

Workshop on Challenges to Transport Theory
for Heavy-Ion Collisions, ECT*, May 20-24, 2019

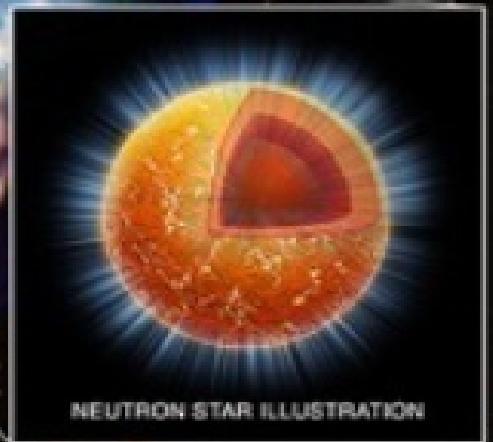
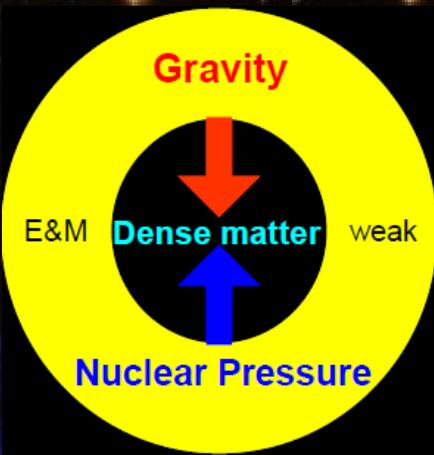
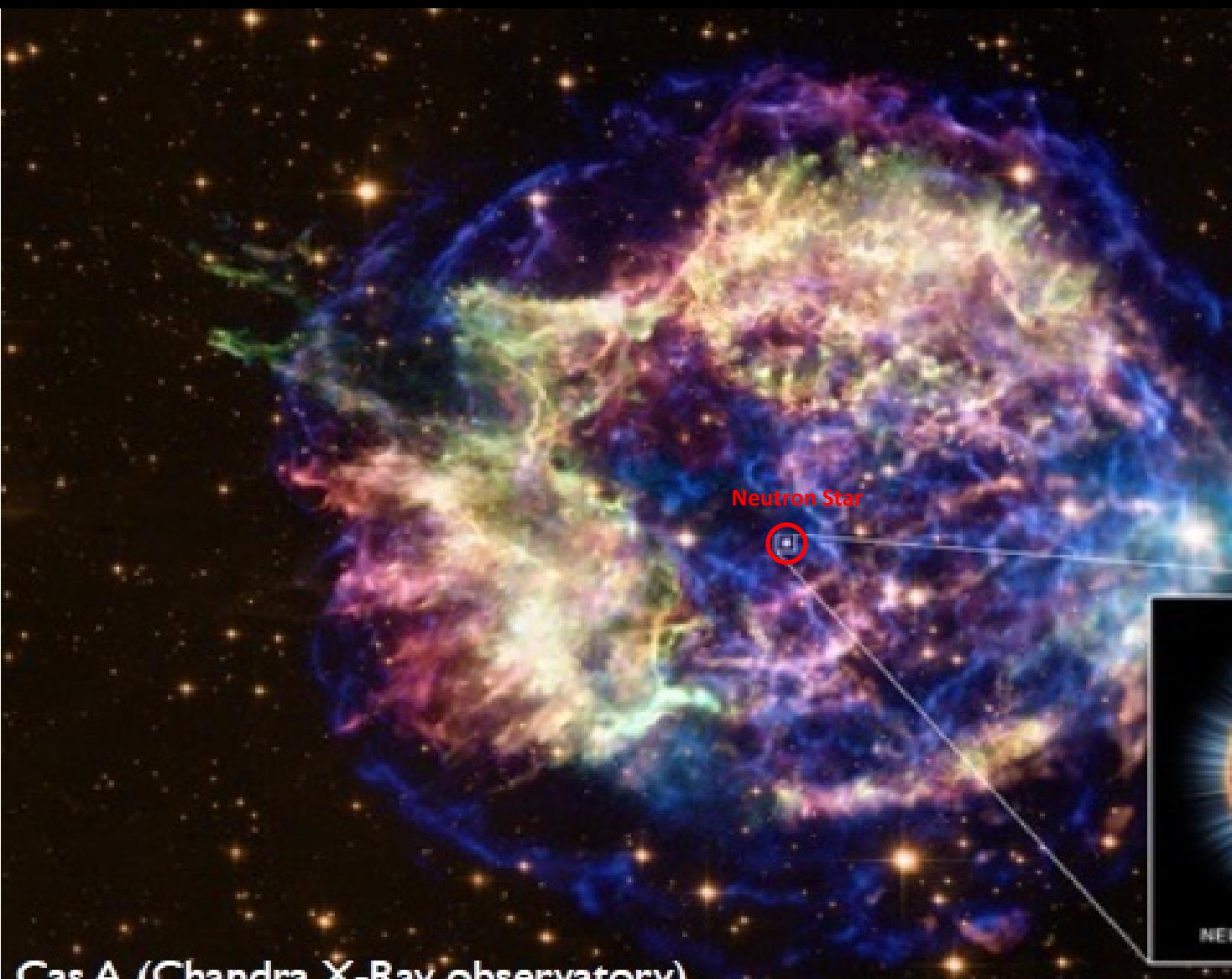
Medium Energy Heavy-Ion Collisions Encounter Short-Range Correlations

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2019/5/21/15:00

Supernova explosion



NEUTRON STAR ILLUSTRATION

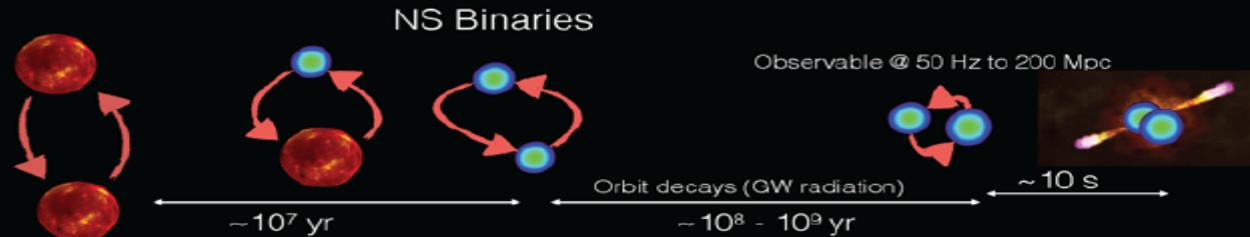
Cas A (Chandra X-Ray observatory)

1	H
2	He
3	Li
4	B
5	Be
6	C
7	N
8	O
9	F
10	Ne
11	Na
12	Mg
13	Al
14	Si
15	P
16	S
17	Cl
18	Ar
19	K
20	Ca
21	Ti
22	V
23	Cr
24	Mn
25	Fe
26	Co
27	Ni
28	Cu
29	Zn
30	Ga
31	Ge
32	As
33	Se
34	Br
35	Kr
36	Xe
37	Rb
38	Sr
39	Y
40	Zr
41	Nb
42	Mo
43	Tc
44	Ru
45	Rh
46	Pd
47	Ag
48	Cd
49	Sn
50	Bi
51	Sb
52	Te
53	I
54	Xe
55	Cs
56	Ba
57	Hf
58	Ta
59	W
60	Re
61	Ov
62	Pt
63	Au
64	Hg
65	Ga
66	Tb
67	Dy
68	Ho
69	Er
70	Tm
71	Yb
72	Lu
73	Ac
74	Th
75	Pa
76	U
77	Pm
78	Em
79	Un
80	Tl
81	Pb
82	Ba
83	Po
84	At
85	Ro
86	Ru
87	Fr
88	Ra

Merging Neutron Stars Exploding Massive Stars
Dying Low Mass Stars Exploding White Dwarfs Big Bang
Exploding White Dwarfs Cosmic Ray Fission

Source: <http://www.astro.vt.edu/~breske/elements.html>

Element Origins



GW170817

Heavy elements production

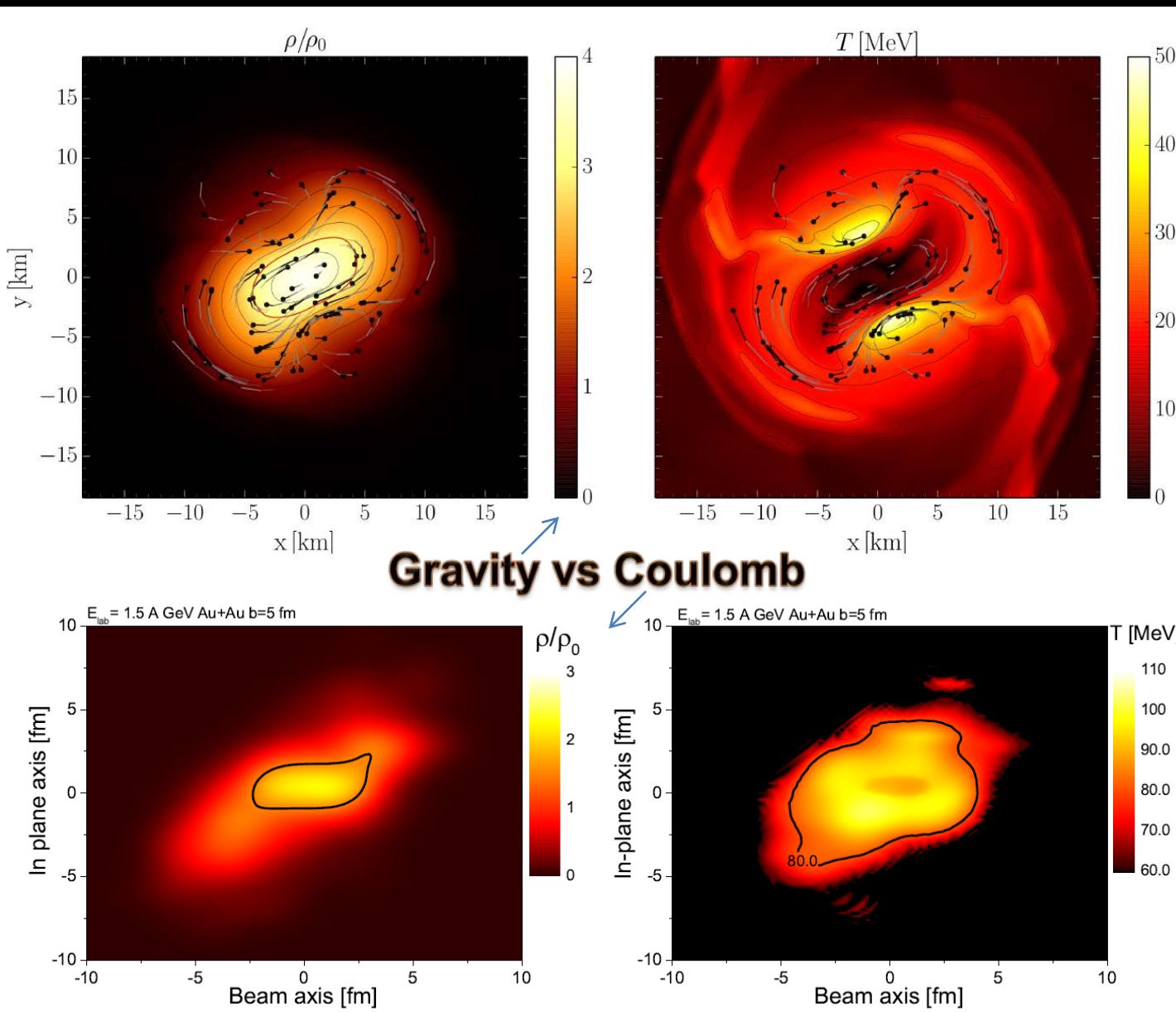
GRB

GW

Neutron star mergers and heavy-ion collisions

density

temperature



M. Hanauske et al.,
J. Phys.: Conf. Ser.
878 012031

n-star merger

EOS

Au +Au
1.5A GeV

Equation of State of nuclear matter

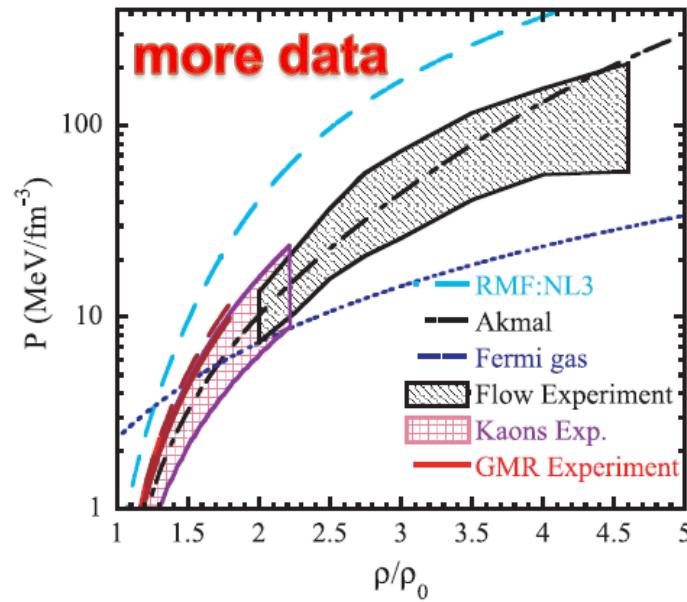
N=Z

$$E(\rho, \delta) \approx E(\rho, \delta = 0) + E_{\text{sym}}(\rho) \delta^2,$$

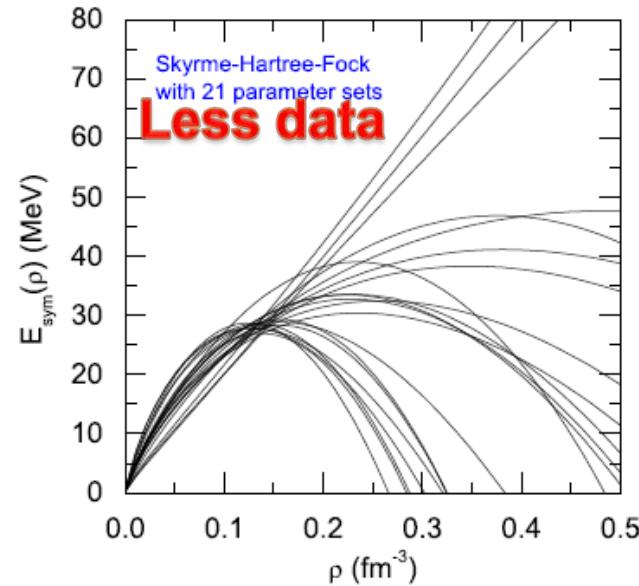
N-rich

$$\delta = \frac{\rho_n - \rho_p}{\rho_n + \rho_p}$$

Constraints at high densities have been extracted



poorly known especially at high densities

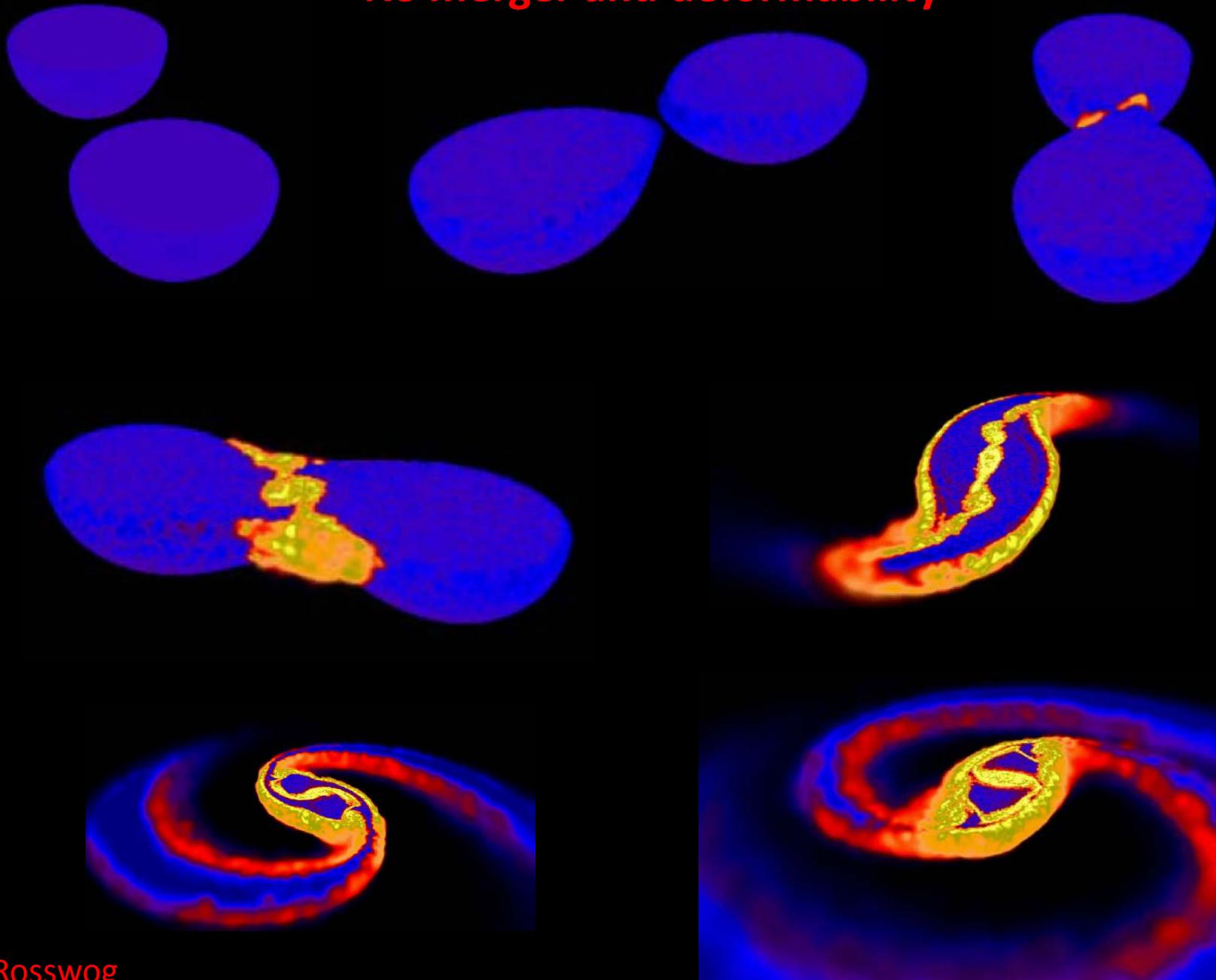


W.G. Lynch, et al., arXiv: 090.0412, [nucl-ex]

