

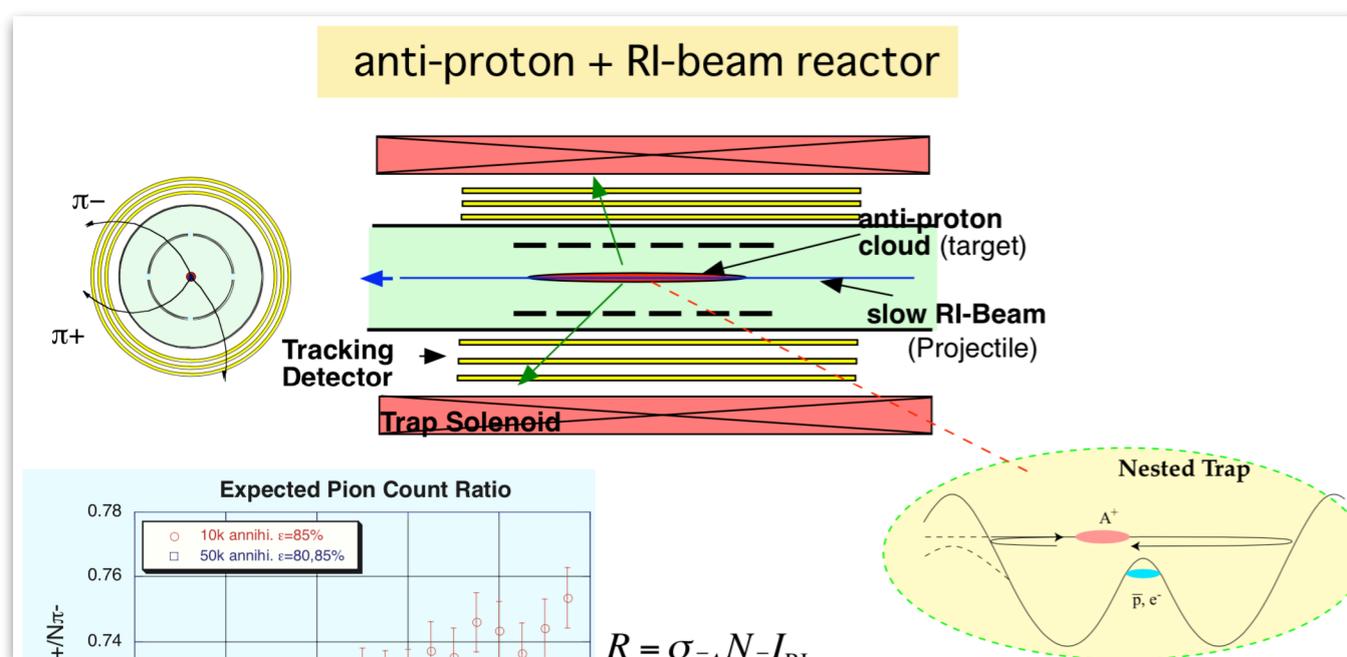
Antiprotonic atoms with short-lived nuclei

— **Obsolete Projects** —

Michiharu Wada, KEK

pbar-RI at SLOWRI/RIKEN 2003

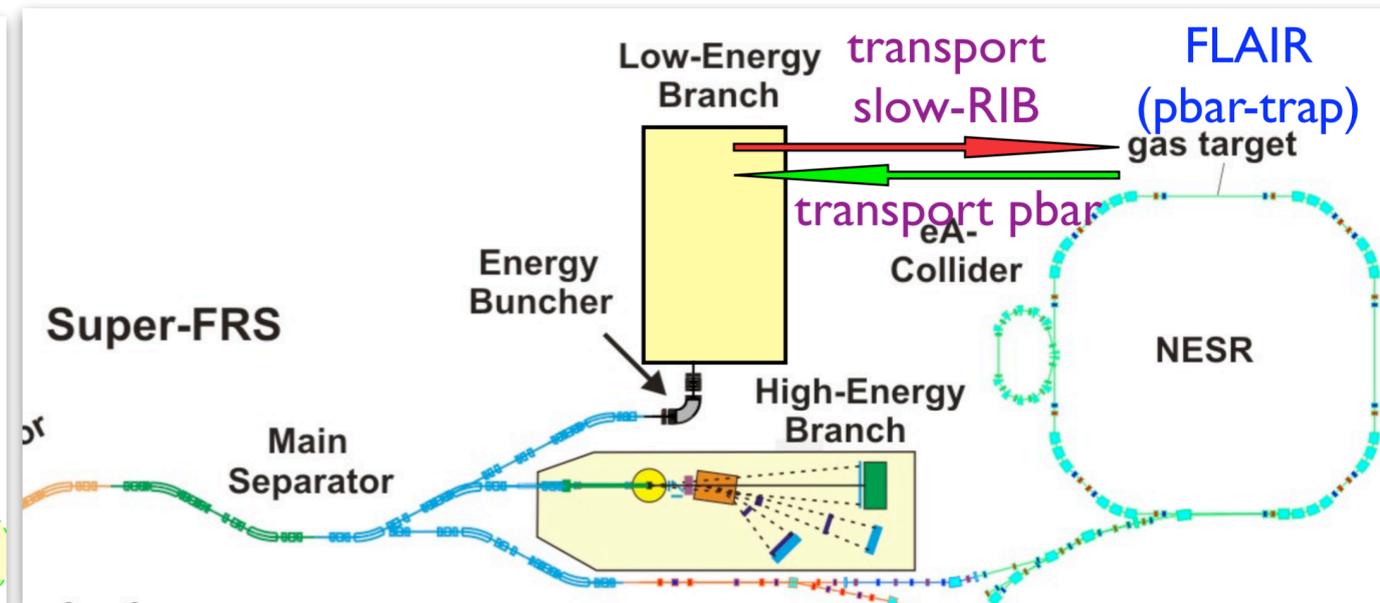
Exo-pbar at FAIR/GSI 2004



M. Wada, Y. Yamazaki NIMB214(2004)196

- ◆ trap pbar @ AD/CERN
- ◆ send RIKEN via Siberian Railway
- ◆ pbar + halo nuclei @ SLOWRI/RIKEN

too ambitious?



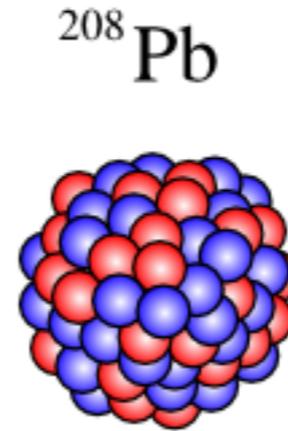
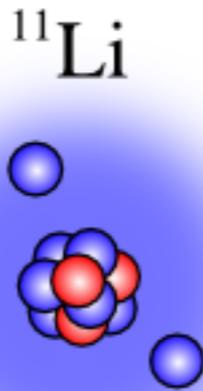
GSI TDR 2005

- ◆ trap pbar @FLAIR
- ◆ RI-beams from Super-FRS LEB
- ◆ pbar + halo nuclei @FAIR

**no facilities have been built,
no collaboration exists**

Motivation

Since 1985, “neutron halo”
becomes sensational in
nuclear physics



More Spectroscopic or
Direct indication?

Interaction
Cross Section



Matter Radius

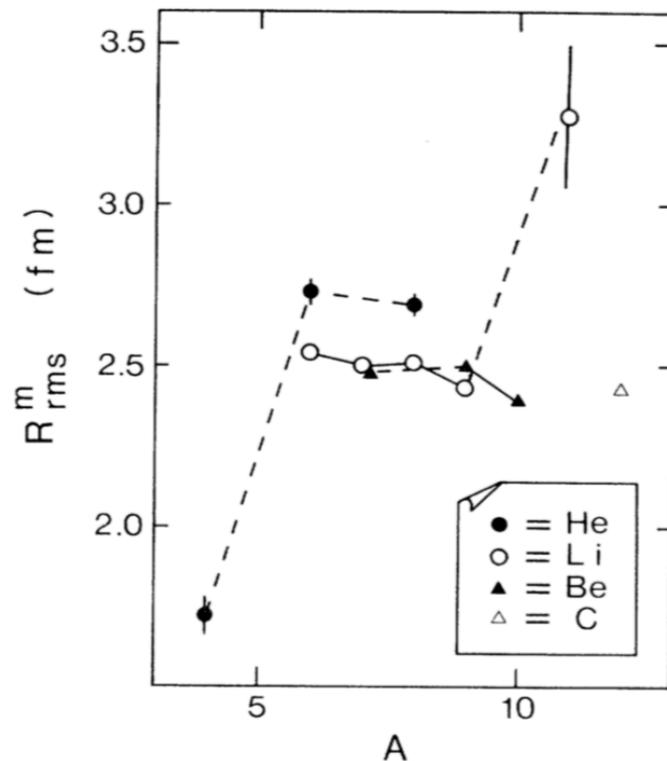


FIG. 3. Matter rms radius R_{rms}^m . Lines connecting isotopes are only guides for the eye. Differences in radii are seen for isobars with $A = 6, 8$, and 9 . The ^{11}Li isotope has a much larger radius than other nuclei.

