Spin Hydrodynamics from the Entropy Principle: flaws and remedies

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— Spin and quantum features of QCD plasma —

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General Introduction

Thermal System with Vorticity A black hole / a QGP has no hair:

Stable black holes are characterized by



[From Forbes]



Thermal System with Vorticity Gunther et al. (2016) **Baryon Density** 300 **Speculated Phase Diagram** 250 $\frac{T}{\mathrm{MeV}}$ 200 Temperature T **Quark-Gluon Plasma** sQGP 150100**Hadronic** Phase Ouarkvonic Regime Liquid-Gas **Crystalline States** 50100 150200 250300 3500 400 **Color Superconductors** $\frac{\mu_B}{\text{MeV}}$ Nuclear Superfluid Chemical Potential UR $\mathcal{S}^{\mu} = su^{\mu} \quad j^{\mu} = nu^{\mu}$ **Ideal hydrodynamics:** $\partial_{\mu}T^{\mu\nu} = 0 \quad \rightarrow \quad \partial_{\mu}\mathcal{S}^{\mu} = 0$

Eliminating $(\partial \cdot u) \rightarrow n\dot{s} - s\dot{n} = 0 \rightarrow s/n = (const.)$

Thermal System with VorticityGrand Canonical Descriptions[EoSs]
$$de = Tds + \mu dn + \omega_{\mu\nu} dJ^{\mu\nu}$$
 $\langle n \rangle = \frac{\partial p}{\partial \mu}$ $dp = sdT + nd\mu + J_{\mu\nu} d\omega^{\mu\nu}$ $\langle J^z \rangle = \frac{\partial p}{\partial \omega^z}$ n : baryon density J : angular mom. μ : baryon chemical pot. \mathcal{J} : angular mom. ω : spin chemical pot. \mathcal{L} : counterparts?

Thermal System with Vorticity

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Ahmed-Cong-Kubiznak-Mann-Visser (2023)

fixed	$(J,\mathcal{V},C):$	$F \equiv E - TS$	Canonical
fixed	$(ilde{\Omega}, \mathcal{V}, C)$:	$W \equiv E - TS - \tilde{\Omega}J$	Grand Canonical
fixed	(J, \mathcal{V}, μ) :	$G \equiv E - TS - \mu C$	



Thermal System with Vorticity

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Ahmed-Cong-Kubiznak-Mann-Visser (2023)

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fixed	(J, \mathcal{V}, μ) :	$G \equiv E - TS - \mu C$	



In the grand canonical case, a first-order Hawking-Page transition is seen.

> Consistent with Chen-Zhang-Li-Hou-Huang (2020) 7



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Thermal System with Vorticity , Merzik, Merzik, Merzik, Merzik, Merzik, Merzik, Merzik, M **EoS calculable from the rotating HRG model Pressure:** ssure: $p_i^{\pm} = \pm \frac{T}{8\pi^2} \sum_{\ell=-\infty}^{\infty} \int_{(\Lambda_{\ell}^{\mathrm{IR}})^2} dk_r^2 \int dk_z \sum_{\nu=\ell}^{\ell+2S_i} J_{\nu}^2(k_r r)$ $\times \log \left\{ 1 \pm \exp[-(\varepsilon_{\ell,i} - \mu_i)/T] \right\}$ = 0.1 GeV1.2 $\omega = 0.2 \text{GeV}$ $\omega = 0.3 \text{GeV}$ 1.0 $\varepsilon_{\ell,i} = \sqrt{k_r^2 + k_z^2 + m_i^2} - (\ell + S_i)\omega$ 0.8 م ا_{0.6} 0.4 $\Lambda_{\ell}^{\rm IR} = \xi_{\ell,1} \omega$ 0.2 0.0 0.03 0.06 0.09 0.12 0.15 0.00 0.18 *r* [GeV⁻¹] Fujimoto-Fukushima-Hidaka (2021)



Power counting??

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∆ (GeV)

Technical Introduction

[Disclaimer]

Instead of an incomplete review, the general review is completely omitted... sorry!

Entropy Principle – review

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Very powerful strategy — Good example from Son-Surowka (2009) (The notation is slightly changed here)

Ideal Hydrodynamics ~ $\mathcal{O}(\partial^0)$

Entropy

$$s = \beta(e+p) - n\alpha$$
 $(\alpha = \beta\mu)$

Entropy Current

$$\mathcal{S}^{\mu}_{(0)} = su^{\mu} = \beta(e+p)u^{\mu} - n\alpha u^{\mu}$$
$$= \beta(u_{\nu}\Theta^{\mu\nu}_{(0)} + pu^{\mu} - \mu j^{\mu}_{(0)}) \quad \rightarrow \text{Generalized later!}$$

 $\partial_{\mu} \mathcal{S}^{\mu}_{(0)} = 0$ is easily concluded (ideal).

