

Energy conservation in pion production

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Challenges to Transport Theory for Heavy-Ion Collisions
ECT*, Trento, 22 May 2019



Elliptic Flow vs. Pion ratios

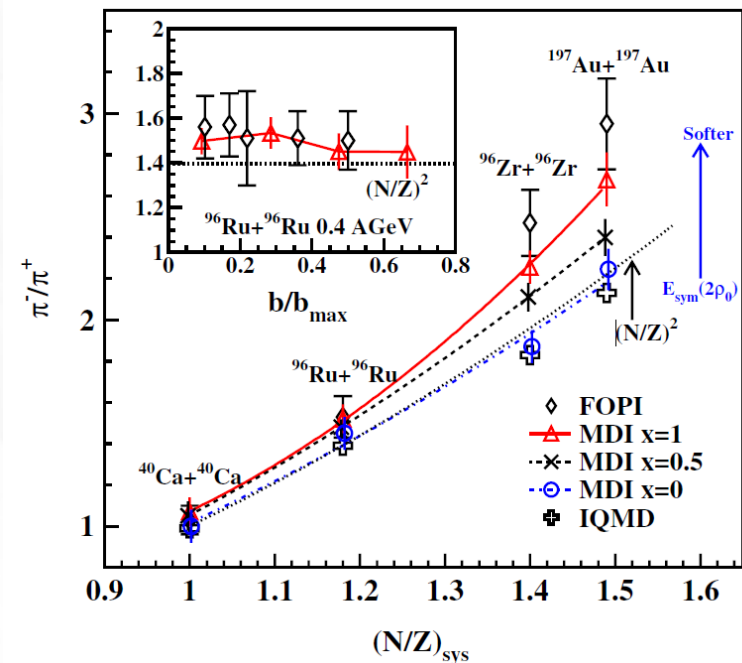
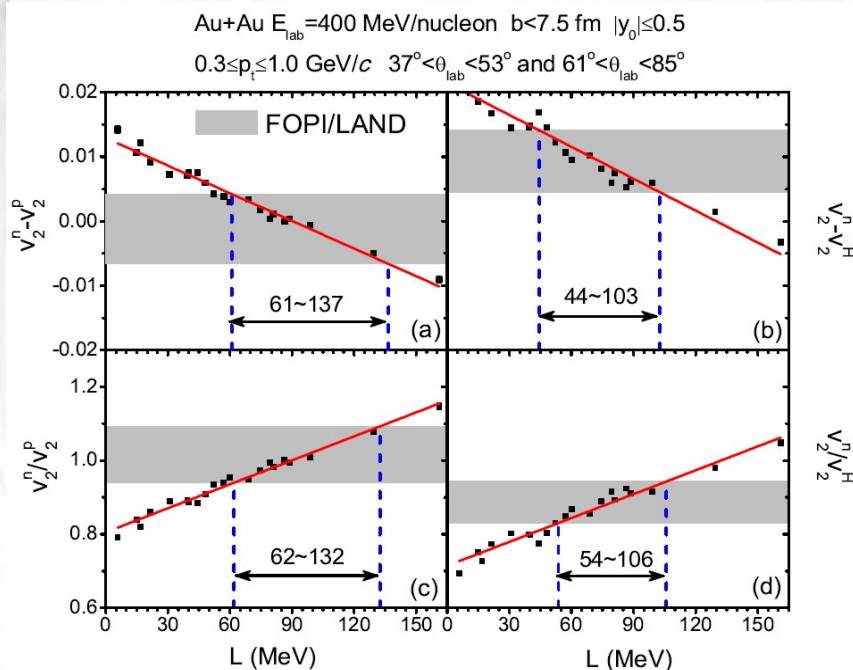
$$\frac{dN}{d\phi} \sim 1 + 2v_1 \cos \phi + 2v_2 \cos 2\phi$$

UrQMD - Y. Wang et al. PRC 89, 044603 (2014)

Isobar model (no symmetry potential)

$$\pi^-/\pi^+ = (5N^2 + NZ)/(5Z^2 + NZ)$$

IBUU - Z.Xiao et al. PRL 102,062502 (2009)



TuQMD - linear/moderately stiff
 M.D. Cozma et al. PRC 88, 044912 (2013)
 UrQMD - linear
 P. Russotto et al. PLB 697, 491 (2011)
 IBUU - linear/moderately stiff
 G.-C. Yong private communication

ImQMD - stiff
 Z.-Q. Feng et al., PLB 683, 140 (2010)
 Boltzmann-Langevin - super-soft
 W.-J. Xie et al., PLB 718, 1510 (2013)
 TuQMD (VEC) - super-soft
 M.D. Cozma, PLB 753, 166 (2016)
 pBUU - no sensitivity to SE
 J. Hong et al. PRC 90, 024605 (2014)

Energy Conservation (QMD type models)

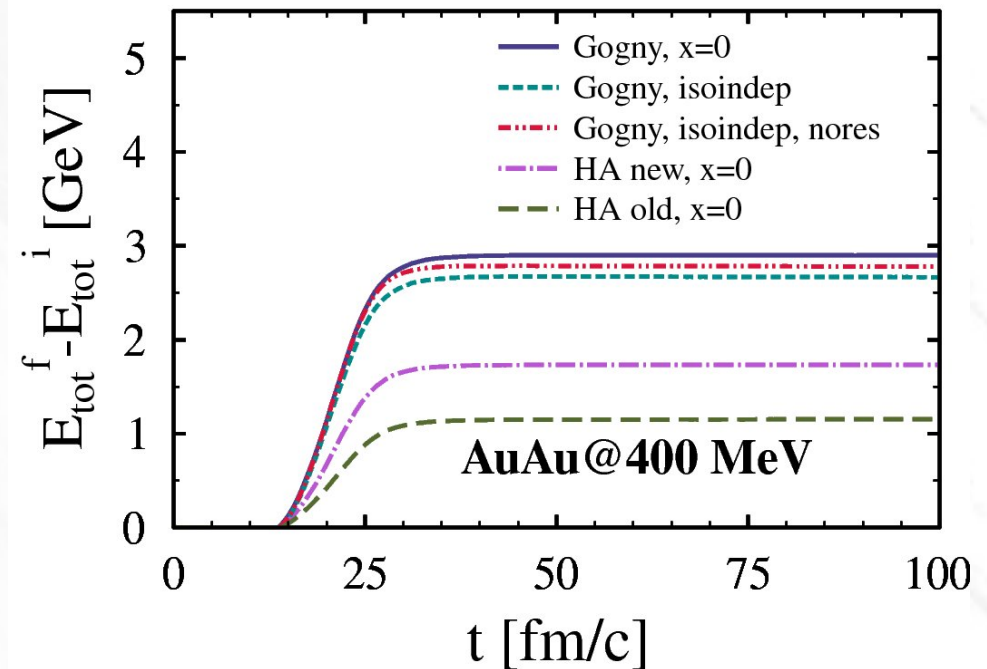
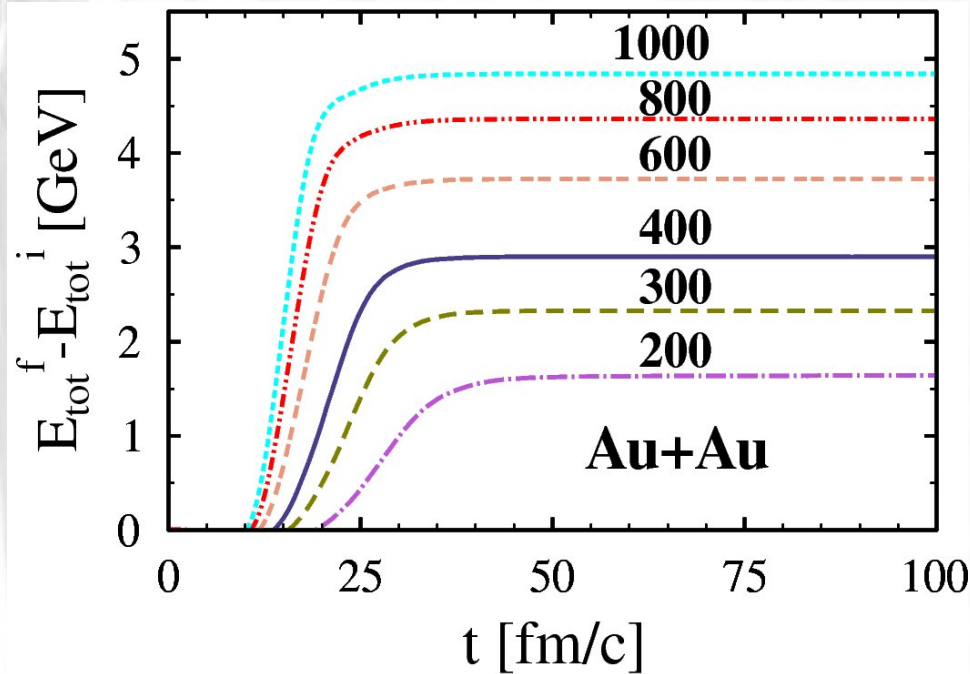
80's transport models – total energy conserved on average (potentials dependent only on density)

