





Time-resolved sensing of electromagnetic fields with single electron interferometry

Gerbold Ménard

Superconducting Devices for Quantum Optics and Quantum Simulations 7th October 2025 - Trento











Research team



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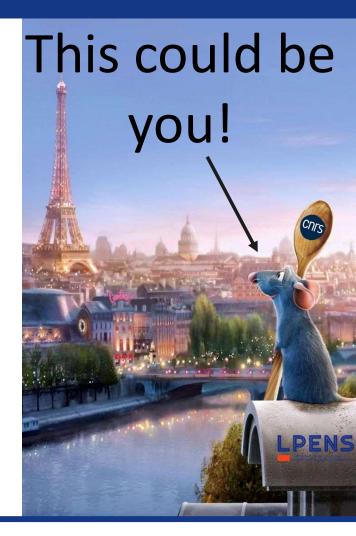






Advertisement

Open post-doc position in Paris at LPENS!!





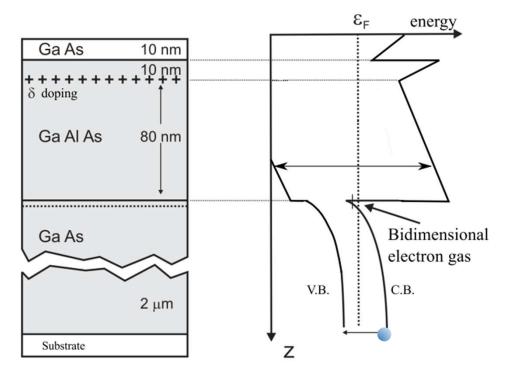
Outline:

- The quantum Hall effect
- Electron quantum optics
- Fabry-Perot interferometry

The quantum Hall effect: GaAs heterostructures

Two-dimensional electron gas in a magnetic field: quantum Hall effect

GaAs based 2D electron gas (2DEG)



Creation of a potential trap for electrons through band engineering

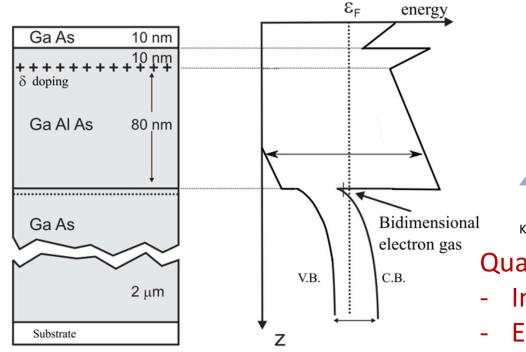
J. Gabelli, PhD thesis (2006)

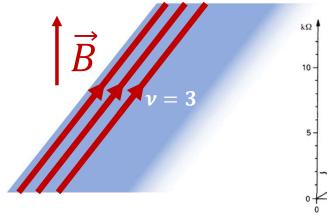


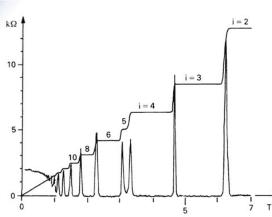
The quantum Hall effect: Edge states

Two-dimensional electron gas in a magnetic field: quantum Hall effect

GaAs based 2D electron gas (2DEG)







K. Von Klitzing, Rev. Mod. Phys. 58, 519 (1986)

Quantum picture:

- Insulating bulk, chiral edge states
- Edge channels with quantized conductance

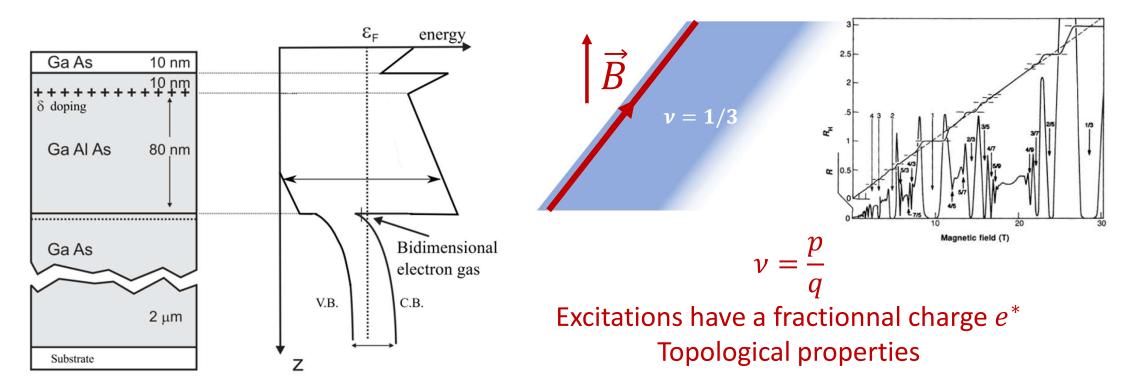
$$\rho = \nu e^2/h$$
, $\nu = h n_{2D}/eB \in \mathbb{N}$



J. Gabelli, PhD thesis (2006)

The quantum Hall effect

High purity 2DEG: fractionnal quantum Hall effect



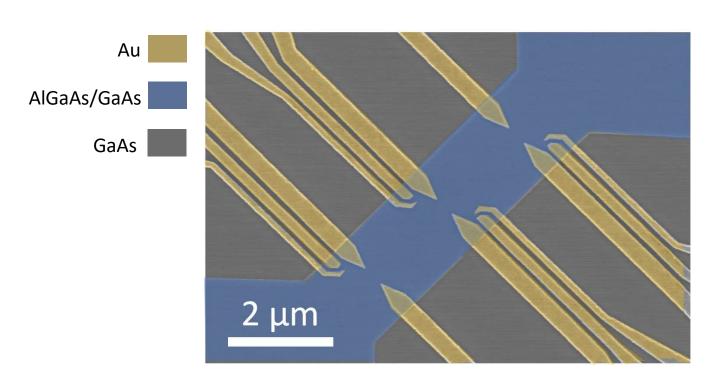
JP Eisenstein and HL Stormer, "The fractional quantum Hall effect," Science 248, 1510-1516 (1990).



Sample fabrication

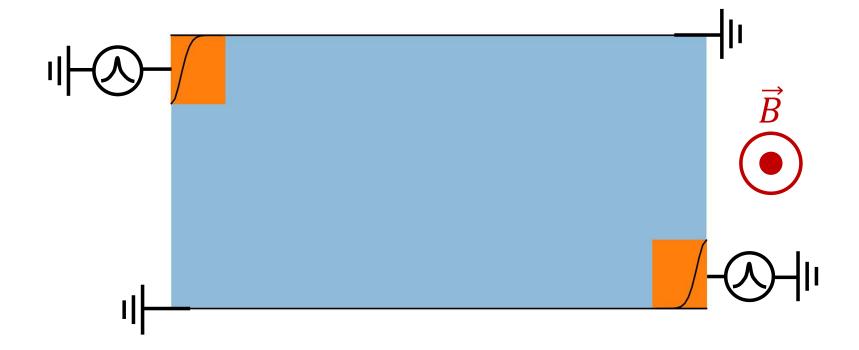
Using electronic lithography we can define gates with \simeq 10 nm resolution or etch away the AlGaAs layer.

