DPS in quarkonium hadroproduction and photoproduction

Matteo Rinaldi
INFN sezione di Perugia

Outline

- Introduction to double parton scattering (DPS) and hadronic Physics
- Data and interpretation
- DPS at the EIC?
- Nuclear DPS at the EIC?

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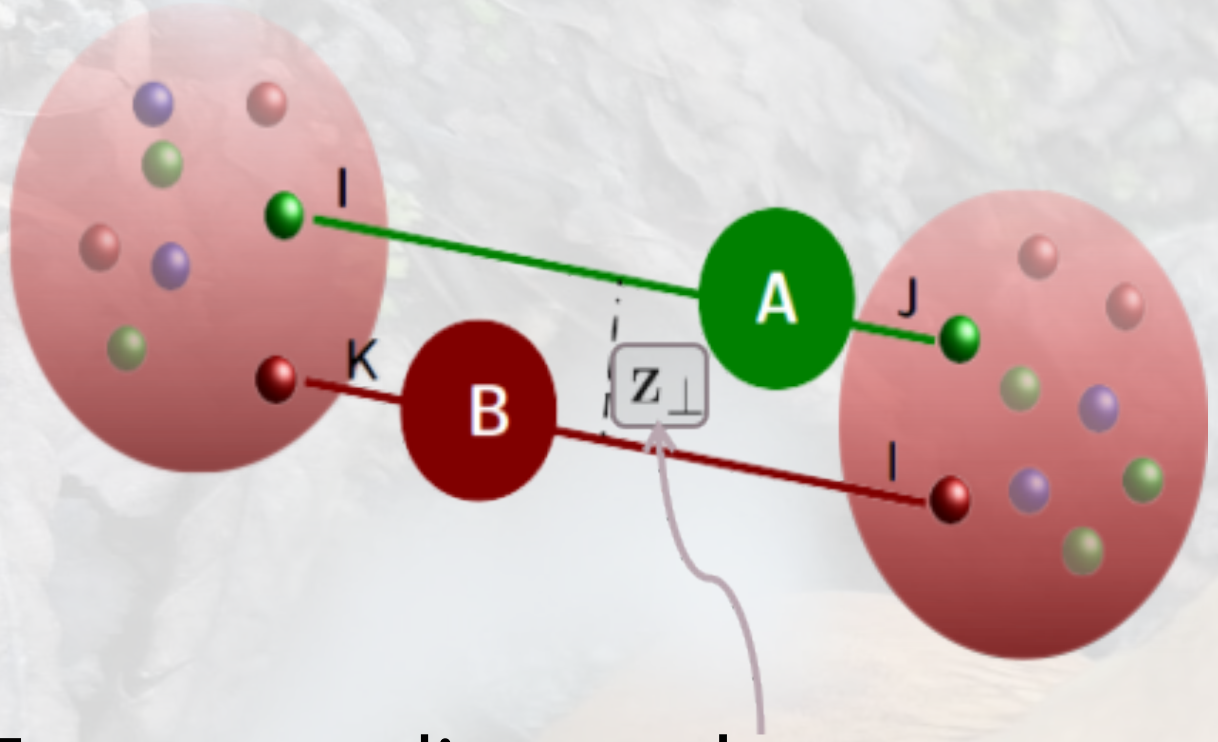
Double Parton Scattering

Multiparton interaction (MPI) can contribute to the, pp, pA and AA cross section @ the LHC:



Transverse distance between two partons

$$d\sigma \propto \int d^2 z_\perp \, F_{ij}(x_1, x_2, z_\perp, \mu_A, \mu_B) F_{kl}(x_3, x_4, z_\perp, \mu_A, \mu_B)$$

Double Parton Distribution (DPD)

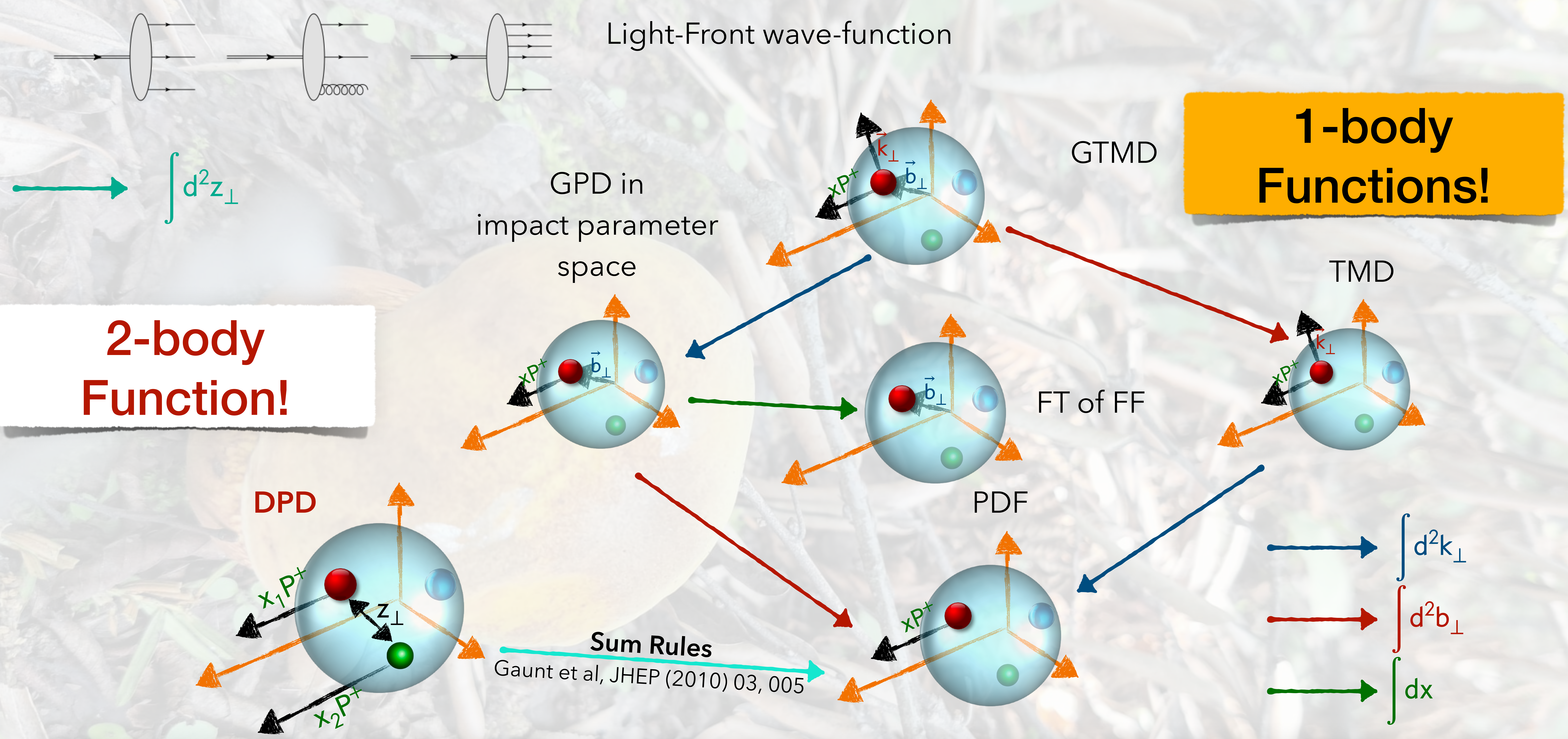
N. Paver and D. Treleani, Nuovo Cimento 70A, 215 (1982)
Mekhfi, PRD 32 (1985) 2371
M. Diehl et al, JHEP 03 (2012) 089

$$F_{ij}^{\lambda_1,\lambda_2}(x_1, x_2, \vec{k}_\perp) = (-8\pi P^+)\frac{1}{2}\sum_\lambda \int d\vec{z}_\perp \, e^{i\vec{z}_\perp\cdot\vec{k}_\perp}$$
$$\times \int \left[\prod_l^3 \frac{dz_l^-}{4\pi}\right] e^{ix_1P^+z_1^-/2} e^{ix_2P^+z_2^-/2} e^{-ix_1P^+z_3^-/2}$$
$$\times \langle \lambda, \vec{P} = \vec{0} | \hat{\mathcal{O}}_i^1 \left(z_1^-\frac{\bar{n}}{2}, z_3^-\frac{\bar{n}}{2} + \vec{z}_\perp\right) \hat{\mathcal{O}}_j^2 \left(z_2^-\frac{\bar{n}}{2} + \vec{z}_\perp, 0\right) | \vec{P} = \vec{0}, \lambda\rangle$$

$$\hat{\mathcal{O}}_i^k(z, z') = \bar{q}_i(z)\hat{O}(\lambda_k)q_i(z')$$

$$\hat{O}(\lambda_k) = \frac{\not{n}}{2}\frac{1 + \lambda_k\gamma_5}{2} .$$

Multidimensional picture of hadrons



How to build up a DPD

$F_{ik}(x_1, x_2, \vec{z}_\perp)$ is unknown. For phenomenology @LHC kinematics (small x and many partons produced)

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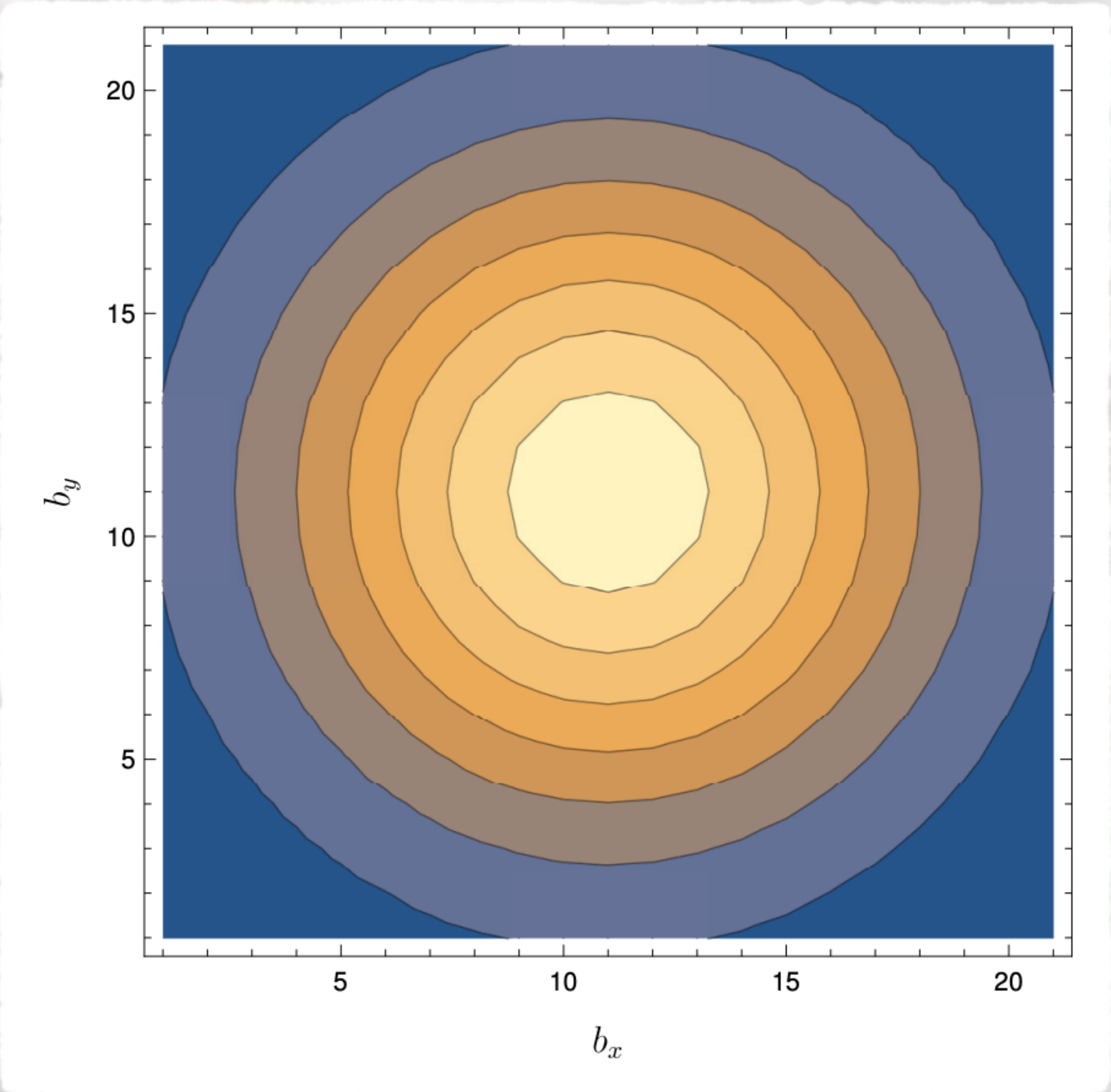
$$F_{ik}(x_1, x_2, \vec{z}_\perp) \sim g(x_1, x_2)\tilde{T}(\vec{z}_\perp)$$

Models can help to grasp general features

M.R., S. Scopetta et al, PRD 87 (2013) 114021
M.R., S, Scopetta et al, JHEP 12 (2014) 028
A. V. Manohar et al, PRD 87 (2013) 3, 034009

$$\langle b_\perp^2 \rangle_{x_1,x_2}^{ij} = \frac{\int d^2 b_\perp b_\perp^2 \tilde{F}_{ij}(x_1, x_2, b_\perp, Q^2)}{\int d^2 b_\perp \tilde{F}_{ij}(x_1, x_2, b_\perp, Q^2)}$$



M.R. and F. A. Ceccopieri, JHEP 09 (2019) 097

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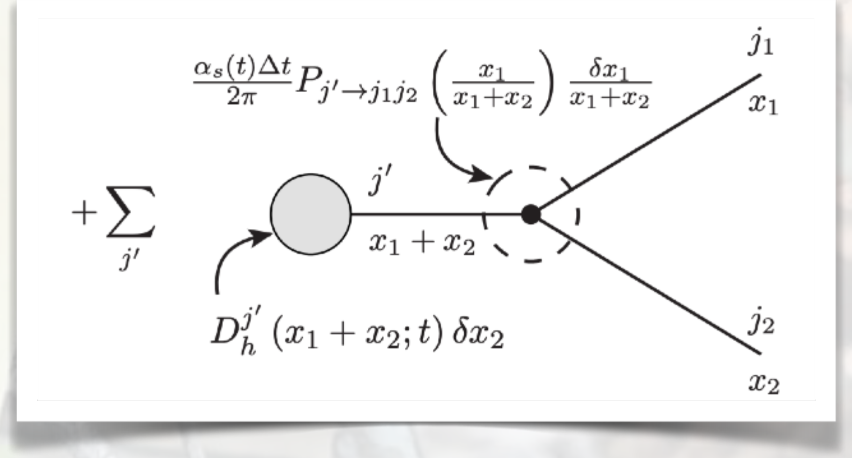
uncorrelated scenario: $$F_{ik}(x_1, x_2, \vec{z}_\perp) \sim g(x_1, x_2)\tilde{T}(\vec{z}_\perp)$$

double PDF

Sum Rules

pQCD evolution

PDF(x_1)*PDF(x_2) uncorrelated scenario



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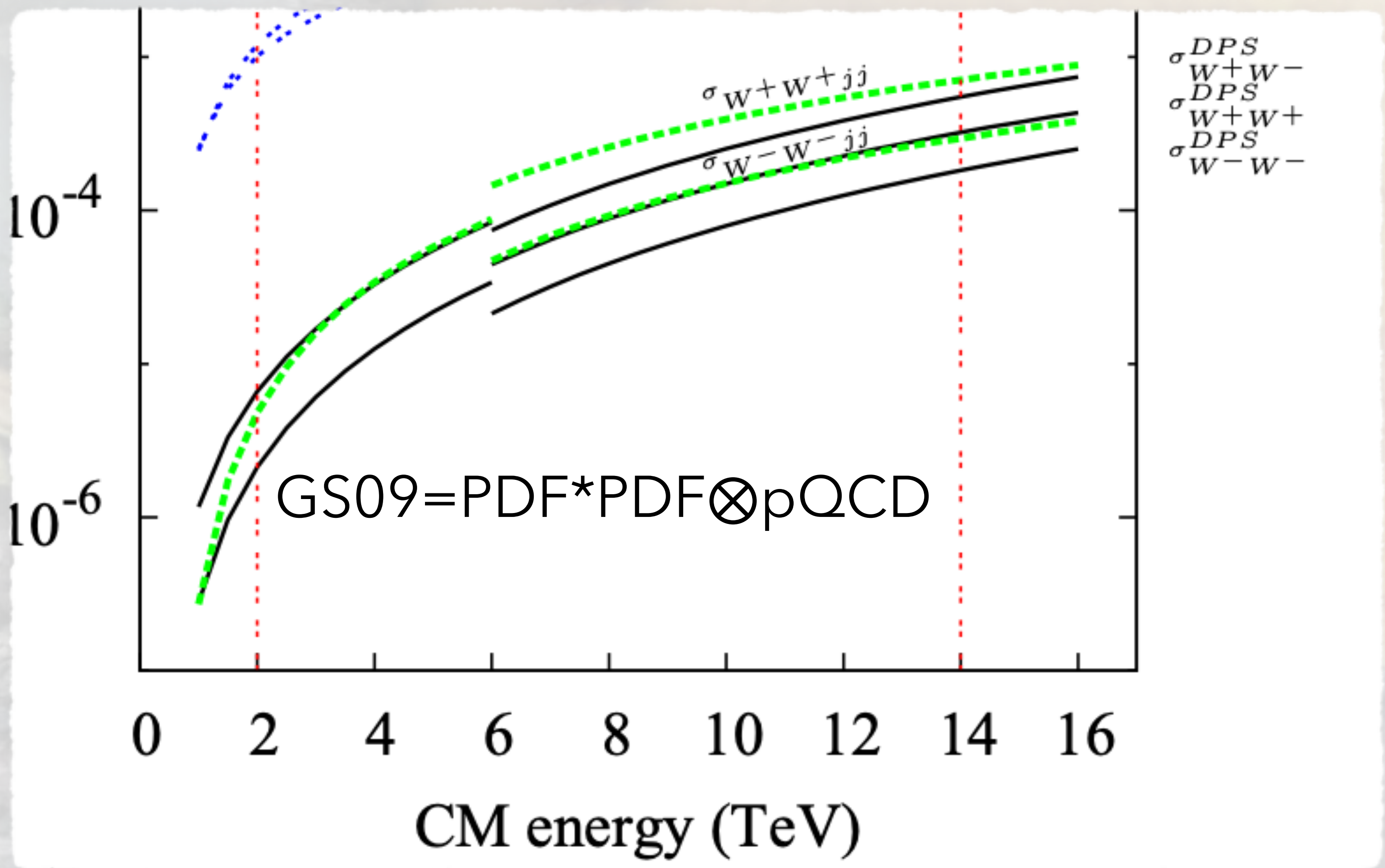
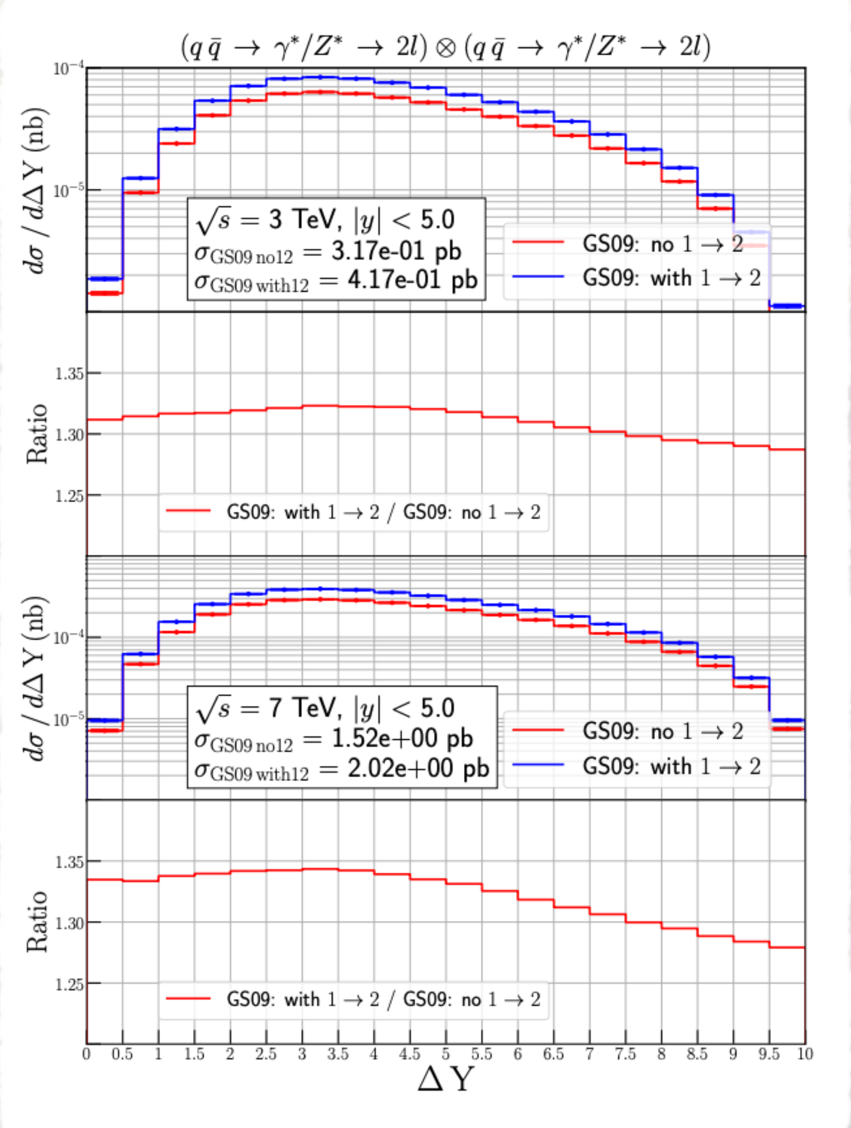
J. R. Gaunt et al, EPJC 69 (2010) 54-65





O. Fedkevych, J. R. Gaunt, JHEP 02 (2023) 090

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$F_{ik}(x_1, x_2, \vec{z}_\perp)$ is unknown. For phenomenology @LHC kinematics (small x and many partons produced)

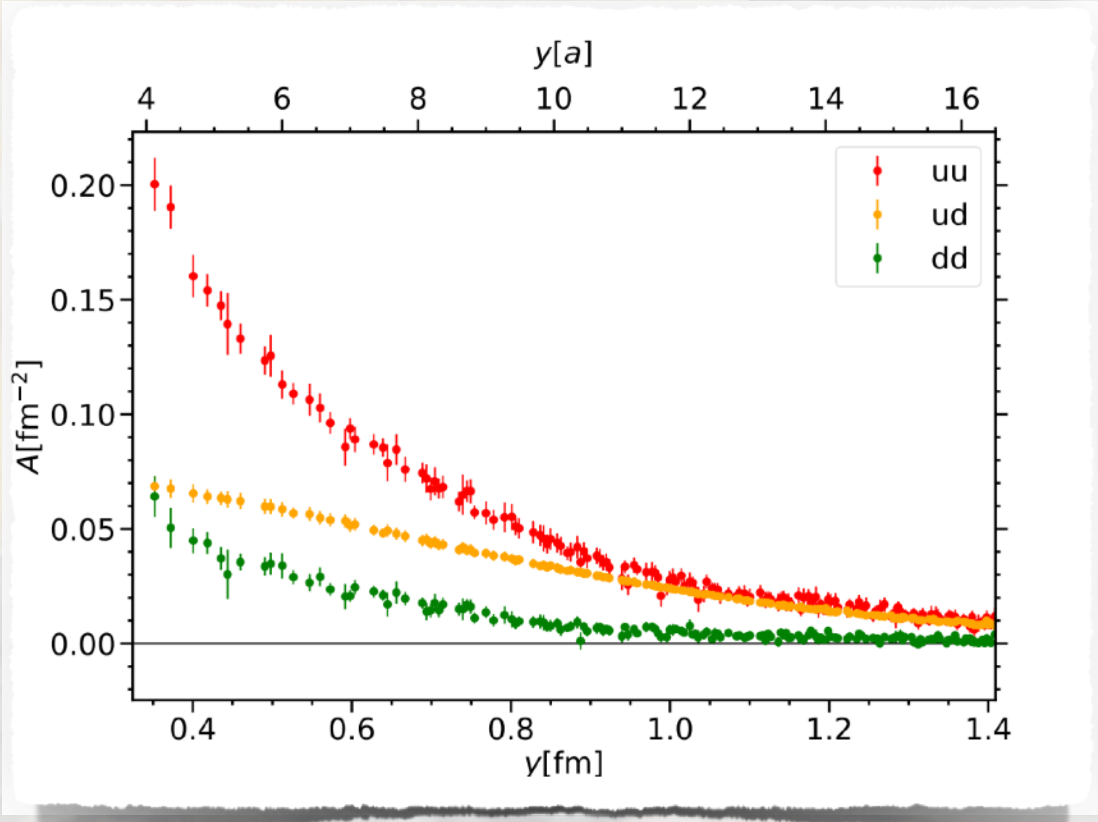
double PDF

uncorrelated scenario: $F_{ik}(x_1, x_2, \vec{z}_\perp) \sim g(x_1, x_2)\tilde{T}(\vec{z}_\perp)$ → Probability distribution of two partons at given distance

Unknown Non perturbative object

Some information from lattice

G. S. Bali et al, JHEP 09 (2021) 121



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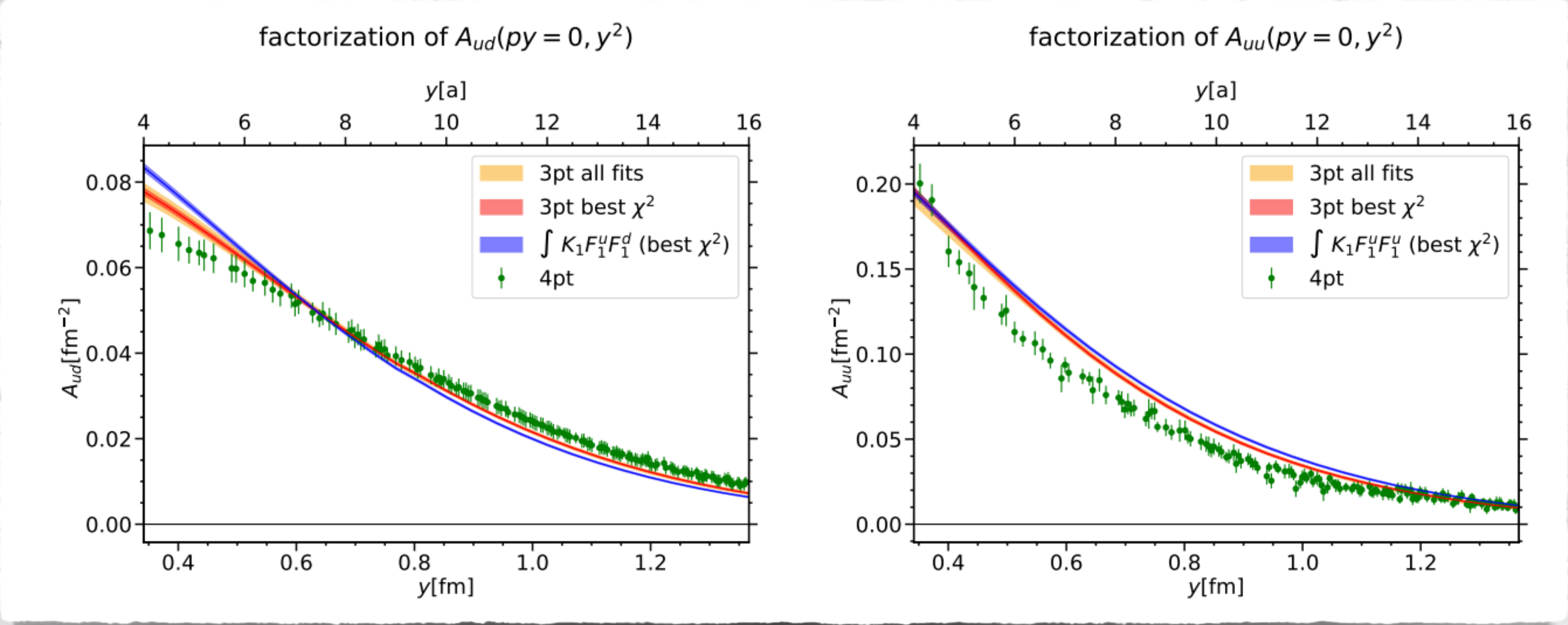

Violation of the usual ansatz:

$$\tilde{T}(\vec{z}_\perp) \sim \int d^2b_1\, f(b_1) f(b_1 - \vec{z}_\perp)$$

Transverse distribution

DPD = GPD x GPD
NOT WELL REPRODUCED!

G. S. Bali et al, JHEP 09, 106 (2021)

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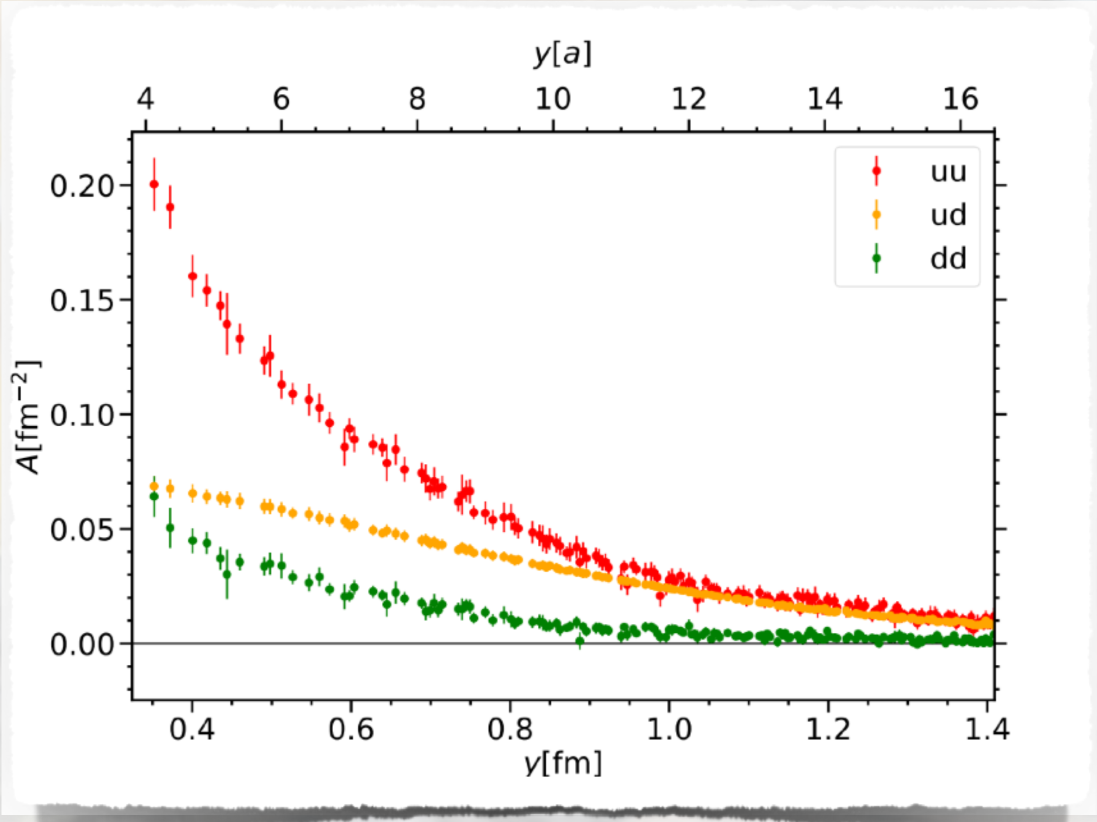
Unknown Non perturbative object

Some information from lattice

Some constraints from data

Some Constraints from general properties

G. S. Bali et al, JHEP 09 (2021) 121



$$\sigma_{\text{eff}}^{-1} = \int d^2 z_\perp \, \tilde{T}(z_\perp)^2$$

Some Data and Effective Cross Section

$$\sigma^{pp}_{eff} = \frac{m}{2}\frac{\sigma^{pp}_{A}\sigma^{pp}_{B}}{\sigma^{pp}_{DPS}}$$

POCKET FORMULA

Differential X-section single parton scattering for the process: $pp \longrightarrow A(B) + X$

Differential X-section double parton scattering for the process: $pp \longrightarrow A + B + X$

Some Data and Effective Cross Section

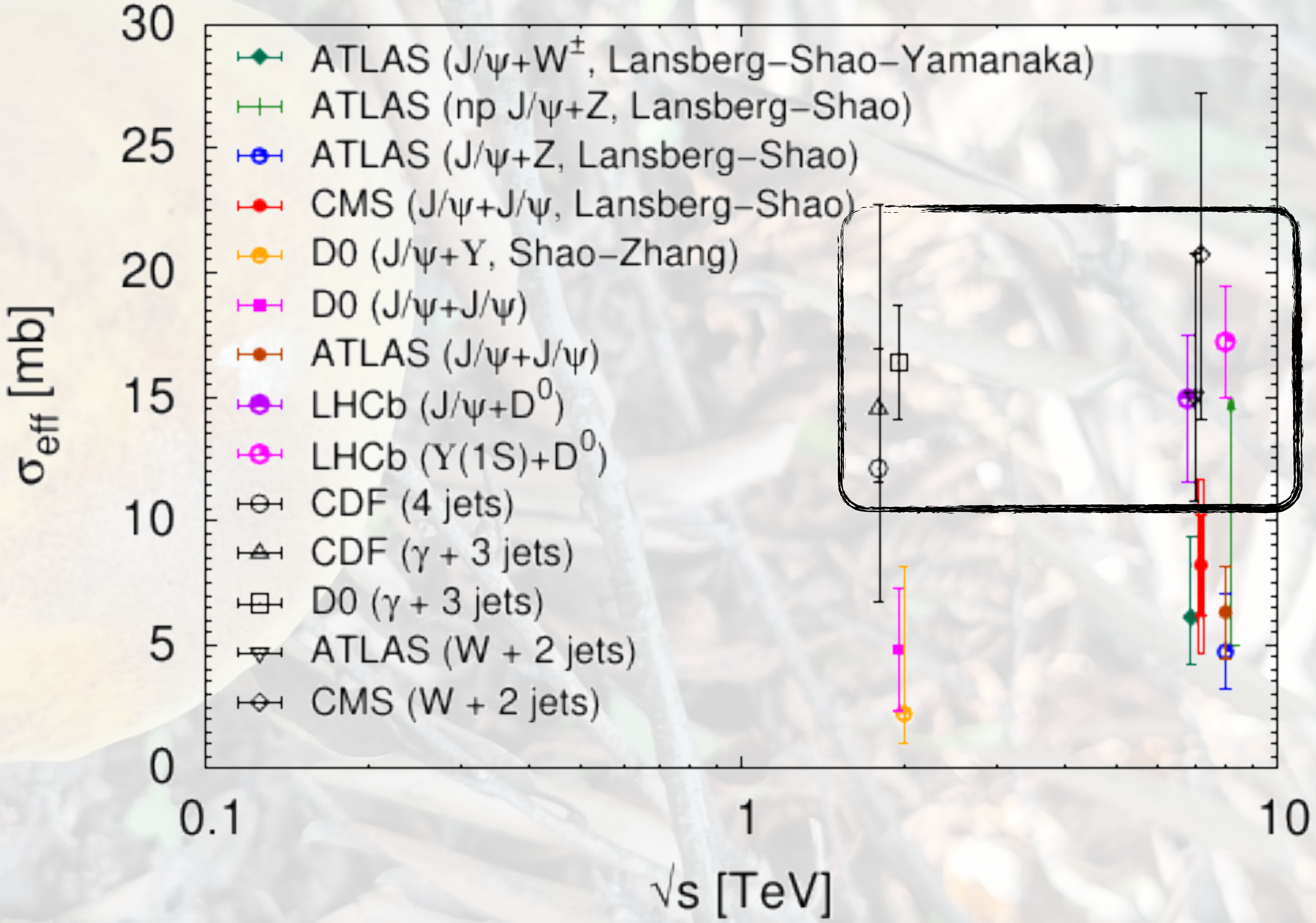
$$\sigma_{\text{eff}}^{\text{pp}} = \frac{m}{2} \frac{\sigma_A^{\text{pp}} \sigma_B^{\text{pp}}}{\sigma_{\text{DPS}}^{\text{pp}}}$$

POCKET FORMULA

Differential X-section single parton scattering for the process: $\text{pp} \longrightarrow A(B) + X$

Differential X-section double parton scattering for the process: $\text{pp} \longrightarrow A + B + X$

Results for W, Jet productions…



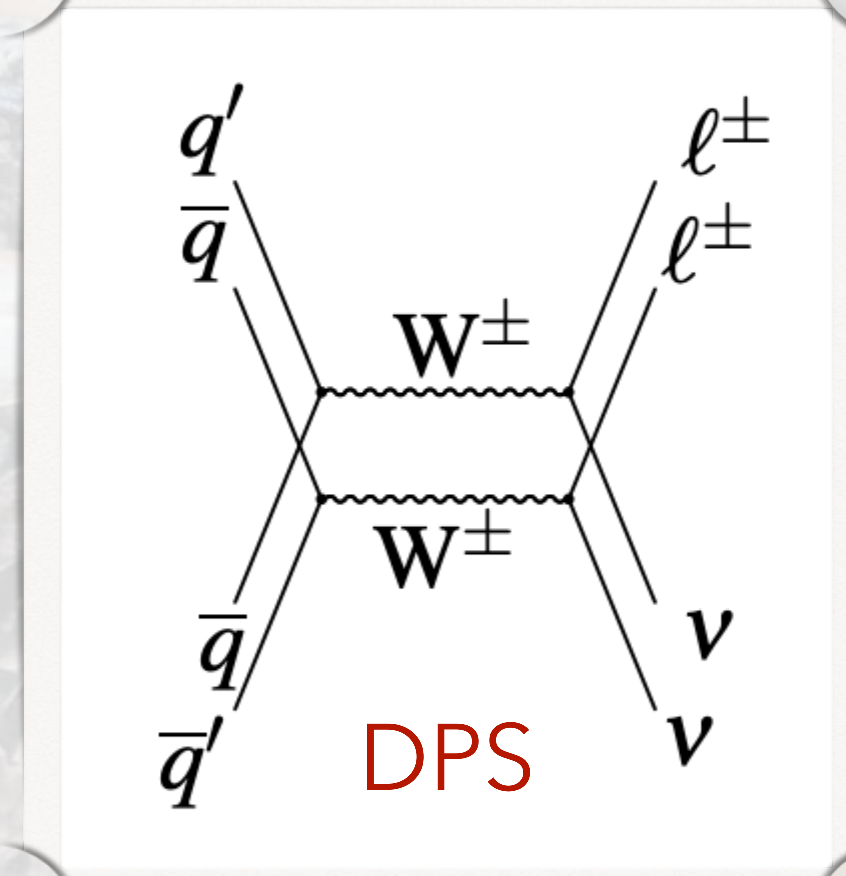
First observation of same sign WW via DPS:

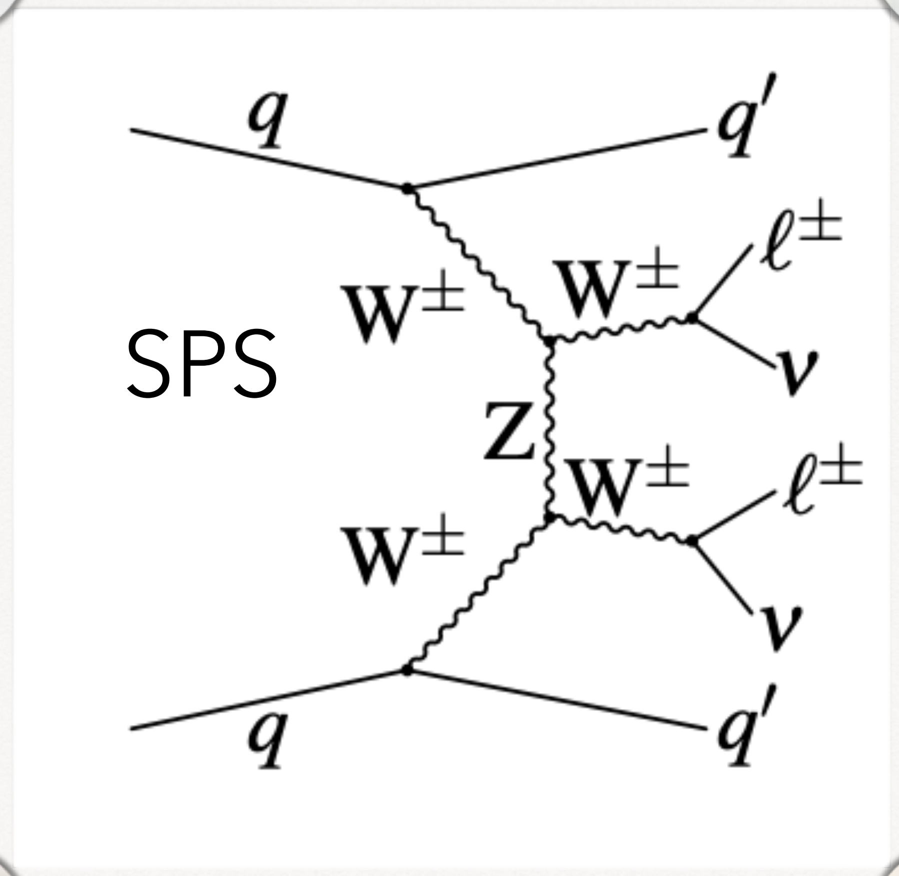
$$\sigma_{\text{eff}} = 12.2^{+2.9}_{-2.2} \quad \text{mb}$$

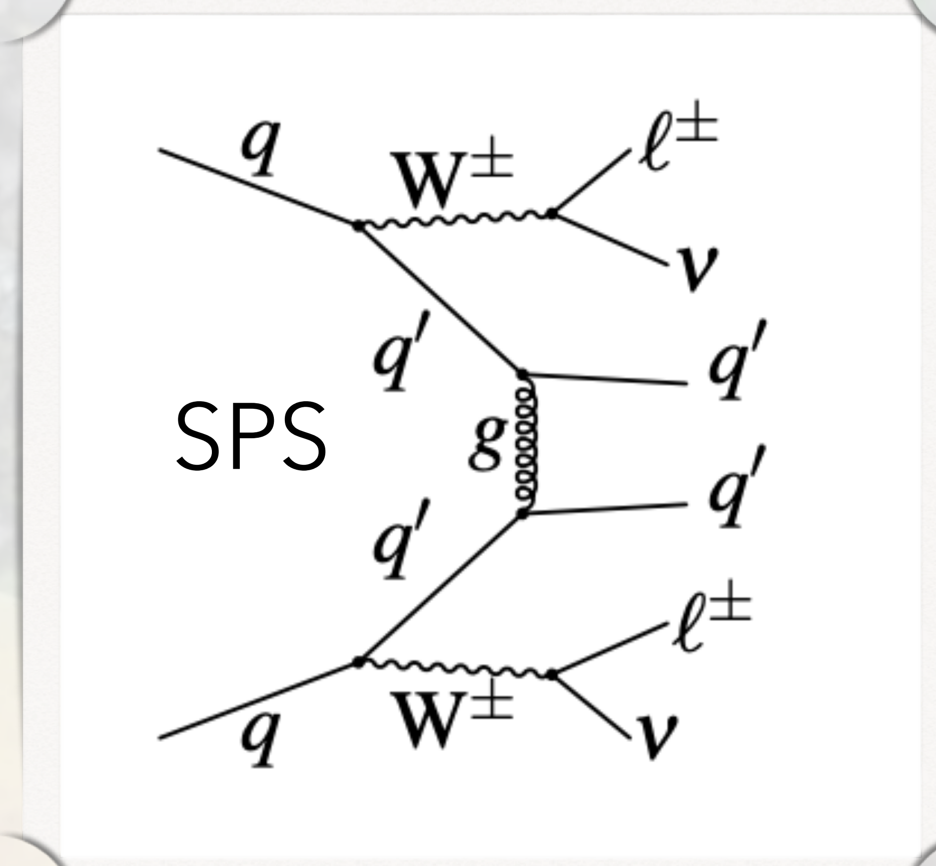
[CMS coll.], PRL 131 (2023) 091803

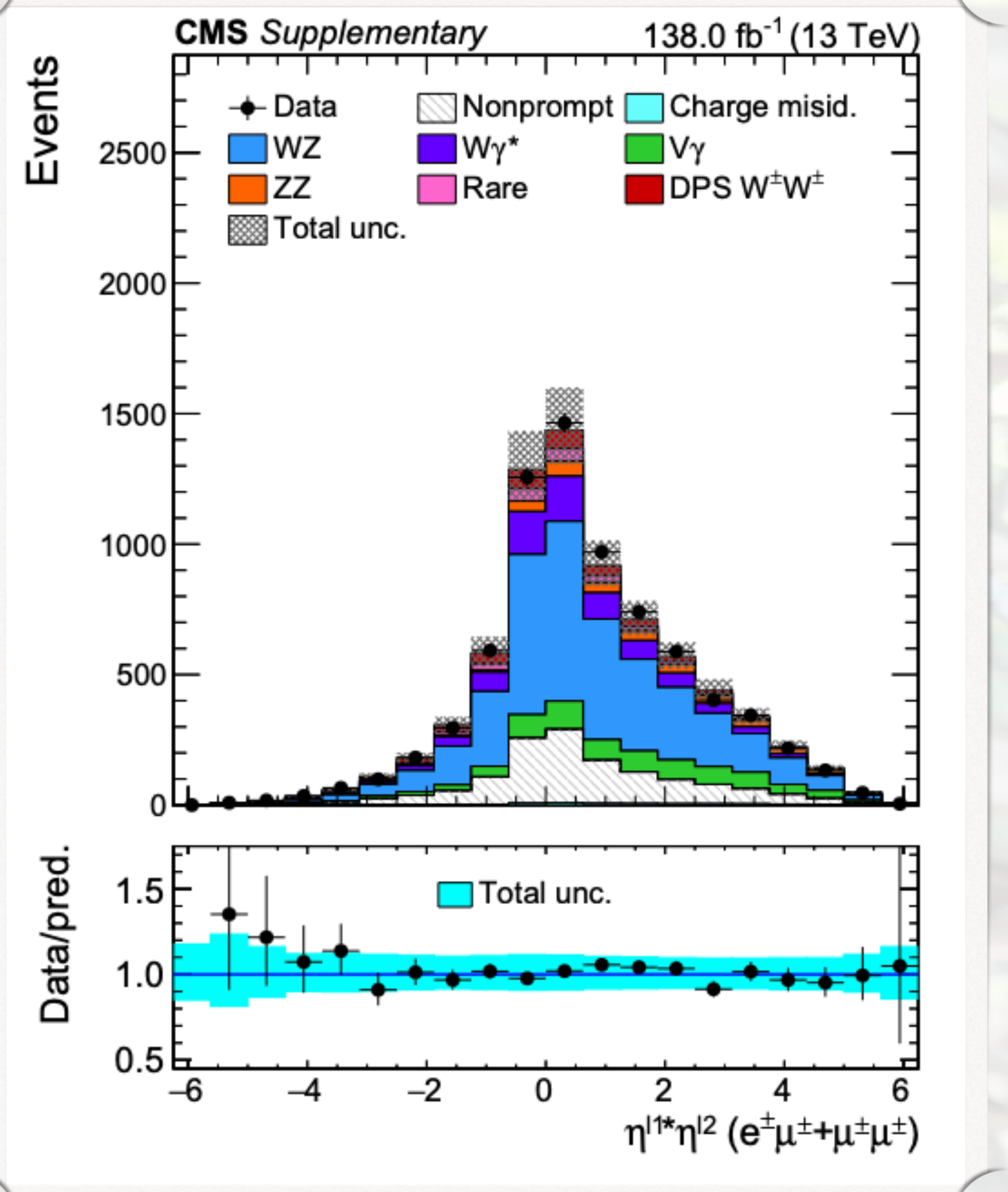
$$\sigma^{\text{DPS}} \sim 6.28 \text{ fb}$$

Some Data and Effective Cross Section - ssWW









$$\sigma_{\text{eff}}^{\text{pp}} = \frac{m}{2} \frac{\sigma_A^{\text{pp}} \sigma_B^{\text{pp}}}{\sigma_{\text{DPS}}^{\text{pp}}}$$

POCKET FORMULA

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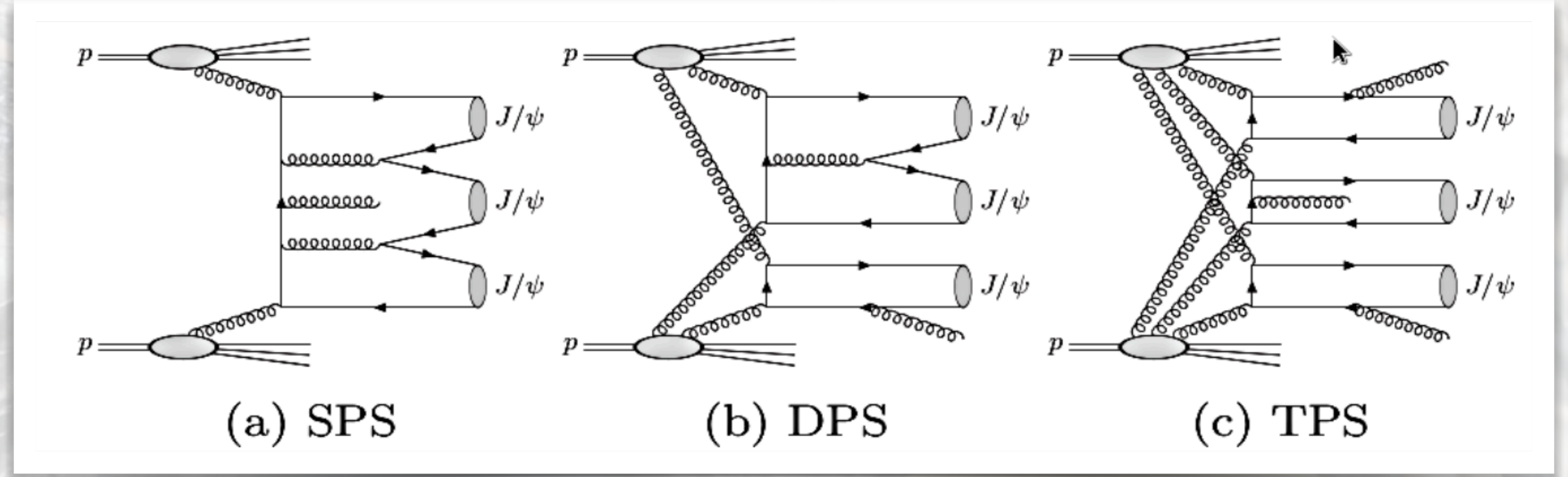
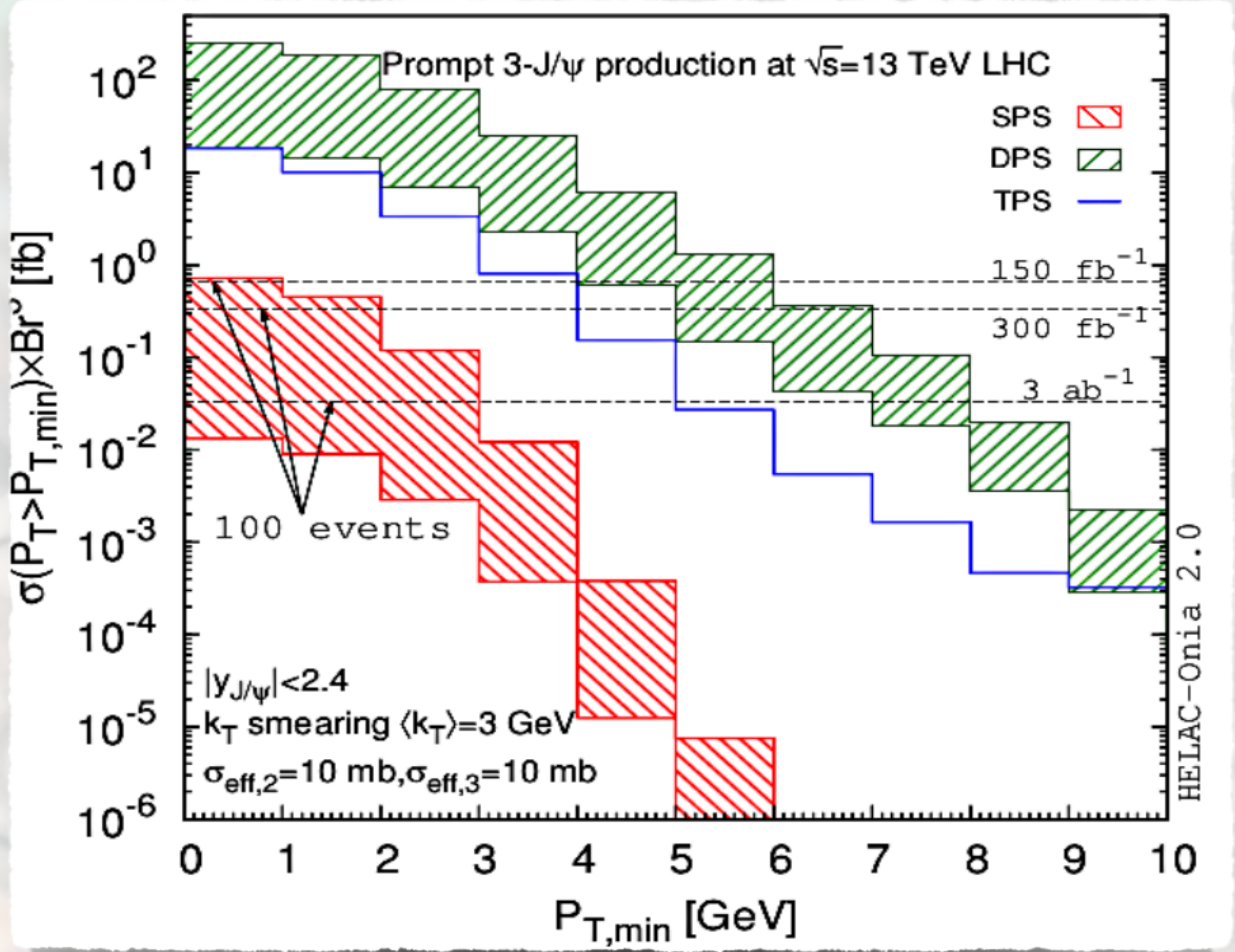
Some Data and Effective Cross Section - 3 J/ψ





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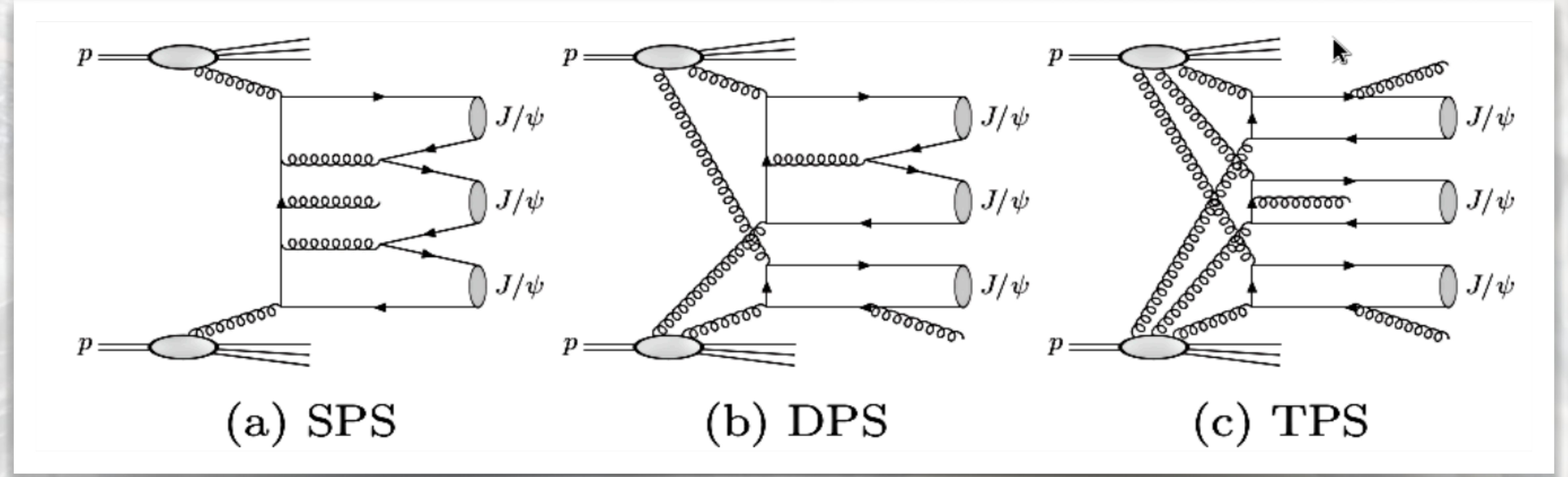
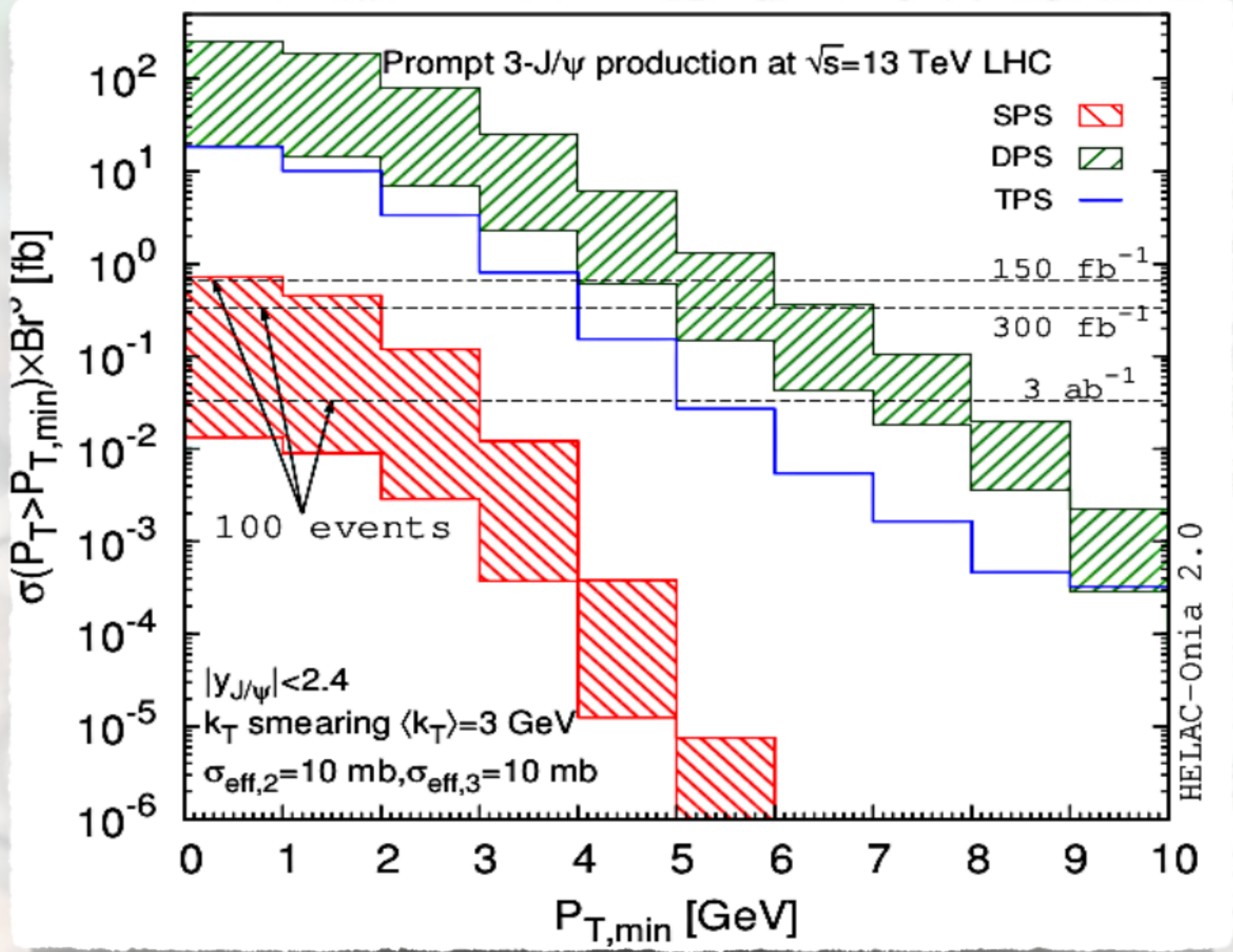
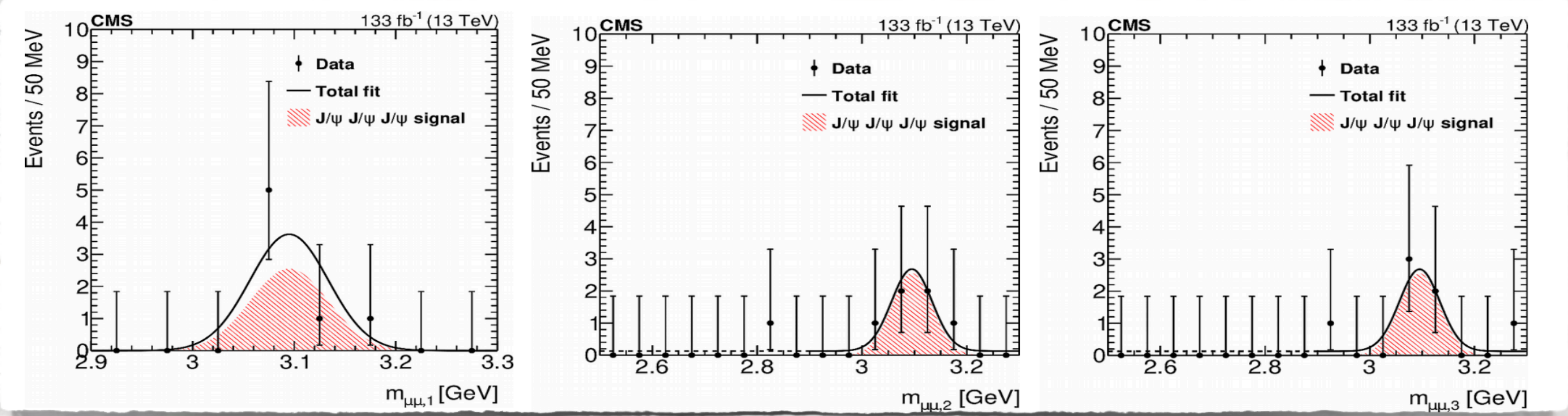


