

Quantum entanglement from local operators

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ECT* Trento, Hamiltonian Lattice Gauge Theories: Status, Novel Developments and Applications

Based on a [WIP](#) with [Marco Panero](#), [Paolo Stornati](#) and [Luca Tagliacozzo](#)



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DI **TORINO**



Istituto Nazionale di Fisica Nucleare

Entanglement in lattice field theories I

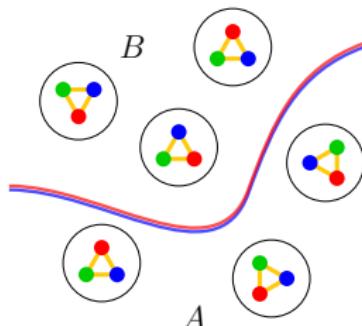
Entanglement in many-body systems —→ applications in different areas of physics:

- Quantum phases of matter and quantum phase transitions [Vidal et al.; PRL 90 (2003)]
- High-energy physics
 - Confinement in gauge theories [Klebanov et al.; Nucl. Phys. B 796 (2008)]
- AdS/CFT and quantum gravity [Ryu, Takayanagi; PRL 96 (2006)]
- Complexity of quantum simulations, thermalization and its violation, scattering processes ...



Dirac Medallists 2024

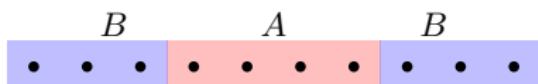
Horacio Casini, Marina Huerta, Shinsei Ryu, Tadashi Takayanagi



For “pioneering contributions to the understanding of quantum entropy in gravity and quantum field theory”

Bipartite entanglement

A hard quantity to handle



$$\rho_A = \sum_{P_{i_1}, \dots, P_{i_N}} c_{i_1 \dots i_N} P_{i_1} \otimes \dots \otimes P_{i_N}$$
$$P \in \{1, X, Y, Z\}$$

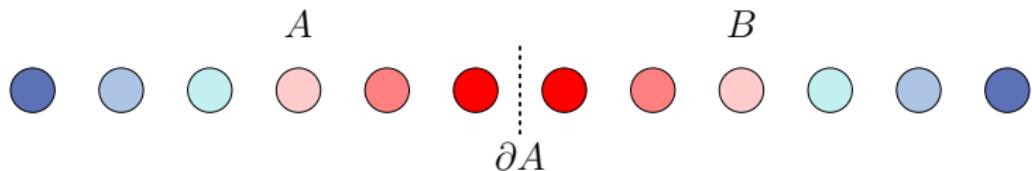
$$\rho_A = \mathrm{Tr}_B \rho$$

$$S = -\mathrm{Tr}_A(\rho_A \log \rho_A)$$

$$S_n = \frac{1}{1-n} \log \mathrm{Tr} \rho_A^n \qquad S_2 \longleftrightarrow \mathrm{Tr} \rho_A^2$$

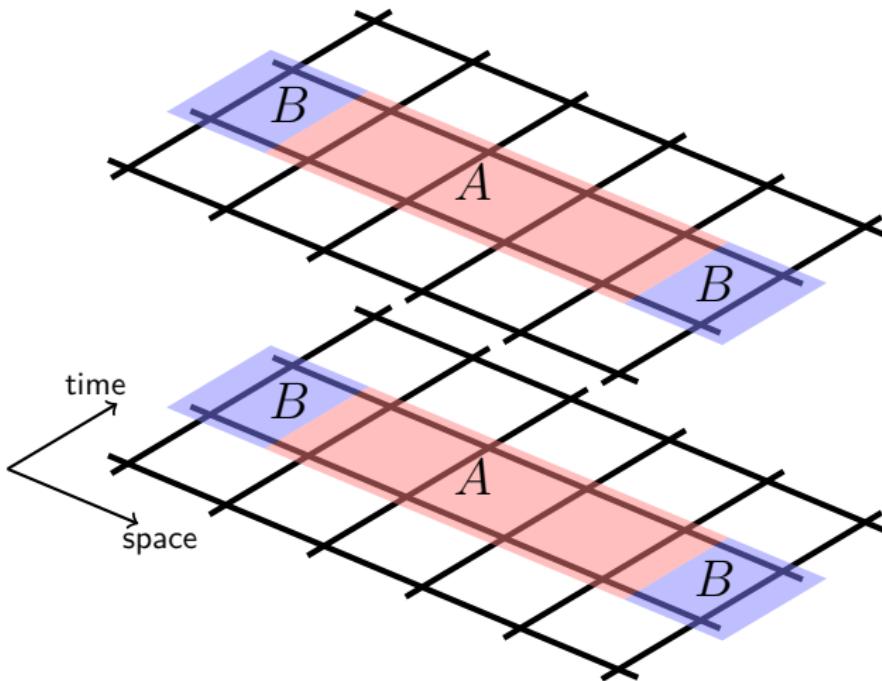
- Computing $c_{i_1 \dots i_N}$ is **exponentially hard**
- Other methods: SWAP operators, random unitaries + classical shadows
...
- In general: progressively **harder** for larger systems

