

SU(2) gauge theory with $N_f=2$
fundamental flavour:
A minimal template for model building ?

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in collaboration with

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Interdisciplinary approach to QCD-like composite dark matter

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**UNIVERSITY OF
PLYMOUTH**

Outline



- Introduction
- Benchmark Results :
 - ★ Setup
 - ★ Spectrum
- New Results :
 - ★ Scattering properties and resonances width
- Summary / Outlook

Introduction

Strongly Interacting Dark Matter

- ♦ Dark Matter is essential to our current understanding of the Universe evolution
- ♦ Properties :
 - ★ Long-lived
 - ★ Electrically neutral
 - ★ Interaction with the Standard Model are suppressed

Why the Lattice ?

- ♦ Lattice simulations provide insights in many strongly coupled theories
- ♦ The lattice can provide information on the dark sector in isolation:
 - ★ Low-lying spectrum
 - ★ Matrix element relevant for direct detection
 - ★ Production cross section ?
 - ★ Self interactions ?
- ♦ Price to pay:
 - ★ The uv completion needs to be fixed.

Motivations to study $SU(2)$ with $N_f=2$

- ◆ Our main original motivation: Composite Higgs / Technicolor framework [Cacciapaglia & Sannino 2014]
- ◆ Light asymmetric Dark Matter [Lewis et al.]
- ◆ SIMP mechanism [Hochberg et al.]

SU(2)_c with N_f=2 fundamental Dirac flavours

- ♦ SU(2) gauge theory with N_f = 2 Dirac fermions in the fundamental representation.

$$\mathcal{L} = -\frac{1}{4} F_{\mu\nu}^a F^{a\mu\nu} + i\bar{U}\gamma^\mu D_\mu U + i\bar{D}\gamma^\mu D_\mu D + \frac{m}{2} Q^T (-i\sigma^2) C EQ + \frac{m}{2} (Q^T (-i\sigma^2) C EQ)^\dagger$$

- ♦ Pseudo-real irrep of SU(2): **global flavour symmetry is upgraded to SU(4)** :

$$Q \equiv \begin{pmatrix} U_L \\ D_L \\ \tilde{U}_L \\ \tilde{D}_L \end{pmatrix} \equiv \begin{pmatrix} U_L \\ D_L \\ -i\sigma_2 C \bar{u}_R^T \\ -i\sigma_2 C \bar{d}_R^T \end{pmatrix}, \quad E = \begin{pmatrix} 0 & 0 & 1 & 0 \\ 0 & 0 & 0 & 1 \\ -1 & 0 & 0 & 0 \\ 0 & -1 & 0 & 0 \end{pmatrix}$$

- ♦ Infinitesimal SU(4) transformation : $Q \longrightarrow \left(1 + i \sum_{n=1}^{15} \alpha^n T^n \right) Q$
- ♦ Generators that leaves the Lagrangian invariant satisfy : $ET^n + T^{nT}E = 0$
- ♦ Chiral symmetry breaking pattern : **SU(4) -> Sp(4)** (**SO(6) -> SO(5)**)
(5 Goldstone bosons)
- ♦ Wess-Zumino-Witten term allowed (3->2 vertex in EFT)

Hadronic operators

♦ Meson operators:

$$\mathcal{O}_{\bar{U}D}^{(\Gamma)} \equiv \bar{U}(x)\Gamma D(x),$$

$$\mathcal{O}_{\bar{D}U}^{(\Gamma)} \equiv \bar{D}(x)\Gamma U(x),$$

$$\mathcal{O}_{\bar{U}U \pm \bar{D}D}^{(\Gamma)} \equiv \frac{1}{\sqrt{2}} \left(\bar{U}(x)\Gamma U(x) \pm \bar{D}(x)\Gamma D(x) \right),$$

♦ Baryon operators:

$$\mathcal{O}_{UD}^{(\Gamma)} \equiv U^T(x)(-i\sigma^2)C\Gamma D(x)$$

$$\mathcal{O}_{DU}^{(\Gamma)} \equiv D^T(x)(-i\sigma^2)C\Gamma U(x)$$

♦ Large euclidean time behaviour :

$$\sum_{\mathbf{x}} \langle O(\mathbf{x}, t) O(\mathbf{0}, 0) \rangle = \sum_n \langle 0|O|n \rangle \langle n|O|0 \rangle e^{-m_n t}$$

What do we know ?

- ♦ Most notably : spectrum investigations at non vanishing fermions mass
 - ★ Goldstone Bosons
 - ★ Vector and Axial mesons
 - ★ Scalar and Pseudo-scalar mesons
 - ★ Pseudo-scalar decay constant

- ♦ In this talk :
 - ★ Review spectrum predictions
 - ★ Goldstone Boson scattering and the vector meson resonance

- ♦ Can we bring useful constraints in the context of :
 - ★ Dark pion produced by through Dark vector mesons
 - ★ SIMP mechanism
 - ★ Self-interacting DM ?

Benchmark results

The setup

R. Lewis, C. Pica, F. Sannino, Phys.Rev. D85 (2012) 014504 [arXiv:1109.3513]

A. Hietanen, C. Pica, R. Lewis, F. Sannino, JHEP 1407 (2014) 116 [arXiv:1404.2794]

A. Hietanen, C. Pica, R. Lewis, F. Sannino [arXiv:1308.4130]

R. Arthur, V.D, A. Hietanen, M. Hansen, C. Pica, F. Sannino [arXiv:1602.06559]

- ♦ Plaquette action + dynamical Wilson Fermions
- ♦ Several volumes $V=L^3 \times T$ \Rightarrow extrapolate to **infinite volume**
- ♦ 4 lattice spacings : a \Rightarrow extrapolate to the **continuum limit**
- ♦ Several fermion masses $m_f \longleftrightarrow m_{PS^2}$ \Rightarrow extrapolate to the **chiral limit**
- ♦ Non-perturbative renormalisation
- ♦ **HiRep** code

L. Del Debbio, A. Patella, C. Pica, Phys.Rev. D81 (2010) 094503

Chiral behavior : GB sector

J. Bijnens and J. Lu, JHEP **11** (2009) 116, [arXiv:0910.5424]

♦ Continuum χ PT:

$$\frac{m_{\text{ps}}^2}{m_f} = 2B \left[1 + \frac{3}{4}x \log \frac{2Bm_f}{\mu^2} + b_M x + \mathcal{O}(x^2) \right]$$
$$f = F \left[1 - x \log \frac{2Bm_f}{\mu^2} + b_F x + \mathcal{O}(x^2) \right]$$
$$x = \frac{2Bm_f}{(4\pi F)^2}$$

♦ Range of applicability: unknown *a priori*

♦ In terms of $\tilde{x} = \frac{m_{\text{ps}}^2}{(4\pi F)^2}$ the expressions are unchanged at NLO

♦ Assess discretisation effects:

✱ Strategy I:

- modelling the discretisation effects and doing a global fit:

$$\frac{m_{\text{ps}}^2}{m_f} = 2B \left[1 - a_M \tilde{x} \log \frac{m_{\text{ps}}^2}{\mu^2} + b_M \tilde{x} + \delta_M \frac{a}{w_0^\chi} + \gamma_M m_{\text{ps}}^2 \frac{a}{w_0^\chi} \right]$$
$$F_{\text{ps}} = F \left[1 - a_F \tilde{x} \log \frac{m_{\text{ps}}^2}{\mu^2} + b_F \tilde{x} + \delta_F \frac{a}{w_0^\chi} + \gamma_F m_{\text{ps}}^2 \frac{a}{w_0^\chi} \right]$$

- Define four data subsets to control fit stability

✱ Strategy II:

