

Towards analog reheating of the universe in the laboratory

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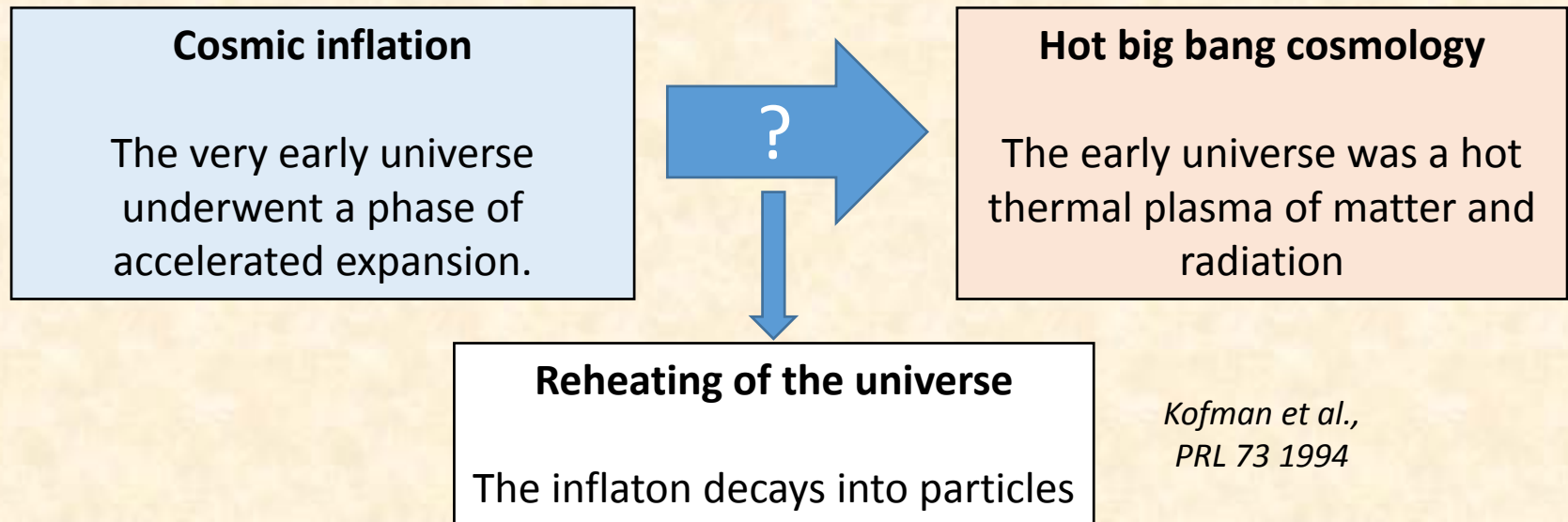


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Synthetic
Quantum
Systems



Reheating of the universe after inflation



What aspects of the reheating dynamics can be simulated with trapped atomic gases?

- Cosmic expansion
- Particle production after inflation
- Thermalization

Cosmic expansion in the lab

- Cosmology:
 - Expanding universe
 - Flat **FRW metrics**: characterized by a scale factor $a(t)$

$$\Delta x_{\text{phys}} = a(t)\Delta x$$

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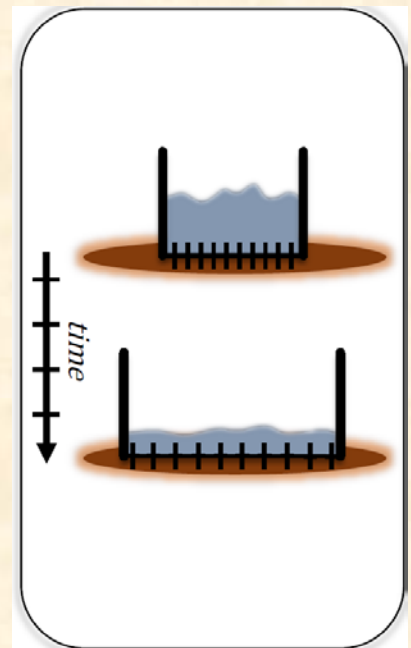
- Expansion in a trapped Bose gas:

$$i\hbar \frac{\partial \psi}{\partial t} = \left(-\frac{\hbar^2 \Delta}{2m} + U_{\text{ext}} \right) \psi + g\psi^\dagger \psi \psi,$$

