



# **Asymmetric collisions in Madgraph5 Impact on nPDF of future D0 and B+ measurements with SMOG2**

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This work has been supported by STRONG-2020 "The strong interaction at the frontier of knowledge: fundamental research and applications" which received funding from the European Union's Horizon 2020 research and innovation programme under grant agreement No 824093.

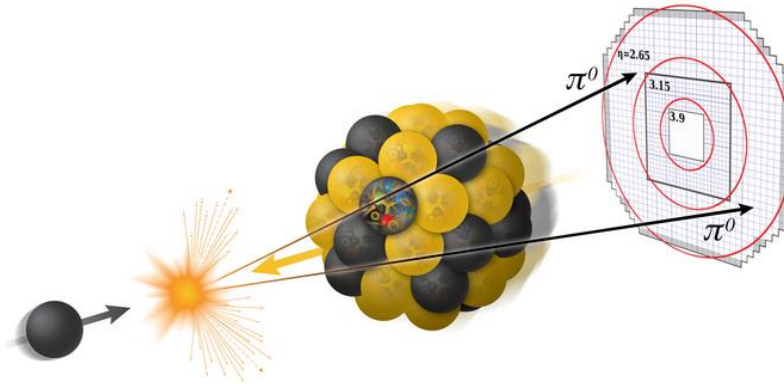
Synergies between LHC and EIC for quarkonium physics  
08-13 July 2024, Trento, Italy

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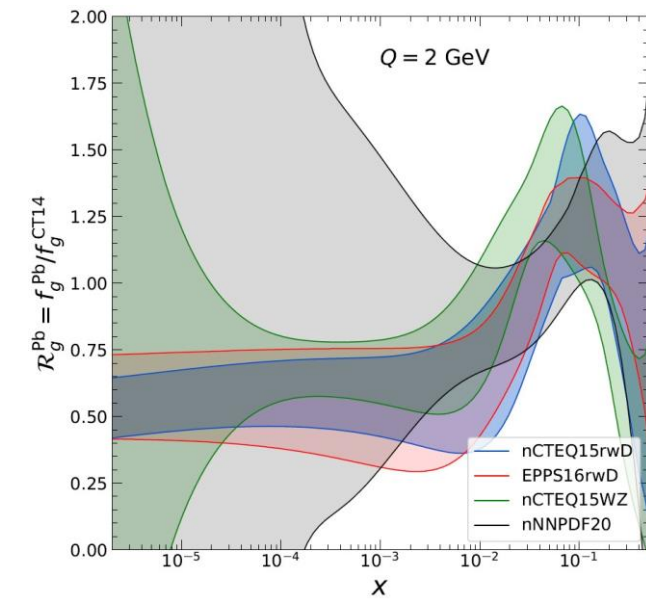
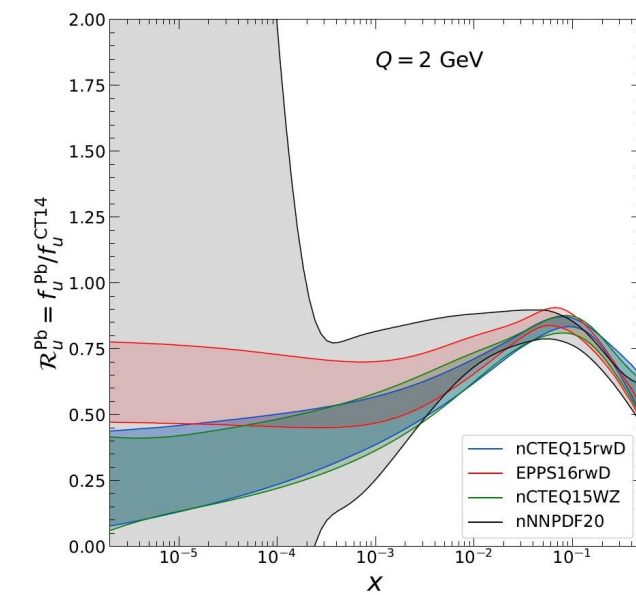
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# Motivation

- Study quark and gluon content of nucleons and nuclei in
  - hadron-hadron scattering,
  - hadron-nucleus scattering,
  - or any asymmetric reactions (nucleus/hadron A + nucleus/hadron B), described by Parton Distribution Functions (PDF)



- Evaluate the baseline for more sophisticated studies, like:
  - new state of matter in heavy-ion collisions,
  - charm and beauty quark production,
  - quarkonium productions and
  - the interpretation of the LHC and RHIC data.
- Develop a reliable and high-precision tool for feasibility studies for future measurements and experimental programs



# Context (State of art)

- There is **no NLO** code for **asymmetric hadron** collisions with **radiative corrections**
  - Modified versions of
    - **MCFM** (Monte Carlo for FeMtobarn processes)
    - **FEWZ** (FEWZ: A Fully Exclusive Numerical Code for QCD and EW Correction to Drell-Yan Process)
- are used by **High Energy Physics** community (they are **private** and **not** properly **validated**)
- It's important to provide such a **tool** for **phenomenology**
  - The **code** should be **interfaceable** even with generators like **Pythia**
  - It should be able to perform calculations at high orders of **QCD (Leading Order, Next-to-Leading Order)** for any types of **Parton Distribution Functions**, or example:
    - **nuclear,**
    - **nucleon,**
    - **pion,**
    - **photon, etc.**

# Framework – Collinear factorization

Cross sections in collinear factorization and perturbative QCD

$$d\sigma = \sum_{a,b} \int dx_1 dx_2 \underbrace{f_a(x_1, \mu_F) f_b(x_2, \mu_F)}_{\text{Parton density functions}} \underbrace{d\hat{\sigma}_{ab \rightarrow K}(\hat{S}, \mu_F, \mu_R)}_{\text{Parton-level (differential) Cross section}}$$

where the **partonic cross section** is calculated using:

$$\hat{\sigma} = \sigma^{Born} \left( 1 + \frac{\alpha_s}{2\pi} \sigma^{(1)} + \left( \frac{\alpha_s}{2\pi} \right)^2 \sigma^{(2)} + \left( \frac{\alpha_s}{2\pi} \right)^3 \sigma^{(3)} + \dots \right)$$

Leading order

Next-to-leading order

Next-to-next-to-leading order

For **charm, beauty, quarkonium** production, the scales are small and  $\alpha_s$  is large (0.15 ~ 0.25), **NLO corrections are very large and cannot be neglected.**

Such processes are usually accompanied with the **largest nuclear corrections** in proton-nucleus and nucleus-nucleus collisions

# Framework - PDFs

**Parton-distribution functions (PDFs):** essential link between hadronic cross sections and perturbatively calculable partonic cross sections

**Challenging situation for PDFs of nucleons inside nuclei (nPDFs):** nuclear data significantly more complex to collect with two additional degrees of freedom (**protons** and **neutrons**)

**nPDFs and PDFs give information on:**

- the **nuclear / hadronic structure** in terms of quarks and gluons;
- the **initial state** of relativistic heavy-ion collisions,  
to use **perturbative probes** of the **Quark Gluon Plasma** to study its properties
- **nPDFs** cannot be computed  
and similarly, to the proton PDFs are **fit to experimental data**.  
*Only the evolution is perturbative*

- Collinear **factorization** in terms of nPDFs is **assumed** and should be tested case by case
- **Automating** computations of cross sections with nPDFs up to **NLO** is highly **desirable**

# Nuclear Modification Factors

For rare/hard probes [ $\sigma_{NN}^{probe} \ll \sigma_{NN}^{inel}$ ]

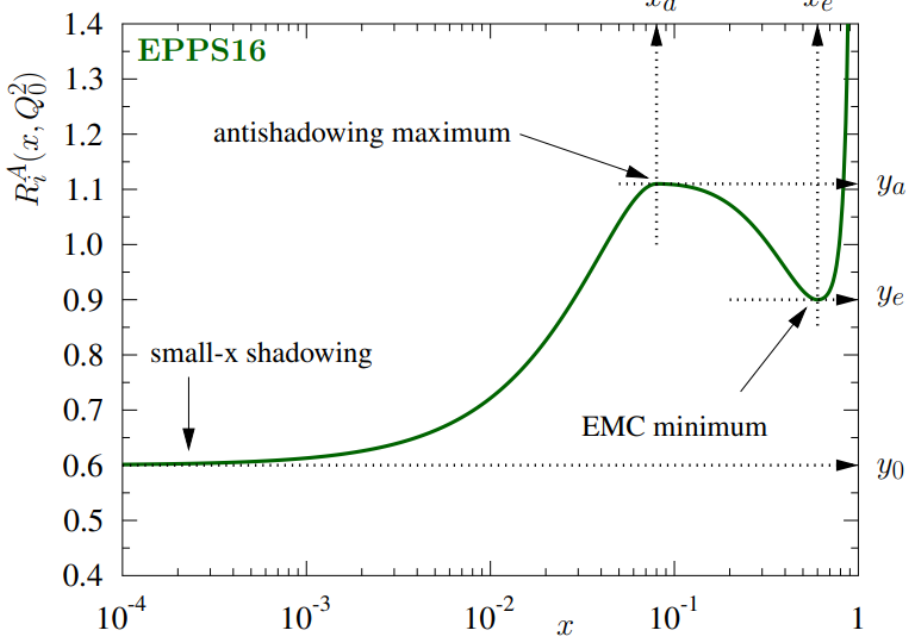
$$\sigma_{AB}^{probe} = A \times B \times \sigma_{NN}^{probe}$$

[Each probe is produced independently]

We can define [nuclear modification factors](#) ( $R_{AA}, R_{pA}$ ):

$$R_{AB} = \frac{\sigma_{AB}}{AB \sigma_{pp}}$$

$$R_{pA} \equiv \frac{\sigma_{pA}}{(1 \times A \times \sigma_{pp})}$$



arXiv:1612.05741v2 [hep-ph]

One expects:

- $R_{pA} > 1$  for  $x \gtrsim 0.8$  (**Fermi-motion region**),
- $R_{pA} < 1$  for  $0.25 \lesssim x \lesssim 0.8$  (**EMC region**),
- $R_{pA} > 1$  for  $0.1 \lesssim x \lesssim 0.25$  (**antishadowing region**)
- $R_{pA} < 1$  for  $x \lesssim 0.1$  (**shadowing region**)
- $R_{pA} \sim 1$ : **absence of nuclear effects**

# nPDFs and MG5

Any **PDFs** can be used in **MG5** up to **NLO** like *proton PDFs* with **LHAPDF** library

Currently only the symmetric mode is implemented

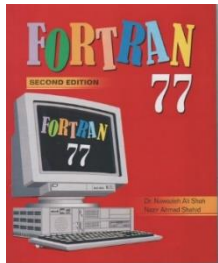
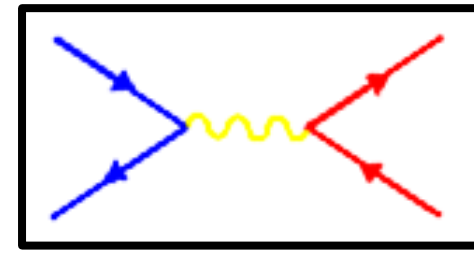
Reminder: we **assume** that

- **the factorization** of the **cross section** even in presence of **nuclear effects**
- all the **nuclear effects** can be accounted by **nPDFs** and thus **can be computed by MG5**.



# MadGraph

- **MG5\_aMC@NLO** is a **metacode**, i.e. a code generating another code
- **Matrix element generator** written in **Python**
- Can compute cross section and generates events at NLO with QCD corrections automatically
- Using **LHAPDF** can compute the cross section for any PDF in it with negligible additional CPU time (but only for **symmetrical** beam species)
- **Scale** and **PDF uncertainties** automatically computed and stored in **Histograms with Uncertainties** (HwU)
- **Output** in multiple formats (**root**, **HwU**, **gnuplot**, etc...)



```
*****
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*           W E L C O M E to
*           M A D G R A P H 5 _ a M C @ N L O
*
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*
*           VERSION 2.7.2           2020-03-17
*
*           The MadGraph5_aMC@NLO Development Team - Find us at
*           https://server06.fynu.ucl.ac.be/projects/madgraph
*           and
*           http://amcatnlo.web.cern.ch/amcatnlo/
*
*           Type 'help' for in-line help.
*           Type 'tutorial' to learn how MG5 works
*           Type 'tutorial aMCatNLO' to learn how aMC@NLO works
*           Type 'tutorial MadLoop' to learn how MadLoop works
*
*****
load MG5 configuration from ../input/mg5_configuration.txt
set collier to /projet/pth/safronov/MG5/MG5_aMC_v2_7_2/HEPTools/lib
set fastjet to /projet/pth/safronov/fastjet-install/bin/fastjet-config
set lhpdf to /projet/pth/safronov/MG5/MG5_aMC_v2_7_2/HEPTools/lhapdf6/bin/lhapdf-config
set ninja to /projet/pth/safronov/MG5/MG5_aMC_v2_7_2/HEPTools/lib
Using default text editor "vi". Set another one in ./input/mg5_configuration.txt
Using default eps viewer "evince". Set another one in ./input/mg5_configuration.txt
Using default web browser "firefox". Set another one in ./input/mg5_configuration.txt
Checking if MG5 is up-to-date... (takes up to 2s)
No new version of MG5 available
Loading default model: sm
INFO: Restrict model sm with file ../models/sm/restrict_default.dat .
INFO: Run "set stdout_level DEBUG" before import for more information.
INFO: Change particles name to pass to MG5 convention
Defined multiparticle p = g u c d s u~ c~ d~ s~
Defined multiparticle j = g u c d s u~ c~ d~ s~
Defined multiparticle l+ = e+ mu+
Defined multiparticle l- = e- mu-
Defined multiparticle vl = ve vm vt
Defined multiparticle vl~ = ve~ vm~ vt~
Defined multiparticle all = g u c d s u~ c~ d~ s~ a ve vm vt e- mu- ve~ vm~ vt~ e+ mu+ t b t~ b~ z w+ h w- ta- ta+
MG5_aMC>
```



# Framework – Collinear factorization

Cross sections in collinear factorization and perturbative QCD

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Leading order

Next-to-leading order

Next-to-next-to-leading order

For **charm, beauty, quarkonium** production, the scales are small and  $\alpha_s$  is large (0.15 ~ 0.25), **NLO corrections are very large and cannot be neglected.**

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# Framework – Collinear factorization

Cross sections in collinear factorization and perturbative QCD

$$\sigma_{h_1 h_2 \rightarrow X} = \sum_{a,b} \int dx_a dx_b \underbrace{f_{a/h_1}(x_a, \mu_F; LHAID_{h_1}) f_{b/h_2}(x_b, \mu_F; LHAID_{h_2})}_{\text{Parton density functions}} \underbrace{\hat{\sigma}_{ab \rightarrow X}(x_a, x_b, \mu_F, \alpha_S(\mu_R; LHAID_{h_2}))}_{\text{Parton-level (differential) Cross section}}$$

Parton density functions

Parton-level  
(differential)  
Cross section

where the **partonic cross section** is calculated using:

$$\hat{\sigma} = \underbrace{\sigma^{Born}}_{\text{Leading order}} \left( 1 + \underbrace{\frac{\alpha_S}{2\pi}}_{\text{Next-to-leading order}} \sigma^{(1)} + \underbrace{\left(\frac{\alpha_S}{2\pi}\right)^2}_{\text{Next-to-next-to-leading order}} \sigma^{(2)} + \left(\frac{\alpha_S}{2\pi}\right)^3 \sigma^{(3)} + \dots \right)$$

Leading order

Next-to-leading order

Next-to-next-to-leading order

For **charm, beauty, quarkonium** production, the scales are small and  $\alpha_S$  is large (0.15 ~ 0.25), **NLO corrections are very large and cannot be neglected.**

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# MadGraph in NLOAccess

**MG5\_aMC@NLO** is now available online with its full NLO version on **NLOAccess** (<https://nloaccess.in2p3.fr/MG5/>)

The asymmetric version of the code is available at:  
<https://github.com/mg5amcnlo/mg5amcnlo.git>



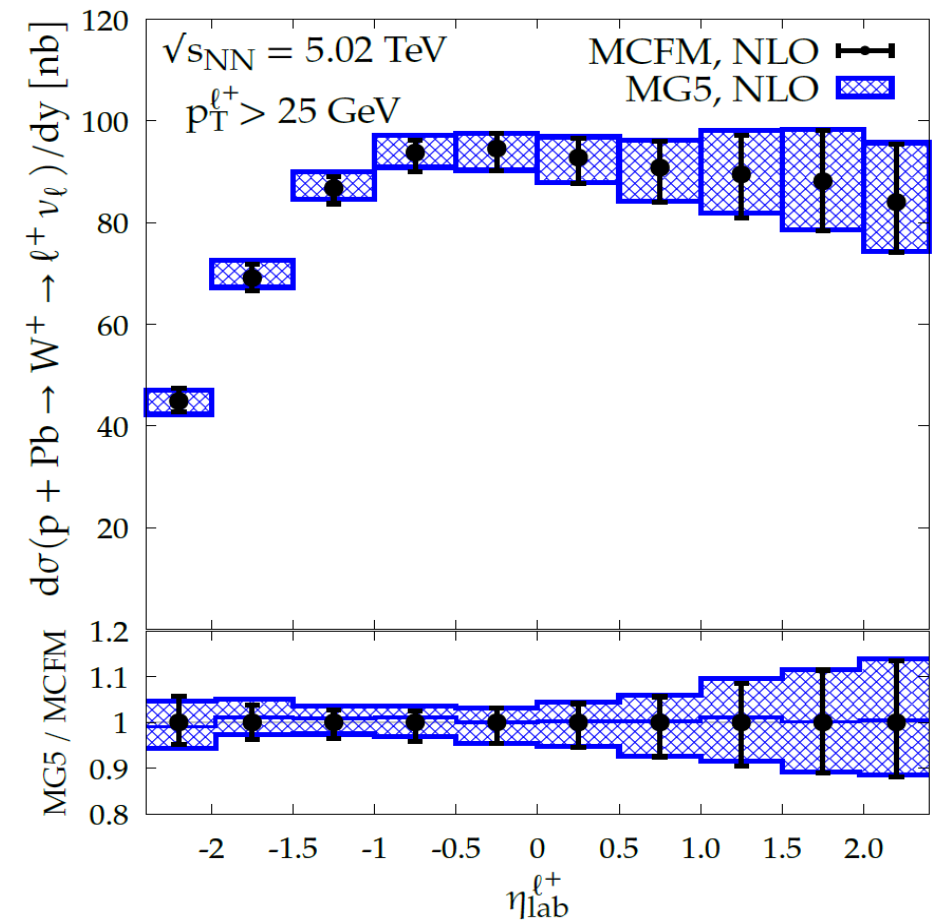
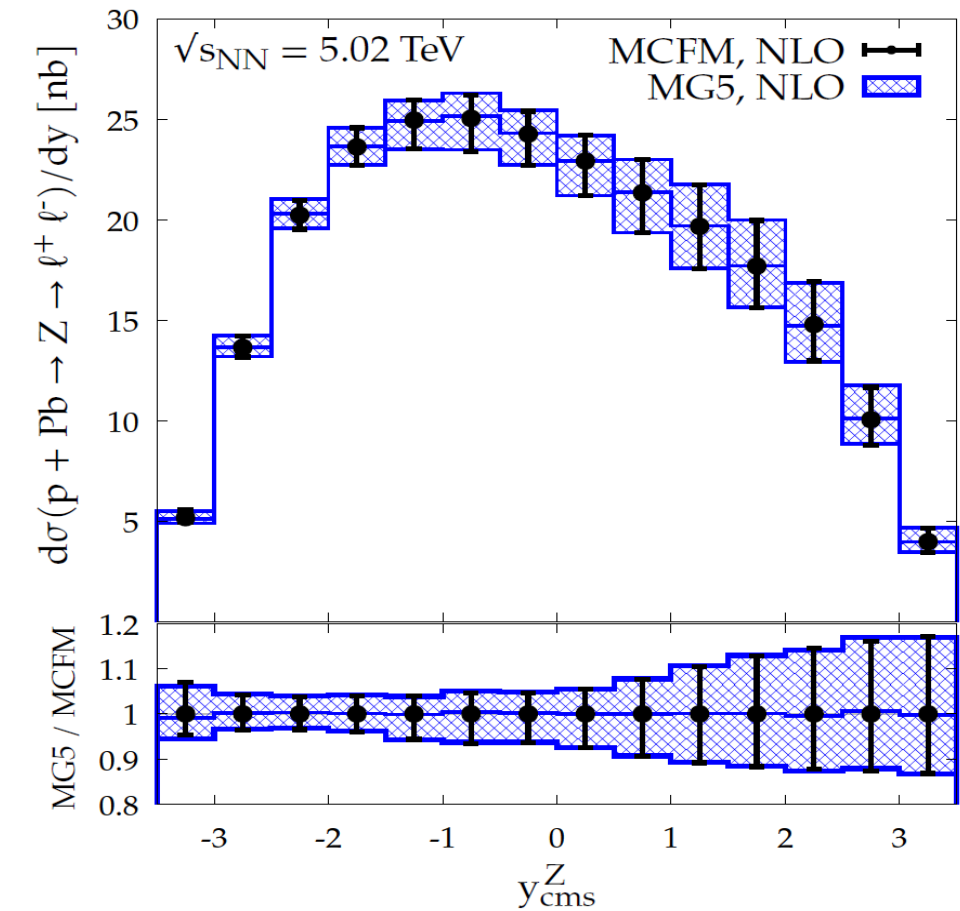
<https://nloaccess.in2p3.fr/>

A screenshot of the GitHub repository page for 'mg5amcnlo'. The repository is public and has 58 stars, 32 forks, and 3 unwatchers. It shows 66 branches and 28 tags. The main branch is ahead of the default branch by 39 commits. A recent commit by AntonSafr96 is shown, dated 01.07.2024. The repository contains folders for .github/workflows, HELAS, MadSpin, PLUGIN, Template, and aloha. The right sidebar shows repository statistics and options like 'Readme', 'View license', and 'Report repository'.

