

INCL for Neutrino Scattering

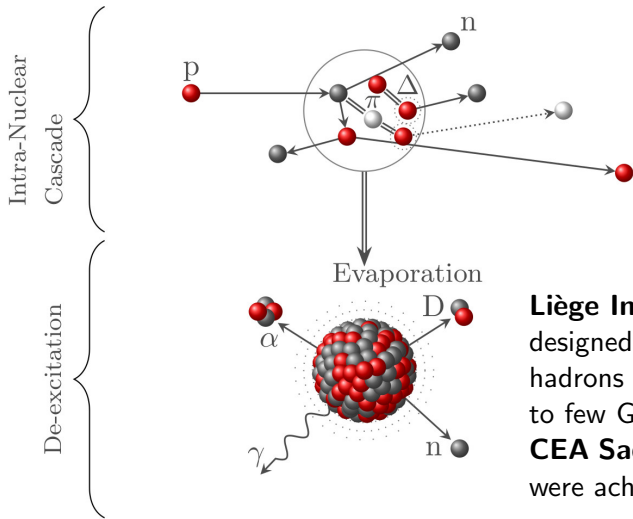
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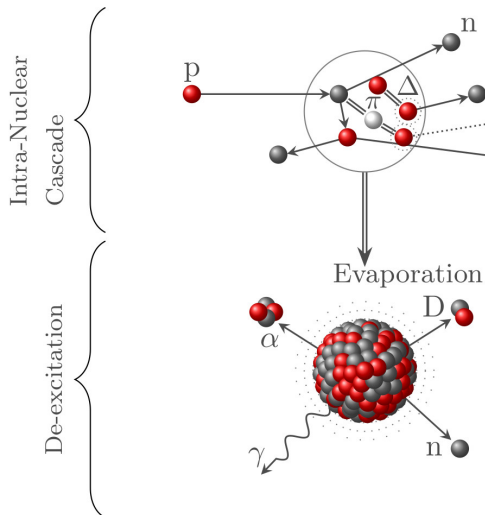



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Liège Intra Nuclear Cascade: is a code initially designed to simulate nuclear reactions induced by hadrons and light nuclei in the few tens of MeV to few GeV range. Its development is ongoing at **CEA Saclay**, and the results presented in this talk were achieved there.



Projectiles: baryons (nucleons, Λ , Σ), mesons (pions and Kaons) or light nuclei ($A \leq 18$). We use neutrino vertex from  **NuWro**

Versatile tool: has been implemented in GEANT4, GENIE, and other

De-excitation: ABLA, SMM, GEMINI

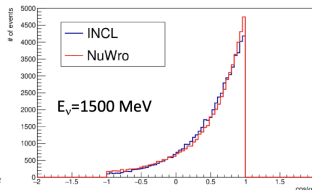
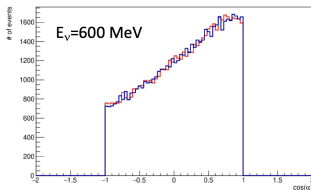
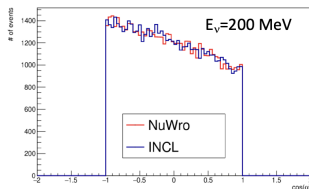
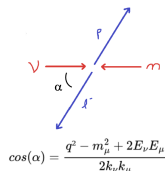
We will use **ABLA**, since it is proved to work for the **light nuclei** (Phys. J. Plus 130, 153 (2015))

Neutrino simulation results:
PRD 106, 3 (2022), PRD 108, 11 (2023)

Llewelyn-Smith formula is **implemented in INCL**:

$$\frac{d\sigma^{\nu/\bar{\nu}}}{d|q^2|} = \frac{M^2 G^2 (f^{CC/NC})^2}{8\pi E_\nu^2} \left[A(q^2) \mp B(q^2) \frac{(s-u)}{M^2} + \frac{C(q^2)(s-u)^2}{M^4} \right]$$

Validation on NuWro calculations for scattering angle in CM



- **Full systematic comparisons** with NuWro are needed before making this version available
- **A paper** is expected to be published next year
- Work of **Antoine Legendre**

Potential

Each nucleon in the nucleus has its **position and momentum** and moves **freely** in a square potential well. Nuclear model is essentially **classical**, with some additional ingredients to mimic quantum effects.

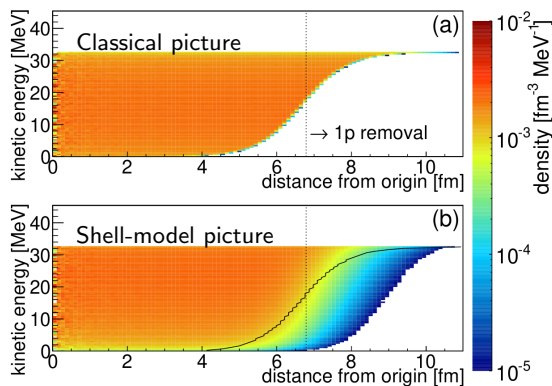
Pauli Blocking

- **strict**: blocked is $p < p_{Fermi}$
- **statistical**: count only nearby nucleons
- strict for the first event and statistical for the subsequent ones

Events inside cascade

- decay/collision
- reflection/transmission with probability to **leave the nucleus as a nuclear cluster**

Space-kinetic-energy density of protons in ^{208}Pb



Phys.Rev.C 91, 034602 (2021)

More details about INCL in **Jean-Christophe's** talk today!

