

# Dynamics in the inner crust: heavy impurity in Fermi superfluid

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Faculty  
of Physics

WARSAW UNIVERSITY OF TECHNOLOGY

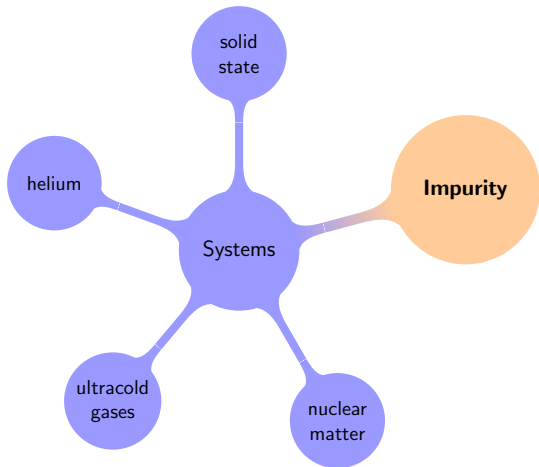


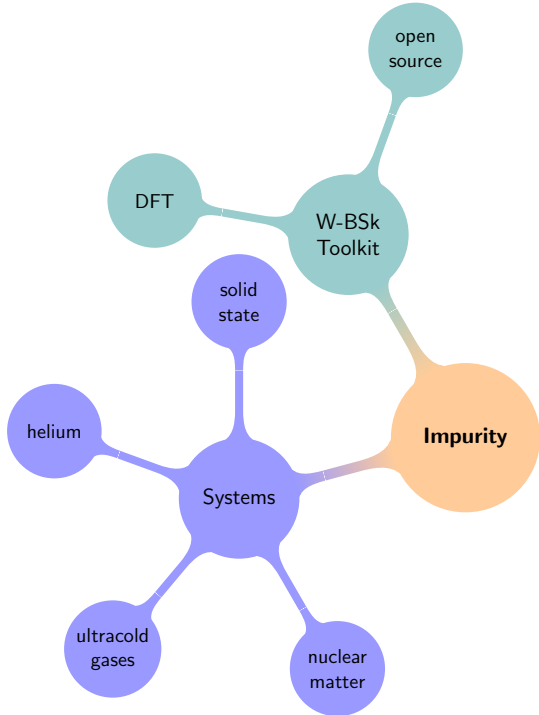
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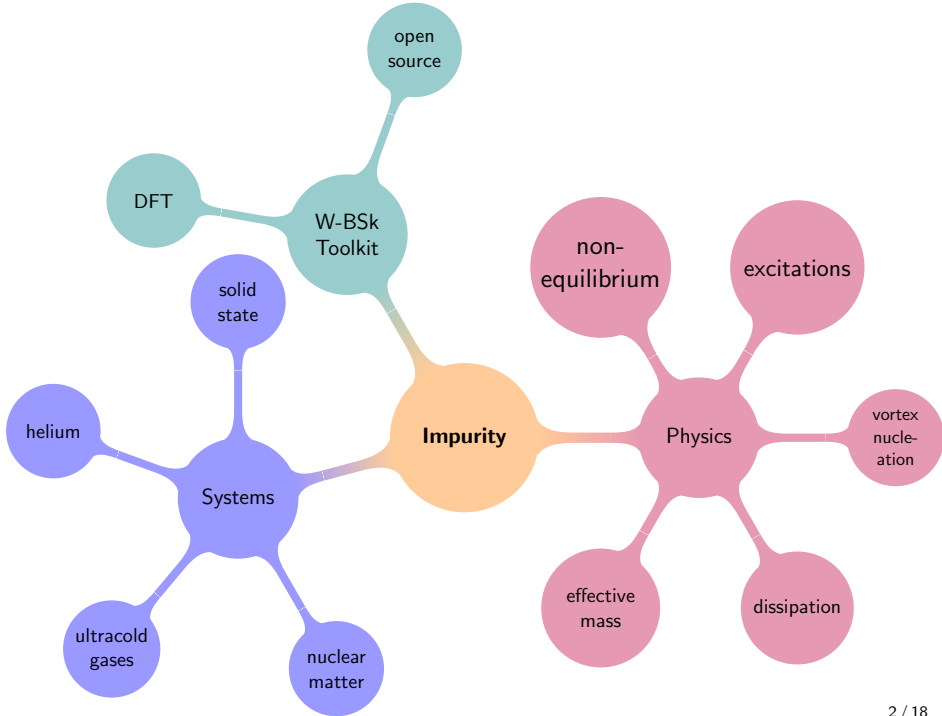


