

Moving up into the 1pi region

Natalie Jachowicz



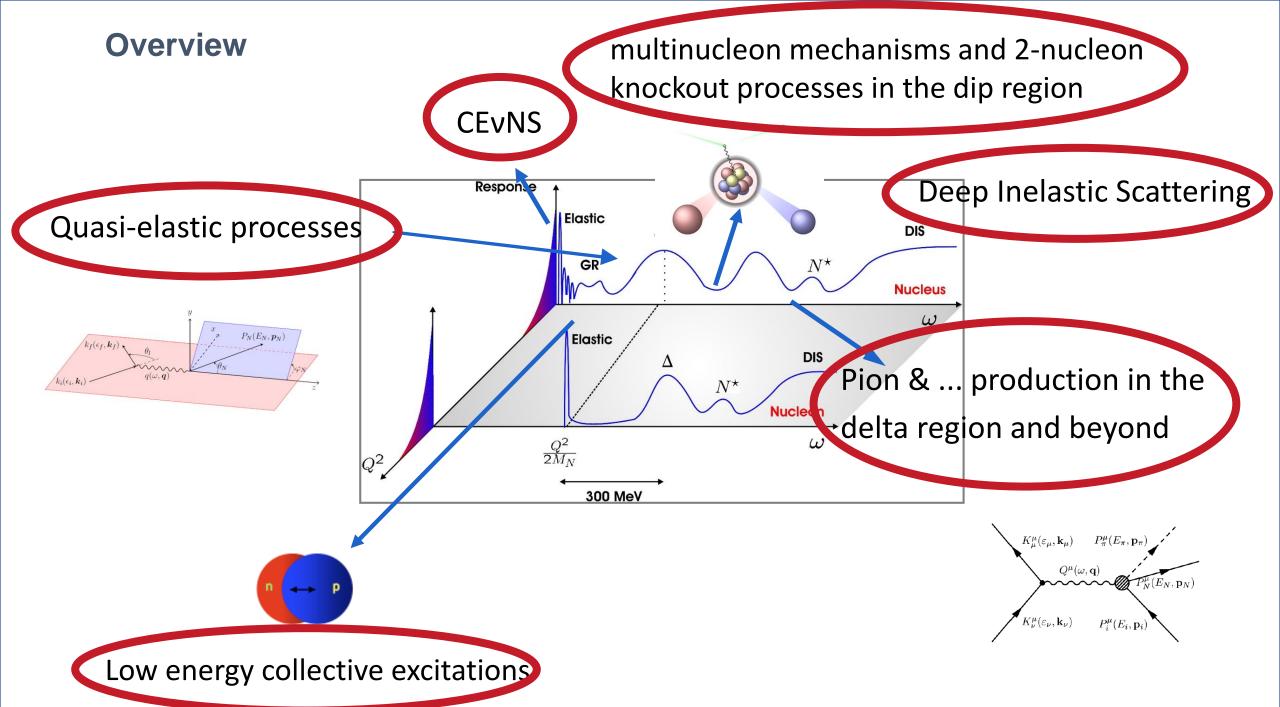


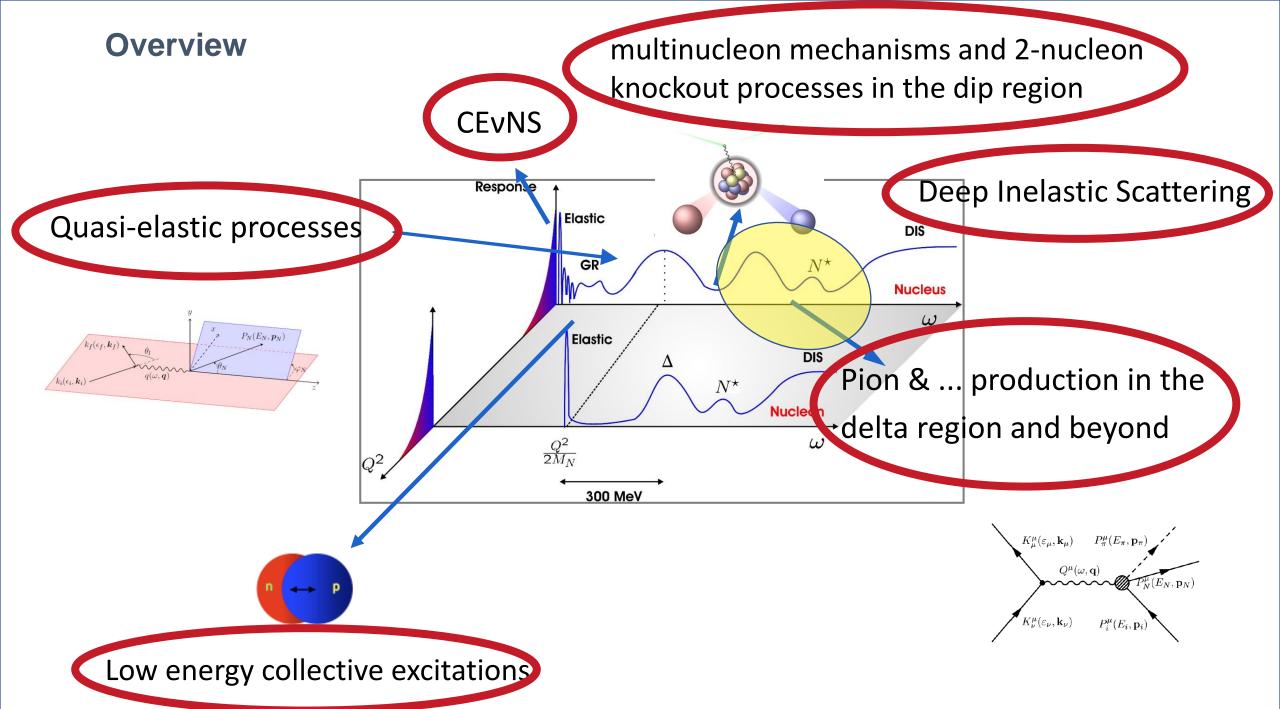
Overvieweof non-CCOpi modelling relevant for

