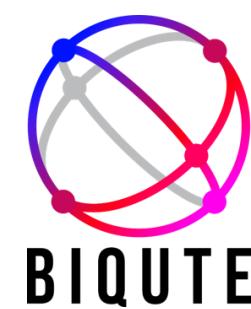




Istituto Nazionale di Fisica Nucleare



Design of an analog quantum simulator with superconducting qubits

Alessandro Cattaneo

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University of Milano-Bicocca

INFN – Milano-Bicocca

Bicocca Quantum Technologies (BiQuTe) Center

Outline

- Introduction
- Physical problem
- Quantum architecture
- Next steps

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□ Introduction <<<

□ Physical problem

□ Quantum architecture

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QUART&T project

QUantum Architectures for Theory & Technology (QUART&T)

Goal

Develop quantum architectures where theoretical models and phenomena of interest to the INFN can be tested.

Motivation

Simulate quantum many-body systems that have a physical related interest. Examples are nuclear reaction and dynamics, lattice quantum chromodynamics.

Approach

Realize an analog quantum simulator dedicated to a specific problem.



Istituto Nazionale di Fisica Nucleare

QUART&T collaboration

INFN Units:

- INFN Bologna
- INFN Ferrara
- INFN Firenze
- INFN Lecce
- INFN Milano
- **INFN Milano-Bicocca**
- INFN Gruppo Collegato di Salerno (INFN Napoli)

INFN National Laboratories:

- INFN Laboratori Nazionali di Frascati (LNF)
- INFN Laboratori Nazionali di Legnaro (LNL)

INFN Research Center:

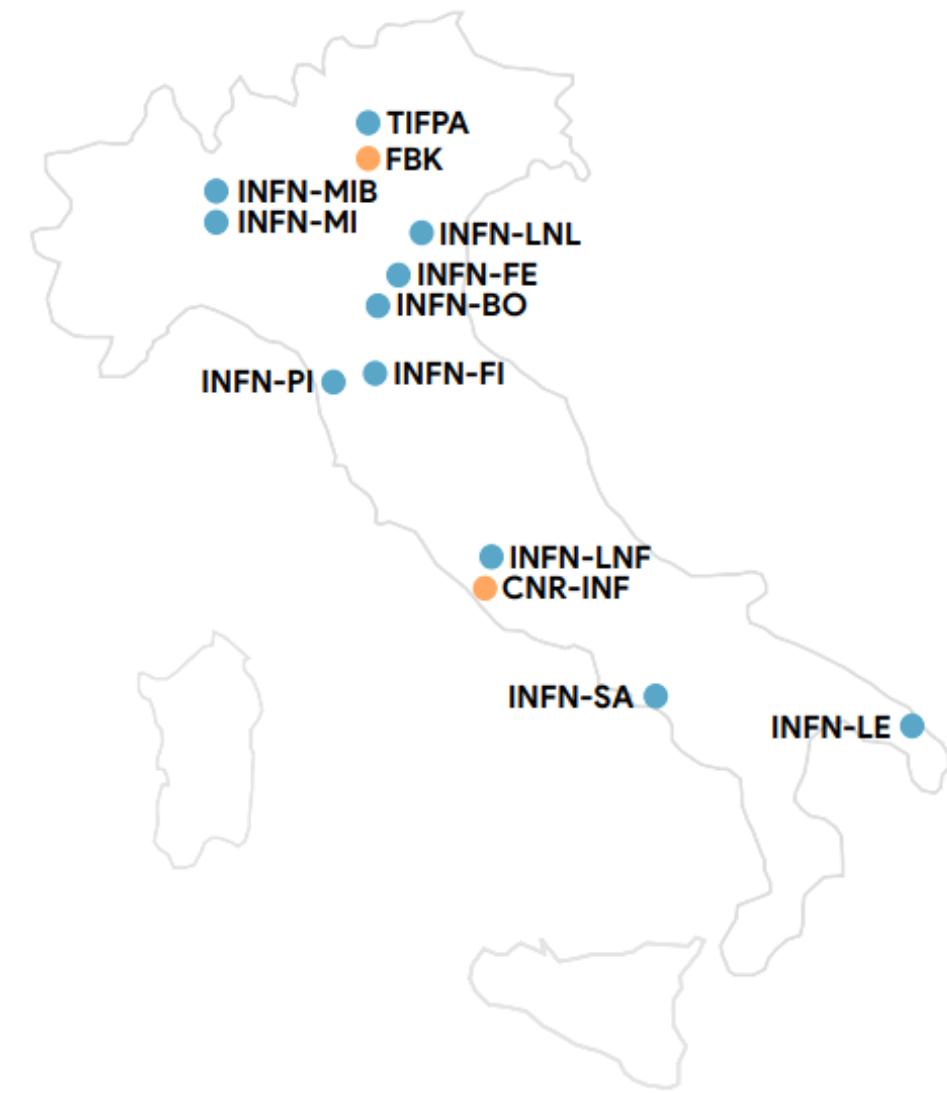
- Trento Institute for Fundamental Physics and Applications (TIFPA)

External Research Centers:

- Fondazione Bruno Kessler (FBK, Trento)
- Istituto di Fotonica e Nanotecnologie (CNR-INFN, Roma)



CSN5
Technological
research



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Neutrinos oscillations in dense environments

- In high-density neutrino environments — such as supernovae, neutron star mergers, or the early universe — neutrino flavor oscillations are affected by neutrino-neutrino interactions.



The Crab Nebula, a remnant of a supernovae.

Neutrinos oscillations in dense environments

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- The Hamiltonian for a **two flavors oscillation** can be written as:

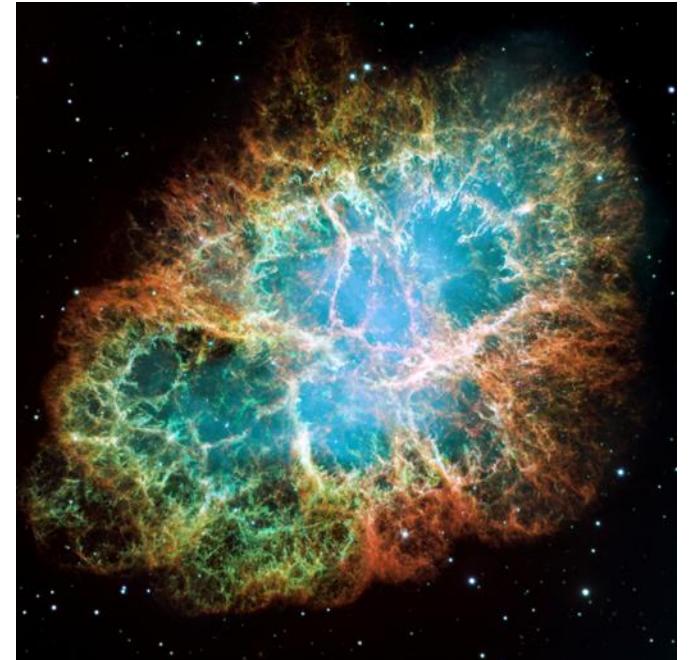
$$H = \sum_{k=1}^N \vec{b} \cdot \vec{\sigma}_k + \sum_{p < q}^N J_{pq} \vec{\sigma}_p \cdot \vec{\sigma}_q \quad \text{with } \vec{\sigma}_k = (\sigma_k^x, \sigma_k^y, \sigma_k^z)$$



1. Vacuum mixing
2. MSW effect



3. Neutrino-neutrino interaction



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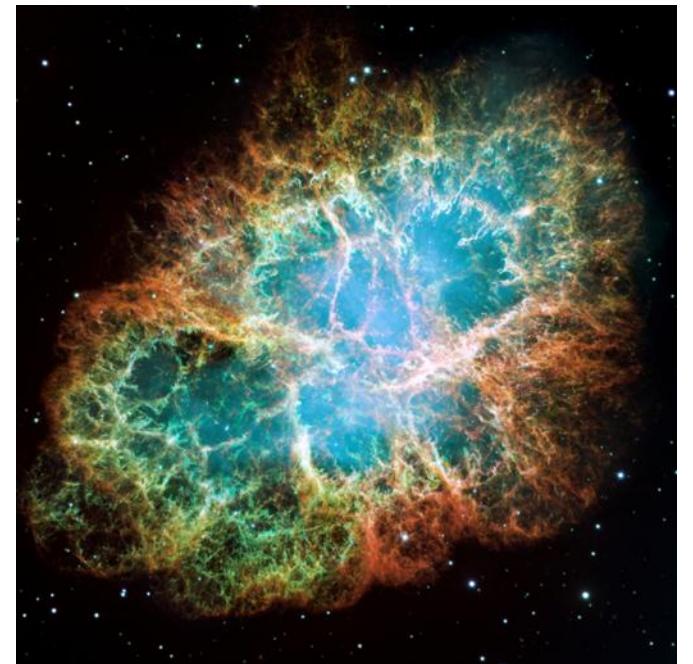


1. Vacuum mixing
2. MSW effect



3. Neutrino-neutrino interaction

- An **exact solution** is necessary to understand the role of quantum correlations.



The Crab Nebula, a remnant of a supernovae.

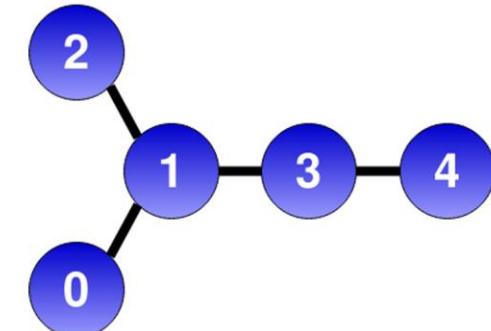
Quantum digital simulations

[[Phys.Rev.D 104 \(2021\) 6, 063009](#)]

- System complexity on a classical computer scales as $\sim e^N$. On a quantum machine $\sim N^k$.

$$e \longrightarrow |0\rangle$$

$$\tau \longrightarrow |1\rangle$$



IBM Quantum Canary Processor Vigo

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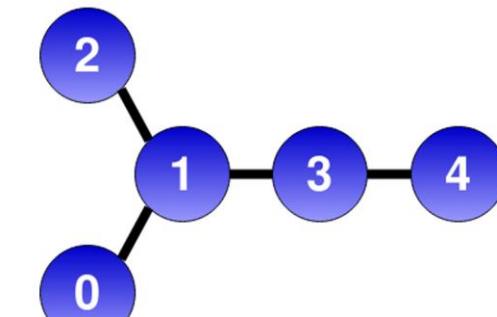
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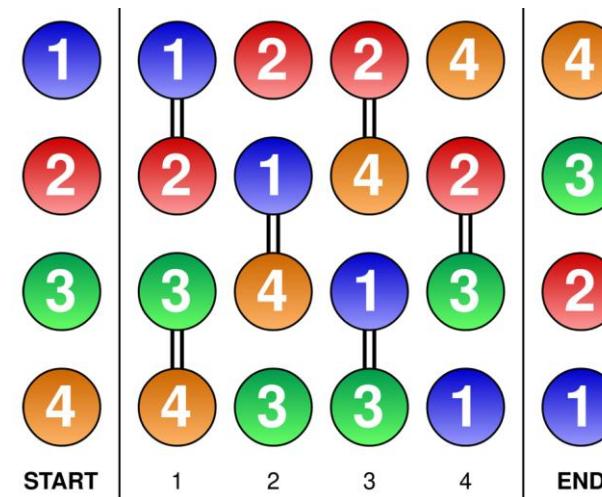
- **First-order Trotter-Suzuki decomposition:**

$$U(t) = \prod_{j=1}^N e^{-it\vec{b}\cdot\vec{\sigma}_j} \prod_{p < q} e^{-itJ_{pq}\vec{\sigma}_p\cdot\vec{\sigma}_q}$$

- **All-to-all connectivity** realized using SWAP gates.