Allow for control of edge states.





Using rf EM fields we can generate plasmonic excitations of the edges called edge magnetoplasmons (EMPs)

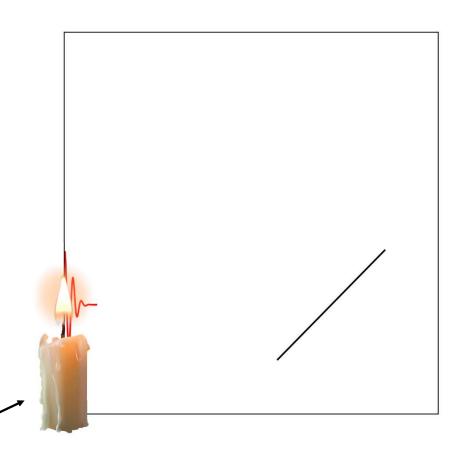




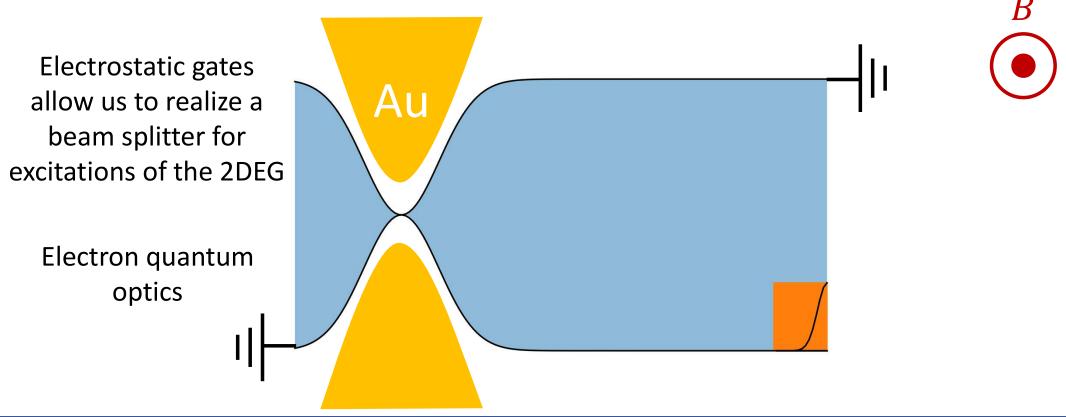
In optics, a beam splitter can be realized by using a semitransparent mirror.

Quantum optics

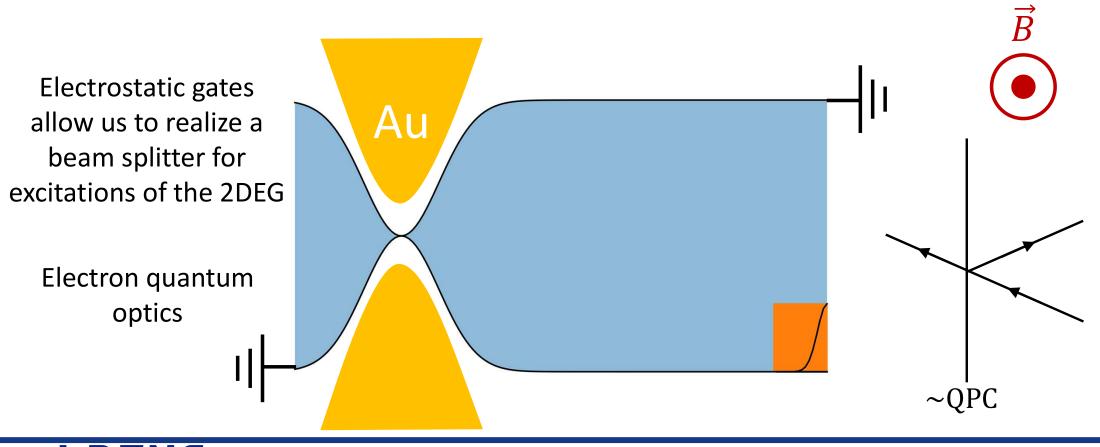
Single photon candle (not sold at Ikea)







