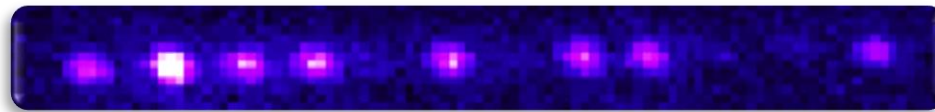


Probing new physics with precision isotope shift spectroscopy



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מכון ויצמן למדע

WEIZMANN INSTITUTE OF SCIENCE

AMO Precision Measurements

- Variation of fundamental constant
- eEDM measurements
- Atomic parity violation
- Test of local Lorentz invariance
- Testing general relativity with atom interferometry
- Probing\bounding new light force-mediators by isotope shift spectroscopy

C. Delaunay, R.Ozeri, G.Perez and Y. Soreq Phys. Rev. D 96, 093001 (2017)

Berengut J. C.; Budker D.; Delaunay C.; Flambaum V. V.; Frugiuele C.; Fuchs E.; Grojean C.; Harnik R.; RO; Perez G.; Soreq Phys. Rev. Lett. 120, 091801 (2018)

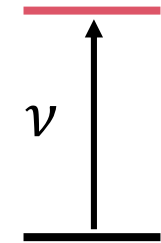
Isotope Shifts

hypothetical new force carriers

A boson that couples to electrons and neutrons

$$V_{\phi}(r) = \frac{-(-1)^s y_e y_n (A - Z) e^{-m_{\phi} r}}{4\pi r}$$

Atomic transition

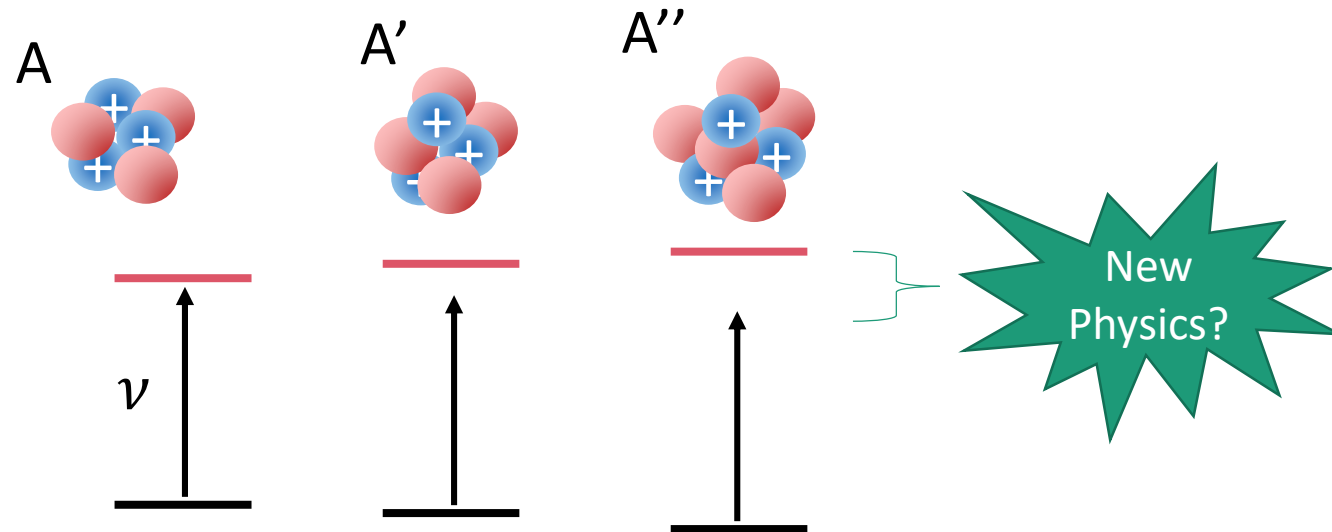

$$\delta\nu/\nu = 10^{-18}$$

An over simplified picture :

Measure different isotope

$$\delta\nu_i^{AA'} \equiv \nu_i^A - \nu_i^{A'}$$

There are **normal** contributions :
Mass shift and **Field shift**



Isotope Mass Shift

Due to the kinetic energy of the nucleus

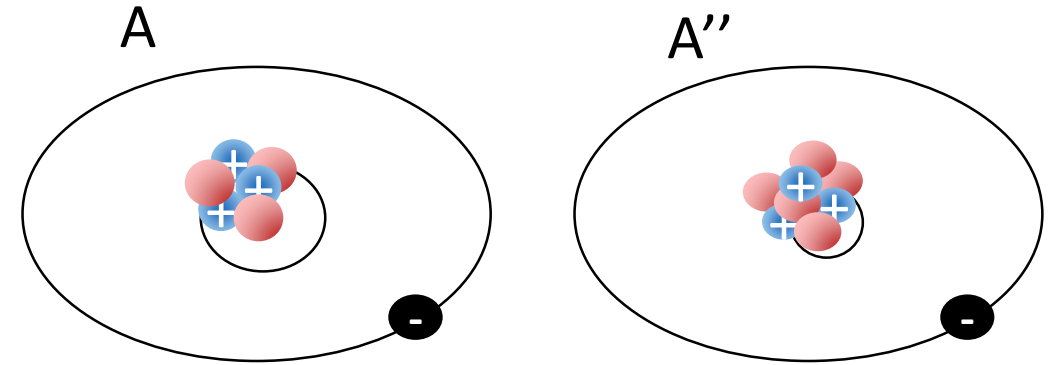
$$\frac{P_n^2}{2M_n} = \frac{(\sum p_e)^2}{2M_n} = \frac{\sum_i p_i^2 + \sum_{i \neq j} p_i p_j}{2M_n}$$

Normal Mass Shift
(Reduced mass)

$$\delta\nu_{NMS}^{AA'} \sim \frac{1000}{A^2} \text{GHz}$$

Specific mass shift

- correlation in the electrons many body wave function.
- Very hard to calculate -> Unknown



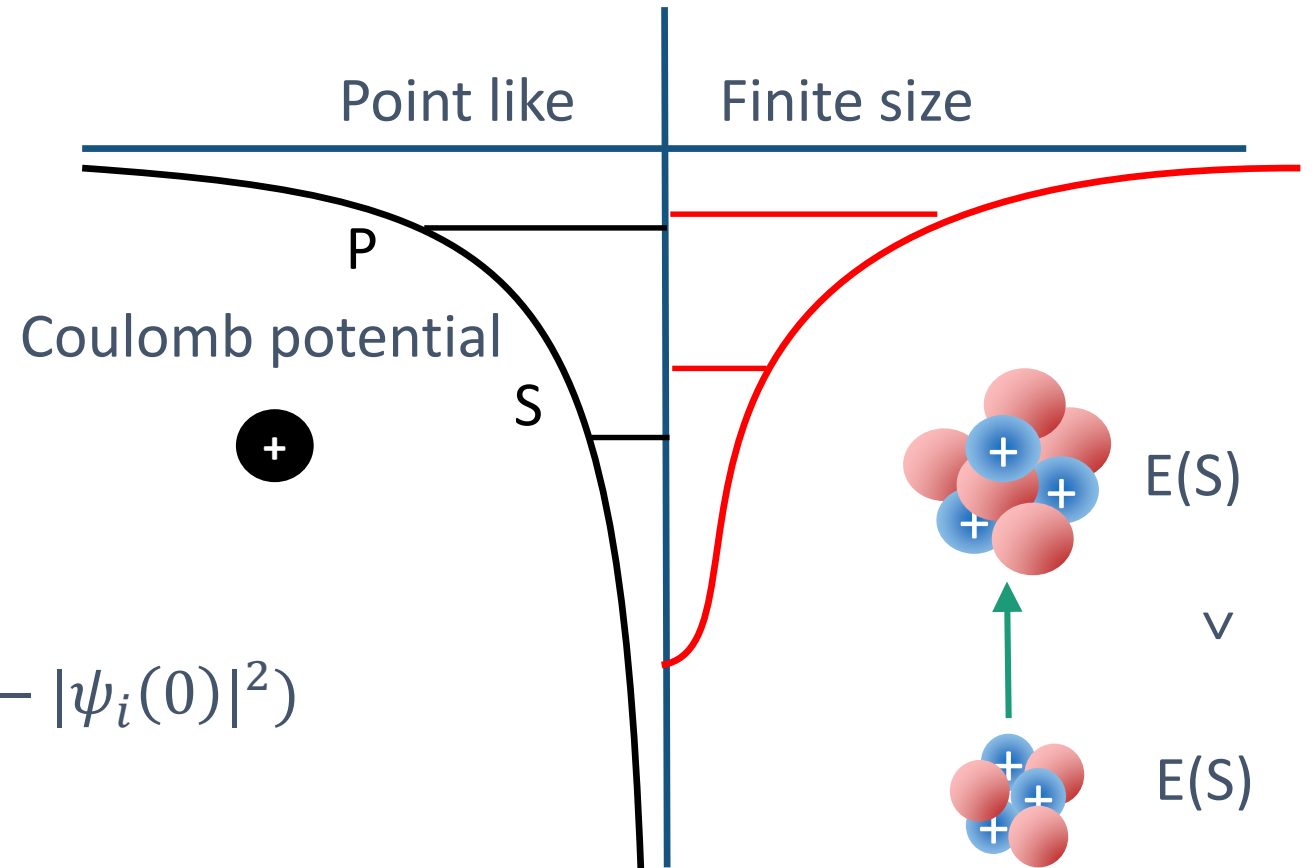
Isotope Field Shift

Due to the finite size of the nucleus
and Charge distribution

Leading order :

$$\delta \nu_{FS}^{AA'} \propto \delta \langle r_c^2 \rangle^{AA'} (|\psi_f(0)|^2 - |\psi_i(0)|^2)$$

