



ALMA MATER STUDIORUM
UNIVERSITÀ DI BOLOGNA

Recent Advances and Future Perspectives in Nuclear Astrophysics at

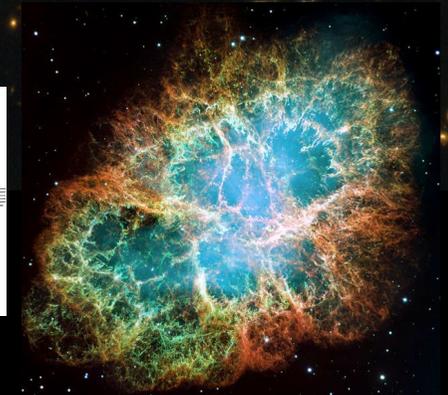
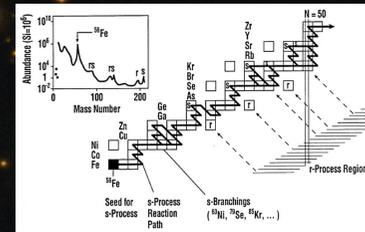
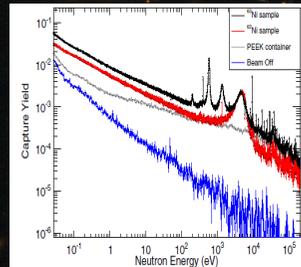
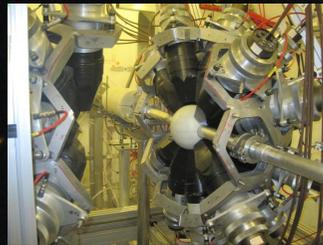


Cristian Massimi

Department of Physics and Astronomy

Outline:

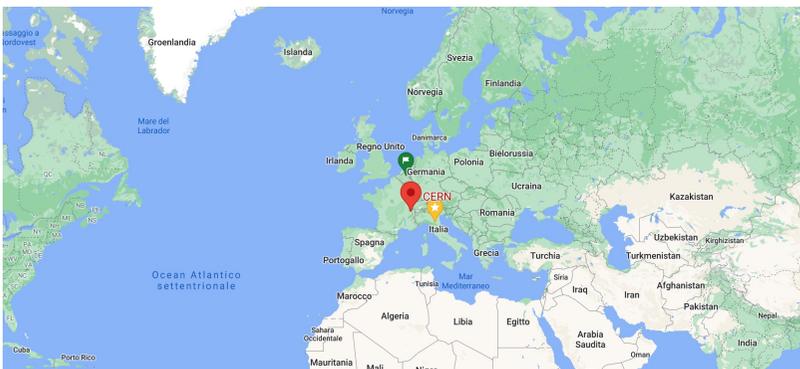
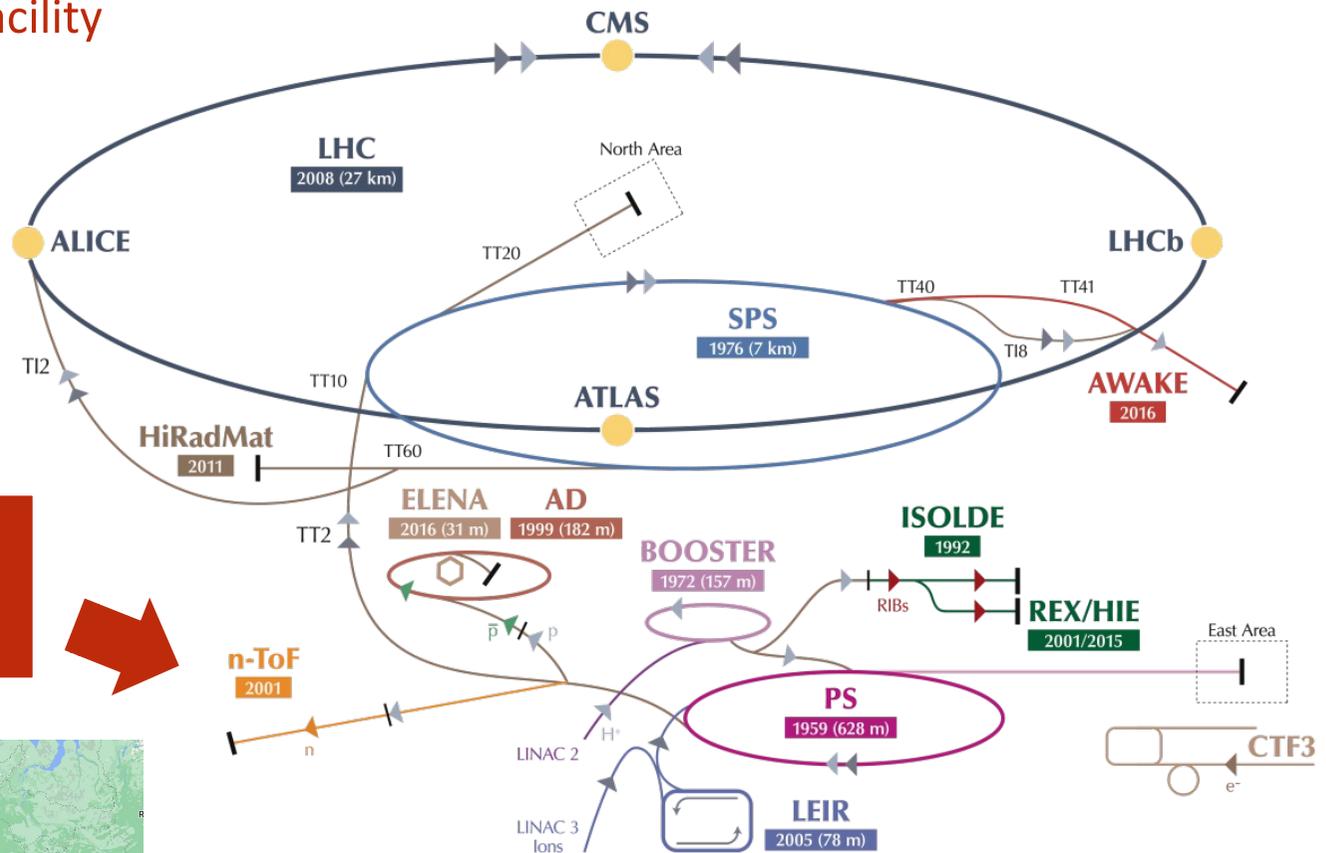
- n_TOF @ CERN 
- **Recent measurements and their impact on Nuclear Astrophysics**
(Branching isotopes, Neutron source of the s process, BBN)
- **Some projects for the near future**



n_TOF: neutron time-of-flight facility @ CERN



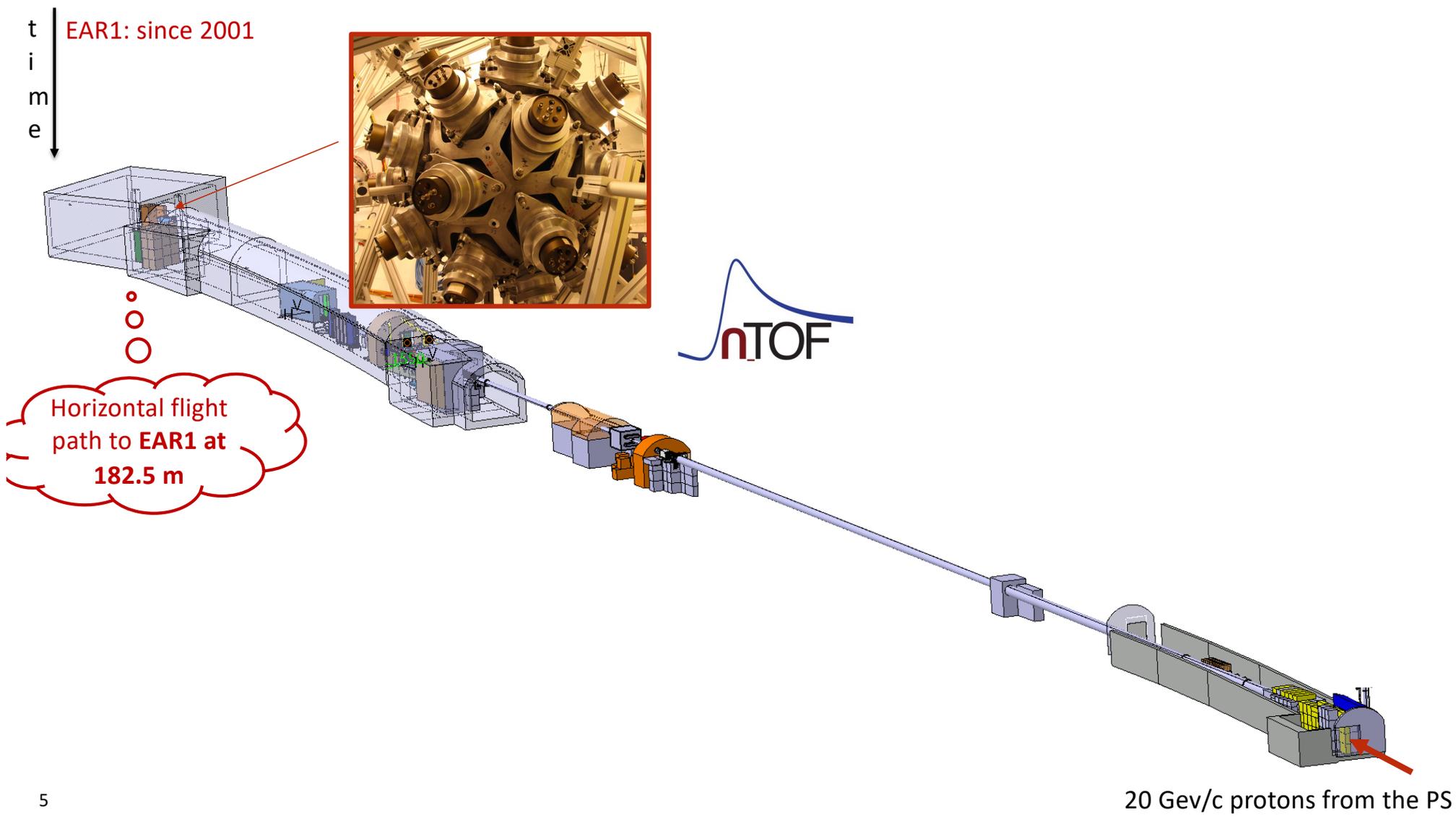
C. Rubbia et al., A high resolution spallation driven facility at the CERN-PS to measure neutron cross sections in the interval from 1 eV to 250 MeV CERN/LHC/98 02(EET) 1998



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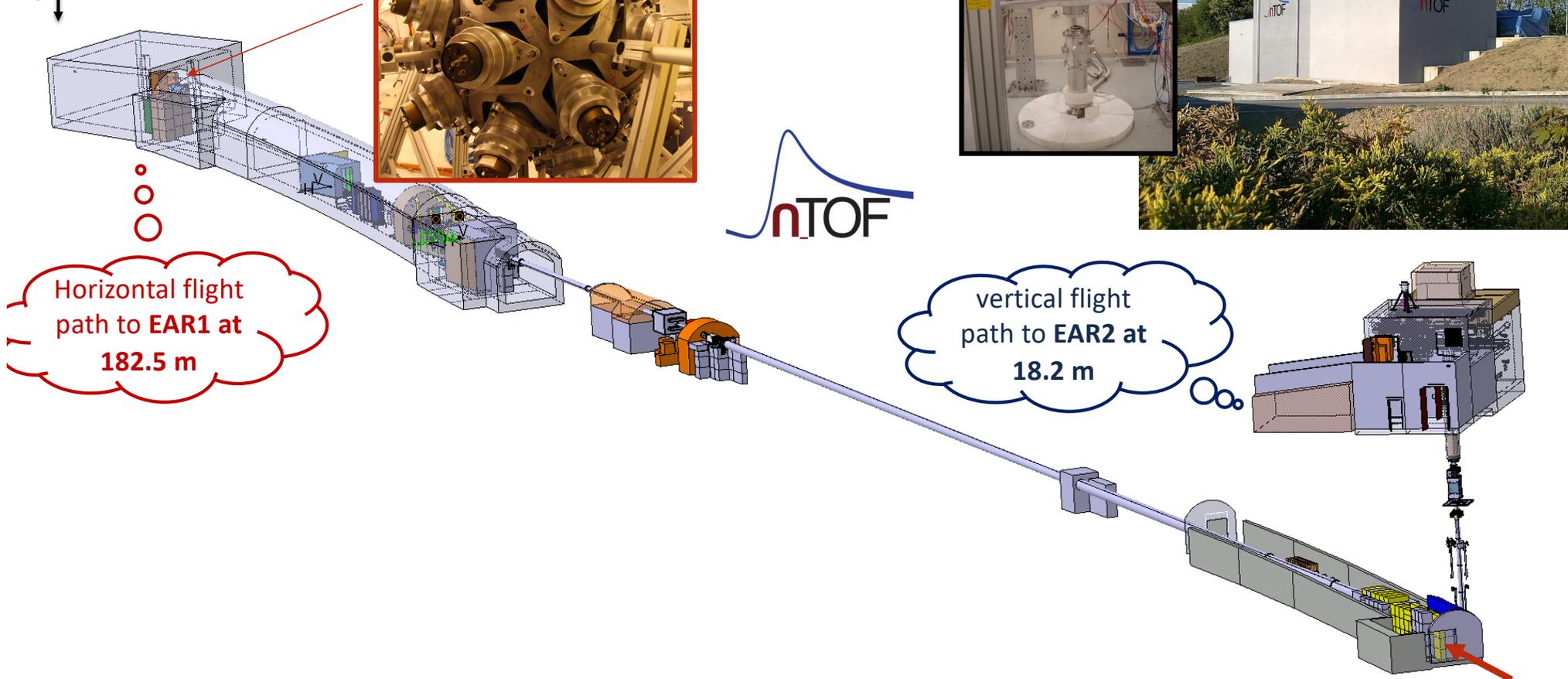
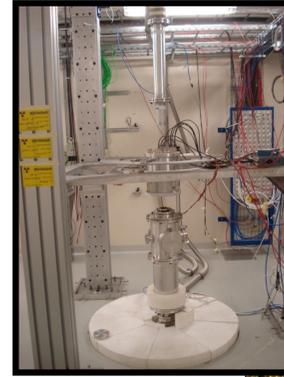
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EAR1: since 2001
EAR2: since 2014

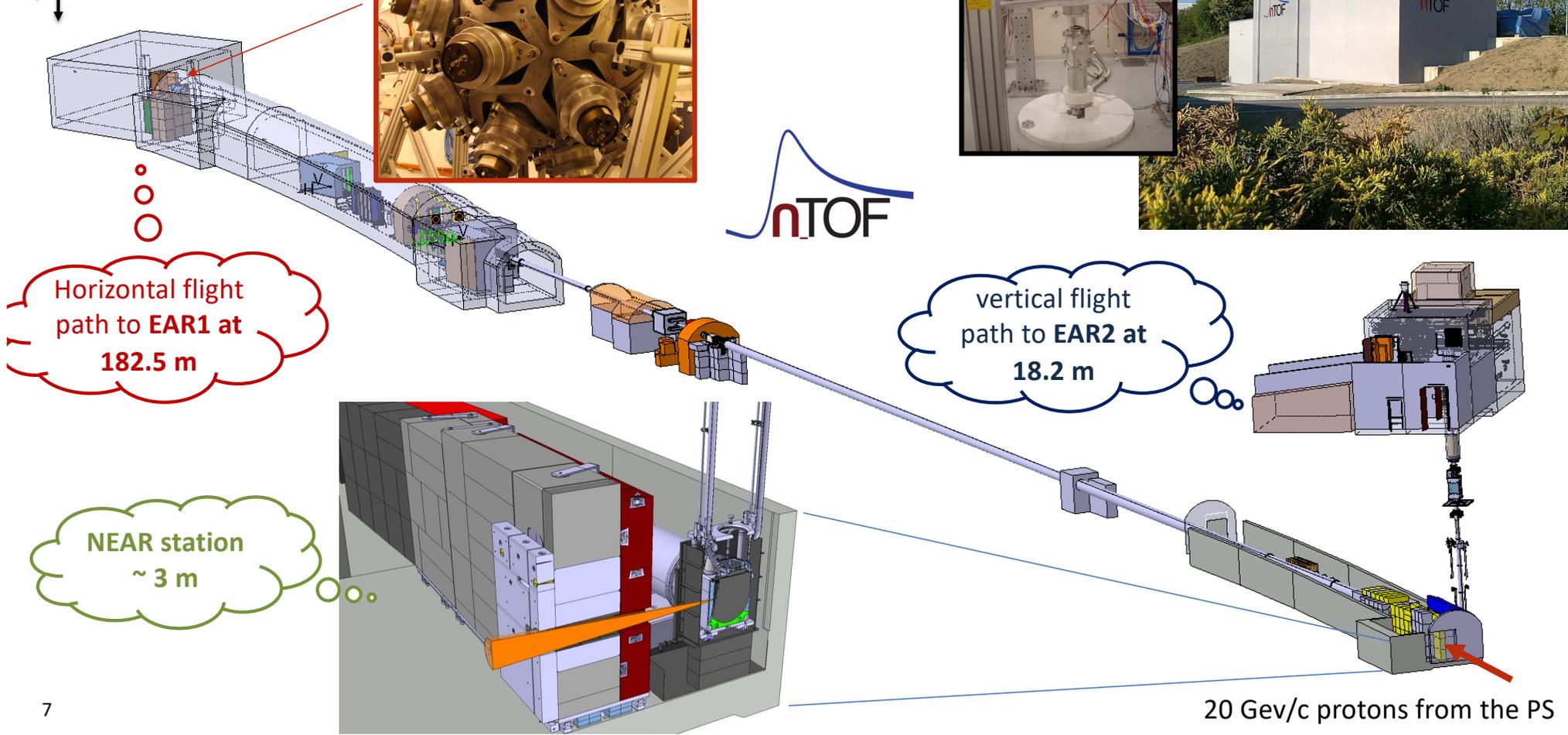
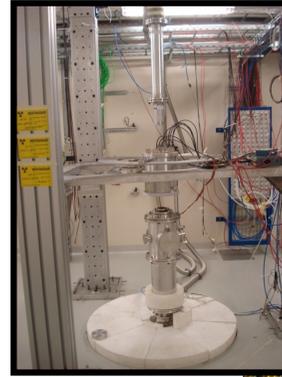


Horizontal flight path to EAR1 at 182.5 m

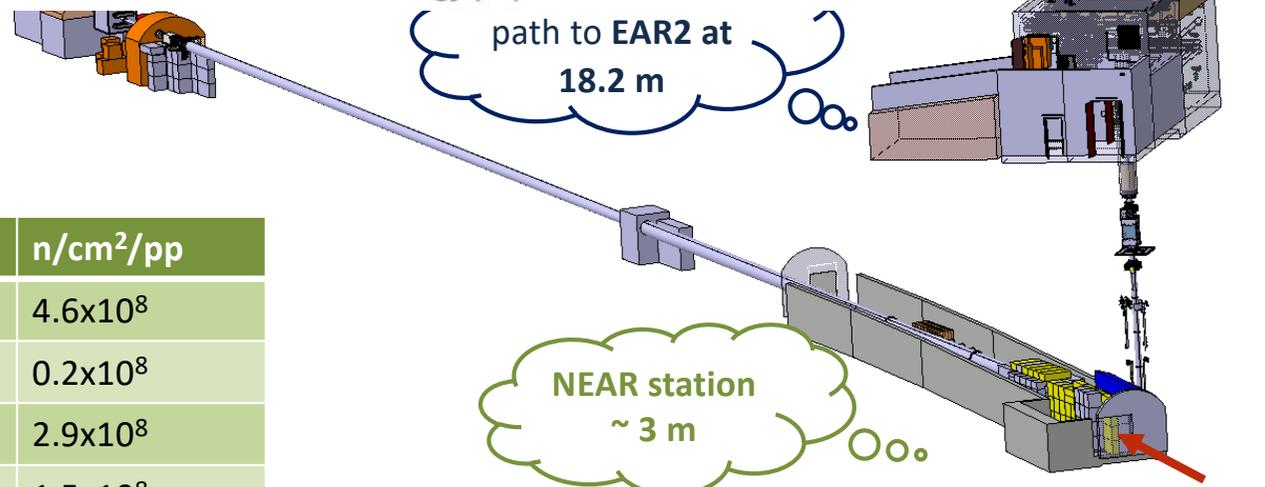
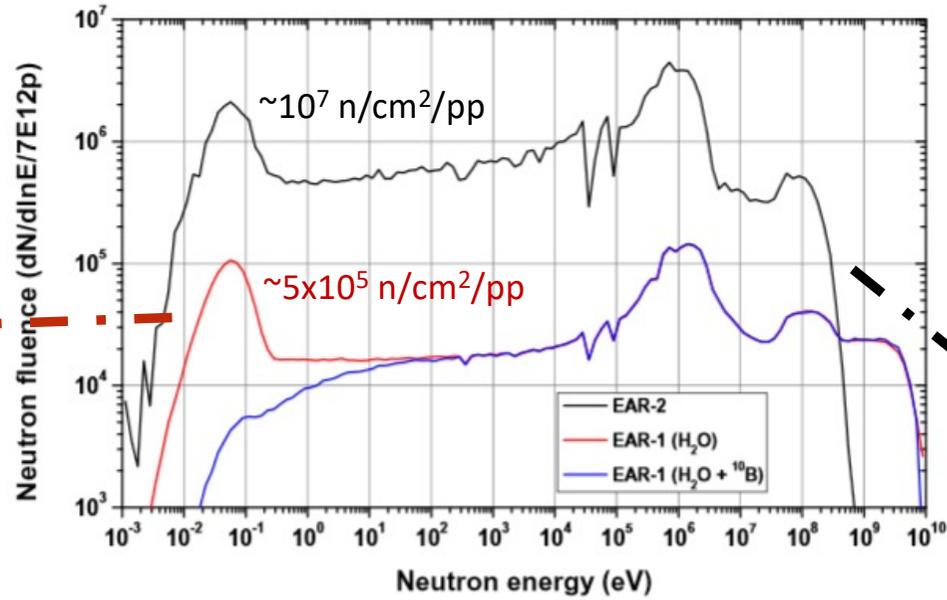
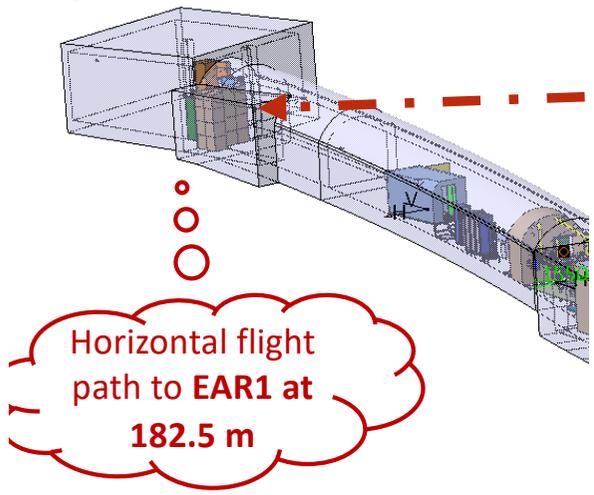
vertical flight path to EAR2 at 18.2 m

20 GeV/c protons from the PS

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m
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EAR1: since 2001
EAR2: since 2014
NEAR: Since 2021



time
 EAR1: since 2001
 EAR2: since 2014
 NEAR: Since 2021



20 GeV/c protons from the PS

Neutron energy region	n/cm ² /pp
ALL energies	4.6×10^8
$E_n < 1$ keV	0.2×10^8
$1 \text{ keV} < E_n < 1 \text{ MeV}$	2.9×10^8
$E_n > 1 \text{ MeV}$	1.5×10^8

n_TOF @ CERN

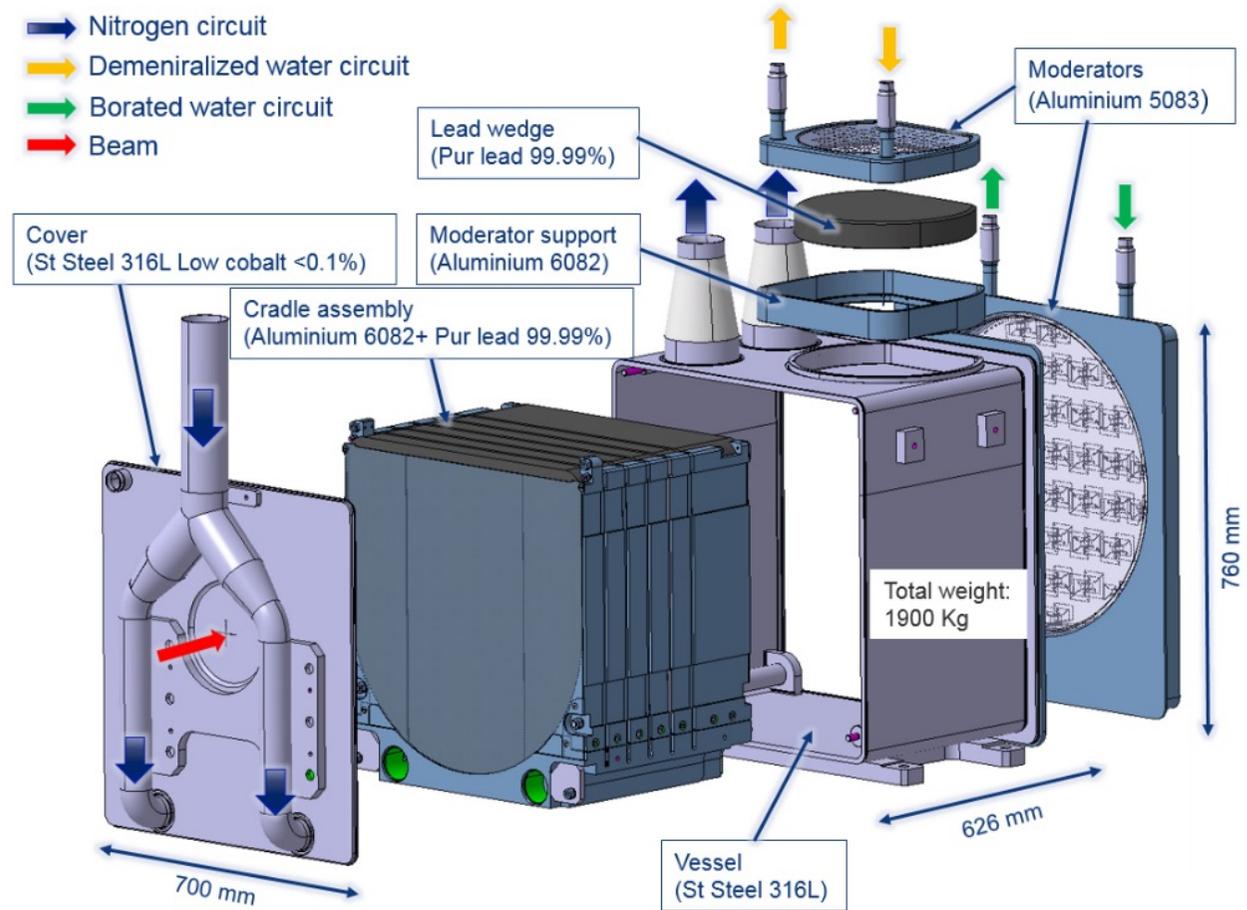
The advantages of n_TOF are a direct consequence of the characteristics of the **PS proton beam: high energy, high peak current, low duty cycle.**

proton beam momentum	20 GeV/c
intensity (dedicated mode)	$\sim 10^{13}$ protons/pulse
repetition frequency	1 pulse/1.2s
pulse width	6 ns (rms)
n/p	300
lead target dimensions	80x80x60 cm ³
cooling & moderation material	N ₂ & (H ₂ O + ¹⁰ B)
moderator thickness in the exit face	5 cm
neutron beam dimension in EAR-1 (capture mode)	2 cm (FWHM)

n_TOF @ CERN

3rd generation spallation target

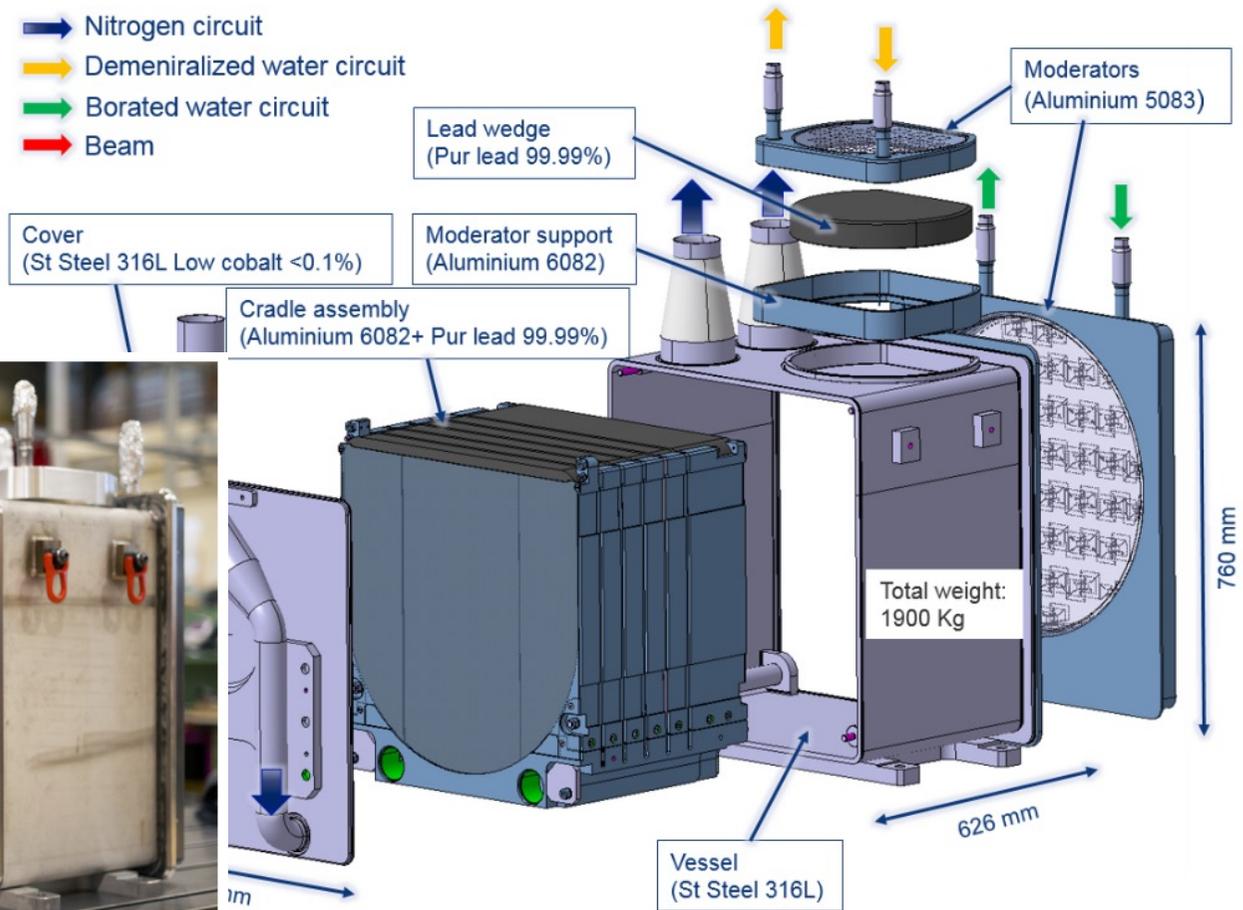
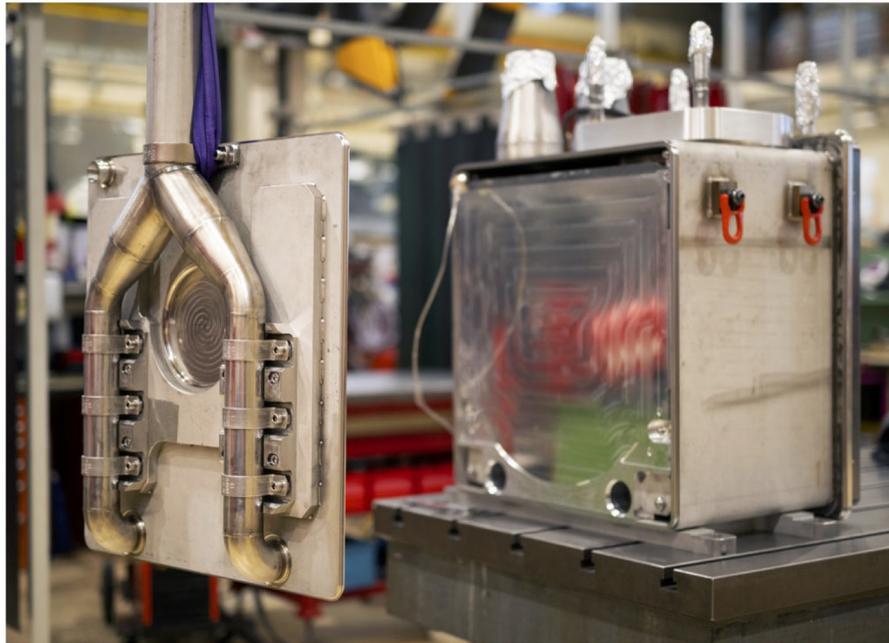
- ❖ pure Pb based
- ❖ N₂-gas cooled, water moderated
- ❖ Several innovations have been introduced



