

# EXPERIMENT AND OPTICAL POTENTIALS: A USER'S POINT OF VIEW

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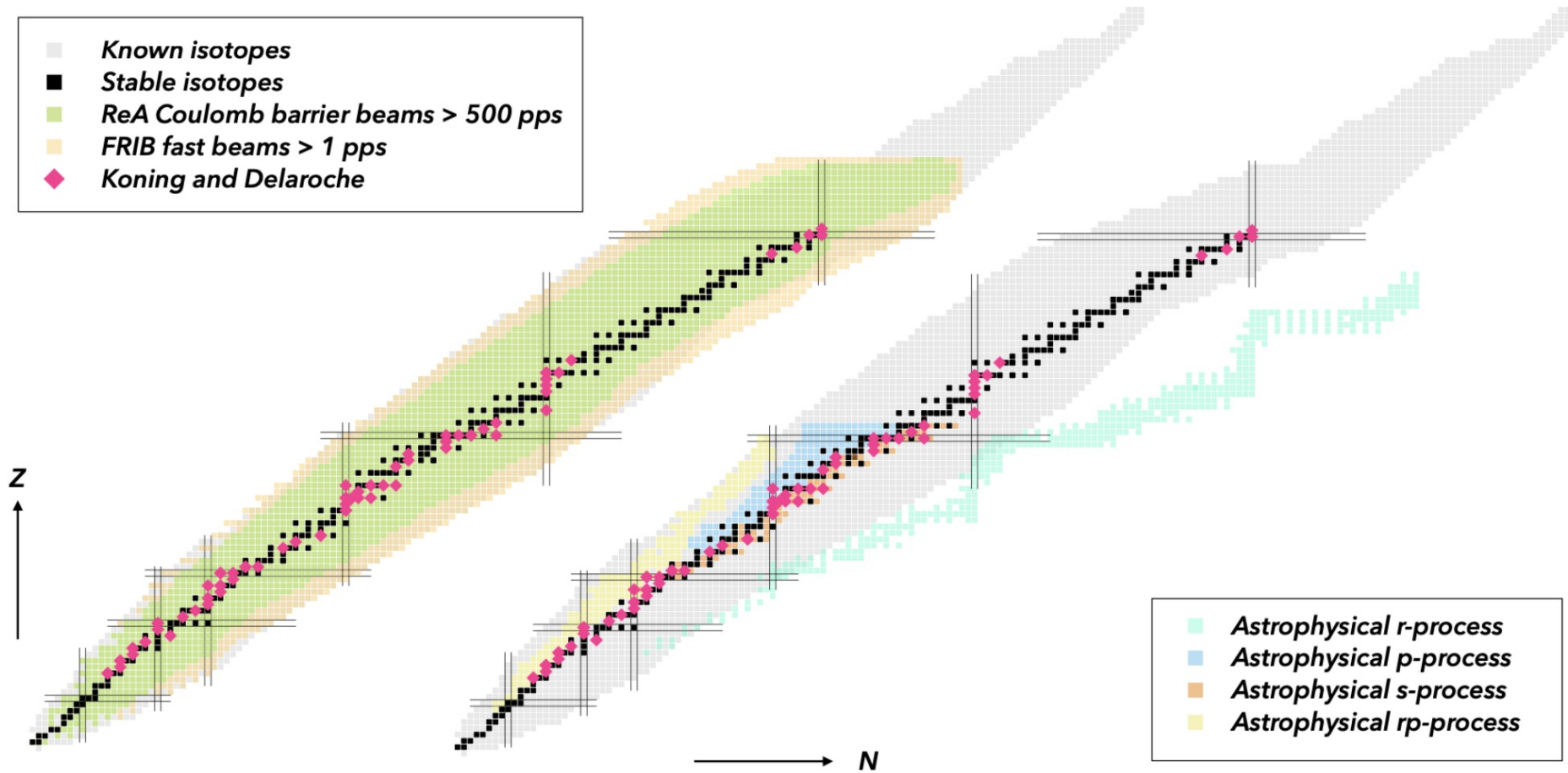
Alexandre Obertelli, TU Darmstadt

# OUTLINE

- **Direct reactions: needs and reality**  
Aumann et al., PPNP (2021); Aumann et al., PRL (2017); Ponnath et al., PLB (2024)
- **One nucleon transfer from  $^{14}\text{O}$ : a sensitivity study**  
Flavigny et al., PRL (2013); Flavigny et al., PRC (2018)
- **Quasifree scattering**  
Enciu et al., PRL (2022)
- **Uncertainties and energy dependence**  
Pohl et al., PRL (2023); Bertulani, PLB (2023); Gomez-Ramos et al., PLB (2023)
- **Stripping from low-energy antiprotons**  
Aumann et al. (PUMA collaboration), EPJA (2022)

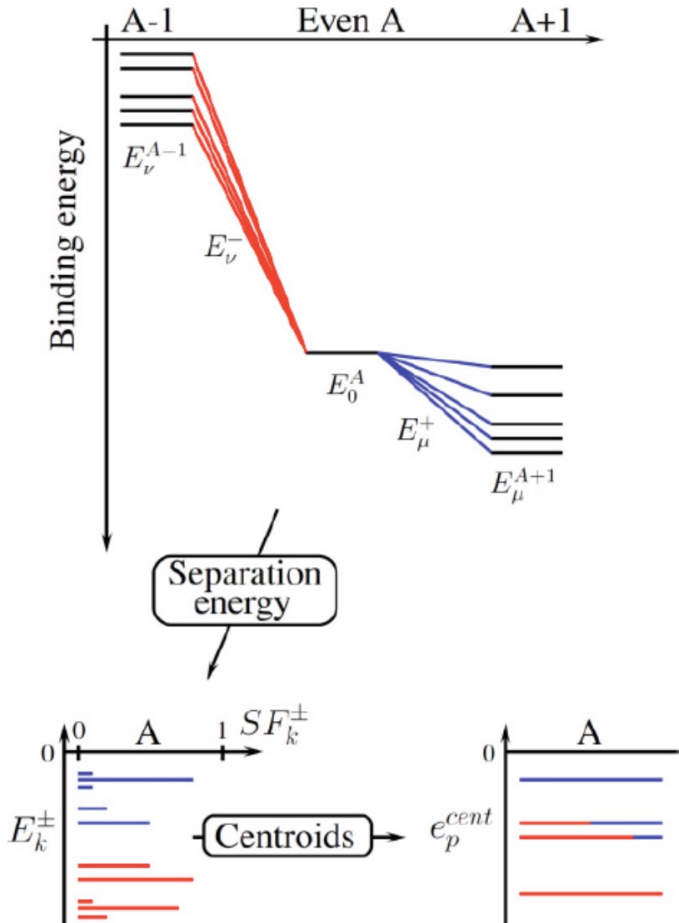


# THE AIM OF DIRECT REACTIONS



Hebborn et al., JPG (2023), shown at Haloweek 2024

# THE AIM OF DIRECT REACTIONS



- One nucleon pickup or stripping used to probe the structure of nuclear states through spectroscopic strength

$$S_k^{n\ell j,+} = |\langle \Psi_k^{A+1} | a_{n\ell j}^\dagger | \Psi_0^A \rangle|^2 \quad S_k^{n\ell j,-} = |\langle \Psi_k^{A-1} | a_{n\ell j} | \Psi_0^A \rangle|^2$$

- Single-particle energies relate to physical state energies via the Baranger relation

$$e_{n\ell j} = \frac{\sum_k S_k^{n\ell j,+} (E_k - E_0) + S_k^{n\ell j,-} (E_0 - E_k)}{\sum_k S_k^{n\ell j,+} + S_k^{n\ell j,-}}$$

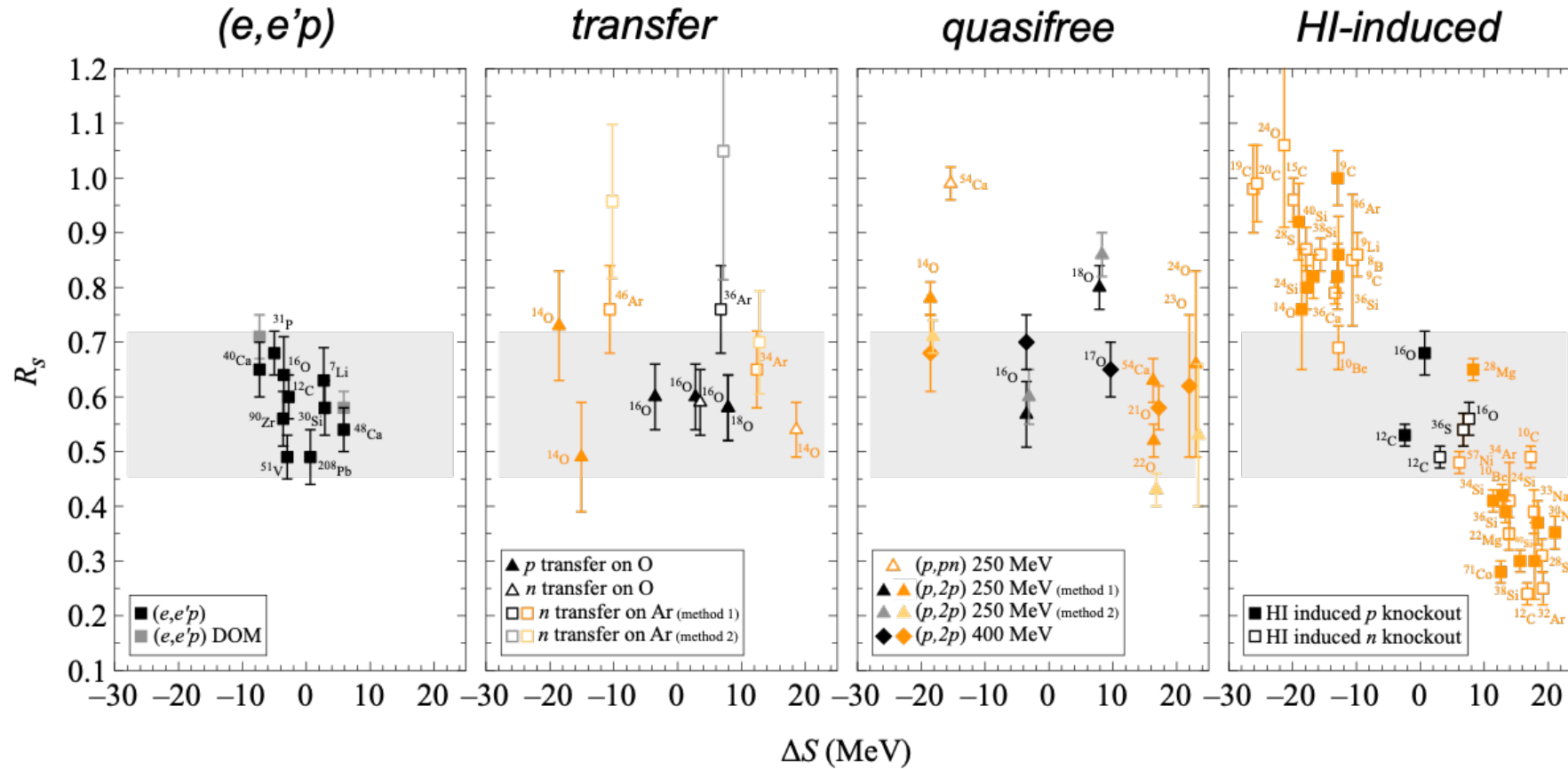
- Limitations:
  - Non observability of single particles and spectroscopic factors
  - Surface or asymptotic sensitivity of direct reactions
  - **Reaction model approximations and undefined uncertainties**
  - Experimental uncertainties and incomplete data sets