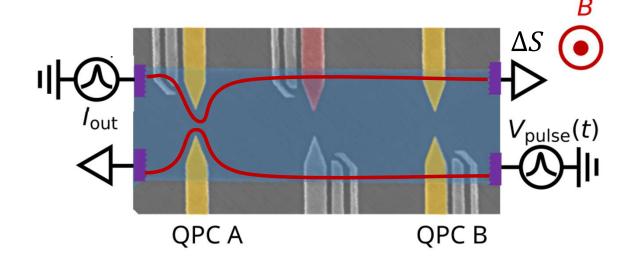




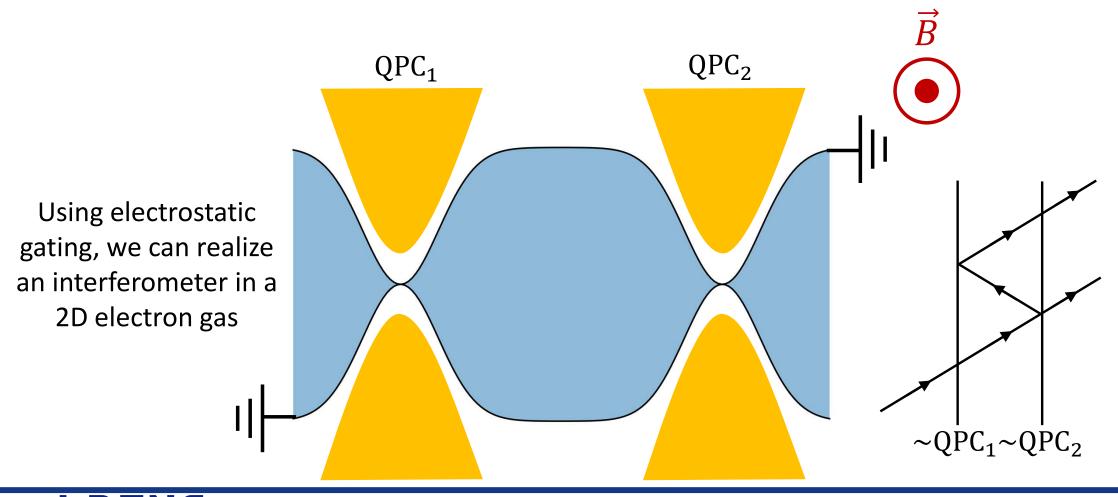


Calibration of pulse: shot noise

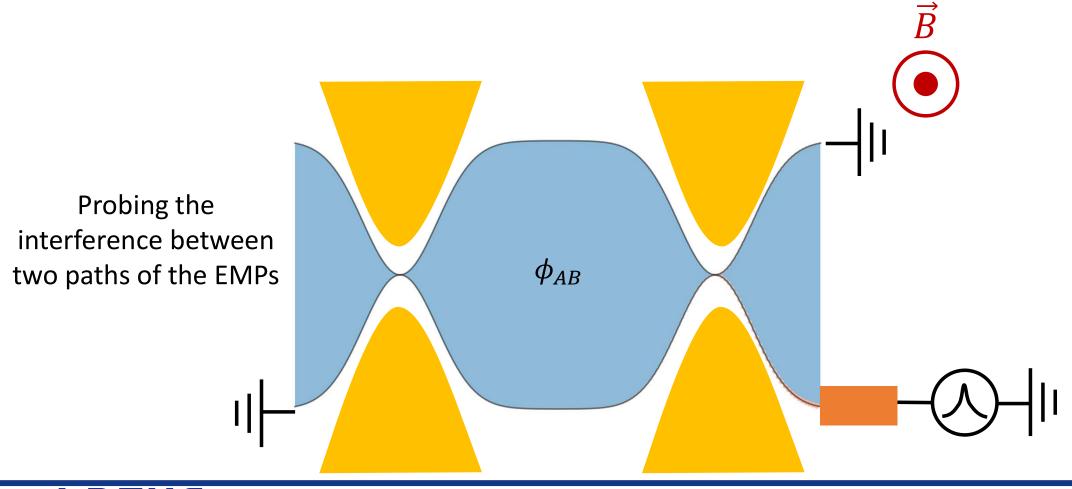
HOM configuration to measure shot noise & characterize Lorentzian pulse