Potential

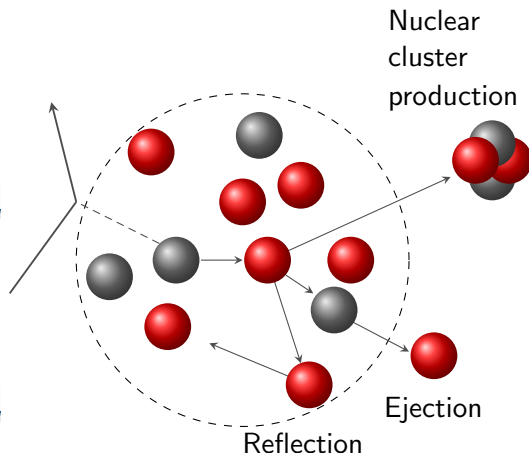
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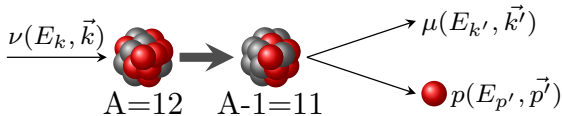
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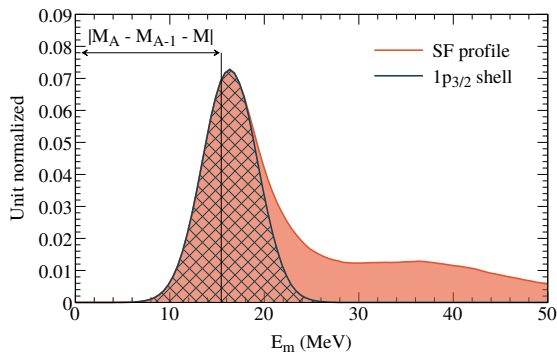


Experimental definition:

$$E_x^{\text{exp}} = E_{\text{missing}} - \underbrace{(M_A - M_{A-1} - M)}_{\text{constant}}$$

- A constant shift of missing energy by ~ 15.4 MeV leads to **non-physical, negative values**
- We use experimental data (J. Phys. G: Nucl. Part. Phys. 16 507 (1999)) to simulate discrete levels

M_{A-1} is the rest mass of the $A - 1$ nucleus
 M_A is the rest mass of the initial A nucleus
 M is the rest mass of the target nucleon
 E_{missing} is the missing energy
 For interaction on carbon,
 $M_A - M_{A-1} - M = 15.4$ MeV



We assume all strength below the peak comes from the symmetric $1p_{3/2}$ shell

Experimental data used for the carbon spectral function

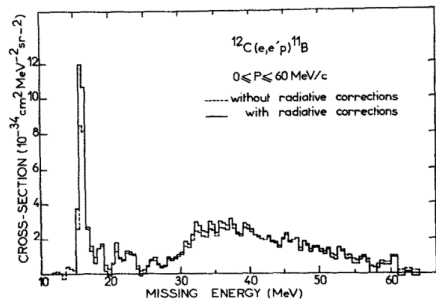


Fig. 7. Energy spectrum of the $^{12}\text{C}(e, e'p)^{11}\text{B}$ reaction before and after the radiative corrections.

Experimental data used to simulate the discrete levels

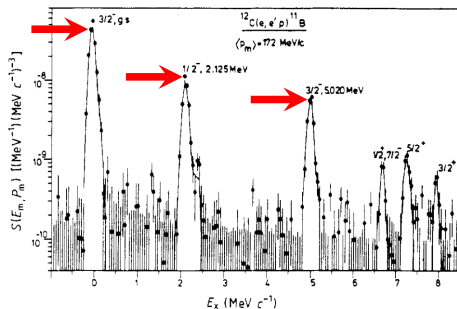
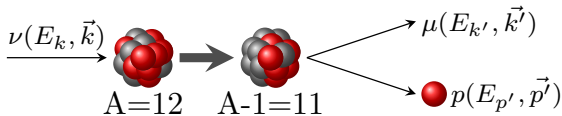


Figure 22. Excitation-energy spectrum of ^{11}B observed in the reaction $^{12}\text{C}(e, e'p)$. Both negative and positive-parity final states are shown.