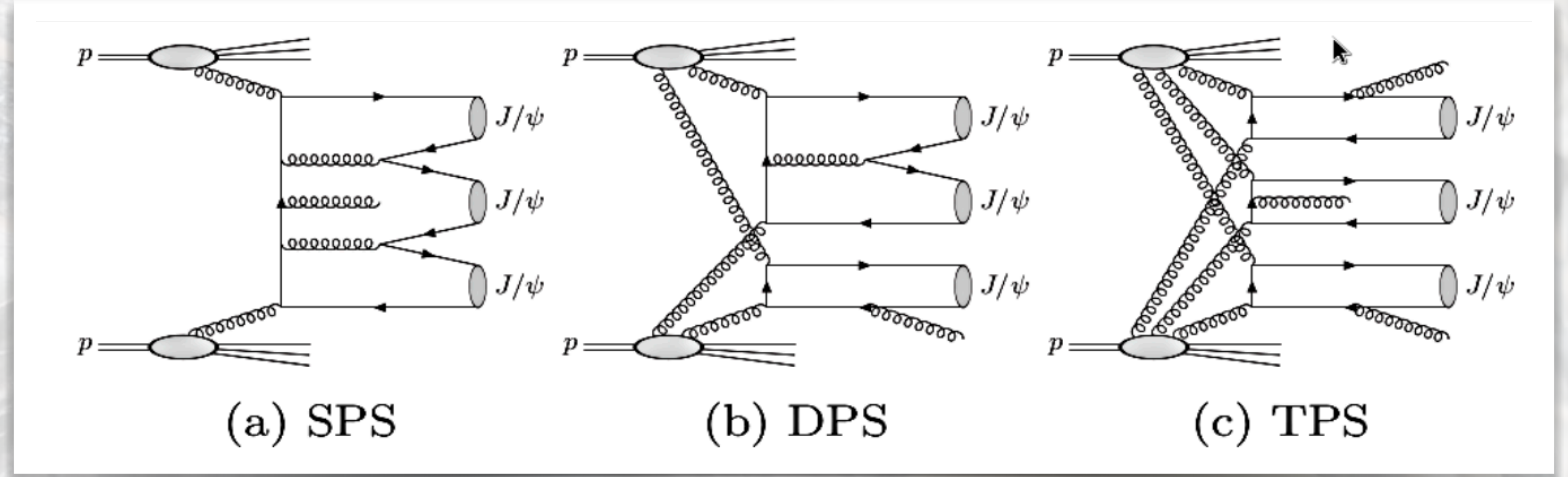


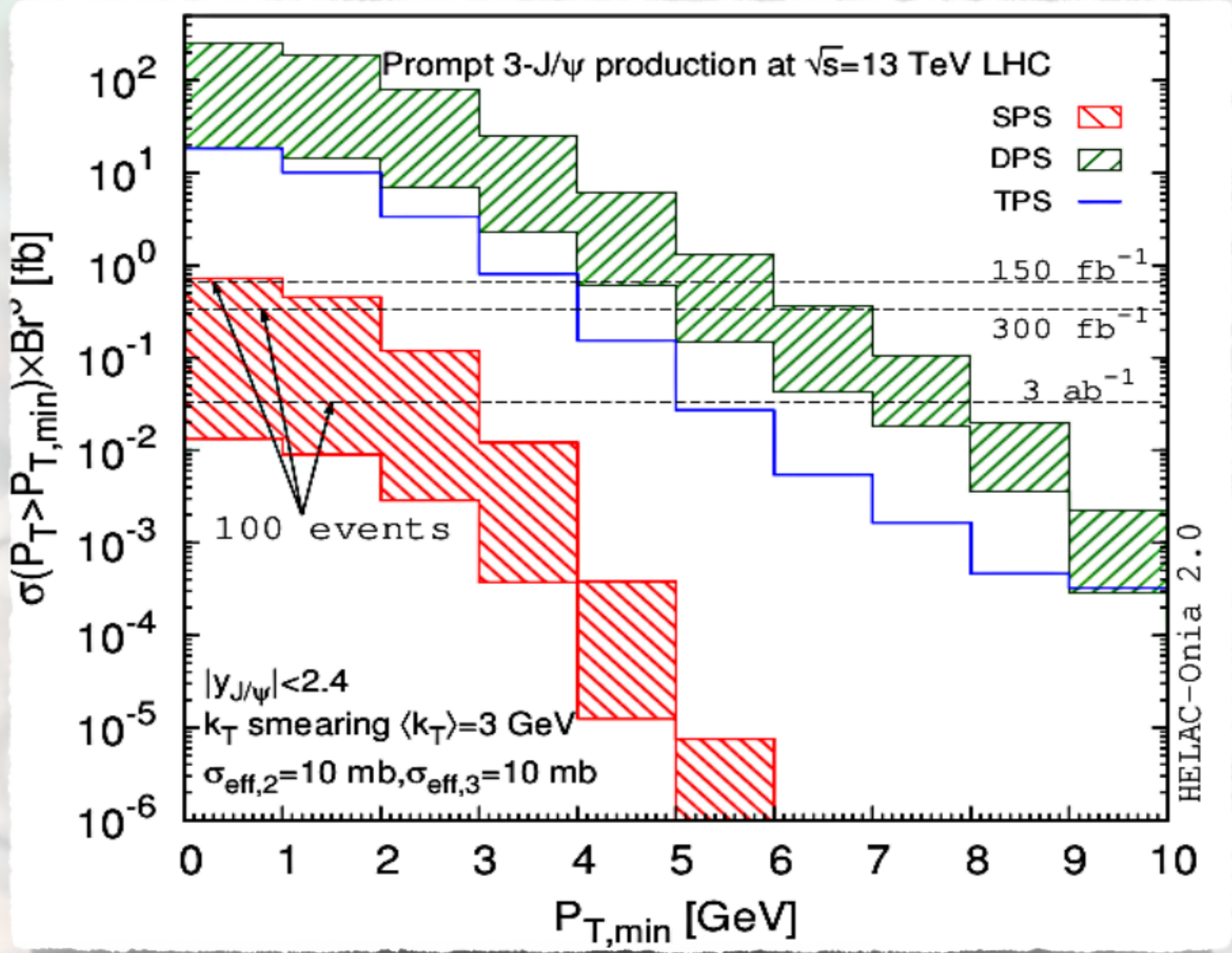


- $\sigma = 272^{+141}_{104}$ (stat) $\pm$ 17 (syst) fb

SPS -> 6%
DPS -> 74%
TPS ->20%

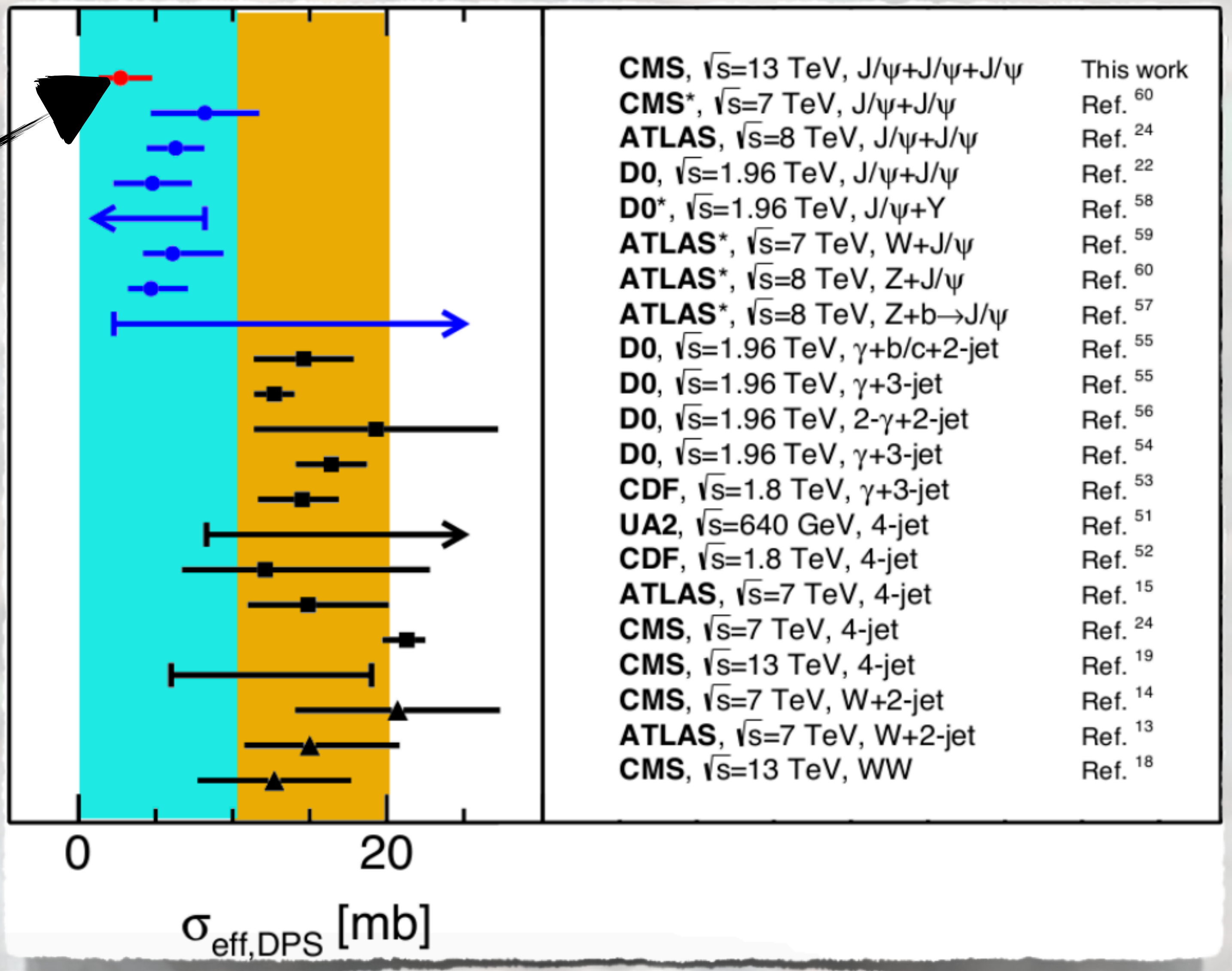
Some Data and Effective Cross Section - 3 J/ψ





Novel way to extract the DPS effective cross-section:

$$\sigma_{\text{eff,DPS}} = 2.7^{1.4}_{-1.0}\ (\text{exp})^{+1.5}_{-1.0}\ (\text{theo})\ \text{mb}$$



Some Data and Effective Cross Section

$$\sigma_{\rm eff}^{\rm pp} = \frac{m}{2}\frac{\sigma_A^{\rm pp}\sigma_B^{\rm pp}}{\sigma_{\rm DPS}^{\rm pp}}$$

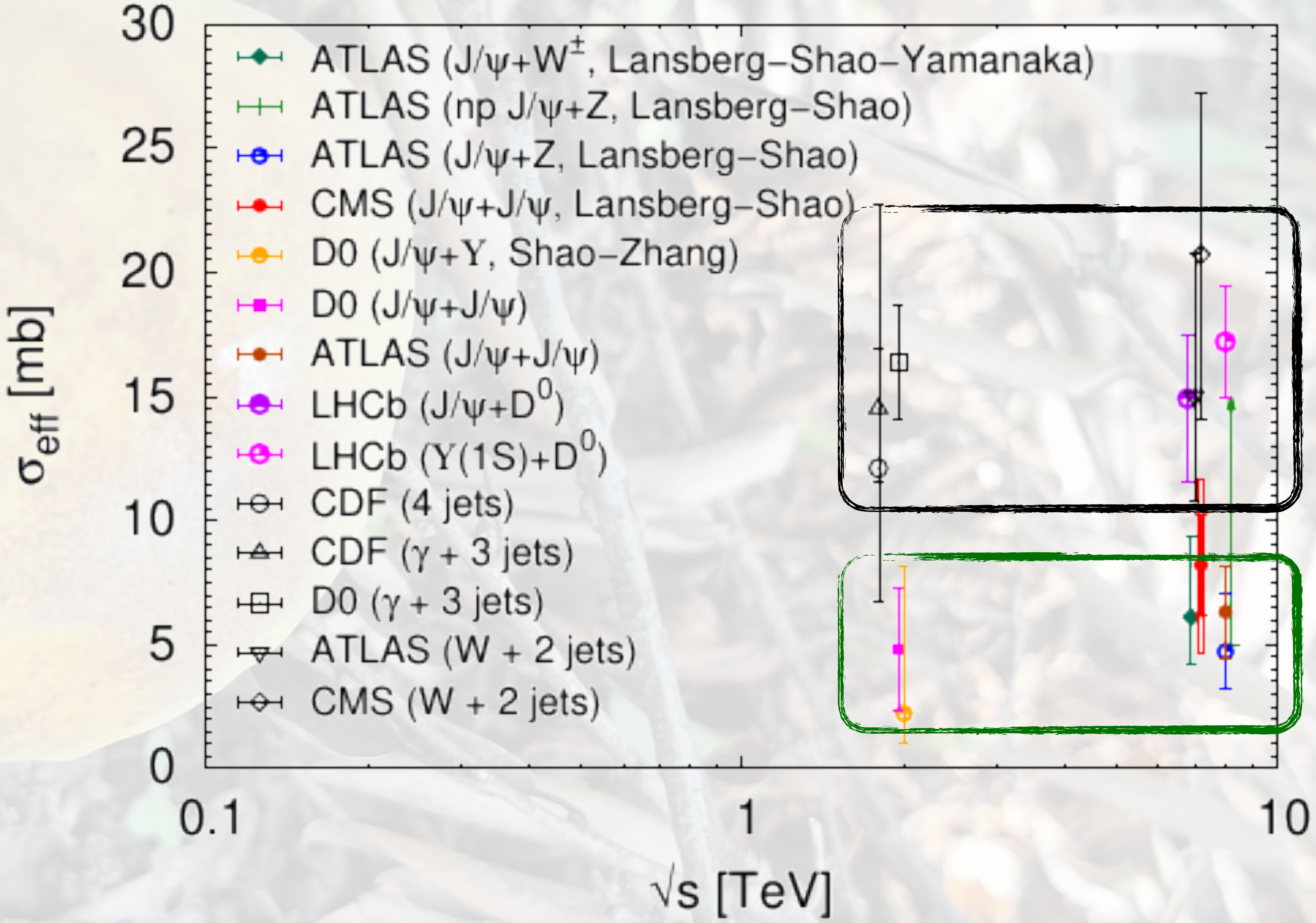
POCKET FORMULA

Differential X-section single parton scattering for the process: $pp \longrightarrow A(B) + X$

Differential X-section double parton scattering for the process: $pp \longrightarrow A + B + X$

Results for W, Jet productions…

Results for quarkonium productions



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Some Data and Effective Cross Section

$$\sigma_{\rm eff}^{\rm pp} = \frac{m}{2}\frac{\sigma_A^{\rm pp}\sigma_B^{\rm pp}}{\sigma_{\rm DPS}^{\rm pp}}$$

POCKET FORMULA

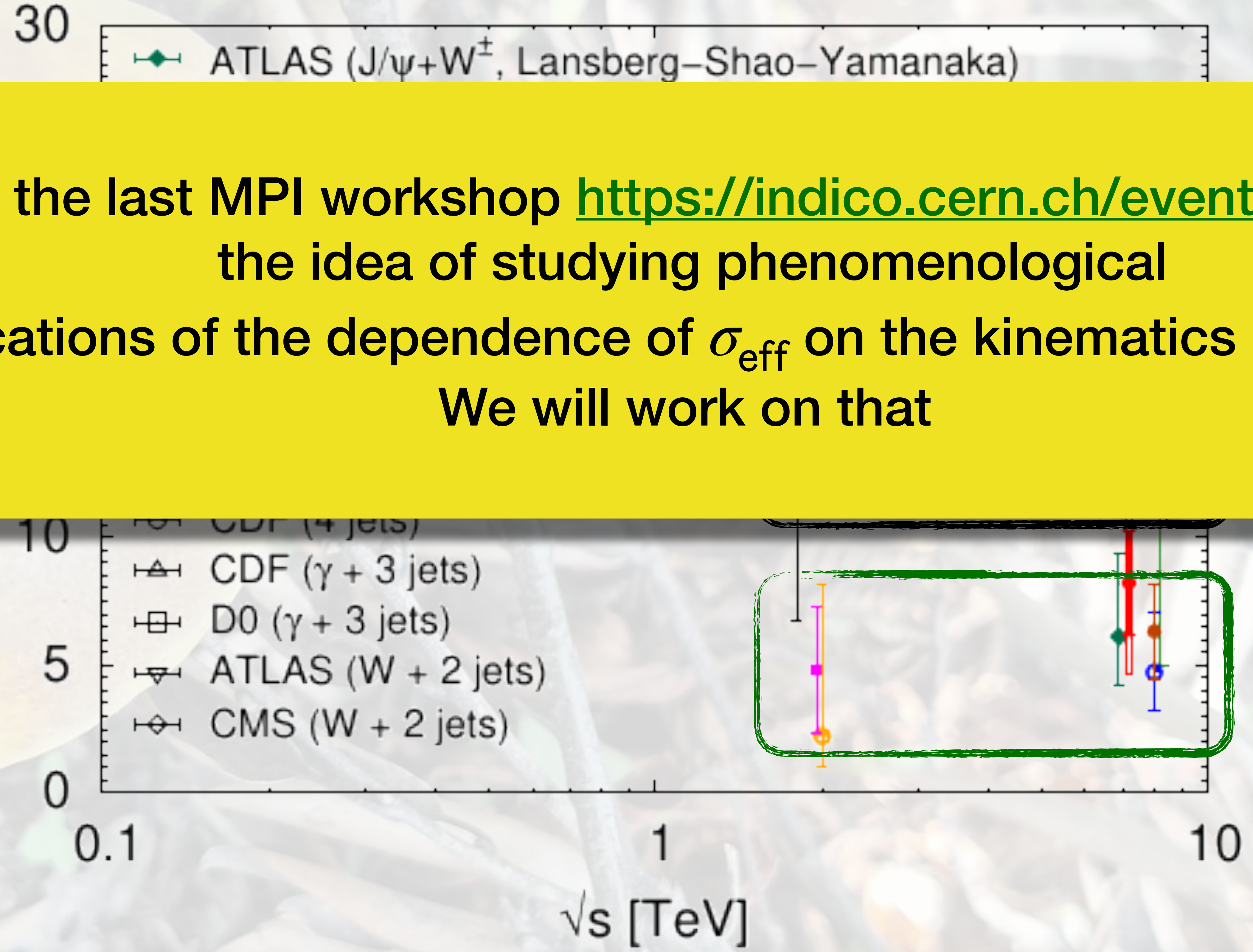
$\sigma_B^{\rm pp}$ → Differential X-section single parton scattering for the process: $pp \longrightarrow A(B) + X$

$\sigma_{\rm DPS}^{\rm pp}$ → Differential X-section double parton scattering for the process: $pp \longrightarrow A + B + X$

1) Process dependent?
2) Sensitive to correlations
3) Sensitive to the inner structure?

predicted by all models!

M.R. et al PLB 752,40 (2016)
M. Traini, M. R. et al, PLB 768, 270 (2017)
M. R. et al, Phys.Rev. D95 (2017) no.11, 114030



From the last MPI workshop https://indico.cern.ch/event/1281679/ the idea of studying phenomenological implications of the dependence of $\sigma_{\rm eff}$ on the kinematics came out!! We will work on that

First observation of WW via DPS:

$2.2^{+2.9}_{-2.2}$ mb

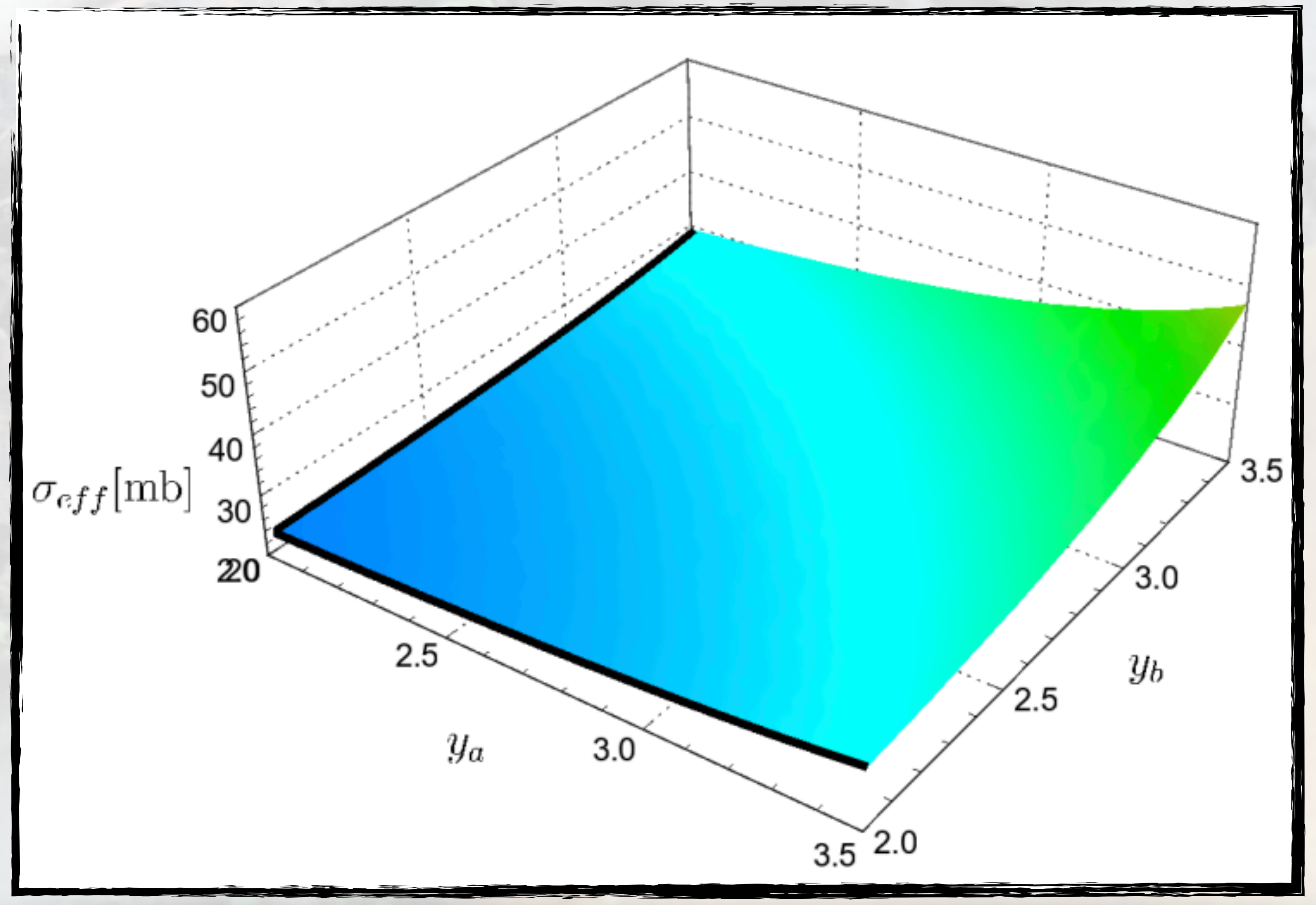
PRL 131 (2023) 091803

$\sigma \sim 6.28$ fb

Some Ideas on "*non-constant*" $\sigma_{\rm eff}$

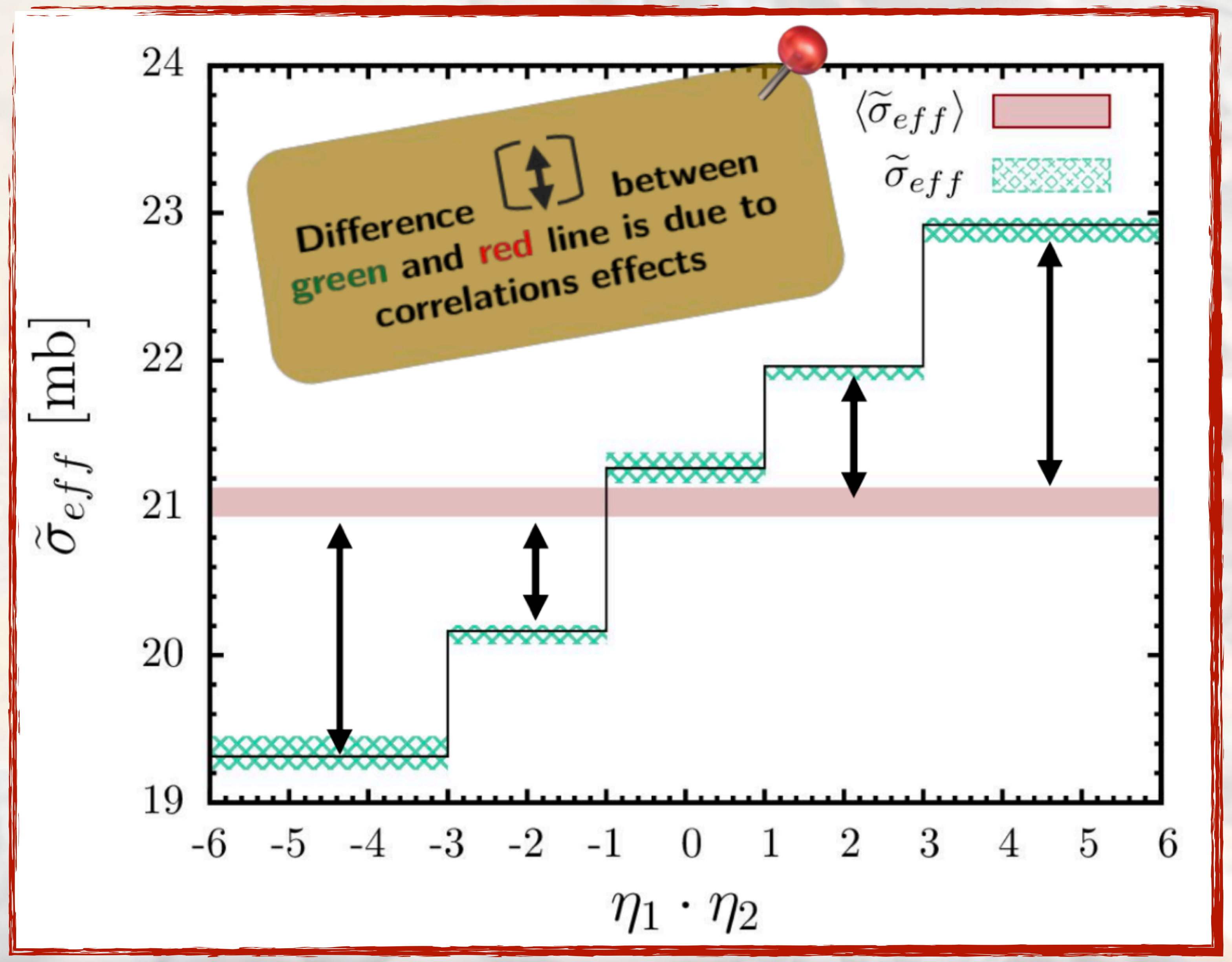
$$\sigma_{\rm eff} = \frac{m}{2}\frac{\sigma_A \sigma_B}{\sigma^{DPS}_{A+B}}$$

$$x_{1,3} = \frac{\sqrt{m_A^2 + k_{T,A}^2}}{\sqrt{s}} e^{\pm y_A}, \quad x_{2,4} = \frac{\sqrt{m_B^2 + k_{T,B}^2}}{\sqrt{s}} e^{\pm y_B}$$



An example for the calculation of $\sigma_{\rm eff}$ from gluon DPDs at high energy scales

M. R. and F. A. Ceccopieri, JHEP 09 (2019) 12, 125003



Predictions from the calculation of same sign W's production at the LHC

F. A. Ceccopieri, M. R. and S. Scopetta, PRD 95 (2017), no.11, 114030

Effective Cross Section and proton structure

If DPDs factorize in terms of PDFs then

$$\sigma_{\text{eff}}^{-1} = \int d^2 z_\perp \, \tilde{T}(z_\perp)^2 = \int \frac{d^2 k_\perp}{(2\pi)^2} T(k_\perp)^2$$

Effective Form Factor (EFF) =
FT of the probability distribution T

As for the standard FF:

$$\langle z_\perp^2 \rangle \propto \frac{d}{k_\perp dk_\perp} T(k_\perp) \Big|_{k_\perp=0}$$

$$T(k_\perp) \propto \int dx_1 dx_2 \, \tilde{F}(x_1, x_2, k_\perp)$$

First moment of DPD

From the asymptotic behavior we got the following relation:

$$\frac{\sigma_{\text{eff}}}{3\pi} \leq \langle z_\perp^2 \rangle \leq \frac{\sigma_{\text{eff}}}{\pi}$$

M. R. and F. A. Ceccopieri, PRD 97, no. 7, 071501 (2018)

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DPD = GPD ⊗ GPD

From the asymptotic behavior we got the following relation:

$$\frac{\sigma_{\text{eff}}}{3\pi} \leq \langle z_\perp^2 \rangle \leq \frac{\sigma_{\text{eff}}}{\pi}$$

Verified in all model calculations:

Constituent quark models for:

proton

M.R. and F. A. Ceccopieri, JHEP 09 (2019) 097

Pion

M.R. EPJC 80 (2020) 7, 678

W. Broniovski and E. R. Arriola PRD 101 (2020), 1, 014019

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M.R. EPJC 80 (2020) 7, 678

M. R. and F. A. Ceccopieri, PRD 97, no. 7, 071501 (2018)

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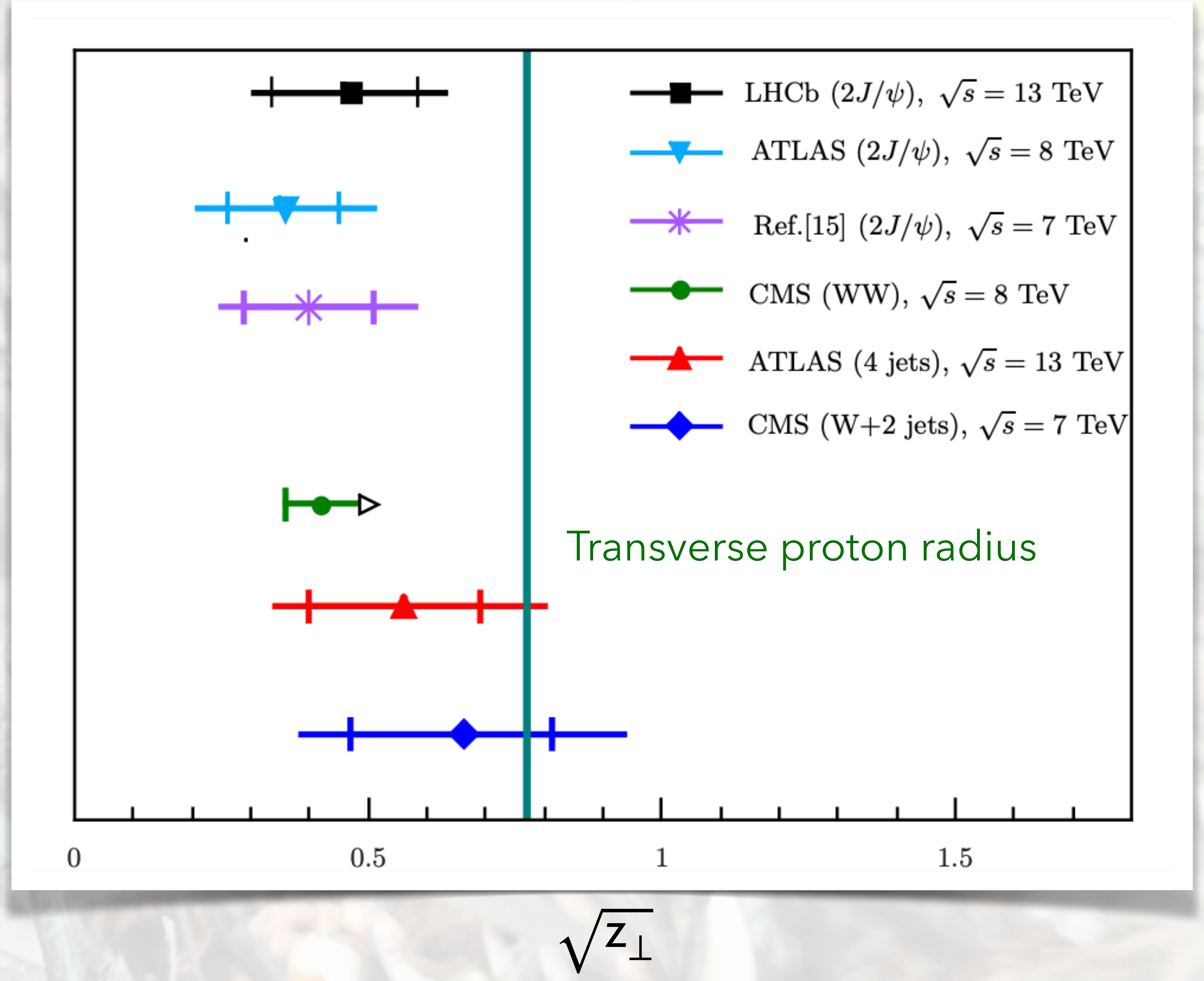

M. R. and F. A. Ceccopieri, PRD 97, no. 7, 071501 (2018)

Effective Cross Section and proton structure

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$$\sigma_{eff}^{-1} = \int d^2z_\perp \, \tilde{T}(z_\perp)^2 = \int \frac{d^2k_\perp}{(2\pi)^2} T(k_\perp)^2$$

1) THE MEAN DISTANCE IS LOWER THEN THE PROTON RADIUS!

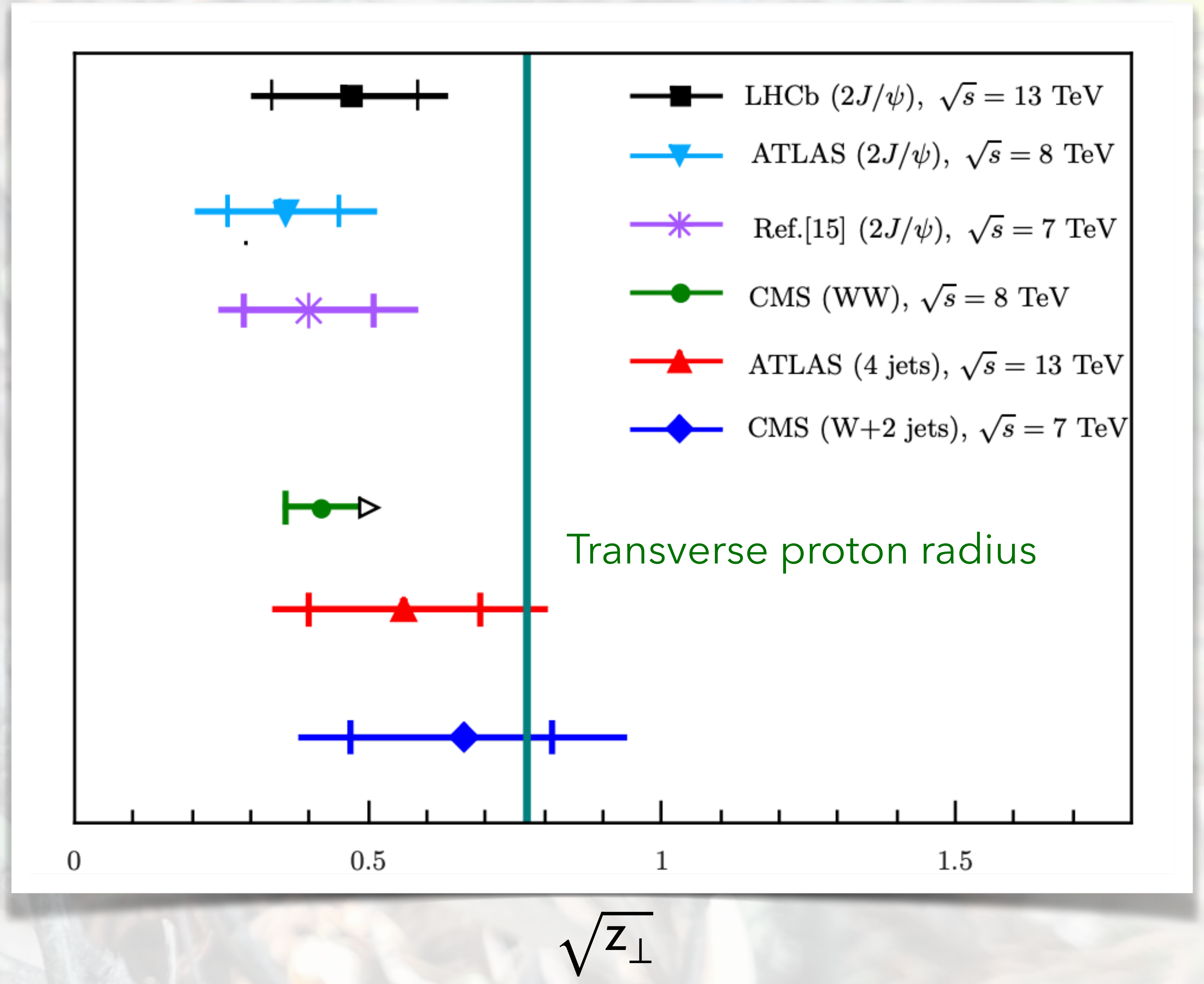
2) in hadron-hadron collisions we do not access directly the distance!

M.R. and F. A. Ceccopieri, JHEP 09 (2019) 097

From the asymptotic behavior we got the following relation:

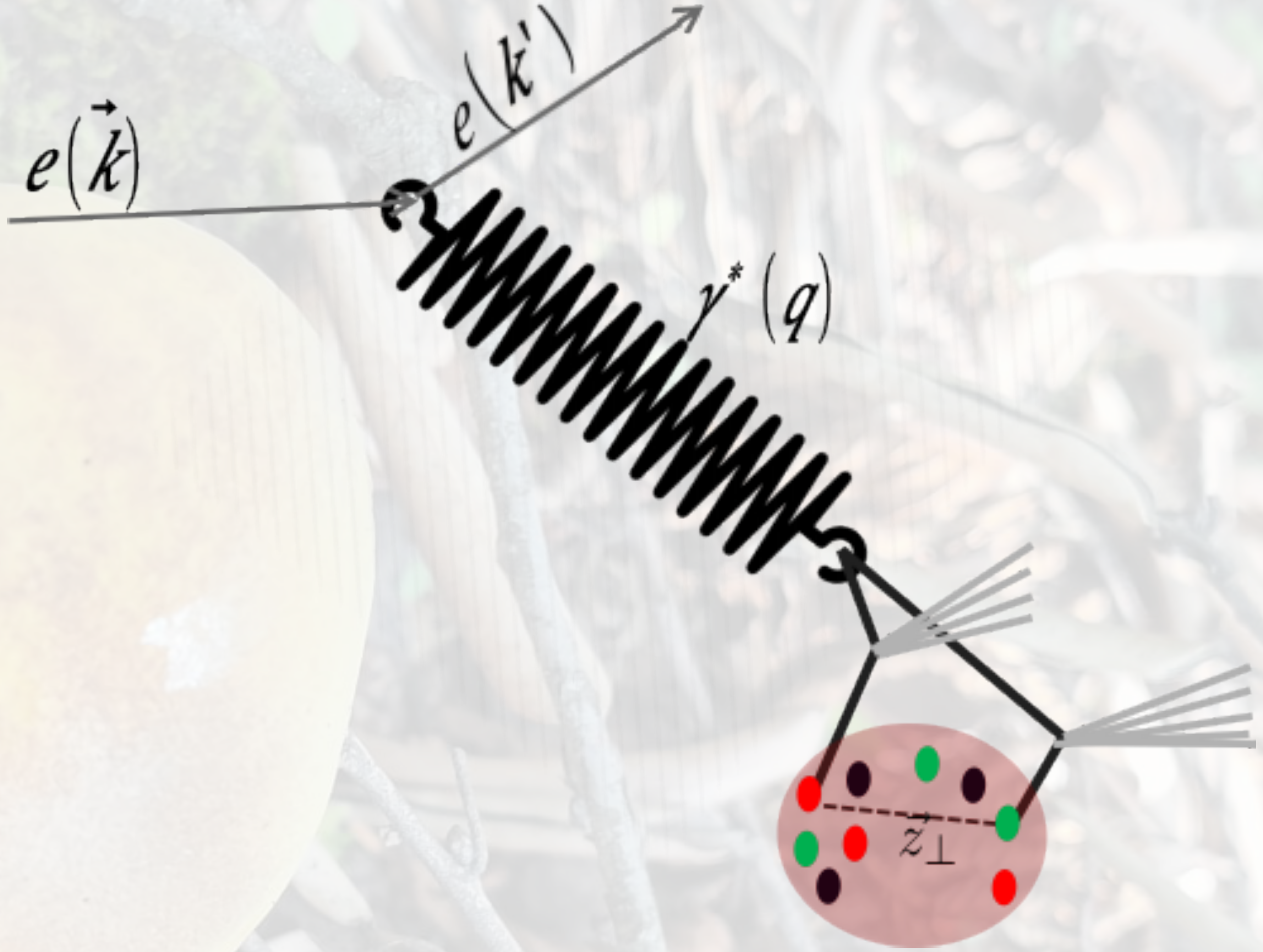
$$\frac{\sigma_{eff}}{3\pi} \leq \langle z_\perp^2 \rangle \leq \frac{\sigma_{eff}}{\pi}$$

M. R. and F. A. Ceccopieri, PRD 97, no. 7, 071501 (2018)



DPS in $\gamma - p$ interactions

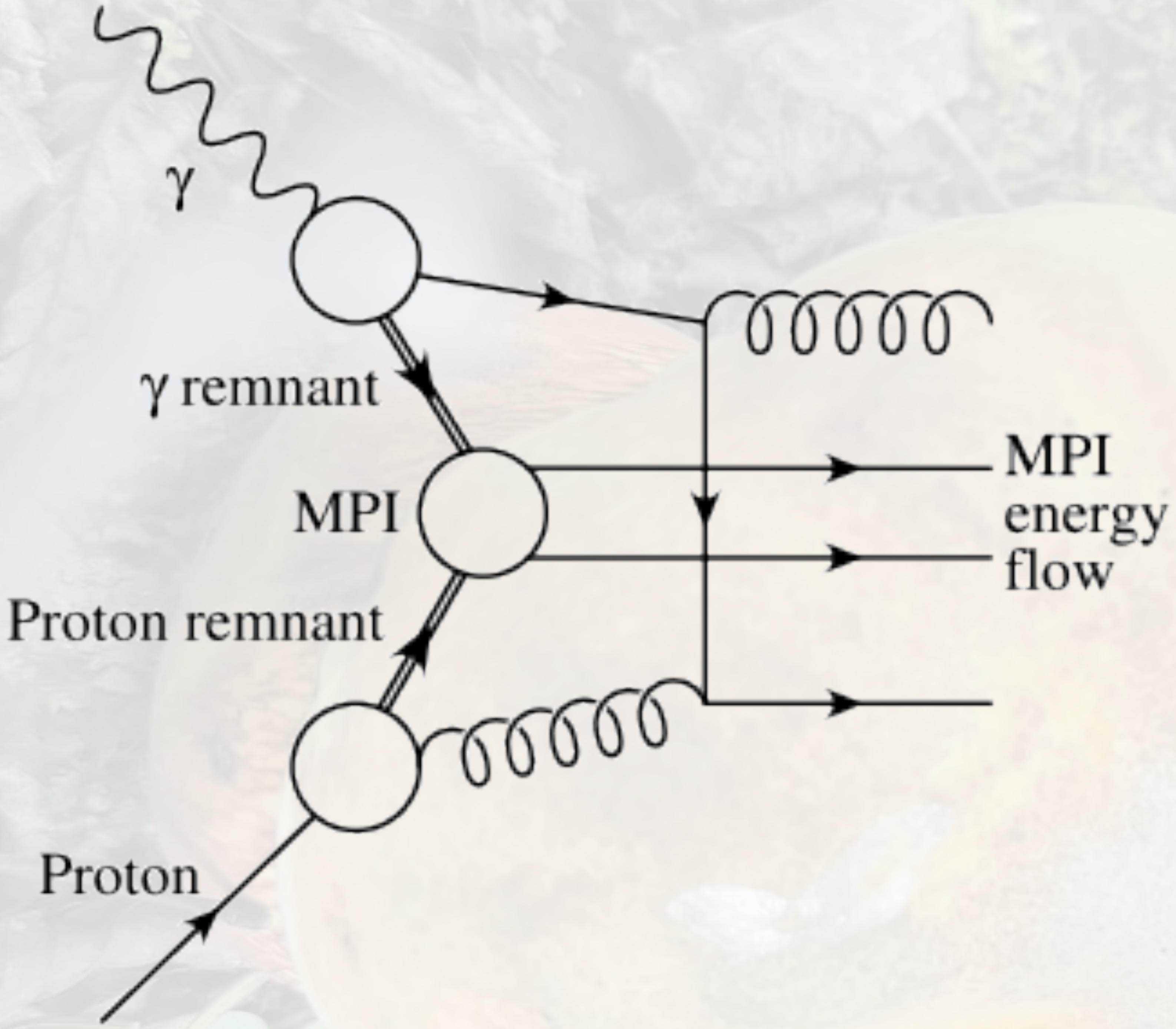
We consider the possibility offered by a DPS process involving a photon FLACTUATING in a quark-antiquark pair interacting with a proton:

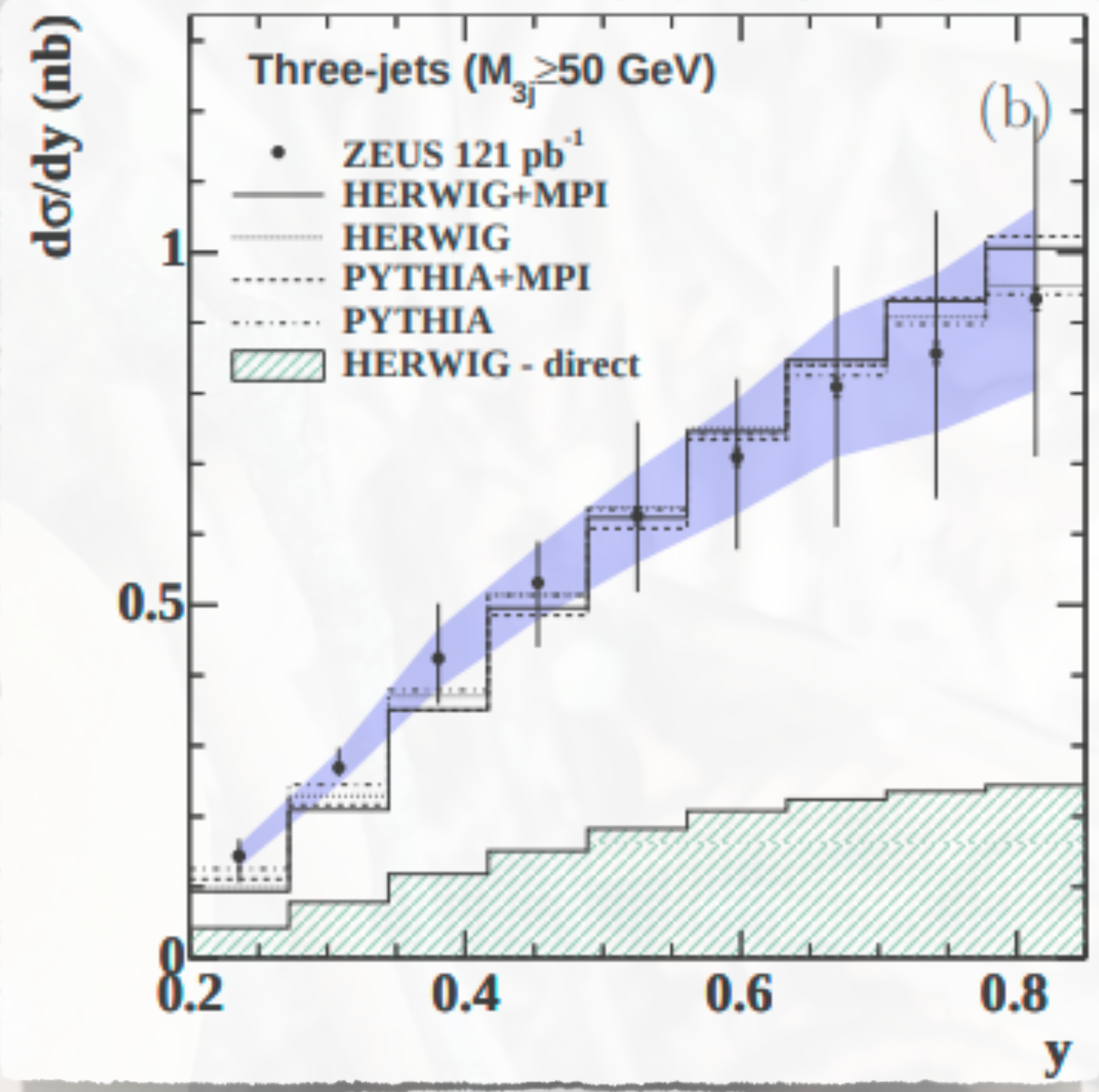


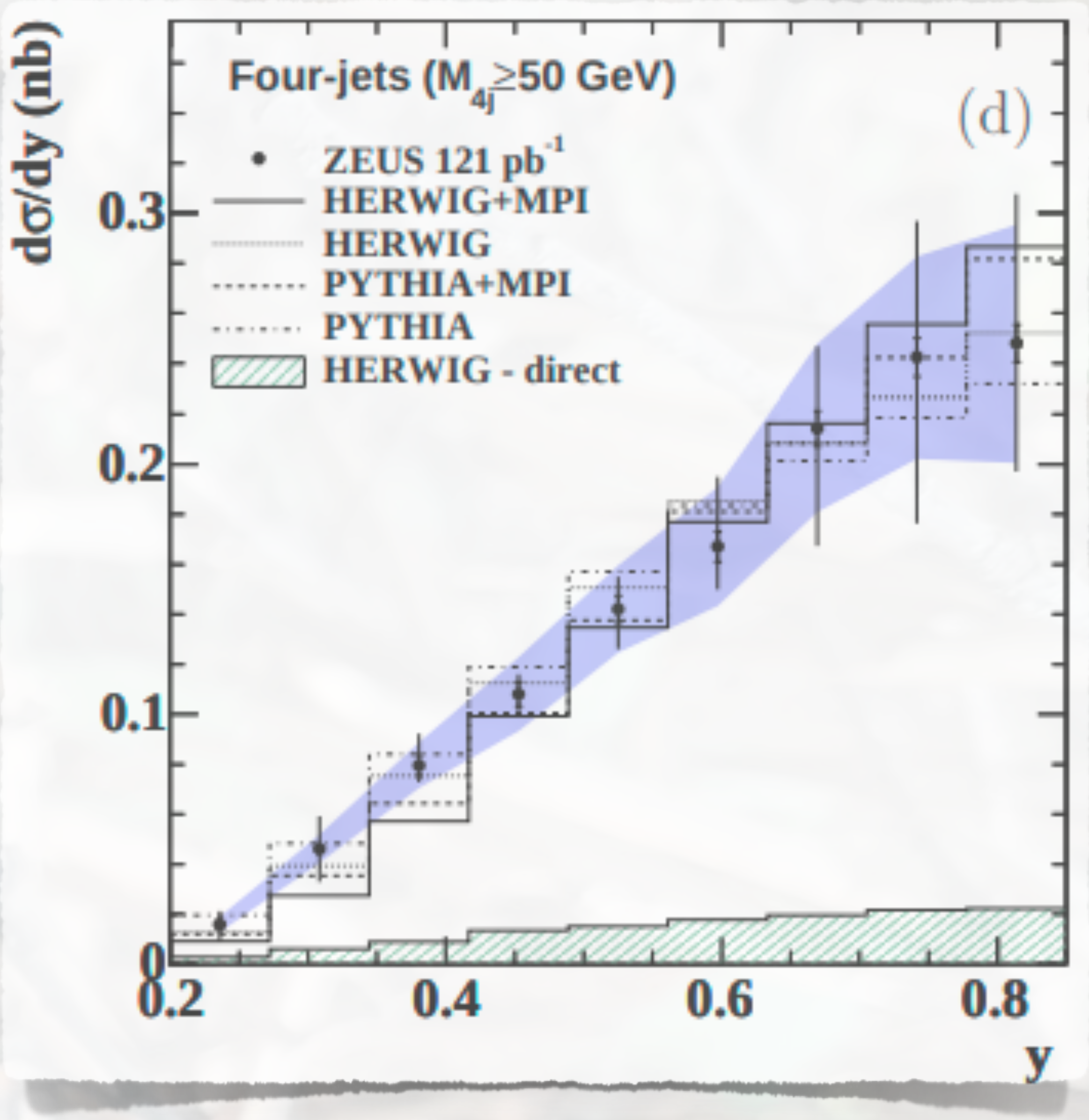
M. R. and F. A. Ceccopieri, PRD 105 (2022) L011501

DPS in $\gamma - p$ interactions

Already at HERA the importance of MPI for the **3,4 jets photo-production** has been addressed:





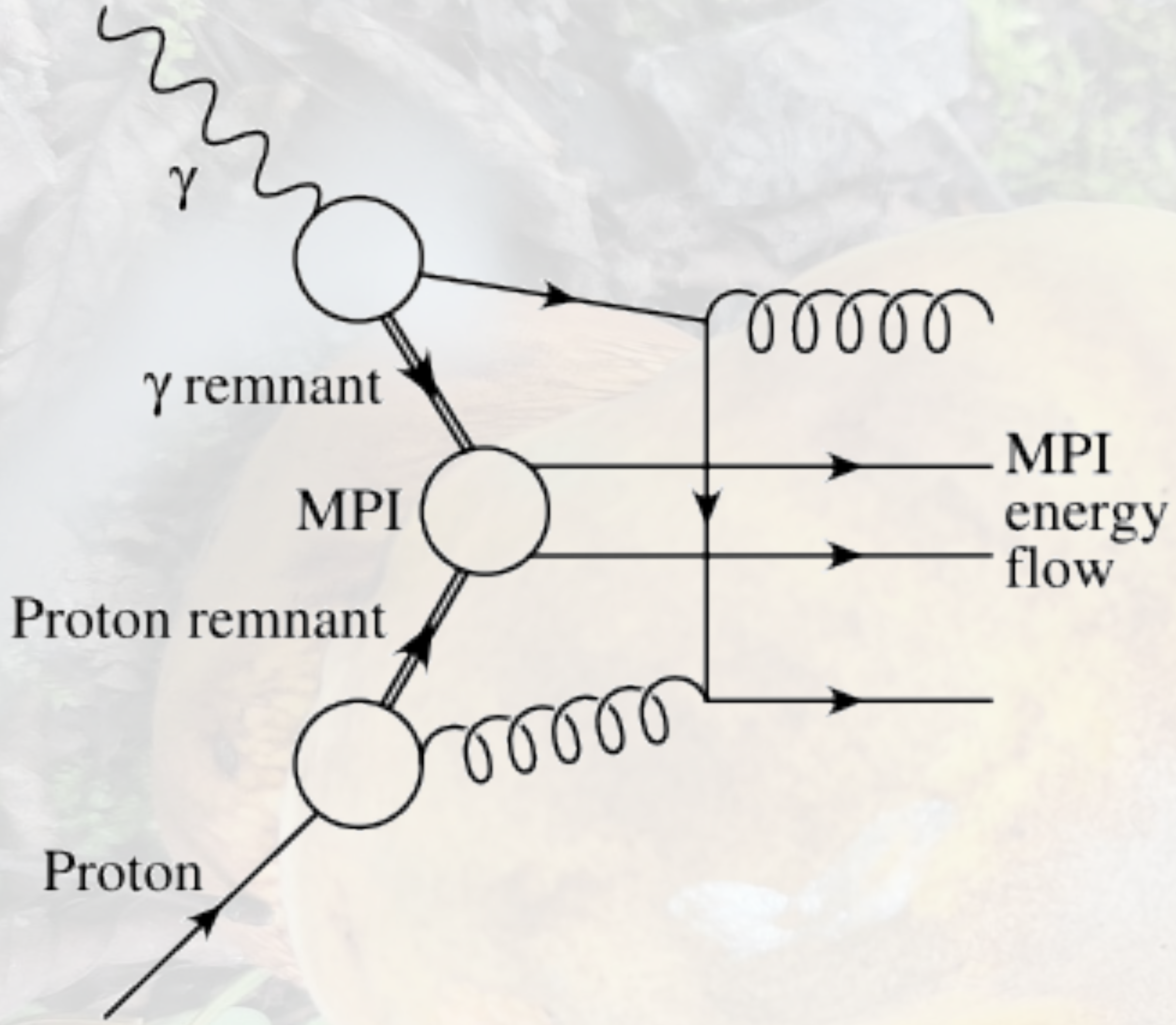


J. R. Forshaw et al, Z phys. C 72, 637

S. Chekanov et al [ZEUS coll.], Nucl. Phys B 792,1 (2008)

DPS in $\gamma - p$ interactions

In order to study the impact of the DPS contribution to a process initiated via photon-proton interactions we evaluated the 4-JET photo-production at HERA (**S. Checkanov et al. (ZEUS), Nucl. Phys B792, 1 (2008)**)



For this first investigation, we make use of the POCKET FORMULA:

$$d\sigma^{4j}_{DPS} = \frac{1}{2}\sum_{ab,cd}\int dy\, dQ^2 \frac{f_{\gamma/e}(y,Q^2)}{\sigma^{\gamma p}_{eff}(Q^2)} \times$$
$$\times \int dx_{p_a} dx_{\gamma_b} f_{a/p}(x_{p_a}) f_{b/\gamma}(x_{\gamma_b}) d\hat{\sigma}^{2j}_{ab}(x_{p_a}, x_{\gamma_b})$$
$$\times \int dx_{p_c} dx_{\gamma_d} f_{c/p}(x_{p_c}) f_{d/\gamma}(x_{\gamma_d}) d\hat{\sigma}^{2j}_{cd}(x_{p_c}, x_{\gamma_d})$$

SPS * SPS

Flux Factor: **P. Nason et al, PLB319**

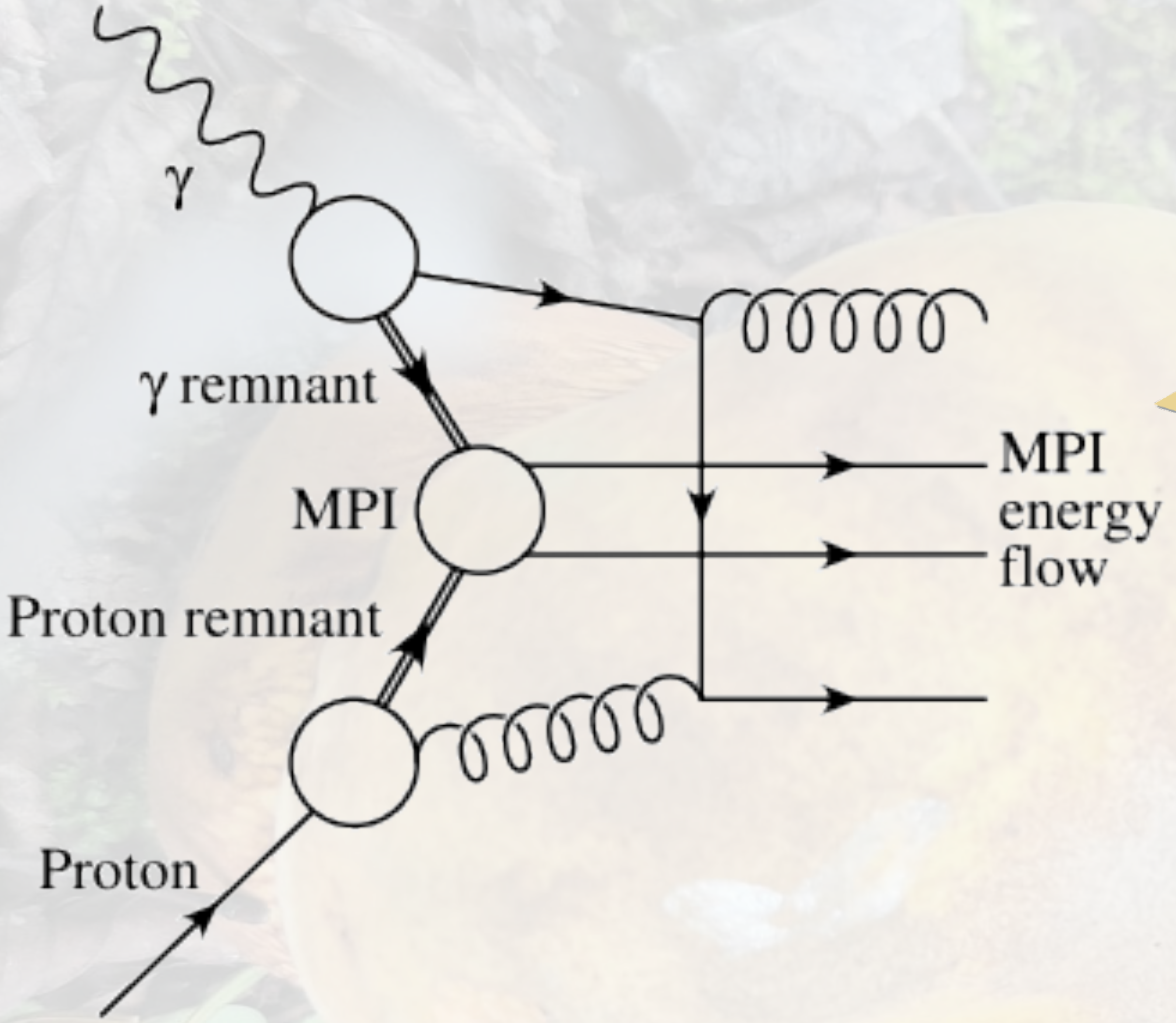
Proton PDF (J. Pumplin et al. JHEP 07, 012 (2002)

Photon PDF (M. Gluck et al. PRD46, 1973 (1992)

DPS in $\gamma - p$ interactions

In order to study the impact of the DPS contribution to a process initiated via photon-proton interactions we evaluated the 4-JET photo-production at HERA (**S. Checkanov et al. (ZEUS), Nucl. Phys B792, 1 (2008)**)

For this first investigation, we make use of the POCKET FORMULA:



(J. Pumplin et al. ...

