



# FRIB

## Nuclear structure-based optical potentials for the era of rare isotope beams

Grigor Sargsyan

FRIB Theory Fellow

Michigan State University  
[sargsyan@frib.msu.edu](mailto:sargsyan@frib.msu.edu)

Towards a consistent approach for nuclear structure and reactions:  
microscopic optical potentials  
ECT\*, Trento, Italy  
19 June 2024



MICHIGAN STATE  
UNIVERSITY

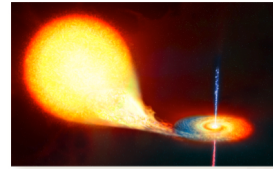
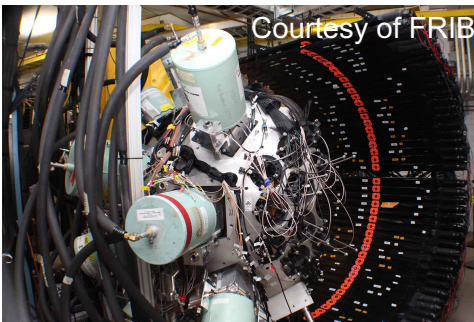
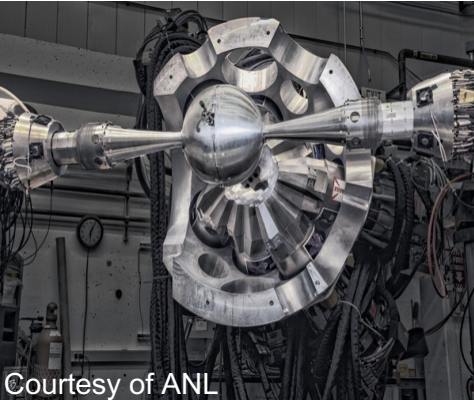


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Science

# The era of rare isotope beams

Experiments at RIB facilities

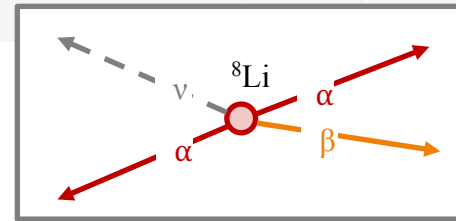
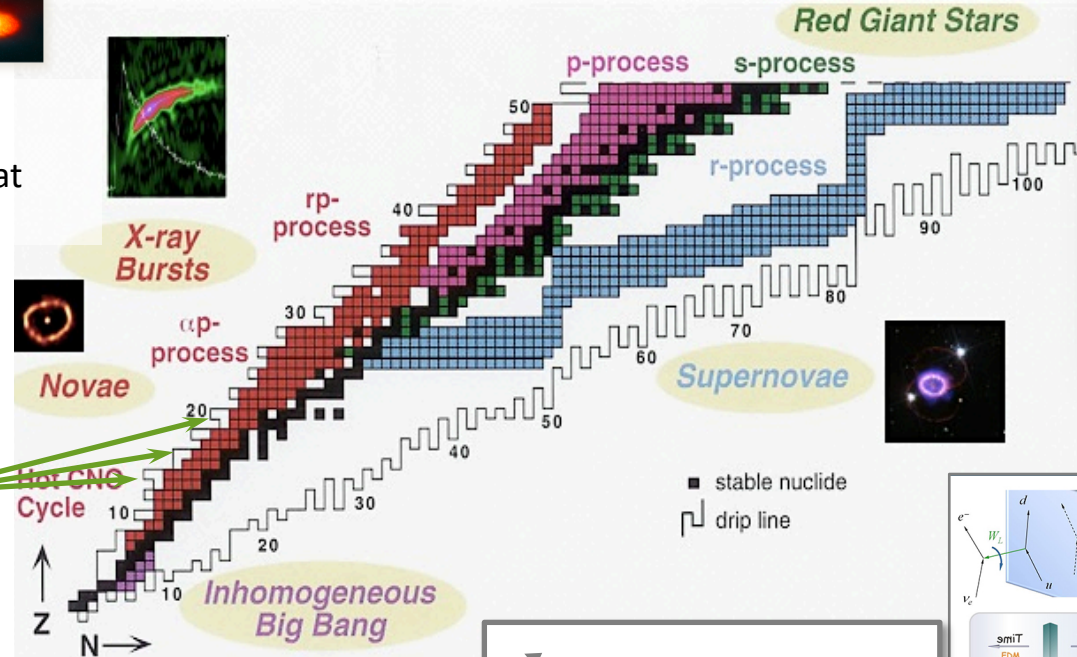


X-ray burst nucleosynthesis

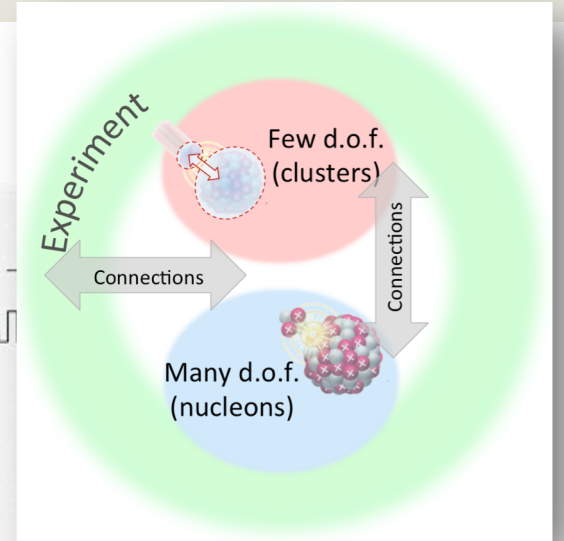
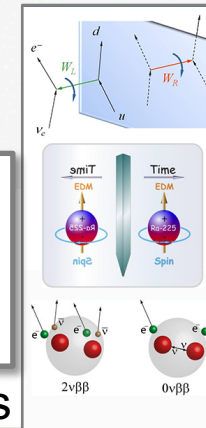
Alpha widths

Reaction rates at low energies

Close proximity of the drip lines



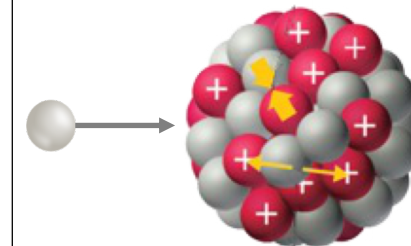
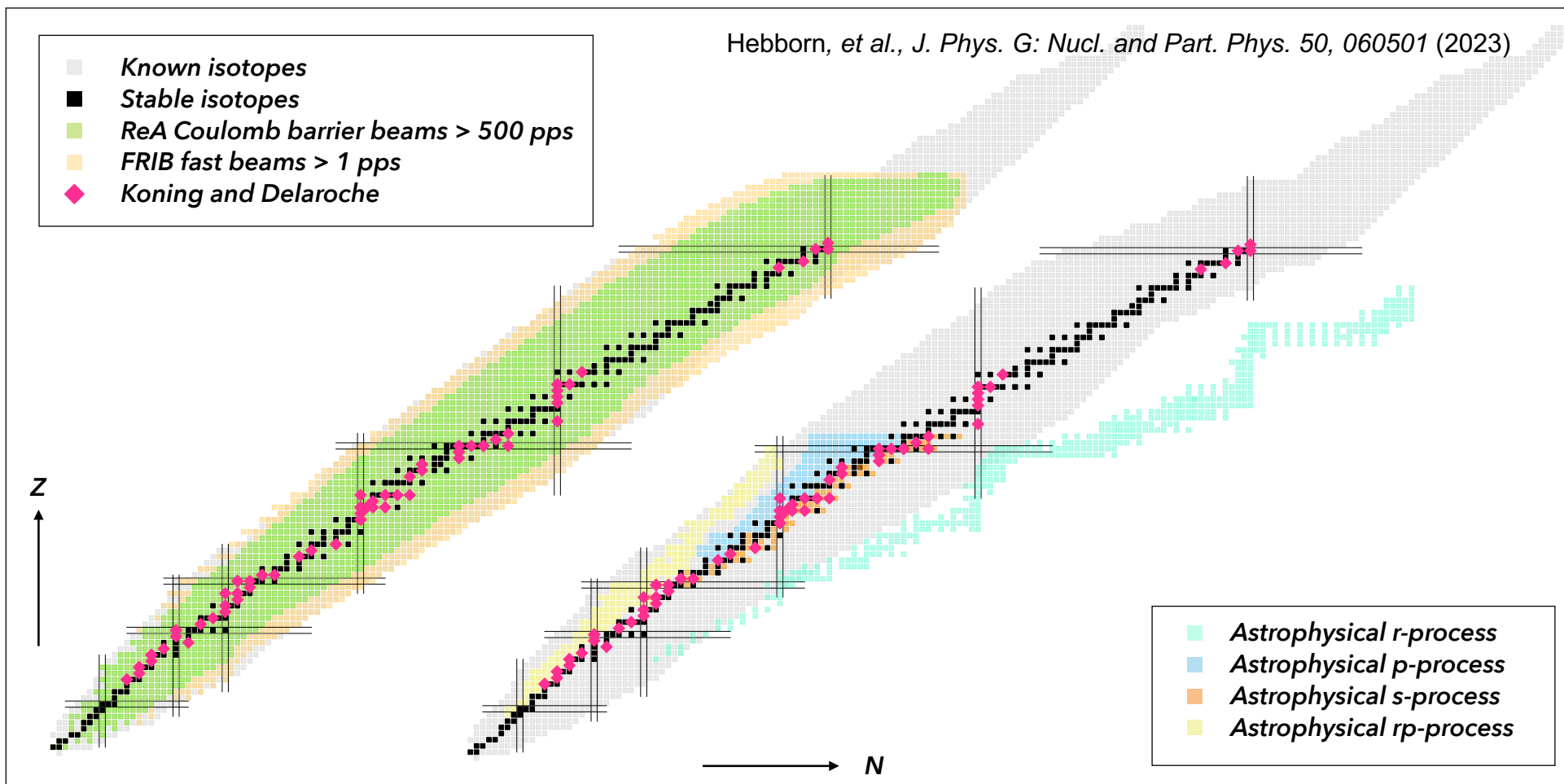
Fundamental symmetries



FRIB Theory Alliance topical program (2018)  
 “From bound states to the continuum”  
 (C. Johnson et al.)

