

# Future prospects of Lambda-proton scattering experiment

Koji Miwa (Tohoku Univ., KEK IPNS, RIKEN)  
on behalf of J-PARC E40, E86, HYPS collaboration

Workshop on Strange hadron as a Precision tool for strongly interacting systems  
May. 13<sup>rd</sup> – 17<sup>th</sup>, 2024



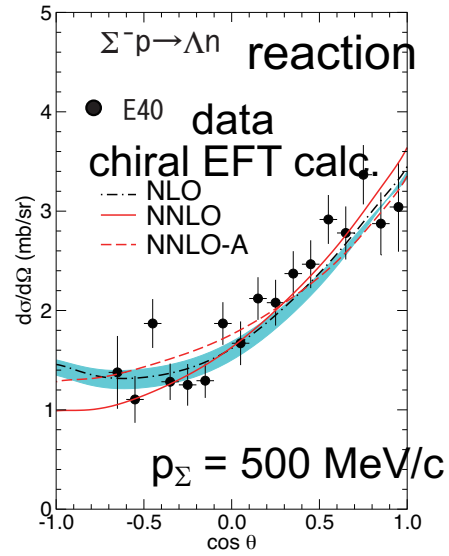
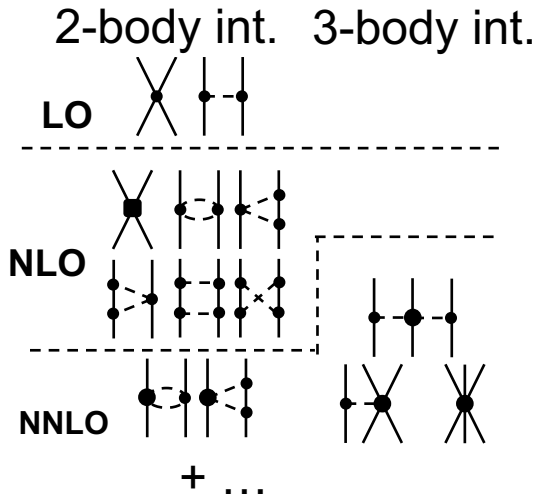
# Contents

- Brief summary of  $\Sigma p$  scattering experiment (J-PARC E40)
- New project of  $\Lambda p$  scattering experiment at SPring-8
- Summary

# Progress of theory & experiment of BB int. study

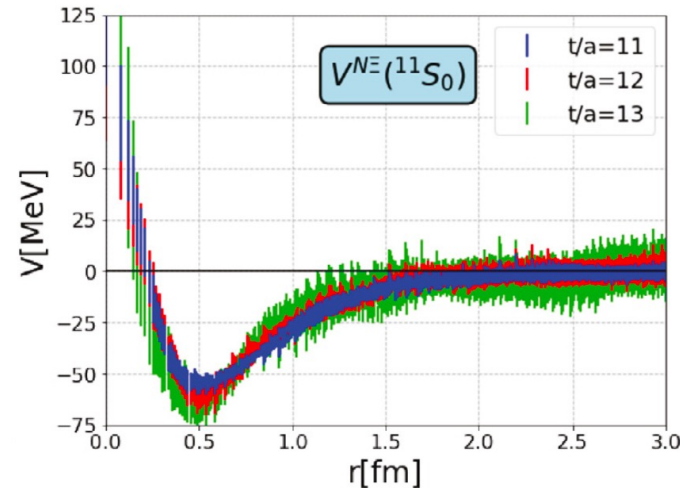
## Theoretical progress

### Hyperon-Nucleon int. w/ chiral effective field theory (J. Haidenbauer et al.)



### Hyperon potential by Lattice QCD

BB interaction at almost physical point for multi-strangeness sector



K. Sasaki et al., Nucl. Phys. A 998 (2020) 121737

## Improving accuracy w/ our new data

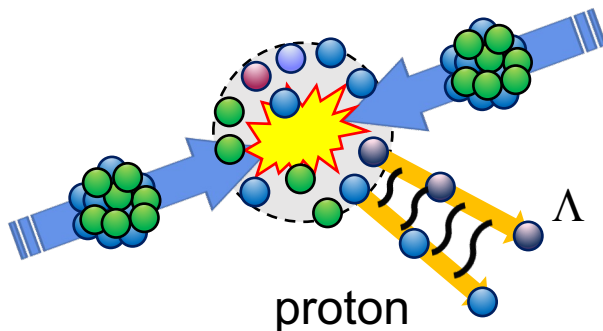
## Experimental progress

### BB interaction from femtoscopy

$$c(k^*) = \int S(r^*) |\Psi(\vec{k}^*, \vec{r}^*)|^2 d^3 r^*$$

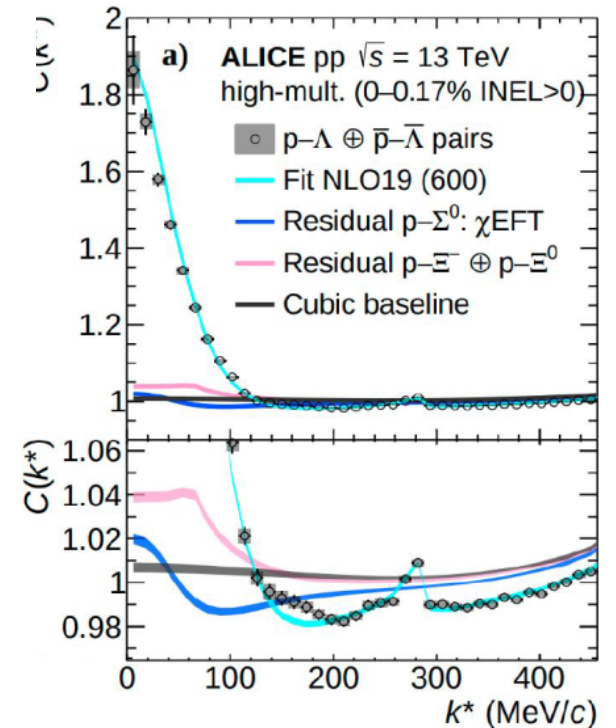
Fix source size ( $S(r^*)$ ) →

Study interaction from wave function ( $\Psi(\vec{k}^*, \vec{r}^*)$ )



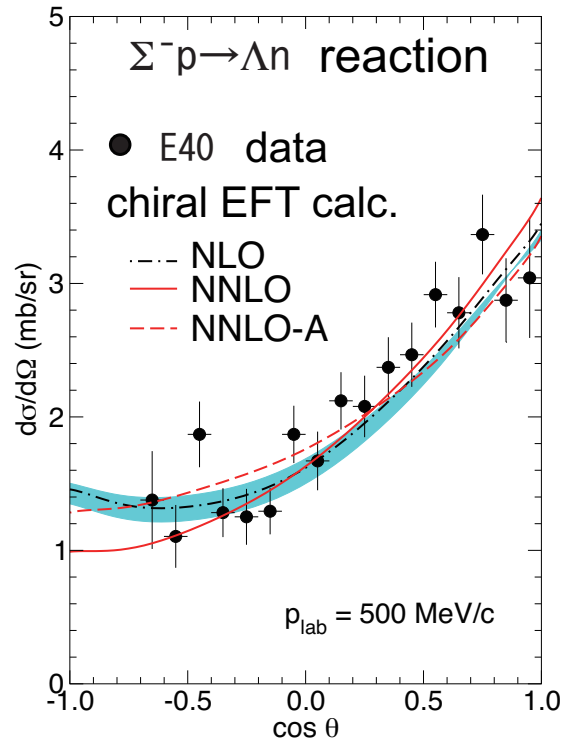
ALICE Collaboration, Phys. Lett. B 833 (2022) 137272

### Particle correlation between $\Lambda$ and p

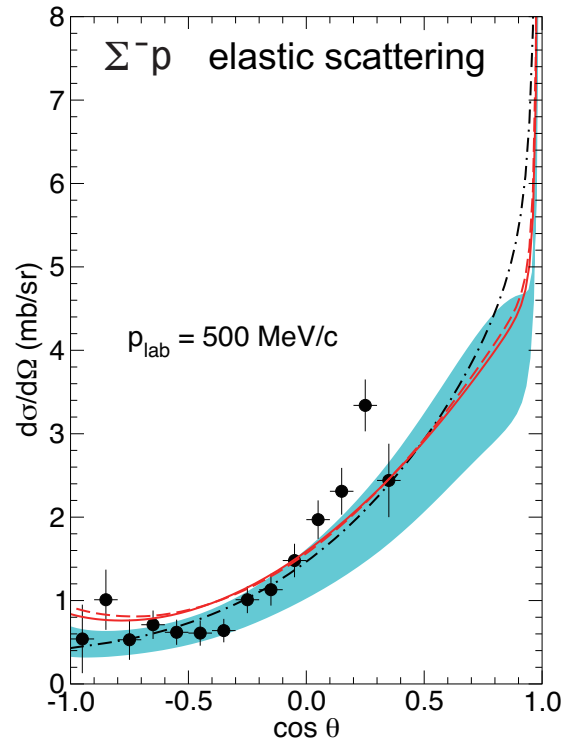


# New $\Sigma p$ scattering data at J-PARC

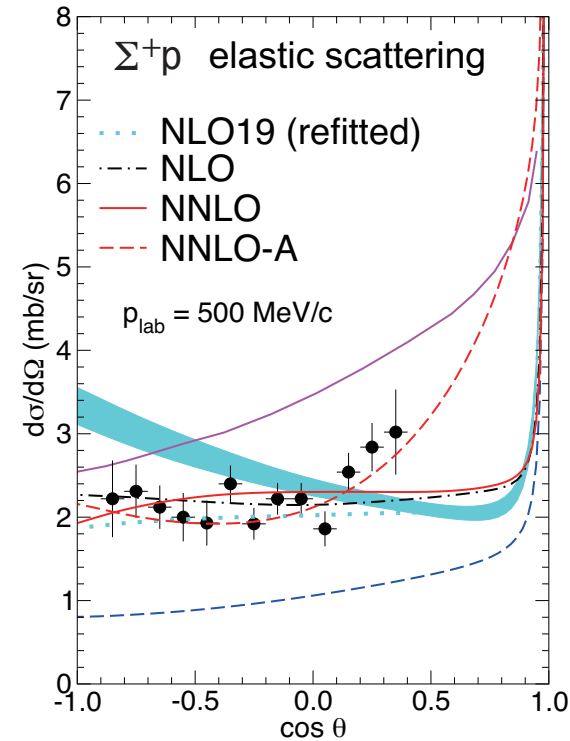
Accurate and systematic data of differential cross sections of  $\Sigma p$  scattering



K. Miwa et al.,  
PRL 128, 072501 (2022)



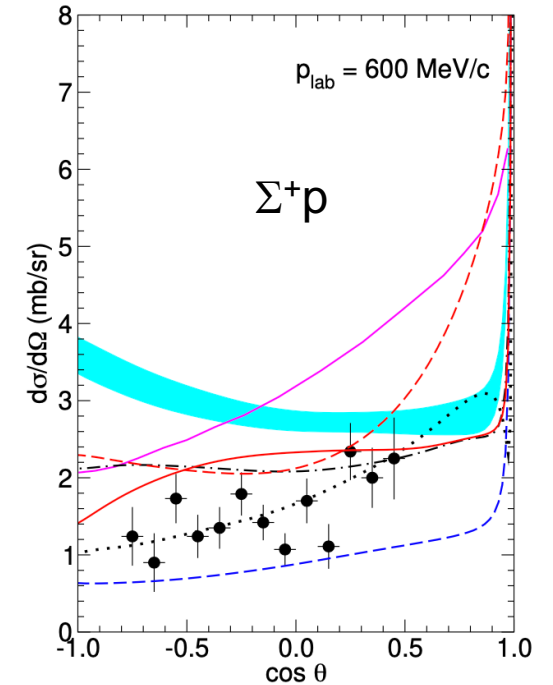
K. Miwa et al.,  
PRC 104, 045204 (2021)



T. Nanamura et al., PTEP 2022 093D01

Difficulty at higher momentum

$\Sigma^+ p \rightarrow \Sigma^+ p$



J. Haidenbauer et al.,  
Eur.Phys.J.A 59 (2023) 3

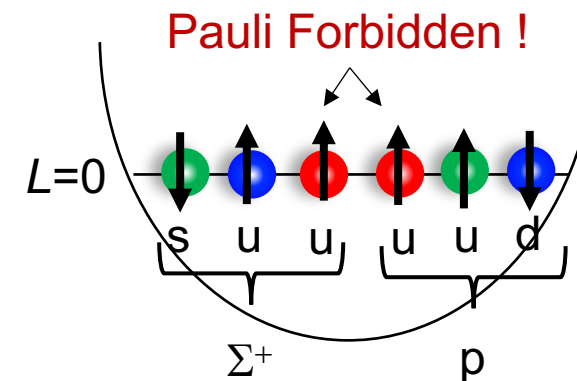
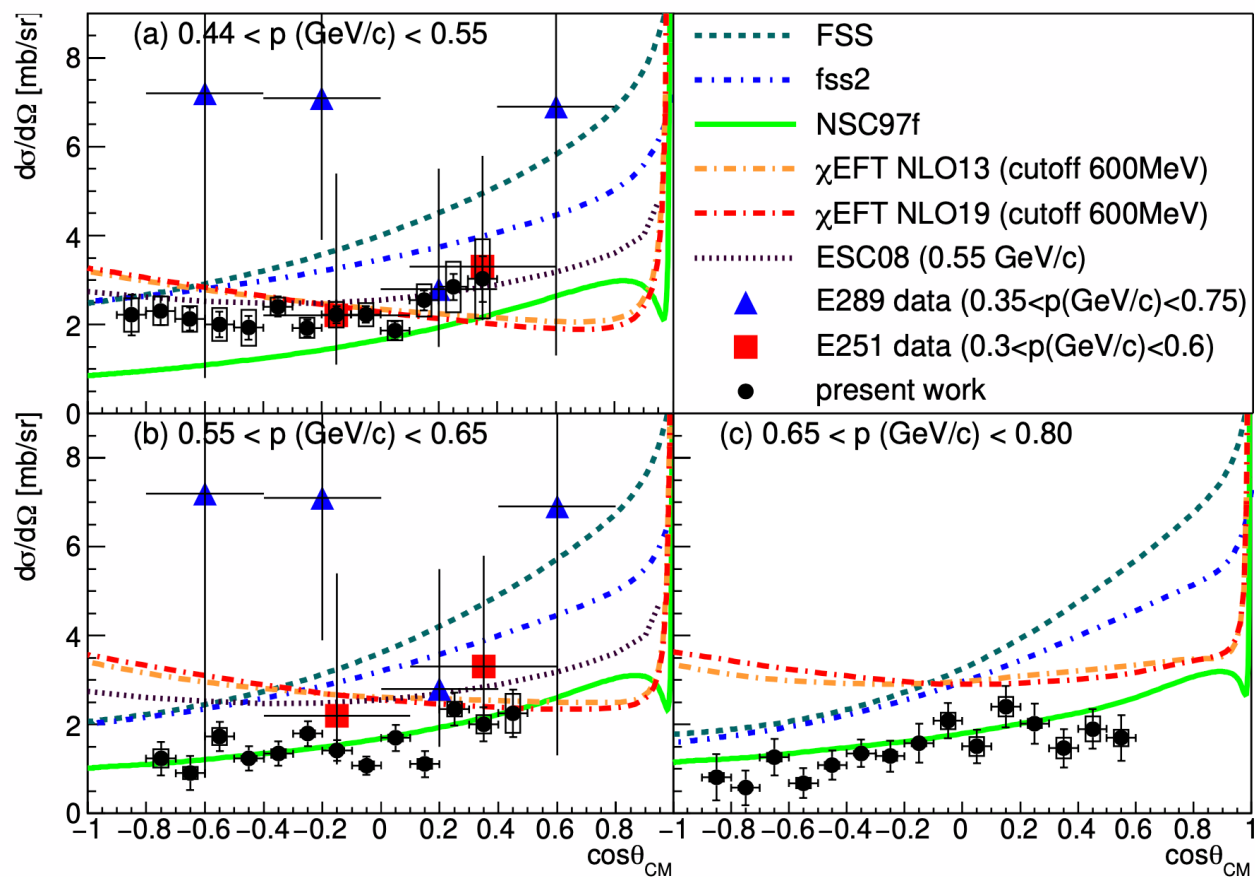
Development of Chiral EFT at NNLO have got started with E40 data

But, the interactions are not uniquely determined yet.

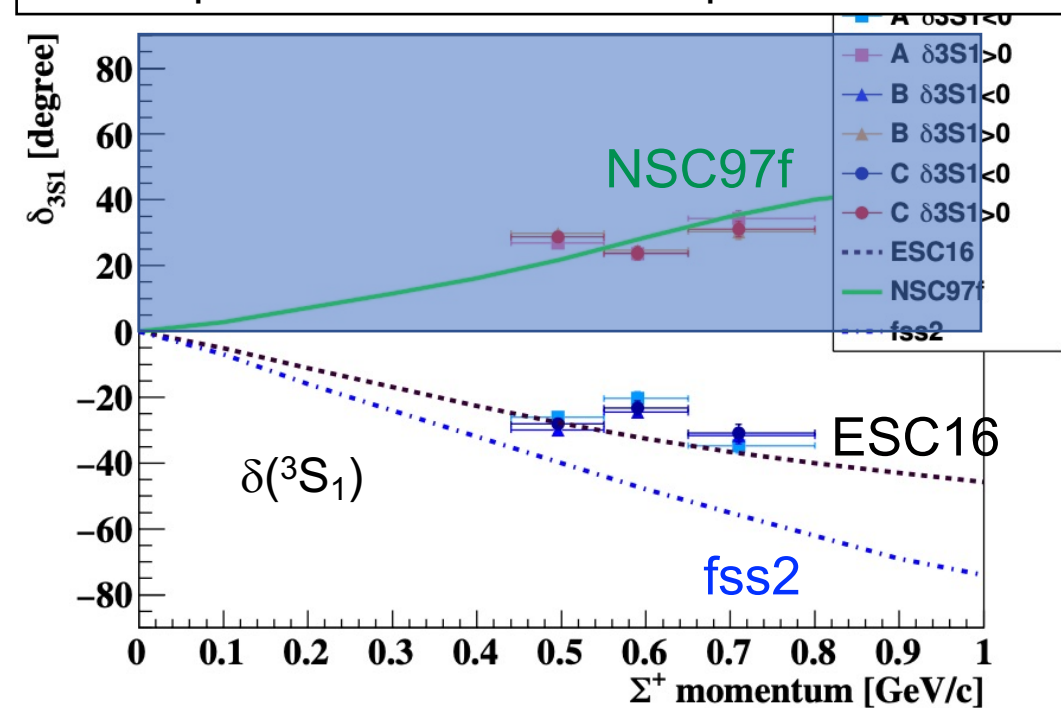
We need more data from additional channels ( $\Lambda p$ , ...) and additional differential observables (polarizations, ...)