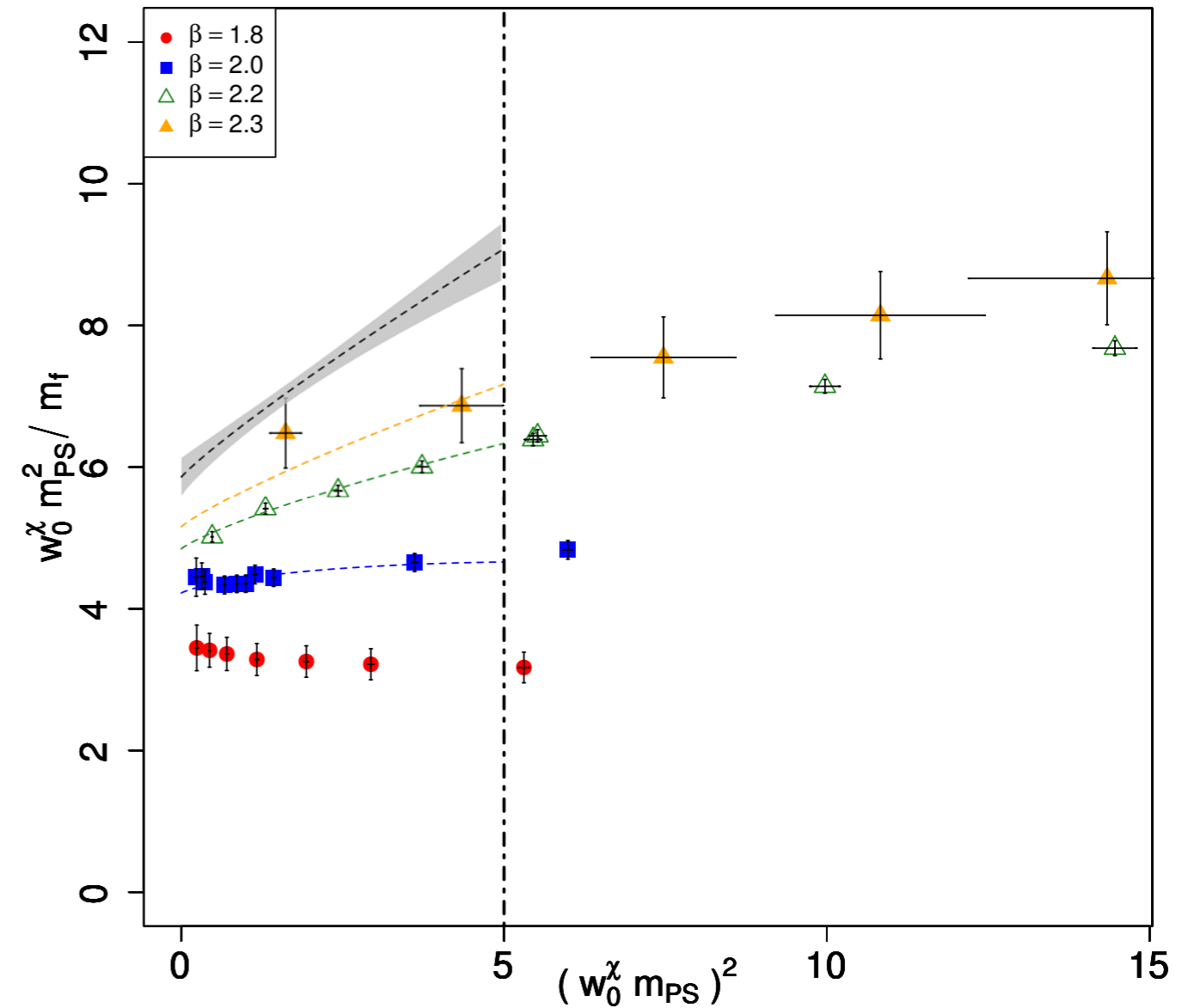
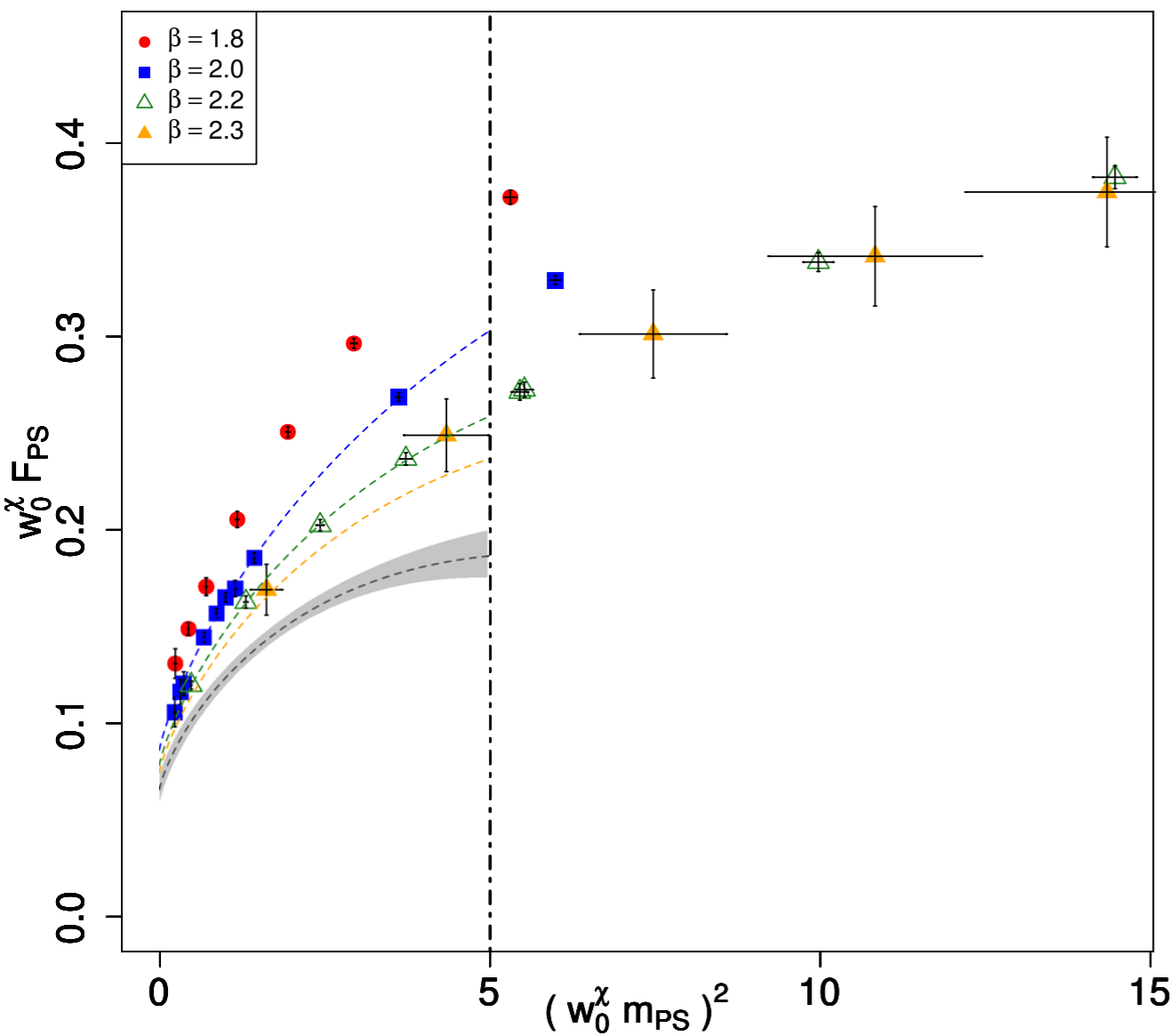
- set $\delta_{M,F}$ & $\gamma_{M,F}$ to zero

- fit each lattice spacings independently

- study fit parameters as a function of the lattice spacing

Goldstone bosons : mass and decay constant

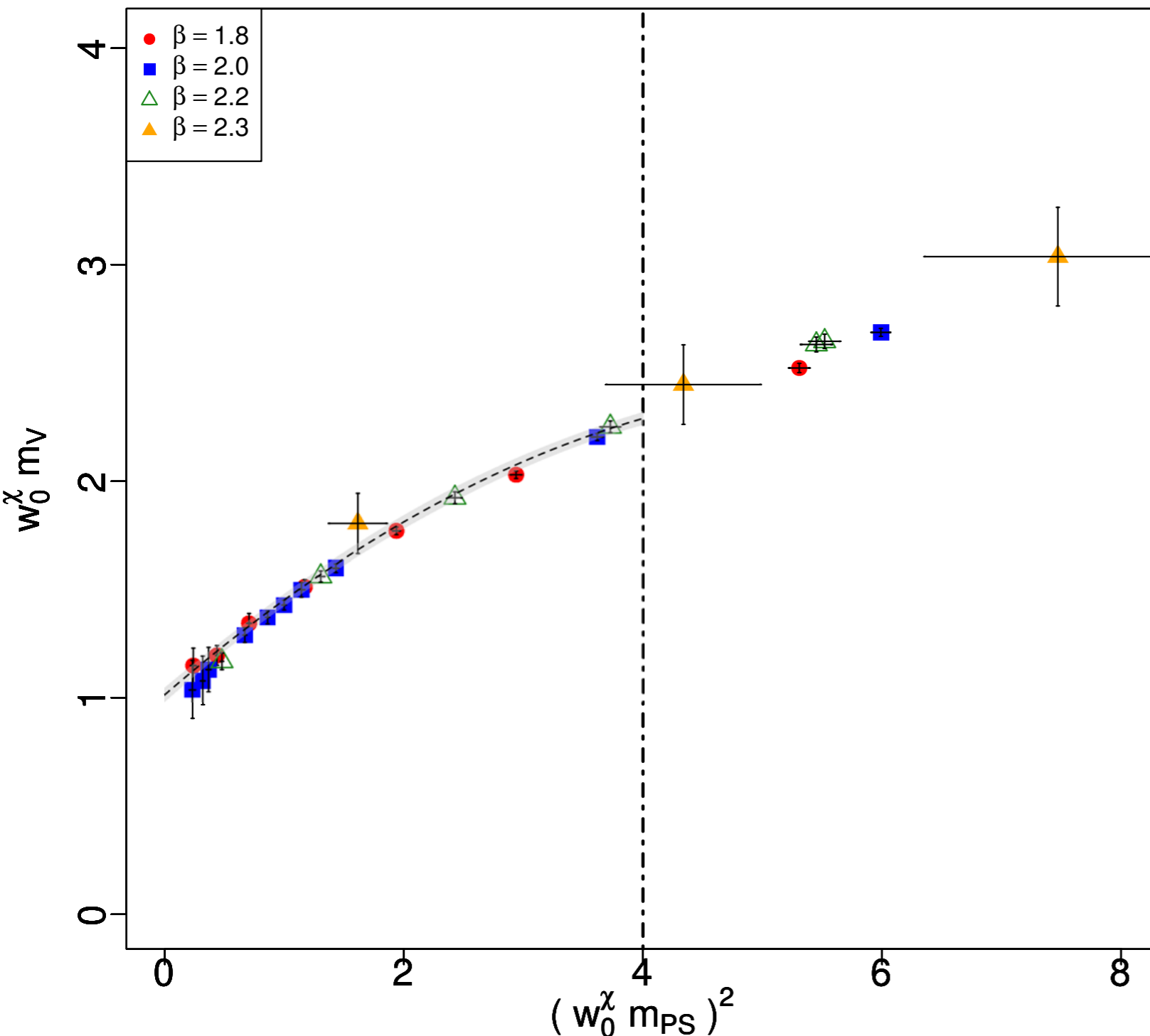
R. Arthur, VD, A. Hietanen, C. Pica, F. Sannino [arXiv:1602.06559]