I. Tanihata et al, PRL55(1985)2676

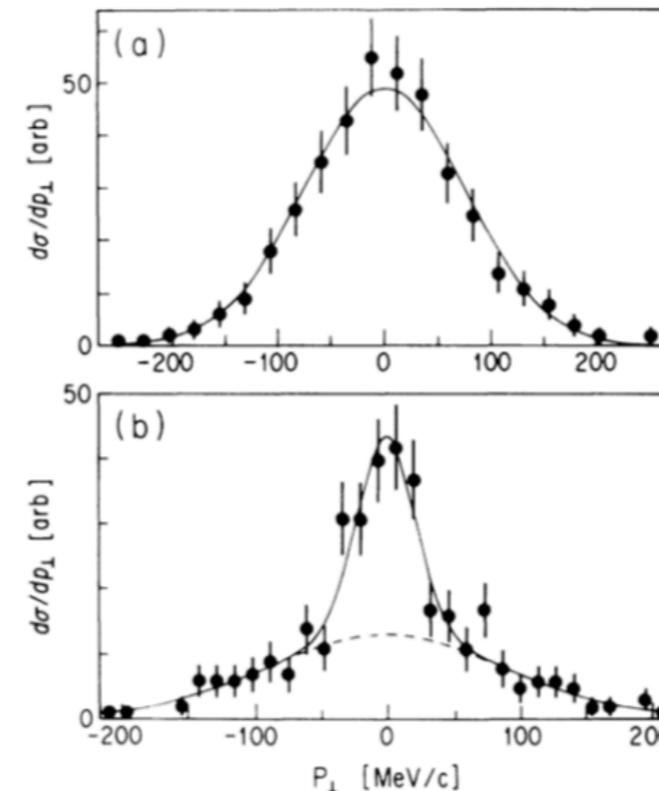


FIG. 1. Transverse-momentum distributions of (a) ^6He fragments from reaction $^8\text{He}+\text{C}$ and (b) ^9Li fragments from reaction $^{11}\text{Li}+\text{C}$. The solid lines are fitted Gaussian distributions. The dotted line is a contribution of the wide component in the ^9Li distribution.

T. Kobayashi et al, PRL60(1988)2599

Narrow Recoil
Momenta



Weakly bound
two neutrons

Optical Spectroscopy of neutron halo

Bohr-Weisskopf effect for valence neutron radius

^{11}Be

Z=4
N=7
(T_{1/2}=13s)

^{10}Be

Z=4
N=6

Core + One loosely bound Neutron

Charge Radius

IS (E0)

r_c

Volume without Charge-less Neutron

Magnetization Radius

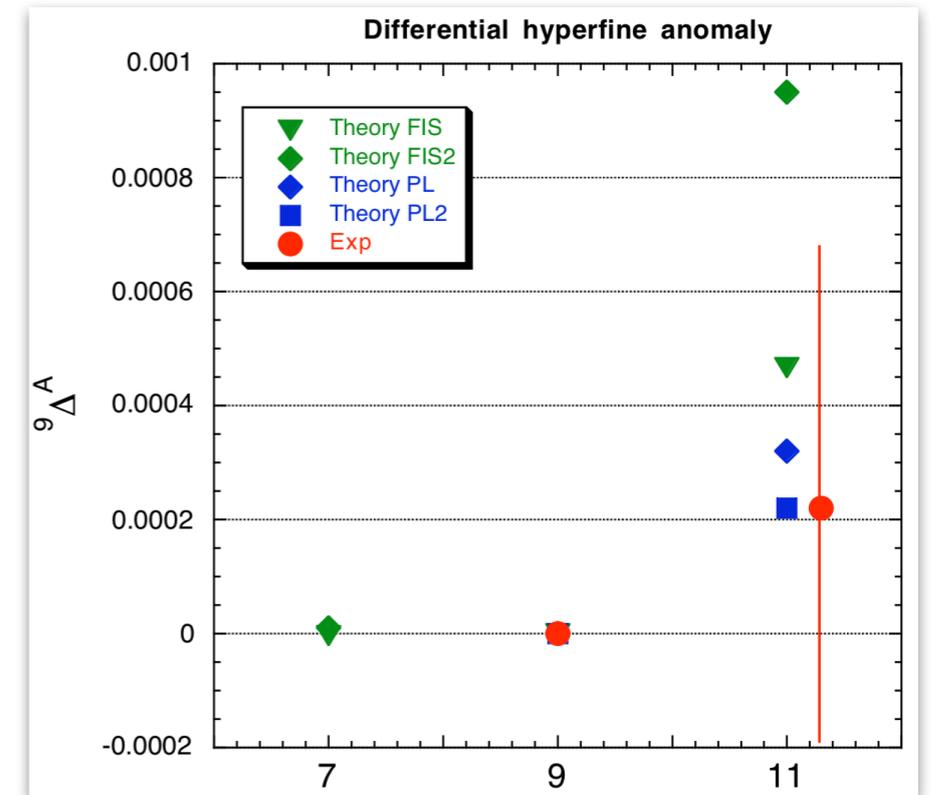
IS (M1)

Distribution of Valence Neutron

r_m

Single Valence Neutron carries most of Magnetization of Nucleus

Cancelled in the Core (core polarization still exist)



$${}^9\Delta_{11} = \frac{A_9/\mu_9}{A_{11}/\mu_{11}} - 1 = 2.2(48) \times 10^{-4}$$

	Be7	Be9	Be11	
HFS constant A (MHz)	-742.7722(4)	-625.0088370529(11)	-2677.308(2)	
Nuclear Mag. Moment (n.m)	[-1.39928(2)]	-1.177432(3)	(-)1.6816(8)*	5x10 ⁻⁴

* W. Geithner et al, PRL 83(1999)3792

More Universal Method?

Neutron halo from Antiproton absorption

Bubble Chamber
experiment in 1973

Evidence for a Neutron Halo in Heavy Nuclei from Antiproton Absorption*

W. M. Bugg, G. T. Condo, and E. L. Hart
The University of Tennessee, Knoxville, Tennessee 37916

and

PRL 31(1973)475

H. O. Cohn and R. D. McCulloch
Oak Ridge National Laboratory, Oak Ridge, Tennessee 37830
(Received 19 April 1973)

From a study of stopping antiprotons in a variety of elements located in a hydrogen bubble chamber, we find evidence for the existence of a neutron fringe in heavy nuclei.

TABLE IV. "Halo factor" analysis.

Element	$N(\pi^-)$ $-N(\pi^+)$	$N(\bar{p}n)$	$N(\bar{p}p)$	$\frac{N(\bar{p}n)}{N(\bar{p}p)}$	$\frac{N(\bar{p}n)}{N(\bar{p}p)} \Big _c$	$\frac{N}{Z}$	Halo factor
C	2302	2586	4089	0.632	1.00	1.00	1.00
Ti	881	1067	1111	0.960	1.52	1.18	1.29 ± 0.21
Ta	1006	1276	931	1.371	2.17	1.48	1.46 ± 0.24
Pb	947	1216	534	2.270	3.59	1.54	2.34 ± 0.50

Charged Pion Ratio

"Calibrate" R_{np} by C-12

$$R_{np} \equiv \sigma_{\bar{p}n} / \sigma_{\bar{p}p} \approx 0.63$$

$$f_{n\text{halo}} = \frac{N(\bar{p}n)}{N(\bar{p}p)} \cdot \frac{Z}{N} \cdot \frac{\sigma_{\bar{p}p}}{\sigma_{\bar{p}n}}$$

