

# Mapping a critical point on pion condensate line in isospin phase diagram from Lee-Yang zeros

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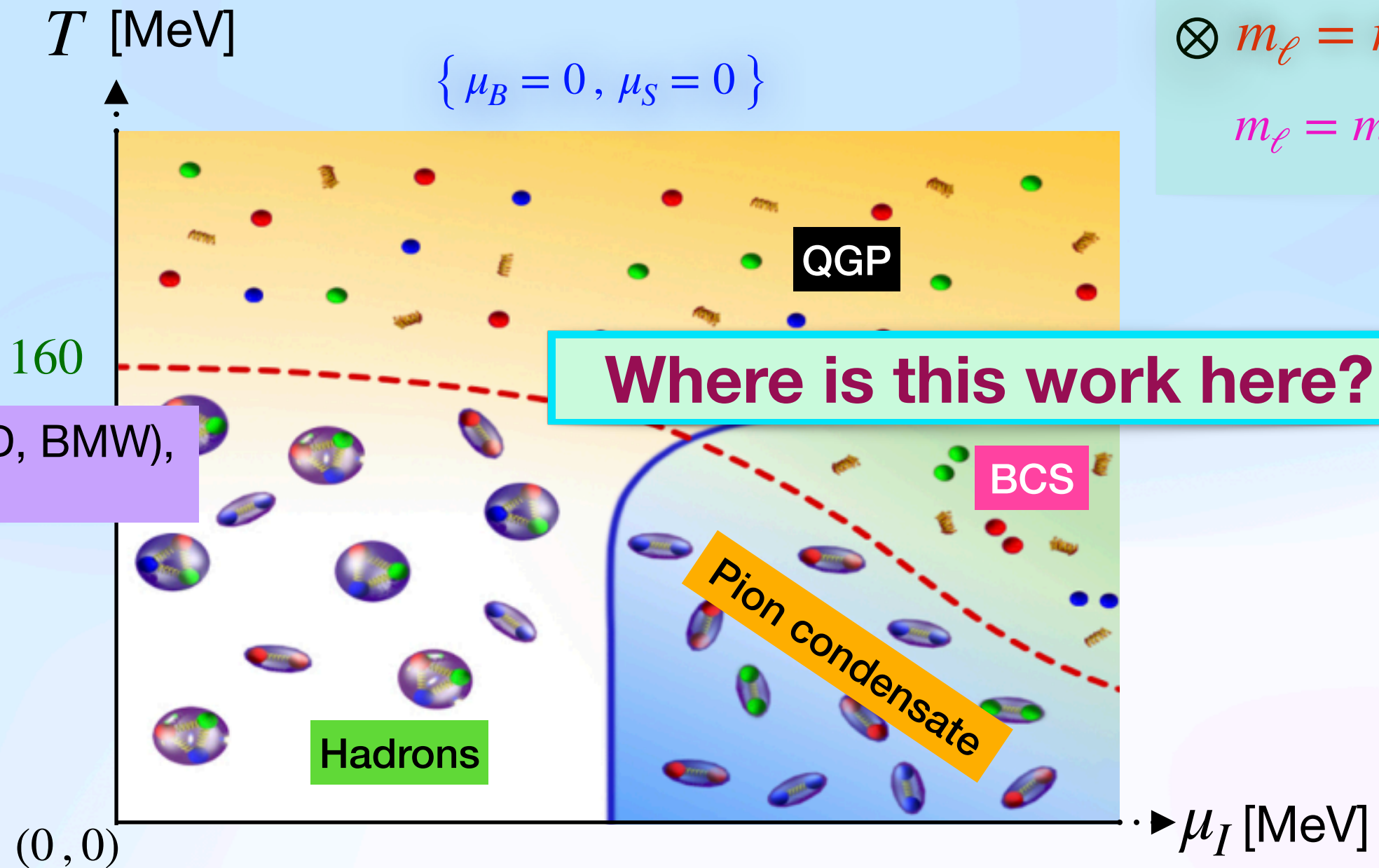
Faculty of Physics, Bielefeld University

Phys.Rev.D 112 (2025) 1, 014511 , arXiv : 2401.14299 [hep-lat]

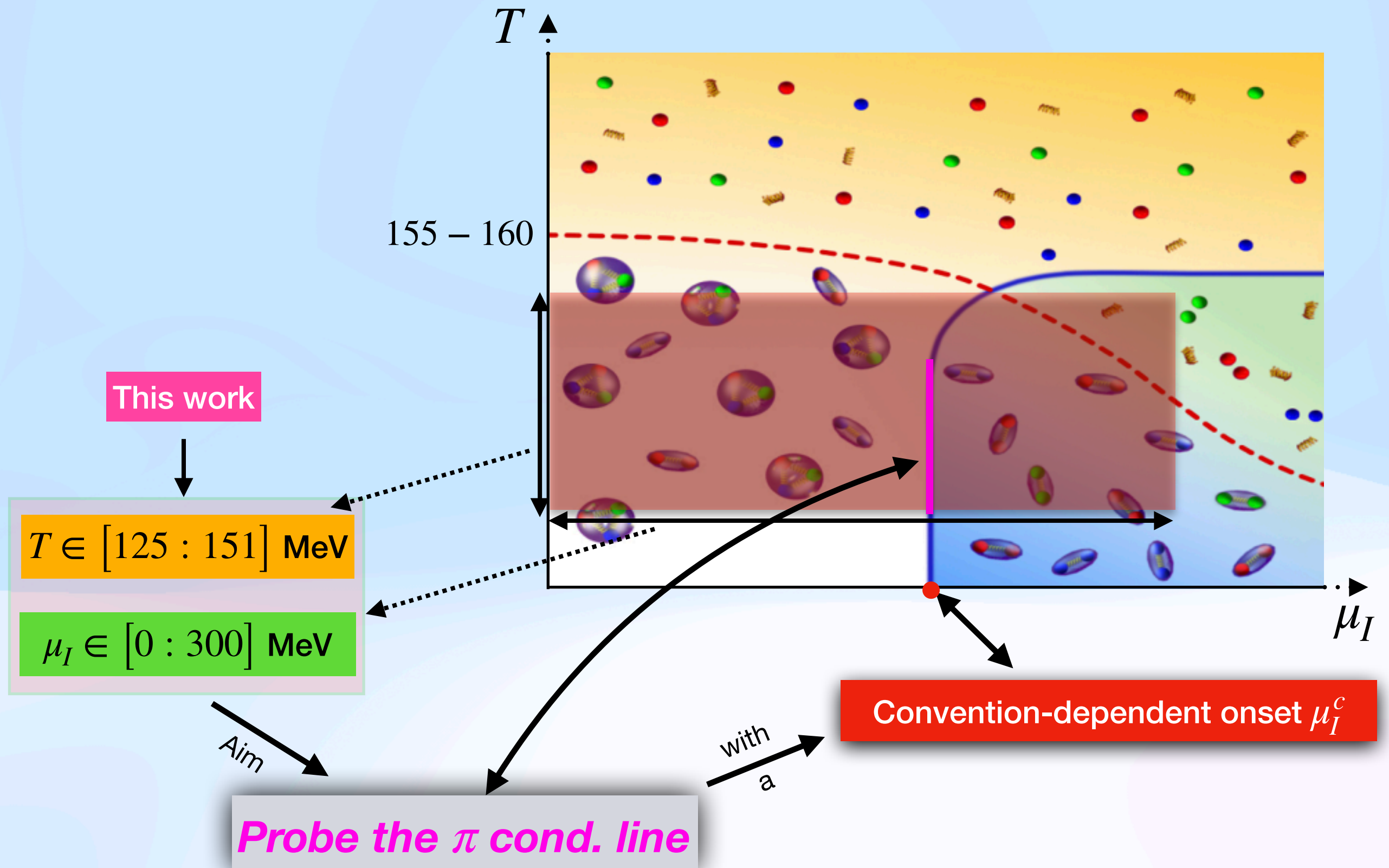
# Plan of talk

- Motivation and domain of this work
- Core methodology and working lattices
- Zeroes, results and their stability
- Mapping a critical point on the critical line
- Radius of convergence and subsequent observations
- Overlap problem and its severity
- Conclusion and Outlook

# Motivation



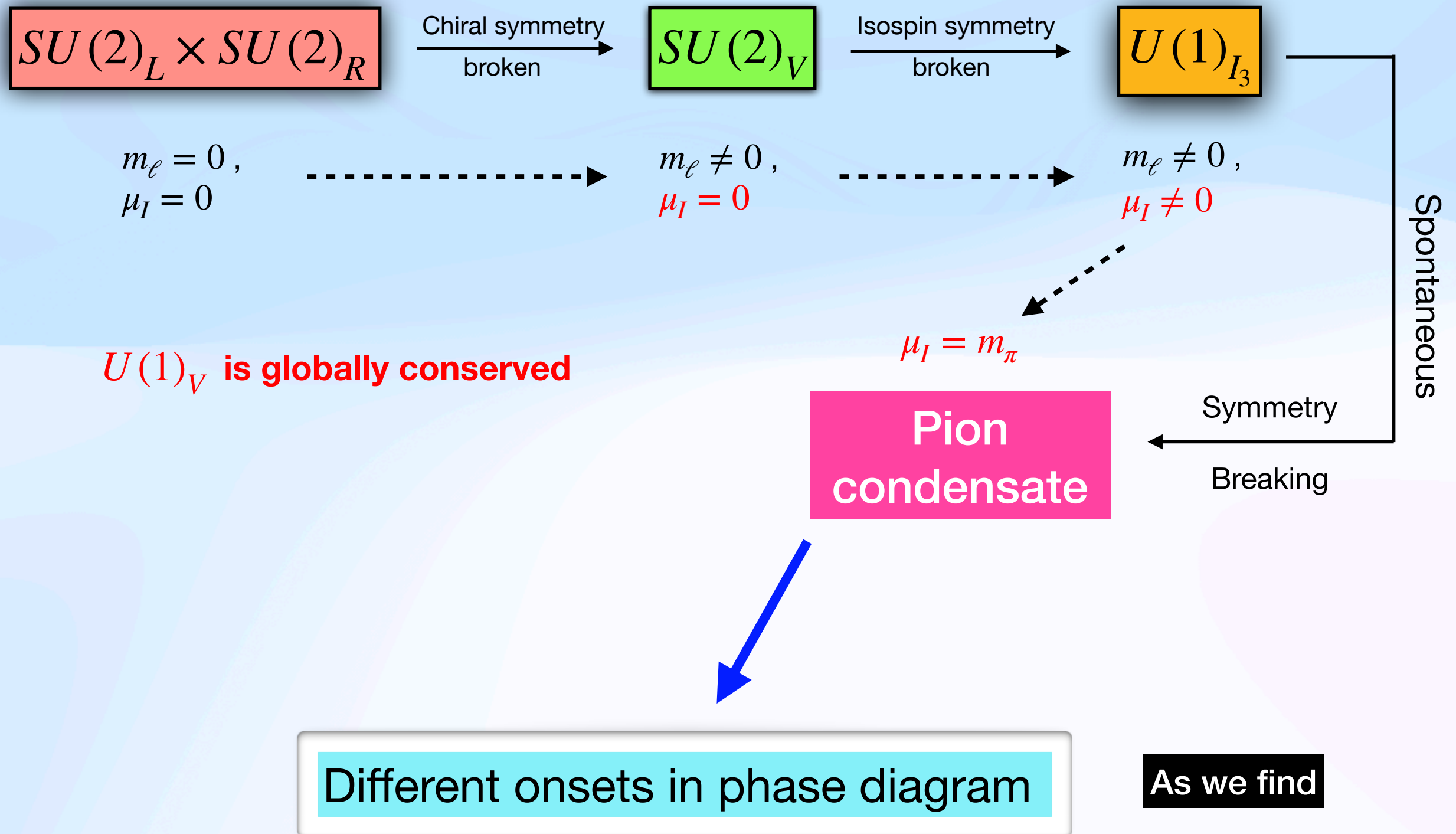
Brandt et. al. : **PRD 97 (2018) 5, 054514**, arXiv : **1712.08190 [hep-lat]**



