

# Resonant Electromagnetic Responses with the Lorentz Integral Transform

- Motivation
- LIT method
- Low-energy continuum observables with LIT method

# Resonant Electromagnetic Responses with the Lorentz Integral Transform

- Motivation
- LIT method: **ab initio** calculation of reactions with full FSI
- Low-energy continuum observables with LIT method

# Resonant Electromagnetic Responses with the Lorentz Integral Transform

- Motivation
- LIT method: **ab initio** calculation of reactions with full FSI
- Low-energy continuum observables with LIT method
  - S-Factor  $S_{12}$  (Coulomb potential !)**
  - Isoscalar  $0^+$  resonance of  $^4\text{He}$**

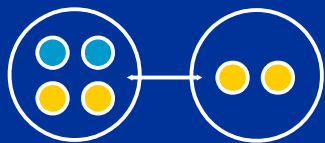
# Motivation

# 6-Body total photodisintegration

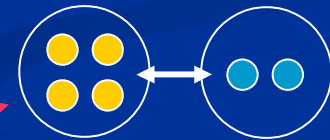
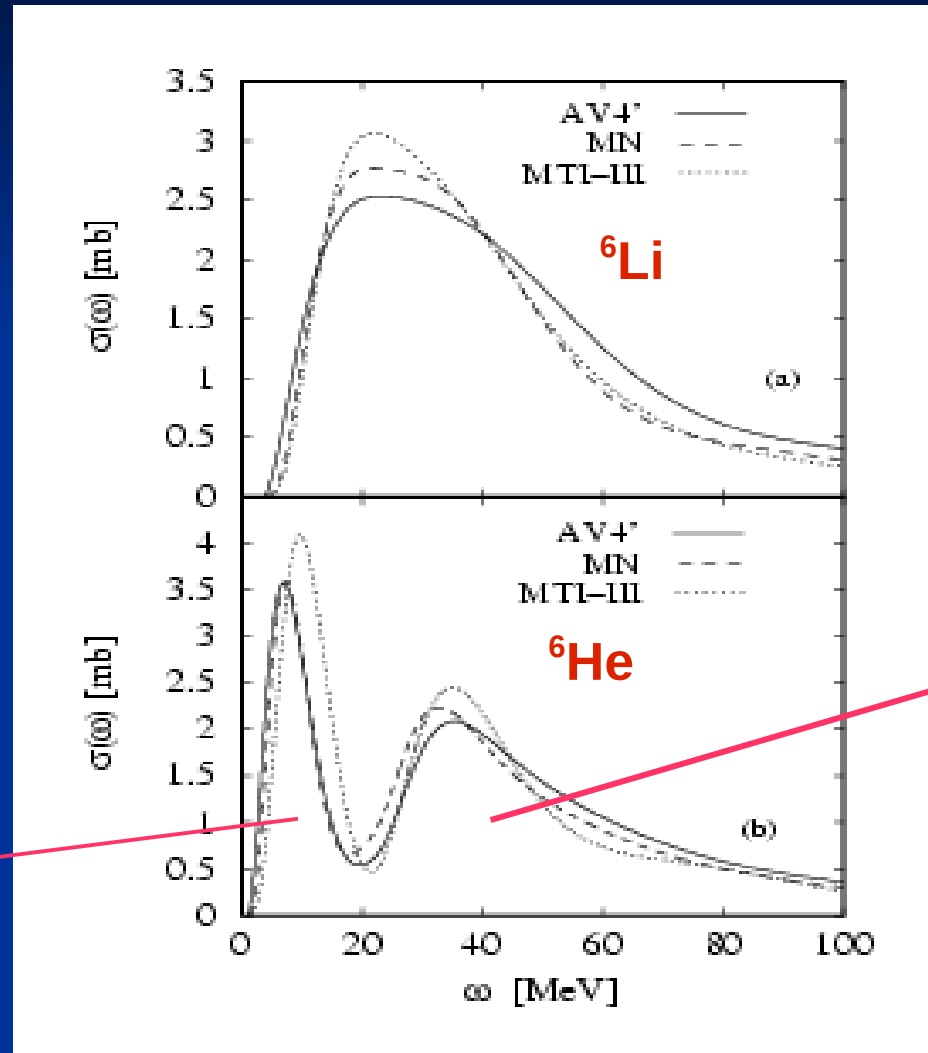
S.Bacca et al.  
PRL 89 (2002) 052502,  
PRC 69 (2004), 057001