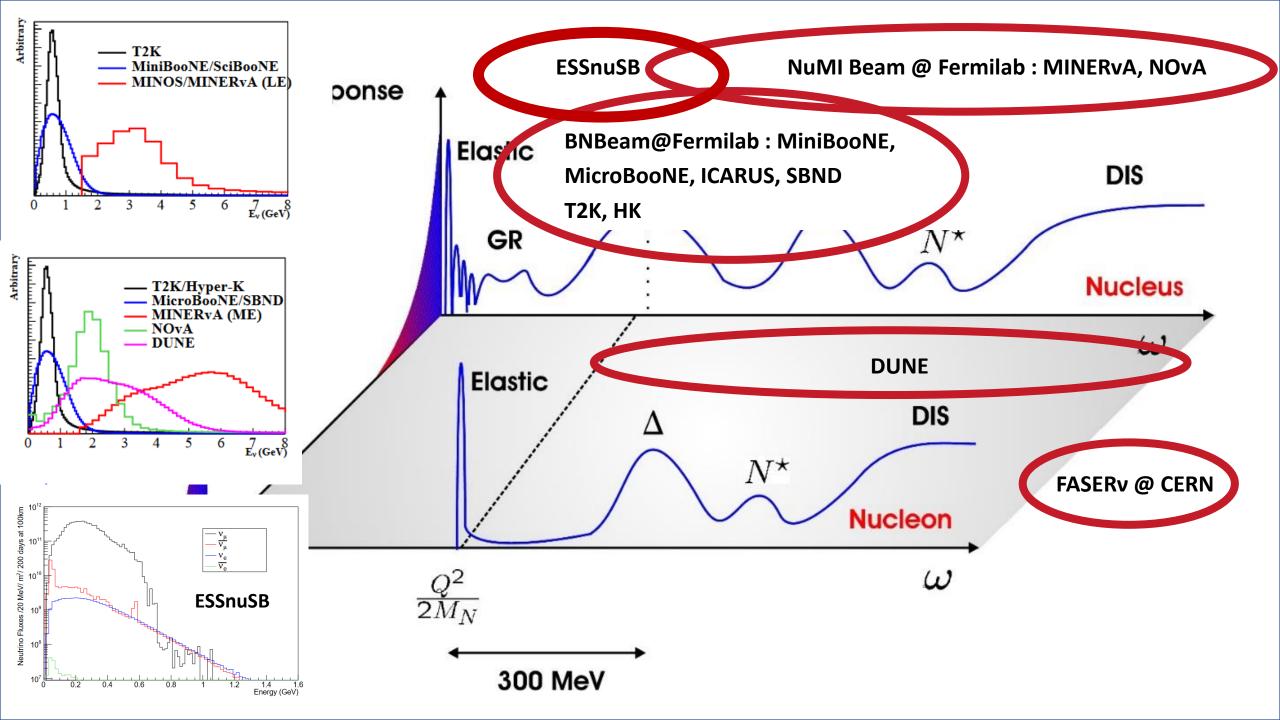
JNE and HK

Natalie Jachowicz

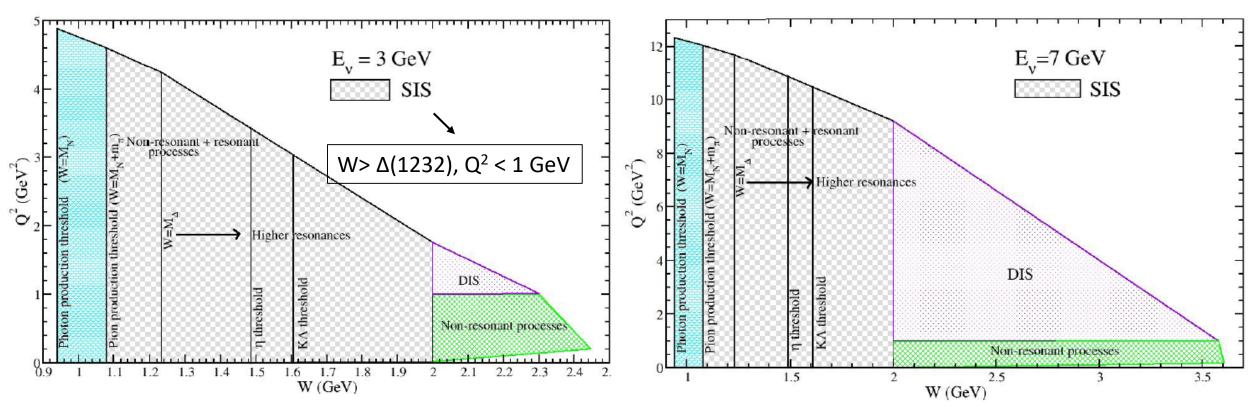








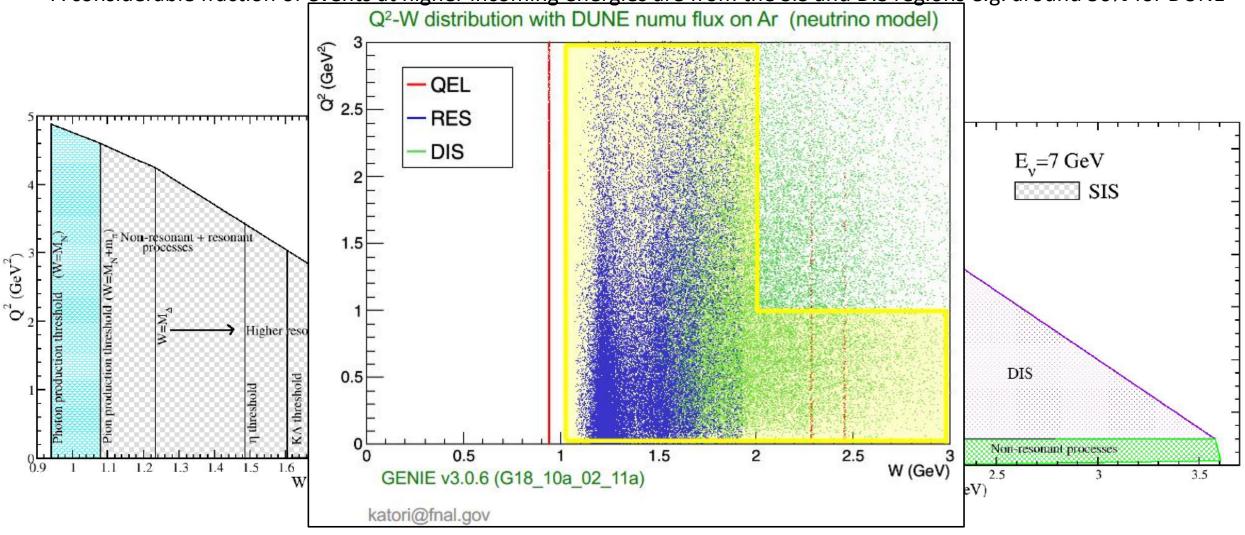
- This kinematic region is not well as well understood as the QE one, both experimentally and theoretically
- A considerable fraction of events at higher incoming energies are from these SIS and DIS regions e.g. around 50% for DUNE



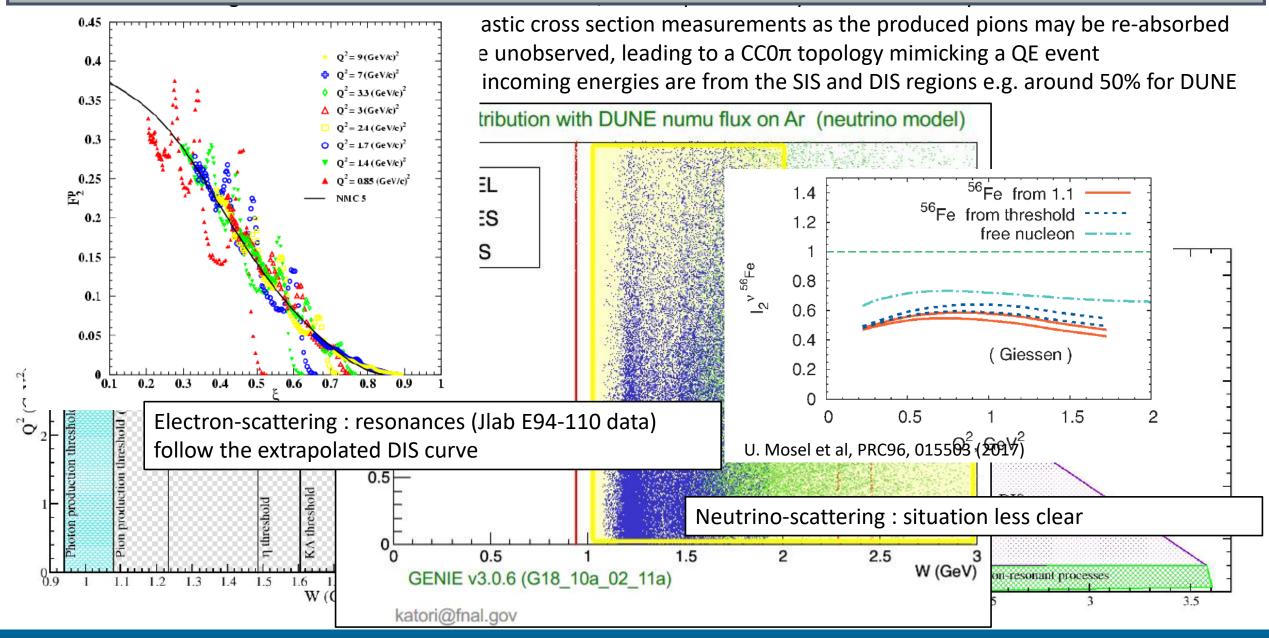
Snowmass WP on theoretical tools for neutrino scattering, L. Alvarez Ruso et al, arXiv:2203.09030

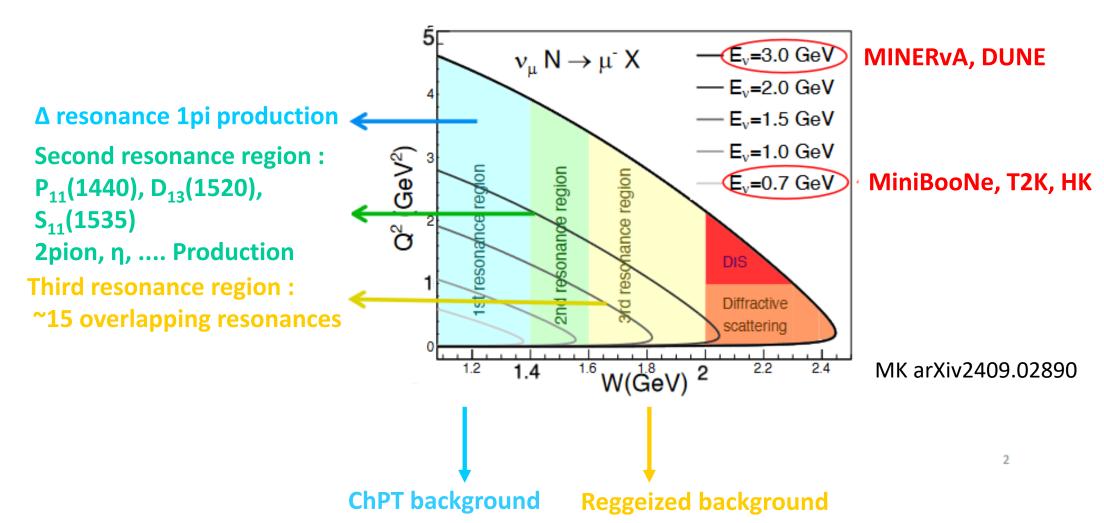
- This kinematic region is not well understood or studied, both experimentally and theoretically
- Important background channel for quasi-elastic cross section measurements as the produced pions may be re-absorbed in the nuclear medium or remain otherwise unobserved, leading to a CC0 π topology mimicking a QE event

• A considerable fraction of events at higher incoming energies are from the SIS and DIS regions e.g. around 50% for DUNE