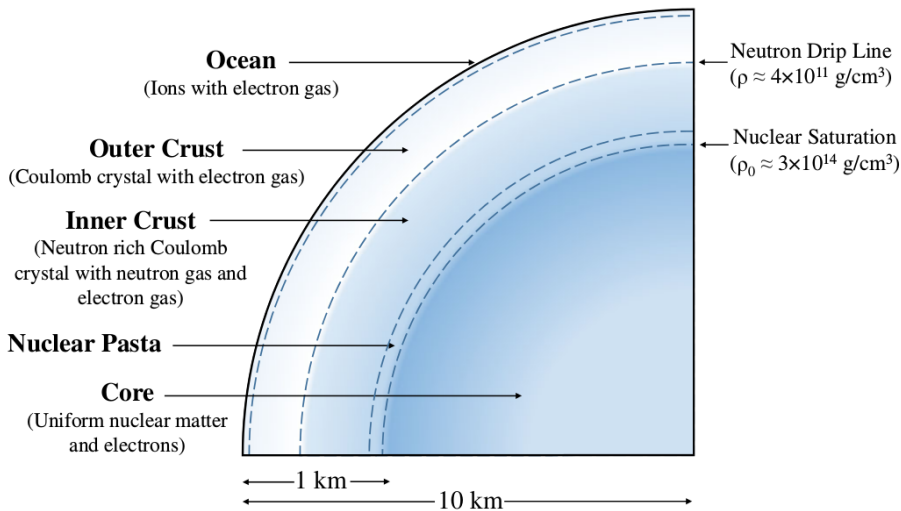
NARODOWE CENTRUM NAUKI

25<sup>th</sup> April 2024, ECT\*, Trento









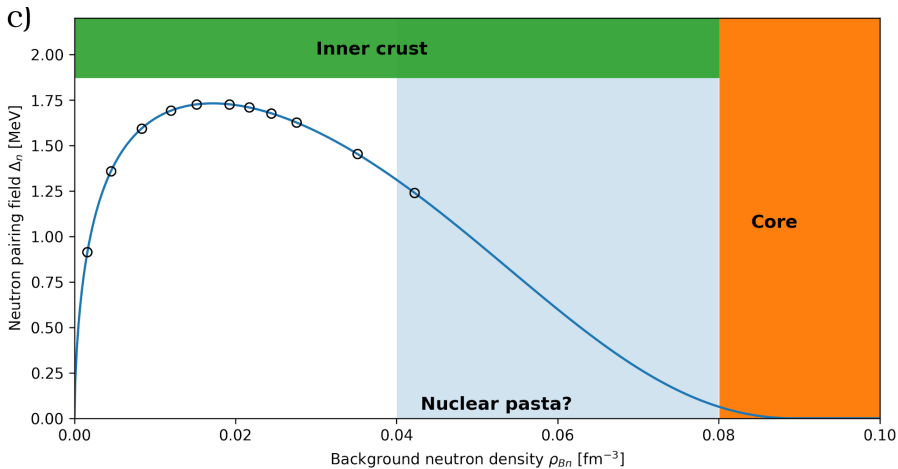
Caplan, M. E., and C. J. Horowitz, *Reviews of Modern Physics* 89, 041002 (2017)

# Crust composition

$\rho_{\max}$ [g cm <sup>-3</sup> ]	Element	Z	N	$R_{\text{cell}}$ [fm]
$8.02 \times 10^6$	<sup>56</sup> Fe	26	30	1404.05
$2.71 \times 10^8$	<sup>62</sup> Ni	28	34	449.48
$1.33 \times 10^9$	<sup>64</sup> Ni	28	36	266.97
$1.50 \times 10^9$	<sup>66</sup> Ni	28	38	259.26
$3.09 \times 10^9$	<sup>86</sup> Kr	36	50	222.66
$1.06 \times 10^{10}$	<sup>84</sup> Se	34	50	146.56
$2.79 \times 10^{10}$	<sup>82</sup> Ge	32	50	105.23
$6.07 \times 10^{10}$	<sup>80</sup> Zn	30	50	80.58
$8.46 \times 10^{10}$	<sup>82</sup> Zn	30	52	72.77
$9.67 \times 10^{10}$	<sup>128</sup> Pd	46	82	80.77
$1.47 \times 10^{11}$	<sup>126</sup> Ru	44	82	69.81
$2.11 \times 10^{11}$	<sup>124</sup> Mo	42	82	61.71
$2.89 \times 10^{11}$	<sup>122</sup> Zr	40	82	55.22
$3.97 \times 10^{11}$	<sup>120</sup> Sr	38	82	49.37
$4.27 \times 10^{11}$	<sup>118</sup> Kr	36	82	47.92

