Flavor evolution of dense neutrinos as an emergent phenomenon

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Flavor Conversion in Core-Collapse Supernovae



Why worry about detailed neutrino transport?





- Explosion mechanism: Shock-wave revival by nu energy deposition
- Nucleosynthesis in neutrino irradiated outflows in SNe and NS-mergers depends on flavor (beta reactions!)
- Signal interpretation of DSNB and next nearby SN
- Collective flavor conversion: interesting theoretical problem in its own right



Kinetic Equation for Neutrino Transport

Flavor-dependent phase-space densities (occupation number matrices)

$$\varrho = \begin{pmatrix} f_{\nu_e} & f_{\langle \nu_e | \nu_\mu \rangle} & f_{\langle \nu_e | \nu_\tau \rangle} \\ f_{\langle \nu_\mu | \nu_e \rangle} & f_{\nu_\mu} & f_{\langle \nu_\mu | \nu_\tau \rangle} \\ f_{\langle \nu_\tau | \nu_e \rangle} & f_{\langle \nu_\tau | \nu_\mu \rangle} & f_{\nu_\tau} \end{pmatrix}$$

Diagonal: Usual occupation numbers Off-diag: Flavor coherence information

and similar for $\overline{\nu}$

Transport equation

$$\left(\partial_t + \vec{v} \cdot \vec{\nabla}_x - \vec{F} \cdot \vec{\nabla}_p\right) \varrho(t, \vec{x}, \vec{p}) = -\mathrm{i} \left[\mathcal{H}(t, \vec{x}, \vec{p}), \varrho(t, \vec{x}, \vec{p})\right] + \mathcal{C}[\varrho(t, \vec{x}, \vec{p})]$$

Streaming

Gravitational forces (redshift, deflection)

Typical approximations in numerical simulations:

(Angular moments, ray-by-ray, ...)

No flavor conversion (large matter effect!)

Reducing 6+1 dimensions

No gravitational deflection

• 3-species transport: v_e, \overline{v}_e, v_x

No muons

Flavor oscillations (vacuum, matter, vv) Collisions

•
$$e^- + p \rightleftharpoons n + v_e$$

• $e^+ + n \rightleftharpoons p + \bar{v}_e$

•
$$e^- + A \rightleftharpoons v_e + A^*$$

- $v + n, p \rightleftharpoons v + n, p$
- $\nu + A \rightleftharpoons \nu + A$
- $v + e^{\pm} \rightleftharpoons v + e^{\pm}$
- $N + N \rightleftharpoons N + N + \nu + \bar{\nu}$
- $e^+ + e^- \rightleftharpoons v + \bar{v}$
- $v_x + v_e, \bar{v}_e \rightleftharpoons v_x + v_e, \bar{v}_e$ $(v_x = v_\mu, \bar{v}_\mu, v_\tau, \text{ or } \bar{v}_\tau)$

• $v_e + \bar{v}_e \rightleftharpoons v_{\mu,\tau} + \bar{v}_{\mu,\tau}$

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Streaming

Gravitational forces (redshift, deflection)

Flavor oscillations (vacuum, matter, vv)

Collisions

Flavor evolution governed by "Hamiltonian matrix" (here for 2 flavors)

$$\mathcal{H} = \frac{\Delta m^2}{4E} \begin{pmatrix} \cos 2\theta & \sin 2\theta \\ \sin 2\theta & -\cos 2\theta \end{pmatrix} + \sqrt{2}G_F \begin{pmatrix} n_e & 0 \\ 0 & 0 \end{pmatrix} + \sqrt{2}G_F \int \frac{d^3\vec{p}}{(2\pi)^3}(\varrho + \overline{\varrho})$$
Vacuum oscillations
MSW effect
Nu-nu interactions,
nus feed back on each other

• Flavor evolution is caused by off-diagonal \mathcal{H} elements (vacuum or nu-nu term) • For $\Delta m^2 = 0$, nu-nu term can still cause run-away modes!

