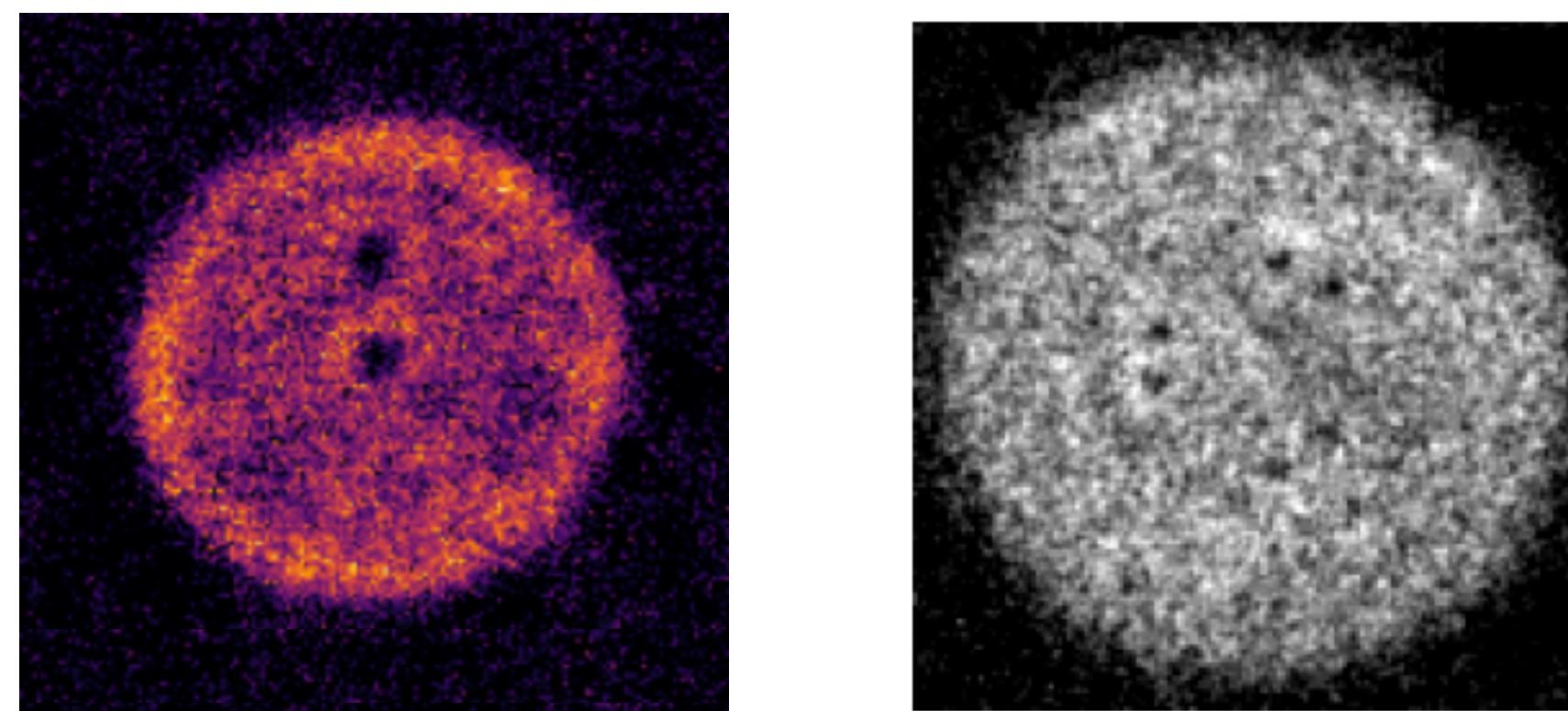


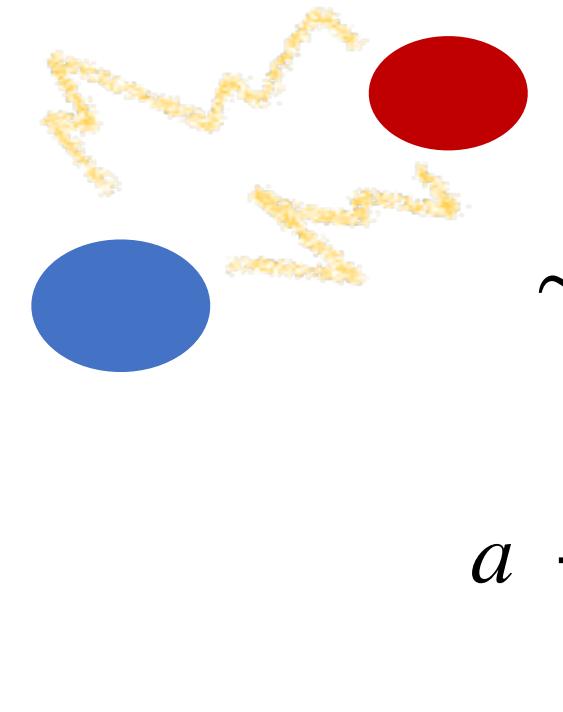
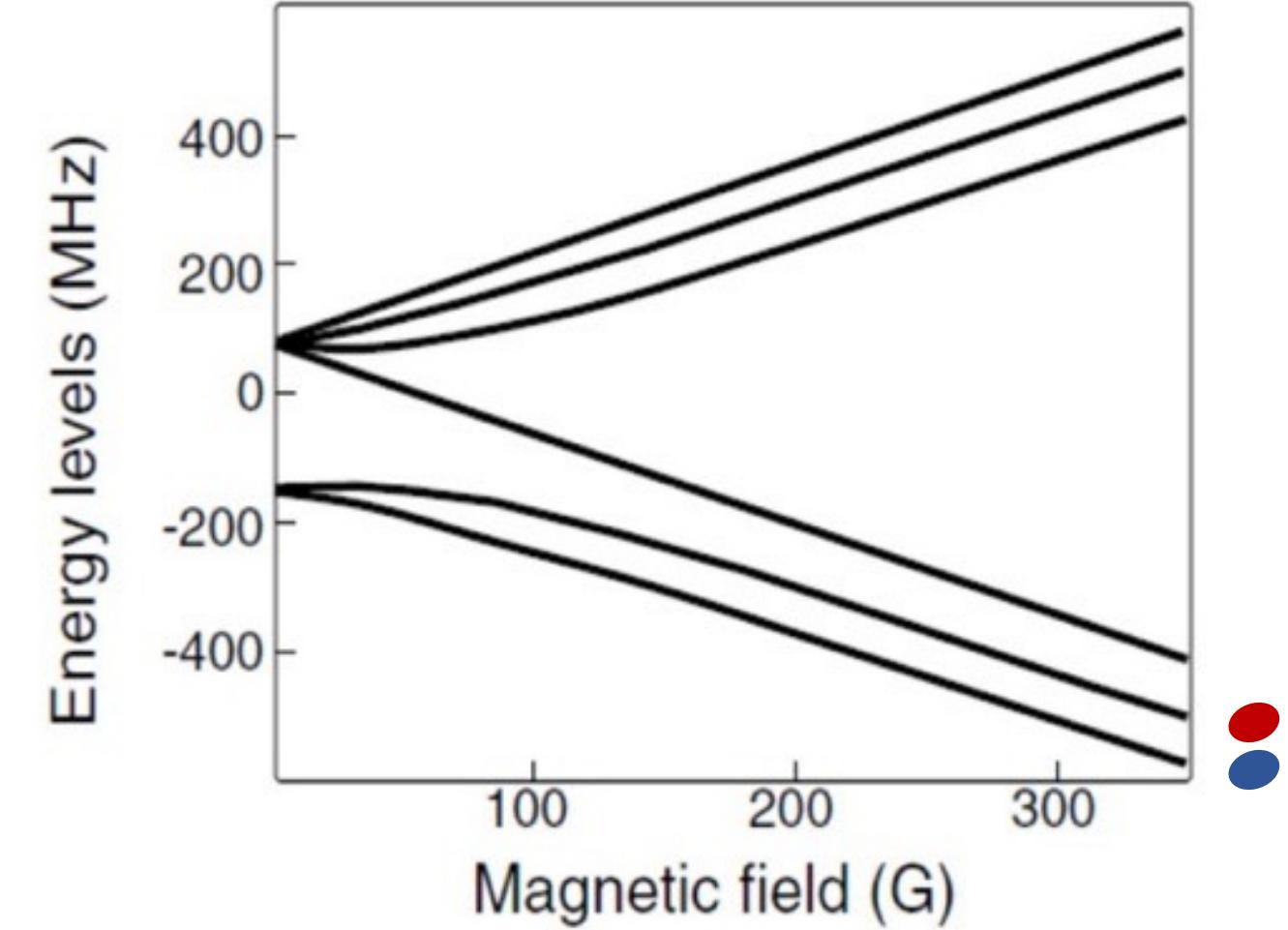
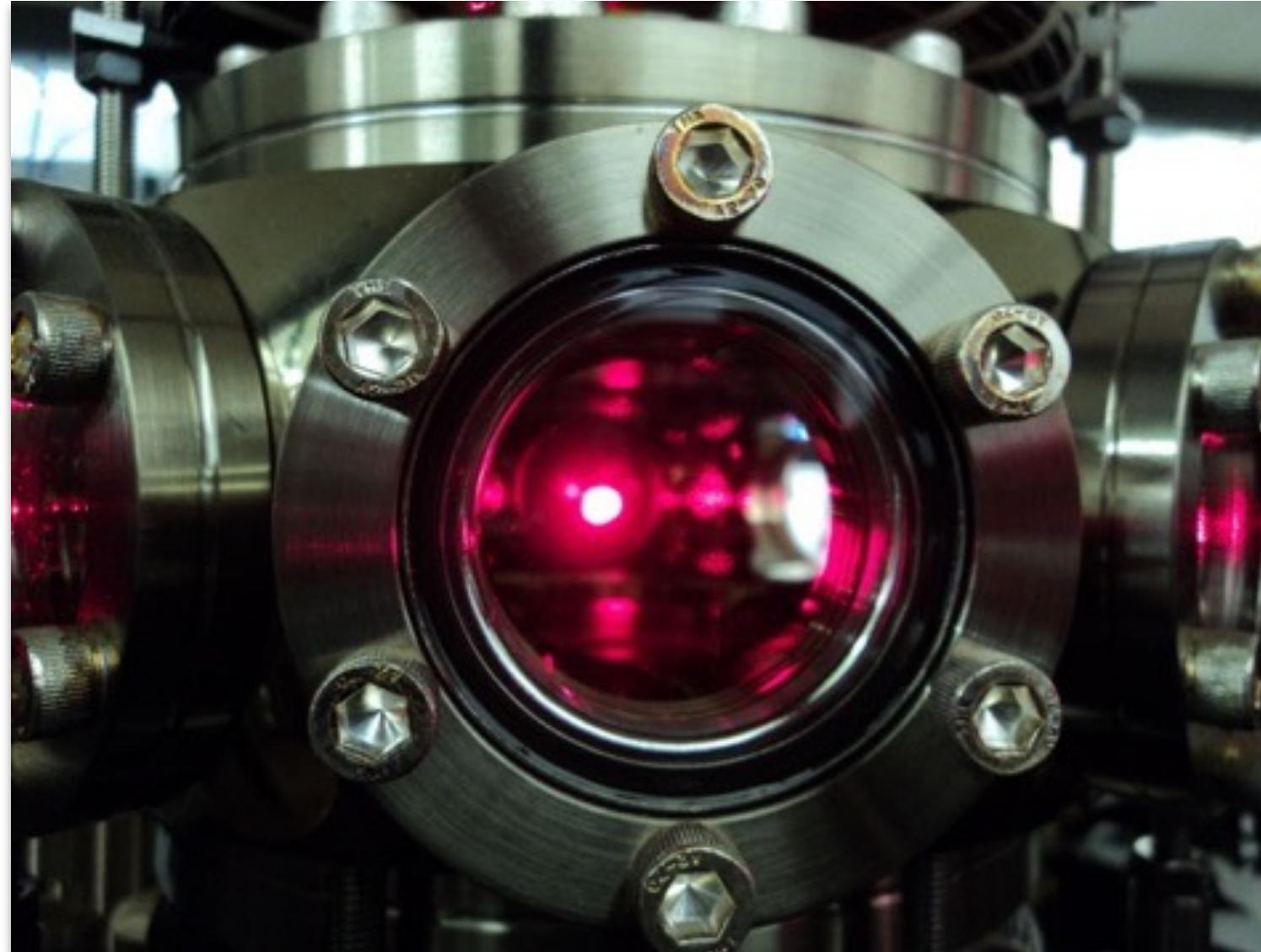
VORTEX DYNAMICS IN STRONGLY-INTERACTING ATOMIC FERMI SUPERFLUIDS: FROM FEW TO MANY

Giulia Del Pace

University of Florence



LITHIUM-6 FERMI SUPERFLUIDS

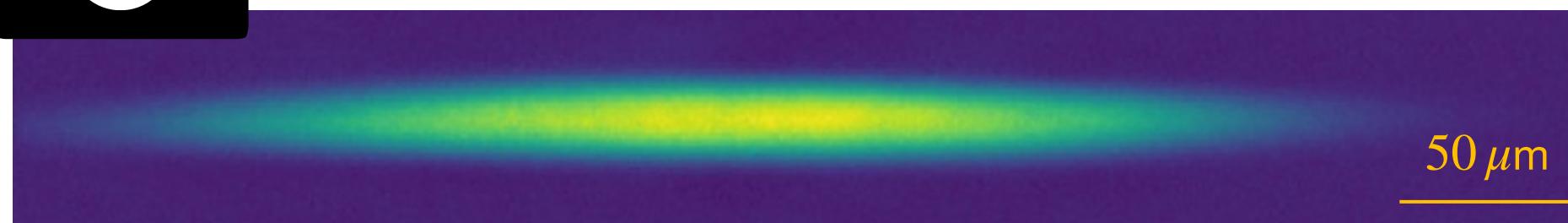
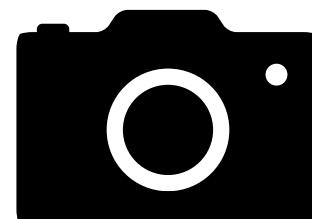


$$\sim \frac{4\pi\hbar^2}{m} a \delta(r)$$

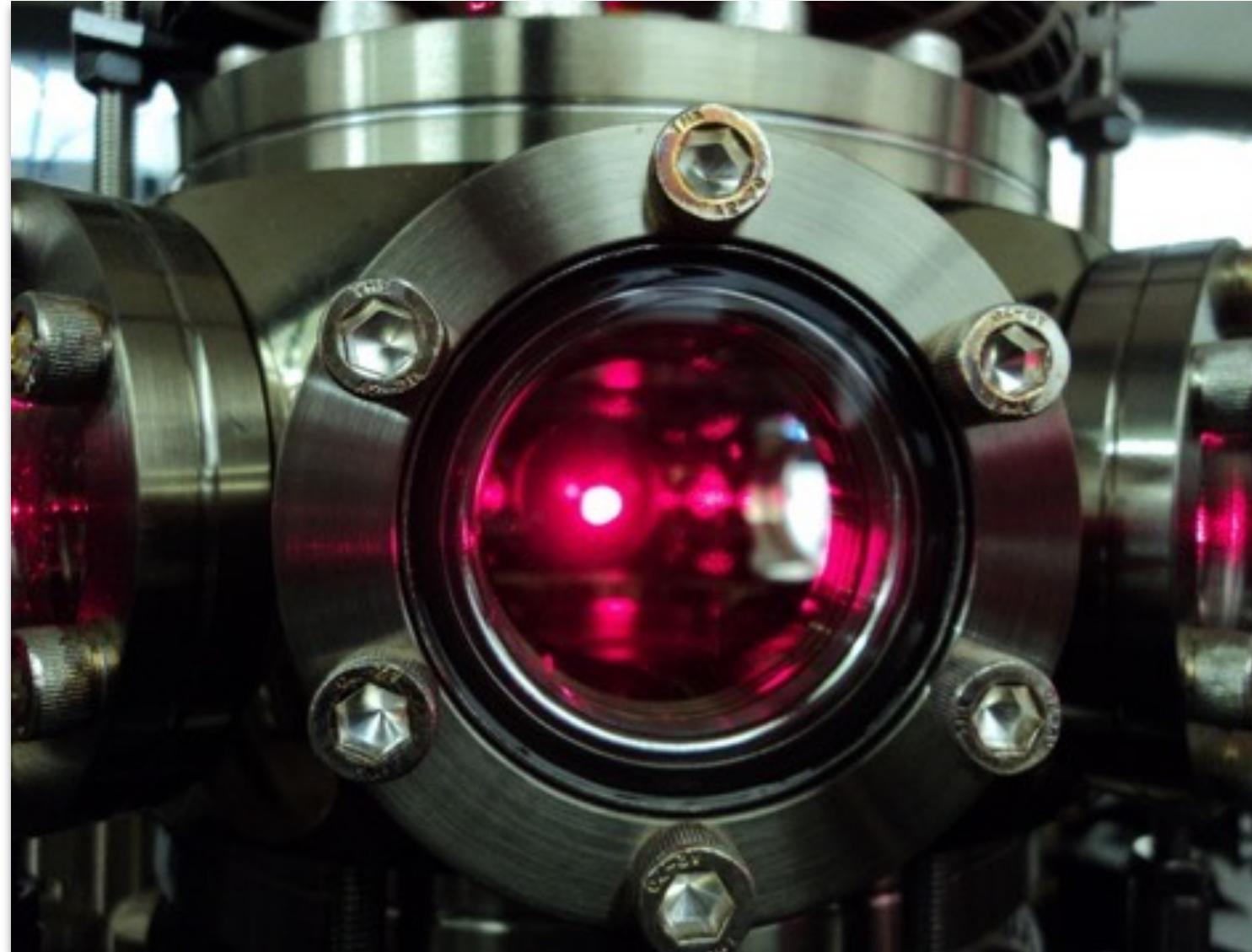
$a \rightarrow$ Scattering length

$N \sim 10^5$ at $T \sim 30$ nK

Magneto optical trap + D₁ molasses
+ Evaporative cooling

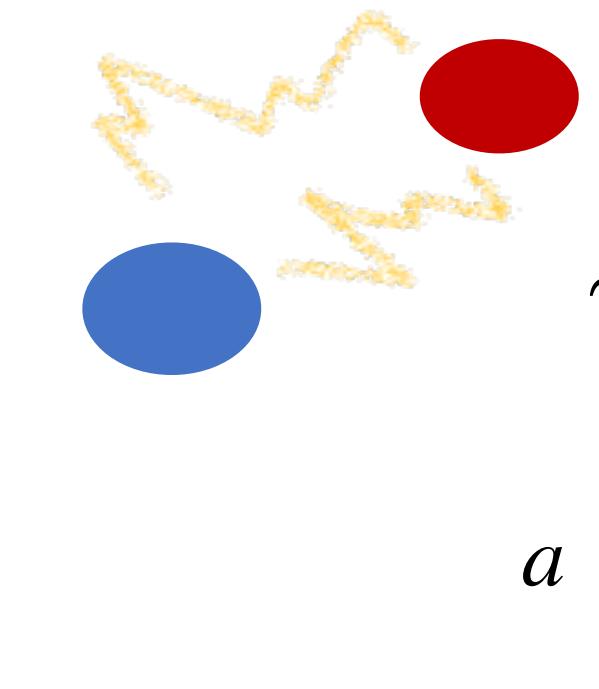
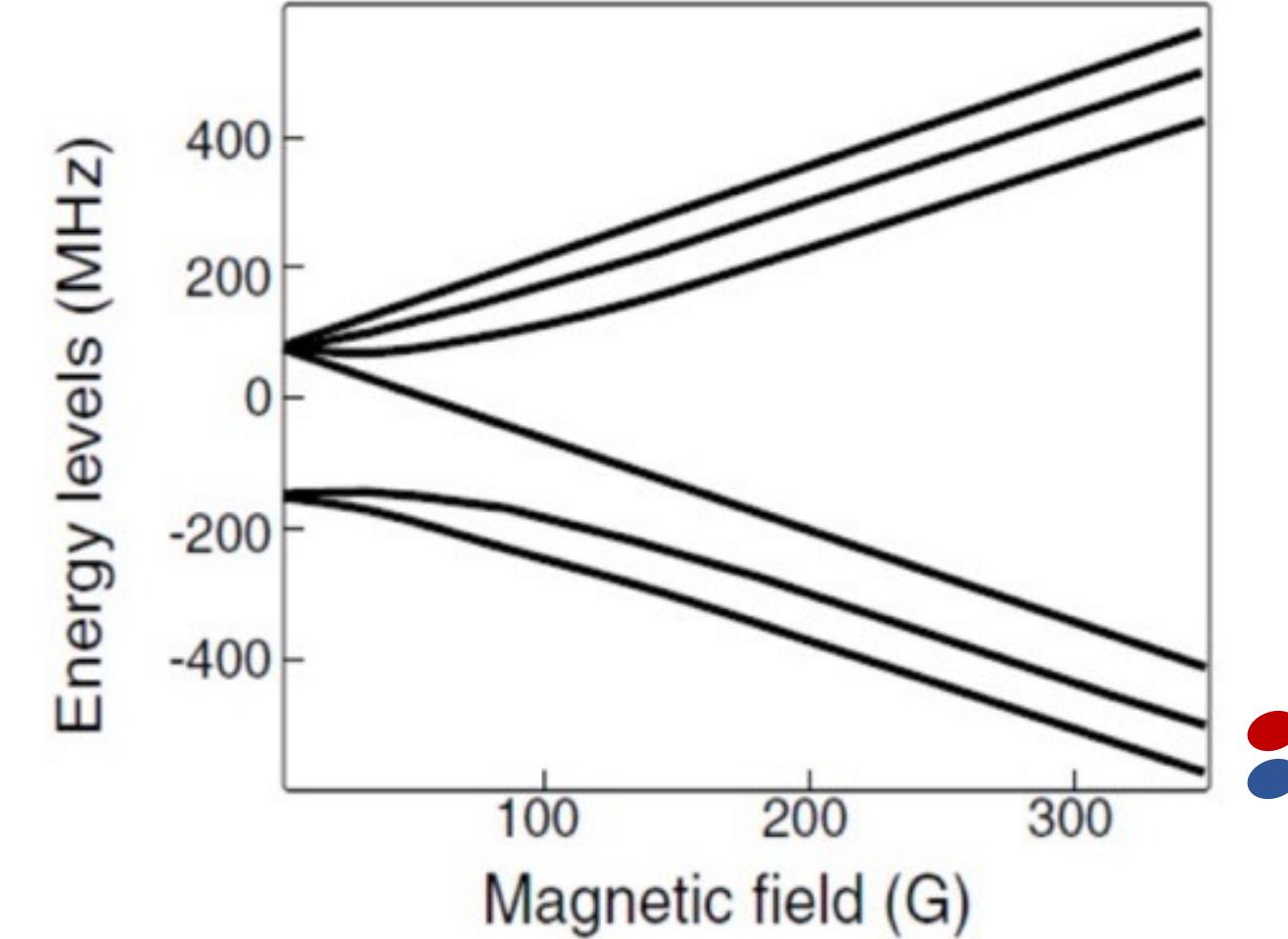
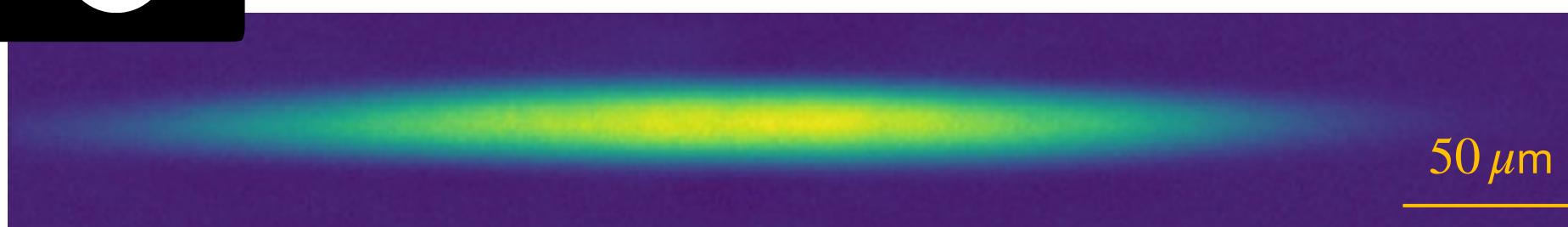
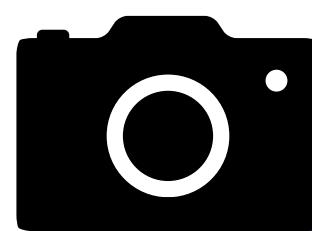