Area law: ground states of local Hamiltonian have few entanglement



$$\mathrm{Tr} \rho_A^2 \simeq \langle \mathcal{O}(\partial A) \rangle ?$$

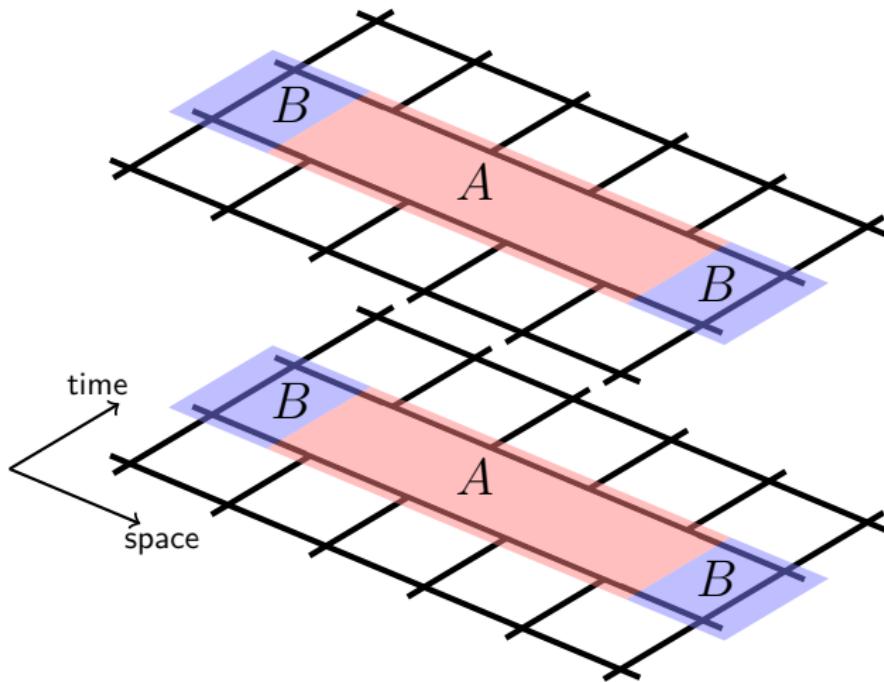
Intuition 2/3: replica trick I



$Z^2 \rightarrow$ Decoupled replicas

[Calabrese, Cardy; J.Stat.Mech. 0406 (2004)]

