# Gravitational Form Factors through the application of novel computational techniques

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Mechanical properties of hadrons: Structure, dynamics, visualization ECT\* Trento, Italy March 3th, 2025





### Outline

I. Motivation and scope of the project:

- Mechanical properties of the proton.
- Accessing GFFs through 3D distributions: Compton form factors.
- 2. Framework:
  - from LQCD.
  - Technical difficulties of using NNets.
- 3. State of the project.
- 4. Deliverables and future directions.

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• Extraction of Gravitational Form Factors using Neural Networks, from experimental data with insights



# Motivation and scope of the project: Mechanical properties of the proton. Extraction of Gravitational Form Factors using Neural Networks, from experimental data with insights from LQCD.







### Mechanical properties of the proton

What are these mechanical properties?

- Internal distribution of forces and pressure.
- Distribution of mass and energy.
- Distribution of angular momentum.
- Physical size as given by the mechanical radius.

There is no direct way to measure the gravitational form factors (GFFs). However, they may be probed indirectly in various exclusive processes like DVCS that simulate the graviton-proton interaction.

The  $2\gamma$  field couples to the EMT as gravity does with a strength many orders of magnitude greater.

The leading contribution to DVCS is described in terms of four generalized parton distributions (GPDs). Two of them give access to the quark GFFs. Therefore it also let to the insight that GPDs can be related to the mechanical properties of nucleons.





$$\xi,t) = A_q(t) + \xi^2 D_q(t),$$

$$x,\xi,t) = B_q(t) - \xi^2 D_q(t),$$







### Mechanical properties of the proton

The Matrix element of the QCD Energy-Momentum tensor encodes key information of the proton including the mass and spin, as the distributions of energy, angular momentum, and various mechanical properties such as, e.g., internal forces inside the system. Momentum



On top of restricting the number of GFFs, Poincaré symmetry imposes additional constraints:

$A(0) = \sum_{q} A_q(0) + A_G(0) = 1,$ $J(0) = \sum_{q} J_q(0) + J_G(0) = \frac{1}{2},$	D(t) is un be deterr
$\frac{1}{2}\Delta\Sigma = \sum_{q}^{q} S_{q}(0),$	D(t)
$\bar{C}(t) = \sum_{q}^{1} \bar{C}_{q}(t) + \bar{C}_{G}(t) = 0,$	The D-te

V. D. Burkert, et. al. Rev. Mod. Phys. 95, 041002

rm is not related to "external properties" of a particle like mass and spin, but to the stress tensor and internal force.

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$$p, \vec{s} \rangle = \overline{u}(p', \vec{s}') \left[ A_a(t) \frac{P^{\mu}P^{\nu}}{M_N} \right]$$

$$\frac{\Delta^{\nu} - g^{\mu\nu}\Delta^2}{4M_N} + \overline{C}_a(t) M_N g^{\mu\nu}$$

$$\frac{\Delta^{\nu} - g^{\mu\nu}\Delta_\lambda}{M_N} + \overline{C}_a(t) \frac{P^{[\mu}i\sigma^{\nu]\lambda}\Delta_\lambda}{M_N} \right] u(p, \vec{s})$$
Pressure

Total angular momentum

Intrinsic angular

Symmetric kinematical variables: P = (p' + p)/2 and  $\Delta = p' - p$ Free Dirac spinor:  $u(p, \vec{s})$ GFFs:  $A_a(t), D_a(t), \overline{C}_a(t), \text{ and } J_a(t)$ 

nknown. It value is not fixed by spacetime symmetries and must mined from experiment.

$$b \qquad \Delta(t) \propto D(t) \propto \int d^3 \mathbf{r} \ p(r) \frac{j_0(r\sqrt{-t})}{t}$$



= ?

### Our goal:

Make a global analysis using NN

Extract GFFs

• Use state-of-the-art ML libraries: Pytorch.

• Handle experimental and LQCD data.





### Accessing GFFs through 3D distributions: Compton form factors

Complications: The actual observables in DVCS are the Compton form factors (CFFs) which are expressed by means of factorization formulae in terms of complex-valued convolution integrals:

> $\operatorname{Re}\mathcal{H}(\xi,t) + i\operatorname{Im}\mathcal{H}(\xi,t)$  $\sum_{q} e_q^2 \int_{-1}^{1} dx \left[ \frac{1}{\xi - x} \right]$

Where  $\xi = x_B/(2 - x_B)$ ,  $x_B$  is the Bjorken scaling variable, and t is the squared momentum transfer to the nucleon.

