

Inclusive quarkonium photoproduction at the LHC

Kate Lynch

Jean-Philippe Lansberg (IJCLab), Charlotte Van Hulse (UAH)
& Ronan McNulty (UCD)

Synergies between LHC and EIC for quarkonium physics
Trento

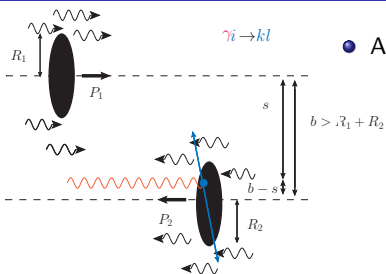


This project is supported by the European Union's Horizon 2020 research and innovation programme under Grant agreement no. 824093

Part I

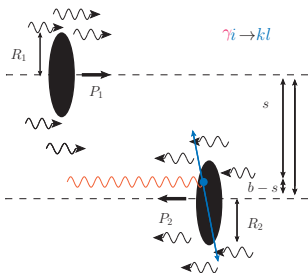
Introduction

Photon-induced interactions @ the LHC



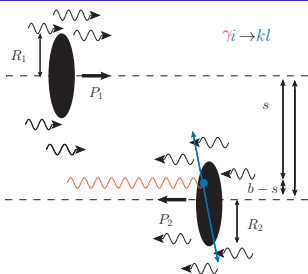
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Photon-induced interactions @ the LHC



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- Photoproduction usually studied in *ep* colliders
→ clean photoproduction environment
- However, the **LHC** is an excellent source of photons
→ can reach extremely large $W_{\gamma p}$

Photon-induced interactions @ the LHC



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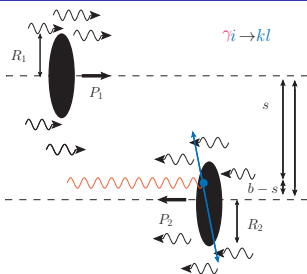
• Energies available at the LHC:

- $pp @ \sqrt{s} = 13 \text{ TeV} \rightarrow W_{\gamma p}^{\max} \approx 5 \text{ TeV} \rightarrow x_{\gamma}^{\max} \approx 0.14$
- $pPb @ \sqrt{s_{NN}} = 8.16 \text{ TeV} \rightarrow W_{\gamma p}^{\max} \approx 1.5 \text{ TeV} \rightarrow x_{\gamma}^{\max} \approx 0.03$

• Energies available at *ep* colliders:

- $W_{\gamma p}^{\max \text{ HERA}} \approx 240 \text{ GeV}$
- $W_{\gamma p}^{\max \text{ EIC}} \approx 100 \text{ GeV}$

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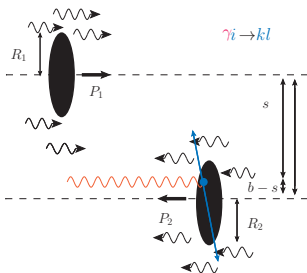
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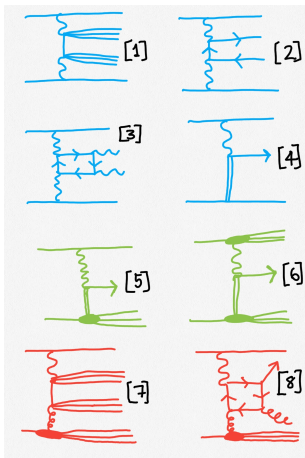
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We will show:

Inclusive quarkonium photoproduction can be measured via UPC at the **LHC**

Photon-induced interactions via UPC @ the LHC

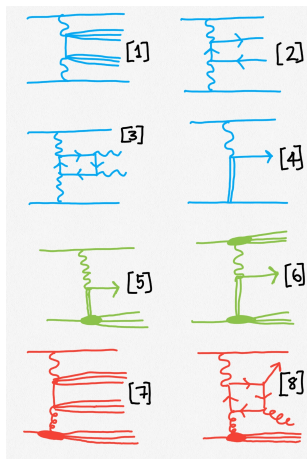
- So far focus of UPCs @ LHC on **exclusive processes** (fully determined final state) [1–4]
- Recently there were photoproduction studies with **nuclear break up** [5] (non-UPC [6*])
- Only published **inclusive** UPC study in PbPb: two-particle azimuthal correlations ATLAS, PRC 104, 014903 (2021)
- Coming soon: **inclusive** photonuclear dijets in PbPb [7]



- [1] **Exclusive dijet**: CMS, PRL 131 (2023) 5, 051901
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Photon-induced interactions via UPC @ the LHC

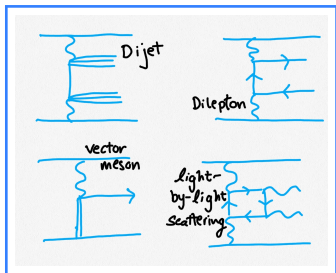
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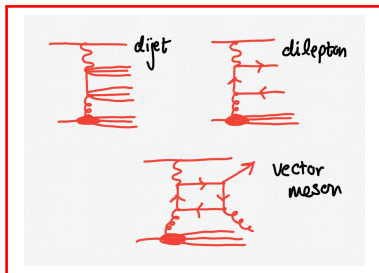
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- [8] **Inclusive quarkonium photoproduction**: **NOT YET MEASURED AT THE LHC!**