LITHIUM-6 FERMI SUPERFLUIDS



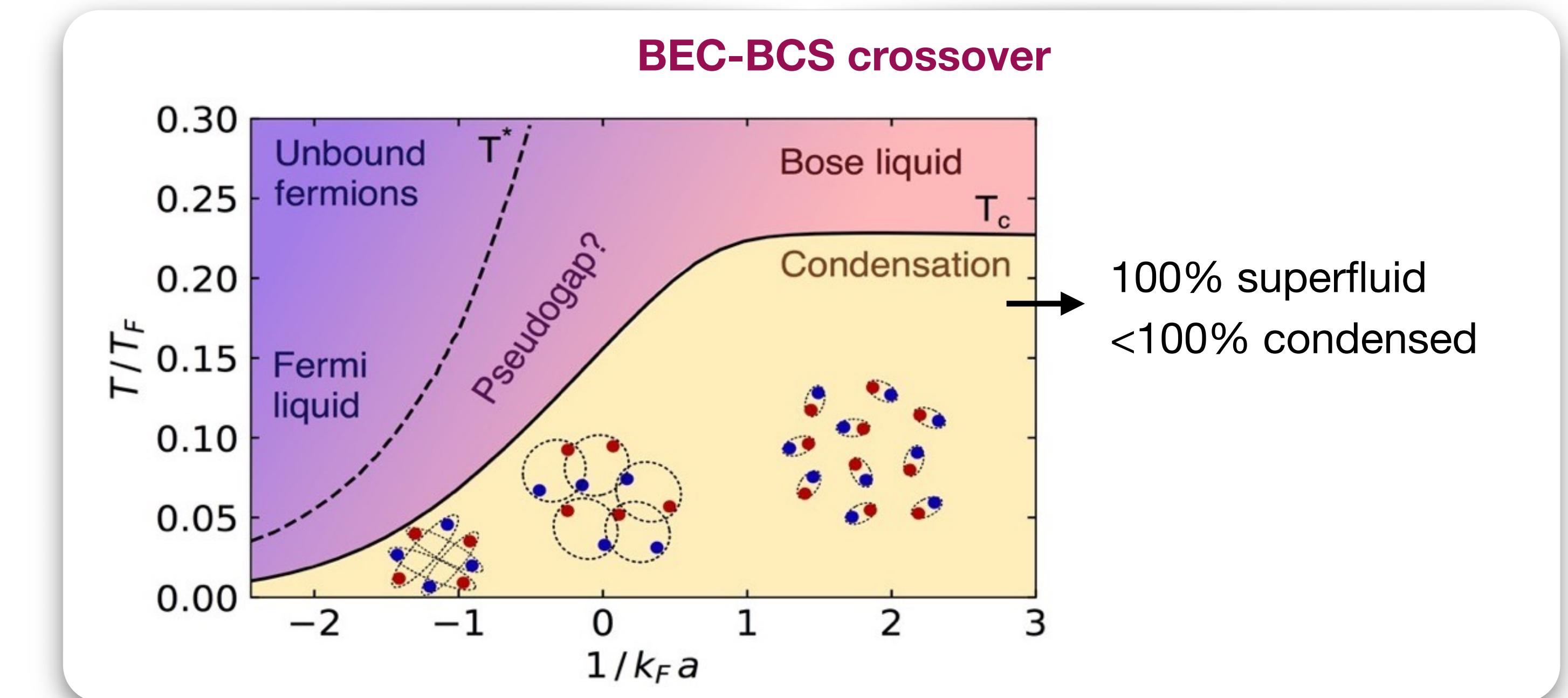
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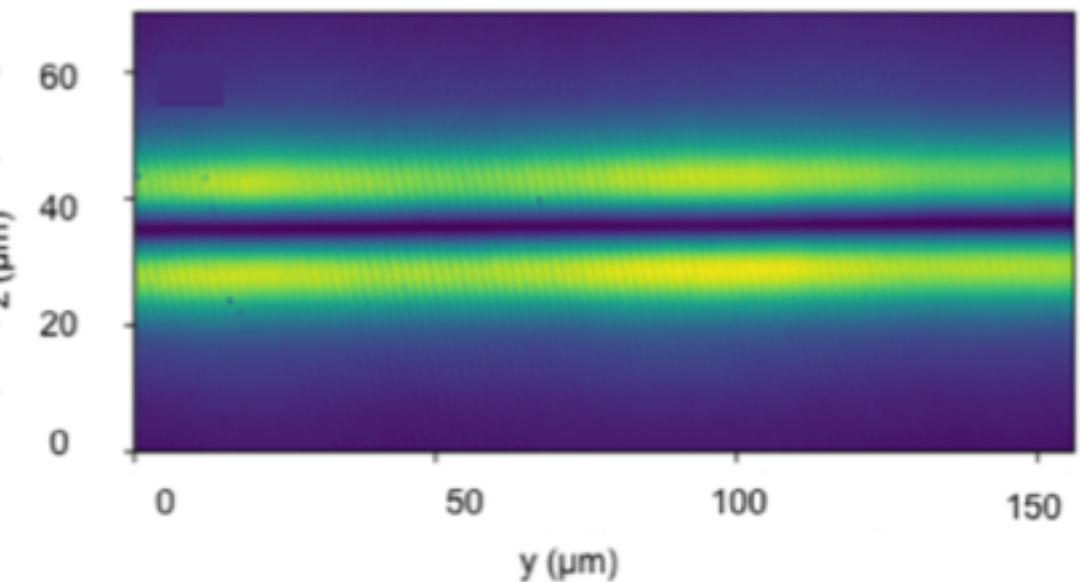


$$\sim \frac{4\pi\hbar^2}{m} a \delta(r)$$

$a \rightarrow$ Scattering length

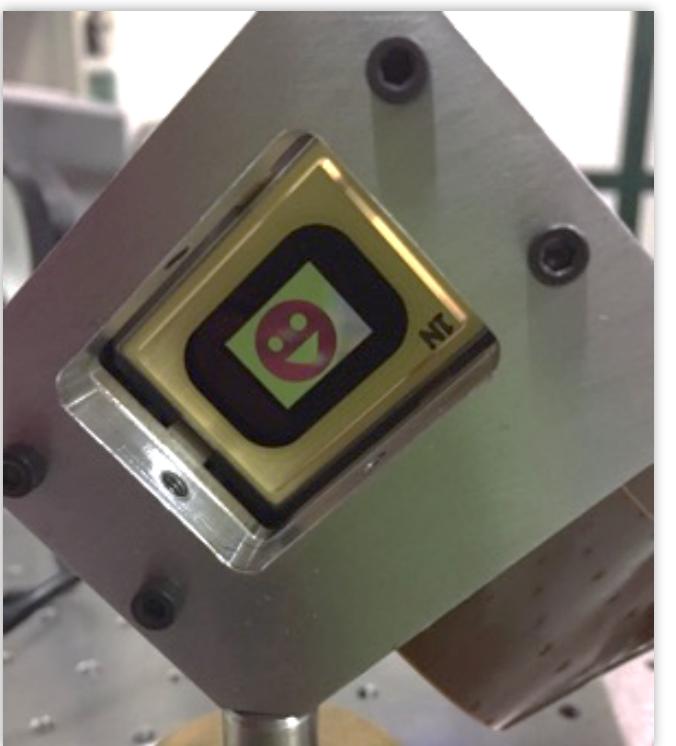


LITHIUM-6 FERMI SUPERFLUIDS @ LENS



TEM_{0,1} – like beam @ 532 nm

Squeeze the cloud along the vertical direction and create from oblate to quasi-2D geometries



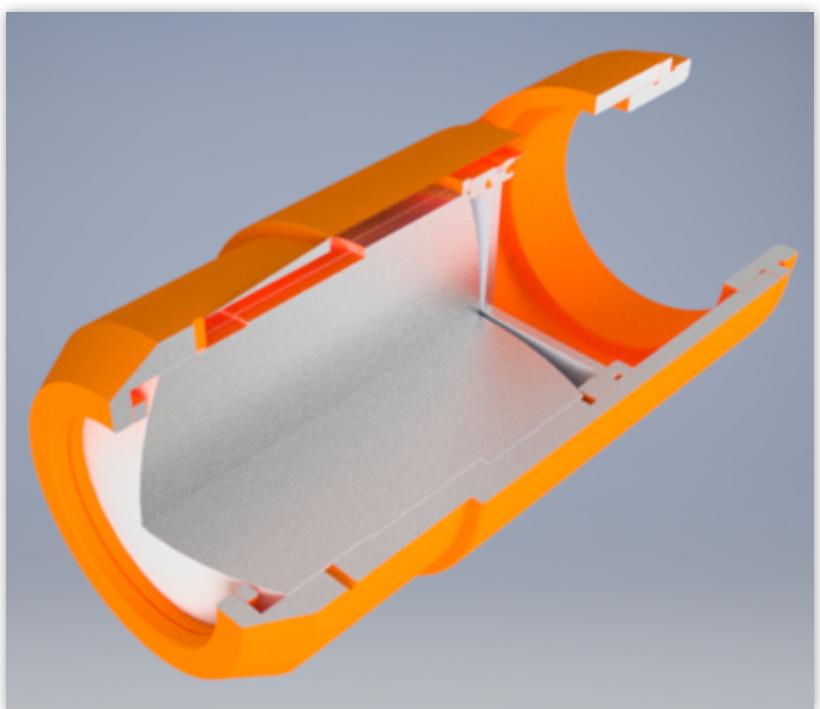
Digital Micromirror Device (DMD)

Vialux V-7000

1024 × 768 mirrors

13.68 μm pitch

Maximum framerate: 22 kHz



High resolution microscope objective:

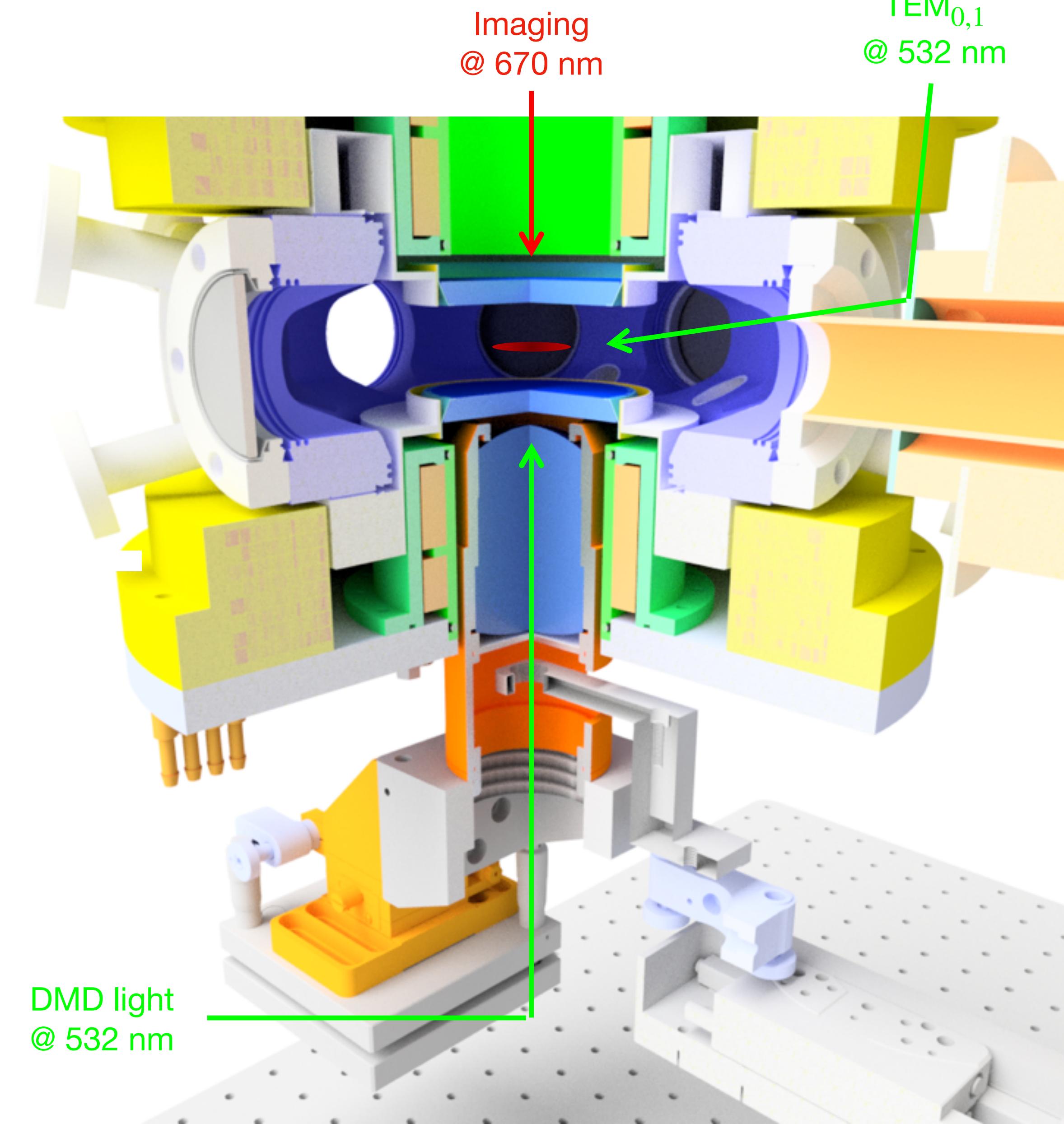
$NA = 0.45$

Effective Focal Length: 47mm

Field Of View: 0.3 mm

Working Distance: 25.1 mm

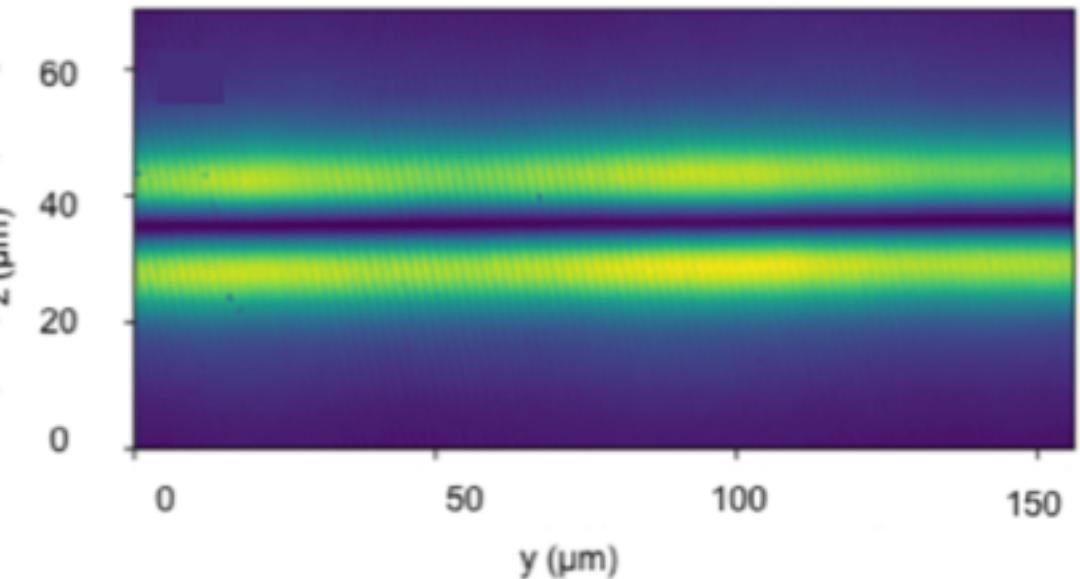
Resolution: 0.9 μm @ 670 nm, 0.7 μm @ 532 nm



TEM_{0,1}

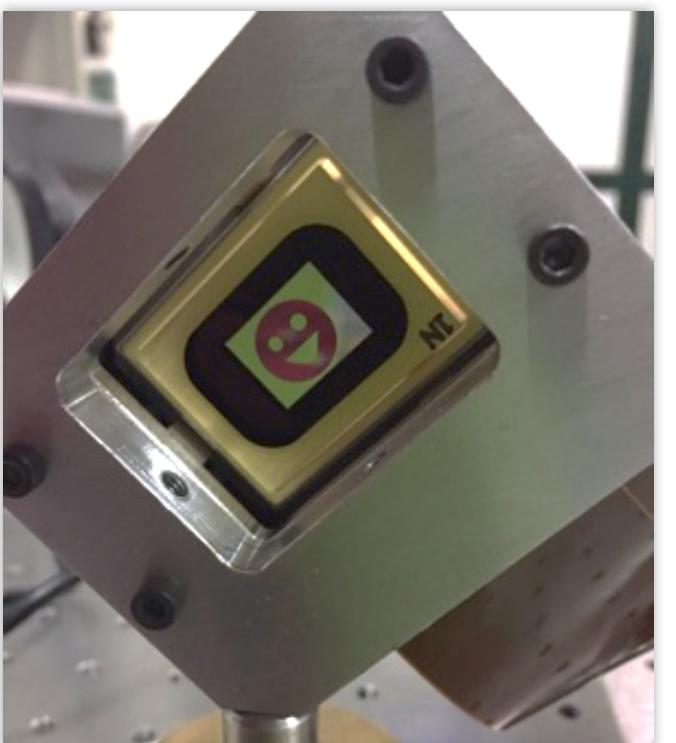
@ 532 nm

LITHIUM-6 FERMI SUPERFLUIDS @ LENS



TEM_{0,1} – like beam @ 532 nm

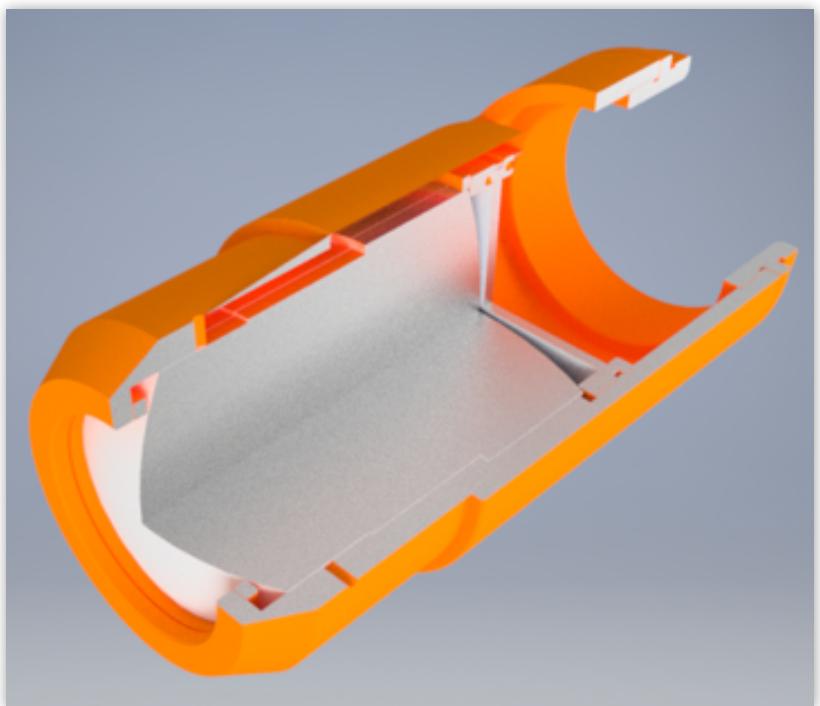
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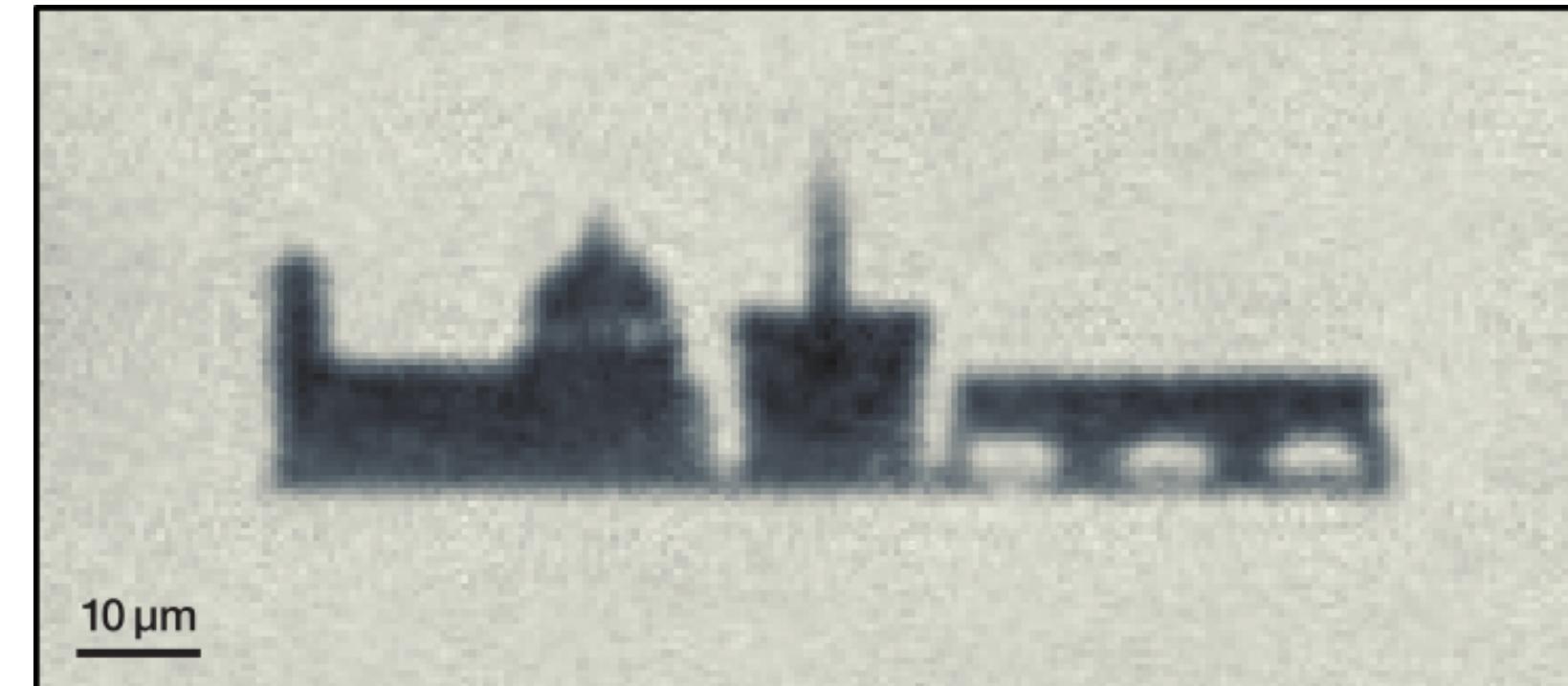
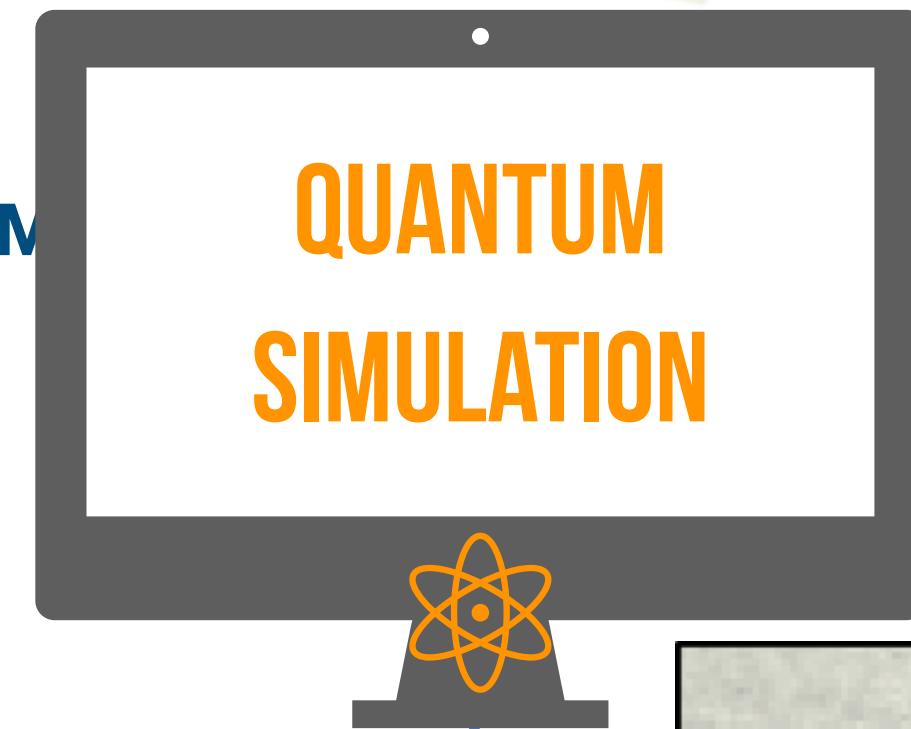
$NA = 0.45$