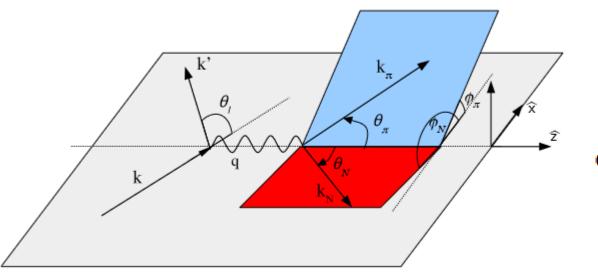


Duality and the transition from nucleon to partonic degrees of freedom





$$d\sigma(E) = \frac{m}{E} \frac{m'}{E'} \frac{d\mathbf{k'}}{(2\pi)^3} \frac{m_N}{E_N} \frac{d\mathbf{p_N}}{(2\pi)^3} \frac{1}{2E_\pi} \frac{d\mathbf{p_\pi}}{(2\pi)^3} \frac{M_B}{E_B} \frac{d\mathbf{p_B}}{(2\pi)^3} \times (2\pi)^4 \delta^4 (K^\mu + P_A^\mu - K'^\mu - P_N^\mu - P_\pi^\mu - P_B^\mu) |\mathcal{M}|^2$$



$$d\sigma(E) = \frac{mm'M_NM_B}{2(2\pi)^8 E E' E_N E_\pi E_B} d\mathbf{k}' d\mathbf{p_N} d\mathbf{p_\pi}$$
$$\delta(E + M_A - E' - E_N - E_\pi - E_B) |\mathcal{M}|^2$$

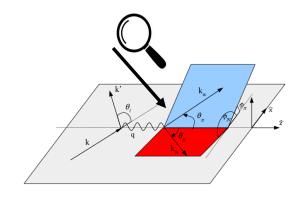
$$\frac{\mathrm{d}^8 \sigma(E)}{\mathrm{d} E' \mathrm{d}\Omega \mathrm{d} E_\pi \mathrm{d}\Omega_\pi \mathrm{d}\Omega_N} = \frac{m_i}{E} \frac{m' k' M_N k_\pi k_N M_B}{2 (2\pi)^8 E_B f_{rec}} |\mathcal{M}|^2$$

$$|\mathcal{M}_W|^2 = \frac{G_F^2 \cos^2 \theta_c}{m'm} L^{\mu\nu} H_{\mu\nu}$$

γ Ημν

Hadron tensor

$$H^{\mu\nu} = \frac{1}{2} \sum_{s: s: t} J^{\mu,*} \left(k_{\pi}, k_{N}, Q, s_{i}, s_{f} \right) J^{\nu} \left(k_{\pi}, k_{N}, Q, s_{i}, s_{f} \right)$$



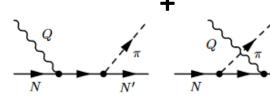
$$J^{\mu}\left(k_{\pi},k_{N},Q,s_{i},s_{f}\right)=\overline{u}\left(k_{N},s_{f}\right)\mathcal{O}^{\mu}\left(k_{\pi},k_{N},k\right)u\left(k_{i},s_{i}\right)$$

Resonances: invariant mass, width Γ , isospin I, spin J, parity P, and for SPP the branching ratio for decay into a pion and nucleon

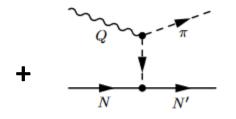
Q π N Res N' N Res N'

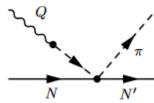
At tree level

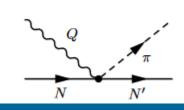
Nucleon pole and cross nucleon pole

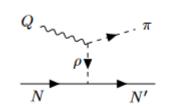


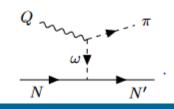
Non-resonant background











+ interferences

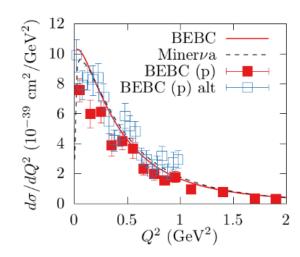
$$J^{\mu} = \sum_{n} \overline{u}\left(k_{N}, s_{f}\right) \mathcal{O}_{n}^{\mu} u\left(k_{i}, s_{i}\right)$$
 $\mathcal{O}_{R_{3/2}}^{\mu} = I_{iso,s} \Gamma_{R\pi N}^{\alpha} \; S_{R,\alpha,\beta}\left(k_{R}\right) \; \Gamma_{QRN}^{\mu\beta}\left(K_{i}, Q\right)$ Vertex function for resonance product

Let's pick only one contribution ... spin 3/2 resonance direct contribution

Vertex function for resonance production

$$\Gamma^{\beta\mu}_{QRN} = \left(\Gamma^{\beta\mu}_{QRN,V} + \Gamma^{\beta\mu}_{QRN,A}\right)\tilde{\gamma}^5$$

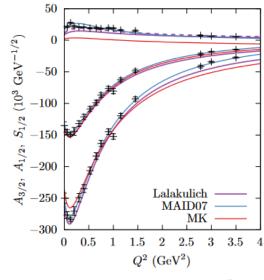
- Vector and axial contribution
- Constraint through CVC and PCAC
- Axial form factors not well constrained,
- Large uncertainties on neutrino data



$$\begin{split} \Gamma_V^{\beta\mu} &= \left[\frac{C_3^V}{M} \left(g^{\beta\mu} \not Q - Q^\beta \gamma^\mu\right) + \frac{C_4^V}{M^2} \left(g^{\beta\mu} Q \cdot k_R - Q^\beta k_R^\mu\right) \right. \\ &+ \left. \frac{C_5^V}{M^2} \left(g^{\beta\mu} Q \cdot k_i - Q^\beta k_i^\mu\right) + C_6^V g^{\beta\mu} \right] \gamma^5, \end{split}$$

$$\begin{split} \Gamma_A^{\beta\mu} &= \frac{C_3^A}{M} \left(g^{\beta\mu} \mathcal{Q} - Q^\beta \gamma^\mu \right) + \frac{C_4^A}{M^2} \left(g^{\beta\mu} Q \cdot k_R - Q^\beta k_R^\mu \right) \\ &+ C_5^A g^{\beta\mu} + \frac{C_6^A}{M^2} Q^\beta Q^\mu, \end{split}$$

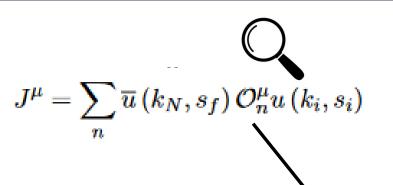
e.g. in Lalakulich parametrization



$$C_3 = \frac{2.13/D_V(Q^2)}{1 + Q^2/4M_V^2}$$

$$C_4 = -1.51/2.13C_3$$

$$C_5 = \frac{0.48/D_V(Q^2)}{1 + Q^2/0.776M_V^2}$$



Let's pick only one contribution ... spin 3/2 resonance direct contribution

$$\mathcal{O}_{R_{3/2}}^{\mu}=I_{iso,s}\Gamma_{R\pi N}^{lpha}\;S_{R,lpha,eta}\left(k_{R}
ight)\;\Gamma_{QRN}^{\mueta}\left(K_{i},Q
ight),$$



Resonance propagator

$$S_3^{\mu\nu} = \frac{k_R + M_R}{k_R^2 - M_R^2 + i M_R \Gamma(W)} \left[g^{\mu\nu} - \frac{1}{3} \gamma^\mu \gamma^\nu - \frac{2}{3} \frac{k_R^\mu k_R^\nu}{M_R^2} + \frac{k_R^\mu \gamma^\nu - k_R^\nu \gamma^\mu}{3 M_R} \right]$$

$$J^{\mu} = \sum_{n} \overline{u}\left(k_{N}, s_{f}\right) \mathcal{O}_{n}^{\mu} u\left(k_{i}, s_{i}\right)$$
 $\mathcal{O}_{R_{3/2}}^{\mu} = I_{iso,s} \Gamma_{R\pi N}^{lpha} \; S_{R,lpha,eta}\left(k_{R}
ight) \; \Gamma_{QRN}^{\mueta}\left(K_{i},Q
ight)$

Let's pick only one contribution ... Spin 3/2 resonance direct contribution



Vertex function for the decay vertex

$$\Gamma^{\mu}_{\pi NR} = rac{\sqrt{2} f_{\pi NR}}{m_{\pi}} k_{\pi}^{\mu} \gamma^5 \tilde{\gamma}^5.$$

Depends on

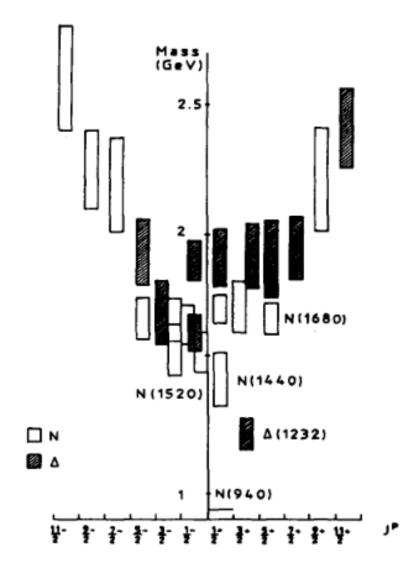
- Coupling constant f_{πNR}
 Resonance width Γ

$$\Gamma_{\pi N}^{3/2}(W) = \frac{I_{iso}}{12\pi} \frac{f_{\pi NR}^2}{m_{\pi}^2} \frac{(k_{\pi}^*)^3}{W} (E_N^* \pm M_N)$$

	$M_R \; ({ m MeV})$	$\Gamma_{exp} \; ({ m MeV})$	$\beta_{\pi N}$	$f_{\pi NR}$
P_{33}	1232	120	1	2.18
S_{11}	1535	150	0.45	0.16
P_{11}	1430	350	0.6	$0.16 \\ 0.49 \\ 1.62$
D_{13}	1515	115	0.6	1.62

$$\Gamma_{\pi N}(M_R) = \beta_{\pi N} \Gamma_{exp}$$

Experimentally determined full width + branching ratio





... and these are only resonances ...