Antiprotonic atoms of stable nuclei by Warsaw Group

Radiochemical identification: Halo factor

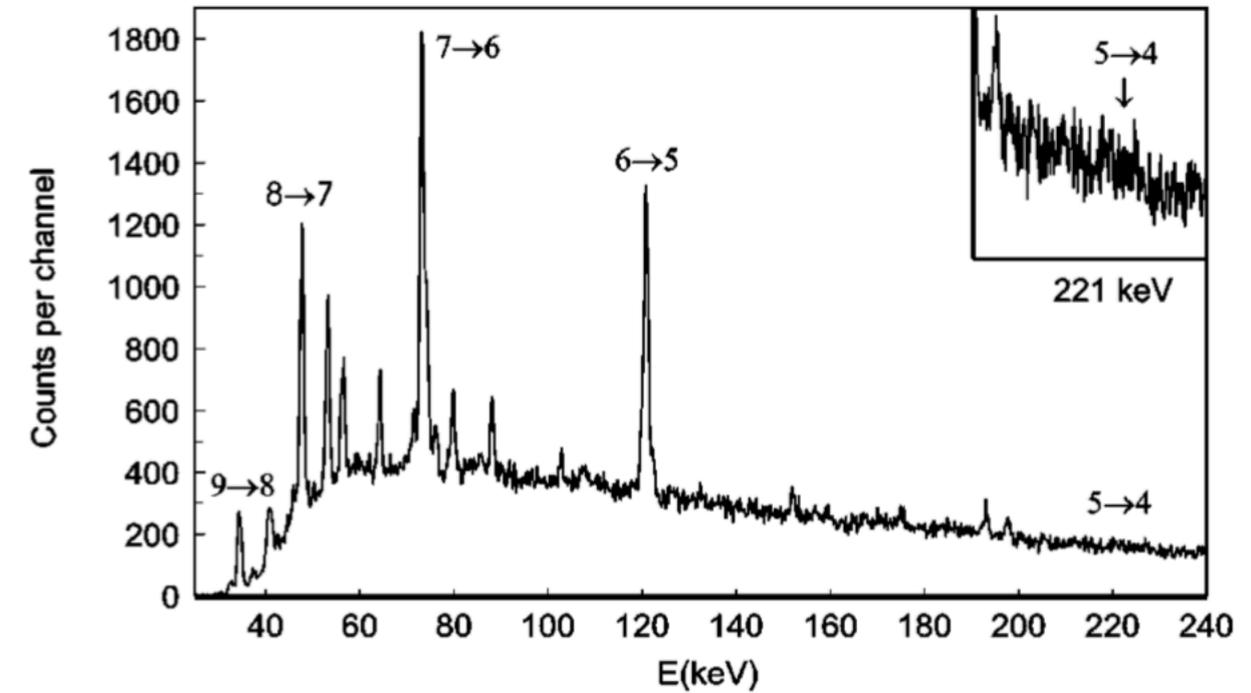
Table 1
Summary of ^{232}Th evaluation

Jastrzebski et al, NPA 558(1993)405c

	Nuclei per 1000 \bar{p}	
	Experiment	Theory, INC [24]
$^{231}\text{Th}(\bar{p}n)$	100 ± 27	19.5
$^{231}\text{Ac}(\bar{p}p)$	12.9 ± 1.6	13.3
	} 113	} 33
$\frac{N(\bar{p}n)}{N(\bar{p}p)}$	7.8 ± 2.2	1.5
$f(\text{peripheral}) \stackrel{\text{def.}}{=} \frac{N(\bar{p}n)}{N(\bar{p}p)} \frac{\text{Ima}_p}{\text{Ima}_n} \frac{Z_t}{N_t} = (7.8 \pm 2.2) \frac{1}{0.63} \frac{90}{142} = 7.8 \pm 2.2$		

Cascade X-rays limit: Nuclear size

Hartmann et al, PRC 65(2001)014306



$r_n - r_p$ vs. $(N - Z)/A$ Trzcinska et al, PRL 87(2001)082501

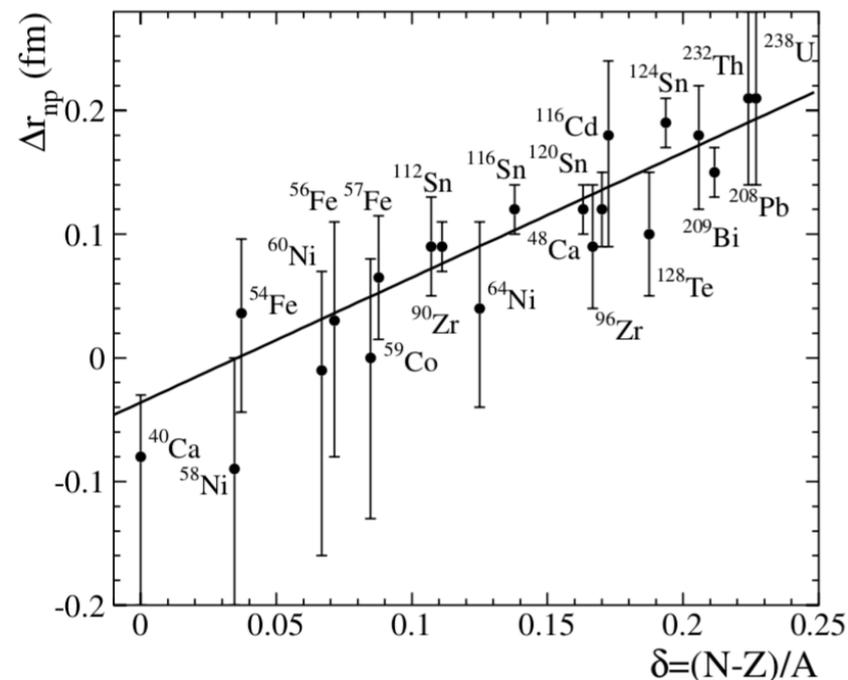


FIG. 4. Difference Δr_{np} between the rms radii of the neutron and proton distributions as deduced from the antiprotonic atom x-ray data, as a function of $\delta = (N - Z)/A$. The proton distributions were obtained from electron scattering data [41] (Sn nuclei) or from muonic atom data [38,42,43] (other nuclei). The

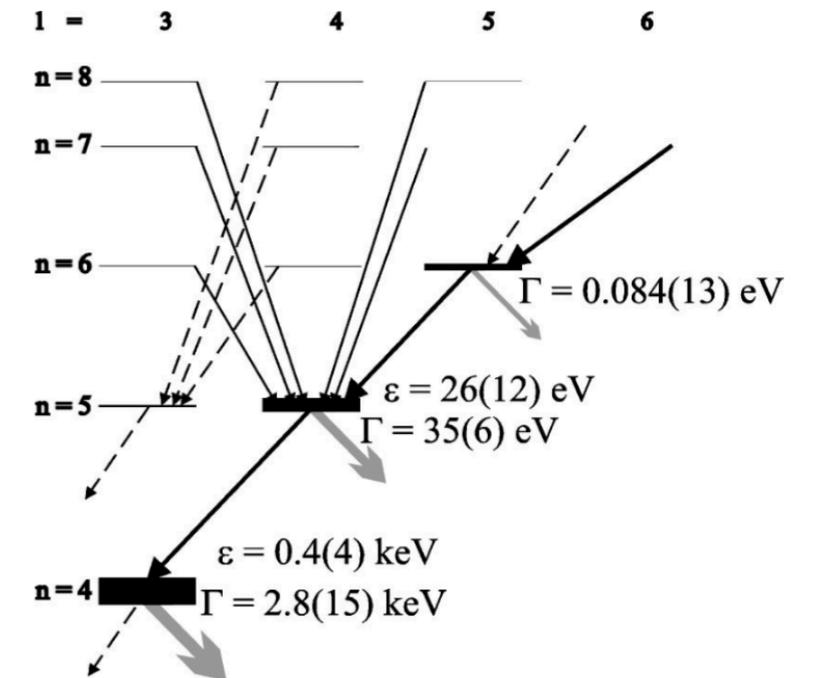
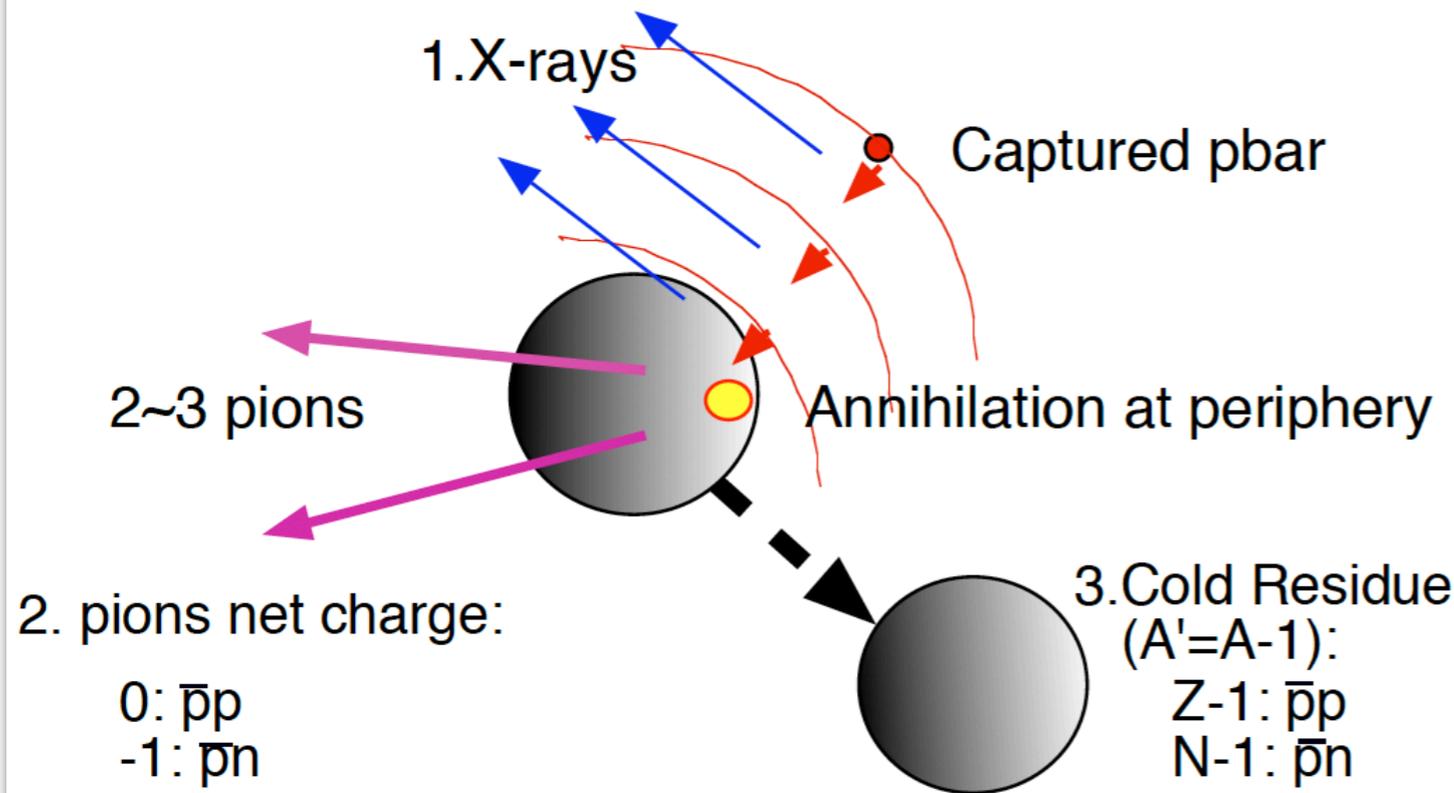