Hagen, Duguet, PRC (2012)

# TODAY'S REALITY



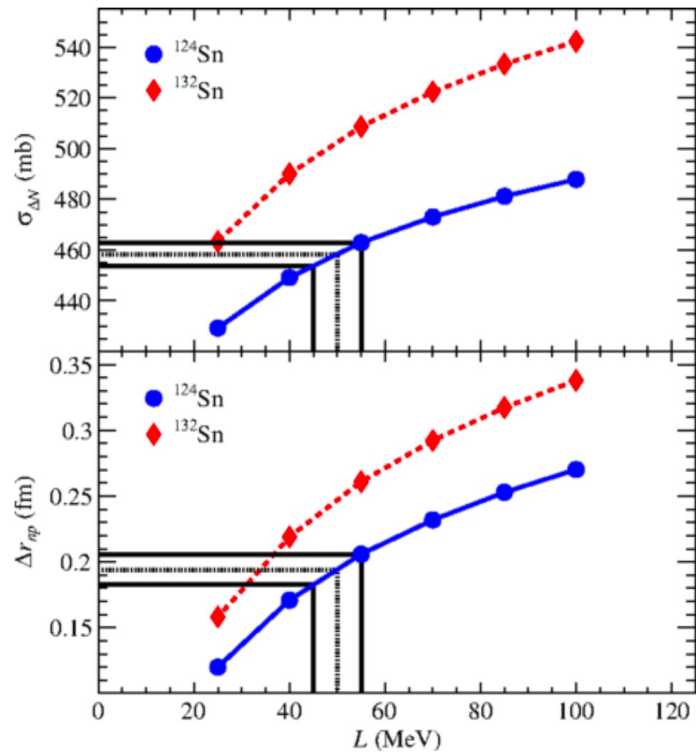
Reaching accuracy and precision of 10% would be a game changer



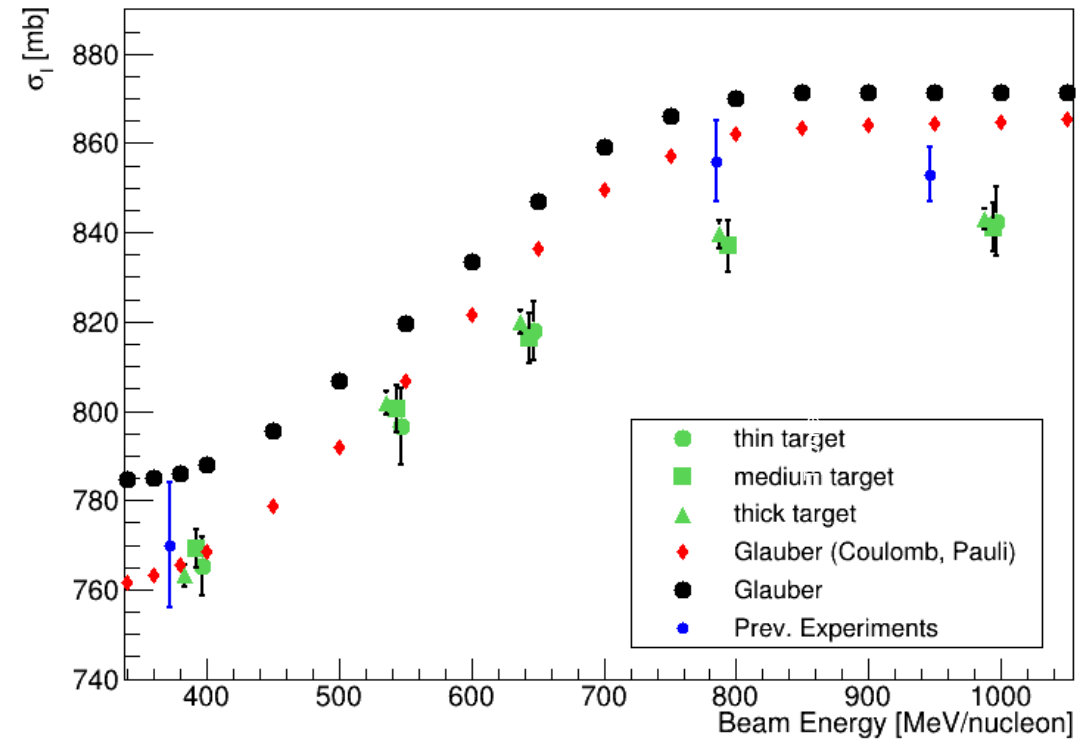
Aumann et al., Prog. Part. Nucl. Phys. (2021)

# SKINS FROM NEUTRON REMOVAL

- Sensitivity to symmetry term of EOS from neutron removal cross sections would require a 1% accuracy
- $^{12}\text{C}$ - $^{12}\text{C}$  interaction cross section reproduced at 3% by Glauber



Aumann et al., PRL (2017)



Ponnath et al., PLB (2024)

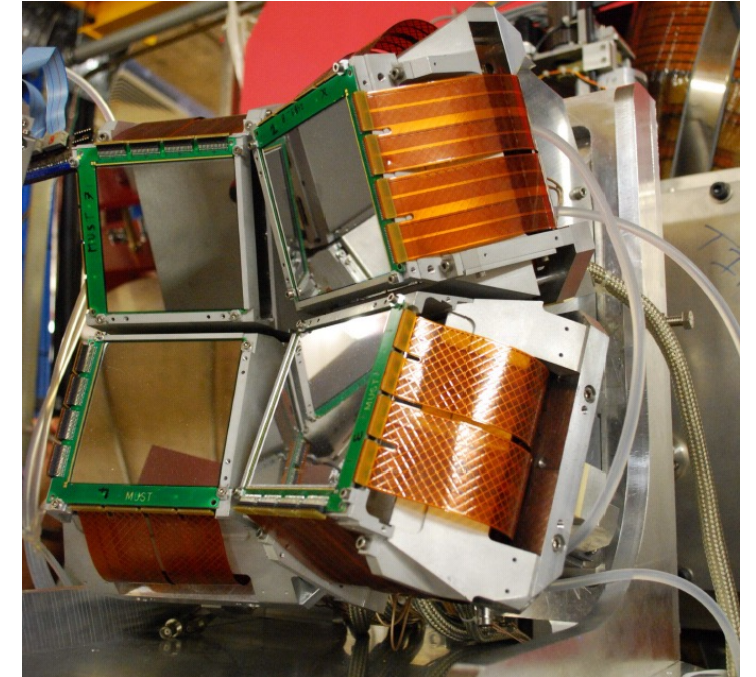
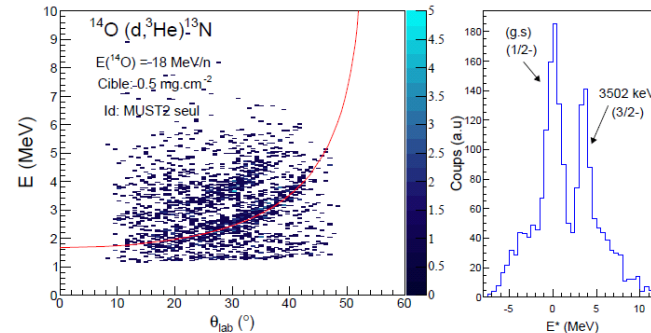
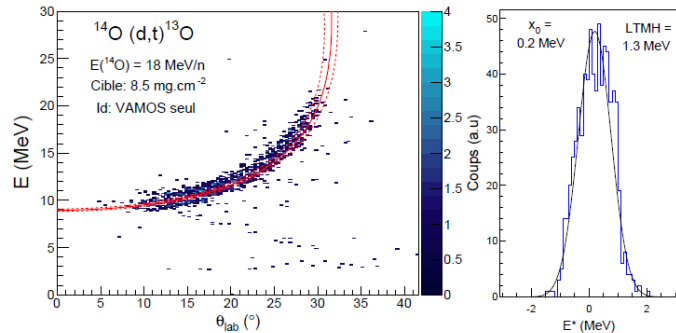
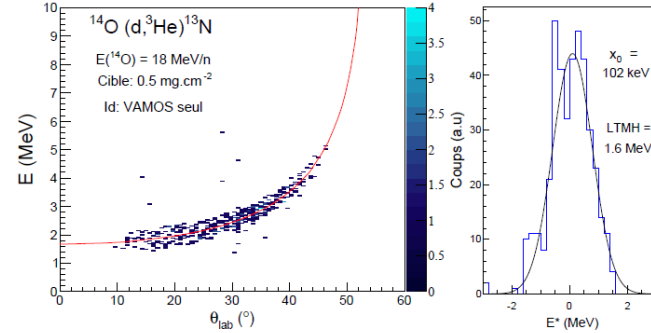
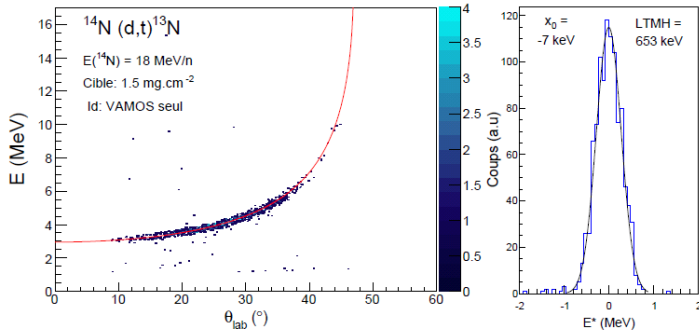
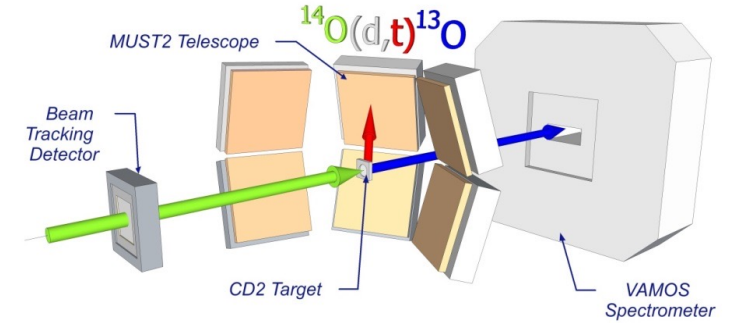
# TRANSFER: AN EXAMPLE