Effective Focal Length: 47mm

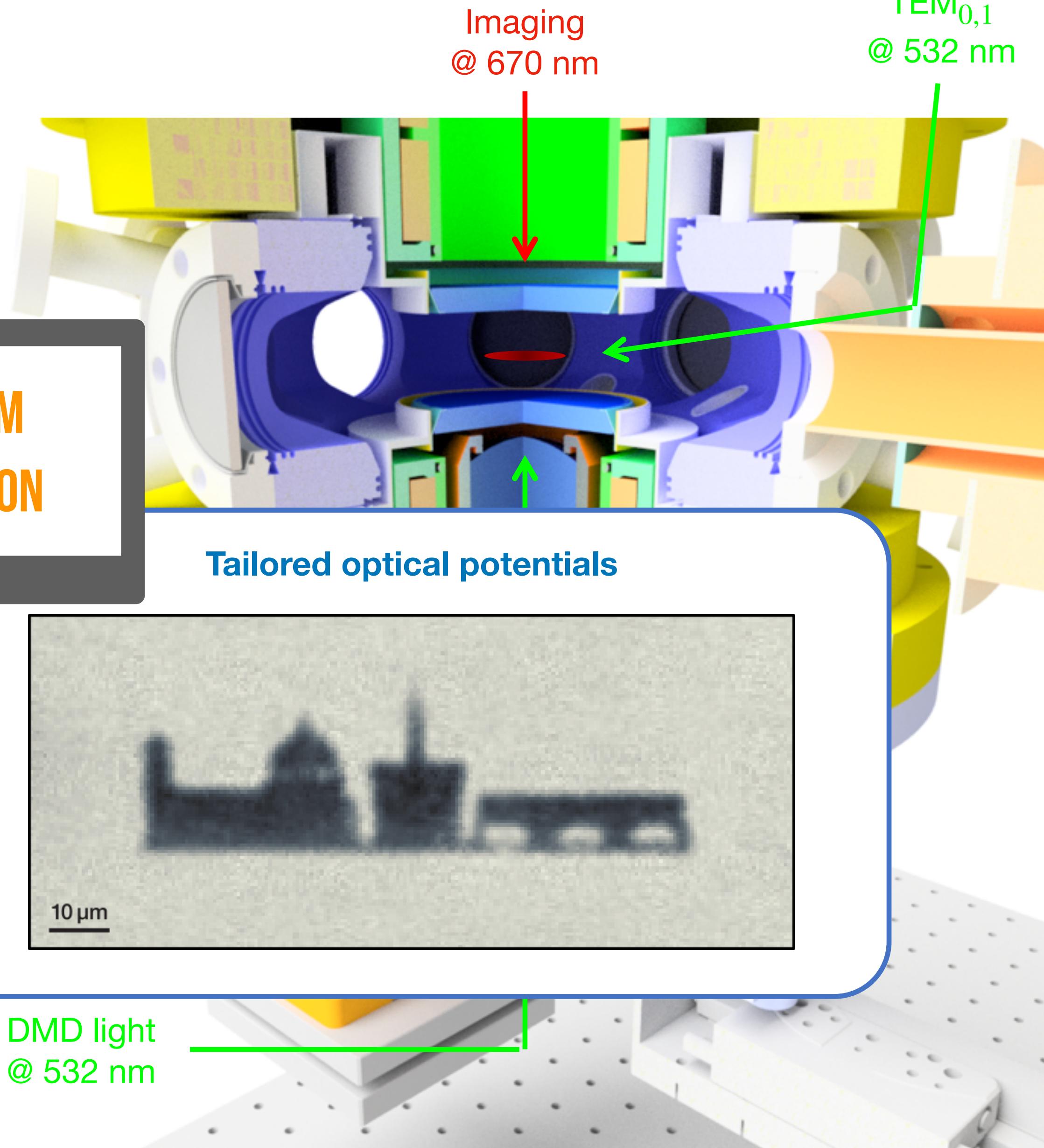
Field Of View: 0.3 mm

Working Distance: 25.1 mm

Resolution: 0.9 μm @ 670 nm, 0.7 μm @ 532 nm

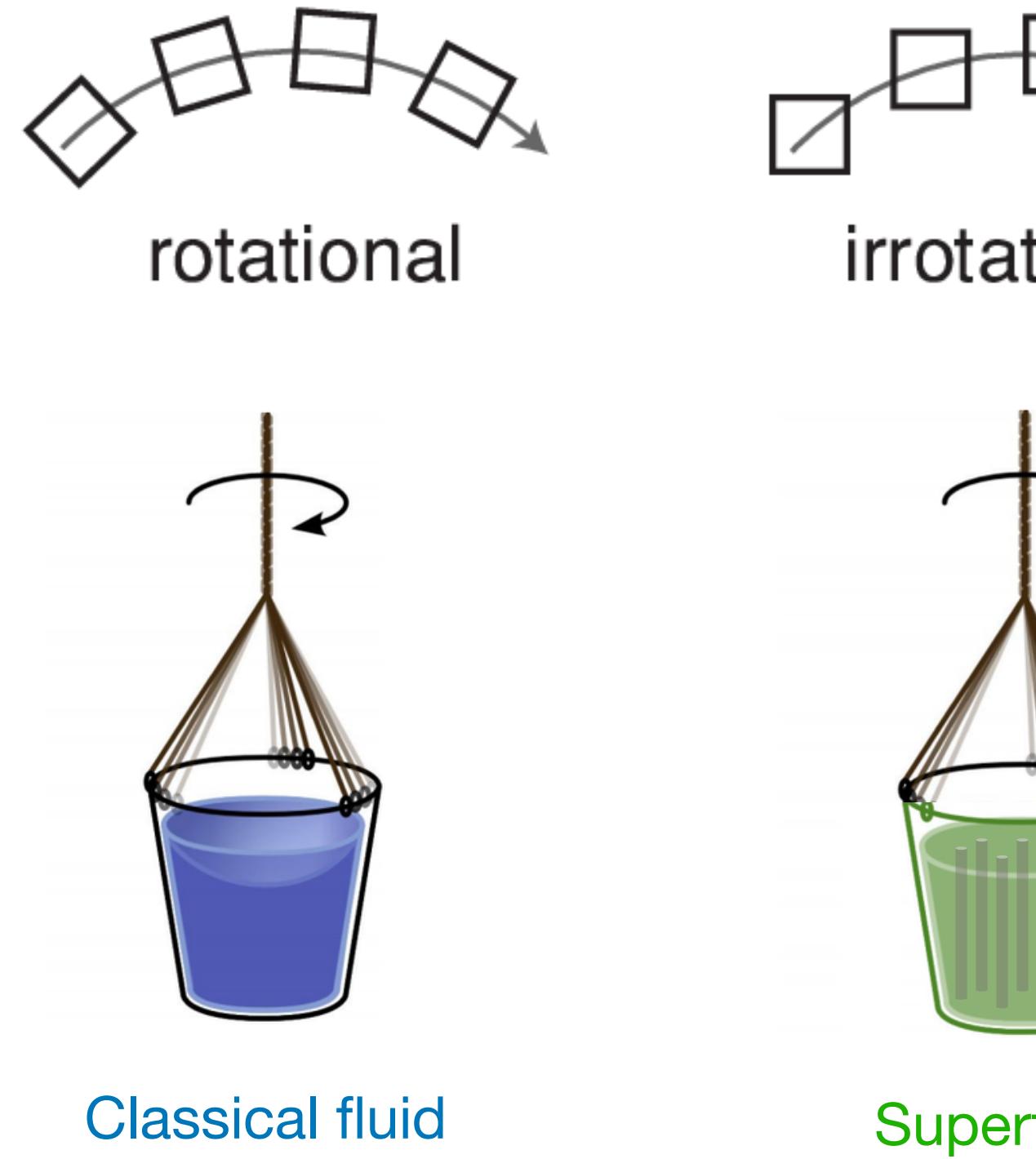


DMD light
@ 532 nm



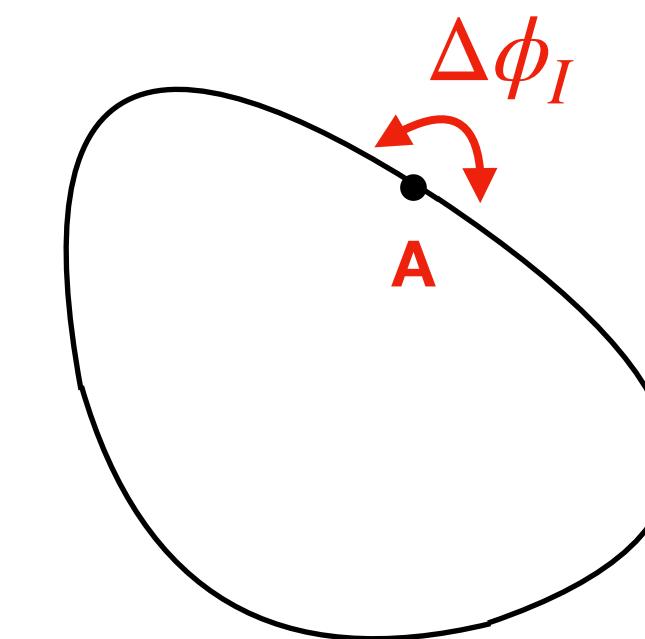
TEM_{0,1}
@ 532 nm

TOPOLOGICAL EXCITATIONS IN ATOMIC SUPERFLUIDS



The motion of an inviscid fluid is **irrotational**

$$\nabla \times v = 0$$



Quantized circulation

$$\Gamma = \oint v \cdot dl = \frac{\hbar}{m} \Delta\phi = \frac{\hbar}{m} 2\pi w$$

For a superfluid:

$$v = \frac{\hbar}{m} \nabla \phi$$

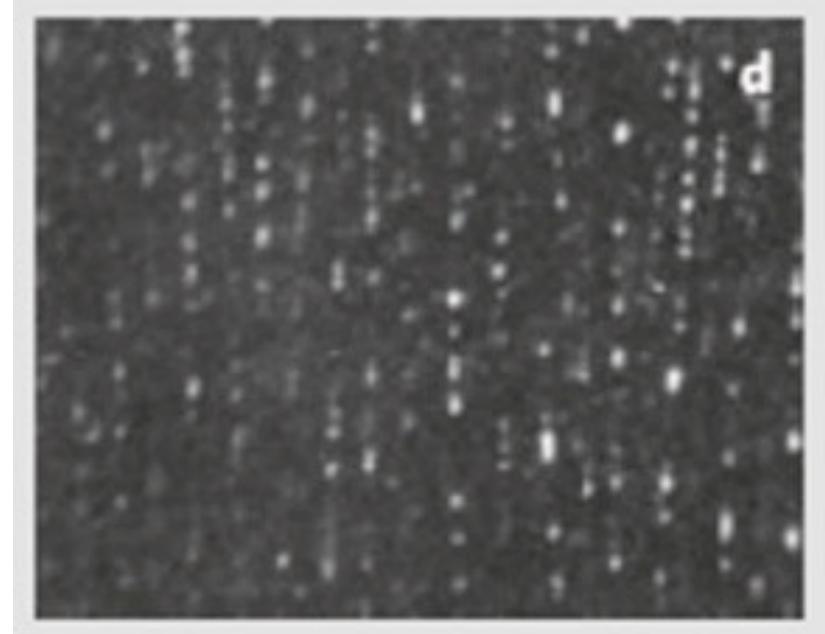
$\Delta\phi = 2\pi w \rightarrow$ single-valuedness of wavefunction

$w \rightarrow$ winding number

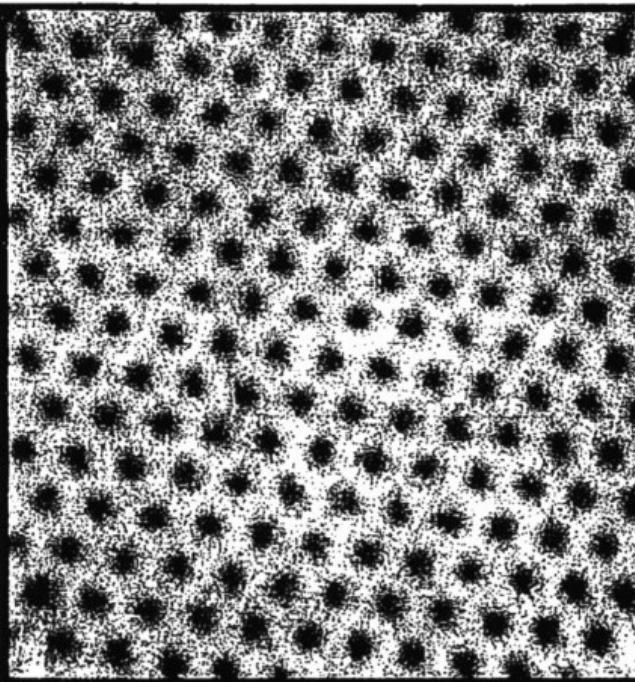
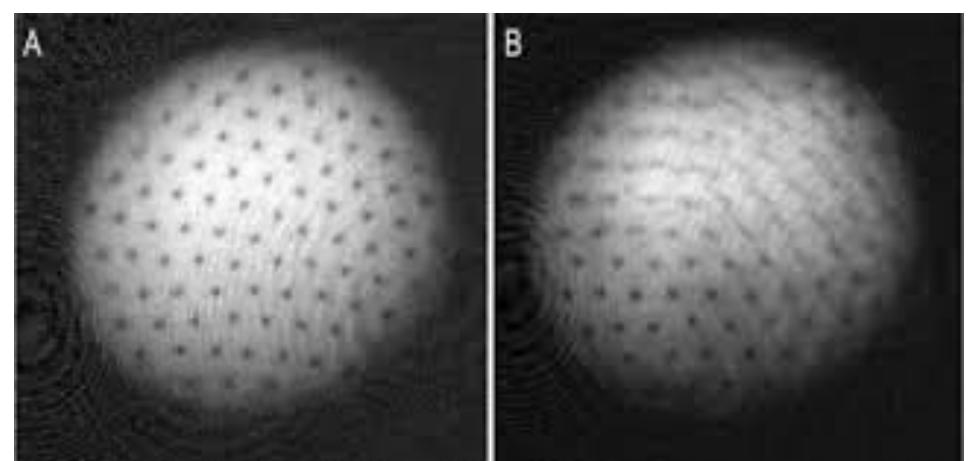
Topological excitations in a multiply-connected geometry
Vortices, (meta)stable currents in rings

TOPOLOGICAL EXCITATIONS IN ATOMIC SUPERFLUIDS

Quantized vortices in superfluids

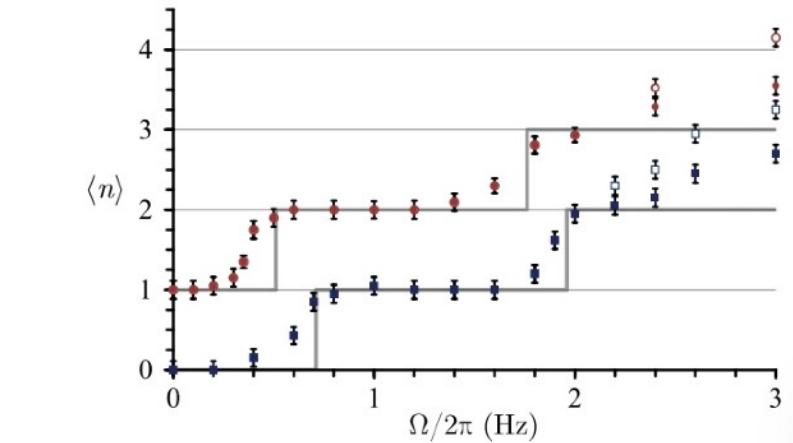
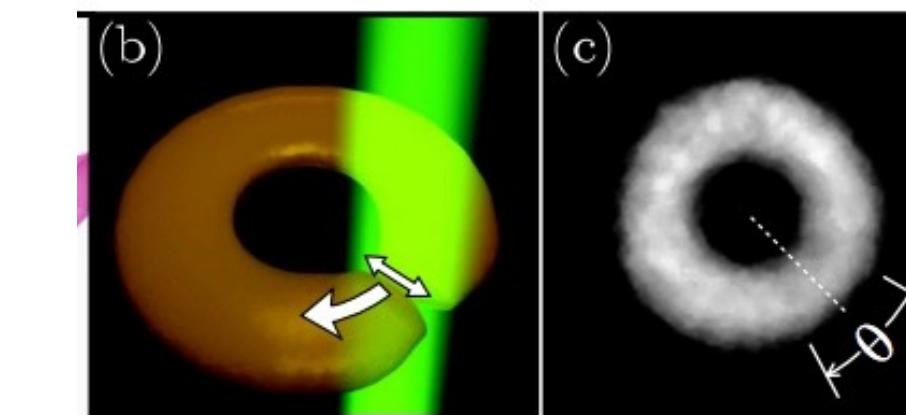
Helium

From: G. Bewley, et al.,
Nature 441, 588 (2006).

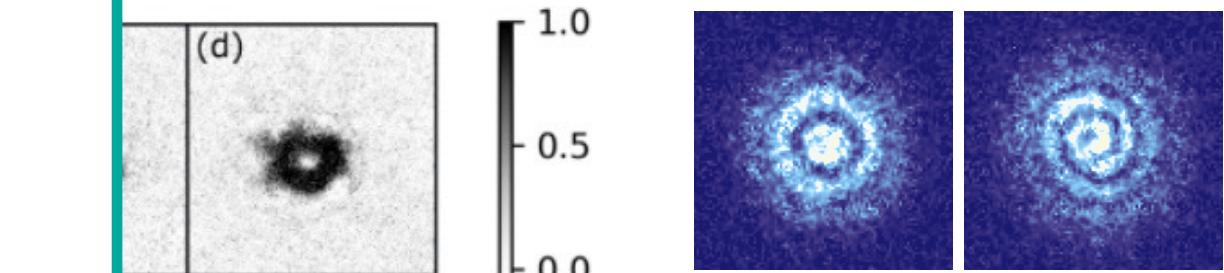
Superconductors**Atomic condensates**

From: J. R. Abo-Shaeer, et al., *Science*
292, Issue 5516, pp. 476-479 (2001)

Quantized circulation in ring geometry

BEC

From: K. Wright, et al. *Phys. Rev. Lett.* 110.2 (2013)
Campbell group (NIST)
Dalibard group (ENS)
Perrin group (Sorbonne- CNRS)
Neely group (University of Queensland)

Fermionic superfluids

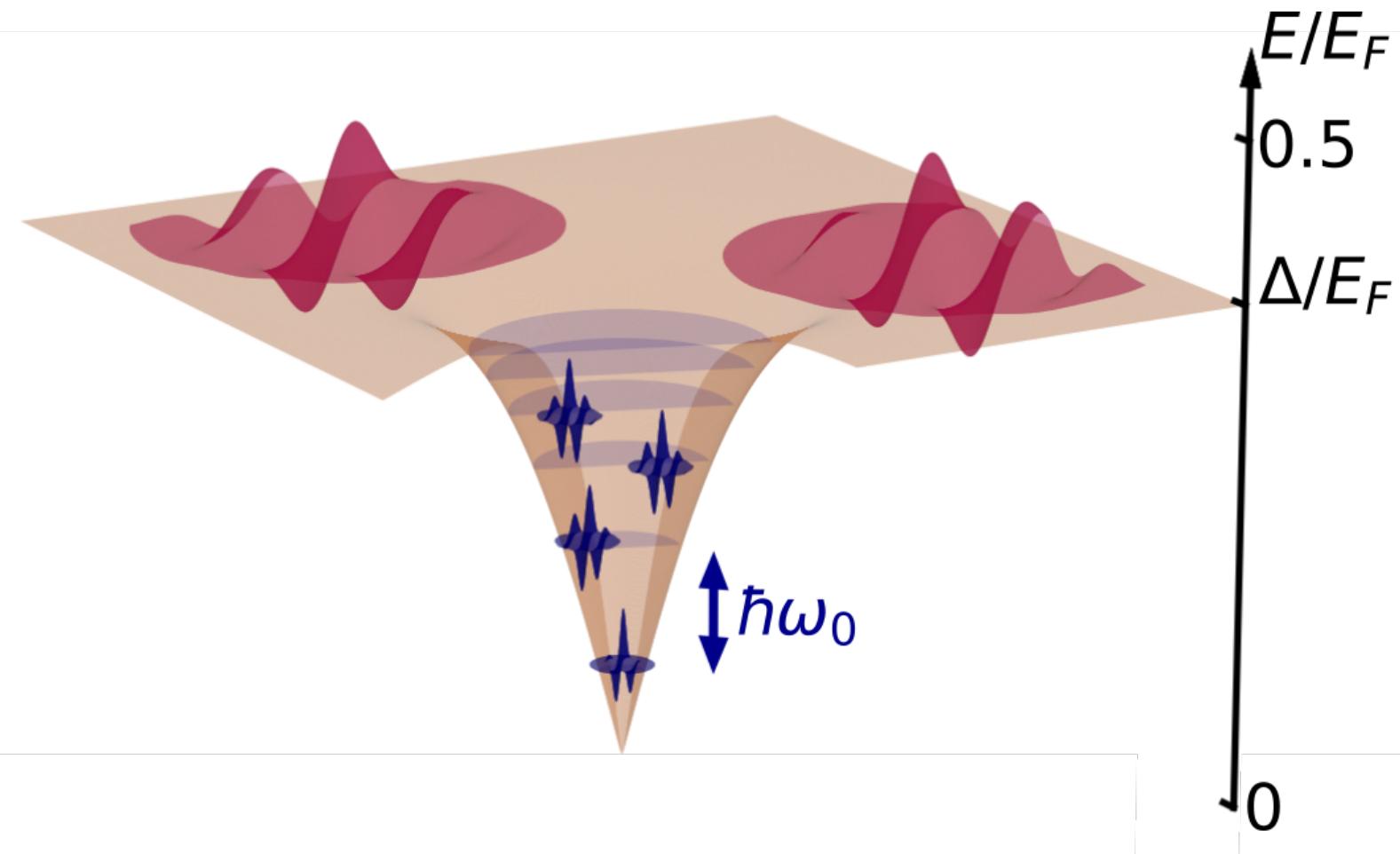
401 (2022).

GDP, et al. *Phys. Rev. X* 12.4
041037 (2022)

How can they **dissipate**?

- ▶ Fundamental physics and quantum simulation
- ▶ Quantum technologies

VORTEX DISSIPATION



Caroli, De Gennes, and Matricon. *Phys. Lett.* 9.4 307-309 (1964)
R. Sensarma, et al., *Phys. Rev. Lett.* 96.9 090403 (2006)



Mutual friction

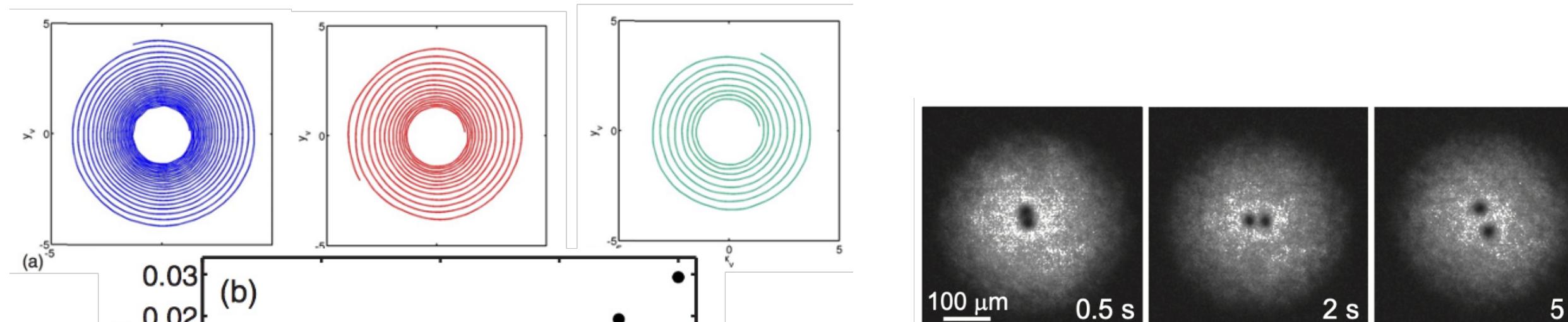
The vortex interacts with **quasiparticle** excitations of the superfluid:

- * Delocalized quasiparticles (sound) \rightarrow bosons, fermions
- * Localized quasiparticle (bound states) \rightarrow fermions

Finite temperature effect

Atomic BEC

Mutual friction determines the motion of single and many vortices

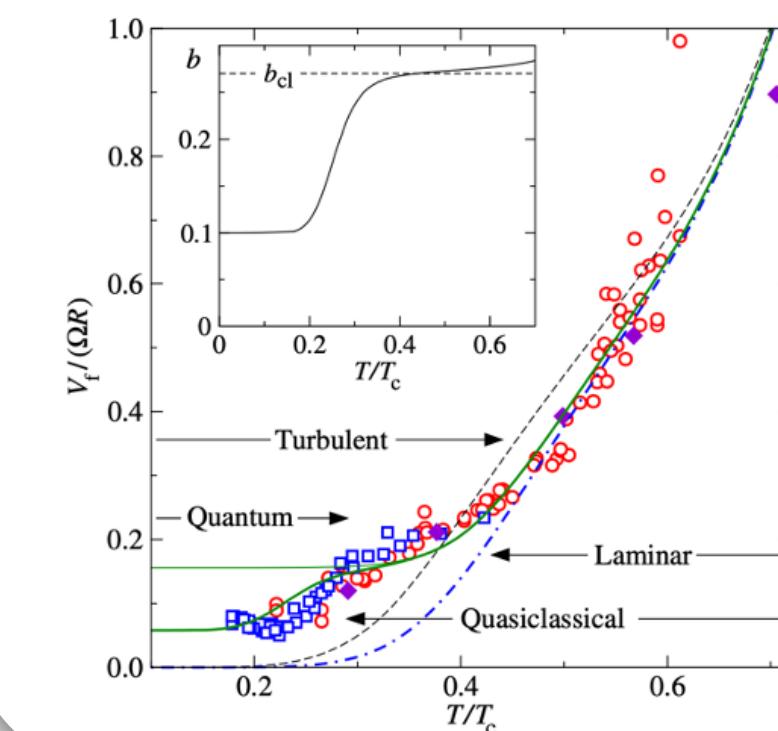


Moon et al., *PRA* 92, 051601(R) (2015)

B. Jackson, et al., *Phys. Rev. A* 79, 053615 (2009)

Superfluid Helium

Mutual friction is introduced as a phenomenological force to account for dissipation

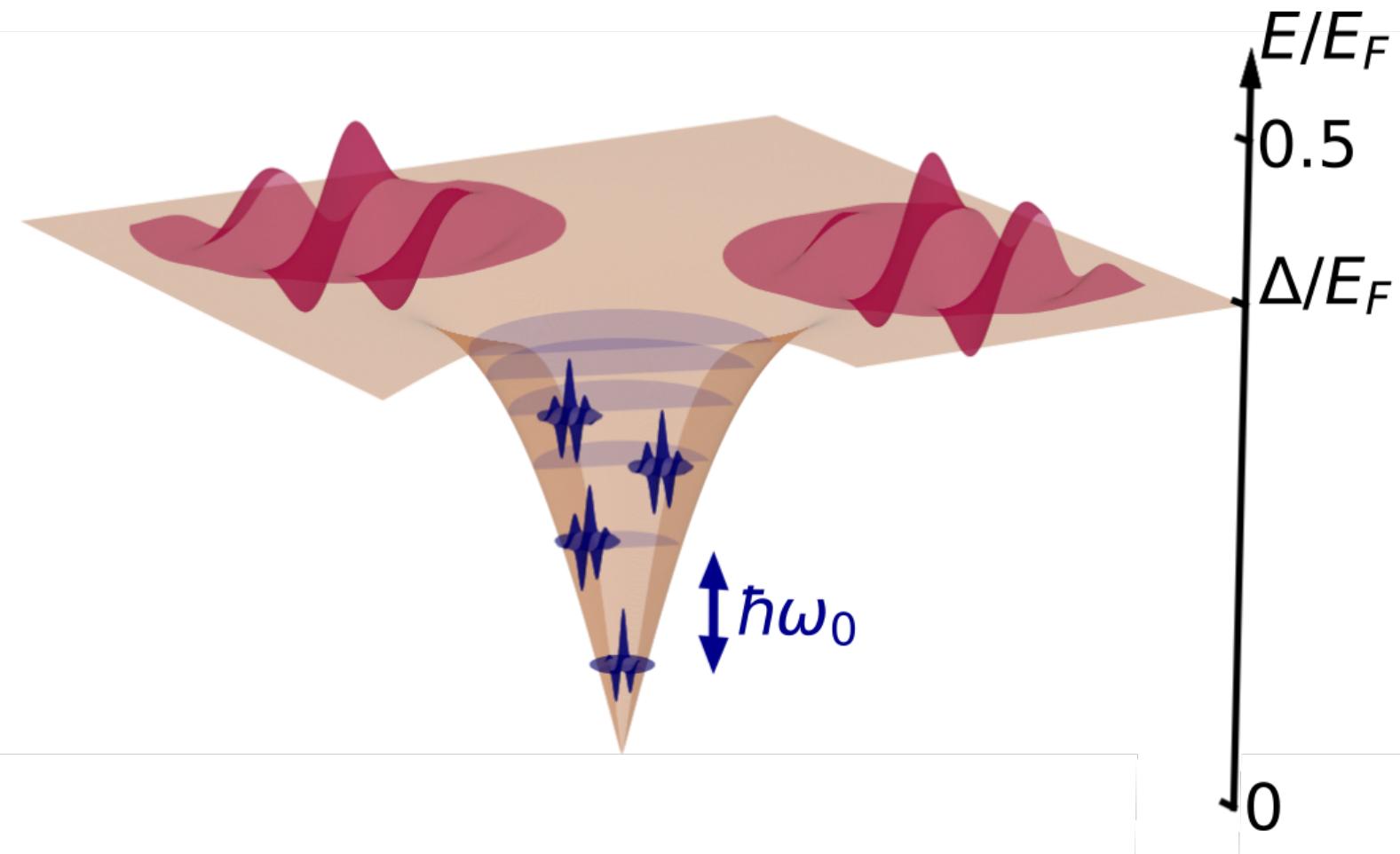


H. E. Hall and W. F. Vinen, *The rotation of liquid helium II – II. The theory of mutual friction in uniformly rotating helium II*, Proc. R. Soc. Lond. 238, 215 (1956).