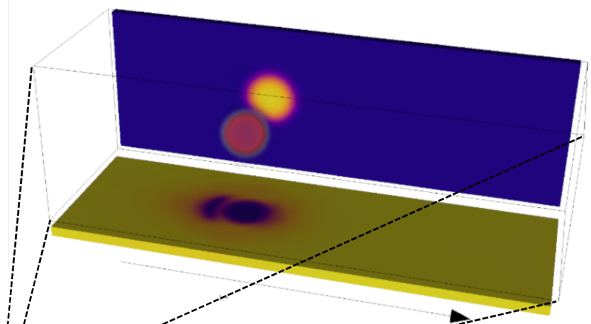
$\rho$ [g cm <sup>-3</sup> ]	Element	Z	N	$R_{\text{cell}}$ [fm]
$4.67 \times 10^{11}$	<sup>180</sup> Zr	40	140	53.60
$6.69 \times 10^{11}$	<sup>200</sup> Zr	40	160	49.24
$1.00 \times 10^{12}$	<sup>250</sup> Zr	40	210	46.33
$1.47 \times 10^{12}$	<sup>320</sup> Zr	40	280	44.30
$2.66 \times 10^{12}$	<sup>500</sup> Zr	40	460	42.16
$6.24 \times 10^{12}$	<sup>950</sup> Sn	50	900	39.32
$9.65 \times 10^{12}$	<sup>1100</sup> Sn	50	1050	35.70
$1.49 \times 10^{13}$	<sup>1350</sup> Sn	50	1300	33.07
$3.41 \times 10^{13}$	<sup>1800</sup> Sn	50	1750	27.61
$7.94 \times 10^{13}$	<sup>1500</sup> Zr	40	1460	19.61
$1.32 \times 10^{14}$	<sup>982</sup> Ge	32	950	14.38

Chamel, Nicolas, and Pawel Haensel. "Physics of neutron star crusts." Living Reviews in relativity 11.1 (2008): 1-182.

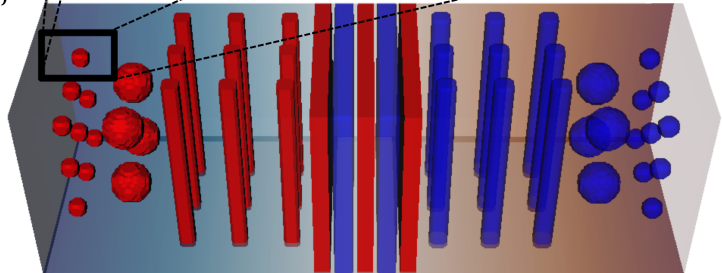


# Nonequilibrium dynamics

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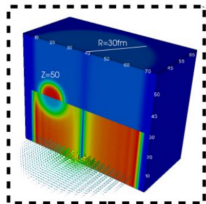


b)

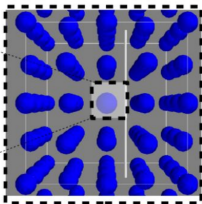




10fm

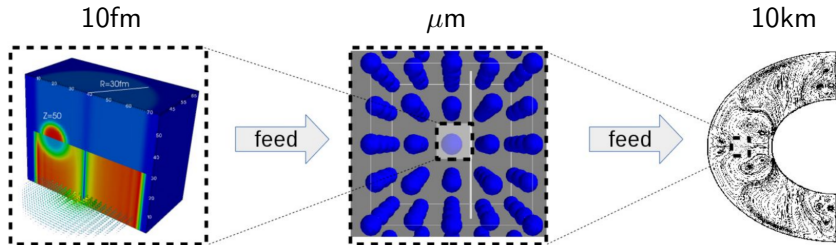


$\mu\text{m}$



10km

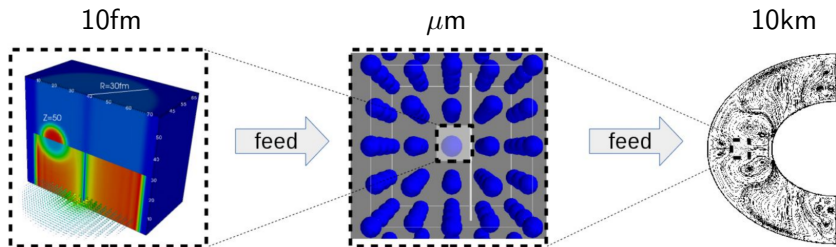




LUMI, Finland (#5 Top 500)



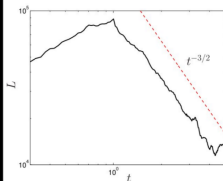
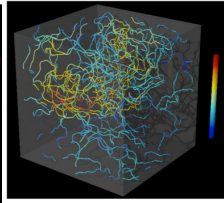
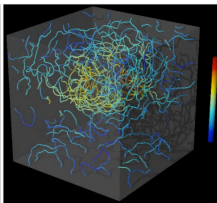
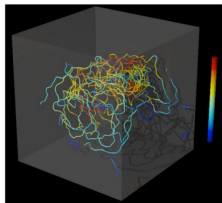
Piz Daint, Switzerland (#37 Top 500)



LUMI, Finland (#5 Top 500)



Piz Daint, Switzerland (#37 Top 500)



# Warsaw University of Technology

# W-SLDA Toolkit W-BSk Toolkit

## W-SLDA Toolkit

Self-consistent solver  
of mathematical problems  
which have structure  
formally equivalent to  
Bogoliubov-de Gennes equations.

