

Calculation of Standard Model β spectra

Leendert Hayen

ECT* Workshop, April 8th 2019

IKS, KU Leuven, Belgium

Trento is enthusiastic for symmetry breaking



ACFI Workshop, Nov 1-3 2018



Table of contents

Introduction

Beta spectrum shape

Current status

Neutron V_{ud} calculation

Challenges

Introduction

β decay context: Colliders

Drought at LHC leaves limited number of viable theories standing

FCC is still rather far away, if it comes



β decay context: (B)SM

Standard Model internal consistency test through CKM unitarity

- Superalowed $0^+ \rightarrow 0^+$ decays
- $T = 1/2$ mirror decays
- Neutron decay (see later)

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BSM: Since EFT entered scene directly compare high and low energy:

- Competitive for scalar & tensor currents
- Complementary for right-handed currents

Cirigliano *et al.*, PPNP 71 (2013) 93

Significant experimental progress, new & improved techniques

- Atomic traps: Jerusalem, TRIUMF, CERN, ANL, TAM, CENPA

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- High-precision spectrum shapes: NCSL, LANL, CENPA, ...

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Great for theory!

Push both in β decay fundamentals and nuclear structure calculations

β decay context: Outside the box

Landscape has changed significantly past 10-20 years

- Neutrino physics: oscillations, reactor anomaly
- Astrophysics: β decay lifetimes, r -process
- Big Bang Nucleosynthesis: mass abundances

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Knowledge of Standard Model β decay spectra plays a significant role in all of them

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Different regimes, overlapping challenges

General Hamiltonian

$$\mathcal{H} = \sum_{j=V,A,S,P,T} \langle f | \mathcal{O}_j | i \rangle \langle e | \mathcal{O}_j [C_j + C_j' \gamma_5] | \nu \rangle + h.c.$$

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QCD influences \rightarrow *induced* currents, influenced through **nuclear structure**?

BSM Observables in β decay

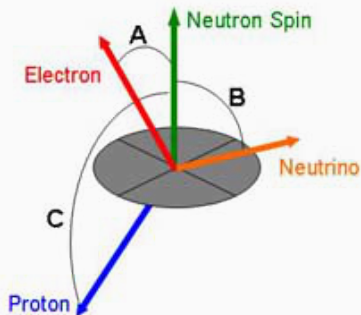
Typical BSM searches through correlations

$$\frac{d\Gamma}{dE_e d\Omega_e d\Omega_\nu} \propto 1 + a_{\beta\nu} \frac{\vec{p}_e \cdot \vec{p}_\nu}{E_e E_\nu} + b_F \frac{m_e}{E_e} + A \frac{\vec{p}_e}{E_e} \langle \vec{l} \rangle + \dots$$

Measure effective correlations

$$\tilde{X} = \frac{X}{1 + b_F \langle \frac{m_e}{E_e} \rangle}$$

Sensitivity to full spectrum!



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Sensitivity comes from b_F

$$b_F = \pm \frac{1}{1 + \rho^2} \left[\text{Re} \left(\frac{C_S + C'_S}{C_V} \right) + \rho^2 \text{Re} \left(\frac{C_T + C'_T}{C_A} \right) \right]$$

because it's **linear** in coupling constants

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because it's **linear** in coupling constants

→ measure β spectrum directly & fit for $1/E_e$

Beta Spectrum Shape

Exploring the Standard Model and Beyond via the allowed β spectrum shape:

$$\frac{dN}{dE_e} \propto 1 + b_{\text{Fierz}} \frac{m_e}{E_e} + b_{WM} E_e$$

b_{Fierz} : Proportional to scalar (Fermi) and tensor (Gamow-Teller) couplings

b_{WM} : Weak Magnetism (main induced current), poorly known for $A > 60$, forbidden decays

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This requires knowledge of the theoretical spectrum shape to $\leq 10^{-3}$ level!