Entropy Principle — review Anomalous Hydrodynamics ~ O(d)

$$\mathcal{S}^{\mu} = \beta(u_{\nu}\Theta^{\mu\nu} + pu^{\mu} - \mu j^{\mu}) \qquad \partial_{\mu}j^{\mu} = C_{\text{anom}}E \cdot B$$
$$\partial_{\mu}\mathcal{S}^{\mu} = \Theta^{\mu\nu}_{(1)}\partial_{\mu}(\beta u_{\nu}) + j^{\mu}_{(1)}(-\partial_{\mu}\alpha + \beta E_{\mu}) - C_{\text{anom}}\alpha E \cdot B$$

How can we make this non-negative?

$$\Theta_{(1)}^{\mu\nu} = 2h^{(\mu}u^{\nu)} + \pi^{\mu\nu} \qquad u \cdot h = (u \cdot \pi)^{\mu} = 0$$

Heat Flow Shear Tensor (Here, we do not care about the frame choice.)

$$\begin{split} & Entropy Principle - review \\ & \Theta_{(1)}^{\mu\nu} = 2h^{(\mu}u^{\nu)} + \pi^{\mu\nu} \qquad u \cdot h = (u \cdot \pi)^{\mu} = 0 \\ & \Delta^{\mu\nu} = g^{\mu\nu} - u^{\mu}u^{\nu} \\ & 2h^{(\mu}u^{\nu)}\partial_{\mu}(\beta u_{\nu}) - j^{\mu}_{(1)}\partial_{\mu}\alpha = (h^{\mu} - \mathcal{H}j^{\mu}_{(1)})(\Delta_{\mu\nu}\partial^{\nu}\beta + \beta\dot{u}_{\mu}) \\ & \mathcal{H}^{-1}\begin{pmatrix} n \\ e + p \end{pmatrix} \Delta_{\mu\nu}\partial^{\nu}\alpha = \Delta_{\mu\nu}\partial^{\nu}\beta + \beta\dot{u}^{\mu} \qquad \beta u^{\nu}\partial_{\nu}u_{\mu} \\ & \pi^{\mu\nu}\partial_{\mu}(\beta u_{\nu}) = \beta\pi^{\mu\nu} \bigg[\partial_{<\mu}u_{\nu>} + \frac{1}{3}\Delta_{\mu\nu}(\partial \cdot u)\bigg] \\ & A^{<\mu\nu>} = \Delta^{(\mu\alpha}\Delta^{\beta\nu)}A_{\alpha\beta} - \frac{1}{3}\Delta^{\mu\nu}\Delta^{\alpha\beta}A_{\alpha\beta} \\ & \text{September 17, 2024 @ ect*, Trento} \end{split}$$

$$\begin{split} &Entropy\ Principle\ -\ review\\ &\partial_{\mu}\mathscr{S}^{\mu} = \Theta_{(1)}^{\mu\nu}\partial_{\mu}(\beta u_{\nu}) + j_{(1)}^{\mu}(-\partial_{\mu}\alpha + \beta E_{\mu}) - C_{\text{anom}}\alpha E \cdot B\\ &= (h^{\mu} - \mathscr{H}j_{(1)}^{\mu})(\Delta_{\mu\nu}\partial^{\nu}\beta + \beta \dot{u}_{\mu})\\ &+ \beta \pi^{\mu\nu} \left[\partial_{<\mu}u_{\nu>} + \frac{1}{3}\Delta_{\mu\nu}(\partial \cdot u)\right] - C_{\text{anom}}E \cdot B \end{split}$$

The last unwanted term is "renormalized" in the currents.

$$\delta j_{(1)}^{\mu} = \xi_V \omega^{\mu} + \xi_B B^{\mu} \qquad \delta S^{\mu} = -\alpha \delta j_{(1)}^{\mu} + D_V \omega^{\mu} + D_B B^{\mu}$$
$$\longleftrightarrow \quad \xi_V = C_{\text{anom}} \left(\mu^2 - \frac{2}{3} \frac{\mu^3}{\mathcal{H}} \right) \qquad \xi_B = C_{\text{anom}} \left(\mu - \frac{1}{2} \frac{\mu^2}{\mathcal{H}} \right)$$
Son-Surowka (2009)

