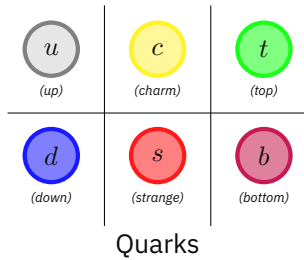


Hadronic Structure and Contour Deformations

Eduardo Ferreira

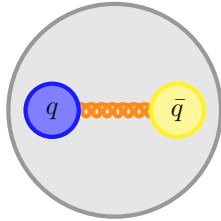
ECT*, Trento – May 26-30, 2025

Hadron Physics



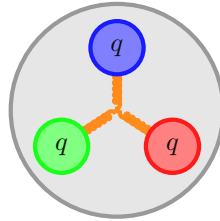
Gluons

Hadron Physics



Mesons

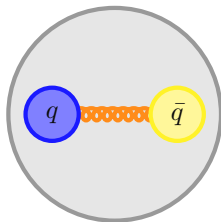
$\pi, \rho, K, \eta, \dots$



Baryons

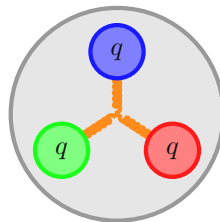
$p, n, \Delta, \Sigma, \Lambda, \dots$

Hadron Physics



Mesons

$\pi, \rho, K, \eta, \dots$



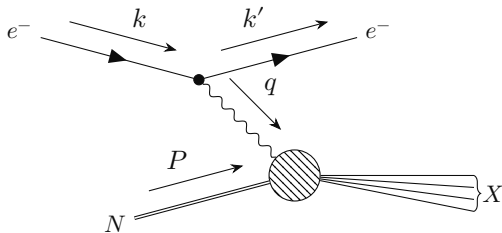
Baryons

$p, n, \Delta, \Sigma, \Lambda, \dots$

- Their dynamics and properties are determined by QCD
 - Atomic nuclei and nuclear stability, EMC effect
 - Radius; spin-, momentum-, charge distributions;
 - Interaction with external currents: $\{e^-, \nu, \dots\}N$ scattering, $N\pi$ scattering, nucleon compton scattering, meson photoproduction

(EIC White Paper), (Eichmann, Sanchis-Alepuz, Williams, Alkofer, Fischer; 2016), (European Muon Collaboration; 1983), ...

Parton model



(PDG Section 18 (Structure Functions))

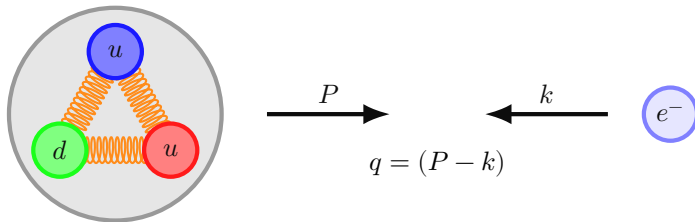
- Common experiment:
Deep inelastic scattering.
- Electron (or neutrino) probe interacts with nucleon via $\gamma, Z, W^\pm$.
- Main idea:
 - Hadrons are *bags* of **partons**.
 - Scattering occurs between the exchanged boson and a parton from the nucleon
- Cross-section separates:

$$\frac{d^2\sigma}{dx dq^2} \propto L_{\mu\nu} W^{\mu\nu}$$

- $W^{\mu\nu}$ described via PDFs.

PDFs/TMDs

- How to they show up?



$$W_{\mu\nu}(p, q, \sigma) \propto \int d^4z \langle p, \sigma | J_\mu^\dagger(z) J_\nu(0) | p, \sigma \rangle = \sum_i \tau(p, q, \sigma)_{\mu\nu}^i f_i(p, q)$$

PDFs/TMDs

- How to they show up?
- Take the collinear, q^2 very large limit.

$$\left| \text{Diagram} \right|^2 = \text{Diagram} \otimes \text{Diagram} + \mathcal{O}\left(\frac{k_{\perp}^2}{q^2}\right)$$

$$\sigma \propto \sum_f \int \frac{dx}{x} H(x, Q^2/\Lambda^2) q_f(x, \Lambda^2) + \mathcal{O}\left(\frac{k_{\perp}^2}{Q^2}\right)$$

$$q_f(x, \Lambda) \propto \int \frac{d^4\ell}{(2\pi)^4} \delta\left(x - \frac{\ell^+}{p^+}\right) \text{Tr} \left[\langle p | \bar{\psi}_f(\ell) \gamma^+ \psi_f(\ell) | p \rangle \right]$$

- Hard scale $Q^2 = -q^2$ introduced allows for factorization – power suppression $\mathcal{O}(Q^{-2})$
- Works also on other processes, like Drell-Yan.

PDFs/TMDs

- How to they show up?
- Take the collinear, q^2 very large limit.

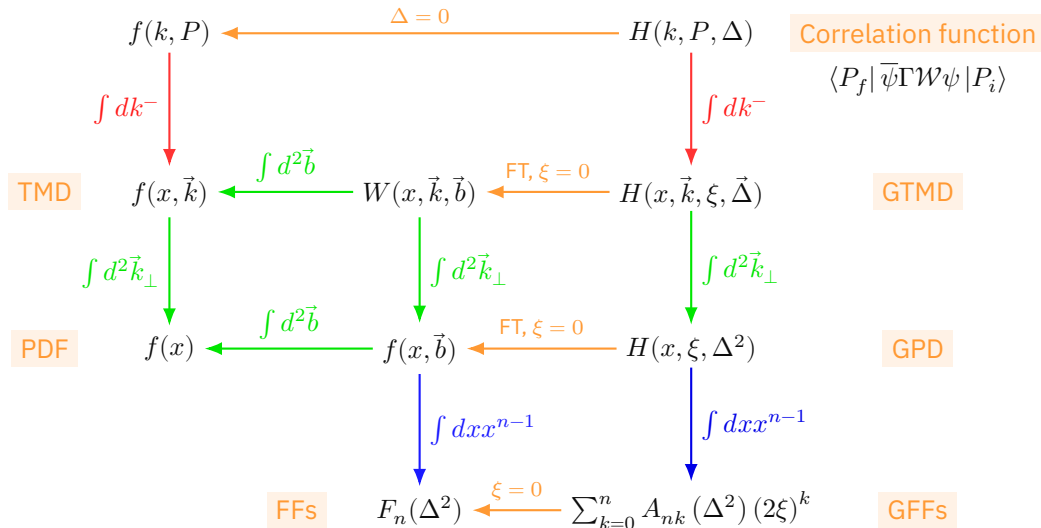
$$\left| \text{Diagram} \right|^2 = \text{Diagram} \otimes \text{Diagram} + \mathcal{O}\left(\frac{k_{\perp}^2}{q^2}\right)$$

