

Neutrino interactions in nuclei

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Atomic nuclei as laboratories for BSM physics,
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National Energy Research
Scientific Computing Center



At "nuclear" energies, understanding neutrino-nucleus interactions very challenging and important!

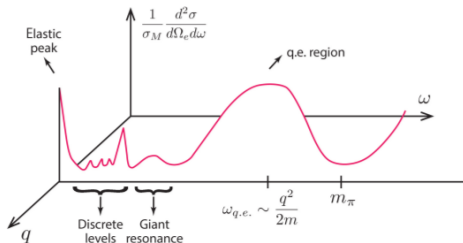
Understanding Nuclei:

- Nuclear interactions and structure
- Exotic nuclei - neutron rich
- Electroweak processes

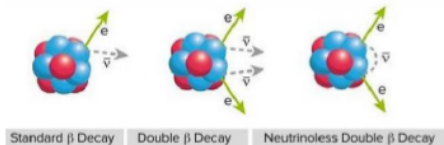
Relevance:

- Neutrino scattering in nuclei (neutrino oscillation experiments)
- Neutrinoless Double Beta Decay
- Neutrino interactions in supernovae and neutron stars, nucleosynthesis

We need a coherent picture of ν -nucleus interactions

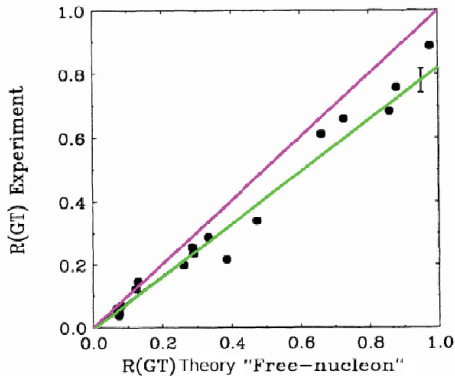
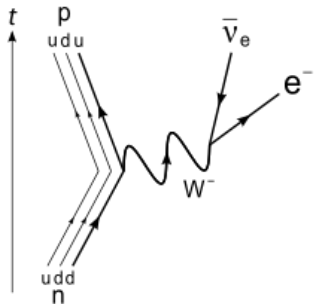


- $\omega \approx \text{few MeV}$, $q \approx 0$: β - and $\beta\beta$ -decays
- $\omega \approx \text{few MeV}$, $q \approx 10^2 \text{ MeV}$: Neutrinoless $\beta\beta$ -decays
- $\omega \leq \text{tens MeV}$: Astrophysics
- $\omega \approx 10^2 \text{ MeV}$: Accelerator neutrinos, ν -nucleus scattering



- The “ g_A quenching” puzzle
- Evidence of the importance of two-body currents
- The nuclear Hamiltonian, currents and QMC method
- β -decay in nuclei
- Electron scattering
- ν -nucleus scattering

The “quenching”- g_A problem



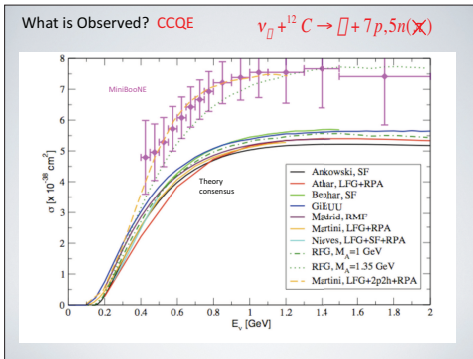
$$g_A^{\text{eff}} \simeq 0.70 g_A$$

Chou et al., PRC 47, 163 (1993)

What's the origin (or is there a **need**) of g_A quenching?

Charge-change quasi-elastic cross-section in ^{12}C

Experimental vs theory disagreement:



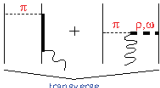
Alvarez-Ruso arXiv:1012.3871

Currents inconsistent with the Hamiltonian.

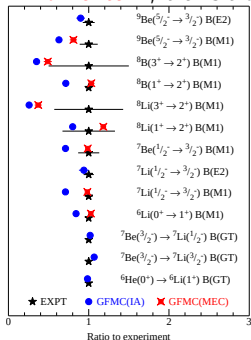
Nucleon-nucleon correlations and two-body processes approximately accounted for.

