

Correlation patterns in a BEC Black Hole Analog with Massive Phonons

Trento Workshop 2019

Richard A. Dudley
dudlra13@wfu.edu

In collaboration with
Paul R. Anderson, Alessandro Fabbri and Roberto Balbinot

September 16, 2019

Introduction

Overview



- ▶ Most work done theoretically and experimentally for BEC analog black holes in 1D systems is with no excitations of the transverse modes for the phonons ([A. Coutant, A. Fabbri, R. Parentani, R. Balbinot, and P.R. Anderson, PRD 86 \(2012\) 064022](#) and [G. Jannes, P. Mai, T.G. Philbin, and G. Rousseaux, PRD 83 \(2011\) 104028](#) are examples of exceptions)

Introduction

Overview



- ▶ Most work done theoretically and experimentally for BEC analog black holes in 1D systems is with no excitations of the transverse modes for the phonons ([A. Coutant, A. Fabbri, R. Parentani, R. Balbinot, and P.R. Anderson, PRD 86 \(2012\) 064022](#) and [G. Jannes, P. Mai, T.G. Philbin, and G. Rousseaux, PRD 83 \(2011\) 104028](#) are examples of exceptions)
- ▶ Transverse excitations generate an effective mass term in the phonon mode equation

Introduction

Overview



- ▶ Most work done theoretically and experimentally for BEC analog black holes in 1D systems is with no excitations of the transverse modes for the phonons ([A. Coutant, A. Fabbri, R. Parentani, R. Balbinot, and P.R. Anderson, PRD 86 \(2012\) 064022](#) and [G. Jannes, P. Mai, T.G. Philbin, and G. Rousseaux, PRD 83 \(2011\) 104028](#) are examples of exceptions)
- ▶ Transverse excitations generate an effective mass term in the phonon mode equation
- ▶ In the analog spacetime this results in a massive minimally coupled scalar field with a potential

Introduction

Overview



- ▶ Previous studies of massive fields in black hole spacetimes show the energy density and pressure are smaller than the massless case

Introduction

Overview



- ▶ Previous studies of massive fields in black hole spacetimes show the energy density and pressure are smaller than the massless case
- ▶ One might expect therefore that structure in the density density correlation function might be similar to the massless case but on a smaller scale

Introduction

Overview



- ▶ Previous studies of massive fields in black hole spacetimes show the energy density and pressure are smaller than the massless case
- ▶ One might expect therefore that structure in the density density correlation function might be similar to the massless case but on a smaller scale
- ▶ Instead we find **fundamental differences** between the massless and massive cases

Introduction

Overview



- ▶ Briefly discuss the inclusion of a mass term in the equations

Introduction

Overview



3

- ▶ Briefly discuss the inclusion of a mass term in the equations

- ▶ Introduce the simple model we've been working with

Introduction

Overview



3

- ▶ Briefly discuss the inclusion of a mass term in the equations
- ▶ Introduce the simple model we've been working with
- ▶ Review results for the two point and density density correlation functions for the case of massless phonons
- ▶ Show the very different results for the massive case

The Model

The Gross-Pitaevskii Equation



- $i\hbar\partial_T \Psi_0 = \left(-\frac{\hbar^2}{2m} \vec{\nabla}^2 + V_{ext} + g |\Psi_0|^2 \right) \Psi_0.$

- Let $\Psi_0 = \sqrt{n} e^{i\theta}$

The Model

The Gross-Pitaevskii Equation



- ▶ $i\hbar\partial_T \Psi_0 = \left(-\frac{\hbar^2}{2m} \vec{\nabla}^2 + V_{ext} + g |\Psi_0|^2 \right) \Psi_0.$
- ▶ Let $\Psi_0 = \sqrt{n} e^{i\theta}$
- ▶ Choose a solution with $n = \text{constant}$.

The Model

The Gross-Pitaevskii Equation



- ▶ $i\hbar\partial_T \Psi_0 = \left(-\frac{\hbar^2}{2m} \vec{\nabla}^2 + V_{ext} + g |\Psi_0|^2 \right) \Psi_0.$
- ▶ Let $\Psi_0 = \sqrt{n} e^{i\theta}$
- ▶ Choose a solution with $n = \text{constant}$.
- ▶ The BEC moves with a constant velocity in the negative x direction.

$$\vec{v} = \frac{\hbar}{m} \vec{\nabla} \theta = -v_0 \hat{x}$$

The Model

The Horizon



► $i\hbar\partial_T \Psi_0 = \left(-\frac{\hbar^2}{2m} \vec{\nabla}^2 + V_{ext} + g |\Psi_0|^2 \right) \Psi_0, \quad \Psi_0 = \sqrt{n} e^{i\theta}$

The Model

The Horizon

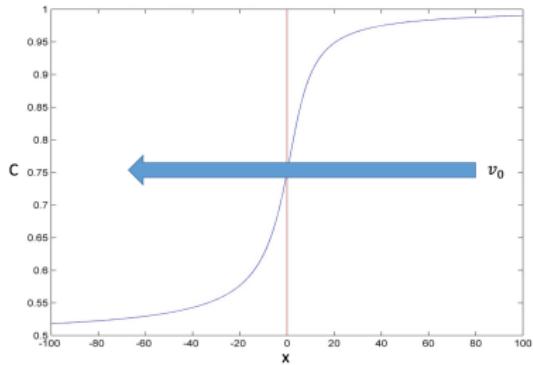


- ▶ $i\hbar\partial_T\Psi_0 = \left(-\frac{\hbar^2}{2m}\vec{\nabla}^2 + V_{ext} + g|\Psi_0|^2\right)\Psi_0, \quad \Psi_0 = \sqrt{n}e^{i\theta}$
- ▶ The sound speed is given by $c = \sqrt{\frac{gn}{m}}$

The Model

The Horizon

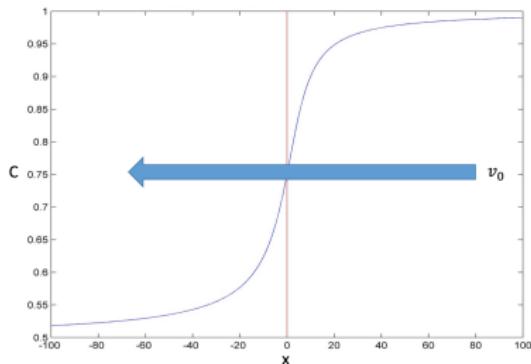
- ▶ $i\hbar\partial_T \Psi_0 = \left(-\frac{\hbar^2}{2m} \vec{\nabla}^2 + V_{ext} + g |\Psi_0|^2 \right) \Psi_0, \quad \Psi_0 = \sqrt{n} e^{i\theta}$
- ▶ The sound speed is given by $c = \sqrt{\frac{gn}{m}}$
- ▶ Let the sound speed vary along the flow.



The Model

The Horizon

- ▶ $i\hbar\partial_T\Psi_0 = \left(-\frac{\hbar^2}{2m}\vec{\nabla}^2 + V_{ext} + g|\Psi_0|^2\right)\Psi_0, \quad \Psi_0 = \sqrt{n}e^{i\theta}$
- ▶ The sound speed is given by $c = \sqrt{\frac{gn}{m}}$
- ▶ Let the sound speed vary along the flow.