$$\sigma \propto \sum_f \int \frac{dx}{x} H(x, Q^2/\Lambda^2) q_f(x, \Lambda^2) + \mathcal{O}\left(\frac{k_{\perp}^2}{Q^2}\right)$$

$$q_f(x, \Lambda) \propto \int \frac{d^4\ell}{(2\pi)^4} \delta\left(x - \frac{\ell^+}{p^+}\right) \text{Tr} \left[\langle p | \bar{\psi}_f(\ell) \gamma^+ \psi_f(\ell) | p \rangle \right]$$

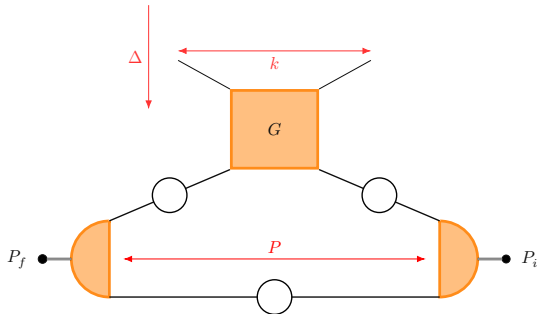
- Focus on the non-perturbative hadronic part.
- Can we build them directly from first-principles **FUN**ctional Methods calculations?

Hadronic quantities



Our Goal

- **Main Goal:** Get partonic distribution functions from hadron-hadron correlations via **FUN**ctional Methods



- G is the four-point quark correlation function, calculated with scattering equation.
- The BSWF is calculated via the meson BSE.
- First scalar toy model, then QCD.

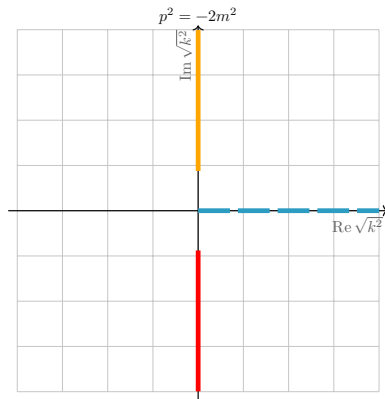
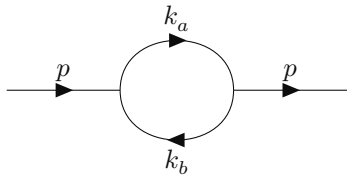
(Mezrag; 2015), (Diehl, Gousset; 1998), (Tiburzi, Miller; 2003),
(Mezrag, Chang, Moutarde, Roberts, Rodríguez-Quintero, Sabatié, Schmidt; 2015),
many others, ...

$$\mathcal{G}^{[\Gamma]}(P, k, \Delta) = \frac{1}{2} \text{Tr} \left[\int dk^- \int \frac{d^4 z}{2\pi^4} e^{ik \cdot z} \langle P_f | \bar{\psi}(z) \mathcal{W} \Gamma \psi(0) | P_i \rangle \right]$$

- Partonic distributions are calculated by integrating the correlator in k^- and taking appropriate traces.

Analytic Structure is Important

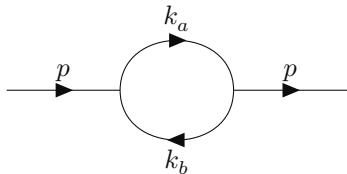
- Calculating quantities in time-like momenta is **complicated** !
 - Analytic structure prevents naïve integration.
 - Poles and branch cuts are present.
- Calculate quantities in the P^2 complex plane.
- **Euclidean** $\Leftrightarrow$ Minkowski



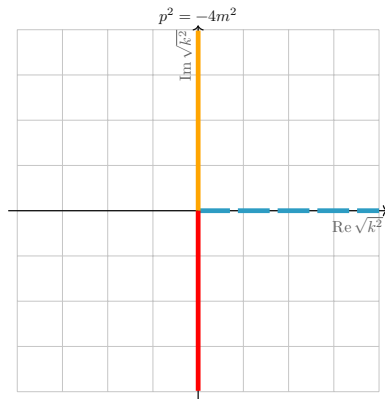
$$I(p^2) = \int d^4k \frac{1}{k_a^2 - m^2 + i\epsilon} \frac{1}{k_b^2 - m^2 + i\epsilon}$$

Analytic Structure is Important

- Calculating quantities in time-like momenta is **complicated** !
 - Analytic structure prevents naïve integration.
 - Poles and branch cuts are present.
- Calculate quantities in the P^2 complex plane.
- **Euclidean** $\Leftrightarrow$ Minkowski

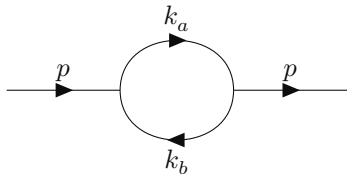


$$I(p^2) = \int d^4k \frac{1}{k_a^2 - m^2 + i\epsilon} \frac{1}{k_b^2 - m^2 + i\epsilon}$$

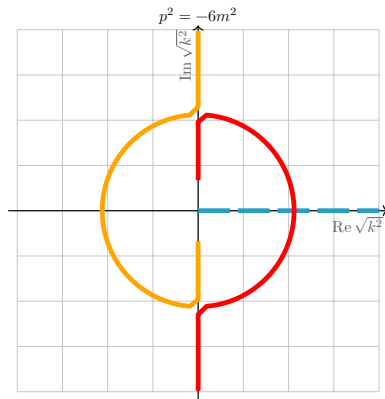


Analytic Structure is Important

- Calculating quantities in time-like momenta is **complicated** !
 - Analytic structure prevents naïve integration.
 - Poles and branch cuts are present.
- Calculate quantities in the P^2 complex plane.
- **Euclidean** $\Leftrightarrow$ Minkowski



$$I(p^2) = \int d^4k \frac{1}{k_a^2 - m^2 + i\epsilon} \frac{1}{k_b^2 - m^2 + i\epsilon}$$



Worked Example

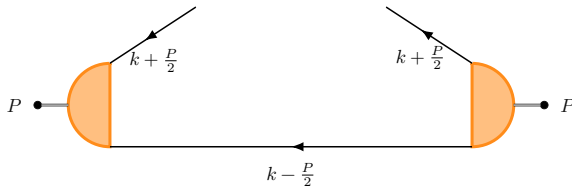
- Simple *tree-level* version of hadronic matrix element.