***So, what happens here??***



# The Theory : Symmetry



# The prevailing notions so far

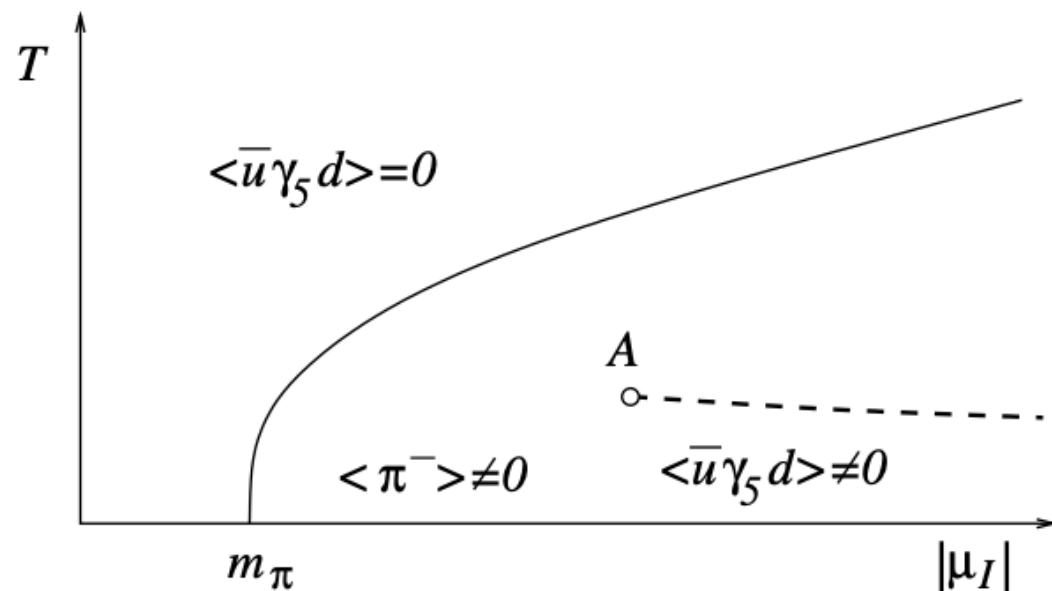
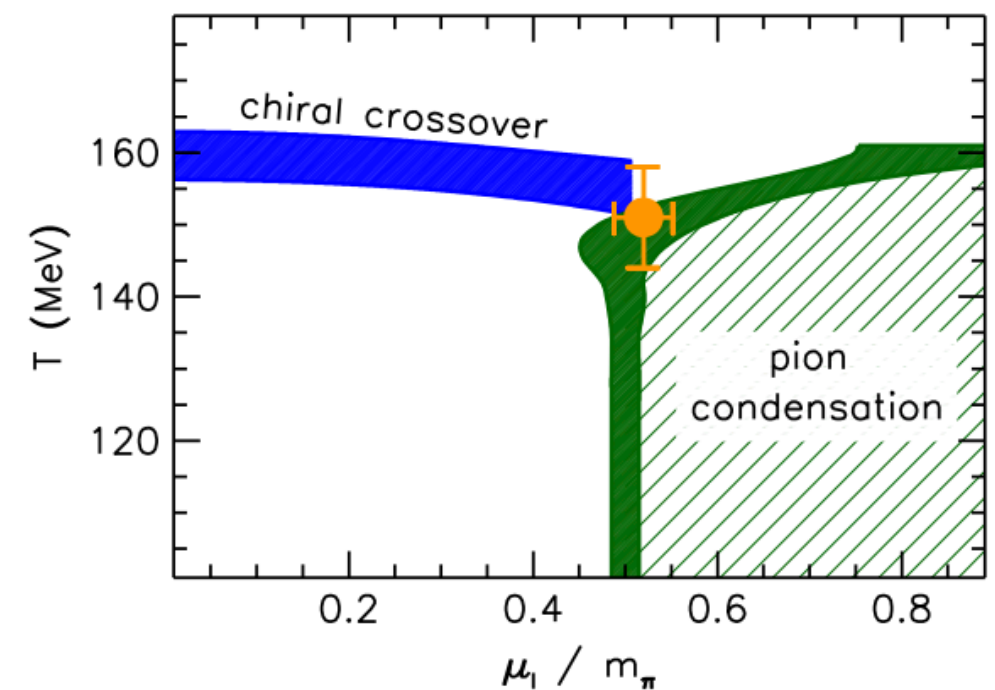


FIG. 1. Phase diagram of QCD at finite isospin density.



Son, Stephanov  
*PRL*

$$\mu_I = \mu_u - \mu_d$$



$$\mu_I^c = m_\pi$$

Adopted here in  
this work

So, what we do here??

Brandt et. al.  
*PRD*

$$\mu_I = \frac{\mu_u - \mu_d}{2}$$



$$\mu_I^c = \frac{m_\pi}{2}$$

On Lattice (Lattice QCD)

In complex isospin densities

Determine the **Lee-Yang zeros**  
of the **QCD partition function**

**Closest zero** and  
**Radius of convergence (RoC)**  
as  $f(\mu_I, T)$

No assumption of the  
nature of phase transition

Regulator-free

See manifestations of this  
**pion condensate onset**

Can RoC be indicator of the  
condensate critical line ???

Implications of this **radius  
of convergence**

WORKING FORMULATION??

# Working formulation

To evaluate the zeros, we use the familiar Newton-Raphson formulation :

$$\mu_{n+1} = \mu_n - \frac{\mathcal{Z}(\mu_n)}{\mathcal{Z}'(\mu_n)}, \quad \mu \equiv \mu_I$$

- $\mathcal{Z}'(\mu_n) = \left[ \partial \mathcal{Z} / \partial \mu \right]_{\mu=\mu_n}$  ( Isospin partition function  $\longrightarrow \mathcal{Z}$  )
- $\mu_n$  ( $n \geq 0 \in \mathbb{Z}^+$ )  $\longrightarrow$  estimate obtained on  $n^{\text{th}}$  iteration . What iteration???
- Start from an initial guess  $\mu_0$  and continue iteratively , unless  $|\mu_{n+1} - \mu_n| \leq \epsilon$

In this work, we choose  $\epsilon = 0.002$

Tolerance

So, how is  $\mathcal{Z}(\mu_I)$  defined here???

$$\mathcal{Z}(\mu_I, V, T) = \int DU e^{-S_g[U]} \left[ \det M(\mu_I, V, T) \right]$$

$$= \left\langle \exp \left[ \sum_{n=1}^N \frac{\mu_I^n}{n!} D_n^{(u)} \right] \right\rangle$$

**Unbiased exponential  
resummation**

**SM**, Hegde ; **PRD 108 (2023) 3, 034502**, arXiv : **2302.06460 [hep-lat]**

$\langle O \rangle$  calculated on gauge ensemble generated at  $\mu_I = 0$

$$D_n = \left. \frac{\partial^n}{\partial \mu_I^n} \ln [\det M] \right|_{\mu_I=0}, \text{ and } D_n^{(u)} \text{ contains } \underline{\text{unbiased corrections}}$$

These corrections are important to reproduce the exact Taylor coefficients order-by-order. (Upto fourth order here)

**SM**, Hegde, Schmidt ; **PRD 106 (2022), 3, 034504**, arXiv : **2205.08517 [hep-lat]**

And now the working lattice ...

# Lattice setup

- (2+1)-flavor Highly Improved Staggered Quark (HISQ) action
- Working lattices  $\longrightarrow 32^3 \times 8$  lattices (higher volumes in future)
- Total of **20K configurations** ( $N_{conf} = 20K$ )
- Working temperatures  $\longrightarrow 125 \leq T \leq 171$  MeV
- Physical light and strange quark masses, for each temperature ( $m_\ell = m_s / 27$ )

**How do we implement the simulations???**



# Implementation

- Choose the initial complex  $\mu_0$  from the complex set

$$S = \{\mu_0 : 0 \leq \text{Re}(\mu_0) \leq 2, 0 \leq \text{Im}(\mu_0) \leq 2\}$$

in steps of 0.1 on each direction, along  $\text{Re}(\mu_0)$  &  $\text{Im}(\mu_0)$

- Choose the tolerance value

$$\epsilon = 0.002$$

- Choose the upper bound of the number of iterations

$$N_{max} = 10^8$$

with  $N_{conf} = 20K$ ,  $N_B = 50$  (bootstrap samples)

And :

# WORKFLOW

Start with a  $\mu_0$  value  $\in S$

$\mu_{NR}^{(1)}$

$\mu_{NR}^{(2)}$

.....

