

Quantum Simulation of Fundamental Physics

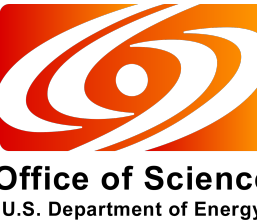


Progress and Thoughts
ECT*, Trento, September 1, 2025

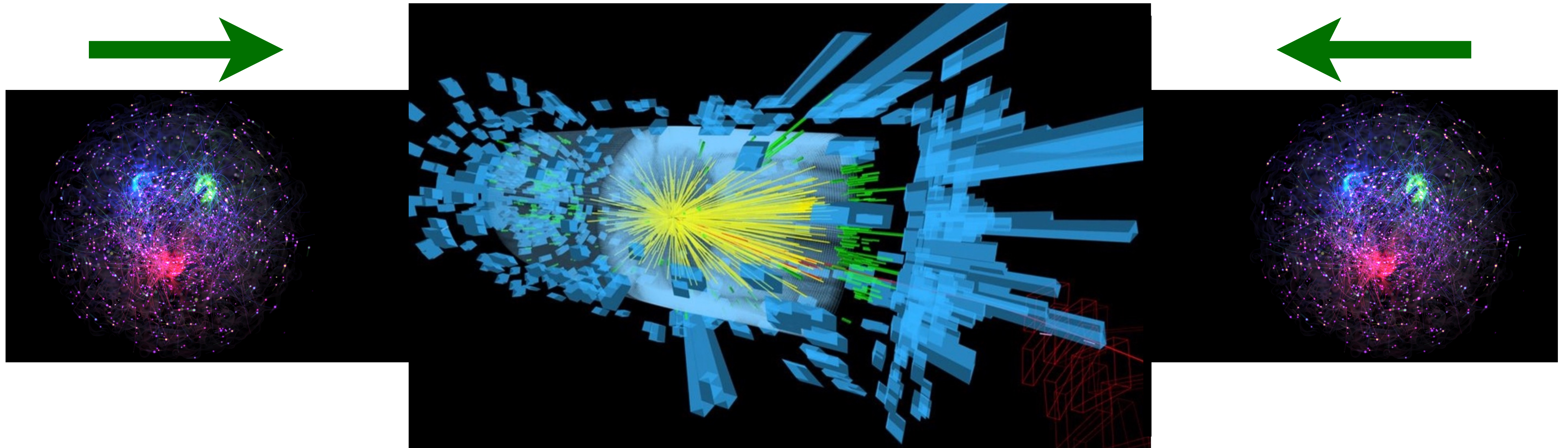
Martin Savage
InQubator for Quantum Simulation (IQuS),
University of Washington



<https://iqus.uw.edu/>



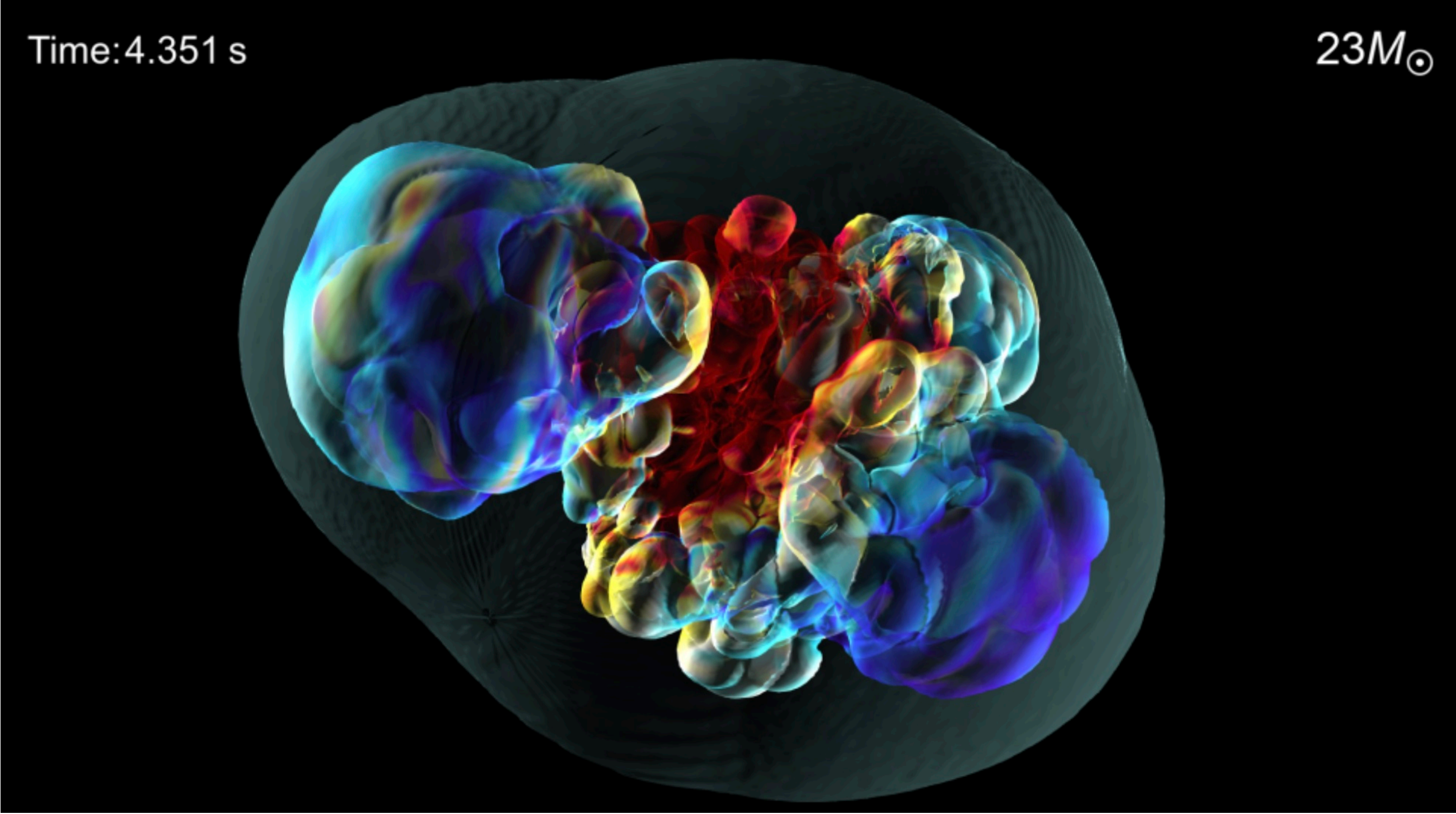
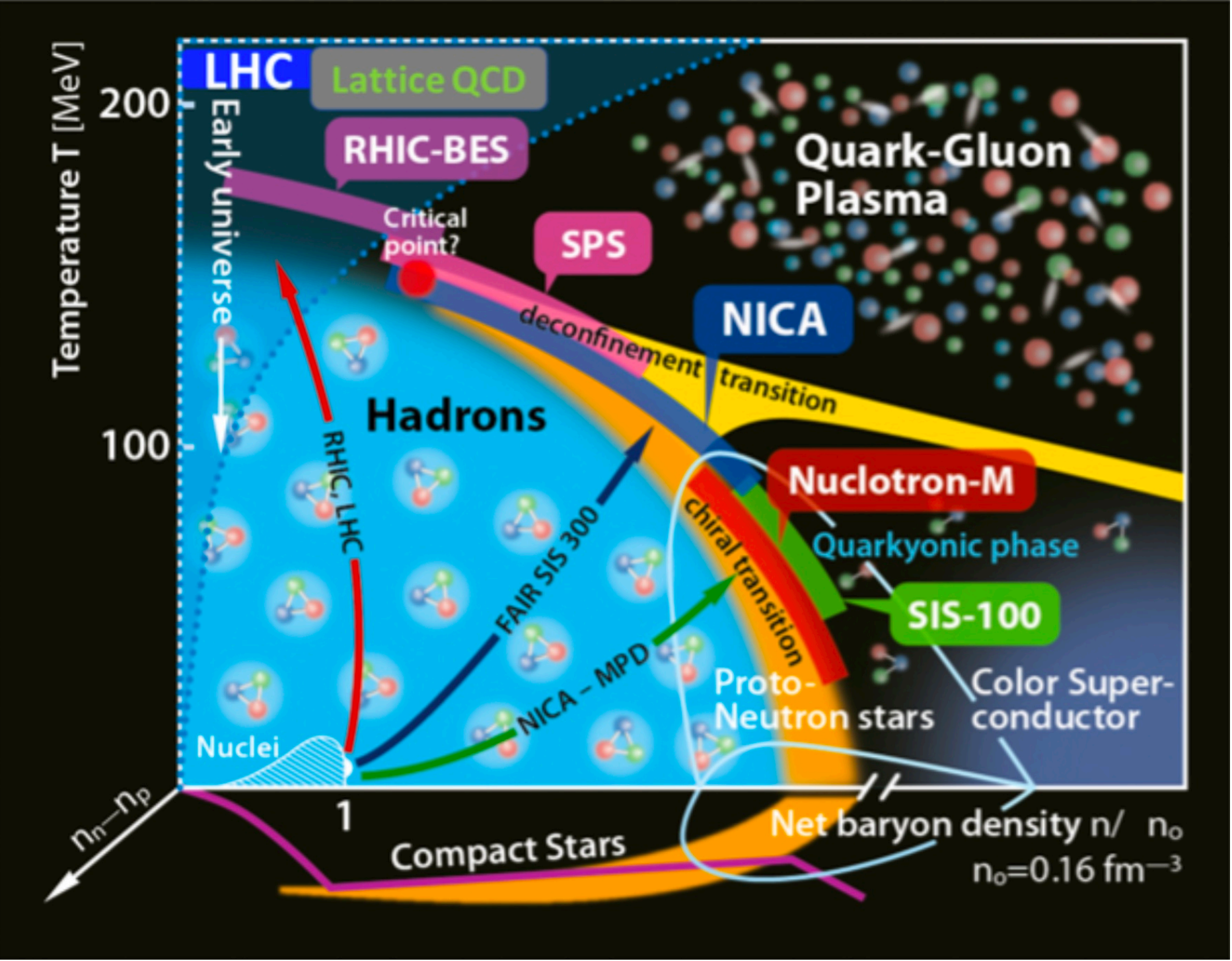
High-Energy Collisions



There is a lot of physics that takes place in collisions of matter

- Production of new particles, e.g., Higgs
- Collisions of nuclei with satellites, energy loss with materials - quantification
- Properties and thermalization of quark-gluon plasma/liquid
- Fragmentation, hadronization, confinement - QCD is now the background
- Transport in strongly interacting, correlated non-equilibrium matter - open quantum system(s)
- Neutrino interactions with matter and properties

Hot and Dense Non-Equilibrium Matter



Possible early universe signals in proton collisions at the Large Hadron Collider

Raghunath Sahoo^{1,2,*} and Tapan Kumar Nayak^{2,3}
¹Department of Physics, Indian Institute of Technology Indore, Simrol, Indore 453552, India
²CERN, CH 1211, Geneva 23, Switzerland and
³School of Physical Sciences, National Institute of Science Education and Research, HBNI, Jatni-752050, India
(Dated: January 4, 2022)

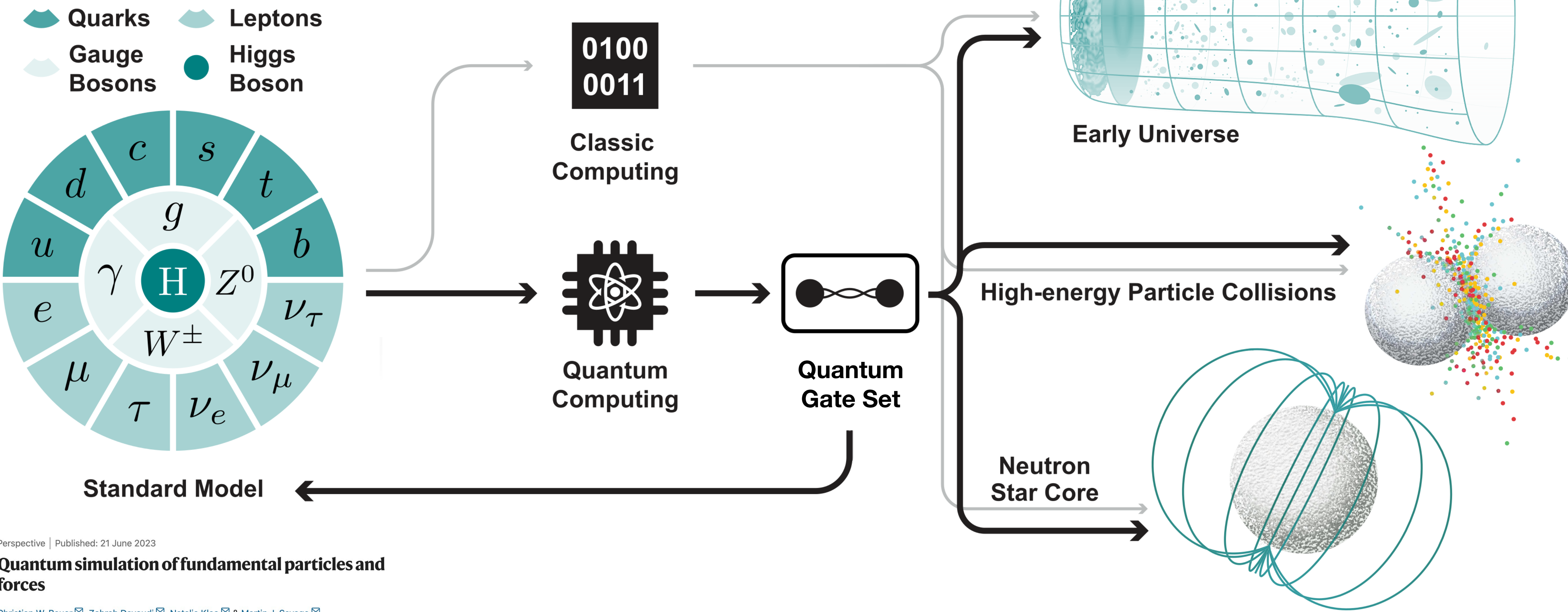
Next-Generation 3D Core-Collapse
Supernova Simulations

PI Adam Burrows, Princeton University
co-PI David Vartanyan, University of California Berkeley
Matt Coleman, Princeton University
Chris White, Princeton University

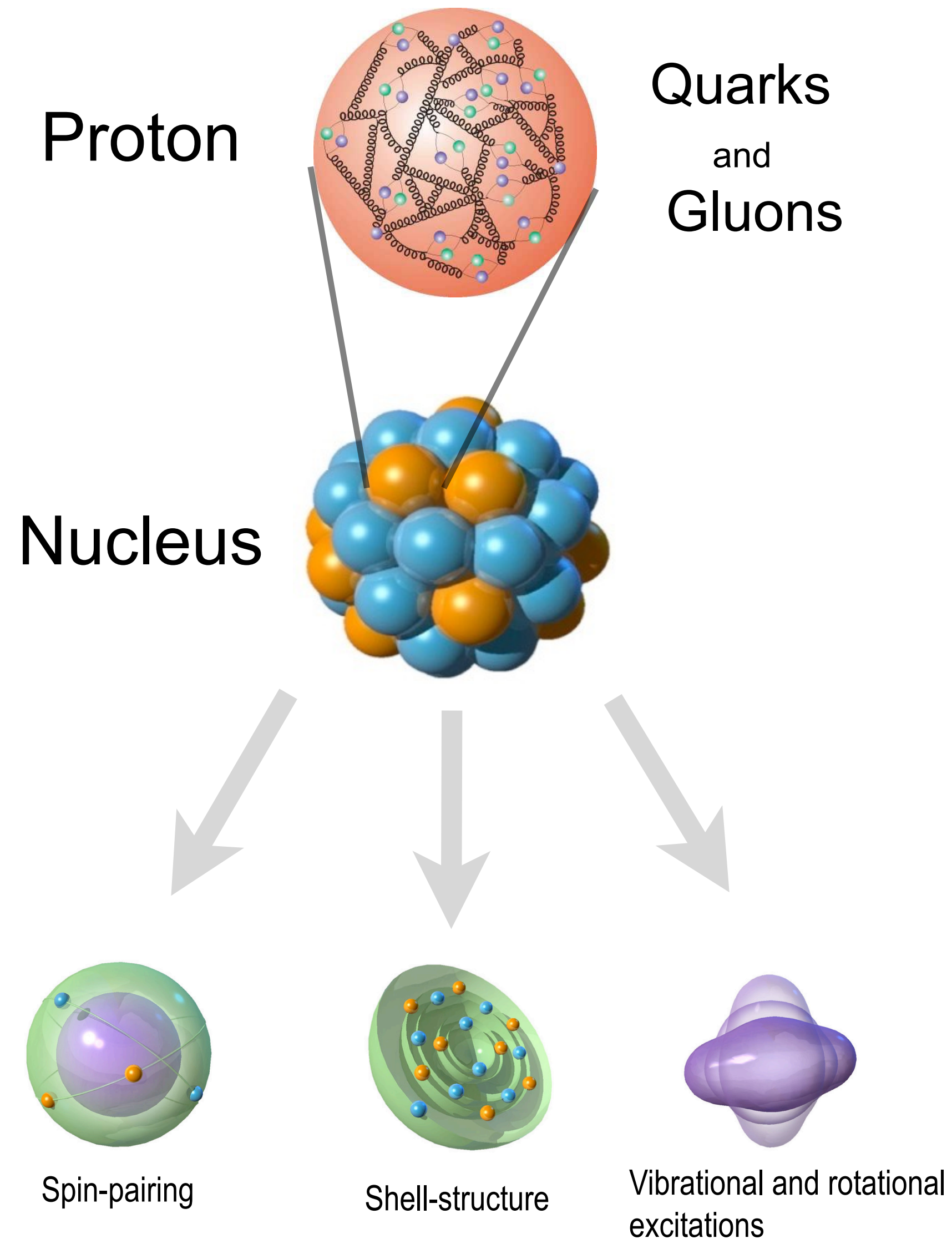
Particles & Interactions

Simulation

Phases & Dynamics of Matter



Ordinary Matter



$$\Lambda_{\text{QCD}}$$

$$\frac{m_u}{\Lambda_{\text{QCD}}} \quad \frac{m_d}{\Lambda_{\text{QCD}}} \quad \frac{m_s}{\Lambda_{\text{QCD}}}$$

$$\alpha_e$$

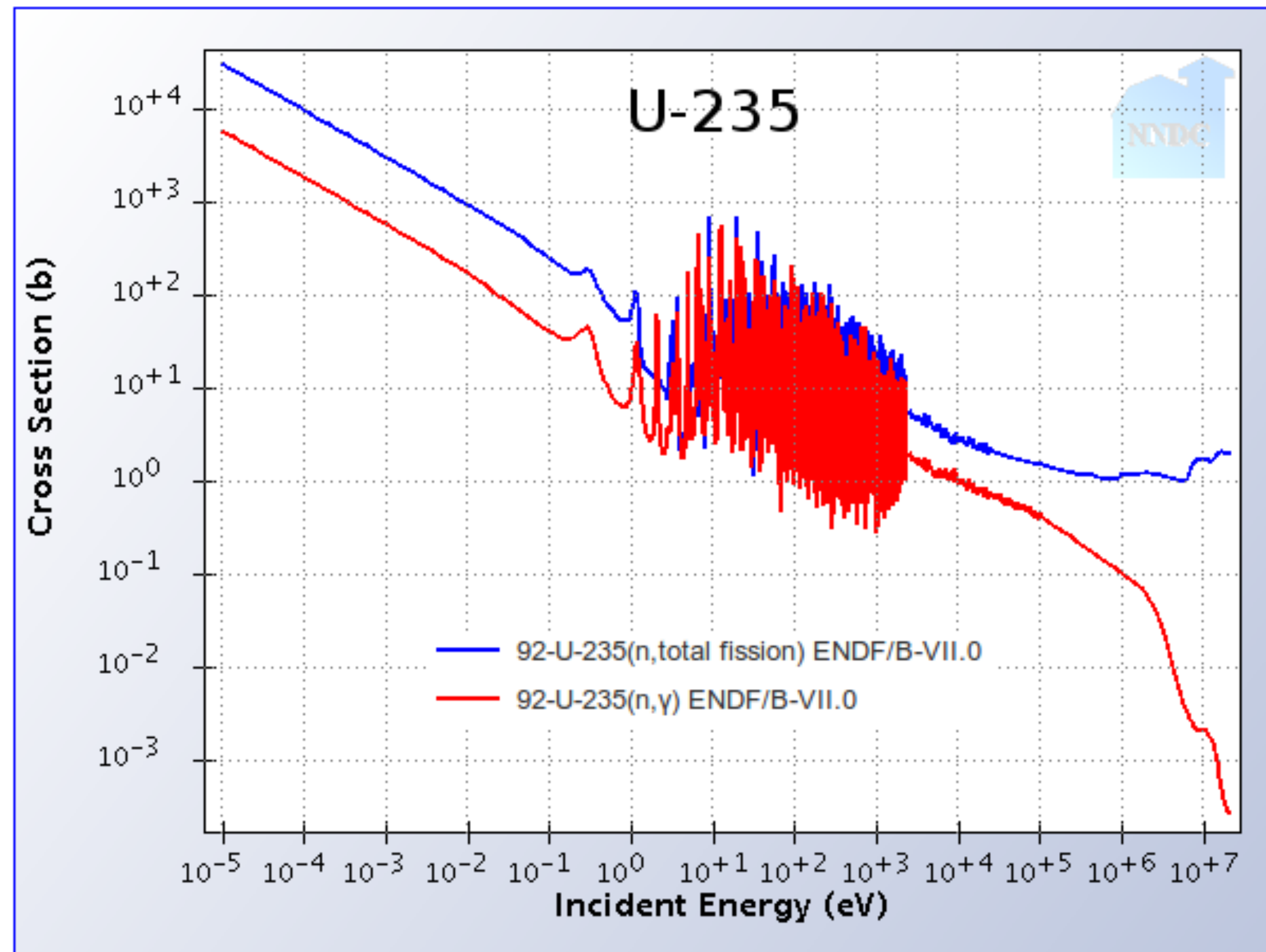
Small number of input parameters responsible for all of strongly interacting matter

Dimensionless plus a scale

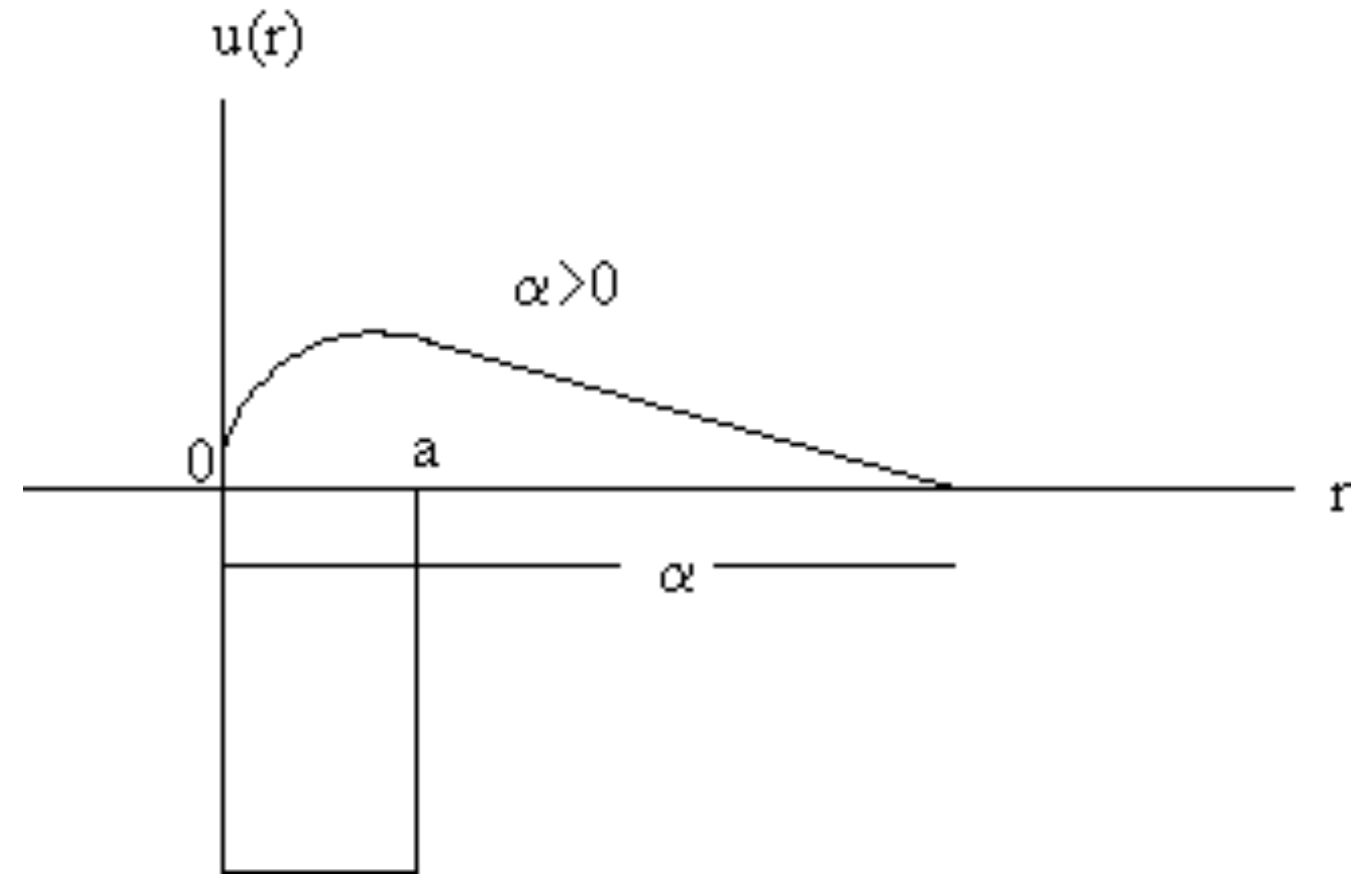
LQCD - dimensionless

Hierarchies

Neutron-Induced Fission



NN wavefunction at $E=0$



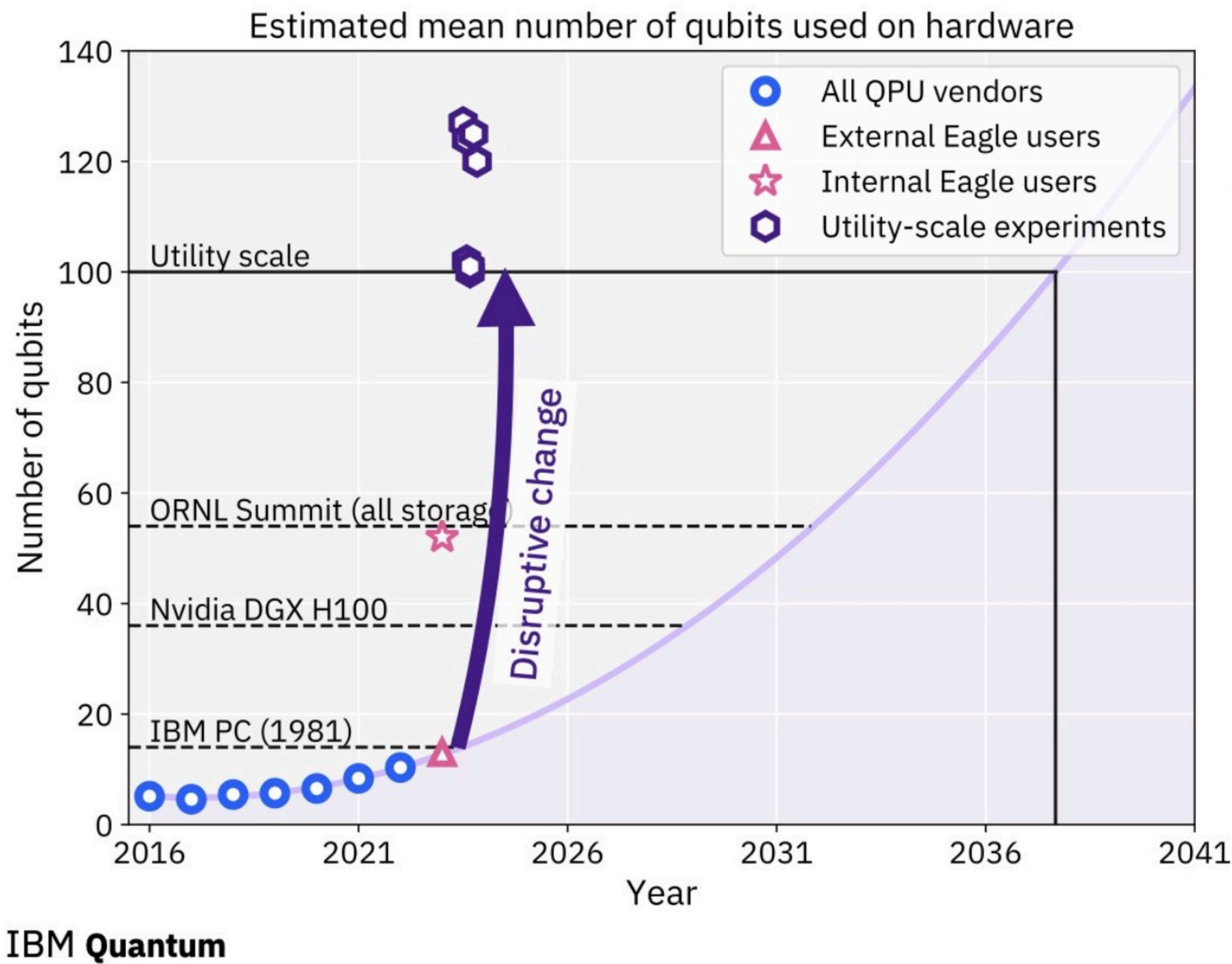
Mass $^{235}\text{U} \sim 220\,900\text{ MeV}$



IBM Quantum Summit - NYC December 2023

Utility-scale experiments

With quantum systems composed of 100+ qubits, researchers are beginning to explore algorithms and applications at scales beyond brute-force classical computation [using IBM Quantum systems](#).



Evidence for the utility of quantum computing before fault tolerance

[127 qubits / 2880 CX gates](#)

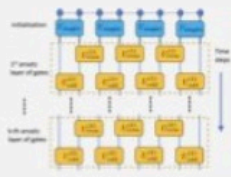
Nature, 618, 500 (2023)



Simulating large-size quantum spin chains on cloud-based superconducting quantum computers

[102 qubits / 3186 CX gates](#)

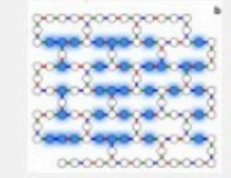
arXiv:2207.09994



Uncovering Local Integrability in Quantum Many-Body Dynamics

[124 qubits / 2641 CX gates](#)

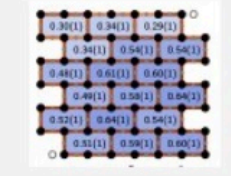
arXiv:2307.07552



Realizing the Nishimori transition across the error threshold for constant-depth quantum circuits

[125 qubits / 429 gates + meas.](#)

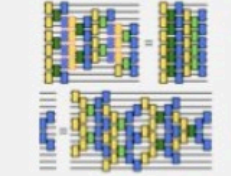
arXiv:2309.02863



Scalable Circuits for Preparing Ground States on Digital Quantum Computers: The Schwinger Model Vacuum on 100 Qubits

[100 qubits / 788 CX gates](#)

arXiv:2308.04481



Efficient Long-Range Entanglement using Dynamic Circuits

[101 qubits / 504 gates + meas.](#)

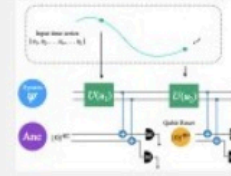
arXiv:2308.13065



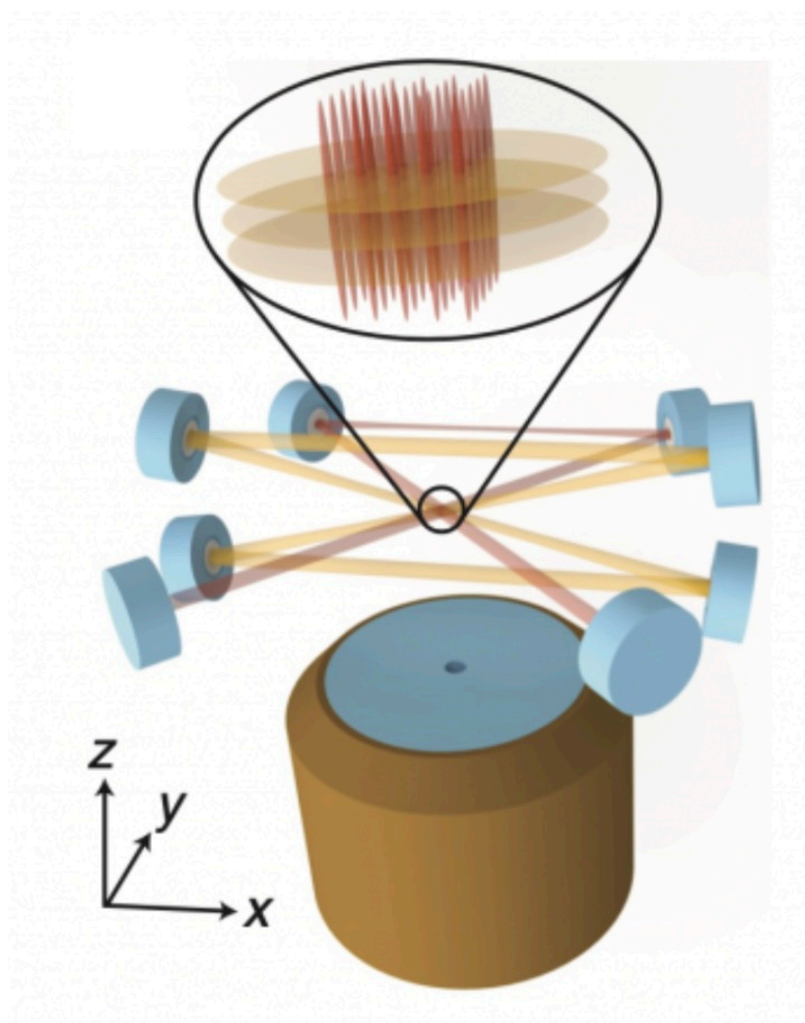
Quantum reservoir computing with repeated measurements on superconducting devices

[120 qubits / 49470 gates + meas.](#)

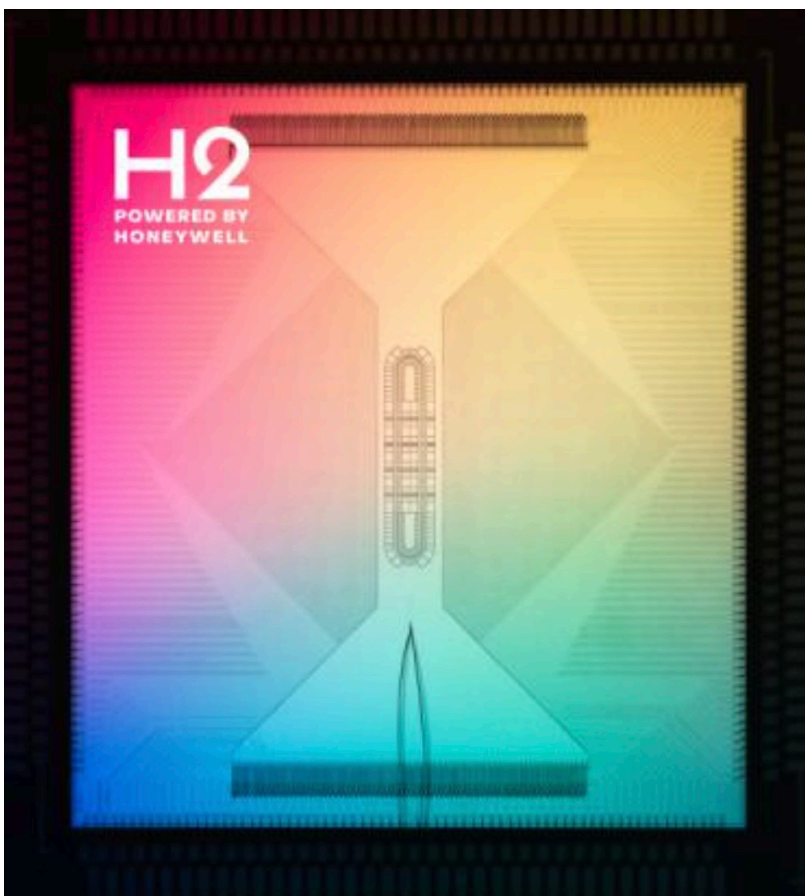
arXiv:2310.06706



Select Recent Advances in Quantum Computing



Cold-Atom arrays with Optical Tweezers



4 Logical Qubits
32-qubit H2-1 trapped ions
(Quantinuum-Microsoft)

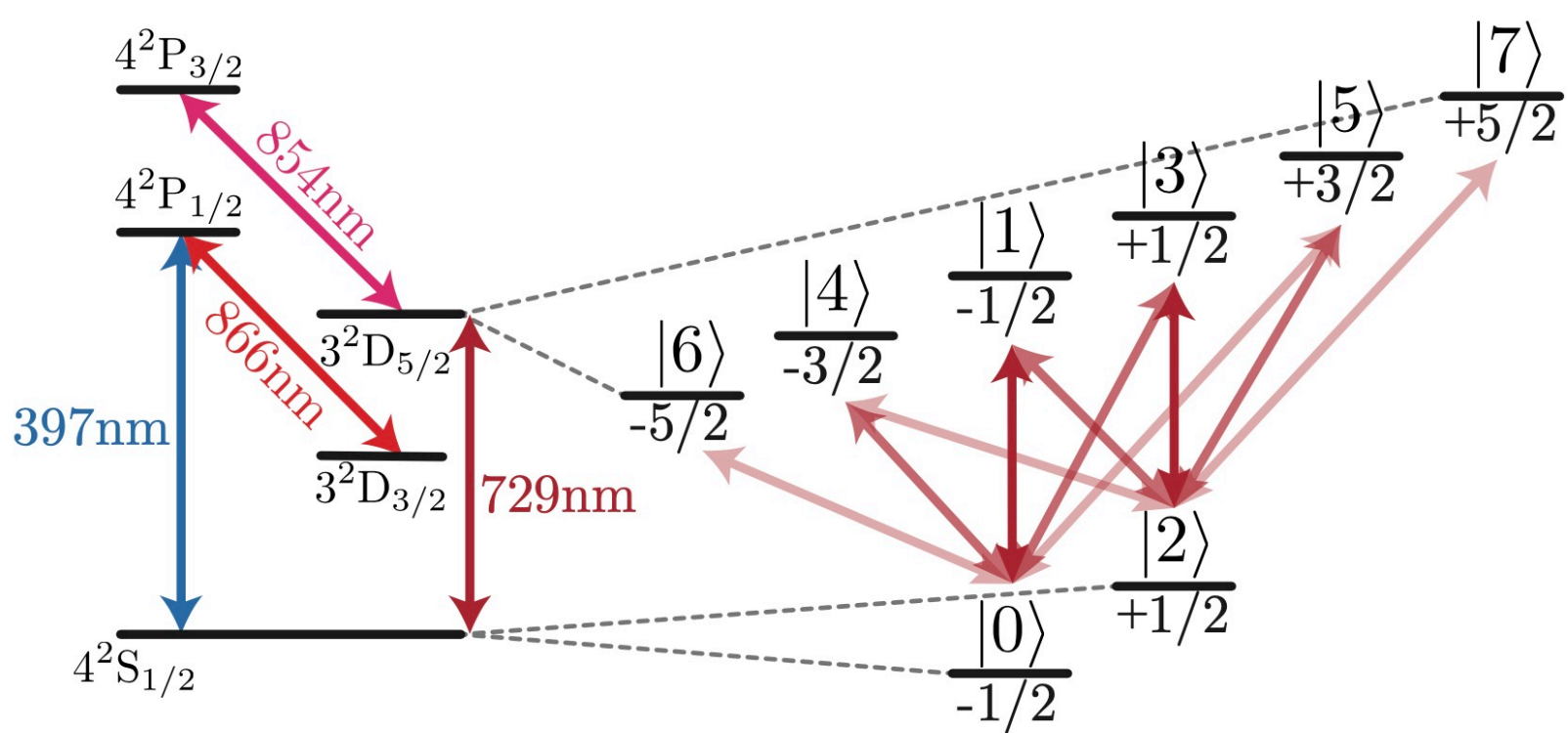
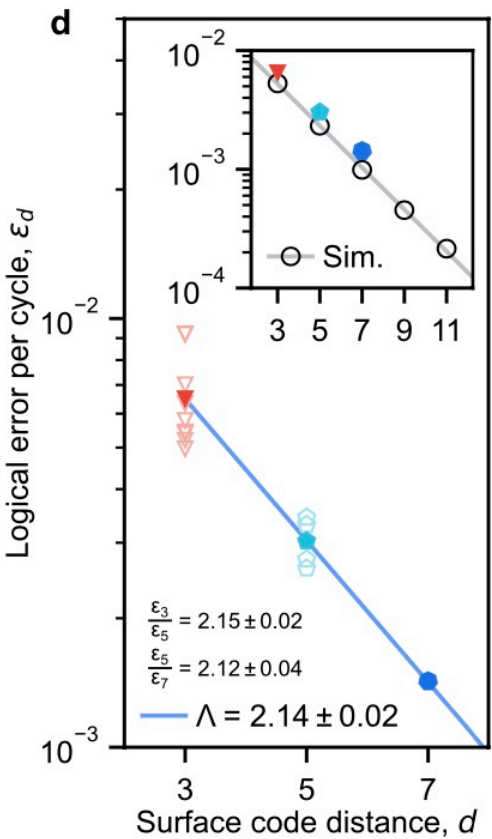
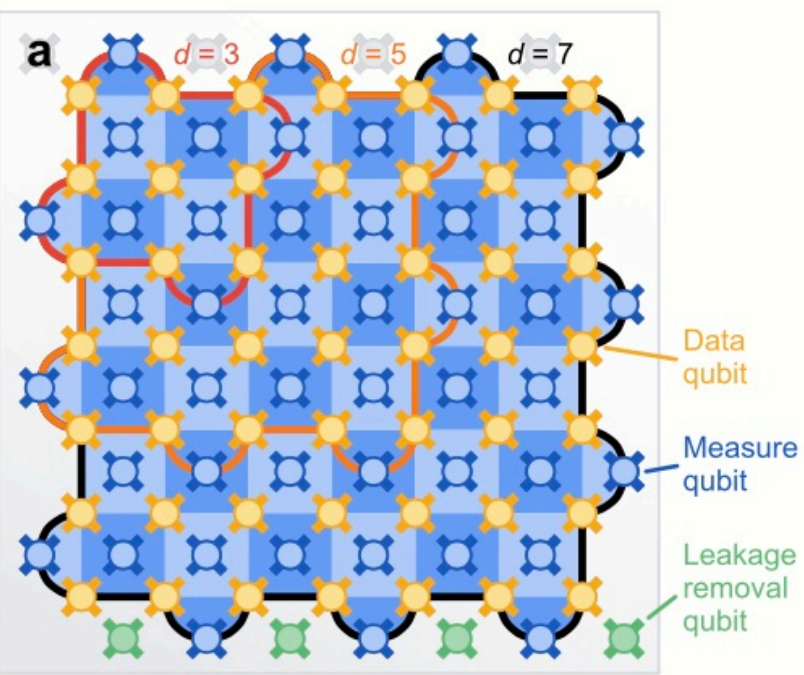
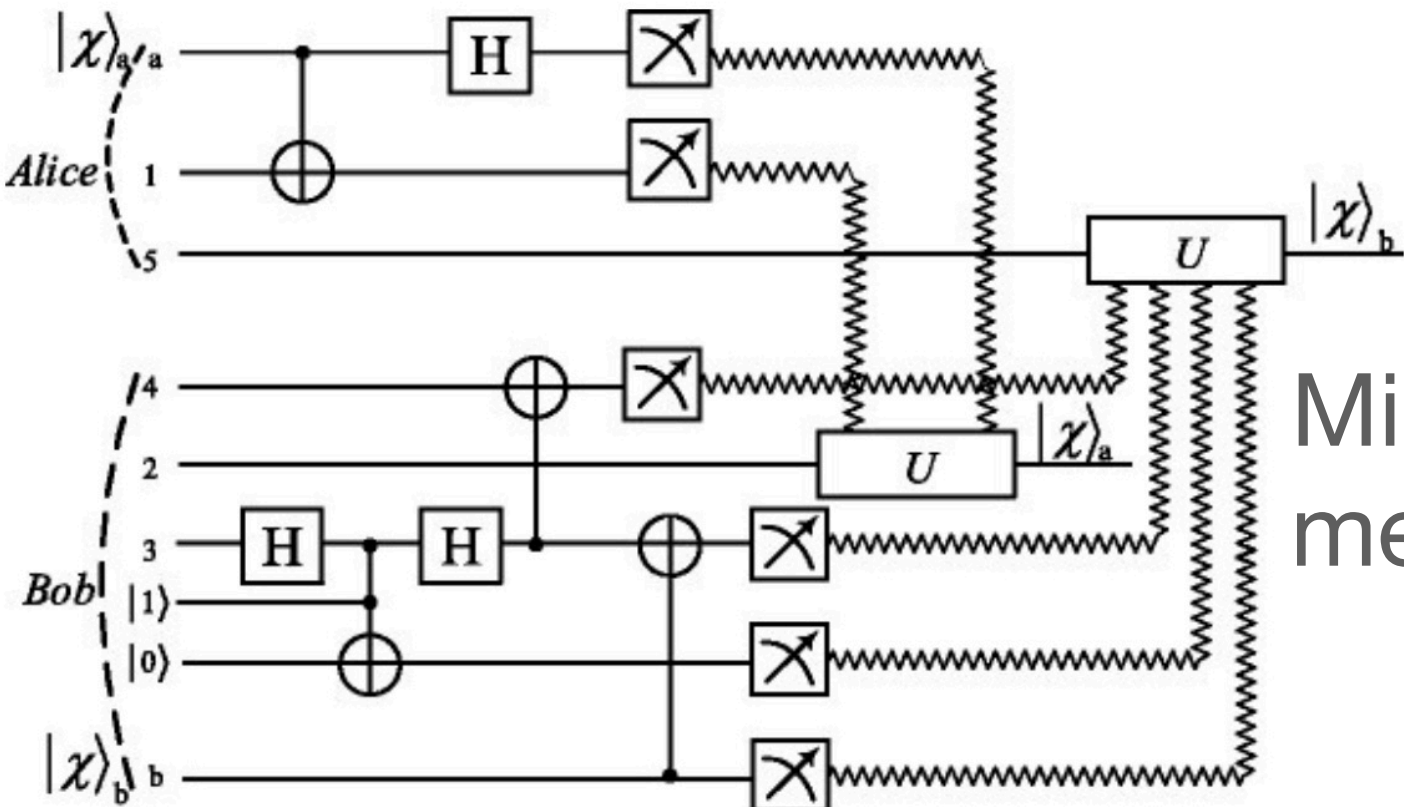


FIG. 1. Level scheme of the $^{40}\text{Ca}^+$ ion.

Qudits with trapped ions

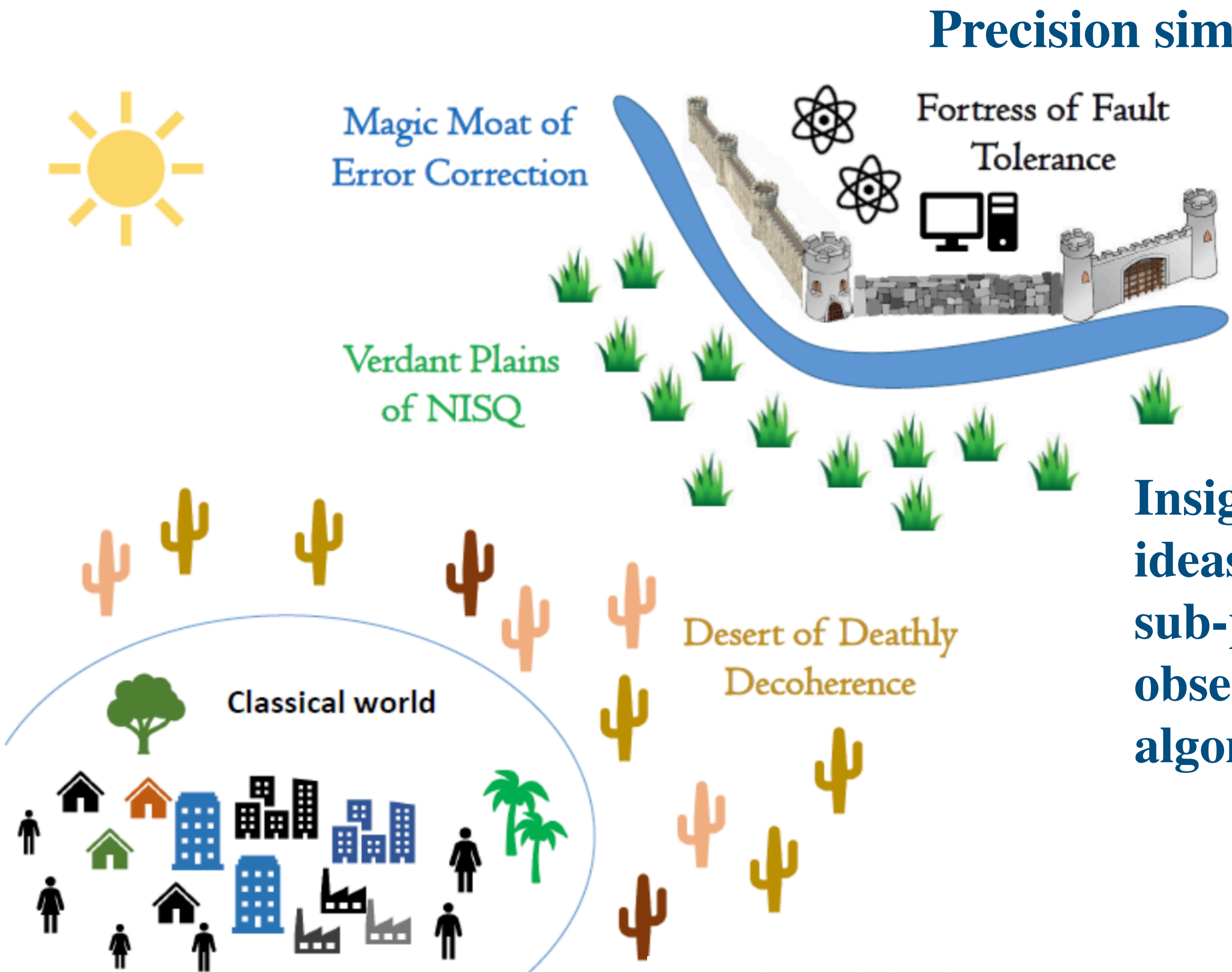


Surface code
>100 superconducting qubits



Mid-circuit measurements

From Classical to Error-Corrected Quantum Computing

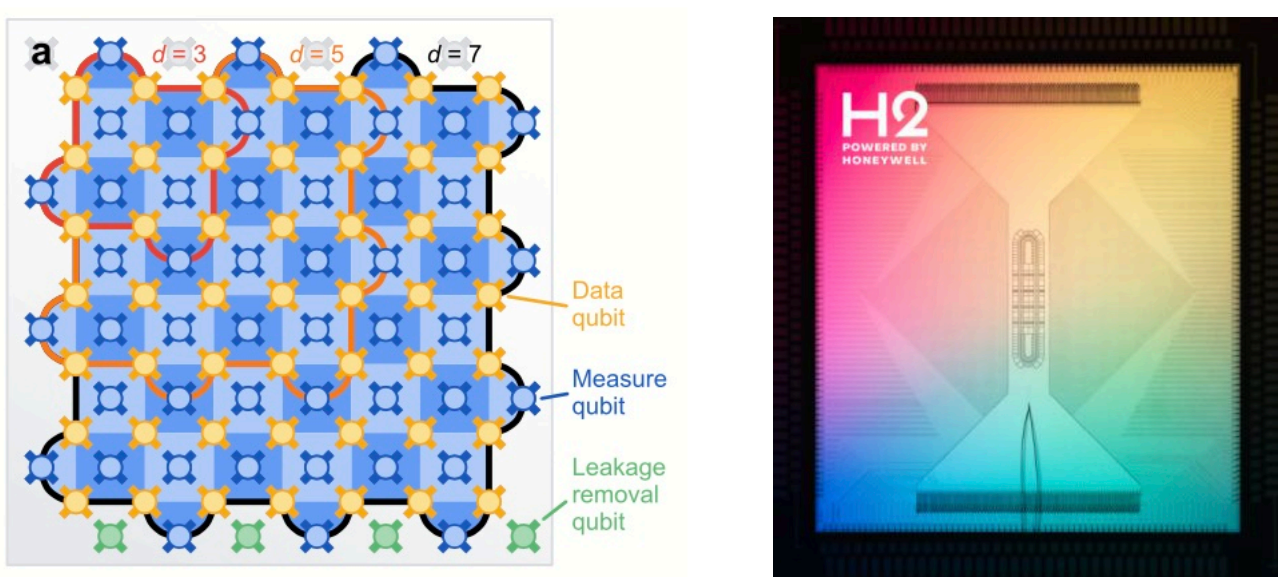


by [Ewan Munro](#), Co-Founder of [Entropica Labs](#).
Landscape of quantum computing from an error correction perspective. Inspired by a [figure](#) by Daniel Gottesman.