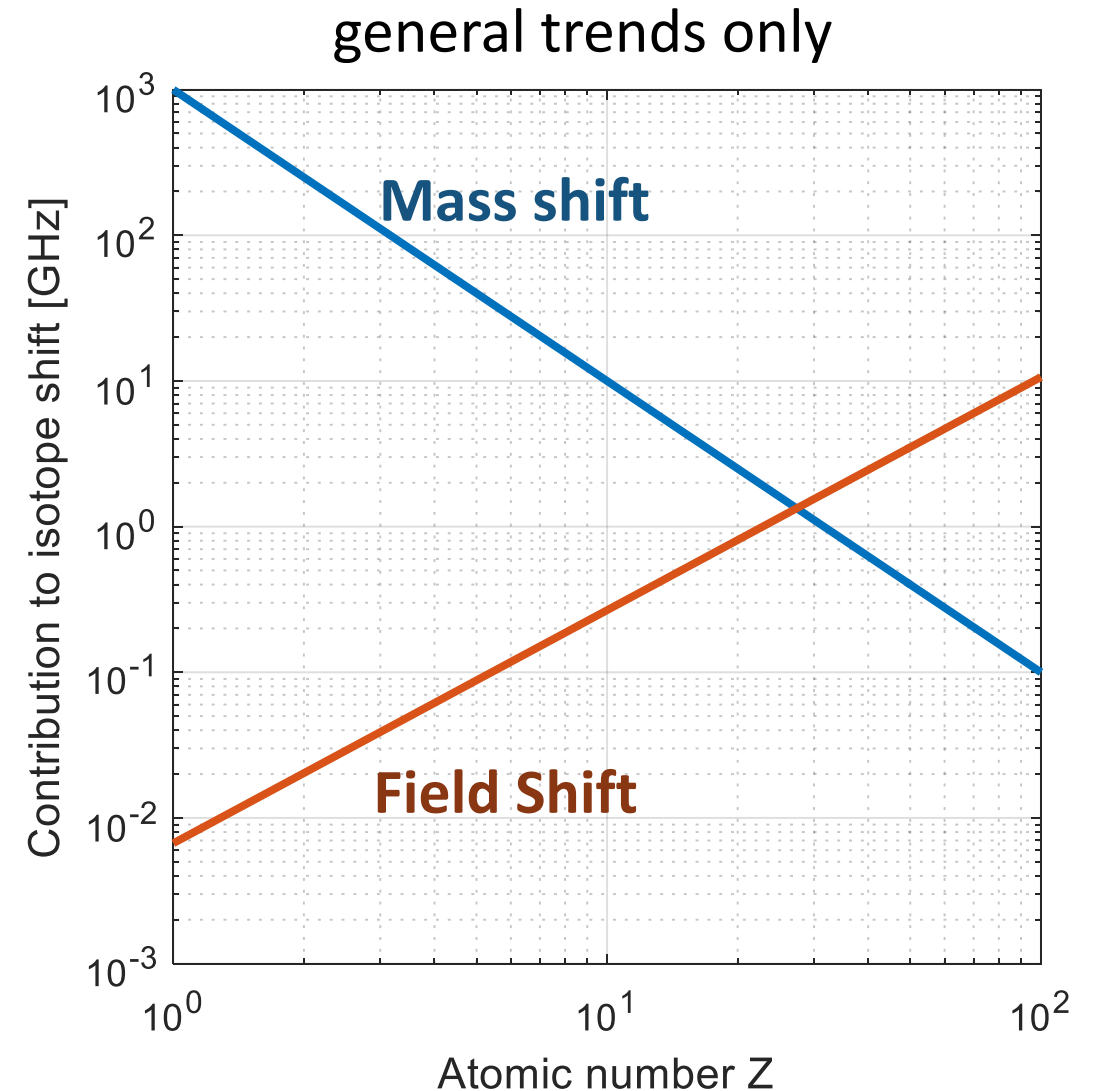
$$\delta \nu_{FS}^{AA'} \sim A^{5/3}$$



Mass shifts and Field Shift

$$\delta\nu^{AA'} = \delta\nu_{MS}^{AA'} + \delta\nu_{FS}^{AA'}$$

- Mass shift dominates in light atoms
- Field shift dominates in heavy atoms
- IS on the order of GHz for $A > 10$
- **Theoretical uncertainty is still poor**



King plot comparison (to the rescue...)

$$\delta\nu^{AA'} = \delta\nu_{MS}^{AA'} + \delta\nu_{FS}^{AA'}$$

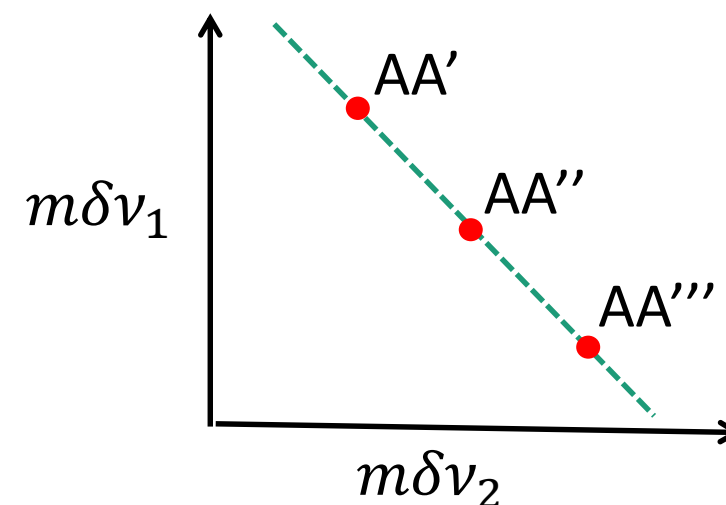
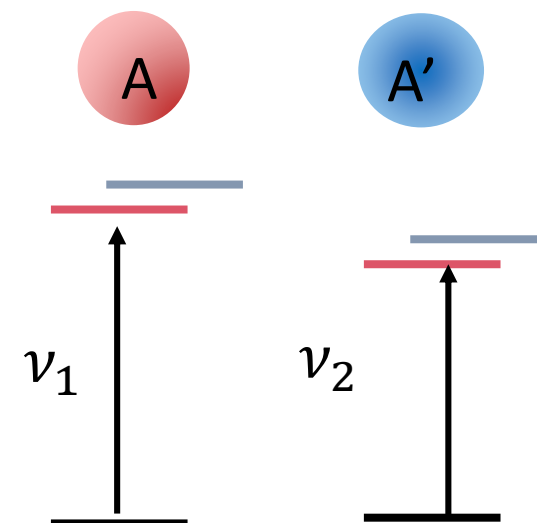
Following King's factorization :

$$\delta\nu_i^{AA'} = K_i\mu + F_i\delta\langle r^2\rangle^{AA'} + X_i\gamma_{AA'}$$

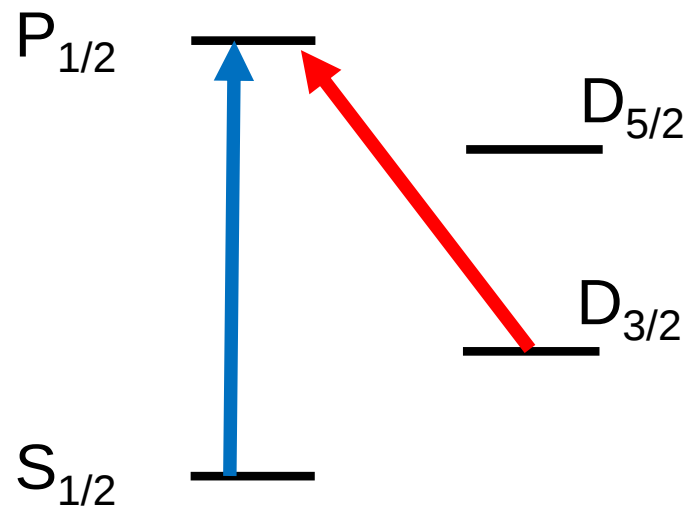
For two transitions $\delta\nu_1^{AA'}$ and $\delta\nu_2^{AA'}$:

$$m\delta\nu_2^{AA'} = K_{21} + F_{21}m\delta\nu_1^{AA'}$$

$$\text{with } m \equiv \frac{M^A M^{A'}}{M^A - M^{A'}} \quad ; F_{21} \equiv \frac{F_2}{F_1} \quad ; K_{21} \equiv K_2 - F_{21}K_1$$

