Impact of ground-state correlations on the nuclear response

ECT* workshop Next generation *ab initio* nuclear theory

Trento, July 16th, 2025

Andrea Porro

Technische Universität Darmstadt



Introduction

- Physics case
- Existing ab initio methods

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IMSRG multipole moments

- Moments of the strength
- IMSRG implementation
- Model-space convergence

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- Interaction sensitivity
- Comparison to experiment
- Comparison to sum rules

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Challenges and opportunities

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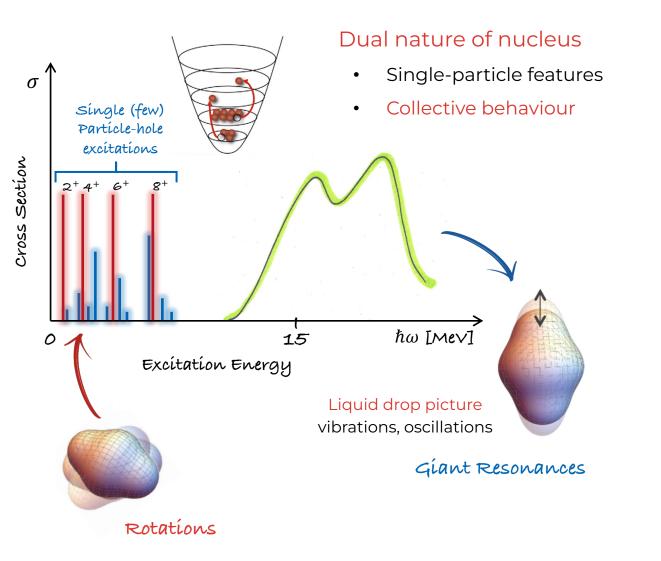
IMSRG multipole moments

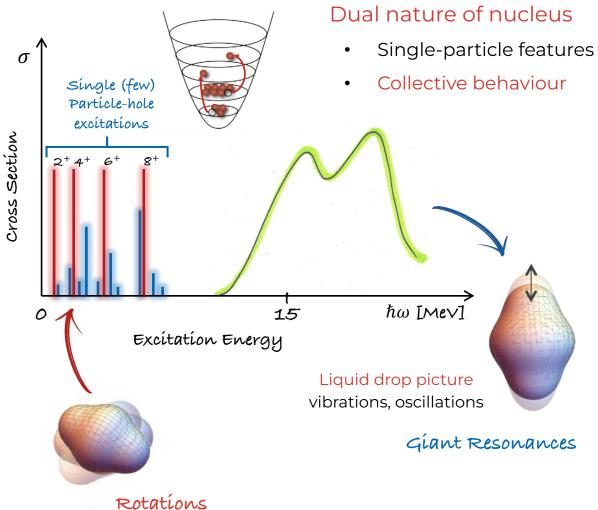
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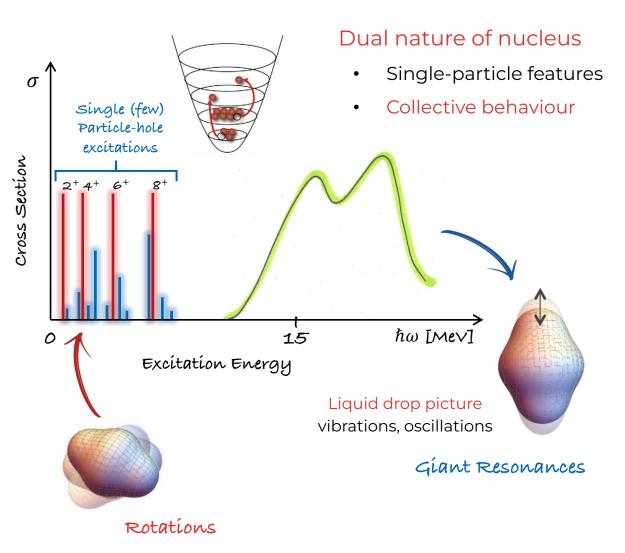
Challenges and opportunities





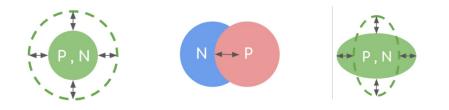
Response function

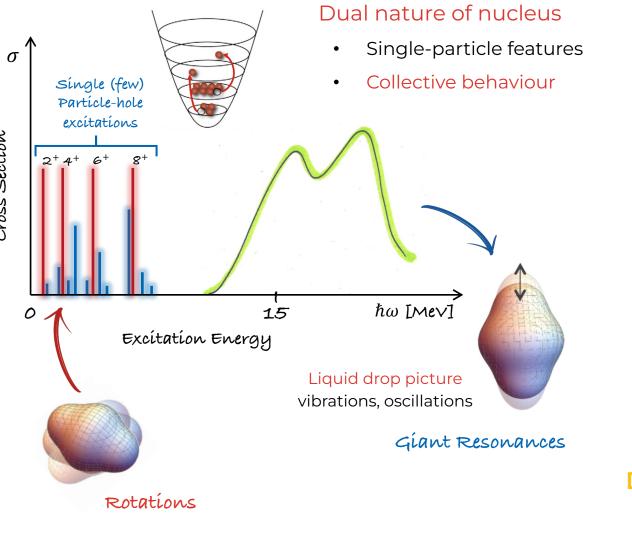
Fully characterise linear response



Response function

Fully characterise linear response

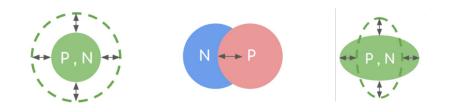




Response function

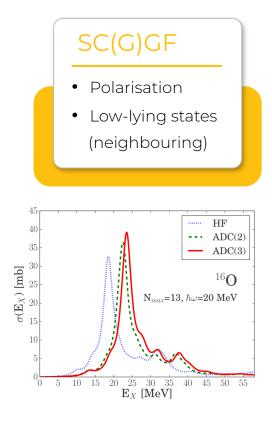
Fully characterise linear response

Studied quantity: multipole response



Different excitations





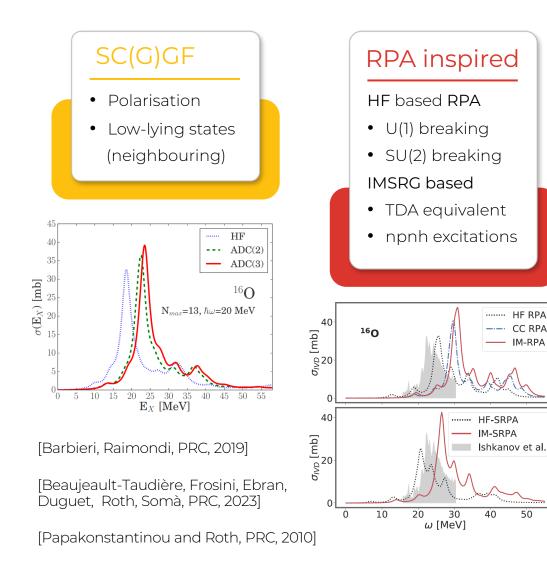
[Barbieri, Raimondi, PRC, 2019]

