

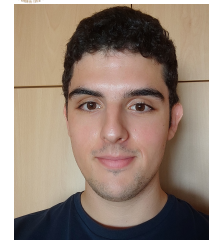
# Neural Quantum States for Quantum Many-Body Physics



*JWT Keeble*



*M. Drissi*



*J. Rozalén Sarmiento*



*A. Rojo-Francàs*



*B. Julià-Díaz*

Dr Arnau Rios Huguet

Institute of Cosmos Sciences  
Universitat de Barcelona

**ECT\* Workshop**  
18 April 2024

## ● Nuclear physics

## ● Neural Quantum States

## ● Deuteron

JWT Keeble & A. Rios, Phys. Lett. B **809**, 135743 (2020)

J Rozalen-Sarmiento, JWT Keeble & A. Rios, Eur. Phys. J. Plus **139**, 189 (2024)

## ● 1D spinless (polarised) systems ( $A > 2$ )

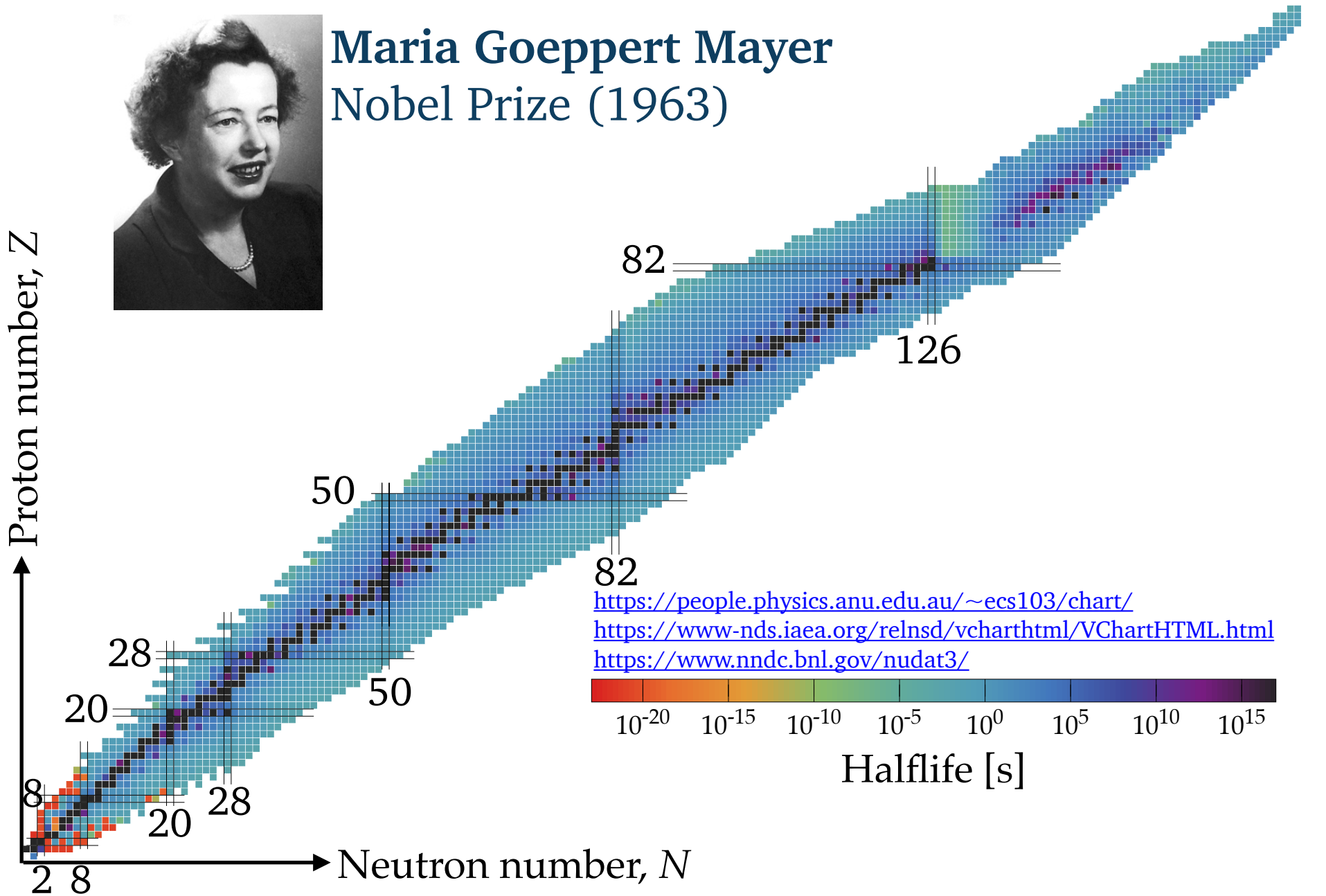
J.W.T. Keeble, M. Drissi, A. Rojo-Francàs, B. Juliá-Díaz & A. Rios, Phys. Rev. A **108**, 063320 (2023),  
arxiv:2304.04725

## ● Optimisation strategy

M. Drissi, J.W.T. Keeble, J. Rozalen-Sarmiento & A. Rios, arxiv:2401.17550



**Maria Goeppert Mayer**  
Nobel Prize (1963)



# Ab initio for isotopes

## $\chi$ -EFT interactions

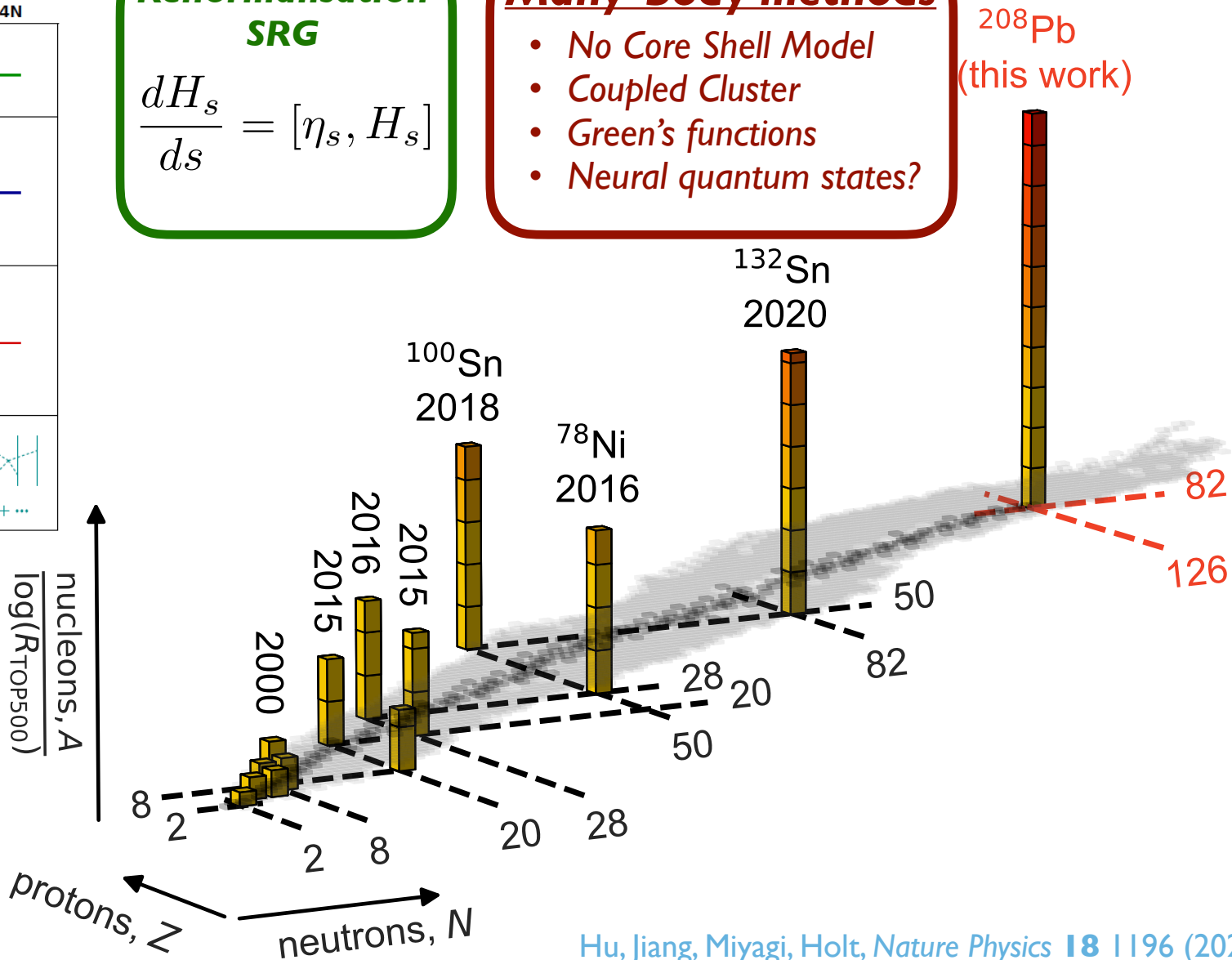
	NN	3N	4N
LO		—	—
NLO		—	—
N <sup>2</sup> LO			—
N <sup>3</sup> LO			

## Renormalisation SRG

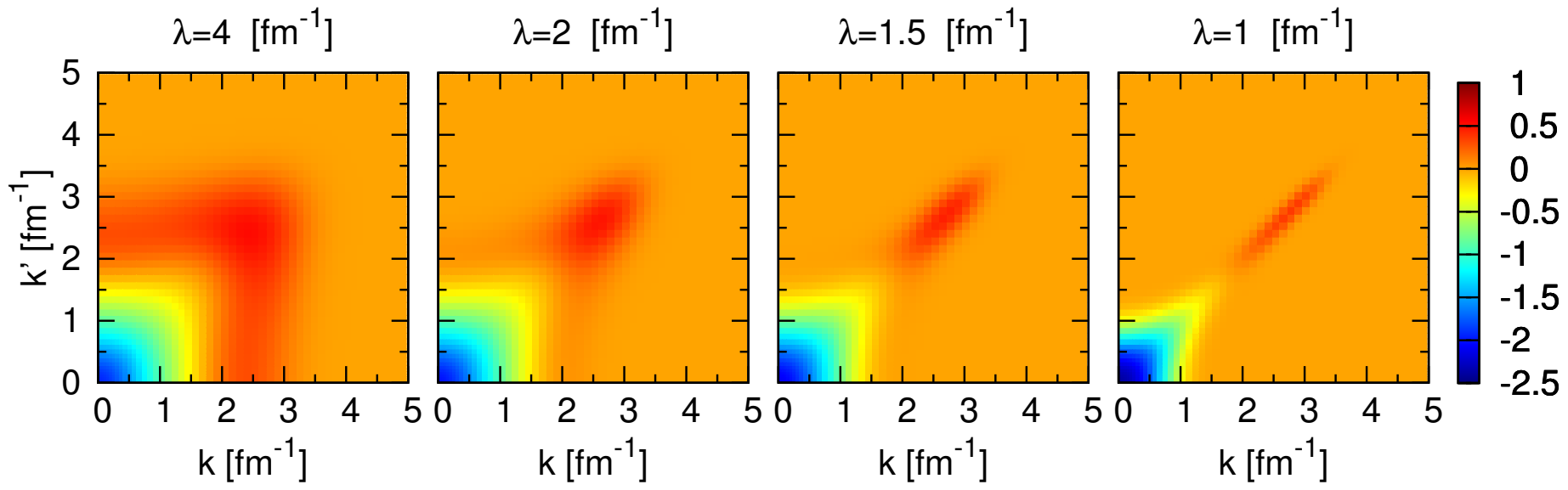
$$\frac{dH_s}{ds} = [\eta_s, H_s]$$

## Many-body methods

- No Core Shell Model
- Coupled Cluster
- Green's functions
- Neural quantum states?



## $^1S_0$ NN matrix elements from N3LO



$$H_s = U(s) H U^\dagger(s) \quad \frac{dH_s}{ds} = [\eta_s, H_s] = [[T_{\text{rel}}, H_s], H_s] \quad \lambda = s^{-1/4}$$

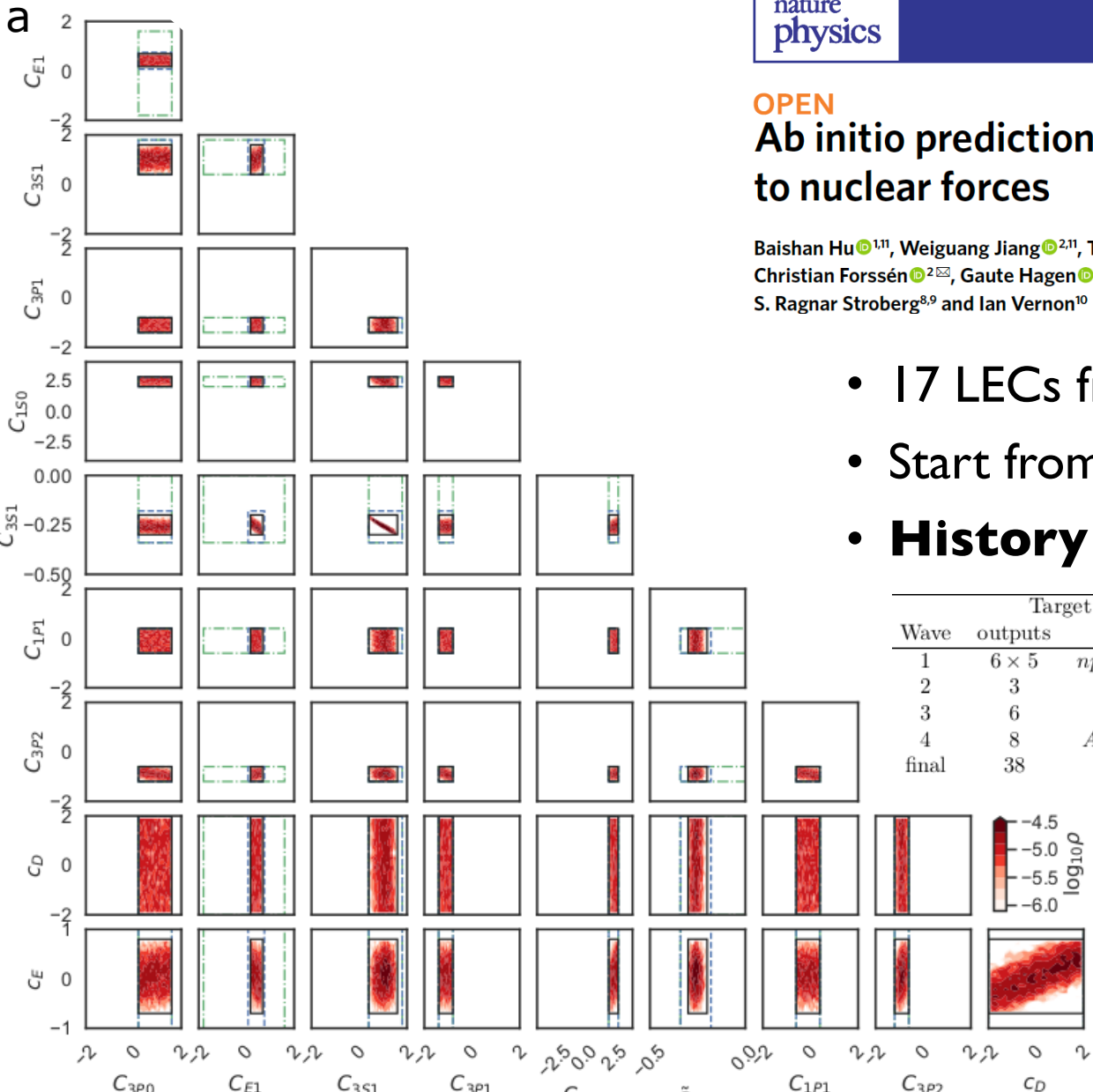
- Series of **unitary** transformation
- Observables **unaltered** at 2-body level, but...
- **Induces** 3-, 4- and up to A-body forces
- **If** these can be treated perturbatively, calculation is **easier**
- **Fixed point?**

**OPEN**  
**Ab initio predictions link the neutron skin of  $^{208}\text{Pb}$  to nuclear forces**

Baishan Hu<sup>1,11</sup>, Weiguang Jiang<sup>2,11</sup>, Takayuki Miyagi<sup>1,3,4,11</sup>, Zhonghao Sun<sup>5,6,11</sup>, Andreas Ekström<sup>2</sup>, Christian Forssén<sup>2,3,4</sup>, Gaute Hagen<sup>1,5,6</sup>, Jason D. Holt<sup>1,7</sup>, Thomas Papenbrock<sup>5,6</sup>, S. Ragnar Stroberg<sup>8,9</sup> and Ian Vernon<sup>10</sup>

- 17 LECs from  $\chi\text{EFT}$
- Start from  **$10^9$**  realisations
- **History matching** reduces to **34**

Wave	Target set $\mathcal{Z}$ outputs	Active inputs	Input samples	NI fraction	Proportion space NI
1	$6 \times 5$ <i>np</i> scattering	5–7	$10^5$ – $10^7$	$10^{-1}$ – $10^{-4}$	$3.2 \cdot 10^{-6}$
2	3 $A = 2$	7	$10^8$	$2.3 \cdot 10^{-4}$	$3.0 \cdot 10^{-7}$
3	6 $A = 2$ –4	13	$10^8$	$3.4 \cdot 10^{-5}$	$9.2 \cdot 10^{-8}$
4	8 $A = 2$ –4, 16	17	$5 \cdot 10^8$	$8.6 \cdot 10^{-6}$	same
final	38 all above	17	4,337	$7.8 \cdot 10^{-3}$	same

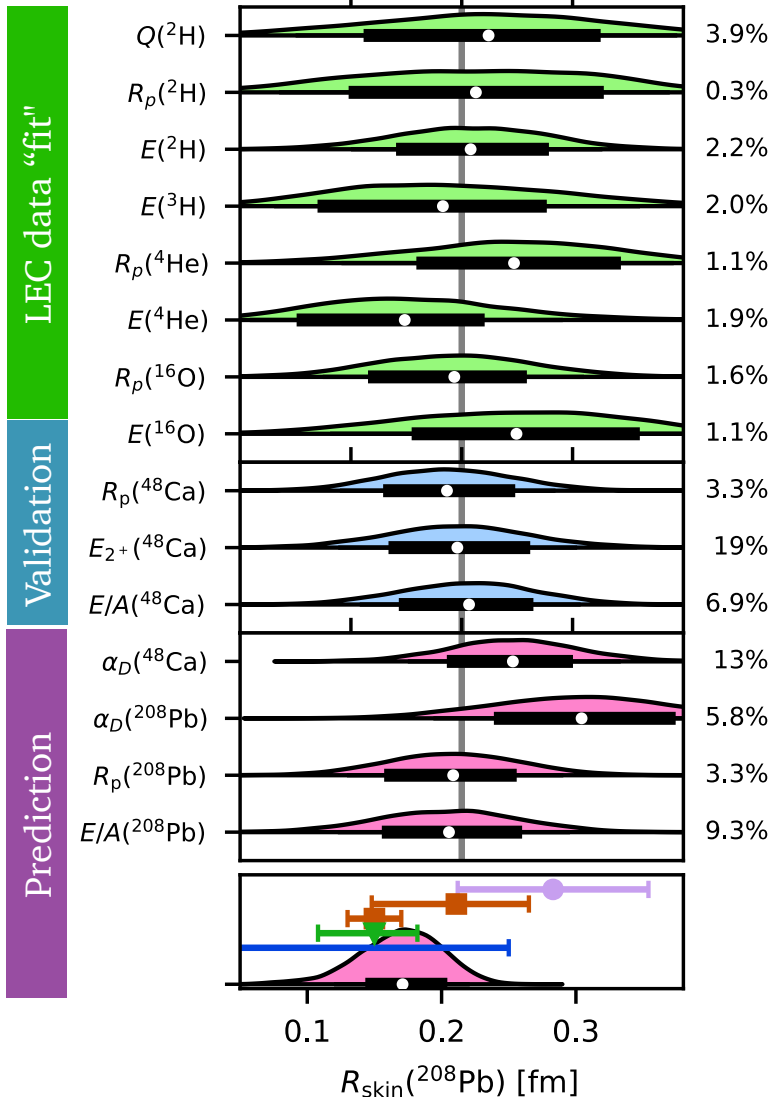


OPEN

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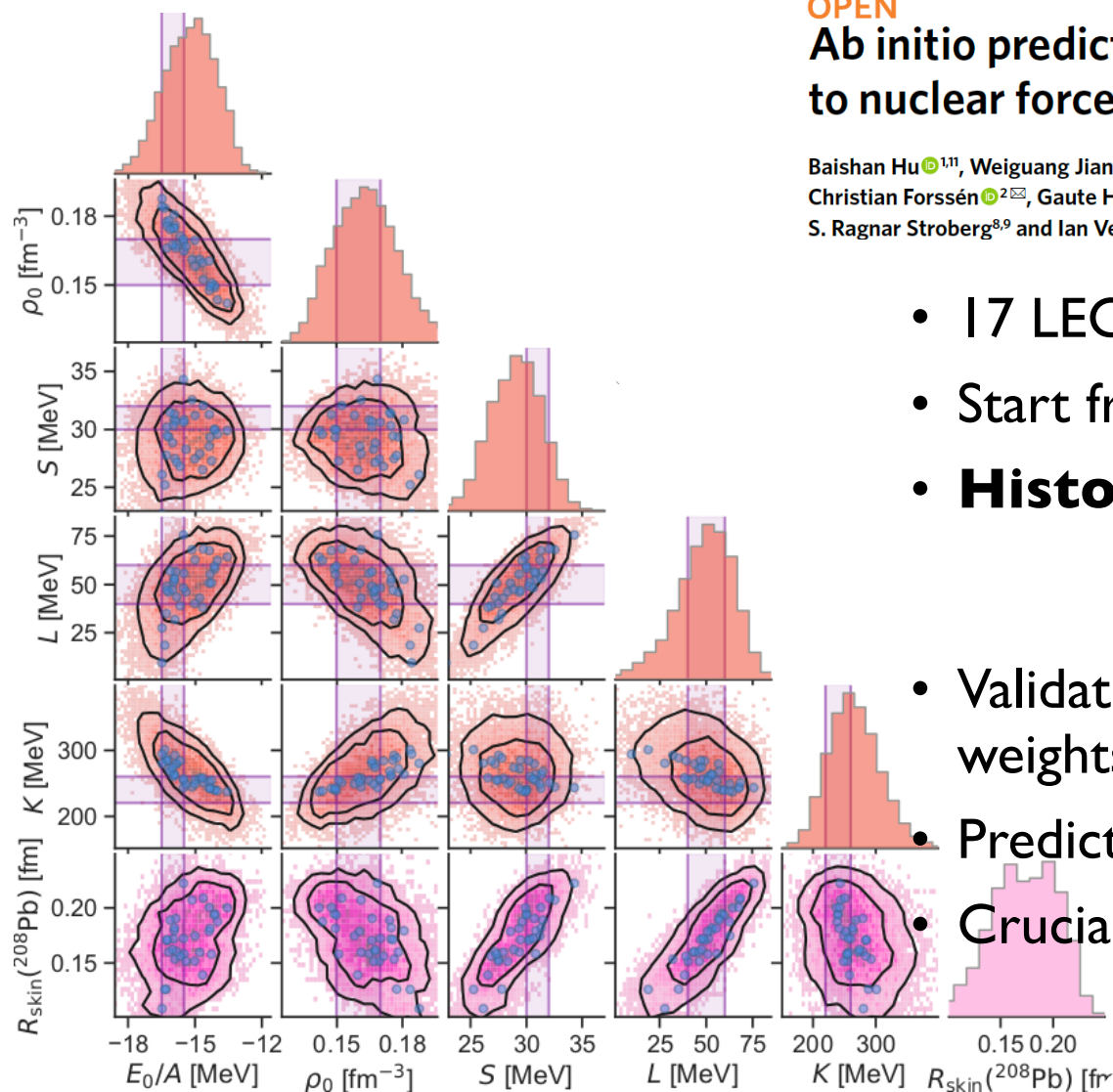
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- Validate data in  $^{48}\text{Ca}$  (importance weights)
- Predict  $^{208}\text{Pb}$  & nuclear matter
- Crucial role of **emulators** (GPs)



OPEN

## Ab initio predictions link the neutron skin of $^{208}\text{Pb}$ to nuclear forces

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● Nuclear physics

● **Neural Quantum States**

● Deuteron

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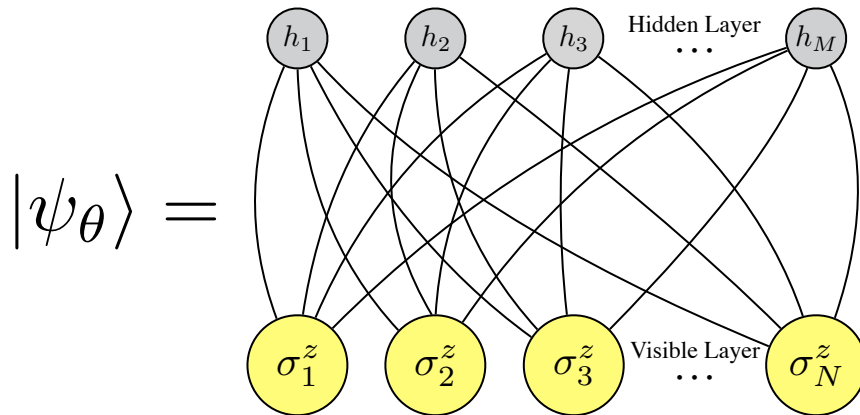
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$$E_{\theta} = \frac{\langle \psi_{\theta} | \hat{H} | \psi_{\theta} \rangle}{\langle \psi_{\theta} | \psi_{\theta} \rangle}$$



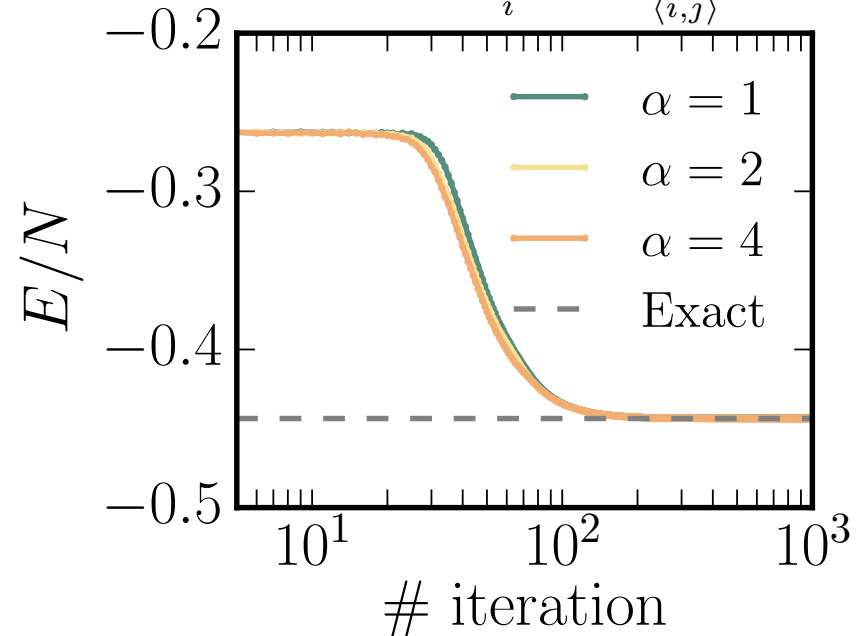
G. Carleo & M. Troyer, *Science* **355**, 602 (2017)