P. Danielewicz, R. Lacey, W.G. Lynch, Science 298 (2002) 1592

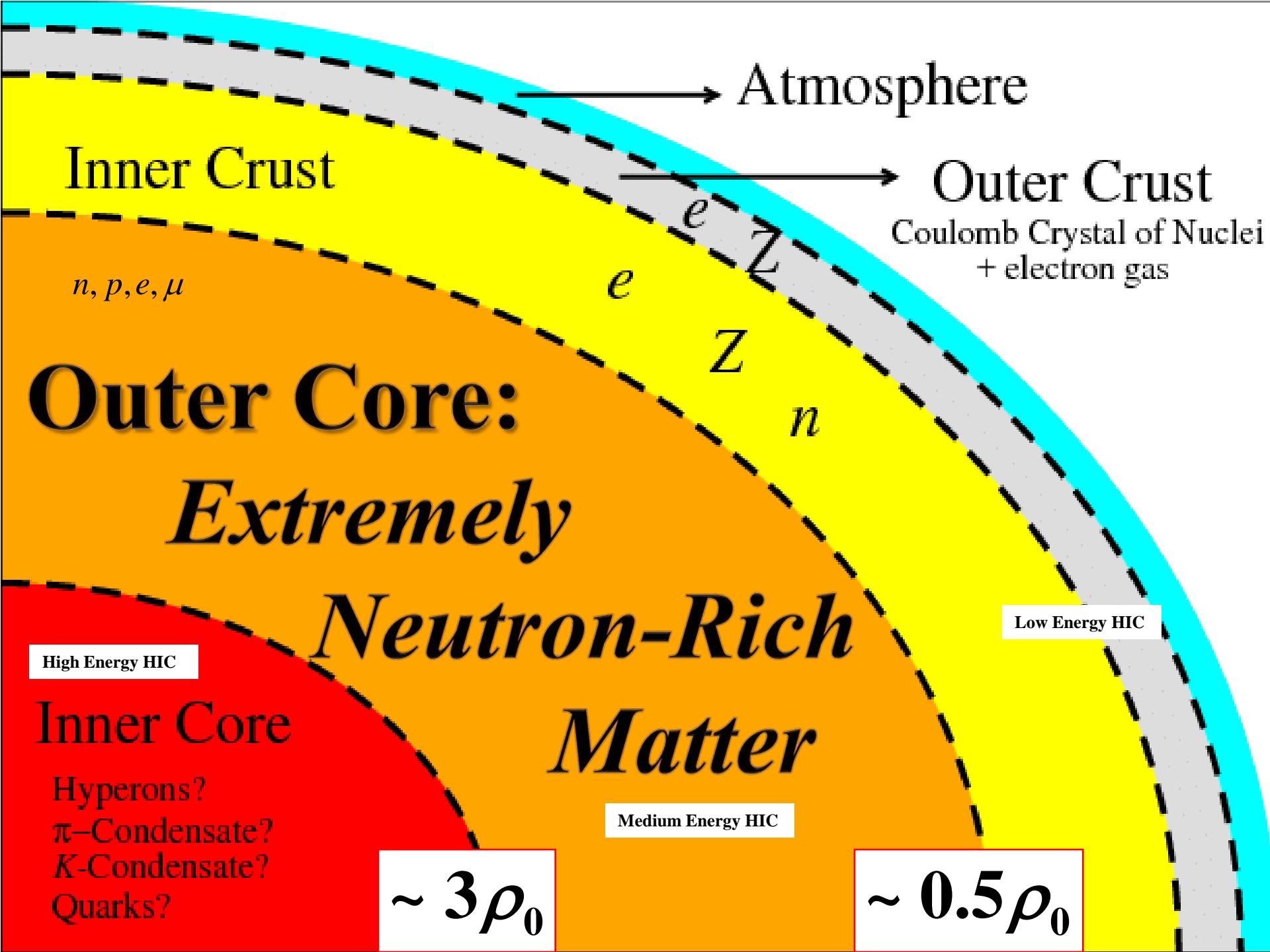
B.A. Li , L.W. Chen, Che Ming Ko, Phys. Rep. 464, 113 (2008)

NS merger and deformability

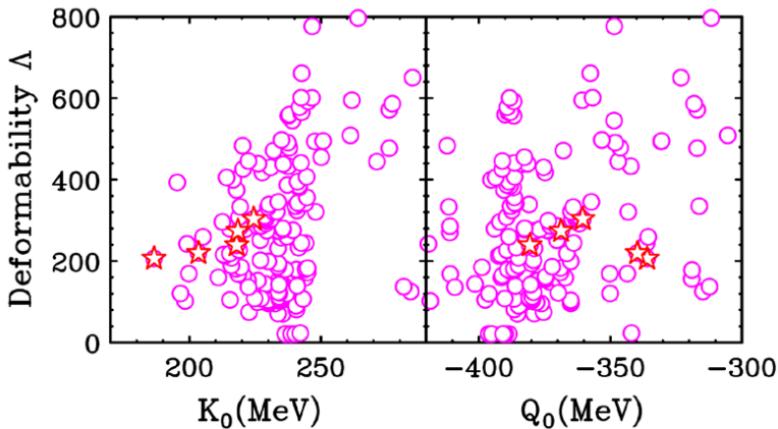


S. Rosswog

<http://compact-merger.astro.su.se/index.html>



Deformability and the symmetry energy



$$\varepsilon(\rho, \delta) = \varepsilon(\rho, \delta=0) + S(\rho)\delta^2 + O(\delta^4) + \dots$$

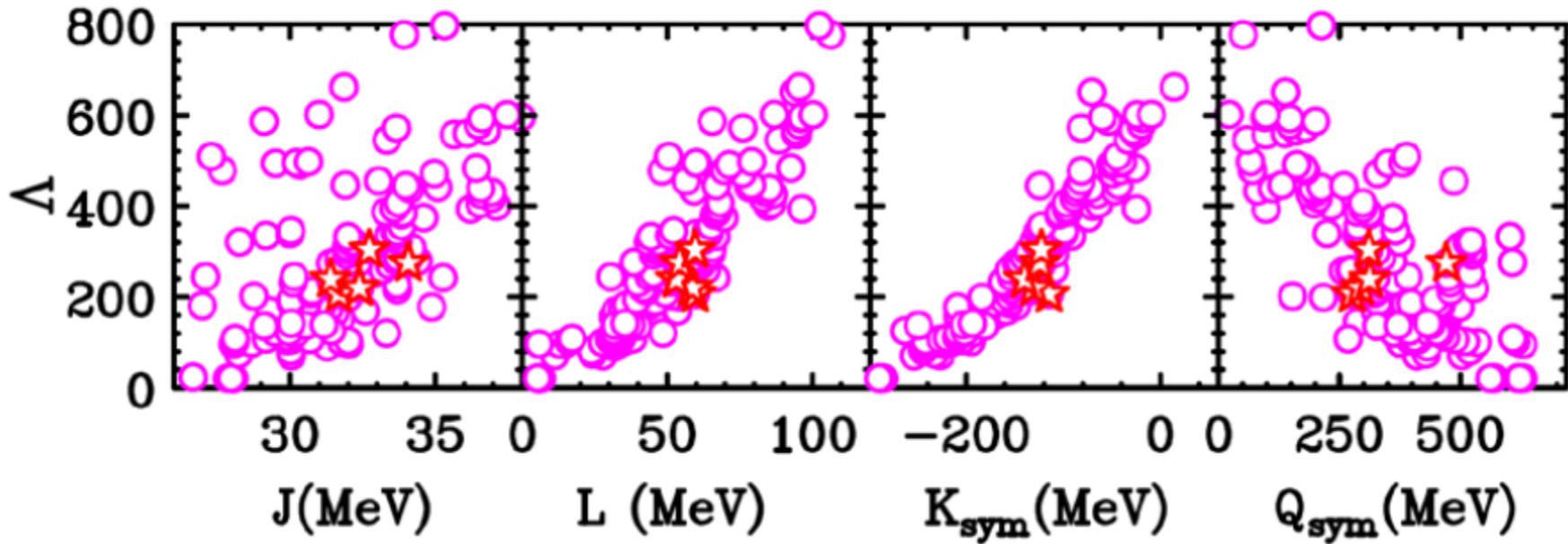
Symmetric EoS:

$$\varepsilon_{SNM}(\rho) = E_0 + \frac{1}{2}K_0 x^2 + \frac{1}{6}Q_0 x^3 + O(x^4)$$

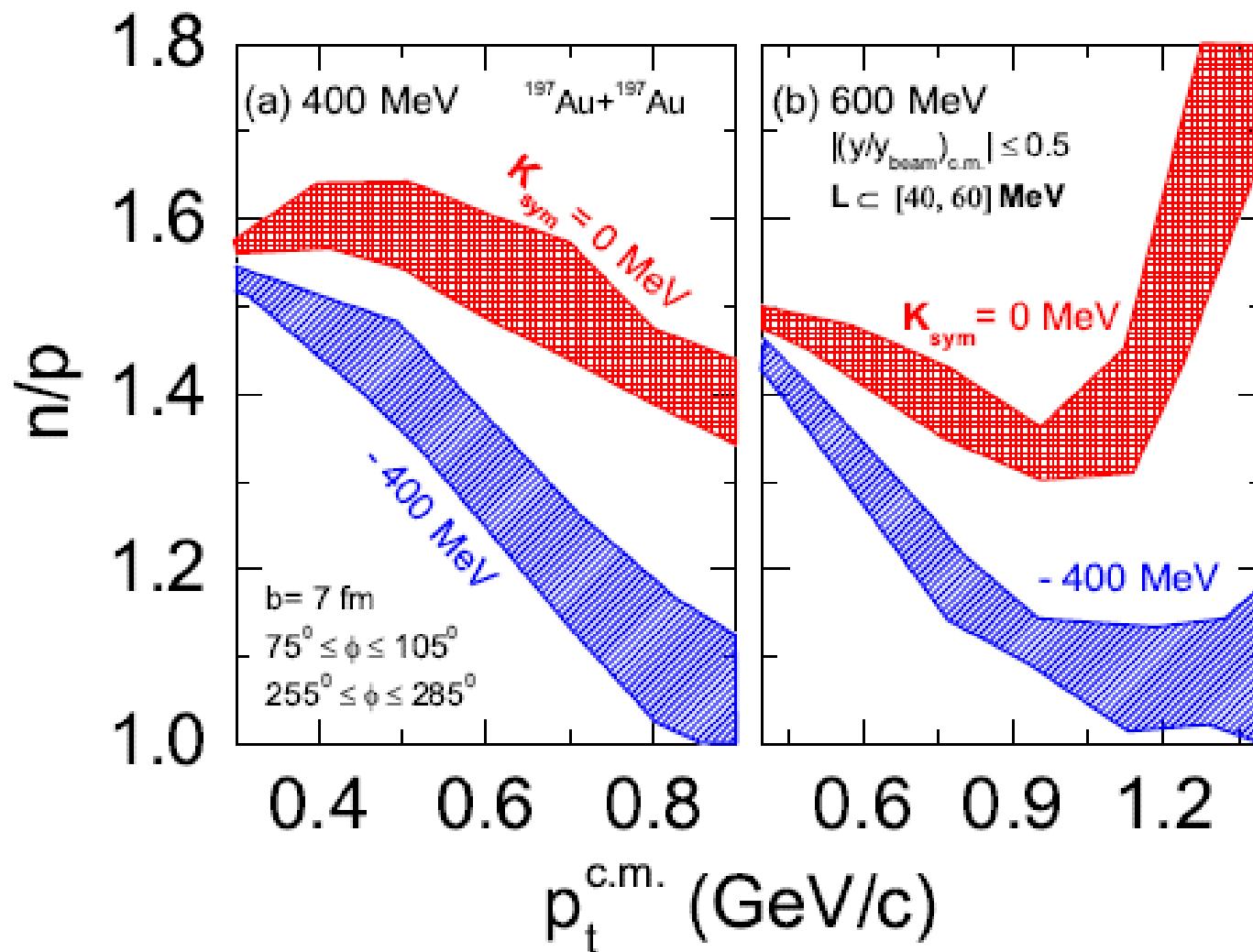
Symmetry energy:

$$S(\rho) = J + Lx + \frac{1}{2}K_{sym}x^2 + \frac{1}{6}Q_{sym}x^3 + O(x^4)$$

$$\Lambda = (2/3)k_2[(c^2/G)(R/m)]^5$$

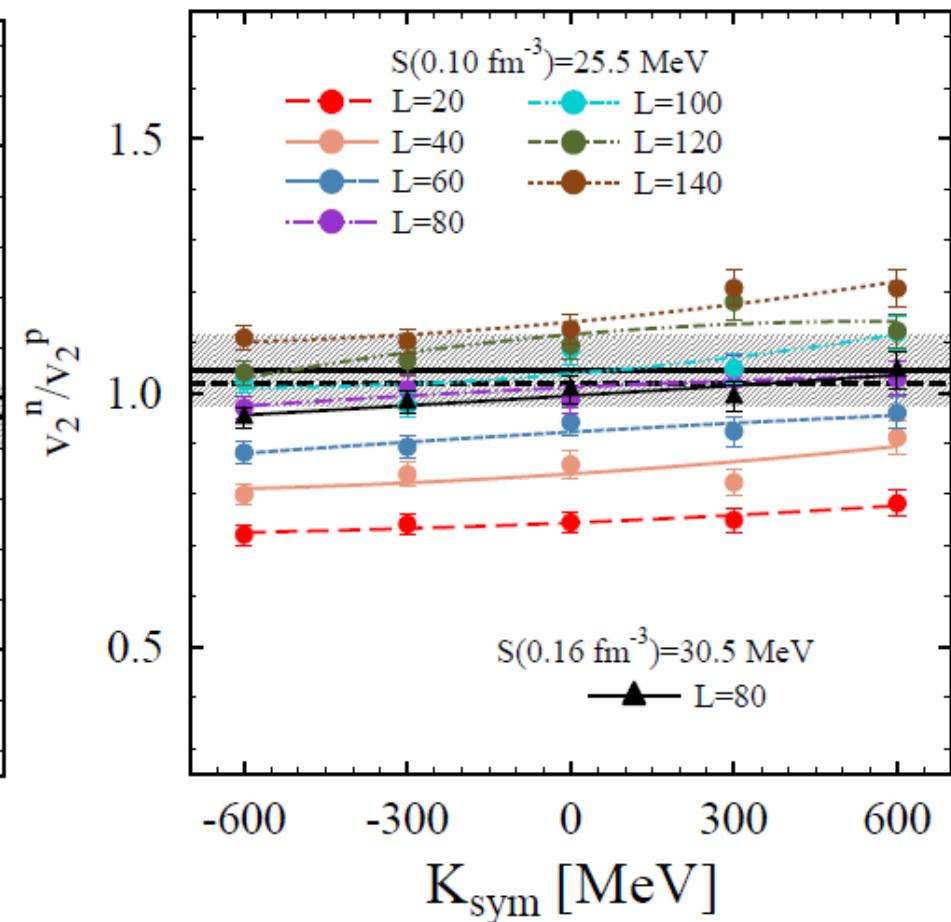
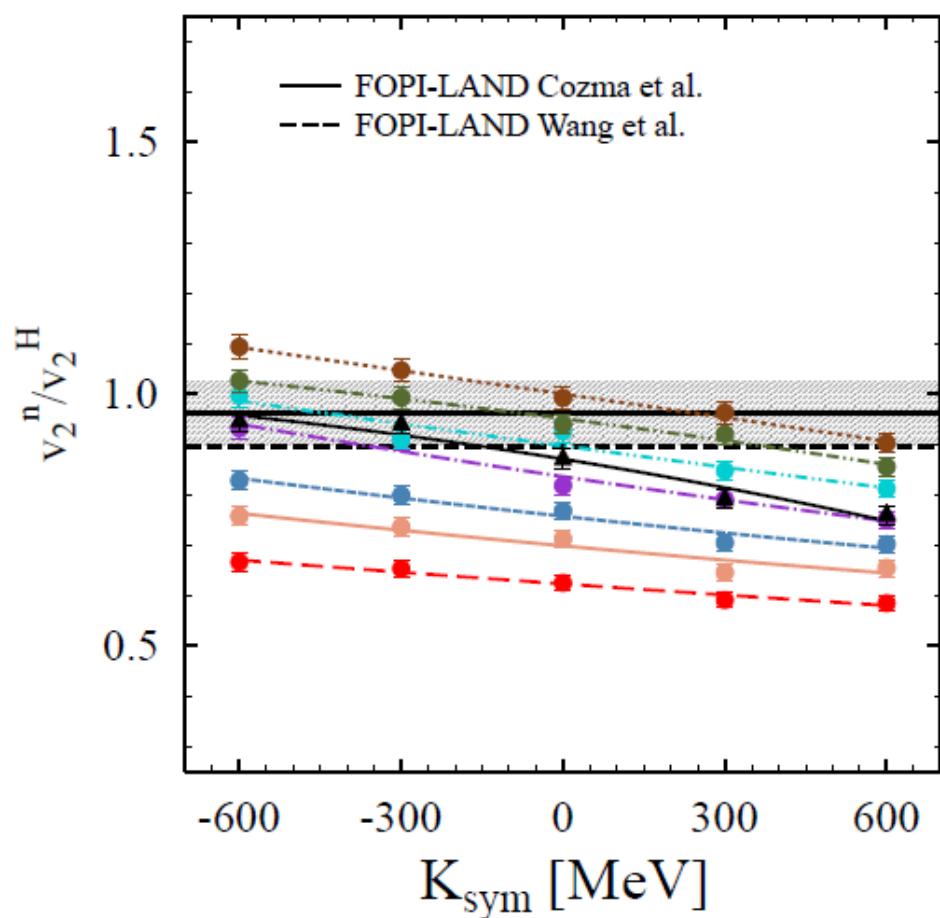


Probing Ksym by HIC



Probing Ksym by HIC

Cozma, M.D. Eur. Phys. J. A (2018) 54: 40



More detailed data are needed.

Challenges of transport model

old days:

- In-medium effects of transport cannot be experimentally well determined
- Non-equilibrium quantum many-body system in finite volume is troublesome

nowadays:

- Pursuing more details/tiny effects
- Nucleon-nucleon Short-range interactions is not well understood

Short-Range-Correlation by Jlab

Final State Interactions

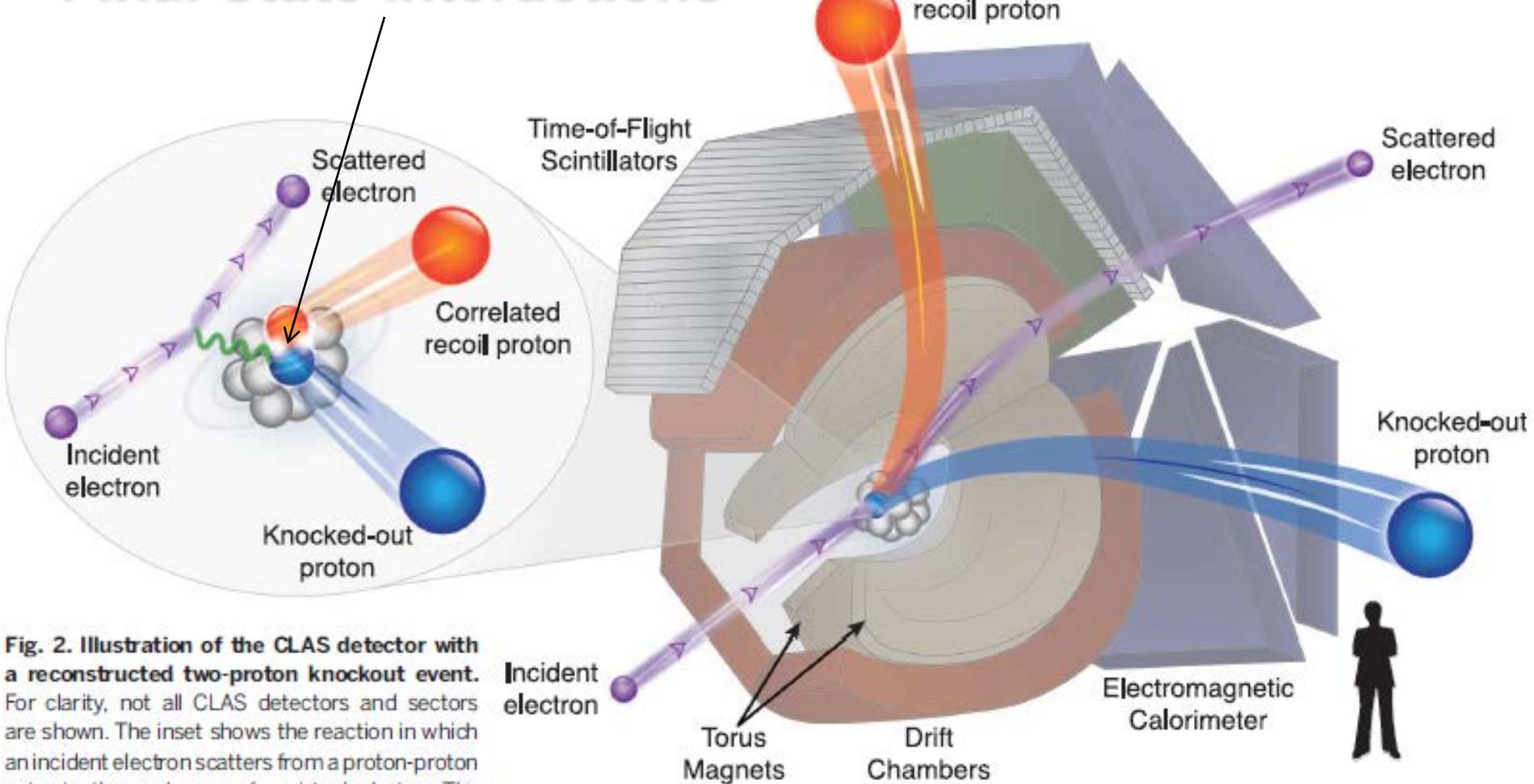


Fig. 2. Illustration of the CLAS detector with a reconstructed two-proton knockout event. For clarity, not all CLAS detectors and sectors are shown. The inset shows the reaction in which an incident electron scatters from a proton-proton pair via the exchange of a virtual photon. The human figure is shown for scale.

O. Hen et al., Science 346, 614 (2014).

What are Short Range Correlations (SRC) in nuclei ?

(Eli Piasetzky)

SRC $\sim R_N$

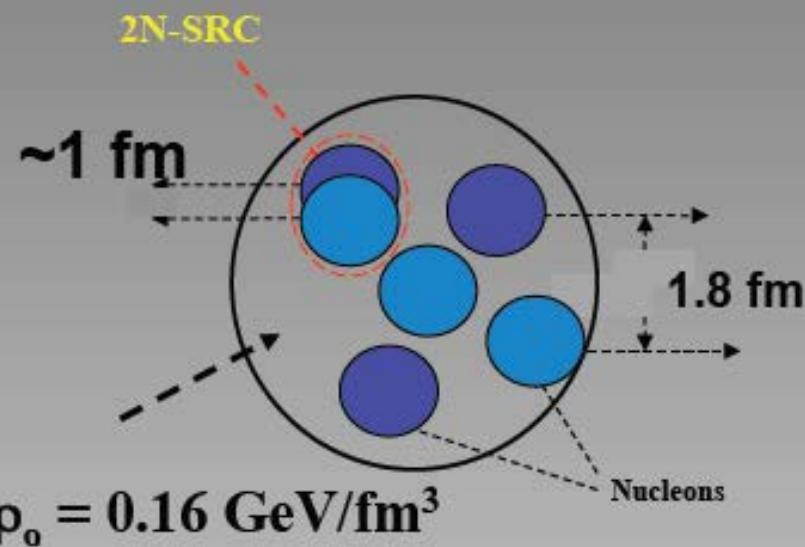
LRC $\sim R_A$

$k_F \sim 250 \text{ MeV/c}$

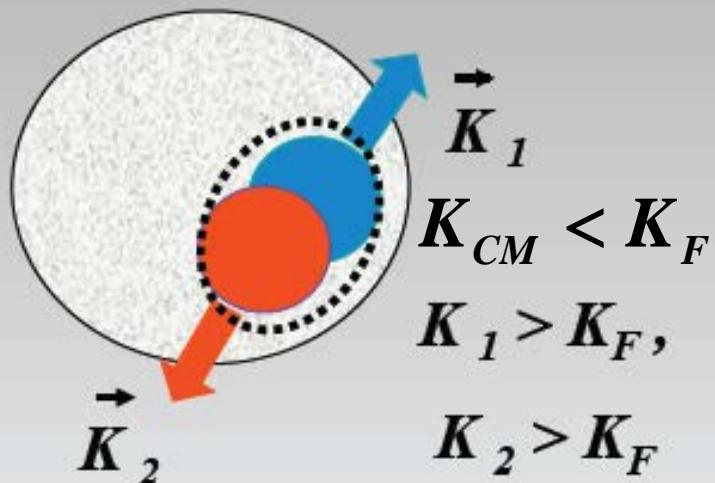
High momentum tail:

$300-600 \text{ MeV/c}$

$1.5 K_F - 3 K_F$

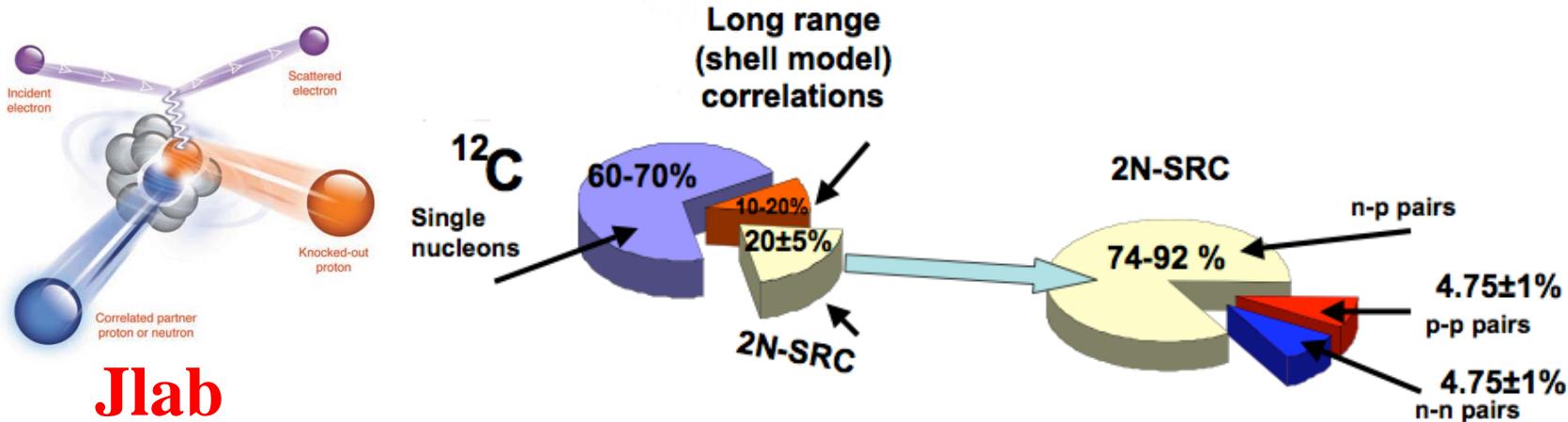


In momentum space:



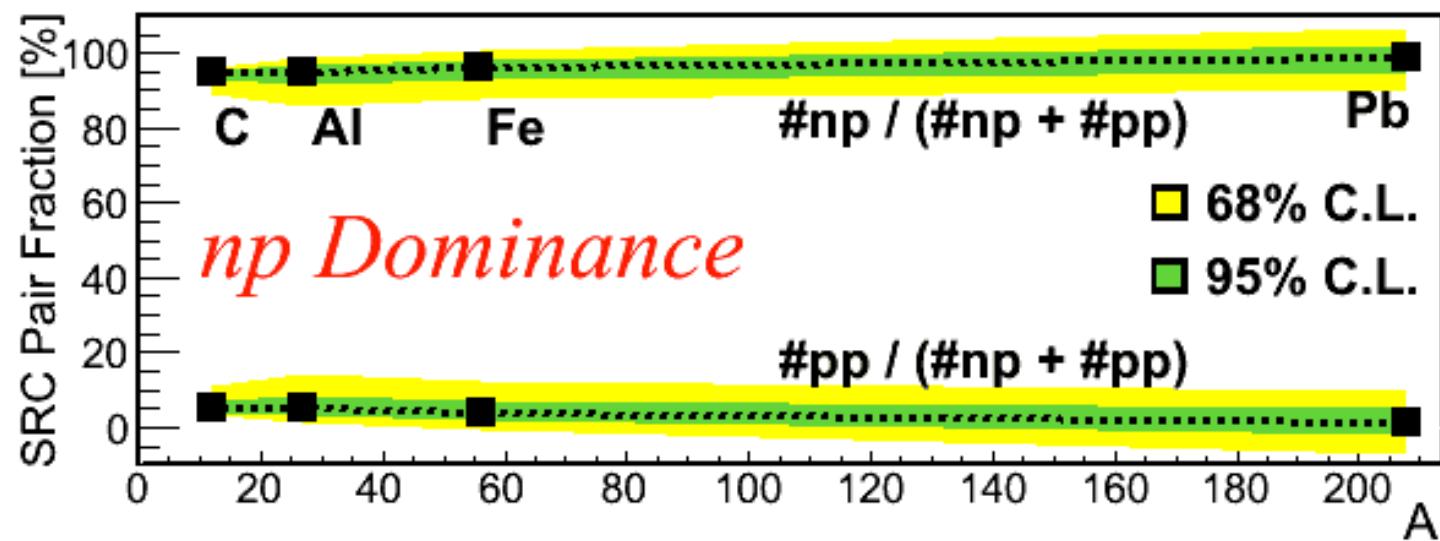
A pair with large relative momentum between the nucleons and small CM momentum.