- Large cut-off effects

Vector meson

R. Arthur, VD, A. Hietanen, C. Pica, F Sannino [arXiv:1602.06559]

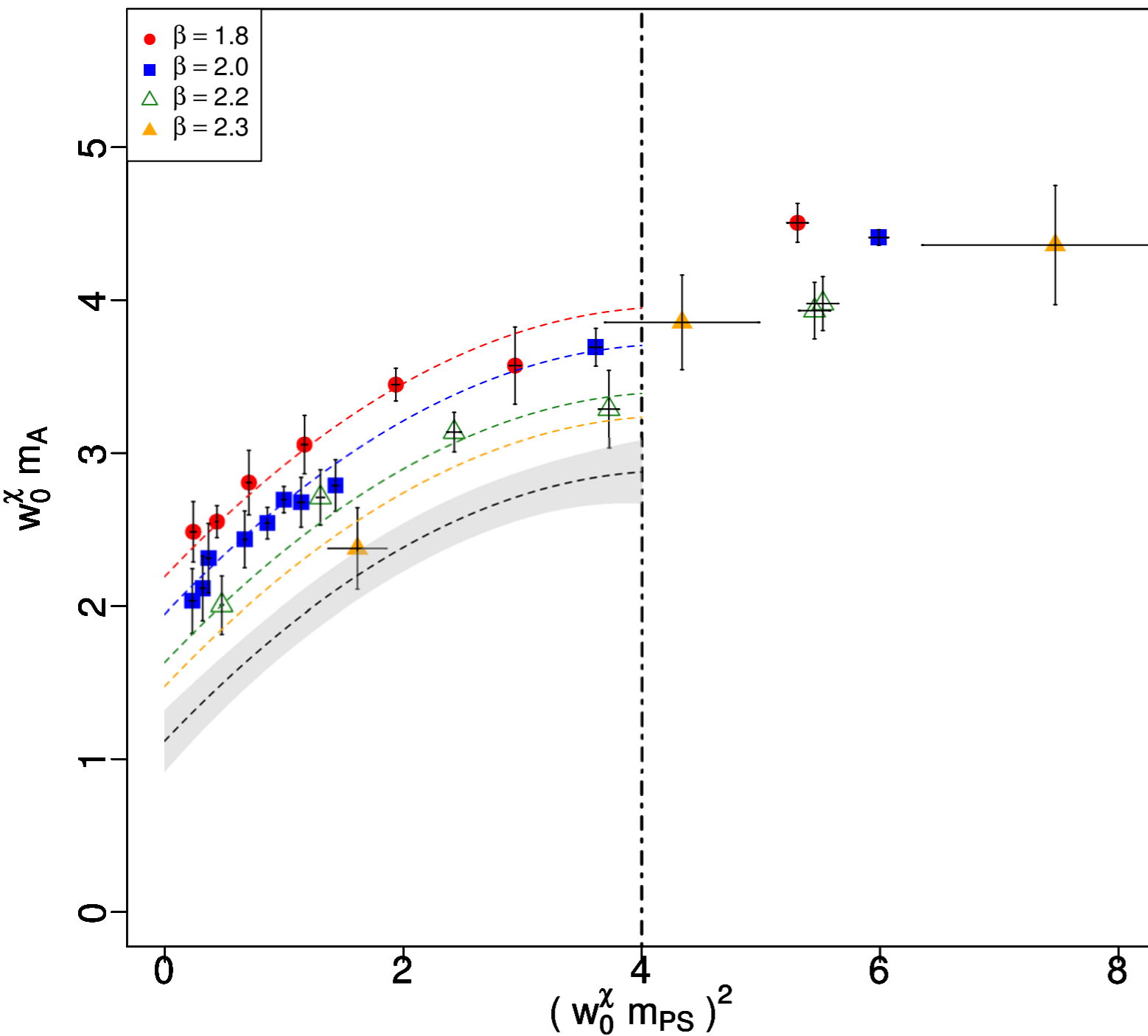


- Vector meson stable in the regime of our simulations ($m_V < 2 m_{PS}$)
- All 4 lattice spacings are consistent
- $m_V = 13.1(2.2) F_{PS}$

$$w_0^\chi m_X = w_0^\chi m_X^\chi + A(w_0^\chi m_{ps})^2 + B(w_0^\chi m_{ps})^4 + C \frac{a}{w_0}$$

Axial meson

R. Arthur, VD, A. Hietanen, C. Pica, F Sannino [arXiv:1602.06559]



- Axial meson stable in the regime of our simulations ($m_A < 2 m_{PS}$)
- Significant cut-off effects
- $m_A = 14.5(3.6) F_{PS}$

General remarks

- ♦ Scalar sector might be the way to discover new physics
 - ♦ In the real world :
 - ◉ $m_\sigma/f_{PS} \sim 5$ and width is large
 - ◉ The ratio depends a lot on the quark mass
 - ♦ Here :
 - ◉ 0^{++} is not light because of symmetry reasons
 - ◉ N_c dependence is not know
 - ♦ Technical issues :
 - ◉ Rigorous treatment of resonances on the lattice is expensive and difficult (see later)
 - ◉ Disconnected contributions are noisy
- Very challenging !**

Masses (II)

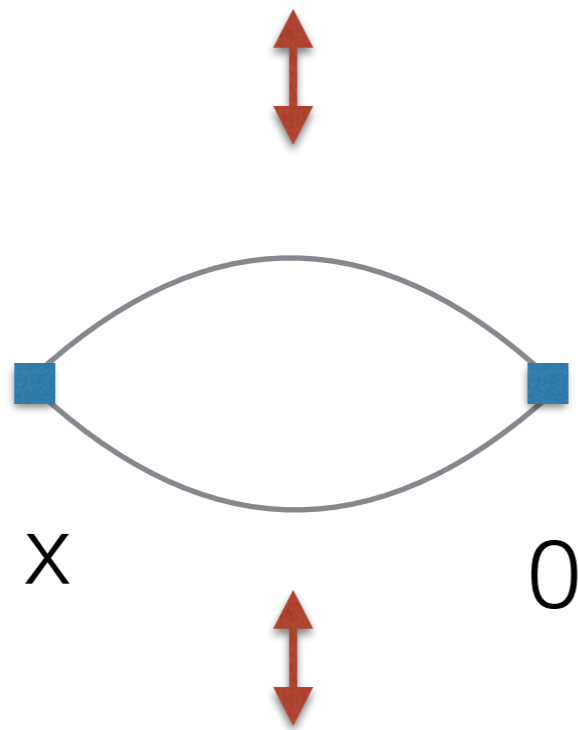
Forgetting about the fact that the σ is unstable

♦ Simplest fermionic interpolating field : $J = \bar{u}u + \bar{d}d$ (iso-singlet, σ)

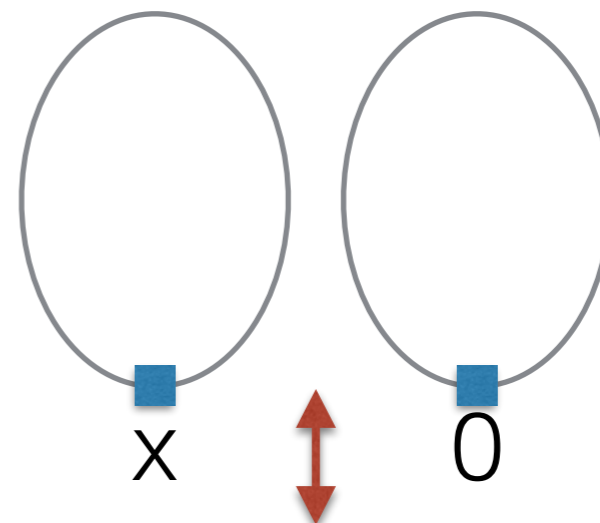
♦ Correlator : $C_{2\text{pts}}(t) = \sum_{\vec{x}} \langle J(t, \vec{x}) \bar{J}(0) \rangle$

♦ After integration over the fermions :

$$C_{2\text{pts}}(t) = \sum_{\vec{x}} \text{tr}\{S(\mathbf{x}, 0)S(0, \mathbf{x})\} + N_f \sum_{\tilde{\mathbf{x}}} \text{tr}\{S(\mathbf{x}, \mathbf{x})\}\text{tr}\{S(0, 0)\}$$



=connected contribution



=disconnected contribution

« infinitely many gluons » exchanged

♦ connected contribution : two-point function $J = \bar{u}u - \bar{d}d$ (iso-vector, a_0)