+

physically expand the trapping potential

$$U_{\text{ext}} \rightarrow U_{\text{ext}}[a(t)]$$

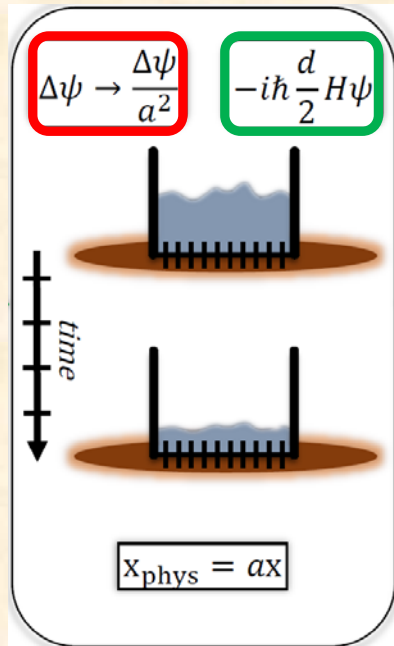


M. Uhlmann et al. New J. Phys. 7, 2005, Fedichev and Fischer PRL 91, 2003

S. Eckel et al. PRX 8 2018

Cosmic expansion in the lab

Alternative approach: identify the (lab) coordinates with co-moving ones:



$$i\hbar \frac{\partial \psi}{\partial t} = \left(-\frac{\hbar^2}{2m} \frac{\Delta}{a^2} + U_{\text{ext}} \right) \psi + g\psi^+ \psi \psi - i\hbar \frac{d}{2} H\psi,$$

Red-shift

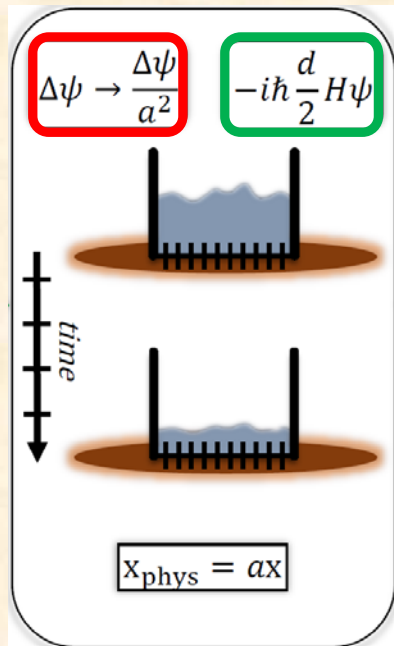
Hubble friction: $H = \frac{\dot{a}}{a}$

Non-relativistic limit of KG equation

$$\frac{1}{c^2} \ddot{\varphi} + \frac{1}{c^2} dH \dot{\varphi} - \frac{\Delta \varphi}{a^2} + \frac{\delta U}{\delta \varphi} = 0,$$

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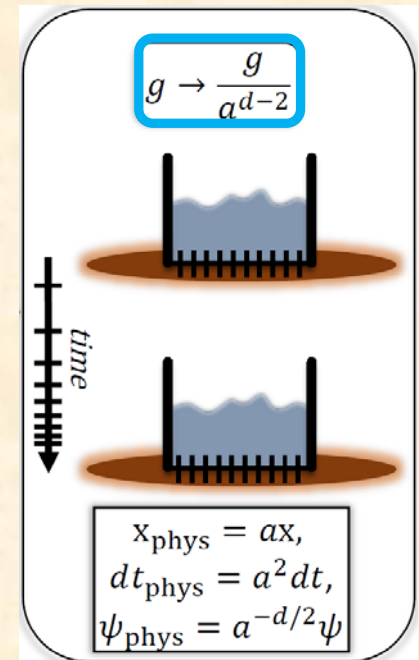
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- Transform to new variables

$$\psi = \tilde{\psi} a^{-d/2}, \quad dt = a^2 d\tau,$$

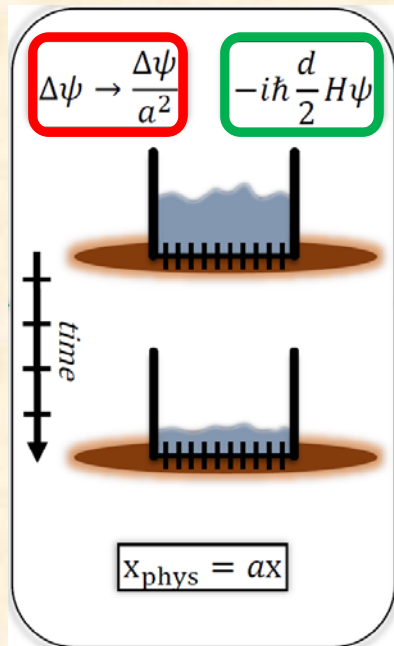
$$i\hbar \frac{\partial \tilde{\psi}}{\partial \tau} = \left(-\frac{\hbar^2}{2m} \Delta + \tilde{U}_{\text{ext}} \right) \tilde{\psi} + \left(\frac{g}{a^{d-2}} \right) \tilde{\psi}^+ \tilde{\psi} \tilde{\psi}.$$