Worked Example

- Simple *tree-level* version of hadronic matrix element.
- Understand Minkowski $\Leftrightarrow$ Euclidean.

Worked Example

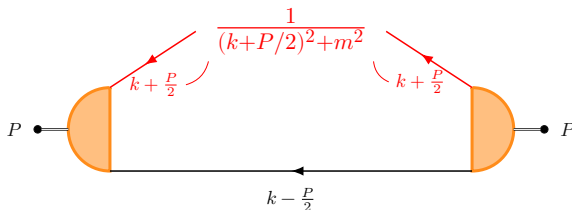
- Simple *tree-level* version of hadronic matrix element.
- Understand Minkowski $\Leftrightarrow$ Euclidean.
- All elements known – **Residue Integration**



$$\mathcal{G}^{[\Gamma]}(P, k, \Delta = 0) = \int dq^- \langle P | \bar{\phi}\left(q + \frac{P}{2}\right) \phi\left(q + \frac{P}{2}\right) | P \rangle \Big|_{q_{\perp} = k_{\perp}, q^+ = \frac{\alpha}{2} P^+}$$

Worked Example

- Simple *tree-level* version of hadronic matrix element.
- Understand Minkowski $\Leftrightarrow$ Euclidean.
- All elements known – **Residue Integration**

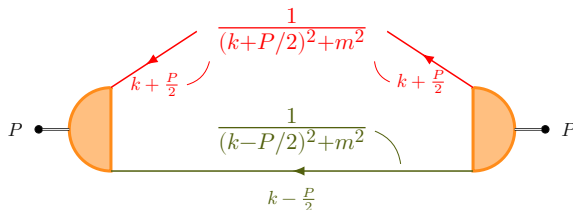


$$\mathcal{G}^{[\Gamma]}(\alpha, k_{\perp}) = \int dq^{-} \left[\frac{1}{\left(q + \frac{P}{2}\right)^2 + m^2} \right]^2$$

$$q^{+} = \frac{\alpha}{2} P^{+}, q_{\perp} = k_{\perp}$$

Worked Example

- Simple *tree-level* version of hadronic matrix element.
- Understand Minkowski $\Leftrightarrow$ Euclidean.
- All elements known – **Residue Integration**

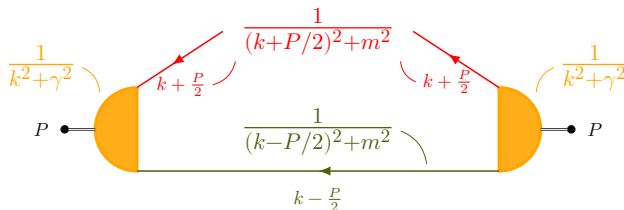


$$\mathcal{G}^{[\Gamma]}(\alpha, k_{\perp}) = \int dq^- \left[\frac{1}{(q + \frac{P}{2})^2 + m^2} \right]^2 \frac{1}{(q - \frac{P}{2})^2 + m^2}$$

$$q^+ = \frac{\alpha}{2} P^+, q_{\perp} = k_{\perp}$$

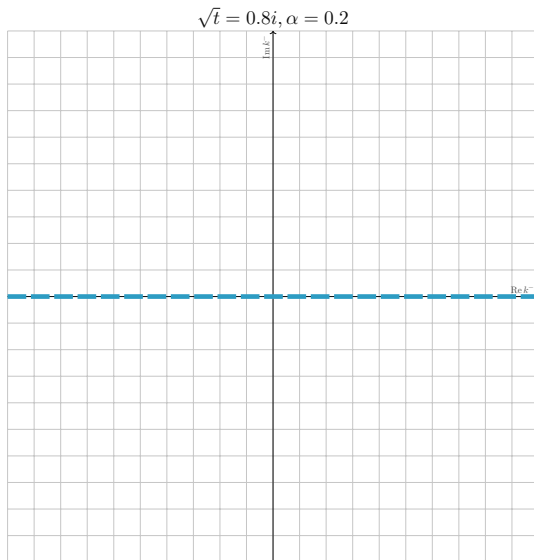
Worked Example

- Simple *tree-level* version of hadronic matrix element.
- Understand Minkowski $\Leftrightarrow$ Euclidean.
- All elements known – **Residue Integration**



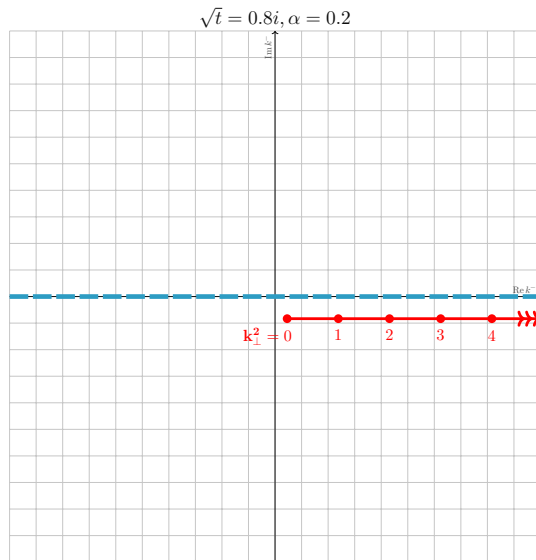
$$\mathcal{G}^{[\Gamma]}(\alpha, k_{\perp}) = \int dq^- \left[\frac{1}{\left(q + \frac{P}{2}\right)^2 + m^2} \right]^2 \frac{1}{\left(q - \frac{P}{2}\right)^2 + m^2} \left[\frac{1}{q^2 + \gamma^2} \right]^2 \Bigg|_{q^+ = \frac{\alpha}{2} P^+, q_{\perp} = k_{\perp}}$$