SPIRAL beam:  $^{14}\text{O}$  at 18 MeV/nucleon

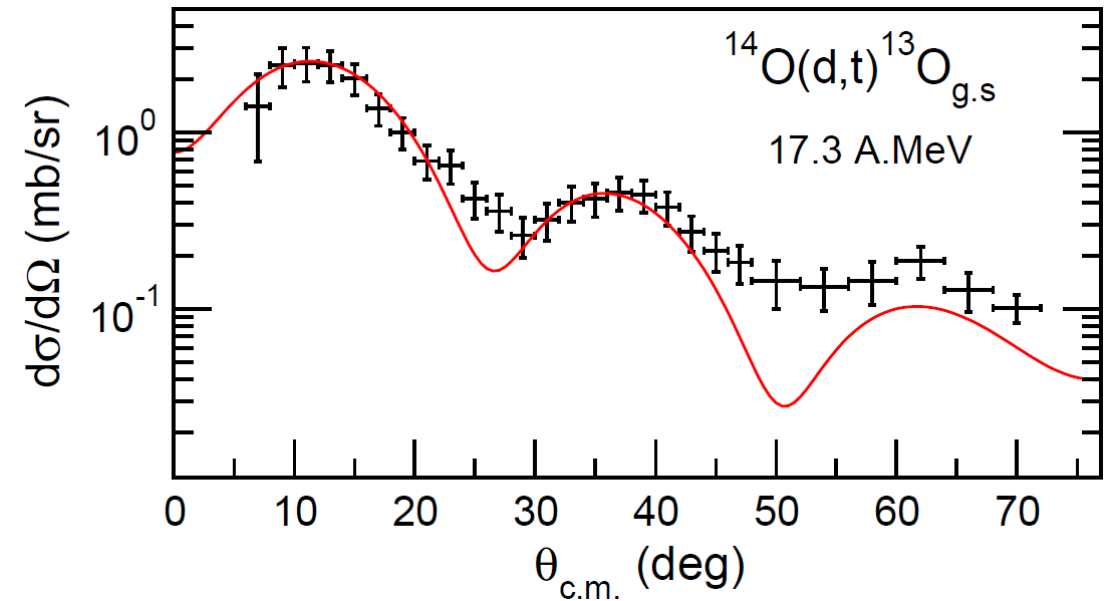
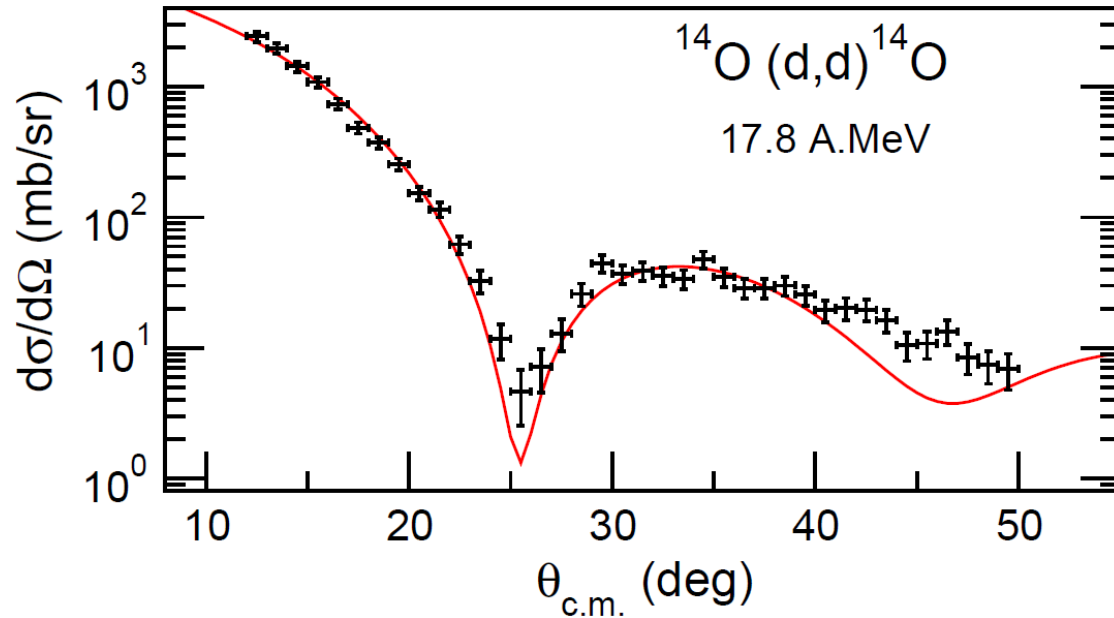
Intensity:  $5 \cdot 10^4$  pps

Targets:  $\text{CD}_2$

Reactions:  $(d,d)$ ,  $(d,t)$  and  $(d,^3\text{He})$ , *fully exclusive measurements*



# TRANSFER AND ELASTIC SCATTERING DIFFERENTIAL CROSS SECTIONS

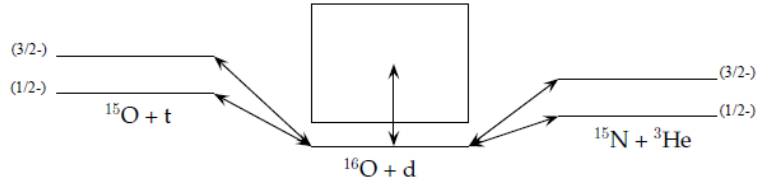


Flavigny et al., PRL (2013)



# ANALYSIS

- Coupled Reaction Channels (CRC) analysis (FRESCO code)



- Potentials

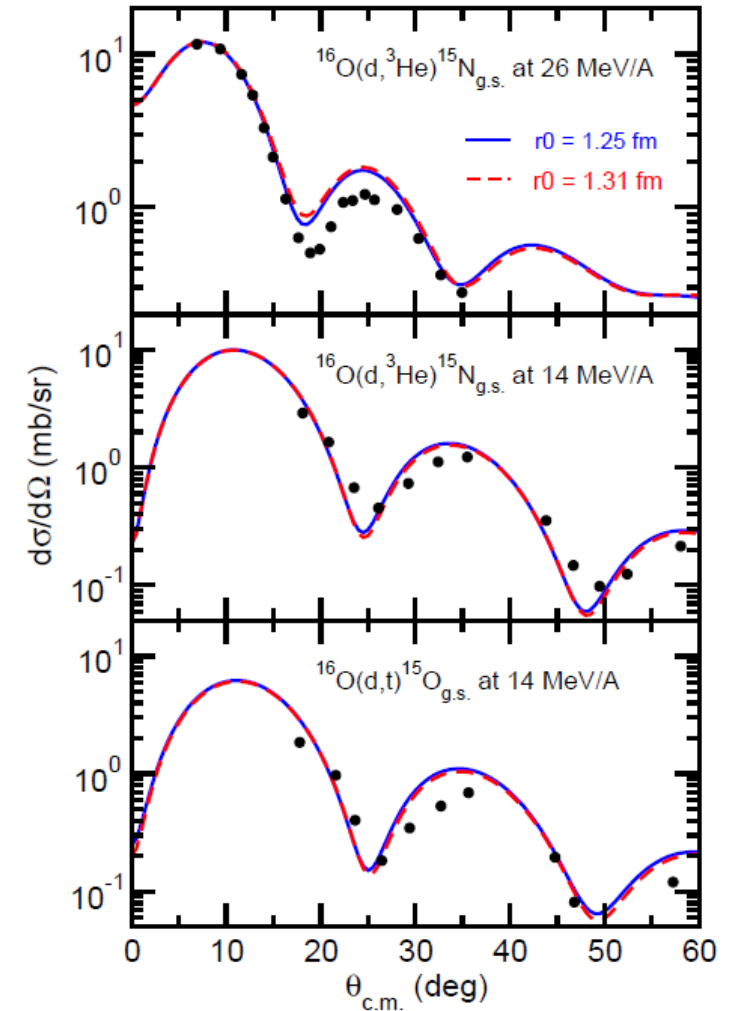
- Entrance: Global N-nucleus + folding *ex. König and Delaroche, NPA (2003)*
- Renormalised entrance potential on elastic scattering ( $V:1.1, W:0.8$ )
- Exit: Global  $t/3\text{He}$  – nucleus *ex. Perey and Perey, ADNDT (1976)*

- Form factors

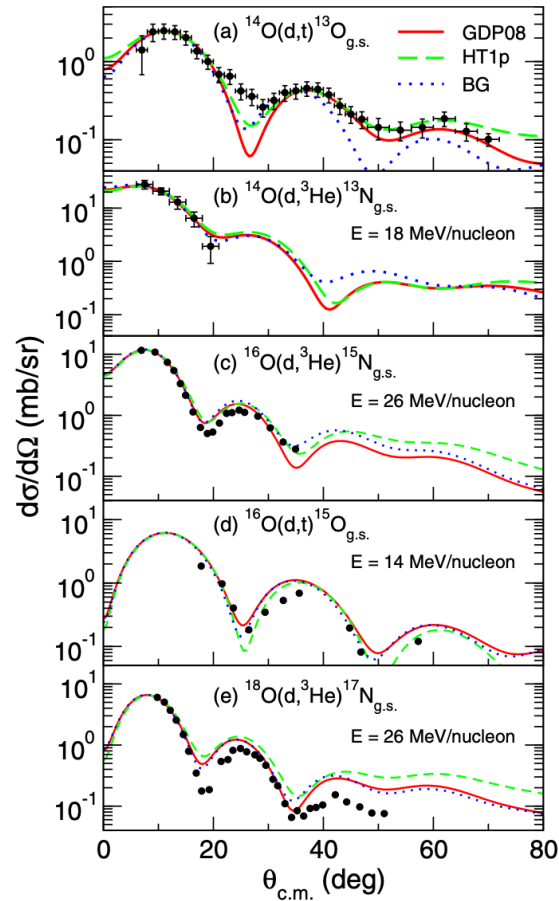
- WS,  $r_0=1.31$  fm from  $(e,e'p)$  or  $r_0=1.25$  fm and  $a_0=0.65$  fm