H. Saito, *J. Phys. Soc. Japan* **86**, 093001 (2017)

H. Saito, *J. Phys. Soc. Japan* **87**, 074002 (2018)

## 1D Heisenberg model

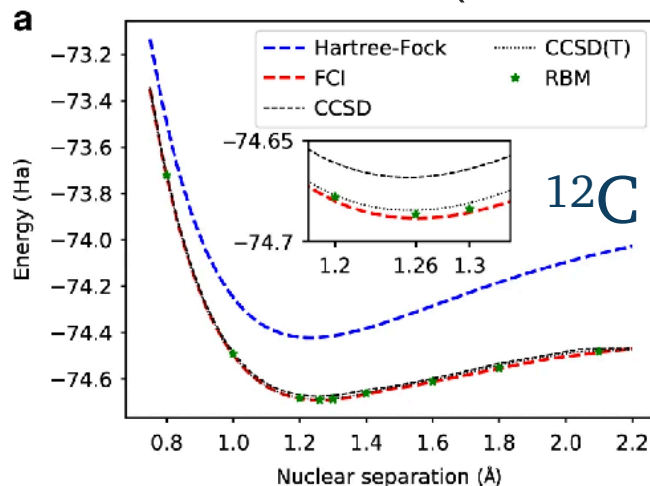
$$\mathcal{H}_{\text{TFI}} = -h \sum_i \sigma_i^x - \sum_{\langle i,j \rangle} \sigma_i^z \sigma_j^z$$



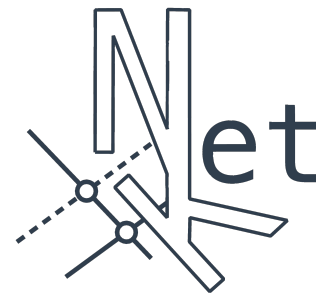
- **Machine learning** techniques to solve the many-body problem
- **Variational Artificial Neural Networks** (now NQS)
  - ANNs **efficiently represent** high-dimensional data (coarse graining?)
  - ANNs are setup for **variational** problems
  - ANNs can be used in **dynamics**

# NQs in quantum chemistry

## NetKet framework (Zurich & Flatiron)

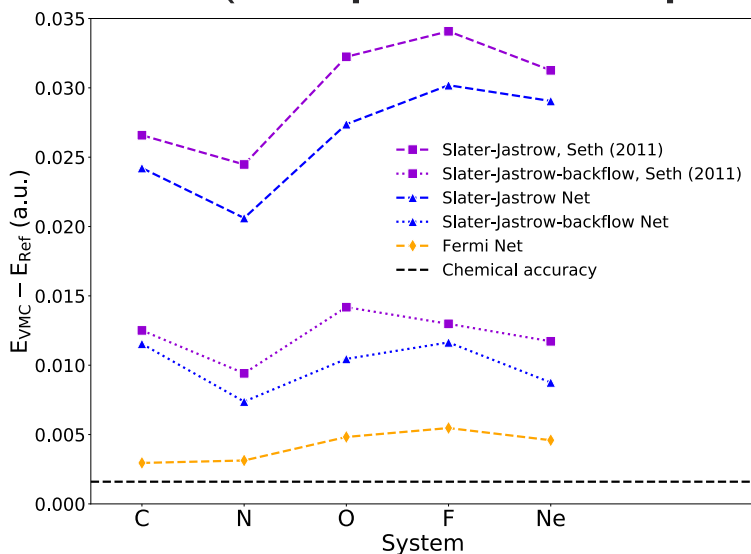


Choo, Carleo et al *Nat Comm* **11** 2368 (2019)



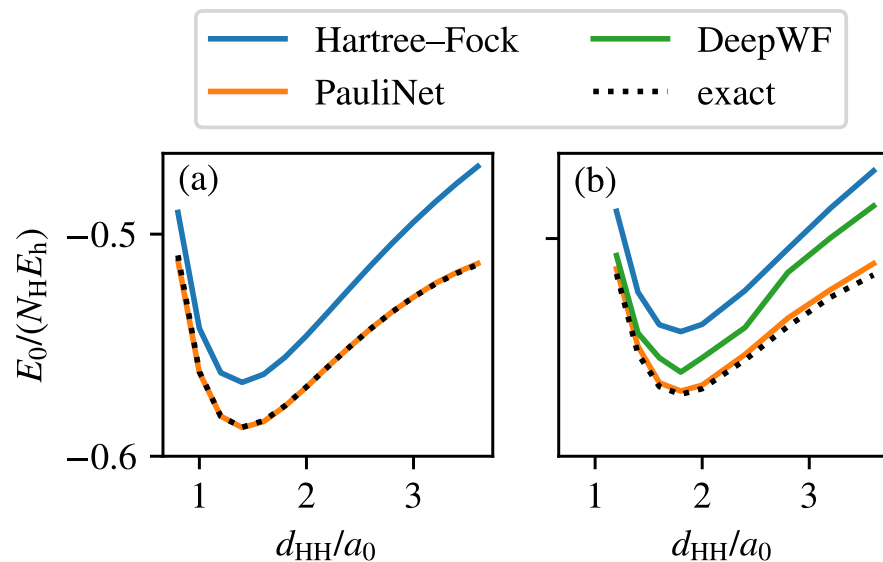
<https://www.netket.org/>

## FermiNet (DeepMind & Imperial)

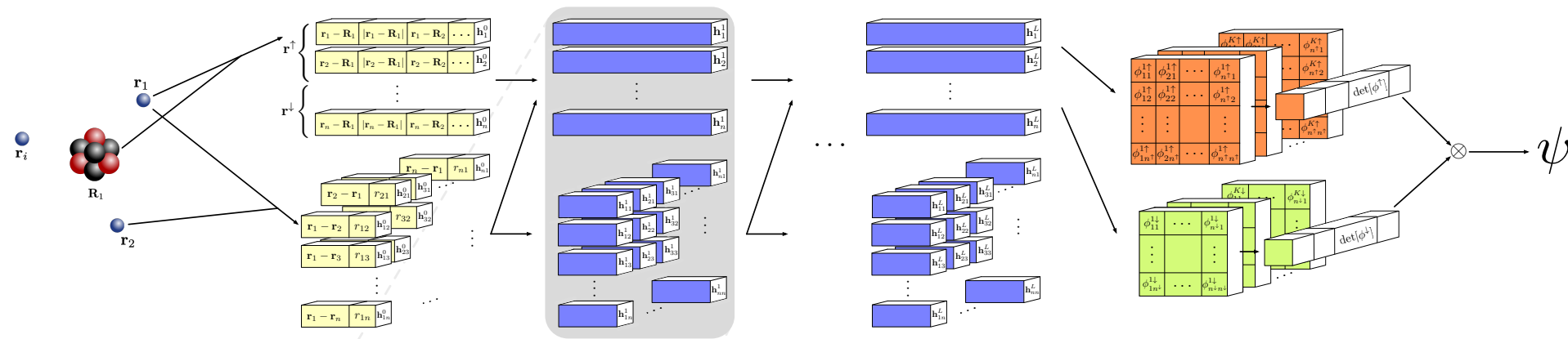


Pfau et al *Phys Rev Res* **2** 033429 (2020)

## PauliNet (FU & TU Berlin)



Hermann et al *Nat Chem* **12** 891 (2020)



**Input:** N-particle positions:  $\mathbf{r}_i - \mathbf{R}$  &  $\mathbf{r}_i - \mathbf{r}_j$

**Network:**

1) Single-particle orbitals (with backflow)

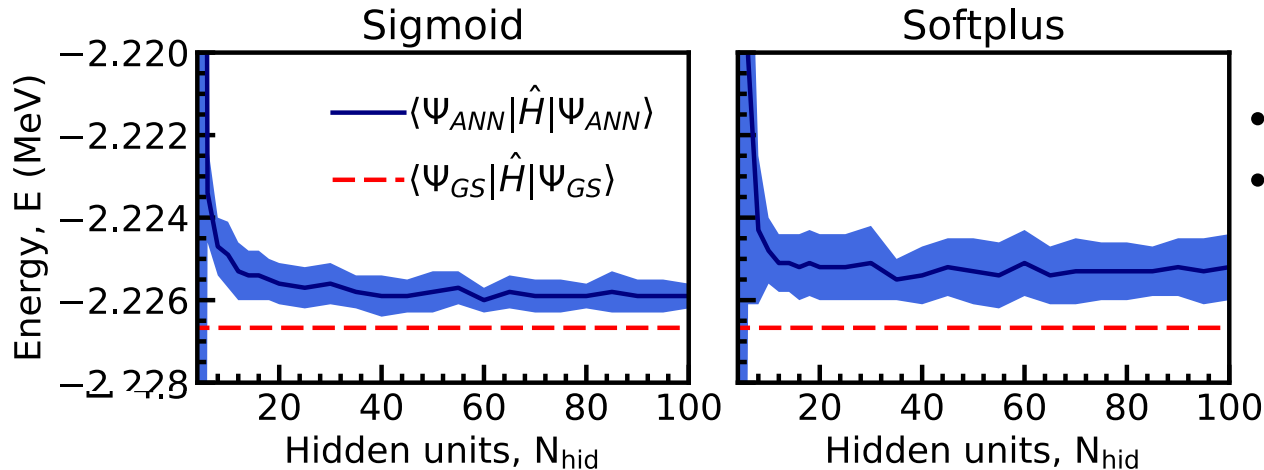
$$\phi_i(\mathbf{r}_j) \rightarrow \phi_i(\mathbf{r}_j; \mathbf{r}_1, \dots, \mathbf{r}_{j-1}, \mathbf{r}_{j+1}, \dots, \mathbf{r}_n) = \phi_i(\mathbf{r}_j; \{\mathbf{r}_{/j}\})$$

$$\phi_i^{k\alpha}(\mathbf{r}_j^\alpha; \{\mathbf{r}_{/j}^\alpha\}; \{\mathbf{r}^{\bar{\alpha}}\}) = (\mathbf{w}_i^{k\alpha} \cdot \mathbf{h}_j^{L\alpha} + g_i^{k\alpha}) \sum_m \pi_{im}^{k\alpha} \exp(-|\Sigma_{im}^{k\alpha}(\mathbf{r}_j^\alpha - \mathbf{R}_m)|)$$

2) N-body wavefunction

$$\psi(\mathbf{r}_1^\uparrow, \dots, \mathbf{r}_{n\downarrow}^\downarrow) = \sum_k \omega_k \left( \det \left[ \phi_i^{k\uparrow}(\mathbf{r}_j^\uparrow; \{\mathbf{r}_{/j}^\uparrow\}; \{\mathbf{r}^\downarrow\}) \right] \det \left[ \phi_i^{k\downarrow}(\mathbf{r}_j^\downarrow; \{\mathbf{r}_{/j}^\downarrow\}; \{\mathbf{r}^\uparrow\}); \right] \right)$$

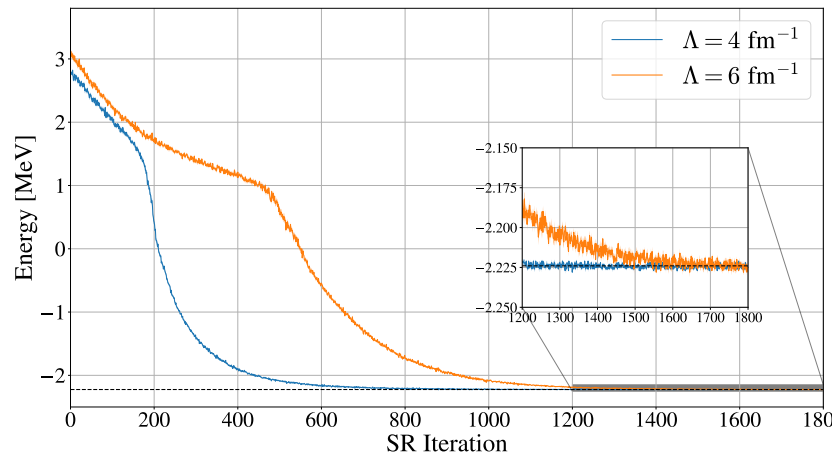
JWT Keeble & A. Rios, Phys. Lett. B **809** 135743 (2020) arxiv:1911.13092  
 J Rozalén, JWT Keeble & A Rios, EPJPlus arXiv:2205.12795



- Momentum basis
- EM N3LO

## • $A < 16$ , Neutron matter

Adams, Carleo, **Lovato** & Rocco, Phys. Rev. Lett. **127** 022502 (2021) arXiv:2007.14282  
 Gnech, Adams, **Lovato** et al, Few-Body Systems **63** 7 (2022) arXiv:2108.06836  
 Lovato, Adams, Carleo & Rocco, Phys. Rev. Research **4** 043178 (2022) arXiv:2206.10021



- $\pi$ -less EFT
- Gaussian int
- Spin?



ECT\*  
EUROPEAN CENTRE  
FOR THEORETICAL STUDIES  
IN NUCLEAR PHYSICS AND RELATED AREAS

## Many-body quantum physics with machine learning

4–8 Sept 2023  
ECT\*  
Europe/Rome timezone



<https://indico.ectstar.eu/event/177/>



● Nuclear physics

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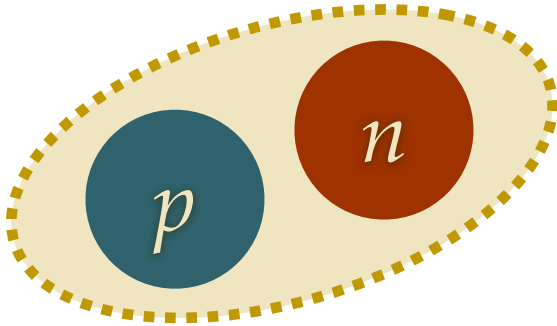
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# Why the deuteron?

- **Simplest** nuclear (np) bound state
- Experimental info **well-known**
- **Computationally** inexpensive benchmark
- Already requires 2 **states coupling**  $^3SD_1$
- Solved in **momentum** space



$$\hat{H} = \hat{H}_{CM} + \hat{H}_r$$

$$\hat{H}_r \psi(q) = E_d \psi(q)$$

$$\hat{H}_r = \frac{q^2}{2\mu} + V(q, q')$$

## Deuteron

$$E_d = 2.224575(9) \text{ MeV}$$

$$r_d = 1.97535(85) \text{ fm}$$

$$Q = 0.2859(3) \text{ fm}^2$$

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[arxiv:1911.13092](https://arxiv.org/abs/1911.13092)

# PYTORCH

- **Scalability** of Neural Networks
- Direct calculation of **gradients** via **Autograd**
- Portability to multiple GPUs



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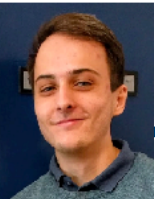
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PYT  RCH

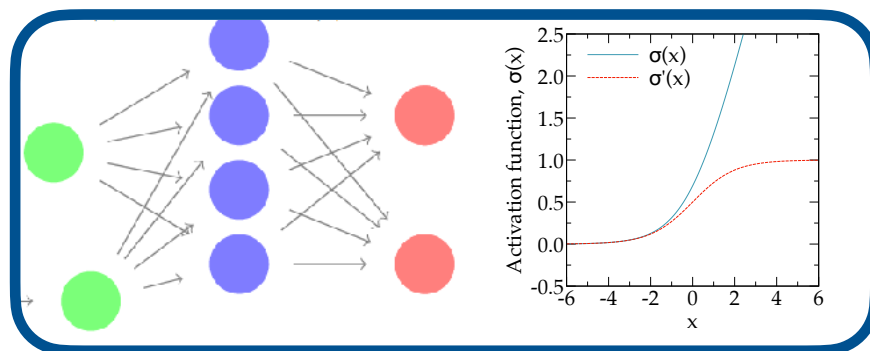
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# NQs three-step process

## I. Network architecture & initialisation



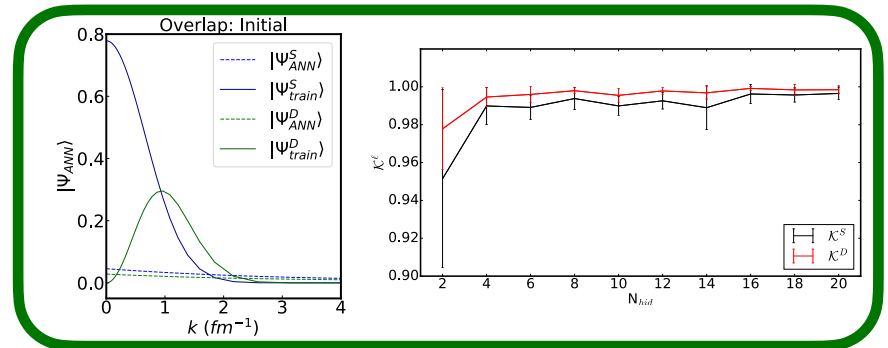
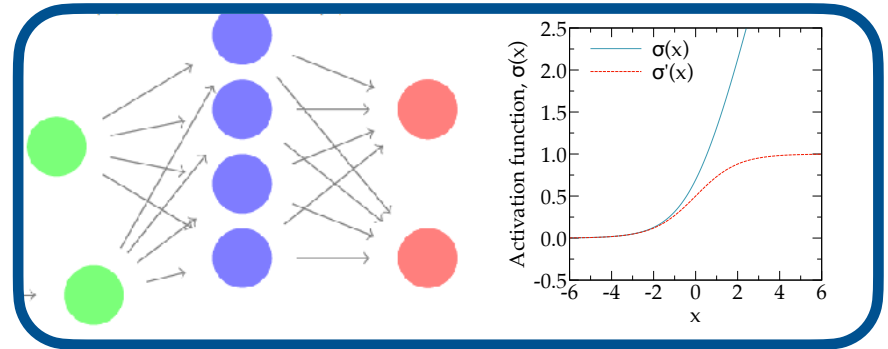


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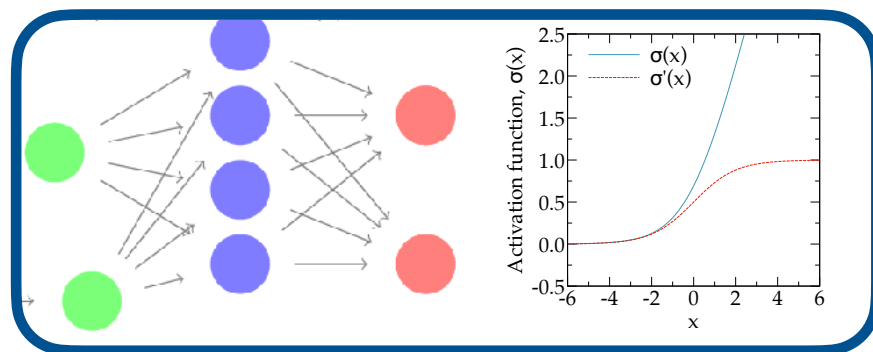
2. Network pre-training



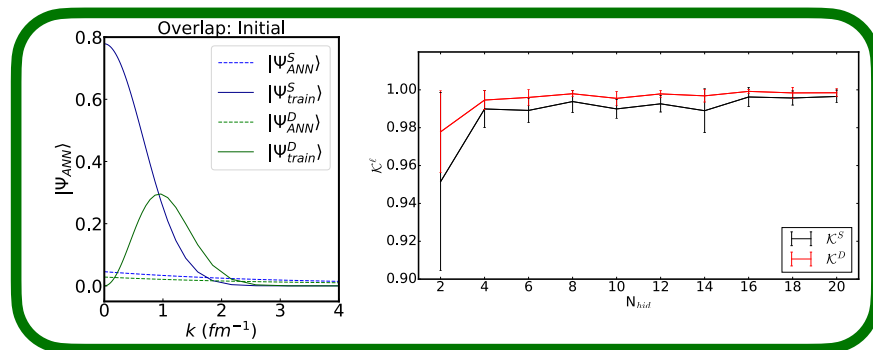


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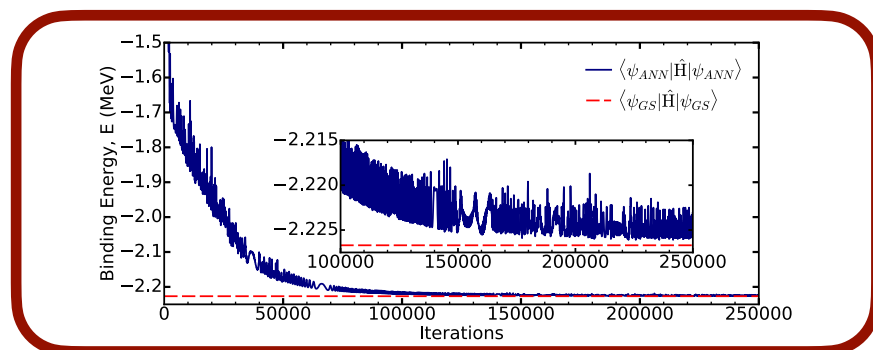
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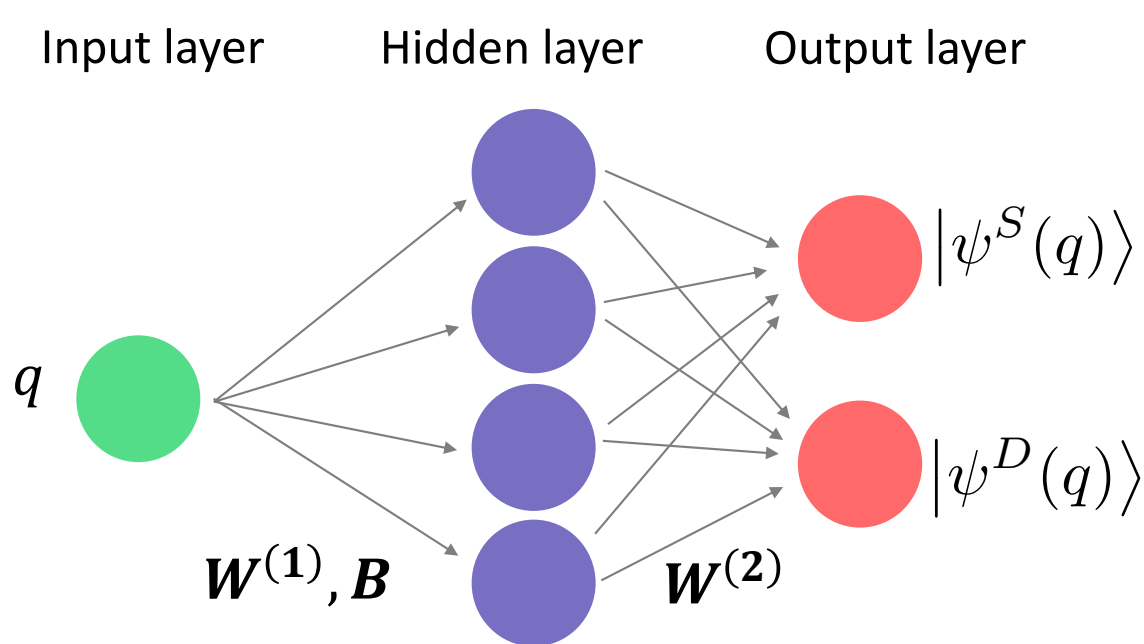


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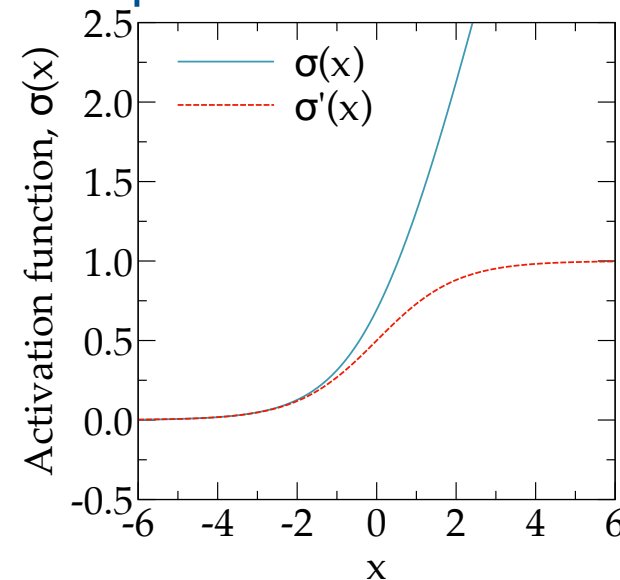


3. Energy minimisation





Softplus  $\sigma(x) = \ln(1 + e^x)$



## Wavefunction ansatz

$$\psi_{\text{ANN}}^L(q) = \sum_{i=1}^{N_{\text{hid}}} \mathcal{W}_{i,L}^{(2)} \sigma \left( \mathcal{W}_i^{(1)} q + b_i \right)$$

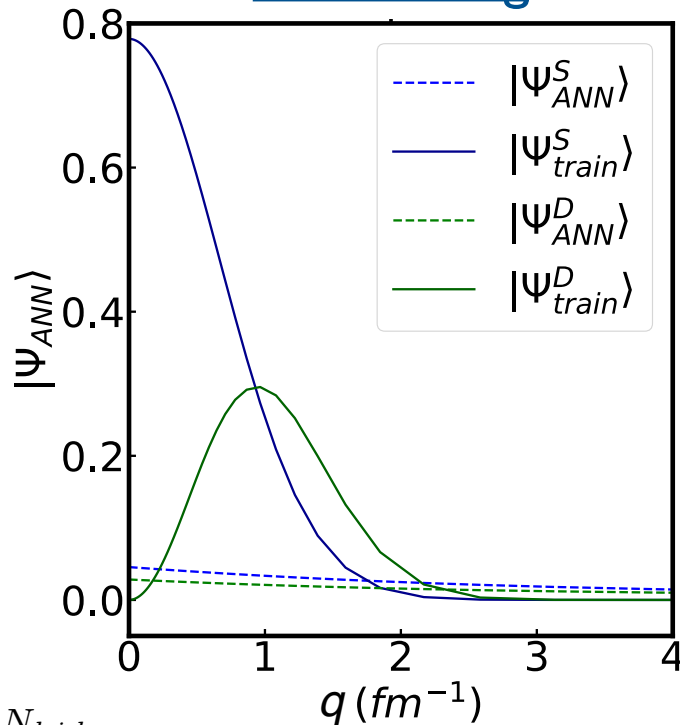
# of parameters

$$N_{\mathcal{W}} = 4 \times N_{\text{hid}}$$

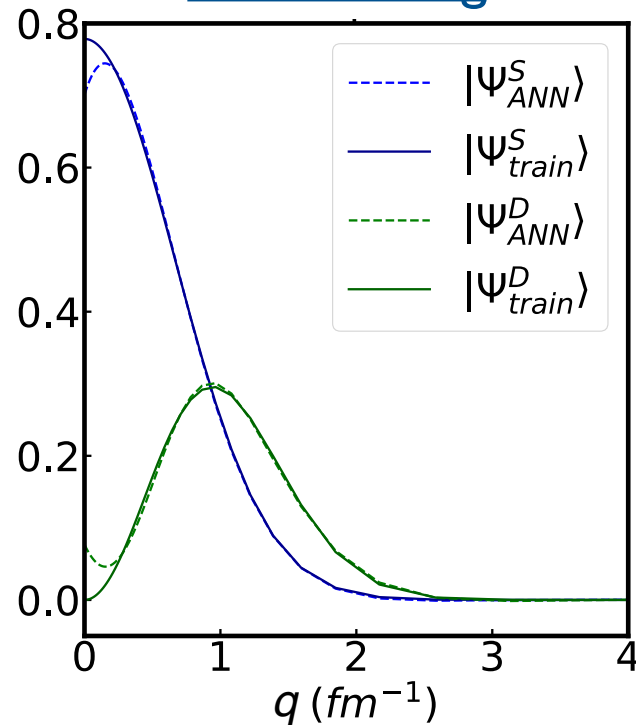
- Hidden layer provides **representability** (UAT)
- Architecture is **minimal**

# Step 2: Network training

Pre-training



Post-training



$$\psi_{\text{ANN}}^L(q) = \sum_{i=1}^{N_{\text{hid}}} \mathcal{W}_{i,L}^{(2)} \sigma(\mathcal{W}_i^{(1)} q + b_i)$$