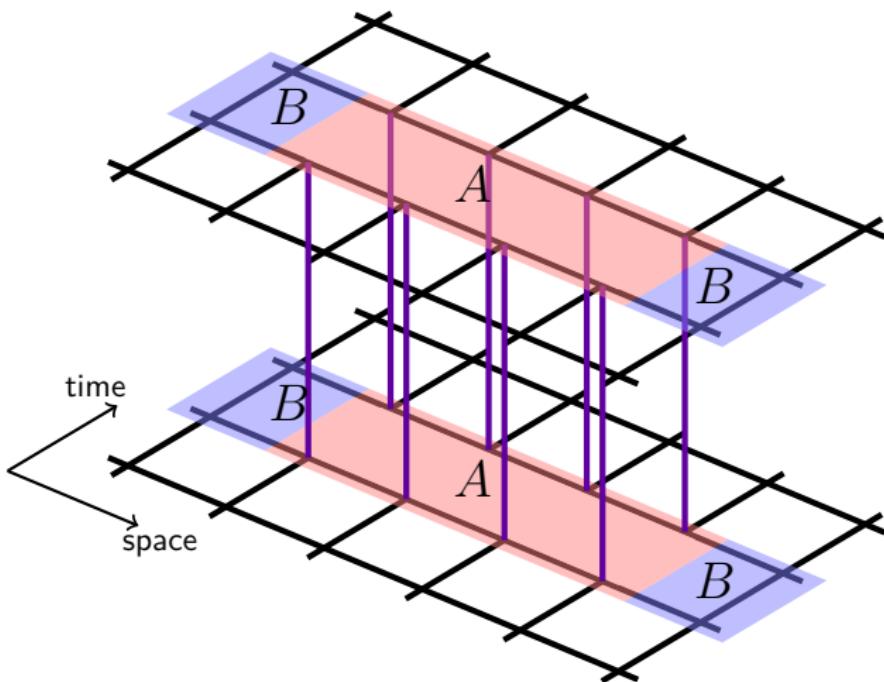
Intuition 2/3: replica trick I



$Z^2 \rightarrow$ Decoupled replicas

[Calabrese, Cardy; J.Stat.Mech. 0406 (2004)]

Intuition 2/3: replica trick I



$Z^2 \rightarrow$ Decoupled replicas

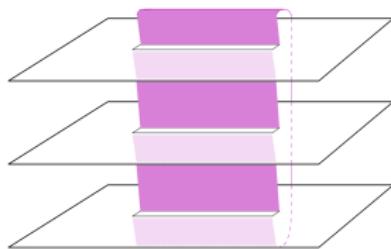
$Z_2 \rightarrow$ Coupled replicas

$$S_2 = -\log \frac{Z_2}{Z^2}$$

[Calabrese, Cardy; J.Stat.Mech. 0406 (2004)]

Replicated manifolds

- Widely used in **CFT calculations**
[Calabrese, Cardy; J.Stat.Mech. 0406 (2004)]

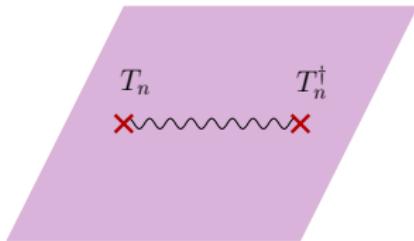


From [Cardy et al.; J.Stat.Phys 130 (2008)]

Twist fields

- The replica trick implies that Rényi entropies can be computed as **correlators of local operators**

$$\text{Tr } \rho_A^n = \langle T_n T_n^\dagger \rangle$$



- Exploited in **Lagrangian simulations**

[Buividovich, Polikarpov; Nucl.Phys.B 802 (2008)]

[Itou et al.; PTEP (2016)]

[Rabenstein et al.; PRD 100 (2019)]

[AB, Panero; JHEP 06 (2024)]

- A lattice characterization is missing...

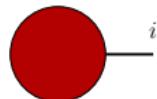
Goals

Lattice characteriza-
tion of twist operators

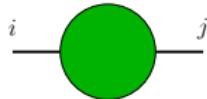
Systematic way to mea-
sure entanglement
with local operators

Quick dictionary: Tensor Networks I

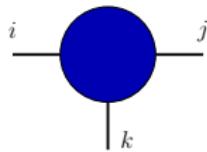
Tensor diagrams



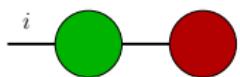
v_i vector



M_{ij} matrix



A_{ijk} rank-3 tensor



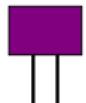
$M_{ij} v_j$ matrix-vector multiplication