Collective phenomena – momentum dependence

Isospin effects – isospin asymmetry dependence

Violation of total energy conservation

Determination of final state kinematics of 2-body collisions neglects medium effects



Energy conservation (in-medium)

$$\sqrt{p_1^2 + m_1^2} + U(p_1) + \sqrt{p_2^2 + m_2^2} + U(p_2) = \sqrt{p_1'^2 + m_1'^2} + U(p_1') + \sqrt{p_2'^2 + m_2'^2} + U(p_2')$$

- **rarely considered** in transport models below 1 AGeV, with a few exceptions:
G. Ferini et al. PRL 97, 202301 (2006), C.Fuchs et al. PRC 55, 411 (1997),
T. Song, C.M. Ko PRC 91, 014901 (2015)
- **Ansatz** for the isospin 3/2 resonance potential motivated by decay channel
– see also S.A. Bass et al., PRC 51, 3343 (1995)
- **imposed** in the CM of the colliding nuclei (not in the Eckart frame)
- **reactions**: $NN \leftrightarrow NR$, $R \leftrightarrow N\pi$ ($R \leftrightarrow N\pi\pi$ not corrected)

$$\begin{aligned} U(\Delta^{++}) &= U^p \\ U(\Delta^+) &= \frac{2}{3}U^p + \frac{1}{3}U^n \\ U(\Delta^0) &= \frac{1}{3}U^p + \frac{2}{3}U^n \\ U(\Delta^-) &= U^n \end{aligned}$$

B.-A. Li, NPA 708 (365) 2002

Reaction	$\Delta U = U^f - U^i$	Effect
$nn \rightarrow p\Delta^-$	$U^p - U^n < 0$	more π^-
$nn \rightarrow n\Delta^0$	$1/3(U^p - U^n) < 0$	more π^-, π^0
$np \rightarrow p\Delta^0$	$1/3(U^p - U^n) < 0$	more π^-, π^0
$np \rightarrow n\Delta^+$	$1/3(U^n - U^p) > 0$	less π^+, π^0
$pp \rightarrow p\Delta^+$	$1/3(U^n - U^p) > 0$	less π^+, π^0
$pp \rightarrow n\Delta^{++}$	$U^n - U^p > 0$	less π^+

Transport Model

Quantum Molecular Dynamics (TuQMD): HIC 0.1-2.0 GeV/A

previously applied to study:

- dilepton emission in HIC: [K.Shekter, PRC 68, 014904 \(2003\)](#); [D. Cozma, PLB640,170 \(2006\)](#); [E.Santini PRC78,03410 \(2008\)](#)
- EoS of symmetric nuclear matter: [C. Fuchs, PRL 86, 1974 \(2001\)](#); [Z.Wang NPA 645, 177 \(1999\)](#)
- In-medium effects and HIC dynamics: [C. Fuchs, NPA 626,987 \(1997\)](#); [U. Maheswari NPA 628,669 \(1998\)](#)

upgrades implemented at IFIN-HH (Bucharest):

- various parametrizations for the EoS: optical potential, symmetry energy [PRC 88, 044912 \(2013\)](#)
- various parametrizations for elastic cross-sections (also in medium ones) [PLB 700, 139 \(2011\)](#)
- threshold effects for baryon resonance & π meson emission/absorption [PLB 753, 166 \(2016\)](#)
- pion optical potential [PRC 95, 014601 \(2017\)](#)
- **planned:** threshold effects for reactions involving strangeness degrees of freedom

Pion production: two step process

- resonance excitation in baryon-baryon collisions
parametrization of the OBE model of
[S.Huber et al., NPA 573, 587 \(1994\)](#)
- resonance decay:
Breit-Wigner shape of the resonance spectral function;
parameters -> [K. Shekhter, PRC 68, 014904 \(2003\)](#)
decay channels:

$$\begin{aligned} R &\rightarrow N\pi, & R &\rightarrow N\pi\pi \\ R &\rightarrow \Delta(1232)\pi, & R &\rightarrow N(1440)\pi \end{aligned}$$

Pion absorption:

- resonance model (all 4* resonances below 2 GeV)
[K. Shekhter, PRC 68, 014904 \(2003\)](#)

Isospin dependence of EoS

momentum dependent – generalization of the Gogny interaction:

MDI potential

Das, Das Gupta, Gale, Li PRC67, 034611 (2003)

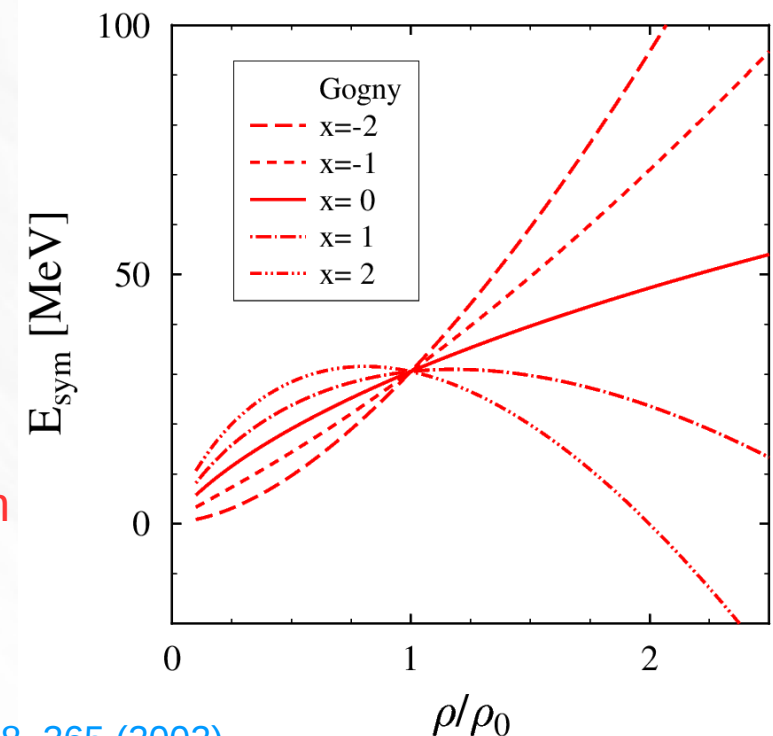
$$\frac{E}{N}(\rho, \beta, x) = \frac{1}{2} A_1 u + \frac{1}{2} A_2(x) u \beta^2 + \frac{B}{\sigma+1} (1 - x \beta^2) + \frac{1}{u \rho_0^2} \sum_{\tau, \tau'} C_{\tau\tau'} \int \int d^3 p d^3 p' \frac{f_{\tau}(p, p') f_{\tau'}(p, p')}{1 + (\vec{p} - \vec{p}')^2 / \Lambda^2}$$

$$A_2(x) = A_2^0 + \frac{2x B}{\sigma+1} \bar{u}^{\sigma-1}$$

$$u = \frac{\rho}{\rho_0}$$

x	L _{sym} [MeV]	K _{sym} [MeV]
-2	152	418
-1	106	127
0	61	-163
1	15	-454
2	-301	-745

$$S(\rho) = J + \frac{L}{3} \frac{\rho - \rho_0}{\rho_0} + \frac{K_{sym}}{18} \frac{(\rho - \rho_0)^2}{\rho_0^2}$$



resonance potential: isovector component unknown

$$V(\Delta^{++}) = V_n = V_s + \frac{3}{2} V_v$$

$$V(\Delta^+) = \frac{2}{3} V_n + \frac{1}{3} V_p = V_s + \frac{1}{2} V_v$$

$$V(\Delta^0) = \frac{1}{3} V_n + \frac{2}{3} V_p = V_s - \frac{1}{2} V_v$$

$$V(\Delta^-) = V_p = V_s - \frac{3}{2} V_v$$

$$V_s \equiv \frac{1}{2} (V_n + V_p)$$

$$\delta \equiv \frac{1}{3} (V_n - V_p)$$

$$V_v = \delta$$

B.-A. Li, NPA 708, 365 (2002)