**EIHH**

Semirealistic NN  
potentials

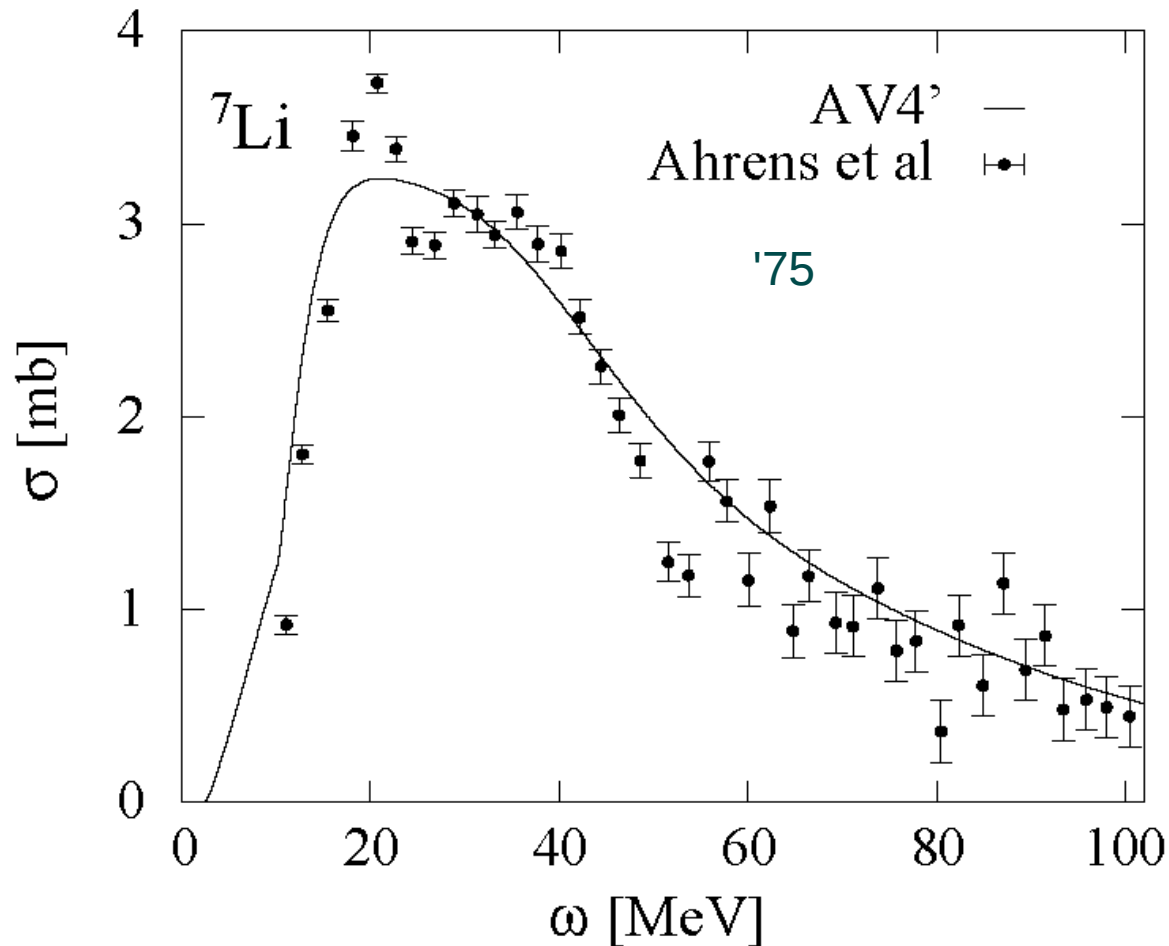


**soft mode**



**classical GT mode**

# 7-Body total photodisintegration

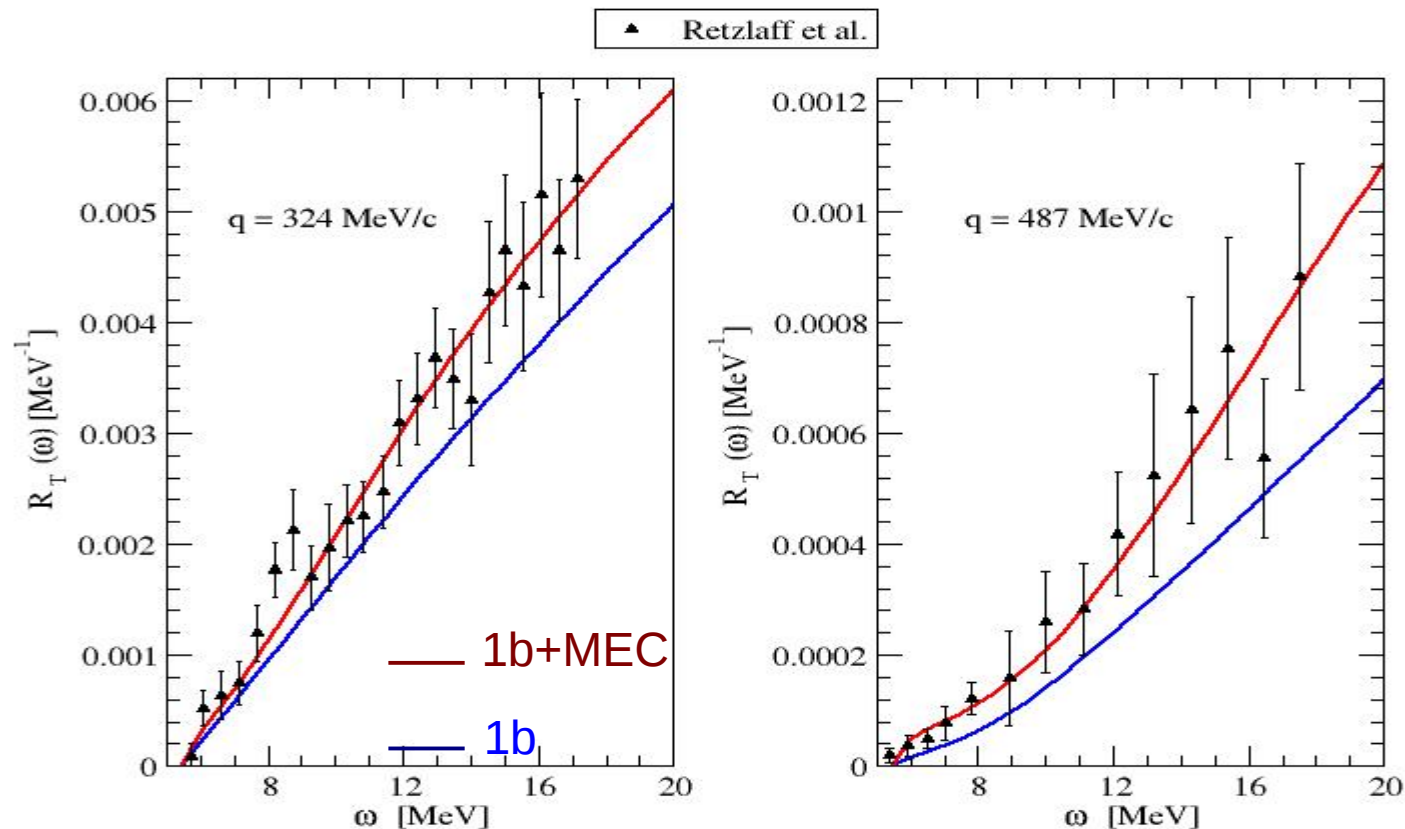


**S.Bacca et al.**  
PLB 603(2004) 159

**EIHH**

Semirealistic NN  
potential

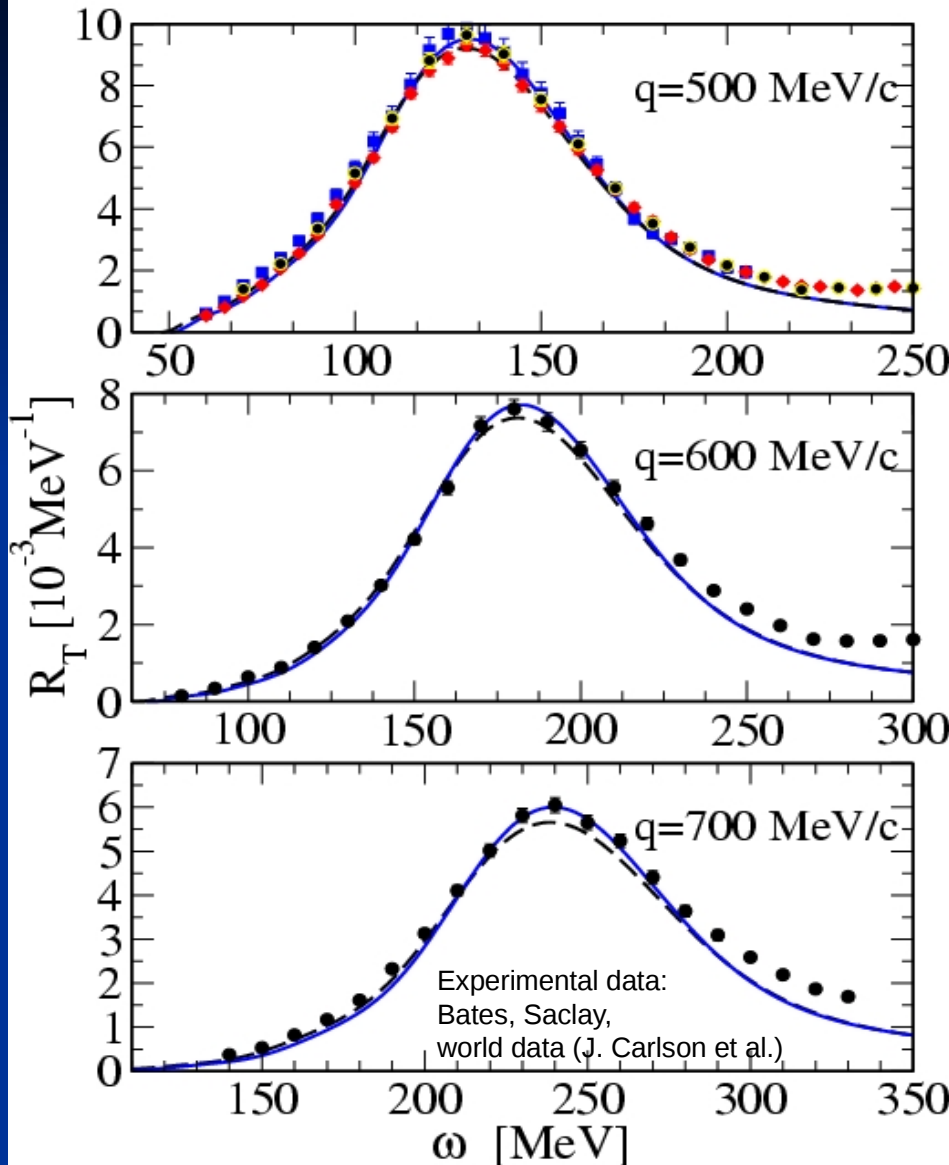
# Inelastic transverse $^3\text{He}$ (e,e') response function $R_T(q,\omega)$