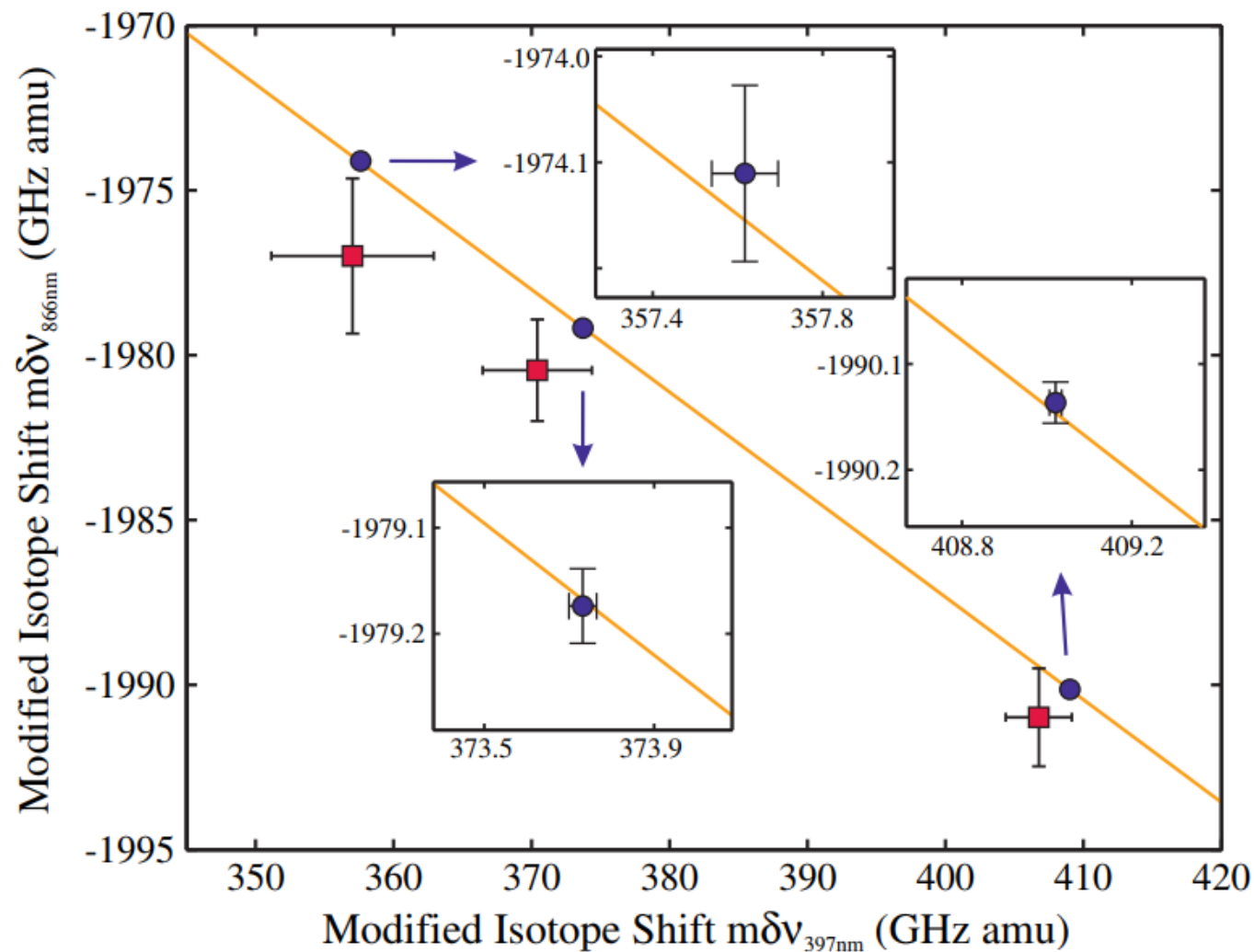


King's plot for **dipole** transitions in Ca^+



Linear at the level of 100 kHz

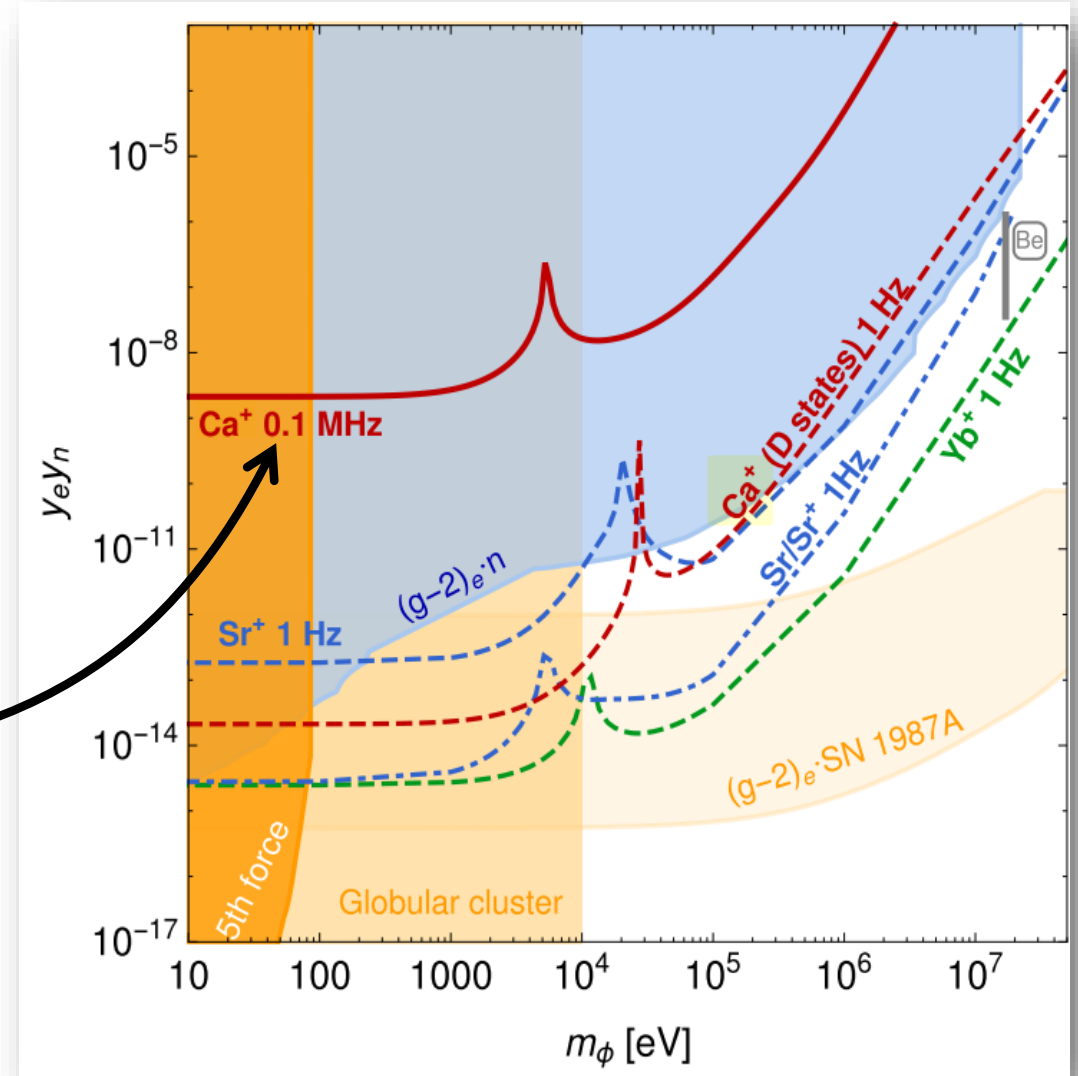
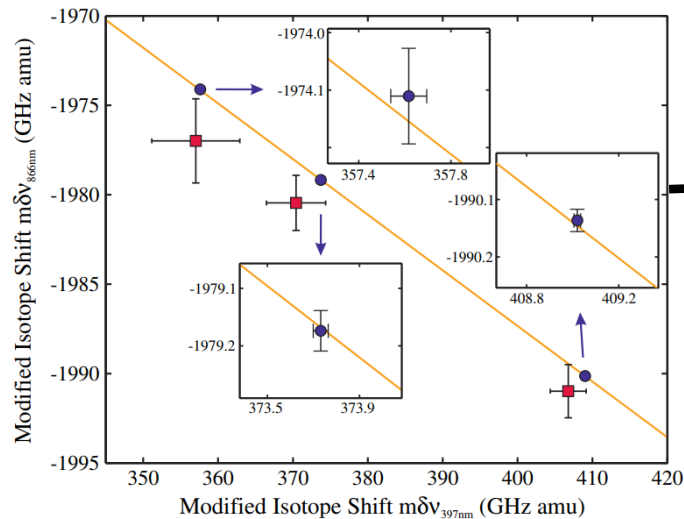
Gebert et al PRL 115, 053003 (2015)



Bounds on new force-mediators

For hypothetical new force carriers of spin $s = 0, 1$ or 2 and mass m_ϕ :

$$V_\phi(r) = \frac{-(-1)^s y_e y_n}{4\pi} \frac{e^{-m_\phi r}}{r}$$



Requirements

- At least four different even isotopes (without hyperfine)
- Two narrow optical transitions (could be neutral and ions)
- Transitions between as different states as possible
- Possible candidates :
 - Ion and neutral : Ca, Yb (Sr)
 - E2 transitions in ions :
Ca⁺, Sr⁺, Ba⁺, Yb⁺
 - E2 and E3 in Yb⁺
- **Small Standard Model nonlinearity**

Nonlinearity in King plots

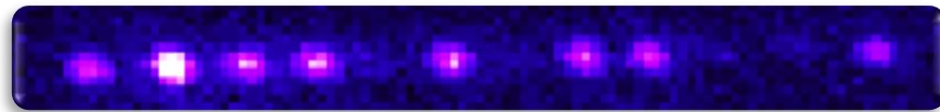
PHYSICAL REVIEW A **97**, 032510 (2018)

Isotope shift, nonlinearity of King plots, and the search for new particles

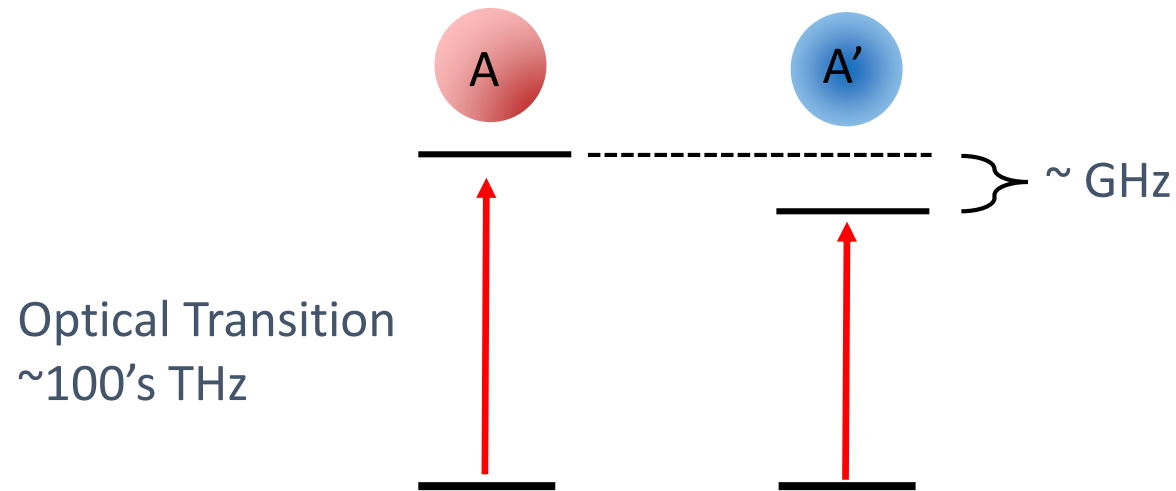
V. V. Flambaum,^{1,2} A. J. Geddes,¹ and A. V. Viatkina²