courtesy of Oliver Aberle and Marco Calviani, CERN

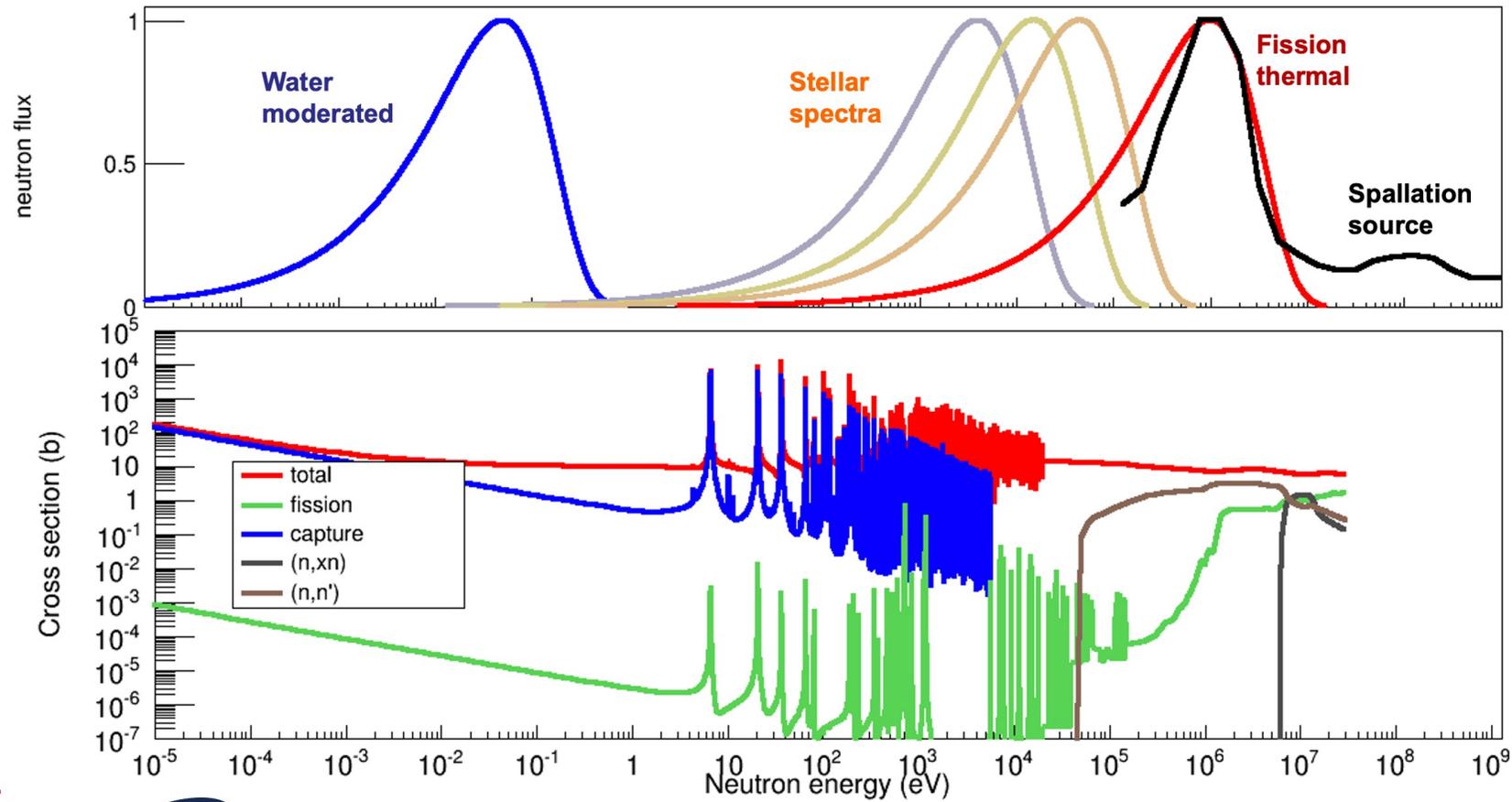
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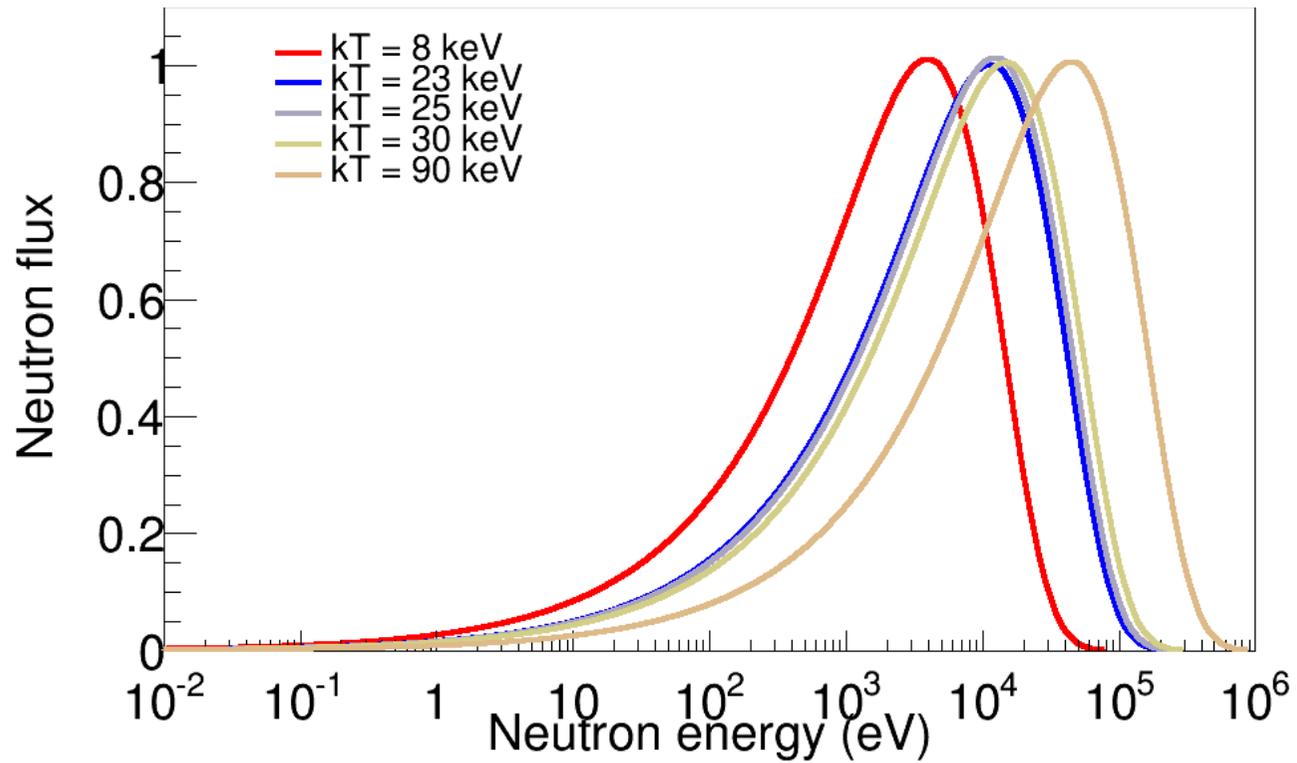
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Neutron-induced reactions



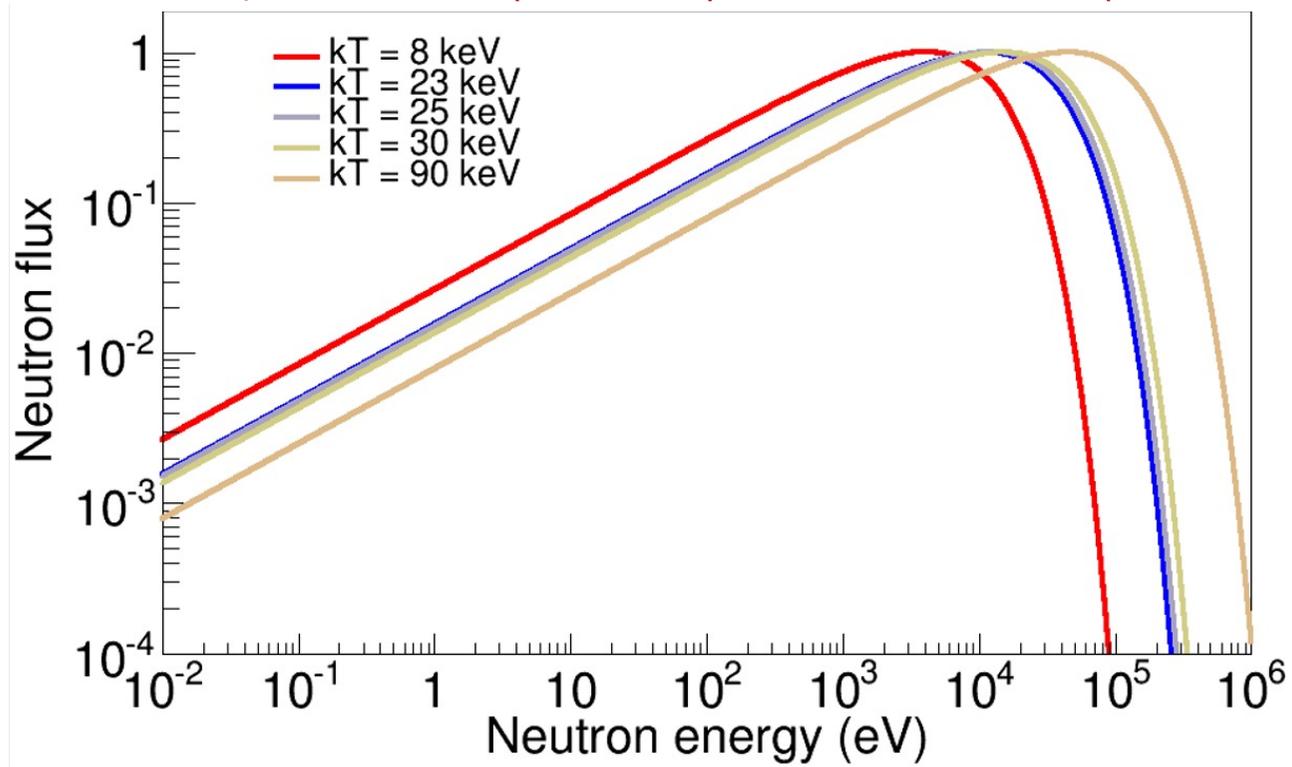
Neutron-induced reactions

Stellar spectra: AGB (8, 23 keV) and Massive stars (25, 90 keV)



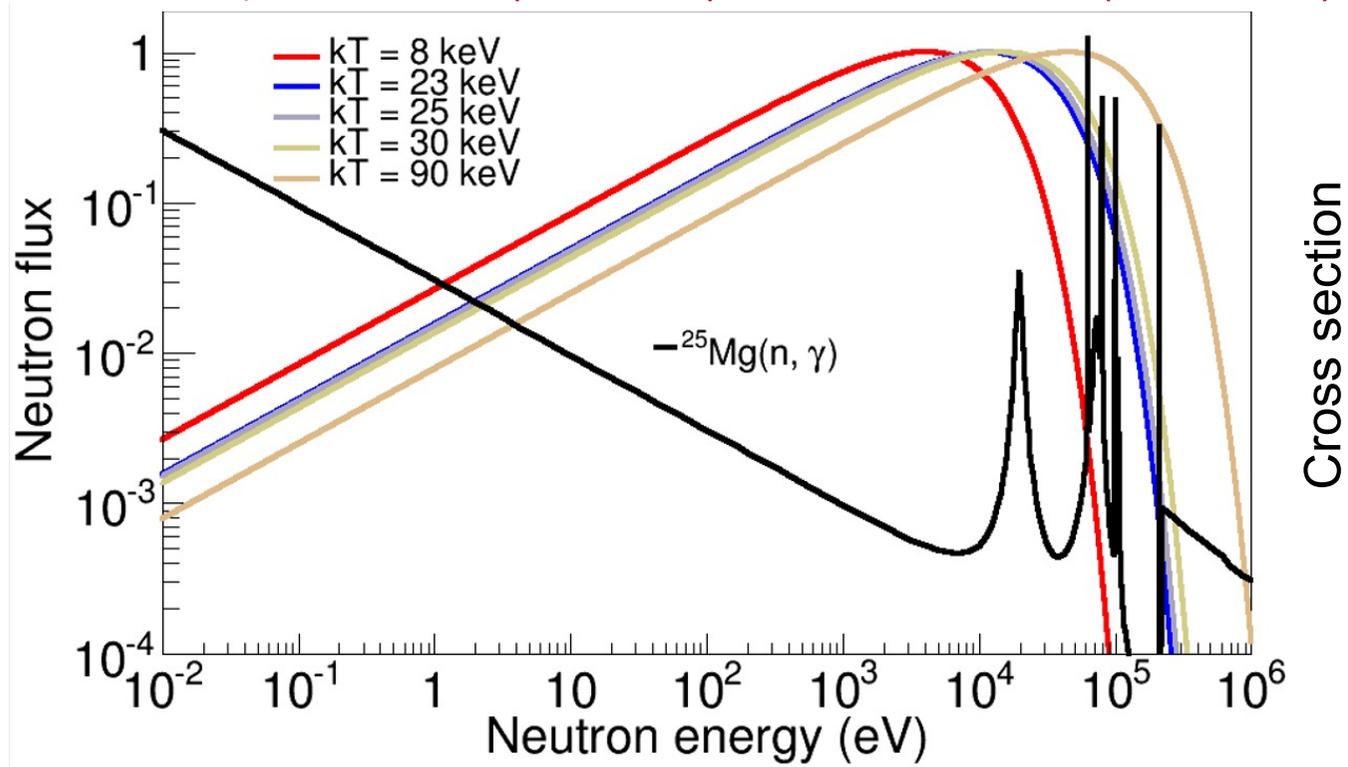
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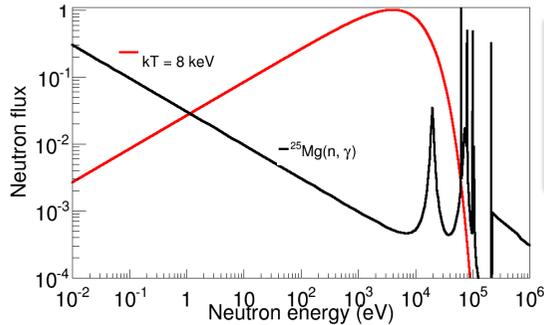


Neutron-induced reactions

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Neutron-induced reactions

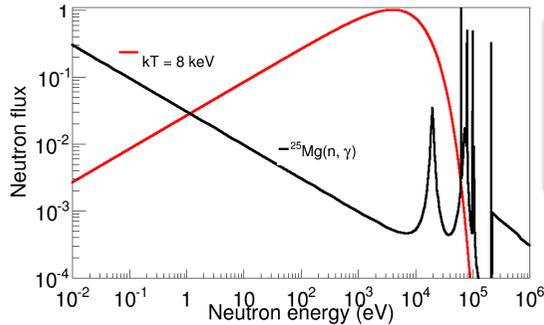


For nuclear astrophysics, what is important is the **Maxwellian Averaged Cross-Sections (MACS)** at various temperatures (kT depends on stellar site).

Reaction rate ($\text{cm}^{-3}\text{s}^{-1}$): $r = N_A N_n v \sigma(v) \Rightarrow r = N_A N_n \langle \sigma \cdot v \rangle$

$$\text{MACS} \equiv \frac{\langle \sigma \cdot v \rangle}{v_T} = \frac{2}{\sqrt{\pi}(kT)^2} \int_0^\infty \sigma(E) E e^{-E/(kT)} dE$$

Neutron-induced reactions



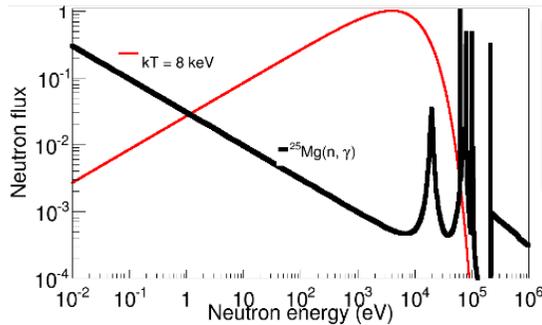
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Two methods to determine MACS:

Neutron-induced reactions



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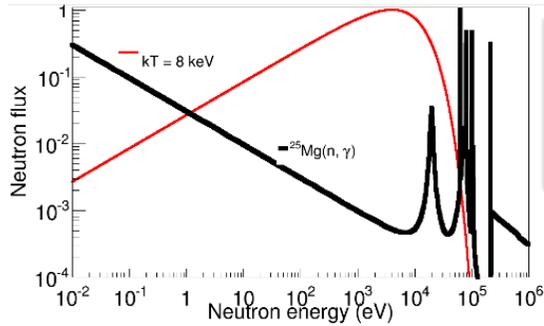
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Two methods to determine MACS:

1. measurement of **energy dependent** neutron capture cross-sections \rightarrow EAR1 & EAR2

Neutron-induced reactions



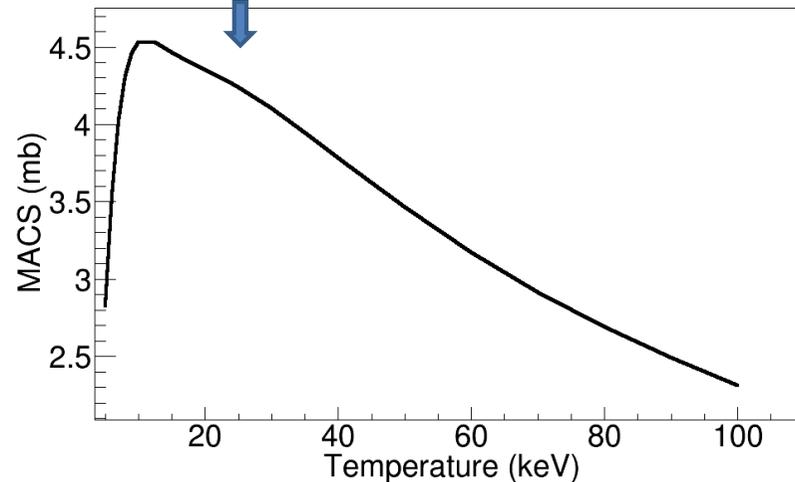
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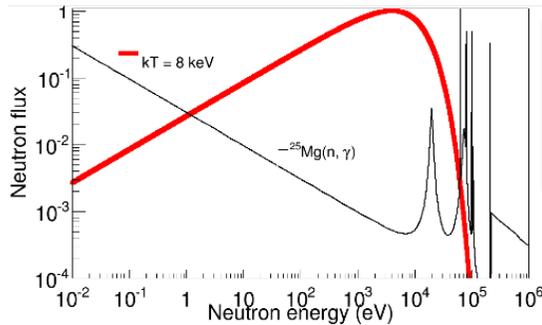
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Neutron-induced reactions



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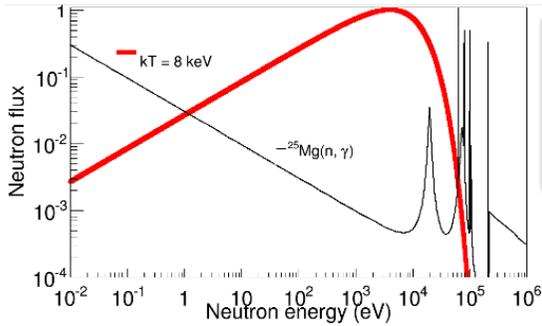
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Two methods to determine MACS:

1. measurement of **energy dependent** neutron capture cross-sections \rightarrow **EAR1 & EAR2**
2. **integral measurement** (energy integrated) using neutron beams with suitable energy \rightarrow **NEAR**

Neutron-induced reactions



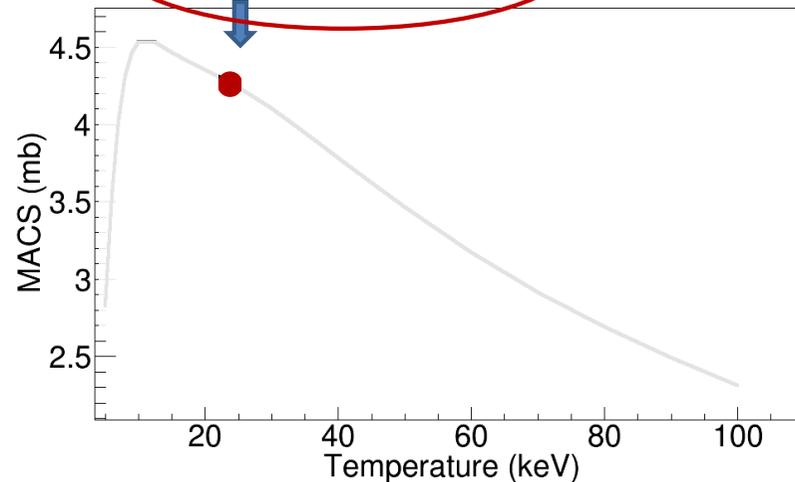
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n_TOF: nuclear data for science (and technology)

... so far

Radiative capture reactions (n,γ)	(89)
Fission reactions (n,f)	(37)
Light particle emission reactions (n,lcp)	(11)
Detector developments	(6)

The n_TOF Collaboration list of publications:
<https://twiki.cern.ch/NTOFPublic/ListOfPublications>



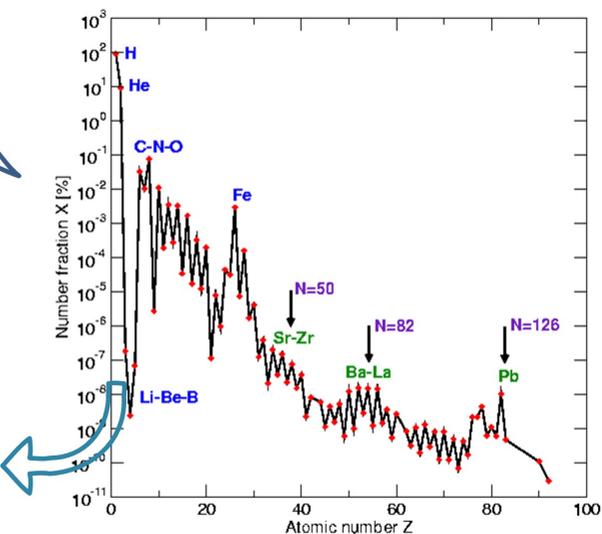
n_TOF: nuclear data for science (and technology)

How chemical elements are synthesized in the Universe?

BIG BANG Nucleosynthesis:

○ ${}^7\text{Be}(n,\alpha)$ and ${}^7\text{Be}(n,p)$ cross section measurement for the Cosmological Lithium Problem.
Later in this presentation

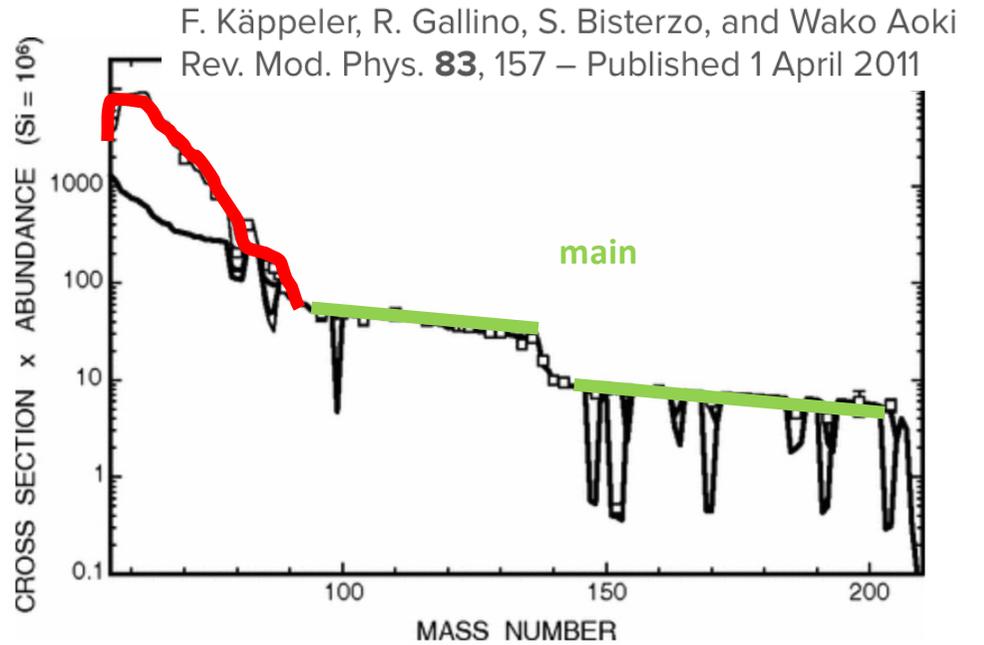
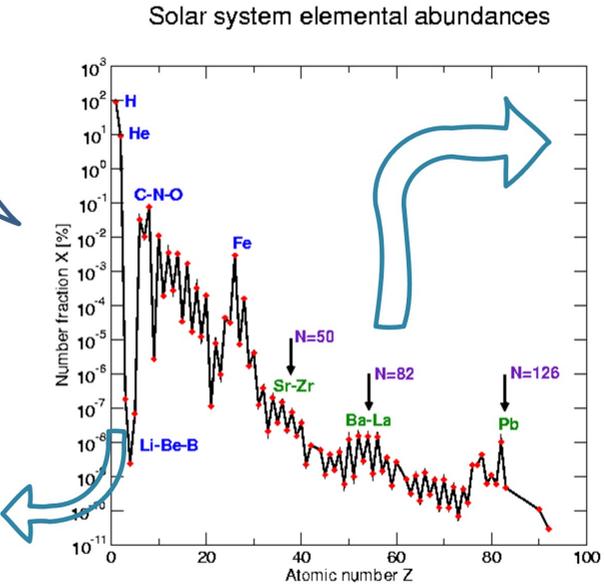
Solar system elemental abundances



n_TOF: nuclear data for science (and technology)

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weak: core He burning in massive stars

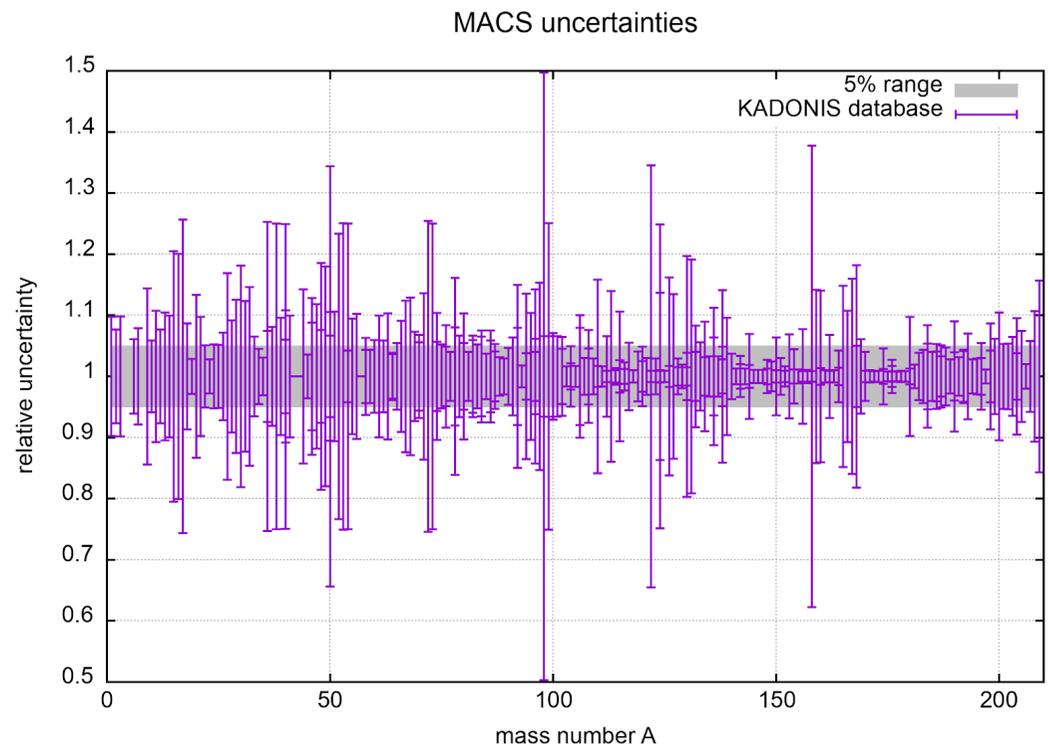
main: He shell flashes in low mass TP-AGB stars

Stellar cross sections (MACS) for the s process

F. Käppeler, R. Gallino, S. Bisterzo, and Wako Aoki
Rev. Mod. Phys. **83**, 157 – Published 1 April 2011

Reducing the uncertainty in the stellar cross section (MACS) is not only a question of better nuclear data: **higher accuracy** in the **reaction rates** opens the possibility to **investigate new astrophysical scenarios**

