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German Research Foundation



Bielefeld Seoul  
International Research Training Group 2235

RMT IN SUB-ATOMIC PHYSICS AND BEYOND

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# **Universal Broadening of Zero Modes: A General Framework and Identification**

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BIELEFELD UNIVERSITY

JOINT WORK WITH  
MARIO KIEBURG AND KIM SPLITTORFF  
[\[arXiv:1902.01733\]](https://arxiv.org/abs/1902.01733)

- ▶ **The Model**
  - ▶ Motivation: Physical Systems
  - ▶ Intuition
  - ▶ Set-up
- ▶ Decoupling of Spectrum
  - ▶ Eigenvalue Equation
  - ▶ Conditions
- ▶ Central Limit Theorem for Matrices
- ▶ Scaling and Applications
- ▶ Conclusions

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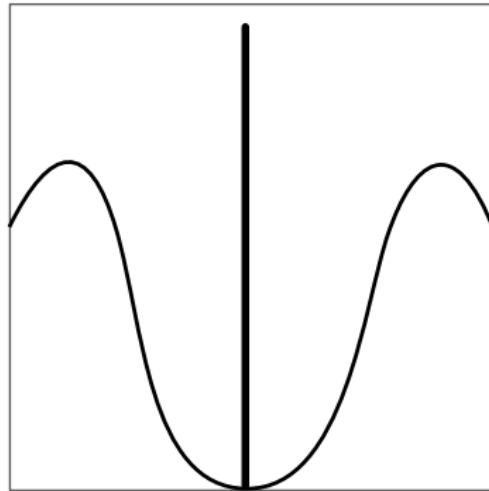
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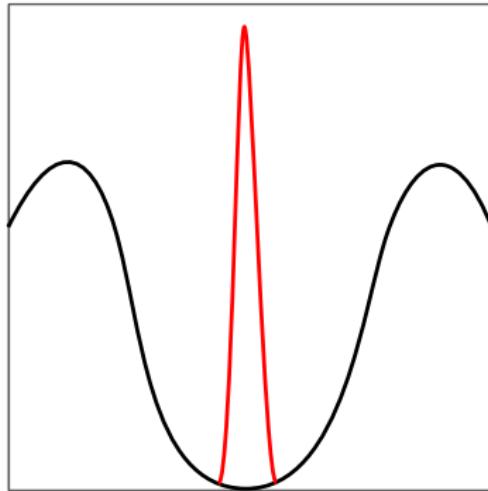
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$$K = \underbrace{A}_{\text{Pure System with zero mode}}$$

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$$K = \underbrace{A}_{\text{Pure System with zero mode}} + \underbrace{\alpha U S U^\dagger}_{\text{Small perturbation}}$$

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## Motivation: Physical Systems

- ▶ Topological modes, but more general framework
- ▶ The behaviour was seen before in several specific systems
  - ▶ Wilson-Dirac Operator in finite-volume lattice-QCD

$$D = \begin{pmatrix} 0 & W \\ W^\dagger & 0 \end{pmatrix} + a \begin{pmatrix} A & 0 \\ 0 & B \end{pmatrix}$$

[Akemann, Damgaard, Splittorff, Verbaarschot [arXiv:1012.0752]]  
[Kieburg, Verbaarschot, Zafeiropoulos [arXiv:1307.7251],[arXiv:1505.01784]]