States as tensors

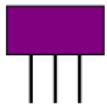
$$|\psi\rangle = \sum_i c_i |i\rangle$$



$$|\psi\rangle = \sum_{ij} c_{ij} |ij\rangle$$



$$|\psi\rangle = \sum_{ijk} c_{ijk} |ijk\rangle$$



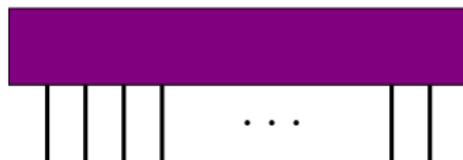
$$|\psi\rangle = \sum_{i_1 i_2 \dots i_N} c_{i_1 i_2 \dots i_N} |i_1 i_2 \dots i_N\rangle$$



Exponential size

- Full representation of a quantum state: **exponential** in system size

$$|\psi\rangle = \sum_{i_1 i_2 \dots i_N} c_{i_1 i_2 \dots i_N} |i_1 i_2 \dots i_N\rangle$$

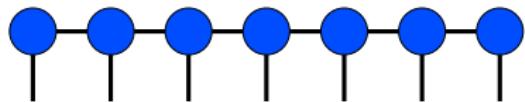


$$\sim \exp(N)$$

MPS

- **Matrix Product States:** **compressed** representation for a *1d* system

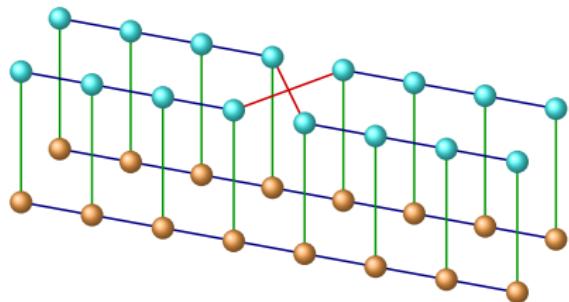
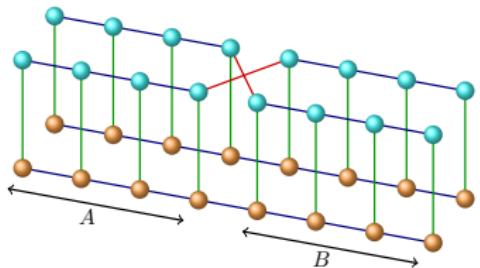
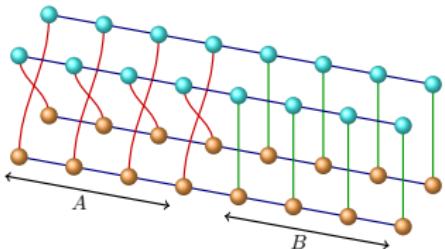
$$c_{i_1 i_2 \dots i_N} \rightarrow [A^1]_{i_1}^{\alpha_1} [A^2]_{i_2}^{\alpha_1 \alpha_2} \dots [A^N]_{i_N}^{\alpha_N}$$



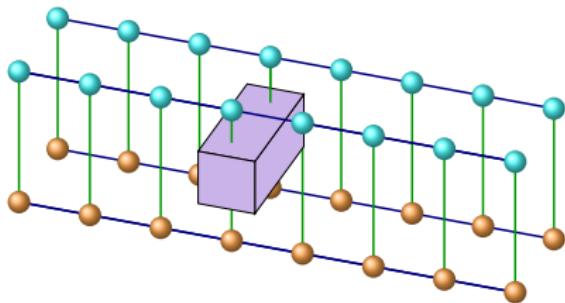
$$\sim \text{poly}(N)$$

Intuition 3/3: MPS

$$\text{Tr } \rho_A^2 =$$

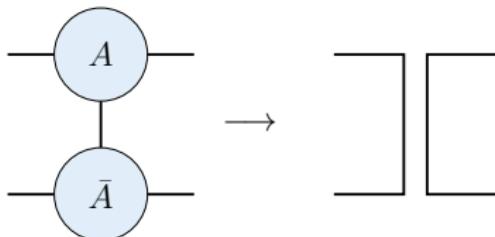


?