Nuclear Physics A262 (1976) 461-492

J.Phys.G: Nucl.Part.Phys. 16 507 (1990)



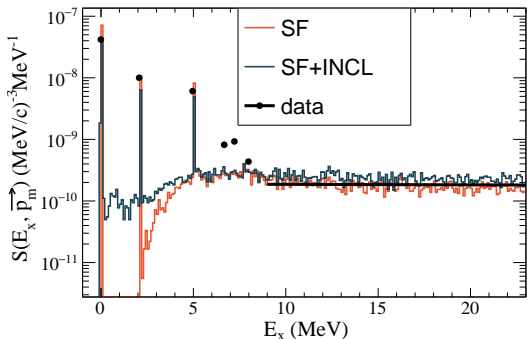
M_R^* is the mass of the excited remnant
 M_R is the rest mass of the remnant
 T_R is the kinetic energy of the excited remnant
 p_m is the missing momentum

For the continuous spectrum part, we can calculate excitation energy as:

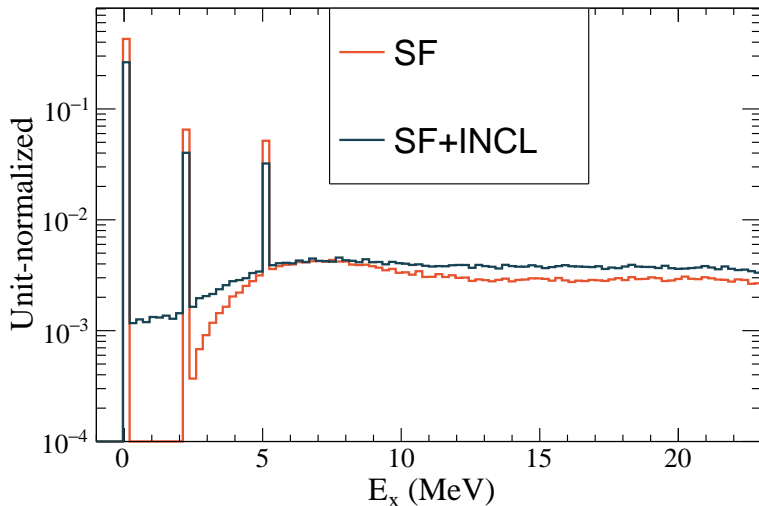
$$E_x = M_R^* - M_R, \text{ where:}$$

$$M_R^* = \sqrt{(E_k + M_A - E_{k'} - E_{p'})^2 - |\vec{p}_m|^2}$$

Otherwise, we model **3 discrete peaks** with strength of 79%, 12%, and 9% (**p-shell**)

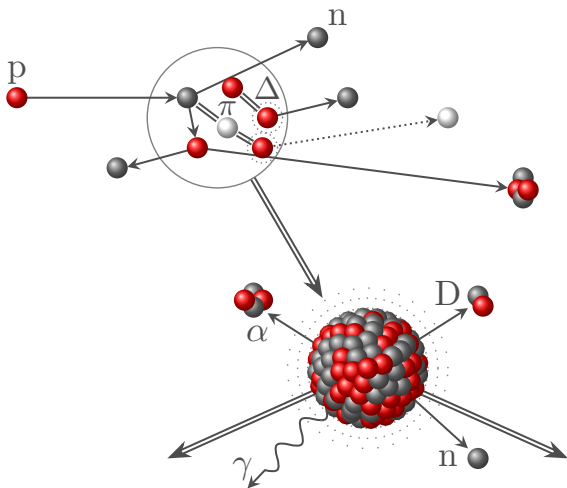


Excitation energy of the nuclear remnant after neutrino interactions for $169.5 < p_m < 174.5$ MeV



Final distribution of the excitation energy

We use INCL+ABLA to simulate **neutrino CCQE** events with the **T2K** energy spectrum. We compare the obtained results to **NuWro**.

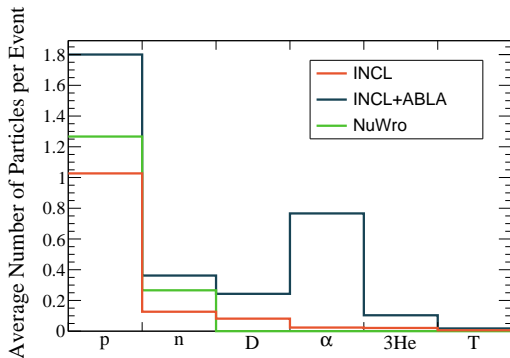


In the paper [Phys.Rev.D 106, 3 \(2022\)](#) we show the **nuclear cluster production for the first time** in FSI.

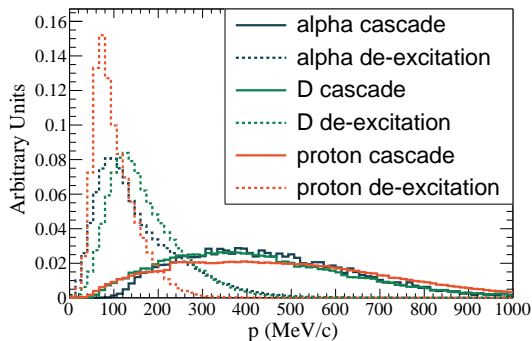
In the paper [Phys.Rev.D 108, 11 \(2023\)](#) we study the impact of the subsequent **de-excitation modelling**, that predicts **more nuclear clusters**.