Beta spectrum shape

Beta Spectrum Shape

Active participation of QED, QCD & WI \rightarrow Complicated system

Weak Hamiltonian is **modified**

1. Emitted β particle immersed in Coulomb field: (electroweak) radiative corrections

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Large scale gap to cross

Quark \rightarrow Nucleon \rightarrow Nucleus \rightarrow Atom \rightarrow Molecule

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Whole slew of **approximations** introduced

Standard Model Calculation: Quark

Starting from the Standard Model $SU(2)_L \times U(1)_Y$ EW sector

$$\mathcal{M} = \frac{g^2}{8} V_{ud} \bar{u} \gamma^\mu (1 - \gamma^5) d \frac{g_{\mu\nu} - q_\mu q_\nu / M_W^2}{q^2 - M_W^2} \bar{e} \gamma^\nu (1 - \gamma^5) \nu$$

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Since $q \ll M_W$, identify Fermi coupling constant

$$\frac{G_F}{\sqrt{2}} = \frac{g^2}{8M_W^2}$$

Standard Model Calculation: Nucleon

Moving to the nucleon system, we face

$$\langle p | \bar{u} \gamma^\mu (1 - \gamma^5) d | n \rangle$$

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Symmetries to the rescue! CVC & PCAC define new *nucleon* currents

$$V^\mu + A^\mu \approx g_V(q^2) \gamma^\mu (1 - \lambda \gamma^5)$$

where $g_V(q^2) \approx 1$ and λ from the lattice

Great progress from lattice QCD, including scalar & tensor charges

Standard Model Calculation: Nucleon

Strong interaction introduces extra terms into the vertex \rightarrow
Construct all Lorentz invariants

$$\begin{aligned}\langle p|V^\mu|n\rangle &= \bar{p} \left[g_V \gamma^\mu + \frac{g_M - g_V}{2M} \sigma^{\mu\nu} q_\nu + i \frac{g_S}{2M} q^\mu \right] n \\ \langle p|A^\mu|n\rangle &= \bar{p} \left[g_A \gamma^\mu \gamma^5 + \frac{g_T}{2M} \sigma^{\mu\nu} q_\nu \gamma^5 + i \frac{g_P}{2M} q^\mu \gamma^5 \right] n\end{aligned}$$

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Introduction of recoil ($\sim q/M$) terms

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Introduction of recoil ($\sim q/M$) terms

CVC requires $g_S = 0$ & $g_M = \mu_p^{an} - \mu_n = 4.7$

Standard Model Calculation: Nucleus

Nucleus is spherical system \rightarrow multipole decomposition,
elementary particle

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Relativistic generalization in Breit frame

$$\langle f | V^0 + A^0 | i \rangle \propto \sum_{LM} (-)^{J_f - M_f} \begin{pmatrix} J_f & L & J_i \\ -M_f & M & M_i \end{pmatrix} (Y_L^M)^* F_L(q^2)$$

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Conservation of angular momentum limits $\#$ terms

Standard Model Calculation: Nucleus

Require transformation from form factors to matrix elements

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Immediately faced with several issues:

- Weak current in strongly bound system?

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Here the going gets rough → severe approximations

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→ **Impulse approximation**, non-interacting nucleons

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- Neglects meson exchange
- Nucleon-nucleon interaction present in many-body methods

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Relativistic nuclear wave functions

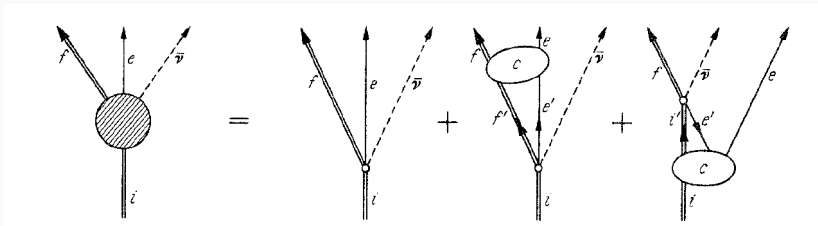
→ **Non-relativistic** nucleons

- expand operator to $\mathcal{O}(v/c)$
- Incomplete wave function basis, core polarization