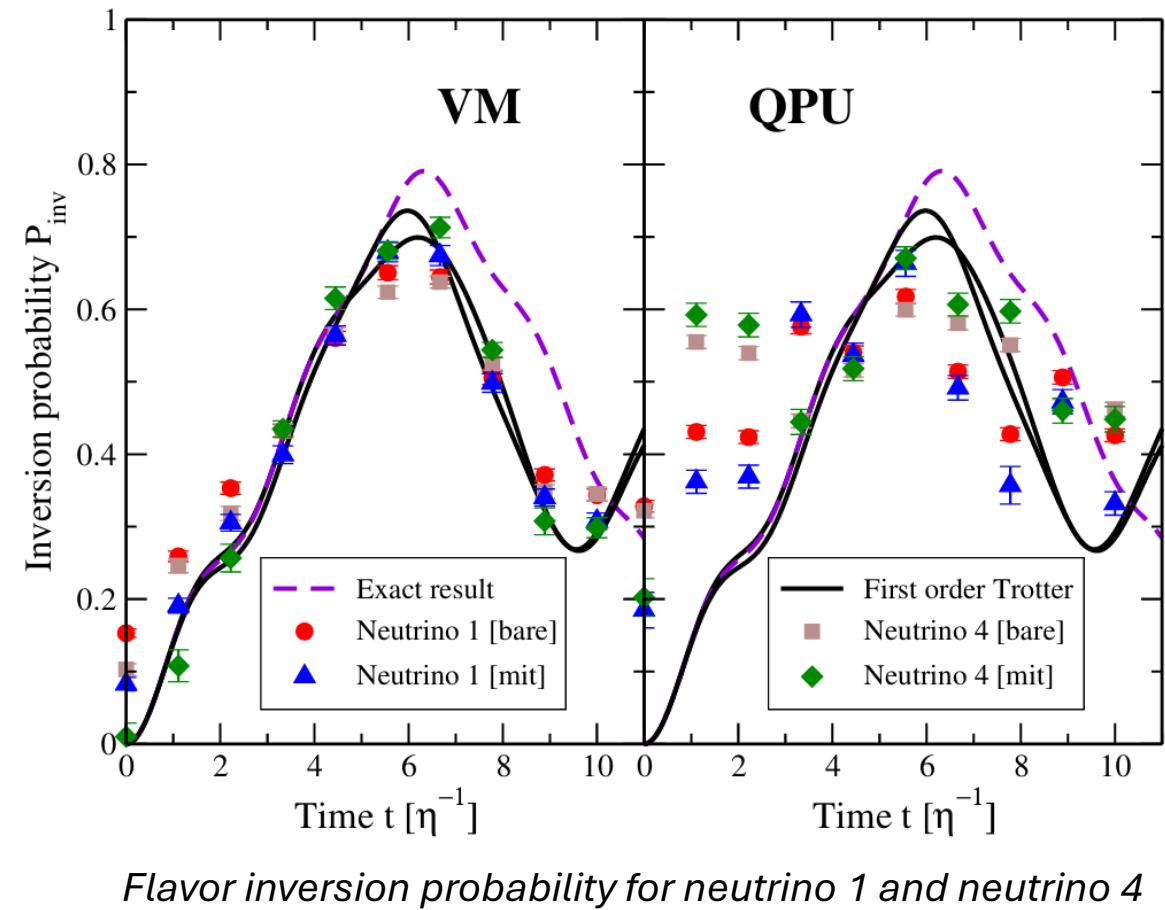
IBM Quantum Canary Processor Vigo



Scheme to realize a single time-step simulation.
Each double line is formed by $3 \binom{N}{2}$ CNOT + $15 \binom{N}{2}$ single qubit gates.

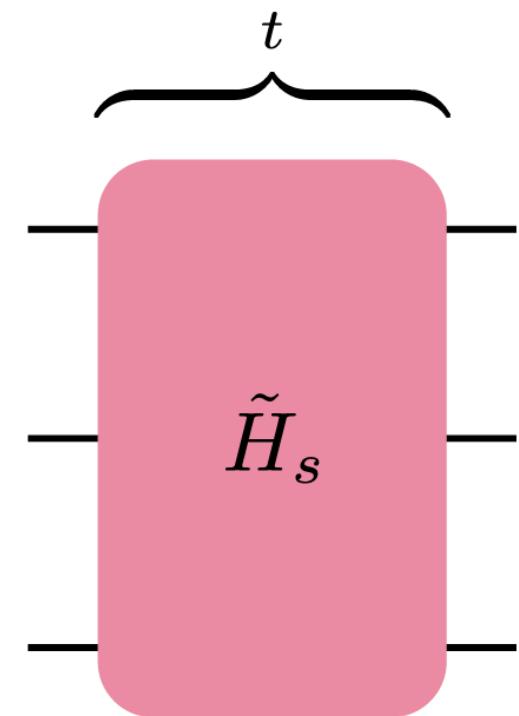
Quantum digital simulation - Limitations

- The **first-order Trotter-Suzuki** decomposition introduce an error of order $O(t^2)$ on the simulation results.
- **N gates** are needed for each simulation step t. This require long run time to reach higher simulated times.
- **Single qubit noise** limits both results accuracy and simulation length.
- **SWAP gate errors** increase as single qubit errors occurs.



Analog approach

- **Analog quantum simulator:** a quantum device whose time evolution emulates the dynamics of a more complex quantum model.

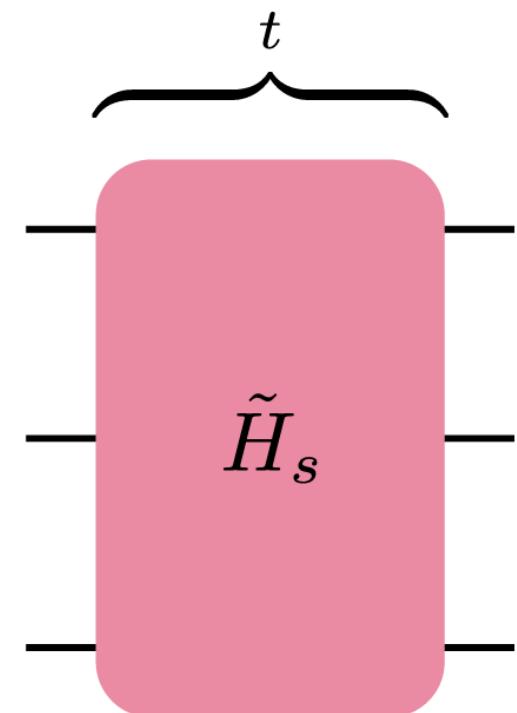


Analog approach

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Main advantages:

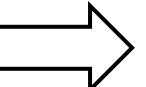
- Each simulated time step require less operations.
- No approximation needed for the time evolution operator.

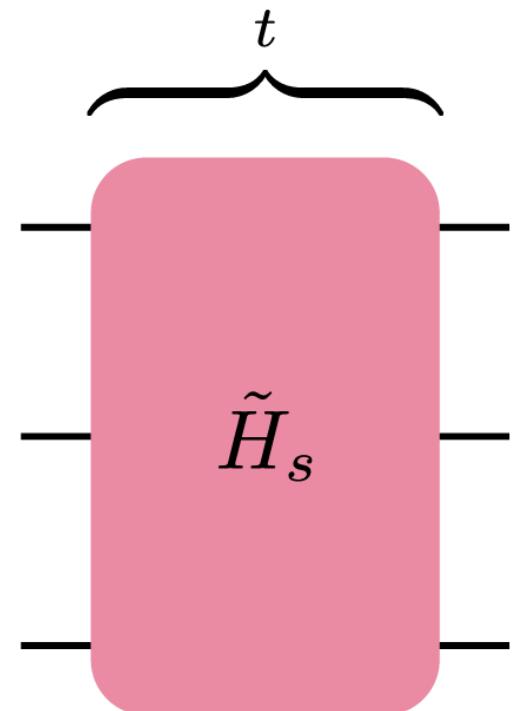


Analog approach

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Main advantages:

- Each simulated time step require less operations.
- No approximation needed for the time evolution operator.
- **2D superconducting platform** 
 - Highly engineerable
 - Parameter tunability



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Hamiltonian

$$H = \sum_{k=1}^N \vec{b} \cdot \vec{\sigma}_k + \sum_{p < q} J_{pq} \vec{\sigma}_p \cdot \vec{\sigma}_q$$

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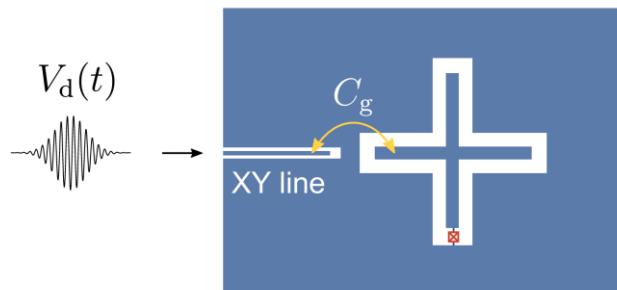
Single qubit gates:

- External drive on single qubits



$$H \sim \Omega V_d(t) (I\sigma_x + Q\sigma_y) \quad I = \cos \phi, Q = \sin \phi$$

