

CORRELATIONS IN RELATIVISTIC HEAVY ION COLLISIONS

Pan-American Few-Body Physics Boot Camp: Fostering Collaboration

22.10.2025

CONTENTS

- 1** Hanbury-Brown-Twiss interferometry
- 2** Heavy ion collisions
- 3** Definition of the correlation function
- 4** Final state interactions

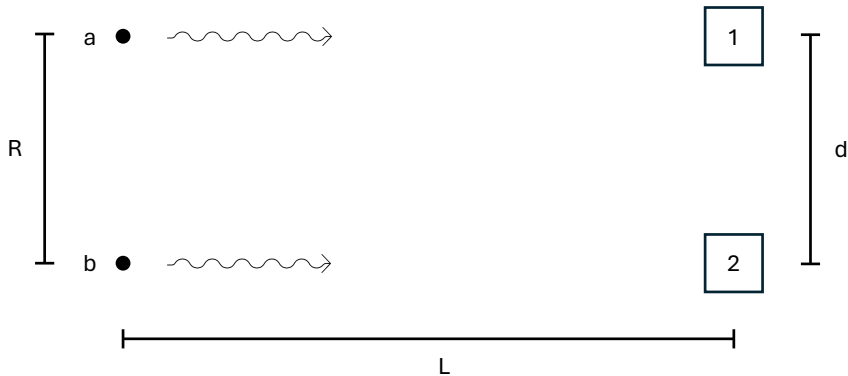
- 5** Nucleon-nucleon correlations
- 6** Toy problems: Coulomb and Square Well
- 7** Summary and Outlook



HANBURY-BROWN-TWISS INTERFEROMETRY

Section 1

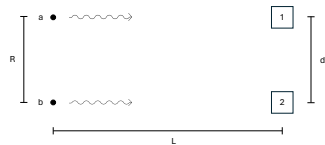
HANBURY-BROWN-TWISS INTERFEROMETRY



HANBURY-BROWN-TWISS INTERFEROMETRY

- Sources a and b emit radiation with amplitudes

$$A_a(\vec{r}) = A \frac{\exp[i(k|\vec{r}-\vec{r}_a|+\phi_a)]}{|\vec{r}-\vec{r}_a|} \text{ and } A_b(\vec{r}) = B \frac{\exp[i(k|\vec{r}-\vec{r}_b|+\phi_b)]}{|\vec{r}-\vec{r}_b|}$$

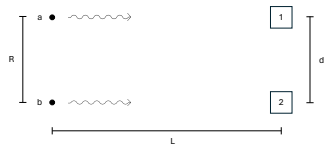


HANBURY-BROWN-TWISS INTERFEROMETRY

- Sources a and b emit radiation with amplitudes

$$A_a(\vec{r}) = A \frac{\exp[i(k|\vec{r}-\vec{r}_a|+\phi_a)]}{|\vec{r}-\vec{r}_a|} \text{ and } A_b(\vec{r}) = B \frac{\exp[i(k|\vec{r}-\vec{r}_b|+\phi_b)]}{|\vec{r}-\vec{r}_b|}$$
- Averaging over the phases ϕ_i we obtain the mean intensities

$$\langle I_1 \rangle = \langle I_2 \rangle \approx \frac{1}{L^2} (|A|^2 + |B|^2)$$



HANBURY-BROWN-TWISS INTERFEROMETRY

- Sources a and b emit radiation with amplitudes

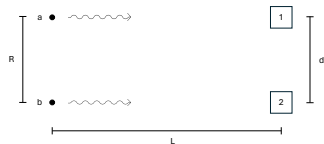
$$A_a(\vec{r}) = A \frac{\exp[i(k|\vec{r}-\vec{r}_a|+\phi_a)]}{|\vec{r}-\vec{r}_a|} \quad \text{and} \quad A_b(\vec{r}) = B \frac{\exp[i(k|\vec{r}-\vec{r}_b|+\phi_b)]}{|\vec{r}-\vec{r}_b|}$$

- Averaging over the phases ϕ_i we obtain the mean intensities

$$\langle I_1 \rangle = \langle I_2 \rangle \approx \frac{1}{L^2} (|A|^2 + |B|^2)$$

- Similar:

$$\langle I_1 I_2 \rangle = \langle I_1 \rangle \langle I_2 \rangle + \frac{2}{L^4} |A|^2 |B|^2 \cos(k(r_{1a} - r_{1b} - r_{2a} + r_{2b}))$$



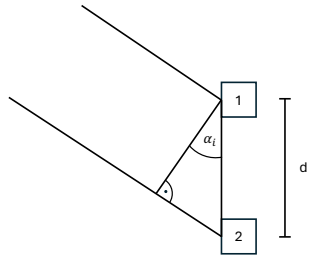
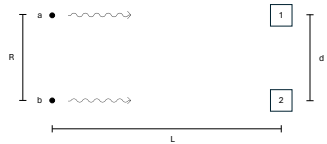
HANBURY-BROWN-TWISS INTERFEROMETRY

- Sources a and b emit radiation with amplitudes

$$A_a(\vec{r}) = A \frac{\exp[i(k|\vec{r}-\vec{r}_a|+\phi_a)]}{|\vec{r}-\vec{r}_a|} \text{ and } A_b(\vec{r}) = B \frac{\exp[i(k|\vec{r}-\vec{r}_b|+\phi_b)]}{|\vec{r}-\vec{r}_b|}$$
- Averaging over the phases ϕ_i we obtain the mean intensities

$$\langle I_1 \rangle = \langle I_2 \rangle \approx \frac{1}{L^2} (|A|^2 + |B|^2)$$
- Similar:

$$\langle I_1 I_2 \rangle = \langle I_1 \rangle \langle I_2 \rangle + \frac{2}{L^4} |A|^2 |B|^2 \cos(k(r_{1a} - r_{1b} - r_{2a} + r_{2b}))$$



HANBURY-BROWN-TWISS INTERFEROMETRY

- Sources a and b emit radiation with amplitudes

$$A_a(\vec{r}) = A \frac{\exp[i(k|\vec{r}-\vec{r}_a|+\phi_a)]}{|\vec{r}-\vec{r}_a|} \quad \text{and} \quad A_b(\vec{r}) = B \frac{\exp[i(k|\vec{r}-\vec{r}_b|+\phi_b)]}{|\vec{r}-\vec{r}_b|}$$

- Averaging over the phases ϕ_i we obtain the mean intensities

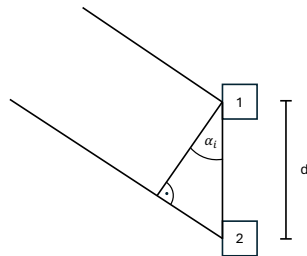
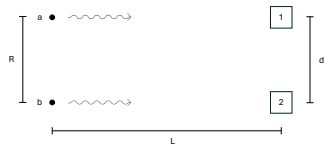
$$\langle I_1 \rangle = \langle I_2 \rangle \approx \frac{1}{L^2} (|A|^2 + |B|^2)$$

- Similar:

$$\langle I_1 I_2 \rangle = \langle I_1 \rangle \langle I_2 \rangle + \frac{2}{L^4} |A|^2 |B|^2 \cos(k(r_{1a} - r_{1b} - r_{2a} + r_{2b}))$$