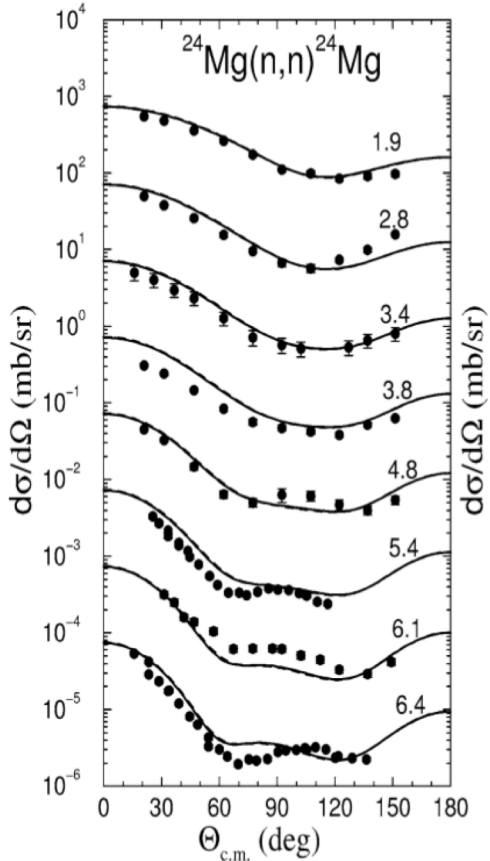


# Our current knowledge of optical potentials (OP) is very limited



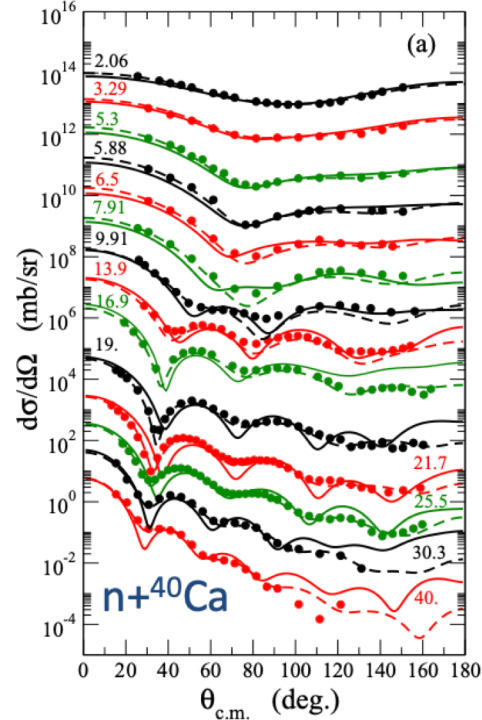
# Different strategies for calculating the nucleon-nucleus optical potential (OP)

phenomenological fit



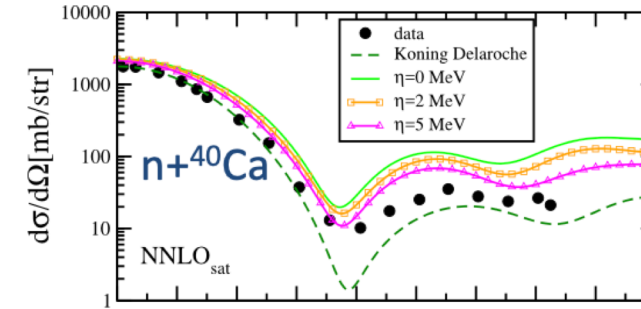
Koning, Delaroche NPA **713** (2003) 231

RPA calculation with added imaginary part



Blanchon *et al.* PRC **91** 014612 (2015)

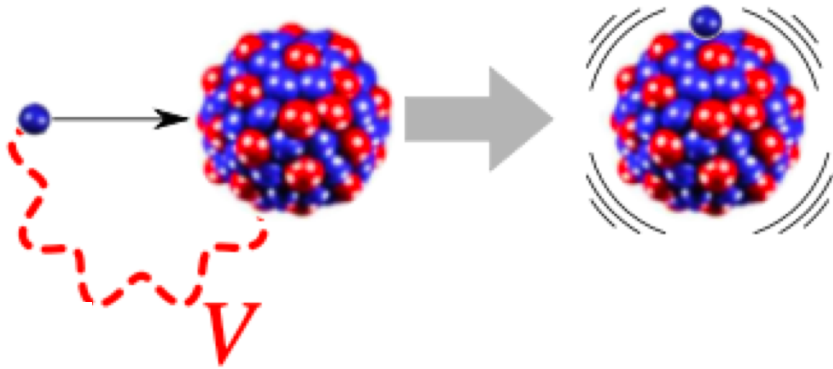
coupled-cluster ab initio with non-zero  $\eta$  parameter



Rotureau *et al.* PRC **98** 044625 (2018)

- Phenomenological fits are widely used, but are disconnected from the structure, and extrapolation away from stability is risky
- Microscopic theories often struggle to get absorption right
- Ab initio approaches are mainly feasible for light or near closed-shell nuclei

# Embedding nuclear structure information within OP



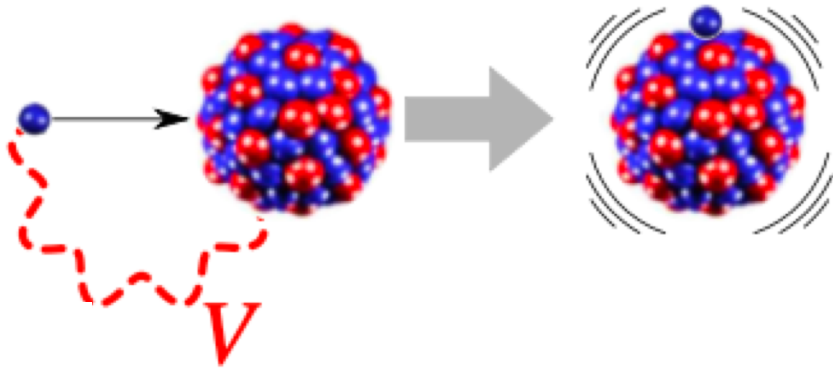
Feshbach formalism

$$\begin{aligned} V(\mathbf{r}, \mathbf{r}', E) &= U_0(\mathbf{r}) + V_{PO}(\mathbf{r}, \mathbf{r}', E - E_i) \\ &= U_0(\mathbf{r}) + \sum_i U_{0i}(\mathbf{r}) G_i(\mathbf{r}, \mathbf{r}', E - E_i) U_{0i}(\mathbf{r}') \end{aligned}$$

Static, energy-independent potential

Polarization potential: Requires input from nuclear structure

# Embedding nuclear structure information within OP



Feshbach formalism

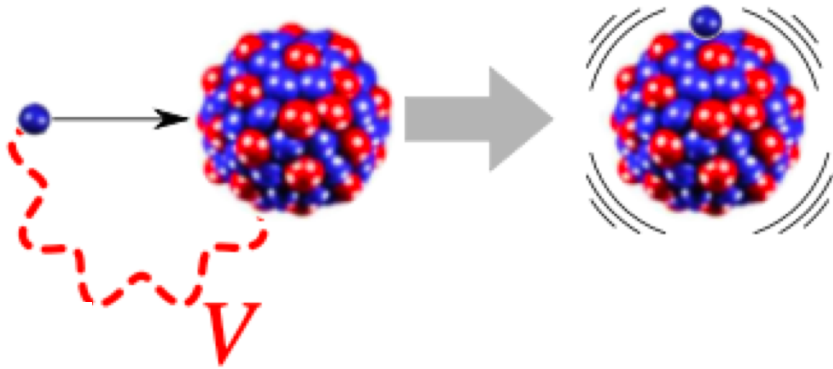
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Static,  
energy-independent  
potential

Polarization potential:  
Requires input from  
nuclear structure

Can be applied to any mass range  
as long as nuclear structure  
calculations are available

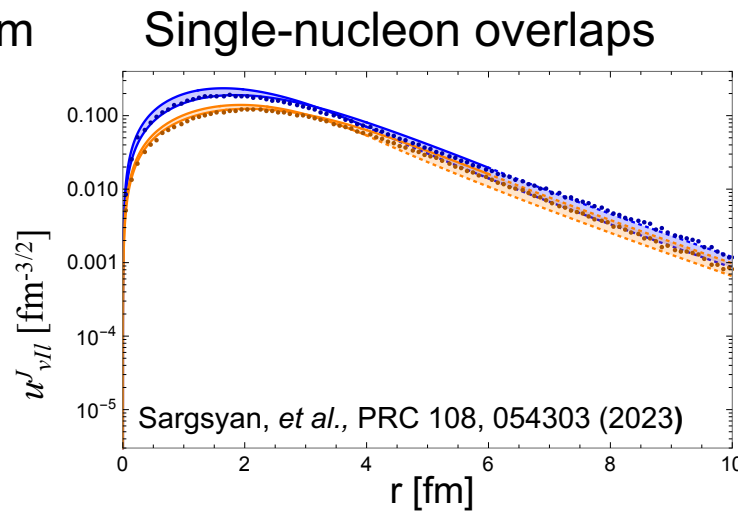
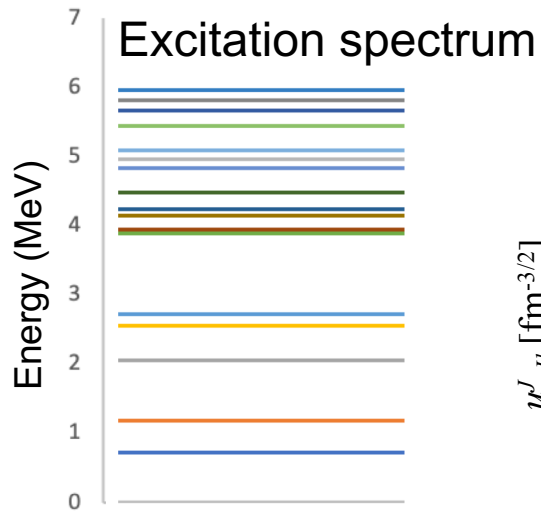
# Embedding nuclear structure information within OP



Feshbach formalism

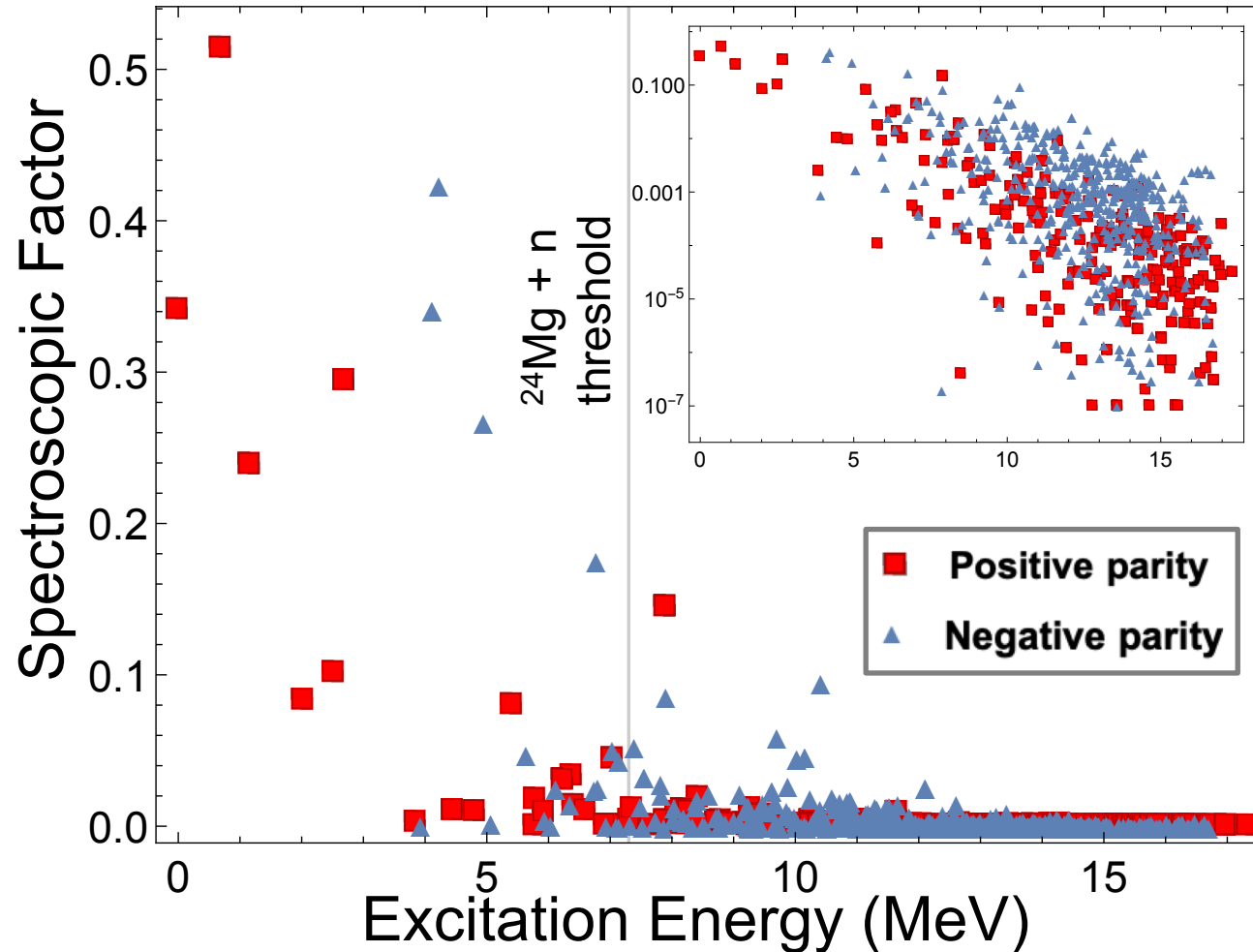
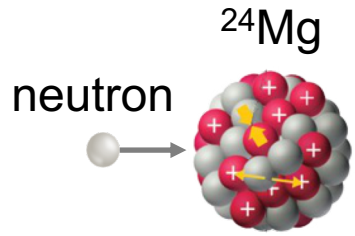
$$\begin{aligned}
 V(\mathbf{r}, \mathbf{r}', E) &= U_0(\mathbf{r}) + V_{PO}(\mathbf{r}, \mathbf{r}', E - E_i) \\
 &= U_0(\mathbf{r}) + \underbrace{\sum_i U_{0i}(\mathbf{r}) G_i(\mathbf{r}, \mathbf{r}', E - E_i) U_{0i}(\mathbf{r}')}_{\text{Polarization potential}}
 \end{aligned}$$