$\mu_{NR}^{(49)}$

$\mu_{NR}^{(50)}$

$\mu_{NR}^{(b)}$   $\longrightarrow$  estimate of the Newton-Raphson root in  $b^{th}$  bootstrap sample

Final estimates :

$$\mu_{NR}(\mu_0) = \frac{1}{50} \sum_{b=1}^{50} \mu_{NR}^{(b)}$$

← Mean

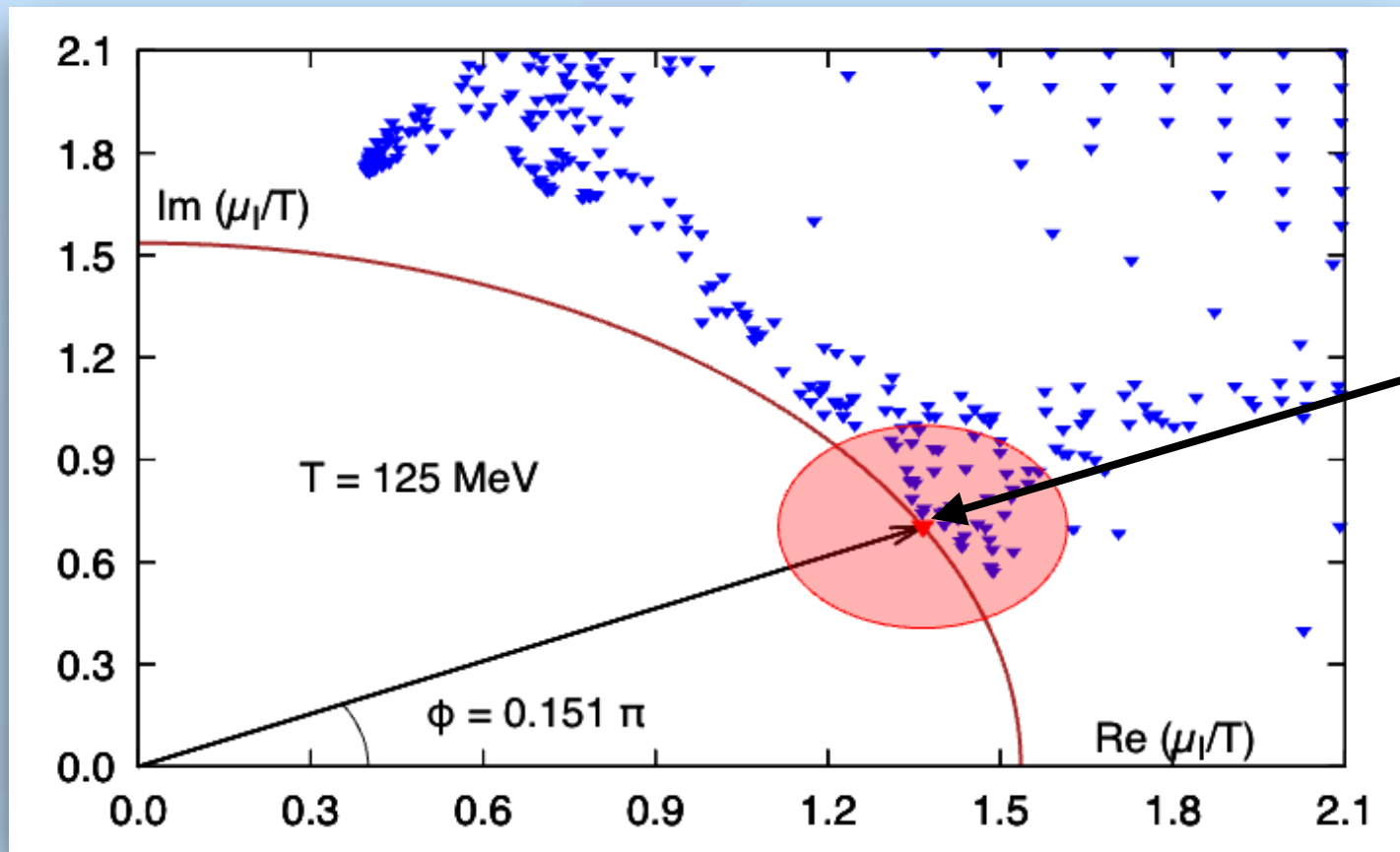
$$\sigma_{NR}^2(\mu_0) = \frac{1}{50} \sum_{b=1}^{50} \left\{ \mu_{NR}^{(b)} - \mu_{NR}(\mu_0) \right\}^2$$

↖ Variance

Follow this for all the other  $\mu_0$

RESULTS ???

# Results



Complex  $\mu_I$  plane

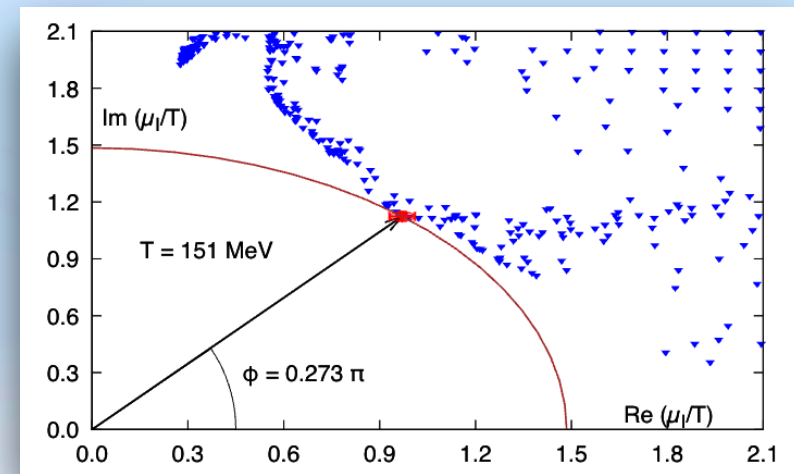
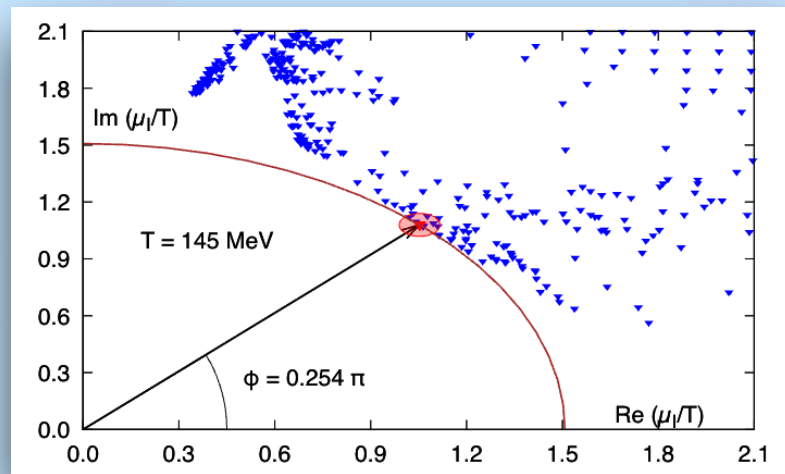
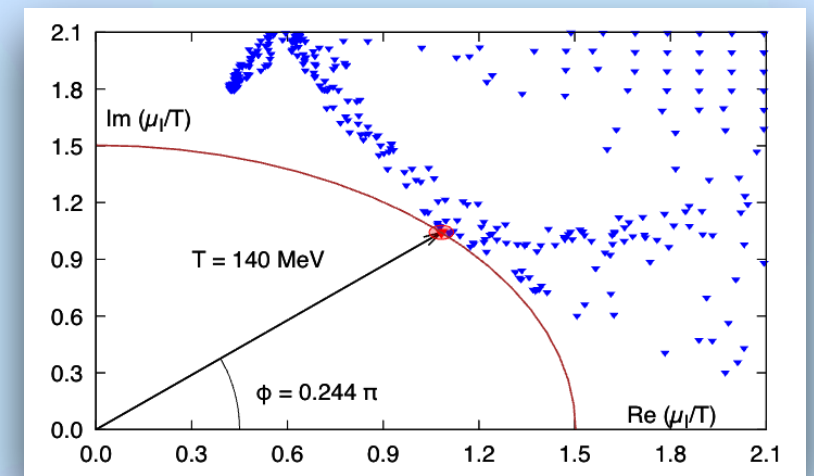
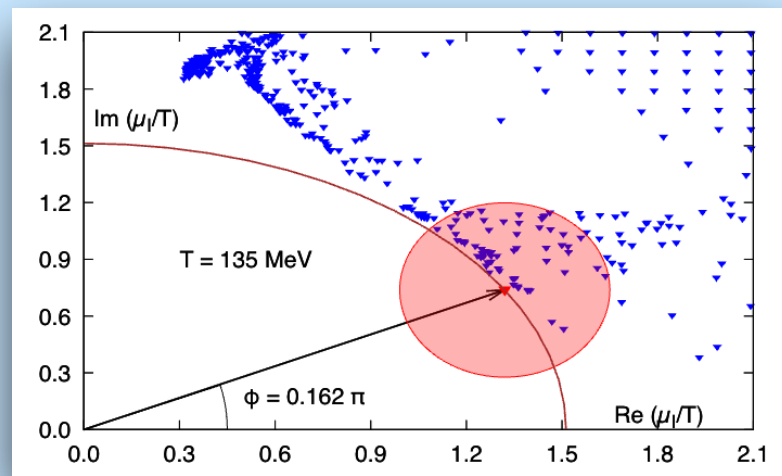
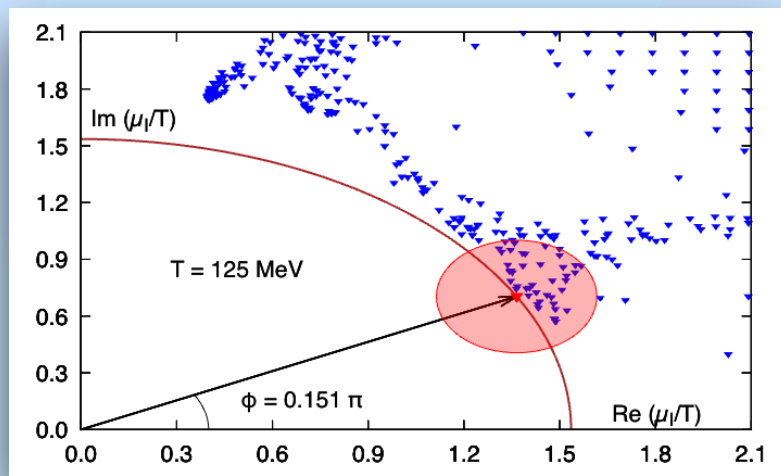
**Nearest** Lee-Yang  
zero  $\mu_I^0$

**Closest** to point of expansion  
→ origin (0, 0) here

- Individual error bars on roots are **only** shown here for  $\mu_I^0$ .
- Elliptical representation of errors : major and minor axes.
- **This line** shows **one quarter** of the **circle of convergence**.
- $\mu_I^0$  makes an **angle**  $\phi$  (radian units) **with the real  $\mu_I$  axis**.