- Approximating yields:

$$r_{1a} - r_{1b} - r_{2a} + r_{2b} = d (\sin(\alpha_a) - \sin(\alpha_b))$$



HANBURY-BROWN-TWISS INTERFEROMETRY

- Sources a and b emit radiation with amplitudes

$$A_a(\vec{r}) = A \frac{\exp[i(k|\vec{r}-\vec{r}_a|+\phi_a)]}{|\vec{r}-\vec{r}_a|} \quad \text{and} \quad A_b(\vec{r}) = B \frac{\exp[i(k|\vec{r}-\vec{r}_b|+\phi_b)]}{|\vec{r}-\vec{r}_b|}$$

- Averaging over the phases ϕ_i we obtain the mean intensities

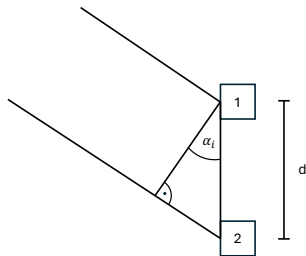
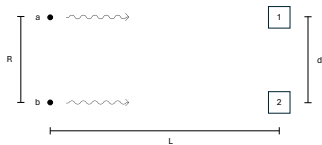
$$\langle I_1 \rangle = \langle I_2 \rangle \approx \frac{1}{L^2} (|A|^2 + |B|^2)$$

- Similar:

$$\langle I_1 I_2 \rangle = \langle I_1 \rangle \langle I_2 \rangle + \frac{2}{L^4} |A|^2 |B|^2 \cos(k(r_{1a} - r_{1b} - r_{2a} + r_{2b}))$$

- Approximating yields:

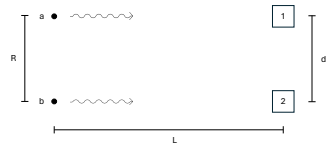
$$\begin{aligned} r_{1a} - r_{1b} - r_{2a} + r_{2b} &= d (\sin(\alpha_a) - \sin(\alpha_b)) \\ &\approx d (\alpha_a - \alpha_b) = d R/L \end{aligned}$$



HANBURY-BROWN-TWISS INTERFEROMETRY

- Sources a and b emit radiation with amplitudes

$$A_a(\vec{r}) = A \frac{\exp[i(k|\vec{r}-\vec{r}_a|+\phi_a)]}{|\vec{r}-\vec{r}_a|} \text{ and } A_b(\vec{r}) = B \frac{\exp[i(k|\vec{r}-\vec{r}_b|+\phi_b)]}{|\vec{r}-\vec{r}_b|}$$



[model: Baym, arXiv preprint nucl-th/9804026, 1998]

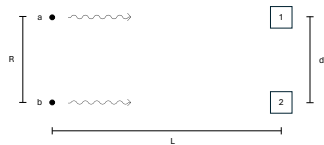
HANBURY-BROWN-TWISS INTERFEROMETRY

- Sources a and b emit radiation with amplitudes

$$A_a(\vec{r}) = A \frac{\exp[i(k|\vec{r}-\vec{r}_a|+\phi_a)]}{|\vec{r}-\vec{r}_a|} \text{ and } A_b(\vec{r}) = B \frac{\exp[i(k|\vec{r}-\vec{r}_b|+\phi_b)]}{|\vec{r}-\vec{r}_b|}$$

- We can define the Correlation function

$$C(d) = \langle I_1 I_2 \rangle / (\langle I_1 \rangle \langle I_2 \rangle)$$



[model: Baym, arXiv preprint nucl-th/9804026, 1998]

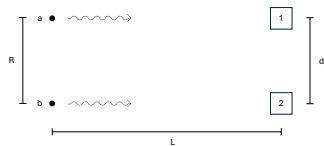
HANBURY-BROWN-TWISS INTERFEROMETRY

- Sources a and b emit radiation with amplitudes

$$A_a(\vec{r}) = A \frac{\exp[i(k|\vec{r}-\vec{r}_a|+\phi_a)]}{|\vec{r}-\vec{r}_a|} \text{ and } A_b(\vec{r}) = B \frac{\exp[i(k|\vec{r}-\vec{r}_b|+\phi_b)]}{|\vec{r}-\vec{r}_b|}$$

- We can define the Correlation function

$$\begin{aligned} C(d) &= \langle I_1 I_2 \rangle / (\langle I_1 \rangle \langle I_2 \rangle) \\ &= 1 + 2 \frac{|A|^2 |B|^2 \cos\left(\frac{kdR}{L}\right)}{(|A|^2 + |B|^2)^2} \end{aligned}$$



[model: Baym, arXiv preprint nucl-th/9804026, 1998]

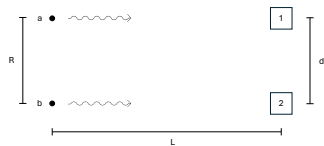
HANBURY-BROWN-TWISS INTERFEROMETRY

- Sources a and b emit radiation with amplitudes

$$A_a(\vec{r}) = A \frac{\exp[i(k|\vec{r}-\vec{r}_a|+\phi_a)]}{|\vec{r}-\vec{r}_a|} \quad \text{and} \quad A_b(\vec{r}) = B \frac{\exp[i(k|\vec{r}-\vec{r}_b|+\phi_b)]}{|\vec{r}-\vec{r}_b|}$$

- We can define the Correlation function

$$\begin{aligned} C(d) &= \langle I_1 I_2 \rangle / (\langle I_1 \rangle \langle I_2 \rangle) \\ &= 1 + 2 \frac{|A|^2 |B|^2 \cos\left(\frac{kdR}{L}\right)}{(|A|^2 + |B|^2)^2} \end{aligned}$$



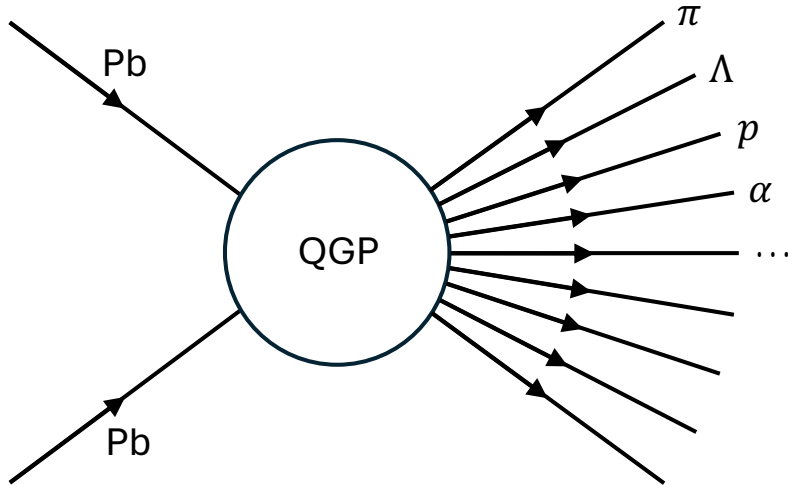
[model: Baym, arXiv preprint nucl-th/9804026, 1998]