Important! Gamma production during de-excitation is just a **small fraction** of possible scenarios

ABLA features a **massive production** of particles with **low momentum**.

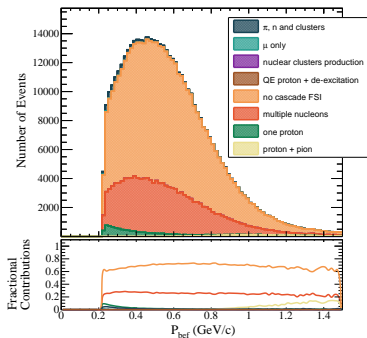


Average number of particles per event produced by INCL, INCL+ABLA, and NuWro.

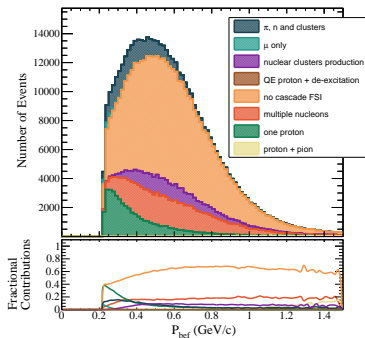


Momentum of nuclear clusters produced during the cascade and de-excitation

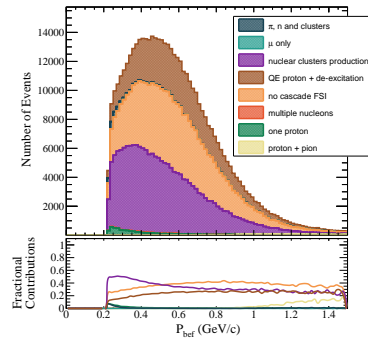
- Large fraction of "no FSI" events (i.e. proton untouched) is now **feature production of other particles** (and nuclear clusters) in the final state due to de-excitation
- Events with **only nucleon production** now feature **nuclear cluster production**



NuWro

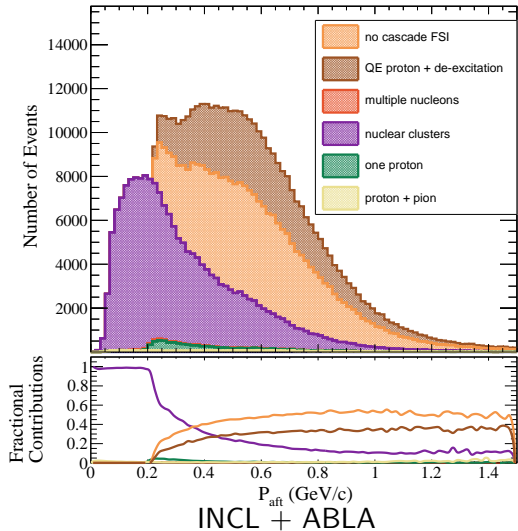
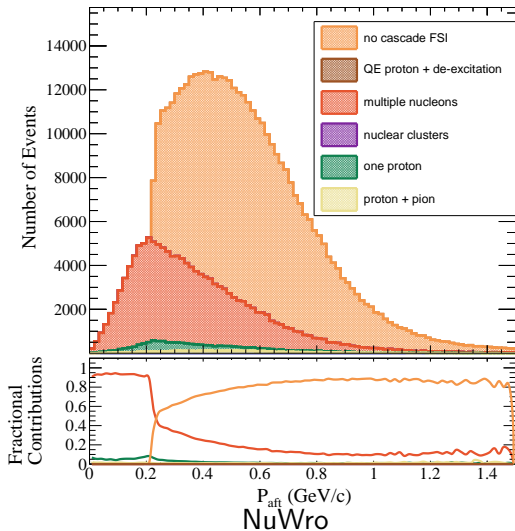


INCL



INCL + ABLA

INCL+ABLA simulation features **massive difference** in nucleon kinematics in comparison to NuWro



We use **Single Transverse Variables (STV)** that allow to disentangle different effects for better FSI estimation. STV are **observable** and **measurable**.

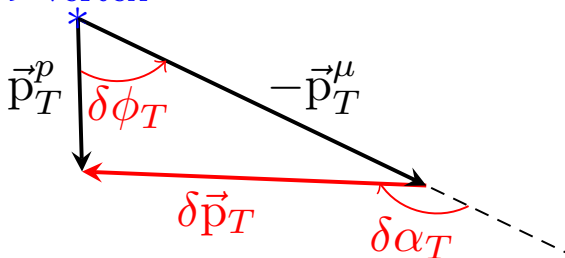
sensitive to FSI: $\delta\alpha_T = \arccos \frac{-\vec{k}'_T \cdot \delta\vec{p}'_T}{k'_T \cdot \delta p'_T}$

sensitive to Fermi Motion:




$$\delta\vec{p}_T = \vec{p}_T^{\vec{p}} + \vec{p}_T^{\vec{\mu}} = \vec{p}_T^{\vec{n}}$$




additional variable: $\delta\phi_T = \arccos \frac{\vec{k}'_T \cdot (\vec{p}_p)_T}{k'_T \cdot (p_p)_T}$