Time-dependent coupling



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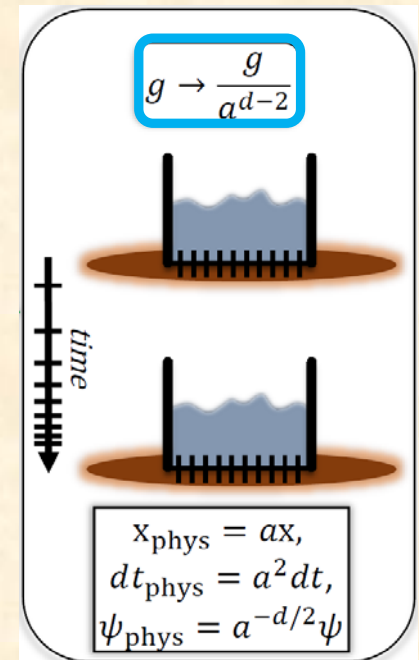
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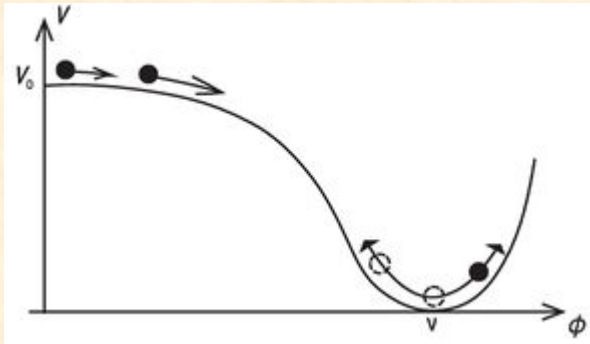


- $d = 2$: scale-invariance: *Pitaevskii and Rosch PRA 55, 1997, Saint-Jalm et al, PRX9, 2019,*
expansion enters only via back-transformation of the variables:

Particle production after inflation

The universe after inflation:

- The inflaton oscillates around the minimum of its potential
- Decays into particles



Ultracold atomic gas:

The total number of atoms is conserved

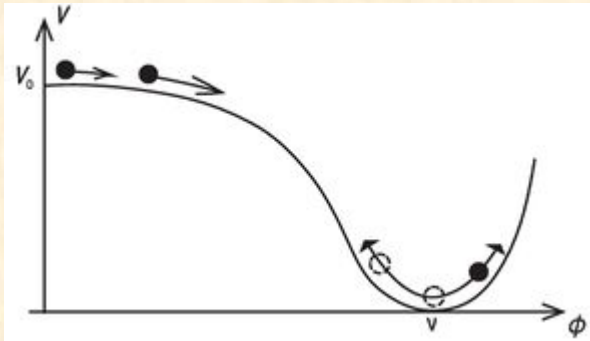


Prevents the decay of the condensate

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Prevents the decay of the condensate

Early stages (preheating):

- Particles/fluctuations live on an effective $\phi(t)$ -dependent potential

$$\delta\ddot{\varphi}_{\mathbf{p}} + M_{\mathbf{p}}^2(\phi(t))\delta\varphi_{\mathbf{p}} = 0$$



Parametric resonance
(explosive particle production)

Sound waves on top of the BEC:

$$\ddot{\Phi}_{\mathbf{p}} + \frac{gn_{\text{cond}}}{m} p^2 \Phi_{\mathbf{p}} = 0$$