Minkowski/LF



$$\text{PDF} \propto \int_{-\infty(1+i\epsilon)}^{\infty(1+i\epsilon)} dk^- \int d^2 k_{\perp} \langle P | \bar{\phi} \phi | P \rangle$$

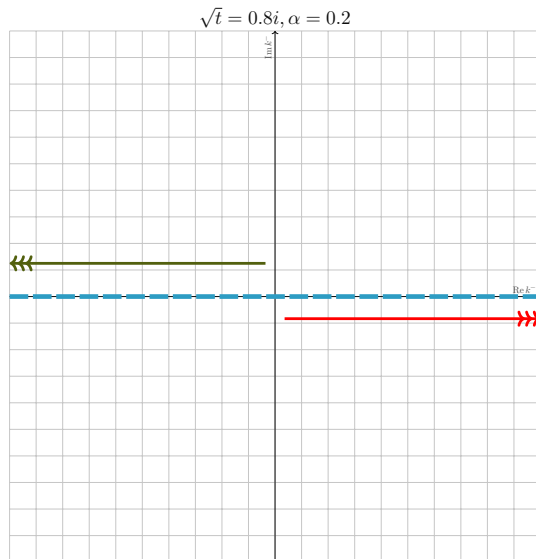
Minkowski/LF



$$\text{PDF} \propto \int_{-\infty(1+i\epsilon)}^{\infty(1+i\epsilon)} dk^- \int d^2k_{\perp} \langle P | \bar{\phi} \phi | P \rangle$$

- Poles and cuts from the elements before: propagators
- $\alpha, P^2 = -4m^2t$ fixed.

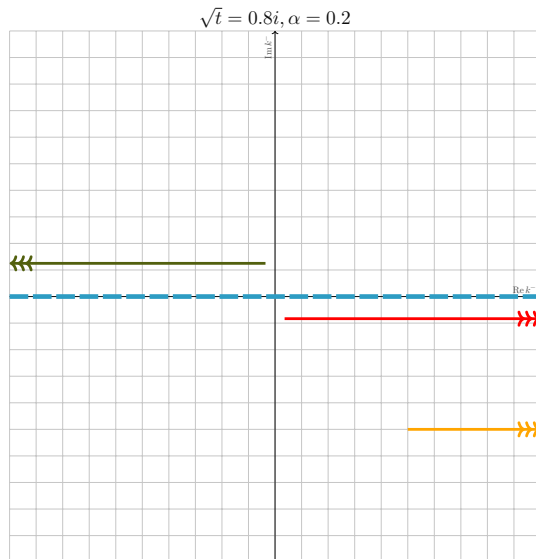
Minkowski/LF



$$\text{PDF} \propto \int_{-\infty(1+i\epsilon)}^{\infty(1+i\epsilon)} dk^- \int d^2k_{\perp} \langle P | \bar{\phi} \phi | P \rangle$$

- Poles and cuts from the elements before: propagators
- $\alpha, P^2 = -4m^2t$ fixed.

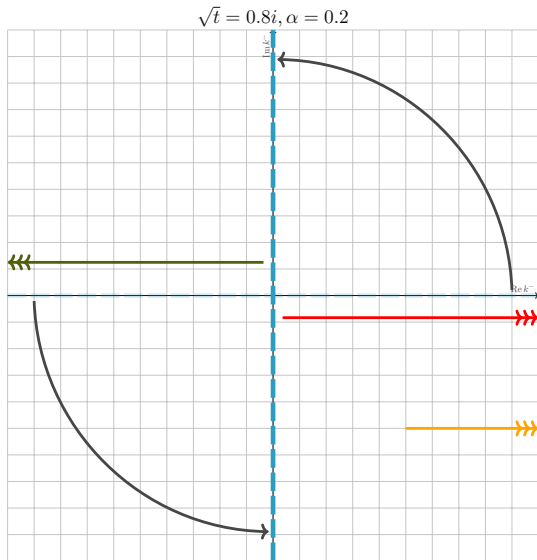
Minkowski/LF



$$\text{PDF} \propto \int_{-\infty(1+i\epsilon)}^{\infty(1+i\epsilon)} dk^- \int d^2k_{\perp} \langle P | \bar{\phi} \phi | P \rangle$$

- Poles and cuts from the elements before: propagators and monopole amplitudes.
- $\alpha, P^2 = -4m^2t$ fixed.
- Result is given by summing residues for poles above the real axis.

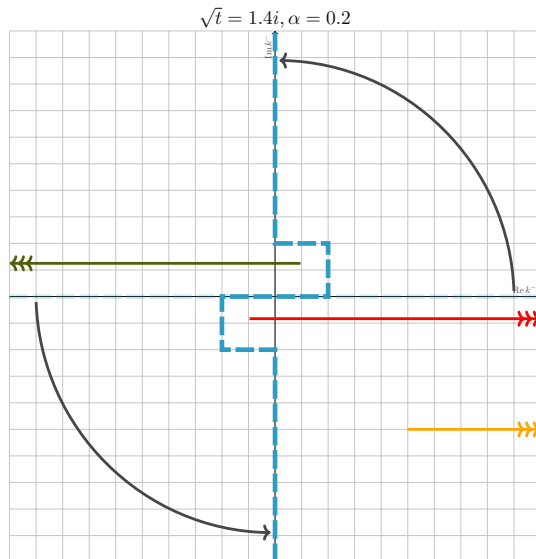
Minkowski/LF



$$\text{PDF} \propto \int_{-\infty(1+i\epsilon)}^{\infty(1+i\epsilon)} dk^- \int d^2k_{\perp} \langle P | \bar{\phi} \phi | P \rangle$$

- Poles and cuts from the elements before: propagators and monopole amplitudes.
- $\alpha, P^2 = -4m^2t$ fixed.
- Result is given by summing residues for poles above the real axis.
- Can Wick Rotate (may need contour deformations).