Polarization potential:  
Requires input from  
nuclear structure



e.g., shell model, RPA,  
ab initio models, ...

# 1<sup>st</sup> ingredient for constructing OP: shell model input



Around 600  
intrinsic states

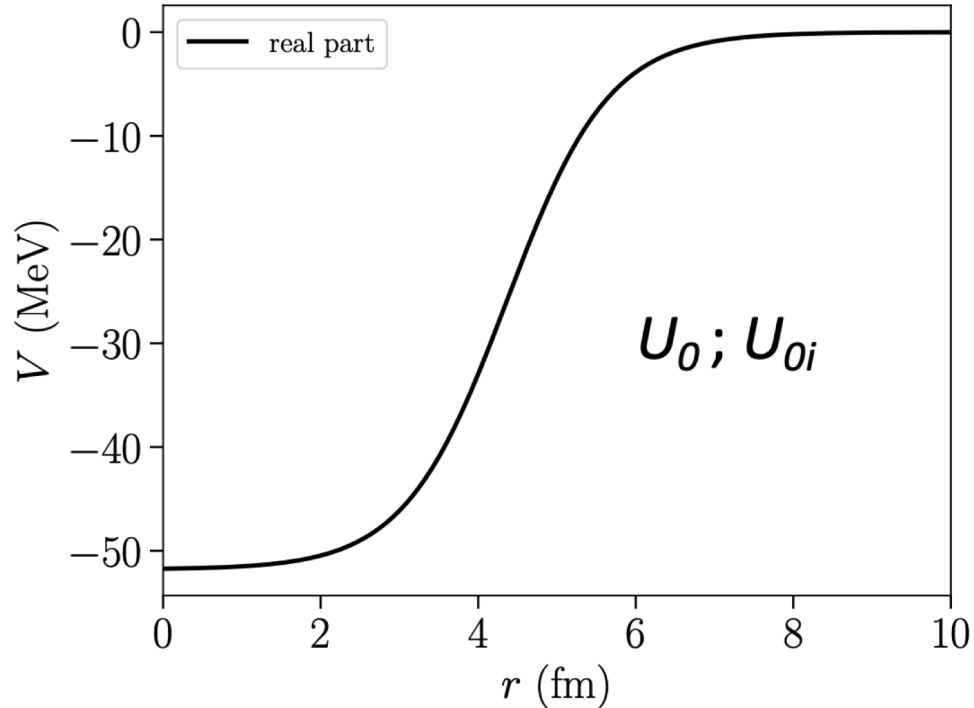
Shell model calculations with  
PSDPF potential M Bouhelal, *et al.*, Nucl. Phys. A 864 (2011)





# 2<sup>nd</sup> ingredient: static potential and couplings

$$V(\mathbf{r}, \mathbf{r}', E) = U_0(\mathbf{r}) + \sum_i U_{0i}(\mathbf{r}) G_i(\mathbf{r}, \mathbf{r}', E - E_i) U_{0i}(\mathbf{r}')$$



- static potential  $U_0$ : real, local Woods-Saxon adjusted to reproduce binding energy of  $^{25}\text{Mg}$
- couplings  $U_{0i}$ : same real Woods-Saxon, but adjusted to each  $E_i$  and multiplied by spectroscopic factor  $S_i$  from shell model



# 3<sup>rd</sup> ingredient: iterative scheme for self consistent OP

$$V(\mathbf{r}, \mathbf{r}', E) = U_0(\mathbf{r}) + \sum_i U_{0i}(\mathbf{r}) G_i(\mathbf{r}, \mathbf{r}', E - E_i) U_{0i}(\mathbf{r}')$$

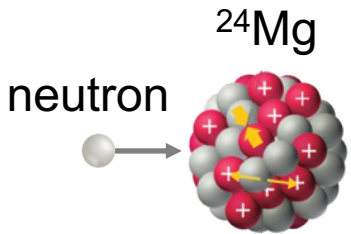
$$G(\mathbf{r}, \mathbf{r}', E) = [E - T - V(\mathbf{r}, \mathbf{r}', E)]^{-1}$$

➤ Start with  $U_0$  and obtain  $V^{(1)}$

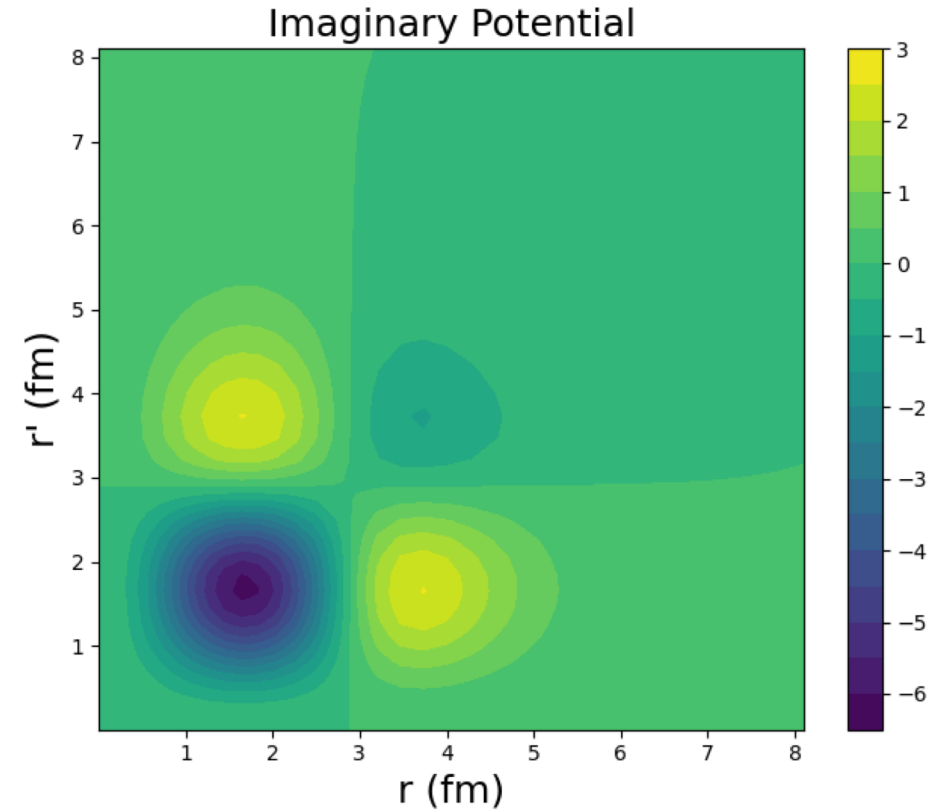
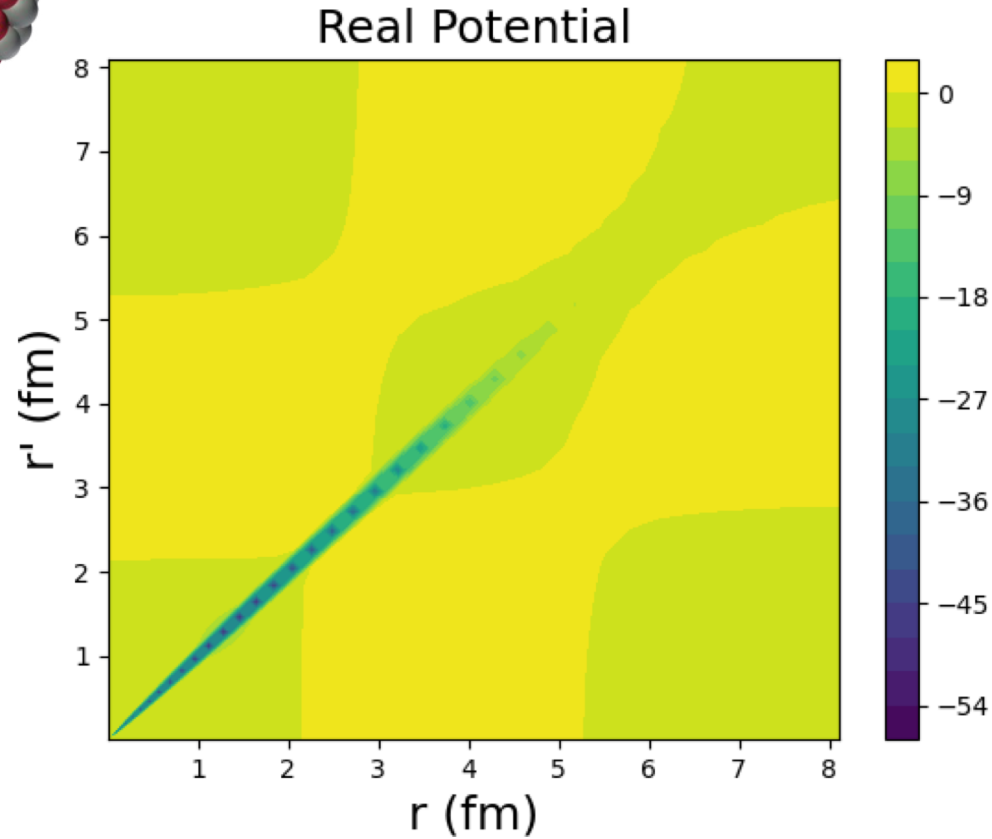
➤ Plug back in  $V^{(1)}$  and obtain  $V^{(2)}$

➤ Repeat until the volume integral converges  $J^{(n)} = \int \mathcal{V}^{(n)}(\mathbf{r}, \mathbf{r}') d\mathbf{r} d\mathbf{r}'$ ,

