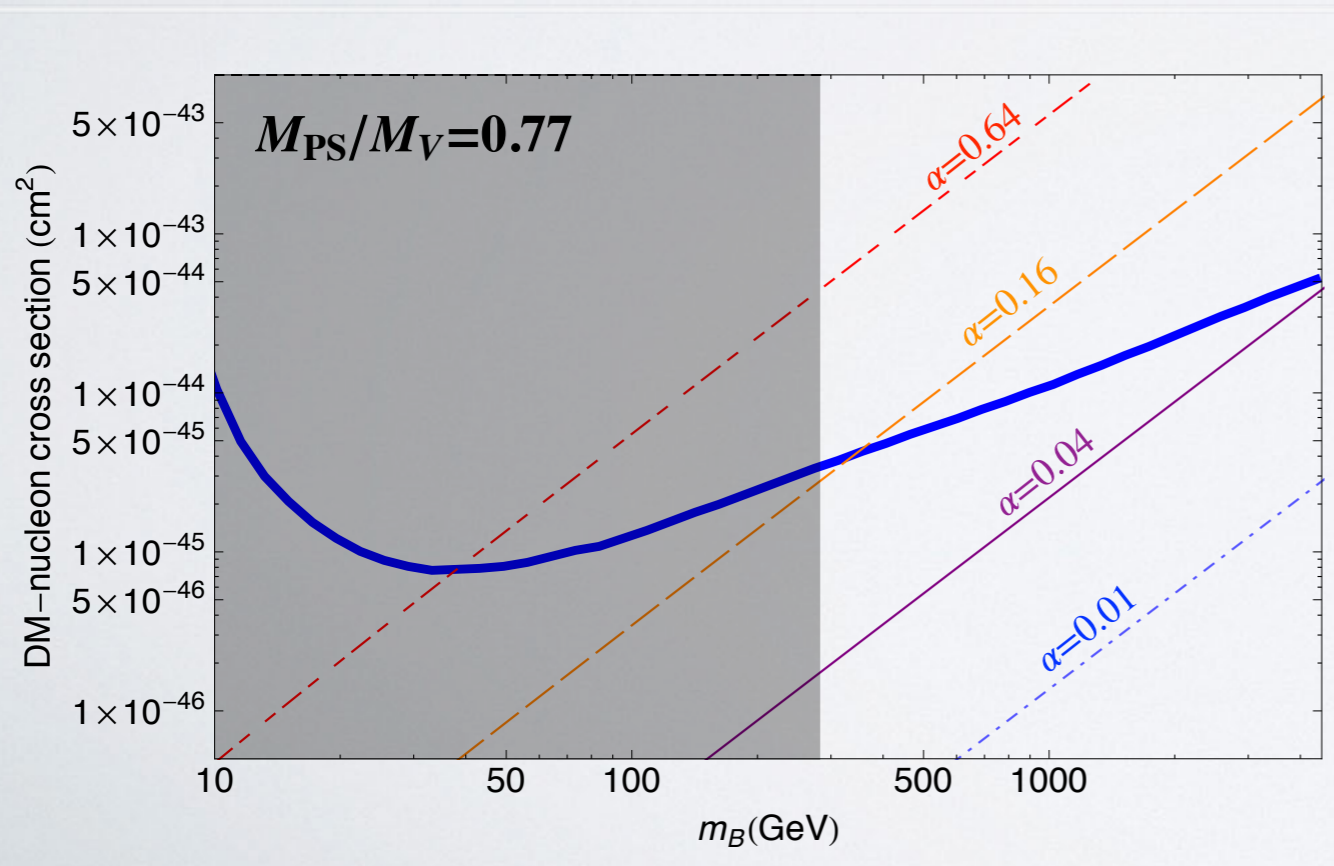
Bounds from Higgs exchange



- ◆ Lattice results for the cross-section are compared to **experimental** bounds
- ◆ Coupling space in specific models can be vastly constrained