The main quantity we have to evaluate is:

$$\sigma_{eff}^{\gamma p}(Q^2)$$

Flux Factor

P. Nason et al, PLB319

$$\ldots \frac{f_{\gamma/e}(y, Q^2)}{\sigma_{eff}^{\gamma p}(Q^2)} \times$$

$$\ldots (x_{\gamma_b}) d\hat{\sigma}_{ab}^{2j}(x_{p_a}, x_{\gamma_b}) \quad \text{SPS}$$

$$*$$

$$\ldots_{\gamma}(x_{\gamma_d}) d\hat{\sigma}_{cd}^{2j}(x_{p_c}, x_{\gamma_d}) \quad \text{SPS}$$

Photon PDF

(M. Gluck et al. PRD46, 1973 (1992))

The $\gamma - p$ effective cross-section

M. R. and F. A. Ceccopieri, PRD 105 (2022) L011501

The expression of this quantity is very similar to the proton-proton collision case and can be formally derived by comparing the product of SPS cross sections and the DPS one obtained in **Gaunt, JHEP 01, 042 (2013)** and describing a DPS from a vector bosons splitting with given Q^2 virtuality

Proton EFF

$$\left[\sigma_{\text{eff}}^{\gamma p}(Q^2)\right]^{-1} = \int \frac{d^2k_\perp}{(2\pi^2)} T_p(k_\perp) T_\gamma(k_\perp; Q^2)$$

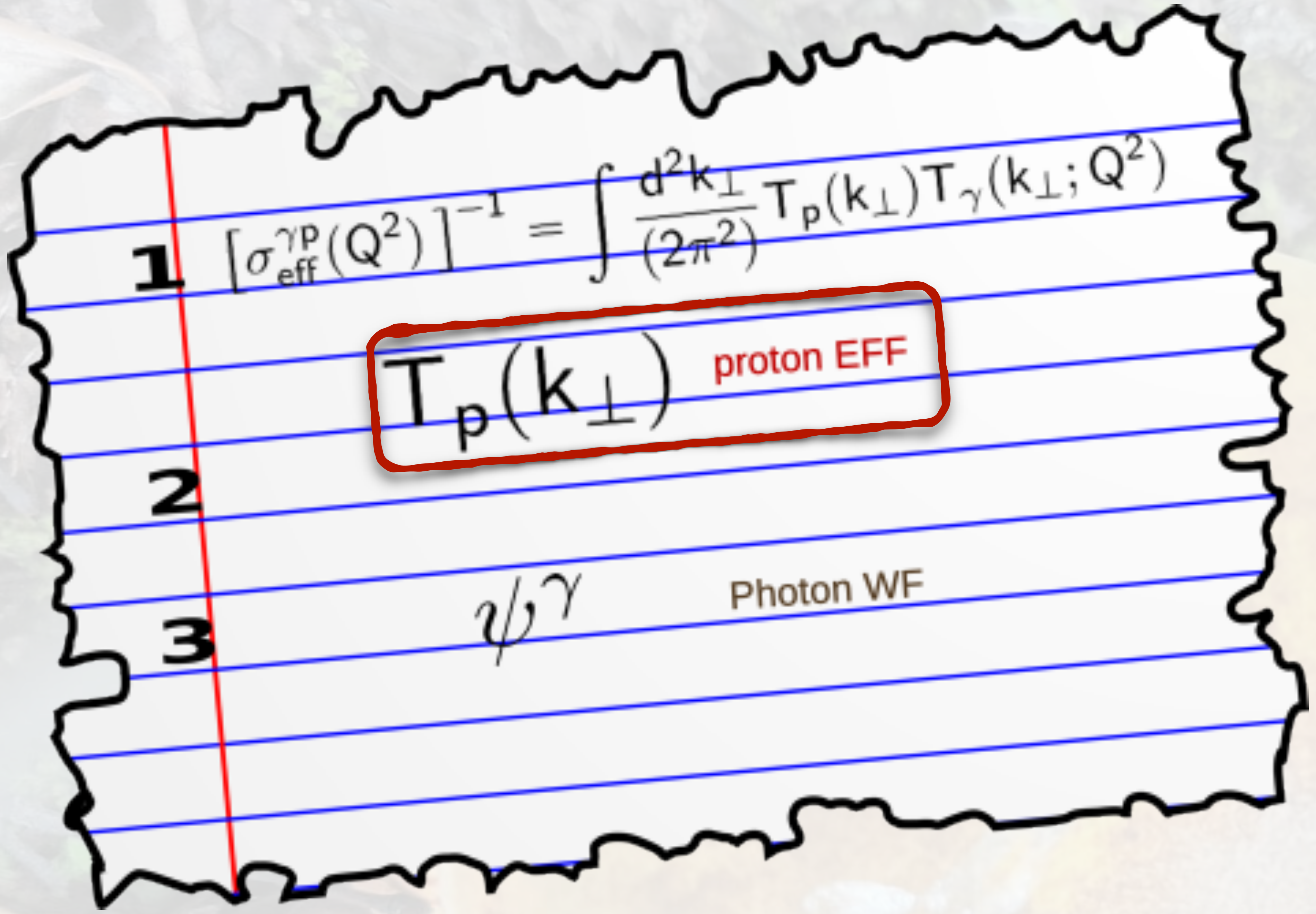
Photon EFF

The full DPS cross section depends on the amplitude of the splitting photon in a $q - \bar{q}$ pair. The latter can be formally described within a Light-Front (LF) approach in terms of LF wave functions

The $\gamma - p$ effective cross-section

M. R. and F. A. Ceccopieri, PRD 105 (2022) L011501

The main ingredients of the calculations:



For the proton EFF use has been made of three choices:

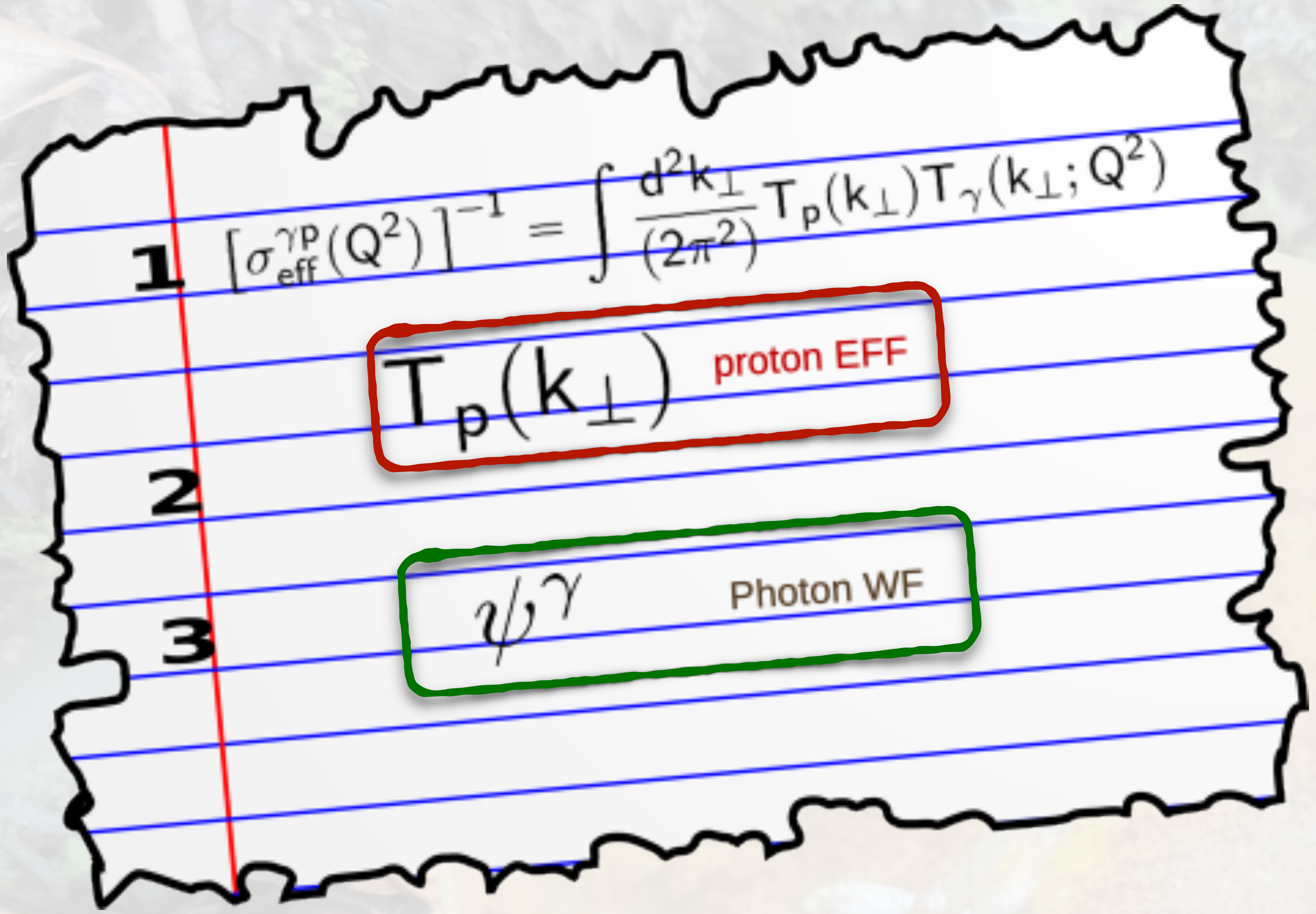
1) G1 $\quad e^{-\alpha_1 k_\perp^2}, \quad \alpha_1 = 1.53 \text{ GeV}^{-2} \Longrightarrow \sigma_{\text{eff}}^{pp} = 15 \text{ mb}$

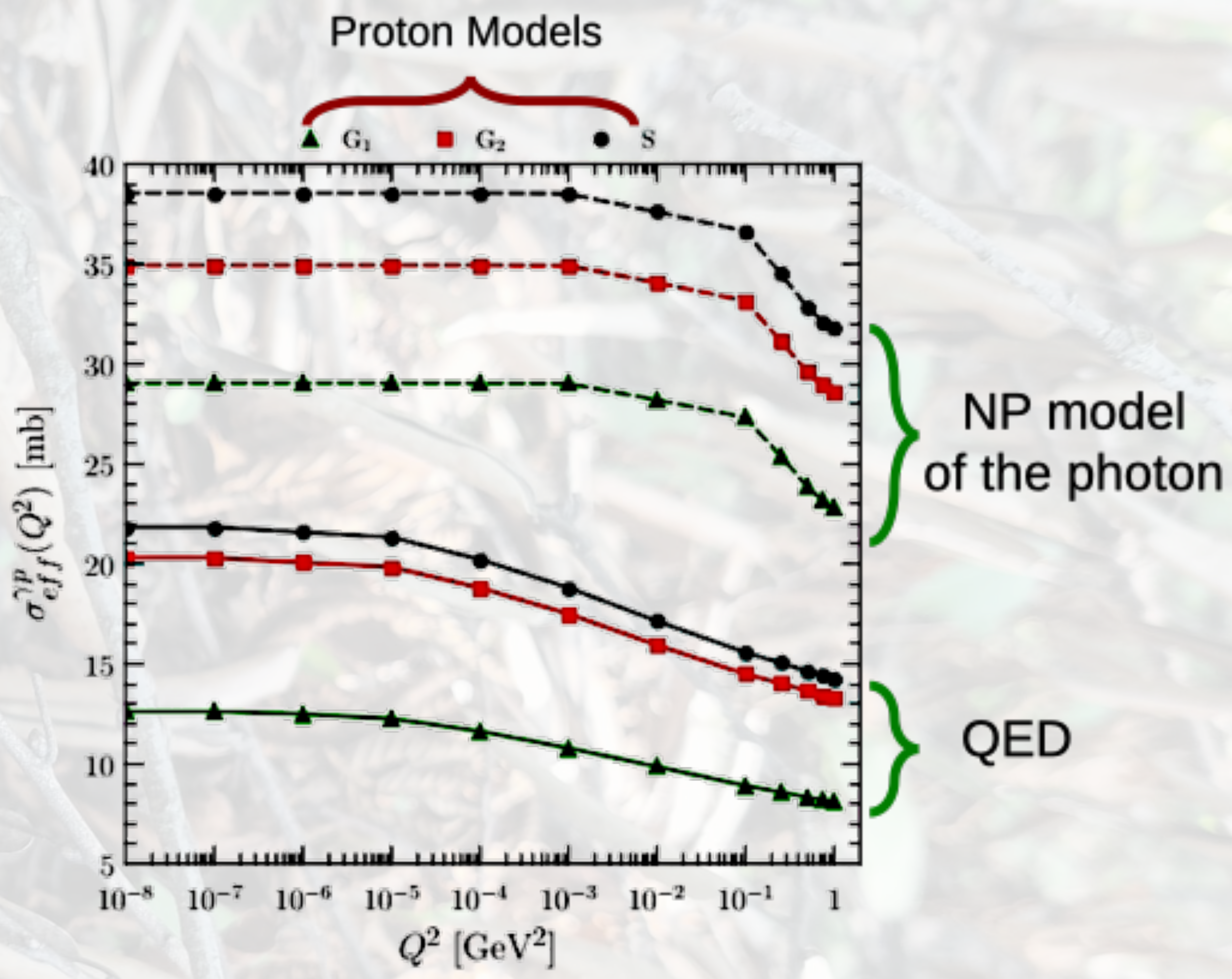
2) G2 $\quad e^{-\alpha_2 k_\perp^2}, \quad \alpha_2 = 2.56 \text{ GeV}^{-2} \Longrightarrow \sigma_{\text{eff}}^{pp} = 25 \text{ mb}$

3) S $\quad \left(1 + \frac{k_\perp^2}{m_g^2}\right)^{-4}, \quad m_g^2 = 1.1 \text{ GeV}^2 \Longrightarrow \sigma_{\text{eff}}^{pp} = 30 \text{ mb}$

The $\gamma - p$ effective cross-section

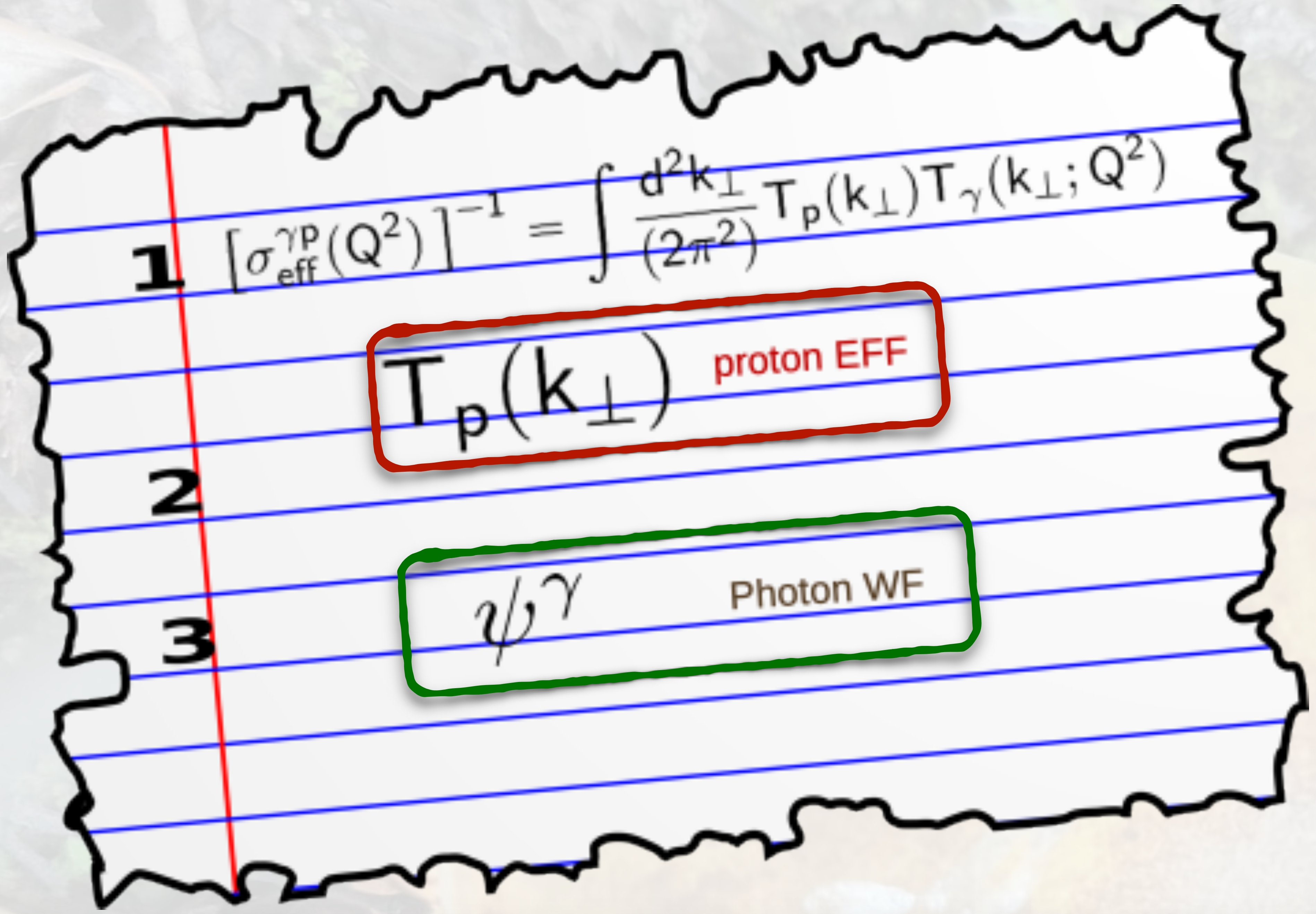
M. R. and F. A. Ceccopieri, PRD 105 (2022) L011501



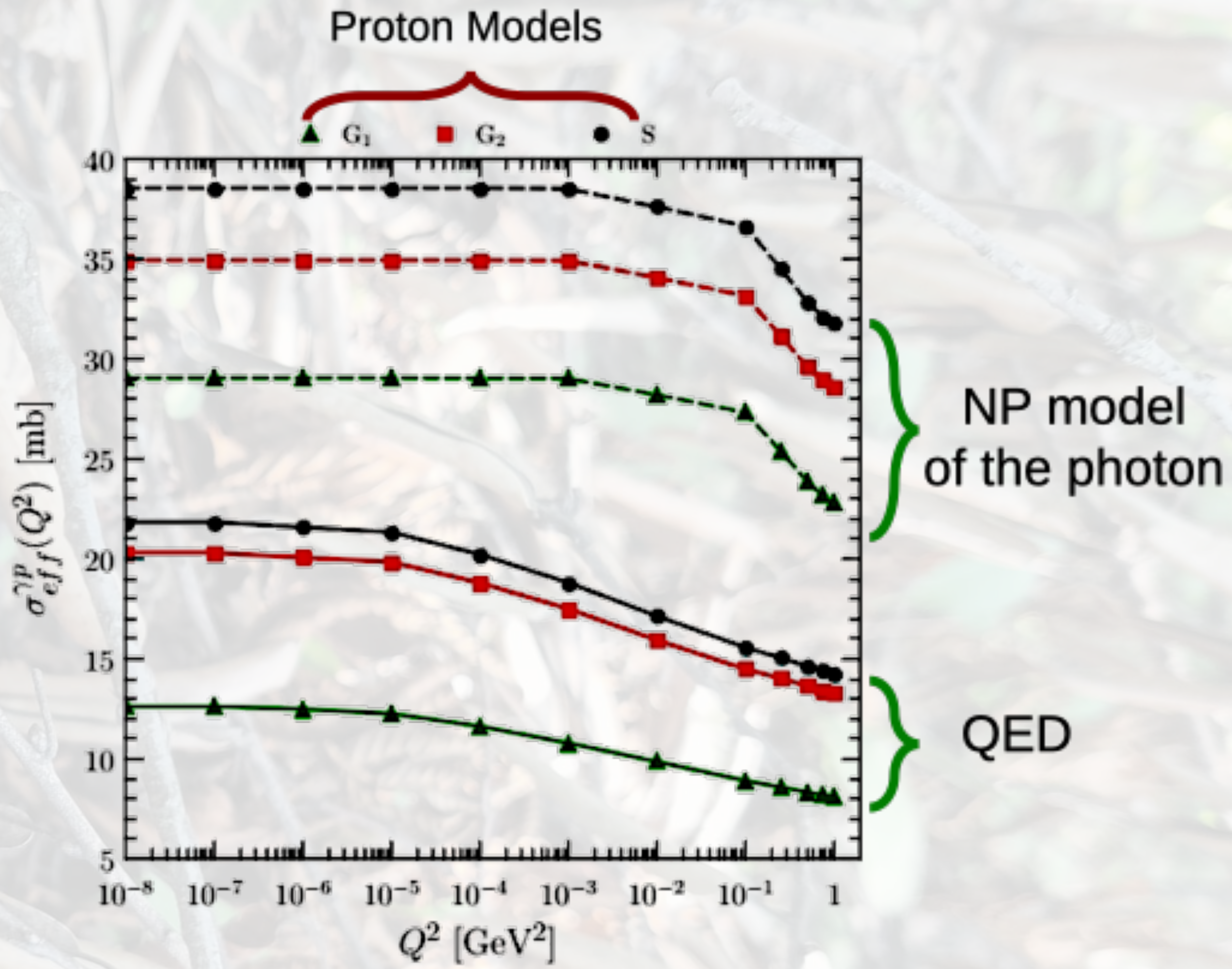


The $\gamma - p$ effective cross-section

M. R. and F. A. Ceccopieri, PRD 105 (2022) L011501



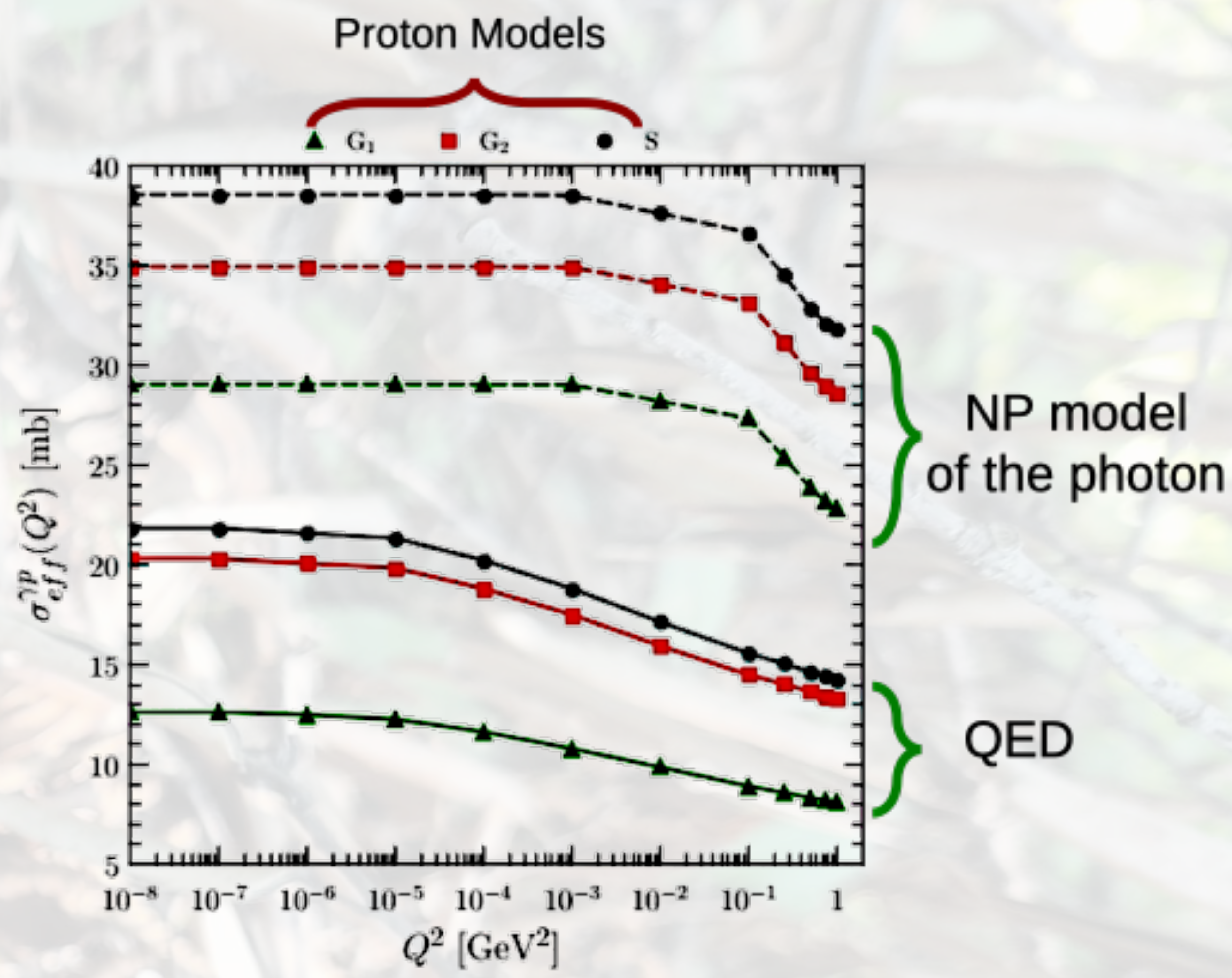
The effective cross-section depends on the photon virtuality! (NEW)



The 4-jets DPS cross-section

M. R. and F. A. Ceccopieri, PRD 105 (2022) L011501

$$d\sigma^{4j}_{DPS} = \frac{1}{2}\sum_{ab,cd}\int dy\, dQ^2\, \frac{f_{\gamma/e}(y,Q^2)}{\sigma^{\gamma p}_{eff}(Q^2)} \times$$
$$\times \int dx_{p_a} dx_{\gamma_b} f_{a/p}(x_{p_a}) f_{b/\gamma}(x_{\gamma_b}) d\hat{\sigma}^{2j}_{ab}(x_{p_a}, x_{\gamma_b})$$
$$\times \int dx_{p_c} dx_{\gamma_d} f_{c/p}(x_{p_c}) f_{d/\gamma}(x_{\gamma_d}) d\hat{\sigma}^{2j}_{cd}(x_{p_c}, x_{\gamma_d})$$



KINEMATICS:

$E_T^{jet} > 6$ GeV

$|\eta_{jet}| < 2.4$

$Q^2 < 1$ GeV^2

$0.2 \leqslant y \leqslant 0.85$

The ZEUS collaboration quoted an integrated total 4-jet cross section of 136 pb
S. Checkanov et al. (ZEUS), Nucl. Phys B792, 1 (2008)

The 4-jets DPS cross-section

M. R. and F. A. Ceccopieri, PRD 105 (2022) L011501

$$d\sigma^{4j}_{DPS} = \frac{1}{2}\sum_{ab,cd}\int \ldots$$

$$\times \int dx_{p_a} dx_{\gamma_b} f_{a/p}(\ldots$$

$$\times \int dx_{p_c} dx_{\gamma_d} f_{c/p}(\ldots$$

KINEMATICS:

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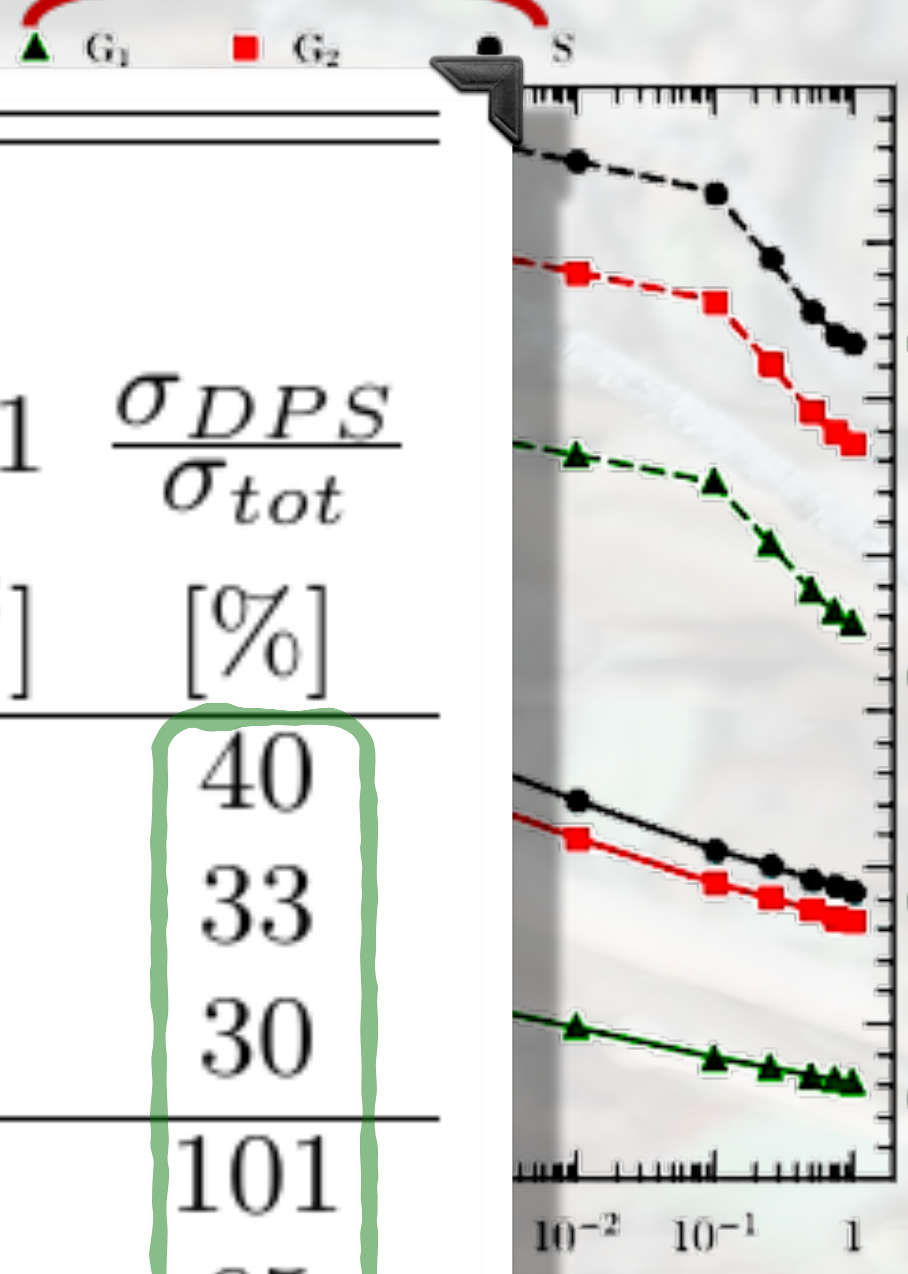
Photon	Proton	σ_{DPS} [pb] $Q^2 \leq 10^{-2}$ [GeV^2]	$10^{-2} \leq Q^2 \leq 1$ [GeV^2]	$Q^2 \leq 1$ [GeV^2]	$\frac{\sigma_{DPS}}{\sigma_{tot}}$ [%]
NP Model	G_1	35.1	18.6	53.7	40
	G_2	29.1	15.2	44.3	33
	S	26.4	13.7	40.1	30
QED	G_1	87.8	54.3	142.1	101
	G_2	54.3	33.4	87.7	65
	S	50.5	31.1	81.6	60



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S. Checkanov et al. (ZEUS), Nucl. Phys B792, 1 (2008)

The 4-jets DPS cross-section

M. R. and F. A. Ceccopieri, PRD 105 (2022) L011501

$$d\sigma^{4j}_{DPS} = \frac{1}{2}\sum_{ab,cd}\int \ldots$$

$$\times \int dx_{p_a} dx_{\gamma_b} f_{a/p}(\ldots$$

$$\times \int dx_{p_c} dx_{\gamma_d} f_{c/p}(\ldots$$

KINEMATICS:

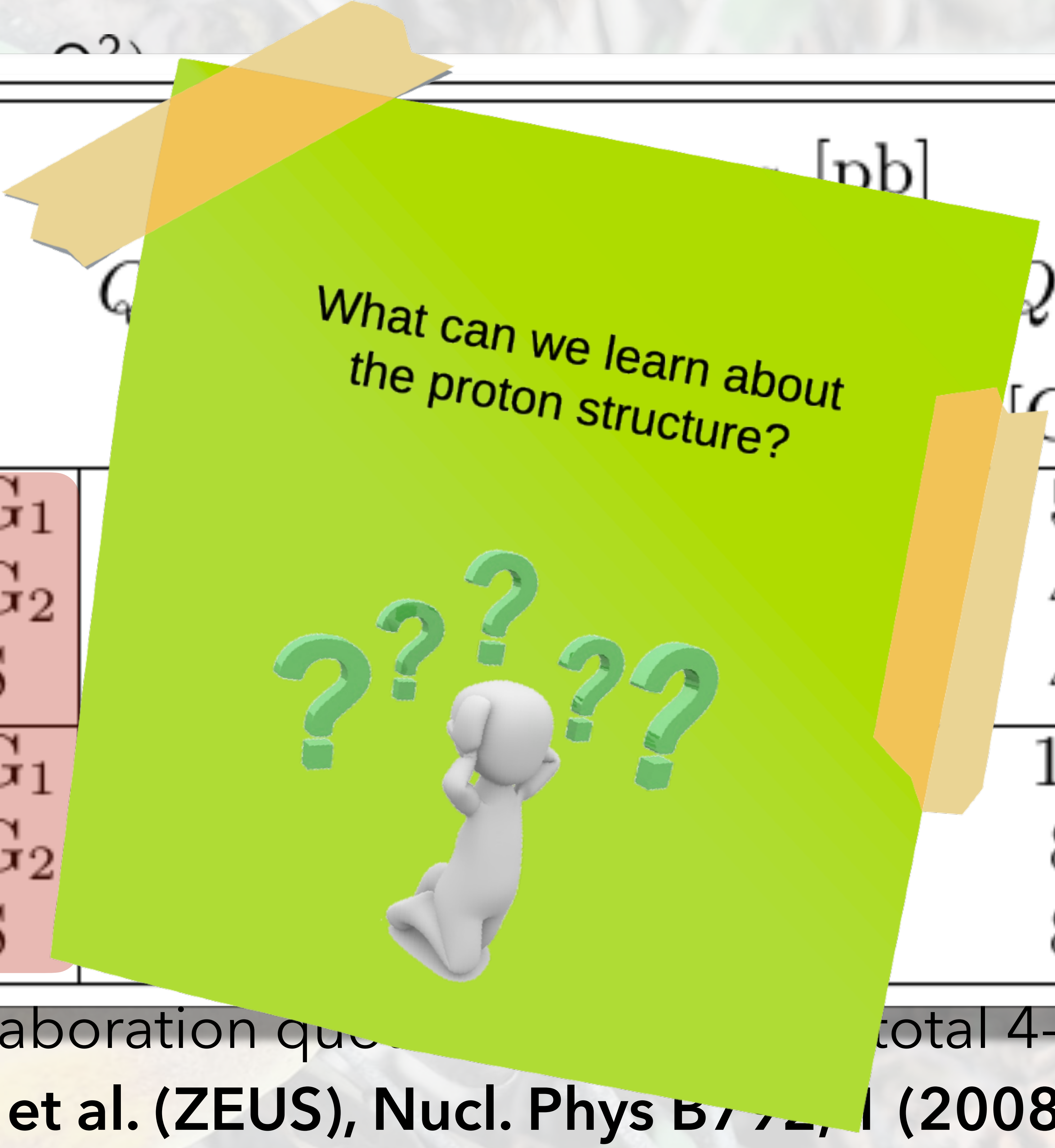
$E_T^{jet} > 6\ \mathrm{GeV}$

$|\eta_{jet}| < 2.4$

$Q^2 < 1\ \mathrm{GeV}^2$

$0.2 \leqslant y \leqslant 0.85$

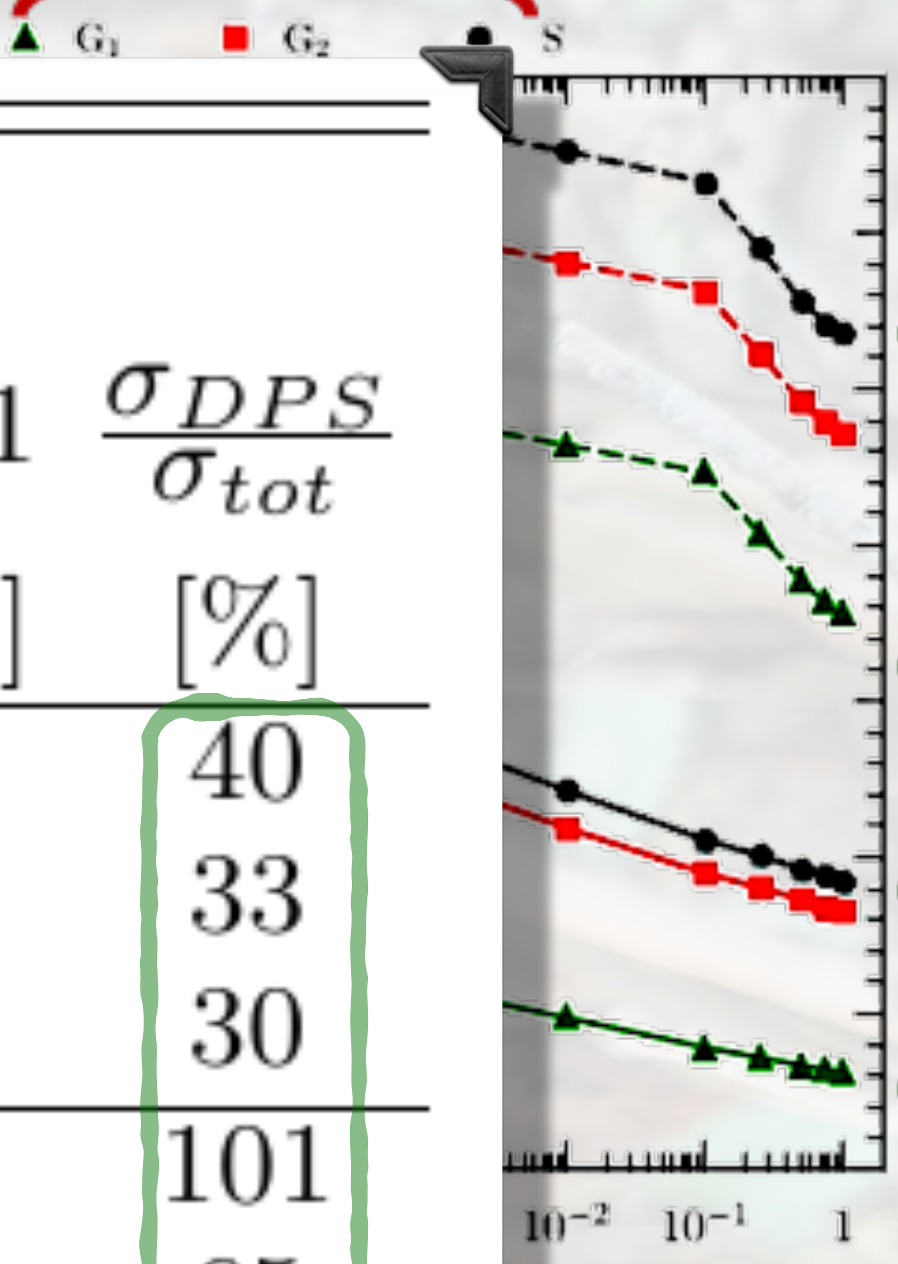
Proton / Photon		$Q^2 \leq 1$ [GeV^2]	$\frac{\sigma_{DPS}}{\sigma_{tot}}$ [%]
NP Model	G_1	53.7	40
	G_2	44.3	33
	S	40.1	30
QED	G_1	142.1	101
	G_2	87.7	65
	S	81.6	60



The ZEUS collaboration qu… total 4-jet cross section of 136 pb

S. Checkanov et al. (ZEUS), Nucl. Phys B… (2008)

Proton Models



NP model of the photon

QED

A key to the proton structure

M. R. and F. A. Ceccopieri, PRD 105 (2022) L011501

The effective cross section can be also written in terms of probability distribution:

$$\left[\sigma_{\rm eff}^{\gamma p}(Q^2)\right]^{-1} = \int d^2 z_\perp \; \tilde{F}_2^p(z_\perp)\tilde{F}_2^\gamma(z_\perp; Q^2)$$

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M. R. and F. A. Ceccopieri, PRD 105 (2022) L011501

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We can expand the distribution related to the photon:

$$\tilde{F}_2^\gamma(z_\perp; Q^2) = \sum_n \boxed{C_n(Q^2)} z_\perp^n$$

Coefficients determined in a given approach describing the photon structure

A key to the proton structure

M. R. and F. A. Ceccopieri, PRD 105 (2022) L011501

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$$\left[\sigma_{\text{eff}}^{\gamma p}(Q^2)\right]^{-1} = \sum_n C_n(Q^2) \langle z_\perp^n \rangle_p$$

Mean value of the transverse distance between two partons in the PROTON

A key to the proton structure

M. R. and F. A. Ceccopieri, PRD 105 (2022) L011501

The effective cross section can be also written in terms of probability distribution:

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Mean value of the transverse distance between two partons in the PROTON

If we could measure $\sigma_{\rm eff}^{\gamma p}(Q^2)$ we could access NEW INFORMATION ON THE PROTON STRUCTURE

A key to the proton structure

M. R. and F. A. Ceccopieri, PRD 105 (2022) L011501

The effective cross section can be also written in terms of probability distribution:

$$\left[\sigma_{eff}^{\gamma p}(Q^2)\right]^{-1} = \int d^2 z_\perp \; \tilde{F}_2^p(z_\perp)\tilde{F}_2^\gamma(z_\perp; Q^2)$$

We can exp... photon:

$\tilde{F}_2^\gamma(z_\perp$ ...

Coefficients determined in a given approach describing the photon structure

$\left[\sigma_{eff}^{\gamma p}(Q$ ... $z_\perp^n\rangle_p$

Mean value of the transverse distance between two partons in the PROTON

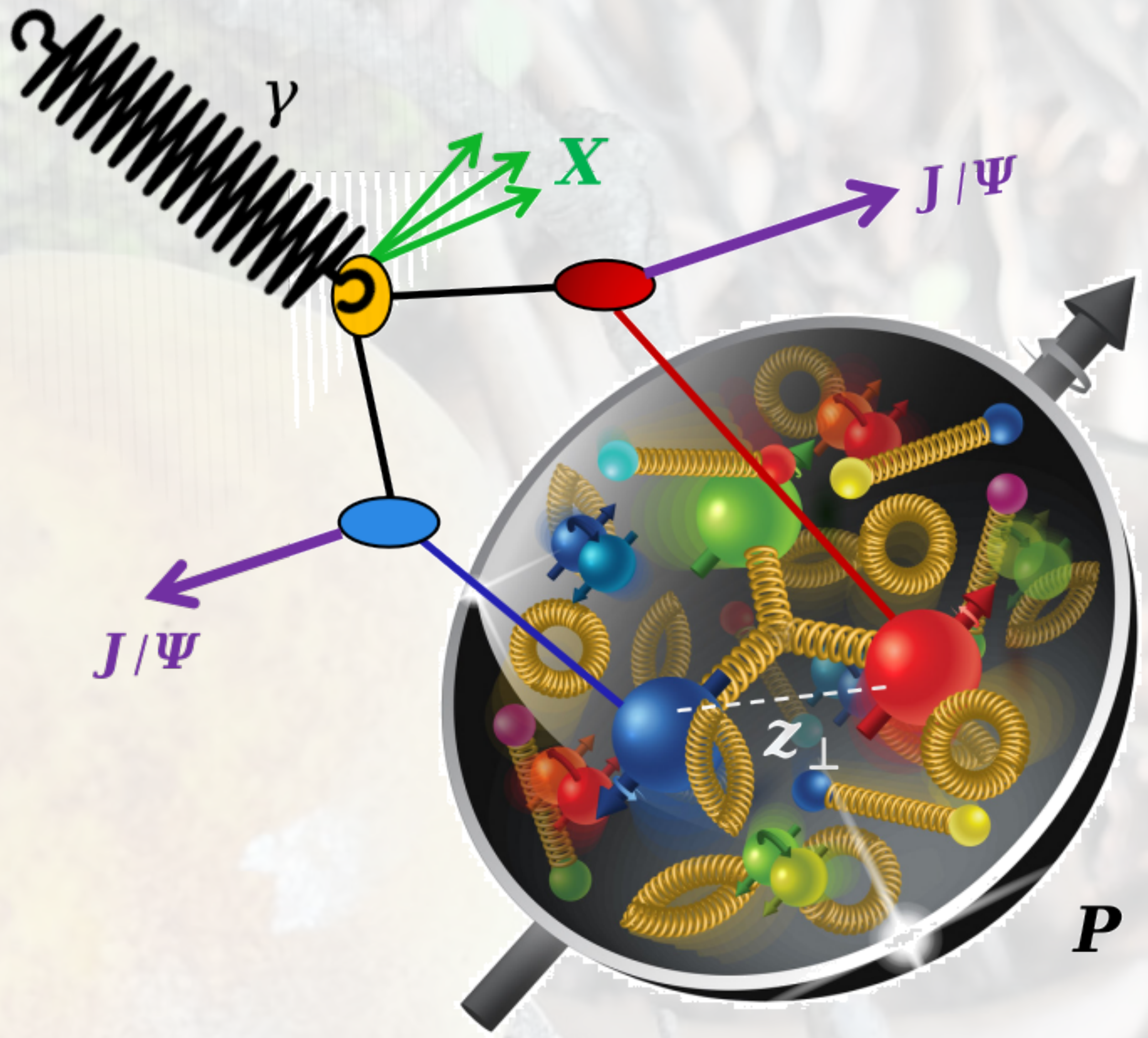
We estimated that with an integrated luminosity of 200 pb^{-1} Q^2 effects can be observed

If we could measure $\sigma_{eff}^{\gamma p}(Q^2)$ we could access NEW INFORMATION ON THE PROTON STRUCTURE

Di J/ψ photo-production@EIC

F. A. Ceccopieri, H. S. Shao, J. P. Lansberg, M. R. and R. Sangem in prep.

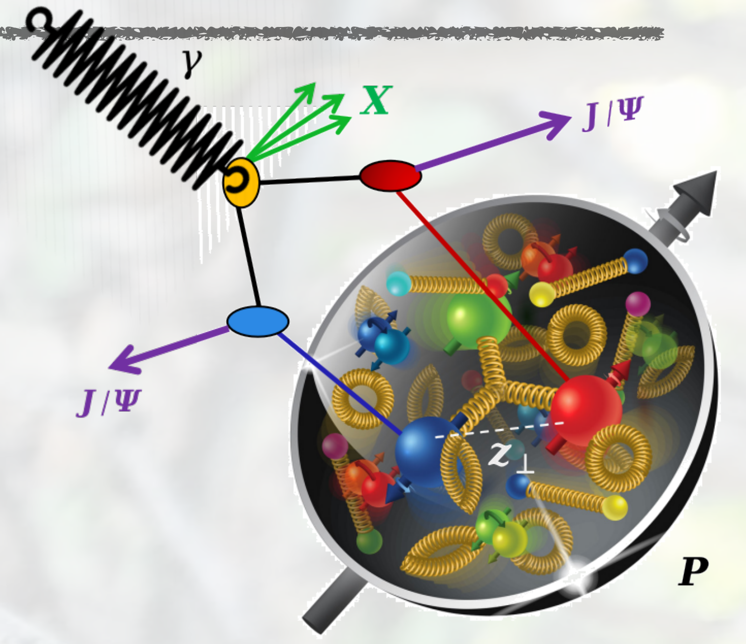
Illustration of DPS for $\gamma + p \to J/\psi + J/\psi + X$



We consider the possibility of **resolved** photon to estimate the DPS cross section in quarkonium-pair photoproduction at the EIC

Di J/ψ photo-production@EIC

F. A. Ceccopieri, H. S. Shao, J. P. Lansberg, M. R. and R. Sangem in prep.



*Slide from R. Sangem

$$\sigma_{SPS}^{(J/\psi,J/\psi)} \propto \sum_{a=g,q} \int dx_{p_a}\, f_{a/p}(x_{p_a},\mu)\, d\hat{\sigma}^{\gamma a \to J/\psi + J/\psi + a}$$

unresolved/direct

$$\sigma_{SPS}^{(J/\psi,J/\psi)} \propto \sum_{a,b=g,q} \int dx_{\gamma_a}\, dx_{p_b}\, f_{a/\gamma}(x_{\gamma_a},\mu)\, f_{b/p}(x_{p_b},\mu)\, d\hat{\sigma}^{ab \to J/\psi + J/\psi}$$

resolved

$$\sigma_{DPS}^{(J/\psi,J/\psi)} \propto \frac{1}{2}\frac{1}{\sigma_{eff}^{\gamma p}} \sum_{a,b,c,d} \int dx_{\gamma_a}\, dx_{p_b}\, f_{a/\gamma}(x_{\gamma_a},\mu)\, f_{b/p}(x_{p_b},\mu)\, d\hat{\sigma}_{SPS}^{ab\to J/\psi}(x_{\gamma_a},x_{p_b}) \times dx_{\gamma_c}\, dx_{p_d}\, f_{c/\gamma}(x_{\gamma_c},\mu)\, f_{d/p}(x_{p_d},\mu)\, d\hat{\sigma}_{SPS}^{cd\to J/\psi}(x_{\gamma_c},x_{p_d})$$

Proton PDF

Photon PDF

Partonic x-sections

Single SPS resolved (namely same partonic cross section as hadroproduction)

Di J/ψ photo-production@EIC

F. A. Ceccopieri, H. S. Shao, J. P. Lansberg, M. R. and R. Sangem in prep.

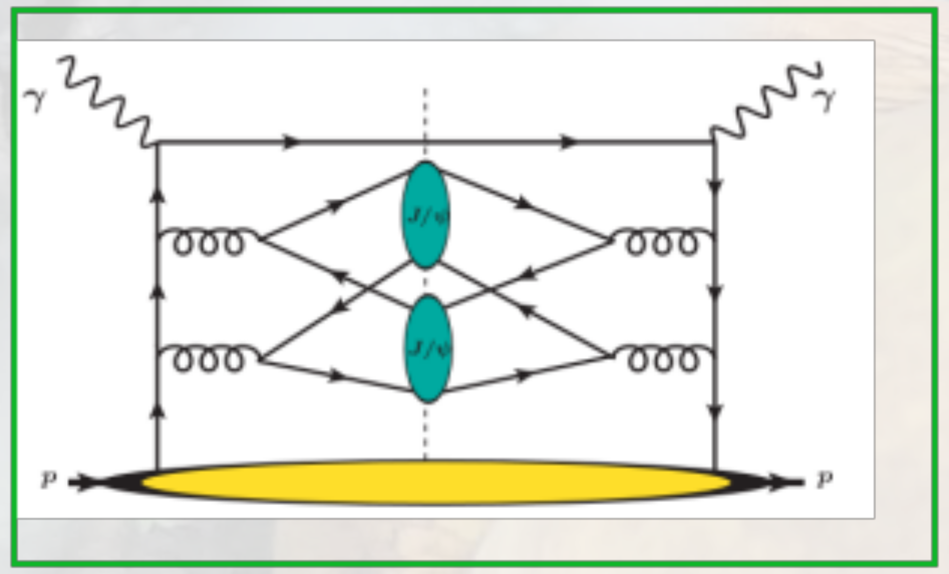
*Slide from R. Sangem

SPS

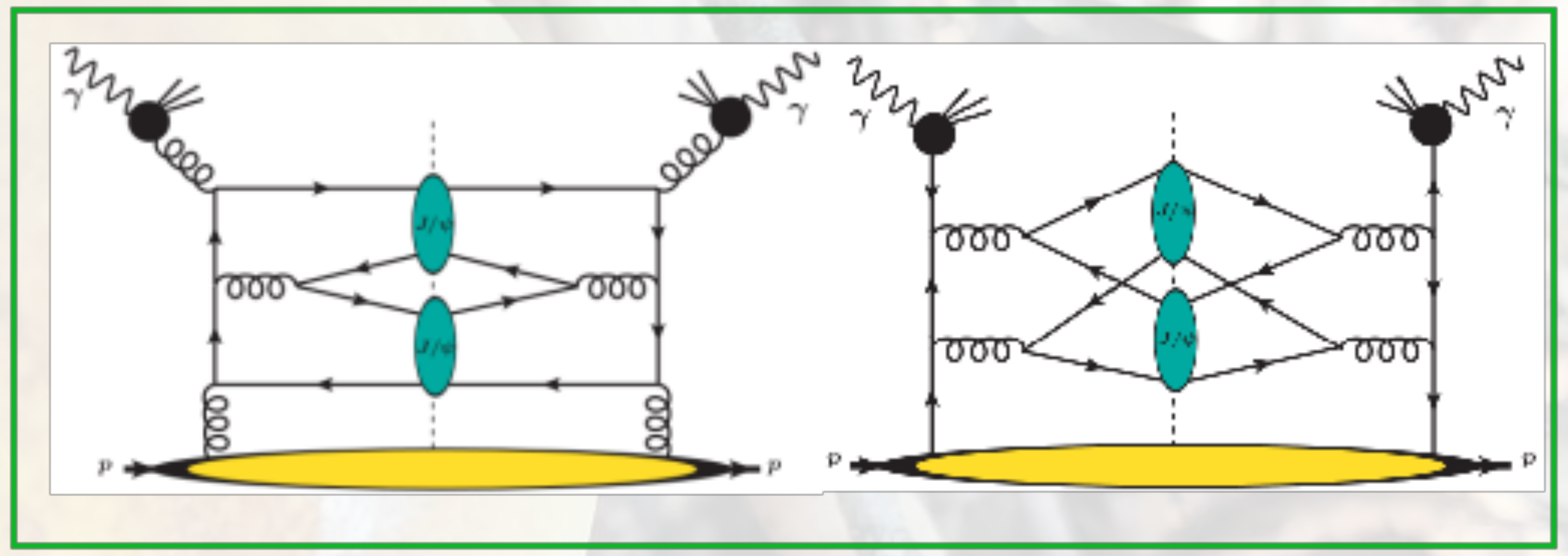
Unresolved | **Resolved**

DPS

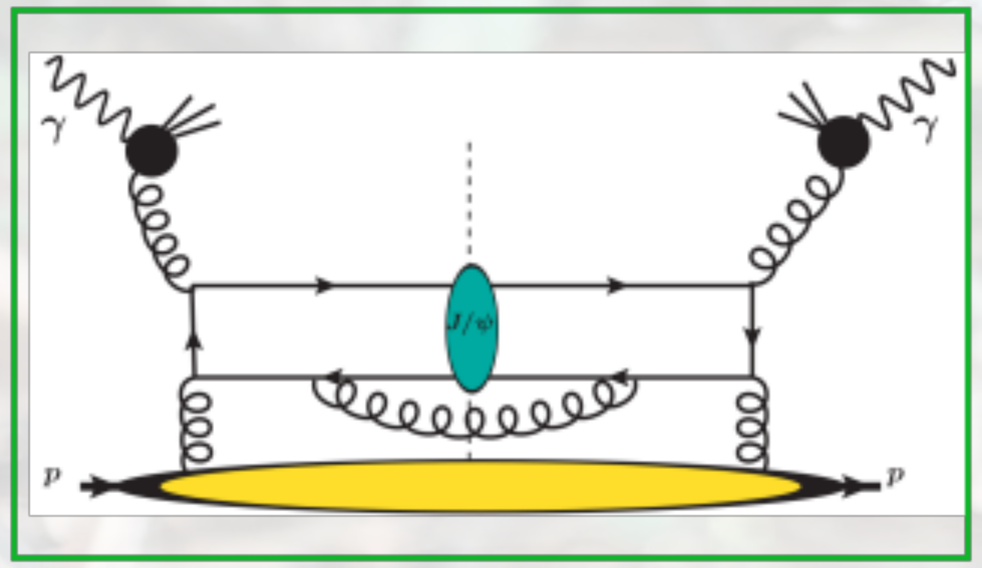
Resolved



$\gamma + q \to J/\psi + J/\psi + q$



$i + j \to J/\psi + J/\psi$ $i, j = g, q$



$|g + g \to J/\psi + g|^2$

- GRV photon PDF is used PRD 46, 1973 (1992) , while CT18NLO PDF for proton T.J. Hou et al., PRD 103, 014013 (2021)
- HELAC-Onia latest version is used for generating matrix elements HS Shao, CPC 184, 2562 (2013), 198, 238 (2016)
- CO LDMEs are taken from M. Butenschoen and B. A. Kniehl, PRD 84, 051501 (2011)
- We expect at least 600 four-muon events with 100 fb^{-1} luminosity

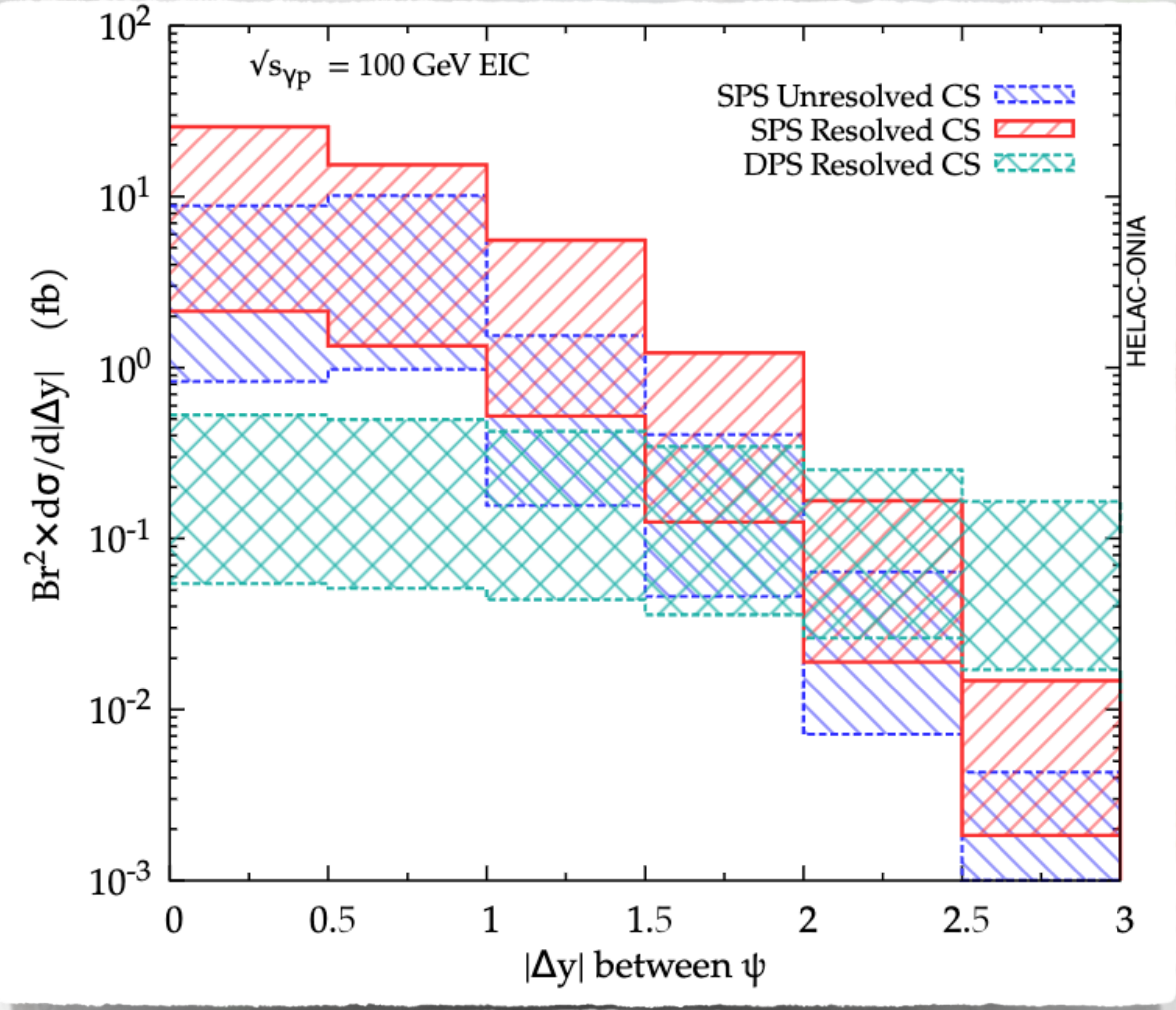
Numerical Results

PRELIMINARY

F. A. Ceccopieri, H. S. Shao, J. P. Lansberg, M. R. and R. Sangem in prep.

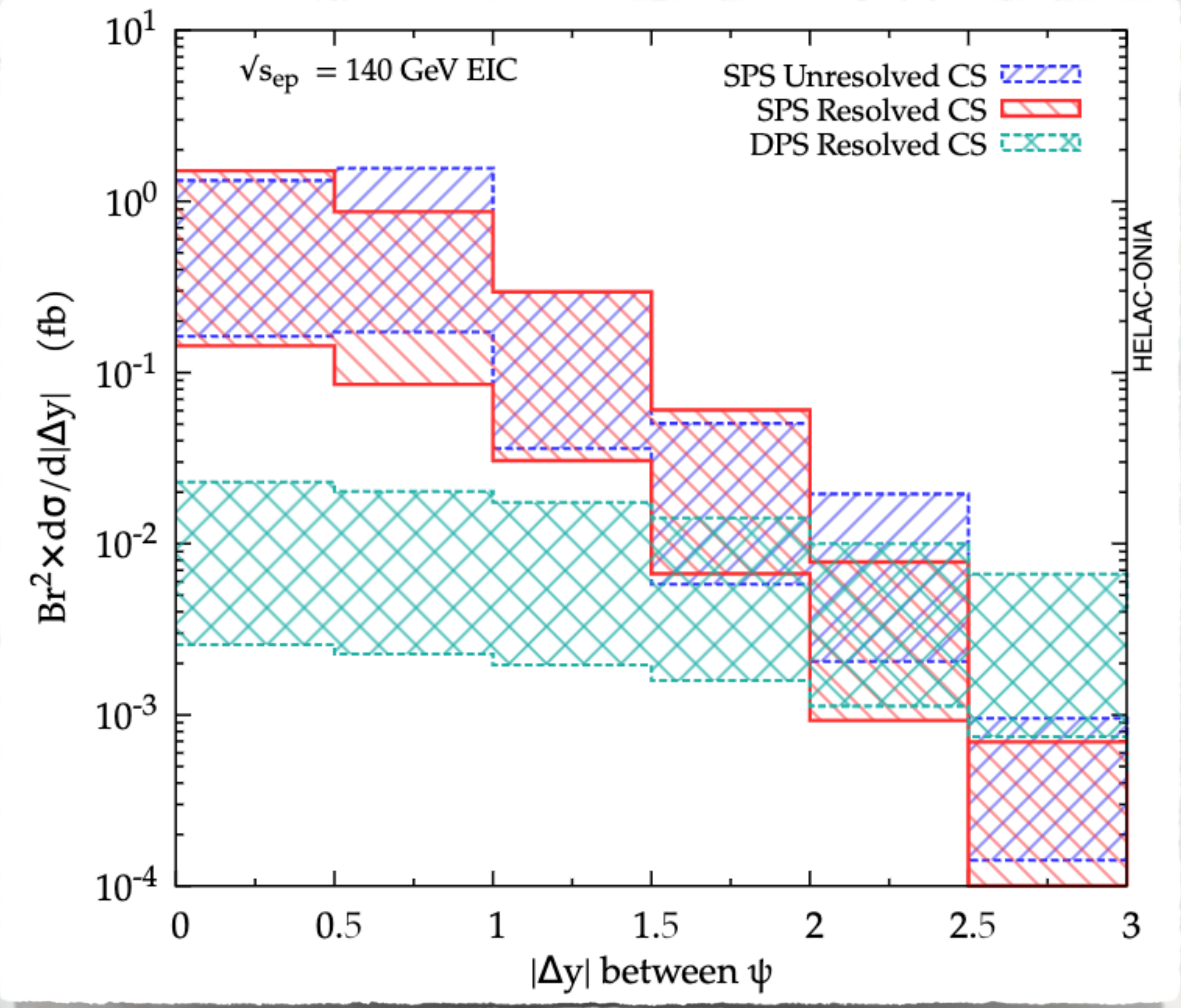
Absolute rapidity difference between the two J/ψ

$\sqrt{s_{\gamma p}} = 100$ GeV



- DPS dominates at high $|\Delta y|$
- DPS is suppressed at low $|\Delta y|$

$\sqrt{s_{\gamma p}} = 140$ GeV



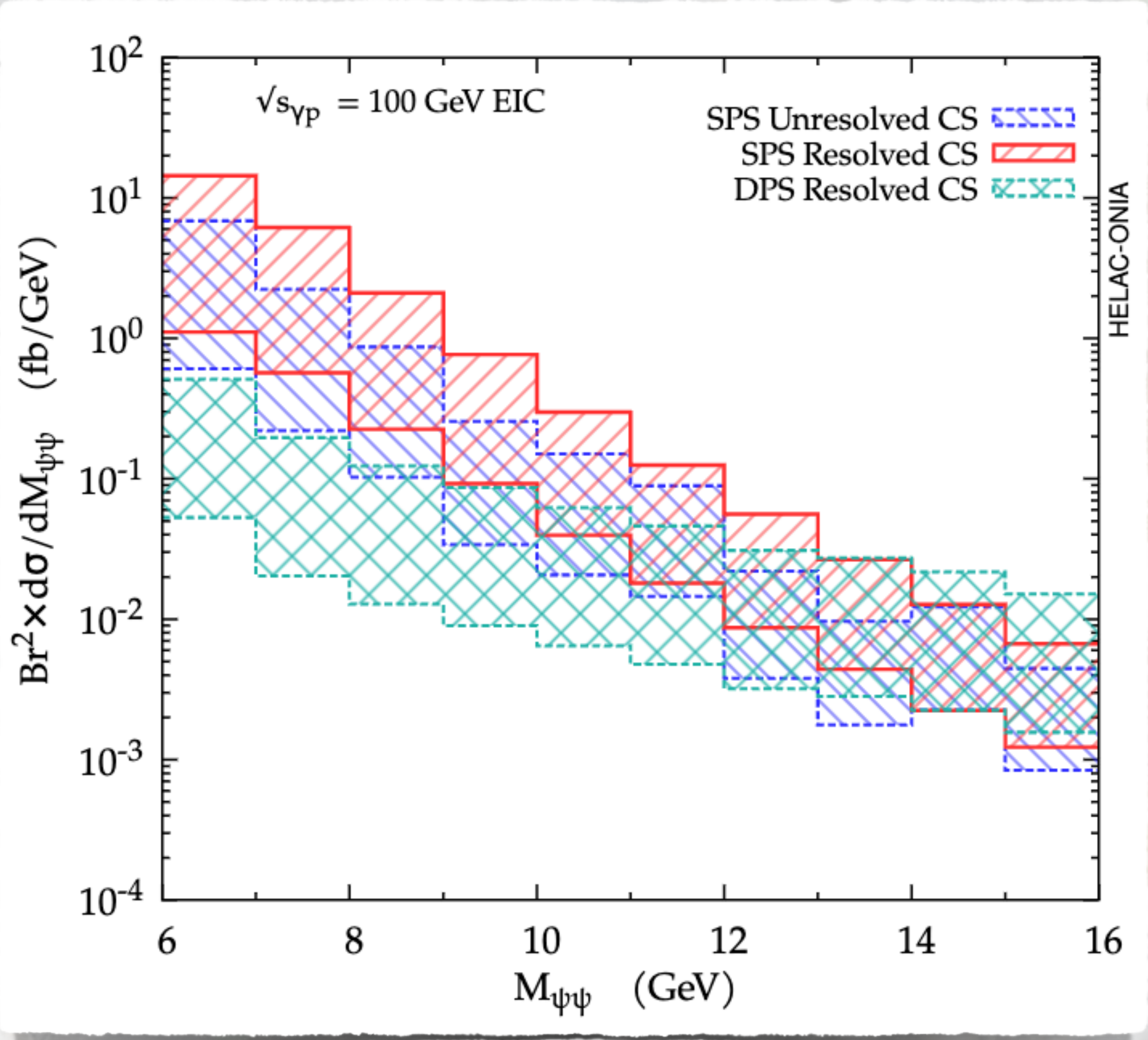
Numerical Results

PRELIMINARY

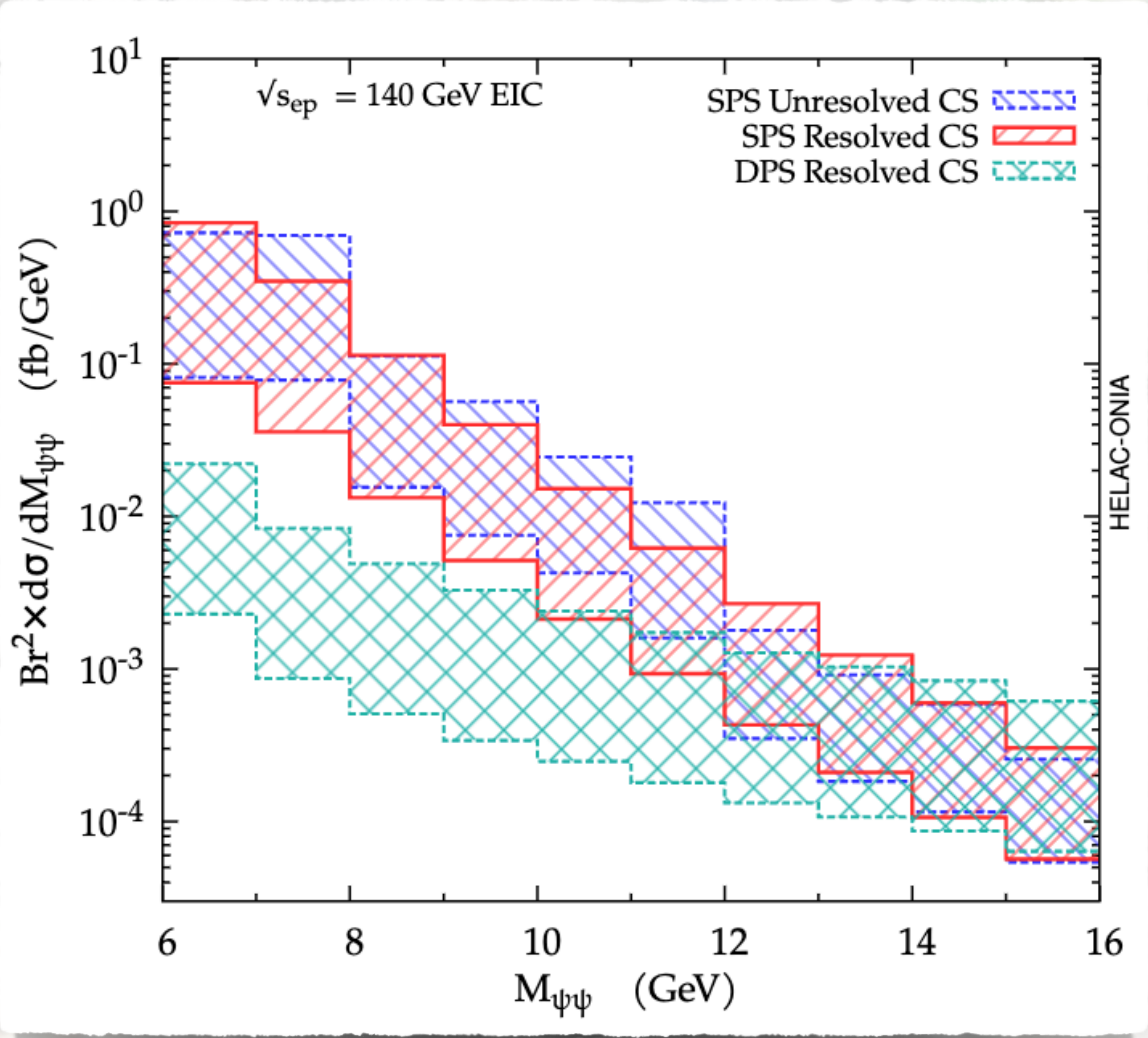
F. A. Ceccopieri, H. S. Shao, J. P. Lansberg, M. R. and R. Sangem in prep.