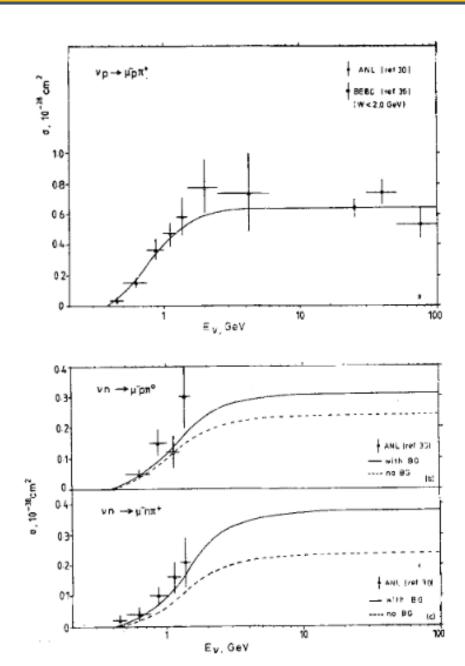
M. Giannini, PPNP 24, 253 (1990)

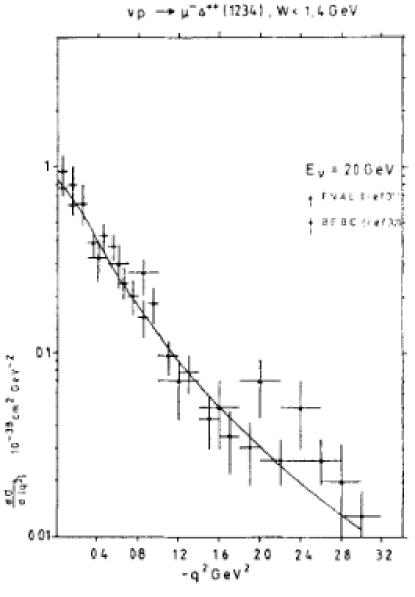
Rein Sehgal model

D. Rein and L.M. Sehgal, Annals of Physics 133, 79 (1981)

This is an attempt to describe all existing data on neutrino production of single pions in the resonance region up to W = 2 GeV in terms of the relativistic quark model of Feynman, Kislinger and Ravndal (FKR). We considered single pion production to be mediated by all interfering resonances below 2 SeV. A simple noninterfering, nonresonant background of isospin 1/2 was added. It improved agreement with experiment, particularly in the ratio of isospin amplitudes in charged current reactions, at the expense of one additional constant. All total cross sections, cross section ratios and W-distributions are well reproduced at low and high energies, with charged and neutral currents (supposing the Salam-Weinberg theory with $\sin^2 \theta_w \approx \frac{1}{4}$ to be correct), and for neutrinos and antineutrinos, giving predictions where data are lacking. New predictions have been made for complex angular distributions in \mathcal{N}_{π} channels exhibiting strong interference between neighbouring resonances. These are sensitive (for 1.1 GeV $\lesssim W \lesssim$ 1.5 GeV) to the sign of the Roper resonance $P_{11}(1450)$ which is controversial in photoproduction experiments.

Rein Sehgal model





D. Rein and L.M. Sehgal, Annals of Physics 133, 79 (1981)

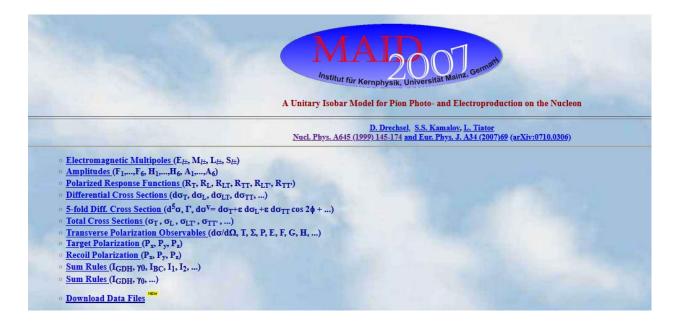
MAID

- Mainz unitary isobar model
- Designed to analyze world data on π photoproduction and (e,e' π)
- Extensions to kaon production, eta production, γππ
- 13 four-star resonances with masses below 2 GeV, unitarized Breit-Wigner shapes
- Background unitarized using K-matrix formalism

https://maid.kph.uni-mainz.de

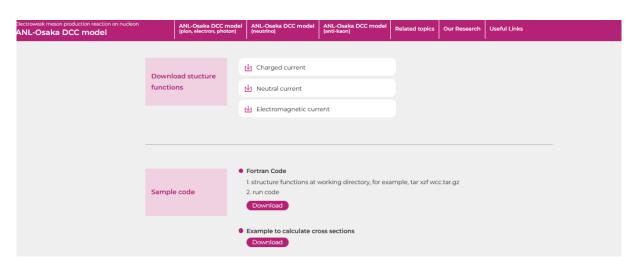
channel	$d\sigma$	Σ	T	P	total
$n\pi^+$	4646	760	645	205	6256
$p\pi^0$	4936	673	353	540	6502
$p\pi^-$	1554	206	94	88	1942

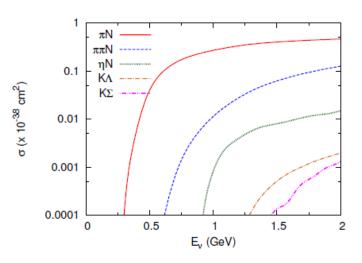
SAID database



ANL-Osaka DCC

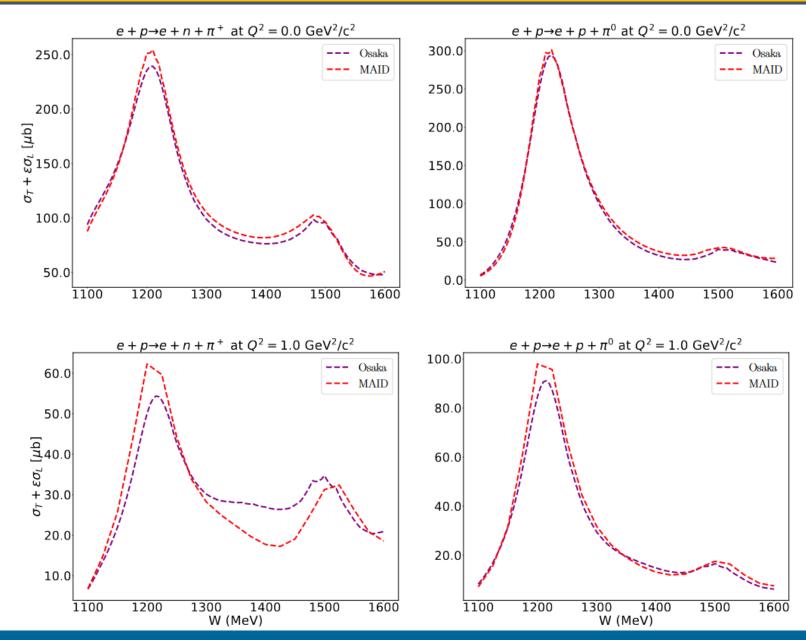
- predicts πN , γN partial wave amplitudes used to determine cross sections for πN , ηN , $K\Lambda$ and $K\Sigma$ scattering, photoproduction of these channels, 2π production off the nucleon, pion electro- and neutrinoproduction
- Up to W~2 GeV, including ~20 N* states
- Fit to ~30.000 data points for foto pion nucleon, hadron and pion production, and electroproduction for different incoming energies,
- Dynamic coupled channel formalism, different reaction mechanisms contribute in a consistent way to the final state



