





Inclusive D^0 photoproduction in UPCs

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1. Inclusive D^0 in ultraperipheral collision (UPC)

- Charm photoproduction in LCPT
- Collinear limit

2. Comparison to CMS data

UPC



UPC



UPC



LO in $\alpha_{\rm s}$



LO in $\alpha_{\rm s}$



• Fock state expansion

$$ig|\gamma_\lambda(q^+,\mathbf{q}_\perp;Q^2)ig
angle_{\mathrm{D}}\propto\sum_{qar{q}}\Psi^{\gamma_\lambda
ightarrow qar{q}}|qar{q}
angle_0+\sum_{qar{q}g}\Psi^{\gamma_\lambda
ightarrow qar{q}g}|qar{q}g
angle_0+\dots$$

Beuf, Hänninen, Lappi, Mulian (2022)∂

• Invariant amplitude from scattering matrix element

$$ig\langle q(m{k}_{0\perp}',k_0'^+)ar{q}(m{k}_{1\perp}',k_1'^+)ig| \hat{S} - 1 \underbrace{ig| \gamma^*(q^+,m{q}_\perp;Q^2)ig
angle}_{ ext{F.T to mixed space}} \propto \mathcal{M}(m{q}_\perp
ightarrow m{k}_{0\perp}',m{k}_{1\perp}')$$

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$$\langle q(\mathbf{k}_{0\perp}', k_{0}'^{+}) \bar{q}(\mathbf{k}_{1\perp}', k_{1}'^{+}) | \hat{S} - 1 \underbrace{|\gamma^{*}(q^{+}, \mathbf{q}_{\perp}; Q^{2})}_{\text{F.T to mixed space}} \propto \mathcal{M}(\mathbf{q}_{\perp} \rightarrow \mathbf{k}_{0\perp}', \mathbf{k}_{1\perp}')$$

$$\mathcal{M}_{\gamma^* \to q\bar{q}} = \int_{\mathbf{x}_{0\perp}\mathbf{x}_{1\perp}} e^{i\mathbf{k}_{0\perp}\mathbf{x}_{0\perp}} e^{i\mathbf{k}_{1\perp}\mathbf{x}_{1\perp}} \left(\left[U_F(\mathbf{x}_{0\perp}) U_F^{\dagger}(\mathbf{x}_{1\perp}) \right]_{\beta_0\beta_1} - \delta_{\beta_0\beta_1} \right) \psi^{\gamma^* \to q_0\bar{q}_1}$$

$$\mathrm{d}\sigma^{\gamma^*+\mathcal{A}\to c\bar{c}+X} = \int \mathrm{d}[\mathrm{P.S}]\langle |\mathcal{M}|^2 \rangle$$

Dominguez, Marquet, Xiao, Yuan (2011)

• $D^0 = ar{u}c \longrightarrow$ we integrate over the antiquark momentum $m{k}_{1\perp}$

• $D^0 = ar{u} c \longrightarrow$ we integrate over the antiquark momentum $oldsymbol{k}_{1\perp}$

Color structure in $\langle |\mathcal{M}|^2 \rangle$:

 $\left\langle \mathrm{Tr} \left(U(\boldsymbol{x}_{0\perp}) U^{\dagger}(\boldsymbol{x}_{1\perp}) U(\boldsymbol{x}_{1\perp}') U^{\dagger}(\boldsymbol{x}_{0\perp}') \right) \right\rangle - \left\langle \mathrm{Tr} \left(U(\boldsymbol{x}_{0\perp}) U^{\dagger}(\boldsymbol{x}_{1\perp}) \right) \right\rangle - \left\langle \mathrm{Tr} \left(U(\boldsymbol{x}_{1\perp}') U^{\dagger}(\boldsymbol{x}_{0\perp}') \right) \right\rangle + N_{c}$

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• Mixed space

$$\frac{\mathrm{d}\sigma^{\gamma^*+A\to c+X}}{\mathrm{d}^2 \boldsymbol{k}_{0\perp} \mathrm{d}y_0} \propto \int\limits_{\substack{\boldsymbol{x}_{0\perp}, \boldsymbol{x}_{1\perp}, \\ \boldsymbol{x}_{0\perp}'}} e^{i\boldsymbol{k}_{0\perp}(\boldsymbol{x}_{0\perp}-\boldsymbol{x}_{0\perp}')} \Big[D_{\boldsymbol{x}_{0\perp}-\boldsymbol{x}_{0\perp}'} - D_{\boldsymbol{x}_{0\perp}-\boldsymbol{x}_{1\perp}} - D_{\boldsymbol{x}_{0\perp}'-\boldsymbol{x}_{1\perp}} + 1 \Big] \psi(\boldsymbol{x}_{0\perp}, \boldsymbol{x}_{1\perp}) \psi^{\dagger}(\boldsymbol{x}_{0\perp}', \boldsymbol{x}_{1\perp})$$