MG5 extension to asymmetric collisions will be included on NLOAccess

# Validations of MG5 in asymmetric collisions

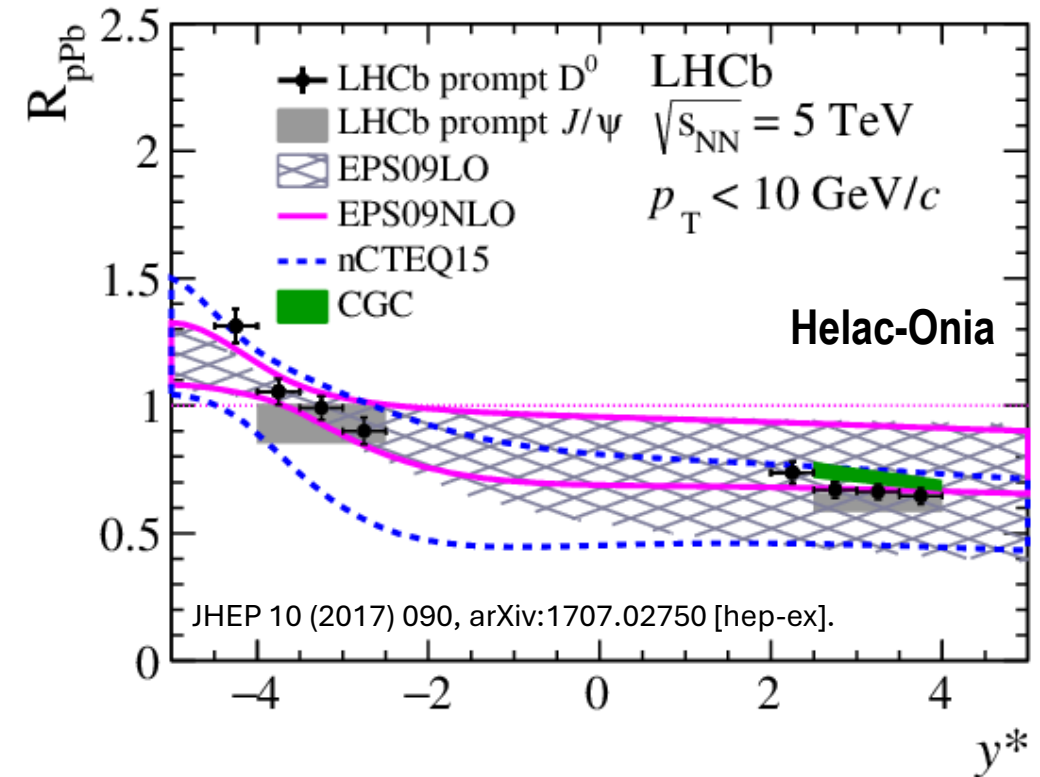
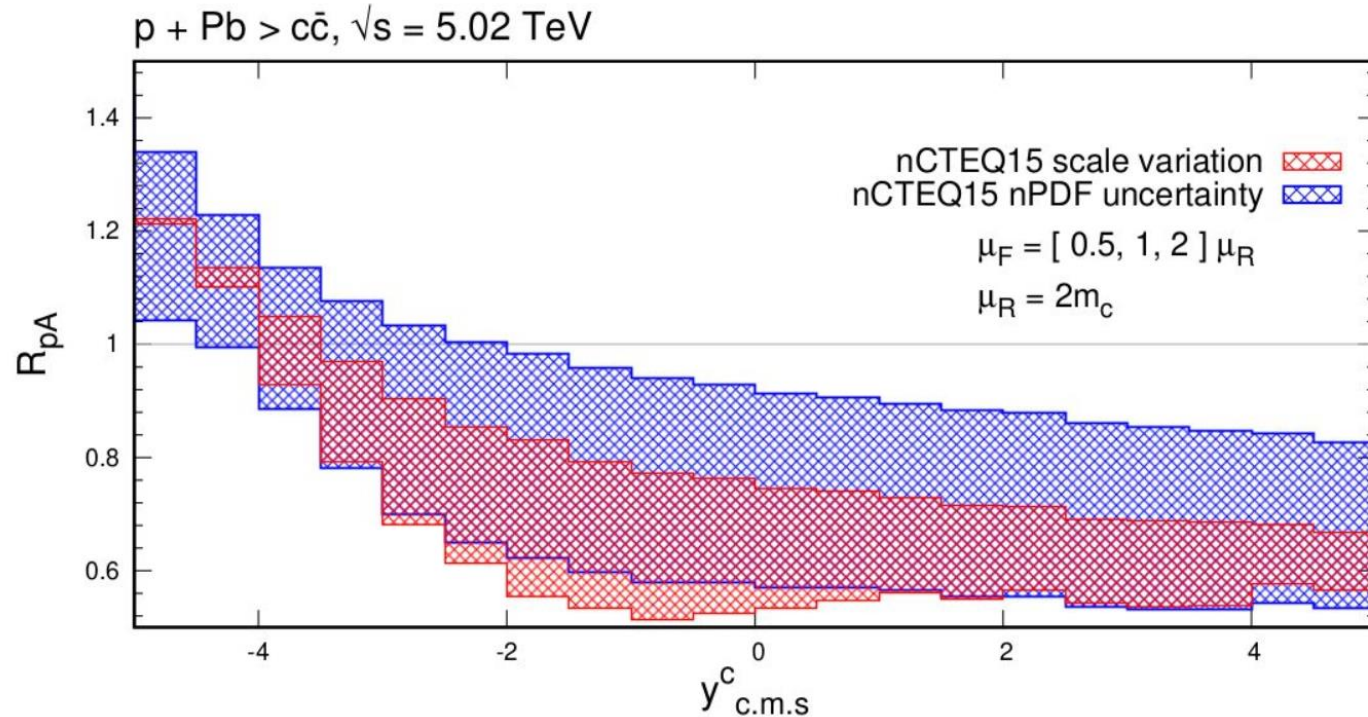
## Validation vs MCFM for W and Z production in proton-lead collisions at NLO



- **Very good agreement** between **MG5** and **MCFM**-based computations both for central value and uncertainties
- **Uncertainties match**, if **MCFM**-based computation done with **asymmetric** error estimation

# Example: $c$ production in $p$ Pb collision at LHC

*Charm quark production*

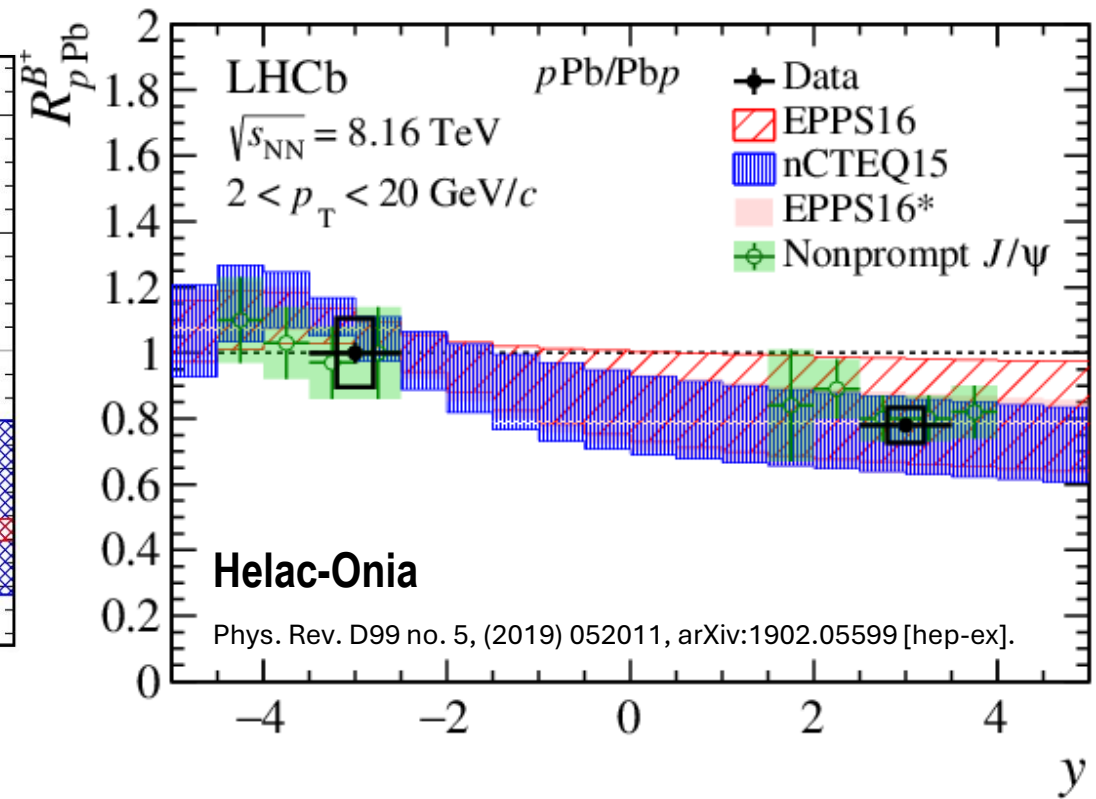
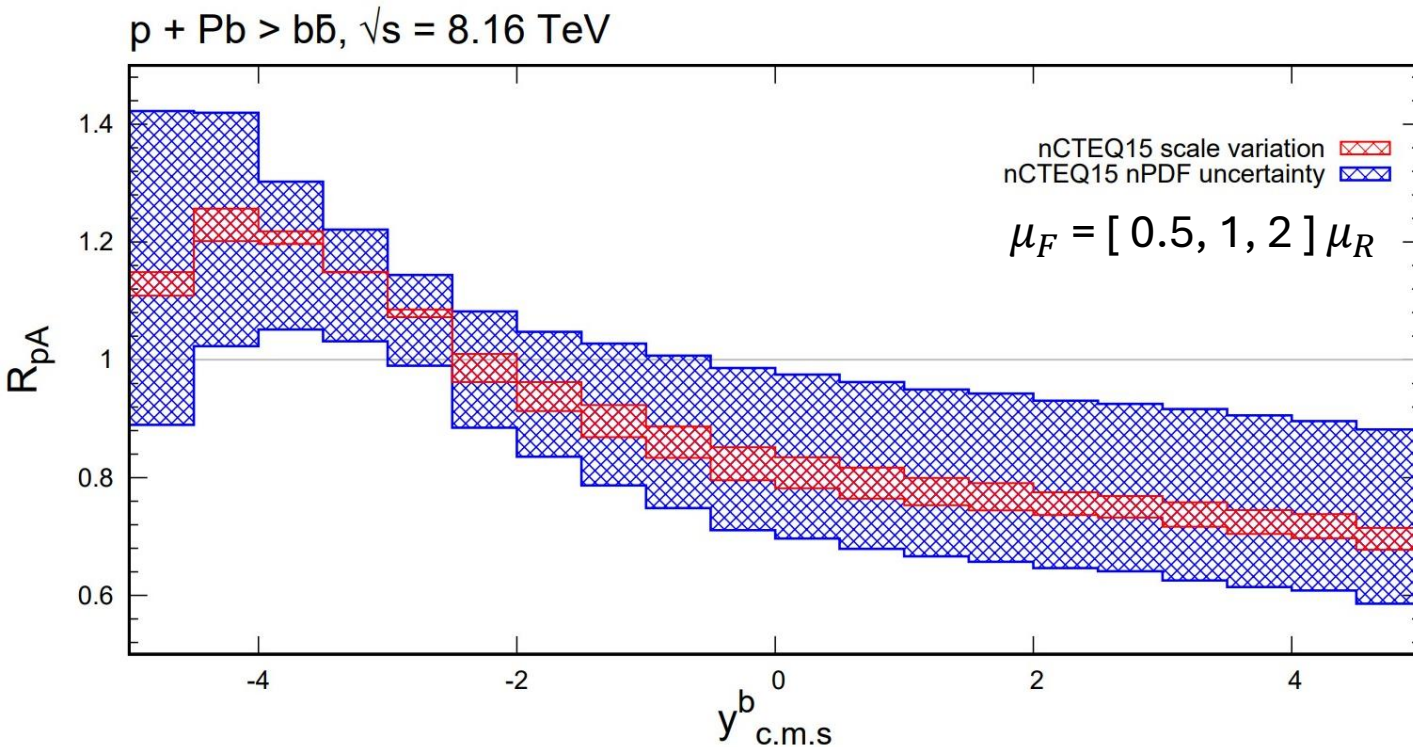


To make this plot, one just needs as input two numbers:  
**LHAPDF IDs** of proton and nCTEQ15 for Lead  
**Scale uncertainty is automatically** computed.

For charm production,  $\mu_F$  uncertainty nearly as **large** as the **nPDF uncertainty**.

# Example: $b$ production in $p$ Pb collision at LHC

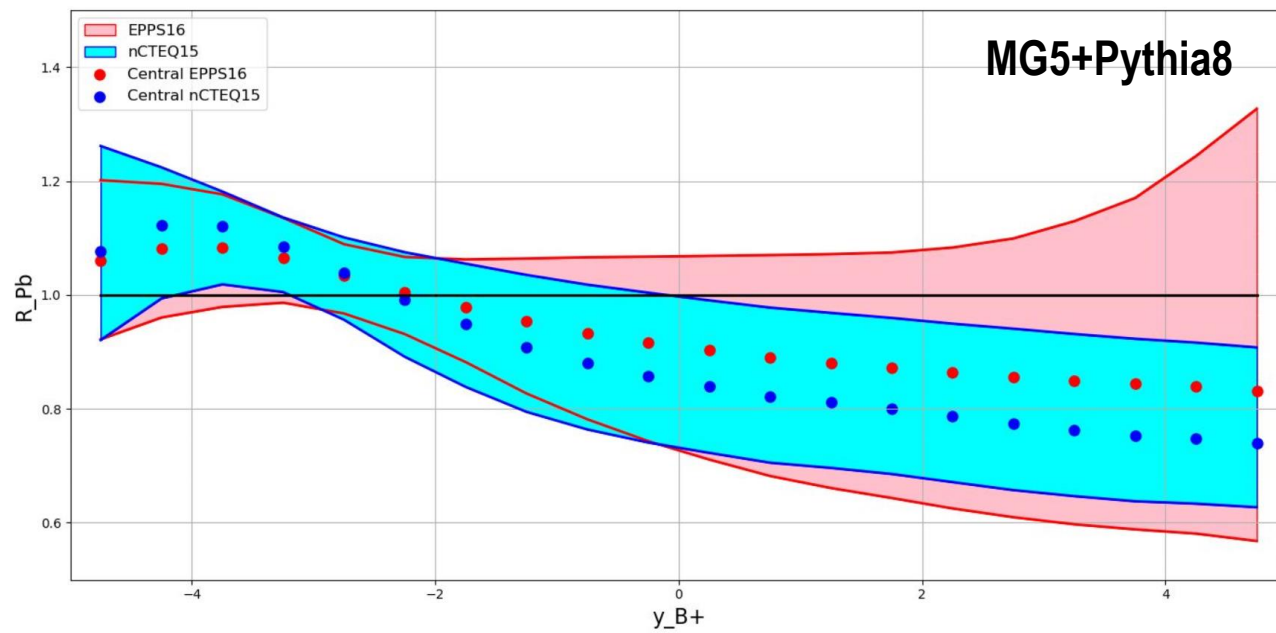
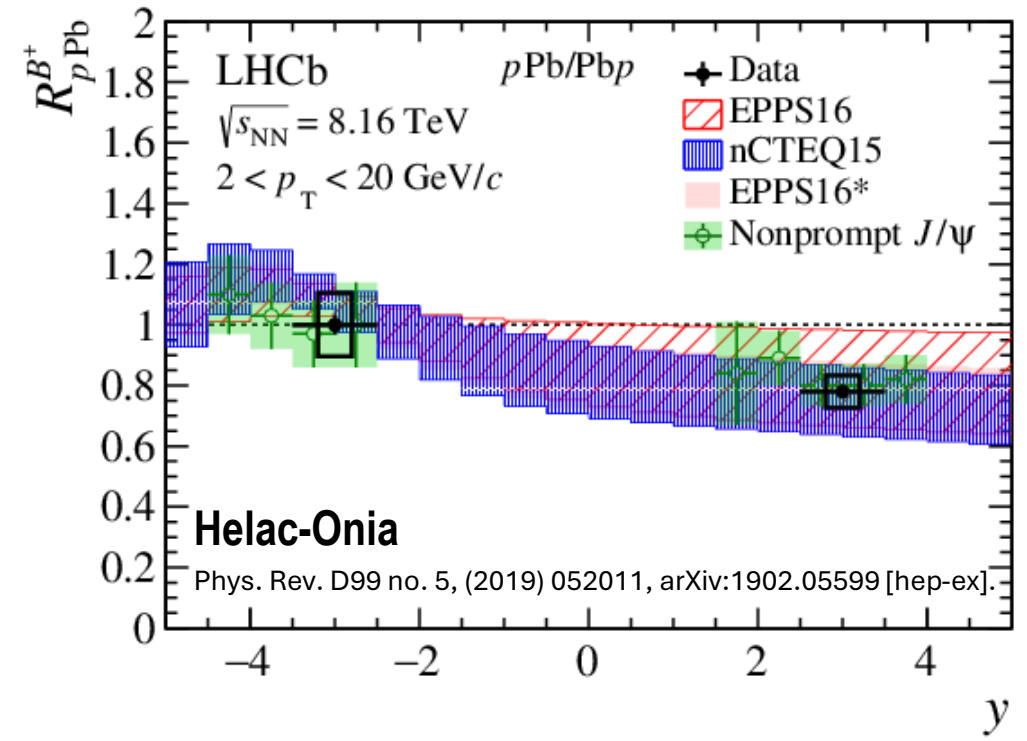
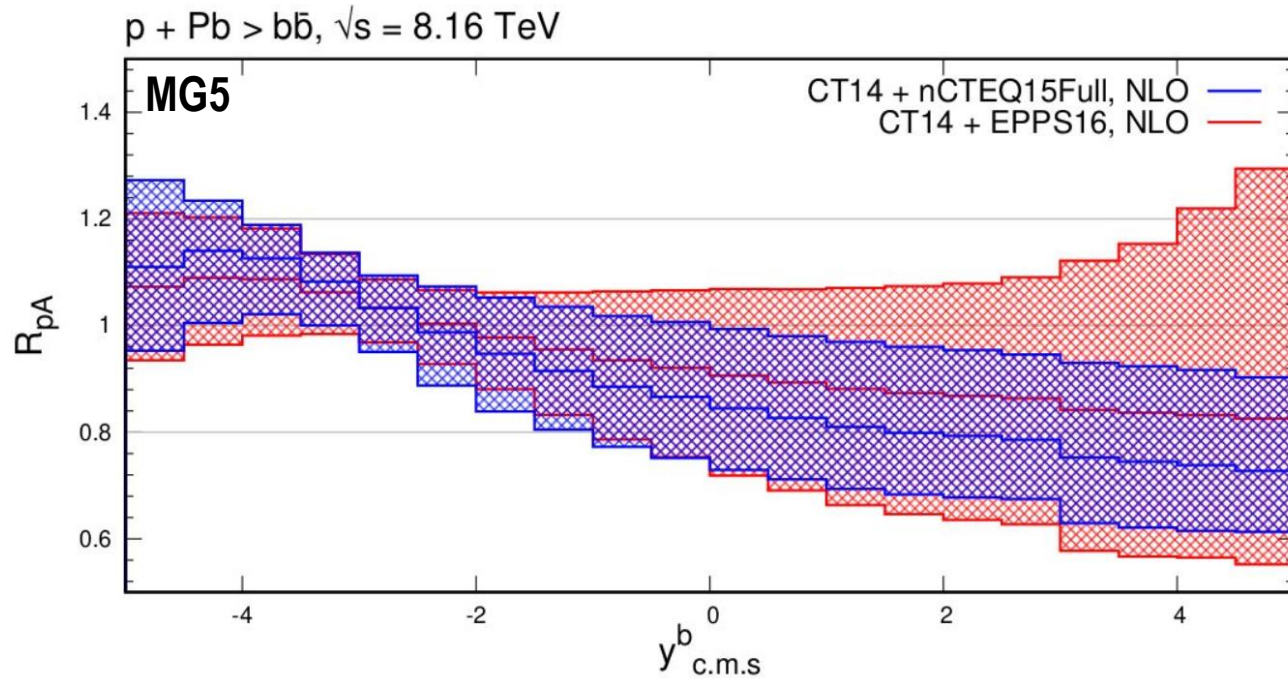
*Bottom quark production*



To make this plot, one just needs as input two numbers:  
**LHAPDF IDs** of proton and nCTEQ15 for Lead  
**Scale uncertainty is automatically** computed.