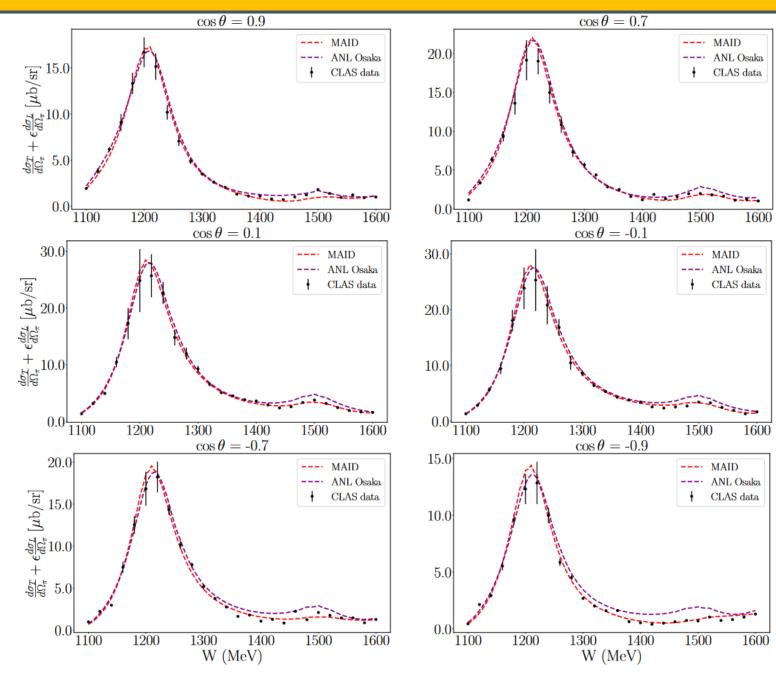


https://www.phy.anl.gov/theory/research/anl-osaka-pwa/ S.X. Nakamura, H. Kamano, T. Sato, PRD92, 074024 (2015) https://www.rcnp.osaka-u.ac.jp/~anl-osk/neutrino/index.html

MAID vs ANL-Osaka

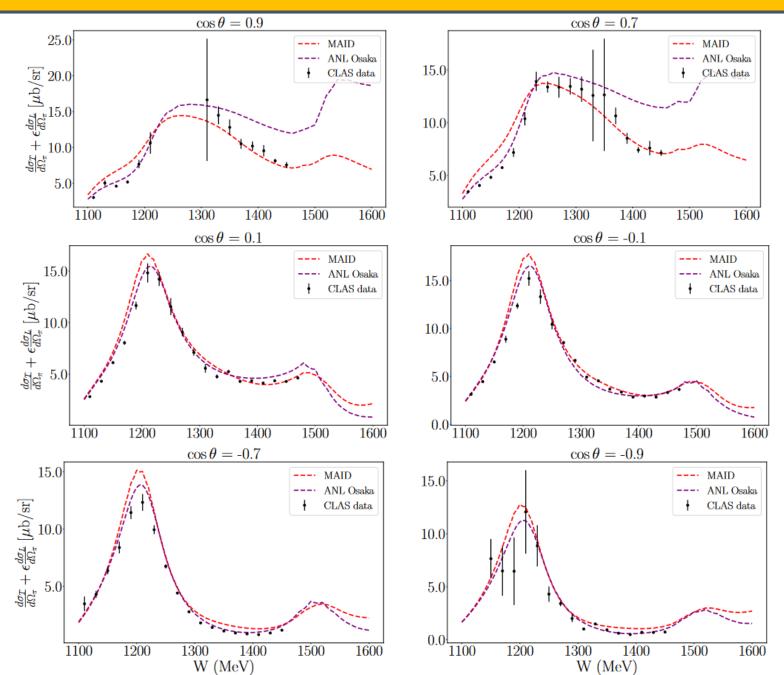


MAID vs ANL-Osaka



$$e + p \rightarrow e + p + \pi^0$$

MAID vs ANL-Osaka



$$e + p \rightarrow e + n + \pi^+$$

MK model

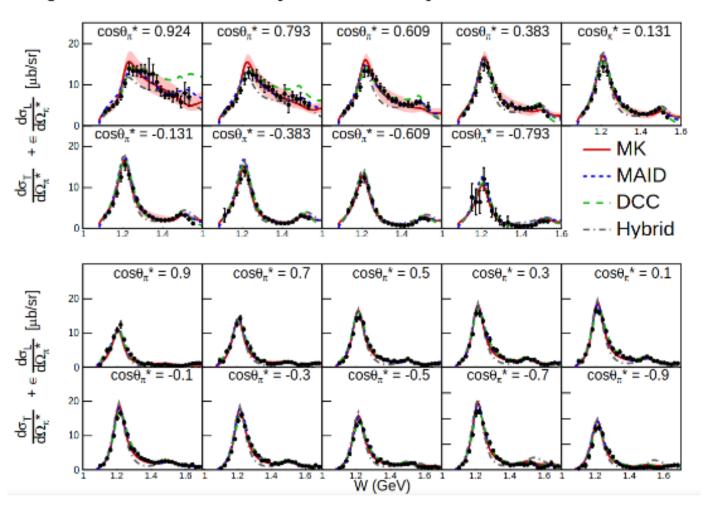
- Single pion production off the nucleon
- developed aiming at describing neutrino induced processes!
- RS with 17 resonances W<2GeV
- Non-resonant background HNV model + Regge
- Fit of a lot of parameters to data
- $d\sigma/dWdQ^2d\Omega_{\pi}$

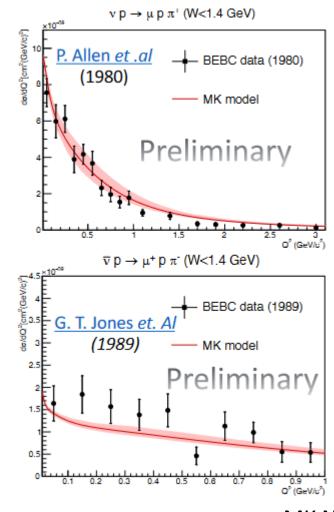
# data point	Photon, electron, pion, Neutrino Channels	Q ² Range (GeV/C) ²	W Range GeV	Form Factors	
≈ 9800	$\gamma \: p \to n + \pi^+ \: , \ \ \gamma p \to p + \pi^0$	0	1.08 - 2.0	Proton	
≈ 31000	$ep \rightarrow en + \pi^+, ep \rightarrow ep + \pi^0$	0.16 - 6.0	1.08 - 2.0	∀ ec	
≈ 2500	$\gamma n \rightarrow p + \pi^-$	0	1.08 - 2.0	Neutron S	
≈ 700	NEW en \rightarrow ep $+ \pi^-$	0.4 - 1.0	1.08 - 1.8		
≈ 400	$\pi^+ p \rightarrow p + \pi^+$, $\pi^- p \rightarrow p + \pi^-$	0	1.08 - 2.0		
<100	$\nu N \rightarrow l^- N + \pi$, $\bar{\nu} N \rightarrow l^+ N + \pi$	Q ² >0 Integrated	1.08 – 2.0 Integrated	Axial-Vector	

MK model

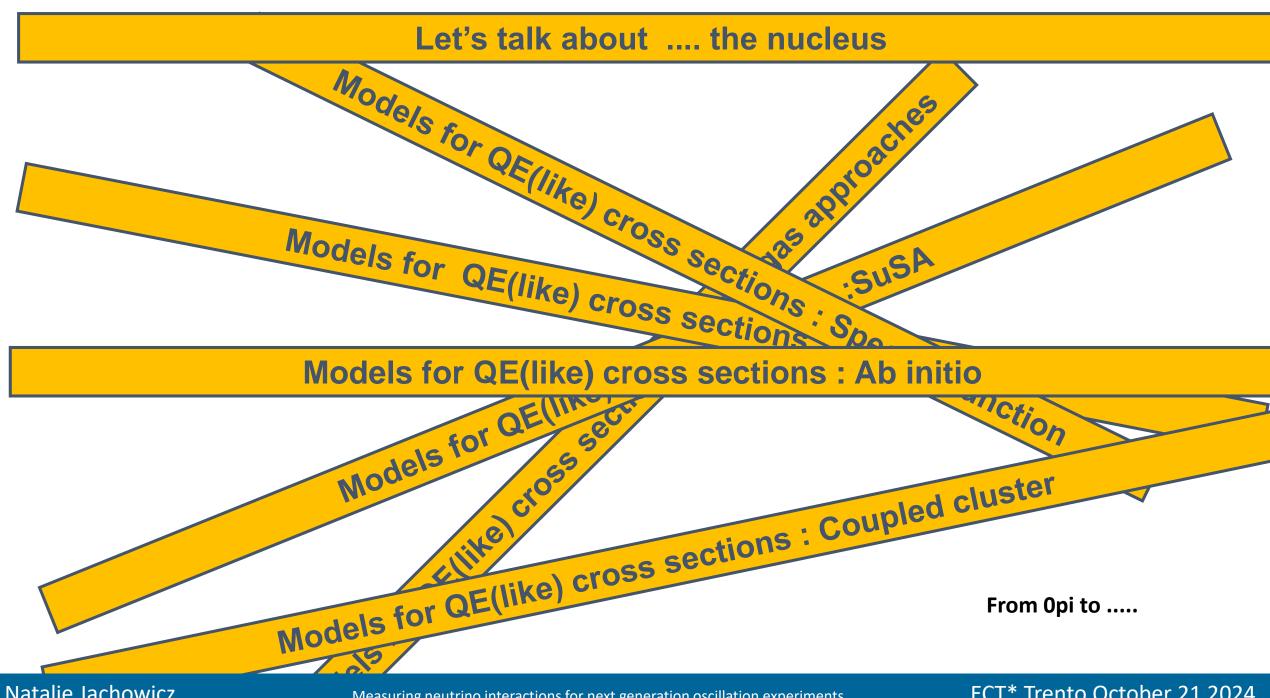
Electron-induced SPP: high quality proton target data