Modulate the scattering length: $g(t)$

*Robertson et al.,
PRD 98, 2018*

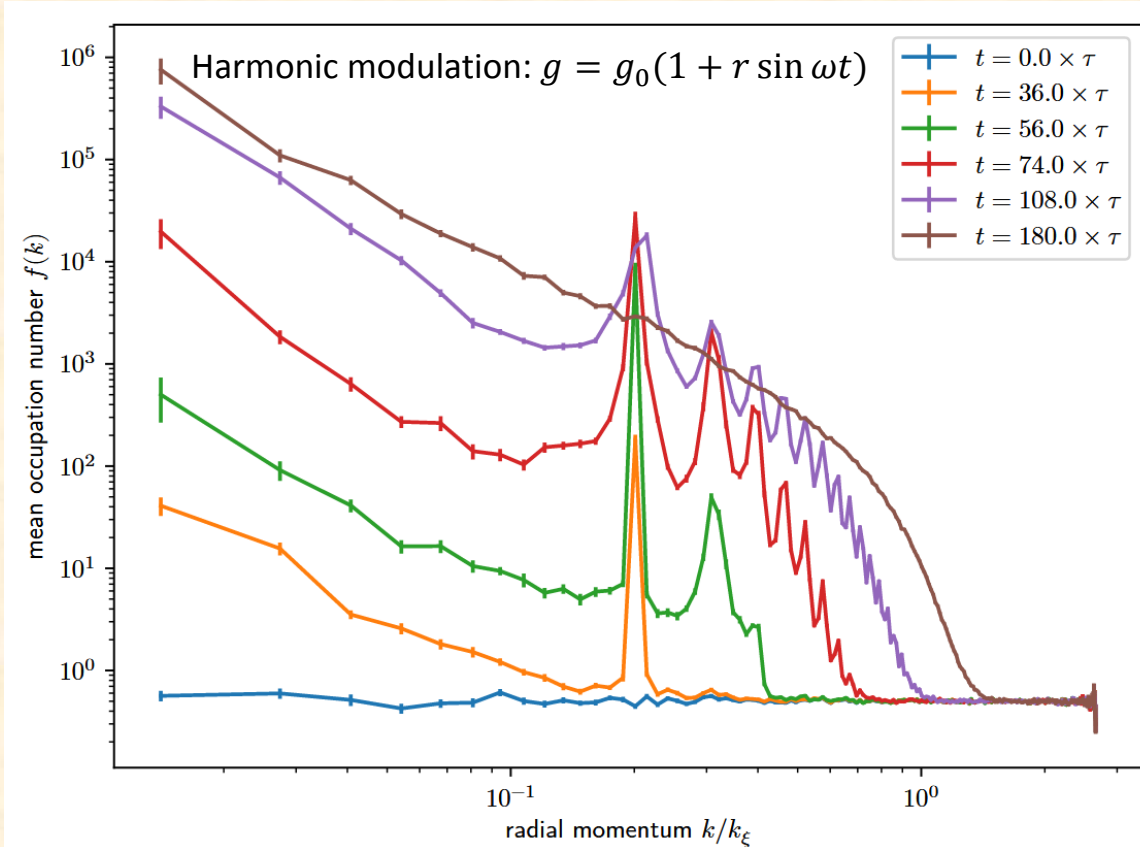
*Pollack et al.,
PRA 81, 2010*

*Vidanović et al.,
PRA 84, 2011*

*Nguyen et al.,
PRX 9, 2019*

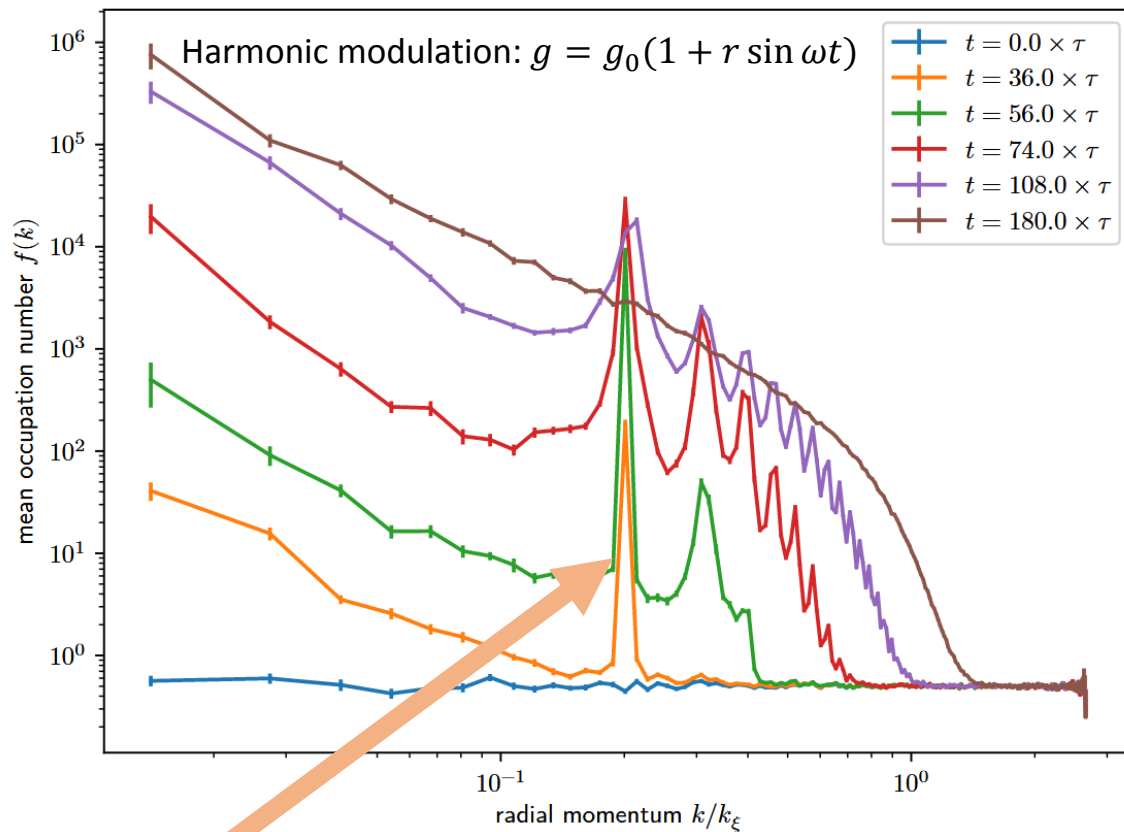