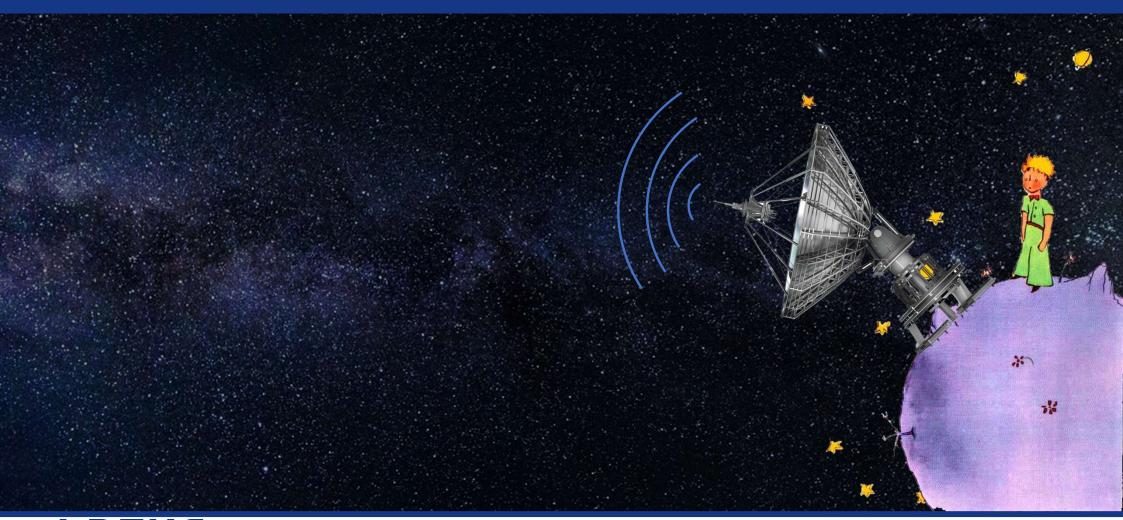








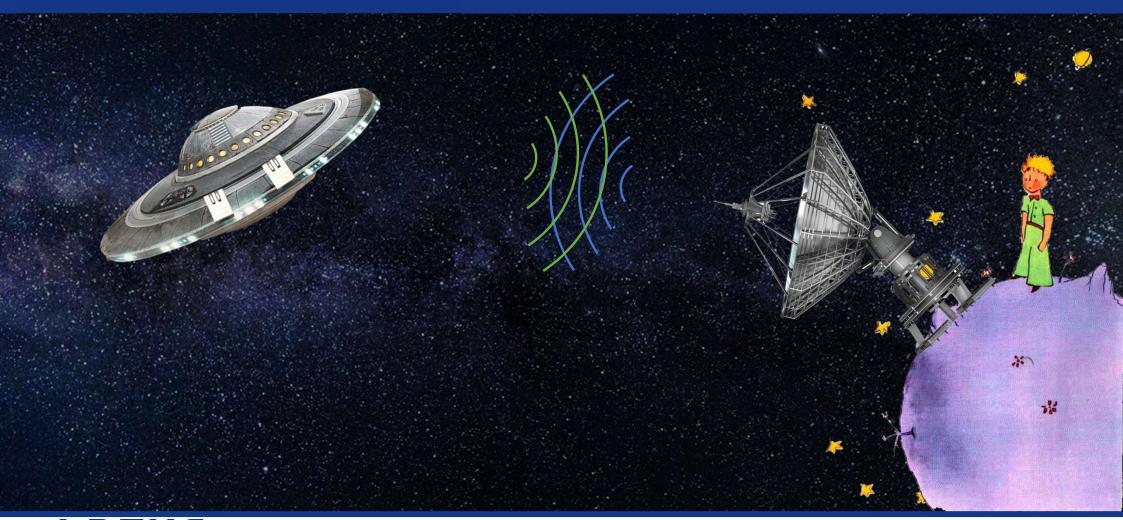




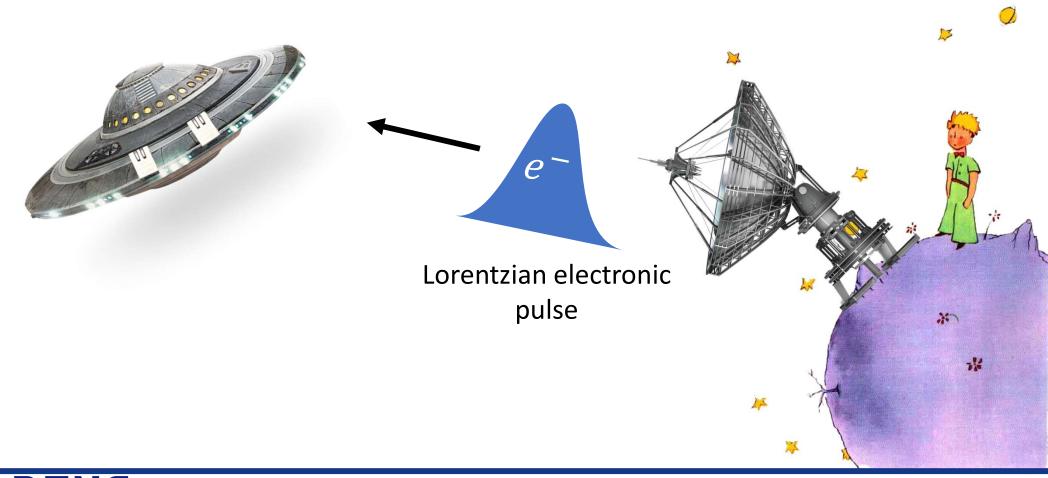




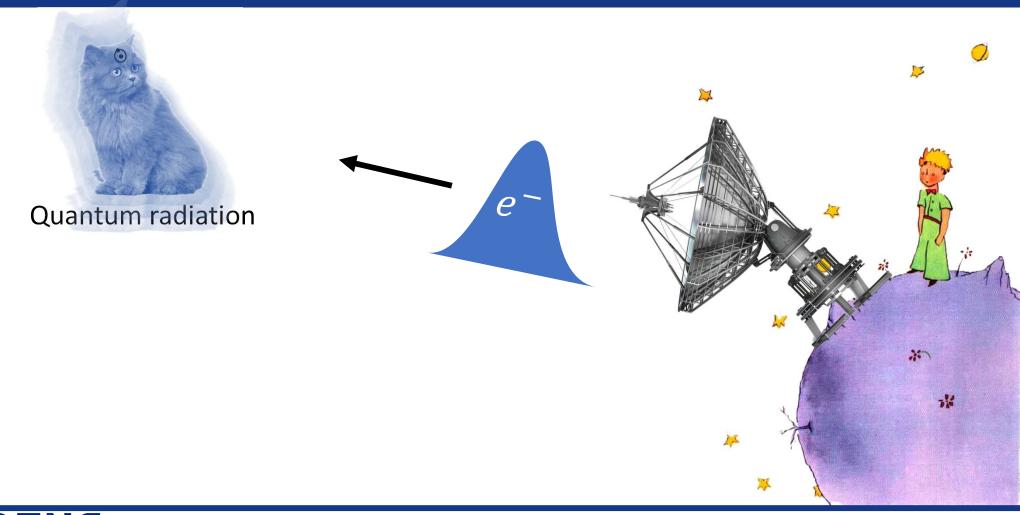




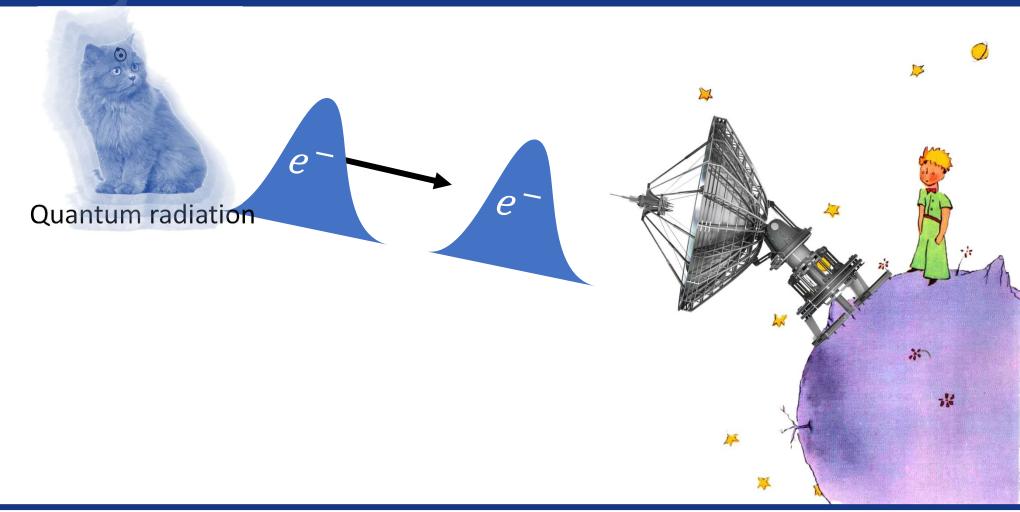




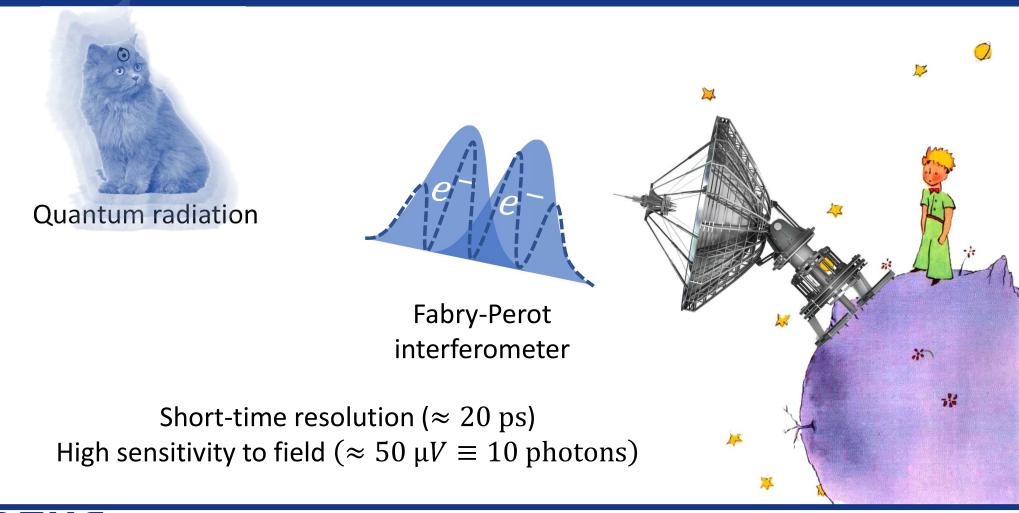




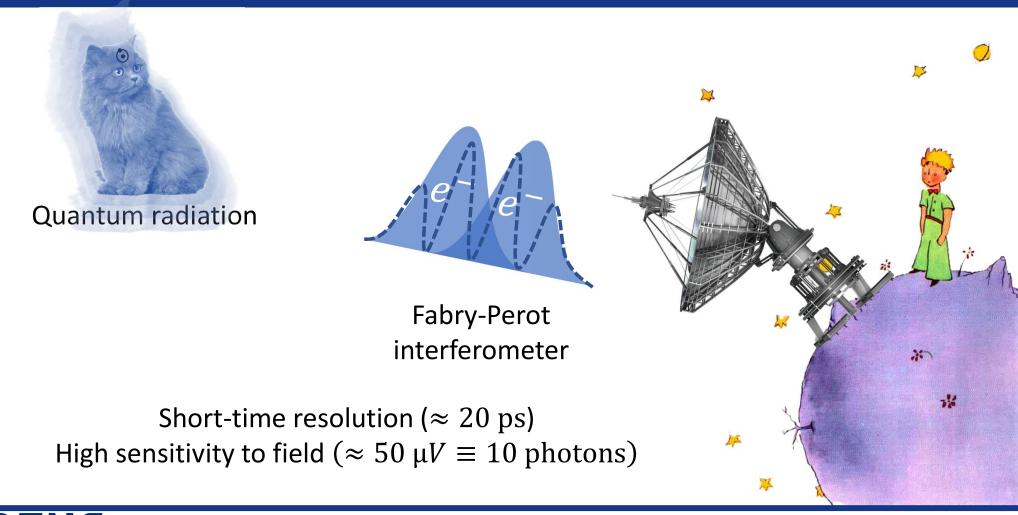






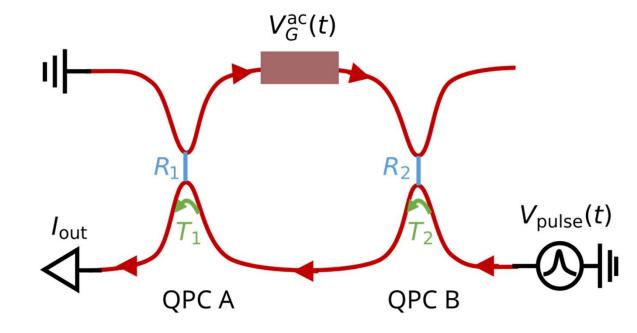






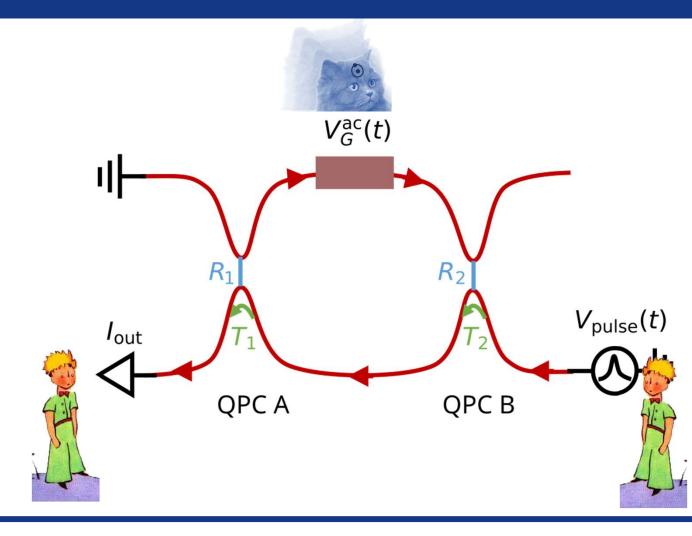


Using Lorentzian shaped pulse to probe classical voltage $V_G^{\rm ac}(t)$