P. M. Walmsley et al., PRL 99, 265302 (2007)
J. T. Mäkinen, and V. B. Eltsov. *Phys. Rev. B* 97.1 014527 (2018)

From: V. B. Eltsov et al., PRL 99, 265301 (2007)

VORTEX DISSIPATION



Sound interaction

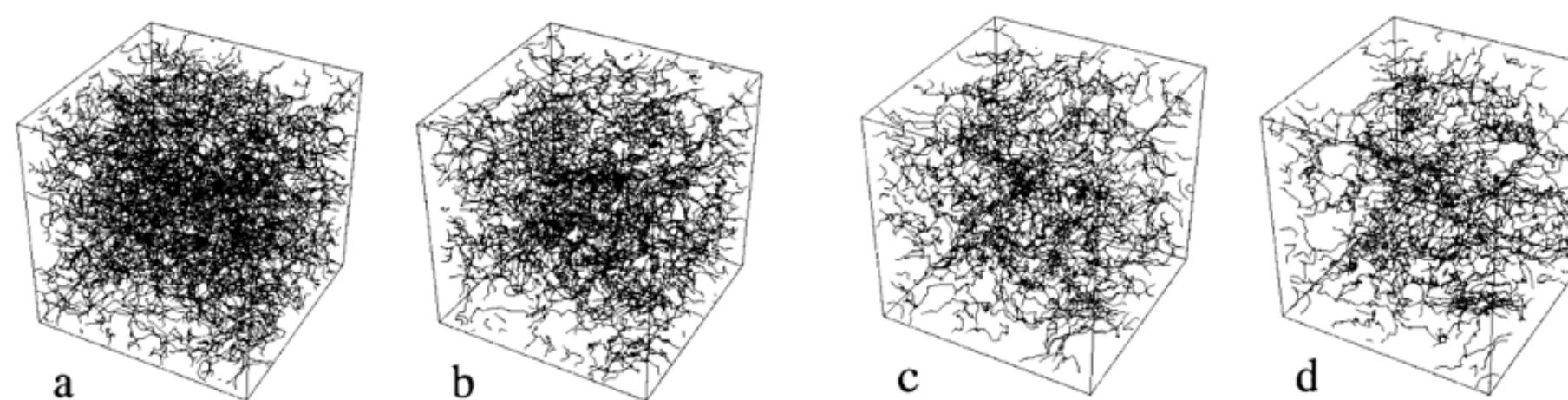
The **incompressible** energy of the vortex can be transformed into **compressible** energy of the superfluid: sound emission

Also at zero temperature

Caroli, De Gennes, and Matricon. *Phys. Lett.* 9.4 307-309 (1964)
R. Sensarma, et al., *Phys. Rev. Lett.* 96.9 090403 (2006)

3D

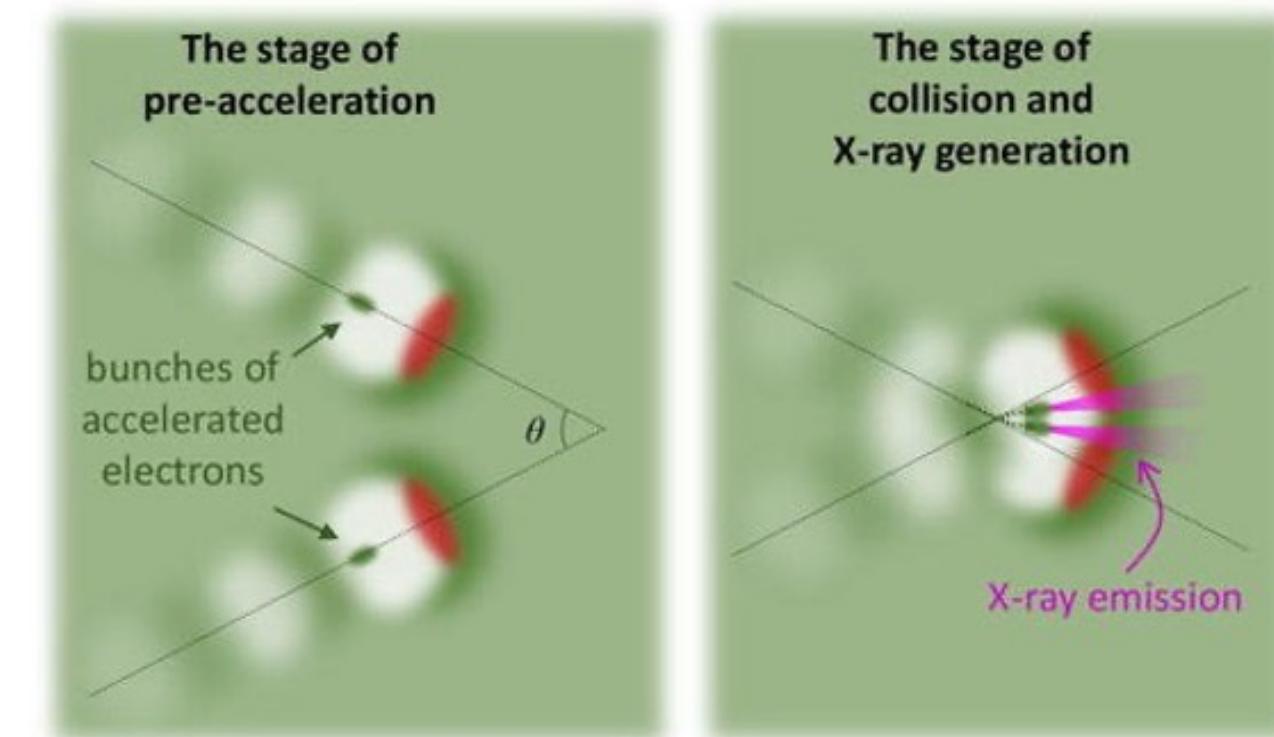
Vortex reconnection, Kelvin waves



Stagg et al., PRA 94, 053632 (2016)
Y. Minowa, et al., Nat. Phys. 21, 233–238 (2025).
M. Tsubota et al., PRB 62, 11751 (2000)

2D

Accelerating vortices \longleftrightarrow accelerating charges?



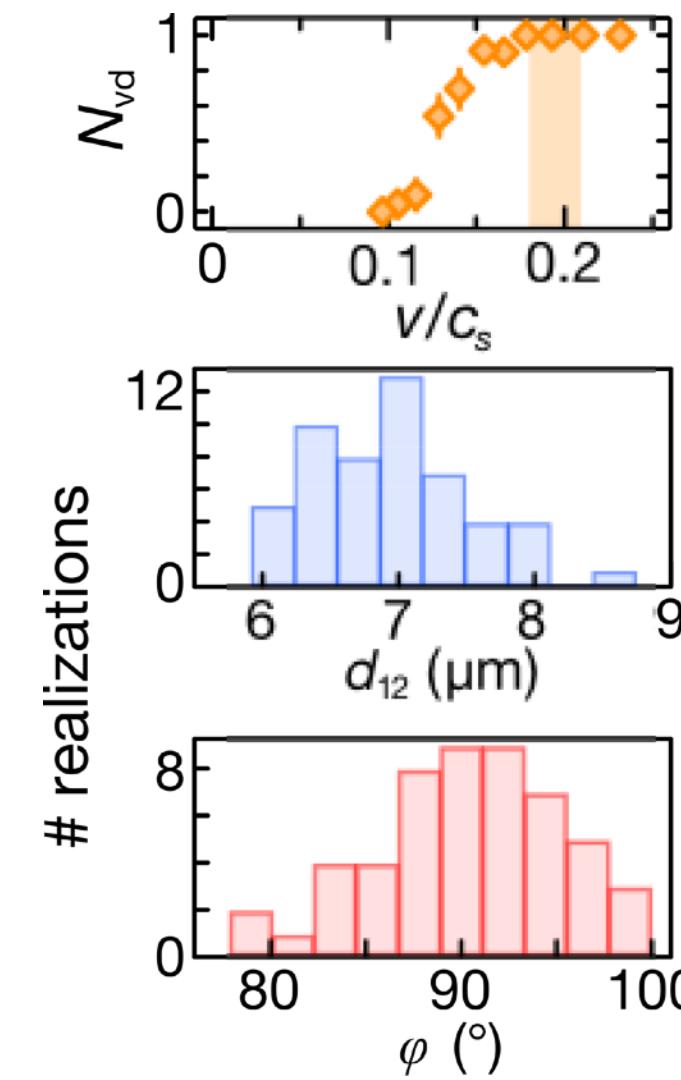
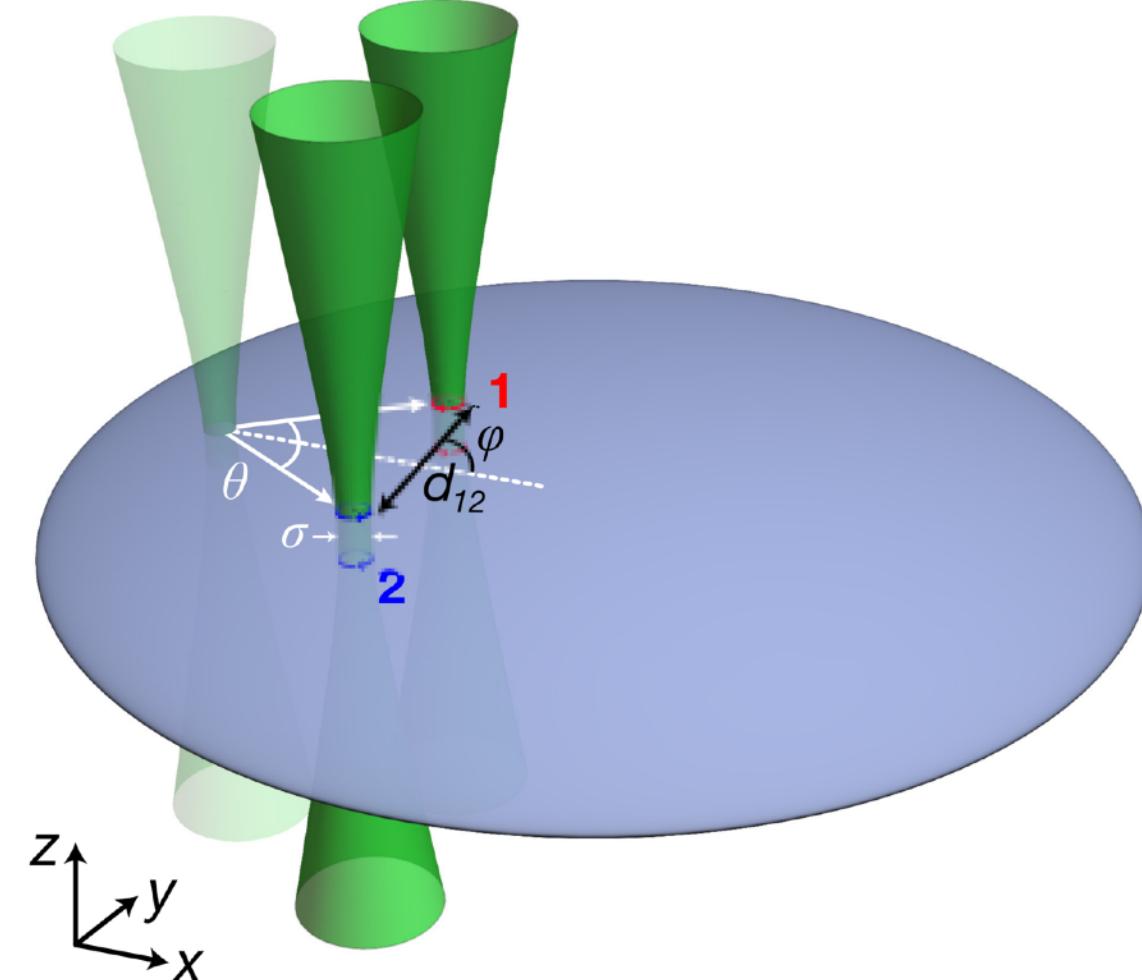
Picture taken from: E. Vallin et al., *Physics of Plasmas* 24, 093101 (2017)

VORTICES IN ATOMIC FERMI SUPERFLUIDS

W. J. Kwon, et al. *Nature* 600, 64–69 (2021).

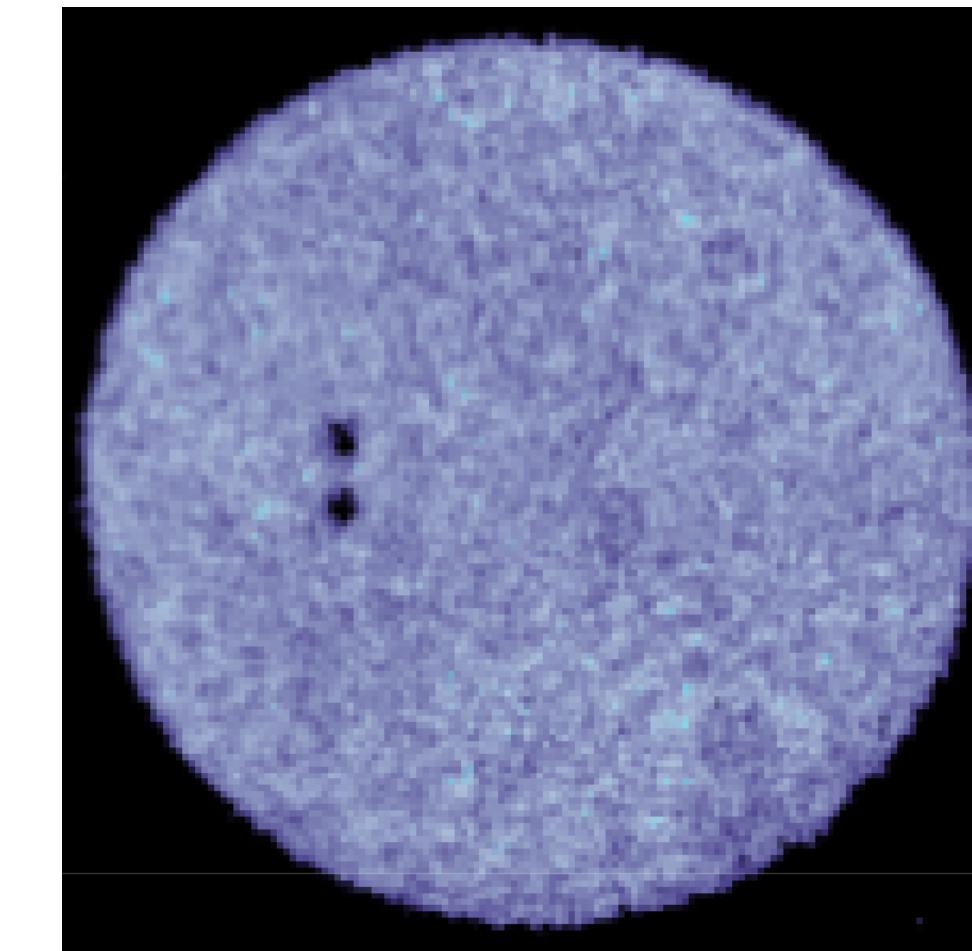
Excitation of vortices

Chopstick method

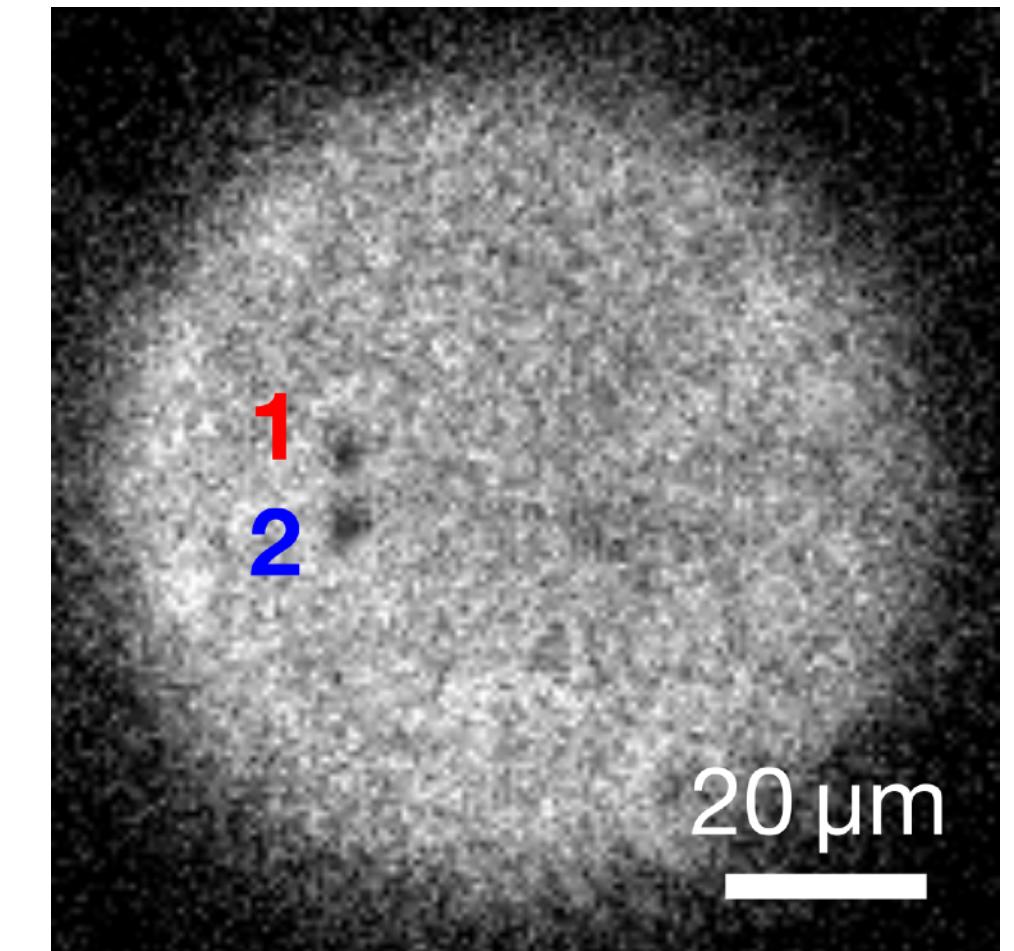


Detection of vortices

Chopstick ON

In situ

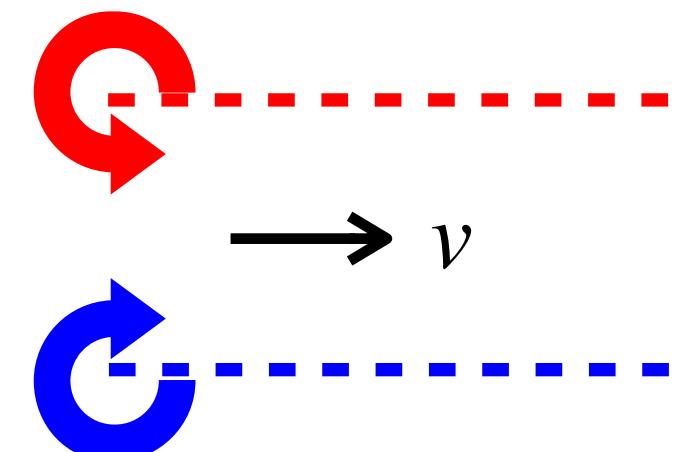
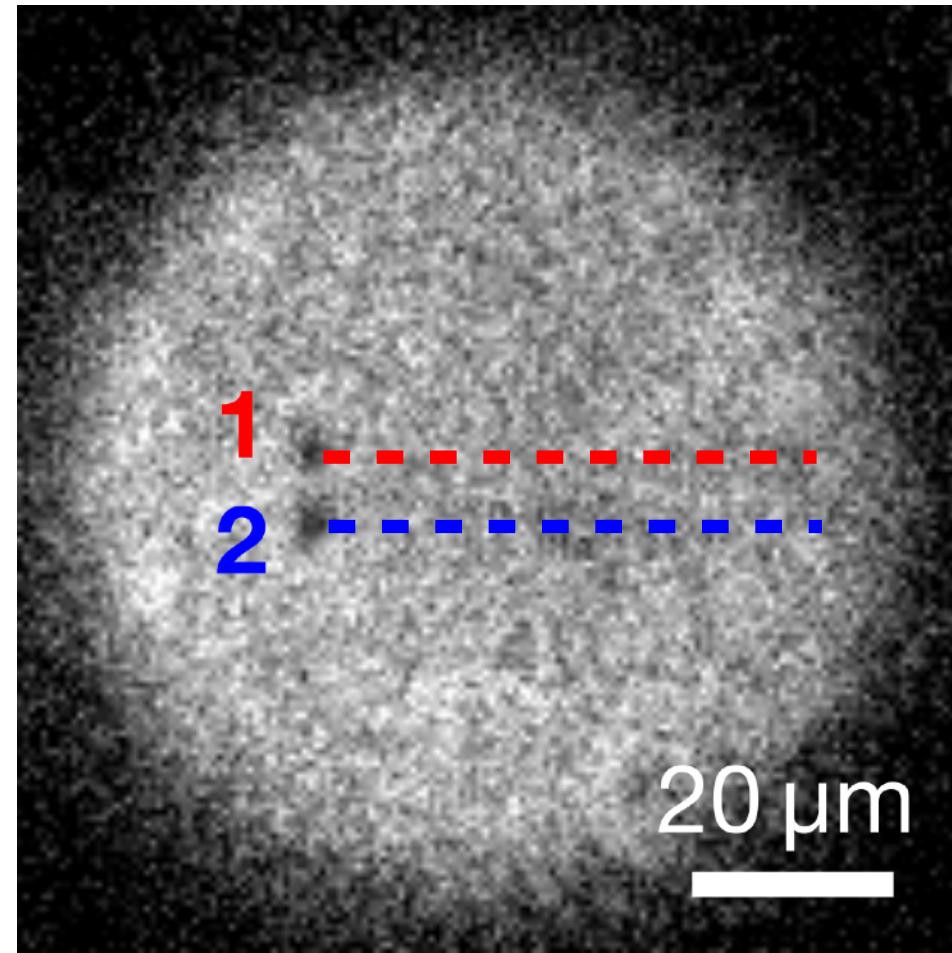
Chopstick OFF

Time-of-flightWith Bose superfluids:E. C. Samson, *Phys. Rev. A* 93, 023603 (2016).
T. Neely et al., arXiv:2402.09920v2

- ▶ Create arbitrary configuration of **vortex dipole**
- ▶ Precisely track the vortex position

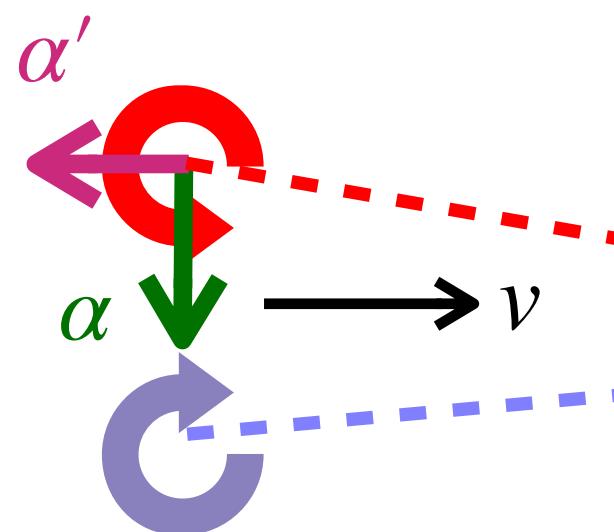
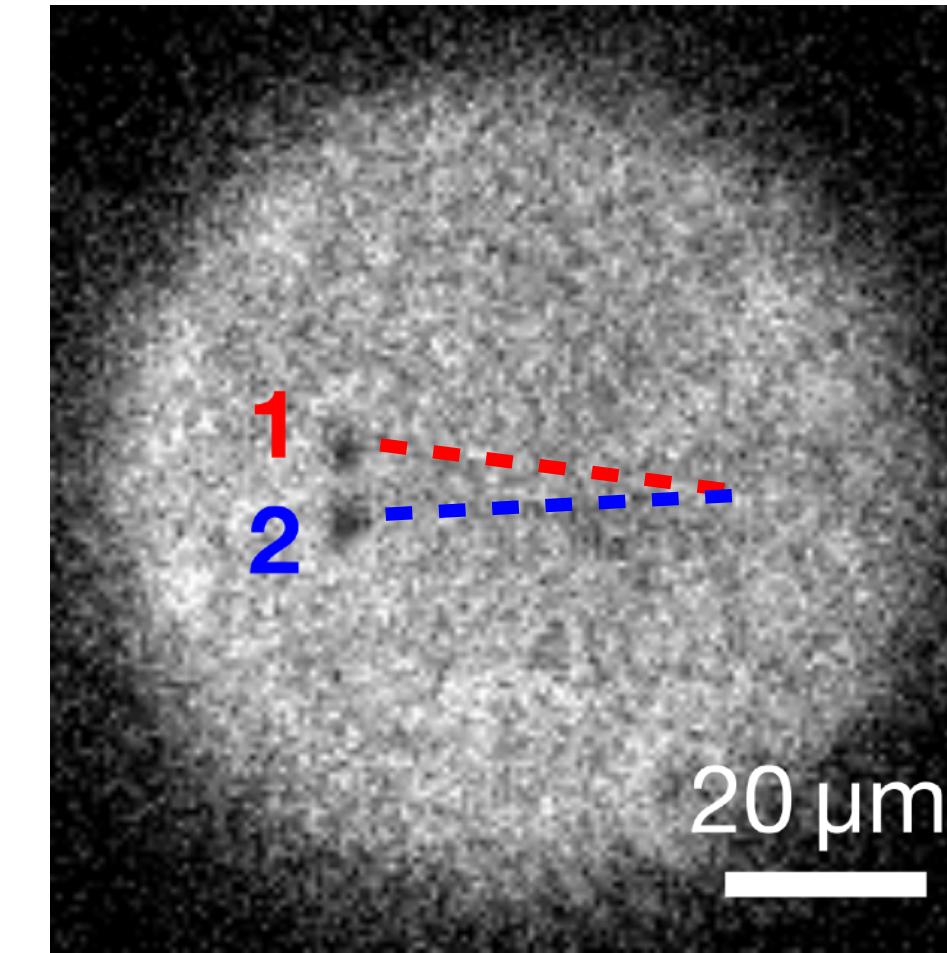
SINGLE DIPOLE DYNAMIC