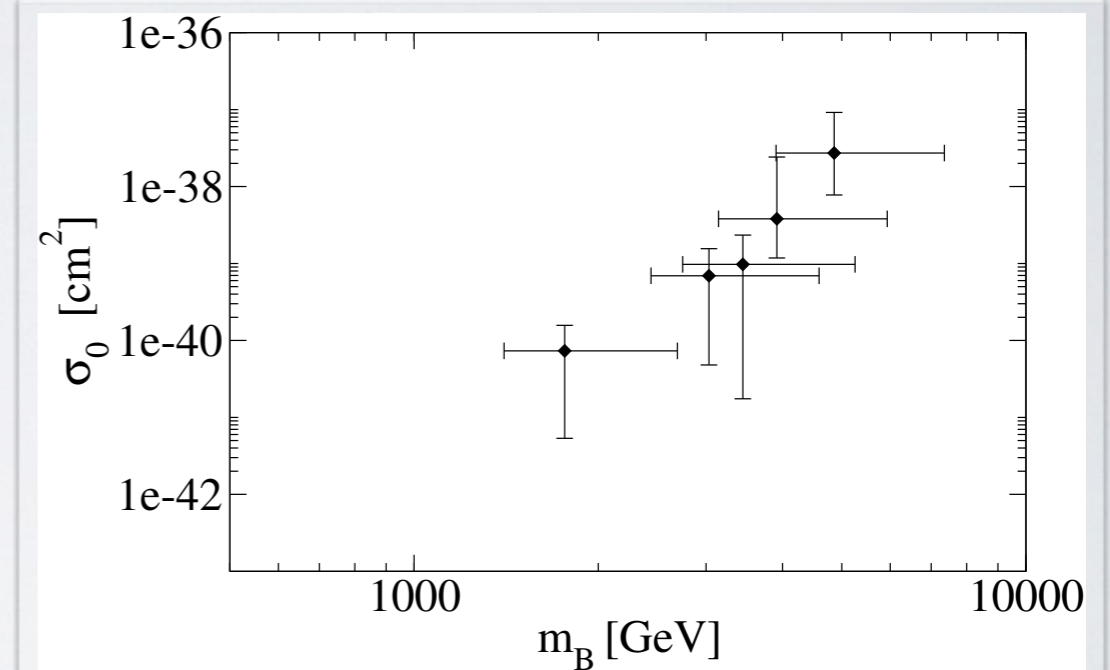
SU(4) $N_f=4$ Stealth DM

[LSD, 1402.6656-1503.04203]



SU(3) $N_f=8$ “technibaryon”

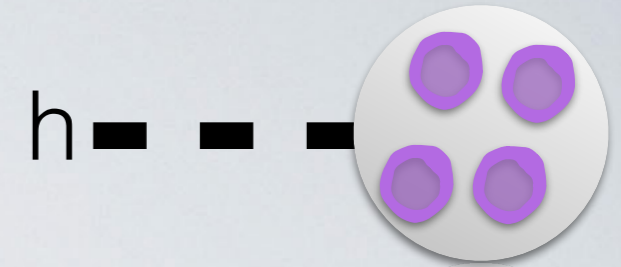
[LatKMI, 1510.07373]



- ◆ Some candidates can be excluded as dominant sources of dark matter
- ◆ There is lattice evidence for universality of dark scalar form factors

[DeGrand et al., 1501.05665]