static problems: st-wslida

$$\begin{pmatrix} h_a(\mathbf{r}) - \mu_a & \Delta(\mathbf{r}) \\ \Delta^*(\mathbf{r}) & -h_b^*(\mathbf{r}) + \mu_b \end{pmatrix} \begin{pmatrix} u_n(\mathbf{r}) \\ v_n(\mathbf{r}) \end{pmatrix} = E_n \begin{pmatrix} u_n(\mathbf{r}) \\ v_n(\mathbf{r}) \end{pmatrix}$$

time-dependent problems: td-wslida

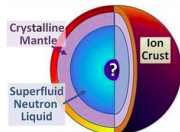
$$i\hbar \frac{\partial}{\partial t} \begin{pmatrix} u_n(\mathbf{r}, t) \\ v_n(\mathbf{r}, t) \end{pmatrix} = \begin{pmatrix} h_a(\mathbf{r}, t) - \mu_a & \Delta(\mathbf{r}, t) \\ \Delta^*(\mathbf{r}, t) & -h_b^*(\mathbf{r}, t) + \mu_b \end{pmatrix} \begin{pmatrix} u_n(\mathbf{r}, t) \\ v_n(\mathbf{r}, t) \end{pmatrix}$$

Extension to nuclear matter  
in neutron stars

Unified solvers for static and  
time-dependent problems

Dimensionalities of  
problems: 3D, 2D and 1D

## Extension to nuclear matter in neutron stars



The W-SLDA Toolkit has been expanded to encompass  
available as the W-BSk Toolkit.

ALL FUNCTIONALITIES →

## Getting the code

 [DOWNLOAD](#)

The W-SLDA & W-BSk Toolkits are free to download. It is published as open source under GNU GPL License. In order to get W-SLDA or W-BSk Toolkit click "Read more" and follow instructions.

[READ MORE →](#)

Integration with VisIt:  
visualization, animation and  
analysis tool



» [Contact us](#)

» [Contributing to W-SLDA](#)

$$\varepsilon(\rho, \vec{\nabla}\rho, \nu, \tau, \mathbf{j}) = \frac{\hbar^2}{2M}\tau + \varepsilon_\rho(\rho) + \varepsilon_\tau(\rho, \tau, \mathbf{j}) + \varepsilon_{\Delta\rho}(\rho, \vec{\nabla}\rho) + \varepsilon_\pi(\rho, \vec{\nabla}\rho, \nu)$$

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$$h(r) = \frac{\delta\varepsilon}{\delta\rho} - \nabla \frac{\delta\varepsilon}{\delta\tau} \nabla - \frac{i}{2} \left\{ \frac{\delta\varepsilon}{\delta\mathbf{j}}, \nabla \right\}$$

$$\Delta(r) = \frac{\delta\varepsilon}{\delta\nu}$$

## Superfluid **L**ocal **D**ensity **A**pproximation

A. Bulgac, Physical Review A **76**, 040502 (2007)

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## Superfluid **Local Density Approximation**

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### Hartree-Fock-Bogoliubov equations

$$\begin{pmatrix} h(r) & \Delta(r) \\ \Delta^*(r) & -h^*(r) \end{pmatrix} \begin{pmatrix} u_k(r) \\ v_k(r) \end{pmatrix} = \epsilon_k \begin{pmatrix} u_k(r) \\ v_k(r) \end{pmatrix}$$



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Quality of results highly depends on the quality of density functional!

### 1. Energy density

Under the assumption of invariance under time reversal, the HFB energy is written as the integral of a purely local energy-density functional

$$E_{\text{HFB}} = \int \mathcal{E}_{\text{HFB}}(\mathbf{r}) d^3\mathbf{r}, \quad (\text{A1})$$

where

$$\begin{aligned} \mathcal{E}_{\text{HFB}}(\mathbf{r}) = & \mathcal{E}_{\text{Sky}}[\rho_n(\mathbf{r}), \nabla\rho_n(\mathbf{r}), \tau_n(\mathbf{r}), \mathbf{J}_n(\mathbf{r}), \rho_p(\mathbf{r}), \nabla\rho_p(\mathbf{r}), \\ & \tau_p(\mathbf{r}), \mathbf{J}_p(\mathbf{r})] + \mathcal{E}_{\text{Coul}}[\rho_p(\mathbf{r})] \\ & + \mathcal{E}_{\text{pair}}[\rho_n(\mathbf{r}), \bar{\rho}_n(\mathbf{r}), \rho_p(\mathbf{r}), \bar{\rho}_p(\mathbf{r})]. \end{aligned} \quad (\text{A2})$$