# $d\sigma/d\Omega$ of $\Sigma^+p$ elastic scattering

T. Nanamura et al., Prog. Theor. Exp. Phys. **2022** 093D01



First experimental derivation of phase shift of  $^3S_1$

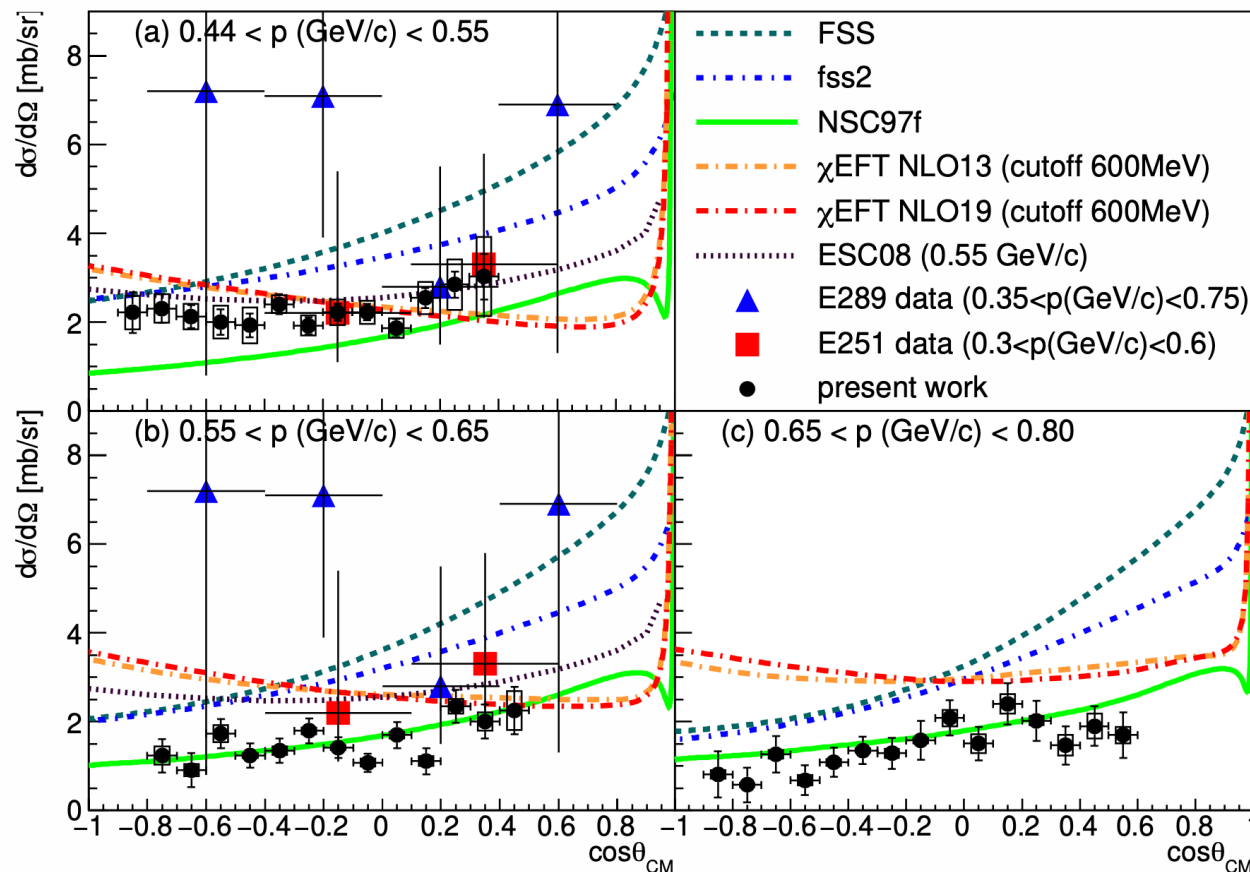


**E40 data : much smaller than fss2 prediction and E289 results**

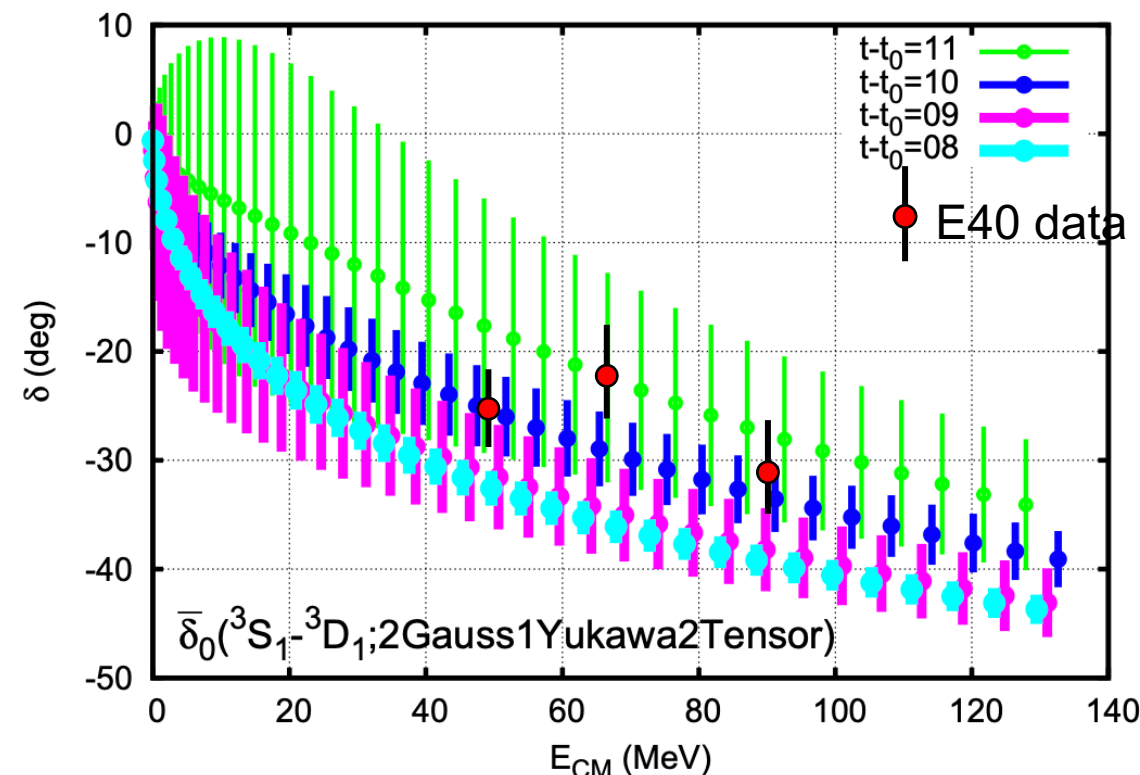
**Derived phase shift suggests that the  $^3S_1$  interaction is moderately repulsive.**

# $d\sigma/d\Omega$ of $\Sigma^+p$ elastic scattering

T. Nanamura et al., Prog. Theor. Exp. Phys. **2022** 093D01



Comparison with HAL QCD  $\Sigma N$  potential



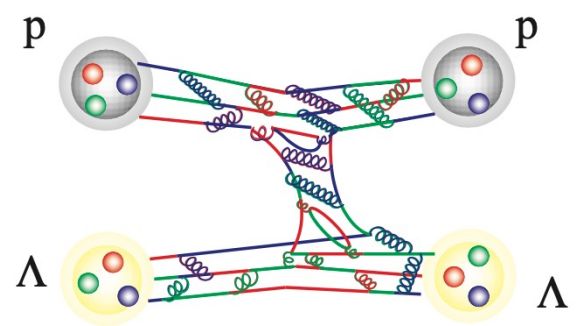
H. Nemura et al., EPJ Web of Conf., 175, 05030 (2018)

**E40 data : much smaller than fss2 prediction and E289 results**

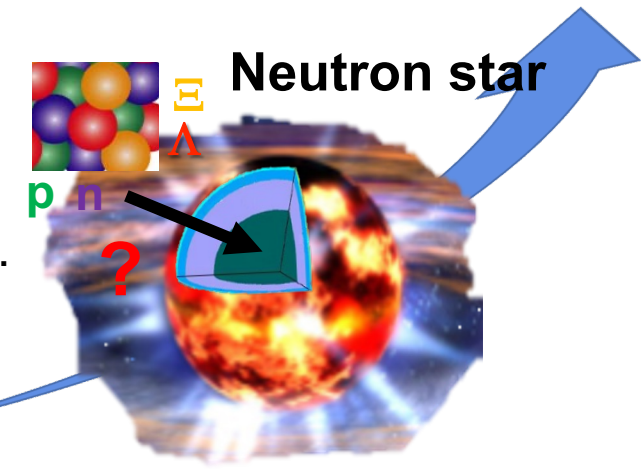
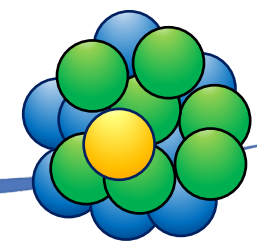
**Derived phase shift suggest that the  $^3S_1$  interaction is moderately repulsive.**

# Toward $\Lambda p$ scattering

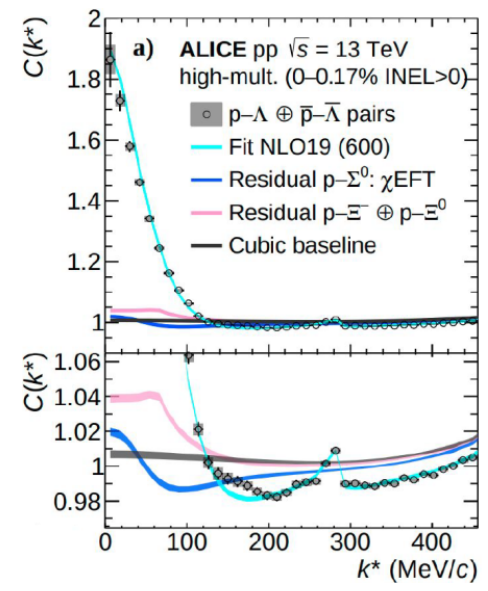
**Reliable  $\Lambda N$  two-body interaction :**  
key to deepen  $\Lambda$  hypernuclear physics



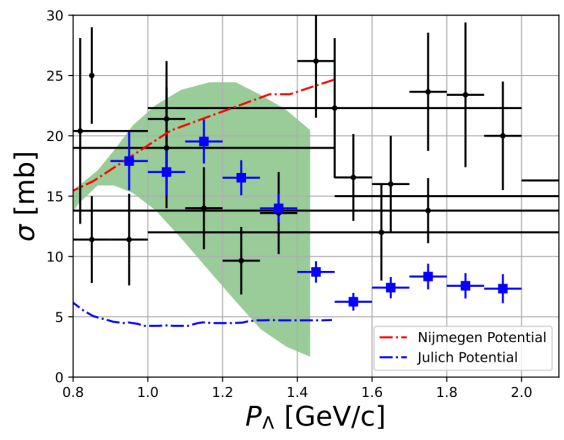
**$\Lambda$  hypernuclei**  
key to reveal  $\Lambda NN$  int.



## Femtoscscopy from HIC

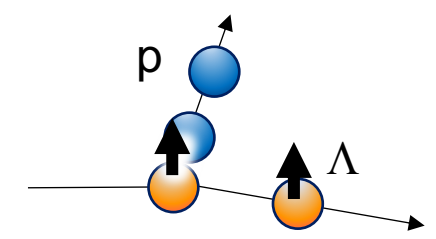


## New cross section data from Jlab CLAS



**New project at SPring-8, J-PARC**

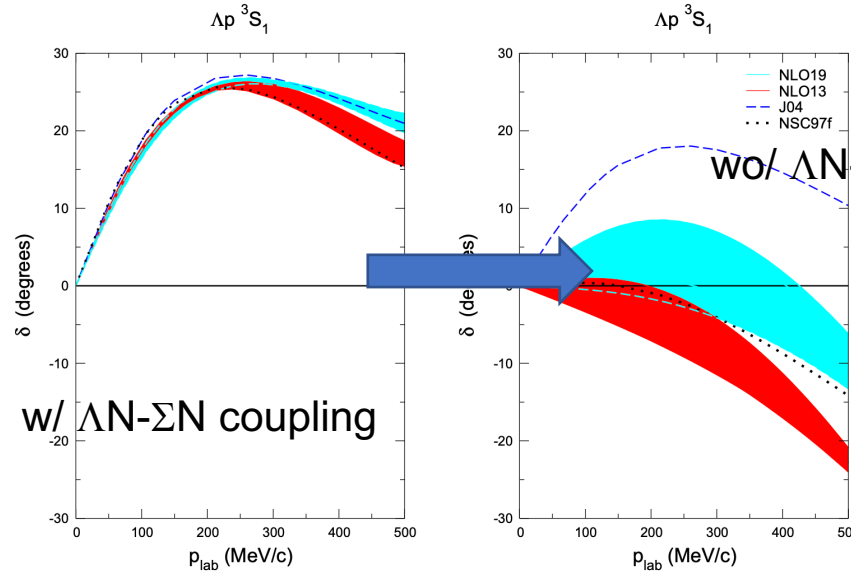
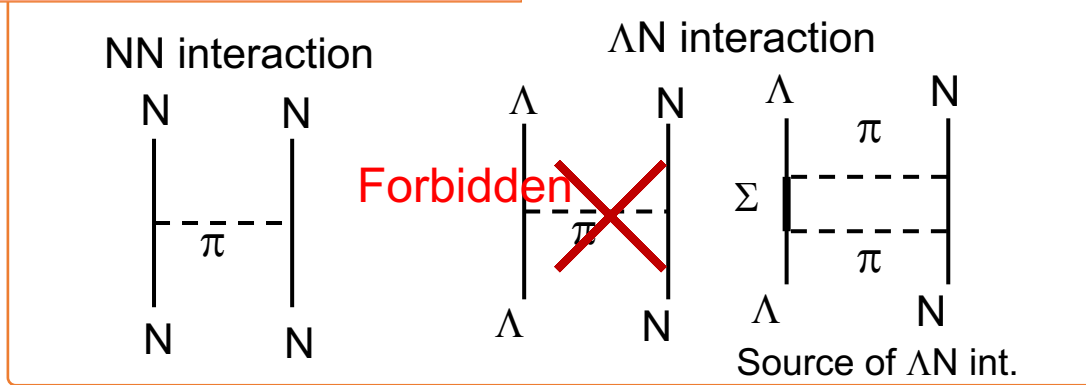
**$\Lambda p$  scattering w/ (polarized)  $\Lambda$**



# Origin of the density dependence of $\Lambda N$ interaction

$\Lambda$  is the only isospin 0 baryon that makes up matter  $\rightarrow$  One-pion exchange is forbidden

## Interaction in free space



$\Lambda N$  interaction becomes less attractive (repulsive) w/o  $\Lambda N$ - $\Sigma N$  coupling

## Interaction in nuclear medium

Two-body NN int. might not be affected very much by other nucleons

**Suppression of two-body  $\Lambda N$  int.**

$\rightarrow$  Less attractive

Due to the Pauli blocking for the intermediate nucleon,  $\Lambda N$ - $\Sigma N$  coupling can be very suppressed

**3 body int. mediated with  $\Sigma$**

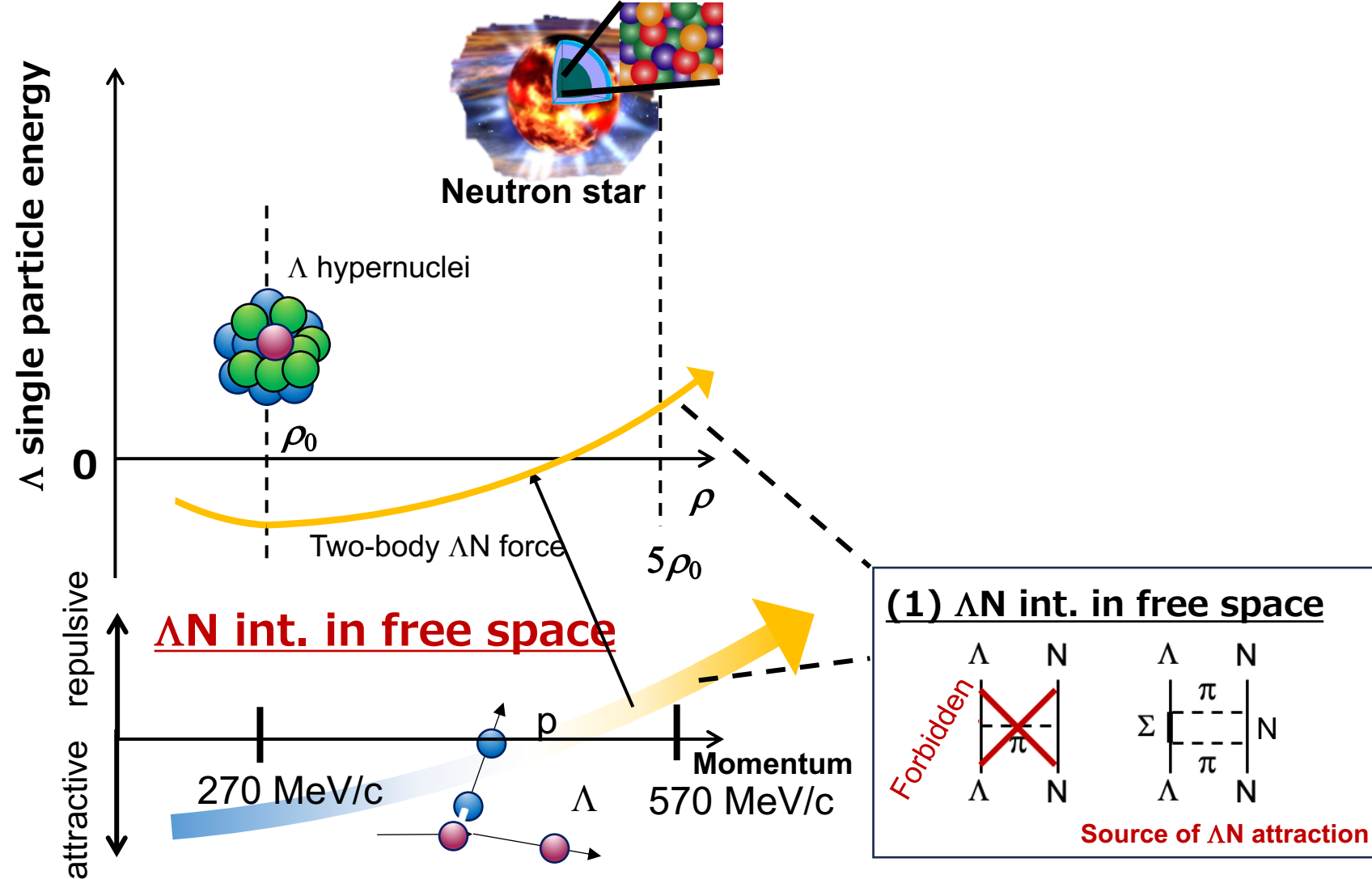
Additional three-body interaction, which does not appear in NN interaction.  
What is the density dependence of this 3-body int. ?