- ▶ A sonic horizon occurs at $x = 0$
For $x > 0, \quad v_0 > c$
For $x < 0, \quad v_0 < c$

The Model

Perturbations



► $i\hbar\partial_T\Psi_0 = \left(-\frac{\hbar^2}{2m}\vec{\nabla}^2 + V_{ext} + g|\Psi_0|^2\right)\Psi_0, \quad \Psi_0 = \sqrt{n}e^{i\theta}$

The Model

Perturbations



- ▶ $i\hbar\partial_T\Psi_0 = \left(-\frac{\hbar^2}{2m}\vec{\nabla}^2 + V_{ext} + g|\Psi_0|^2\right)\Psi_0, \quad \Psi_0 = \sqrt{n}e^{i\theta}$
- ▶ Looking at perturbations of the form $\hat{n} = n + \hat{n}_1, \hat{\theta} = \theta + \hat{\theta}_1$

The Model

Perturbations



- ▶ $i\hbar\partial_T\Psi_0 = \left(-\frac{\hbar^2}{2m}\vec{\nabla}^2 + V_{ext} + g|\Psi_0|^2\right)\Psi_0, \quad \Psi_0 = \sqrt{n}e^{i\theta}$
- ▶ Looking at perturbations of the form $\hat{n} = n + \hat{n}_1, \hat{\theta} = \theta + \hat{\theta}_1$
- ▶ Considering scales much larger than the healing length(hydrodynamic approximation)

$$\left[-\frac{c}{c^2 - v_0^2}\partial_T^2 + c\partial_x\left(\frac{c^2 - v_0^2}{c^2}\partial_x\right) + c(\partial_y^2 + \partial_z^2)\right]\hat{\theta}_1 = 0 \quad (1)$$

The Model

Perturbations



- ▶ $i\hbar\partial_T\Psi_0 = \left(-\frac{\hbar^2}{2m}\vec{\nabla}^2 + V_{ext} + g|\Psi_0|^2\right)\Psi_0, \quad \Psi_0 = \sqrt{n}e^{i\theta}$
- ▶ Looking at perturbations of the form $\hat{n} = n + \hat{n}_1, \hat{\theta} = \theta + \hat{\theta}_1$
- ▶ Considering scales much larger than the healing length(hydrodynamic approximation)

$$\left[-\frac{c}{c^2 - v_0^2}\partial_T^2 + c\partial_x\left(\frac{c^2 - v_0^2}{c^2}\partial_x\right) + c(\partial_y^2 + \partial_z^2)\right]\hat{\theta}_1 = 0 \quad (1)$$

- ▶ Excitation of a mode in the transverse direction yields an effective mass term.

$$e^{i(k_y y + k_z z)} \rightarrow m^2 = k_y^2 + k_z^2$$

(F. Chevy, V. Bretin, P. Rosenbusch, K.W. Madison, and J. Dalibard, PRL 88 (2002) 250402)

The Model

Analog Spacetime



- ▶ It is possible to show this is equivalent to a massive minimally coupled scalar field propagating in the 2D spacetime

$$ds^2 = \frac{c^2 - v_0^2}{c} dt^2 + \frac{c}{c^2 - v_0^2} dx^2 = -(c^2 - v_0^2)(dt^2 - dx^{*2})$$

The Model

Analog Spacetime



- ▶ It is possible to show this is equivalent to a massive minimally coupled scalar field propagating in the 2D spacetime

$$ds^2 = \frac{c^2 - v_0^2}{c} dt^2 + \frac{c}{c^2 - v_0^2} dx^2 = -(c^2 - v_0^2)(dt^2 - dx^{*2})$$

- ▶ The wave equation is $[-\partial_t^2 + \partial_{x_*}^2 - m_{\text{eff}}^2 + V_{\text{eff}}] \hat{\theta}_2 = 0$

The Model

Analog Spacetime



- ▶ It is possible to show this is equivalent to a massive minimally coupled scalar field propagating in the 2D spacetime

$$ds^2 = \frac{c^2 - v_0^2}{c} dt^2 + \frac{c}{c^2 - v_0^2} dx^2 = -(c^2 - v_0^2)(dt^2 - dx^{*2})$$

- ▶ The wave equation is $[-\partial_t^2 + \partial_{x_*}^2 - m_{\text{eff}}^2 + V_{\text{eff}}] \hat{\theta}_2 = 0$
- ▶

$$m_{\text{eff}} \equiv m^2(c^2 - v_0^2)$$

$$V_{\text{eff}} \equiv \frac{c^2 - v_0^2}{c} \left[\frac{1}{2} \frac{d^2 c}{dx^2} \left(1 + \frac{v_0^2}{c^2} \right) - \frac{1}{4c} \left(\frac{dc}{dx} \right)^2 + \frac{5v_0^2}{4c^3} \left(\frac{dc}{dx} \right)^2 \right].$$

The Model



► $[-\partial_t^2 + \partial_{x^*}^2 - m_{\text{eff}} + V_{\text{eff}}] \hat{\theta}_2 = 0$

$$m_{\text{eff}} \equiv m^2(c^2 - v_0^2)$$

$$V_{\text{eff}} \equiv \frac{c^2 - v_0^2}{c} \left[\frac{1}{2} \frac{d^2 c}{dx^2} \left(1 + \frac{v_0^2}{c^2} \right) - \frac{1}{4c} \left(\frac{dc}{dx} \right)^2 + \frac{5v_0^2}{4c^3} \left(\frac{dc}{dx} \right)^2 \right].$$

where

The Model



► $[-\partial_t^2 + \partial_{x^*}^2 - m_{\text{eff}} + V_{\text{eff}}] \hat{\theta}_2 = 0$

$$m_{\text{eff}} \equiv m^2(c^2 - v_0^2)$$

$$V_{\text{eff}} \equiv \frac{c^2 - v_0^2}{c} \left[\frac{1}{2} \frac{d^2 c}{dx^2} \left(1 + \frac{v_0^2}{c^2} \right) - \frac{1}{4c} \left(\frac{dc}{dx} \right)^2 + \frac{5v_0^2}{4c^3} \left(\frac{dc}{dx} \right)^2 \right].$$

where

► Note that m_{eff} vanishes at the horizon.

The Model



► $[-\partial_t^2 + \partial_{x^*}^2 - m_{\text{eff}} + V_{\text{eff}}] \hat{\theta}_2 = 0$

$$m_{\text{eff}} \equiv m^2(c^2 - v_0^2)$$

$$V_{\text{eff}} \equiv \frac{c^2 - v_0^2}{c} \left[\frac{1}{2} \frac{d^2 c}{dx^2} \left(1 + \frac{v_0^2}{c^2} \right) - \frac{1}{4c} \left(\frac{dc}{dx} \right)^2 + \frac{5v_0^2}{4c^3} \left(\frac{dc}{dx} \right)^2 \right].$$

where

- Note that m_{eff} vanishes at the horizon.
- V_{eff} vanishes at the horizon and at $x = \pm\infty$ if $c \rightarrow \text{constant}$.

Our Model



- ▶ Set $V_{\text{eff}} = 0$

Our Model



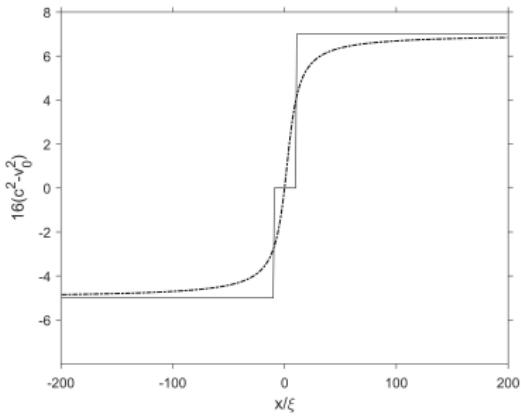
- ▶ Set $V_{\text{eff}} = 0$
- ▶ Let

$$c^2 - v_0^2 \rightarrow c_R^2 - v_0^2 \Theta(x^* - R x_0^*), \quad x > 0$$
$$c^2 - v_0^2 \rightarrow c_L^2 - v_0^2 \Theta(x^* - L x_0^*), \quad x < 0$$