Spin Hydro Introduction

Spin Hydro – review ಜನ್ ಸಚೆದನ್ ಸಚೆದನ್ ಸಚೆದನ್ ಸಚೆದನ್ ಸಚೆದ ಸಚೆದನ್ ಸಚೆದನ್ ಸಚೆದನ್ ಸಚೆದನ್ Spin Hydrodynamics (Hattori-Hongo-Huang-Matsuo-Taya 2019) Noether current from rotational symmetry $J^{\lambda\mu\nu} = x^{\mu}T^{\lambda\nu} - x^{\nu}T^{\lambda\nu} + \Sigma^{\lambda\mu\nu}$ **Orbital Spin** $\partial_{\mu}T^{\mu\nu} = \partial_{\lambda}J^{\lambda\mu\nu} = 0 \longrightarrow \partial_{\lambda}\Sigma^{\lambda\mu\nu} = -2\Theta^{[\mu\nu]}$

Anti-symmetric part of the energy-momentum tensor is the source of the spin (from the orbital part).

Spin Hydro – review

Spin Hydrodynamics (Hattori-Hongo-Huang-Matsuo-Taya 2019)

$$S^{\mu} = \beta(u_{\nu}\Theta^{\mu\nu} + pu^{\mu} - \mu j^{\mu} - \omega_{\rho\sigma}S^{\rho\sigma}u^{\mu})$$

$$\Sigma^{\lambda\mu\nu} = u^{\lambda}S^{\mu\nu} + \Sigma^{\lambda\mu\nu}_{(1)}$$
Induced by the spin potential
$$\Theta^{\mu\nu} = \Theta^{(\mu\nu)} + 2q^{[\mu}u^{\nu]} + \phi^{\mu\nu}$$
 $u \cdot q = (u \cdot \phi)^{\mu} = 0$
Anti-symmetric part

$$\partial_{\mu}\mathcal{S}^{\mu} = \cdots (2q^{[\mu}u^{\nu]} + \phi^{\mu\nu})\partial_{\mu}(\beta u_{\nu}) - \partial_{\mu}(u^{\mu}S^{\rho\sigma})\beta\omega_{\rho\sigma}$$

Spin Hydro — review Spin Hydrodynamics (Hattori-Hongo-Huang-Matsuo-Taya 2019) $\partial_{\mu}S^{\mu} = \dots + (2q^{[\mu}u^{\nu]} + \phi^{\mu\nu})\partial_{\mu}(\beta u_{\nu}) - \partial_{\mu}(u^{\mu}S^{\rho\sigma})\beta\omega_{\rho\sigma}$

$$= \dots + q^{\mu} (\partial_{\mu} \beta - \beta \dot{u}_{\mu} + 4\beta \omega_{\mu\nu} u^{\nu}) \qquad 2\Theta^{[\rho\sigma]} + \beta \phi^{\mu\nu} (2\omega_{\mu\nu} + \partial_{\mu} u_{\nu})$$

$$q^{\mu} = \lambda \Delta^{\mu\nu} (\partial_{\nu}\beta - \beta \dot{u}_{\nu} + 4\beta \omega_{\nu\rho} u^{\rho})$$

$$\phi^{\mu\nu} = \gamma \Delta^{\mu\rho} \Delta^{\nu\sigma} (2\omega_{\rho\sigma} + \partial_{[\rho} u_{\sigma]})$$

Same conclusion in the Belinfante symmetric form Fukushima-Pu (2020)

Spin Hydro Revisited

Spin Hydro – puzzle

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$$\partial_{\lambda} \Sigma^{\lambda\mu\nu} = -2\Theta^{[\mu\nu]}$$
$$\Sigma^{\lambda\mu\nu} = u^{\lambda} S^{\mu\nu} + \Sigma^{\lambda\mu\nu}_{(1)}$$

$$\Theta^{\mu\nu} = \Theta^{(\mu\nu)} + 2q^{[\mu}u^{\nu]} + \phi^{\mu\nu}$$

$$q^{\mu} = -\frac{1}{2} u_{\nu} \partial_{\lambda} \Sigma^{\lambda \mu \nu}$$

This current is already solved in terms of spin!?

In the whole formulation, the d.o.f. are redundant. Hongo-Huang-Kaminski-Stephanov-Yee (2021)

Spin Hydro – puzzle

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Should we really keep 6 generators? — No!

3 rotations + 3 boosts

Hermitian non-Hermitian

Lorentz group has no hermitian representation with finite dimensions.