- Continuous virtual Z gate $\rightarrow H \sim \frac{\delta\omega}{2}\sigma_z$



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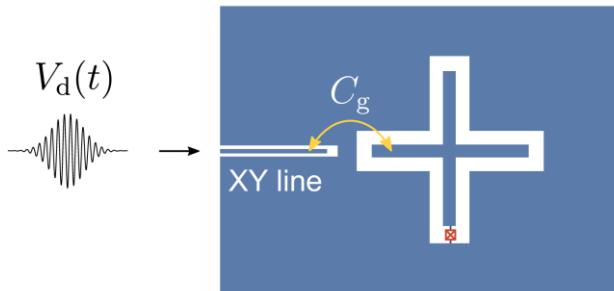
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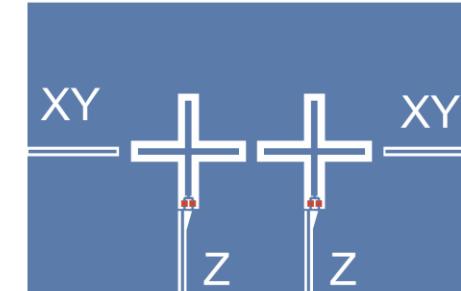
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- Continuous virtual Z gate → $H \sim \frac{\delta\omega}{2} \sigma_z$



Two qubit interactions:

- Capacitive coupling → $H \sim \frac{g}{2} (\sigma_x \sigma_x + \sigma_y \sigma_y)$
- Higher states energy shift → $H \sim \frac{\zeta}{4} (\sigma_z \sigma_z)$
- Parametric couplings



Two qubit interaction – Capacitive coupling

- Direct capacitive coupling between transmon qubits is described by:

$$H = -g([\sigma^+ - \sigma^-] \otimes [\sigma^+ - \sigma^-])$$

- Under **Rotating Wave Approximation** (valid if $\omega_{q1}, \omega_{q2} \gg g$) fast oscillating terms disappear leaving:

$$H = g(\sigma^+ \sigma^- + \sigma^- \sigma^+) = \boxed{\frac{g}{2}(\sigma_x \sigma_x + \sigma_y \sigma_y)}$$

Interaction we are looking for!

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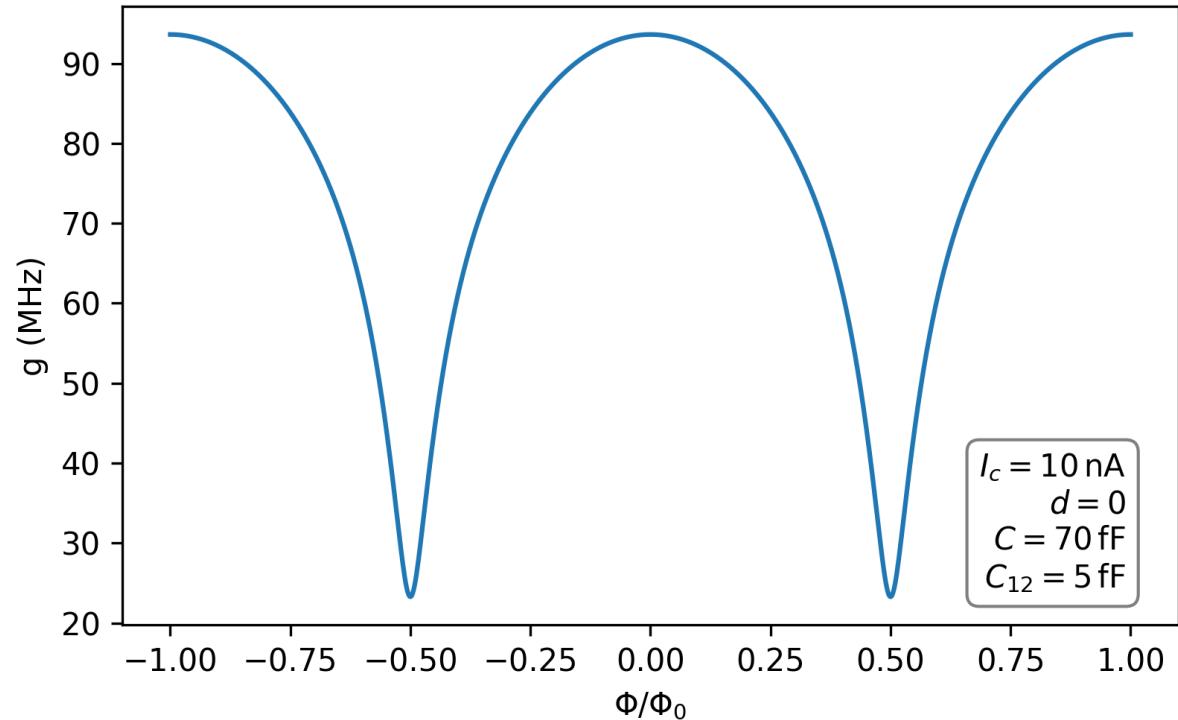
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Interaction we are looking for!

- g has a **qubit frequency dependency**:

$$g \sim \frac{1}{2} \frac{e^2}{\hbar} \frac{C_{12}}{\sqrt{C_1 C_2}} \sqrt{\omega_1 \omega_2}$$

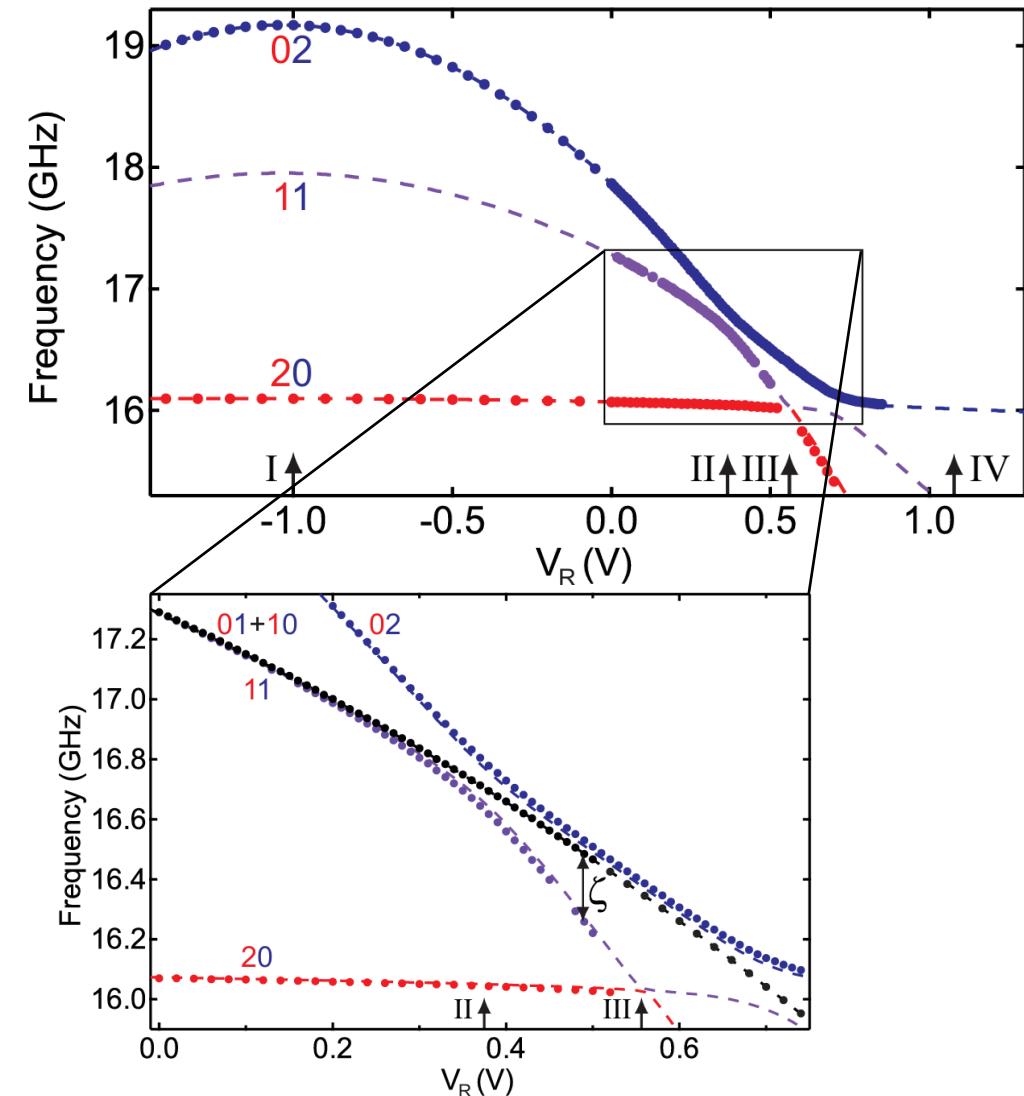


Two qubit interaction - ZZ coupling

[Nature 460, 240–244 (2009)]

- Idea: use the **energy repulsion** between $|20\rangle$, $|02\rangle$ and $|11\rangle$ to realize an effective ZZ interaction on the computational subspace;

