

Stability of three neutrino flavor conversion in supernovae

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Together with: Rasmus S. L. Hansen and Manfred Lindner

14.5.2019 @Trento SN neutrinos at the crossroads

based on: arXiv:1905.03647

Content

- Motivation
- Neutrinos in dense media
- Three flavor stability analysis
- Application to toy-systems
- Analysis in basis of propagting states
- Conclusion

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- Extend analysis to three flavors and search for instabilities

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 $\begin{array}{ccc}
u_x &
u_e \\
&
u_y \end{array}$

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 ρ_i

Y

• Ensemble of neutrinos is described by a $\begin{pmatrix} \nu_e \\ \nu_x \\ \nu_\mu \end{pmatrix} := R_{23}^{\dagger} \begin{pmatrix} \nu_e \\ \nu_\mu \\ \nu_{\tau} \end{pmatrix}$

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$$i \partial_t \rho_i(t, \mathbf{v}) = [H_i(t, \mathbf{v}), \rho_i(t, \mathbf{v})]$$

 $H_i(t, \mathbf{v}) = H_i^{\text{vac}} + H^{\lambda} + H_i^{\nu\nu}(\mathbf{v})$







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ee e $\begin{array}{ccc}
 \nu_x & \nu_e \\
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 \end{array}$ $H_i^{\nu\nu}(\mathbf{v})$ ρ_i $\sqrt{2}G_F \operatorname{diag}(N_e, N_\mu, N_\tau)$

0.2 sec e e^{-} e 108 1.0 sec ν_e ν_y (km ⁻¹) 2.0 sec ν_x 106 5.0 sec 7.0 sec Ar. Hr ρ_i 104 ωH 10 10² 104 105 103 r (km) $\sqrt{2}G_F \operatorname{diag}(N_e, N_\mu, N_\tau)$ Mirizzi et al: arXiv:1508.00785]

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$$\mathcal{S}^{i}_{\alpha\beta}(t,\mathbf{v}) = Q^{i}_{\alpha\beta}e^{-i\Omega t}$$

$$\alpha \neq \beta$$

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 $\begin{bmatrix} \omega_i \underline{\underline{A}}(\theta_{12}, \theta_{13}) + \underline{\underline{B}}(\theta_{13}) \eta \omega_i \end{bmatrix} \begin{pmatrix} Q_{ex}^i \\ Q_{ey}^i \\ Q_{xy}^i \end{pmatrix}$

 $\Omega\left(egin{matrix}Q^i_{ex}\Q^i_{ey}\Q^i_r\end{matrix}
ight)$

• Time-Derivative:

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• (II) Matter-Term:

 $\begin{bmatrix} \omega_{i}\underline{\underline{A}}(\theta_{12},\theta_{13}) + \underline{\underline{B}}(\theta_{13})\eta\omega_{i} \end{bmatrix} \begin{pmatrix} Q_{ex}^{i}\\ Q_{ey}^{i}\\ Q_{xy}^{i} \end{pmatrix}$ $\begin{pmatrix} \lambda & 0 & 0\\ 0 & \lambda & 0\\ 0 & 0 & 0 \end{pmatrix} \begin{pmatrix} Q_{ex}^{i}\\ Q_{ey}^{i}\\ Q_{ey}^{i} \\ Q_{ey}^{i} \end{pmatrix}$

 $\Omega \begin{pmatrix} Q_{ex}^i \\ Q_{ey}^i \\ Q^i \end{pmatrix}$

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• (III) Nu-Nu-Term:

$$\begin{bmatrix} \omega_{i}\underline{A}(\theta_{12},\theta_{13}) + \underline{B}(\theta_{13})\eta\omega_{i} \end{bmatrix} \begin{pmatrix} Q_{ex}^{i}\\ Q_{ey}^{i}\\ Q_{xy}^{i} \end{pmatrix}$$
$$\begin{pmatrix} \lambda & 0 & 0\\ 0 & \lambda & 0\\ 0 & 0 & 0 \end{pmatrix} \begin{pmatrix} Q_{ex}^{i}\\ Q_{ey}^{i}\\ Q_{xy}^{i} \end{pmatrix}$$
$$3\mu \sum_{j=1}^{N} g_{j}(1 - \mathbf{v}^{i}\mathbf{v}^{j}) \begin{pmatrix} (Q_{ex}^{i} - Q_{ex}^{j})\\ (Q_{ey}^{i} - Q_{ey}^{j})\\ 0 \end{pmatrix}.$$

 Q_{ex}^i Q_{ey}^i

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 $\Omega \begin{pmatrix} Q_{ex}^i \\ Q_{ey}^i \\ O^i \end{pmatrix}$

<u>Aim</u>: Solving the eigenvalue equation and looking for imaginary omega



 $egin{array}{ccc} ar{
u}_y & ar{
u}_e & \ ar{
u}_x & ar{Q}_{ar{L}} \end{array}$ $egin{array}{ccc}
u_x &
u_e \
u_y &
u_y \end{array}$ \underline{Q}_R



[Chakraborty et al: arXiv:1602.00698]



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 $A_{\pm} := \frac{1}{2}(L \pm R)$

 $\bar{A}_{\pm} := \frac{1}{2}(\bar{L} \pm \bar{R})$





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$$H'_{\text{diag}} \approx \begin{pmatrix} \lambda + \omega_i (s_{12}^2 c_{13}^2 + \eta s_{13}^2) & 0 & 0 \\ 0 & \omega_i c_{12}^2 & 0 \\ 0 & 0 & \omega_i (s_{12}^2 s_{13}^2 + \eta c_{13}^2) \end{pmatrix}$$

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Basis of propagating states

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$$\omega_{\text{eff}} = \omega(-c_{12}^2 + c_{13}^2 s_{12}^2 + \eta s_{13}^2)$$



[Duan et al:arXiv:astro-ph/0606616]

[Chakraborty et al: arXiv:1507.07569]

Summary

- For high neutrino densities as in a CcSN neutrino-neutrino interactions must be considered
- We studied collective oscillations in a large matter background including all three flavors and search for flavor instabilities
- Besides known flavor instabilities for IO, we find instabilities also for NO in simple toy models
- These new instabilities can be understood as originating from an effective two-neutrino system with IO