The first term here, the energy density for the Skyrme force of this paper, is given by

$$\begin{aligned} \mathcal{E}_{\text{Sky}} = & \sum_{q=n,p} \frac{\hbar^2}{2M_q} \tau_q + \frac{1}{2} t_0 \left[ \left(1 + \frac{1}{2} x_0\right) \rho^2 - \left(\frac{1}{2} + x_0\right) \sum_{q=n,p} \rho_q^2 \right] + \frac{1}{4} t_1 \left\{ \left(1 + \frac{1}{2} x_1\right) \left[ \rho\tau + \frac{3}{4} (\nabla\rho)^2 \right] \right. \\ & \left. - \left(\frac{1}{2} + x_1\right) \sum_{q=n,p} \left[ \rho_q \tau_q + \frac{3}{4} (\nabla\rho_q)^2 \right] \right\} + \frac{1}{4} t_2 \left\{ \left(1 + \frac{1}{2} x_2\right) \left[ \rho\tau - \frac{1}{4} (\nabla\rho)^2 \right] + \left(\frac{1}{2} + x_2\right) \right. \\ & \left. \times \sum_{q=n,p} \left[ \rho_q \tau_q - \frac{1}{4} (\nabla\rho_q)^2 \right] \right\} + \frac{1}{12} t_3 \rho^\alpha \left[ \left(1 + \frac{1}{2} x_3\right) \rho^2 - \left(\frac{1}{2} + x_3\right) \sum_{q=n,p} \rho_q^2 \right] \\ & + \frac{1}{4} t_4 \left\{ \left(1 + \frac{1}{2} x_4\right) \left[ \rho\tau + \frac{3}{4} (\nabla\rho)^2 \right] - \left(\frac{1}{2} + x_4\right) \sum_{q=n,p} \left[ \rho_q \tau_q + \frac{3}{4} (\nabla\rho_q)^2 \right] \right\} \rho^\beta \\ & + \frac{\beta}{8} t_4 \left[ \left(1 + \frac{1}{2} x_4\right) \rho (\nabla\rho)^2 - \left(\frac{1}{2} + x_4\right) \nabla\rho \cdot \sum_{q=n,p} \rho_q \nabla\rho_q \right] \rho^{\beta-1} + \frac{1}{4} t_5 \left\{ \left(1 + \frac{1}{2} x_5\right) \left[ \rho\tau - \frac{1}{4} (\nabla\rho)^2 \right] \right. \\ & \left. + \left(\frac{1}{2} + x_5\right) \sum_{q=n,p} \left[ \rho_q \tau_q - \frac{1}{4} (\nabla\rho_q)^2 \right] \right\} \rho^\gamma - \frac{1}{16} (t_1 x_1 + t_2 x_2) J^2 + \frac{1}{16} (t_1 - t_2) \sum_{q=n,p} J_q^2 \\ & - \frac{1}{16} (t_4 x_4 \rho^\beta + t_5 x_5 \rho^\gamma) J^2 + \frac{1}{16} (t_4 \rho^\beta - t_5 \rho^\gamma) \sum_{q=n,p} J_q^2 + \frac{1}{2} W_0 \left( \mathbf{J} \cdot \nabla\rho + \sum_{q=n,p} \mathbf{J}_q \cdot \nabla\rho_q \right). \end{aligned} \quad (\text{A3})$$

# Brussels-Montreal (BSk) functional

## **Experimental data**

- atomic masses
- nuclear charge radii
- symmetry energy
- incompressibility

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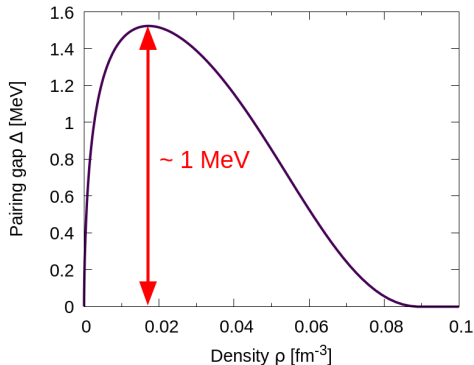
## N-body calculations

- EoS of pure neutron matter
- $^1S_0$  pairing gaps in nuclear matter
- effective masses in nuclear matter

Chamel et al., Phys. Rev. C **80**, 065804 (2009)