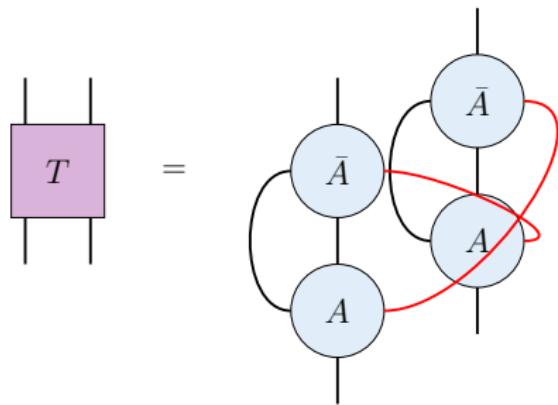


Inverting the MPS

$$\chi \xrightarrow{\quad} \begin{matrix} A \\ d \end{matrix} = \chi \xrightarrow{\quad} \begin{matrix} U \\ S \\ V^\dagger \\ d \end{matrix}$$



$$\chi \xrightarrow{\quad} \begin{matrix} \bar{A} \\ d \end{matrix} = \chi \xrightarrow{\quad} \begin{matrix} V \\ S^{-1} \\ U^\dagger \\ d \end{matrix}$$



Invert the MPS

- If one inverts tensor by tensor we get

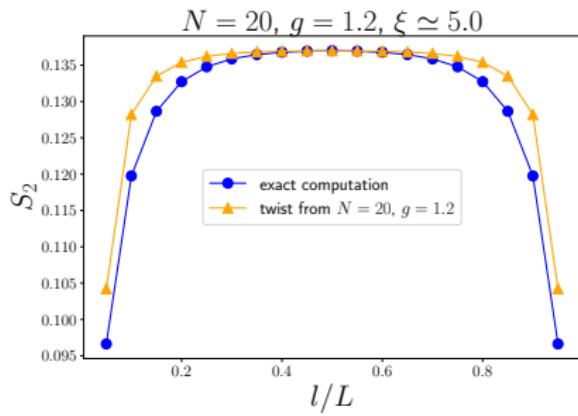
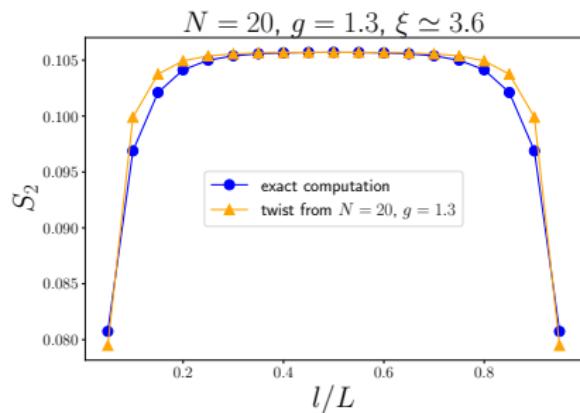
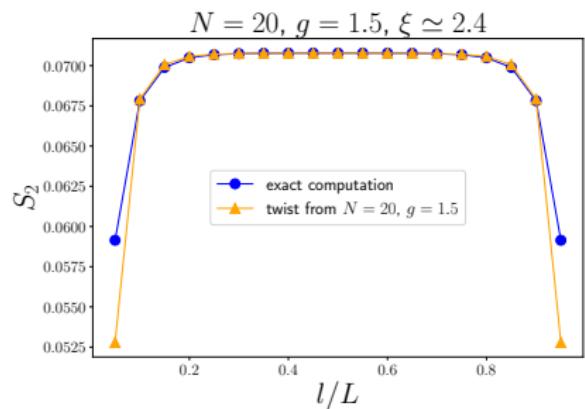
$$\text{Tr } \rho_A^2 = \langle T \rangle$$

- **However:** in general T depends on the site
- Unclear how to derive it, but likely it implies the **knowledge of** S_2

What can we learn?