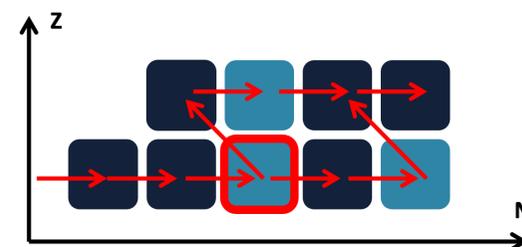
[nuclear clocks, constrains on the BBN, AGB modelling, nucleosynthesis conditions in explosive scenarios, meteoritic grains, others]



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Sample	Half-life (yr)	Q value (MeV)	Comment
^{63}Ni	100.1	β^- , 0.066	TOF work in progress (Couture, 2009), sample with low enrichment
^{79}Se	2.95×10^5	β^- , 0.159	Important branching, constrains s -process temperature in massive stars
^{81}Kr	2.29×10^5	EC, 0.322	Part of ^{79}Se branching
^{85}Kr	10.73	β^- , 0.687	Important branching, constrains neutron density in massive stars
^{95}Zr	64.02 d	β^- , 1.125	Not feasible in near future, but important for neutron density low-mass AGB stars
^{134}Cs	2.0652	β^- , 2.059	Important branching at $A = 134, 135$, sensitive to s -process temperature in low-mass AGB stars, measurement not feasible in near future
^{135}Cs	2.3×10^6	β^- , 0.269	So far only activation measurement at $kT = 25$ keV by Patronis <i>et al.</i> (2004)
^{147}Nd	10.981 d	β^- , 0.896	Important branching at $A = 147/148$, constrains neutron density in low-mass AGB stars
^{147}Pm	2.6234	β^- , 0.225	Part of branching at $A = 147/148$
^{148}Pm	5.368 d	β^- , 2.464	Not feasible in the near future
^{151}Sm	90	β^- , 0.076	Existing TOF measurements, full set of MACS data available (Abbondanno <i>et al.</i> , 2004a; Wisshak <i>et al.</i> , 2006c)
^{154}Eu	8.593	β^- , 1.978	Complex branching at $A = 154, 155$, sensitive to temperature and neutron density
^{155}Eu	4.753	β^- , 0.246	So far only activation measurement at $kT = 25$ keV by Jaag and Käppeler (1995)
^{153}Gd	0.658	EC, 0.244	Part of branching at $A = 154, 155$
^{160}Tb	0.198	β^- , 1.833	Weak temperature-sensitive branching, very challenging experiment
^{163}Ho	4570	EC, 0.0026	Branching at $A = 163$ sensitive to mass density during s process, so far only activation measurement at $kT = 25$ keV by Jaag and Käppeler (1996b)
^{170}Tm	0.352	β^- , 0.968	Important branching, constrains neutron density in low-mass AGB stars
^{171}Tm	1.921	β^- , 0.098	Part of branching at $A = 170, 171$
^{179}Ta	1.82	EC, 0.115	Crucial for s -process contribution to ^{180}Ta , nature's rarest stable isotope
^{185}W	0.206	β^- , 0.432	Important branching, sensitive to neutron density and s -process temperature in low-mass AGB stars
^{204}Tl	3.78	β^- , 0.763	Determines $^{205}\text{Pb}/^{205}\text{Tl}$ clock for dating of early Solar System

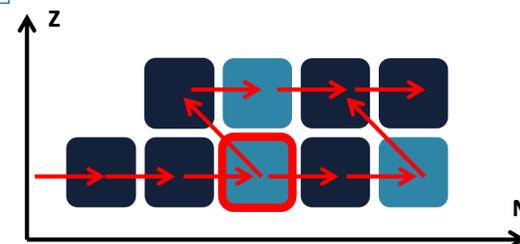


Stellar cross sections (MACS) for the s process

U. Abbondanno, et al. (The n_TOF Collaboration), *Phys. Rev. Lett.* **94** (2004) 161103
 C. Lederer, et al. (The n_TOF Collaboration), *Phys. Rev. Lett.* **110** (2013) 022501
 C. Guerrero, et al. (The n_TOF Collaboration), *Phys. Rev. Lett.* **125** (2020) 142701
 A. Casanovas-Hoste, et al., (The n_TOF Collaboration) *Phys. Rev. Lett.* **133** (2024) 052702

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¹⁴⁷ Pm	2.6234	β^- , 0.225	Part of branching at A = 147/148
¹⁴⁸ Pm	5.368 d	β^- , 2.464	Not feasible in the near future
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¹⁷⁹ Ta	1.82	EC, 0.115	Crucial for s-process contribution to ¹⁸⁰ Ta, nature's rarest stable isotope
¹⁸⁵ W	0.206	β^- , 0.432	Important branching, sensitive to neutron density and s-process temperature in low-mass AGB stars
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Stellar cross sections (MACS) for the s process

Some cross sections measured in 2001 - 2024

- ❖ Branching point isotopes:

^{151}Sm , ^{63}Ni , ^{147}Pm , ^{171}Tm , ^{203}Tl , ^{79}Se

- ❖ Abundances in presolar grains:

$^{28,29,30}\text{Si}$, $^{91,92,93,94,96}\text{Zr}$, $^{94,96}\text{Mo}$, ^{146}Nd

- ❖ Magic Nuclei and end-point:

^{139}La , ^{140}Ce , ^{90}Zr , ^{89}Y , ^{88}Sr , $^{204,206,207,208}\text{Pb}$, ^{209}Bi

- ❖ Seeds isotopes:

$^{54,56,57}\text{Fe}$, $^{58,60,62,64}\text{Ni}$, $^{59}\text{Ni}(n,\alpha)$

- ❖ Isotopes of special interest:

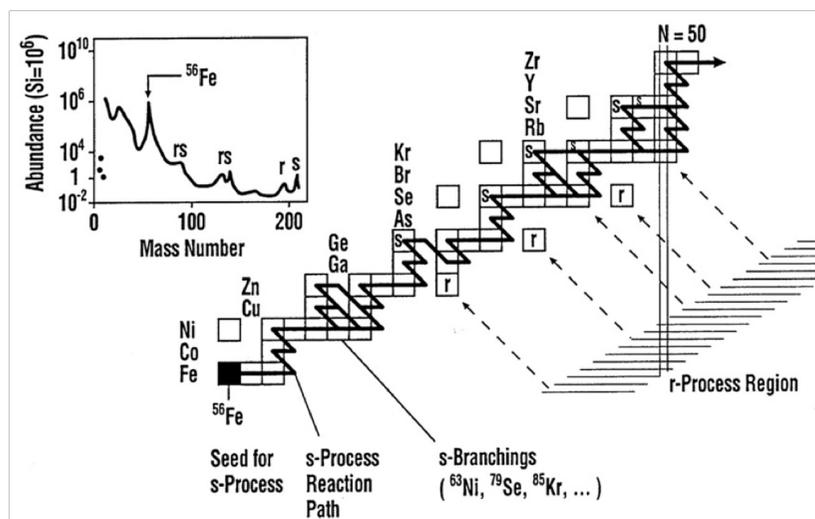
$^{186,187,188}\text{Os}$ (cosmochronometer), ^{197}Au (reference cross section), $^{24,25,26}\text{Mg}$, $^{33}\text{S}(n,\alpha)$, $^{14}\text{N}(n,p)$, $^{35}\text{Cl}(n,p)$,

$^{26}\text{Al}(n,p)$, $^{26}\text{Al}(n,\alpha)$ (neutron poison), ^{154}Gd (s-only isotopes), $^{40}\text{K}(n,p)$, $^{40}\text{K}(n,\alpha)$, $^{63,65}\text{Cu}$, $^{93,94}\text{Nb}$, ^{68}Zn ,

$^{69,71}\text{Ga}$, $^{70,72,73,74,76}\text{Ge}$, $^{77,78,80}\text{Se}$ (weak component), $^{155,157,160}\text{Gd}$, $^7\text{Li}(n,p)$, $^7\text{Li}(n,\alpha)$ BBN

- ❖ Neutron Sources $^{22}\text{Ne}(\alpha,n)^{25}\text{Mg}$ and $^{13}\text{C}(\alpha,n)^{16}\text{O}$:

$n+^{25}\text{Mg}$, $n+^{16}\text{O}$



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❖ Isotopes of special interest:

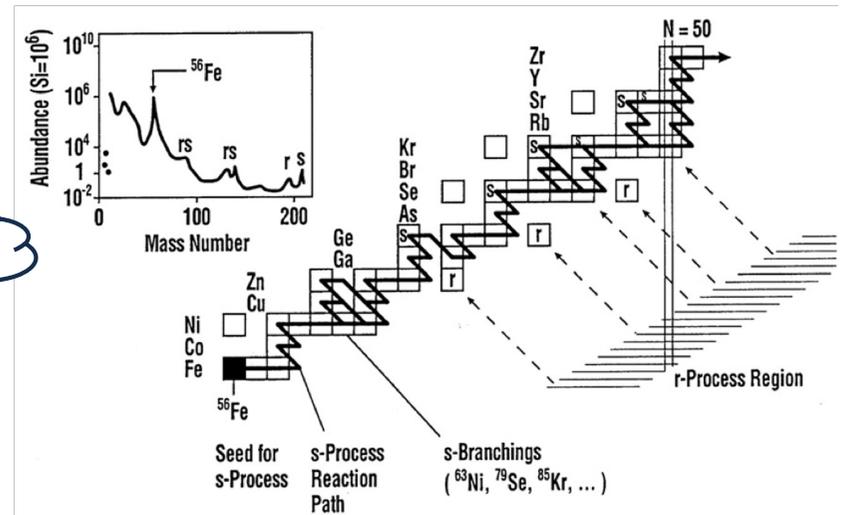
$^{186,187,188}\text{Os}$ (cosmochronometer), ^{197}Au (reference cross section), $^{24,25,26}\text{Mg}$, $^{33}\text{S}(n,\alpha)$, $^{14}\text{N}(n,p)$, $^{35}\text{Cl}(n,p)$,

$^{69,71}\text{Ga}$, $^{73,75,76}\text{Ge}$, $^{77,78,80}\text{Se}$ (weak component), $^{155,157,160}\text{Gd}$, $^{7}\text{Li}(n,p)$, $^{7}\text{Li}(n,\alpha)$ BBN

❖ Neutron Sources $^{22}\text{Ne}(\alpha,n)^{25}\text{Mg}$ and $^{13}\text{C}(\alpha,n)^{16}\text{O}$:

$n+^{25}\text{Mg}$, $n+^{16}\text{O}$

backup slides



$^{26}\text{Al}(n,p)$, $^{26}\text{Al}(n,\alpha)$

backup slides

Some recent results

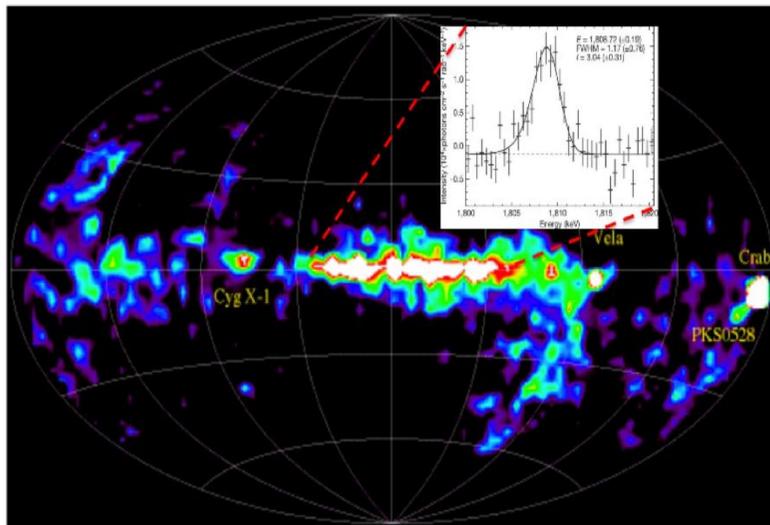
➤ $^{26}\text{Al}(n,p)$, $^{26}\text{Al}(n,\alpha)$

➤ $^{140}\text{Ce}(n,\gamma)$



The cosmic γ -ray emitter ^{26}Al

INTEGRAL Measured abundance 2.8(8) Solar Masses
 [R. Diehl, *Nature* **439**, 45(2006)]



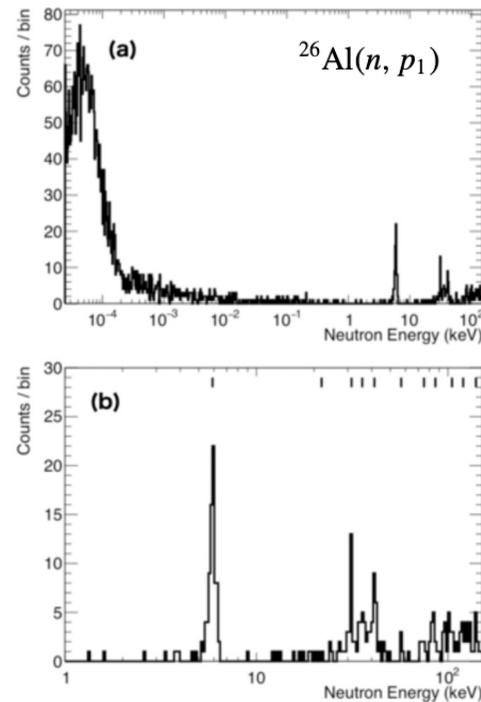
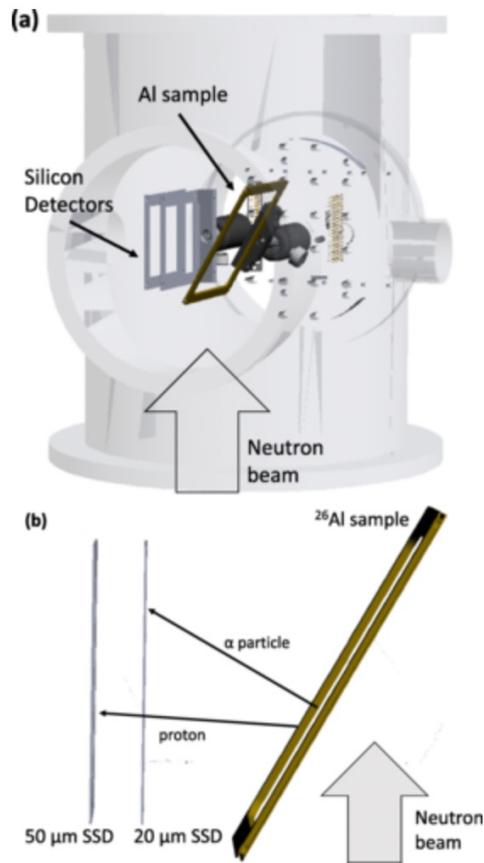
C Illiadis et al., *Ast. J. Supp.* **193**, 16 (2011) Sensitivity study of ^{26}Al abundance in Massive stars

FACTOR CHANGES OF FINAL $^{26}\text{Al}^g$ ABUNDANCE RESULTING FROM REACTION RATE VARIATIONS FOR CONVECTIVE SHELL C/NE BURNING^a, ASSUMING FIVE SPECIES OF ^{26}Al

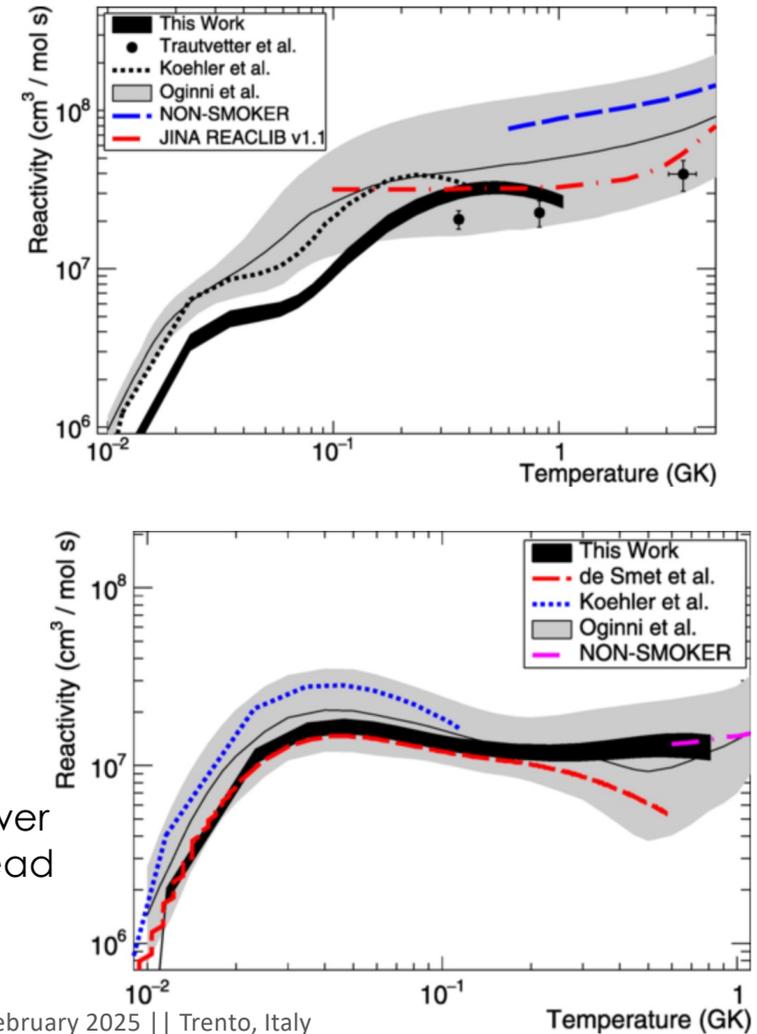
Reaction ^b	Rate multiplied by						Source ^c	Uncertainty ^d
	100	10	2	0.5	0.1	0.01		
$^{26}\text{Al}^g(\text{n,p})^{26}\text{Mg}$	0.017	0.16	0.63	1.3	1.9	2.0	present	
$^{25}\text{Mg}(\text{p},\gamma)^{26}\text{Al}^g$	2.9	5.4	1.5	0.63	0.35	0.29	il10	5%
$^{25}\text{Mg}(\text{p},\gamma)^{26}\text{Al}^m$	6.7	3.0	0.75	0.71	il10	6%
$^{26}\text{Al}^g(\text{n},\alpha)^{23}\text{Na}$	0.12	0.54	present	
$^{26}\text{Al}^m(\text{n,p})^{26}\text{Mg}$	0.58	present	

→ $^{26}\text{Al}(\text{n,p})$ and $^{26}\text{Al}(\text{n},\alpha)$ reaction rates represent critical uncertainties for ^{26}Al material processed by explosive and convective burning in massive stars and ejected into the ISM by core collapse supernovae

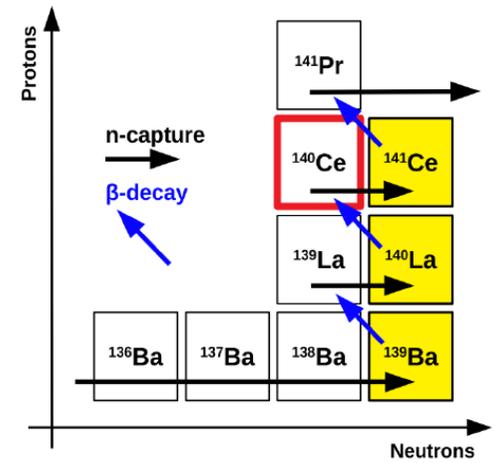
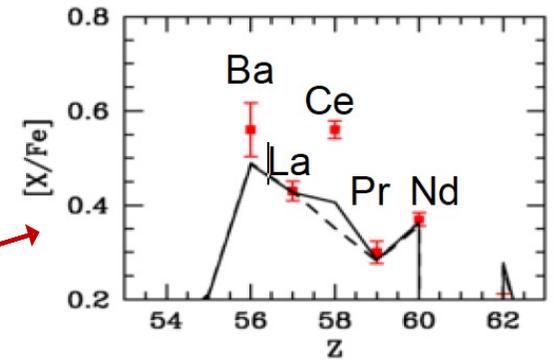
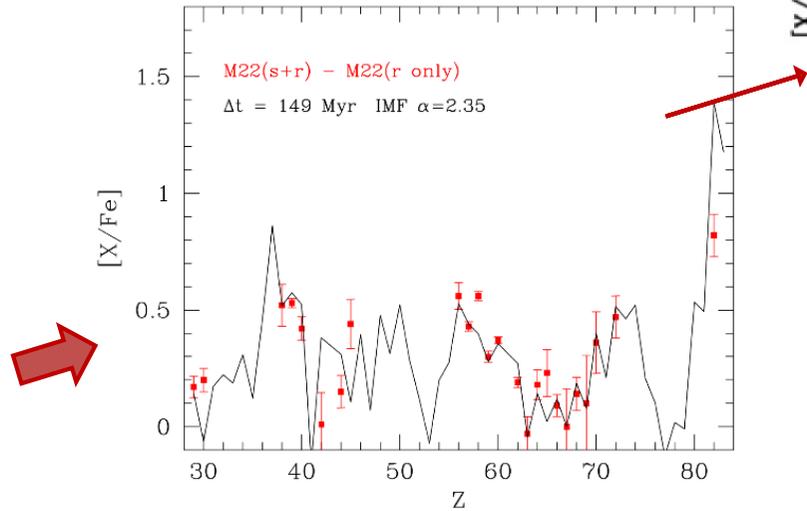
The cosmic γ -ray emitter ^{26}Al



The astrophysical reactivities are lower than literature data, which would lead to a higher destruction of ^{26}Al



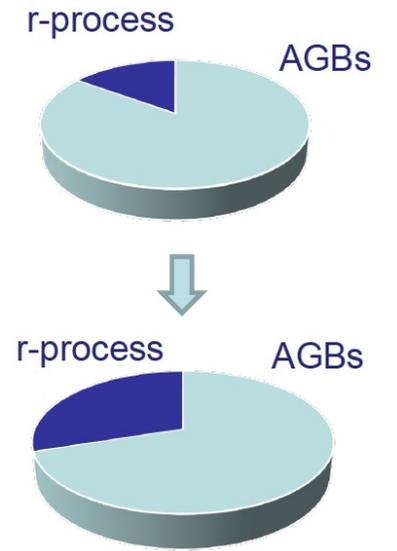
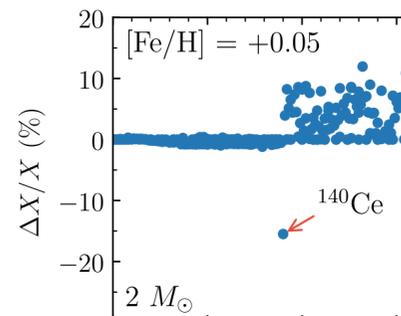
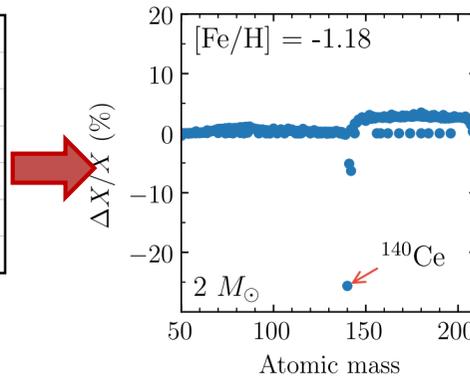
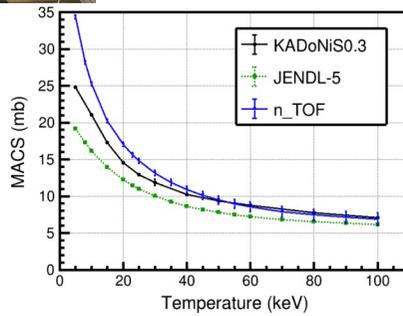
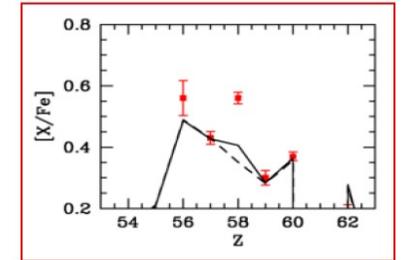
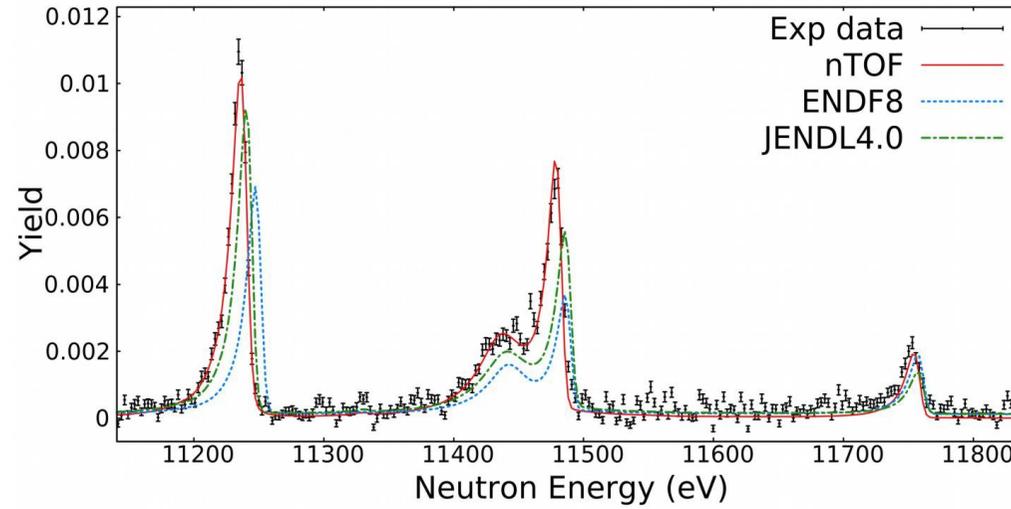
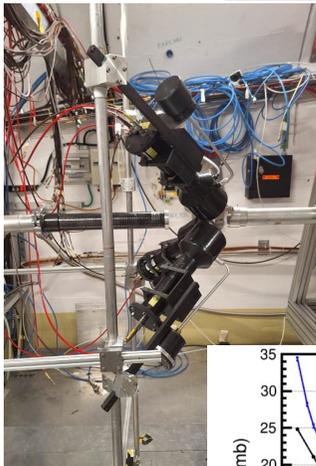
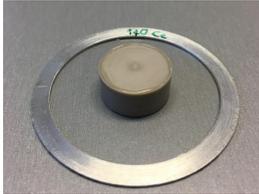
$^{140}\text{Ce}(n,\gamma)$ and the M22 globular cluster



- Abundances in s-process enhanced stars in M22 show a strong overproduction of Ce wrt AGB-models (Strainero, 2014)
- Ce abundance should be dominated by ^{140}Ce , 80% s process (Prantzos et al. 2020)
- Neutron-magic (very low CS) implies that previous CS determinations may be affected by neutron-sensitivity effects

$^{140}\text{Ce}(n,\gamma)$ and the M22 globular cluster

CeO₂ powder
 ^{140}Ce (99.4%)
 10 grams.



ALMA MATER STUDIORUM
 UNIVERSITÀ DI BOLOGNA

One of the ongoing projects: RAMEN



rame + n = *ramen*
Copper is "rame" in italian

➤ $^{63,65}\text{Cu}(n,\gamma)$ and (n,tot)





$^{63,65}\text{Cu}(n,\gamma)$ and $^{63,65}\text{Cu}(n,\text{tot})$

How copper was produced?

Not clear! Candidates:

1. **Weak s process (Massive stars)**
2. **Main s process (AGB)**
3. **SNe Ia**
4. **SNe II**

^{61}Zn 1.48 m β^+	^{62}Zn 9.19 h β^+	^{63}Zn 38.47 m β^+	^{64}Zn 48.63 59 mb	^{65}Zn 243.63 d 162 mb, β^+	^{66}Zn 27.9 35 mb	^{67}Zn 4.1 153 mb
^{60}Cu 23.70 m β^+	^{61}Cu 3.33 h β^+	^{62}Cu 9.67 m β^+	^{63}Cu 69.17 94 mb	^{64}Cu 12.70 h β^+	^{65}Cu 30.83 41 mb	^{66}Cu 5.12 m β^-
^{59}Ni 75.99 ka 87 mb, β^+	^{60}Ni 26.223 30 mb	^{61}Ni 1.14 82 mb	^{62}Ni 3.634 22.3 mb	^{63}Ni 100.11 a 31 mb, β^-	^{64}Ni 0.926 8.7 mb	^{65}Ni 2.52 h β^-
^{58}Co 70.86 d β^+	^{59}Co 100 38 mb	^{60}Co 5.27 a β^-	^{61}Co 1.65 h β^-	^{62}Co 1.50 m β^-	^{63}Co 27.40 s β^-	^{64}Co 300.00 ms β^-
^{57}Fe 2.119 40 mb	^{58}Fe 0.282 12.1 mb	^{59}Fe 44.50 d β^-	^{60}Fe 1.50 Ma β^-	^{61}Fe 5.98 m β^-	^{62}Fe 1.13 m β^-	^{63}Fe 6.01 s β^-



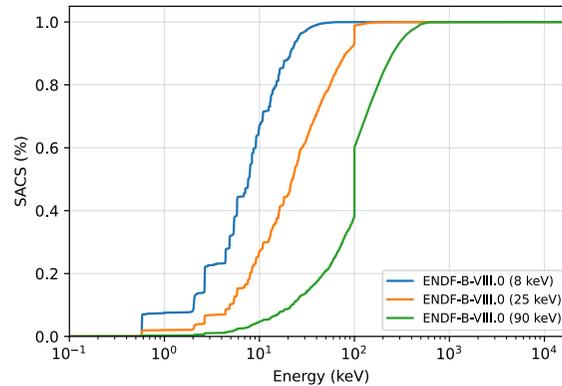
$^{63,65}\text{Cu}(n,\gamma)$ and $^{63,65}\text{Cu}(n,\text{tot})$

How copper was produced?