V. Efros et al., FBS 47, 157 (2010)

# Inelastic transverse $^3\text{He}$ (e,e') response function $R_T(q,\omega)$

L. Yuan et al., PLB 706, 90 (2011)

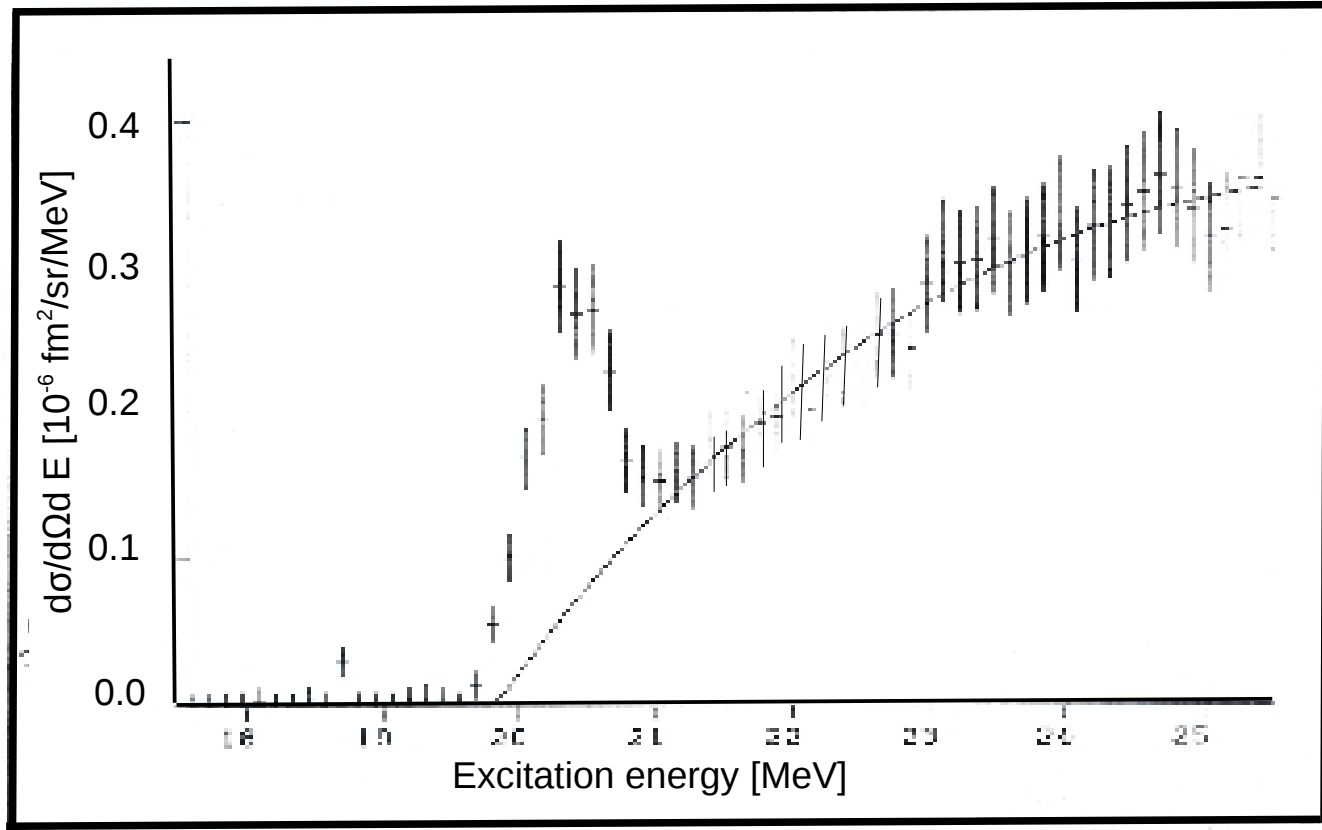




# Resonances

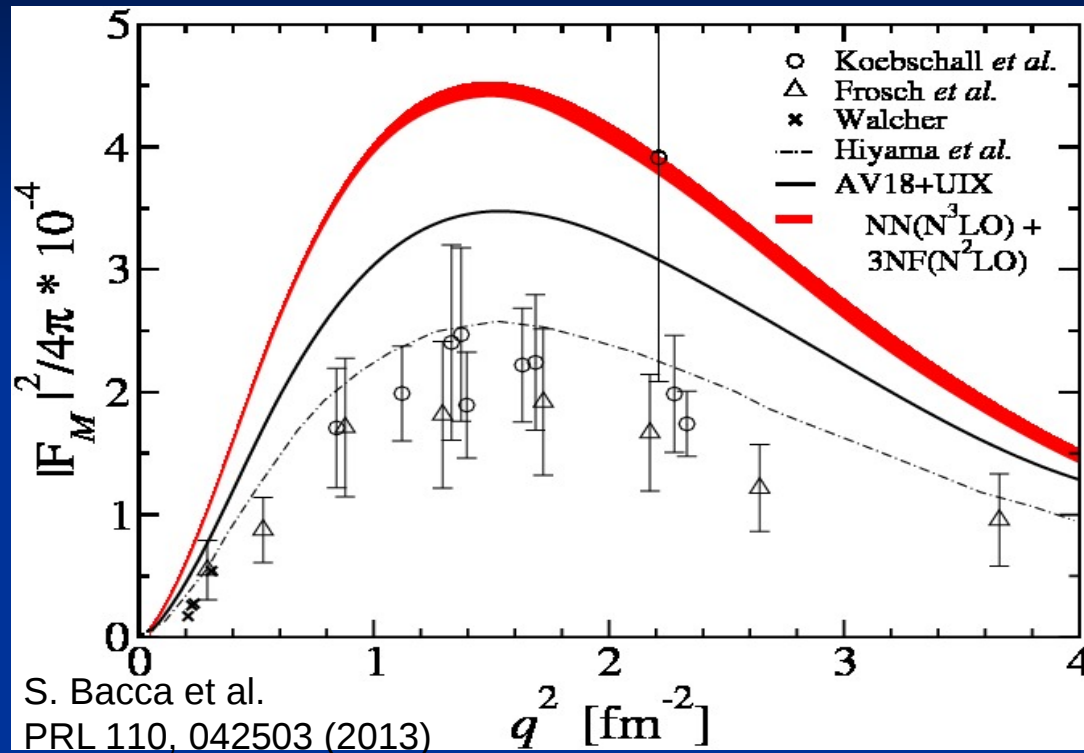
# $0^+$ Resonance in the $^4\text{He}$ compound system

*G. Köbschall et al./ Quasi bound state in  $^4\text{He}$  - Nucl. Phys. A405, 648 (1983)*



Resonance at  $E_R = -8.2 \text{ MeV}$ , i.e. above the  $^3\text{H}$ -p threshold. **Strong evidence** in electron scattering off  $^4\text{He}$ ,  $\Gamma = 270 \pm 50 \text{ keV}$

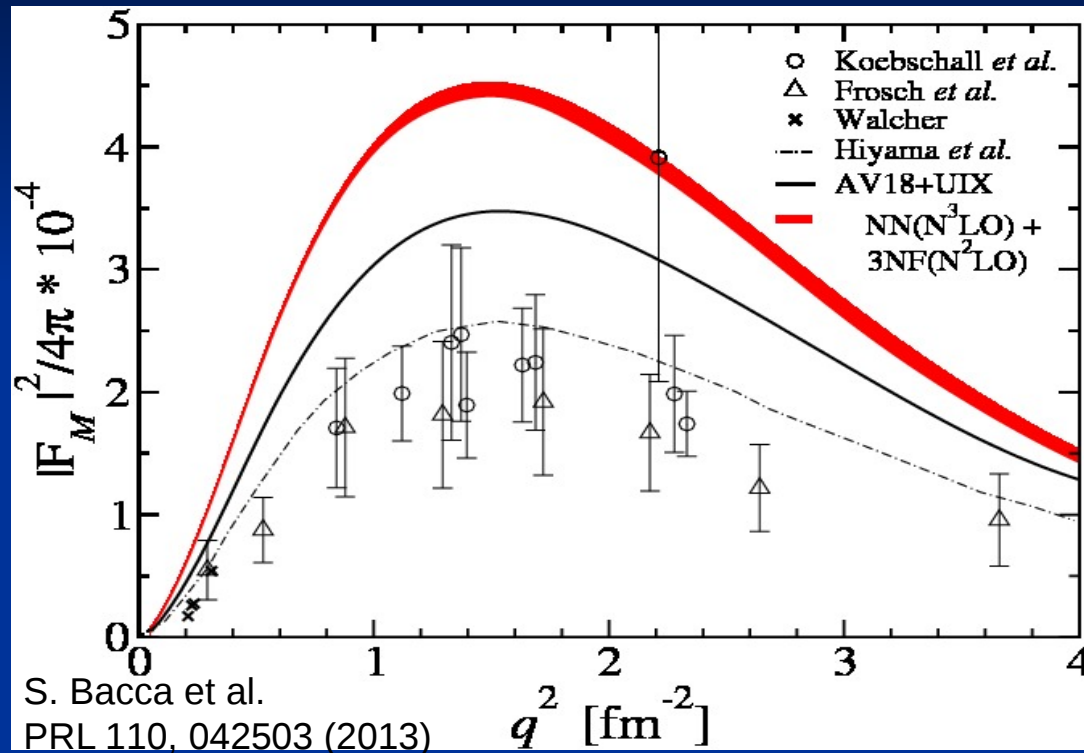
# Comparison to experimental results



LIT/EIHH Calculation for AV18+UIX and Idaho-N3LO+N2LO

dash-dotted: AV8' + central 3NF (Hiyama et al.)