Need of g_A “unquenching” ???

Two-body processes

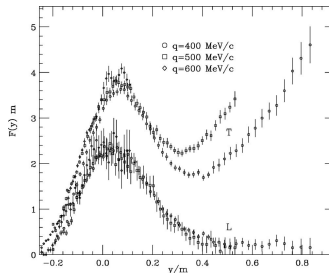
$$j = j^{(1)} + j^{(2)}(v) + j^{(3)}(v^{2\pi})$$


Low-momentum, transitions:



Pastore et al, PRC 2014

High-momentum, e^- scattering:
rescaled longitudinal vs transverse
electromagnetic response in ^{12}C

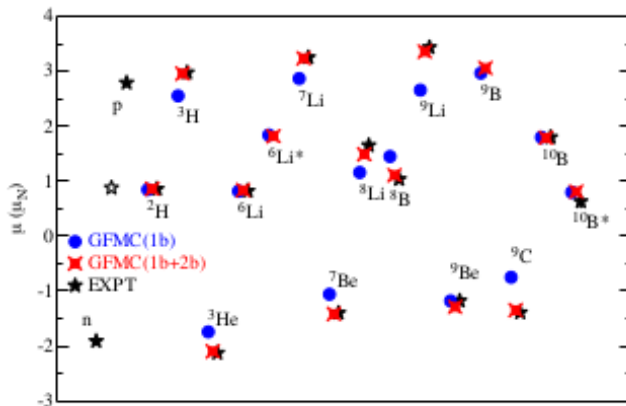


Benhar, Day, Sick, RMP (2008)

Without two-body processes, the
longitudinal and transverse response
is about the same

Electromagnetic two-body processes

Magnetic moments in light nuclei:



Pastore et al, PRC 2013

Nuclear Hamiltonian (only pions)

	2N force	3N force	4N force
LO		—	—
NLO		—	—
N ² LO			—
N ³ LO			

Expansion in powers of Q/Λ , $Q \sim 100$ MeV, $\Lambda \sim 1$ GeV.

Long-range physics given by pion-exchanges (no free parameters).

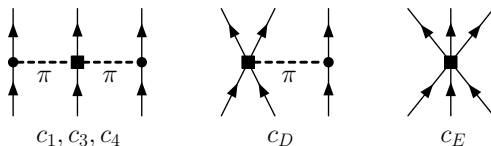
Short-range physics: contact interactions (LECs) to fit.

Operators need to be regulated $\rightarrow$ **cutoff dependency!**

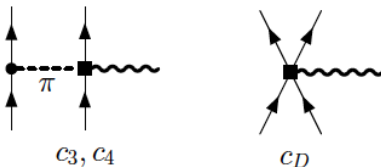
Order's expansion provides a way to quantify uncertainties!

Three-body forces and currents

Chiral three-body forces at N²LO:



Chiral two-body currents:



NN and NNN often use different regulators and cutoffs (local vs non-local, or a mix).

Hamiltonian-currents consistency?

Nuclear Hamiltonian

Model: non-relativistic nucleons strongly interacting with a nucleon-nucleon (NN) and three-nucleon interaction (TNI).

$$H = -\frac{\hbar^2}{2m} \sum_{i=1}^A \nabla_i^2 + \sum_{i<j} v_{ij} + \sum_{i<j<k} V_{ijk}$$

v_{ij} NN fitted on scattering data and TNI to properties of light nuclei.

Quantum Monte Carlo methods used to solve the many-body Schroedinger equation in imaginary time t :

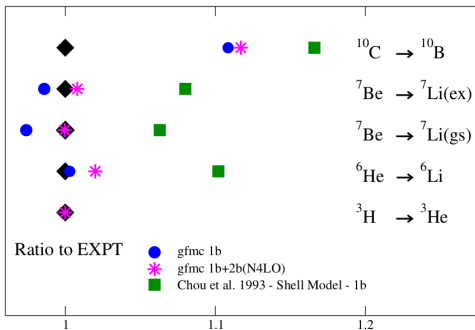
$$H \psi(\vec{r}_1 \dots \vec{r}_N) = E \psi(\vec{r}_1 \dots \vec{r}_N) \quad \psi(t) = e^{-Ht} \psi(0)$$

Ground-state extracted in the limit of $t \rightarrow \infty$.