Our Model

- Set $V_{\text{eff}} = 0$
- Let

$$\begin{aligned}c^2 - v_0^2 &\rightarrow c_R^2 - v_0^2 \Theta(x^* - R x_0^*), & x > 0 \\c^2 - v_0^2 &\rightarrow c_L^2 - v_0^2 \Theta(x^* - L x_0^*), & x < 0\end{aligned}$$



Correlation functions



► $[-\partial_t^2 + \partial_{x_*}^2 - m_{\text{eff}}] \hat{\theta}_2 = 0$

Correlation functions



- ▶ $[-\partial_t^2 + \partial_{x_*}^2 - m_{\text{eff}}^2] \hat{\theta}_2 = 0$
- ▶ The two point function $\langle \{\hat{\theta}_2(T, x) \hat{\theta}_2(T' x')\} \rangle$

Correlation functions

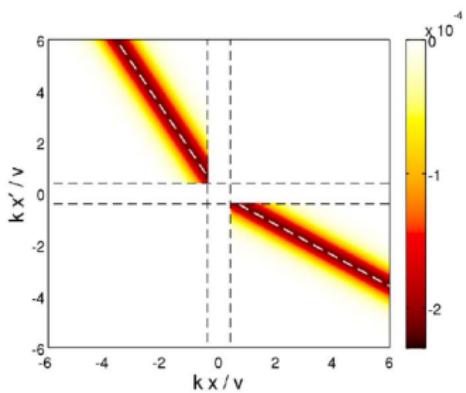
- ▶ $[-\partial_t^2 + \partial_{x_*}^2 - m_{\text{eff}}^2] \hat{\theta}_2 = 0$
- ▶ The two point function $\langle \{\hat{\theta}_2(T, x) \hat{\theta}_2(T' x')\} \rangle$
- ▶ The D-D correlation function can be written in terms of the two point function as follows:

$$G_2(T, x; T', x') = \frac{\hbar n}{2m l_\perp c(x)^2 c(x')^2} \lim_{T \rightarrow T'} D \sqrt{c(x)c(x')} \langle \{\hat{\theta}_2(T, x) \hat{\theta}_2(T' x')\} \rangle \quad (2)$$

$$D = \partial_T \partial_{T'} - v_0 \partial_T \partial_{x'} - v_0 \partial_x \partial_{T'} + v_0^2 \partial_x \partial_{x'}. \quad (3)$$

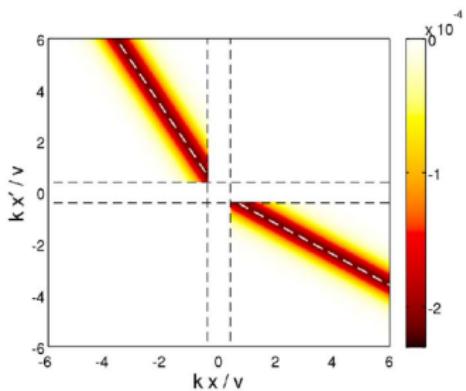
The Massless Case

- DD correlation function from R. Balbinot, A. Fabbri, S. Fagnocchi, A. Recati, and I. Carusotto, PRA **78** (2008) 021603



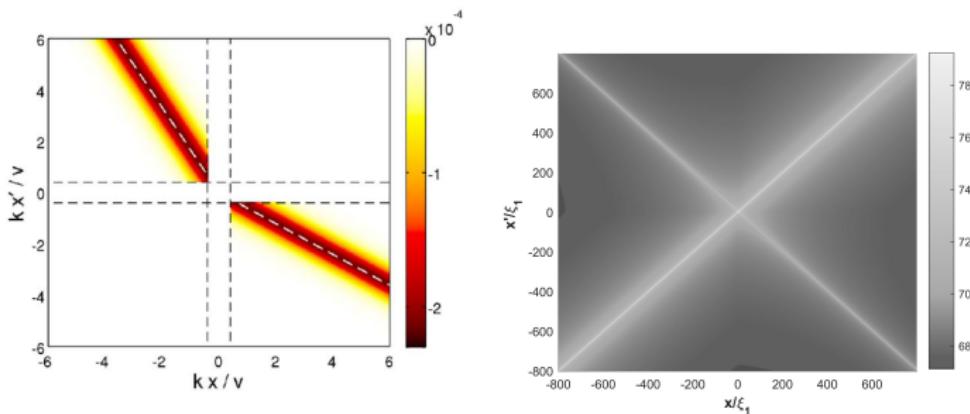
The Massless Case

- DD correlation function from R. Balbinot, A. Fabbri, S. Fagnocchi, A. Recati, and I. Carusotto, PRA **78** (2008) 021603
 - $V_{\text{eff}} = 0, m_{\text{eff}} = 0$



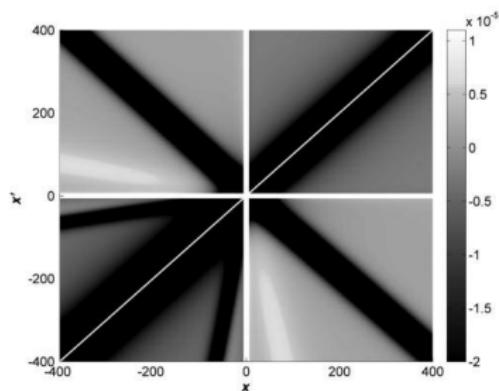
The Massless Case

- ▶ DD correlation function from R. Balbinot, A. Fabbri, S. Fagnocchi, A. Recati, and I. Carusotto, PRA **78** (2008) 021603
 - ▶ $V_{\text{eff}} = 0, m_{\text{eff}} = 0$
- ▶ Two -Point function for current speed of sound profile



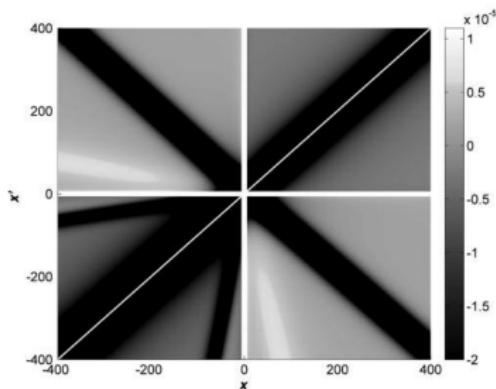
The Massless Case

- DD correlation function from P.R. Anderson, R. Balbinot, A. Fabbri, and R. Parentani [PRD 87 \(2013\) 124018](#)



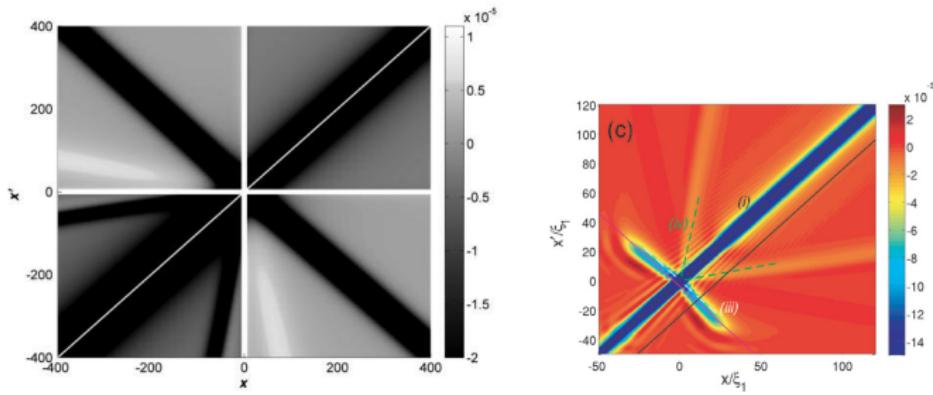
The Massless Case

- ▶ DD correlation function from P.R. Anderson, R. Balbinot, A. Fabbri, and R. Parentani [PRD 87 \(2013\) 124018](#)
 - ▶ $V_{\text{eff}} \neq 0, m_{\text{eff}} = 0$