Ion	Z	A	A_1	A_2	A_3	Pair of transitions	Method 4	Method 5	Nonlinearity (Hz)		QMS
									Quadratic term inc. MB		
									Without α_p	With α_p	
Ca ⁺	20	40	42	44	48	$3p^6 4s^2 S_{1/2} \rightarrow 3p^6 3d^2 D_{3/2}$	3.0×10^{-4}	-6.6×10^{-2}	$\pm 2.9 \times 10^{-3}$	$\pm 2.7 \times 10^{-3}$	± 3.0
						$3p^6 4s^2 S_{1/2} \rightarrow 3p^6 3d^2 D_{5/2}$					
Sr ⁺	38	84	86	88	90	$4p^6 5s^2 S_{1/2} \rightarrow 4p^6 4d^2 D_{3/2}$	1.1×10^{-2}	-2.6	± 0.23	± 0.25	± 9.0
						$4p^6 5s^2 S_{1/2} \rightarrow 4p^6 4d^2 D_{5/2}$					
Ba ⁺	56	132	134	136	138	$5p^6 6s^1 2S_{1/2} \rightarrow 5p^6 5d^2 D_{3/2}$	-3.9×10^{-2}	7.4	∓ 2.0	∓ 1.9	∓ 1.8
						$5p^6 6s^1 2S_{1/2} \rightarrow 5p^6 5d^2 D_{5/2}$					
Yb ⁺	70	168	170	172	176	$4f^{14} 6s^2 S_{1/2} \rightarrow 4f^{13} 6s^2 2F_{7/2}^o$	-3.1	39	± 12260	± 12130	± 28
						$4f^{14} 6s^2 S_{1/2} \rightarrow 4f^{14} 5d^2 D_{3/2}$					
						$4f^{14} 6s^2 S_{1/2} \rightarrow 4f^{14} 5d^2 D_{3/2}$	3.1	-18	± 392	± 386	± 1.1
						$4f^{14} 6s^2 S_{1/2} \rightarrow 4f^{14} 5d^2 D_{5/2}$					
Hg ⁺	80	196	198	200	204	$5d^{10} 6s^2 S_{1/2} \rightarrow 5d^9 6s^2 2D_{3/2}$	3.0	-13	± 2406	± 2382	± 0.38
						$5d^{10} 6s^2 S_{1/2} \rightarrow 5d^9 6s^2 2D_{5/2}$					

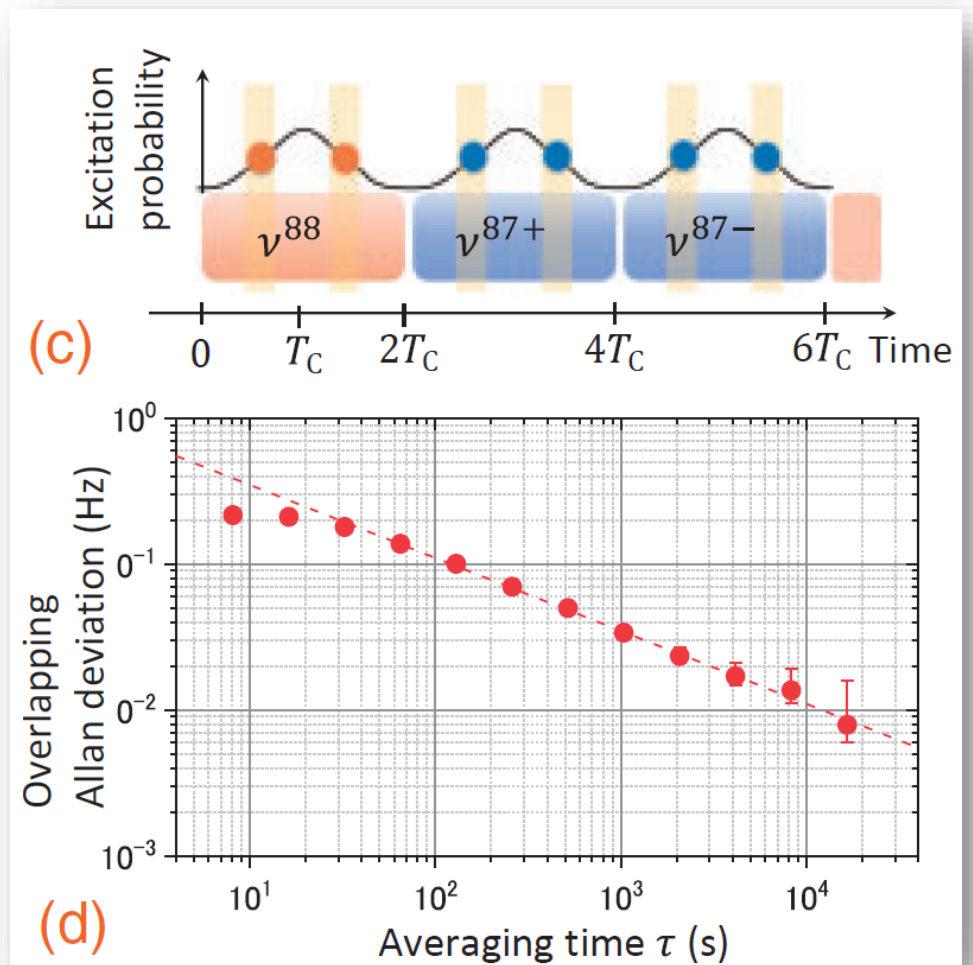
Precision isotope shift spectroscopy in trapped ions