Self-Induced Flavor Conversion

Flavor conversion (vacuum or MSW) for a neutrino of given momentum p

 Requires lepton flavor violation by masses and mixing

$$\nu_{\boldsymbol{e}}(p) \to \nu_{\boldsymbol{\mu}}(p)$$

$$\frac{\Delta m_{\rm atm}^2}{2E} = 10^{-10} \,\rm eV = 0.5 \,\rm km^{-1}$$

Pair-wise flavor exchange by $\nu - \nu$ refraction (forward scattering)

- No net flavor change of pair
- Requires dense neutrino medium (collective effect of interacting neutrinos)
- Can occur without masses/mixing (and then does not depend on $\Delta m^2/2E$)
- Familiar as neutrino pair process $\mathcal{O}(G_F^2)$ Here as coherent refractive effect $\mathcal{O}(G_F)$

$$\begin{split} \nu_{e}(p) + \overline{\nu}_{e}(k) &\rightarrow \nu_{\mu}(p) + \overline{\nu}_{\mu}(k) \\ \nu_{e}(p) + \nu_{\mu}(k) &\rightarrow \nu_{\mu}(p) + \nu_{e}(k) \end{split}$$

$$\sqrt{2}G_{\rm F}n_{\nu} = 10^{-5}{\rm eV} = 0.5~{\rm cm}^{-1}$$

E = 12.5 MeV R = 80 km $L_{\nu} = 40 \times 10^{51} \text{erg/s}$

Transport of Particles and Flavor Coherence

How to interpret the Liouville operator (advection term)? $\oint \mathcal{D} \varrho(t, \vec{x}, \vec{p}) = -i \left[\mathcal{H}(t, \vec{x}, \vec{p}), \varrho(t, \vec{x}, \vec{p}) \right] + \mathcal{C}[\varrho(t, \vec{x}, \vec{p})]$

• Flavor evolution along trajectories (of individual neutrinos or wave packets)

$$\frac{d}{ds}\varrho(\underbrace{t,\vec{x}},\vec{p}) = -i \left[\mathcal{H}(t,\vec{x},\vec{p}), \varrho(t,\vec{x},\vec{p})\right] + \mathcal{C}[\varrho(t,\vec{x},\vec{p})]$$

$$\downarrow$$
Time and space coordinates related by neutrino velocity along trajectory

• Flavor evolution in space and time (of neutrino radiation field)

$$\partial_t \varrho + \frac{1}{2} \{ \underbrace{\vec{\nabla}_p \mathcal{H}}_{\downarrow}, \vec{\nabla}_x \varrho \} - \frac{1}{2} \{ \vec{\nabla}_x \mathcal{H}, \vec{\nabla}_p \varrho \} = -i [\mathcal{H}, \varrho] + \mathcal{C}[\varrho]$$

Velocity operator (matrix in flavor space)

Neutrino quantum field obeys the Dirac equation Kinetic equation for ϱ is always an approximation / simplification / expansion Different intuitive pictures should give the same answer for observables

Velocity in Advection Term

- Advection term after "separation of scales" between gradients and neutrino $ec{p}$
- Wigner function $\varrho(t, \vec{x}, \vec{p})$ makes sense as a phase-space density

$$\partial_t \varrho + \frac{1}{2} \{ \underbrace{\vec{\nabla}_p \mathcal{H}}_{\downarrow}, \vec{\nabla}_r \varrho \} - \frac{1}{2} \{ \vec{\nabla}_x \mathcal{H}, \vec{\nabla}_p \varrho \} = -i [\mathcal{H}, \varrho] + \mathcal{C}[\varrho]$$

$$\downarrow$$
Velocity operator (matrix in flavor space) $\vec{\nabla}_p E_{kin} = \vec{\nabla}_p \sqrt{\vec{p}^2 + m^2} = \frac{\vec{p}}{E} = \vec{v}$

Ultrarelativistic limit and ignore gravity or refractive deflection

Leading-order effect of ν masses on the rhs in the usual oscillation term

Stirner, Sigl & Raffelt: *Liouville term for neutrinos: Flavor structure and wave interpretation*, JCAP 1805 (2018) 016 [arXiv:1803.04693]

Georg Raffelt, MPI Physics, Munich

Kinetic equations for flavor transport

Rudszky: *Kinetic equations for neutrino spin- and type-oscillations in a medium,* Astrophys. Space Sci. 165 (1990) 65