Invariant mass of the J/ψ pair

$\sqrt{s_{\gamma p}} = 100$ GeV



$\sqrt{s_{\gamma p}} = 140$ GeV



Numerical Results

PRELIMINARY

F. A. Ceccopieri, H. S. Shao, J. P. Lansberg, M. R. and R. Sangem in prep.

Invariant mass of the J/ψ pair

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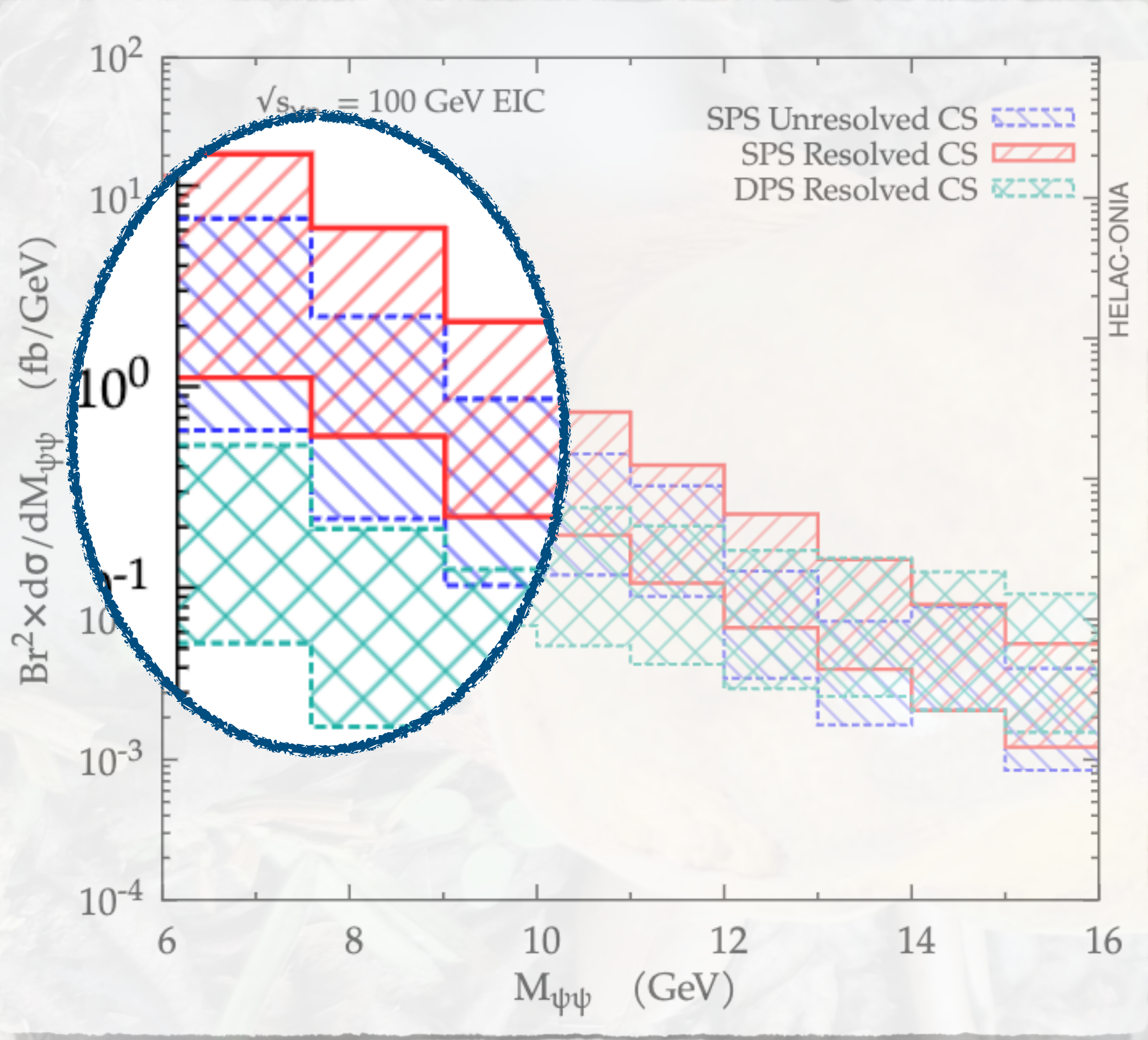
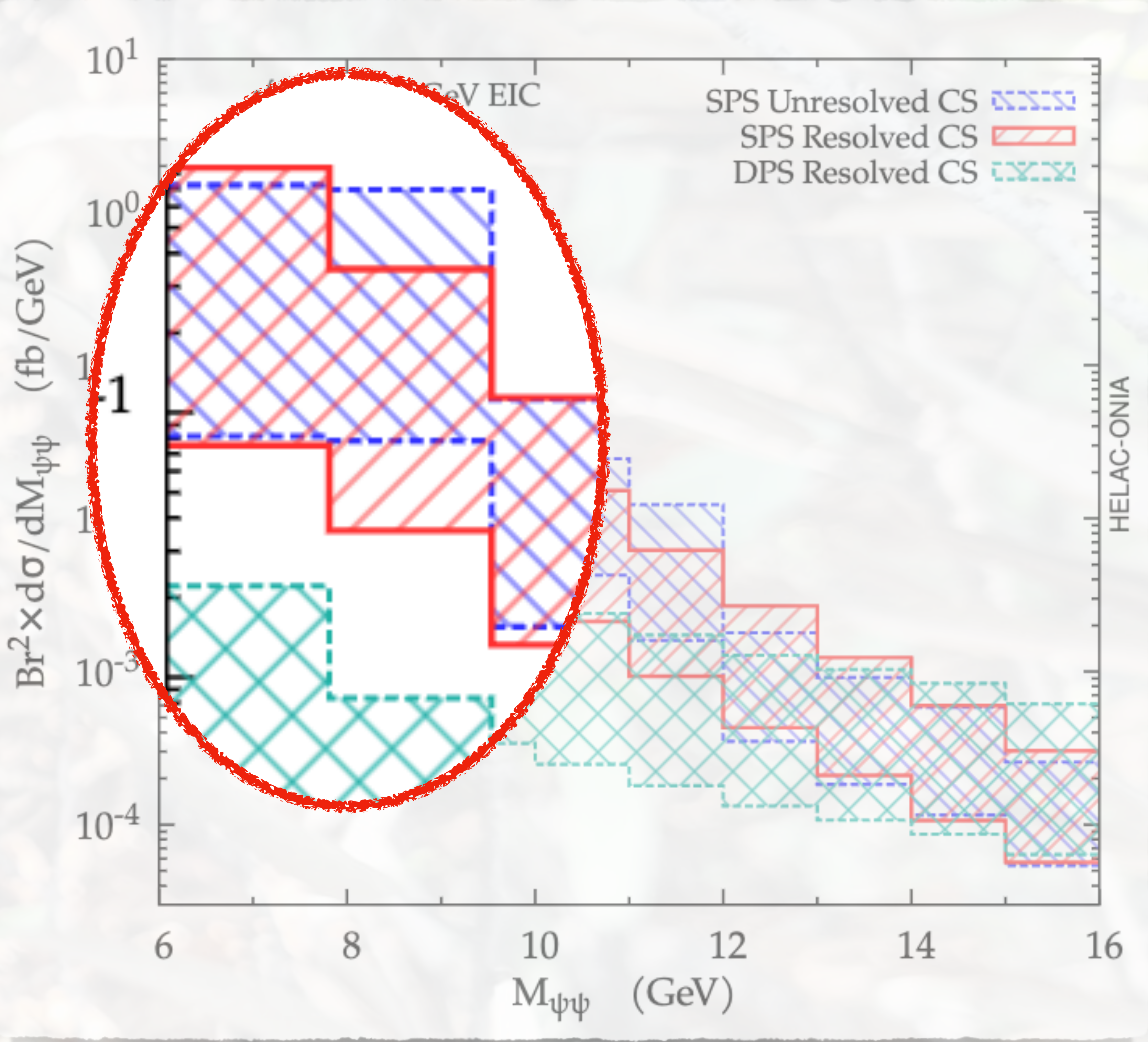


a) at low invariant mass:

- DPS smaller then SPS, but not negligible
- DPS negligible

$\sqrt{s_{\gamma p}} = 140$ GeV



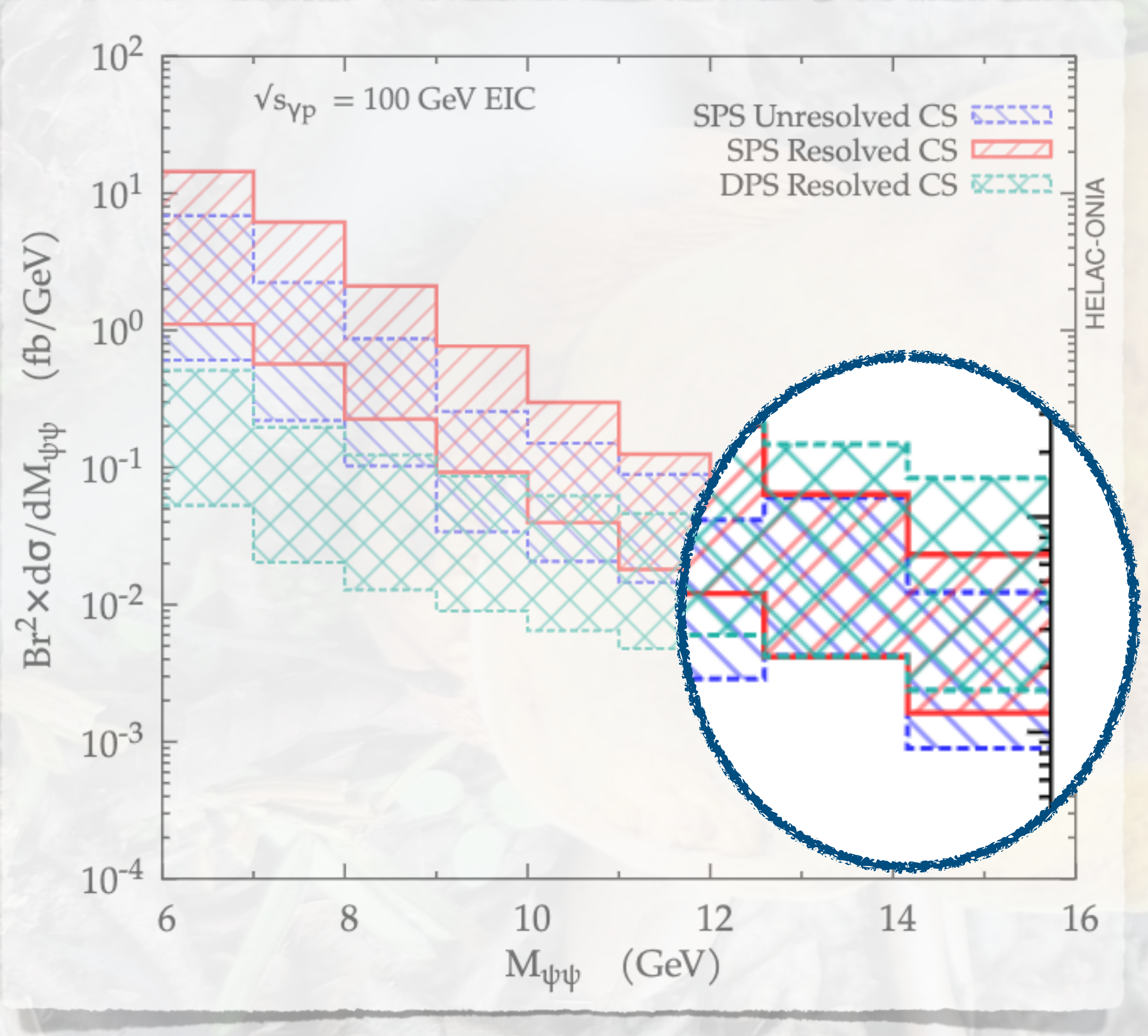
Numerical Results

PRELIMINARY

F. A. Ceccopieri, H. S. Shao, J. P. Lansberg, M. R. and R. Sangem in prep.

Invariant mass of the J/ψ pair

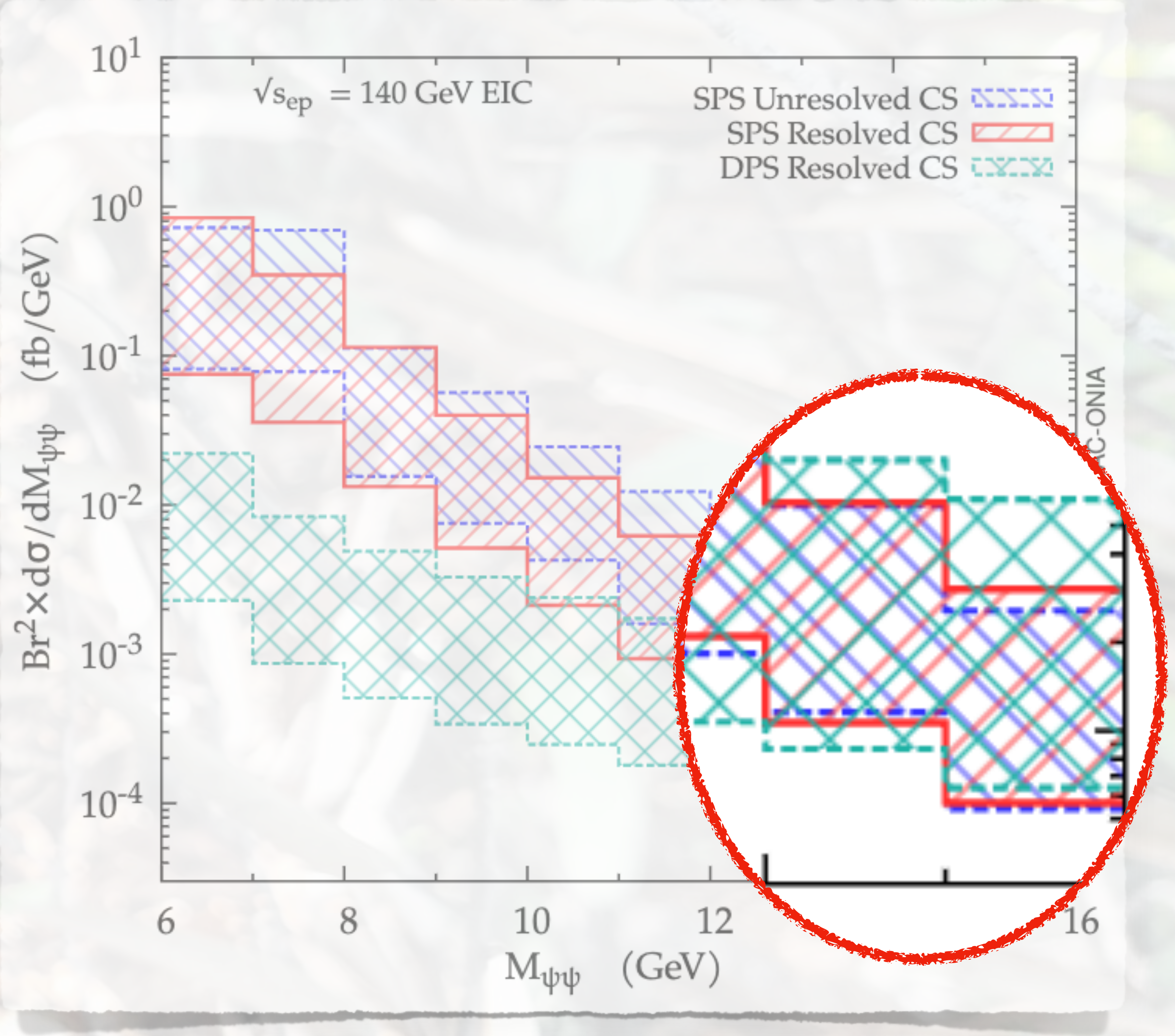
$\sqrt{s_{\gamma p}} = 100$ GeV

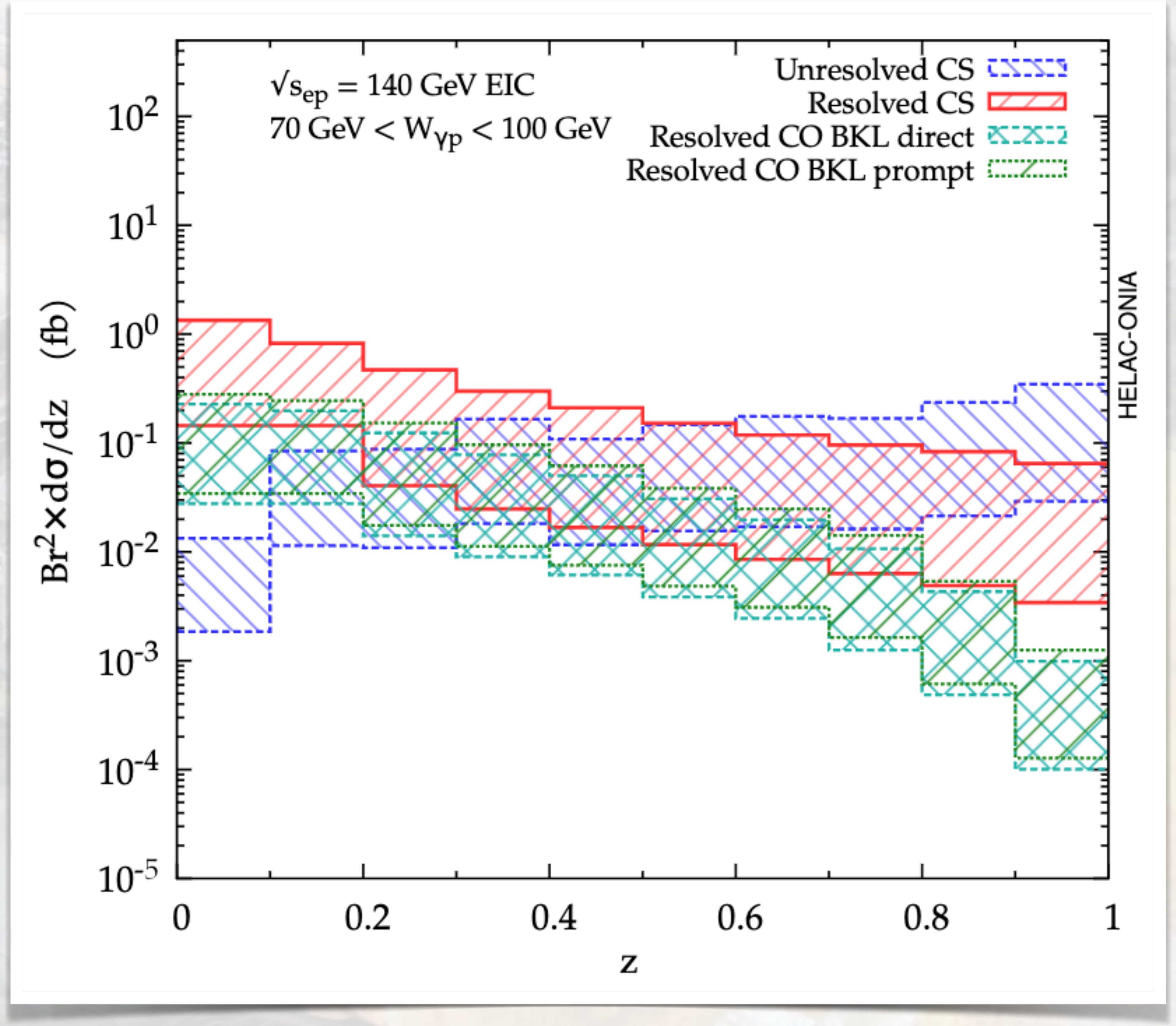


a) at low invariant mass:
- DPS smaller then SPS, but not negligible
- DPS negligible

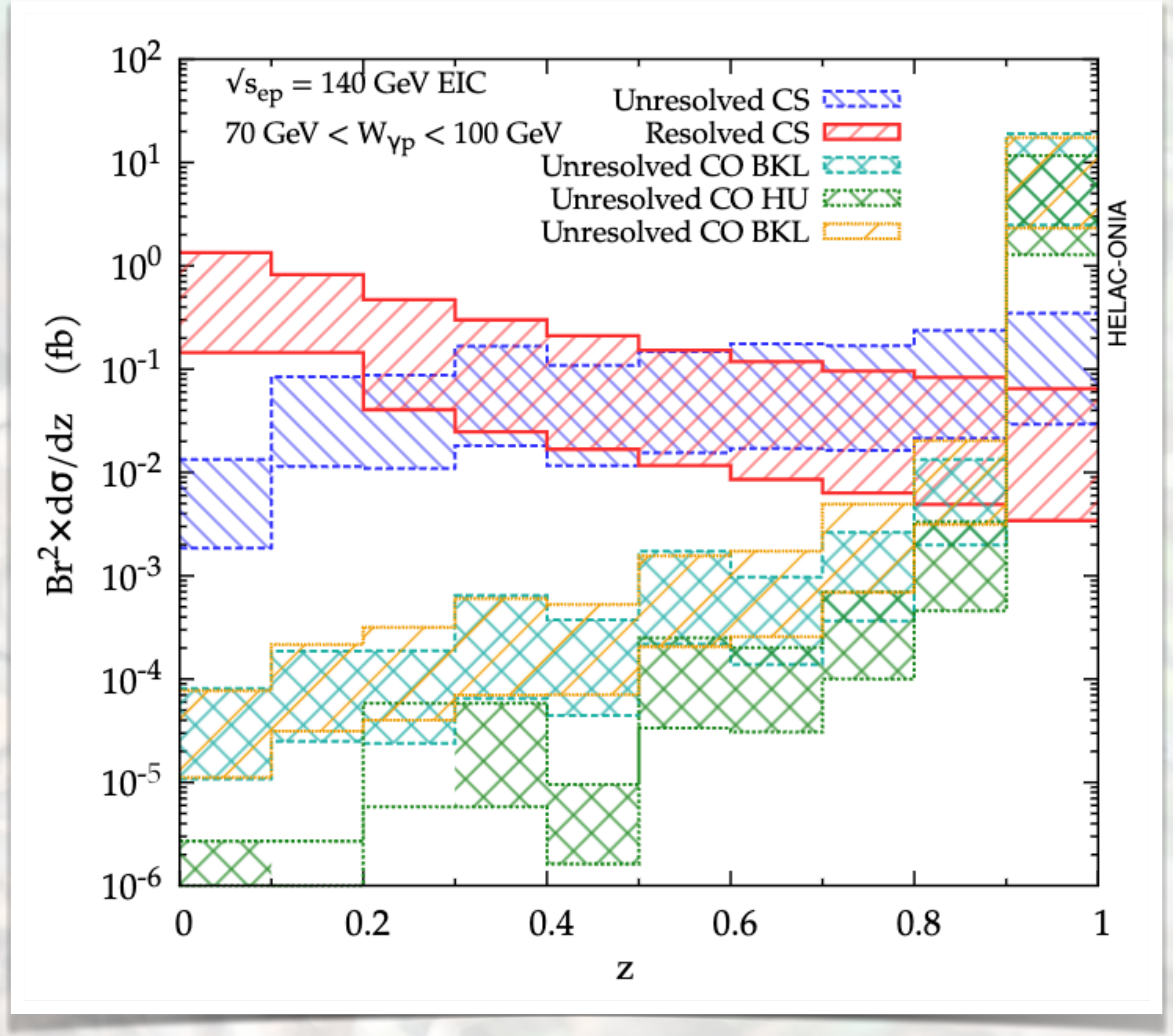
b) at low invariant mass:
- DPS bigger then SPS
- DPS similar to SPS

$\sqrt{s_{\gamma p}} = 140$ GeV





$$\mu_0 = \frac{H_T}{2} = \frac{\sum_i \sqrt{p_{Ti}^2 + m_i^2}}{2}$$



CO negligible if we cut z<0.9 (to be checked)

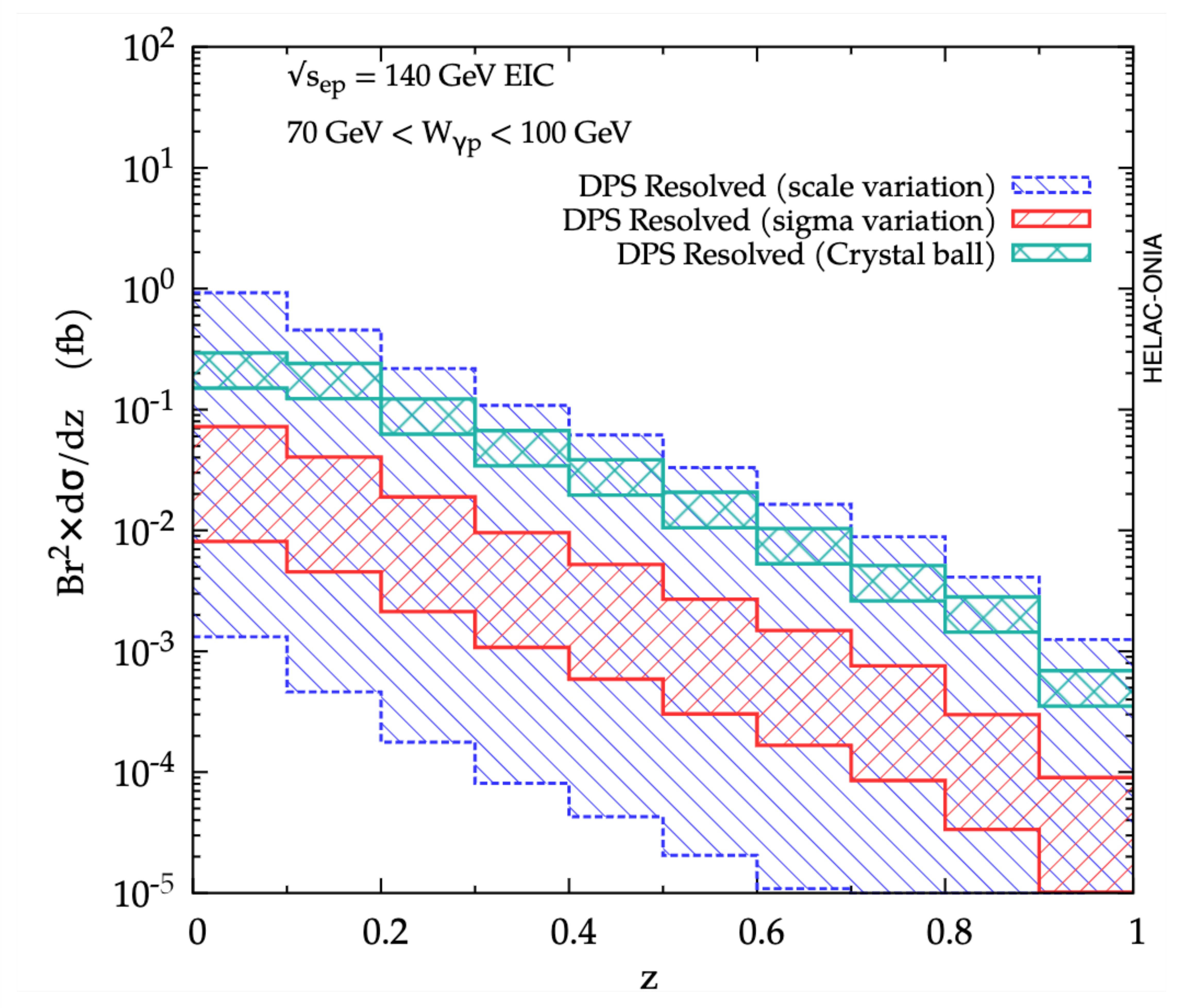
Numerical Results

PRELIMINARY

F. A. Ceccopieri, H. S. Shao, J. P. Lansberg, M. R. and R. Sangem in prep.

We also considered to use the Crystal Ball parametrization of the square of the amplitude $gg \longrightarrow \mathcal{Q} + X$

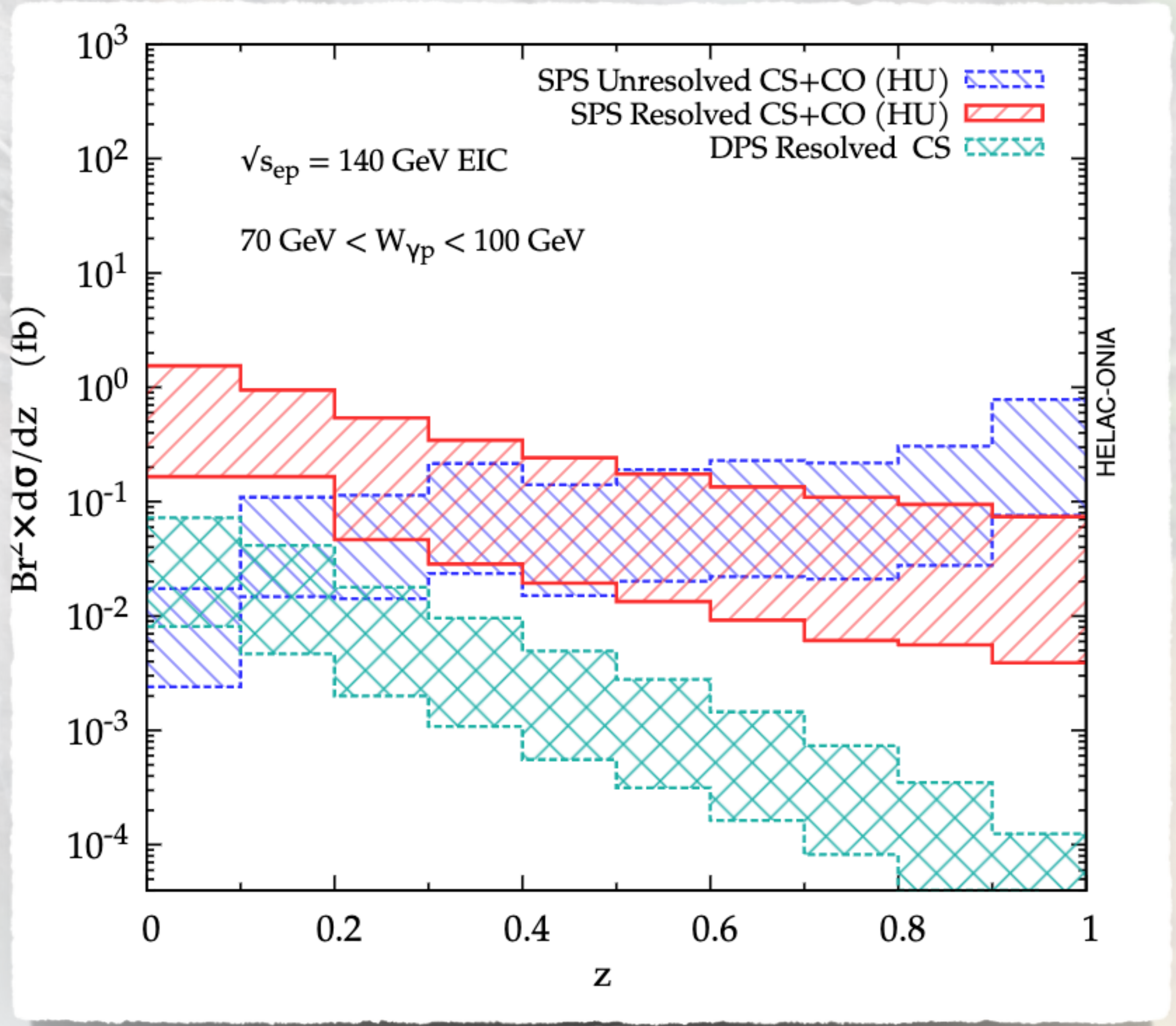
$$\begin{cases} \frac{\lambda^2 \kappa \hat{s}}{M_{\mathcal{Q}}^2} \exp(-\kappa \frac{P_T^2}{M_{\mathcal{Q}}^2}) & \text{when } P_T \leq \langle P_T \rangle \\ \frac{\lambda^2 \kappa \hat{s}}{M_{\mathcal{Q}}^2} \exp(-\kappa \frac{\langle P_T \rangle^2}{M_{\mathcal{Q}}^2}) \left(1 + \frac{\kappa}{n} \frac{P_T^2 - \langle P_T \rangle^2}{M_{\mathcal{Q}}^2}\right)^{-n} & \text{when } P_T > \langle P_T \rangle \end{cases}$$

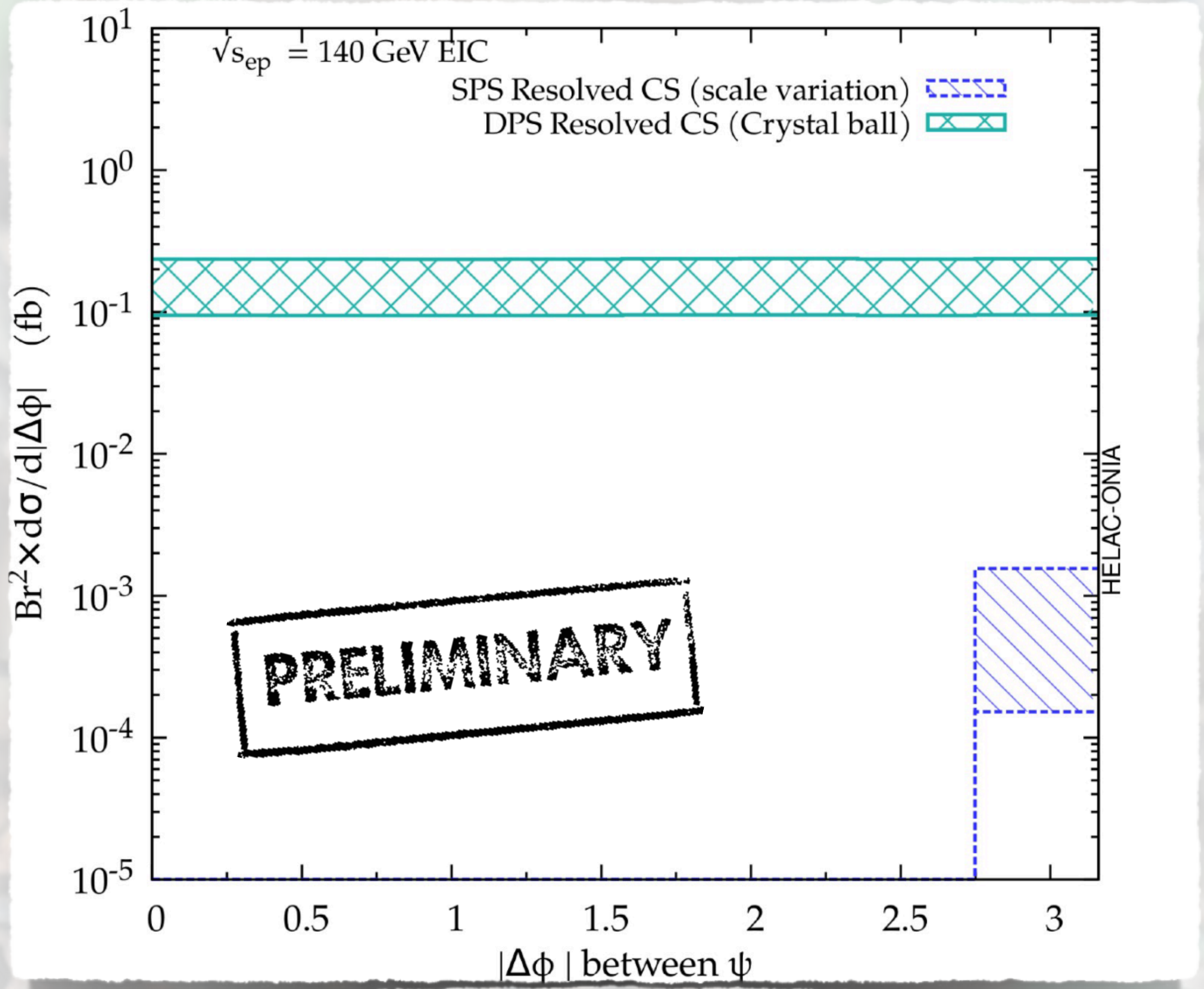


Numerical Results

PRELIMINARY

F. A. Ceccopieri, H. S. Shao, J. P. Lansberg, M. R. and R. Sangem in prep.



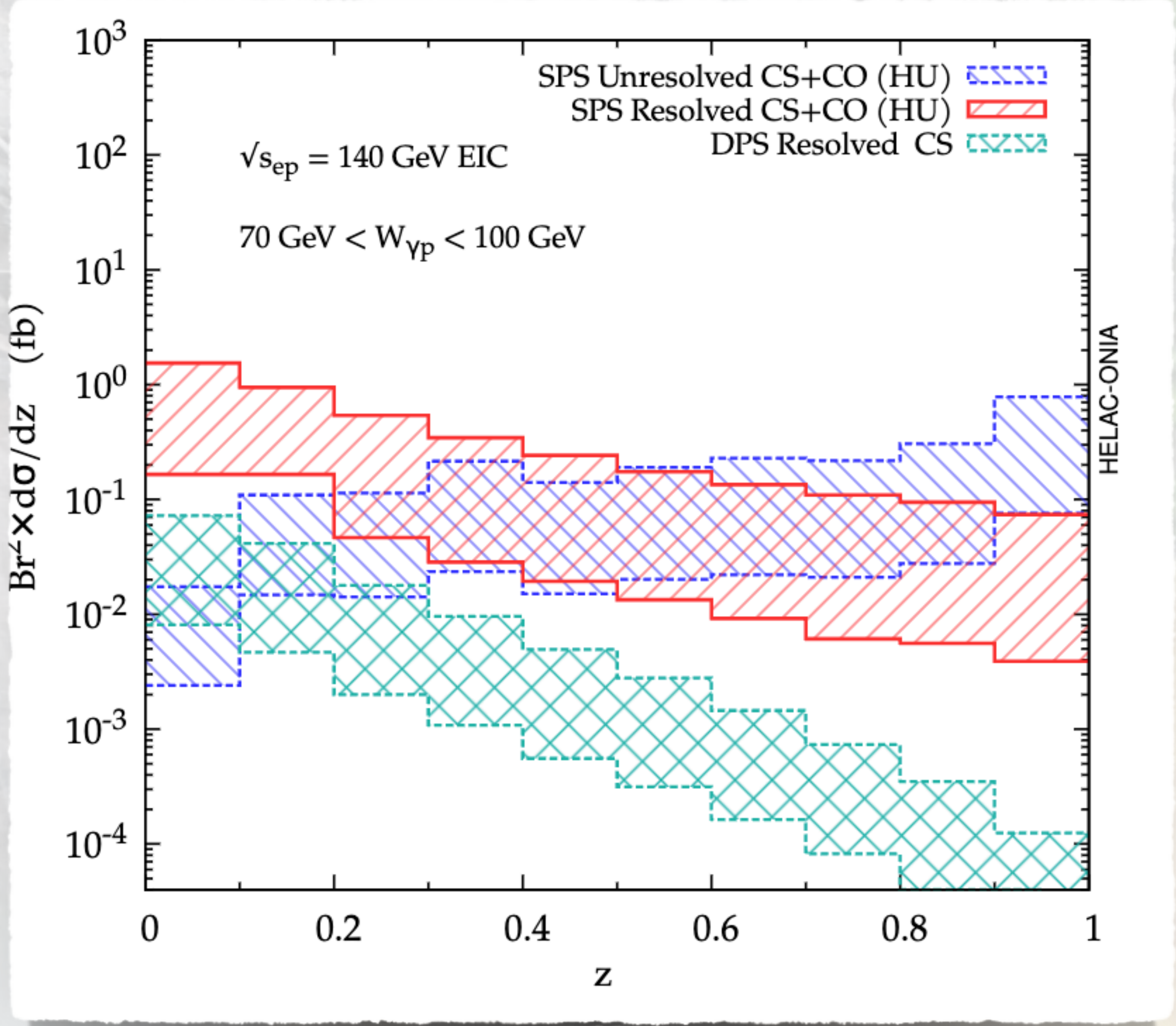


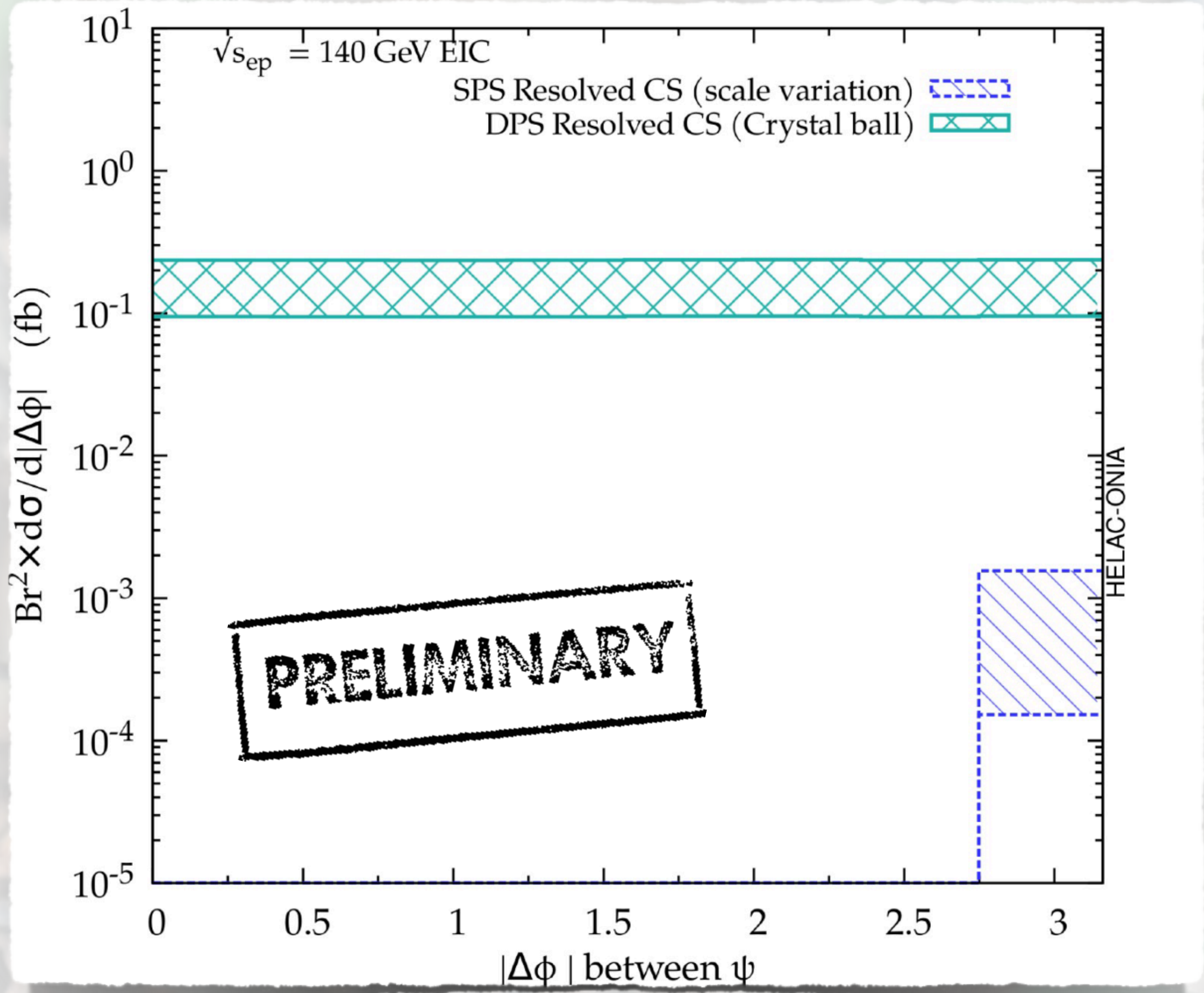
* for z<0.1, SPS resolved dominates → unique opportunity to investigate the PHOTON structure

Numerical Results

PRELIMINARY

F. A. Ceccopieri, H. S. Shao, J. P. Lansberg, M. R. and R. Sangem in prep.



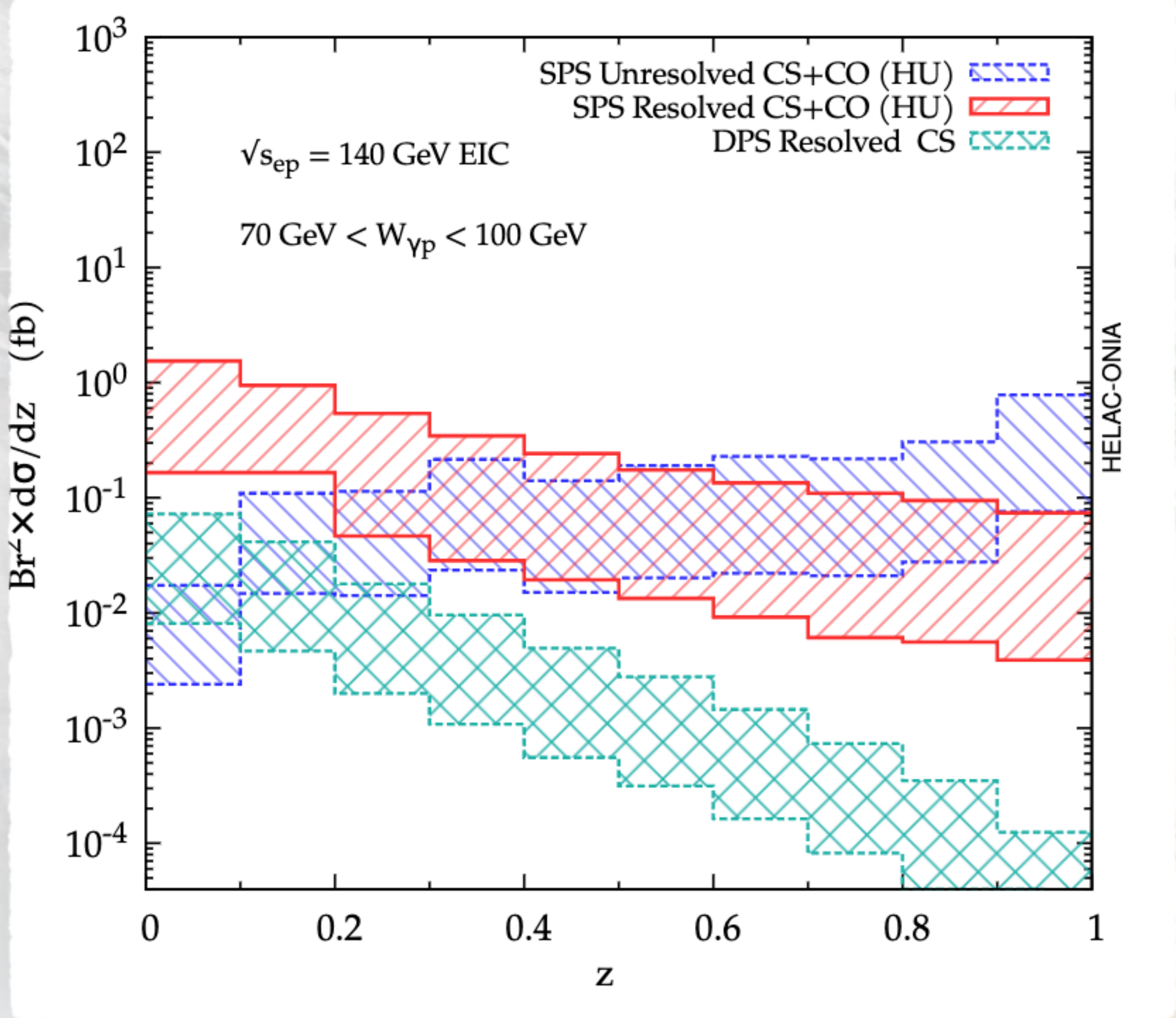


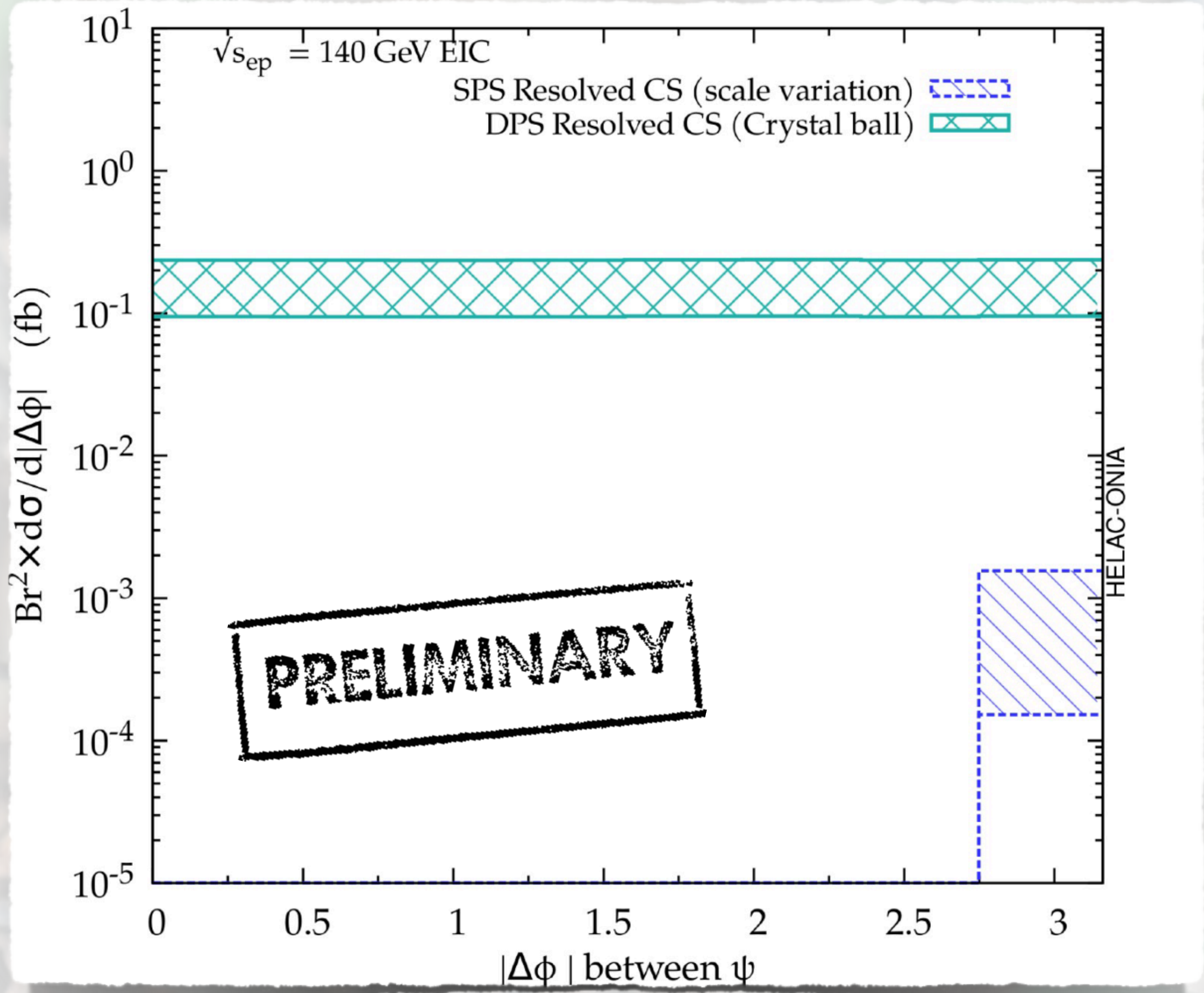
- for $z<0.1$, SPS resolved dominates → unique opportunity to investigate the PHOTON structure
- for high z, the direct SPS contribution dominates → we test the quarkonia production via direct photoproduction

Numerical Results

PRELIMINARY

F. A. Ceccopieri, H. S. Shao, J. P. Lansberg, M. R. and R. Sangem in prep.





- for $z<0.1$, SPS resolved dominates → unique opportunity to investigate the PHOTON structure
- for high z, the direct SPS contribution dominates → we test the quarkonia production via direct photoproduction
- as for DPS studies @LHC, the cross-section dependence on the relative azimuthal angle is relevant to access the DPS contribution!

Numerical Results

PRELIMINARY

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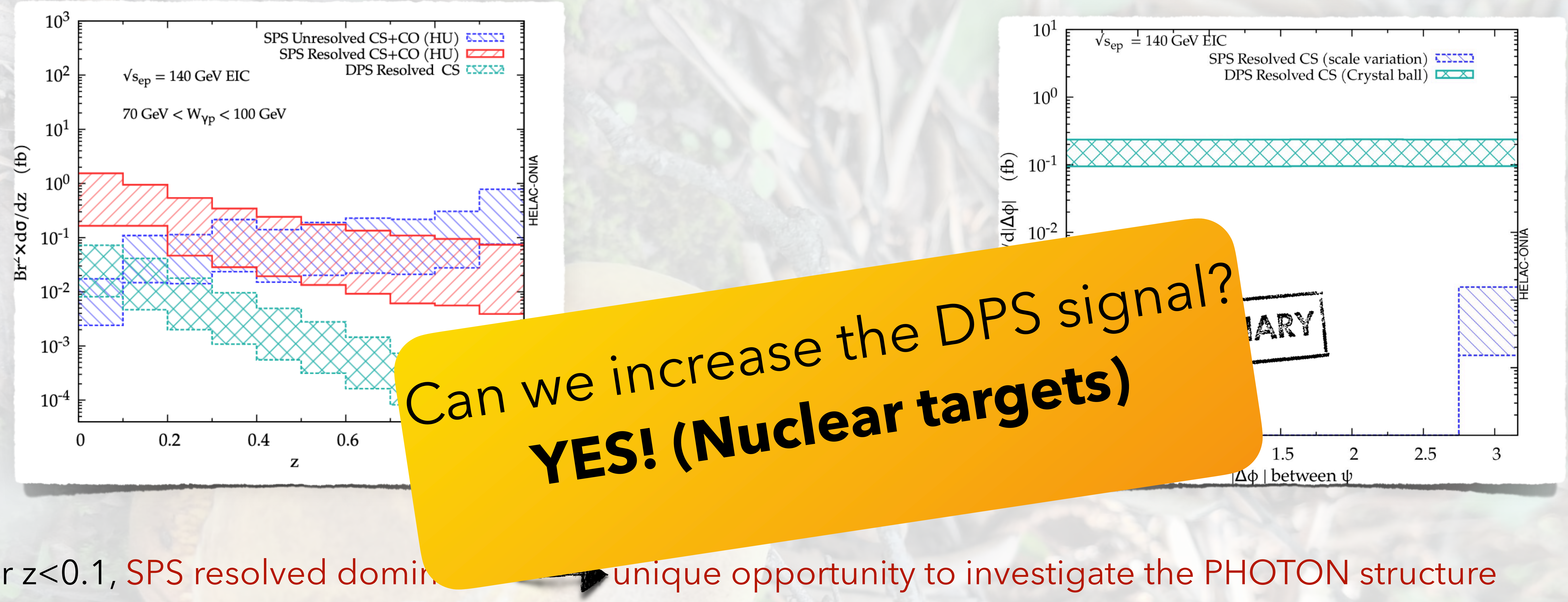

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- as for DPS studies @LHC, the cross-section dependence on the relative azimuthal angle is relevant to access the DPS contribution!