Nucleon-nucleon Short-range-Correlations



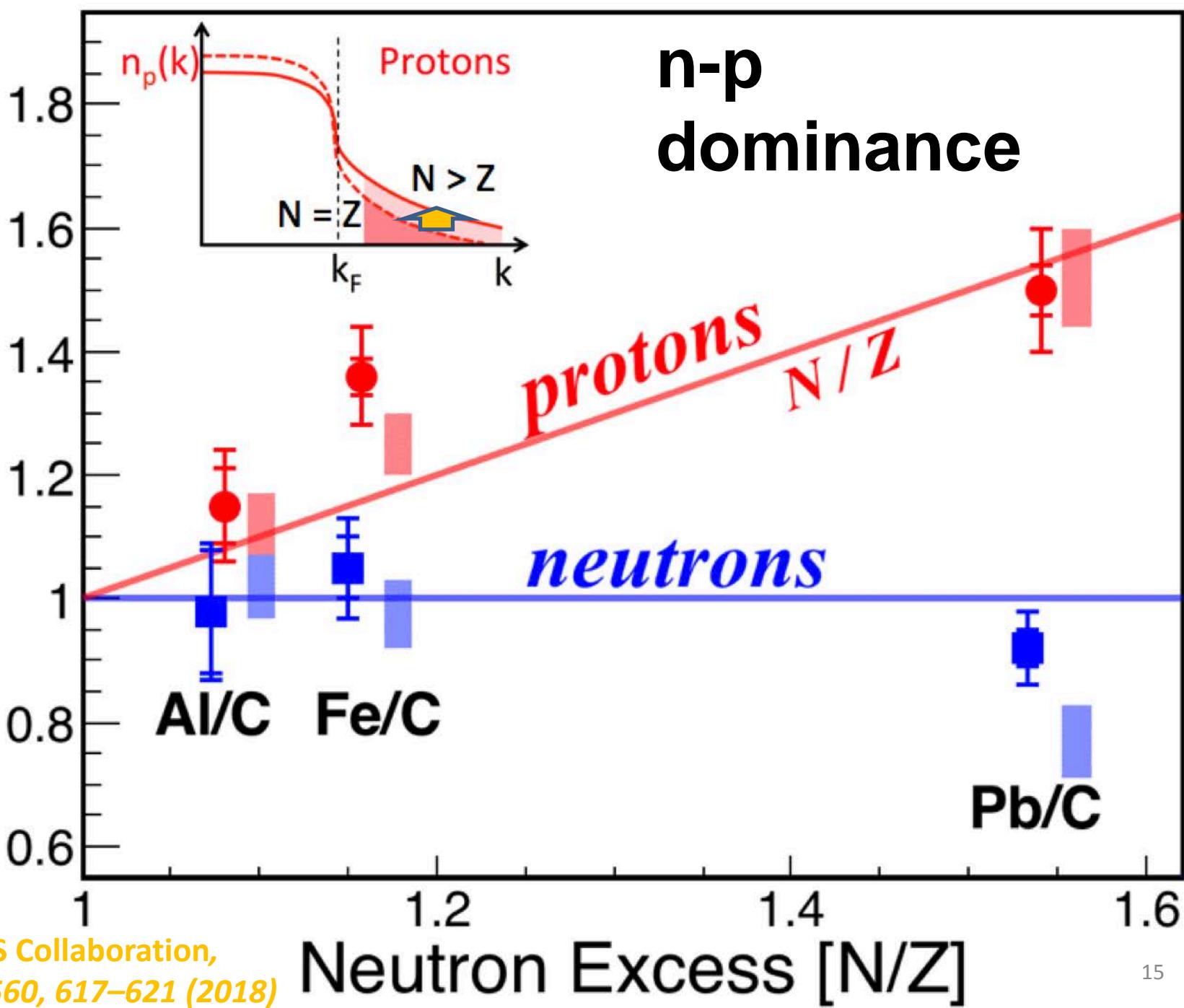
Jlab

R. Subedi et al., Science 320, 1476 (2008).

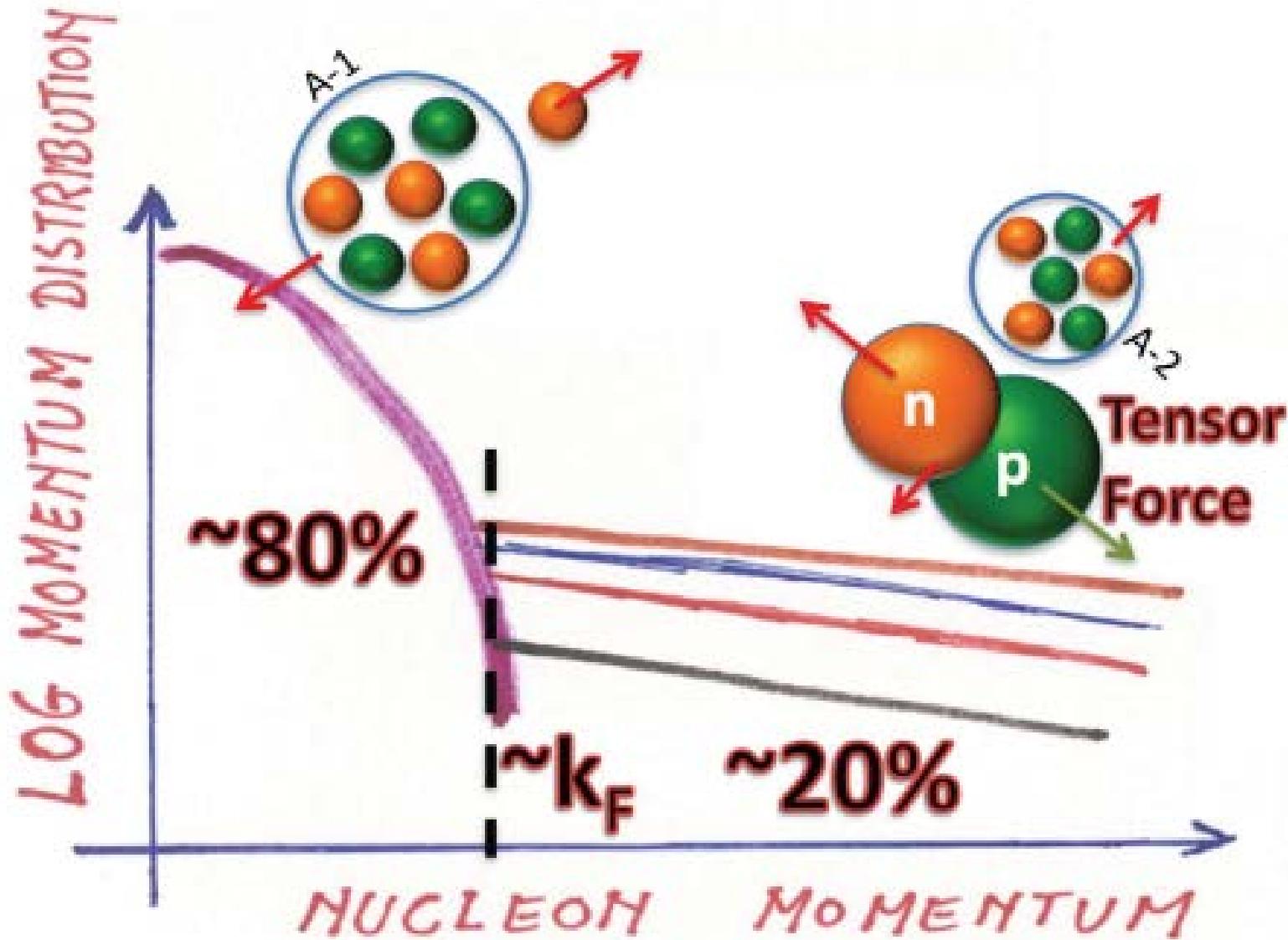


O. Hen et al., Science 346, 614 (2014).

High-Momentum Fraction

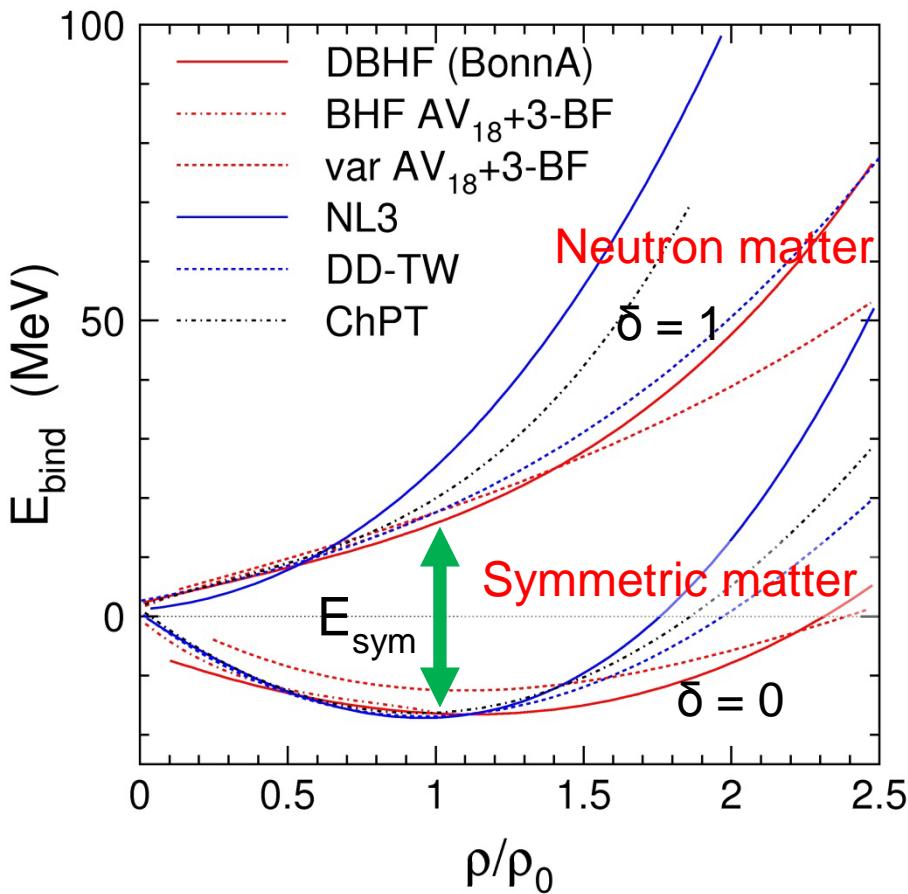


Dominant features of nucleon mom. Dis.



How SRC affect symmetry energy

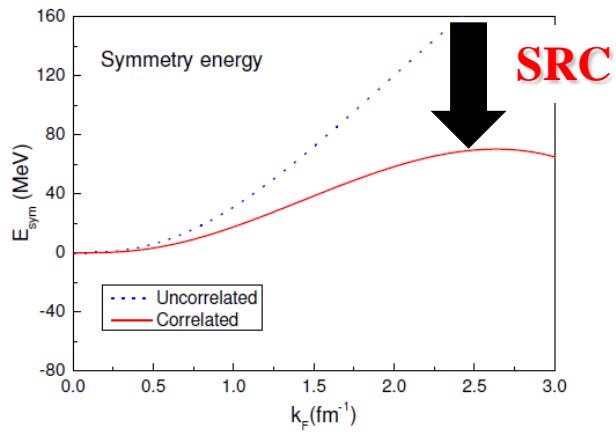
Fuchs and Wolter, EPJA 30 (2006)



$$E(\rho, \delta) \approx E(\rho, \delta = 0) + E_{\text{sym}}(\rho)\delta^2, \quad \delta = \frac{\rho_n - \rho_p}{\rho_n + \rho_p}$$

$$E_{\text{sym}} \propto E_{\text{nentron}} - E_{\text{proton}}$$

n-p SRC:
En-Ep↓ Esym↓



The Esym is decided by the difference of En and Ep

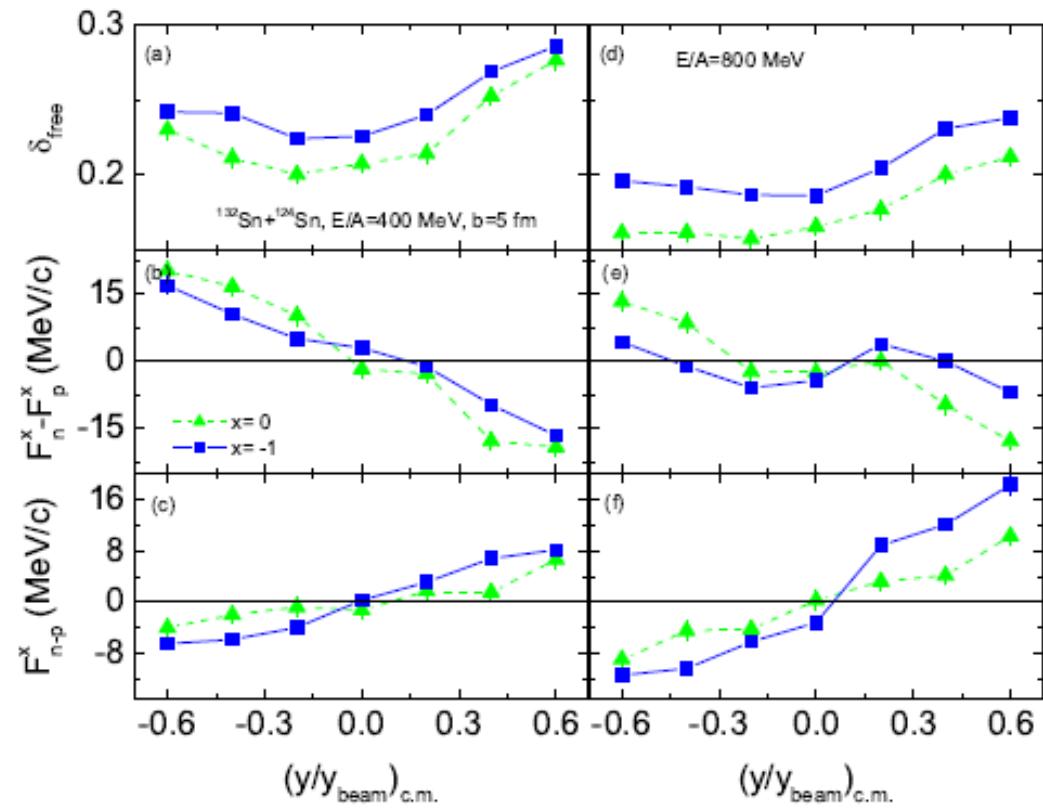
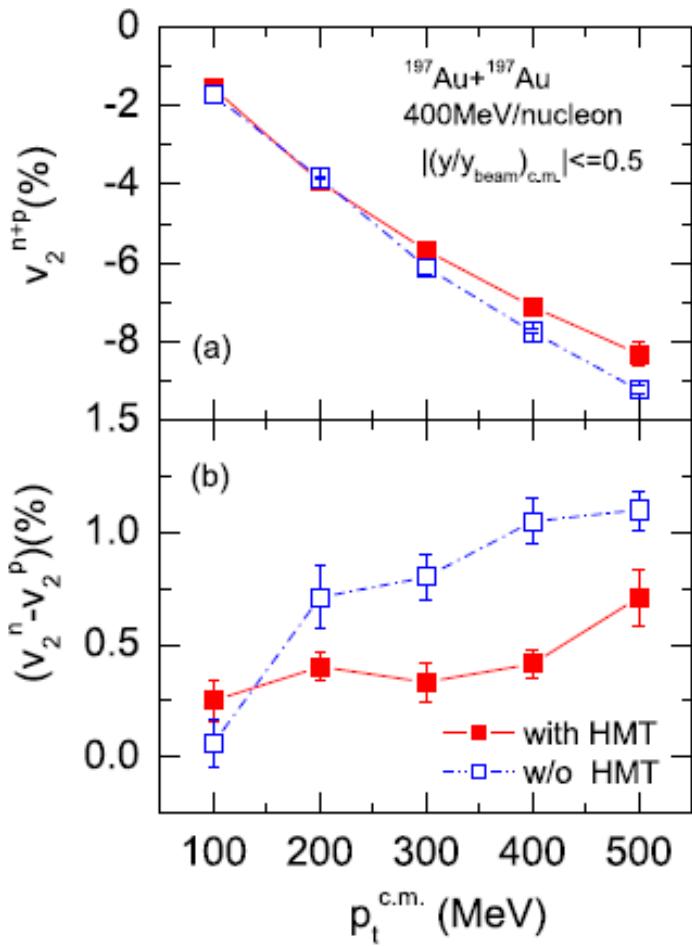
Xu and Li, arXiv:1104.2075