Precision simulations

Insights,
ideas,
sub-parts,
observables,
algorithms

Physics Output
Optimization



Logical Qubits

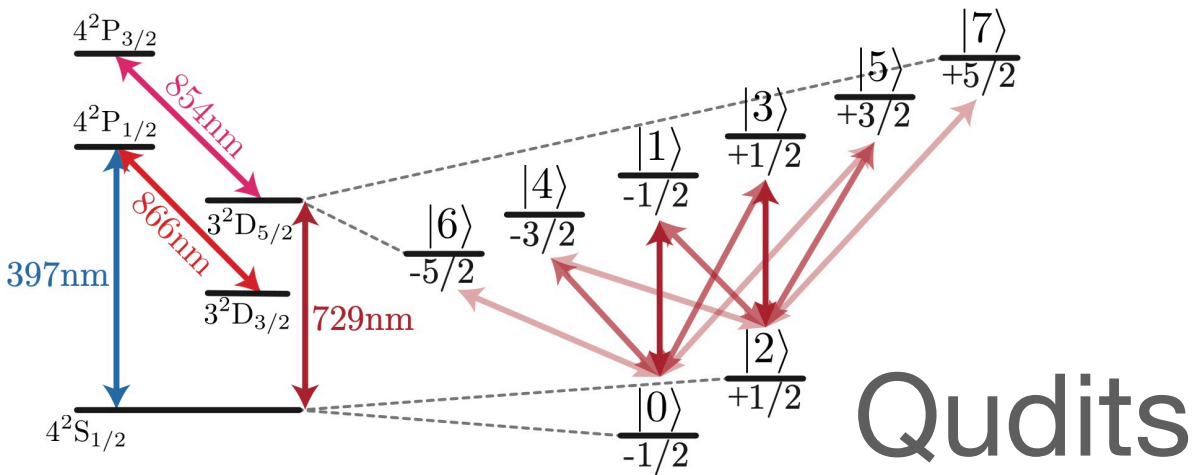
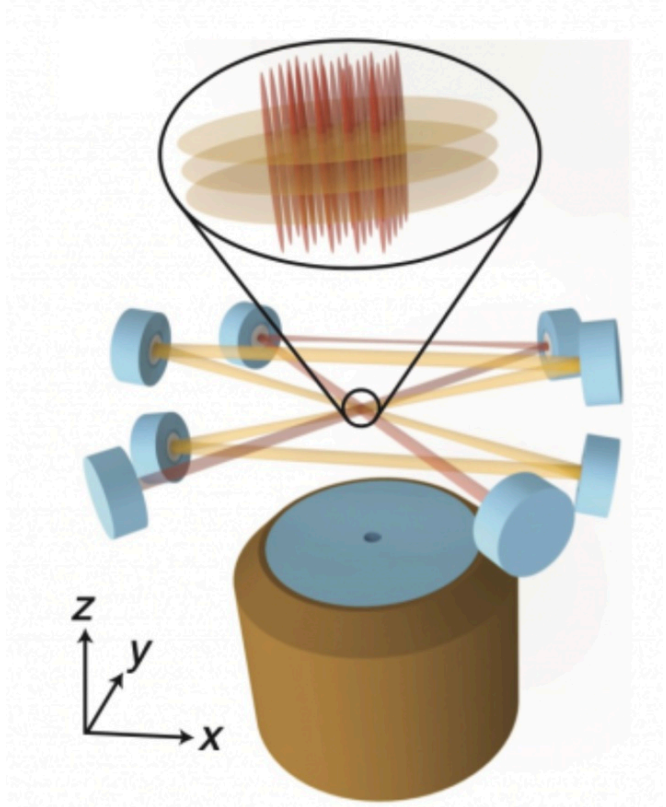


FIG. 1. Level scheme of the $^{40}\text{Ca}^+$ ion.



Cold
Atoms

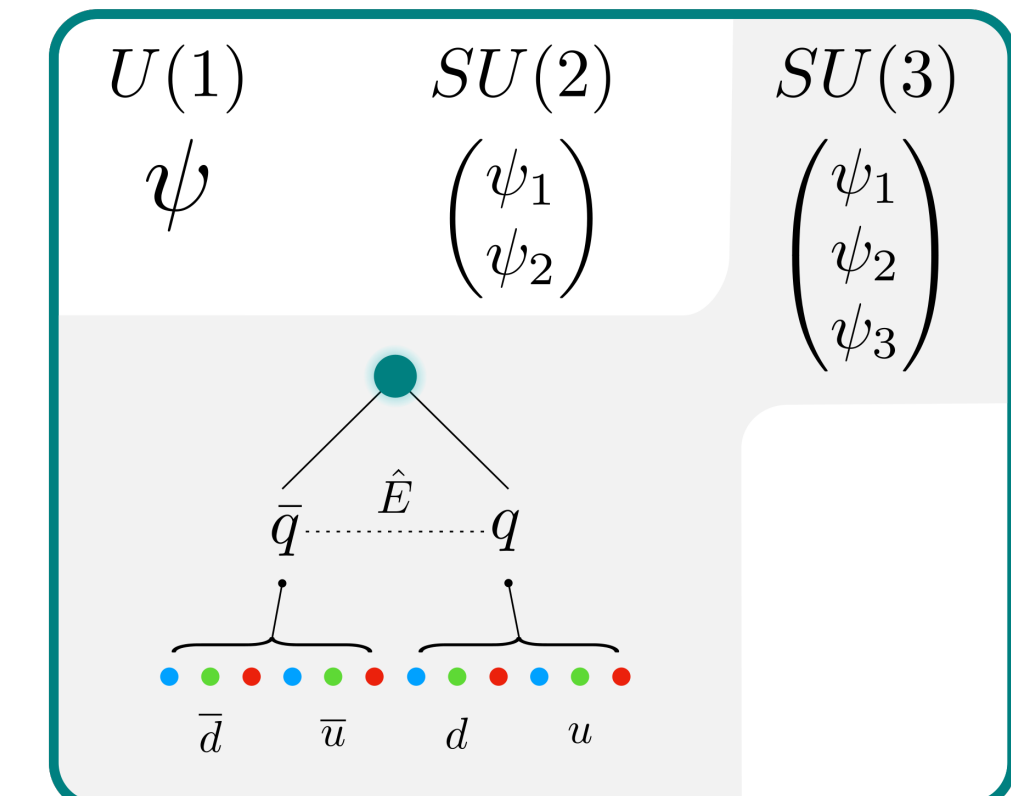
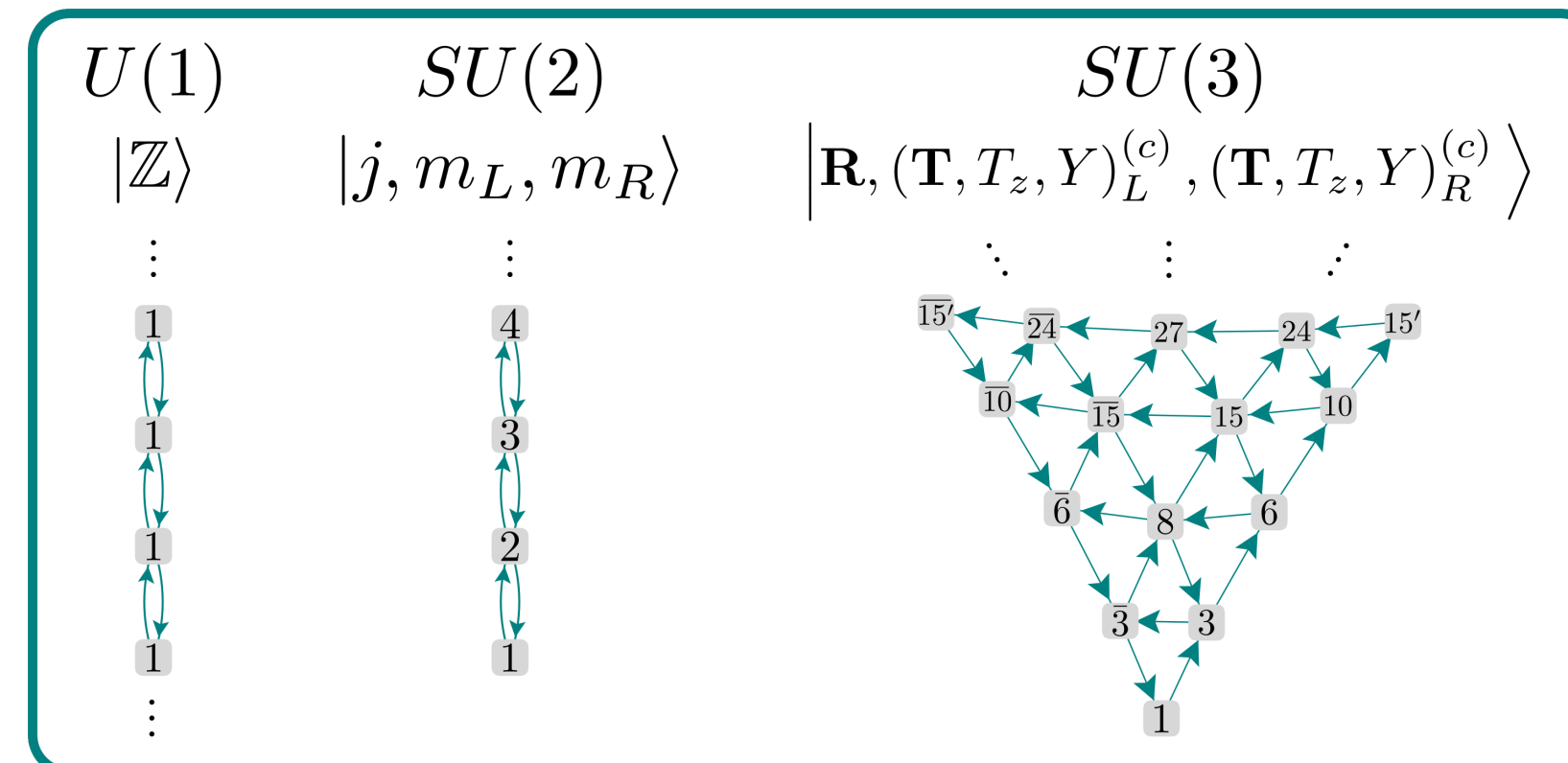
Simulating Lattice Gauge Field Theories

Hamiltonian
Kogut-Susskind
1970's

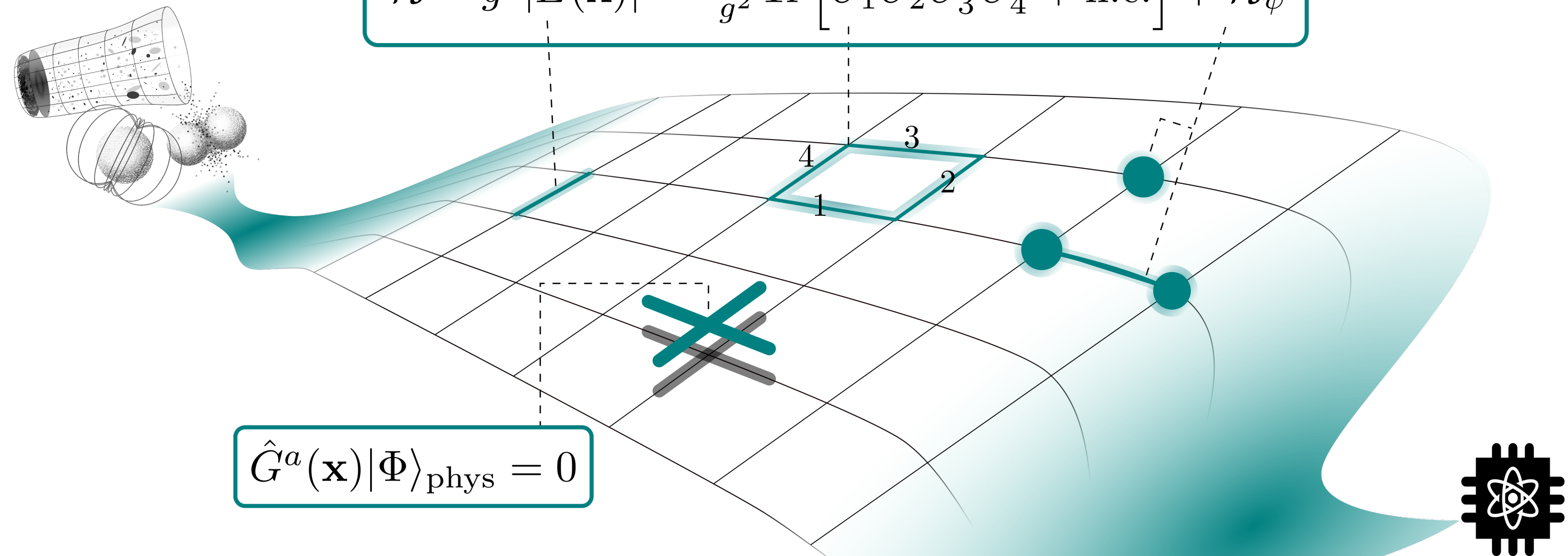
Yang-Mills:
Byrnes-Yamamoto
2005

SU(N):
Zohar et al
(2013)

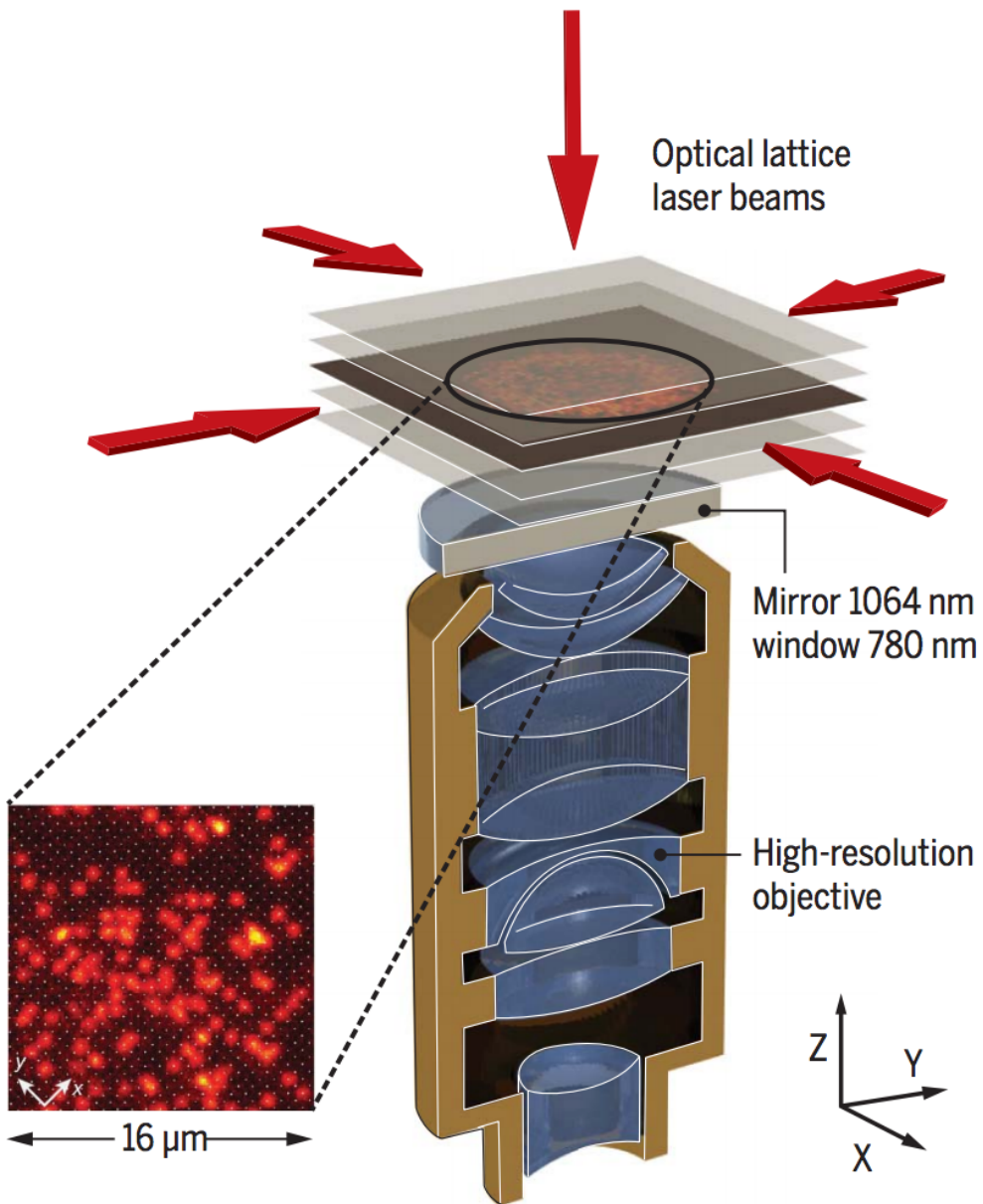
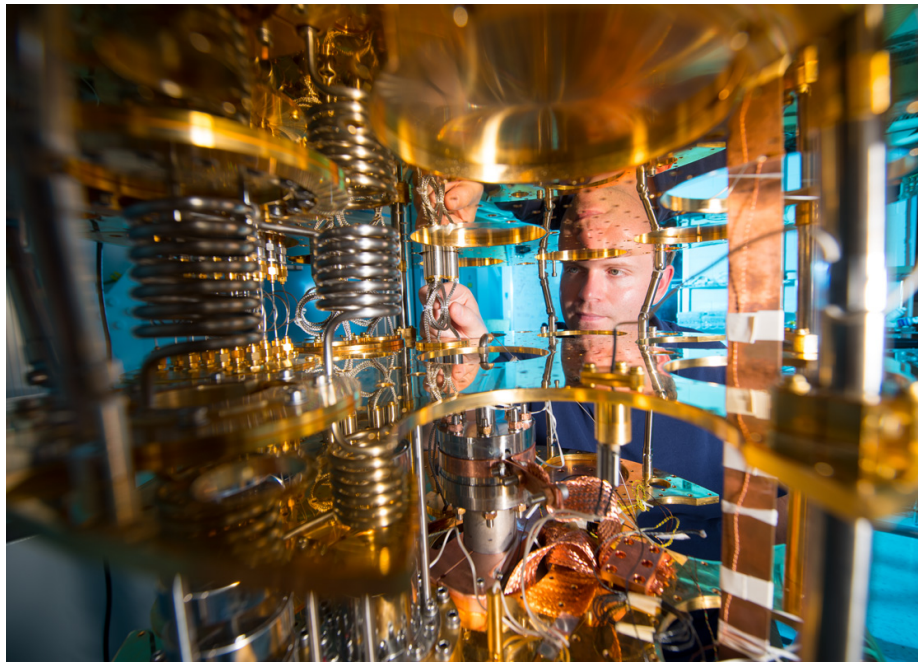
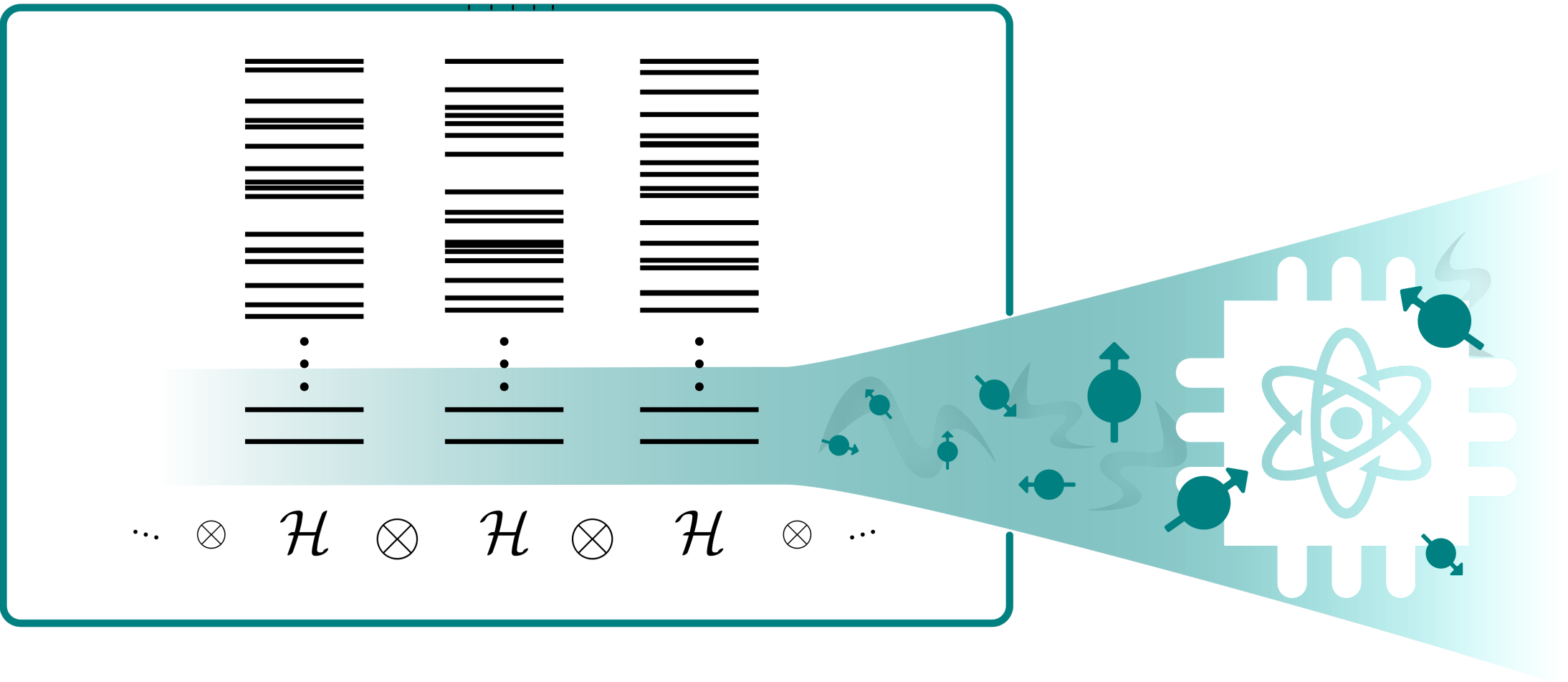
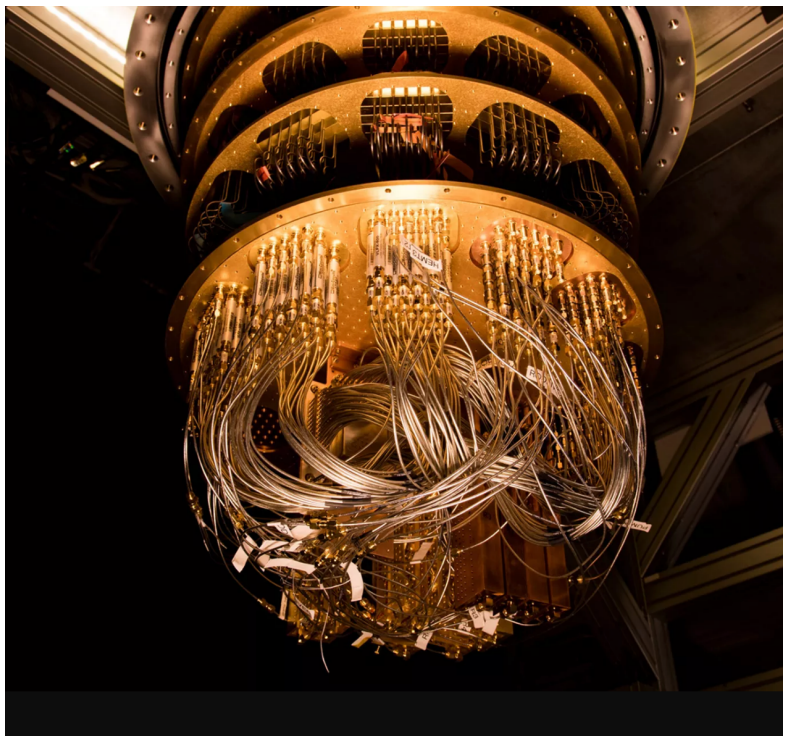
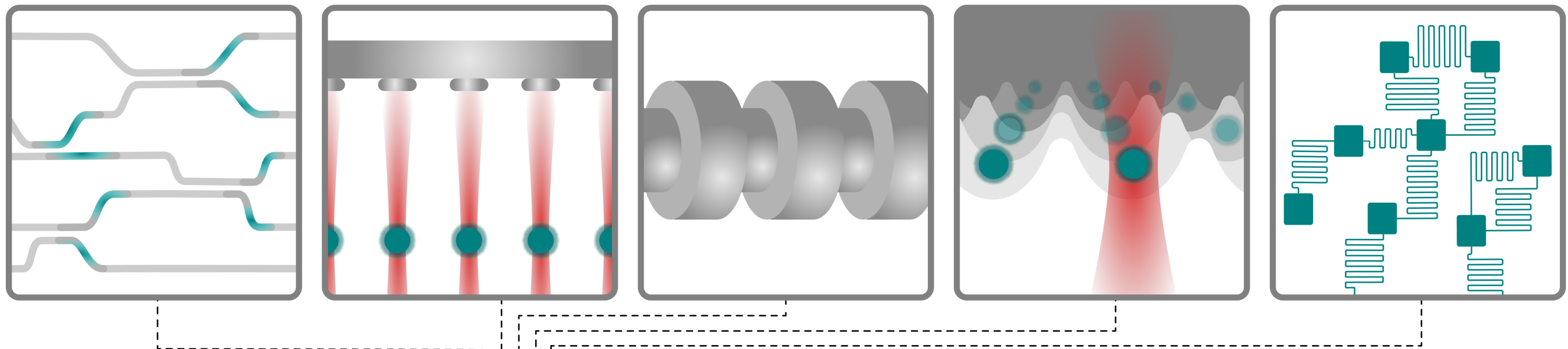
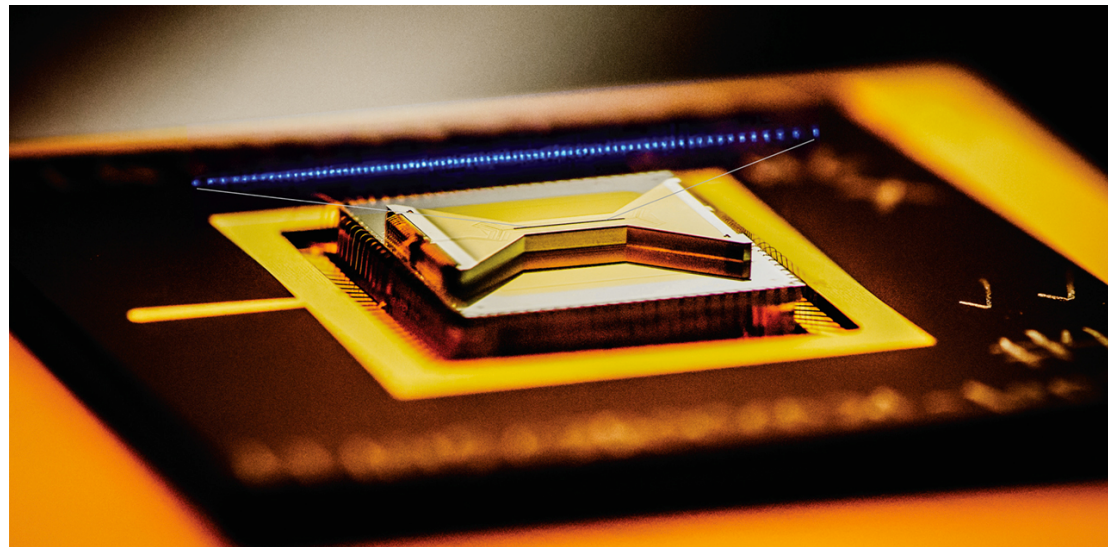
QLM
Banerjee et al
Tagliacozzo et al
(2013)



$$\hat{\mathcal{H}} \sim g^2 |\hat{E}(\mathbf{x})|^2 - \frac{1}{g^2} \text{Tr} \left[\hat{U}_1 \hat{U}_2 \hat{U}_3^\dagger \hat{U}_4^\dagger + \text{h.c.} \right] + \hat{\mathcal{H}}_\psi$$



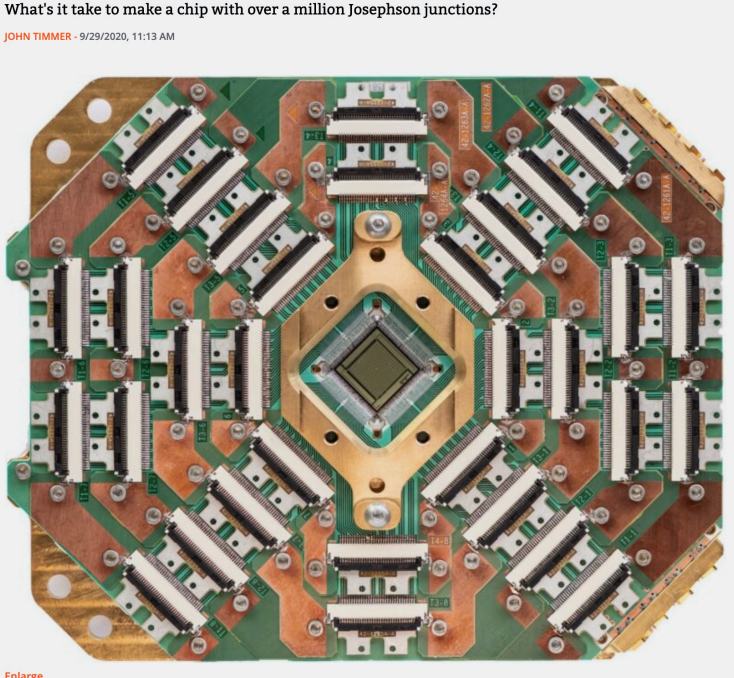
Encoding Systems in Multi-Hilbert Spaces Embedded in Large HPC systems



Map scalar, fermion and vector systems

Optimize for target observables - Physics Aware

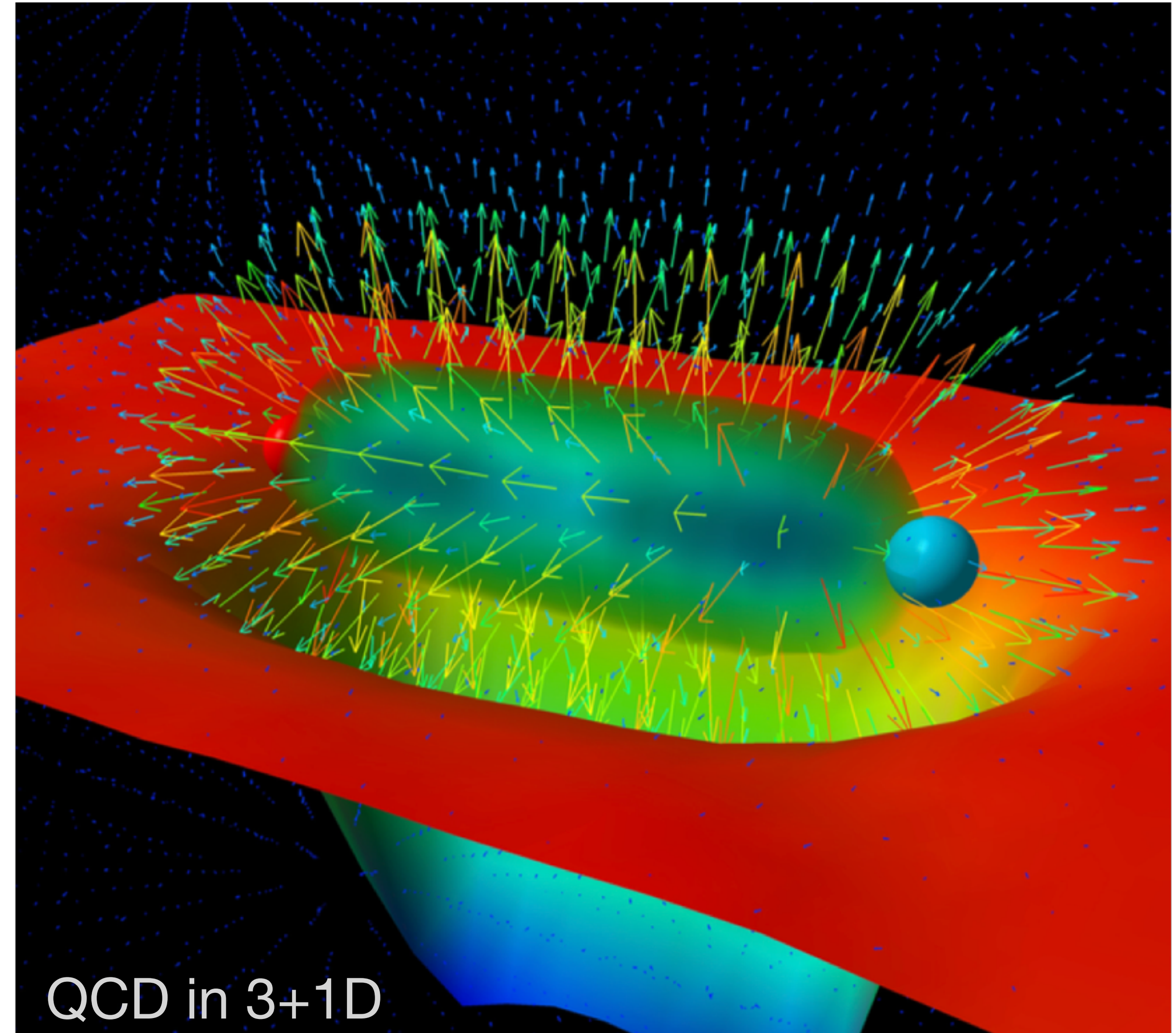
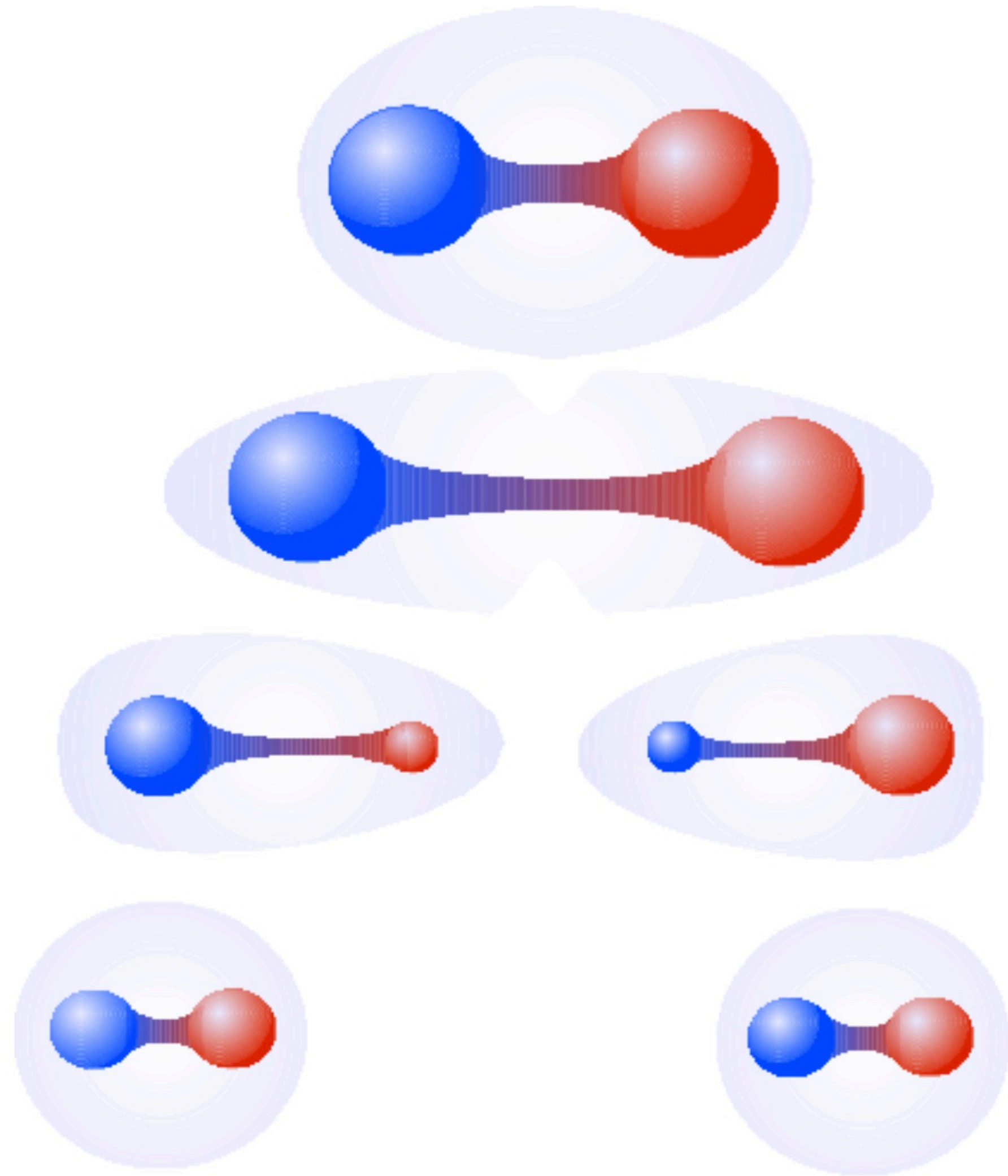
Human-intensive co-design exploration



Low-Dimensional Models:

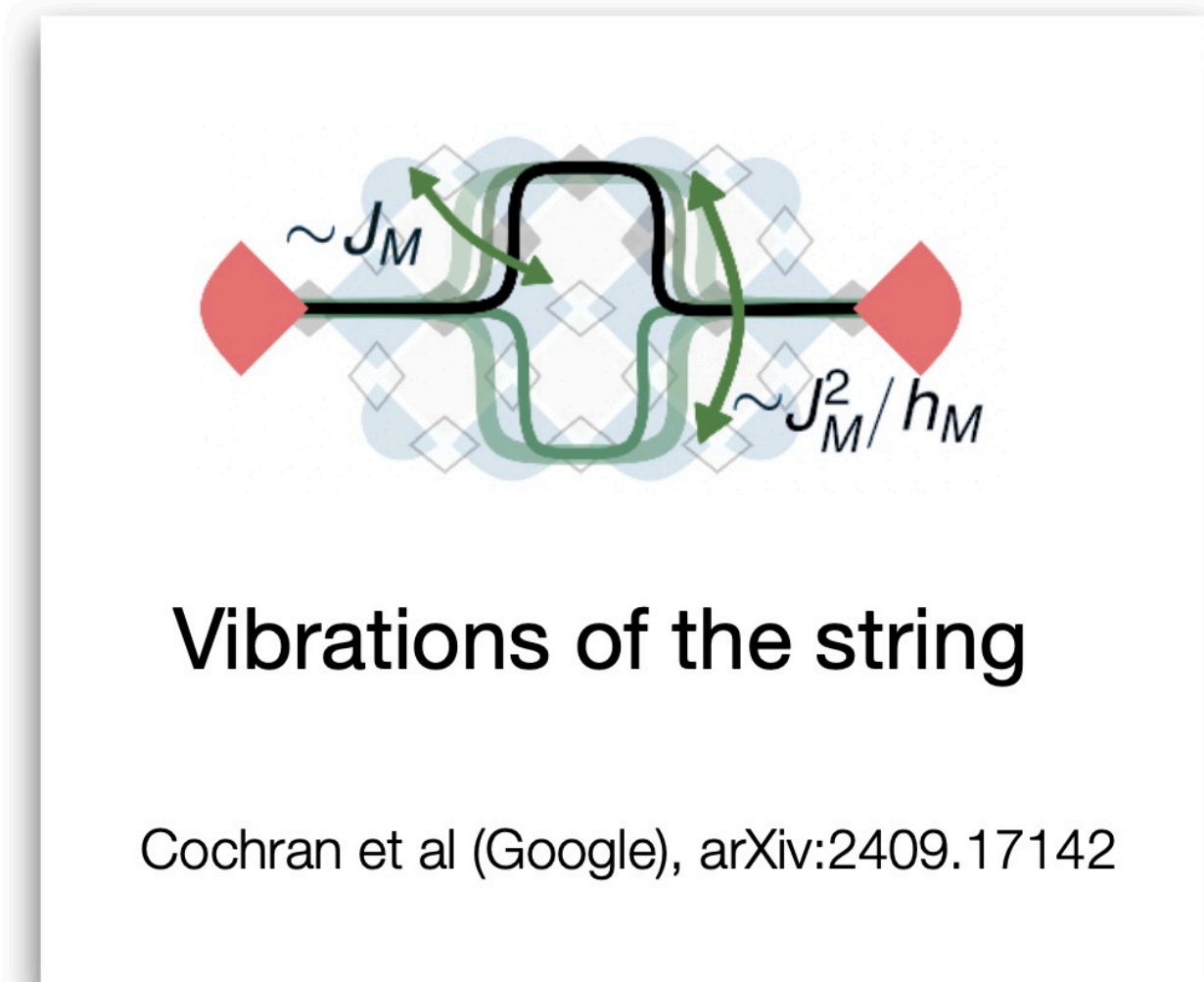
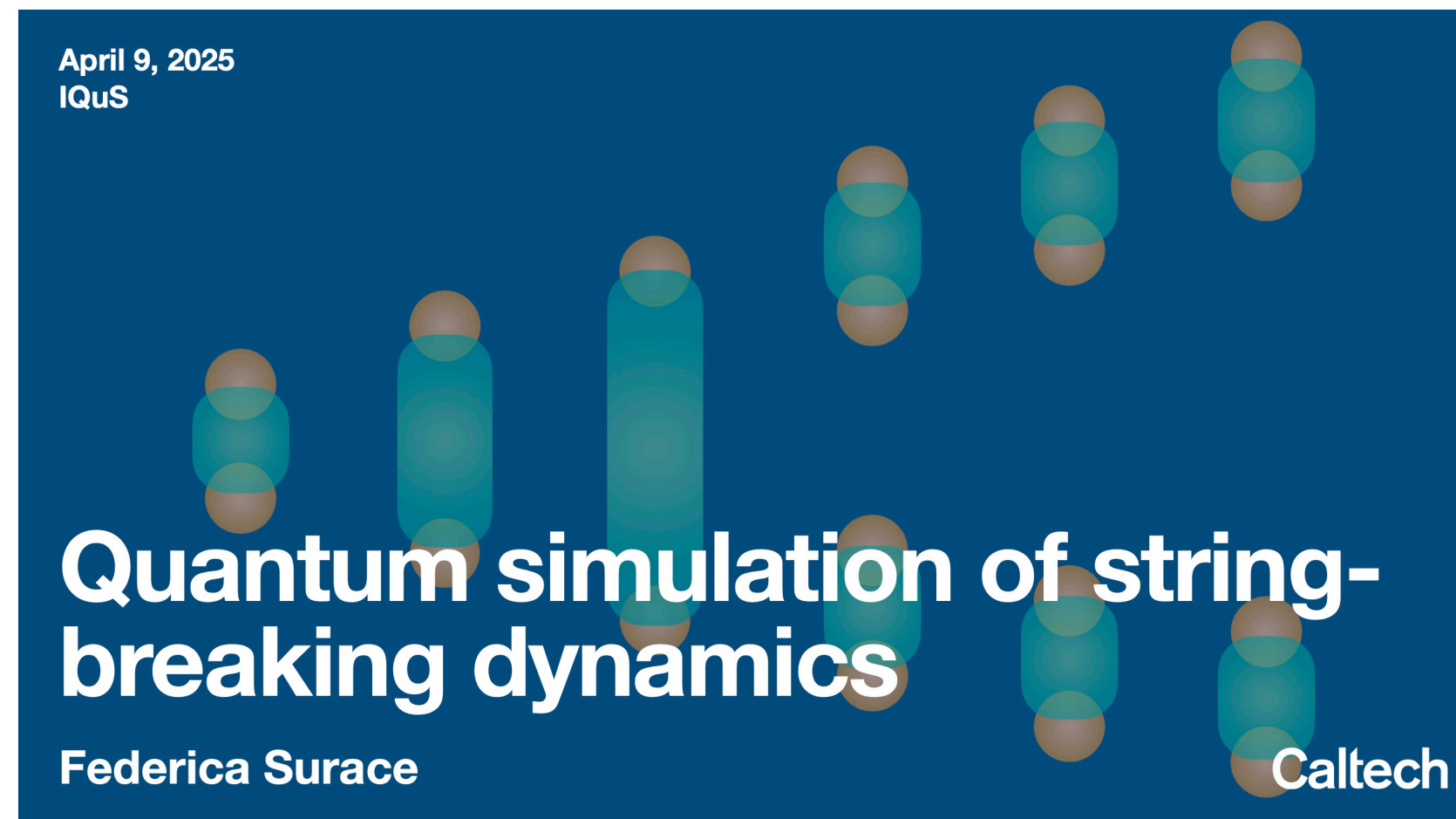
e.g., Quantum Electromagnetism in 1 Space and 1 Time Dimensions

2+1D simulations are starting - Zn-like



This model is being used by several groups pursuing quantum simulations

String Breaking as an Example



Quantum simulation of string breaking: summary

Models

1+1D

- Z_2 LGT, trapped ions
- $SU(2)$ LGT, superconducting qubits
- $U(1)$ LGT, atoms in optical lattices

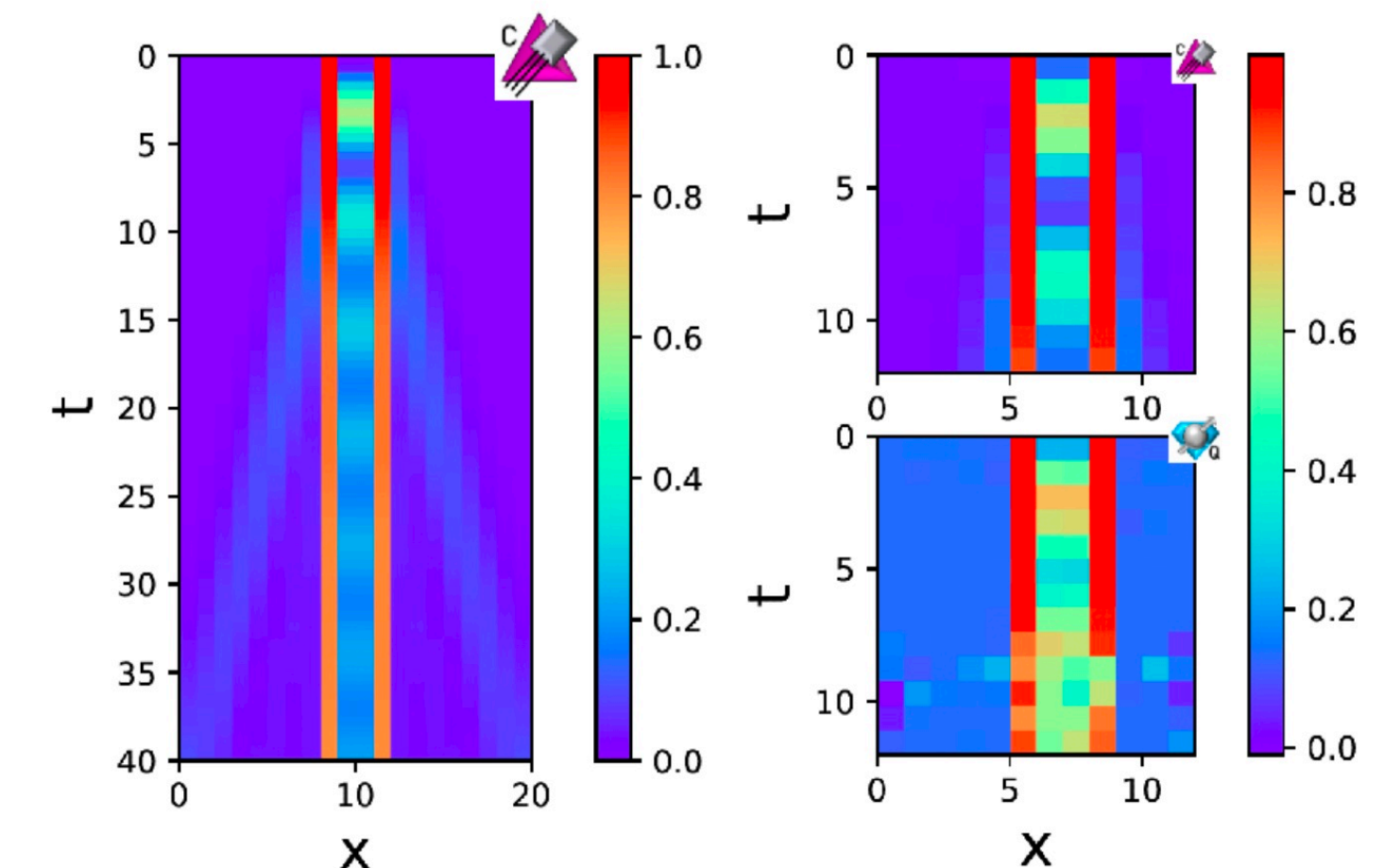
2+1D

- Z_2 LGT, superconducting qubits
- $U(1)$ LGT, Rydberg atom arrays

Phenomenology

- Ground state
- Quench dynamics
 - Resonant string breaking
- Dynamics of strings in 2+1D
- **Ramp dynamics**