DPS in pA collisions

For DPS in pA and AA collisions the following references were missing:

1)Same-sign WW production in proton-nucleus collisions at the LHC as a signal for double parton scattering
D. d'E. & A. Snigirev, PLB 718 (2013) 1395-1400

2)Enhanced J/ΨJ/\PsiJ/Ψ-pair production from double parton scatterings in nucleus-nucleus collisions at the Large Hadron Collider
D. d'E. & A. Snigirev, PLB 727 (2013) 157-162

3)Pair production of quarkonia and electroweak bosons from double-parton scatterings in nuclear collisions at the LHC
D. d'E. & A. Snigirev, Nucl. Phys. A 931 (2014) 303-308

and for TPS:

Triple-parton scatterings in proton–nucleus collisions at high energies
D. d'E. & A. Snigirev, EPJC 78 (2018) 5, 359

DPS in pA collisions

$$F_{a_1a_2}(x_1, x_2, \mathbf{y}_\perp) = 2p^+ \int \frac{dz_1^-}{2\pi} \frac{dz_2^-}{2\pi} dy^- e^{i(x_1 z_1^- + x_2 z_2^-)p^+} \times \langle A | \mathcal{O}_{a_2}(0, z_2) \mathcal{O}_{a_1}(y, z_1) | A \rangle$$

In this case we have two mechanisms that contribute:

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In this case we have two mechanisms that contribute:

B. Blok et al, EPJC (2013) 73:2422

DPS 1: The two partons belong to the SAME nucleon in the nucleus!

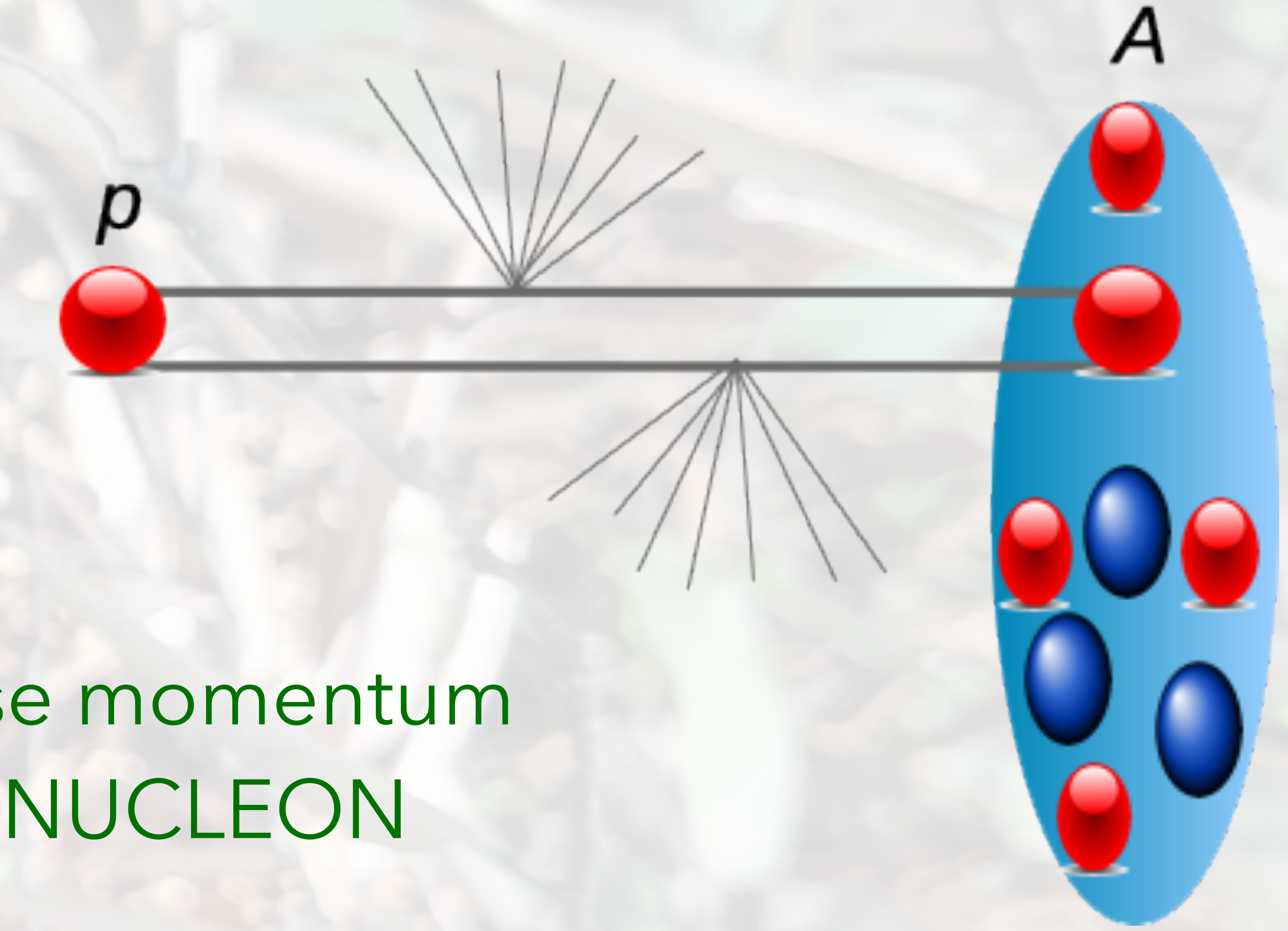
DPD of te nucleon inside the nucleus

$$\tilde{F}^1_{a_1a_2}(x_1,x_2,k_\perp) = \sum_{N=p,n}\int \frac{1}{\xi}\tilde{F}^N_{a_1a_2}\left(\frac{x_1}{\xi},\frac{x_2}{\xi},k_\perp\right)\rho^N_A(\xi,p_{t,N})\frac{d\xi}{\xi}\,d^2p_{t,N}$$

Momentum fraction carried by a NUCLEON

Light-Cone Momentum Distribution

Transverse momentum of the NUCLEON



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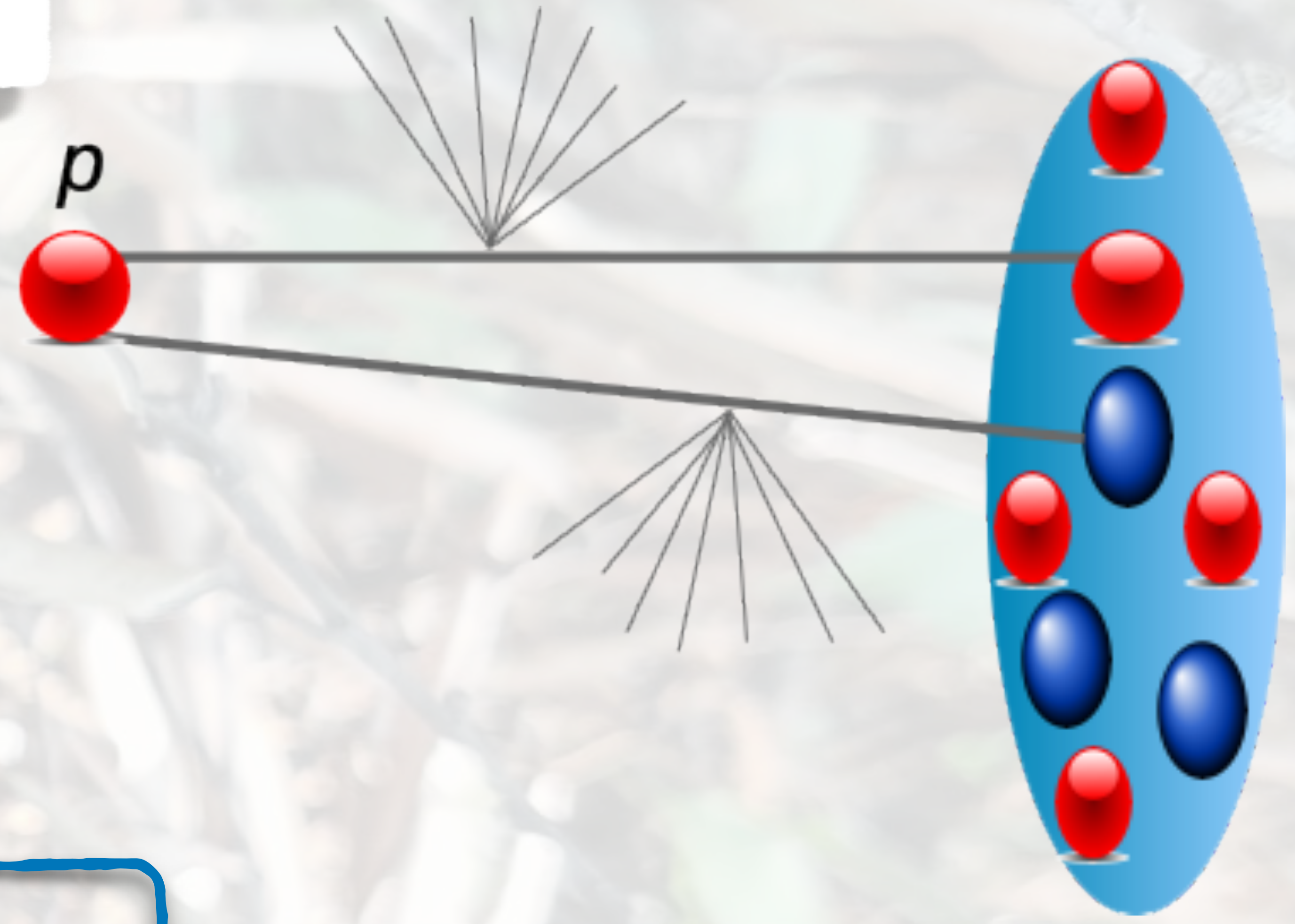
B. Blok et al, EPJC (2013) 73:2422

DPS 2: The two partons belong to the DIFFERENT nucleons in the nucleus!



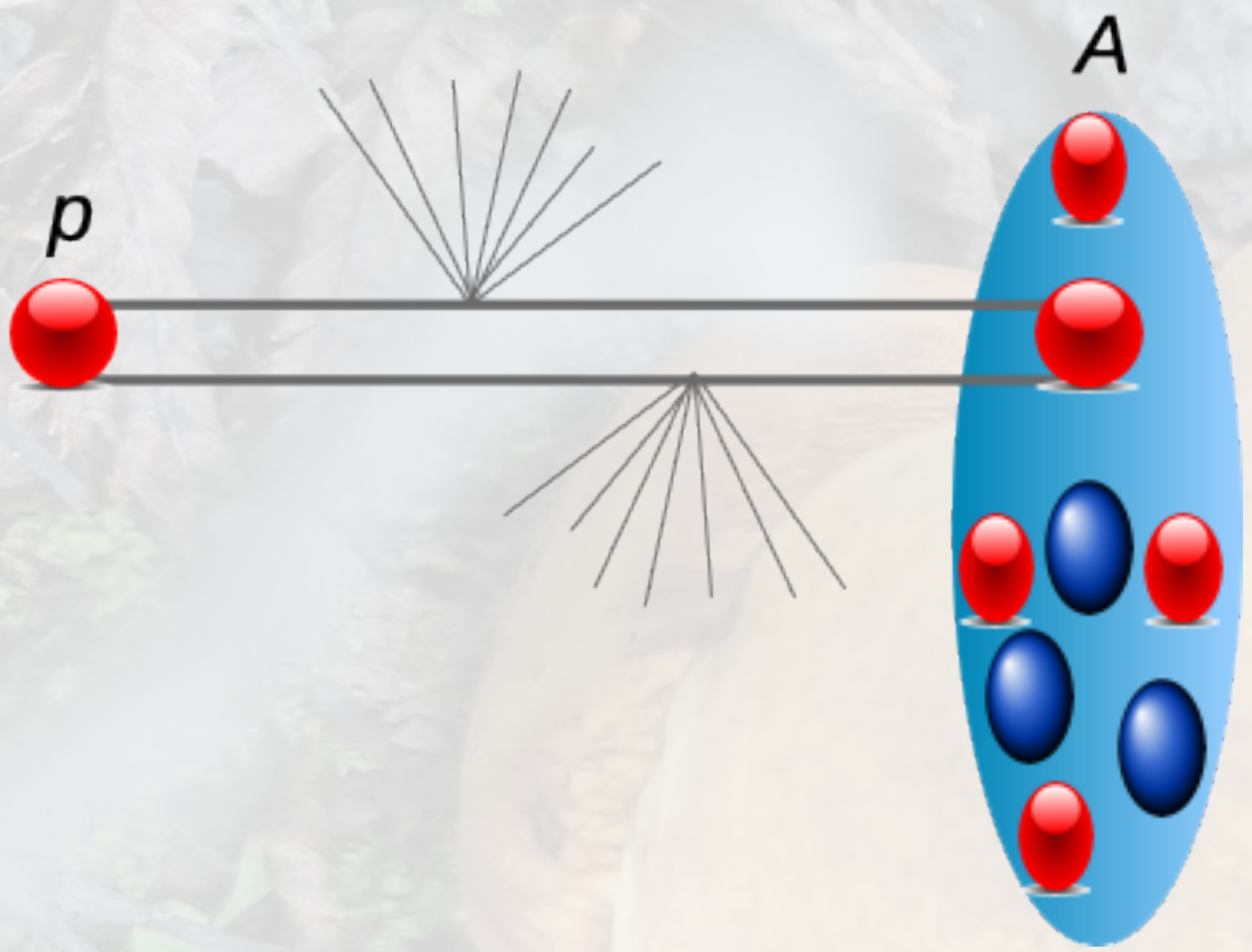
$$\tilde{F}^2_{a_1a_2}(x_1, x_2, \vec{k}_\perp) \propto \int \frac{1}{\xi_1\xi_2} \prod_{i=1}^{i=A} \frac{d\xi_i d^2 p_{ti}}{\xi_i} \delta\left(\sum_i \xi_i - A\right) \delta^{(2)}\left(\sum_i \mathbf{p}_{ti}\right) \psi_A^*(\xi_1, \xi_2, p_{t1}, p_{t2}, \ldots)$$

$$\times \psi_A\left(\xi_1, \xi_2, p_{t1} + \vec{k}_\perp, p_{t2} - \vec{k}_\perp, \ldots\right) G_{a_1}^{N_1}\left(x_1/\xi_1, |\vec{k}_\perp|\right) G_{a_2}^{N_2}\left(x_2/\xi_2, |\vec{k}_\perp|\right)$$

Nucleus wf

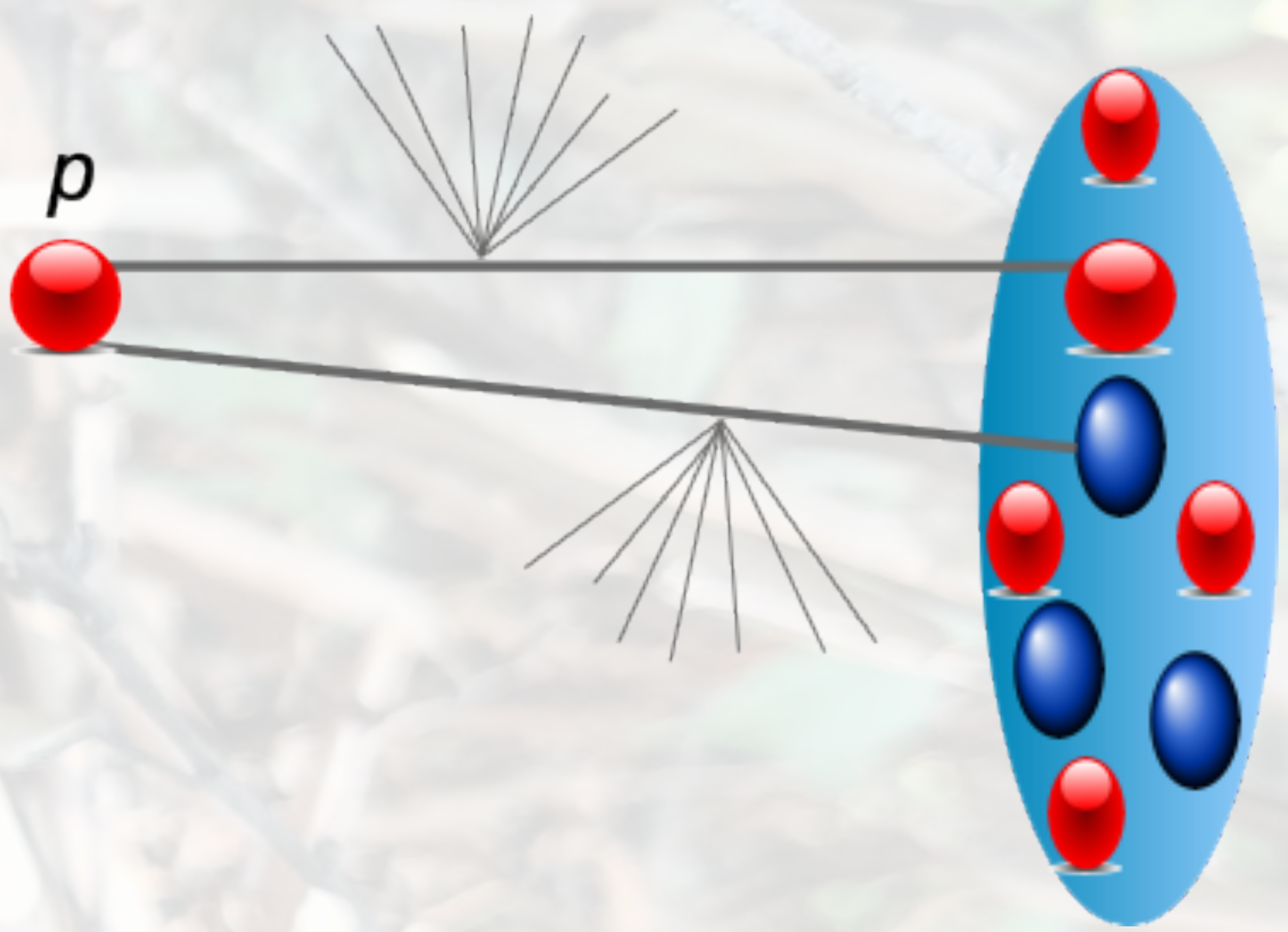
Nucleon GPD

DPS in pA collisions



$$\sigma_{DPS2} \sim A^{1/3} \sigma_{DPS1}$$

$$\sigma_{DPS1} \sim A \sigma_{DPS}^{pp}$$



DPS in γA collisions with light nuclei?

M.R. in progress

In p-Pb collisions there are some difficulties (personal view):

1) both cross-sections (DPS1 and DPS2) depends on proton DPD (still almost unknown) therefore both mechanisms are very important → could be difficult to extract some information on the proton DPD

2) for heavy nuclei is difficult to perform calculation with wave-function obtained from realistic potentials

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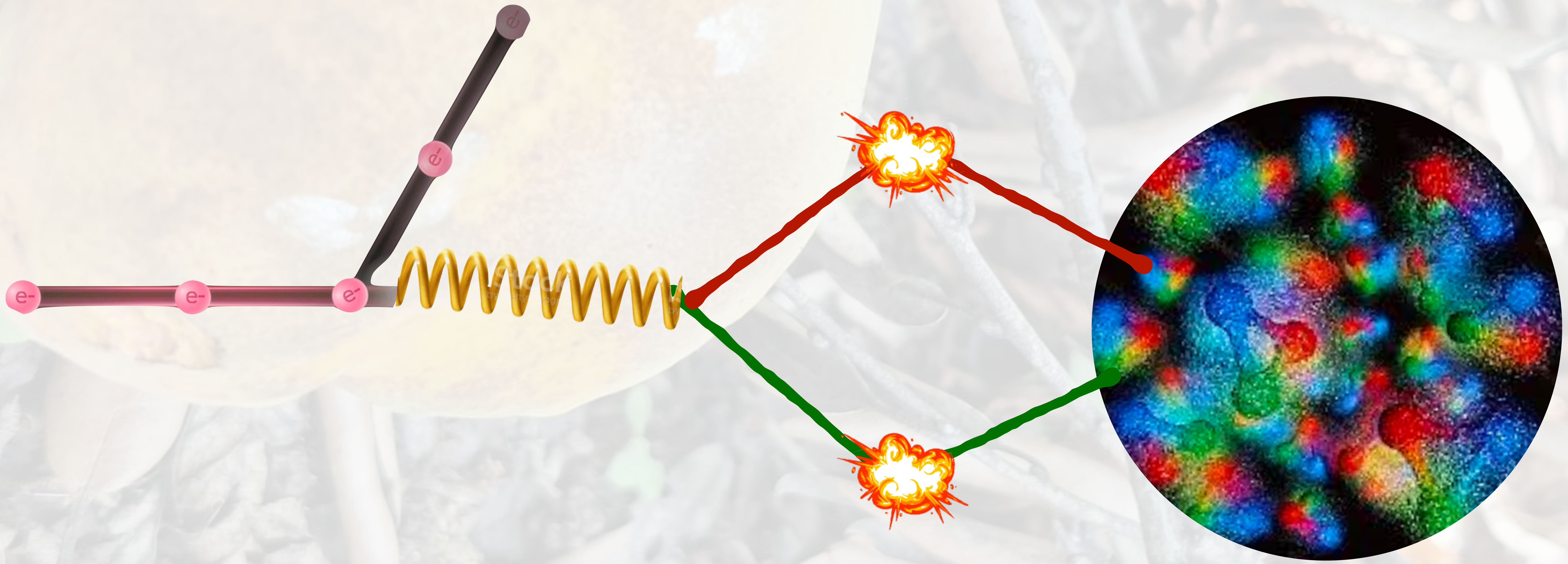
POSSIBLE SOLUTION?



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POSSIBLE SOLUTION?

1) In γA the DPS2 will not contain any DPD of the proton → this mechanism can now be viewed as a background that can be evaluated if we properly treat the photon (as previously discussed) and the Nuclear geometry

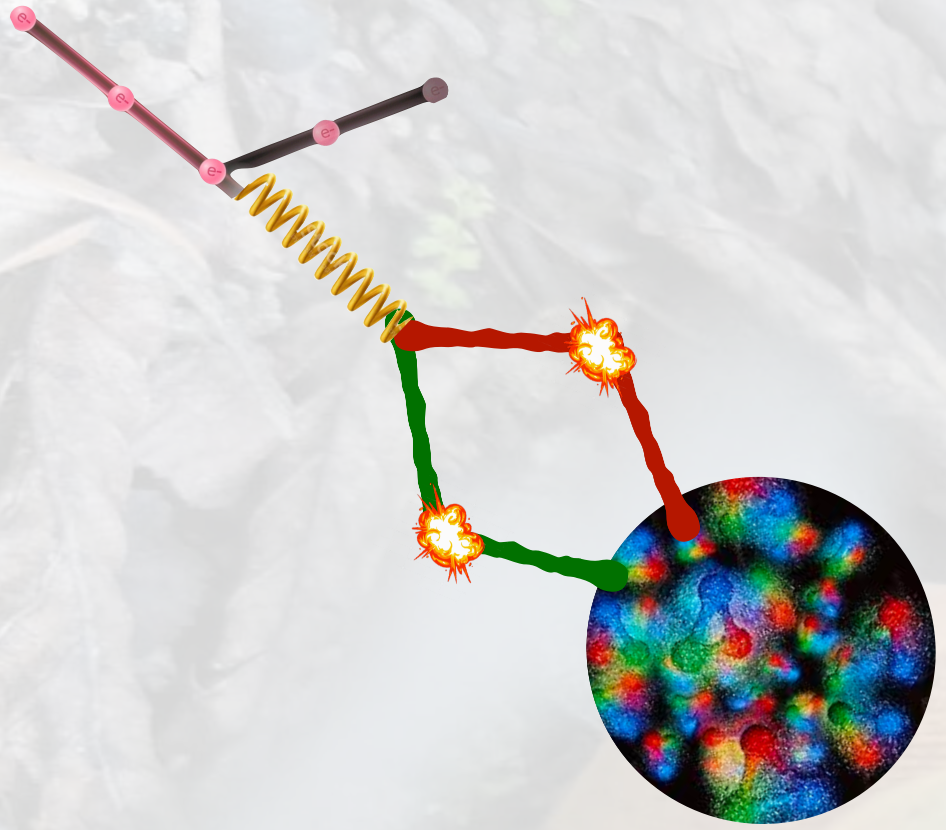
2) For light nuclei these calculations can be done starting from realistic wave-function (Av18 or chiral potential)!

Could we access the DPD of bound nucleons? Double EMC effect?

DPS1 in γA collisions with light nuclei?

M.R. in progress

For example in DPS1:



$$\tilde{F}^1_{a_1 a_2}(x_1, x_2, k_\perp) = \sum_{N=p,n} \int \frac{1}{\xi} \tilde{F}^N_{a_1 a_2}\left(\frac{x_1}{\xi}, \frac{x_2}{\xi}, k_\perp\right) \boxed{\rho^N_A(\xi, p_{t,N})} \frac{d\xi}{\xi} d^2 p_{t,N}$$

The nuclear light-cone distribution can be evaluated with realistic wave-function (from Av18 +UIV potential) in a fully relativistic and Poincaré covariant approach for:

1) H^2 in **E. Pace and G. Salmé, TNPI2000 (2001), arXiv:nucl-th/0106004**

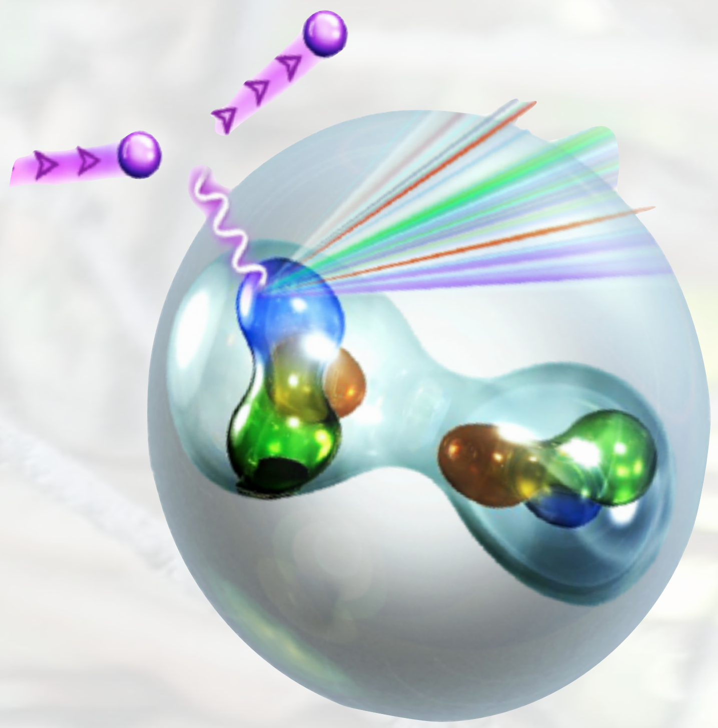
2) He^3 in e.g. **A. Del Dotto et al, PRC 95, 014001 (2017), M.R. et al, PLB 839 (2023), 137810**

3) He^4 from **F. Fornetti, E. Pace, M. R., G. Salmé, S. Scopetta and M. Viviani, PLB 851 (2024) 138587**

The nuclear EMC effect

M.R. in progress

Let us consider DIS processes of nuclei at high energies and evaluate the ratio of cross-sections of different nuclei (A = generic nucleus with A nucleons, 2 = deuteron):



$$R = \frac{\sigma^A}{\sigma^2}$$

since the energy binding per nucleon is few MeV

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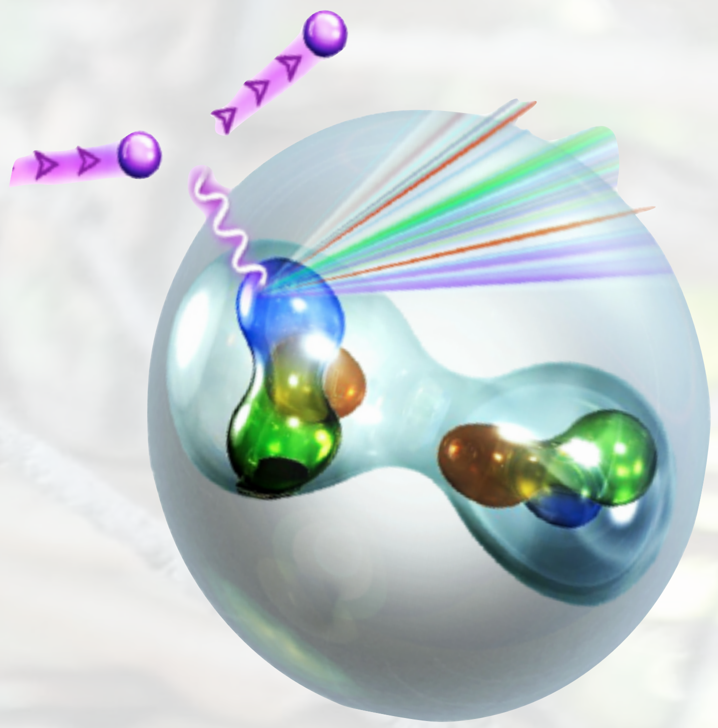

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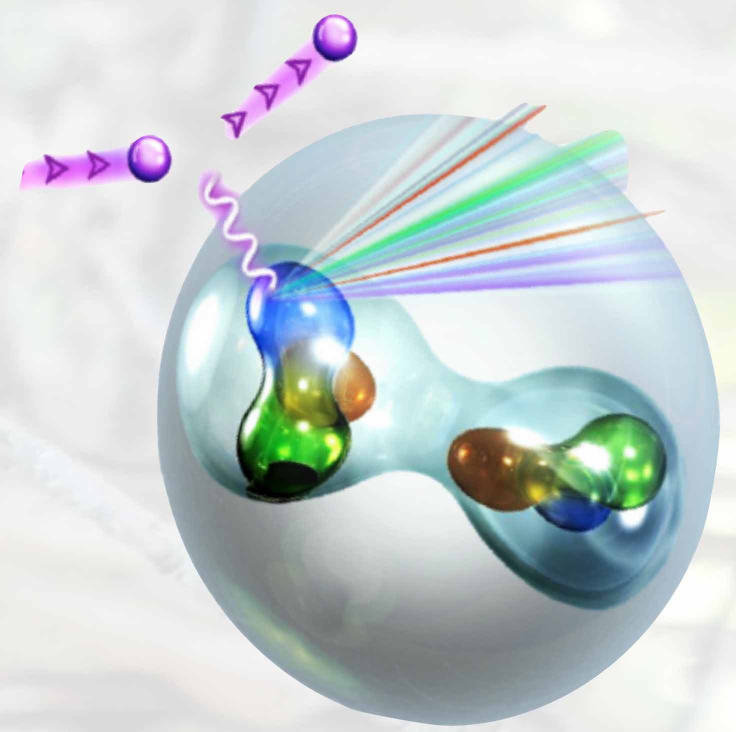


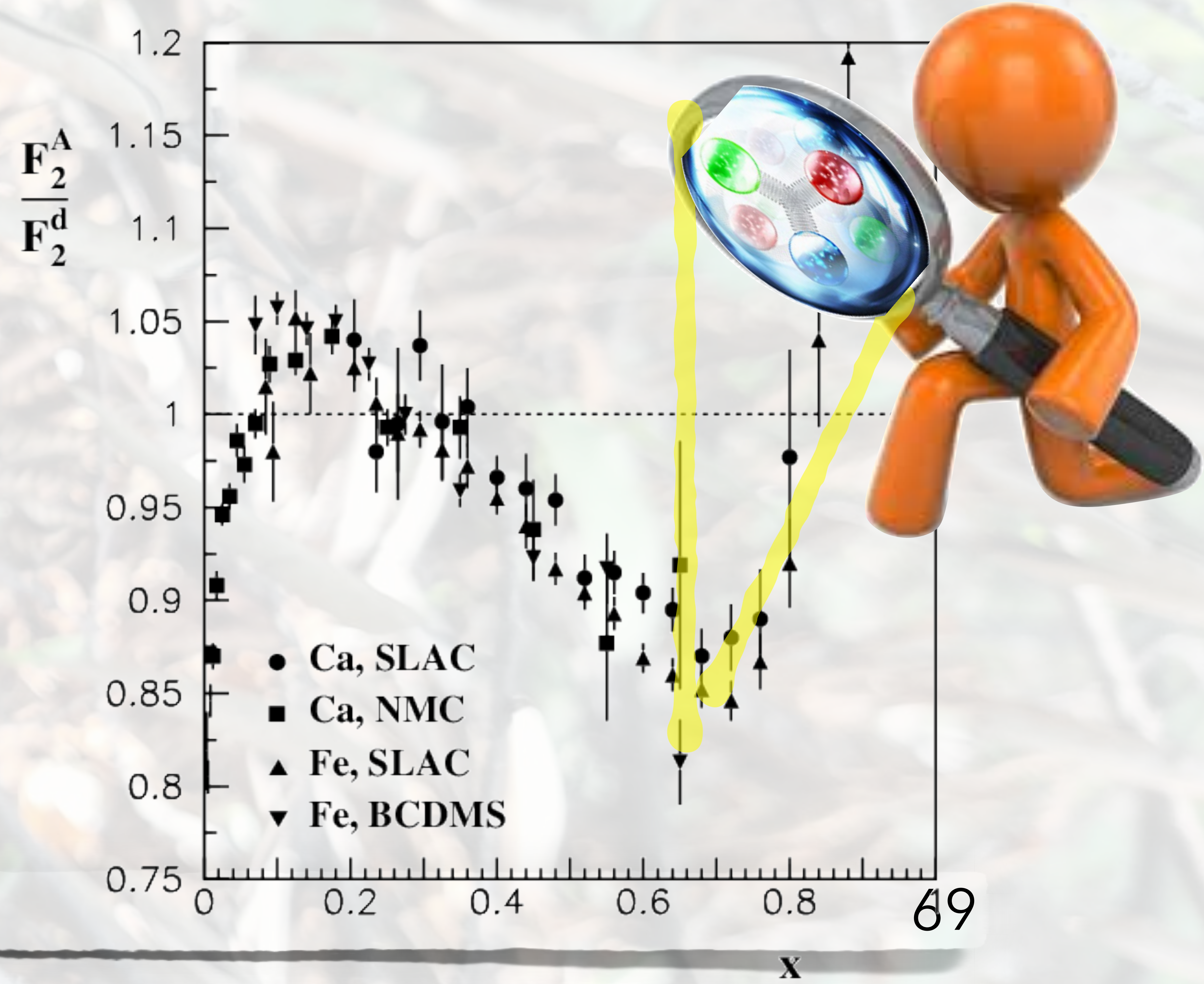
$$R = \frac{\sigma^A}{\sigma^2} \sim \frac{F_2^A(x)}{F_2^2(x)}$$

Nucleare Structure Function

since the energy binding per nucleon is few MeV

$$R \sim 1$$

We found a 10% effect! Why?



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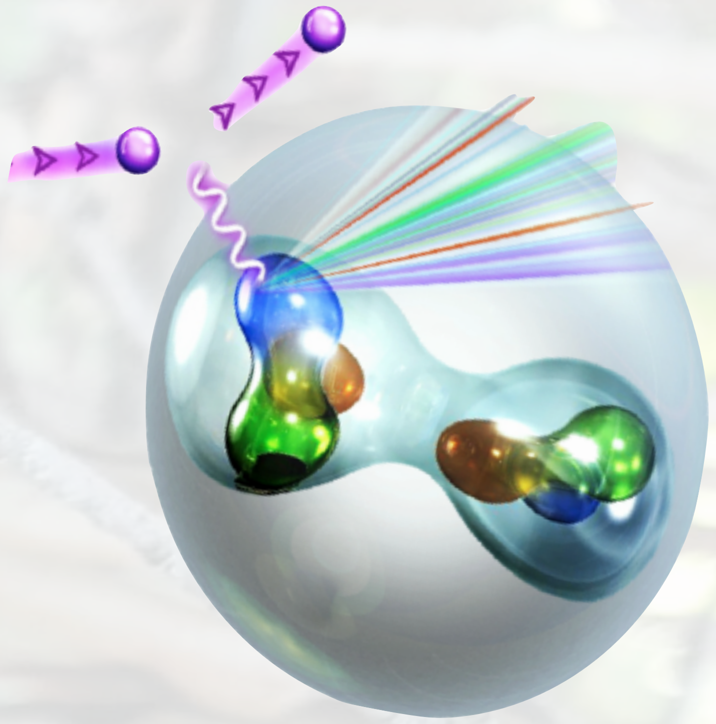

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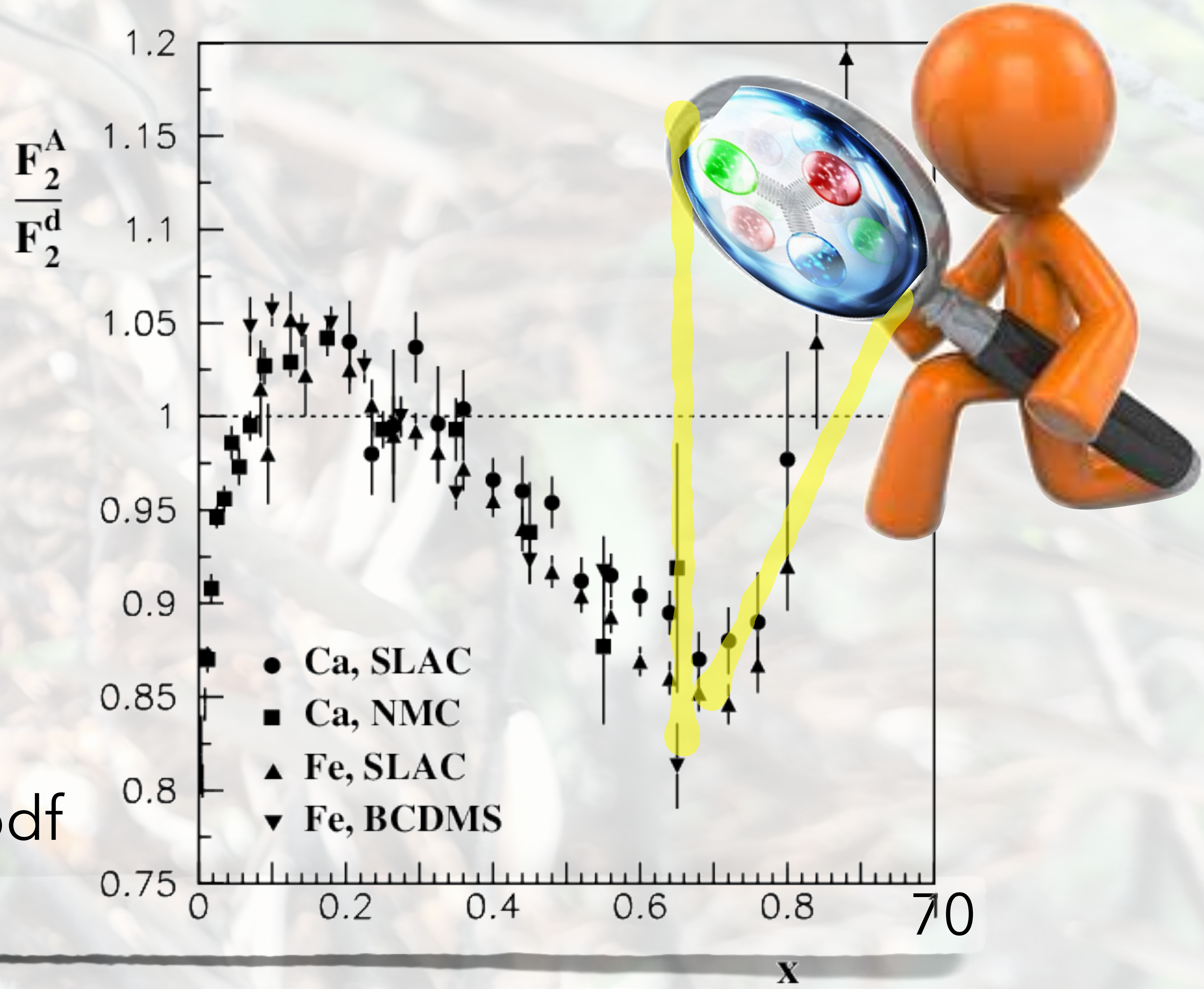
EMC = European Muon Collaboration

Everyone's Model is Cool

that "EMC means Everyone's Model is Cool". It is interesting to note that none of the earliest models were that concerned with the role of two-nucleon correlations, except in relation to six-quark bags.

https://cds.cern.ch/record/1734943/files/vol53-issue4-p035-e.pdf



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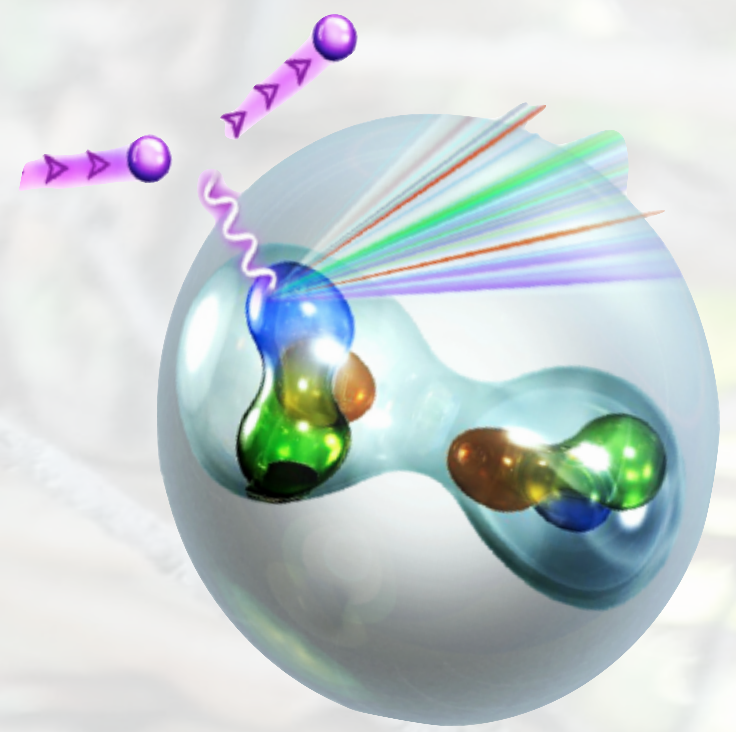
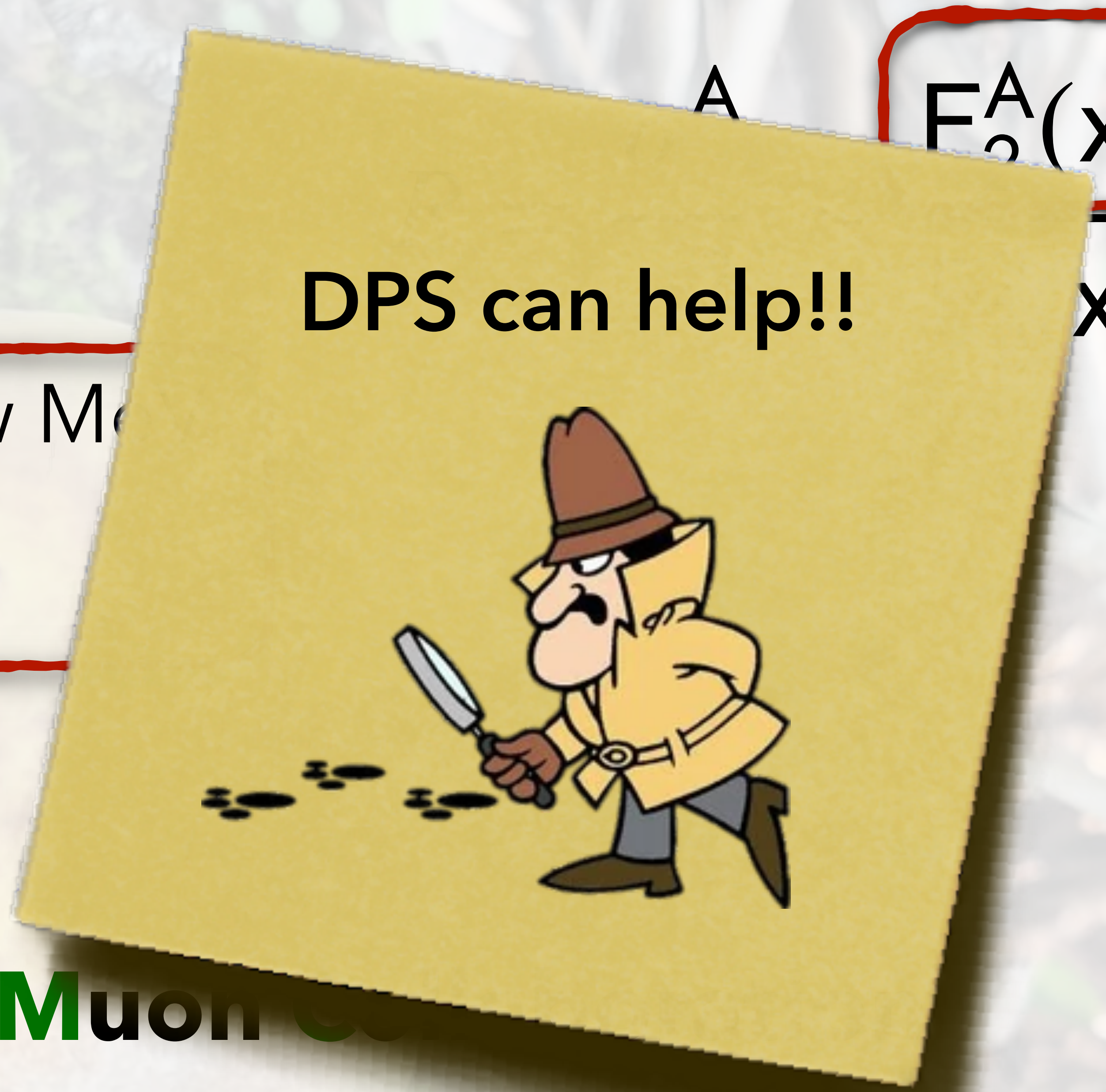


A $F_2^A(x)$ Nucleare Structure Function

x)



since the energy binding per nucleon is few M

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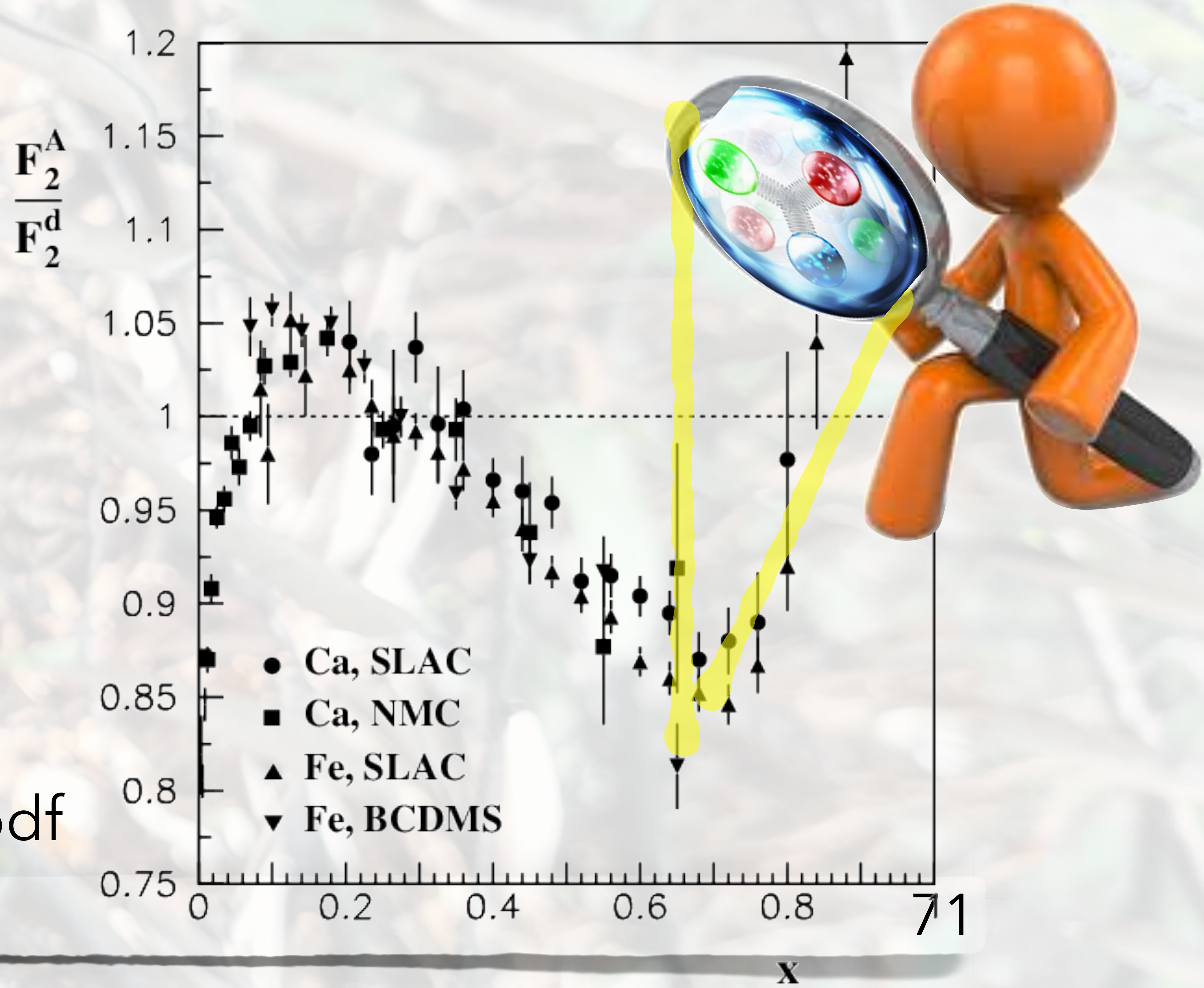
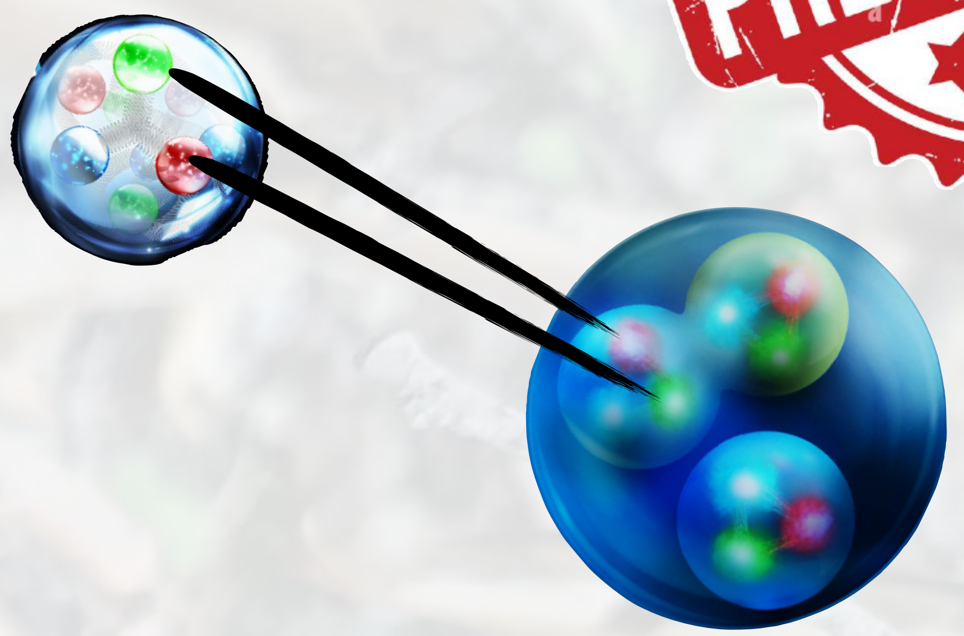


DPS1 and double EMC effect

Nuclear DPD (in momentum space) corresponding to the DPS1 mechanism:



$$\tilde{F}^{1}_{a_1 a_2}(x_1, x_2, k_\perp) = \sum_{N=p,n} \int \frac{1}{\xi} \tilde{F}^{N}_{a_1 a_2}\left(\frac{x_1}{\xi}, \frac{x_2}{\xi}, k_\perp\right) \boxed{\rho^N_A(\xi, p_{t,N})} \frac{d\xi}{\xi} d^2 p_{t,N}$$

The nuclear light-cone distribution can be evaluated with realistic wave-function (from Av18 +UIV potential) for light nuclei and modeled for heavy ions.

We can define the double structure functions (dSF): $F^{2,A}(x_1, x_2) \equiv \sum_{ij} e_i^2 e_j^2 x_1 x_2 \tilde{F}^1_{ij}(x_1, x_2, 0)$

We can generalize the EMC ratio: $R^A_{EMC}(x) = \frac{F^A_2(x)}{A} \frac{2}{\boxed{F^2_2(x)}}$ (^2H) $\longrightarrow$ $R^A_{2EMC}(x_1, x_2) = \frac{F^{2,A}(x_1, x_2)}{A} \frac{2}{F^{2,2}(x_1, x_2)}$

DPS1 and double EMC effect

PRELIMINARY

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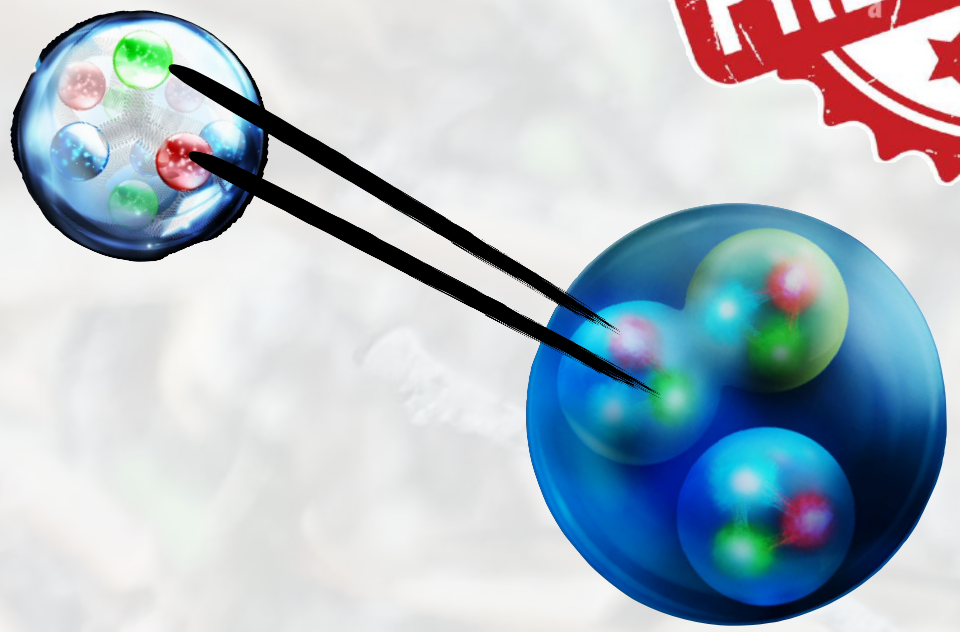

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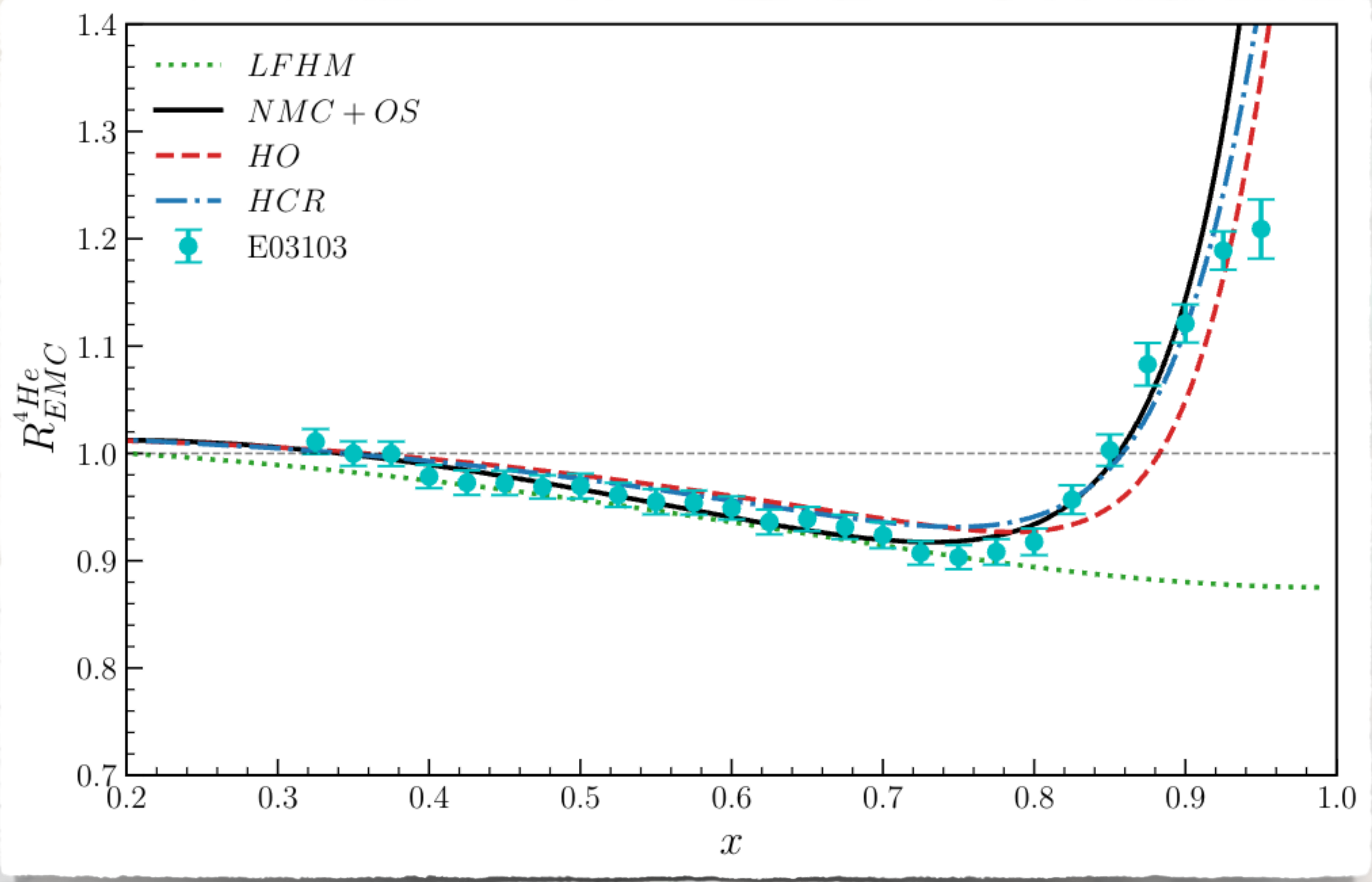
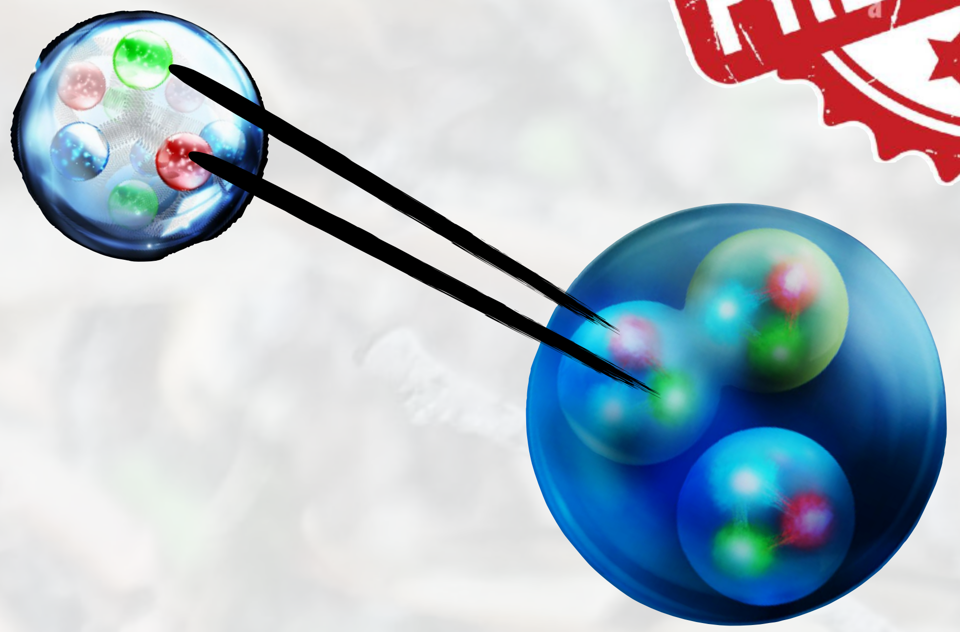
$$R^A_{EMC}(x)=\frac{F^A_2(x)}{A}\frac{2}{F^2_2(x)}$$



DPS1 and double EMC effect

Nuclear DPD (in momentum space) corresponding to the DPS1 mechanism:



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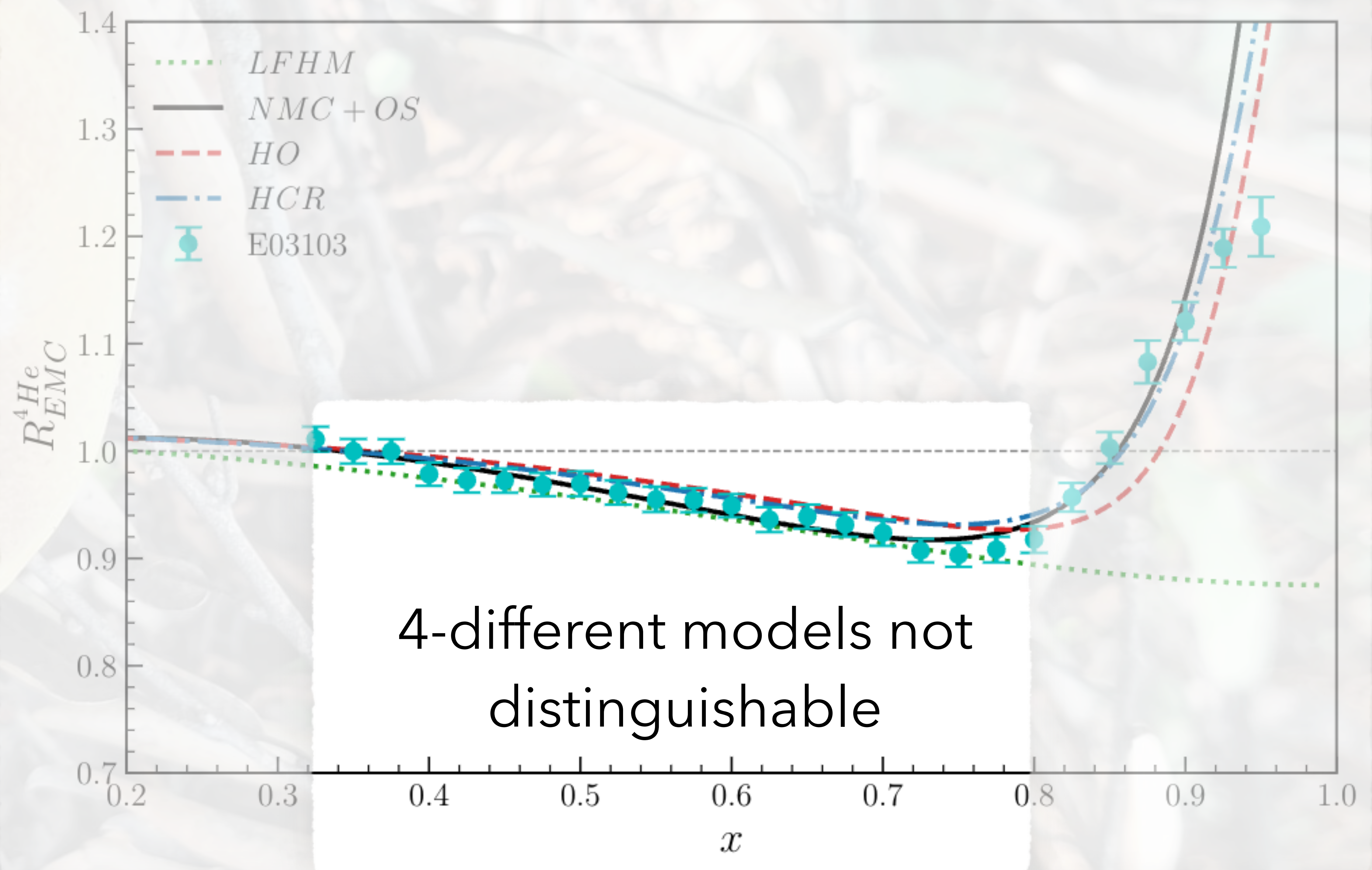

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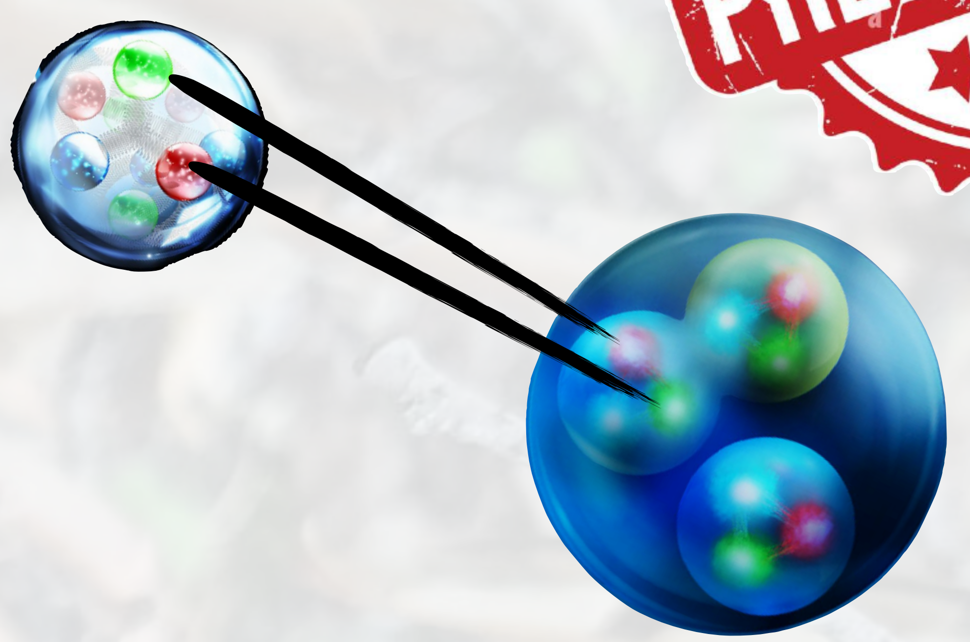

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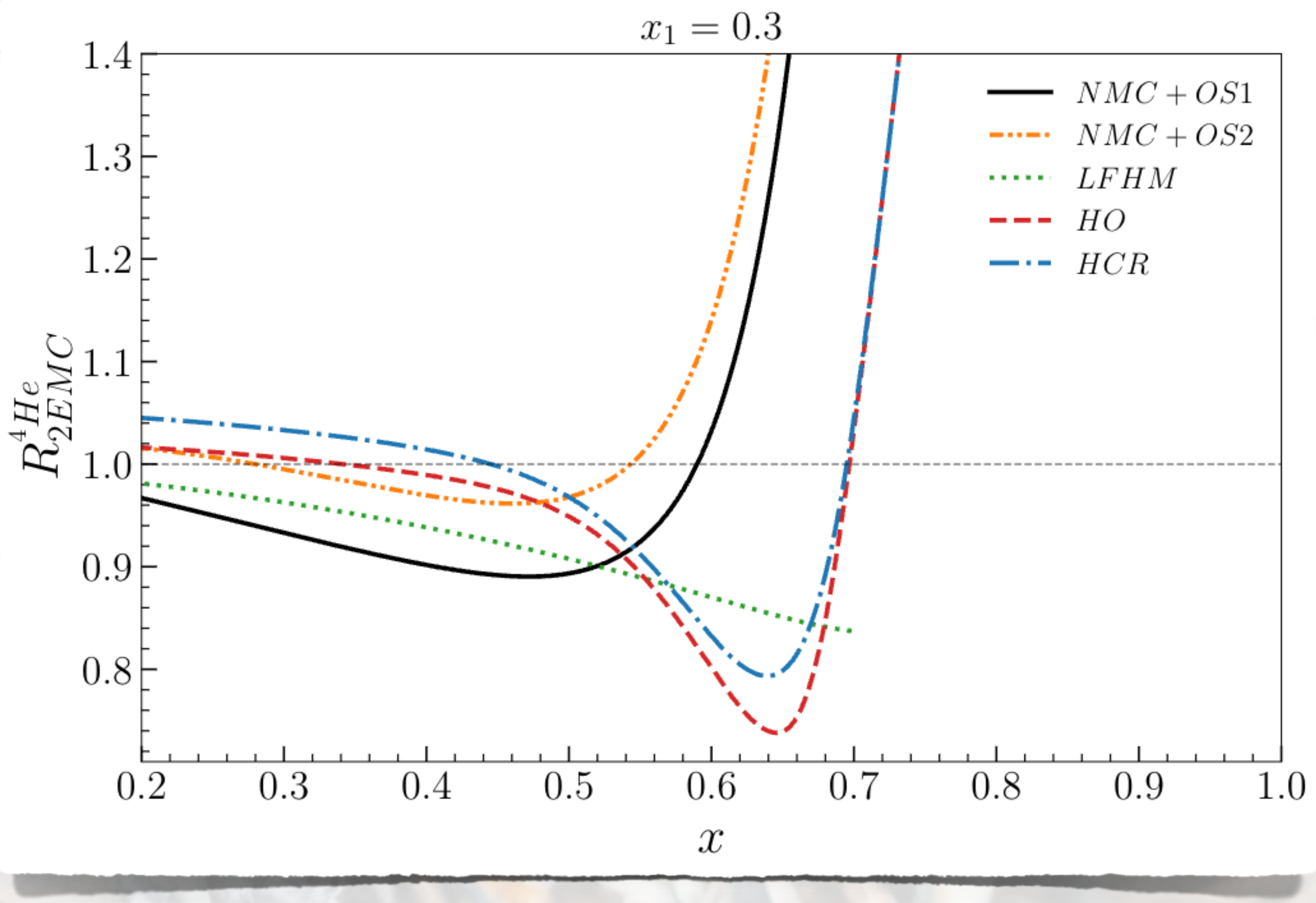
$$R^A_{2EMC}(x_1, x_2) = \frac{F^{2,A}(x_1, x_2)}{A} \frac{2}{F^{2,2}(x_1, x_2)}$$

4-different models leading different results!!