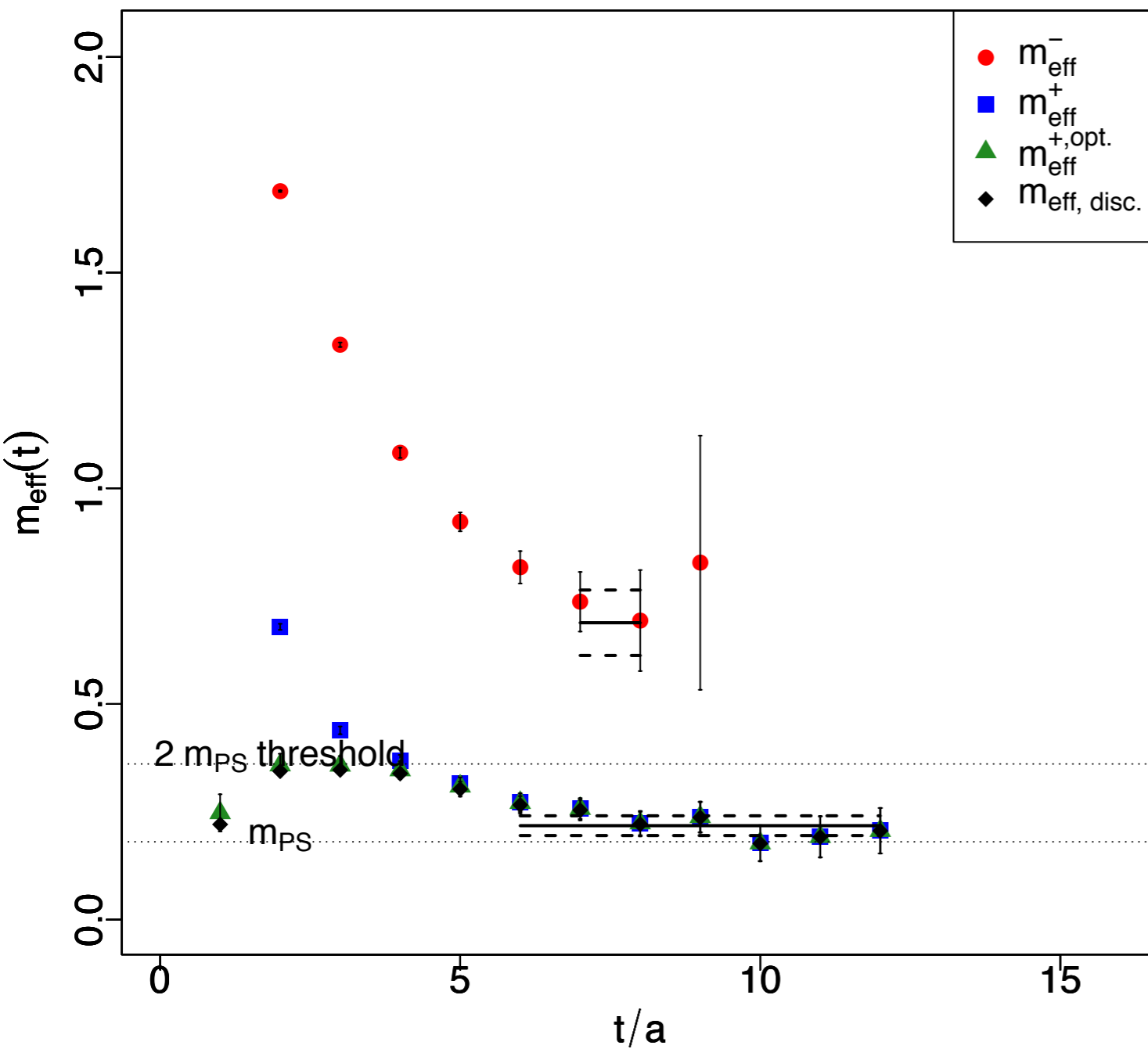
Masses (II)

- ♦ Disconnected contribution estimated using stochastic estimators (64 volume sources/conf.)
- ♦ Same calculation can be used to compute η'
- ♦ Interpolating field $J = \bar{u}\gamma_5 u + \bar{d}\gamma_5 d$ (iso-scalar)
- ♦ *not* a GB because of the $U(1)_A$ anomaly

results on 3 new states of the theory : σ, a_0, η'

Effective mass : scalar channel

R. Arthur, VD, A. Hietanen, C. Pica, F. Sannino [arXiv:1607.06654]

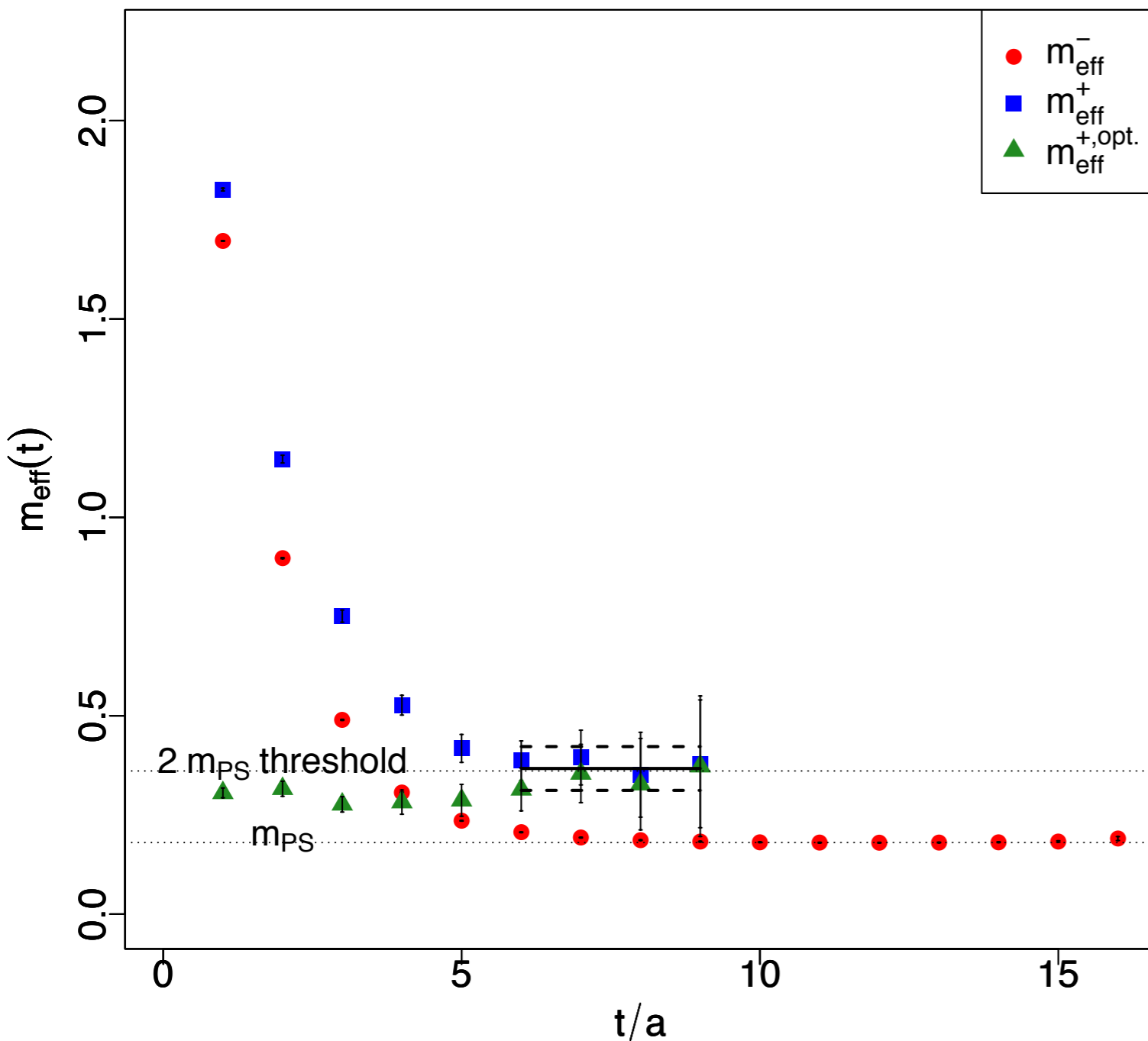


- Most chiral run: $\beta=2.0$, $m=-0.958$, $V=32^4$
- σ is stable in our setup
- $m_{\sigma} \sim m_{\text{PS}}$!

No need to consider the resonance analysis for the σ

Effective mass : pseudoscalar channel

R. Arthur, VD, A. Hietanen, C. Pica, F. Sannino [arXiv:1607.06654]

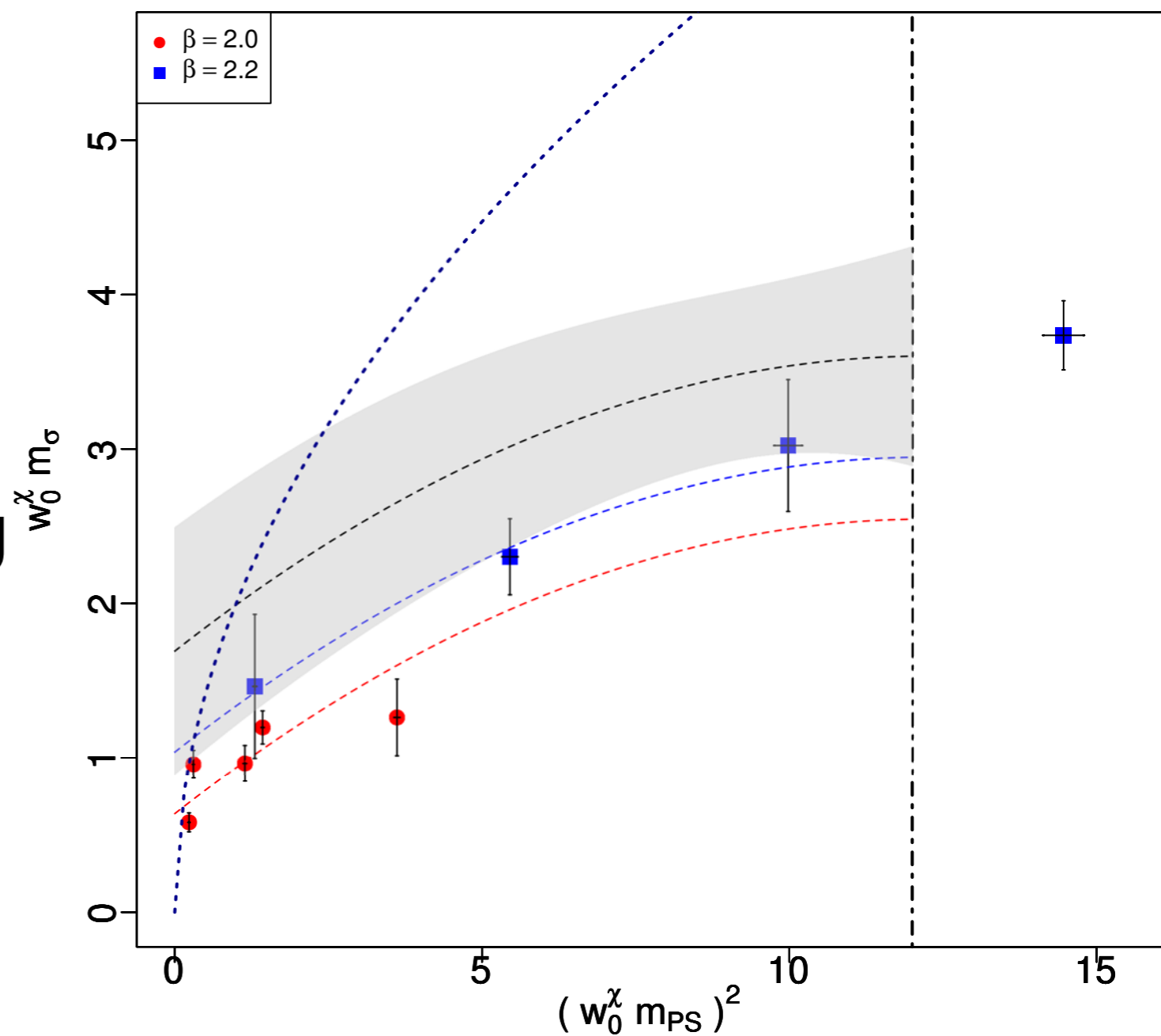


- Most chiral run: $\beta=2.0$, $m=-0.958$, $V=32^4$
- η' signal

Contribution from the disconnected diagrams is sizeable only for our lightest masses