- ▶ Coupled Chiral Systems  
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- ▶ Difficult to distinguish topological and non-topological modes  
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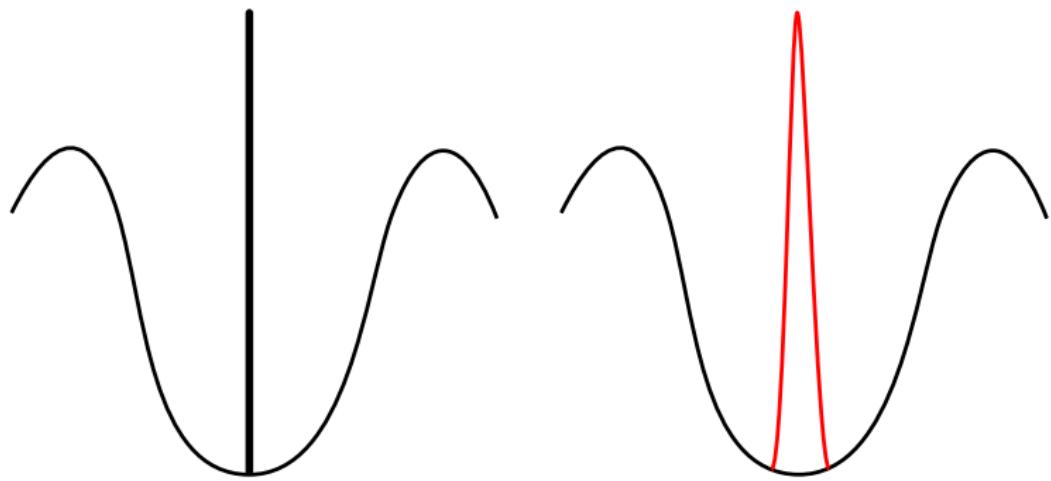
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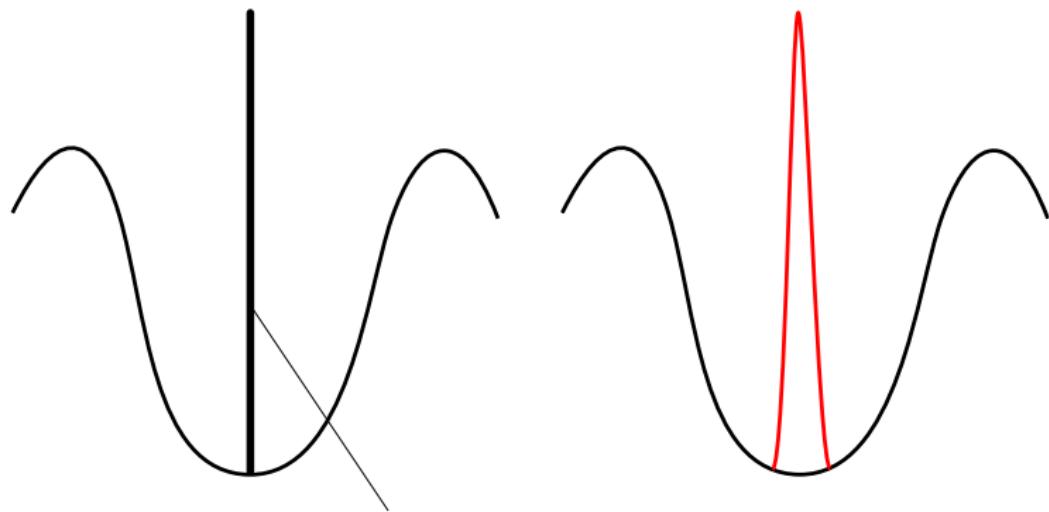
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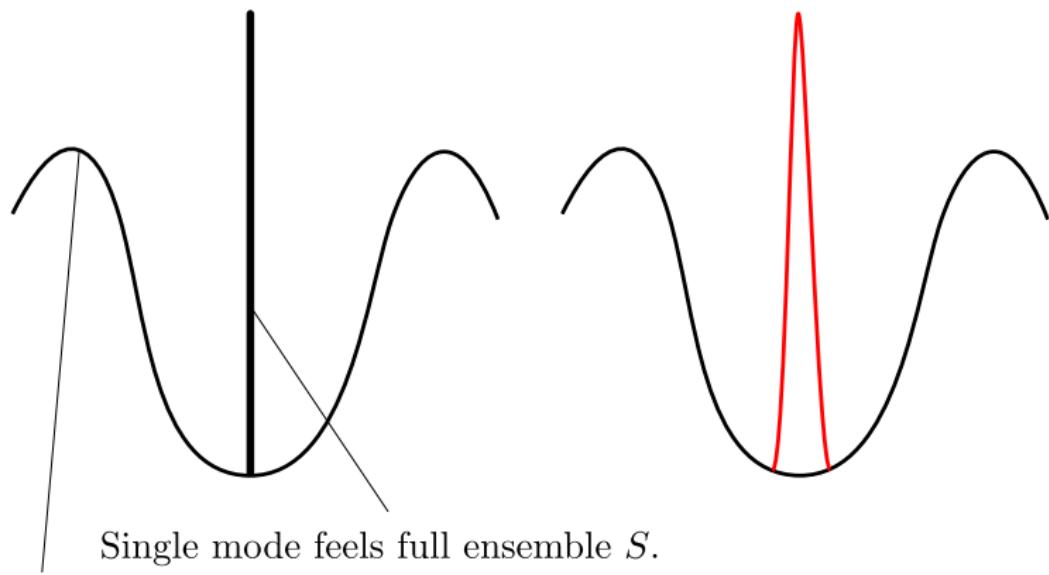
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Single mode feels full ensemble  $S$ .

