Inference of hyperon–nucleon interactions from light hypernuclei

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SPICE: Strange hadrons as a Precision tool for strongly InteraCting systEms

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Introduction & motivation

Strangeness physics

- Interdisciplinary field connecting particle physics, nuclear physics, and astrophysics
- One of its major goals is to understand the elusive interaction of hyperons with nucleons and the nuclear medium

Theoretical analysis of hypernuclei

- Using 'effective' YN interaction models & mean-field / shell-model approaches – successful but difficult to link with the underlying free-space YN interaction, limited predictive power
- Using 'realistic' (free-space) YN interaction models ...

Constraining YN interactions

- YN scattering 'pure' but very difficult to realize, sparse database with large uncertainties (J-PARC)
- Heavy-ion collisions production and decays of light hypernuclei, correlation femtoscopy (HADES, ALICE, STAR)
- Final-state interactions in hyperon photoproduction (CLAS)
- ► Lattice QCD (HAL QCD, NPLQCD)
- Hypernuclei precise spectroscopy of hypernuclear energy levels

Introduction & motivation

Theoretical analysis of hypernuclei using realistic YN interactions

- Combines modern developments of YN interactions based on *x*EFT and ab initio few- and many-body approaches
- Computationally demanding
- \blacktriangleright Can reveal deficiencies of existing YN interaction models \rightarrow calibration?

Calibration of YN interaction models using hypernuclei requires

- Advanced 'ab initio' computational methods
- Quantified method uncertainties, σ_{method} associated with the solution of the many-body problem
- Quantified model uncertainties, σ_{model} associated with the choice of the nuclear interaction
- Overcoming the computational demands large number of evaluations
- Sensitivity analysis hypernuclear spectra might not be sensitive to certain parameters (LECs) of the YN interaction models
- Simultaneous fitting of other observables

Ab initio calculations of light hypernuclei

Ab initio calculations of light hypernuclei

Ab initio methods aim to solve the (hyper)nuclear many-body problem starting from realistic (free-space) interactions exactly or with controlled approximations

Ab initio no-core shell model

Quasi-exact method to solve the few- and many-body Schrödinger equation

$$\bigg(\sum \frac{\hat{\boldsymbol{p}}_i^2}{2m_i} + \sum \hat{\boldsymbol{V}}_{NN;ij} + \sum \hat{\boldsymbol{V}}_{NNN;ijk} + \sum \hat{\boldsymbol{V}}_{YN;ij}\bigg) \Psi = \boldsymbol{E} \Psi$$

[Navrátil et al., JPG 36, 083101 (2009); DG et al., FBS 55, 857 (2014); Wirth et al., PRL 113, 192502 (2014); Le et al., EPJA 56, 301 (2020)]

 Wave function is expanded and Hamiltonian is diagonalized in a finite A-particle harmonic oscillator (HO) basis

$$\Psi(\mathbf{r}_{1},\ldots,\mathbf{r}_{A}) = \sum_{N \leq N_{max}} \Phi_{N,\omega}^{HO}(\mathbf{r}_{1},\ldots,\mathbf{r}_{A})$$

Converges to exact results for $N_{max} \rightarrow \infty$

- Input NN+NNN and YN interactions derived from χ EFT
 - ► The NNLO_{sim} family at NNLO [Carlsson et al., PRX 6, 011019 (2016)]
 - ► Jülich YN at LO [Polinder et al., NPA 779, 244 (2006)]

Ab initio calculations of light hypernuclei: method uncertainties

Method uncertainties associated with convergence of the solution of the many-body problem



- NCSM-calculated energies typically exhibit undesired dependence on the HO basis frequency ħω and truncation N_{max}
- Convergence properties of observables calculated in finite HO bases are rather well understood [Wendt et al., PRC 91, 061391 (2015)]
 - NCSM model-space parameters (N_{max}, ħω) recast into infrared (IR) and ultraviolet (UV) scales (L_{IR}, Λ_{UV})
 - In a regime with negligible UV corrections, IR corrections are universal