On HIC-hadron probe

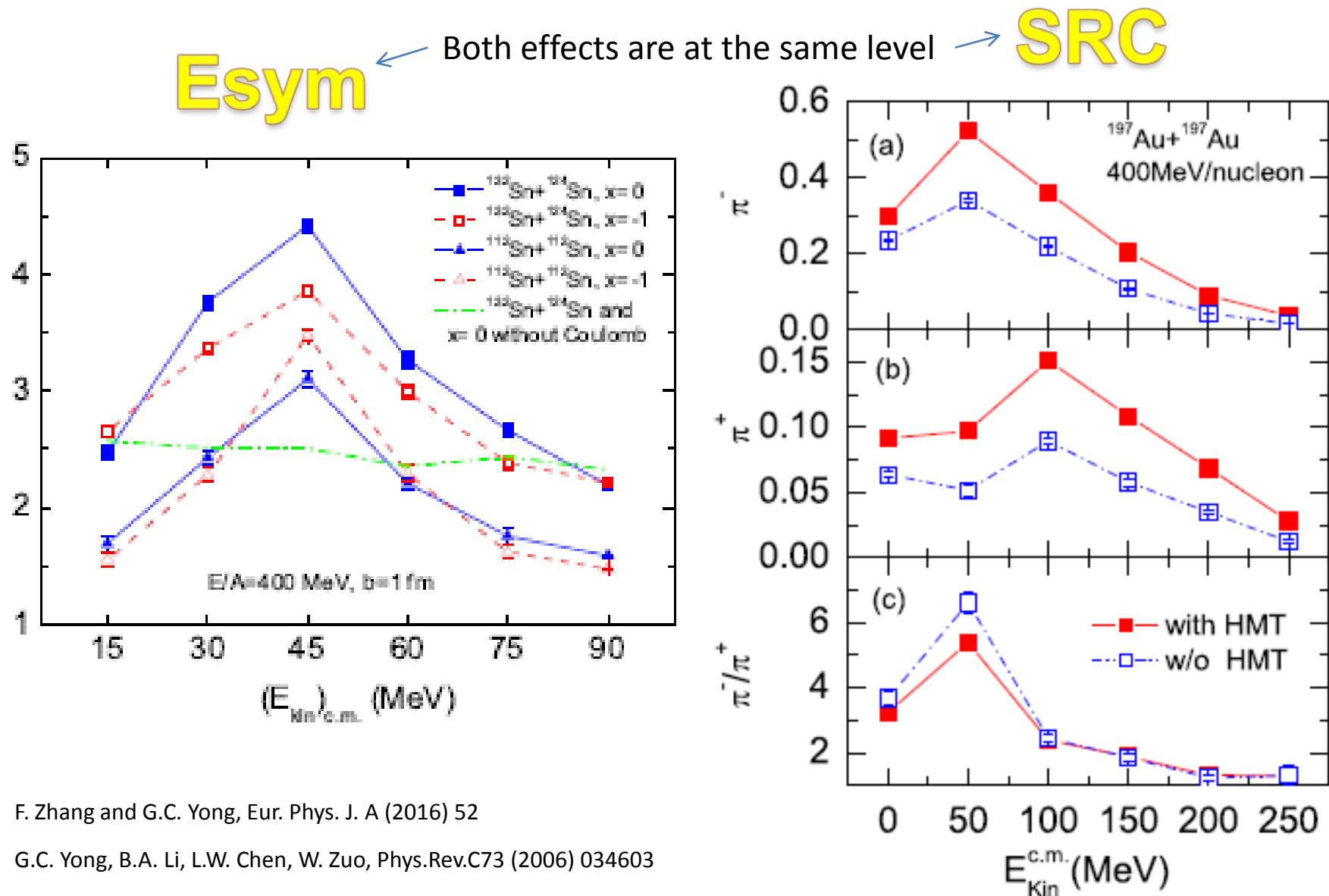
SRC

Both effects are at the same level

Esym



On HIC-hadron probe



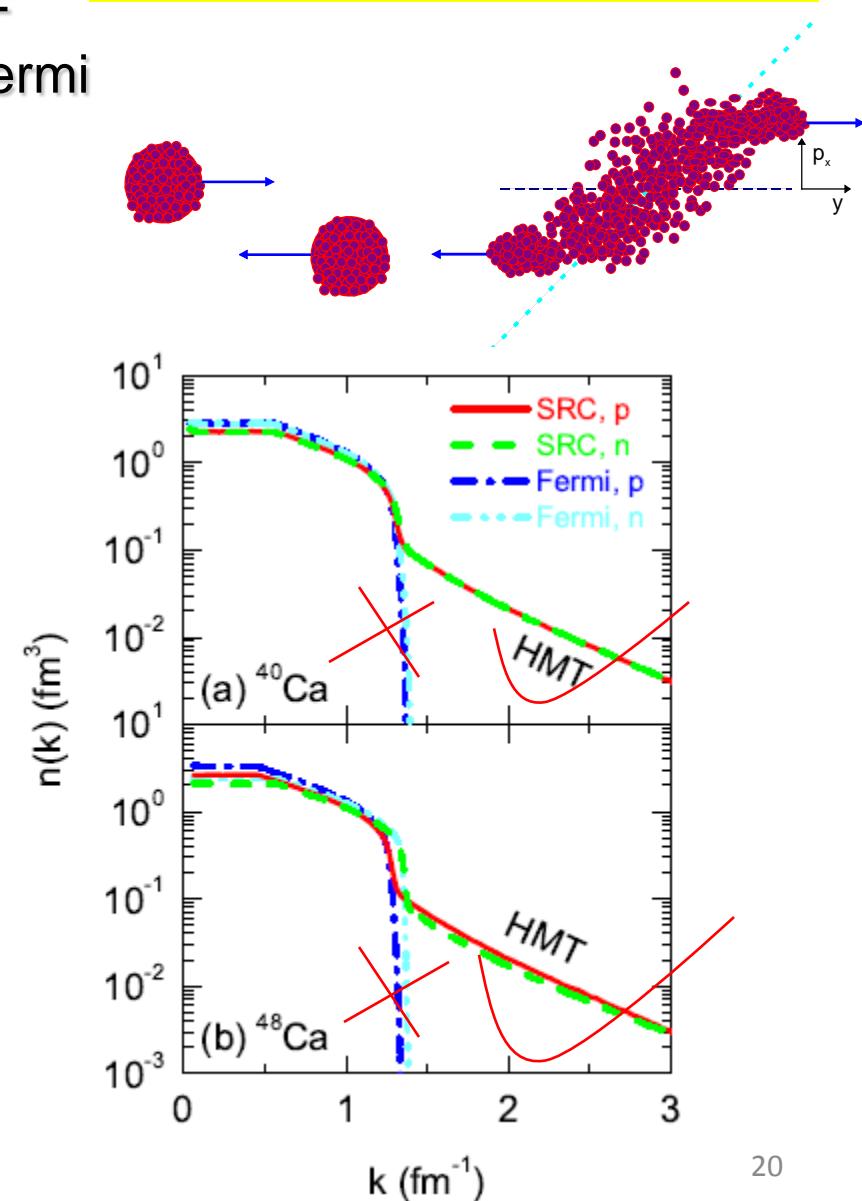
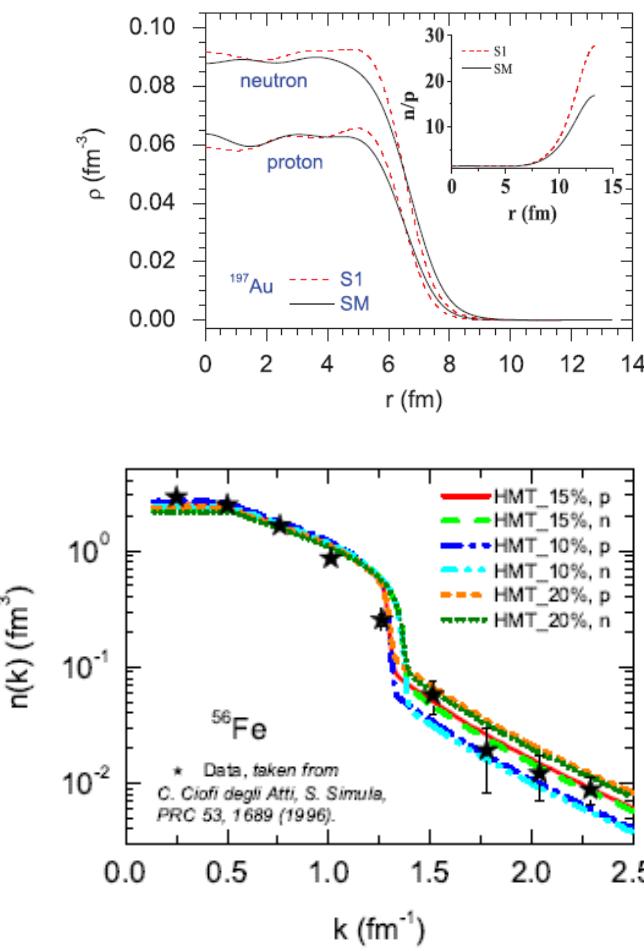
F. Zhang and G.C. Yong, Eur. Phys. J. A (2016) 52

G.C. Yong, B.A. Li, L.W. Chen, W. Zuo, Phys. Rev. C73 (2006) 034603

(1) Initialization:

- (1) Nucleons in coordinates, RMF, SHF
- (2) Momentum-space: localThomas-Fermi

Reproduce HMT with equal
Numbers of n and p



Formula used to reproduce SRC/HMT

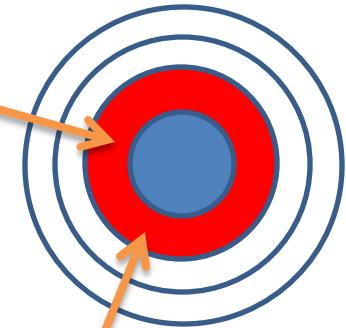
$$n(k) = \begin{cases} C_1, k \leq k_F \\ C_2/k^4, k_F < k \leq \lambda k_F \end{cases}$$

$$K_{F_{n,p}}(r) = [3\pi^2 \rho(r)_{n,p}]^{1/3}$$

$$\int_0^{\lambda k_F} n(k) k^2 dk = 1$$

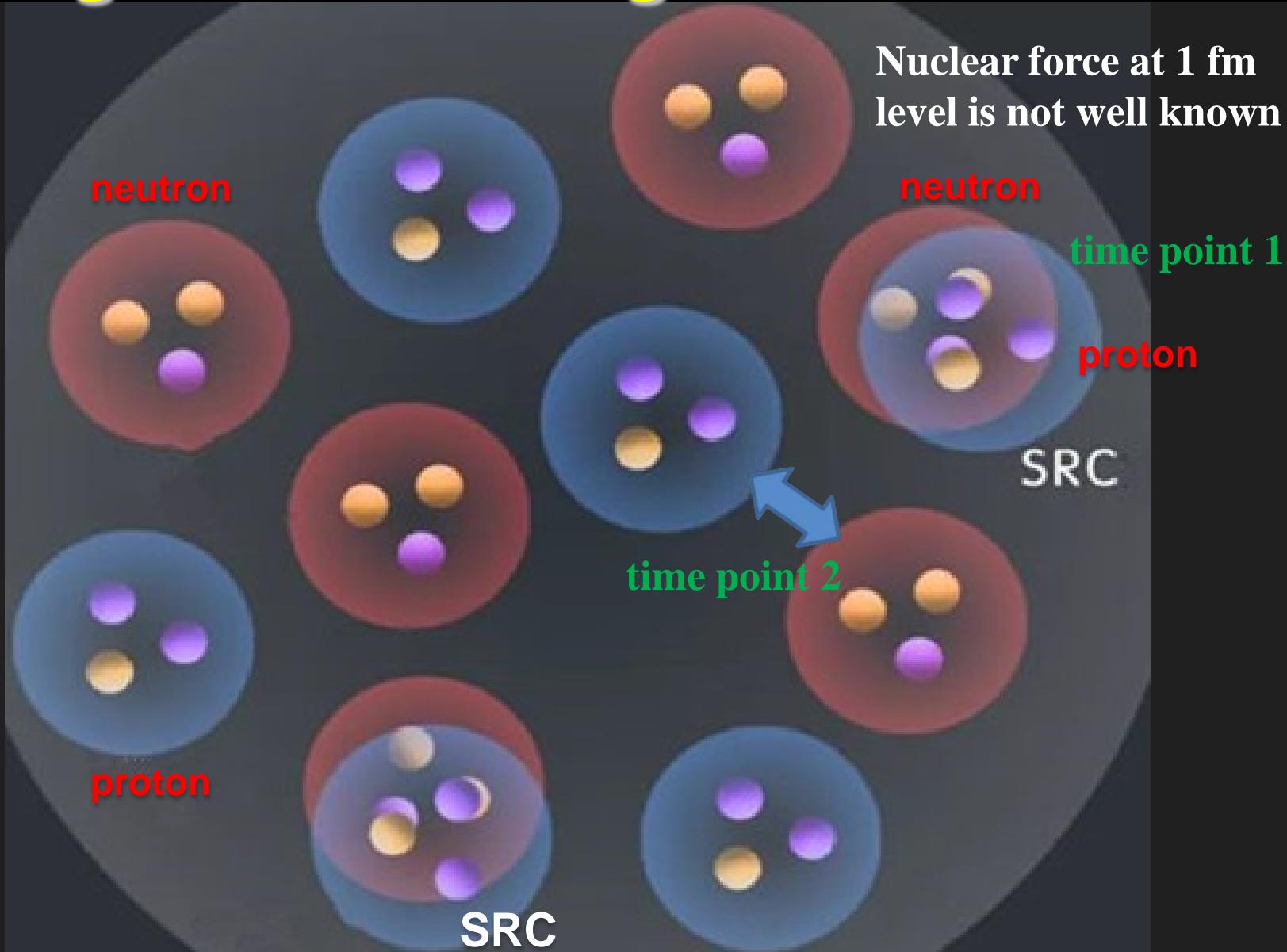
$$\frac{n(k)}{\text{nucleus}} = \frac{1}{N, Z} \int_0^{r_{\max}} d^3 r \rho(r)_{n,p} n[k, K_{F_{n,p}}(r)]$$

Nucleus = many Spherical shells

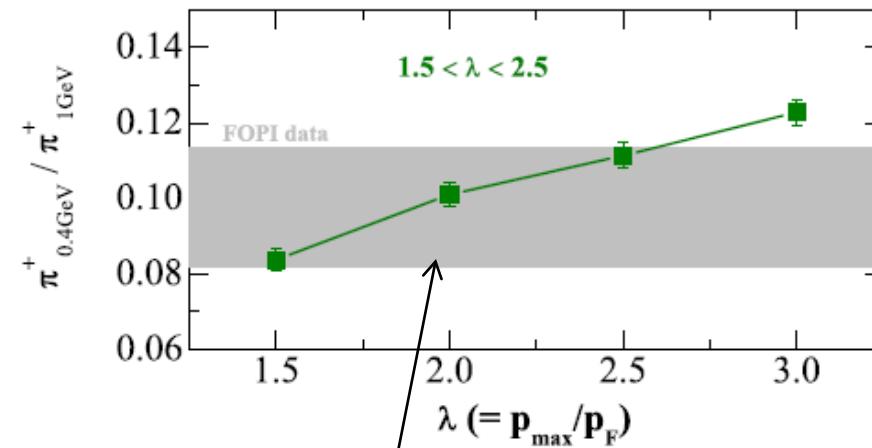
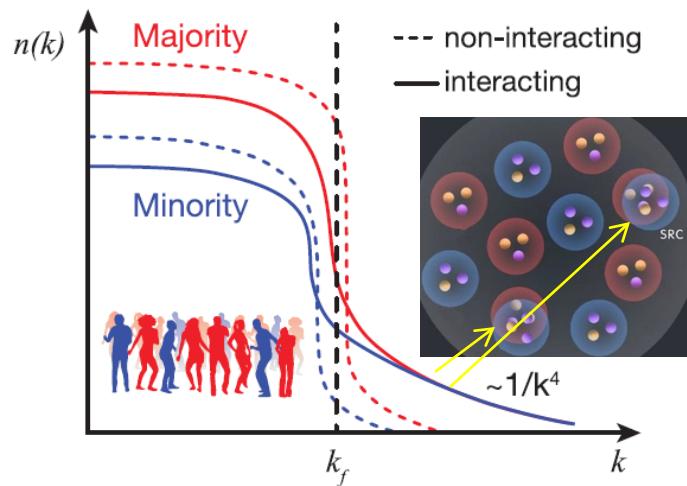


Nuclear spherical shell

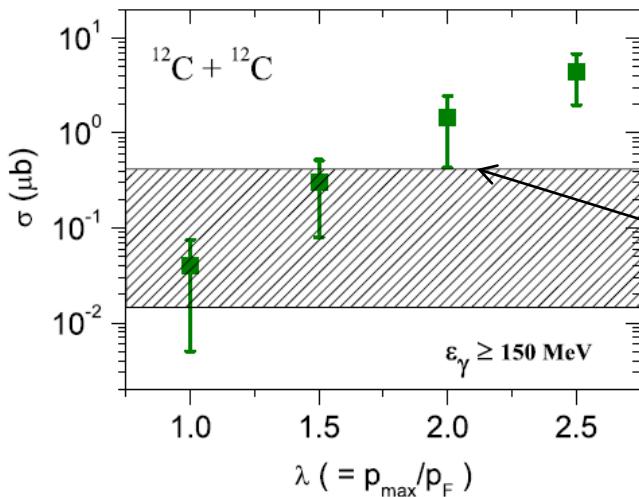
Imagined SRC-image of nucleus



Constrain nucleon motion from HIC



O. Hen et al., Science 346, 614 (2014).



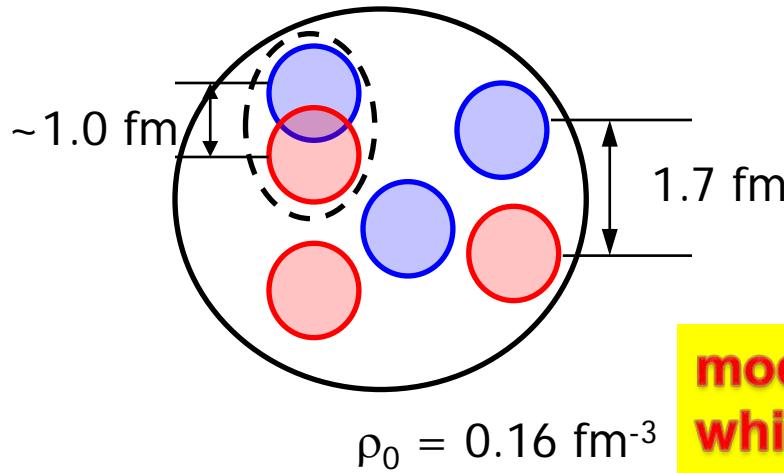
In nuclei:
There is a HMT cutoff

$$p_{\max} \leq 2 p_F$$

$$p_{\max} \leq 2.75 p_F \Leftarrow \text{Exp.} \leftarrow \text{deuteron}$$

(2) Mean-field:

Yong, PRC93, 044610 (2016)



$$U(\rho, \delta, \vec{p}, \tau) = A_u(x) \frac{\rho_{\tau'}}{\rho_0} + A_l(x) \frac{\rho_\tau}{\rho_0}$$

$$+ B \left(\frac{\rho}{\rho_0} \right)^\sigma (1 - x\delta^2) - 8x\tau \frac{B}{\sigma + 1} \frac{\rho^{\sigma-1}}{\rho_0^\sigma} \delta \rho_{\tau'}$$

$$+ \frac{2C_{\tau,\tau}}{\rho_0} \int d^3 \vec{p}' \frac{f_\tau(\vec{r}, \vec{p}')}{1 + (\vec{p} - \vec{p}')^2 / \Lambda^2}$$

$$+ \frac{2C_{\tau,\tau'}}{\rho_0} \int d^3 \vec{p}' \frac{f_{\tau'}(\vec{r}, \vec{p}')}{1 + (\vec{p} - \vec{p}')^2 / \Lambda^2}$$

Considering n-p correlations:

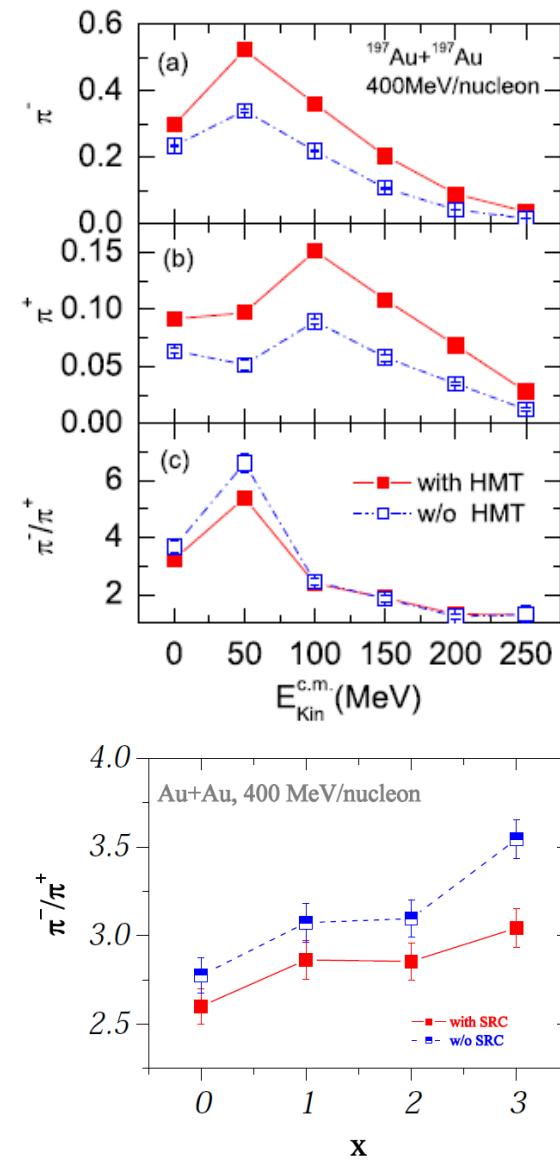
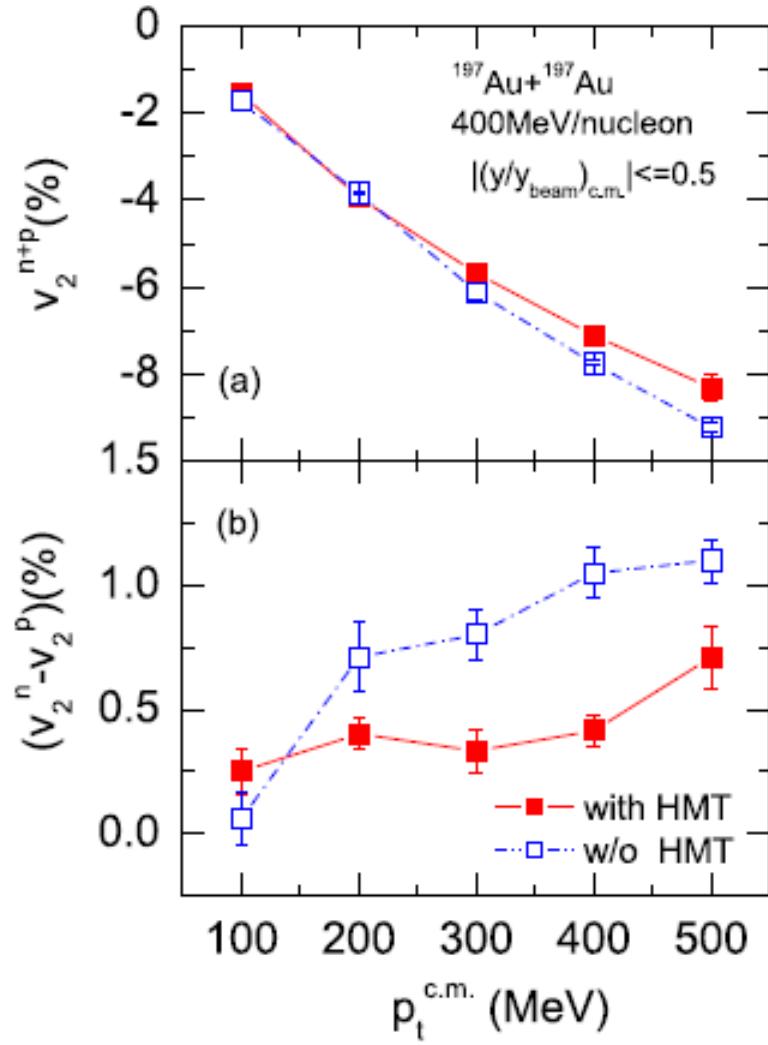
The parameters A, B, C, et al., are readjusted

$$U = U_0 + U_{\text{sym}}$$

$$E_{\text{sym}}^{\text{kin}} \approx 0$$

$$E_{\text{sym}}^{\text{pot}} \approx 31.6(\text{MeV})$$

SRC effects in HIC



Transition momentum:

$$n_p^{\text{HMT}}(k)/n_n^{\text{HMT}}(k) \simeq \rho_n/\rho_p$$

$$\int_{k_F}^{\lambda k_F} n^{\text{HMT}}(k) k^2 dk / \int_0^{\lambda k_F} n(k) k^2 dk \simeq 20\%$$

$$\int_0^{\lambda k_F} n(k) k^2 dk = 1$$

respective – transition – momentum

A)

$$n(k) - n = \begin{cases} C_1, k \leq k_{F_n} \\ \downarrow \\ C_2/k^4, k_{F_n} < k \leq \lambda k_{F_n} \end{cases}$$

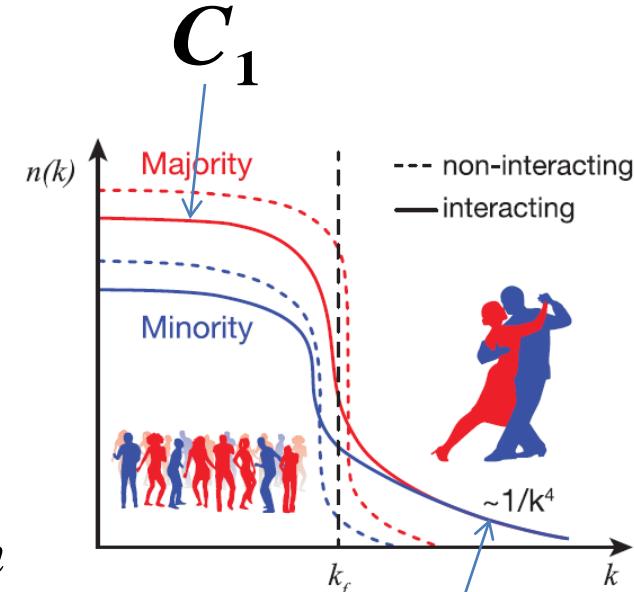
$$n(k) - p = \begin{cases} C_1, k \leq k_{F_p} \\ \downarrow \\ C_2/k^4, k_{F_p} < k \leq \lambda k_{F_p} \end{cases}$$

majority – transition – momentum

B)

$$n(k) - n = \begin{cases} C_1, k \leq k_{F_n} \\ \downarrow \\ C_2/k^4, k_{F_n} < k \leq \lambda k_{F_n} \end{cases}$$

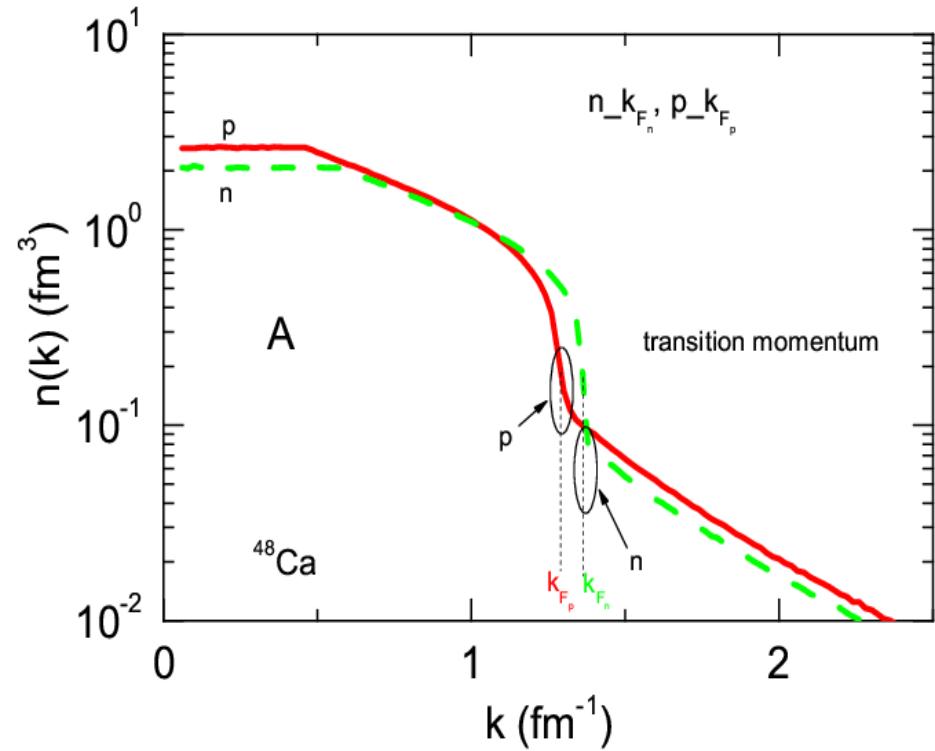
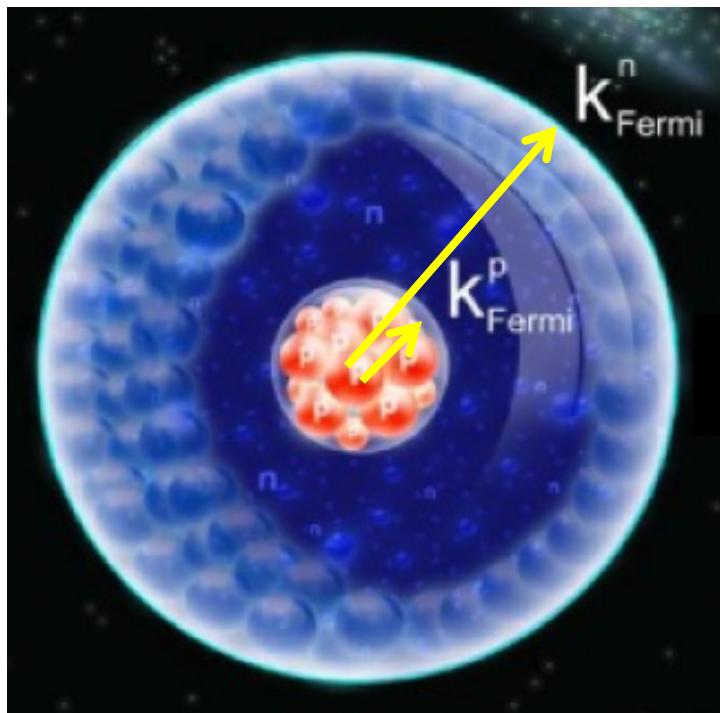
$$n(k) - p = \begin{cases} C_1, k \leq k_{F_p} \\ \downarrow \\ C_2/k^4, k_{F_p} < k \leq \lambda k_{F_p} \end{cases}$$



Transition momentum

Neutron-star matter

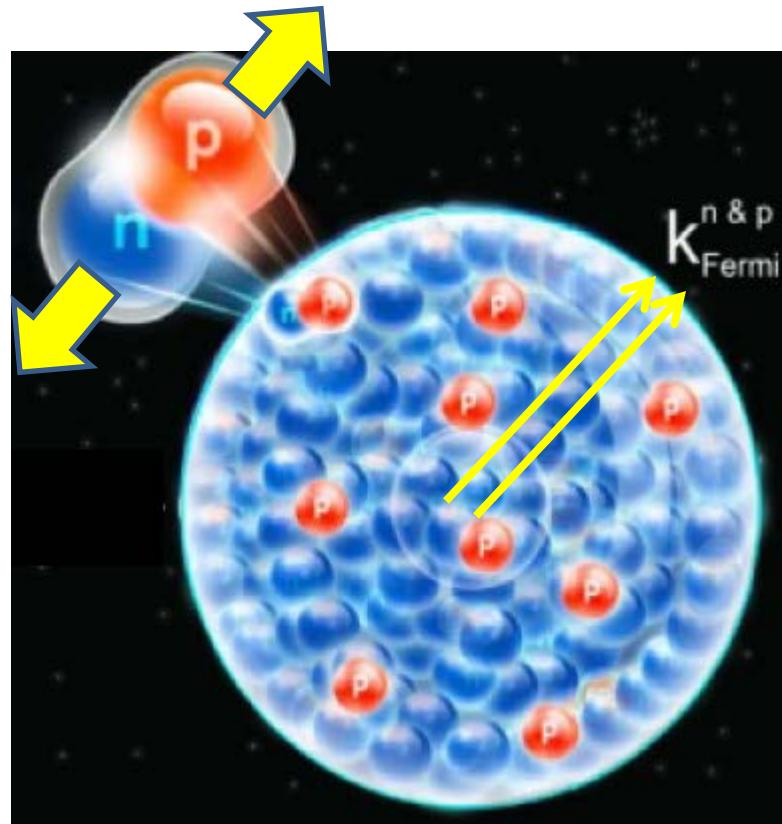
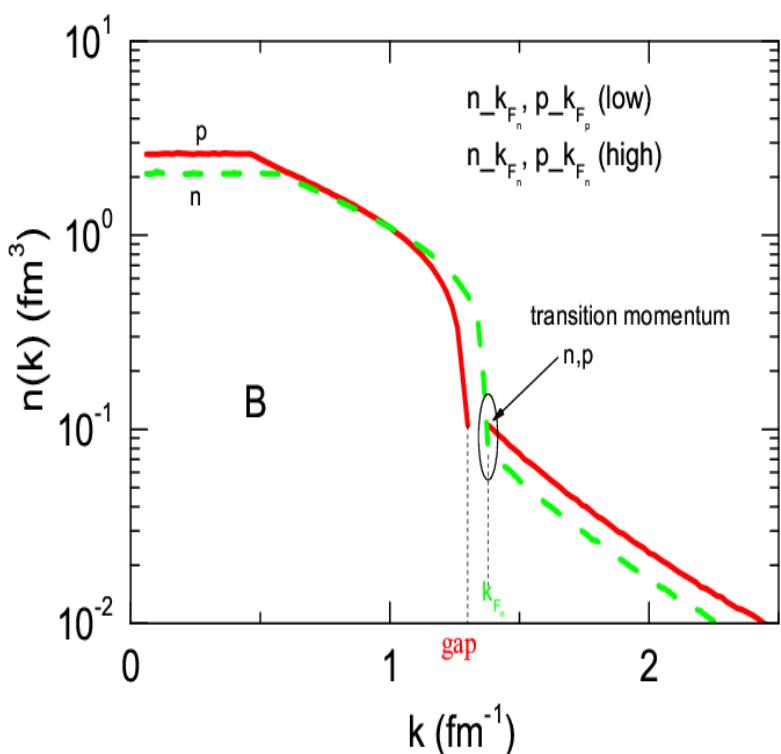
$$k_{F_{n,p}}(r) = [3\pi^2 \hbar^3 \rho(r)_{n,p}]^{\frac{1}{3}}$$



k_{proton} : not matched

Transition mom. = Respective Fermi mom.