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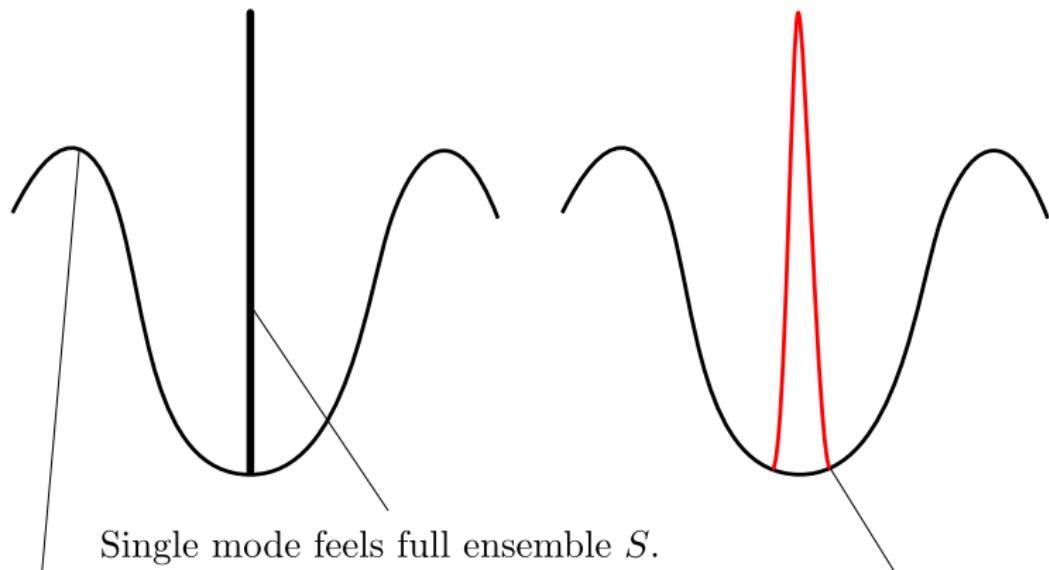