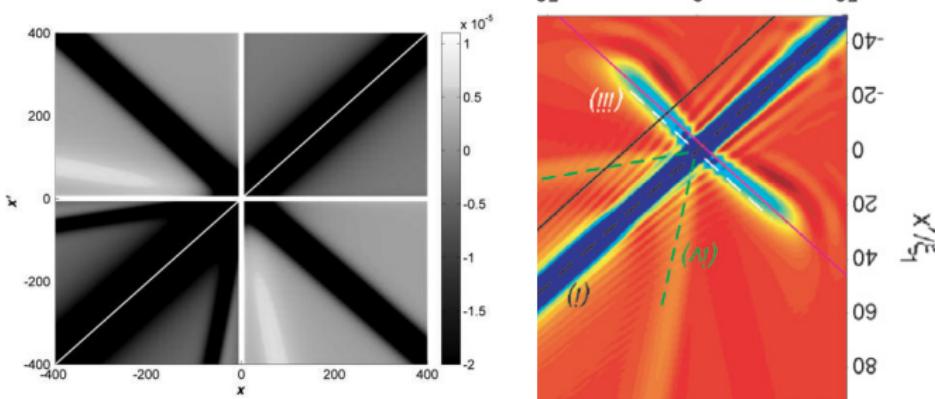
The Massless Case

- ▶ DD correlation function from P.R. Anderson, R. Balbinot, A. Fabbri, and R. Parentani [PRD 87 \(2013\) 124018](#)
 - ▶ $V_{\text{eff}} \neq 0, m_{\text{eff}} = 0$
- ▶ 2008 condensed matter calculation. I. Carusotto, S. Fagnocchi, A. Recati, R. Balbinot, and A. Fabbri, [NJP 10 \(2008\) 103001](#)



The Massless Case

- ▶ DD correlation function from P.R. Anderson, R. Balbinot, A. Fabbri, and R. Parentani [PRD 87 \(2013\) 124018](#)
 - ▶ $V_{\text{eff}} \neq 0, m_{\text{eff}} = 0$
- ▶ 2008 condensed matter calculation. I. Carusotto, S. Fagnocchi, A. Recati, R. Balbinot, and A. Fabbri, [NJP 10 \(2008\) 103001](#)



Results

Massive Two-Point Function



13

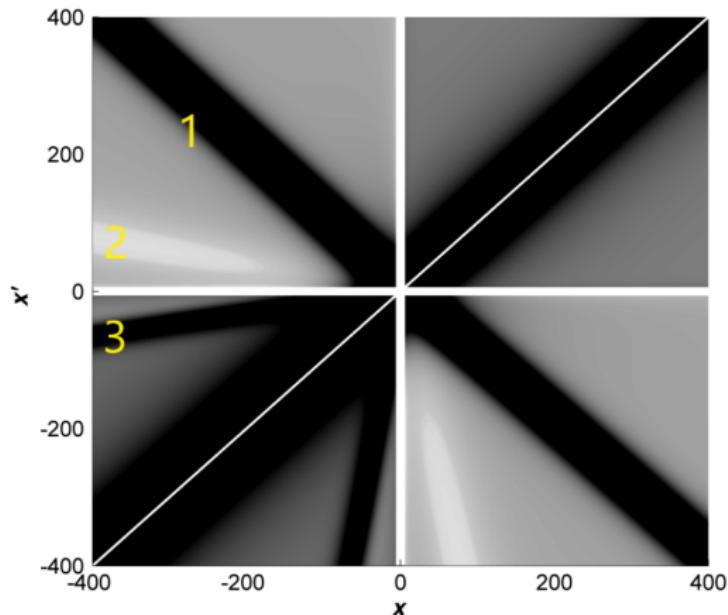


Figure: PRD 2013 Density Density results.

The Massless Case

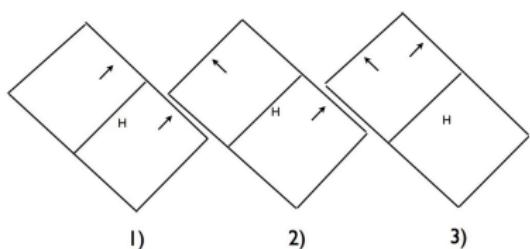
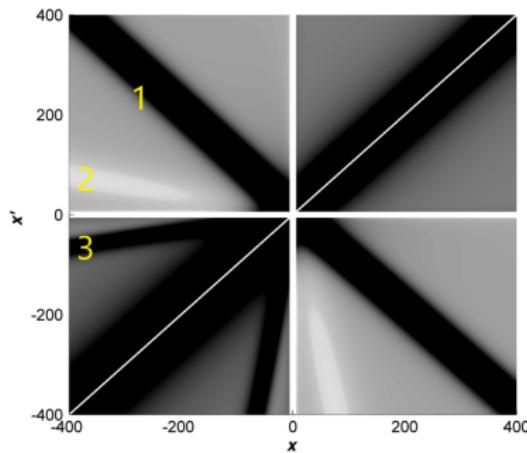


Figure: Origin of peaks in massless case.

Results

Massive Two-Point Function

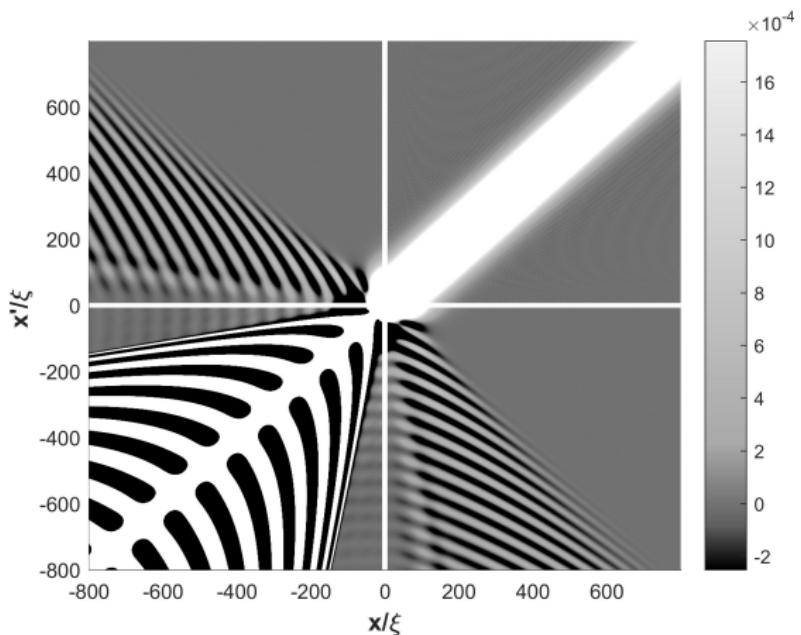


Figure: Massive two-point function correlation function for 1+1D BEC BH analog with masslike term, where $m = 4 \times 10^{-2} m_a$.

Results

Massive Two-Point Function with massless peaks superimposed

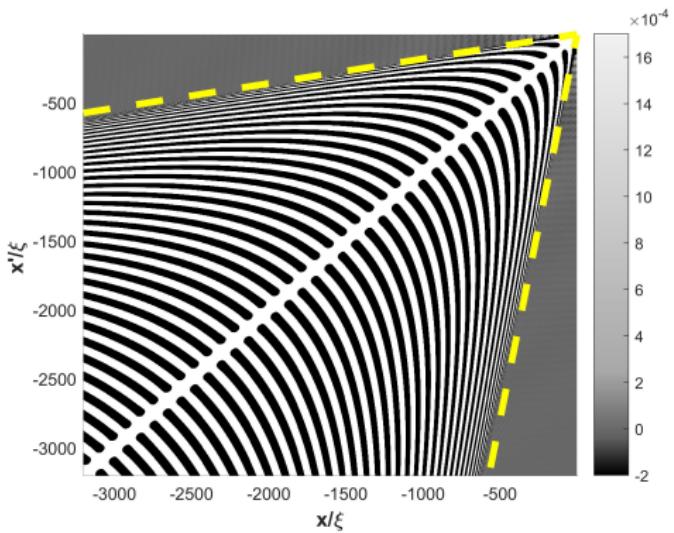


Figure: massive two-point function correlation function for 1+1D BEC BH analog with masslike term where $m = 4 \times 10^{-2} m_a$ with both points in the interior of the analog BH. The yellow dashed lines are the locations of the peaks found in the massless case.