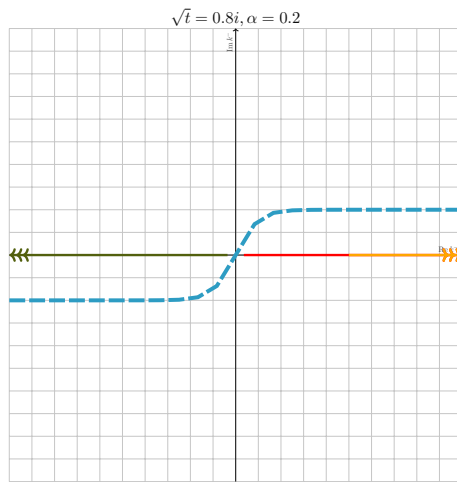
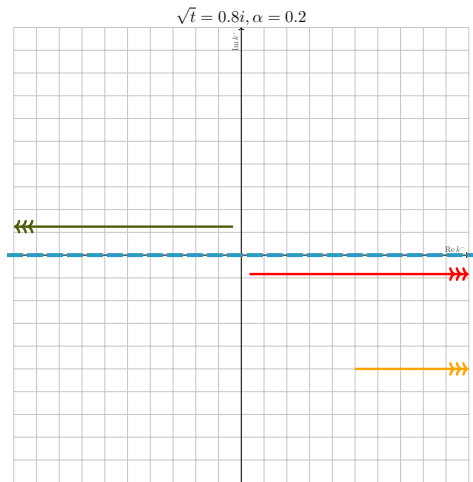
Minkowski/LF



$$\text{PDF} \propto \int_{-\infty(1+i\epsilon)}^{\infty(1+i\epsilon)} dk^- \int d^2k_{\perp} \langle P | \bar{\phi} \phi | P \rangle$$

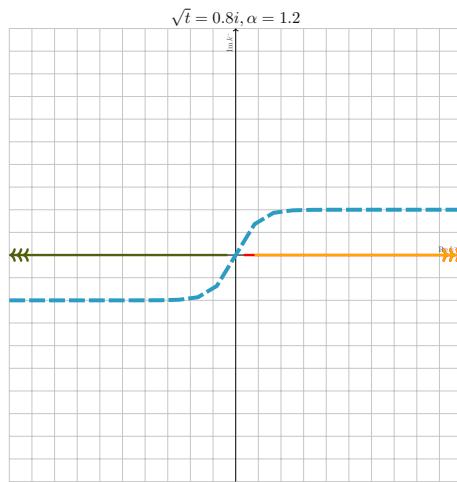
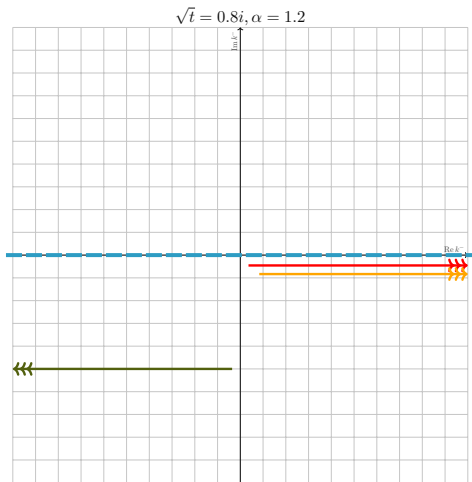
- Poles and cuts from the elements before: propagators and monopole amplitudes.
- $\alpha, P^2 = -4m^2t$ fixed.
- Result is given by summing residues for poles above the real axis.
- Can Wick Rotate (may need contour deformations).

Imaginary time boundary conditions vs $i\epsilon$ prescription



- Both match on $|\alpha| \leq 1$. Only complex boundaries give analytic function in α .

Imaginary time boundary conditions vs $i\epsilon$ prescription



- Both match on $|\alpha| \leq 1$. Only complex boundaries give analytic function in α .

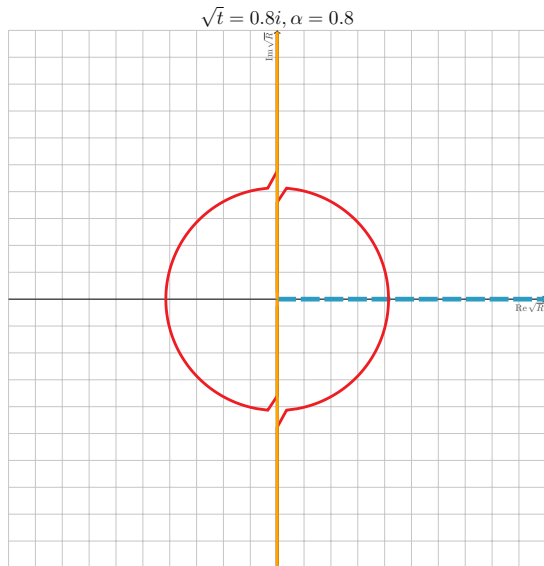
Euclidean

$$k_{\perp}^2 = R \quad k^- \propto \hat{k} \cdot \hat{P} = Z \quad P^2 = 4mt^2$$

$$\text{PDF} \propto \int_0^{\infty} dR \int_{-\infty}^{\infty} dZ \langle P | \bar{\phi} \phi | P \rangle$$

- Use only Lorentz invariants.
- Hyperspherical coordinates in k .
- We check cuts on $k_{\perp}^2 = R$.
 - Avoiding cuts on R means selecting the correct residues in Minkowski.

Euclidean

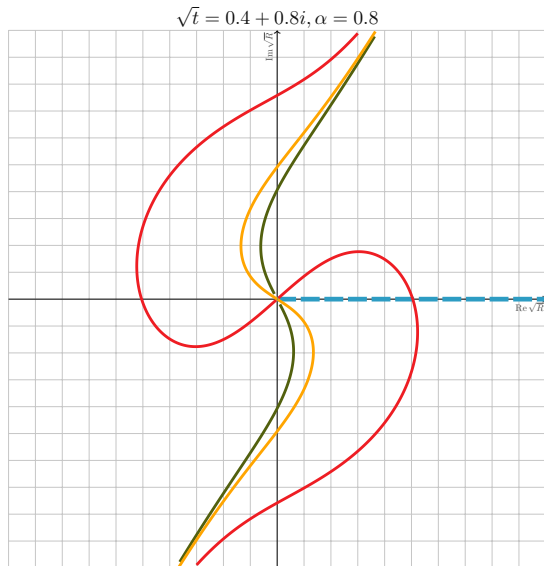


$$k_{\perp}^2 = R \quad k^{-} \propto \hat{k} \cdot \hat{P} = Z \quad P^2 = 4mt^2$$

$$\text{PDF} \propto \int_0^{\infty} dR \int_{-\infty}^{\infty} dZ \langle P | \bar{\phi} \phi | P \rangle$$

- Use only Lorentz invariants.
- Hyperspherical coordinates in k .
- We check cuts on $k_{\perp}^2 = R$.
 - Avoiding cuts on R means selecting the correct residues in Minkowski.