! In the following: V_v expressed in units of $[\delta]$

Approximations

PLB 753, 166 (2016)

Elastic scattering:

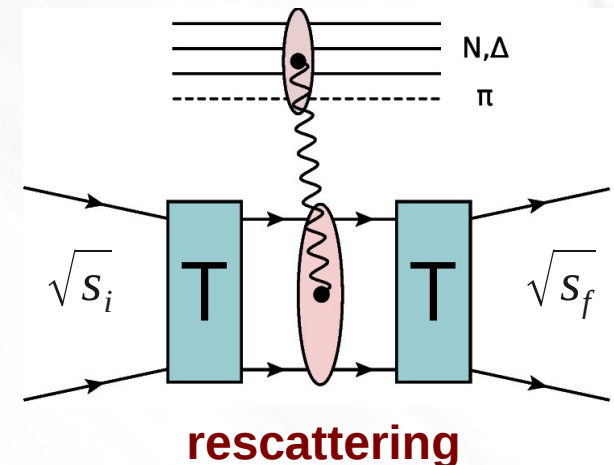
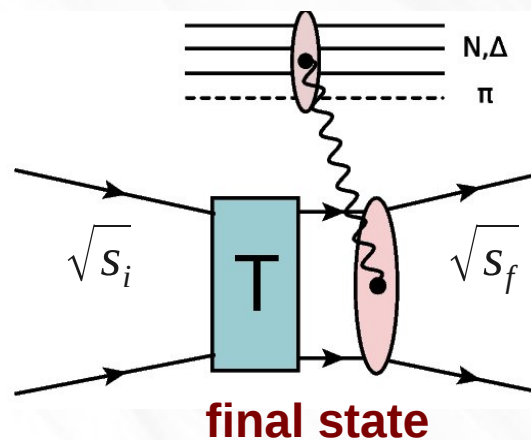
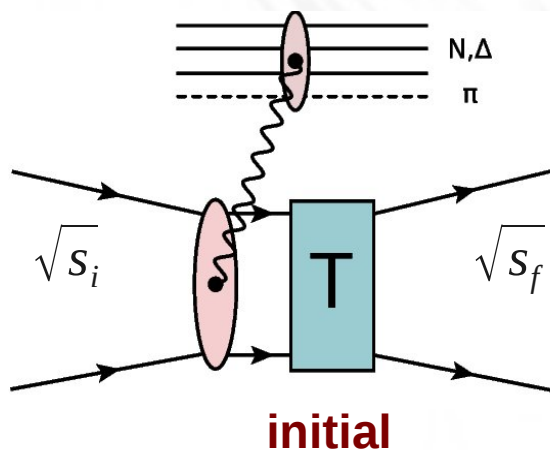
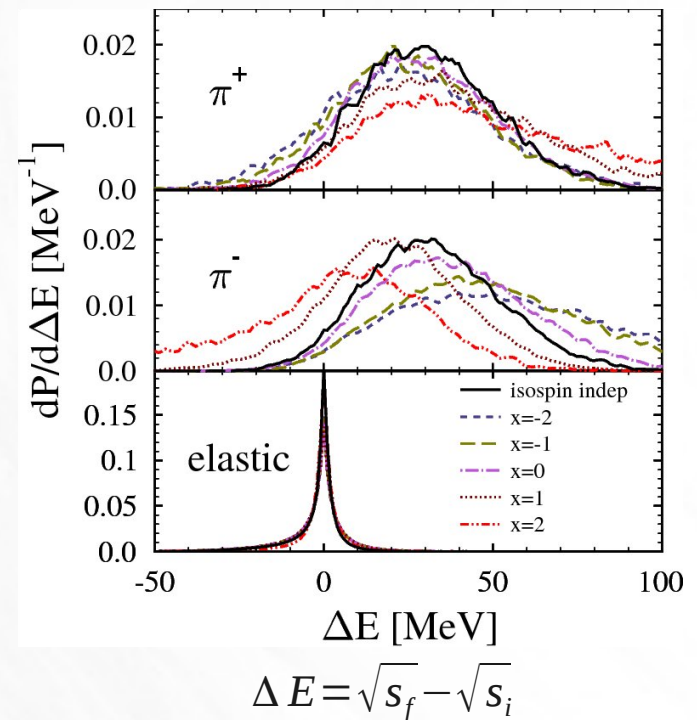
$$\sqrt{s_f} \approx \sqrt{s_i}$$

$$\sqrt{s^*} = 0.5(\sqrt{s_f} + \sqrt{s_i})$$

Resonance excitation:

$$\sqrt{s_f} - \sqrt{s_i} \approx 25 \text{ MeV}$$

$$\sqrt{s^*} = \sqrt{s_f}$$



Different “scenarios”

VEC – vacuum energy conservation constraint

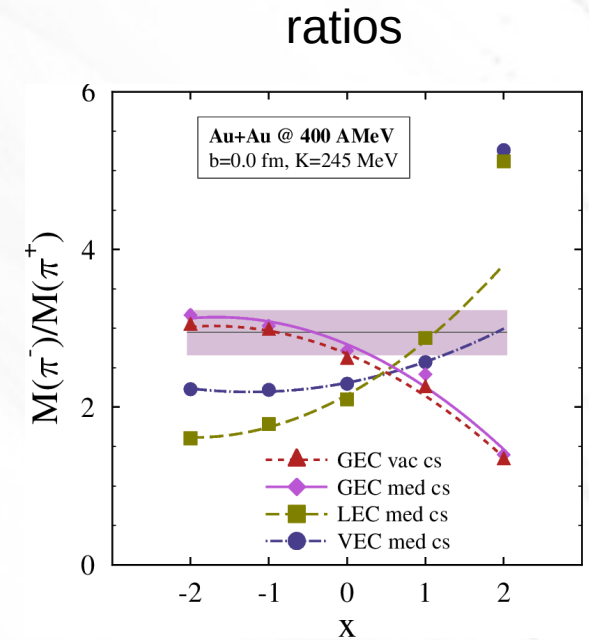
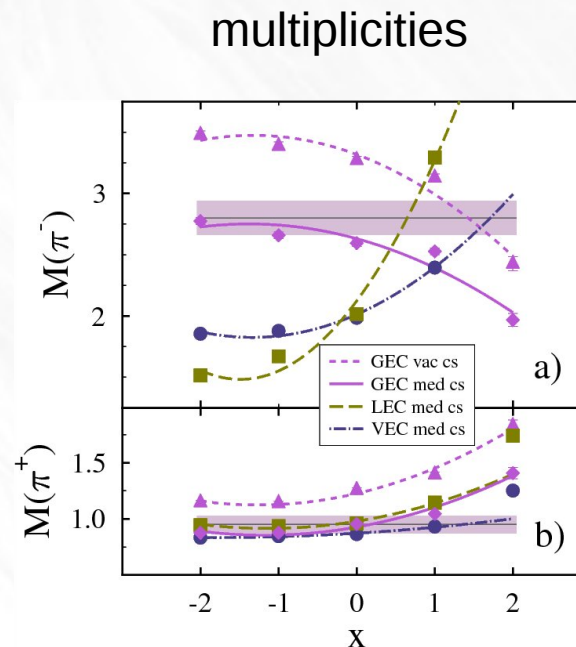
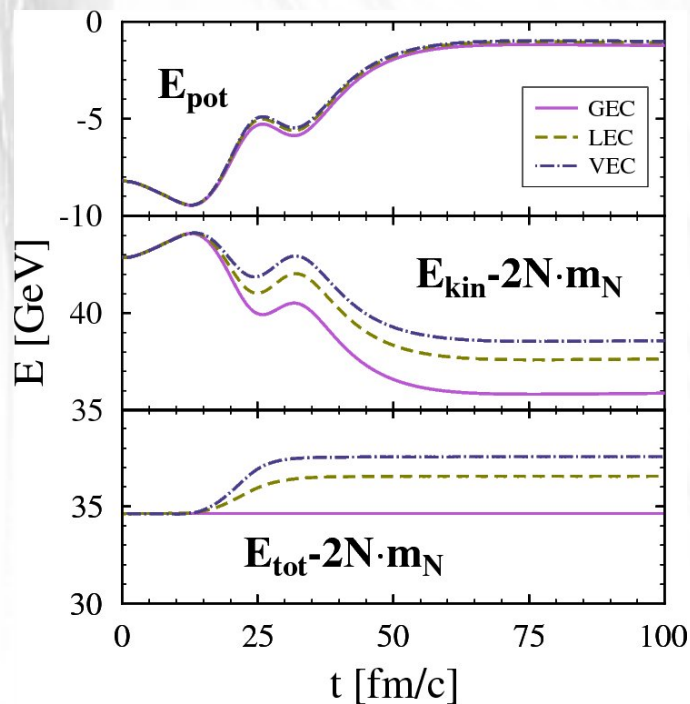
LEC - “local” energy conservation – limited impact on multiplicities and ratios