σ resonance : $0(0^+)$

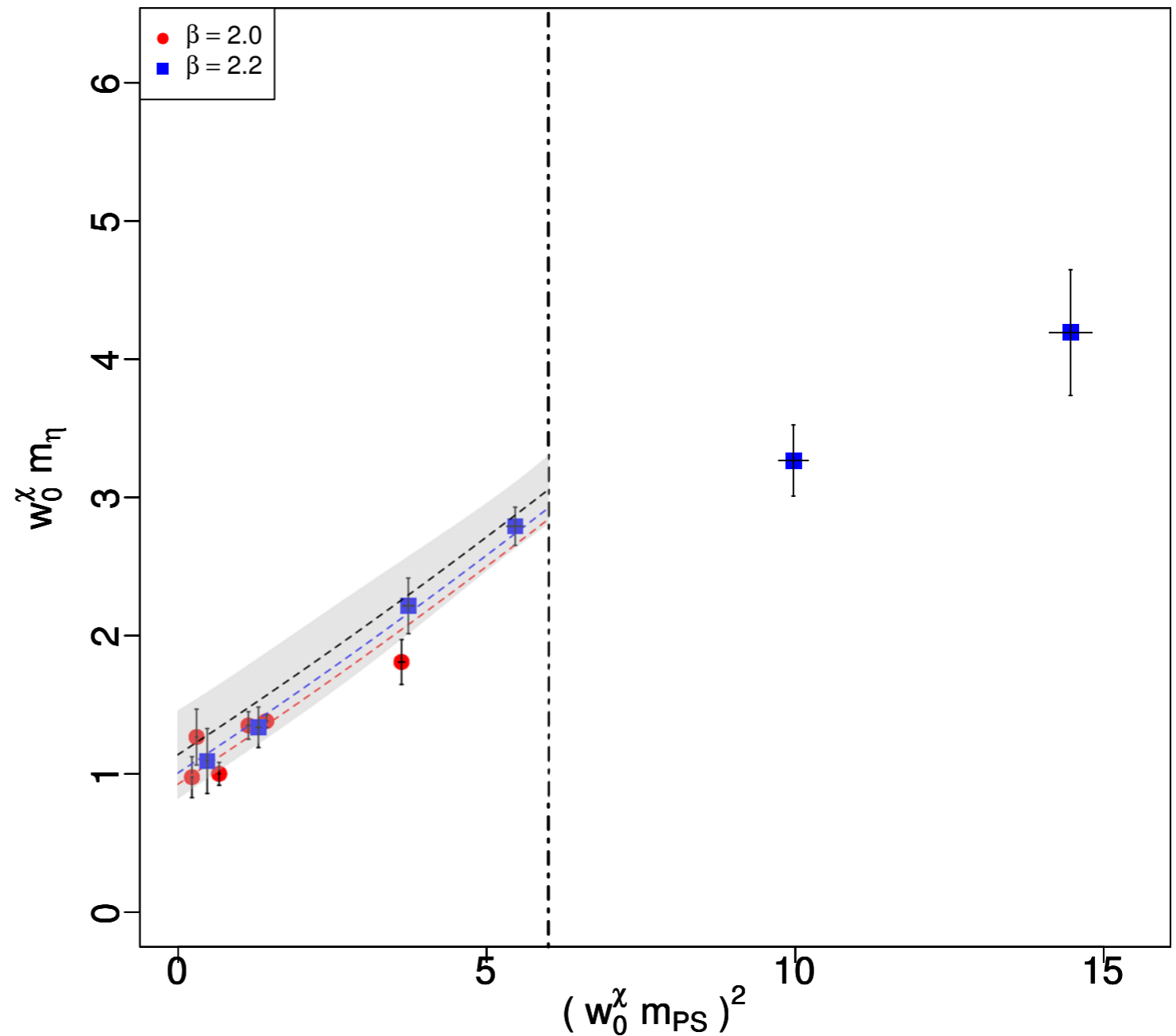
- σ stable in our simulation
- Polynomial global fit including 2 lattice spacing
- $m_\sigma = 21.7(10.8) F_{PS}$



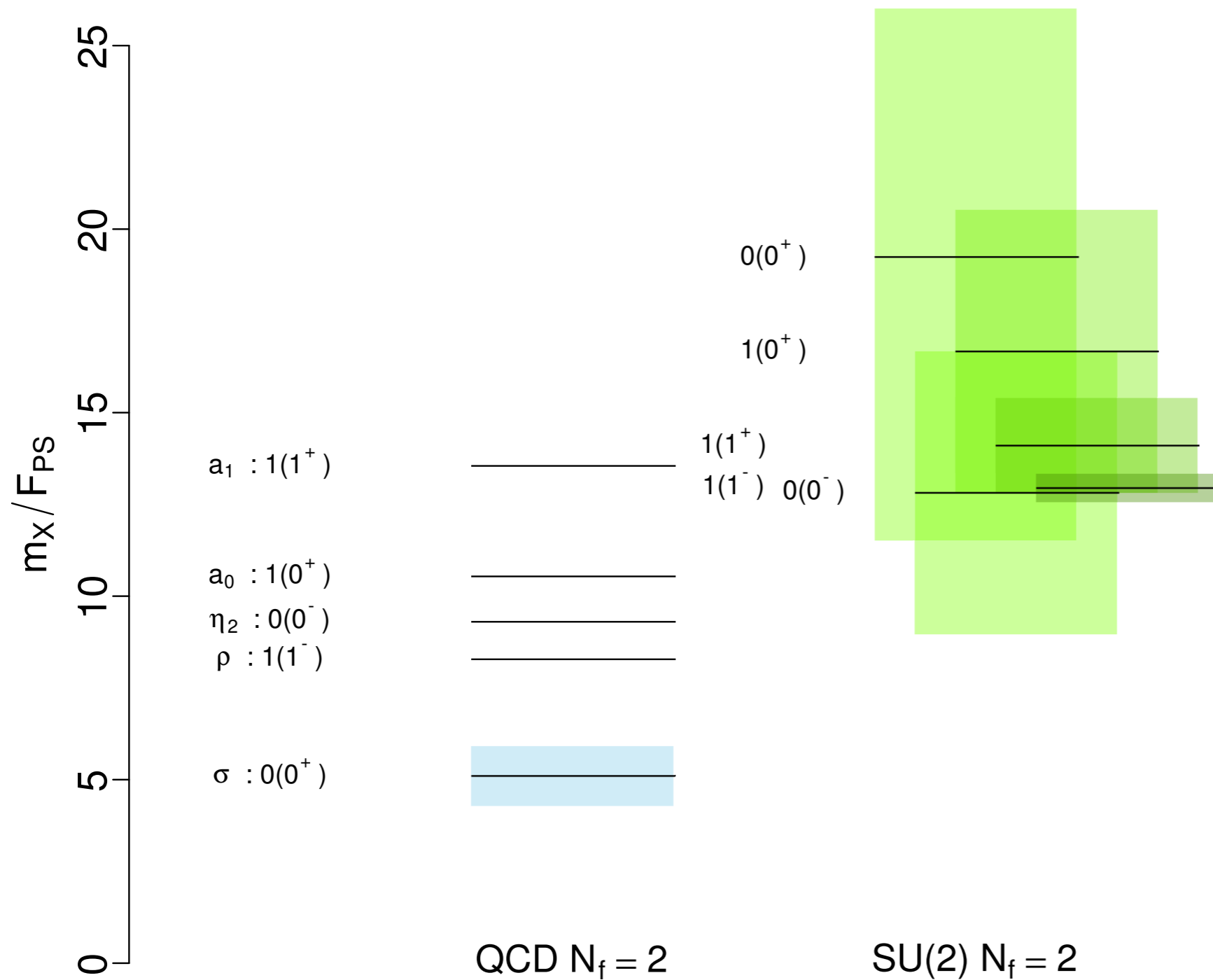
$$w_0^\chi m_X = w_0^\chi m_X^\chi + A(w_0^\chi m_{ps})^2 + B(w_0^\chi m_{ps})^4 + C \frac{a}{w_0}$$

$$\eta' : 0(0^-)$$

- 2 lattice spacings
- $m_\eta > m_{ps}$ only at our lightest fermion masses.
- no cutoff effects are seen.
- $m_\eta = 14.6(4.7) F_{PS}$



Summary and comparison with $N_c=3$



- Shifted upward compared to QCD ?
- Neglects the decay of resonances

New Results: improved setup

Scattering

- ♦ In the TC framework the Goldstone boson equivalence theorem : at large energies, external vector bosons states equivalent to Goldstone Boson states
- ♦ e.g : $\rho \rightarrow \pi^+ \pi^-$ corresponds to vector resonance in W^+W^- scattering