$SU(2)$ lattice gauge theory 1+1D





Confinement and Scalable Circuits

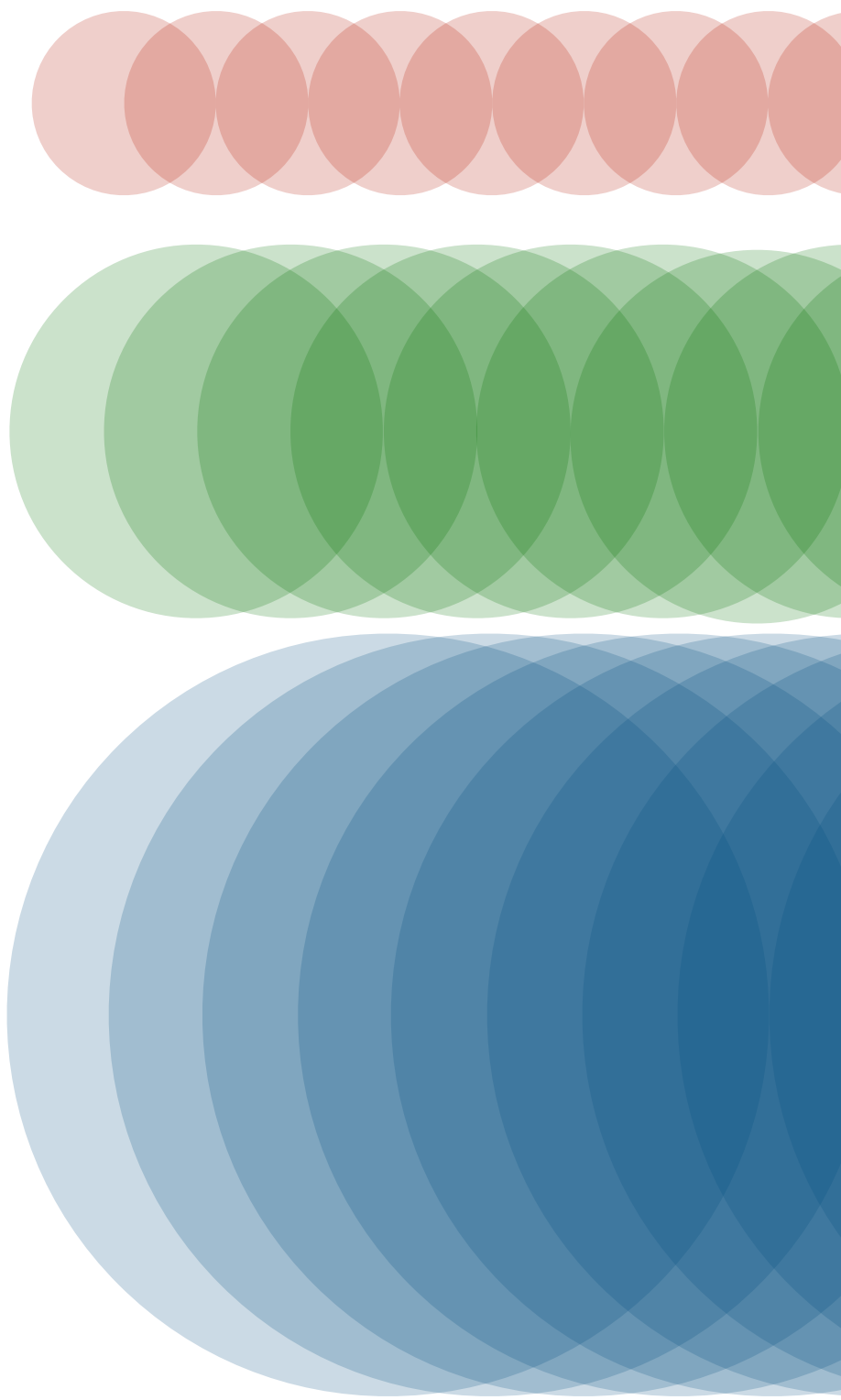
(2023-)



Roland Farrell, Marc Illa,
Anthony Ciavarella and MJS

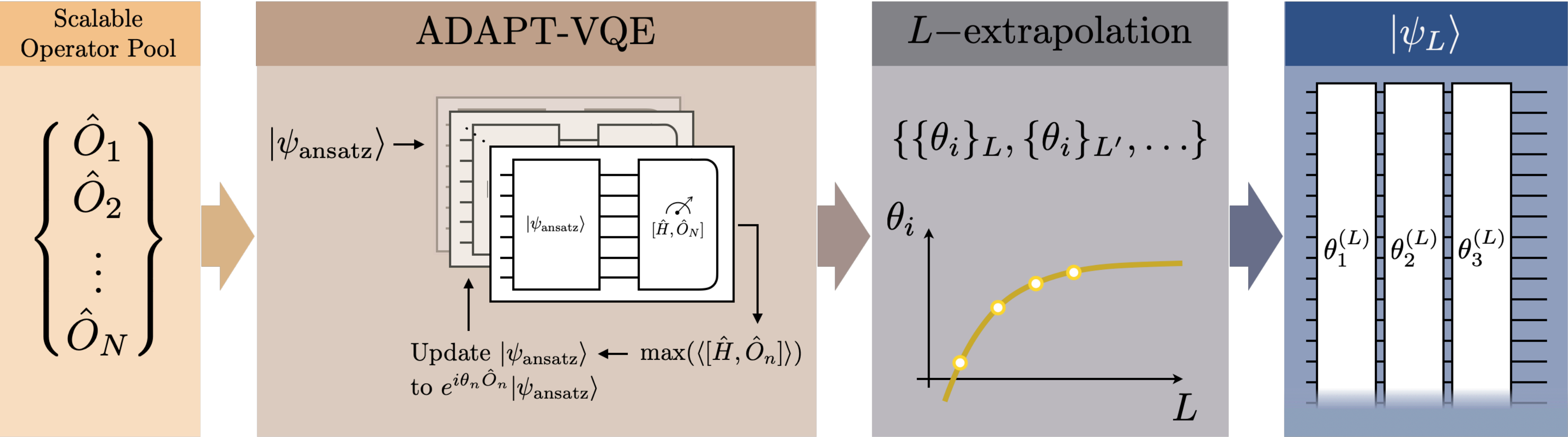
$$\hat{H} = \hat{H}_m + \hat{H}_{kin} + \hat{H}_{el} = \frac{m}{2} \sum_{j=0}^{2L-1} \left[(-1)^j \hat{Z}_j + \hat{I} \right] + \frac{1}{2} \sum_{j=0}^{2L-2} \left(\hat{\sigma}_j^+ \hat{\sigma}_{j+1}^- + \text{h.c.} \right) + \frac{g^2}{2} \sum_{j=0}^{2L-2} \left(\sum_{k \leq j} \hat{Q}_k \right)^2$$

Local Nearest Neighbor Non-local



Symmetries and Confinement

Classical Extrapolations



Classical Optimization Quantum Implementation

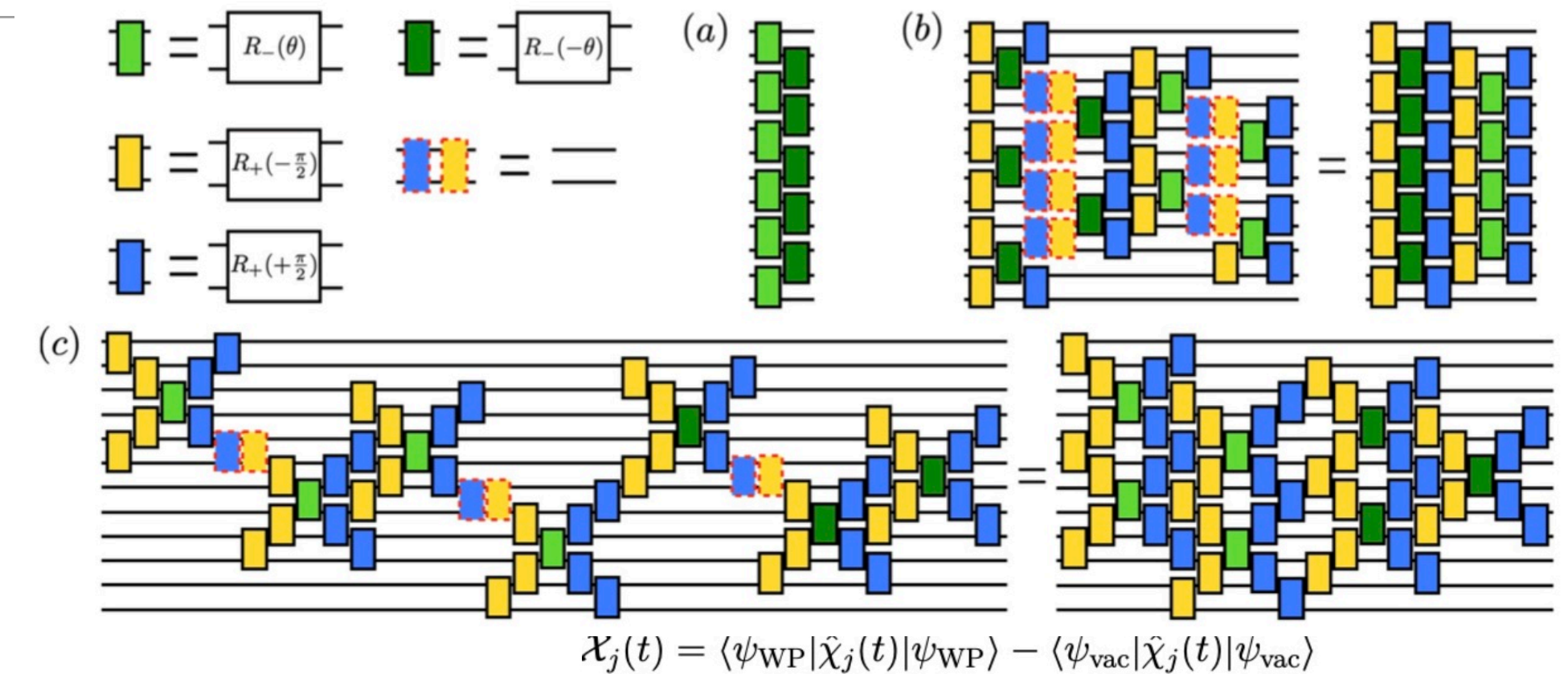
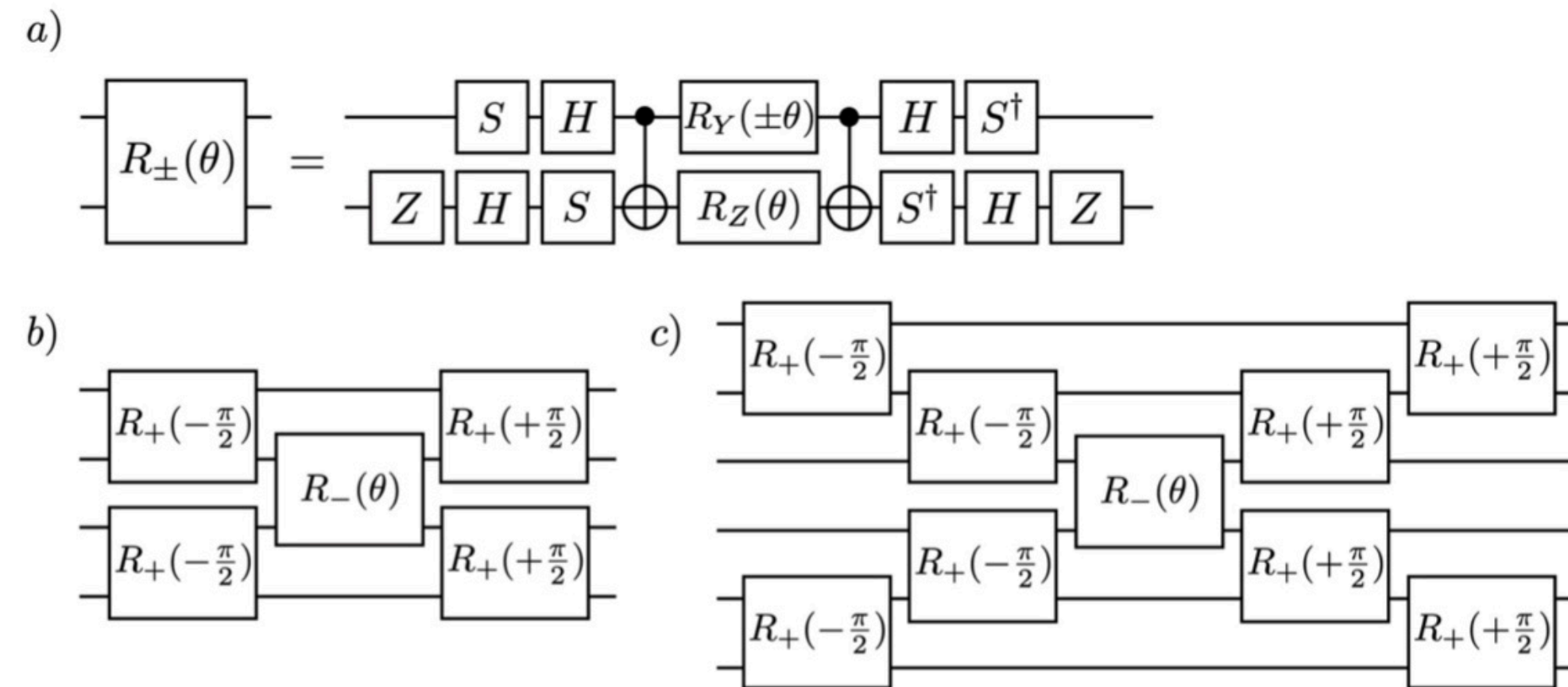
Scalable Circuits for Preparing Ground States on Digital Quantum Computers: The Schwinger Model Vacuum on 100 Qubits
Roland C. Farrell, Marc Illa, Anthony N. Ciavarella, and Martin J. Savage
PRX Quantum **5**, 020315 – Published 18 April 2024

Quantum simulations of hadron dynamics in the Schwinger model using 112 qubits
Roland C. Farrell, Marc Illa, Anthony N. Ciavarella, and Martin J. Savage
Phys. Rev. D **109**, 114510 – Published 10 June 2024

Builds upon ADAPT-VQE
by Sophia Economou *et al.*

Production using IBM's QPU Torino

(The largest quantum simulation that had been performed)

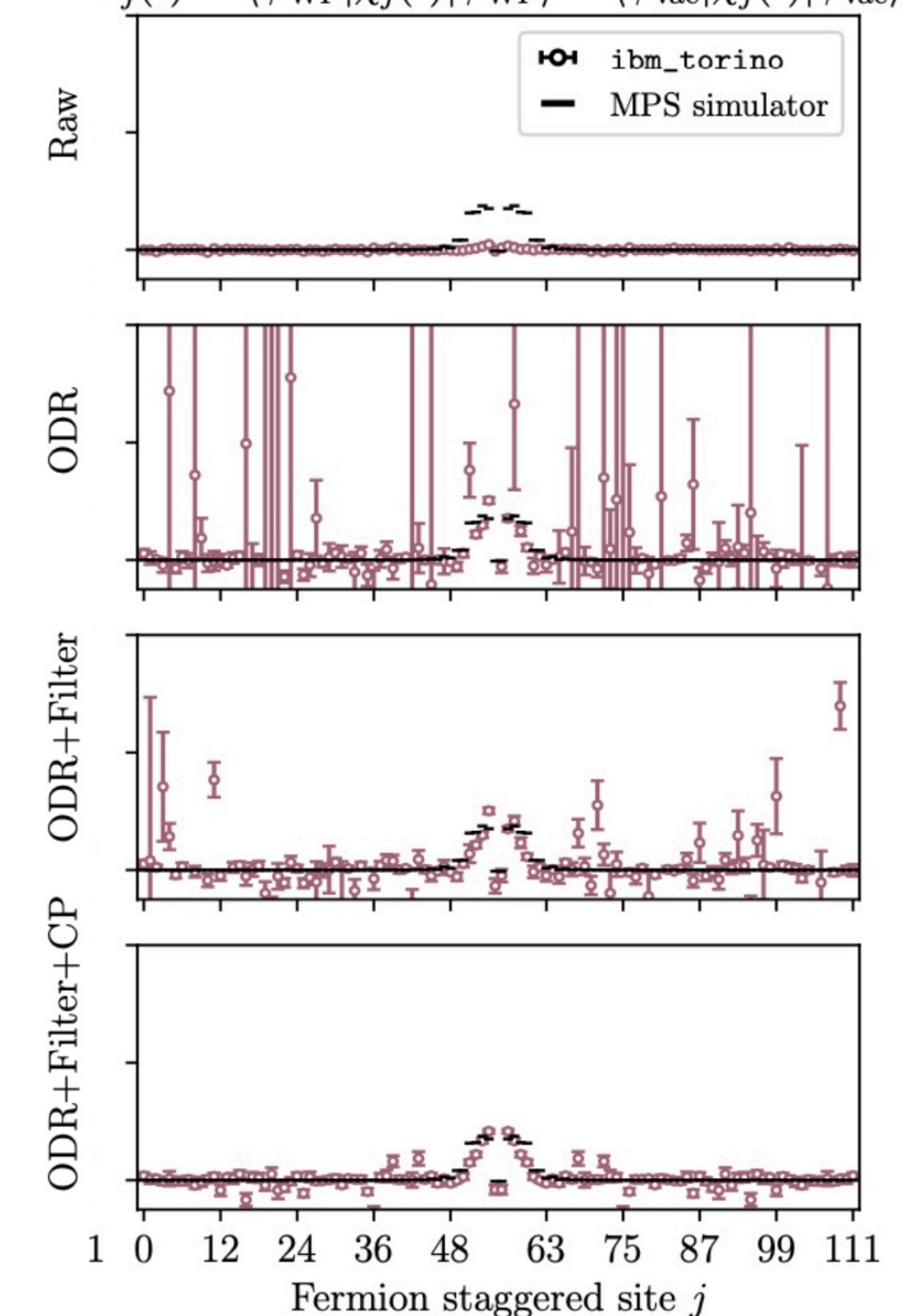
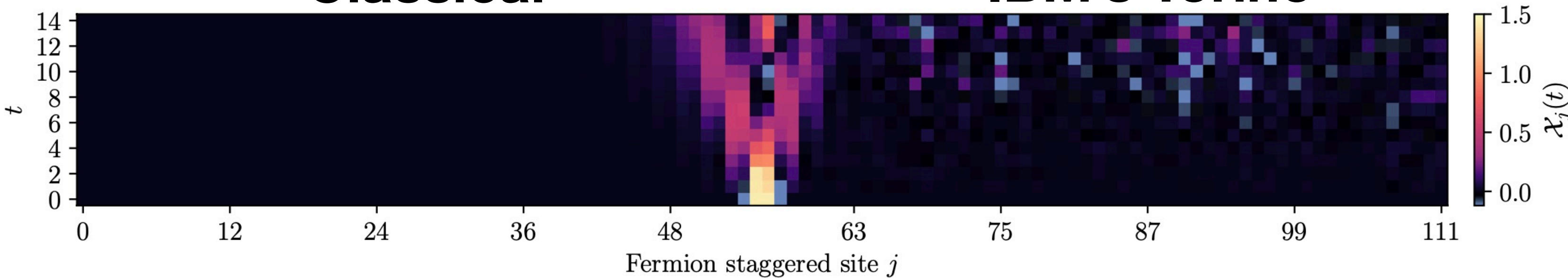


Production highlights

- 14K CNOTs for 14 Trotter steps
- 1.05 Trillion total CNOTs applied
- 154 Million shots
- 112 qubits x 370 depth

Classical

IBM's Torino



Decoherence Renormalization

Mitigating Depolarizing Noise on Quantum Computers with Noise-Estimation Circuits

Miroslav Urbanek, Benjamin Nachman, Vincent R. Pascuzzi, Andre He, Christian W. Bauer, and Wibe A. de Jong
Phys. Rev. Lett. **127**, 270502 – Published 27 December 2021

Self-mitigating Trotter circuits for SU(2) lattice gauge theory on a quantum computer

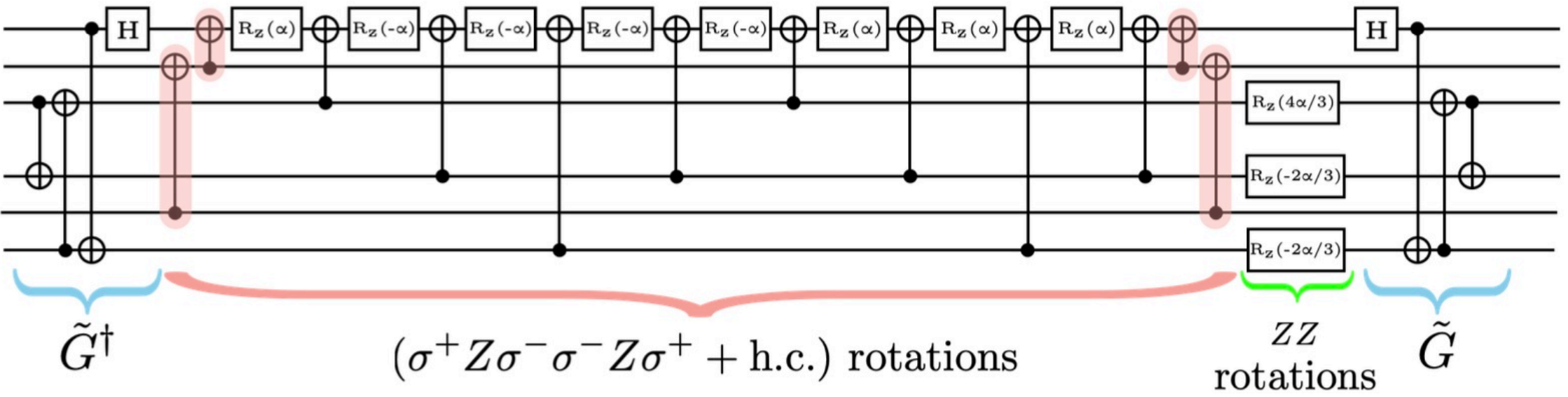
Sarmed A Rahman, Randy Lewis, Emanuele Mendicelli, and Sarah Powell
Department of Physics and Astronomy, York University,
Toronto, Ontario, Canada, M3J 1P3

(Dated: May 2022. Updated: October 2022.)

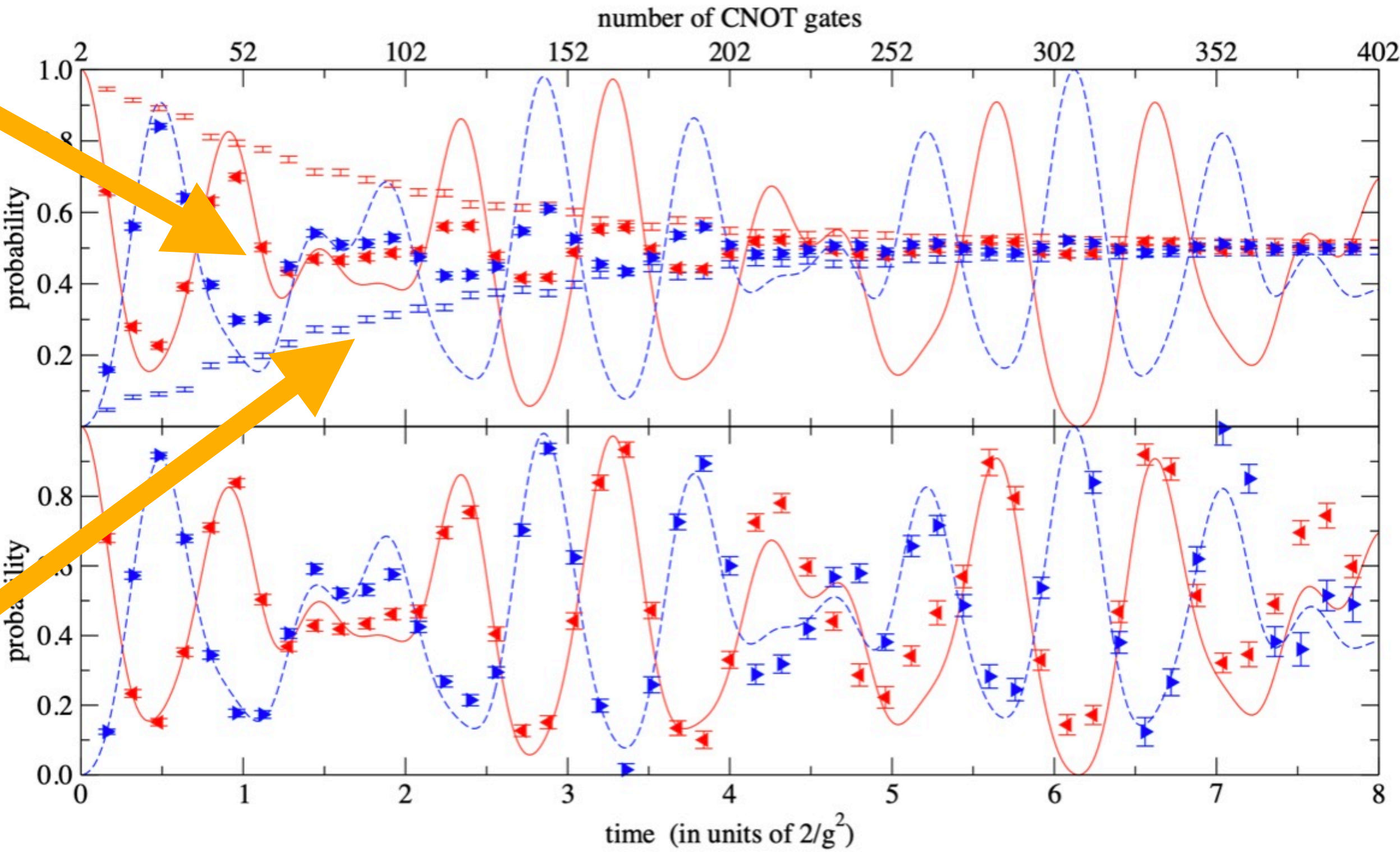
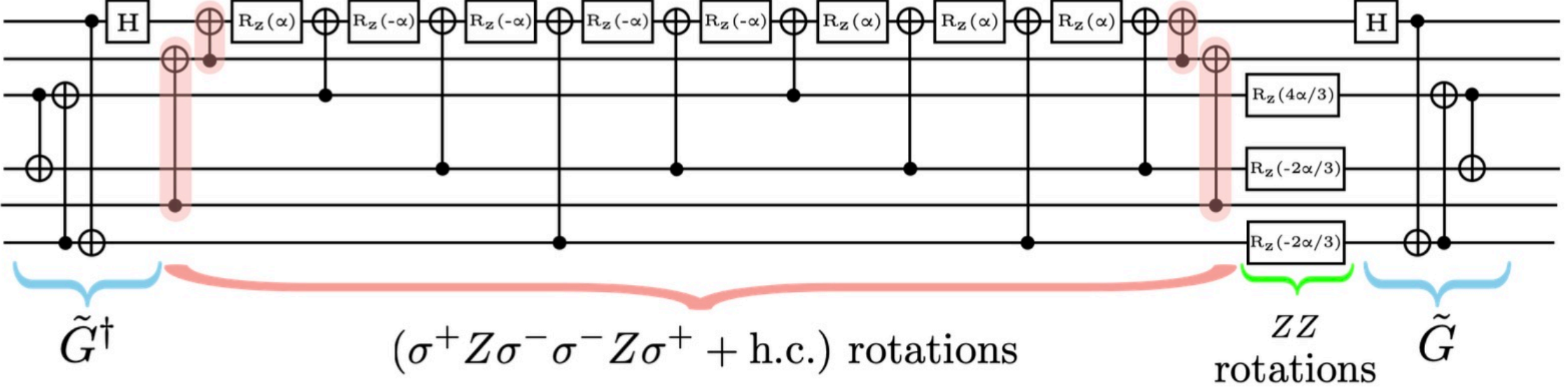
The device is approaching a classical, depolarized set of qubits as time goes by.

Mitigation methods are essential and effective

“Physics circuit”

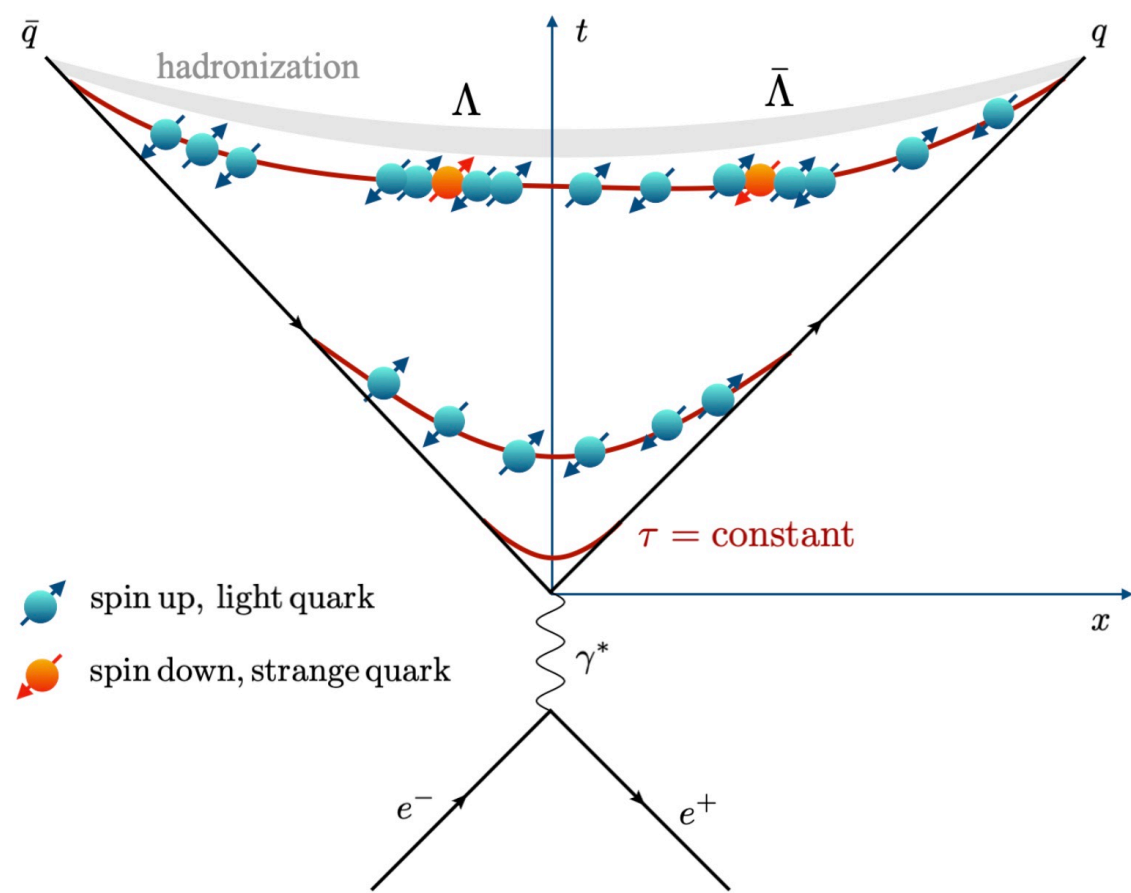


“Mitigation circuit” - all angles set to zero (e.g.)

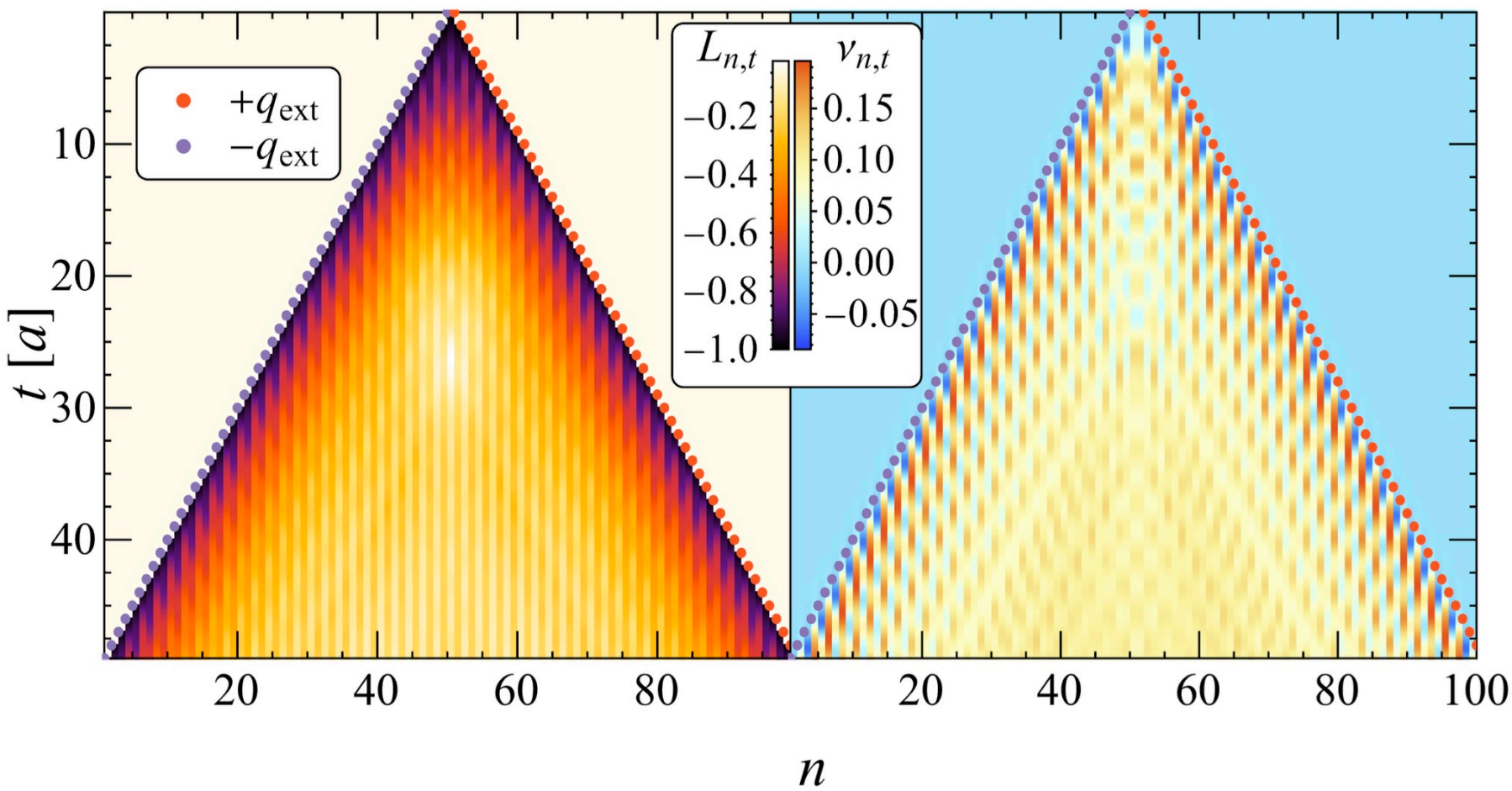


Hadronization and Fragmentation

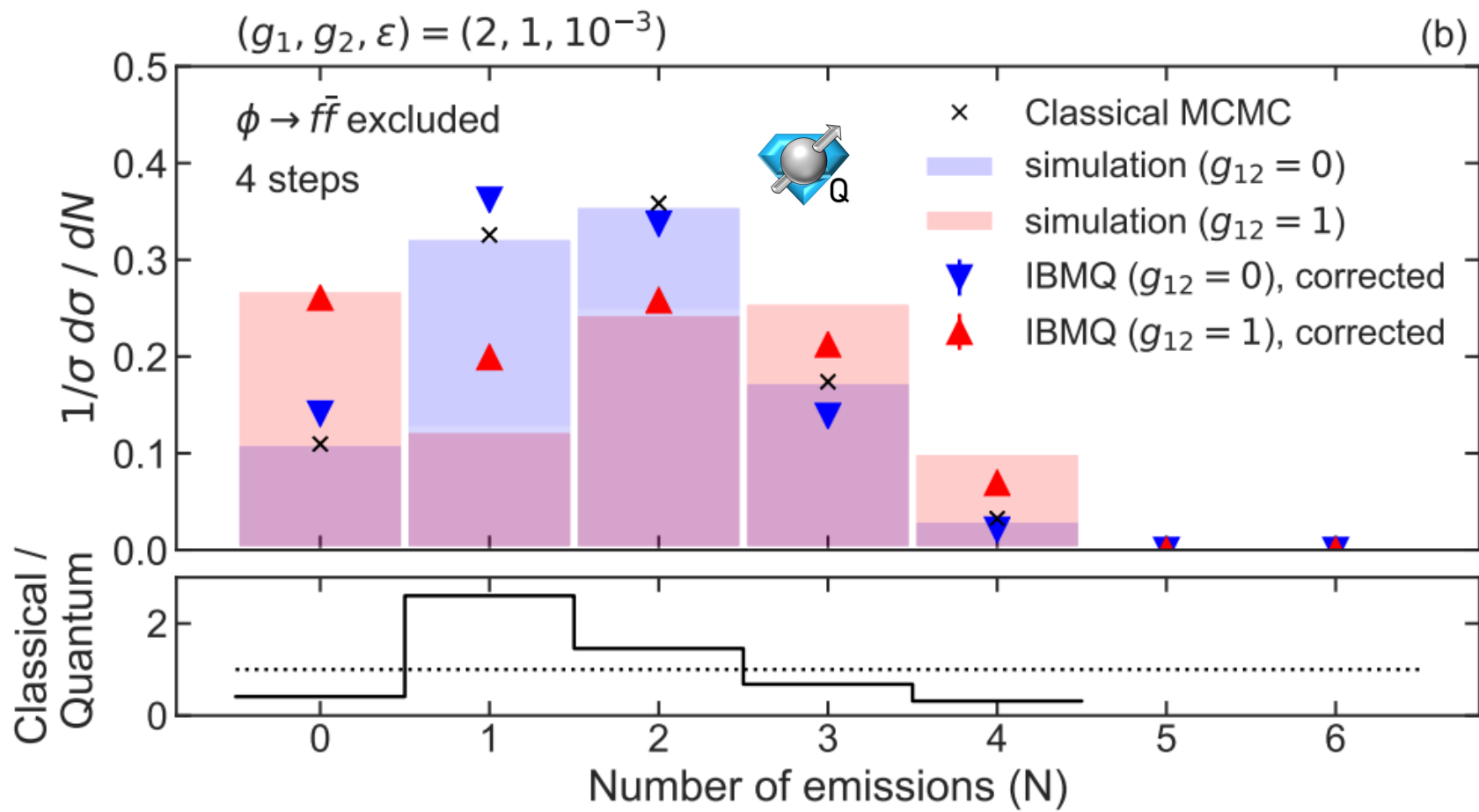
Baryon Entanglement, Barata *et al*, Gong *et al*



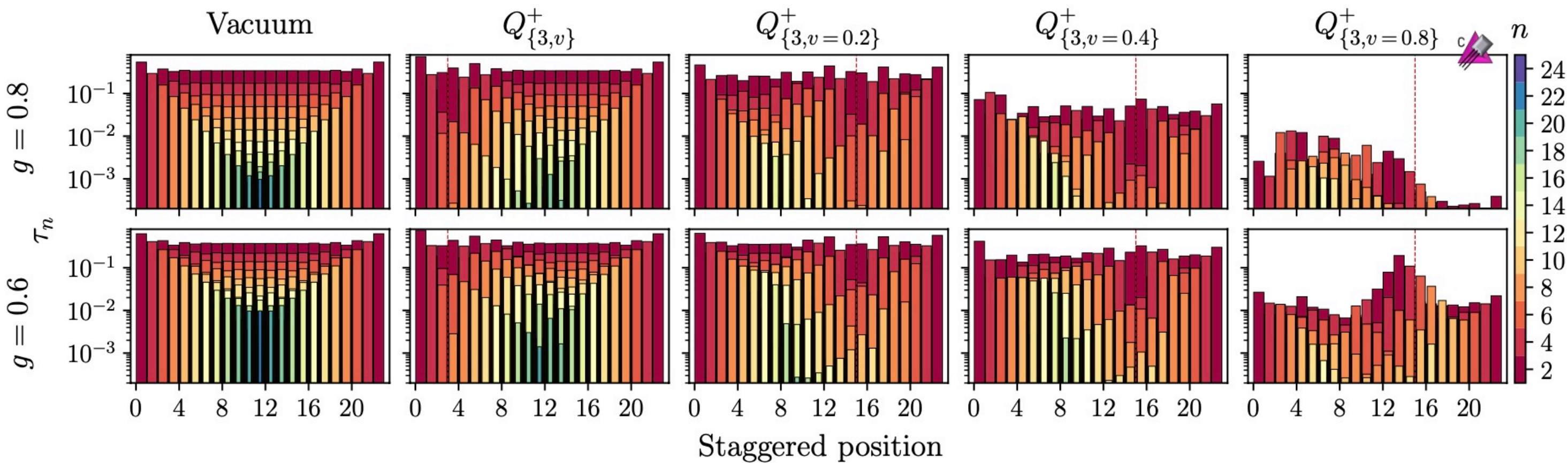
Hadronization, Florio *et al*



Fragmentation, Bauer *et al*

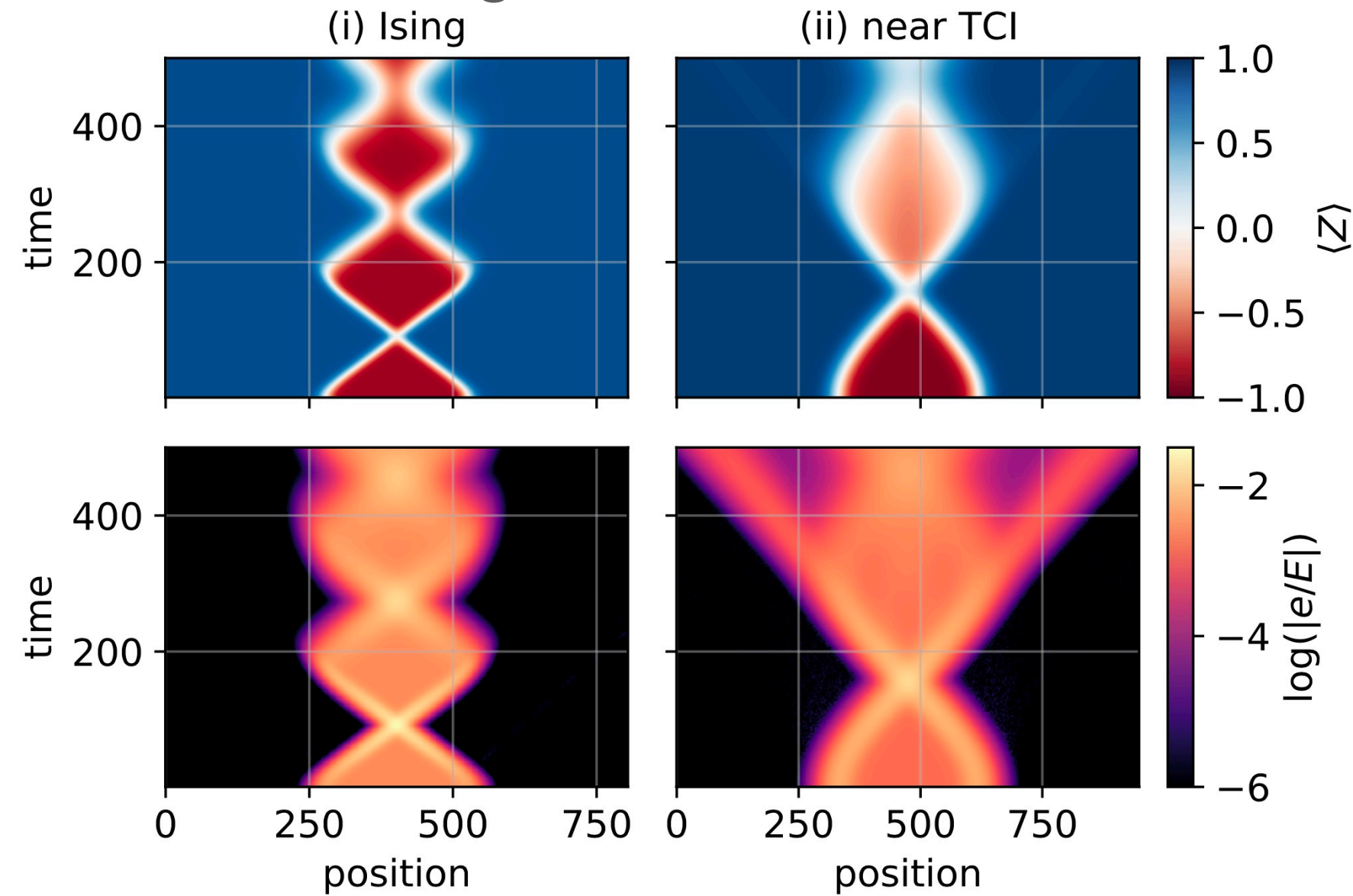


Energy Loss and Multi-partite Entanglement, Farrell *et al*



1+1D Preparing for the Future

Ising+ , Milsted et al



Adopted

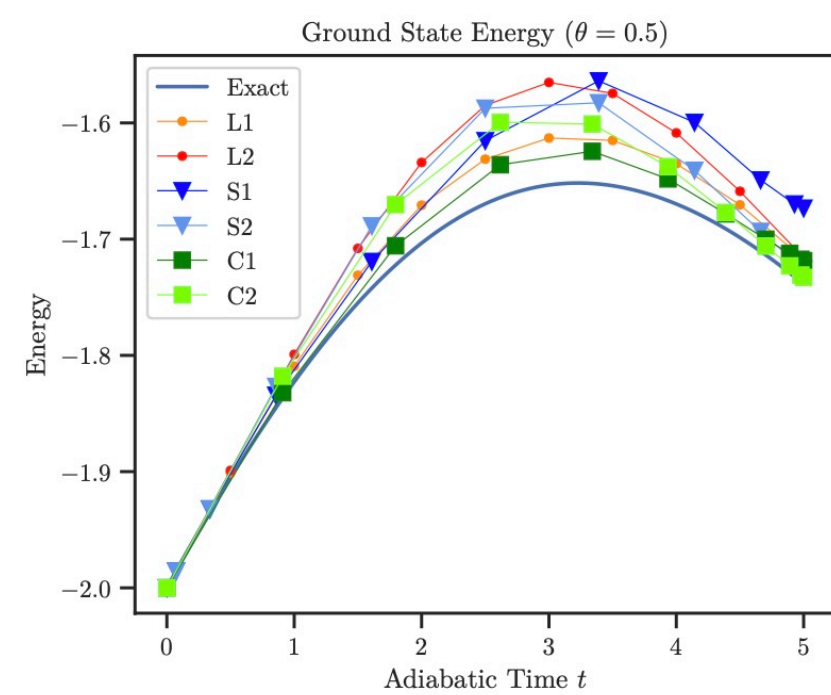
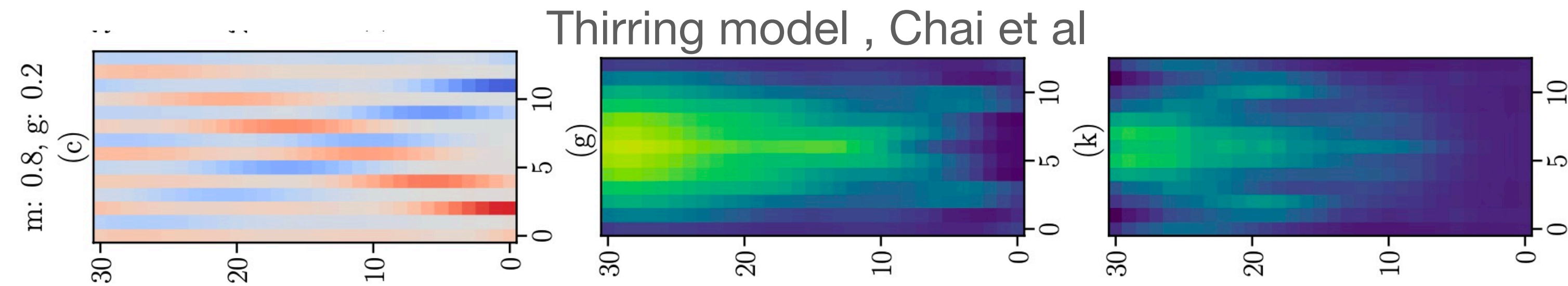
$$\Psi(k) = \mathcal{N}_\Psi \exp(-ik\mu) \exp\left(-\frac{(k-k_0)^2}{4\sigma^2}\right)$$

$$b_\Psi^\dagger = \sum_k \Psi(k) b_k^\dagger$$

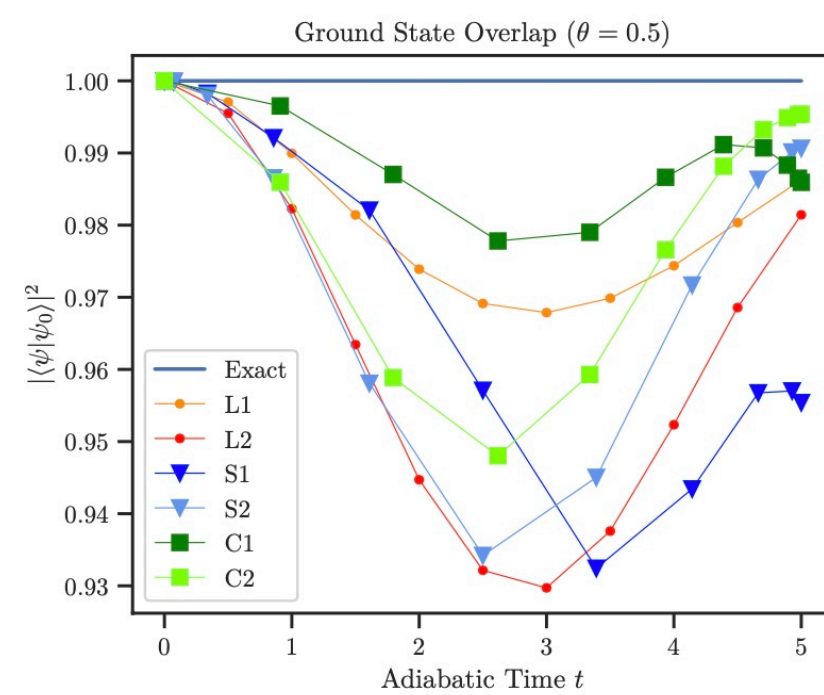
Optimized

$$\eta(p, q) = \mathcal{N}_\eta \exp\left(\frac{i(p-q)\mu_k^A}{2}\right) \exp\left(-\frac{(p-q)^2}{4\sigma_k^2}\right)$$

Wavepackets , Davoudi et al

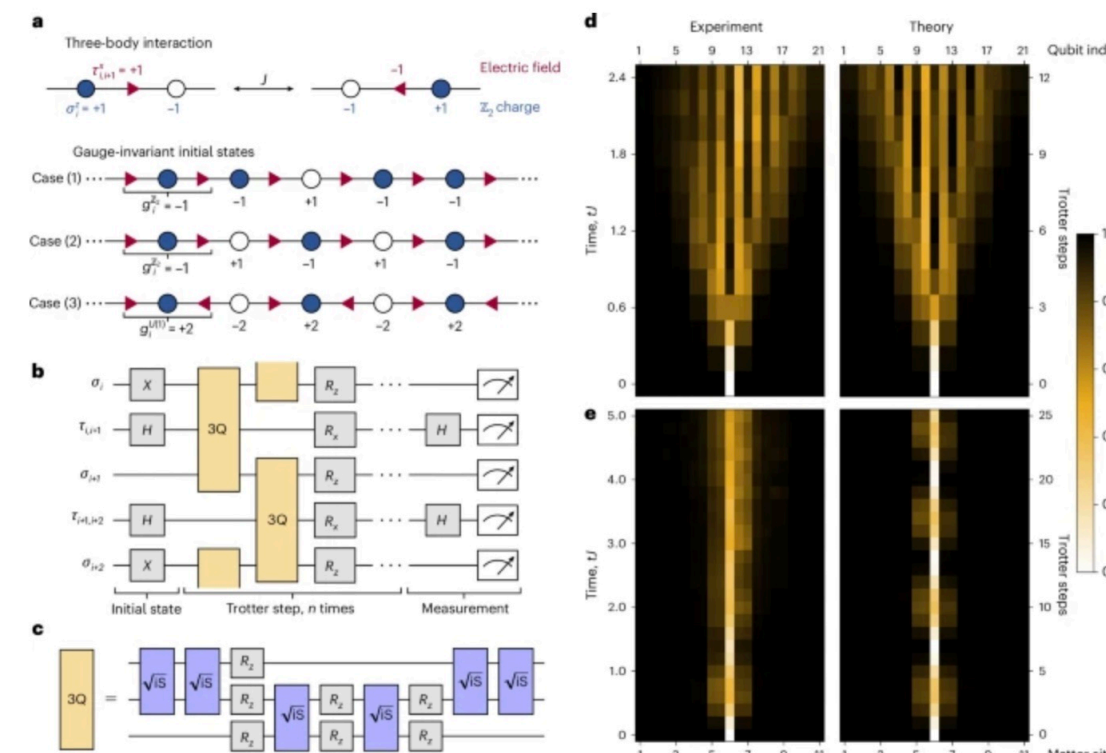


(a) QAOA, Rodeo , Pederiva et al



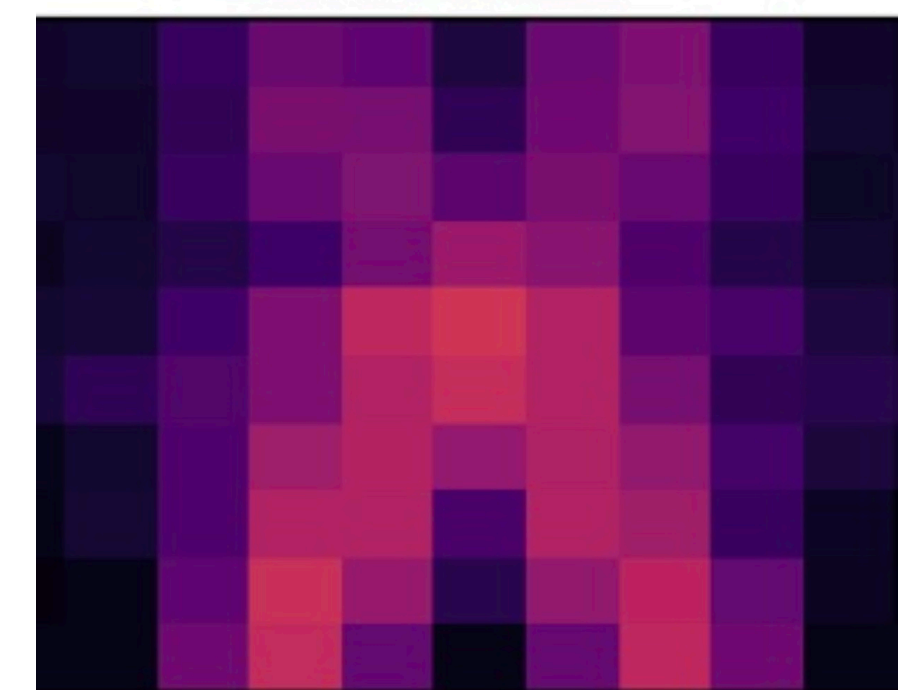
(b)

Fig. 1: Simulating confinement in a \mathbb{Z}_2 LGT on a digital superconducting chip.