Standard Model Calculation: Nucleus

Final state interactions

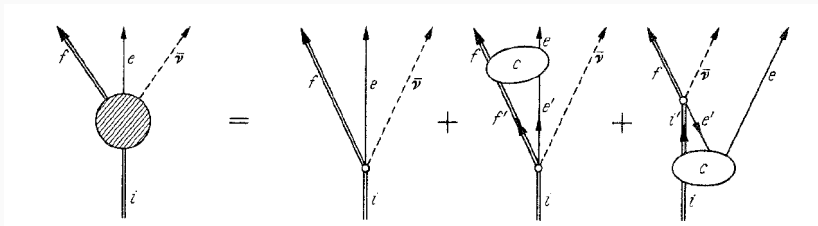
1. Coulomb interaction



Standard Model Calculation: Nucleus

Final state interactions

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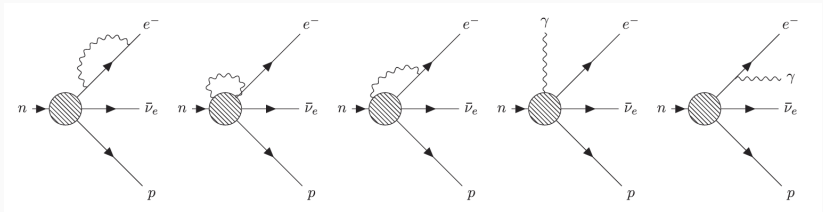


→ Fermi function, *induced Coulomb* terms

Standard Model Calculation: Nucleus

Final state interactions

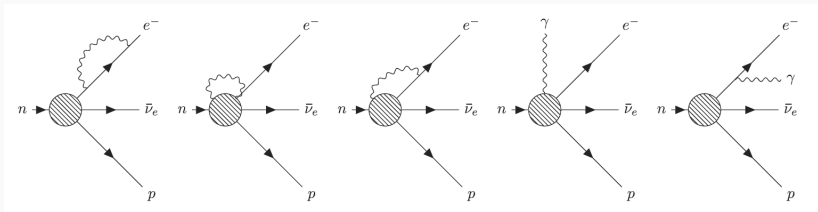
2. EW Radiative corrections



Standard Model Calculation: Nucleus

Final state interactions

2. EW Radiative corrections



+ higher orders, γW boxes: talks by M. Gorshteyn, C. Y. Seng, M. Ramsey-Musolf

Standard Model Calculation: Atom

Must consider total nuclear + atomic Hamiltonian

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Changes

- Available phase space
- Final state interactions
- Opens new decay modes (bound & exchange)

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Changes

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Require *atomic* wave functions

- Central & static potential
- Sudden approximation

Standard Model Calculation: Molecule

Similar as atomic system, but changes

- Available phase space
- Molecular excitation, ionization
- Recoil correction & distribution

Standard Model Calculation: Molecule

Similar as atomic system, but changes

- Available phase space
- Molecular excitation, ionization
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Enter quantum chemistry

- Born-Oppenheimer approximation
- MOLCAO

Current status

Beta Spectrum Shape

Significant effort to rediscover & renew formalisms

Beta spectrum of unique first-forbidden decays as a novel test for fundamental symmetries

Ayala Glick-Magid^a, Yonatan Mishnayot^{a,b,c}, Ish Mukul^b, Michael Hass^b,
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PHYSICAL REVIEW C **95**, 024327 (2017)

Spectrum-shape method and the next-to-leading-order terms of the β -decay shape factor

M. Haaranen, J. Kotila, and J. Suhonen

Department of Physics, University of Jyväskylä, P.O. Box 35 (YFL), FI-40014 Jyväskylä, Finland

(Received 10 November 2016; published 28 February 2017)

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Relativistic Theory and Ab Initio Simulations of Electroweak Decay Spectra in Medium-Heavy Nuclei and of Atomic and Molecular Electronic Structure

Tommaso Morresi, Simone Taioli✉, Stefano Simonucci

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$$\begin{aligned} N(W)dW = & \frac{G_V^2 V_{ud}^2}{2\pi^3} F_0(Z, W) L_0(Z, W) U(Z, W) R_N(W, W_0, M) \\ & \times Q(Z, W, M) R(W, W_0) S(Z, W) X(Z, W) r(Z, W) \\ & \times C(Z, W) D_C(Z, W, \beta_2) D_{FS}(Z, W, \beta_2) \\ & \times pW(W_0 - W)^2 dW \end{aligned}$$