$$\zeta = E_{|11\rangle} - E_{|01\rangle} - E_{|10\rangle}$$



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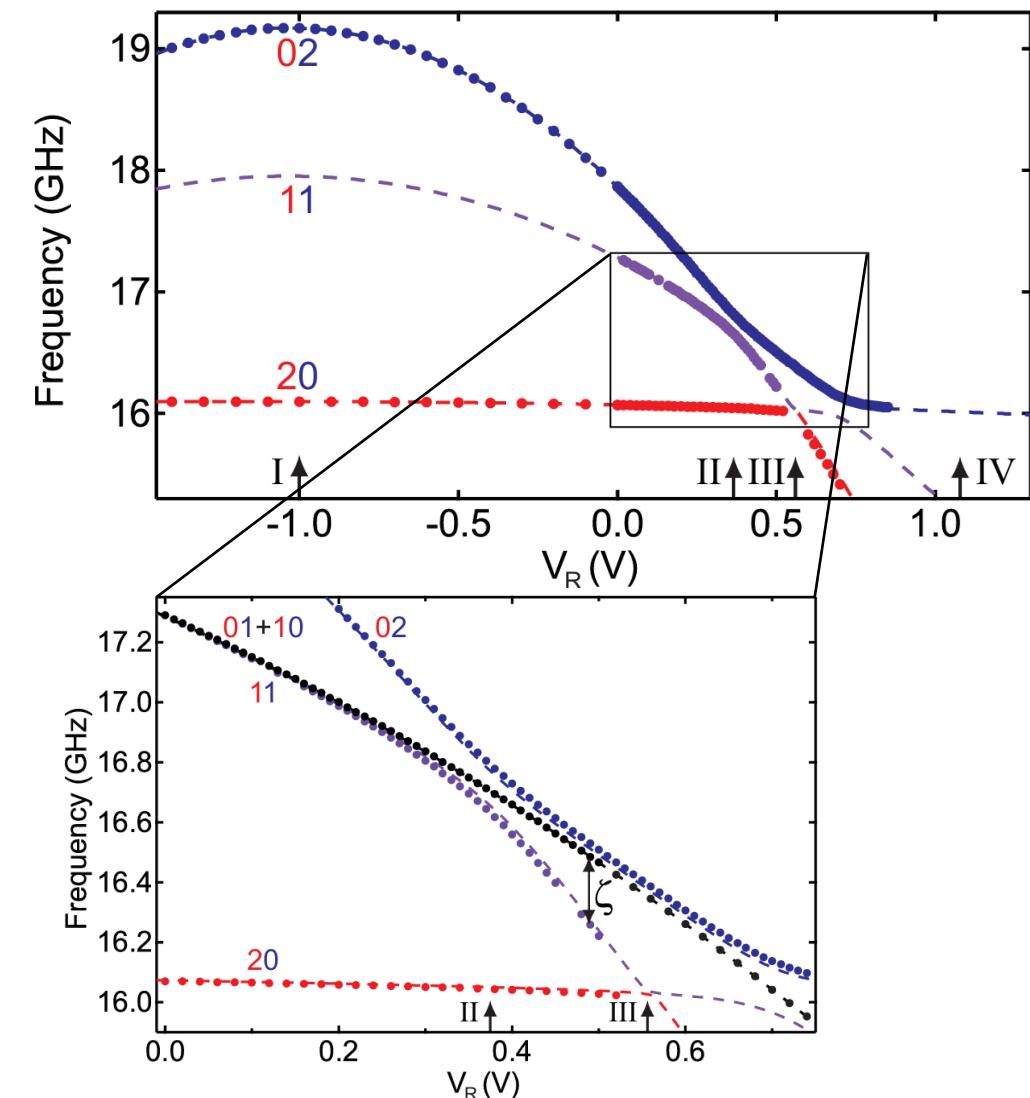
$$\zeta = E_{|11\rangle} - E_{|01\rangle} - E_{|10\rangle}$$

- The interaction hamiltonian on the computational subspace is:

$$H \sim \frac{\zeta}{4} \sigma_z \otimes \sigma_z - A \sigma_z \otimes \mathbb{I} - B \mathbb{I} \otimes \sigma_z$$

Interaction we are looking for! Can be used for single qubit gates!

- Same type of interaction used in digital architectures to realize a CPHASE gate.



Two qubit interaction – ZZ coupling

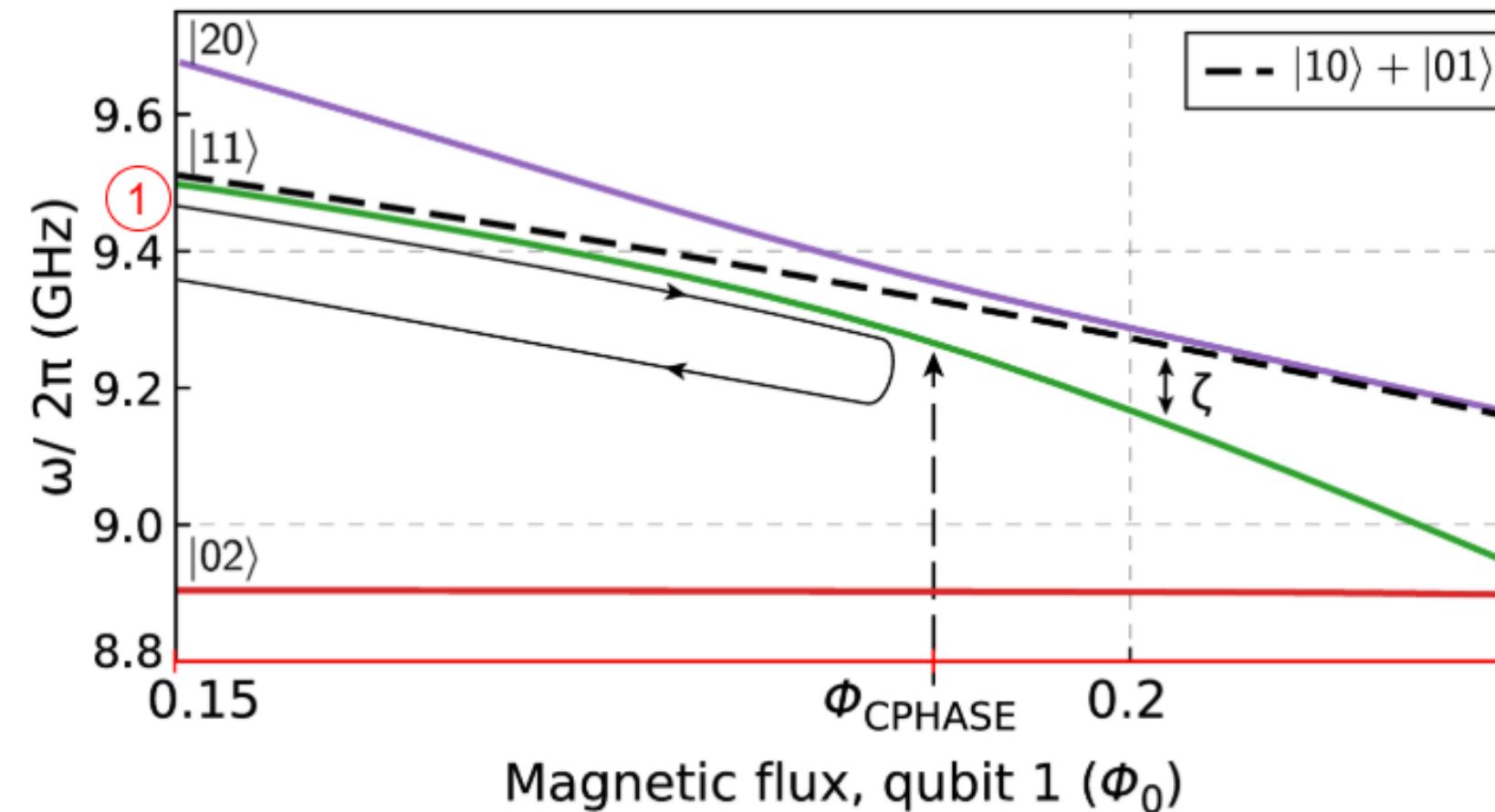
How to properly create the ZZ effective interaction?

[\[Appl. Phys. Rev. 1 2019\]](#)