	Notre analyse		$(e,e'p)[3]$
	$r_0 = 1.25$ fm	$r_0 = 1.31$ fm	$r_0 = 1.31$ fm
$C^2 S_{exp}$	1.5(3)	0.94(30)	1.17(7)
$C^2 S_{th}$	1.51	1.51	1.51
$R_s$	1.0(2)	0.62(20)	0.77(5)

- Microscopic from SCGF theory

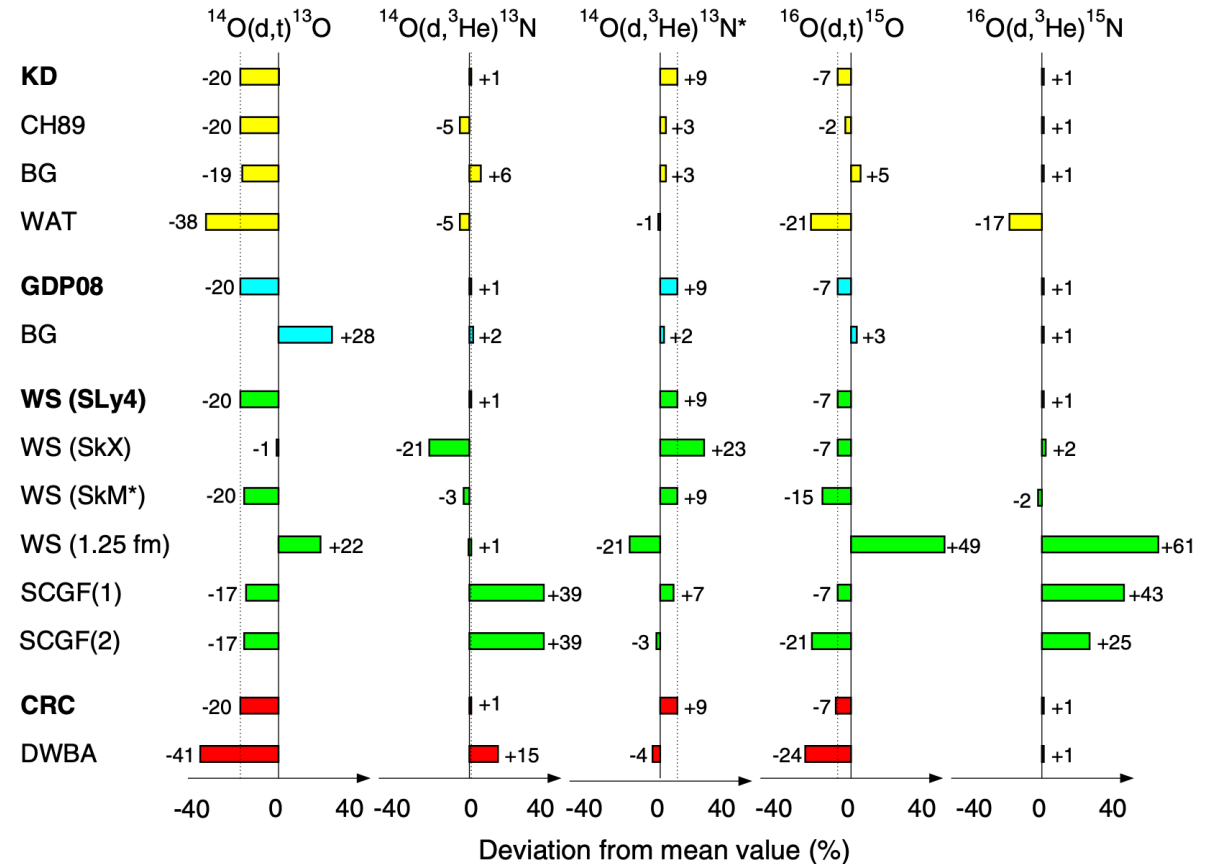


# SENSITIVITY STUDY



Flavigny et al., PRC (2018)

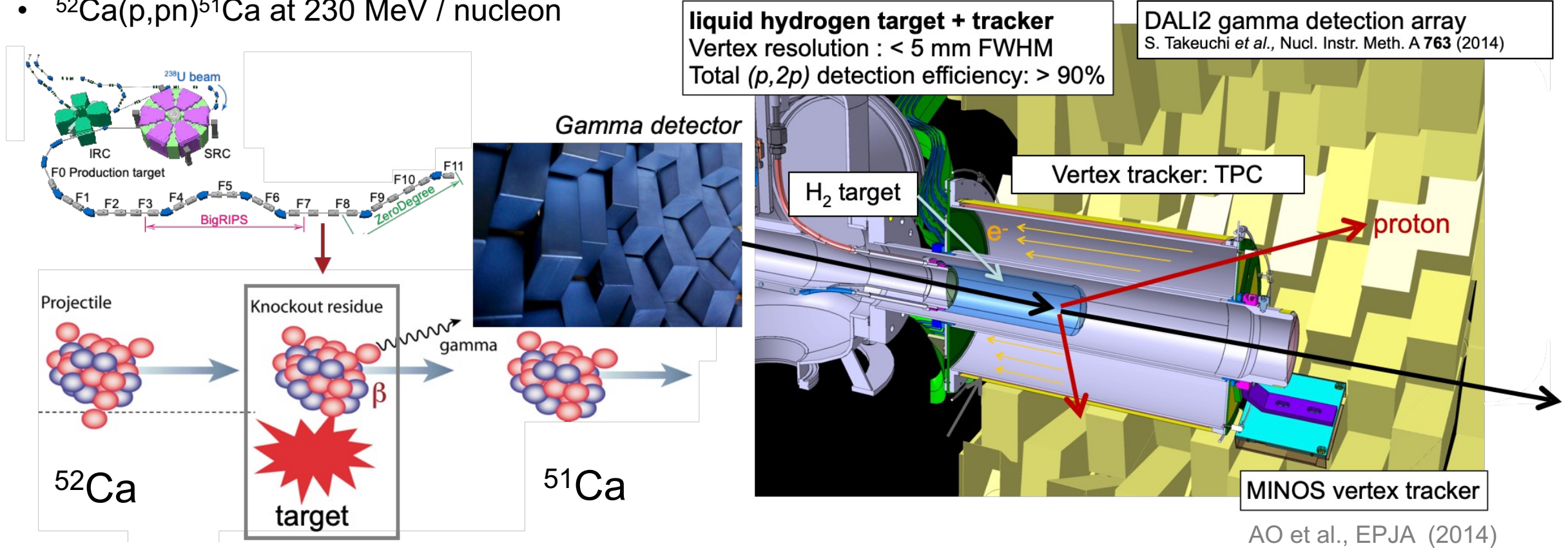
König & D., NPA (2003)  
 Varner et al., PR (1991)  
 Becchetti & G., PR (1969)  
 Watson et al., PR (1969)  
 Pang et al., PRC (2009)



Global potentials fitted on data from stable nuclei, leading to non-estimated systematic errors from to isospin / binding energy dependence

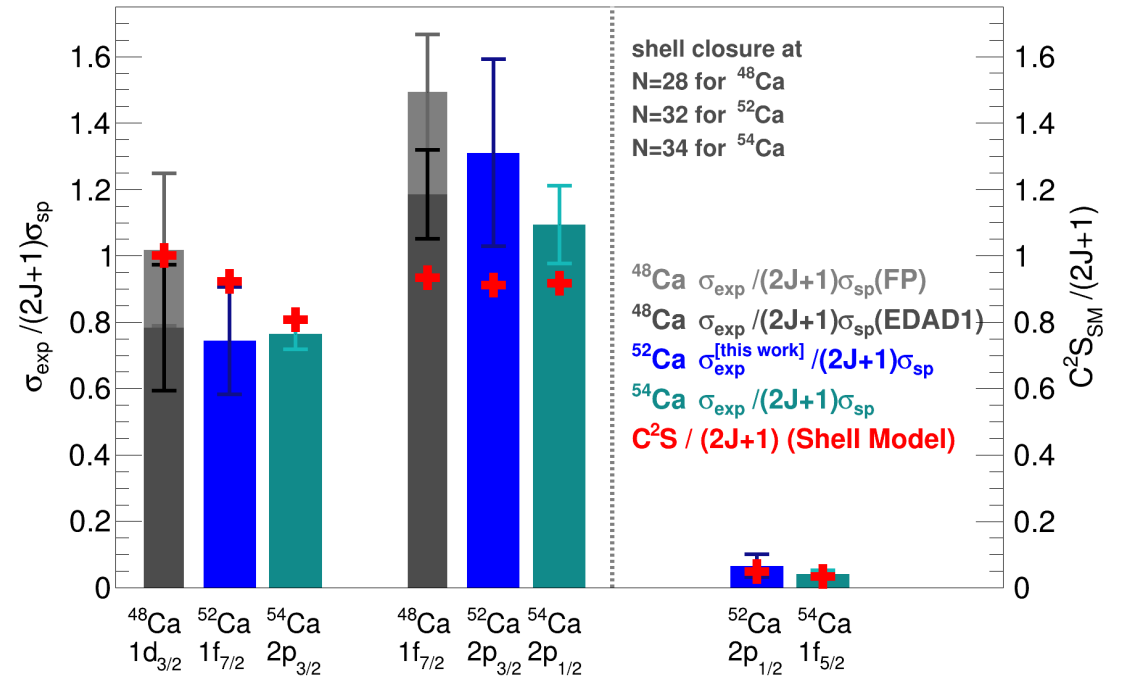
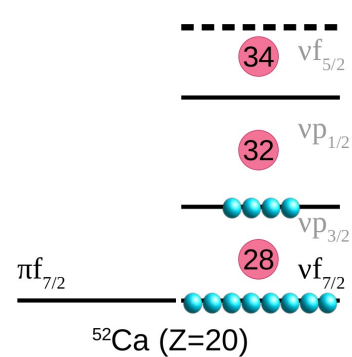
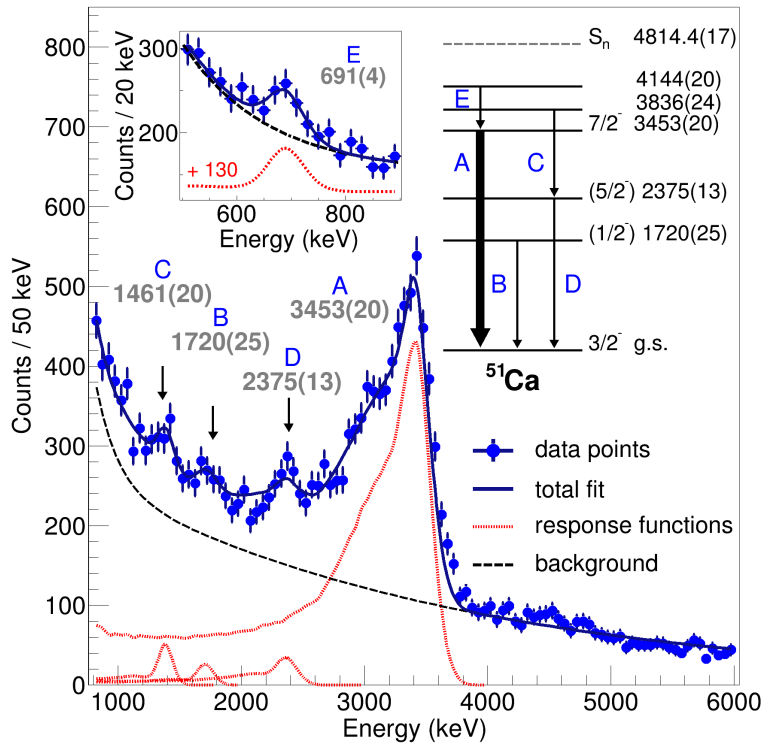
# THE EXAMPLE OF $^{52}\text{Ca}(p,pn)^{51}\text{Ca}$