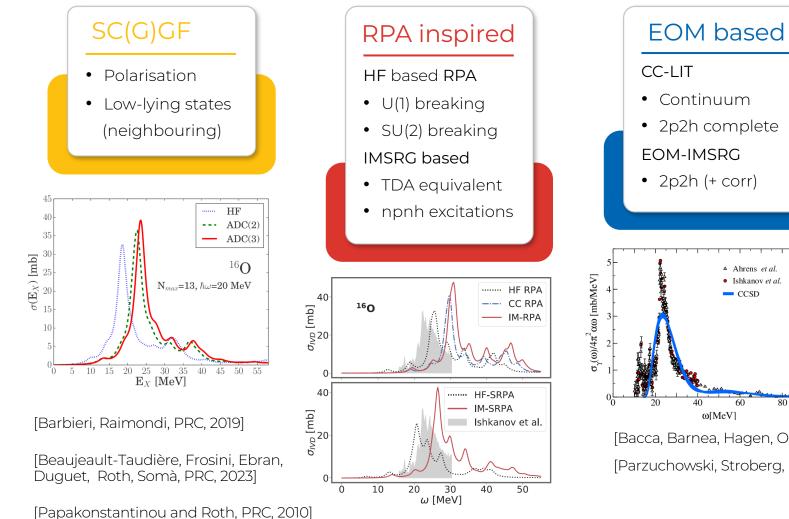
HF RPA

CC RPA

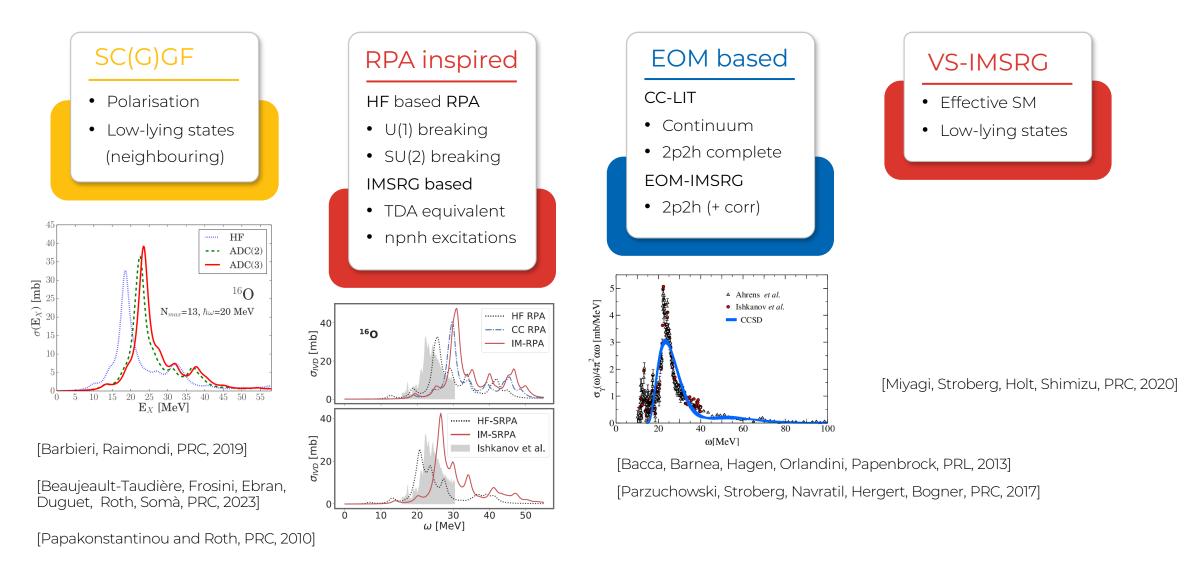
IM-RPA

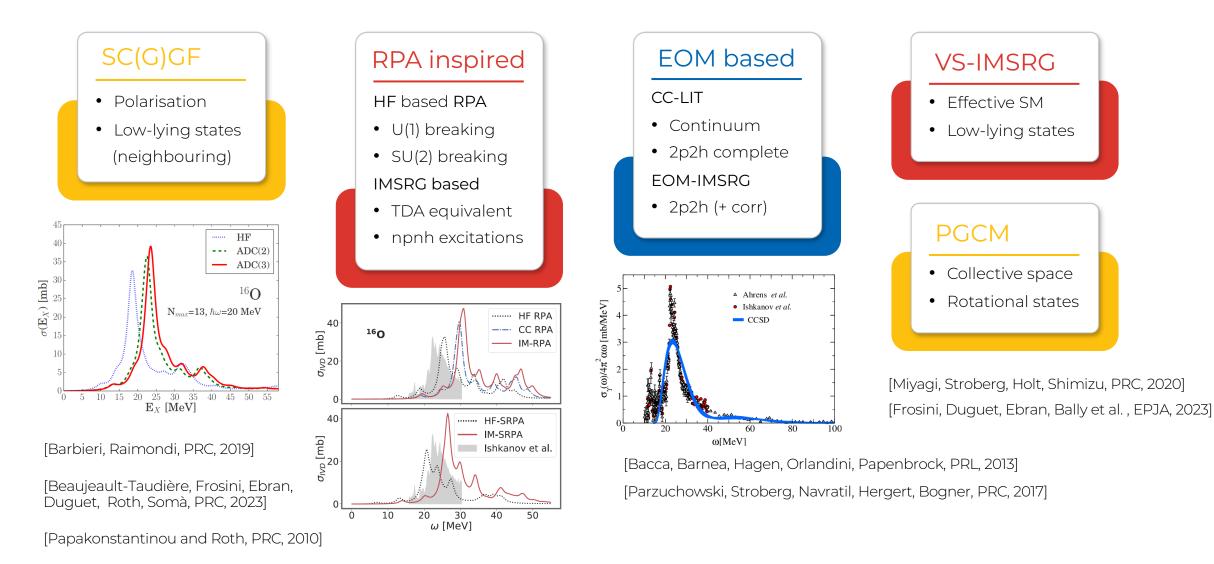
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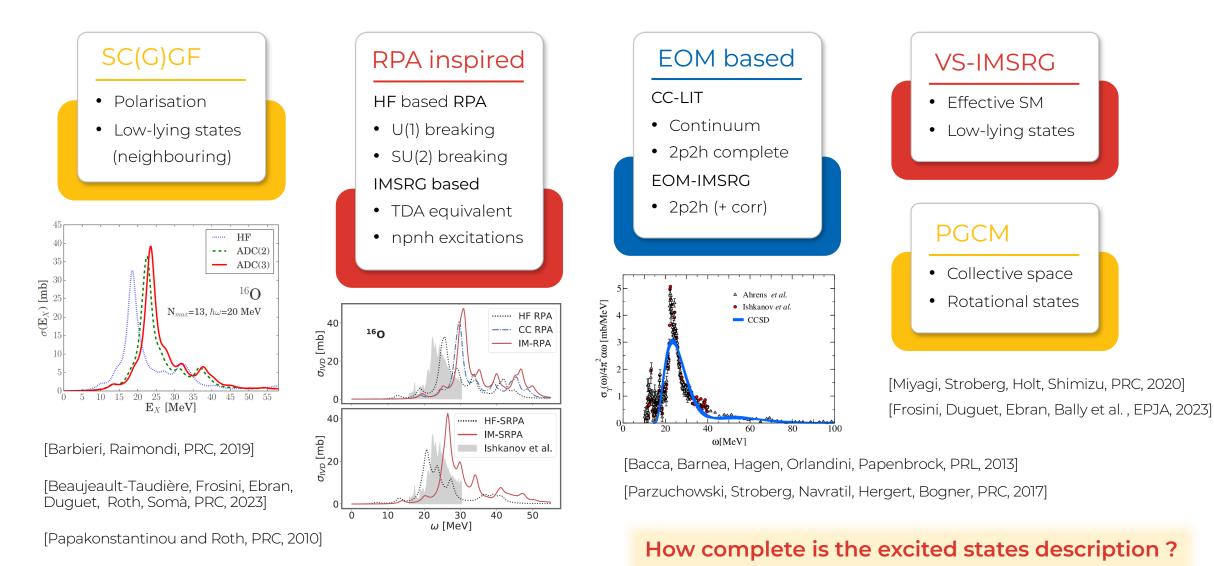




(a) Ahrens et al. (a) Ahrens et al. (b) Ahrens et al. (c) CCSD (c) CCSD







Uncertainty quantification / benchmark

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$$S(Q_{\lambda}, E) \equiv \sum_{\mu\nu} |\langle \Psi_{\nu} | Q_{\lambda\mu} | \Psi_{0} \rangle|^{2} \,\delta(E_{\nu} - E_{0} - E)$$

Studied quantity: multipole response

$$S(Q_{\lambda}, E) \equiv \sum_{\mu\nu} |\langle \Psi_{\nu} | Q_{\lambda\mu} | \Psi_{0} \rangle|^{2} \,\delta(E_{\nu} - E_{0} - E)$$

Related moments

$$m_k(Q_\lambda) \equiv \int_0^\infty E^k S(Q_\lambda, E) dE$$
$$= \sum_{\mu\nu} (E_\nu - E_0)^k |\langle \Psi_\nu | Q_{\lambda\mu} | \Psi_0 \rangle|^2$$

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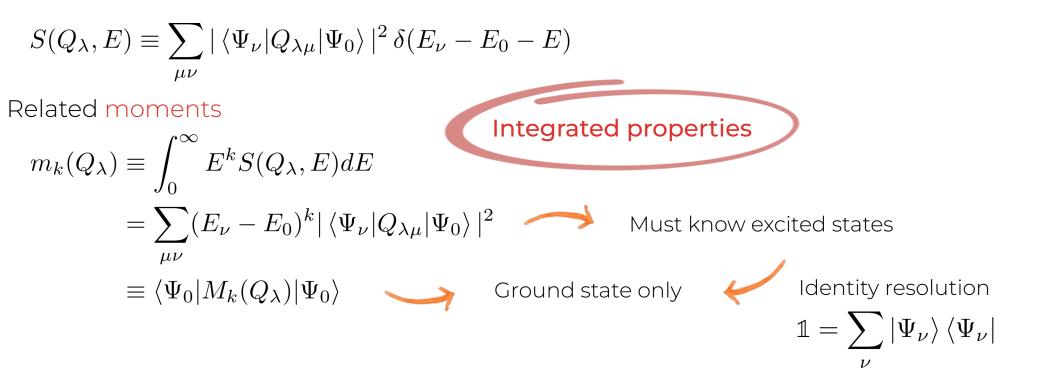
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Studied quantity: multipole response

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 ν

Complexity shifted to operator structure

$$M_0(Q_\lambda) \equiv \sum_{\mu} (-1)^{\mu} Q_{\lambda,-\mu} Q_{\lambda\mu}$$
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• Exact treatment for exc states

 ν

• Many-body truncation only GS

"Exact sum rules with approximate ground states"

[Johnson et al., JPG, 2020]