Figures from M. Kabirnezhad [arxiv:2203.15594]





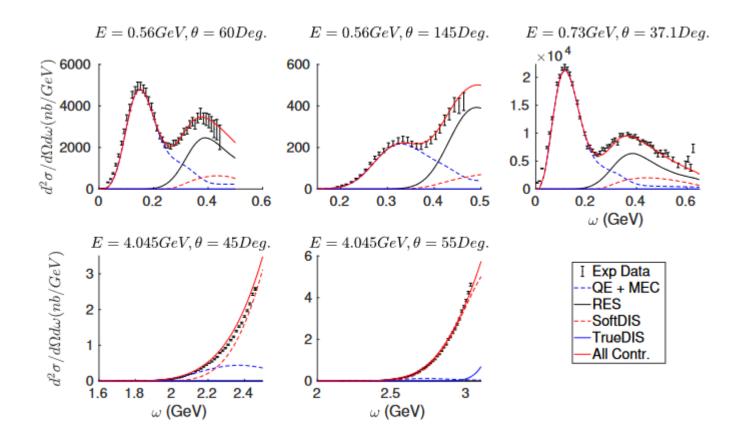
MK NuInt 2024



Let's talk about the nucleus : DCC implementation in SuSA

Incorporating:

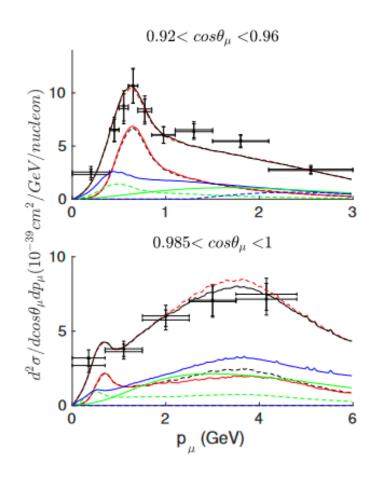
- ANL-Osaka DCC results in SuSAv2
- Inelastic contributions: Bodek-Ritchie, Bosted-Christy or parton distribution functions

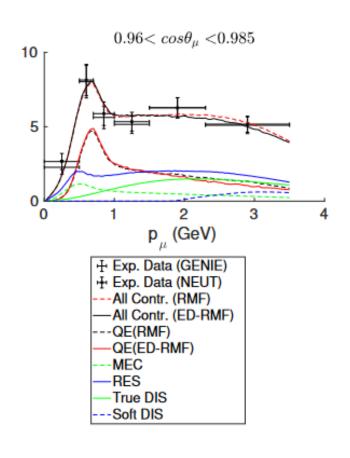


Inclusive e-12C

Gonzalez-Rosa et al.PRD 108, 113008 (2023).

Let's talk about the nucleus : DCC implementation in SuSA

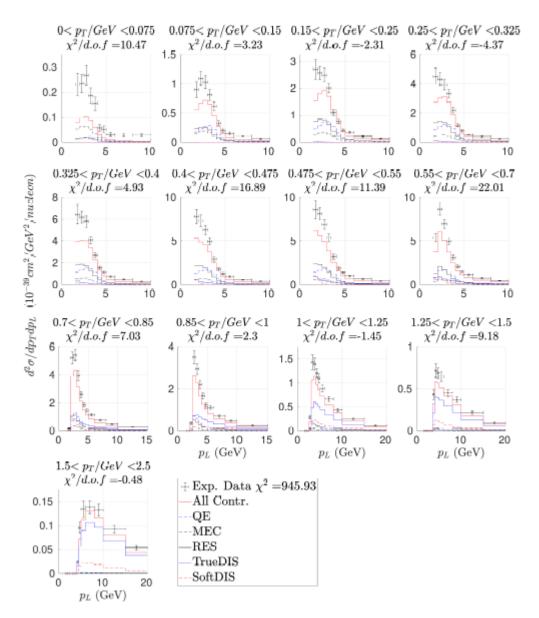




T2K inclusive flux-folded CC

J. Gonzalez-Rosa et al. PRD 108, 113008 (2023).

Let's talk about the nucleus : DCC implementation in SuSA

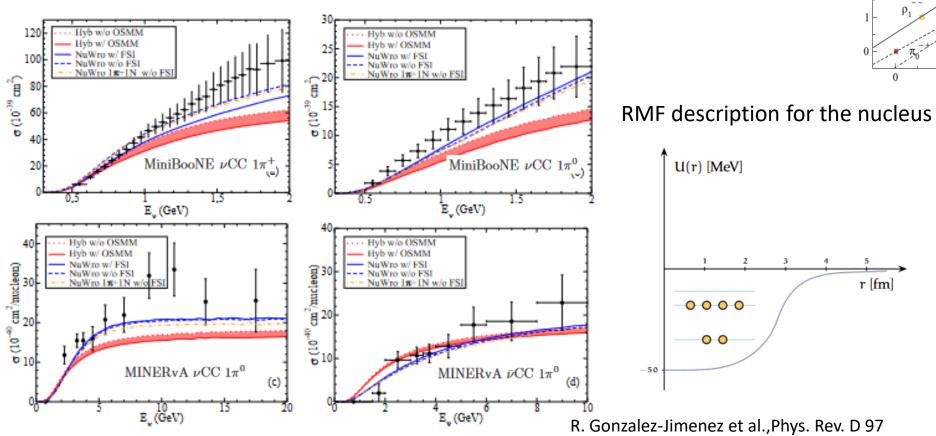


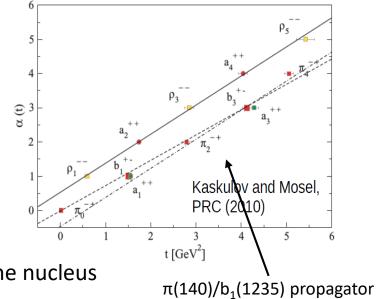
 $1.5 < p_L/GeV < 2$ $2 < p_L/GeV < 2.5$ $2.5 < p_L/GeV < 3$ $3 < p_L/GeV < 3.5$ $\chi^2/d.o.f = 114.84$ $\chi^2/d.o.f = 127.67$ $\chi^2/d.o.f = -0.01$ $\chi^2/d.o.f = 125.47$ 0.5 0.5 0.5 0.5 $4.5 < p_L/GeV < 5$ $3.5 < p_L/GeV < 4$ $4 < p_L/GeV < 4.5$ $5 < p_L/GeV < 6$ $(10^{-39}cm^2/GeV^2/nucleon)$ $\chi^2/d.o.f = 60.46$ $\chi^2/d.o.f = 10.5$ $\chi^2/d.o.f = 5.9$ $\chi^2/d.o.f = -118.41$ 0.5 $6 < p_L/GeV < 7$ $7 < p_L/GeV < 8$ $8 < p_L/GeV < 9$ $9 < p_L/GeV < 10$ $\chi^2/d.o.f = -6.78$ $\chi^2/d.o.f = 38.14$ $\chi^2/d.o.f = 31.2$ $\chi^2/d.o.f = 50.08$ 0.6 0.5 $15 < p_L/GeV < 20$ $10 < p_L/GeV < 15$ $20 < p_L/GeV < 40$ $40 < p_L/GeV < 60$ 0.1 $\chi^2/d.o.f = 54.25$ $\chi^2/d.o.f = 9.46$ $\chi^2/d.o.f = 122.91$ $\chi^2/d.o.f = 95.89$ 0.2 0.004 0.02 0.05 0.1 0.002 $p_T(GeV)$ $p_T(GeV)$ $p_T(GeV)$ $p_T(GeV)$

Minerva longitudinal and transverse muon momentum bins

Ghent Hybrid model

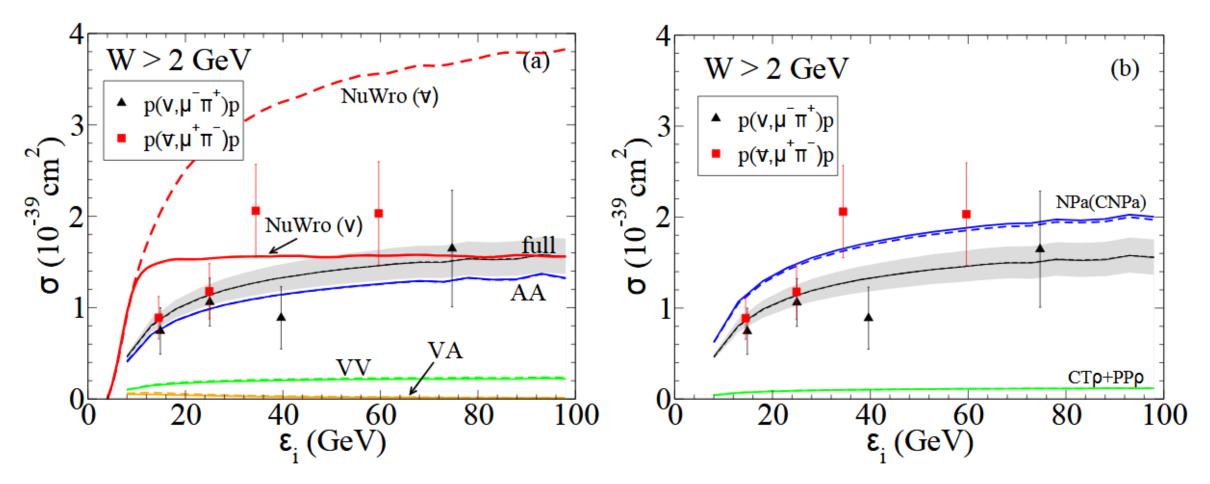
- Designed aiming at the description of pion production off the nucleus
 - Low energy model: resonances + background from ChPT in HNV description
 - High energy model: Reggeized background to overcome shortcomings of the low energy description
- Aiming at fully exclusive description