No more details! Ask if interested...

β -decays in light nuclei

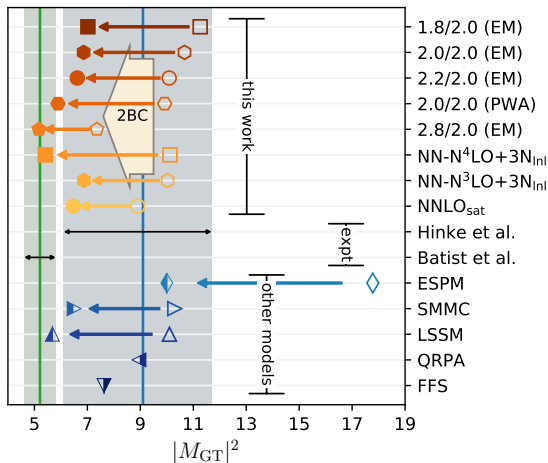
QMC calculations using a correlated wave function compared to shell-model calculations using the AV18+IL7 Hamiltonian and chiral currents.



Pastore, et al., PRC 97, 022501 (2018).

The effect of correlations in the nuclear wave function is critical!

β -decay in ^{100}Sn



ESPM: Extreme
Single Particle Model

SMMC: Shell Model
MC

LSSM: Large Space
Shell Model

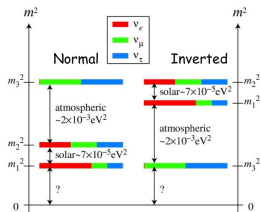
QRPA: quasiparticle
random phase
approximation

FFS: finite Fermi
systems

Gysbers et al., Nature Physics (2019).

Motivation

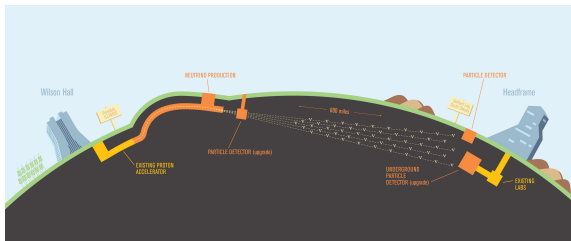
DUNE - Deep Underground Neutrino Experiment - to measure neutrino oscillations and CP violation



Simplified 2 flavors evolution (CP violation non included):

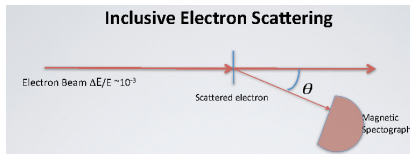
$$P_{\alpha \rightarrow \beta} = \sin^2(2\theta_{\alpha\beta}) \sin^2 \left(1.267 \frac{\Delta m_{\alpha\beta}^2 L}{E} \frac{\text{GeV}}{\text{eV}^2 \text{km}} \right)$$

Need to know E !

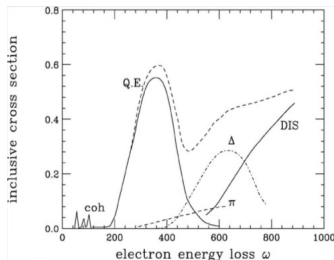


Introduction: electron energy and cross-section

Electron energy easy to know:



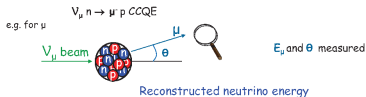
Electron scattering in nuclei:



Benhar, Day, Sick, RMP (2008)

Introduction: neutrino energy and cross-section

E_ν difficult to reconstruct. Example: CCQE process

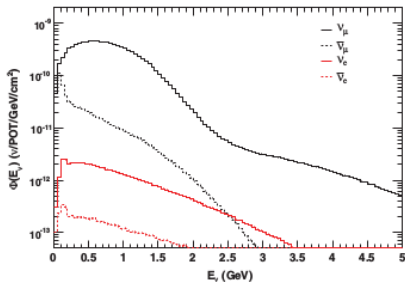


$$\overline{E_\nu} = \frac{E_\mu - m_\mu^2/(2M)}{1 - (E_\mu - P_\mu \cos \theta)/M}$$

via two-body kinematics

Neutral current process even more difficult.