$$\phi = \tan^{-1} \left[ \frac{\mu_{I,i}^T}{\mu_{I,r}^T} \right]$$

Other  $T's??$

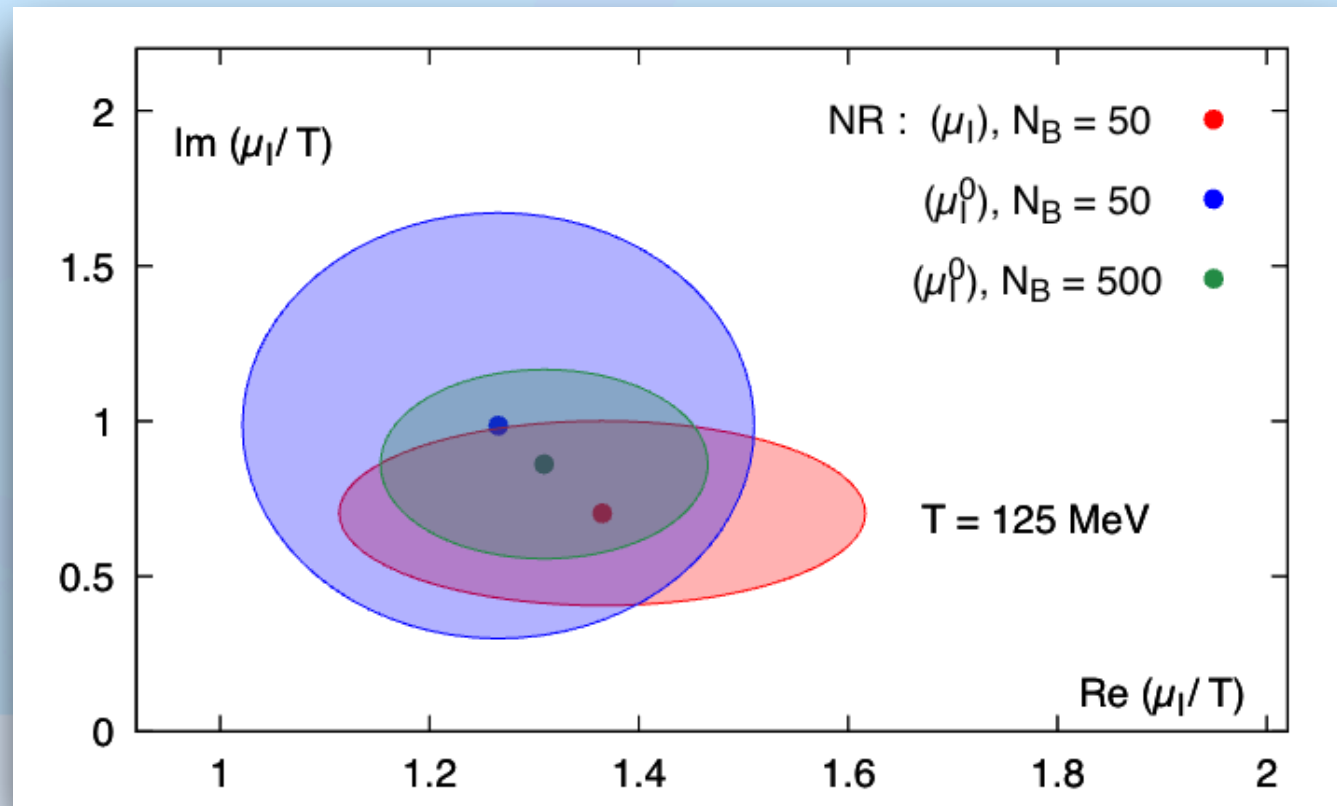


- $\mu_I^0$  approaching real axis with reducing  $T$  (reducing angle  $\phi$  )
- Expecting to find **real**  $\mu_I^0 \Rightarrow$  **genuine critical point** at **lower**  $T$

How stable (reliable)  
are these results ??

However

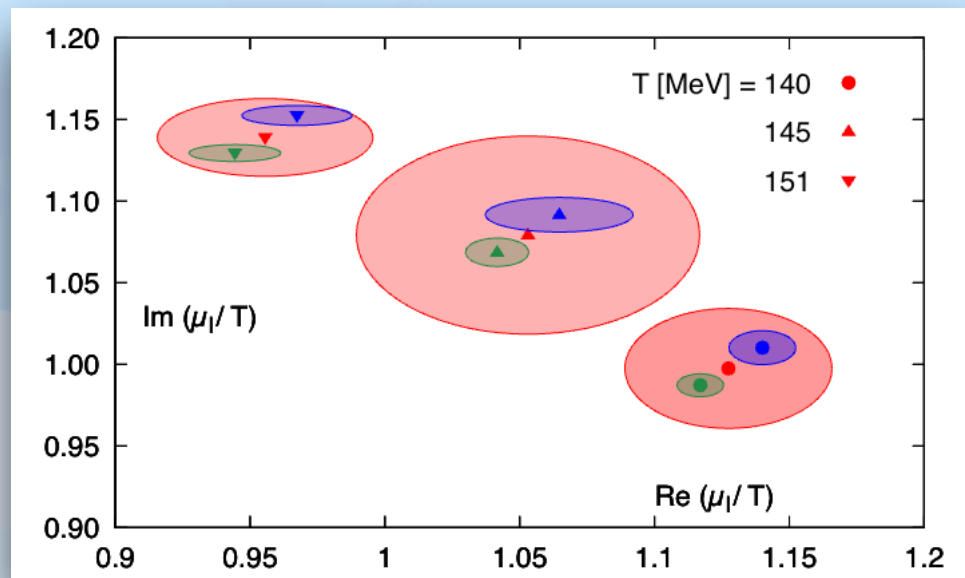
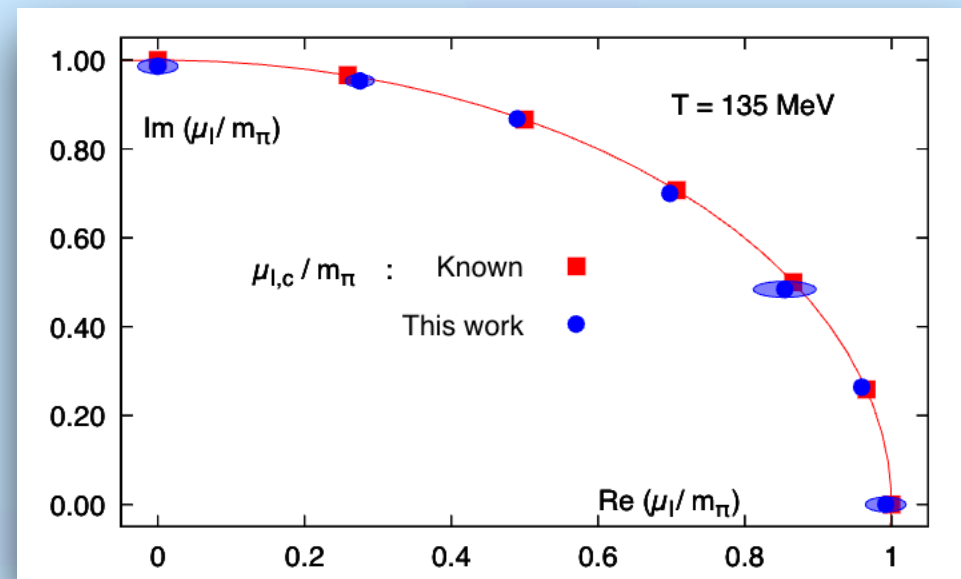
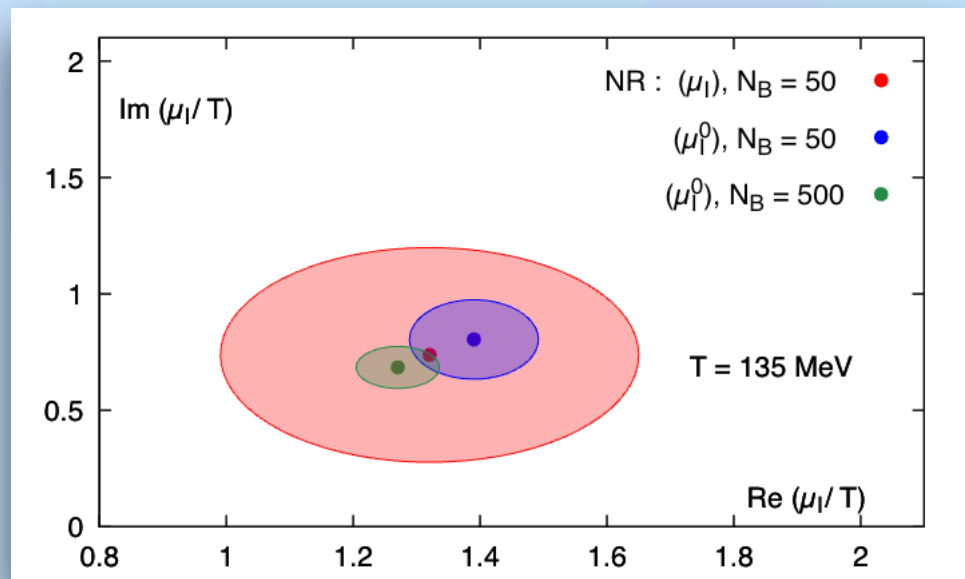
# Stability of results



New estimates ← using  
 $\mu_I^0$  as initial guesses

- **Good agreement (overlap)** with the **old**  $\iff$  **new estimates** of the roots / zeros
- The present estimates of zeros **are reliable !!!**

What about other  $T$  ??



All possible **critical points** for  
 $T = 135$  MeV as **initial guesses**

$$\mu_0 = (m_\pi \cos \phi, m_\pi \sin \phi), \quad \phi = \left\{ \frac{n\pi}{12}, 1 \leq n \leq 6 \right\}$$

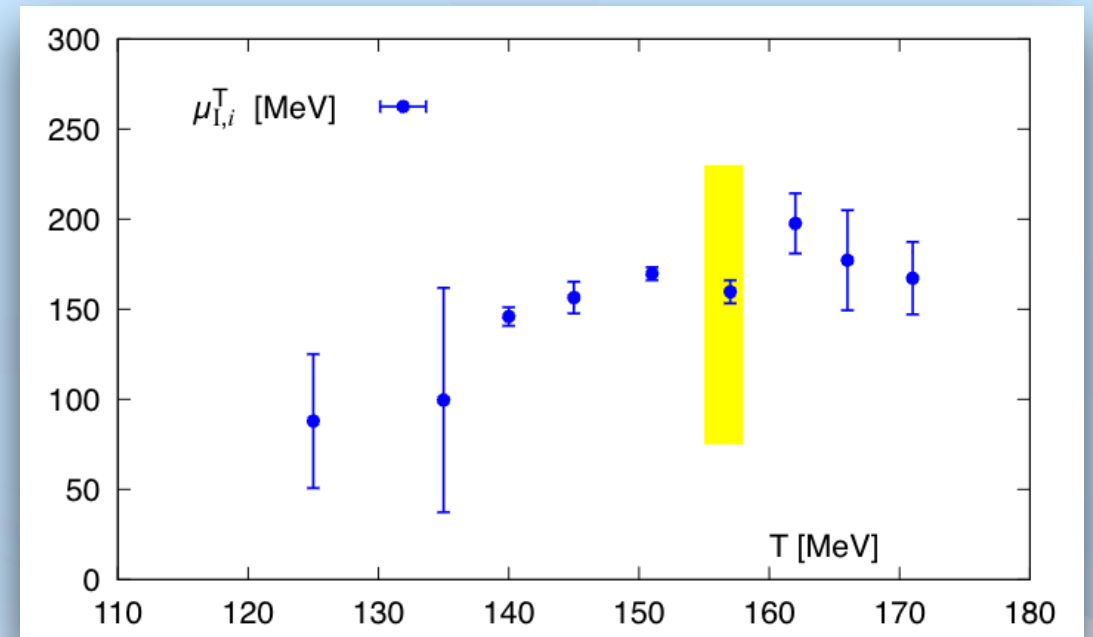
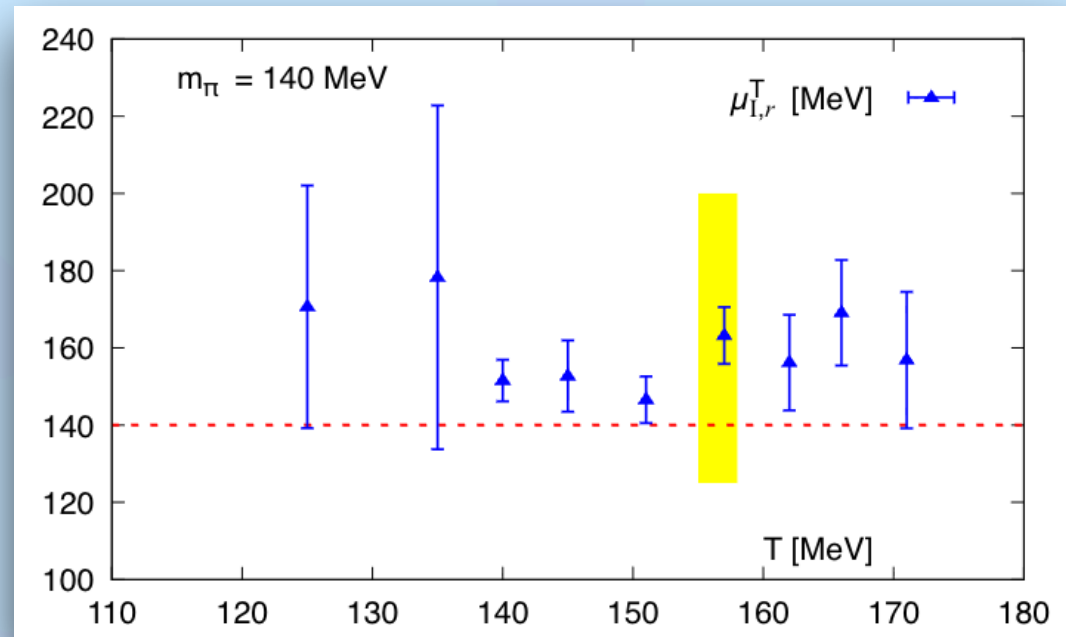
**Commendable agreement** upto  $T = 151$  MeV