LH *et al.*, Rev. Mod. Phys. 90 (2018) 015008

LH, Severijns, Comp. Phys. Comm. 10.1016/j.cpc.2019.02.012



Order of magnitude estimates

Nuclear structure sensitivity in shape factor

$$C(Z, W) \sim 1 \pm \frac{4}{3} \frac{W}{M_N} \frac{\mathbf{b}}{A_c} \pm \frac{4\sqrt{2}}{21} \alpha ZWR\Lambda - \frac{1}{3WM_c} (\pm 2\mathbf{b} + \mathbf{d})$$

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Fill in typical numbers to obtain

| Matrix element | Name | Slope (% MeV ⁻¹) |
|----------------|----------------------|------------------------------|
| b | Weak Magnetism | 0.5 |
| d | Induced Tensor | 0.1 |
| Λ | Induced Pseudoscalar | 0.1 |

Weak magnetism is generally more stable than others

→ **essential** to get this right

Nuclear matrix elements

Overview

$$b = A(g_M \mathcal{M}_{GT} + g_V \mathcal{M}_L)$$

$$d = A(g_A \mathcal{M}_{\sigma L} + g_T \mathcal{M}_{GT})$$

$$\Lambda \propto \mathcal{M}_{1y}$$

Nuclear matrix elements

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where

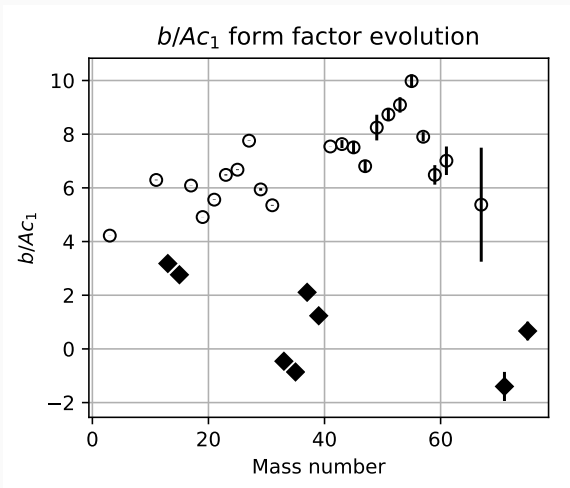
$$\mathcal{M}_L = \langle \beta | \sum \tau_i \vec{l}_i | \alpha \rangle$$

$$\mathcal{M}_{\sigma L} = \langle \beta | \sum \tau_i i \vec{\sigma}_i \times \vec{l}_i | \alpha \rangle$$

$$\mathcal{M}_{1y} \propto \langle \beta | \sum \tau_i r^2 C_{121}^{nn'k} \sigma_{i,n} Y_2^{n'}(\hat{r}_i) | \alpha \rangle$$

Weak magnetism

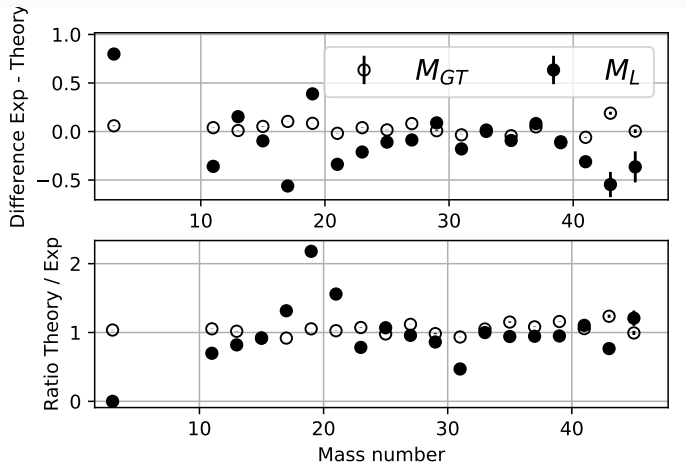
Mirror nuclei have CVC-determined WM



open: $I + 1/2$, closed: $I - 1/2$

Weak magnetism

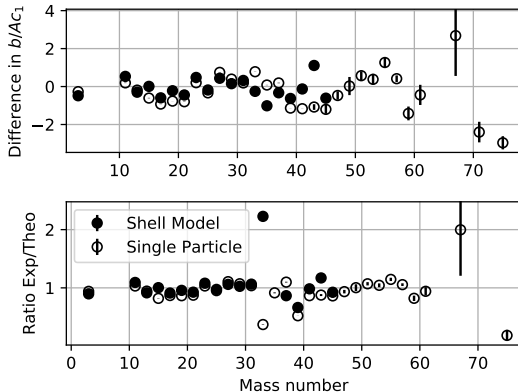
How does shell model perform right now?