Bounds from Higgs exchange

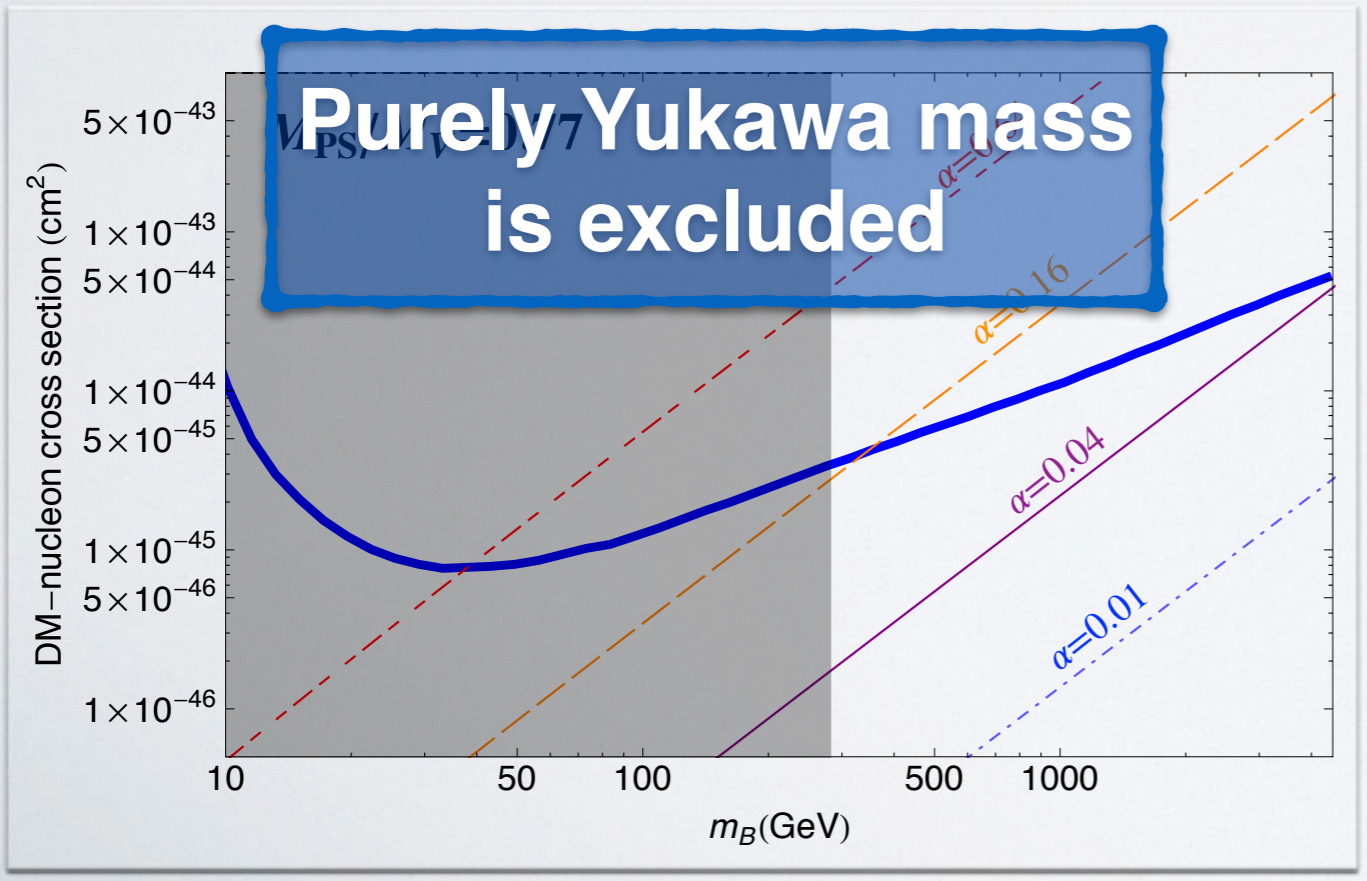


- ◆ Lattice results for the cross-section are compared to **experimental** bounds
- ◆ Coupling space in specific models can be vastly constrained

SU(4) $N_f=4$ Stealth DM

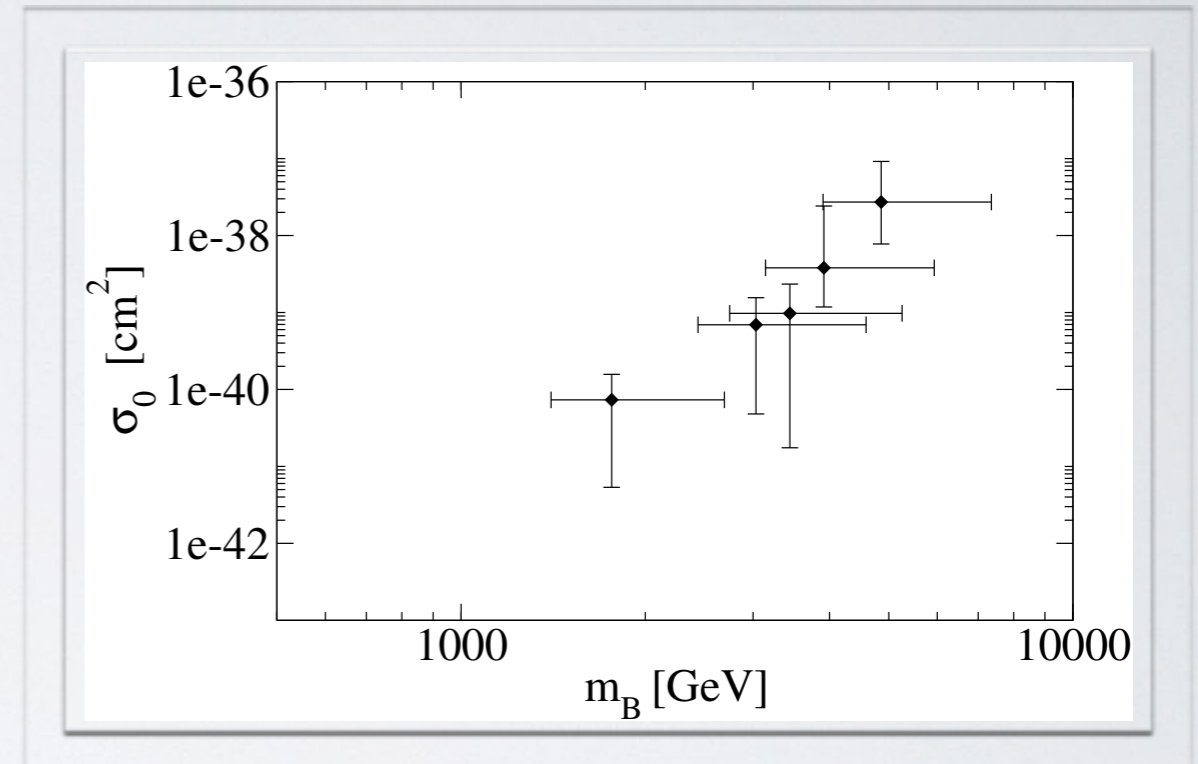
[LSD, 1402.6656-1503.04203]