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Lattice simulations ($2D$, $N \sim 10^8$ atoms) based on TWA



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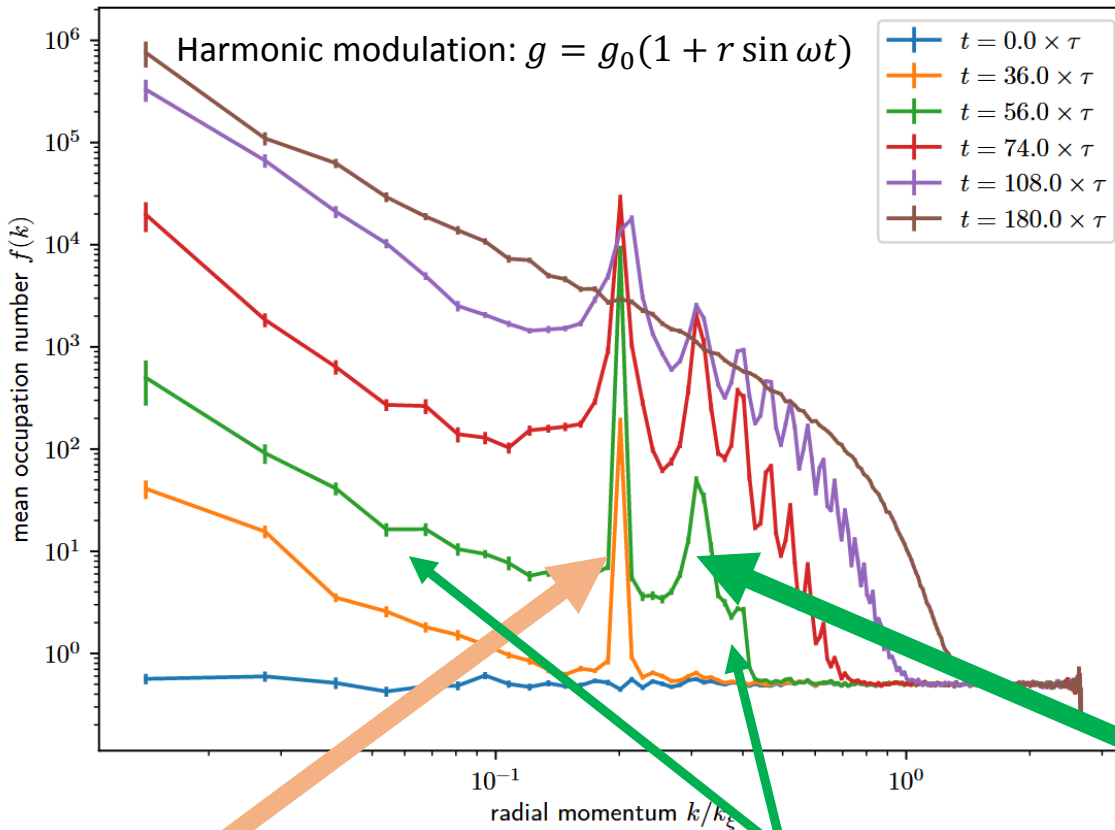


Linear regime

Parametric resonance
peak at $\omega_1 = \omega/2$.