# Toward the elucidation of the density dependence of the $\Lambda N$ interaction

Ingredients of  $\Lambda N$  interaction  
in nuclear medium

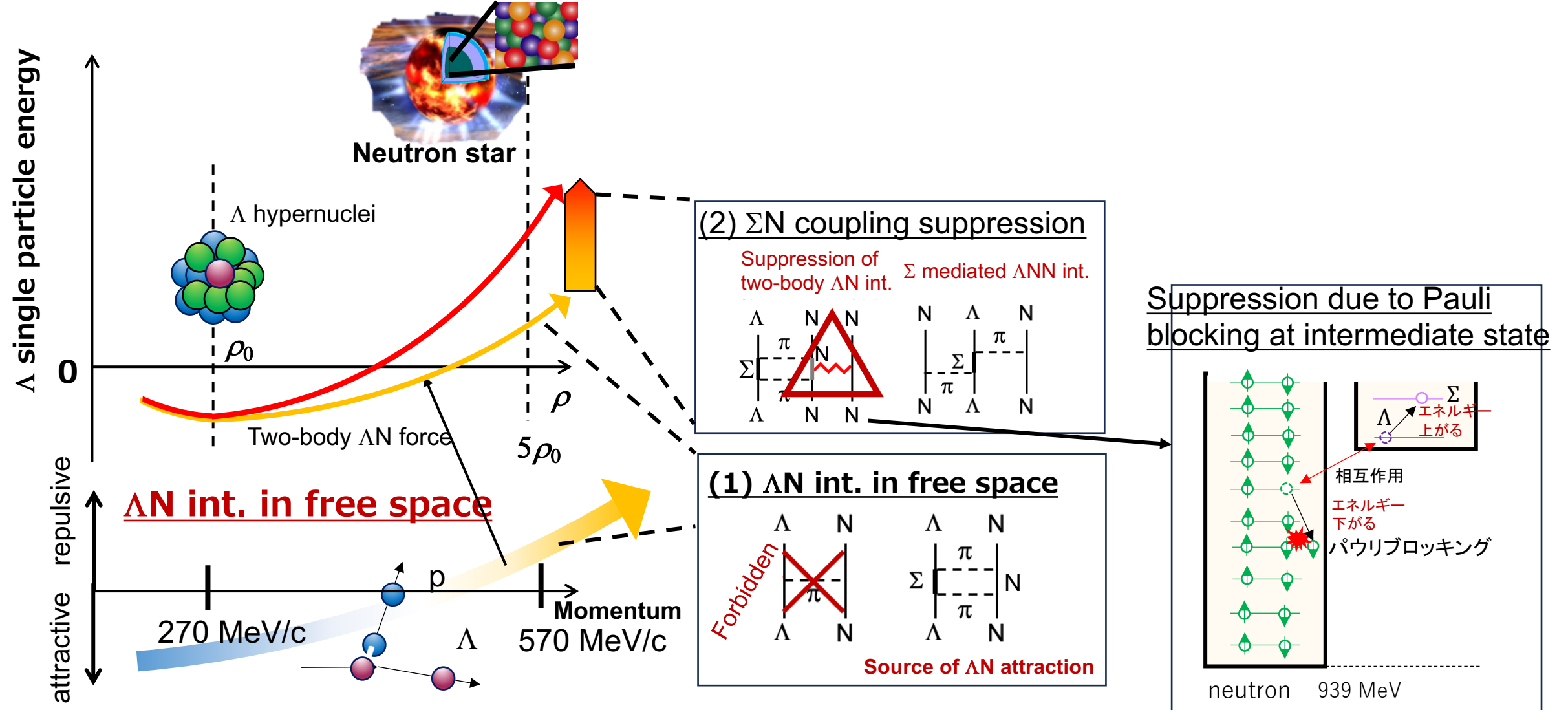
## $\Lambda N$ int. in nuclear medium



# Toward the elucidation of the density dependence of the $\Lambda N$ interaction

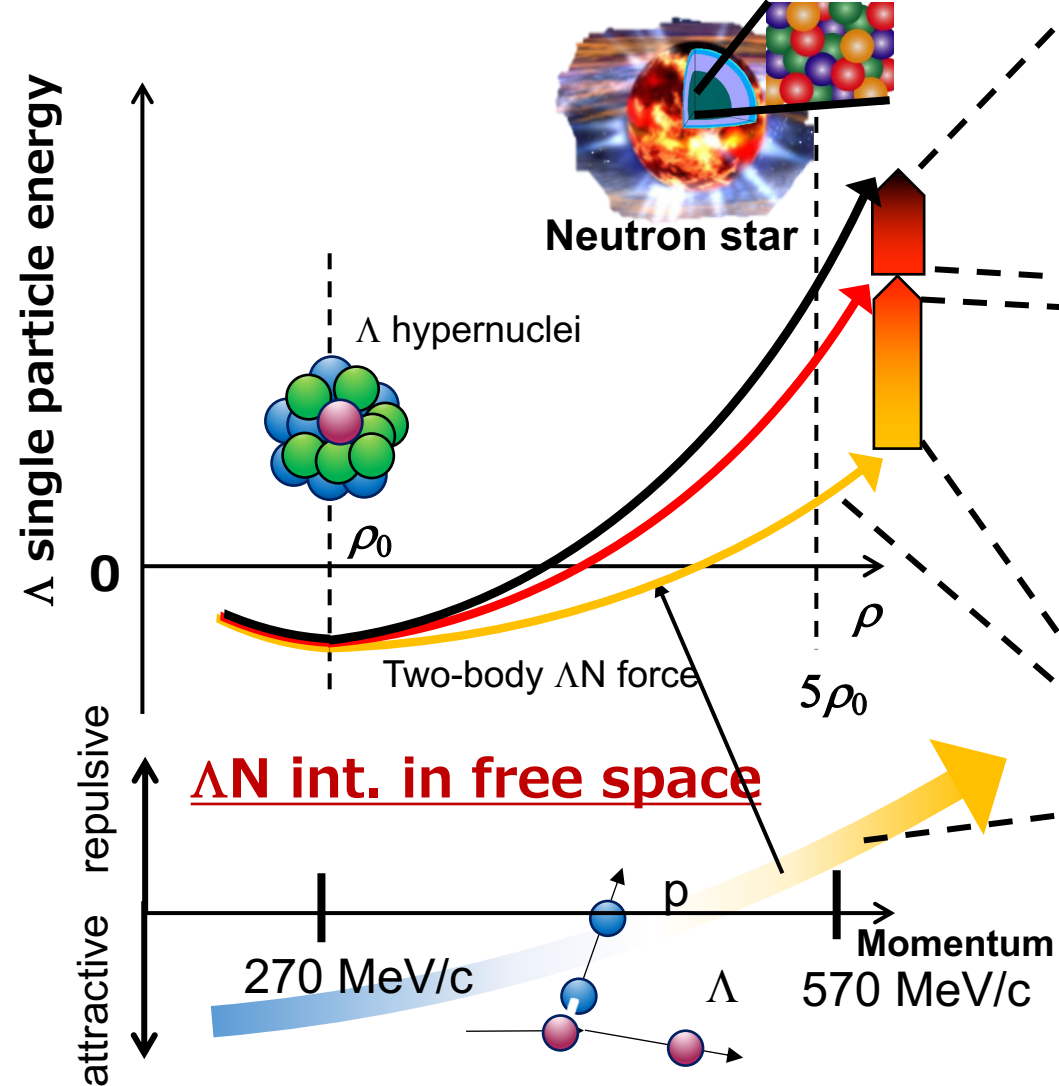
Ingredients of  $\Lambda N$  interaction  
in nuclear medium

## $\Lambda N$ int. in nuclear medium



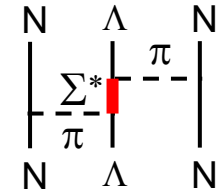
# Toward the elucidation of the density dependence of the $\Lambda N$ interaction

## $\Lambda N$ int. in nuclear medium



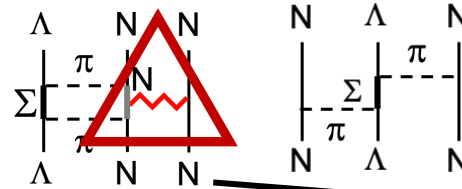
Ingredients of  $\Lambda N$  interaction in nuclear medium

$\Lambda NN$  int. mediated with  $\Sigma^*$

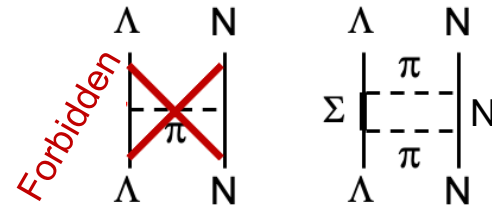


(2)  $\Sigma N$  coupling suppression

Suppression of two-body  $\Lambda N$  int.  $\Sigma$  mediated  $\Lambda NN$  int.

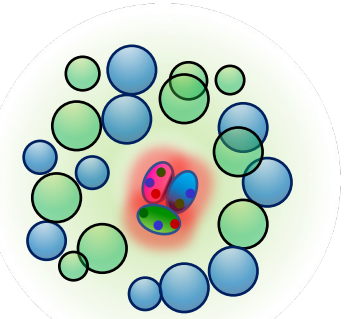


(1)  $\Lambda N$  int. in free space

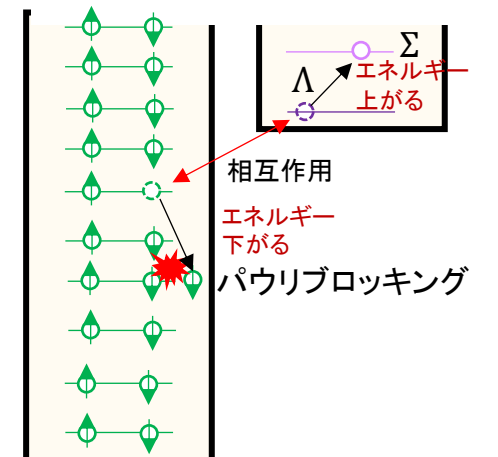


Source of  $\Lambda N$  attraction

Repulsive force by  $\Lambda NN$  force

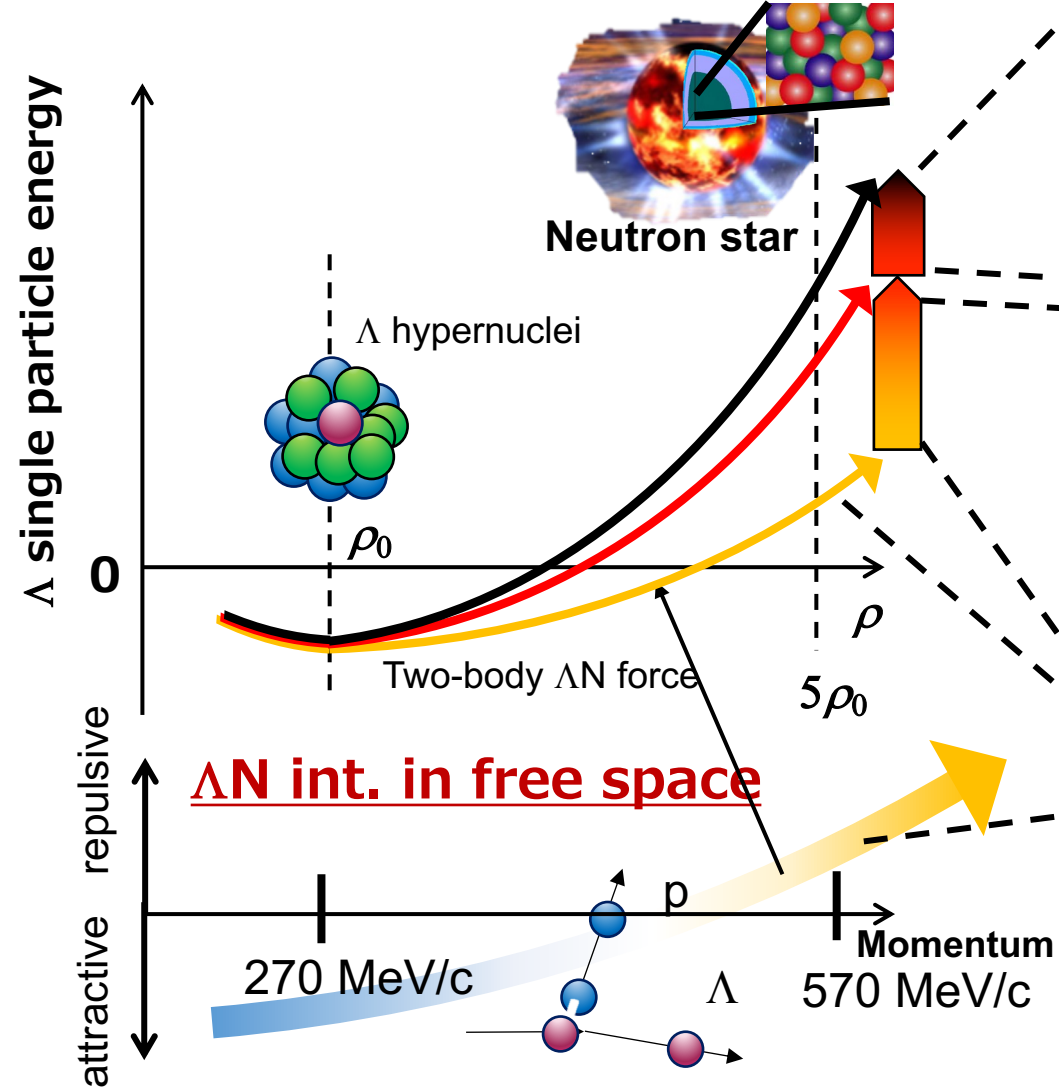


Suppression due to Pauli blocking at intermediate state

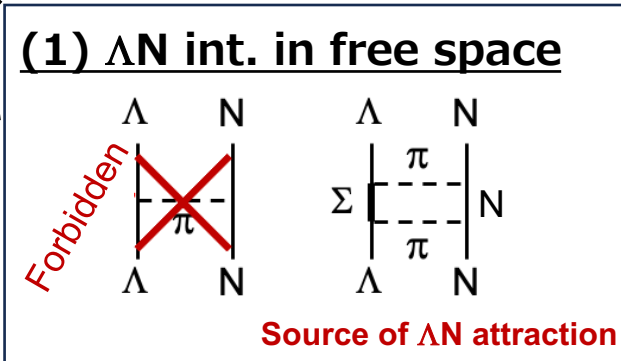
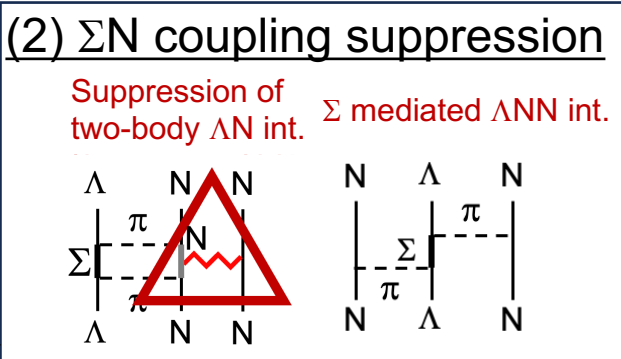
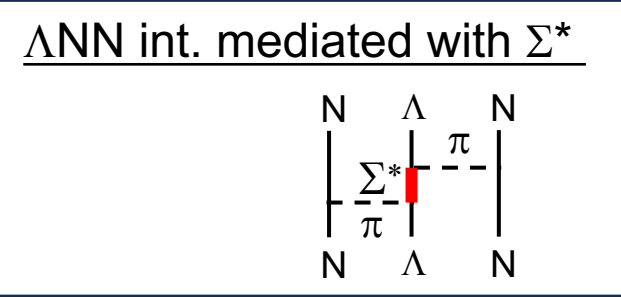


# Toward the elucidation of the density dependence of the $\Lambda N$ interaction

## $\Lambda N$ int. in nuclear medium



Ingredients of  $\Lambda N$  interaction in nuclear medium



Determination of  $\Sigma^*$  mediated  $\Lambda NN$  force

Precise  $\Lambda$  hypernuclear spectroscopy

Determination of  $\Lambda N$ - $\Sigma N$  coupling

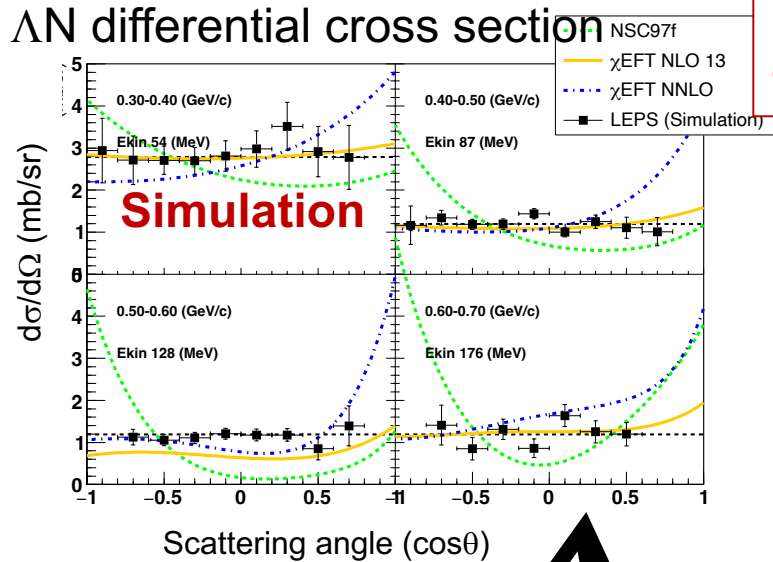
$\Sigma N$  cusp measurement  
Lattice QCD

Realistic  $\Lambda N$  interaction

$\Lambda N$  scattering data  
Lattice QCD

# Collaborative research regarding the two-body $\Lambda N$ , $\Sigma N$ int.

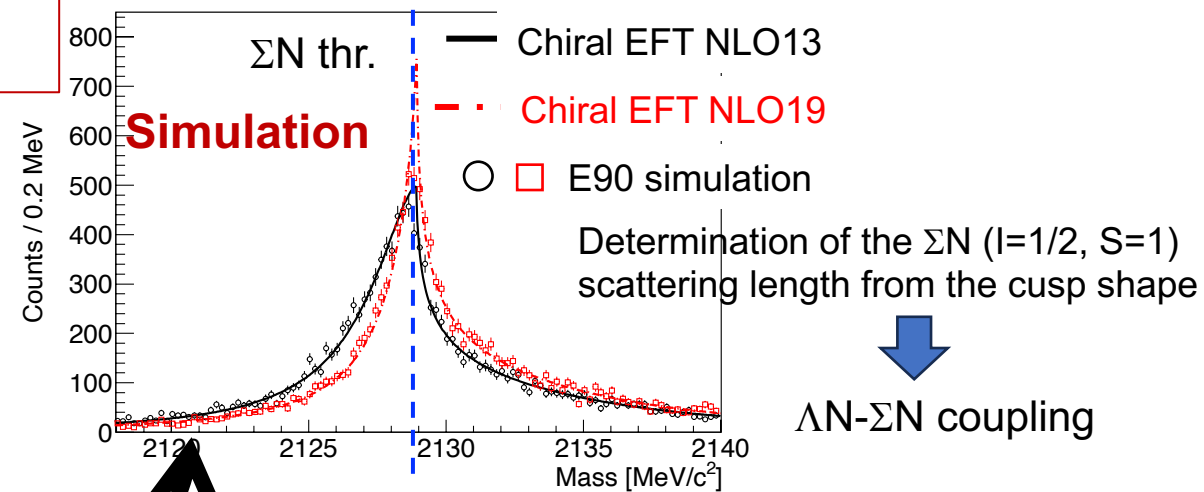
## (1) $\Lambda p$ scattering experiment (Koji Miwa)



The differential cross section and cusp structure must be represented with the same  $\Lambda N$ - $\Sigma N$  coupling

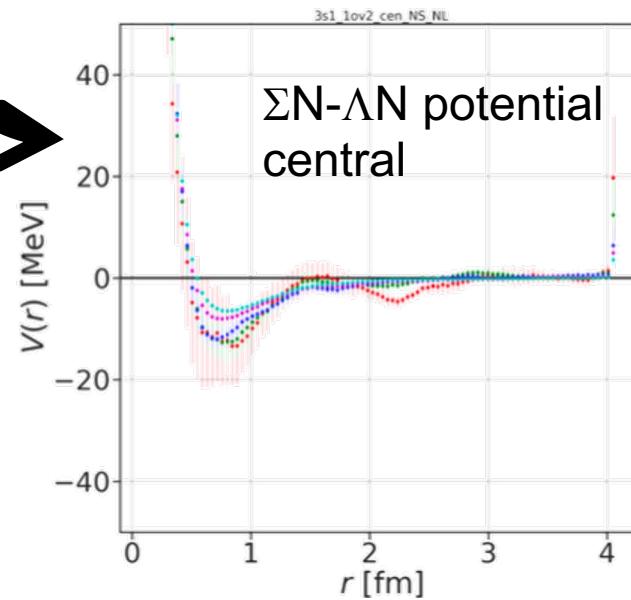
Momentum dependence of  $d\sigma/d\Omega$