Neutron-star matter

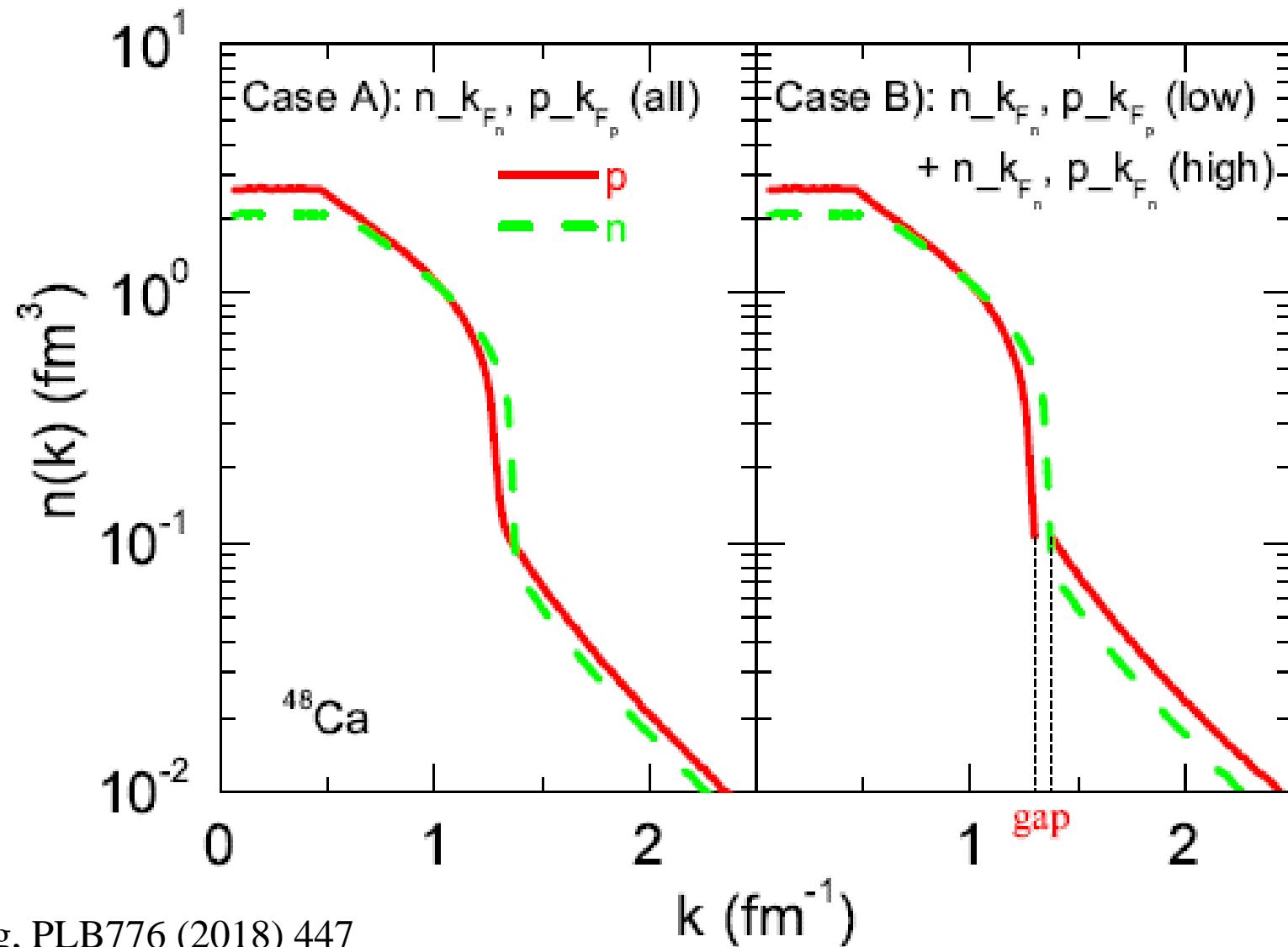


k_proton: matched

Minority transition mom.
= majority transition mom.

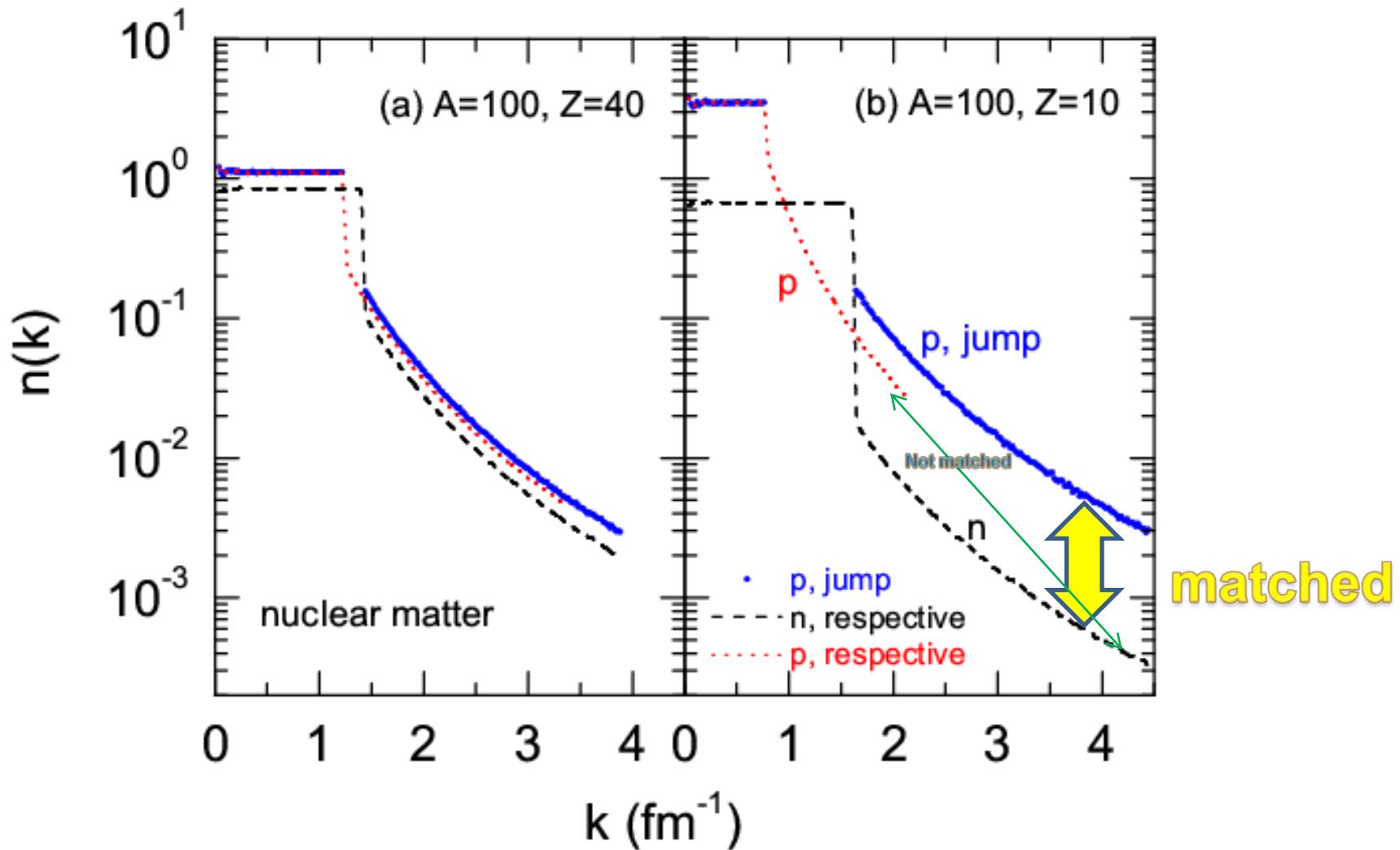
Which one is correct?

$$nucleus : 20\% \times (1 - \delta^2)$$

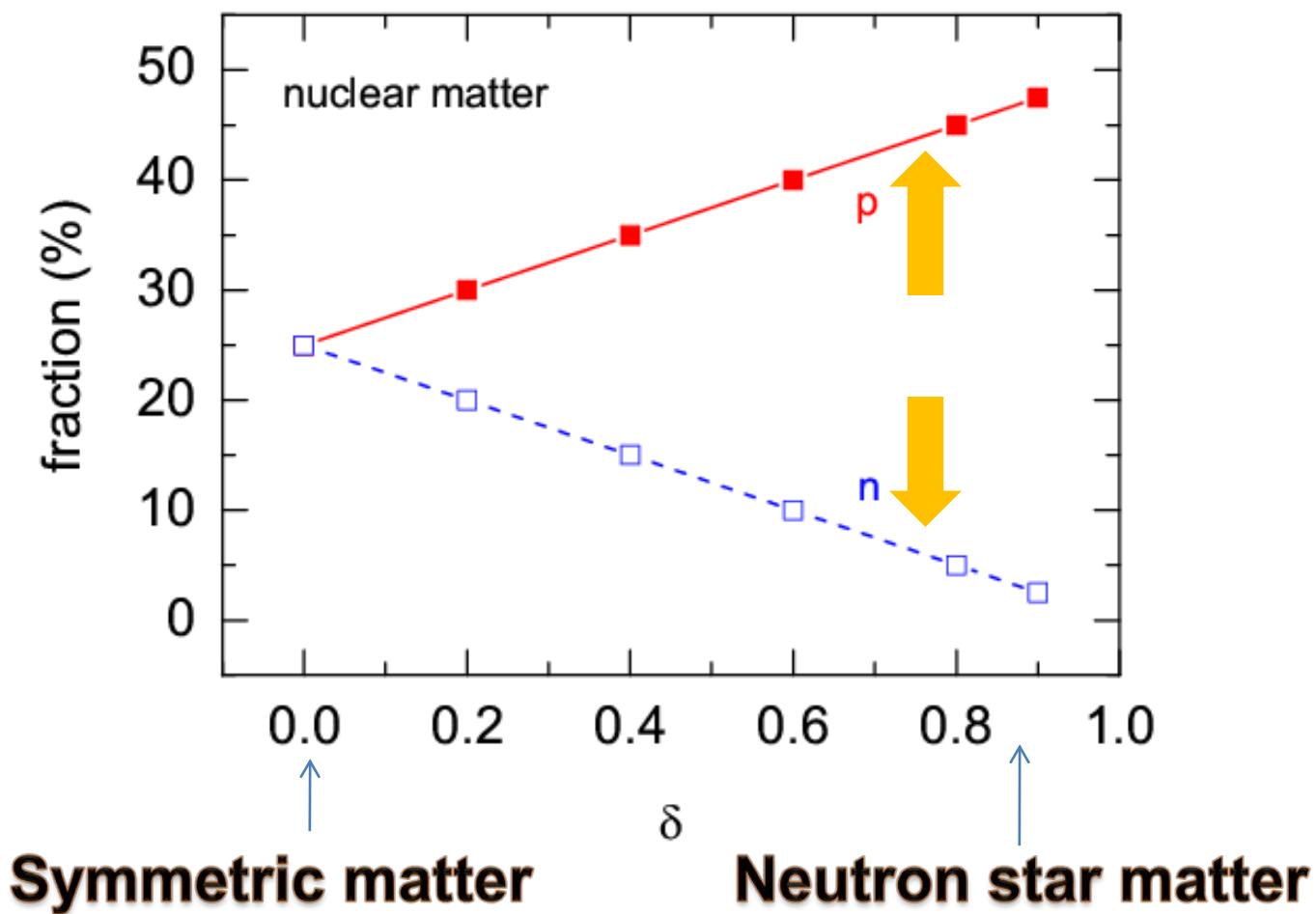


Which one is correct?

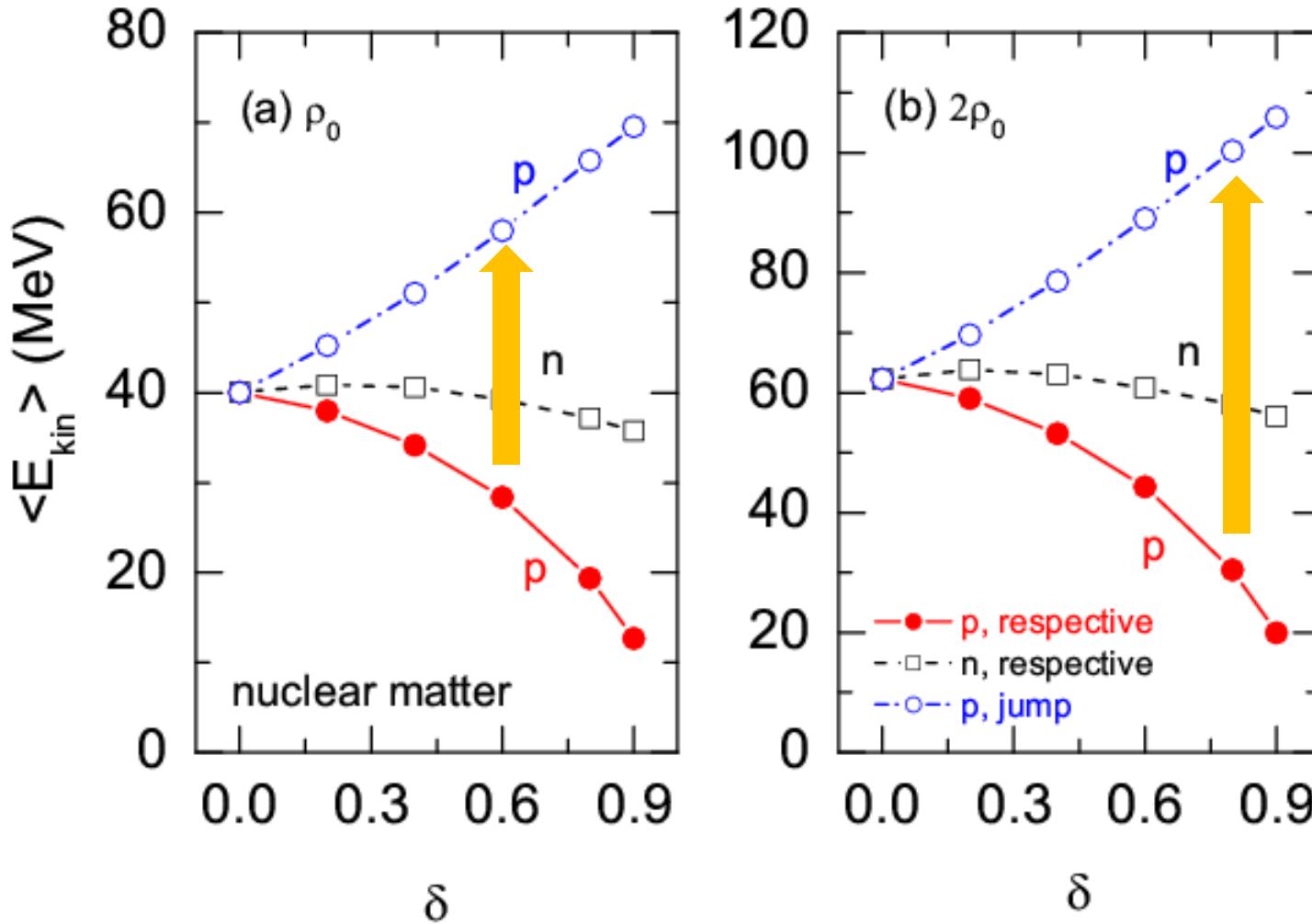
$$\text{matter} : 25\% \times (1 - \delta^2)$$