This convolution integral cannot be inverted in a model-independent way to yield GPDs.

Thanks to the DVCS CFFs it is possible to obtain the fixed-t dispersion relation (DR) where the subtraction constant is related to the so-called D-term.

$$\operatorname{Re}\mathcal{H}(\xi,t) = \mathcal{C}_{\mathcal{H}}(t) + \frac{1}{\pi} \operatorname{P.V.} \int_{0}^{1} \mathrm{d}\xi' \left[\frac{1}{\xi - \xi'} - \frac{1}{\xi + \xi'}\right] \operatorname{Im}\mathcal{H}(\xi',t)$$

Subtraction constant

The CFFs are linked to measurable quantities, such as differential cross sections and beam, charge, and target polarization asymmetries.

$$(t) = rac{1}{x - i\epsilon} - rac{1}{\xi + x - i\epsilon} \bigg] H_q(x, \xi, t)$$



# Accessing GFFs

Considering some assumptions like:

I. At LO, gluons don't contribute to DVCS much, so one can assume that the subtraction constant is dominantly coming from quarks, but for the moment we neglect also strange and heavier quarks.

2. To talk about "quark pressure" one needs to neglect energy-momentum flow from quarks to gluons described by EMT form factor C(t).

2. In the asymptotic limit  $d_n^q(t)$  vanish for n>1. Where  $d_1^q(t)$  determine the symptomatic form of the GPDs.

3. Considerate the dominance of the flavor singlet combination.

data remains compatible with zero within large uncertainties. experimental data and enhance the results with lattice calculations!

$$\mathcal{C}_{\mathcal{H}}(t) = 2\sum_{q} e_{q}^{2} \int_{-1}^{1} \mathrm{d}z \, rac{D_{\mathrm{term}}^{q}(z,t)}{1-z}, 
onumber \ d_{1}^{u} = d_{1}^{d} = d_{1}^{Q}/2$$

$$D_{\text{term}}^q(z,t) = (1-z^2) \sum_{\text{odd } n} d_n^q(t) C_n^{3/2}(z)$$

$$D_q(t) = \frac{4}{5} d_1^q(t) = \int_{-1}^1 dz \, z \, D_{\text{term}}^q(z, t)$$

$$\Delta(t)=4\left(rac{4}{9}d_1^u(t)+rac{1}{9}d_1^d(t)
ight)$$

$$D^Q(t) = rac{18}{25} \Delta(t) \ .$$

- To date, a more conservative extraction of the subtraction constant from the currently available experimental
- By utilizing this DR, along with other phenomenological assumptions, we can extract the CFFs and GFFs from



### Matching theory with experiment

interference with the Bethe-Heitler radiation occurs

The BH process gives access to elastic form factors, it does not see the 3D structure of hadrons.

$$\frac{d^5\sigma}{dx_B dQ^2 d|t| d\phi d\phi_S} = \frac{\alpha^3 x_B}{16\pi^2 Q^4 \sqrt{1+\epsilon^2}} \left|\mathcal{T}\right|^2$$

$$|\mathcal{T}|^2 = |\mathcal{T}_{\rm BH} + \mathcal{T}_{\rm DVCS}|^2 = |\mathcal{T}_{\rm BH}|^2 + |\mathcal{T}_{\rm DVCS}|^2 +$$

The DVCS amplitude can be decomposed either in helicity amplitudes or, equivalently, in complex valued CFFs.