FIG. 3. Mean widths and shifts of all levels with measurable strong interaction effects. The weight of the different calcium isotopes is given in the caption of Fig. 1.

Antiprotonic radioactive atoms (Exo+pbar)

A new probe for nuclear structure study of unstable nuclei

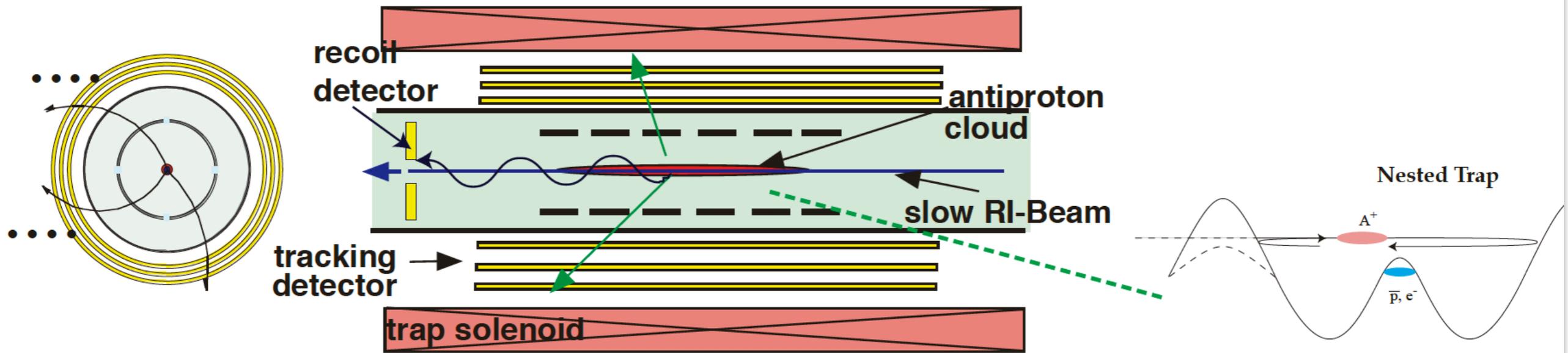


- * annihilation at $\rho \approx \rho_0 1/1000$
- * pbar distinguishes p and n

physical quantity	observable	method	for RIB*	previous works for stable nuclei
nuclear size	X-ray (min. n,l)		?	Trzcinska et al, PRL87(2001) 82501
p,n abundance at nuclear surface	pion net charge	calorimetric		Bugg et al, bubble chamber exp. for C,Ti,Ta,Pb. PRL31(1973)475
		statistical	○	
	cold residue	gamma-ray		Jastrebski et al, PRC47(1993)216
		recoil momentum PI	○	
surface nucleon's momentum	cold residue	recoil momentum	○	

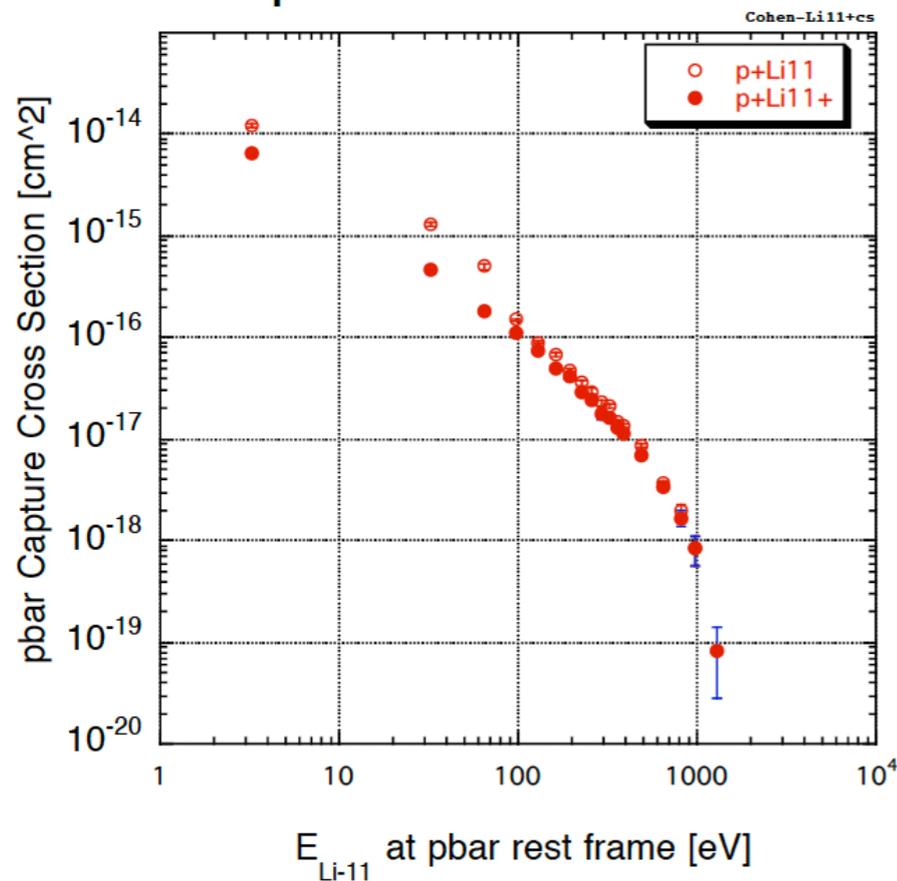
* radioactive nuclei require efficient detection methods

Production of pbar-RI atom



M.Wada, Y.Yamazaki, NIM B214 (2004) 196-200

Capture Cross Section



Cohen PRA69(2004)022501

Yield estimation:

5M pbar (target) + 1 k atom/s Li-II, 10ms cycle

$$\begin{aligned}
 R &= \sigma_{\bar{p}A} N_{\bar{p}} I_{RI} \\
 &= 4 \times 10^{-16} [\text{cm}^2] \cdot 5 \times 10^8 [\text{cm}^{-2}] \cdot 10 [\text{atoms}] \cdot 5 \times 10^3 [1/10\text{ms}] \\
 &= 1 \times 10^{-2} [\bar{p}RI/10\text{ms}]
 \end{aligned}$$

$$\bar{p}^{11}\text{Li}^{2+} : 1 \text{ atom/s}$$

Measurement Methods

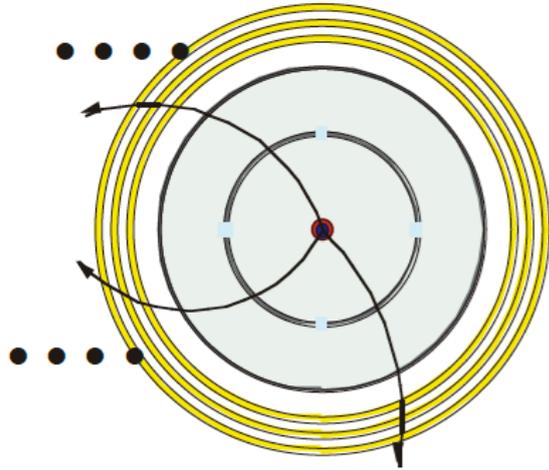
1. Pion's net charge $\bar{p}p:0$ $\bar{p}n:-1$

a) Calorimetric detection of charged pions

- O 1 event distinguishes pbar-p or pbar-n annihilations
- X Detection efficiency must be unity.