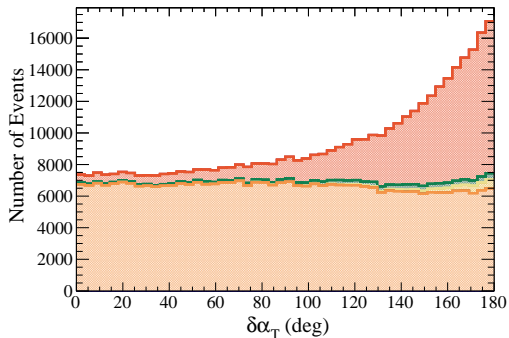
ν vertex



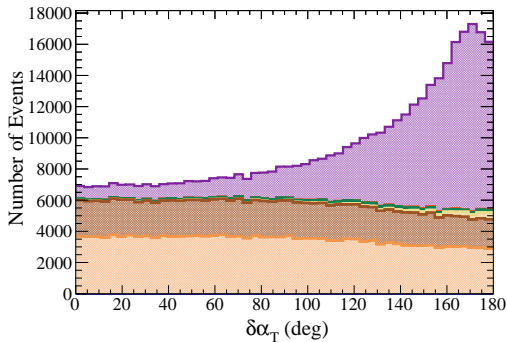
High $\delta\alpha_T$ strongly depends on FSI and is affected by de-excitation and Pauli blocking

-  no cascade FSI
-  QE proton + de-excitation
-  multiple nucleons

-  nuclear clusters production
-  one proton
-  proton + pion



NuWro

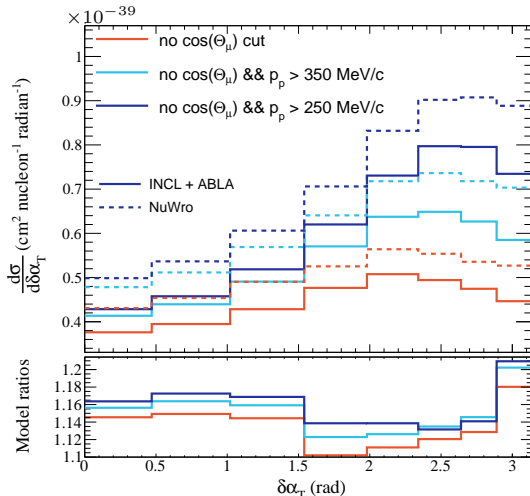
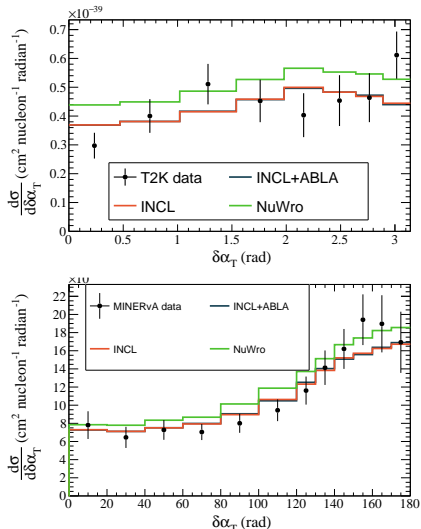


INCL+ABLA

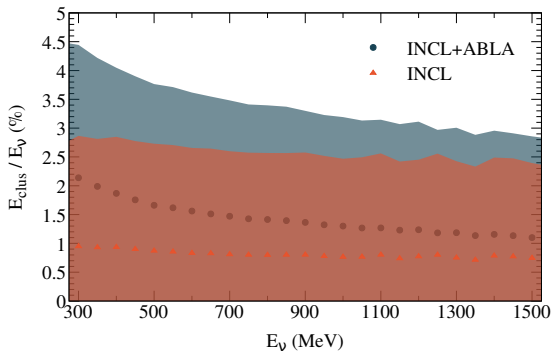
Current detector threshold in ND280 and MINERvA scintillators is **too large**, so we **cannot see the difference** between INCL and

Lower threshold provides better sensitivity to distinguish models

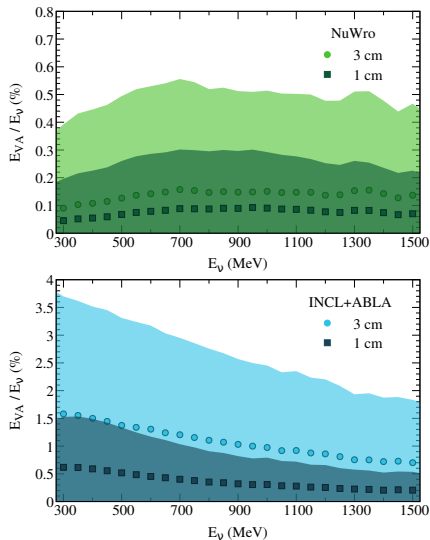
NuWro



The **actual fraction** of neutrino energy going to the kinetic energy of the subleading hadrons is **non-negligible**.