Purely Yukawa mass is excluded



SU(3) $N_f=8$ "technibaryon"

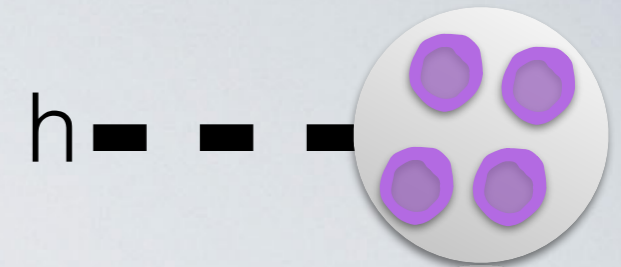
[LatKMI, 1510.07373]



- ◆ Some candidates can be excluded as dominant sources of dark matter
- ◆ There is lattice evidence for universality of dark scalar form factors

[DeGrand et al., 1501.05665]

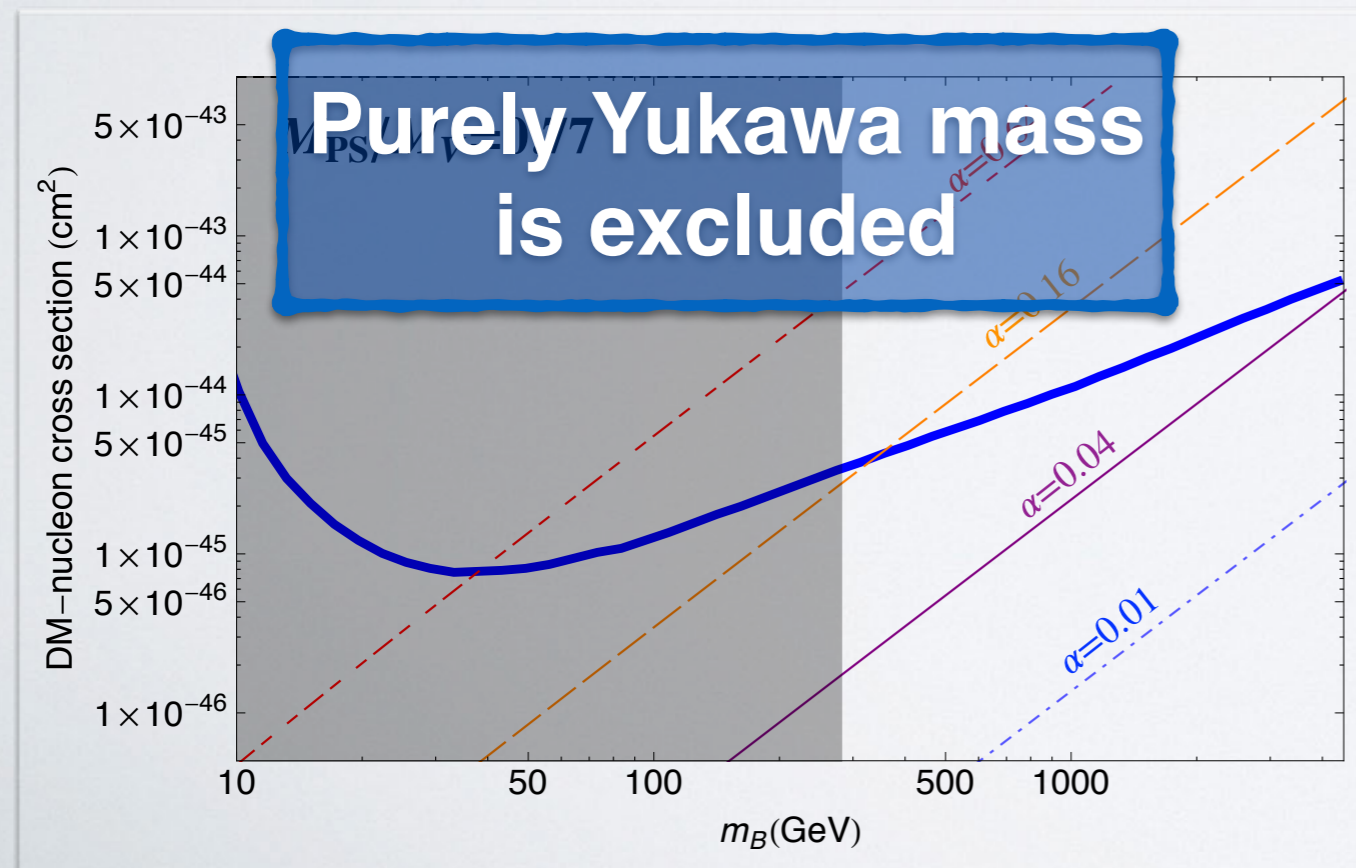
Bounds from Higgs exchange



- ◆ Lattice results for the cross-section are compared to **experimental** bounds
- ◆ Coupling space in specific models can be vastly constrained