Fully anti-symmetrized $\Sigma^{\lambda\mu\nu}$ should correspond to the physical observables, i.e., $\Sigma^{00i} = 0$

$$\Sigma^{\lambda\mu\nu} = u^{\lambda}S^{\mu\nu} + u^{\mu}S^{\nu\lambda} + u^{\nu}S^{\lambda\mu} + \Sigma^{\lambda\mu\nu}_{(1)}$$

Spin Hydro — challenge Shuo Fang-KF-SP-Dong-Lin Wang (soon!) Can we really construct the hydro? — Yes! (Discussions did not converge yet.) Differences:

 q^{μ} is already solved as $q^{\mu} = -\frac{1}{2}u_{\nu}\partial_{\lambda}\Sigma^{\lambda\mu\nu} = \frac{1}{2}(S^{\mu\nu}\dot{u}_{\nu} + \Delta^{\mu}_{\nu}\partial_{\lambda}S^{\nu\lambda})$

Some extra terms from

Simplification

$$-\partial_{\mu}(u^{\mu}S^{\rho\sigma})\beta\omega_{\rho\sigma} \neq 2\Theta^{[\rho\sigma]}\beta\omega_{\rho\sigma}$$

 $(u \cdot S)^{\mu} = 0$ $(u \cdot \omega)^{\mu} = 0$

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Spin Hydro — challenge Shuo Fang-KF-SP-Dong-Lin Wang (soon!)





Maybe..., these do not appear, if the entropy current is chosen better.

Spin Hydro — challenge Shuo Fang-KF-SP-Dong-Lin Wang (soon!)

$$\begin{split} &\frac{1}{2}(S^{\mu\nu}\dot{u}_{\nu} + \Delta^{\mu}_{\nu}\partial_{\lambda}S^{\nu\lambda})(\partial_{\mu}\beta \bigoplus \beta\dot{u}_{\mu}) \\ &= -\frac{1}{2}(S^{\mu\nu}\dot{u}_{\nu} + \Delta^{\mu}_{\nu}\partial_{\lambda}S^{\nu\lambda})(\partial_{\mu}\beta \bigoplus \beta\dot{u}_{\mu}) \\ &+ (S^{\mu\nu}\dot{u}_{\nu} + \Delta^{\mu}_{\nu}\partial_{\lambda}S^{\nu\lambda})\partial_{\mu}\beta \\ &= -\frac{1}{2}(\Delta^{\mu}_{\nu}\partial_{\lambda}S^{\nu\lambda} + \beta^{-1}S^{\mu\nu}\partial_{\nu}\beta)(\partial_{\mu}\beta + \beta\dot{u}_{\mu}) \\ &+ \Delta^{\mu}_{\nu}\partial_{\lambda}S^{\nu\lambda}\partial_{\mu}\beta \end{split} \qquad \begin{array}{l} \text{Renormalized in} \\ \text{the heat current.} \end{array}$$



Spin Hydro – results ik. Mengik. Mengik. Mengik. Mengik. Men Mengik. Mengik. Mengik. Mengik. Mengik. Mengik. **Shuo Fang-KF-SP-Dong-Lin Wang (soon!)** $\delta h^{\mu} = -\frac{1}{2} \Delta^{\mu}_{\nu} \partial_{\lambda} S^{\nu\lambda} - \frac{1}{2} \beta^{-1} S^{\mu\nu} \partial_{\nu} \beta$ $$\begin{split} \delta \pi^{\mu\nu} &= S^{\lambda < \rho} \partial_{\lambda} u^{\sigma >} \\ \delta \Pi &= \frac{1}{3} S^{\mu\nu} \partial_{\mu} u_{\nu} - 2 \left(\frac{\partial p}{\partial e} \right)_{n, S^{\mu\nu}} S^{\mu\nu} \omega_{\mu\nu} \end{split}$$ $\phi^{\mu\nu} = \gamma \Delta^{\mu\rho} \Delta^{\nu\sigma} (2\beta\omega_{\rho\sigma} + \partial_{\rho}u_{\sigma}) + \beta^{-1}\dot{\beta}S^{\mu\nu}$

We are still checking the calculations... Coefficients may be wrong, but the strategy is correct.

Spin Hydro — puzzle again Shuo Fang-KF-SP-Dong-Lin Wang (soon!)

Tensor decomposition and renormalization seem to be not unique...???

The present demonstration is just one simple example... we found a family of non-negative decomposition of tensors.

Any principle still overlooked???

Summary

- Spin hydrodynamics in the canonical (and Belinfante) form have some redundancy.
 - \square Spin has not 6 but 3 charge observables.
 - □ Symmetrized form should be physical.
- Extra terms can be renormalized so that the second law of thermodyanmics is satisfied.
 - \square Overconstrained problem is resolved.
 - **Tensor decomposition seems to be non-unique.**
 - □ A family of decompositions with unfixed parameters was found to satisfy the entropy principle... more physical conditions?