Two qubit interaction - ZZ coupling

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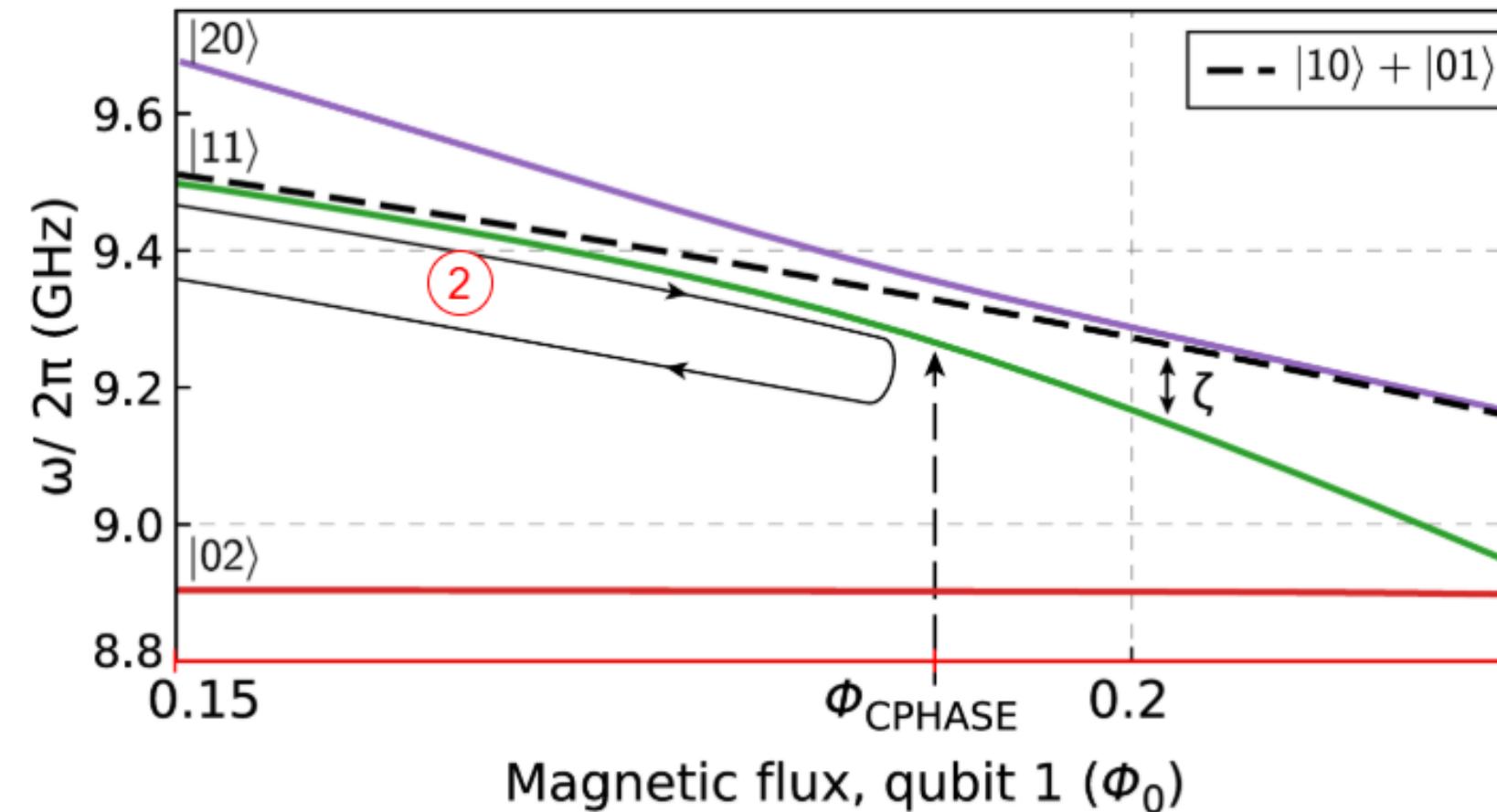
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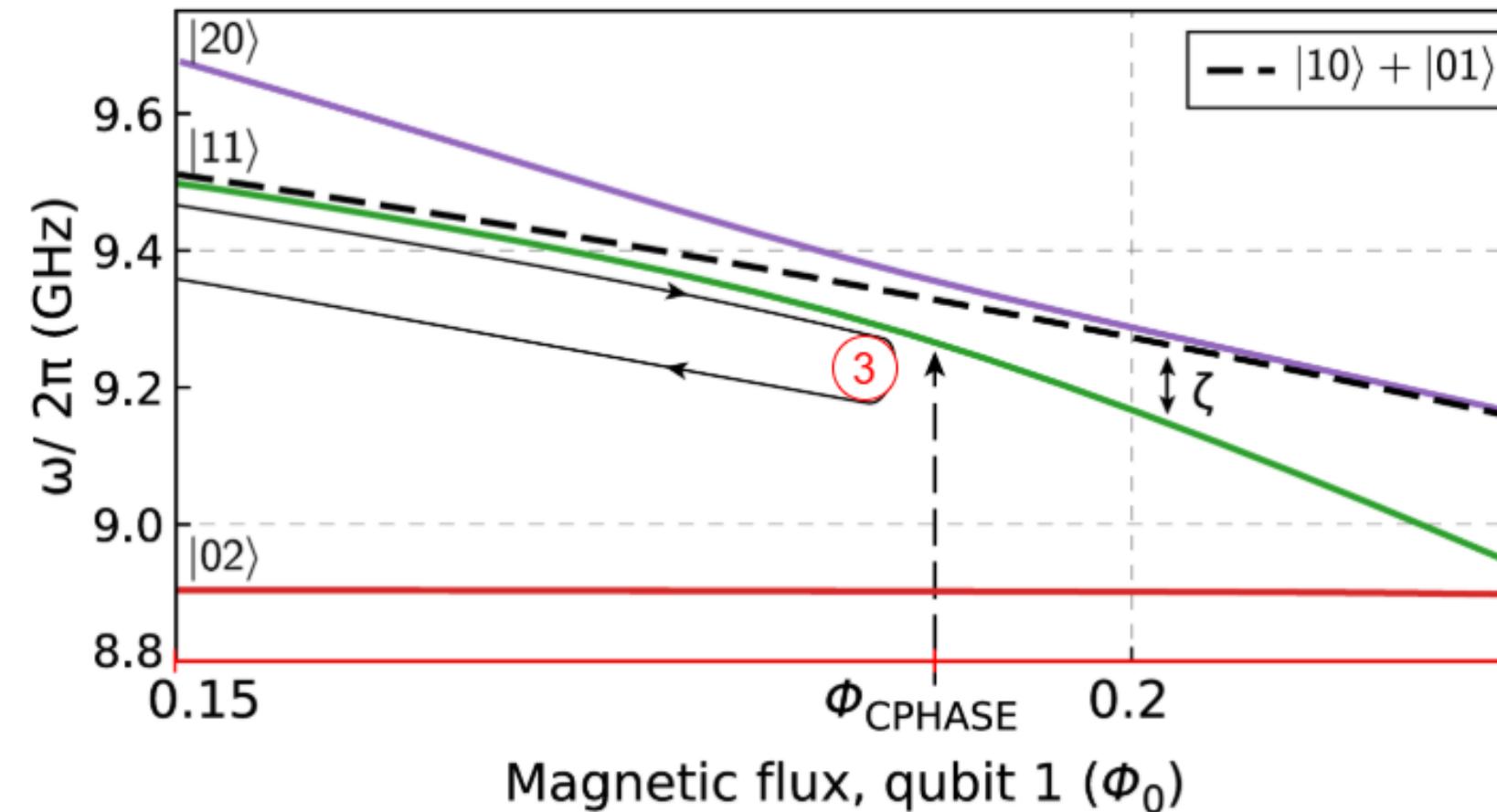
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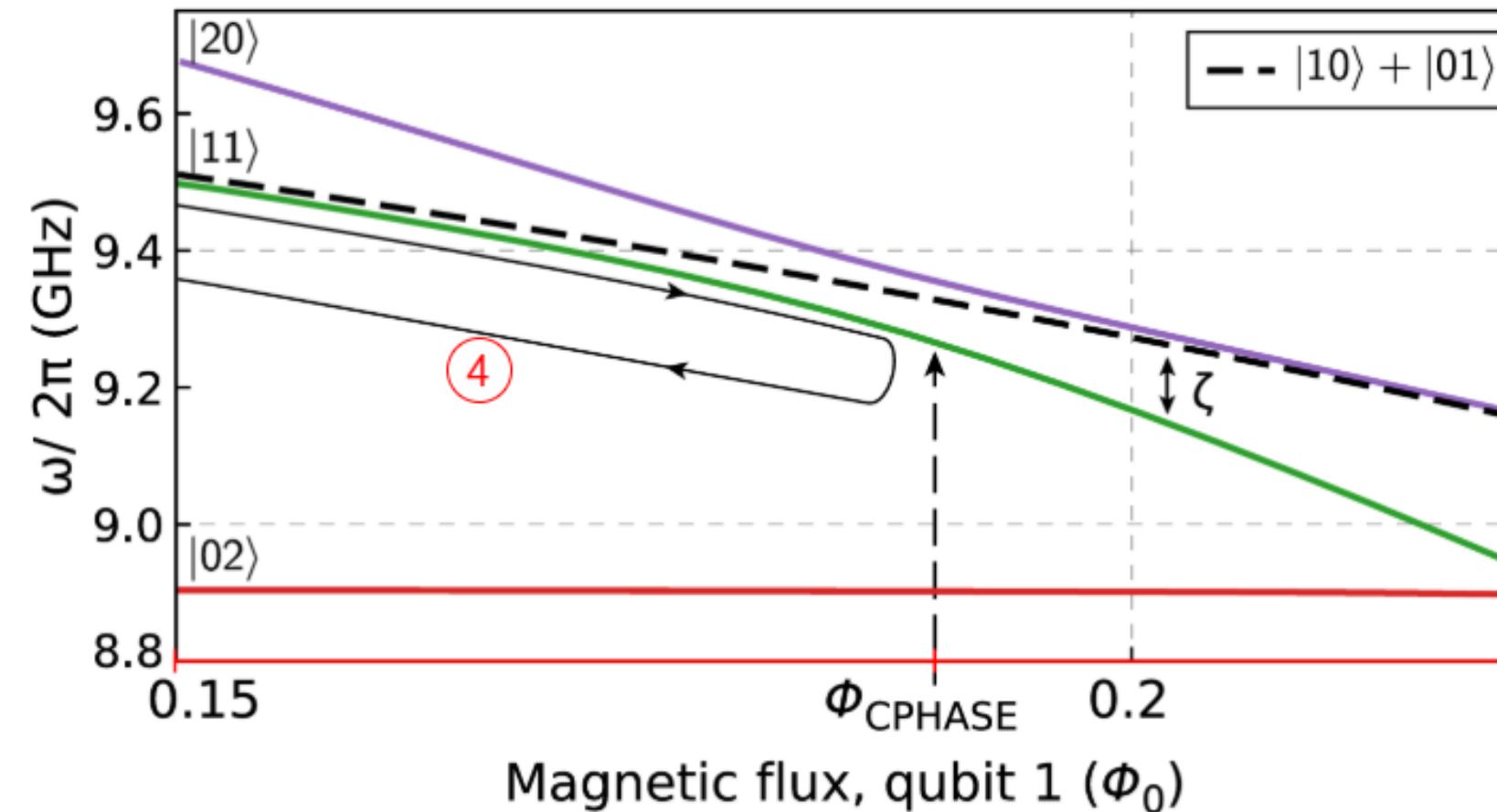
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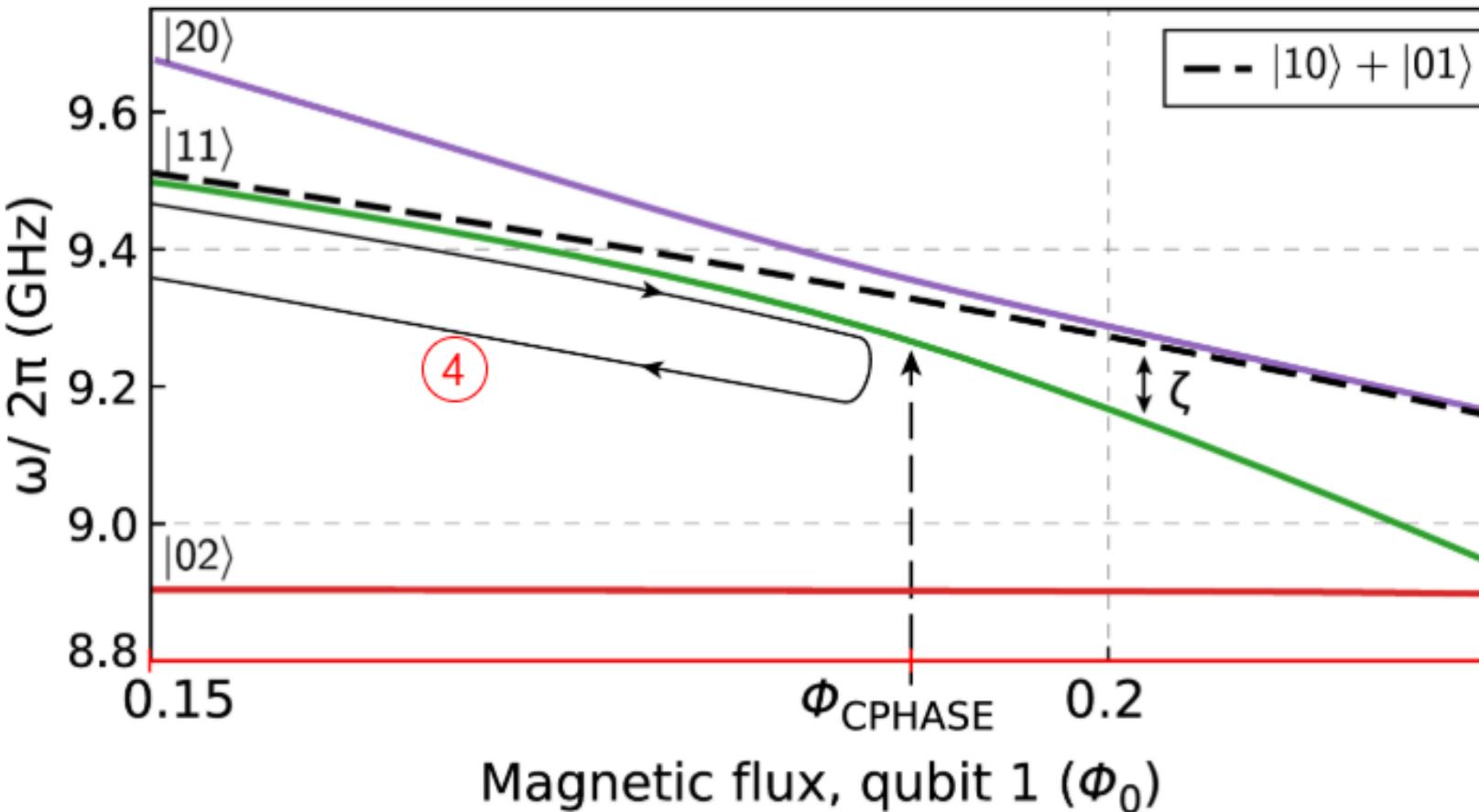
[Appl. Phys. Rev. 1 2019]



Two qubit interaction - ZZ coupling

How to properly create the ZZ effective interaction?