Studied quantity: multipole response

Exact implementation up to m₁

$$S(Q_{\lambda}, E) \equiv \sum_{\mu\nu} |\langle \Psi_{\nu} | Q_{\lambda\mu} | \Psi_{0} \rangle|^{2} \delta(E_{\nu} - E_{0} - E)$$
Effective two-body Hamiltonian
$$H = H^{[1]} + H^{[2]}$$
Spherical tensor operators
$$m_{k}(Q_{\lambda}) \equiv \int_{0}^{\infty} E^{k} S(Q_{\lambda}, E) dE$$

$$= \sum_{\mu\nu} (E_{\nu} - E_{0})^{k} |\langle \Psi_{\nu} | Q_{\lambda\mu} | \Psi_{0} \rangle|^{2}$$
Must know excited states
$$\equiv \langle \Psi_{0} | M_{k}(Q_{\lambda}) | \Psi_{0} \rangle$$
Ground state only
$$1 = \sum_{\nu} |\Psi_{\nu} \rangle \langle \Psi_{\nu}|$$
Complexity shifted to operator structure
$$Effective two-body Hamiltonian
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Spherical tensor operators
$$Q_{\lambda\mu}^{\dagger} = (-1)^{\mu}Q_{\lambda,-\mu}$$

$$1 = \sum_{\nu} |\Psi_{\nu} \rangle \langle \Psi_{\nu}|$$

$$M_0(Q_\lambda) \equiv \sum_{\mu} (-1)^{\mu} Q_{\lambda,-\mu} Q_{\lambda\mu}$$
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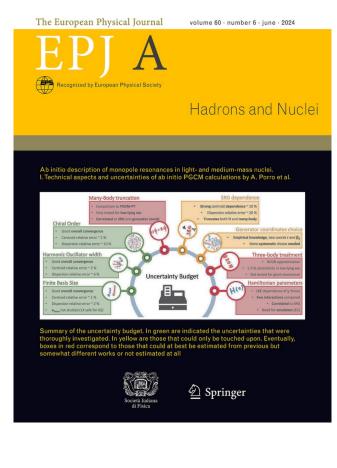
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[Johnson et al., JPG, 2020]

Previous PGCM study



- I. [EPJA (2024) 60, 133]
- II. [EPJA (2024) 60, 134]
- III. [EPJA (2024) 60, 155]
- IV. [EPJA (2024) 60, 233]

| Eur. Phys. J. A (2024) 60:155 | |
|---|--|
| https://doi.org/10.1140/epja/s10050-024-01377-5 | |

Regular Article - Theoretical Physics

Ab initio description of monopole resonances in light- and medium-mass nuclei

III. Moments evaluation in ab initio PGCM calculations

A. Porro^{1,2,3,a}, T. Duguet^{3,4}, J.-P. Ebran^{5,6}, M. Frosini⁷, R. Roth^{1,8}, V. Somà³



Previous PGCM study





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- [EPJA (2024) 60, 133] Ι. 11. [EPJA (2024) 60, 134]
- |||. [EPJA (2024) 60, 155]
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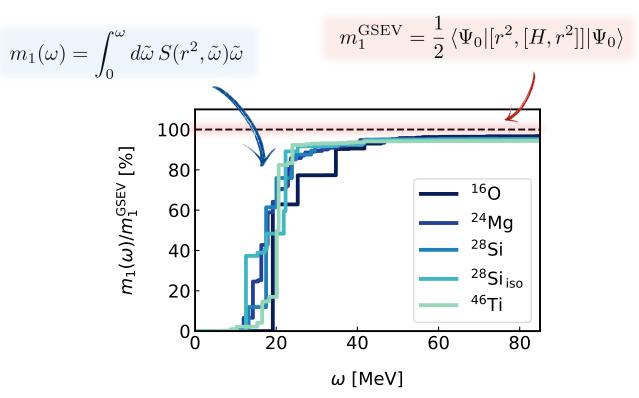
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Previous PGCM study



Ab initio description of monopole resonances in light- and medium-mass nuclei. I. Technical aspects and uncertainties of ab initio PGCM calculations by A. Porro et al



thoroughly investigated. In yellow are those that could only be touched upon. Eventually, boxes in red correspond to those that could at best be estimated from previous but somewhat different works or not estimated at all



- I. [EPJA (2024) 60, 133] II. [EPJA (2024) 60, 134]
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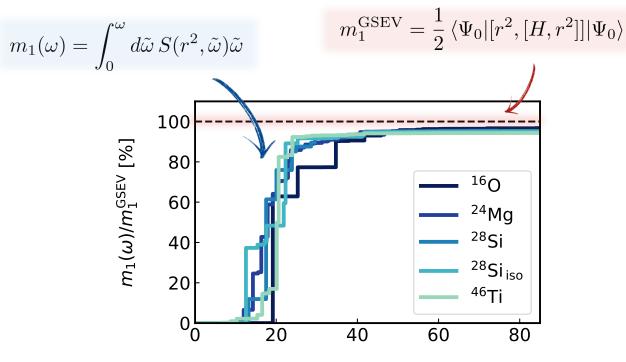
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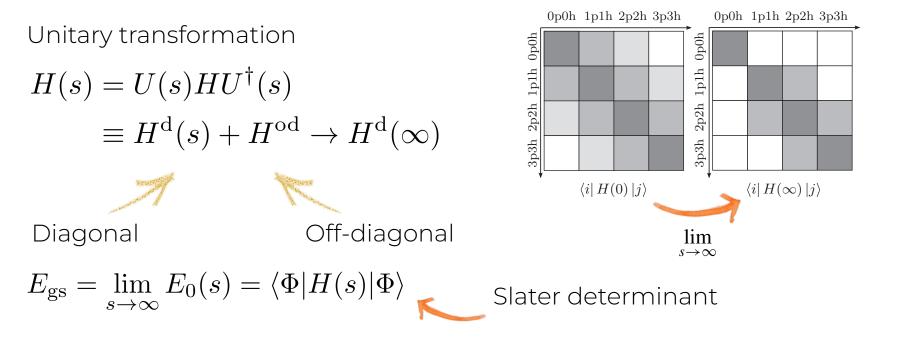
updates



 ω [MeV]

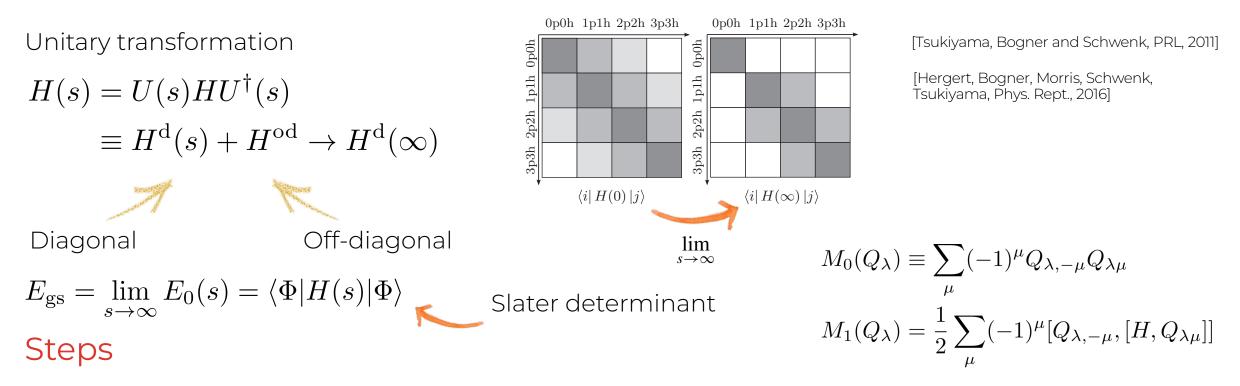
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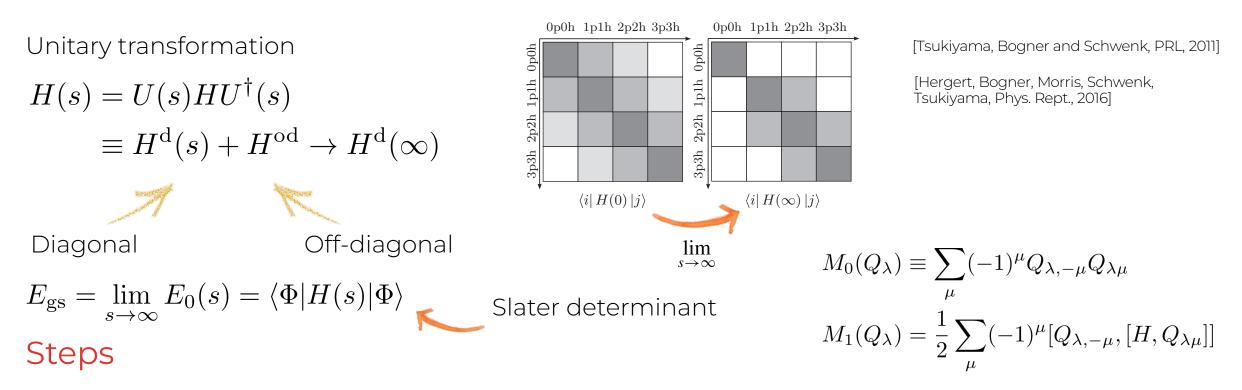


[Tsukiyama, Bogner and Schwenk, PRL, 2011]

[Hergert, Bogner, Morris, Schwenk, Tsukiyama, Phys. Rept., 2016]