## (2) $\Sigma N$ cusp measurement (Yudai Ichikawa)



Independent determinations of  $\Lambda N$ - $\Sigma N$  coupling through the scattering length measurement

## (3) $\Lambda N$ , $\Sigma N$ Lattice QCD potential (Takahiro Doi)



$\Lambda N$ ,  $\Sigma N$  and  $\Lambda N$ - $\Sigma N$  coupling potentials by HAL QCD

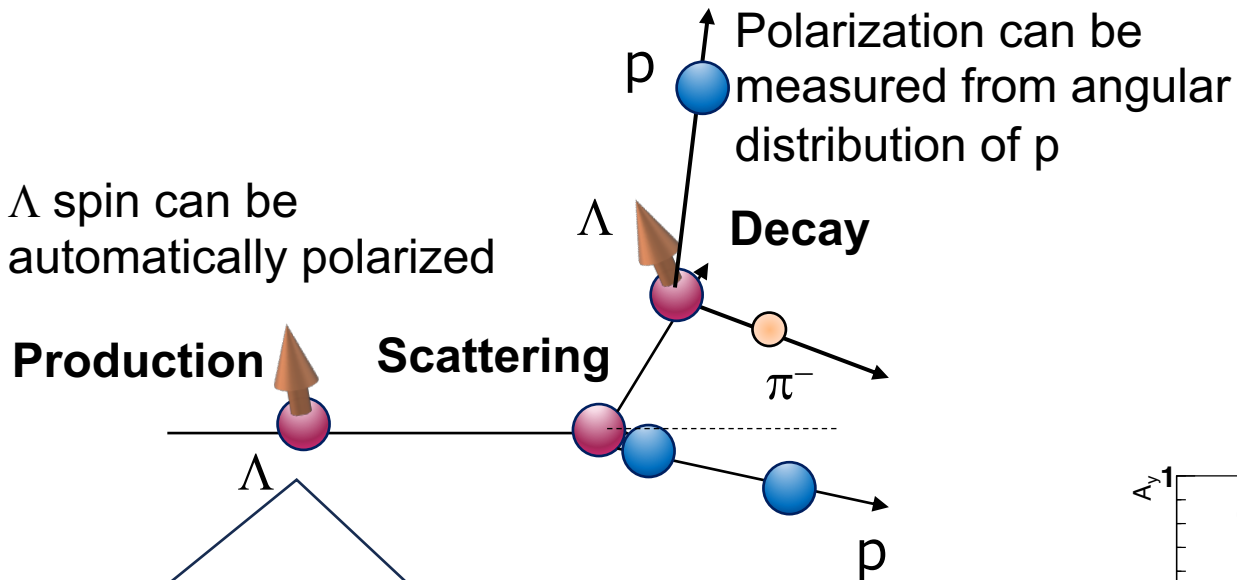
$\Lambda N$ - $\Sigma N$  coupling

[T. Doi, presentation at 3rd J-PARC HEF-ex WS](#)

Universal understanding of  $\Lambda N$  interaction from scattering and Lattice QCD

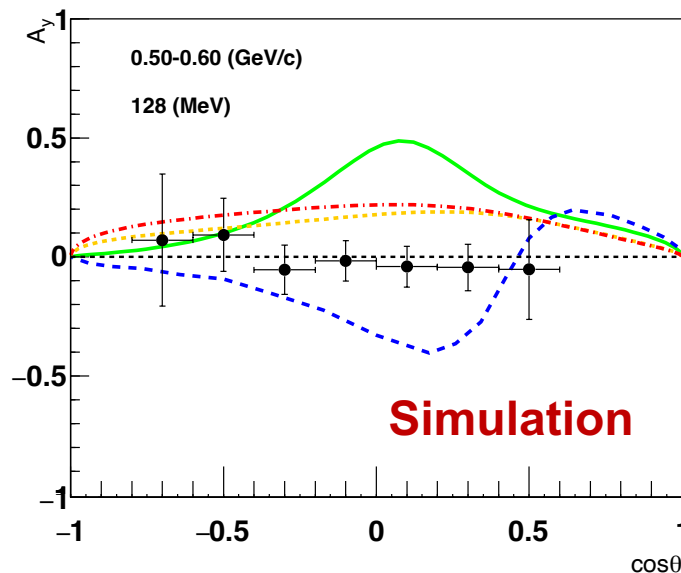
# $\Lambda p$ scattering experiment with polarized $\Lambda$ beam (J-PARC E86)

**Advantage of scattering experiment: Spin observables can be measured thanks to self polarimeter of hyperon**

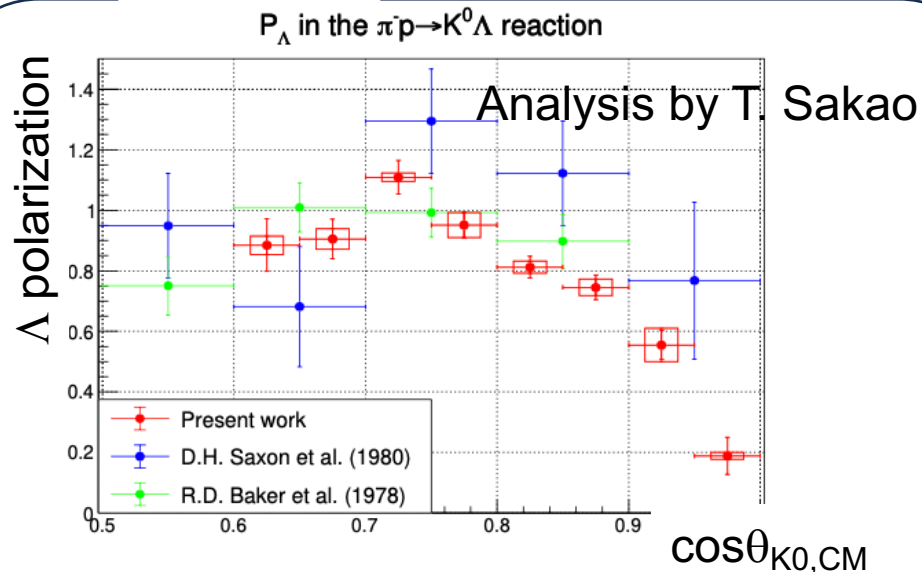
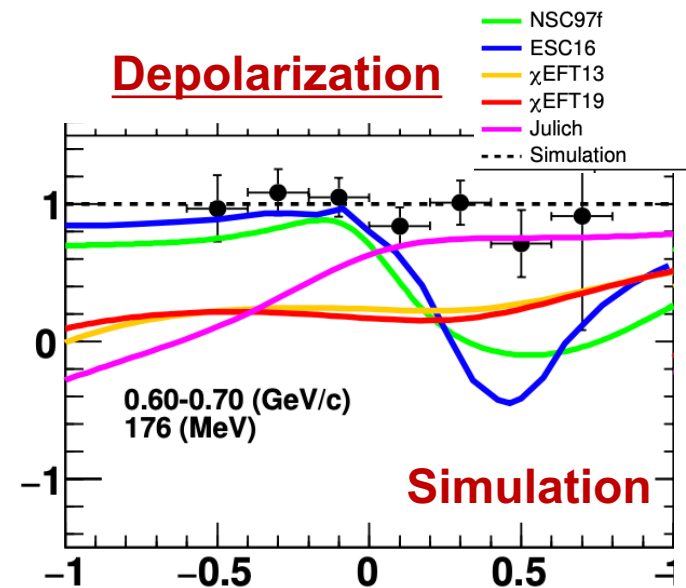


- Left-Right asymmetry of  $\Lambda p$  scattering (Analyzing power)
  - spin-orbit interaction
- Polarization change before and after the scattering (Depolarization)
  - spin-spin interaction, tensor interaction

## Analyzing power



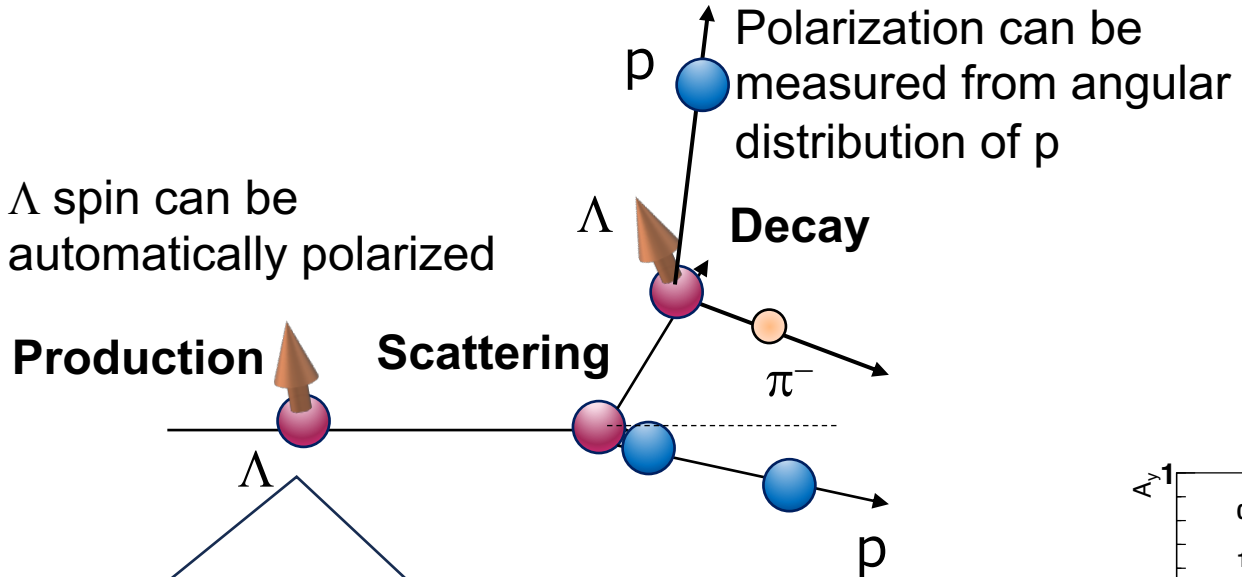
## Depolarization



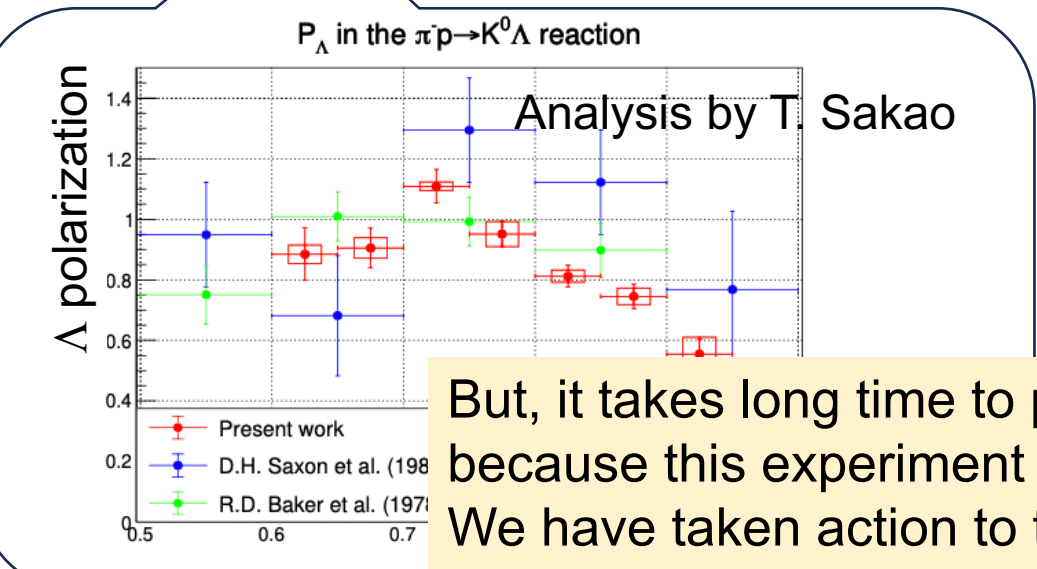
Essential constraint to determine spin-dependent  $\Lambda N$  interaction

# $\Lambda p$ scattering experiment with polarized $\Lambda$ beam (J-PARC E86)

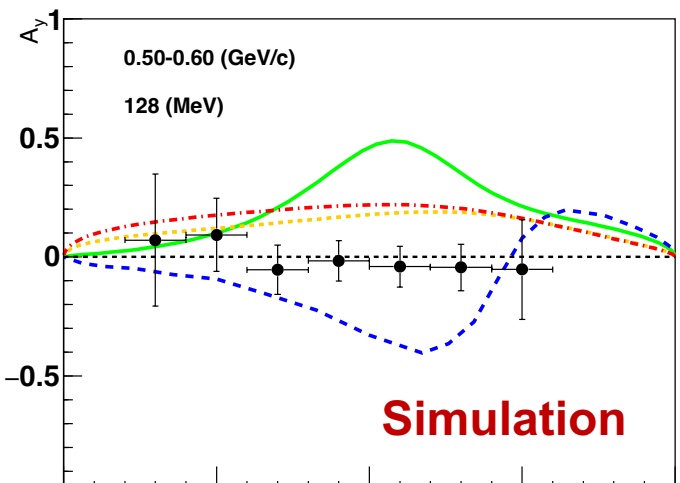
**Advantage of scattering experiment: Spin observables can be measured thanks to self polarimeter of hyperon**



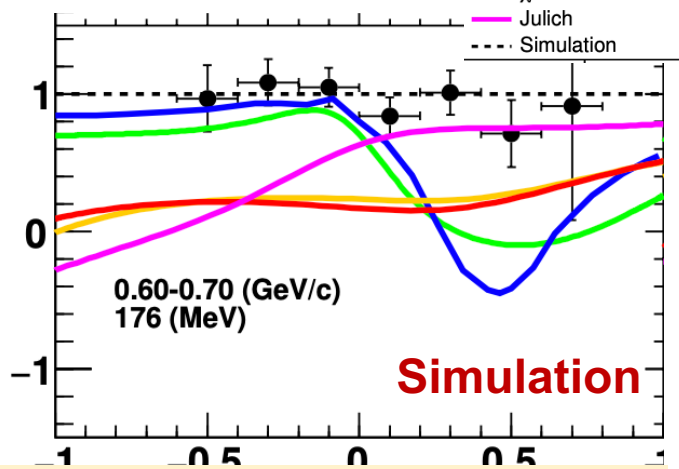
- Left-Right asymmetry of  $\Lambda p$  scattering (Analyzing power)
  - spin-orbit interaction
- Polarization change before and after the scattering (Depolarization)
  - spin-spin interaction, tensor interaction



## Analyzing power



## Depolarization



But, it takes long time to perform this experiment, because this experiment needs a construction of new K1.1 beam line. We have taken action to take data using photon beam at SPring-8 to measure  $d\sigma/d\Omega$ .

# $\Lambda p$ scattering experiment using photo-produced $\Lambda$ at SPring-8 (HYPS project)

This project is performed as RIKEN-TOHOKU project



# Building $\Lambda N$ interaction from $\Lambda N$ scattering experiment using photo-produced $\Lambda$

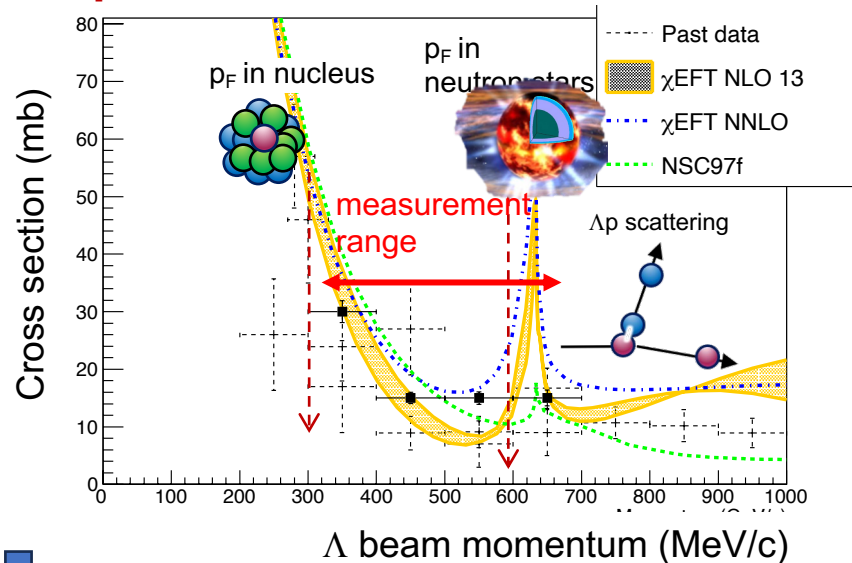
## Purpose of research

Building the realistic  $\Lambda N$  interaction by providing  $\Lambda N$  scattering data to chiral EFT theory

$\Lambda N$  interaction is still uncertain due to lack of scattering data, although the interaction is essential to describe many-body system with  $\Lambda$  such as hypernuclei and neutron stars.