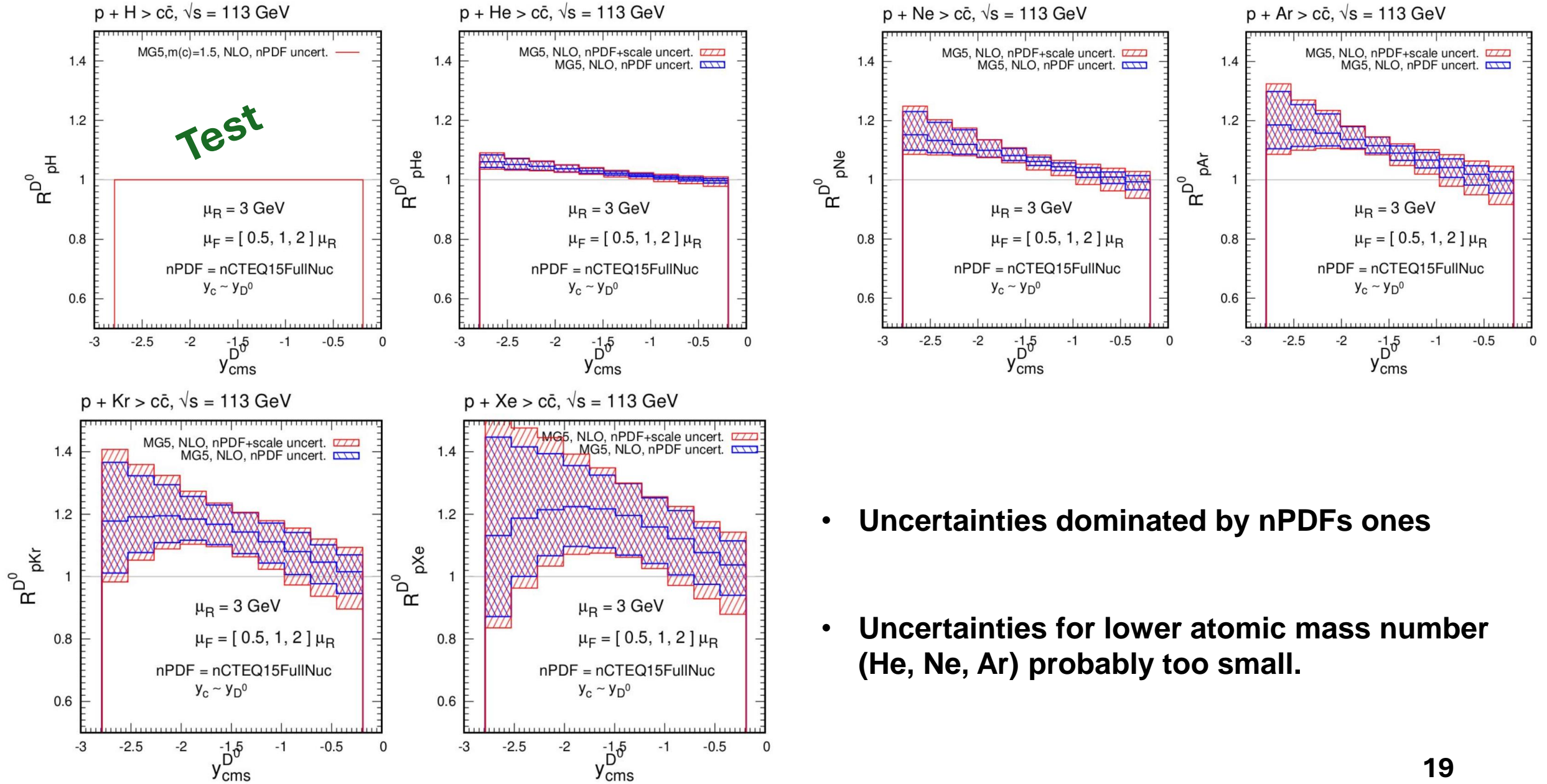
# Example: $B^+$ production in $pPb$ collision at LHC



**Impact on nPDF of  
future D0 and B+  
measurements with  
SMOG2**

**Directed by  
Cynthia Hadjidakis  
(IJClab)**

# $R^{D^0}_{pA}$ for different nuclei (H, He, Ne, Ar, Kr, Xe): rapidity dependence

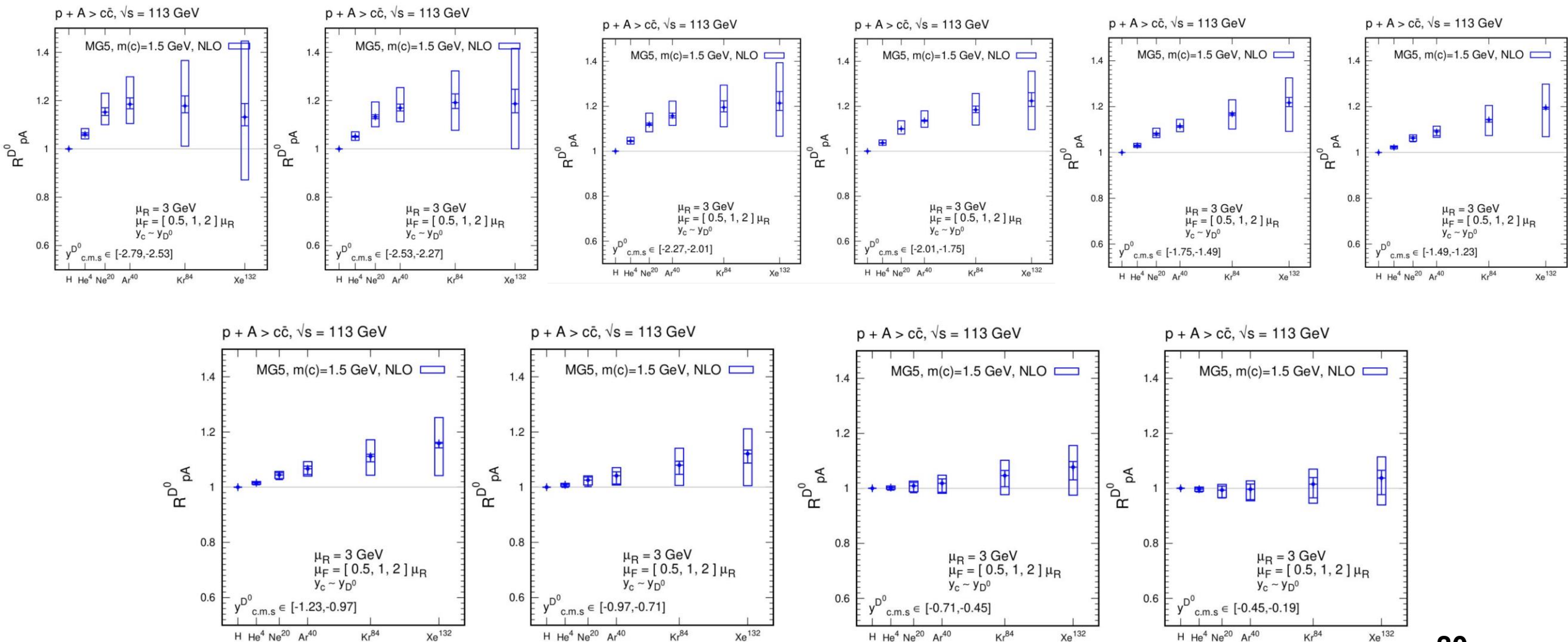


- **Uncertainties dominated by nPDFs ones**
- **Uncertainties for lower atomic mass number (He, Ne, Ar) probably too small.**

# $R^{D^0}_{pA}$ for different nuclei (H, He, Ne, Ar, Kr, Xe): A dependence

Vertical bar – scale uncertainties

Boxes – nPDF uncertainties



# Production of the D0 pseudo-data of SMOG2 LHCb using MG5 output

LHCb-PUB-2018-015  
04/12/2018

$$\sqrt{s} = 113 \text{ GeV}$$

$$2.0 < y_{\text{lab}} < 4.6$$

$$-2.79 < y_{\text{cms}} < -0.19$$

System	$\sqrt{s_{NN}}$ (GeV)	< pressure> ( $10^{-5}$ mbar)	$\rho_S$ ( $\text{cm}^{-2}$ )	$\mathcal{L}$ ( $\text{cm}^{-2}\text{s}^{-1}$ )	Rate (MHz)	Time (s)	$\int \mathcal{L}$ ( $\text{pb}^{-1}$ )
<u>pH<sub>2</sub></u>	115	4.0	$2.0 \times 10^{13}$	$6 \times 10^{31}$	4.6	$2.5 \times 10^6$	150
<u>pD<sub>2</sub></u>	115	2.0	$1.0 \times 10^{13}$	$3 \times 10^{31}$	4.3	$0.3 \times 10^6$	9
<u>pAr</u>	115	1.2	$0.6 \times 10^{13}$	$1.8 \times 10^{31}$	11	$2.5 \times 10^6$	45
<u>pKr</u>	115	0.8	$0.4 \times 10^{13}$	$1.2 \times 10^{31}$	12	$2.5 \times 10^6$	30
<u>pXe</u>	115	0.6	$0.3 \times 10^{13}$	$0.9 \times 10^{31}$	12	$2.5 \times 10^6$	22
<u>pHe</u>	115	2.0	$1.0 \times 10^{13}$	$3 \times 10^{31}$	3.5	$3.3 \times 10^3$	0.1
<u>pNe</u>	115	2.0	$1.0 \times 10^{13}$	$3 \times 10^{31}$	12	$3.3 \times 10^3$	0.1
<u>pN<sub>2</sub></u>	115	1.0	$0.5 \times 10^{13}$	$1.5 \times 10^{31}$	9.0	$3.3 \times 10^3$	0.1
<u>pO<sub>2</sub></u>	115	1.0	$0.5 \times 10^{13}$	$1.5 \times 10^{31}$	10	$3.3 \times 10^3$	0.1
PbAr	72	8.0	$4.0 \times 10^{13}$	$1 \times 10^{29}$	0.3	$6 \times 10^5$	0.060
<u>PbH<sub>2</sub></u>	72	8.0	$4.0 \times 10^{13}$	$1 \times 10^{29}$	0.2	$1 \times 10^5$	0.010
<u>pAr</u>	72	1.2	$0.6 \times 10^{13}$	$1.8 \times 10^{31}$	11	$3 \times 10^5$	5

40

40

$$\sigma_{D_0 + \bar{D}_0} = 2 \times \sigma_c \times \varepsilon_{eff} \times Br \times f(c \rightarrow D_0) \times \mathcal{L}$$

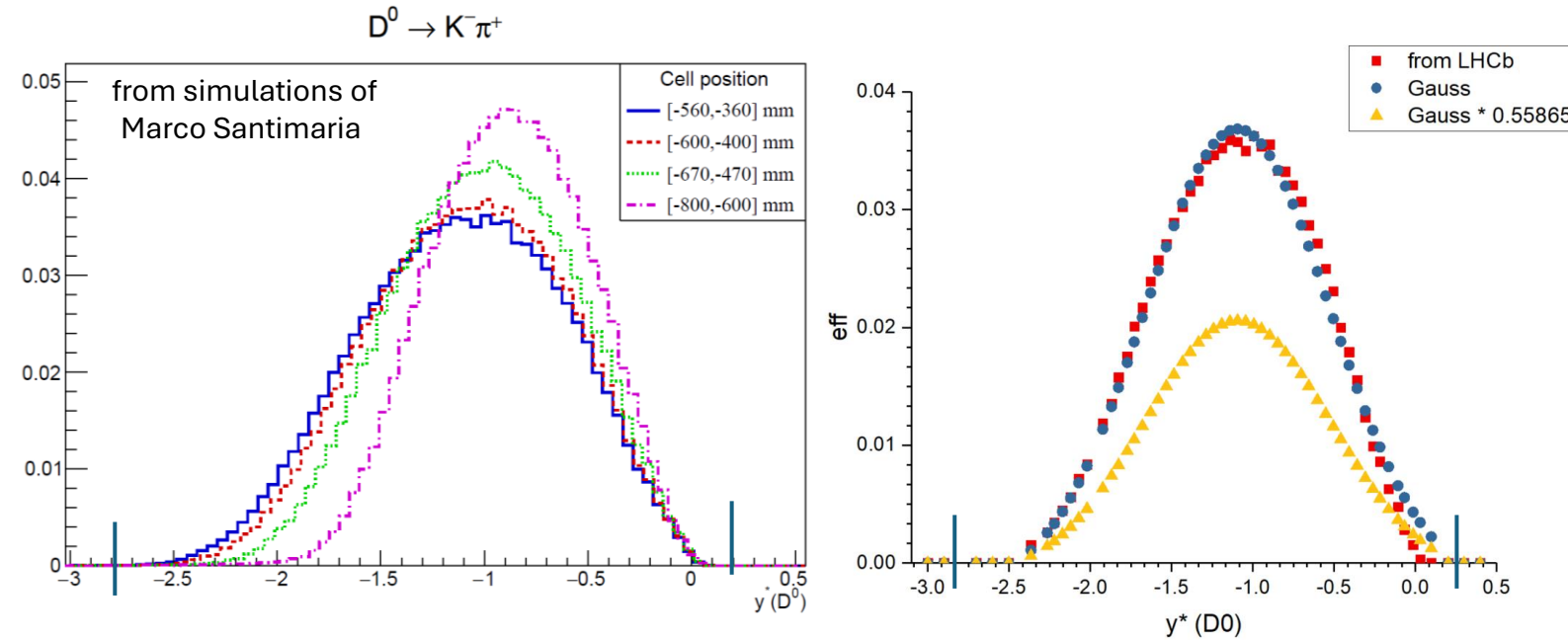
$$Br(D^0 \rightarrow K^- \pi^+) = (3.9 \pm 0.09 \pm 0.12)\%$$

$$f(c \rightarrow D^0) = 0.542 \pm 0.024$$

$$f = \frac{A}{B} \quad \sigma_f^2 \approx f^2 \left[ \left( \frac{\sigma_A}{A} \right)^2 + \left( \frac{\sigma_B}{B} \right)^2 - 2 \frac{\sigma_{AB}}{AB} \right]^{[17]}$$

We have less statistics  
(limited by) for  $H_2$

# Calculation of the AccxEff and uncertainties



Acceptance efficiency integrated over  $y$  is taken as **0.5-1%** (we took **1%**)  
(priv. comm. Emilie Maurice)

To achieve this, we have to scale Gaussian from the LHCb by factor **0.56** (yellow distribution)

Uncertainties are considered as uncorrelated over  $y$ !

Systematic uncertainty estimations are based on pPb systematics for  $D^0$  production:

arXiv:1707.02750v1 (5 TeV)

arXiv:2205.03936v4 (8.16 TeV)