[Appl. Phys. Rev. 1 2019]



Adiabatic motion is fundamental to **always stay in one dressed state.**



Prevent the system from evolving into $|20\rangle$.

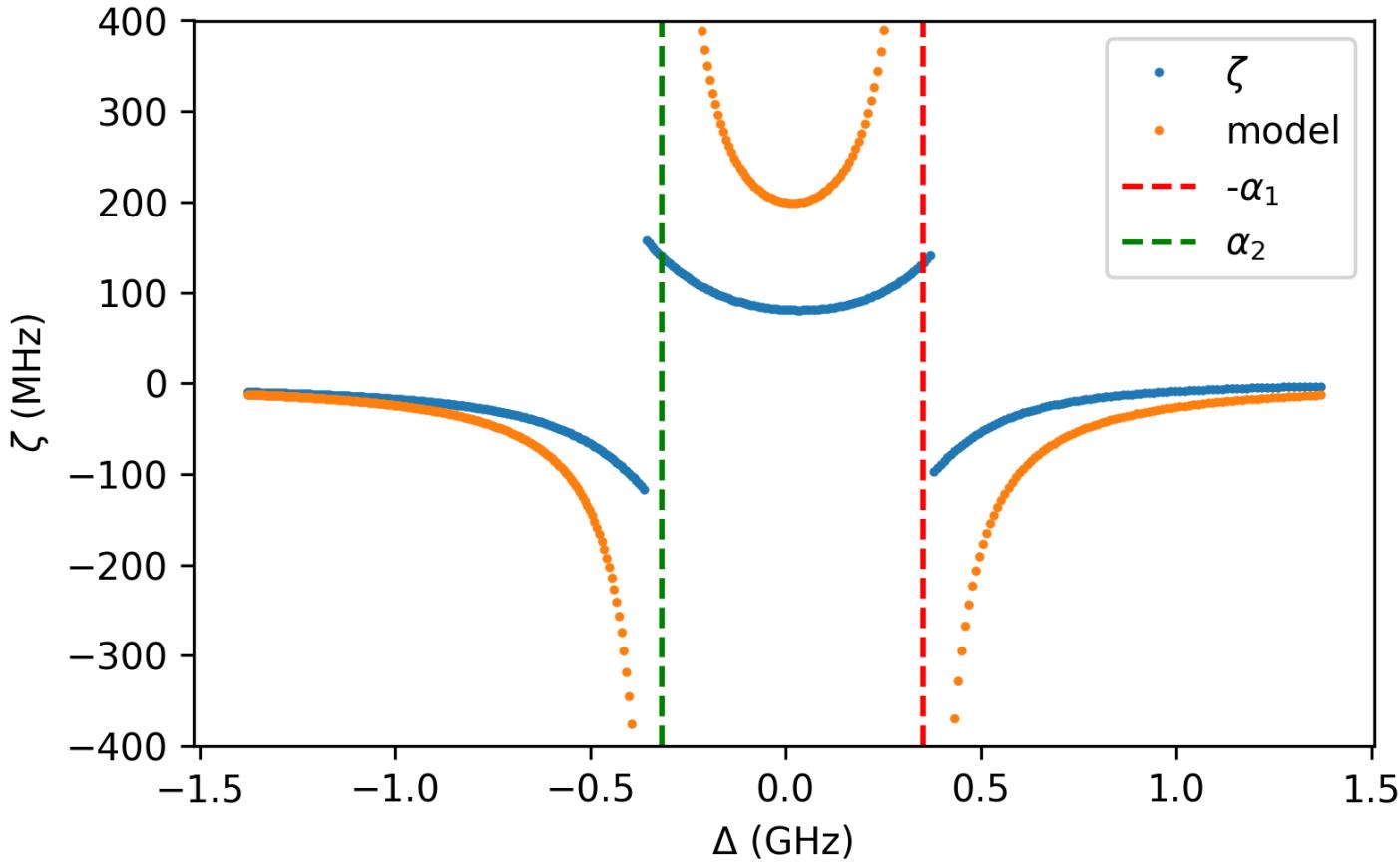
Adiabatic condition: $\left| \frac{d\omega}{dt} \right| \ll 8g^2$

Two qubit interaction - ZZ coupling strength

- There exist an analytical model valid in the dispersive regime ($|\omega_1 - \omega_2| \gg g$):

$$\zeta \approx 2g^2 \left(\frac{1}{\Delta - \alpha_2} - \frac{1}{\Delta + \alpha_1} \right)$$

- **Numerical approximations** are needed for resonance regimes.

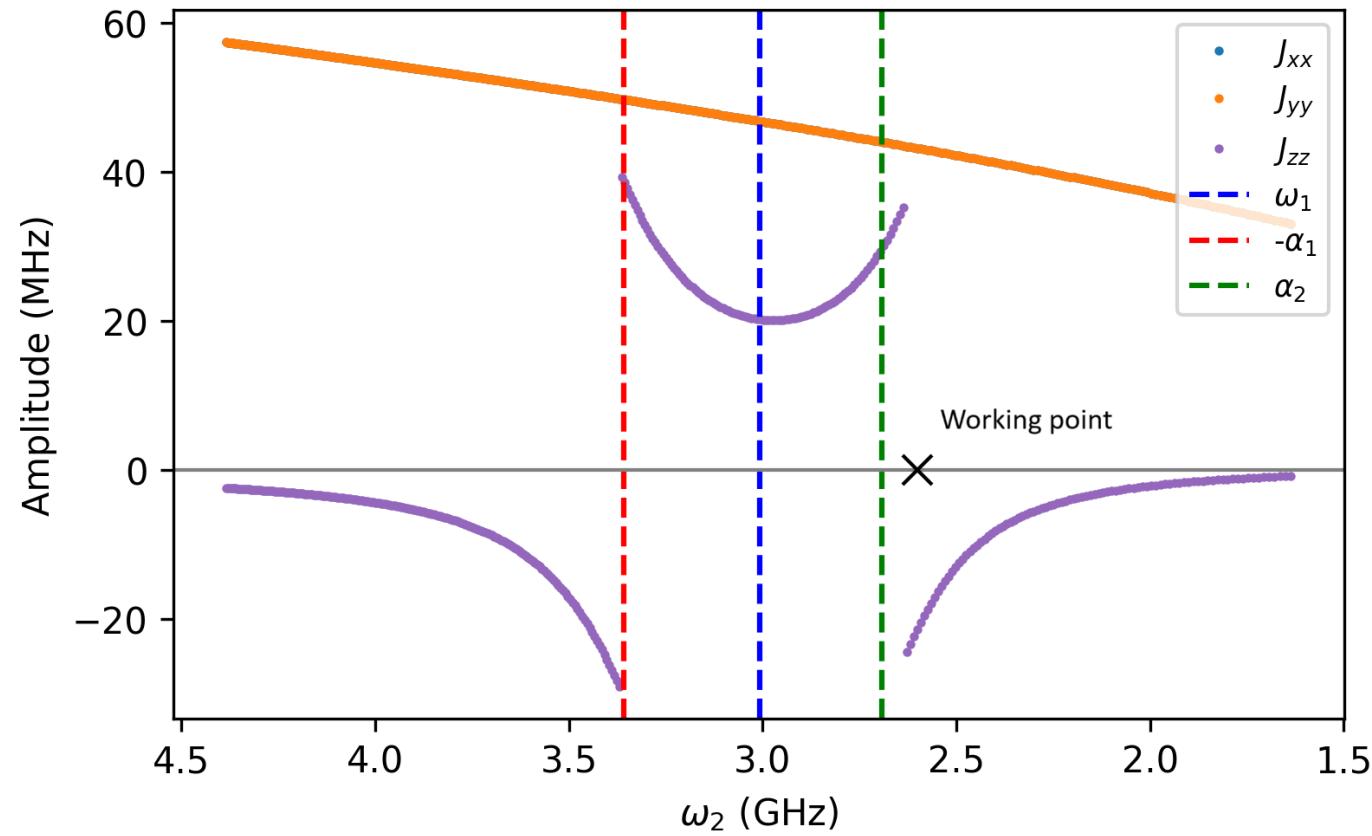


Two qubit interaction – Strength comparison

- At the working point - coupling strength parameters:

$$J_{xx} = 43 \text{ MHz} \quad J_{yy} = 43 \text{ MHz} \quad J_{zz} = -21 \text{ MHz}$$

$$\Delta_{J_{xx}} \sim 20 \text{ MHz} \quad \Delta_{J_{zz}} \sim 65 \text{ MHz}$$



