

Chiral effective field theory for nuclear forces and the dense matter equation of state

Isak Svensson

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Outline

- ① Chiral EFT (χ EFT) and nuclear forces
- ② Bayesian uncertainty quantification (UQ) in ab initio nuclear theory
- ③ Many-body perturbation theory (MBPT) calculations of nuclear matter
- ④ Posterior predictive distributions for nuclear matter (preliminary!)
- ⑤ New χ EFT bands for dense matter equation of state (EOS)

Making predictions in nuclear theory

The time-independent Schrödinger equation:

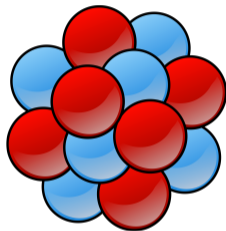
$$(\hat{H}_0 + \hat{V}) |\Psi\rangle = E |\Psi\rangle$$

We need:

- 1 a model for the interaction potential \hat{V} :
chiral effective field theory
- 2 a many-body method for solving the S.E.

2-body scattering: Solving Lippmann-Schwinger equation

Many-body methods: NCSM, QMC, ... (light systems, $A \lesssim 16$), CC, IMSRG, MBPT, ... (not-so-light systems)



χ EFT and nuclear forces

χ EFT:

- Systematic expansion in low momenta:
 $(Q/\Lambda_b)^k$
- Power counting: assigns each contribution to an order k^a
- Orders designated leading order (LO), next-to-leading order (NLO), N^2 LO, N^3 LO, ...
- Many-body forces enter consistently at sub-leading orders

^aNo contributions for $k = 1$.

(Weinberg, van Kolck, Kaplan, Savage, Wise, Bernard, Epelbaum, Kaiser, Meißner, ...)

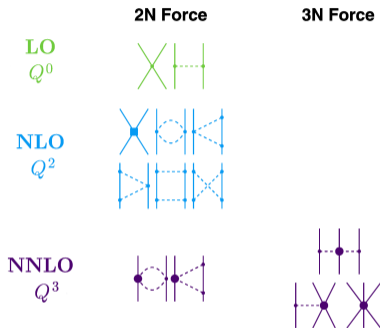


Figure adapted from Entem et al., Phys. Rev. C 96 (2017).

χ EFT and nuclear forces

- Two main variants: Δ -less (previous slide) and Δ -full
- Degrees of freedom in Δ -less: nucleons and pions
- Additionally in Δ -full: $\Delta(1232)$ -isobar

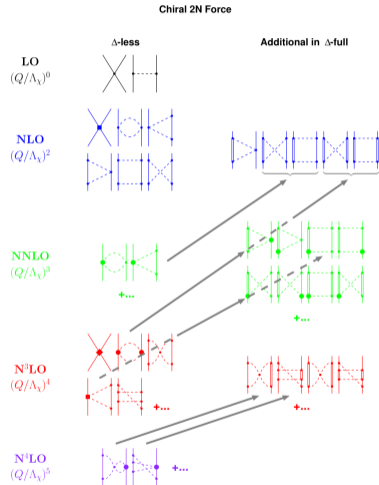


Figure from Machleidt & Entem,
Phys. Rept. 503 (2011).

χ EFT and nuclear forces - UQ

χ EFT (in principle) enables **uncertainty quantification**:

- Each order suppressed by $\sim (Q/\Lambda_b) < 1$ (gives a handle on **truncation errors**)
- Short-range physics accounted for by **unknown low-energy constants (LECs)**
- Number of LECs grows with order: ~ 15 at $N^2\text{LO}$, ~ 30 at $N^3\text{LO}$
- LECs ($\vec{\alpha}$) fitted to scattering and other nuclear observables

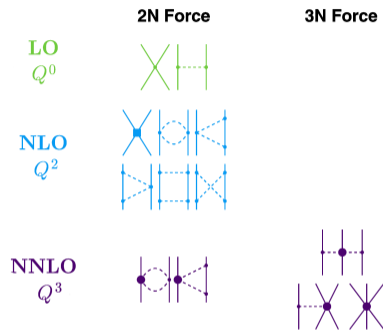


Figure adapted from Entem et al., Phys. Rev. C 96 (2017).

Bayesian UQ for χ EFT pioneered by the BUQEYE collaboration: Furnstahl, Melendez, Phillips, Wesolowski ...