Raffelt & Sigl: *General kinetic description of relativistic mixed neutrinos*, NPB 406 (1993) 423

Sirera & Perez: *Relativistic Wigner function approach to neutrino propagation in matter,* PRD 59 (1999) 125011 [hep-ph/9810347]

Yamada: **Boltzmann equations for neutrinos with flavor mixings**, PRD 62 (2000) 093026 [astro-ph/0002502]

Cardall: *Liouville equations for neutrino distribution matrices,* PRD 78 (2008) 085017 [arXiv:0712.1188]

Vlasenko, Fuller & Cirigliano: *Neutrino quantum kinetics*, PRD 89 (2014) 105004 [arXiv:1309.2628]

Serreau & Volpe: *Neutrino-antineutrino correlations in dense anisotropic media,* PRD 90 (2014) 125040 [arXiv:1409.359]

Hansen & Smirnov: *The Liouville equation for flavour evolution of neutrinos and neutrino wave packets,* JCAP 1612 (2016) 019 [arXiv:1610.00910]

Stirner, Sigl & Raffelt: *Liouville term for neutrinos: Flavor structure and wave interpretation,* JCAP 1805 (2018) 016 [arXiv:1803.04693]

Correlated Trajectories vs. Field of Flavor Coherence

Assume globally spherically symmetric neutrino emission from SN core



- Every ν meets every other ν at most once
- Nonlinear feedback on flavor evolution?



- Oscillating (or unstable) field $\varrho(r)$ of flavor coherence, acting back upon itself
- Do not worry about individual neutrinos

Evolution of the Questions

Bulb model of neutrino emission:

- Nu-nu interaction determined by aspect ratio of emission surface
- Instability as a function of radius adiabatic conversion possible
- Flavor pendulum, spectral splits, multi-angle matter effect, three-flavor effects, ...
- Spurious instabilities (need many angle bins in numerical studies)
- Instability in the transverse direction: Spontaneous symmetry breaking



Evolution of the Questions

Halo effect:

- Small re-scattered flux, much larger angular leverage
- Impact on collective effects?
- How to deal with backward flux?



Evolution of the Questions

Non-stationarity:

- Time-variation (of SN emission) in source region?
- Self-induced time variation, "pulsating modes" more unstable than stationary ones?



Fast Flavor Conversion

Flavor evolution governed by "Hamiltonian matrix" (here for 2 flavors)

$$\mathcal{H}_{\vec{p}} = \frac{\Delta m^2}{4E} \begin{pmatrix} \cos 2\theta & \sin 2\theta \\ \sin 2\theta & -\cos 2\theta \end{pmatrix} + \mathcal{H}_{mat} + \sqrt{2}G_F \int \frac{d^3\vec{p}'}{(2\pi)^3} (1 - \vec{v} \cdot \vec{v}') \begin{pmatrix} f_{\nu_e} & f_{\langle \nu_e | \nu_x \rangle} \\ f_{\langle \nu_x | \nu_e \rangle} & f_{\nu_x} \end{pmatrix}$$

Flavor conversion caused by off diagonals

Energy scales of the problem:

 $\begin{array}{ll} \mu = \sqrt{2} \ G_{\rm F} \ n_{\nu \overline{\nu}} & \mbox{Required for any collective effects} \\ \omega_E = \Delta m^2 / 2E & \mbox{Vacuum oscillation frequency} \\ \lambda = \sqrt{2} \ G_{\rm F} \ n_{\rm e} & \mbox{Matter effect} \end{array}$

Slow modes:

Require $\omega_E \neq 0$ Possible growth rate: $\kappa \sim \sqrt{\mu \omega_E}$, requires "crossing" of ω_E distribution

Fast modes:

Dynamical even for $\omega_E = 0$ Growth rate: $\kappa \sim \sqrt{\mu \omega_E}$ slow growth $\kappa \sim \mu$ fast growth, requires "crossing" of angle distribution

Neutrino gas in the near-free streaming regime



Linearisation for Fast Flavor Modes

Evolution equation: $iv^{\alpha}\partial_{\alpha}\varrho = [\mathcal{H}, \varrho]$ with $v^{\alpha} = (1, \vec{v})$