$$\begin{split} |\mathcal{T}_{\rm DVCS}|^2 &= \frac{2\left(2-2y+y^2\right)}{y^2 Q^2 \left(2-x_{\rm B}\right)^2} \bigg[ 4\left(1-x_{\rm B}\right) \left(|\mathcal{H}|^2+|\widetilde{\mathcal{H}}|^2\right) - \left(x_{\rm B}^2+(2-x_{\rm B})^2 \frac{\Delta^2}{4M^2}\right) |\mathcal{E}|^2 \\ &- x_{\rm B}^2 \left(\mathcal{H}\mathcal{E}^*+\mathcal{E}\mathcal{H}^*+\widetilde{\mathcal{H}}\widetilde{\mathcal{E}}^*+\widetilde{\mathcal{E}}\widetilde{\mathcal{H}}^*\right) - x_{\rm B}^2 \frac{\Delta^2}{4M^2} |\widetilde{\mathcal{E}}|^2 \bigg], \quad y = \frac{Q^2}{xs}. \end{split}$$

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#### Measurements of DVCS are mostly realized via the process of leptoproduction of a real photon, where also an



Thursday's talk by Marija Cuic







# Experimental observables

the Bethe-Heitler (BH) process in the denominator of the DVCS amplitude.

For example, the first sine harmonic of the beam spin asymmetry (BSA), as measured e.g. in Jefferson Lab is defined as:

$$A_{LU}^{-,\sin\phi} \equiv \frac{1}{\pi} \int_{-\pi}^{\pi} d\phi \, \sin\phi \, A_{LU}^{-}(\phi)$$

This harmonics are then proportional to linear combinations of CFFs:

$$A_{LU}^{-,\sin\phi} \propto \mathfrak{Im}\left(F_1\mathcal{H} -$$

If cross-sections are measured, using weighted Fourier integrals can help to cancel the strongly oscillating Bethe-Heitler propagators and interference terms improving the convergence of harmonic terms

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# For observables like asymmetries, it is possible to still access more or less linear CFFs using the dominance of

$$-rac{t}{4M^2}F_2\mathcal{E}+rac{x_B}{2}(F_1+F_2) ilde{\mathcal{H}}
ight)$$















# •Framework: Extraction of Gravitational Form Factors using Neural Networks, from experimental data. •Parallel projects: Database. Technical difficulties of using NNets.





# Framework

We are building upon existing tools: Geoderic

- Python software framework dedicated to the study of GPDs.
- scattering at next-to-leading order and beyond, Nuclear Physics B, Volume 794, Issues 1–2, 2008

To obtain a more reliable estimate of uncertainties with reduced model dependence, we decided to utilize Artificial Neural Networks to parameterize the CFFs.

- Use State-of-the-art standard ML libraries. 
   **O PyTorch**
- A tool to fit all existing data using NN to extract GFFs.

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• K. Kumerički, D. Müller, K. Passek-Kumerički, Towards a fitting procedure for deeply virtual Compton