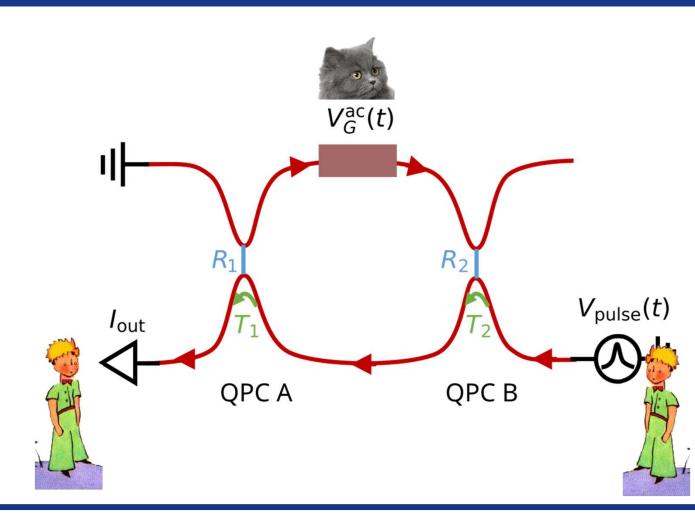


Using Lorentzian shaped pulse to probe classical voltage $V_G^{\rm ac}(t)$





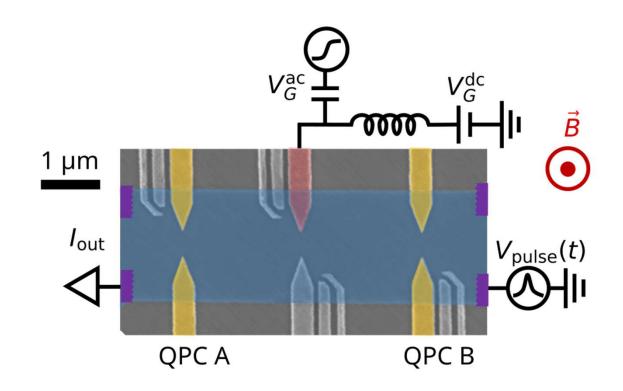
Using Lorentzian shaped pulse to probe **classical** voltage $V_G^{\rm ac}(t)$





AlGaAs/GaAs based 2DEG

2 QPCs & 1 excitation gate





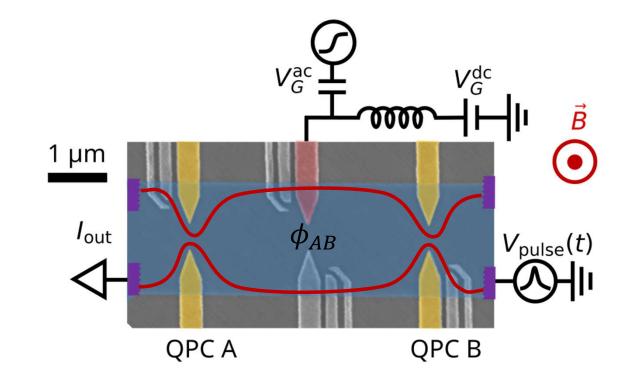
AlGaAs/GaAs based 2DEG

2 QPCs & 1 excitation gate

Device can be used in Fabry-Perot mode

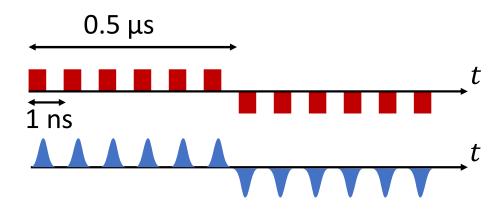
Typical time:

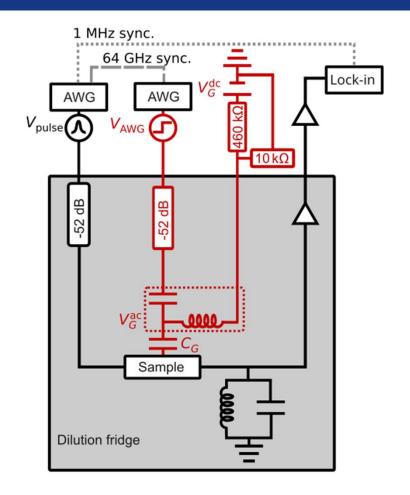
$$au_L =$$
 30 ps $au_e = 35 - 113$ ps





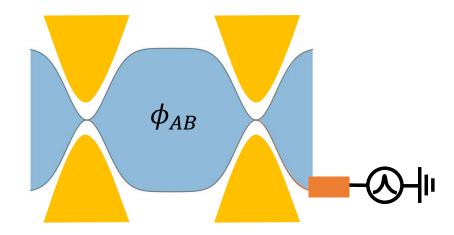
10 GHz pulse/drive period 1 MHz cycle Detection through tank circuit

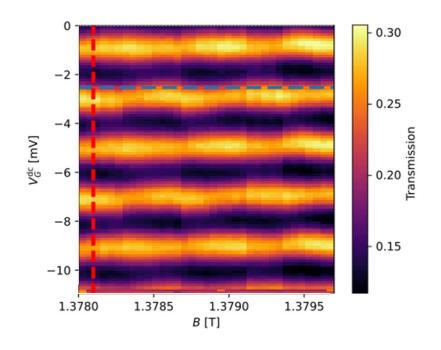






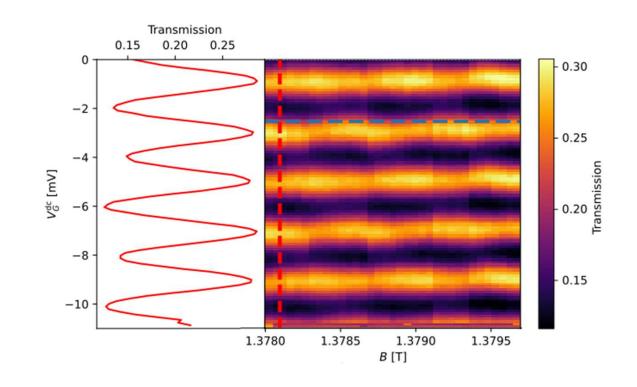
Oscillating behavior as a function of B and $V_G^{\rm dc}$ ($V_G^{\rm ac}=0$)







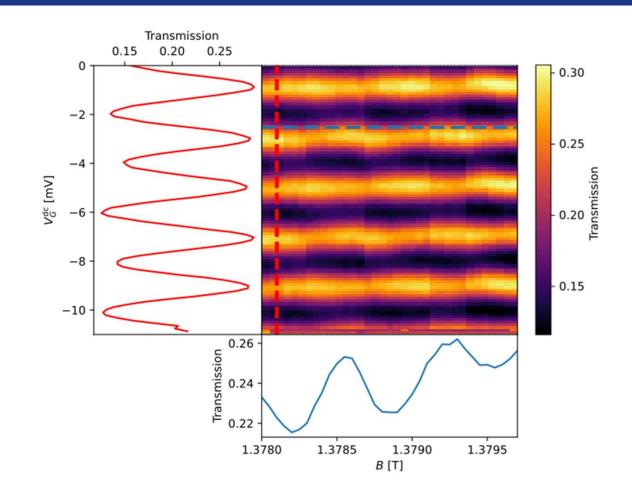
Oscillating behavior as a function of B and $V_G^{\rm dc}$ ($V_G^{\rm ac}=0$)





Oscillating behavior as a function of B and $V_G^{\rm dc}$ ($V_G^{\rm ac}=0$)

Amplitudes are much stronger along $V_G^{ ext{dc}}$ axis due to Coulomb effect





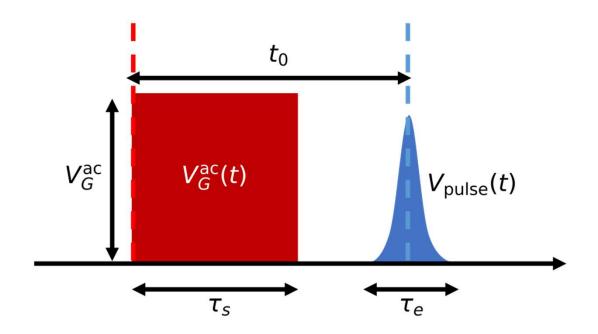
Square drive probed by Lorentzian pulse.