## Initialisation

$$\mathcal{W}^{(1)} \in [-1, 0), \quad \mathcal{W}^{(2)} \in [0, 1)$$

$$b^{(1)} \in [-1, 1)$$

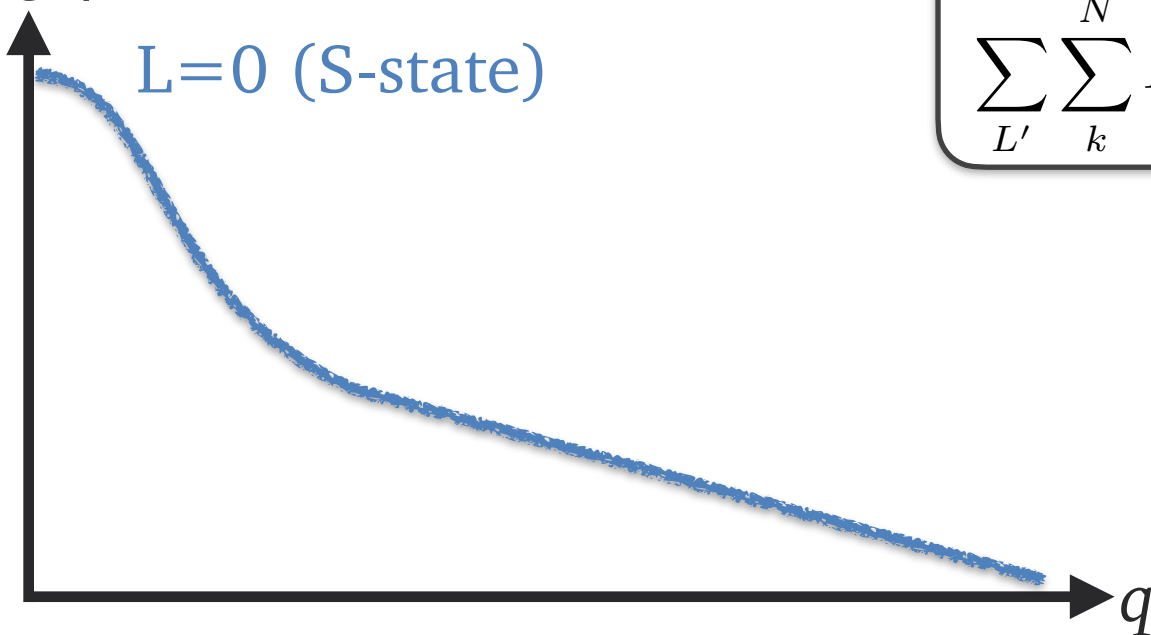
## **Training**

## Target wavefunction

$$\psi_{\text{targ}}^L(q) = \mathcal{N}_L q^L e^{-\frac{b^2 q^2}{2}}$$

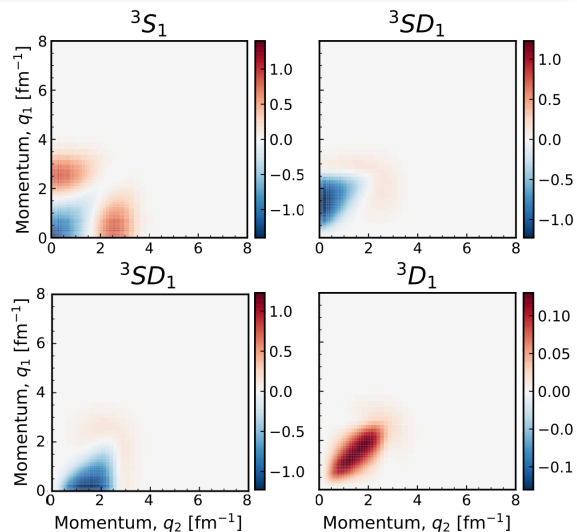
$\log \psi$

$L=0$  (S-state)



**Effective 1-body eq**

$$\sum_{L'} \sum_k^N H_{ik}^{LL'} |\psi^L(q_k)\rangle = E |\psi^L(q_i)\rangle$$

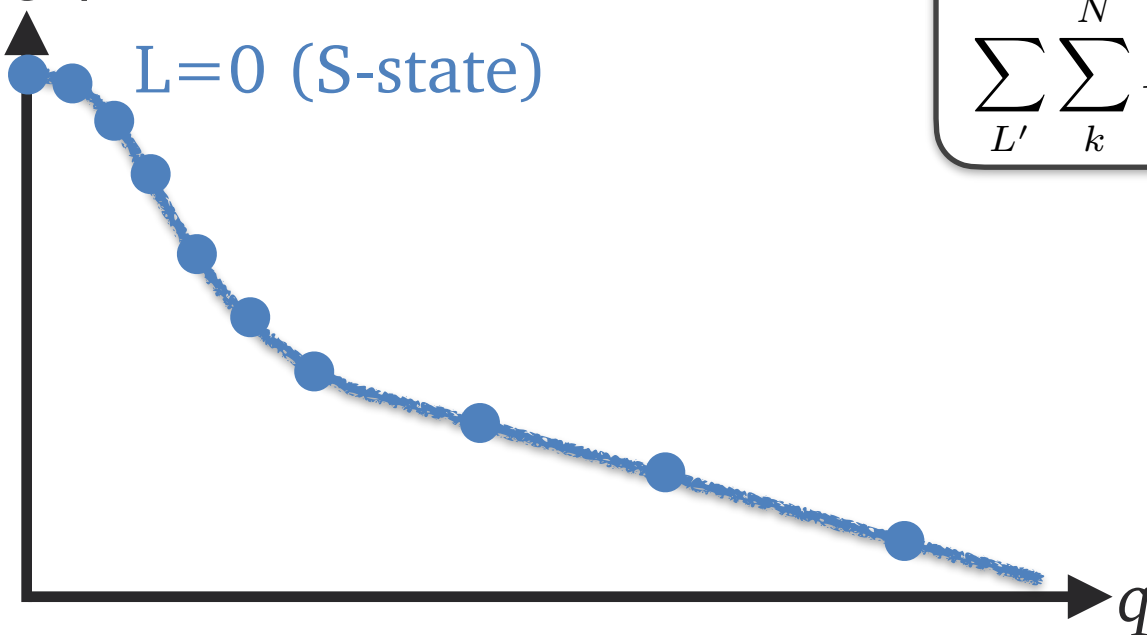


$$\langle \psi | K | \psi \rangle = \sum_L \int_0^\infty dq q^2 \frac{q^2}{2\mu} |\psi^L(q)|^2$$

- **No** need for **derivatives** in momentum space
- Tangential transform on  $N=64$  Gauss-Legendre mesh
- Training & validation **set**:  $N=64$  points

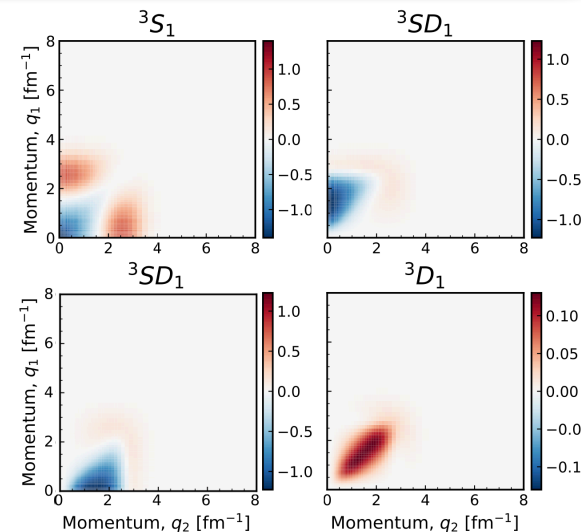
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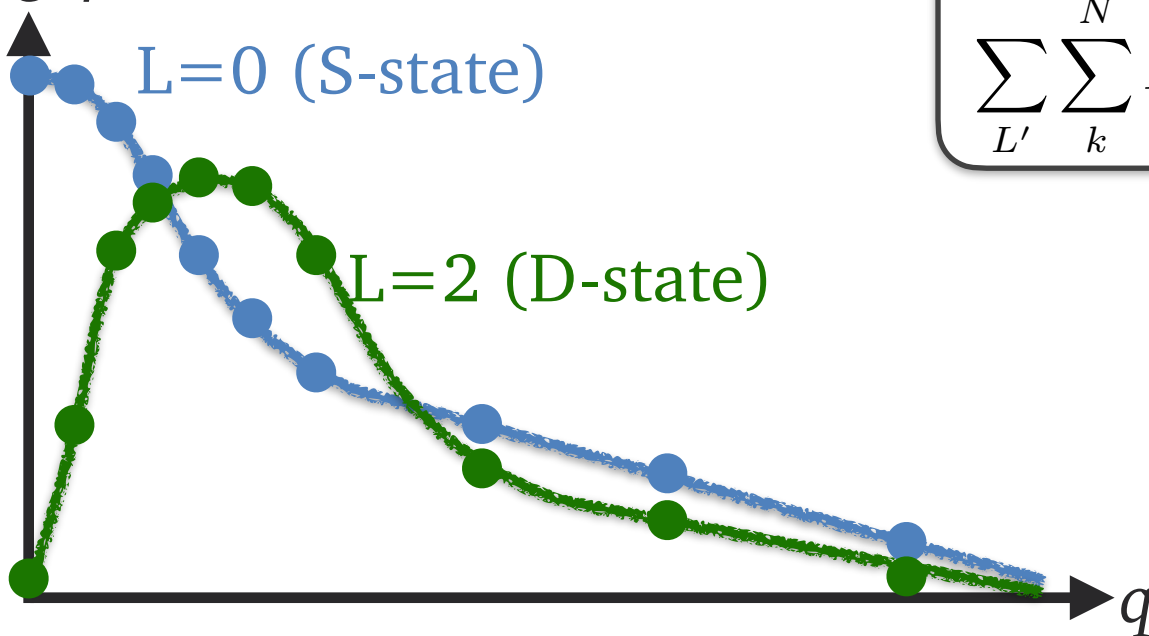


$$\langle \psi | K | \psi \rangle = \sum_L \int_0^\infty dq q^2 \frac{q^2}{2\mu} |\psi^L(q)|^2 \approx \sum_L \sum_i^N w_i q_i^2 \frac{q_i^2}{2\mu} |\psi^L(q_i)|^2$$

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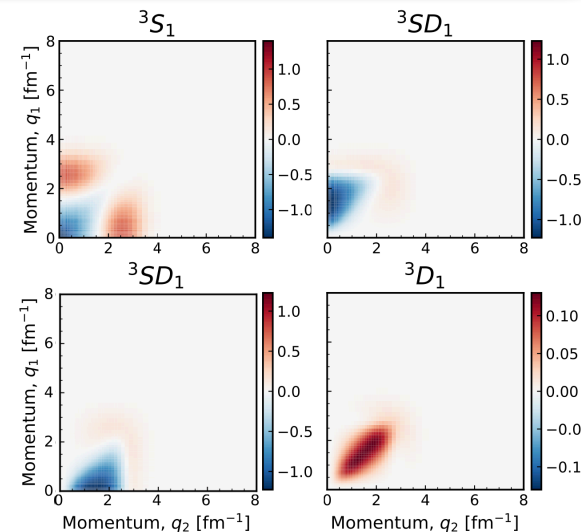


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# Energy minimisation - movie

$$\frac{\partial E}{\partial \mathcal{W}} = \frac{2}{\langle \psi | \psi \rangle} \left\{ \left\langle \frac{\partial \psi}{\partial \mathcal{W}} \middle| \hat{H} \middle| \psi \right\rangle - E \left\langle \frac{\partial \psi}{\partial \mathcal{W}} \middle| \psi \right\rangle \right\}$$

## RMS prop

$$\mathcal{W}_{t+1} = \mathcal{W}_t - \frac{\alpha}{\sqrt{S_t + \epsilon}} \frac{\partial \mathcal{C}}{\partial \mathcal{W}_t}$$

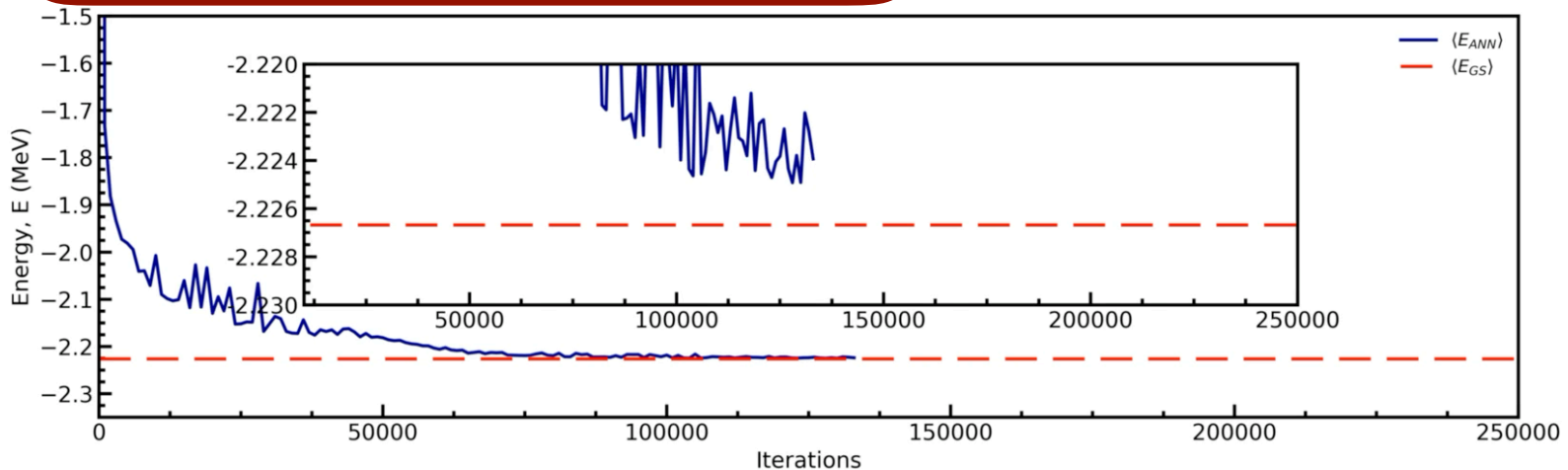
$$S_t = \gamma S_{t-1} + (1 - \gamma) \left( \frac{\partial \mathcal{C}}{\partial \mathcal{W}_t} \right)^2$$

## Hyperparameters

$$\alpha = 10^{-2}$$

$$\epsilon = 10^{-8}$$

$$\gamma = 0.9$$



# Energy minimisation - movie

$$\frac{\partial E}{\partial \mathcal{W}} = \frac{2}{\langle \psi | \psi \rangle} \left\{ \left\langle \frac{\partial \psi}{\partial \mathcal{W}} \middle| \hat{H} \middle| \psi \right\rangle - E \left\langle \frac{\partial \psi}{\partial \mathcal{W}} \middle| \psi \right\rangle \right\}$$

## RMS prop

$$\mathcal{W}_{t+1} = \mathcal{W}_t - \frac{\alpha}{\sqrt{S_t + \epsilon}} \frac{\partial \mathcal{C}}{\partial \mathcal{W}_t}$$

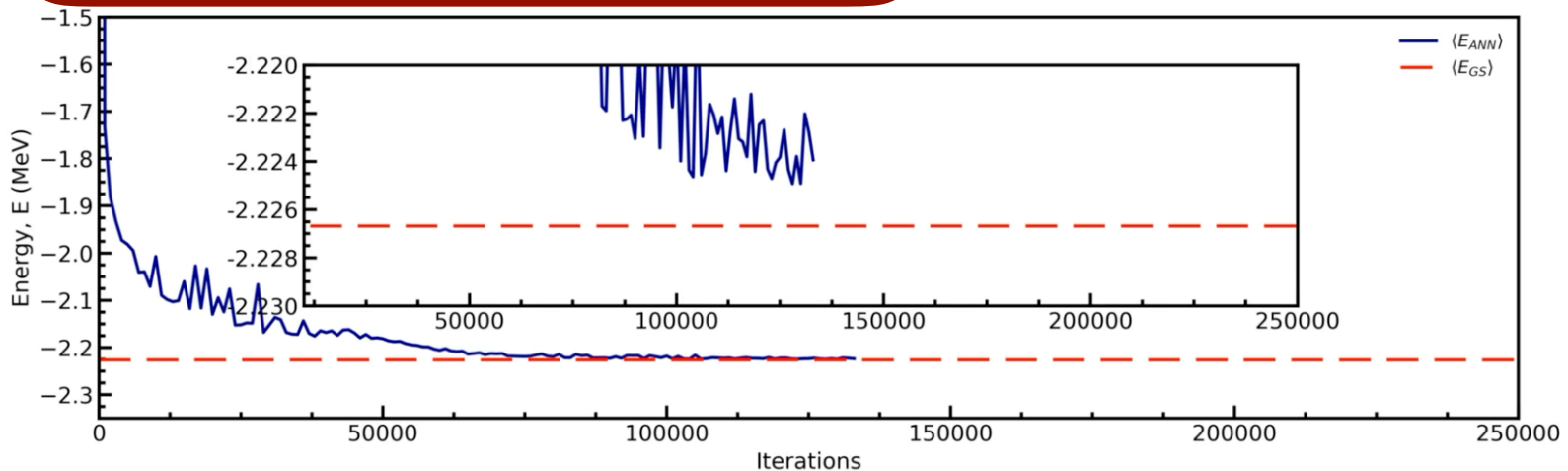
$$S_t = \gamma S_{t-1} + (1 - \gamma) \left( \frac{\partial \mathcal{C}}{\partial \mathcal{W}_t} \right)^2$$

## Hyperparameters

$$\alpha = 10^{-2}$$

$$\epsilon = 10^{-8}$$

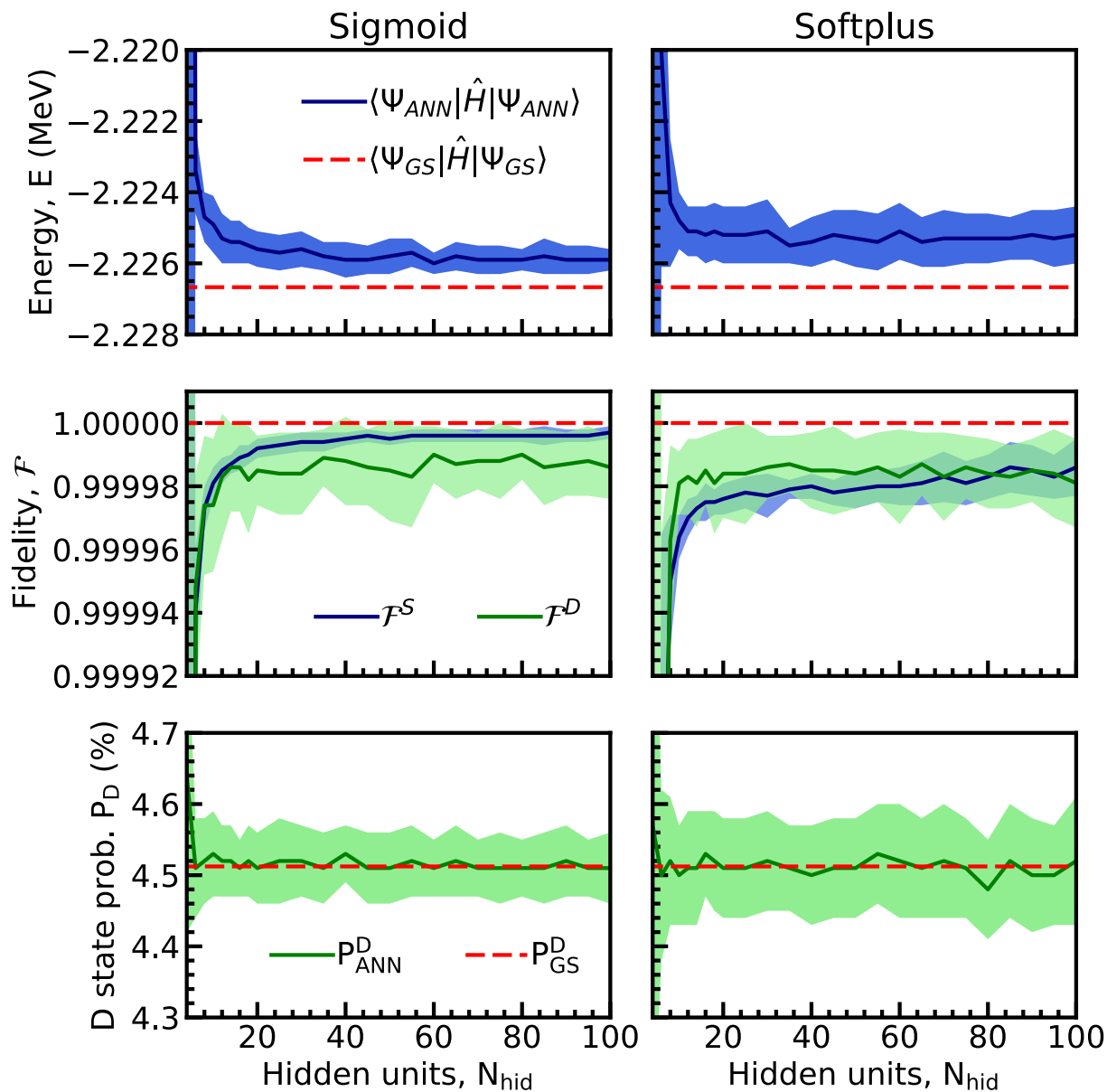
$$\gamma = 0.9$$



# Results: dependence on $N_{hid}$

## Fidelity

$$\mathcal{F}^L = \frac{\langle \psi_{GS}^L | \psi_{ANN}^L \rangle^2}{\langle \psi_{GS}^L | \psi_{GS}^L \rangle \langle \psi_{ANN}^L | \psi_{ANN}^L \rangle}$$

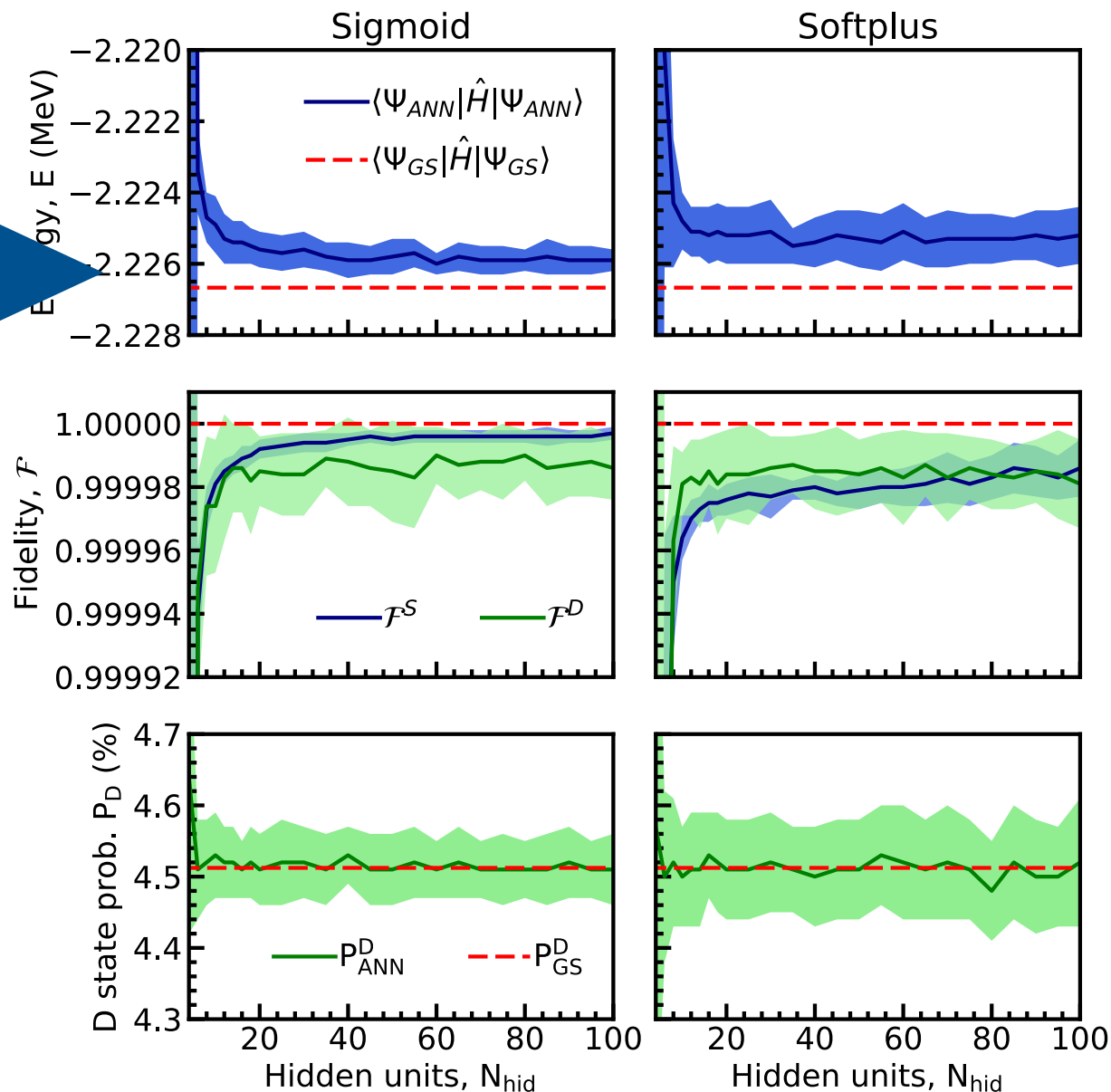


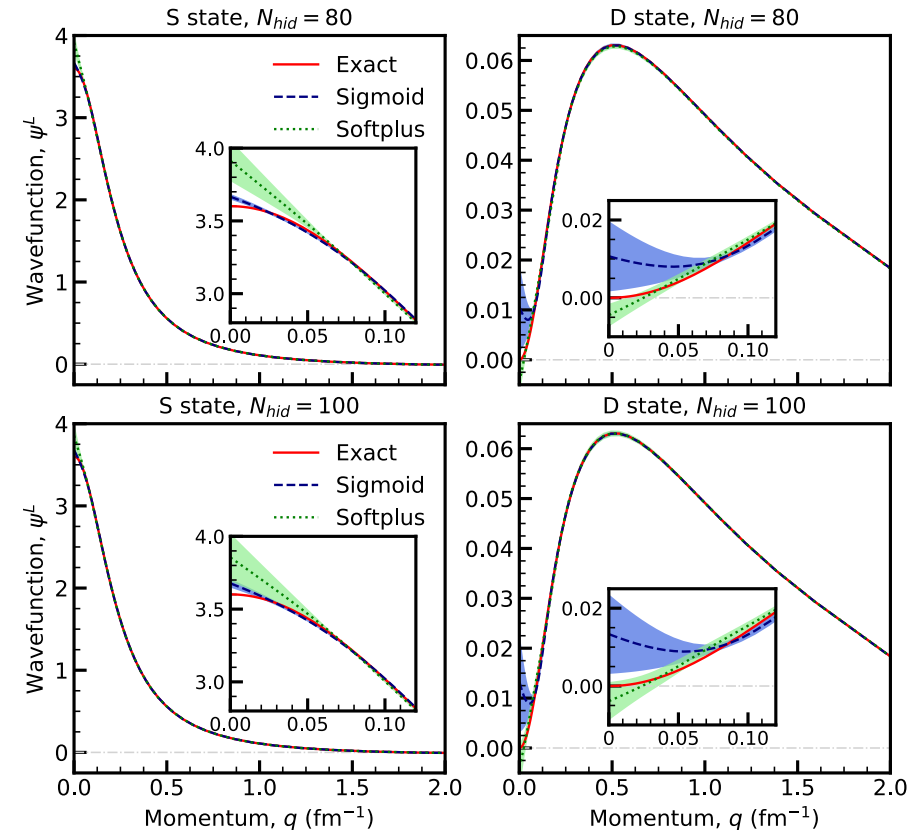
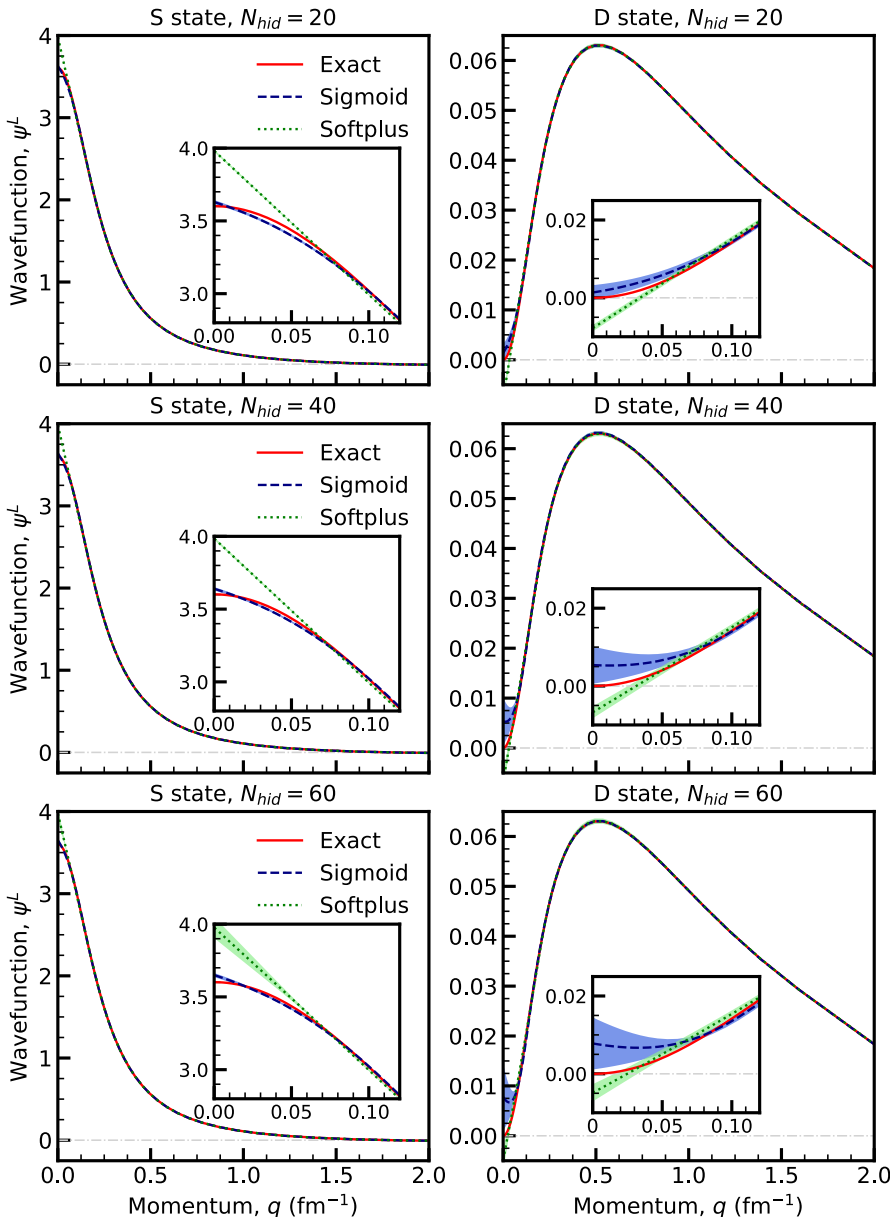
# Results: dependence on $N_{hid}$

keV accuracy  
No bias-variance

## Fidelity

$$\mathcal{F}^L = \frac{\langle \psi_{GS}^L | \psi_{ANN}^L \rangle^2}{\langle \psi_{GS}^L | \psi_{GS}^L \rangle \langle \psi_{ANN}^L | \psi_{ANN}^L \rangle}$$