Not clear! Candidates:

1. Weak s process (Massive stars)
2. Main s process (AGB)
3. SNe Ia
4. SNe II

The s process requires
MACS @ $kT=8, 25, 90$ keV



^{61}Zn 1.48 m β^+	^{62}Zn 9.19 h β^+	^{63}Zn 38.47 m β^+	^{64}Zn 48.63 59 mb	^{65}Zn 243.63 d 162 mb, β^+	^{66}Zn 27.9 35 mb	^{67}Zn 4.1 153 mb
^{60}Cu 23.70 m β^+	^{61}Cu 3.33 h β^+	^{62}Cu 9.67 m β^+	^{63}Cu 69.17 94 mb	^{64}Cu 12.70 h β^+	^{65}Cu 30.83 41 mb	^{66}Cu 5.12 m β^-
^{59}Ni 75.99 ka 87 mb, β^+	^{60}Ni 26.223 30 mb	^{61}Ni 1.14 82 mb	^{62}Ni 3.634 22.3 mb	^{63}Ni 100.11 a 31 mb, β^-	^{64}Ni 0.926 8.7 mb	^{65}Ni 2.52 h β^-
^{58}Co 36 d β^+	^{59}Co 100 38 mb	^{60}Co 5.27 a β^-	^{61}Co 1.65 h β^-	^{62}Co 1.50 m β^-	^{63}Co 27.40 s β^-	^{64}Co 300.00 ms β^-
^{57}Fe 19 mb	^{58}Fe 0.282 12.1 mb	^{59}Fe 44.50 d β^-	^{60}Fe 1.50 Ma β^-	^{61}Fe 5.98 m β^-	^{62}Fe 1.13 m β^-	^{63}Fe 6.01 s β^-

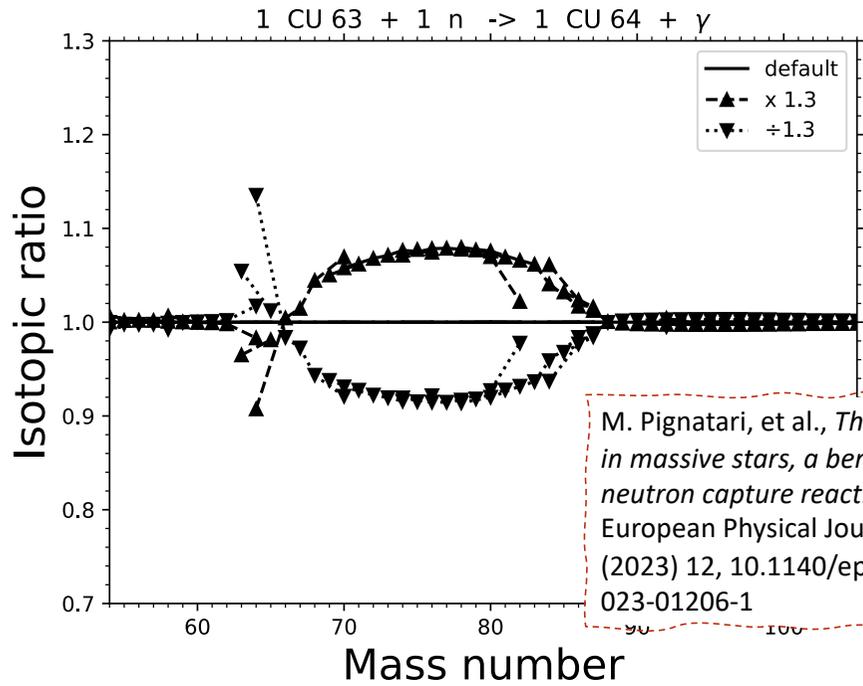
With accurate determination of Cu MACS, it would be possible to clarify what is the s-process contribution to Cu. Once this is done, it will be possible to constrain the Cu production by other nucleosynthesis processes, where stellar and nuclear uncertainties are much larger.

$^{63,65}\text{Cu}(n,\gamma)$ and $^{63,65}\text{Cu}(n,\text{tot})$

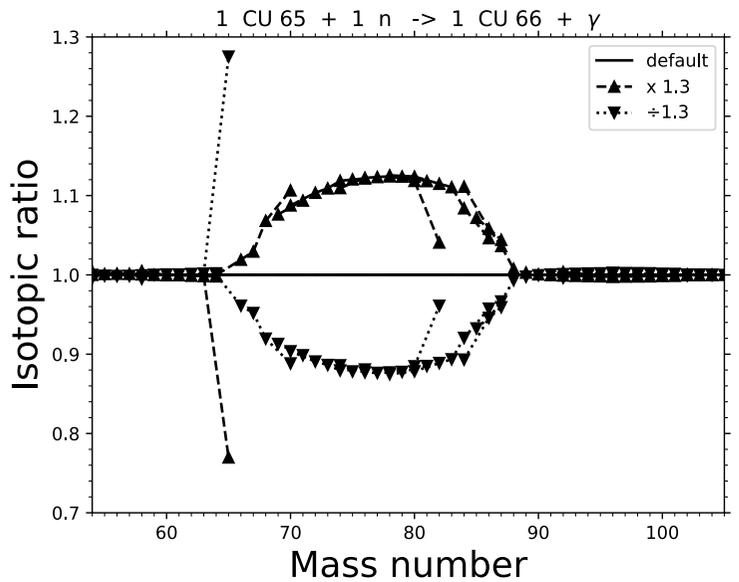
Propagation effect in the weak s-process nucleosynthesis

Impact of $^{63,65}\text{Cu}(n,\gamma)$ cross sections on the efficiency for the production of elements heavier than Cu

^{61}Zn 1.48 m β^+	^{62}Zn 9.19 h β^+	^{63}Zn 38.47 m β^+	^{64}Zn 48.63 59 mb	^{65}Zn 243.63 d 162 mb, β^+	^{66}Zn 27.9 35 mb	^{67}Zn 4.1 153 mb
^{60}Cu 23.70 m β^+	^{61}Cu 3.33 h β^+	^{62}Cu 9.67 m β^+	^{63}Cu 69.17 94 mb	^{64}Cu 12.70 h β^+	^{65}Cu 30.83 41 mb	^{66}Cu 5.12 m β^-
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						^{64}Co 300.00 ms β^-
						^{63}Fe 6.01 s β^-



M. Pignatari, et al., *The s process in massive stars, a benchmark for neutron capture reaction rates*, European Physical Journal A **59** (2023) 12, 10.1140/epja/s10050-023-01206-1



$^{63,65}\text{Cu}(n,\gamma)$ and $^{63,65}\text{Cu}(n,\text{tot})$

▼ **Recommended MACS30** (Maxwellian Averaged Cross Section @ 30keV)

$^{63}\text{Cu}(n,\gamma)^{64}\text{Cu}$

Total MACS at 30keV: 60.1 ± 6.2 mb

Cross sections do not include stellar enhancement factors!

▼ **History**

Version	Total MACS [mb]	Partial to gs [mb]	Partial to isomer [mb]
1.0	60.1 ± 6.2	-	-
0.3	55.6 ± 2.2	-	-
0.0	94 ± 10	-	-

(Version 0.0 corresponds to Bao et al.)

▼ **Comment**

New rec. value is from [HKU08](#), renormalized by $632 \text{ mb}/586 \text{ mb} = 1.0785$, and recalculated with normalized energy dependencies of [tendl15](#), [endfb71](#), [jendl40](#). Uncertainty is the deviation between different evaluations plus 4% exp. uncertainty from [HKU08](#). **Note the large deviation between the activation measurement and the TOF measurements. More investigation needed!**

Last review: April 2017

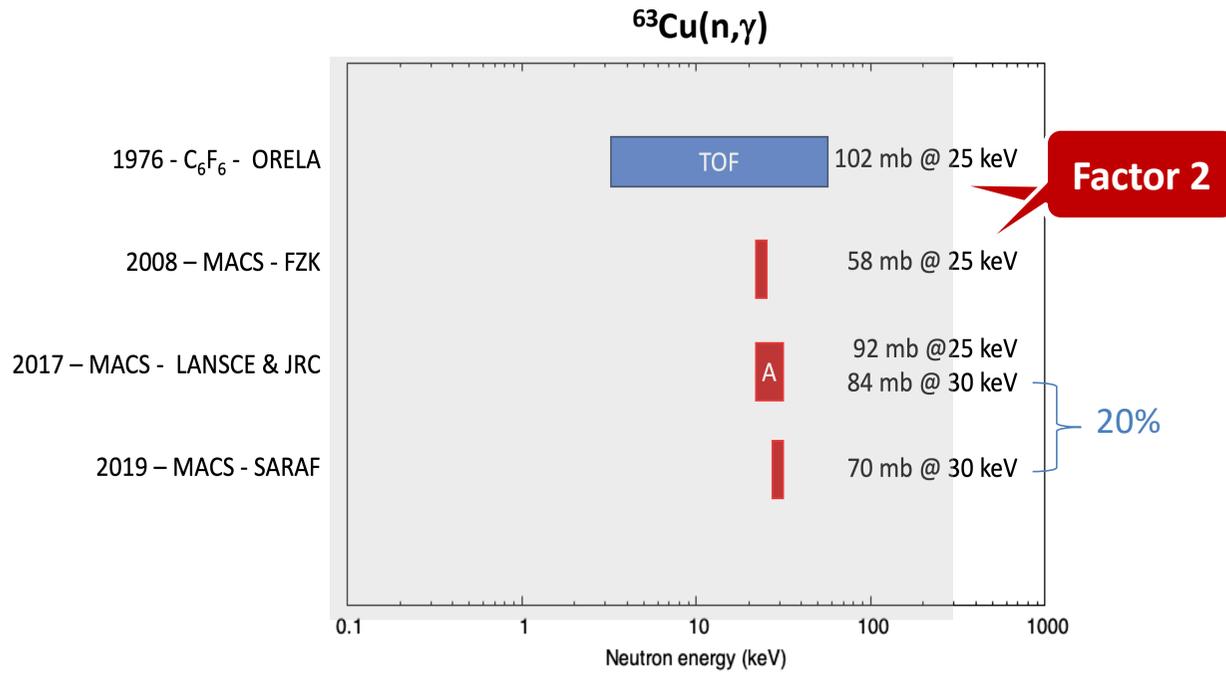
View Maxwellian-Averaged (n,g) Cross Section

Isotope

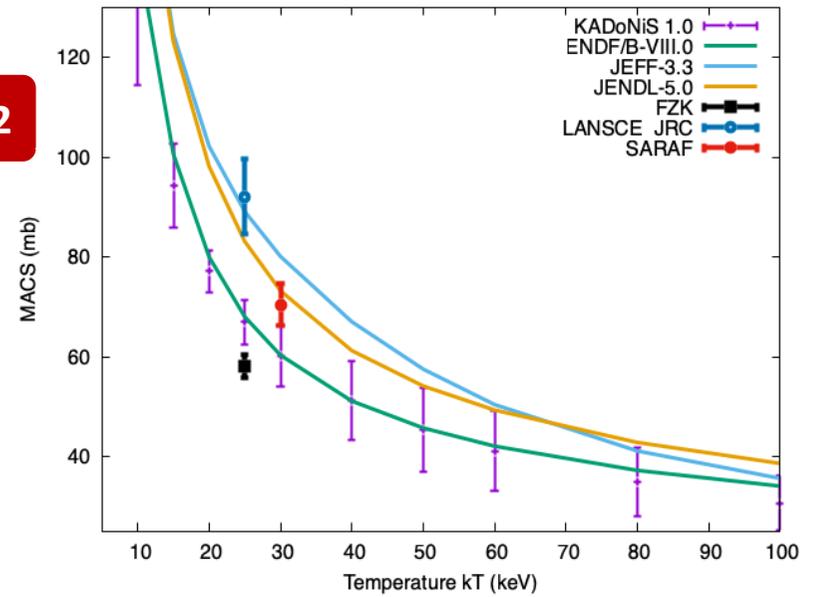
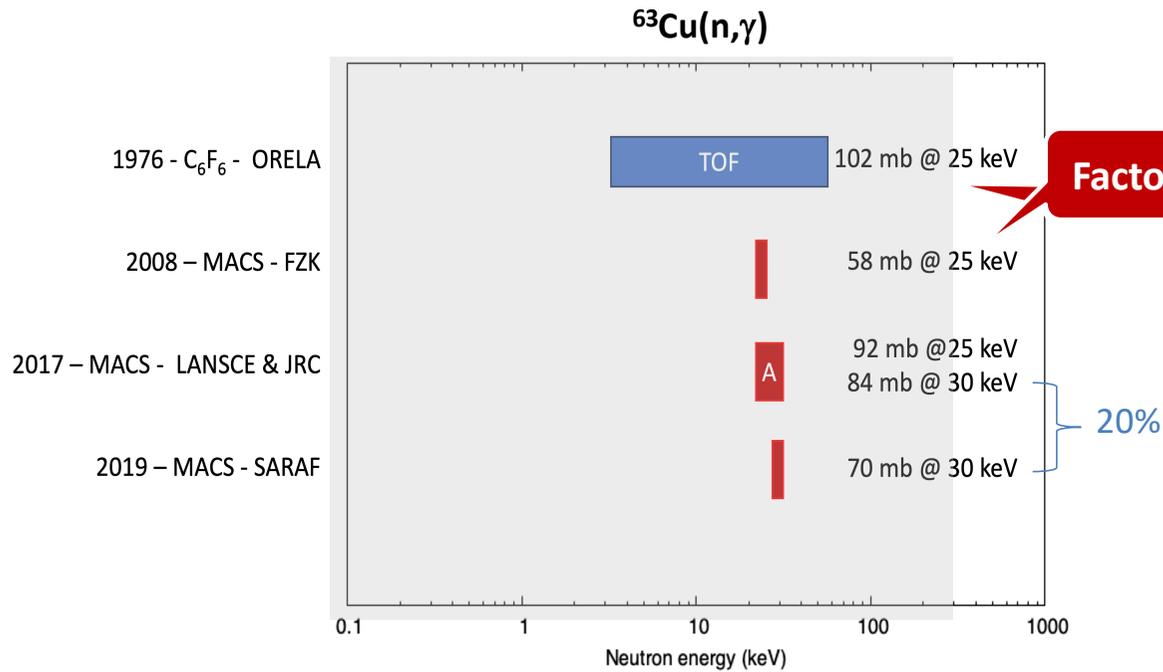
(Examples: Ba138, Ta180m, Se.)



$^{63,65}\text{Cu}(n,\gamma)$ and $^{63,65}\text{Cu}(n,\text{tot})$

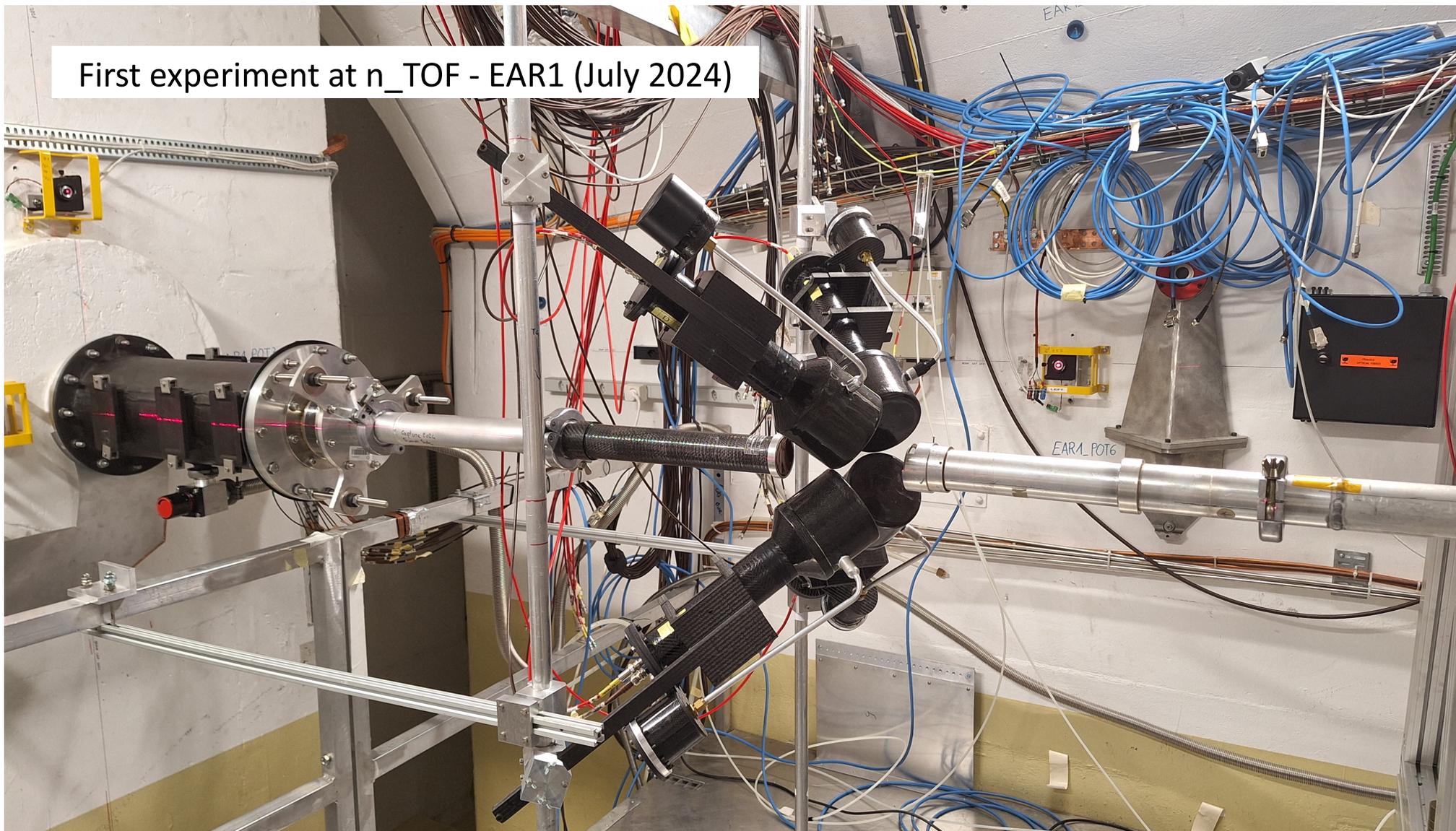


$^{63,65}\text{Cu}(n,\gamma)$ and $^{63,65}\text{Cu}(n,\text{tot})$



- Time-of-flight data needed for $E_n < 300$ keV
- n_TOF can accurately cover this energy region

First experiment at n_TOF - EAR1 (July 2024)



First experiment at n_TOF - EAR1 (July 2024)

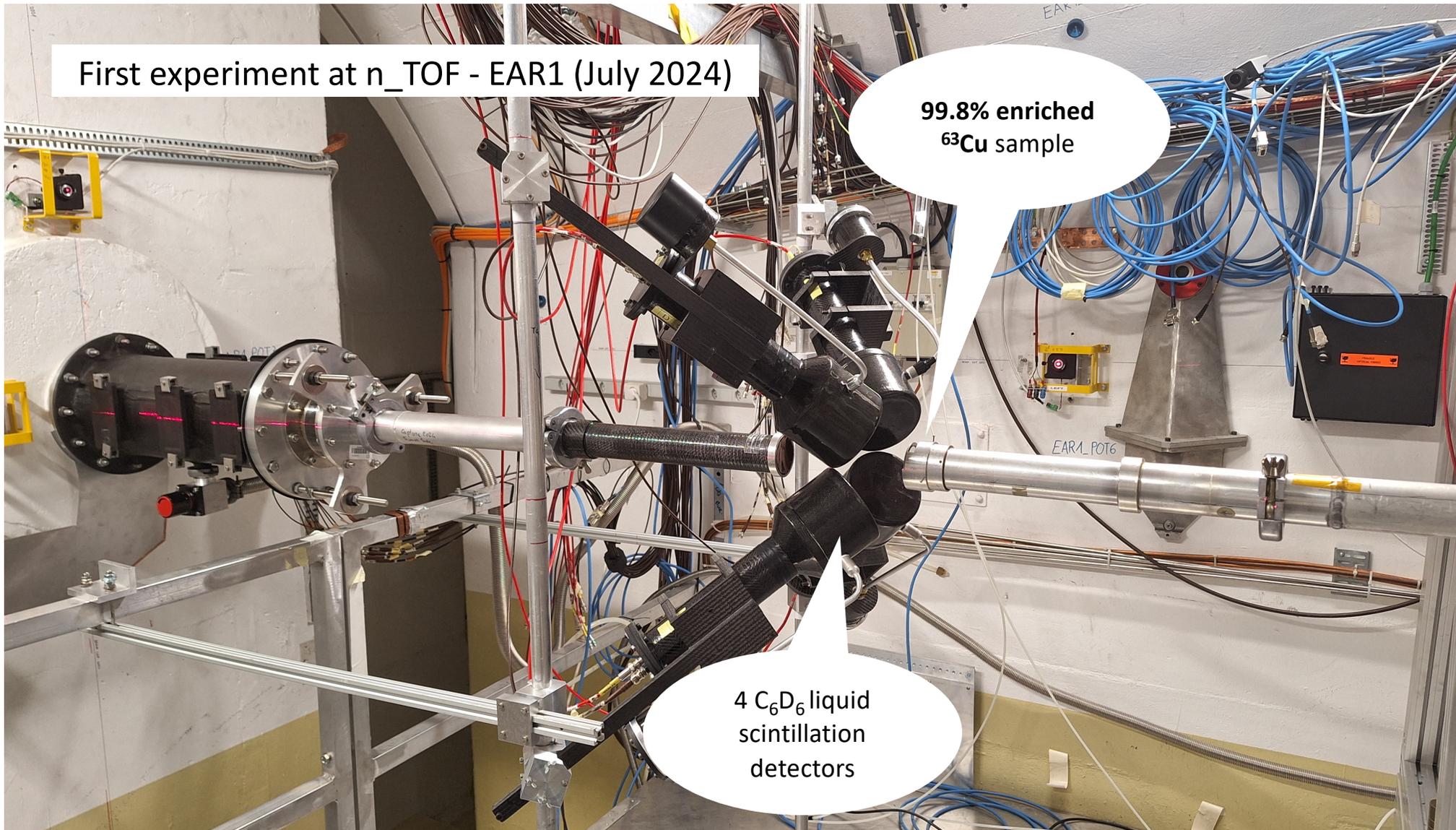
99.8% enriched
 ^{63}Cu sample



First experiment at n_TOF - EAR1 (July 2024)

99.8% enriched ^{63}Cu sample

4 C_6D_6 liquid scintillation detectors



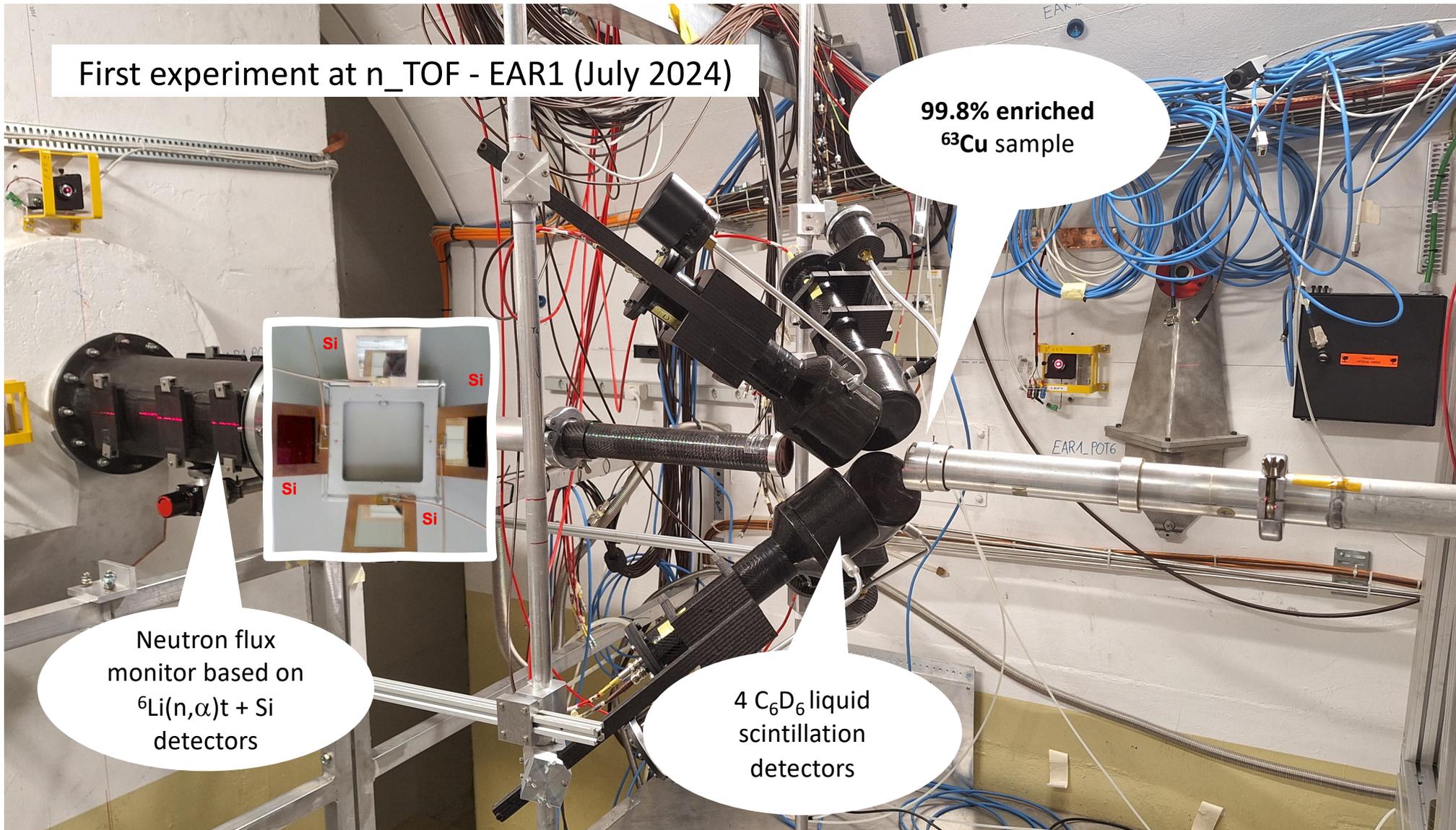
First experiment at n_TOF - EAR1 (July 2024)

99.8% enriched ^{63}Cu sample



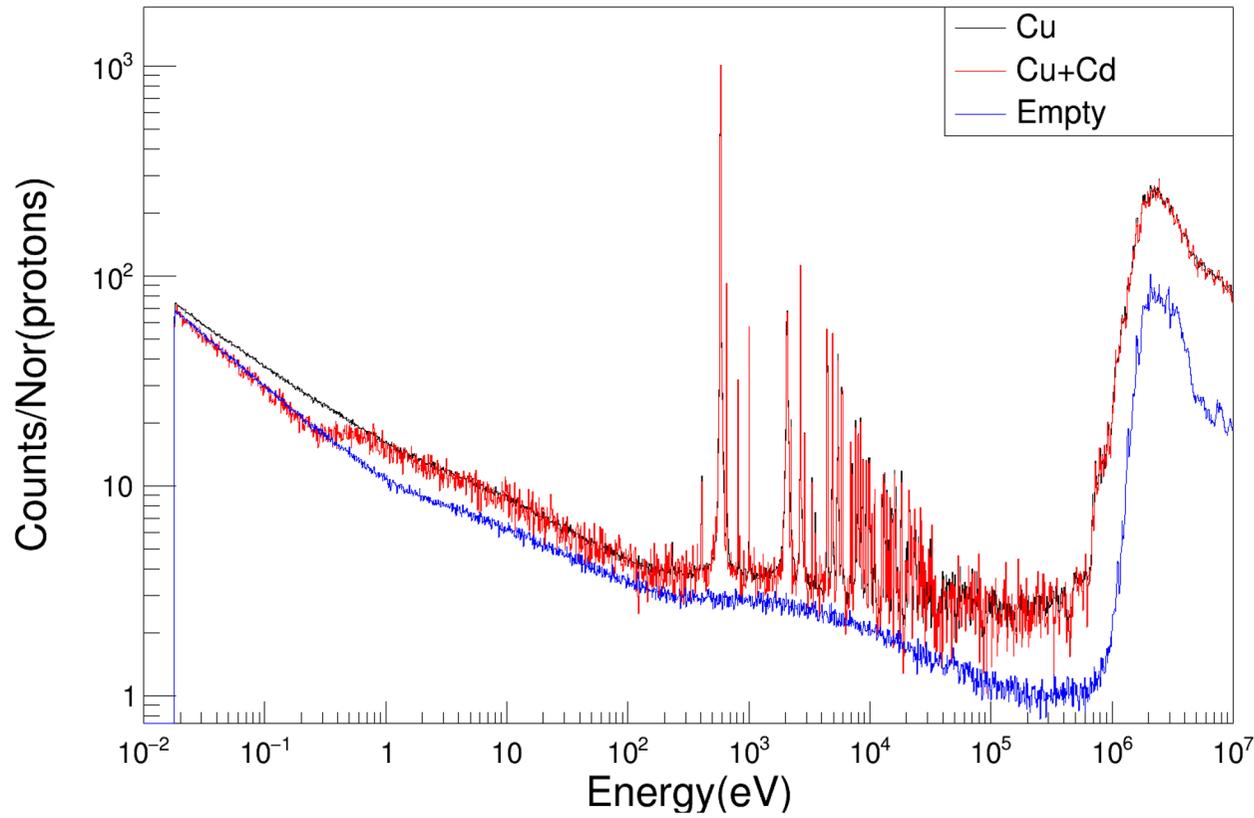
Neutron flux monitor based on $^6\text{Li}(n,\alpha)t + \text{Si}$ detectors

4 C_6D_6 liquid scintillation detectors



$^{63,65}\text{Cu}(n,\gamma)$ and $^{63,65}\text{Cu}(n,\text{tot})$

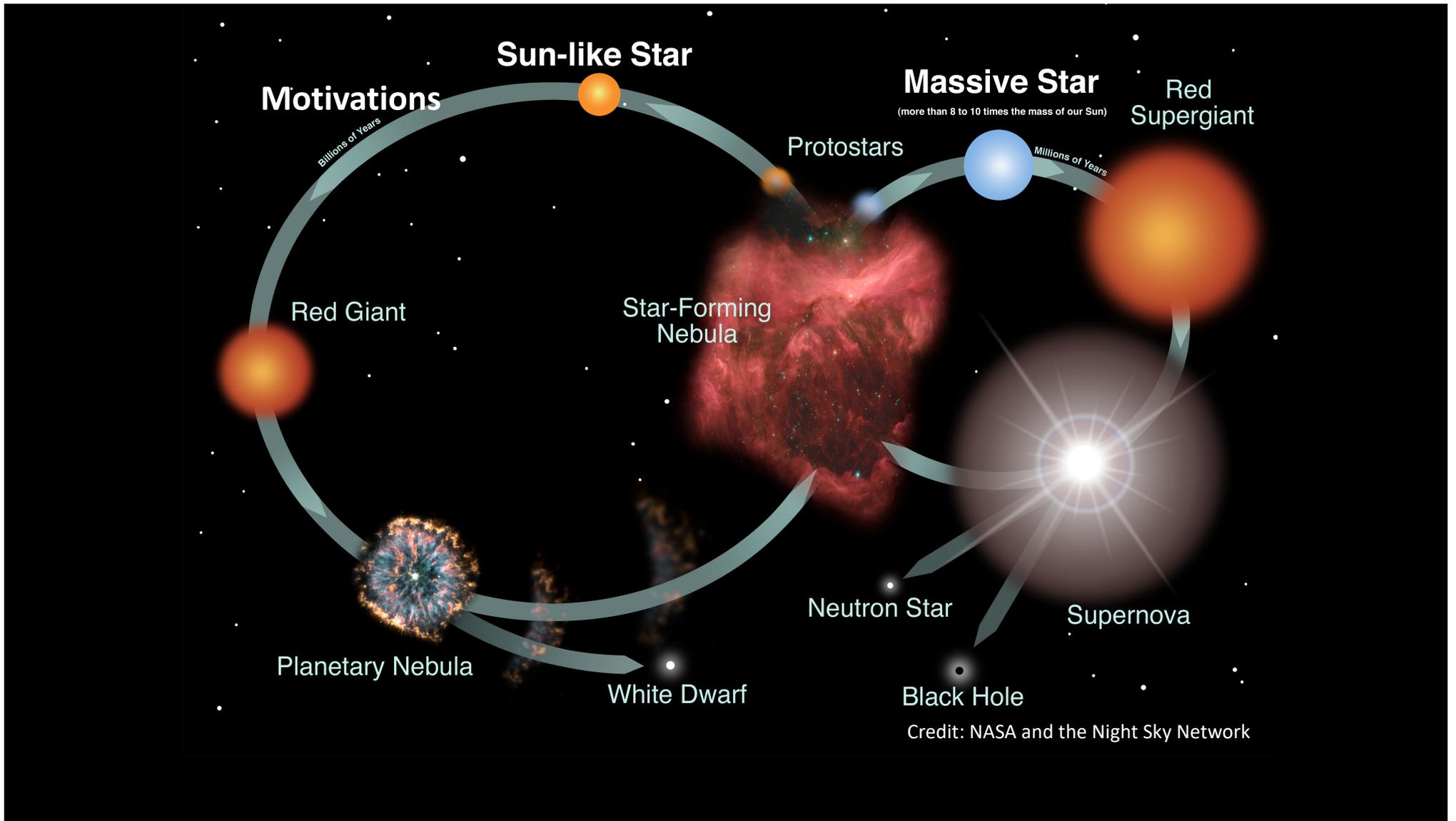
C_6D_6 counts normalized to neutron intensity