Exclusive vs. inclusive photoproduction at the LHC

Exclusive: fully determined final state

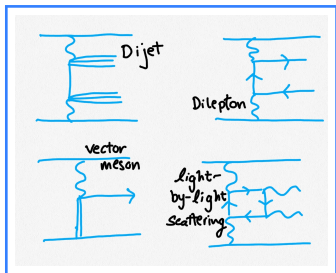


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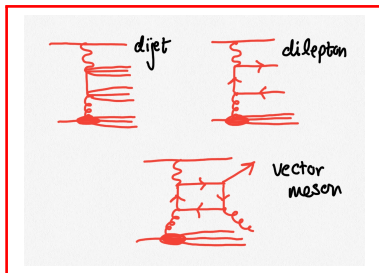
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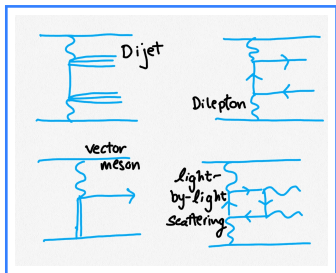
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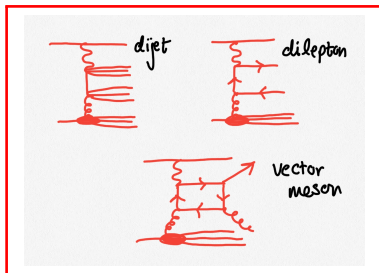
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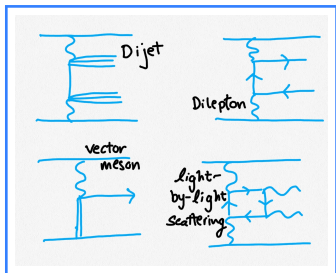
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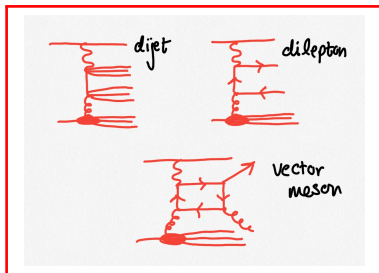
Exclusive vs. inclusive photoproduction at the LHC

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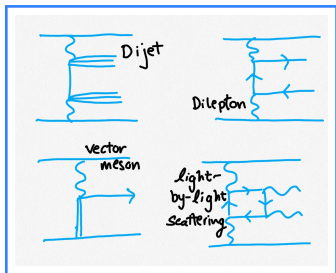
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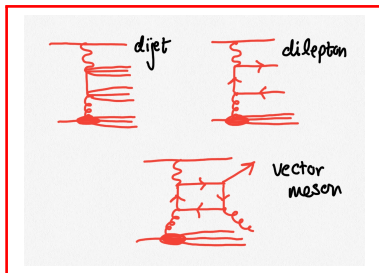
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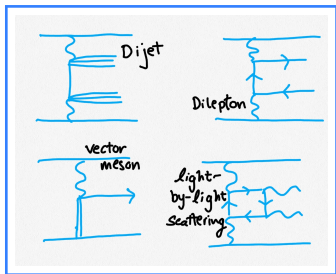
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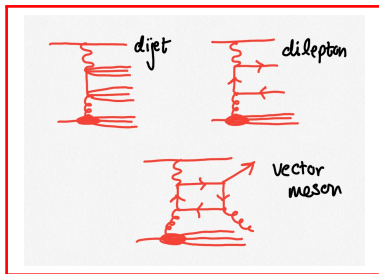
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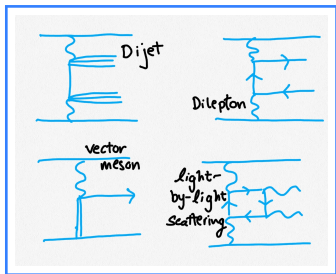
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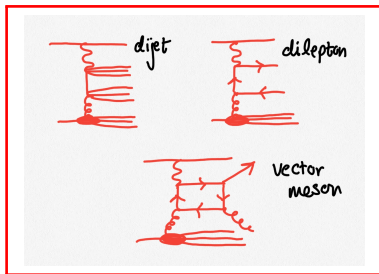
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- Measured at the LHC

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- Probe **Parton Distribution Functions**
- Colourful exchange
- Challenging: large backgrounds
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- Initial state kinematics **partially** determined by the final state
- Can and should be measured at the LHC

Quarkonium production status

- Discovered 50 years ago quarkonia are bound states of heavy quarks
- To date there is **no theoretical mechanism** that can **describe all of the data**
- Different models make different assumptions of the hadronisation
 - **Colour Evaporation model**: 1 free parameter per meson
 - × fails to describe di- J/ψ data
 - **Colour Singlet model**: no free parameters
 - × tends to undershoot large p_T data
 - **Colour Octet mechanism** (extension to CSM via non-relativistic QCD): free parameters
 - × cannot simultaneously describe the photoproduction and polarisation data

Maxim Nefedov, QaT 2023

LDME fit	J/ψ hadropr.	J/ψ photopr.	J/ψ polar.	η_c hadropr.
Butenschön et al.	✓ ($p_T > 3$ GeV)	✓	✗	✗
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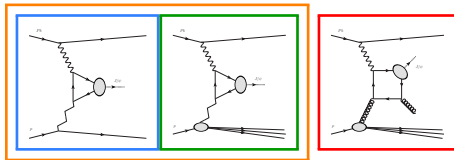
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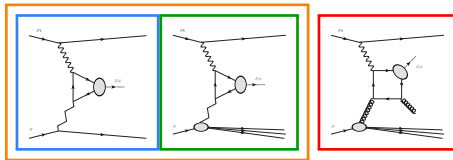
More inclusive photoproduction data → possible at ~~EIC~~ in 10 years LHC today!