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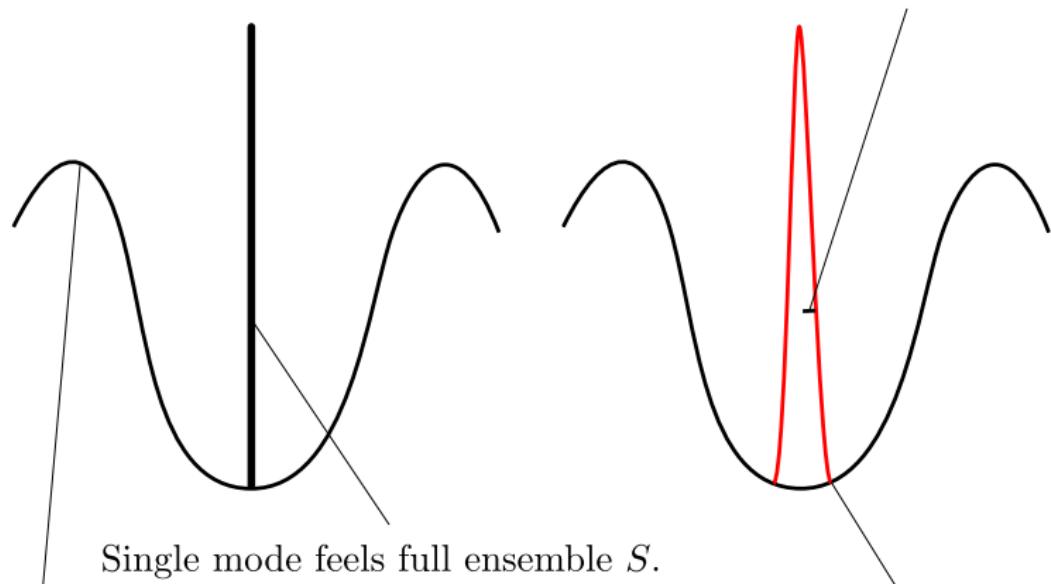
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Width proportional to  $\alpha$ .



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## Set-up

- ▶  $A$  and  $S$  deterministic.
- ▶ Broadening comes from averaging over change of basis.

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## Eigenvalue Equation

- ▶ Choice of basis

$$A = \left( \begin{array}{c|c} A' = \text{diag}(\lambda_{\nu+1}, \dots, \lambda_N) & 0_{(N-\nu) \times \nu} \\ \hline 0_{\nu \times (N-\nu)} & 0_{\nu \times \nu} \end{array} \right)$$
$$USU^\dagger = \left( \begin{array}{c|c} S_1 & S_2 \\ \hline S_2^\dagger & S_3 \end{array} \right), \quad S_3 \text{ corresponds to zero modes}$$
$$U = \begin{pmatrix} U_1 \\ U_2 \end{pmatrix}, \quad U_1 \text{ is } (N - \nu) \times N, \quad U_2 \text{ is } \nu \times N$$

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$$\begin{aligned}\det(K - \lambda \mathbf{1}_N) &= \det \left( \begin{pmatrix} A' - \alpha S_1 & \alpha S_2 \\ \alpha S_2^\dagger & \alpha S_3 \end{pmatrix} - \lambda \mathbf{1}_N \right) \\ &= \det(A' + \alpha S_1 - \lambda \mathbf{1}_{N-\nu}) \\ &\quad \times \det(\alpha S_3 - \lambda \mathbf{1}_\nu - \alpha^2 S_2^\dagger (A' + \alpha S_1 - \lambda \mathbf{1}_{N-\nu})^{-1} S_2) \\ &\stackrel{\alpha \ll 1}{=} \det(A' + \alpha S_1 - \lambda) \det(\alpha S_3 - \lambda \mathbf{1}_\nu)\end{aligned}$$

Former zero modes are determined by  $S_3$ .

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## Conditions

- ▶ Perturbation small enough

$$\alpha = o\left(\frac{1}{\|S\|_{\text{op}}} \sqrt{\frac{N}{\text{Tr}(A')^{-2}}}\right)$$

- ▶ Centred

$$\text{Tr} S = 0$$

- ▶ Sufficient mixing for limit

$$\lim_{N \rightarrow \infty} q = \infty,$$

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- ▶ Possible for all Altland-Zirnbauer classes
- ▶ Result for non-chiral classes is

$$\lim_{N \rightarrow \infty} p(S_3) \propto \exp \left[ -\frac{\gamma N^2 \text{Tr} S_3^2}{2\alpha^2 \text{Tr} S^2} \right]$$