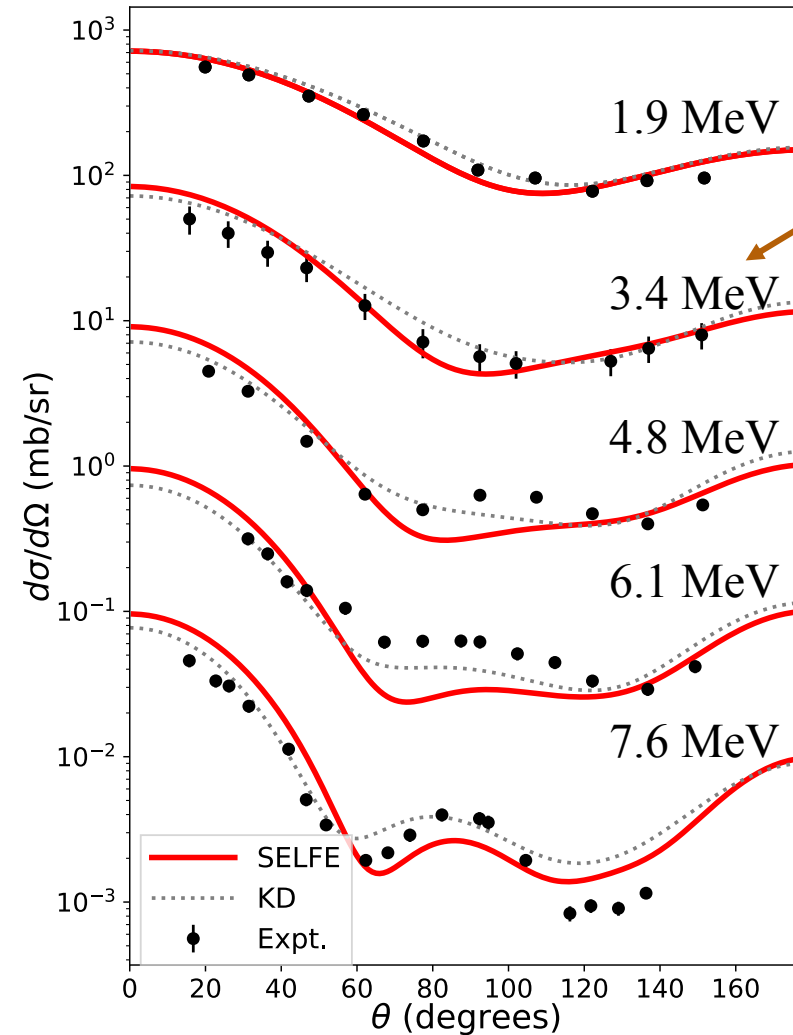
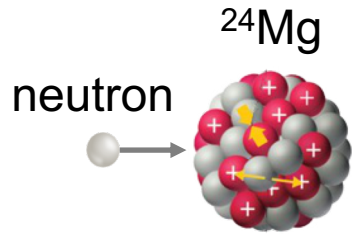
# OP is complex, energy-dependent, dispersive, and non-local



3.4 MeV projectile energy



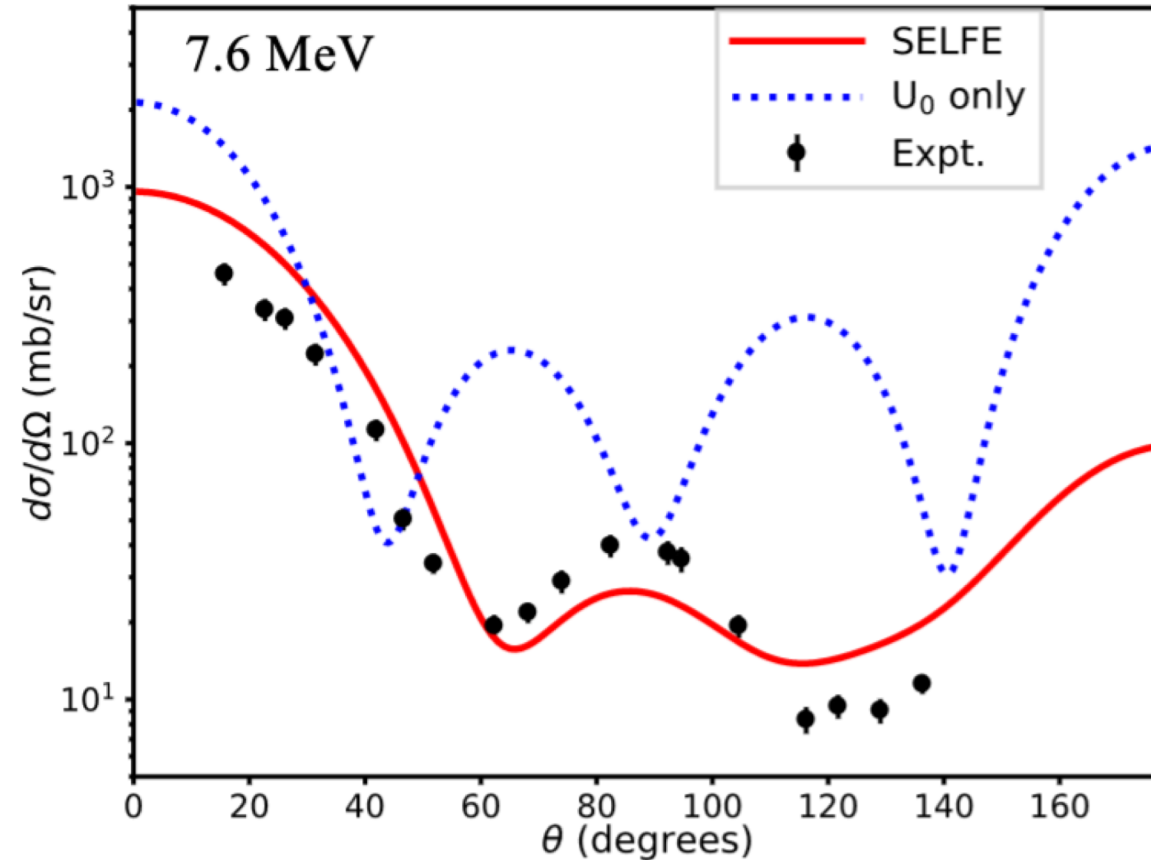
# Accurate prediction without parameters fitted to experimental scattering data!