Existing J/ψ photoproduction measurements from HERA



- Data exists for **diffractive** (exclusive and proton-dissociative) & **inclusive/inelastic** photoproduction @ HERA $\sqrt{s} = 320$ GeV

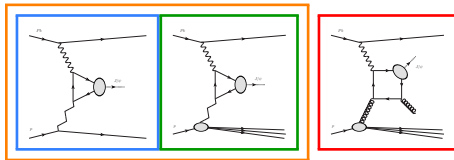
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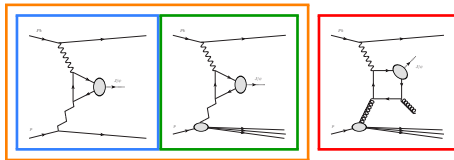
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- HERA result: $\sigma_{\text{exclusive}}^{\text{HERA}} \simeq \sigma_{\text{dissociative}}^{\text{HERA}} \simeq \sigma_{\text{inclusive}}^{\text{HERA}}$
- Expectation: $\sigma_{\text{exclusive}}^{\text{LHC}} \simeq \sigma_{\text{dissociative}}^{\text{LHC}} \simeq \sigma_{\text{inclusive}}^{\text{LHC}} \rightarrow$ only difference is photon flux!
- **Exclusive** and **proton-dissociative** photoproduction have been measured @ LHC
- Expect that **inclusive yield** is sufficiently large we will demonstrate this

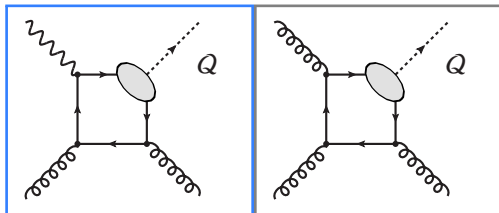
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- Expect that **inclusive yield** is sufficiently large we will demonstrate this
- Measuring **inclusive** quarkonium photoproduction to **understand the quarkonium hadronisation**

Is it feasible to measure inclusive quarkonium photoproduction at the LHC?

- Anticipate sizeable **photoproduction** yield
- Large hadronic background must be shown to be suppressed



Proton-lead is the ideal collision system

- Enhanced photon flux w.r.t. pp : $\propto Z^2$
- No ambiguity as to the photon emitter: reconstruction of z and $W_{\gamma p}$
- Less pileup than pp

Part II

Methodology

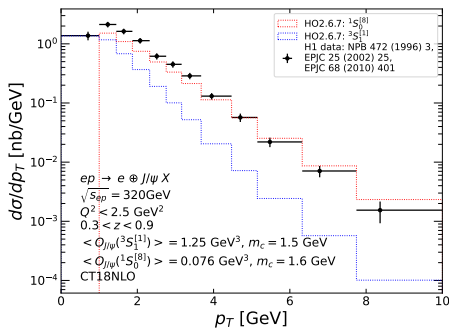
Building a Monte Carlo sample

We must:

- 1 Evaluate yield & P_T reach: need reliable Monte Carlo (MC) sample

Problem:

- Only **LO MC** for quarkonia + QCD corrections are large!
 - **LO CS** **undershoots** undershoots large P_T data
 - **LO CO** **captures** large P_T data



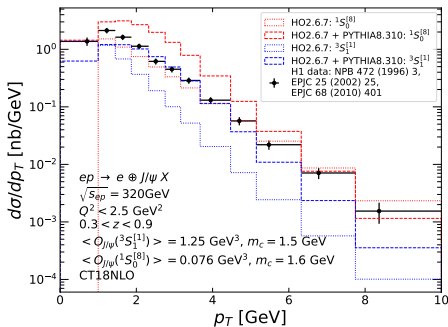
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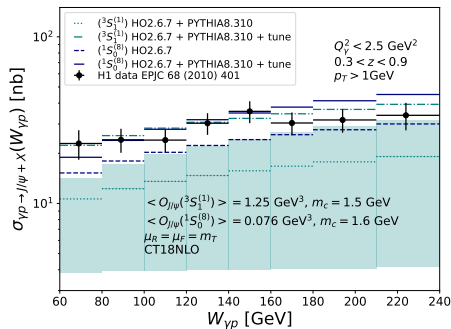
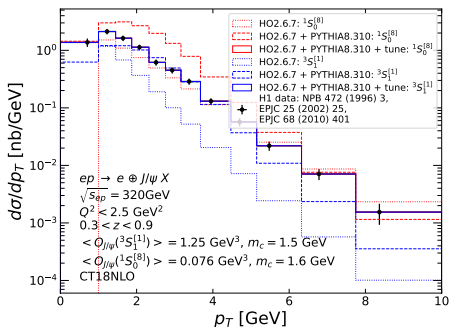
- Only **LO MC** for quarkonia + QCD corrections are large!
 - **LO CS + PS undershoots** improved but still undershoots
 - **LO CO + PS captures** overshoots low P_T data



Building a Monte Carlo sample

We must:

- Evaluate yield & P_T reach: need reliable Monte Carlo (MC) sample
Solution: perform tune in P_T to HERA data + keep \sqrt{s} and y dependence from photon flux



- Reject background: reliable background MC + background reduction strategy

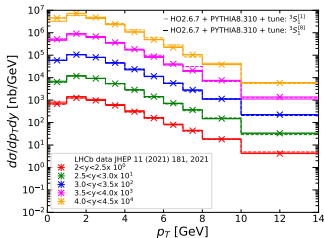
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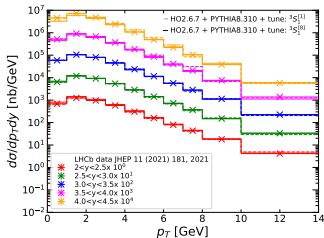
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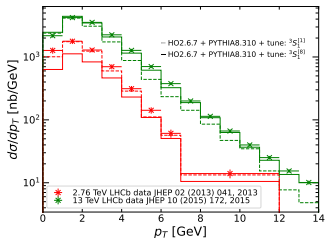
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 - 2 Tuning is \sqrt{s} independent