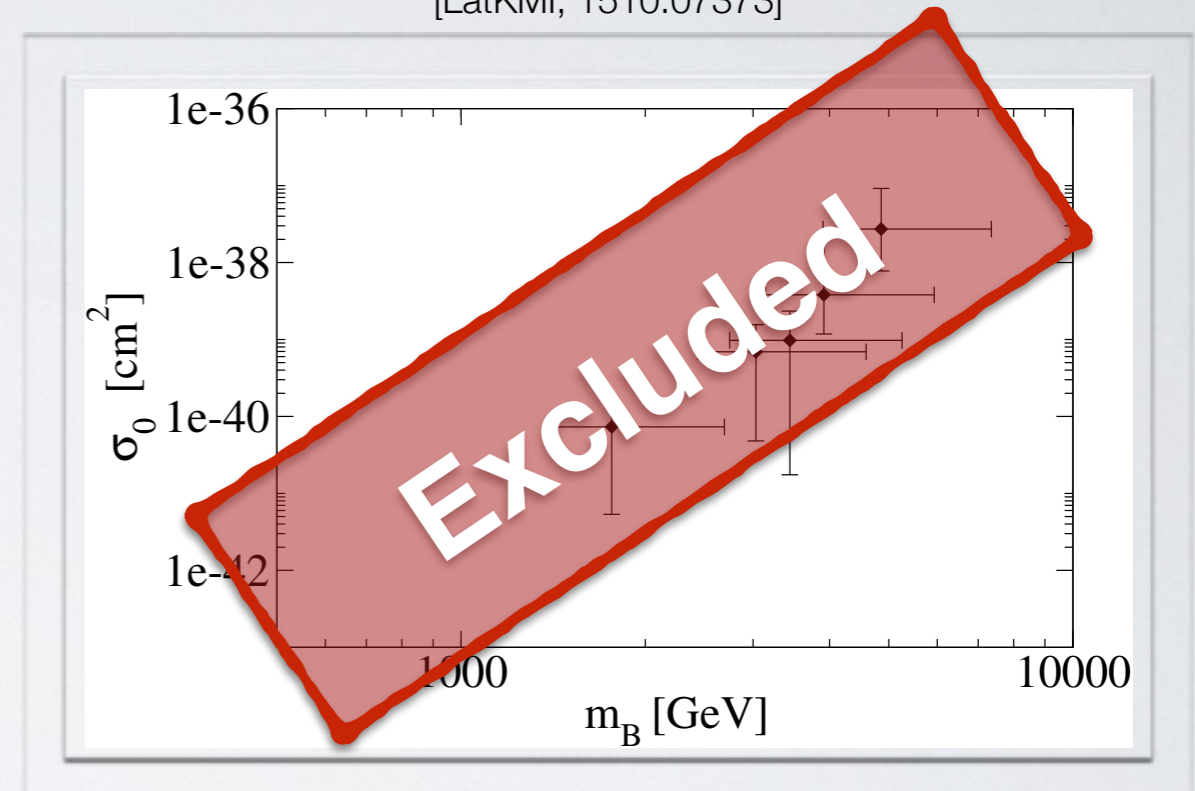
SU(4) $N_f=4$ Stealth DM

[LSD, 1402.6656-1503.04203]



SU(3) $N_f=8$ "technibaryon"

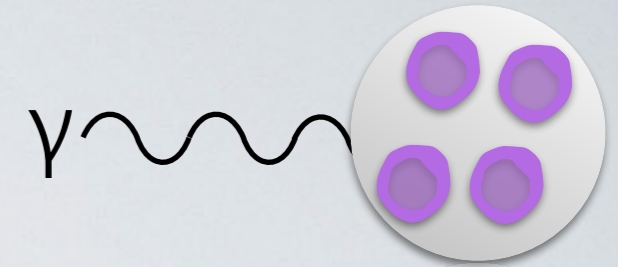
[LatKMI, 1510.07373]



- ◆ Some candidates can be excluded as dominant sources of dark matter
- ◆ There is lattice evidence for universality of dark scalar form factors

[DeGrand et al., 1501.05665]