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First non-linear corrections set in when $N_{exc}(t) \sim \sqrt{N}$.

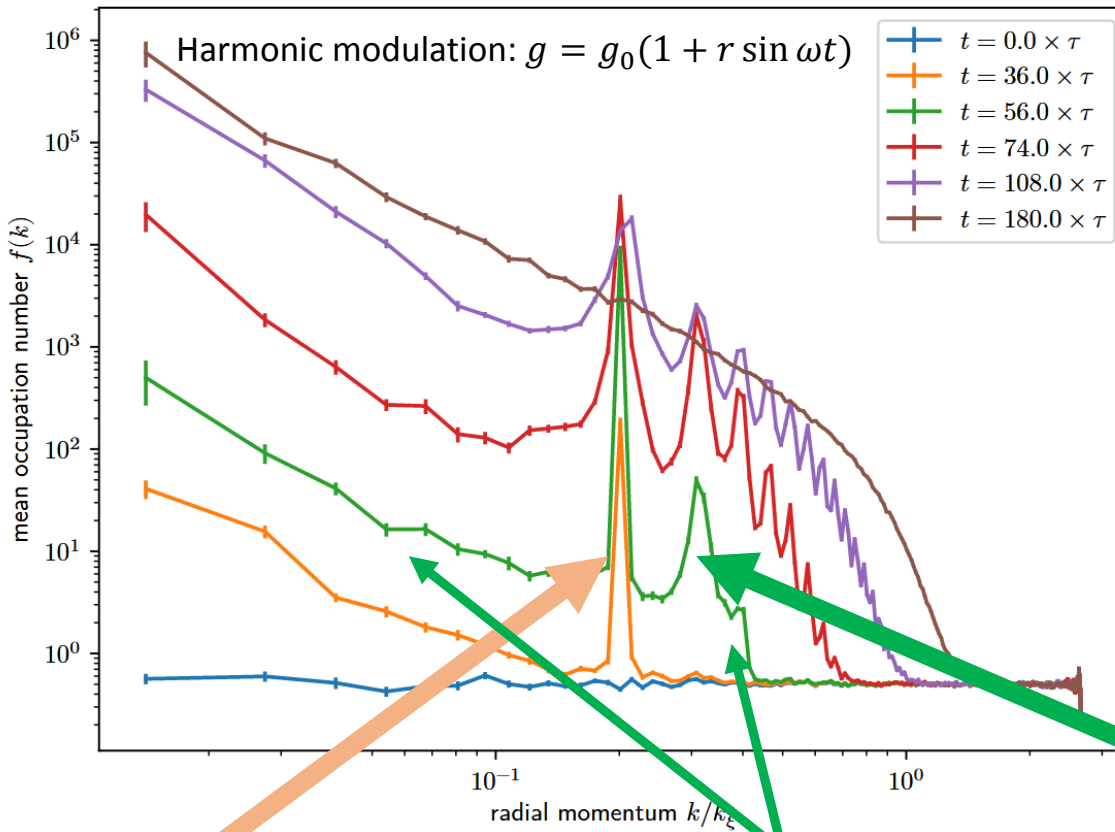
Parametric resonance peak at $\omega_1 = \omega/2$.

Broad band from off-shell re-scattering at $0 < p < 2p_1$.

Narrow peak from on-shell re-scattering at $\omega_2 = 2\omega_1$.