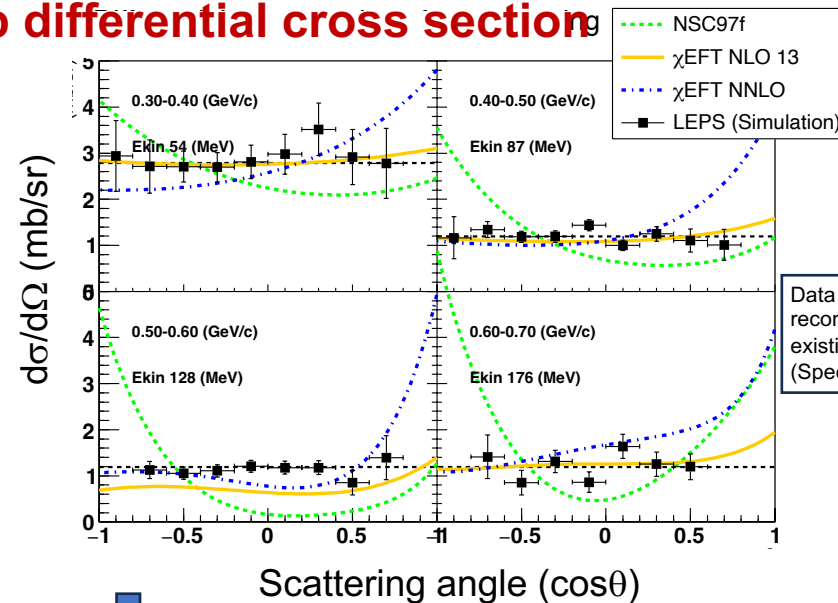
We plan to perform  $\Lambda p$  scattering experiment at BL33LEP

## $\Lambda p$ total cross section



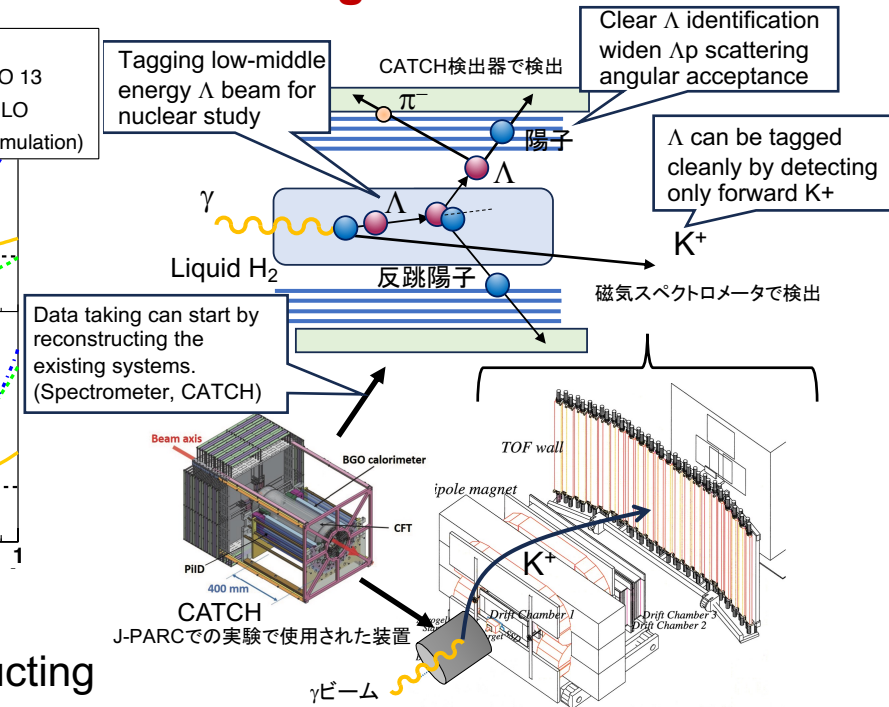
Density (radial) dependence of  $\Lambda N$  interaction

## $\Lambda p$ differential cross section



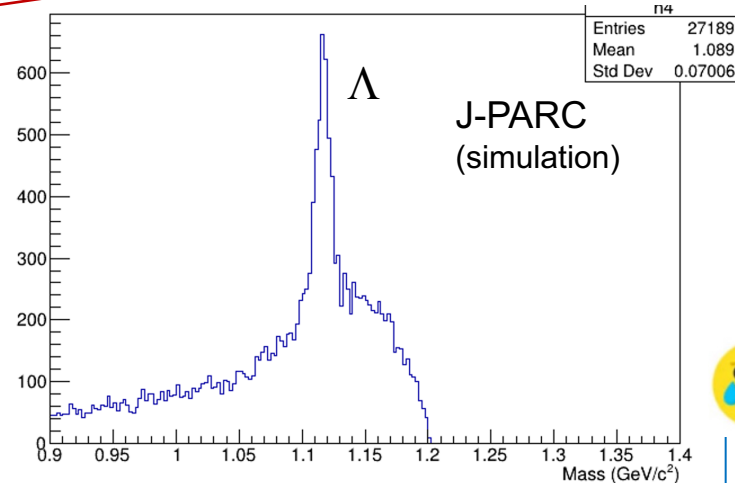
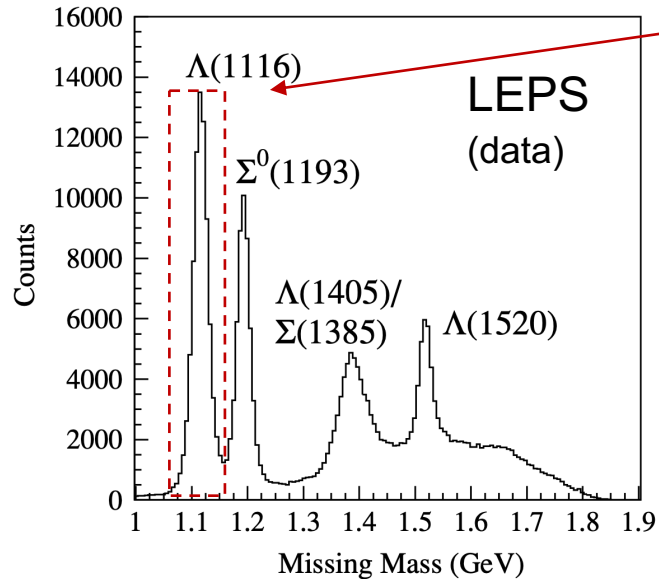
Essential input for constructing realistic  $\Lambda N$  interaction

## Advantage at BL33LEP



# Why BL33LEP?

**Advantage of  $\gamma$  beam:**  $\Lambda$  production can be identified most clearly by detecting  $K^+$



$K^0$  identification

→ Some background remains due to identification method



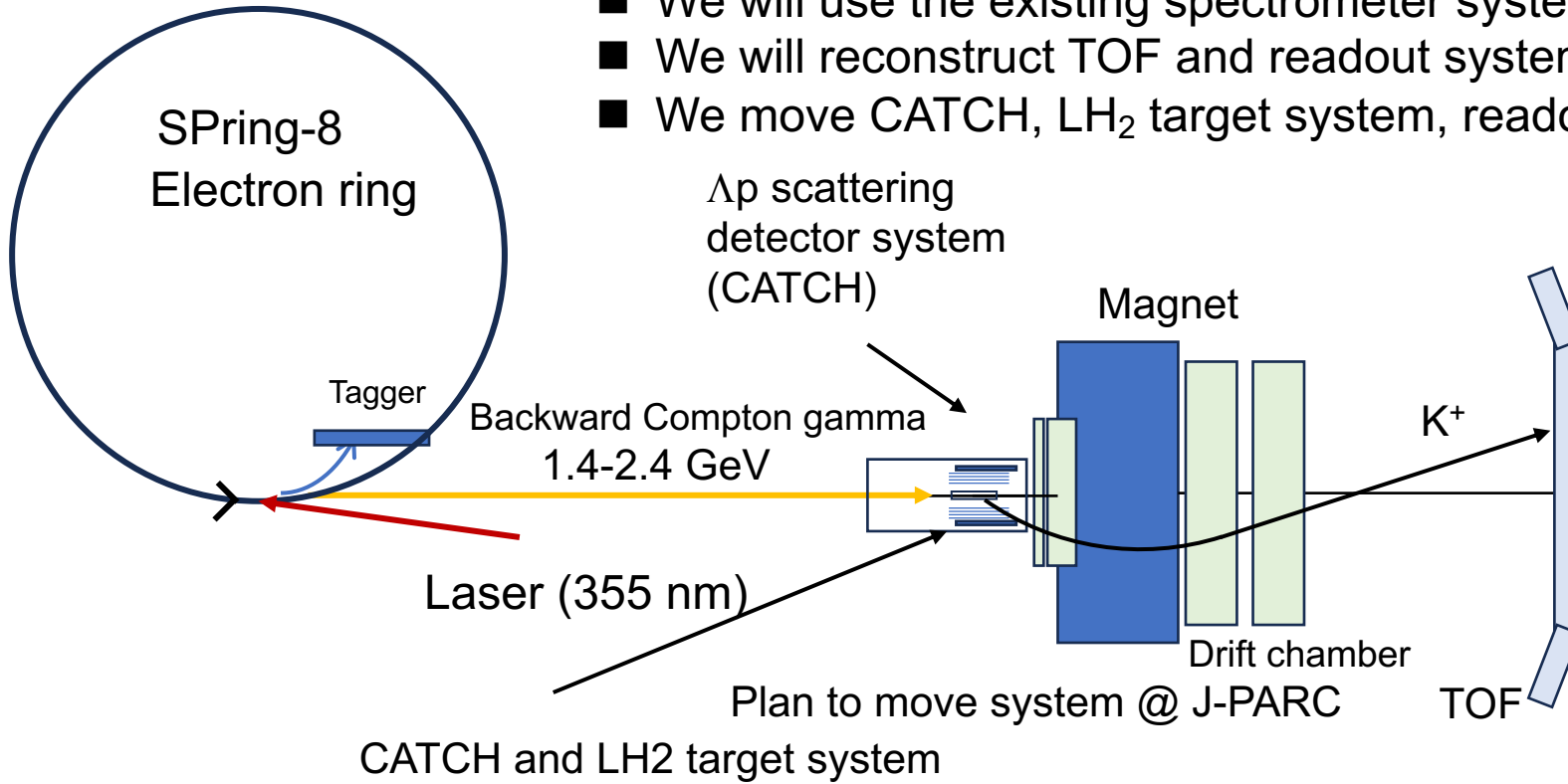
In the scattering analysis, we need to require two protons. This limits angular acceptance of  $\Lambda p$  scattering

**Advantage of backward Compton  $\gamma$  beam:** A forward spectrometer can be placed, making it possible to tag low-energy  $\Lambda$  with small momentum transfer for the first time .

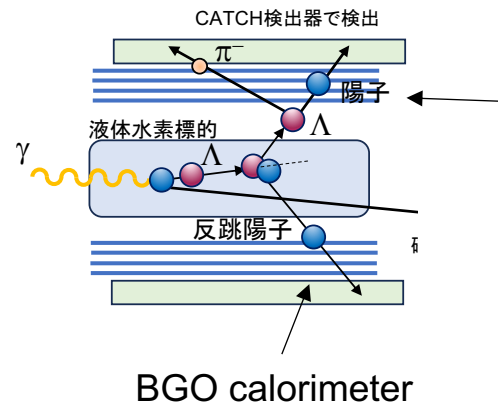
We can get  $\Lambda$  beam from 0.3  $\text{GeV}/c$  (for nuclear study) to 0.6  $\text{GeV}/c$  (for neutron star)

# Experimental setup of $\Lambda p$ scattering experiment at BL33LEP

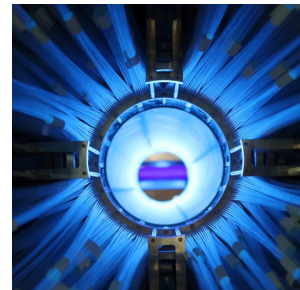
- We will use the existing spectrometer system.
- We will reconstruct TOF and readout systems
- We move CATCH, LH<sub>2</sub> target system, readout system to SPring-8 from J-PARC



CATCH and LH<sub>2</sub> target system



Cylindrical fiber tracker



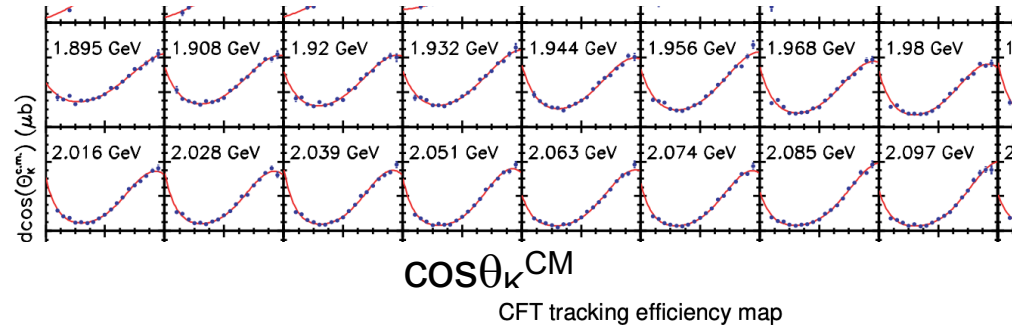
# Feasibility study at BL33LEP : Acceptance of $\Lambda$ beam detection

We estimated essential parameters w/ simulation study

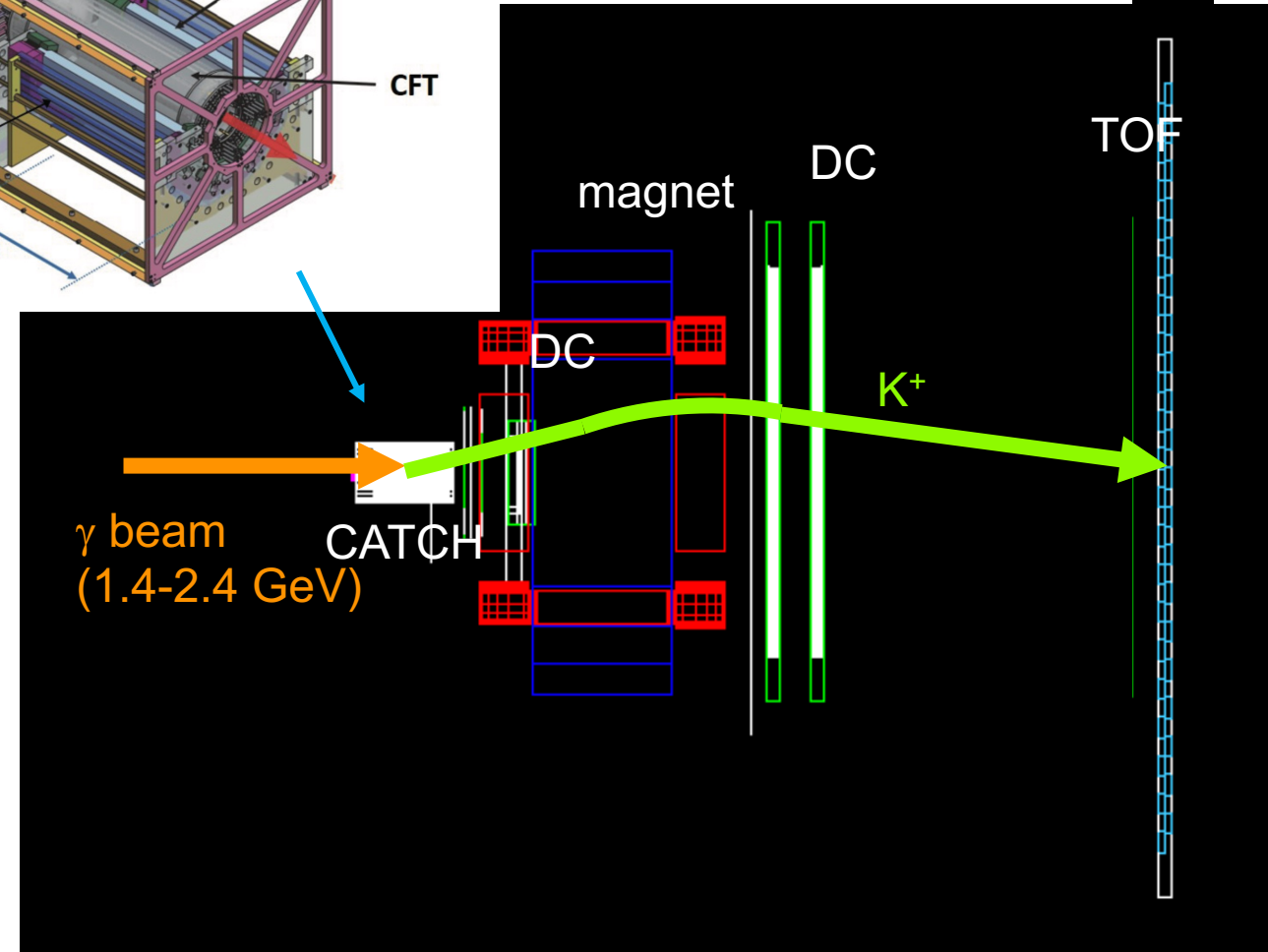
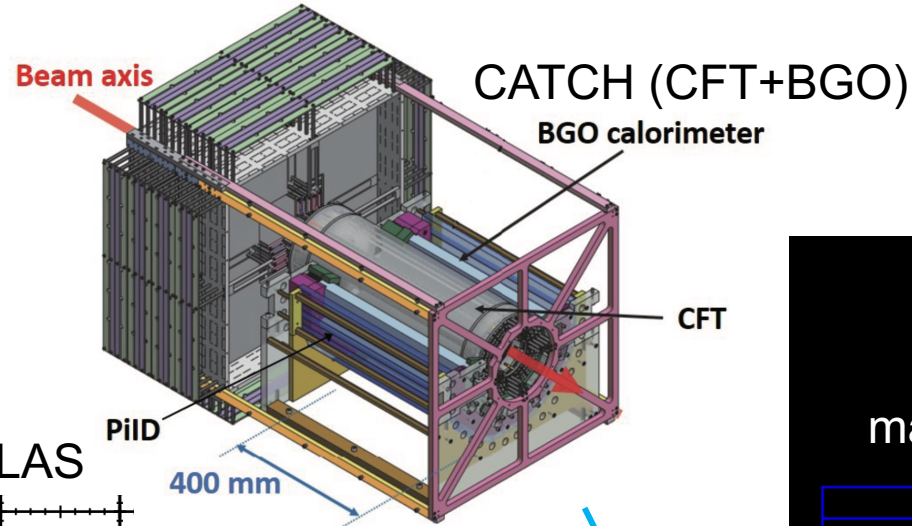
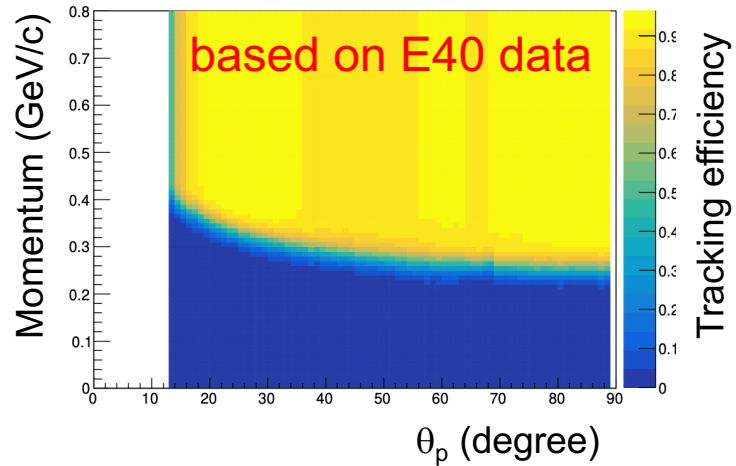
- Acceptance for  $K^+$
- $\Lambda$  beam momentum distribution
- $\Lambda p$  scattering identification quality