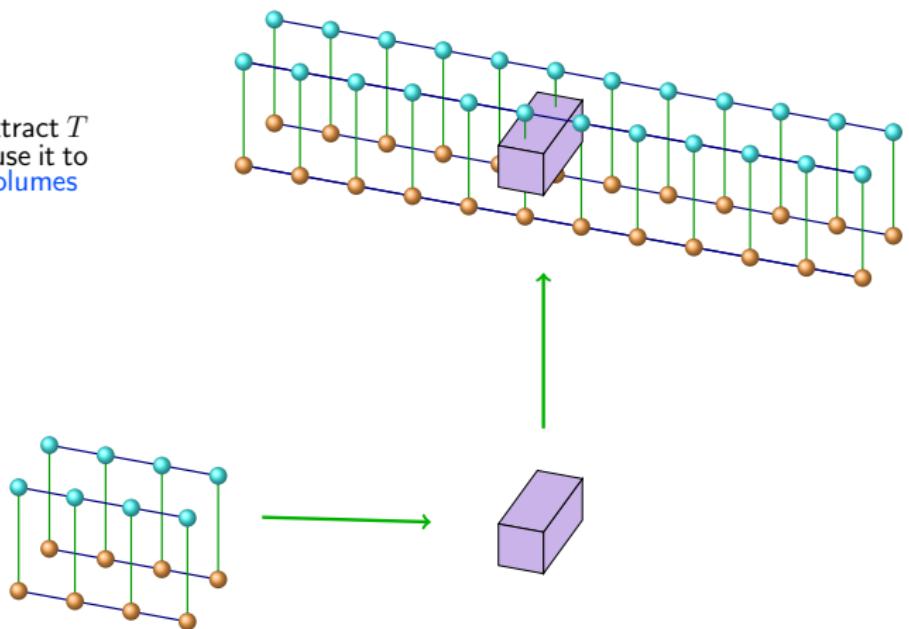
- An **exact** scheme might be too costly
- Can we design an **approximate** yet **systematically improvable** scheme?
- A first idea: use the **same operator** to compute S_2 in the whole system
- Benchmark: transverse field Ising model with OBC

Strategies to compute S_2 I

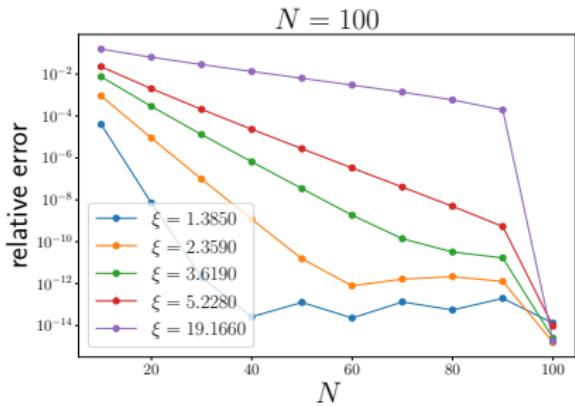
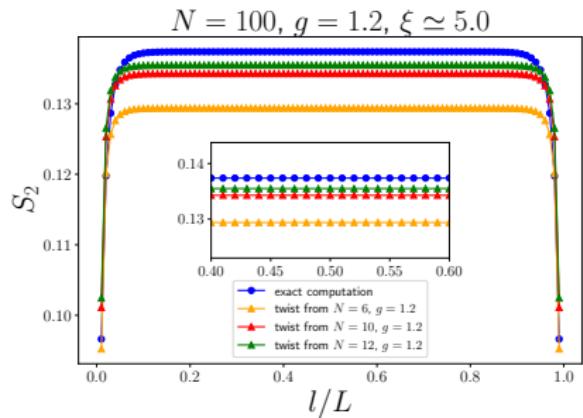
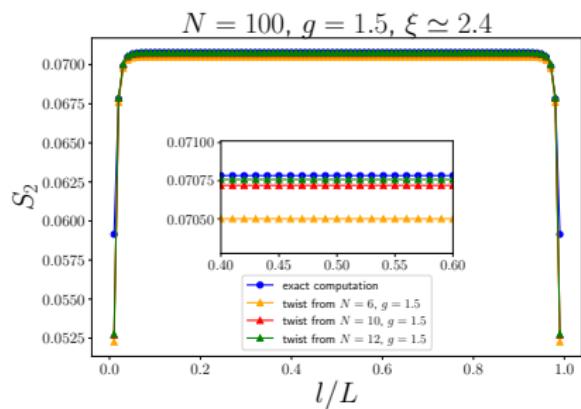


Strategies to compute S_2 II

Another strategy: extract T for **small volumes** and use it to compute S_2 in **large volumes**