Results

Massive Two-Point Function with massless peaks superimposed

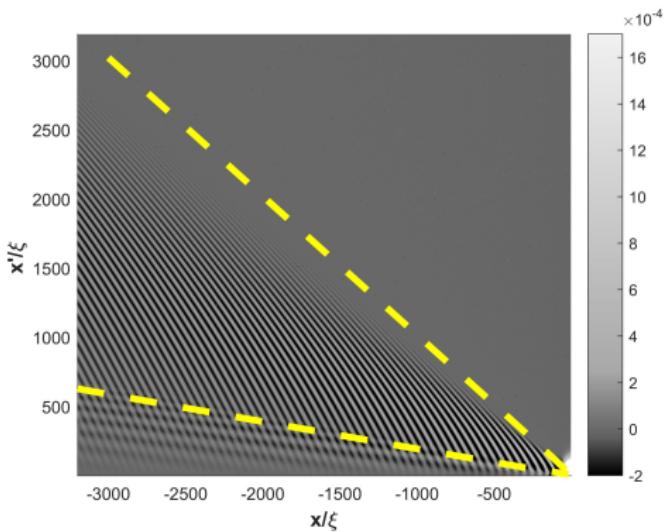


Figure: massive two-point function correlation function for 1+1D BEC BH analog with masslike term where $m = 4 \times 10^{-2} m_a$ with one point in the interior and 1 point in the exterior of the analog BH. The yellow dashed lines are the locations of the peaks found in the massless case.