Euclidean

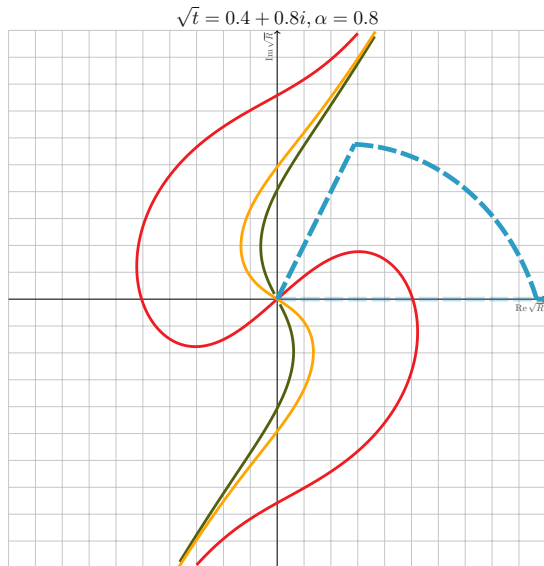


$$k_{\perp}^2 = R \quad k^- \propto \hat{k} \cdot \hat{P} = Z \quad P^2 = 4mt^2$$

$$\text{PDF} \propto \int_0^{\infty} dR \int_{-\infty}^{\infty} dZ \langle P | \bar{\phi} \phi | P \rangle$$

- Use only Lorentz invariants.
- Hyperspherical coordinates in k .
- We check cuts on $k_{\perp}^2 = R$.
 - Avoiding cuts on R means selecting the correct residues in Minkowski.
- We need slightly complex masses.

Euclidean

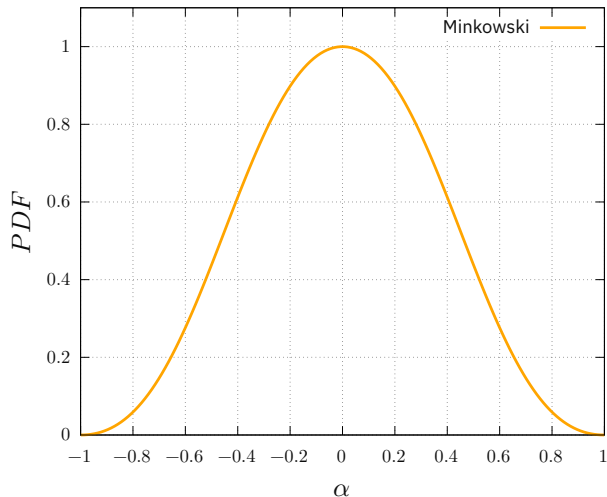


$$k_{\perp}^2 = R \quad k^{-} \propto \hat{k} \cdot \hat{P} = Z \quad P^2 = 4mt^2$$

$$\text{PDF} \propto \int_0^{\infty} dR \int_{-\infty}^{\infty} dZ \langle P | \bar{\phi} \phi | P \rangle$$

- Use only Lorentz invariants.
- Hyperspherical coordinates in k .
- We check cuts on $k_{\perp}^2 = R$.
 - Avoiding cuts on R means selecting the correct residues in Minkowski.
- We need slightly complex masses.

Euclidean

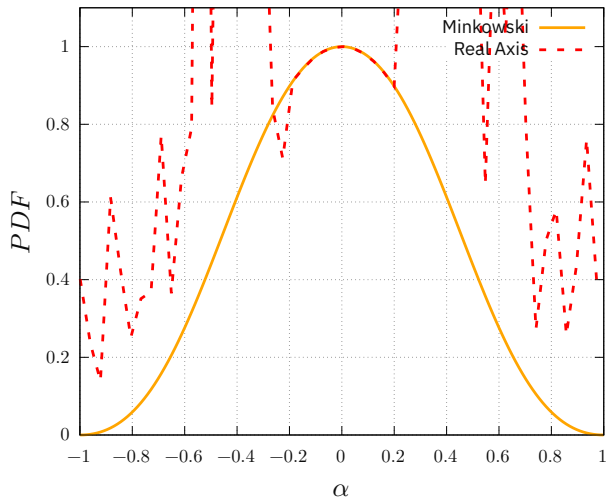


$$k_{\perp}^2 = R \quad k^- \propto \hat{k} \cdot \hat{P} = Z \quad P^2 = 4mt^2$$

$$\text{PDF} \propto \int_0^{\infty} dR \int_{-\infty}^{\infty} dZ \langle P | \bar{\phi} \phi | P \rangle$$

- Use only Lorentz invariants.
- Hyperspherical coordinates in k .
- We check cuts on $k_{\perp}^2 = R$.
 - Avoiding cuts on R means selecting the correct residues in Minkowski.
- We need slightly complex masses.

Euclidean

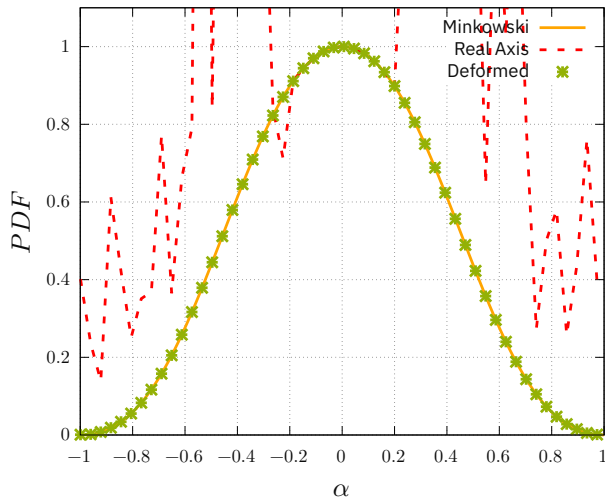


$$k_{\perp}^2 = R \quad k^- \propto \hat{k} \cdot \hat{P} = Z \quad P^2 = 4mt^2$$

$$\text{PDF} \propto \int_0^{\infty} dR \int_{-\infty}^{\infty} dZ \langle P | \bar{\phi} \phi | P \rangle$$

- Use only Lorentz invariants.
- Hyperspherical coordinates in k .
- We check cuts on $k_{\perp}^2 = R$.
 - Avoiding cuts on R means selecting the correct residues in Minkowski.
- We need slightly complex masses.