Without dissipation



$$E_i = E_f$$

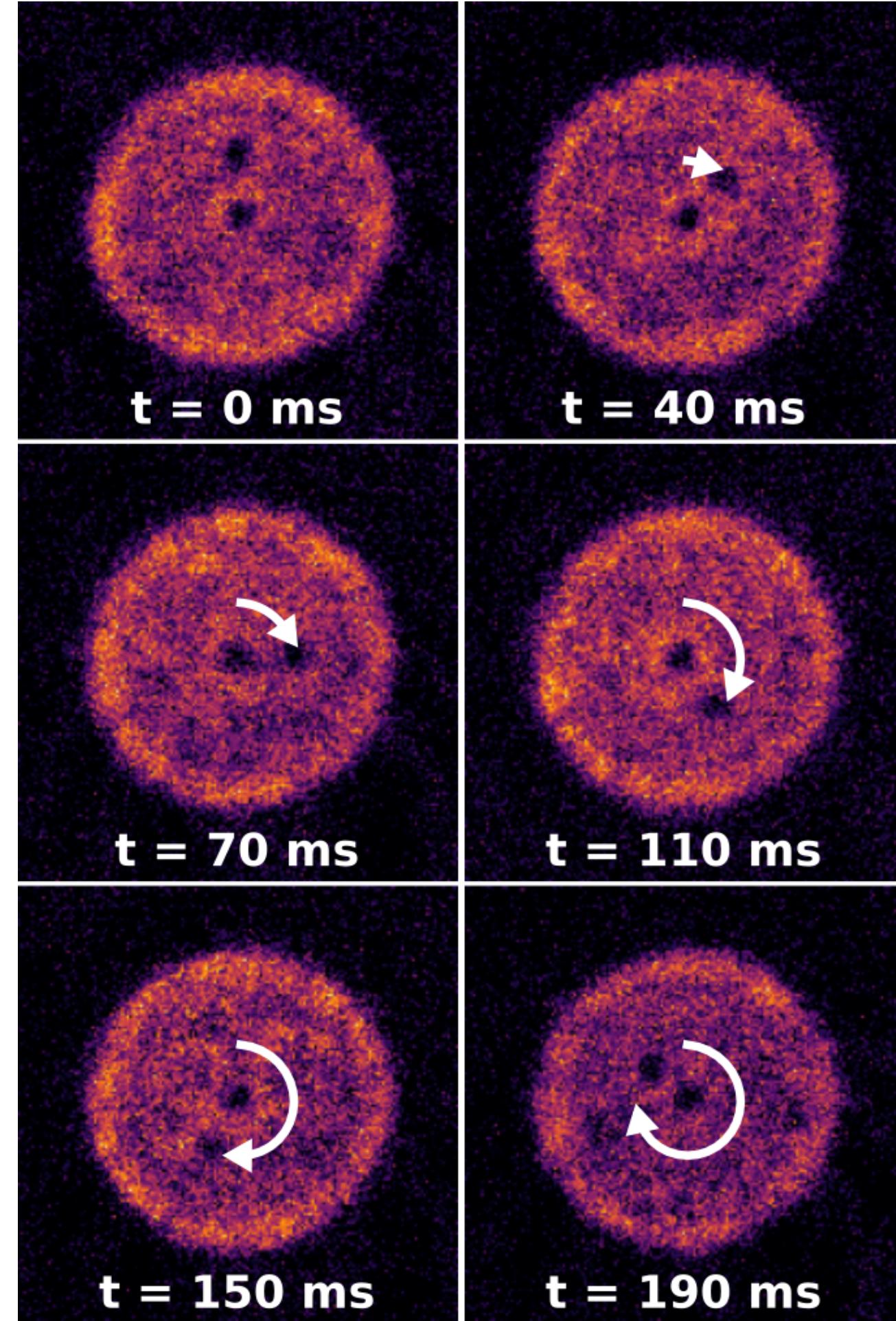
With dissipation



$$E_i > E_f$$

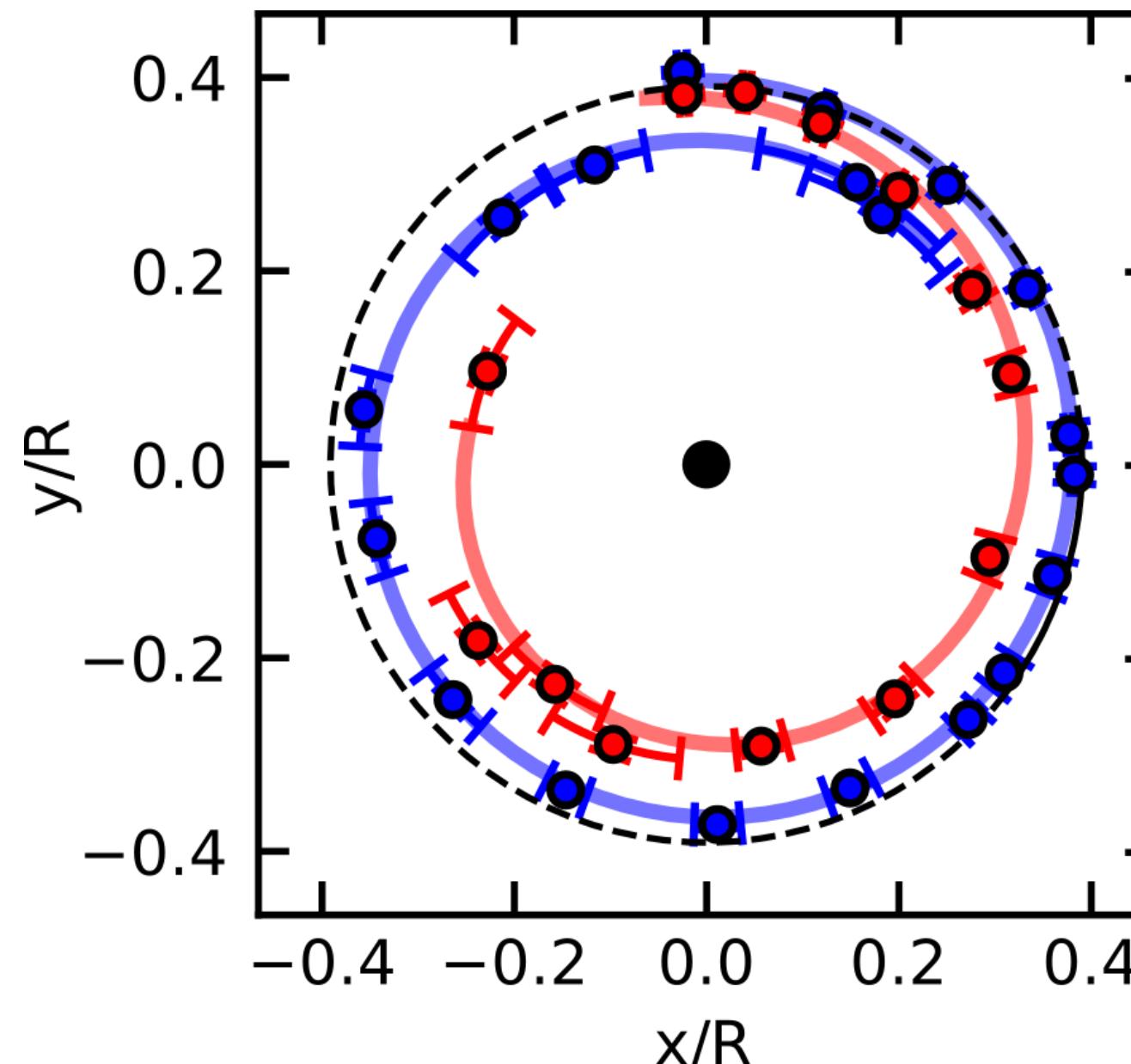
$$E \propto \log d$$
$$v \propto 1/d$$

Experimental realization



MUTUAL FRICTION IN A VORTEX DIPOLE

N. Grani, et al., arXiv:2503.21628 (2025)



Point Vortex Model (PVM)

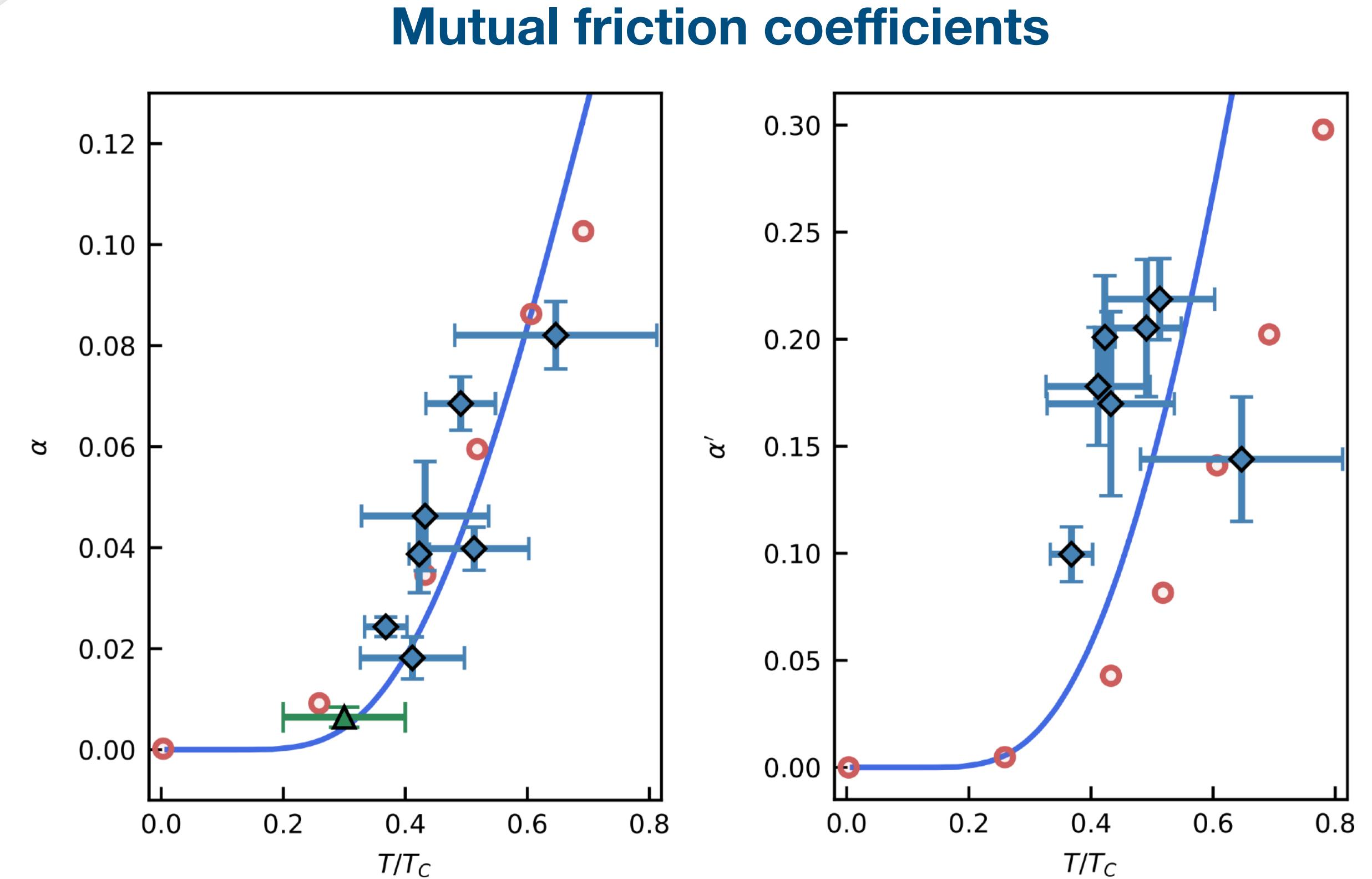
$$\frac{d\vec{r}}{dt} = (1 - \alpha')\vec{v}_s^0 - \alpha\sigma\hat{z} \times \vec{v}_s^0$$

$\sigma \rightarrow$ Circulation

$v_s^0 \rightarrow$ Superfluid velocity

$$T/T_c = 0.36(4)$$

$$T/T_c = 0.50(6)$$



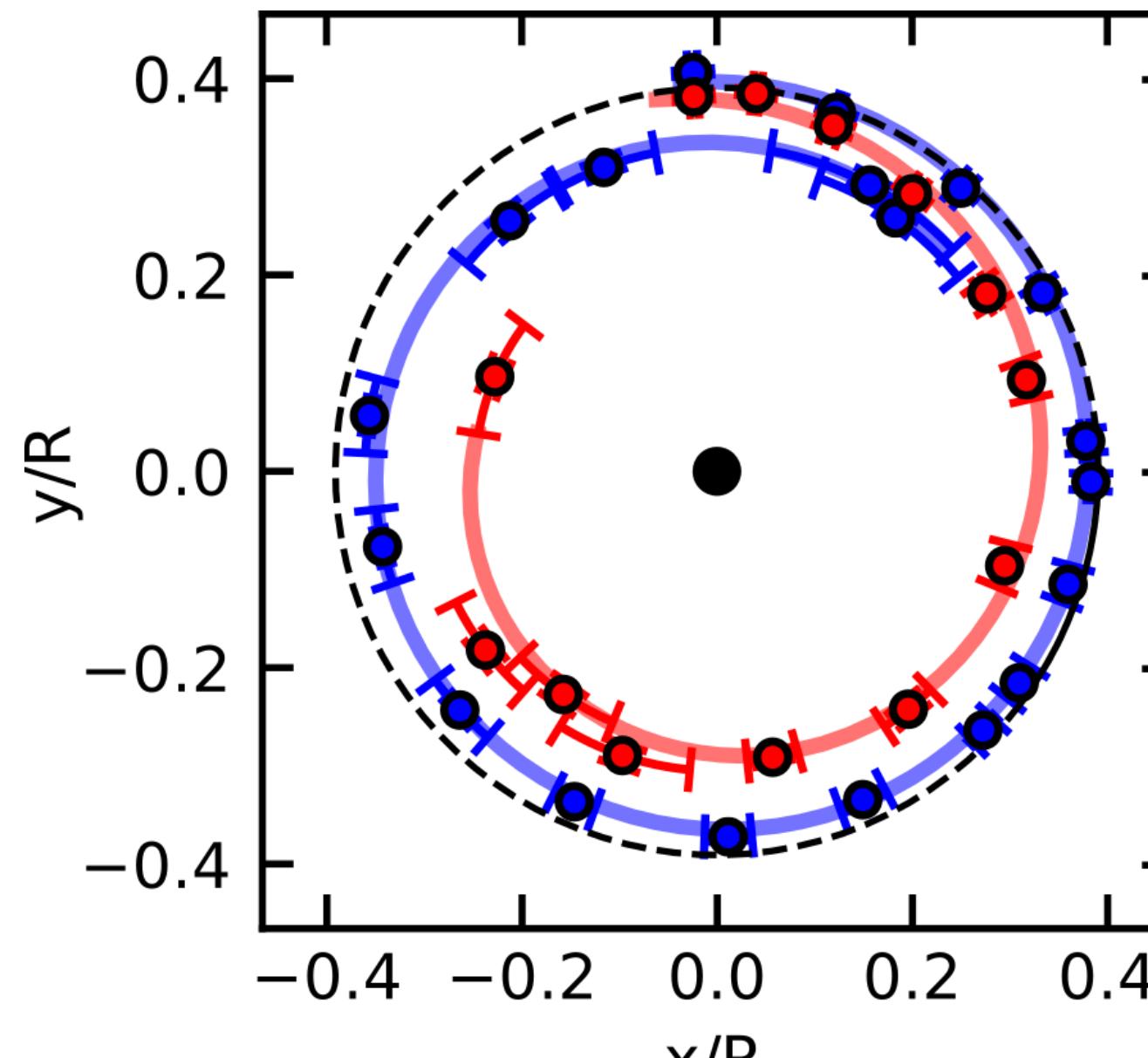
- SLDA simulation [1]
- Analytical model including localized QP [2]

[1] A. Bulgac, Phys. Rev. A 76, 040502 (2007)

[2] N. B. Kopnin, Reports on Progress in Physics 65, 1633 (2002)

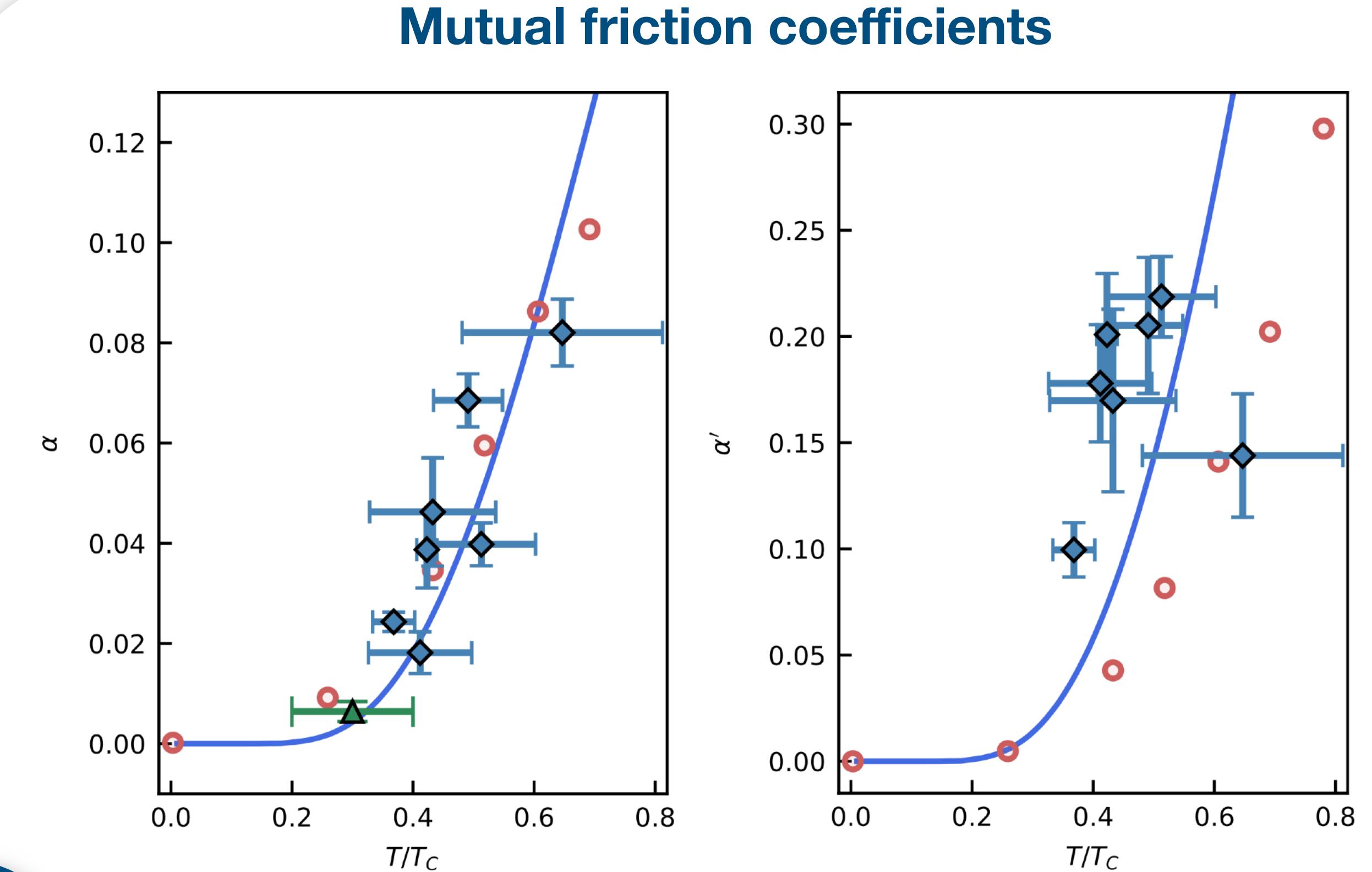
MUTUAL FRICTION IN A VORTEX DIPOLE

N. Grani, et al., arXiv:2503.21628 (2025)



Point Vortex Model (PVM)

- ▶ First experimental observation of **nonzero** α' in ultra cold atoms
- ▶ Connection with **microscopic mechanism** of dissipation



- SLDA simulation [1]
- Analytical model including localized QP [2]

[1] A. Bulgac, Phys. Rev. A 76, 040502 (2007)

[2] N. B. Kopnin, Reports on Progress in Physics 65, 1633 (2002)

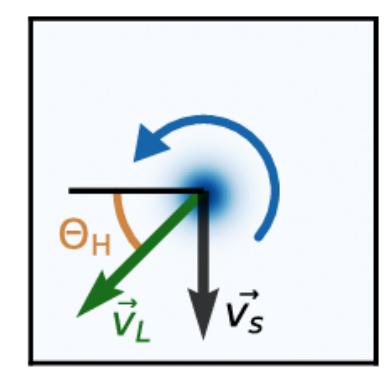
MICROSCOPIC DISSIPATION

N. Grani, et al., arXiv:2503.21628 (2025)

Can we extract information on **microscopic quantities** given the **phenomenological** parameters α and α' ?

- ▶ Hall angle: direction of the vortex motion with respect to the superfluid velocity

$$\Theta_H = \tan^{-1} \frac{1 - \alpha'}{\alpha} = \omega_0 \tau$$



- ▶ Localized quasiparticle lifetime

$$\tau = \Theta_H / \omega_0$$

$$\hbar \omega_0 = |\Delta|^2 / E_F$$

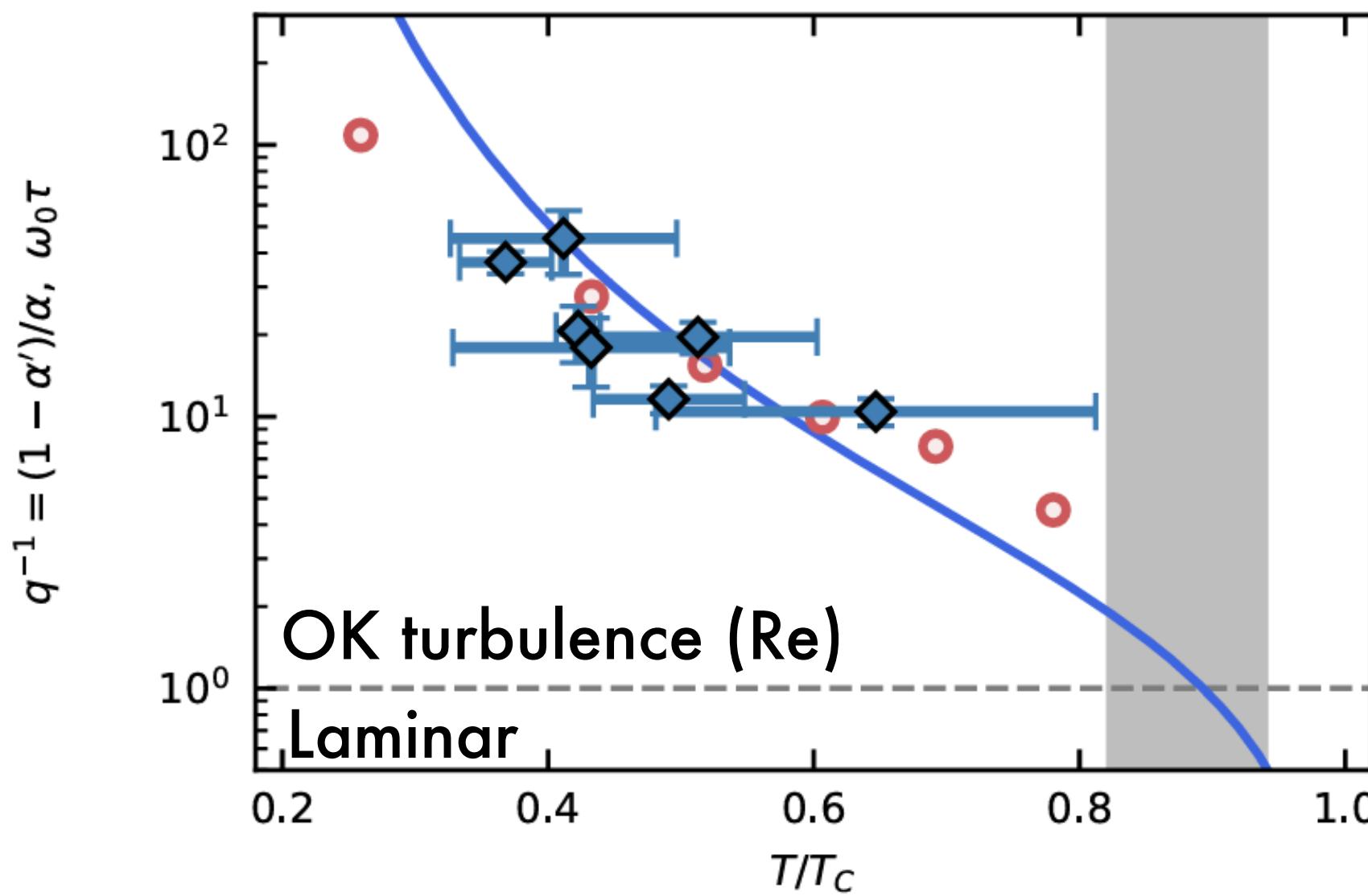
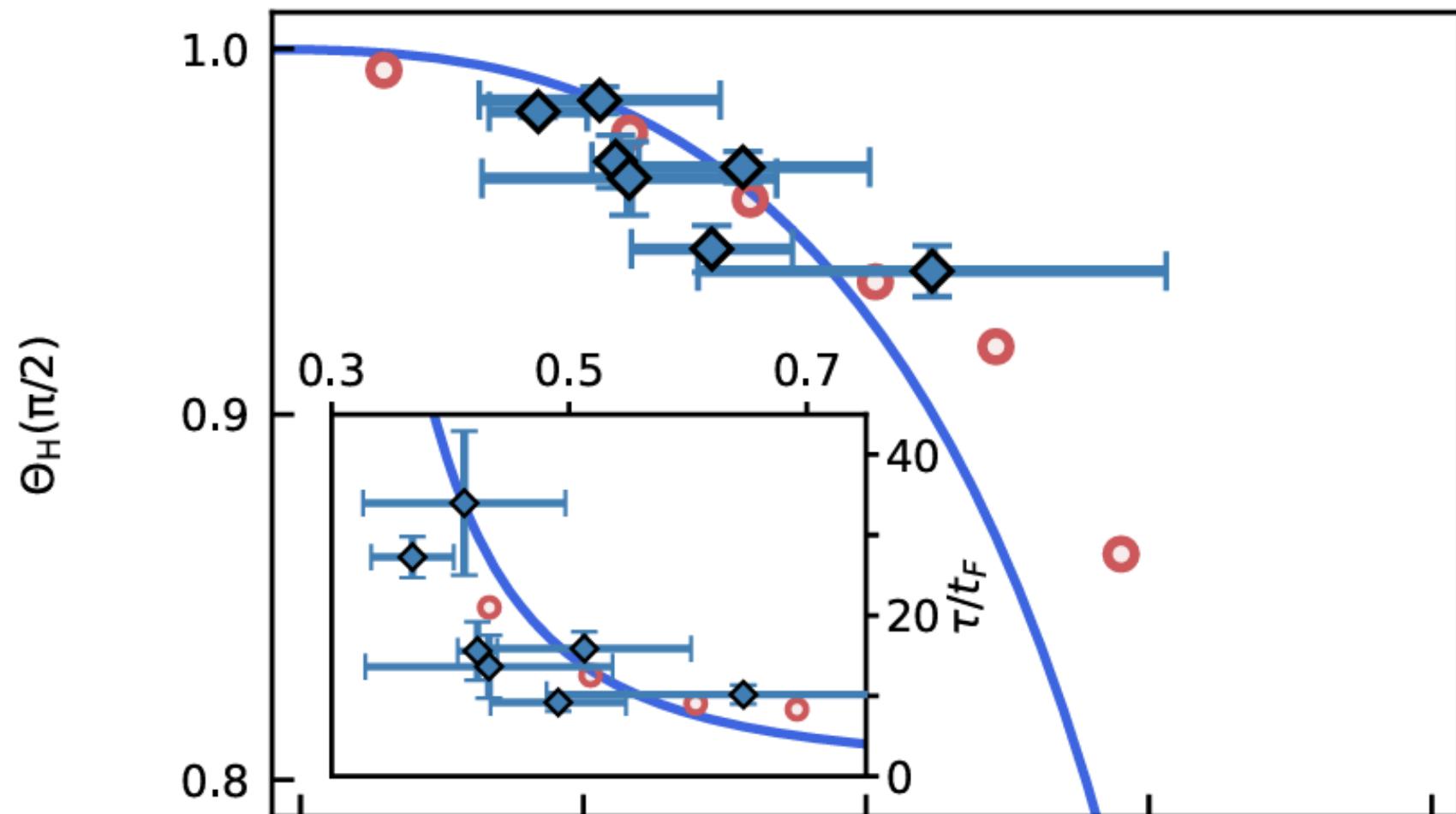
- ▶ Intrinsic superfluid parameter: regime for superfluid turbulence