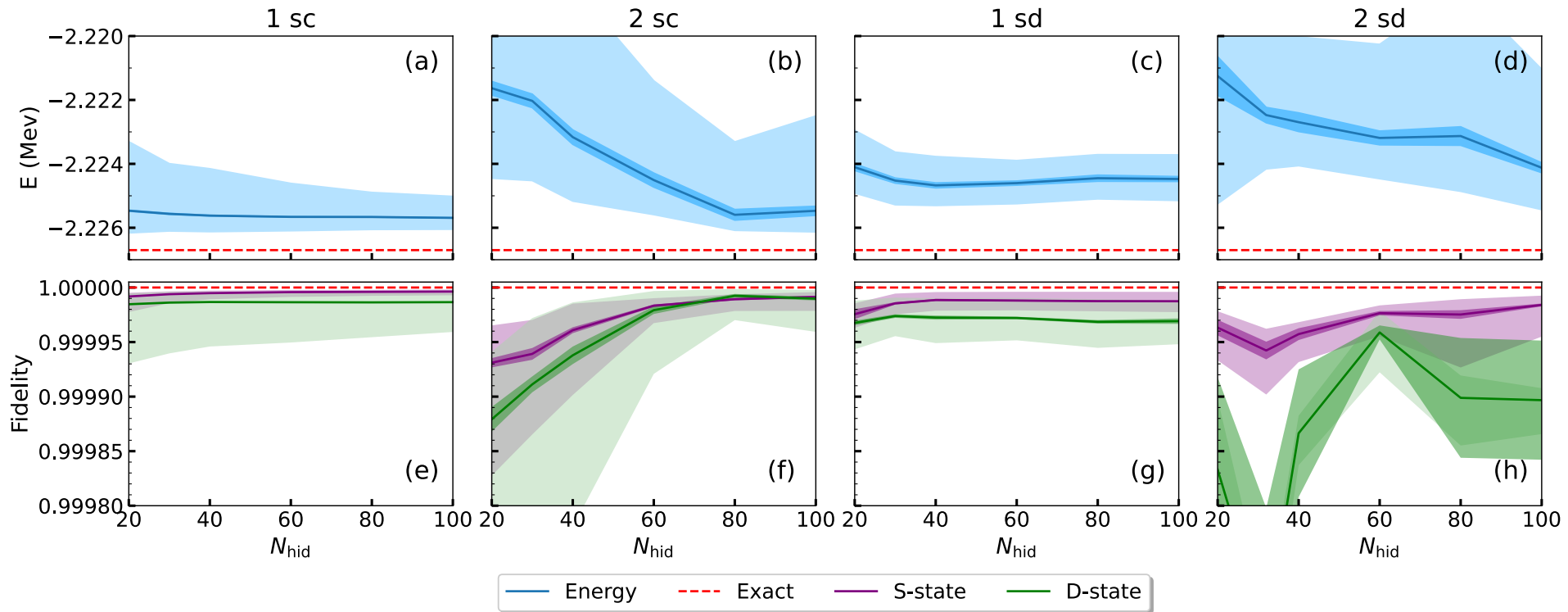
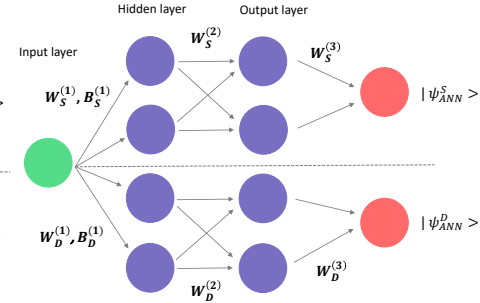
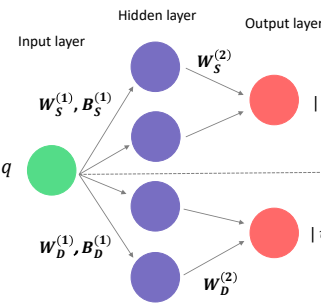
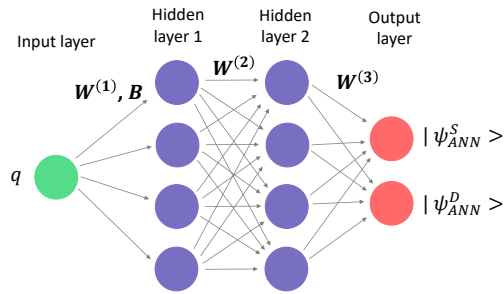
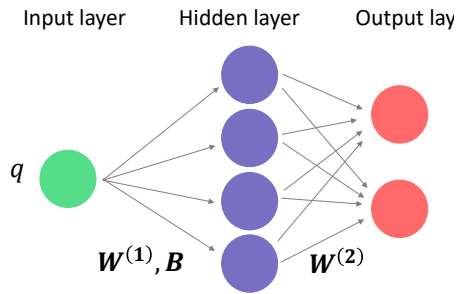




- **Variance increases with  $N_{hid}$**
- **But in costless regions**

$$\langle \psi | K | \psi \rangle = \sum_L \int_0^\infty dq q^2 \frac{q^2}{2\mu} |\psi^L(q)|^2$$

# Different architectures



## ● Nuclear physics

## ● Neural Quantum States

## ● Deuteron

JWT Keeble & A. Rios, Phys. Lett. B **809**, 135743 (2020)

J Rozalen-Sarmiento, JWT Keeble & A. Rios, Eur. Phys. J. Plus **139**, 189 (2024)

## ● **1D spinless (polarised) systems ( $A > 2$ )**

J.W.T. Keeble, M. Drissi, A. Rojo-Francàs, B. Juliá-Díaz & A. Rios, Phys. Rev. A **108**, 063320 (2023),  
arxiv:2304.04725

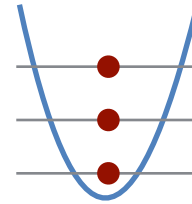
## ● Optimisation strategy

M. Drissi, J.W.T. Keeble, J. Rozalen-Sarmiento & A. Rios, arxiv:2401.17550

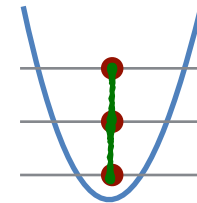


## Deep-VMC approach

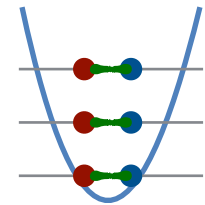
1. **Free** spinless fermions in 1D ✓



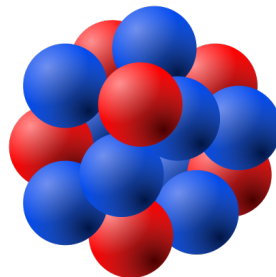
2. **Interacting** spinless fermions in 1D ✓



3. **Interacting spin-full** fermions in 1D ✗



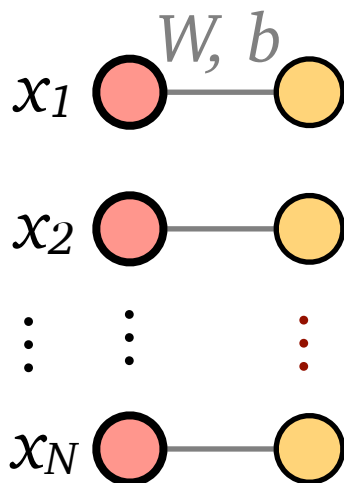
4. **3D nuclei** ✗



# Antisymmetry? Equivariance!

$$f_i(x_{\pi(1)}, \dots, x_{\pi(N)}) = f_{\pi(i)}(x_1, \dots, x_N)$$

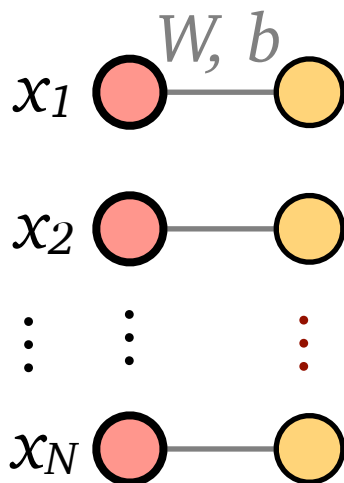
$$\begin{pmatrix} f_1 \\ f_2 \\ \vdots \\ f_N \end{pmatrix} = \begin{pmatrix} W & 0 & \cdots & 0 \\ 0 & W & \cdots & 0 \\ \vdots & \vdots & \ddots & \vdots \\ 0 & 0 & \cdots & W \end{pmatrix} \begin{pmatrix} x_1 \\ x_2 \\ \vdots \\ x_N \end{pmatrix} + \begin{pmatrix} b \\ b \\ \vdots \\ b \end{pmatrix} = \begin{pmatrix} Wx_1 + b \\ Wx_2 + b \\ \vdots \\ Wx_N + b \end{pmatrix}$$



# Antisymmetry? Equivariance!

$$f_i(x_{\pi(1)}, \dots, x_{\pi(N)}) = f_{\pi(i)}(x_1, \dots, x_N)$$

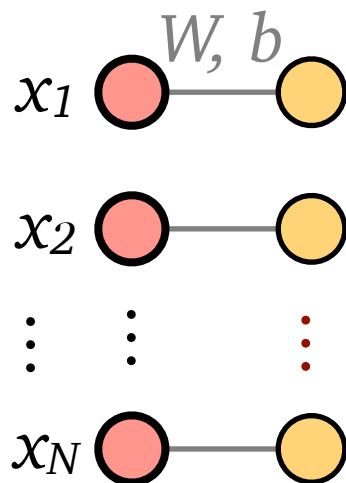
$$\begin{pmatrix} f_1 \\ f_2 \\ \vdots \\ f_N \end{pmatrix} = \begin{pmatrix} W & 0 & \cdots & 0 \\ 0 & W & \cdots & 0 \\ \vdots & \vdots & \ddots & \vdots \\ 0 & 0 & \cdots & W \end{pmatrix} \begin{pmatrix} x_1 \\ x_2 \\ \vdots \\ x_N \end{pmatrix} + \begin{pmatrix} b \\ b \\ \vdots \\ b \end{pmatrix} = \begin{pmatrix} Wx_1 + b \\ Wx_2 + b \\ \vdots \\ Wx_N + b \end{pmatrix}$$



# Antisymmetry? Equivariance!

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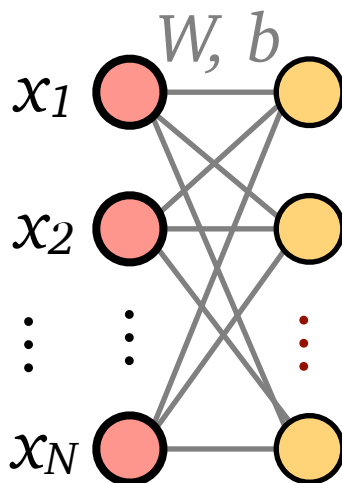
$$\begin{pmatrix} f_1 \\ f_N \\ \vdots \\ f_2 \end{pmatrix} = \begin{pmatrix} W & 0 & \cdots & 0 \\ 0 & W & \cdots & 0 \\ \vdots & \vdots & & \vdots \\ 0 & 0 & & W \end{pmatrix} \begin{pmatrix} x_1 \\ x_N \\ \vdots \\ x_2 \end{pmatrix} + \begin{pmatrix} b \\ b \\ \vdots \\ b \end{pmatrix} = \begin{pmatrix} Wx_1 + b \\ Wx_N + b \\ \vdots \\ Wx_2 + b \end{pmatrix}$$



# Antisymmetry? Equivariance!

$$f_i(x_{\pi(1)}, \dots, x_{\pi(N)}) = f_{\pi(i)}(x_1, \dots, x_N)$$

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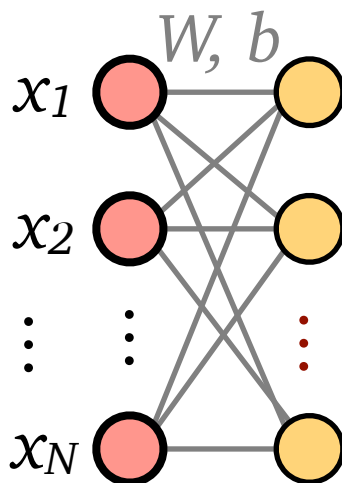


# Antisymmetry? Equivariance!

$$f_i(x_{\pi(1)}, \dots, x_{\pi(N)}) = f_{\pi(i)}(x_1, \dots, x_N)$$

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$$Wx_i + V \sum_{j \neq i} x_j + b = (W - V)x_i + V \sum_j x_j + b$$



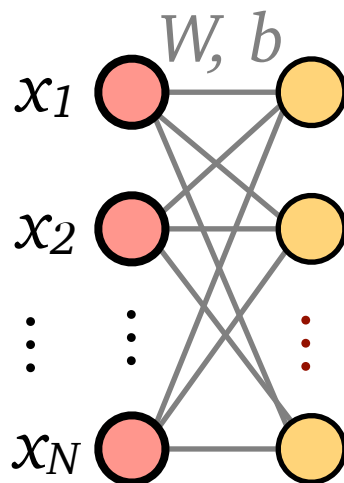
# Antisymmetry? Equivariance!

$$f_i(x_{\pi(1)}, \dots, x_{\pi(N)}) = f_{\pi(i)}(x_1, \dots, x_N)$$

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$$Wx_i + V \sum_{j \neq i} x_j + b = (\cancel{W}^{W'} V)x_i + V \sum_j x_j + b$$

$\pi$ -invariant



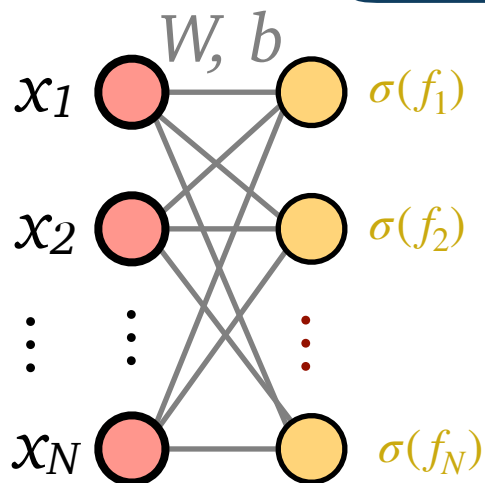
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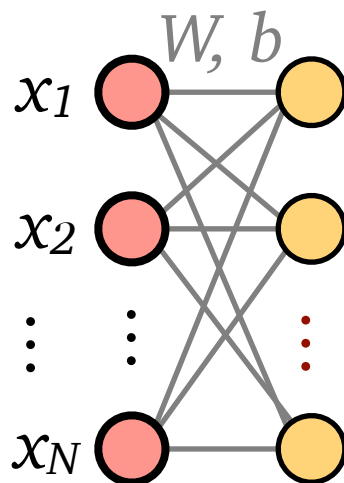
# Antisymmetry? Equivariance!

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$$Wx_i + V \sum_{j \neq i} x_j + b = (\cancel{W}^{W'} V)x_i + V \sum_j x_j + b$$

*$\pi$ -invariant*



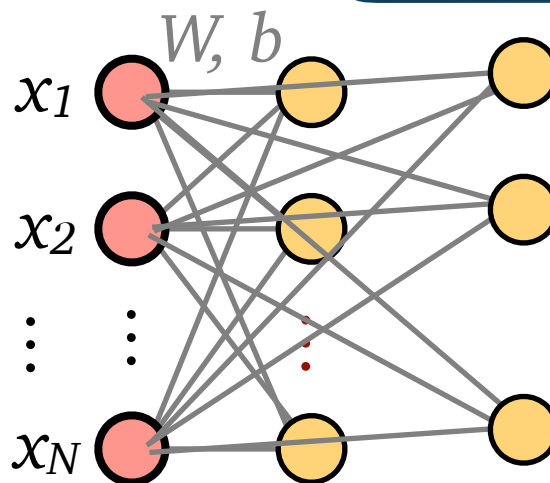
# Antisymmetry? Equivariance!

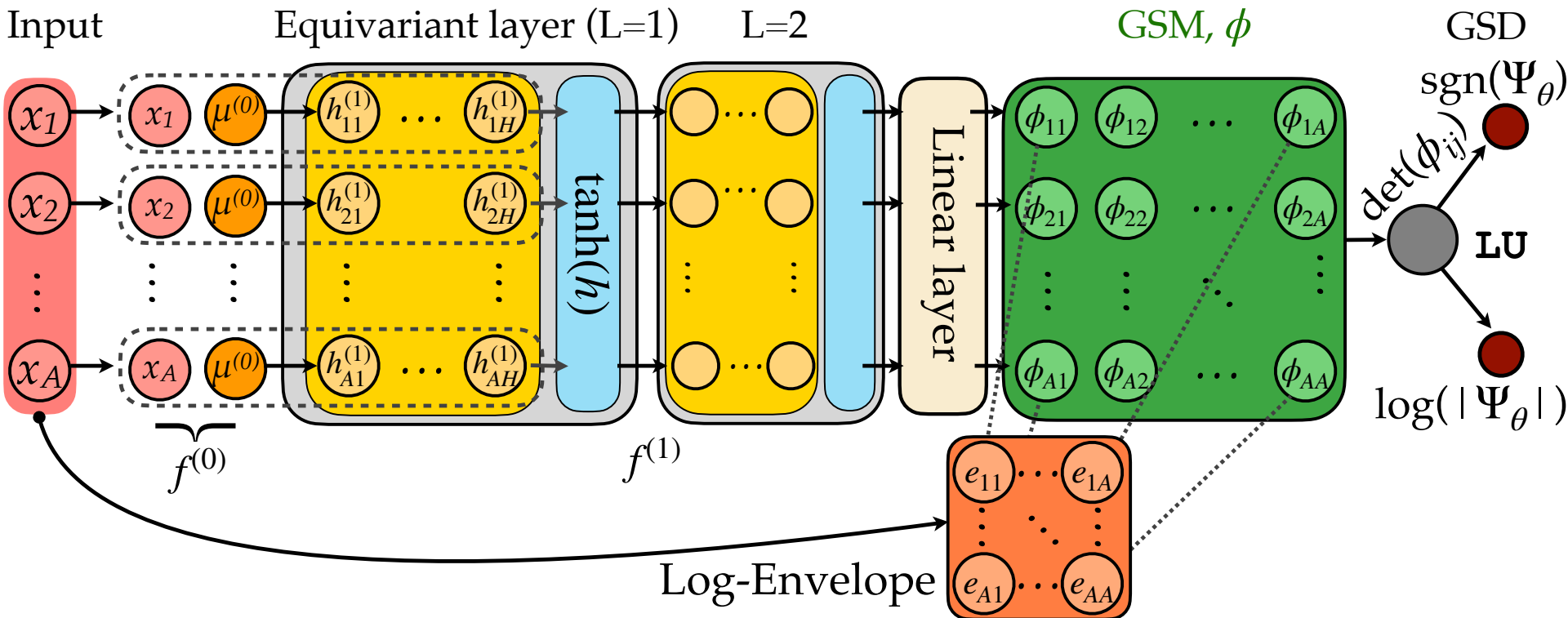
$$f_i(x_{\pi(1)}, \dots, x_{\pi(N)}) = f_{\pi(i)}(x_1, \dots, x_N)$$

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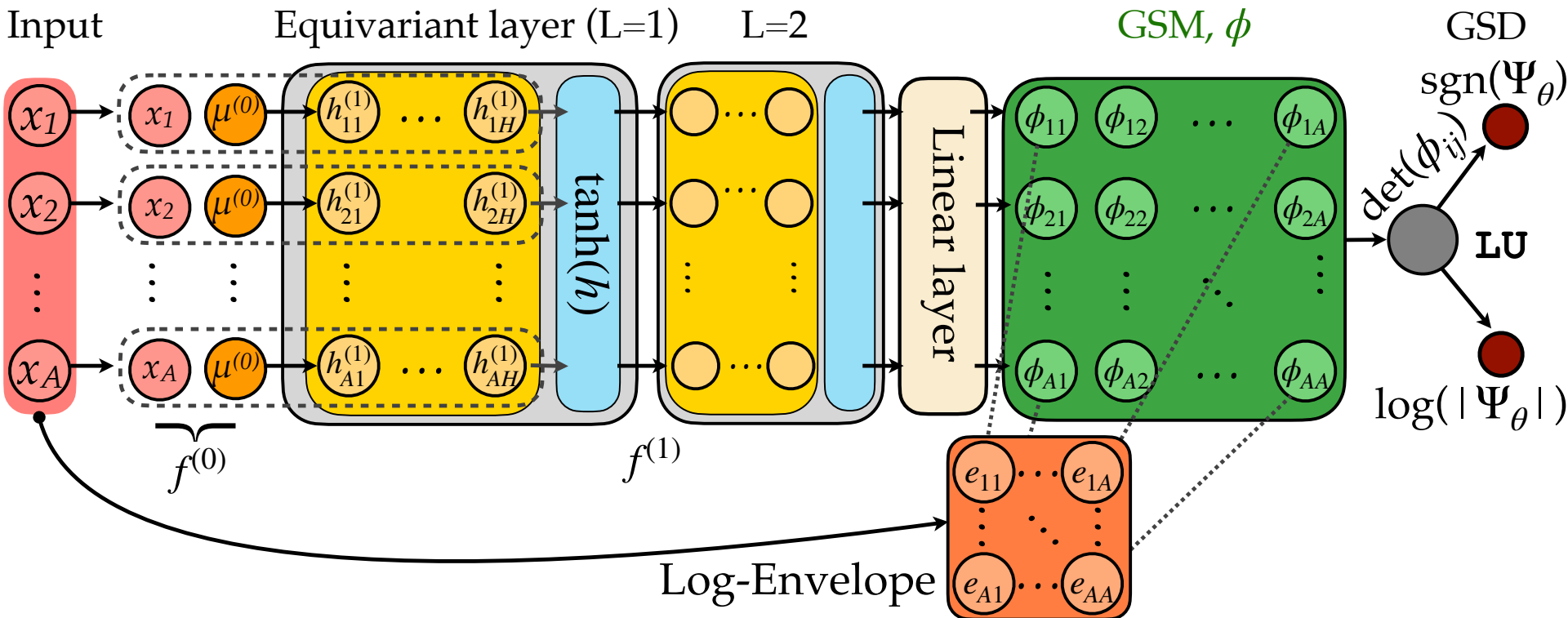
$$Wx_i + V \sum_{j \neq i} x_j + b = (\cancel{W}^{W'} V)x_i + V \sum_j x_j + b$$