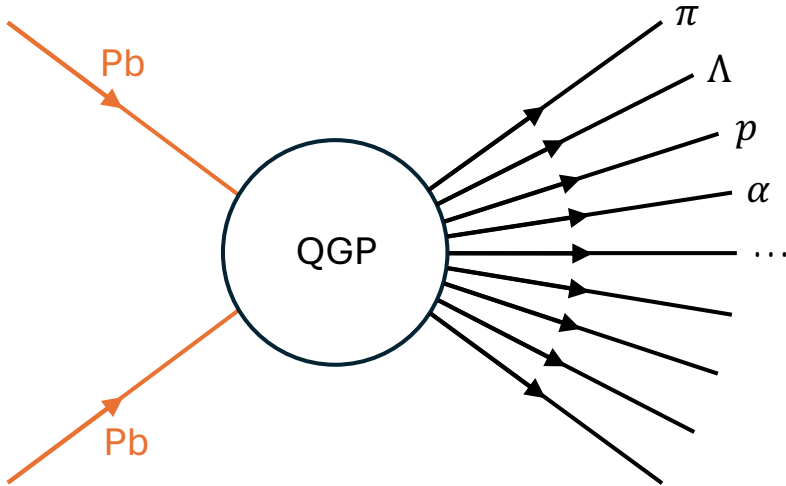
$\Rightarrow$ if k and L are known we can extract R by differing d and measuring the correlated signal

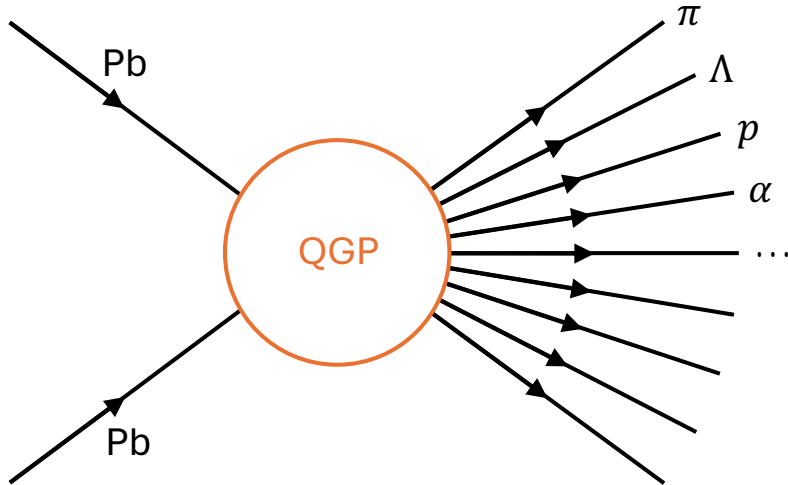


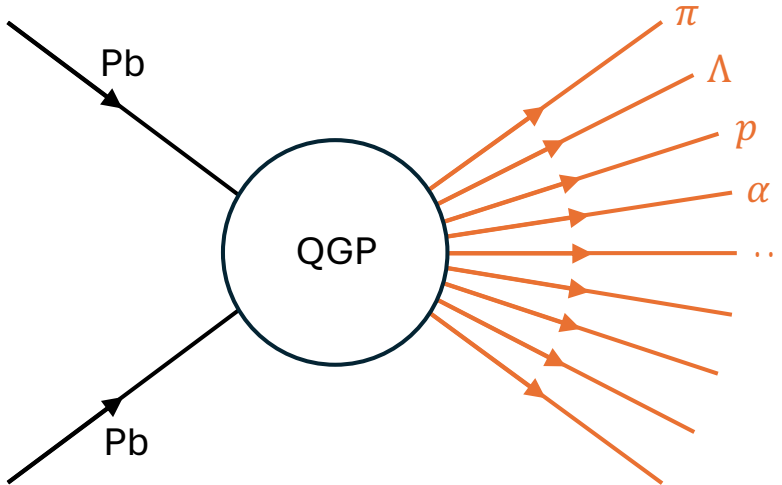
HEAVY ION COLLISIONS

Section 2









DEFINITION OF THE CORRELATION FUNCTION

Section 3

DEFINITION OF THE CORRELATION FUNCTION

- Start from the i -particle distribution

$$N_i(\vec{p}_1, \vec{p}_2, \dots, \vec{p}_i) = \left\langle a_{\vec{p}_1}^\dagger a_{\vec{p}_2}^\dagger \dots a_{\vec{p}_i}^\dagger a_{\vec{p}_i} \dots a_{\vec{p}_2} a_{\vec{p}_1} \right\rangle$$

DEFINITION OF THE CORRELATION FUNCTION

- Start from the i -particle distribution

$$N_i(\vec{p}_1, \vec{p}_2, \dots, \vec{p}_i) = \left\langle a_{\vec{p}_1}^\dagger a_{\vec{p}_2}^\dagger \dots a_{\vec{p}_i}^\dagger a_{\vec{p}_i} \dots a_{\vec{p}_2} a_{\vec{p}_1} \right\rangle$$

- Define the correlation function [Wiedemann Heinz, Physics Reports 319, 145–230 (1999)]

$$C(\vec{p}_1, \vec{p}_2) = \frac{N_2(\vec{p}_1, \vec{p}_2)}{N_1(\vec{p}_1)N_1(\vec{p}_2)}$$

DEFINITION OF THE CORRELATION FUNCTION

- Start from the i -particle distribution

$$N_i(\vec{p}_1, \vec{p}_2, \dots, \vec{p}_i) = \left\langle a_{\vec{p}_1}^\dagger a_{\vec{p}_2}^\dagger \dots a_{\vec{p}_i}^\dagger a_{\vec{p}_i} \dots a_{\vec{p}_2} a_{\vec{p}_1} \right\rangle$$

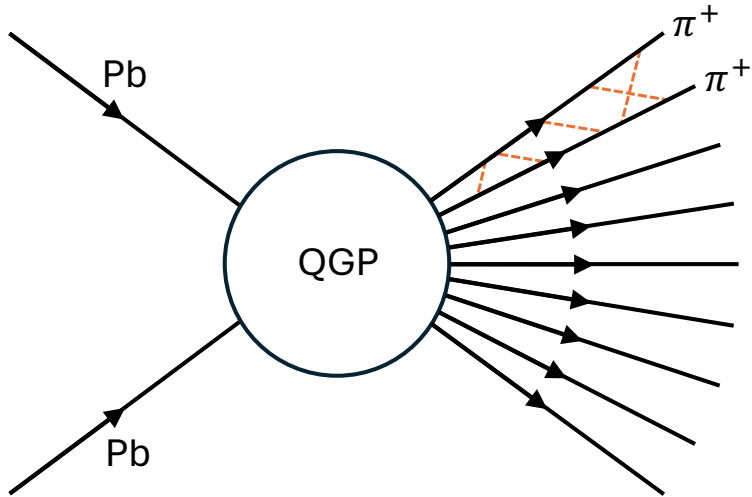
- Define the correlation function [Wiedemann Heinz, Physics Reports 319, 145–230 (1999)]

$$C(\vec{p}_1, \vec{p}_2) = \frac{N_2(\vec{p}_1, \vec{p}_2)}{N_1(\vec{p}_1)N_1(\vec{p}_2)}$$

- If particles tend to turn up in pairs we observe correlations above unity for small relative momenta