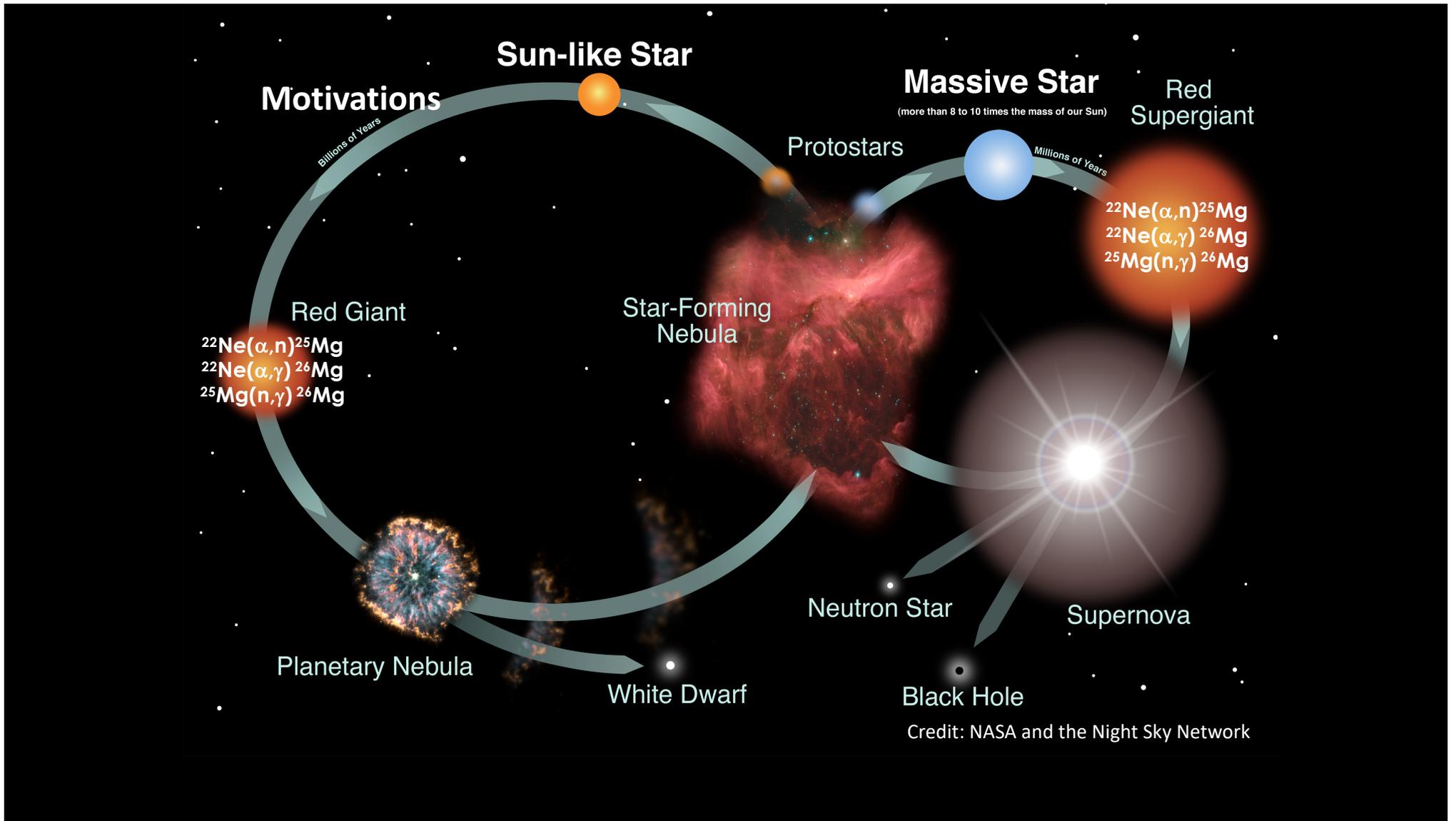


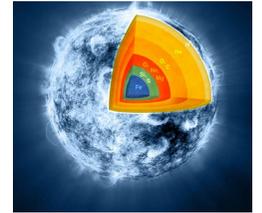
Analysis ongoing

Some future perspectives

- $^{25}\text{Mg}(n,\gamma)$
- (cyclic) Activation @ NEAR







$^{25}\text{Mg}(n,\gamma)$ for neutron source reaction in stars

➤ NEUTRON POISON:

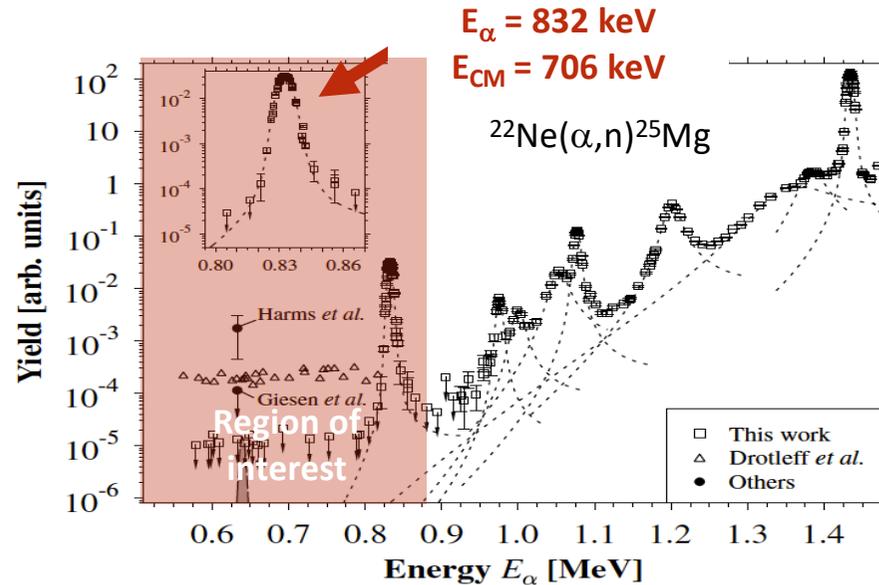
- $^{25,26}\text{Mg}$ are the most important neutron poisons due to neutron capture on Mg stable isotopes, i.e. $^{25,26}\text{Mg}(n,\gamma)$, in competition with neutron capture on ^{56}Fe (the basic s-process seed for the production of heavier isotopes).

➤ CONSTRAINTS for $^{22}\text{Ne}(\alpha,n)^{25}\text{Mg}$ and $^{22}\text{Ne}(\alpha,\gamma)^{26}\text{Mg}$:

- $^{22}\text{Ne}(\alpha,n)^{25}\text{Mg}$ is one of the most important neutron source in Red Giant stars. Its reaction rate is very uncertain because of the poorly known property of the states in ^{26}Mg . From neutron measurements the energy, J^π and **energy** of ^{26}Mg states can be deduced, in addition to Γ_γ and Γ_n .

$^{25}\text{Mg}(n,\gamma)$

Data in the literature



➤ CONSTRAINTS for $^{22}\text{Ne}(\alpha,n)^{25}\text{Mg}$ and $^{22}\text{Ne}(\alpha,\gamma)^{26}\text{Mg}$:

- $^{22}\text{Ne}(\alpha,n)^{25}\text{Mg}$ is one of the most important neutron source in Red Giant stars. Its reaction rate is very uncertain because of the poorly known property of the states in ^{26}Mg . From neutron measurements the energy, J^π and **energy** of ^{26}Mg states can be deduced, in addition to Γ_γ and Γ_n .

$^{25}\text{Mg}(n,\gamma)$ for neutron source reaction in stars



Reevaluation of the $^{22}\text{Ne}(\alpha,\gamma)^{26}\text{Mg}$ and $^{22}\text{Ne}(\alpha,n)^{25}\text{Mg}$ reaction rates

Philip Adsley^{1,2,3,*}, Umberto Battino^{4,1}, Andreas Best^{5,6}, Antonio Cacioli^{7,8}, Alessandra Guglielmetti⁹, Gianluca Imbriani^{5,6}, Heshani Jayatissa¹⁰, Marco La Cognata¹¹, Livio Lamia^{12,11,13} *et al.*

Show more

Phys. Rev. C **103**, 015805 – Published 19 January, 2021

DOI: <https://doi.org/10.1103/PhysRevC.103.015805>

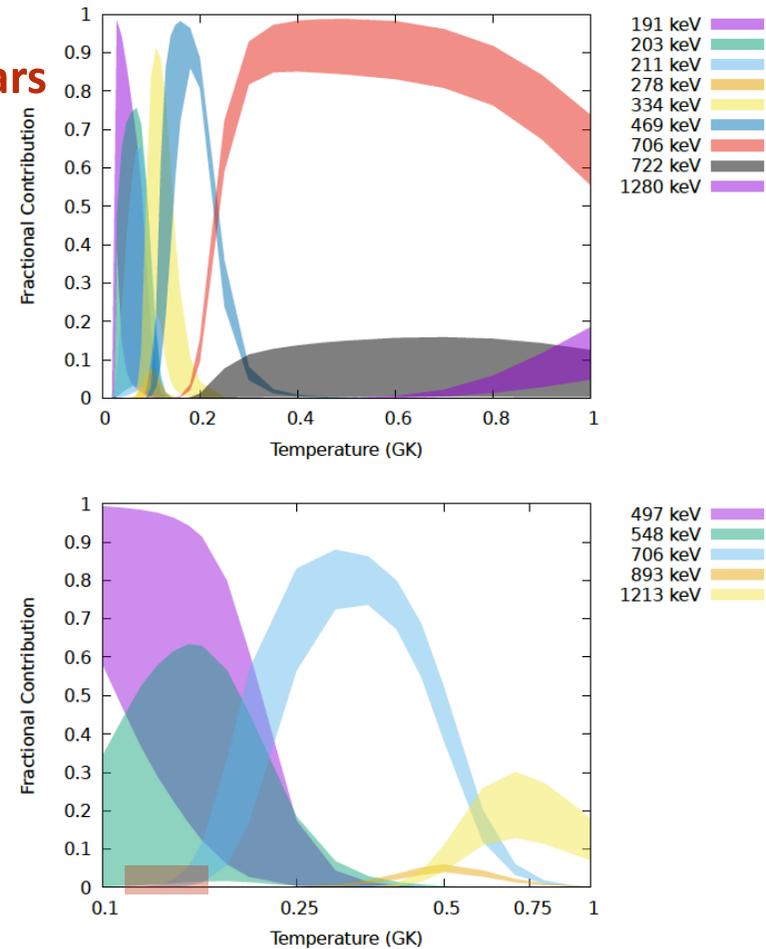


FIG. 1. Fractional contributions of selected resonances to the (top) $^{22}\text{Ne}(\alpha,\gamma)^{26}\text{Mg}$ and (bottom) $^{22}\text{Ne}(\alpha,n)^{25}\text{Mg}$ reaction rates. These fractional contributions are for the recommended

$^{25}\text{Mg}(n,\gamma)$ for neutron source reaction in stars



Reevaluation of the $^{22}\text{Ne}(\alpha,\gamma)^{26}\text{Mg}$ and $^{22}\text{Ne}(\alpha,n)^{25}\text{Mg}$ reaction rates

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Show more

Phys. Rev. C **103**, 015805 – Published 19 January, 2021

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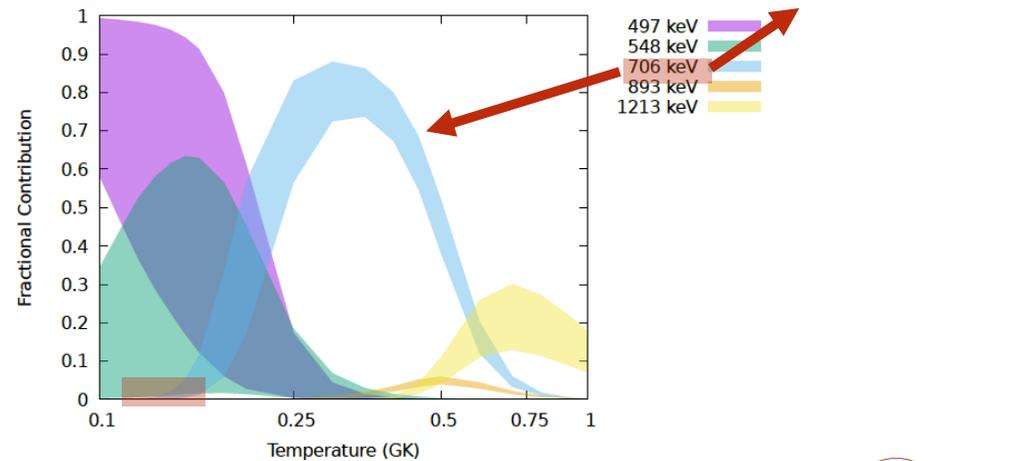
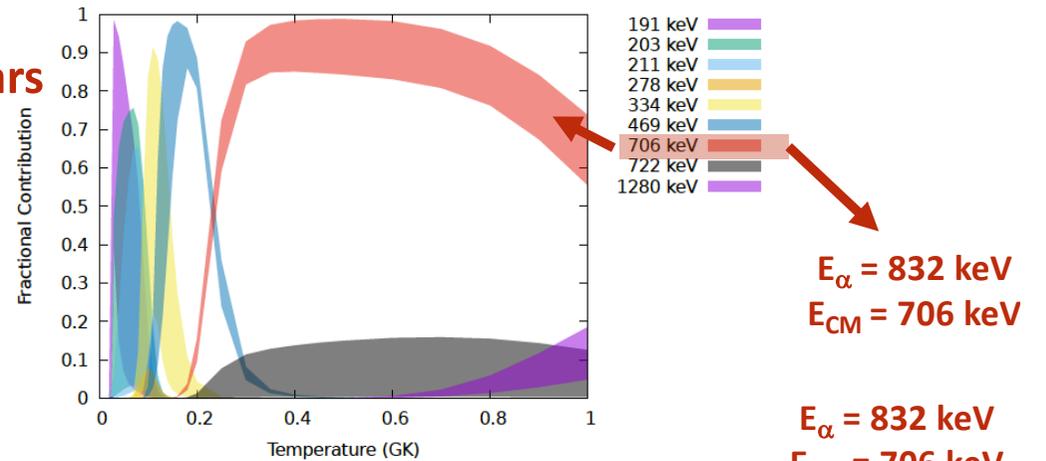
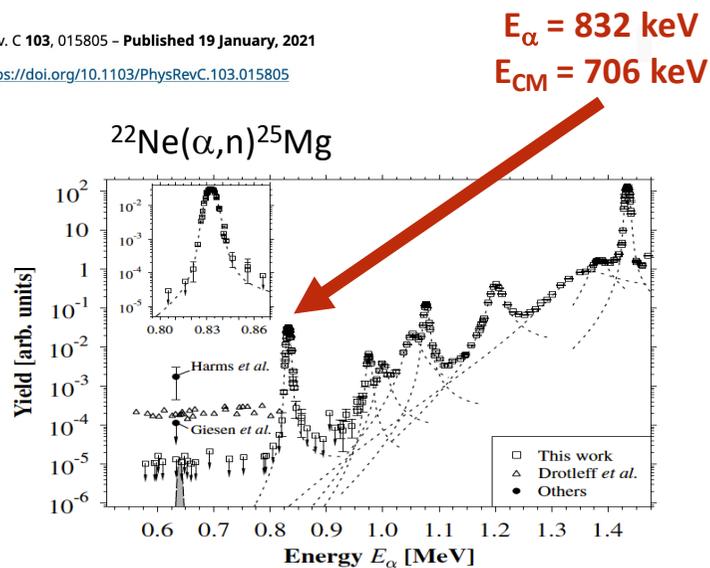


FIG. 1. Fractional contributions of selected resonances to the (top) $^{22}\text{Ne}(\alpha,\gamma)^{26}\text{Mg}$ and (bottom) $^{22}\text{Ne}(\alpha,n)^{25}\text{Mg}$ reaction rates. These fractional contributions are for the recommended



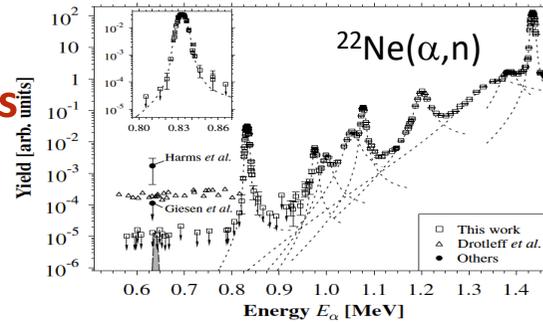
$^{25}\text{Mg}(n,\gamma)$ for neutron source reaction in stars

Resonance strength $^{22}\text{Ne}(\alpha,n)^{25}\text{Mg}$:

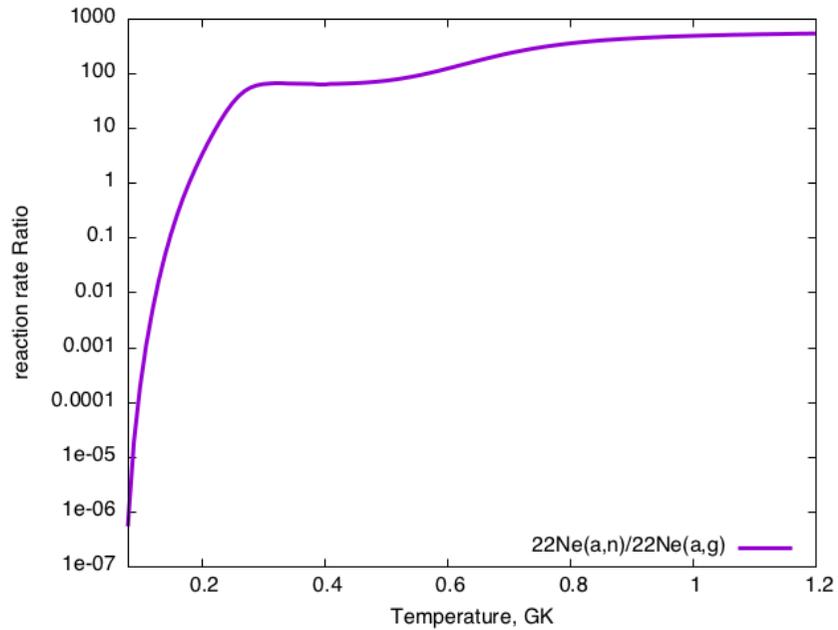
$$\omega_\alpha = g \Gamma_\alpha \Gamma_n / (\Gamma_\alpha + \Gamma_\gamma + \Gamma_n)$$

Resonance strength $^{22}\text{Ne}(\alpha,\gamma)^{26}\text{Mg}$:

$$\omega_\gamma = g \Gamma_\alpha \Gamma_\gamma / (\Gamma_\alpha + \Gamma_\gamma + \Gamma_n)$$



$^{22}\text{Ne}(\alpha,n) / ^{22}\text{Ne}(\alpha,\gamma)$



$^{25}\text{Mg}(n,\gamma)$ for neutron source reaction in stars

Resonance strength $^{22}\text{Ne}(\alpha,n)^{25}\text{Mg}$:

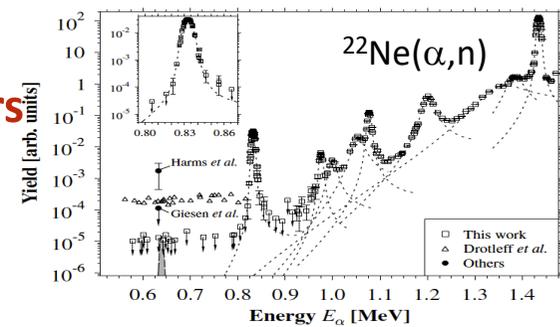
$$\omega_\alpha = g \Gamma_\alpha \Gamma_n / (\Gamma_\alpha + \Gamma_\gamma + \Gamma_n)$$

Resonance strength $^{22}\text{Ne}(\alpha,\gamma)^{26}\text{Mg}$:

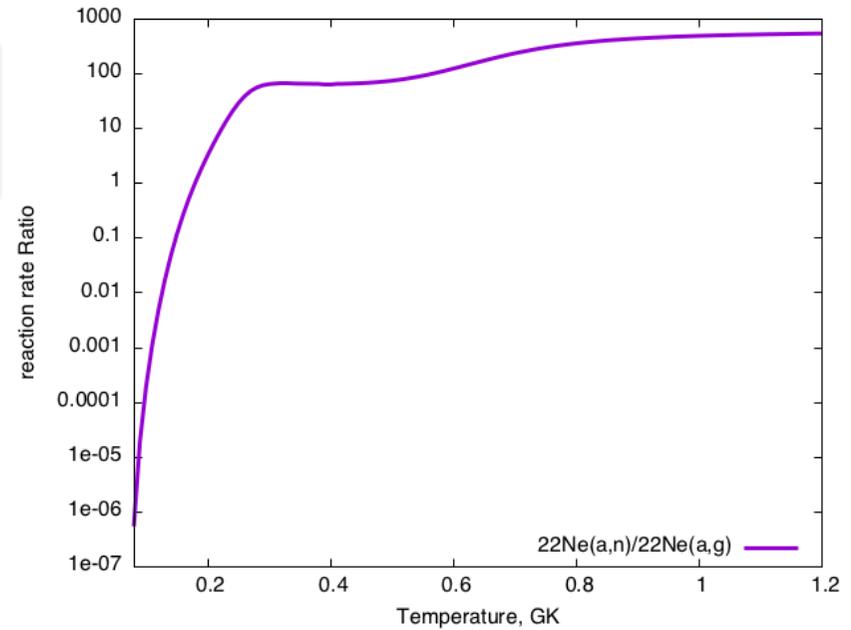
$$\omega_\gamma = g \Gamma_\alpha \Gamma_\gamma / (\Gamma_\alpha + \Gamma_\gamma + \Gamma_n)$$

$$\frac{\omega_\alpha}{\omega_\gamma} = \frac{\Gamma_n}{\Gamma_\gamma}$$

Resonance strength ratio and Reaction rate ratio independent of Γ_α



$^{22}\text{Ne}(\alpha,n) / ^{22}\text{Ne}(\alpha,\gamma)$



$^{25}\text{Mg}(n,\gamma)$ for neutron source reaction in stars

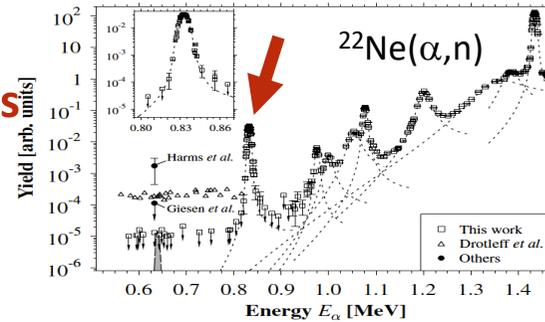
Resonance strength $^{22}\text{Ne}(\alpha,n)^{25}\text{Mg}$:

$$\omega_\alpha = g \Gamma_\alpha \Gamma_n / (\Gamma_\alpha + \Gamma_\gamma + \Gamma_n)$$

Resonance strength $^{22}\text{Ne}(\alpha,\gamma)^{26}\text{Mg}$:

$$\omega_\gamma = g \Gamma_\alpha \Gamma_\gamma / (\Gamma_\alpha + \Gamma_\gamma + \Gamma_n)$$

$$\frac{\omega_\alpha}{\omega_\gamma} = \frac{\Gamma_n}{\Gamma_\gamma}$$



$^{22}\text{Ne}(\alpha,n) / ^{22}\text{Ne}(\alpha,\gamma)$

Publication	YEAR	Result	comment
Shahina, PRC	2024	$\Gamma_n / \Gamma_\gamma = 2.85(71)$	ω_α res. strength
M. Wiescher, EPJA	2023	$\Gamma_n = 0.4 - 1.0$ eV $\Gamma_\gamma = 1.33$ eV	Re-evaluation
Y. Chen, PRC	2021	$\Gamma_n = 0.4$ eV $\Gamma_\gamma = 1.33$ eV	$^{25}\text{Mg}(d,p)^{26}\text{Mg}$ transfer
S. Ota, PLB	2020	$\Gamma_n / \Gamma_\gamma = 1.14(26)$	transfer

$^{25}\text{Mg}(n,\gamma)$ for neutron source reaction in stars

Resonance strength $^{22}\text{Ne}(\alpha,n)^{25}\text{Mg}$:

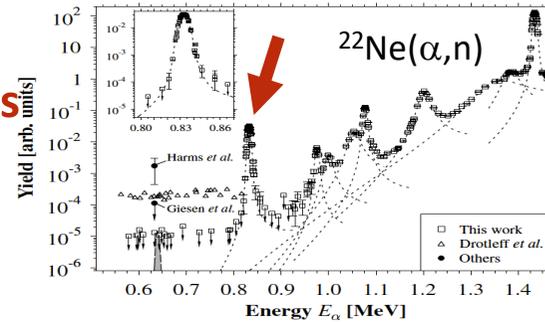
$$\omega_\alpha = g \Gamma_\alpha \Gamma_n / (\Gamma_\alpha + \Gamma_\gamma + \Gamma_n)$$

Resonance strength $^{22}\text{Ne}(\alpha,\gamma)^{26}\text{Mg}$:

$$\omega_\gamma = g \Gamma_\alpha \Gamma_\gamma / (\Gamma_\alpha + \Gamma_\gamma + \Gamma_n)$$

$$\frac{\omega_\alpha}{\omega_\gamma} = \frac{\Gamma_n}{\Gamma_\gamma}$$

Neutron width Γ_n and γ -ray width Γ_γ can be deduced from $n + ^{25}\text{Mg}$ experiments

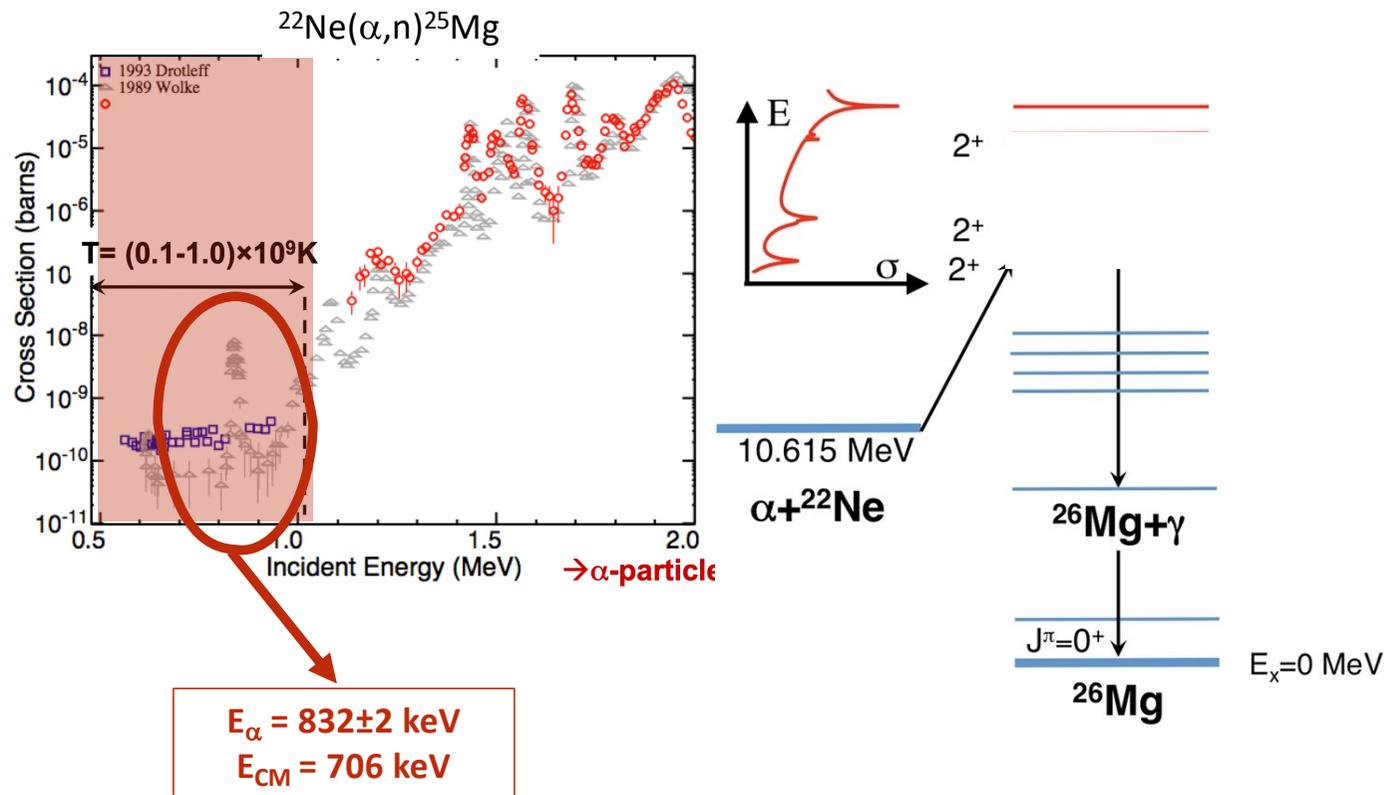


$$^{22}\text{Ne}(\alpha,n) / ^{22}\text{Ne}(\alpha,\gamma)$$

$$\sigma_\gamma(E_n) = g \frac{\pi}{k_n^2} \frac{\Gamma_n \Gamma_\gamma}{(E_n - E_R)^2 + (\Gamma/2)^2}$$