## Input

Angular distribution of  $\Lambda$  production by CLAS

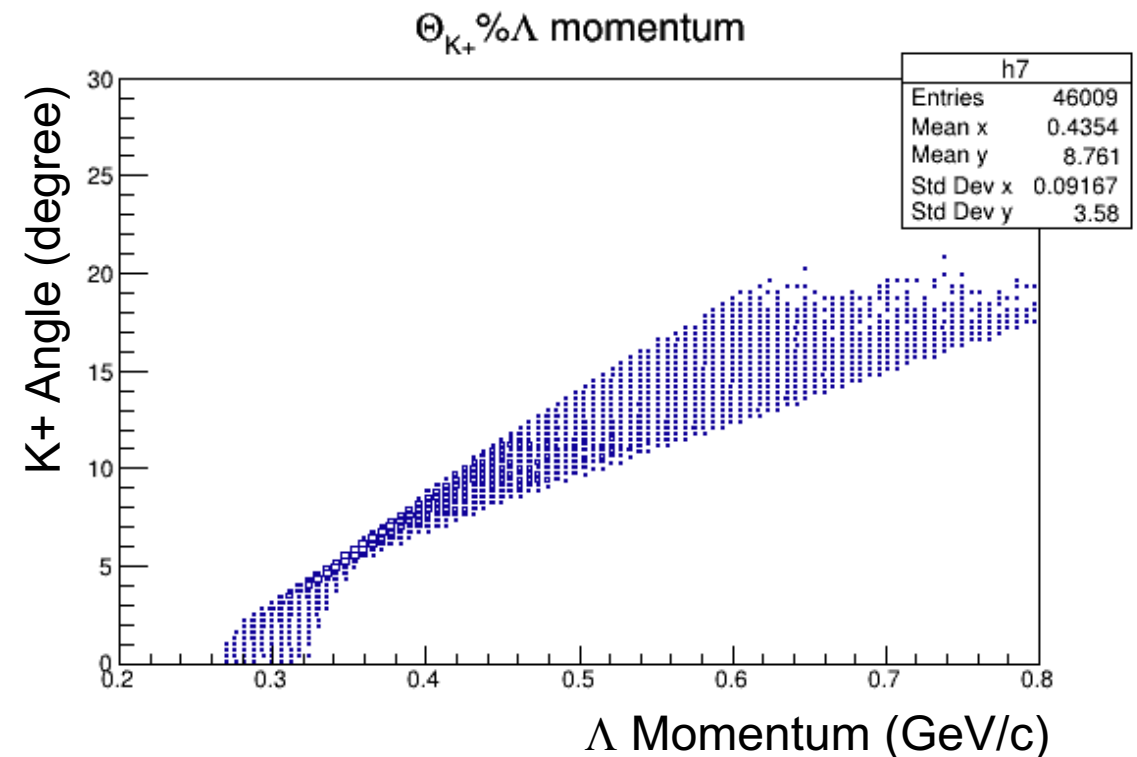
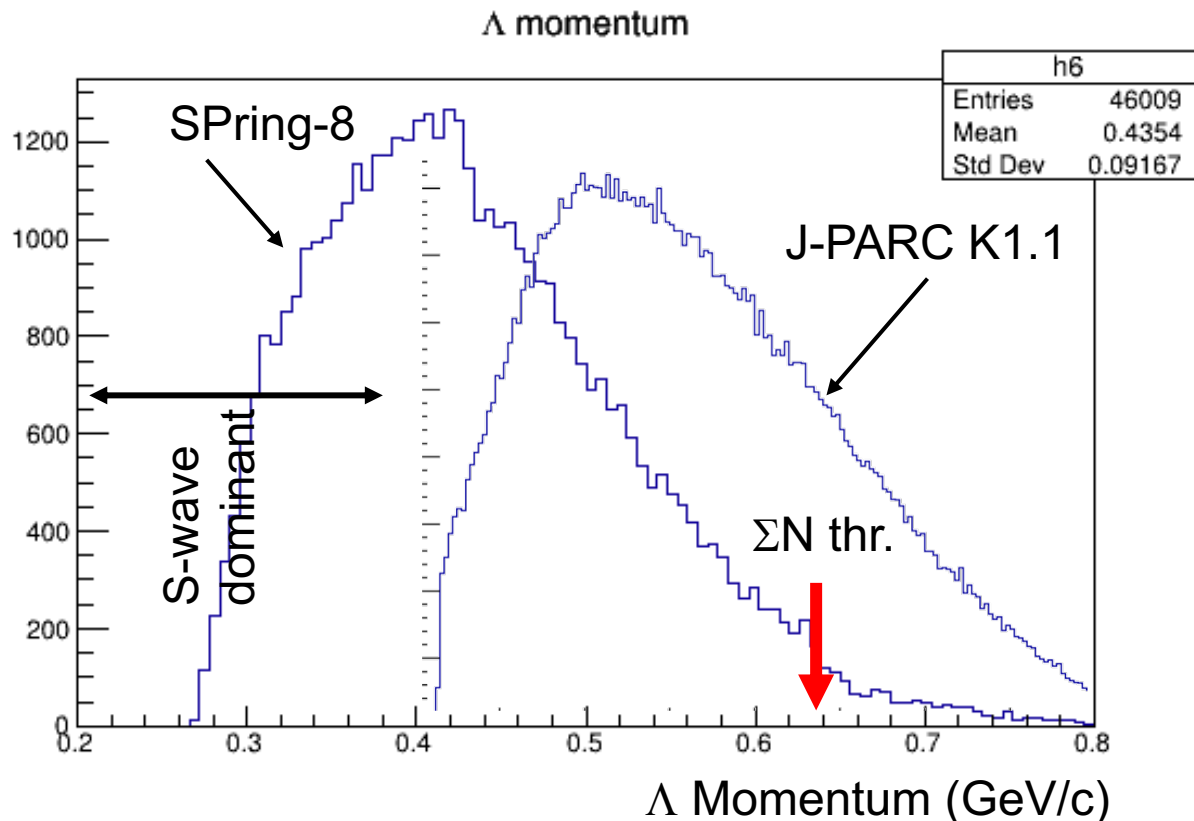


CATCH efficiency



# $\Lambda$ beam identification (Acceptance, Momentum)

- Acceptance for mimicked setup for HYPS :  $\sim 10\%$ 
  - Corresponding  $K^+$  momentum : 1 ~ 2 GeV/c
- $\Lambda$  momentum range : 0.3 ~ 0.55 GeV/c
  - Cover lower momentum region. Close relationship with hypernuclear physics.
  - Good complementary with K1.1



# $\Lambda$ beam yield estimation

Items	Estimated values
$\gamma$ beam intensity	2 MHz
$\Lambda$ production cross section	1.5 $\mu\text{b}$
Liquid H <sub>2</sub> target thickness, number	30 cm $\rightarrow$ 12.7 x 10 <sup>23</sup> [1/cm <sup>2</sup> ]
K <sup>+</sup> acceptance	0.11
K <sup>+</sup> survival rate	0.69 (for $p_{K^+}=1.5$ GeV/c, L=3.7 m)
DAQ, analysis efficiency	0.9 (assumption)
Tagged $\Lambda$ yield per second	0.281 [1/s]
Tagged $\Lambda$ yield per day	2.42 x 10 <sup>4</sup> [1/day]

We need to accumulate  $10^7$   $\Lambda$  beams for 10% order accuracy:

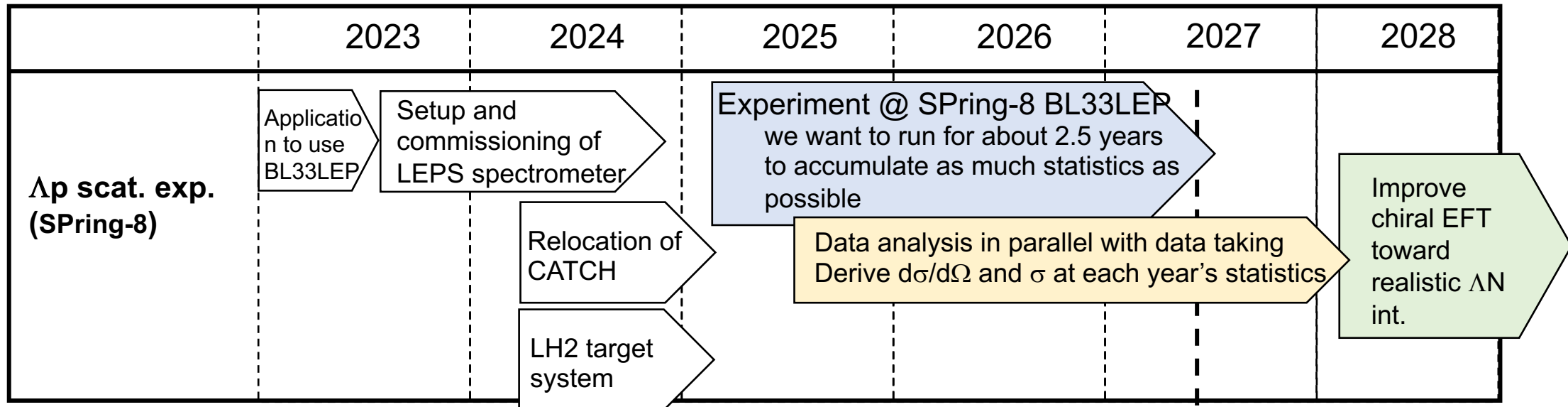
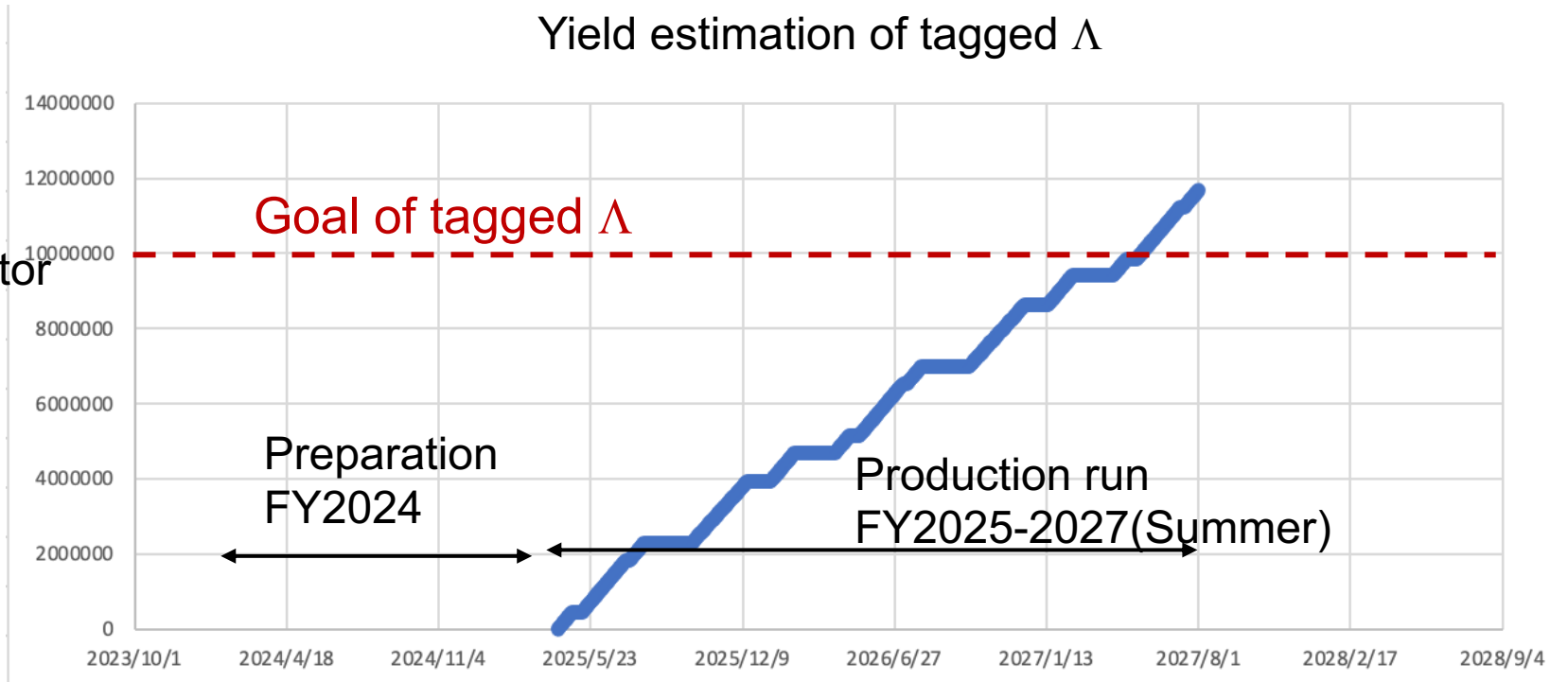
c.f.  $\Sigma^-$ -p scattering (E40)  $1.7 \times 10^7$   $\Sigma^-$  beam



Beam time of ~400 days is necessary

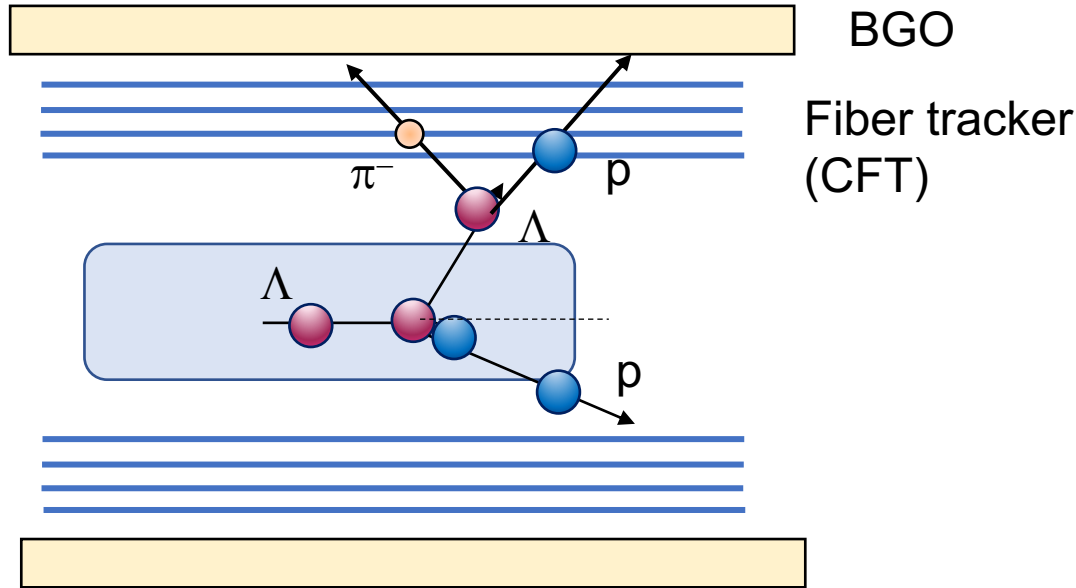
# Timeline

FY2024 : Commissioning of detector  
 FY2025: Start data taking  
 until Aug. 2027



Beam stop at 2027 August ?

# $\Lambda$ p scattering identification



Proton can be stopped in BGO

→ Proton's direction and energy information

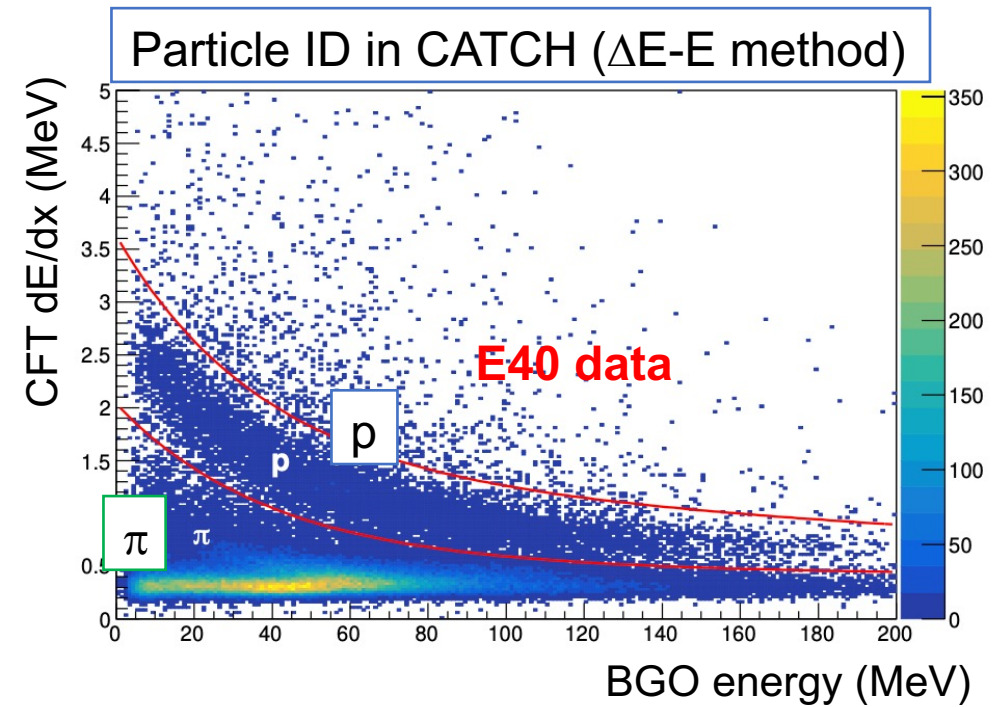
Pion cannot be stopped in BGO for many cases

→ Only direction information

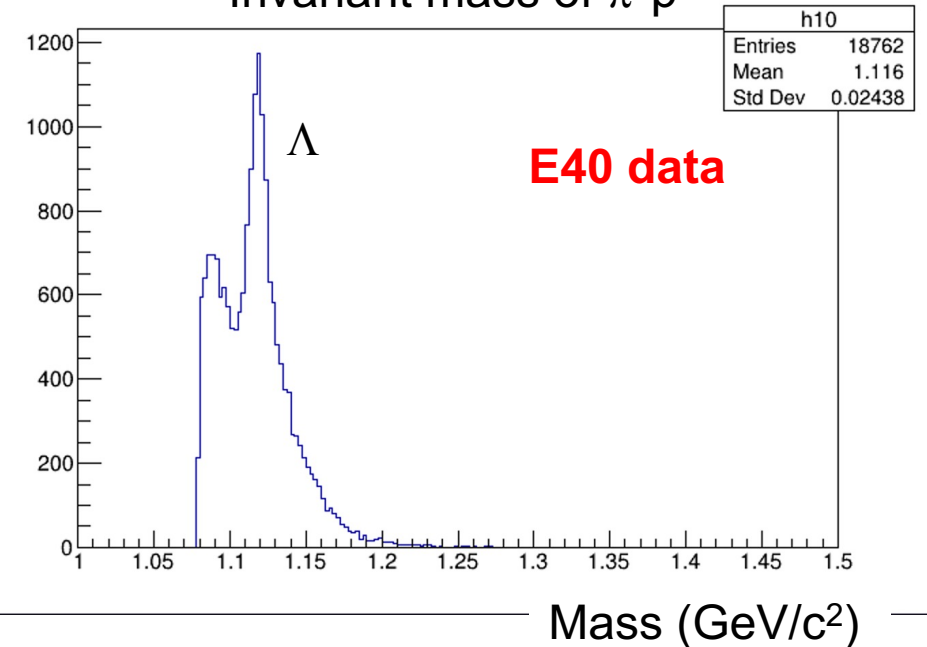
But,  $\pi^-$  from  $\Lambda$  decay has low momentum ( $\sim 150$  MeV/c)

→ many of  $\pi^-$  can be stopped (resolution is not so good)

$\pi^-$  energy is calculated from  $\Lambda$ 's decay kinematics.