♦ Effective Lagrangian

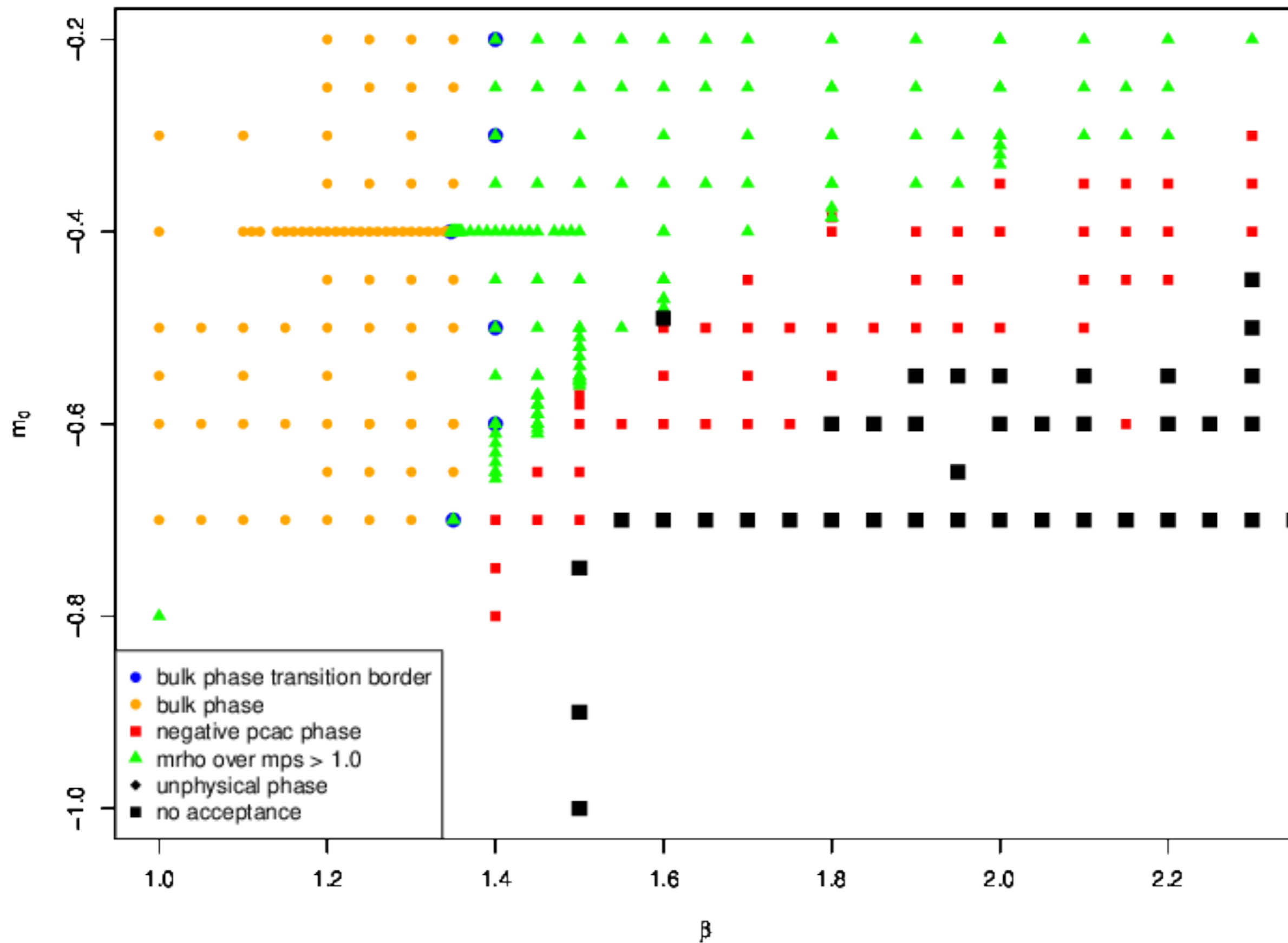
$$\mathcal{L}_{eff} = g_{\rho\pi\pi} \rho_{[ij]}^\mu \partial_\mu \pi_i \pi_j$$

Wilson Clover fermions + Symanzik improved gauge

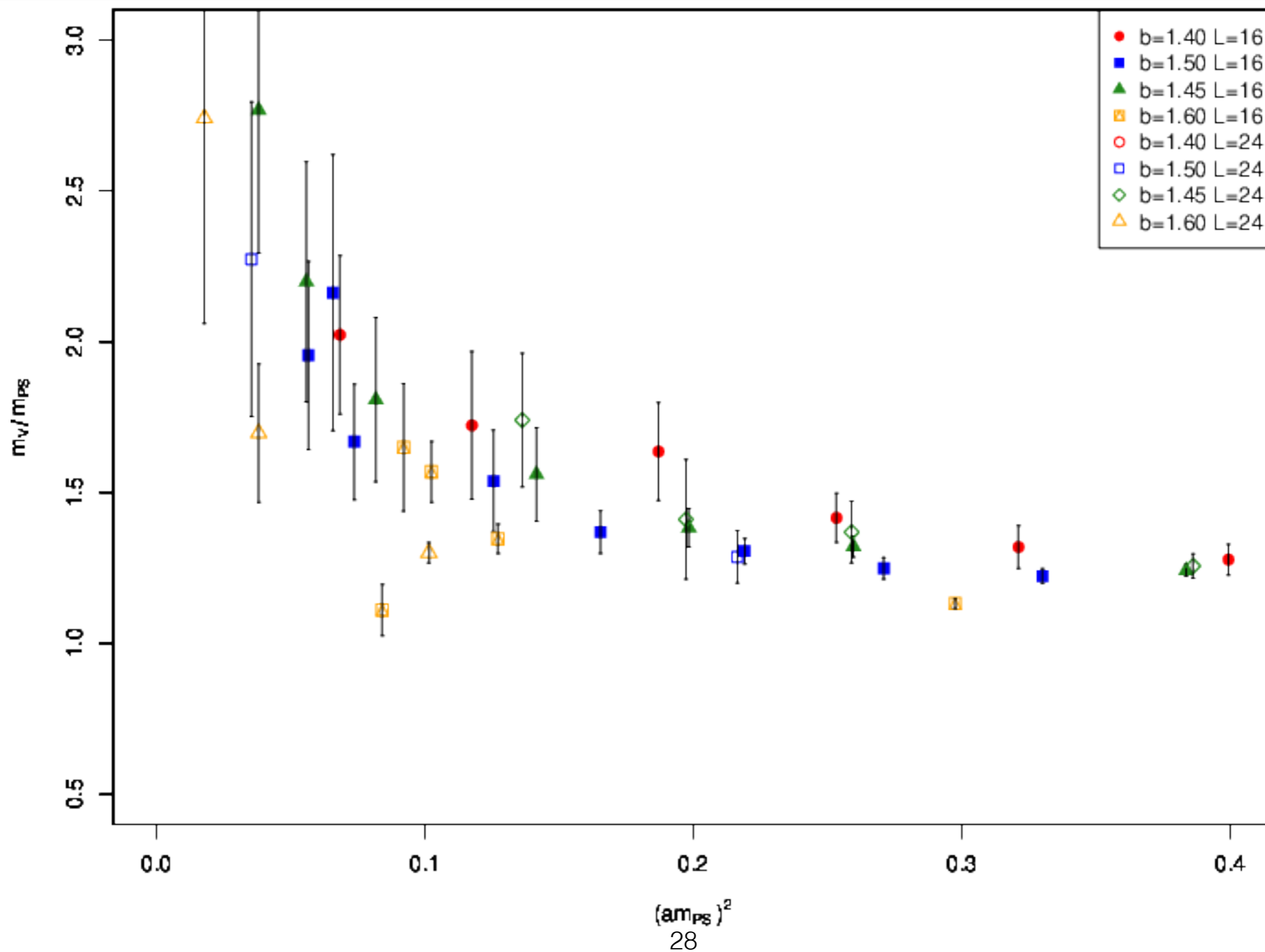
b	m0	L	T	Blk_size	stat	ampac	dampac	amps	damps	afps	dafps	amV	damV
1.25	-0.6000	16	32	100	6500	0.2851	0.0004	1.303	0.002	0.233	0.0024	1.468	0.007
1.30	-0.6000	16	32	100	6500	0.2321	0.0005	1.154	0.002	0.203	0.0025	1.324	0.009
1.35	-0.6000	16	32	100	6500	0.1701	0.0006	0.956	0.003	0.168	0.0025	1.129	0.013
1.35	-0.7000	16	32	10	1666	0.0357	0.0009	0.423	0.006	0.105	0.0035	0.655	0.081
1.40	-0.6000	16	32	10	1500	0.0984	0.0006	0.686	0.003	0.124	0.0022	0.862	0.020
1.40	-0.6100	16	32	10	1392	0.0846	0.0008	0.632	0.005	0.118	0.0032	0.808	0.031
1.40	-0.6200	16	32	10	1218	0.0696	0.0009	0.567	0.006	0.109	0.0034	0.748	0.039
1.40	-0.6300	16	32	10	1621	0.0552	0.0007	0.503	0.005	0.102	0.0024	0.713	0.040
1.40	-0.6400	16	32	10	1795	0.0413	0.0007	0.433	0.005	0.093	0.0025	0.708	0.070
1.40	-0.6500	16	32	10	1585	0.0250	0.0010	0.343	0.010	0.081	0.0040	0.591	0.082
1.40	-0.6570	16	32	10	2779	0.0137	0.0008	0.262	0.010	0.067	0.0036	0.529	0.066
1.45	-0.5500	16	32	10	2000	0.0898	0.0005	0.619	0.003	0.105	0.0016	0.769	0.012
1.45	-0.5700	16	32	20	4156	0.0631	0.0005	0.510	0.003	0.093	0.0013	0.673	0.017
1.45	-0.5800	16	32	10	1791	0.0493	0.0007	0.445	0.005	0.085	0.0019	0.616	0.028
1.45	-0.5900	16	32	10	1602	0.0346	0.0009	0.376	0.007	0.077	0.0027	0.587	0.058
1.45	-0.6000	16	32	10	1766	0.0197	0.0010	0.286	0.010	0.064	0.0033	0.517	0.076
1.45	-0.6050	16	32	10	2371	0.0110	0.0011	0.236	0.013	0.051	0.0049	0.520	0.090
1.45	-0.6100	16	32	20	3424			0.195	0.015			0.541	0.083
1.50	-0.5000	16	32	10	1500	0.0869	0.0006	0.575	0.004	0.091	0.0019	0.703	0.013
1.50	-0.5100	16	32	10	1500	0.0731	0.0007	0.521	0.005	0.084	0.0019	0.650	0.017
1.50	-0.5200	16	32	10	1513	0.0600	0.0007	0.468	0.005	0.078	0.0018	0.611	0.019
1.50	-0.5300	16	32	10	1500	0.0459	0.0008	0.407	0.006	0.072	0.0019	0.557	0.027
1.50	-0.5400	16	32	10	1325	0.0330	0.0010	0.354	0.009	0.062	0.0026	0.545	0.058
1.50	-0.5500	16	32	10	1750	0.0190	0.0010	0.272	0.013	0.050	0.0028	0.453	0.048
1.50	-0.5525	16	32	10	2214	0.0141	0.0010	0.256	0.012	0.041	0.0028	0.555	0.115
1.50	-0.5550	16	32	20	4117	0.0120	0.0011	0.238	0.017	0.038	0.0035	0.465	0.066
1.60	-0.4000	16	32	10	1500	0.0993	0.0005	0.546	0.004	0.071	0.0012	0.617	0.008
1.60	-0.4500	16	32	10	1810	0.0407	0.0007	0.357	0.007	0.048	0.0013	0.480	0.014
1.60	-0.4700	16	32	10	2432	0.0162	0.0011	0.320	0.015	0.022	0.0015	0.502	0.023
1.60	-0.4800	16	32	10	2106	0.0018	0.0016	0.304	0.027	0.002	0.0020	0.501	0.046
1.60	-0.4900	16	32	10	2866	-0.0163	0.0006	0.290	0.012	-0.033	0.0020	0.322	0.021
1.80	-0.2500	16	32	5	500	0.1393	0.0004	0.566	0.004	0.053	0.0009	0.602	0.005
1.80	-0.3000	16	32	5	500	0.0893	0.0004	0.501	0.006	0.045	0.0010	0.548	0.007
1.80	-0.3500	16	32	10	2500	0.0333	0.0003	0.451	0.006	0.019	0.0003	0.464	0.008
1.80	-0.3850	16	32	10	1917	-0.0062	0.0005	0.429	0.007	-0.004	0.0003	0.461	0.007
2.00	-0.2000	16	32	5	500	0.1266	0.0002	0.533	0.004	0.048	0.0009	0.578	0.004
2.00	-0.2500	16	32	5	500	0.0757	0.0003	0.331	0.008	0.038	0.0011	0.365	0.008
2.00	-0.3000	16	32	5	500	0.0250	0.0006	0.234	0.013	0.034	0.0014	0.314	0.020
1.45	-0.5500	24	48	10	1828	0.0899	0.0003	0.621	0.002	0.106	0.0016	0.781	0.025
1.45	-0.5700	24	48	10	1684	0.0628	0.0004	0.509	0.002	0.093	0.0018	0.697	0.052
1.45	-0.5800	24	48	10	1485	0.0487	0.0004	0.444	0.003	0.085	0.0023	0.627	0.088
1.45	-0.5900	24	48	10	1020	0.0340	0.0009	0.369	0.007	0.078	0.0044	0.643	0.081
1.50	-0.5200	24	48	10	1225	0.0603	0.0005	0.465	0.004	0.079	0.0023	0.599	0.040
1.50	-0.5560	24	48	10	1673	0.0100	0.0007	0.188	0.009	0.048	0.0029	0.428	0.096
1.60	-0.4500	24	48	10	2693	0.0404	0.0002	0.319	0.002	0.051	0.0008	0.414	0.010
1.60	-0.4700	24	48	10	1991	0.0151	0.0005	0.195	0.007	0.036	0.0012	0.331	0.043
1.60	-0.4800	24	48	10	1758			0.134	0.012			0.367	0.085