Ghent Hybrid model

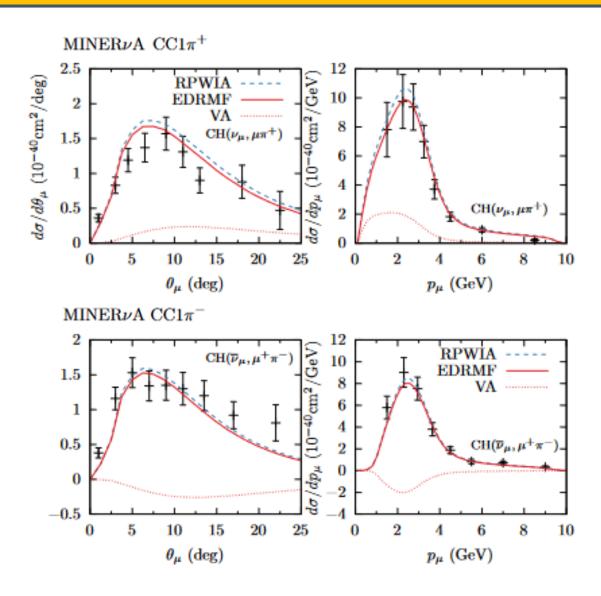
Going beyond W>2 GeV

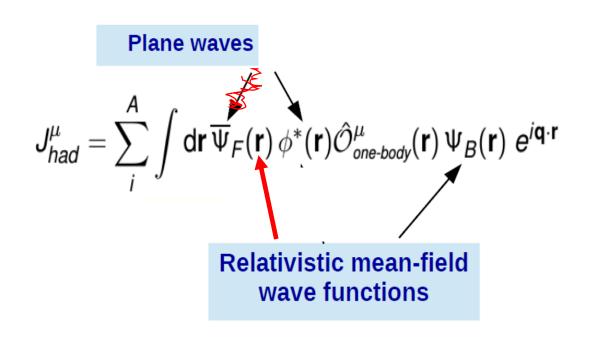


R. Gonzalez-Jimenez et al., Phys. Rev. D 97

Data : P. Allen et al. Nucl. Phys. B264, 221 (1986).

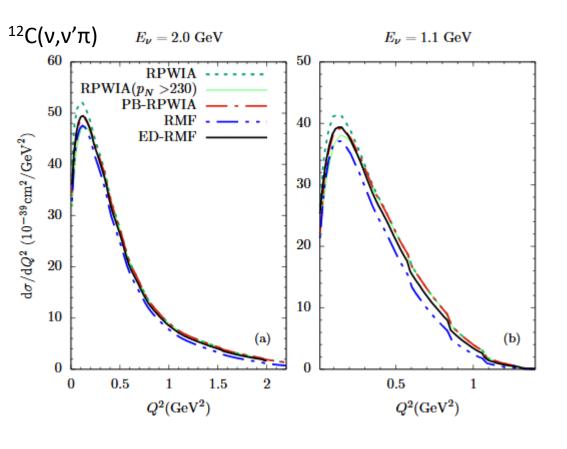
Let's talk about The nucleus

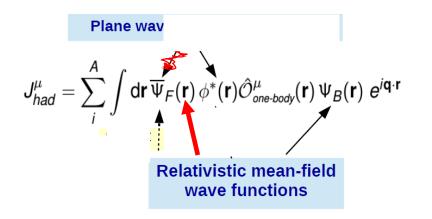


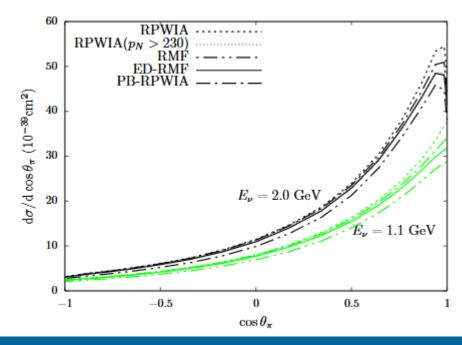


A. Nikolakopoulos, PhD UGent 2021

Let's talk about The nucleus

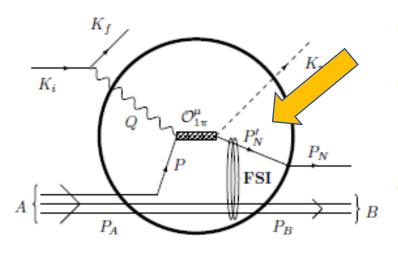




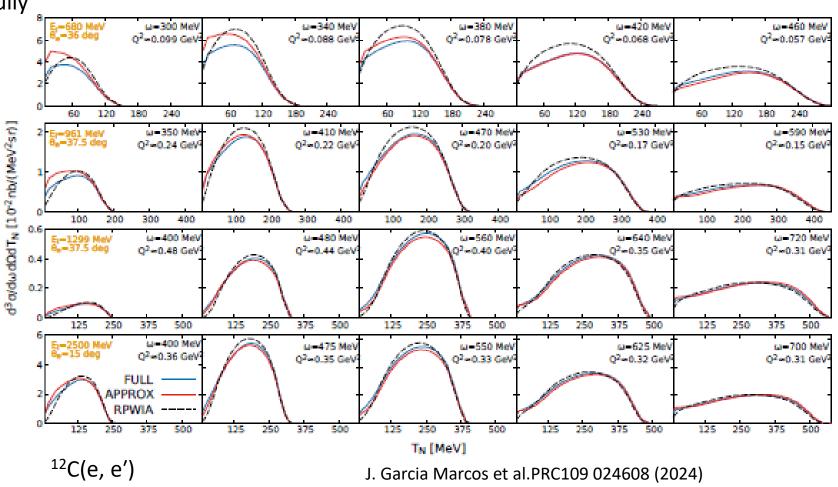


Let's talk about The nucleus

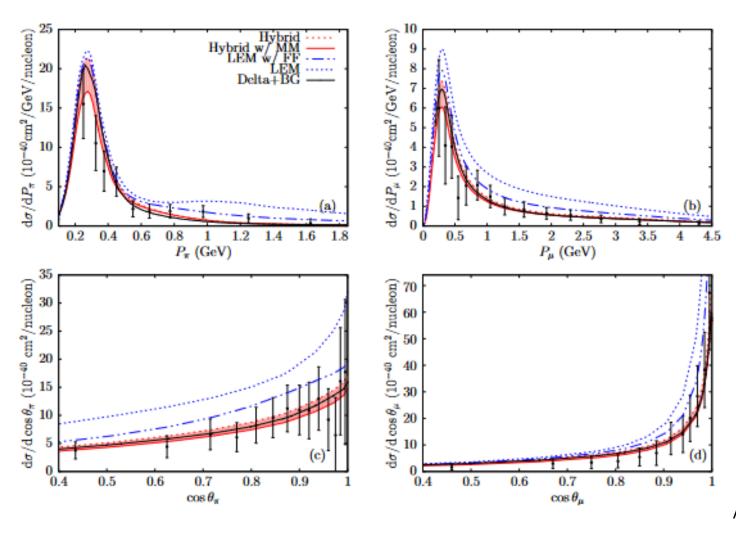
Beyond the asymptotic approximation: fully microscopic calculation of nucleon momentum in the nuclear medium