GEC- “global” energy conservation – conserve energy of the entire system

-in-medium cross-sections for the inelastic channels

$$\sigma^{*(12 \rightarrow 23)} = \left[\frac{m_1^* m_2^* m_3^* m_4^*}{m_1 m_2 m_3 m_4} \right]^{1/2} \sigma^{(12 \rightarrow 34)}$$

AuAu@400 MeV, b=0 fm



Experimental data: [W. Reisdorf et al. \(FOPI\) NPA 848, 366 \(2010\)](#)

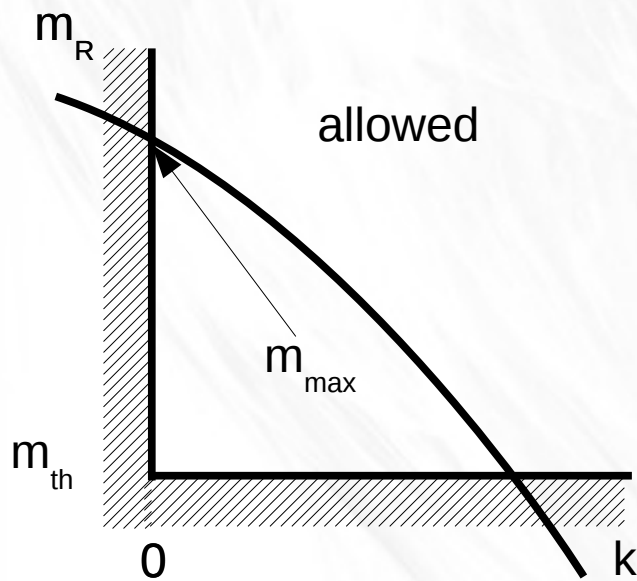
softer

Final state kinematics

final state phase space in NN->NR

VEC

$$\sqrt{s} = \sqrt{m_N^2 + k^2} + \sqrt{m_R^2 + k^2}$$



$$m_{th} \leq m_R \leq m_{max}$$

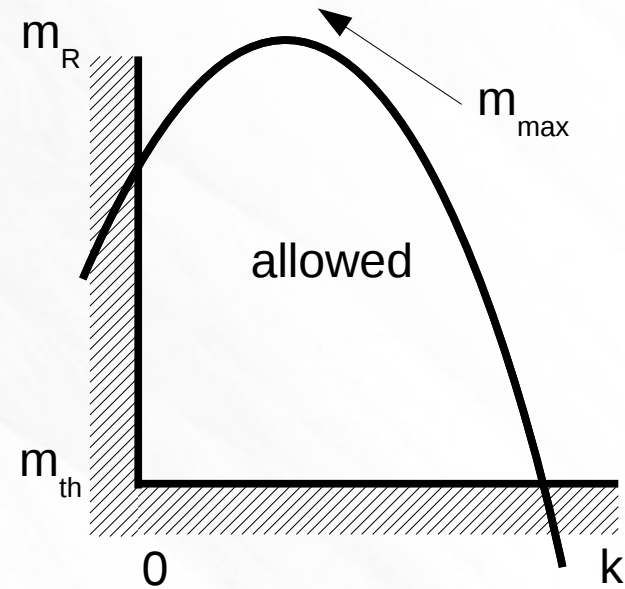
$$0 \leq \theta \leq \pi \quad \text{consistent with } \frac{d\sigma}{d\Omega}$$

$$0 \leq \Phi \leq 2\pi$$

LEC, GEC

$$\sqrt{s_{fin}} = \sqrt{m_N^2 + k^2} + \sqrt{m_R^2 + k^2}$$

$$\sqrt{s_{fin}} = \sqrt{s_{ini}} + \delta \sqrt{s}$$



$$m_{th} \leq m_R \leq m_{max}$$

$$\theta_{min} \leq \theta \leq \theta_{max}$$

$$\Phi_{min} \leq \Phi \leq \Phi_{max}$$

$$m_{th} \leq m_R \leq m_R(k=0)$$

approximate to vacuum case

Detailed Balance (GEC)

Resonance absorption: detailed balance

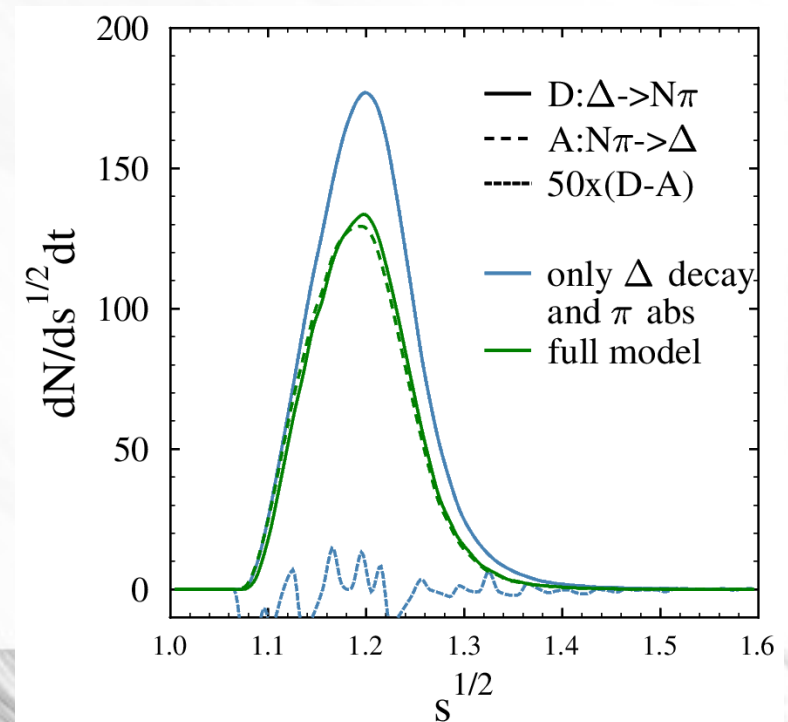
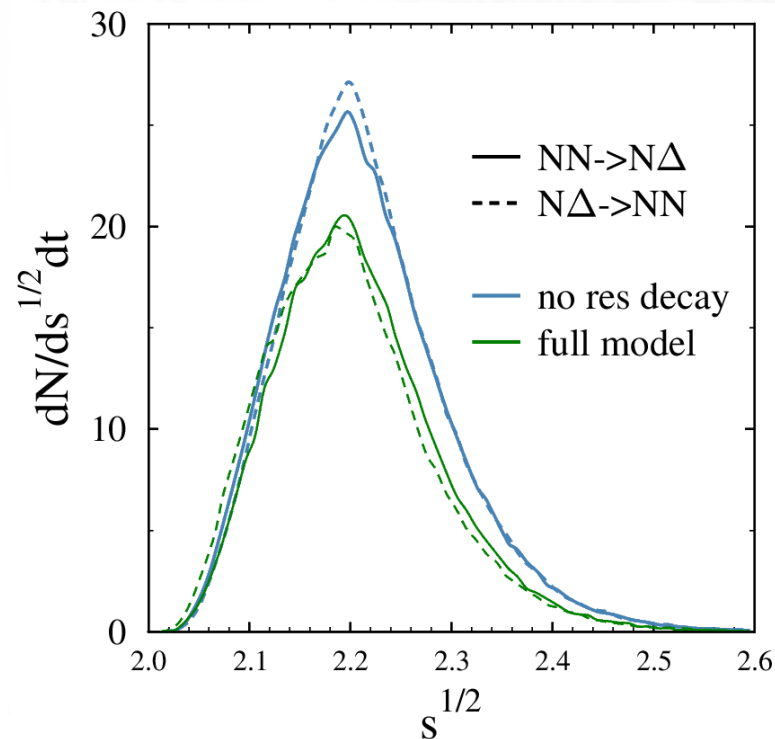
P. Danielewicz et al., NPA 533, 712 (1992)