FINAL STATE INTERACTIONS

Section 4



FINAL STATE INTERACTIONS

- Correlation function in presence of final state interactions

$$C(\vec{q}) = \int d^3r S(\vec{r}) |\psi_{\vec{q}}(\vec{r})|^2$$

[S. E. Koonin, Physics Letters B 70, 43–47 (1977); S. Pratt, Physical Review D 33, 72–79 (1986)]

FINAL STATE INTERACTIONS

- Correlation function in presence of final state interactions

$$C(\vec{q}) = \int d^3r S(\vec{r}) |\psi_{\vec{q}}(\vec{r})|^2$$

[S. E. Koonin, Physics Letters B 70, 43–47 (1977); S. Pratt, Physical Review D 33, 72–79 (1986)]

Assumptions

[M. A. Lisa, Annual Review of Nuclear and Particle Science 55, 357–402 (2005); E. Epelbaum, arXiv preprint arXiv:2504.08631 (2025)]

FINAL STATE INTERACTIONS

- Correlation function in presence of final state interactions

$$C(\vec{q}) = \int d^3r S(\vec{r}) |\psi_{\vec{q}}(\vec{r})|^2$$

[S. E. Koonin, Physics Letters B 70, 43–47 (1977); S. Pratt, Physical Review D 33, 72–79 (1986)]

Assumptions

- Nonrelativistic treatment

[M. A. Lisa, Annual Review of Nuclear and Particle Science 55, 357–402 (2005); E. Epelbaum, arXiv preprint arXiv:2504.08631 (2025)]

FINAL STATE INTERACTIONS

- Correlation function in presence of final state interactions

$$C(\vec{q}) = \int d^3r S(\vec{r}) |\psi_{\vec{q}}(\vec{r})|^2$$

[S. E. Koonin, Physics Letters B 70, 43–47 (1977); S. Pratt, Physical Review D 33, 72–79 (1986)]

Assumptions

- Nonrelativistic treatment
- Chaotic source assumption

[M. A. Lisa, Annual Review of Nuclear and Particle Science 55, 357–402 (2005); E. Epelbaum, arXiv preprint arXiv:2504.08631 (2025)]

FINAL STATE INTERACTIONS

- Correlation function in presence of final state interactions

$$C(\vec{q}) = \int d^3r S(\vec{r}) |\psi_{\vec{q}}(\vec{r})|^2$$

[S. E. Koonin, Physics Letters B 70, 43–47 (1977); S. Pratt, Physical Review D 33, 72–79 (1986)]

Assumptions

- Nonrelativistic treatment
- Chaotic source assumption
- Emission function S is independent of the momenta of the emitted particles

[M. A. Lisa, Annual Review of Nuclear and Particle Science 55, 357–402 (2005); E. Epelbaum, arXiv preprint arXiv:2504.08631 (2025)]

FINAL STATE INTERACTIONS

- Correlation function in presence of final state interactions

$$C(\vec{q}) = \int d^3r S(\vec{r}) |\psi_{\vec{q}}(\vec{r})|^2$$

[S. E. Koonin, Physics Letters B 70, 43–47 (1977); S. Pratt, Physical Review D 33, 72–79 (1986)]

Assumptions

- Nonrelativistic treatment
- Chaotic source assumption
- Emission function S is independent of the momenta of the emitted particles
- Equal time approximation

[M. A. Lisa, Annual Review of Nuclear and Particle Science 55, 357–402 (2005); E. Epelbaum, arXiv preprint arXiv:2504.08631 (2025)]

FINAL STATE INTERACTIONS

- Correlation function in presence of final state interactions

$$C(\vec{q}) = \int d^3r S(\vec{r}) |\psi_{\vec{q}}(\vec{r})|^2$$

[S. E. Koonin, Physics Letters B 70, 43–47 (1977); S. Pratt, Physical Review D 33, 72–79 (1986)]

Assumptions

- Nonrelativistic treatment
- Chaotic source assumption
- Emission function S is independent of the momenta of the emitted particles
- Equal time approximation
- No interaction between the residual system and correlated particles

[M. A. Lisa, Annual Review of Nuclear and Particle Science 55, 357–402 (2005); E. Epelbaum, arXiv preprint arXiv:2504.08631 (2025)]

FINAL STATE INTERACTIONS

- Correlation function in presence of final state interactions

$$C(\vec{q}) = \int d^3r S(\vec{r}) |\psi_{\vec{q}}(\vec{r})|^2$$

[S. E. Koonin, Physics Letters B 70, 43–47 (1977); S. Pratt, Physical Review D 33, 72–79 (1986)]

Assumptions

- Nonrelativistic treatment
- Chaotic source assumption
- Emission function S is independent of the momenta of the emitted particles
- Equal time approximation
- No interaction between the residual system and correlated particles
- model dependence of the source

[M. A. Lisa, Annual Review of Nuclear and Particle Science 55, 357–402 (2005); E. Epelbaum, arXiv preprint arXiv:2504.08631 (2025)]

NUCLEON-NUCLEON CORRELATIONS

Section 5

NUCLEON-NUCLEON CORRELATIONS

- The absolute value squared of the s-wave function for NN scattering is given by

$$|\psi_{NN}(\vec{q}, \vec{r})|^2 = \frac{1}{4} |\psi_{1S_0}(\vec{q}, \vec{r})|^2 + \frac{3}{4} |\psi_{3S_1}(\vec{q}, \vec{r})|^2$$

NUCLEON-NUCLEON CORRELATIONS

- The absolute value squared of the s-wave function for NN scattering is given by

$$|\psi_{NN}(\vec{q}, \vec{r})|^2 = \frac{1}{4} |\psi_{1S_0}(\vec{q}, \vec{r})|^2 + \frac{3}{4} |\psi_{3S_1}(\vec{q}, \vec{r})|^2$$

- The asymptotic wave functions can be written as

$$\psi_{S_{0/1}}(\vec{q}, \vec{r}) = e^{iqz} + f_q^{(0/1)} \frac{e^{iqr}}{r}$$

NUCLEON-NUCLEON CORRELATIONS

- The absolute value squared of the s-wave function for NN scattering is given by

$$|\psi_{NN}(\vec{q}, \vec{r})|^2 = \frac{1}{4} |\psi_{1S_0}(\vec{q}, \vec{r})|^2 + \frac{3}{4} |\psi_{3S_1}(\vec{q}, \vec{r})|^2$$

- The asymptotic wave functions can be written as

$$\psi_{S_{0/1}}(\vec{q}, \vec{r}) = e^{iqz} + f_q^{(0/1)} \frac{e^{iqr}}{r}$$

- Approximating the scattering amplitude in s-wave approximation using ERE yields

$$f_q(\theta) = f_q = \frac{1}{-\frac{1}{a_0} + \frac{1}{2}r_0q^2 - iq}$$

NUCLEON-NUCLEON CORRELATIONS

- Frequently used is a Gaussian source which can be parametrized by

$$S(r) = \frac{1}{\sqrt{2\pi}^3 \sigma^3} e^{-\frac{r^2}{2\sigma^2}}$$

NUCLEON-NUCLEON CORRELATIONS

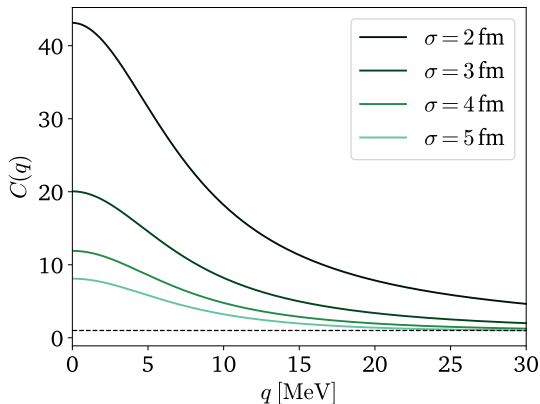
- Frequently used is a Gaussian source which can be parametrized by