Results

Massive Two-Point Function Decreasing Mass



18

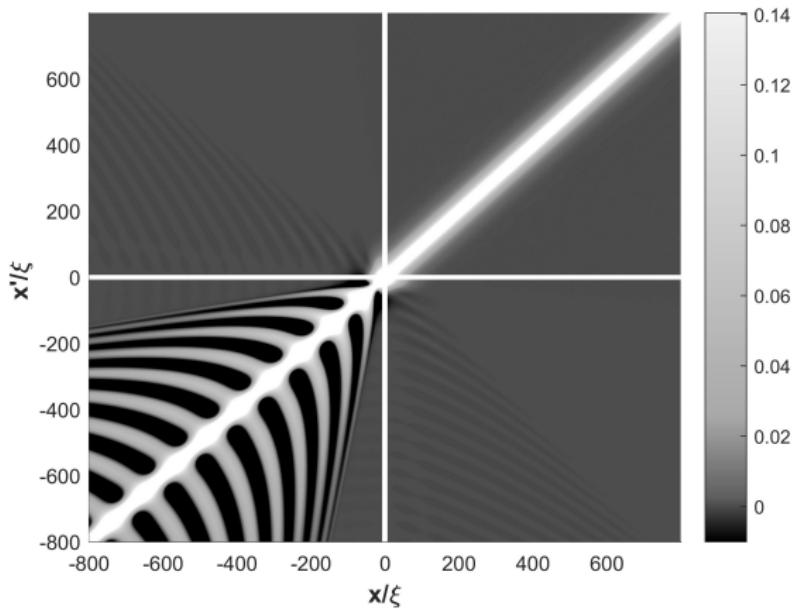


Figure: $m = 4 \times 10^{-2} m_a$

Results

Massive Two-Point Function Decreasing Mass

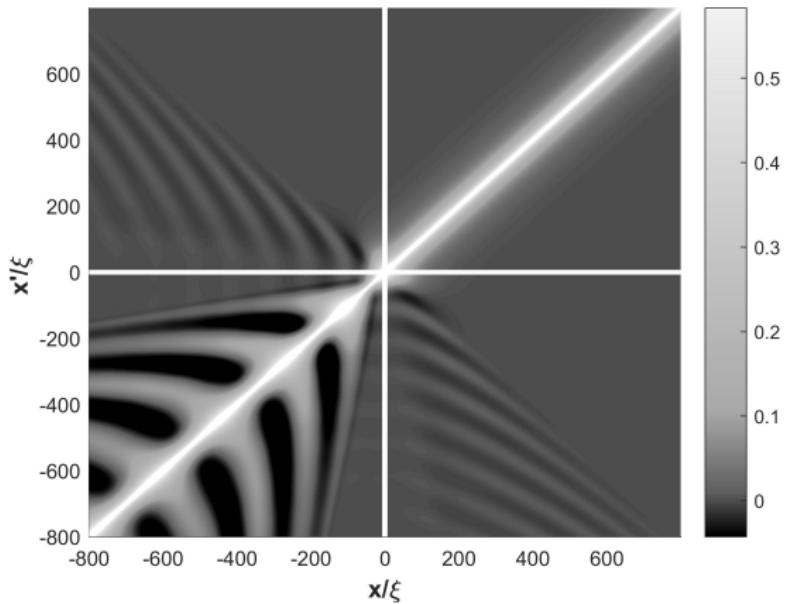


Figure: $m = 2 \times 10^{-2} m_a$

Results

Massive Two-Point Function Decreasing Mass

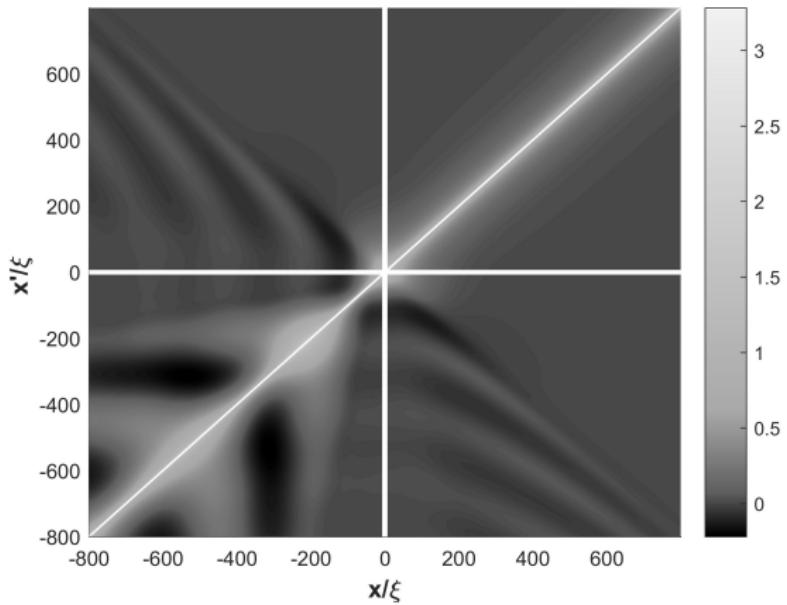


Figure: $m = 1 \times 10^{-2} m_a$

Results

Massive Two-Point Function Decreasing Mass



21

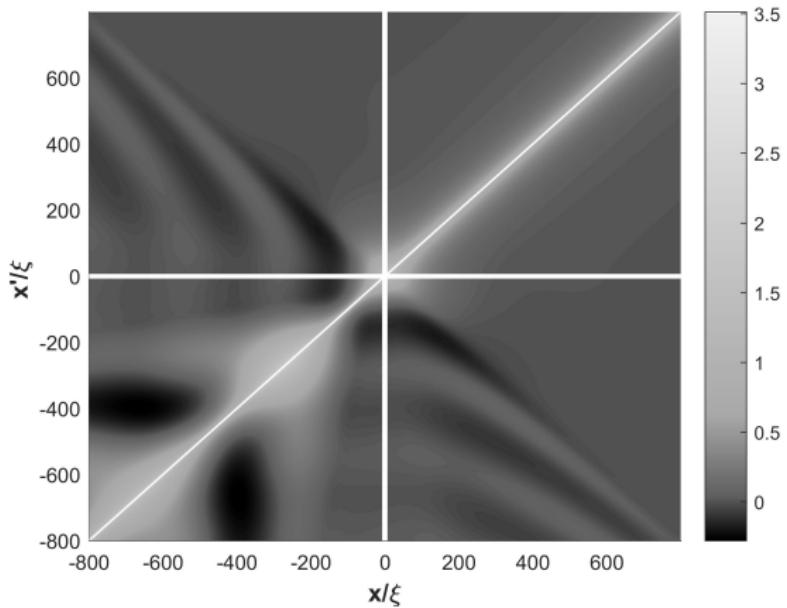


Figure: $m = 8 \times 10^{-3} m_a$

Results

Massive Two-Point Function Decreasing Mass



22

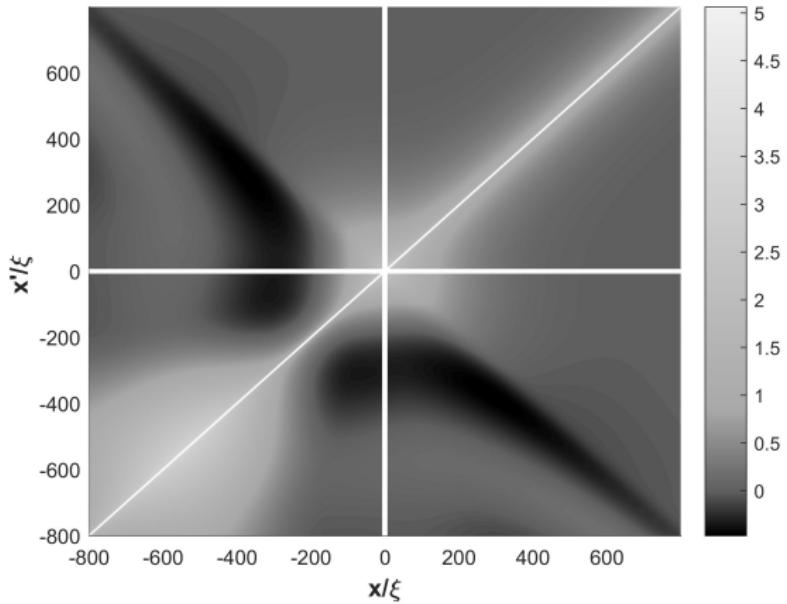


Figure: $m = 6 \times 10^{-3} m_a$

Results

Massive Two-Point Function Decreasing Mass

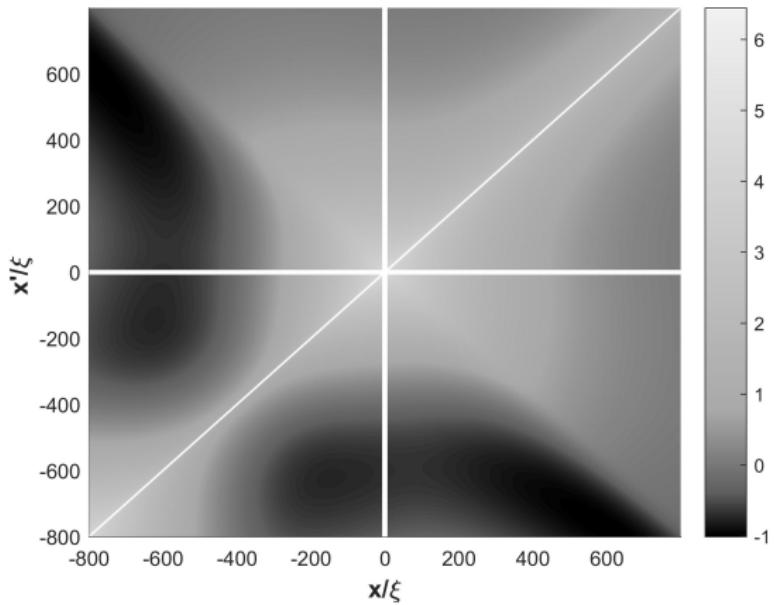


Figure: $m = 4 \times 10^{-3} m_a$

Results

Massive Two-Point Function Decreasing Mass



24

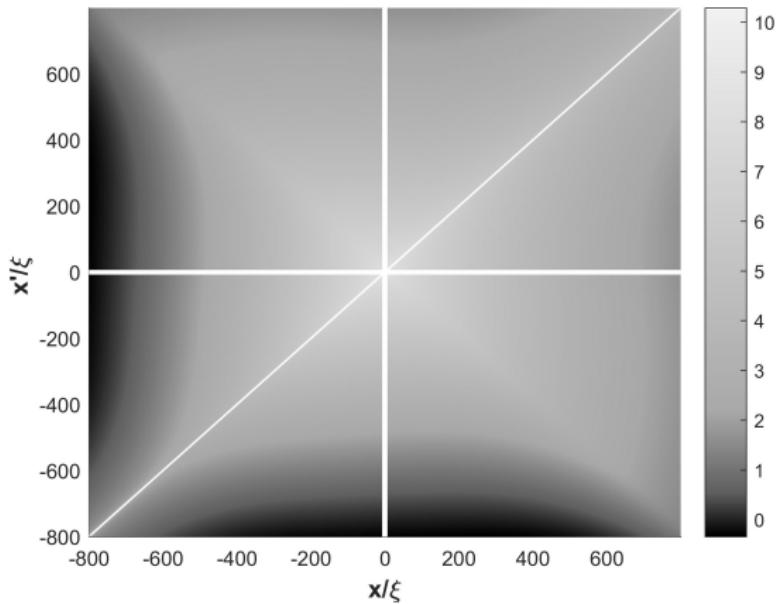


Figure: $m = 2 \times 10^{-3} m_a$

Results