$$q = \frac{\alpha}{1 - \alpha'}$$

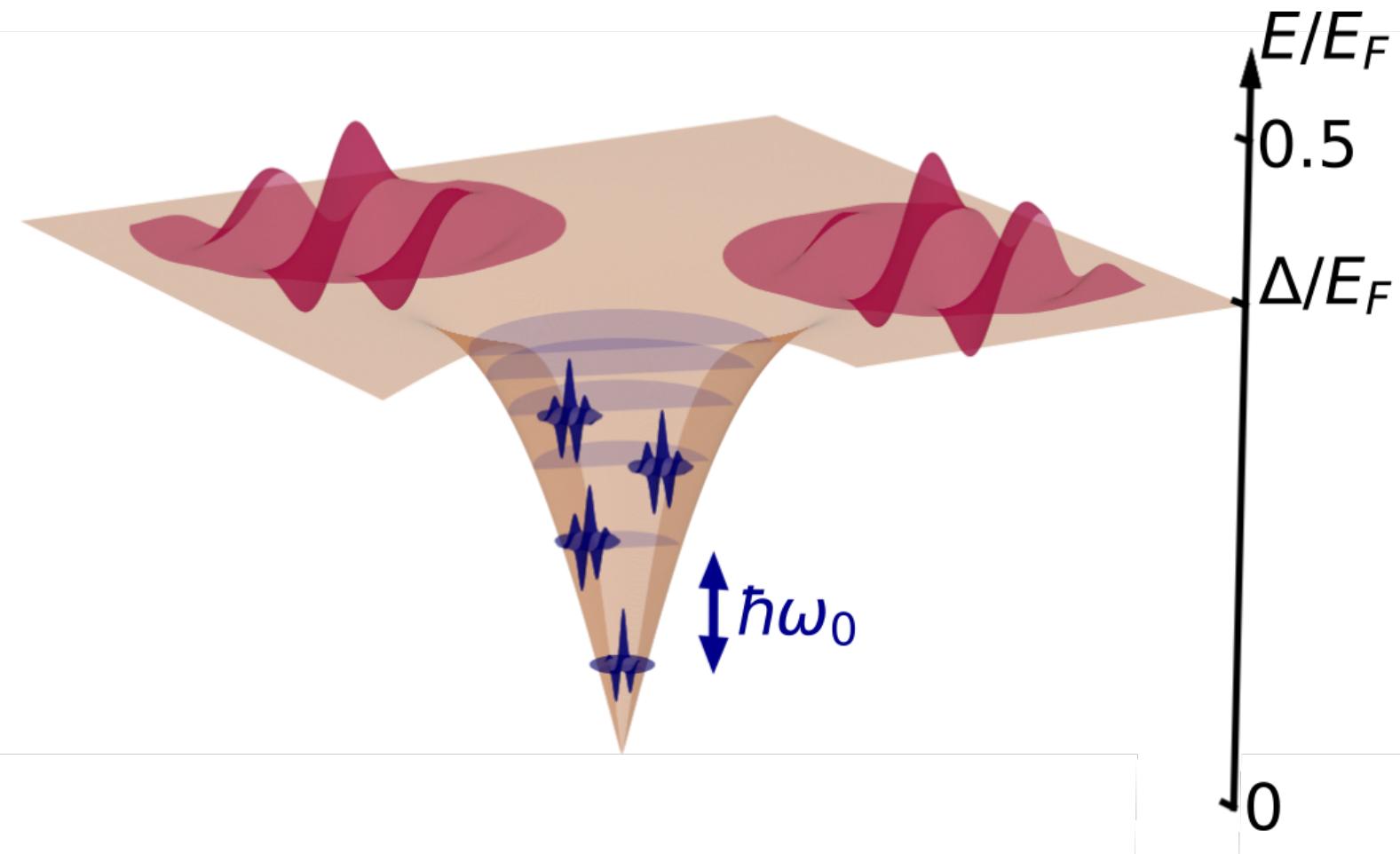
M. Heyl, et al., Nat. Comm. 13, 6986 (2022).

N. Kopnin and M. Salomaa, Phys. Rev. B 44, 9667 (1991).

A. Finne, et al., Nature 424, 1022 (2003).



SOUND INTERACTION



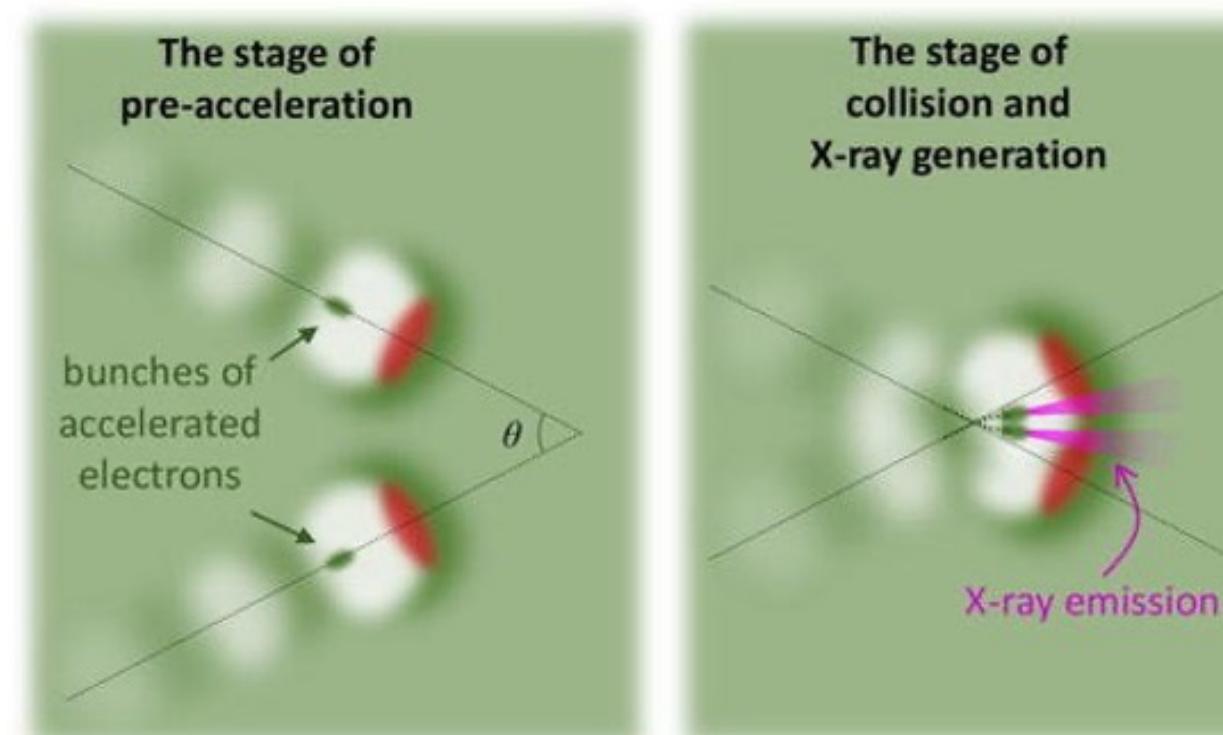
Sound interaction

The **incompressible** energy of the vortex can be transformed into **compressible** energy of the superfluid: sound emission

Also at zero temperature

Caroli, De Gennes, and Matricon. *Phys. Lett.* 9.4 307-309 (1964)

R. Sensarma, et al., *Phys. Rev. Lett.* 96.9 090403 (2006)



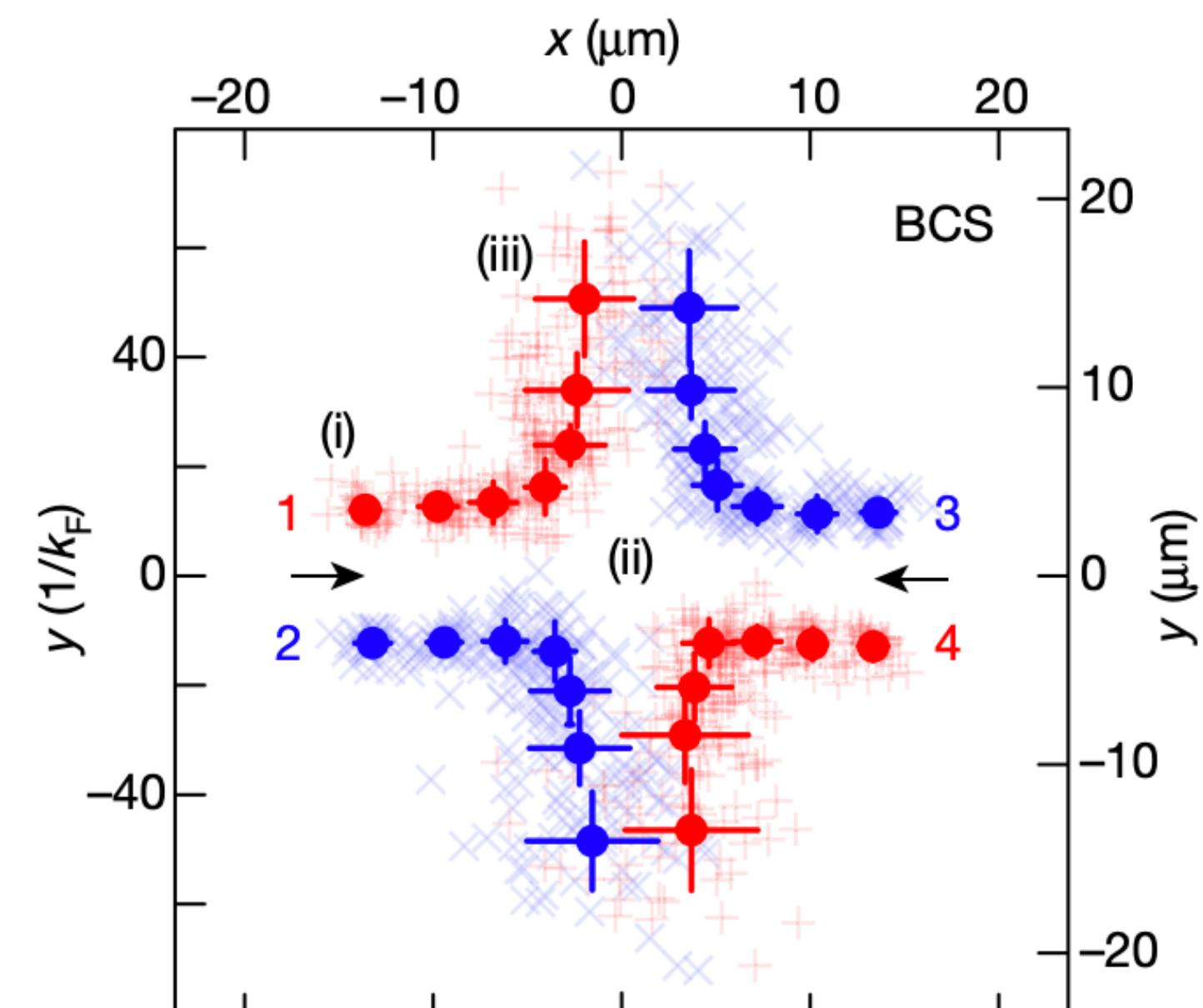
- ▶ Vortices at short distance
- ▶ Acceleration —> sound emission
- ▶ Vortex-dipole collisions

Picture taken from: E. Vallin et al., *Physics of Plasmas* 24, 093101 (2017)

DIPOLE-DIPOLE COLLISIONS

W. J. Kwon, et al. *Nature* 600, 64–69 (2021).

To promote **sound dissipation** we work at T=0 and we engineer collisions between short-distance dipole



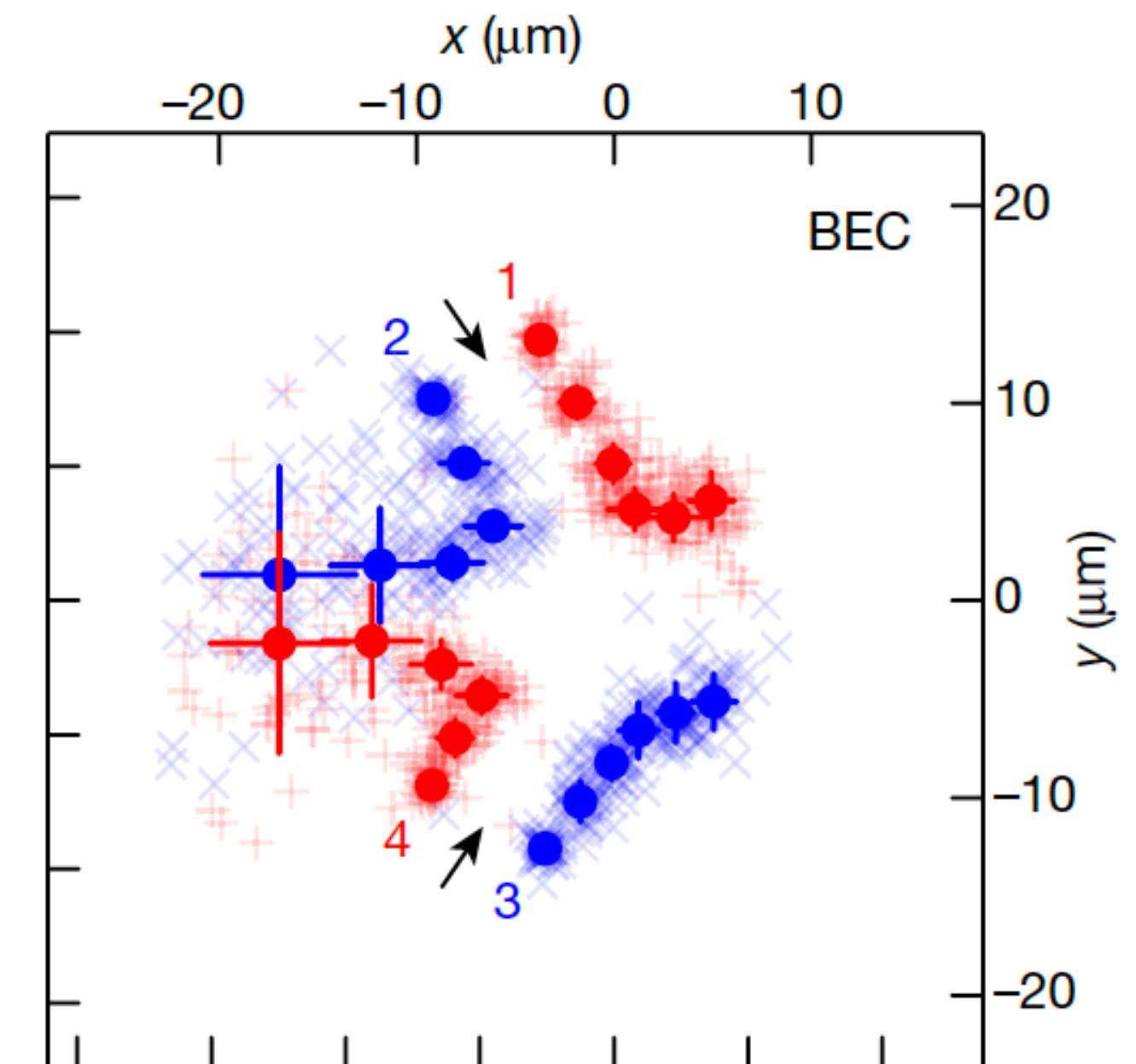
Head-on collision

Dissipation

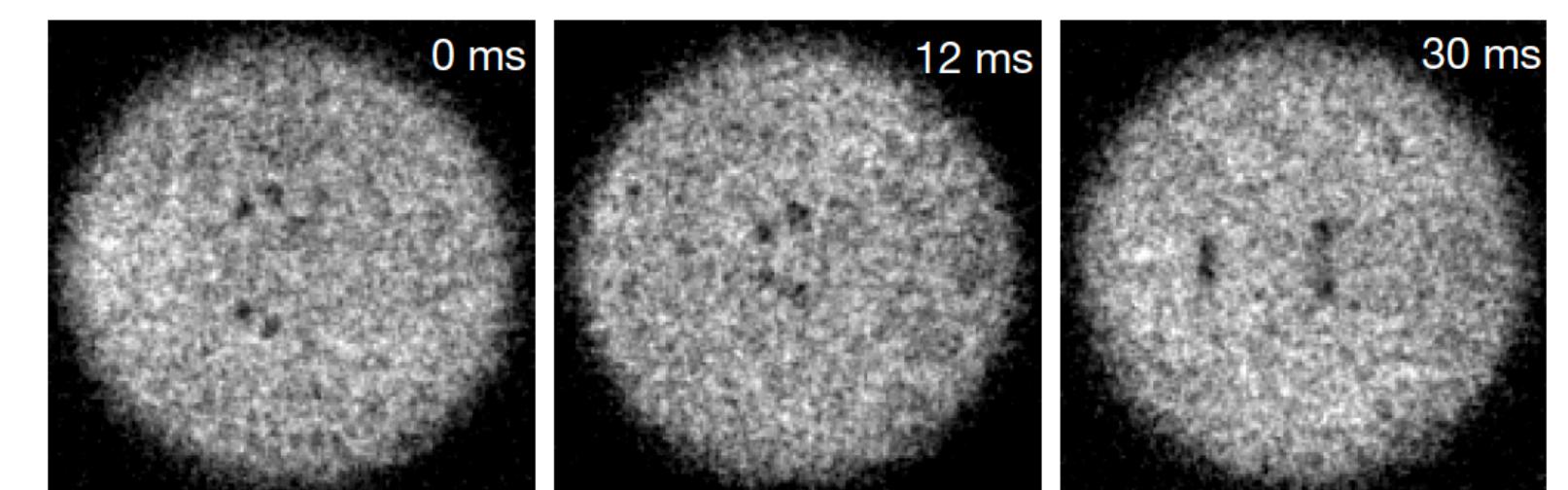
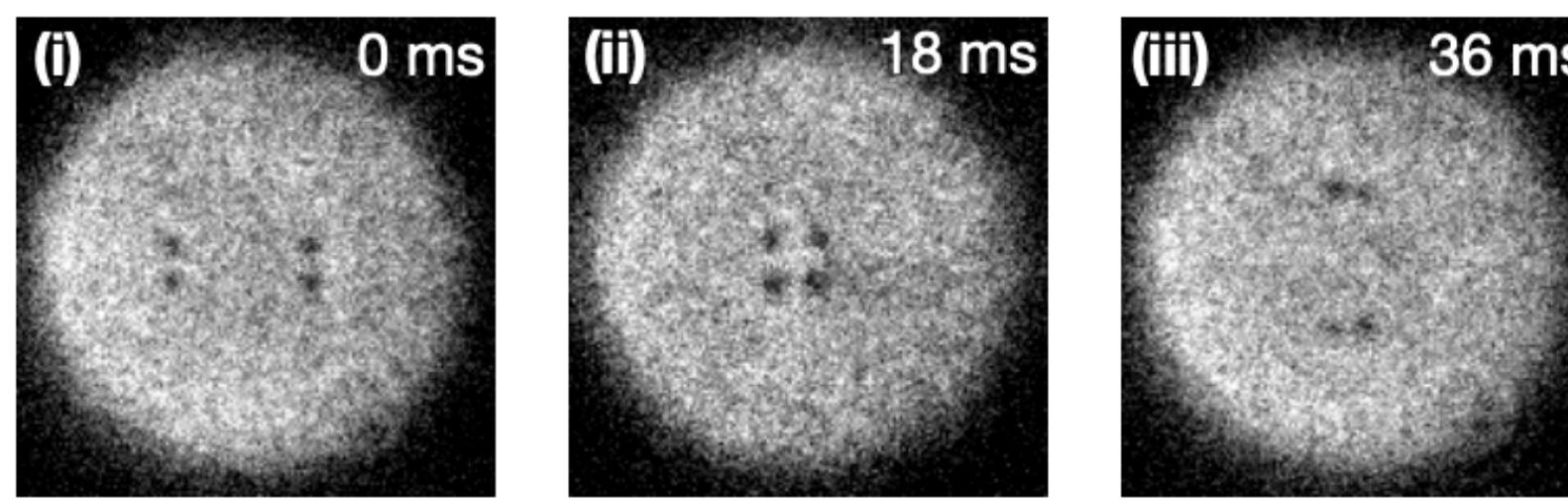
$$E \propto \log d$$

Dipole size ratio after/before the collision

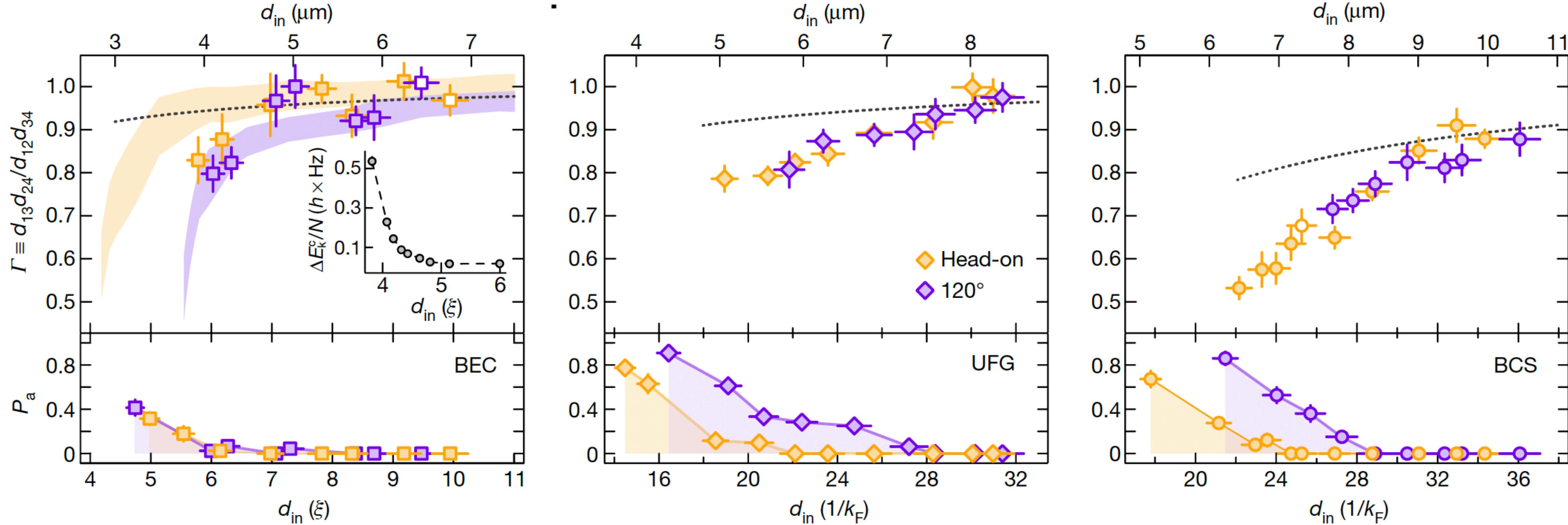
$$\Gamma = \frac{d_{13}d_{24}}{d_{12}d_{34}}$$