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$$p(S') = \int_{\mathcal{K}_\nu} d\mu(U_2) \delta(S' - \kappa U_2 S U_2^\dagger)$$

- ▶ Rewrite Haar-measure

$$\int_{\mathcal{K}_\nu} d\mu(U_2) f(U_2) = \frac{\int_{\mathcal{G}_\nu} dU_2 f(U_2) \delta(\mathbf{1}_\nu - U_2 U_2^\dagger)}{\int_{\mathcal{G}_\nu} dU_2 \delta(\mathbf{1}_\nu - U_2 U_2^\dagger)}.$$

- ▶ Rewrite  $\delta$ -function

$$\delta(X) \propto \lim_{\epsilon \rightarrow 0} \int_{\mathcal{H}_\nu} dH \exp[-\epsilon \text{Tr} H^2 + i \text{Tr} XH]$$

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$$p(S') \propto \lim_{\epsilon \rightarrow 0} \int_{\mathcal{G}_\nu} dU_2 \int_{\mathcal{P}_\nu} dP \int_{\mathcal{H}_\nu} dH \exp[-\epsilon \text{Tr} H^2 + i \text{Tr} (\overbrace{S' - \kappa U_2 S U_2^\dagger}^{\text{From spectrum}}) H] \\ \times \exp[\epsilon \gamma N \text{Tr}(\mathbf{1}_\nu - iP)^2 + \gamma N \text{Tr}(\underbrace{\mathbf{1}_\nu - U_2 U_2^\dagger}_{\text{From Haar-measure}})(\mathbf{1}_\nu - iP)]$$

- ▶ Interchange integrals and perform  $U_2$ -integral

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- ▶ Expanding the determinant

$$\ln \det^{-\gamma} \left[ \mathbf{1}_{N\nu} + i \frac{S}{\gamma \sqrt{\text{Tr} S^2}} \otimes H \right] = \gamma \sum_{j=1}^{\infty} \frac{1}{j} \text{Tr} \left( -i \frac{S}{\gamma \sqrt{\text{Tr} S^2}} \right)^j \text{Tr} H^j$$

- ▶ Higher orders vanish

$$\left| \frac{\text{Tr}(S)^j}{(\text{Tr} S^2)^{j/2}} \right| \leq \frac{\|S\|_{\text{op}}^{j-2} \text{Tr} S^2}{(\text{Tr} S^2)^{j/2}} = \frac{1}{q^{j-2}} \xrightarrow{N \rightarrow \infty} 0$$

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$$\lim_{N \rightarrow \infty} p(S') \propto \lim_{N \rightarrow \infty} \lim_{\epsilon \rightarrow 0} \int_{\mathcal{H}_\nu} dH \exp[-\epsilon \text{Tr} H^2 + i \text{Tr} S' H] \\ \times \det^{-\gamma} [\mathbf{1}_{N\nu} + i\gamma^{-1} S/\sqrt{\text{Tr} S^2} \otimes H]$$

- ▶ Expanding the determinant

$$\ln \det^{-\gamma} \left[ \mathbf{1}_{N\nu} + i \frac{S}{\gamma \sqrt{\text{Tr} S^2}} \otimes H \right] = \gamma \sum_{j=1}^{\infty} \frac{1}{j} \text{Tr} \left( -i \frac{S}{\gamma \sqrt{\text{Tr} S^2}} \right)^j \text{Tr} H^j$$

- ▶ Higher orders vanish

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# Central Limit Theorem for Matrices

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$$\lim_{N \rightarrow \infty} p(S') = \frac{\int_{\mathcal{H}_\nu} dH \exp[-\text{Tr}H^2/(2\gamma) + i\text{Tr}S'H]}{\int_{\mathcal{H}_\nu} d\bar{S} \int_{\mathcal{H}_\nu} dH \exp[-\text{Tr}H^2 - \text{Tr}\bar{S}^2/4]} = \frac{\exp[-\gamma \text{Tr}S'^2/2]}{\int_{\mathcal{H}_\nu} d\bar{S} \exp[-\gamma \text{Tr}\bar{S}^2/2]}$$