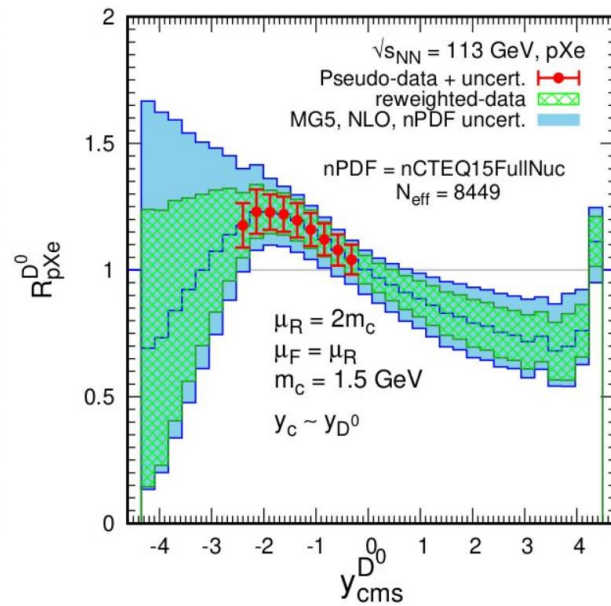
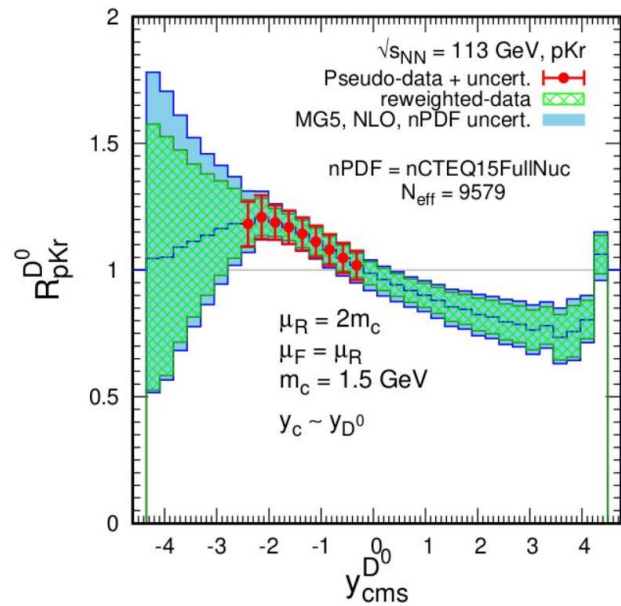
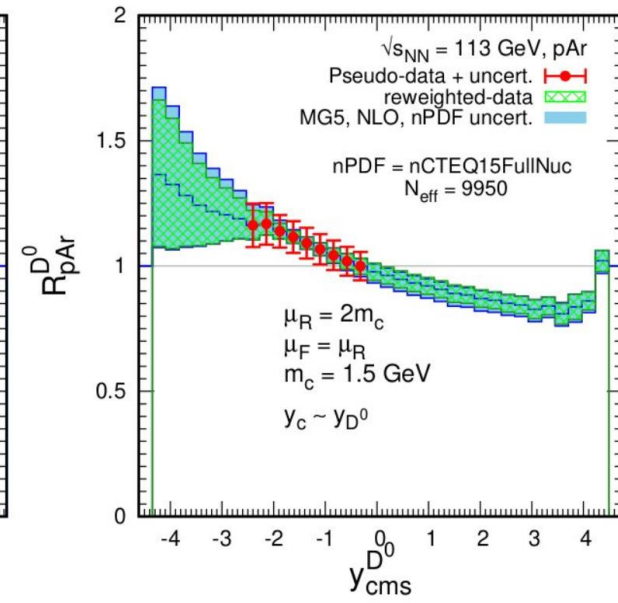
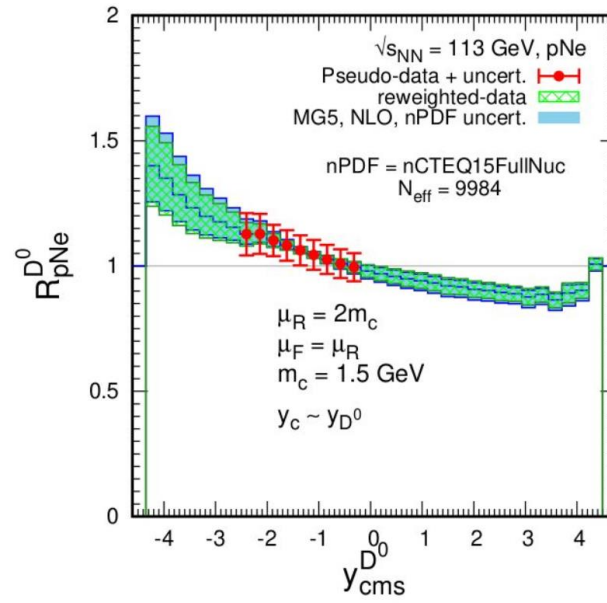
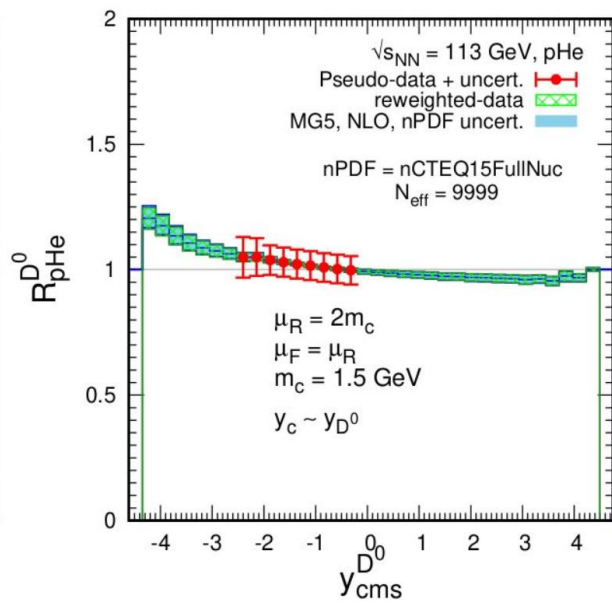
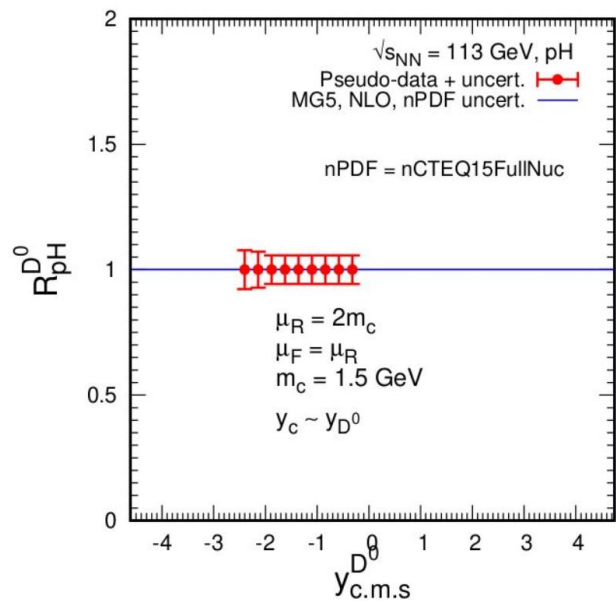
PID	Lum	Tracking	Signal	Prompt	Sim.sample	BR
1%	2%	3%	1%	0.1%	1%	0.8%

$$\frac{\delta\sigma_{syst}}{\sigma_{syst}} = 4\%$$

Except for bins  $-2.5 < y_{cms} < -2.0$

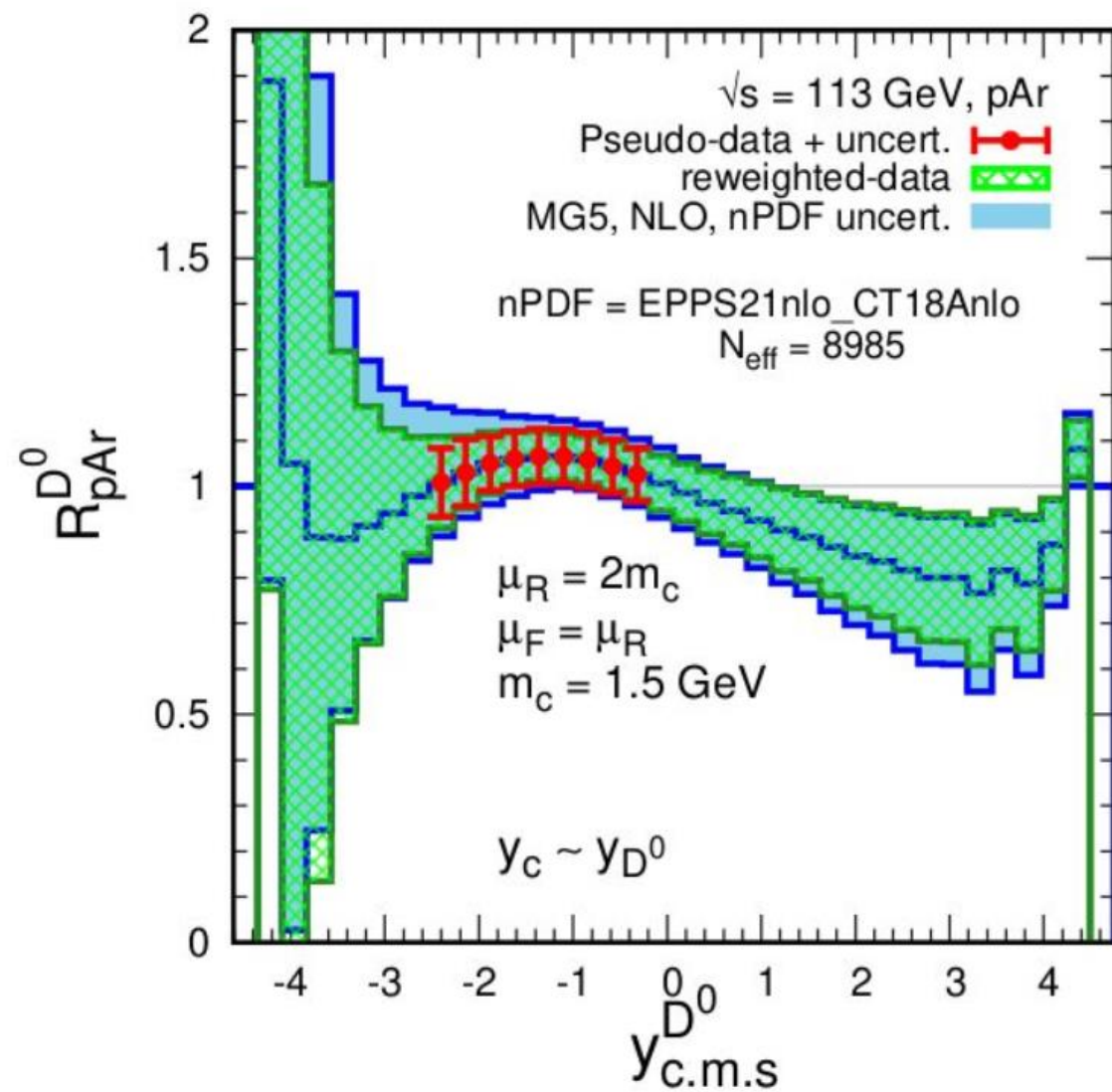
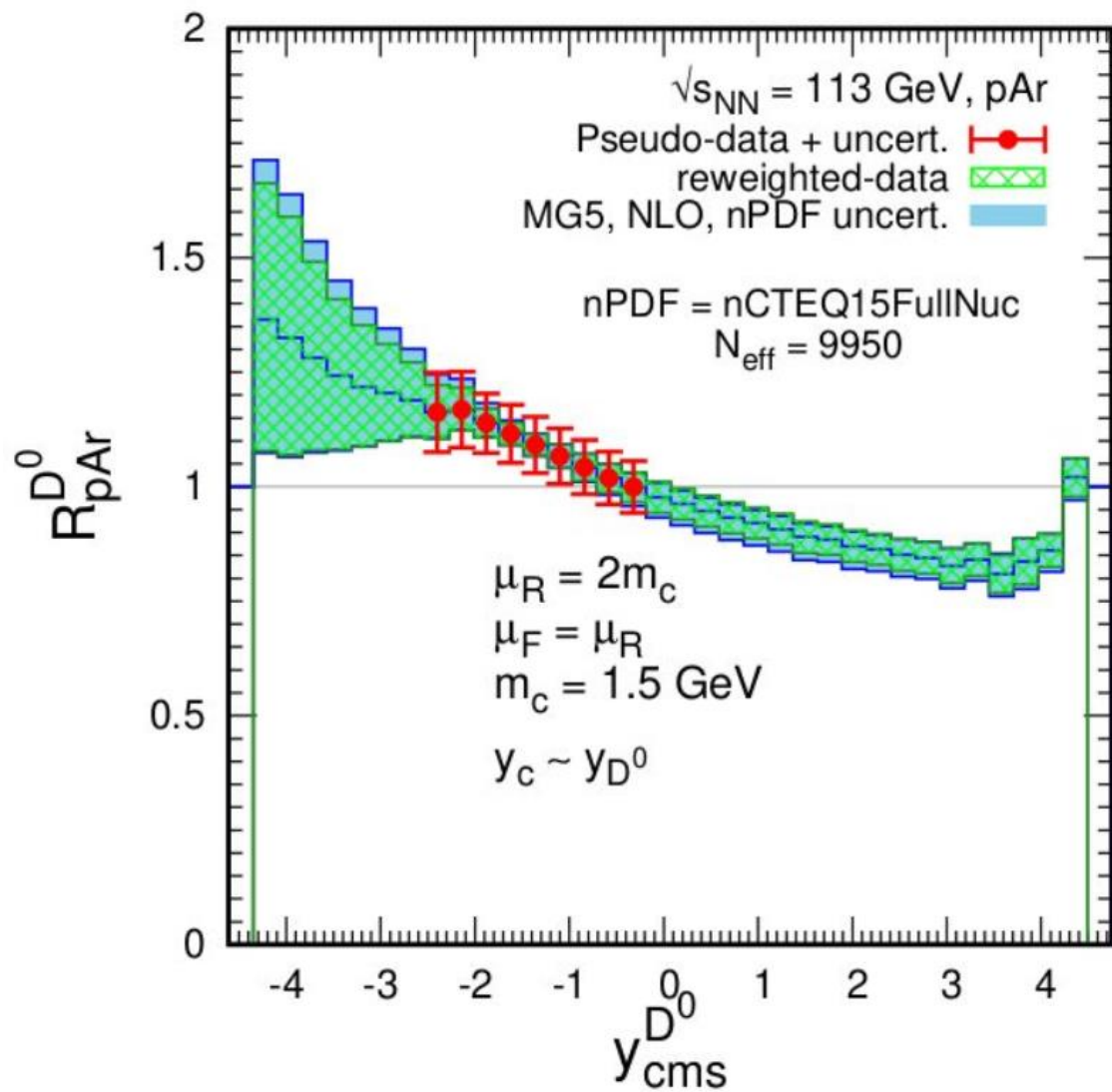
$$\frac{\delta\sigma_{syst}}{\sigma_{syst}} = 5\%$$

# $R_{pA}^{D^0}$ for different nuclei (H, He, Ne, Ar, Kr, Xe): reweighing of the nPDFs



Steps for the reweighting were taken from  
[arXiv:1610.02925v2](https://arxiv.org/abs/1610.02925v2)

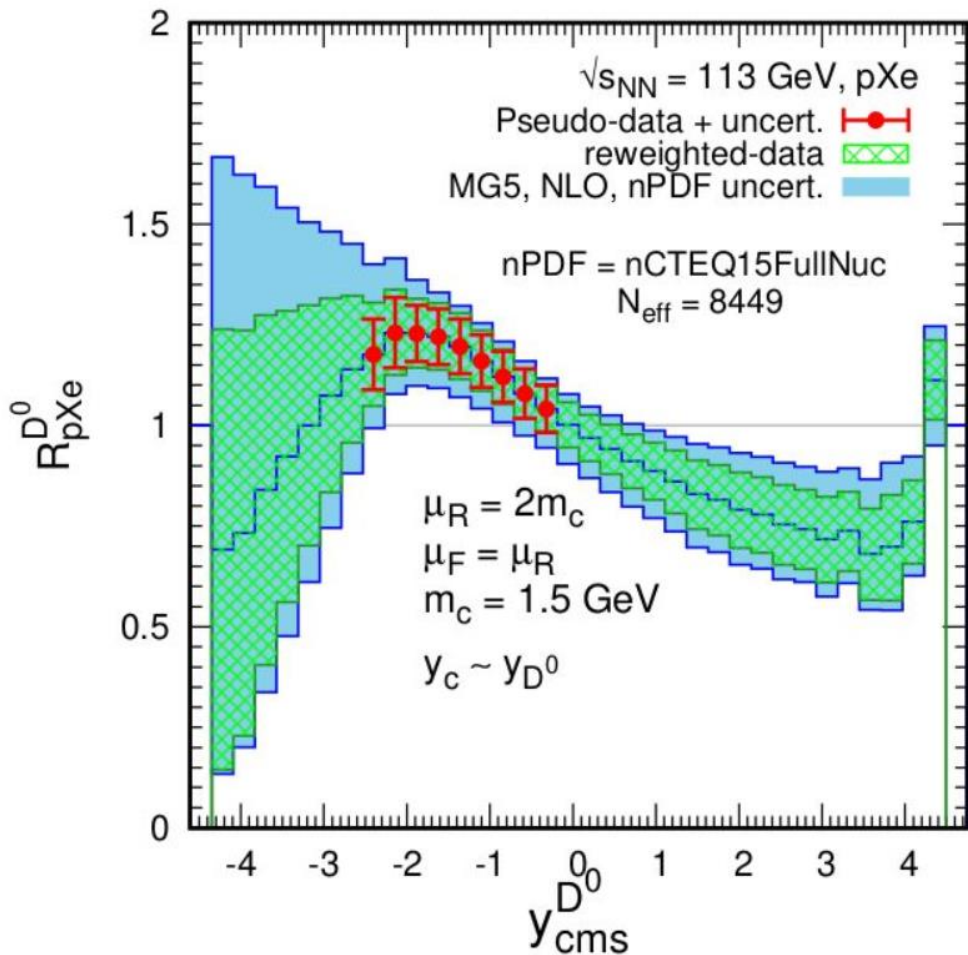
# $R_{pA}^{D^0}$ for Ar: reweighing of the nPDFs, nCTEQ vs EPPS21





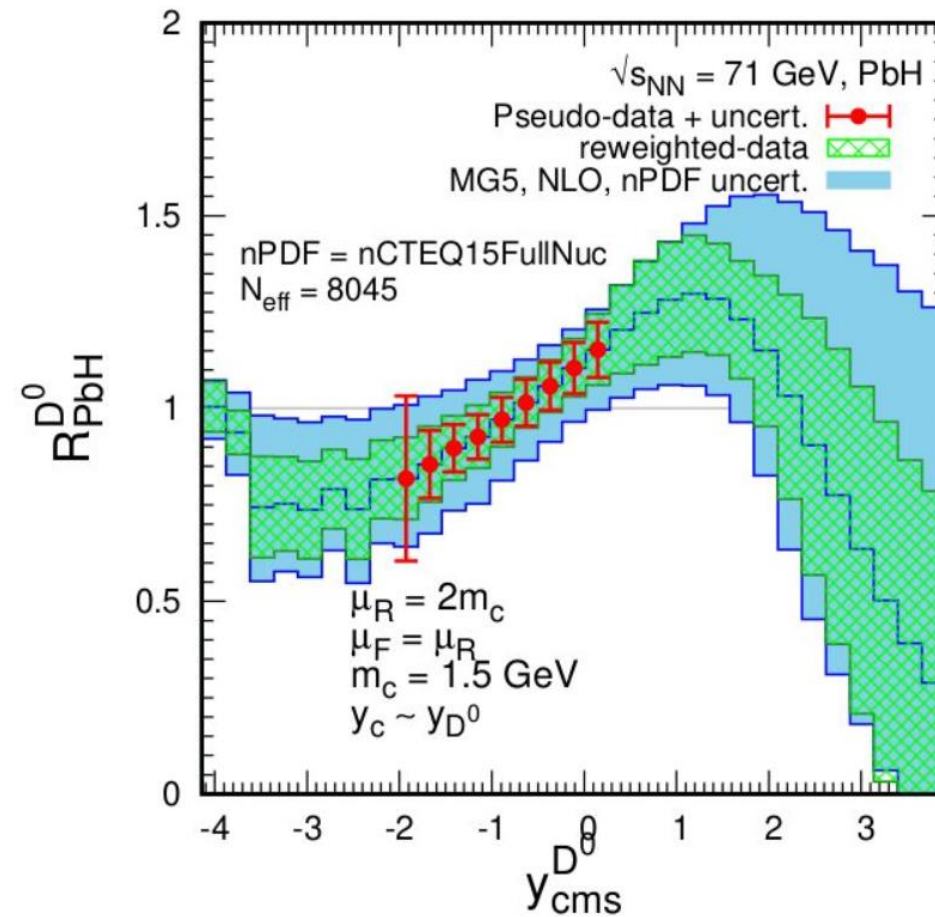
# Reweighting of the nPDFs: proton beam on Xe target vs Pb beam on H target