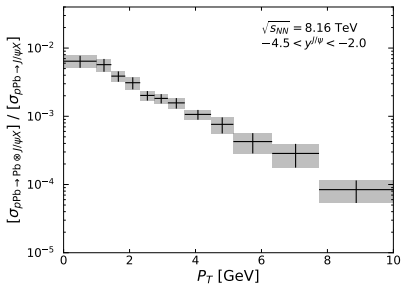
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Validation 2: tune vs. 13- and 2.76 TeV data.

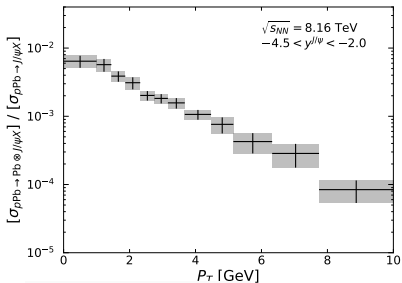


Background-reduction techniques

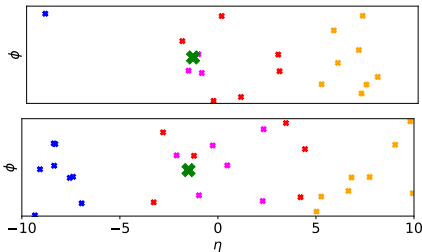
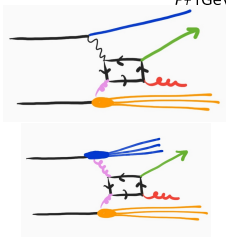


- Large yields but huge **background**!
- **Background** reduction critical at large P_T
- **Hadroproduced** J/ψ are associated with more detector activity than **photoproduced** J/ψ

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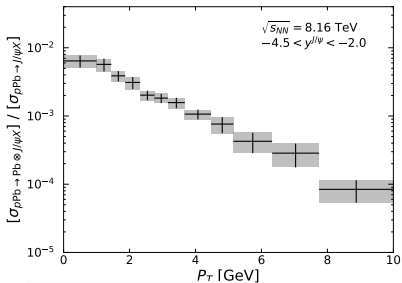


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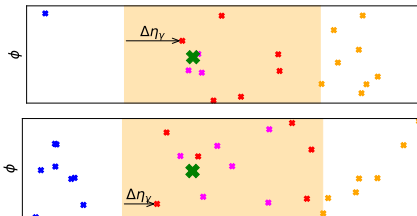
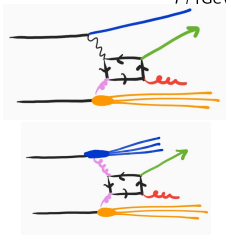


- 3 background-reduction techniques based on different detector acceptances

Background-reduction techniques

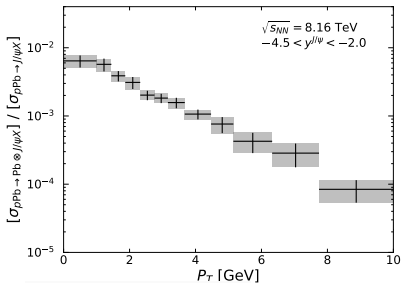


- Large yields but huge **backgrounds!**
- **Background** reduction critical at large P_T
- **Hadroproduced** J/ψ are associated with more detector activity than **photoproduced** J/ψ

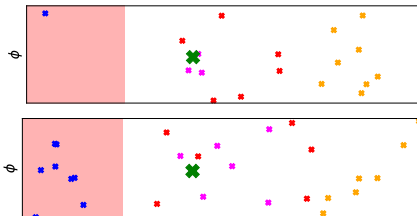
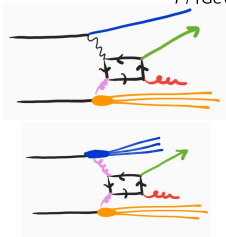


- 3 background-reduction techniques based on different detector acceptances:
| central $\Delta\eta_\gamma$: distance in rapidity between main detector on photon-going side and closet particle activity

Background-reduction techniques

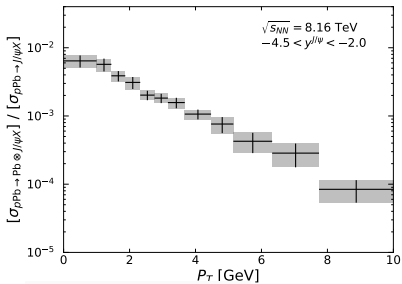


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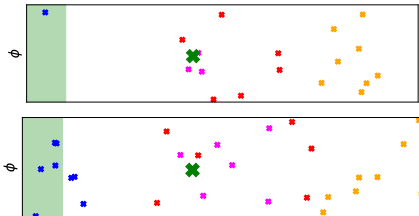
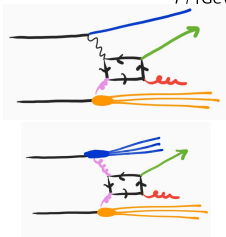


- 3 background-reduction techniques based on different detector acceptances:
I **central** II **forward**