120° collision



SOUND DISSIPATION IN VORTEX COLLISION

 W. J. Kwon, et al. *Nature* **600**, 64–69 (2021).


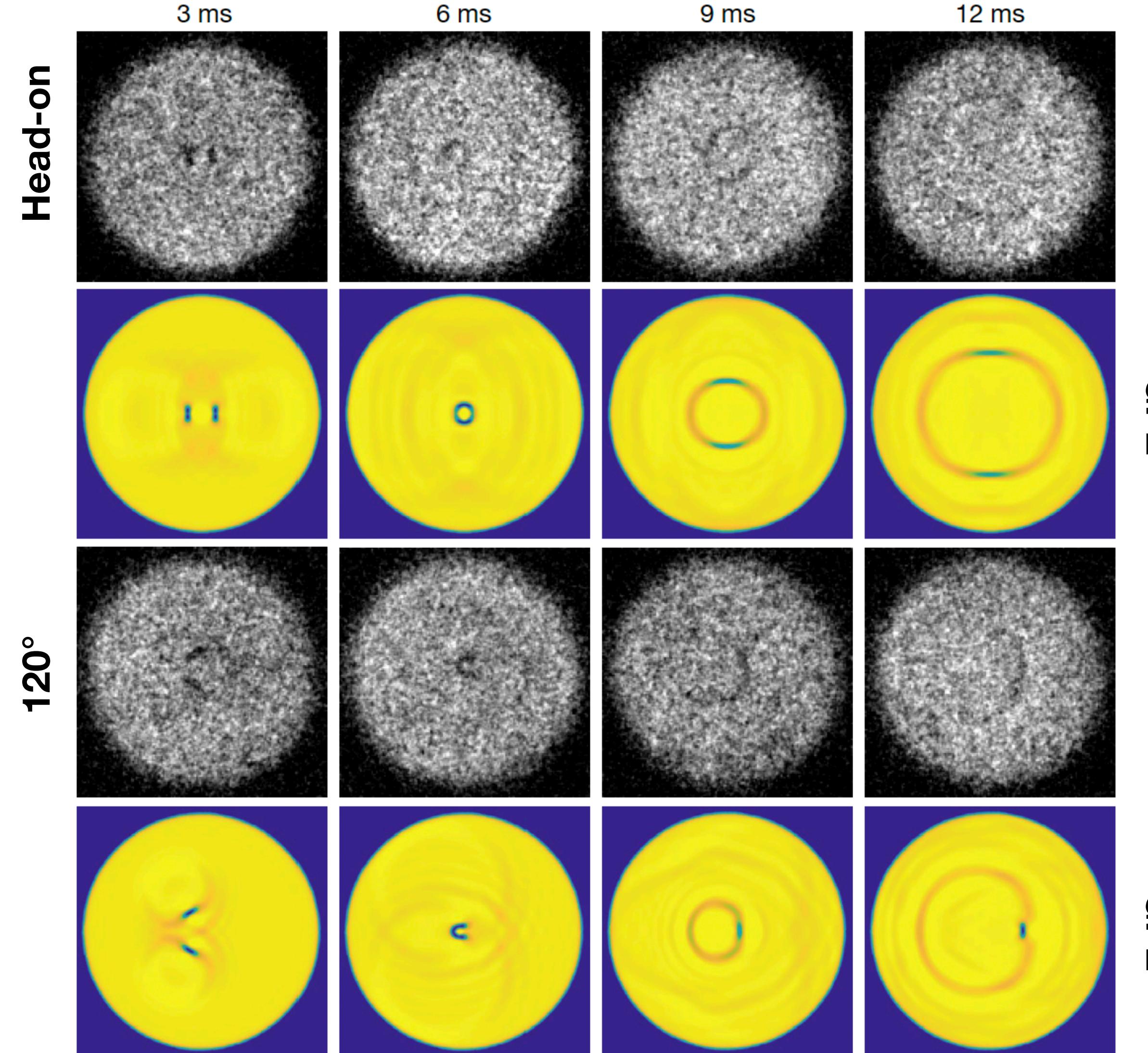
..... PVM

GPE

- ▶ The disagreement with PVM signal the presence of a different source of dissipation at short distance: **sound**
- ▶ The sound dissipation increases for **fermionic** superfluids
- ▶ For shorter dipole the probability of **annihilation** increases

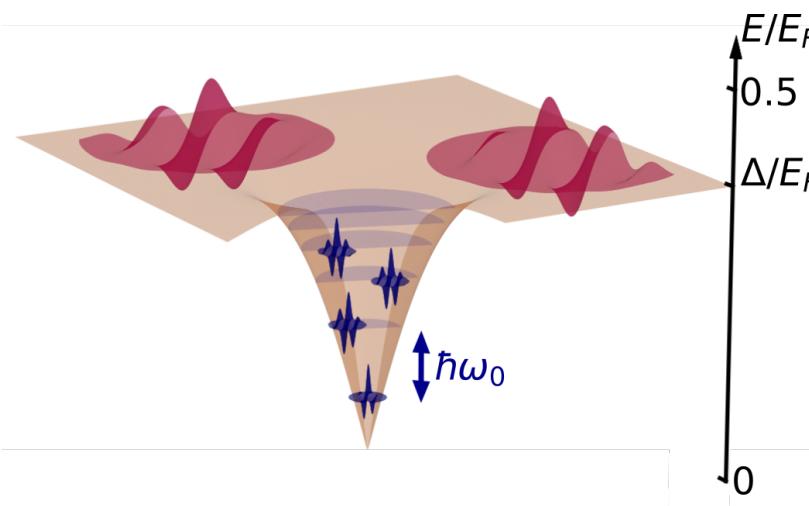
VORTEX - ANTIVORTEX ANNIHILATION

W. J. Kwon, et al. *Nature* 600, 64–69 (2021).

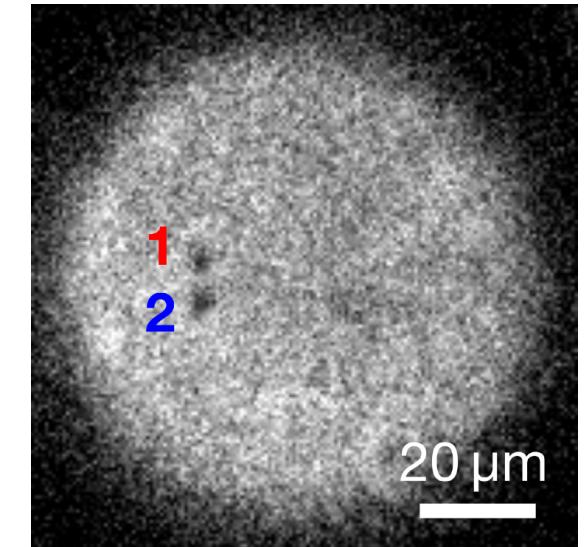


- ▶ Vortex-antivortex annihilation via **sound emission**: the incompressible energy of the dipole is fully converted in compressible sound energy
- ▶ The sound emission is non-symmetric (rarefaction pulse)

CONCLUSIONS

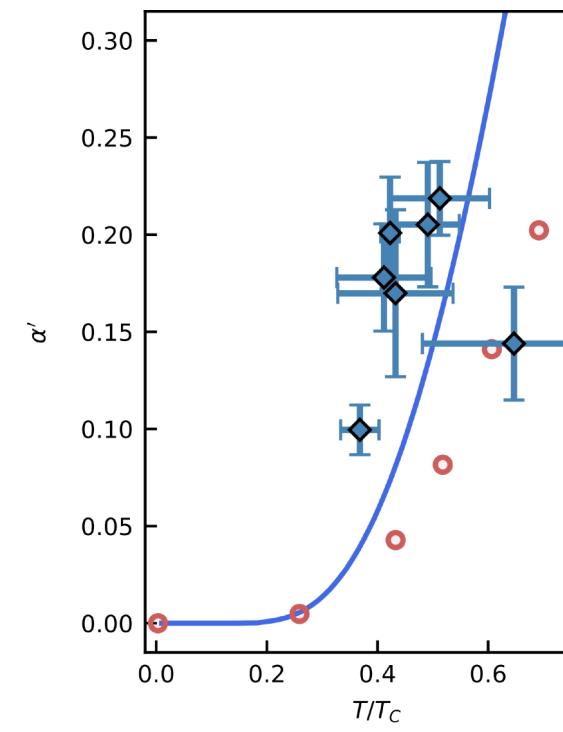
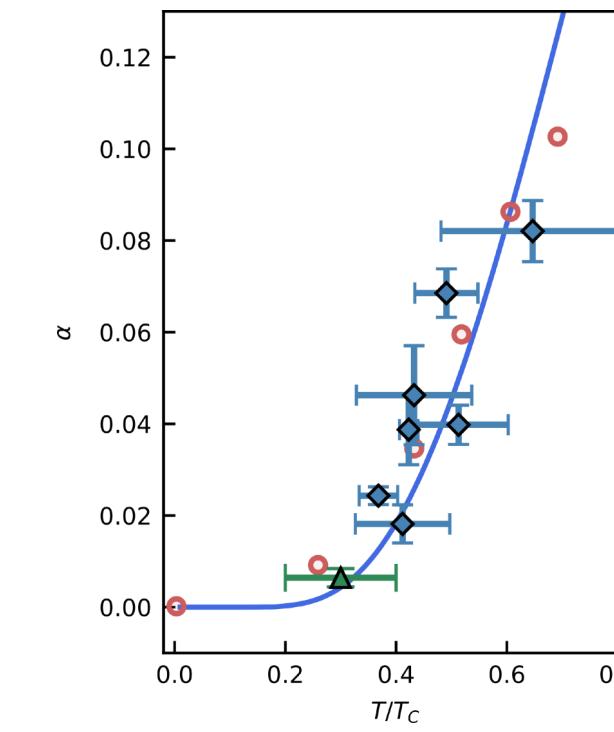
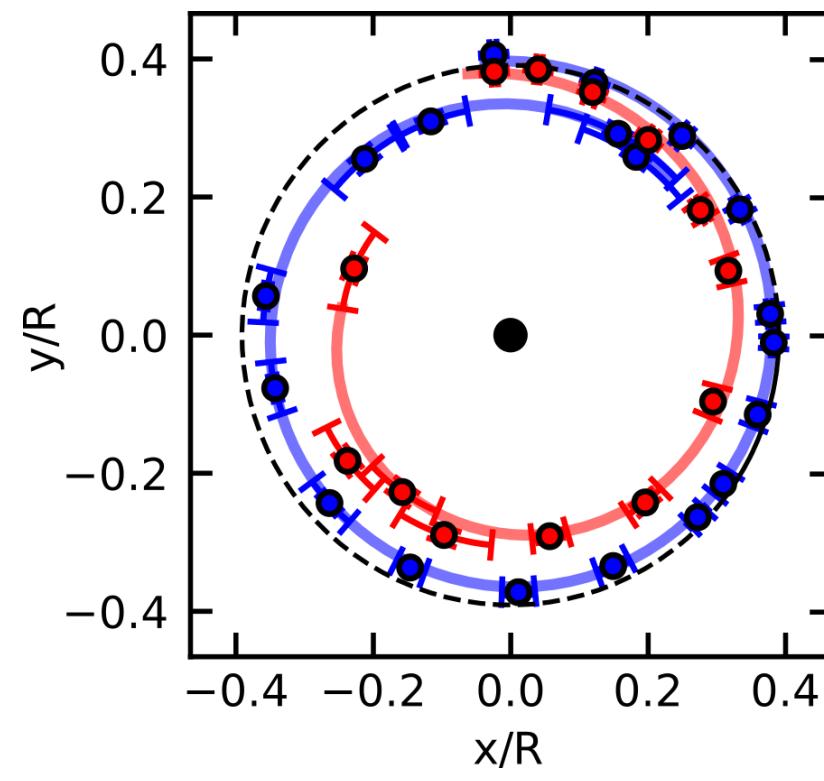


- ▶ Ultracold atoms are an ideal platform for **studying vortex dynamic in a bottom-up approach**: create on-demand vortex configuration, detect with micron-scale resolution
- ▶ We reconstruct the vortex configuration dynamics by following the **vortex trajectories**



Mutual friction

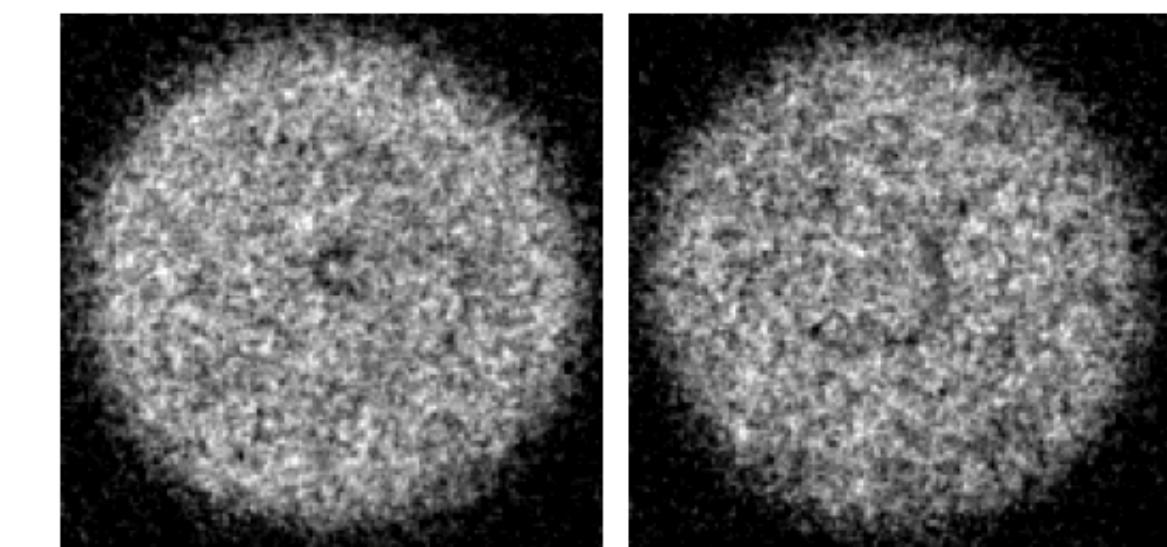
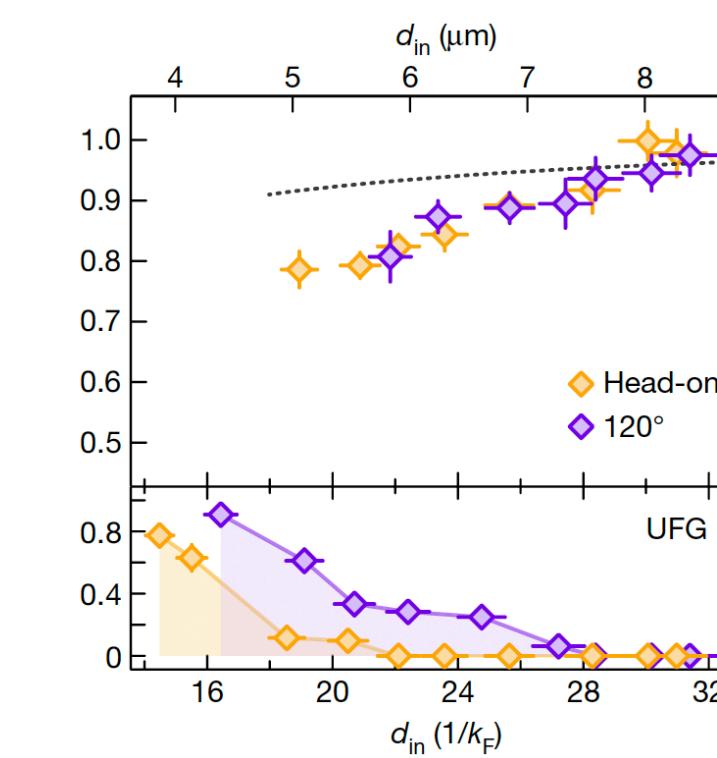
- ▶ We observe an **increase of mutual friction** for increasing temperature, and extract the mutual friction coefficients
- ▶ We identify a role of **localized quasiparticles** in dissipation of vortices in Fermi superfluids



N. Grani, et al., arXiv:2503.21628 (2025)

Sound interaction

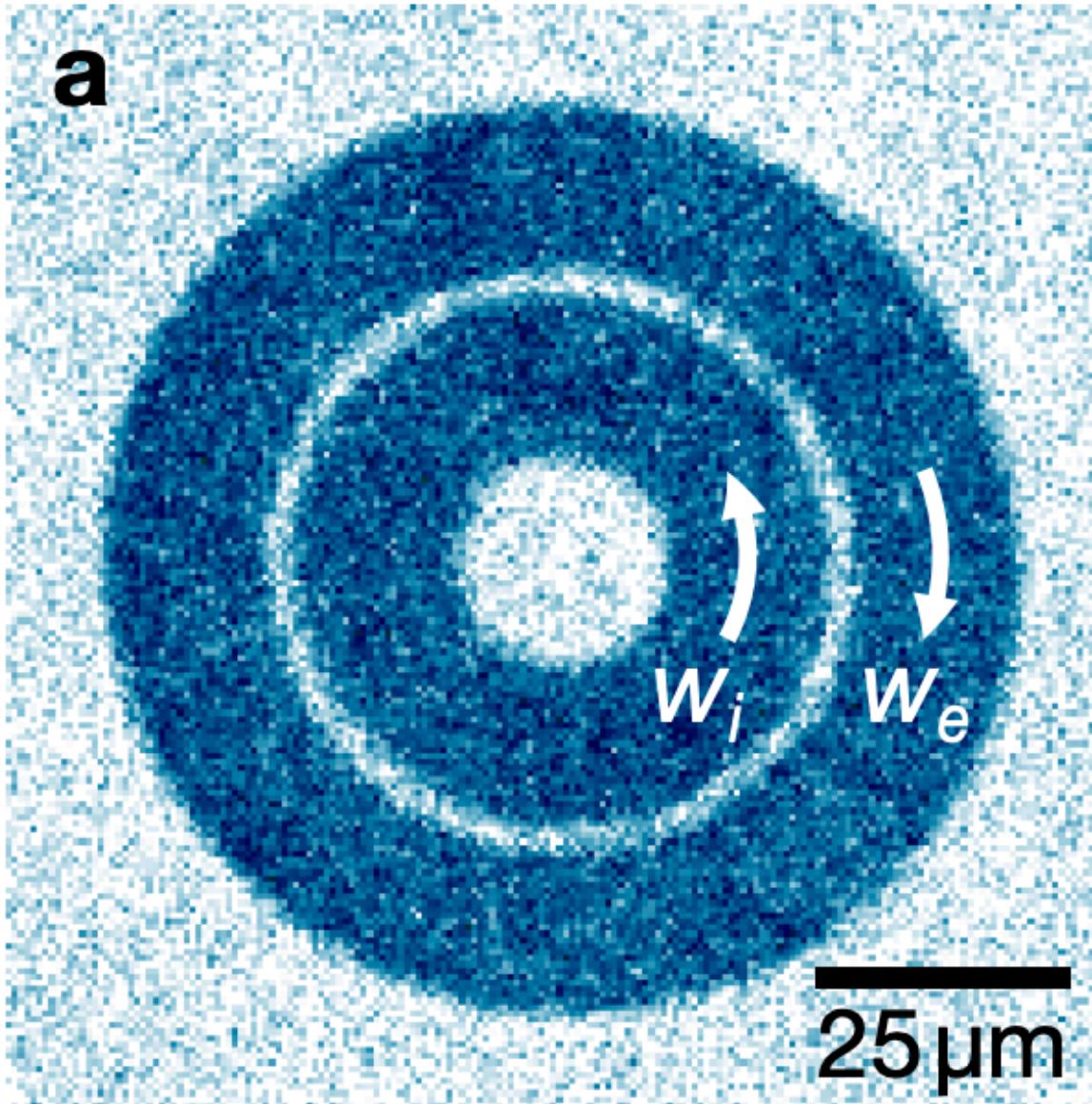
- ▶ In colliding dipoles at short distance we observe an excess dissipation due to **sound interaction** (incompressible → compressible energy)
- ▶ We directly visualize **vortex-antivortex annihilation** via sound emission



W. J. Kwon, et al. Nature 600, 64–69 (2021).

MORE VORTICES?

D. Hernández-Rajkov, et al. *Nat. Phys.* **20**, 939–944 (2024).

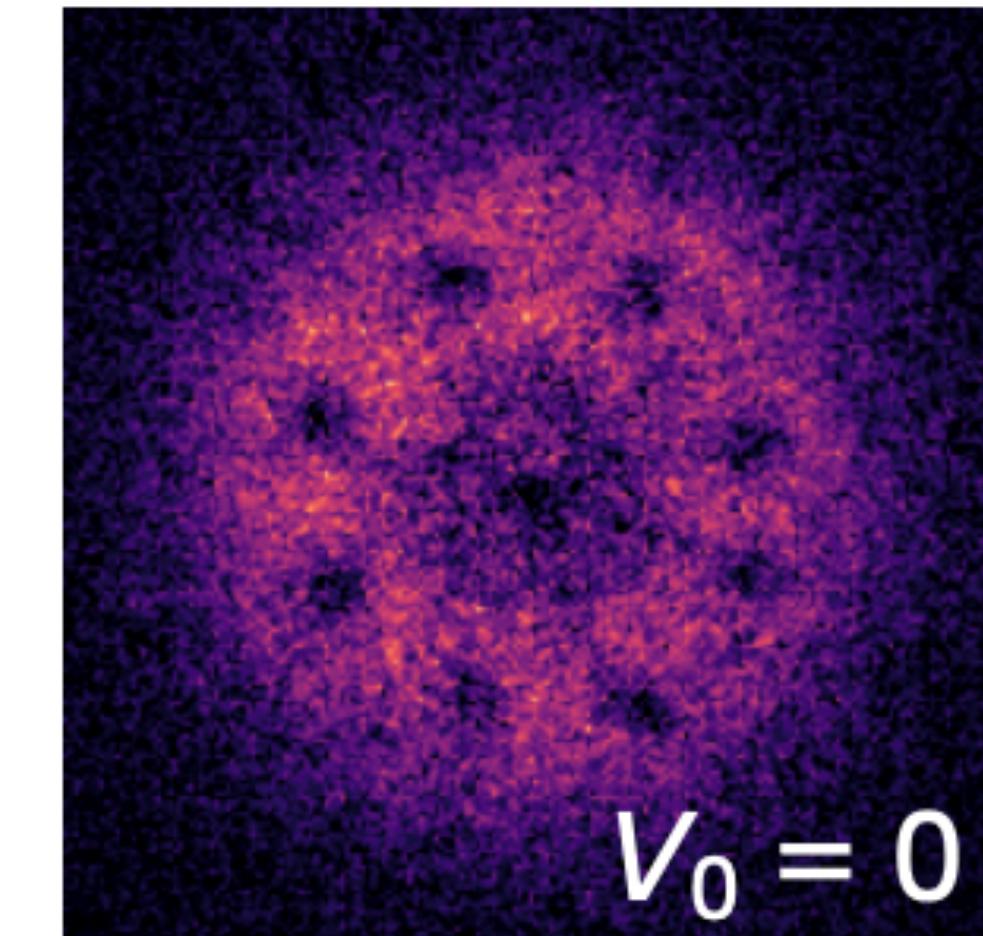
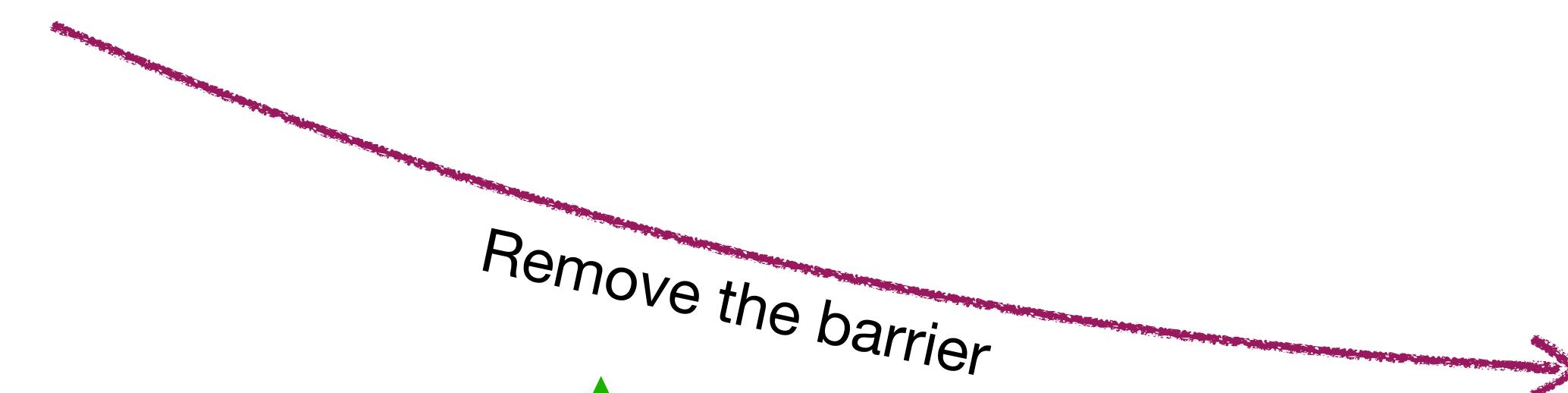


Double concentric ring geometry

We imprint opposite winding $w_i = -w_e$ in the two rings

The relative velocity at the interface of radius R_0 is:

$$\Delta v = \frac{\hbar(w_e - w_i)}{mR_0}$$



See Diego Hernandez
Rajkov talk on Thursday!