What can be actually seen in the detector (Birks quenching applied):

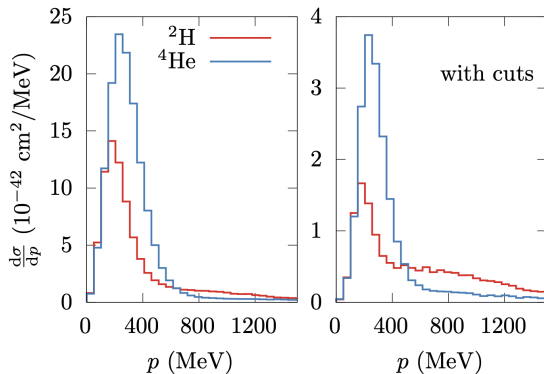


Main goal: describe semi-inclusive one-nucleon knockout on argon for the kinematics of the MicroBooNE experiment

- consistent **vertex simulation:** RPWIA, EDRMF
- **no** initial interaction **excitation** calculation: not enough data
- generators: ACHILLES, NEUT, NuWro, INCL

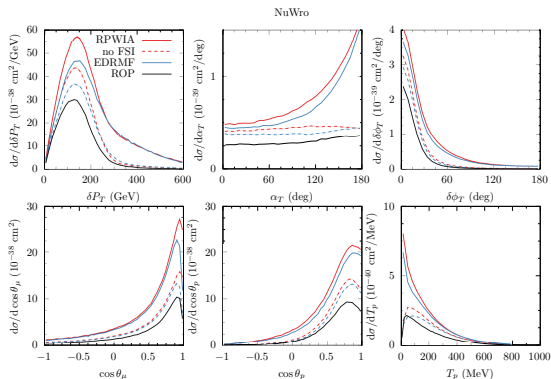
MicroBooNE cuts:

- $100 \text{ MeV} < p_{\mu} < 1200 \text{ MeV}$
- only one proton with momentum (300; 1000) MeV
- no π^0 , no π^{\pm} with $p > 70 \text{ MeV}$