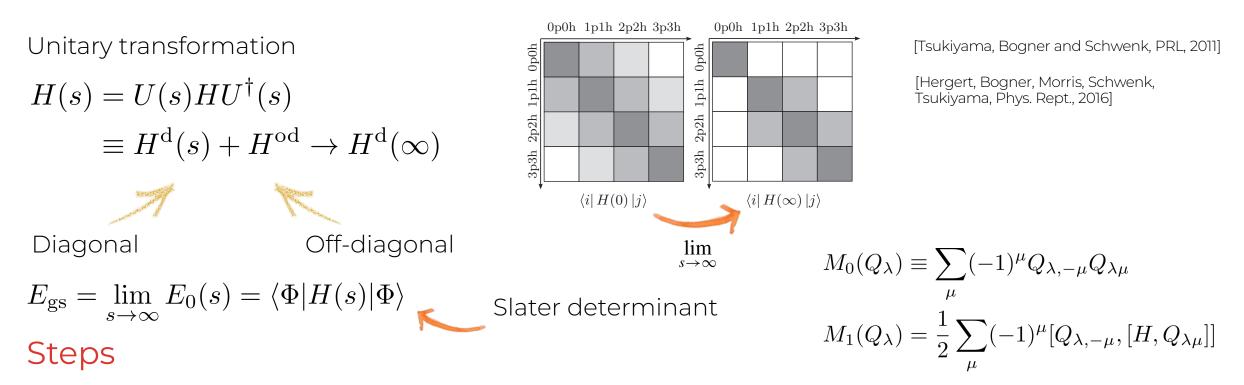


Start from the moment operator in the HO basis



Start from the moment operator in the HO basis

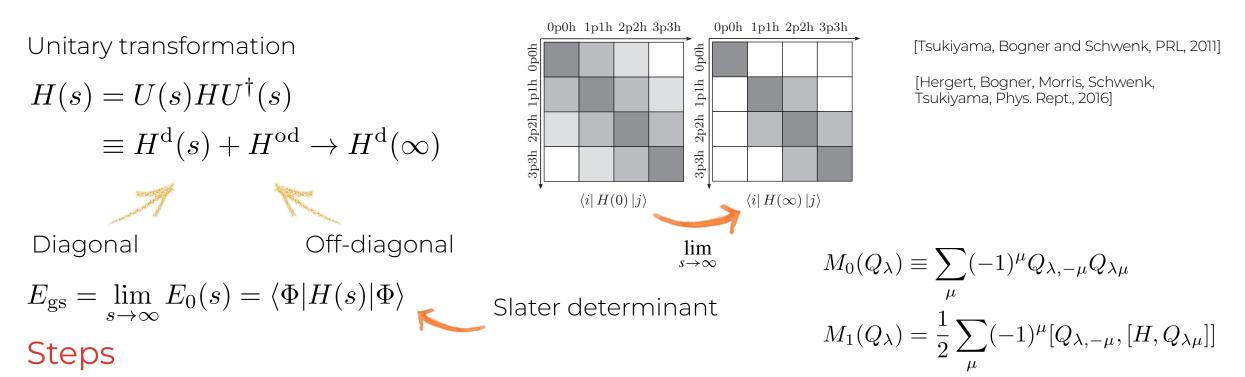
J-scheme expressions of m_0 and m_1 [Lu and Johnson, PRC 97 (2018) 3, 034330]



Start from the moment operator in the HO basis

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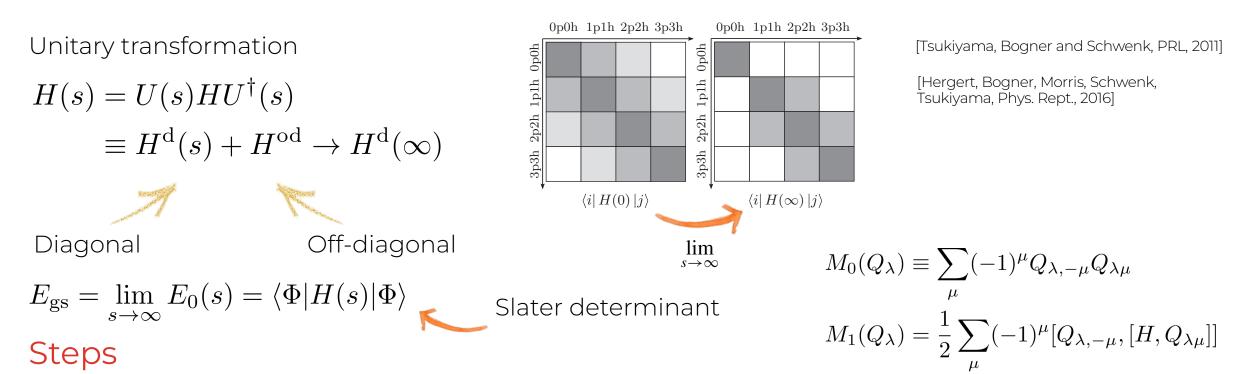
Implemented within imsrg++ code
[github.com/ragnarstroberg/imsrg]



- Start from the moment operator in the HO basis
- Perform an IMSRG(2) calculation

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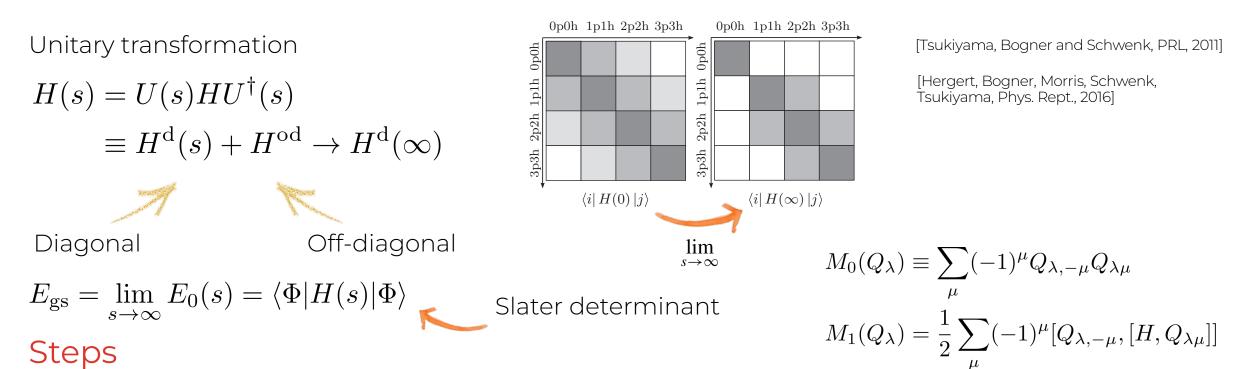


- Start from the moment operator in the HO basis
- Perform an IMSRG(2) calculation
- Evolve moment operators using Magnus $U(s)\equiv e^{\Omega(s)}$

J-scheme expressions of m_0 and m_1 [Lu and Johnson, PRC 97 (2018) 3, 034330]

Implemented within imsrg++ code
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Strategy in the IMSRG framework



- Start from the moment operator in the HO basis
- Perform an IMSRG(2) calculation
- Evolve moment operators using Magnus $U(s)\equiv e^{\Omega(s)}$

Benchmarks

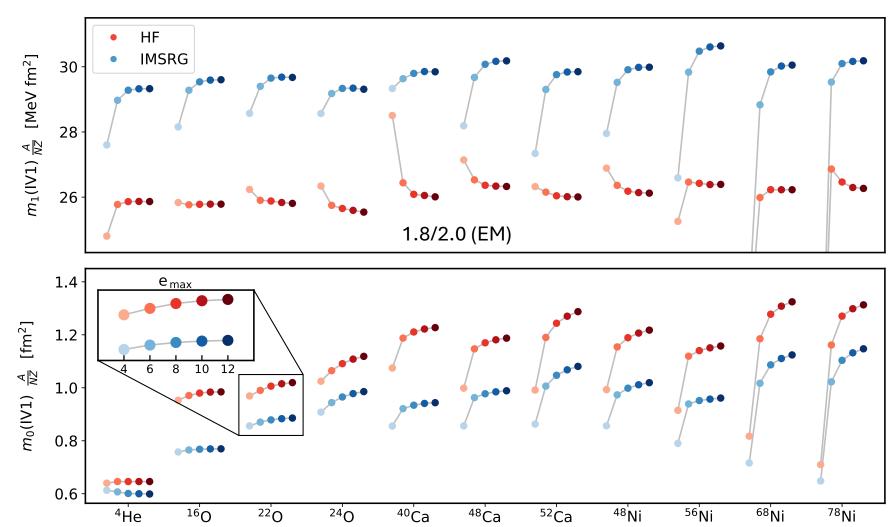
- HF value of m₀ against TDA
- HF value of m₁ against RPA

J-scheme expressions of m_0 and m_1 [Lu and Johnson, PRC 97 (2018) 3, 034330]

Implemented within imsrg++ code
[github.com/ragnarstroberg/imsrg]

$$Q_{1\mu}^{\rm IV} = \frac{N}{A} \sum_{i=1}^{Z} r_i Y_{1\mu}(\hat{r}_i) - \frac{Z}{A} \sum_{i=1}^{N} r_i Y_{1\mu}(\hat{r}_i)$$

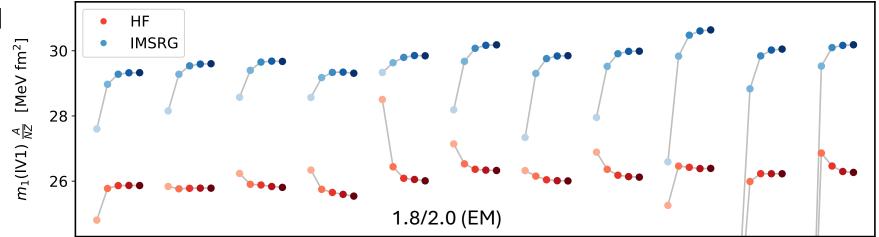
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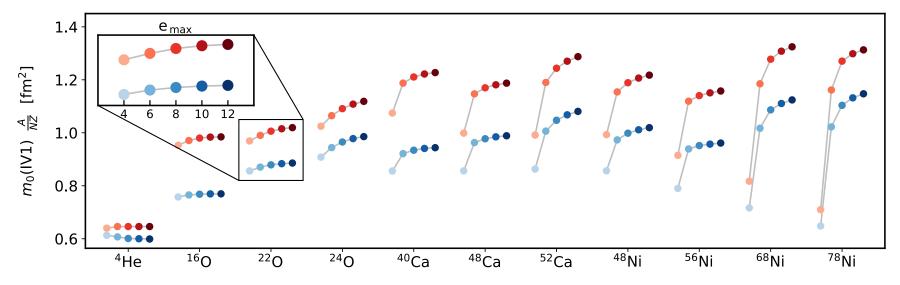


$$M_1(Q_{\lambda}) = \frac{1}{2} \sum_{\mu} (-1)^{\mu} [Q_{\lambda,-\mu}, [H, Q_{\lambda\mu}]]$$