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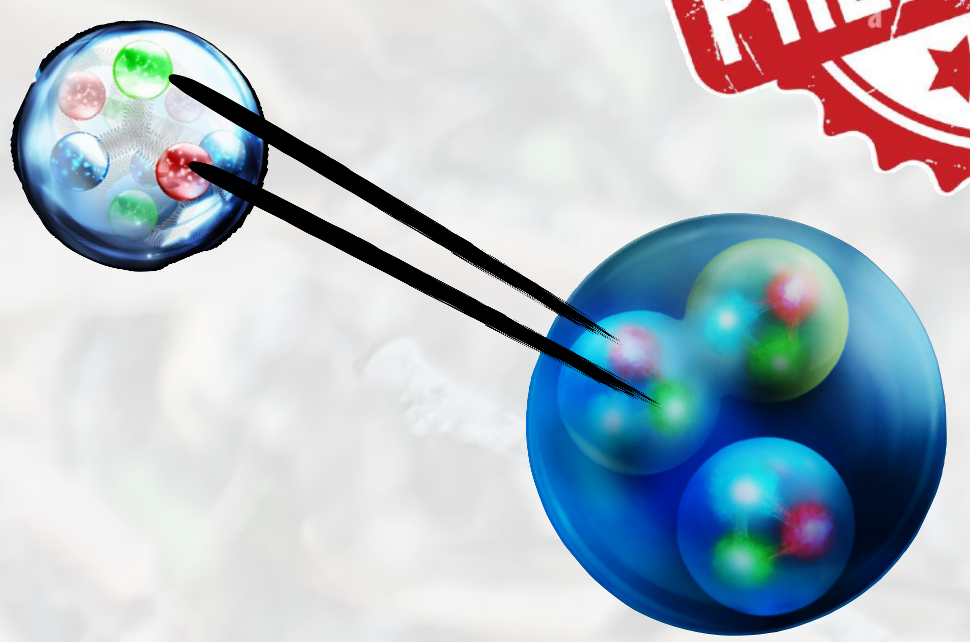

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Deep inside the nucleon DPDs

Usually, DPDs are studied for Physics of the low *x* region, i.e.:

$$\tilde{F}_{ij}(x_1, x_2, 0) \sim f_i(x) f_j(x_2)$$

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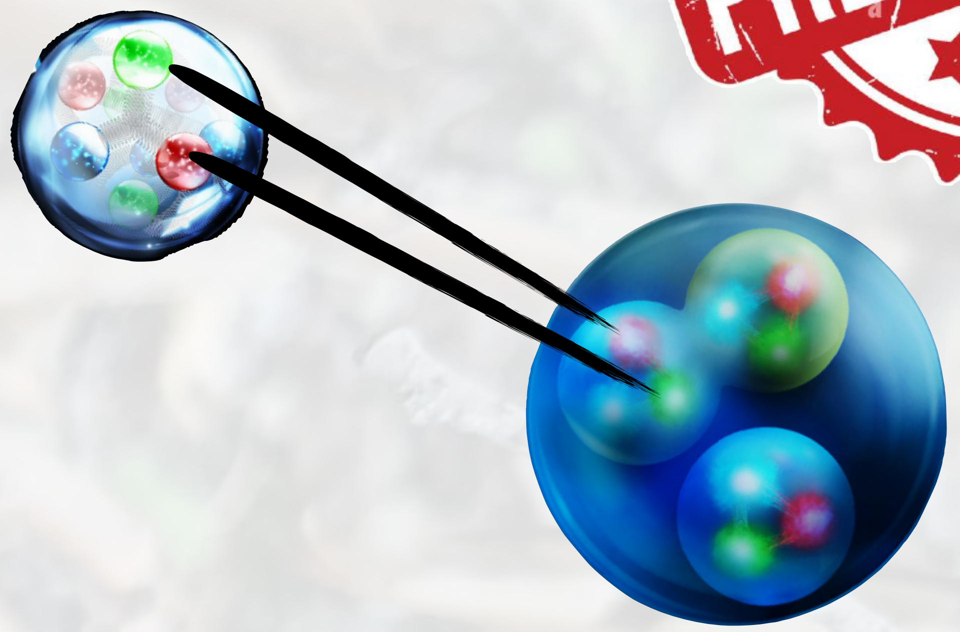

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what happens when $x_1 + x_2 \sim 1$? How fast DPDs go to zero? **We do not know!**

DPS1 and double EMC effect

PRELIMINARY

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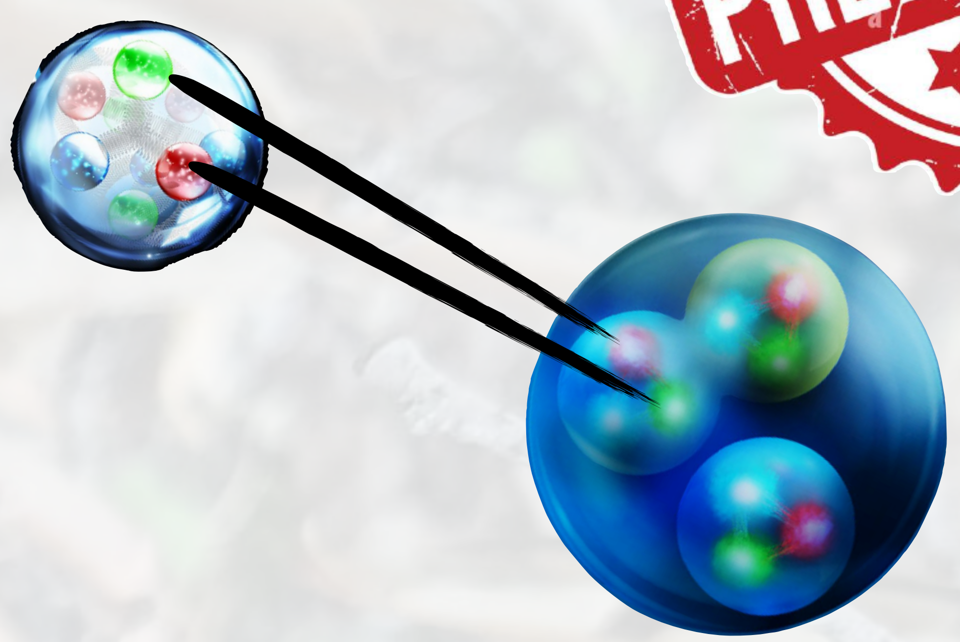

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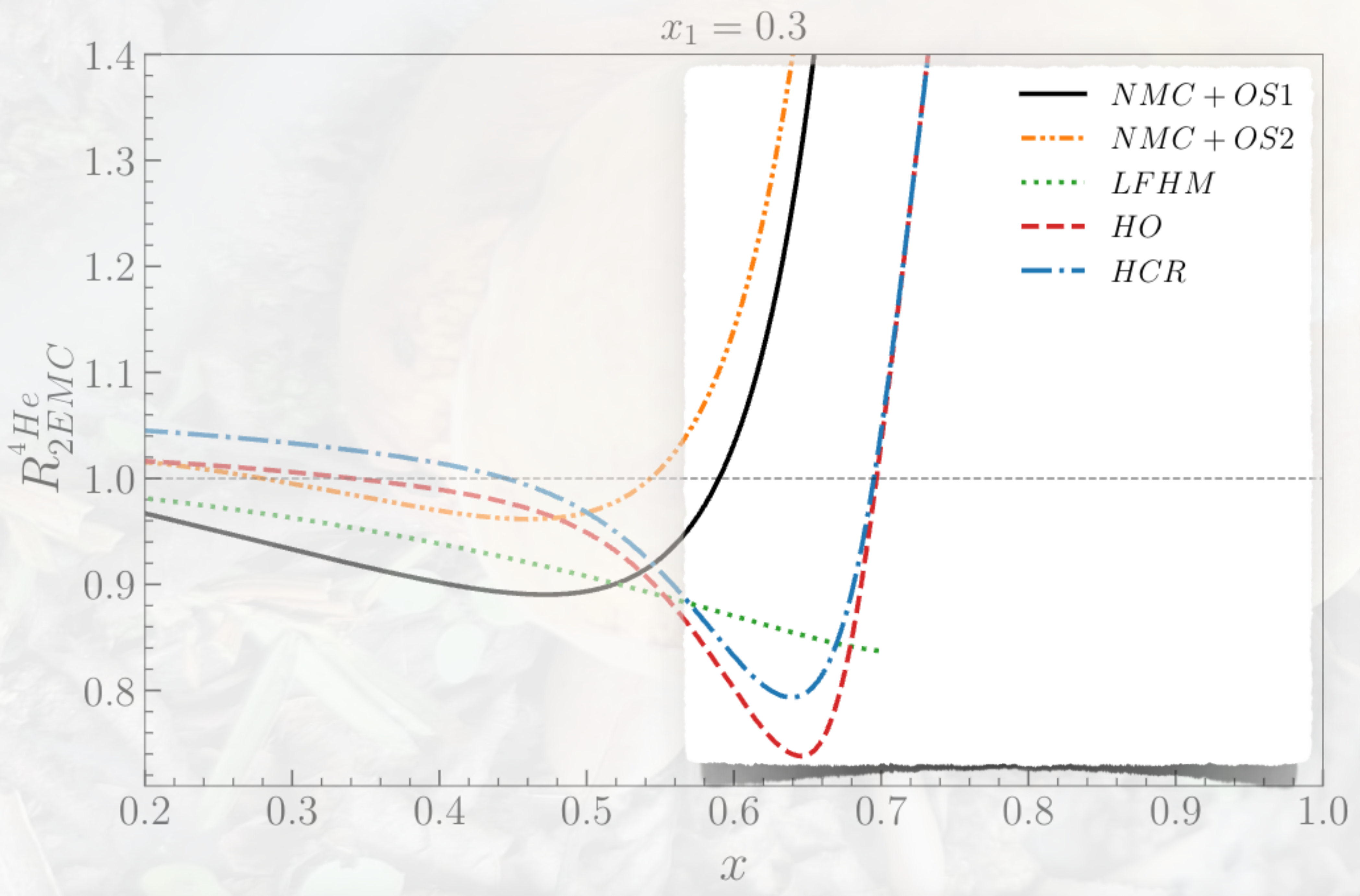

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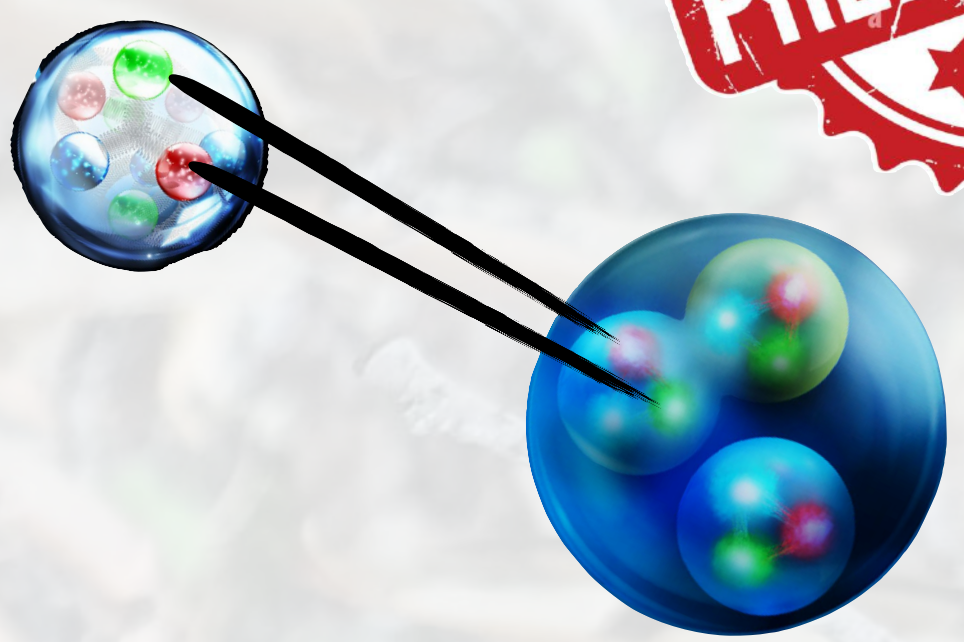
what happens when $x_1 + x_2 \sim 1$? How fast DPDs go to zero?

This information could be extracted from this ratio! Its rise depends on how proton and neutron DPDs go to zero

DPS1 and double EMC effect

Nuclear DPD (in momentum space) corresponding to the DPS1 mechanism:



$$\tilde{F}^1_{a_1 a_2}(x_1, x_2, k_\perp) = \sum_{N=p,n} \int \frac{1}{\xi} \tilde{F}^N_{a_1 a_2}\left(\frac{x_1}{\xi}, \frac{x_2}{\xi}, k_\perp\right) \boxed{\rho^N_A(\xi, p_{t,N})} \frac{d\xi}{\xi} d^2 p_{t,N}$$

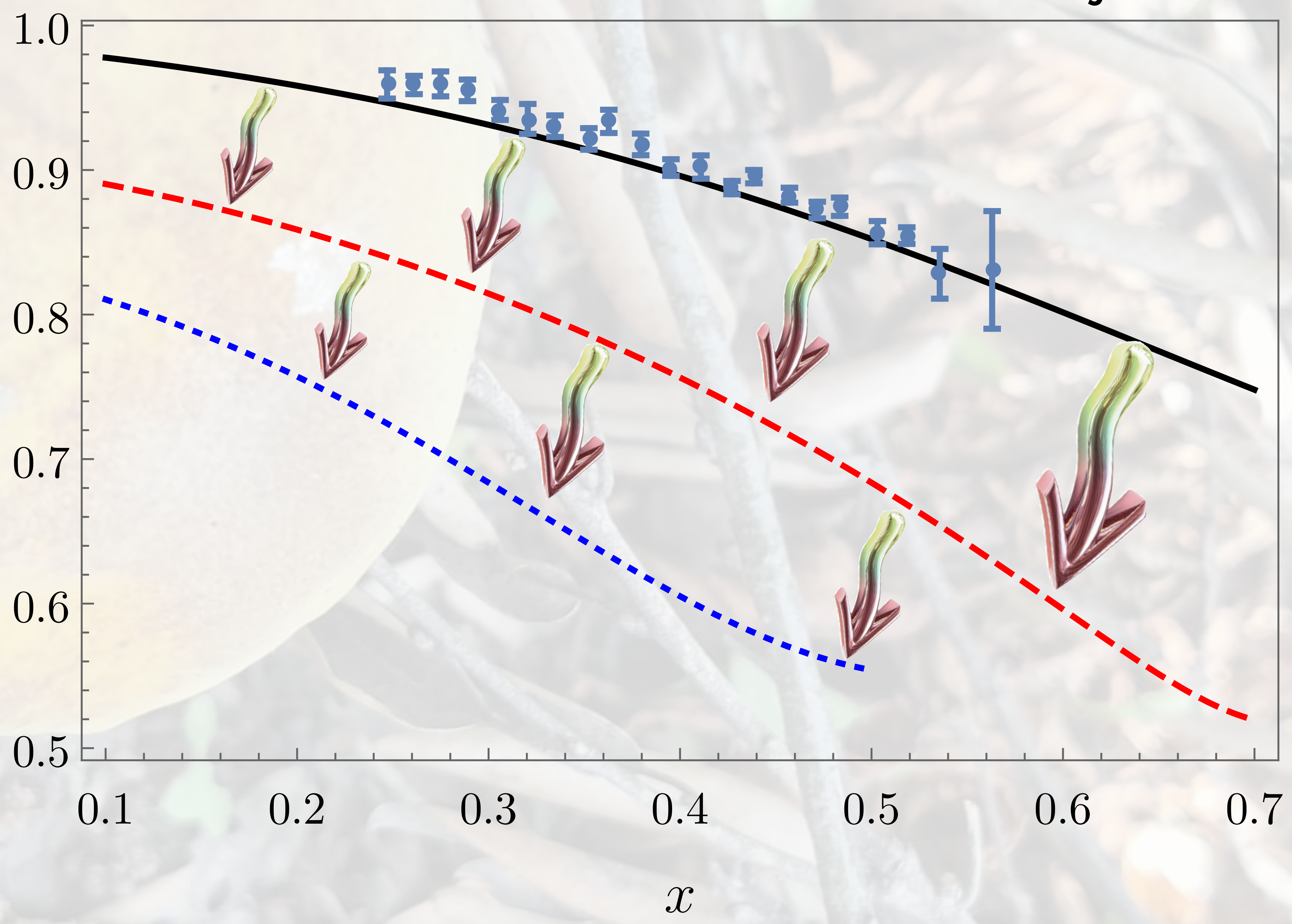
We can define the double structure functions (dSF): $$F^{2,A}(x_1, x_2) \equiv \sum_{ij} e_i^2 e_j^2 x_1 x_2 \tilde{F}^1_{ij}(x_1, x_2, 0)$$

We can generalize the EMC ratio:

$$R^A_{EMC}(x) = \frac{F^A_2(x)}{A} \frac{2}{F^2_2(x)}$$

$$R^A_{2EMC}(x_1, x_2) = \frac{F^{2,A}(x_1, x_2)}{A} \frac{2}{F^{2,2}(x_1, x_2)}$$

Pb



$R^{Pb}_{EMC}(x)$

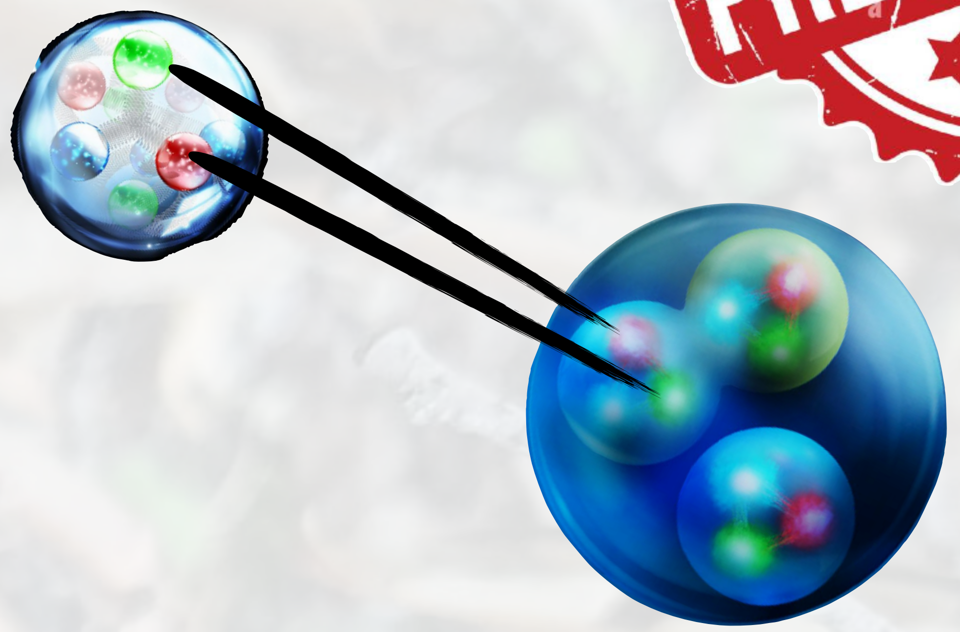
$R^{Pb}_{2EMC}(x, 0.3)$

$R^{Pb}_{2EMC}(x, 0.5)$

Calculations with the model of:

D. N. Kim and G. A. Miller, PRC 106, no.5, 055202 (2022)

DPS1 and double EMC effect



Nuclear DPD (in momentum space) corresponding to the DPS1 mechanism:

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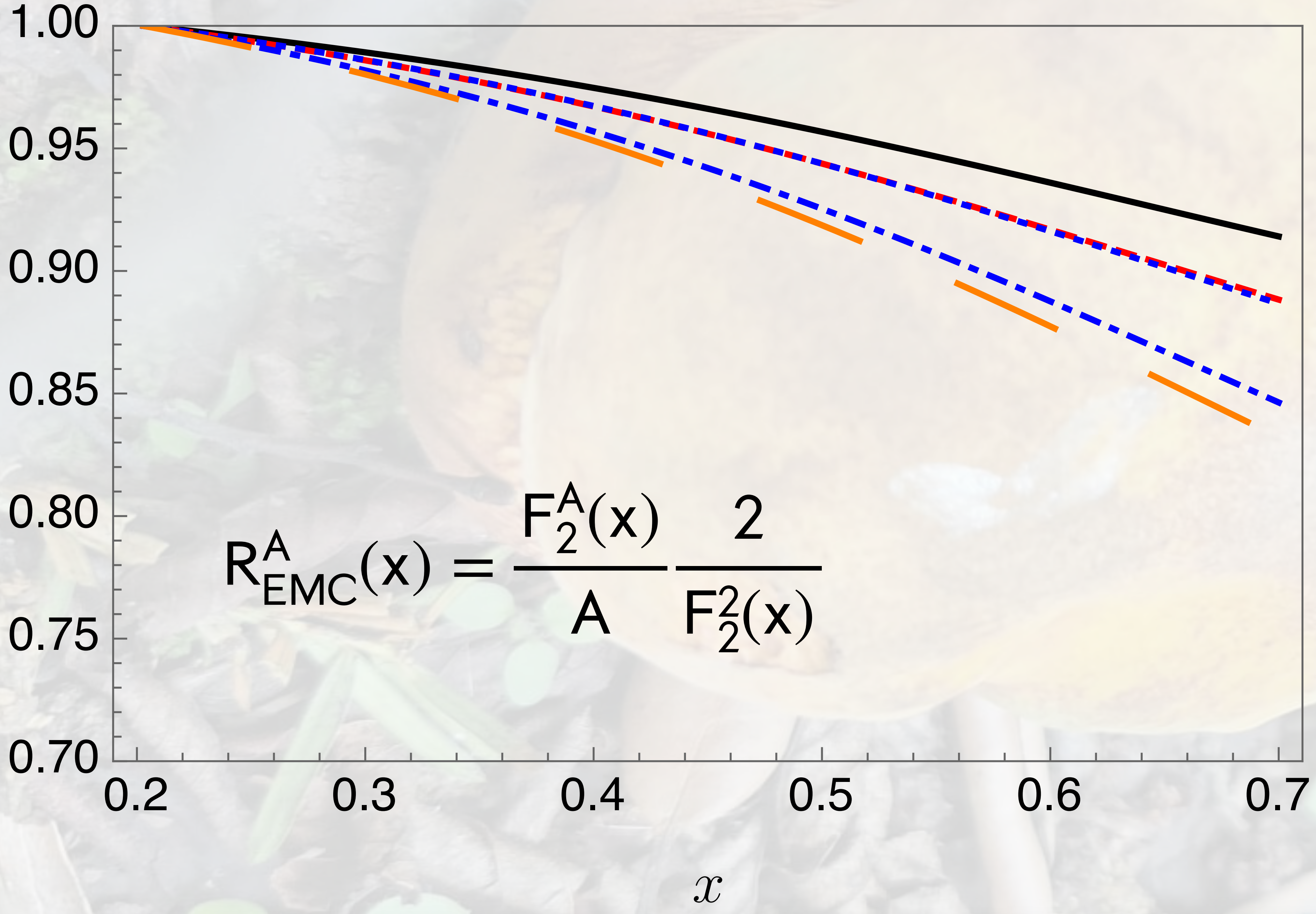
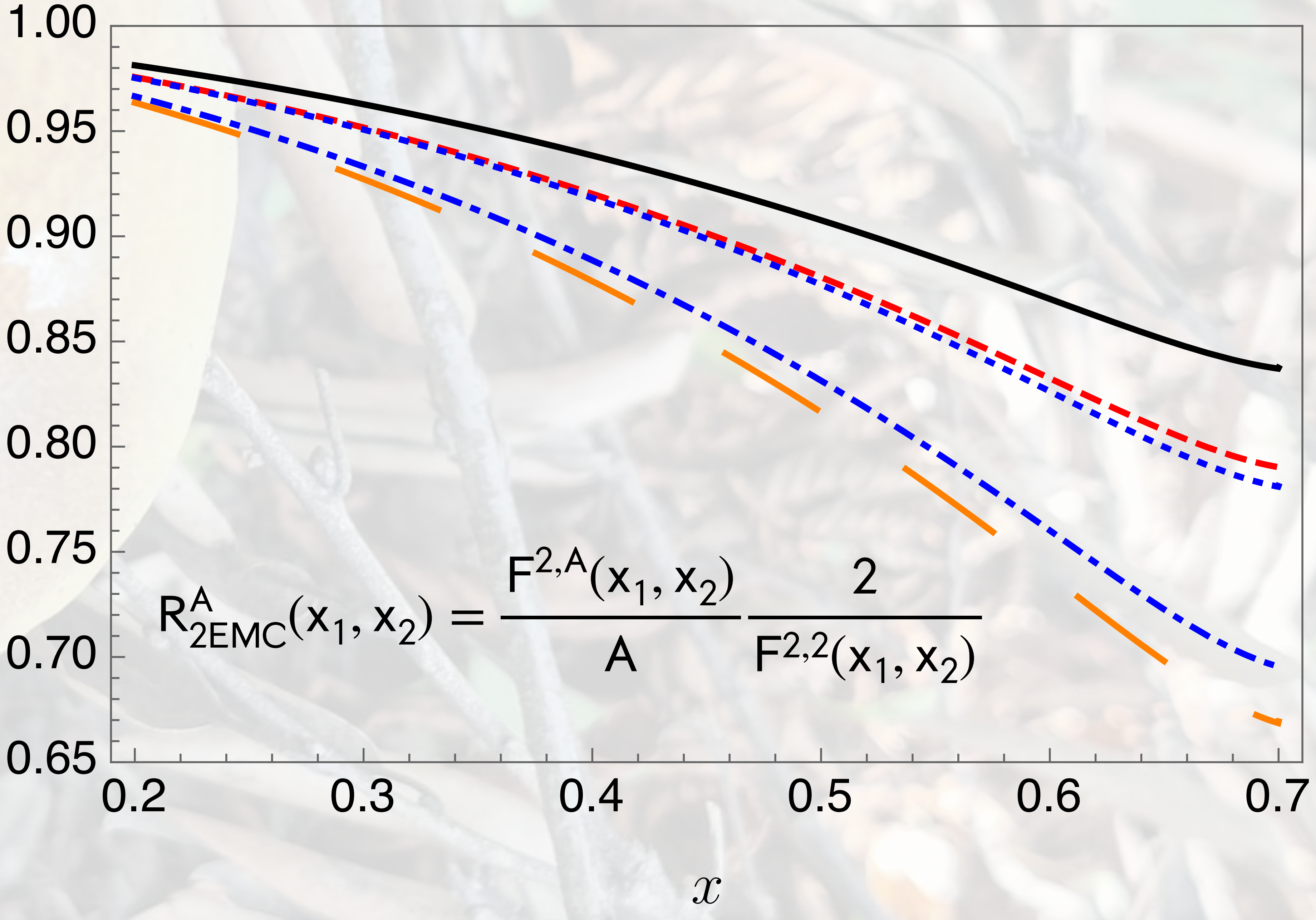
Calculations with the model of:

D. N. Kim and G. A. Miller, PRC 106, no.5, 055202 (2022)

Stronger dependence on A in DPS!





^{4}He

^{12}C

^{63}Cu

^{197}Au

^{208}Pb

DPS2 in γA collisions with light nuclei?

M.R. in progress

For example in DPS2:

$$\tilde{F}^2_{a_1a_2}(x_1, x_2, \vec{k}_\perp) \propto \int \frac{1}{\xi_1\xi_2} \prod_{i=1}^{i=A} \frac{d\xi_i d^2 p_{ti}}{\xi_i} \delta\left(\sum_i \xi_i - A\right) \delta^{(2)}\left(\sum_i \mathbf{p}_{ti}\right) \psi^*_A(\xi_1, \xi_2, p_{t1}, p_{t2}) \psi_A\left(\xi_1, \xi_2, p_{t1} + \vec{k}_\perp, p_{t2} - \vec{k}_\perp\right)$$

$$\times G^{N_1}_{a_1}\left(\frac{x_1}{\xi_1}, |\vec{k}_\perp|\right) G^{N_2}_{a_2}\left(\frac{x_2}{\xi_2}, |\vec{k}_\perp|\right);$$

$$\underset{\xi_i \sim 1}{\sim} G^{N_1}_{a_1}\left(x_1, |\vec{k}_\perp|\right) G^{N_2}_{a_2}\left(x_2, |\vec{k}_\perp|\right)$$

$$\times \left[\int \frac{1}{\xi_1\xi_2} \prod_{i=1}^{i=A} \frac{d\xi_i d^2 p_{ti}}{\xi_i} \delta\left(\sum_i \xi_i - A\right) \delta^{(2)}\left(\sum_i \mathbf{p}_{ti}\right) \psi^*_A(\xi_1, \xi_2, p_{t1}, p_{t2}) \psi_A\left(\xi_1, \xi_2, p_{t1} + \vec{k}_\perp, p_{t2} - \vec{k}_\perp\right)\right]$$

DPS2 in γA collisions with light nuclei?

M.R. in progress

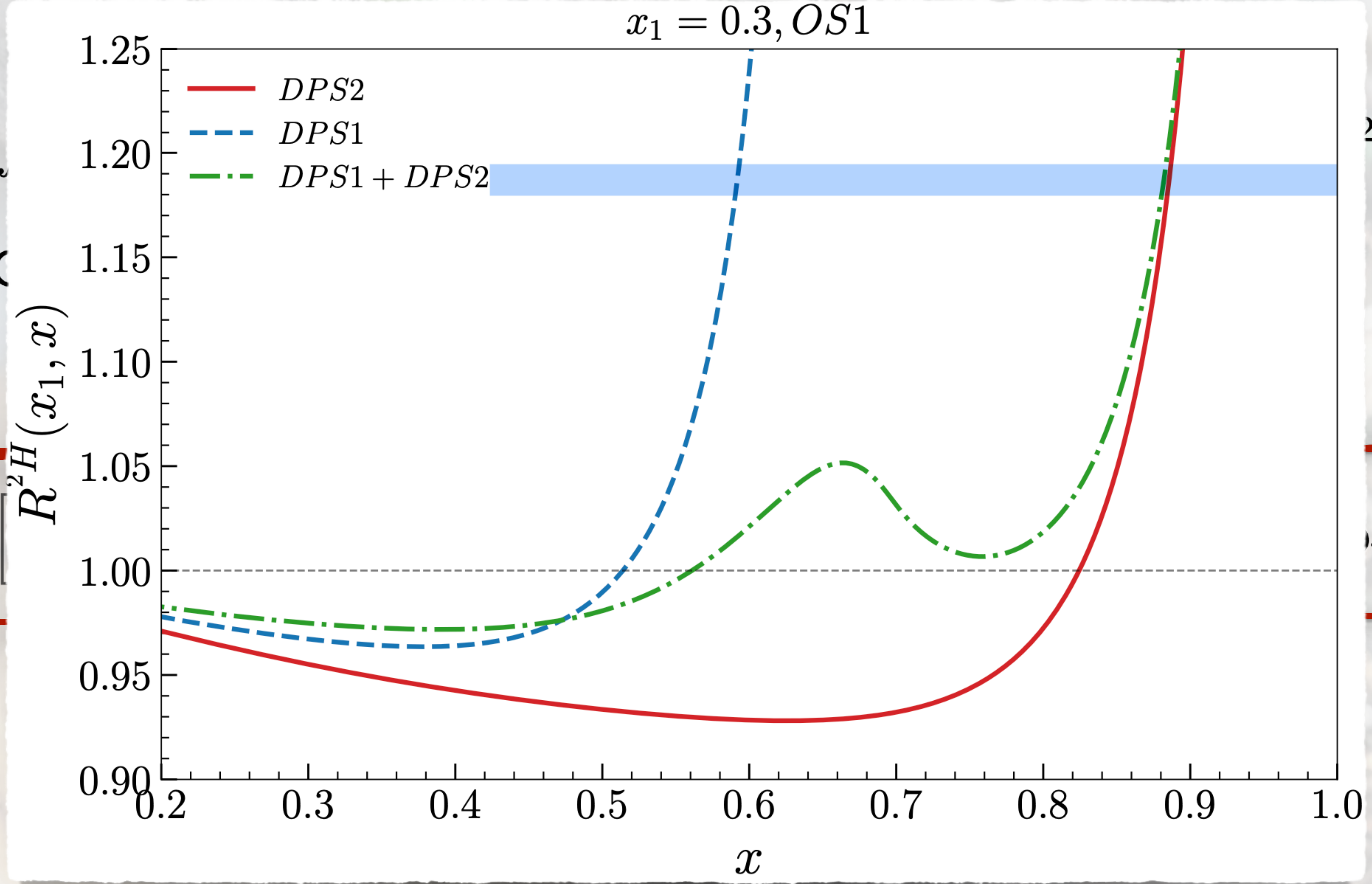
For example in DPS2:

$\tilde{F}^2_{a_1 a_2}(x_1, x_2, \vec{k}_\perp) \propto$ … $)\psi_A\left(\xi_1, \xi_2, p_{t1} + \vec{k}_\perp, p_{t2} - \vec{k}_\perp\right)$

$\times$ C…

$\xi_i \tilde{\sim} 1$

$\times\Big[$ … $_{t2})\psi_A\left(\xi_1, \xi_2, p_{t1} + \vec{k}_\perp, p_{t2} - \vec{k}_\perp\right)\Big]$



DPS2 in γA collisions with light nuclei?

M.R. in progress

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$$\times G^{N_1}_{a_1}\left(\frac{x_1}{\xi_1},|\vec{k}_\perp|\right)G^{N_2}_{a_2}\left(\frac{x_2}{\xi_2},|\vec{k}_\perp|\right);$$

$$\underset{\xi_i\sim 1}{\sim} G^{N_1}_{a_1}\left(x_1,|\vec{k}_\perp|\right)G^{N_2}_{a_2}\left(x_2,|\vec{k}_\perp|\right)$$

$$\times\left[\int \frac{1}{\xi_1\xi_2}\prod_{i=1}^{i=A}\frac{d\xi_i d^2p_{ti}}{\xi_i}\delta\left(\sum_i \xi_i - A\right)\delta^{(2)}\left(\sum_i \mathbf{p}_{ti}\right)\psi^*_A(\xi_1,\xi_2,p_{t1},p_{t2})\psi_A\left(\xi_1,\xi_2,p_{t1}+\vec{k}_\perp,p_{t2}-\vec{k}_\perp\right)\right]$$

Nuclear 2-body form factor $F_2(\vec{k}_\perp,-\vec{k}_\perp)$

Calculated $F_2(\vec{k}_2,\vec{k}_1)$ for ^{3}He and ^{4}He in:

V. Guzey, M.R., S. Scopetta, M. Strikman and M. Viviani et al, "*Coherent J/Ψ electroproduction on He4 and He3 at the EIC: probing Nuclear shadowing one nucleon at a time*", PRL 129 (2022) 24, 242503

DPS2 in γA collisions with light nuclei?

M.R. in progress

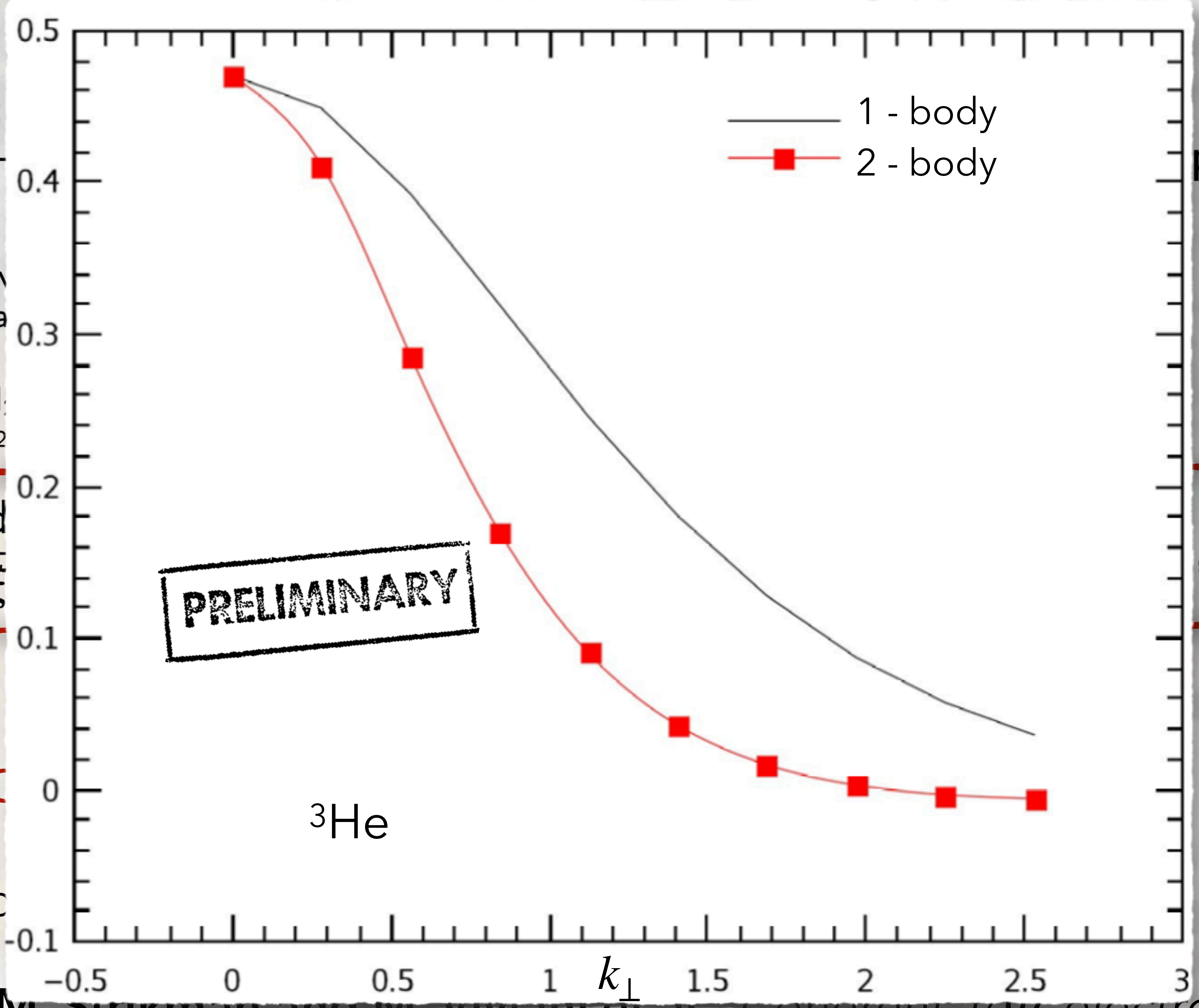
For example in DPS2:

$$\tilde{F}^2_{a_1a_2}(x_1, x_2, \vec{k}_\perp) \propto \int \frac{1}{\xi_1\xi_2} \prod_{i=1}^{i=A} \frac{d\xi_i d^2}{\xi_i} \quad p_{t1} + \vec{k}_\perp, p_{t2} - \vec{k}_\perp\Big)$$

$$\times G^{N_1}_{a_1}\left(\frac{x_1}{\xi_1}, |\vec{k}_\perp|\right) G^{N}_{a}$$

$$\underset{\xi_i \sim 1}{\sim} G^{N_1}_{a_1}\left(x_1, |\vec{k}_\perp|\right) G^{N}_{a_2}$$

$$\times \left[\int \frac{1}{\xi_1\xi_2} \prod_{i=1}^{i=A} \frac{d\xi_i d}{\xi} \quad , p_{t1} + \vec{k}_\perp, p_{t2} - \vec{k}_\perp\Big)\right]$$

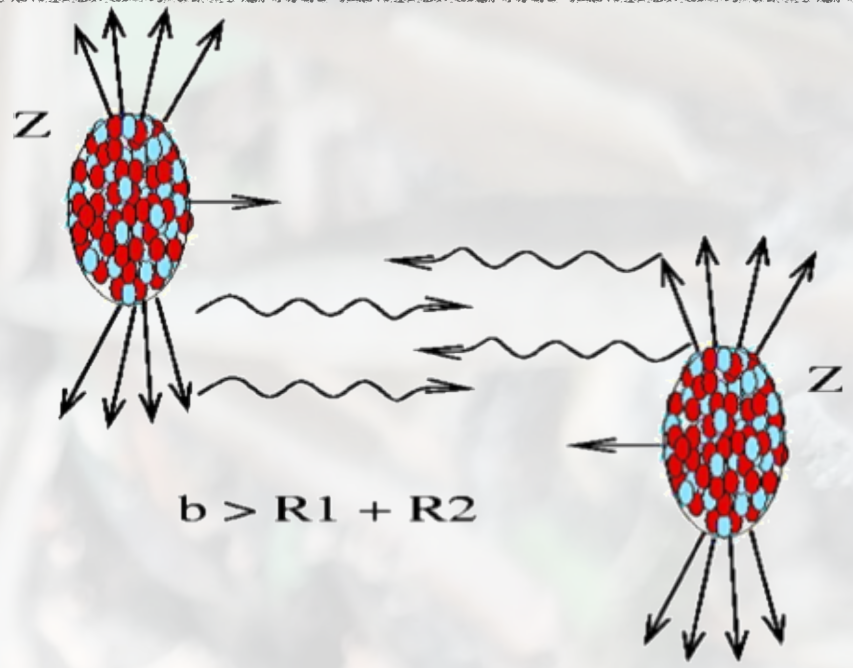


Nuc $\vec{k}_\perp)$

Calculated $F_2(\vec{k}_2, \vec{k}_1)$ fo

V. Guzey, M.R., S. Scopetta, M. Strikman and M. Viviani et al, *"Coherent J/Ψ electroproduction on He4 and He3 at the EIC: probing Nuclear shadowing one nucleon at a time"*, PRL 129 (2022) 24, 242503

J/ψ electroproduction on light-nuclei



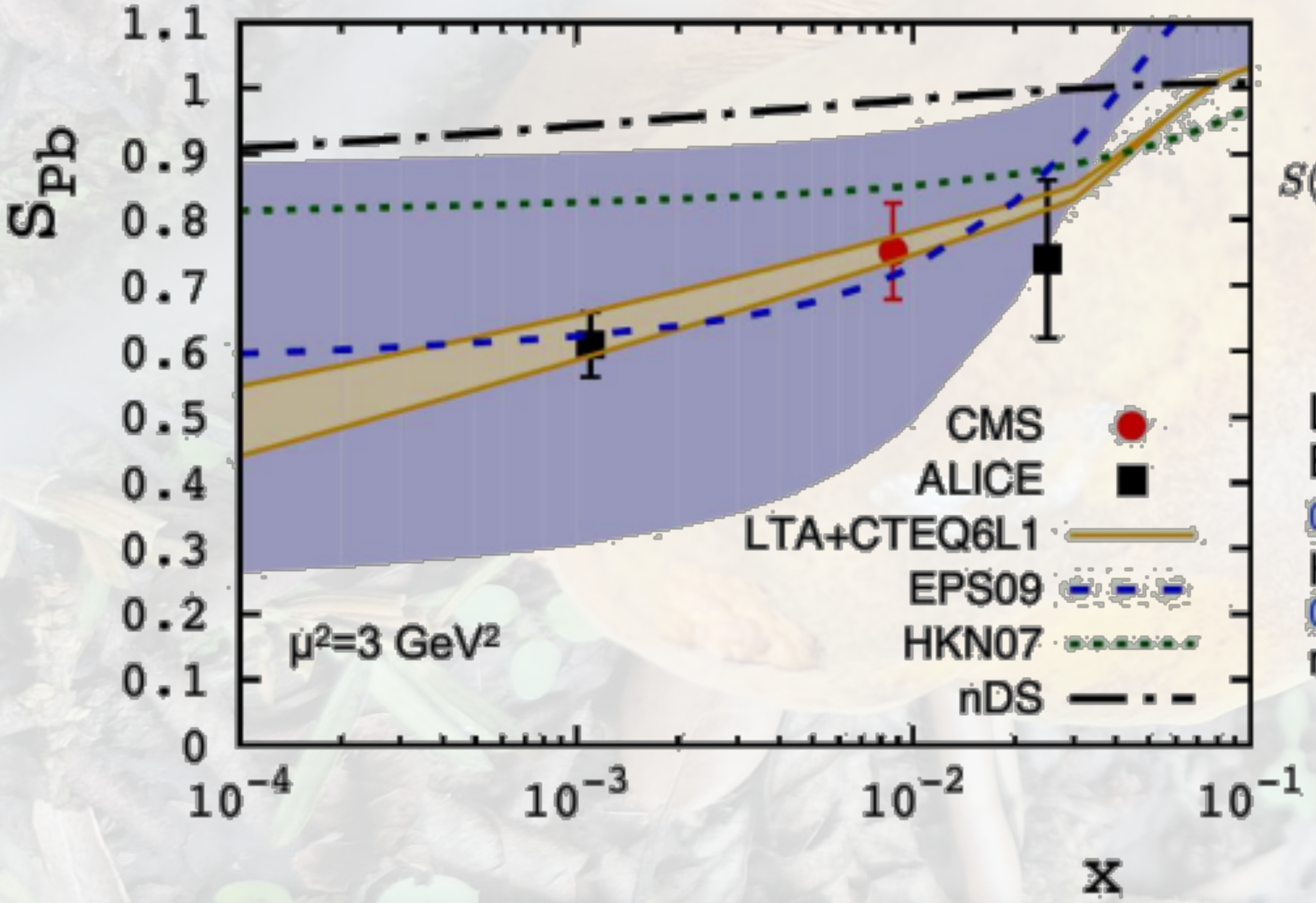
Gluon shadowing in UPC collisions @ LHC

Large (up to 40%) Leading twist (LT) shadowing in:

$\gamma + \mathrm{Pb/Au} \to \rho(\mathrm{J}/\Psi) + \mathrm{Pb/Au}$ Explained/predicted

(**Frankfurt, Guzey, Strikman Phys. Rep. 512 (2012) 255**)

Abbas *et al.* [ALICE], EPJ C 73 (2013) 2617; CMS Collab., PLB 772 (2017) 489 → suppression factor S_{Pb}



$$S(W_{\gamma p}) = \left[\frac{\sigma_{\gamma Pb\to J/\psi Pb}}{\sigma^{IA}_{\gamma Pb\to J/\psi Pb}}\right]^{1/2} = \kappa_{A/N}\frac{G_A(x,\mu^2)}{AG_N(x,\mu^2)}$$

LTA: Guzey, Zhalov JHEP 1310 (2013) 207
EPS09: Eskola, Paukkunen, Salgado, JHEP 0904 (2009) 065
HKN07: Hirai, Kumano, Nagai, PRC 76 (2007) 065207
nDS: de Florian, Sassot, PRD 69 (2004) 074028

Introduction. Studies of nuclear shadowing have a long history [1–5]. In quantum mechanics and in the eikonal limit, it is manifested in the total hadron-nucleus cross section being smaller than the sum of individual hadron-nucleon cross sections. In essence, this is due to simultaneous interactions of the projectile with $k \geq 2$ nucleons of the nuclear target, leading to a reduction (shadowing) of the total cross section. In this frame-

Learning from light nuclei - I

- Problem:
  @ EIC/LHC it is challenging to measure coherent scattering at $t \neq 0$ for $A \simeq 200$; Large coherence length: information on interactions with many nucleons, in average

- Solution:
  use the lightest nuclei, especially ^{3}He and ^{4}He, to study coherent effects for interactions with exactly 2 nucleons in the range of $0 < -t < 0.5$ GeV2 .

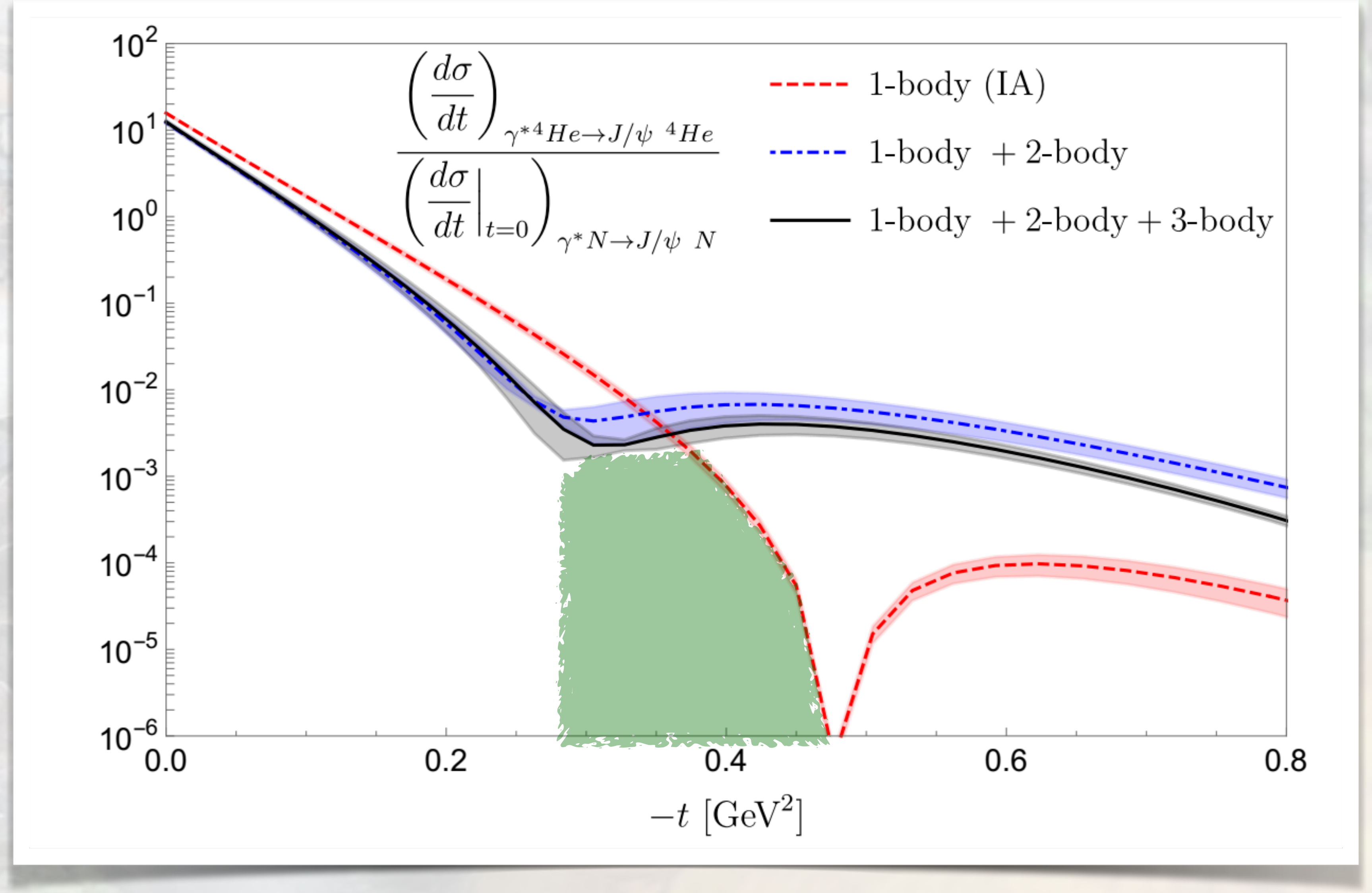
Complementary measurements with light ion beams @ the EIC:

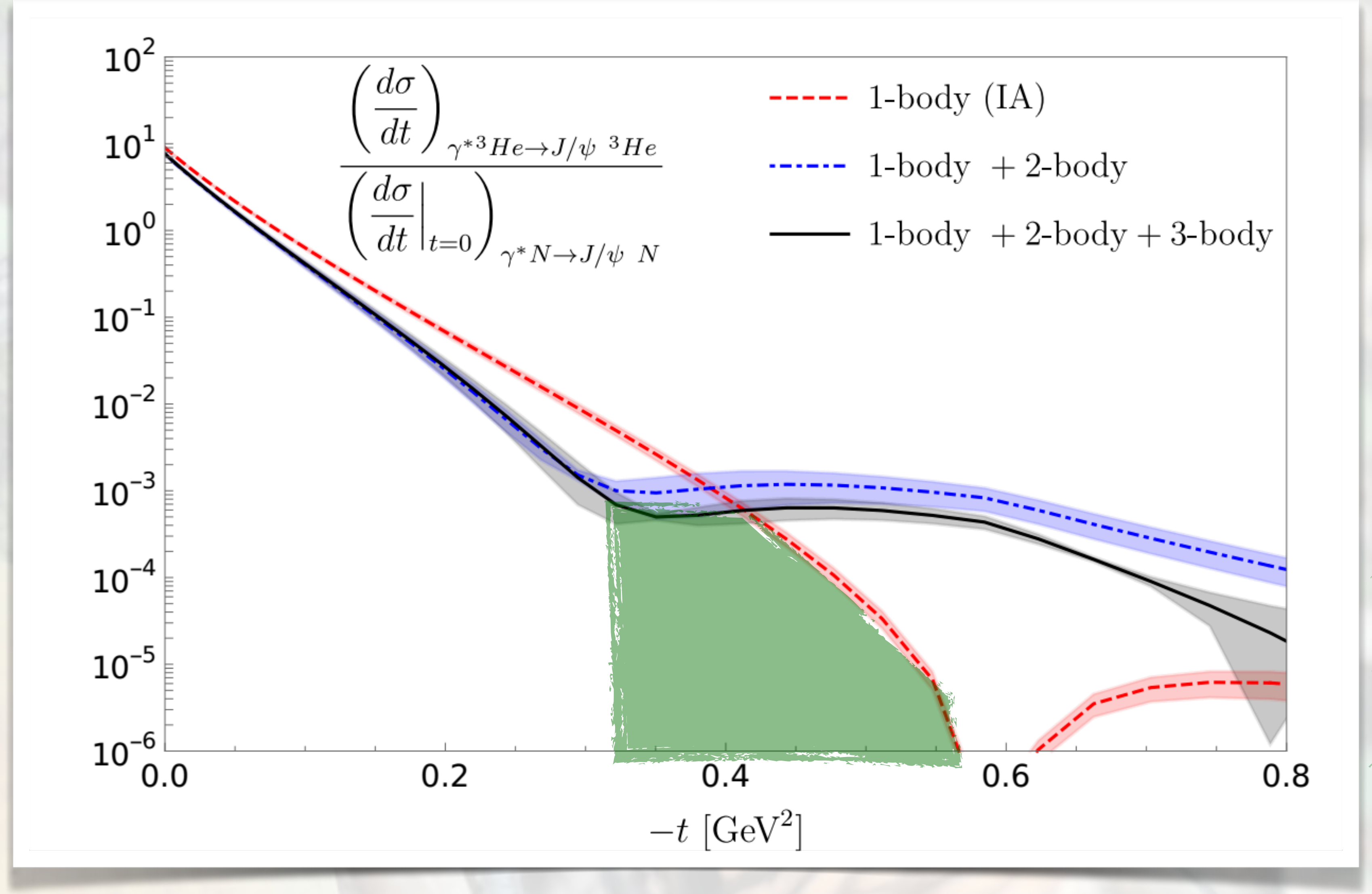
- Scattering off 2 and 3 nucleons can be separately probed
- no excited states -> easy to select coherent events

Here:

Results on J/Ψ diffractive electro-production off ^{3}He – ^{4}He
V. Guzey, M. R., S. Scopetta, M. Strikman and M. Viviani, PRL 129 (2022) 24, 24503

Results for J/Ψ exclusive production @EIC: xB ≈ 10-3





Error bars account:

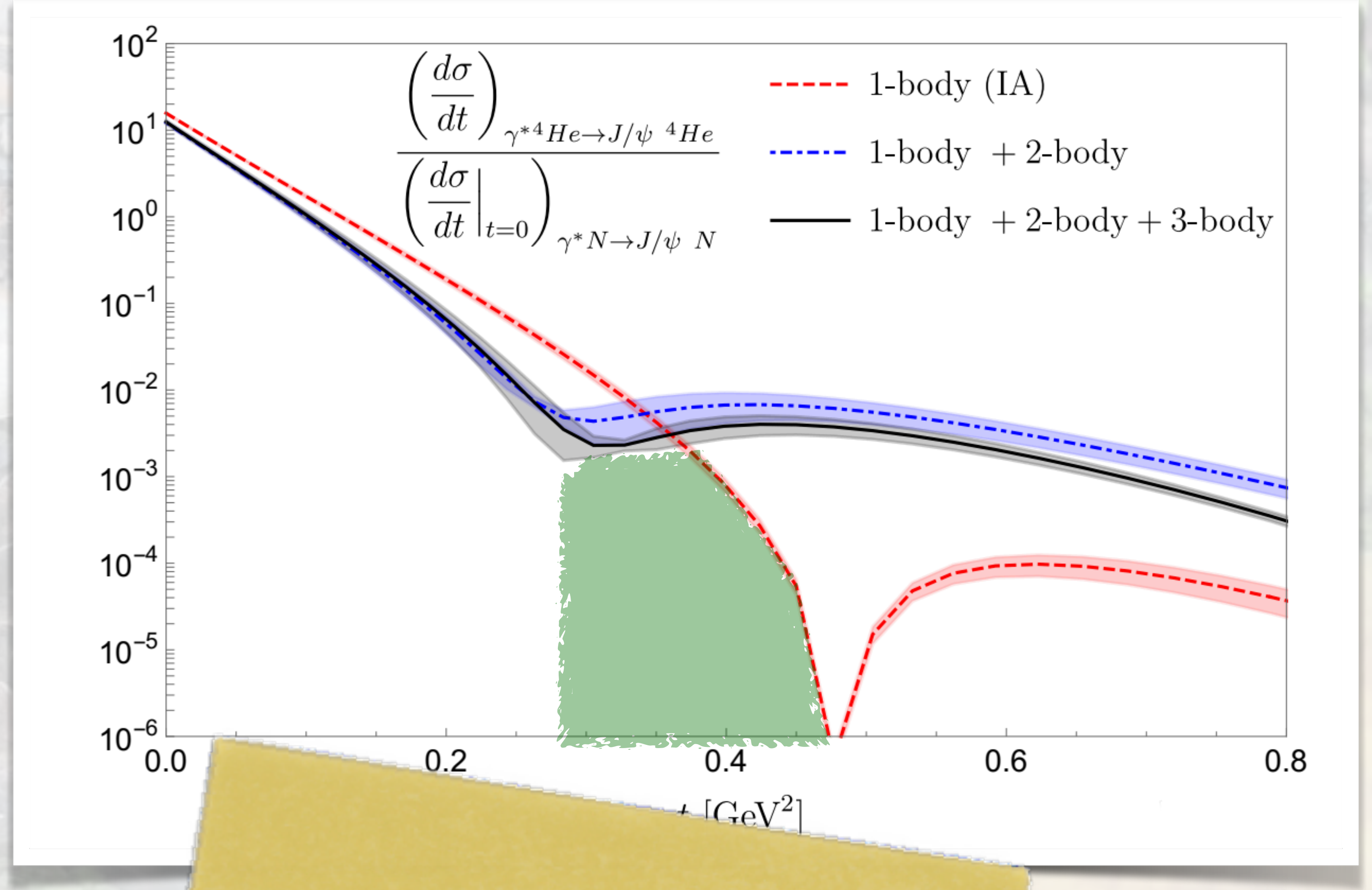
-10% of variation for B_0

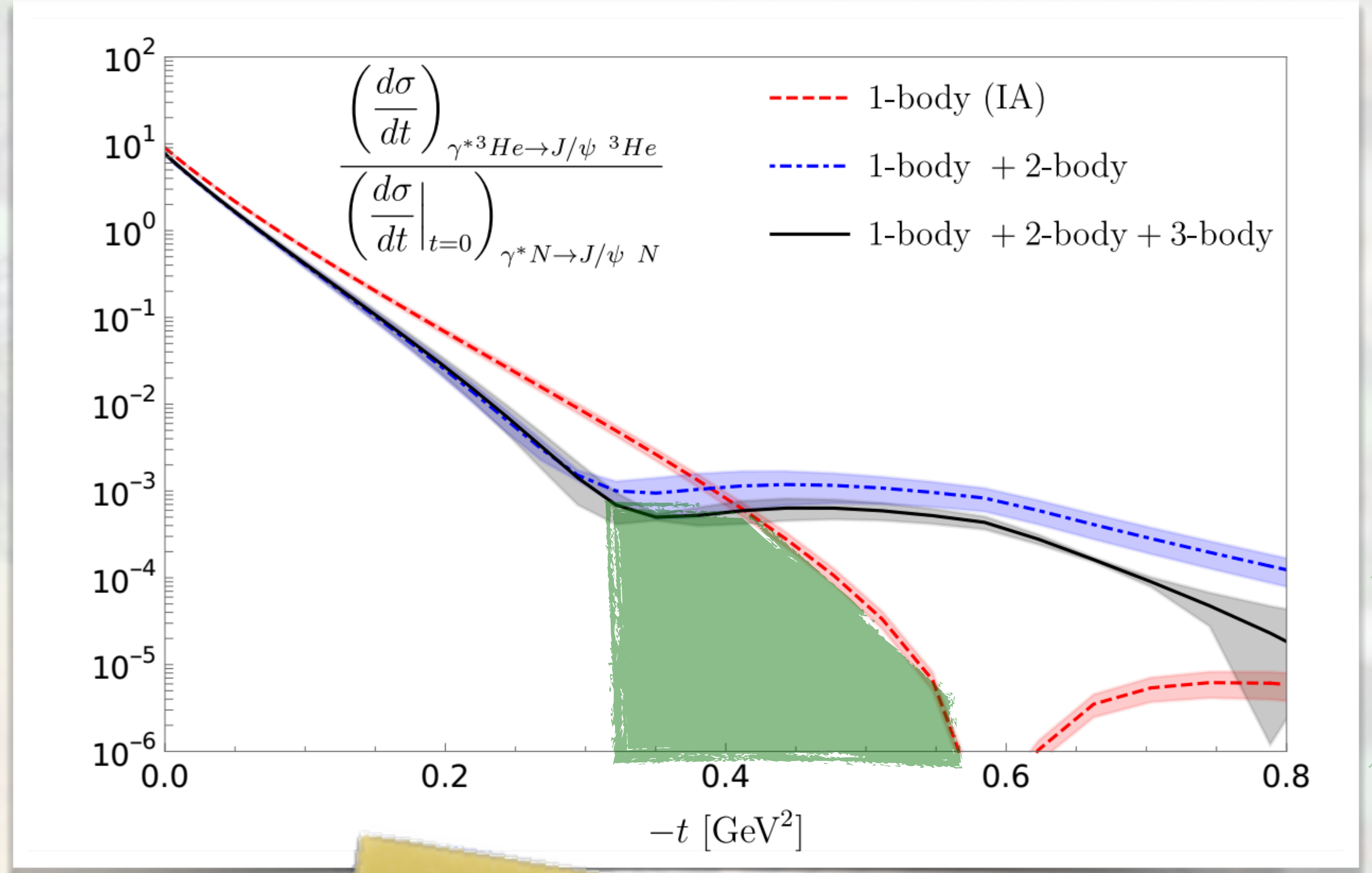
-15 of variation in < σ^2 >

=Multi parton interactions effects

- ✔ 1-body + 2-body re-scatterings dominate the cross-sections shift of the minimum due to 2-body dynamics
- ✔ 1-body dynamics under theoretical control: very good chances to disentangle
- ✔ 2-body dynamics (LT gluon shadowing)
- ✔ unique opportunity to access the real part of the scattering amplitudes in a wide range of t
- ✔ The position of the minimum is extremely sensitive to dynamics and the structure!