‘Easy’ matrix elements only accurate to 10-20%

Weak magnetism

How does shell model perform right now?



$$\Delta b/Ac = 1 \rightarrow 0.1\% \text{ MeV}^{-1}$$

Still large discrepancies for d/Ac

PHYSICAL REVIEW C **95**, 035501 (2017)

2_1^+ to 3_1^+ γ width in ^{22}Na and second class currents

S. Triambak,^{1,2,*} L. Phuthu,¹ A. García,³ G. C. Harper,³ J. N. Orce,¹ D. A. Short,³ S. P. R. Steininger,³ A. Diaz Varela,⁴
R. Dunlop,⁴ D. S. Jamieson,⁴ W. A. Richter,¹ G. C. Ball,⁵ P. E. Garrett,⁴ C. E. Svensson,⁴ and C. Wrede^{3,6}

$$21(6) \geq d/Ac \geq 3(6)$$

Factor 7 differences depending on shell model results \rightarrow killer!

Immediate response last workshop (Alex Brown)

| 3+ to 2+ | USDB | USDA | USD |
|----------------------|-------|-------|-------|
| M(s-tau) (c_1) | 0.042 | 0.012 | 0.027 |
| M(l-tau) (part of b) | -1.07 | -1.00 | -1.00 |
| M(d-tau) | 0.062 | 0.081 | 0.066 |

Relative phases look robust but s-tau is not very uncertain
so we should look at b/d (not b/c and d/c)

Alex Brown, ND2013, NYC, March 4, 2013

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General lesson: when looking at allowed transitions, make sure \mathcal{M}_{GT} is large and stable

Neutron V_{ud} calculation

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Neutron is extremely well-studied system, ideal system for V_{ud}

$$|V_{ud}|^2 \tau_n (f_V + 3f_A \lambda^2) = \frac{2\pi^3}{G_F^2 m_e^5 g_V^2} \frac{1}{1 + RC}$$

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From β decay perspective, need 3 things

- Neutron lifetime

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- λ
- Theory calculations for $f_{V,A}$ and RC

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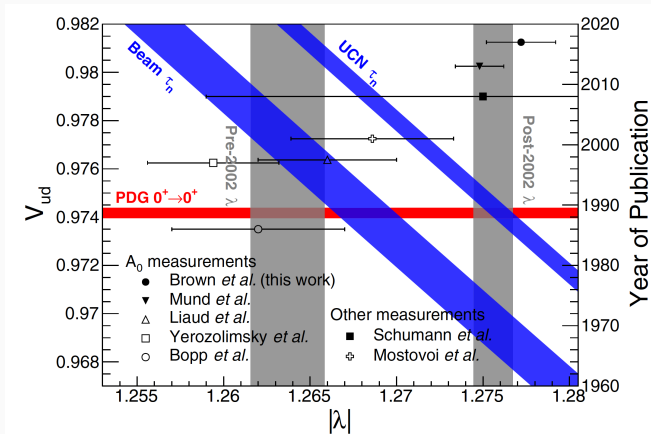
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Clearly, all *trivial* things

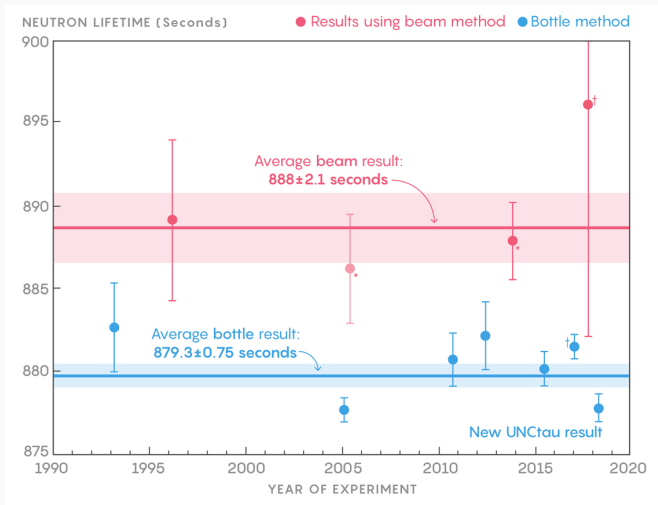
Neutron V_{ud} calculation

Major decades-long community efforts



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Neutron V_{ud} calculation

Well, at least $f_{V,A}$ are well-known, right? RIGHT?