b) Comparison of statistical sum of detected pions

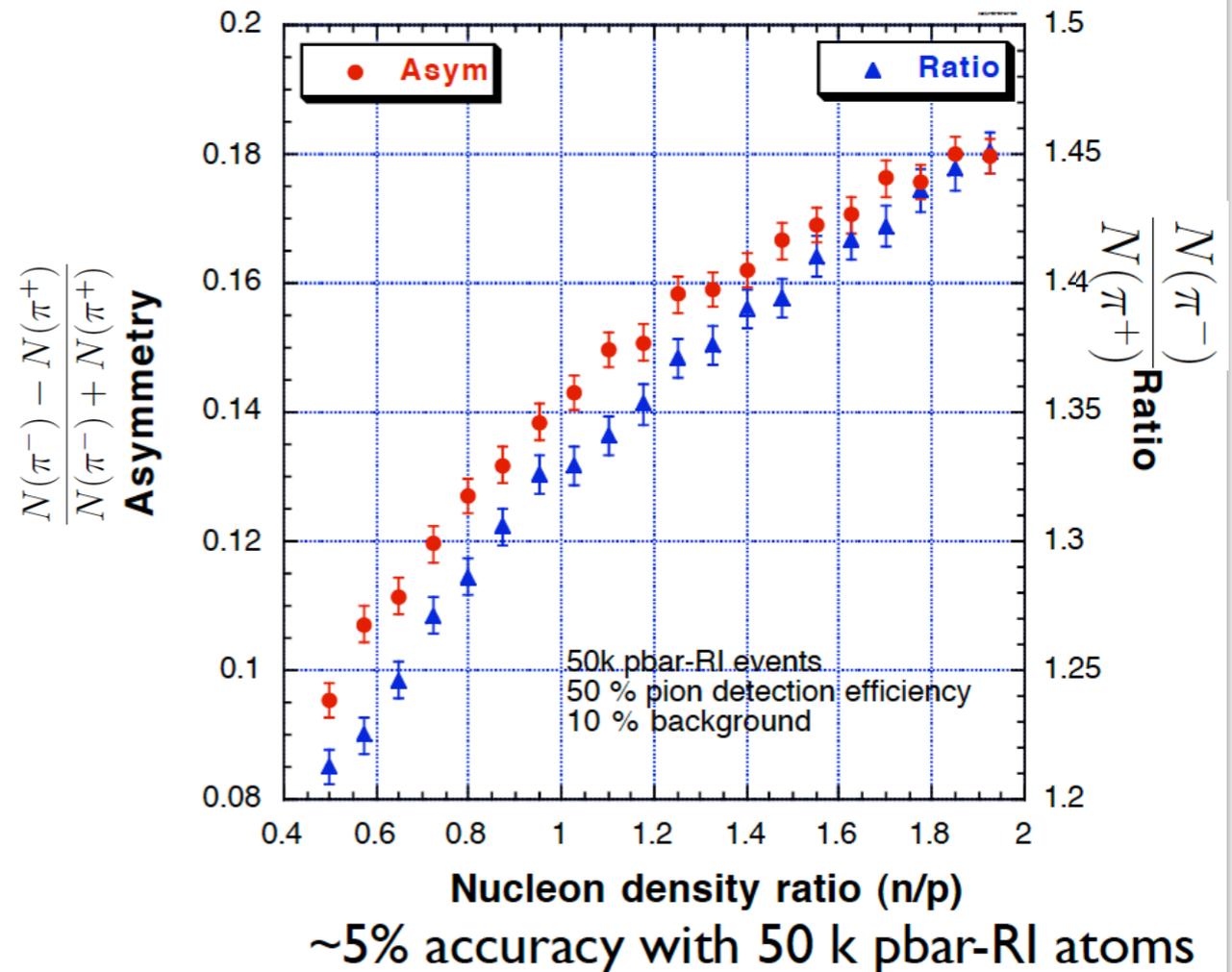
$$\frac{N(\pi^-) - N(\pi^+)}{N(\pi^-) + N(\pi^+)} \text{ or } \frac{N(\pi^-)}{N(\pi^+)} \Rightarrow \frac{N(\bar{p}n)}{N(\bar{p}p)} \Rightarrow \frac{\rho(n)}{\rho(p)} \Big|_{@ \text{ annihilation}}$$



- * Universal method
- * High statistics is needed
- * Background pion should be avoided (e.g. coincidence with recoil nucleus)

		N 11	N 12	N 13	N 14	N 15	N 16	N 17
			0.011s	9.965m	99.634	0.366	7.13s	4.173s
C 8	C 9	C 10	C 11	C 12	C 13	C 14	C 15	C 16
	0.1265s	19.25s	20.39m	98.89	1.11	5730v	2.449s	0.747s
	B 8	B 9	B 10	B 11	B 12	B 13	B 14	B 15
	0.77s	8.5e-19s	19.9	80.1	0.0202s	0.01736s	0.0138s	0.0105s
	Be 7	Be 8	Be 9	Be 10	Be 11	Be 12		Be 14
	3.12d	6.7e-17s	100	1.51e+06s	13.81s	0.0215s		0.00484s
	Li 6	Li 7	Li 8	Li 9	Li 10	Li 11		
	7.5	92.5	0.838s	0.1783s		0.1085s		
He 4		He 6		He 8				
99.9999		0.81s		0.119s				

cold residues are often particle unbound



More detailed statistical treatment including Multiplicity

TABLE 3. Simulated histogram of the multiplicity of charged pions (row) and the total charge per event (column) in a typical experimental condition: $R_{n/p} = 1.0$, $\varepsilon_+ = \varepsilon_- = 0.6$, $\omega_+ = \omega_- = 0.1$, $\lambda_+ = \lambda_- = 0.1$. The number of total annihilation event is 10^5 .

$M \backslash \Sigma_c$	-5	-4	-3	-2	-1	0	+1	+2	+3	+4	
0	0	0	0	0	0	0	0	0	0	0	(11386)*
1	0	0	0	0	17223	0	11233	0	0	0	28456
2	0	0	0	7530	0	21437	0	2844	0	0	31811
3	0	0	1029	0	11901	0	6591	0	179	0	19700
4	0	44	0	1904	0	4394	0	519	0	5	6866
5	1	0	99	0	979	0	451	0	13	0	1543
6	0	2	0	75	0	133	0	14	0	0	224
7	0	0	1	0	7	0	3	0	0	0	11
8	0	0	0	1	0	1	0	0	0	0	2
9	0	0	0	0	1	0	0	0	0	0	1
	1	46	1129	9510	30111	25965	18278	3377	192	5	88612

* unobservable events

ω_-, ω_+ : absorption of π^-, π^+ with residue
 λ_-, λ_+ : charge exchange π^-, π^+ with residue