Isotope Shift Spectroscopy

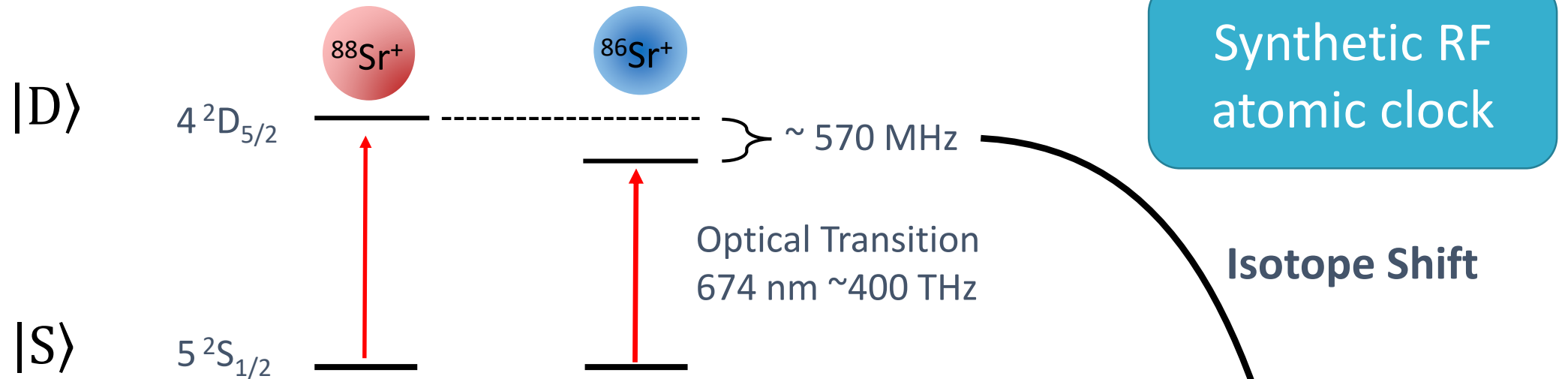


$$\Delta\nu_{S_0, P_0}^{88, 87} = 62,188,134.004(10) \text{ Hz}$$



Isotope Shift Spectroscopy In DFS

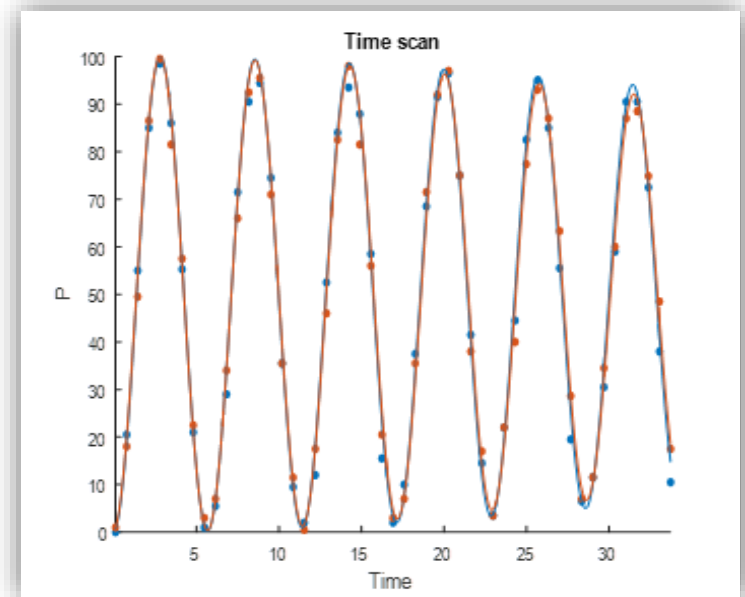
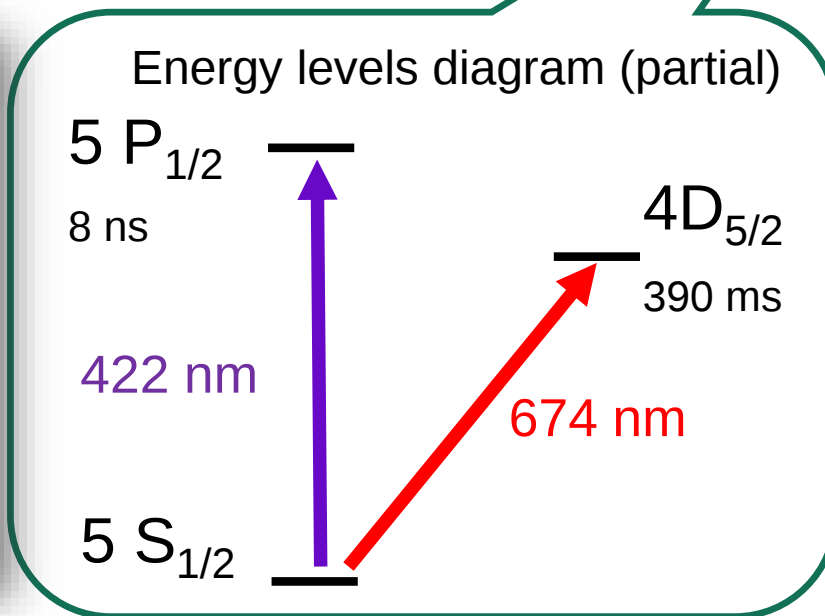
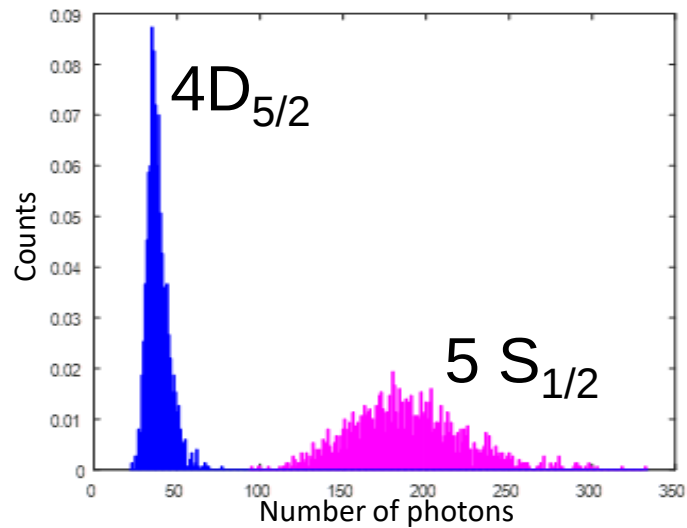
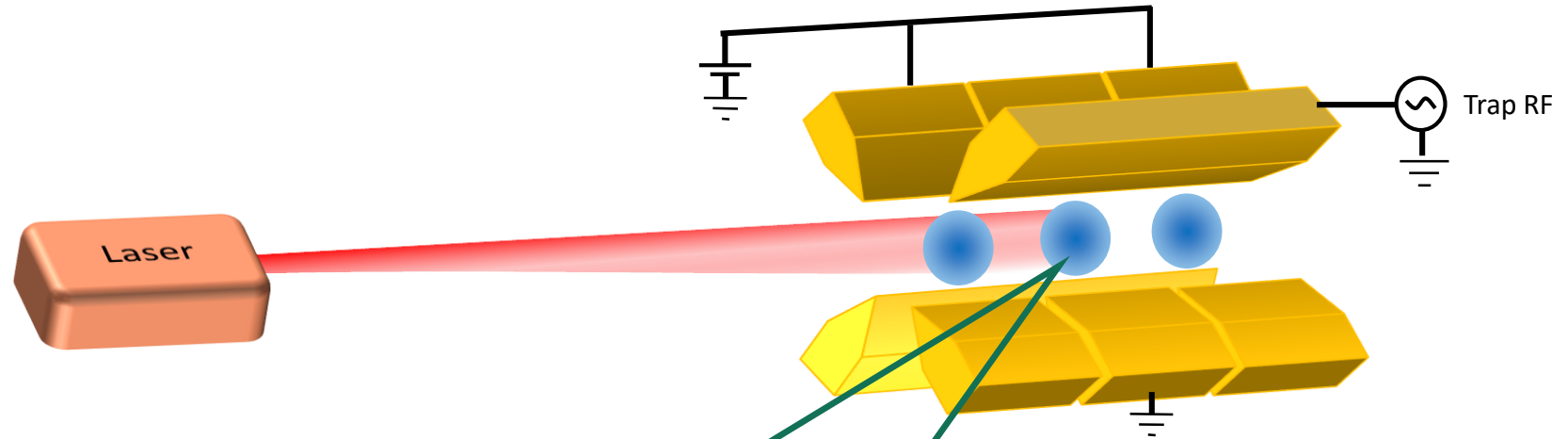
The Quadruple transition in Sr⁺ ion



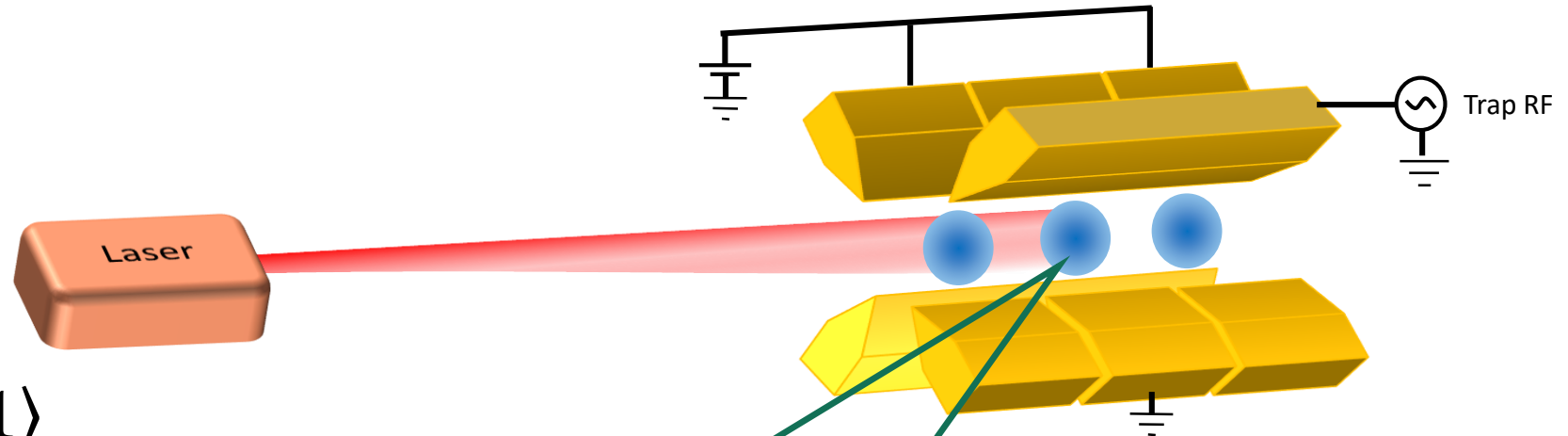
Decoherence Free Subspace

$$\psi_{init} = |S\rangle|D'\rangle + |D\rangle|S'\rangle \rightarrow \psi(T) = |S\rangle|D'\rangle + e^{i2\pi\delta\nu T}|D\rangle|S'\rangle$$