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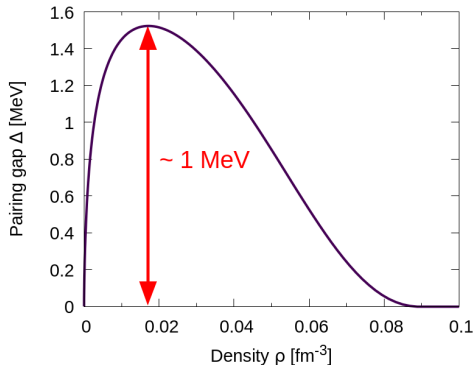
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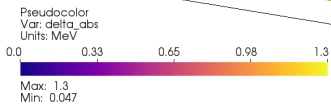
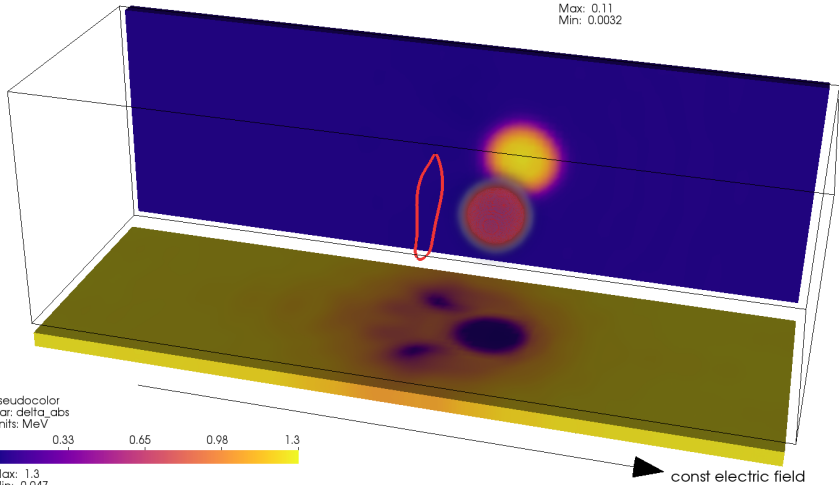
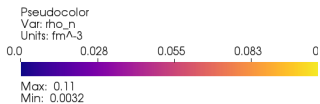
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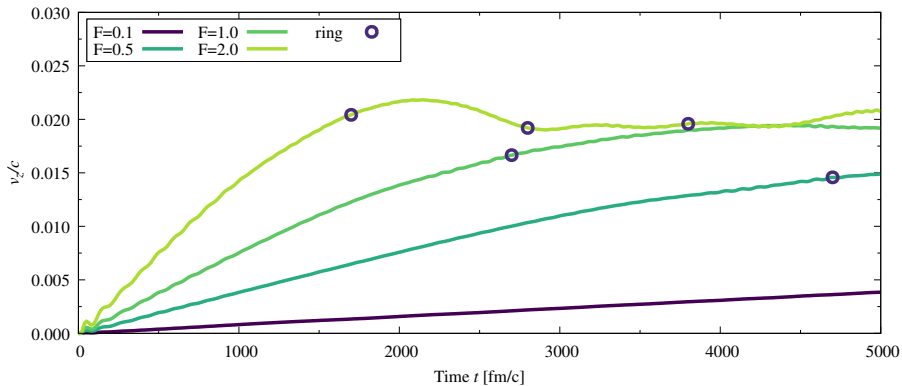
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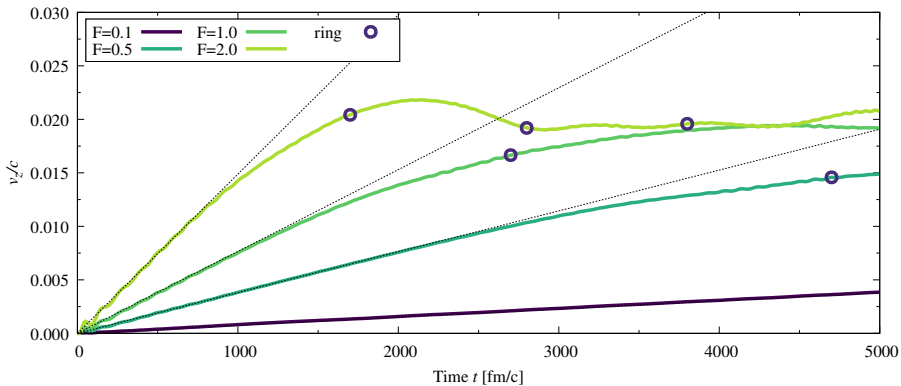
Goriely et al., Phys. Rev. C **93**, 034337 (2016)