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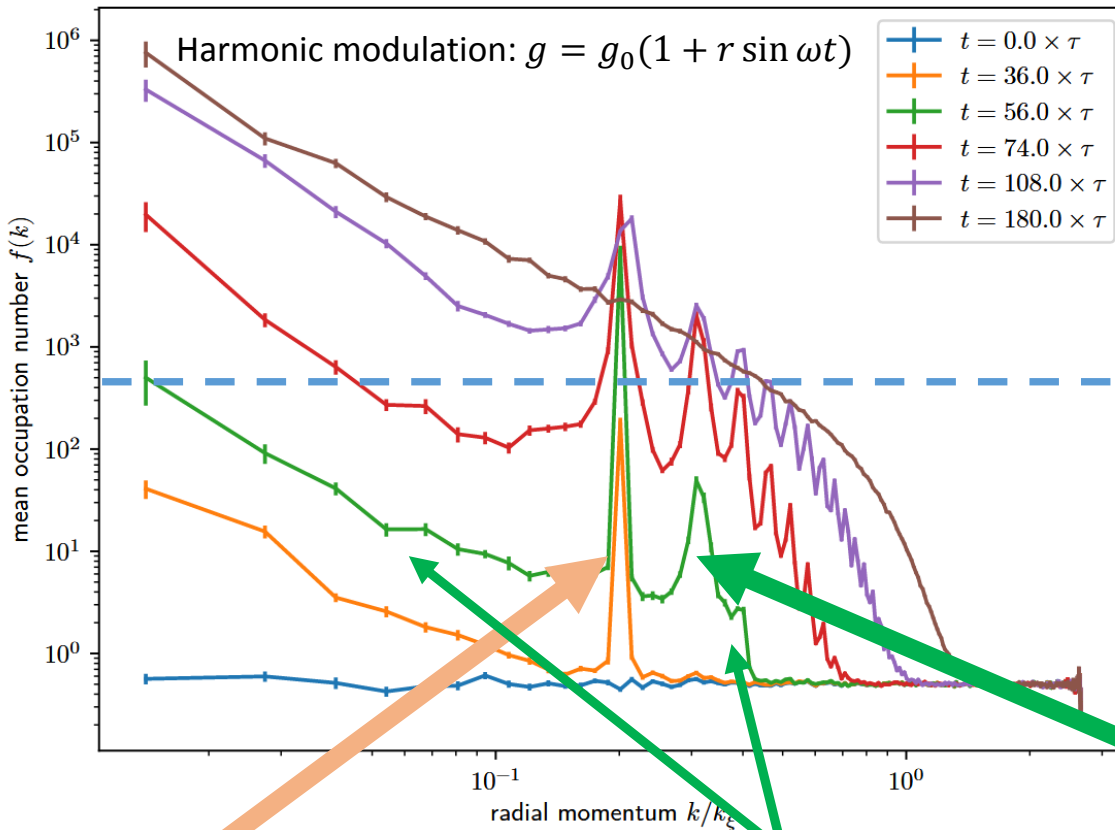
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Non-linearities: analytics

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- Statistical propagator:

$$\tilde{n}(x, y) = \frac{1}{2} \langle \{ \hat{\psi}(x) \hat{\psi}^+(y) \} \rangle - \Psi(x) \Psi^*(y), \quad \tilde{n}(t, t, \mathbf{p}) = f(t, \mathbf{p})$$



$$\Psi(x) = \langle \hat{\psi}(x) \rangle.$$

$$\tilde{m}(x, y) = \frac{1}{2} \langle \{ \hat{\psi}(x) \hat{\psi}(y) \} \rangle - \Psi(x) \Psi(y).$$

$$\left[i\partial_t - \frac{\mathbf{p}^2}{2m} - 2g(t)|\Psi(t)|^2 \right] \tilde{n}(t, t', \mathbf{p}) - g(t)\Psi^2(t)\tilde{m}^*(t, t', \mathbf{p}) = 0$$

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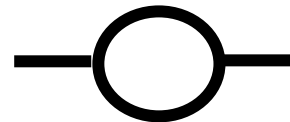
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Linear regime

$$+icg(t)g(t') \int_{\mathbf{q}} \left[4\Psi(t)\Psi^*(t') \left(\tilde{n}(t, t', \mathbf{q})\tilde{n}^*(t, t', \mathbf{p} - \mathbf{q}) + \text{similar terms} \right) \right]$$

+...

First non-linear correction



Zache et al., PRA 95, 2017
Berges et al., PRL 91 2003

Thermalization

Once a sufficient amount of atoms is produced, switch off the modulation

How does the universe approach thermal equilibrium?