- RIBF, RIKEN
- $^{70}\text{Zn}$  primary beam at 345 MeV/nucleon, 200 pA
- $^{52}\text{Ca}(p,pn)^{51}\text{Ca}$  at 230 MeV / nucleon



# THE EXAMPLE OF $^{52}\text{Ca}(p,pn)^{51}\text{Ca}$

- DWIA single-particle cross sections: K. Ogata, K. Yoshida (see talk by K. Yoshida)
- Effect of OP for p-nucleus interaction estimated to 15% at 150 MeV/n ( $^{48}\text{Ca}$ ), less than 5% above 200 MeV/n
- Folding potential with Melbourne G matrix: Amos et al., ANP (2000), DIRAC potential: Cooper et al., PRC (1993)



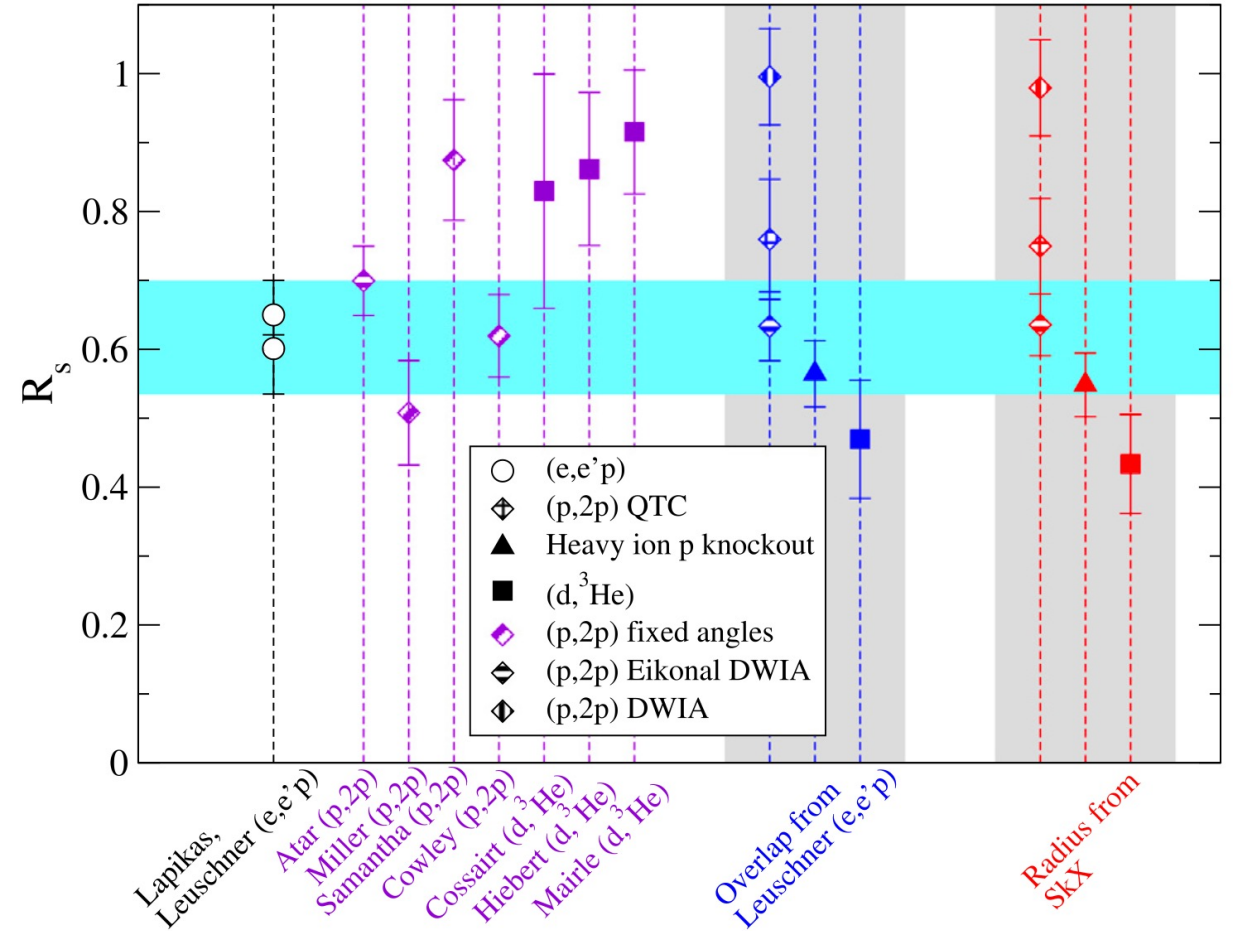
Enciu et al., PRL (2022)

# SENSITIVITY STUDY

- One proton removal from  $^{16}\text{O}$  (inclusive to 1/2- ground state and 3/2- state in  $^{15}\text{N}$ )
- Comparison of analysis from (e,e'p), transfer, quasifree scattering and light-ion induced knockout
- DWIA, eikonal DWIA and QTC (Quantum Transfer to the Continuum) compared: variations of  $R_s$  from 0.6 to 1, i.e. ~40%
- Effect of pA optical potential (QTC) : ~10%

DWIA  $T_{p,pN} = \sqrt{S(lj)} \langle \chi_{\mathbf{k}'_p}^{(-)} \chi_{\mathbf{k}_N}^{(-)} | \tau_{pN} | \chi_{\mathbf{k}_p}^{(+)} \psi_{jlm} \rangle$

QTC  $\mathcal{T}_{if}^{3b} = \langle \Psi_f^{3b(-)} | V_{pN} + U_{pC} - U_{pA} | \phi_{CA} \chi_{pA}^{(+)} \rangle$



Aumann et al., Prog. Part. Nucl. Phys. (2021)