$$S(r) = \frac{1}{\sqrt{2\pi}^3 \sigma^3} e^{-\frac{r^2}{2\sigma^2}}$$

- With this we can calculate

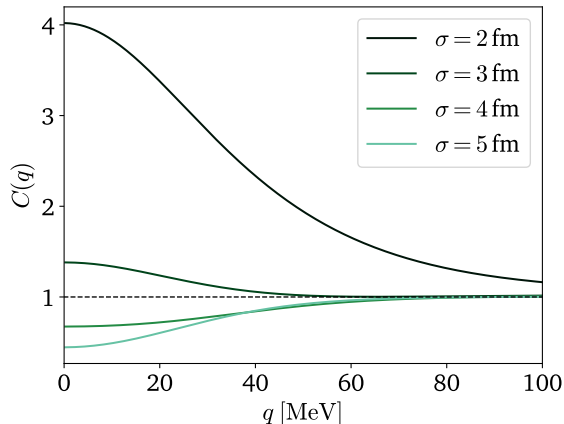
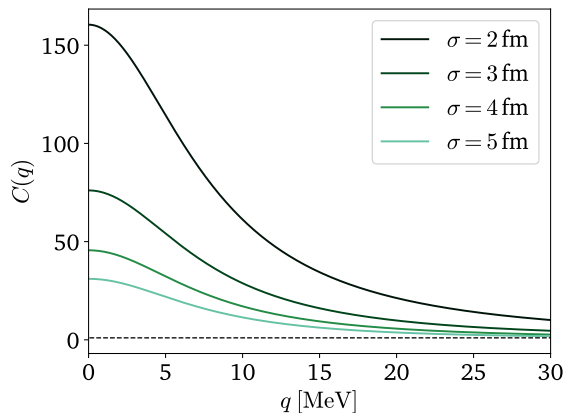
$$C_{NN}(q) = 1 + \frac{e^{-2\sigma^2 q^2} |f_q|^2}{\sigma^2 q} \left[q + \operatorname{Erfi}(\sqrt{2}\sigma q) \left(-\frac{1}{a_0} + \frac{1}{2} r_0 q^2 \right) \right]$$

Correlation of nucleon-nucleon scattering



Parameters: $a_{1S_0} = -23.71$ fm, $r_{1S_0} = 2.70$ fm, $a_{3S_1} = 5.42$ fm and $r_{3S_1} = 1.76$ fm

[V. A. Babenko and N. M. Petrov, Physics of Atomic Nuclei 70, 669–675 (2007)]

Separate correlations for 1S_0 and 3S_1 channels

Parameters: $a_{^1S_0} = -23.71$ fm, $r_{^1S_0} = 2.70$ fm, $a_{^3S_1} = 5.42$ fm and $r_{^3S_1} = 1.76$ fm

[V. A. Babenko and N. M. Petrov, Physics of Atomic Nuclei 70, 669–675 (2007)]

NUCLEON-NUCLEON CORRELATIONS

- The limit $q \rightarrow 0$ can be calculated to be

$$\lim_{q \rightarrow 0} C_{\text{NN}}(q) = 1 + \frac{2a_0^2}{\sigma^2} - \frac{2^{5/2}a_0}{\sqrt{\pi}\sigma}$$

NUCLEON-NUCLEON CORRELATIONS

- The limit $q \rightarrow 0$ can be calculated to be

$$\lim_{q \rightarrow 0} C_{\text{NN}}(q) = 1 + \frac{2a_0^2}{\sigma^2} - \frac{2^{5/2}a_0}{\sqrt{\pi}\sigma}$$

- Two modifications can be checked:

NUCLEON-NUCLEON CORRELATIONS

- The limit $q \rightarrow 0$ can be calculated to be

$$\lim_{q \rightarrow 0} C_{\text{NN}}(q) = 1 + \frac{2a_0^2}{\sigma^2} - \frac{2^{5/2}a_0}{\sqrt{\pi}\sigma}$$

- Two modifications can be checked:
 - Change of the correlation function for full instead of asymptotic solutions

NUCLEON-NUCLEON CORRELATIONS

- The limit $q \rightarrow 0$ can be calculated to be

$$\lim_{q \rightarrow 0} C_{\text{NN}}(q) = 1 + \frac{2a_0^2}{\sigma^2} - \frac{2^{5/2}a_0}{\sqrt{\pi}\sigma}$$

- Two modifications can be checked:
 - Change of the correlation function for full instead of asymptotic solutions
 - Incorporating higher partial waves

NUCLEON-NUCLEON CORRELATIONS

- The limit $q \rightarrow 0$ can be calculated to be

$$\lim_{q \rightarrow 0} C_{\text{NN}}(q) = 1 + \frac{2a_0^2}{\sigma^2} - \frac{2^{5/2}a_0}{\sqrt{\pi}\sigma}$$

- Two modifications can be checked:
 - Change of the correlation function for full instead of asymptotic solutions
 - Incorporating higher partial waves

⇒ Discuss systems with Coulomb and square-well potentials

TOY PROBLEMS: COULOMB AND SQUARE WELL

Section 6

COULOMB SCATTERING PROBLEM

- The Schrödinger equation with attractive Coulomb potential and $\kappa = \alpha m/k$ is given by

$$\left(\nabla^2 + \frac{2\kappa k}{r} + k^2 \right) \psi = 0$$

COULOMB SCATTERING PROBLEM

- The Schrödinger equation with attractive Coulomb potential and $\kappa = \alpha m/k$ is given by

$$\left(\nabla^2 + \frac{2\kappa k}{r} + k^2 \right) \psi = 0$$

- One can derive the scattering solution proportional to e^{ikz} for $r \rightarrow \infty$ to be

[Landau and Lifshitz, Quantum mechanics: non-relativistic theory]

$$\psi(\vec{r}) = e^{\pi\kappa/2} \Gamma(1 - i\kappa) e^{ikr \cos \theta} {}_1F_1(i\kappa, 1, ikr(1 - \cos \theta))$$

COULOMB SCATTERING PROBLEM

- The Schrödinger equation with attractive Coulomb potential and $\kappa = \alpha m/k$ is given by

$$\left(\nabla^2 + \frac{2\kappa k}{r} + k^2 \right) \psi = 0$$

- One can derive the scattering solution proportional to e^{ikz} for $r \rightarrow \infty$ to be

[Landau and Lifshitz, Quantum mechanics: non-relativistic theory]