Projectile energy

Sargsyan, *et al.*, in preparation

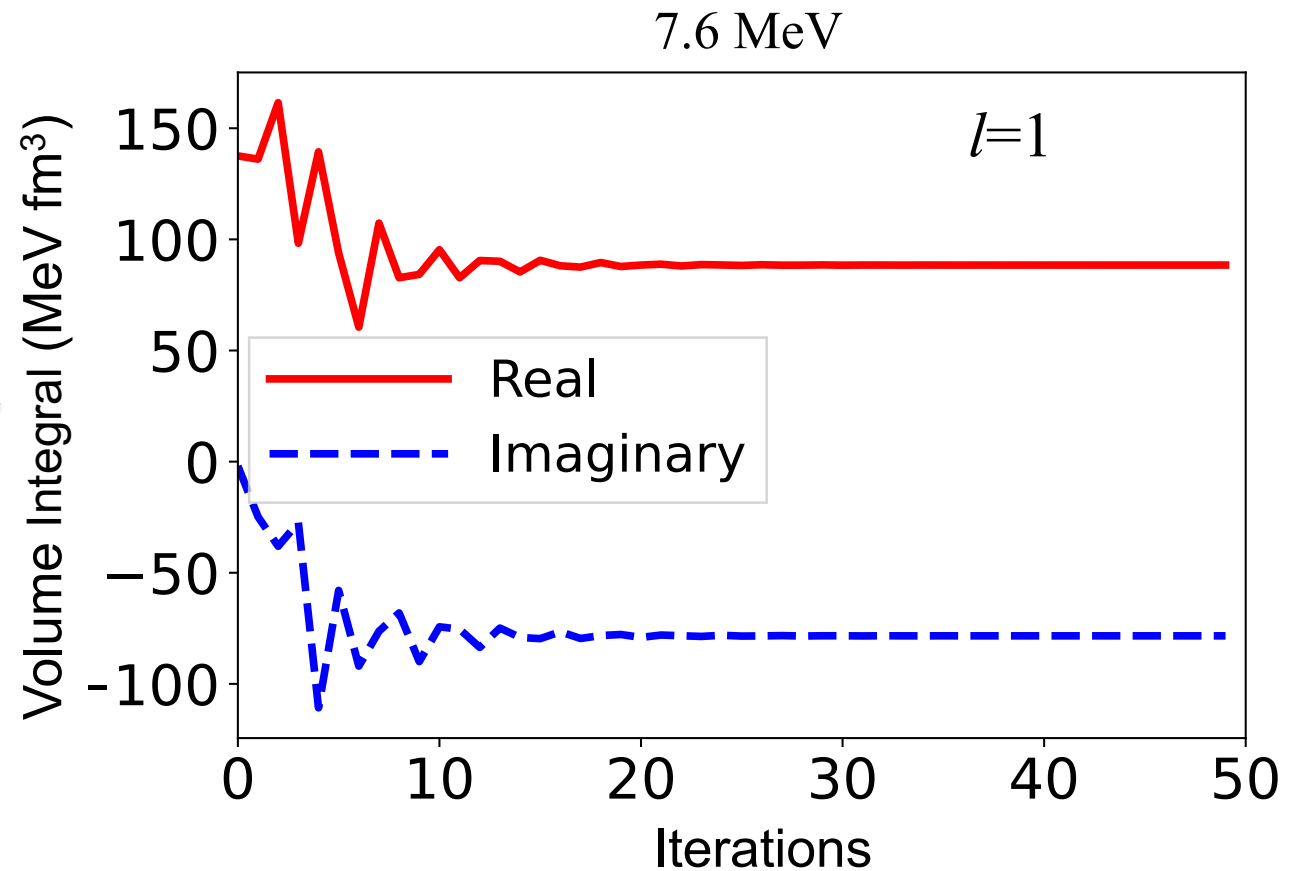
# No phenomenological imaginary terms



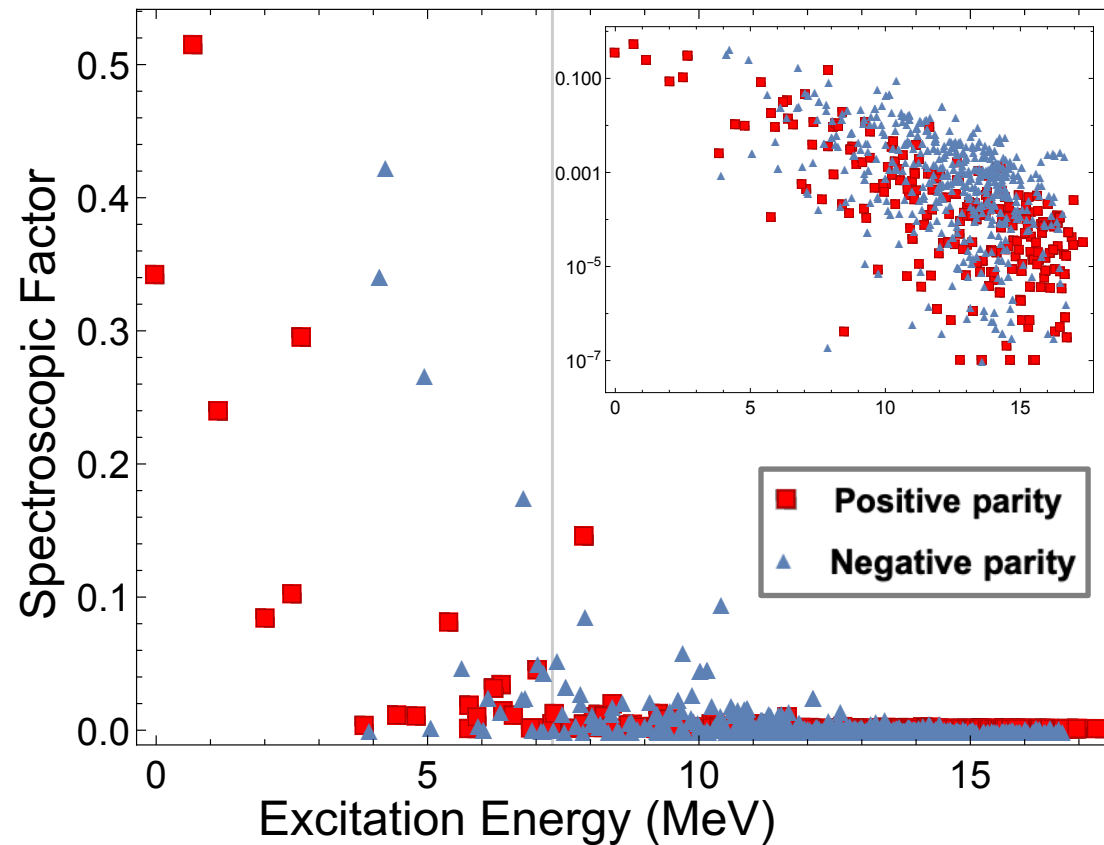
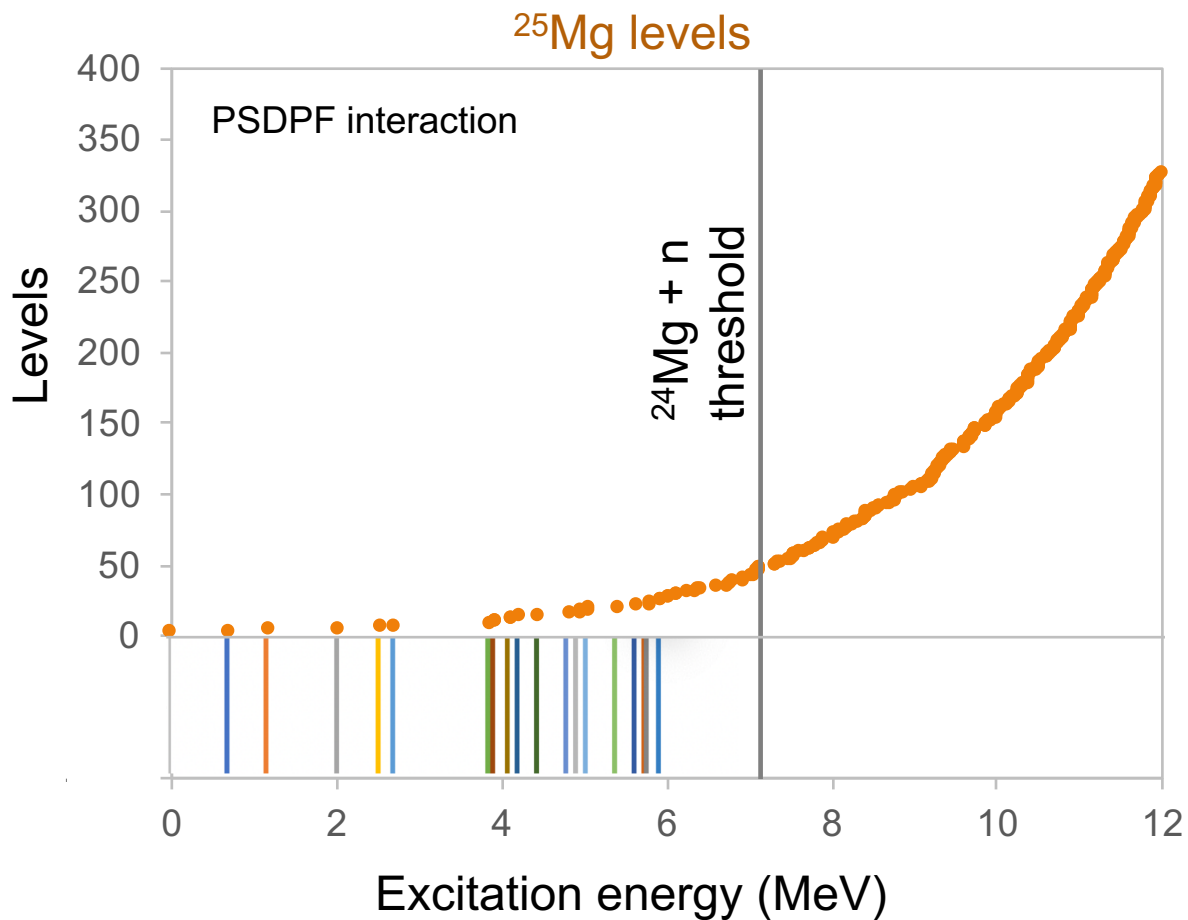
$$V(\mathbf{r}, \mathbf{r}', E) = U_0(\mathbf{r}, E) + V_{PO}(\mathbf{r}, \mathbf{r}', E) = U_0(\mathbf{r}, E) + \sum_i U_{0i}(\mathbf{r}) G_i(\mathbf{r}, \mathbf{r}', E) U_{0i}(\mathbf{r}')$$

# Potential volume integral convergence

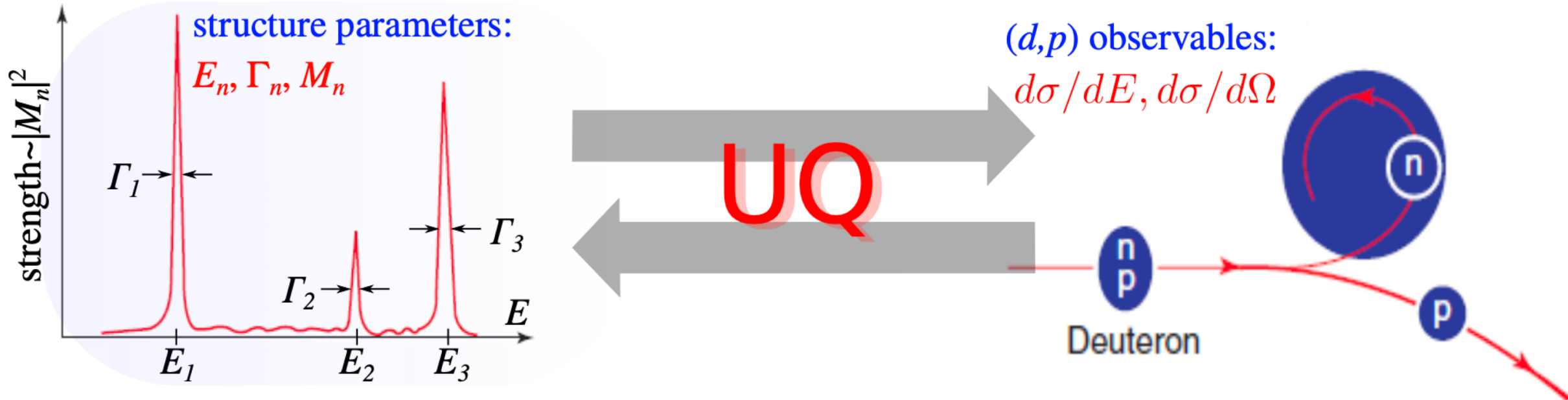
$$J^{(n)} = \int \mathcal{V}^{(n)}(\mathbf{r}, \mathbf{r}') d\mathbf{r} d\mathbf{r}',$$



# Ingredients for constructing neutron+<sup>24</sup>Mg OP



# Quantify uncertainties in the structure parameters that define OP

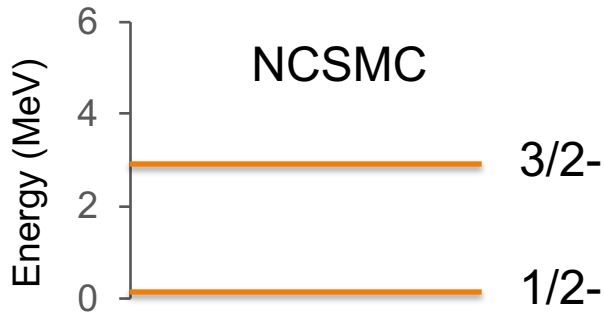


Can also be used to constrain the underlying chiral forces if we use ab initio inputs

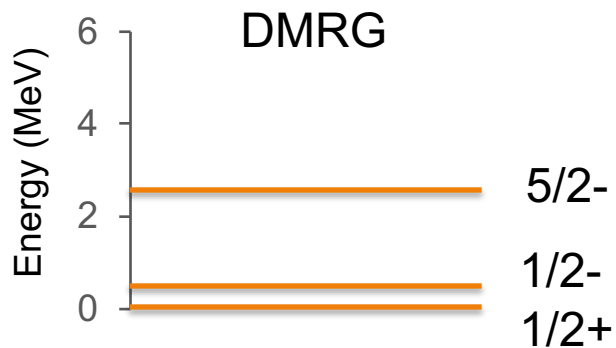


# Elusive ground state of ${}^9\text{He}$

${}^9\text{He}$  -- Nucleus with the highest neutron to proton ratio!

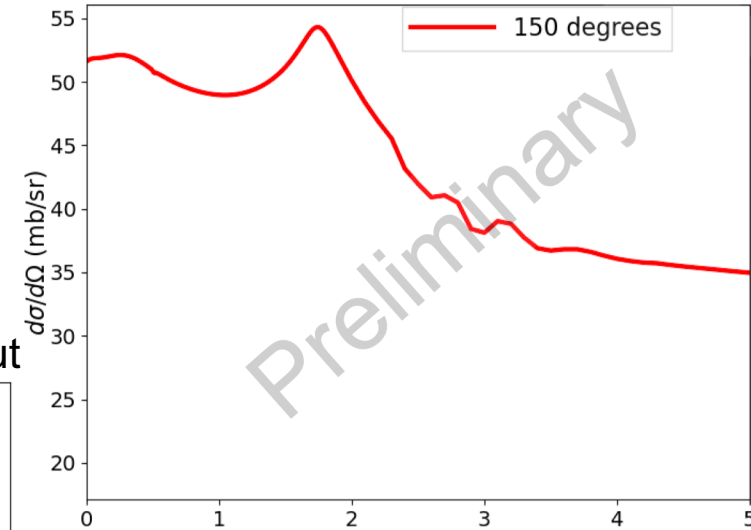


Vorabbi *et al.*, PRC **97**, 034314 (2018)



Fossez, *et al.*, PRC **98**, 061302(R) (2018)