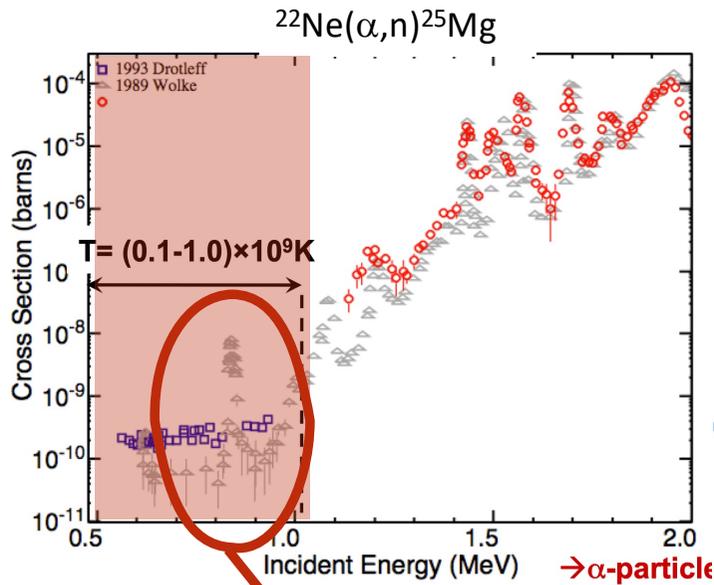
$^{25}\text{Mg}(n,\gamma)$ cross section

$^{25}\text{Mg}(n,\gamma)$ for neutron source reaction in stars

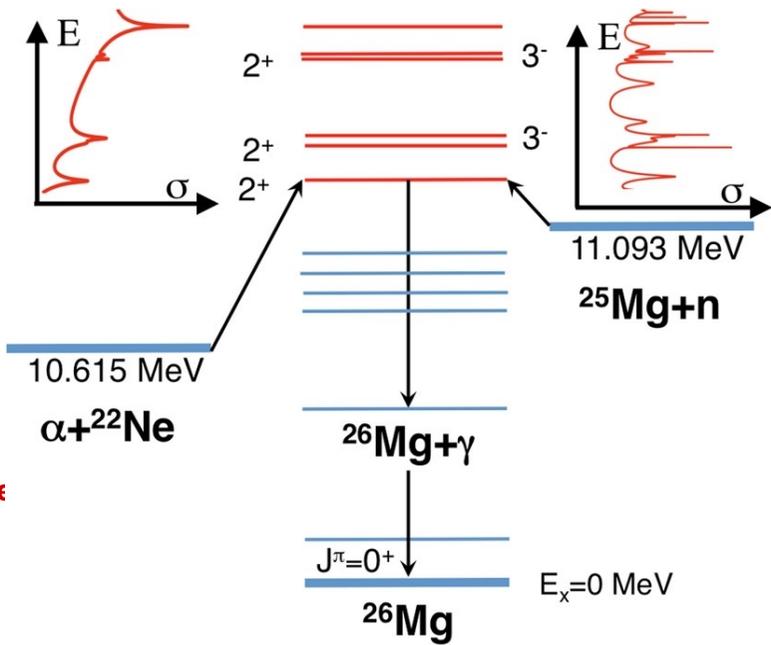


$^{25}\text{Mg}(n,\gamma)$ for neutron source reaction in stars

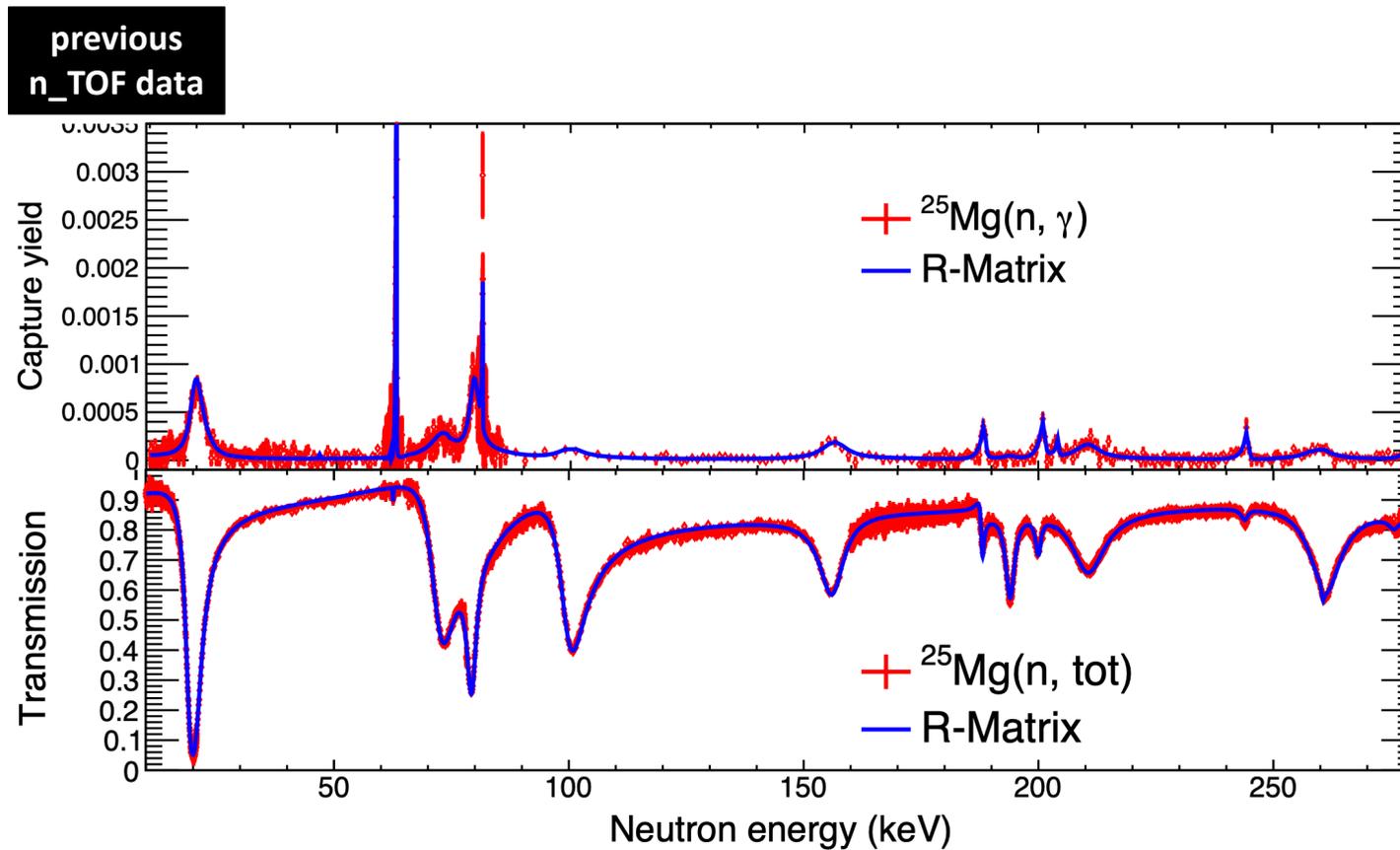
^{26}Mg levels via $n + ^{25}\text{Mg}$:



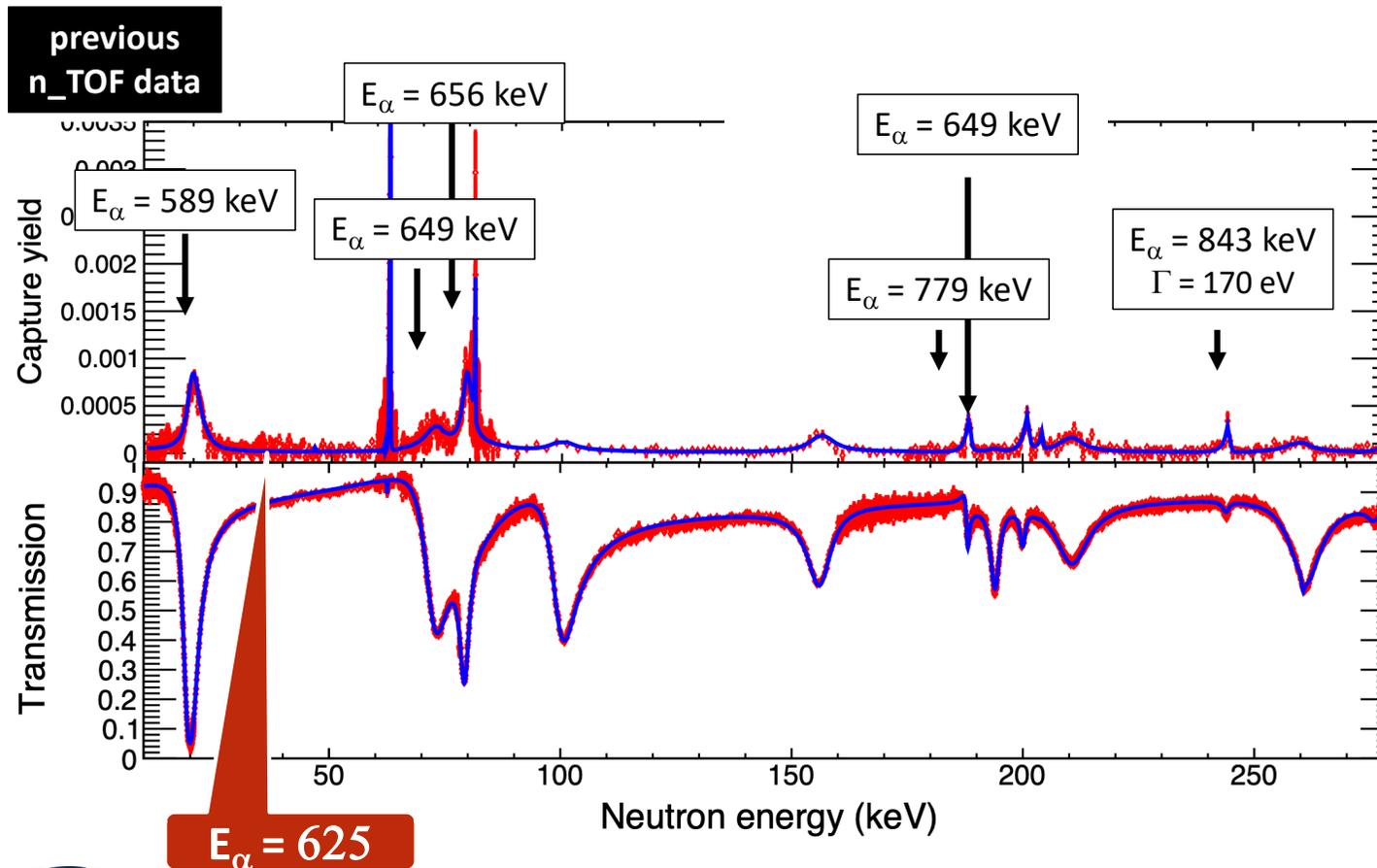
$E_\alpha = 832 \pm 2 \text{ keV}$
 $E_{\text{CM}} = 706 \text{ keV}$



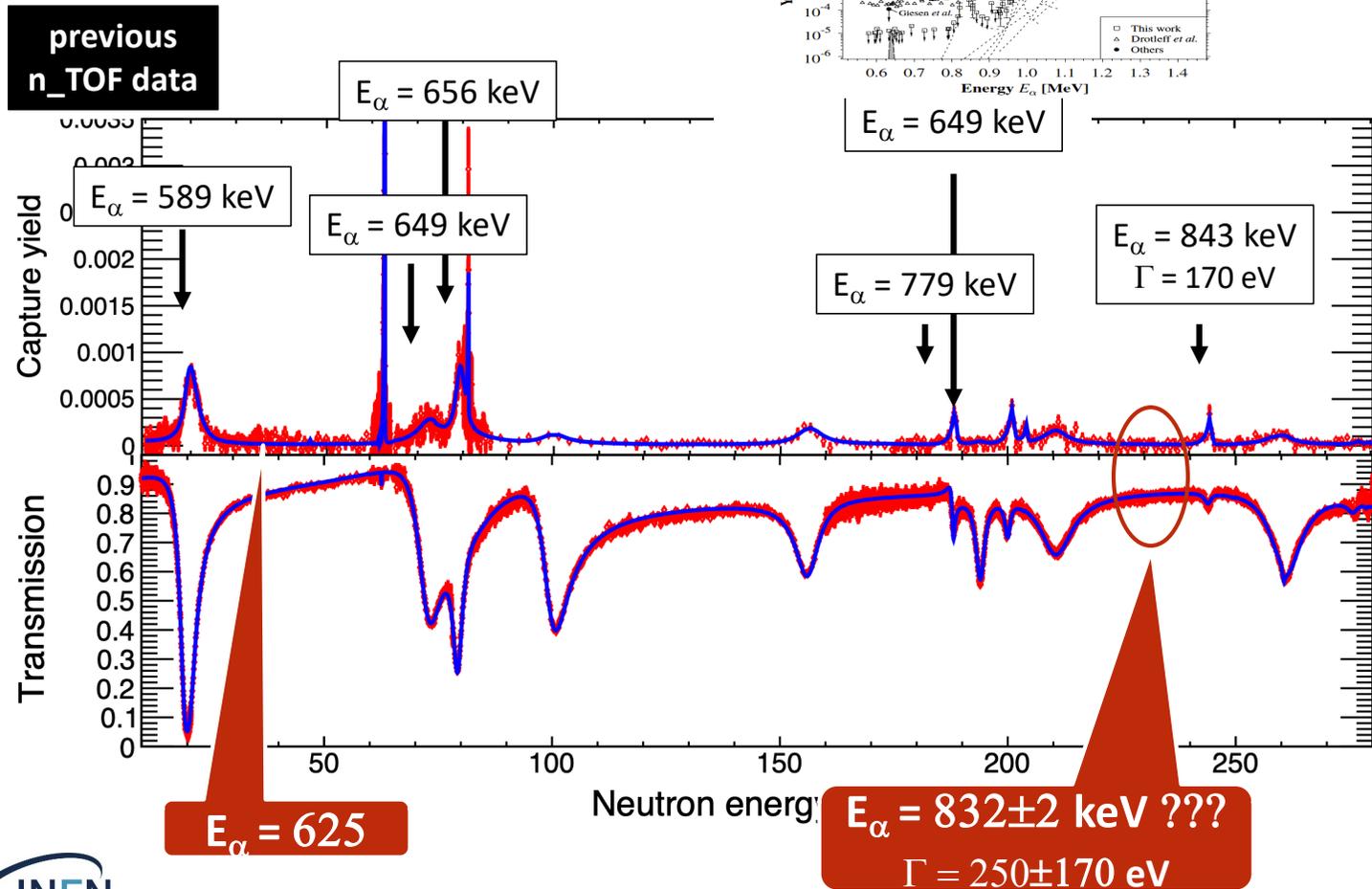
$^{25}\text{Mg}(n,\gamma)$ for neutron source reaction in stars



$^{25}\text{Mg}(n,\gamma)$ for neutron source reaction in stars



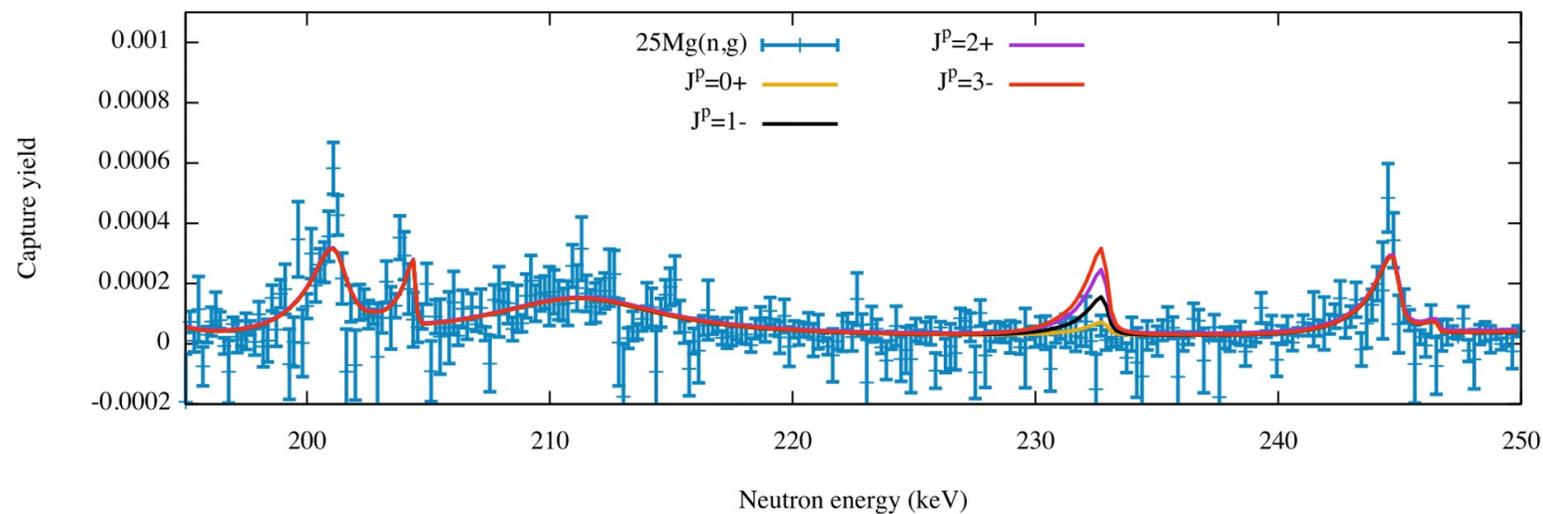
$^{25}\text{Mg}(n,\gamma)$ for neutron source reaction in stars



Proposal: $^{25}\text{Mg}(n,\gamma)^{26}\text{Mg}$ @ n_TOF

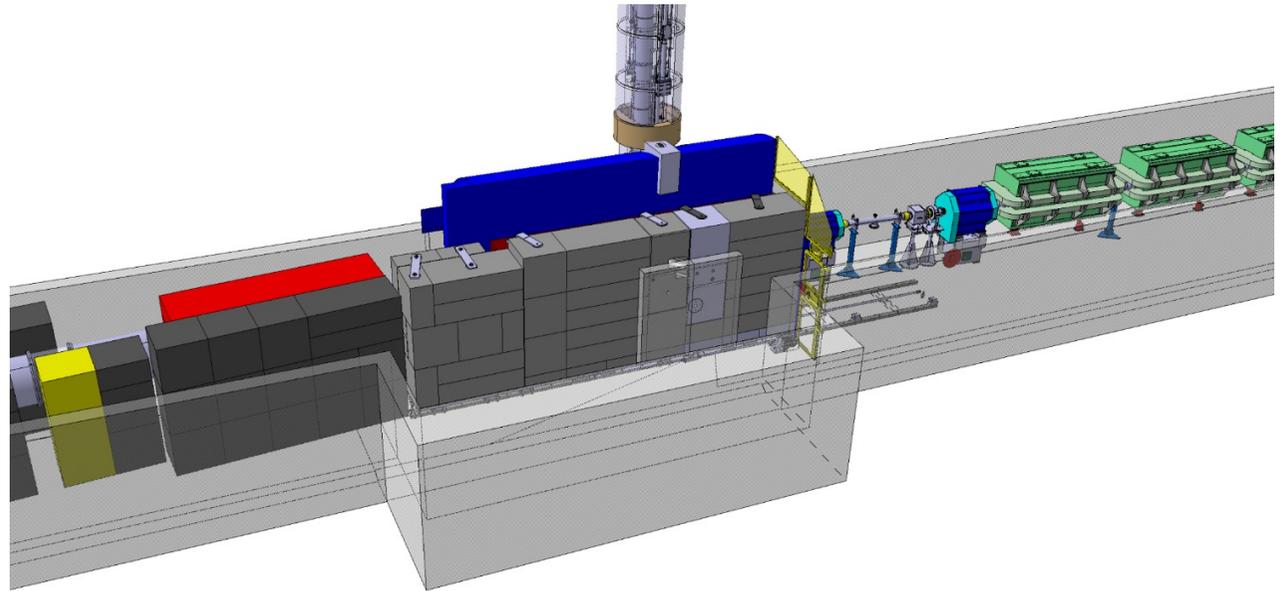
Our proposal is to **repeat the measurement in EAR1** with a factor 4 higher statistics and with some improvements:

- Combined use of C_6D_6 and LaBr detectors
- use of a thicker enriched ^{25}Mg sample
- combine with a capture measurement in EAR2



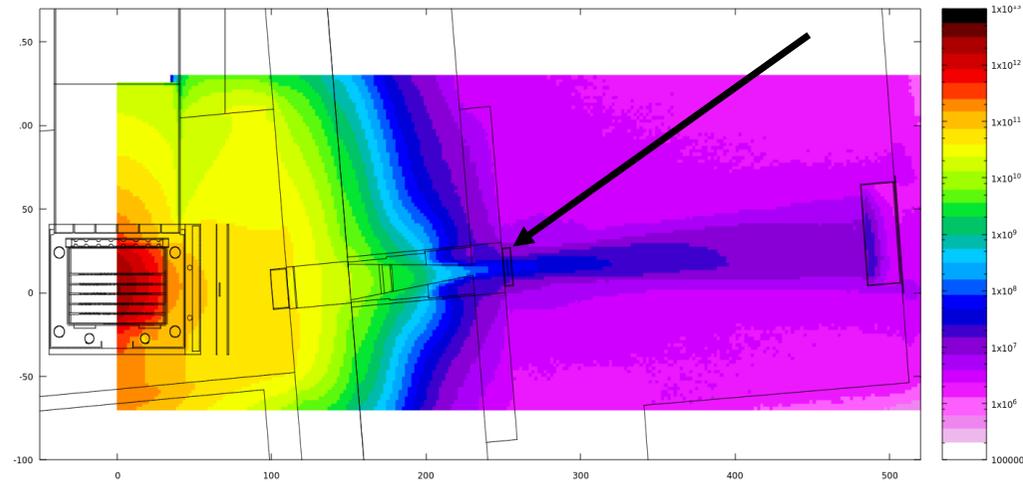
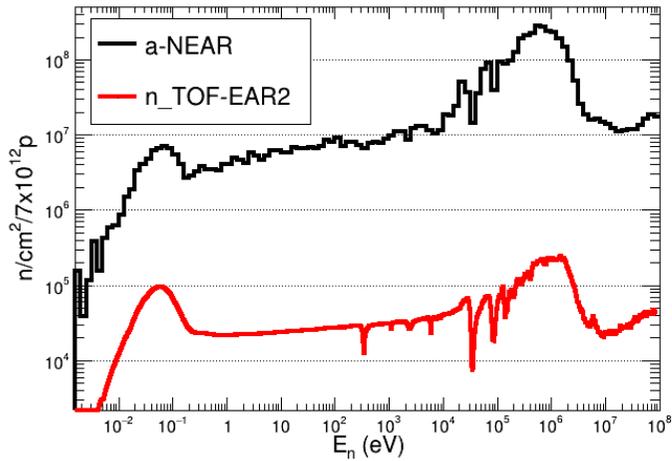
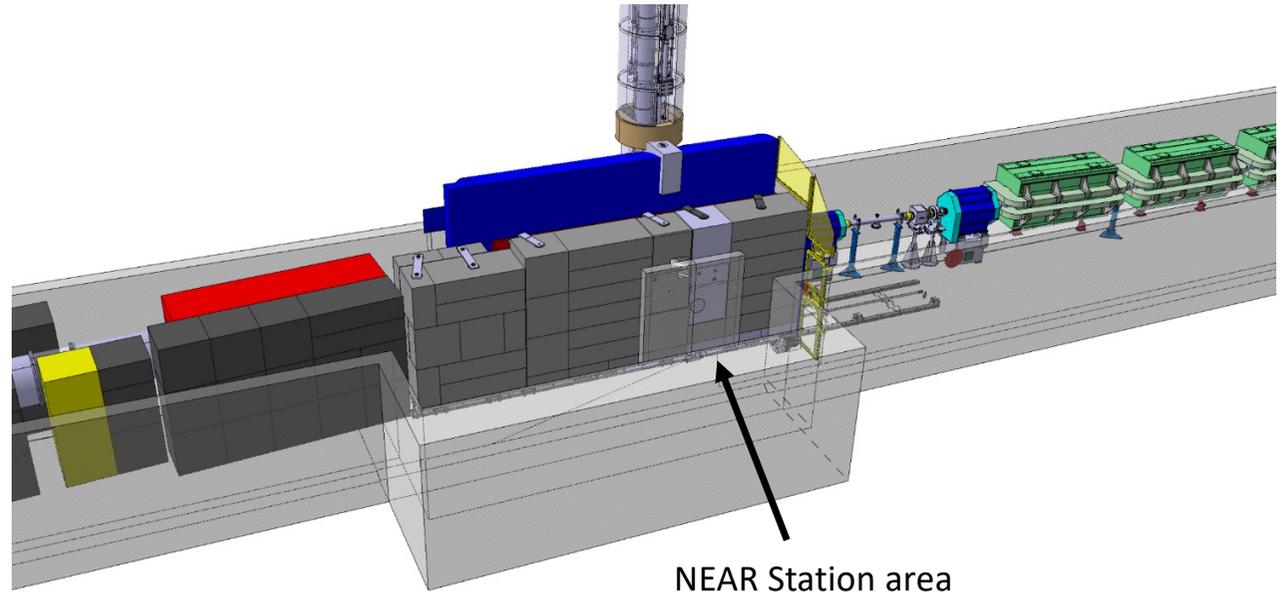
Activation at NEAR

During the design studies of the new shielding around the neutron-target station the opportunity for a new near-target experimental area appeared (NEAR station)



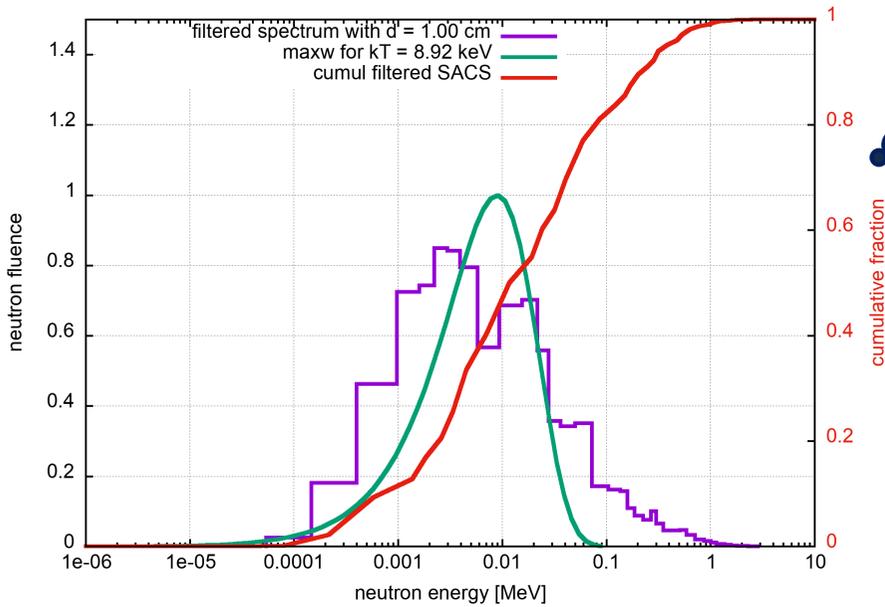
Activation at NEAR

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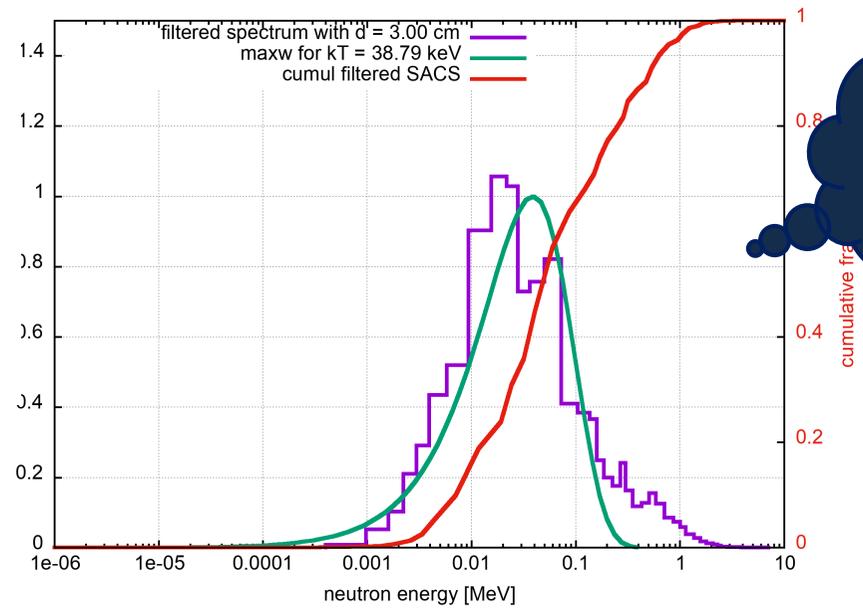
Activation at NEAR

SACS for Au197 (ENDF/B-VIII.0 data)



Filtered neutron spectrum
kT = 9 keV

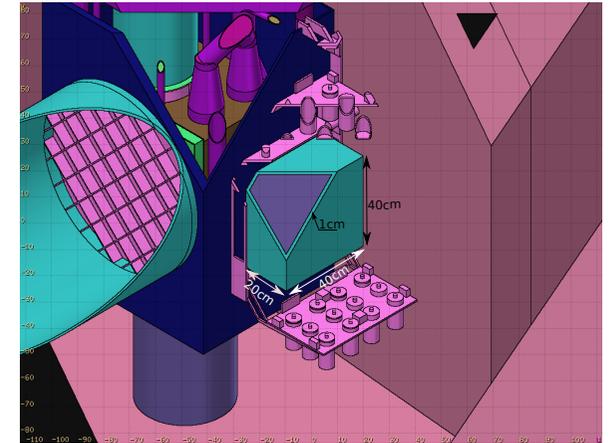
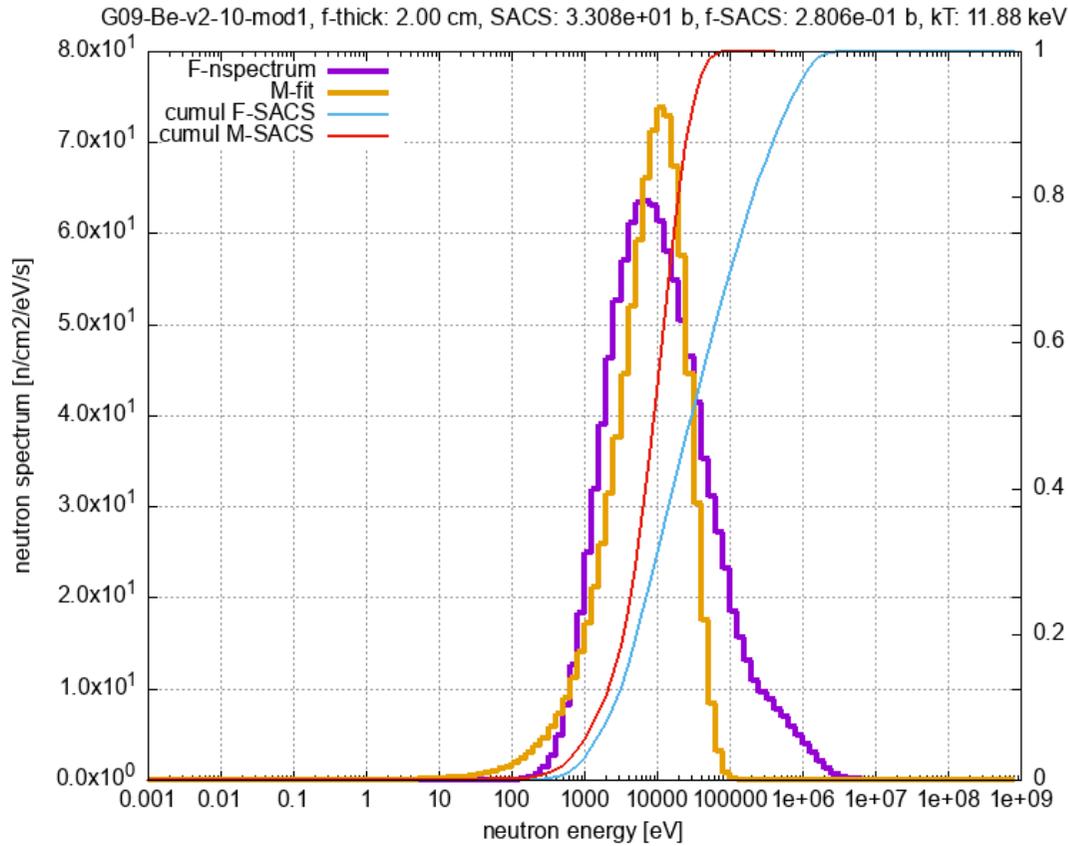
SACS for Au197 (ENDF/B-VIII.0 data)