Bayesian UQ in ab initio nuclear theory

Based on my PhD work in collaboration with Andreas Ekström and Christian Forssén (and BUQEYE): Svensson et al., Phys. Rev. C 105 (2022), Phys. Rev. C 107 (2023), arXiv:2304:02004; Wesolowski et al., Phys. Rev. C 104 (2021)

Bayesian UQ in ab initio nuclear theory

Most common approach to fitting LECs: optimization to (mainly 2-body) nuclear observables/phaseshifts

Has yielded many accurate interactions. To mention a few: Entem & Machleidt Phys. Rev. C 68 (2003), Hebeler et al. Phys. Rev. C 83 (2011), Carlsson et al. Phys. Rev. X 6 (2016), Jiang et al. Phys. Rev. C 102 (2020)

But rigorous UQ is lacking

Bayesian inference and predictions

We make predictions of y using a **posterior predictive distribution (PPD)**:

$$\text{pr}(y|D, I) = \int \text{pr}(y|\vec{\alpha}, I) \text{pr}(\vec{\alpha}|D, I) d\vec{\alpha}$$

For this we need the joint **posterior** for the LECs $\text{pr}(\vec{\alpha}|D, I)$. **Bayes' theorem**:

$$\underset{\text{Posterior}}{\text{pr}(\vec{\alpha}|D, I)} \propto \underset{\text{Likelihood}}{\text{pr}(D|\vec{\alpha}, I)} \times \underset{\text{Prior}}{\text{pr}(\vec{\alpha}|I)}$$

We include experimental errors and truncation errors¹ in our analyses. Our priors are grounded in EFT.

¹Both uncorrelated and (in the latest paper) correlated.

(Breaking) the curse of dimensionality

Problem: $\text{pr}(\vec{\alpha}|D, I)$ is multidimensional ($\sim 15\text{-}30$ parameters). Must use Markov chain Monte Carlo (MCMC).

Even with MCMC, sampling $\text{pr}(\vec{\alpha}|D, I)$ is very challenging due to (i) the dimensionality and (ii) computational cost of calculating observables.

Our approach: use **Hamiltonian Monte Carlo**² (HMC), which is uniquely suited to high-dimensional problems

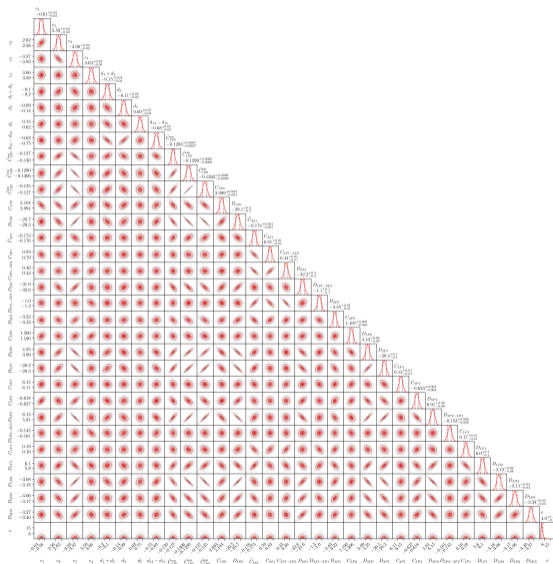
HMC uses **gradients** of the posterior to increase sampling efficiency.

We have found that HMC is \sim **5 times more efficient** than the popular Emcee³ in our application.

²Duane et al., Phys. Lett. B **195**(2) (1987)

³Foreman-Mackey et al., PASP **125** (2013)

Demonstrating the sampling capabilities of HMC



31-dimensional $N^3\text{LO}$
LEC posterior [Phys.
Rev. C 107 (2023)].

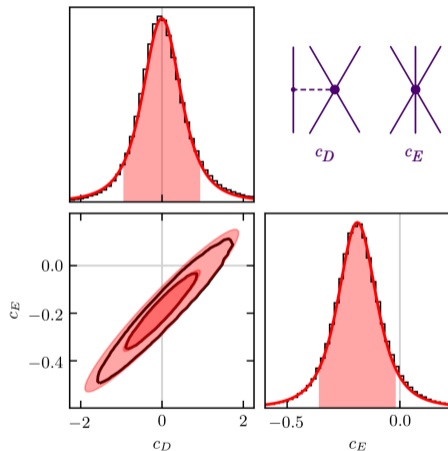
I can zoom in if you
want to see details!