 $E(L_{IR}) = E_{\infty} + a_0 \exp(-2\kappa_{\infty}L_{IR}) + \cdots$

Ab initio calculations of light hypernuclei: method uncertainties

 Infrared extrapolation formulated as a Bayesian inference problem

$$\begin{split} \mathsf{E}(\mathsf{L}_{\mathsf{IR}}) &= \mathsf{E}_{\infty} + \Delta \mathsf{E}_{\mathsf{IR}} \exp(-2\kappa_{\infty}\Delta\mathsf{L}_{\mathsf{IR}}) \\ &\times \left(1 + \frac{\epsilon_{\mathsf{NLO}}}{\kappa_{\infty}(\mathsf{L}_{\mathsf{IR},\,\mathsf{max}} + \Delta\mathsf{L}_{\mathsf{IR}})}\right), \end{split}$$

with data $\mathcal{D} = \{E(L_{IR,i})\}$ calculated in different model spaces and $\vec{\epsilon}_{NLO} \sim N(0, \Sigma(\bar{\epsilon}, \rho))$ providing a stochastic model for the NLO energy correction [DG, Htun, Forssén, PRC 106, 054001 (2022)]

• Validation for $^{3}_{\Lambda}H$



6

		B^{Exp}_{Λ} (MeV)	B^{th}_{Λ} (MeV)	
			median	68 % Cl _{method}
 Method uncertainty quantified by 68 % credible interval for the extrapolated 	³ Η	0.164(43)	0.166	[-0.001, +0.001]
	_Λ H	2.157(77)	2.78	[-0.01, +0.01]
	⁴ ∕he	2.39(3)	2.76	[-0.01, +0.01]
	Λ ⁵ He	3.12(2)	6.03	[-0.28, +0.18]
energy E_∞	${}^{4}_{\Lambda}$ H; 1 $^{+}$	1.067(80)	1.75	[-0.12, +0.10]
	$^4_{\Lambda}$ He; 1 $^+$	0.984(50)	1.71	[-0.13, +0.10]

- Dominating source of uncertainty of hypernuclear observables likely comes from the underlying YN interaction $\leftarrow \chi$ EFT truncation, regulator artifacts, calibration data uncertainties
- Energy levels of light hypernuclei are also sensitive to details of the employed nuclear NN+NNN interactions
- One can naively expect that calculated Λ separation energies should be insensitive to the choice of nuclear interaction, $B_{\Lambda} = E({}^{A}Z) E({}^{A}_{\Lambda}Z)$
- ► A rather weak residual dependence of B_A was found using a limited set of phenomenological [Nogga et al., PRL 88, 172501 (2002)] and *χ*EFT [Le et al., EPJA 56, 301 (2020)] NN interactions

Ab initio calculations of light hypernuclei: model uncertainties

- ► To expose the magnitude of systematic model uncertainties in B_A we employed [DG, Htun, Forssén, PRC 106, 054001 (2022)] the NNLO_{sim} family of 42 different nuclear NN+NNN interactions [Carlsson et al., PRX 6, 011019 (2016)]
 - NNLO_{sim} LECs fitted to reproduce simultaneously πN, NN, and NNN low-energy observables
 - Experimental data uncertainties propagate into the LECs
- Model uncertainty connected to the choice of nuclear Hamiltonian quantified by variance, σ²(NNLO_{sim}), of predictions for B_Λ

For LO YN:

 $\frac{{}^{3}_{\Lambda}H}{\sigma_{model} (keV)} \frac{{}^{3}_{\Lambda}H}{20} \frac{{}^{4}_{\Lambda}H}{80} \frac{{}^{4}_{\Lambda}He}{{}^{4}_{\Lambda}He} \frac{{}^{4}_{\Lambda}He}{{}^{4}_{\Lambda}He_{1^{+}}} \frac{{}^{5}_{\Lambda}He}{{}^{5}_{\Lambda}He}$