Neutron V_{ud} calculation

Well, at least $f_{V,A}$ are well-known, right? RIGHT?

Seminal work by Wilkinson in 1982, exhaustively listed all corrections: found $\Delta f_{V,A} \simeq 10^{-6}$, $f_V = 1.6887(2)$

Neutron V_{ud} calculation

Well, at least $f_{V,A}$ are well-known, right? RIGHT?

Seminal work by Wilkinson in 1982, exhaustively listed all corrections: found $\Delta f_{V,A} \simeq 10^{-6}$, $f_V = 1.6887(2)$

One particular case appears forgotten, however...

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Recap:

$$\langle p|V^\mu|n\rangle = \bar{p} \left[g_V \gamma^\mu + \frac{g_M - g_V}{2M} \sigma^{\mu\nu} q_\nu + i \frac{g_S}{2M} q^\mu \right] n$$

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gives rise to spectrum shape contribution

$$\begin{aligned} \left(\frac{dN}{dW_e} \right)^{\text{wm}} &\propto \frac{4}{3M} \frac{g_M}{g_A \mathcal{M}_{GT}} p_e W_e (W_0 - W_e)^2 \\ &\times \left(W_e - \frac{W_0}{2} - \frac{m_e^2}{2W_e} \right) \end{aligned}$$

represents vector-axial vector spacelike cross term

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Except...

Weinberg, Phys Rev 115 (1959) 481

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There is one more thing: Coulomb corrections on weak magnetism
gives non-negligible terms $\mathcal{O}(\alpha Z/MR)$ besides expected
 $\mathcal{O}(\alpha Z(q/M)qR)$

$$\frac{f_A}{f_V} = 1 + \frac{4}{5} \frac{\alpha Z}{MR} \frac{g_M}{g_A} = 1.0040(2)$$

Plot twist!

Wilkinson Nucl Phys A 377 (1982) 474; Bottino *et al.* Phys Rev C 9 (1974) 2052;

Holstein Phys Rev C 10 (1974) 1215

Neutron V_{ud} calculation

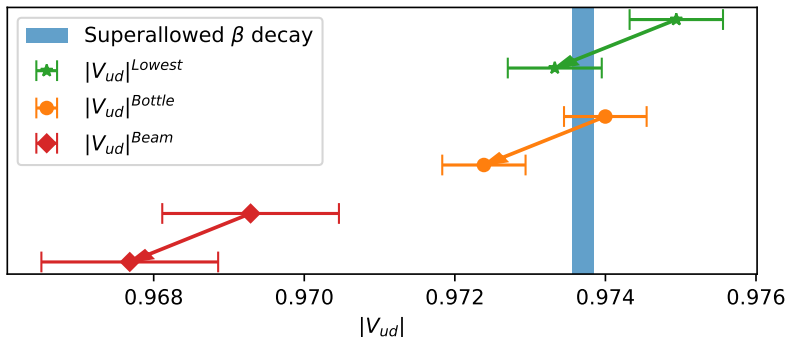
Using $\Delta_R = 0.02467(22)$, $\lambda = 1.27510(66)$

$$|V_{ud}|^2 \tau_n \left(1 + 3 \frac{f_A}{f_V} \lambda^2 \right) = 4903.5(1.1) \text{ s}$$