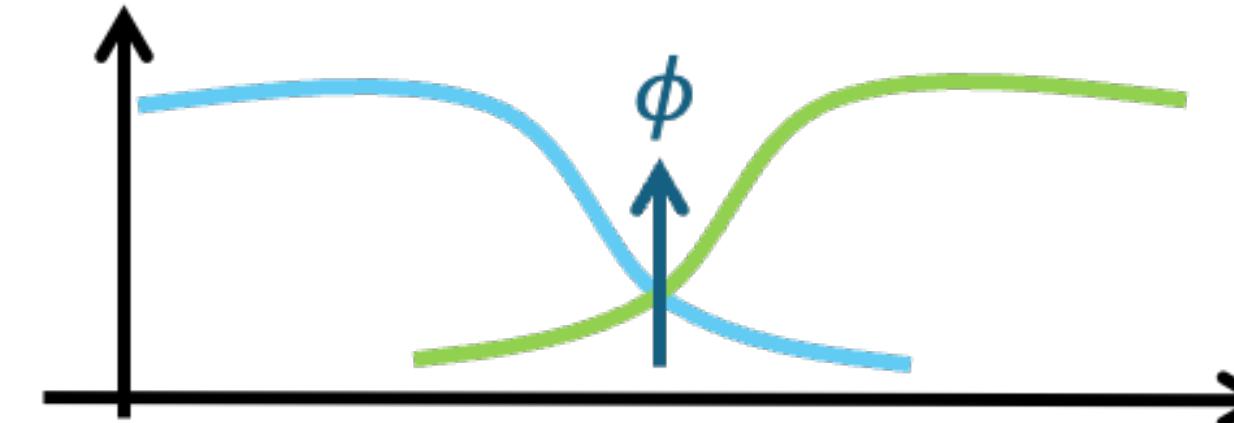
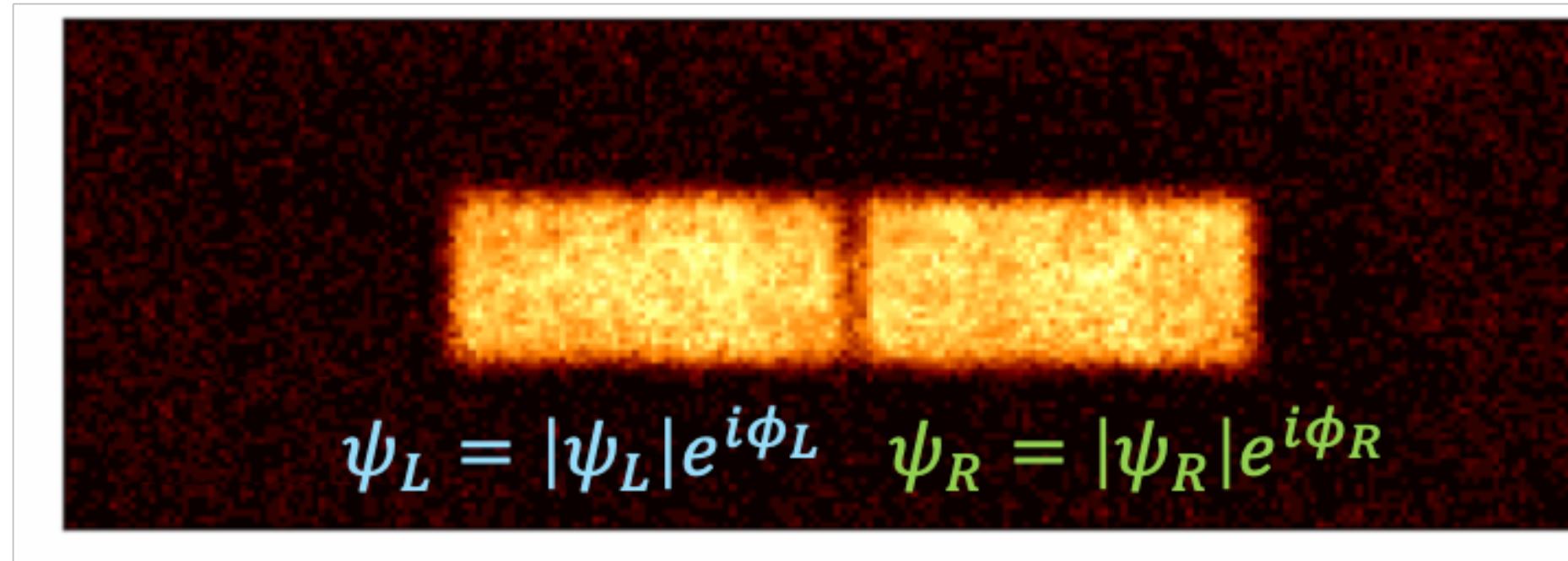
Vortex necklace

With a number of vortices $N_v = \Delta w$

MORE FROM LILAB?

W. J. Kwon, et al. Science 369.6499: 84-88 (2020)

G. Del Pace, et al., arXiv preprint arXiv:2409.03448 (2024)

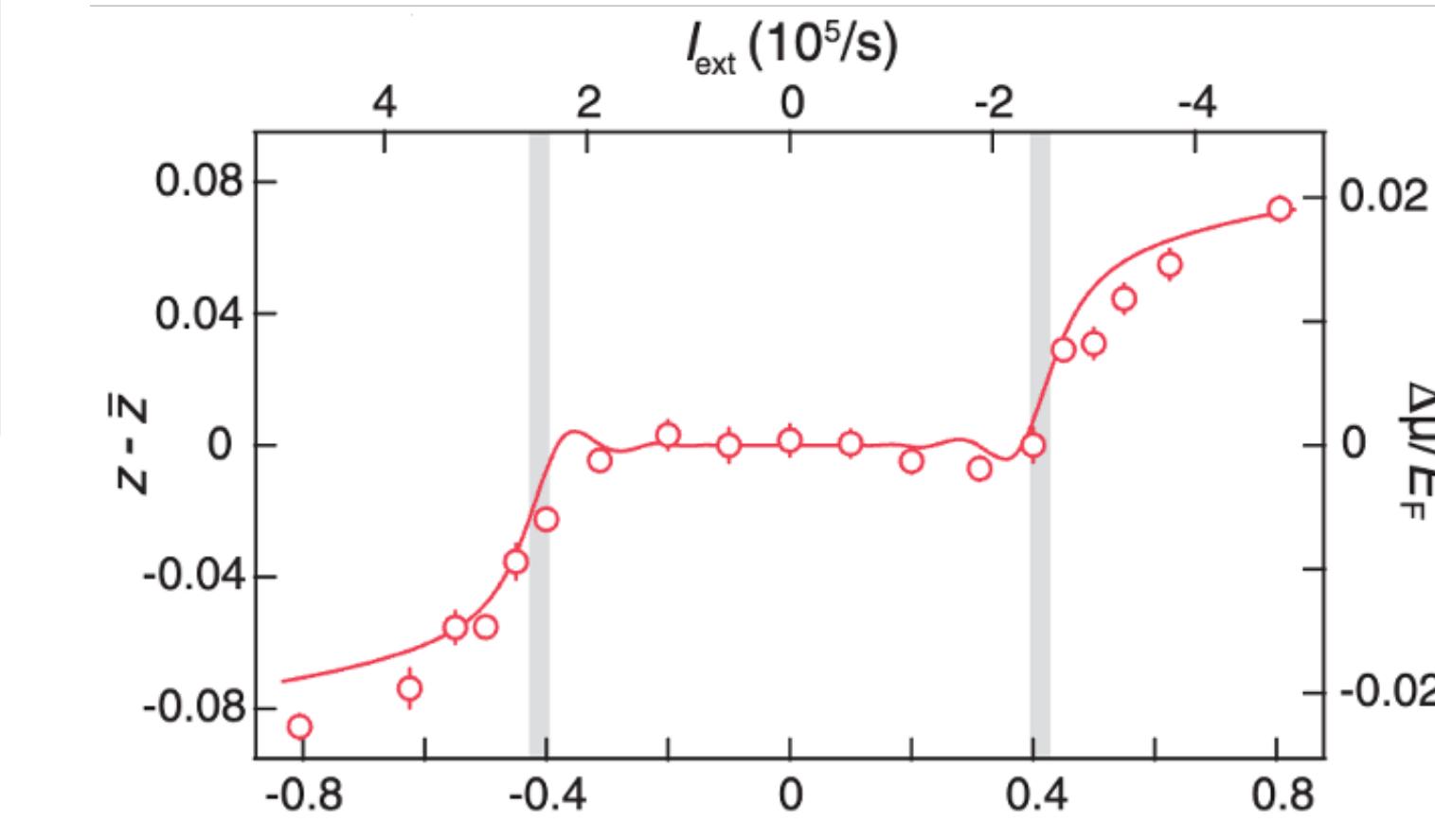


Josephson-Anderson equations:

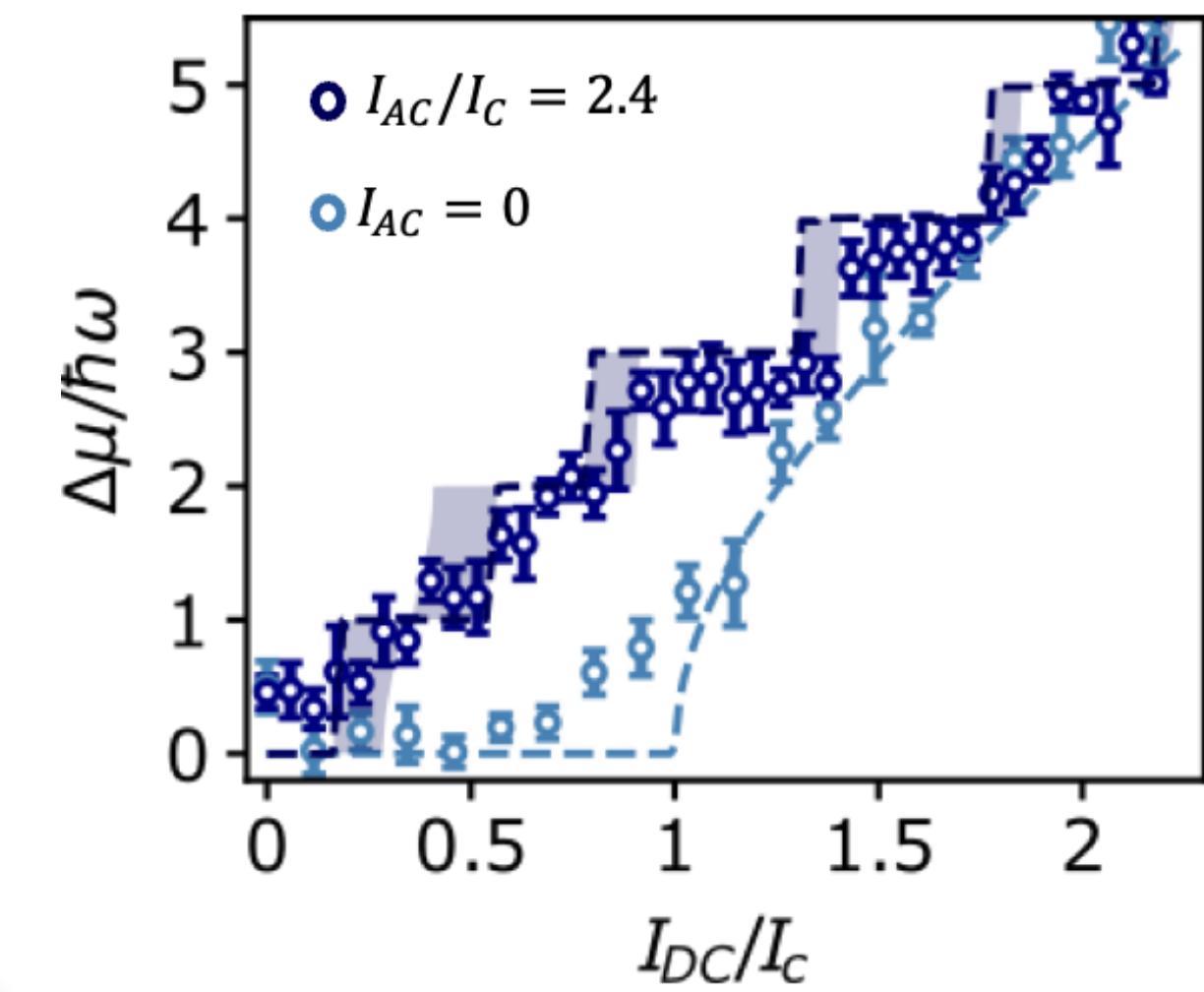
$$I = I_c \sin(\phi)$$

$$\dot{\phi} = -\frac{\Delta\mu}{\hbar}$$

DC Josephson effect

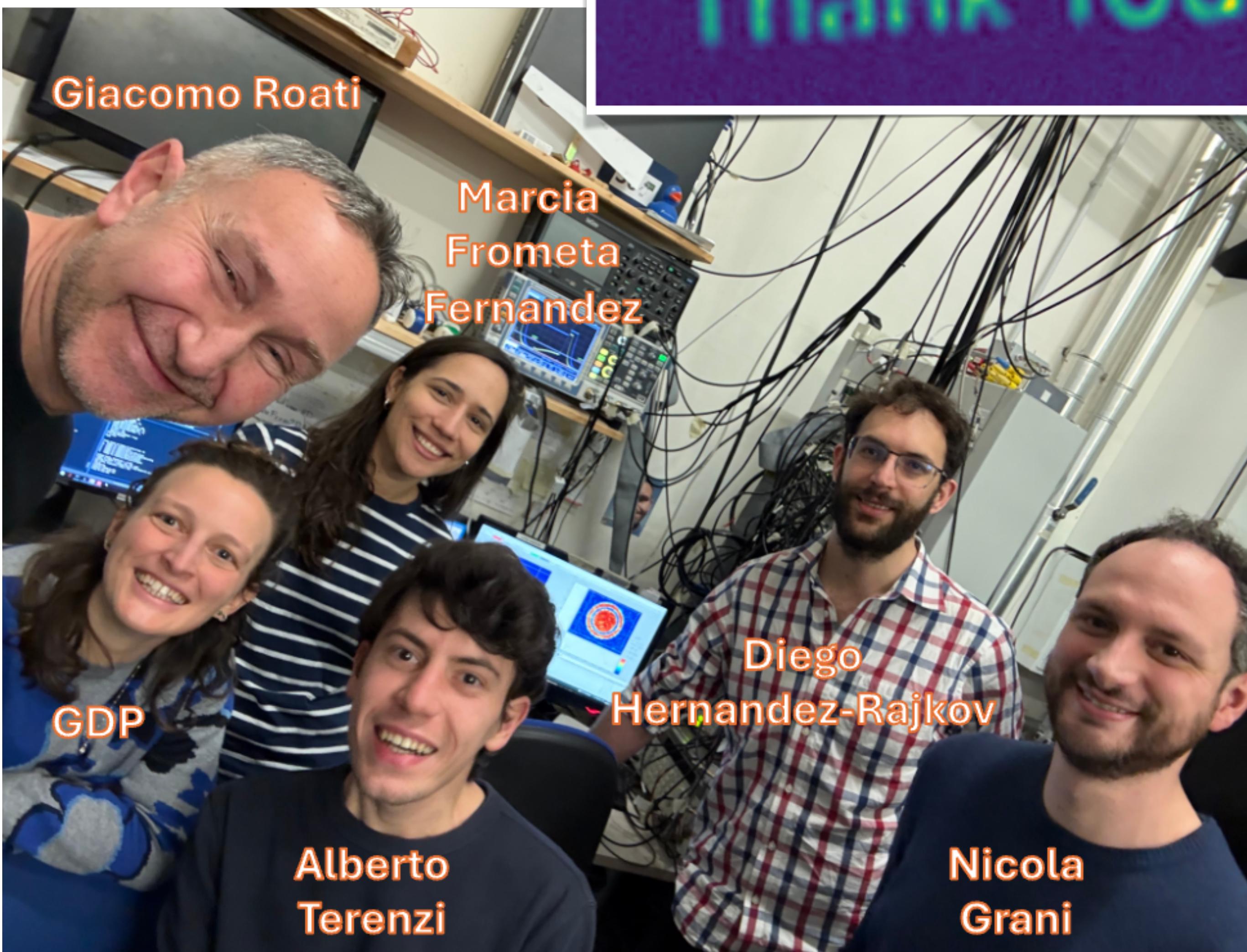


Shapiro steps under AC drive



See Marcia Frometa
Fernandez poster this
afternoon!

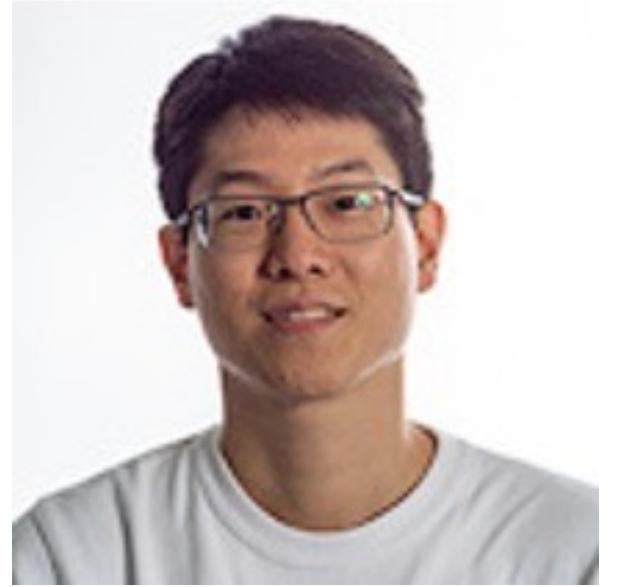
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Keljdja Xhani, Politecnico di Torino
Luca Galantucci, IAC-CNR

FORCES ON VORTICES

$$\vec{F}_N = D(\vec{v}_n - \vec{v}_L) + D' \hat{z} \times (\vec{v}_n - \vec{v}_L),$$

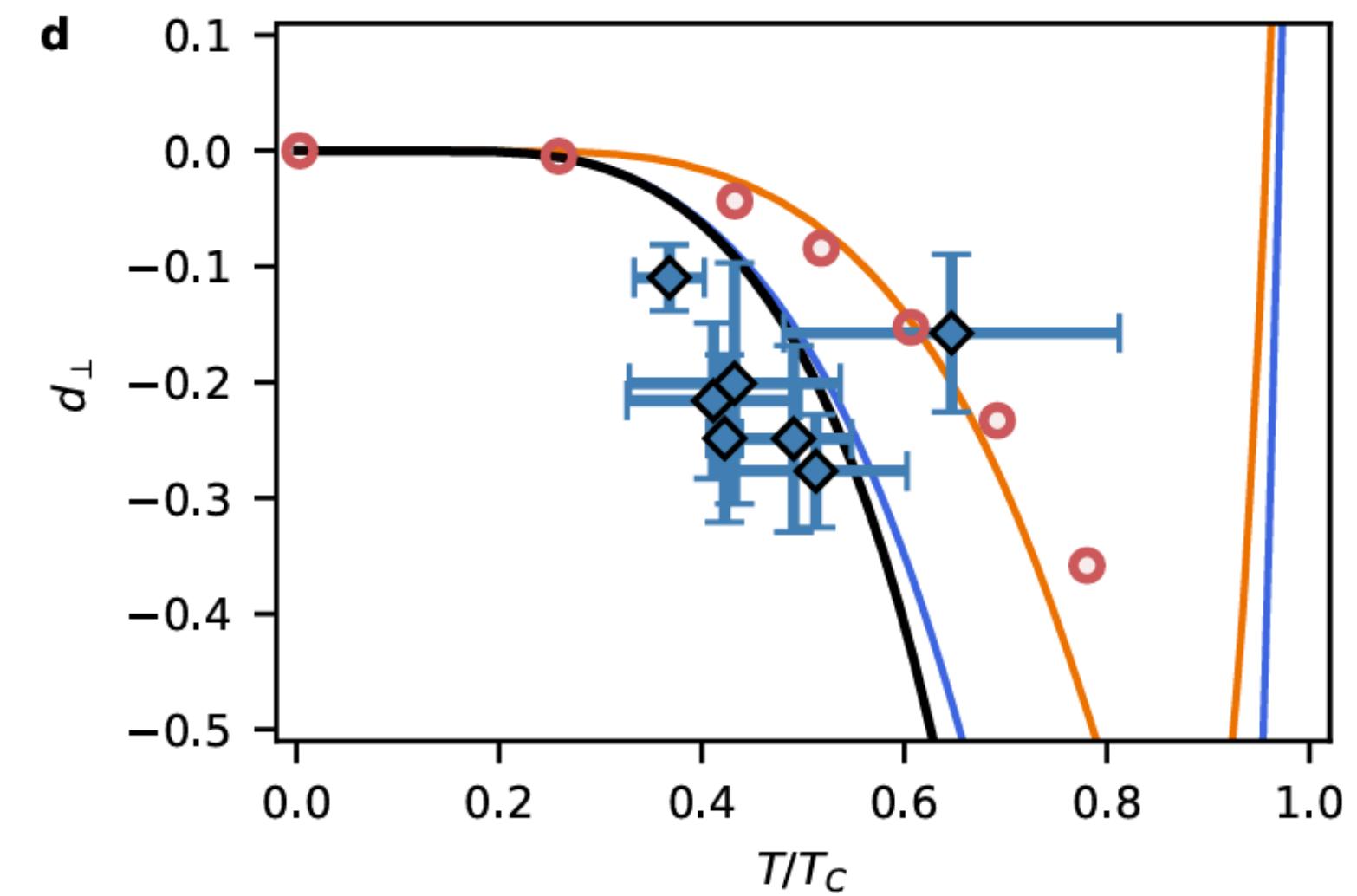
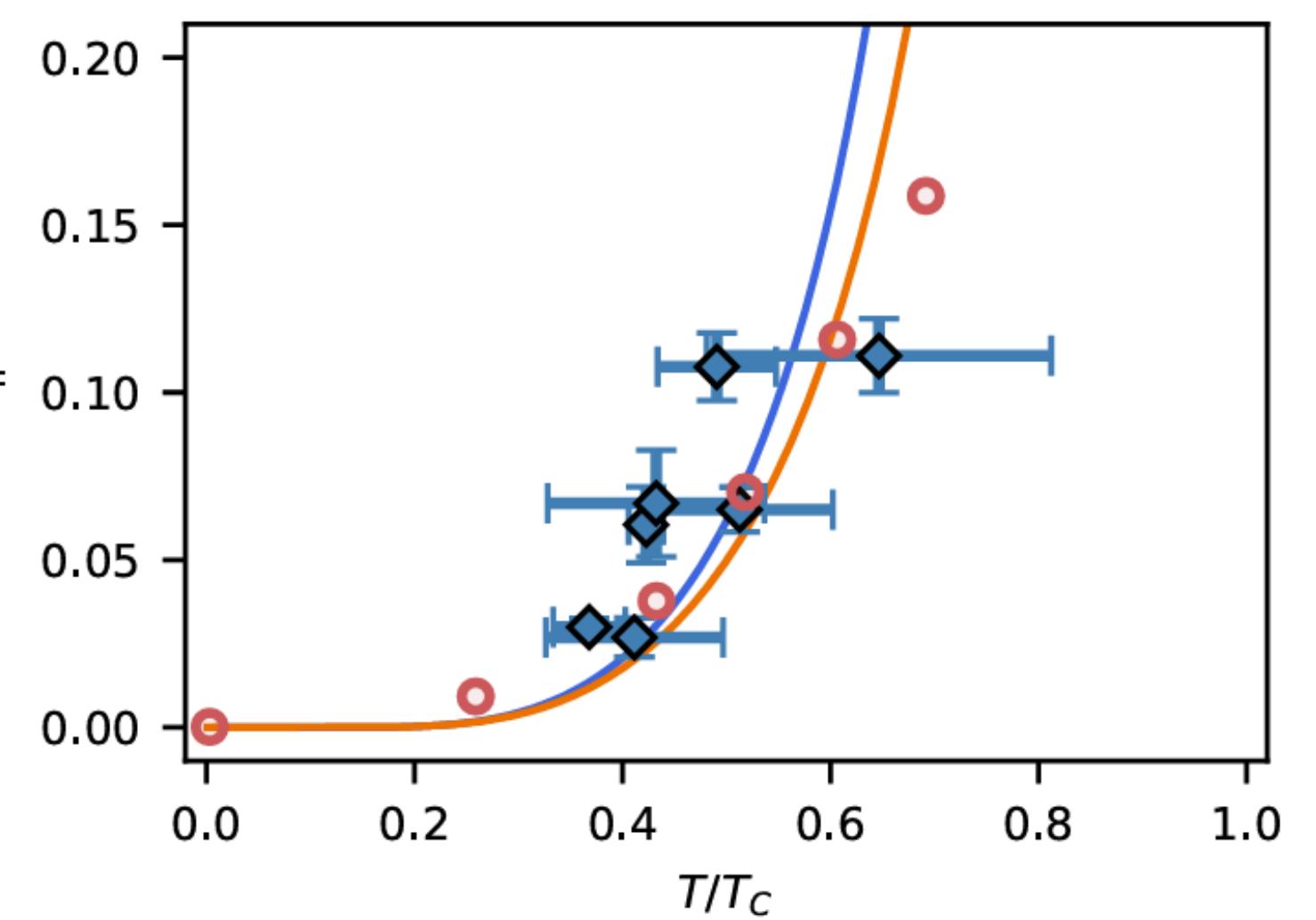
$$\vec{F}_M = \kappa \rho_s (\vec{v}_s - \vec{v}_L) \times \hat{z}$$

$$\vec{F}_N + \vec{F}_M = 0,$$

$$\alpha = \frac{d_{||}}{d_{||}^2 + (1 - d_{\perp})^2},$$

$$\alpha' = 1 - \frac{1 - d_{\perp}}{d_{||}^2 + (1 - d_{\perp})^2},$$

$$d_{||} = \frac{D}{\kappa_c \rho_s}, \quad d_{\perp} = \frac{D'}{\kappa_c \rho_s}.$$



$$\alpha = \frac{\rho_s}{\rho} \left(\omega_0 \tau \tanh \frac{|\Delta|}{2k_B T} \right)^{-1},$$

$$\alpha' = 1 - \frac{\rho_s}{\rho} \left(\tanh \frac{|\Delta|}{2k_B T} \right)^{-1},$$

$$d_{||} = \frac{\rho}{\rho_s} \frac{\omega_0 \tau}{1 + \omega_0^2 \tau^2} \tanh \left(\frac{\Delta}{2k_B T} \right),$$

$$d_{\perp} = 1 - \frac{\rho}{\rho_s} \frac{\omega_0^2 \tau^2}{1 + \omega_0^2 \tau^2} \tanh \left(\frac{\Delta}{2k_B T} \right),$$