Euclidean



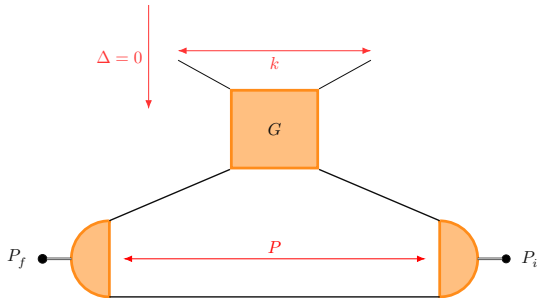
$$k_{\perp}^2 = R \quad k^- \propto \hat{k} \cdot \hat{P} = Z \quad P^2 = 4mt^2$$

$$\text{PDF} \propto \int_0^{\infty} dR \int_{-\infty}^{\infty} dZ \langle P | \bar{\phi} \phi | P \rangle$$

- Use only Lorentz invariants.
- Hyperspherical coordinates in k .
- We check cuts on $k_{\perp}^2 = R$.
 - Avoiding cuts on R means selecting the correct residues in Minkowski.
- We need slightly complex masses.

Re-stating

- Get partonic distribution functions from hadron-hadron correlations via **Functional Methods**



- TMDs/PDFs – $\Delta = 0$.
- G is the four-point quark correlation function – added later.
- Tree level propagators.
- First scalar toy model, then QCD.

(Mezrag; 2015), (Diehl, Gousset; 1998), (Tiburzi, Miller; 2003),
(Mezrag, Chang, Moutarde, Roberts, Rodríguez-Quintero, Sabatié, Schmidt; 2015),
many others, ...

- Partonic distributions are calculated by integrating the correlator in k^- and taking appropriate traces.
- This means: integral in $(-\infty, \infty)$ in $\hat{k} \cdot \hat{P}$ – **Analytic Continuation needed**

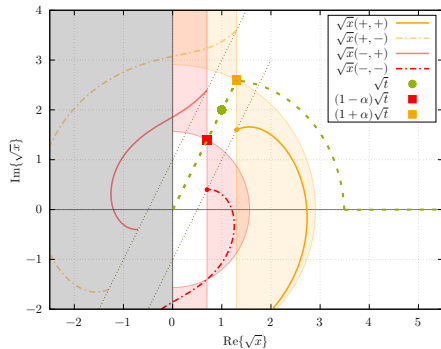
(Eichmann, EF, Stadler; 2022)

Analytic Structure of BSE $\text{Im} \sqrt{t} > \min \left(\frac{1}{1 \pm \alpha} \right)$

$$\mathbf{G}_0^{-1} = (l_+^2 + m^2)(l_-^2 + m^2)$$

- Branch cuts in complex $\sqrt{l}$ plane

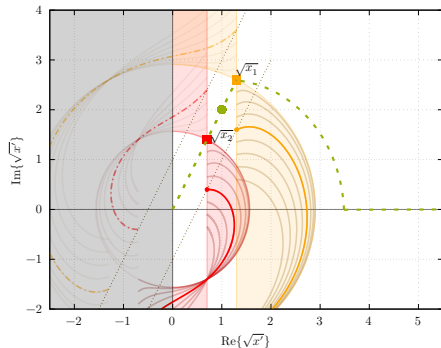
$$\sqrt{l}_\pm^\lambda = \mp(1 \pm \alpha)\sqrt{t} \left[z_l + i\lambda \sqrt{1 - z_l^2 + \frac{1}{(1 \pm \alpha)^2 t}} \right]$$



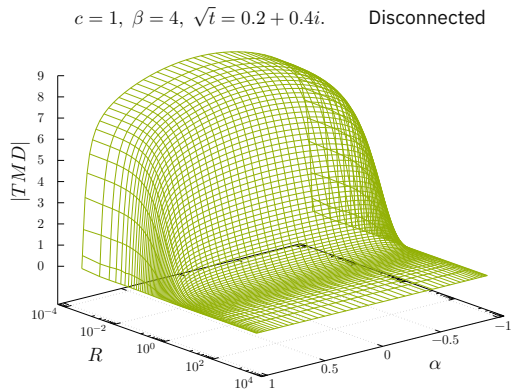
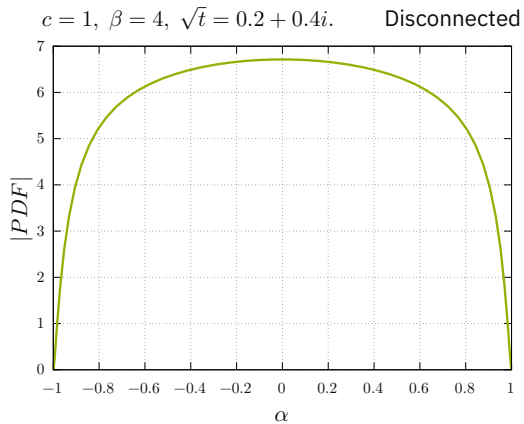
$$\mathbf{K}^{-1} = (q - l)^2 + \mu^2$$

- Branch cuts in complex $\sqrt{l}$ plane depend on path taken:

$$\sqrt{l} = \sqrt{\rho} \left(\Omega \pm i \sqrt{1 - \Omega^2 + \frac{\beta^2}{\rho}} \right)$$



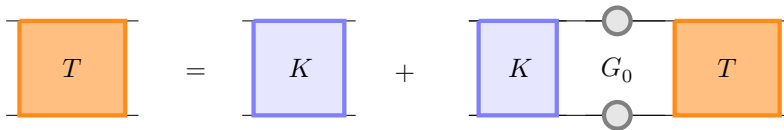
First Results Without 4-point function



4-point function

- 4- point function determined from scattering equation:

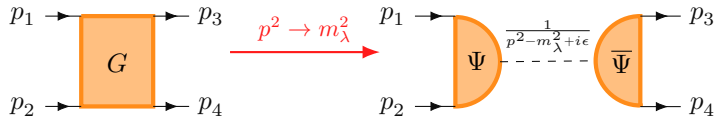
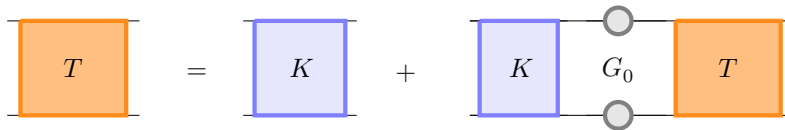
$$\mathbf{G} = \mathbf{G}_0 + \mathbf{G}_0 \mathbf{T} \mathbf{G}_0 \implies \mathbf{T} = \mathbf{K} + \mathbf{K} \mathbf{G}_0 \mathbf{T} \implies \mathbf{T} = (\mathbb{1} - \mathbf{K} \mathbf{G}_0)^{-1} \mathbf{K}.$$



4-point function

- 4-point function determined from scattering equation:

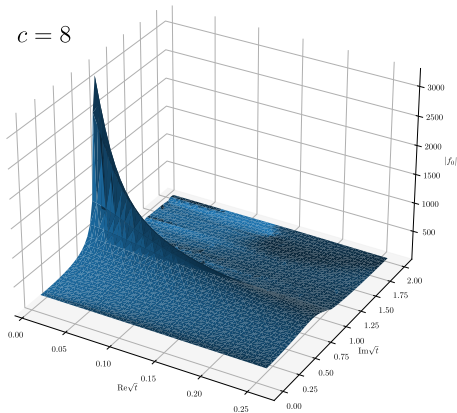
$$G = G_0 + G_0 T G_0 \Rightarrow T = K + K G_0 T \Rightarrow T = (\mathbb{1} - K G_0)^{-1} K.$$



- Same G_0 and K as in the BSE
- All dynamics of 2 ϕ particles:
 - Must produce bound state poles dynamically!**

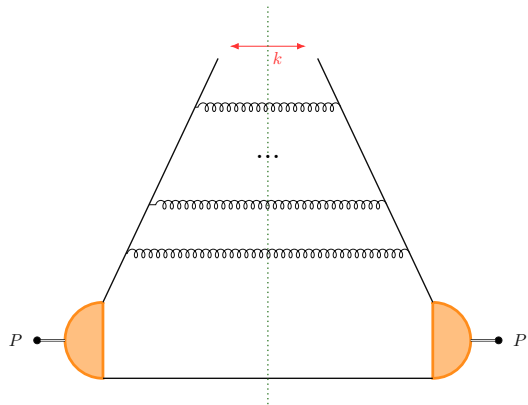
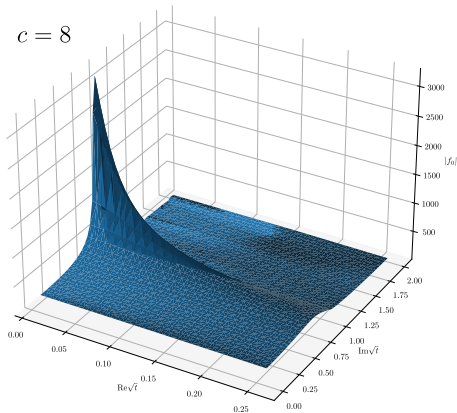
4-point function

- Partial-wave expansion shows bound-state pole in the first Riemann sheet



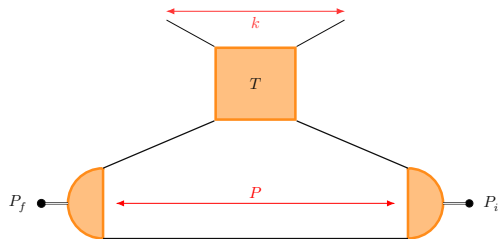
4-point function

- Partial-wave expansion shows bound-state pole in the first Riemann sheet



- Describes both long-range and short-range qq dynamics.
- Further Fock states?

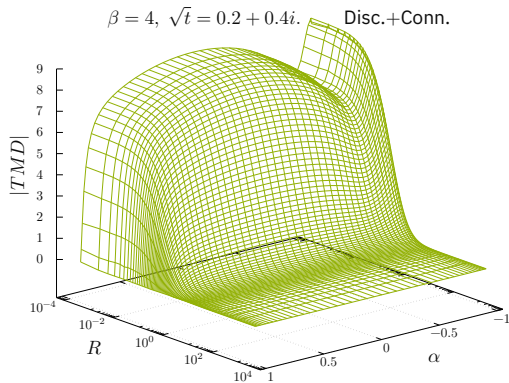
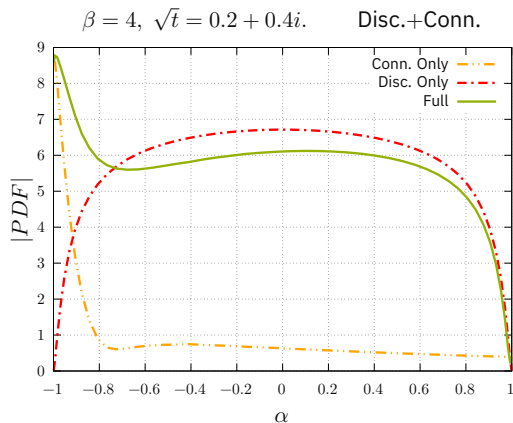
Put everything together



1. Calculate 4-point function
2. Calculate BSE
3. Do the loop integration
4. Project to LF

- We solve one triangle for each point in the R, α grid – **HPC Needed!**
 - $N_\alpha \times N_R \times N_Z$ 4-point functions!
- Obtained PDFs/TMDs for a system of two bound scalars ϕ .

Results *With 4-point function*



■ Publication coming soon!

Conclusions

- Framework for calculation of PDFs/TMDs.
- Applications to hadronic structure calculations starting from self-consistent first principles calculations.
- Calculation of $q\bar{q}$ scattering equation – rich dynamics possible.
- Future applications to realistic QCD models soon.