Simulation of neutrino energy distribution:



MiniBooNE Coll., PRD (2009)

Knowledge of cross-section
+ near detector
= determination of E_ν

Inclusive scattering

Electron scattering:

$$\left(\frac{d^2\sigma}{d\epsilon' d\Omega} \right)_{\nu/\bar{\nu}} = \left(\frac{d\sigma}{d\Omega} \right)_M \left[\frac{Q^4}{q^4} R_L(q, \omega) + \left(\frac{Q^2}{2q^2} + \tan^2 \frac{\theta}{2} \right) R_T(q, \omega) \right]$$

R_T and R_L transverse and longitudinal response functions.

Neutrino scattering:

$$\left(\frac{d^2\sigma}{d\epsilon' d\Omega} \right)_{\nu/\bar{\nu}} = \frac{G^2}{2\pi^2} k' \epsilon' \cos^2 \frac{\theta}{2} \left[R_{00}(q, \omega) + \frac{\omega^2}{q^2} R_{zz}(q, \omega) - \frac{\omega}{q} R_{0z}(q, \omega) + \left(\tan^2 \frac{\theta}{2} + \frac{Q^2}{2q^2} \right) R_{xx+yy}(q, \omega) \mp \tan \frac{\theta}{2} \sqrt{\tan^2 \frac{\theta}{2} + \frac{Q^2}{q^2}} R_{xy}(q, \omega) \right]$$

R_{00} , R_{zz} , R_{0z} , R_{xx+yy} , and R_{xy} neutrino response functions.

R_{xy} is important for ν vs $\bar{\nu}$ processes.

$$\begin{aligned} R(q, \omega) &= \sum_n \langle \Psi | j^\dagger(q) | n \rangle \langle n | j(q) | \Psi \rangle \delta(\omega - E_n + E_0) \\ &= \int dt \langle \Psi | j^\dagger(q) \exp[i(H - \omega)t] j(q) | \Psi \rangle \\ &= \int dt E(q, \tau) \end{aligned}$$

Using QMC we can calculate **exactly** $E(q, \tau)$ and then reconstruct $R(q, \omega)$.

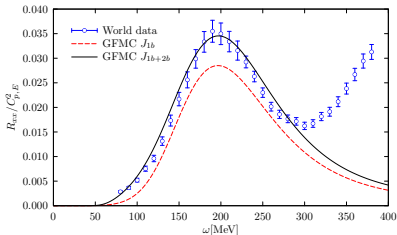
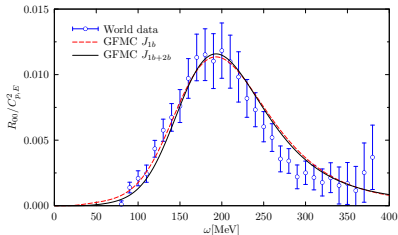
Ingredients:

- Hamiltonian H
- Ground-state Ψ (H)
- Currents described by the electroweak operators $\mathbf{j}(q)$, constructed consistently with H .

Response functions

Using the maximum entropy method, we can reconstruct the response functions.

Longitudinal and transverse response functions of ^4He ($q=600$ MeV)

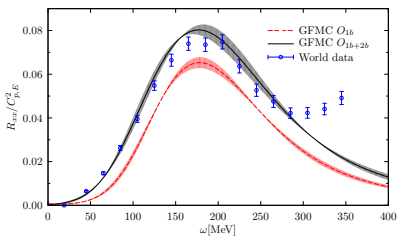
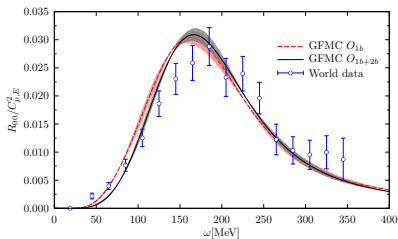


Lovato, Gandolfi, Carlson, Pieper, Schiavilla, PRC (2015)

Similar agreement also with other kinematics, $q=400$, 500 , and 700 MeV.