Background-reduction techniques



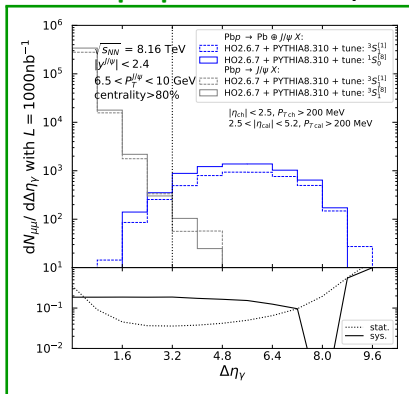
- Sufficient yields but huge backgrounds!
- Background reduction becomes more critical at larger P_T
- **Hadroproduced** J/ψ are associated with more detector activity than **photoproduced** J/ψ



- 3 background-reduction techniques based on different detector acceptances:
I central II forward III far-forward

Method I: Rapidity gaps in LHC detectors

General purpose detector [ATLAS, CMS]



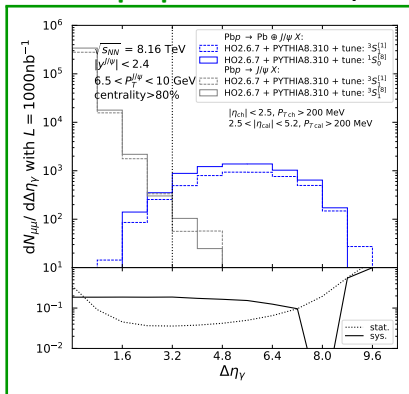
Broad rapidity coverage:

CMS/ATLAS 10 units

clean separation between photoproduction
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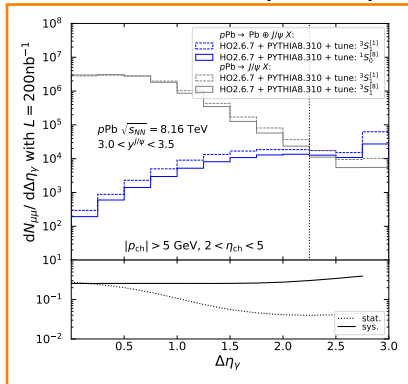


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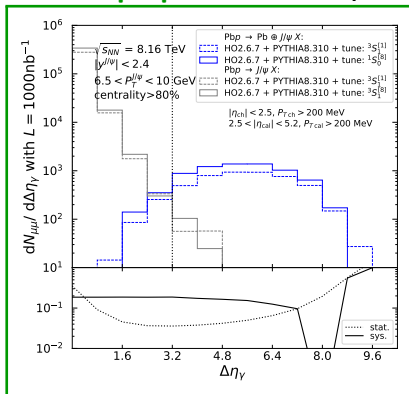
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LHCb 3 units, ALICE 1.8 units

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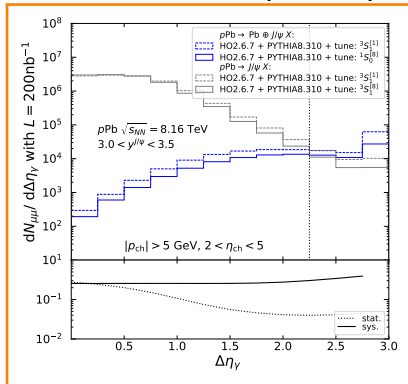


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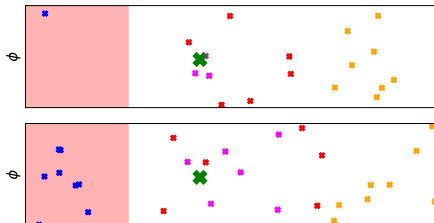
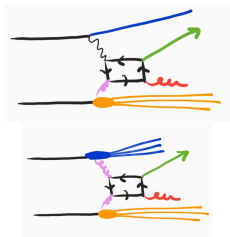
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- Selecting a cut value that minimises that statistical uncertainty:
 - removes $\mathcal{O}(99.99\%)$ ($\mathcal{O}(99.9\%)$) of background events → $S/B \gtrsim \mathcal{O}(1)$

Background-reduction techniques

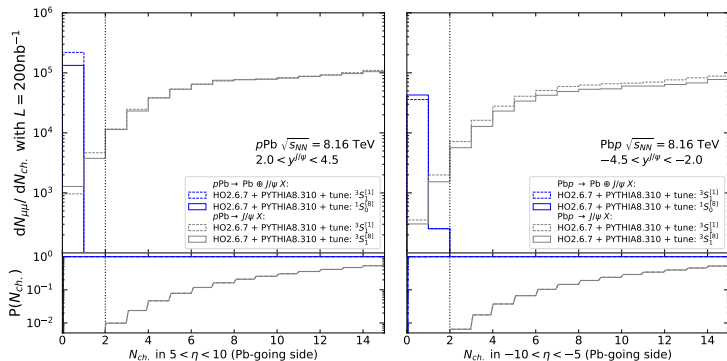


Method II: forward activity with **HeRSChel** at LHCb

- forward scintillator sensitive to **charged particle activity** in the region $5 < |\eta| < 10$
- Photoproduction events identified with **no HeRSChel activity**

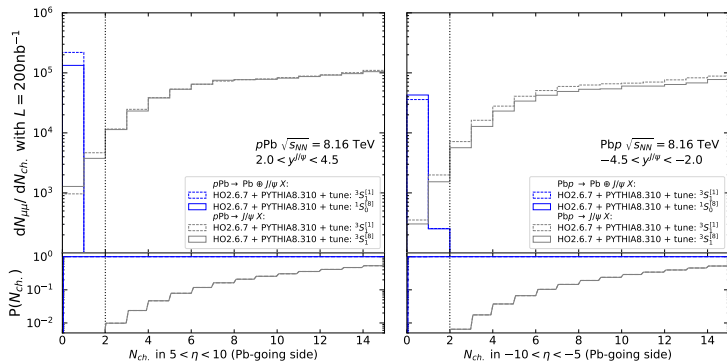
Selecting events based on activity in HeRSChel