# Comparison to experimental results



Observable is strongly dependent on potential model

Breathing Mode? (S. Bacca et al., PRC 91, 024303 (2015))

- Resonance strength determined with LIT method
- But resonance width undetermined
- Why ?

# LIT Method

# LIT method

The LIT of a function  $R(E)$  is defined as follows

$$\Rightarrow L(\sigma) = \int dE \mathcal{L}(E, \sigma) R(E) ,$$

where the kernel  $\mathcal{L}$  is a Lorentzian,

$$\Rightarrow \mathcal{L}(E, \sigma) = \frac{1}{(E - \sigma_R)^2 + \sigma_I^2}$$

For inclusive reactions the LIT  $L(\sigma)$  is calculated by solving an equation of the form

$$(H - \sigma) \tilde{\Psi} = S ,$$

where  $H$  is the Hamiltonian of the system under consideration and  $S$  is an asymptotically vanishing source term related to the operator inducing the specific reaction.

The solution  $\tilde{\Psi}$  is localized and the LIT is given by

$$L(\sigma) = \langle \tilde{\Psi} | \tilde{\Psi} \rangle .$$

Alternative way:

$$L(\sigma) = -\frac{1}{\sigma_I} \text{Im}(\langle S | \frac{1}{\sigma_R + i\sigma_I - H} | S \rangle) .$$

The source term  $S$  for inclusive reactions has the form

$$\Rightarrow |S\rangle = \theta|0\rangle,$$

where the operator  $\theta$  induces a specific electroweak reaction.

The corresponding response function is given by

$$\Rightarrow R(E_f) = \int dE_f |\langle f|\theta|0\rangle|^2 \delta(E_f - E_0 - \omega)$$

Ingredients of the solution of the LIT equation via an expansion on a basis of dimension  $N$ :

$N$  eigenstates with eigenenergies

$$\phi_n$$

$$E_n$$

and strength

**LIT states**

$$S_n = |\langle \phi_n | \theta | 0 \rangle|^2 \quad \Leftarrow \text{strength to a LIT state}$$

leading to the following LIT

$$\Rightarrow L(\sigma) = \sum_{i=1}^N \frac{S_n}{(\sigma_R - E_n)^2 + \sigma_I^2}$$



# main point of the LIT :

Schrödinger-like equation with a source

$$(H - E_0 - \omega_0 + i\Gamma) \tilde{\Psi} = S$$

The  $\tilde{\Psi}$  solution is unique and has **bound state like** asymptotic behavior



one can apply **bound state methods**

# Hyperspherical Harmonics (HH)

Expansion of ground-state wave function and LIT states on HH basis

HH:  $(A-1)$  Jacobi coordinates  $\rightarrow$  1 hyperradius +  $3A-4$  hyperangles

Acceleration of Convergence

with two-body correlations: CHH

or with effective interaction (Lee-Suzuki approach): EIHH

# Inversion of the LIT

- LIT is calculated for a fixed  $\sigma_l$  in many  $\sigma_r$  points
- Express the searched response function formally on a basis set with  $M$  basis functions  $f_m(E)$  and open coefficients  $c_m$  with correct threshold behaviour for the  $f_m(E)$  (e.g.,  $f_m = f_{thr}(E) \exp(-\alpha E/m)$ )
- Make a LIT transform of the basis functions and determine coefficients  $c_m$  by a fit to the calculated LIT
- Increase  $M$  up to the point that a sufficient convergence is obtained (uncontrolled oscillations should not be present)

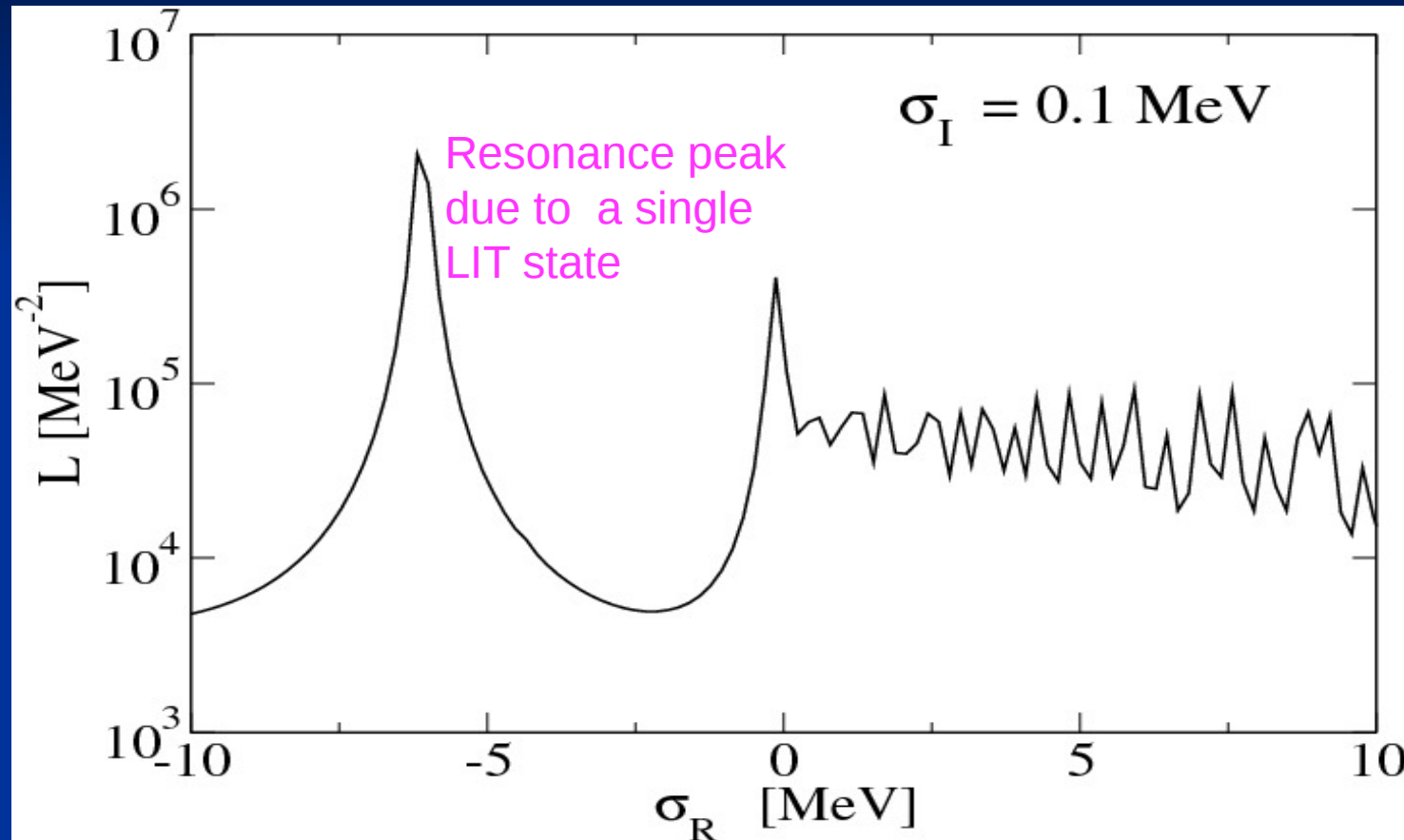
**A regularization method is needed for the inversion**

Why were we unable to determine the width of the  $^4\text{He}$   
isoscalar monopole resonance?