Reliable therefore, at least for this work

$\therefore$  *Let's analyse these 0's*



# Real and imaginary parts of $\mu_I^0$



- $\mu_{I,r}^T \sim \text{Re}(\mu_I^0)$  close to  $m_\pi$  for **lower values of  $T$**

- **Monotonically reducing**  $\mu_{I,i}^T \sim \text{Im}(\mu_I^0)$  with **decreasing  $T$**

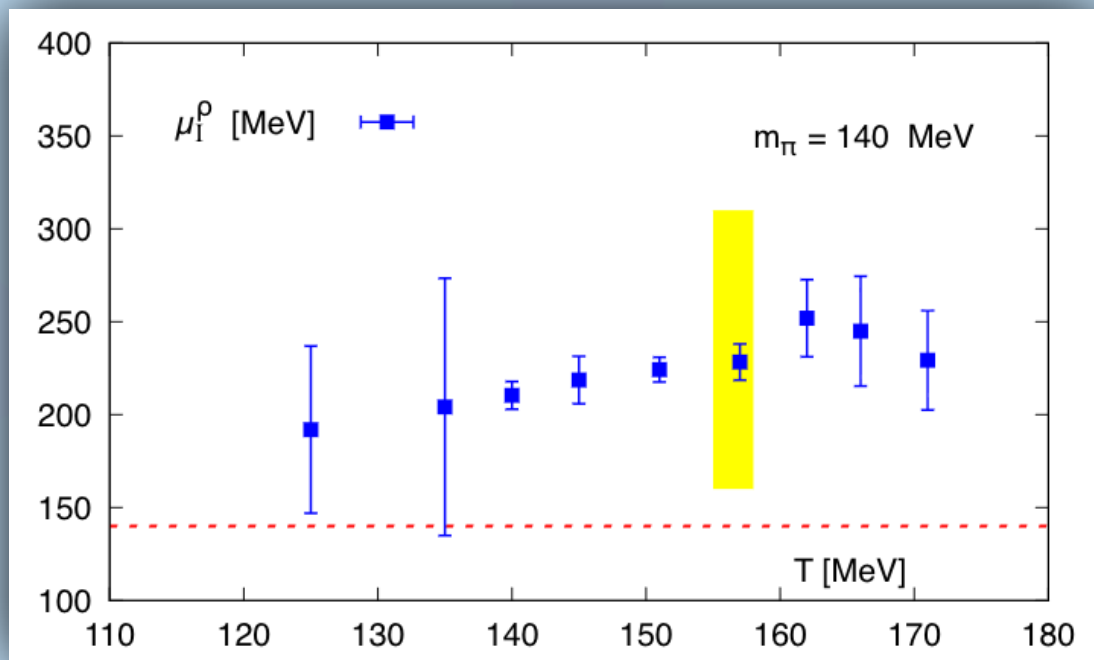
Also, **Increasing errors** for  $\mu_{I,r}^T$  and  $\mu_{I,i}^T$  for  $T < 140$  MeV



signs of possible “**vicinity of critical point(s)**”

**What about the coveted RoC ???**

# Radius of convergence (RoC) :



( Note,  $|\mu_I^0| = \mu_I^\rho / T$  )

$$\mu_I^\rho = \sqrt{\left(\mu_{I,r}^T\right)^2 + \left(\mu_{I,i}^T\right)^2}$$

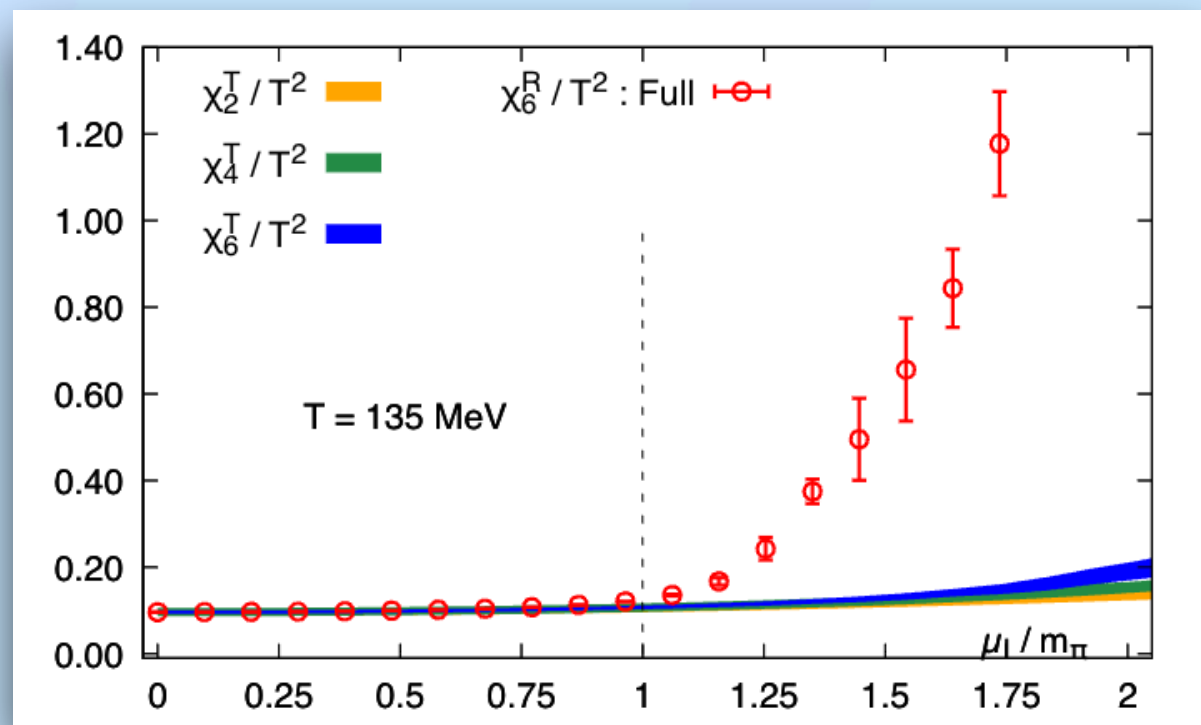
Dimensional ( in MeV units )

- **Monotonic reduction** :  $\mu_I^\rho$  decrease with  $T$  for  $125 \leq T \leq 151$  MeV

Qualitatively consistent with  $\chi PT$  predictions

- $\mu_I^\rho \sim m_\pi$  , within errors for  $T = 135$  MeV

Can it indicate possible critical point / line???



- Resummed susceptibility  $\chi_6^R$  **deviates sharply** from Taylor counterpart results for  $\mu_I > m_\pi$  (**good agreement** for  $\mu_I < m_\pi$ )

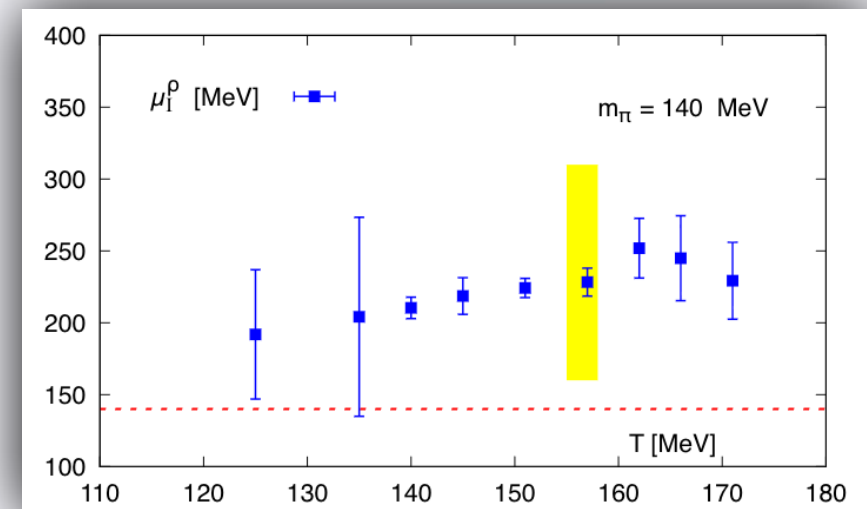
similar  
to

Borsanyi et. al. : **PRD 109, 054509 (2024)**,  
arXiv : **2308.06105 [hep-lat]**