Massive Two-Point Function Decreasing Mass



25

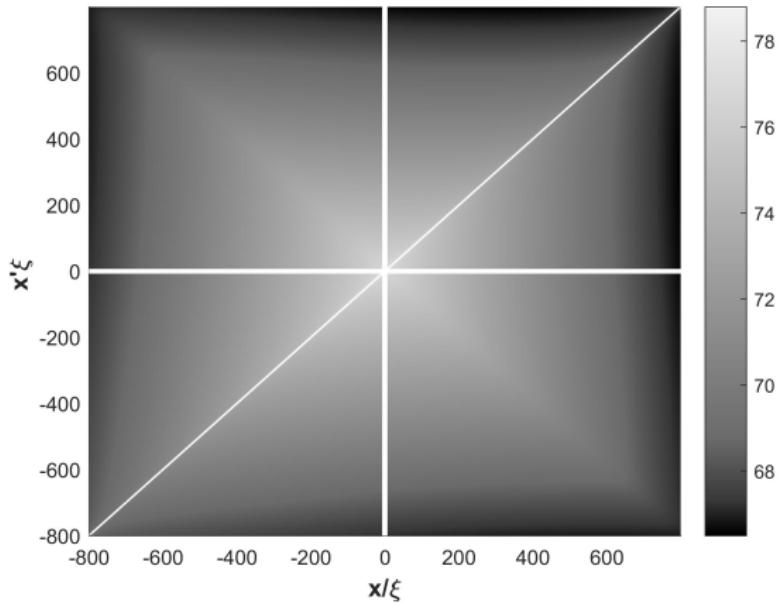


Figure: $m = 1 \times 10^{-4} m_a$

Results

Massive Two-Point Function Low Mass Comparison to Massless Result

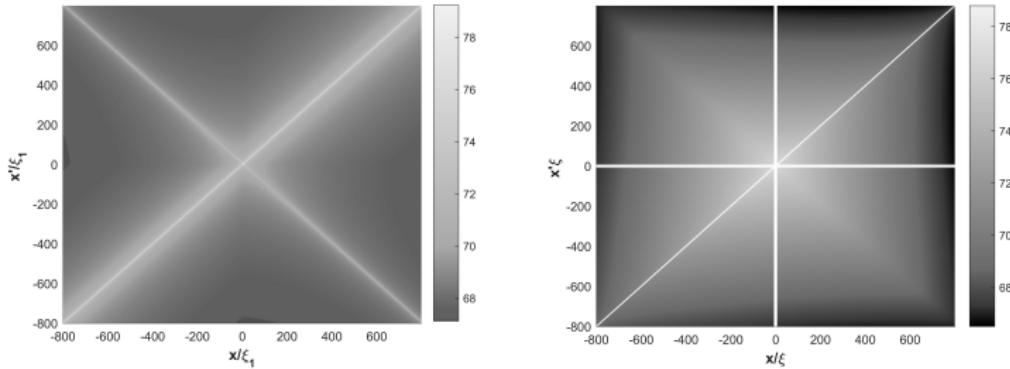


Figure: Left: Analytic solution for the massless two point function. Right: Massive two-point function correlation function for 1+1D BEC BH analog with masslike term where $m = \times 10^{-4} m_a$.

Results

Massive Two-Point Function Low Mass Result



27

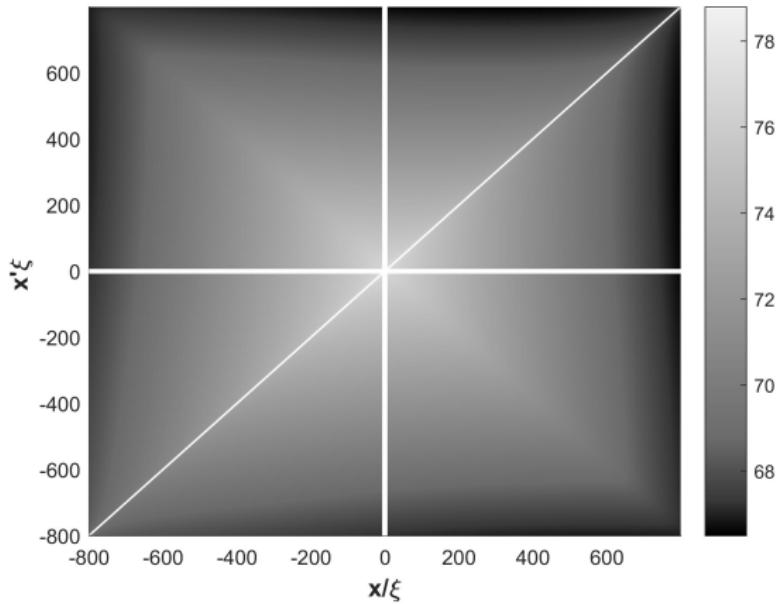


Figure: $m = 1 \times 10^{-4} m_a$

Results

Massive Two-Point Function Low Mass Result



28

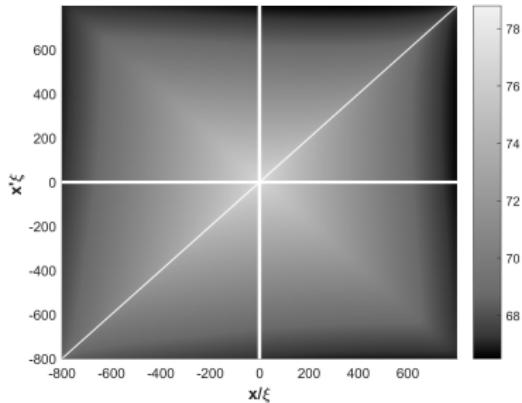


Figure: $m = 1 \times 10^{-4} m_a$

Results

Massive Two-Point Function Low Mass Result

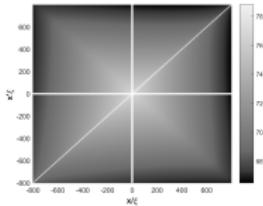


Figure: $m = 1 \times 10^{-4} m_a$

Results

Massive Two-Point Function Low Mass Result



30

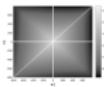
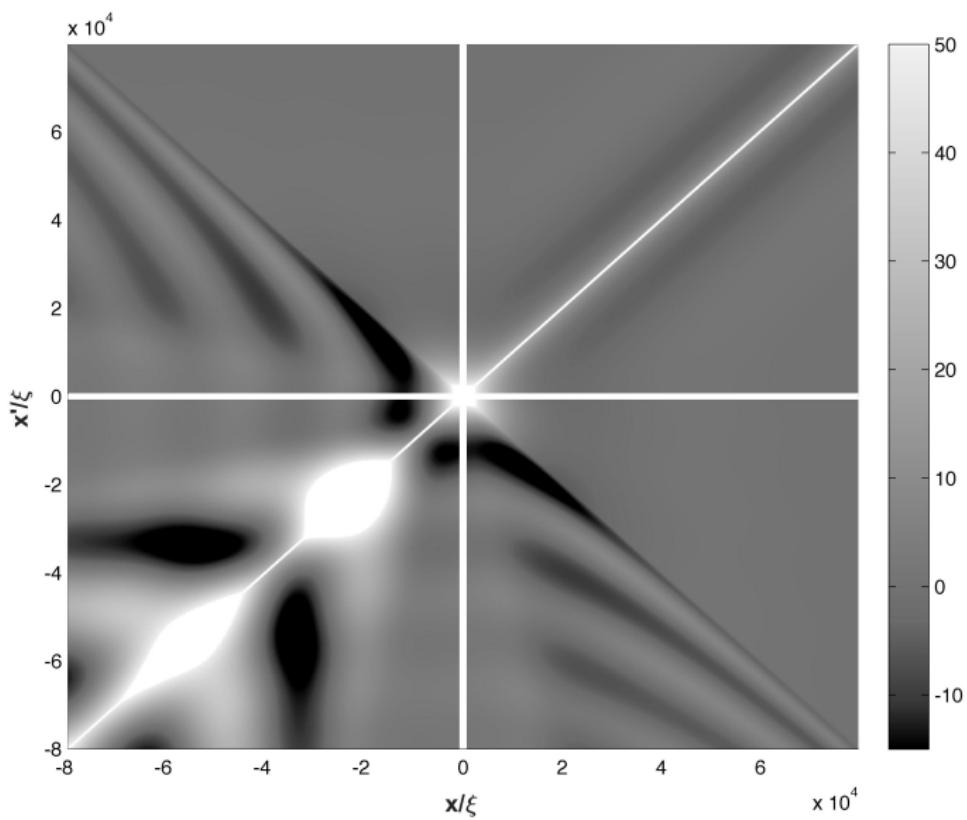