Inferring three-nucleon forces

Three-nucleon forces play an essential role in the description of many-body systems.

In collaboration with BUQEYE we have inferred the two leading $3N$ LECs (c_D, c_E).

Practically usable data are rather lacking as many observables provide degenerate constraints.



Posterior for three-nucleon force LECs c_D, c_E [Wesolowski, IS, et al., Phys. Rev. C 104 (2021)]

Some take-aways

Fully Bayesian UQ is now possible in nuclear theory.

but

Much work remains on accurate error modeling⁴. The fixed-LEC interactions mentioned earlier provide more reliable results.

⁴See, e.g., BUQEYE: Millican et al., 2402.13165.

Goal: combine EOS calculations with Bayesian UQ

MBPT calculations of nuclear matter EOS by Keller et al.:

PHYSICAL REVIEW LETTERS **130**, 072701 (2023)


Nuclear Equation of State for Arbitrary Proton Fraction and Temperature Based on Chiral Effective Field Theory and a Gaussian Process Emulator

J. Keller^{1,2,*}, K. Hebeler^{1,2,3,†} and A. Schwenk^{1,2,3,‡}

¹*Technische Universität Darmstadt, Department of Physics, 64289 Darmstadt, Germany*

²*ExtreMe Matter Institute EMMI, GSI Helmholtzzentrum für Schwerionenforschung GmbH,
64291 Darmstadt, Germany*

³*Max-Planck-Institut für Kernphysik, Saupfercheckweg 1, 69117 Heidelberg, Germany*

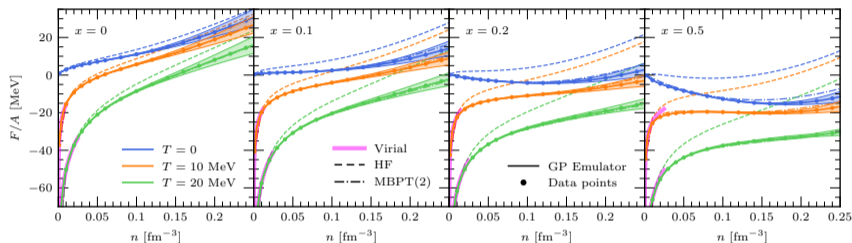
 (Received 5 May 2022; revised 9 December 2022; accepted 12 January 2023; published 17 February 2023)

See also previous work by Christian Drischler

MBPT calculations of nuclear matter EOS at N³LO

Pure neutron matter

Symmetric nuclear matter



Energy per particle as a function of number density for temperature $T = 0, 10, 20$ MeV and proton fraction $x = 0.0, 0.1, 0.2, 0.5$.

Uncertainty bands using the **EKM prescription**⁵ (i.e., not a Bayesian approach):

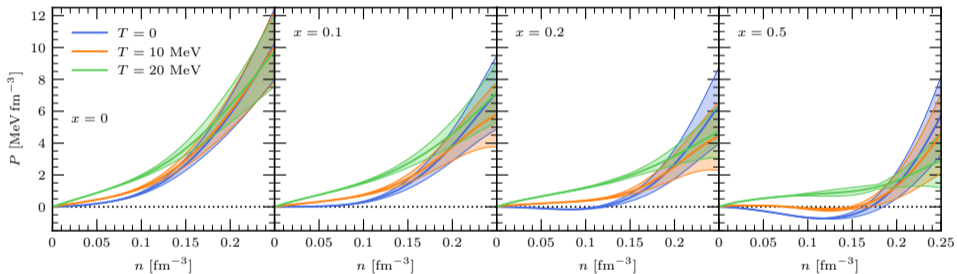
$$\Delta y^{(k)} = \frac{Q}{\Lambda_b} \max \left(|y^{(k)} - y^{(k-1)}|, \Delta y^{(k-1)} \right)$$

⁵Epelbaum et al., Eur. Phys. J. A 51 (2015)

MBPT calculations of nuclear matter EOS at N³LO

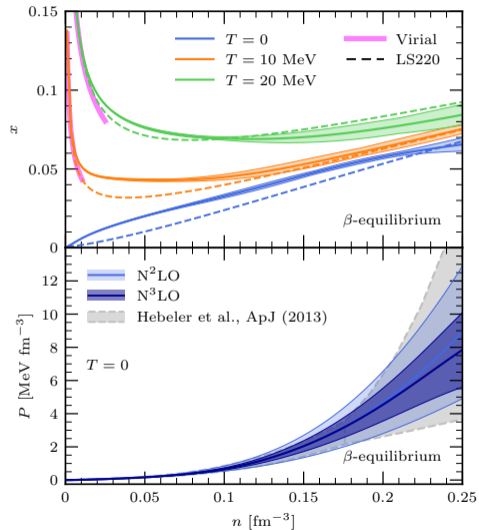
Pure neutron matter

Symmetric nuclear matter



Pressure as a function of number density for temperature $T = 0, 10, 20$ MeV and proton fraction $x = 0.0, 0.1, 0.2, 0.5$.