Electromagnetic response functions of ^{12}C

Electromagnetic longitudinal and transverse response functions of ^{12}C
($q=570$ MeV)

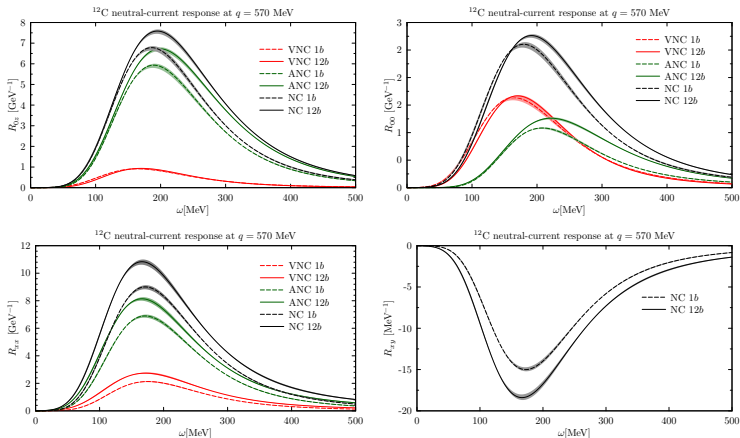


Lovato et al., PRL (2016).

Role of two-nucleon currents very important (as expected).

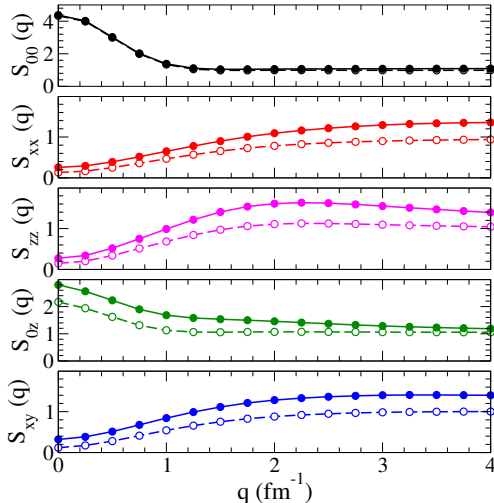
Electroweak response functions of ^{12}C

Transverse vector, axial, and neutral current of ^{12}C ($q=570$ MeV)



Lovato, et al., PRC 97, 022502 (2018)

Electroweak sum-rules in ^{12}C

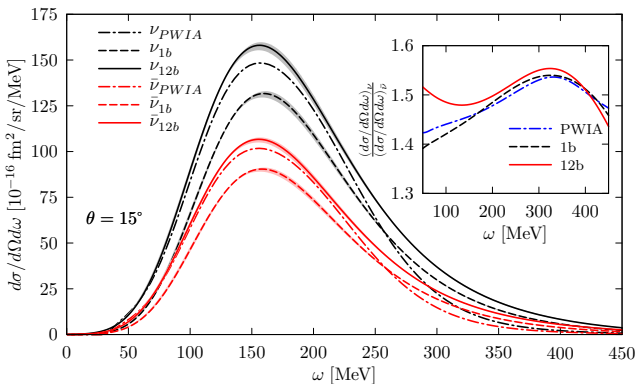


Lovato, Gandolfi, Carlson, Pieper, Schiavilla, PRL (2014).

Two-body operators enhance sum-rules up to 50%.

Electroweak cross-section of ^{12}C

From the response functions, we can calculate the cross-section:



Lovato, et al., PRC 97, 022502 (2018)

Summary and future work

Conclusions:

- “Quenching” of g_A *maybe* understood. Two-body currents and nuclear correlations very important.
- Electron scattering in ^{12}C calculated using GFMC. Good agreement with experiments. One- and two-body vector currents tested.
- Two-body axial currents show a similar enhancement in response functions and sum rules.

In progress/future work:

- Calculation of charge changing weak currents.
- Extension to larger nuclei.

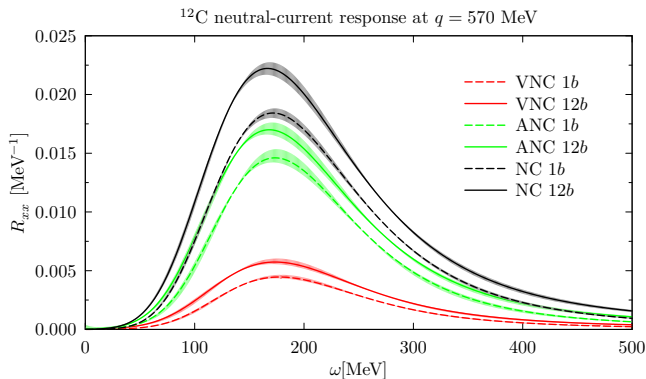
Acknowledgments:

- Joe Carlson (LANL)
- S. Pastore (WUSTL)
- A. Lovato, S. Pieper (ANL)
- R. Schiavilla (Jlab/ODU)

Extra slides

Euclidean electroweak response functions of ^{12}C

Transverse vector, axial, and neutral current of ^{12}C ($q=570$ MeV)

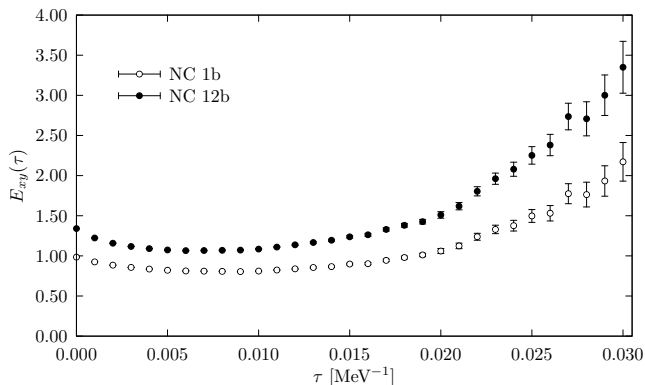


Axial currents give the largest contribution.

Role of axial form factor?

Euclidean electroweak response functions of ^{12}C

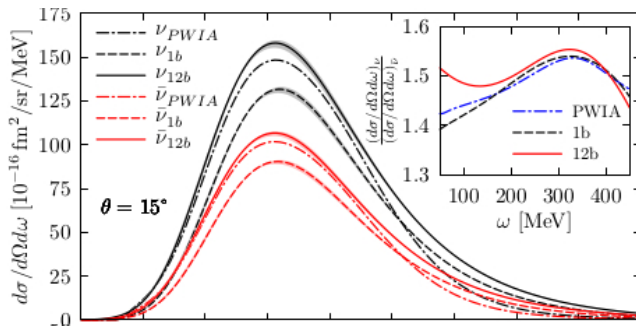
R_{xy} term responsible for ν vs $\bar{\nu}$ response. ^{12}C , $q=570$ MeV



Lovato, Gandolfi, Carlson, Pieper, Schiavilla, PRC (2015)

Electroweak cross-section of ^{12}C

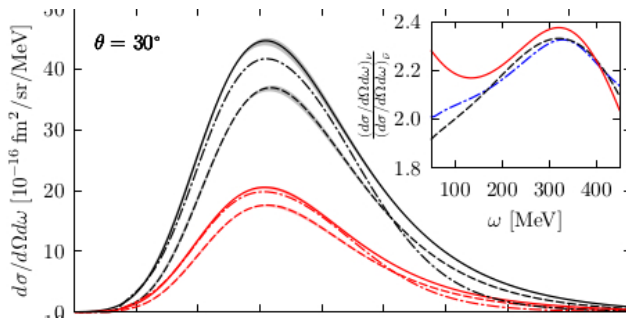
From the response functions, we can calculate the cross-section:



Lovato, et al., PRC 97, 022502 (2018)

Electroweak cross-section of ^{12}C

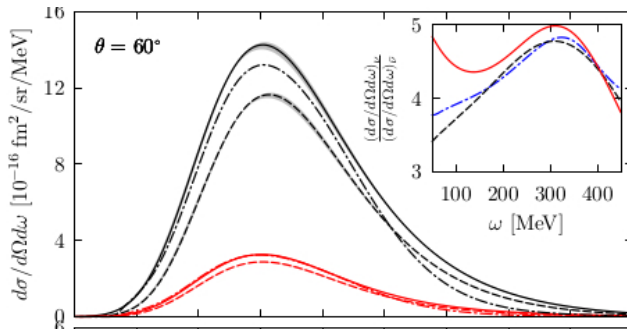
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Electroweak cross-section of ^{12}C