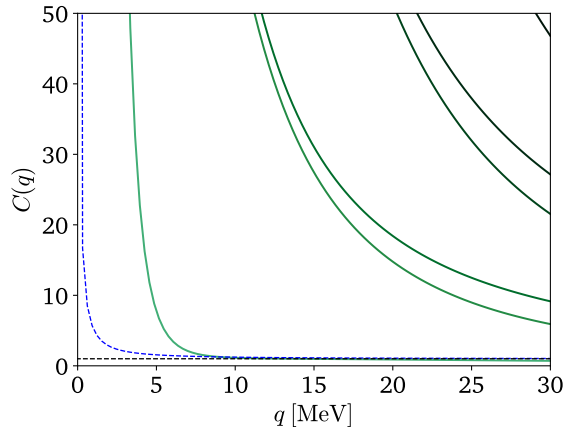
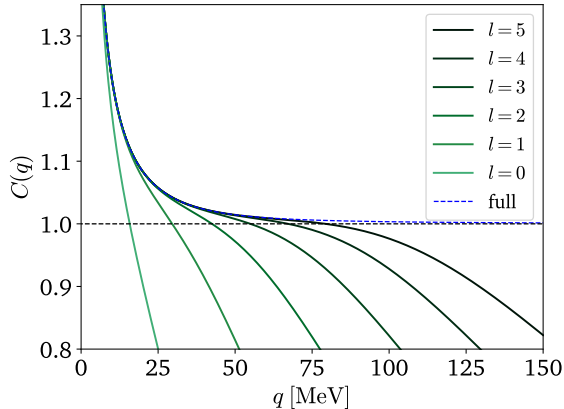
$$\psi(\vec{r}) = e^{\pi\kappa/2} \Gamma(1 - i\kappa) e^{ikr \cos \theta} {}_1F_1(i\kappa, 1, ikr(1 - \cos \theta))$$

- The decomposition into partial waves is given by [A. Messiah, Quantum mechanics]

$$\psi(\vec{r}) = \sum_{l=0}^{\infty} (2l+1) i^l e^{i\sigma_l} R^l(kr) P_l(\cos \theta)$$

with $\sigma_l = \arg(\Gamma(l+1 - i\kappa))$ and $R^l(kr) = \frac{2^l e^{\kappa\pi/2} |\Gamma(l+1 - i\kappa)|}{(2l+1)!} e^{ikr} (kr)^l {}_1F_1(l+1 - i\kappa, 2l+2, -2ikr)$

Coulomb correlation functions (pion-proton system), on the right for asymptotic partial waves



Parameters: $\sigma = \sqrt{2} \cdot 3 \text{ fm}$

SQUARE WELL

- The square well potential in three dimensions is given by

$$V(r) = \begin{cases} -V_0 & r < b \\ 0 & r > b \end{cases}$$

SQUARE WELL

- The square well potential in three dimensions is given by

$$V(r) = \begin{cases} -V_0 & r < b \\ 0 & r > b \end{cases}$$

- The asymptotic s-wave solution to this is given by

$$\psi_{\text{asymptotic}}(r) = A \frac{\sin(kr + \delta)}{kr}$$

SQUARE WELL

- The square well potential in three dimensions is given by

$$V(r) = \begin{cases} -V_0 & r < b \\ 0 & r > b \end{cases}$$

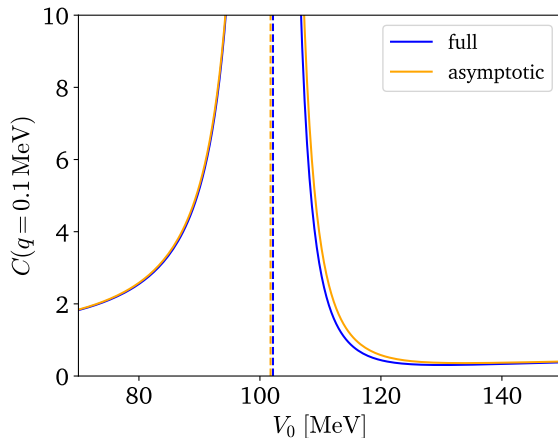
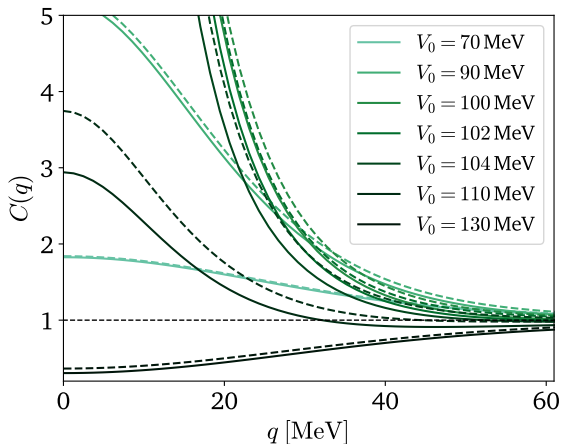
- The asymptotic s-wave solution to this is given by

$$\psi_{\text{asymptotic}}(r) = A \frac{\sin(kr + \delta)}{kr}$$

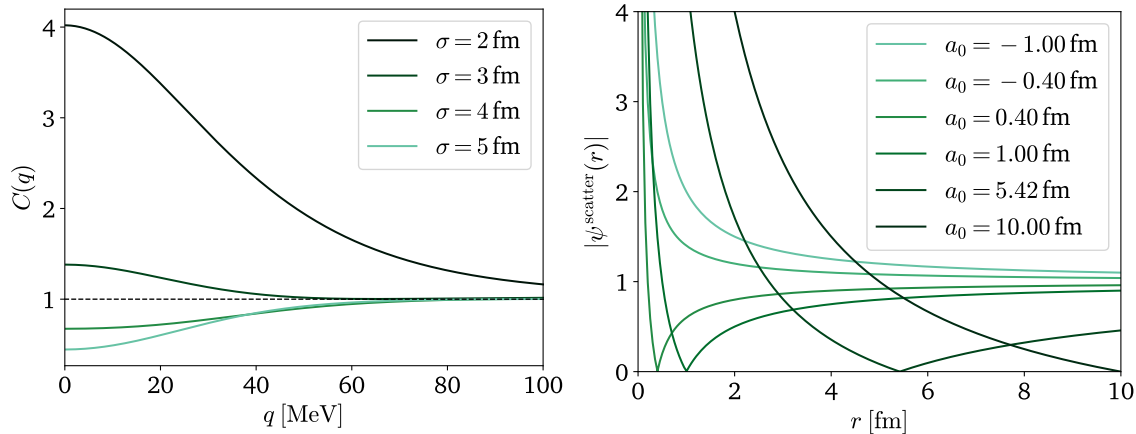
with the norm A , the modified wave number $k' = \sqrt{k^2 + 2V_0}$ (with $m = 1$) and

$$\cot \delta = (k \sin(kr) \sin(k'r) + k' \cos(kr) \cos(k'r)) / (k \cos kr \sin(k'r) - k' \sin(kr) \cos(k'r))$$

Square well correlation functions



Parameters: $\sigma = \sqrt{2} \cdot 3 \text{ fm}$, $b = 1 \text{ fm}$, $\mu = 938/2 \text{ MeV}$

Nucleonic 3S_1 correlation and wave functions

Parameters: (left plot:) $\sigma = \sqrt{2} \cdot 3$ fm, $a_{3S_1} = 5.42$ fm, (both plots:) $r_{3S_1} = 1.76$ fm, (right plot:) $q = 0.1$ MeV

[V. A. Babenko and N. M. Petrov, Physics of Atomic Nuclei 70, 669–675 (2007)]