$$\operatorname{Re}\mathcal{H}(\xi,t) = \mathcal{C}_{\mathcal{H}}(t) + \frac{1}{\pi} \operatorname{P.V.} \int_{0}^{1} \mathrm{d}\xi' \left[\frac{1}{\xi - \xi'} - \frac{1}{\xi + \xi'}\right] \operatorname{Im}\mathcal{H}(\xi,t) = \mathcal{C}_{\mathcal{H}}(t) + \frac{1}{\pi} \operatorname{P.V.} \int_{0}^{1} \mathrm{d}\xi' \left[\frac{1}{\xi - \xi'} - \frac{1}{\xi + \xi'}\right] \operatorname{Im}\mathcal{H}(\xi,t) = \mathcal{C}_{\mathcal{H}}(t) + \frac{1}{\pi} \operatorname{P.V.} \int_{0}^{1} \mathrm{d}\xi' \left[\frac{1}{\xi - \xi'} - \frac{1}{\xi + \xi'}\right] \operatorname{Im}\mathcal{H}(\xi,t) = \mathcal{C}_{\mathcal{H}}(t) + \frac{1}{\pi} \operatorname{P.V.} \int_{0}^{1} \mathrm{d}\xi' \left[\frac{1}{\xi - \xi'} - \frac{1}{\xi + \xi'}\right] \operatorname{Im}\mathcal{H}(\xi,t) = \mathcal{C}_{\mathcal{H}}(t) + \frac{1}{\pi} \operatorname{P.V.} \int_{0}^{1} \mathrm{d}\xi' \left[\frac{1}{\xi - \xi'} - \frac{1}{\xi + \xi'}\right] \operatorname{Im}\mathcal{H}(\xi,t) = \mathcal{C}_{\mathcal{H}}(t) + \frac{1}{\pi} \operatorname{P.V.} \int_{0}^{1} \mathrm{d}\xi' \left[\frac{1}{\xi - \xi'} - \frac{1}{\xi + \xi'}\right] \operatorname{Im}\mathcal{H}(\xi,t) = \mathcal{C}_{\mathcal{H}}(t) + \frac{1}{\pi} \operatorname{P.V.} \int_{0}^{1} \mathrm{d}\xi' \left[\frac{1}{\xi - \xi'} - \frac{1}{\xi + \xi'}\right] \operatorname{Im}\mathcal{H}(\xi,t) = \mathcal{C}_{\mathcal{H}}(t) + \frac{1}{\pi} \operatorname{P.V.} \int_{0}^{1} \mathrm{d}\xi' \left[\frac{1}{\xi - \xi'} - \frac{1}{\xi + \xi'}\right] \operatorname{Im}\mathcal{H}(\xi,t) = \mathcal{C}_{\mathcal{H}}(t) + \frac{1}{\pi} \operatorname{P.V.} \int_{0}^{1} \mathrm{d}\xi' \left[\frac{1}{\xi - \xi'} - \frac{1}{\xi + \xi'}\right] \operatorname{Im}\mathcal{H}(\xi,t) = \mathcal{C}_{\mathcal{H}}(t) + \frac{1}{\pi} \operatorname{P.V.} \int_{0}^{1} \mathrm{d}\xi' \left[\frac{1}{\xi - \xi'} - \frac{1}{\xi + \xi'}\right] \operatorname{Im}\mathcal{H}(\xi,t) = \mathcal{C}_{\mathcal{H}}(t) + \frac{1}{\pi} \operatorname{P.V.} \int_{0}^{1} \mathrm{d}\xi' \left[\frac{1}{\xi - \xi'} - \frac{1}{\xi + \xi'}\right] \operatorname{Im}\mathcal{H}(\xi,t) = \mathcal{C}_{\mathcal{H}}(t) + \frac{1}{\pi} \operatorname{P.V}(t) + \frac$$

It's essentially a least-squares fit of a complex, multi-parameter function, implying no theoretical bias.



# Database

uncertainties.

We developed a publicly accessible database format capable of storing current and future data as well as pseudodata.

- Stores data and pseudodata coming from various elastic and exclusive measurements and lattice-QCD results.
- Interfaces are provided in the main programming languages: Python and C++.
- Interface with GEPARD and Parsons.
- Accessible by anyone in the world using Github.

arxiv.org/abs/2503.18152

https://opengpd.github.io/gpddatabase/index.html

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#### Currently, there was no unified common database storing all available data which is necessary to reduce the

Database for studying Generalised Parton Distributions (GPDs

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available data adding new data

acknowledgements

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View PyPI repository:

View reference article:

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#### Available datasets

uuid	collaboration	reference	type	pseudo	observables	C
EqbtDRkv	HallA	https://arxiv.org/ pdf/1703.09442.pdf	DVCS	False	CrossSectionDifferenceLU	N
PusMstKs	CLAS	https://arxiv.org/ pdf/2211.11274	DVCS	False	ALL	N
vGAKAf7P	CLAS	https://arxiv.org/ pdf/1501.07052.pdf	DVCS	False	ALU AUL ALL	N
75ueQoQw	COMPASS	https://arxiv.org/ abs/1802.02739	DVCS	False	CrossSectionUUVirtualPhotoProduction TSlope	N
ob8hLTm2	CLAS	https://arxiv.org/ pdf/1810.02110	DVCS	False	CrossSectionUU	D C ex e
3gYp9R4P	HERMES	https://arxiv.org/ pdf/hep- ex/0605108.pdf	DVCS	False	AcCos1Phi	N
bmTzHHvg	HallA	https://arxiv.org/ pdf/2109.02076	DVCS	False	CrossSectionUU	B in cı
RQncbKtk	CLAS	https://arxiv.org/ pdf/1504.02009v1.pdf	DVCS	False	CrossSectionUU CrossSectionDifferenceLU	N
AtY8o7Ej	HallA	https://arxiv.org/ pdf/1504.05453v1.pdf	DVCS	False	CrossSectionUU CrossSectionDifferenceLU	N

#### In collaboration with P. Sznajder and C. Mezrag





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# Database

uncertainties.