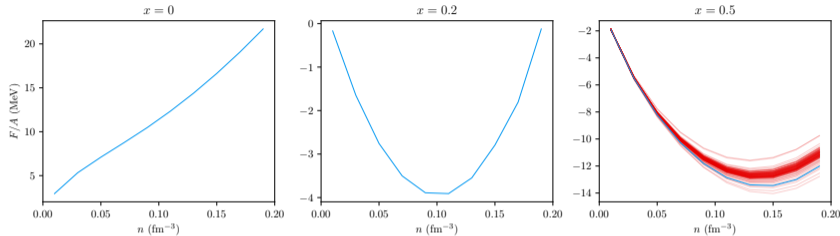
MBPT calculations of nuclear matter EOS at N³LO



Top: Proton fraction in β -equilibrium as a function of density

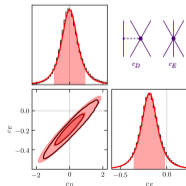
Bottom: Pressure in β -equilibrium as a function of density
N²LO and N³LO bands up to $1.5n_0$
Hebel et al. up to $1.1n_0$, then a piecewise polytrope high-density parametrization

Preliminary: PPD for nuclear matter EOS



Red: PPD for the energy per particle for symmetric nuclear matter at zero temperature.
Blue: results from Keller et al. 2023.

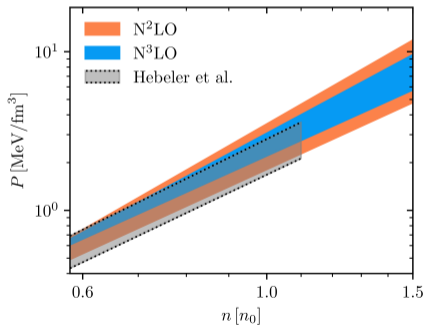
Ongoing work with Achim Schwenk, Kai Hebeler, Hannah Götting, Alex Tichai: PPDs for nuclear matter EOS including LEC variations, correlated truncation errors, MBPT method error. Arbitrary proton fraction and temperature.



Constraining the dense matter equation of state with new NICER mass-radius measurements and new chiral effective field theory inputs

NATHAN RUTHERFORD ¹, MELISSA MENDES ^{2,3,4}, ISAK SVENSSON ^{2,3,4}, ACHIM SCHWENK ^{2,3,4}, ANNA L. WATTS ⁵,
KAI HEBELER ^{2,3,4}, JONAS KELLER ^{2,3}, CHANDA PRESCOD-WEINSTEIN ¹, DEVARSHI CHOUDHURY ⁵,
GEERT RAALJMAKERS ⁶, TUOMO SALMI ⁵, PATRICK TIMMERMAN ⁵, SERENA VINCIGUERRA ⁵,
SEBASTIEN GUILLOT ^{7,8} AND JAMES M. LATTIMER ⁹

Submitted to ApJL a few weeks ago; see Melissa's talk tomorrow



Pressure as a function of density for matter in β -equilibrium.

- New bands include muons in addition to electrons and neutrons/protons
- We trust χ EFT to higher density ($1.5n_0$)
- New bands calculated directly in β -equilibrium; Hebeler bands use an empirical parametrization
- Plan: map bands to LECs

Outlook

- Simultaneous Bayesian inference for 2- and 3-body forces
- Improved modeling of errors—lots of work remains
- Improved UQ for nuclear matter calculations with correlated truncation errors using Gaussian processes (talk to Hannah Göttling!)
- Improved inferences of neutron star properties as new data become available (see Melissa's talk)

Thank you! Collaborators:

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Jordan Melendez

Melissa Mendes
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Nathan Rutherford
Achim Schwenk
Alex Tichai
Anna L. Watts
Sarah Wesolowski
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