Institute for Nuclear Physics | P. Quoss

SUMMARY AND OUTLOOK

- The interferometric technique of Hanbury-Brown Twiss interferometry can be adapted to the setting of heavy-ion collisions

SUMMARY AND OUTLOOK

- The interferometric technique of Hanbury-Brown Twiss interferometry can be adapted to the setting of heavy-ion collisions
- This results in the definition of the correlation function that can be directly measured in experiments

SUMMARY AND OUTLOOK

- The interferometric technique of Hanbury-Brown Twiss interferometry can be adapted to the setting of heavy-ion collisions
- This results in the definition of the correlation function that can be directly measured in experiments
- Bound states play an important role in determining the shape of the correlation function

SUMMARY AND OUTLOOK

- The interferometric technique of Hanbury-Brown Twiss interferometry can be adapted to the setting of heavy-ion collisions
- This results in the definition of the correlation function that can be directly measured in experiments
- Bound states play an important role in determining the shape of the correlation function
- Higher order partial waves become increasingly more important for higher relative momenta

SUMMARY AND OUTLOOK

- The interferometric technique of Hanbury-Brown Twiss interferometry can be adapted to the setting of heavy-ion collisions
- This results in the definition of the correlation function that can be directly measured in experiments
- Bound states play an important role in determining the shape of the correlation function
- Higher order partial waves become increasingly more important for higher relative momenta
- Looking forward we will combine the femtoscopic approach with an EFT calculation for the hypertriton

THANK YOU FOR YOUR ATTENTION!

CORRELATION FUNCTION OF IDENTICAL BOSONS

- Define the source function (also Wigner Function) to be

$$S(x, K) = \langle N \rangle_P \int d^4r d^4p d^4y \rho(r, p) e^{iy(K-p)} J_0^* \left(x + \frac{1}{2}y \right) J_0 \left(x - \frac{1}{2}y \right)$$

CORRELATION FUNCTION OF IDENTICAL BOSONS

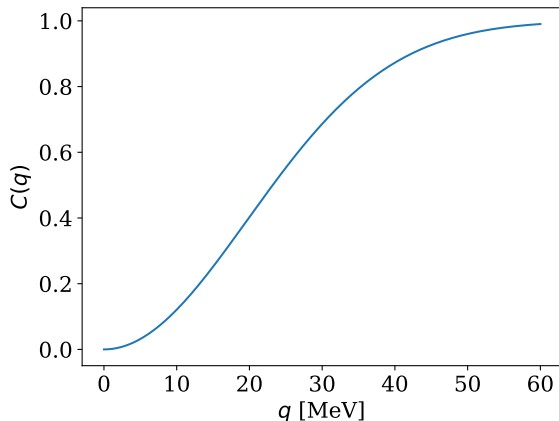
- Define the source function (also Wigner Function) to be

$$S(x, K) = \langle N \rangle_P \int d^4r d^4p d^4y \rho(r, p) e^{iy(K-p)} J_0^* \left(x + \frac{1}{2}y \right) J_0 \left(x - \frac{1}{2}y \right)$$

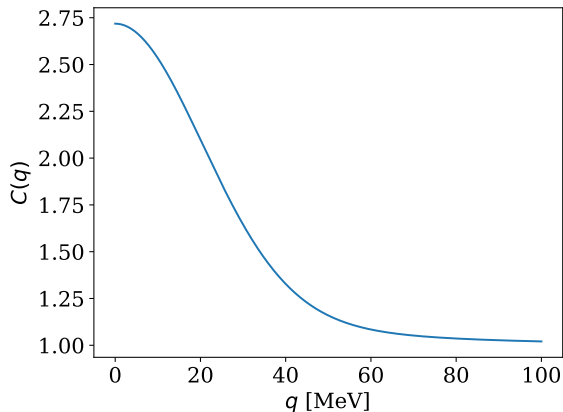
- With this we can calculate

$$c(\vec{q}, \vec{K}) = 1 + \frac{|\int d^4x S(x, K) e^{-iqx}|^2}{|\int d^4x S(x, K)|^2}$$

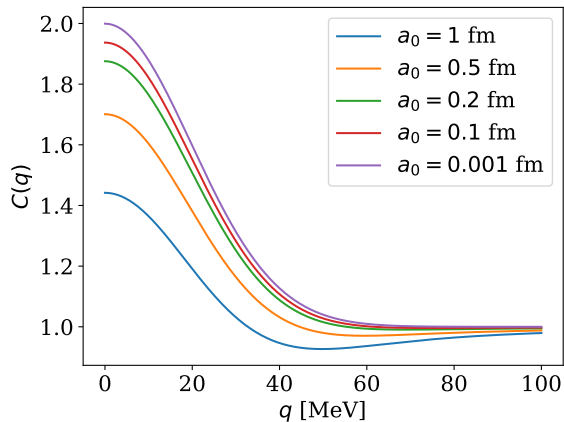
Correlation of Fermions

Parameters: $\sigma = 5$ fm

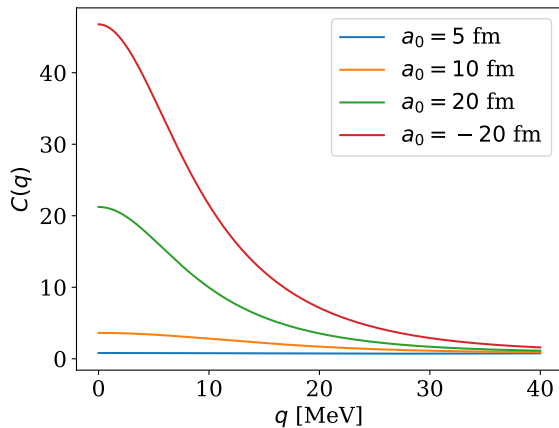
Correlation of Bosons

Parameters: $\sigma = 5$ fm, $a_0 = -1$ fm, $r_0 = 0$ fm

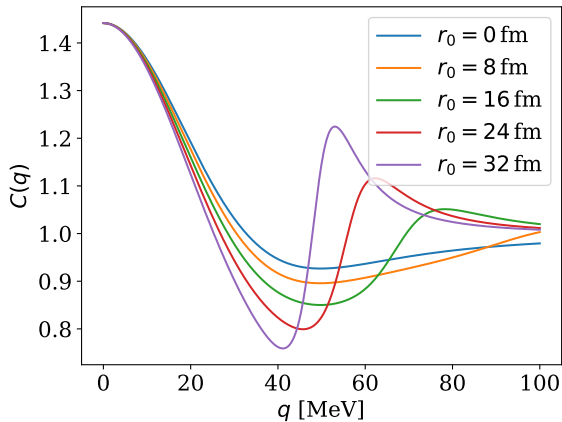
Correlation of Bosons

Parameters: $\sigma = 5$ fm, $r_0 = 0$ fm

Correlation of Bosons

Parameters: $\sigma = 5$ fm, $r_0 = 0$ fm

Correlation of Bosons

Parameters: $\sigma = 5$ fm and $a_0 = 1$ fm

RESULTS FOR GAUSSIAN MODEL

- Model the emission function by

$$S(x, K) = \delta(t) \frac{1}{(2\sigma\sqrt{\pi})^3} \exp\left[-\frac{\vec{x}^2}{4\sigma^2}\right]$$