We developed a publicly accessible database format capable of storing current and future data as well as pseudodata.

- Get the list of available data files.
- Load a given data file
- Get metadata.
- Accessible by anyone in the world using Github.

arxiv.org/abs/2503.18152

https://opengpd.github.io/gpddatabase/index.html

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#### **Available datasets**

uuid	collaboration	reference	type	pseudo	observables	C
EqbtDRkv	HallA	https://arxiv.org/ pdf/1703.09442.pdf	DVCS	False	CrossSectionDifferenceLU	N
PusMstKs	CLAS	https://arxiv.org/ pdf/2211.11274	DVCS	False	ALL	N
vGAKAf7P	CLAS	https://arxiv.org/ pdf/1501.07052.pdf	DVCS	False	ALU AUL ALL	N
75ueQoQw	COMPASS	https://arxiv.org/ abs/1802.02739	DVCS	False	CrossSectionUUVirtualPhotoProduction TSlope	N
ob8hLTm2	CLAS	https://arxiv.org/ pdf/1810.02110	DVCS	False	CrossSectionUU	D C e
3gYp9R4P	HERMES	https://arxiv.org/ pdf/hep- ex/0605108.pdf	DVCS	False	AcCos1Phi	N
bmTzHHvg	HallA	https://arxiv.org/ pdf/2109.02076	DVCS	False	CrossSectionUU	B ir c
RQncbKtk	CLAS	https://arxiv.org/ pdf/1504.02009v1.pdf	DVCS	False	CrossSectionUU CrossSectionDifferenceLU	N
AtY8o7Ej	HallA	https://arxiv.org/ pdf/1504.05453v1.pdf	DVCS	False	CrossSectionUU CrossSectionDifferenceLU	N

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# Technical difficulties of using NNets

representations.

- This shows how as the number of models increases, the mean residual prediction tends to stabilize.
- The standard deviation plot shows how the uncertainty in our predictions increases as the ensemble shrinks



A larger ensemble tends to 'average out' individual model noise, improving the predictive stability results and uncertainty estimation of our model.

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#### To solve instability in results we use ensemble analysis by averaging over multiple learned



# State of the project



# Experimental observables

Experimental data we have used:

- We started with a specific experiment.
- Describe 2007 and 2015 CLAS data to compare with previous results.
- Add the rest of the CLAS data at 6 GeV
- Add CLAS data at ~ || GeV.

Collab	Year	Observable	1.b.e.	$x_B$	$Q^2[GeV^2]$	t [Ge
CLAS	2001	$A_{LU}^{sin(\phi)}$	4.25	0.19	1.25	0.19
CLAS	2006	$A_{UL}^{sin(\phi)}$	5.7	0.2-0.4	1.82	0.15-0.
CLAS	2007	$A_{LU}$	5.77	0.13-0.35	1.1-3	0.1-0.
CLAS	2009	$A_{LU}$	4.8	0.12-0.48	1.0-2.8	0.09-1
CLAS	2015	$A_{LU}(\phi), A_{UL}(\phi), A_{LL}(\phi)$	5.93	0.18-0.4	1.6-3.2	0.1-0.4
CLAS	2015	$\sigma(\phi), \Delta\sigma(\phi)$	5.75	0.1-0.58	1-4.6	0.09-0.
CLAS	2018	$\sigma(\phi)$	5.88	0.12-0.5	1.1-4	0.1-1.
CLAS	2023	A <sub>LL</sub>	10.2	0.09-0.45	1.3-6	0.1-2.
CLAS	2023	A <sub>LL</sub>	10.6	0.09-0.62	1.1-7.2	0.1-2.

We made cuts at t=0.51 GeV<sup>2</sup>, and Q<sup>2</sup><1.5 GeV<sup>2</sup> or  $-t/Q^2$ <0.25

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Lepton-beam-energy (l.b.e.)





### State of the project

#### Description of 2015 CLAS data through the NN









# State of the project

• First, we used 2008 beam spin asymmetry and 2015 cross-section CLAS data to reproduce previous results.