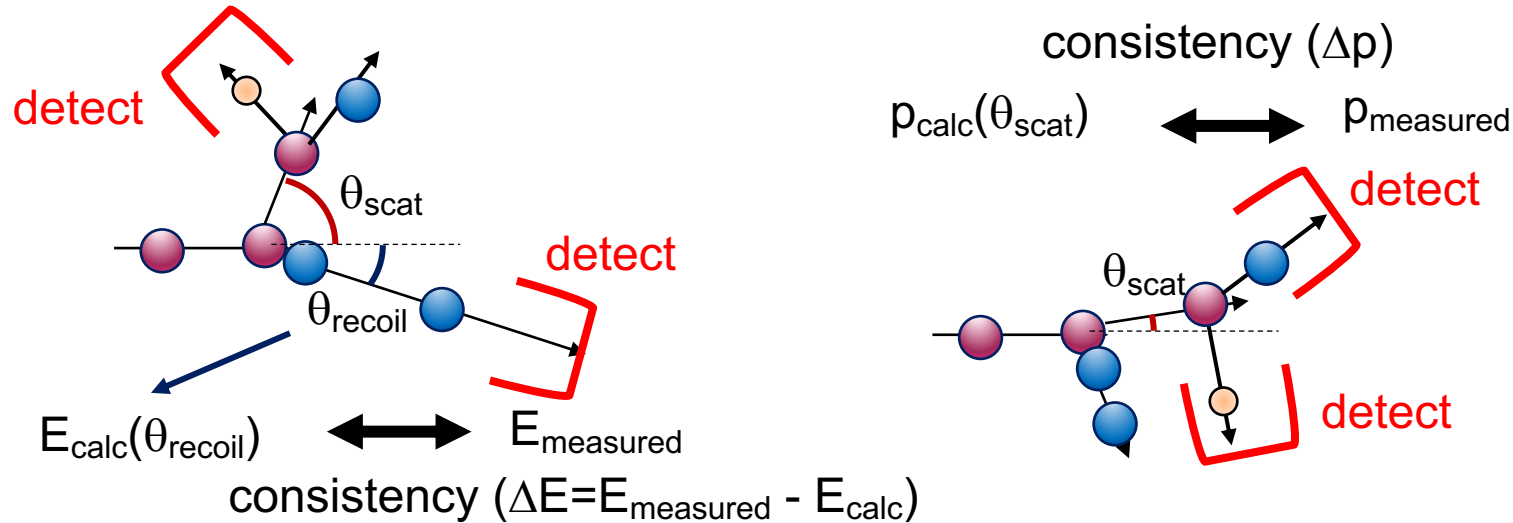


Invariant mass of  $\pi^-p$

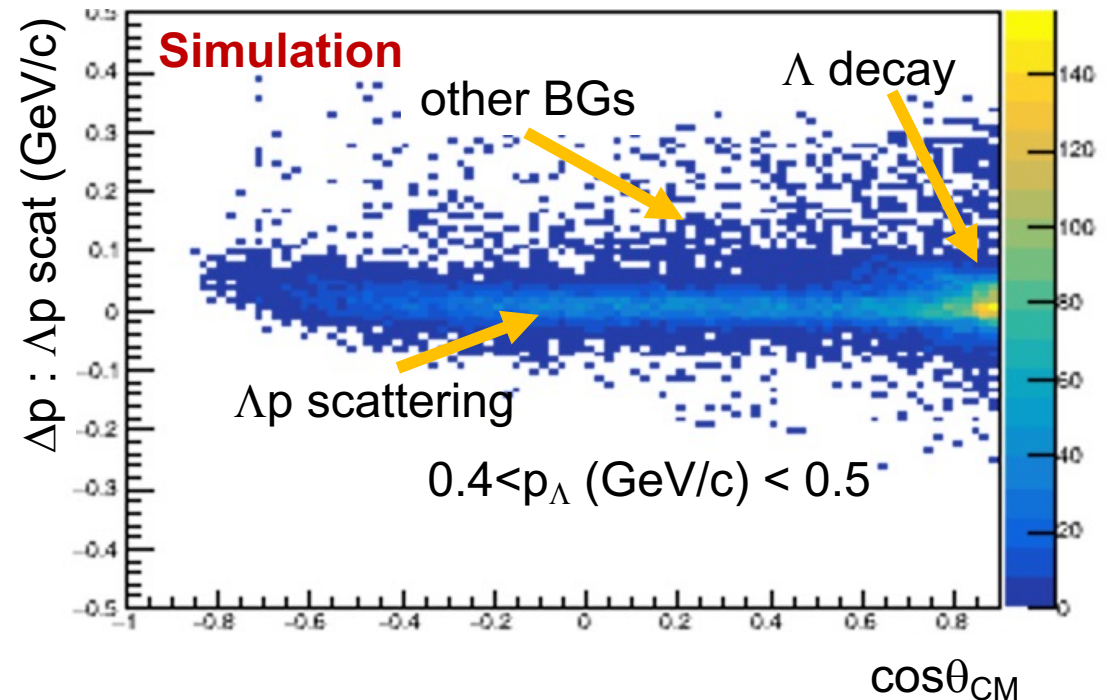
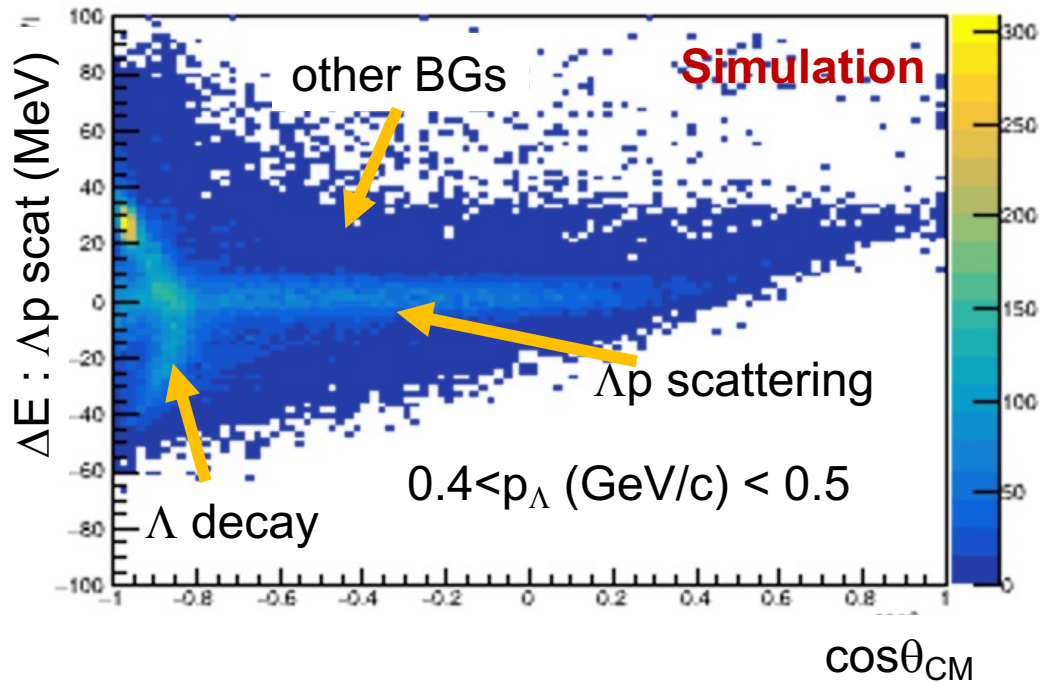




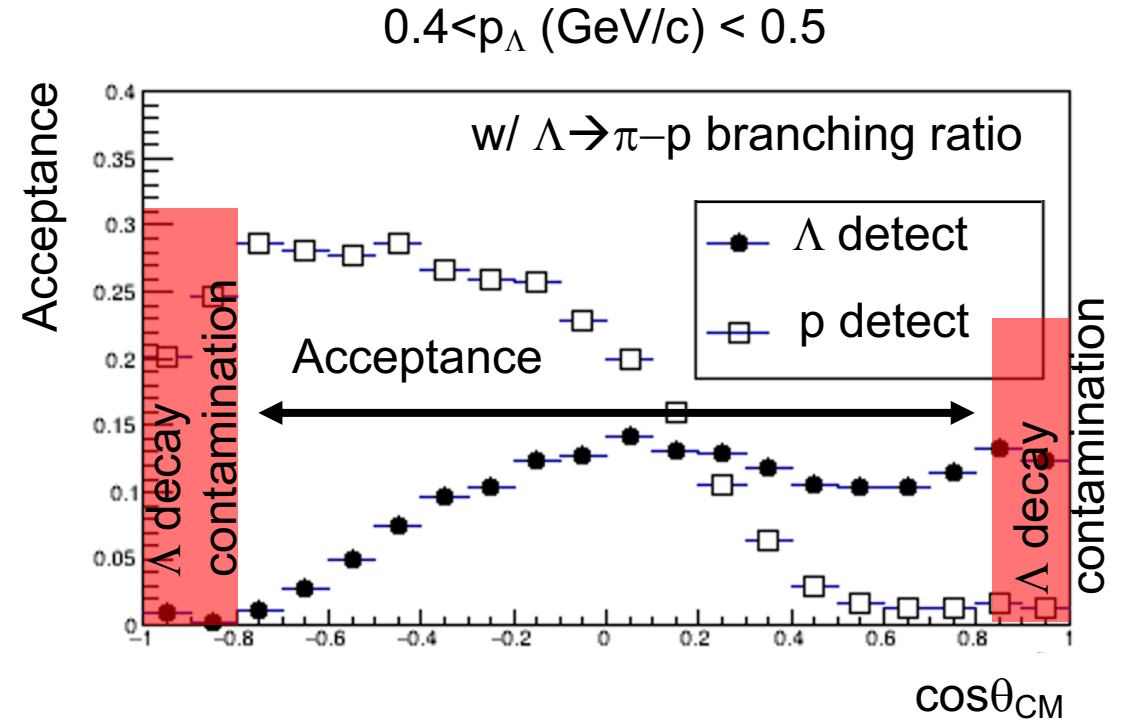
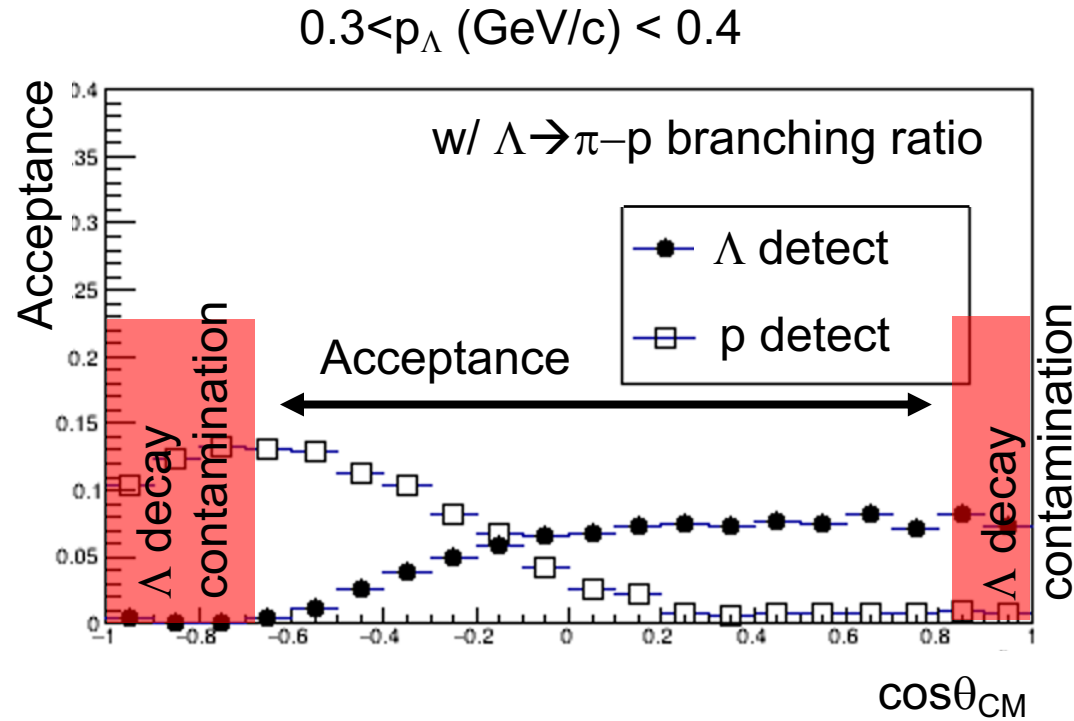
# Angular acceptance of $\Lambda p$ scattering by CATCH



$\Lambda p$  cross section : 30 mb (assumed)  
 $pp, \pi^- p$  : realistic cross section  
 $\Lambda$  decay : 1/100



# Angular acceptance of $\Lambda p$ scattering by CATCH



Forward scattering angle : covered by  $\Lambda$  detection

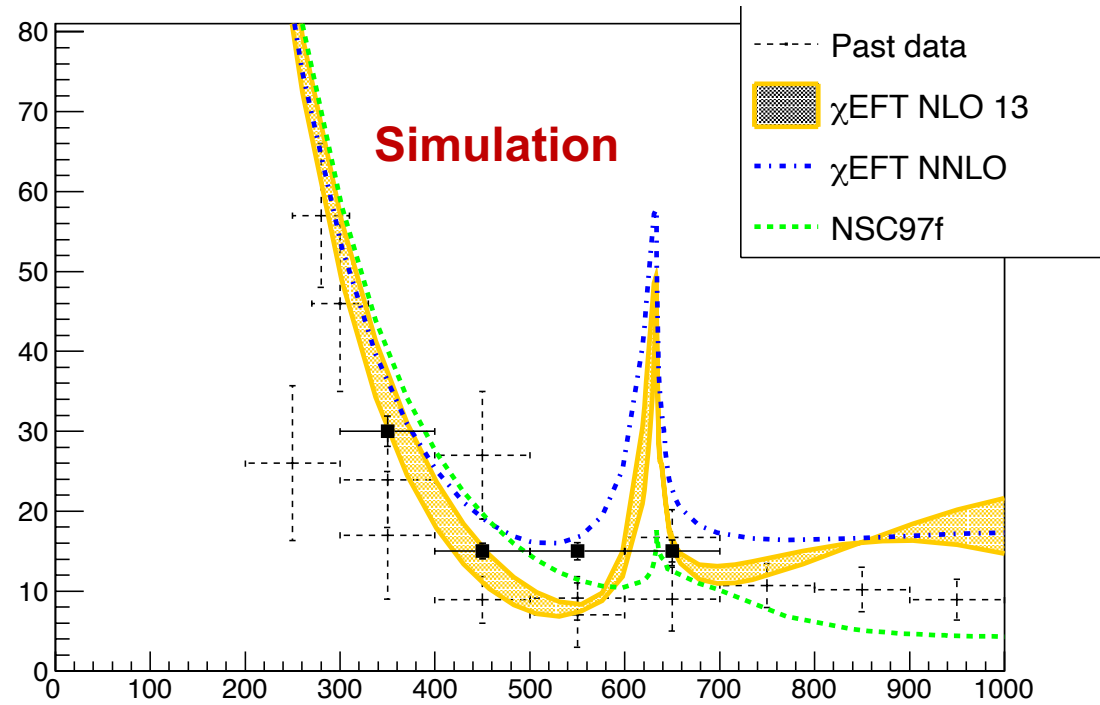
Backward scattering angle : covered by proton detection

But, very forward and very backward regions might be hard due to  $\Lambda$  decay contamination

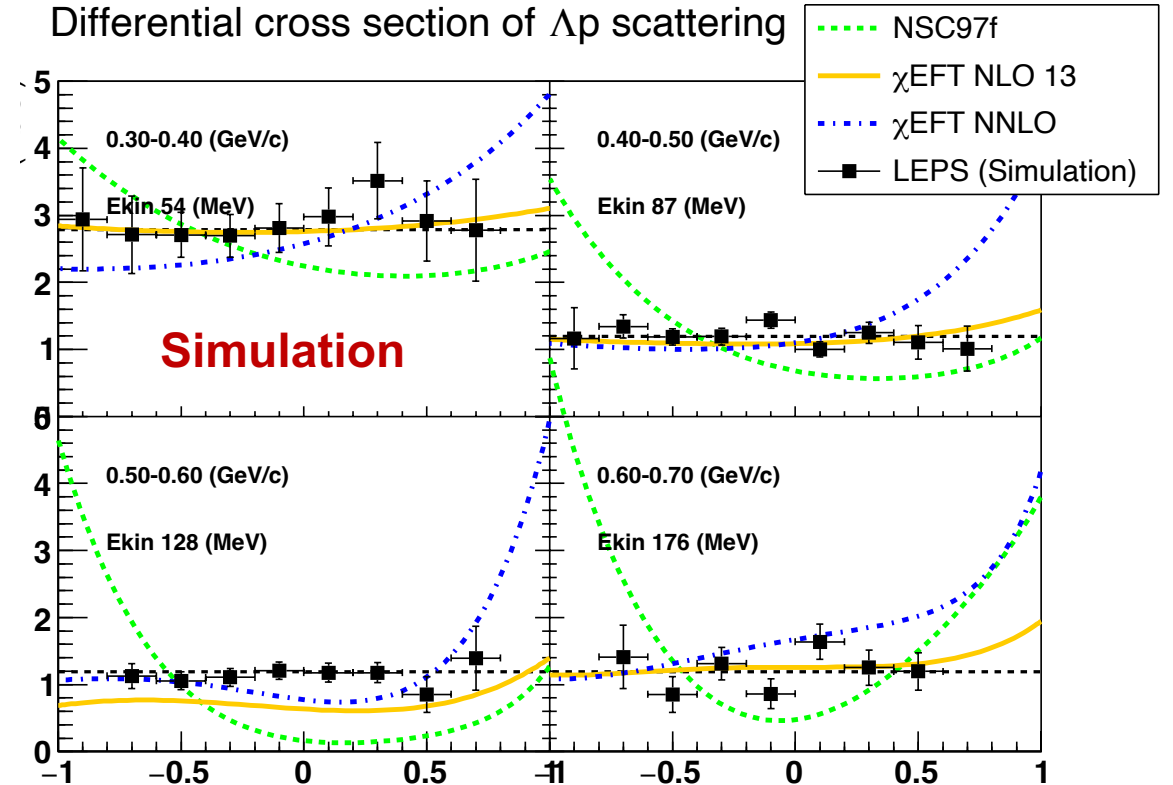
Even though, rather wide acceptance can be obtained.

# Expected results

Momentum dependence of total cross section



Differential cross section of  $\Lambda p$  scattering



Accurate  $d\sigma/d\Omega$  data and total cross section can put strong constraint on  $\Lambda N$  interaction theory.

Chiral NNLO  $\Lambda N$  interaction shows rather attractive nature

- Larger cross section around  $p_\Lambda \sim 500$  MeV/c
- Deeper  $U(\Lambda)$  potential (-35~-37 MeV)

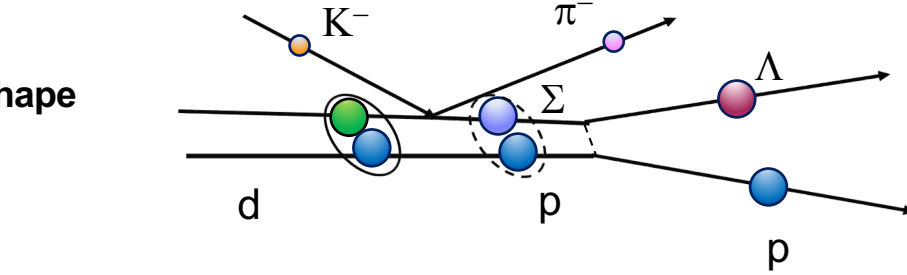
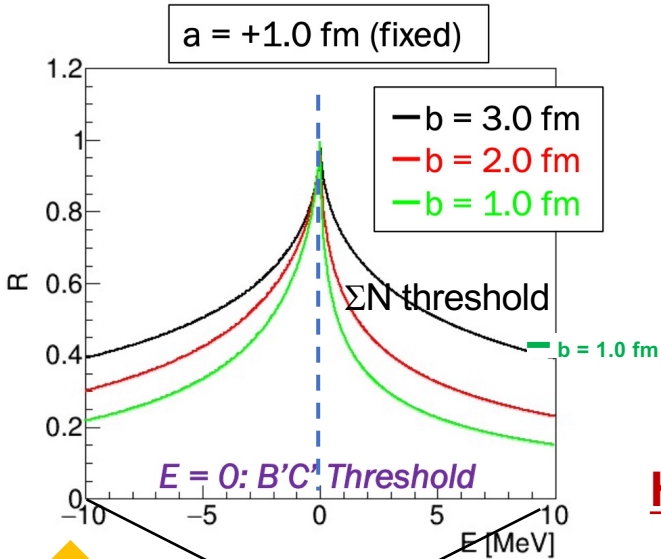
J. Haidenbauer et al.,  
Eur.Phys.J.A 59 (2023) 3

# Summary

- Many progresses have been obtained in the BB interactions study.
  - Lattice QCD, Chiral EFT, ...
  - Femtoscopy is successfully used for the hadron-hadron interaction study.
  - YN scattering experiment gets possible!
- New collaborative project regarding the two-body  $\Lambda N$ ,  $\Sigma N$  interactions
  - $\Lambda p$  scattering experiment with photo-produced  $\Lambda$
  - Precise measurement of  $\Sigma N$  cusp shape with S-2S
  - Lattice QCD potential of  $\Lambda N$ ,  $\Sigma N$ ,  $\Lambda N$ - $\Sigma N$  potentials
- New experimental project will begin at SPring-8 to measure  $\Lambda p$  scattering cross section
  - $\Lambda$  particle ( $300 < p_\Lambda < 600$  MeV/c) can be identified cleanly by  $\gamma p \rightarrow K^+ \Lambda$  reaction.
  - Experimental technology developed at J-PARC will be introduced.
  - Our goal
    - Total cross section measurement better than 10%
    - First derivation of the differential cross section for  $\Lambda p$  elastic scattering