RESULTS FOR GAUSSIAN MODEL

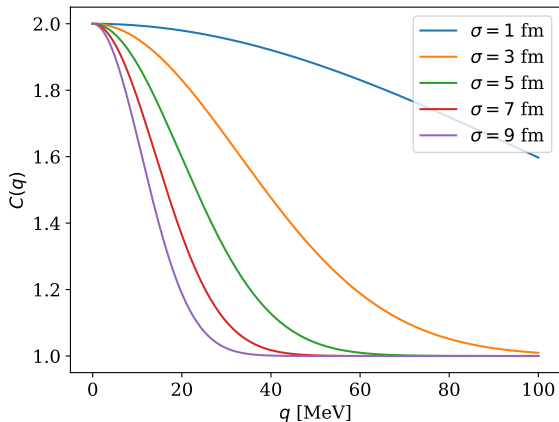
- Model the emission function by

$$S(x, K) = \delta(t) \frac{1}{(2\sigma\sqrt{\pi})^3} \exp\left[-\frac{\vec{x}^2}{4\sigma^2}\right]$$

- This leads to the correlation function

$$c(\vec{q}, \vec{K}) = 1 + \exp[-2\sigma^2 \vec{q}^2]$$

RESULTS FOR GAUSSIAN MODEL



RESULTS: IDENTICAL FERMIONS AND BOSONS

- For identical particles we need to consider an explicitly (anti-) symmetrized wave function

$$\psi_{\text{asymptotic}}(\vec{q}, \vec{r}) = \frac{1}{\sqrt{2}} \left(e^{iqz} \pm e^{-iqz} + [f_q(\theta) \pm f_q(\pi - \theta)] \frac{e^{iqr}}{r} \right)$$

RESULTS: IDENTICAL FERMIONS AND BOSONS

- For identical particles we need to consider an explicitly (anti-) symmetrized wave function

$$\psi_{\text{asymptotic}}(\vec{q}, \vec{r}) = \frac{1}{\sqrt{2}} \left(e^{iqz} \pm e^{-iqz} + [f_q(\theta) \pm f_q(\pi - \theta)] \frac{e^{iqr}}{r} \right)$$

- In S-wave approximation the scattering amplitude can be expressed as

$$f_q(\theta) = f_q = \frac{1}{-\frac{1}{a_0} + \frac{1}{2}r_0q^2 - iq}$$

RESULTS: IDENTICAL FERMIONS AND BOSONS

- For identical particles we need to consider an explicitly (anti-) symmetrized wave function

$$\psi_{\text{asymptotic}}(\vec{q}, \vec{r}) = \frac{1}{\sqrt{2}} \left(e^{iqz} \pm e^{-iqz} + [f_q(\theta) \pm f_q(\pi - \theta)] \frac{e^{iqr}}{r} \right)$$

- In S-wave approximation the scattering amplitude can be expressed as

$$f_q(\theta) = f_q = \frac{1}{-\frac{1}{a_0} + \frac{1}{2}r_0q^2 - iq}$$

- Model the source function by

$$S(r) = \frac{1}{\sqrt{2\pi}^3 \sigma^3} e^{-\frac{r^2}{2\sigma^2}}$$

RESULT: IDENTICAL FERMIONS AND BOSONS

- With this we obtain for fermions in S-wave approximation

$$C_{\text{fermions}}(q) = 1 - e^{-2\sigma^2 q^2}$$

RESULT: IDENTICAL FERMIONS AND BOSONS

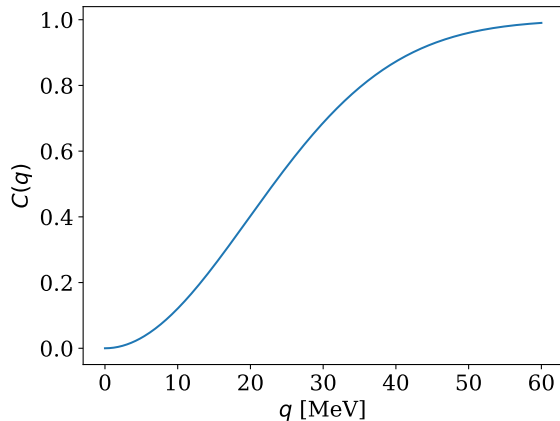
- With this we obtain for fermions in S-wave approximation

$$C_{\text{fermions}}(q) = 1 - e^{-2\sigma^2 q^2}$$

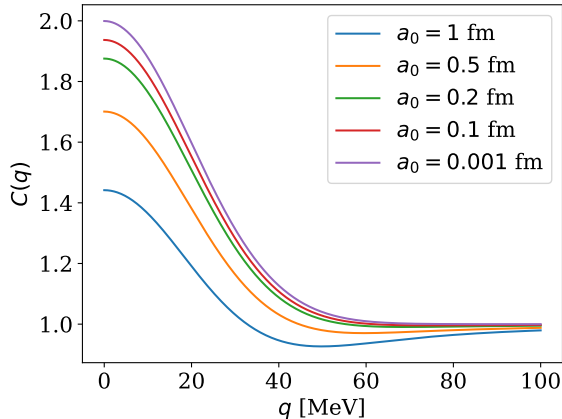
- While for Bosons

$$C_{\text{bosons}}(q) = 1 + e^{-2\sigma^2 q^2} + \frac{2e^{-2\sigma^2 q^2} |f_q|^2}{\sigma^2 q} \left[q + \text{Erfi}(\sqrt{2}\sigma q) \left(-\frac{1}{a_0} + \frac{1}{2}r_0 q^2 \right) \right]$$

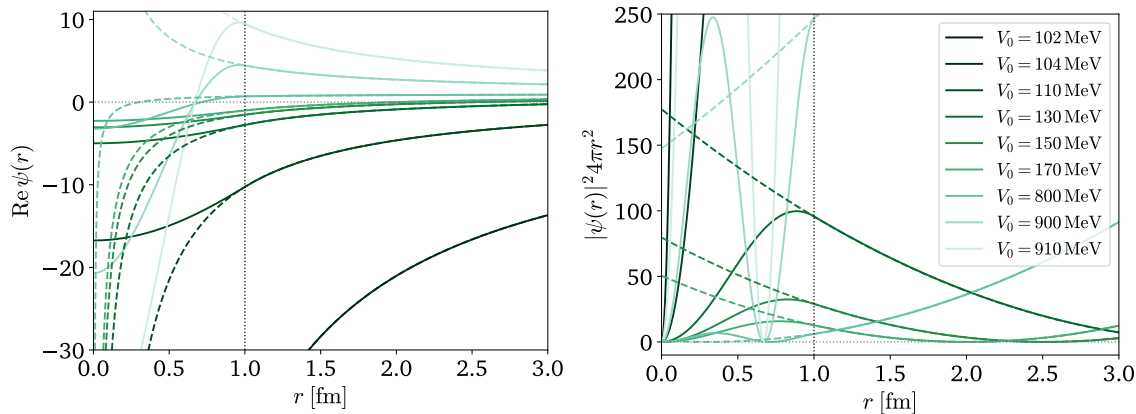
Correlation of Fermions

Parameters: $\sigma = 5$ fm

Correlation of Bosons

Parameters: $\sigma = 5$ fm, $r_0 = 0$ fm

Square well wave functions

Parameters: $b = 1$ fm, $q = 0.1$ MeV