- Large correlation impact
- Relative difference ~0.2%
- Similar error for $\hbar\omega$ variations

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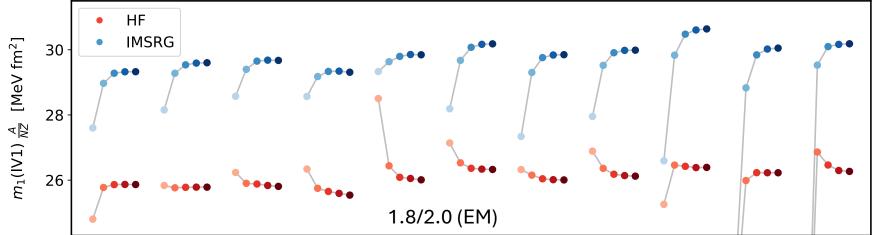
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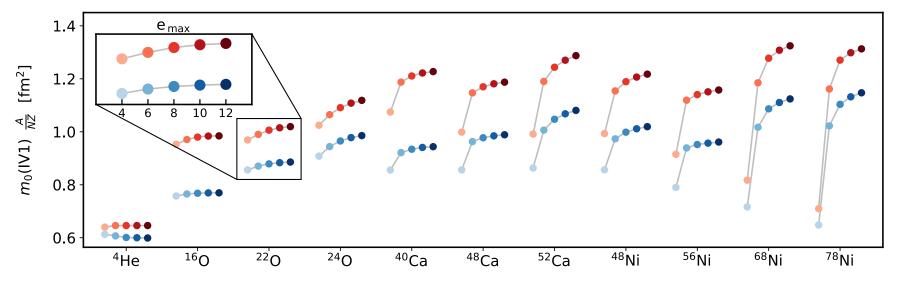
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$$M_0(Q_\lambda) \equiv \sum_{\mu} (-1)^{\mu} Q_{\lambda,-\mu} Q_{\lambda\mu}$$

- Slower convergence
- Relative difference ~1.3%
- 2% error for $\hbar\omega$ variations

$$Q_{1\mu}^{\rm IV} = \frac{N}{A} \sum_{i=1}^{Z} r_i Y_{1\mu}(\hat{r}_i) - \frac{Z}{A} \sum_{i=1}^{N} r_i Y_{1\mu}(\hat{r}_i)$$





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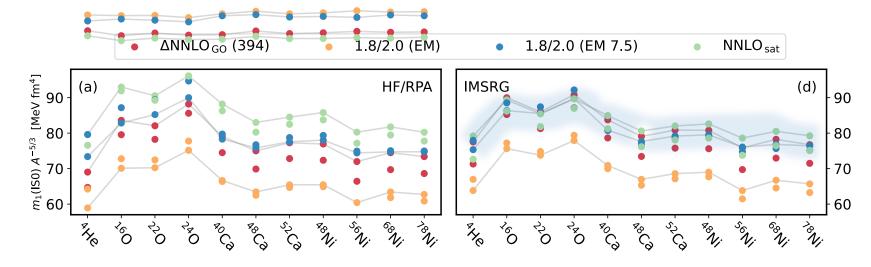
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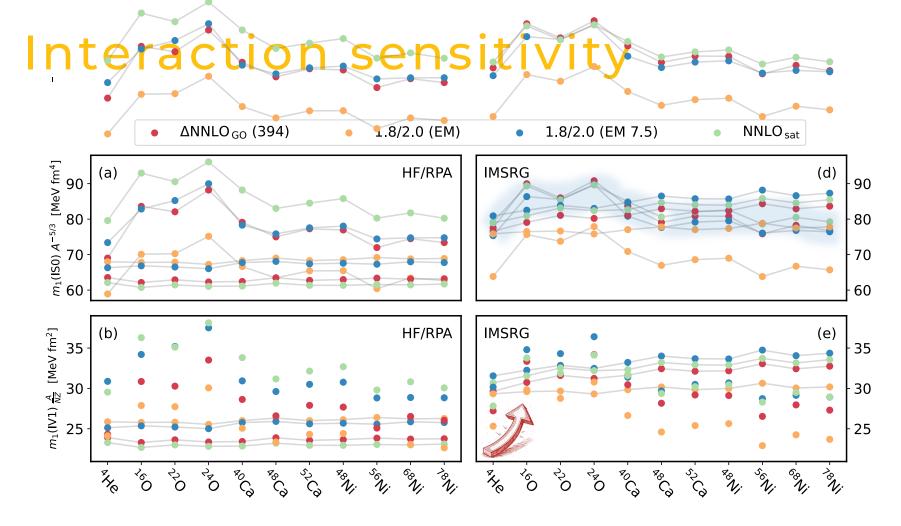
Challenges and opportunities

Interaction sensitivity



Monopole

- Reduced spread
- ~5% correlations effect



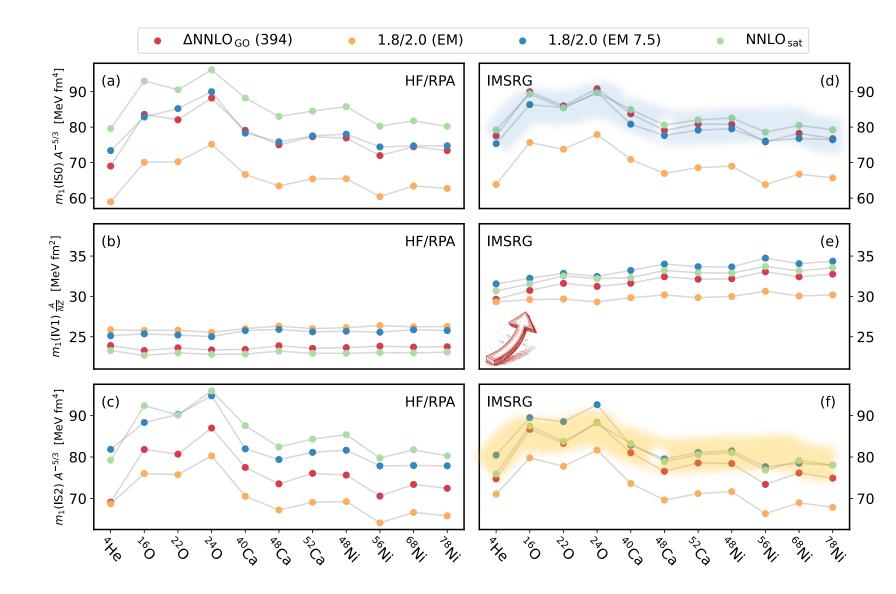
Monopole

- Reduced spread
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Dipole

- Increase up to 40%
- 2% spread (w/o 1.8/2.0(EM))

Interaction sensitivity



Monopole

- Reduced spread
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Dipole

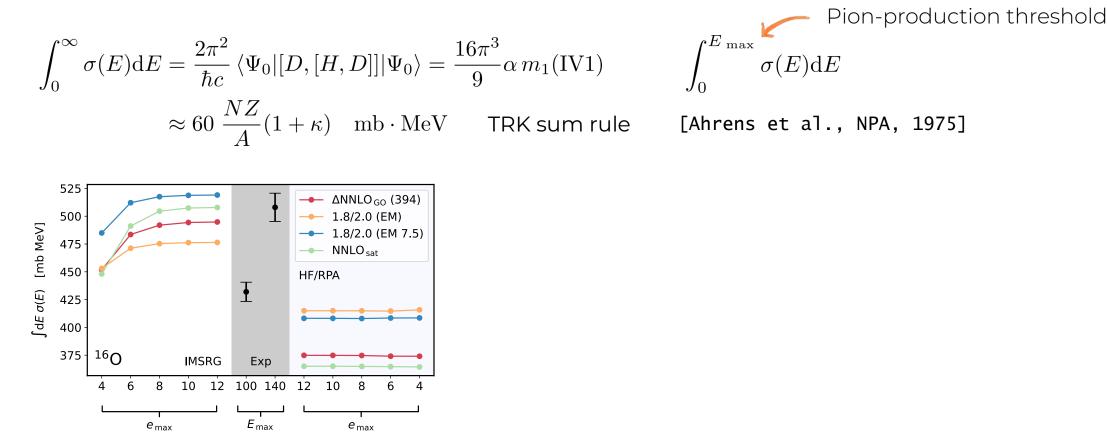
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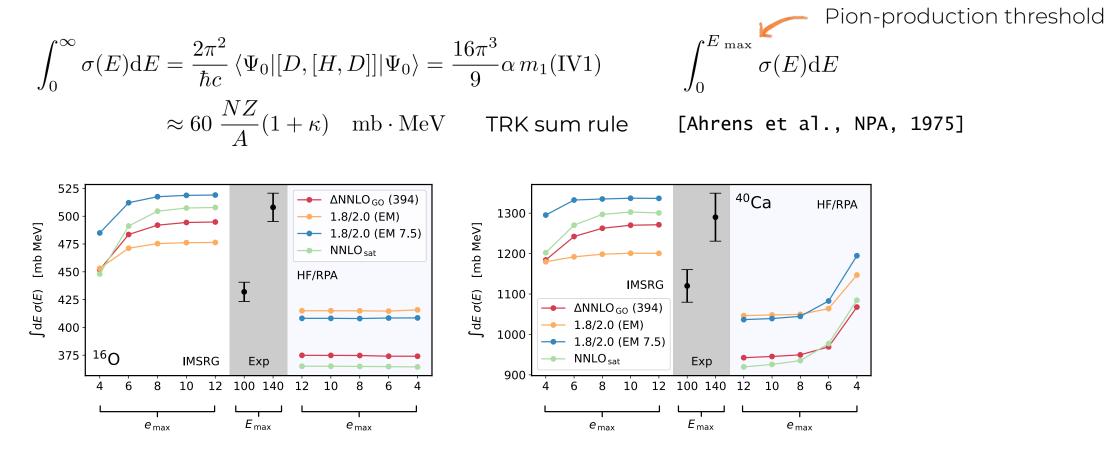
Quadrupole