Filtered neutron spectrum
kT = 39 keV

Activation at NEAR

Need for a moderator

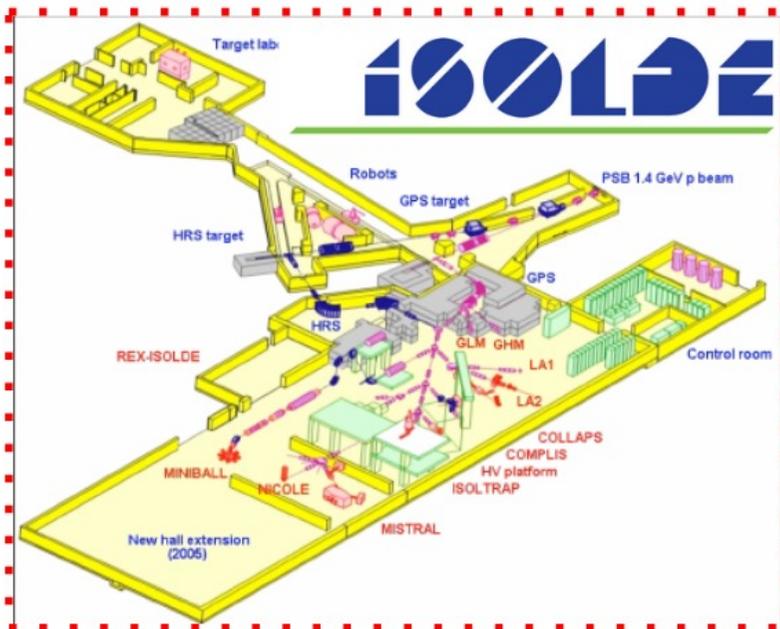


Filtered moderated neutron spectrum
kT = 12 keV

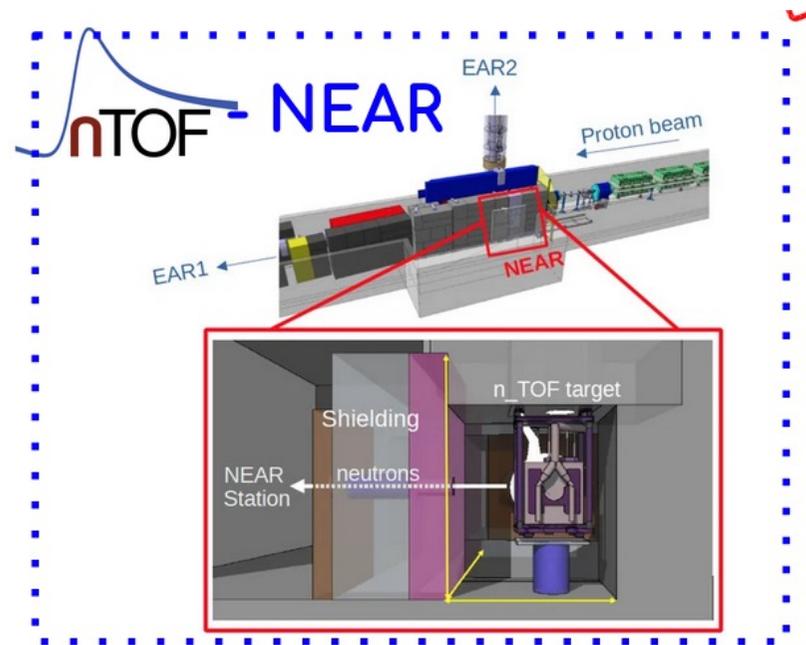
Activation at NEAR

Sample production at ISOLDE
Rare and/or short-lived isotopes

(n, γ) activation at n_TOF

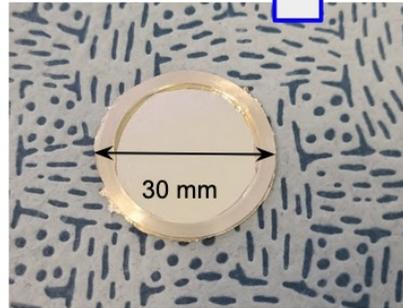


+



Examples: ^{59}Fe , ^{94}Nb , ^{125}Sb , ^{134}Cs , ^{135}Cs , ^{144}Ce , ^{148}Pm , ^{154}Eu , ^{155}Eu , ^{160}Tb , ^{170}Tm , ^{171}Tm , and ^{181}Hf (s-process), ^{137}Cs , ^{66}Ni , ^{72}Zn (i-process)

Activation at NEAR



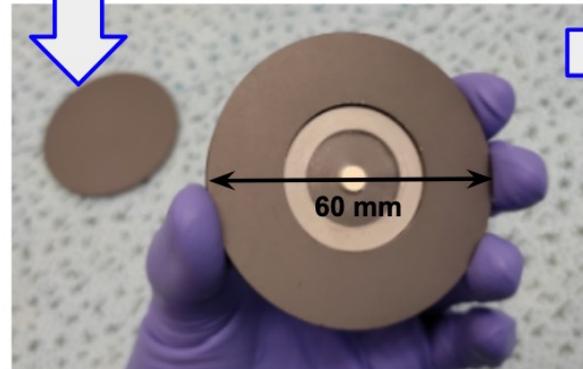
1) ^{135}Cs sample:
30mm diam Al frame

**Expected dates:
Beginning 2025**

2) Irradiation at NEAR: Week(s)-long irradiations



3) Access to NEAR (6h) + manual
transport to decay station (~h)



4) Decay measurement at GEAR: HPGe decay station



$T_{1/2} \sim \text{seconds}$

Activation at NEAR (cyclic ?)



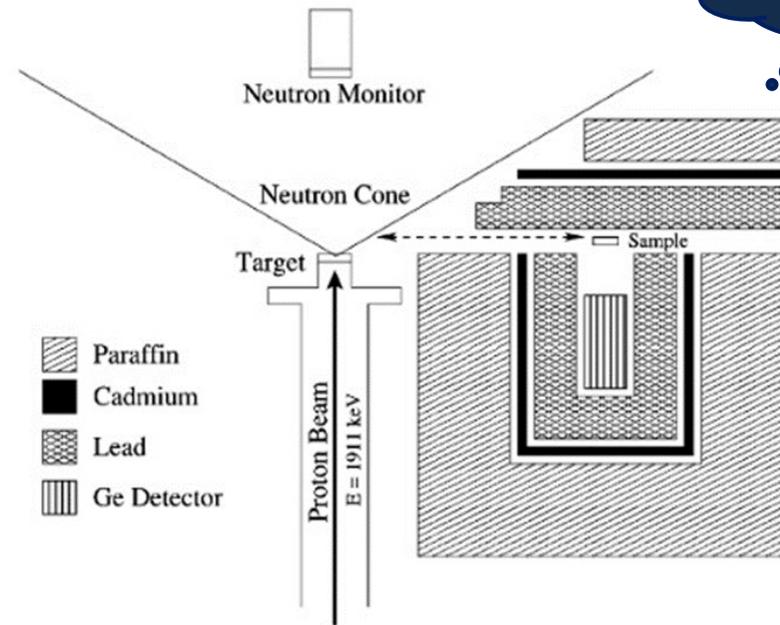
Nuclear Instruments and Methods in Physics
Research Section A: Accelerators, Spectrometers,
Detectors and Associated Equipment



Volume 337, Issues 2-3, 1 January 1994, Pages 492-503

The fast cyclic neutron activation technique at the Karlsruhe 3.75 MV Van de Graaff accelerator and the measurement of the $^{107,109}\text{Ag}(n, \gamma)^{108,110}\text{Ag}$ cross sections at $kT = 25 \text{ keV}$

Hermann Beer ^a, G. Rupp ^a, G. Walter ^b, F. Voss ^a, F. Käppeler ^a



(n,γ) activation with low-lived products

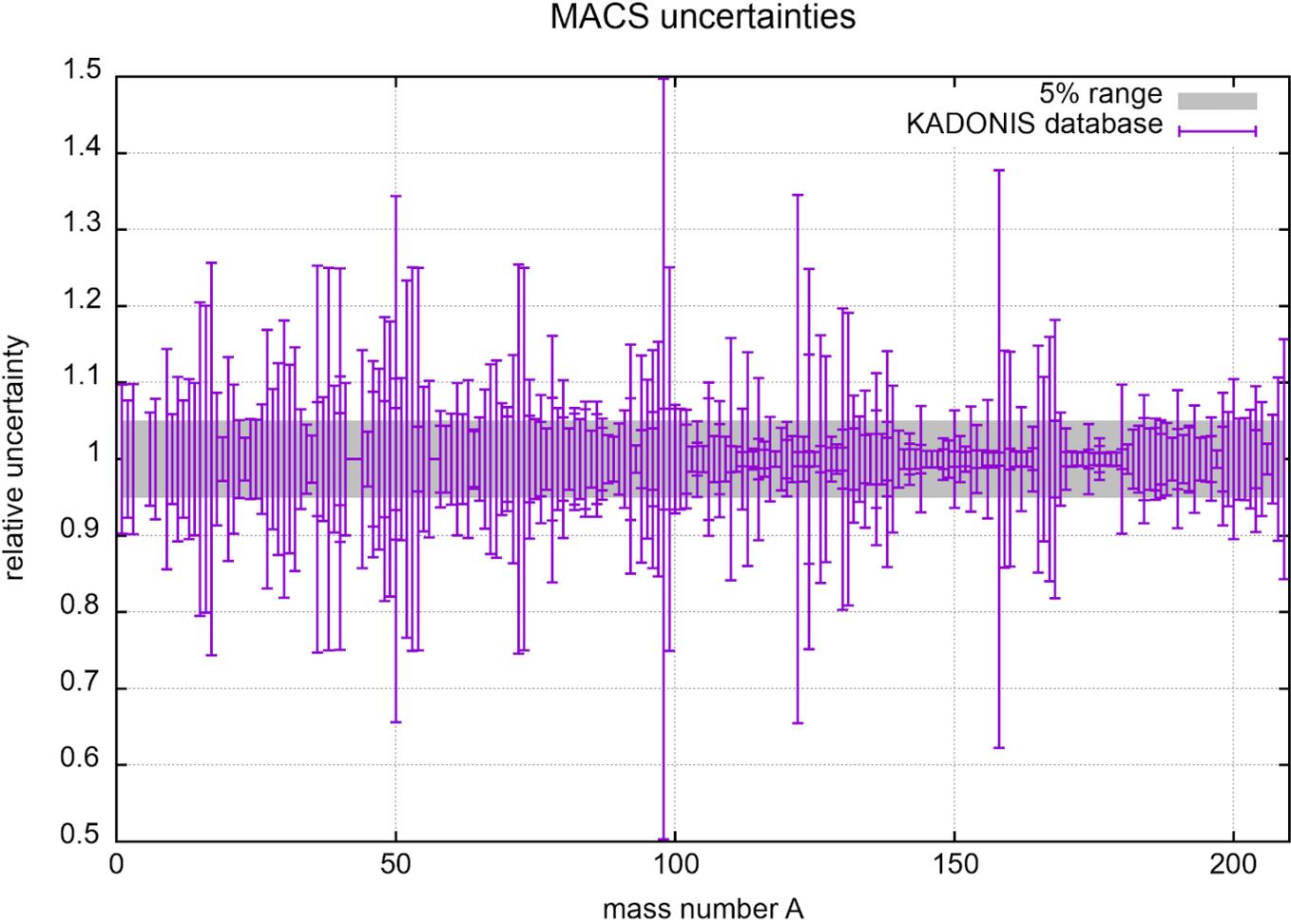
Requisites

Beam period: Rep. rate of n_TOF (max 0.8 Hz) is well suited for short lived (seconds)

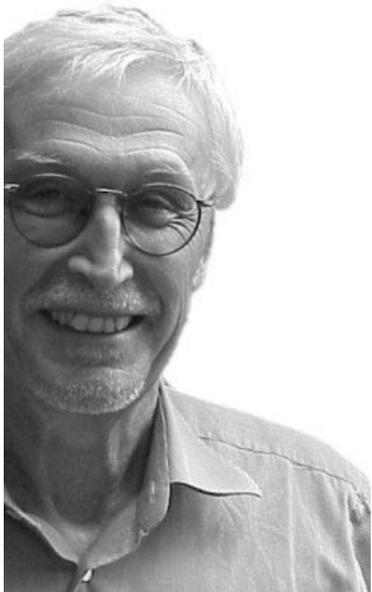
Operate a high resolution γ -ray detector in **the harsh radiation environment in the NEAR bunker**

Conclusions

Conclusions

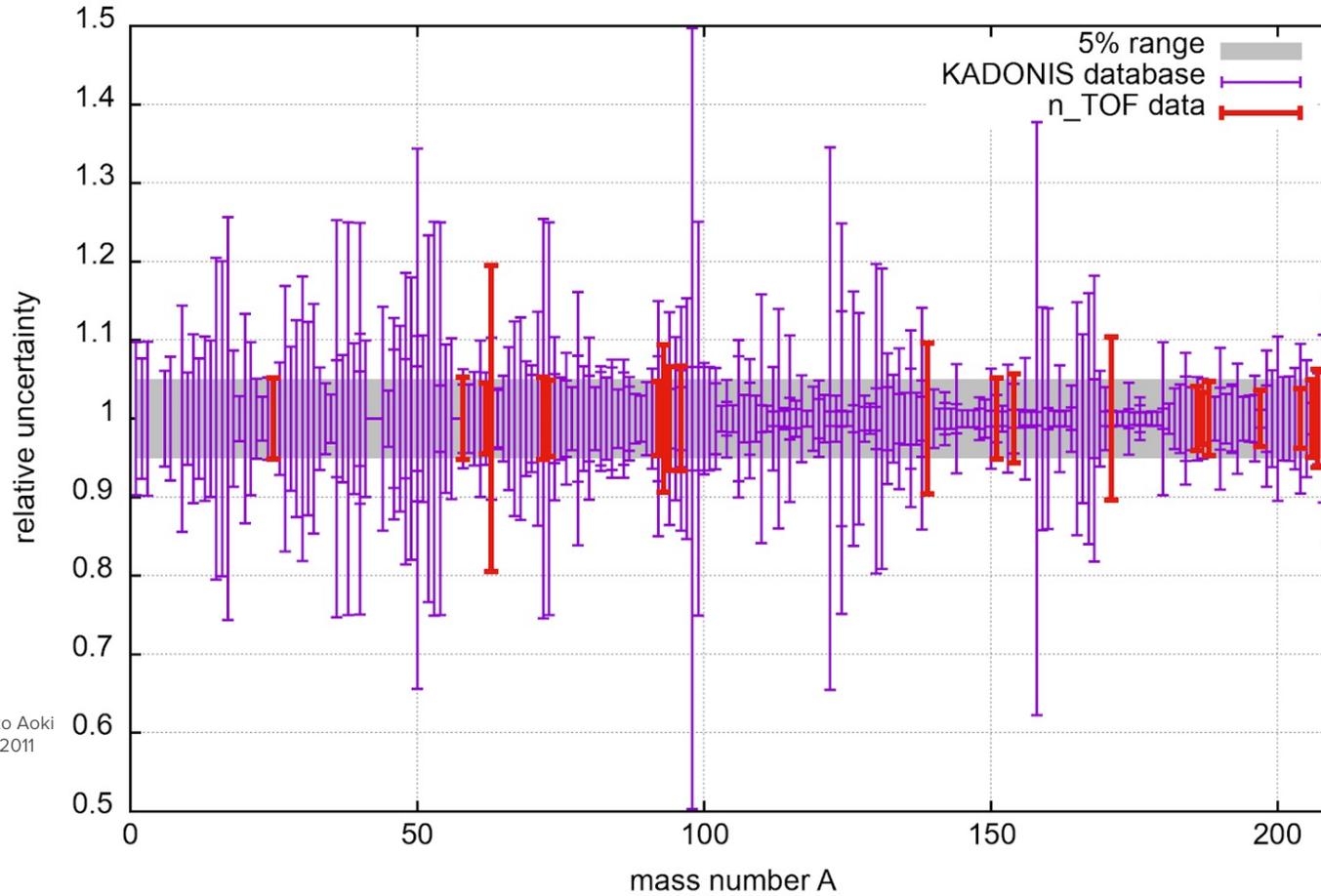


Conclusions



F. Käppeler, R. Gallino, S. Bisterzo, and Wako Aoki
Rev. Mod. Phys. **83**, 157 – Published 1 April 2011

MACS uncertainties



Conclusions

- Efforts in neutron-beam facilities (EAR2, NEAR, etc.), detection systems, and sample production techniques (ILL, PSI, ISOLDE) have enabled significant progress in the measurement of challenging neutron-capture cross-sections.
- However, there are still many neutron-capture cross-sections whose accuracy needs to be improved to the 5% level or lower. Achieving this should be possible for stable isotopes with state-of-the-art instrumentation and facilities.
- For unstable nuclei, further developments in detection systems and facilities will be required.

Thank you for your attention

Thanks to the organizers



Many thanks to the
n_TOF Collaboration





ALMA MATER STUDIORUM
UNIVERSITÀ DI BOLOGNA

Cristian Massimi

Department of Physics and Astronomy

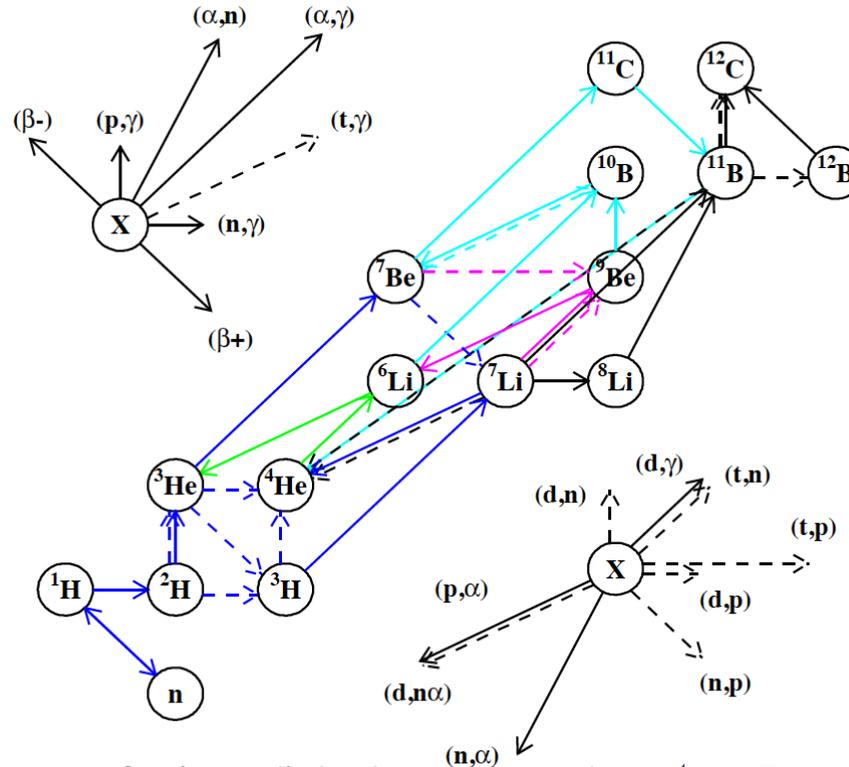
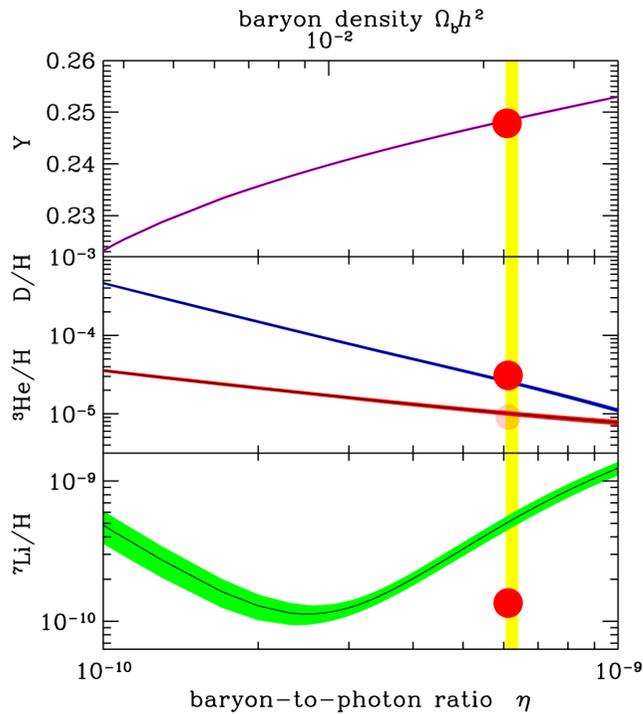
cristian.massimi@unibo.it

www.unibo.it

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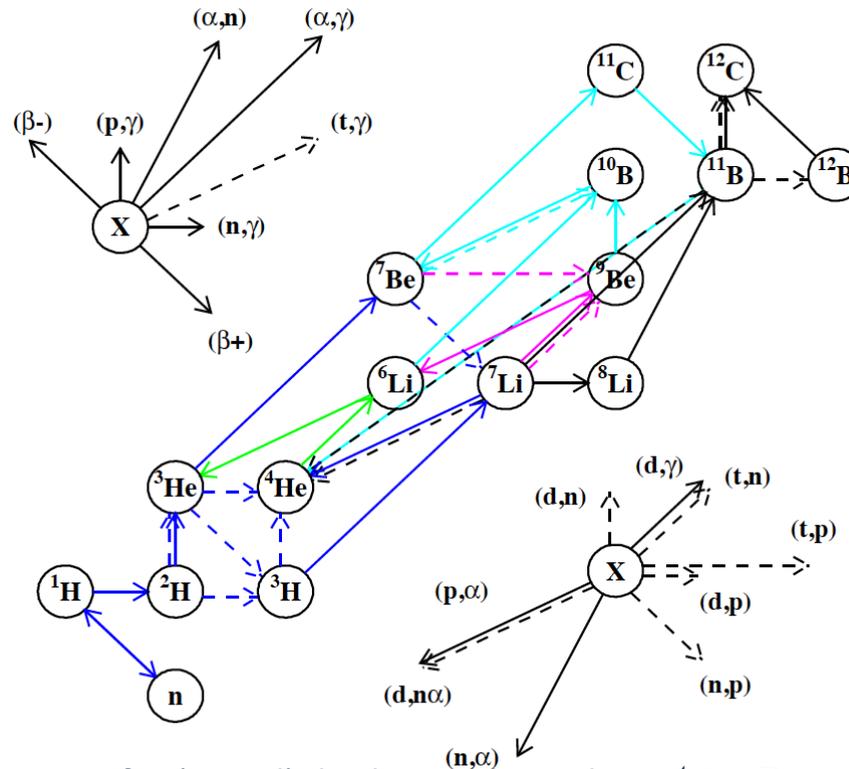
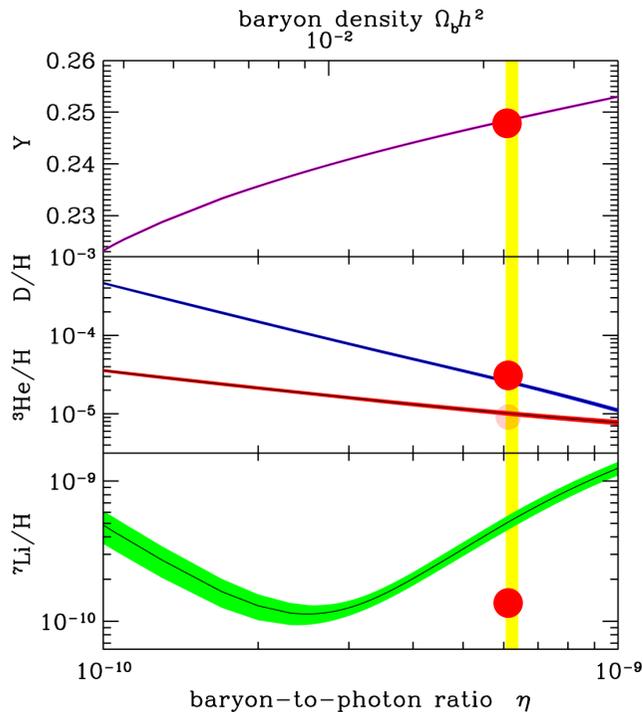
Cosmological lithium problem and ${}^7\text{Be}$



BBN successfully predicts the abundances of primordial elements such as ${}^4\text{He}$, D and ${}^3\text{He}$. Large **discrepancy** for ${}^7\text{Li}$, which is produced from electron capture decay of ${}^7\text{Be}$



Cosmological lithium problem and ⁷Be



~ 95% of ⁷Li is produced by the decay of ⁷Be ($T_{1/2}=53.2$ d)

↓

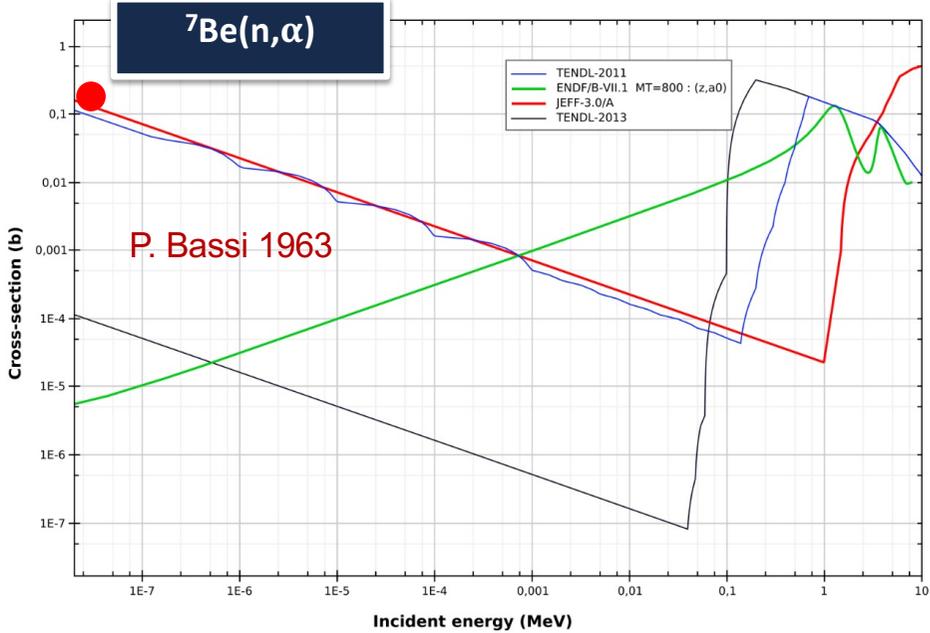
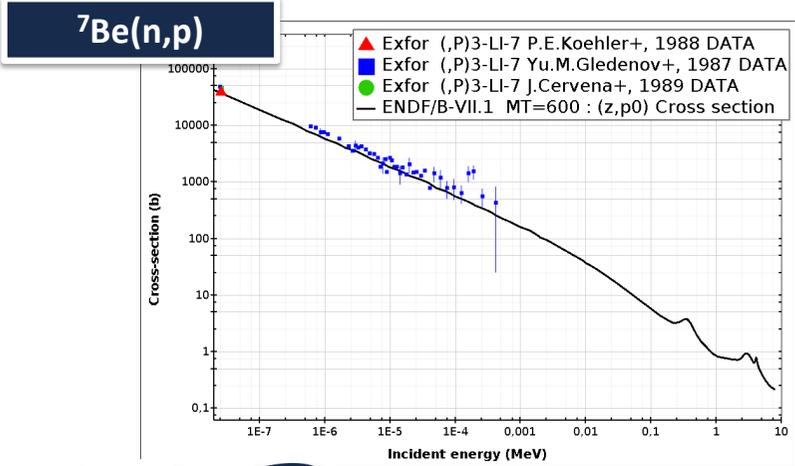
⁷Be is the key

BBN successfully predicts the abundances of primordial elements such as ⁴He, D and ³He. Large **discrepancy** for ⁷Li, which is produced from electron capture decay of ⁷Be



Cosmological lithium problem and ^7Be

- ^7Be is **destroyed** by:
- $(n,p) \approx 97\%$
 - $(n,\alpha) \approx 2.5\%$
- With a **10 times higher destruction rate** of ^7Be the cosmological lithium problem could be solved (**nuclear solution**)



Date in the literature: scarce and uncertain

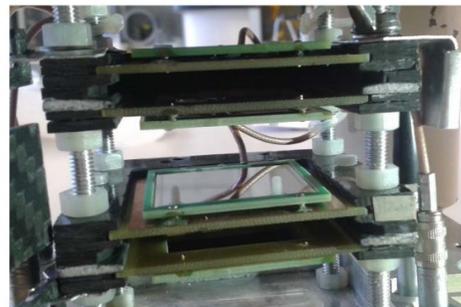
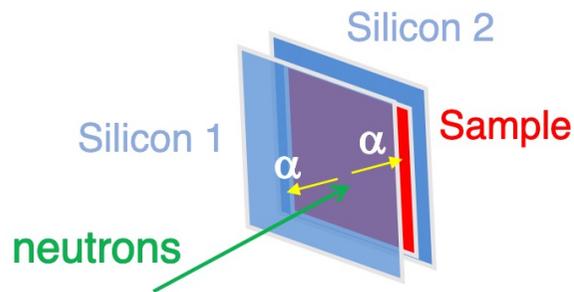
Cosmological lithium problem and ^7Be

The (n,α) reaction produces **two α -particles** emitted back-to-back with **several MeV energy** ($Q\text{-value}=19\text{ MeV}$)