Results for J/Ψ exclusive production @EIC: xB ≈ 10-3





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=Multi parton interactions effects

PRECISE MEASUREMENTS SUCH THOSE @EIC ARE NECESSARY!

A NEW LINK BETWEEN NUCLEAR DPS PHYSICS AND DIFFRACTIVE PROCESSES

- ✔ 1 … nate the cross-sections shift of the …
- ✔ … control: very good chances to dise…
- ✔ … ving)
- ✔ … eal part of the scattering amplitu…
- ✔ … xtremely sensitive to dynamics a…

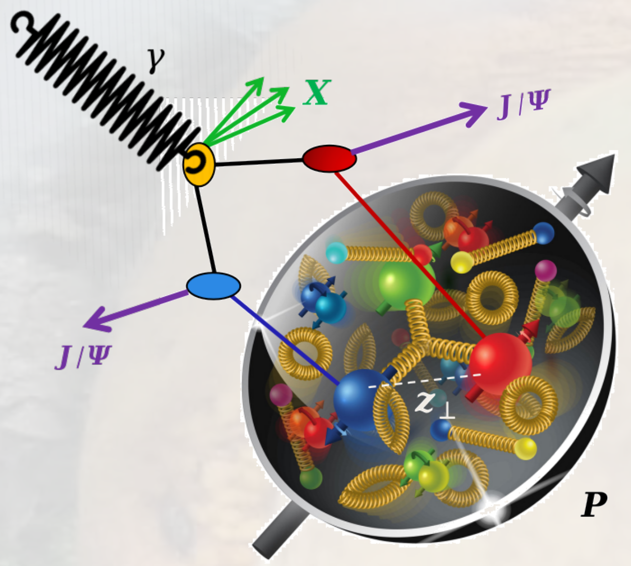
CONCLUSIONS

1) We demonstrated DPS represents a new way to access new information of hadrons

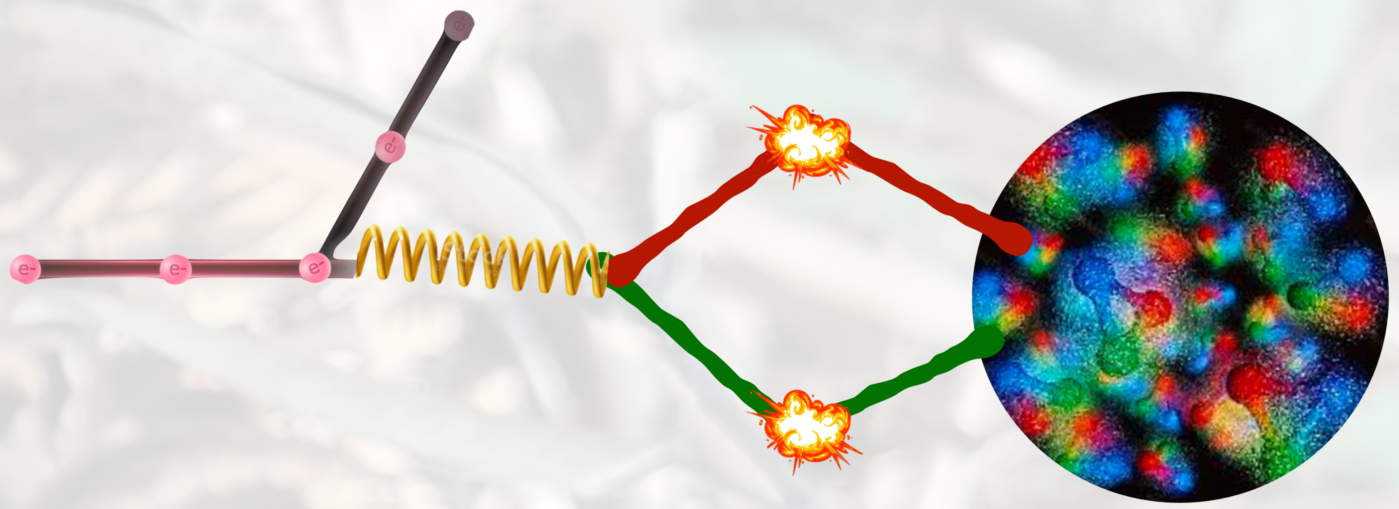
2) Several experimental analyses and theoretical developments are on going

3) We proposed to consider DPS initiated via photon-proton interactions:

a) DPS@EIC



b) Nuclear DPS@EIC

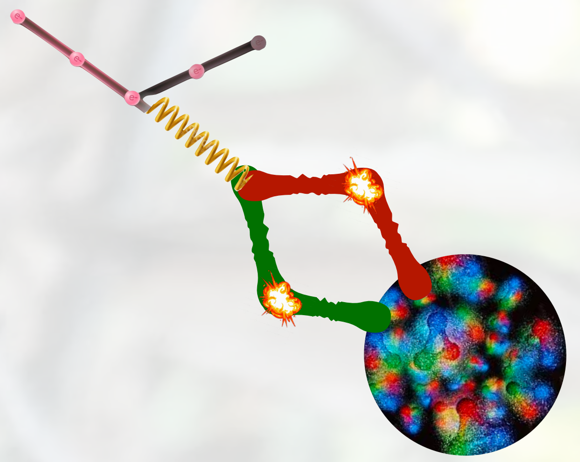


a) DPS contributes, in particular in the 4-jets photoproduction
b) We have estimated SPS and DPS cross sections for quarkonium-pair photoproduction at the EIC using the NRQCD framework
c) The dependence of $\sigma_{\rm eff}^{\gamma p}(Q^2)$ on Q^2 can unveil the mean distance of partons in the proton
d) Quarkonium-pair photoproduction is a promising channel to probe the gluonic content of the photon structure

DPS in γA collisions with light nuclei?

M.R. in progress

For example in DPS1: $\tilde{F}^{1}_{a_1a_2}(x_1, x_2, k_\perp) = \sum_{N=p,n} \int \frac{1}{\xi} \tilde{F}^{N}_{a_1a_2}\left(\frac{x_1}{\xi}, \frac{x_2}{\xi}, k_\perp\right) \boxed{\rho^N_A(\xi, p_{t,N})} \frac{d\xi}{\xi} d^2p_{t,N}$



Let us check sum rules:

$$\int_0^1 dx_1 \int_0^{1-x_1} dx_2 F_{i_1i_2}(x_1, x_2, k_\perp = 0) = \begin{cases} N_{i_1}N_{i_2} & \text{for } i_1 \neq i_2 \\ \left(N_{i_1} - 1\right) N_{i_2} & \text{for } i_1 = i_2 \end{cases}$$

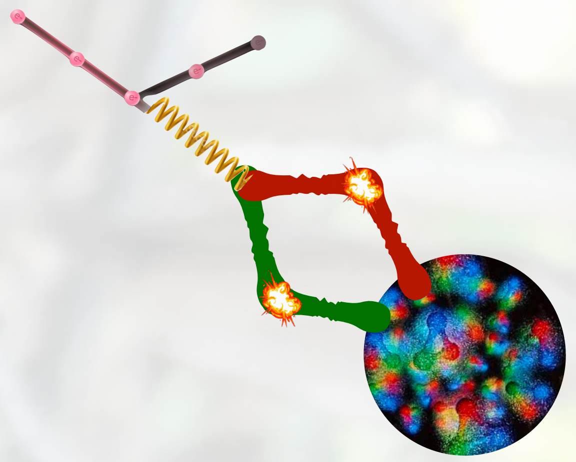
Gaunt's sum rules

J. R. Gaunt and W. J. Stirling, JHEP 03, 005 (2010)

DPS in γA collisions with light nuclei?

M.R. in progress

For example in DPS1: $\tilde{F}^{1}_{a_1 a_2}(x_1, x_2, k_\perp) = \sum_{N=p,n} \int \frac{1}{\xi} \tilde{F}^{N}_{a_1 a_2}\left(\frac{x_1}{\xi}, \frac{x_2}{\xi}, k_\perp\right) \boxed{\rho^N_A(\xi, p_{t,N})} \frac{d\xi}{\xi} d^2 p_{t,N}$



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Gaunt's sum rules

J. R. Gaunt and W. J. Stirling, JHEP 03, 005 (2010)

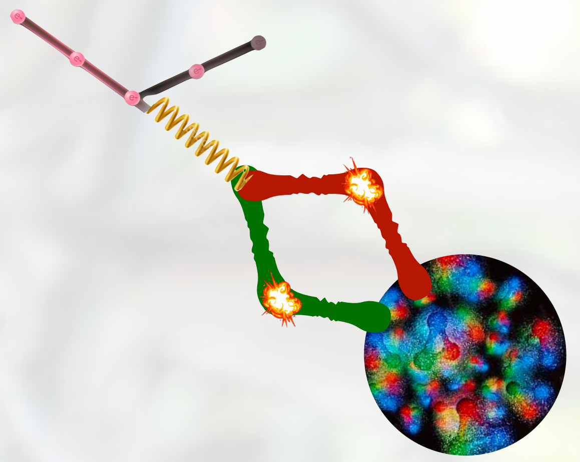
However for the nuclear case one needs also the DPS2

Thus we can introduce approximated partial sum rules (APSR)

DPS in γA collisions with light nuclei?

M.R. in progress

For example in DPS1: $$\tilde{F}^1_{a_1a_2}(x_1,x_2,k_\perp) = \sum_{N=p,n}\int \frac{1}{\xi}\tilde{F}^N_{a_1a_2}\left(\frac{x_1}{\xi},\frac{x_2}{\xi},k_\perp\right)\rho^N_A(\xi,p_{t,N})\frac{d\xi}{\xi}d^2p_{t,N}$$



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Gaunt's sum rules

J. R. Gaunt and W. J. Stirling, JHEP 03, 005 (2010)

APSR: Since $f^A_n(\xi) = \int d^2p_{t,N}\,\rho^N_A(\xi,p_{t,N})$ is peaked around 1/A

$$\int_0^A dx_1\int_0^{A-x_1}dx_2\tilde{F}^{A,1}_{i_1i_2}(x_1,x_2,0) \sim \sum_{n=N,P}\int d\xi f^A_n(\xi)\begin{cases}\left(N^n_{i_1}-1\right)N^n_{i_2} & i_1=i_2\\ N^n_{i_1}N^n_{i_2} & i_1\neq i_2\end{cases}$$

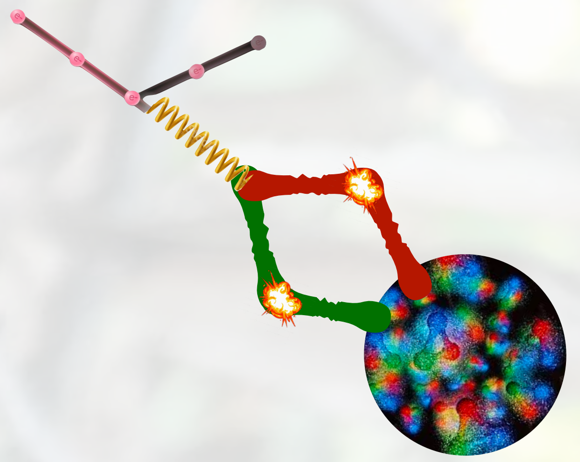
Normalized to 1

Gaunt's sum rules for the nucleon DPD: numbers of quarks with given flavor i in the nucleon n

DPS in γA collisions with light nuclei?

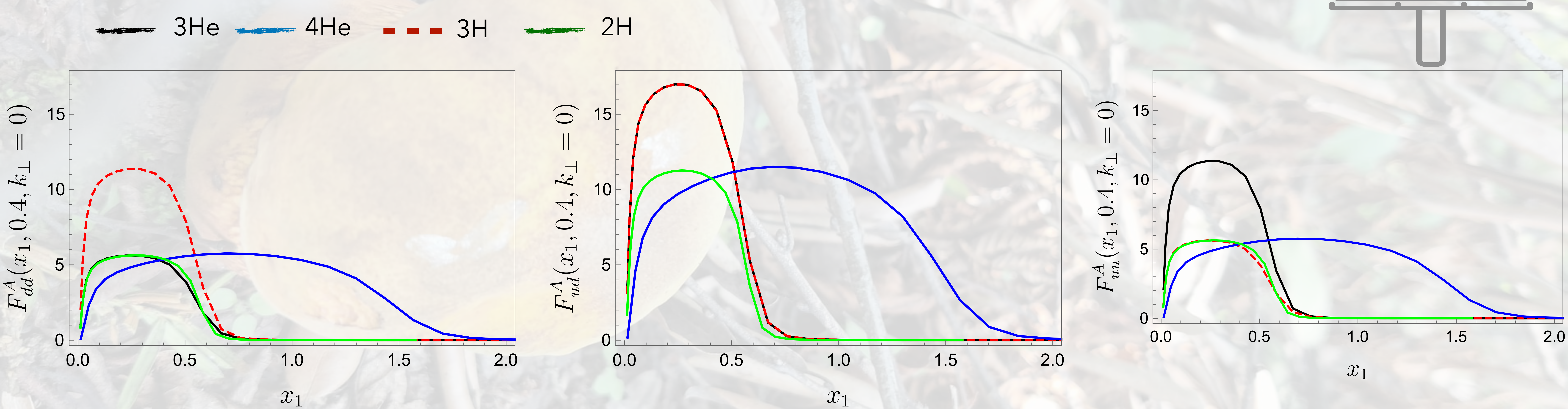
M.R. in progress

For example in DPS1: $$\tilde{F}^{1}_{a_1 a_2}(x_1, x_2, k_\perp) = \sum_{N=p,n} \int \frac{1}{\xi} \tilde{F}^{N}_{a_1 a_2}\left(\frac{x_1}{\xi}, \frac{x_2}{\xi}, k_\perp\right) \rho^{N}_{A}(\xi, p_{t,N}) \frac{d\xi}{\xi} d^2 p_{t,N}$$



- Only for valence quarks. We used a relativistic constituent quark model for the nucleon DPD
- verified approximated partial sum rules (numerically)

$0 < x_i < A$



DPS in γA collisions with light nuclei?

M.R. in progress

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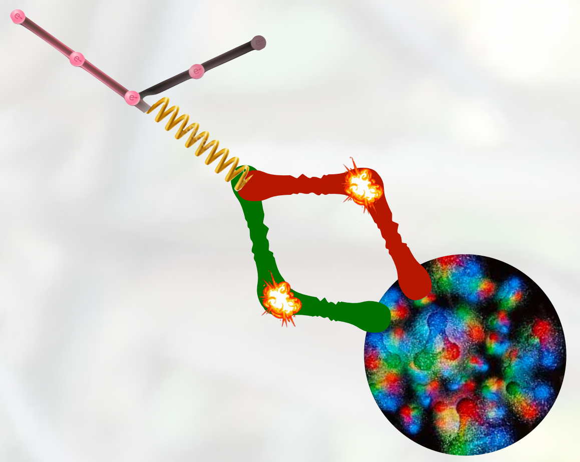

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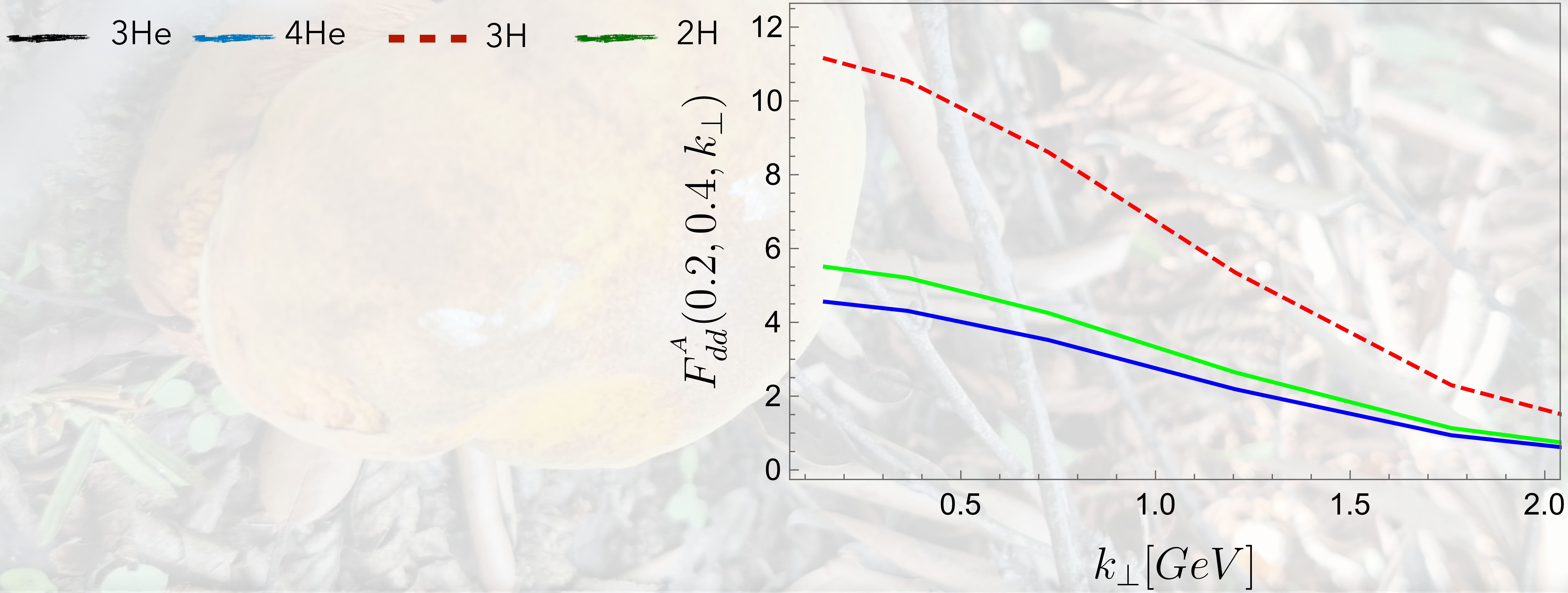


CONCLUSIONS

Backup - $\sigma_{\text{eff}}^{\gamma p}(Q^2 \to \infty)$

1) we show that high virtual behavior of the effective cross sections correctly follows the result in **J.R. Gaunt JHEP 01, 042 (2013)**, i.e.:

$$\sigma_{eff}^{\gamma p}(Q^2 \to \infty) = \sigma_{1v2}^{pp} = \left[\int \frac{d^2 k_\perp}{(2\pi)^2} T_p(k_\perp) \right]^{-1}$$

2) In Ref. **M.Rinaldi and F.A: Ceccopieri JHEP 09, 097 (2019)**, we prove, in a general framework:

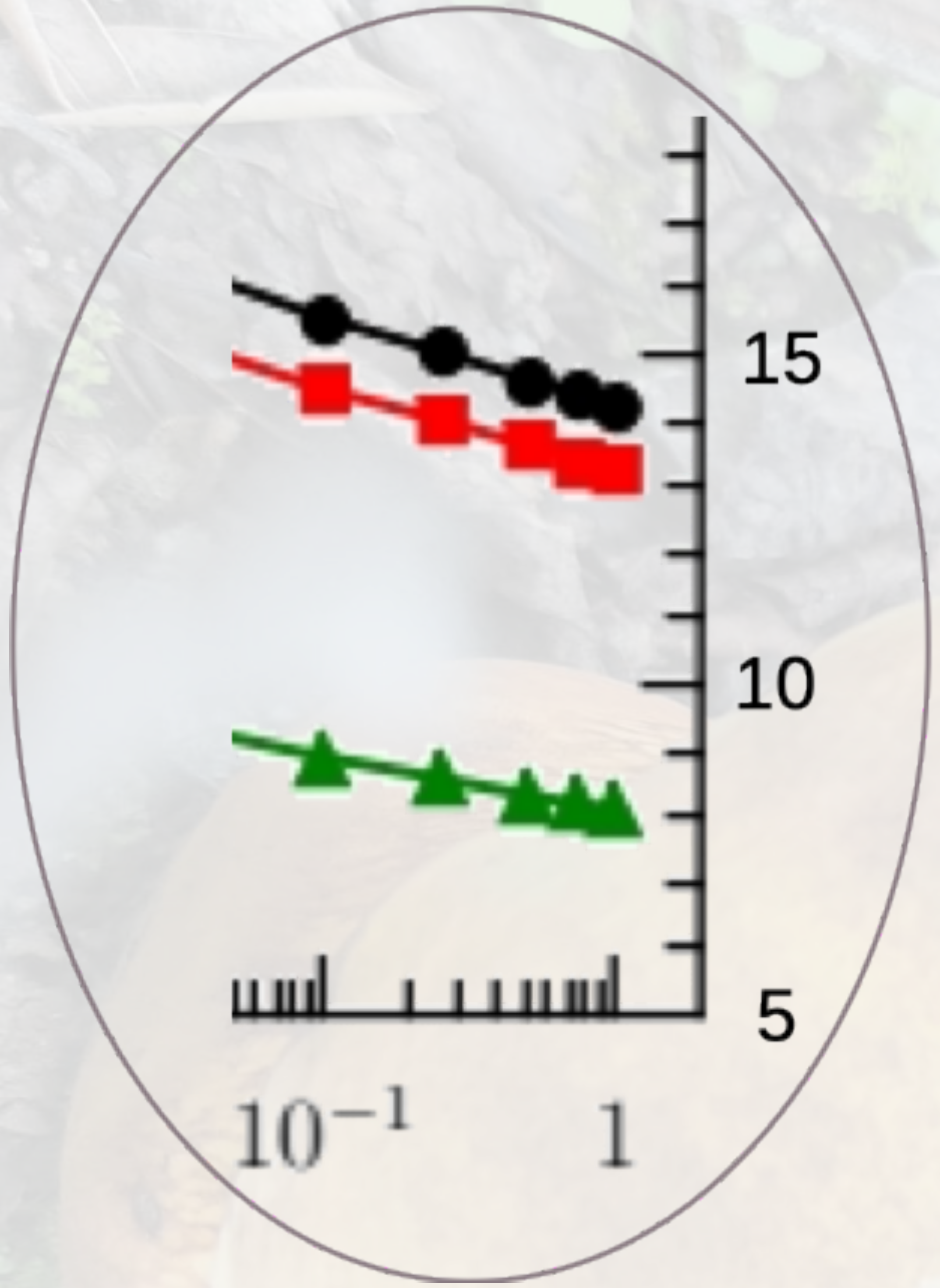
$$\frac{\pi}{2} \langle b^2 \rangle \leq \sigma_{eff}^{\gamma p}(Q^2 \to \infty) \leq 2\pi \langle b^2 \rangle$$

Being: $\sigma_{\text{eff}}^{\gamma p}(Q^2 \to \infty) = \sigma_{\text{eff}}^{2v1}$

Extracted from data

$$\frac{\sigma_{eff}^{pp}}{6} \leq \sigma_{eff}^{\gamma p}(Q^2 \to \infty) \leq 2\sigma_{eff}^{pp}$$

Backup - $\sigma_{\text{eff}}^{\gamma p}(Q^2 \to \infty)$



~ 30/2 mb

~ 25/2 mb

~ 15/2 mb

$$[\sigma_{eff}^{\gamma p}(Q^2)]^{-1} = \int \frac{d^2 k_\perp}{(2\pi)^2} T_p(k_\perp) T_\gamma(k_\perp; Q^2)$$

$$[\sigma_{eff}^{\gamma p}(Q^2)]^{-1} \underset{Q^2>>1}{\sim} \int \frac{d^2 k_\perp}{(2\pi)^2} T_p(k_\perp) \times 1$$

For the proton models we have used:

$$\int \frac{d^2 k_\perp}{(2\pi)^2} T_p(k_\perp) \sim 2 \int \frac{d^2 k_\perp}{(2\pi)^2} T_p(k_\perp)^2$$

$$\sigma_{eff}^{\gamma p}(Q^2 >> 1 \text{ GeV}^2) \sim \sigma_{eff}^{pp}/2$$

Thus for QED: $Q^2 >> 1 \text{ GeV}^2$ almost approximates the asymptotic

DPS in pA collisions

The DPS cross-section

$$d\sigma^{ML}_{DPS} = \frac{m}{2} \sum_{i,j,k,l} d\hat{\sigma}^{M}_{ik} d\hat{\sigma}^{L}_{jl} \int d^2 b_\perp \, F^{ij}_{p}(x_1, x_2, \vec{b}_\perp) \int d^2 B \Bigg\{$$

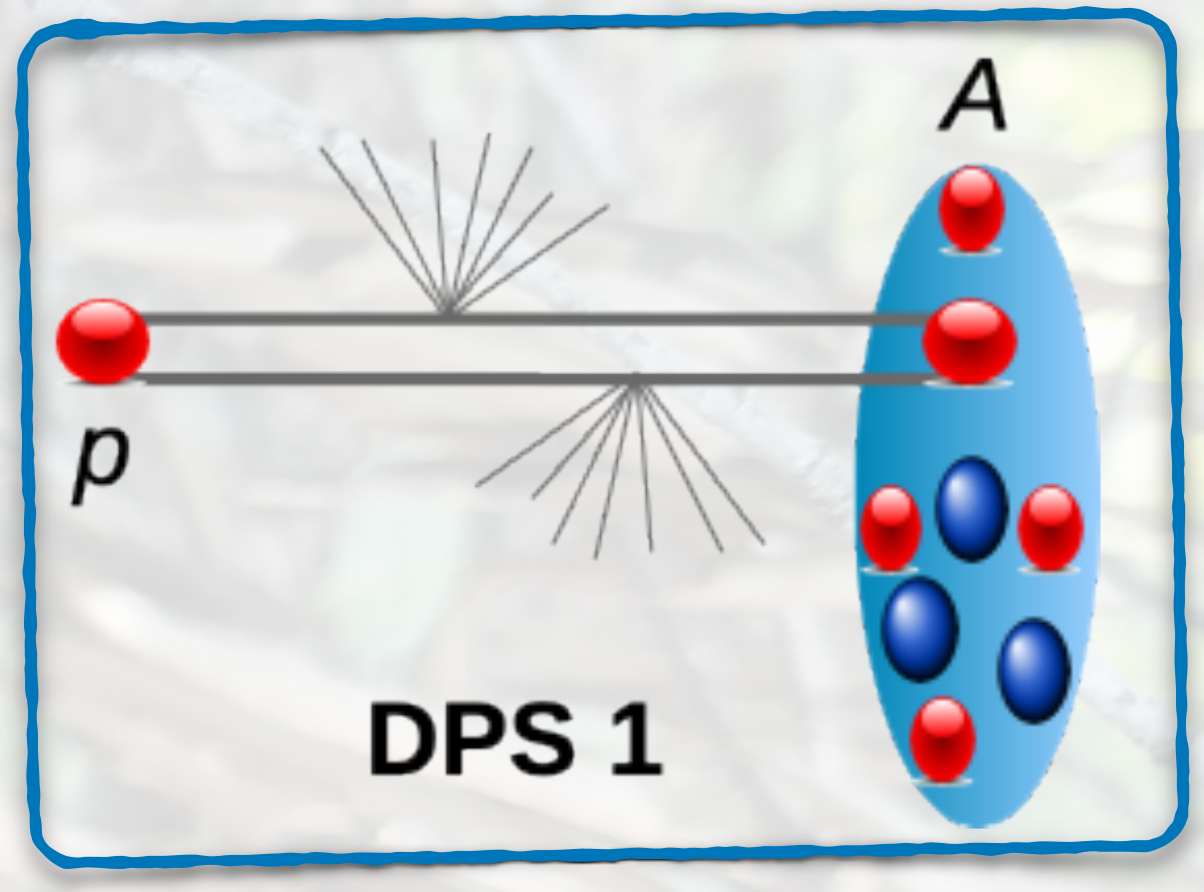
the thickness function as a function of the impact parameter B

$$\bar{T}(\vec{b}_\perp + \vec{B}) \sim \bar{T}(\vec{B})$$

$$\bar{T}_N(B) = \int dz \, \underbrace{\rho_N}(\sqrt{B^2 + z^2})$$

Wood-Saxon distribution for pb normalized to A

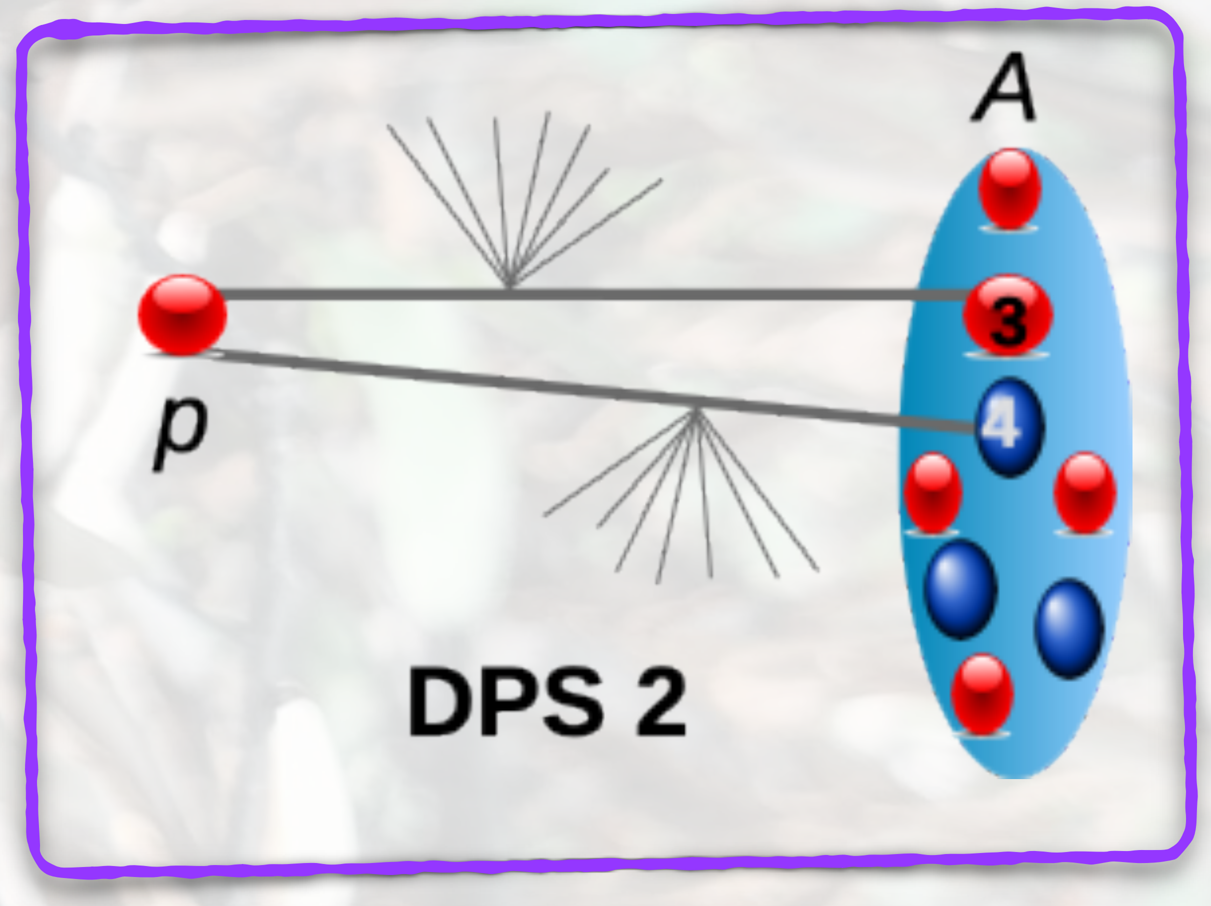
$$\sum_{N=p,n} F^{kl}_{N}(x_3, x_4, \vec{b}_\perp) \bar{T}_N(B)$$



$$+$$

$$\sum_{N_3, N_4 = p,n} f^{k}_{N_3/A}(x_3) f^{l}_{N_4/A}(x_4) T_{N_3}(B) T_{N_4}(B)$$

$$\Bigg\}$$

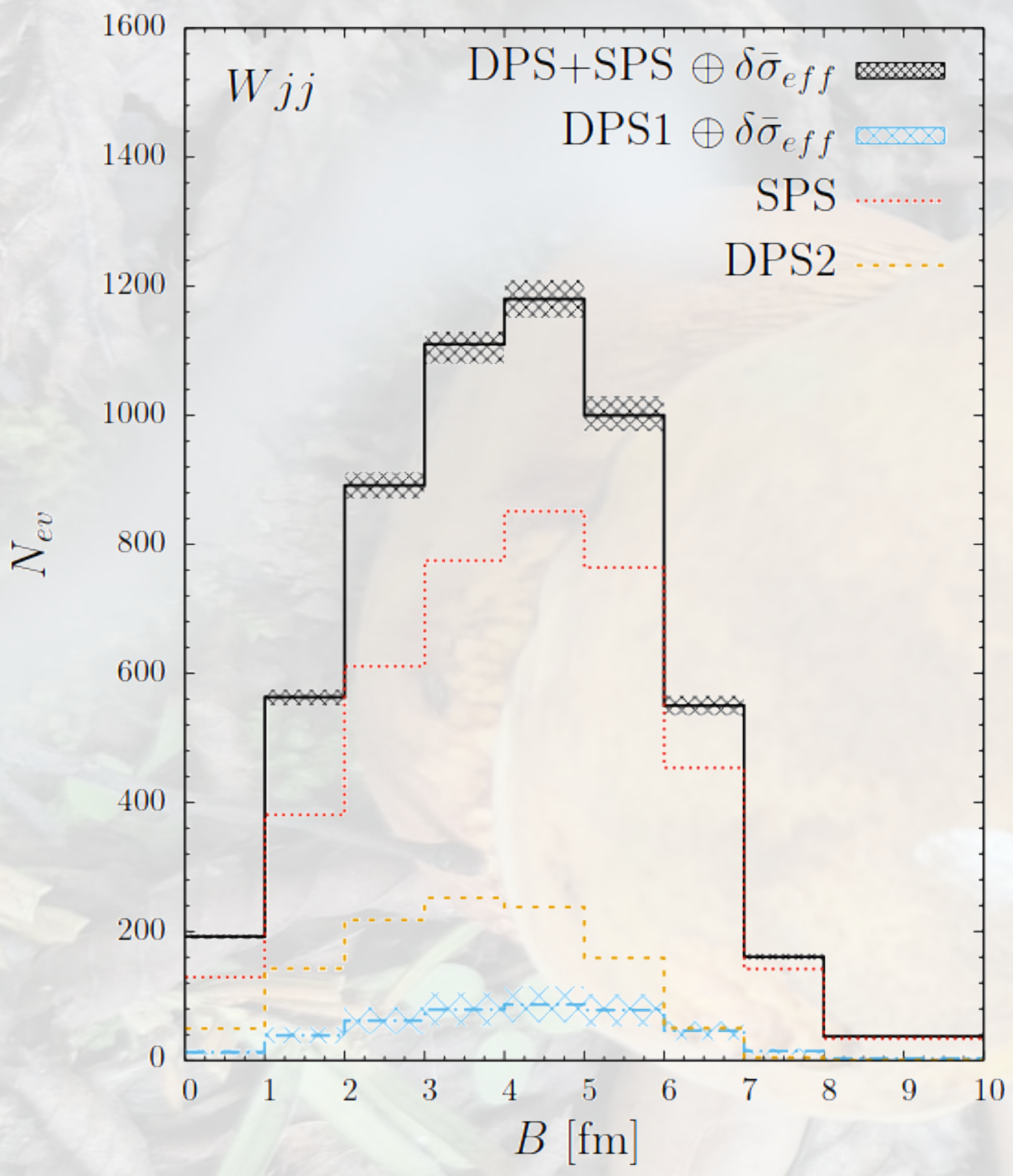


DPS in pA collisions

Some examples of predictions:

W+di-jets

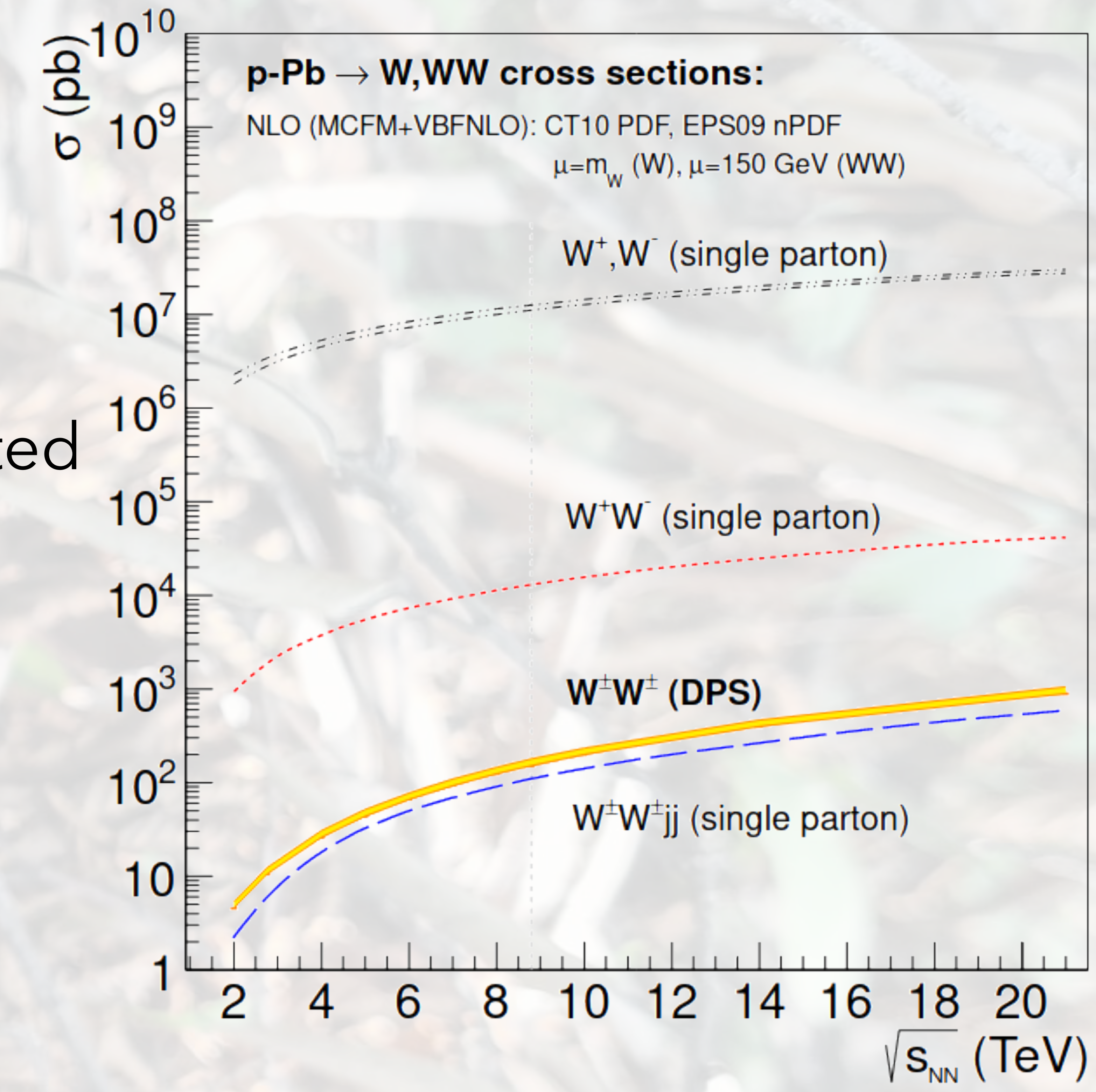
B. Blok and F. A. Ceccopieri EPJC (2020) 80, 278



- SPS dominant
- DPS2 bigger then DPS1 has expected

Same sign WW

D. D'Enterria and Snigirev, PLB 718 (2013) 1395-1400



DPS in γA collisions with light nuclei?

M.R. in progress

For example in DPS2:

$$\tilde{F}^2_{a_1a_2}(x_1, x_2, \vec{k}_\perp) \propto \int \frac{1}{\xi_1\xi_2} \prod_{i=1}^{i=A} \frac{d\xi_i d^2p_{ti}}{\xi_i} \delta\left(\sum_i \xi_i - A\right) \delta^{(2)}\left(\sum_i \mathbf{p}_{ti}\right) \psi^*_A(\xi_1, \xi_2, p_{t1}, p_{t2}) \psi_A\left(\xi_1, \xi_2, p_{t1} + \vec{k}_\perp, p_{t2} - \vec{k}_\perp\right)$$
$$\times\, G^{N_1}_{a_1}\left(\frac{x_1}{\xi_1}, |\vec{k}_\perp|\right) G^{N_2}_{a_2}\left(\frac{x_2}{\xi_2}, |\vec{k}_\perp|\right);$$

if we approximate: $\xi_i \sim 1$ we get:

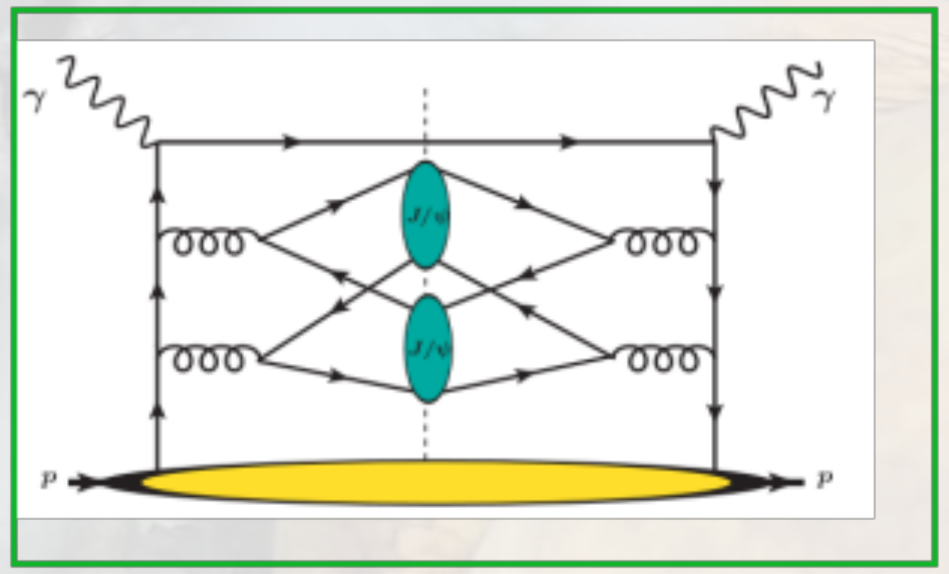
Di J/ψ photo-production@EIC

F. A. Ceccopieri, H. S. Shao, J. P. Lansberg, M. R. and R. Sangem in prep.

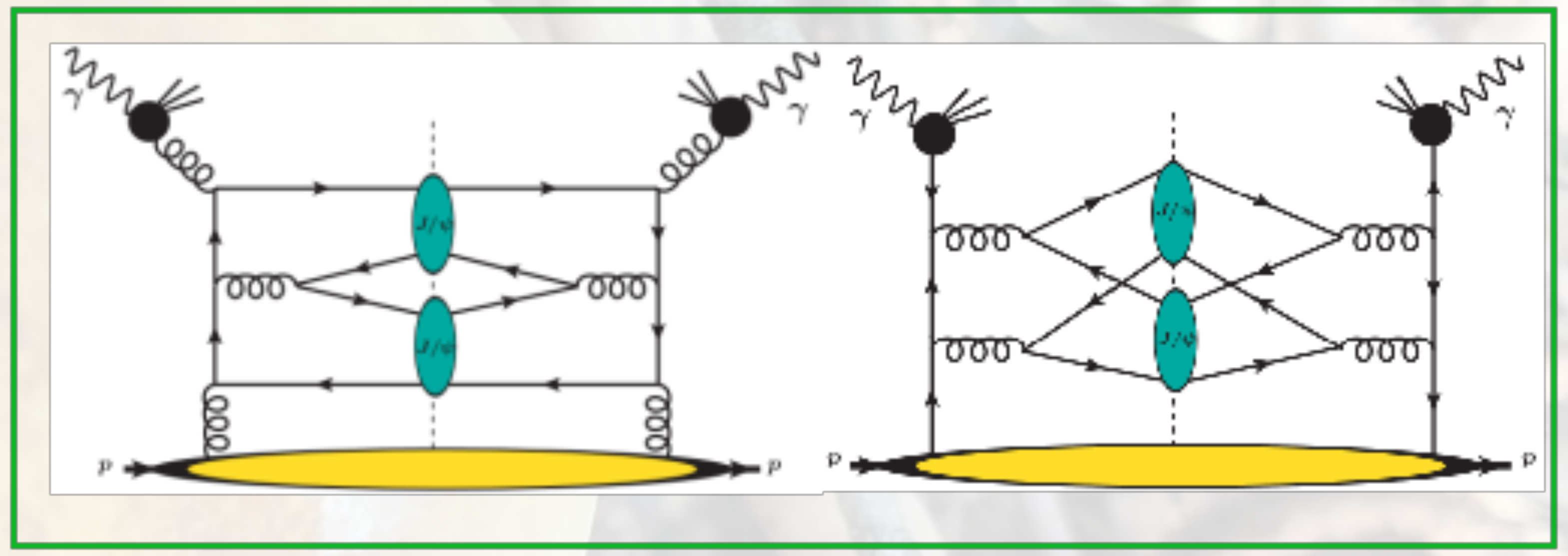
*Slide from R. Sangem

SPS → **Unresolved**, **Resolved**

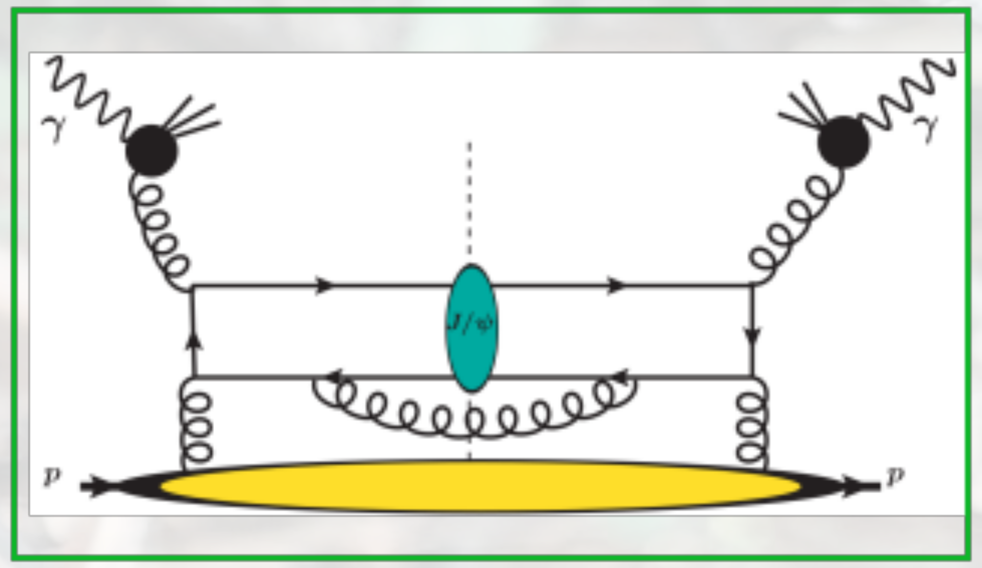
DPS → **Resolved**



$\gamma + q \to J/\psi + J/\psi + q$



$i + j \to J/\psi + J/\psi$ $\quad i, j = g, q$



$|g + g \to J/\psi + g|^2$

Range of cross sections in CSM =100 GeV

$\sigma_{SPS}^{(J/\psi,J/\psi)} \times Br^2 = 4 - 30$ fb

$\sigma_{DPS}^{(J/\psi,J/\psi)} \times Br^2 = 0.2 - 5$ fb

(Resolved) $\sigma_{eff}^{\gamma p} = 10$ mb for DPS

$\sigma_{SPS}^{(J/\psi,J/\psi)} \times Br^2 = 2 - 12$ fb (Unresolved)

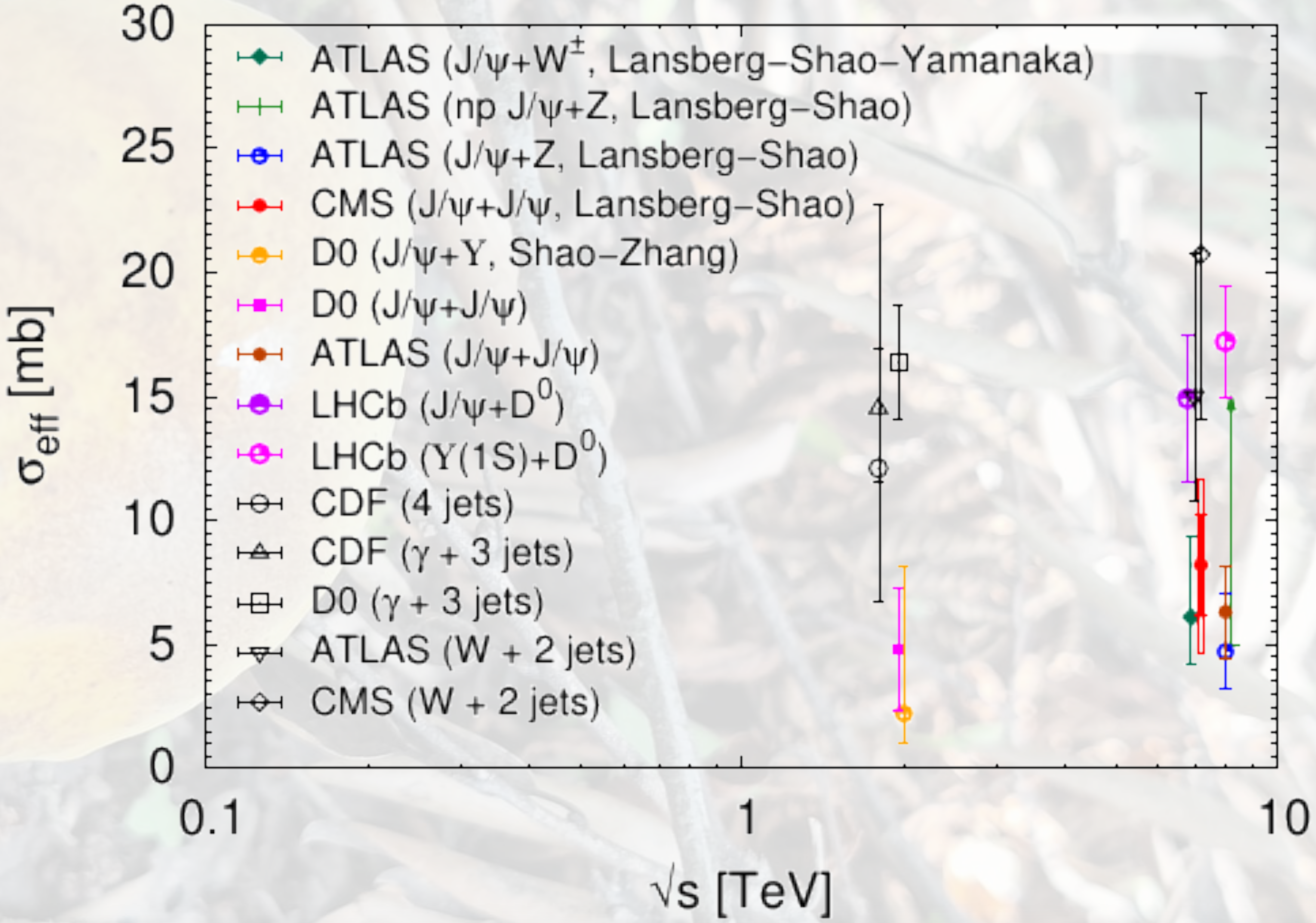
Some Data and Effective Cross Section

$$\sigma_{\text{eff}}^{pp} = \frac{m}{2}\frac{\sigma_A^{pp}\sigma_B^{pp}}{\sigma_{DPS}^{pp}}$$

POCKET FORMULA

Differential X-section single parton scattering for the process: $pp \longrightarrow A(B) + X$

Differential X-section double parton scattering for the process: $pp \longrightarrow A + B + X$



First observation of same sign WW via DPS:

$$\sigma_{\text{eff}} = 12.2^{+2.9}_{-2.2} \quad \text{mb}$$

[CMS coll.], PRL 131 (2023) 091803

$$\sigma^{DPS} \sim 6.28 \text{ fb}$$

Some Data and Effective Cross Section

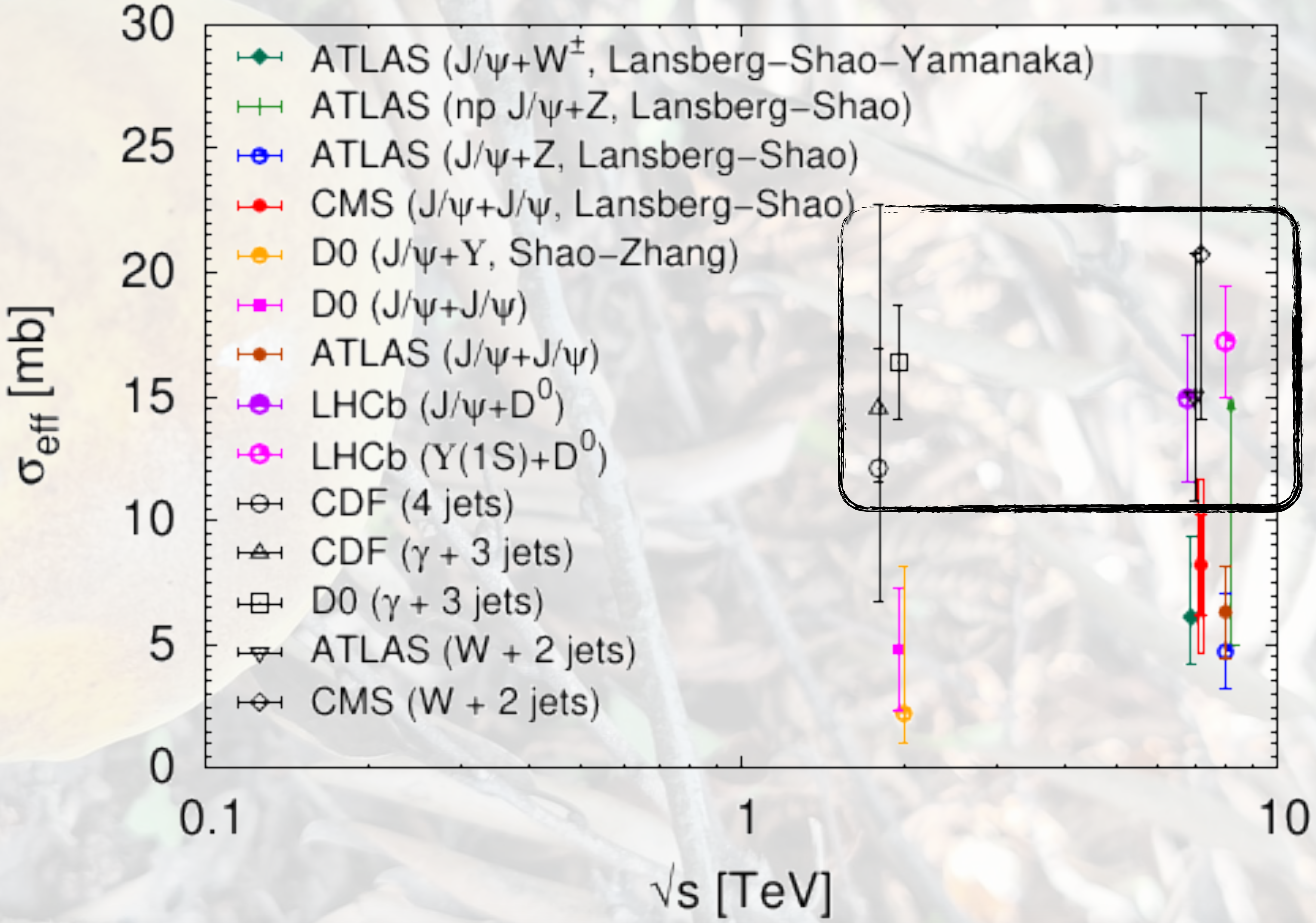
$$\sigma_{eff}^{pp} = \frac{m}{2}\frac{\sigma_A^{pp}\sigma_B^{pp}}{\sigma_{DPS}^{pp}}$$

POCKET FORMULA

Differential X-section single parton scattering for the process: $pp \longrightarrow A(B) + X$

Differential X-section double parton scattering for the process: $pp \longrightarrow A + B + X$

Results for W, Jet productions…



First observation of same sign WW via DPS:

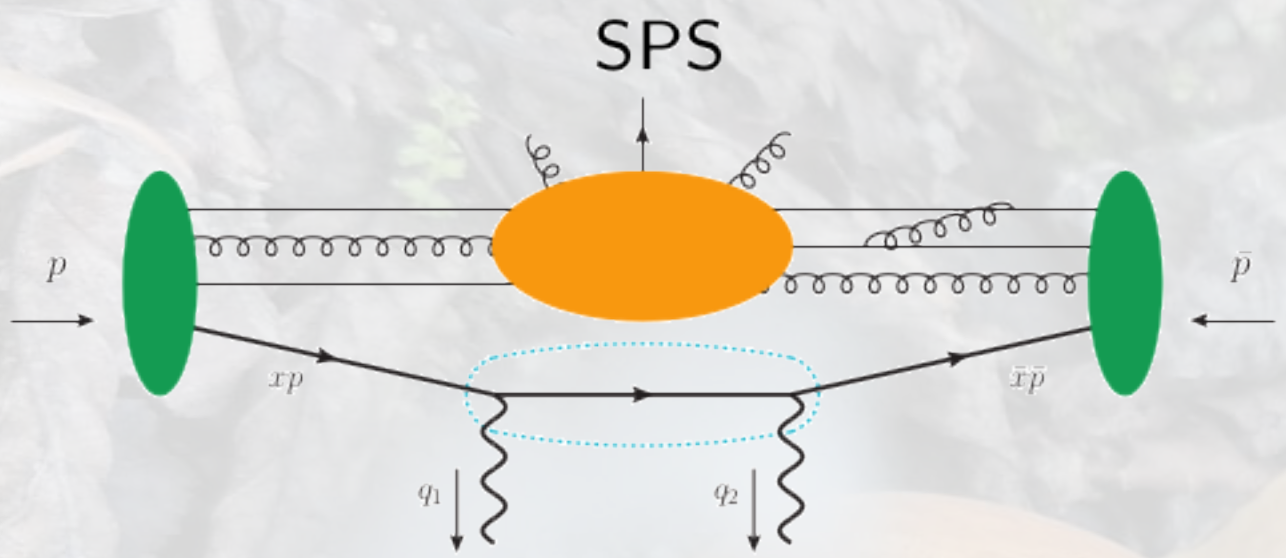
$$\sigma_{eff} = 12.2^{+2.9}_{-2.2} \quad \text{mb}$$

[CMS coll.], PRL 131 (2023) 091803

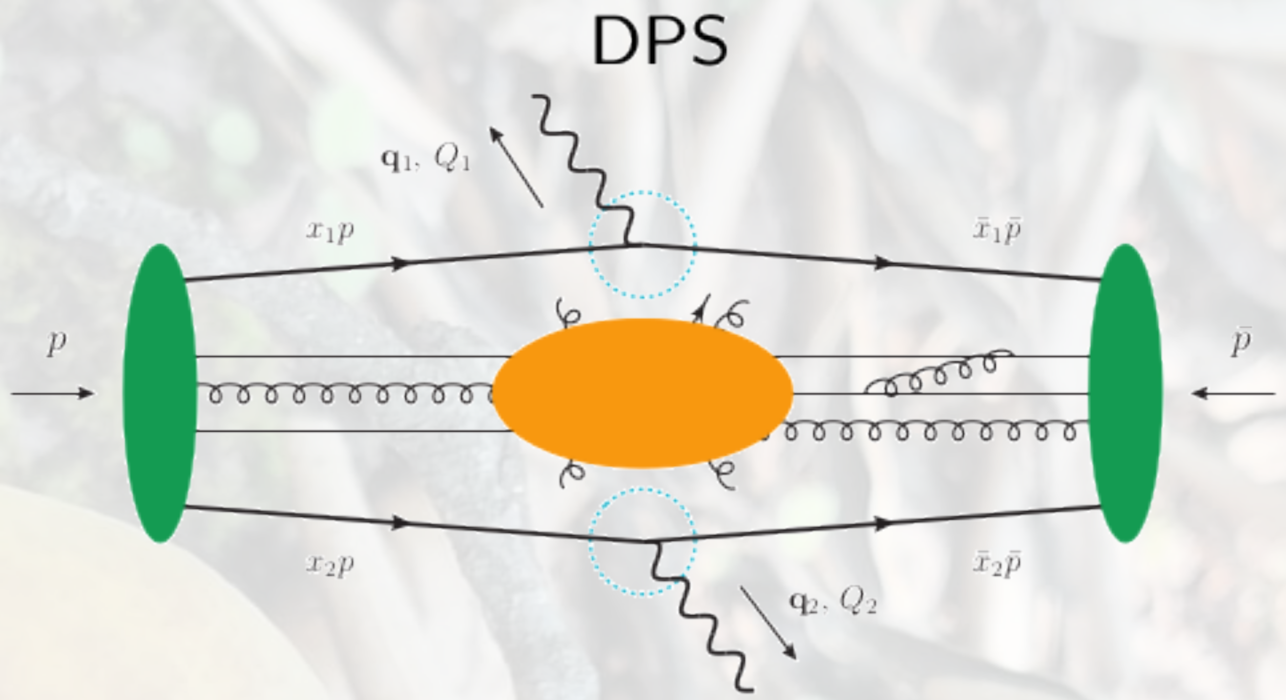
$$\sigma^{DPS} \sim 6.28 \text{ fb}$$

Double Parton Scattering scales

Scale analysis of SPS and DPS processes



$$|\mathbf{q}_1^\perp + \mathbf{q}_2^\perp| \sim \Lambda \ll Q$$



$$|\mathbf{q}_1^\perp| \sim \Lambda \ll Q$$
$$|\mathbf{q}_2^\perp| \sim \Lambda \ll Q$$

where:
- $Q = min(Q_1, Q_2)$
- Λ *transverse momentum scale*
- $\Lambda_{QCD} << \Lambda << Q$

Usually:

$$\frac{d^2\sigma_{SPS}}{d^2q_1\ d^2q_2} \sim \frac{d^2\sigma_{DPS}}{d^2q_1\ d^2q_2}$$

$$\frac{\sigma_{DPS}}{\sigma_{SPS}} \sim \mathcal{O}\left(\frac{\Lambda^2}{Q^2}\right)$$

First appearance in theory studies:
Politzer
Paver, Treleani
Mekhfi

Other ground-setting works:
Gaunt, Stirling
Blok et al.
Diehl et al.
Manohar, Waalewijn
Ryskin, Snigierev
. . .