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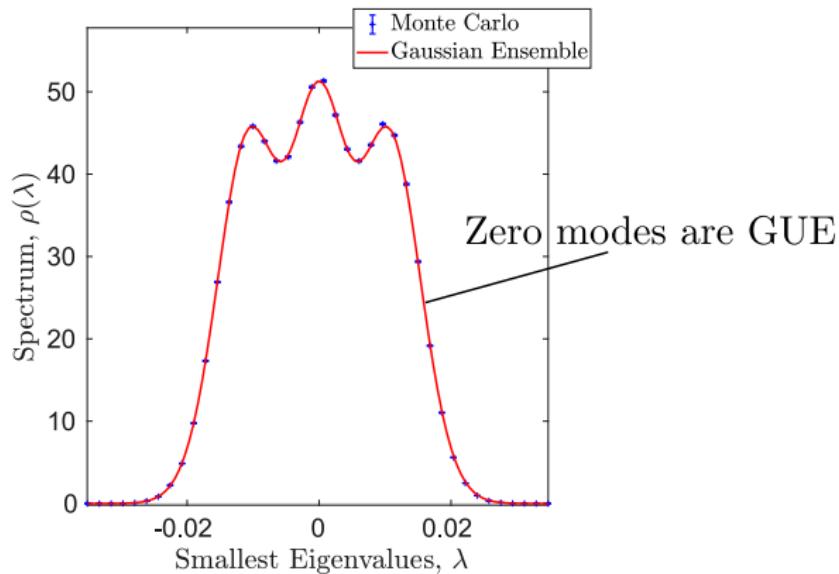
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- ▶ Assume  $\text{Tr}S^2 \sim N$  and  $\alpha$  fixed, the broadened modes scale as  $\sqrt{N}$
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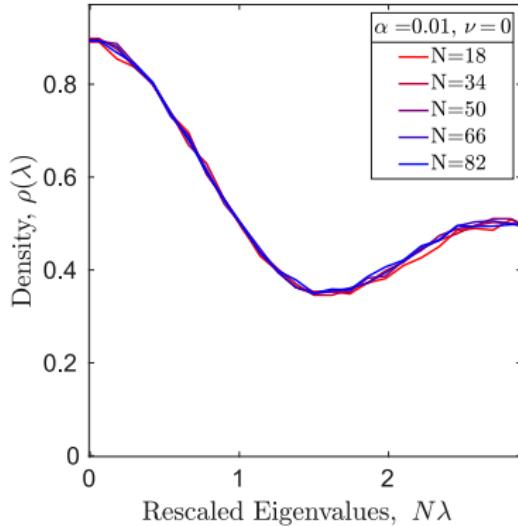
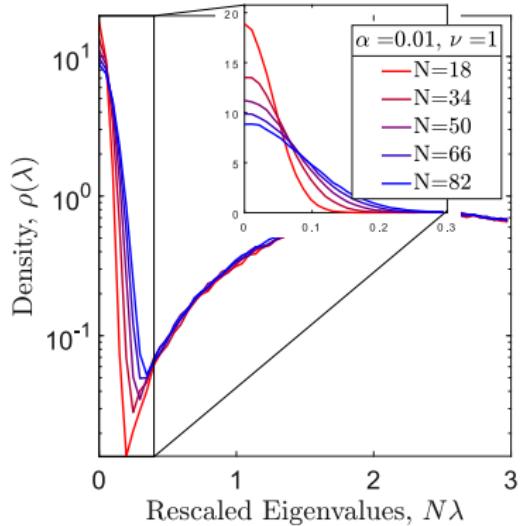
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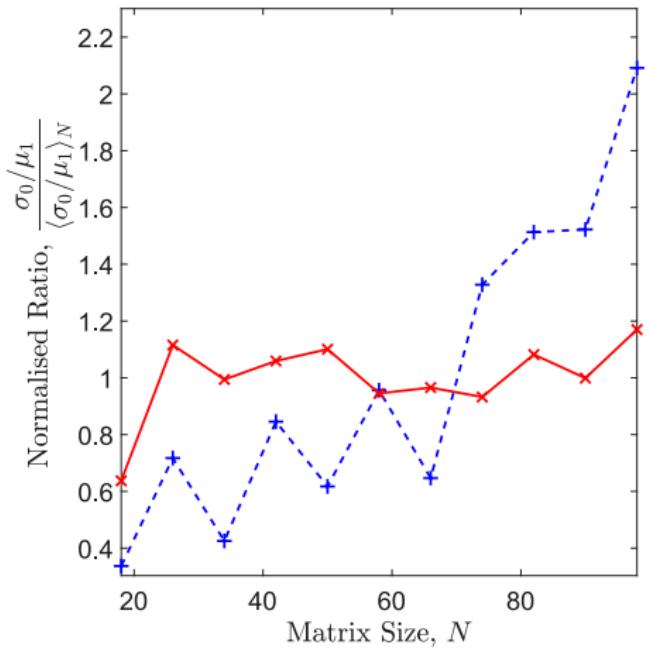
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$$K = \begin{pmatrix} iM & 0 \\ 0 & -iM \end{pmatrix} + \alpha O \begin{pmatrix} 0 & iW \\ -iW^T & 0 \end{pmatrix} O^T$$