 τ_e : pulse width

 t_0 : time delay

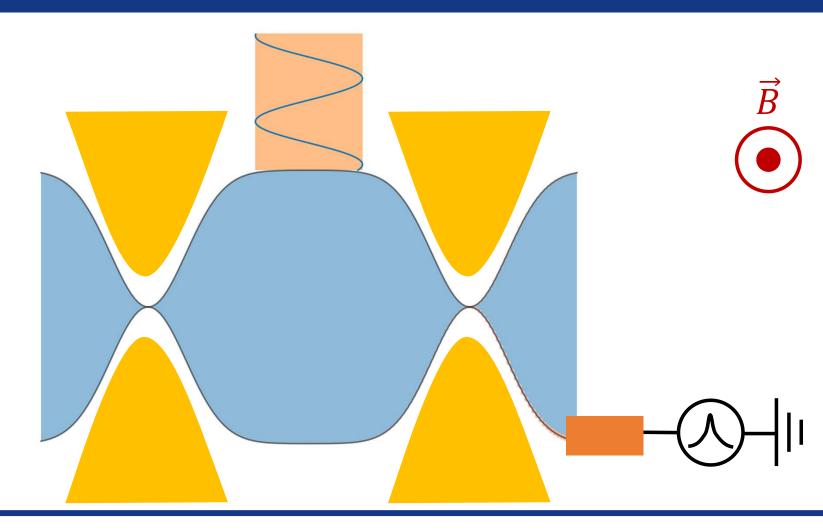
 au_L : time cavity round-trip

 $V_G^{\rm ac}$: Amplitude drive



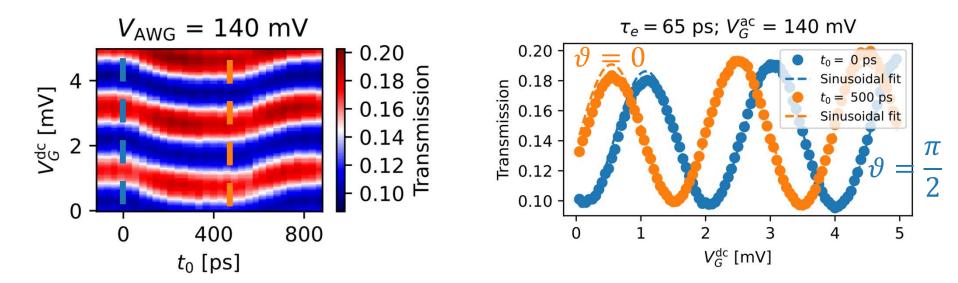


Cavity comparable to EMP size





Fabry-Perot interferometry: results



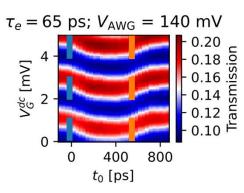
A phase ϑ and contrast C are extracted from sinusoidal fits of the transmission at fixed values of t_0

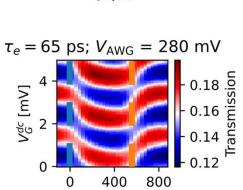


Fabry-Perot interferometry: results

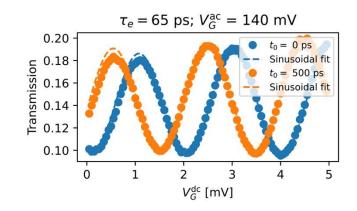
The phase-shift depends on the amplitude of the drive V_G^{ac} .

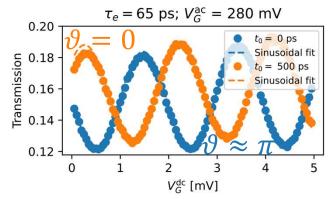
140 mV $\rightarrow \pi/2$ shift 280 mV $\rightarrow \pi$ shift





 t_0 [ps]

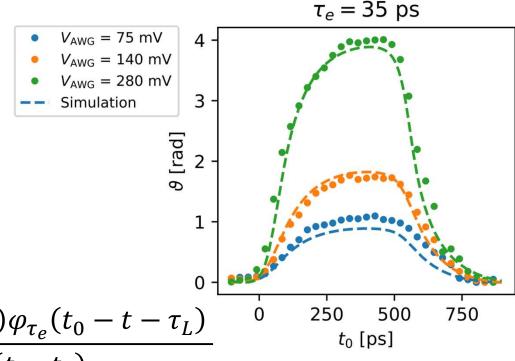






Effect of the drive amplitude:

- phase shift is proportionnal to $V_{
 m AWG}$ with very good agreement
- Theory follows contrast qualitatively but is underestimated. Reproduces well contrast dips.

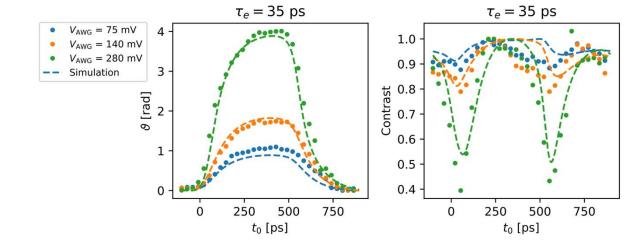


$$C(t_0)e^{i\vartheta(t_0)} = \frac{\int dt \ e^{2i \ \frac{C_G}{e}V_G^{\rm ac}(t)} \varphi_{\tau_e}^*(t_0 - t) \varphi_{\tau_e}(t_0 - t - \tau_L)}{\int dt \ \varphi_{\tau_e}^*(t) \varphi_{-\tau_e}(t - t_L)} \stackrel{0}{\longrightarrow} \frac{250 \ 500 \ 750}{t_0 \ [ps]}$$



Effect of the drive amplitude:

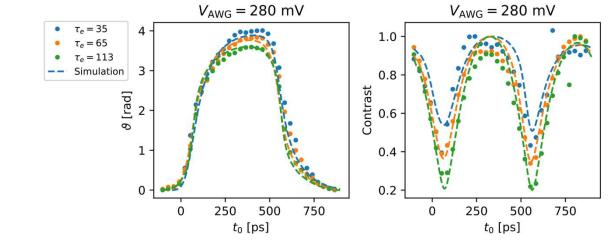
- phase shift is proportionnal to $V_{
 m AWG}$ with very good agreement
- Theory follows contrast qualitatively but is underestimated. Reproduces well position of contrast dips.



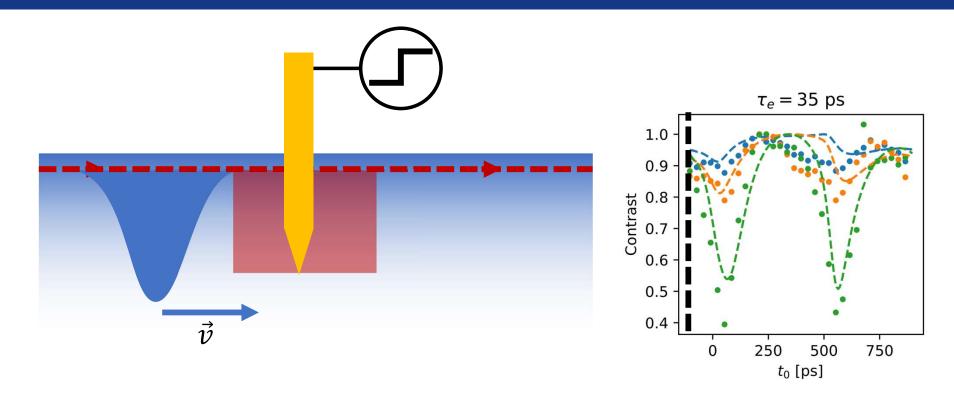


Effect of the pulse width:

- Much smaller effect on both phase
 & amplitude
- Longer pulse diminishes contrast
- Resolution of $50 \mu V$.