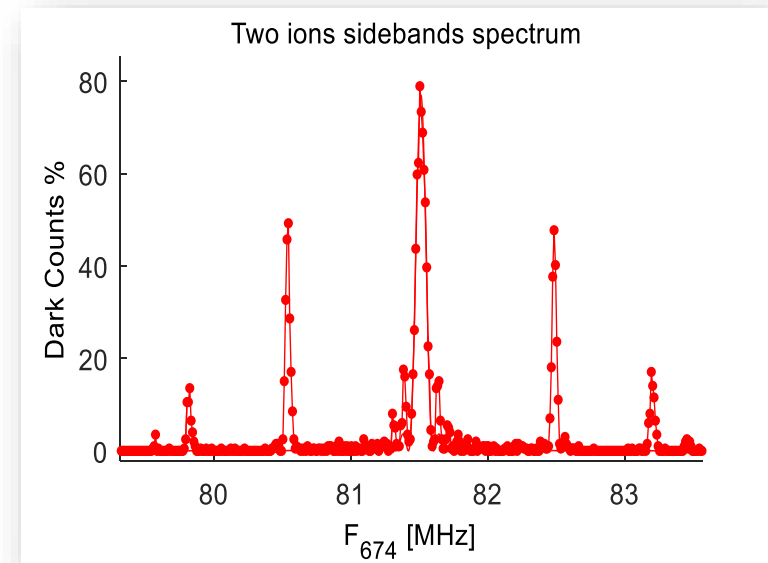
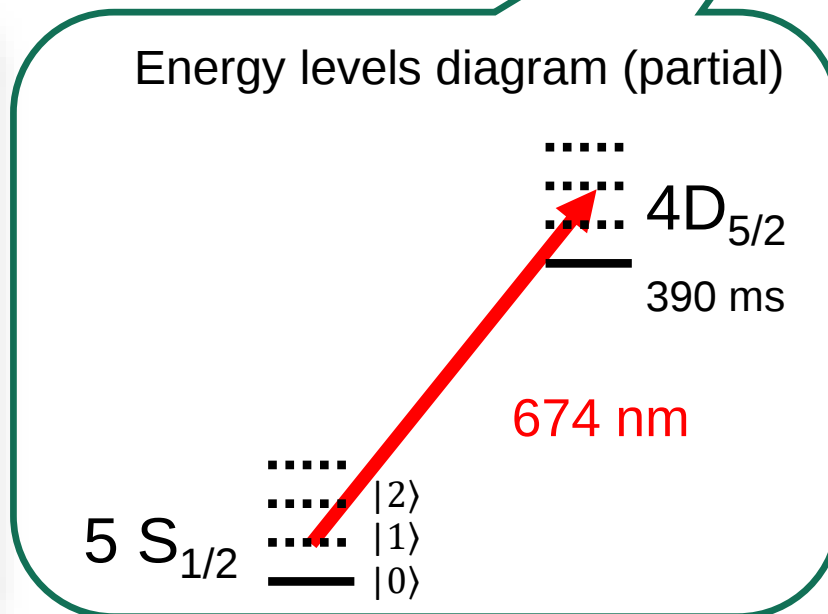
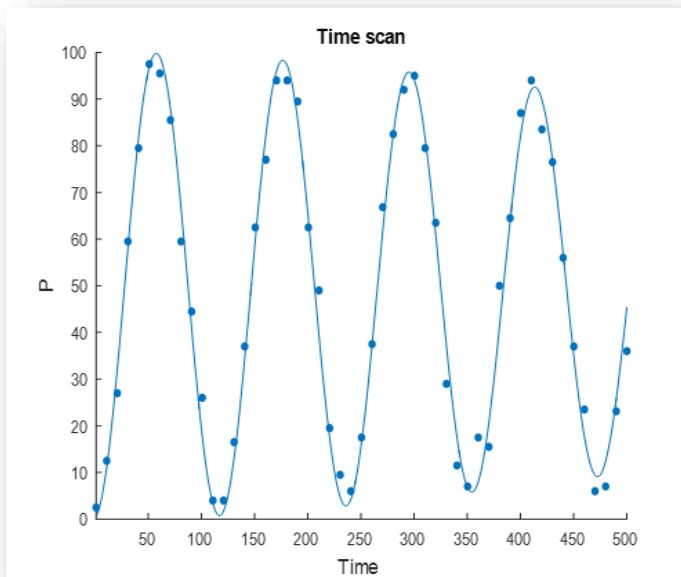
Trapped Ions



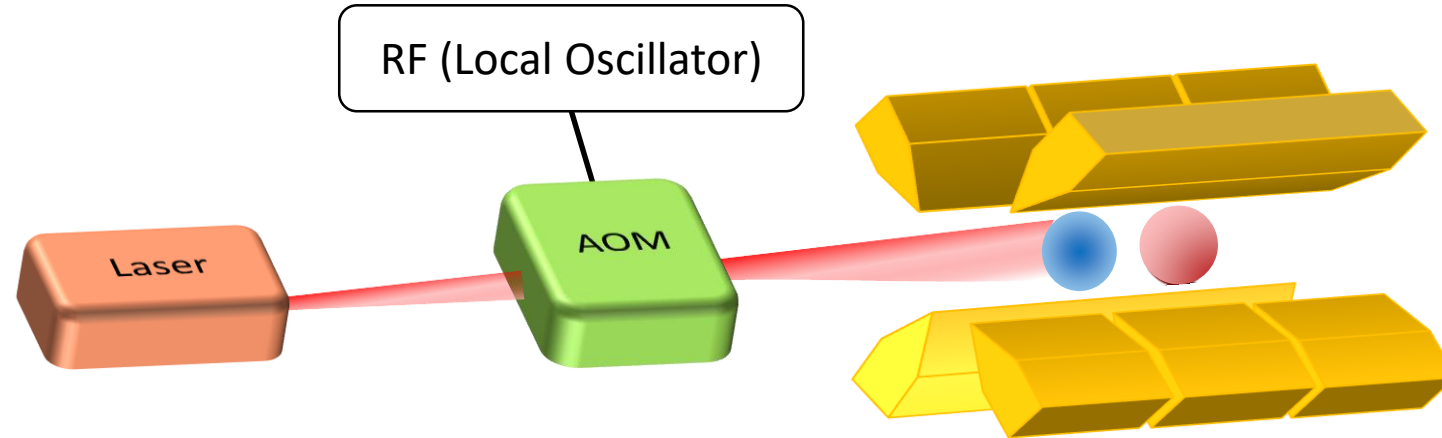
Trapped Ions



$$|S\rangle|0\rangle \leftrightarrow |D\rangle|1\rangle$$

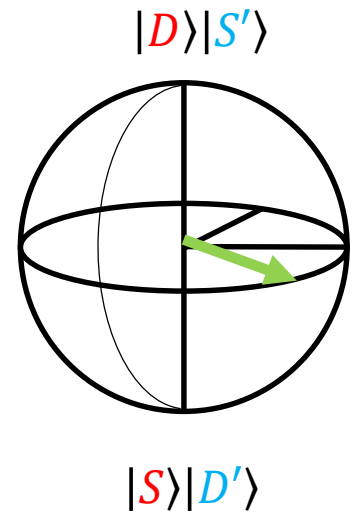


Isotope Shift Spectroscopy In DFS

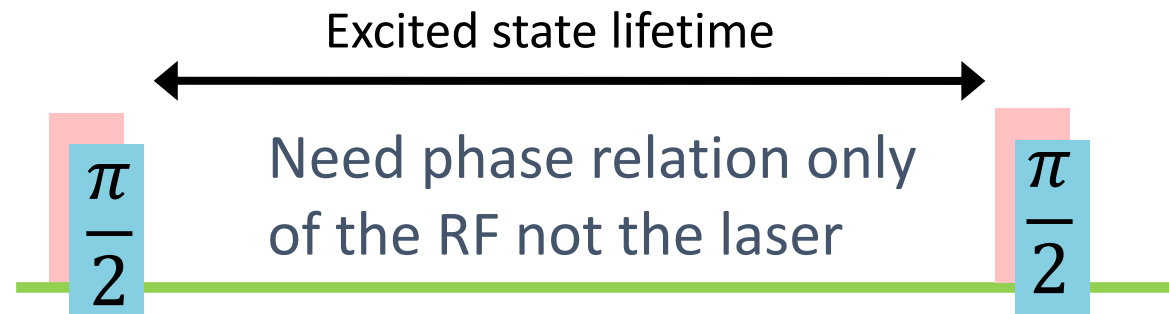


Measuring the Isotope shift in decoherence free subspace

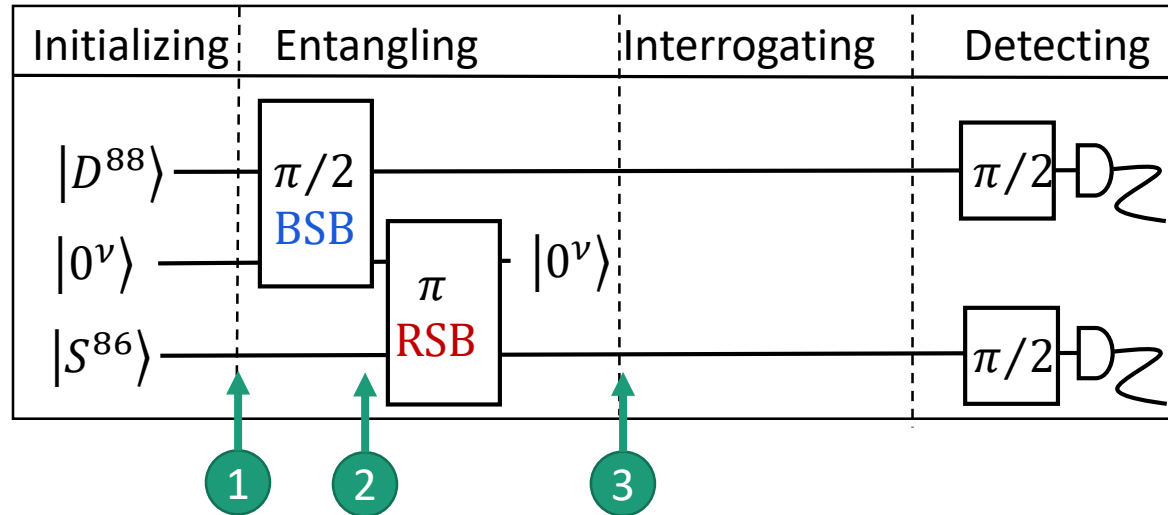
$$\psi(T) = |S\rangle|D'\rangle + e^{i2\pi\delta\nu T}|D\rangle|S'\rangle$$



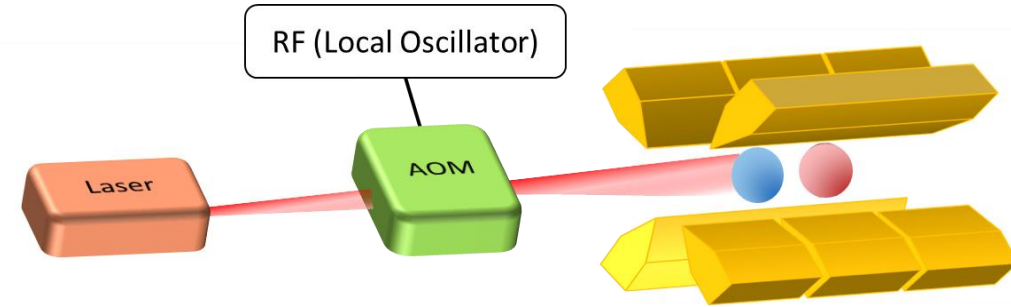
Ramsey sequence:



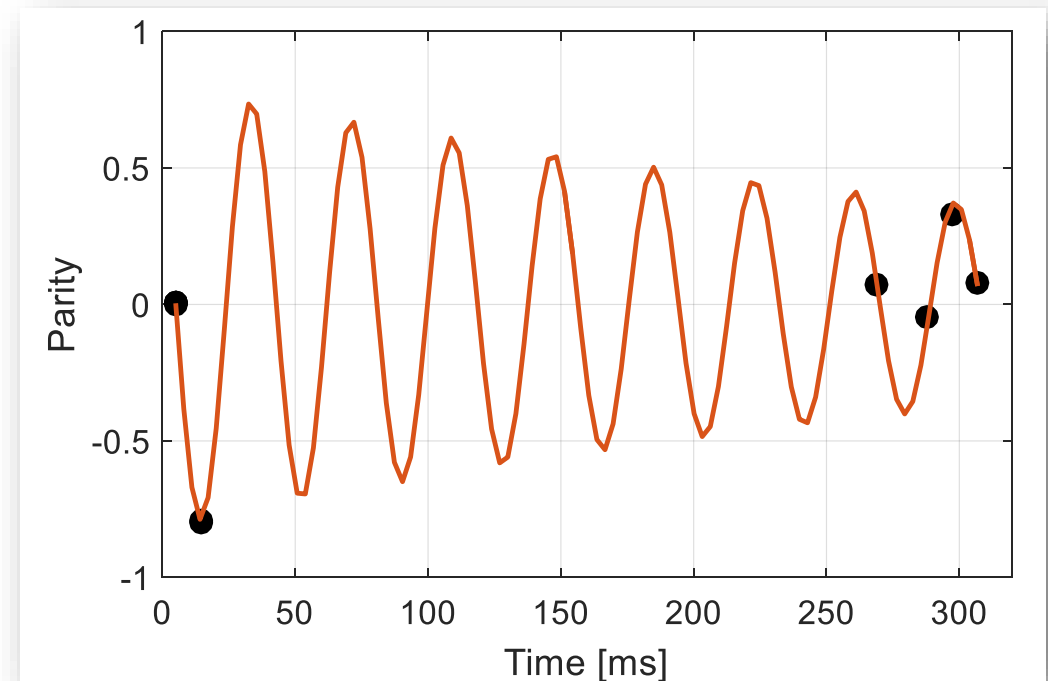
The Experiment Sequence - Entangled



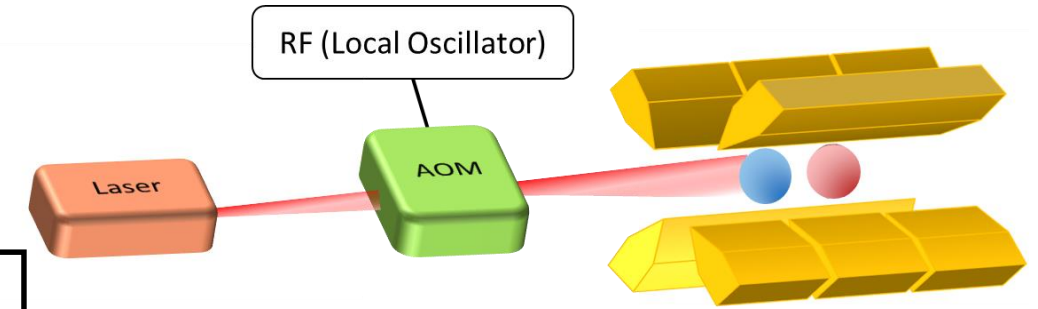
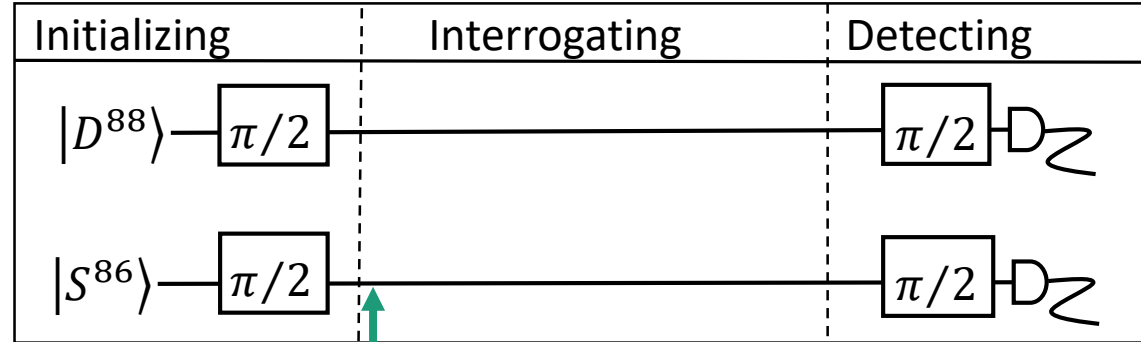
- 1 $\psi(T) = (|S^{88}\rangle|D^{86}\rangle|0\rangle)$
- 2 $\psi(T) = 1/\sqrt{2}(|S^{88}\rangle|0\rangle + |D^{88}\rangle|1\rangle)|S^{86}\rangle$
- 3 $\psi(T) = 1/\sqrt{2}(|S^{88}\rangle|D^{86}\rangle + e^{i2\pi\delta\nu T}|D^{88}\rangle|S^{86}\rangle)$



$$Parity = P_{SS} + P_{DD} - (P_{DS} + P_{SD})$$



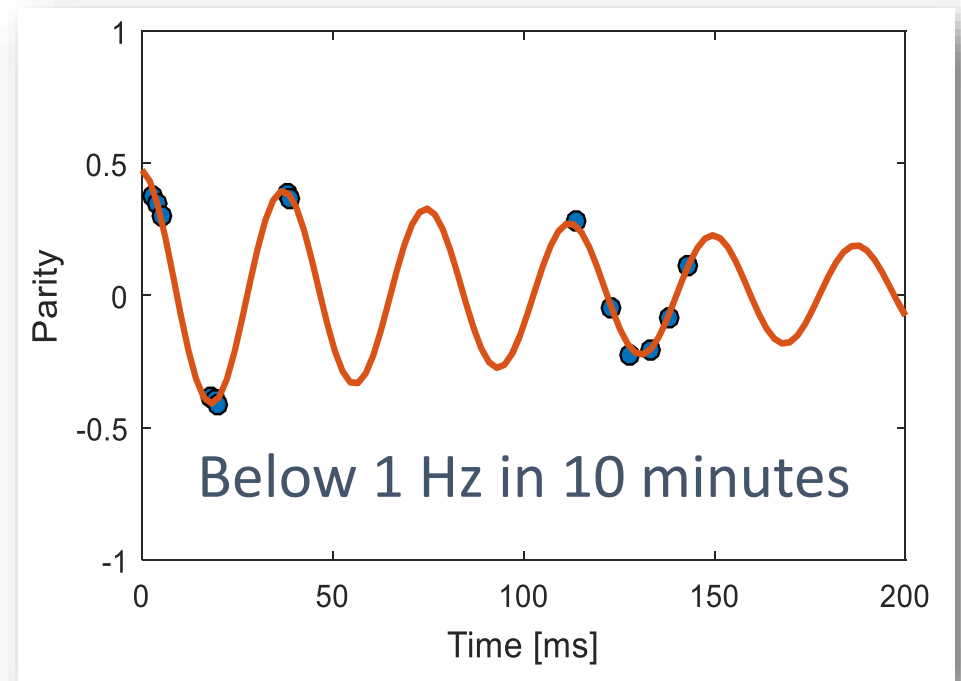
The Experimental Sequence - Separable



$$\text{Parity} = P_{SS} + P_{DD} - (P_{DS} + P_{SD})$$

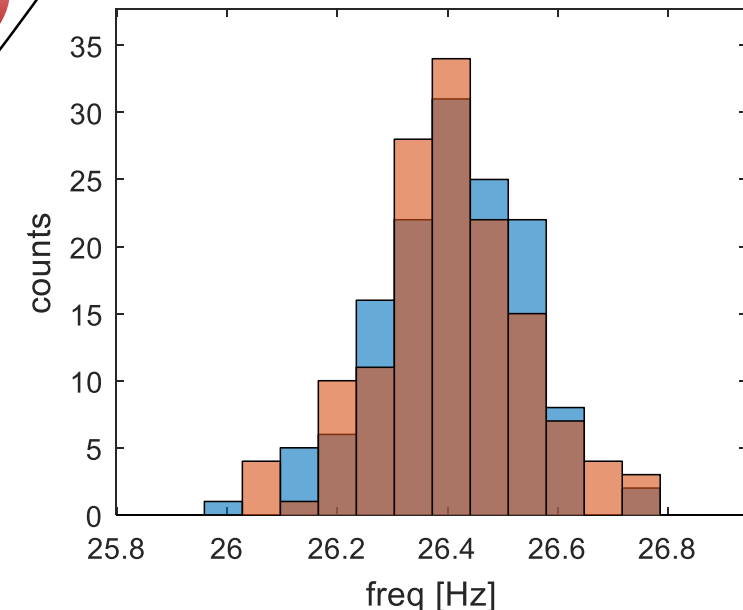
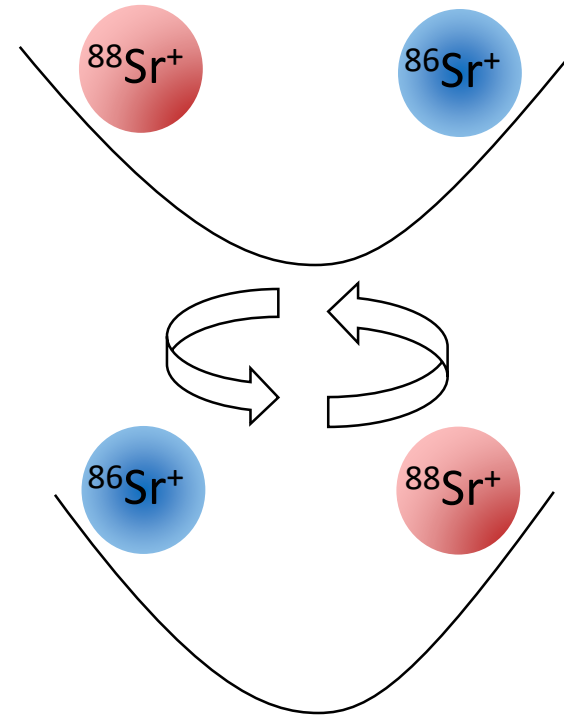
$$\psi(T) = 1/2 (|S^{88}\rangle|D^{86}\rangle + e^{i2\pi\delta\nu T}|D^{88}\rangle|S^{86}\rangle + e^{-i\phi(T)}|S^{88}\rangle|S^{86}\rangle + e^{i\phi(T)}|D^{88}\rangle|D^{86}\rangle)$$