To answer this let us check our very first LIT calculation from 1997:

$^4\text{He}(e,e')$  inelastic longitudinal response function  
with a central NN potential

Unpublished result from a CHH calculation with the TN potential (V. Efros, WL, G. Orlandini, PRL 78,432 (1997))



To study the problem better let us consider first instead of a four-body reaction a **simpler three-body reaction**:



at low energies

LIT calculation with central MTI/III NN potential in unretarded dipole approximation

**Aim:** Increase low-energy density of LIT states

To study the problem better let us consider first instead of a four-body reaction a **simpler three-body reaction**:



at low energies

LIT calculation with central MTI/III NN potential in unretarded dipole approximation

**How:** Increase number of basis states, both, hyperradial and hyperspherical

To study the problem better let us consider first instead of a four-body reaction a **simpler three-body reaction**:



at low energies

LIT calculation with central MTI/III NN potential in unretarded dipole approximation

**Also note:** hyperradial basis states consist in an expansion on Laguerre polynomials times a spatial cutoff  $\exp(-\rho/b)$   
Increase of **b** shifts spectrum to lower energies



To study the problem better let us consider first instead of a four-body reaction a **simpler three-body reaction**:

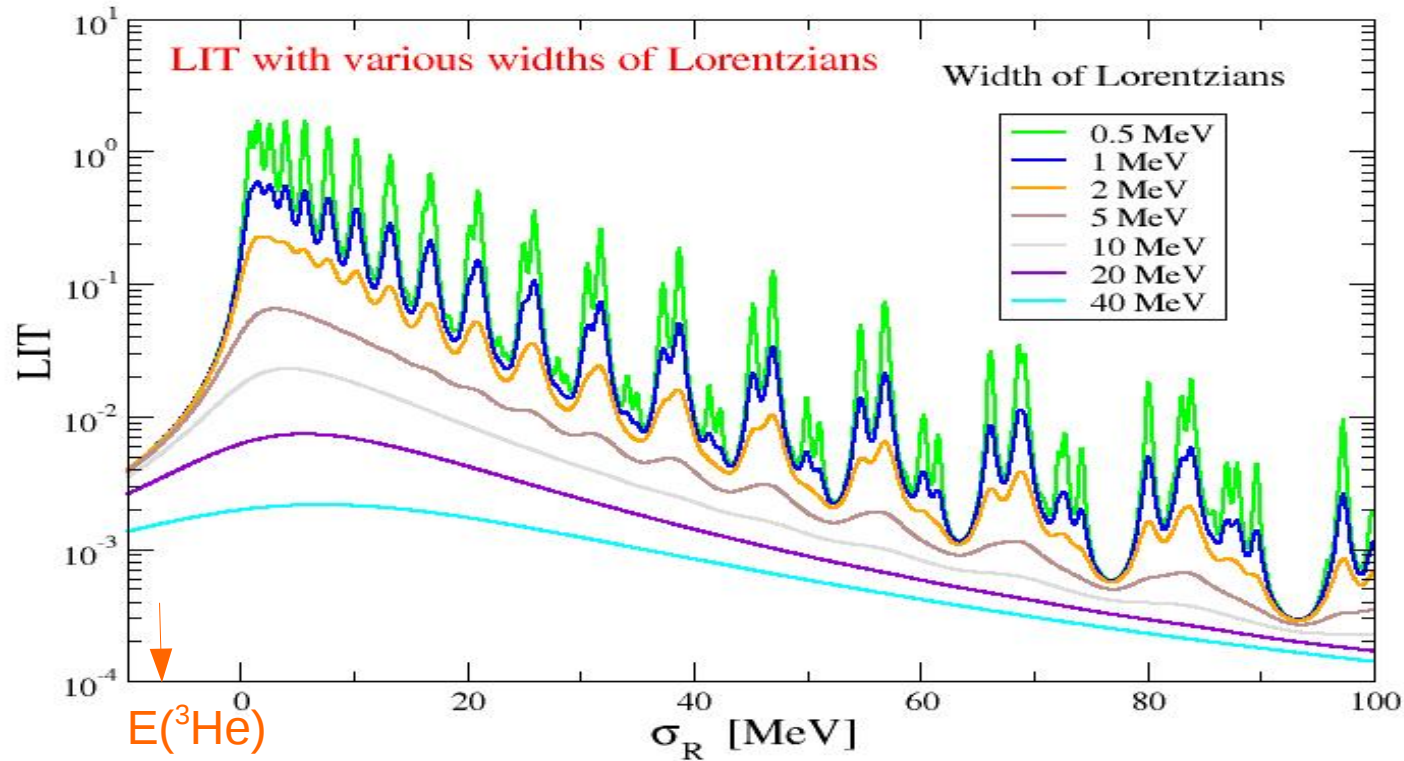


at low energies

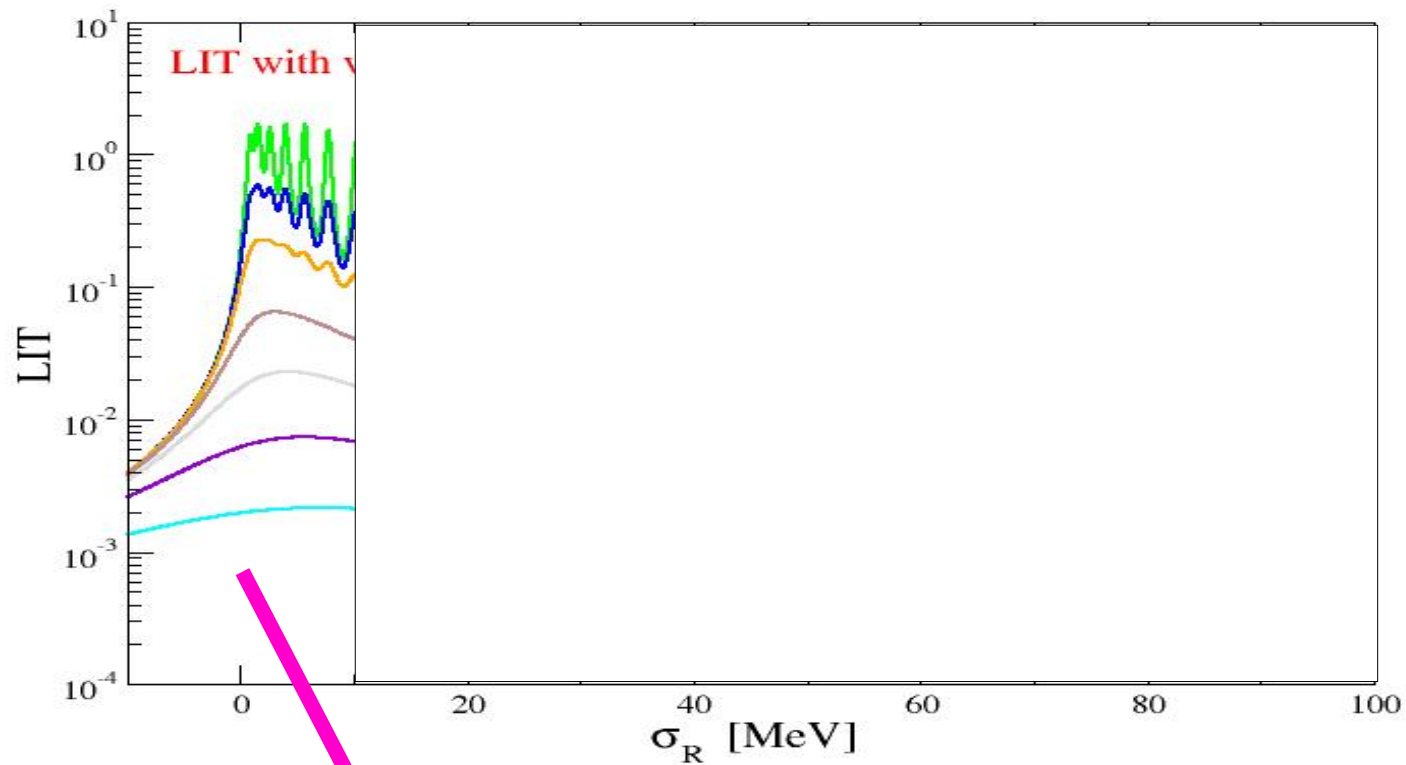
LIT calculation with central MTI/III NN potential in unretarded dipole approximation

**Next slide:** LIT with 30 hyperspherical and 31 hyperradial basis functions  $\Rightarrow$  **930 basis states** with  $b = 0.6$  fm