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• Momentum space

$$\begin{aligned} \frac{\mathrm{d}\sigma^{\gamma^* + A \to c + X}}{\mathrm{d}^2 \mathbf{k}_{0\perp} \mathrm{d}y_0} \propto \int \frac{\mathrm{d}^2 \boldsymbol{\ell}_{\perp}}{(2\pi)^2} D(\boldsymbol{\ell}_{\perp}) \{\dots\} \\ \{\dots\} &= 2zm^2 \bigg[\frac{1}{m^2 + (\mathbf{k}_{0\perp} - \boldsymbol{\ell}_{\perp})^2} - \frac{1}{m^2 + \mathbf{k}_{0\perp}^2} \bigg]^2 + 4z \Big(z^2 + (1-z)^2 \Big) \bigg[\frac{\mathbf{k}_{0\perp} - \boldsymbol{\ell}_{\perp}}{m^2 + (\mathbf{k}_{0\perp} - \boldsymbol{\ell}_{\perp})^2} - \frac{\mathbf{k}_{0\perp}}{m^2 + \mathbf{k}_{0\perp}^2} \bigg]^2 \\ &- \text{Soft scale: } \boldsymbol{\ell}_{\perp}^2 \sim Q_{\mathrm{S}}^2 \\ &- \text{Quasi-real photon: } Q^2 \to 0 \end{aligned}$$

Full expression:
$$I_{\ell} = \int \frac{d^2 \ell_{\perp}}{(2\pi)^2} D(\ell_{\perp}) \{\dots\}$$

$$\begin{array}{ll} {\sf Full \ \ expression: \ I_\ell = \int \frac{{\rm d}^2 \ell_\perp}{(2\pi)^2} {\cal D}(\ell_\perp) \{\dots\}} \\ {\sf onst., \quad \ell_\perp \ll {\cal Q}_{\rm s}} \end{array}$$

$$D(\ell_{\perp}) \sim egin{cases} ext{const.}, & \ell_{\perp} \ll Q_{ ext{s}} \ 1/\ell_{\perp}^4, & \ell_{\perp} \gg Q_{ ext{s}} \end{cases}$$

$$\begin{split} \mathsf{Full} \ \ \mathsf{expression:} \ \ \mathbf{I}_\ell &= \int \frac{\mathrm{d}^2 \boldsymbol{\ell}_\perp}{(2\pi)^2} D(\boldsymbol{\ell}_\perp) \{\dots\} \\ D(\boldsymbol{\ell}_\perp) &\sim \begin{cases} \mathsf{const.}, \quad \boldsymbol{\ell}_\perp \ll Q_{\mathrm{s}} \\ 1/\boldsymbol{\ell}_\perp^4, \quad \boldsymbol{\ell}_\perp \gg Q_{\mathrm{s}} \end{cases} \qquad \qquad \{\dots\} \sim \begin{cases} \boldsymbol{\ell}_\perp^2, \qquad \boldsymbol{\ell}_\perp \ll k_{0\perp} \ , \\ \mathsf{const.}, \quad \boldsymbol{\ell}_\perp \gg k_{0\perp} \ . \end{cases} \end{split}$$

$$\begin{array}{ll} \text{(A)} \quad \ell_{\perp} \ll Q_{\text{s}} \ll k_{0\perp} : \int_{0}^{Q_{\text{s}}} \mathrm{d}\ell_{\perp}\ell_{\perp}^{3} \\ \text{(B)} \quad Q_{\text{s}} \ll \ell_{\perp} \ll k_{0\perp} : \int_{Q_{\text{s}}}^{k_{0\perp}} \mathrm{d}\ell_{\perp}\frac{1}{\ell_{\perp}} \\ \text{(C)} \quad Q_{\text{s}} \ll k_{0\perp} \ll \ell_{\perp} : \int_{k_{0\perp}}^{\infty} \mathrm{d}\ell_{\perp}\frac{1}{\ell_{\perp}^{3}} \end{array}$$

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• In the limit $\boldsymbol{\ell}_\perp \ll \boldsymbol{k}_{0\perp}$:

$$\Pi_{\ell_{\perp}}^{
m collinear} = \int^{m k_{0\perp}} rac{{
m d}^2 m \ell_{\perp}}{(2\pi)^2} D(\ell_{\perp}) m \ell_{\perp}^2 \Biggl\{ 2z rac{2m^2 k_{0\perp}^2 + [z^2 + (1-z)^2](k_{0\perp}^2 + m^4)}{(k_{0\perp}^2 + m^2)^4} \Biggr\}$$

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• Numerical comparison with the full expression

$$\mathrm{I}_{\ell} = \int \frac{\mathrm{d}^2 \boldsymbol{\ell}_{\perp}}{(2\pi)^2} D(\boldsymbol{\ell}_{\perp}) \{\dots\}$$

Collinear limit, ℓ_{\perp} -integrand



Collinear limit, cross section



- Small $k_{D\perp}$: collinear result suppressed
 - Large $k_{D\perp}$: matching up to some normalization

Comparison to CMS data

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- Rederived inclusive charm cross section for D^{0} in UPC
- Collinear limit vs full expression: matching at large $k_{\rm D\perp}$
- Comparison to CMS data: good agreement at small $k_{\mathrm{D}\perp}$
- To improve: NLO, fit BK initial condition including massive quarks *Casuga, Karhunen, Mäntysaari (2024) ⊗*