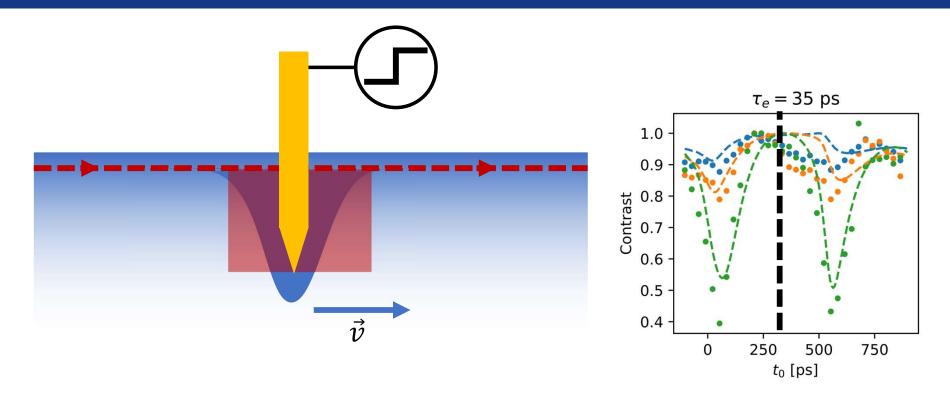






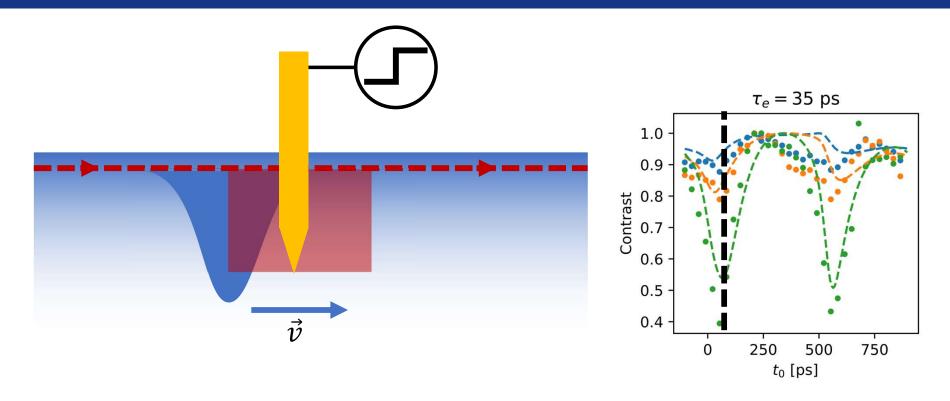
For small t_0 , no overlap between wave function and drive





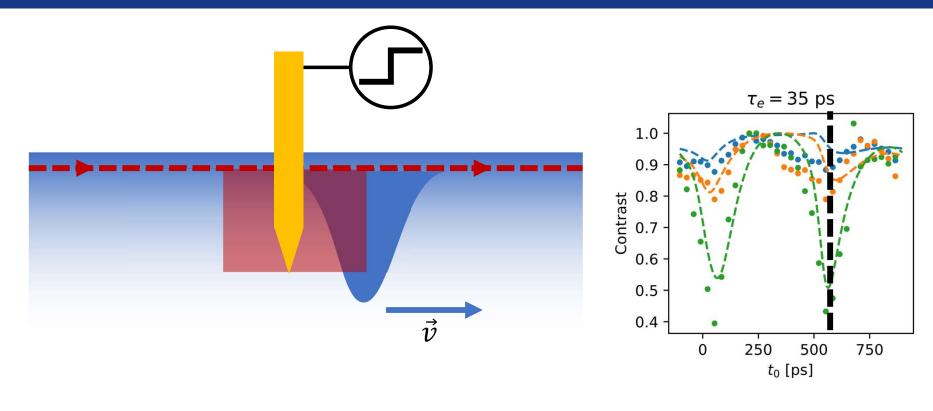
When both drive and pulse coincide, the contrast is roughly identical to that in absence of drive





When pulse only partially overlaps with drive, wavepacket spread reduces contrast depending on position of the electron.

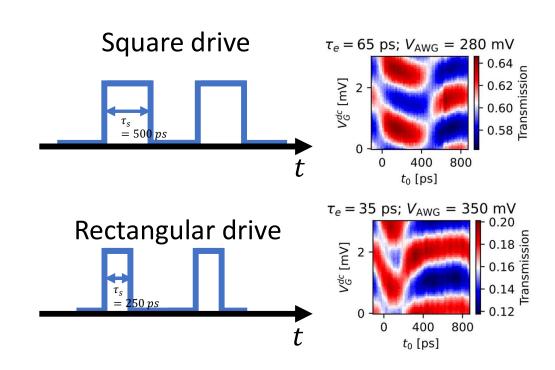




Same on other side

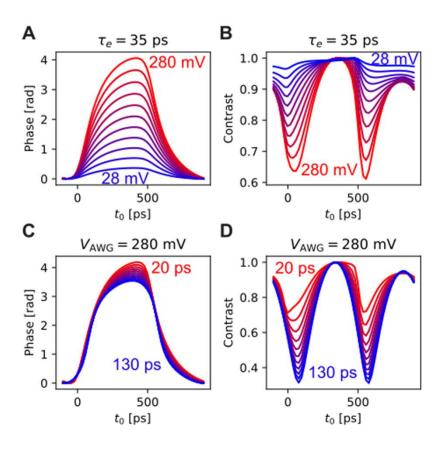


Changing shape of drive from square ($\tau_s = 500 \text{ ps}$) to rectangular ($\tau_s = 250 \text{ ps}$).

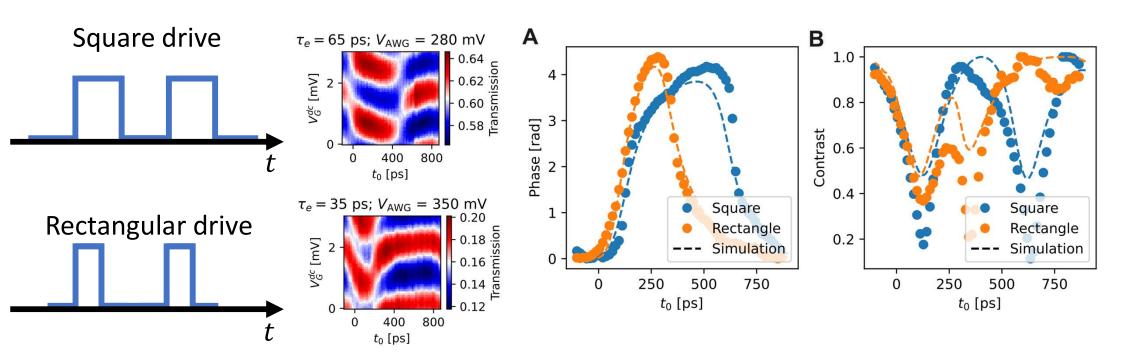




Varying amplitude of and width of pulse







Good agreement on the phase (limited by τ_r). Contrast has quantitative difference but good position of dips.



Conclusions

- Fast sensing probe of time-dependent voltage
- Few 10s of ps time-resolution
- Limited by gate polarization rise-time
- Sensitivity ≈ 10 photons but could be easily improved
- Potential application to detection of squeezing, EMP resonance, field tomography...







Thank you for your attention!