$$\sigma^{NR \rightarrow NN} = \frac{1}{4} \frac{m_R p_{NN}^2}{p_{NR}} \sigma^{NN \rightarrow NR} \times \left[\frac{1}{4\pi} \int d\Omega \frac{1}{2\pi} \int_{m_N+m_\pi}^{M_{\max}(\theta, \phi)} dM M p'_{NR} A_R(M) \right]^{-1}$$

$\sqrt{s_i} \rightarrow p_{NR}, \quad \sqrt{s_f} \rightarrow p_{NN}$

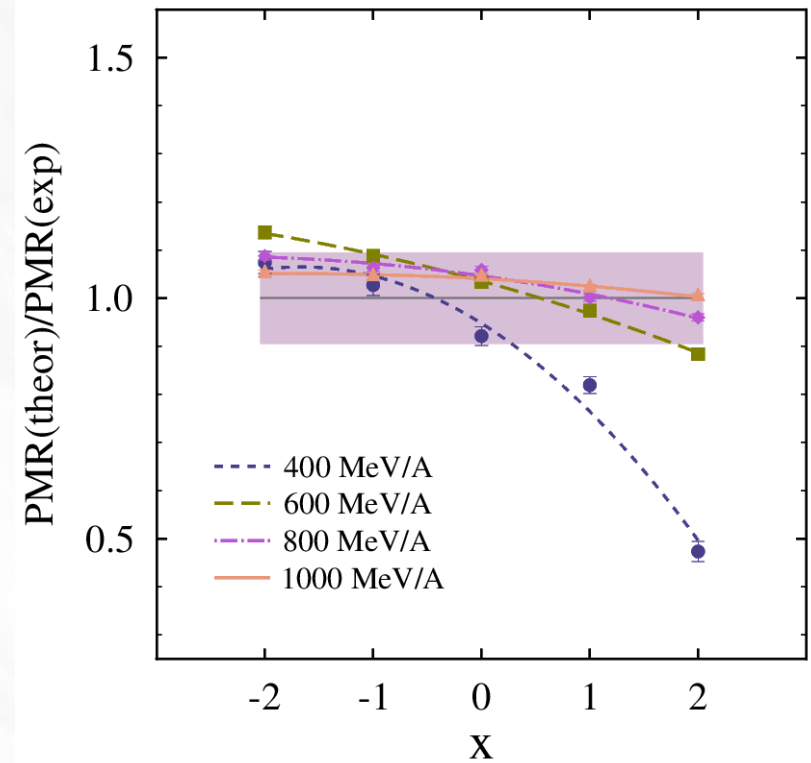
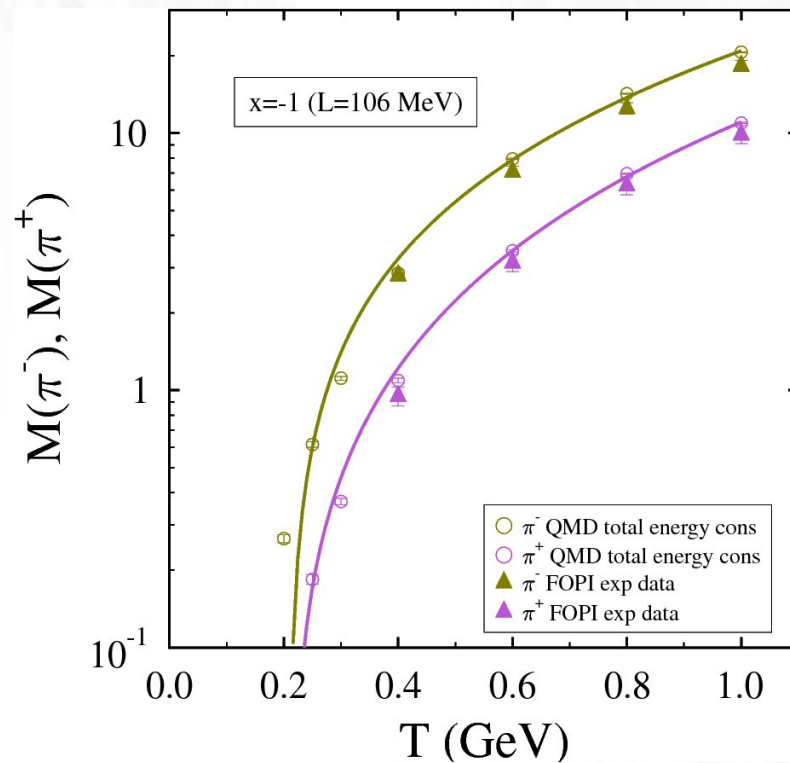
average over solid angle not considered in [PLB 753, 166 \(2016\)](#):

Box calculation: $\beta=0.0$, $T=60$ MeV



Energy dependence

FOPI data: W.Reisdorf et al., NPA 781, 459 (2007)



Impact of the $\Delta(1232)$ potential

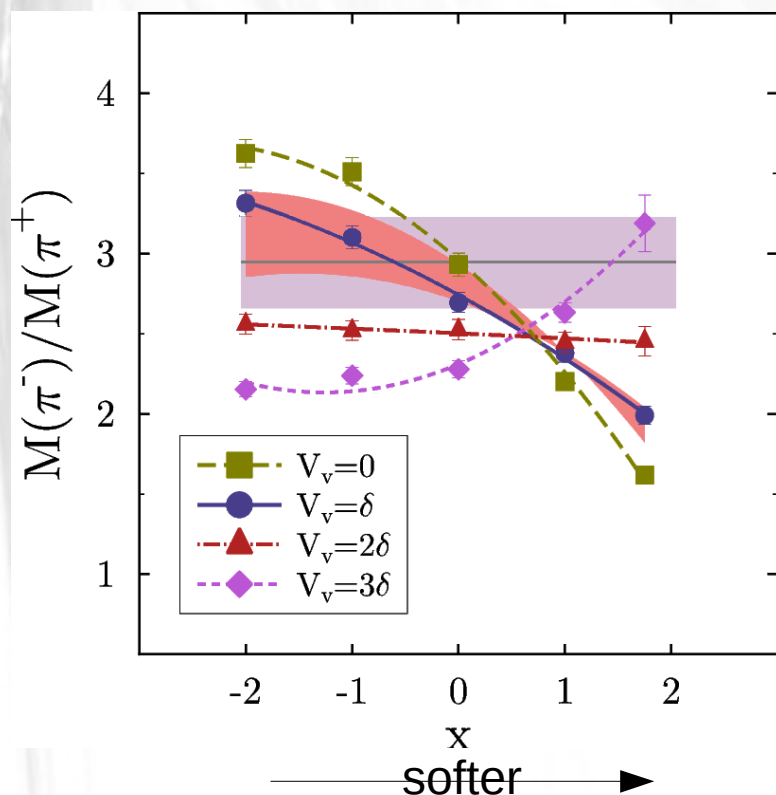
Phenomenology – inclusive electron nucleus scattering (He,C,Fe) **attractive**
 - Δ -nucleus potential deeper than the nucleon-nucleus potential