#### CFFs using NN and dispersion relation



Qualitative good agreement with Kumericki's previous results. 

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#### ImH and ReH CFFs

Jefferson Lab







![](_page_21_Picture_7.jpeg)

![](_page_21_Figure_9.jpeg)

# State of the project: Inputs from Lattice QCD

#### Electromagnetic form factors from Lattice QCD

• The obtained fit will be used as a constraint for the neural network.

![](_page_22_Figure_3.jpeg)

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I.J. High Energy. Phys. 2024, 162 (2024)

![](_page_22_Picture_8.jpeg)

# State of the project: Inputs from Lattice QCD

- calculated in Lattice across a wide range of t.

![](_page_23_Figure_4.jpeg)

![](_page_23_Picture_9.jpeg)

### State of the project: D-term

Finally, the subtraction constant obtained from the NN leads to the following D-term values:

![](_page_24_Figure_2.jpeg)

$$D^Q(t)=rac{18}{25}\Delta(t)$$

#### • More data needs to be added to reduce uncertainties as well as a bigger ensemble for the NNets

![](_page_24_Picture_8.jpeg)

# Deliverables and future directions

![](_page_25_Picture_1.jpeg)

# Double DVCS (DDVCS)

#### Another privilege process for GPDs extraction is the Double DVCS

![](_page_26_Picture_2.jpeg)

Possibility of directly measuring GPDs for  $x \neq \xi$ 

![](_page_26_Picture_4.jpeg)

Providing a potential solution to deconvolution challenges.

![](_page_26_Picture_6.jpeg)

Better estimations of the moments of GPDs.

![](_page_26_Picture_8.jpeg)

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#### *PoS* DIS2024 (2025) 239

![](_page_26_Figure_11.jpeg)

Monday's talk by Victor Martínez-Fernández K. Deja, et. al, Acta Phys. Polon. Supp. 16 (2023) 7, 7-A24

![](_page_26_Picture_13.jpeg)

![](_page_26_Picture_15.jpeg)

![](_page_26_Picture_16.jpeg)

![](_page_26_Picture_17.jpeg)

# Deliverables and future directions

- Understand physical constraints from LQCD and implement them in the network architecture.
- Preliminary CFFs and D-term results for all CLAS data have been extracted.

#### Future directions on Lattice QCD computation:

- Calculate the isoscalar GFFs using Lattice QCD. I.J. High Energy. Phys. 2024, 162 (2024)
- Implement physical constraints from LQCD in the network architecture.

#### Future directions on the fitting framework:

- Implement DDVCS process into Gepard.
- Optimization of NNets of Gepard, for instance, implementation of the possibility for use GPU device.
- (TMDs & GPDs).
- Global analysis by steps: add HallA data and the rest of DVCS data.

Completed implementation of DR with PyTorch into Gepard
A tool to fit all existing data using NN to extract GFFs capable of handling both experimental and LQCD data.

Implement NN to which is a C++ framework for the phenomenology of the 3D structure of the nucleon

![](_page_27_Picture_20.jpeg)

# Thank you all for your attention

![](_page_28_Picture_1.jpeg)

# Other details

Data used thus far:

- Therefore new 2023 BSA CLASI2 data at 10.6 and 10.2 GeV (Phys. Rev. Lett. 130, 211902) was implemented to raise kinematic restrictions, and then the rest of CLAS data was added.
- Planning to use new 2022 Hall A data (Phys. Rev. Lett. 128, 252002)
- Planning to use upcoming cross-section CLASI2 data.

#### Neural Network construction in Gepard:

- previously necessary.
- Trained 10 neural nets per model with an ensemble of at least 15 trainings.
- •The architecture of nets parametrizing imaginary part of CFF is:
  - 2 input neurons (for  $x_B$  and t)
  - I7 and 25 neurons in hidden layers for reproducing previous results.
  - 20 and 25 neurons in hidden layers for the new CLAS data
  - constant.

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• First, we used 2008 beam spin asymmetry and 2015 cross-section CLAS data to reproduce previous results.

• Fits are performed in harmonic space, therefore harmonic analysis (Fourier transform) of the data was

I neuron in the output layer for each required imaginary part of CFF as well as the subtraction

![](_page_29_Picture_19.jpeg)