Confinement in Z2, Mildenerger et al

ibm fez



Scalar FT, Zemlevskiy

Error Correction in Gauge Theories

Abelian \mathbb{Z}_n models

Quantum error correction with gauge symmetries

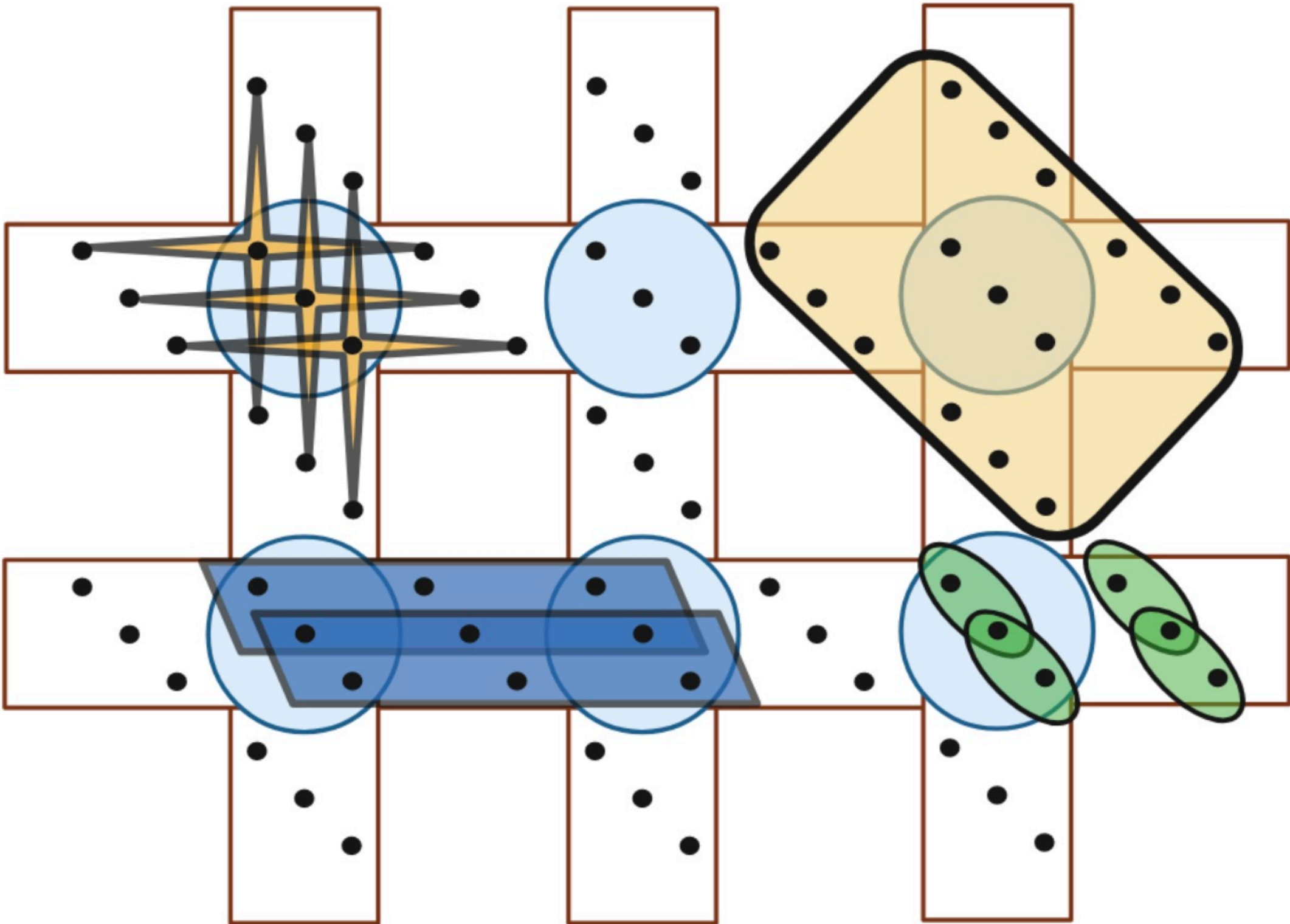
[Abhishek Rajput](#) ✉, [Alessandro Roggero](#) ✉ & [Nathan Wiebe](#) ✉

[npj Quantum Information](#) **9**, Article number: 41 (2023) | [Cite this article](#)

Fault-tolerant simulation of Lattice Gauge Theories with gauge covariant codes

Luca Spagnoli* and Alessandro Roggero†
*Dipartimento di Fisica, University of Trento, via Sommarive 14, I-38123, Povo, Trento, Italy and
INFN-TIFPA Trento Institute of Fundamental Physics and Applications, Trento, Italy*

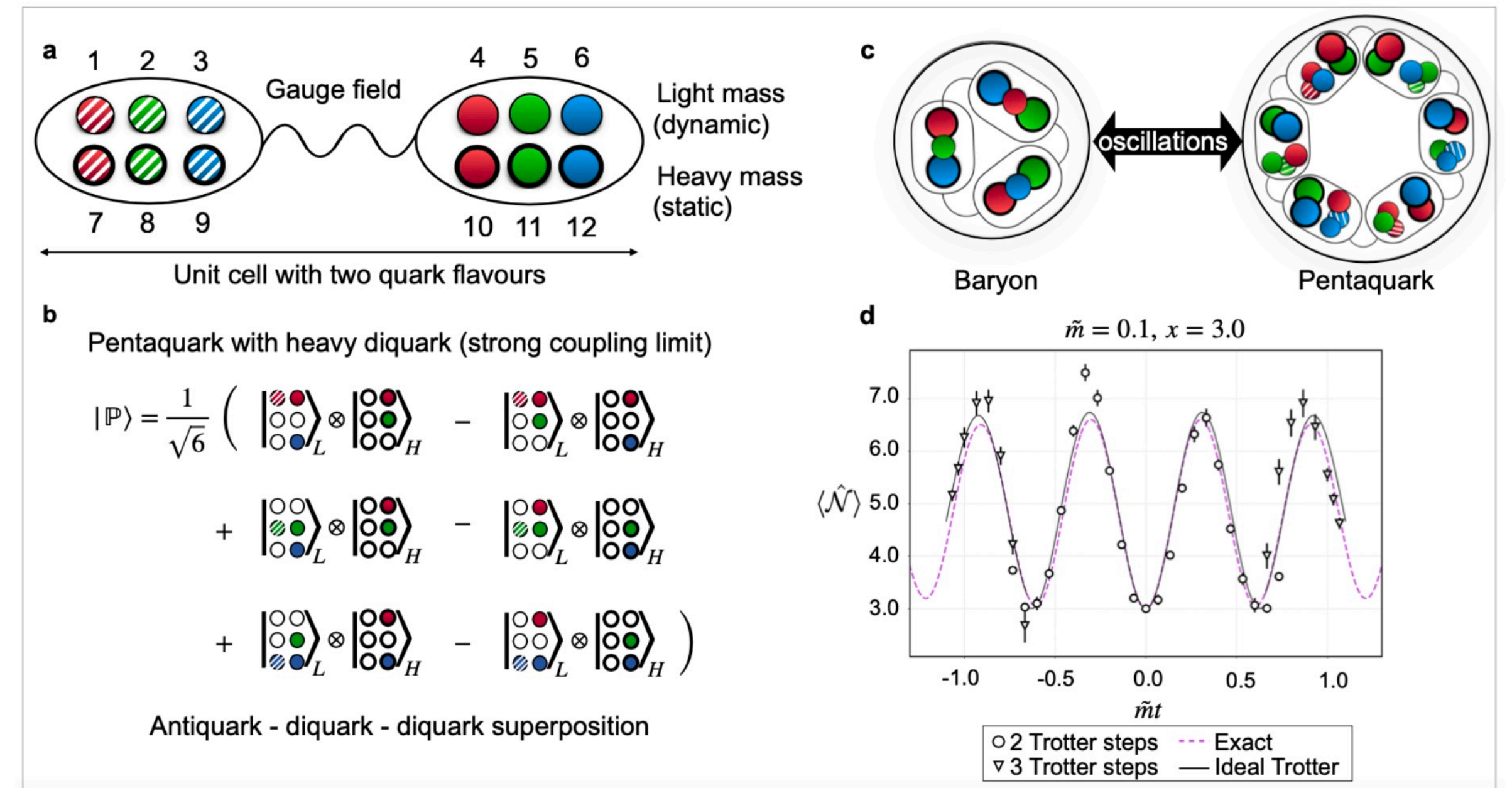
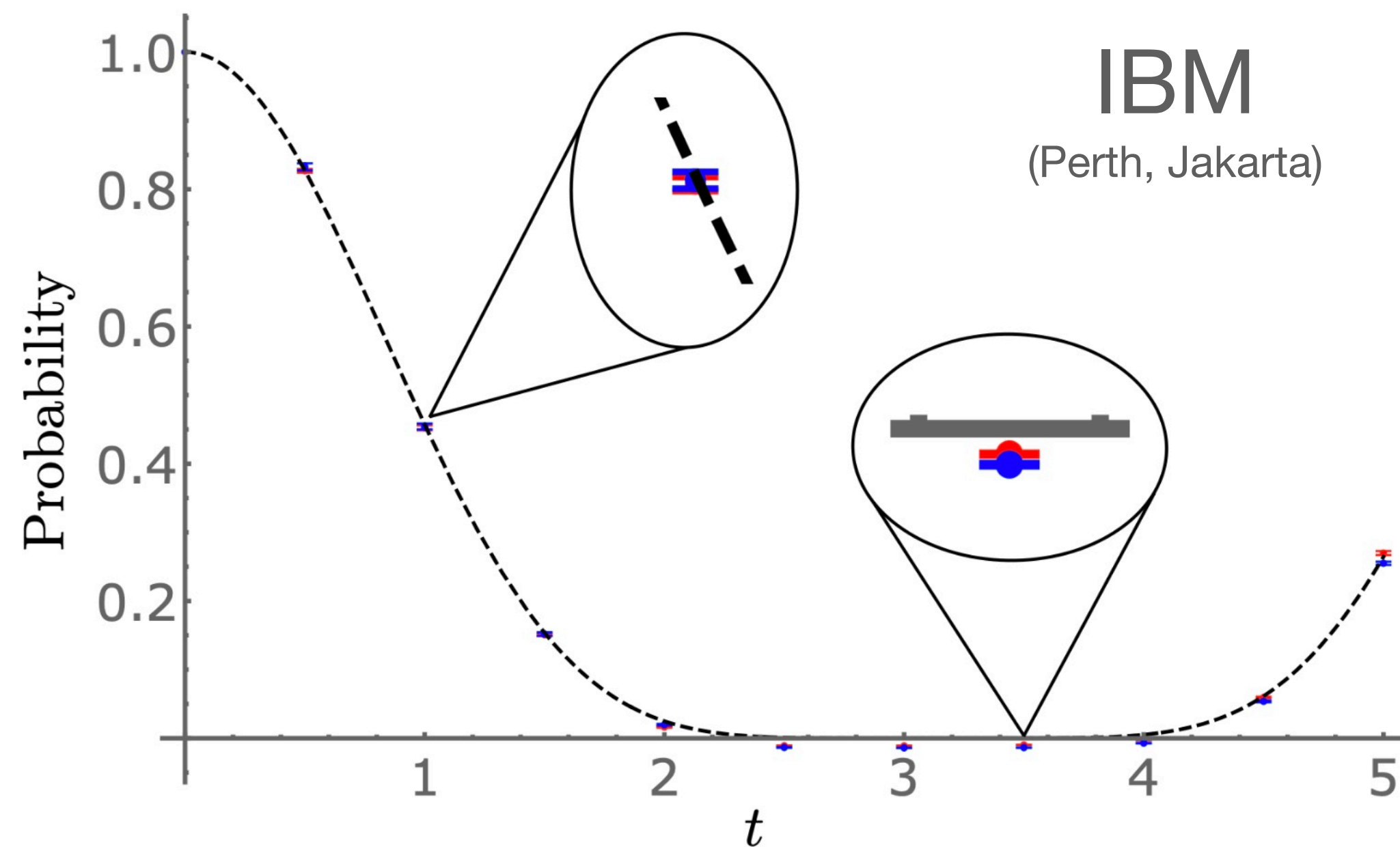
Nathan Wiebe‡
*Department of Computer Science, University of Toronto, Toronto, ON M5S 2E4, Canada
Pacific Northwest National Laboratory, Richland, WA 99354, USA and
Department of Physics, University of Washington, Seattle, WA 98195, USA*
(Dated: October 7, 2024)



Gauss’s law  Stabilizers

1+1D Quantum Chromodynamics

Trivial Vacuum-to-Vacuum



Preparations for quantum simulations of quantum chromodynamics in 1 + 1 dimensions. I. Axial gauge

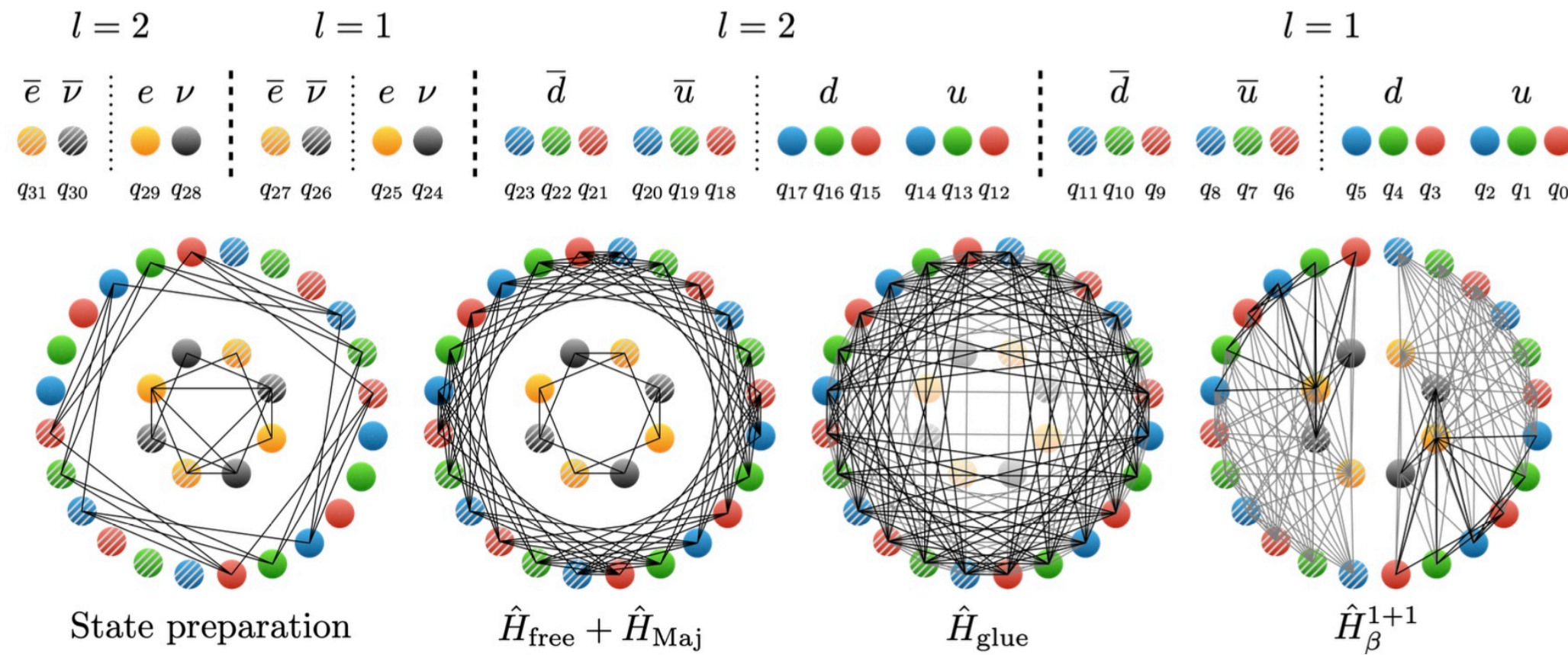
Roland C. Farrell, Ivan A. Chernyshev, Sarah J. M. Powell, Nikita A. Zemlevskiy, Marc Illa, and Martin J. Savage
Phys. Rev. D **107**, 054512 – Published 30 March 2023

Simulating one-dimensional quantum chromodynamics on a quantum computer:
Real-time evolutions of tetra- and pentaquarks

Yasar Y. Atas*,^{1,2,†} Jan F. Haase*,^{1,2,3,‡} Jinglei Zhang,^{1,2,§} Victor Wei,^{1,4}
Sieglinde M.-L. Pfaendler,⁵ Randy Lewis,⁶ and Christine A. Muschik^{1,2,7}

SU(3) and higher is quite different from U(1) and SU(2)

Ivan A. Chernyshev ¹ Roland C. Farrell ² Marc Illa ³ Martin J. Savage ^{3,*} Andrii Maksymov ⁴ Felix Tripier ⁴ Miguel Angel Lopez-Ruiz ⁴ Andrew Arrasmith ⁴ Yvette de Sereville ⁴ Aharon Brodutch ⁴ Claudio Grotto ⁴ Ananth Kaushik ⁴ and Martin Roetteler ^{4,†}

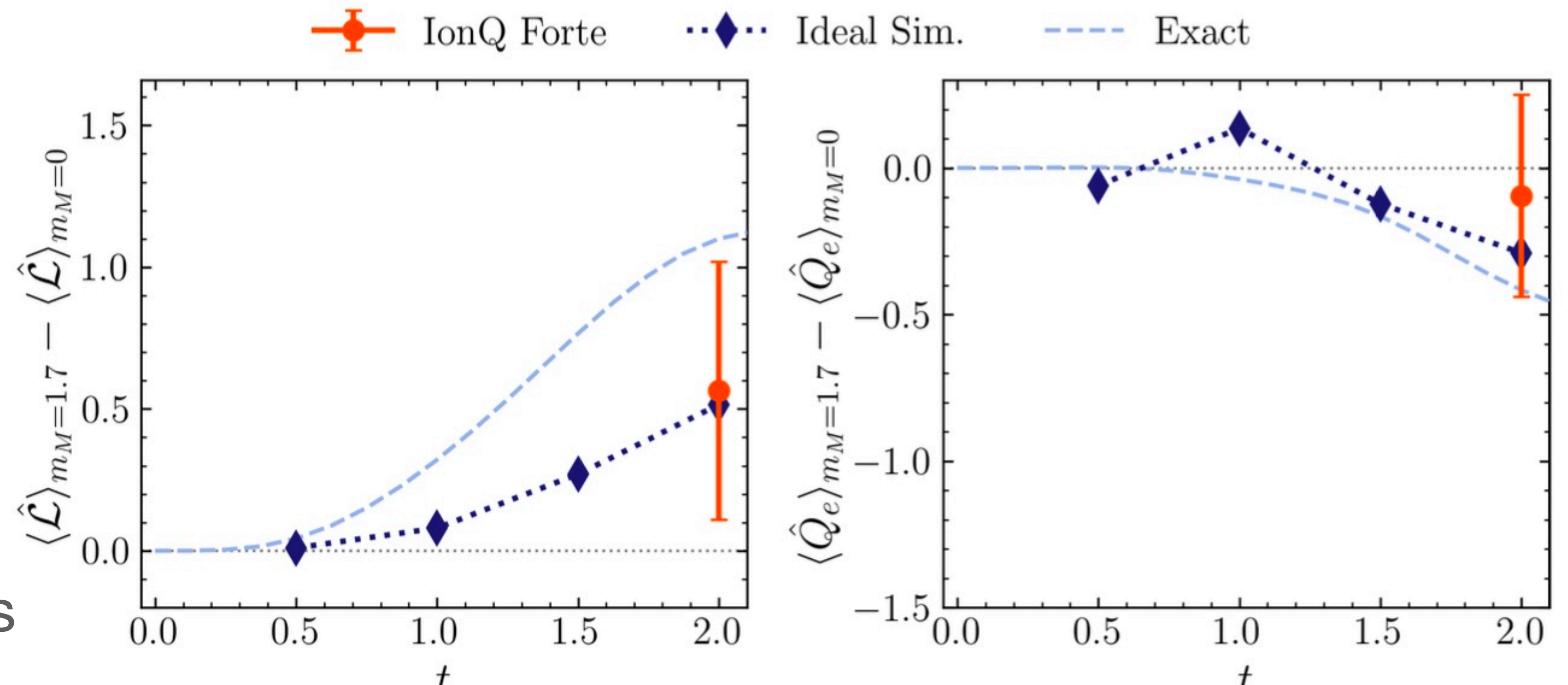


$$\hat{H}_{\beta}^{1+1} = \frac{G}{\sqrt{2}} \sum_{\text{spatial sites}} (\bar{\psi}_u \gamma^{\mu} \psi_d \bar{\psi}_e \gamma_{\mu} \mathcal{C} \psi_{\nu} + \text{h.c.})$$

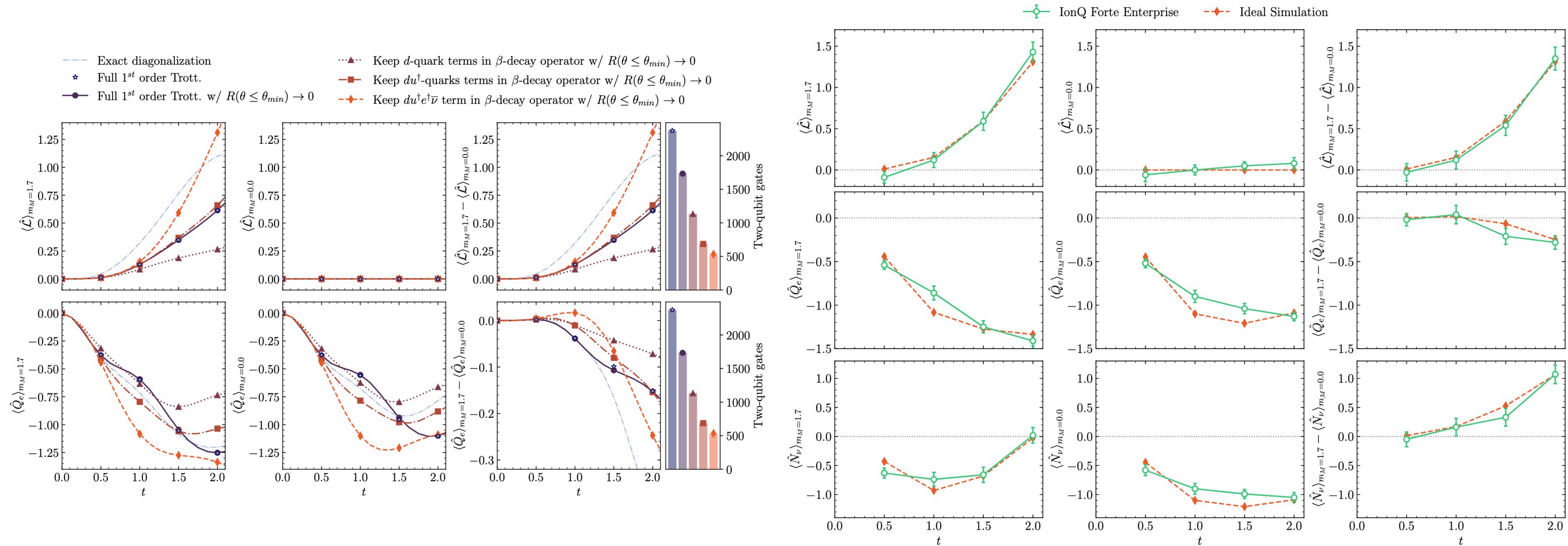
$$\approx \frac{G}{\sqrt{2}} \sum_{n \text{ even}} \left(\phi_n^{(u)\dagger} \phi_n^{(d)} \phi_n^{(e)\dagger} \phi_{n+1}^{(\nu)} + \text{h.c.} \right),$$

$$\hat{H}_{\text{Maj}} = \frac{1}{2} m_M \sum_{n \text{ even}} \left(\phi_n^{(\nu)} \phi_{n+1}^{(\nu)} + \text{h.c.} \right)$$

>2300 entangling gates



Ivan A. Chernyshev ¹ Roland C. Farrell ² Marc Illa ³ Martin J. Savage ^{3,*} Andrii Maksymov ⁴ Felix Tripier ⁴ Miguel Angel Lopez-Ruiz ⁴ Andrew Arrasmith ⁴ Yvette de Sereville ⁴ Aharon Brodutch ⁴ Claudio Grotto ⁴ Ananth Kaushik ⁴ and Martin Roetteler ^{4,†}



>450 entangling gates

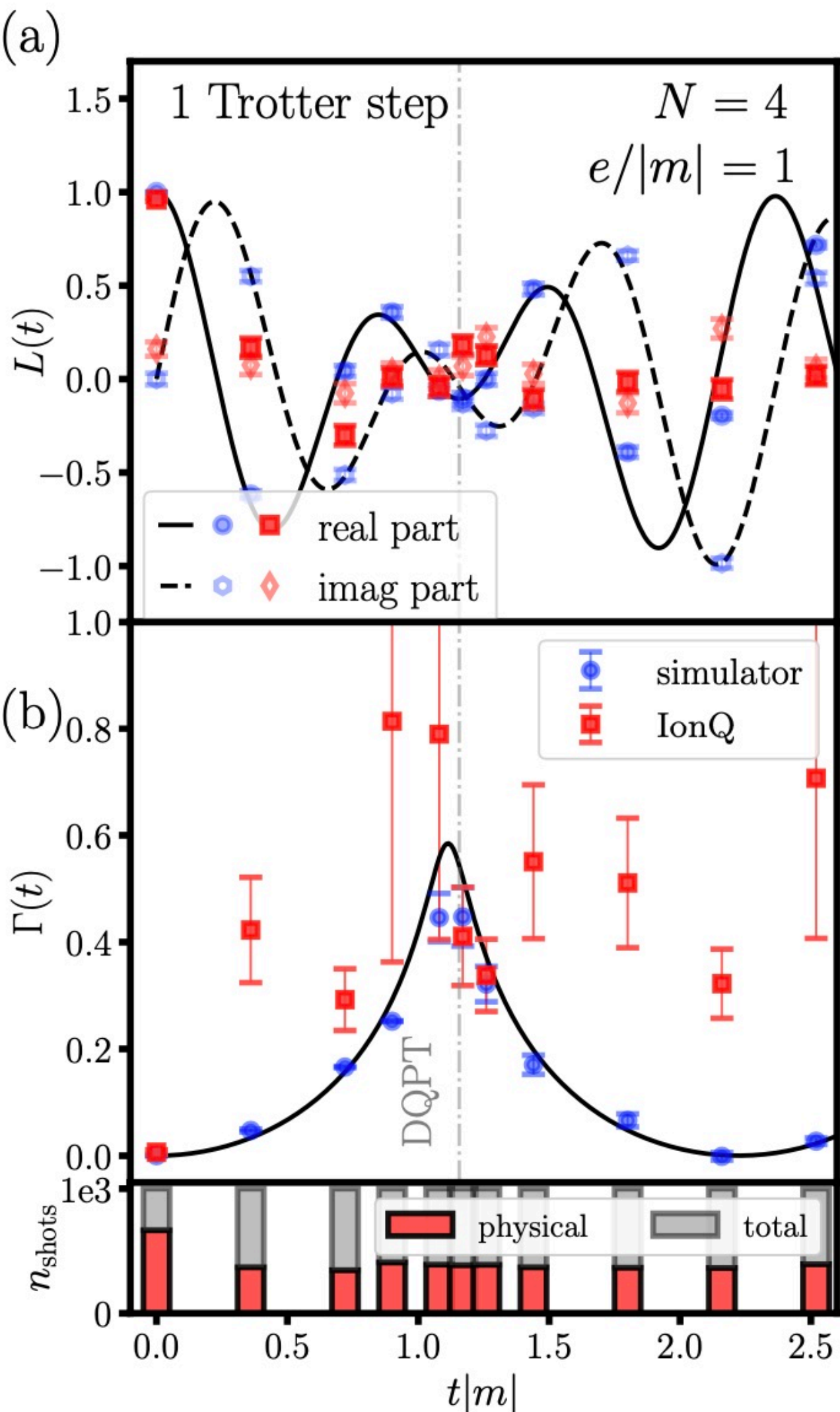
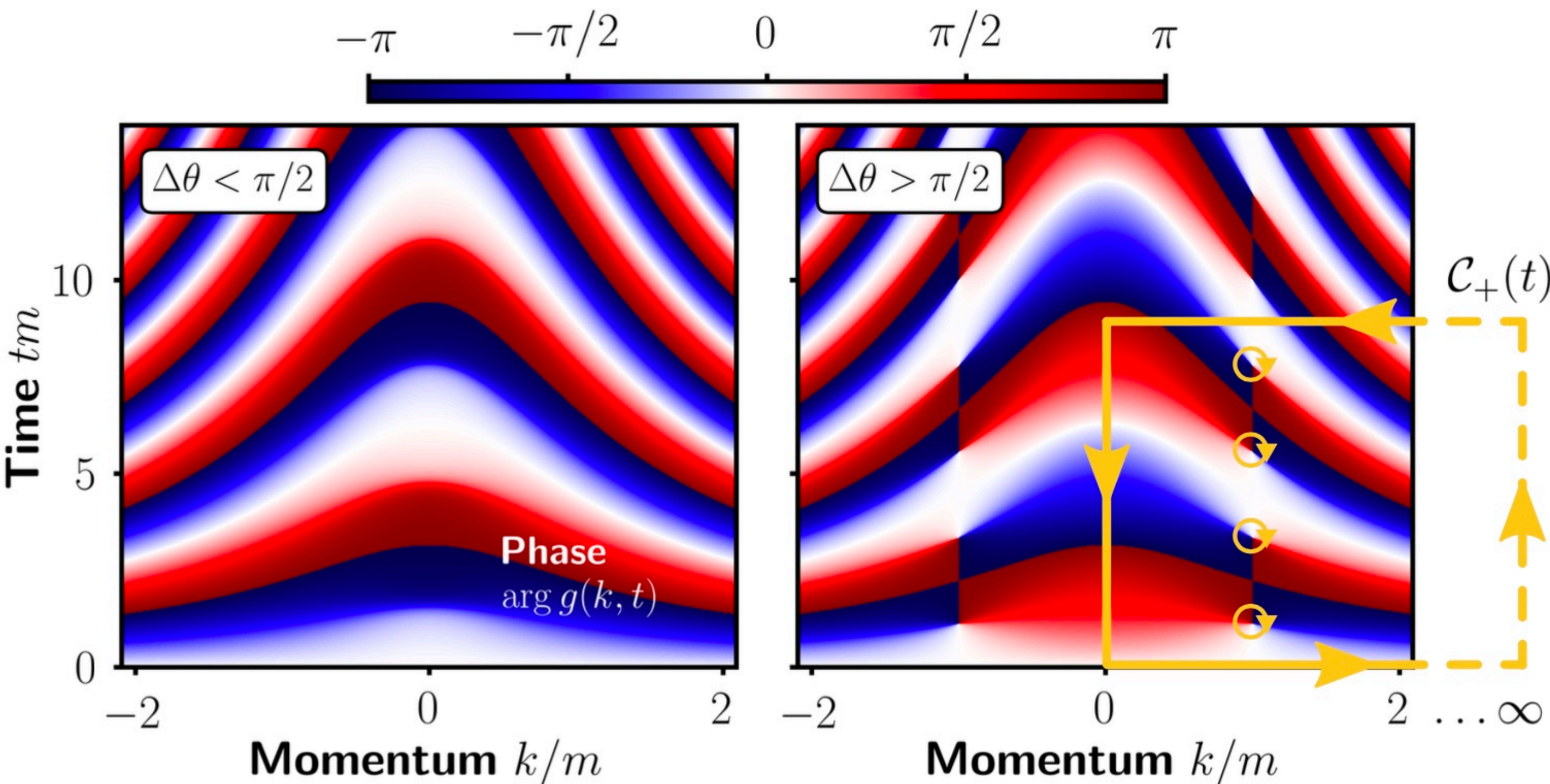
Dynamical Quantum Phase Transitions

Dynamical topological transitions in the massive Schwinger model with a θ -term

T. V. Zache,^{1,*} N. Mueller,² J. T. Schneider,¹ F. Jendrzejewski,³ J. Berges,¹ and P. Hauke^{1,3}

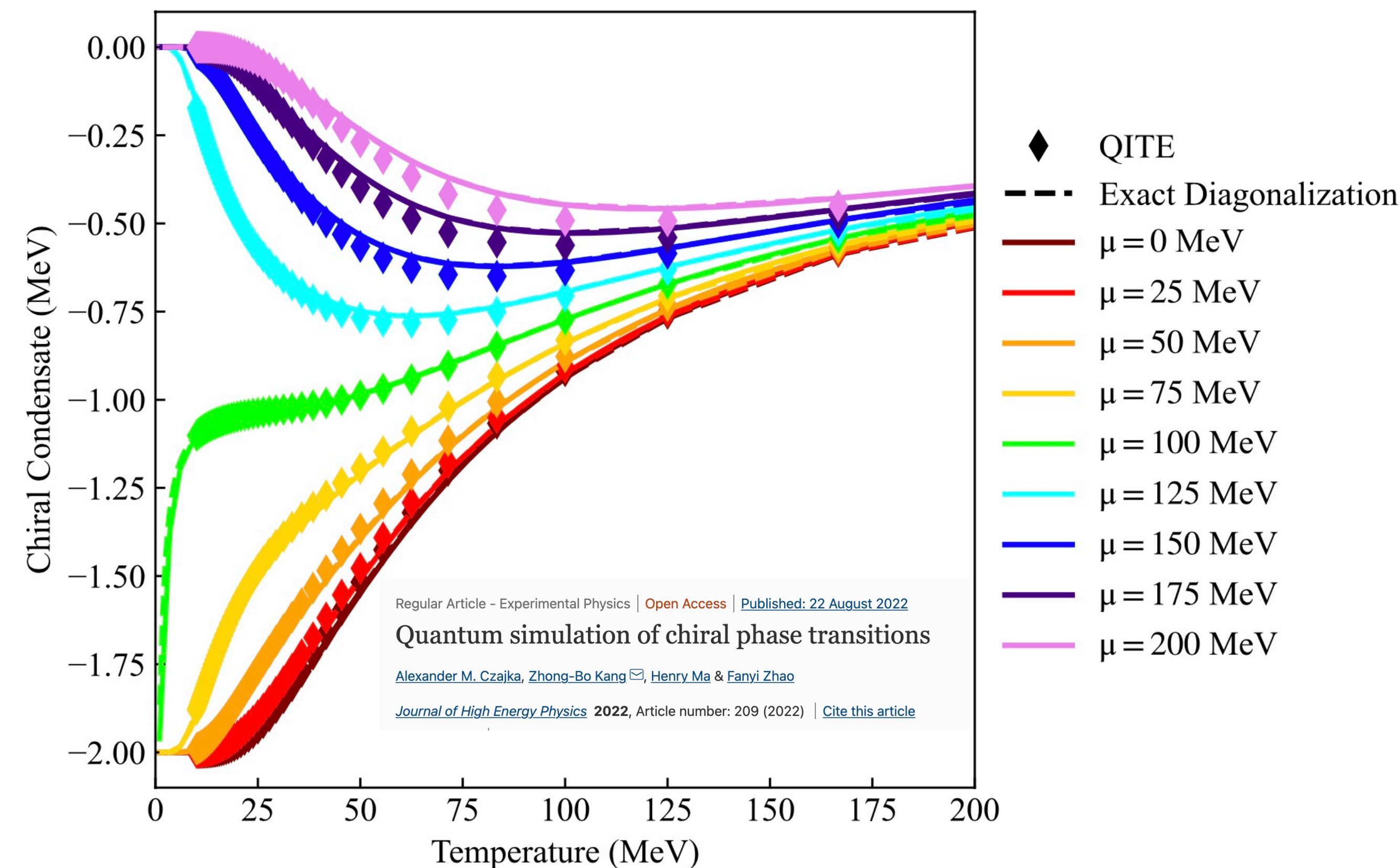
Quantum computation of dynamical quantum phase transitions and entanglement tomography in a lattice gauge theory

Niklas Mueller,^{1,2,3,*} Joseph A. Carolan,⁴ Andrew Connelly,⁵
Zohreh Davoudi,^{1,6,†} Eugene F. Dumitrescu,^{7,‡} and Kübra Yeter-Aydeniz⁸

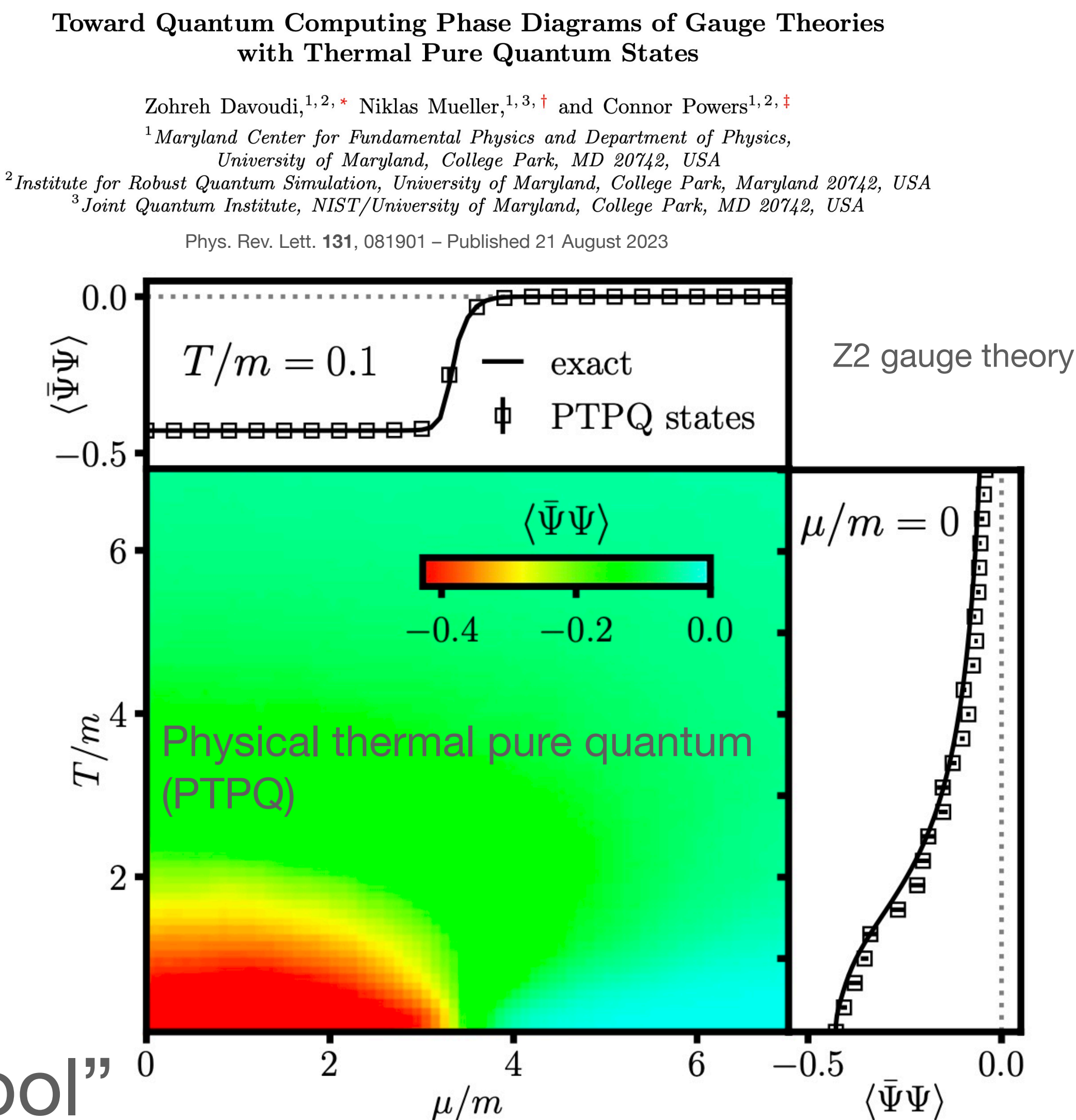


Modeling the QCD Phase Diagram

$$\mathcal{H} = \bar{\psi}(i\gamma_1\partial_1 + m)\psi - g(\bar{\psi}\psi)^2 - \mu\bar{\psi}\gamma_0\psi$$

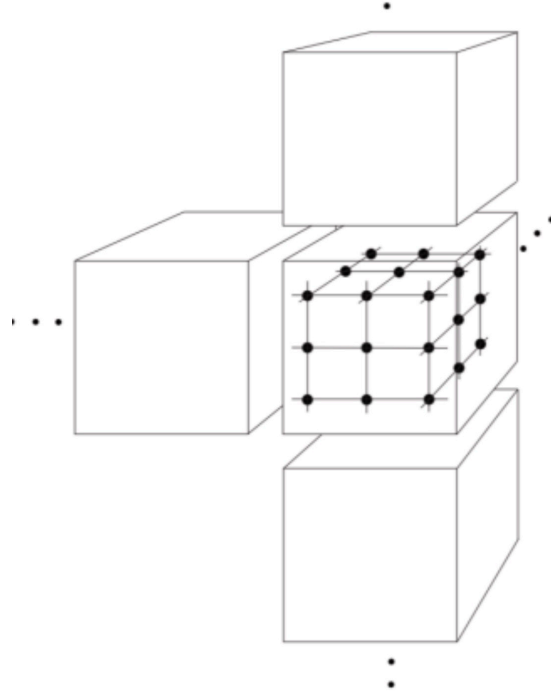


QITE algorithms to “cool”



Dynamical Gauge Fields - Yang-Mills

Byrnes-Yamamoto — Kogut-Susskind

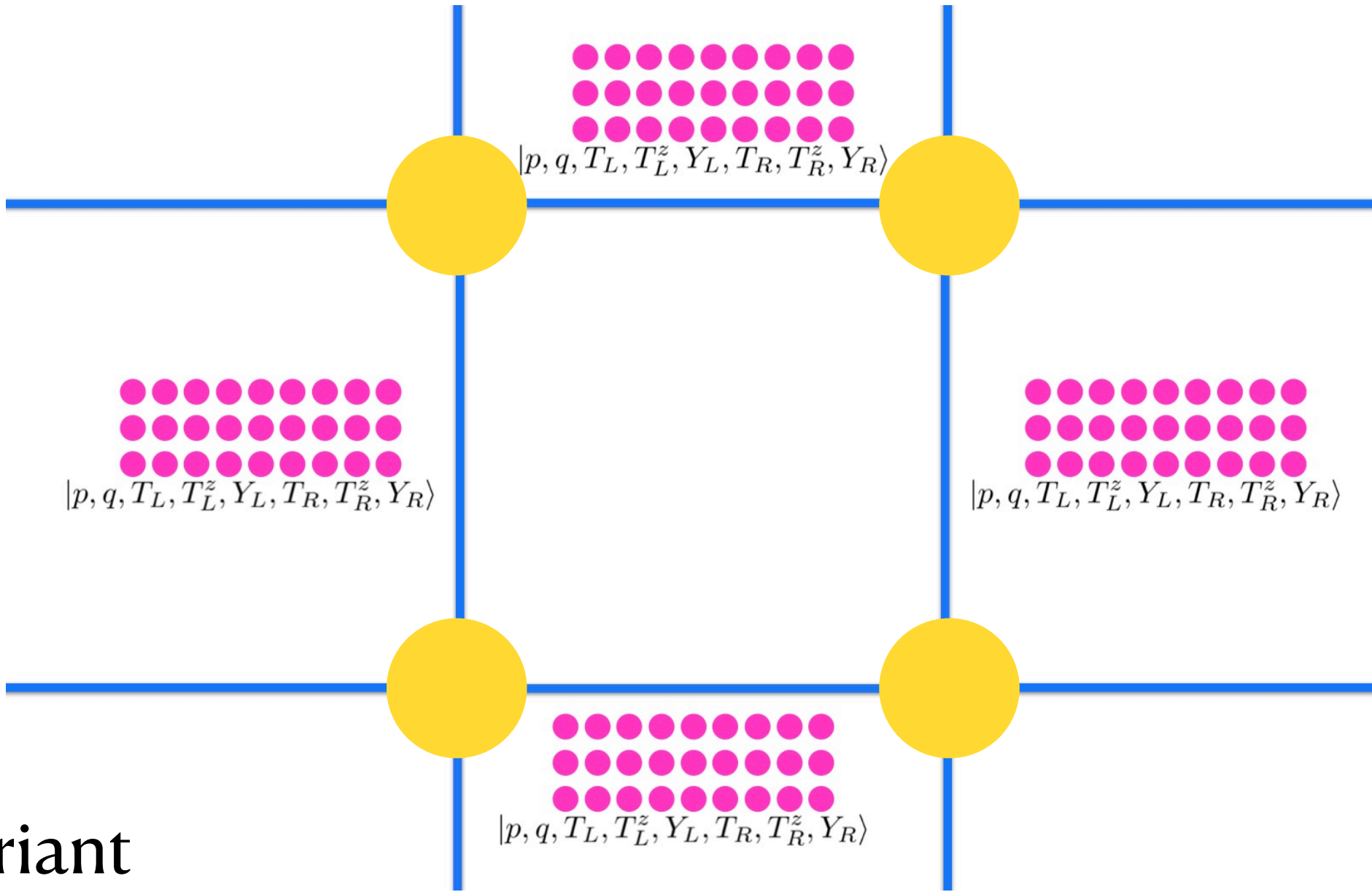
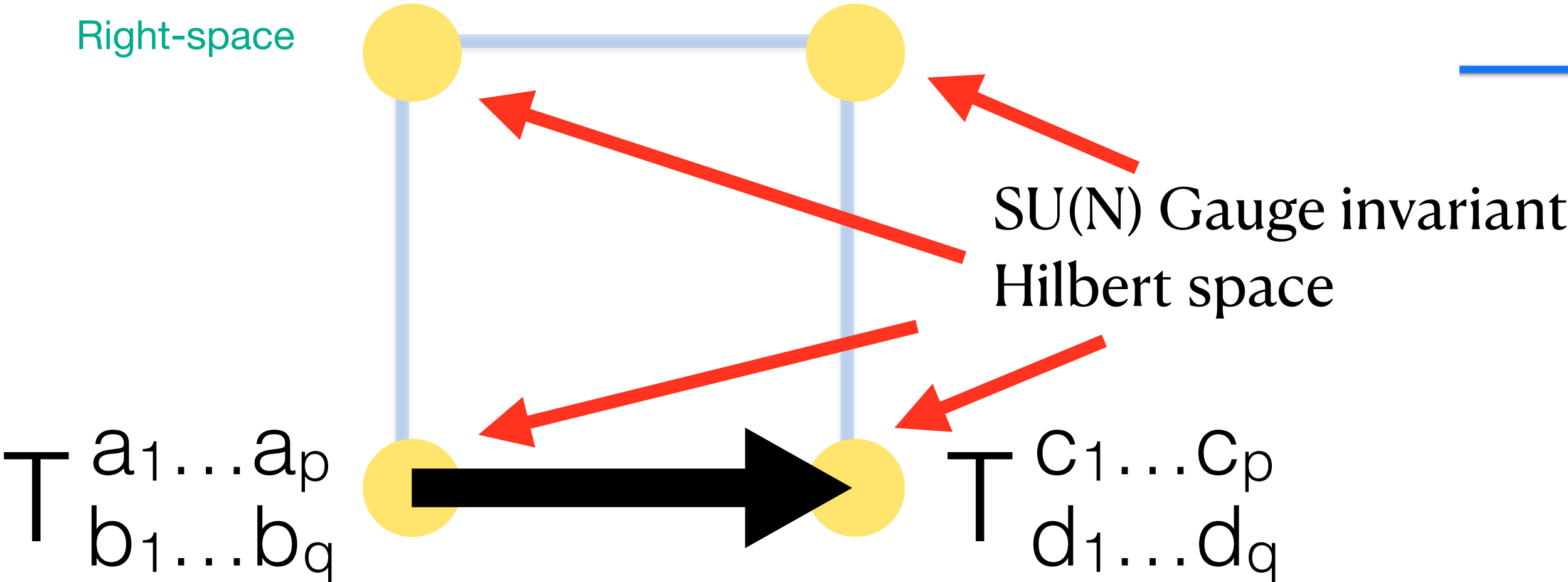