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Pseudo-Gauge Sym. in Energy-Momentum Tensor Fukushima-Pu (2020)

 $\mathcal{T}^{\mu\nu} = \Theta^{\mu\nu} + \partial_{\lambda} K^{\lambda\mu\nu} \quad \text{(conserved current redefined)}$ $K^{\lambda\mu\nu} = \frac{1}{2} \left(\Sigma^{\lambda\mu\nu} - \Sigma^{\mu\lambda\nu} + \Sigma^{\nu\mu\lambda} \right)$

$$\mathcal{T}^{\mu\nu} = \Theta^{\mu\nu} + \frac{1}{2} \partial_{\lambda} (u^{\lambda} S^{\mu\nu} - u^{\mu} S^{\lambda\nu} + u^{\nu} S^{\mu\lambda})$$
$$= \Theta^{\mu\nu}_{(s)} + \frac{1}{2} \left[\partial_{\lambda} (u^{\mu} S^{\nu\lambda} + u^{\nu} S^{\mu\lambda}) \right]$$

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Pseudo-Gauge Sym. in Energy-Momentum Tensor Fukushima-Pu (2020)

$$\mathcal{T}^{\mu\nu} = \Theta^{\mu\nu} + \frac{1}{2} \partial_{\lambda} (u^{\lambda} S^{\mu\nu} - u^{\mu} S^{\lambda\nu} + u^{\nu} S^{\mu\lambda})$$
$$= \Theta^{\mu\nu}_{(s)} + \frac{1}{2} \left[\partial_{\lambda} (u^{\mu} S^{\nu\lambda} + u^{\nu} S^{\mu\lambda}) \right]$$

Spin induced terms are "renormalized" as

$$2h^{(\mu}u^{\nu)} + \pi^{\mu\nu} + \frac{1}{2} \left[\partial_{\lambda} (u^{\mu}S^{\nu\lambda} + u^{\nu}S^{\mu\lambda}) \right] \\ = \delta e u^{\mu}u^{\nu} + 2 \left(h^{(\mu} + \delta h^{(\mu)})u^{\nu)} + \pi^{\mu\nu} + \delta \pi^{\mu\nu} \right)$$

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Pseudo-Gauge Sym. in Energy-Momentum Tensor Fukushima-Pu (2020)

$$\delta e = u_{\rho} \partial_{\sigma} S^{\rho\sigma}$$

$$\delta h^{\mu} = \frac{1}{2} \Big[\Delta^{\mu}_{\sigma} \partial_{\lambda} S^{\sigma\lambda} + u_{\rho} S^{\rho\lambda} \partial_{\lambda} u^{\mu} \Big]$$

$$\delta \pi^{\mu\nu} = \partial_{\lambda} (u^{<\mu} S^{\nu>\lambda}) + \delta \Pi \Delta^{\mu\nu}$$

$$\delta \Pi = \frac{1}{3} \partial_{\lambda} (u^{\sigma} S^{\rho\lambda}) \Delta_{\rho\sigma}$$

An electric current $j \propto \nabla \times S$ is implied... Spin Vorticity Effect

Pseudo-Gauge Sym. in Energy-Momentum Tensor Fukushima-Pu (2020)

$$\partial_{\mu}S^{\mu} = \dots + \frac{1}{2} \Big[\partial_{\lambda}(u^{\mu}S^{\nu\lambda} + u^{\nu}S^{\mu\lambda}) \Big] \partial_{\mu}\frac{u_{\nu}}{T} - \frac{\omega_{\rho\sigma}}{T} \partial_{\lambda}(u^{\lambda}S^{\rho\sigma}) \\ = \frac{1}{2} \partial_{\mu} \Big[\partial_{\lambda}(u^{\lambda}S^{\mu\nu} + u^{\mu}S^{\nu\lambda} + u^{\nu}S^{\mu\lambda}) \frac{u_{\nu}}{T} \Big] \\ \hline \mathbf{Total \ derivative} \\ - \frac{1}{2} \Big[\partial_{\lambda}(u^{\lambda}S^{\mu\nu}) \Big] \partial_{\mu}\frac{u_{\nu}}{T} - \frac{\omega_{\rho\sigma}}{T} \partial_{\lambda}(u^{\lambda}S^{\rho\sigma}) \\ \mathbf{Absorbed \ in \ the \ entropy, \\ \mathbf{then \ it \ is \ just \ canonical!}} \\ \mathbf{Canonical \ results} \\ \end{bmatrix}$$