- Reduced spread
- ~5% correlations effect

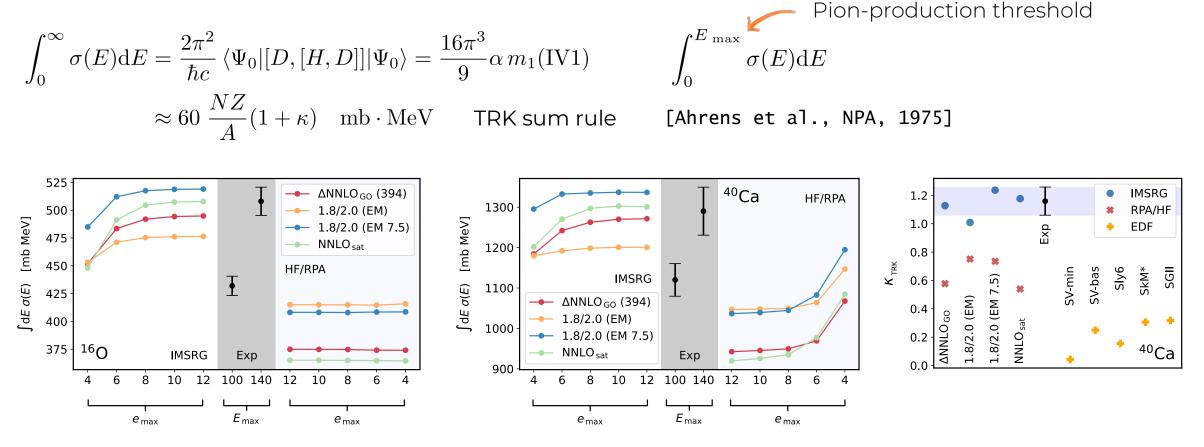
$$\int_0^\infty \sigma(E) dE = \frac{2\pi^2}{\hbar c} \langle \Psi_0 | [D, [H, D]] | \Psi_0 \rangle = \frac{16\pi^3}{9} \alpha \, m_1(\text{IV1})$$
$$\approx 60 \, \frac{NZ}{A} (1+\kappa) \quad \text{mb} \cdot \text{MeV} \qquad \text{TRK sum rule}$$

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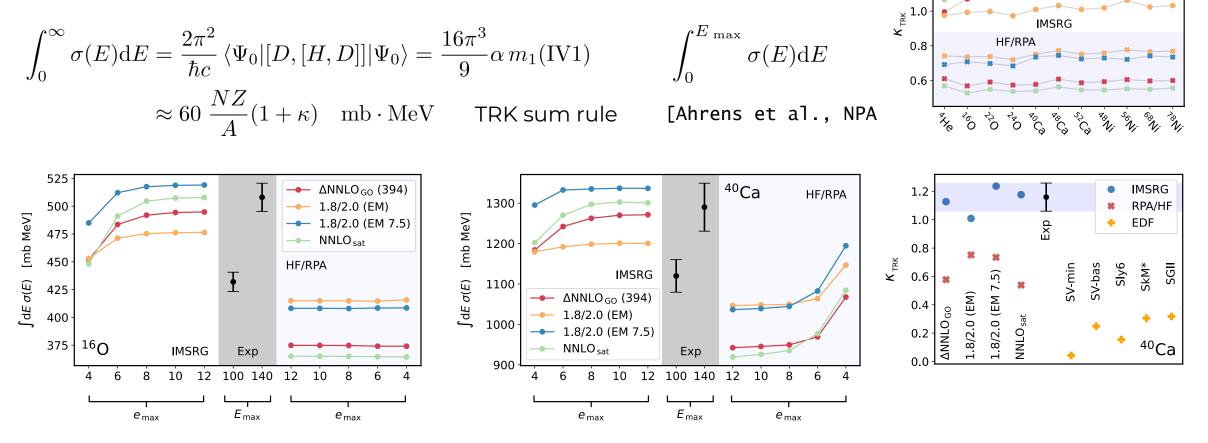
Comparison to exp only makes sense for integrated quantities



Comparison to EDF calculations

[Courtesy of P.-G. Reinhard]

Comparison to exp only makes sense for integrated quantities



Comparison to EDF calculations

• * ΔNNLO_{GO} (394)

• * 1.8/2.0 (EM)

1.4

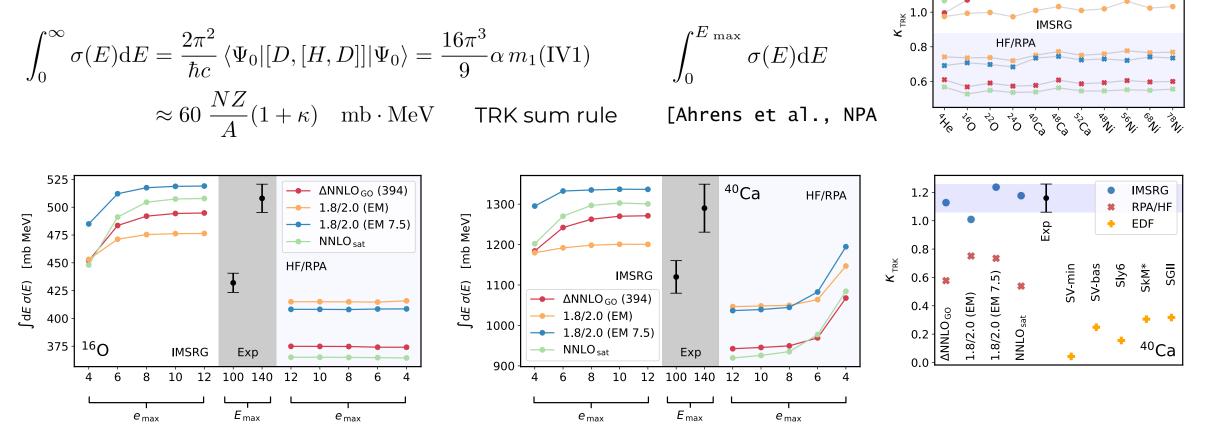
1.2

• * 1.8/2.0 (EM 7.5)

NNLO_{sat}

[Courtesy of P.-G. Reinhard]

Comparison to exp only makes sense for integrated quantities



Both needed for consistent description

- Ground-state correlations
- Commutator expression generates 2-body currents

Comparison to EDF calculations

• * ΔNNLO_{GO} (394)

• * 1.8/2.0 (EM)

1.4

1.2

14

• * 1.8/2.0 (EM 7.5)

NNLO_{sat}

[Courtesy of P.-G. Reinhard]

Sum rules extensively studied in the past

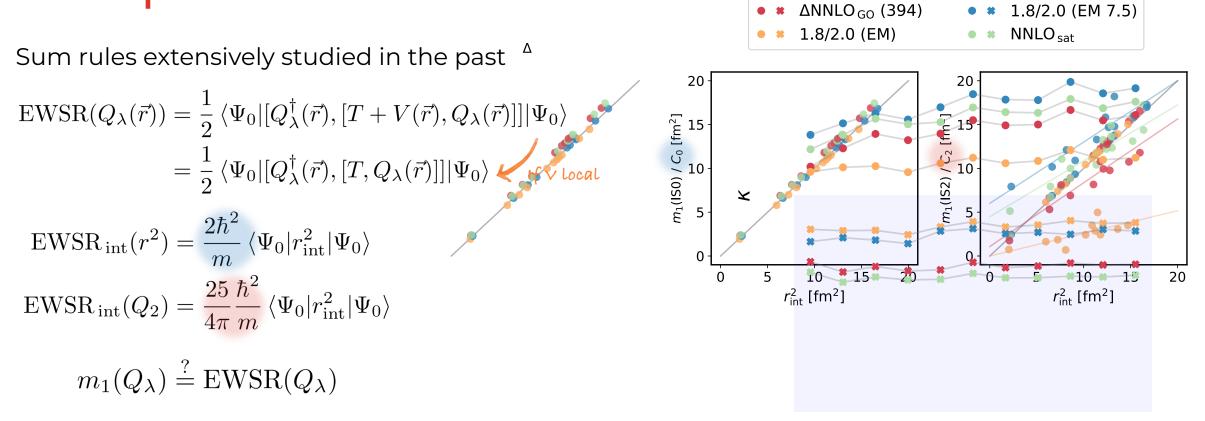
$$\begin{split} \mathrm{EWSR}(Q_{\lambda}(\vec{r})) &= \frac{1}{2} \langle \Psi_{0} | [Q_{\lambda}^{\dagger}(\vec{r}), [T + V(\vec{r}), Q_{\lambda}(\vec{r})]] | \Psi_{0} \rangle \\ &= \frac{1}{2} \langle \Psi_{0} | [Q_{\lambda}^{\dagger}(\vec{r}), [T, Q_{\lambda}(\vec{r})]] | \Psi_{0} \rangle \checkmark_{\mathrm{if}} \mathsf{v} \mathsf{local} \end{split}$$