*Micha and Tkachev
PRD 70 2004*

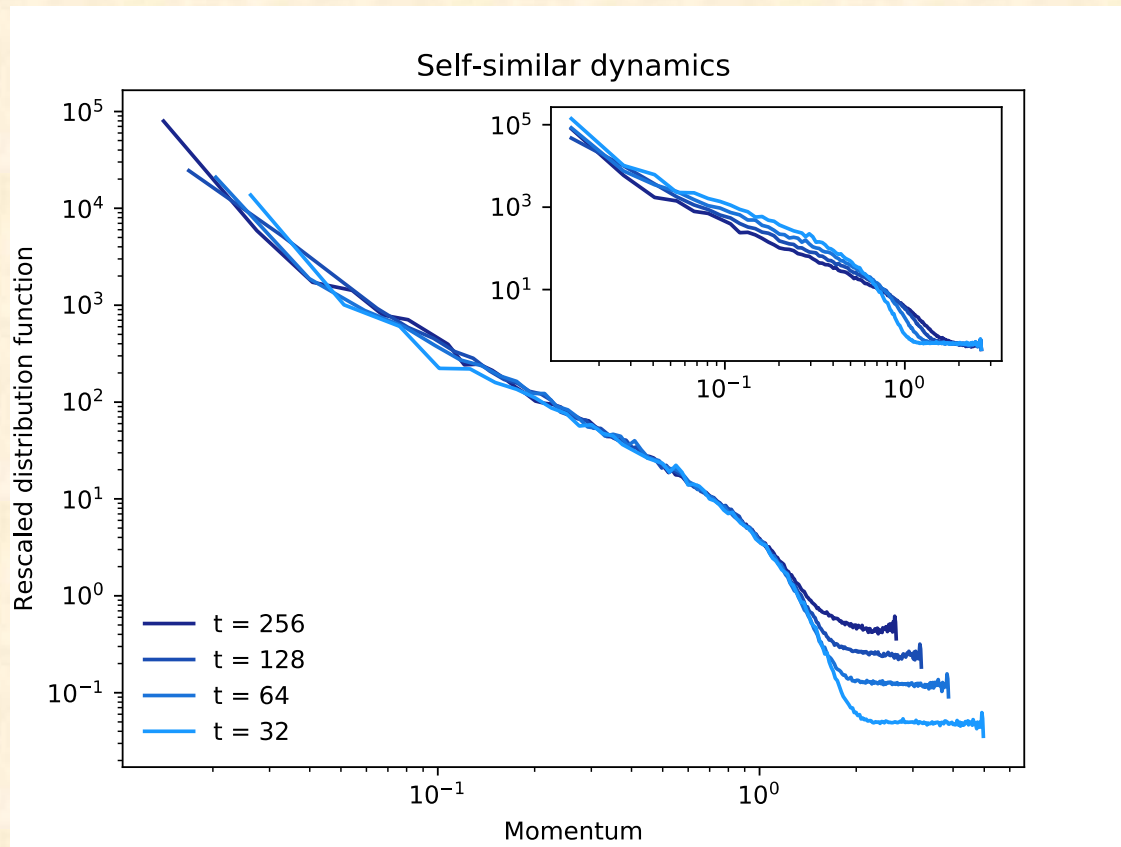
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Energy transport with
self-similar scaling
behavior:

$$f(t, p) = s^\alpha f_S(s^\beta p)$$

where $s = t/t_S$



Non-thermal fixed point

Numerically extracted
exponents:

$$\alpha \approx 1, \beta \approx 0.3$$

Thermalization vs expansion

- When $f_{\text{char}} \sim 1$, TWA becomes inapplicable



Relaxation to a Bose-Einstein distribution

Thermalization vs expansion

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Relaxation to a Bose-Einstein distribution

- Interaction rates decrease with expansion
 - Can the particles achieve thermal equilibrium?

*Kolb Turner,
Front.Phys. 69 (1990)*

$\Gamma \gg H$: thermalization possible

$\Gamma \ll H$: freeze-out

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$\Gamma \ll H$: freeze-out

- These effects are present in the Bose gas

• $d = 2 \Rightarrow g = \text{const} \Rightarrow$

Freeze-out happens at $\tau(t_{\text{phys}} = \infty)$,
if $\tau_{\text{therm}} > \tau(t_{\text{phys}} = \infty)$

• $d = 3 \Rightarrow g \propto a^{-1} \Rightarrow$

Freeze-out can happen
before $\tau(t_{\text{phys}} = \infty)$

My collaborators

Paper: in preparation

Theory + Future implementation in “Oberthaler lab”

Aleksandr
Chatrchyan



Jürgen
Berges



Kevin
Geier



Philipp
Hauke

Markus
Oberthaler



Helmut
Strobel



Celia
Viermann



Maurus
Hans

Thank you for your attention!