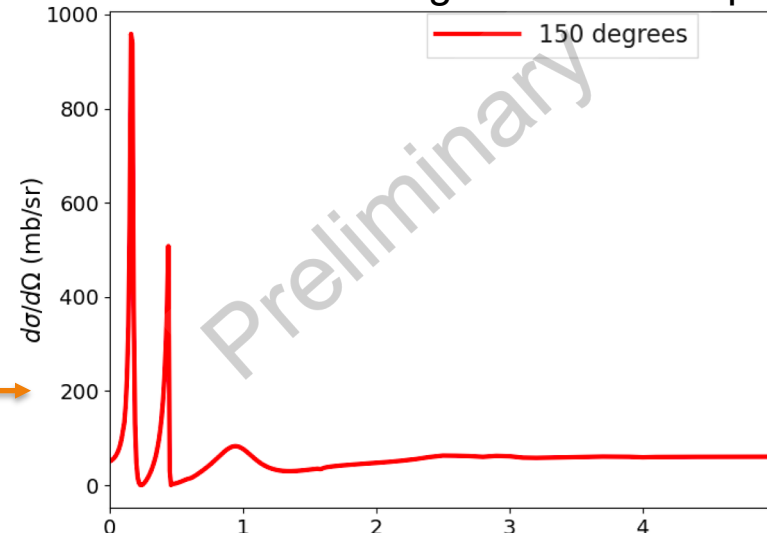
$n+{}^8\text{He}$  elastic scattering with NCSMC input



Neutron Energy (MeV)

Calculate (d,p) observables with and without  $1/2+$  state and propose an experiment at FRIB to measure  ${}^8\text{He}(d,p)$

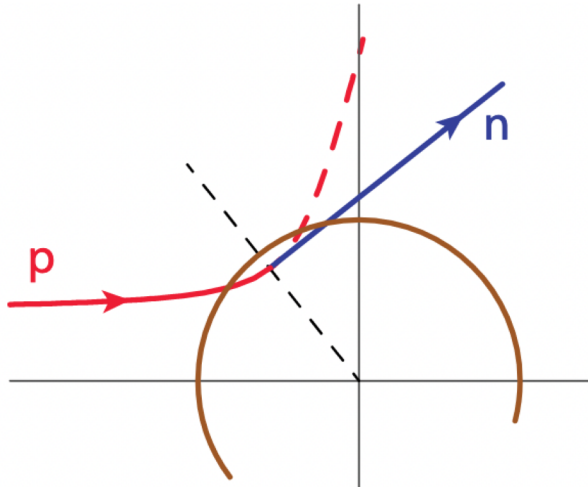
$n+{}^8\text{He}$  elastic scattering with DMRG input



Neutron Energy (MeV)

See Gregory's talk tomorrow for GF transfer!

# Reactions for the studies of charge-exchange processes



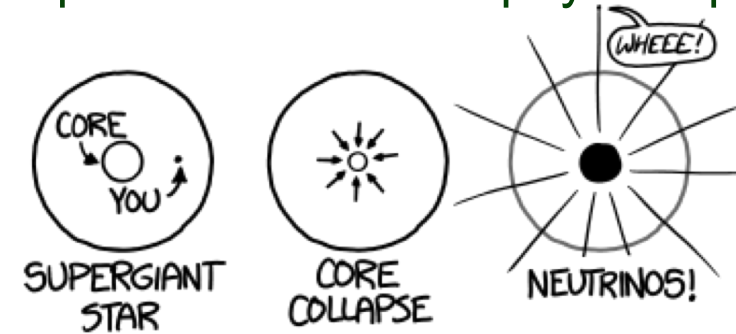
Danielewicz *et al.*, Nuc. Phys. A, 2017

- Charge-exchange reactions have been used to constrain (double)  $\beta$ -decay rates and neutrino-induced reactions

$$\frac{d\sigma}{d\Omega}(q \approx 0) \sim B(GT)$$

Taddeucci *et al.*, Nucl. Phys. A469 (1987) 125-172

- Electron capture rates in astrophysical processes



- FRIB PAC experiment  
“Constraining electron-capture rates in and near the N=20 island of inversion”

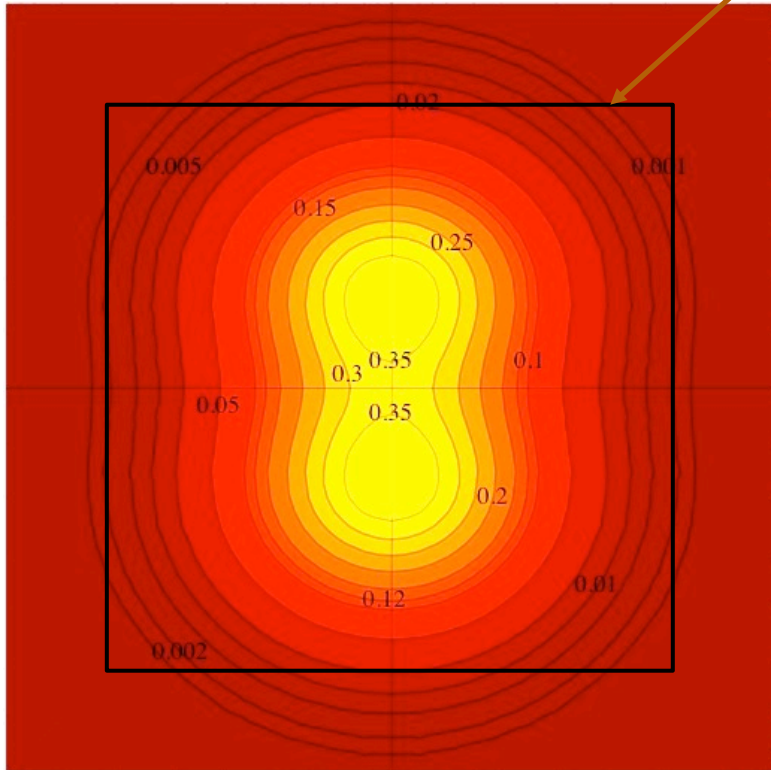


- **Need isospin dependent optical potentials**
- Use ab initio symmetry-adapted no-core shell model to provide input

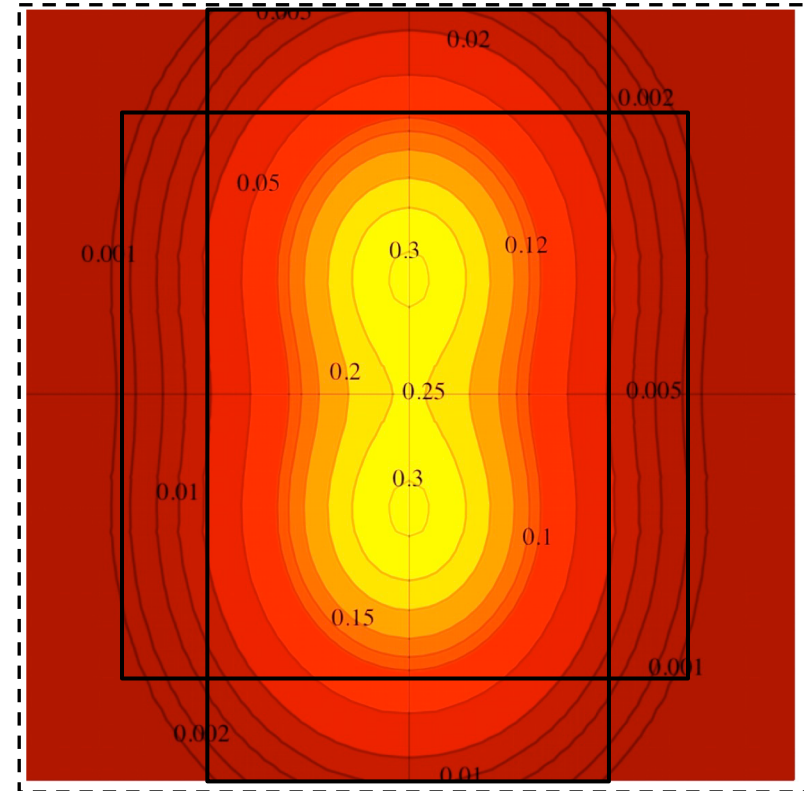
# Symmetry-adapted no-core shell model (SA-NCSM)