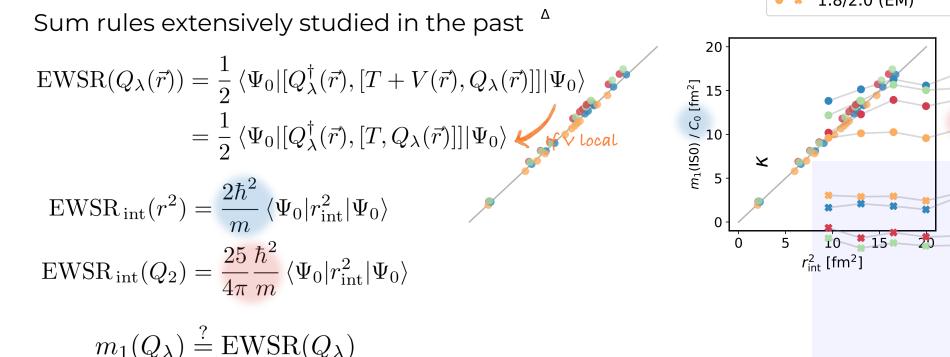
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$$\begin{aligned} \operatorname{EWSR}(Q_{\lambda}(\vec{r})) &= \frac{1}{2} \langle \Psi_{0} | [Q_{\lambda}^{\dagger}(\vec{r}), [T + V(\vec{r}), Q_{\lambda}(\vec{r})]] | \Psi_{0} \rangle \\ &= \frac{1}{2} \langle \Psi_{0} | [Q_{\lambda}^{\dagger}(\vec{r}), [T, Q_{\lambda}(\vec{r})]] | \Psi_{0} \rangle \checkmark \mathsf{from local} \end{aligned}$$
$$\begin{aligned} \operatorname{EWSR}_{\mathrm{int}}(r^{2}) &= \frac{2\hbar^{2}}{m} \langle \Psi_{0} | r_{\mathrm{int}}^{2} | \Psi_{0} \rangle \\ \operatorname{EWSR}_{\mathrm{int}}(Q_{2}) &= \frac{25}{4\pi} \frac{\hbar^{2}}{m} \langle \Psi_{0} | r_{\mathrm{int}}^{2} | \Psi_{0} \rangle \end{aligned}$$

Sum rules extensively studied in the past

$$\begin{aligned} \operatorname{EWSR}(Q_{\lambda}(\vec{r})) &= \frac{1}{2} \langle \Psi_{0} | [Q_{\lambda}^{\dagger}(\vec{r}), [T + V(\vec{r}), Q_{\lambda}(\vec{r})]] | \Psi_{0} \rangle \\ &= \frac{1}{2} \langle \Psi_{0} | [Q_{\lambda}^{\dagger}(\vec{r}), [T, Q_{\lambda}(\vec{r})]] | \Psi_{0} \rangle \checkmark \mathsf{free}_{\mathsf{Ifvecal}} \end{aligned}$$
$$\begin{aligned} \operatorname{EWSR}_{\mathrm{int}}(r^{2}) &= \frac{2\hbar^{2}}{m} \langle \Psi_{0} | r_{\mathrm{int}}^{2} | \Psi_{0} \rangle \\ \operatorname{EWSR}_{\mathrm{int}}(Q_{2}) &= \frac{25}{4\pi} \frac{\hbar^{2}}{m} \langle \Psi_{0} | r_{\mathrm{int}}^{2} | \Psi_{0} \rangle \\ m_{1}(Q_{\lambda}) \stackrel{?}{=} \operatorname{EWSR}(Q_{\lambda}) \end{aligned}$$

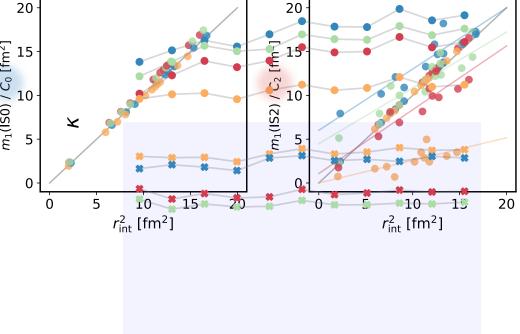


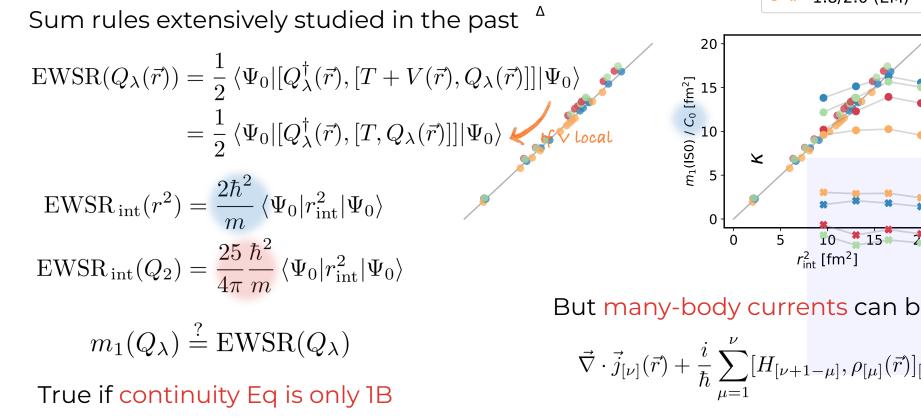


True if continuity Eq is only 1B

$$\vec{\nabla} \cdot \vec{j}(\vec{r}) + \frac{i}{\hbar} [H, \rho(\vec{r})] = 0$$

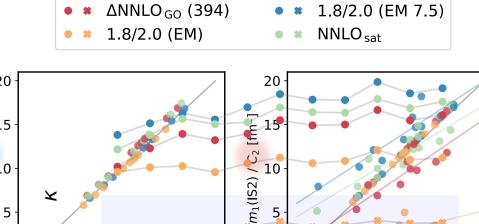
$$\vec{\nabla} \cdot \vec{j}_{[1]}(\vec{r}) + \frac{i}{\hbar} [T, \rho_{[1]}(\vec{r})] = 0$$
If v local





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If \vee local



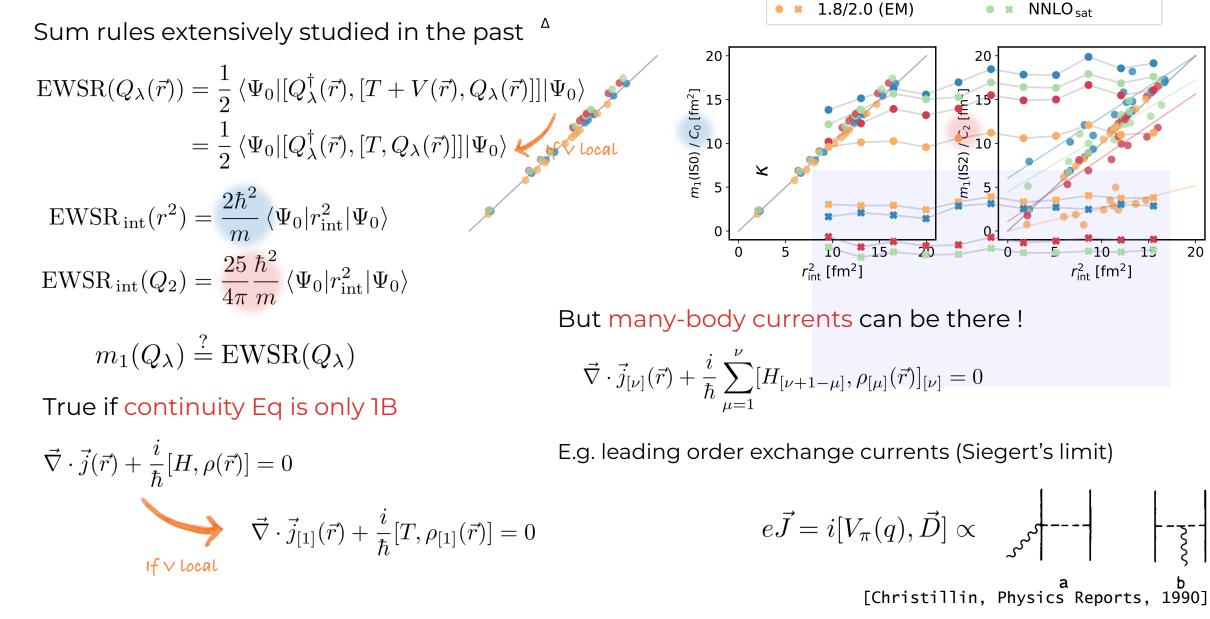
But many-body currents can be there !

$$\vec{\nabla} \cdot \vec{j}_{[\nu]}(\vec{r}) + \frac{i}{\hbar} \sum_{\mu=1}^{\nu} [H_{[\nu+1-\mu]}, \rho_{[\mu]}(\vec{r})]_{[\nu]} = 0$$

*10 × 15

 $r_{\rm int}^2$ [fm²]