Manifesting

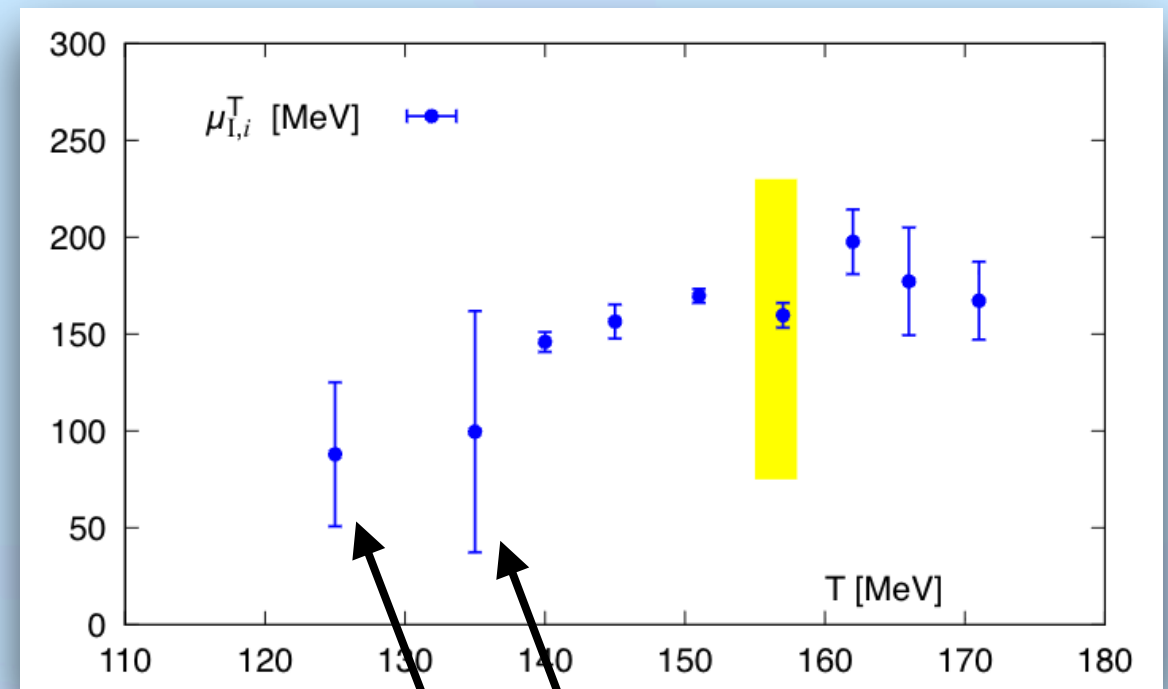
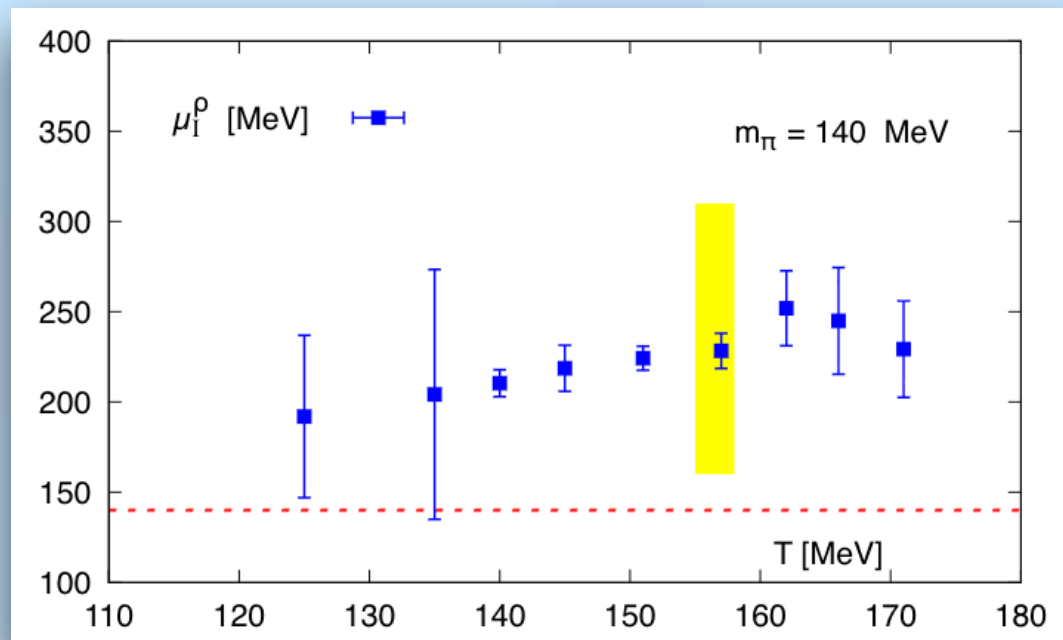
**divergence**

**Possible phase transition** signatures



RoC :  $\mu_I^\rho \longrightarrow$  **good indicator** of  $\pi$ -boundary , at least in low  $T$

**However**



$$\mu_{I,i}^T \neq 0$$

for  $T \leq 135$  MeV

Non-zero **imaginary** part

$(\mu_I^c, T_I^c) = (m_\pi, 135)$  is  
**NOT**  
**A CRITICAL POINT**

So, analyse lower  $T...$

But

No further data for  $T < 125$  MeV  
available on these lattices



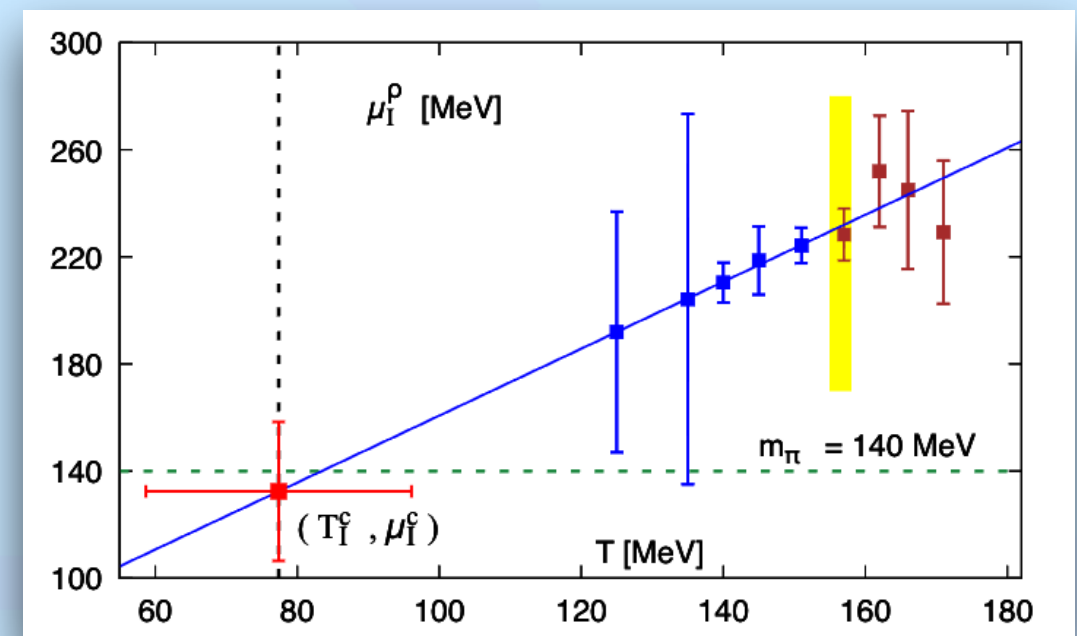
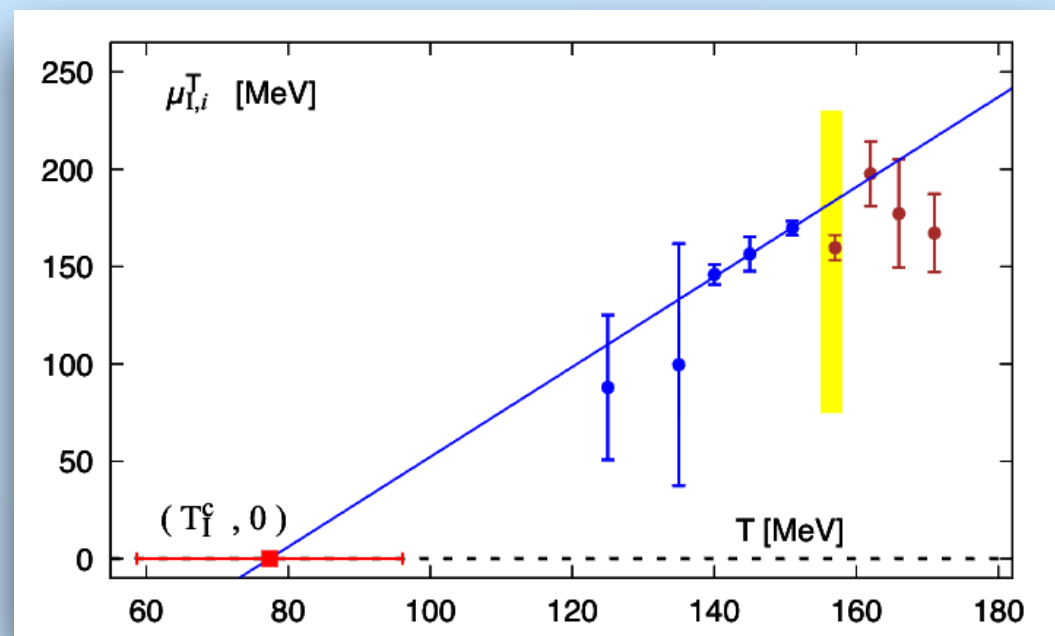
**Extrapolate**  $\mu_{I,i}^T$  and  $\mu_I^\rho$  in  $T$ , to lower  $T$



1. Find :  $T = T_I^c \ni \mu_{I,i}^T = 0$  , and hence,  $\mu_I^0(T_I^c) \in \mathbb{R}$
2. Find :  $\mu_I^\rho$  at  $T = T_I^c$  :  $\Rightarrow \mu_I^c$  { since, at  $T = T_I^c$  ,  $\mu_I^\rho = \text{Re}(\mu_I^0)$  }



Thus, obtain an estimate of a critical point  $(\mu_I^c, T_I^c)$



**Linear extrapolation** of  $\mu_{I,i}^T$  and  $\mu_I^\rho$

+

only in  $T \in [125 : 151]$  MeV

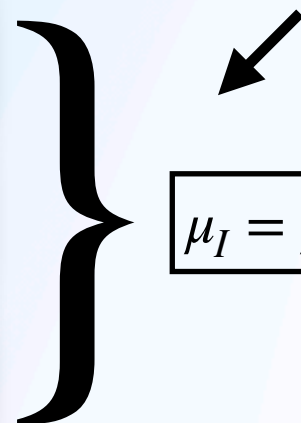
$$(\mu_I^c, T_I^c) := (132.40 \pm 25.91, 77.38 \pm 18.71) \text{ MeV}$$

This lies on the present state-of-the-art pion condensate line in the QCD isospin phase diagram.  $\iff$  Nomenclature with  $\mu_I^c = m_\pi$  (One of the achievements).