*$\pi$ -invariant*

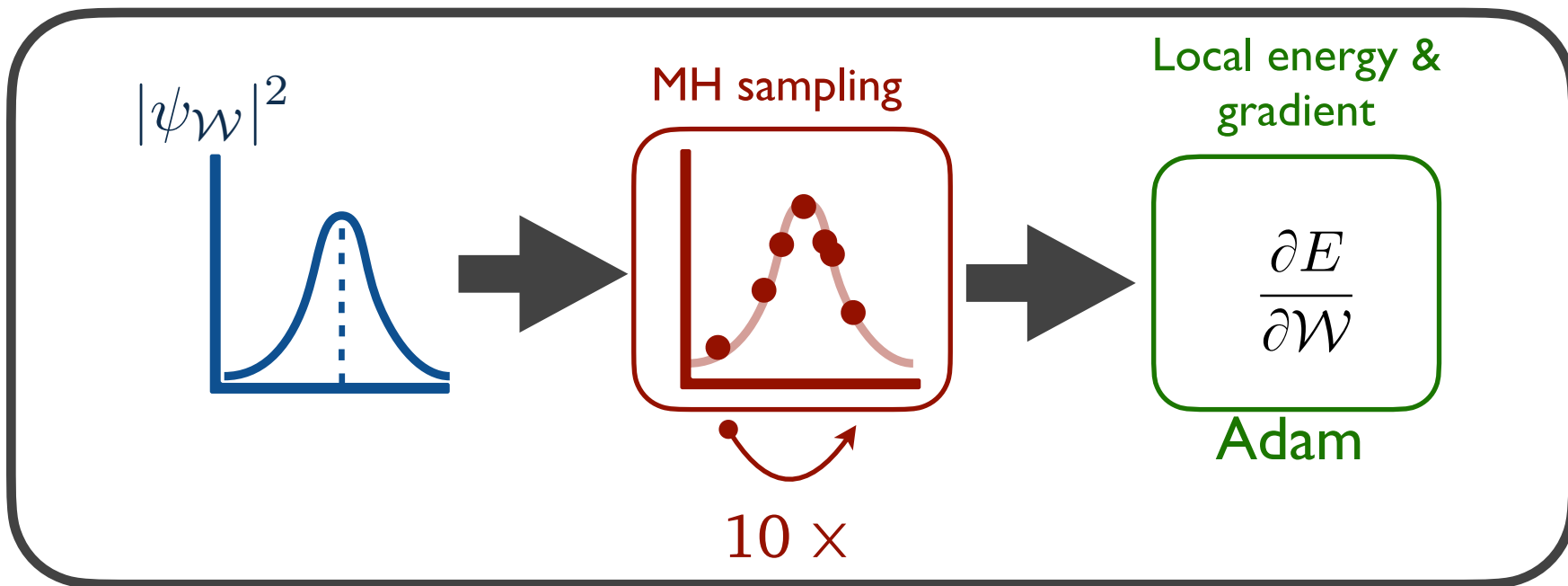




- Permutation **equivariance**
- $D=1$  determinants
- Fixed **latent width**  $H=64$
- $L=2$  layers
- Envelope layer (boundaries)
- **Linear increase** of # parameters increase in  $A$



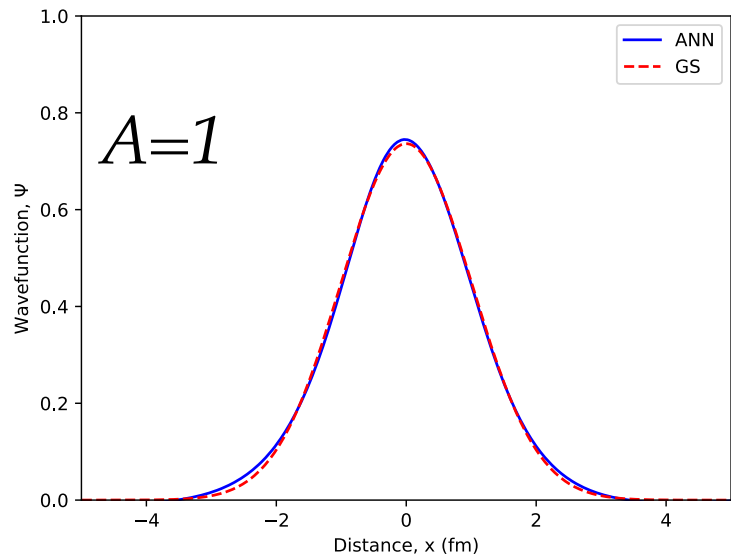
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- $D=1$  determinants
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- $L=2$  layers
- Envelope layer (boundaries)
- **Linear increase** of # parameters increase in  $A$



- 4096 walkers in **real** space
- **Local energy & gradient** every 10 MH epochs
- Total of **100,000** epochs
- Energy computed from  $10^4$  batches  $\times$  4096 walkers
- Width of Metropolis-Hastings adapted on-the-fly
- For KE, laplacian wrt input needed

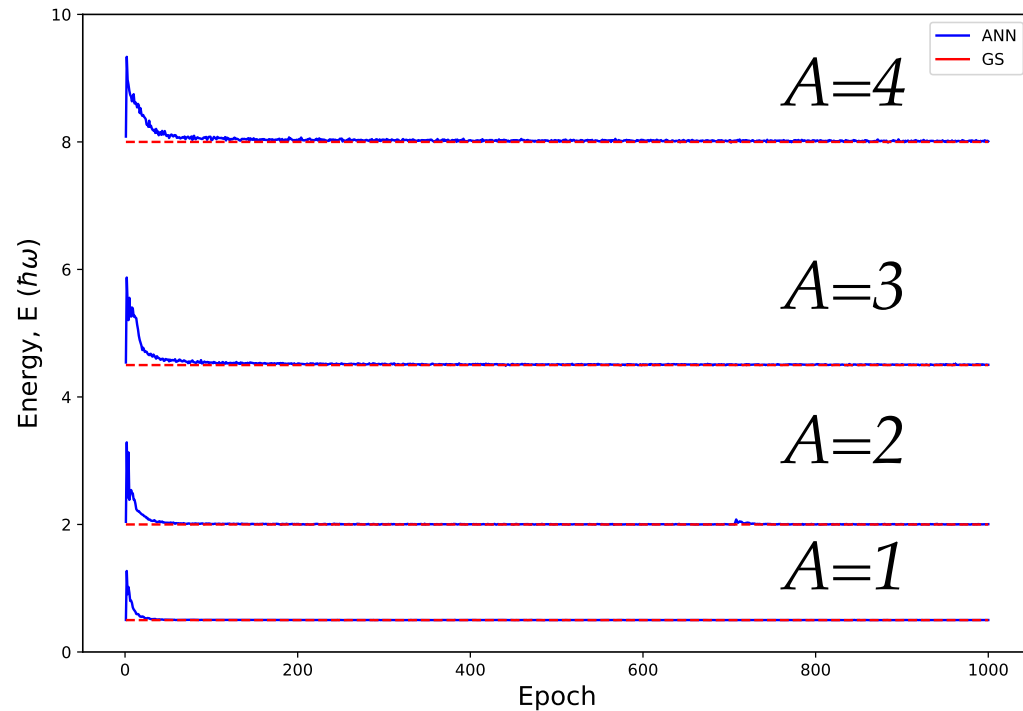
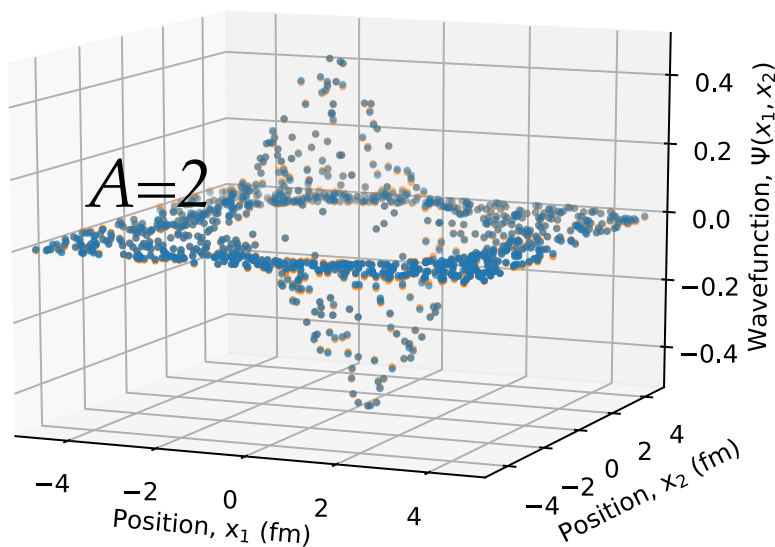
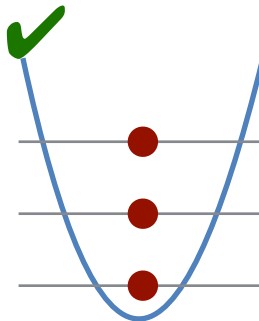
$$- \sum_i \frac{\hbar^2}{2m} \frac{\partial^2}{\partial x_i^2} \Psi(x_1, \dots, x_A)$$

# Trapped spinless fermions in 1D HO

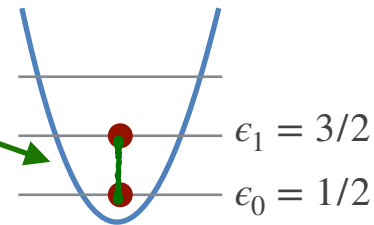


- **Antisymmetrization** ✓

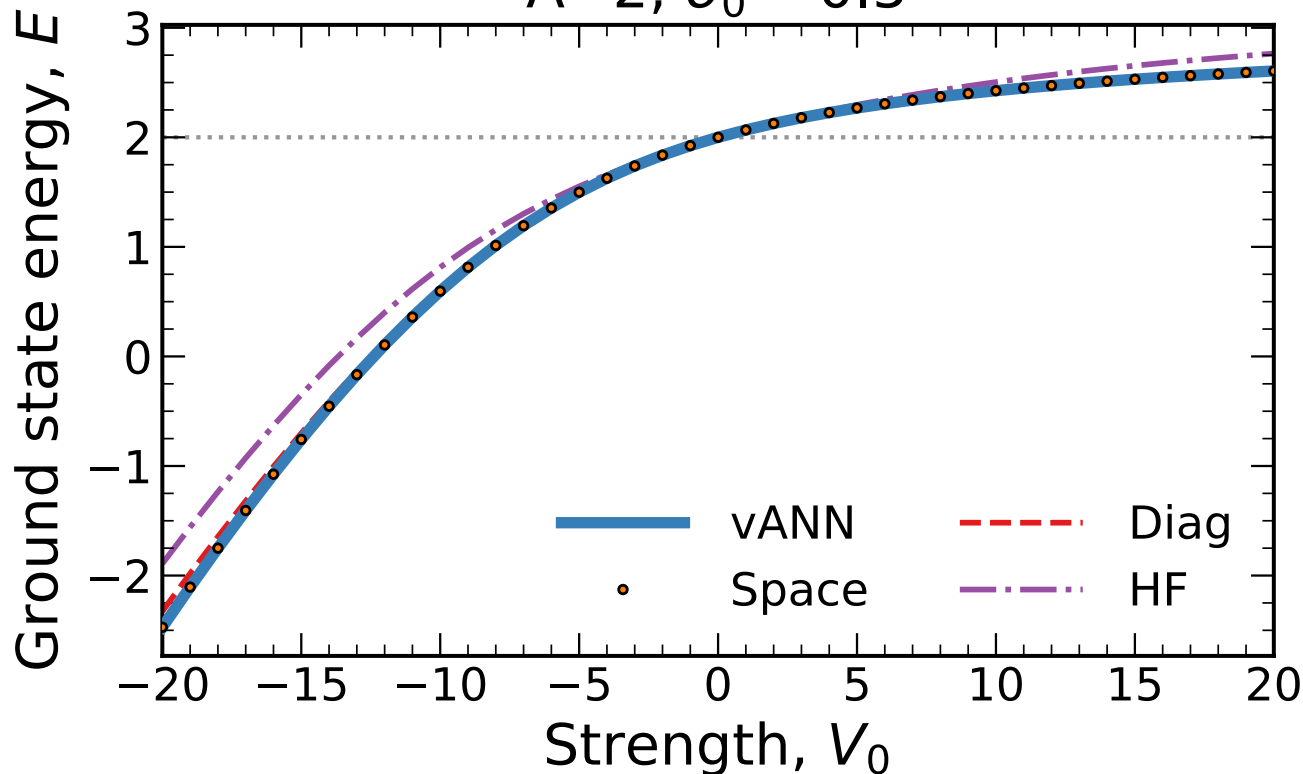
$$E = \frac{A^2}{2}$$



$$\hat{H} = - \sum_i \frac{\partial_{x_i}^2}{2} + \sum_i \frac{x_i^2}{2} + \sum_{i < j} \frac{V_0}{\sqrt{2\pi\sigma_0}} \exp\left(-\frac{(x_i - x_j)^2}{2\sigma_0^2}\right)$$



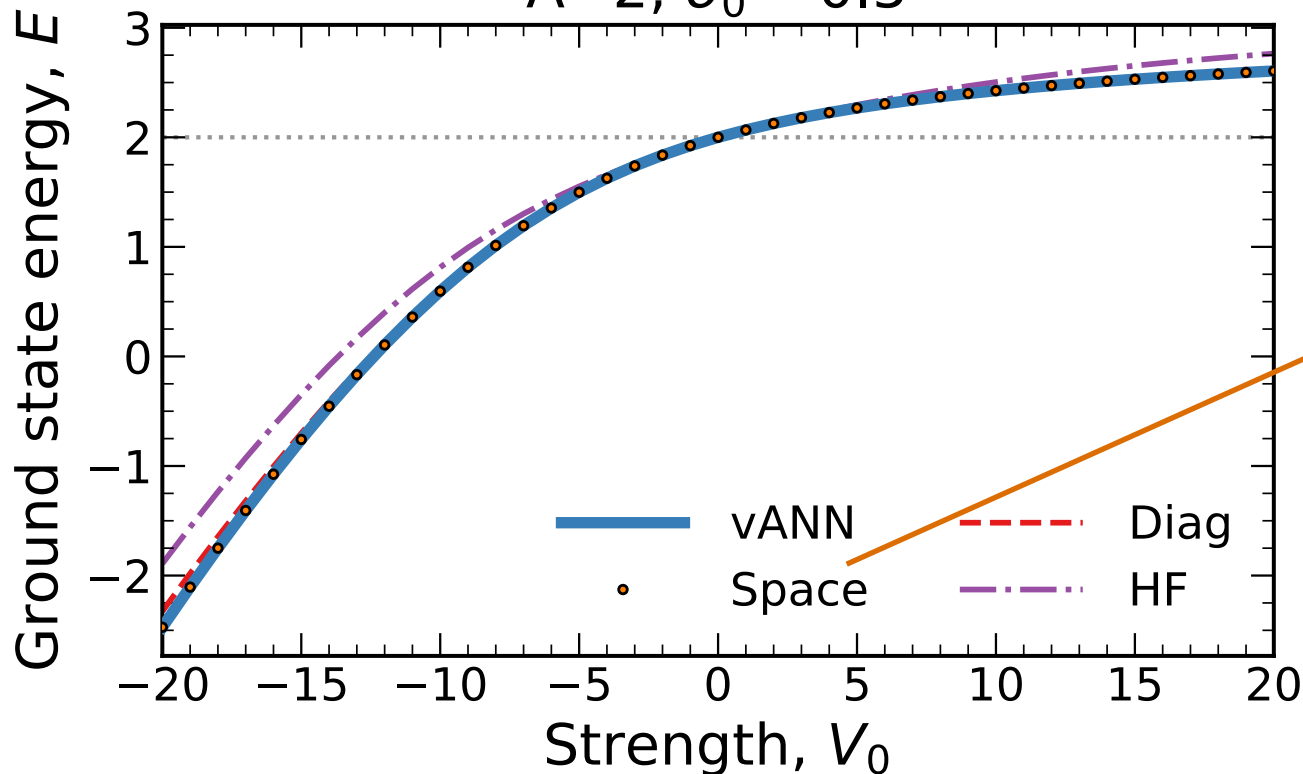
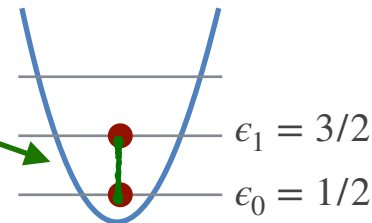
$$A=2, \sigma_0 = 0.5$$



$$\Psi(x_1, x_2) = \begin{vmatrix} \phi_1(x_1, x_2) & \phi_1(x_2, x_1) \\ \phi_2(x_1, x_2) & \phi_2(x_2, x_1) \end{vmatrix}$$

$$\hat{H} = - \sum_i \frac{\partial_{x_i}^2}{2} + \sum_i \frac{x_i^2}{2} + \sum_{i < j} \frac{V_0}{\sqrt{2\pi\sigma_0}} \exp\left(-\frac{(x_i - x_j)^2}{2\sigma_0^2}\right)$$

$$A=2, \sigma_0 = 0.5$$



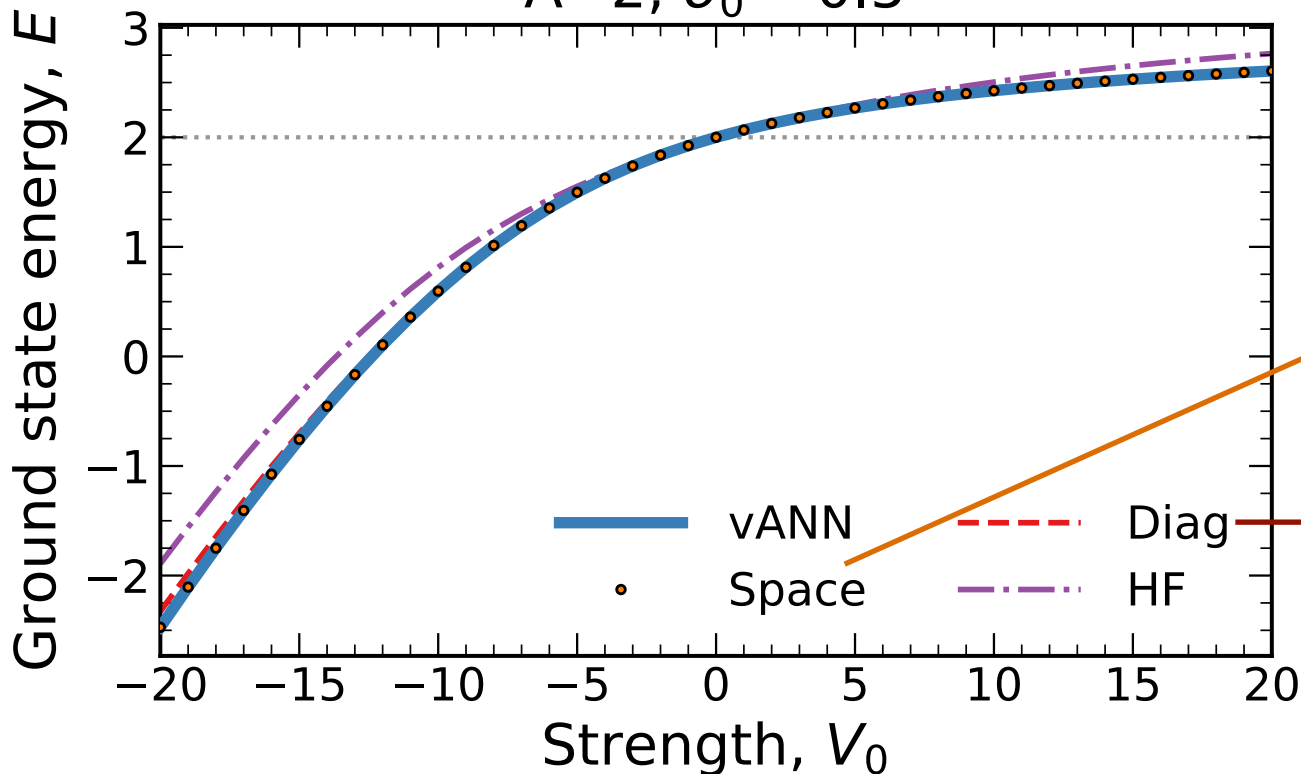
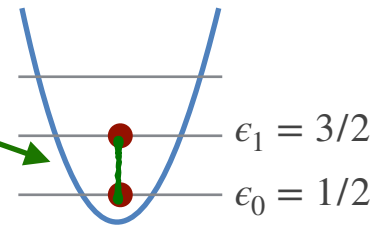
$\Psi(x_1, x_2) = \psi(X)\phi(r)$   
 1.  $\psi$  is HO  
 2. Solve  $\phi$  on 1D space grid

$$\Psi(x_1, x_2) = \begin{vmatrix} \phi_1(x_1, x_2) & \phi_1(x_2, x_1) \\ \phi_2(x_1, x_2) & \phi_2(x_2, x_1) \end{vmatrix}$$



$$\hat{H} = - \sum_i \frac{\partial_{x_i}^2}{2} + \sum_i \frac{x_i^2}{2} + \sum_{i < j} \frac{V_0}{\sqrt{2\pi\sigma_0}} \exp\left(-\frac{(x_i - x_j)^2}{2\sigma_0^2}\right)$$

$$A=2, \sigma_0 = 0.5$$



$\Psi(x_1, x_2) = \psi(X)\phi(r)$

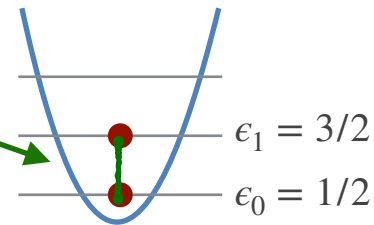
1.  $\psi$  is HO
2. Solve  $\phi$  on 1D space grid

Full CI on HO basis

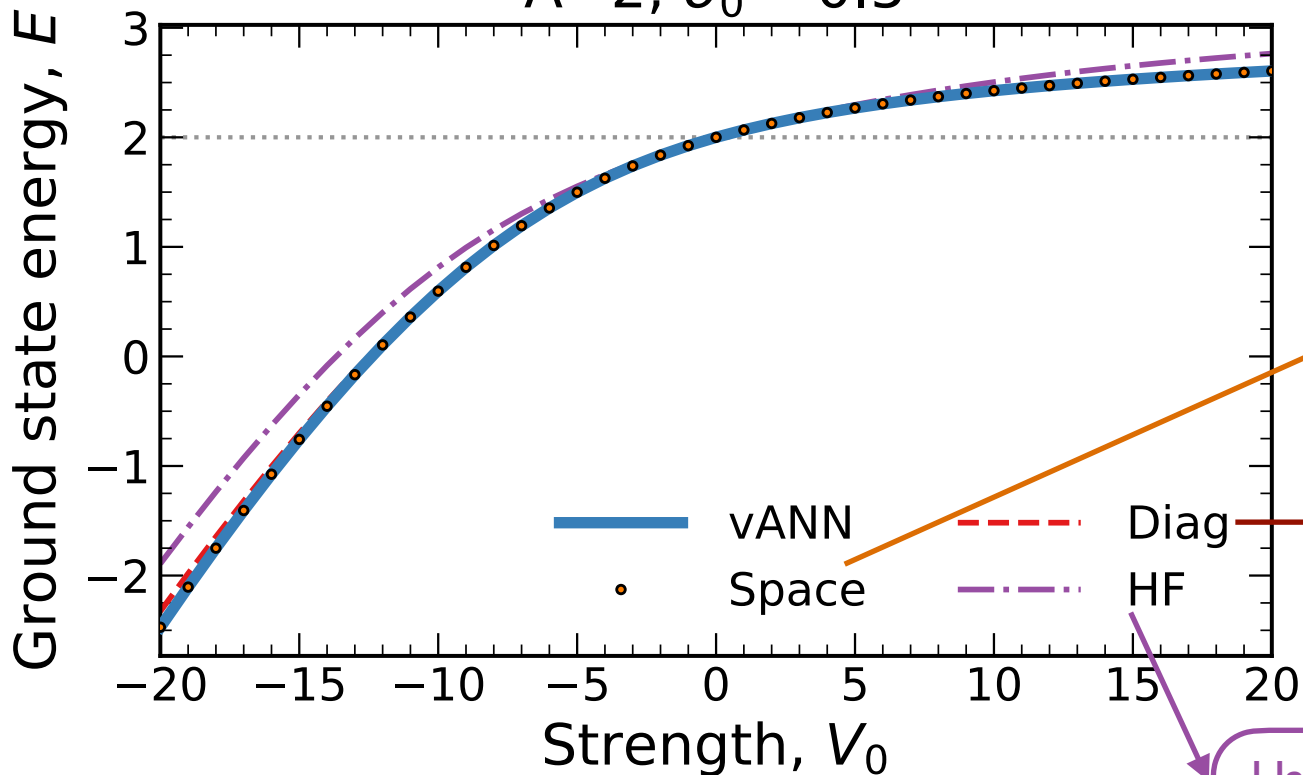
$$\Psi(x_1, x_2) = \begin{vmatrix} \phi_1(x_1, x_2) & \phi_1(x_2, x_1) \\ \phi_2(x_1, x_2) & \phi_2(x_2, x_1) \end{vmatrix}$$

# A=2, finite range

$$\hat{H} = - \sum_i \frac{\partial_{x_i}^2}{2} + \sum_i \frac{x_i^2}{2} + \sum_{i < j} \frac{V_0}{\sqrt{2\pi\sigma_0}} \exp\left(-\frac{(x_i - x_j)^2}{2\sigma_0^2}\right)$$



A=2,  $\sigma_0 = 0.5$



$\Psi(x_1, x_2) = \psi(X)\phi(r)$

1.  $\psi$  is HO
2. Solve  $\phi$  on 1D space grid

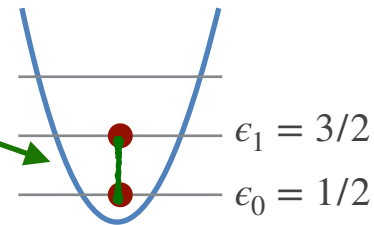
Full CI on HO basis

Hartree-Fock:  
 | Slater determinant  
 No backflow

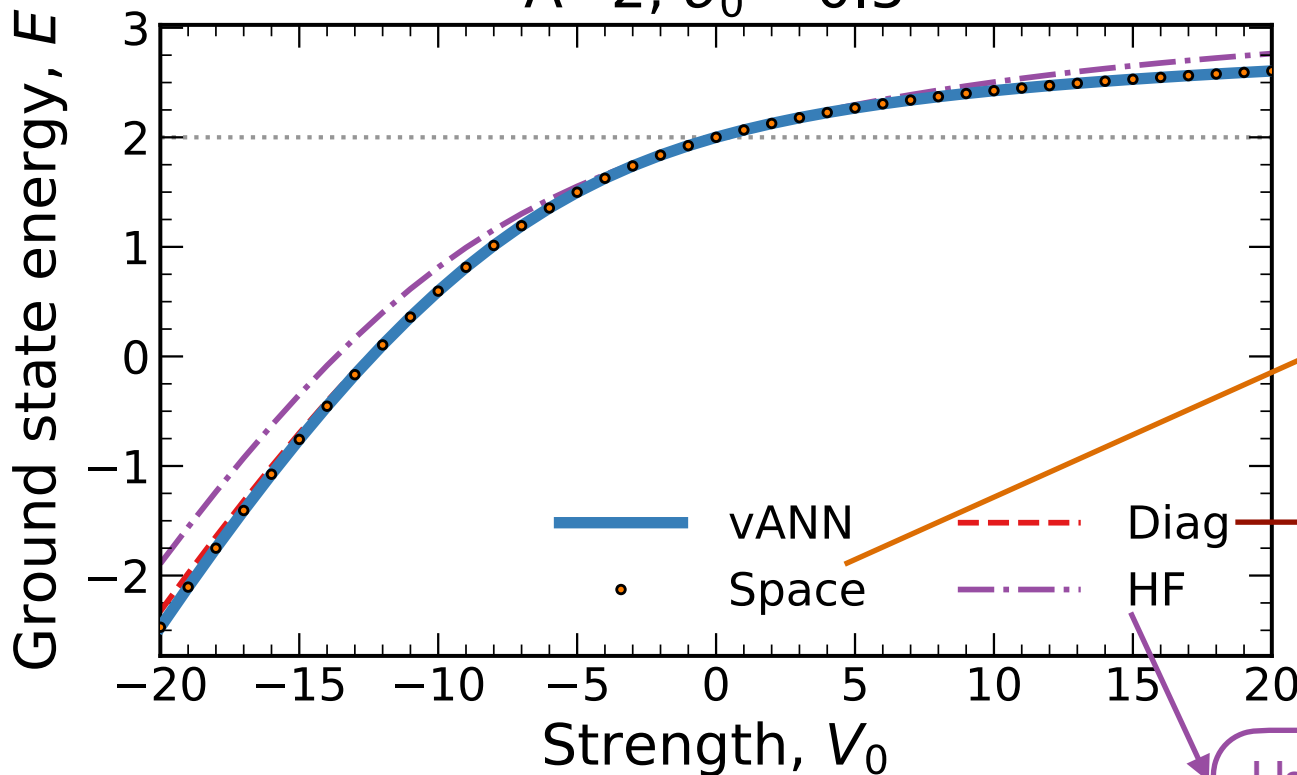
$$\Psi(x_1, x_2) = \begin{vmatrix} \phi_1(x_1, x_2) & \phi_1(x_2, x_1) \\ \phi_2(x_1, x_2) & \phi_2(x_2, x_1) \end{vmatrix}$$