- Differential yield w.r.t. the number of charged particles on the γ -emitter side within $5 < \eta < 10$ for photo- and hadroproduced J/ψ



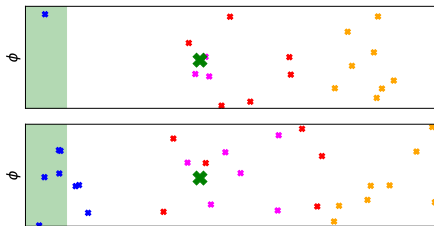
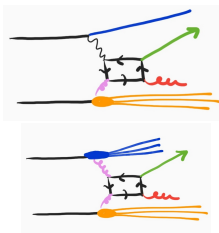
Selecting events based on activity in HeRSChel

- Differential yield w.r.t. the number of charged particles on the γ -emitter side within $5 < \eta < 10$ for **photo-** and hadroproduced J/ψ



- We anticipate a clear distinction between **photo-** and hadroproduction
- Necessary to perform a full detector simulation to include HeRSChel response

Background-reduction techniques



Method III: far-forward activity with **zero-degree calorimeter** at ALICE, ATLAS, & CMS

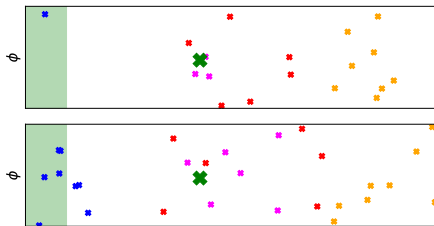
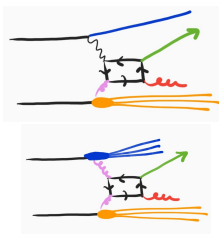
- Detector close to the beam pipe ($|\eta| \gtrsim 8$) sensitive to **neutral particles**
- UPCs identified as most peripheral events (80 – 100% centrality)

[Already done in p Pb collisions: ALICE, JHEP 02 (2021) 002]

- Selecting events with **0 neutrons** in ZDC can further enhance signal purity

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- This would **not** be possible in PbPb where there is a non-negligible photoproduction cross section with neutron emissions $\mathcal{O}(50\%)$

- disentangling the photon emitter CMS, Phys.Rev.Lett. 131 (2023) 26, 262301, PRC 93, 055206 (2016)

Part III

Results

One of the advantages of $p\text{Pb}$ over pp is the significantly reduced pile-up. However, given the possibility of a $p\text{Pb}$ run with a sizeable μ value we should consider the efficacy of methods I–III under such conditions:

- Method I: **rapidity gaps**
 - **Calorimeter** based rapidity-gap definitions not possible
 - **Only** rapidity-gap definitions based on **charged tracks** possible
 - Reduced $\Delta\eta$ reach for ATLAS and CMS $10 \rightarrow 5$ units
- Method II: **HeRSChEL**
 - Timing is insufficient
- Method III: **ZDC**
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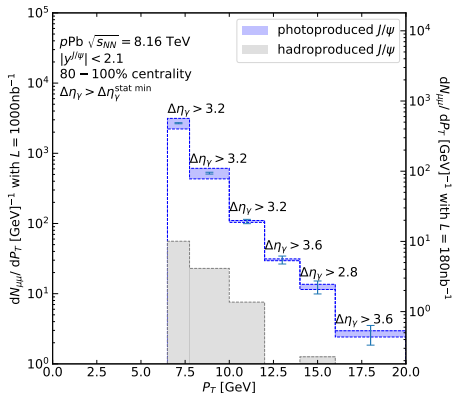
Comment on methods I–III in HL run

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The same comments apply to exclusive UPCs

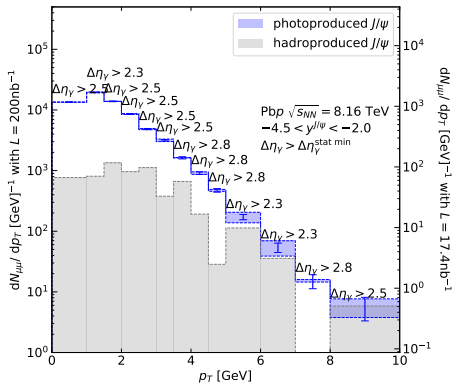
Photoproduction yields: ATLAS & CMS



- Possible to **isolate photoproduction** with CMS and ATLAS using methods I & III
- Possible to further enhance signal purity by selecting 0n events
- With Run3+4 lumi, possible to extend the P_T reach from 10 GeV (HERA data) to **20 GeV**

detector	CMS	LHCb	CMS	LHCb
	<u>Run 2 lumi:</u>		<u>Run 3+4 lumi:</u>	
yield	$\mathcal{O}(10^3 - 10^5)$	$\mathcal{O}(10^3 - 10^4)$	$\mathcal{O}(10^4 - 10^6)$	$\mathcal{O}(10^4 - 10^5)$
P_T reach	14 GeV	8 GeV	20 GeV	14 GeV

Photoproduction yield: LHCb



- Possible to **isolate photoproduction** at LHCb using method I alone
- Combining with HeRSChEL information (method II) will improve background removal
- Expect ψ' yield to be $\sim 1/20$ of J/ψ yield no P_T differential data from HERA!

