# Constraining BSM models by nonperturbative Higgs physics

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- same number of states also seen in lattice computations [Shrock'85-88, Maas'12, Maas&Mufti'13]
- BUT! Mismatch to lattice computations for SU(3) [Maas&Törek'16,'18]

## BEH physics - SU(3) with fundamental scalars

- Conventional perspective:  $SU(3) \rightarrow SU(2)$
- Mismatch conventional analysis vs lattice spectrum:



### elementary fields

[Maas&Törek '16,'18]

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- $\Rightarrow$  Rethinking of the particle spectrum!

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- 4. Perform standard perturbation theory on the right-hand side  $\langle (\phi^{\dagger}\phi)(x) (\phi^{\dagger}\phi)(y) \rangle = v^2 \langle h(x)h(y) \rangle + \langle h(x)h(y) \rangle^2 + \mathcal{O}(g,\lambda)$

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  - Confirmed on the lattice for SU(2)-Higgs theory [Maas'12]

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- Mapping of local to global multiplets in the Standard Model  $\rightarrow$  phenomenological implications

[Fröhlich et al '80, '81] Trento, 09.04.2019 6 / 11

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- Scale dependent field transformations:

$$\begin{split} \Gamma_{k} &= \int_{X} \left[ \frac{1}{4} F_{i\mu\nu} F_{i}^{\mu\nu} + (D_{\mu}\phi)^{\dagger} D^{\mu}\phi + U(\phi^{\dagger}\phi) \right. \\ &+ h_{\Phi} \Phi_{k} \phi^{\dagger}\phi + h_{V} V_{\mu} \mathrm{i} \phi^{\dagger} D^{\mu}\phi + \cdots \right] \\ &\partial_{t} \Phi_{k} = (\phi^{\dagger}\phi) \partial_{t} A_{k}, \qquad \partial_{t} V_{k} = (\mathrm{i} \phi^{\dagger} D^{\mu}\phi) \partial_{t} B_{k} \end{split}$$

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$$\partial_t \Phi_k = (\phi^{\dagger} \phi) \partial_t A_k, \qquad \quad \partial_t V_k = (i \phi^{\dagger} D^{\mu} \phi) \partial_t B_k$$

• Compute properties of *n*-point functions on both sides of the FMS relation

R. Sondenheimer (KFU Graz)

# Nonperturbative checks of FMS Preliminary!



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$J^P$	SU(N) invariant	SU(N-1) invariant	el. Fields	Multiplicity
0+	$\phi^{\dagger}\phi$	h	h	1
	$\phi^{\dagger} D^2 \phi$	$h, (A_{\rm s}^2), ( A_{\rm f} ^2)$		
1-	d <sup>†</sup> D <sup>µ</sup> d	Λμ	Δμ	$(N - 1)^2 - 1$
1	$\varphi^{,} D^{,} \varphi$	$\mathcal{A}_{\mathrm{S}}^{\prime}$	$A_a^{\dagger}$	(N = 1) = 1
			$ A_{\rm f}^{\mu},A_{\rm f}^{\mu} $	2(N-1)
			$A_{\rm s}^{\mu}$	1

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0+	$\phi^{\dagger}\phi$	h	h	1
	$\phi^{\dagger} D^2 \phi$	$h_{\rm r} (A_{\rm s}^2), ( A_{\rm f} ^2)$		
	$\mathrm{tr}F^{2}$	$(trF_{a}^{2}), ( F_{f} ^{2}), (F_{s}^{2})$		
1-	$\phi^{\dagger}D^{\mu}\phi$	$A^{\mu}_{ m s}$	$A^{\mu}_{ m a}$	$(N-1)^2 - 1$
	$\phi^{\dagger}F^{\mu u}D_{ u}\phi$	$(F_{\mathrm{f}}^{\mu u}A_{\mathrm{f} u})$ , $(F_{\mathrm{s}}^{\mu u}A_{\mathrm{s} u})$	$A^{\mu}_{ m f}$ , $A^{\dagger \mu}_{ m f}$	2(N-1)
			$A^{\mu}_{ m s}$	1

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- Nonperturbative check for *N* = 3 on the lattice [Maas, Törek'16,'18]
- Also mismatches between spectra for other representations or gauge groups

$$\mathrm{SU}(5) \xrightarrow{\langle \Sigma \rangle \sim w} \mathrm{SU}(3) \times \mathrm{SU}(2) \times \mathrm{U}(1) \xrightarrow{\langle \phi \rangle \sim v} \mathrm{SU}(3) \times \mathrm{U}(1) \quad (w \gg v)$$

$J^P$	SU(5) invariant	${ m SU}(3) imes { m U}(1)$ invariant	el. Fields
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• Similar considerations for SO(10), E(6),  $\cdots$ 

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	<i>O</i> <sub>0-</sub>	$h, \cdots$	$+18$ fields with $m \sim w$
	$O_{\pm 1+}$	states $\sim$ w	
	$O_{\pm 1-}$	states $\sim$ w	
1-	$O_{0+}^{\mu}$	$(Z^{\mu}), (A^{\mu}), \cdots$	$A^{\mu}$
	$O_{0-}^{\mu}$	$Z^{\mu}$ , $(A^{\mu})$ , $\cdots$	$W^{\pm\mu}$
	$O_{\pm 1+}^{\mu}$	states $\sim$ w	$Z^{\mu}$
	$O_{\pm 1-}^{\mu}$	states $\sim$ w	$G^{\mu}(8)$
			$ $ +12 Leptoquarks with $m \sim w$

- Similar considerations for  $\mathrm{SO}(10),\,\mathrm{E}(6),\,\cdots$
- Rethinking of GUT construction!

### Conclusions & Outlook

- Observable spectrum must be gauge invariant
- Non-Abelian gauge theory: composite operator
- FMS mechanism provides a mapping of local to global multiplets
- Same results in leading order for the standard model
- BSM model building can be affected

## Thank you for your attention!