nuisance effects can be evaluatable by including multiplicity

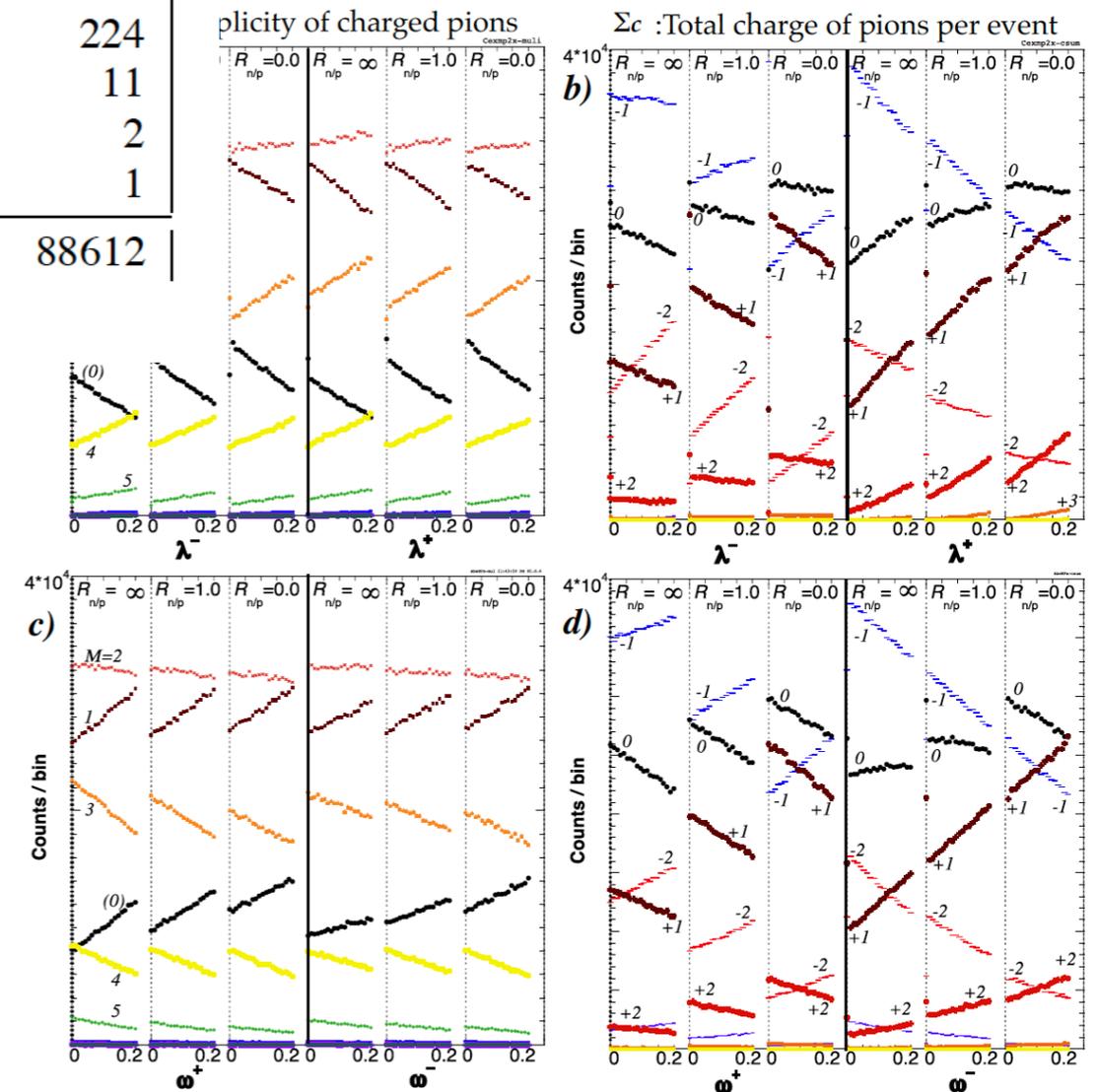
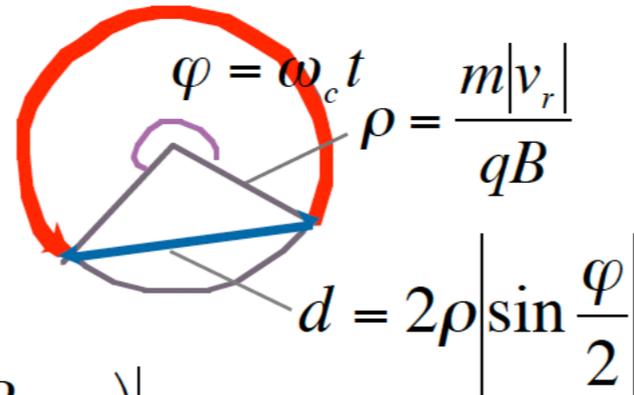
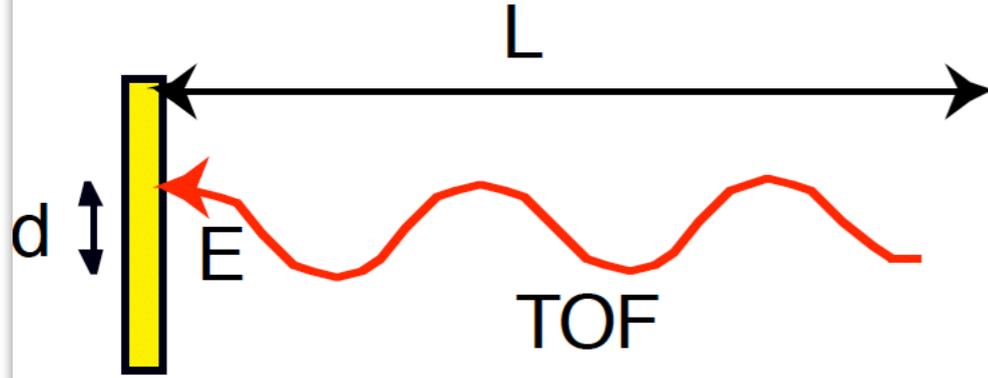
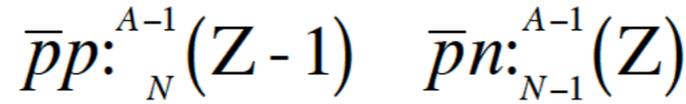


FIGURE 5. Simulation of the dependence of the detected charged pion's multiplicity (left (a, c)), the total charge per event (right (b, d)) on the charge exchange probability λ_{\pm} (top (a, b)) and on the absorption probability ω_{\pm} (bottom (c, d)). For a) and b), λ_- and λ_+ are scanned from 0.0 to 0.2 with $\lambda_+ = 0.1$ for left three scans ($\lambda_- = 0.1$ for right three scans). For c) and d), ω_- and ω_+ are scanned from 0.0 to 0.2 with $\omega_+ = 0.1$ for left three scans ($\omega_- = 0.1$ for right three scans). Other parameters are fixed: $\varepsilon_{\pm} = 0.6$, $\omega_{\pm} = 0.1$ and $\lambda_{\pm} = 0.1$.

Measurement Methods

2. Recoiled Cold Residue



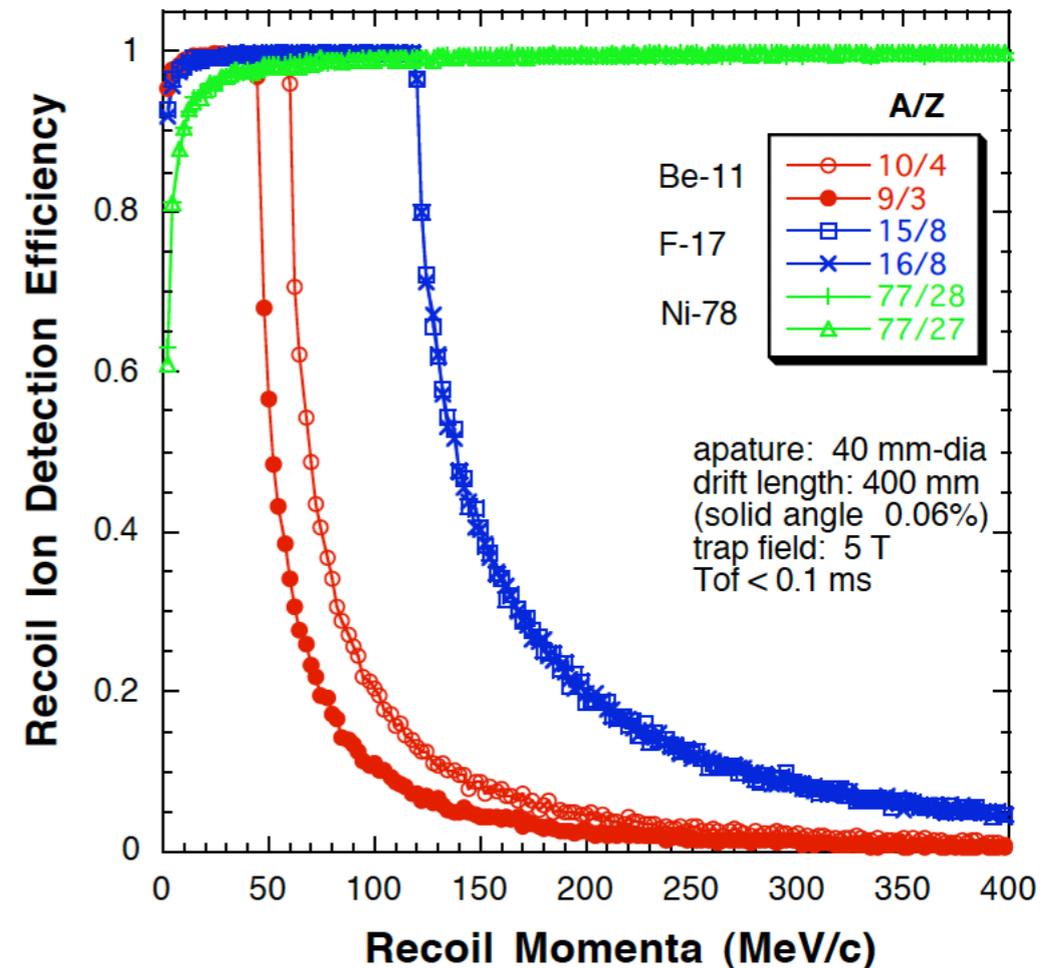
$$d = 2 \frac{m}{qB} \sqrt{\frac{2E}{m} - \left(\frac{L}{tof}\right)^2} \left| \sin\left(\frac{qB}{2m} tof\right) \right|$$

Particle Identification & Momentum measurement

direct deduction of pbar-p, pbar-n event ratio

$$\frac{N\left({}_{N-1}^{A-1}Z\right)}{N\left({}_N^{A-1}(Z-1)\right)} \approx \frac{N(\bar{p}n)}{N(\bar{p}p)}$$

- * Higher sensitivity
- * Ambiguity due to nuclear excitation, re-scattering etc
- * Recoil momentum $\approx \mathbf{P}$ vanished nucleon + other effects



Possible Experimental Locations

1. @CERN AD
ISOLDE (RI-Beam) → L.E. Beam Transport → AD (pbar)

PUMA2 ?