O'Connell, Sealock PRC 42, 2290 (1990)

Ab initio calculations – Argonne v_{28} interaction (BBG)=> **repulsive** Δ potential

Baldo, Ferreira NPA 569, 645 (1994)
 Malfliet, de Jong PRC 46, 2567 (1992)

- 3D reduction of Bethe-Salpeter equation similar (DB)
 - strong dominant repulsive contribution from T=2 sector



$$V(\Delta^{++}) = V_s + \frac{3}{2} V_v$$

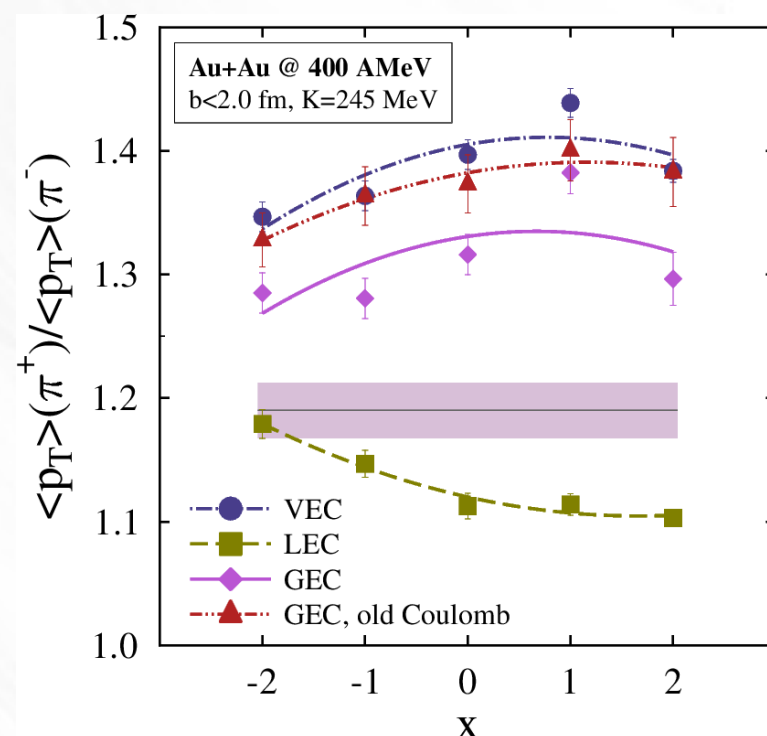
$$V(\Delta^+) = V_s + \frac{1}{2} V_v$$

$$V(\Delta^0) = V_s - \frac{1}{2} V_v$$

$$V(\Delta^-) = V_s - \frac{3}{2} V_v$$

$$V_s = \frac{1}{2}(V_n + V_p)$$

$$\delta = \frac{1}{3}(V_n - V_p)$$



DC, PRC 95, 14601 (2017)

M.D. Cozma, PLB 753, 166 (2016)

Exp. data: W. Reisdorf et al., NPA 781, 459 (2007)

Pion optical potential

Sources:

Theoretical models: Effective Hadronic Models

J. Nieves et al., NPA 554, 509; 554 (1993)
M. Doring et al., PRC 77, 024602 (2008)

pion self-energy starting from basic interaction terms πNN , $\pi N\Delta$, $\pi NN^*(1440)$

Chiral Perturbation Theory

N. Kaiser et al., PLB 512, 283 (2001)
C. Baru et al., NPA 872, 69 (2011)
W. Weise, Acta. Phys. Pol, B31, 2715 (2000)

Experiment: Pionic Atoms

T. Yamazaki et al., Phys. Rep. 514, 1 (2012)
E. Friedman, A. Gal, Phys. Rep. 452, 89 (2007)

Pion-Nucleus Scattering

R. Seki et al, PRC 27, 2799; 2817 (1983)

Ericson-Ericson parametrization (M. Ericson et al. Ann. Phys. 36, 323 (1966))

$$V_{\text{opt}}(r) = \frac{2\pi}{\mu} \left[-q(r) + \vec{\nabla} \frac{\alpha(r)}{1 + 4/3\pi\lambda\alpha(r)} \vec{\nabla} \right]$$

$$q(r) = \epsilon_1 (\bar{b}_0 \rho + \bar{b}_1 \beta \rho) + \epsilon_2 B_0 \rho^2$$

$$\alpha(r) = \epsilon_1^{-1} (c_0 \rho + c_1 \beta \rho) + \epsilon_2^{-1} (C_0 \rho^2 + C_1 \beta \rho^2)$$

$$\epsilon_1 = 1 + m_\pi / m_N \quad \epsilon_2 = 1 + 2m_\pi / m_N$$

	$\bar{b}_0 [m_\pi^{-1}]$	$\bar{b}_1 [m_\pi^{-1}]$	$\text{Re } B_0 [m_\pi^{-4}]$	$\text{Im } B_0 [m_\pi^{-4}]$	λ	$c_0 [m_\pi^{-3}]$	$c_0 [m_\pi^{-3}]$	$\text{Re } C_0 [m_\pi^{-6}]$	$\text{Im } C_0 [m_\pi^{-6}]$
SM-1	-0.0283	-0.120	0.0	0.042	1	0.223	0.250	0.0	0.10
SM-2	0.030	-0.143	-0.150	0.046	1	0.210	0.180	0.11	0.09
Batty-1	-0.017	-0.130	-0.048	0.0475	1	0.255	0.170	0.0	0.09
Batty-2	-0.023	-0.085	-0.021	0.049	1	0.210	0.089	0.118	0.058
Konijn-2	0.025	-0.094	-0.265	0.0546	1	0.273	0.184	-0.140	0.105

subset from K. Itahashi et al., PRC 62, 025202 (2000)

Symmetry Energy Constraints

Au+Au@400 MeV/A, $b < 2.0$ fm

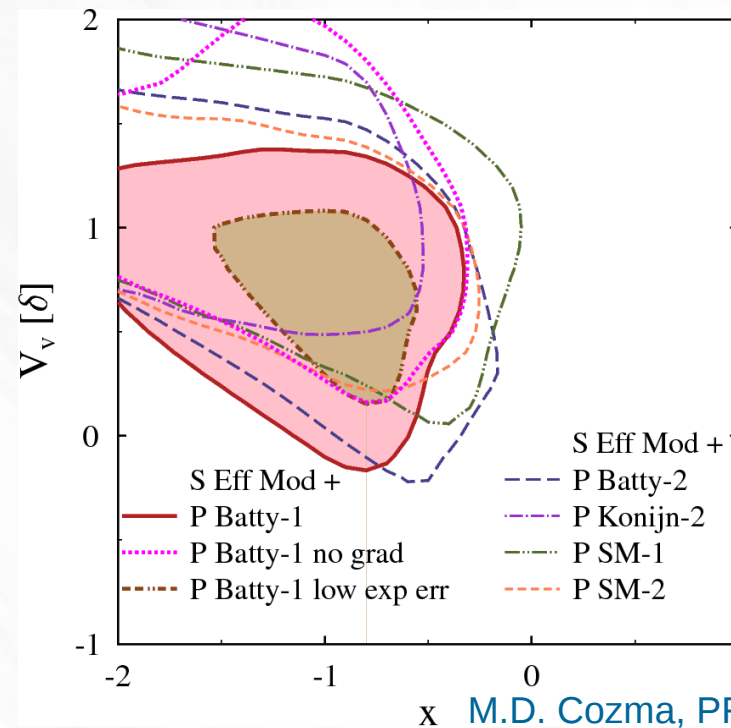
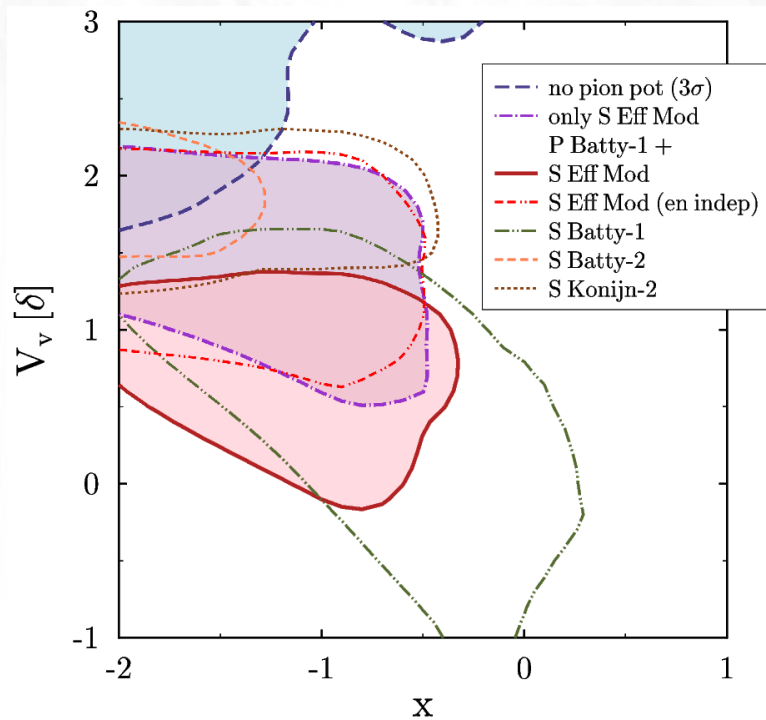
S-wave impact