Fraction of SRC-nucleon

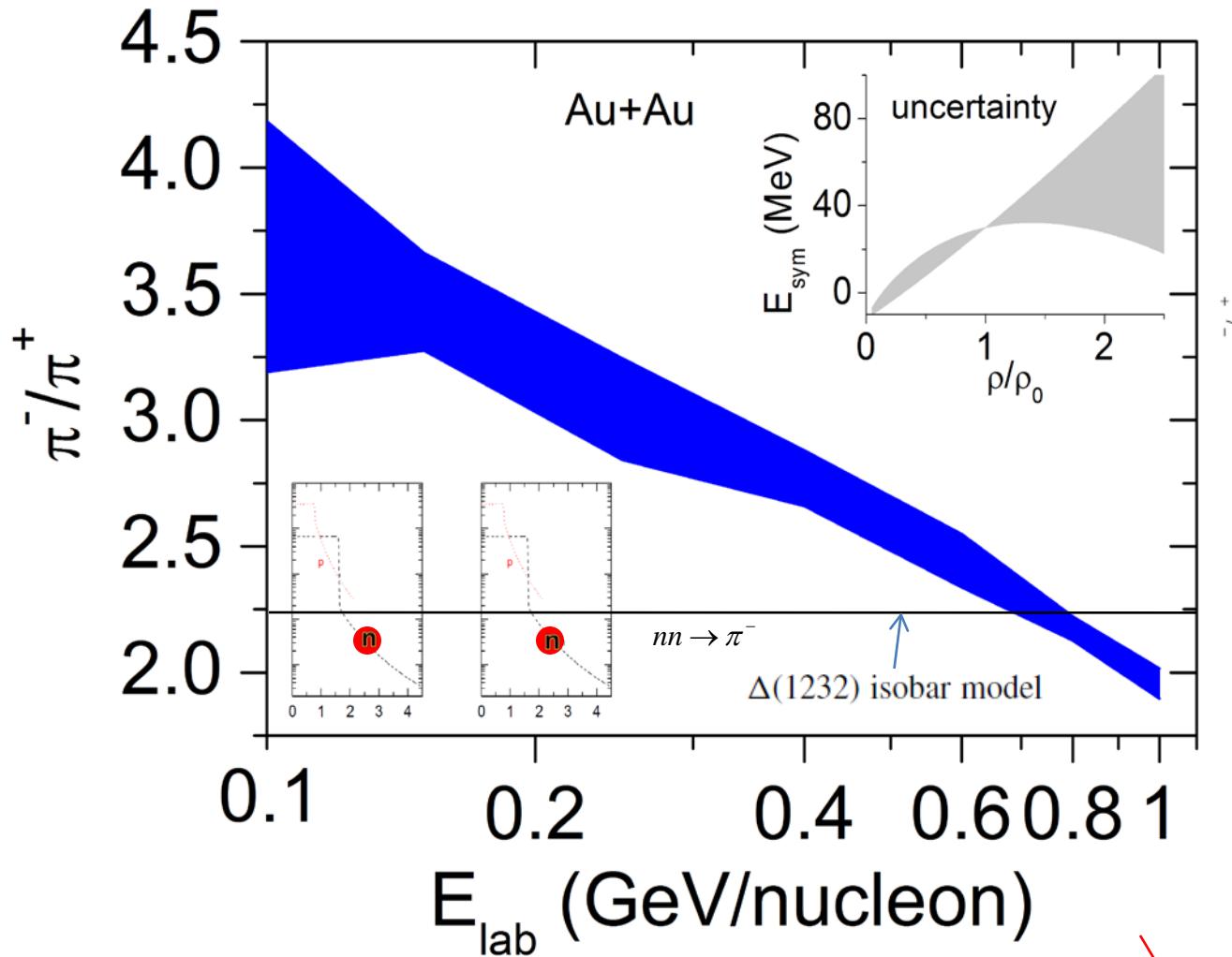


On Nucleon kinetic energy



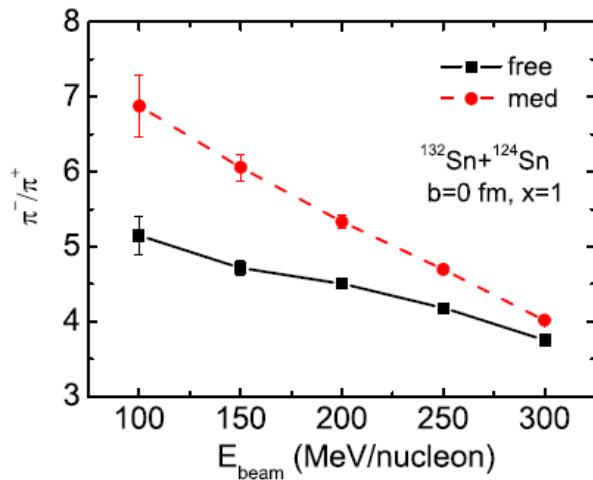
Proton ekin is enhanced significantly

SRC effects on pion ratio in HIC



$\Delta(1232)$ isobar model R. Stock, Phys. Rep. 135, 259 (1986)

$$R_{\text{isob}} \equiv (\pi^-/\pi^+)_{\text{res}} \equiv (5N^2 + NZ)/(5Z^2 + NZ) \approx (N/Z)_{\text{dense}}^2$$

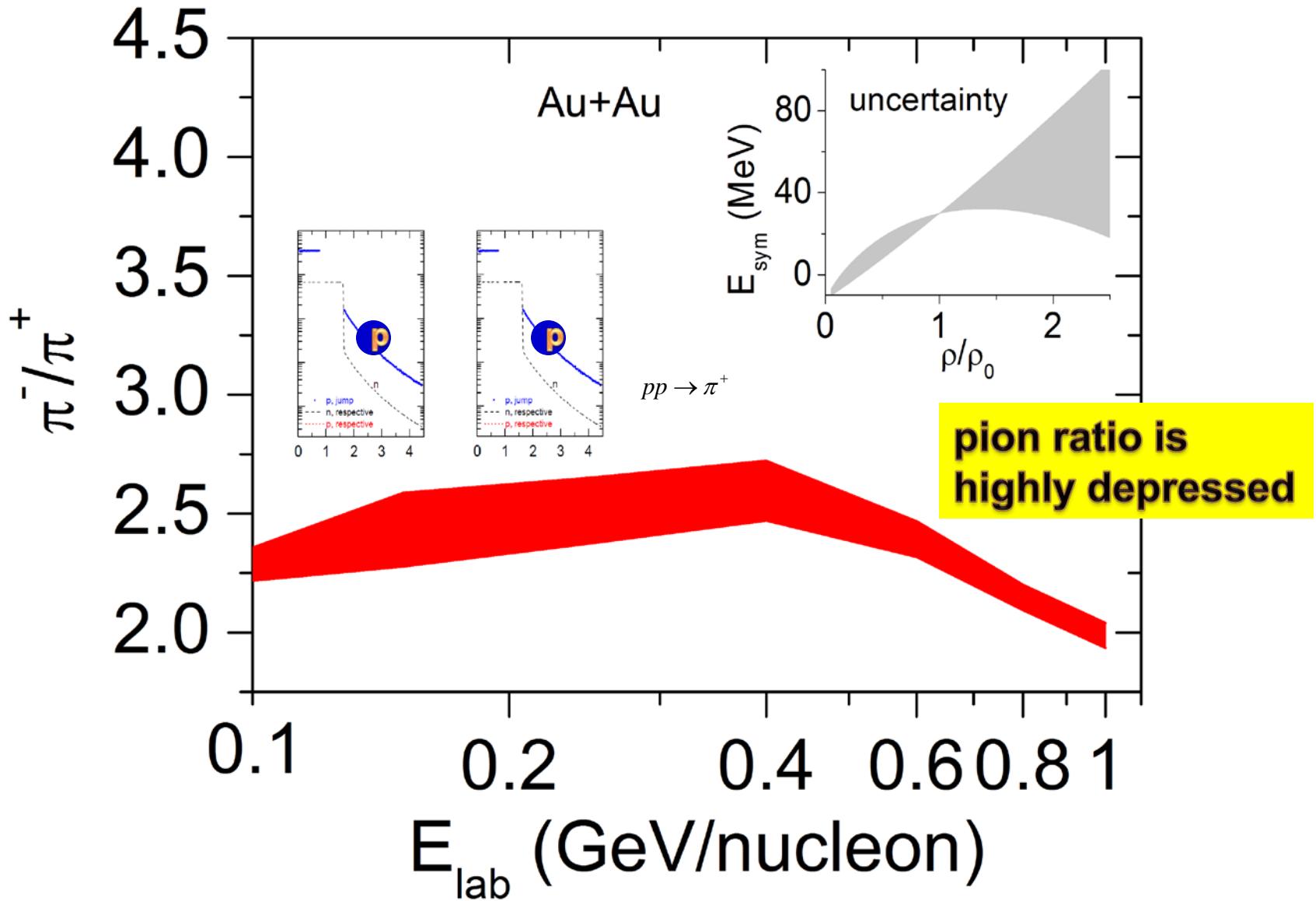


Without SRC

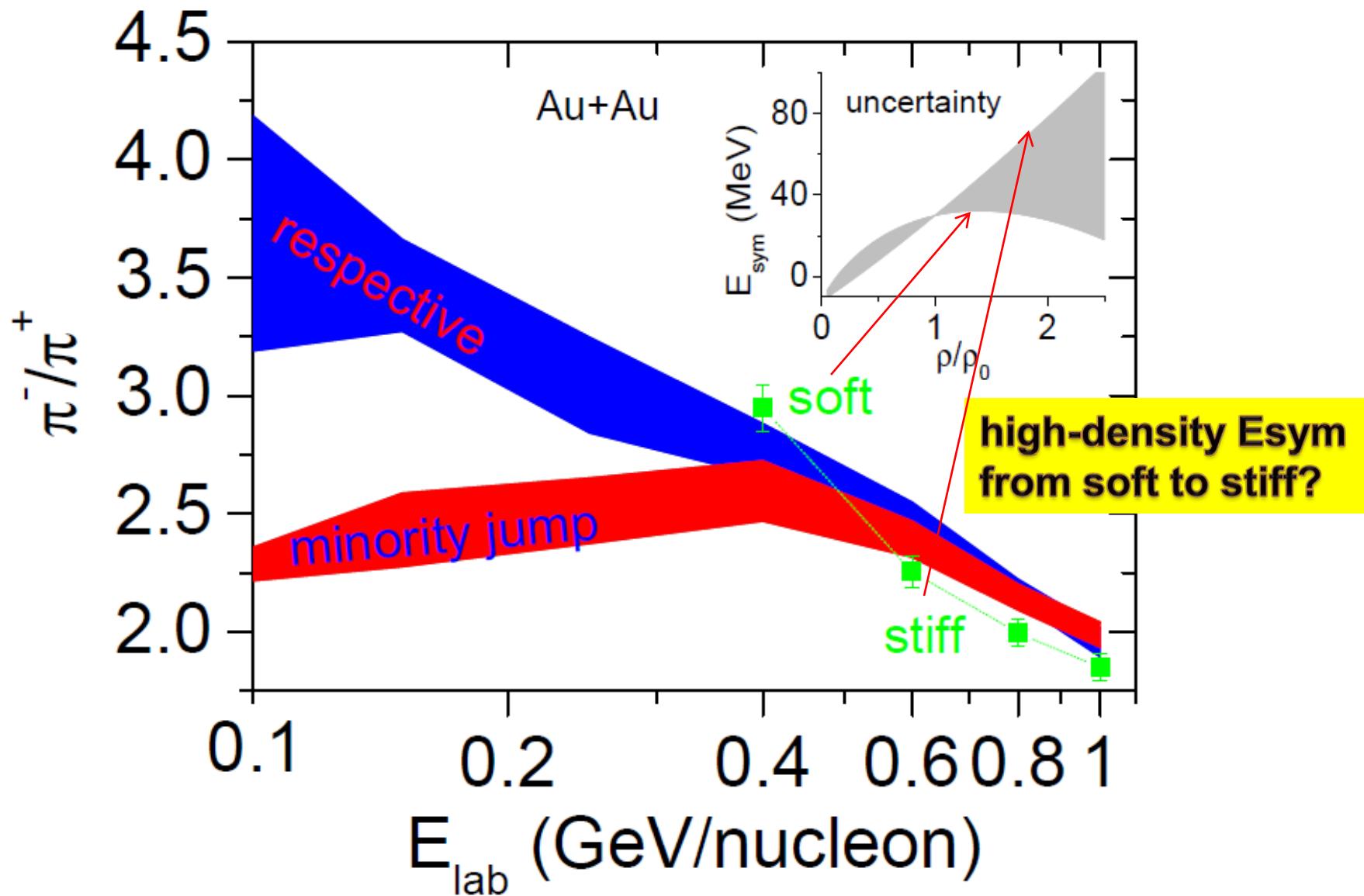
Guo, Yong, Zuo,
PRC90, 044605(2014)

**Isobar model
is always wrong**

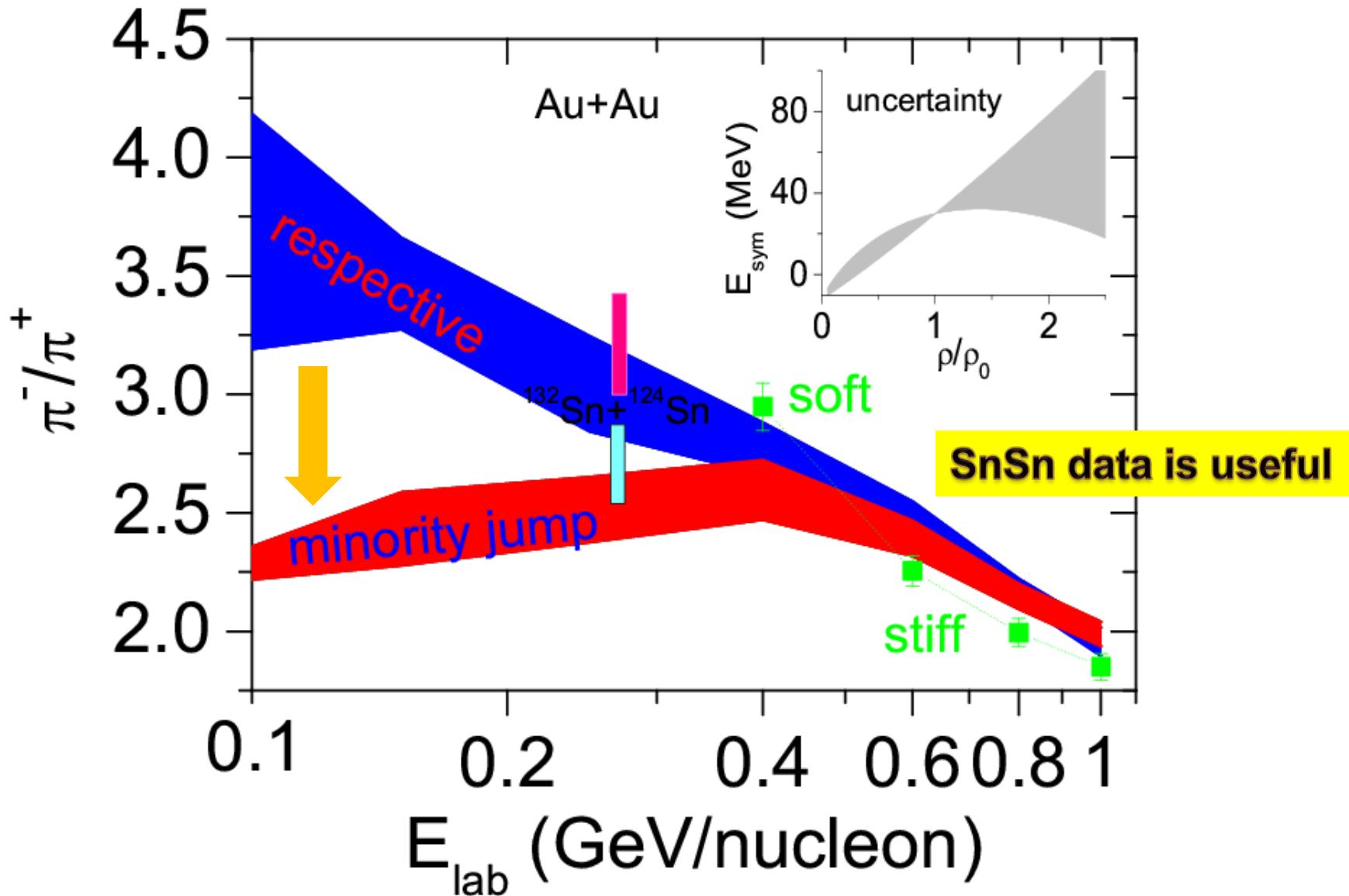
SRC effects on pion ratio in HIC



SRC effects on pion ratio in HIC



SRC effects on pion ratio in HIC



Summary&Comments

- It is not possible that SRC does not affect Esym sensitive probes.
- The picture of SRC is not clear (most many-body approaches fails to fit experimental SRC data) although there are some solid data.
- To check the SRC effects, 100-400 MeV beam energy HICs are favored.
- A perfect SRC consideration in the transport model seems hard to do, but we can first consider the main factors.
- Beam energy scans from 100-600 MeV beam energy n-rich HIC, more data information (such as pion-,pion+, their mom. Distribution, etc.) are welcomed (everyone can fit well if the experimental information is less).

Welcome your comments or suggestions !