- Reduce cutoff effects
- Or simulate at coarser lattice spacing with larger physical volumes...

Phase diagram: bare parameters



Decaying vector meson !



$\pi\pi$ scattering

- ◆ Consider the process :

$$\pi^a \pi^b \longrightarrow \pi^c \pi^d$$

- ◆ In (2 flavour) QCD pion's belong to "3", and the two pion operators can be classified according to :

$$3 \times 3 = 1 + 3 + 5$$

- ◆ In our case GB belong to "5" dimensional irrep of $SP(4)$:

$$5 \times 5 = 1 + 10 + 14$$

- ◆ There are still 3 channels

Phase shift & cross section

- ◆ Differential cross-section in terms of scattering amplitude:

$$\frac{d\sigma}{d\Omega} = |A(p, \Omega)|^2$$

- ◆ Partial wave decomposition:

$$A(p, \Omega) = \sum_{l,m} 4\pi A_l(p) Y_{lm}^*(\theta_p, \phi_p) Y_{lm}(\theta, \phi)$$

- ◆ Phase shift definition: $A_l(p) = \frac{e^{2i\delta_l(p)} - 1}{2ip} = \frac{e^{i\delta_l(p)} \sin \delta_l(p)}{p}$

- ◆ Relation with effective coupling and width

$$\tan \delta_1 = \frac{g_{\rho\pi\pi}^2}{6\pi} \frac{p^3}{E_{CM}(m_\rho^2 - E_{CM}^2)}, \quad p = \sqrt{E_{CM}^2/4 - m_\pi^2}$$

$$\Gamma_\rho = \Gamma_R(s) \Big|_{s=m_\rho^2} = \frac{g_{\rho\pi\pi}^2}{6\pi} \frac{p_\rho^3}{m_\rho^2}, \quad p_\rho = \sqrt{m_\rho^2/4 - m_\pi^2}.$$

Lattice calculation in a nutshell

[Luscher, Rummukainen-Gottlieb,...]

- ◆ Non interacting pions with momenta \mathbf{p} and $\mathbf{0}$:

$$E_1^0 = \sqrt{m_{\text{PS}}^2 + \mathbf{p}^2} + m_{\text{PS}}$$

$$E_2^0 = \sqrt{m_{\text{V}}^2 + \mathbf{p}^2}$$

- ◆ In a finite box the energy levels are shifted by the interaction. The energy shift is related to the two-pion scattering phase shift δ in infinite volume.

- ◆ Total momentum: $P = \frac{2\pi}{L} \mathbf{e}_z$

$$\tan \delta_1(E_{CM}) = \frac{\gamma \pi^{3/2} q}{\mathcal{Z}_{00}^{\mathbf{d}}(1; q^2) + (2q^{-2}/\sqrt{5}) \mathcal{Z}_{20}^{\mathbf{d}}(1; q^2)} \quad E_{CM} = 2\sqrt{m^2 + \mathbf{p}^{*2}} \quad \mathbf{p}^{*2} \equiv \left(q \frac{2\pi}{L} \right)^2$$

Lattice calculation in a nutshell

- ◆ Lattice approach, compute:

$$C(t) = \begin{pmatrix} \langle 0 | (\pi\pi)^\dagger(t) (\pi\pi)(t_S) | 0 \rangle & \langle 0 | (\pi\pi)^\dagger(t) \rho_3(t_S) | 0 \rangle \\ \langle 0 | \rho_3^\dagger(t) (\pi\pi)(t_S) | 0 \rangle & \langle 0 | \rho_3^\dagger(t) \rho_3(t_S) | 0 \rangle \end{pmatrix}$$

with $(\pi\pi)(t) = \frac{1}{\sqrt{2}} \left(\pi^-(\mathbf{p}, t) \pi^+(\mathbf{0}, t) - \pi^+(\mathbf{p}, t) \pi^-(\mathbf{0}, t) \right)$

- ◆ Large time behaviour:

$$C_{ij}(t) \equiv \langle 0 | O_i^\dagger(t) O_j(0) | 0 \rangle = \sum_{n,m} \langle 0 | O_i^\dagger | n \rangle (e^{-E_n t} \delta_{mn}) \langle m | O_j | 0 \rangle$$

- ◆ Solve the generalised eigenvalue problem:

$$C_{ij}^{-1}(t_0) C_{jk}(t) = V_{in}^{-1} \text{diag} \left(e^{-E_n(t-t_0)} \right)_{nm} V_{mj}$$

Wick contractions

$$C_{11}(t) =$$

The expression for $C_{11}(t)$ is represented by a sum of nine Feynman diagrams. The first row contains three diagrams: a bubble with two external legs labeled 'p', a bubble with one external leg labeled 'p' and one labeled '0', and a square with four external legs labeled 'p', '0', 'p', '0'. The second row contains three diagrams: a bubble with two external legs labeled '0', a bubble with one external leg labeled '0' and one labeled 'p', and a square with four external legs labeled '0', 'p', '0', 'p'. The third row contains three diagrams: a square with four external legs labeled 'p', '0', 'p', '0', a square with four external legs labeled 'p', 'p', '0', '0', and a square with four external legs labeled 'p', 'p', '0', '0'. The diagrams are separated by plus and minus signs as indicated in the equation.

$$C_{12}(t) = -C_{21}^*(t) =$$

The expression for $C_{12}(t) = -C_{21}^*(t)$ is represented by two Feynman diagrams. Both are triangles with three external legs labeled 'p', 'p', and '0'. The first diagram has the '0' leg at the bottom-left, and the 'p' legs at the top-left and right. The second diagram has the '0' leg at the bottom-left, and the 'p' legs at the top-left and right. The two diagrams are separated by a minus sign.

$$C_{22}(t) =$$

The expression for $C_{22}(t)$ is represented by a single Feynman diagram: a bubble with two external legs labeled 'p'.