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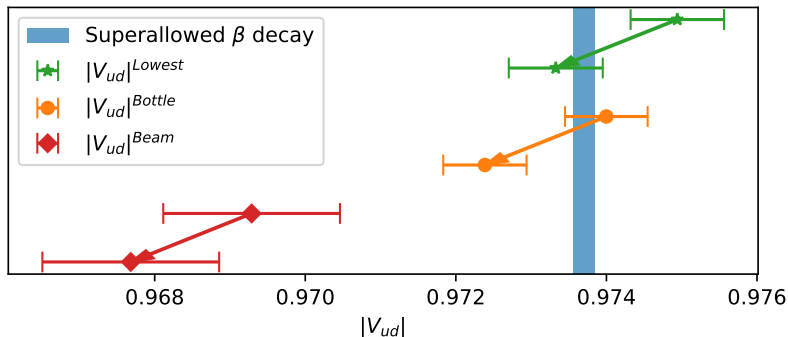
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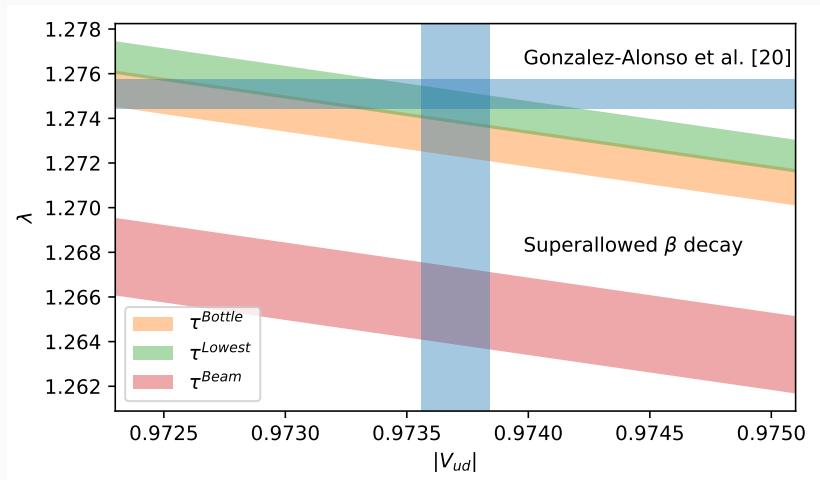
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Yikes!

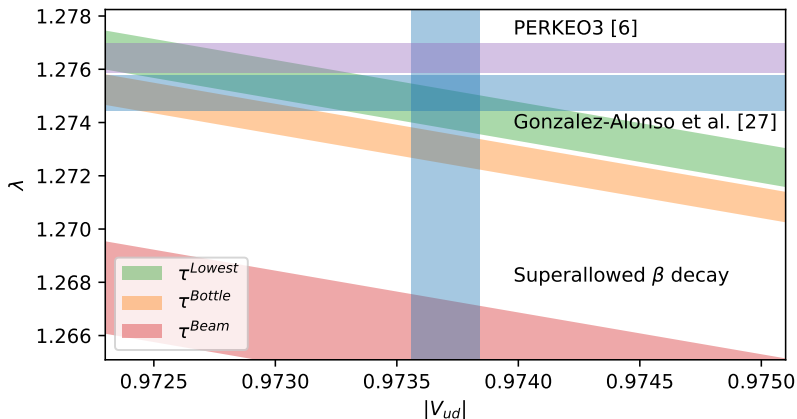
Neutron V_{ud} calculation

Using the new f_A/f_V



Neutron V_{ud} calculation

Using the new f_A/f_V , including latest PERKEO3



Neutron V_{ud} calculation

Assume superallowed V_{ud} , predict 'Standard Model' τ_n

$$\lambda = 1.27510(66) \longrightarrow \tau_n^{SM} = 877.0(8) \text{ s}$$

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Also impacts Big Bang Nucleosynthesis, helium mass abundance

$$\frac{\Delta Y_p}{Y_p} \approx 0.73 \frac{\Delta \tau_n}{\tau_n}.$$

Shift in neutron lifetime leads to

$$\Delta Y_p = -5.1 \cdot 10^{-4}$$

3 σ shift!

Challenges

At $\mathcal{O}(10^{-3})$, nuclear structure is main culprit

- Nuclear matrix elements only precise to 10-20%
- Generally: large meson exchange corrections on induced currents
- Isospin multiplet decays are way to go: WM from CVC, induced tensor = 0
- Major ab initio efforts underway

At $\leq \mathcal{O}(10^{-4})$, *everything* breaks

At $\leq \mathcal{O}(10^{-4})$, *everything* breaks , but not in the same place!

- Low energy: Atomic & Molecular effects (exchange)
- Endpoint: Final state interactions, excitations
- Radiative corrections: higher order, model dependence
- Low Z : recoil corrections to matrix elements
- High Z : everything electromagnetic

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Exciting (unnerving?) developments happening in β decay fundamentals