P-wave: Batty-1

1σ CL contour plots

P-wave impact

S-wave: effective model



M.D. Cozma, PRC 95, 14601 (2017)

No pion potential – unable to describe exp. PMR and PAPTR simultaneously

S-wave - energy dependence impacts V_v
 - density dependence – impact both x (60 MeV on L) and V_v

P-wave - moderate impact on both x (25 MeV on L) and V_v

Summary / Conclusions

- QMD transport models:**
- momentum/isospin dependent interactions lead to total energy conservation violation (scattering term)
 - **solution:** account for potential energies when determining final state kinematics in 2-body collisions
 - **threshold shifts** induced by imposing energy conservation
 - **local energy conservation** (roughly 40% reduction of the magnitude of energy conservation violation)
 - **global energy conservation:** does the job by definition

- Open issues:**
- no momentum transfer
 - only applicable in regions where cross-sections depend strongly on invariant mass
 - not a Lorentz covariant approach (+retardation effects may not be negligible)

- Accomplishments:**
- description of experimental pion multiplicities in HIC in the 0.4-1.0 AGeV range
 - compatible constraints for symmetry energy stiffness from pionic and elliptic flow observables

S-wave pion potential

	$b_0[m_\pi^{-1}]$	$b_1[m_\pi^{-1}]$
Exp	-0.0001 ± 0.0021	-0.0885 ± 0.0021
ChPT	0.0075 ± 0.0031	-0.0861 ± 0.0009
WT	0.0	-0.0790

Energy dependence

inferred from exp. pion-nucleus scattering

$$\bar{b}_0^{eff} = \bar{b}_0 + \rho^{eff} \text{Re } B_0$$

R. Seki et al., PRC 27, 2799 (1983)

C. Garcia-Recio et al., PRC 40, 1308 (1989)

Free-space & ChPT

R.A. Arndt et al., PRC 74, 045205 (2006)

$$\frac{db_0}{d\omega} = -0.00053 m_\pi^{-1} / \text{MeV}$$

$$\frac{db_1}{d\omega} = 0$$

Effective Model

$$b_0(\omega) = -0.010 - 0.00016 \omega$$

$$\bar{b}_1(\rho) = -0.088 \left(1 + \frac{0.6116}{b_1} \frac{\rho}{\rho_0} \right)$$

$$\bar{b}_0(\rho, \omega) = b_0 - \frac{3}{2\pi} (b_0^2 + 2b_1^2) \left(\frac{3\pi^2}{2} \rho \right)^{\frac{1}{3}}$$

M. Krell et al. NPB 11, 521 (1969)

double scattering

$$\bar{b}_0 = b_0 - \frac{3}{2\pi} (b_0^2 + 2b_1^2) \left(\frac{3\pi^2}{2} \rho \right)^{\frac{1}{3}}$$

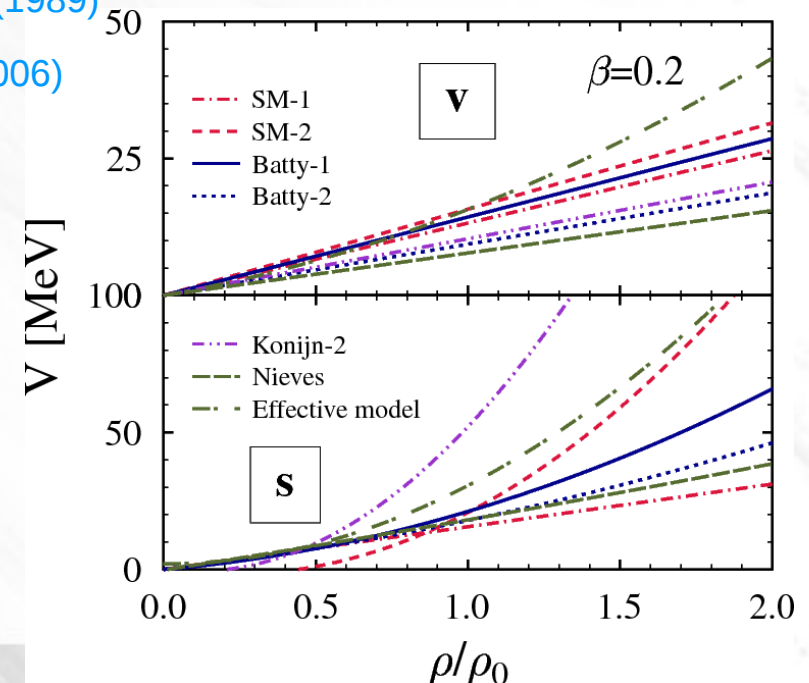
$$b_0 = 0.0 \quad b_1 = \frac{-m_\pi}{8\pi(1+m_\pi/m_N) f_\pi^2}$$

ChPT

$$f_\pi^2(\rho) = f_\pi^2(0) - \frac{\sigma \rho}{m_\pi^2}$$

W. Weise, NPA 690, 98 (2001)

$$b_1(\rho) = \frac{b_1}{1 - \frac{\sigma \rho}{m_\pi^2 f_\pi^2}} \simeq \frac{b_1}{1 - 2.3\rho}$$



S-wave parametrizations: subset from K. Itahashi et al., PRC 62, 025202 (2000)

P-wave pion potential

Energy dependence:

extrapolation of pionic atoms results using a local approximation of the delta-hole model that describes pion-nucleus scattering up to $\omega=300$ MeV

C. Garcia-Recio et al., NPA 526, 685 (1991)

$$f(p^2) = \frac{1 - p_{\text{eff}}^2 / \Lambda_1^2 + p_{\text{eff}}^4 / \Lambda_2^4}{1 - p^2 / \Lambda_1^2 + p^4 / \Lambda_2^4}$$