Nucleus in model space

Conventional Shell Model

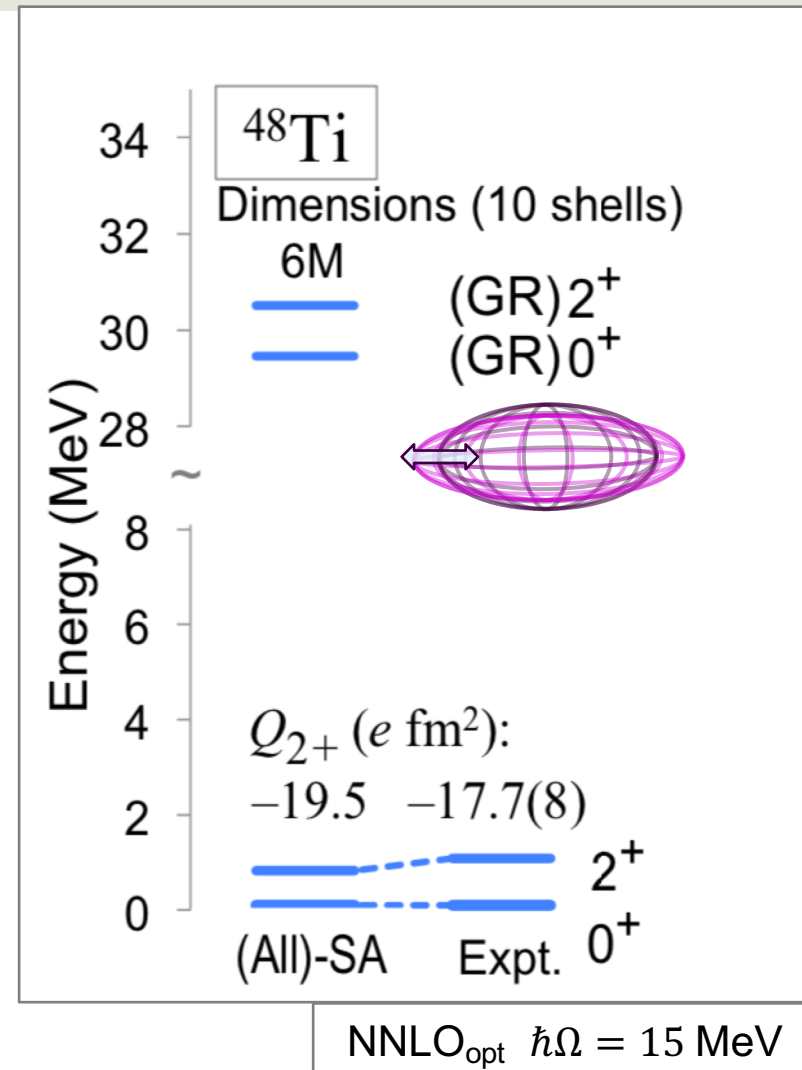
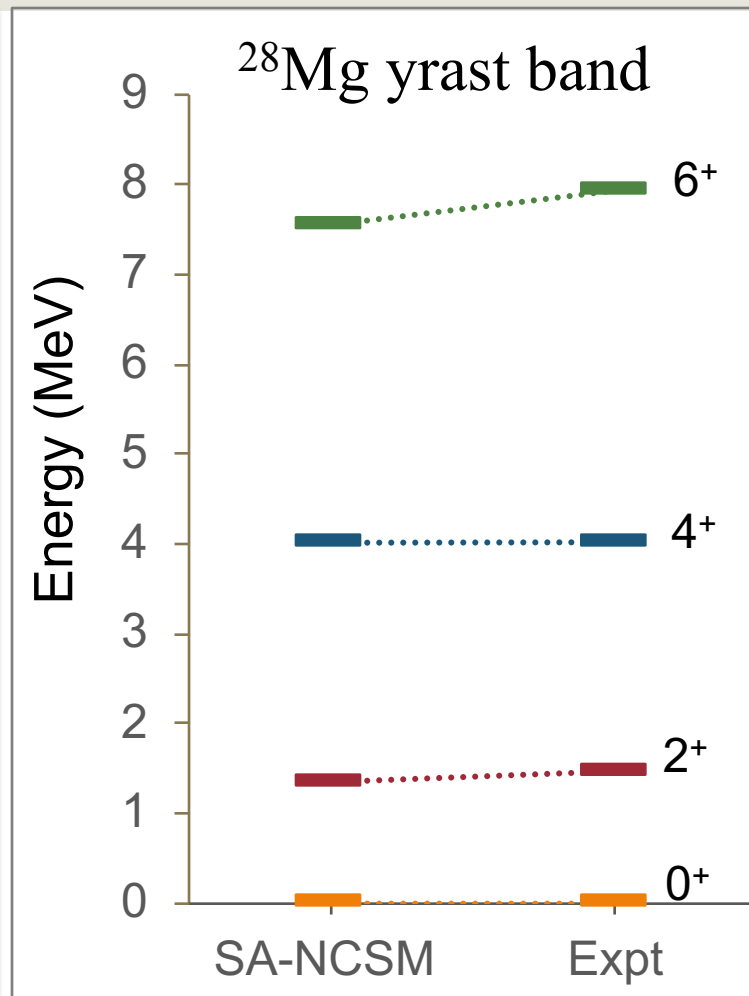
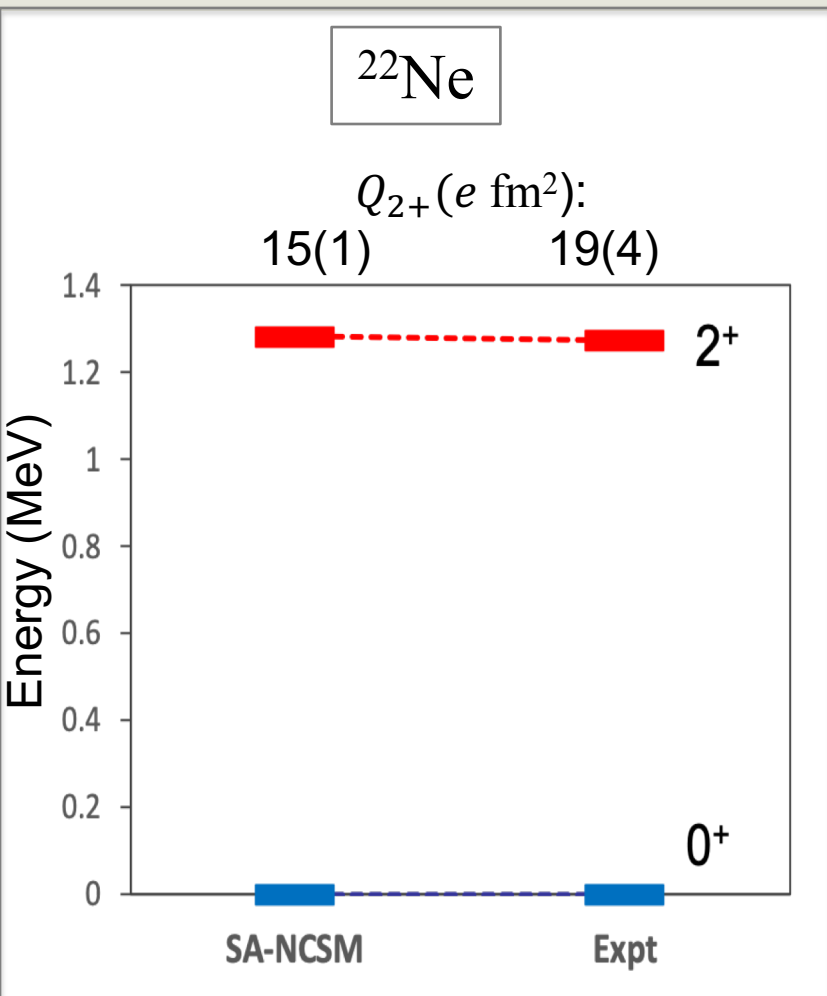


*Ab initio* Symmetry-adapted No-core Shell Model (SA-NCSM)

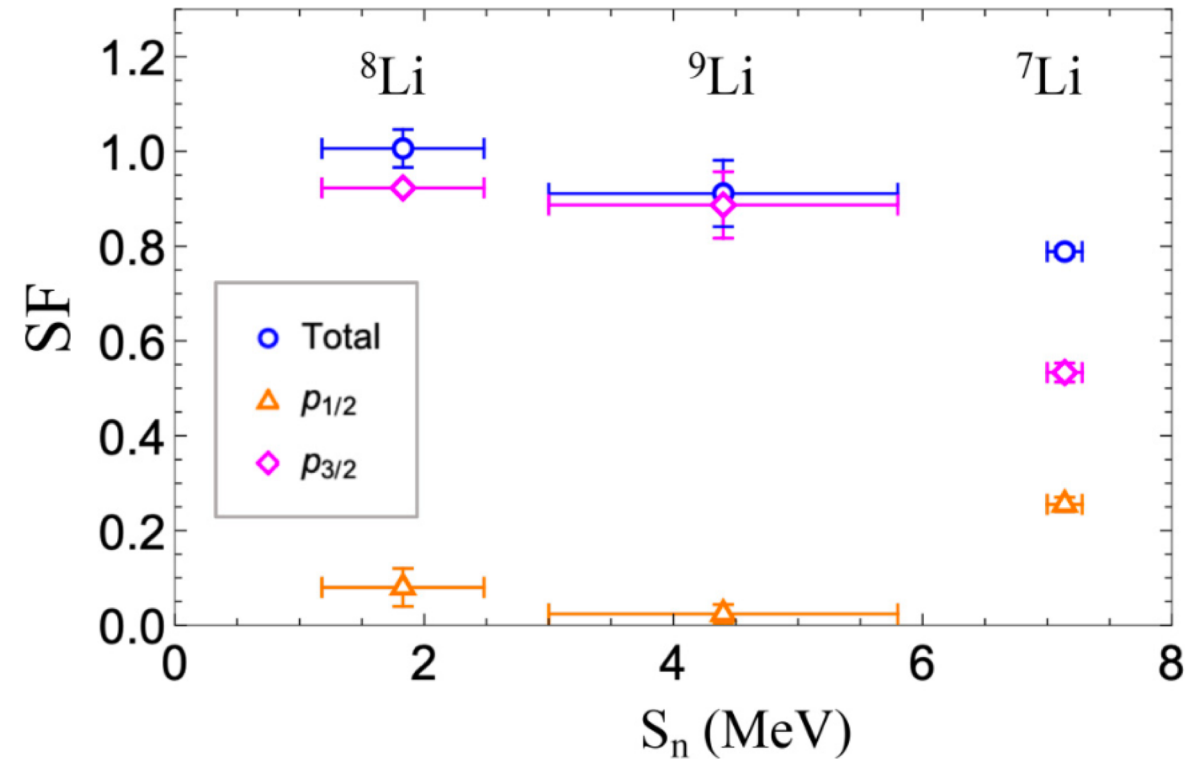
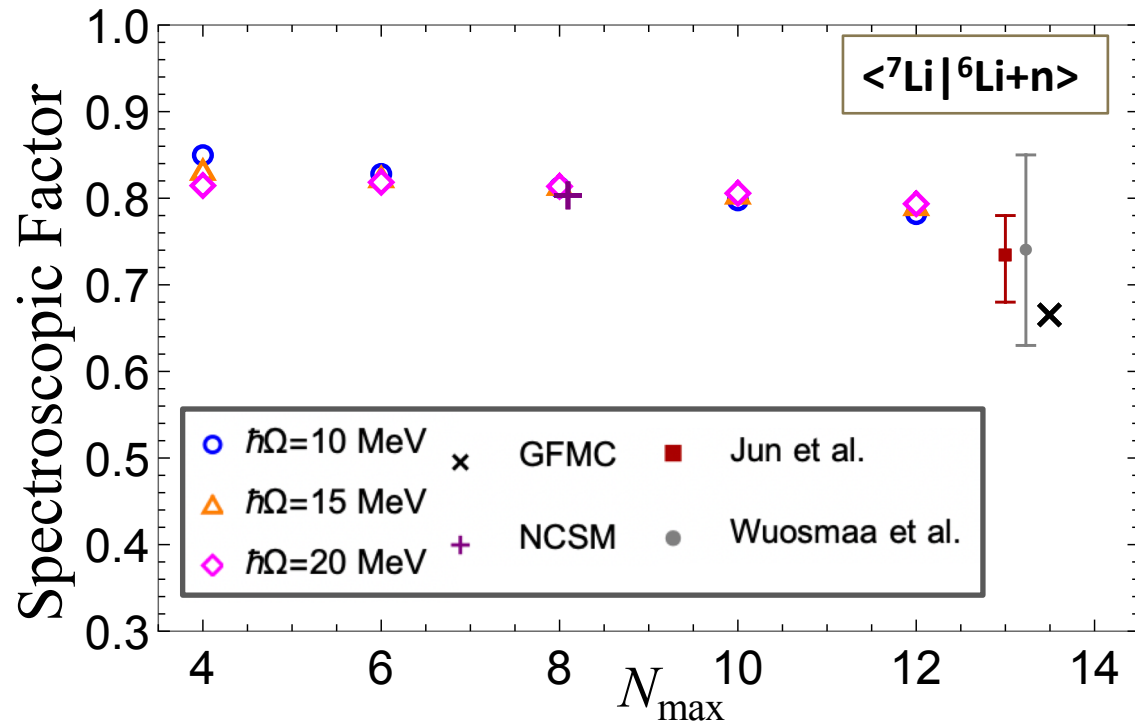


**SU(3) and symplectic symmetry**

# SA-NCSM can reach intermediate mass nuclei



# Ab initio SA-NCSM can provide input for OP



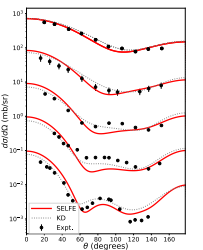
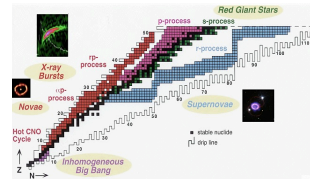
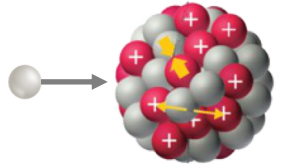
Sargsyan, *et al.* PRC **108**, 054303 (2023)



Facility for Rare Isotope Beams  
 U.S. Department of Energy Office of Science | Michigan State University  
 640 South Shaw Lane • East Lansing, MI 48824, USA  
 frib.msu.edu

# Summary

- We develop a new code to build nucleon-nucleus optical potentials (OPs) for reliable calculations of nuclear reactions
- The method can be applied to any mass range as long as structure calculations are available
- First scattering results for  $^{24}\text{Mg}$  based on shell model structure input are in good agreement with measurements
- We can use ab initio structure input to propagate nucleon-nucleon interaction uncertainties to scattering observables
- Calculations of  $n+^8\text{He}$  scattering with different structure inputs can shed light on the possible parity inversion in  $^9\text{He}$  ground state
- We aim to extend the framework for charge-exchange reaction studies



# Acknowledgements

Gregory Potel, Kostas Kravvaris, Jutta Escher



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Kristina Launey



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U.S. DEPARTMENT OF  
**ENERGY**