Linearisation to find (propagating or unstable) collective modes:

$$\varrho = \begin{pmatrix} f_{\nu_e} & f_{\langle \nu_e | \nu_x \rangle} \\ f_{\langle \nu_x | \nu_e \rangle} & f_{\nu_x} \end{pmatrix} = \frac{f_{\nu_e} + f_{\nu_x}}{2} + \frac{f_{\nu_e} - f_{\nu_x}}{2} \begin{pmatrix} s & S \\ S^* & -s \end{pmatrix}$$
 Field of flavor coherence

Linearised EOM for field of flavor coherence

$$iv^{\alpha}\partial_{\alpha}S_{\vec{p}} = \underbrace{\left(\frac{\Delta m^{2}}{2E} + v^{\alpha}\Lambda_{\alpha}\right)}_{QE}S_{\vec{p}} - \mu v^{\alpha}\int \underbrace{\frac{d^{3}\vec{p}'}{(2\pi)^{3}}v_{\alpha}'\left(g_{\vec{p}'}S_{\vec{p}'} - \overline{g}_{\vec{p}'}\overline{S}_{\vec{p}'}\right)}_{QE}$$

Ignore for fast modes

Same for all *E* and ν and $\overline{\nu}$

Angle distribution of electron lepton number (ELN) carried by neutrinos

$$G_{\vec{v}} = \int \frac{dE \ E^2}{2\pi^2} \frac{f_{\nu_e,\vec{p}} - f_{\overline{\nu}_e,\vec{p}} - f_{\nu_x,\vec{p}} + f_{\overline{\nu}_x,\vec{p}}}{2}$$

Linearised EOM for field of flavor coherence

$$iv^{\alpha}(\partial_{\alpha} + i\Lambda_{\alpha})S_{\vec{v}} = -\mu v^{\alpha} \int \frac{d\vec{v}'}{4\pi} v_{\alpha}' G_{\vec{v}'} S_{\vec{v}'}$$

Matter effect, "rotate away" by including it in derivative if medium is homogeneous and stationary

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Dispersion Relation for Fast Flavor Modes

Linearised EOM for field of flavor coherence – a wave equation

$$iv^{\alpha}\partial_{\alpha}S_{\vec{v}} = -\mu v^{\alpha} \int \frac{d\vec{v}'}{4\pi} v_{\alpha}' G_{\vec{v}'} S_{\vec{v}'}$$

Plane-wave ansatz

$$S_{\vec{v}}(t,\vec{r}) = Q_{\vec{v}}(\Omega,\vec{K}) e^{-i(\Omega t - \vec{K} \cdot \vec{r})}$$

EOM in Fourier space

$$\left(\Omega - \vec{v} \cdot \vec{K}\right)Q_{\vec{v}} = -\mu \int \frac{d\vec{v}'}{4\pi} (1 - \vec{v} \cdot \vec{v}') G_{\vec{v}'}Q_{\vec{v}'}$$

Non-collective solutions:

 (Ω, \vec{K}) real and "below the light cone" $(\Omega - \vec{v} \cdot \vec{K}) = 0$ for some mode \vec{v}

Collective solutions:

$$(\Omega - \vec{v} \cdot \vec{K}) \neq 0$$
 for all \vec{v} modes

 (Ω, \vec{K}) real and "outside the light cone" or imaginary part

Dispersion relation:

 $\Omega = \vec{v} \cdot \vec{K}$ for every \vec{K} continuos infinity of frequencies

Eigenfunctions
$$Q_{\vec{v}} \propto \frac{\vec{a} \cdot \vec{v} + b}{\Omega - \vec{v} \cdot \vec{K}}$$

Dispersion relation: det $\Pi = 0$
 $\Pi^{\mu\nu} = \eta^{\mu\nu} + \int \frac{d\vec{v}}{4\pi} G_{\vec{v}} \frac{v^{\mu}v^{\nu}}{\Omega - \vec{v} \cdot \vec{K}}$

Fast Pairwise Conversion of Supernova Neutrinos: A Dispersion Relation Approach

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Classification of instabilities of "flavor waves" (Two-beam model)



Classification of instabilities of plasma waves (Two-beam model)