Son, Stephanov :  
*PRL*

Andersen et. al.  
(2024) : *PRD*

Andersen et. al.  
(2025) : *PRD*

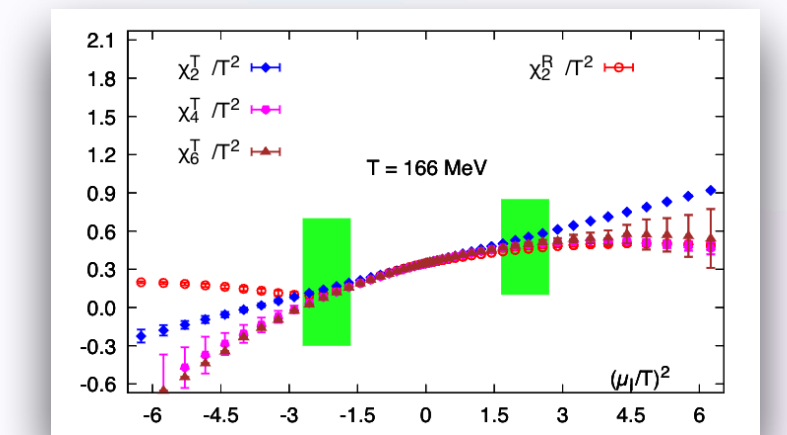
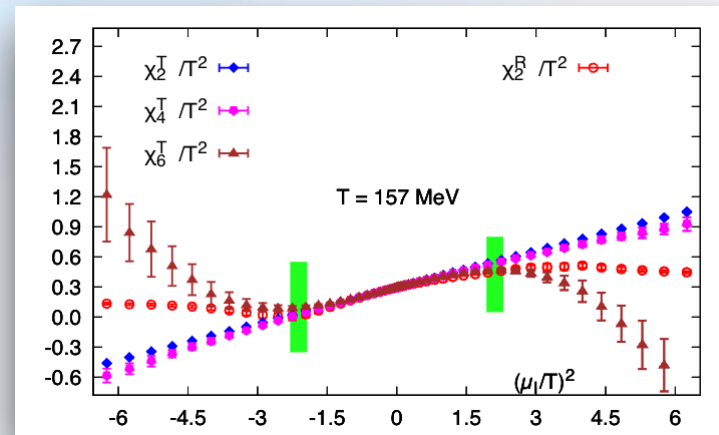
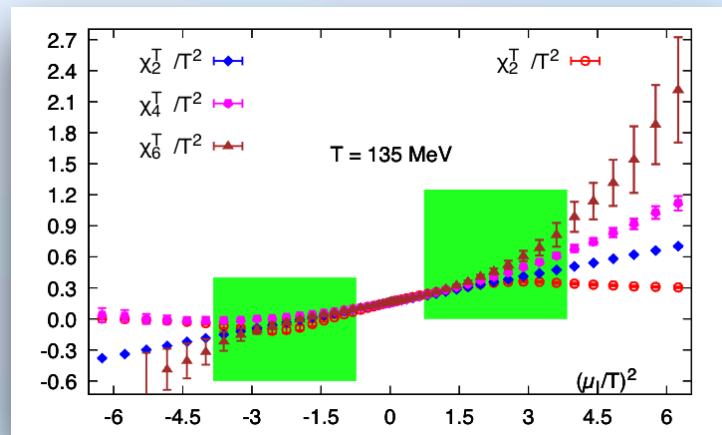
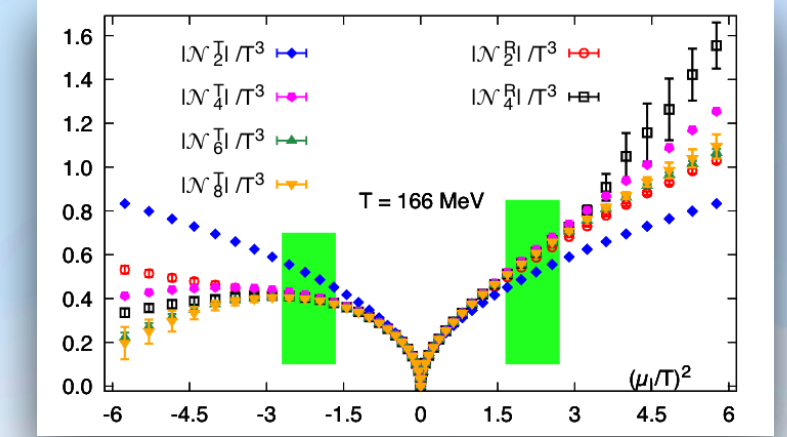
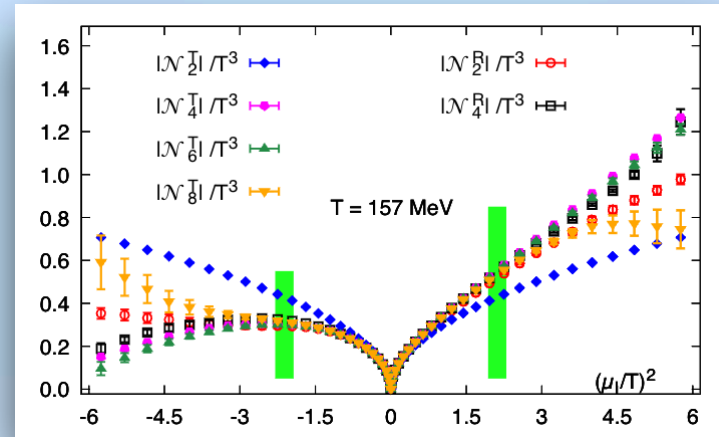
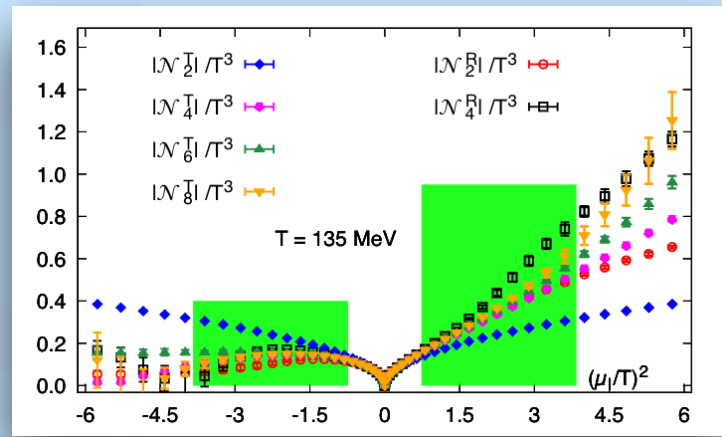
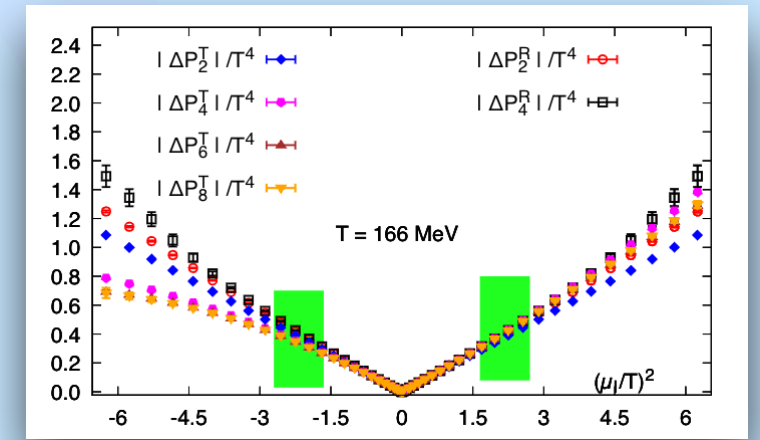
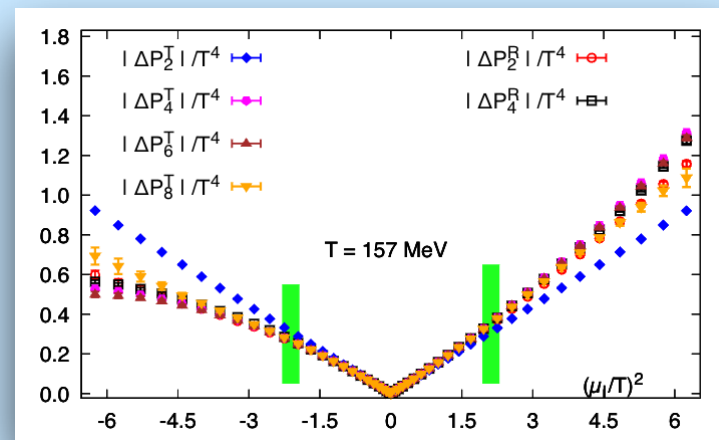
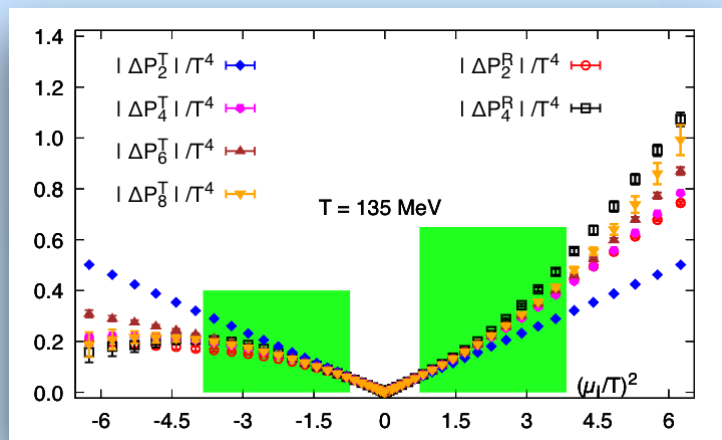


$$\mu_I = \mu_u - \mu_d$$

**within  
error bars**

Some observations related to this RoC ...