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# **Thank you for your time!**

# Extra Slides

## Eigenvalue Equation

$$\begin{aligned} 0 &= \det(K - \lambda \mathbf{1}_N) \\ &= \det \left( \begin{array}{c|c} A' + \alpha S_1 - \lambda \mathbf{1}_{N-\nu} & \alpha S_2 \\ \hline \alpha S_2^\dagger & \alpha S_3 - \lambda \mathbf{1}_\nu \end{array} \right) \\ &= \det(A' - \lambda \mathbf{1}_{N-\nu}) \det \left( \begin{array}{c|c} \mathbf{1}_{N-\nu} + \alpha(A' - \lambda \mathbf{1}_{N-\nu})^{-1} S_1 & \alpha(A' - \lambda \mathbf{1}_{N-\nu})^{-1} S_2 \\ \hline \alpha S_2^\dagger & \alpha S_3 - \lambda \mathbf{1}_\nu \end{array} \right) . \\ &= \det(A' - \lambda \mathbf{1}_{N-\nu}) \det \left( \mathbf{1}_{N-\nu} + \alpha(A' - \lambda \mathbf{1}_{N-\nu})^{-1} U_1 S U_1^\dagger \right) \\ &\quad \times \det \left[ \alpha U_2 S U_2^\dagger - \lambda \mathbf{1}_\nu - \alpha U_2 S U_1^\dagger \left( \mathbf{1}_{N-\nu} + \alpha(A' - \lambda \mathbf{1}_{N-\nu})^{-1} U_1 S U_1^\dagger \right)^{-1} \right. \\ &\quad \left. \times \alpha(A' - \lambda \mathbf{1}_{N-\nu})^{-1} U_1 S U_2^\dagger \right] \end{aligned}$$

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- ▶ Write as Neumann sum

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## Eigenvalue Equation

- ▶ Consider the squared norm of  $\alpha U_1^\dagger (A')^{-1} U_1 S |\chi\rangle$

$$\begin{aligned}\int_{\mathcal{K}} d\mu(U) \alpha^2 \langle \chi | S U_1^\dagger (A')^{-2} U_1 S |\chi\rangle &= \frac{\alpha^2 \text{Tr}(A')^{-2}}{N} \langle \chi | S^2 |\chi\rangle \\ &\leq \frac{\alpha^2 \text{Tr}(A')^{-2} \|S\|_{\text{op}}^2}{N} \xrightarrow{\text{Condition on } \alpha} 1\end{aligned}$$

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