# Precise $\Sigma N$ cusp measurement with $K^-d \rightarrow \Lambda p \pi^-$ reaction (J-PARC E90)

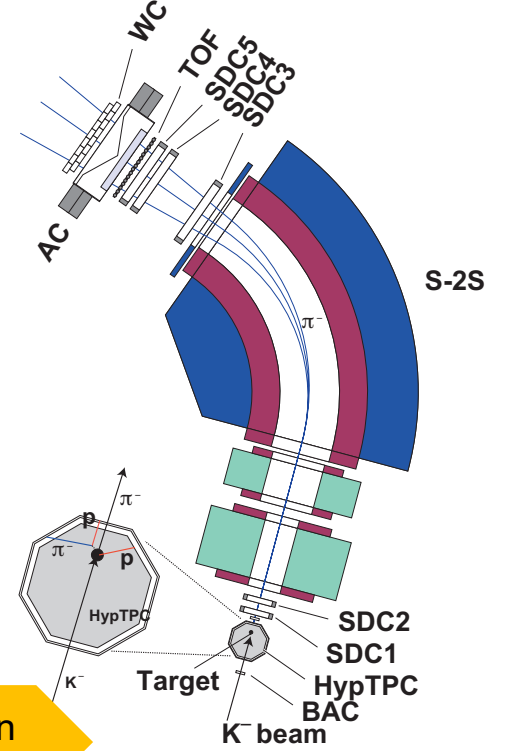
$\Lambda N$ - $\Sigma N$  Coupling dependence of cusp shape



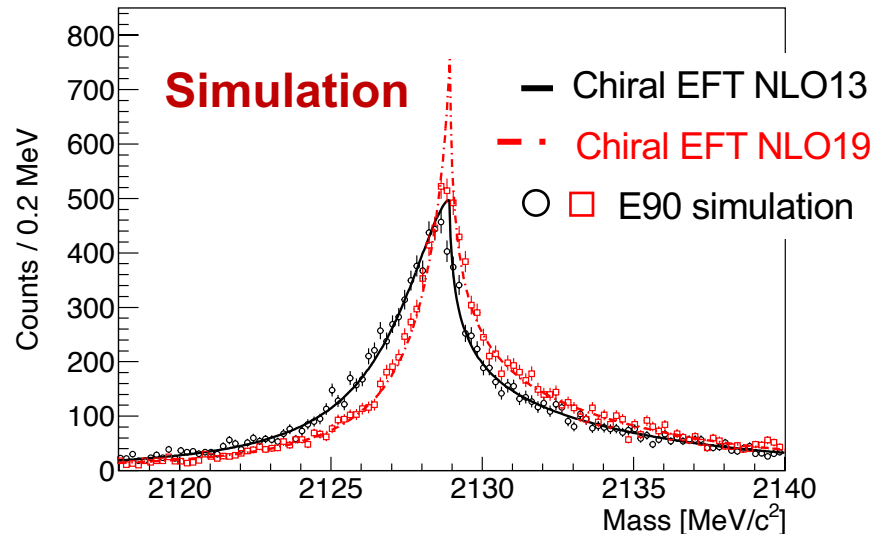
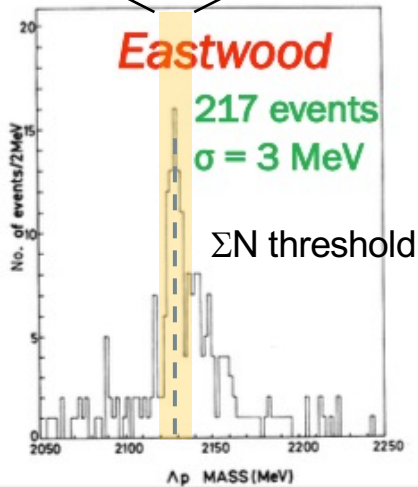
The cusp shape is represented by the scattering length  $A=a + ib$  of  $\Sigma N$  ( $I=1/2, S=1$ ) channel

- Real part  $a$  : Attraction of the  $\Sigma N$  interaction at low energy
- Imaginary part  $b$  :  $\Lambda N$ - $\Sigma N$  coupling

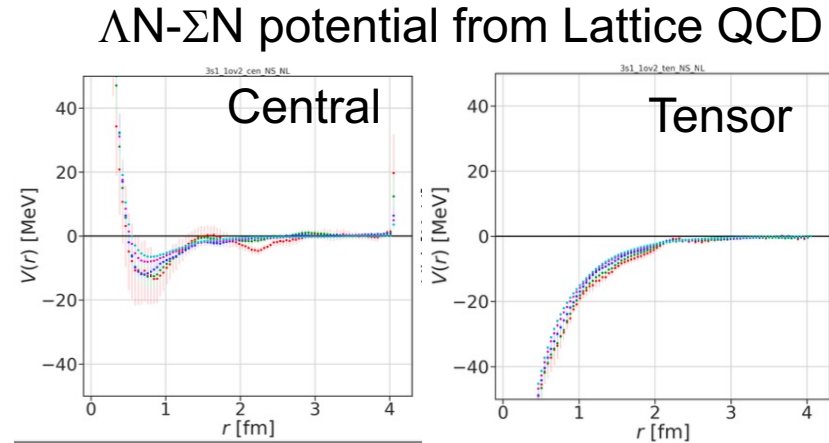
High resolution cusp measurement of  $\sigma=0.4$  MeV w/ S-2S



Past measurement

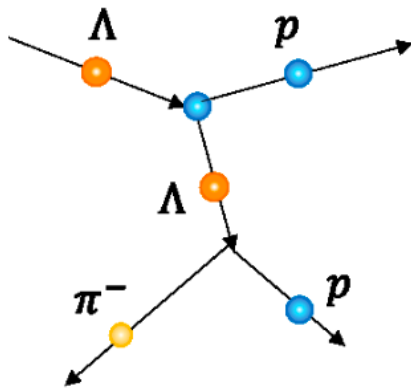


Direct comparison by scattering length

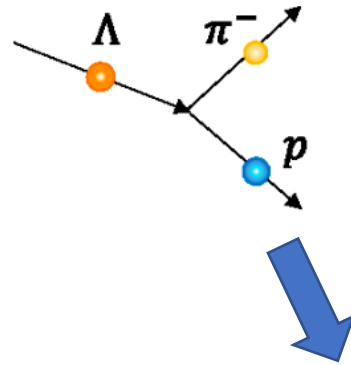


# Background of $\Lambda p$ scattering

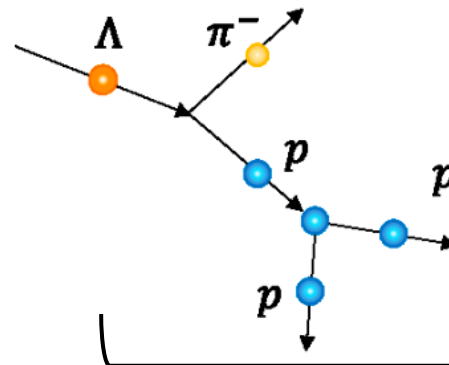
$\Lambda p$  scattering



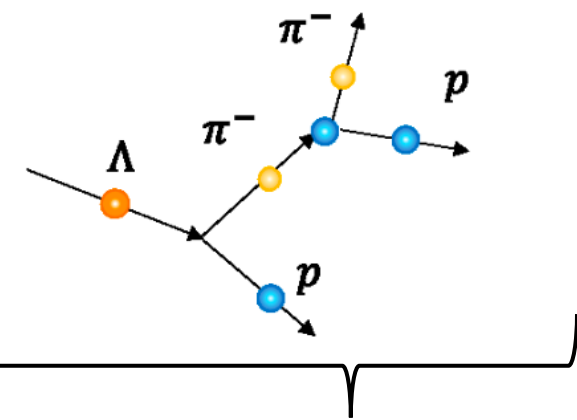
$\Lambda$  decay



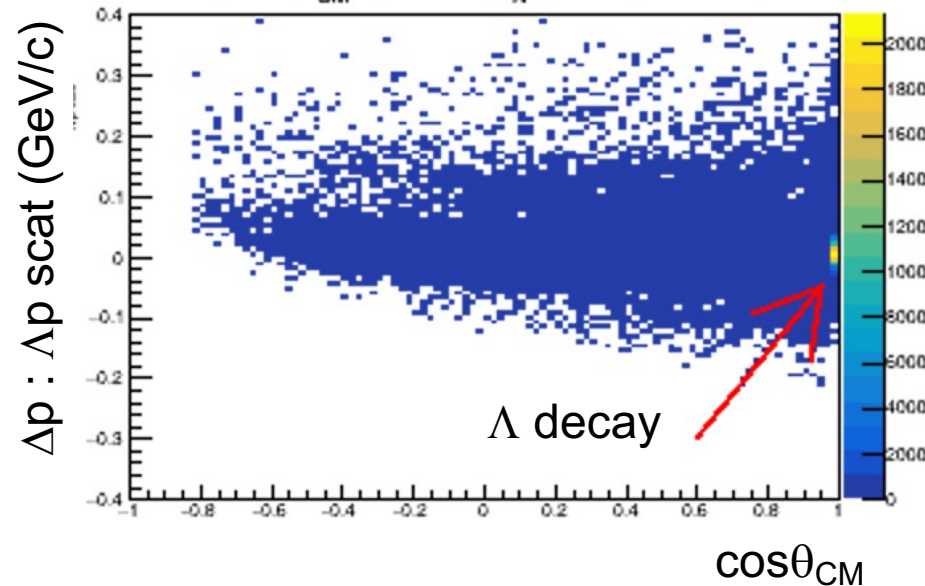
$pp$  scattering from  $\Lambda$  decay



$\pi^- p$  scattering from  $\Lambda$  decay



$\Delta p : \cos(\theta_{CM})$  ( $0.300 < p_{\Lambda} \text{ (GeV/c)} < 0.400$ )



To suppress background,

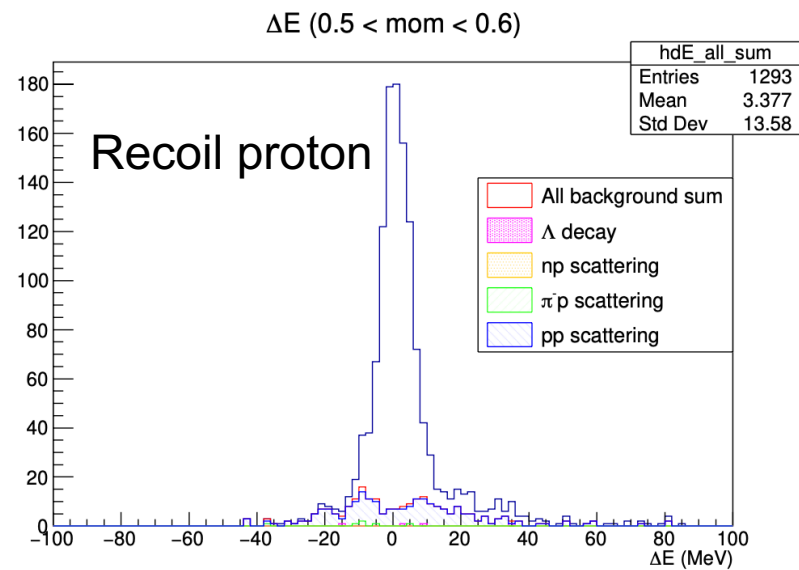
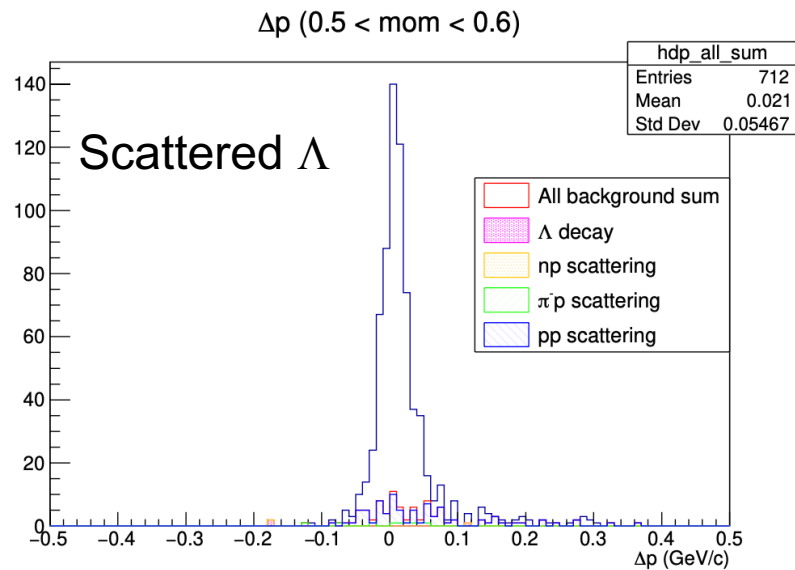
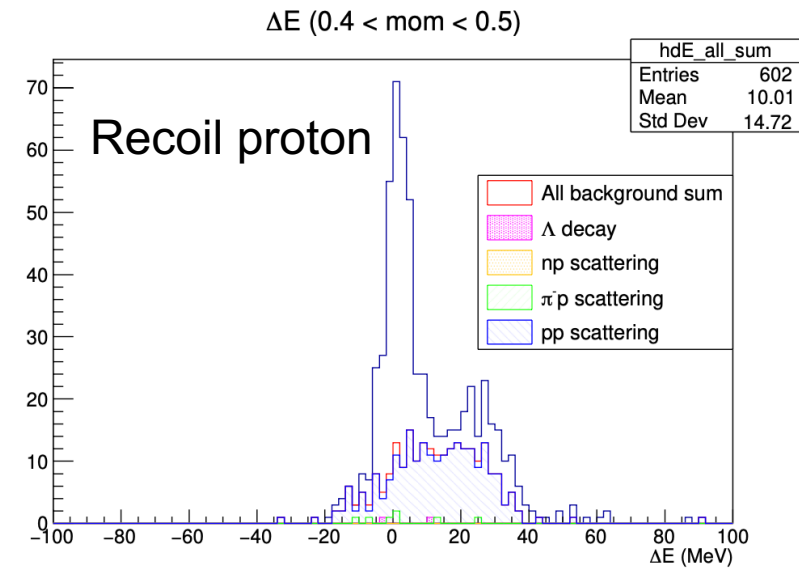
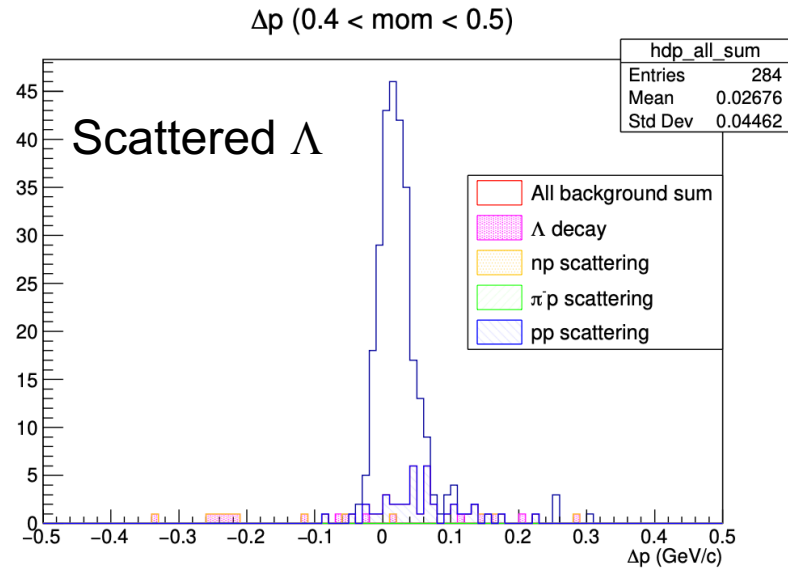
- one proton and one  $\pi^-$
- two protons

are essential

These rescattering events after  $\Lambda$  decay can be partially rejected kinematically

$\Lambda$  decay can be identified as scattering at 0 degree

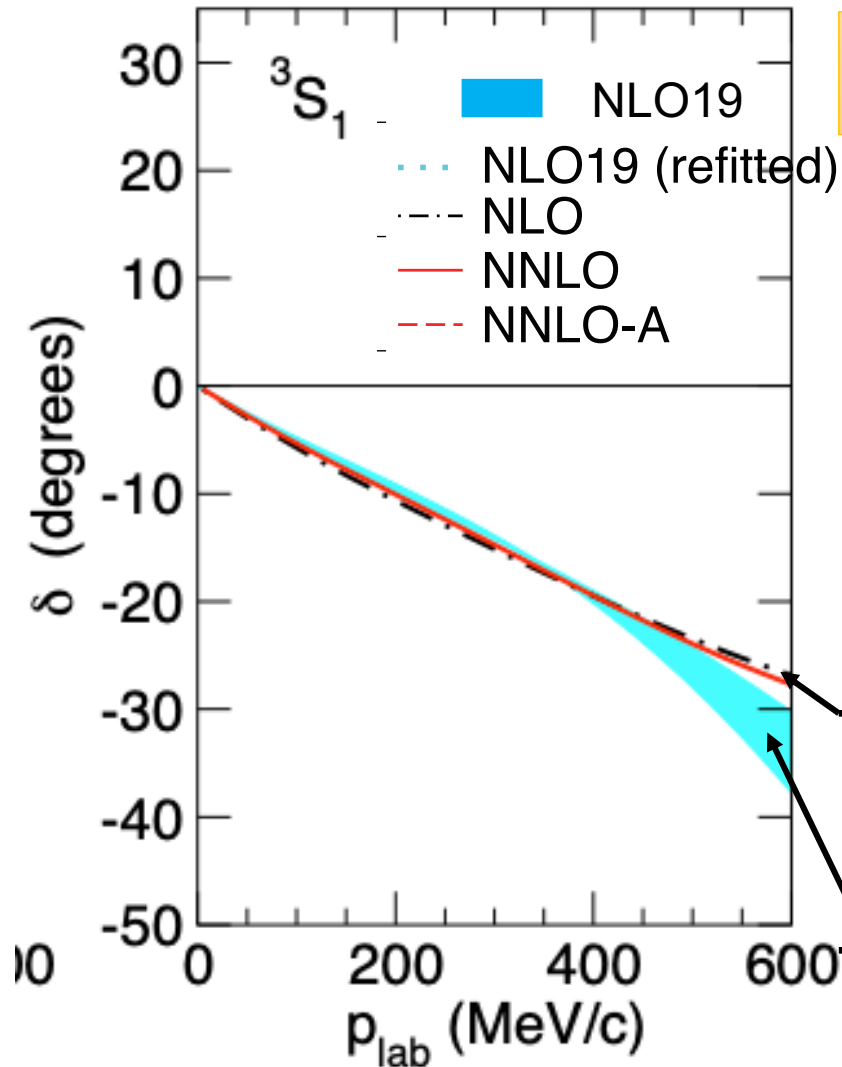
# Event ID for $10^7 \Lambda$







# Phase shift in Chiral EFT NNLO and $U_\Sigma$



J. Haidenbauer et al.,  
arXiv:2301.00722

Based on the E40  $\Sigma^+p$  phase shift,  
 $U_\Sigma$  becomes less repulsive.

${}^3_\Lambda\text{H}$

$YN$ potential	$B_\Lambda$ [MeV]	$E$ [MeV]	$P_\Sigma$ [%]	$U_\Lambda(0)$	$U_\Sigma(0)$
SMS $N^2\text{LO}(500)$	0.147	-2.371	0.25	-33.1	6.4
SMS $N^2\text{LO}(550)^a$	0.139	-2.362	0.25	-38.5	2.5
SMS $N^2\text{LO}(550)^b$	0.125	-2.348	0.24	-35.9	2.5
SMS $N^2\text{LO}(600)$	0.172	-2.395	0.22	-37.8	0.1
NLO13(600)	0.090	-2.335	0.25	-21.6	17.1
NLO19(600)	0.091	-2.336	0.21	-32.6	16.9

# $\Sigma N$ ( $I=3/2$ ) phase shift in chiral EFT