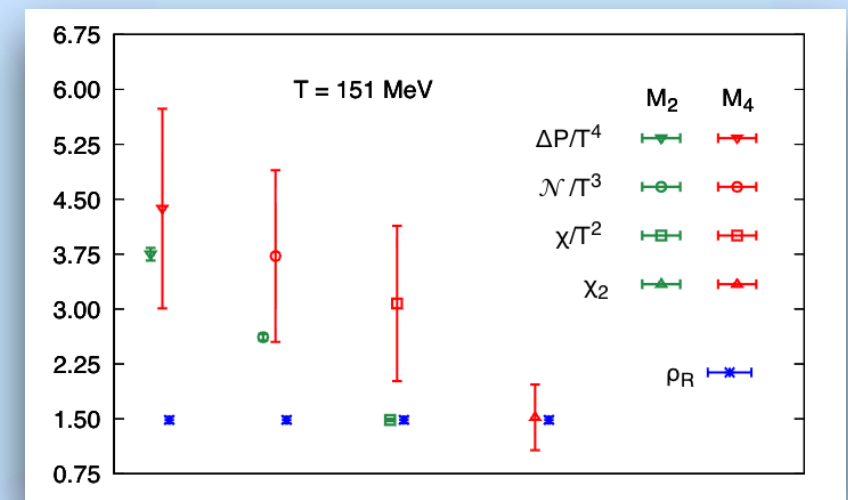
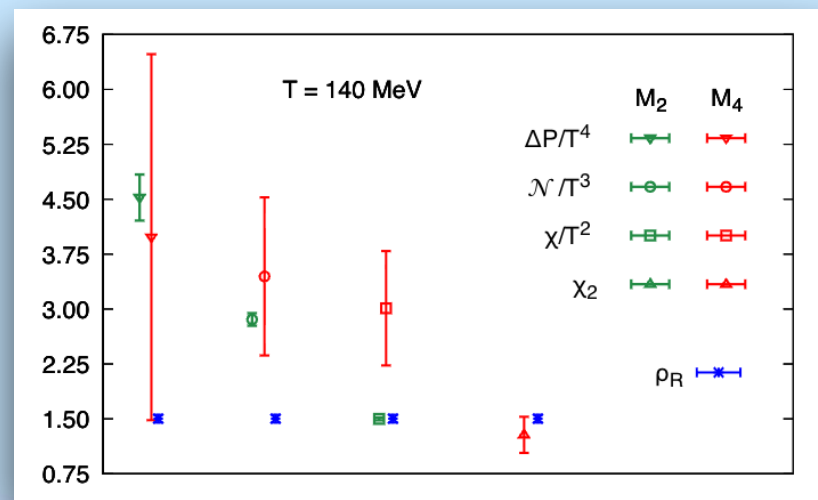
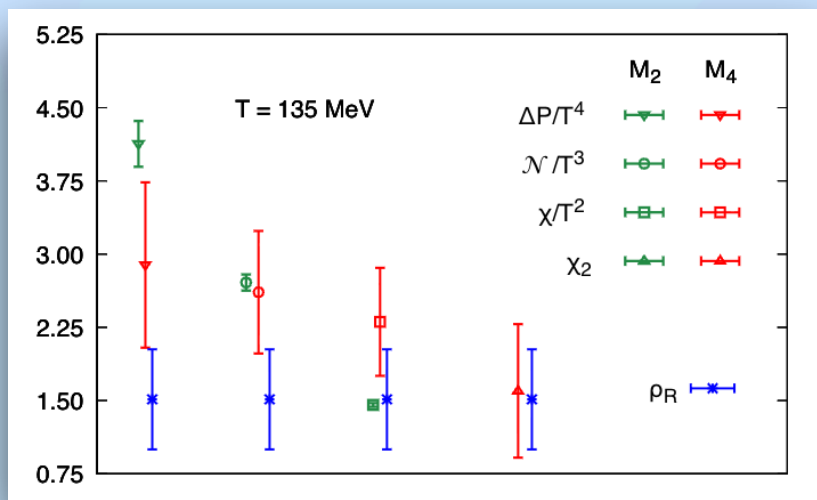




- Order-by-order deviations **beyond**  $\mu_I^\rho$
- Qualitatively, deviations  $\rightarrow \chi > \mathcal{N} > \Delta P$

for **different**  $T$

for **both**  $\text{Re}(\mu_I)$  and  $\text{Im}(\mu_I)$



$$\rho_R = \left| \mu_I^0 \right|$$

$$M_n = \left| \frac{c_{2n+2} c_{2n-2} - c_{2n}^2}{c_{2n+4} c_{2n} - c_{2n+2}^2} \right|^{1/2}$$

$$\frac{\Delta P}{T^4} = \sum_{n=1}^N c_n \left( \frac{\mu_I}{T} \right)^n$$

More order-by-order stable → Mercer-Roberts estimates  $M_n$

- $M_n$  estimates **approach**  $\rho_R = \mu_I^\rho / T$

For all the three  $T$

For higher order  $\mu_I$  derivatives of  $\Delta P$

# Overlap problem

Quantified  
by

Kurtosis  $\kappa$

$$\kappa = M_4 / \sigma^4$$

$$M_4 = \frac{1}{N} \sum_{i=1}^N (x_i - \bar{x})^4$$

Here,  $N = N_{conf}$

$$\sigma = \left( \frac{1}{N} \sum_{i=1}^N (x_i - \bar{x})^2 \right)^{1/2}$$

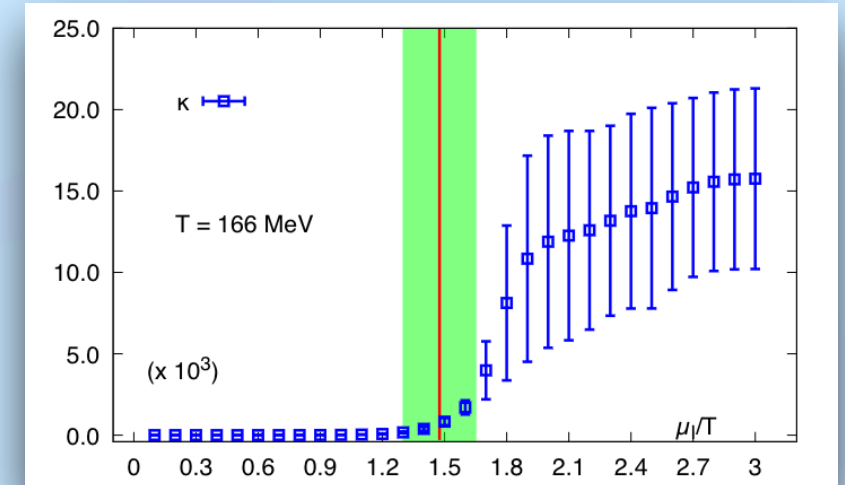
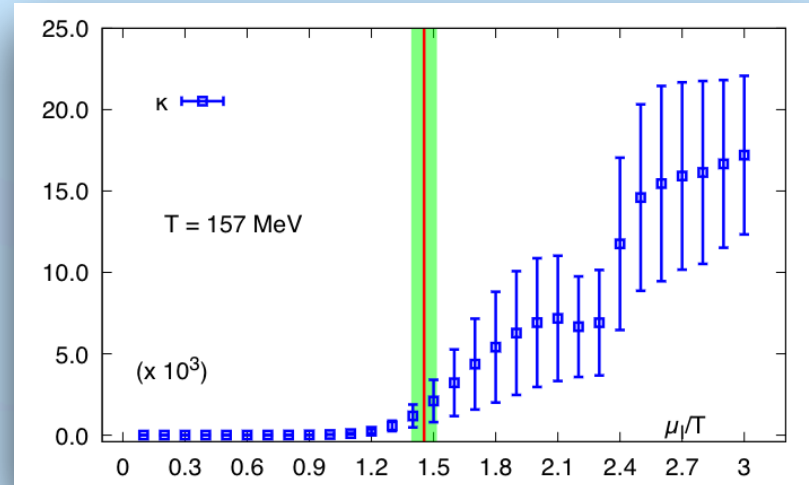
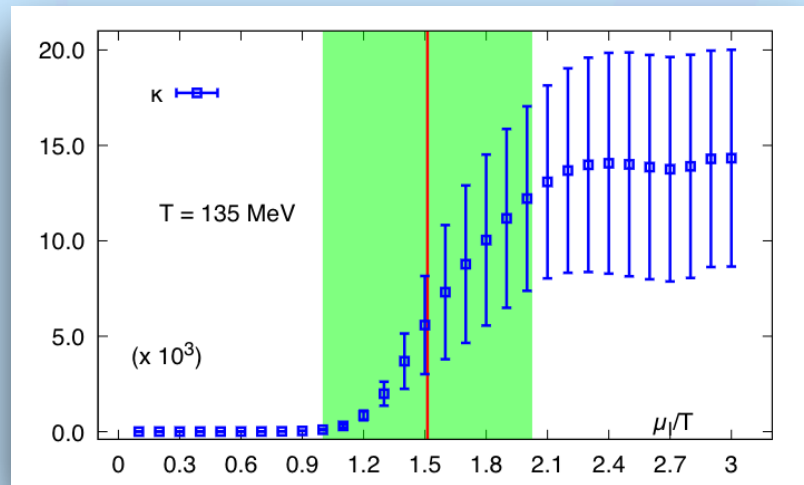
$$\bar{x} = \frac{1}{N} \sum_{i=1}^N x_i$$

Results???

describing distribution  $D$

$$D = \left\{ R_i(\mu_I) : i = 1, 2, \dots, N_{conf} \right\}, \quad R(\mu_I) = \det M(\mu_I) / \det M(0)$$

# Overlap problem



- **Controllable** until **Radius of convergence** (drastic after that)  $\rightarrow$  for all the three  $T$ .
- Indicates the efficacy of  $\mu_I = 0$  extrapolations to **determine finite  $\mu_I$  observables**.
- This truly **breaks down** and is **unreliable** beyond the **Radius of convergence**  $\mu_I^\rho$ .

# Summary

- Present estimates of Lee-Yang zeros  $\rightarrow$  map to a **critical point**  $\Rightarrow$  **exists on** the **present state-of-the-art  $\pi$ -cond. line** (**within error bars**).
- **Early indications of a phase transition** here, from **monotonic reduction of  $\text{Im}(\mu_I^0)$  with reducing  $T$ , while  $\text{Re}(\mu_I^0)$  approaching  $m_\pi$** .
- The root estimate results so far, **show stability** and well-defined convergence.
- **Radius of convergence (RoC)** is a **good indicator** of this  $\pi$ -condensate line, with **divergences** manifesting **beyond this RoC**, for different observables.
- Good agreement : **RoC**  $\iff$  **MR** estimates for higher order observables.
- **Overlap problem** becoming noticeably **severe beyond this RoC**.

With lot of future things to do...

# Future works and Outlook

- Simulations at **higher lattice volumes** ( larger  $N_\sigma$  for same  $N_\tau$  )

Finite- $V$  analysis

if these zeros truly **become real**  $\rightarrow$  confirm a **phase transition** , and

**NOT a crossover**

- Also, how do the  $\text{Im}(\mu_0)$  behave or, **scale** with volumes ?

$1/V$   
scaling

**First order**

or

Critical-exponent  
dependent scaling

**Second order**

- Estimate of  $T_I^c$  in  
thermodynamic limit

- Universality class  
determination, therefore

**From a LY zero perspective**





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EUROPEAN CENTRE  
FOR THEORETICAL STUDIES  
IN NUCLEAR PHYSICS AND RELATED AREAS

**ECT\* workshop**

**“Analytic structure of QCD and Yang-Lee edge singularity”**

**08-12** Sep. 2025

ECT\* Villa Tambosi, Villazzano

**THANK YOU FOR YOUR ATTENTION**

**Backup slides**

## The nomenclature of isospin

$$B = \frac{1}{3} (N_u + N_d + N_s) , \quad S = -N_s , \quad I = \frac{1}{2} (N_u - N_d)$$

$$\mu_B = \frac{3}{2} (\mu_u + \mu_d) \qquad \mu_I = (\mu_u - \mu_d)$$

$$B_u = B_d = B_s = 1/3 \quad (\text{Baryons are 3-quark systems})$$

$$|I_p| = |I_n| = 1/2, \text{ since } 2I + 1 = 2 \quad (\text{proton and neutron})$$

$$\text{Thus, } I_u = -I_d \text{ and } |I_u| = |I_d| = 1/2$$