2. @CERN ISOLDE
CERN AD(pbar) Portable Trap → ISOLDE (RI-Beam)



Loading/Unloading ~2 hours

+ 10 hours ?

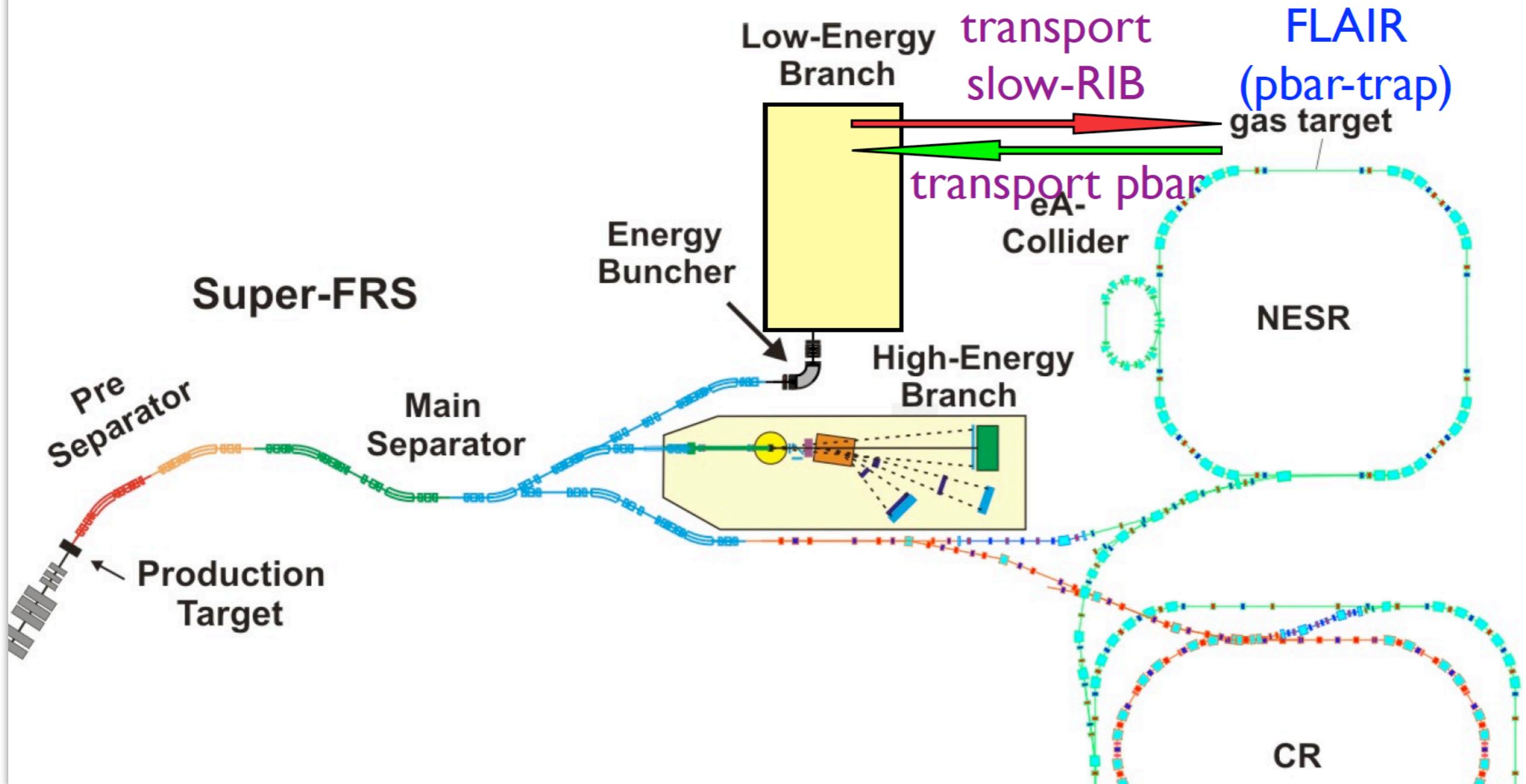


PUMA

3. @RIKEN
CERN AD(pbar) Portable Trap → RIKEN-RIBF (RI-Beam)

4. @GSI FAIR
SuperFRS-LEB (RI-Beam) BTL → FLAIR (pbar)