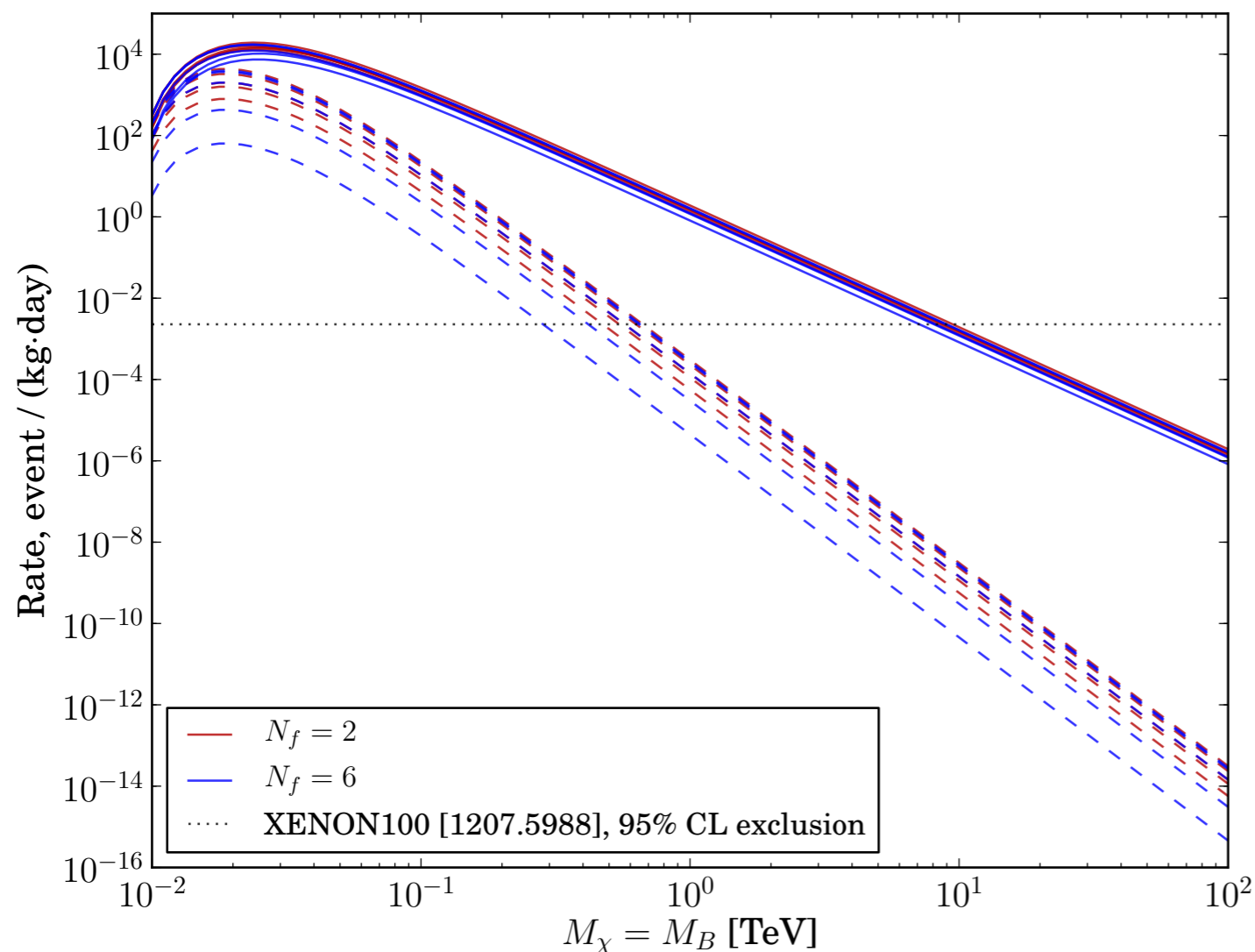
Bounds from EM moments



Mesonic and Baryonic EM form factors
directly from lattice simulations

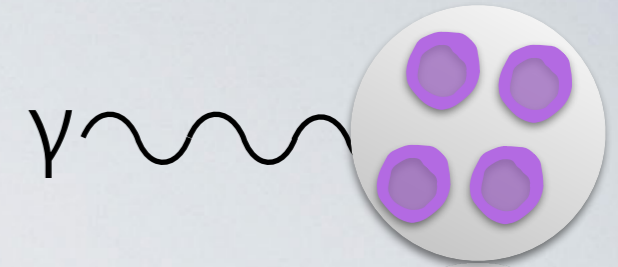
SU(3) $N_f=2,6$ dark fermionic baryon

[LSD, 1301.1693]