$$\varepsilon(\rho_q, \vec{\nabla}\rho_q, \nu_q, \tau_q, \mathbf{j}_q)$$



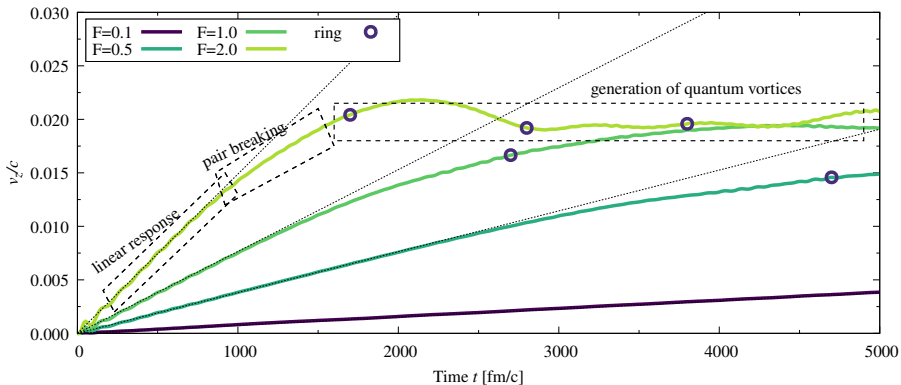




$$v_z = a_z t$$

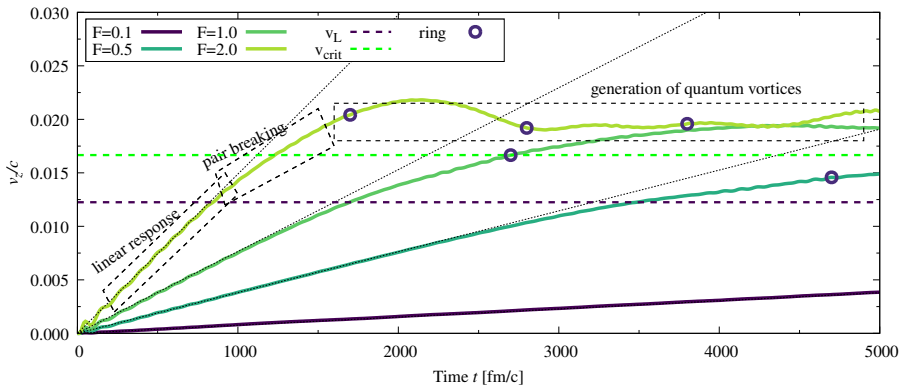
$$M_{eff}^{(d)} = \frac{F}{a_z}$$





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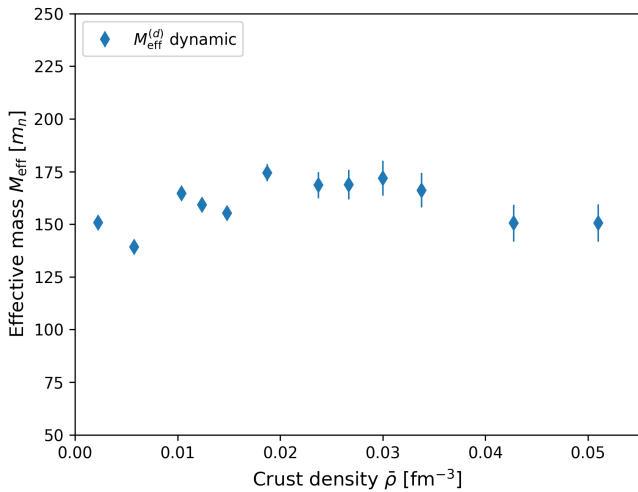


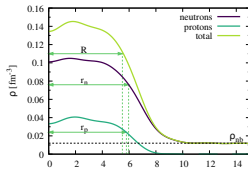
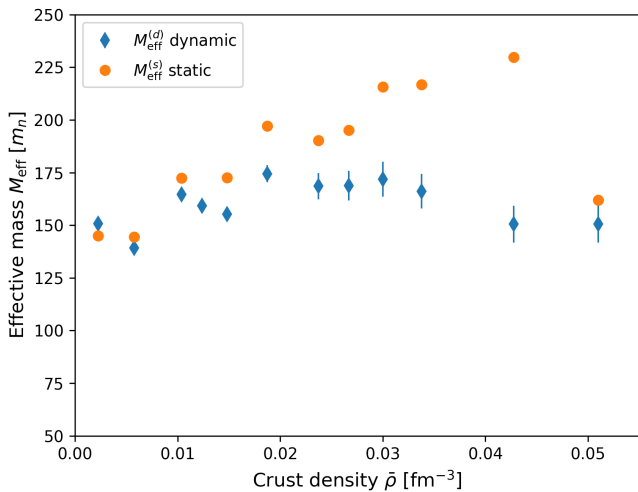
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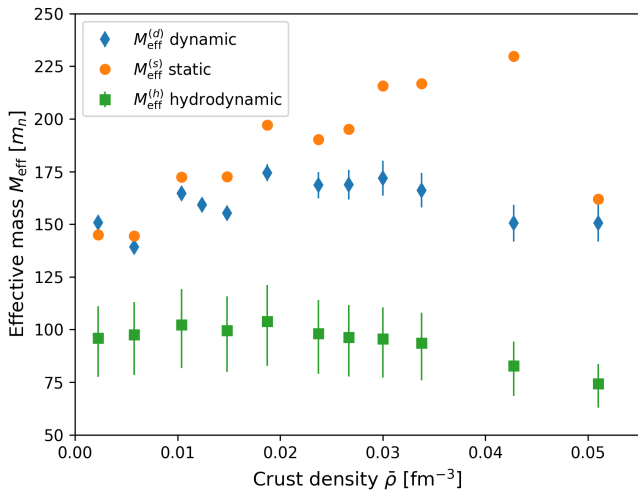
$$M_{eff}^{(d)} = \frac{F}{a_z}$$

$$v_L = \frac{\Delta}{\hbar k_F}$$

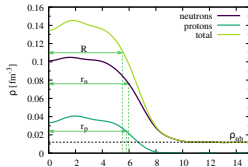
$$v_{crit} = \frac{e \Delta}{2 \hbar k_F}$$