Let's talk about the Δ in the medium

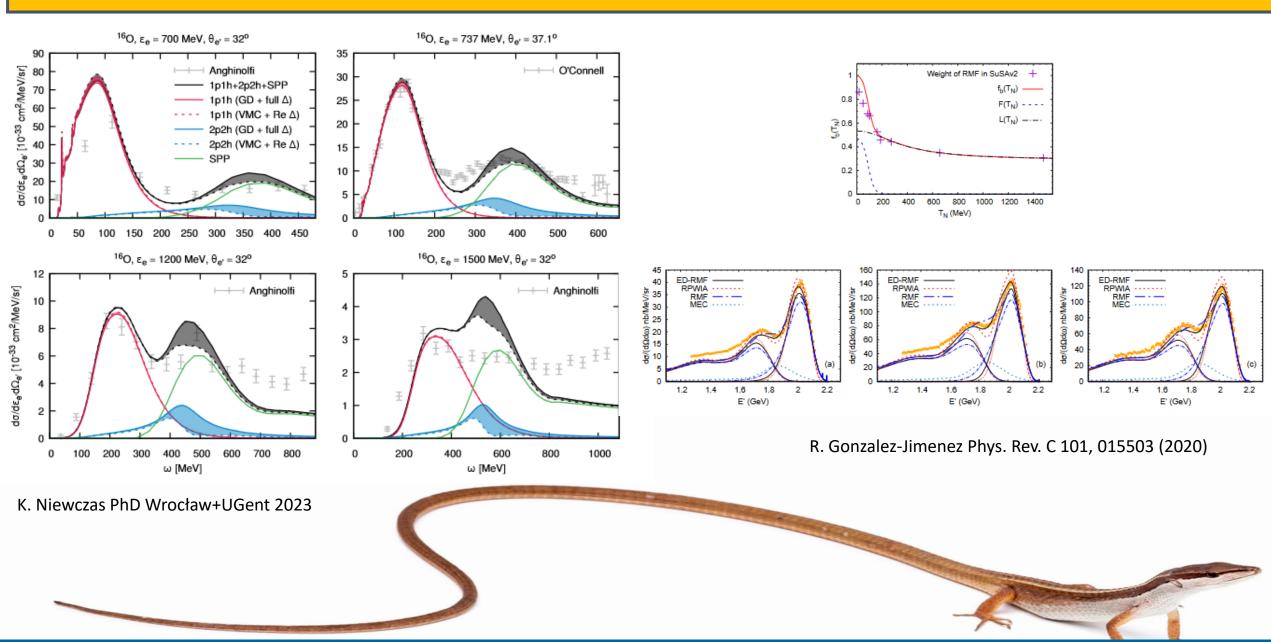


Most commonly used

- Oset-Salcedo medium modifications
- Evaluated in FG model

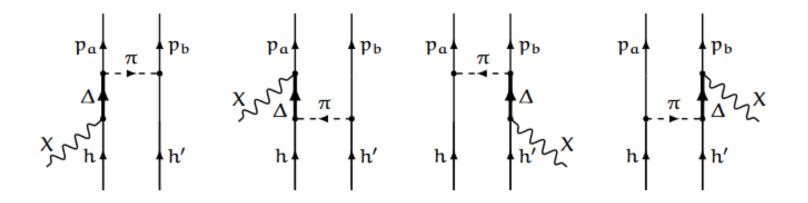
A. Nikolakopoulos, PhD thesis UGent 2021

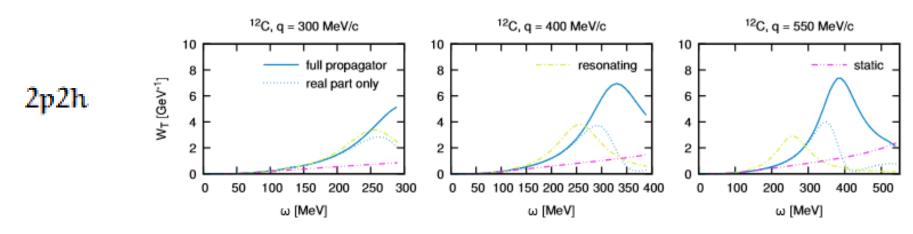
Let's talk about tails



Let's talk about consistency

Consistency of Delta description with MECs description?

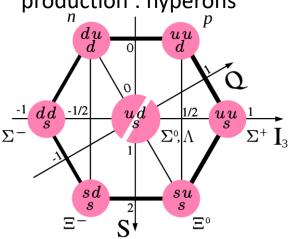


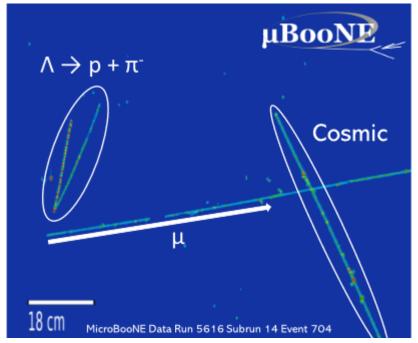


K. Niewczas, PhD thesis 2023

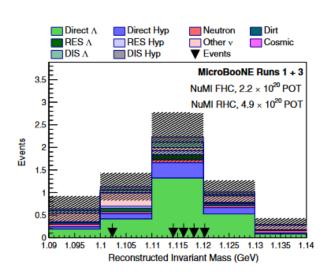
Let's talk about strange things

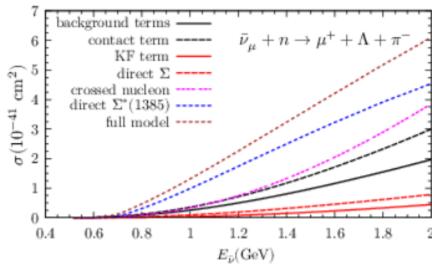
Neutrino strangeness production : hyperons

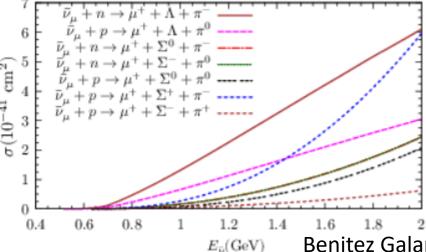




MicroBooNe Phys. Rev. Lett. 130, 231802



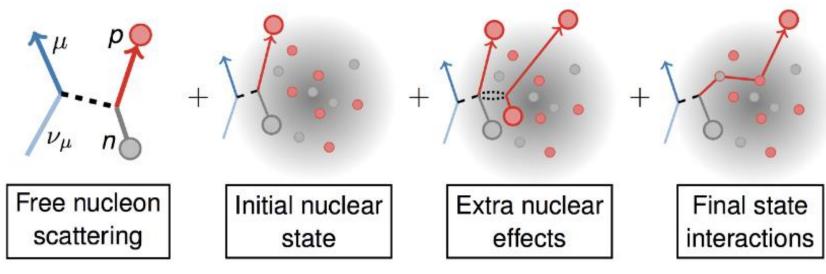




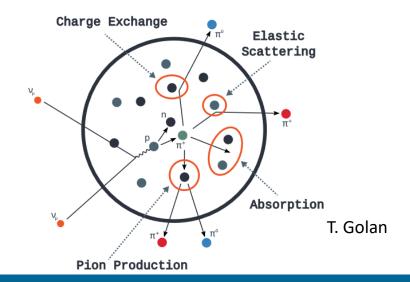
Pion production at low energies!

Benitez Galan, Rafi Alam, Ruiz Simo, PRD104 (2021)

What is going on between initial and final state ...?



K. Niewczas @ NuFACT2021



- Partially taken into account by some of the nuclear models
- MC Generators used for oscillation analyses tend to rely on efficient but approximate models

Event generators: bridges between theory and data

Challenges in generator developments:

• Microscopic input for vertex – exclusive description

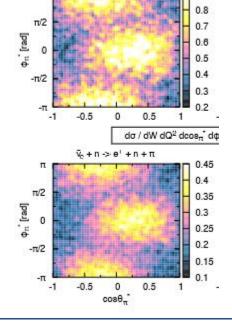
V_c + p -> e + p + π¹

- o Default NuWro
- Free nucleon
- Fixed kinematics:

$$E = 1 \text{ GeV}$$

$$Q^2 = 0.1 \text{ GeV}^2$$

$$W = 1230 \text{ MeV}$$

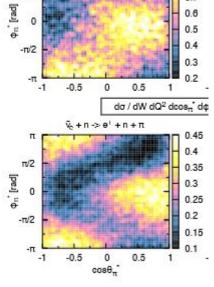


- o Ghent LEM
- Free nucleon
- Fixed kinematics:

$$E = 1 \text{ GeV}$$

$$Q^2 = 0.1 \text{ GeV}^2$$

$$W = 1230 \text{ MeV}$$

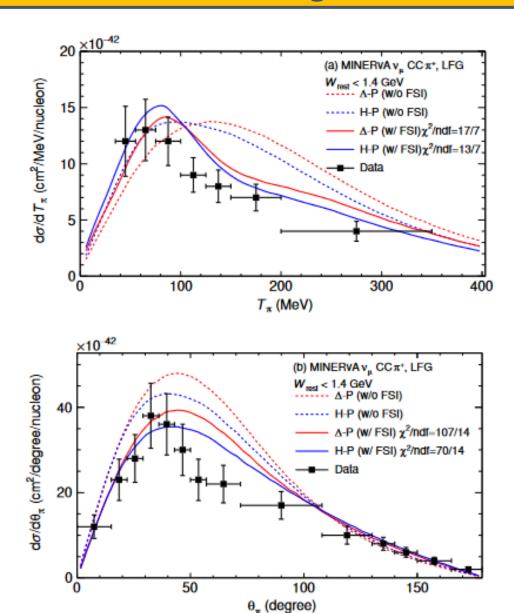


V_c + p -> e + p + π

Need to go beyond isotropic approaches in MC generators and include detailed exclusive microscopic calculation results in the simulation

K. Niewczas et al. PRD 103, 053003 (2021)

Event generators: bridges between theory and data



Integration of Ghent hybrid model in NuWro for pion production off the nucleus

Q. Yan et al. arXiv 2405.05212

Summarizing ...

- Single pion production in the Delta region is relatively well understood
- At higher kinematics, for more complex final states and especially for W>2GeV:
 we'll have to go on opening