20



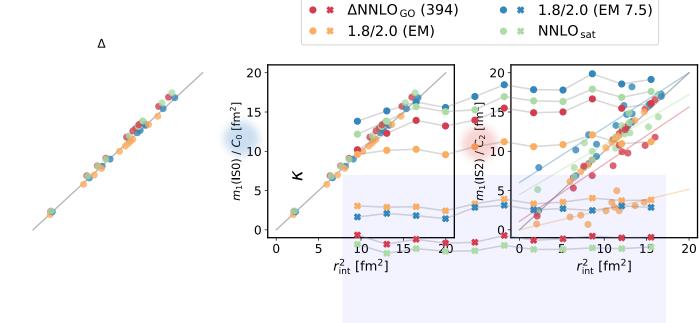
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1.8/2.0 (EM 7.5)

 $\Delta NNLO_{GO}$ (394)

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Let's look more closely



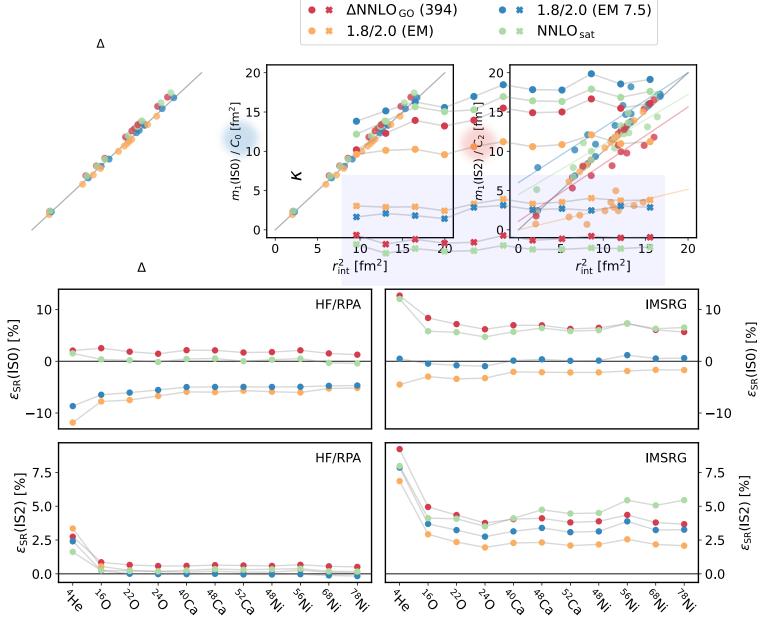
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 $m_1(Q_\lambda) \stackrel{?}{=} \mathrm{EWSR}(Q_\lambda)$

Let's look more closely

Relative difference

$$\varepsilon_{\rm SR}(Q_{\lambda})[\%] \equiv \frac{m_1(Q_{\lambda}) - \mathrm{EWSR}(Q_{\lambda})}{\mathrm{EWSR}(Q_{\lambda})} \times 100$$



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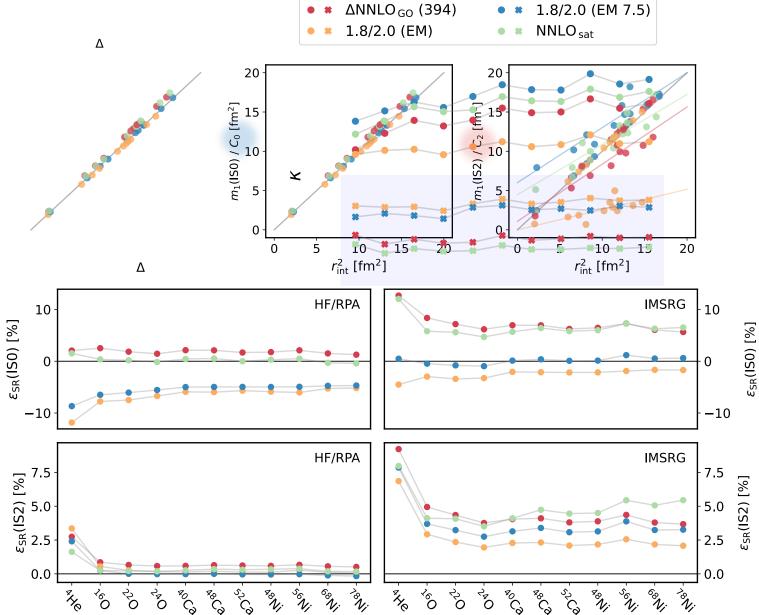
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Differences from EWSR

- Nonlocalities
- Two-body currents



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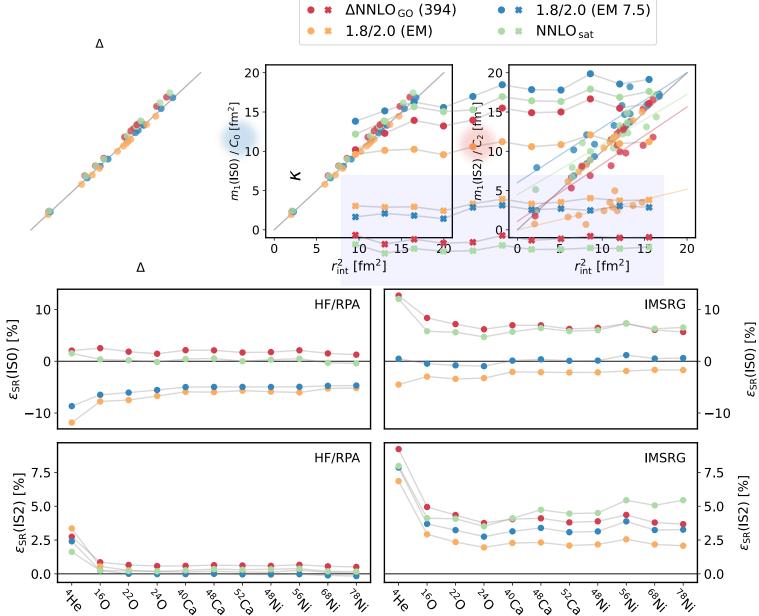
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Differences from EWSR

- Nonlocalities
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m1 from moments is better!



Introduction

- Physics case
- Existing ab initio methods

IMSRG multipole moments

- Moments of the strength
- IMSRG implementation
- Model-space convergence

Numerical results

- Interaction sensitivity
- Comparison to experiment
- Comparison to sum rules

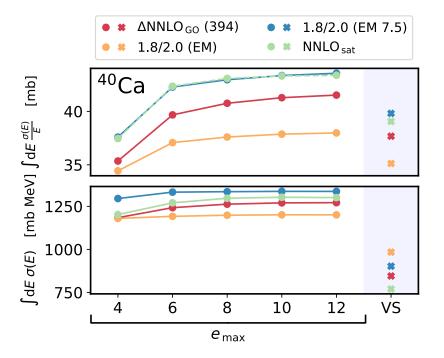
Challenges and opportunities

Going open-shell

Going open-shell

Comparison to VS calculation for ⁴⁰Ca with ²⁸Si core

- Large uncertainties for m_1 and m_0
- Two-step decoupling
- Is the core well described ? (deformation)



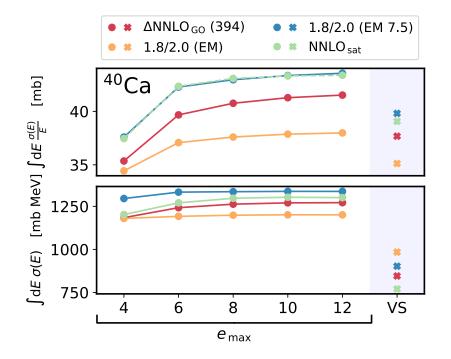
Going open-shell

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Other possibilities within the IMSRG

- Multi-reference formulation
- Symmetry-breaking calculations



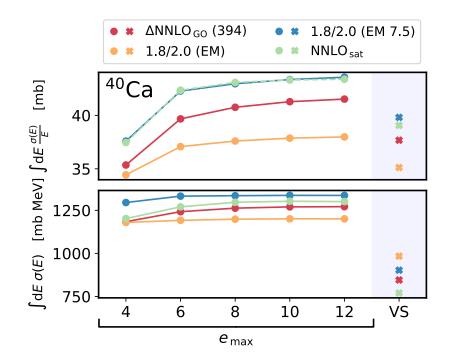
Going open-shell

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Benchmarks and uncertainty quantification

No limitation on the many-body method (and operators) of choice

- Exact treatment of excited states
- Can benchmark response calculations going through exc states EOM, LIT etc. (can send matrix elements)



Going open-shell

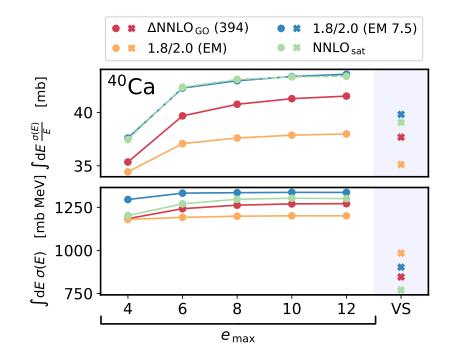
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Systematic studies of H properties



Thank you for the attention



Robert Roth Achim Schwenk Alexander Tichai



Thomas Duguet Jean-Paul Ebran Mikael Frosini Vittorio Somà



Francesca Bonaiti

Summary

- Moment method useful for uncertainty quantification /benchmark of exc states
- Can be used for systematic studies of H properties
- Small finite-basis uncertainty
- Correlation effect ~5% in monopole and quadrupole
- Qualitative change (~40%) for dipole
- Better agreement with data and smaller interaction spread
- Comparison to EWSR diagnostic of non-localities and two-body currents