Landau & Lifshitz, Vol.10, Physical Kinetics Chapter VI, Instability Theory

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SN Neutrinos at the Cross Roads, Trento, 13–17 May 2019

Dispersion Relation for Neutrino Flavor Waves



Fast Flavor Waves



Non-collective modes: Infinitely many $\Omega = \vec{v} \cdot \vec{K}$

- Flavor coherence carried by every neutrino mode separately
- Quick kinematical decoherence

Collective modes: $\Omega(\vec{K})$ according to collective dispersion relation • Flavor wave (or wave packet) propagates and/or grows

Dispersion Relation for Isotropic Case



Dispersion Relation for Isotropic Case



Fast Flavor Waves

Non-collective modes:



Collective modes



- Infinitely many neutrino velocity projections on \vec{K}
- Each carries along its initial flavor coherence
- Kinematical decoherence of initial wave packet (Does not happen in two-beam model)
- Fast dissipation of any initial wave packet

- Infinitely many neutrino velocity projections on \vec{K}
- Move through wave packet (here taken with vanishing central wave number)
- Wave packet moves in neutrino gas,
- independently of velocities of neutrino "beams"

Wave packet of flavor coherence

Dispersion Relation vs. Eigenvalues of Hamiltonian



Dispersion relation:

For fixed μ find $\Omega(K)$ from

$$\left(\Omega - \vec{v} \cdot \vec{K}\right)Q_{\vec{v}} = -\mu \int \frac{d\vec{v}'}{4\pi} (1 - \vec{v} \cdot \vec{v}') G_{\vec{v}'}Q_{\vec{v}'}$$

Eigenvalues of Hamiltonian:

For fixed *K* find eigenvalues $w = \Omega(\mu)/K$ of \mathcal{H}

$$i\partial_t S_{\vec{v}}(t,\vec{K}) = \mathcal{H}(S_{\vec{v}})$$

$$\mathcal{H}(S_{\vec{v}}) = \vec{v} \cdot \vec{K} \, S_{\vec{v}} - \mu \int \frac{d\vec{v}'}{4\pi} (1 - \vec{v} \cdot \vec{v}') \, G_{\vec{v}'} S_{\vec{v}'}$$



Bound vs Scattering States



Continuum limit: Box size $\rightarrow \infty$

Continuum of scattering states

Continuum of scattering states (with phase shifts) + Bound state

Non-collective modes~ Scattering statesCollective modes~ Bound states

No spectral crossing



"Weak" spectral crossing



"Strong" spectral crossing



Spectral crossing – Continuous Limit

Energy levels





- Solutions with complex eigenvalues appear as merging of two real eigenvalues
- With increasing μ must emerge below the light cone

Continuous limit of vanishing mode spacing:

- Critical points at $w = \cos(\theta_0)$, i.e. at crossing where $G(\cos \theta_0) = 0$
- Interaction strength $\mu_{1,2}$ of critical points follow easily from $G(\cos \theta)$
- Single crossing: Complex solution guaranteed
- Several crossings: Not guaranteed

 \rightarrow See Tobias Stirner's talk

Summary

- Neutrino-neutrino interactions lead to emergence of collective modes of flavor coherence (propagating or unstable)
- Need not exist for every \vec{K}/μ (dispersion relations can end)
- Co-exist with non-collective modes
- "Wave packet of flavor coherence" dissipates by kinematical decoherence between non-collective modes
- Contains non-dissipating (propagating or growing) projection for sufficiently strong nu-nu interaction effect
- Explicit formulation of eigenfunctions for non-collective modes leads to simple identification of critical points → Tobias Stirner's talk
- Stable collective modes "peel off" from the light cone and exist only outside
- Unstable collective modes begin/end under the light cone from coalescence of non-collective modes

Capozzi, Raffelt & Stirner, work in progress (2019)

Many Open Questions

Flavor evolution in dense neutrino flows still on the level of simplified toy models and parametric studies

- Realistic normal-mode analysis without symmetry assumptions?
- Realistic triggering of stable or unstable flavor waves?
- Do tachyonic modes really lead to flavor equilibration? (Going beyond linearised stability analysis)
- Realistic impact on SN explosion and nucleosynthesis?



It is only the beginning. A lot more work ahead of us ...