Many ways to map/distribute the field(s) in the UV (lattice spacing)
 Consider the Kogut-Susskind basis = electric basis

$$\hat{H} = \frac{g^2}{2} \sum_{\text{links}} \hat{E}^2 - \frac{1}{2g^2} \sum_{\square} \left(\hat{\square} + \hat{\square}^\dagger \right)$$

Electric Field Casimir operator
 $|p, q, T_L, T_L^z, Y_L, T_R, T_R^z, Y_R\rangle$
 Irrep Left-space Right-space

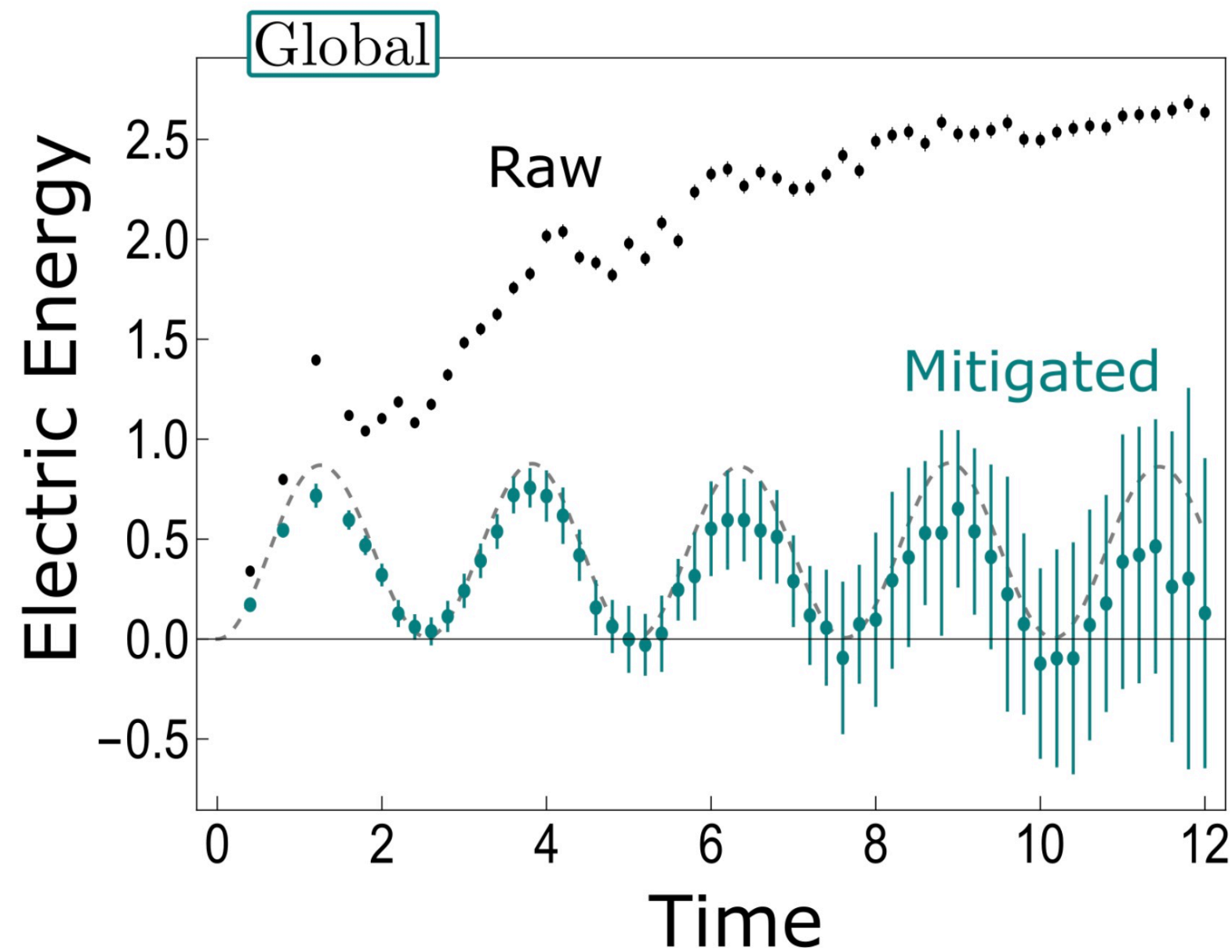
Magnetic Field operator
 Off-diagonal on electric basis



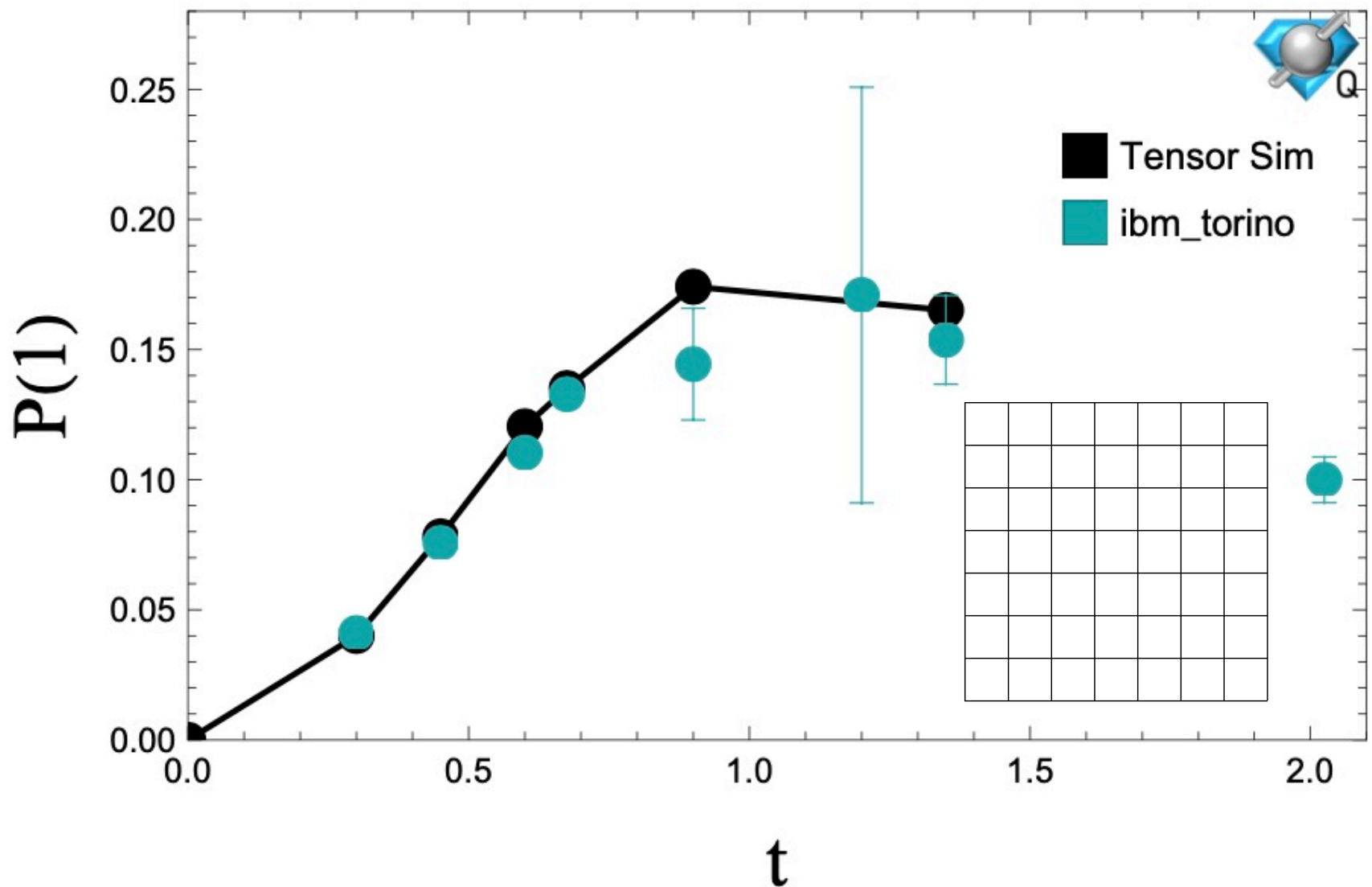
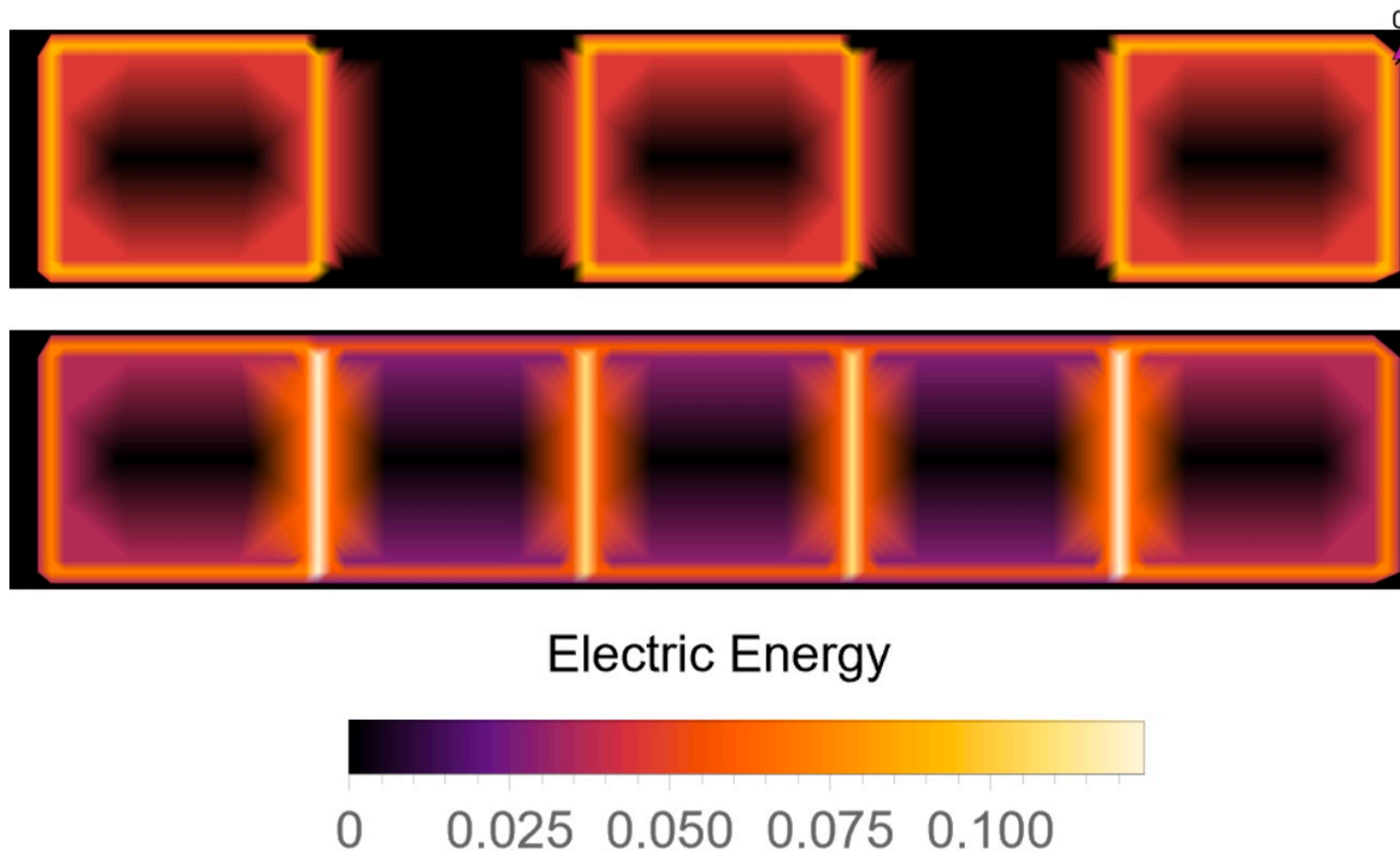
Truncations in irrep space !!!!!

SU(3) Yang-Mills Plaquettes

IBM



One Plaquette



A Trailhead for Quantum Simulation of SU(3) Yang-Mills Lattice Gauge Theory in the Local Multiplet Basis

Anthony Ciavarella^{1,*} Natalie Klco^{2,†} and Martin J. Savage^{1,‡}

Quantum Circuits for SU(3) Lattice Gauge Theory

Praveen Balaji¹, Cianán Conefrey-Shinozaki¹, Patrick Draper^{*1}, Jason K. Elhaderi¹, Drishti Gupta¹, Luis Hidalgo¹, Andrew Lytle¹, and Enrico Rinaldi²

Preparation of the SU(3) Lattice Yang-Mills Vacuum with Variational Quantum Methods

Anthony N Ciavarella^{1,*} and Ivan A Chernyshev^{1,†}

Quantum Simulation of SU(3) Lattice Yang Mills Theory at Leading Order in Large

Anthony N. Ciavarella^{1,*} and Christian W. Bauer^{1,2,†}

Hot Starts

Perturbative Expansions
and
Effective Field Theories

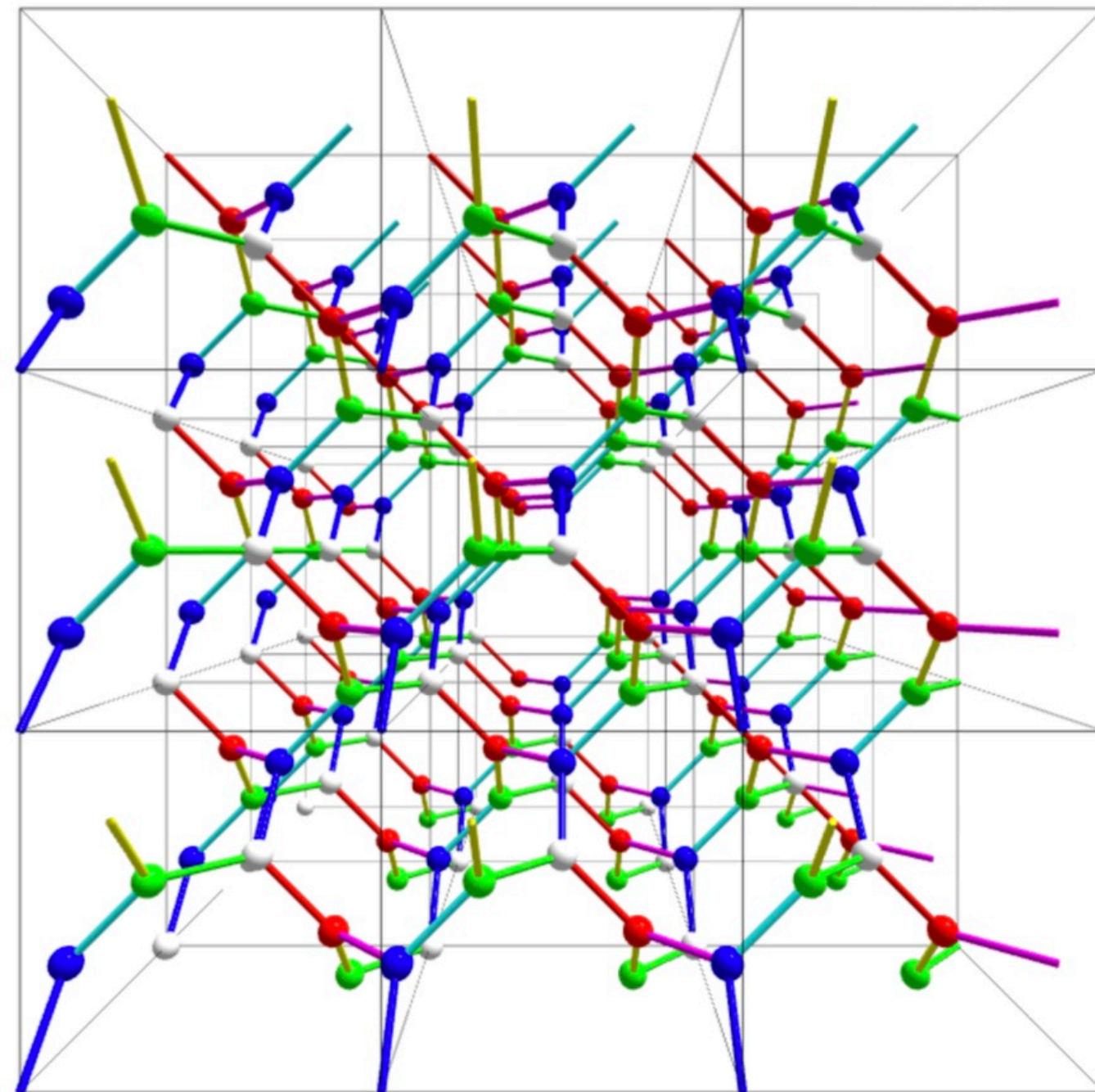
Reducing Group Space Overheads in non-Abelian LGT

reducing the number of links per vertex

From square plaquettes to triamond lattices for $SU(2)$ gauge theory

Ali H. Z. Kavaki^{*} and Randy Lewis[†]

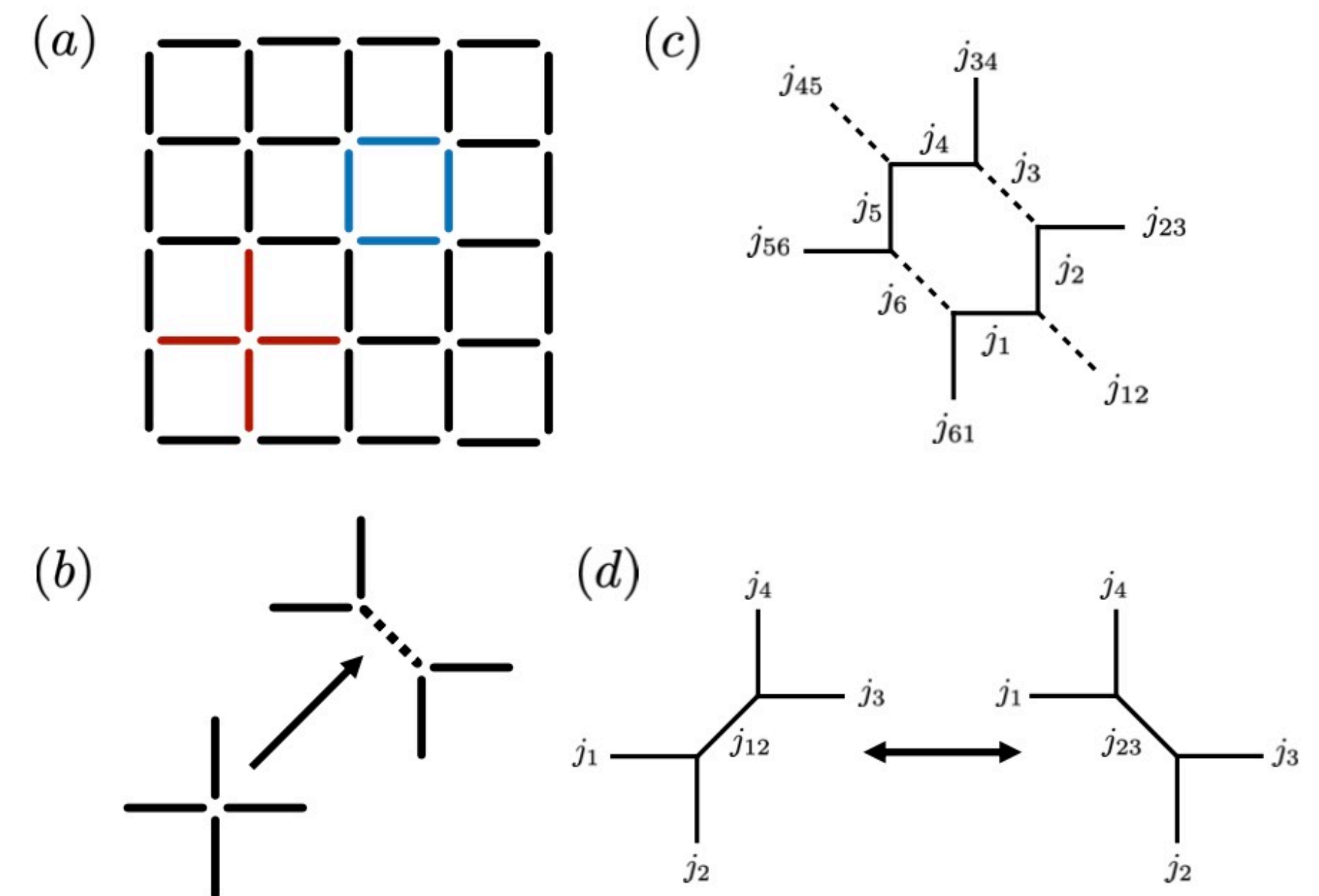
2024



Quantum and classical spin network algorithms for q -deformed Kogut-Susskind gauge theories

Torsten V. Zache,^{*} Daniel González-Cuadra, and Peter Zoller

2023





Transport Properties

Shear Viscosity in 2+1D SU(2)

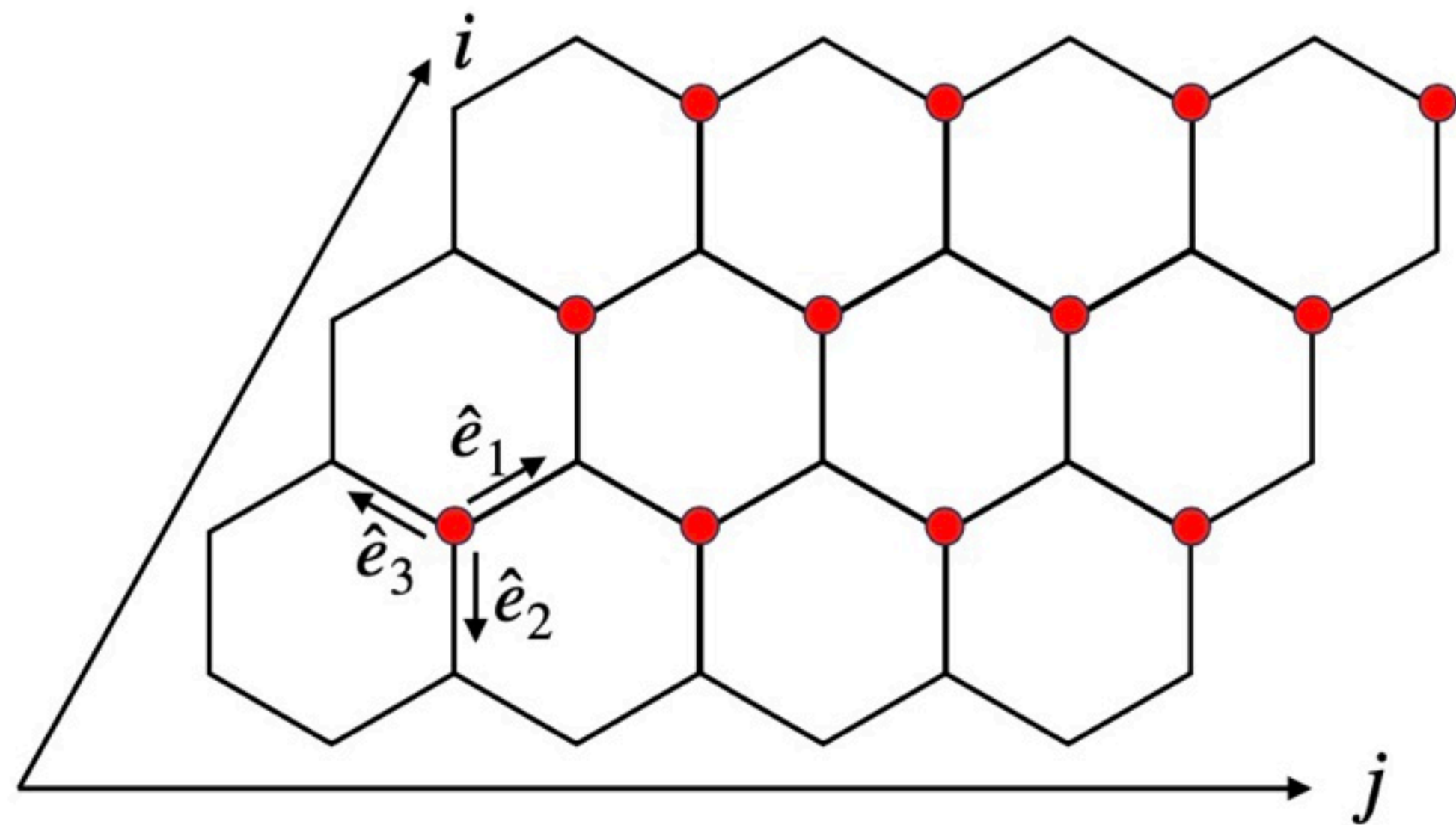
Editors' Suggestion

Open Access

Classical and quantum computing of shear viscosity for $(2 + 1)D$ SU(2) gauge theory

Berndt Mueller and Xiaojun Yao

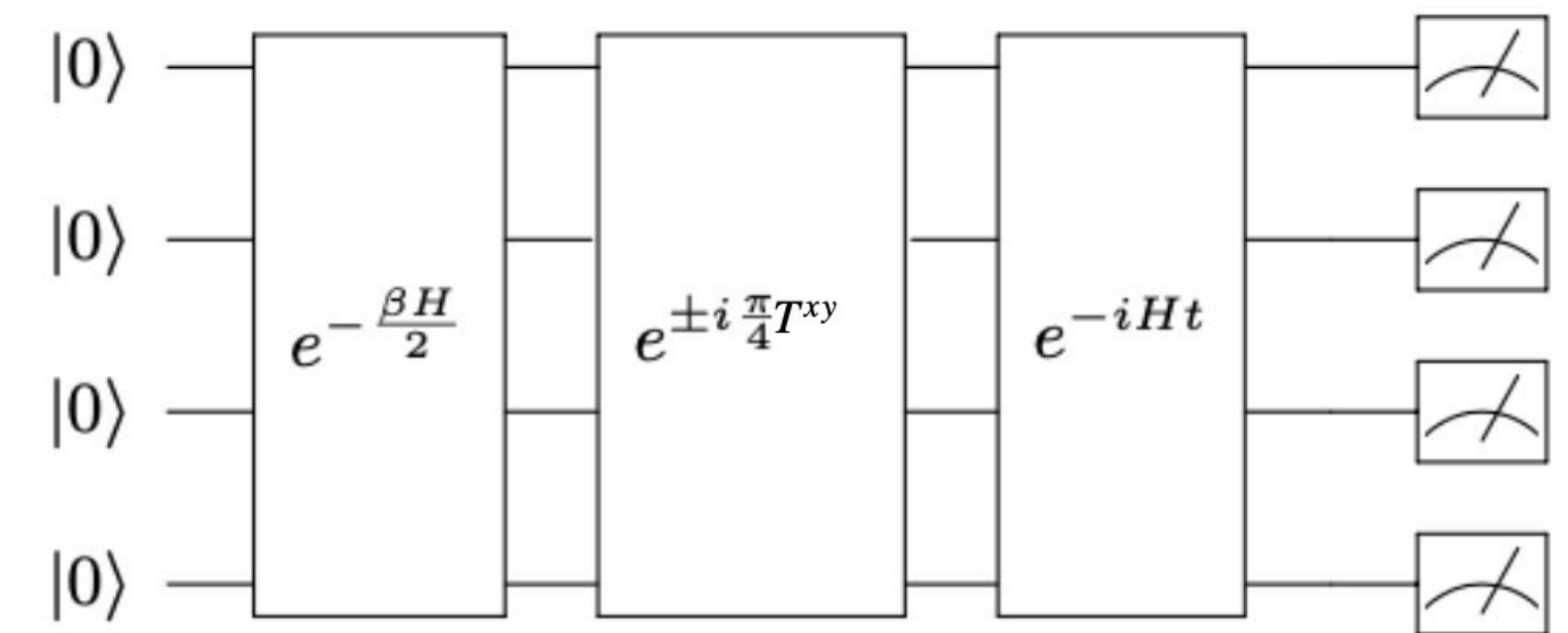
Francesco Turro, Anthony Ciavarella, and Xiaojun Yao
Phys. Rev. D **109**, 114511 – Published 13 June 2024



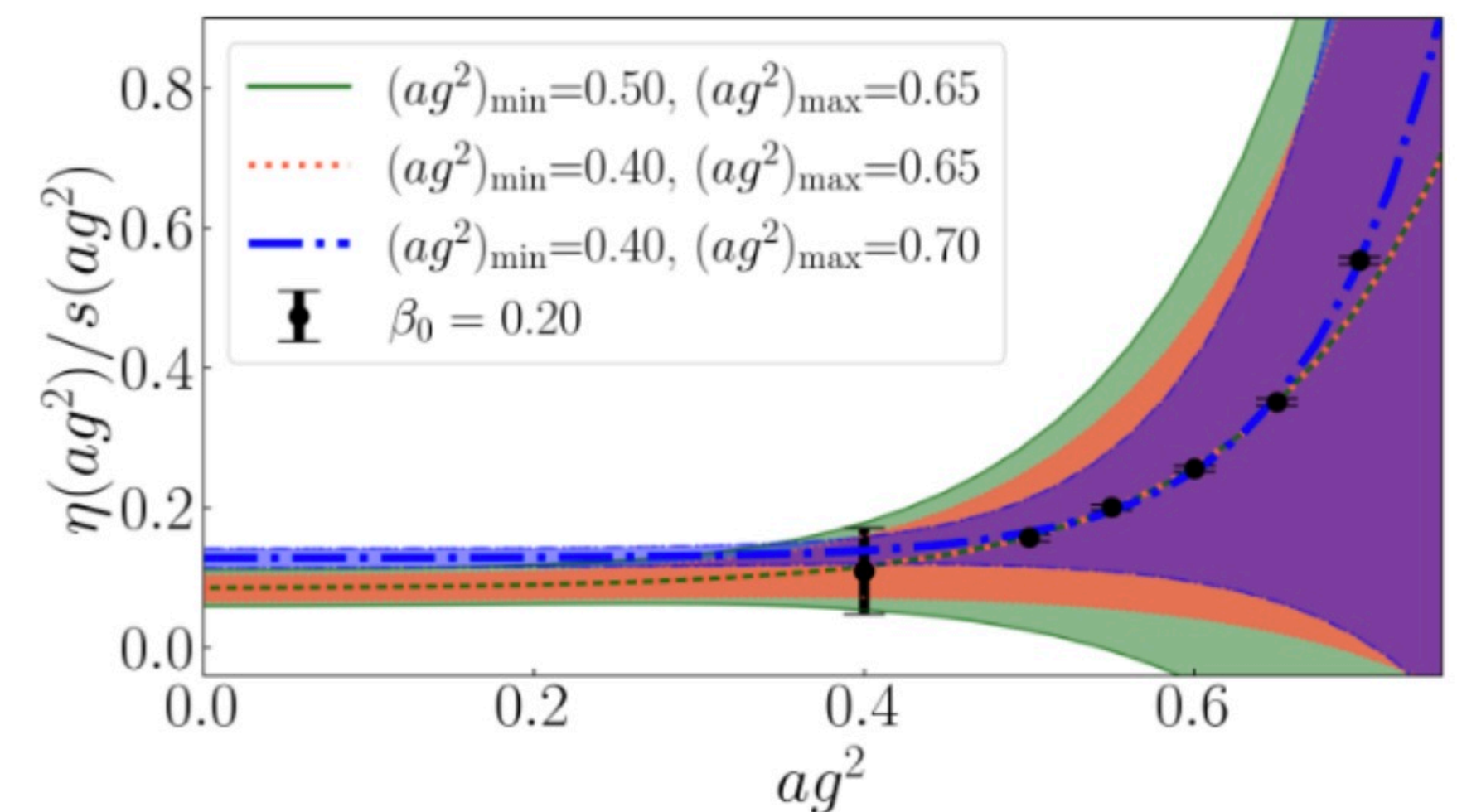
$$H = \frac{3\sqrt{3}g^2}{4} \sum_{\text{links}} E_i^a E_i^a - \frac{4\sqrt{3}}{9g^2 a^2} \sum_{\text{plaqs}} \text{Hexagon}$$

$$T^{xy} = -\frac{g^2}{\sqrt{3}a^2} ((E_1^a)^2 - (E_3^a)^2)$$

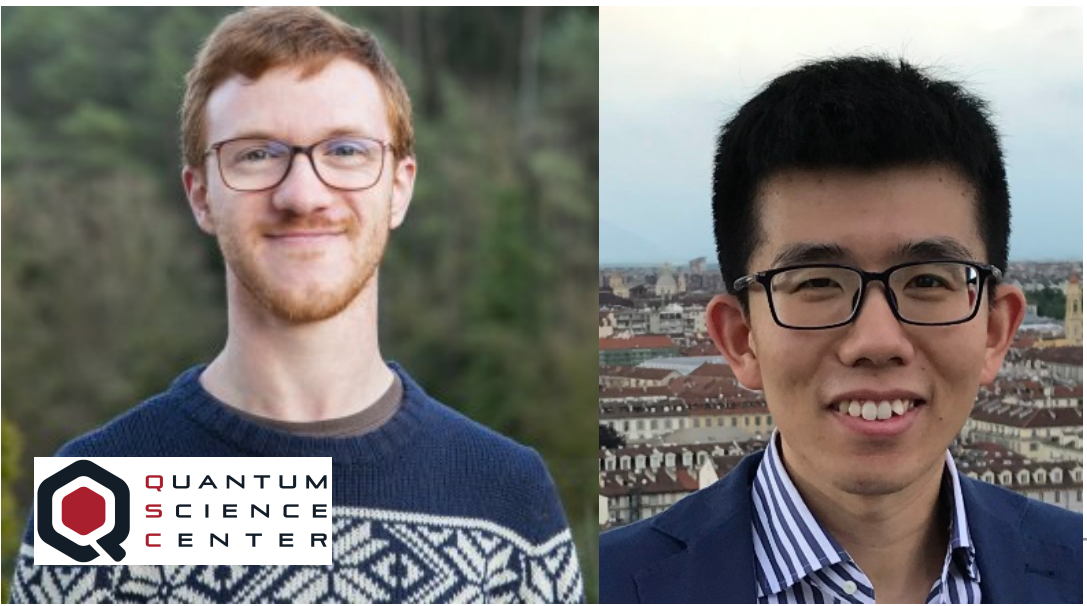
Quantum algorithm for G_r^{xy}



On 4×4 lattice w/ $j_{\text{max}} = 0.5$



At the Quantum Limit, same as liquid created in heavy-ion collisions



Improved Hamiltonian for Honeycomb lattices non-Abelian LGTs in 2+1

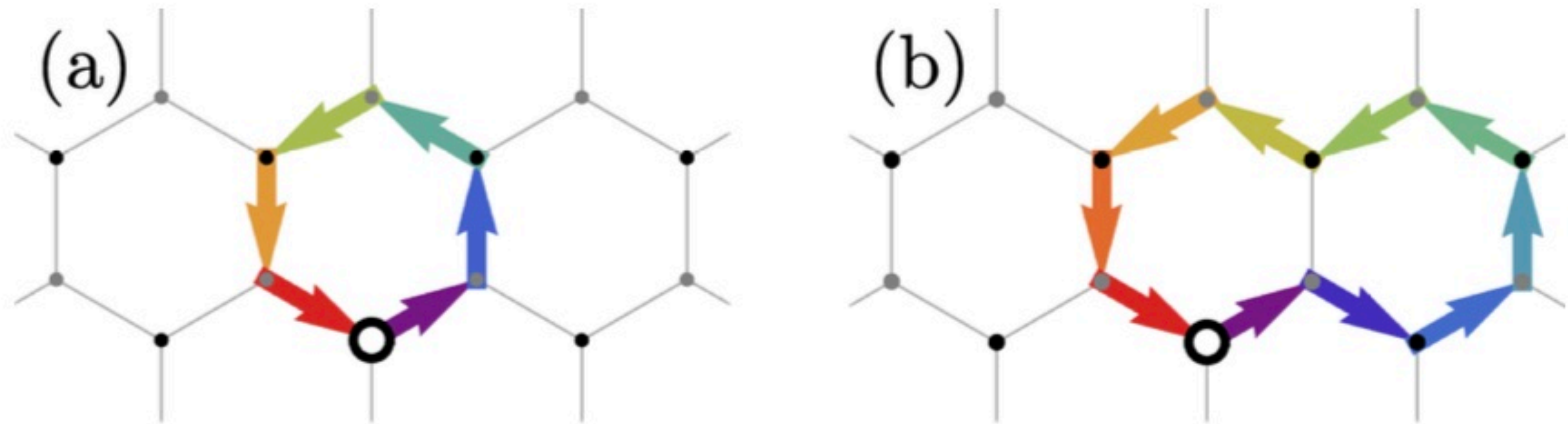
Improved honeycomb and hyperhoneycomb lattice
Hamiltonians for quantum simulations of non-Abelian
gauge theories

Marc Illa^{*}, Martin J. Savage^{†,‡}, and Xiaojun Yao[§]

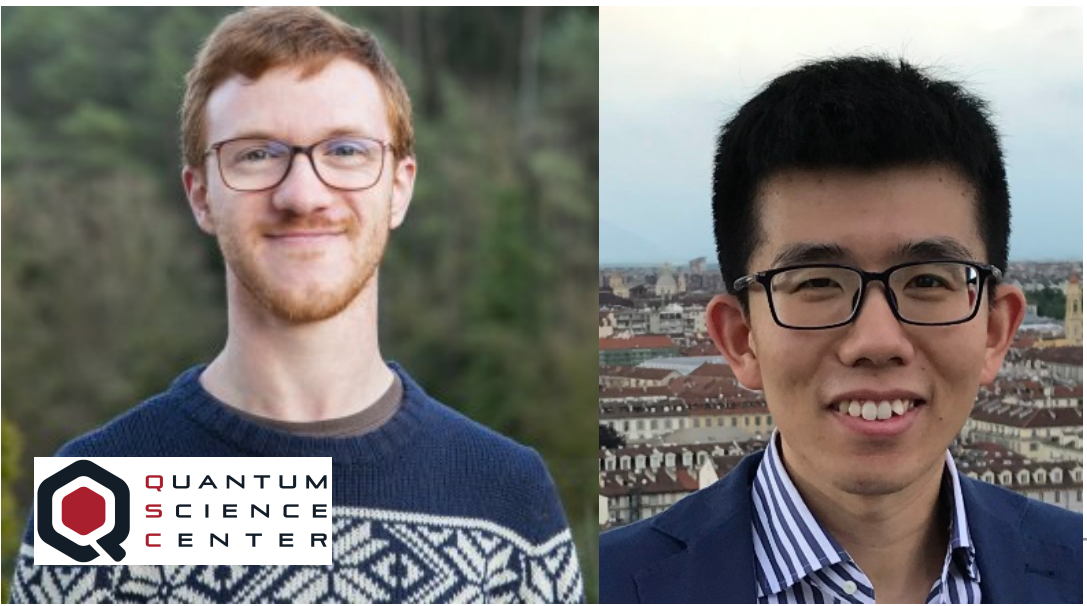
Show more

Phys. Rev. D **111**, 114520 – Published 26 June, 2025

DOI: <https://doi.org/10.1103/3rwf-f844>



$$\Gamma_{6,10}^{(\text{HC})} = \frac{14}{9}\Gamma_6^{(\text{HC})} - \frac{5}{36}\Gamma_{10}^{(\text{HC})} = \frac{g^2}{4}S_{\text{hex}}^2 \sum_a [B^a(\boldsymbol{x})]^2 + \mathcal{O}(b^7)$$



Improved Hamiltonian for Hyper-Honeycomb lattices non-Abelian LGTs in 3+1

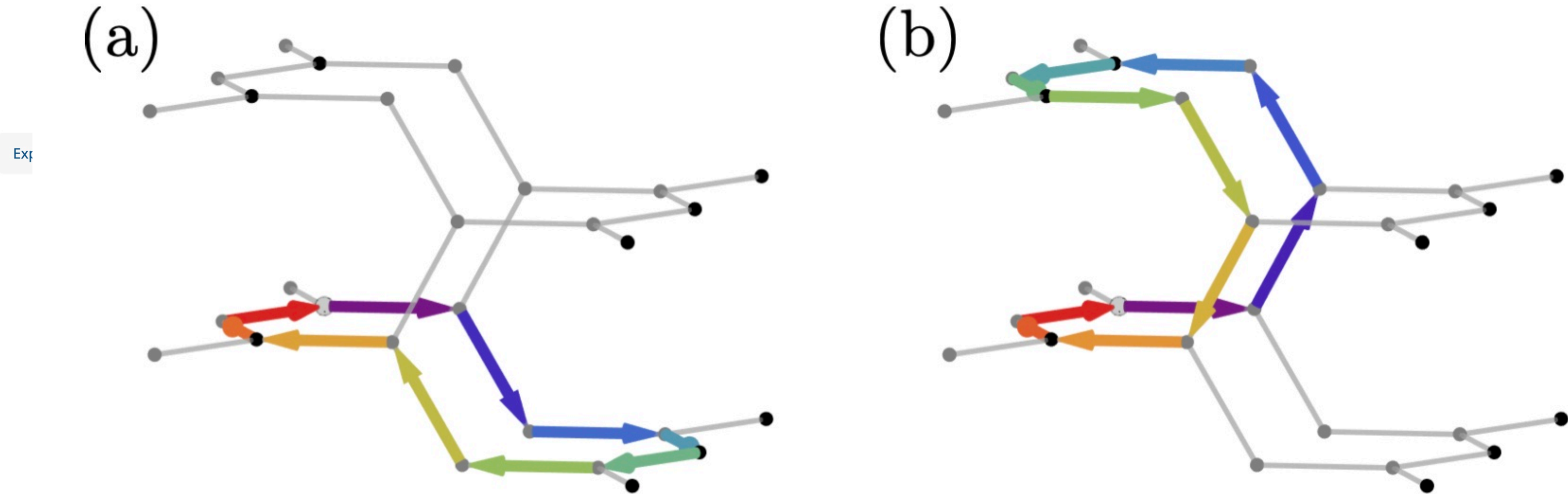
Improved honeycomb and hyperhoneycomb lattice Hamiltonians for quantum simulations of non-Abelian gauge theories

Marc Illa^{*}, Martin J. Savage^{†,‡}, and Xiaojun Yao[§]

Show more

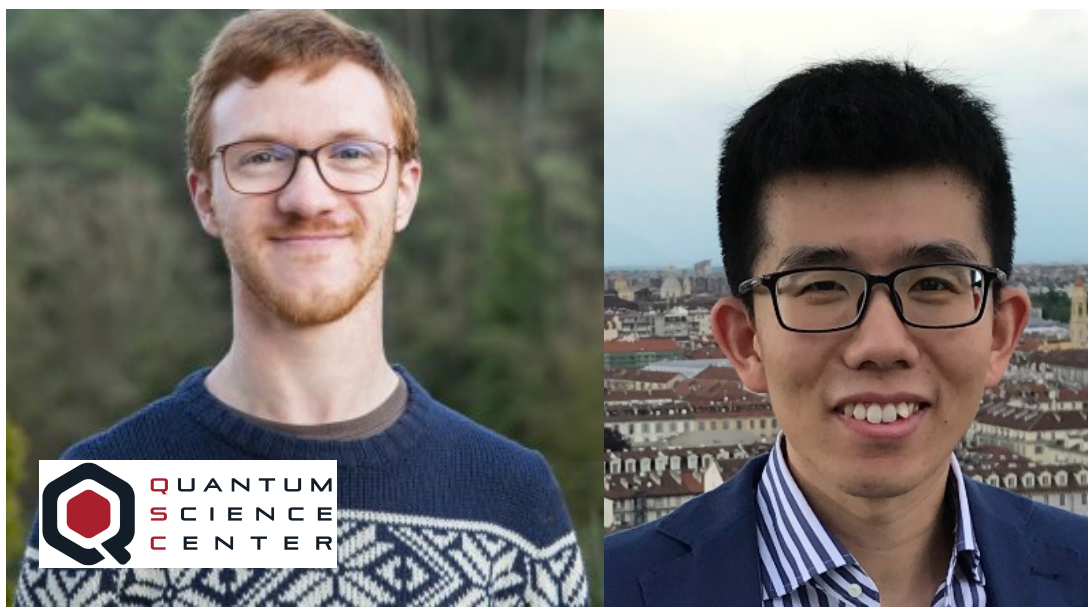
Phys. Rev. D **111**, 114520 – Published 26 June, 2025

DOI: <https://doi.org/10.1103/3rwf-f844>



Defines Leading-Order Magnetic Contribution

$$\Gamma_{10,12}^{(\text{HHC})} = \Gamma_{10}^{(\text{HHC})} + \frac{5}{4}\Gamma_{12}^{(\text{HHC})} = \frac{3g^2bV}{8} \sum_{a,i} [B_i^a(\boldsymbol{x})]^2 + \mathcal{O}(b^6)$$

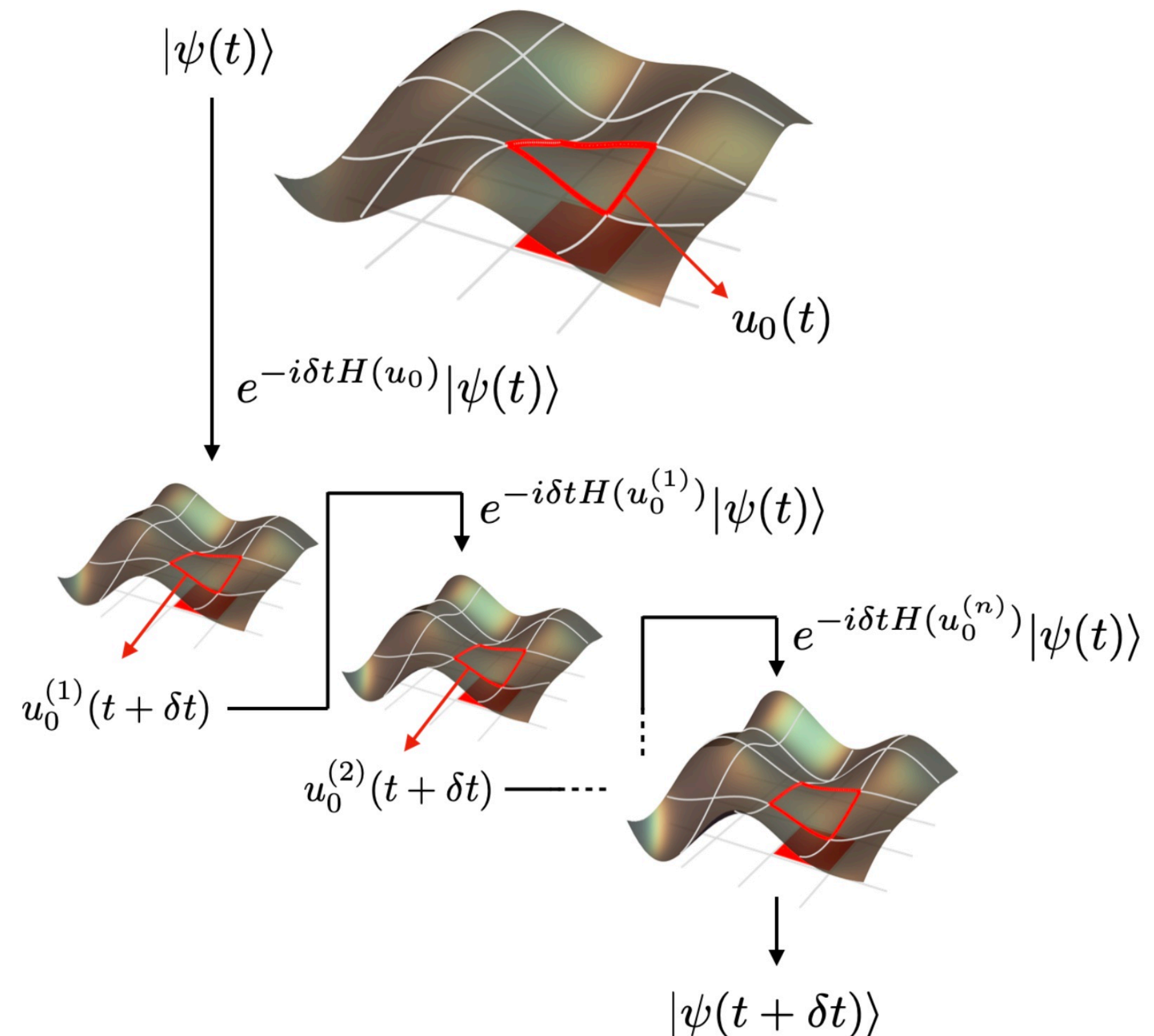


Improved Hamiltonian - Tadpoles

Dynamical Local Tadpole-Improvement in Quantum Simulations of Gauge Theories

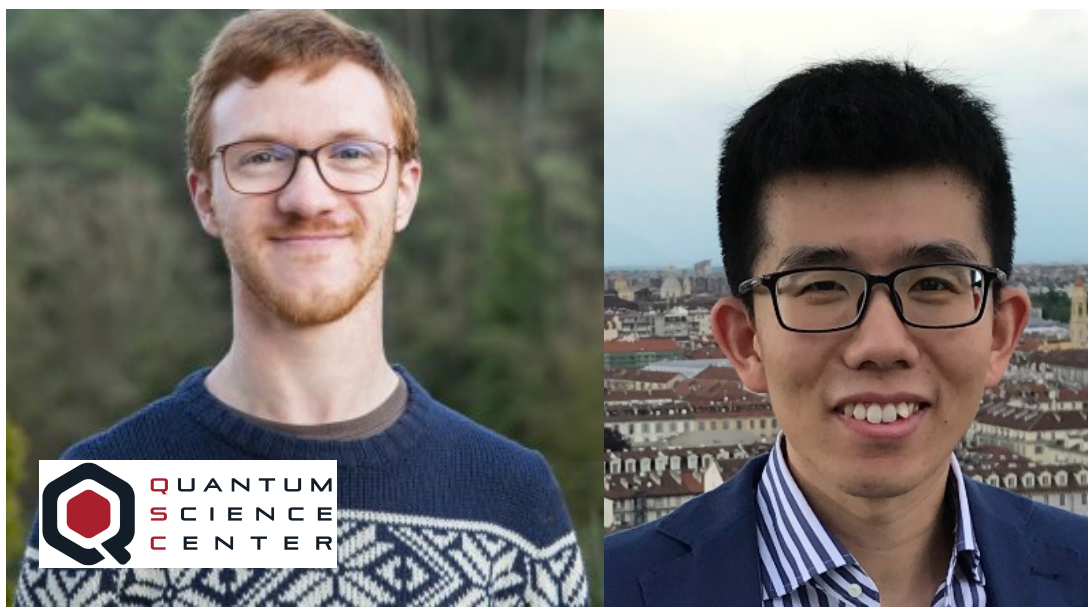
Marc Illa ,* Martin J. Savage ,† and Xiaojun Yao