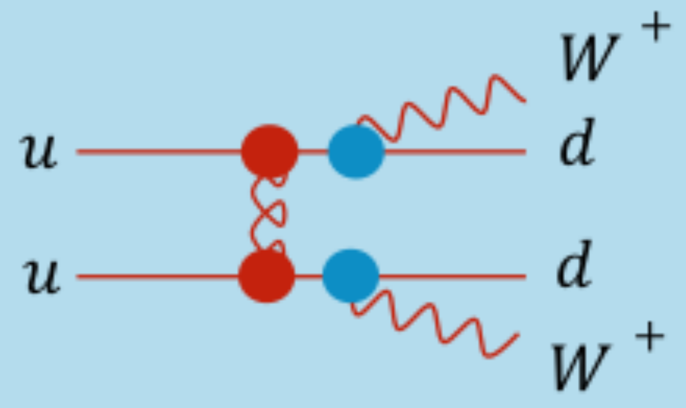
Nagar's slides MPI 2021

Where and Why DPS?

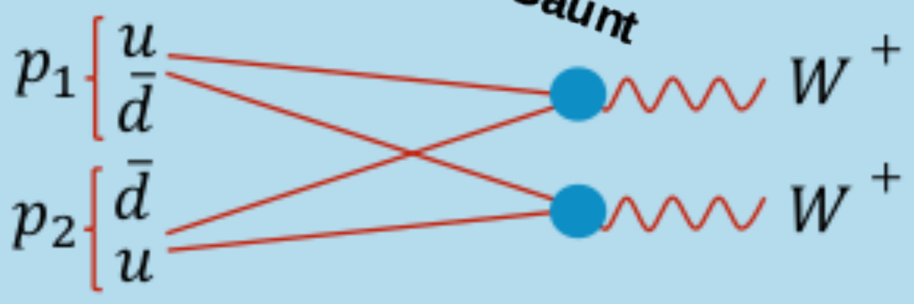
DPS can give a significant contribution to processes where SPS is suppressed by small/multiple coupling constants:

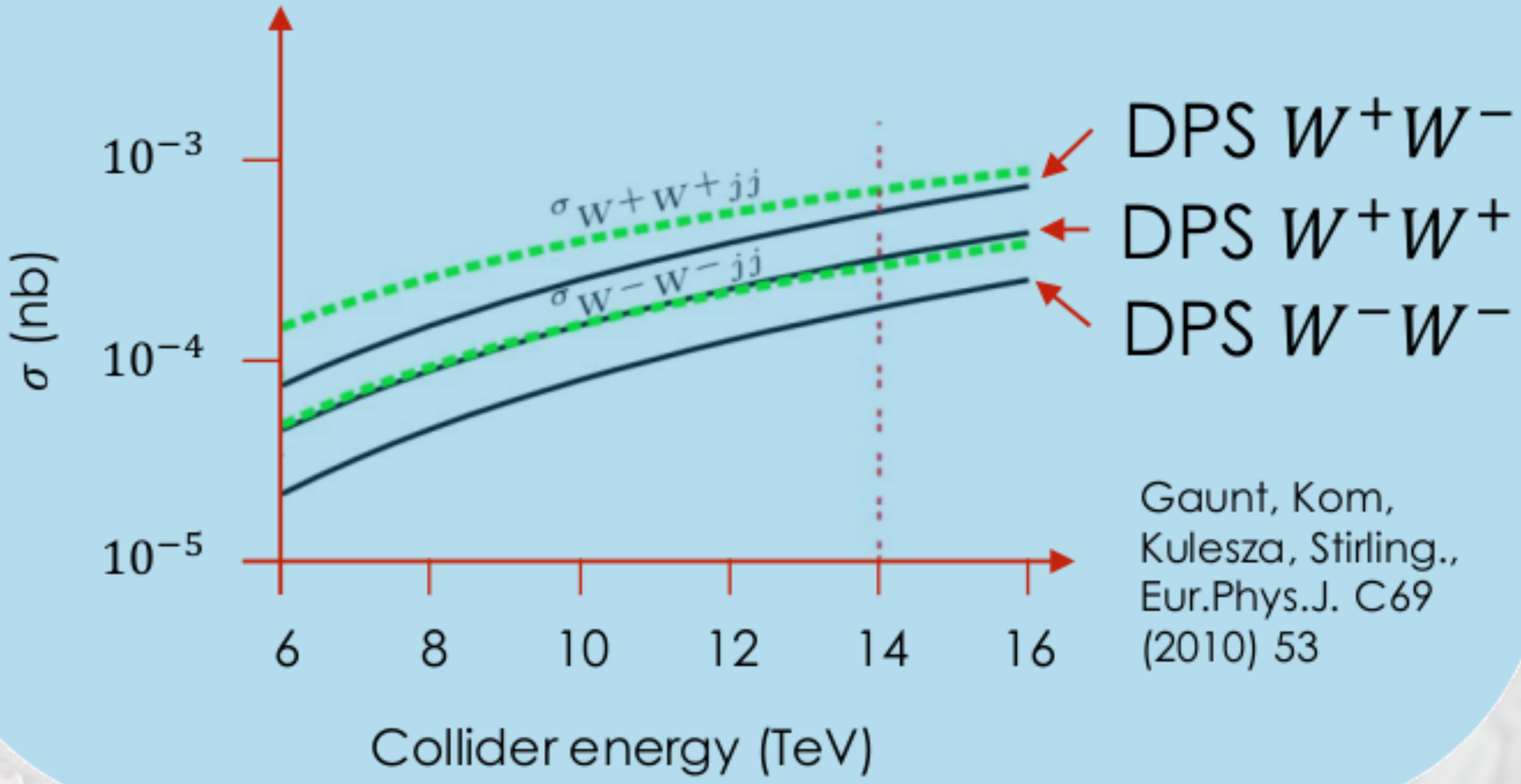
Slide from J. Gaunt

SPS:



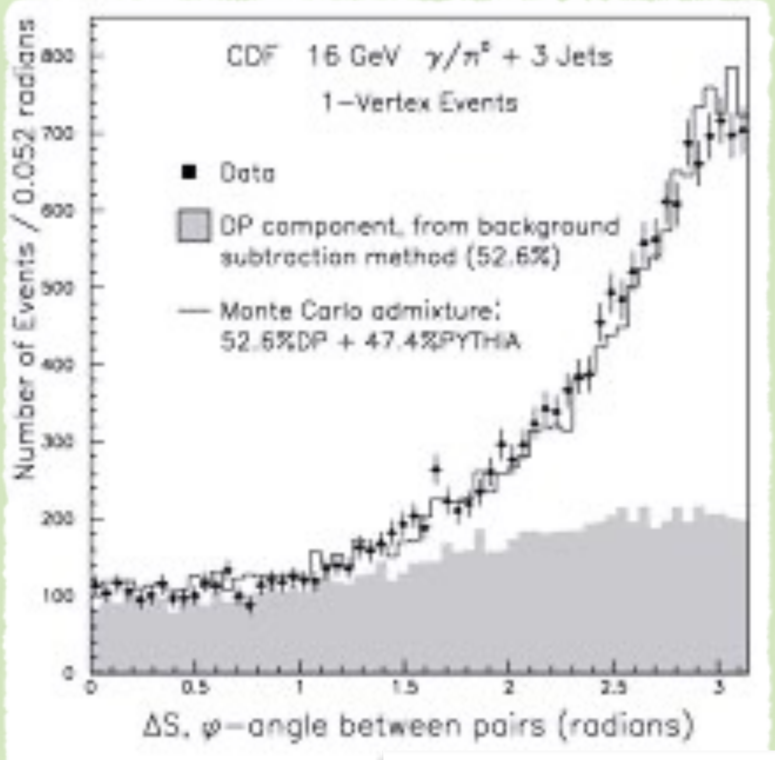
DPS:





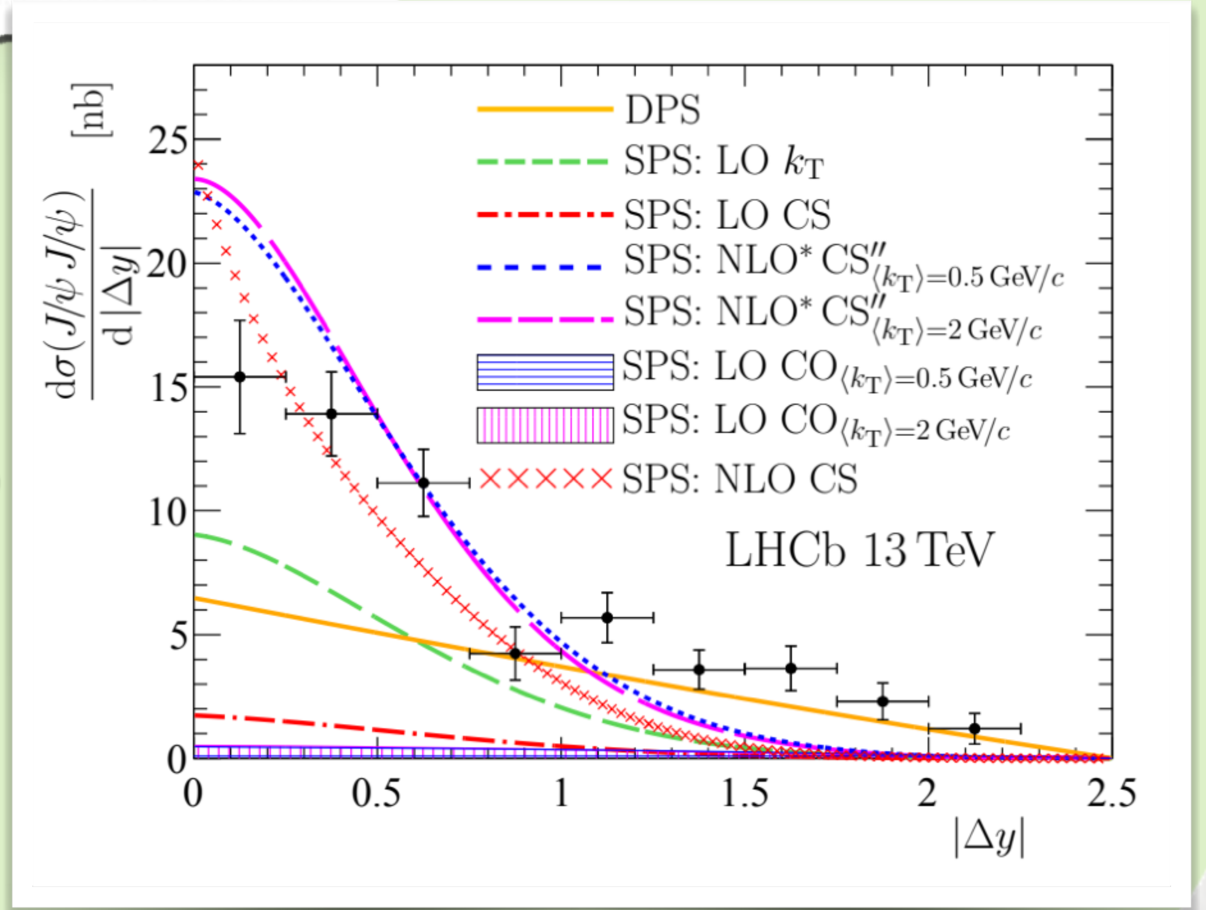
Gaunt, Kom, Kulesza, Stirling., Eur.Phys.J. C69 (2010) 53

…or in certain phase space regions



CDF, $\gamma + 3j$, Phys.Rev. D56 (1997) 3811-3832

LHCb, double J/ψ, JHEP 06, 047, (2017)



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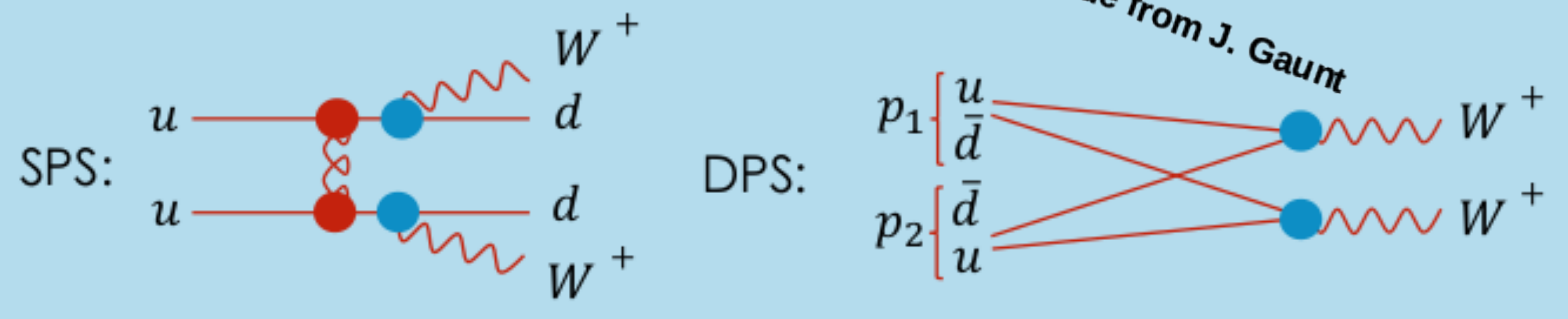
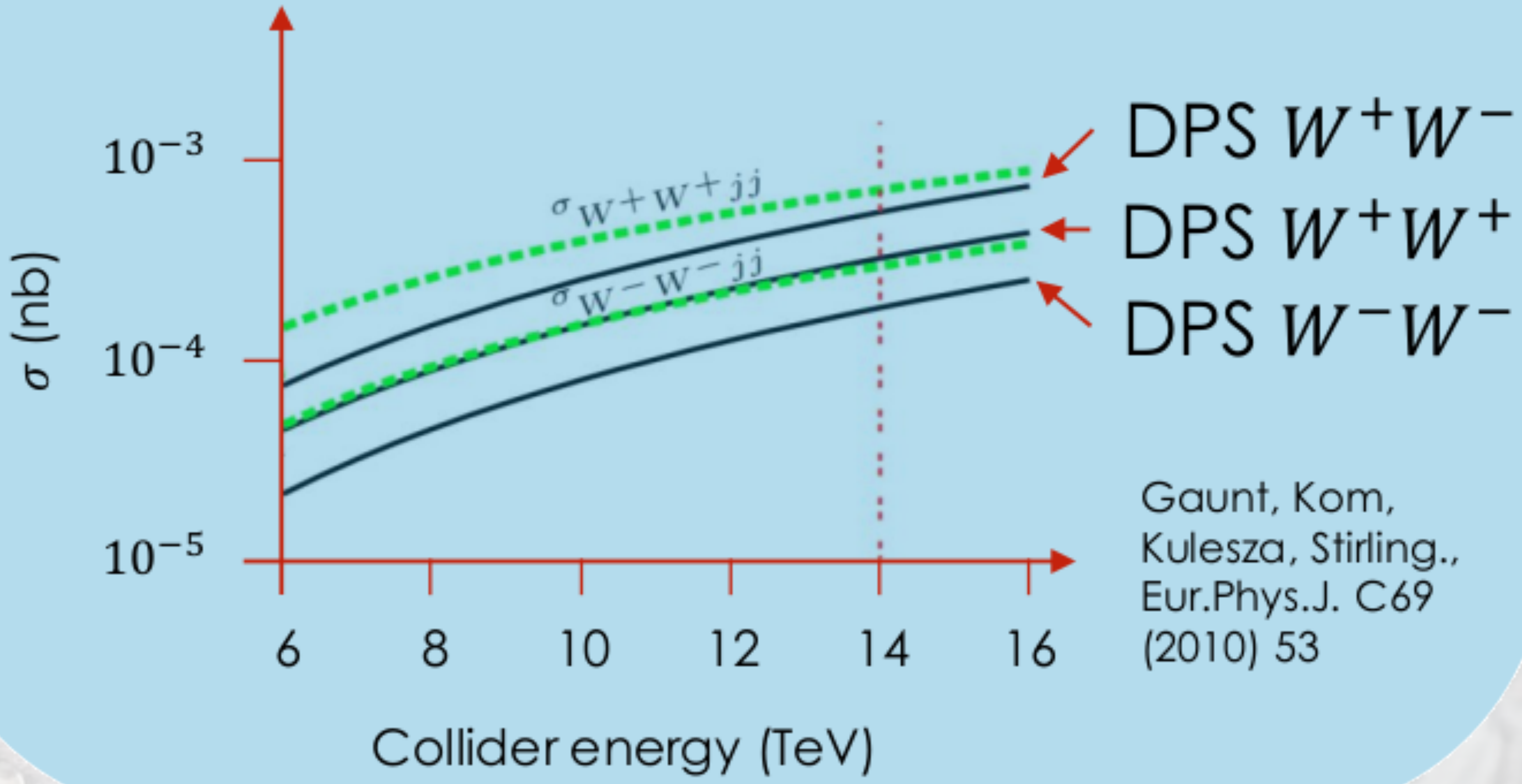
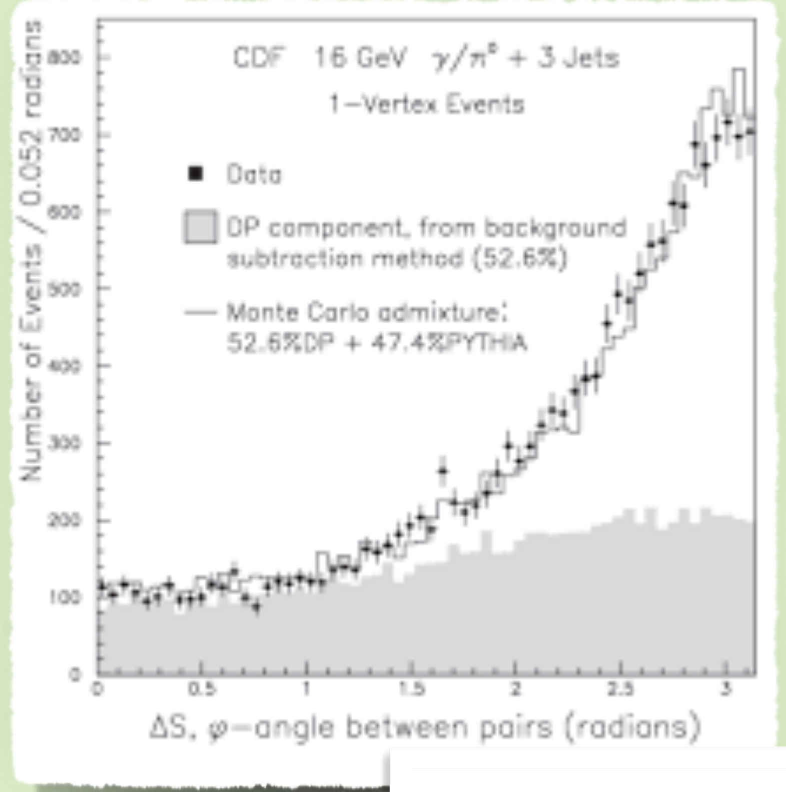
Slide from J. Gaunt





Gaunt, Kom, Kulesza, Stirling., Eur.Phys.J. C69 (2010) 53

…or in certain phase space regions



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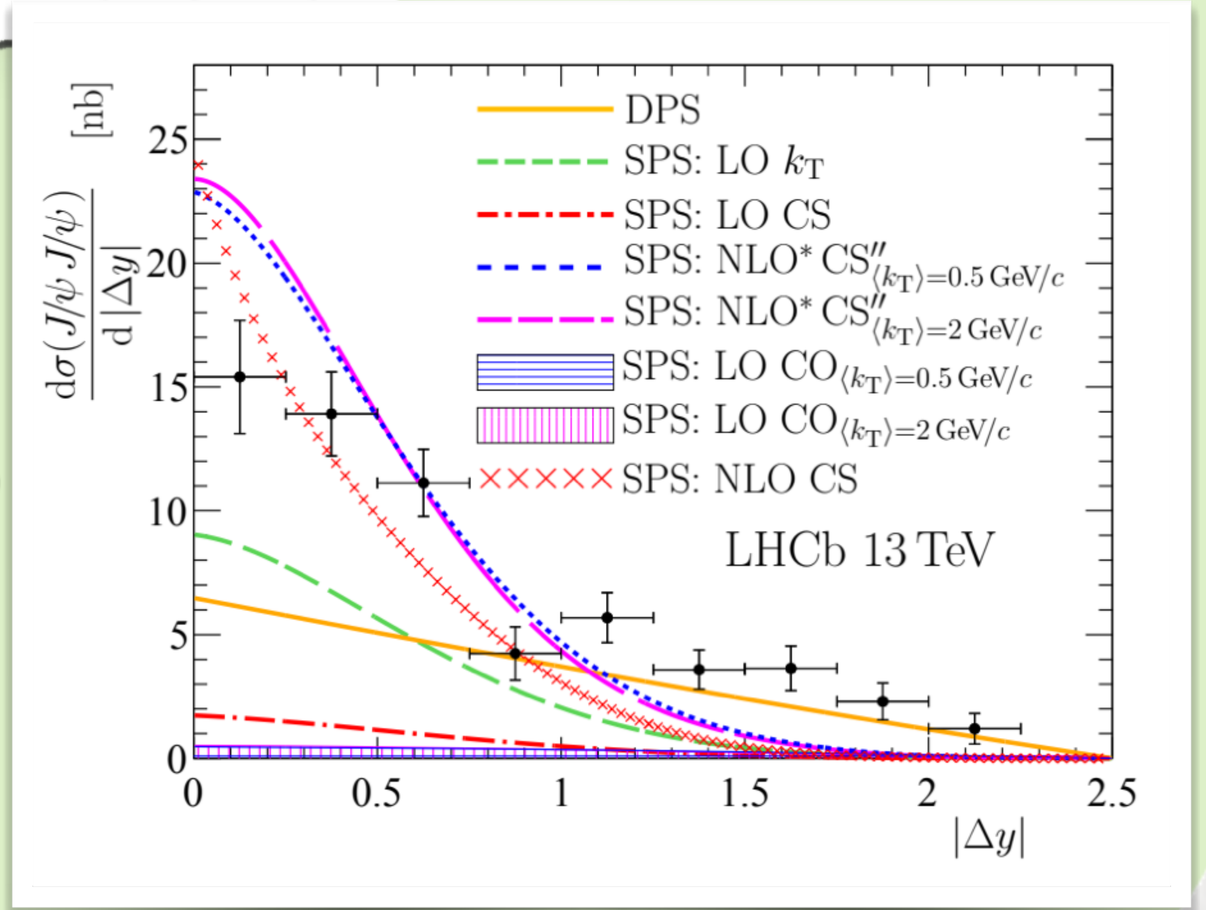
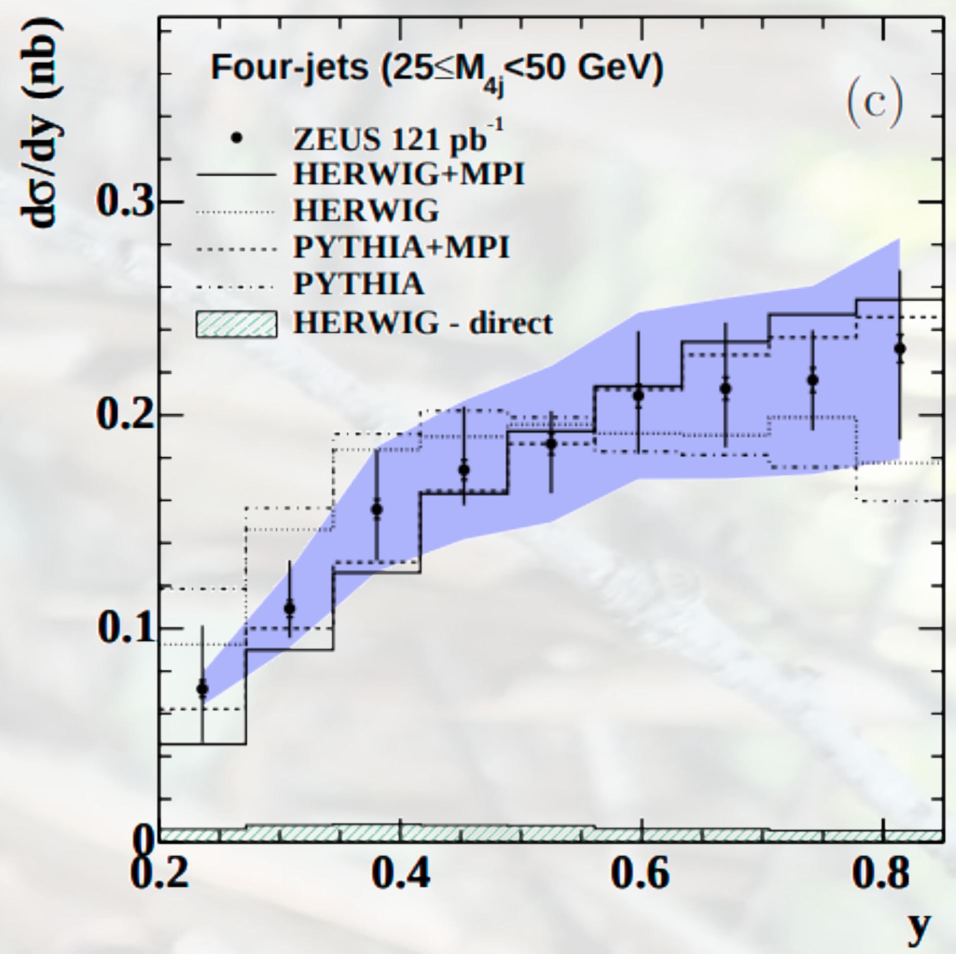


in ep Colliders?



HERA data, ZEUS coll, Nucl. Phys. B 729, 1 (2008)

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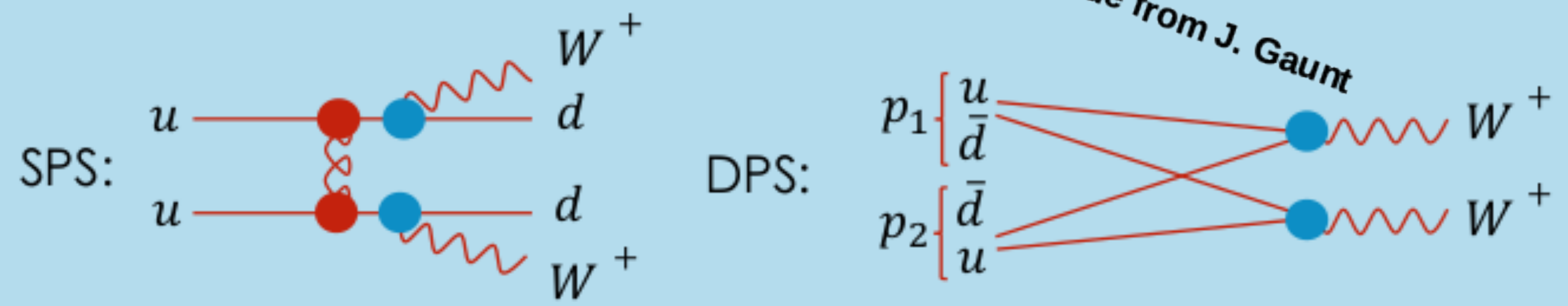

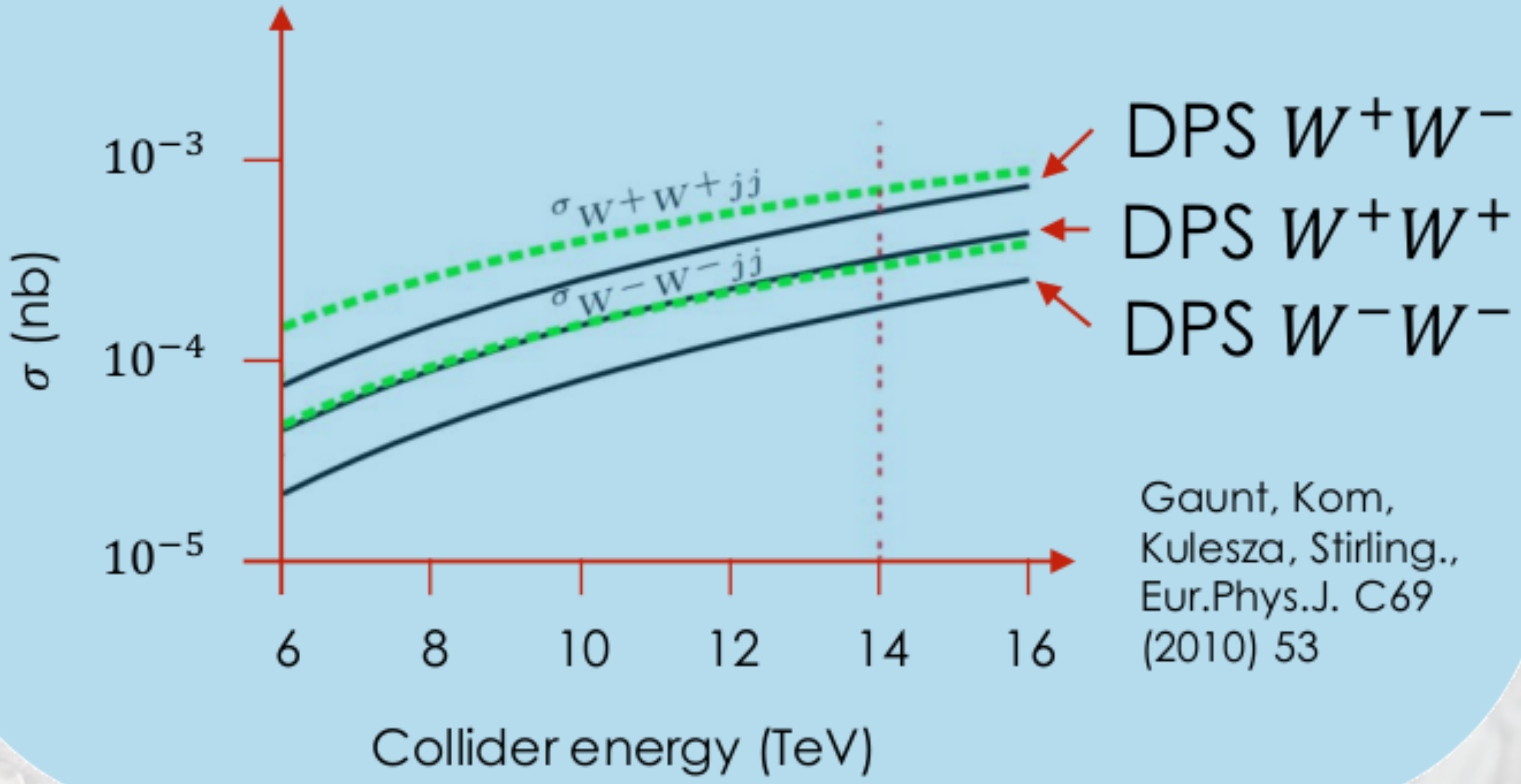

Gaunt, Kom, Kulesza, Stirling., Eur.Phys.J. C69 (2010) 53

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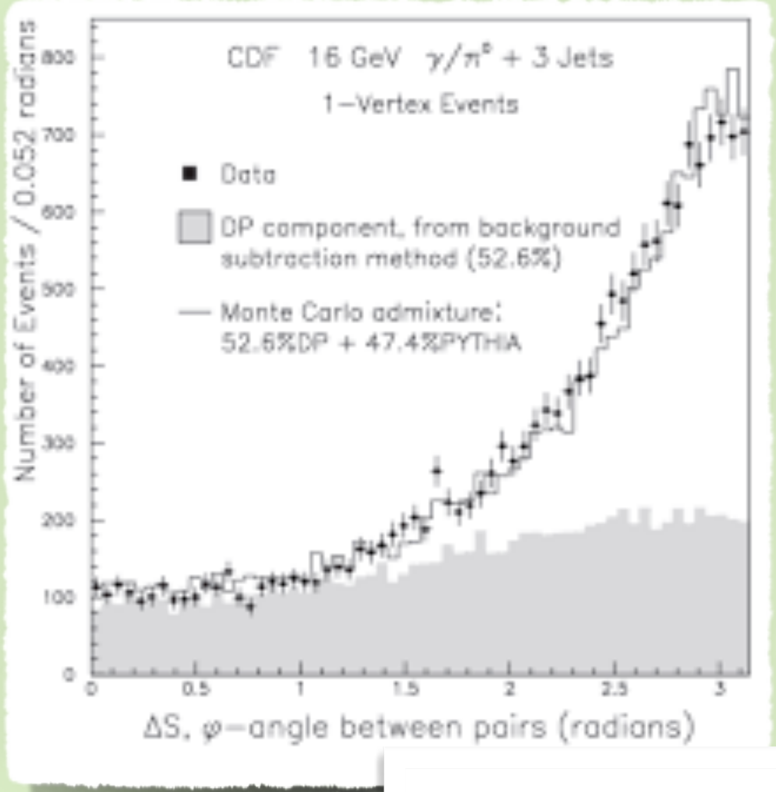
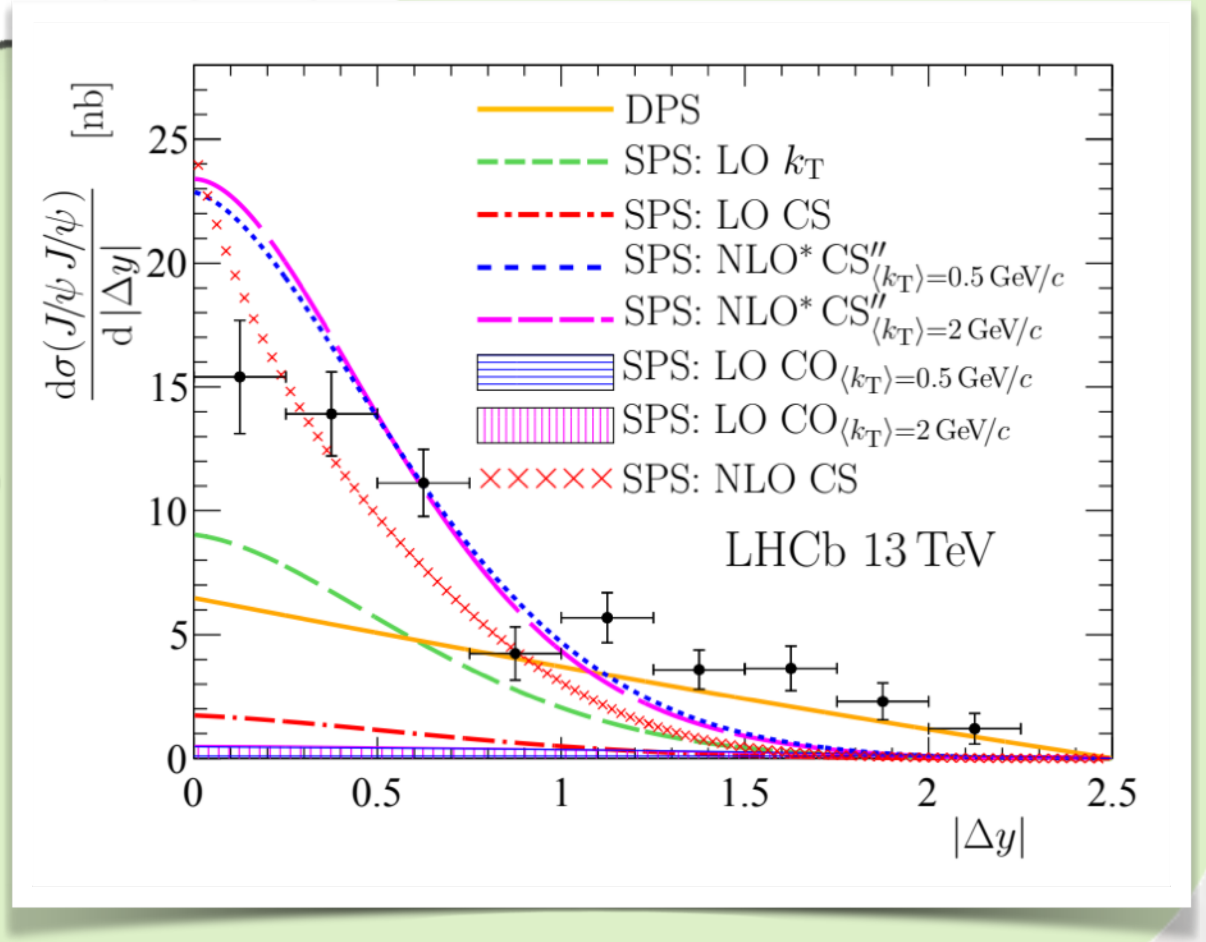
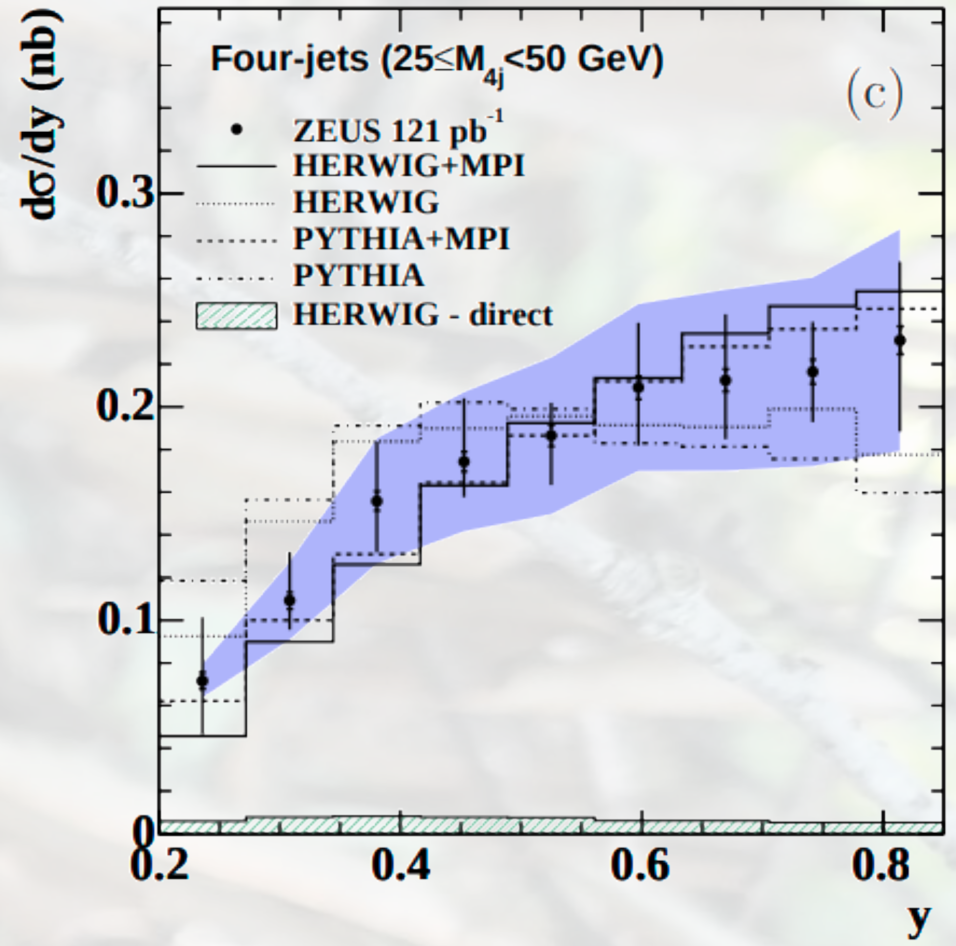


CDF, $\gamma + 3j$, Phys.Rev. D56 (1997) 3811-3832



LHCb, double J/ψ, JHEP 06, 047, (2017)

in ep Colliders?

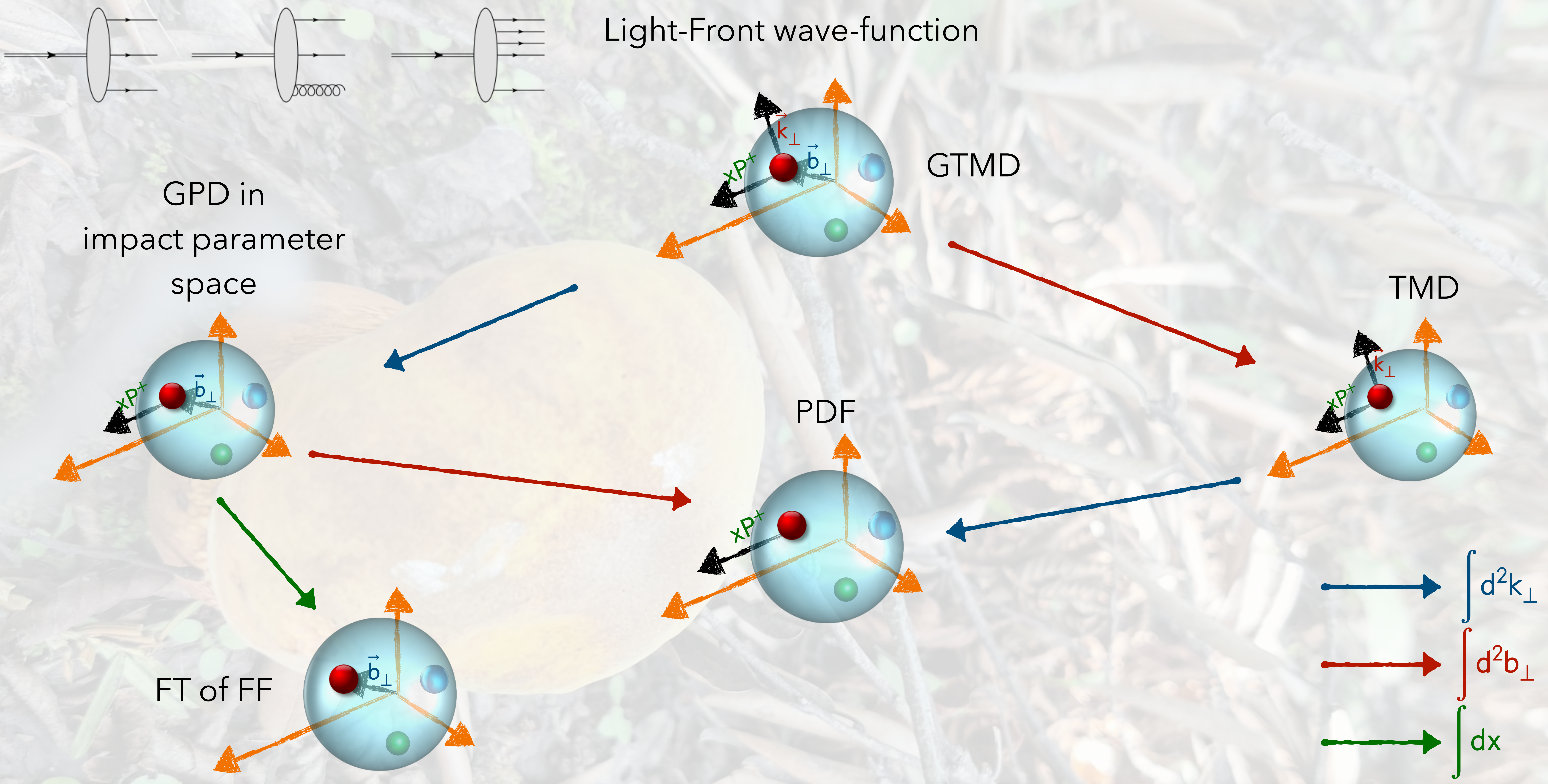


HERA data, ZEUS coll, Nucl. Phys. B 729, 1 (2008)

Access to:
- double parton correlations
- the transverse distance distribution of partons!!

all **UNKNOWN**

Multidimensional picture of hadrons



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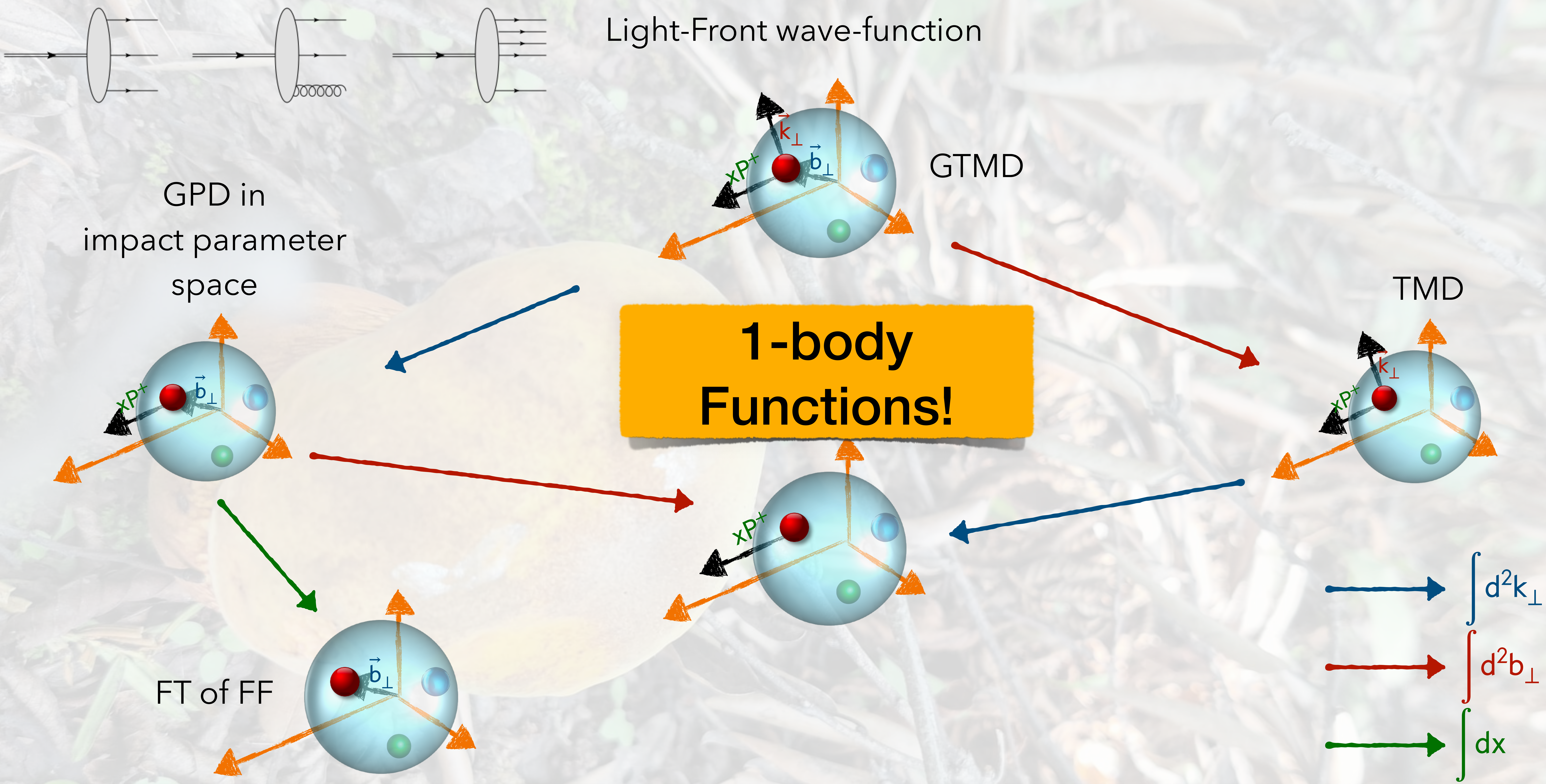
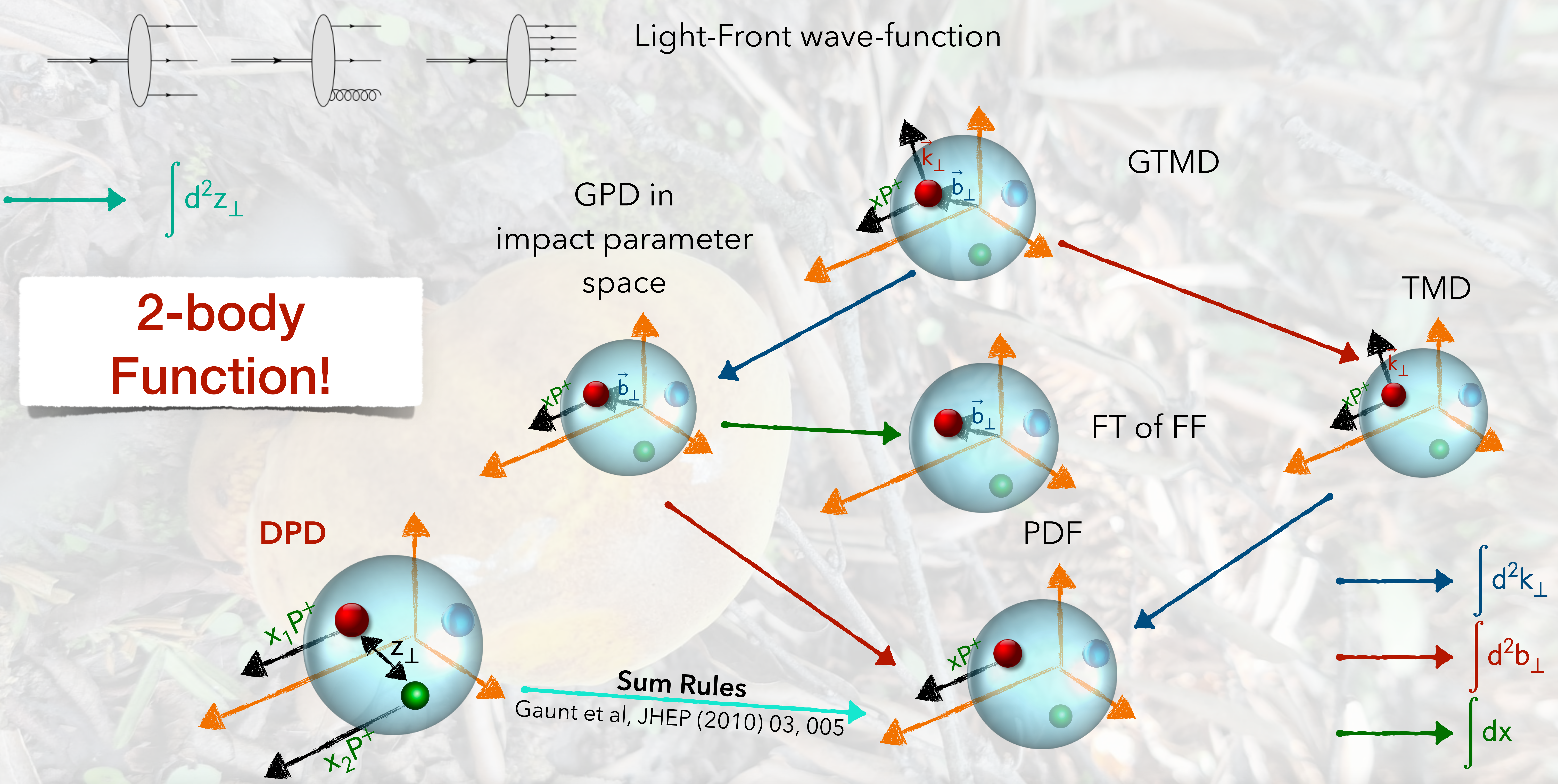


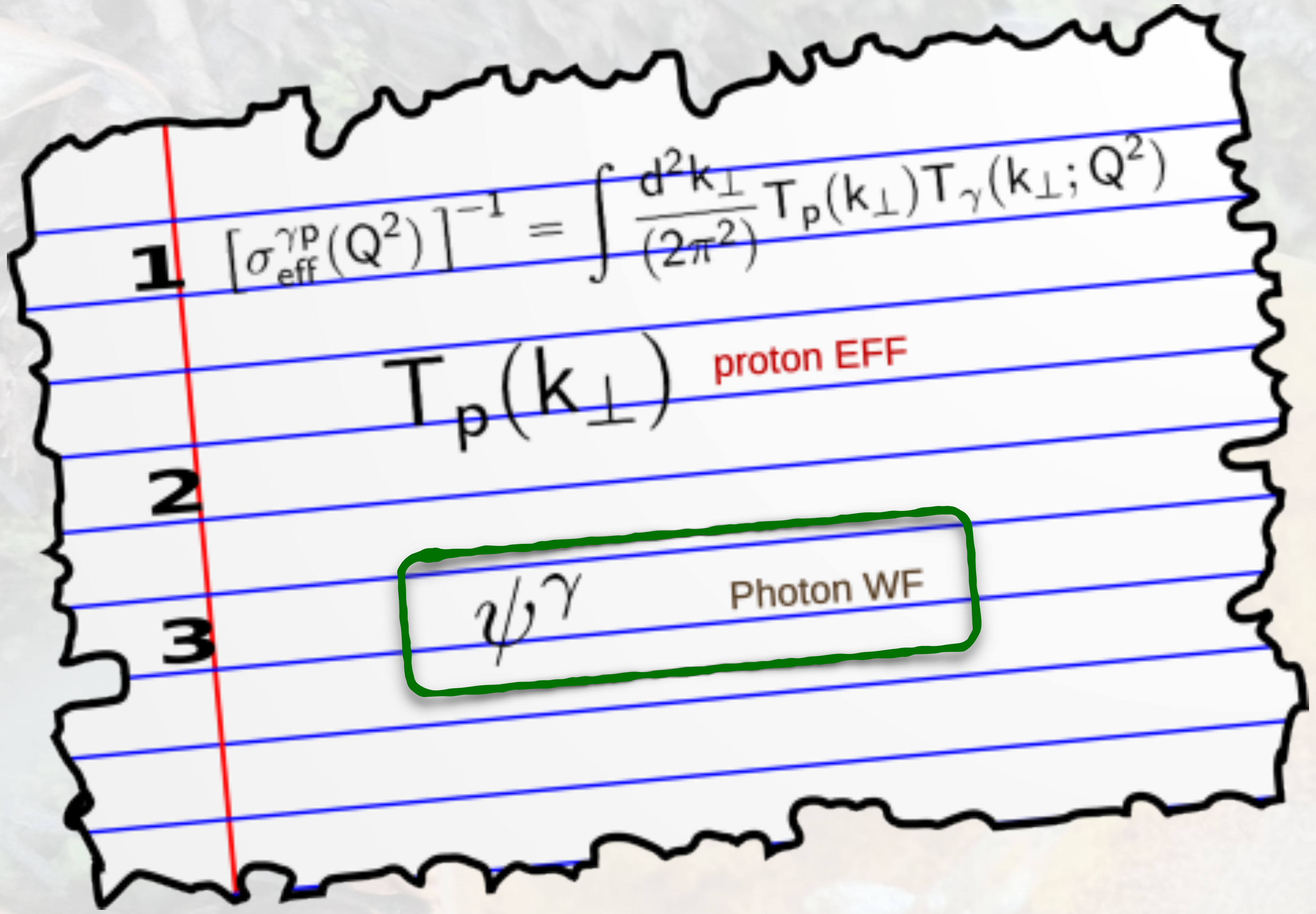
Multidimensional picture of hadrons



The $\gamma - p$ effective cross-section

M. R. and F. A. Ceccopieri, PRD 105 (2022) L011501

The main ingredients of the calculations:



For the photon W.F. use has been made of two choices representing two extreme cases:

1) QED at LO **(S.J. Brodsky et al. PRD50, 3134 (1994))**:

$$\psi_{q,\bar{q}}^{\lambda=\pm}(x, k_{1\perp}; Q^2) = -e_f \frac{\bar{u}_q(k)\ \gamma \cdot \varepsilon^\lambda\ v_{\bar{q}}(q-k)}{\sqrt{x(1-x)}\left[Q^2 + \frac{k_{1\perp}^2 + m^2}{x(1-x)}\right]}$$

2) Non-Pertubative (NP) effects **(E.R.Arriola et al, PRD74,054023 (2006))**

$$\psi_A^\gamma(x, k_{\perp 1}; Q^2) = \frac{6(1 + Q^2/m_\rho^2)}{m_\rho^2\left(1 + 4\frac{k_{\perp 1}^2 + Q^2 x(1-x)}{m_\rho^2}\right)^{5/2}}$$