- ★ baryon similar to QCD neutron
- ★ dark quarks with $Q=Y$
- ★ calculate connected 3pt
- ★ scale set by DM mass
- ★ magnetic moment dominates
- ★ results independent of N_f

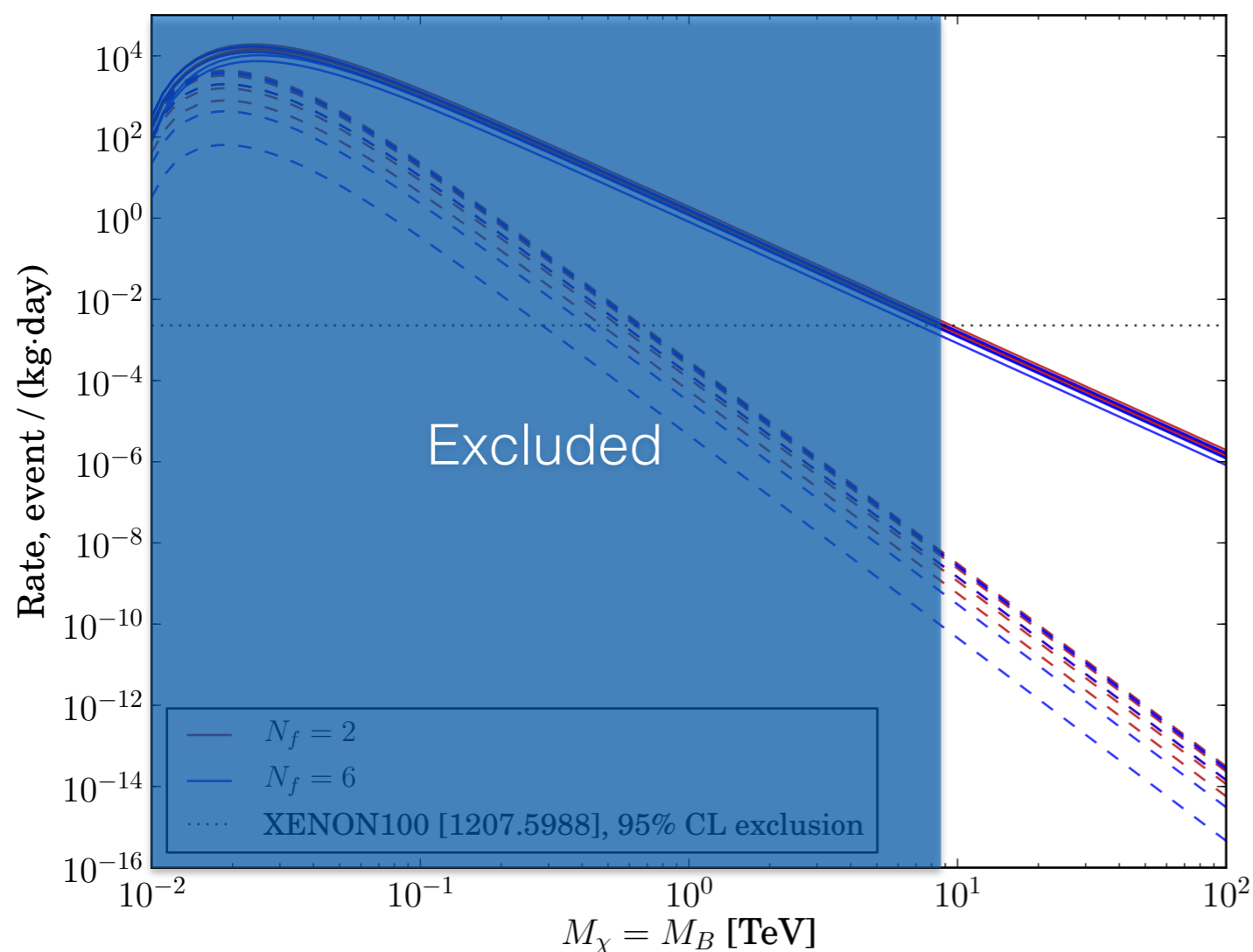
Bounds from EM moments



Mesonic and Baryonic EM form factors
directly from lattice simulations

SU(3) $N_f=2,6$ dark fermionic baryon

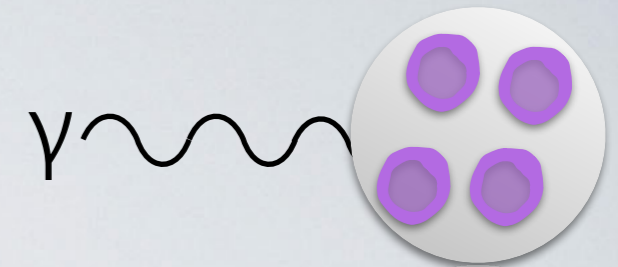
[LSD, 1301.1693]