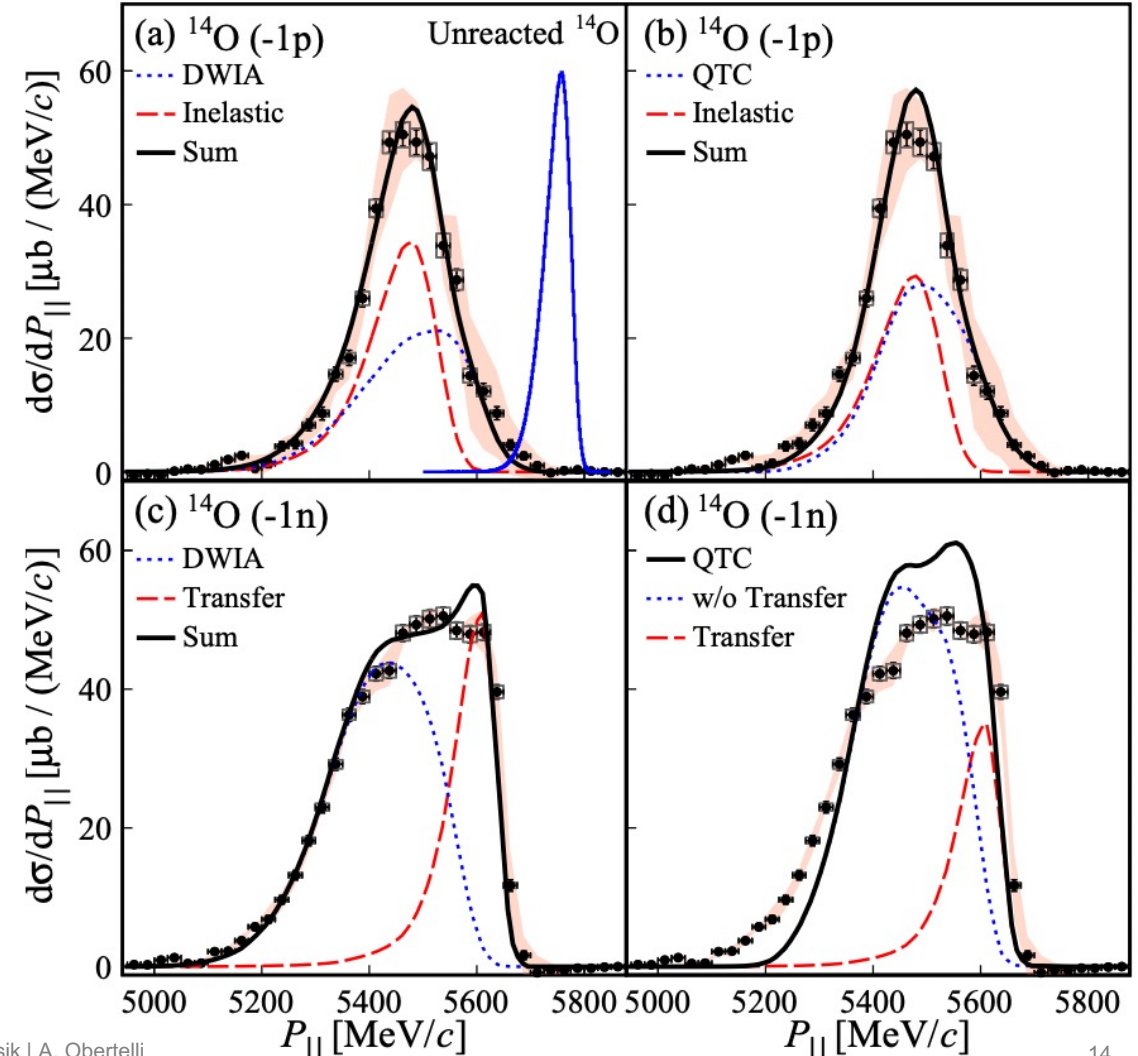
# PROTON-INDUCED REMOVAL AT 94 MEV/U

- Nucleon removal from  $^{14}\text{O}$  on hydrogen, 94 MeV/nucleon.  $^{14}\text{O}$ :  $S_n = 23$  MeV,  $S_p = 4.6$  MeV.
- Interpreted with Quantum Transfer to the Continuum (QTC) and eikonal DWIA.

Residue	$J^\pi$	$\sigma_{\text{exp}}$ [mb]	SF	Theory	$\sigma_{\text{sp}}$ [mb]	$\sigma_{\text{th}}$ [mb]	$R_s$
$^{13}\text{N}_{\text{g.s.}}$	$1/2^-$	10.7(16)	1.58	DWIA	5.2	8.8	1.22(18)
				Inelastic	-	9	
				Sum	-	17.8	0.60(9)
				QTC	7.0	11.9	0.90(13)
$^{13}\text{O}_{\text{g.s.}}$	$3/2^-$	16.7(24)	3.42	DWIA	6.3	23.2	0.72(10)
				Transfer	3	11	
				Sum	-	34.2	0.49(7)
				QTC	10.2	37.6	0.44(6)
				w/o transfer	13.5	49.7	0.34(5)

- See also: Bonnacorso, Brink, PRC (1991), Flavigny et al., PRL (2013)

Pohl et al., PRL (2023)



# CORE DESTRUCTION IN KNOCKOUT

Letter

Isospin dependence in single-nucleon removal cross sections explained through valence-core destruction effects

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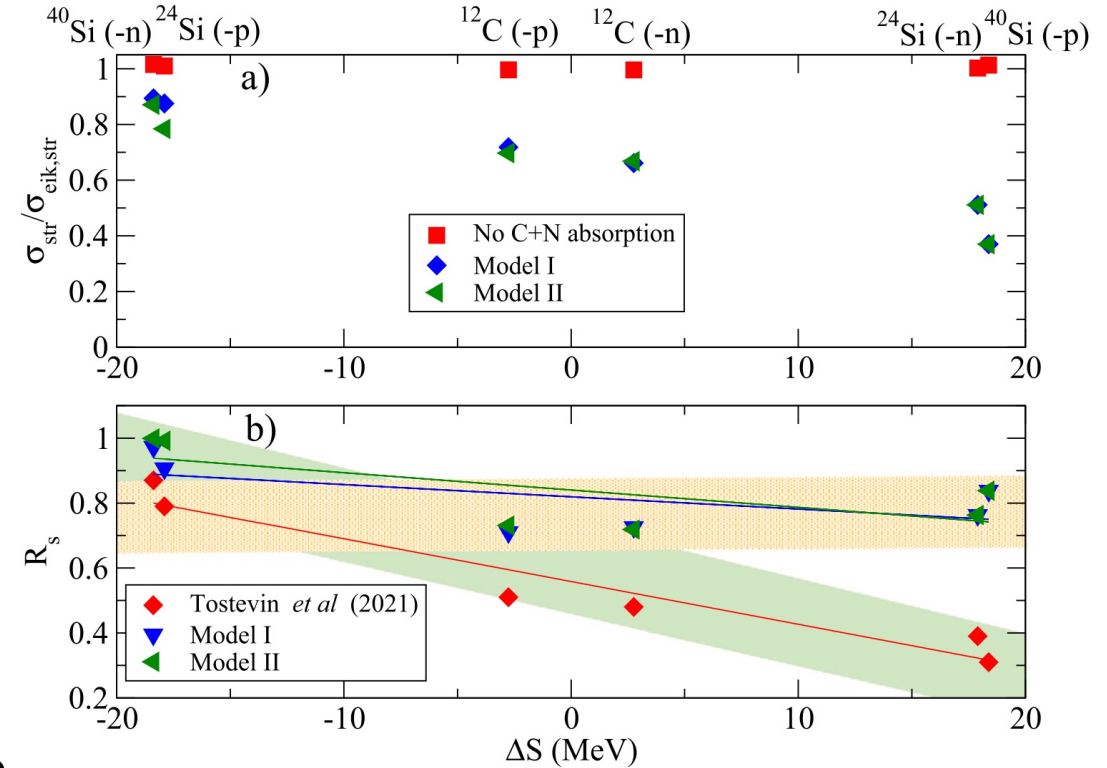
<sup>b</sup> Centro Nacional de Aceleradores (U. Sevilla, J. Andalucía, CSIC), Tomás Alva Edison, 7, 41092 Sevilla, Spain

<sup>c</sup> Instituto Interuniversitario Carlos I de Física Teórica y Computacional (iC1), Apdo. 1065, E-41080 Sevilla, Spain

$$P_{\text{str}}(\vec{b}) \simeq \int dx dy \rho^{(2)\text{eff}}(x, y) \times |S_{CT}^0(b_{CT})|^2 (1 - |S_{VT}^0(b_{VT})|^2)$$

$$\rho^{(2)\text{eff}}(x, y) = \int d^3\vec{r}_1 \int d^3\vec{r}_2 \langle \vec{r}_2 | \rho_f | \vec{r}_1 \rangle \phi_g^*(\vec{r}_2) \phi_g(\vec{r}_1) \times \delta\left(x - \frac{x_1 + x_2}{2}\right) \delta\left(y - \sqrt{\frac{y_1^2 + y_2^2}{2}}\right),$$

- Stripping from light ion target (<sup>9</sup>Be, <sup>12</sup>C)
- Nucleon – core interaction from rescaled global energy-dependent dispersive potential Morillon et al., PRC (2007)
- **large impact on cross section and isospin dependence**
- Conclusions at variance with Bertulani, PLB (2023)
- See also: Louchart et al., PRC (2011), Hebborn et al., PRC (2023)



Gomez-Ramos et al., PLB (2023)

# CORE DESTRUCTION IN KNOCKOUT

(SLIDE ADDED AFTER TALK)

- Eikonal formalism
- **Reduction as much as 9.5 % found**
- **No nucleon-binding-energy dependence found**

**Table 1**

Theoretical values for one-nucleon-removal stripping cross sections for selected single-particle states.  $\Delta S = S_n - S_p$  for neutron removal or  $\Delta S = S_p - S_n$  for proton removal.

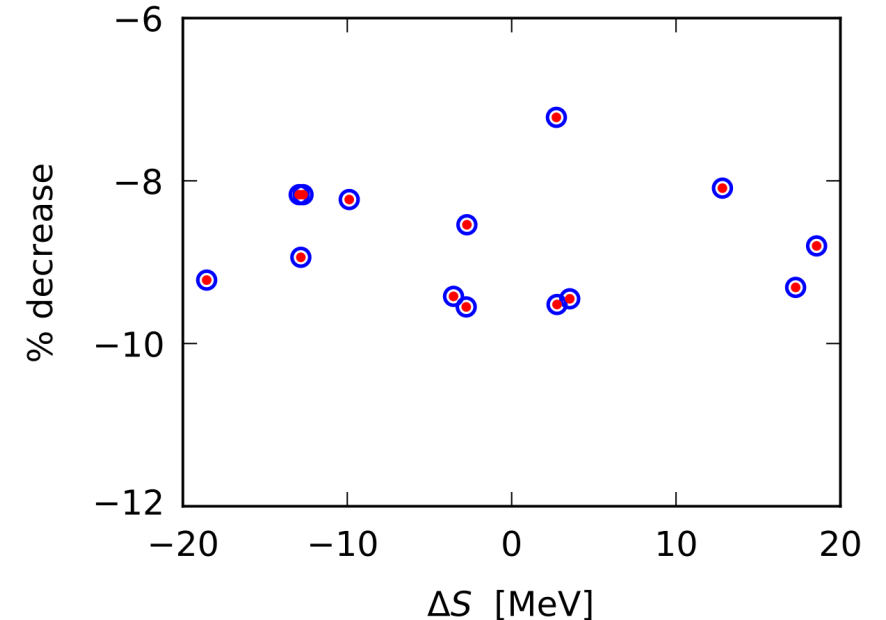
Reaction	$E_{beam}$ MeV/nucleon	$S_p[S_n]$ MeV	$\Delta S$ MeV	$J^\pi(j)$ mb	$\sigma_{str}$ mb	$\sigma_{str}^{rescatt}$ mb	% change
${}^9\text{B}({}^7\text{Li}, {}^6\text{He})$	80	9.98	2.723	$0^+(3/2)$	20.49	19.01	-7.22
${}^9\text{B}({}^7\text{Li}, {}^6\text{Li})$	120	7.25	-2.723	$1^+(3/2)$	24.23	22.16	-8.54
${}^{12}\text{C}({}^8\text{B}, {}^7\text{Be})$	285	0.137	-12.690	$3/2^-(3/2)$	42.49	39.02	-8.17
${}^{12}\text{C}({}^9\text{C}, {}^8\text{B})$	78	1.3	-12.925	$2^+(3/2)$	40.14	36.86	-8.17
${}^{12}\text{C}({}^9\text{Li}, {}^8\text{Li})$	100	4.06	-9.882	$2^+(3/2)$	40.32	37.00	-8.23
${}^9\text{Be}({}^{10}\text{Be}, {}^9\text{Li})$	80	19.64	12.824	$3/2^-(3/2)$	35.33	32.47	-8.09
${}^9\text{Be}({}^{10}\text{Be}, {}^9\text{Be})$	120	6.812	-12.824	$3/2^-(3/2)$	77.62	70.68	-8.94
${}^9\text{Be}({}^{10}\text{C}, {}^9\text{C})$	120	21.28	17.277	$3/2^-(3/2)$	44.57	40.50	-9.31
${}^{12}\text{C}({}^{12}\text{C}, {}^{11}\text{B})$	250	15.95	-2.764	$3/2^-(3/2)$	64.68	58.70	-9.55
${}^{12}\text{C}({}^{12}\text{C}, {}^{11}\text{C})$	250	18.72	2.764	$3/2^-(3/2)$	74.16	67.10	-9.52
${}^{12}\text{C}({}^{14}\text{O}, {}^{13}\text{N})$	305	1.531	-18.552	$1/2^-(1/2)$	37.45	33.99	-9.22
${}^9\text{Be}({}^{14}\text{O}, {}^{13}\text{O})$	53	3.234	18.552	$3/2^-(3/2)$	25.57	23.32	-8.80
${}^{12}\text{C}({}^{16}\text{O}, {}^{15}\text{N})$	2100	22.04	-3.537	$3/2^-(3/2)$	46.90	42.48	-9.42
${}^{12}\text{C}({}^{16}\text{O}, {}^{15}\text{O})$	2100	22.04	3.537	$3/2^-(3/2)$	44.46	40.26	-9.45