Figure: $m = 1 \times 10^{-4} m_a$

Results

Massive Two-Point Function Low Mass Result

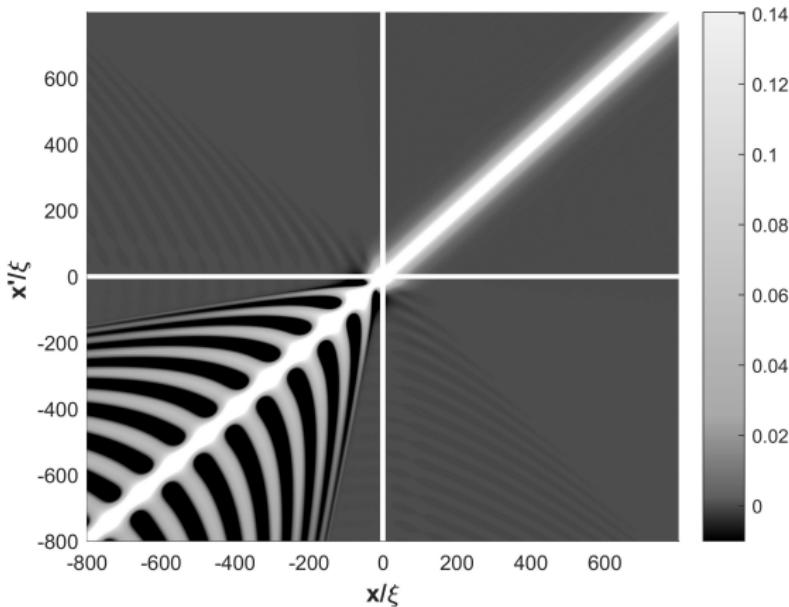


Results

Massive Two-Point Function Undulations

Undulations exist in the massive case. ([A. Coutant, A. Fabbri, R. Parentani, R. Balbinot, and P.R. Anderson, PRD 86 \(2012\) 064022](#))

$$\cos(p_U^m x) \cos(p_U^m x') \quad \text{where } p_U \propto m$$



Results

Massive Two-Point Function Undulations

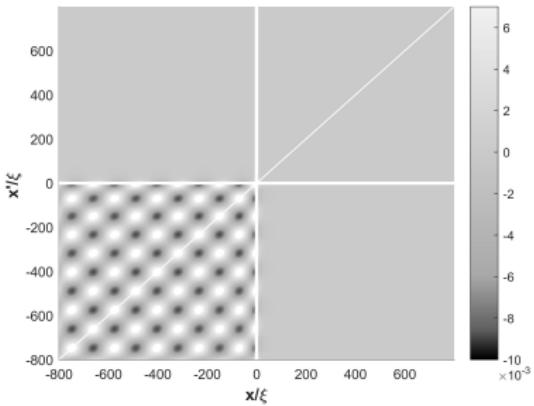


Figure: Massive two-point function correlation function for 1+1D BEC BH analog with masslike term, where $m = 4 \times 10^{-2} m_a$. The plot only includes the frequency range $0 < \omega < m/100$.

Results

Massive Two-Point Function Undulations

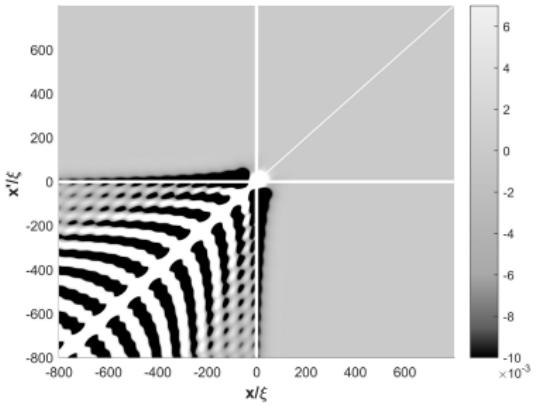


Figure: Massive two-point function correlation function for 1+1D BEC BH analog with masslike term, where $m = 4 \times 10^{-2} m_a$. The plot only includes the frequency range $0 < \omega < m/10$.

Results

Massive Two-Point Function Undulations



35

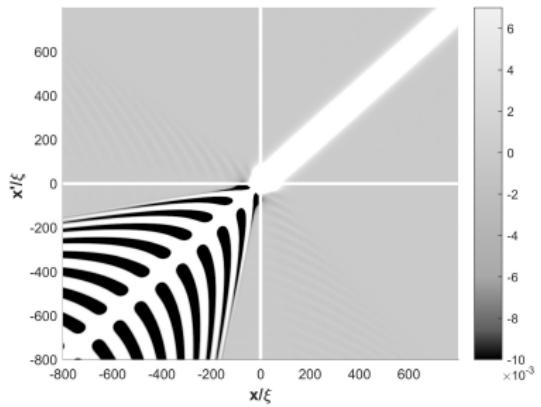


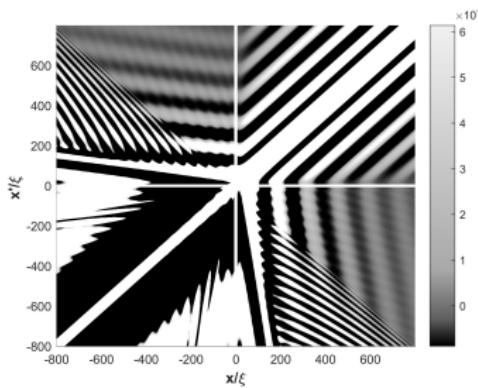
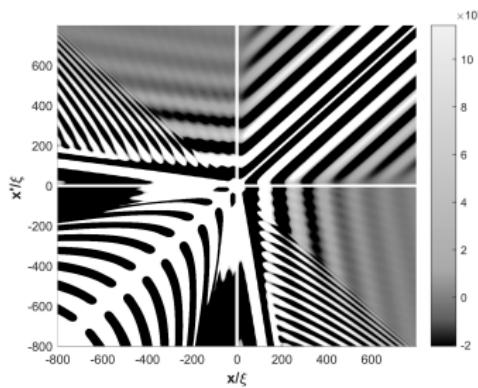
Figure: Massive two-point function correlation function for 1+1D BEC BH analog with masslike term, where $m = 4 \times 10^{-2} m_a$.

Thank you for your time!

Results

Massive Two-Point Function Cancellations

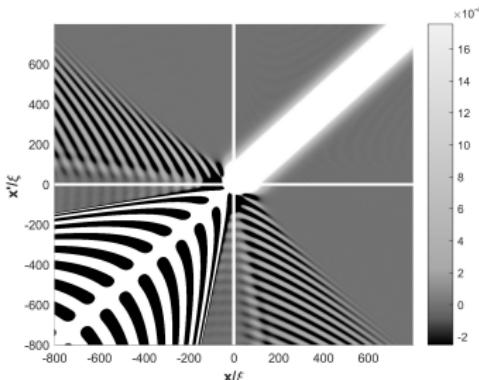
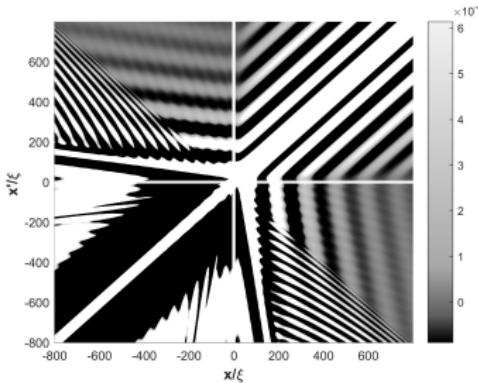
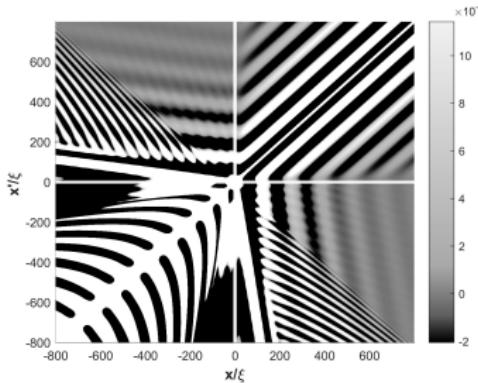
► Left: H_- contribution. Right: i_- contribution.



Results

Massive Two-Point Function Cancellations

► Left: H_- contribution. Right: i_- contribution.



Results

Massive Two-Point Function Stationary Phase Approximation



► Left: H_- contribution. Right: i_- contribution.