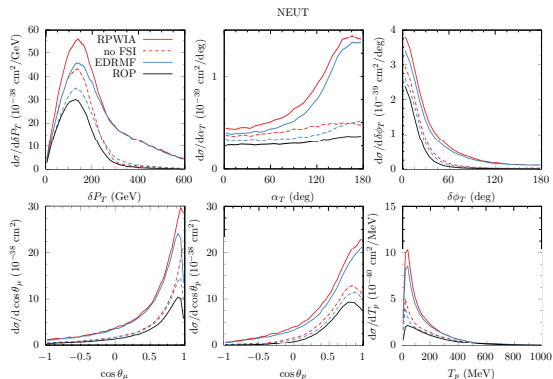


Momenta of α -particles and deuterium produced in INCL+ABLA

Comparison of intranuclear cascade models and the relativistic optical potential



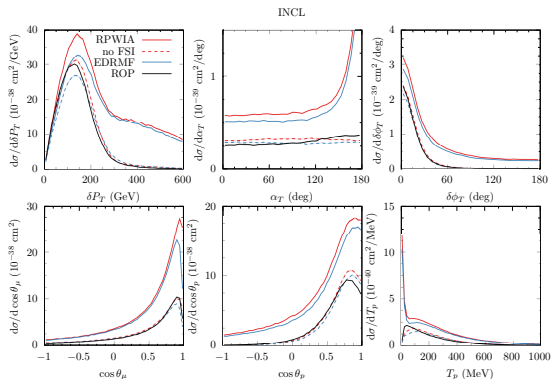
NuWro



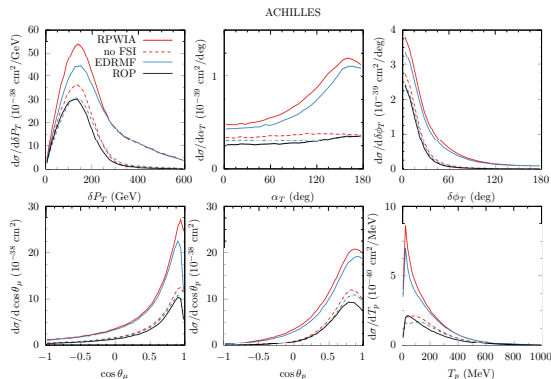
NEUT

Space-like cascades: **NuWro** and **NEUT** compare **well** to ROP

Comparison of intranuclear cascade models and the relativistic optical potential

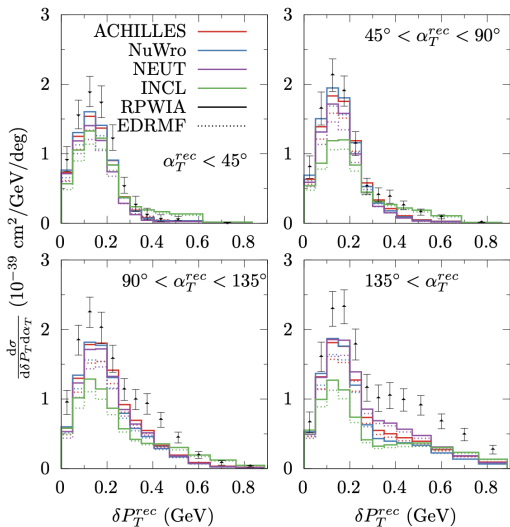


INCL

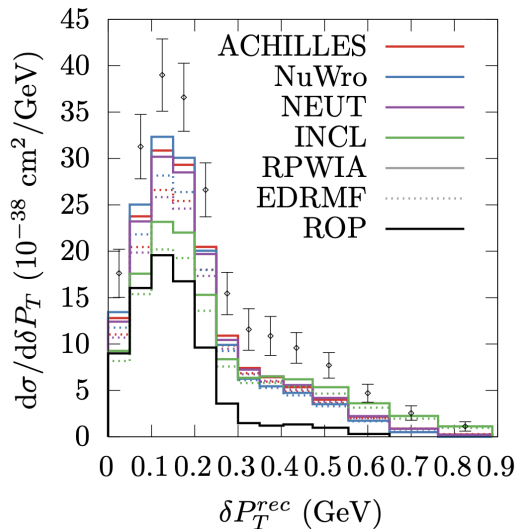


ACHILLES

Time-like cascades: INCL and ACHILLES show remarkable agreement with ROP



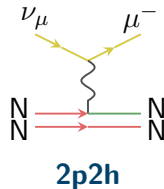
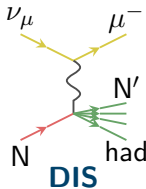
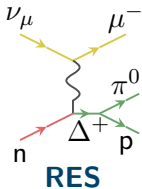
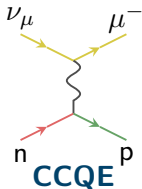
Double differential cross sections



Single differential cross section

- the effort is ongoing to integrate **CCQE** reaction in INCL with subsequent implementation in **Geant4**
- "transparent events" are **no always** transparent: nuclear clusters may be produced
- INCL+ABLA simulation features **important difference** in nucleon kinematics in comparison to NuWro (and the other similar generator used in neutrino scattering)
- An essential novelty of this study is the **simulation of nuclear cluster production** during cascade and de-excitation. It is important for the understanding of the **vertex activity** and calorimetric method of ν **energy reconstruction**
- According to INCL prediction, **nuclear clusters** can be **measurable** in LArTPCs

BACK UP

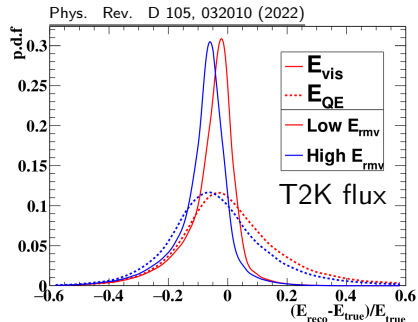


Energy reconstruction using only muon kinematics
(works well for **quasi-elastic reaction**):

$$E_{\nu}^{QE} = \frac{m_p^2 - (m_n - E_B)^2 - m_{\mu}^2 + 2(m_n - E_B)E_{\mu}}{2((m_n - E_B) - E_{\mu} + p_{\mu} \cos \theta_{\mu})}$$

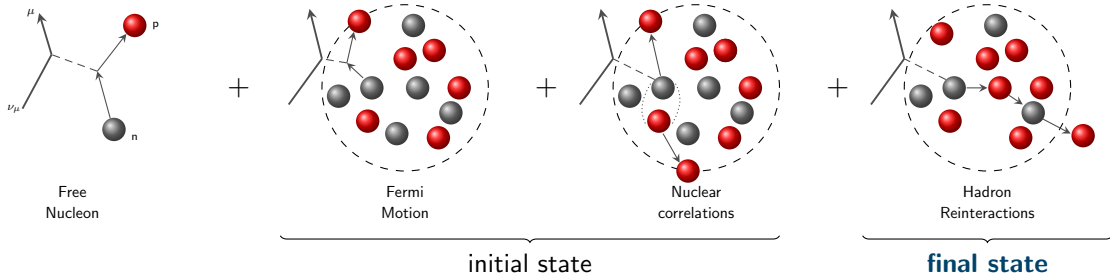
Energy reconstruction using **muon and kinetic energy of the nucleon**:

$$E_{\nu}^{vis} = E_{\mu} + T_N$$

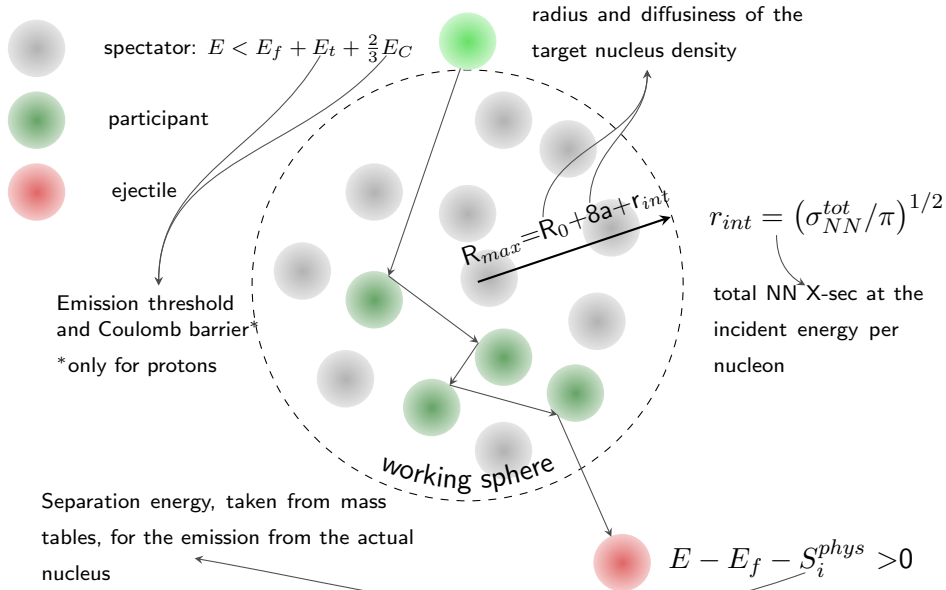


E_{ν}^{vis} , dashed line — QE formula
solid line — $\mu + N$ formula

$\mu + N$ formula gives us more **opportunities**, but also it creates more **challenges** for modelling and we need to **understand better nuclear effects** also on neutrons and protons.

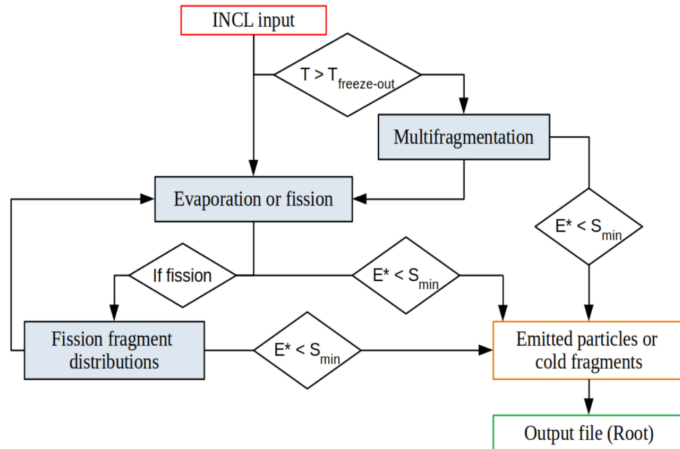


We will focus on **CCQE** ν reaction channel and the **Final State Interactions (FSI)** that are described by **cascade models** and on the nuclear **excitation energy**.



Phys. Rev. C 105, 014623 (2022)

The ablation model **ABLA** describes the **de-excitation** of an excited nuclear system through the emission of γ -rays, neutrons, light-charged particles, and intermediate-mass fragments (or fission in case of hot and heavy remnants).



$$T_{freeze-out} = \max \left[5.5, 9.33e^{(-2.82 \times 10^{-3} A_{rem})} \right]$$

S_{min} —minimum particle separation energy

Using $\mu + p$ is **better** than using muon only, but here we show that we gain even **higher precision** by using all subleading particles

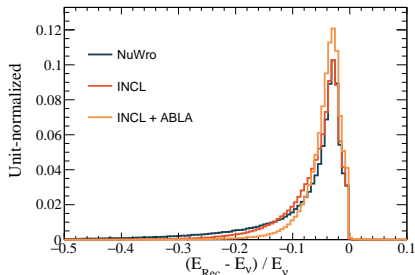
proton only:

$$E_{rec} = E_{\mu} + T_p$$

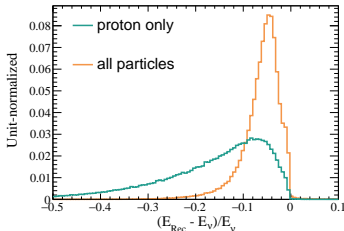


all particles (including clusters)

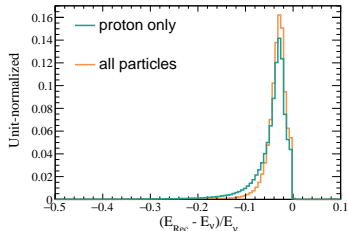
$$E_{rec} = E_{\mu} + \sum_i T_i$$



"all particles" reconstruction



INCL+ABLA cascade FSI



INCL+ABLA no cascade FSI