SMOG2  
probes different  
regions of the rapidity



$$x_2 \in [0.032, 0.33]$$

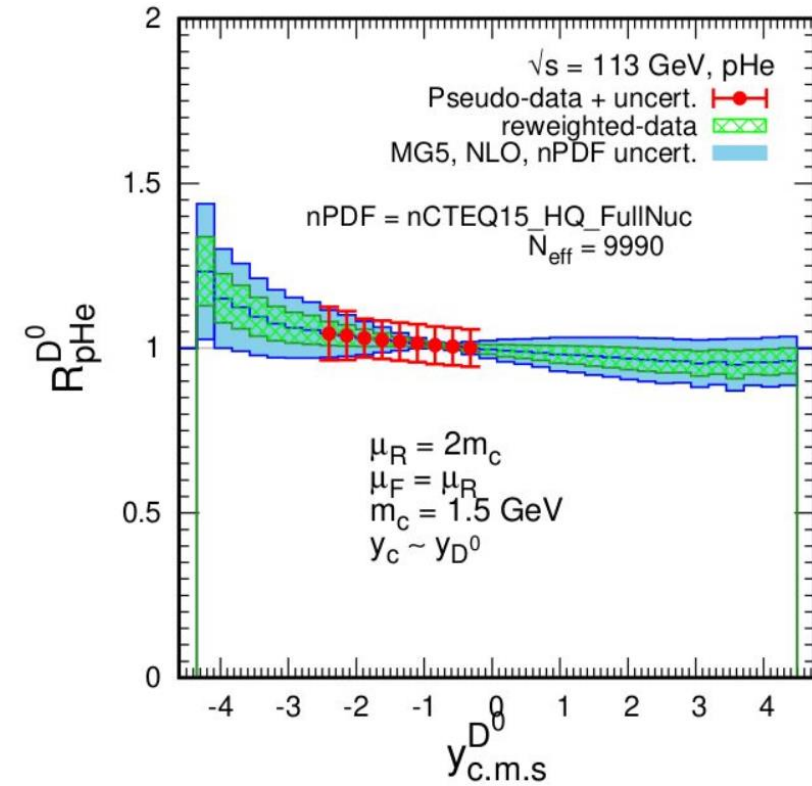
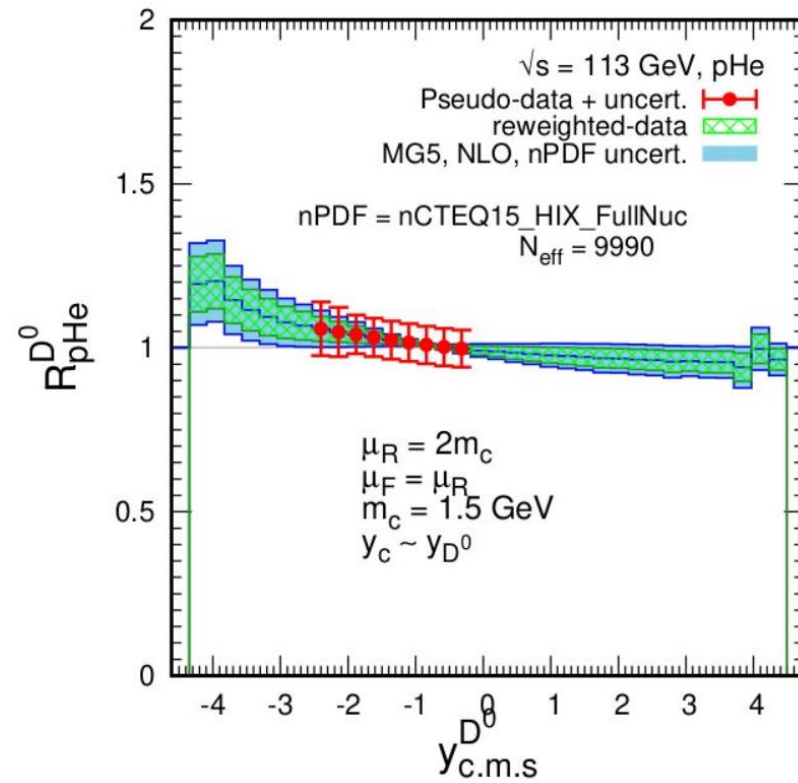
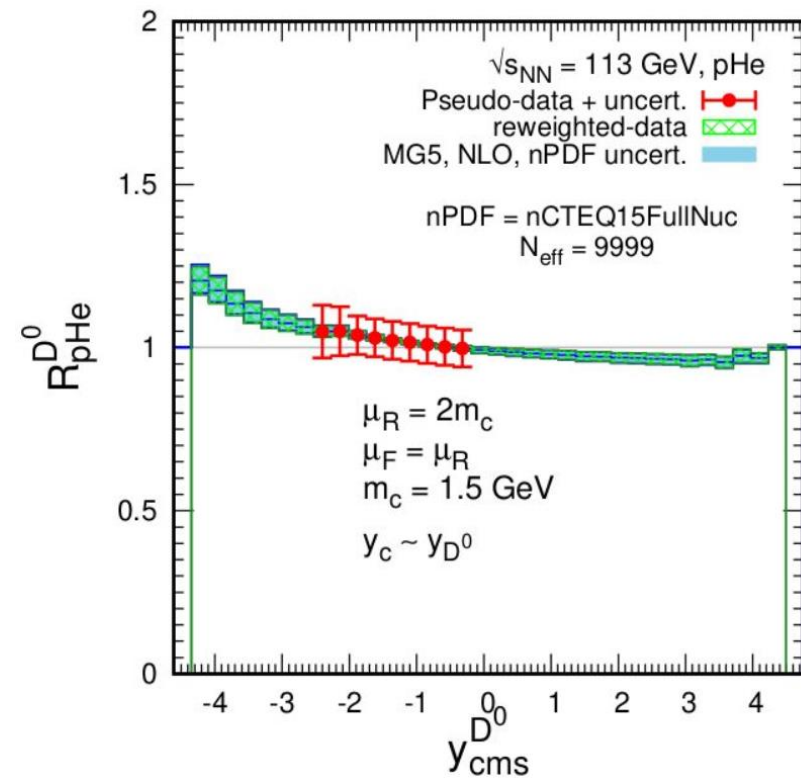
$$x_2 = \frac{2m_c}{113} e^{-y}$$



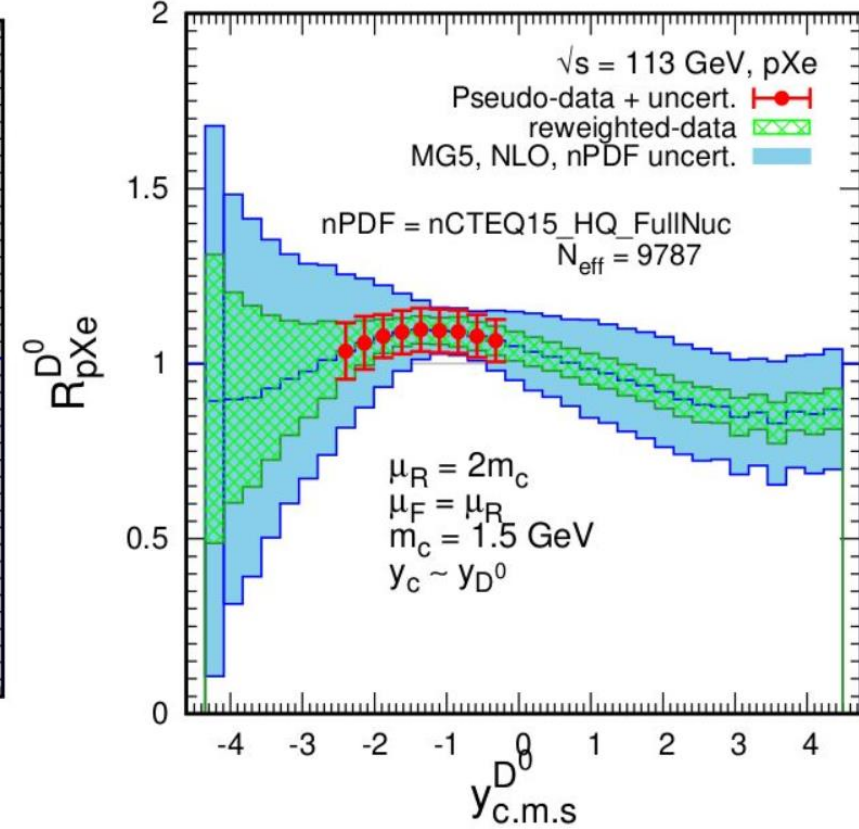
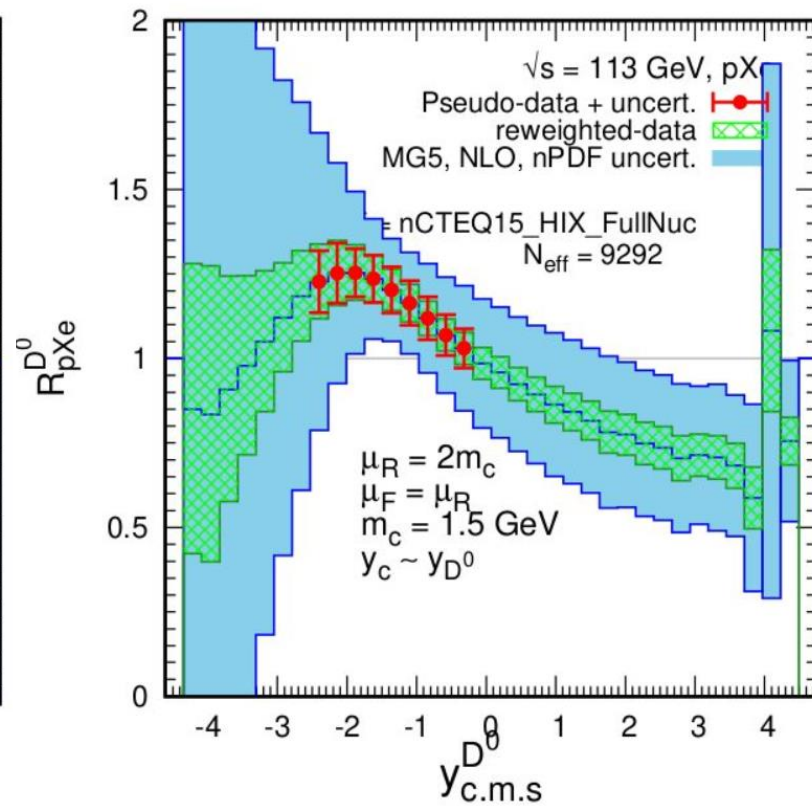
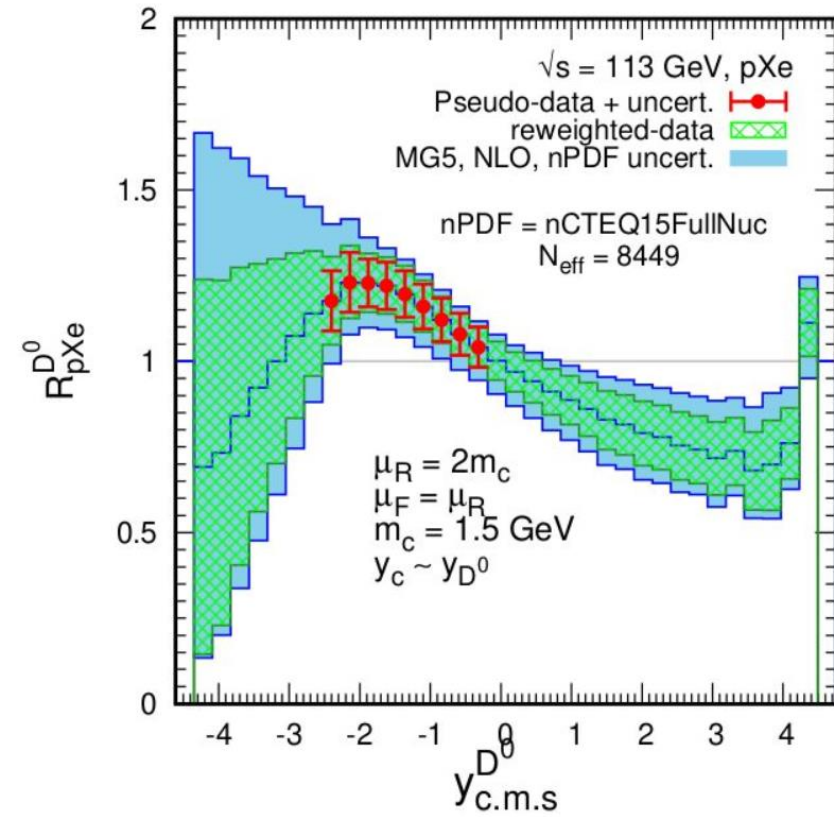
$$x_1 \in [0.0053, 0.056]$$

$$x_1 = \frac{2m_c}{71} e^y$$

# $R^{D^0}_{pA}$ for He and newer nPDFs: reweighting



# $R_{pA}^{D^0}$ for Xe for newer nPDFs: reweighting



# Production of the B<sup>+</sup> pseudo-data of SMOG2 LHCb using MG5 output

LHCb-PUB-2018-015  
04/12/2018

$$\sqrt{s} = 113 \text{ GeV}$$

$$2.0 < y_{\text{lab}} < 4.5$$

$$-2.79 < y_{\text{cms}} < -0.29$$

System	$\sqrt{s_{NN}}$ (GeV)	< pressure > (10 <sup>-5</sup> mbar)	$\rho_s$ (cm <sup>-2</sup> )	$\mathcal{L}$ (cm <sup>-2</sup> s <sup>-1</sup> )	Rate (MHz)	Time (s)	$\int \mathcal{L}$ (pb <sup>-1</sup> )
pH <sub>2</sub>	115	4.0	$2.0 \times 10^{13}$	$6 \times 10^{31}$	4.6	$2.5 \times 10^6$	150
pD <sub>2</sub>	115	2.0	$1.0 \times 10^{13}$	$3 \times 10^{31}$	4.3	$0.3 \times 10^6$	9
pAr	115	1.2	$0.6 \times 10^{13}$	$1.8 \times 10^{31}$	11	$2.5 \times 10^6$	45
pKr	115	0.8	$0.4 \times 10^{13}$	$1.2 \times 10^{31}$	12	$2.5 \times 10^6$	30
pXe	115	0.6	$0.3 \times 10^{13}$	$0.9 \times 10^{31}$	12	$2.5 \times 10^6$	22
pHe	115	2.0	$1.0 \times 10^{13}$	$3 \times 10^{31}$	3.5	$3.3 \times 10^3$	0.1
pNe	115	2.0	$1.0 \times 10^{13}$	$3 \times 10^{31}$	12	$3.3 \times 10^3$	0.1
pN <sub>2</sub>	115	1.0	$0.5 \times 10^{13}$	$1.5 \times 10^{31}$	9.0	$3.3 \times 10^3$	0.1
pO <sub>2</sub>	115	1.0	$0.5 \times 10^{13}$	$1.5 \times 10^{31}$	10	$3.3 \times 10^3$	0.1
PbAr	72	8.0	$4.0 \times 10^{13}$	$1 \times 10^{29}$	0.3	$6 \times 10^5$	0.060
PbH <sub>2</sub>	72	8.0	$4.0 \times 10^{13}$	$1 \times 10^{29}$	0.2	$1 \times 10^5$	0.010
pAr	72	1.2	$0.6 \times 10^{13}$	$1.8 \times 10^{31}$	11	$3 \times 10^5$	5

40  
40

We have less statistics  
(limited by)  
H<sub>2</sub> and Xe

$$\sigma_{B^+} = \sigma_b \times \varepsilon_{eff} \times Br \times f(b \rightarrow B^+) \times \mathcal{L}$$

$$Br(B^+ \rightarrow J/\psi K^+) = (1.020 \pm 0.019) \times 10^{-3}$$

$$Br(J/\psi \rightarrow \mu^+ \mu^-) = (5.961 \pm 0.033)\%$$

$$f(b \rightarrow B^0) = 40.4 \pm 0.6 \%$$

$$f = \frac{A}{B} \quad \sigma_f^2 \approx f^2 \left[ \left( \frac{\sigma_A}{A} \right)^2 + \left( \frac{\sigma_B}{B} \right)^2 - 2 \frac{\sigma_{AB}}{AB} \right]^{[17]}$$

# Production of the pseudo-data of SMOG2 LHCb using MG5 output

A Fixed-Target Programme at the LHC:  
Physics Case and Projected Performances for Heavy-Ion, Hadron, Spin and  
Astroparticle Studies

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I. Schienbein<sup>f,2</sup>, J. Seixas<sup>g,h,i,2</sup>, H.S. Shao<sup>l,2</sup>, A. Signori<sup>k,c,l,2</sup>, B. Trzeciak<sup>m,n,2</sup>, S.J. Brodsky<sup>o</sup>, G. Cavoto<sup>p</sup>,  
C. Da Silva<sup>q</sup>, F. Donato<sup>r</sup>, E.G. Ferreira<sup>s,4</sup>, I. Hrivnáčová<sup>a</sup>, A. Klein<sup>q</sup>, A. Kurepin<sup>u</sup>, C. Lorcé<sup>v</sup>, F. Lyonnet<sup>w</sup>,  
Y. Makdisi<sup>x</sup>, S. Porteboeuf Houssais<sup>y</sup>, C. Quintans<sup>h</sup>, A. Rakotozafindrabe<sup>z</sup>, P. Robbe<sup>a</sup>, W. Scandale<sup>aa</sup>,  
N. Topilskaya<sup>u</sup>, A. Uras<sup>ab</sup>, J. Wagner<sup>ac</sup>, N. Yamanaka<sup>aa,af,ad,ae</sup>, Z. Yang<sup>ag</sup>, A. Zelenski<sup>x</sup>

arXiv:1807.00603v3 [hep-ex] 28 Jan 2021

Target			LHCb							
			Proton beam ( $\sqrt{s_{NN}} = 115$ GeV)				Pb beam ( $\sqrt{s_{NN}} = 72$ GeV)			
			$\mathcal{L}$ [cm <sup>-2</sup> s <sup>-1</sup> ]	$\sigma_{inel.}$ [mb]	$\Gamma_{inel.}$ [kHz]	$\int \mathcal{L}$ [pb <sup>-1</sup> ]	$\mathcal{L}$ [cm <sup>-2</sup> s <sup>-1</sup> ]	$\sigma_{inel.}$ [mb]	$\Gamma_{inel.}$ [kHz]	$\int \mathcal{L}$ [nb <sup>-1</sup> ]
Internal gas target	Gas-Jet	H <sup>†</sup>	$4.3 \times 10^{30}$	39	168	43	$5.6 \times 10^{26}$	1.8	1	$5.6 \times 10^{-1}$
		H <sub>2</sub>	$1.0 \times 10^{33}$	39	40000	$1.0 \times 10^4$	$1.2 \times 10^{29}$	1.8	212	$1.2 \times 10^2$
		D <sup>†</sup>	$4.3 \times 10^{30}$	72	309	43	$5.6 \times 10^{26}$	2.2	1	$5.6 \times 10^{-1}$
		<sup>3</sup> He <sup>†</sup>	$3.4 \times 10^{32}$	117	40000	$3.4 \times 10^3$	$4.7 \times 10^{28}$	2.5	118	47
		Xe	$3.1 \times 10^{31}$	1300	40000	$3.1 \times 10^2$	$2.3 \times 10^{28}$	6.2	186	23
	Storage Cell	H <sup>†</sup>	$9.2 \times 10^{32}$	39	35880	$9.2 \times 10^5$	$1.2 \times 10^{29}$	1.8	212	$1.2 \times 10^2$
		H <sub>2</sub>	$1.0 \times 10^{33}$	39	40000	$1.0 \times 10^4$	$1.2 \times 10^{29}$	1.8	212	$1.2 \times 10^2$
		D <sup>†</sup>	$5.6 \times 10^{32}$	72	40000	$5.6 \times 10^3$	$8.8 \times 10^{28}$	2.2	194	88
		<sup>3</sup> He <sup>†</sup>	$1.3 \times 10^{33}$	117	40000	$1.3 \times 10^4$	$8.3 \times 10^{28}$	2.5	206	83
		Xe	$3.1 \times 10^{31}$	1300	40000	$3.1 \times 10^2$	$3.0 \times 10^{28}$	6.2	186	30
Internal solid target on beam halo	Wire Target	C (500 $\mu$ m)	$2.8 \times 10^{30}$	271	760	28	$5.6 \times 10^{26}$	3.3	2	$5.6 \times 10^{-1}$
		Ti (500 $\mu$ m)	$1.4 \times 10^{30}$	694	972	14	$2.8 \times 10^{26}$	4.7	1	$2.8 \times 10^{-1}$
		W (500 $\mu$ m)	$1.6 \times 10^{30}$	1700	2720	16	$3.1 \times 10^{26}$	6.9	2	$3.1 \times 10^{-1}$
Beam splitting	E1039	NH <sub>3</sub> <sup>†</sup>	$7.2 \times 10^{31}$	420	30240	$7.2 \times 10^2$	$1.4 \times 10^{28}$	19	259	14
		ND <sub>3</sub> <sup>†</sup>	$7.2 \times 10^{31}$	519	37368	$7.2 \times 10^2$	$1.4 \times 10^{28}$	22	314	14
	Unpolarised solid target	C (5 mm)	$2.8 \times 10^{31}$	271	7600	$2.8 \times 10^2$	$5.6 \times 10^{27}$	3.3	18	5.6
		Ti (5 mm)	$1.4 \times 10^{31}$	694	9720	$1.4 \times 10^2$	$2.8 \times 10^{27}$	4.7	13	2.8
		W (5 mm)	$1.6 \times 10^{31}$	1700	27200	$1.6 \times 10^2$	$3.1 \times 10^{27}$	6.9	21	3.1

$$\sqrt{s} = 113 \text{ GeV}$$

$$2.0 < y_{lab} < 4.5$$

$$-2.79 < y_{cms} < -0.29$$

# Calculation of the AccxEff and uncertainties

Table 3: Measured  $\varepsilon_{\text{tot}}$  of  $B^\pm$  events at 13 TeV, in the bins of  $B^\pm p_T$  and  $y$ . The efficiencies and uncertainties are in percent.