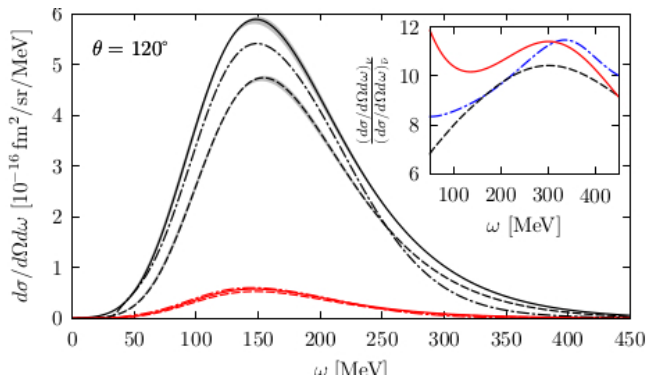
From the response functions, we can calculate the cross-section:



Lovato, et al., PRC 97, 022502 (2018)

Electroweak cross-section of ^{12}C

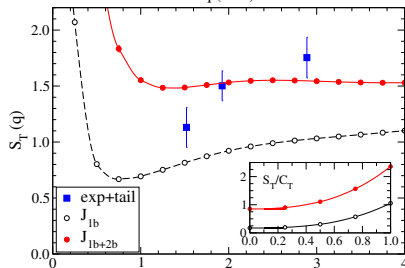
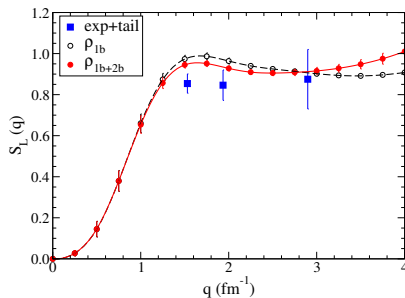
From the response functions, we can calculate the cross-section:



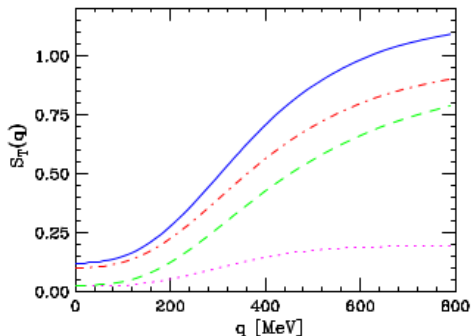
Lovato, et al., PRC 97, 022502 (2018)

Electromagnetic sum-rules in ^{12}C

Sum rules: $S_{L,T}(q) = C_{L,T} \int R_{L,T}(\omega, q) d\omega$



Transverse sum rule of ^{12}C

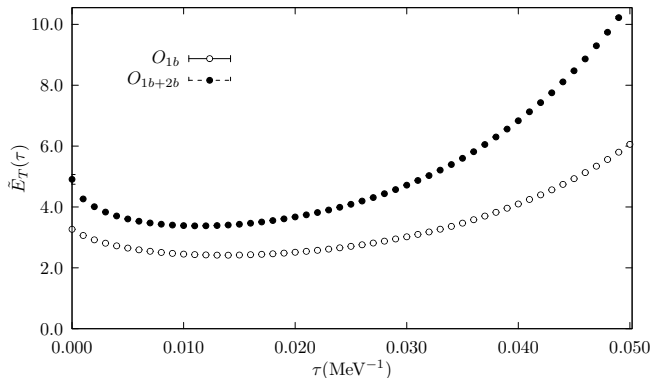


- Full response,
 $\langle \Psi | (j_1 + j_2)^\dagger (j_1 + j_2) | \Psi \rangle$
- $\langle \Psi | j_1^\dagger j_1 + j_2^\dagger j_2 | \Psi \rangle$
- $\langle \Psi | j_1^\dagger j_1 | \psi \rangle$
- interference

Benhar, Lovato, Rocco, PRC (2015)

Euclidean response

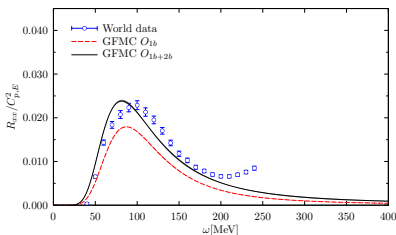
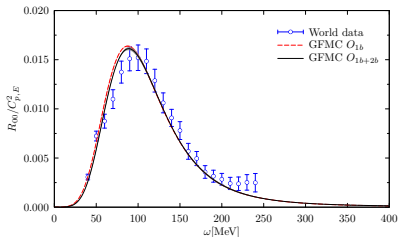
Transverse electromagnetic (euclidean) response functions of ^4He ($q=500$ MeV)