# A=2, finite range

$$\hat{H} = - \sum_i \frac{\partial_{x_i}^2}{2} + \sum_i \frac{x_i^2}{2} + \sum_{i < j} \frac{V_0}{\sqrt{2\pi\sigma_0}} \exp\left(-\frac{(x_i - x_j)^2}{2\sigma_0^2}\right)$$



A=2,  $\sigma_0 = 0.5$



$\Psi(x_1, x_2) = \psi(X)\phi(r)$

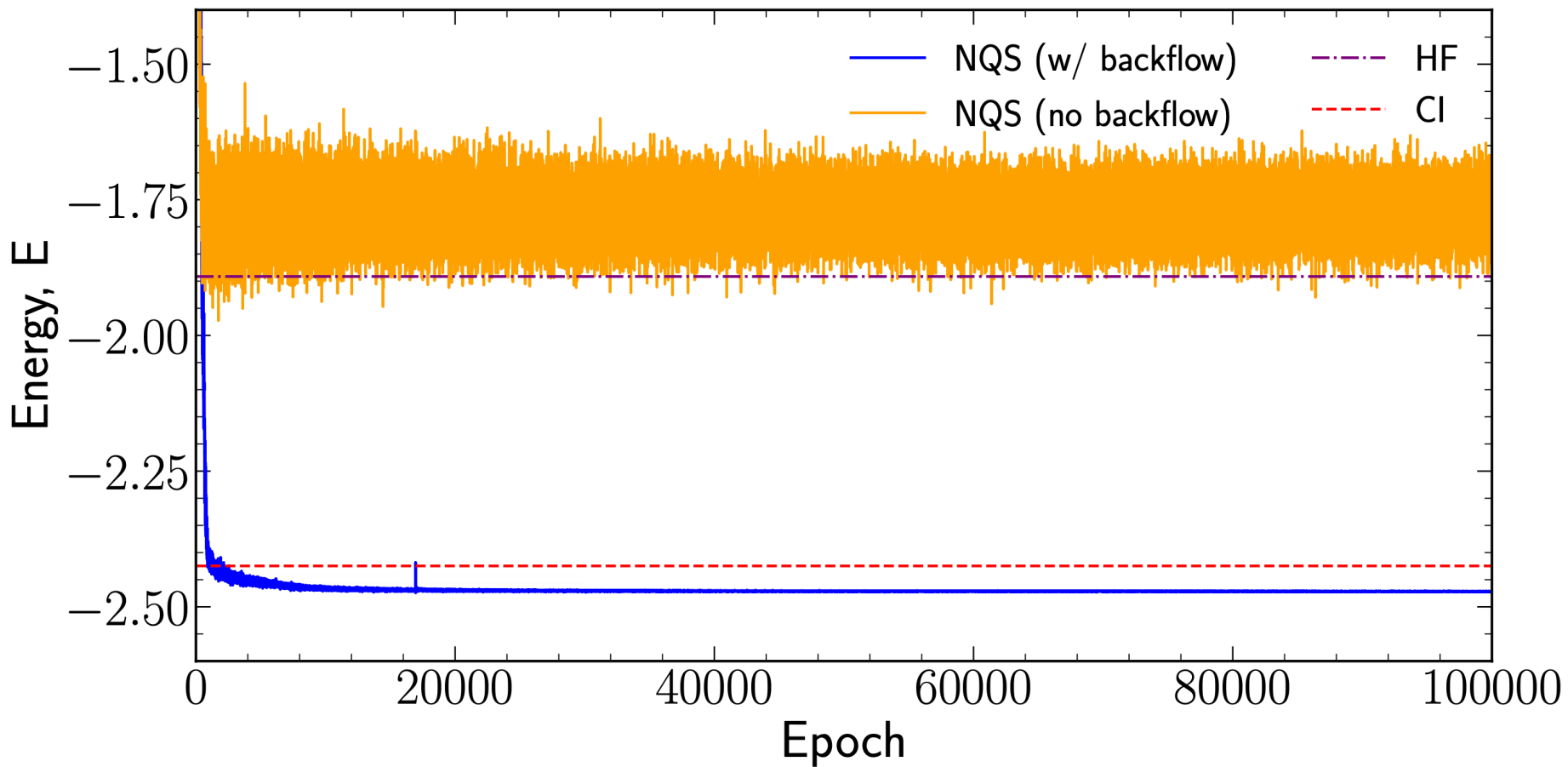
- $\psi$  is HO
- Solve  $\phi$  on 1D space grid

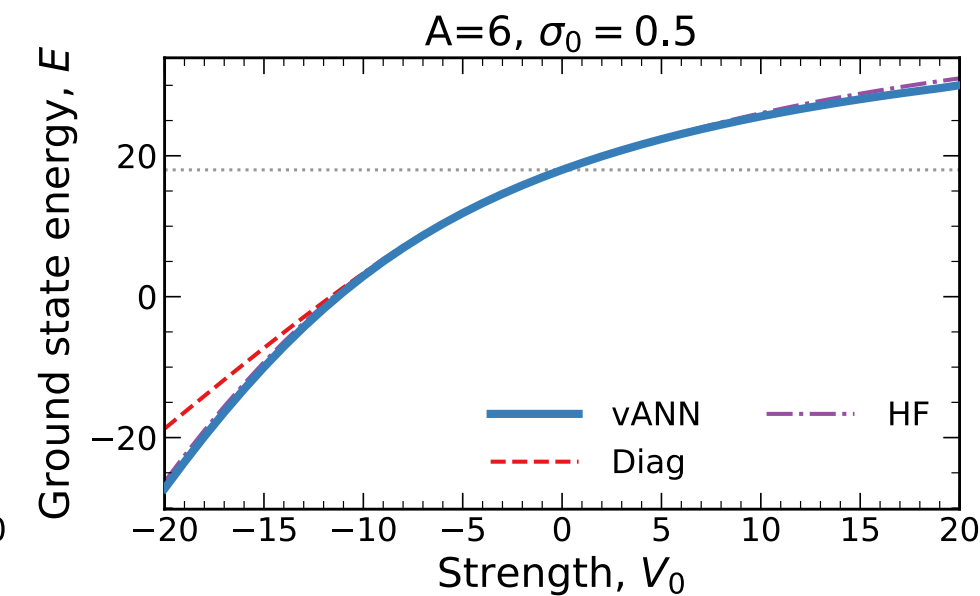
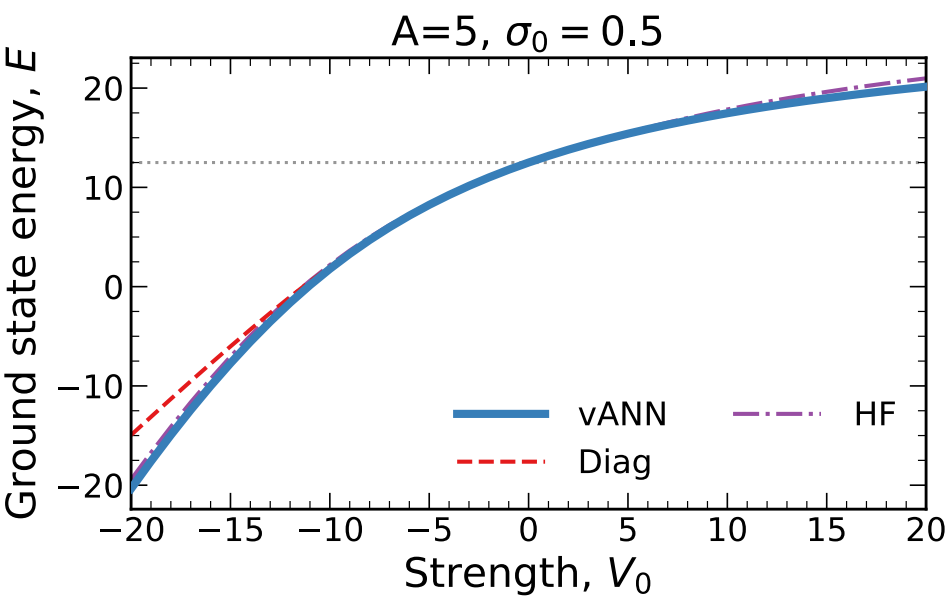
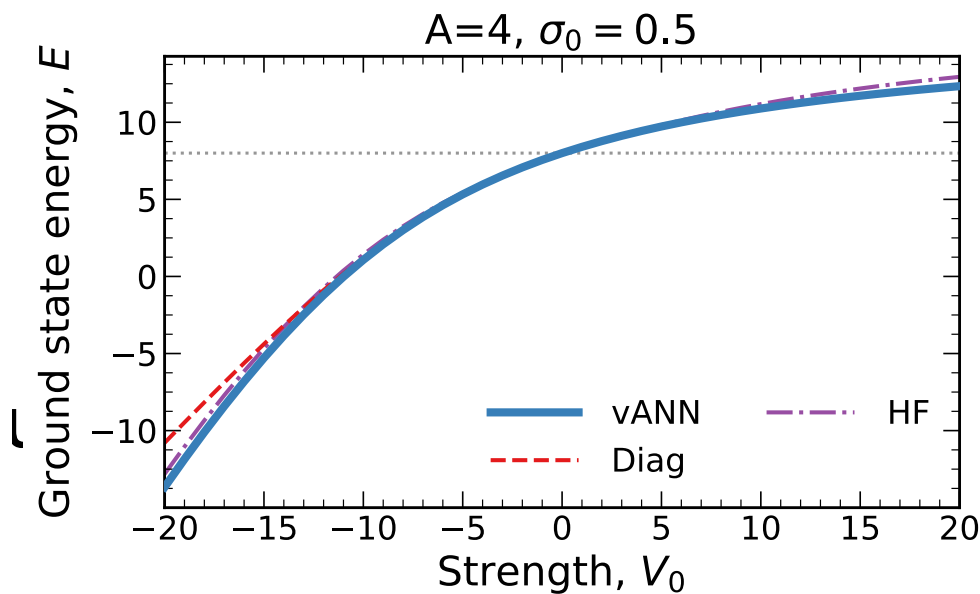
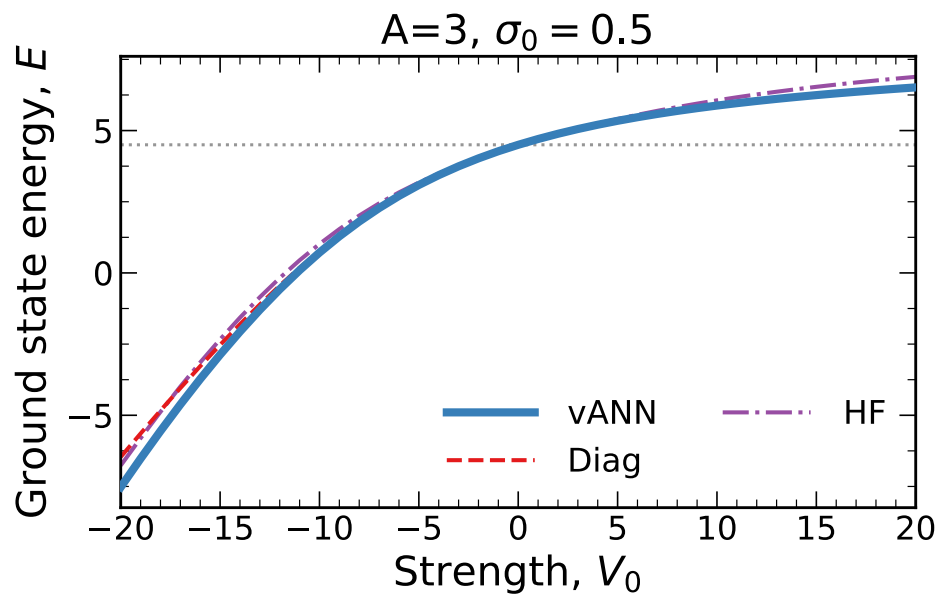
Full CI on HO basis

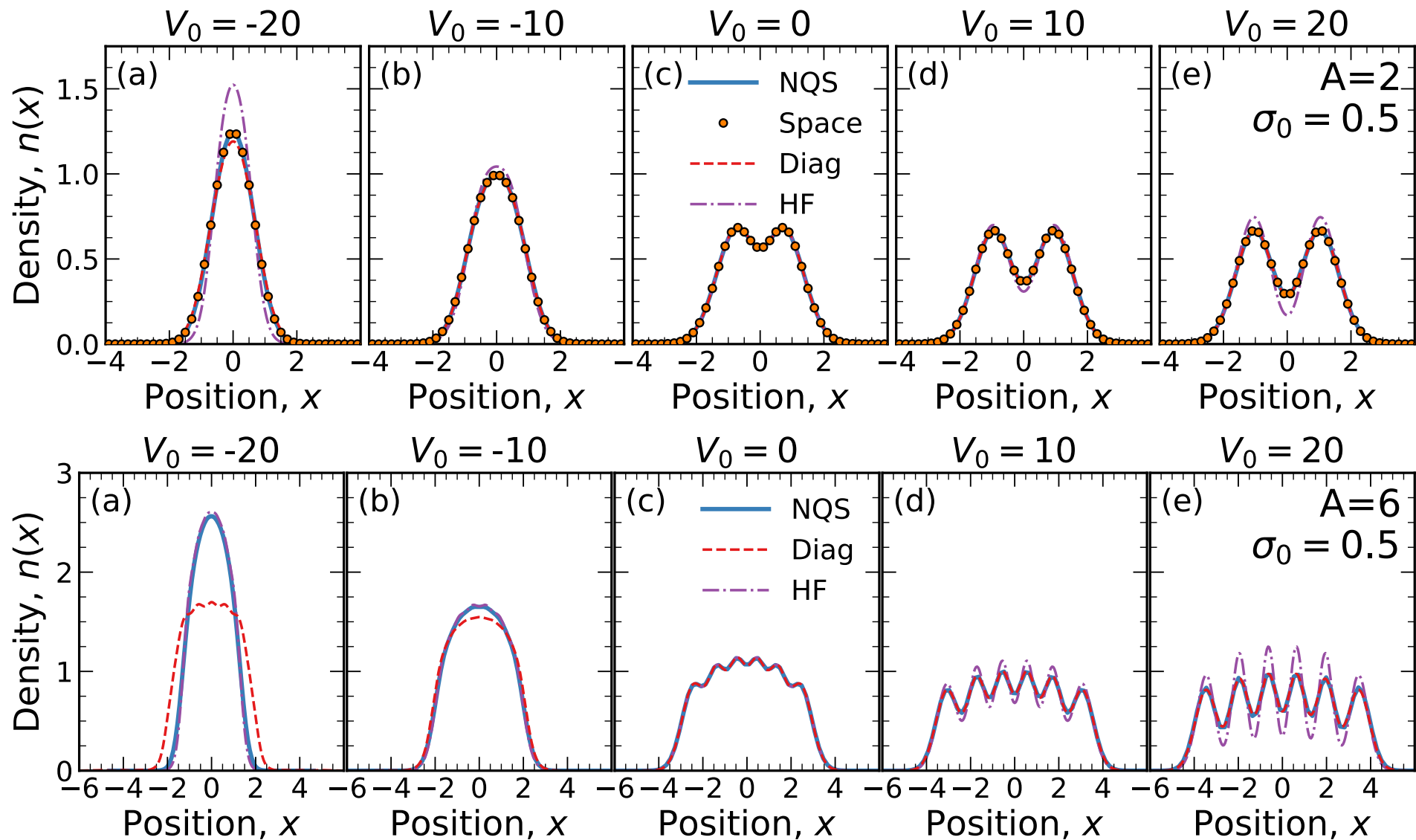
Hartree-Fock:  
 | Slater determinant  
 No backflow

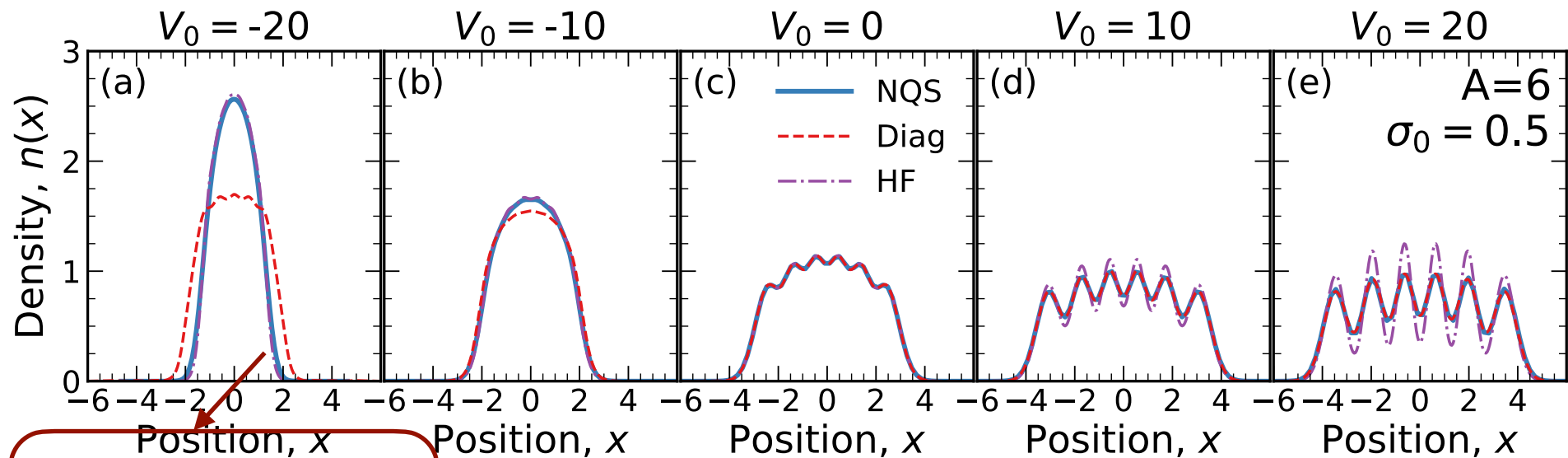
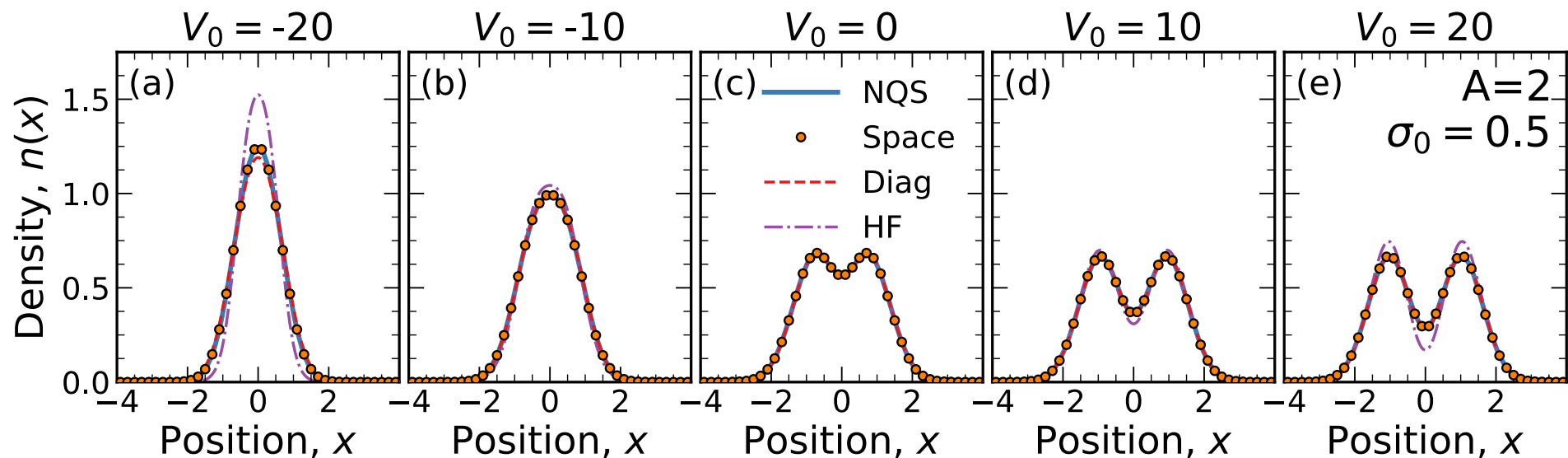
$$\Psi(x_1, x_2) = \begin{vmatrix} \phi_1(x_1, \blacksquare) & \phi_1(x_2, \blacksquare) \\ \phi_2(x_1, \blacksquare) & \phi_2(x_2, \blacksquare) \end{vmatrix}$$

$$A=2, \sigma_0=0.5, V_0=-20$$

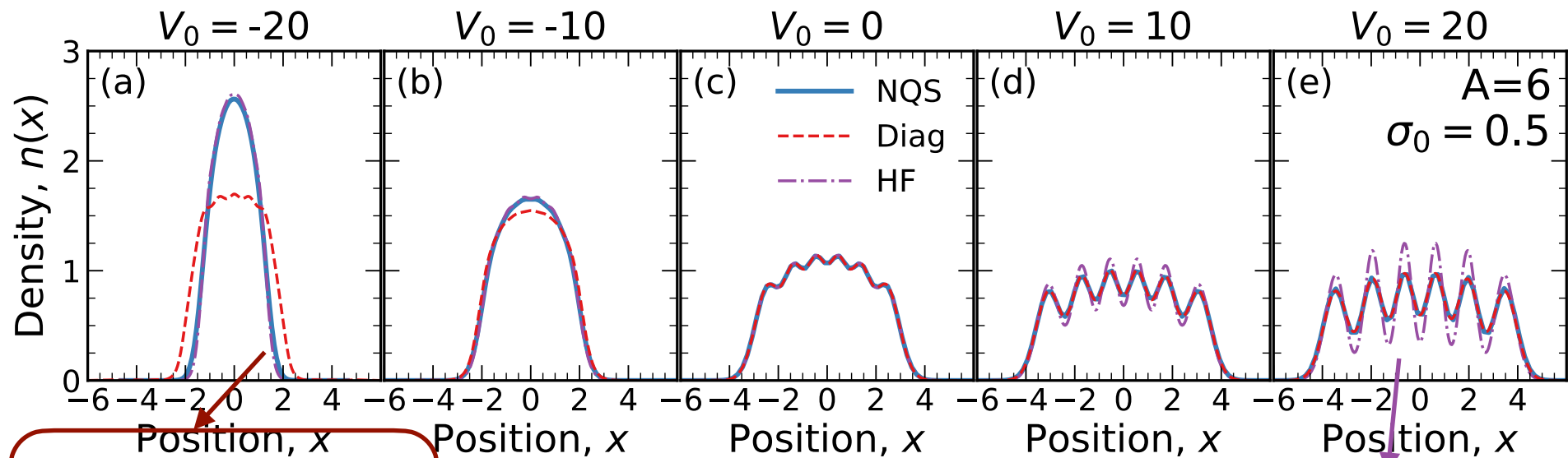
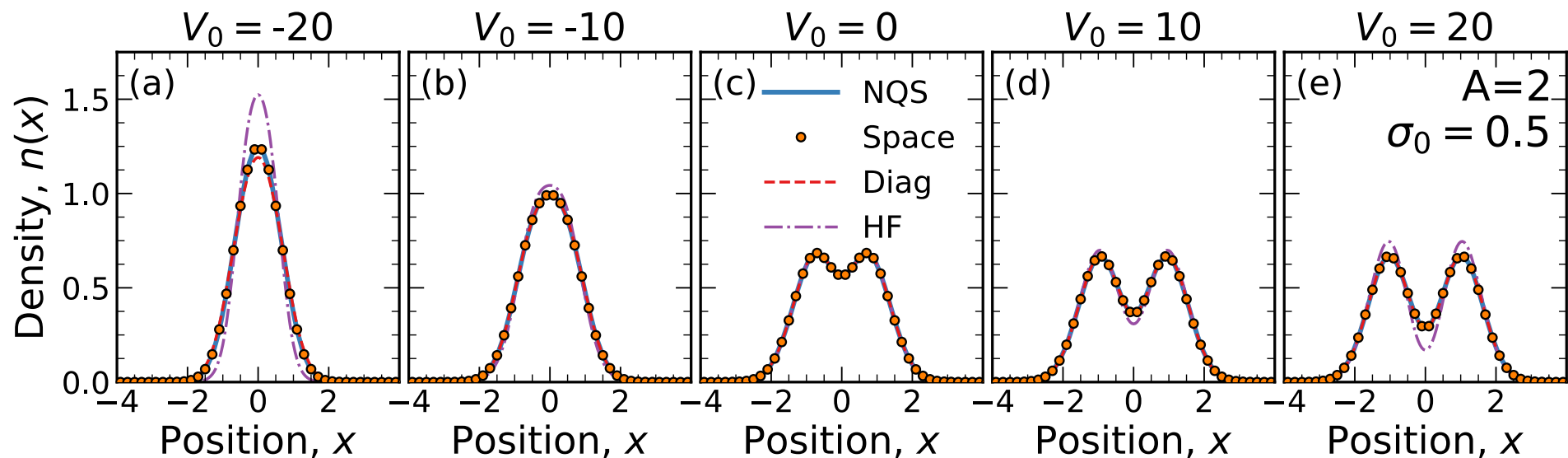