Experimental Locations @FAIR, GSI



Beam Transport between Super FRS LEB and FLAIR

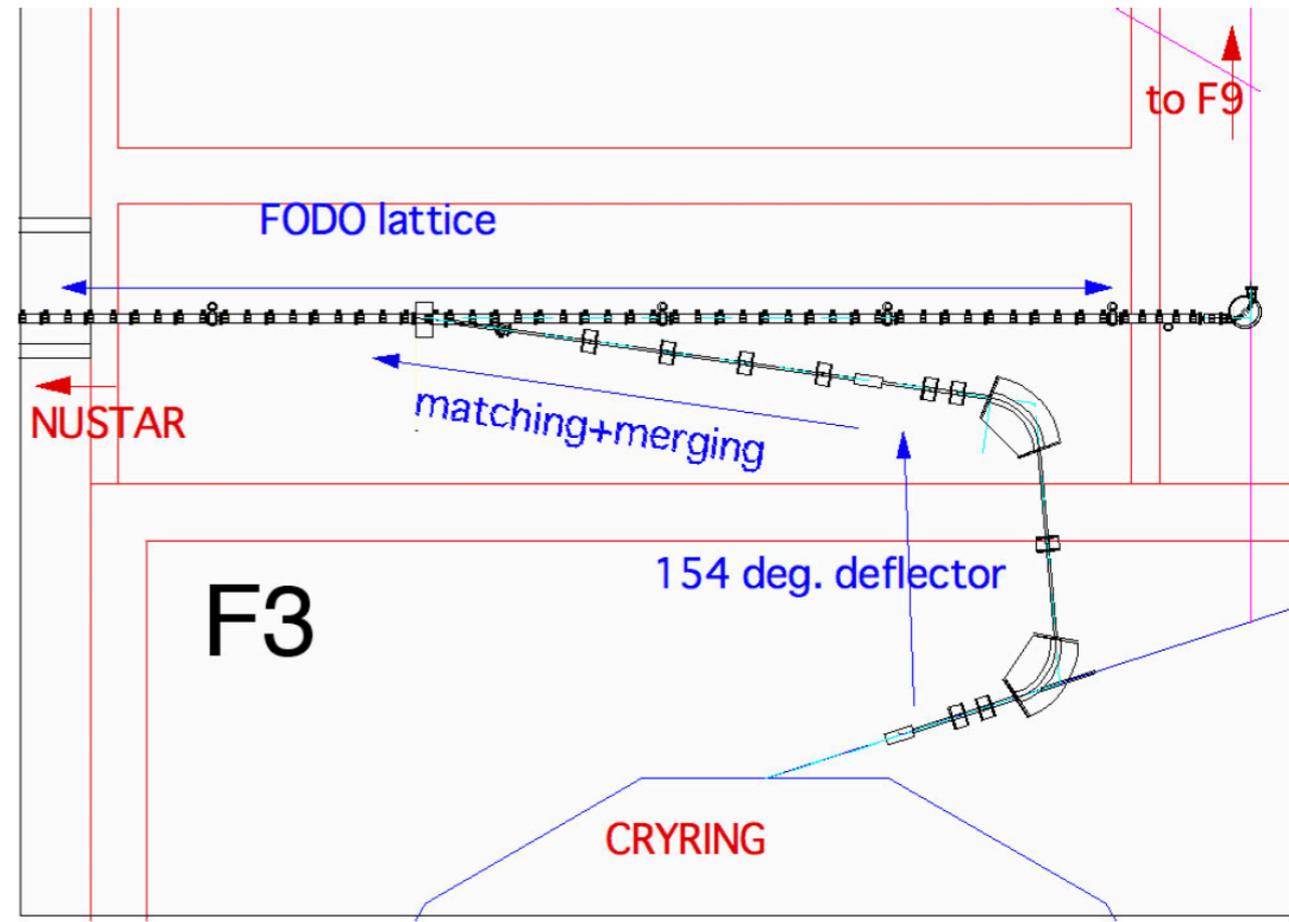
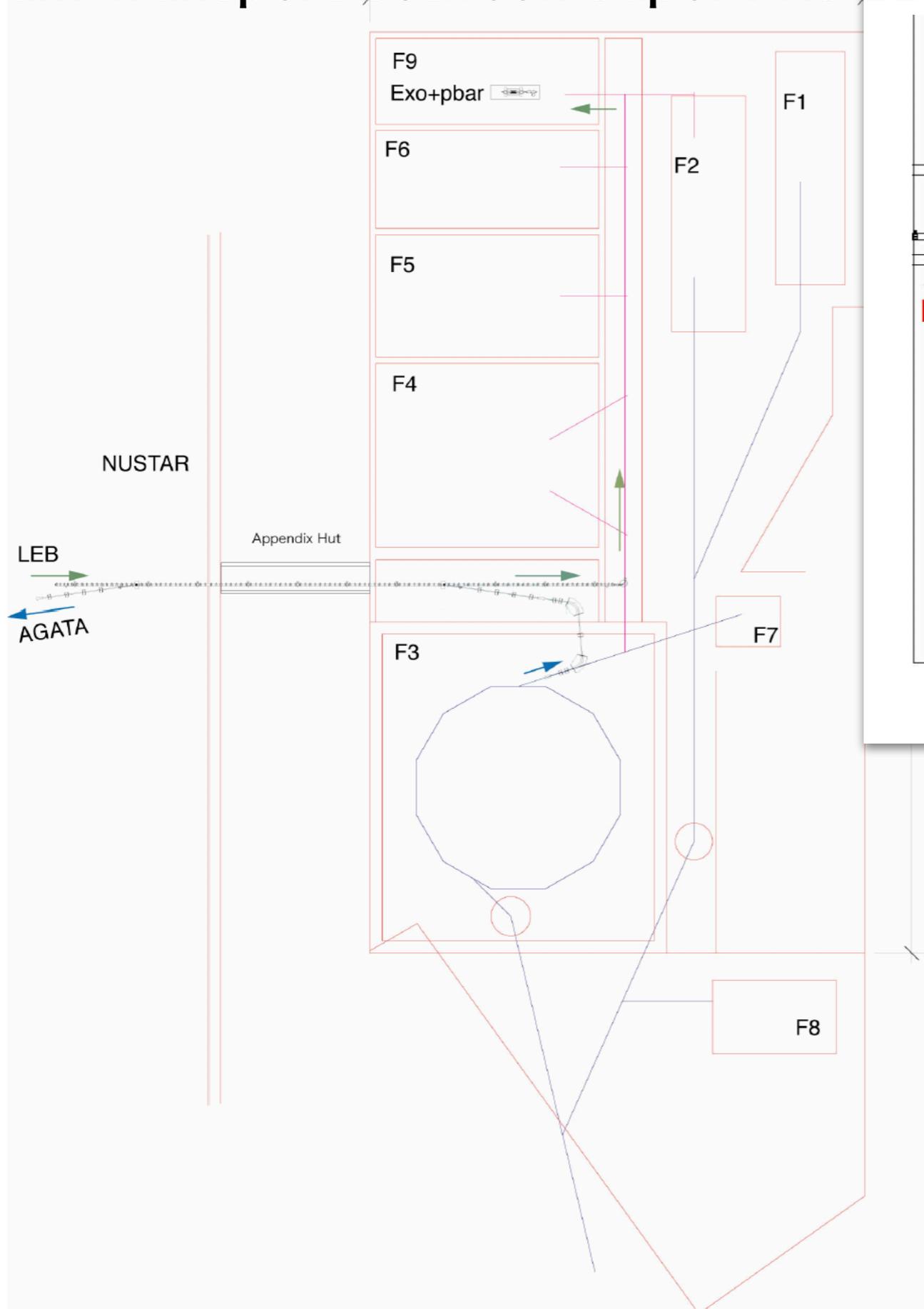


Fig. 89: Core part of LBTL in FLAIR building

GSJ TDR 2005

Fig. 88: Long low-energy beam transport line (LBTL) between NUSTAR and FLAIR

Optics of bi-directional BTL for Heavy Ions and Antiprotons

Table 28: Characteristics of the beams to be transported

	Energy	Emittance	Magnetic Rigidity	Electrostatic Rigidity
Low energy RIB M^{1+}	30 keV	$\sim 50 \pi$ mm mrad	~ 0.2 Tm	60 kV
Medium energy RIB $^{200}\text{Au}^{79+}$	$\sim 6A$ MeV	$< 5 \pi$ mm mrad	~ 0.9 Tm	~ 30 MV

Heavy Ion

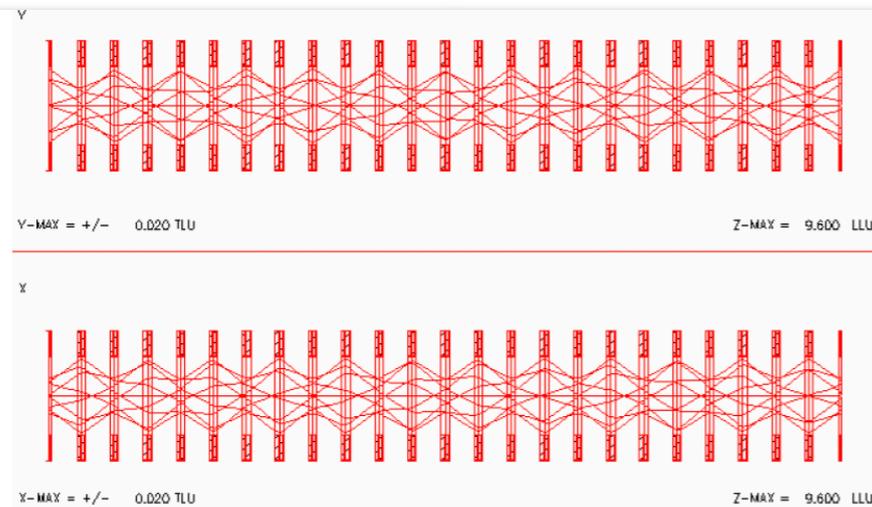


Fig. 90: Part of the beam optics design for low-energy RIB (30 keV, A^+ , 60π mm mrad). Electrostatic quadrupoles with an aperture diameter of 24 mm and a length of 100 mm are placed with a period of 400 mm.

Antiproton

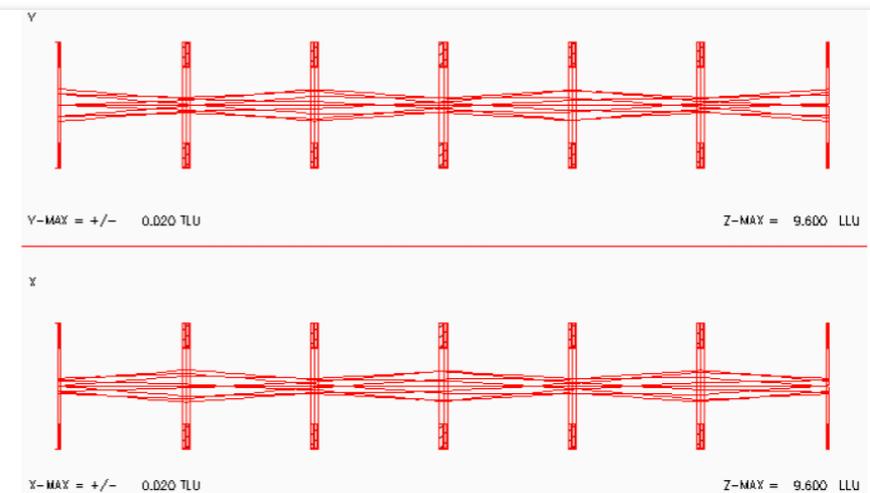


Fig. 91: Part of the beam optics design for medium-energy RIB (6 MeV/u, $A=200$, $Z=79$, Electrostatic Rigidity = 30 MV, 5π mm mrad) from CRYRING. The same periodic structure of electrostatic quadrupoles for low-energy RIB is used, but only one out of 4 elements is activated with a typical voltage of +17 kV while three elements in between are turned off.

Comparison of proposals

Exo+pbar (pbar-RI) TRAP

pbar-A RING

Measurable Quantity	n/p at surface	n/p at rms radius matter rms radii **
Drip-line Nuclides	○ pion detection	× cold residue - unbound
Short-lived Nuclides	○ 10 ms (Intensity > 10 ³ /s)	× beam cooling
Recoil Momentum	△ requires accurate position measurements	○ inverse kinematics
Feasibility	need evaluation, developments	established method

** matter rms radii can be measured with (normal) protons

Yell to

PUMA

Best & Unique Location:

CERN AD + CERN ISOLDE

Strong Collaboration:

Lead by Alexandre

Possible X-ray measurements:

Xenon TPC ?