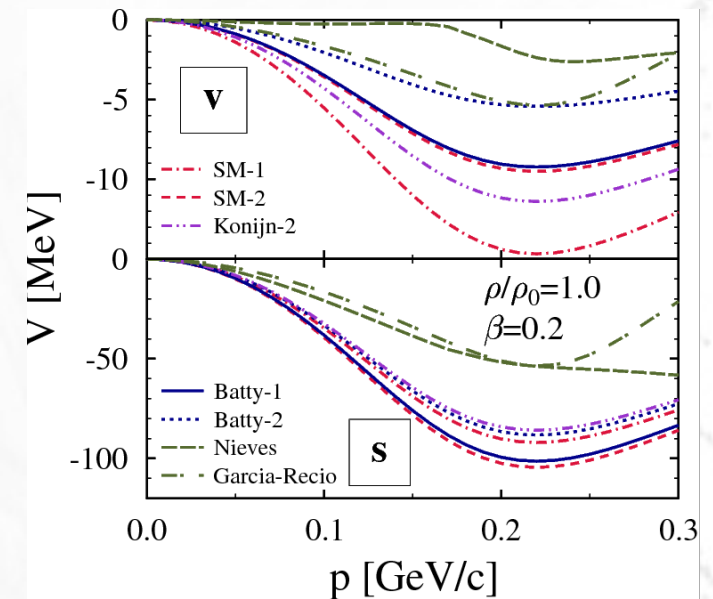
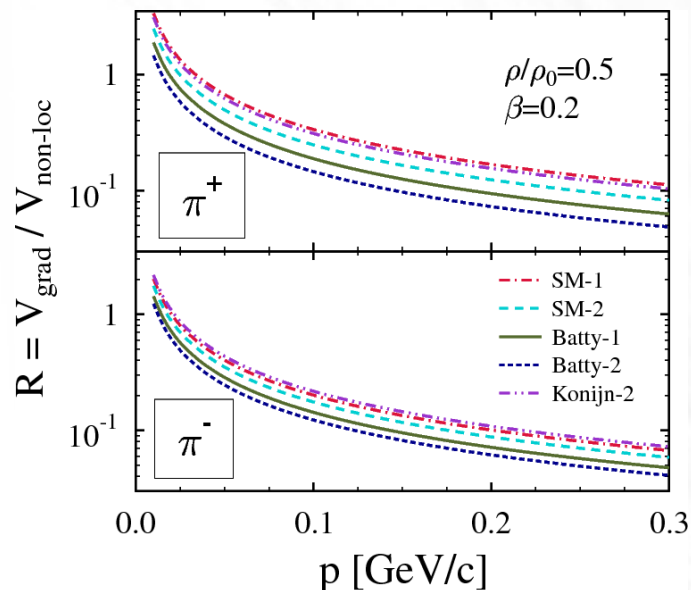
$$\Lambda_1 = 0.55 \text{ GeV} \quad \Lambda_2 = 0.22 \text{ GeV}$$

$$p_{\text{eff}} = 0.05 \text{ GeV}$$

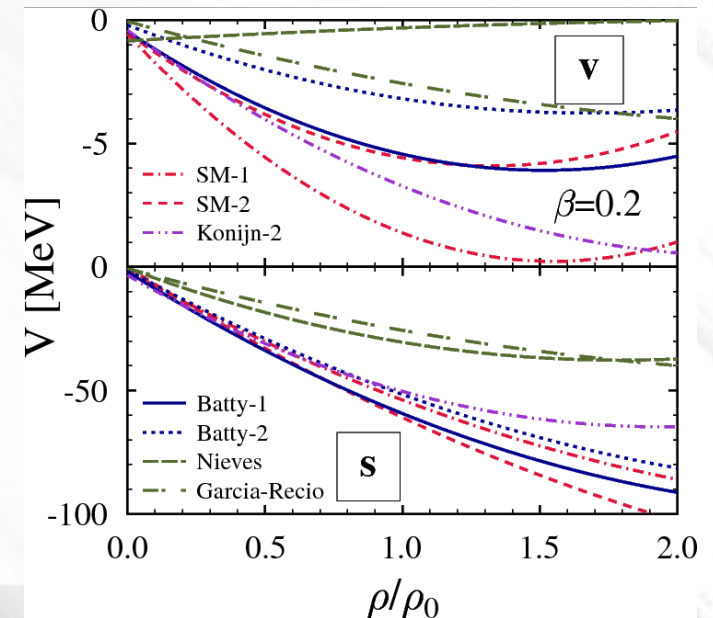
Gradient terms:

$$V_{\text{opt}}^p(r) = \frac{2\pi}{\mu} \vec{\nabla} \frac{\alpha(r)}{1 + 4/3 \pi \lambda \alpha(r)} \vec{\nabla}$$

$$\Rightarrow \text{terms} \sim \vec{p} \cdot \vec{\nabla} \rho \quad \vec{p} \cdot \vec{\nabla} \beta$$



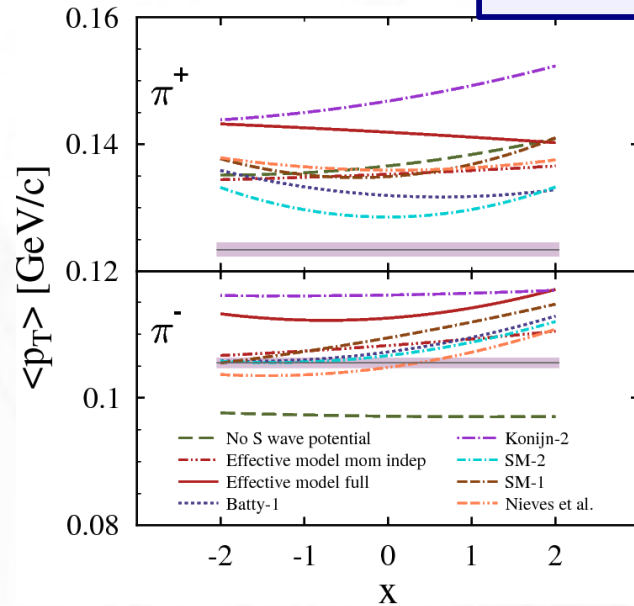
Density dependence



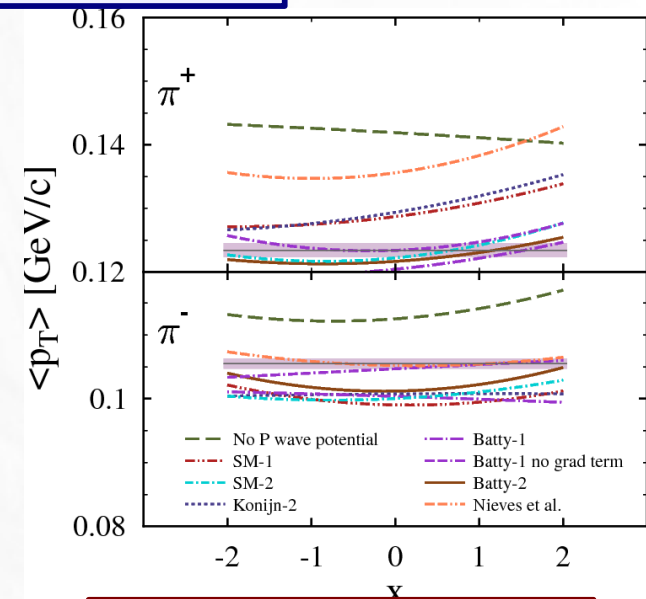
P-wave parametrizations: subset from K. Itahashi et al., PRC 62, 025202 (2000)

Impact on pion observables

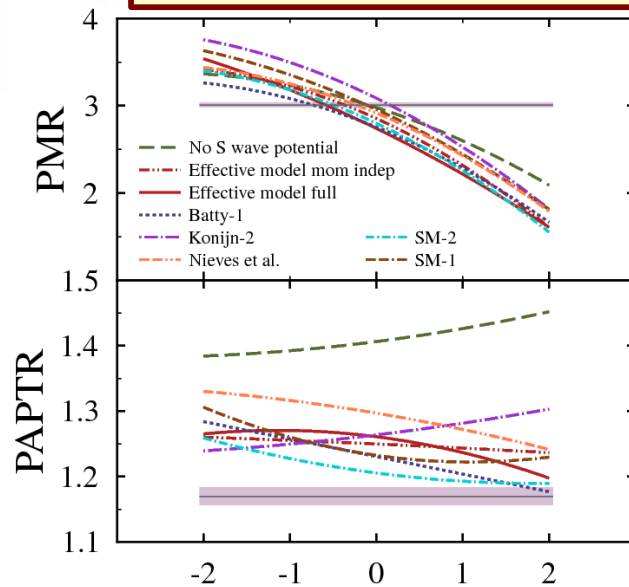
Au+Au@400 MeV/A, 3.35 fm < b < 6.0 fm



$|y/y_p| < 1.75$
 $0 < p_T < 0.33$ GeV/c
 exp data:
 W.Reisdorf et al. (FOPI),
 private communication



S-wave pion potential



P-wave pion potential