Strategies to compute S_2 III



Matrix inversion

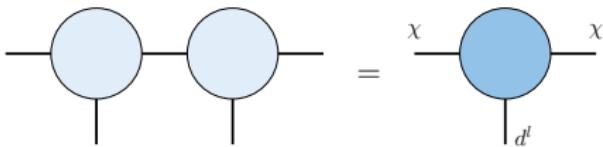
- Larger ξ requires larger χ
- The "inversion" progressively **fails** to give access to the virtual legs



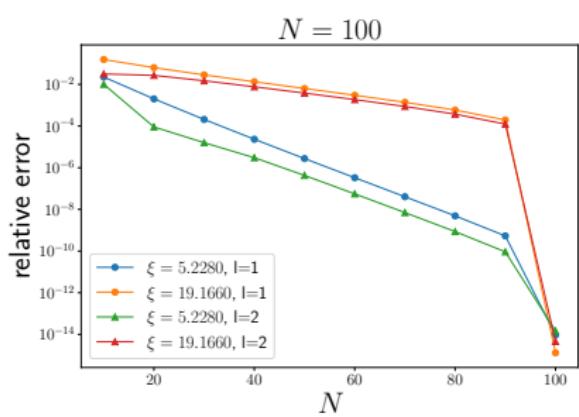
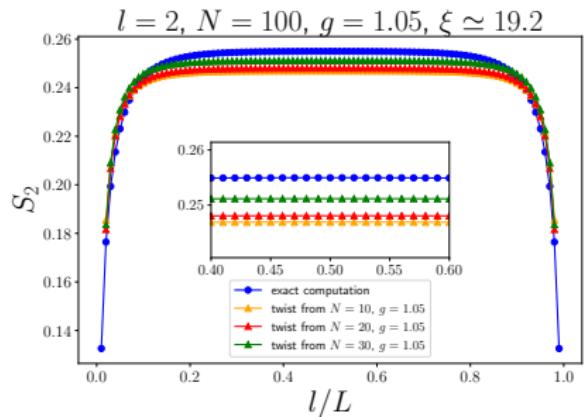
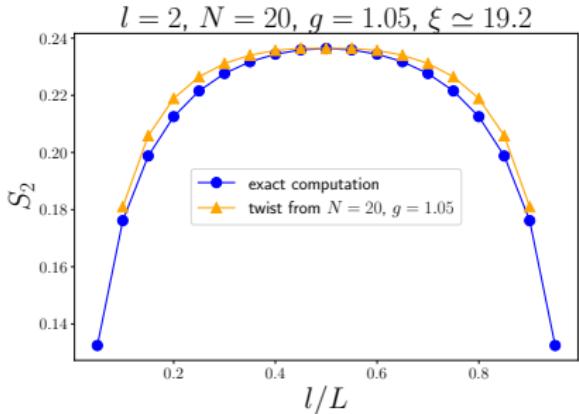
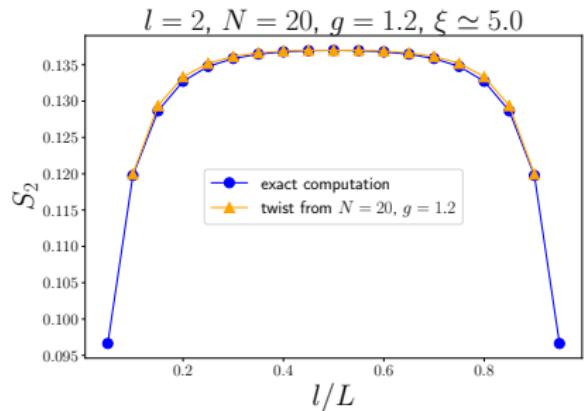
$$\chi^2 > d$$

Larger tensors

- Contract together l tensors
- **Trade-off:** larger matrices provide a better inversion but are more difficult to measure



Larger twists II



Conclusions and outlook

Conclusions

- Local operators inspired to twist fields can provide a measure of entanglement
- Approximate and systematically improvable scheme at least in gapped systems
- Larger twist provide a progressively better approximation

Outlook

- Tomography of such operators?
- Precise lattice definition of twist fields from CFT?
- How to deal with large ξ ?
- MPO version of large twists?