$p_T$ [GeV/c]	$2.0 < y < 2.5$	$2.5 < y < 3.0$	$3.0 < y < 3.5$	$3.5 < y < 4.0$	$4.0 < y < 4.5$
0.0 – 0.5	$0.8 \pm 0.1$	$4.7 \pm 0.1$	$7.6 \pm 0.2$	$6.8 \pm 0.2$	$2.3 \pm 0.1$

**AccxEff is  
Different for each of the y  
region**

**Systematic uncertainty estimations are based  
on pPb and pp systematics for B+ production:**

**arXiv:1902.05599v2 (8.16 TeV)**

**arXiv:1710.04921v2 (13 TeV)**

PID	Lum	BR	Binning	mass-fits	Acc	Rec	Track	Trigger
0.4%	2%	2.8%	2.5%	2%	0.2%	0.1%	2%	3%

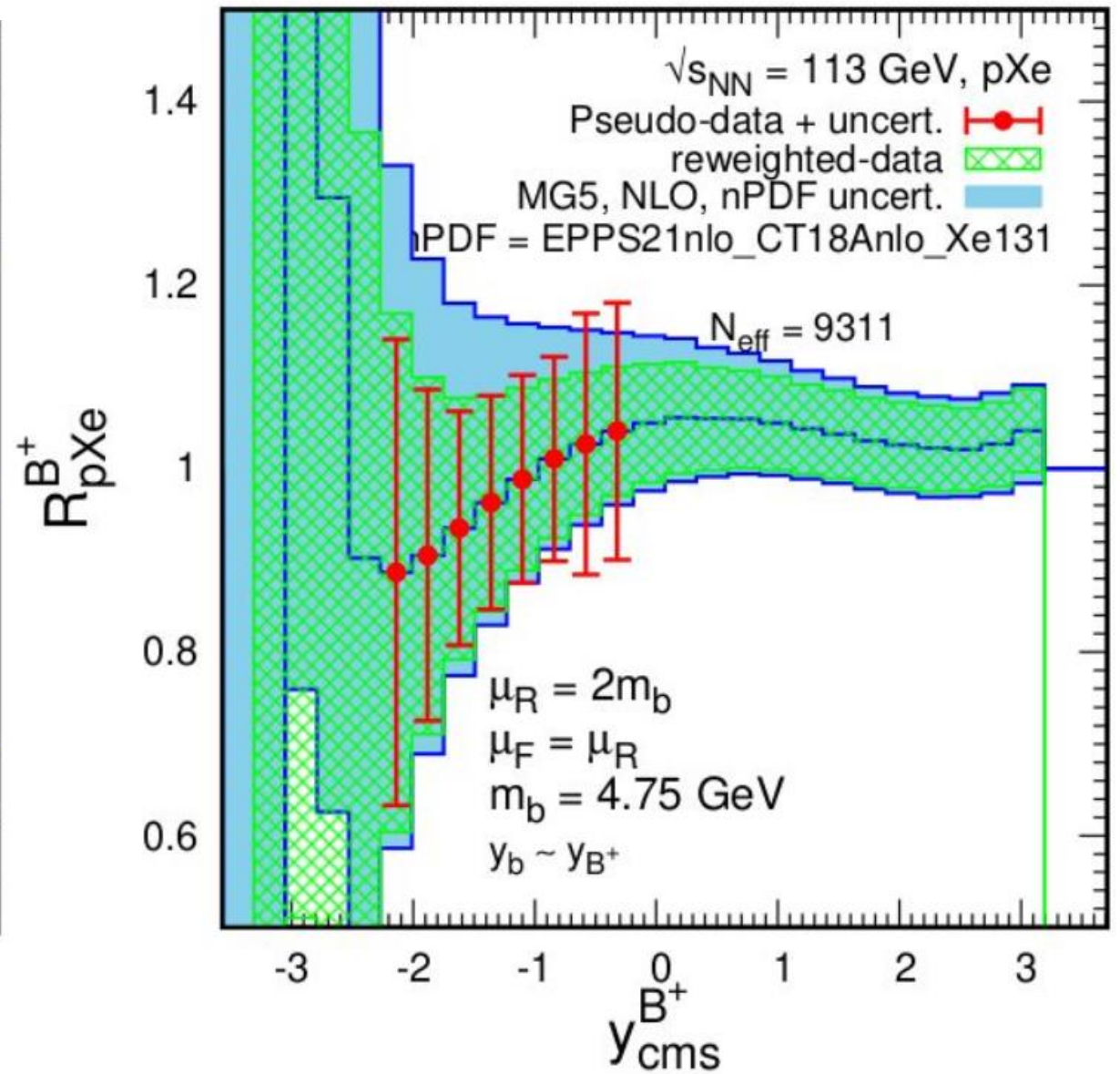
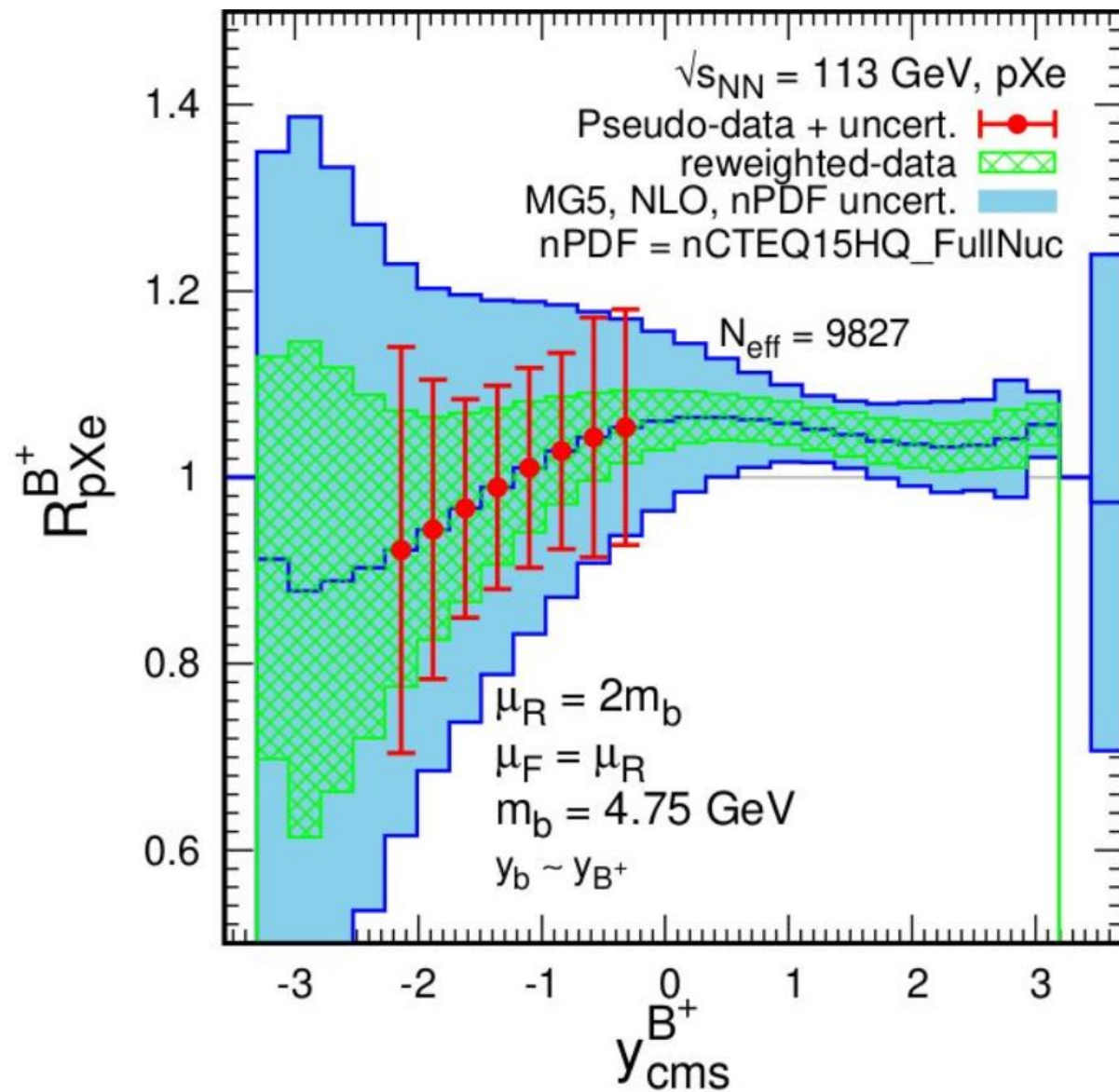
GEC Selection Weighting

0.7% 1% 0.2%

$$\frac{\delta\sigma_{\text{syst}}}{\sigma_{\text{syst}}} = 6.1\%$$

**Uncertainties are  
considered as uncorrelated  
over y!**

# $R_{pA}^{B^+}$ for Xe nPDF: reweighing



# Conclusions

- **Asymmetric collisions in MadGraph5 have been implemented and successfully validated**
- **The code is available at GitHub right now!**
- **Also, it will be a part of NLOAccess and official MadGraph5**
  
- **Reweighting is working with chosen D0-meson pseudo-data with SMOG2:**
  - **Pseudo-data uncertainties dominated by the systematic ones**
  - **If data-point uncertainties lower than nPDF uncertainties (Xe target), data can constrain gluon nPDF in and outside the probed rapidity interval**
  
- **Pseudo-data has the strong impact on the nPDF, for**
  - Xe target ( $x_2 \in [0.032, 0.33]$ )**
  - and**
  - Pb-beam ( $x_1 \in [0.0053, 0.056]$ )**
  
- **For lighter target it's not clear how reliable are uncertainties for nPDFs.**
  
- **Very similar situation with B+-meson pseudo-data**
- **Except that statistic uncertainties are much larger, due to the low luminosities of the Hydrogen and a target**
  
- **Some pseudo-data systematic uncertainties can be correlated over  $y$  assuming some experimental correlation could improve the reweighting picture.**



# Backup

# Quark nPDFs

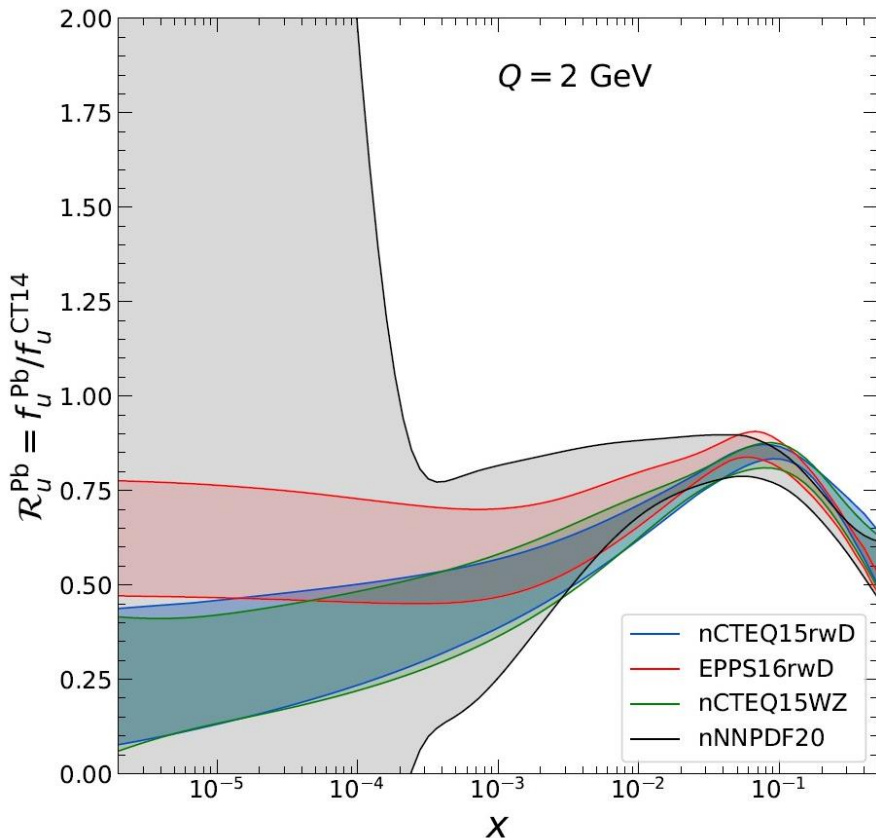
Since the early 1980s, from the ratio of structure functions  $F_2$ , we know that the **nuclei** are **not** a simple collection of **free nucleons**.

In other words, **nPDFs deviate** from a simple **sum of nucleon PDFs**. To **study** such deviations, it is customary to rely on **NMFs**, like:

$$R[F_2^{\ell A}] = \frac{F_2^{\ell A}}{(ZF_2^{\ell p} + (A - Z)F_2^{\ell n})}$$

$$R_i^A(x, \mu_F) = \frac{Zf_i^{p/A} + (A - Z)f_i^{n/A}}{Zf_i^p + (A - Z)f_i^n}$$

arXiv:1712.07024v2 [hep-ph]

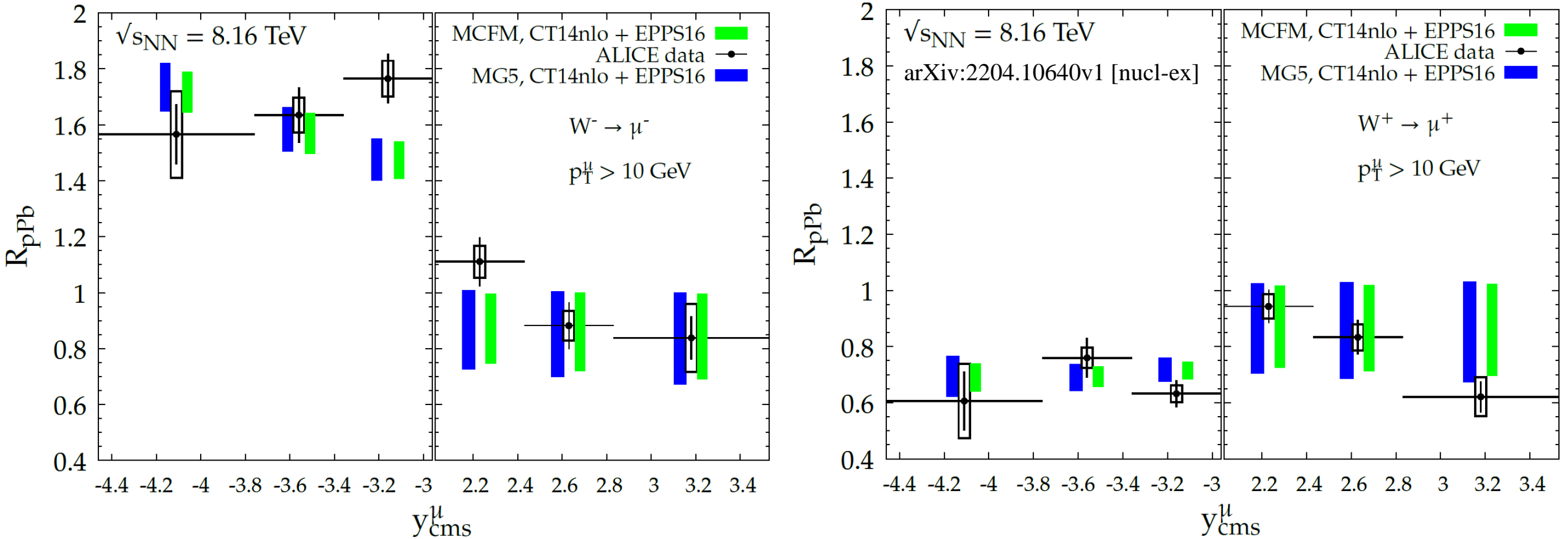


One expects:

- $R_q^A > 1$  for  $x \gtrsim 0.8$  (Fermi-motion region),
- $R_q^A < 1$  for  $0.25 \lesssim x \lesssim 0.8$  (EMC region),
- $R_q^A > 1$  for  $0.1 \lesssim x \lesssim 0.25$  (antishadowing region)
- $R_q^A < 1$  for  $x \lesssim 0.1$  (shadowing region)
- $R_q^A \sim 1$ : absence of nuclear effects

# Validations of MG5 in asymmetric collisions

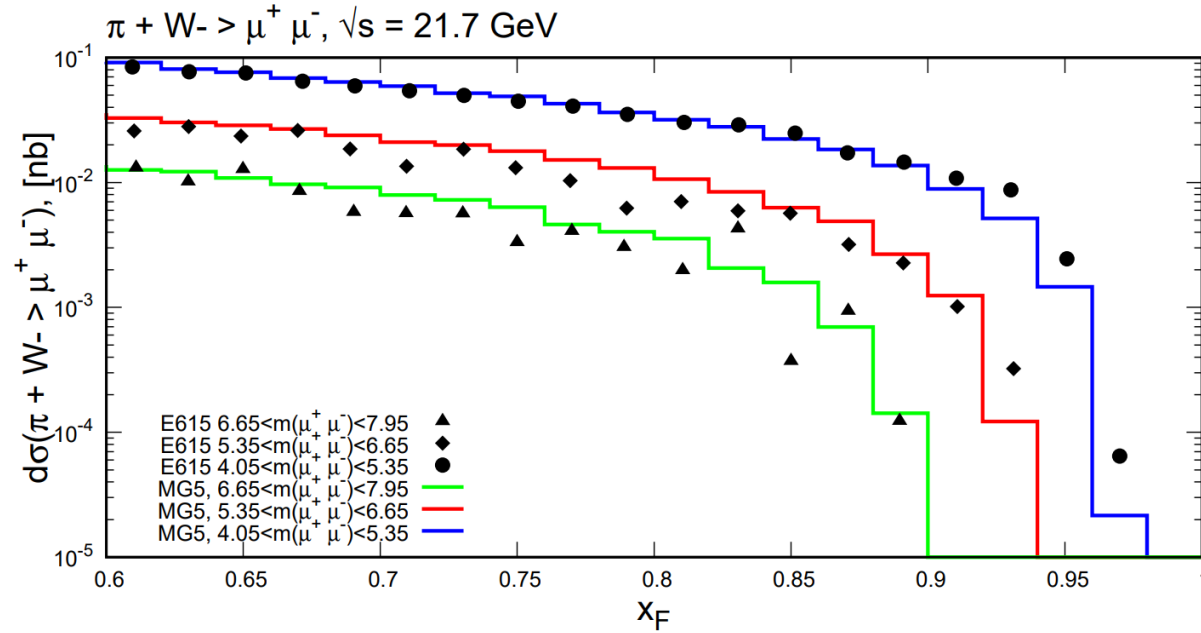
## Validation vs MCFM for CT14 + EPPS16 for W production at NLO



- **Good agreement** between **MG5** and **MCFM**-based computations for **EPPS16**
- **Good agreement** between **MG5** and experimental data
- Slight difference in the uncertainty since MCFM-based computation done with **symmetric uncertainties**