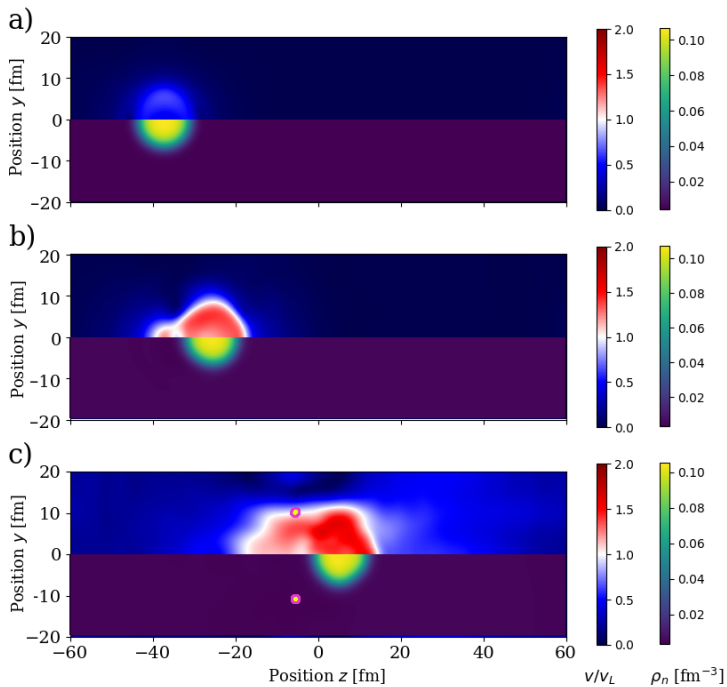




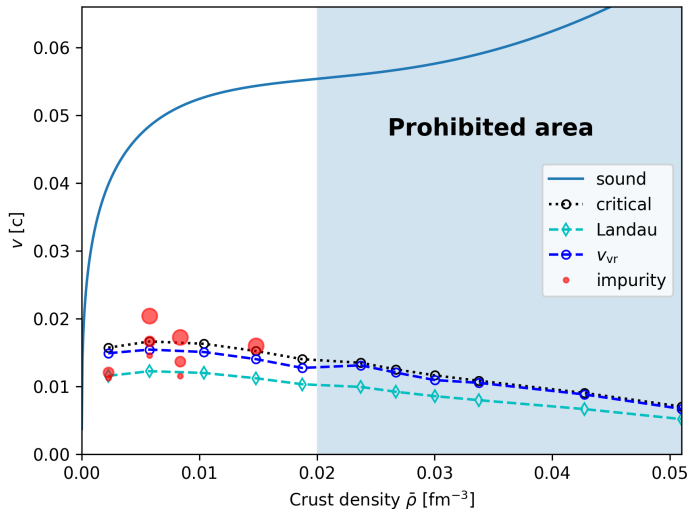


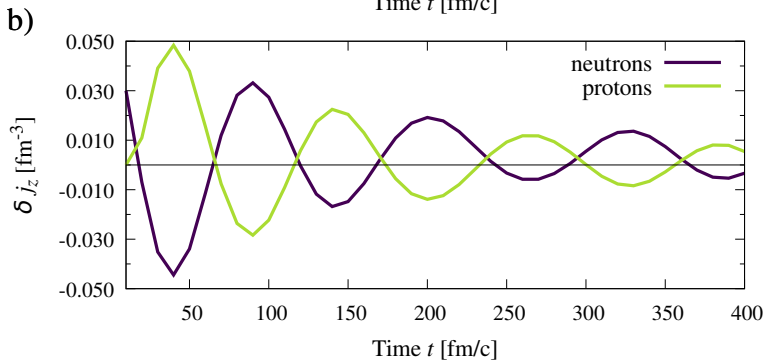
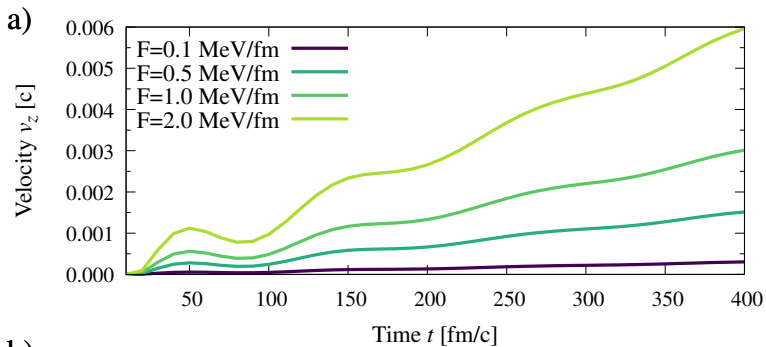
$$M_h = \frac{4}{3} \pi R^3 M \frac{(\rho_{\text{in}} - \rho_{\text{out}})^2}{\rho_{\text{in}} + 2\rho_{\text{out}}}$$





$$v_{vr} = \frac{1}{2\pi R} \frac{\hbar c}{M_n c^2} \left( \ln \frac{8R}{a_{core}} - \alpha \right)$$



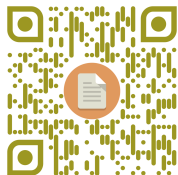




# Summary

- fully self-consistent 3D (TD)HFB calculations
- BSk31 Energy Density Functional
- effective parameters can be extracted
- effective mass
- dissipation channels
- creating vortex rings
- giant dipole resonance

More details



arxiv:2403.17499

# Thank you!