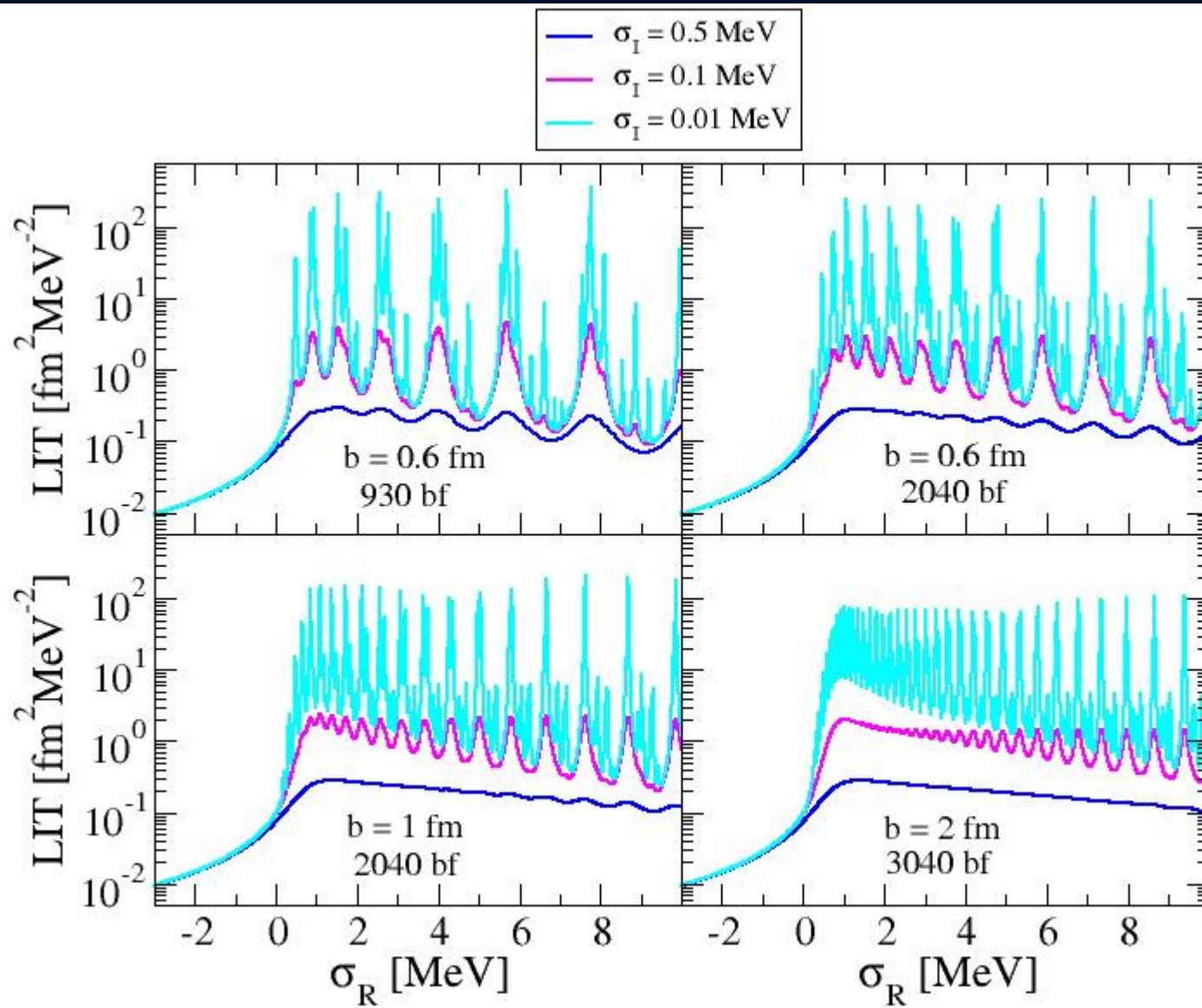
## LIT with various widths of Lorentzians



30 hyperspherical and 31 hyperradial basis functions  
 $\Rightarrow$  930 basis states  
 $b = 0.6$  fm



Increase LIT state density and **zOOM** in



# Observation

The LIT is a method with a controlled resolution

**But not a single LIT state** below three-body breakup threshold  
In present LIT calculation! Similar problem as in the previous four-body case

Solution: use instead of the HH basis a somewhat modified basis

# New A-body basis

Note one of the (A-1) Jacobi vectors can be written in the following form:

$$\eta = \mathbf{r}_A - \mathbf{R}_{\text{cm}}(1,2,\dots,A-1)$$

This is the coordinate one would use for the scattering of a nucleon with a (A-1)-nucleon system. In other words the relevant coordinate for a two-body breakup. Therefore

A-body HH basis  $\longrightarrow$  (A-1)-body HH basis times expansion on  $\eta$   
radial part: Laguerre polynomials  
angular part:  $Y_{\text{LM}}(\theta_\eta, \phi_\eta)$

Four-body system: HH for 3 particles plus 4-th particle coordinate  $\eta$

# New A-body basis

Note one of the (A-1) Jacobi vectors can be written in the following form:

$$\eta = \mathbf{r}_A - \mathbf{R}_{\text{cm}}(1,2,\dots,A-1)$$

This is the coordinate one would use for the scattering of a nucleon with a (A-1)-nucleon system. In other words the relevant coordinate for a two-body breakup. Therefore

A-body HH basis  $\longrightarrow$  (A-1)-body HH basis times expansion on  $\eta$   
radial part: Laguerre polynomials  
angular part:  $Y_{\text{LM}}(\theta_\eta, \phi_\eta)$

Three-body system: pair coordinate for two particles plus 3rd particle coordinate  $\eta$