Average to zero in the Parity signal

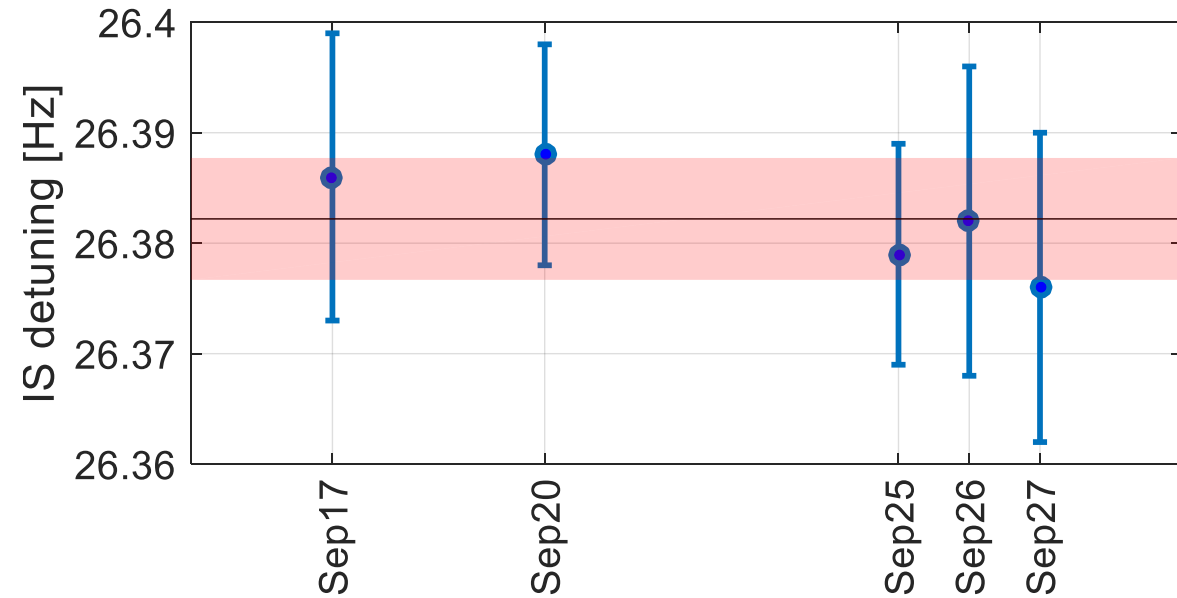
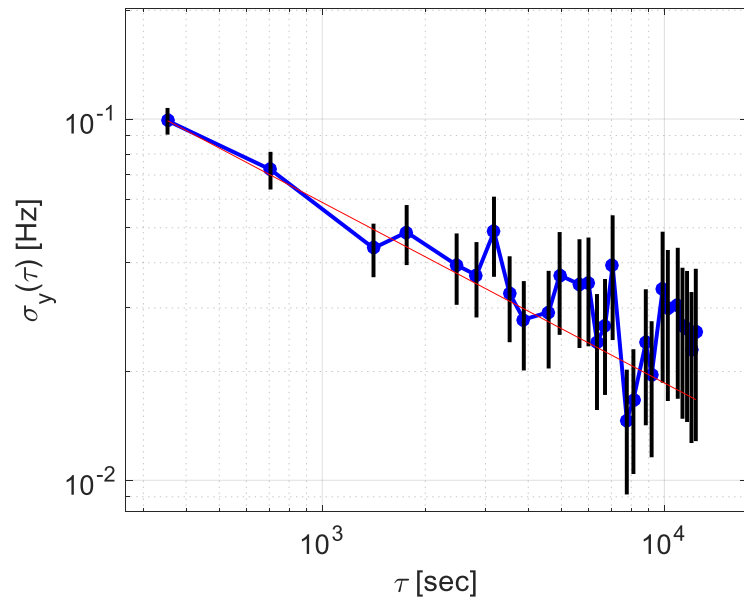
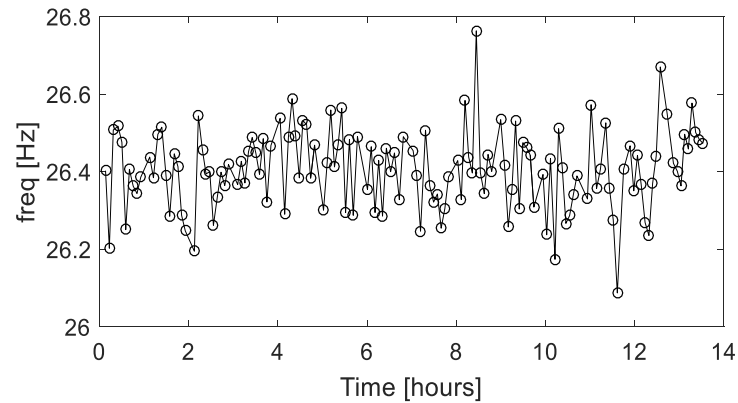


Systematic Uncertainties

- Common mode rejection:
 - Common Magnetic field noise
 - Quadruple shift (for only two ions)
 - Blackbody radiation
- Spatial Inhomogeneity :
 - Magnetic field gradient
 - Micromotion :
 - Second order Doppler
 - ac stark shift
- Light shift (laser light leakage)
- g-factor ?



Preliminary Result

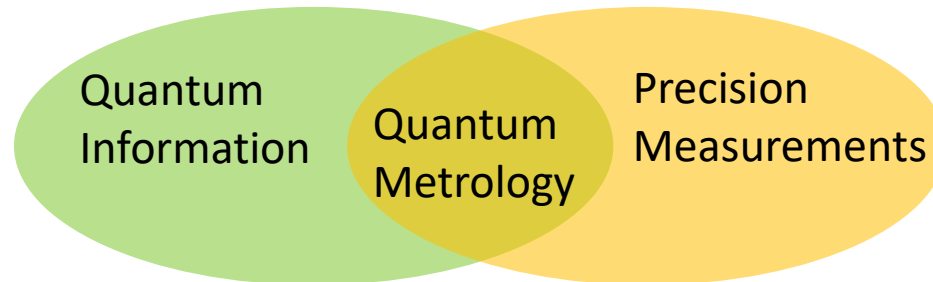


$$\Delta\nu_{S_{1/2}, D_{5/2}}^{88,86} = 570,264,116.382 (\pm 0.005 \text{ stat.}) [\text{Hz}]$$

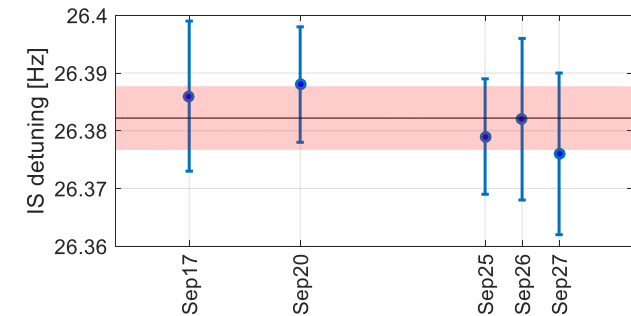
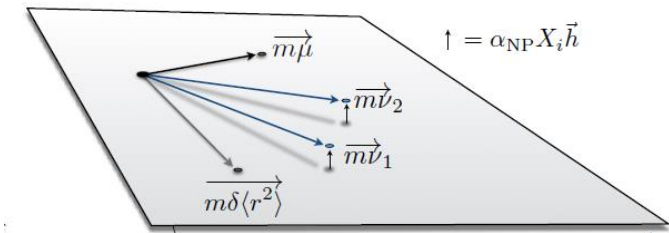
systematic ?

Summary

- King plot linearity has the potential to bound new physics
- Isotope shift can be measured with very high precision (relatively easy in trapped ions)



- Future plans:
 - Measuring King plot for the two E2 transitions is Yb^+ (Ca^+)
 - For Strontium we need the fourth isotope ^{90}Sr
 - Measuring the isotope g-factor



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Precision mass measurements: 10^{-10}



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The most precise atomic mass measurements in Penning traps

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Table 10

Atomic masses of the most abundant isotopes of strontium and ytterbium measured at FSU [109].

Atom	FSU mass (u)	σ_m/m (ppt)
^{86}Sr	85.909 260 730 9(91)	105
^{87}Sr	86.908 877 497 0(91)	105
^{88}Sr	87.905 612 257 1(97)	110
^{170}Yb	169.934 767 241(18)	105
^{171}Yb	170.936 331 514(19)	110
^{172}Yb	171.936 386 655(18)	105
^{173}Yb	172.938 216 213(18)	105
^{174}Yb	173.938 867 539(18)	105
^{176}Yb	175.942 574 702(22)	125