detector	CMS	LHCb	CMS	LHCb
	<u>Run 2 lumi:</u>		<u>Run 3+4 lumi:</u>	
yield	$\mathcal{O}(10^3 - 10^5)$	$\mathcal{O}(10^3 - 10^4)$	$\mathcal{O}(10^4 - 10^6)$	$\mathcal{O}(10^4 - 10^5)$
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Kinematic reconstruction: $W_{\gamma p}$ and z

We have shown that it is possible to measure P_T -differential inclusive photoproduction cross sections at the LHC without waiting for the EIC

- What about $d\sigma/dz$ and as a function of $W_{\gamma p}$?
 - Fully equivalent to ep measurements

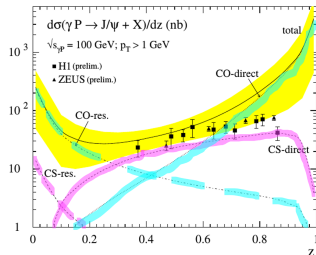
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octet vs. singlet

Kramer, hep-ph/016120



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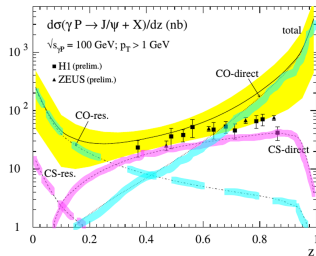
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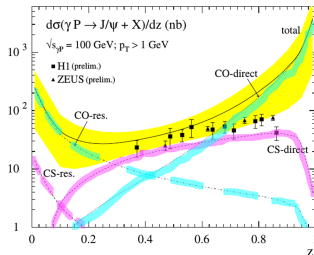
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Kramer, hep-ph/016120



- Let us reconstruct the photon kinematics from the final state :

$$\text{Pb}(P_{\text{Pb}}) + p(P_p) \xrightarrow{\gamma(P_\gamma)} \text{Pb}(P'_{\text{Pb}}) + J/\psi(P_\psi) + X(P_X) \text{ thus } P_\gamma = P_\psi + P_X - P_p$$

- $W_{\gamma p} \simeq \underbrace{(2(P_\psi + P_X - P_p) \cdot P_p)^{1/2}}_{P_\gamma} \quad \& \quad z = \frac{P_p \cdot P_\psi}{P_p \cdot (P_\psi + P_X - P_p)}$

- We only need to measure $(P_\psi \cdot P_p)$ & $(P_X \cdot P_p)$ or equivalently $P_X^- = E_X - P_{X,z}$

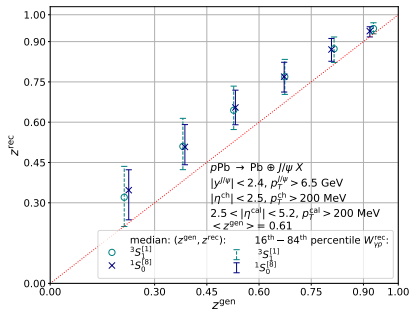
- NB: In the exclusive case, $P_X \simeq P'_p \Rightarrow P_\gamma + P'_p = P_\psi + P'_p$ and $W_{\gamma p} \simeq M_\psi e^{-y_\psi}$

Kinematic reconstruction: results

- Limited detector coverage $\Rightarrow P_{\text{reconstructed}}^- < P_{\text{generated}}^-$

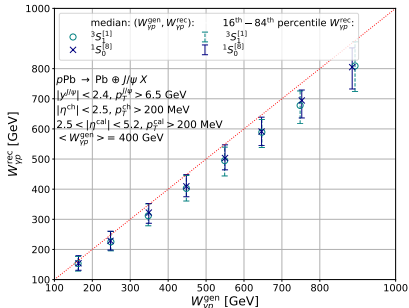
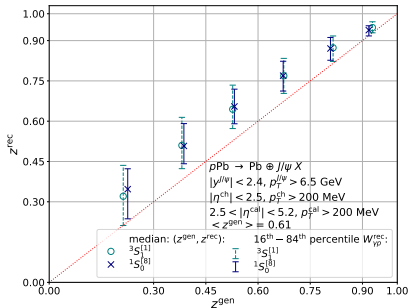
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- This results in the following biases;
 - $Z_{\text{rec}} > Z_{\text{gen}}$



Kinematic reconstruction: results

- Limited detector coverage $\Rightarrow P^-_{\text{reconstructed}} < P^-_{\text{generated}}$
- This results in the following biases;
 - $z_{\text{rec}} > z_{\text{gen}}$ & $W_{\gamma p}^{\text{rec}} < W_{\gamma p}^{\text{gen}}$



- For CMS and ATLAS: z reconstruction allows for $\mathcal{O}(5 - 6)$ bins (similar to HERA) improves with increasing values of z
- $W_{\gamma p}$ reconstruction allows for $\mathcal{O}(10)$ bins improves for decreasing values of $W_{\gamma p}$