Boson-Fermion duality

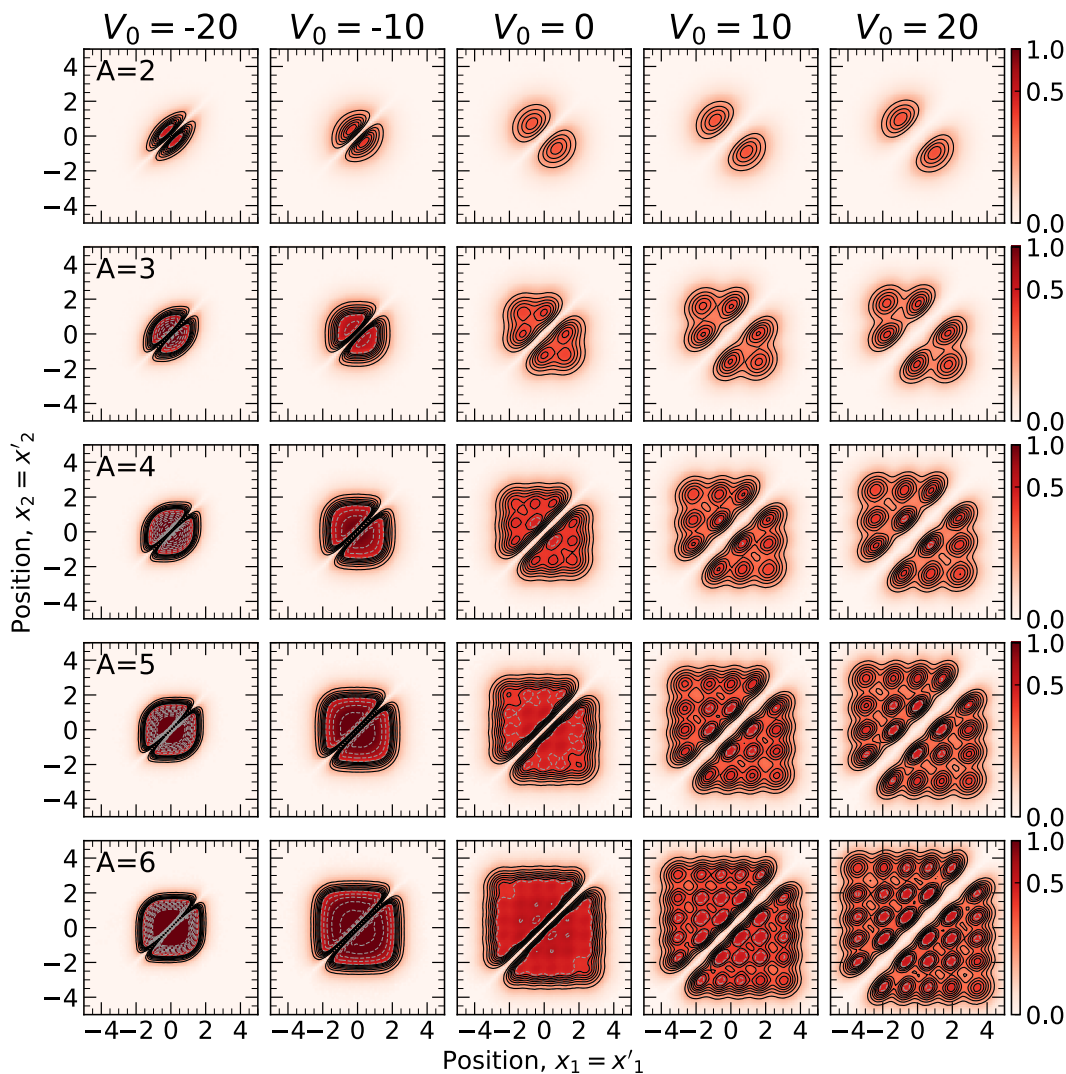


Boson-Fermion duality

Wigner crystallisation

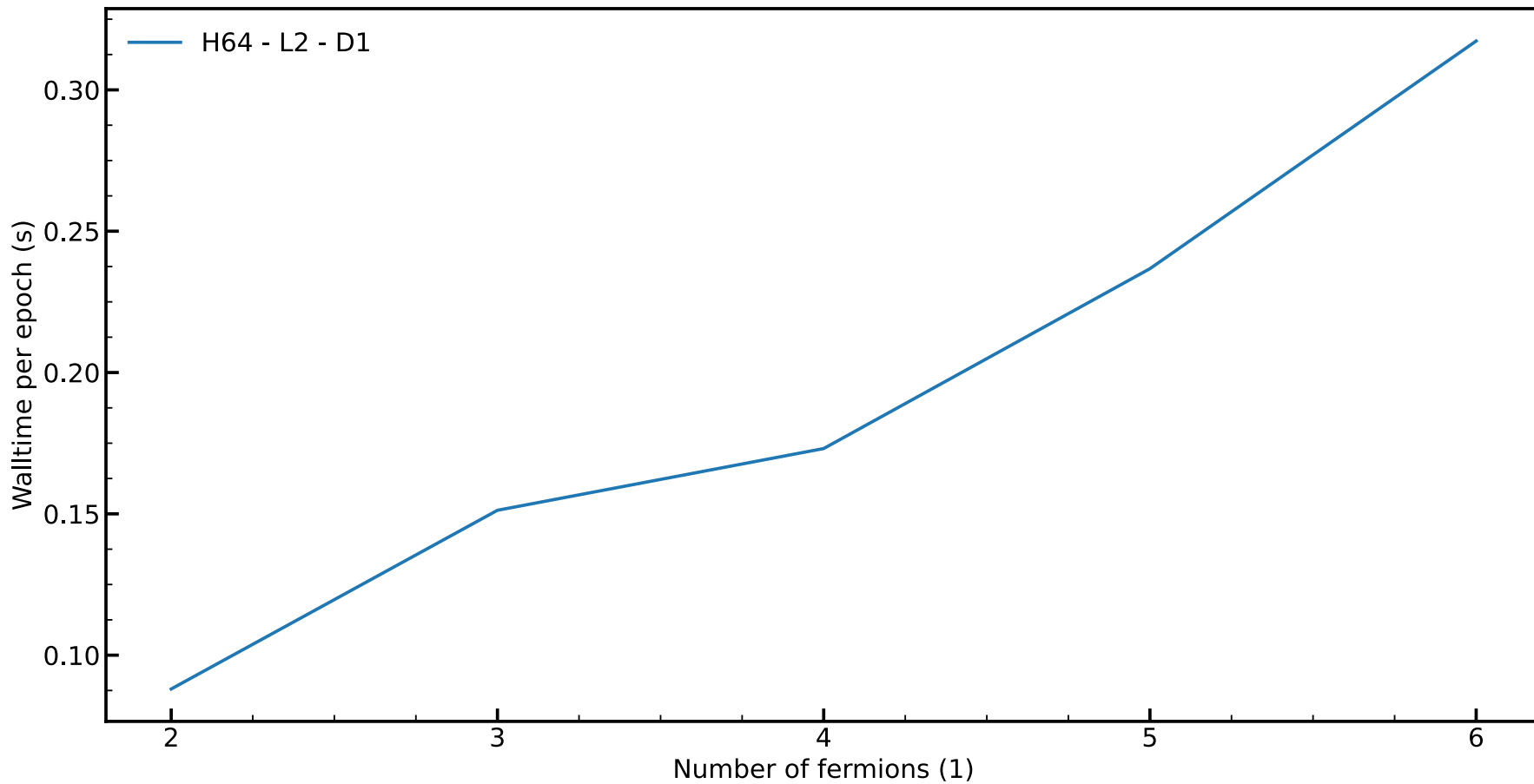


# Pair distribution function



Free case

$$g(x_1, x_2) = \frac{e^{-x_1^2 - x_2^2}}{\pi} (x_1 - x_2)^2 \mathcal{G}^{(A)}(x_1, x_2)$$



- Not so steep in practice!

## ● Nuclear physics

## ● Neural Quantum States

## ● Deuteron

JWT Keeble & A. Rios, Phys. Lett. B **809**, 135743 (2020)

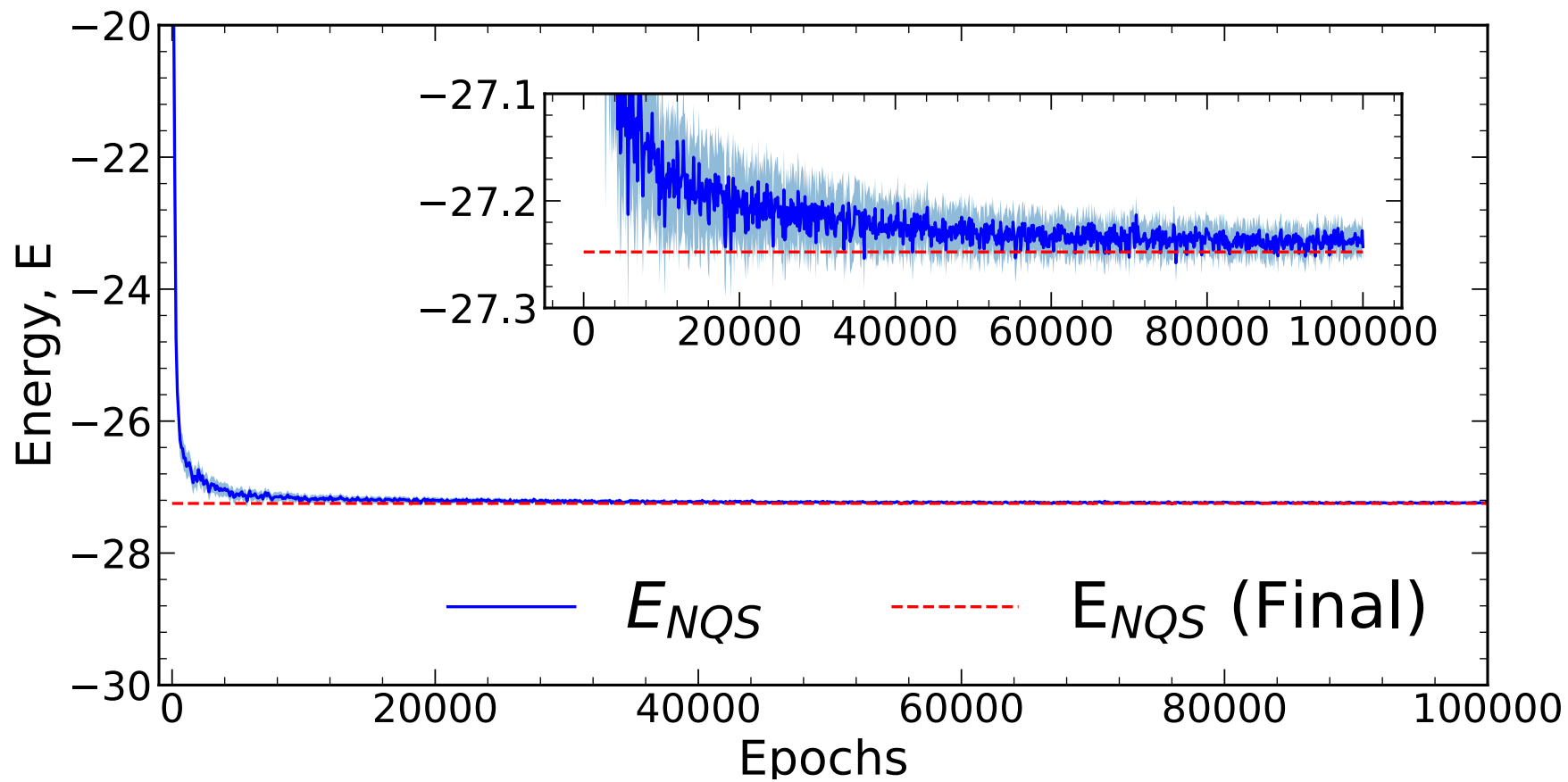
J Rozalen-Sarmiento, JWT Keeble & A. Rios, Eur. Phys. J. Plus **139**, 189 (2024)

## ● 1D spinless (polarised) systems ( $A > 2$ )

J.W.T. Keeble, M. Drissi, A. Rojo-Francàs, B. Juliá-Díaz & A. Rios, Phys. Rev. A **108**, 063320 (2023),  
arxiv:2304.04725

## ● Optimisation strategy

M. Drissi, J.W.T. Keeble, J. Rozalen-Sarmiento & A. Rios, arxiv:2401.17550



- **Cost function:**  $E(\theta)$

- **Goal**

$$E^* = \min_{\theta \in \mathbb{R}^D} E(\theta)$$

$$\theta^* = \operatorname{argmin}_{\theta \in \mathbb{R}^D} E(\theta)$$

- **Problems**

- Dimensionality of the problem  $D \sim 10^5$
- $E(\theta)$  highly nonlinear

- **Solution**

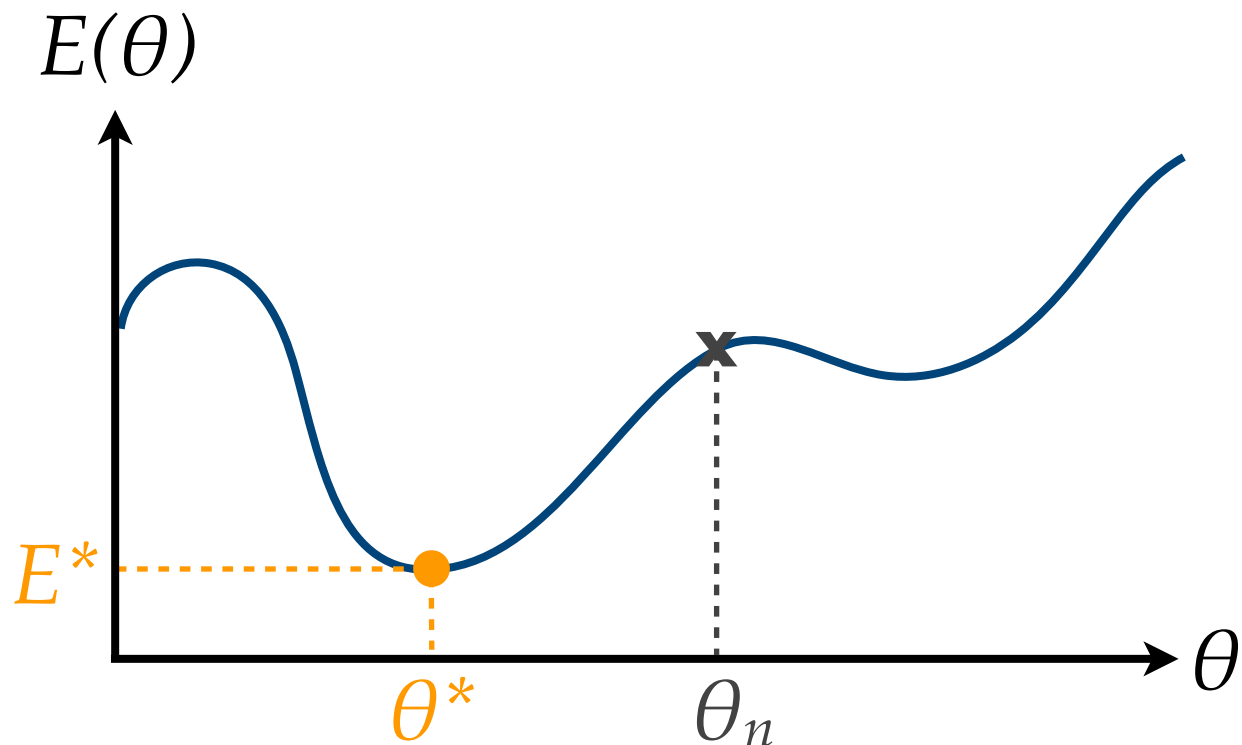
- Solve a sequence of linear/quadratic optimizations

- Epoch  $n$

$$E_{n+1} = E(\theta_{n+1})$$

$$\theta_{n+1} = \theta_n + \operatorname{argmin}_{\delta \in T_n} M_n(\delta)$$

- with **local model**  $M_n(\delta)$  & **trust region**  $T_n$

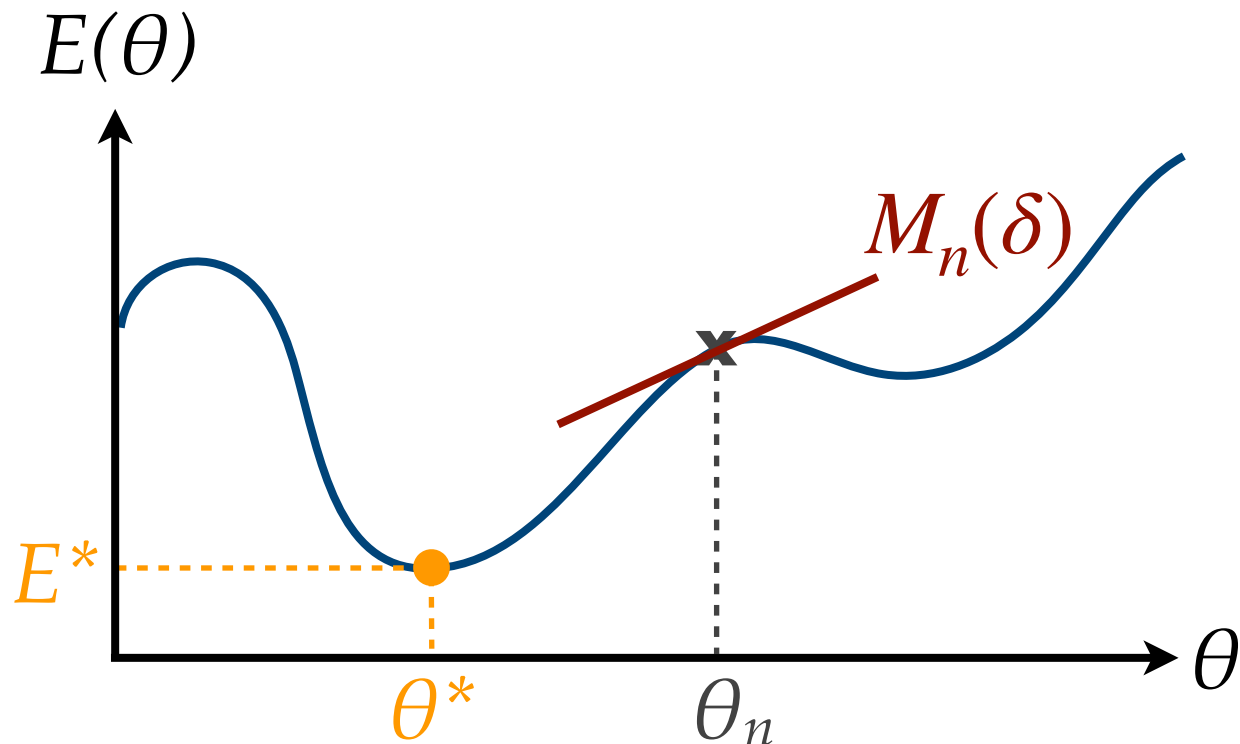


- Epoch  $n$

$$E_{n+1} = E(\theta_{n+1})$$

$$\theta_{n+1} = \theta_n + \operatorname{argmin}_{\delta \in T_n} M_n(\delta)$$

- with **local model**  $M_n(\delta)$  & **trust region**  $T_n$

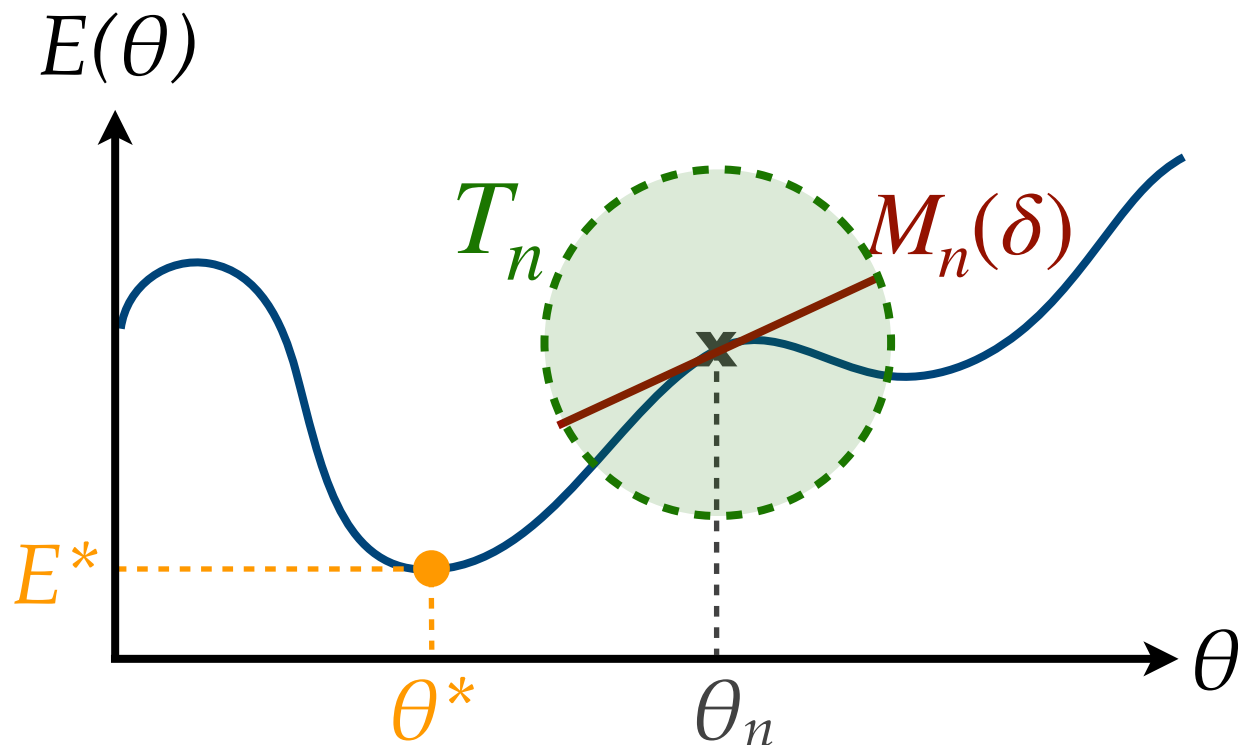


- Epoch  $n$

$$E_{n+1} = E(\theta_{n+1})$$

$$\theta_{n+1} = \theta_n + \operatorname{argmin}_{\delta \in T_n} M_n(\delta)$$

- with **local model**  $M_n(\delta)$  & **trust region**  $T_n$



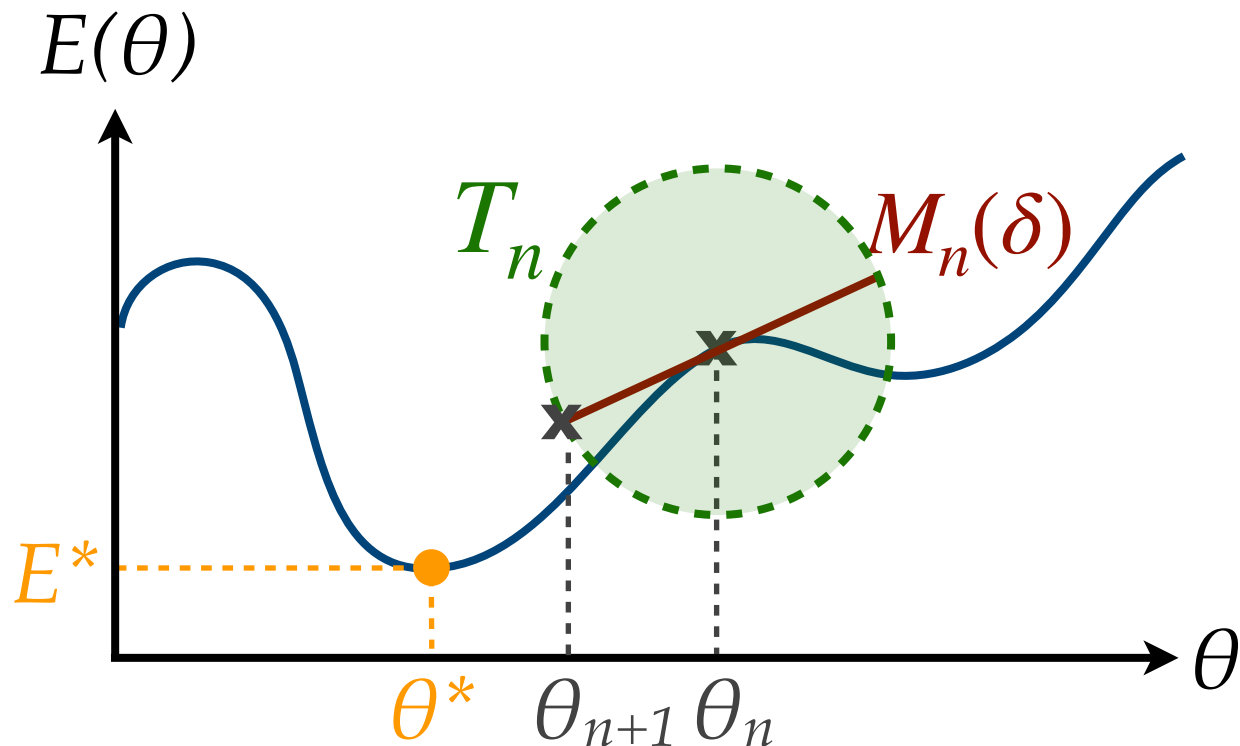


- Epoch  $n$

$$E_{n+1} = E(\theta_{n+1})$$

$$\theta_{n+1} = \theta_n + \operatorname{argmin}_{\delta \in T_n} M_n(\delta)$$

- with **local model**  $M_n(\delta)$  & **trust region**  $T_n$

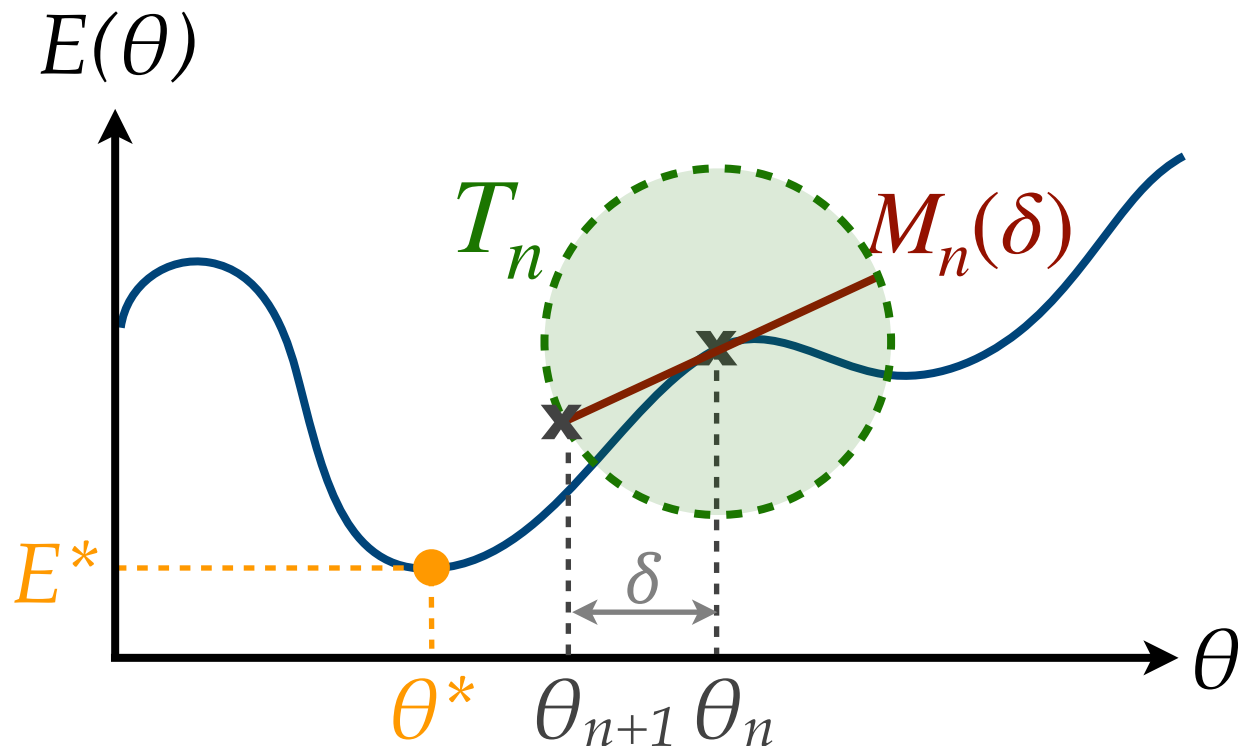


- Epoch  $n$

$$E_{n+1} = E(\theta_{n+1})$$

$$\theta_{n+1} = \theta_n + \operatorname{argmin}_{\delta \in T_n} M_n(\delta)$$

- with **local model**  $M_n(\delta)$  & **trust region**  $T_n$



# Regularized quadratic model

- Model
- Regularization

$$M_n(\delta) = \frac{1}{2} \delta^T Q \delta + L^T \delta + C$$

$$M_n(\delta) \leftarrow M_n(\delta) + \frac{1}{2} \delta^T R_n \delta$$

## Gradient descent

$$M_n(\delta) = \nabla E(\theta_n)^T \delta + E(\theta_n)$$

$$T_n \equiv \{ \delta : \|\delta\|_2 \leq \alpha \|\nabla E(\theta_n)\|_2 \}$$

$\alpha \equiv$  Learning rate

$$\delta_n = -\alpha \nabla E(\theta_n)$$

## Natural gradient descent

Hessian matrix

$$H_{ij} = \partial_{\theta_i} \partial_{\theta_j} E(\theta)$$

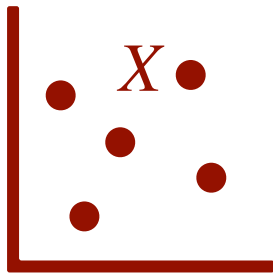
$$M_n(\delta) = \frac{1}{2} \delta^T H \delta + \nabla E(\theta)^T \delta + C$$

$$T_n \equiv \{ \delta : \delta^T H(\theta_n) \delta \leq r^2 \}$$

$$\delta_n = -H^{-1}(\theta_n) \nabla E(\theta_n)$$

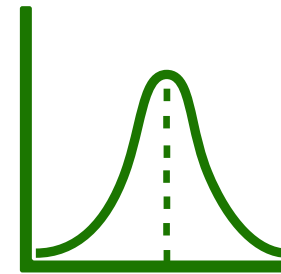
## Data

$q$  = target distribution



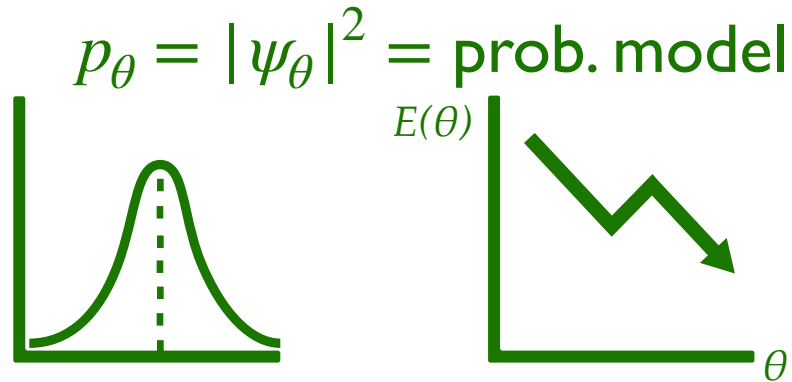
## Model

$p_\theta$  = prob. model



- **Cost function:** cross entropy  $L(\theta) = -\mathbb{E}_{X \sim q} [\ln p_\theta(X)]$
- Kullback-Leibler divergence  $D_{KL}(p_{\theta_1}, p_{\theta_2}) = -\mathbb{E}_{X \sim p_{\theta_1}} [\ln p_{\theta_1}(X) - \ln p_{\theta_2}(X)]$
- KL near minimum  $D_{KL}(p_\theta, p_{\theta+\delta\theta}) \approx \frac{1}{2} \delta^T F(\theta) \delta + O(\delta^3)$
- **Fisher** information matrix/“metric”  $F_{ij} = \mathbb{E} \left[ \partial_{\theta_i} \ln p_\theta \partial_{\theta_j} \ln p_\theta \right]$
- Update rule: natural gradient descent  $\delta_n = -F^{-1}(\theta_n) \nabla E(\theta_n)$