Benchmark ensemble

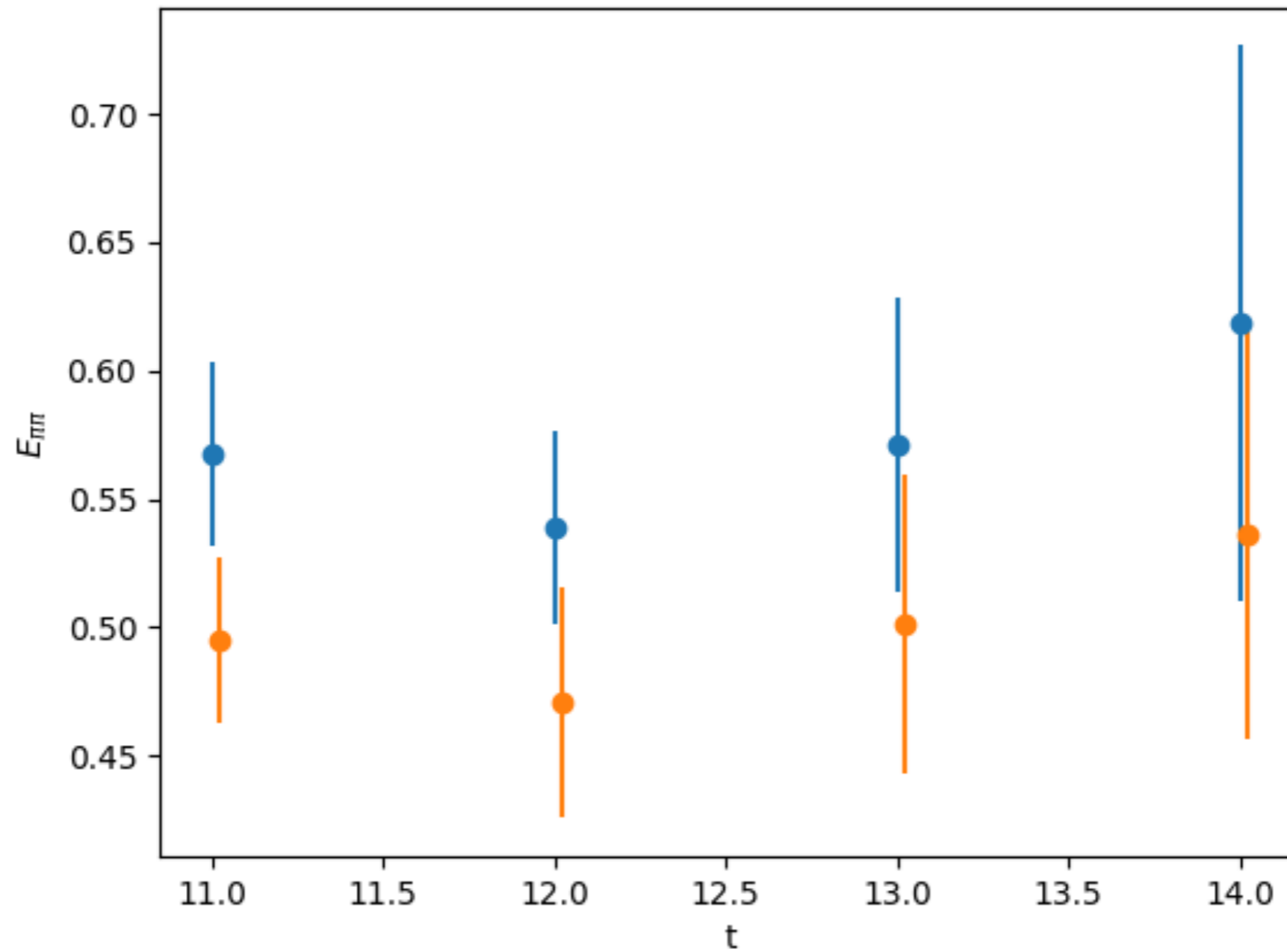
β	1.45
m_0	-0.6050
c_{sw}	1.0
<hr/>	
am_π	0.2114(8)
am_ρ^{naive}	0.444(9)
am_{pcac}^0	0.01110(7)
af_π	0.0564(3)
<hr/>	
# trajectories	1600
# analysed	140

- ♦ For the time being, we use $p = \frac{2\pi}{L} \mathbf{e}_z$
- ♦ We plan to use one or two more moving frames

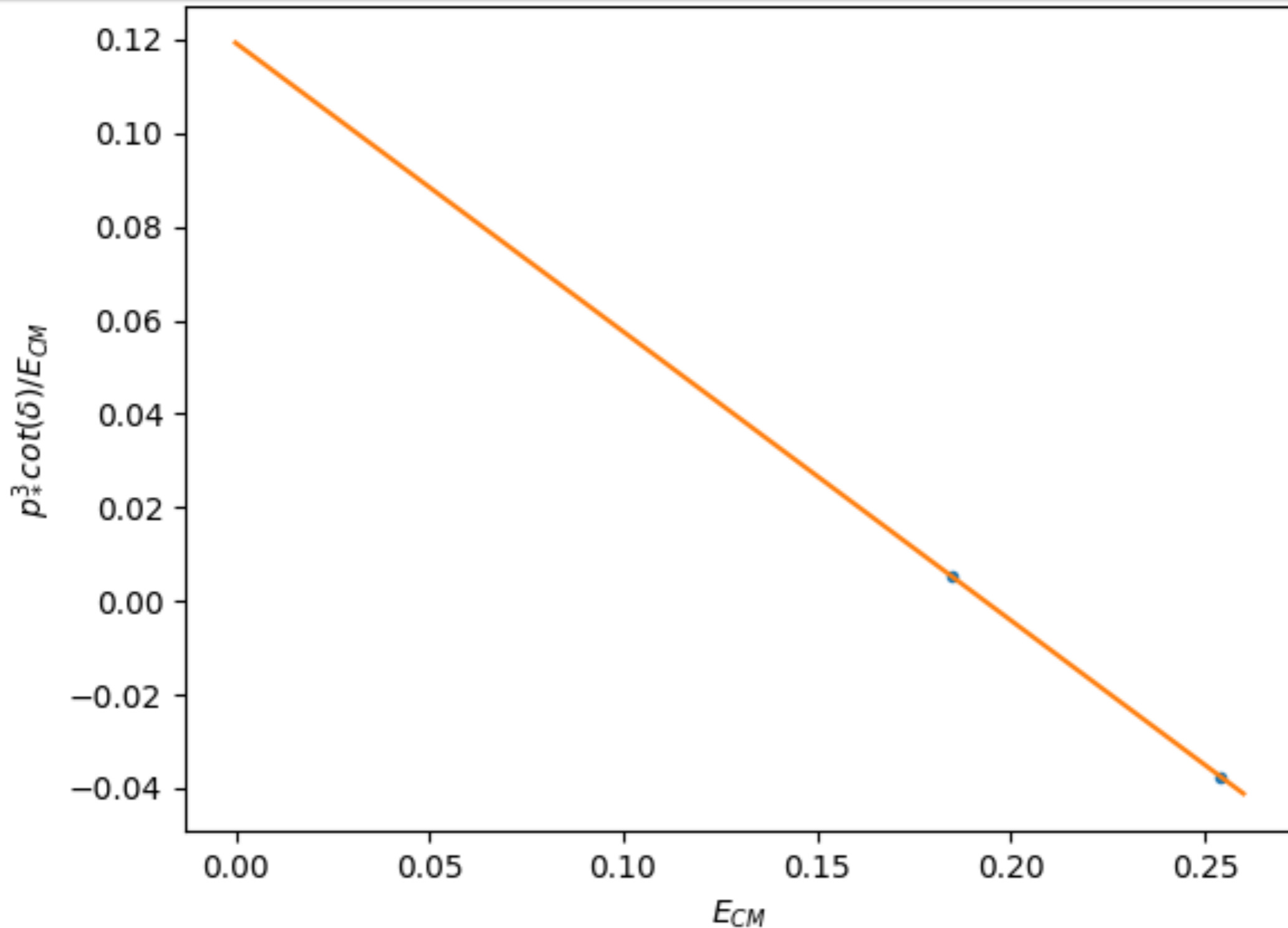
Energy levels

Eigenvalue $\lambda_i(t) = \exp(-E_i(t - t_0))$, $t_0 = 10$.

Effective mass $E_i(t) = \ln \lambda_i(t) / \lambda_i(t + 1)$



Preliminary: phase shift



Fitting central values only (for now) **PRELIMINARY**:

$$g_{\rho\pi\pi} = 5.527$$

$$M_\rho = 0.440$$

Summary & Conclusions

- Lattice as a laboratory to explore non-perturbative dynamics of gauge theories
- Prediction of the spectrum for a range of underlying fermion mass
- Decay channels are often closed for kinematic reasons (scalar, vector, axial are stable in most of our simulations)
- Benchmark predictions in the chiral limit for spin-1 resonances:
 - ★ $m_V/F_{PS} = 13.1(2.2)$; $m_A/F_{PS} = 14.5(3.6)$
- Preliminary results in the spin-0 sector :
 - ★ $m_\sigma/F_{PS} = 21.7(10.8)$
 - ★ $m_\eta/F_{PS} = 14.6(4.7)$
- Ongoing/Outlook :
 - ★ Update results on the spectrum
 - ★ Preliminary results on the vector meson resonance
 - ★ Generalisation to other resonances?