# Thank you!



Theory Alliance  
FACILITY FOR RARE ISOTOPE BEAMS

Grigor Sargsyan

# Back up slide zone

## Iterative scheme for self consistent $V_{PO}$

$$\mathcal{V}^{(0)} = V_{00},$$

$$\mathcal{V}^{(1)} = V_{00} + \lim_{\eta \rightarrow 0} \sum_i V_{0i}(r_n) \left( E - T - \mathcal{V}^{(0)}(E_i; \mathbf{r}_n, \mathbf{r}'_n) + i\eta \right)^{-1} V_{i0}(r'_n),$$

...

$$\mathcal{V}^{(n+1)} = V_{00} + \lim_{\eta \rightarrow 0} \sum_i V_{0i}(r_n) \left( E - T - \mathcal{V}^{(n)}(E_i; \mathbf{r}_n, \mathbf{r}'_n) + i\eta \right)^{-1} V_{i0}(r'_n),$$

$$J^{(n)} = \int \mathcal{V}^{(n)}(\mathbf{r}_n, \mathbf{r}'_n) d\mathbf{r} d\mathbf{r}',$$

$$\varepsilon = \left| \frac{J^{(n+1)} - J^{(n)}}{J^{(n+1)} + J^{(n)}} \right| \ll 1.$$

Volume integral  
convergence condition



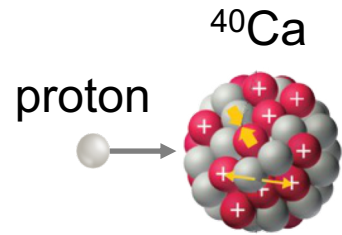
# Elastic and absorption cross sections can be calculated from the OP

$$V(\mathbf{r}, \mathbf{r}', E) = U_0(\mathbf{r}) + V_{PO}(\mathbf{r}, \mathbf{r}', E - E_i)$$

$$(E - T - V(\mathbf{r}, \mathbf{r}', E))\phi = 0 \quad \longrightarrow \quad \text{elastic scattering cross sections from phase shifts}$$

$$\sigma_{abs} \sim \langle \phi | \text{Im}(V_{PO}) | \phi \rangle = 0 \quad \longrightarrow \quad \text{Absorption cross section from imaginary part of the polarization potential}$$

# $^{40}\text{Ca} + p$ elastic scattering with 10 states



Parameters for levels in  $^{40}\text{Ca}$

$\lambda_n^\pi$	$1^-$	$2^+$	$2^+$	$2^+$	$3^-$	$3^-$	$4^+$	$4^+$	$5^-$	$5^-$
$E_n$ (MeV)	18.0	3.9	8.0	16.0	3.73	15.73	8.0	20.0	4.48	16.48
$\beta_\lambda(n)$	0.087	0.143	0.309	0.250	0.354	0.380	0.254	0.457	0.192	0.653

Rao, et al., Nuclear Physics A207 (1973) 182-208.

$$V(\mathbf{r}, \mathbf{r}', E) = V_{00}(\mathbf{r}, E) + V_{PO}(\mathbf{r}, \mathbf{r}', E)$$

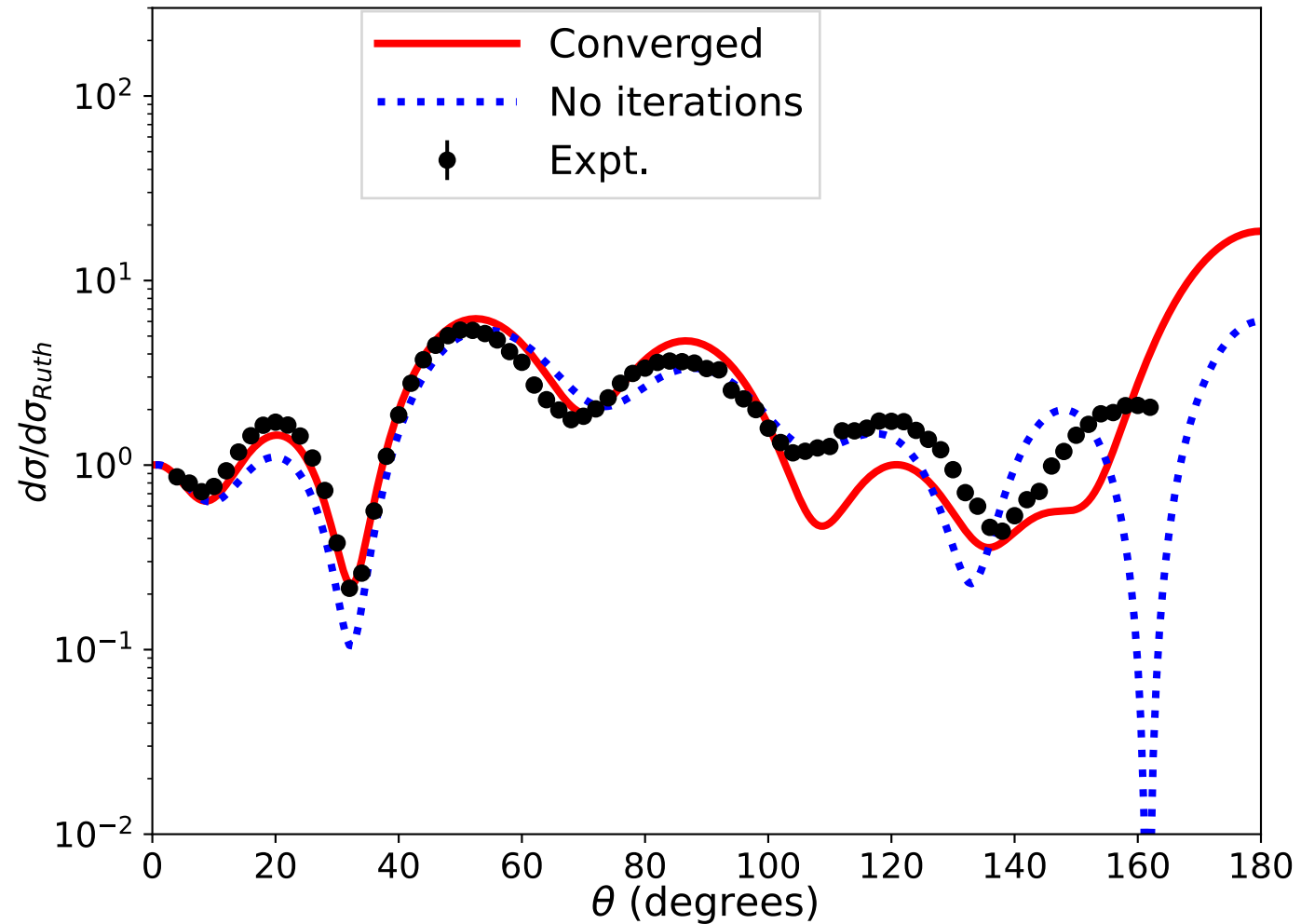
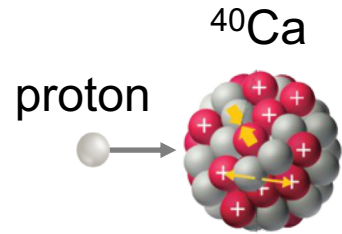


Woods-Saxon:  
 $V_0 = -50$  MeV  
 $R_0 = 1.2A^{1/3}$   
 $a = 0.65$



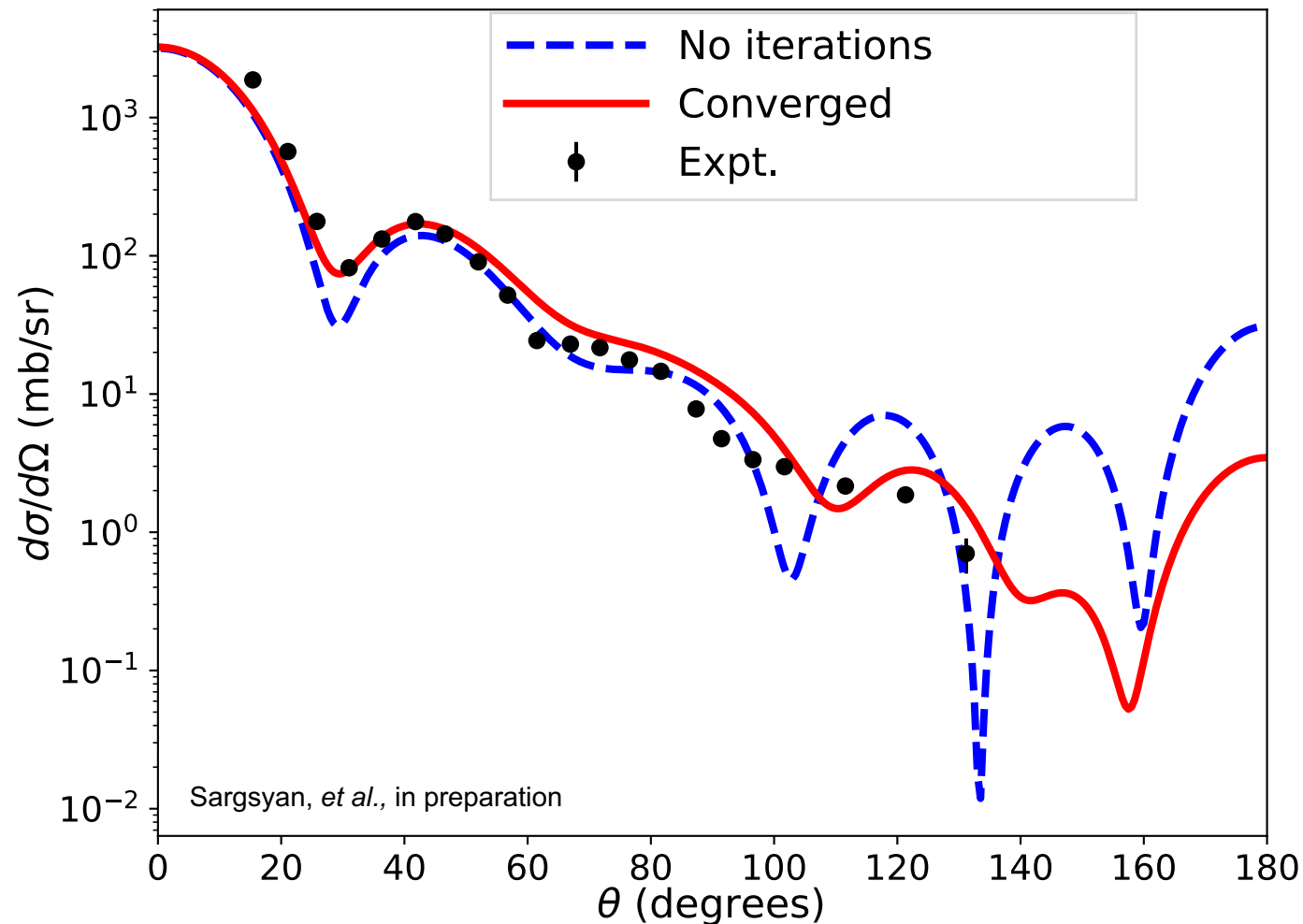
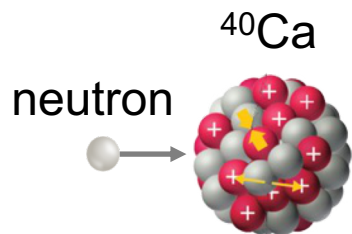
10 intrinsic  
 states

# $^{40}\text{Ca} + p$ elastic scattering at 30 MeV



10 intrinsic states

# Neutron elastic scattering over $^{40}\text{Ca}$ at 30 MeV



10 intrinsic states