Space-time dependent tadpole corrections



$$\hat{H} = \frac{g^2}{2a^{d-2}} \sum_{b, \text{links}} |\hat{\mathbf{E}}^{(b)}|^2 + \frac{1}{2a^{4-d}g^2} \sum_i \left[2N_c - \frac{1}{u_{0,i}^4} \left(\hat{\square}_i + \hat{\square}_i^\dagger \right) \right]$$

$$u_{0,i} = \left(1 + \frac{1}{2N_c} \langle \psi | \hat{\square}_i + \hat{\square}_i^\dagger | \psi \rangle \right)^{1/4}$$



Improved Hamiltonian - Tadpoles

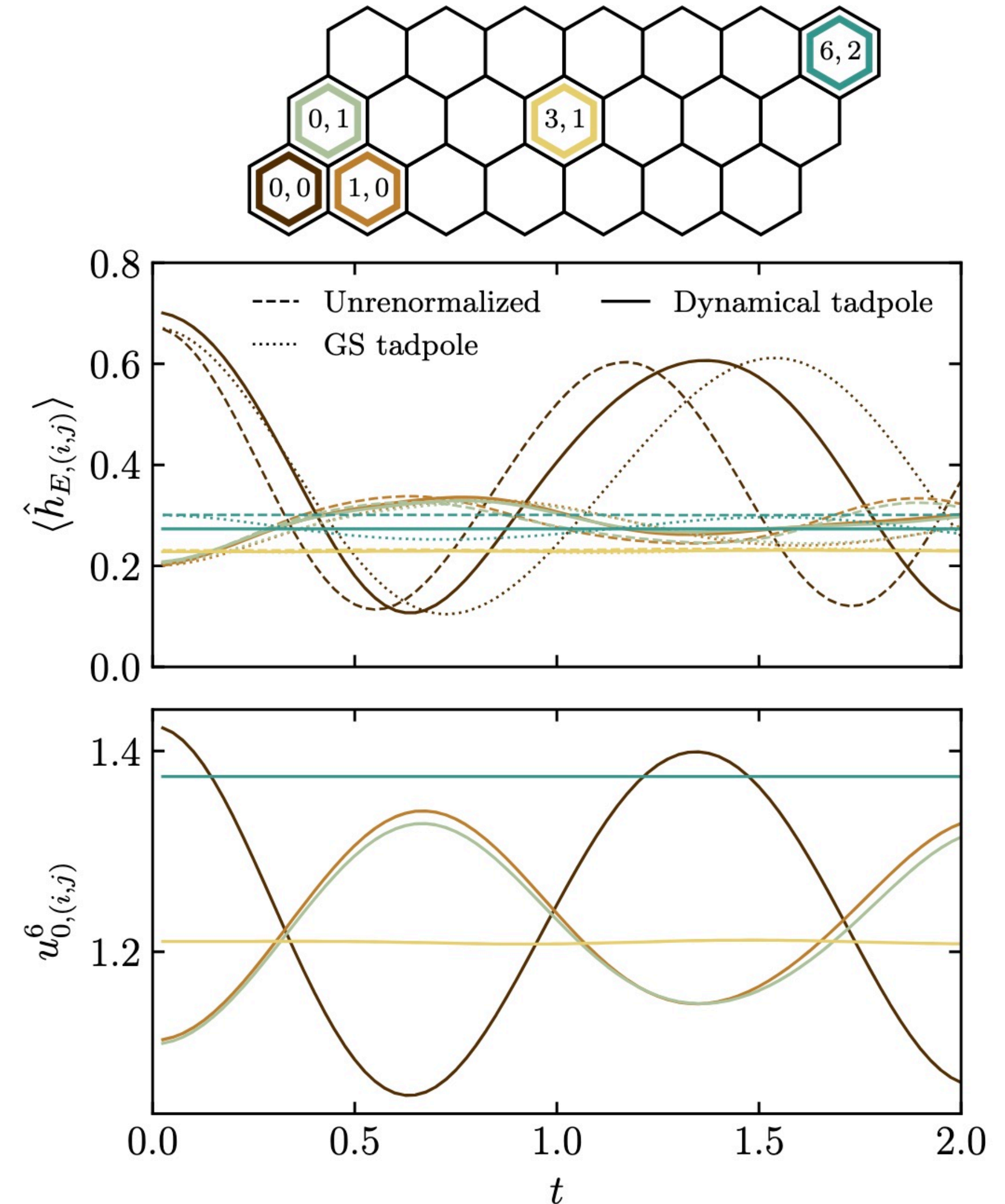
Dynamical Local Tadpole-Improvement in Quantum Simulations of Gauge Theories

Marc Illa ,* Martin J. Savage ,† and Xiaojun Yao

Scalar Field Theory with bounded, smooth wave packet on interacting vacuum

$$\langle \Psi | \phi_\alpha^2 | \Psi \rangle \rightarrow \frac{1}{L} \sum_k \frac{1}{2E_k} + \left| \sum_k \frac{g_k V_{k\alpha}}{\sqrt{E_k}} \right|^2$$

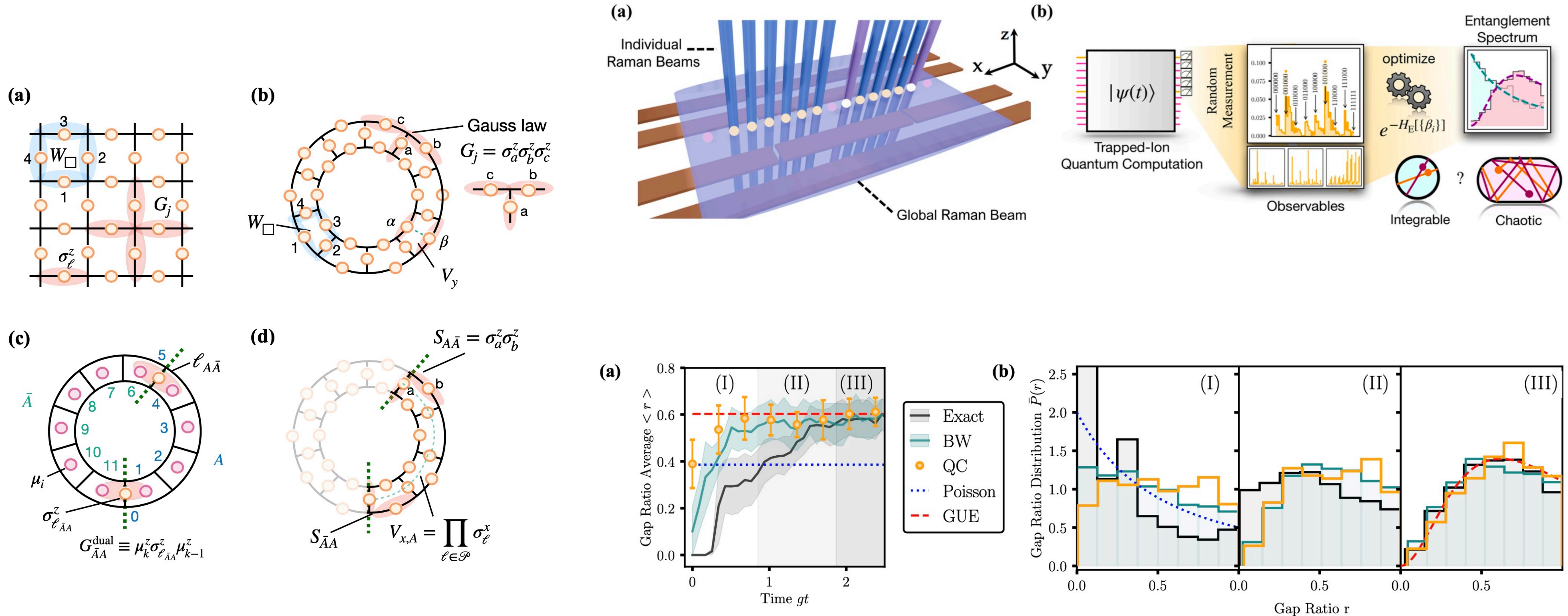
Vacuum tadpole correction plus Lattice spacing finite density term



Entanglement and Thermalization

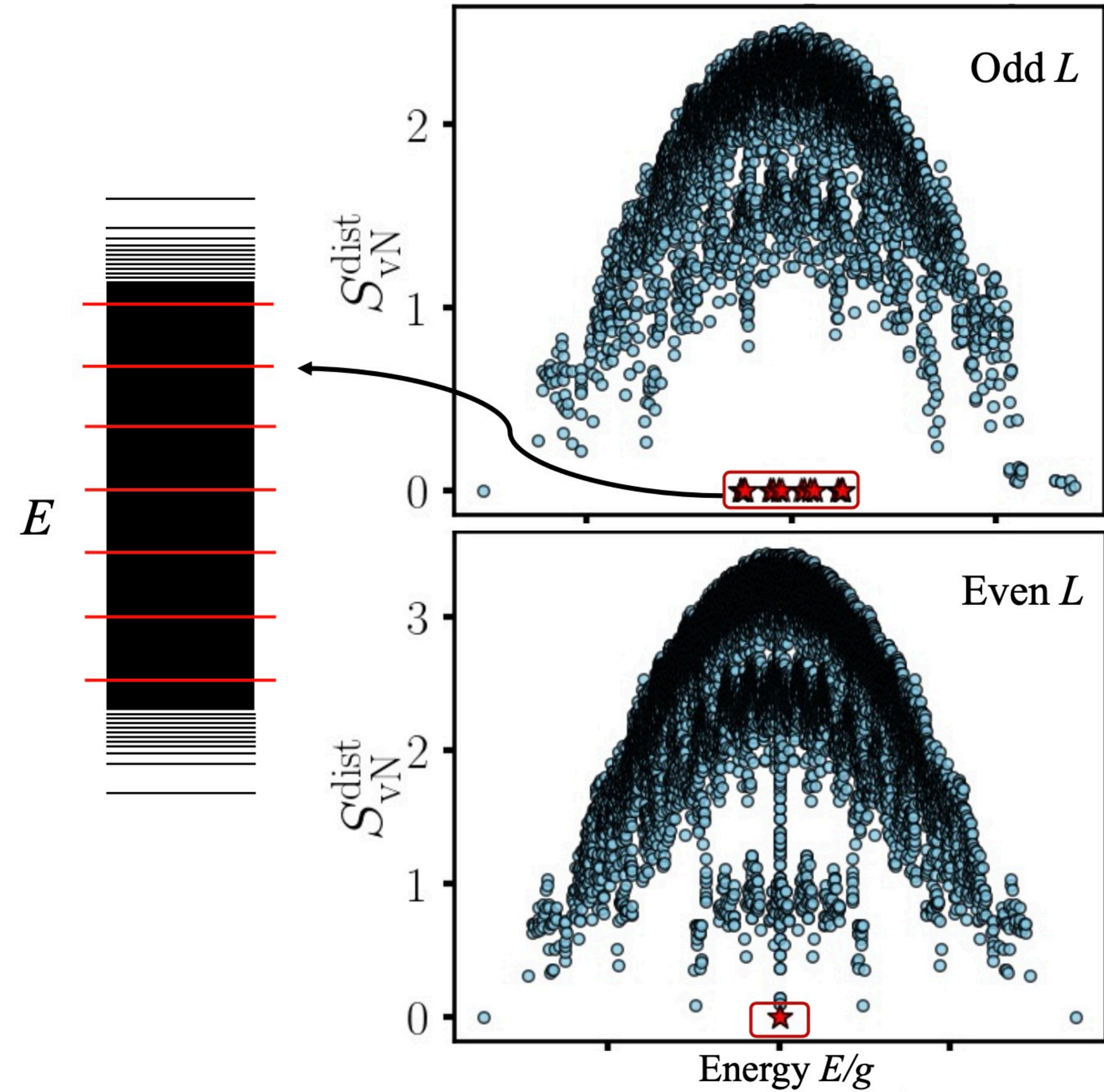
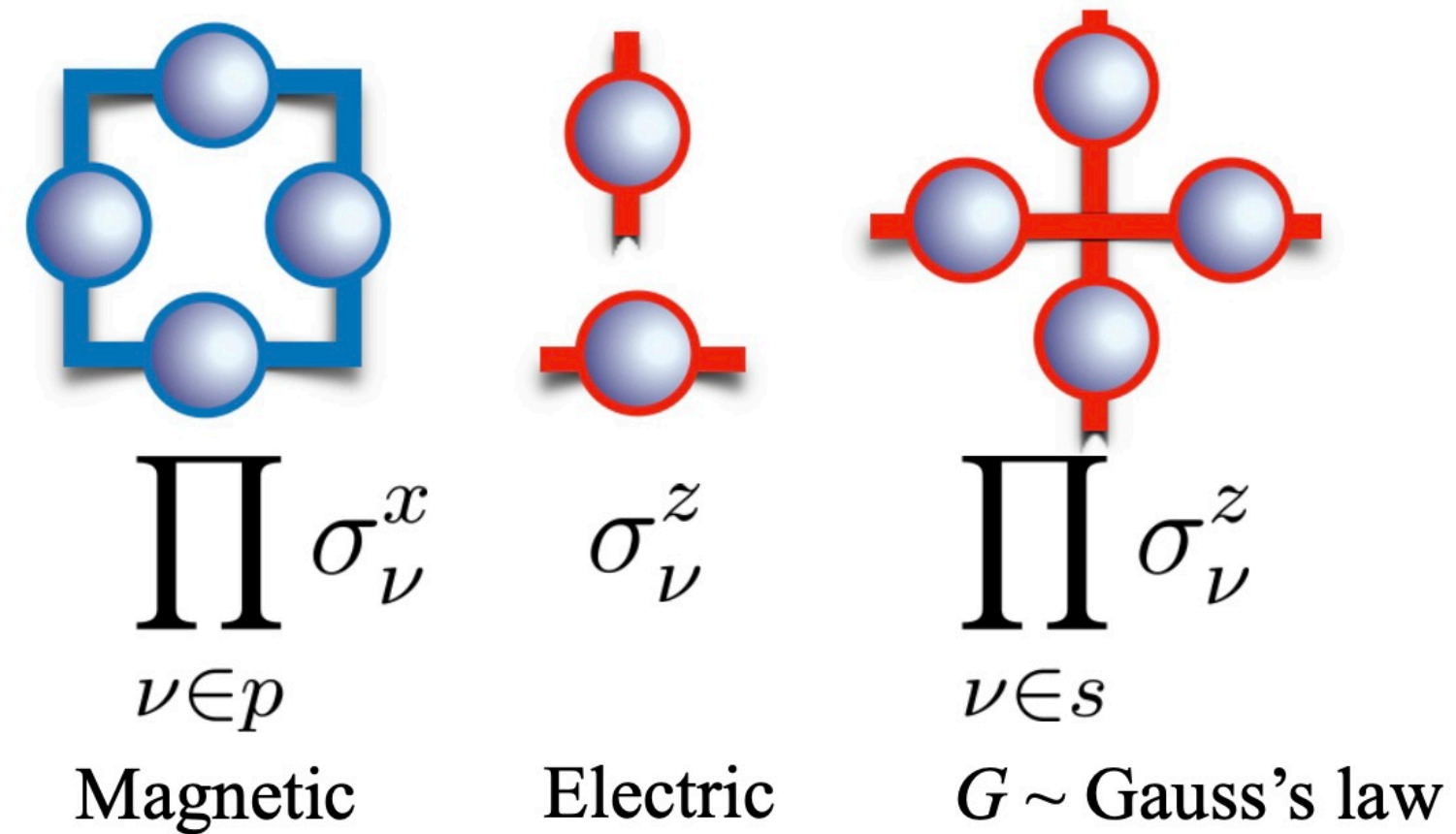
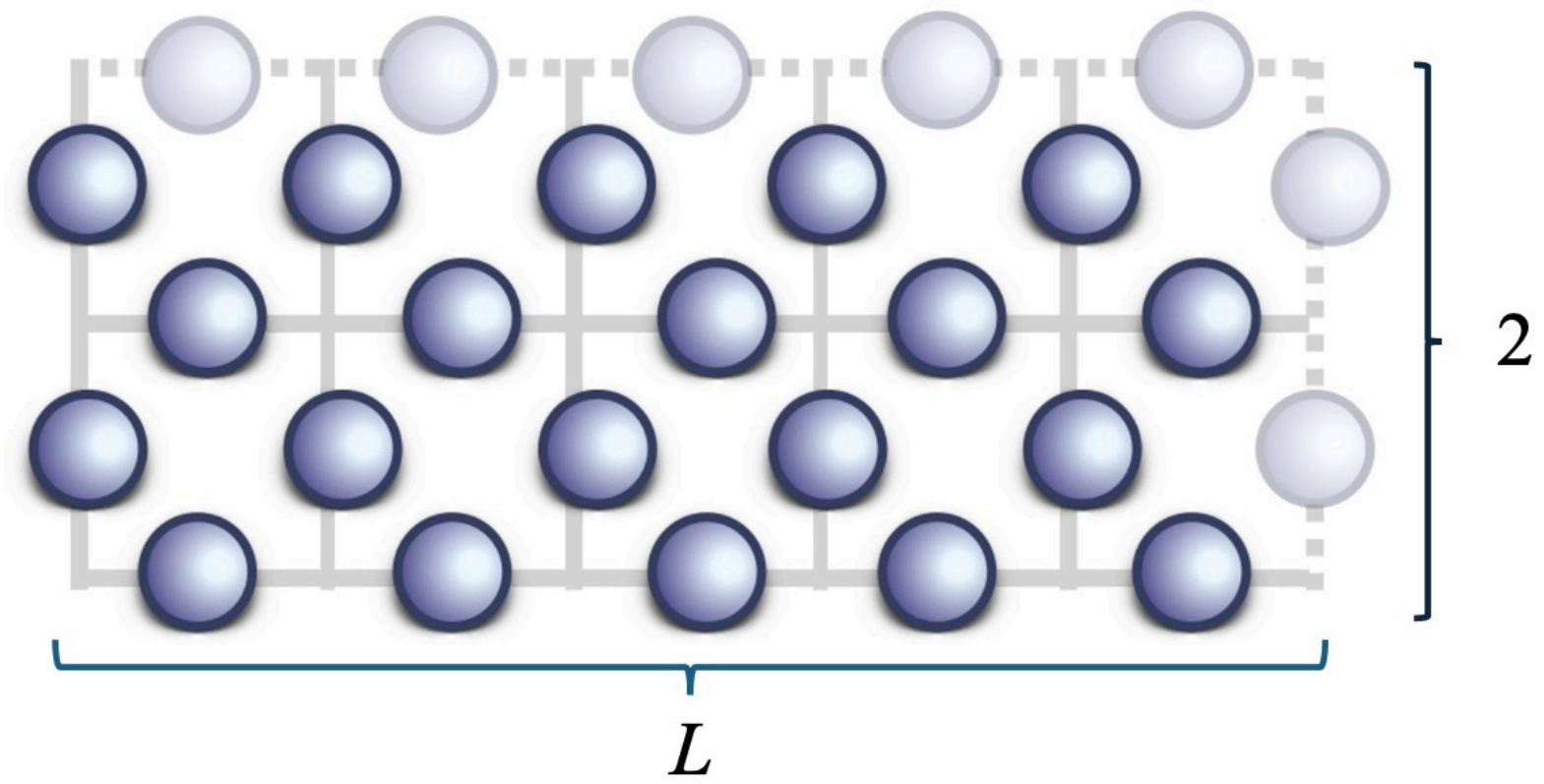
Quantum Computing Universal Thermalization Dynamics in a (2+1)D Lattice Gauge Theory

Niklas Mueller,^{1,*} Tianyi Wang,^{2,3,4} Or Katz,^{3,5,6} Zohreh Davoudi,^{7,8,4,9} and Marko Cetina^{2,3,5,4}



Entanglement and Thermalization

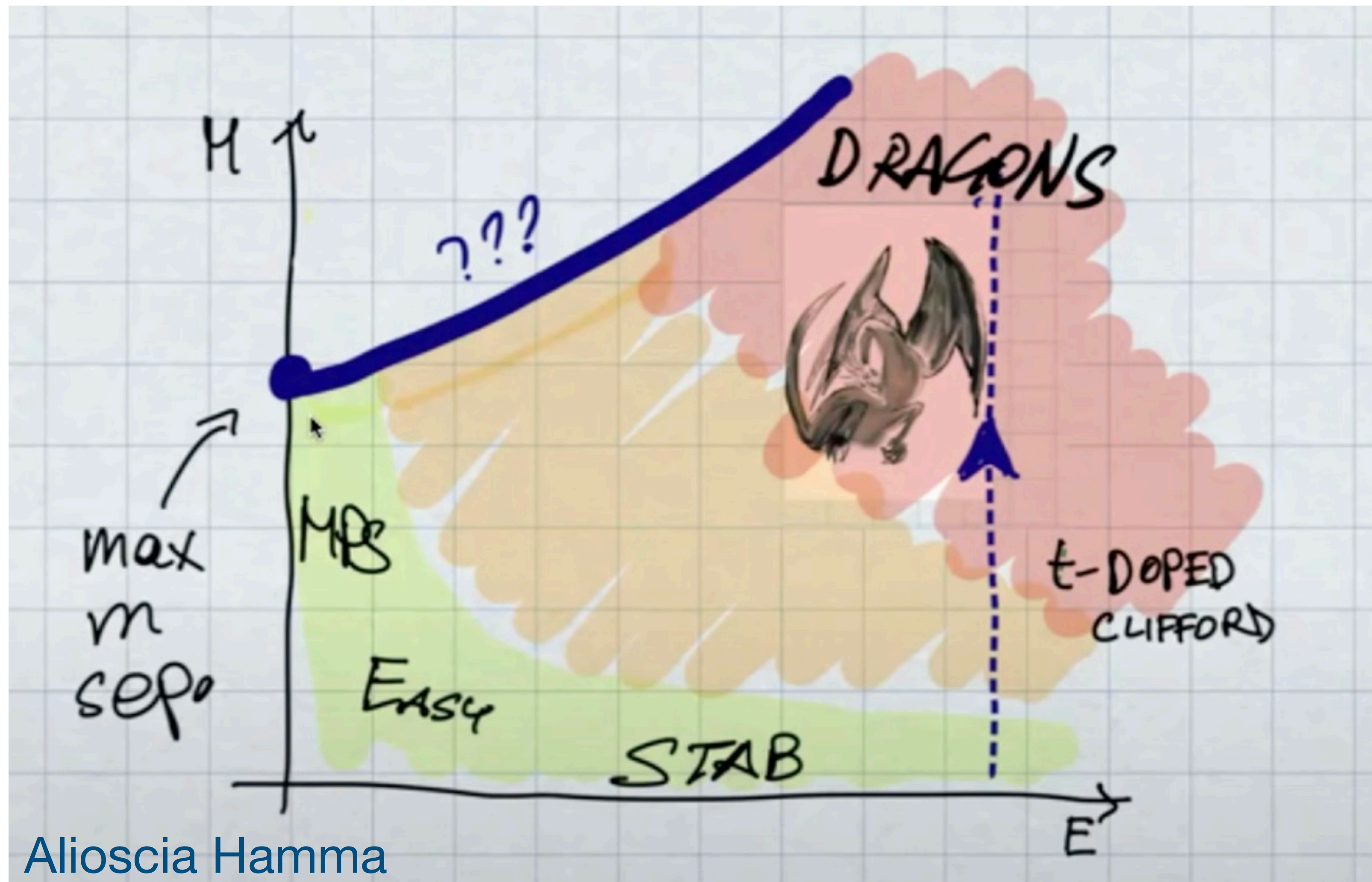
Stabilizer Scars



When is a Quantum Computer Required?

Quantum Complexity of Physical Systems?

Magic

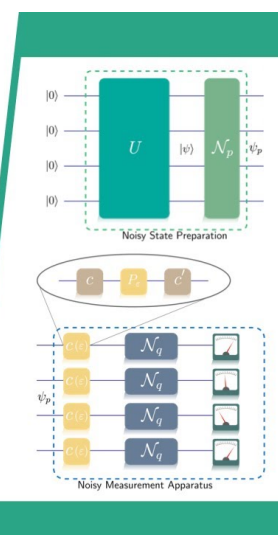


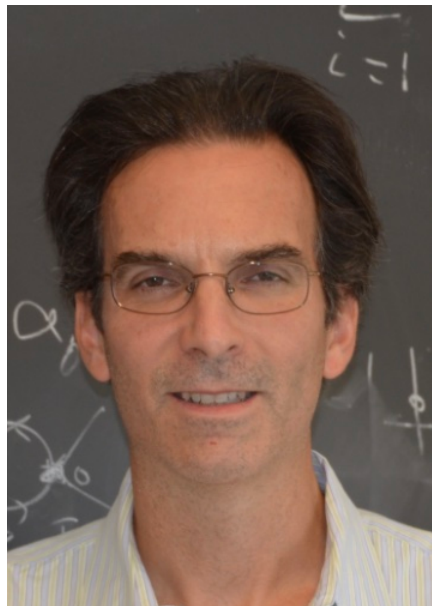
Entanglement

qmts #4

Why should you care about Magic

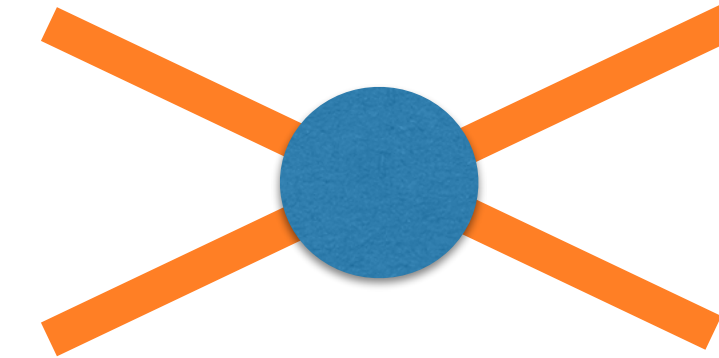
Prof. Alioscia Hamma
Dr. Salvatore F.E. Oliviero





Entanglement Suppression and Emergent Symmetries of Strong Interactions

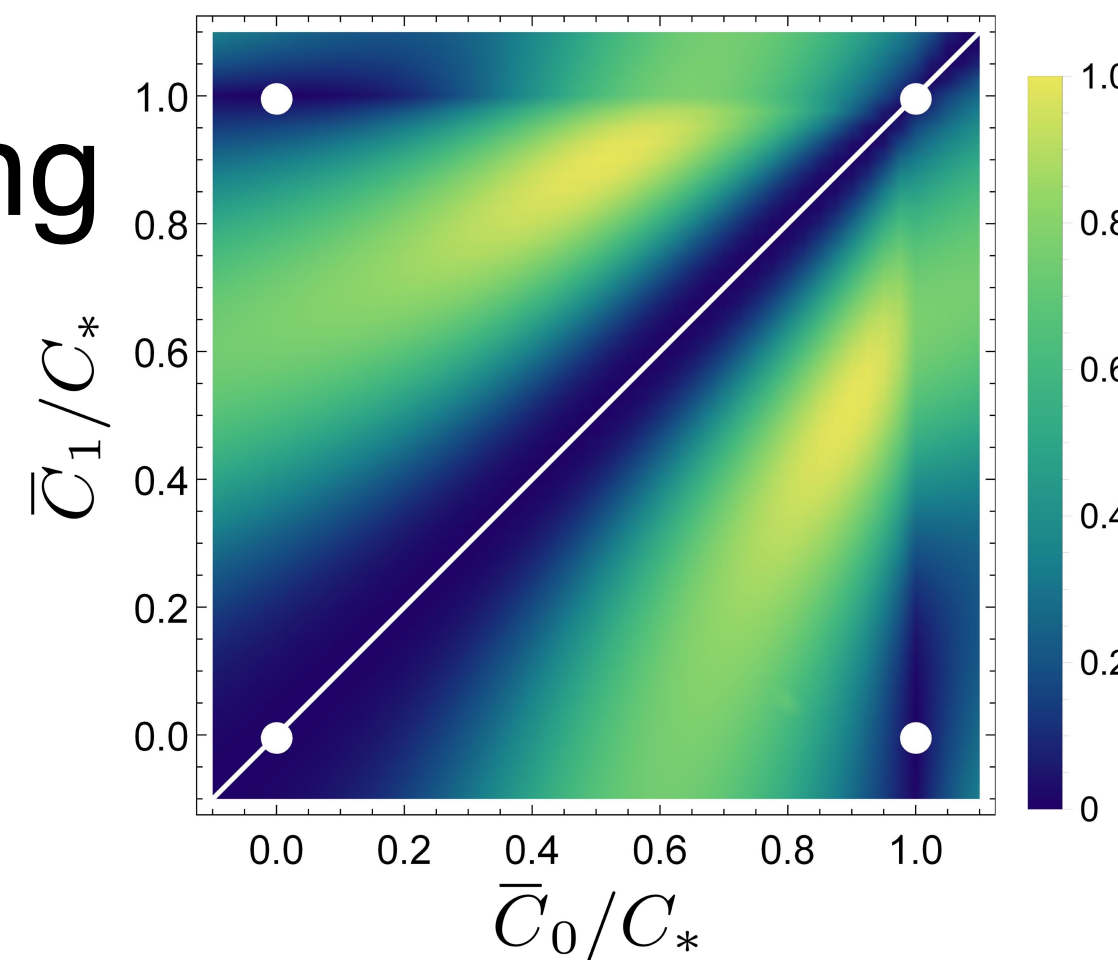
Silas R. Beane, David B. Kaplan, Natalie Klco, and Martin J. Savage
Phys. Rev. Lett. **122**, 102001 – Published 14 March 2019



Low-energy QCD: NN and YN S-Wave Scattering

$$\hat{\mathbf{S}}_\sigma = \frac{1}{4} (3e^{i2\delta_3} + e^{i2\delta_1}) \hat{\mathbf{1}} + \frac{1}{4} (e^{i2\delta_3} - e^{i2\delta_1}) \hat{\boldsymbol{\sigma}} \cdot \hat{\boldsymbol{\sigma}}$$

$$e_p(\hat{A}) \longrightarrow \mathcal{E}(\hat{\mathbf{S}}_\sigma) = \frac{1}{6} \sin^2(2(\delta_3 - \delta_1))$$

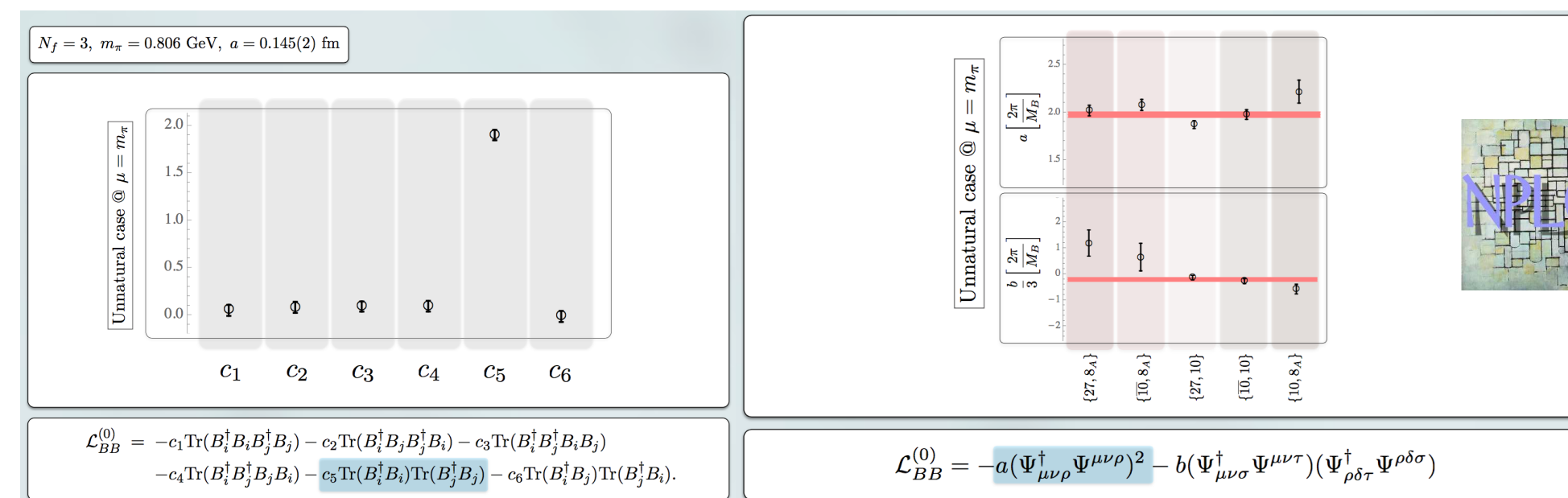


conformal points



Wigner symmetry

SU(4) for 2 flavors and **SU(16)** for 3 flavors
- more symmetry than large- N_c , [SU(4) and SU(6)]

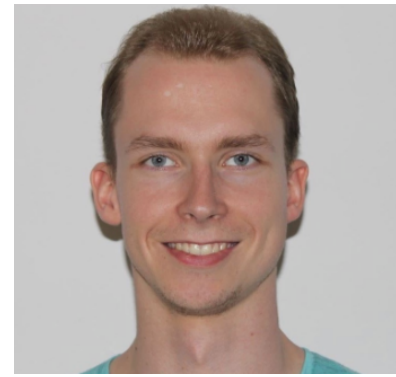


Emergent approximate symmetries in nuclear systems



Suppressed fluctuations in entanglement

Suppressed sign problems in classical simulations



Entanglement Rearrangement, Effective Model Spaces and Symmetries

[Home](#) > [The European Physical Journal A](#) > [Article](#)

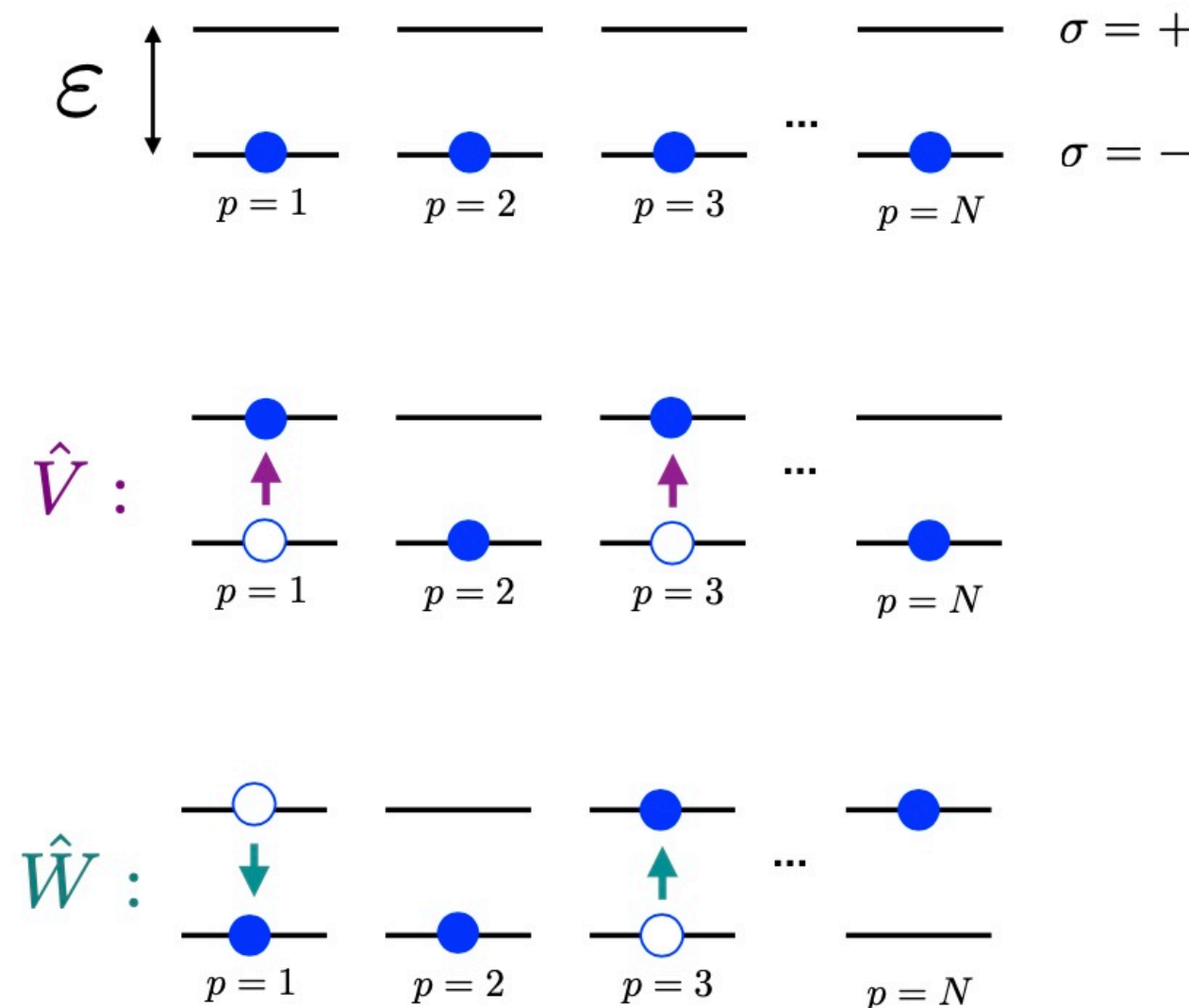
Multi-body entanglement and information rearrangement in nuclear many-body systems: a study of the Lipkin–Meshkov–Glick model

Regular Article – Theoretical Physics | Published: 17 October 2023

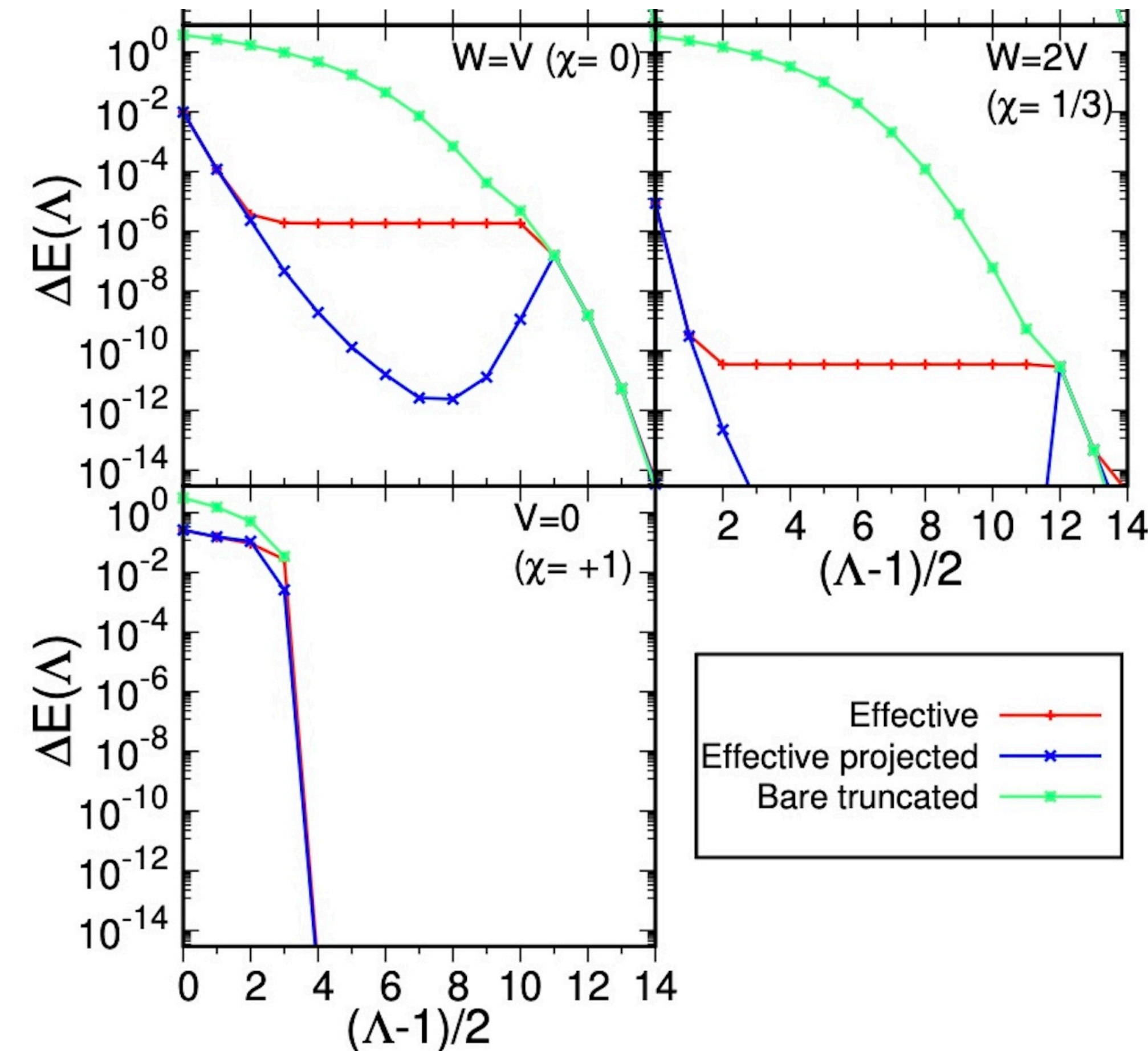
Volume 59, article number 231, (2023) [Cite this article](#)

S. Momme Hengstenberg, Caroline E. P. Robin & Martin J. Savage

$$\begin{aligned}\hat{H} &= \frac{\varepsilon}{2} \sum_{\sigma p} \sigma \hat{c}_{p\sigma}^\dagger \hat{c}_{p\sigma} - \frac{V}{2} \sum_{pq\sigma} \hat{c}_{p\sigma}^\dagger \hat{c}_{q\sigma}^\dagger \hat{c}_{q-\sigma} \hat{c}_{p-\sigma} \\ &\quad - \frac{W}{2} \sum_{pq\sigma} \hat{c}_{p\sigma}^\dagger \hat{c}_{q-\sigma}^\dagger \hat{c}_{q\sigma} \hat{c}_{p-\sigma} , \\ &= \varepsilon \hat{J}_z - \frac{V}{2} (\hat{J}_+^2 + \hat{J}_-^2) - \frac{W}{2} (\hat{J}_+ \hat{J}_- + \hat{J}_- \hat{J}_+ - \hat{N}) ,\end{aligned}$$



$$\begin{pmatrix} \hat{c}_{p+}(\beta) \\ \hat{c}_{p-}(\beta) \end{pmatrix} = \begin{pmatrix} \cos(\beta/2) & -\sin(\beta/2) \\ \sin(\beta/2) & \cos(\beta/2) \end{pmatrix} \begin{pmatrix} \hat{c}_{p+} \\ \hat{c}_{p-} \end{pmatrix}$$





The Magic Power of the S-Matrix - Fluctuations in Magic

$$\overline{\mathcal{M}}(\hat{\mathbf{S}}) \equiv \frac{1}{\mathcal{N}_{ss}} \sum_{i=1}^{\mathcal{N}_{ss}} \mathcal{M}(\hat{\mathbf{S}}|\Psi_i\rangle)$$

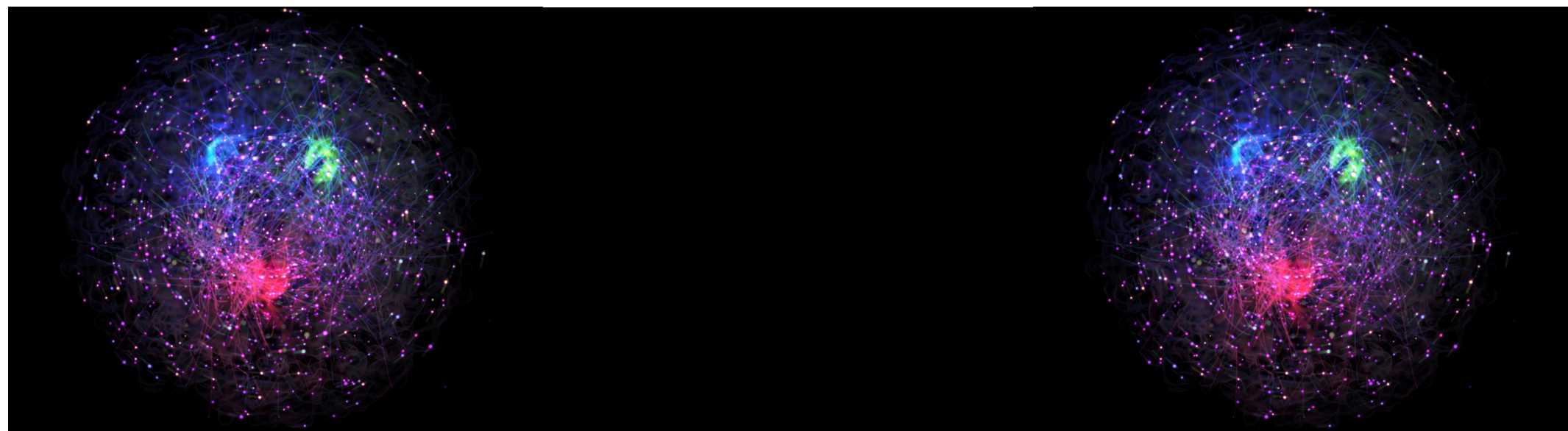
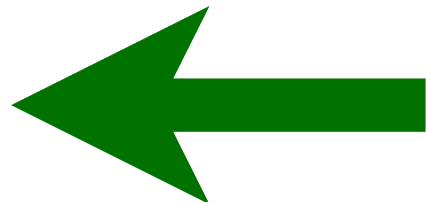
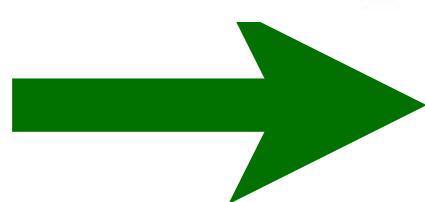
arXiv > nucl-th > arXiv:2405.10268

Nuclear Theory

[Submitted on 16 May 2024 (v1), last revised 20 May 2024 (this version, v2)]