Note: results multiplied by $\exp(\tau q^2/2m)$

Longitudinal and transverse response response

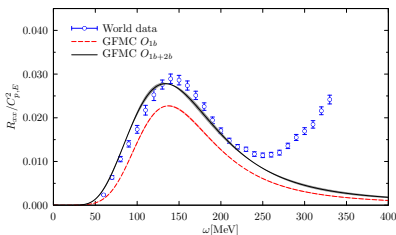
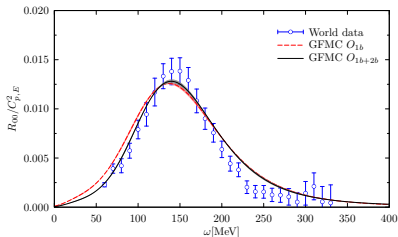
Longitudinal and transverse electromagnetic response functions of ^4He ($q=400$ MeV)



Note: results multiplied by $\exp(\tau q^2/2m)$

Longitudinal and transverse response response

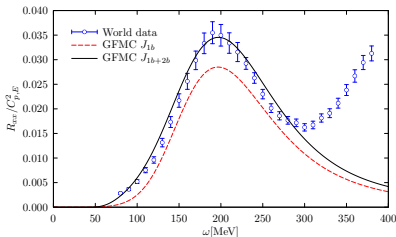
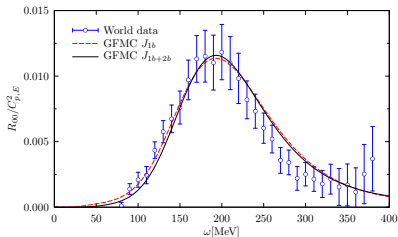
Longitudinal and transverse electromagnetic response functions of ^4He ($q=500\text{ MeV}$)



Note: results multiplied by $\exp(\tau q^2/2m)$

Longitudinal and transverse response response

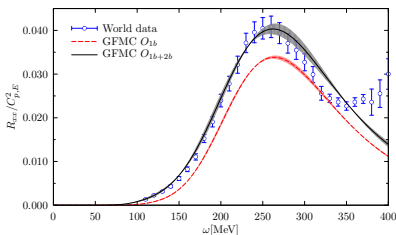
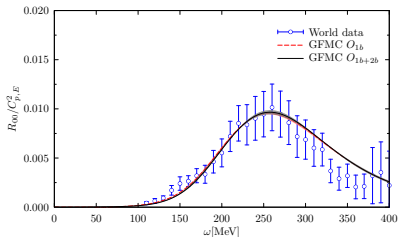
Longitudinal and transverse electromagnetic response functions of ^4He ($q=600\text{ MeV}$)



Note: results multiplied by $\exp(\tau q^2/2m)$

Longitudinal and transverse response response

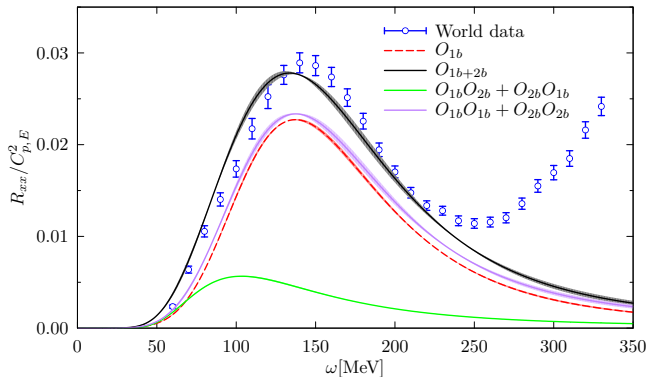
Longitudinal and transverse electromagnetic response functions of ^4He
($q=700$ MeV)



Note: results multiplied by $\exp(\tau q^2/2m)$

Transverse response response

Transverse electromagnetic response functions of ^4He ($q=500$ MeV).
Role of the interference:



Electromagnetic sum-rules in ^{12}C

Sum rules: $S_{L,T}(q) = C_{L,T} \int R_{L,T}(\omega, q) d\omega$