- ★ baryon similar to QCD neutron
- ★ dark quarks with $Q=Y$
- ★ calculate connected 3pt
- ★ scale set by DM mass
- ★ magnetic moment dominates
- ★ results independent of N_f

$M_B > 10$ TeV

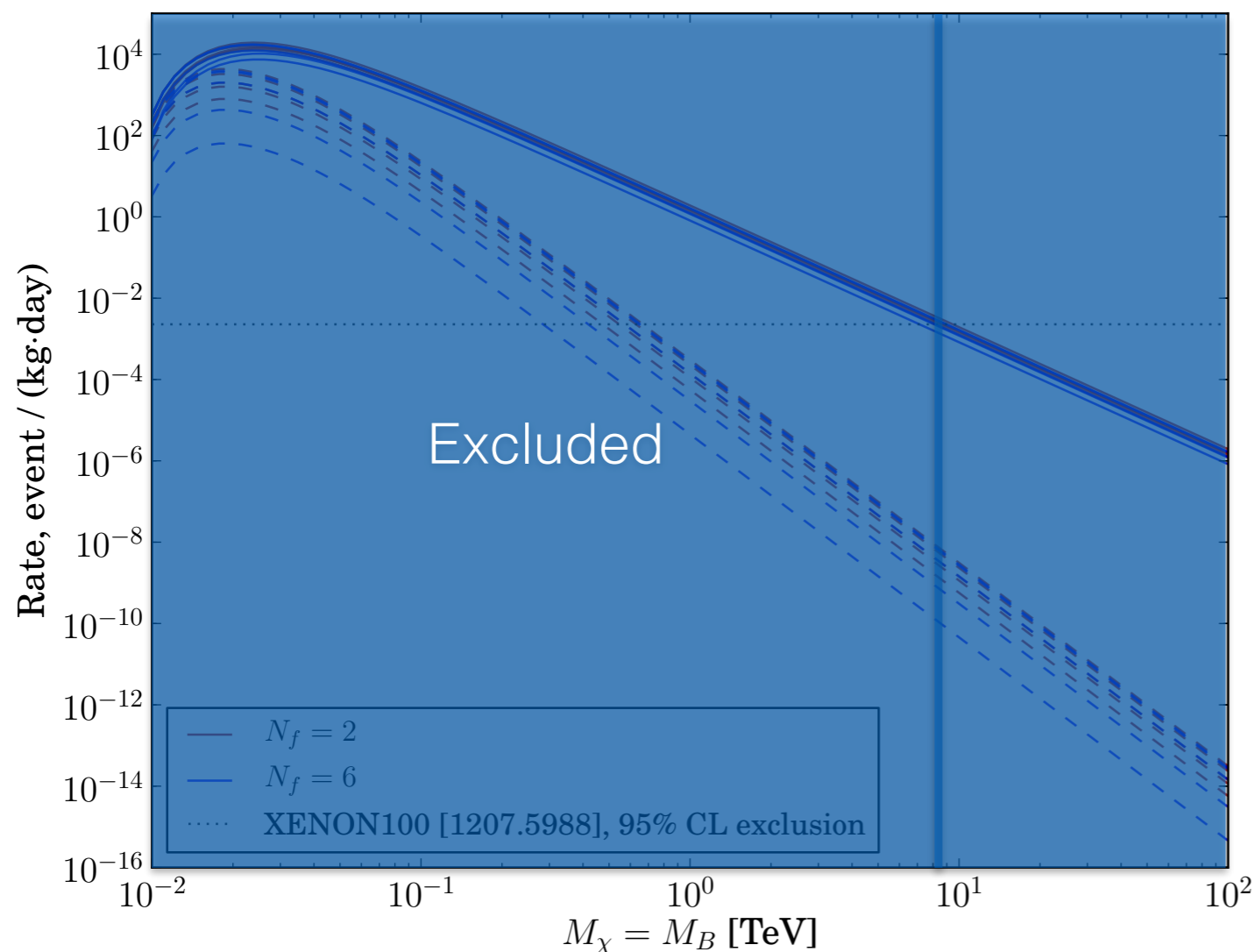
Bounds from EM moments



Mesonic and Baryonic EM form factors
directly from lattice simulations

SU(3) $N_f=2,6$ dark fermionic baryon

[LSD, 1301.1693]



- ★ baryon similar to QCD neutron
- ★ dark quarks with $Q=Y$
- ★ calculate connected 3pt
- ★ scale set by DM mass
- ★ magnetic moment dominates
- ★ results independent of N_f

$M_B > 10$ TeV

pushed to > 100 TeV
with new LUX