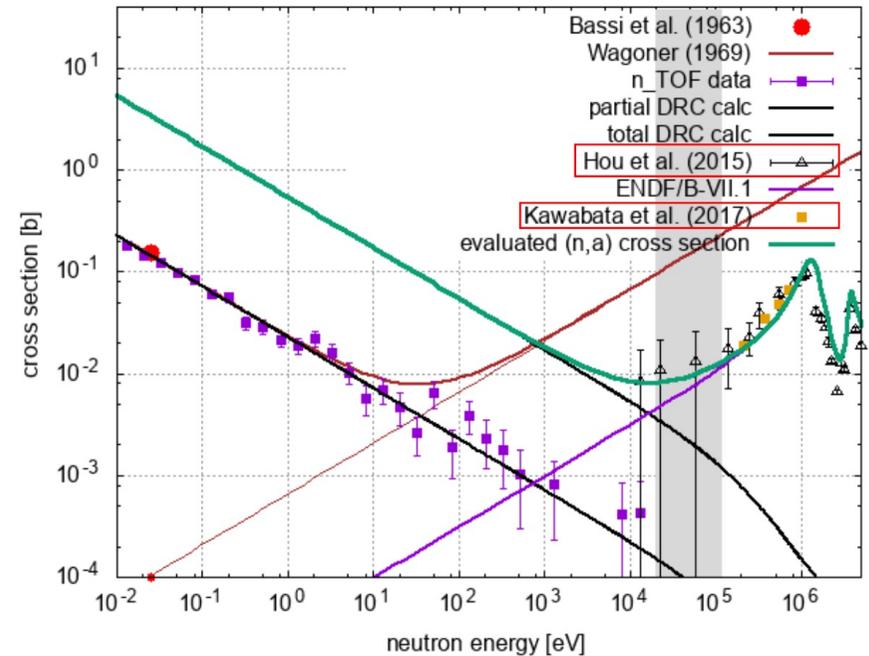
2 Sandwiches of **silicon detector** (140 mm, $3\times 3\text{cm}^2$) with ^7Be sample in between **directly inserted in the neutron beam**

Coincidence technique: strong background rejection

$^7\text{Be}(n,\alpha)$



M. Barbagallo *et al.* (The n_TOF Collaboration), *Phys. Rev. Lett.* **117** 152701 (2016)



- 4 **Electrodeposition on a 5- μm -thick Al foil**
- 3
- 2
- 1



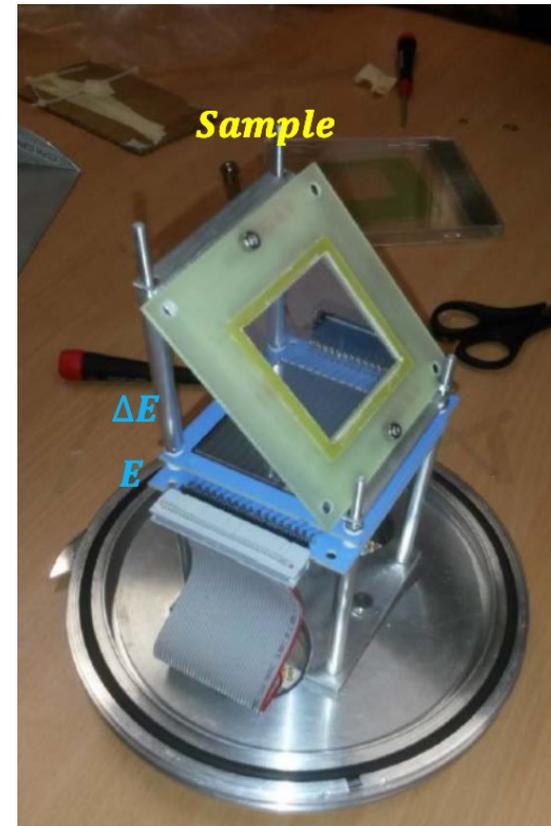
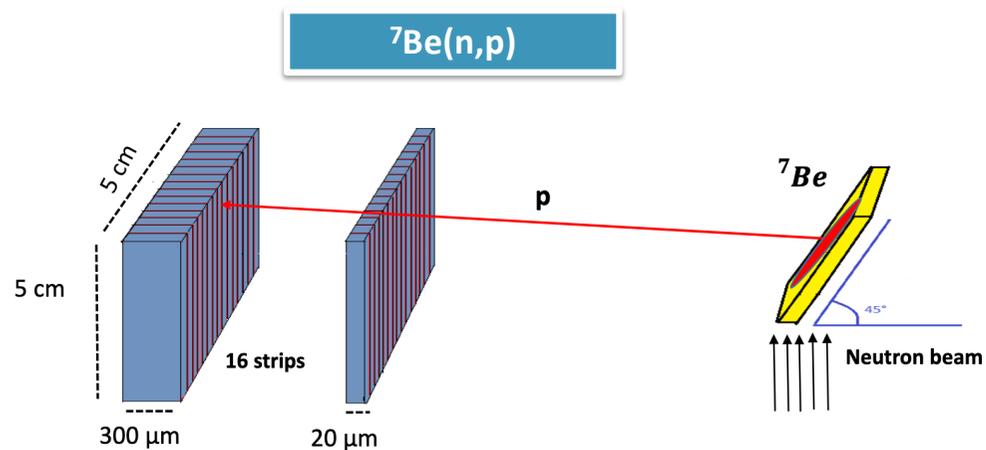
Cosmological lithium problem and ${}^7\text{Be}$

The (n,p) reaction cross section is very high

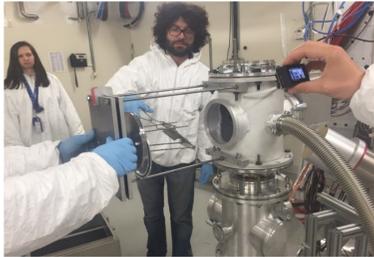
Q-value = 1.6 MeV

Silicon counter telescope $\Delta E-E$

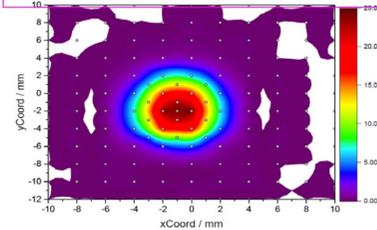
A few ng of 100% **enriched** sample is needed.



Cosmological lithium problem and ^7Be



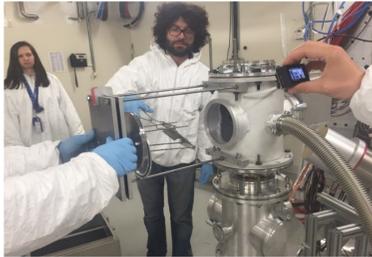
Sample characterization @PSI
(Gaussian profile 0.5 cm FWHM)



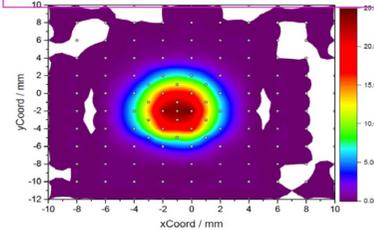
- 200 GBq of ^7Be extracted from the cooling water of the SINQ spallation source at PSI
- Transported to ISOLDE at CERN and installed in the ion source to produce 30 keV ion beam.
- ^7Be beam separated by means of a magnetic dipole, and implanted on a 20 m thick aluminum backing.
- Sample of **1 GBq** ^7Be (~80 ng) transported to EAR2@n_TOF and placed in the neutron beam.

$^7\text{Be}(n,p)$

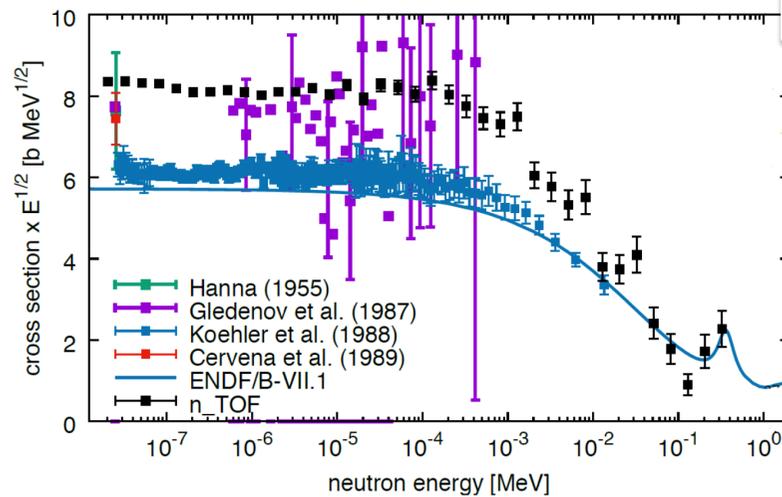
Cosmological lithium problem and ^7Be



Sample characterization @PSI
(Gaussian profile 0.5 cm FWHM)



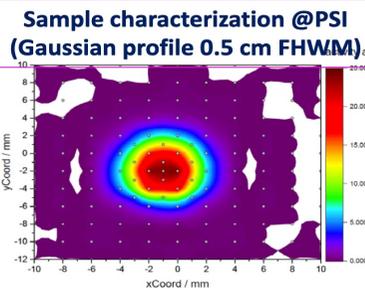
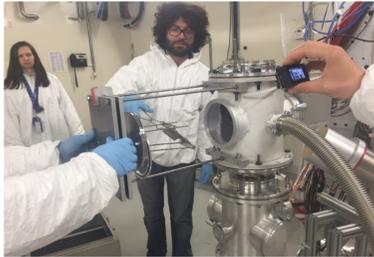
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$^7\text{Be}(n,p)$

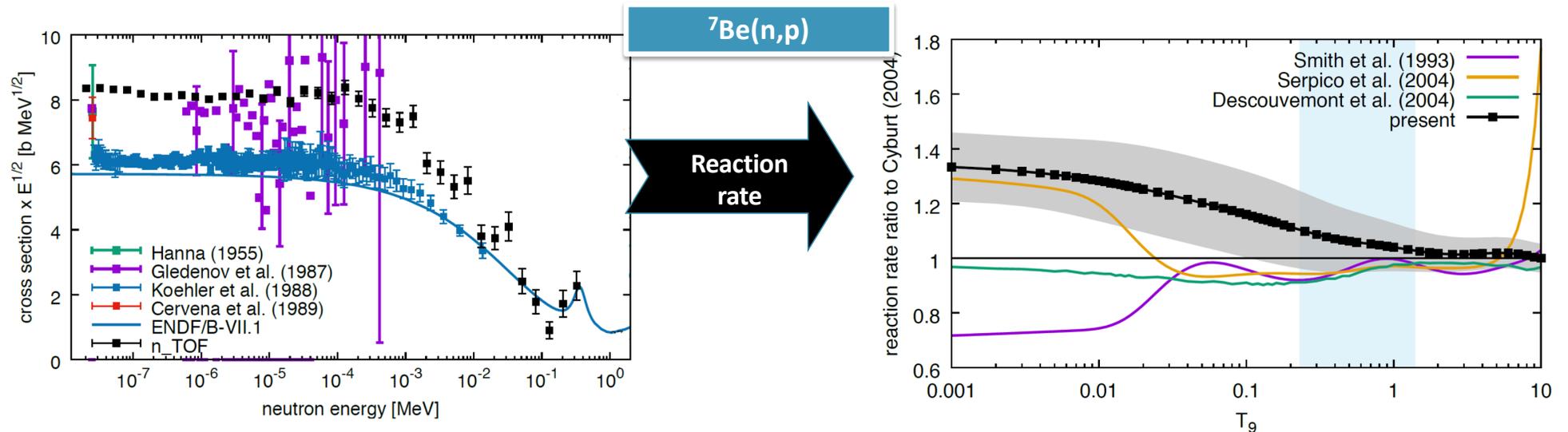


Cosmological lithium problem and ^7Be



$^7\text{Be}(n,\alpha)$ and $^7\text{Be}(n,p)$ results exclude these channels as a solution for the problem

- 200 GBq of ^7Be extracted from the cooling water of the SINQ spallation source at PSI
- Transported to ISOLDE at CERN and installed in the ion source to produce 30 keV ion beam.
- ^7Be beam separated by means of a magnetic dipole, and implanted on a 20 m thick aluminum backing.
- Sample of **1 GBq** ^7Be (~80 ng) transported to EAR2@n_TOF and placed in the neutron beam.



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Extra slides

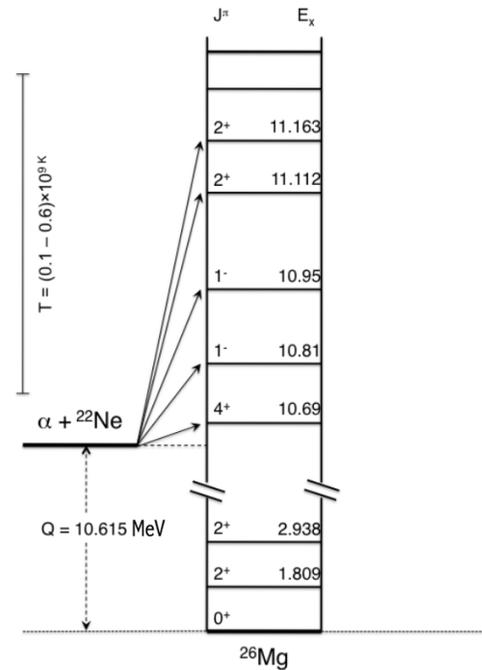
Constraints for the $^{22}\text{Ne}(\alpha, n)^{25}\text{Mg}$ reaction

Element	Spin/ parity
^{22}Ne	0^+
^4He	0^+

$$\vec{J} = \underbrace{\vec{I} + \vec{i}} + \vec{\ell}$$

$$\vec{J} = \mathbf{0} + \vec{\ell}$$

Only **natural-parity** (0^+ , 1^- , 2^+ , 3^- , 4^+ , ...) **states in ^{26}Mg** can participate in the $^{22}\text{Ne}(\alpha, n)^{25}\text{Mg}$ reaction



Extra slides

Constraints for the $^{22}\text{Ne}(\alpha, n)^{25}\text{Mg}$ reaction

Element	Spin/parity
^{25}Mg	$5/2^+$
n	$1/2^+$

$$\vec{J} = \vec{I} + \vec{i} + \vec{\ell}$$

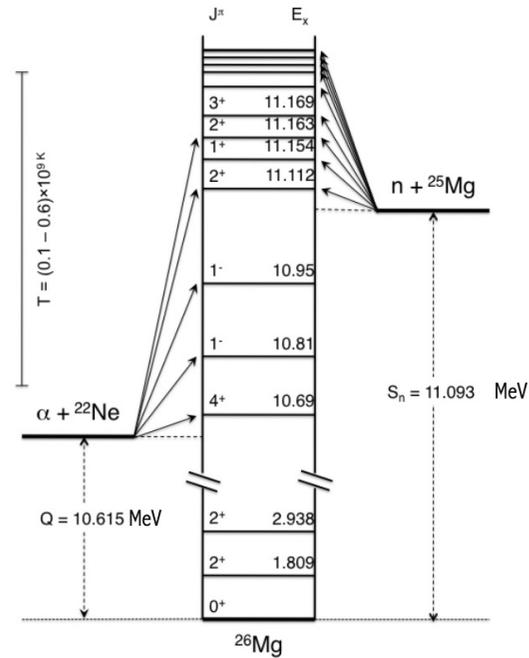
$$\vec{J} = \underline{2} + \vec{\ell} \quad \vec{J} = \underline{3} + \vec{\ell}$$

s-wave $\rightarrow J^\pi = \underline{2}^+, 3^+$

p-wave $\rightarrow J^\pi = \underline{1}^-, 2^-, \underline{3}^-, 4^-$

d-wave $\rightarrow J^\pi = \underline{0}^+, 1^+, \underline{2}^+, 3^+, \underline{4}^+, 5^+$

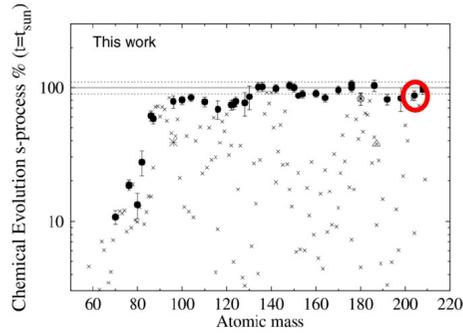
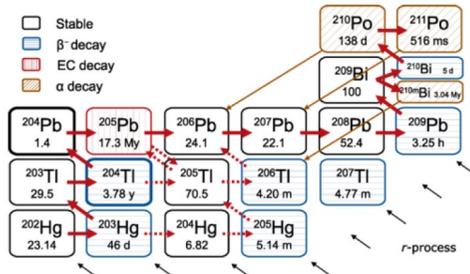
States in ^{26}Mg populated by $^{25}\text{Mg}(n, \gamma)$ reaction



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C. Domingo Pardo
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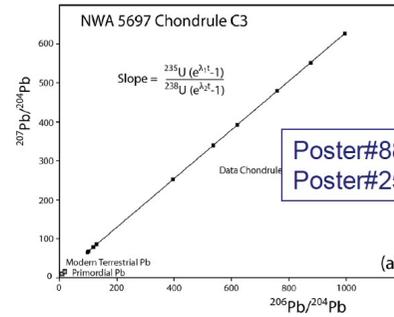
Origin of the heaviest s-only isotope ^{204}Pb : Neutron capture on ^{204}Tl (3.78y)



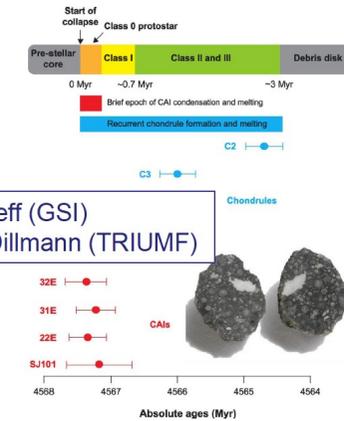
A Casanovas-Hoste *et al.* (The n_{TOF} Collaboration)
 Physical Review Letters **133**, 052702 (2024)
 DOI: [10.1103/PhysRevLett.133.052702](https://doi.org/10.1103/PhysRevLett.133.052702)

- Pure s-process origin
 - No r-process contribution (shielded)
 - No radiogenic contribution from U/Th
- Benchmark for AGB- and GCE-models (Travaglio 2004, Bisterzo 2014)
- Primordial abundance preserved; the reference for Pb-Pb clock (Connelly 2012)
- Its nucleosynthesis is impacted (dominated) by the branching at ^{204}Tl
- Thus far only theoretical estimates existed for the latter (x2 variations)!
- γ -process contribution (Pignatari 2016) and/or fractionation in ESS (González 2014)?

$$\left(\frac{^{204}\text{Pb}}{^{208}\text{Pb}}\right)_{\text{today}} = \left(\frac{^{204}\text{Pb}}{^{208}\text{Pb}}\right)_{\text{at}} + \left(\frac{^{235}\text{U}, ^{232}\text{Th}}{^{204}\text{Pb}}\right)_{\text{today}} (e^{\lambda t} - 1)$$



J.N. Connelly *et al.* Science **338** 2012
 DOI: [10.1126/science.1226919](https://doi.org/10.1126/science.1226919)

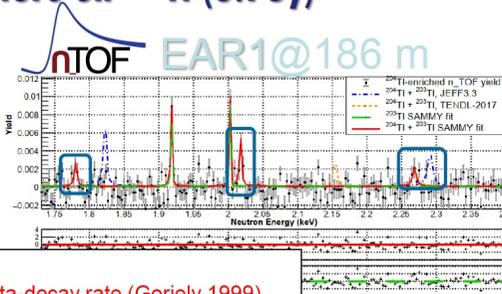
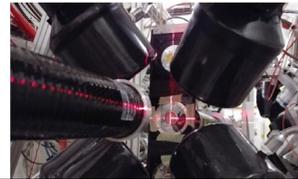
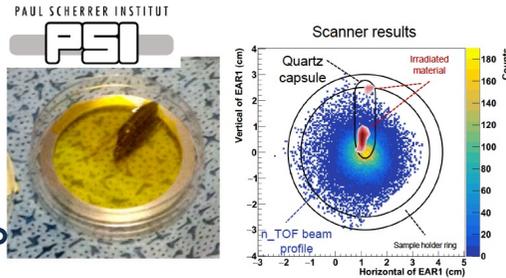


Poster#88 T. Neff (GSI)
 Poster#254 I. Dillmann (TRIUMF)

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NPA - 2024

Origin of the heaviest s-only isotope ^{204}Pb : Neutron capture on ^{204}Tl (3.78y)



→ Remaining open questions:
- Thermal dependency of the beta-decay rate (Goriely 1999)
- Strength of the $^{22}\text{Ne}(\alpha, n)$ source

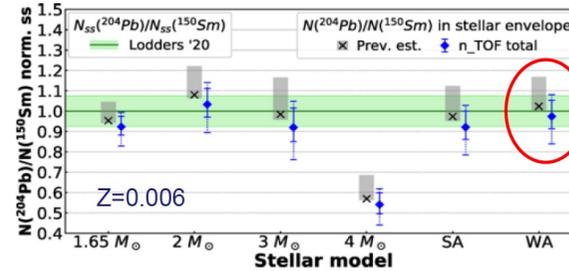
Pellet containing 225 mg of ^{203}Tl , 99.5% isotopic purity, produced at PSI by machine pressing and inserted into quartz container

Irradiated at ILL reactor with thermal neutrons for 55 days: 9 mg of ^{204}Tl produced

180 GBq of β activity plus radioactive impurities and bremsstrahlung



A Casanovas-Hoste *et al.* (The n_TOF Collaboration)
Physical Review Letters **133**, 052702 (2024)
DOI: [10.1103/PhysRevLett.133.052702](https://doi.org/10.1103/PhysRevLett.133.052702)

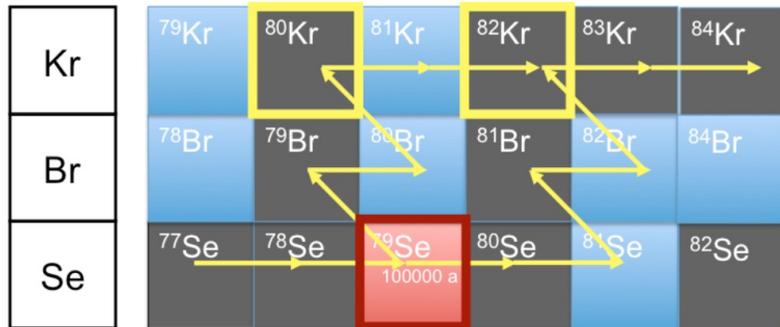


The uncertainty arising from the $^{204}\text{Tl}(n, \gamma)$ cross section on the s-process abundance of ^{204}Pb has been reduced from ~30% down to +8%/-6%, and the s-process calculations are in agreement with K. Lodders in 2021.

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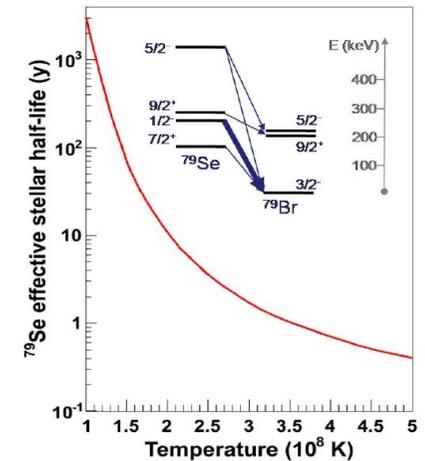


Some ongoing activities: $^{79}\text{Se}(n,\gamma)$



s-process
branching at ^{79}Se

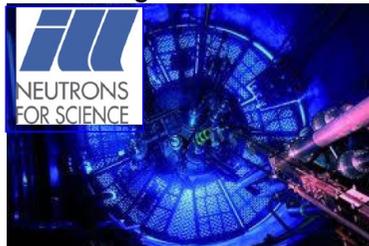
- ^{80}Kr and ^{82}Kr are s-only isotopes
- Kr isotopic ratios measured in SiC grains
- branching ratio provides information on the thermodynamical conditions of the star



Timeline:

- 2018 Sample produced @ ILL (France)
- 2019 Sample characterize at PSI (Switzerland)
- 2019 test of an innovative (n, γ) detector i-TED¹
- 2021 Time-of-flight measurement @ EAR1

ILL: ~3 mg of ^{79}Se via ^{78}Se n-activation



6MBq γ -ray emitters
3 mg of ^{79}Se
1.6 MBq of ^{60}Co
5 MBq of ^{75}Se

PSI: $^{208}\text{Pb}^{78}\text{Se}$ alloy



2.8 g of ^{208}Pb
1.0 g of ^{78}Se



Preparation of PbSe targets for ^{79}Se neutron capture cross section studies

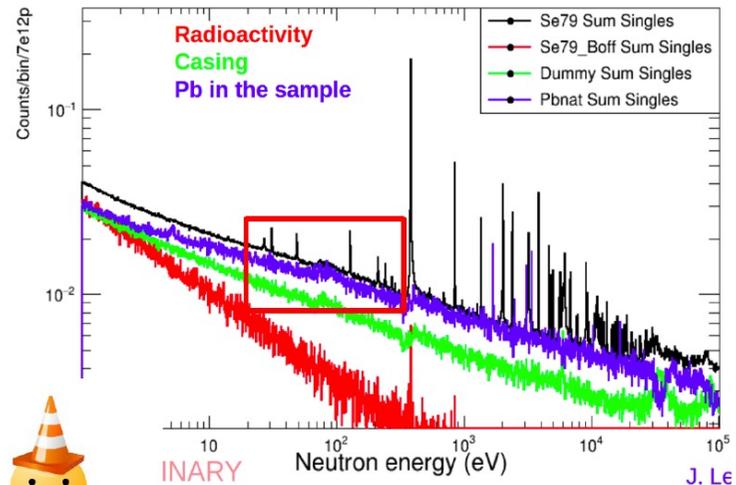
Nadine M. Chieira ^{a,*}, Emilio Andrea Maugeri ^a, Ivan Danilov ^b, Javier Balibrea-Correa ^b, Cesar Domingo-Pardo ^b, Ulli Köster ^c, Jorge Lerendegui-Marco ^b, Mario Vesicht ^{a,d}, Ivan Zivadinovic ^{a,e}, Dorothea Schumann ^a, the n_TOF collaboration

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^dÉcole Polytechnique, Université de Lausanne, Switzerland
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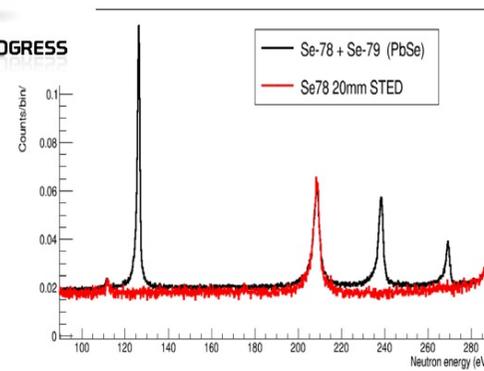


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Some ongoing activities: $^{79}\text{Se}(n,\gamma)$



WORK IN PROGRESS



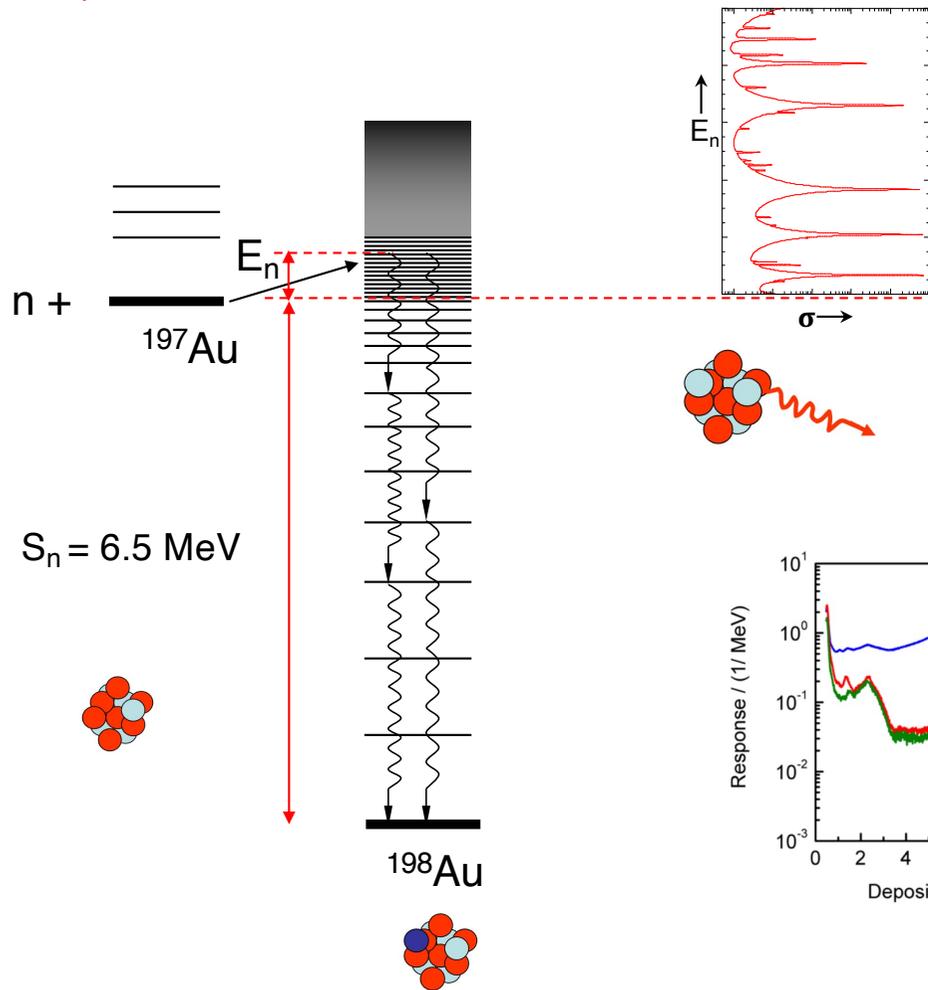
J. Lerendegui-Marco et al. (analysis in progress)



backup slides



(n, γ)

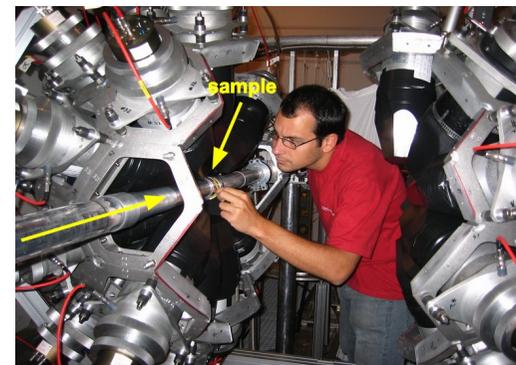
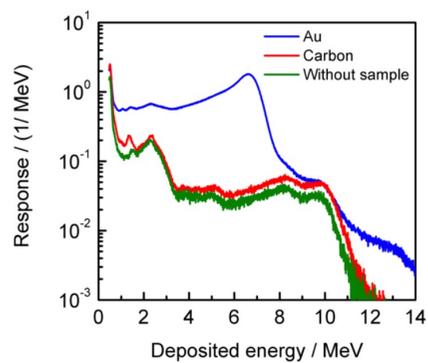


The C_6D_6 Total Energy Detectors (TED)

C_6D_6 scintillators
at 135°



The BaF_2 Total γ -ray Absorption Detector



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Spallation-based neutron target for direct studies of neutron-induced reactions in inverse kinematics

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We discuss the possibility to build a neutron target for nuclear reaction studies in inverse kinematics utilizing a storage ring and radioactive ion beams. The proposed neutron target is a specially designed spallation target surrounded by a large moderator of heavy water (D_2O). We present the resulting neutron spectra and their properties as a target. We discuss possible realizations at different experimental facilities.

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