$$(1 - |S_n|^2)|S_c|^2 \rightarrow (1 - |S_n|^2)|S_c|^2 \left(1 - \langle |S_{nc}|^2 \rangle\right),$$

with

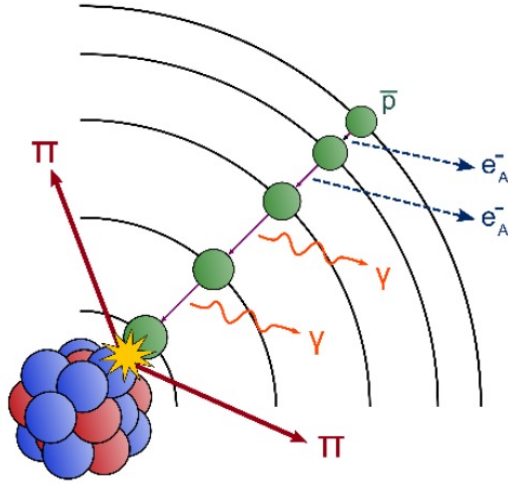
$$\langle |S_{nc}|^2 \rangle = \frac{1}{\sigma_{NN}^{el}} \int d\Omega \frac{d\sigma_{NN}^{el}(\theta)}{d\Omega} |S_{nc}(b_{nc}(\theta, \phi))|^2,$$

Bertulani, PLB (2023)

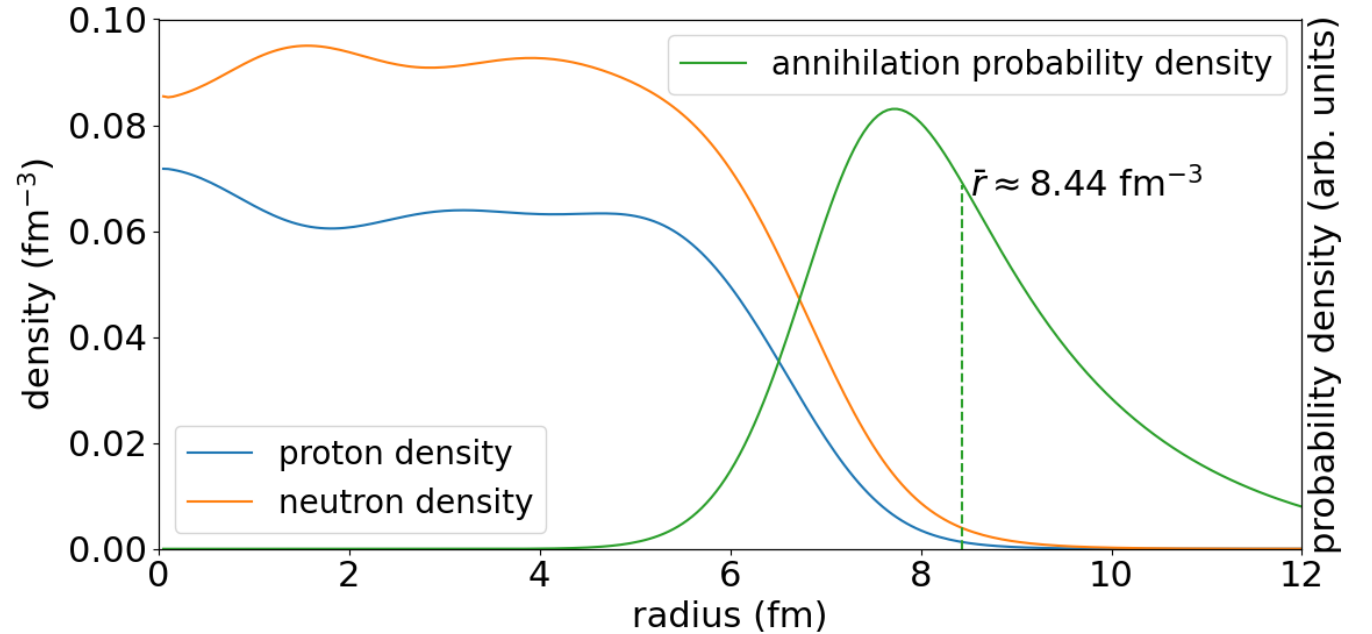




# STRIPPING FROM ANNIHILATION



- $\bar{p}$  captured in antiprotonic orbital ( $\sim$ QED)
- then annihilate in tail  $\rho_{n/p}(r)$  (QCD)

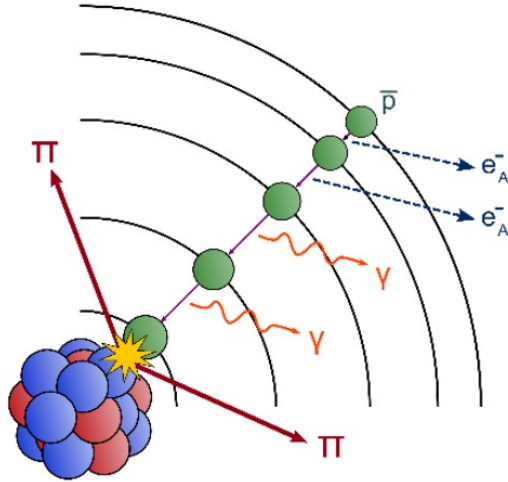


Two main methods applied in the past (LEAR at CERN) for nuclear structure:

- X rays from lower orbitals
- pions from annihilation

The interpretation of both methods depends on  $\bar{p}$  - nucleus optical potential

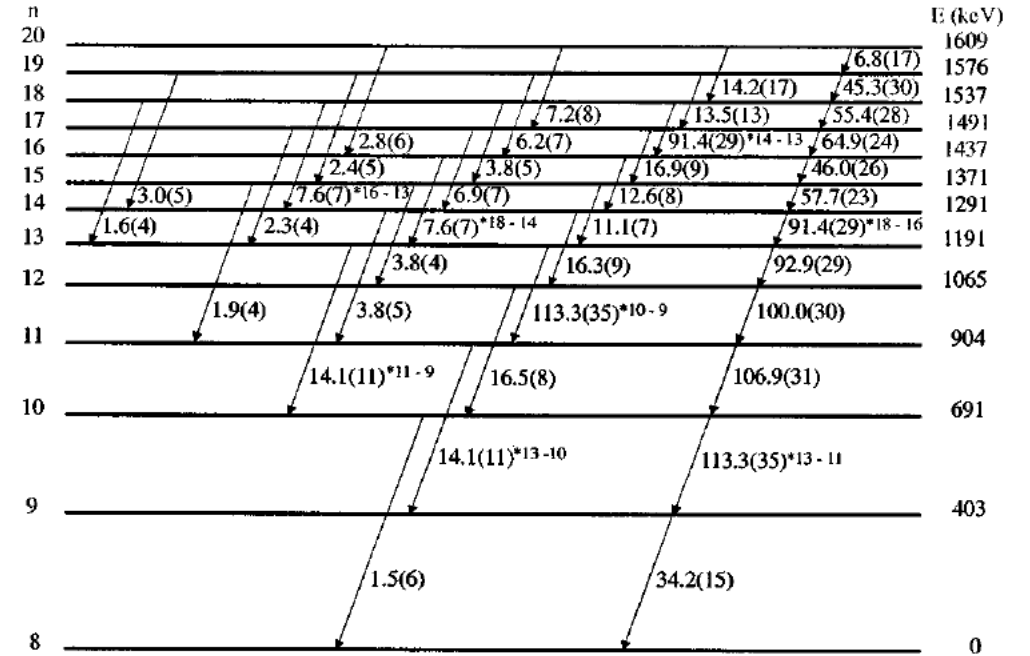
# STRIPPING FROM ANNIHILATION



- $\bar{p}$  captured in antiprotonic orbital ( $\sim$ QED)
- Energy shifts and decay widths of inner levels sensitive to the matter distribution

$$\Gamma_{ne} = \int \text{Im } V(r) |\Psi_{ne}(r)|^2 r^2 dr$$

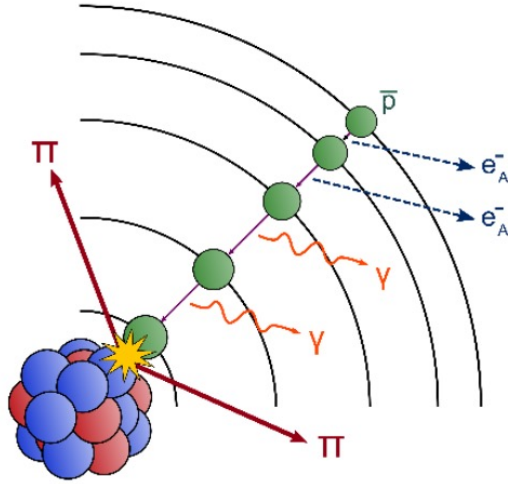
with  $V(r) = \frac{2\pi}{\mu} a\rho(r)$