## Stochastic reconfiguration



- **Cost function:** min distance in imaginary time step

$$|\phi\rangle = e^{-\epsilon\hat{H}}|\psi_\theta\rangle \approx (1 - \epsilon\hat{H})|\psi_\theta\rangle \Rightarrow L(\theta) = \min_\theta D_{FS}(\phi, \psi_\theta)$$

- Fubini-Study distance

$$D_{FS}(\psi_{\theta_1}, \psi_{\theta_2}) = \arccos \sqrt{\frac{\langle \psi_{\theta_1} | \psi_{\theta_2} \rangle \langle \psi_{\theta_2} | \psi_{\theta_1} \rangle}{\langle \psi_{\theta_1} | \psi_{\theta_1} \rangle \langle \psi_{\theta_2} | \psi_{\theta_2} \rangle}}$$

- FS near minimum

$$D_{FS}(\psi_\theta, \psi_{\theta+\delta\theta}) \approx \frac{1}{2} \delta^T G(\theta) \delta + O(\delta^3)$$

- Quantum geometric tensor

$$G_{ij} = \left\langle \partial_{\theta_i} \psi_\theta \middle| \partial_{\theta_i} \psi_\theta \right\rangle - \left\langle \partial_{\theta_i} \psi_\theta \middle| \psi \right\rangle \left\langle \psi_\theta \middle| \partial_{\theta_j} \psi \right\rangle$$

- Update rule

$$\delta_n = -G^{-1}(\theta_n) \nabla E(\theta_n)$$

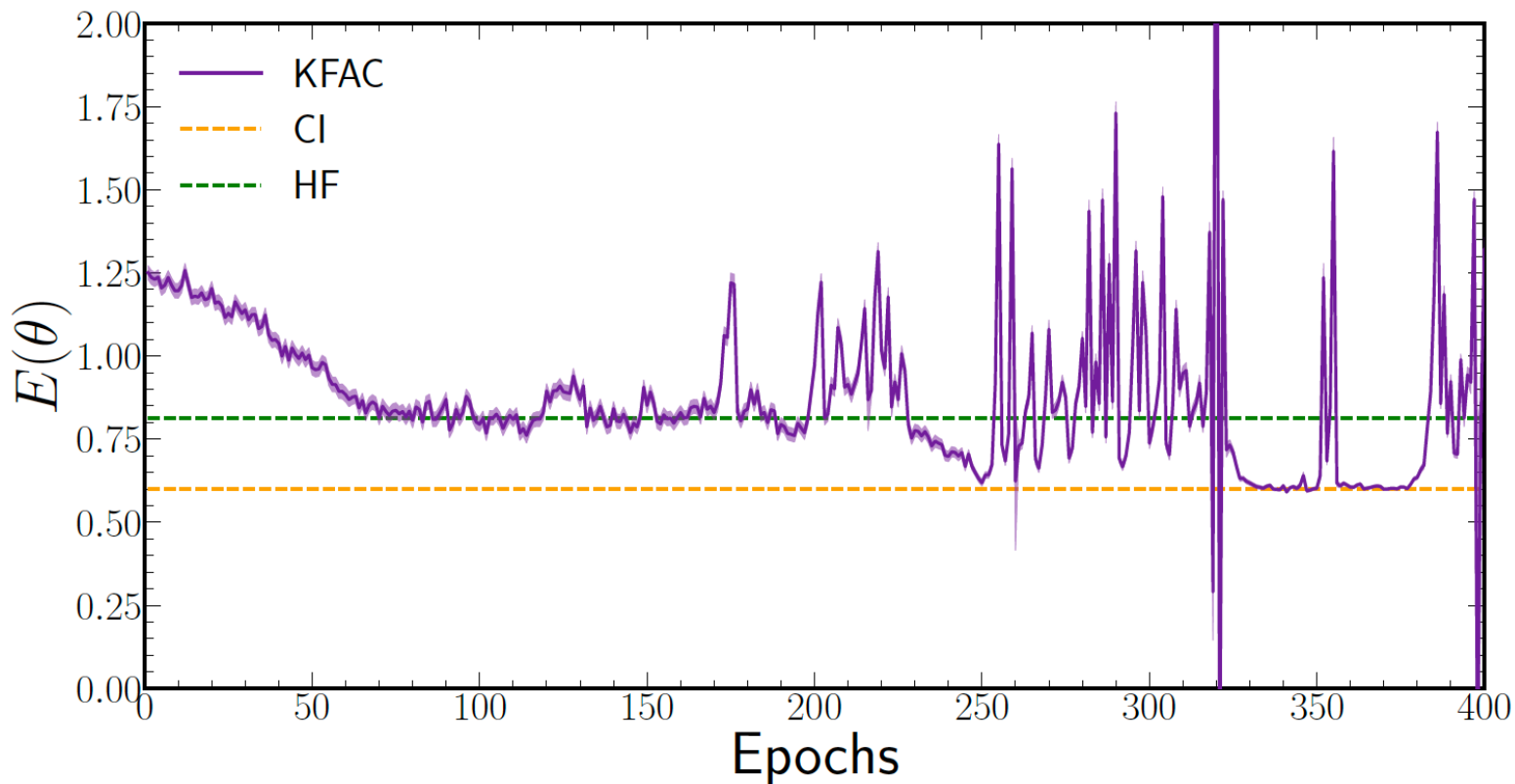
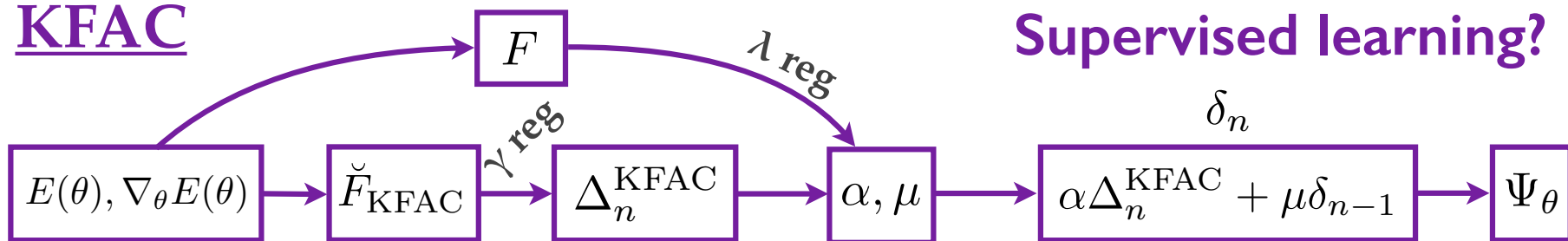
- Update rule for  $\delta$  sometimes **too crude**. Call it  $\Delta_n$
- Add **scaling factor**  $\alpha$  + **momentum**  $\mu$  & minimise:

$$M_n(\alpha\Delta_n + \mu\delta_{n-1})$$

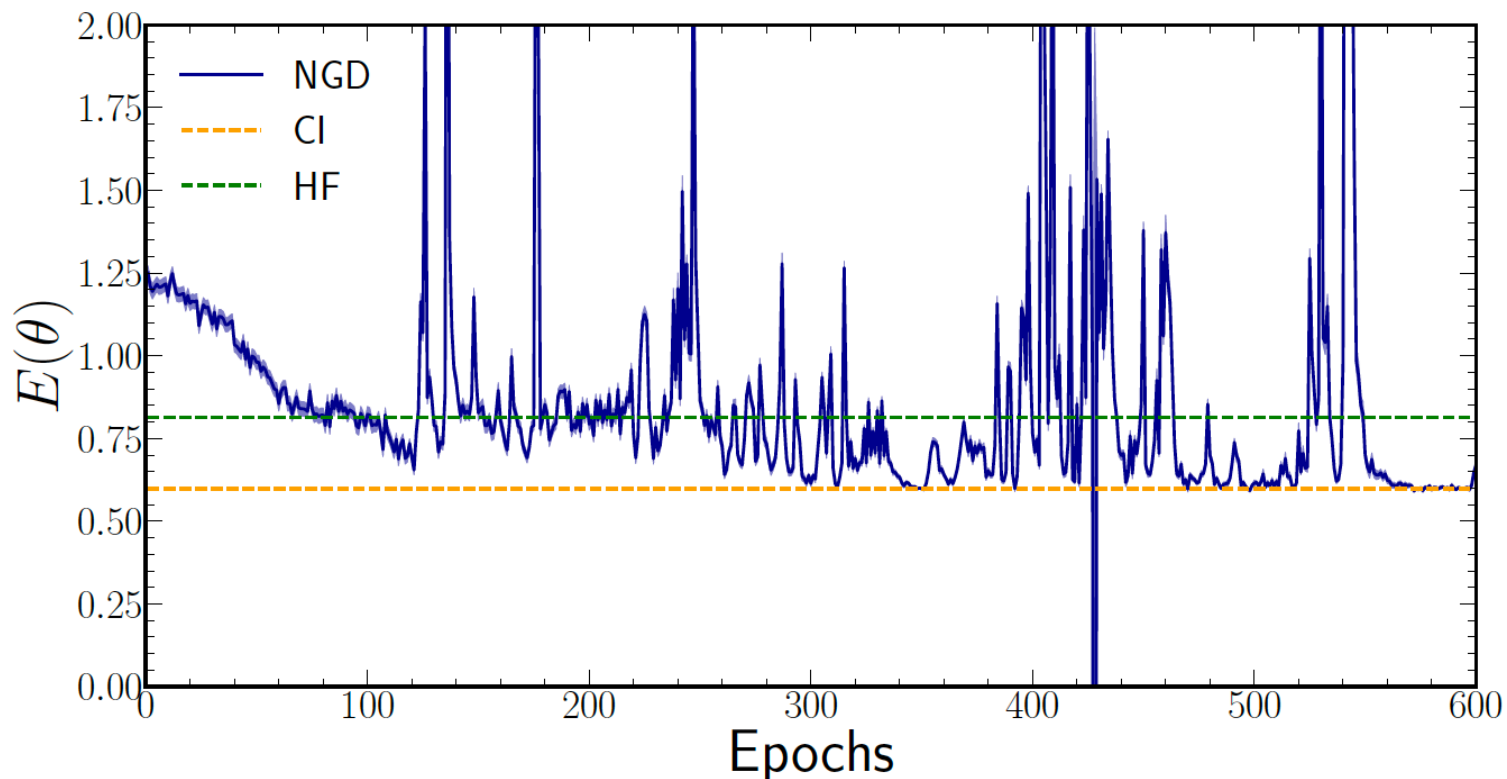
- Optimal values:

$$\begin{pmatrix} \alpha^* \\ \mu^* \end{pmatrix} = - \begin{pmatrix} \Delta_n^T F(\theta_n) \Delta_n & \Delta_n^T F(\theta_n) \delta_{n-1} \\ \delta_{n-1}^T F(\theta_n) \Delta_n & \delta_{n-1}^T F(\theta_n) \delta_{n-1} \end{pmatrix}^{-1} \begin{pmatrix} \nabla E(\theta_n)^T \Delta_n \\ \nabla E(\theta_n)^T \delta_{n-1} \end{pmatrix}$$

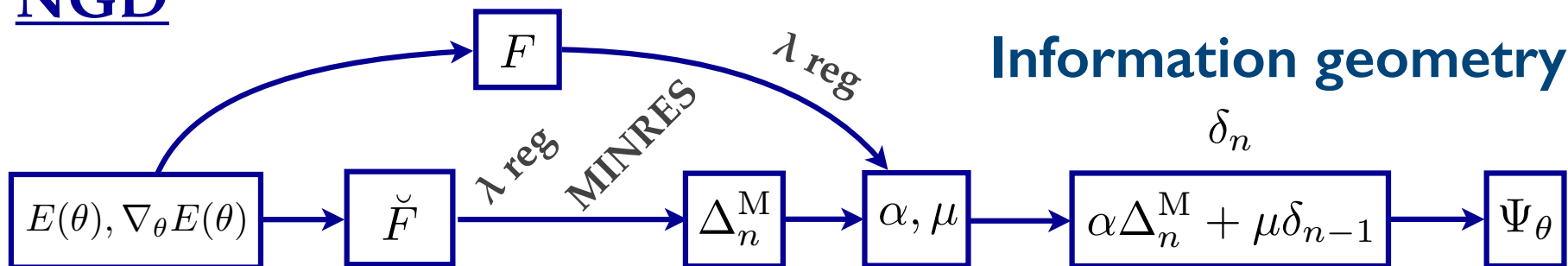
- Here, we can use **exact** Fisher (no need for inversion)
- Approximate inverse Fisher by KFAC version

KFAC:  $A = 2, V_0 = -10$ KFAC

Natural gradient Descent (NGD):  $A = 2, V_0 = -10$



**NGD**





- **Cost function** from score  $\mathcal{L}(\theta) \equiv -\mathbb{E}_{X \sim p_\theta} [\mathcal{S}(X, p_\theta)]$
- **Scoring rule:**  $\mathcal{S}(X, p_\theta)$  is any integrable function that takes as inputs a probability distribution,  $p$ , and a sample,  $X$ , and returns a real valued quantity.
- **Expected score:**  $S(p, q) \equiv \mathbb{E}_{X \sim p} [\mathcal{S}(X, q)]$
- **Proper** scoring rules:  $S(p, p) \leq S(p, q)$

Dawid, *Ann. Inst. Stat. Math.* 59, 77–93 (2007);

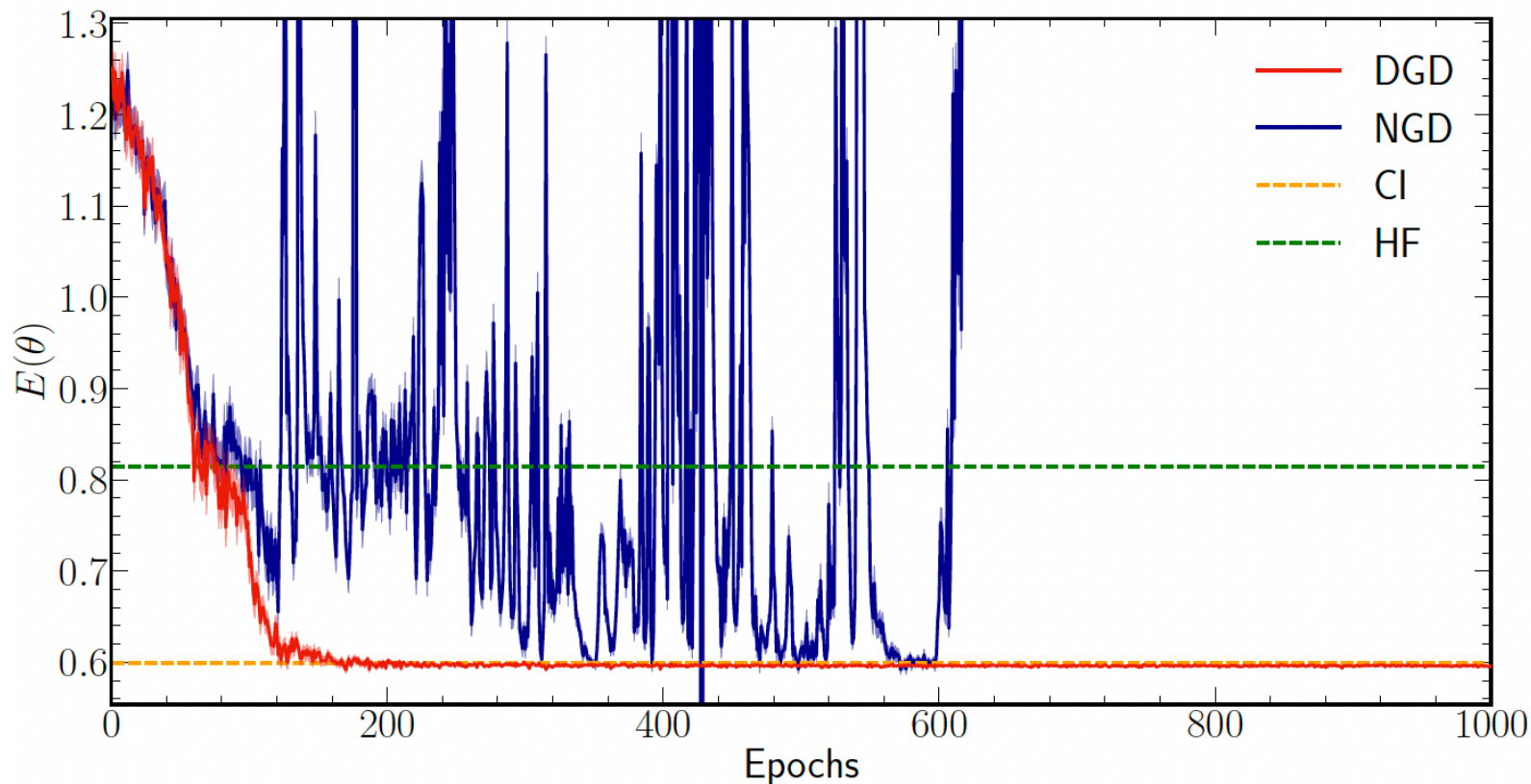
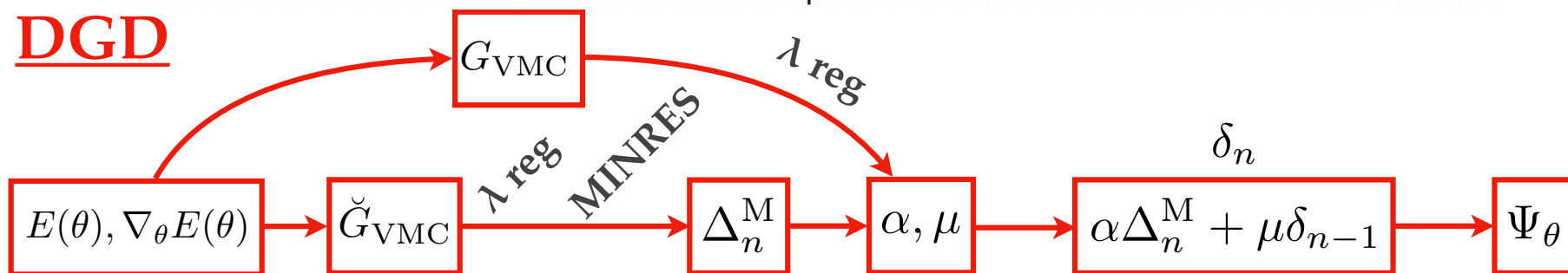
Gneiting, Raftery, *J. Am. Stat. Assoc.* 102, 359–378 (2007)

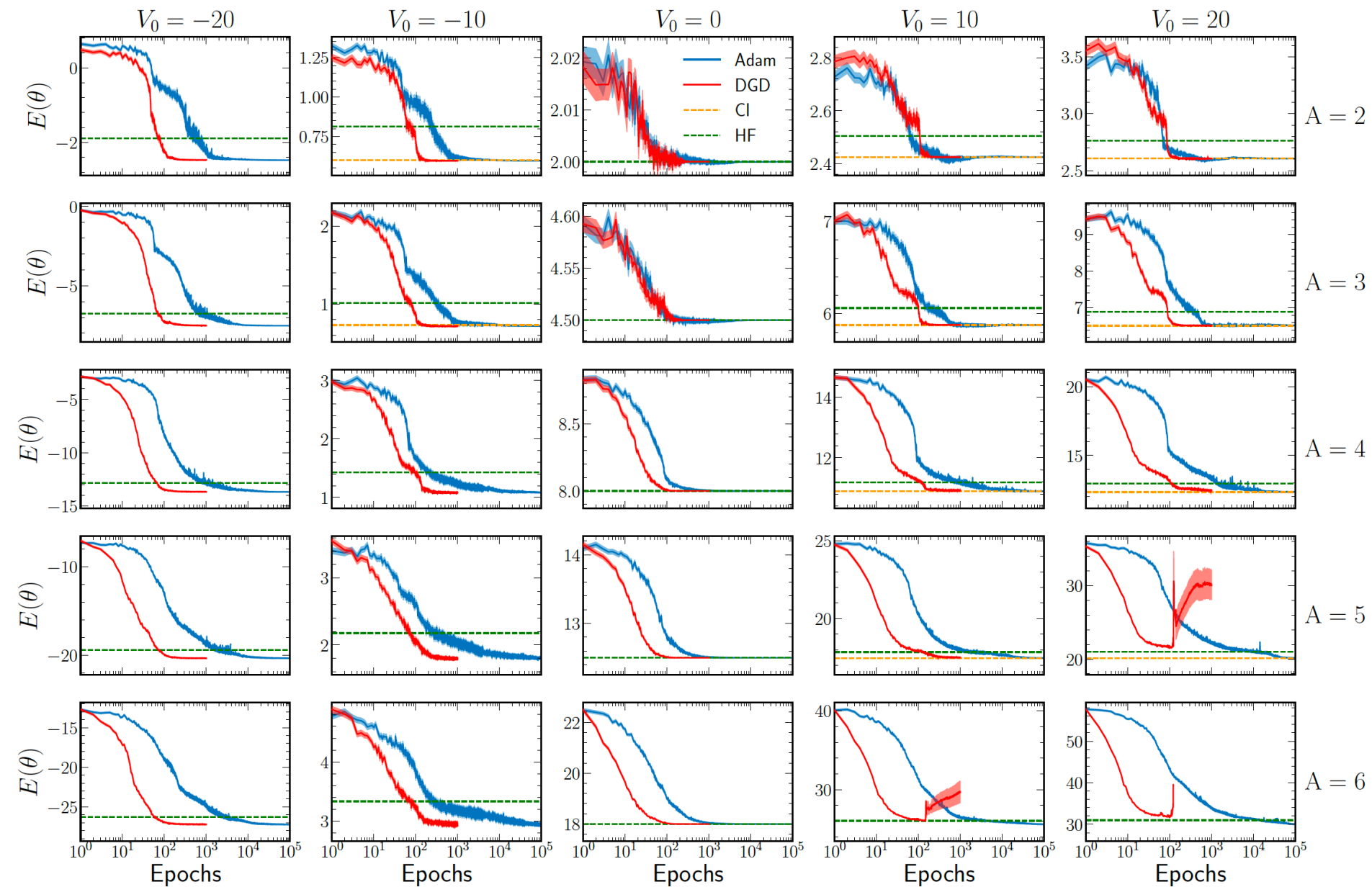
- **Generalized entropy & cross-entropy:**  $S(p, p) \leq S(p, q)$
- **Divergence:**  $D_S(p, q) = S(p, q) - S(p, p)$
- **Divergence near minimum:**  $D_S(p_\theta, p_{\theta+\delta}) = \frac{1}{2} \delta^T G_S(\theta) \delta + O(\delta^3)$
- **G is semipositive definite.**
- **Update rule: decision gradient descent**  $\delta_n = -G^{-1}(\theta_n) \nabla \mathcal{L}(\theta_n)$

- **Information geometry:**  $\mathcal{S}(X, p_\theta) = -\ln p_\theta(X)$  ✗
- **VMC scoring rule:**  $\mathcal{S}_{\text{VMC}}(X, p_\theta) = -E_{L,\theta}(X)$  ✓
- **Cost function** from score  $\mathcal{L}(\theta) = -\mathbb{E}_{X \sim p_\theta} [\mathcal{S}(X, p_\theta)] = E(\theta)$
- **Local energy:**  $E_{L,\theta}(X) \equiv \Psi_\theta(X)^{-1} H \Psi_\theta(X)$   

$$= -\frac{1}{2} \sum_{i=1}^A \left[ \partial_{x_i}^2 \ln |\Psi_\theta(X)| + (\partial_{x_i} \ln |\Psi_\theta(X)|)^2 \right] + V(X)$$

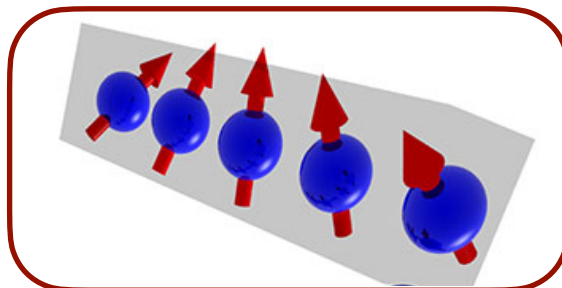
Hyvärinen score!
- **Metric:**  $G_{\text{VMC}}(\theta)_{\theta_i \theta_j} \equiv \frac{1}{4} \sum_{k=1}^A \mathbb{E}_{X \sim p_\theta} \left[ \partial_{\theta_i} \partial_{x_k} \ln p_\theta(X) \partial_{\theta_j} \partial_{x_k} \ln p_\theta(X) \right]$
- **Update rule:**  $\delta_n = -G_{\text{VMC}}^{-1}(\theta_n) \nabla E(\theta_n)$

NGD vs DGD:  $A = 2, V_0 = -10$ **DGD**



- **Machine Learning** techniques can solve **nuclear & condensed matter** many-body problems
- **Potential**: same as VMC, possibly less expensive
- **Challenges**
  - **Antisymmetry** (dealt with)
  - **Symmetries in general** (not dealt just yet)
  - **Numerical scaling** in **different** systems
  - **Optimization strategies**

$$P\Psi \rightarrow \Psi$$





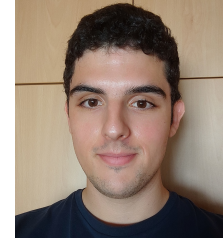
JWT Keeble



M. Drissi

UNIVERSITAT DE  
BARCELONA

J. Rozalén  
J. Gómez Micó  
A. Rojo-Francás  
B. Juliá-Díaz



[arnau.rios@ub.edu](mailto:arnau.rios@ub.edu) [twitter: @riosarnau](https://twitter.com/riosarnau)  
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