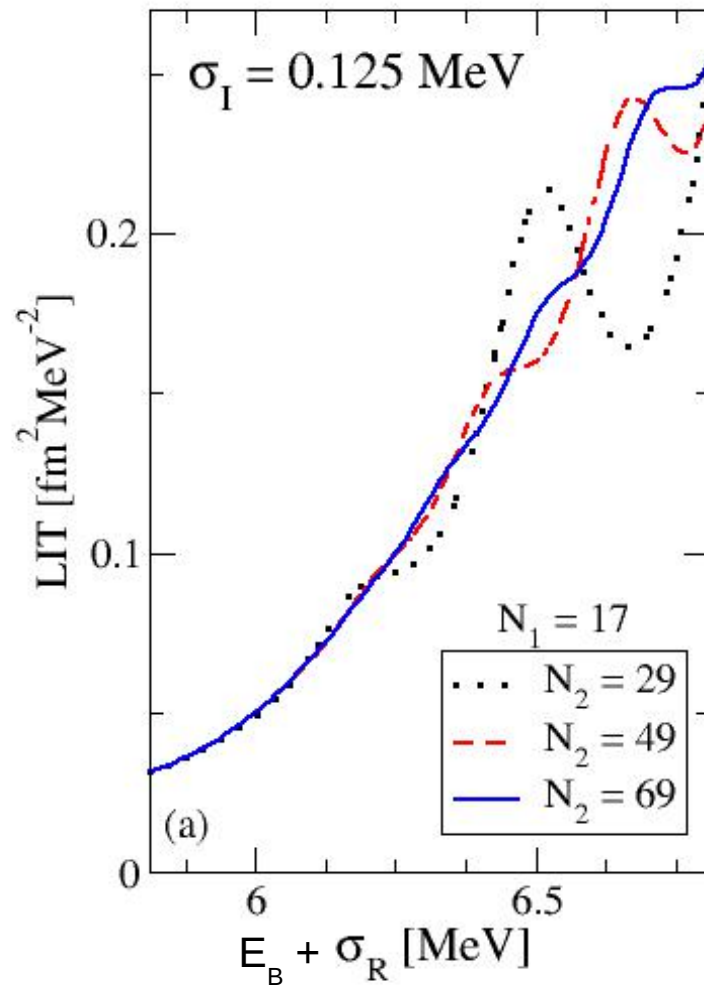
## First three-body case



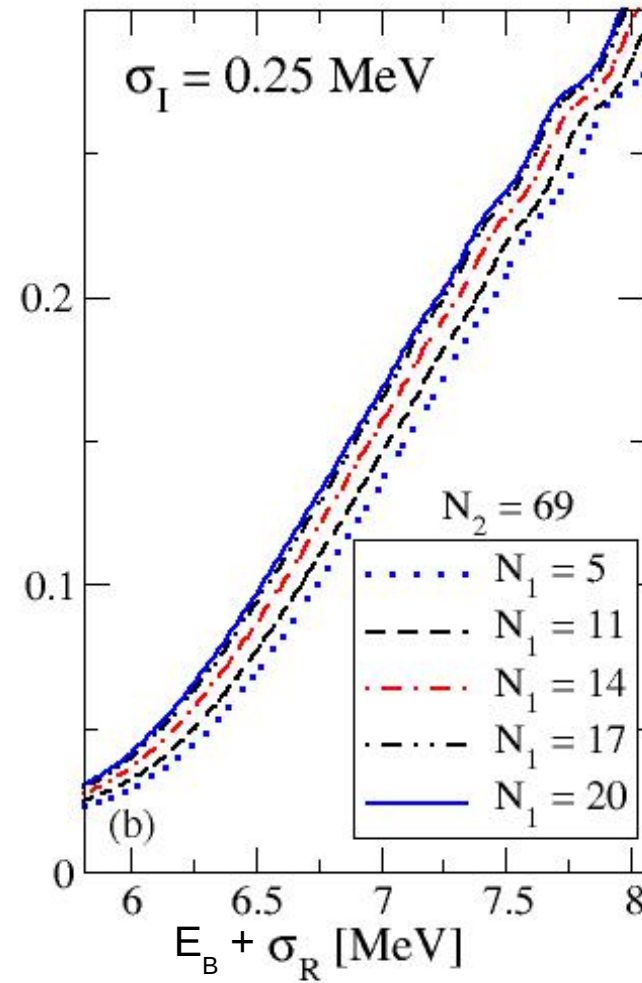
With convergence for expansions in pair and third particle coordinate



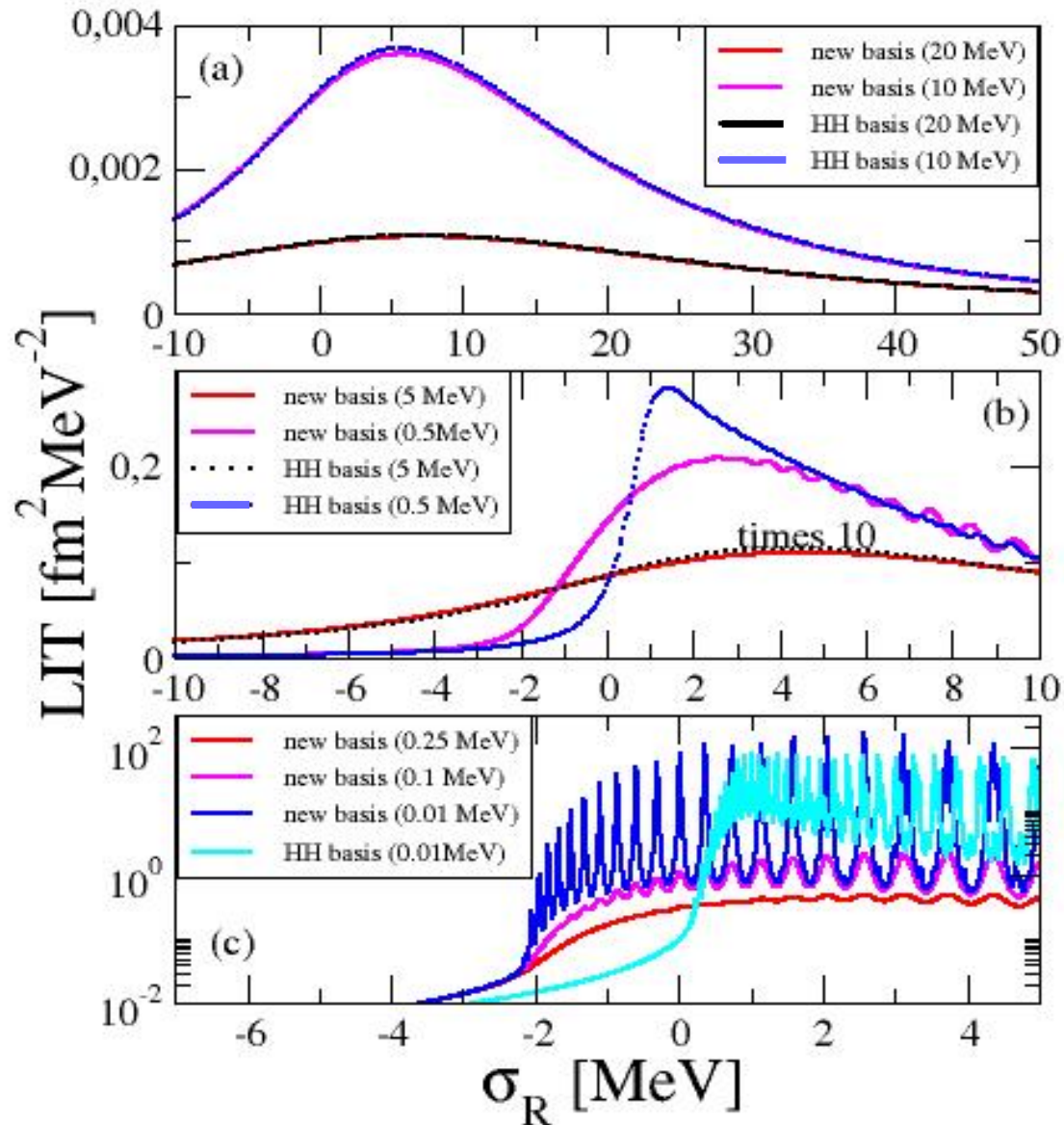
Third particle coordinate



Pair coordinate



# LIT results with HH and new basis

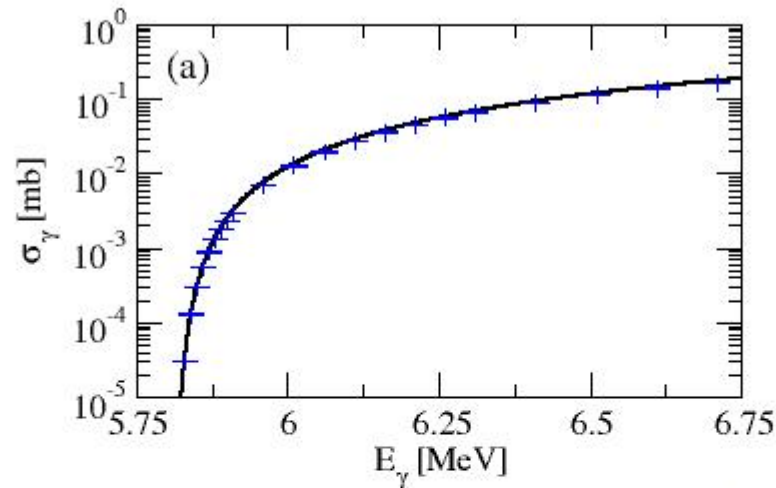


# Inversions

Implement correct threshold behaviour for  ${}^3\text{He} + \gamma \rightarrow \text{d} + \text{p}$

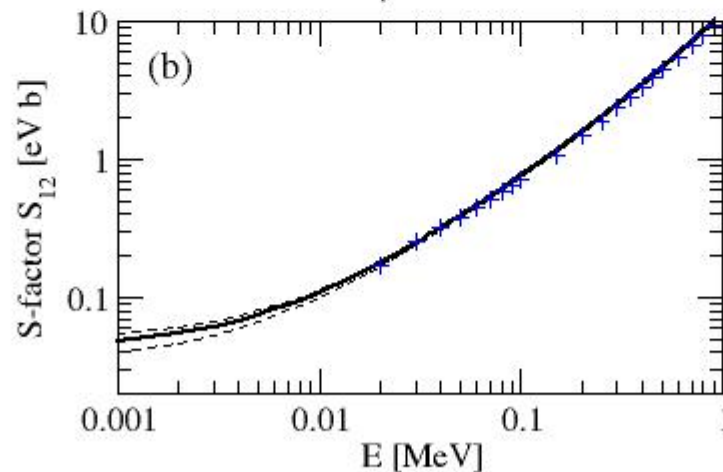
Due to Coulomb potential: usual Gamow factor