The Magic in Nuclear and Hypernuclear Forces

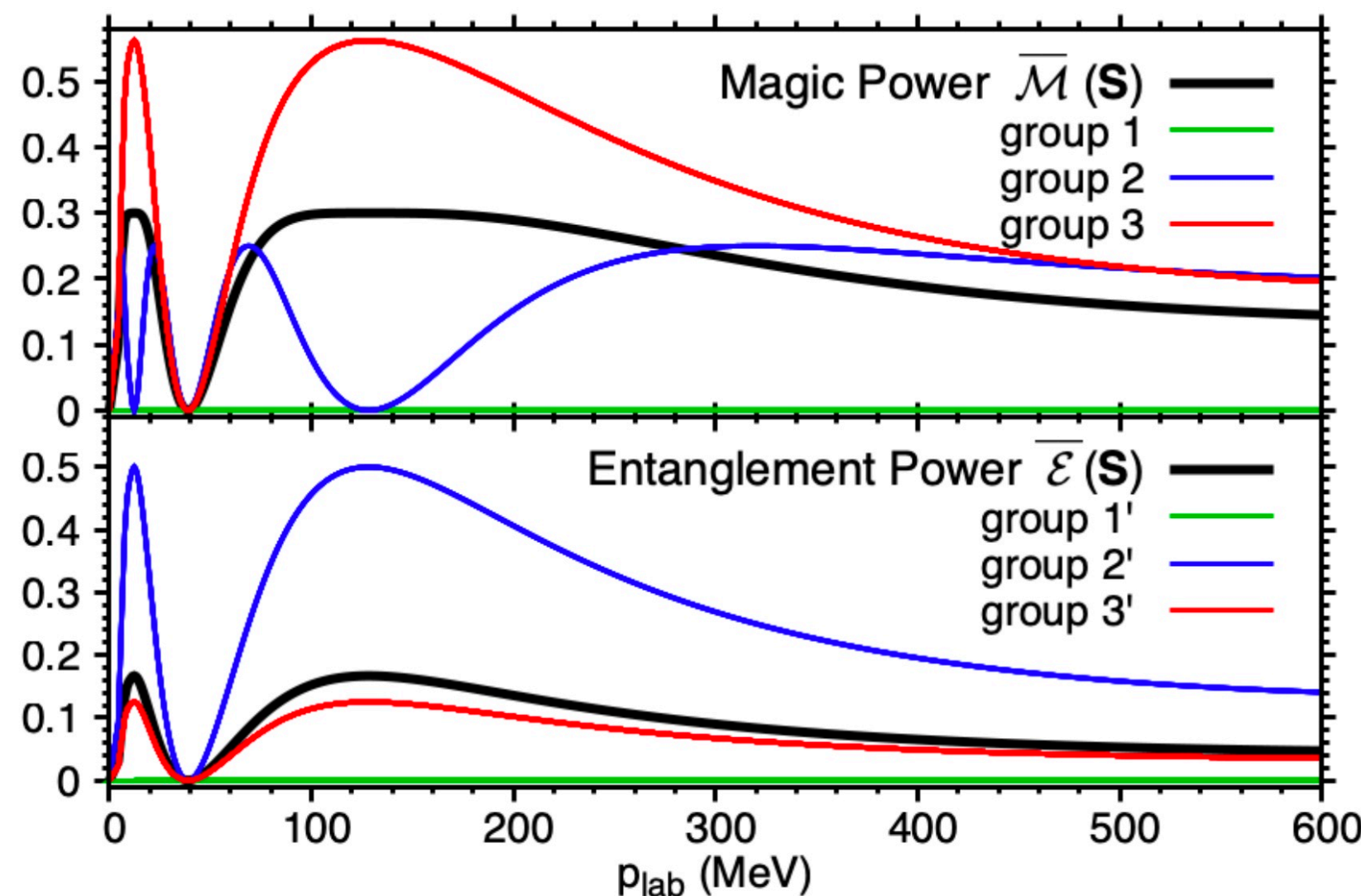
Caroline E. P. Robin, Martin J. Savage



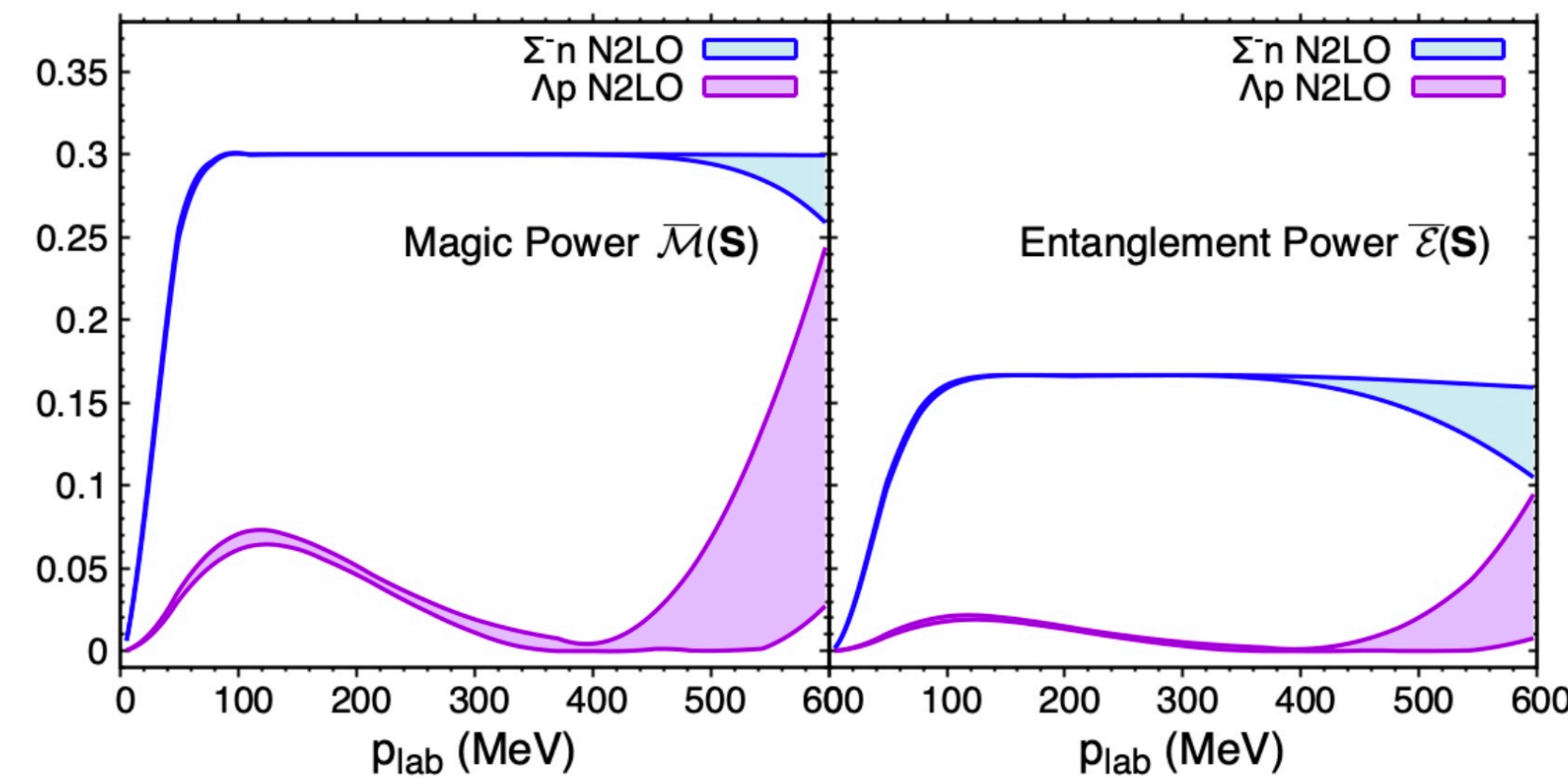
$$\hat{\mathbf{S}}_{\sigma} = \frac{1}{4} (3e^{i2\delta_3} + e^{i2\delta_1}) \hat{\mathbf{1}} + \frac{1}{4} (e^{i2\delta_3} - e^{i2\delta_1}) \hat{\boldsymbol{\sigma}} \cdot \hat{\boldsymbol{\sigma}}$$

$$\overline{\mathcal{M}}(\hat{\mathbf{S}}) = \frac{3}{20} (3 + \cos(4 \Delta\delta)) \sin^2(2 \Delta\delta) ,$$

$$\overline{\mathcal{E}}(\hat{\mathbf{S}}) = \frac{1}{6} \sin^2(2 \Delta\delta) ,$$



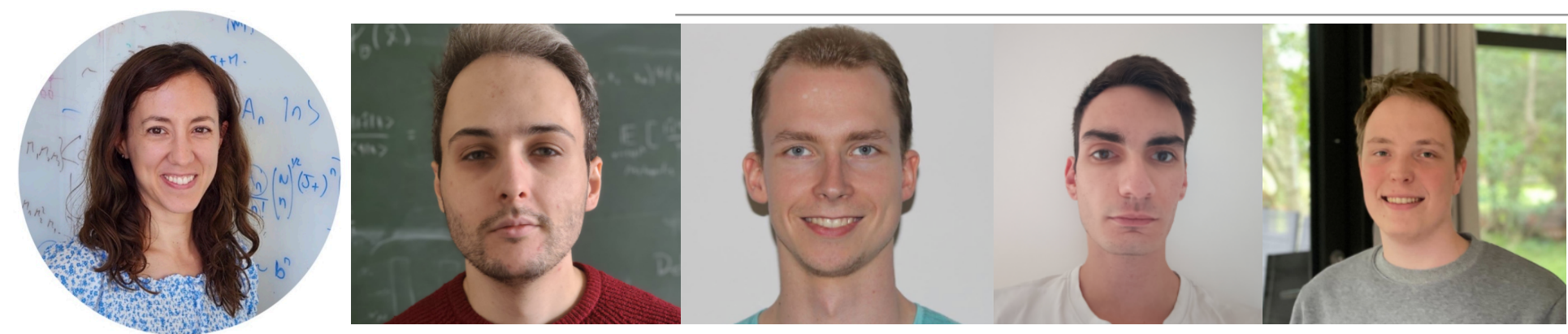
Neutron-Proton



Hyperon-Nucleon

Indicates the fluctuations in quantum magic due to nuclear forces

Magic and Multi-Partite Entanglement in Nuclei



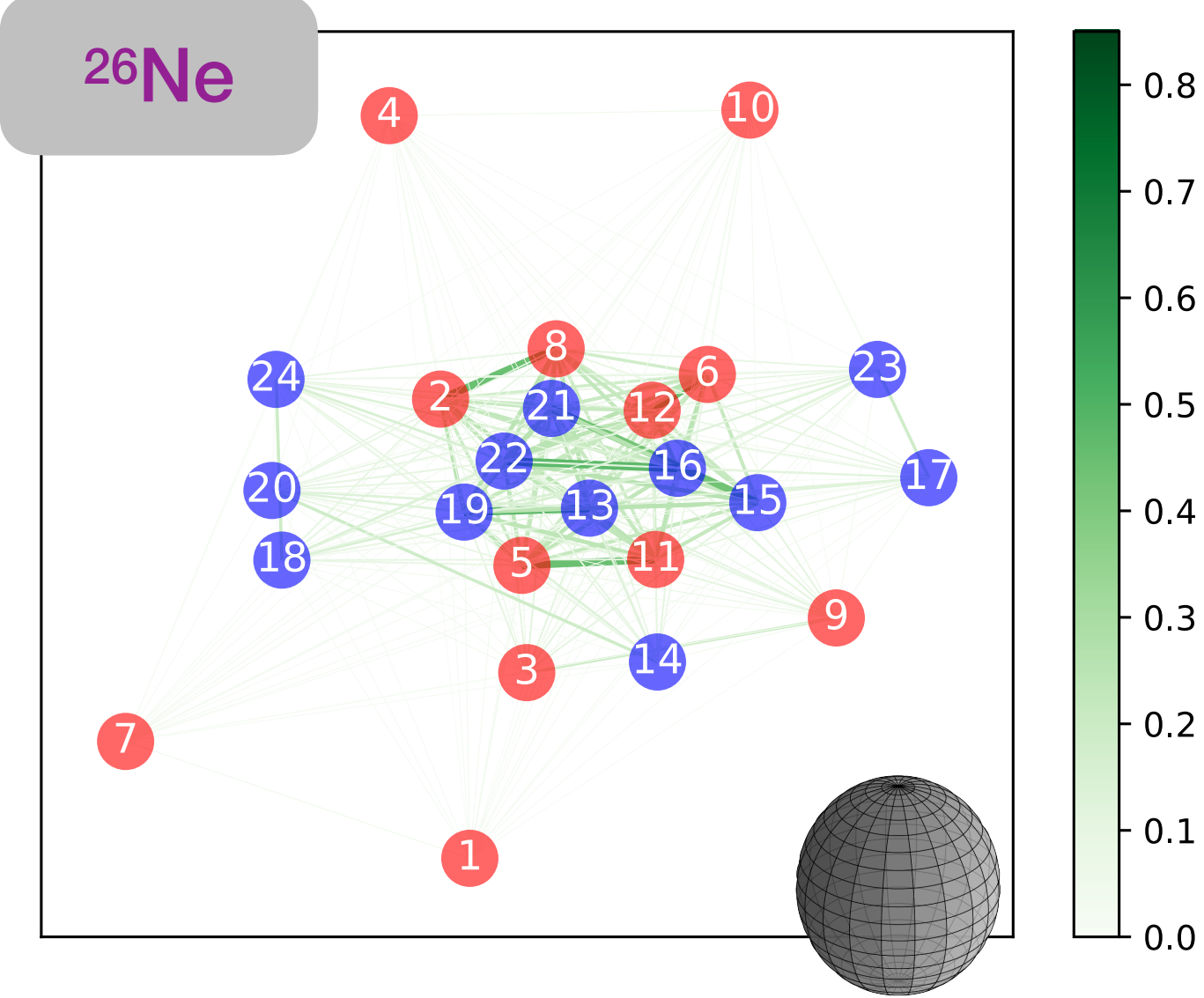
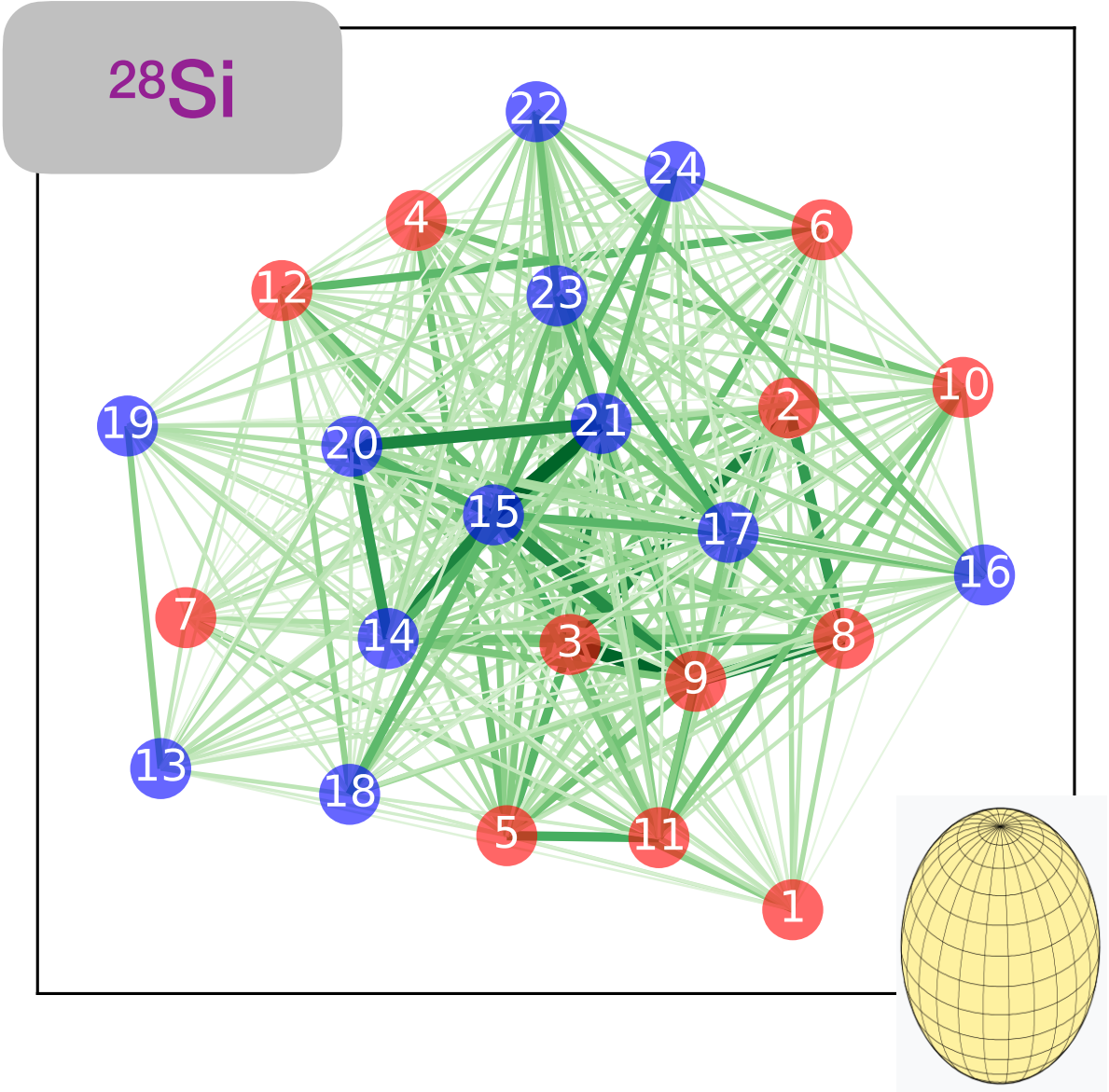
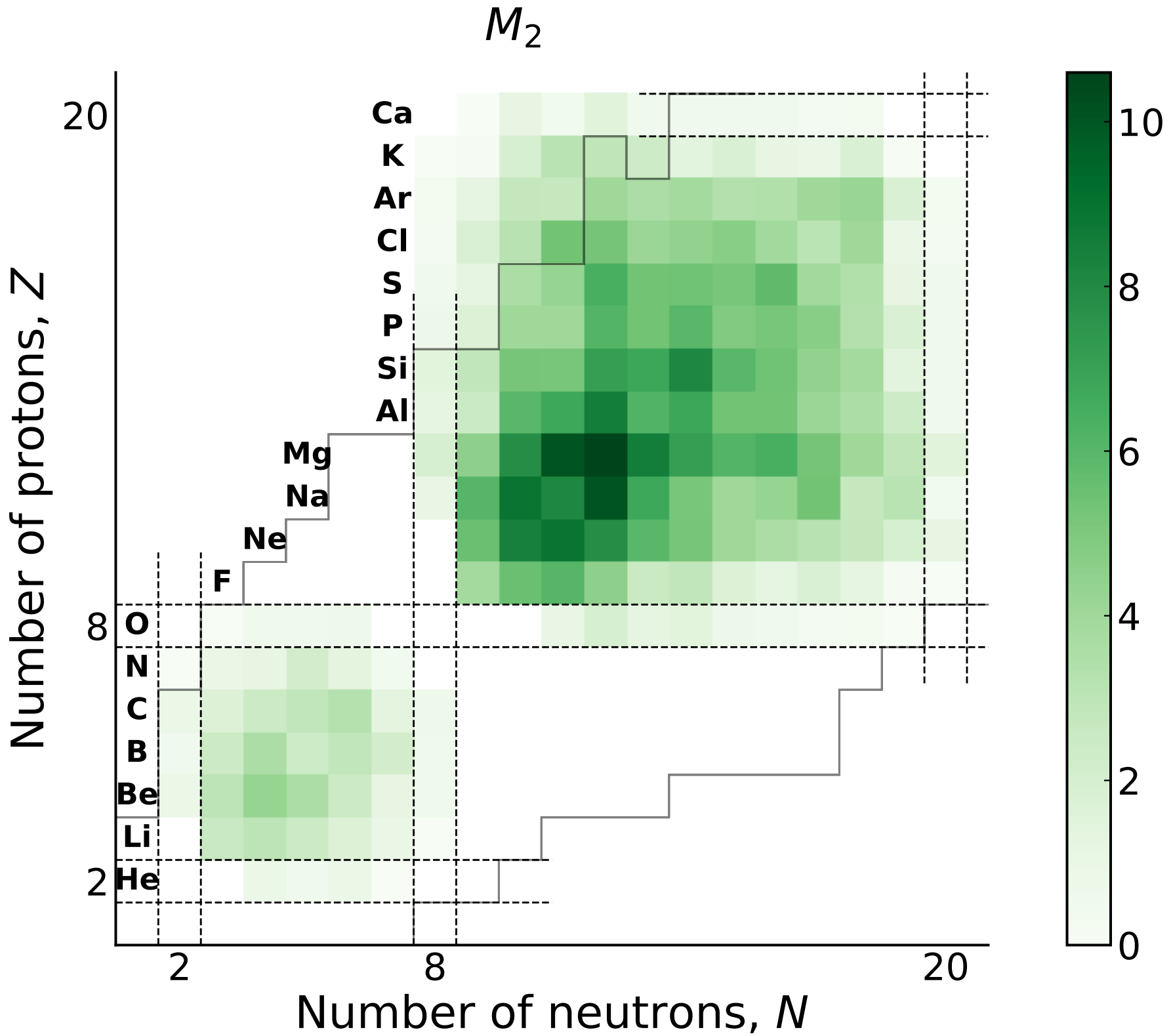
Quantum magic and multipartite entanglement in the structure of nuclei

[Florian Brökemeier](#) ¹, [S. Momme Hengstenberg](#) ¹, [James W. T. Keeble](#) ¹, [Caroline E. P. Robin](#) ^{1,2,*}, [Federico Rocco](#) ¹, and [Martin J. Savage](#) ^{3,†}

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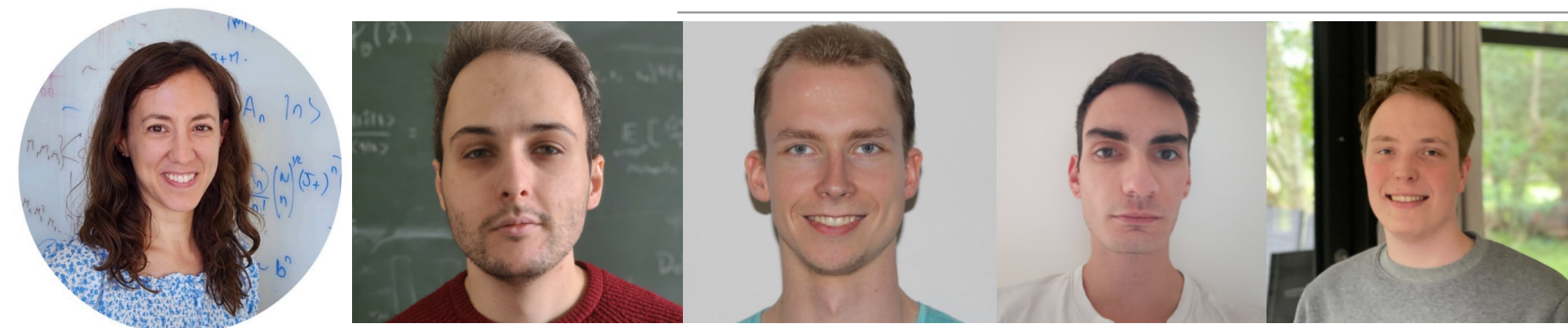
Phys. Rev. C **111**, 034317 – Published 11 March, 2025

DOI: <https://doi.org/10.1103/PhysRevC.111.034317>



8-tangles

Magic and Multi-Partite Entanglement in Nuclei



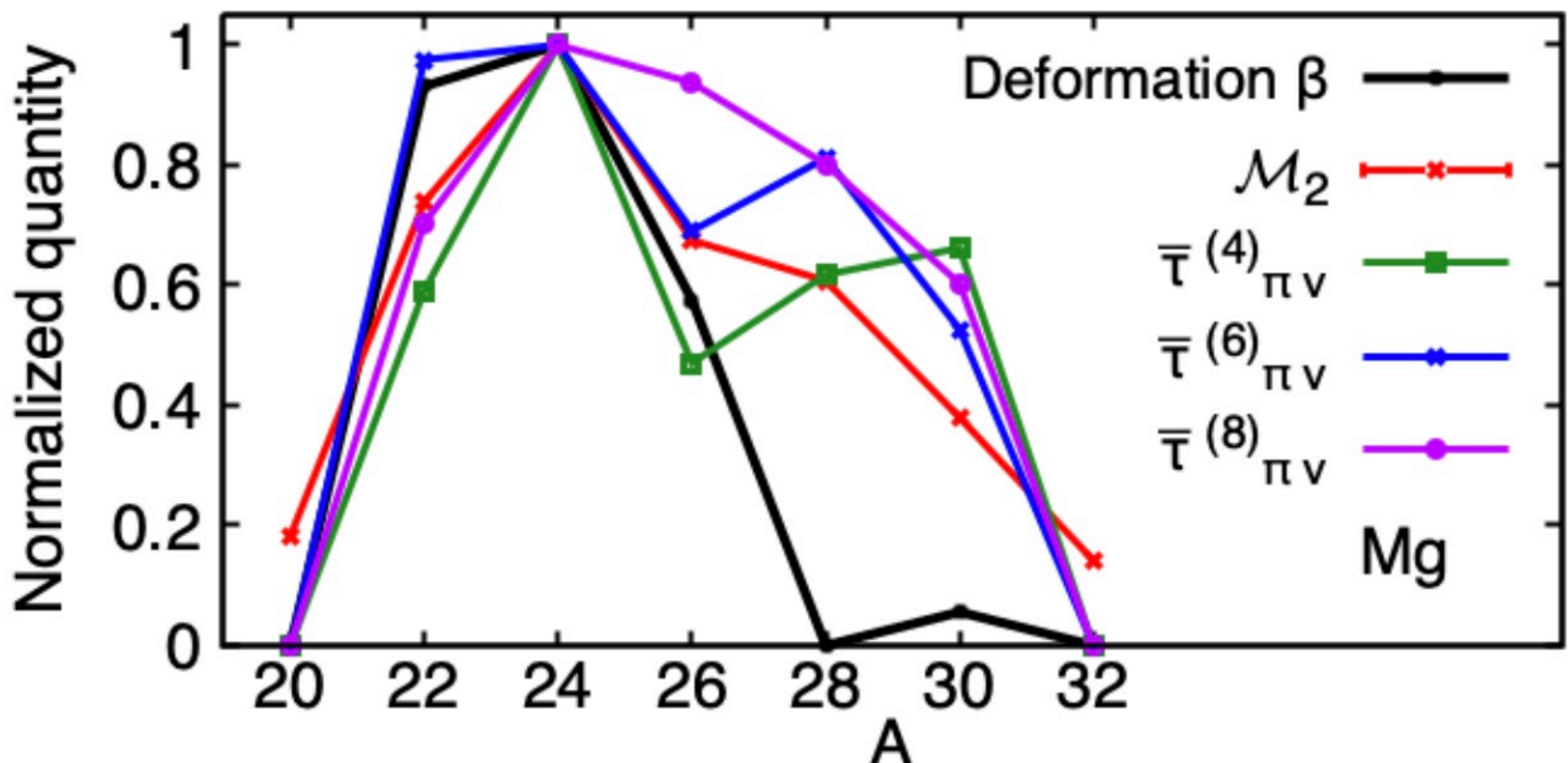
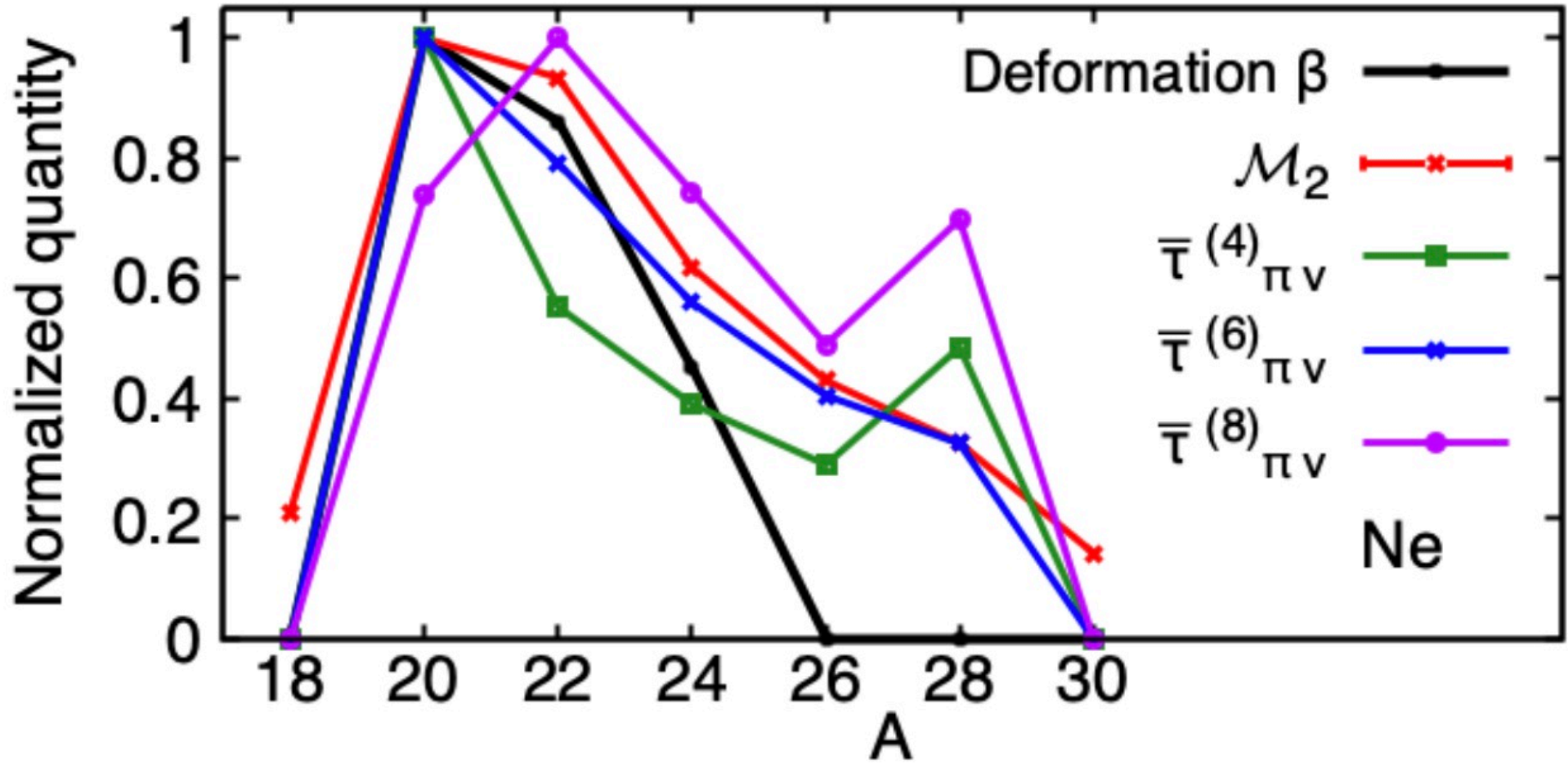
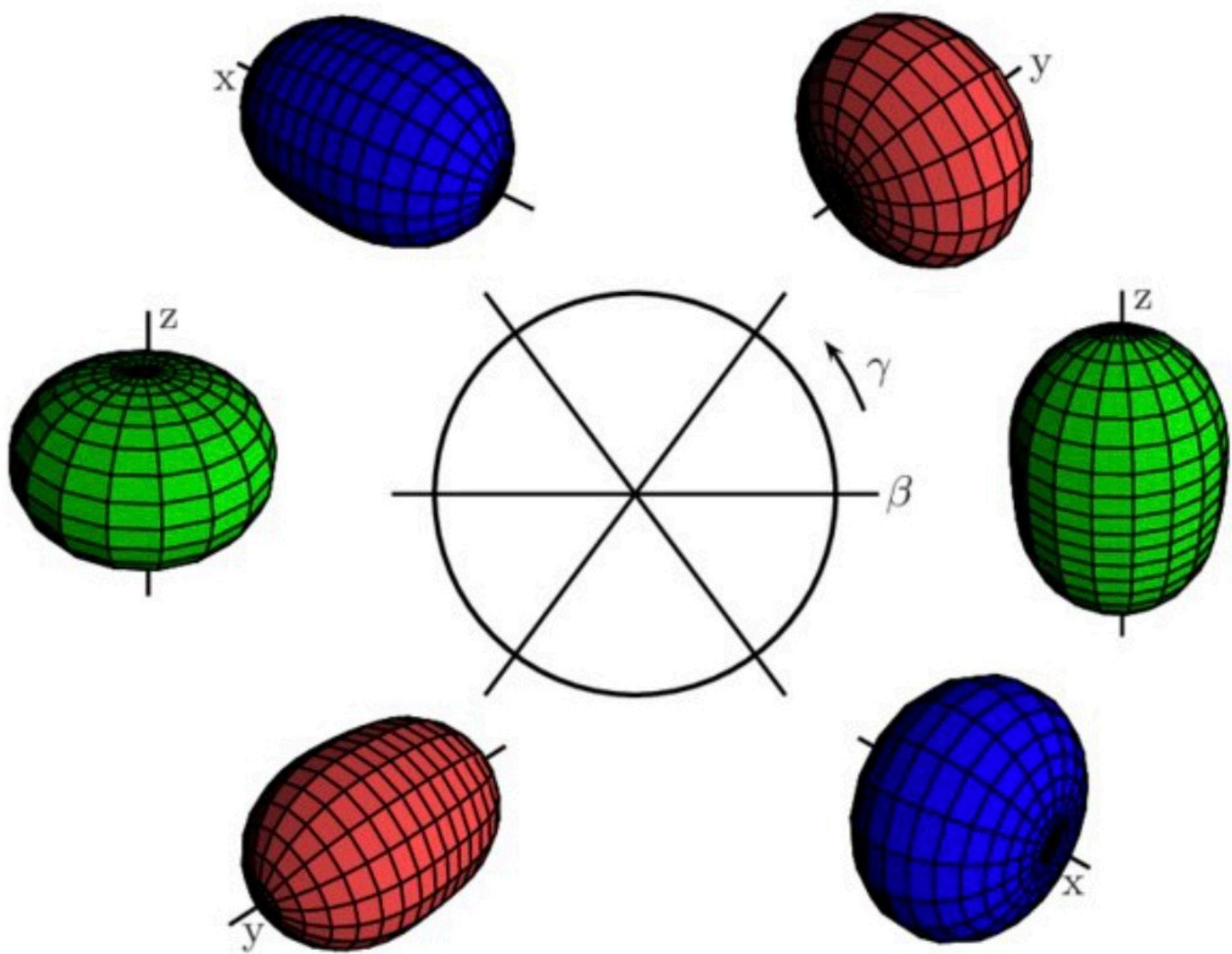
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Phys. Rev. C 111, 034317 – Published 11 March, 2025

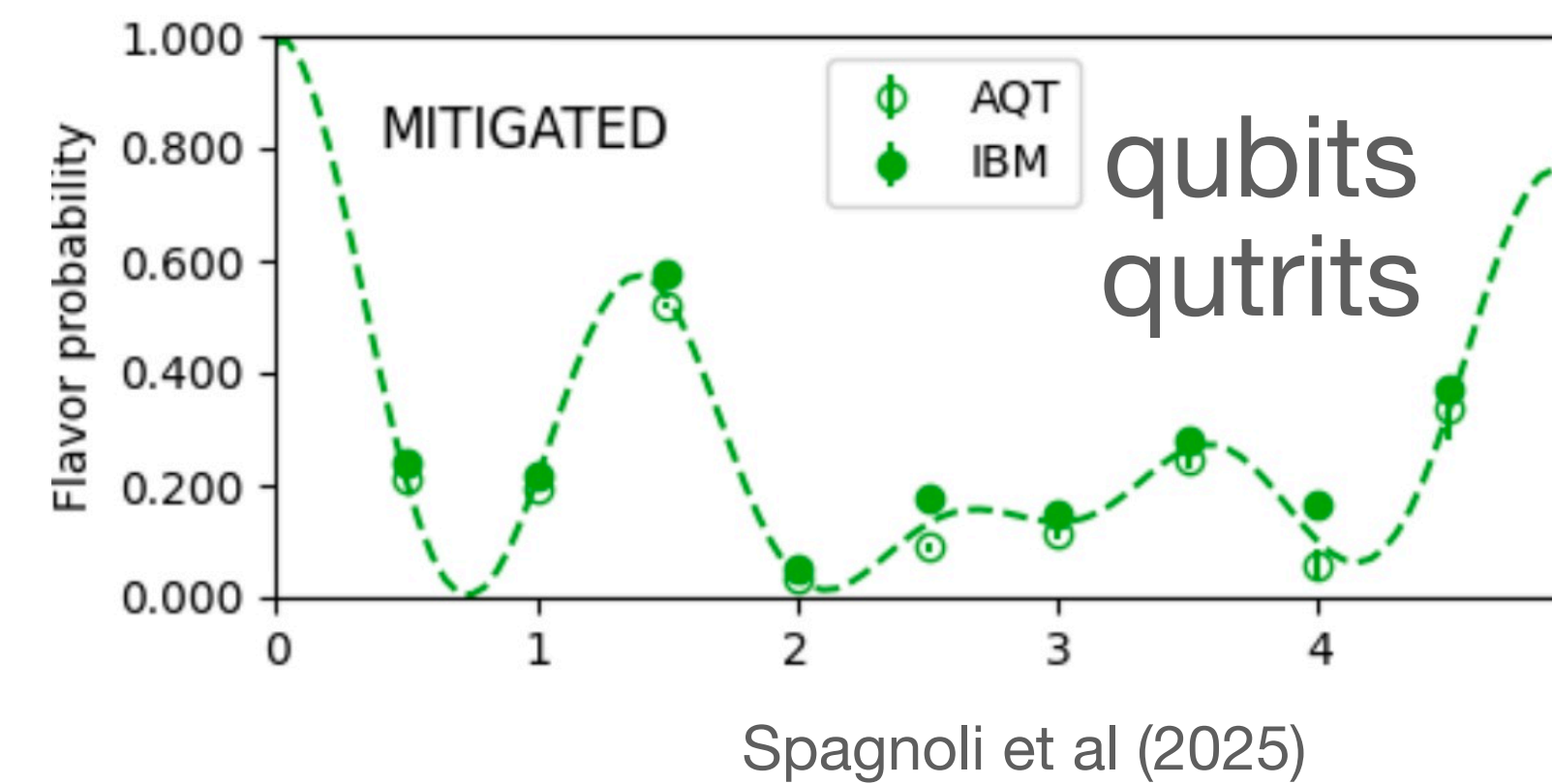
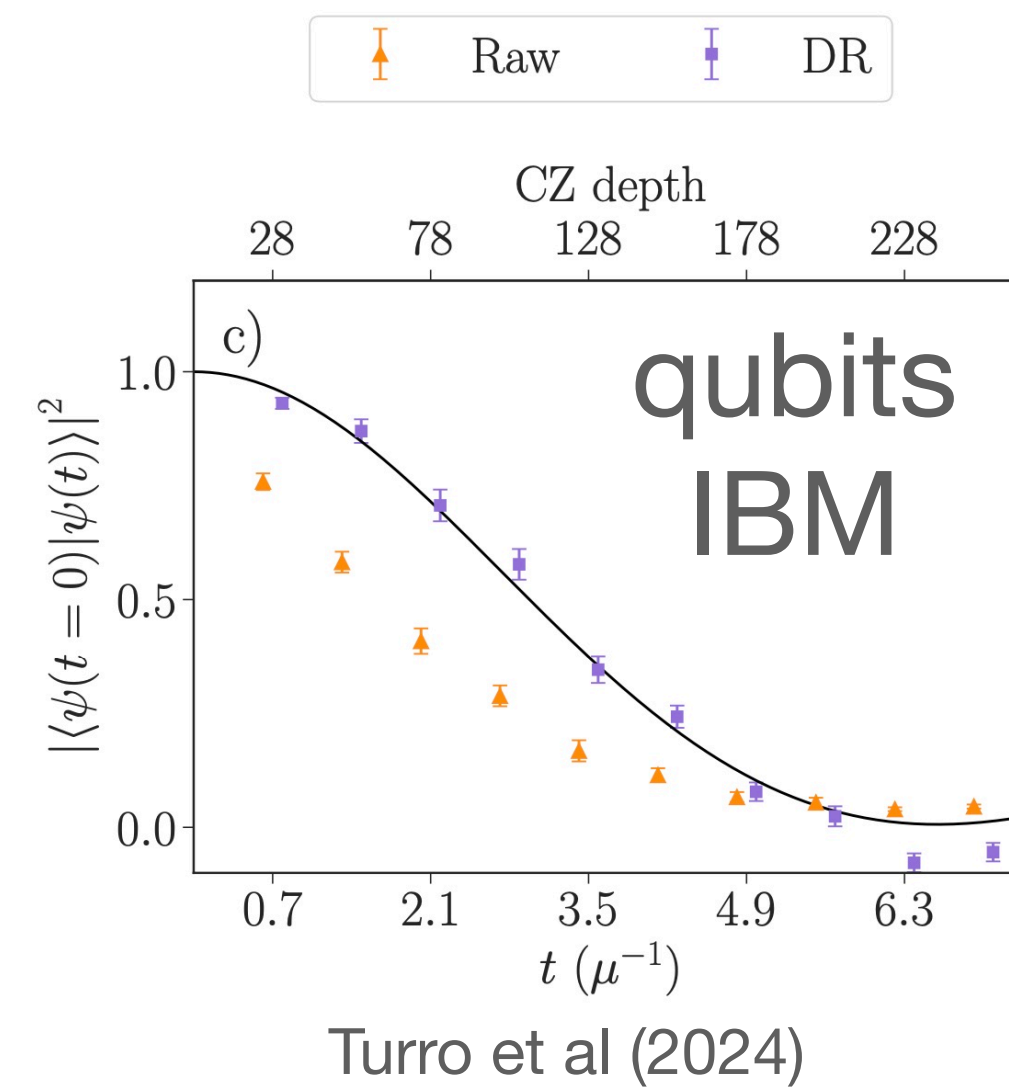
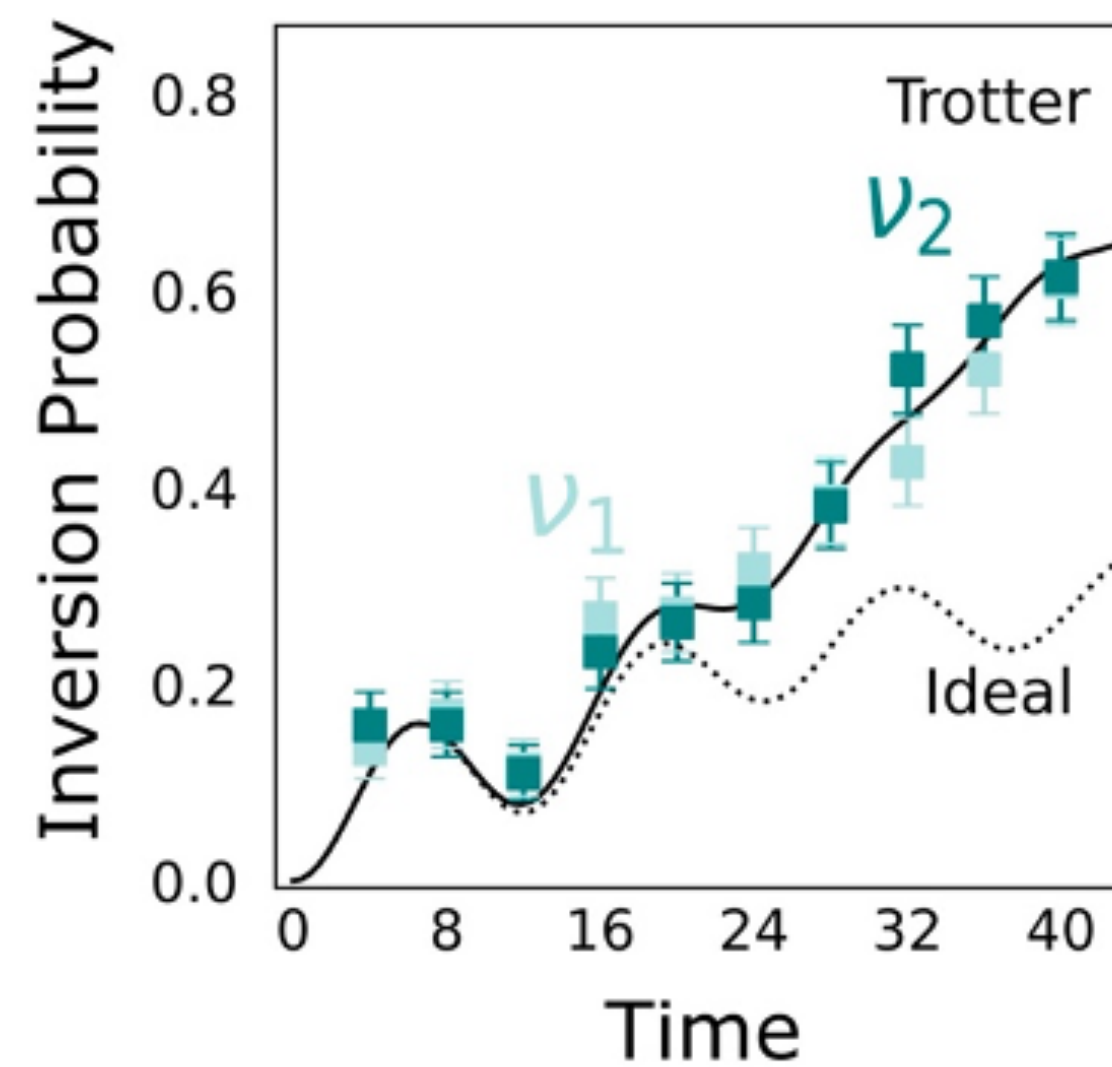
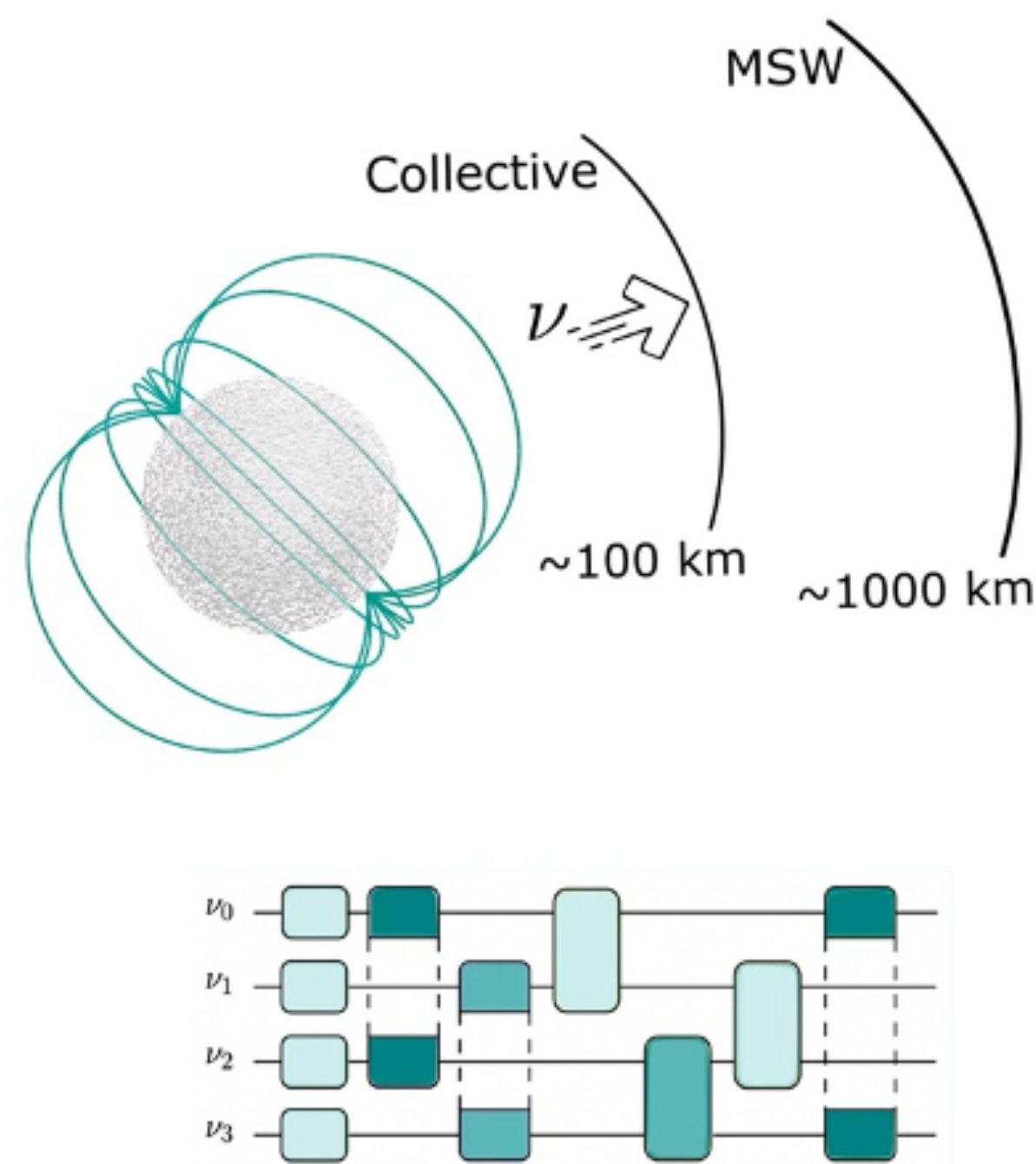
DOI: <https://doi.org/10.1103/PhysRevC.111.034317>



$$H_{FS} = - \sum_{k=1}^N \frac{\omega_k}{2} \sigma_k^z + \frac{\mu}{2N} \sum_{i < j}^N \mathcal{J}_{ij} \vec{\sigma}_i \cdot \vec{\sigma}_j$$

Neutrino Flavor Dynamics in Supernova

Simulations with qubits and qutrits





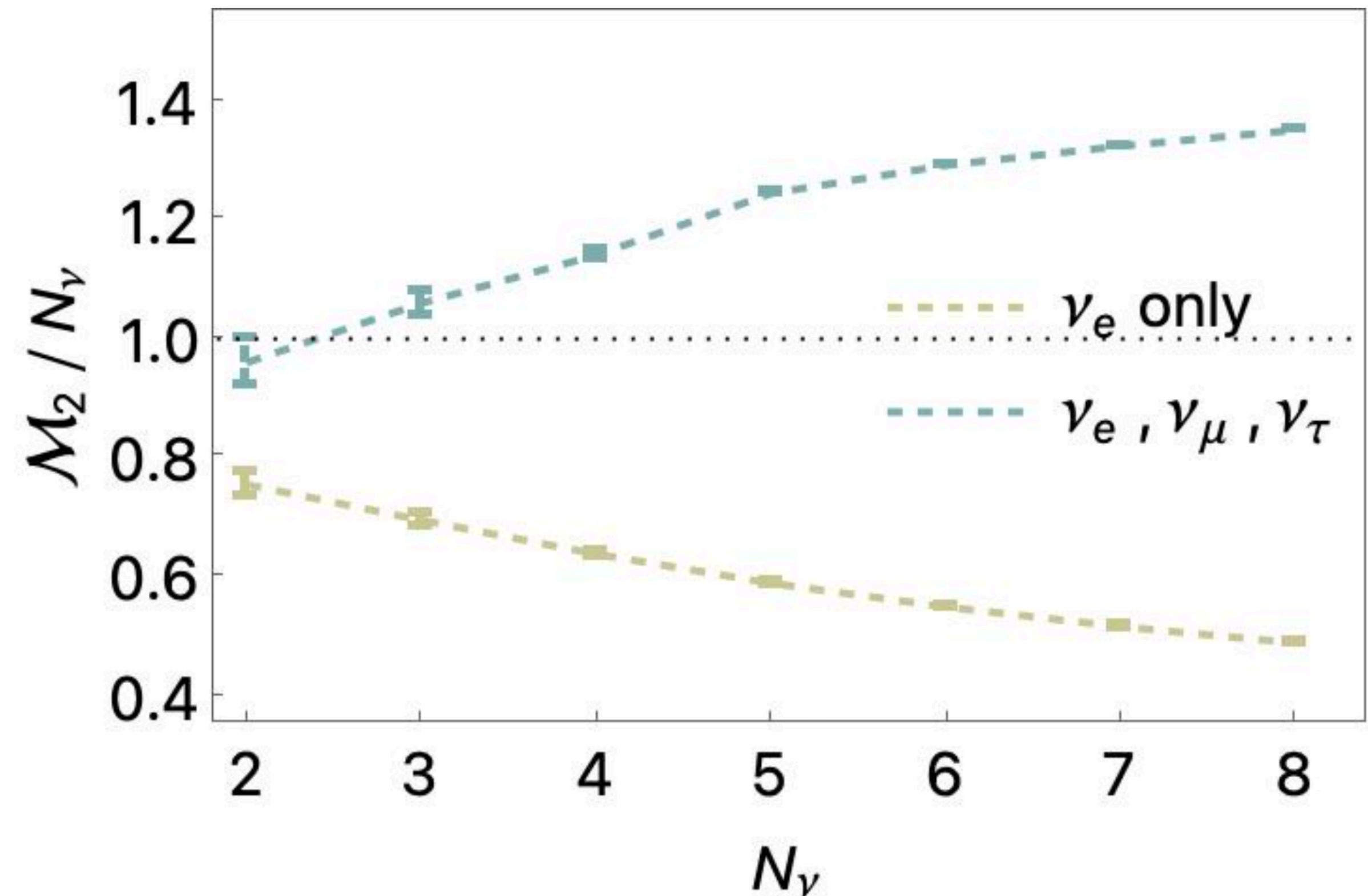
Quantum magic and computational complexity in the neutrino sector

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DOI: <https://doi.org/10.1103/PhysRevResearch.7.023228>

while entangled states can support a maximum value of $M_2 = 2.23379$.