# Example: *Drell-Yan* production in $\pi W$ collision

$4.05 \text{ GeV}/c < m_{\mu\mu} < 8.55 \text{ GeV}/c$

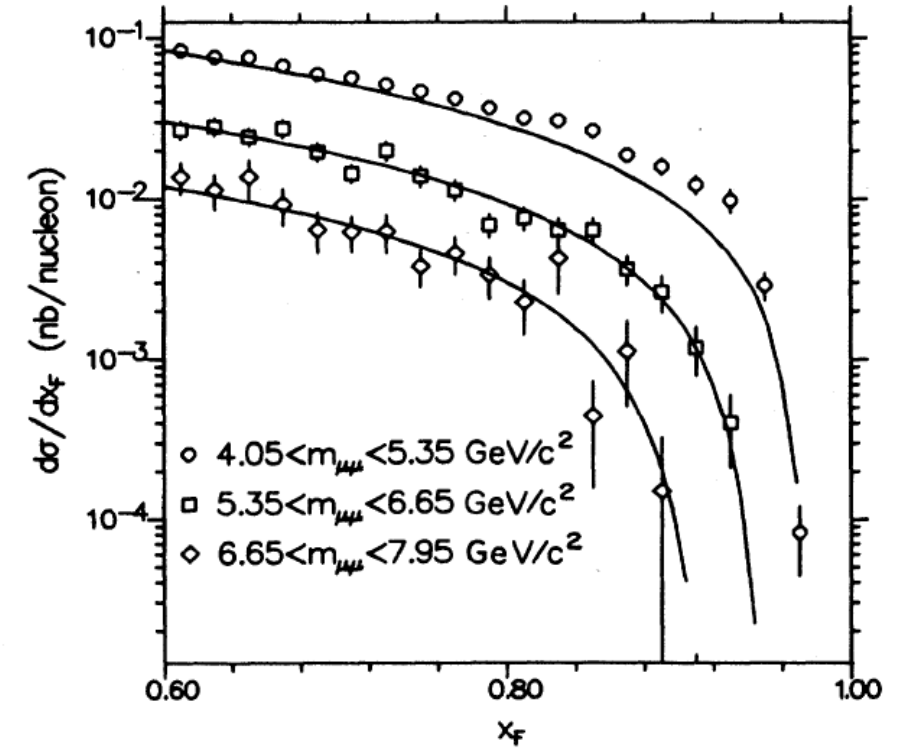


$$\tau = \frac{m_{\mu\mu}^2}{s}$$

$$x_F = x_{\text{pion}} - x_{\text{nucleus}}$$

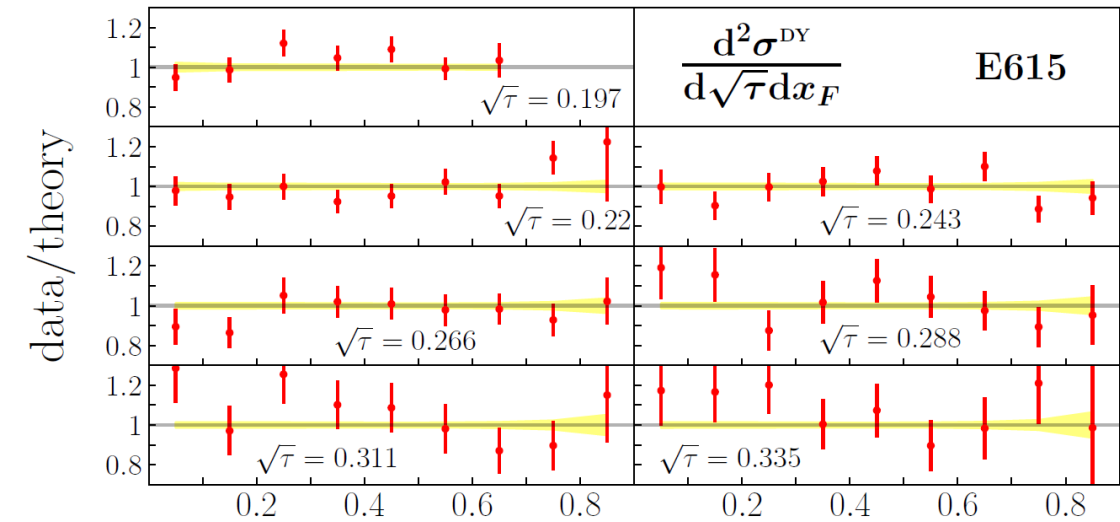
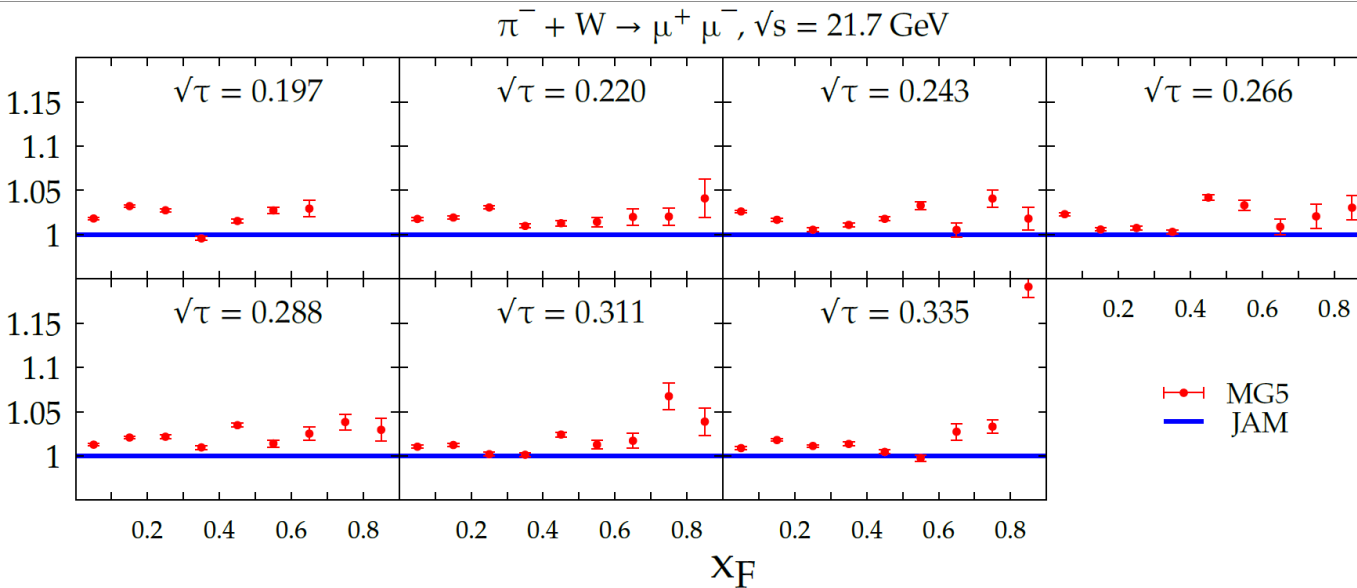
$$x_{\text{pion}} = \sqrt{\tau} e^Y$$

$$x_{\text{nucleus}} = \sqrt{\tau} e^{-Y}$$



*Phys.Rev.D* 39 (1989) 92-122

# Example: *Drell-Yang* production in $\pi W$ collision



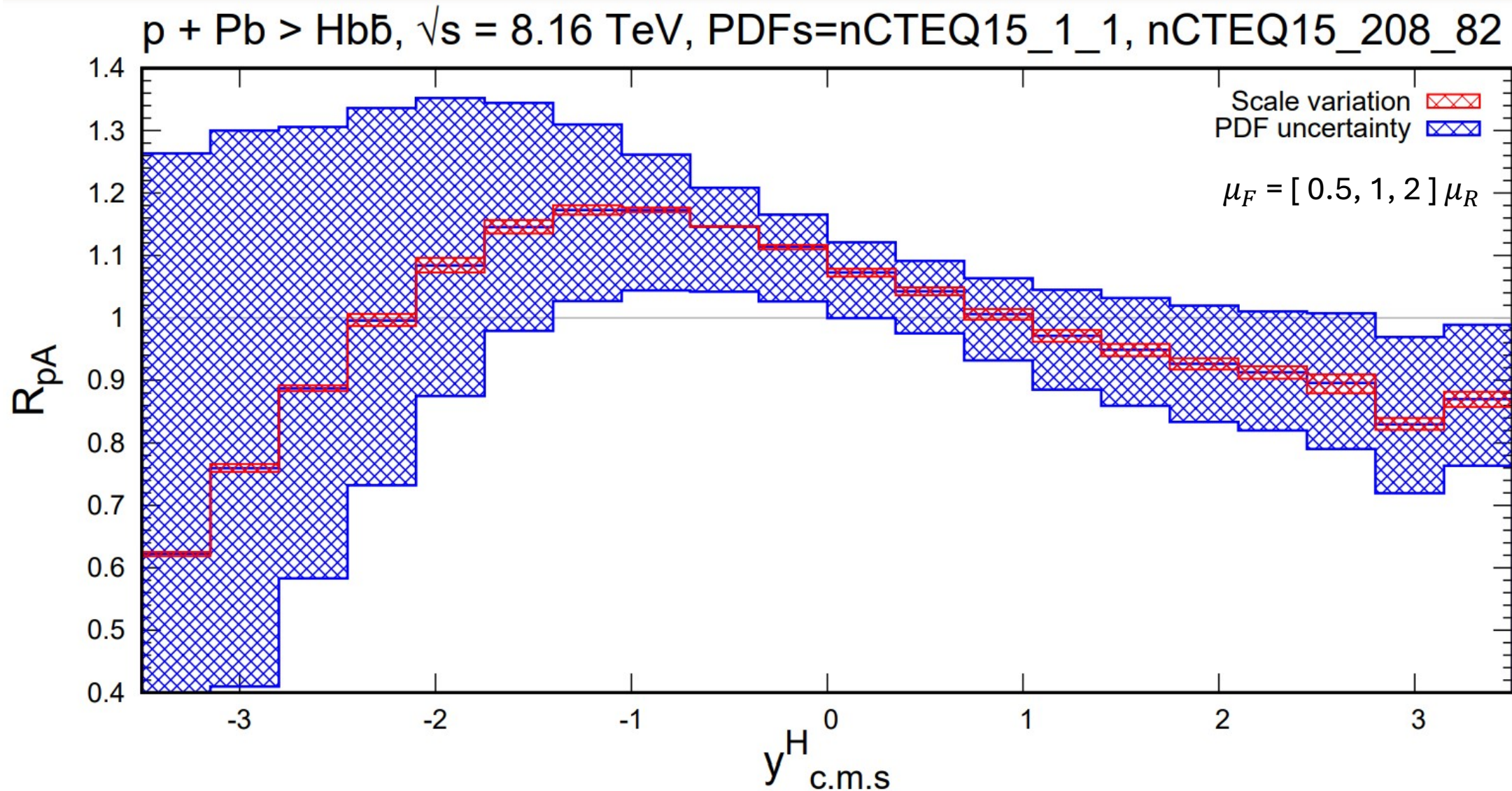
arXiv:2103.02159v2 [hep-ph] 30 Nov 2021

*Phys.Rev.D* 39 (1989) 92-122

- For all  $\sqrt{\tau}$  regions and bins, differences do not exceed 5-10% percent range
- Results **match** well those produced by the **JAM** collaboration
- Small **differences**, could **arise** from **instabilities** that relate to **Monte-Carlo algorithms** and very **narrow regions** of invariant masses of muon pairs

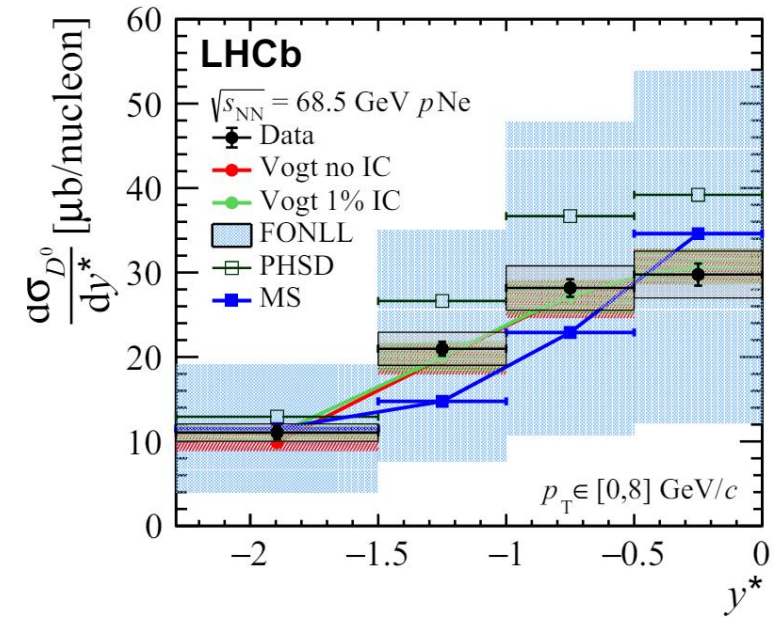
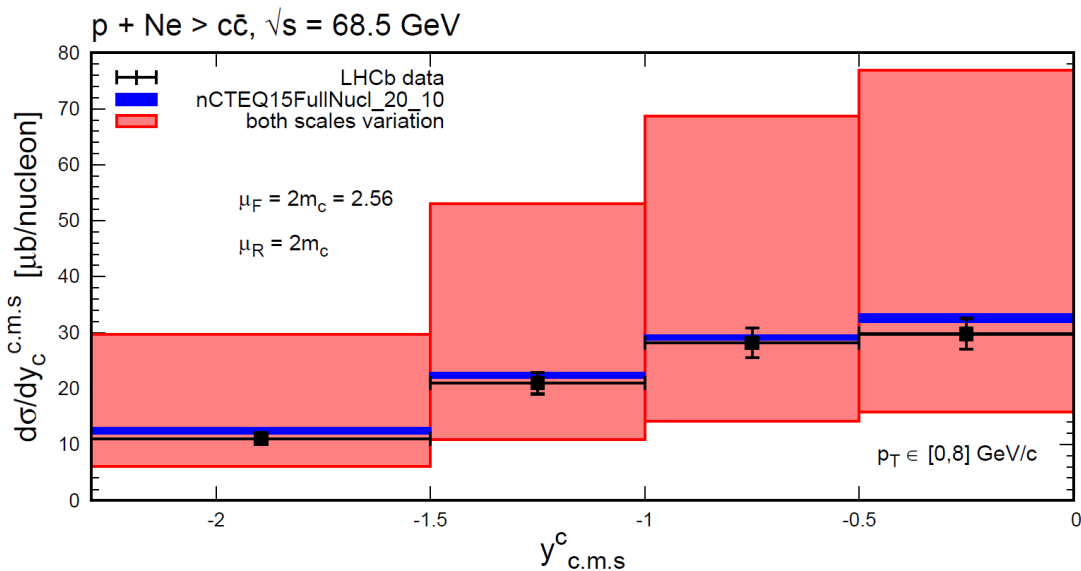
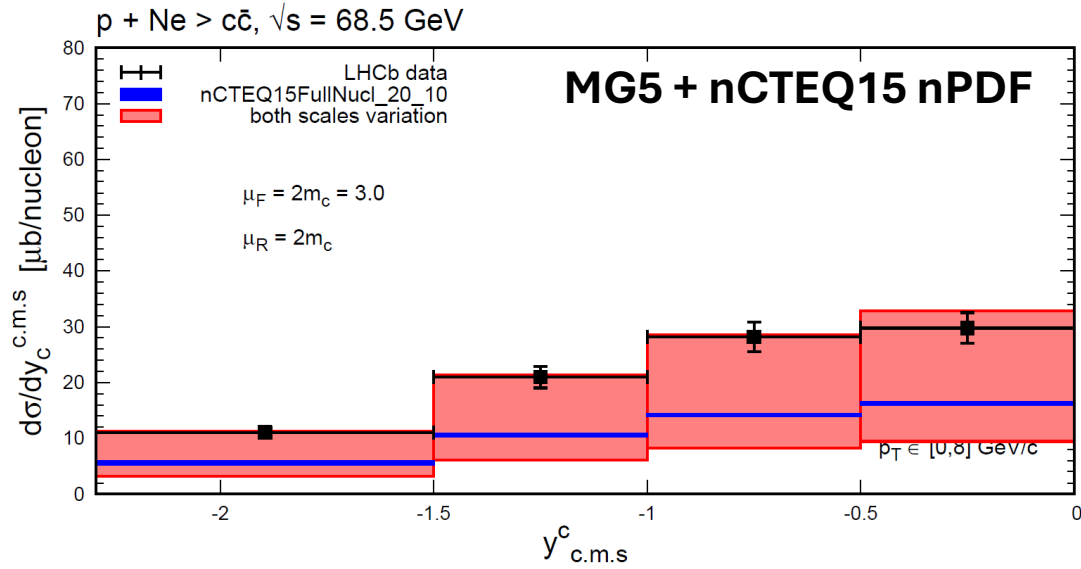
$$\tau = \frac{m_{ll}^2}{s}$$

# “Un”real prediction: Higgs+ $b\bar{b}$ at NLO



- Rapidity dependence for other particles can be obtained by changing a single line in the analysis file

# $D^0$ production for the proton-Ne collisions, for two different masses for $m_c$



Assumption:  $y_c \sim y_{D^0}$

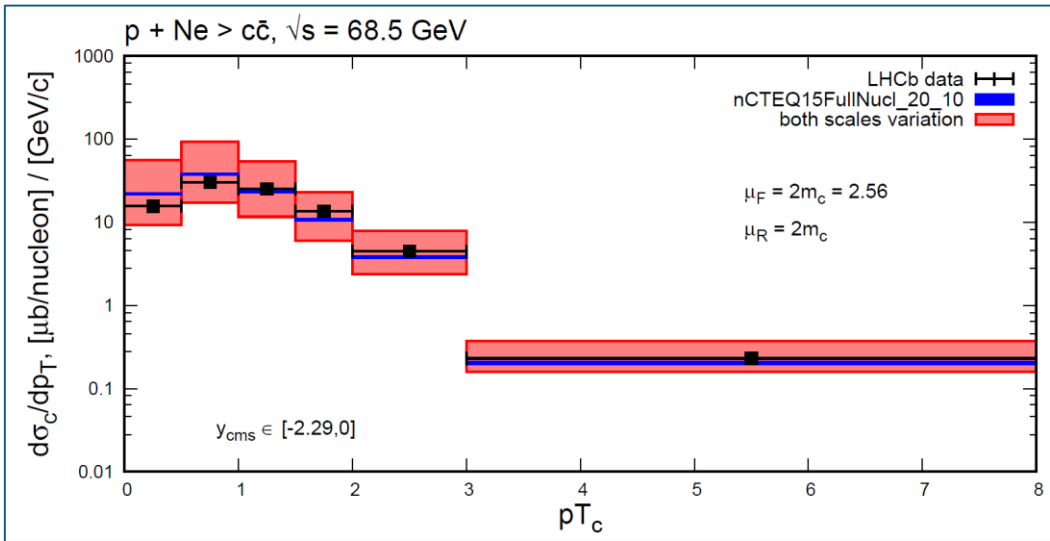
On the plots we have a cross-section of c-quark production,

multiplied by the factor:  $2 * f(c \rightarrow D^0)$

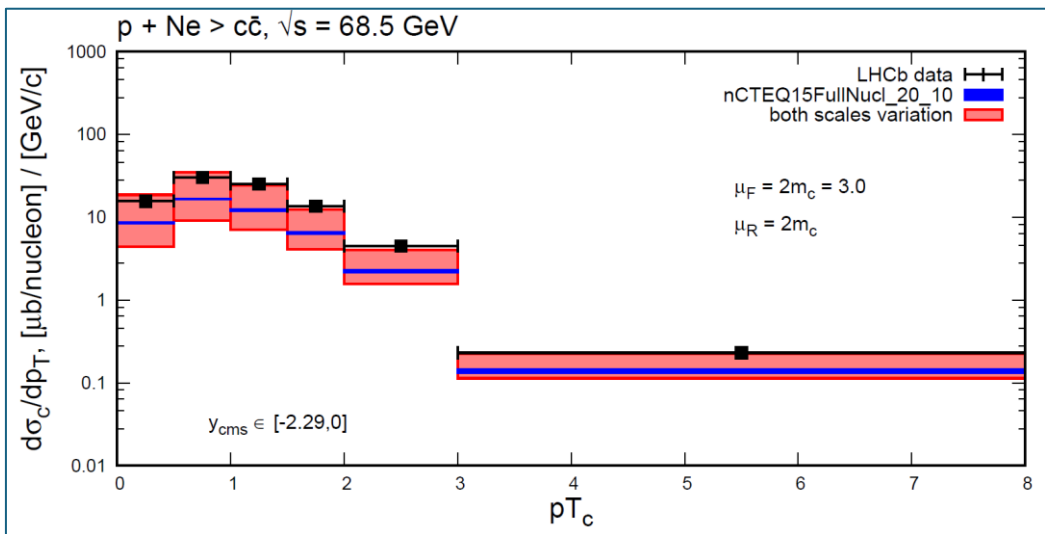
$$f(c \rightarrow D^0) = 0.542 \pm 0.024$$