## Comparison with explicit calculation of continuum state



Cross section  
 ${}^3\text{He} + \gamma \longrightarrow \text{d} + \text{p}$

LIT: full curves  
 cont. wf: +



S-factor  
 $\text{d} + \text{p} \longrightarrow {}^3\text{He} + \gamma$

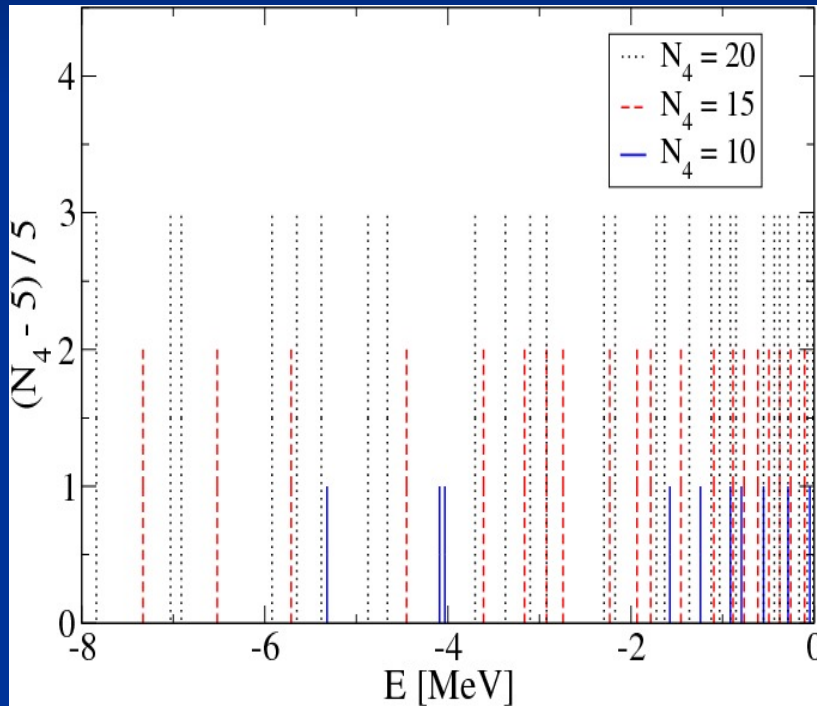
Error due to  
 inversion: dashed

(Standard deviation from inversions with 11-18 basis functions) S. Deflorian, V. Efros, WL, FBS 58:3 (2017)

# Back to the $^4\text{He}$ resonance

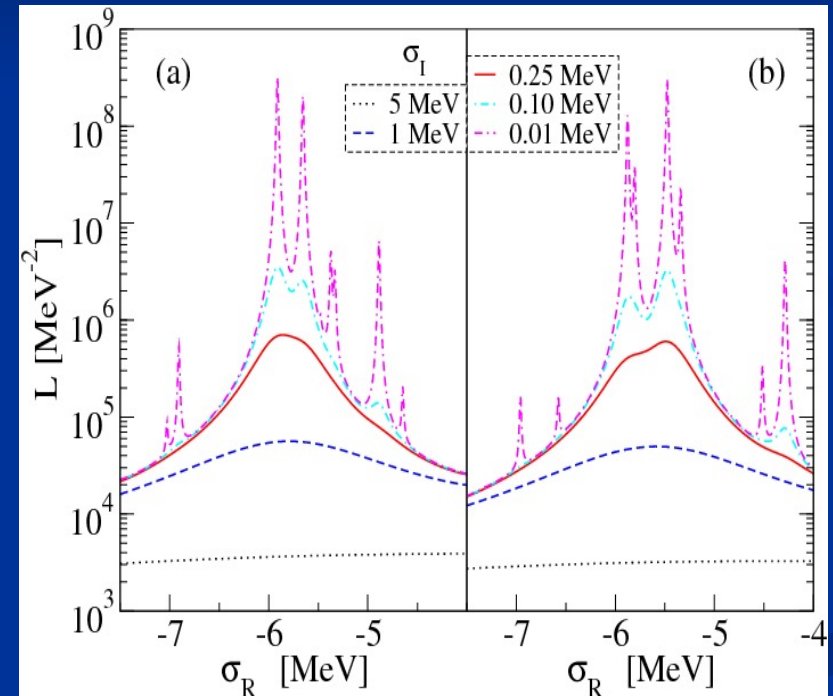
# Results with new basis

Number of basis functions in 4-th particle coordinate



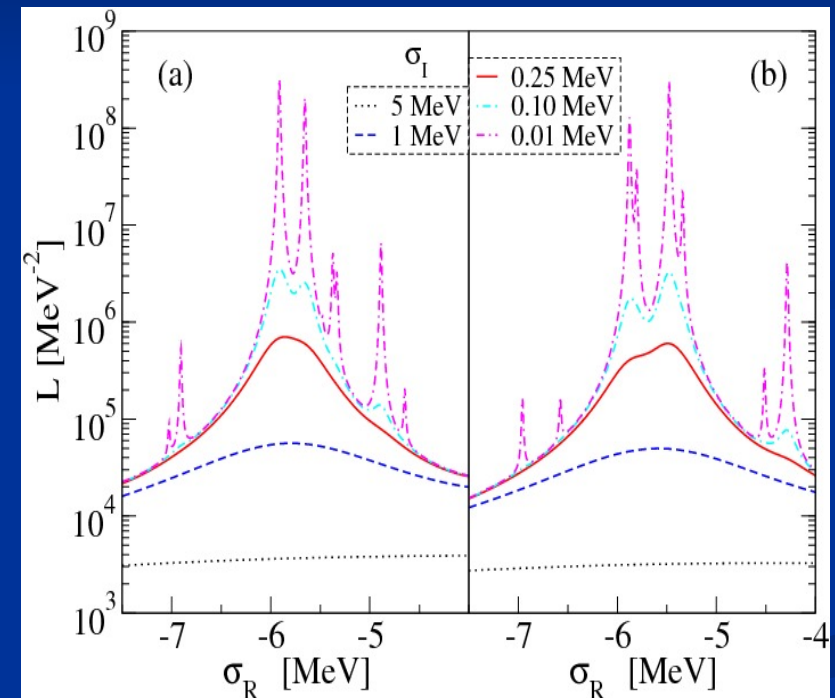
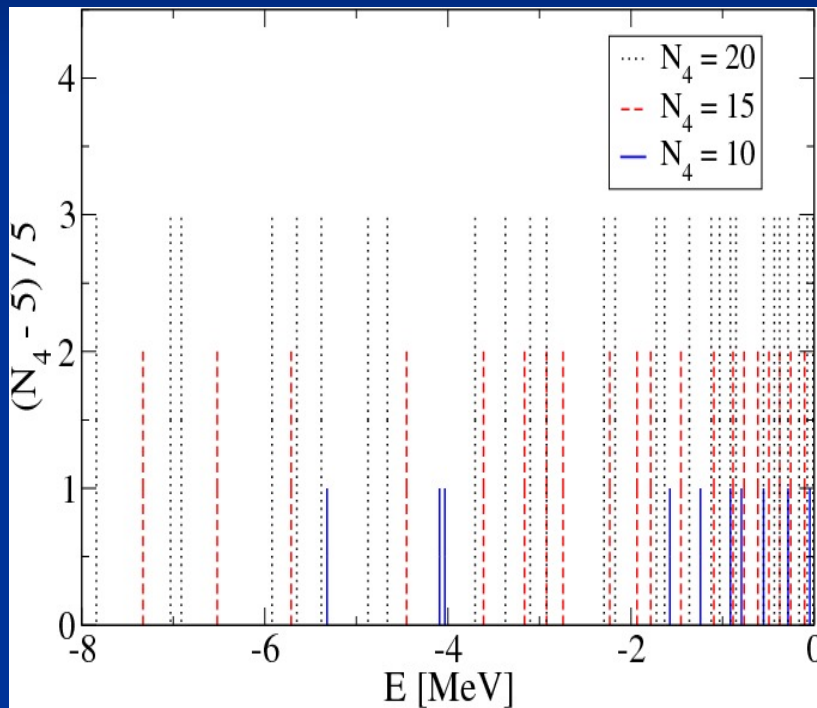
$N_4 = 20$

$N_4 = 21$



LIT

# Results with new basis



**Inversion:  $\Gamma = 180(70)$  keV**

WL, PRC 91, 054001 (2015)

# Summary and Outlook

- LIT method: ab initio method for reaction cross sections
- Extension to calculations for  $A > 4$  with realistic nuclear interaction possible
  - HH but not beyond  $A = 10$
  - other ab initio methods (coupled cluster)
  - Halo EFT