 $E_{sep}^{th} \left({}_{\Lambda}^{4}\mathrm{H},\mathrm{He;}\,0^{+}
ight) \left(\mathrm{MeV}
ight)$ F 6.5 $_{\rm sep}^{\rm th}({}^{4}_{\Lambda}{\rm H},{\rm He};1^{+})~({\rm MeV})$ 1.6 450500 550600 Λ_{NN} (MeV)

 T_{Lab}^{max} (MeV)

150

250

► A smaller NN+NNN-model dependence was found for NLO and NNLO YN interactions [Le et al., EPJA 60, 3 (2024)]

Ab initio calculations of hypernuclei: the curse of dimensionality

- Ab initio methods provide a reliable link between the properties of hypernuclei and the underlying hyperon-nucleon interactions
- Is it possible to directly incorporate them in optimization of hyperon-nucleon forces which require a large number of model evaluations?



This is not feasible given their computational cost

 Reoptimization of 2 LECs to the p-shell hypernuclei Λ separation energies [Knoll, Roth, PLB 846, 138258 (2023)] Emulating ab initio NCSM calculations: eigenvector continuation

Emulating ab initio NCSM calculations: eigenvector continuation

Eigenvector continuation is based on the fact that when a Hamiltonian depends smoothly on some real-valued control parameter(s), any eigenvector is a smooth function of that parameter(s) and its trajectory is confined to a very low-dimensional subspace

[Frame et al., PRL 121, 032501 (2018); König et al., PLB 810, 135814 (2020)]



- ► Write the Hamiltonian in a **linearized** form $H(\vec{c}) = H_0 + \sum c_i H_i$
- Select 'training' points $\{\vec{c}_i\}$ and solve the exact problem $H(\vec{c}_i) |\psi_i\rangle = E_i |\psi_i\rangle$
- ► Project the Hamiltonian onto the subspace of training eigenvectors {|ψ_i⟩} and diagonalize the generalized eigenvalue problem

$$\tilde{\mathsf{H}}(\vec{\mathsf{C}}_{+}) \ket{\tilde{\psi}} = \tilde{\mathsf{E}}_{+} \tilde{\mathsf{N}} \ket{\tilde{\psi}},$$

where $\tilde{H}_{ij} = \langle \psi_i | H(\vec{c}_+) | \psi_j \rangle$, $\tilde{N}_{ij} = \langle \psi_i | \psi_j \rangle$ and \tilde{E}_+ approximates E_+

Emulating ab initio NCSM calculations: eigenvector continuation

Hypernuclear Hamiltonian with LO YN interactions can be linearized,

 $H = H_0 + \frac{C_{27}V_{27}}{V_{27}} + \frac{C_{10*}V_{10*}}{V_{10*}} + \frac{C_{10}V_{10}}{V_{10}} + \frac{C_{8a}V_{8a}}{V_{8a}} + \frac{C_{8s}V_{8s}}{V_{8s}},$

where C_is are the 5 independent $SU_f(3)$ LECs and H_0 contains the kinetic energy, NN+NNN interactions, and hypernuclear meson-exchange and Coulomb interactions



- ► ${}^{5}_{\Lambda}$ He; ${}^{1}_{2}$, model space truncation N_{max} = 12
- Vary one LEC, C₂₇, within ±100 % relative variation with respect to the nominal LOYN(Λ_{YN}=600 MeV) value
- Select 1, 2, 4 exact NCSM eigenvectors to construct the emulators
- Accurate and lighting-fast emulation of ab initio NCSM calculations
- ► Continued eigenvectors stay within the same $(N_{max}, \hbar\omega)$ model space \rightarrow extrapolation of observables to infinite model space is still necessary

Emulating ab initio NCSM calculations: cross validation

- Select 2, 4, 8, 16, 32 points in the 5-dimensional space of LOYN LECs using the Latin hypercube space-filling design in a ±20 % domain around the nominal values to train the emulators
- ► Select randomly 256 exact NCSM calculations within the same domain of LECs



► We can achieve relative accuracy of |δ_{rel}| < 1, 0.1, 0.002 % using 8, 16, 32 training points</p>