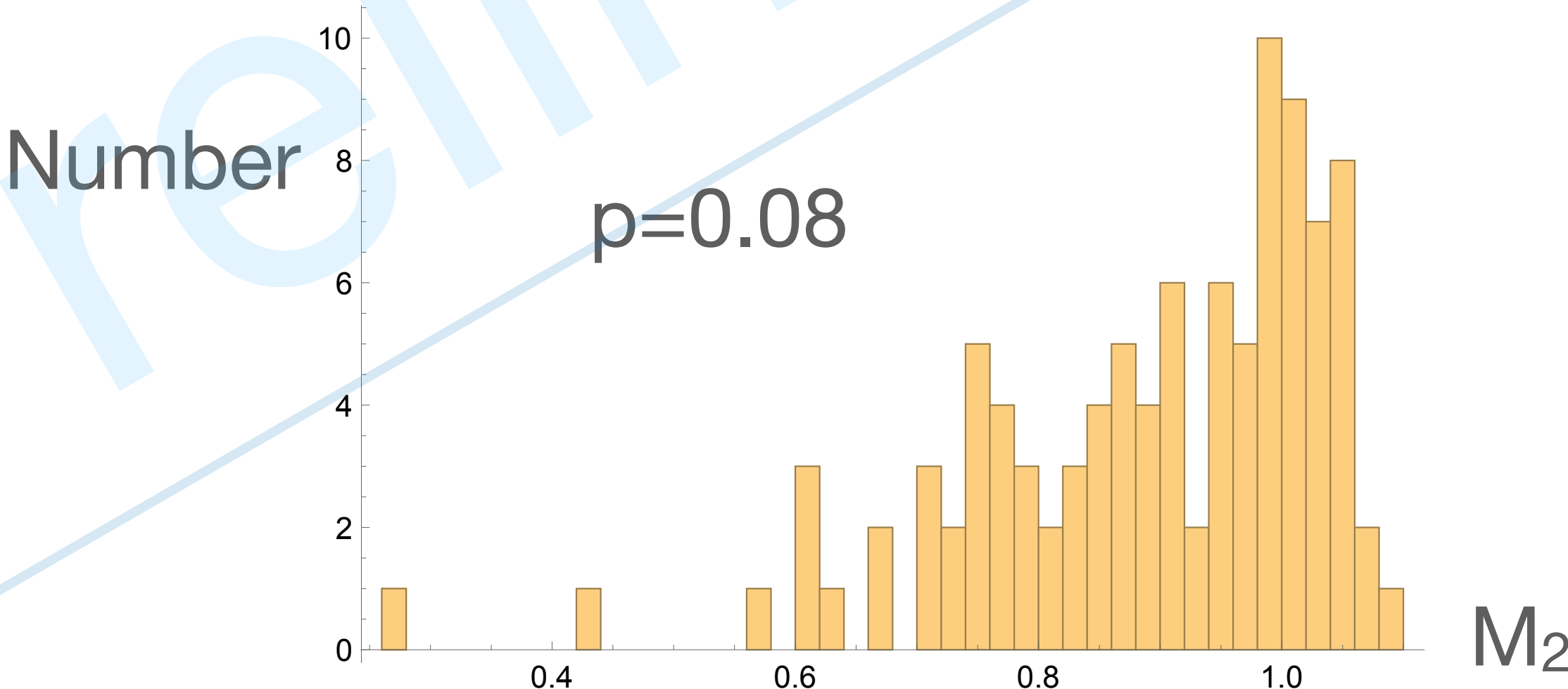
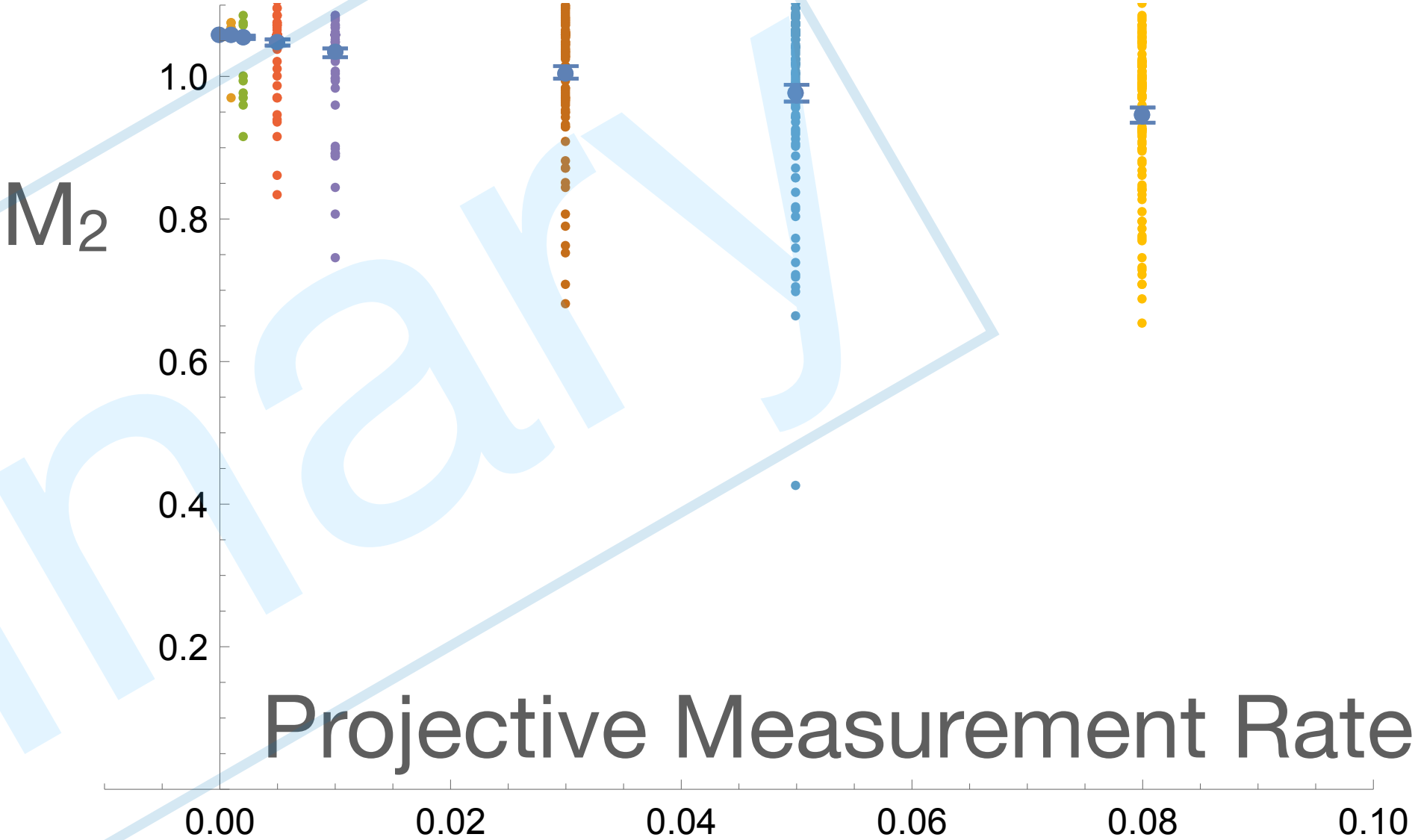
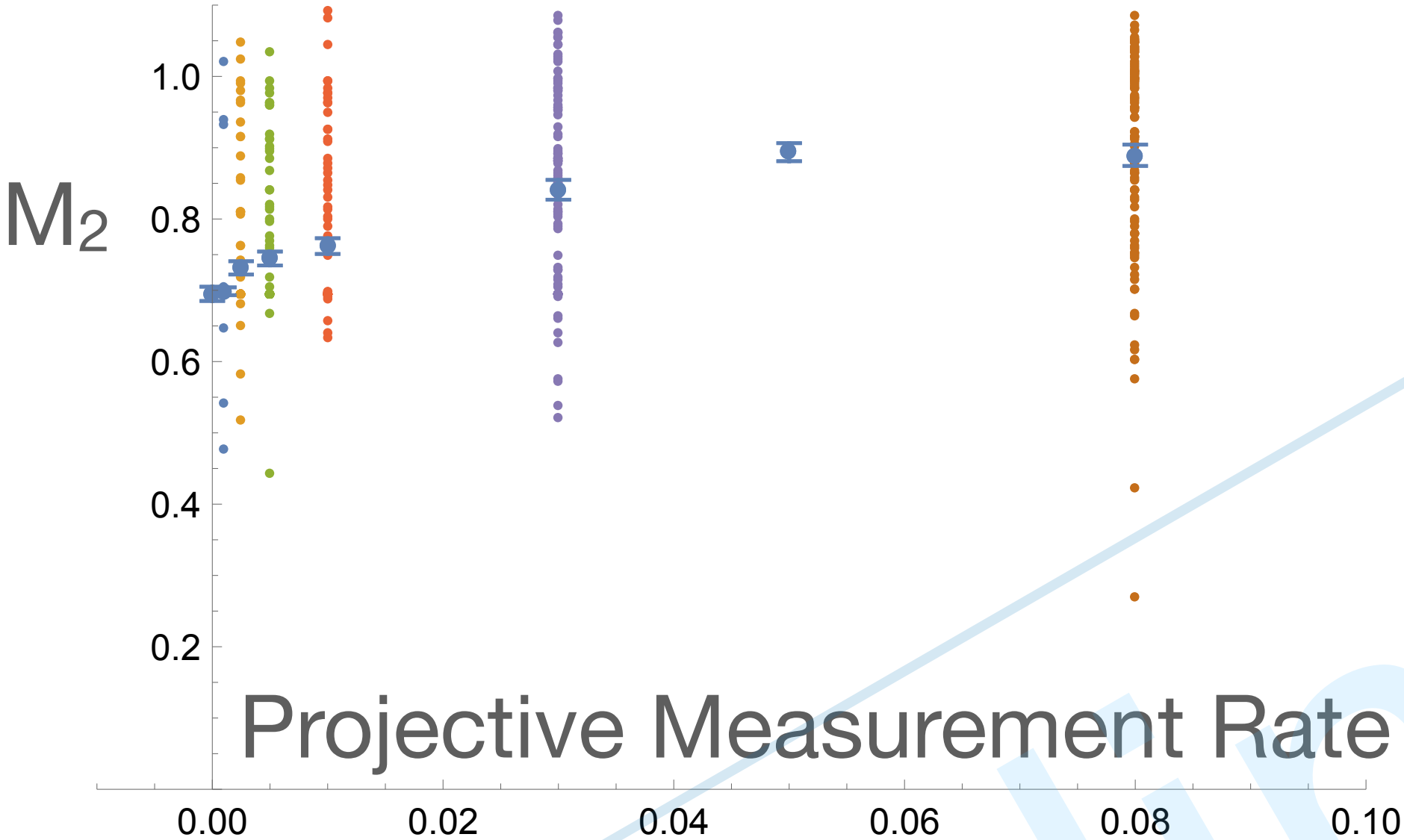
$$\hat{X}|j\rangle = |j+1\rangle \quad , \quad \hat{Z}|j\rangle = \omega^j |j\rangle \quad , \quad \omega = e^{i2\pi/3}$$

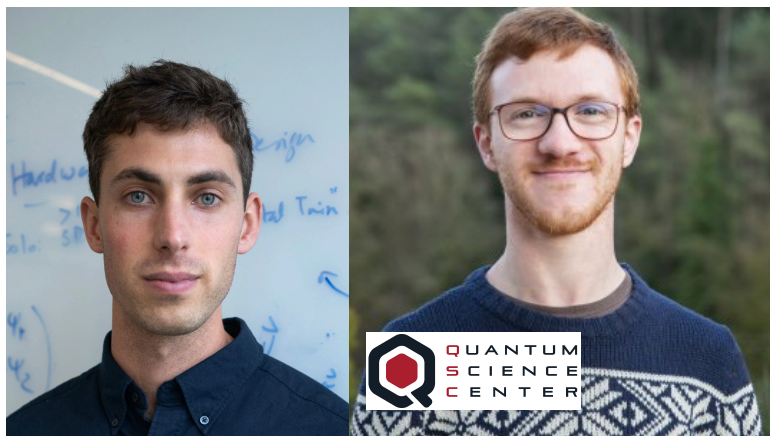




Marc Illa and Caroline Robin

Quantum Complexity in Neutrinos





Lorentz Violation by Lattice Spacing

Steps toward quantum simulations of hadronization and energy loss in dense matter

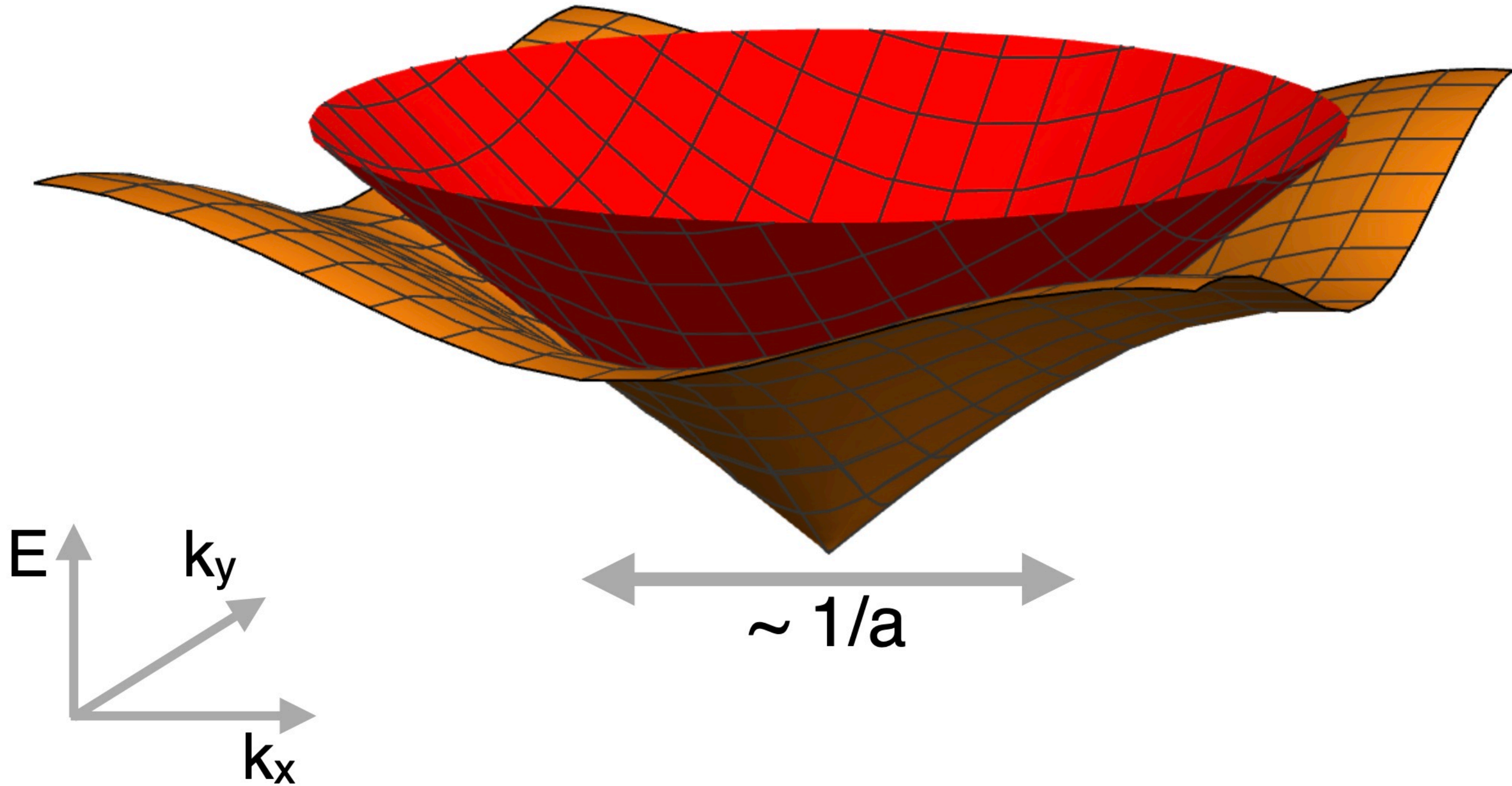
[Roland C. Farrell](#) ^{1,2,*}, [Marc Illa](#) ^{1,†}, and [Martin J. Savage](#) ^{1,‡,§}

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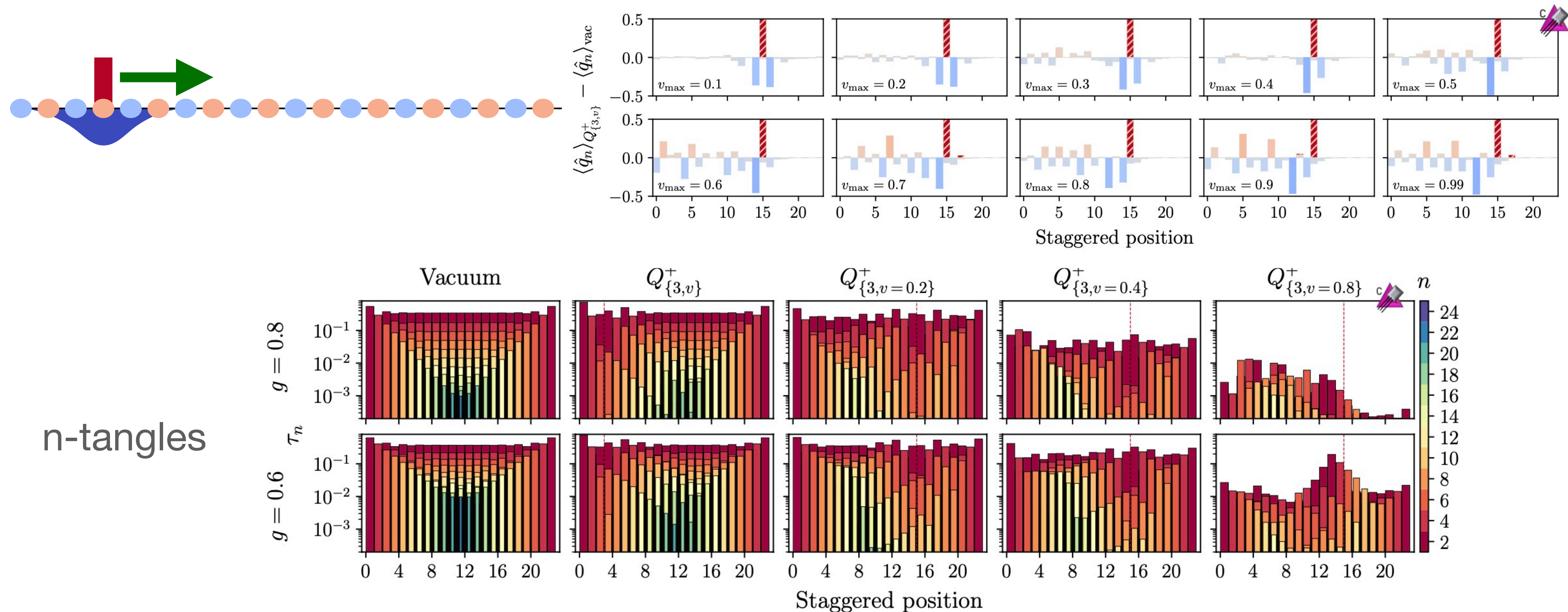
Phys. Rev. C **111**, 015202 – Published 14 January, 2025

DOI: <https://doi.org/10.1103/PhysRevC.111.015202>

Dispersion Relation

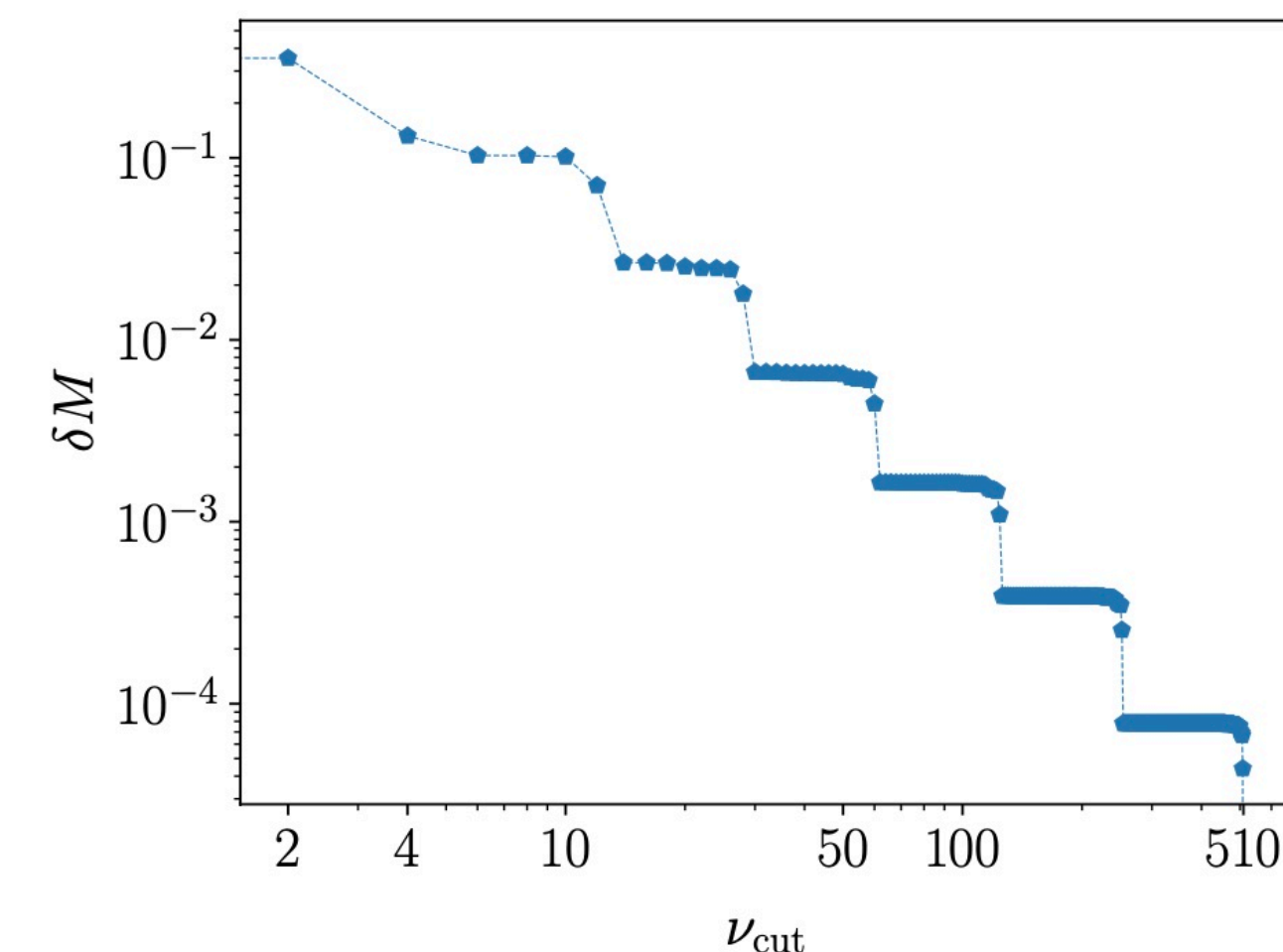
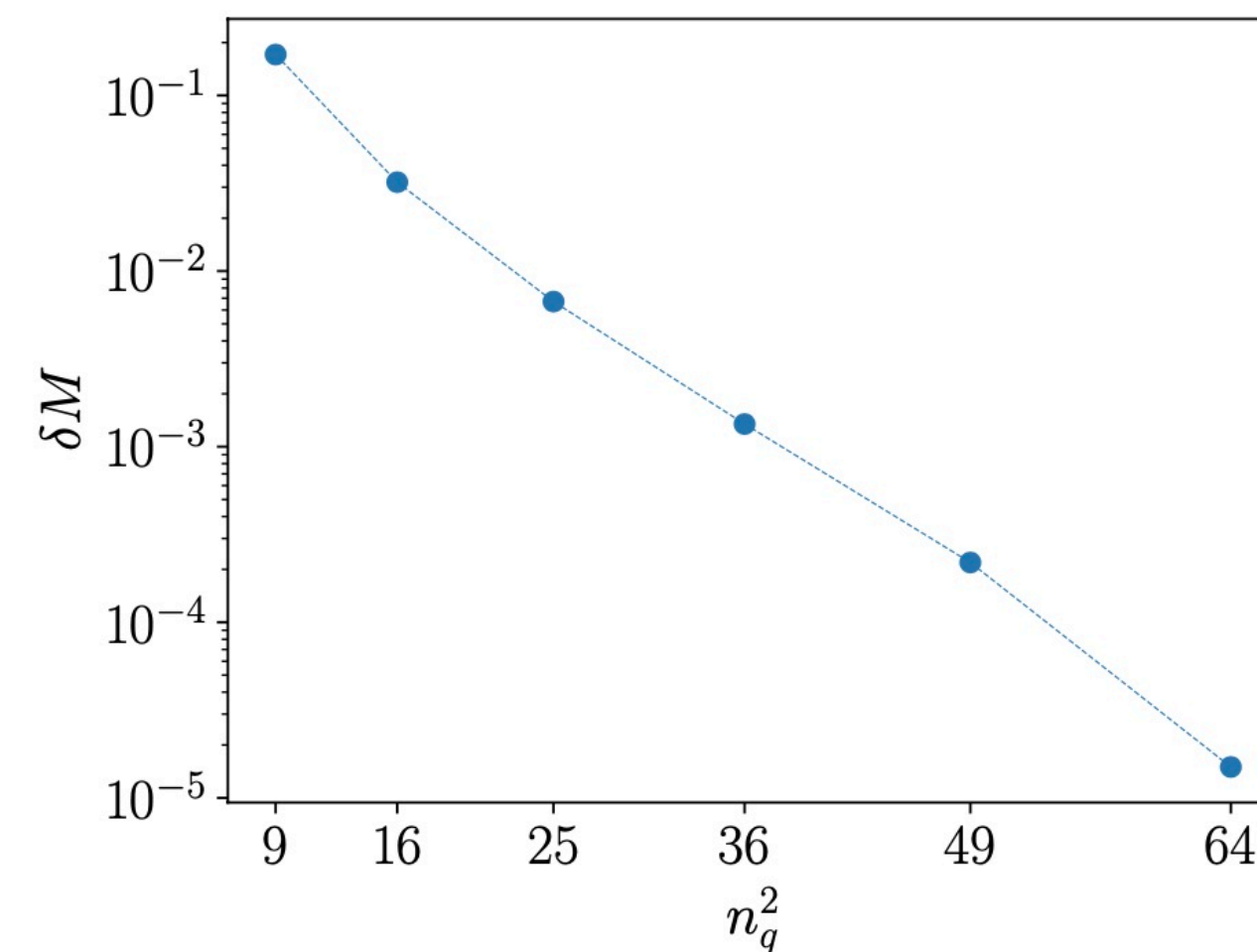
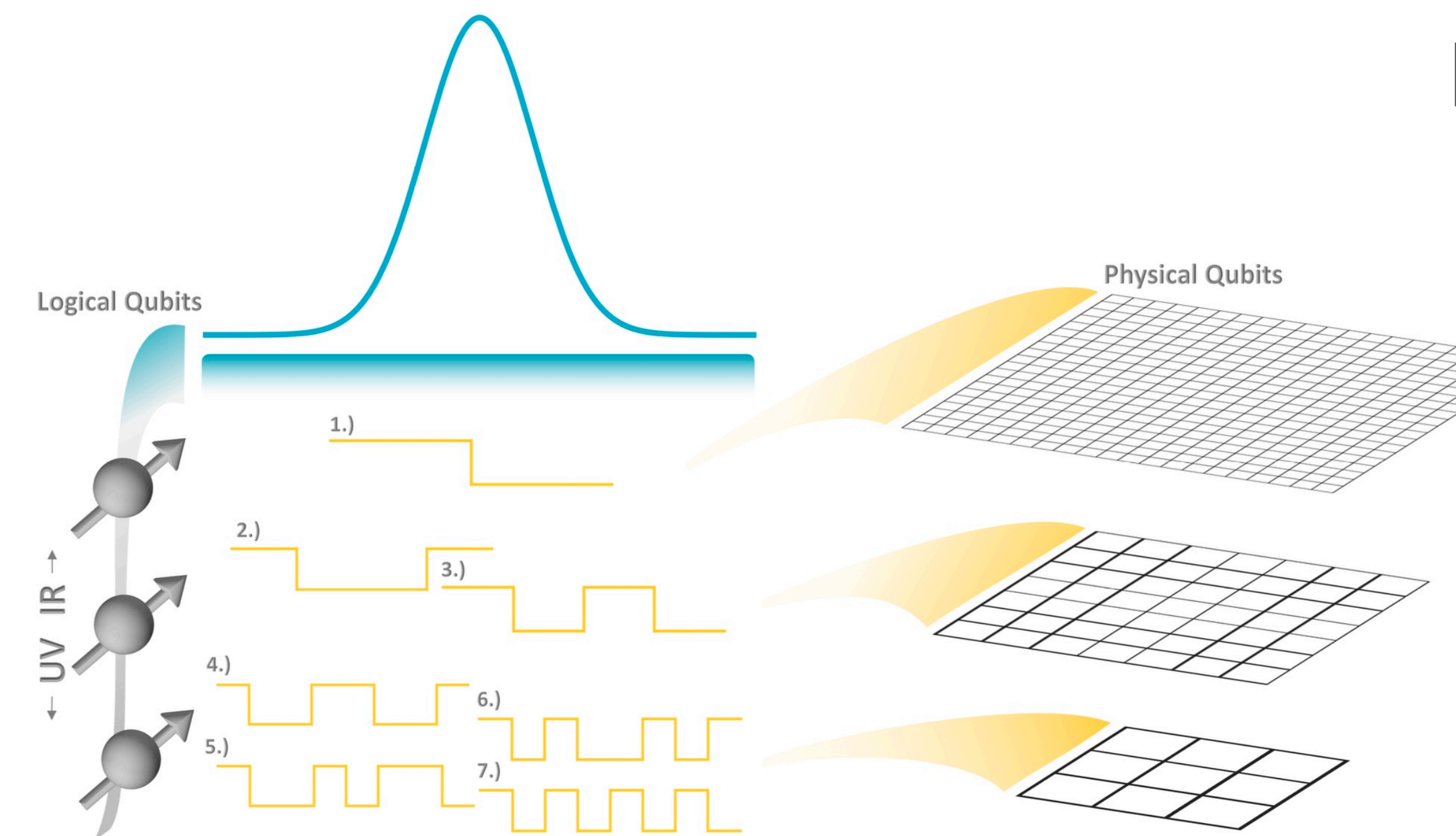


Lorentz Violation by the Lattice Spacing

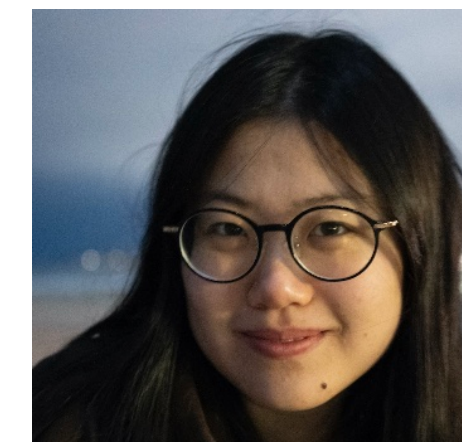


From NISQ to Fault Tolerant Sequency Hierarchies

Magic in a Gaussian: $\mathcal{M}_{\text{lin}} = 0.362007$



UV versus IR qubits



Sequency Hierarchy Truncation (SeqHT) for Adiabatic State Preparation and Time Evolution in Quantum Simulations

Hierarchical qubit maps and hierarchically implemented quantum error correction

Natalie Klco and Martin J. Savage
Phys. Rev. A **104**, 062425 – Published 15 December 2021

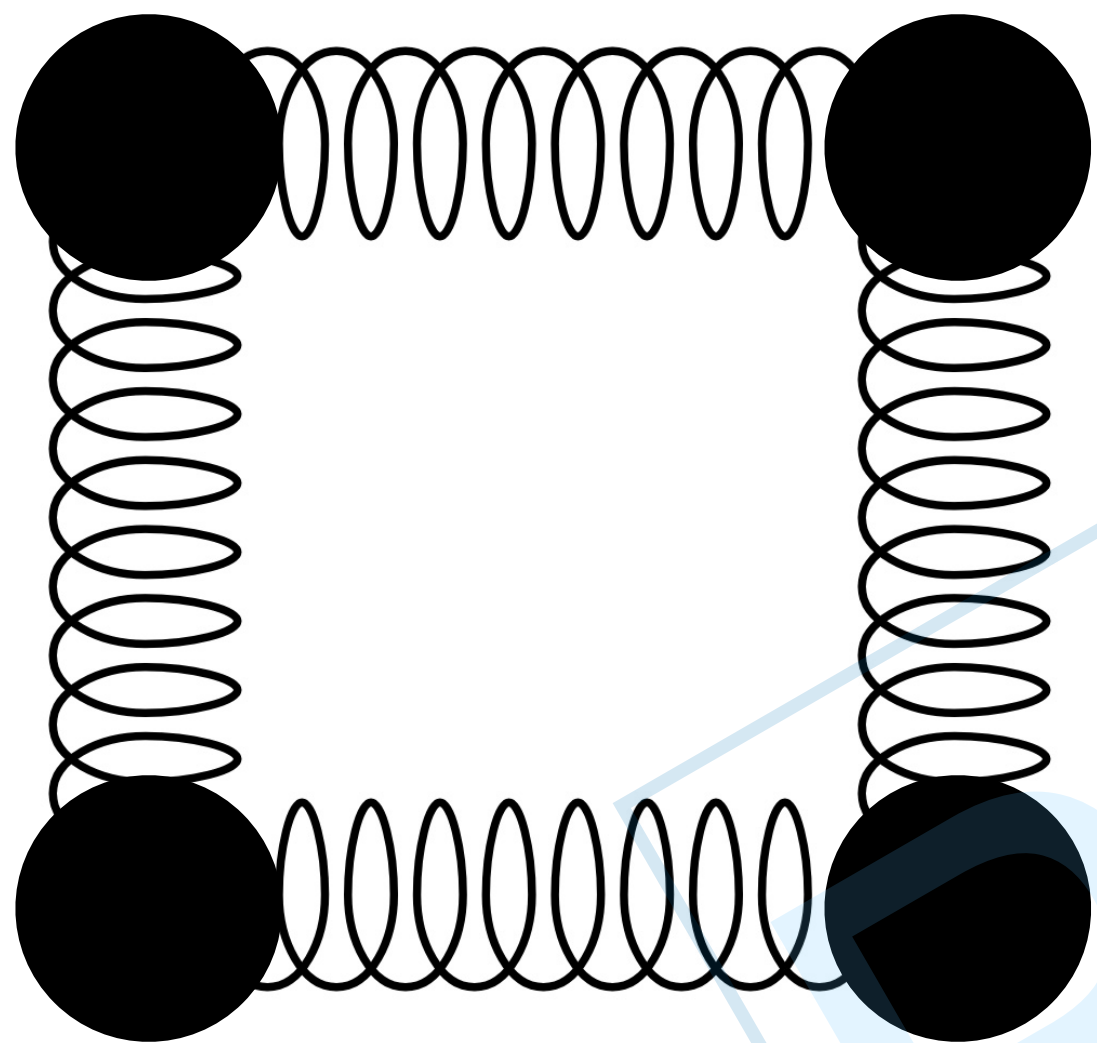
• e-Print: [2407.13835](https://arxiv.org/abs/2407.13835) [quant-ph] Zhiyao Li ,* Dorota M. Grabowska ,† and Martin J. Savage ,‡



1-Plaquette of SU(2) Yang-Mills Magic to Guide Truncations (?)

Marc Illa + MJS

$$\hat{H} = \frac{g^2}{2} \sum_{\mathbf{a}, \text{links}} |\mathbf{E}^{\mathbf{a}}|^2 + \frac{1}{2g^2} \left(4 - \hat{\square} - \hat{\square}^\dagger \right) \quad H_{j,j'} = \frac{1}{2} g^2 j(j+1) \delta_{j,j'} + \frac{1}{g^2} (2\delta_{j,j'} - \delta_{j+1,j'} - \delta_{j-1,j'})$$



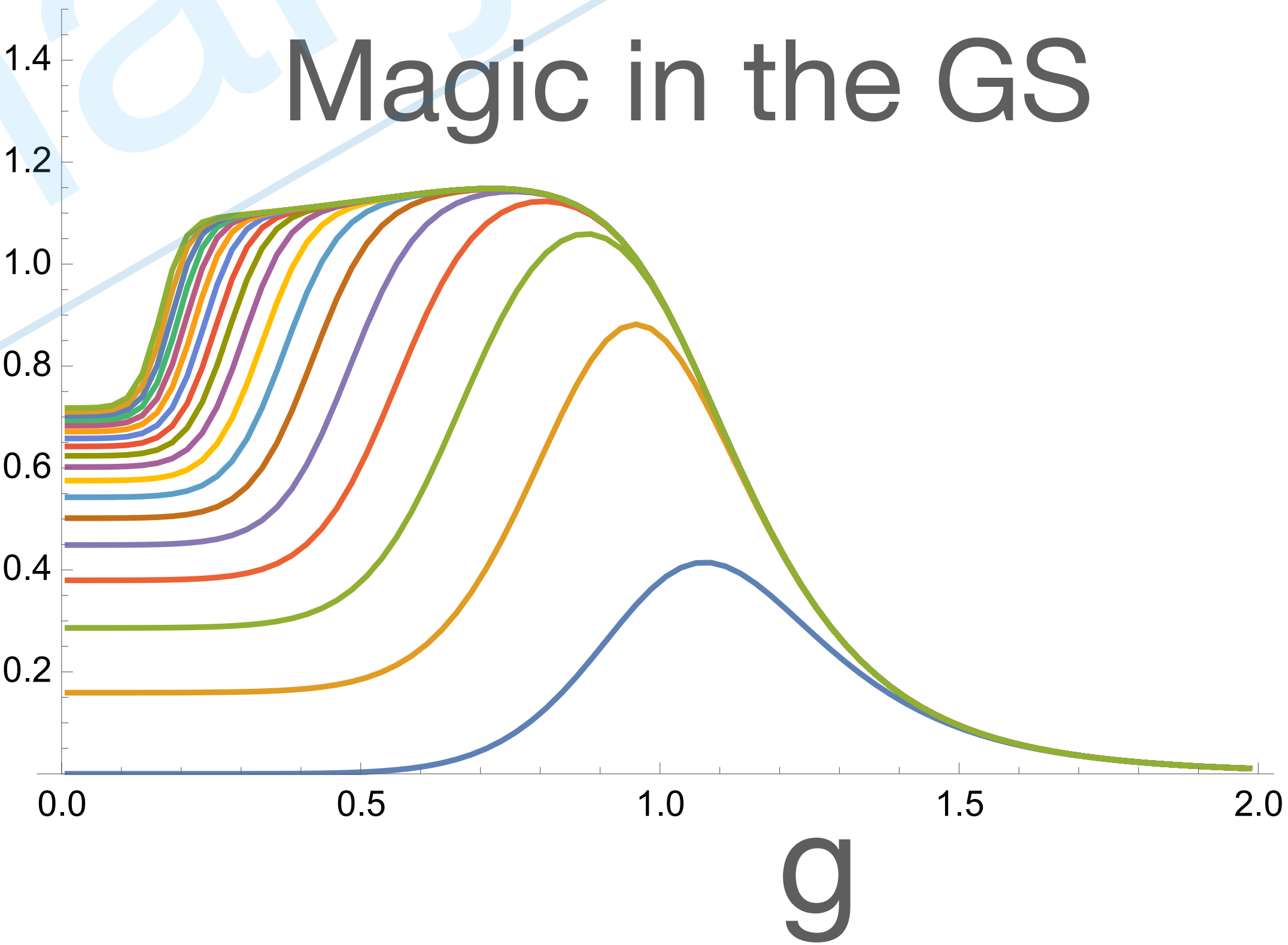
qudit



$j_{\text{Max}} = 1/2, 1, 3/2, 2, \dots$

M_2

Magic in the GS



Summary and Outlook

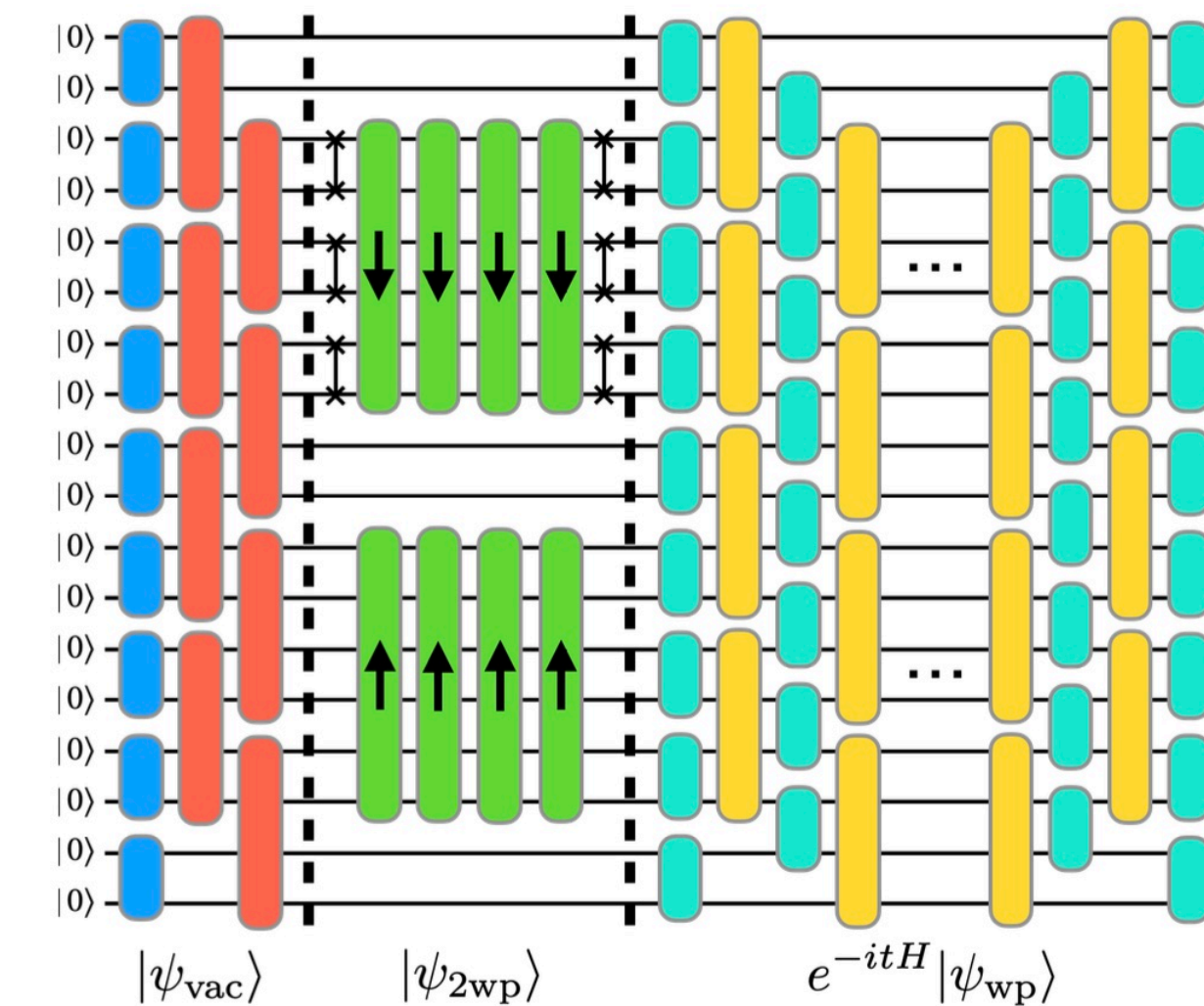
A unique time in the history of computing !!



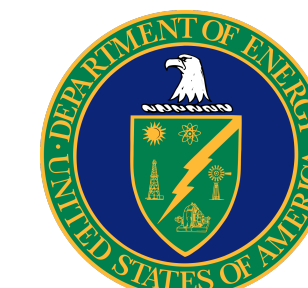
Quantum simulations focused on fundamental physics are advancing

Evolution from NISQ to fault tolerant/error correcting

Understanding how to organize quantum complexity

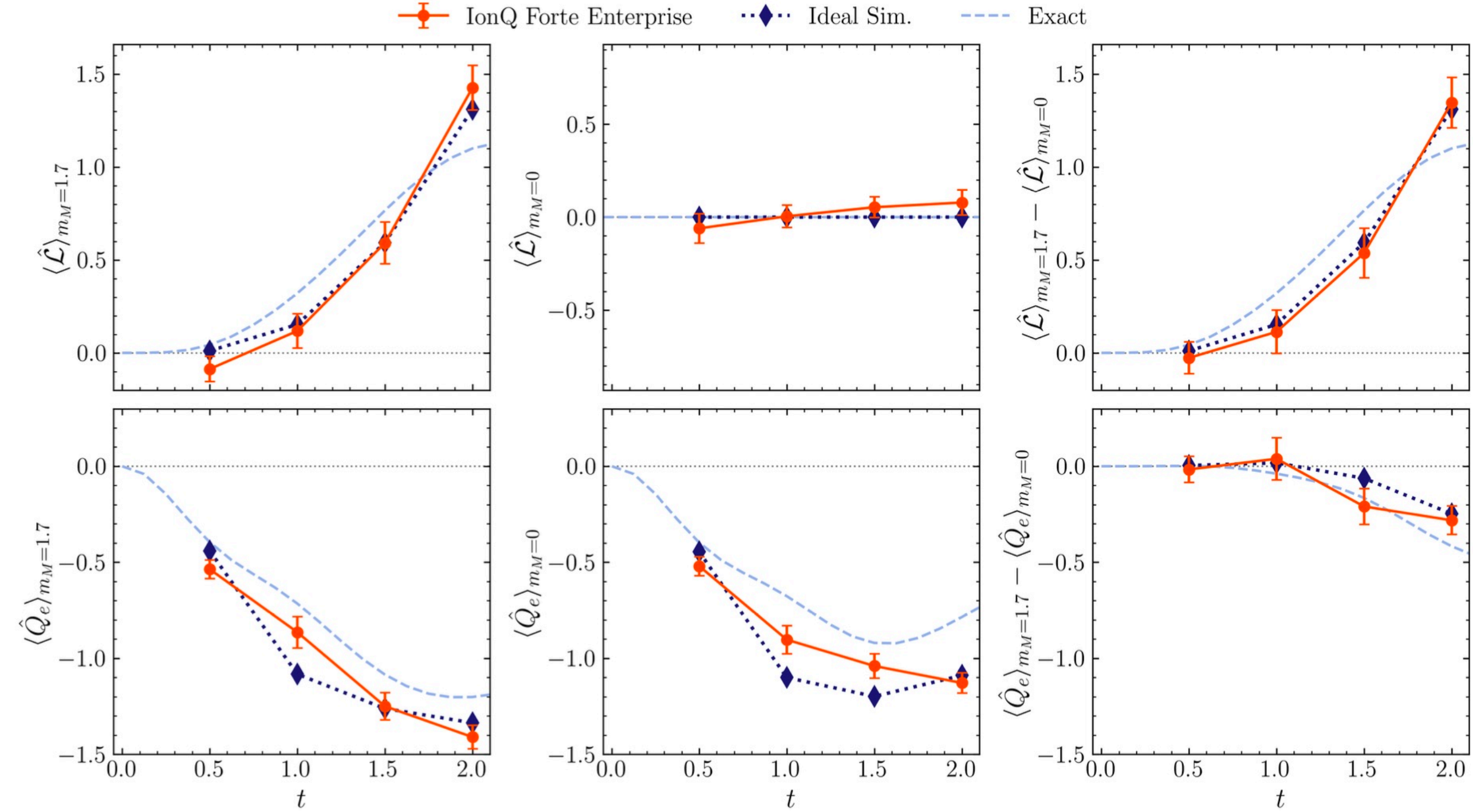
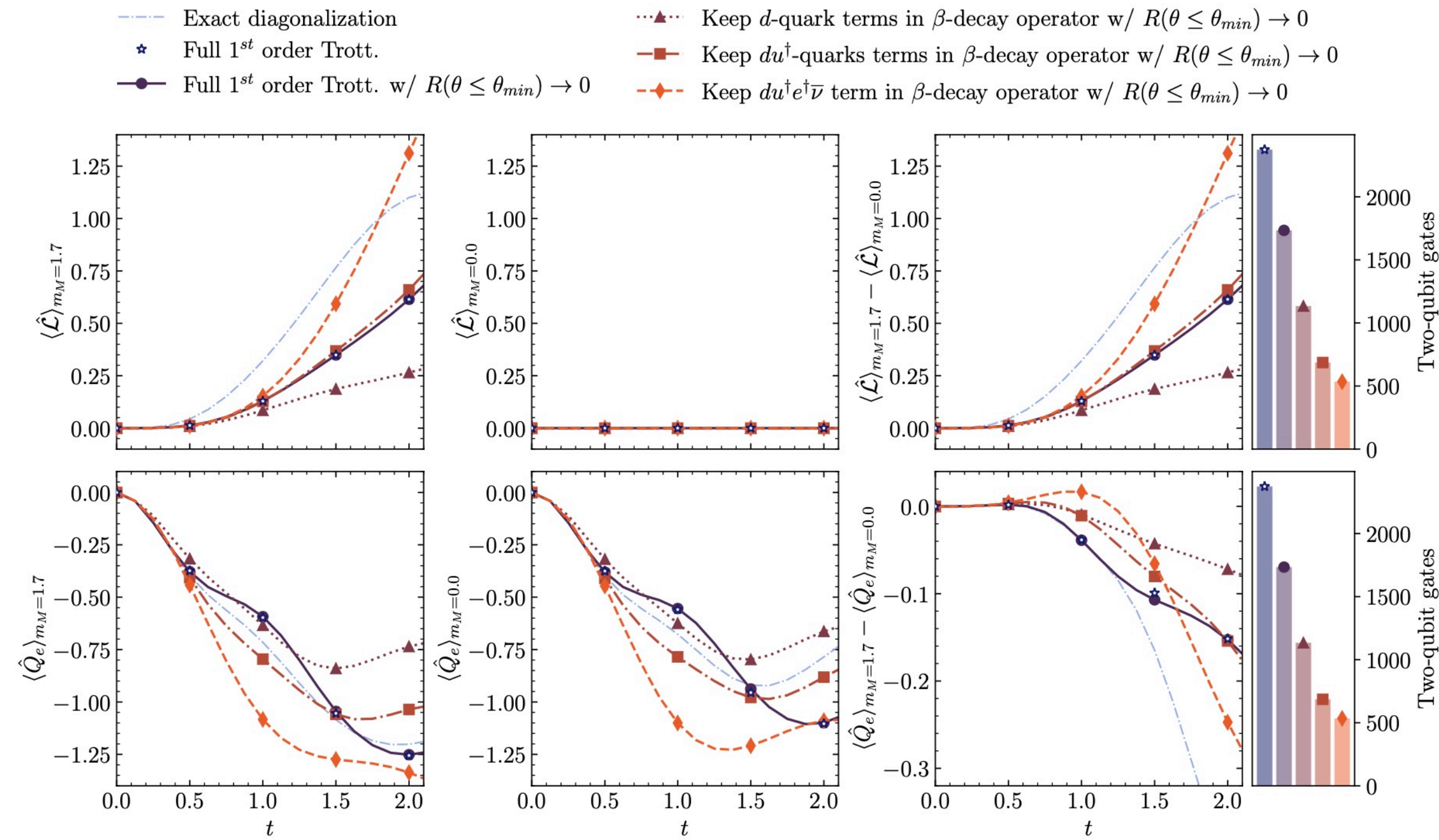


(Multiple) focused teams with access to forefront quantum hardware are essential



FIN

Ivan A. Chernyshev ¹ Roland C. Farrell ² Marc Illa ³ Martin J. Savage ^{3,*} Andrii Maksymov ⁴ Felix Tripier ⁴ Miguel Angel Lopez-Ruiz ⁴ Andrew Arrasmith ⁴ Yvette de Sereville ⁴ Aharon Brodutch ⁴ Claudio Grotto ⁴ Ananth Kaushik ⁴ and Martin Roetteler ^{4,†}



>450 entangling gates