**Charm mass  
significant for the scale variation**

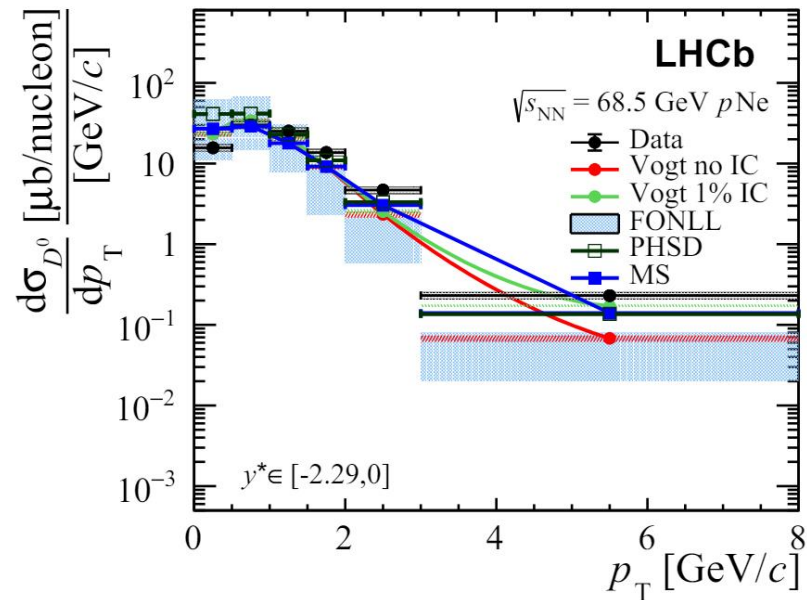
# $D^0$ production for the proton-Ne collisions, for two different masses of $m_c$



**MG5 + nCTEQ15 nPDF**



LHCb data: arXiv:2211.11633v3 [hep-ex] 20 Feb 2024



On the plots we have a cross-section of c-quark  
production, multiplied by the factor:

$$2 * f(c \rightarrow D^0)$$

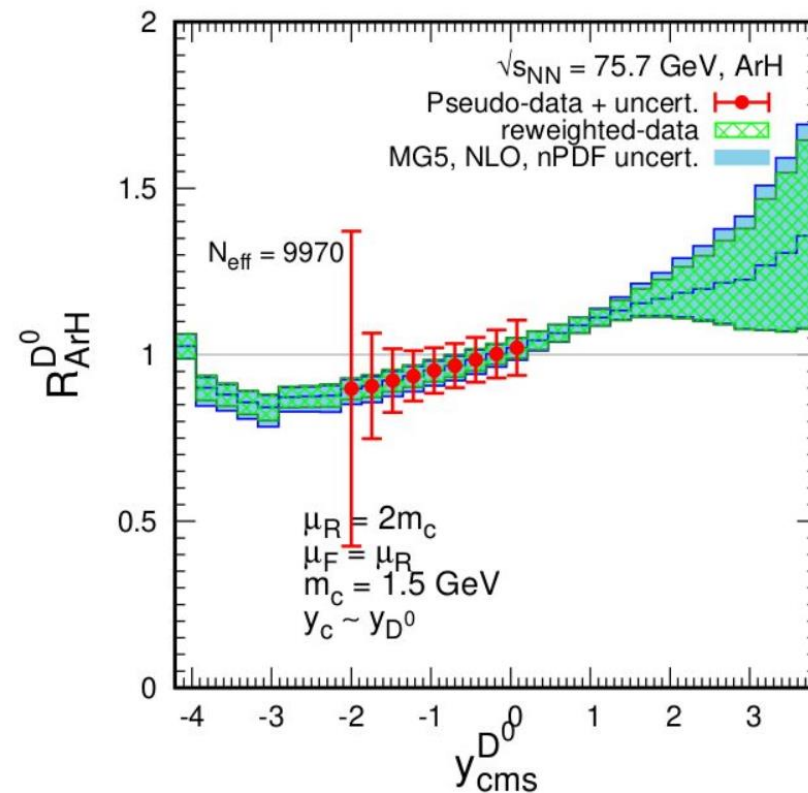
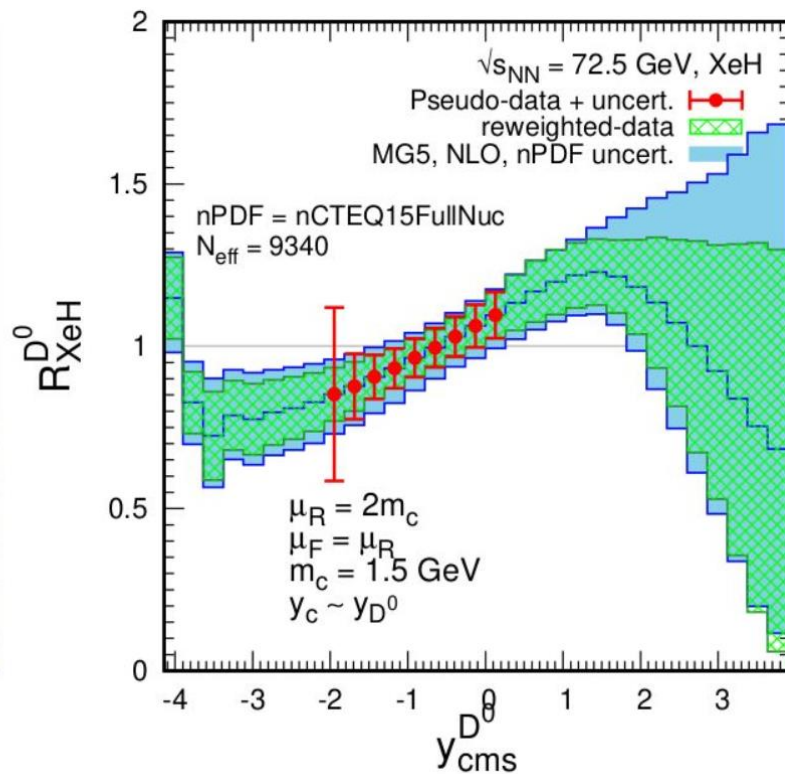
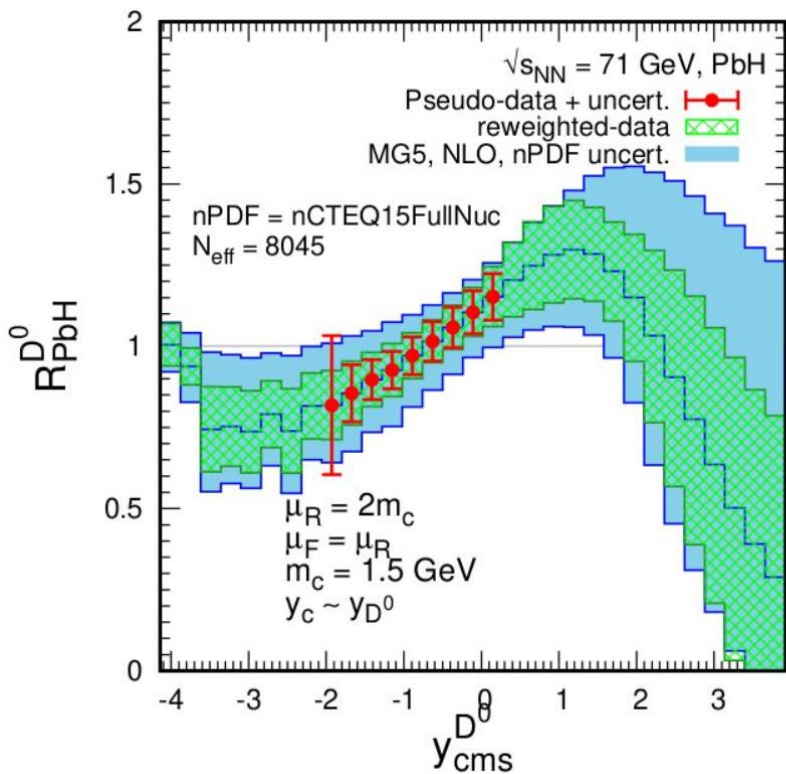
$$f(c \rightarrow D^0) = 0.542 \pm 0.024$$

Assumption that  $pT_c \sim pT_{D^0}$   
not correct for **large**  $pT$ :

one can try to use **Pythia** to compute decay  
kinematics



# Reweighting of the nPDFs: Heavy-Ion beams



## Reweighting process

**Def. of PDF replicas  $f_k$ :**

$$f_k = f_0 + \sum_{i=1}^N \frac{f_i^{(+)} - f_i^{(-)}}{2} R_{ki}$$

$f_0$  – best fit (central) PDF

$f_i^{(+)}$ ,  $f_i^{(-)}$  – are the plus and minus error PDFs corresponding to the eigenvector direction  $i$

**Def. of weight:**

$$w_k = \frac{e^{-\frac{1}{2}\chi_k^2/T}}{\frac{1}{N_{\text{rep}}} \sum_i^{N_{\text{rep}}} e^{-\frac{1}{2}\chi_i^2/T}}$$

$T$  – is the tolerance criterion used when defining Hessian error PDFs  
 $\chi_k^2 - \chi^2$  for a given replica  $k$ .

**Def. of  $\chi^2$  :**

$$\chi_k^2 = \sum_j^{N_{\text{data}}} \frac{(D_j - T_j^k)^2}{\sigma_j^2}$$

$j$  – runs over all data points in the data set(s)

$N_{\text{data}}$  – total number of data points,

$D_j$  – the experimental measurement at point  $j$

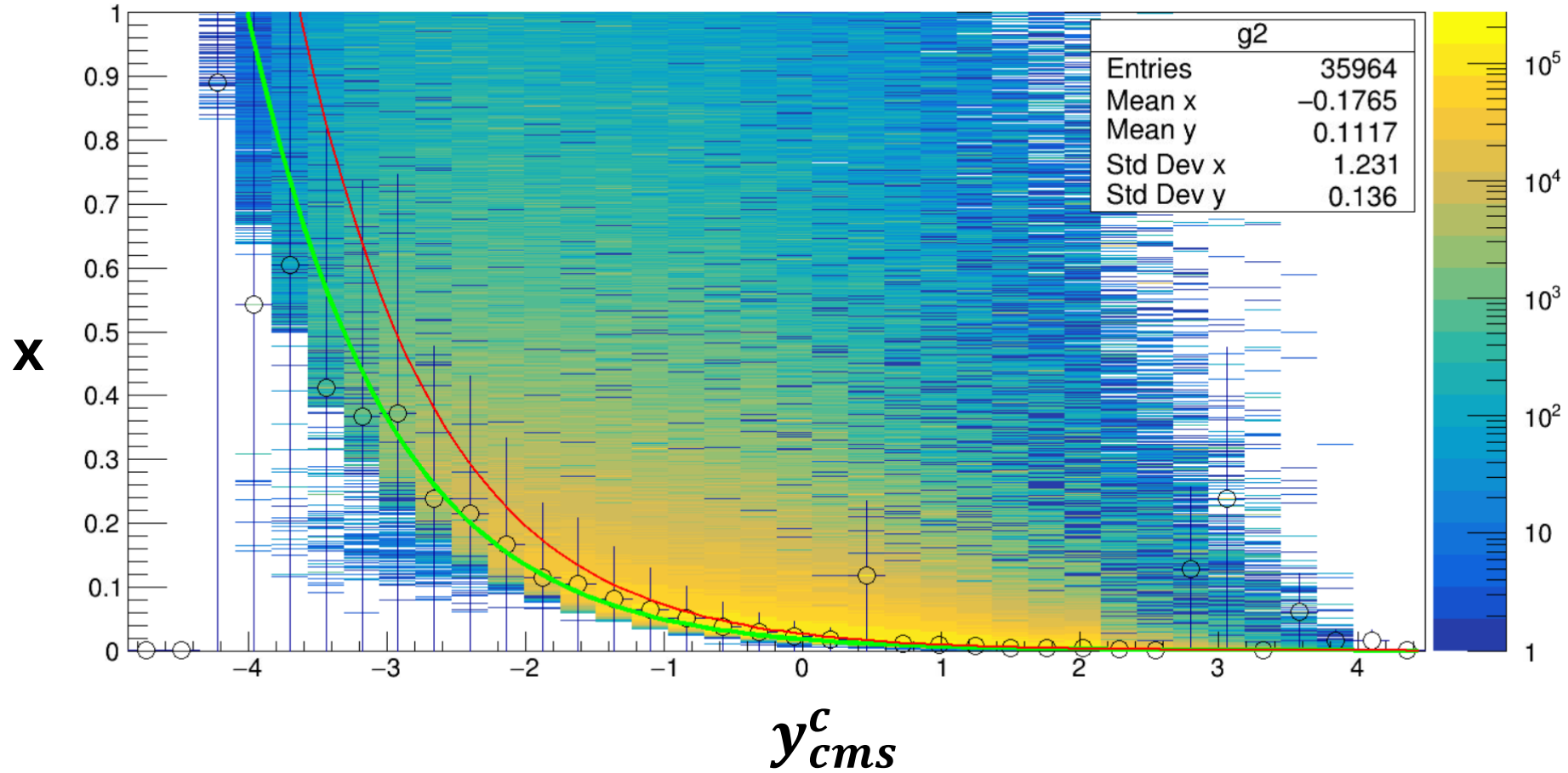
$\sigma_j$  – experimental uncertainty

$T_k^j$  – theoretical prediction calculated with PDFs given by replica  $k$ .

$$\langle \mathcal{O} \rangle_{\text{new}} = \frac{1}{N_{\text{rep}}} \sum_{k=1}^{N_{\text{rep}}} w_k \mathcal{O}(f_k),$$

$$\delta \langle \mathcal{O} \rangle_{\text{new}} = \sqrt{\frac{1}{N_{\text{rep}}} \sum_{k=1}^{N_{\text{rep}}} w_k (\mathcal{O}(f_k) - \langle \mathcal{O} \rangle_{\text{new}})^2}$$

# $x_2$ vs $y$ for $pXe$



Red:  $m_c = 1.5 \text{ GeV}$

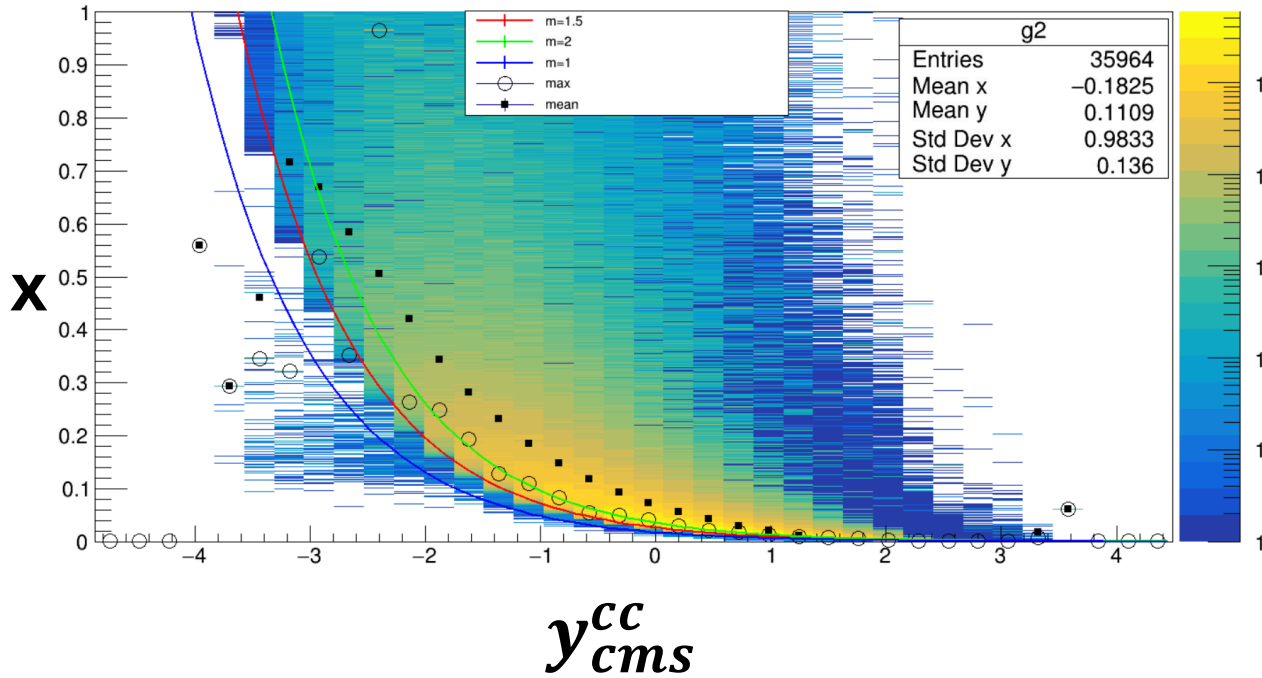
Green:  $m_c = 1 \text{ GeV}$  (fit)

Open-points: **maximum** of the **distribution**  
per bin of  $y$

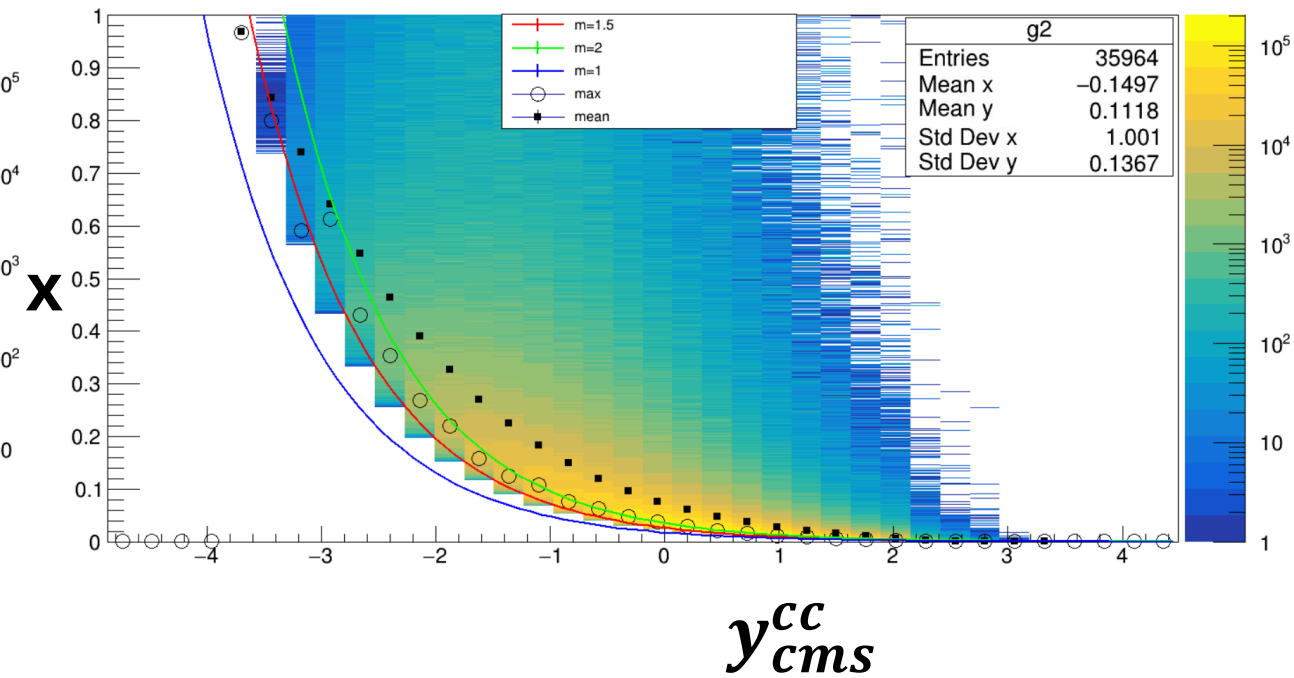
$$x_2 = \frac{2m_c}{113} e^{-y}$$

# $x_2$ vs $y$ for $pXe$

## NLO



## LO



Red:  $m_c = 1.5 \text{ GeV}$

Green:  $m_c = 1 \text{ GeV}$  (fit)

Open-points: **maximum** of the **distribution**  
per bin of  $y$

$$x_2 = \frac{2m_c}{113} e^{-y}$$