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Global sensitivity analysis

- Addresses the question of how variance of the output of a model can be apportioned to variances of the model inputs [Saltelli et al., CPC 181, 259 (2010)]
- ► Allows to identify the most influential LECs of *χ*EFT YN interactions which determine the hypernuclear energy spectra
- For an output Y = f(a) of a model f, we decompose the total variance as

$$\label{eq:Var} \text{Var}\left[Y\right] = \sum_{i=1}^d V_i + \sum_{i < j=1}^d V_{ij} + \cdots \,,$$

where

$$\begin{split} & \mathsf{V}_{i} = \mathsf{Var}\left[\mathsf{E}_{\vec{\alpha} \sim (\alpha_{i})}[\mathsf{Y}|\alpha_{i}]\right], \\ & \mathsf{V}_{ij} = \mathsf{Var}\left[\mathsf{E}_{\vec{\alpha} \sim (\alpha_{i},\alpha_{j})}[\mathsf{Y}|\alpha_{i},\alpha_{j}]\right] - \mathsf{V}_{i} - \mathsf{V}_{j}, \end{split}$$

are variances of conditional expectation of Y

- The variance integrals are computed by using quasi-MC sampling, including 95 % confidence intervals
- The first-, second-, and higher-order (Sobol') sensitivity indices

$$S_i = \frac{V_i}{Var\left[Y\right]}, \quad S_{ij} = \frac{V_{ij}}{Var\left[Y\right]}, \quad \cdot \cdot$$

Total effect

$$S_{Ti} = S_i + S_{ij} + \cdots$$

Identify the most influential LECs:

Y = Λ separation energies of ${}_{\Lambda}^{3}H_{\frac{1}{2}^{+}}$, ${}_{\Lambda}^{4}H_{0^{+}}$, ${}_{\Lambda}^{4}H_{0^{+}}$, ${}_{\Lambda}^{4}H_{1^{+}}$, ${}_{\Lambda}^{4}H_{1^{+}}$, ${}_{\Lambda}^{5}He_{\frac{1}{2}^{+}}$, $\vec{\alpha} =$ the 5 LECs of the LOYN interaction; independent and uniformly distributed within $\pm 2\%$ ($\pm 20\%$) variation around the nominal values of LOYN(Λ_{YN} =600 MeV) for ${}_{\Lambda}^{3}H$ (A = 4, 5)



- $\label{eq:static} \blacktriangleright \ S_i \approx S_{Ti} \rightarrow energies \\ are additive in all \\ LECs$
- C₂₇ is responsible for most of the variation in energy

$$\begin{split} C_{1S_{0}}^{\Lambda\Lambda} &= \frac{1}{10}(9\mathsf{C}_{27} + \mathsf{C}_{8s}) \\ C_{3S_{1}}^{\Lambda\Lambda} &= \frac{1}{2}(\mathsf{C}_{10^{*}} + \mathsf{C}_{8a}) \\ C_{3S_{1}}^{\Sigma\Sigma} &= \mathsf{C}_{10} \end{split}$$

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- Simultaneous fitting of bound-state and scattering observables is inevitable
- ► Can we improve the description of A separation energies in light hypernuclei with a small variation of LO YN LECs?



- Proof-of-principle simple least-squares optimization
 - LECs restricted up to ± 40 % variation around the nominal values of LOYN(A_{YN}=600 MeV)
- ► Theoretical precision σ²_{th} = σ²_{method} + σ²_{model}

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Summary & outlook

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Ab initio calculations of light hypernuclei

► Hypernuclear observables, such as A separation energies in light hypernuclei, suffer from sizable theoretical uncertainties associated with the choice of nuclear interaction (with LO *x*EFT YN interaction)

Emulating ab initio NCSM

- Eigenvector continuation provides fast and accurate emulation of ab initio calculations of light hypernuclei
- ► Global sensitivity analysis identifies **the most influential LECs** of χ EFT YN interactions which **determine the energy spectra** of light hypernuclei
- A significantly better description of energy levels of light hypernuclei can be achieved with a relatively small variation of the LECs

Outlook

Simultaneous optimization of YN interactions using bound-state and scattering observables with accompanying uncertainty quantification Thank you!