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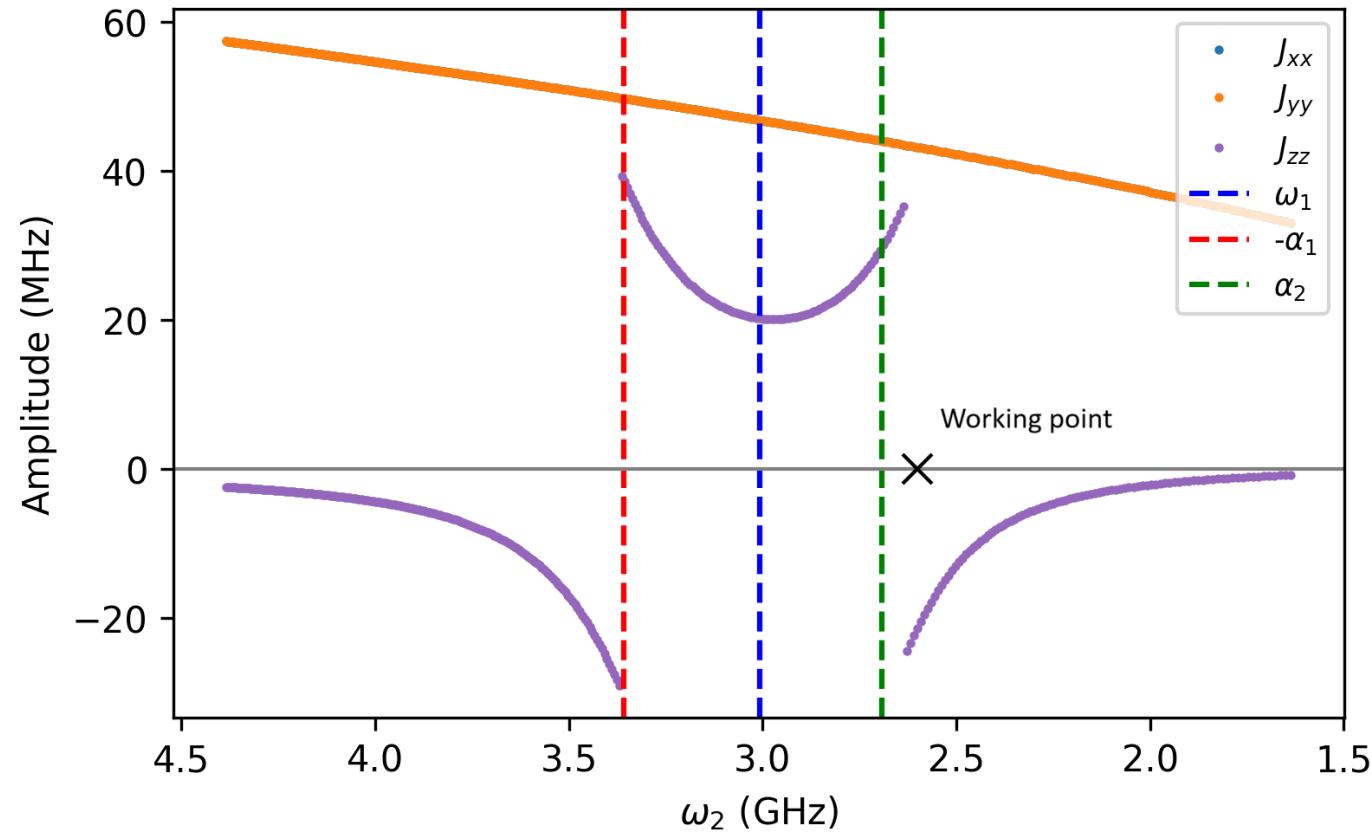
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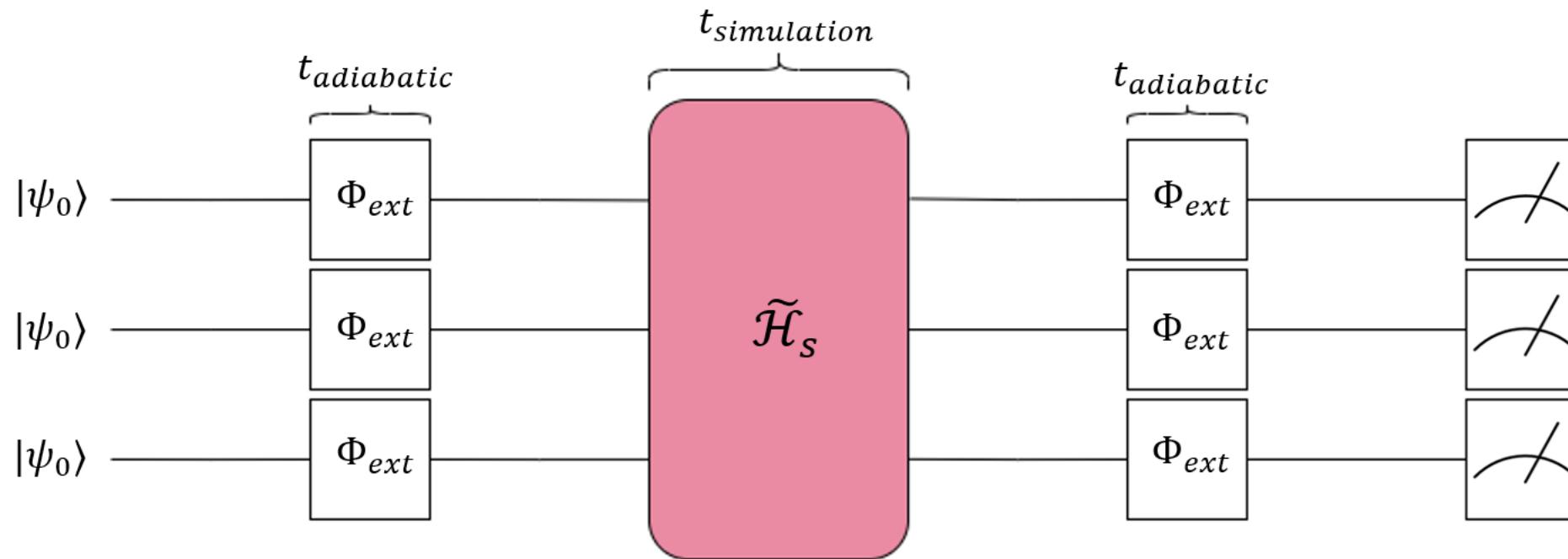
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Future developments

- Engineering design parameters to obtain **coupling strength with the same order of magnitude.**
- Interaction between more qubits.



How to use – Simulation setup



- 1) Prepare the initial state in a idling point
- 2) Change qubit frequencies to reach the working point
- 3) Let the system evolve
- 4) Goes back to the idling point
- 5) Measure

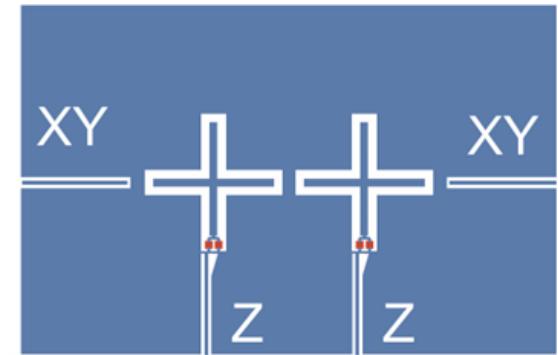
- Process the data considering the **unwanted time evolution** caused by moving between idling and working points.
- Single qubit rotations can be turned on sending electromagnetic pulses during H_S time evolution.

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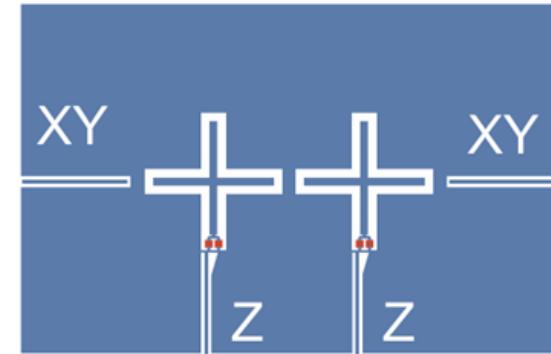
Scaling to more qubits

- **Directly capacitive coupling:**
 - Limited to a low number.
 - Difficult to design precise coupling capacitances C_{ij} .

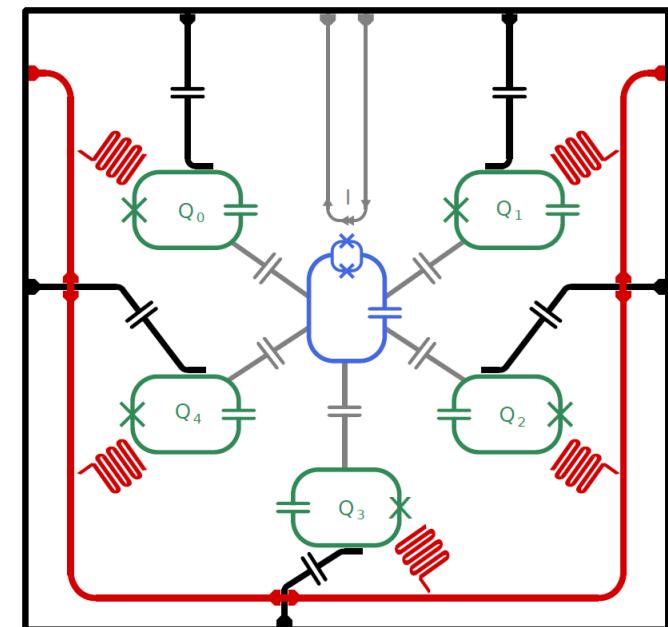


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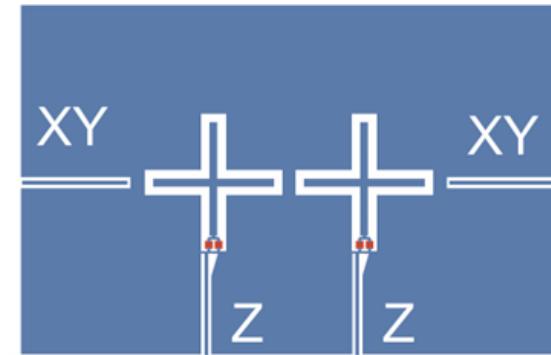


- Use one «central» resonator to connect many qubits:
 - Hamiltonian interaction terms do not connect directly data qubits.
 - Frequency overcrowding in the resonator.

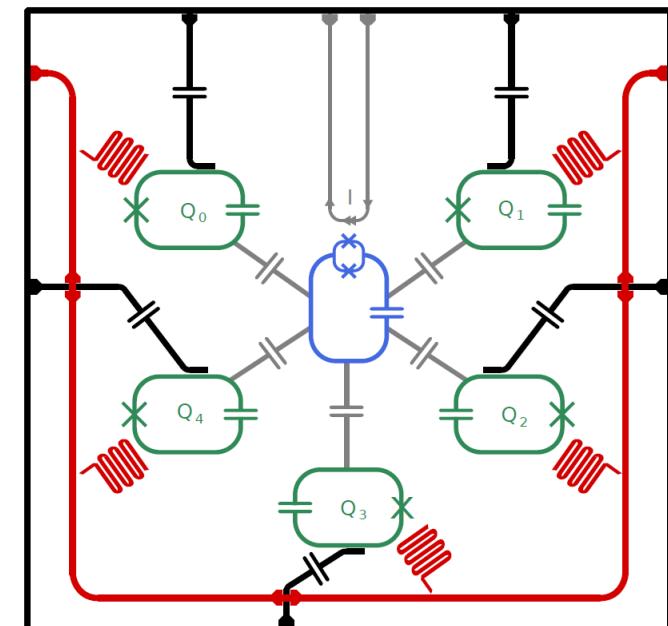


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- Use one «central» resonator to connect many qubits:
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 - Frequency overcrowding in the resonator.
- The presence of many qubits can alter the energy spectrum
→ **change the influence from higher states** → harder to obtain ZZ interaction.