Summary and outlook

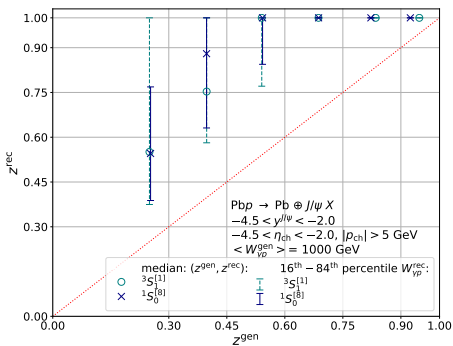
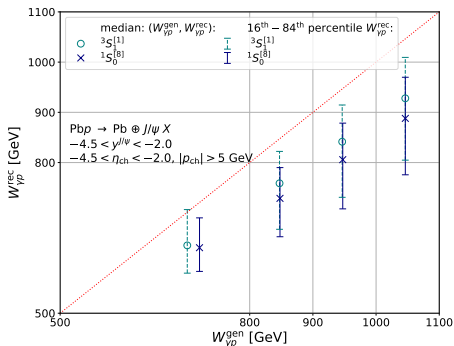
- A **proton-lead** collision system allows the LHC to be used as a **photon-nucleon** collider
 - Feasible to measure inclusive J/ψ , ψ' and Υ photoproduction at the LHC
 - Complementary to HERA measurements with a doubled P_T reach
 - It can be done now $\mathcal{O}(10)$ years before the EIC
- CMS and ATLAS are the **most favourable** experiments with the largest P_T reach and broadest pseudorapidity coverage

(CMS has additional advantage of measuring $P_T \rightarrow 0$ GeV)
- Possible to make measurements at ALICE and LHCb too!
- Despite the impossibility to measure the intact Pb ion,
possible to reconstruct z and $W_{\gamma p}$
 - Binning competitive with HERA, confirms the reach in $W_{\gamma p}$ up to 1 TeV!
 - Possibility to isolate resolved-photon contributions

Backup

Reconstruction of kinematic variable in LHCb

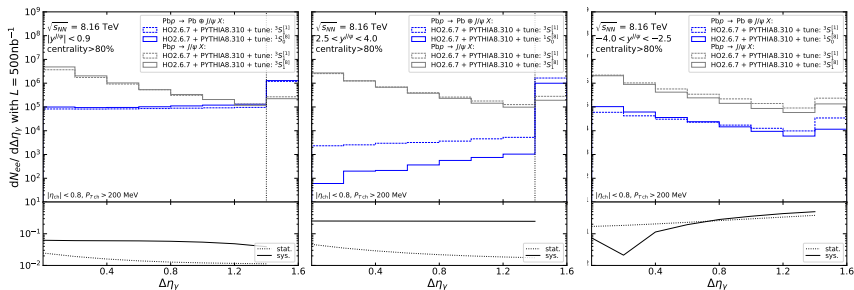
Owing to the narrow pseudorapidity gap coverage in LHCb ($2 < \eta < 5$), reconstruction of kinematics at LHCb is not possible.



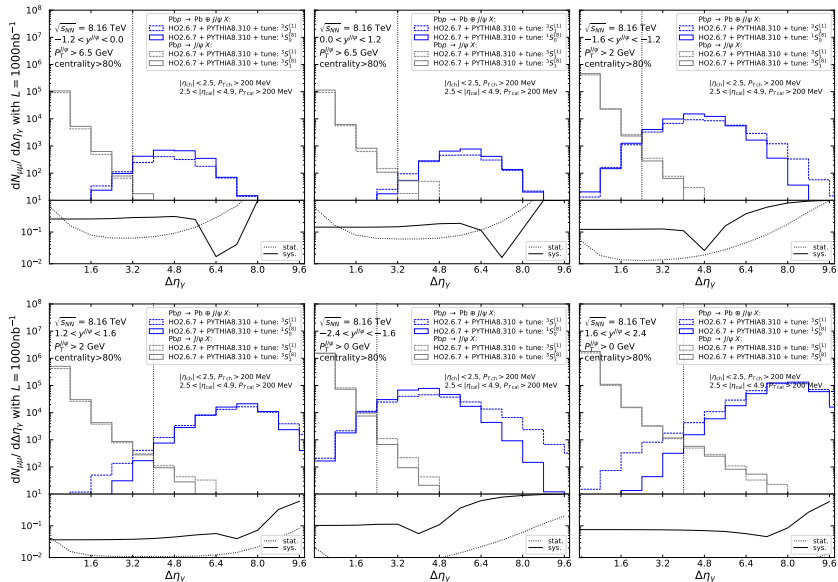
- When $z^{\text{rec}} = 1$ only the J/ψ is captured!

Rapidity gap distributions in ALICE

$\Delta\eta_\gamma$ -differential yield for J/ψ in the ALICE acceptance.



$\Delta\eta_{\gamma}$ - y -diff. yield for J/ψ in CMS low- P_T acceptance

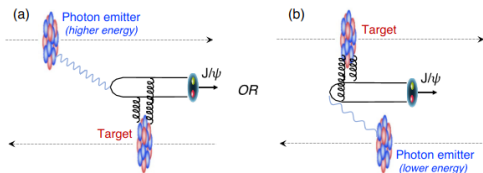


Neutron emission: disentangling the photon emitter

- For exclusive vector meson production in PbPb collisions there is an ambiguity as to which Pb ion is the photon emitter
- At a given rapidity either:

(a) $x_\gamma = \frac{m_{T J/\psi}}{\sqrt{s}} e^{+y_{J/\psi}}$, $x_{\text{Pb}} = \frac{m_{T J/\psi}}{\sqrt{s}} e^{-y_{J/\psi}}$ or (b) $x_\gamma = \frac{m_{T J/\psi}}{\sqrt{s}} e^{-y_{J/\psi}}$, $x_{\text{Pb}} = \frac{m_{T J/\psi}}{\sqrt{s}} e^{+y_{J/\psi}}$

ALICE, JHEP 10 (2023) 119; CMS, Phys.Rev.Lett. 131 (2023) 26, 262301 PRC 03 055206 (2016)



- Neutron emissions (detected with ZDCs) serve as an impact parameter filter
- Larger photon energies are associated with smaller impact parameters
- $0nXn$ and $XnXn$ select smaller impact parameter and larger x_γ compared to $0n0n$