172Y @ CERN, R. Schmidt et al., PRC (1998)

Nuclear potential induces shift of inner level energies of few 100 eV (X ray energies: 100 keV)

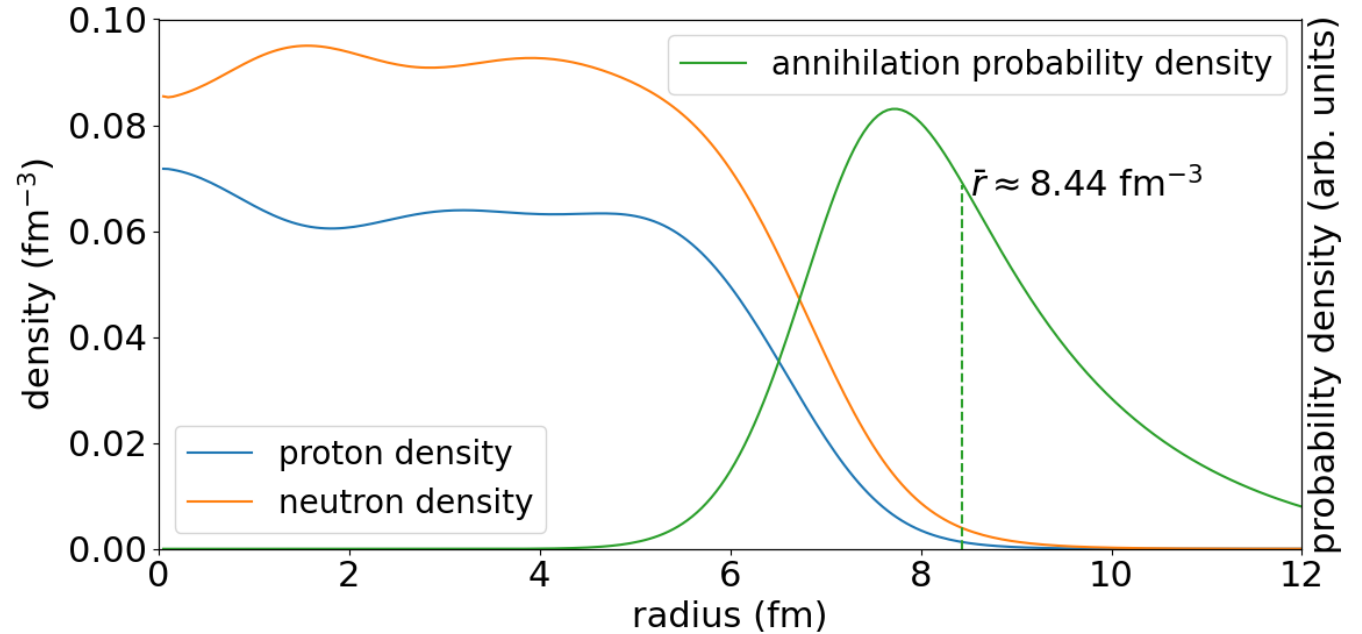
# STRIPPING FROM ANNIHILATION



- $\bar{p}$  captured in antiprotonic orbital ( $\sim$ QED)
- then annihilate in tail  $\rho_{n/p}(r)$  (QCD)
- Conservation of total charge

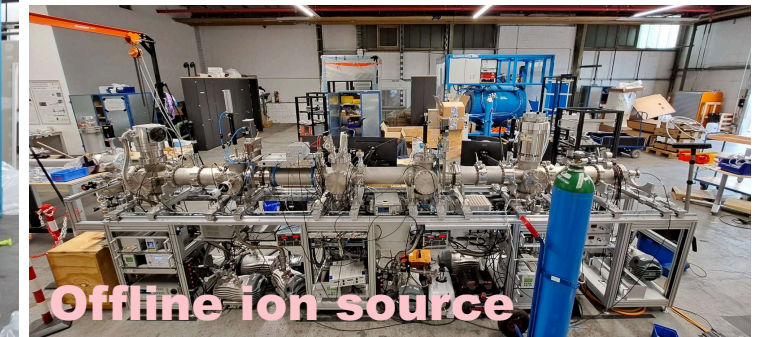
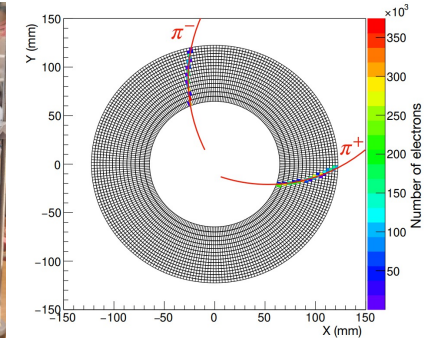
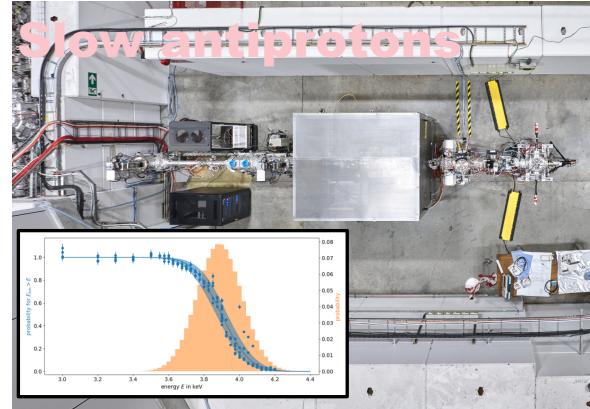
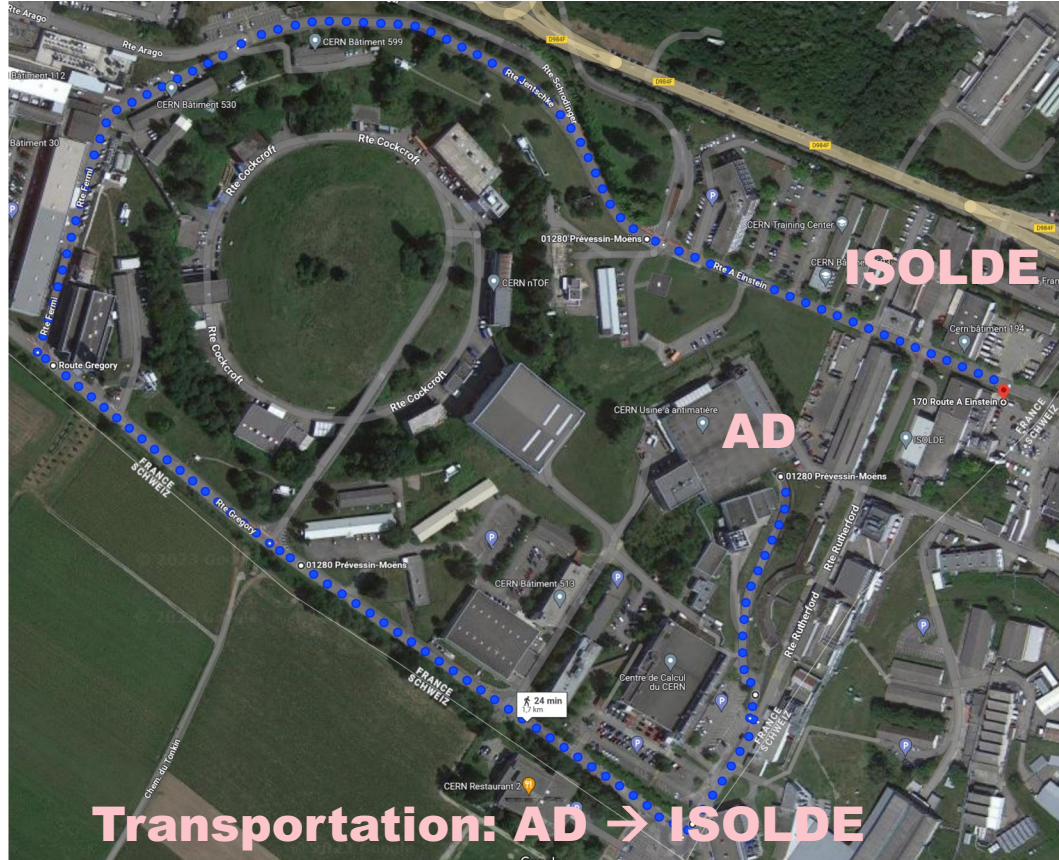
$$\sum_{\pi} q_{\pi} = \begin{cases} 0 & \text{for } \bar{p}p \\ -1 & \text{for } \bar{p}n \end{cases}$$

- Collaboration with Lazauskas, Hupin et al. (theory)
- See talks by: **Duerinck and Finelli**



- First application of method: Bugg et al., PRL (1973)
- Application to rare isotopes first proposed by: Wada and Yamazaki, NIM B (2004)
- First experiments at PUMA, CERN, in 2025

# THE PUMA EXPERIMENT



Aumann et al. (PUMA collaboration), EPJA (2022)



# SUMMARY

- Direct reactions are unique to probe the wavefunction of nuclear states.
- Optical potentials are major components for the analysis and interpretation of direct reaction cross sections  
Accuracy and precision of 10% or less for cross sections would allow quantitative statements
- In **transfer reactions**, uncertainties from isospin dependence and separation energy are lacking  
Variations of cross sections of  $\sim O(10)$ - $O(100)$  % due to optical potentials
- **Nucleon quasifree scattering**: optical potentials lead to variations of cross sections of  $\sim O(10)$ %  
Cluster knockout (deuteron, ... not discussed in this talk) is expected to be more dependent on the OP  
Systematic uncertainties from reaction formalism and related approximations are large ( $> 20\%$ )
- Inclusive stripping reactions at or below  $\sim 150$  MeV / nucleon may deviate significantly from eikonal and favour larger uncertainties than at higher energies (Pauli blocking, core excitation from FSI, transfer)
- **Low-energy antiprotons** sensitive to nuclear surface; soon applied to RI at CERN

# SCOPE AND AIMS

- Optical potentials from ab initio theory, Dispersive optical model (DOM) and other microscopic methods
- Consistency among structure and reaction mechanisms
- Uncertainties
- Implementation of microscopic nucleon-nucleus OP in actual reactions
- Applications
  - structure of nuclei away from stability investigated at facilities at various energy regimes
  - antiproton-nucleus optical potentials
  - deformed nuclei
  - astrophysics