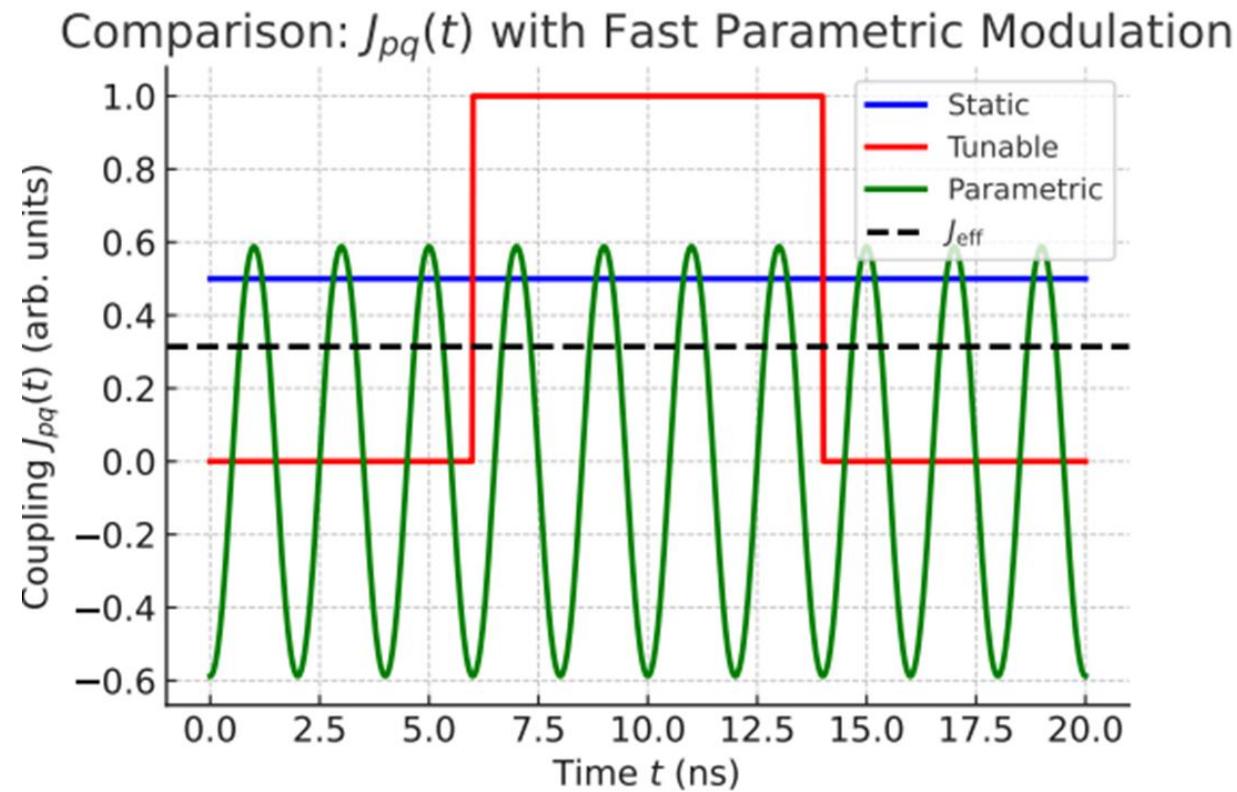
Different approaches – Parametric coupling

- **Key idea:** using an **AC pulse** to modulate the qubit frequency. By rapidly varying the frequency qubit 1 at the detuning $\Delta = \omega_1 - \omega_2$ we can activate a SWAP interaction with qubit 2.

- Time-dependent coupling:

$$H_{int} \sim J(t) (\sigma^- \sigma^+ + \sigma^+ \sigma^-)$$

- ZZ coupling can be activated similarly.
Exploiting higher energy levels.



[10.48550/arXiv.2305.02907]

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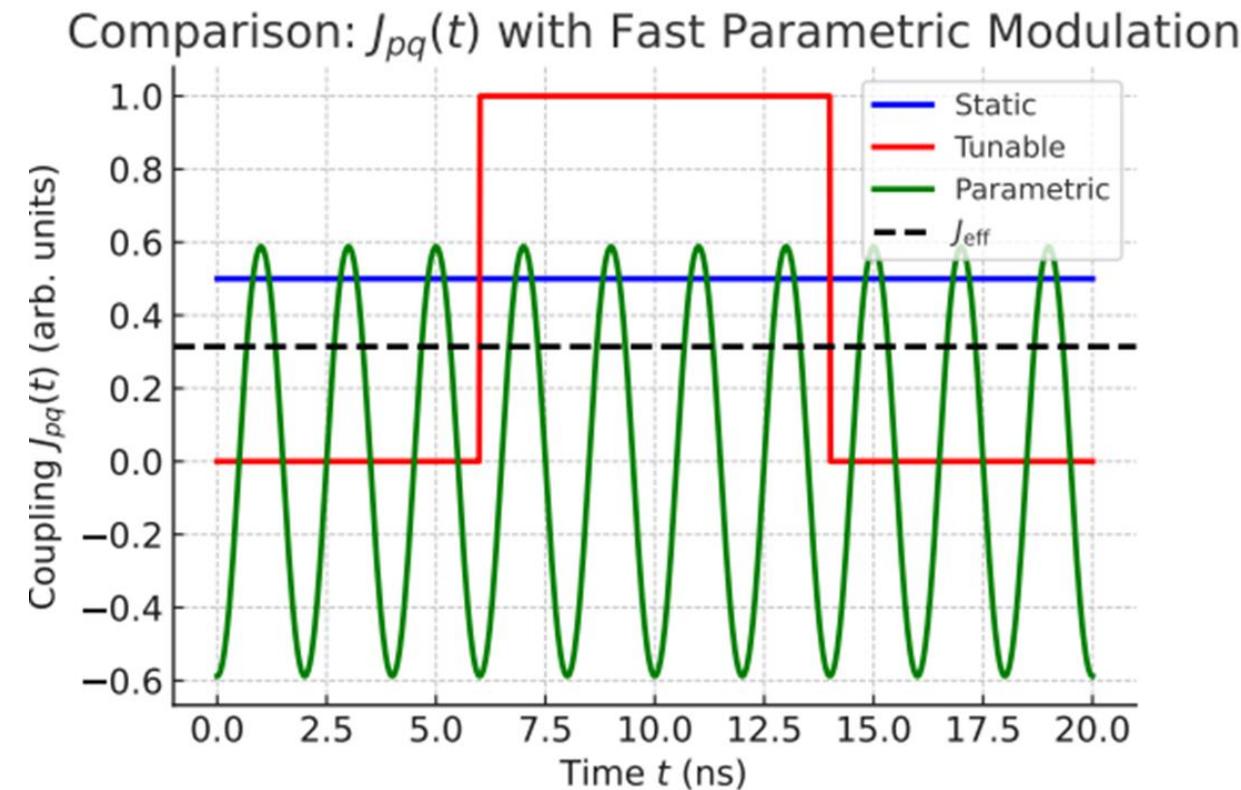
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Advantages:

- High control over coupling strength amplitude.
- Qubit frequency independent strength → no frequency overcrowding.



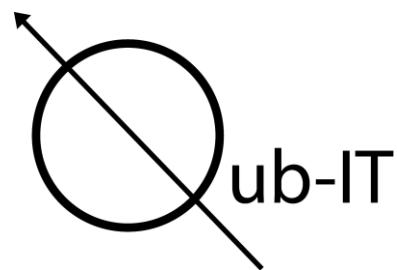
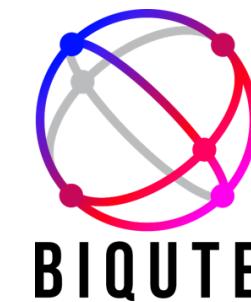
[10.48550/arXiv.2305.02907]

Thank you for your attention!

UNIMIB/INFN-MIB group: Matteo Borghesi, Pietro Campana, Lucia Canonica, Rodolfo Carobene, Alessandro Cattaneo, Hervè Corti, Marco Faverzani, Elena Ferri, Sara Gamba, Andrea Giachero, Marco Gobbo, Danilo Labranca, Roberto Moretti, Angelo Nucciotti, Luca Origo



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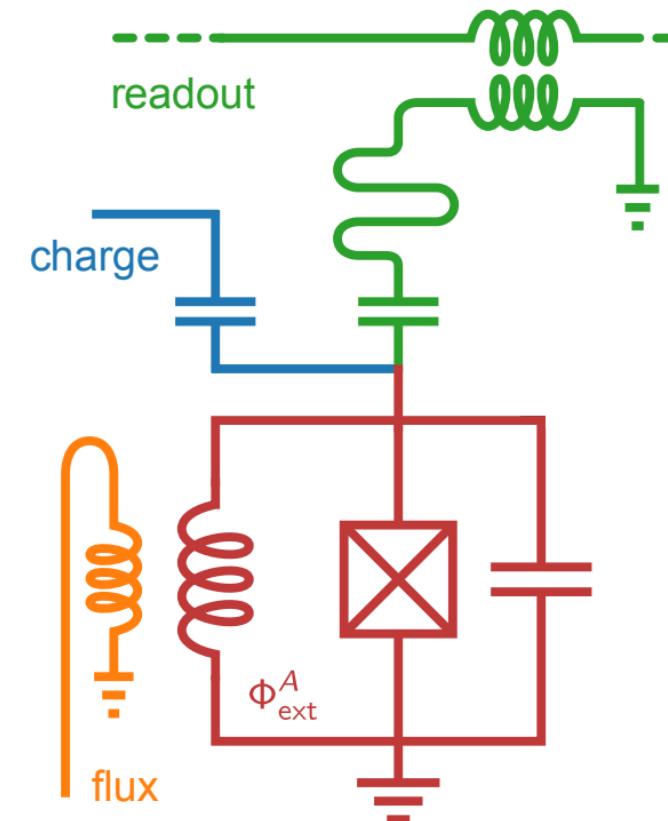


Different approaches – Fluxonium

- The **fluxonium** is superconducting qubit which offers an **high control on anharmonicity**. Typical values are between $0.5 - 1.5 \text{ GHz}$ and can be both positive and negative.
- High anharmonicity \rightarrow High ZZ interaction.
- It is composed of a Josephson Junction, a large inductance (array of JJ), and a capacitor.

Complications:

- Fabrication becomes more complicated;
- Higher flux sensitivity;
- Require more space \rightarrow harder to scale;
- Read out is more complicated.



[PhysRevLett.129.010502]